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Stability in the energy space for chains of solitons of the one-dimensional Gross-Pitaevskii equation

Fabrice Béthuel ¹, Philippe Gravejat ², Didier Smets ³

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Abstract

We establish the stability in the energy space for sums of solitons of the one-dimensional Gross-Pitaevskii equation when their speeds are mutually distinct and distinct from zero, and when the solitons are initially well-separated and spatially ordered according to their speeds.

1 Introduction

In this paper, we focus on the one-dimensional Gross-Pitaevskii equation

$$i\partial_t\Psi + \partial_x^2\Psi + \Psi(1 - |\Psi|^2) = 0, \quad (\text{GP})$$

for a function $\Psi : \mathbb{R} \times \mathbb{R} \rightarrow \mathbb{C}$, supplemented with the boundary condition at infinity

$$|\Psi(x, t)| \rightarrow 1, \quad \text{as } |x| \rightarrow +\infty. \quad (1)$$

The Gross-Pitaevskii equation is a defocusing nonlinear Schrödinger equation. In one spatial dimension, solutions with finite Ginzburg-Landau energy

$$\mathcal{E}(\Psi) := \frac{1}{2} \int_{\mathbb{R}} |\partial_x \Psi|^2 + \frac{1}{4} \int_{\mathbb{R}} (1 - |\Psi|^2)^2,$$

are globally defined, and the energy is conserved along the flow.

In 1972, V.E. Zakharov and A.B. Shabat [20] have shown that the Gross-Pitaevskii equation supplemented with condition (1) is integrable by means of an inverse scattering transform. As a consequence, they exhibited exact soliton and multi-soliton solutions. It is generally believed that such solutions play a central role in the long-time asymptotics of the flow, and sometimes conjectured that any initial datum with finite Ginzburg-Landau energy should eventually resolve into a multi-soliton solution plus a vanishing (in an appropriate sense) dispersive tail. Detailed asymptotic expansions for smooth and fast decaying solutions, based on their scattering data at the initial time, have actually been obtained by A.H. Vartanian in [19], using reductions to Riemann-Hilbert problems. To our knowledge, for initial data in the energy space, there is no rigorous evidence supporting the aforementioned behaviour further than the orbital stability of

¹Laboratoire Jacques-Louis Lions, Université Pierre et Marie Curie, Boîte Courrier 187, 75252 Paris Cedex 05, France. E-mail: bethuel@ann.jussieu.fr

²Centre de Mathématiques Laurent Schwartz, École Polytechnique, 91128 Palaiseau Cedex, France. E-mail: gravejat@math.polytechnique.fr

³Laboratoire Jacques-Louis Lions, Université Pierre et Marie Curie, Boîte Courrier 187, 75252 Paris Cedex 05, France. E-mail: smets@ann.jussieu.fr

a single soliton (see e.g. [12, 1, 2, 10]). The main goal of this paper is to provide a partial justification through the proof of the orbital stability in the energy space of some finite sums of solitons along the Gross-Pitaevskii flow. Our approach does not make any use of the integrability of the equation, and is largely influenced by the corresponding study of the generalized Korteweg-de Vries equation (see [15]) and of the nonlinear Schrödinger equation with vanishing data at infinity (see [16]) by Y. Martel, F. Merle and T.-P. Tsai. The main ingredient is a monotonicity formula for the momentum (see Subsection 3.2).

1.1 Soliton solutions to the Gross-Pitaevskii equation

A soliton with speed c is a solution to (GP) which takes the form

$$\Psi(x, t) := U_c(x - ct).$$

Its profile U_c is a solution to the ordinary differential equation

$$-ic\partial_x U_c + \partial_{xx}^2 U_c + U_c(1 - |U_c|^2) = 0. \quad (2)$$

The solutions to (2) with finite energy are classified and indeed explicit. For $|c| \geq \sqrt{2}$, all of them are identically constant. When $|c| < \sqrt{2}$, there exist non-constant solutions with finite energy. Up to the invariances of the problem, i.e. the multiplication by a constant of modulus one and the invariance by translation, they are uniquely given by the expression

$$U_c(x) = \sqrt{\frac{2-c^2}{2}} \operatorname{th}\left(\frac{\sqrt{2-c^2}x}{2}\right) + i\frac{c}{\sqrt{2}}. \quad (3)$$

Notice that the soliton U_c does not vanish on \mathbb{R} when $c \neq 0$. As a matter of fact, our stability analysis below requires non-vanishing solutions. Indeed, it is performed on a reformulation of (GP) which only makes sense for such solutions, and which we introduce next.

1.2 Hydrodynamical form of the Gross-Pitaevskii equation

Provided a solution Ψ to (GP) does not vanish, it may be written, at least formally, as

$$\Psi := \varrho \exp i\varphi,$$

where $\varrho := |\Psi|$. Still on a formal level, the functions $\eta := 1 - \varrho^2$ and $v := \partial_x \varphi$ are then solutions to an hydrodynamical form of (GP), namely

$$\begin{cases} \partial_t \eta = \partial_x (2v - 2\eta v), \\ \partial_t v = \partial_x \left(\eta - v^2 - \partial_x \left(\frac{\partial_x \eta}{2(1-\eta)} \right) + \frac{(\partial_x \eta)^2}{4(1-\eta)^2} \right). \end{cases} \quad (\text{HGP})$$

For $k \geq 0$, we introduce the Hilbert spaces $X^k(\mathbb{R}) := H^{k+1}(\mathbb{R}) \times H^k(\mathbb{R})$, equipped with the norm

$$\|(\eta, v)\|_{X^k}^2 := \|\eta\|_{H^{k+1}}^2 + \|v\|_{H^k}^2,$$

and their open subsets ¹

$$\mathcal{NV}^k(\mathbb{R}) := \{(\eta, v) \in X^k(\mathbb{R}), \text{ s.t. } \max_{x \in \mathbb{R}} \eta(x) < 1\},$$

¹The notation \mathcal{NV} stands for non-vanishing.

consisting of states with non-vanishing corresponding Ψ . In the sequel, we set $X(\mathbb{R}) := X^0(\mathbb{R})$ and $\mathcal{NV}(\mathbb{R}) := \mathcal{NV}^0(\mathbb{R})$.

The Ginzburg-Landau energy ²

$$E(\eta, v) := \int_{\mathbb{R}} e(\eta, v) := \frac{1}{8} \int_{\mathbb{R}} \frac{(\partial_x \eta)^2}{1 - \eta} + \frac{1}{2} \int_{\mathbb{R}} (1 - \eta)v^2 + \frac{1}{4} \int_{\mathbb{R}} \eta^2, \quad (4)$$

and the momentum

$$P(\eta, v) := \int_{\mathbb{R}} p(\eta, v) := \frac{1}{2} \int_{\mathbb{R}} \eta v,$$

are well-defined and smooth functionals on $\mathcal{NV}(\mathbb{R})$. In view of (4), the space $\mathcal{NV}(\mathbb{R})$ is the energy space for states (η, v) satisfying the non-vanishing condition $\max_{x \in \mathbb{R}} \eta(x) < 1$.

Concerning the Cauchy theory for (HGP), we have

Theorem 1 ([18]). *Let $k \in \mathbb{N}$ and $(\eta^0, v^0) \in \mathcal{NV}^k(\mathbb{R})$. There exists a maximal time $T_{\max} > 0$ and a unique solution $(\eta, v) \in \mathcal{C}^0([0, T_{\max}), \mathcal{NV}^k(\mathbb{R}))$ to equation (HGP) with initial datum (η^0, v^0) . The maximal time T_{\max} is continuous with respect to the initial datum (η^0, v^0) , and is characterized by*

$$\lim_{t \rightarrow T_{\max}} \max_{x \in \mathbb{R}} \eta(x) = 1 \quad \text{if} \quad T_{\max} < +\infty.$$

The flow map $(\eta^0, v^0) \mapsto (\eta, v)$ is locally well-defined and continuous from $\mathcal{NV}^k(\mathbb{R})$ to $\mathcal{C}^0([0, T], \mathcal{NV}^k(\mathbb{R}))$ for any $T < T_{\max}$. Moreover, the energy E and the momentum P are constant along the flow.

We refer to [18] for a proof. It is a consequence of the Cauchy theory for (GP) in Zhidkov spaces $Z^k(\mathbb{R})$ (see e.g. [21, 8, 9]), and the fact that the mapping $\Psi \mapsto (\eta, v)$ is locally Lipschitz from $Z^k(\mathbb{R})$ to $X^k(\mathbb{R})$ provided that $(\eta, v) \in \mathcal{NV}(\mathbb{R})$. The uniqueness part requires slightly more care.

1.3 Statement of the main result

When $c \neq 0$, the solitons U_c described in Subsection 1.1 do not vanish and may thus be written under the form

$$U_c := \varrho_c \exp i\varphi_c.$$

In view of formula (3), the maps $\eta_c := 1 - \varrho_c^2$ and $v_c := \partial_x \varphi_c$ are given by ³

$$\eta_c(x) = \frac{2 - c^2}{2\text{ch}\left(\frac{\sqrt{2-c^2}x}{2}\right)^2} \quad \text{and} \quad v_c(x) = -\frac{c\eta_c(x)}{2(1 - \eta_c(x))} = -\frac{c(2 - c^2)}{2\left(2\text{ch}\left(\frac{\sqrt{2-c^2}x}{2}\right)^2 - 2 + c^2\right)}. \quad (5)$$

²When Ψ may be written as $\Psi = \varrho \exp i\varphi$, the quantities $\mathcal{E}(\Psi)$ and $E(\eta, v)$, with $\eta = 1 - \varrho^2$ and $v = \partial_x \varphi$, are exactly equal.

³A striking point in formula (5) is the fact that, up to some scaling, the map η_c is exactly the soliton of the Korteweg-de Vries equation

$$\partial_t u + \partial_x^3 u + u\partial_x u = 0. \quad (\text{KdV})$$

Like the Gross-Pitaevskii equation, the Korteweg-de Vries equation indeed owns a family of solitons u_σ , given, up to translations, by

$$u_\sigma(x) := \frac{3\sigma}{\text{ch}\left(\frac{\sqrt{\sigma}x}{2}\right)^2},$$

for some positive speed σ . Actually, the link between the two equations does not reduce to the level of the solitons. Both of them are integrable by means of the inverse scattering transform (see e.g. [17] for the Korteweg-de Vries equation), and moreover, as was pointed out in [2, 3] (see also [5]), the Korteweg-de Vries equation appears as the limit of system (HGP) in some long-wave regime.

In the sequel, we set

$$Q_{c,a} := (\eta_{c,a}, v_{c,a}) := (\eta_c(\cdot - a), v_c(\cdot - a)),$$

for $0 < |c| < \sqrt{2}$ and $a \in \mathbb{R}$. We introduce the set of admissible speeds as

$$\text{Adm}_N := \{\mathbf{c} := (c_1, \dots, c_N) \in (-\sqrt{2}, \sqrt{2})^N, \text{ s.t. } c_k \neq 0, \forall 1 \leq k \leq N\},$$

and, for $\mathbf{c} \in \text{Adm}_N$ and $\mathbf{a} := (a_1, \dots, a_N) \in \mathbb{R}^N$, we define

$$R_{\mathbf{c},\mathbf{a}} := (\eta_{\mathbf{c},\mathbf{a}}, v_{\mathbf{c},\mathbf{a}}) = \sum_{k=1}^N Q_{c_k, a_k} = \left(\sum_{k=1}^N \eta_{c_k, a_k}, \sum_{k=1}^N v_{c_k, a_k} \right) = \left(\sum_{k=1}^N \eta_{c_k}(\cdot - a_k), \sum_{k=1}^N v_{c_k}(\cdot - a_k) \right).$$

In the sequel, we refer to small perturbations of the functions $R_{\mathbf{c},\mathbf{a}}$ as chains of solitons. We also denote

$$\nu_{\mathbf{c}} := \min \left\{ \sqrt{2 - c_k^2}, 1 \leq k \leq N \right\}.$$

Our main result is

Theorem 2. *Let $\mathbf{c}^* \in \text{Adm}_N$ be such that*

$$c_1^* < \dots < c_N^*. \quad (6)$$

There exist positive numbers α^ , L^* and A^* , depending ⁴ only on \mathbf{c}^* , such that the following holds. Assume that $(\eta^0, v^0) \in \mathcal{NV}(\mathbb{R})$ satisfies*

$$\alpha^0 := \left\| (\eta^0, v^0) - R_{\mathbf{c}^*, \mathbf{a}^0} \right\|_X \leq \alpha^*, \quad (7)$$

for some points $\mathbf{a}^0 = (a_1^0, \dots, a_N^0) \in \mathbb{R}^N$ such that

$$L^0 := \min \{ a_{k+1}^0 - a_k^0, 1 \leq k \leq N-1 \} \geq L^*. \quad (8)$$

Then the unique solution (η, v) to (HGP) with initial datum (η^0, v^0) is globally defined on $[0, +\infty)$, and there exists a function $\mathbf{a} = (a_1, \dots, a_N) \in \mathcal{C}^1([0, +\infty), \mathbb{R}^N)$ such that

$$\sum_{k=1}^N |a'_k(t) - c_k^*| \leq A^* \left(\alpha^0 + \exp \left(- \frac{\nu_{\mathbf{c}^*} L^0}{33} \right) \right), \quad (9)$$

and

$$\left\| (\eta(\cdot, t), v(\cdot, t)) - R_{\mathbf{c}^*, \mathbf{a}(t)} \right\|_X \leq A^* \left(\alpha^0 + \exp \left(- \frac{\nu_{\mathbf{c}^*} L^0}{33} \right) \right), \quad (10)$$

for any $t \in [0, +\infty)$.

In different terms, chains of solitons with strictly ordered speeds are preserved along the Gross-Pitaevskii flow in positive time. Since the initial speeds are strictly ordered, the mutual distances between the individual solitons appearing in the chain increase proportionally with time, so that the interactions between solitons decrease to zero. As a matter of fact, integrating (9) yields

$$a_{k+1}(t) - a_k(t) \geq \left(c_{k+1}^* - c_k^* - 2A^* \left(\alpha^0 + \exp \left(- \frac{\nu_{\mathbf{c}^*} L^0}{33} \right) \right) \right) t + a_{k+1}(0) - a_k(0).$$

⁴A close inspection of the proof actually shows that α^* , L^* and A^* only depend on a lower bound on $\nu_{\mathbf{c}^*}$, on $\min\{c_j^* - c_k^*, 1 \leq j < k \leq N\}$ and on $\min\{c_k^*, 1 \leq k \leq N\}$.

A typical example of chain of solitons is given by the exact multi-solitons constructed by Zakharov and Shabat [20] (see also [7]). In particular, Theorem 2 provides the orbital stability of multi-solitons in the energy space in positive time.

In view of the time reversibility of the Gross-Pitaevskii equation, Theorem 2 also provides the stability in negative time of chains of solitons with reversely ordered speeds. To our knowledge, the question of the stability in the energy space of chains of solitons for both positive and negative times remains open when the chain includes at least two solitons. For the generalized Korteweg-de Vries equations, results of this type have recently been obtained by Martel and Merle [13, 14].

Finally, we believe that our arguments still apply to non-integrable versions of the nonlinear Schrödinger equation with non-vanishing condition at infinity, provided individual solitons are themselves orbitally stable. The latter condition has recently been treated in details in a work by Chiron [4], where numerous examples are presented.

1.4 Elements in the proofs

We will present a detailed description of the elements in the proof of Theorem 2 in Subsections A to G below. We precede this description by a voluntarily oversimplified overview in the case of a two-soliton, in the hope that this may ease the reading of the detailed description.

The solitons Q_c may be characterized as minimizers of some variational problems. In terms of $Q_c = (\eta_c, v_c)$, equation (2) reduces to

$$E'(Q_c) + cP'(Q_c) = 0. \quad (11)$$

Equation (11) is the Euler-Lagrange equation for the minimization of the Ginzburg-Landau energy E under fixed scalar momentum P , and the speed c is the corresponding Lagrange multiplier. As a matter of fact, we have

$$E(Q_c) = \min \left\{ E(\eta, v), (\eta, v) \in X(\mathbb{R}) \text{ s.t. } \eta < 1 \text{ and } P(\eta, v) = P(Q_c) \right\}, \quad (12)$$

for any $c \in (-\sqrt{2}, \sqrt{2}) \setminus \{0\}$ (see e.g. [1]), and Q_c is the only minimizer up to the invariances of the problem. This minimizing property, combined with the conservation of E and P by the Gross-Pitaevskii flow, yields the orbital stability of single solitons. Instability of a *chain* of solitons, should it happen, could thus only arise from their mutual interactions. Assume for simplicity that at some fixed time the solution at hand is a chain of two solitons

$$(\eta, v) \simeq Q_{c_1, a_1} + Q_{c_2, a_2},$$

where $a_2 - a_1 \gg 1$. In particular, the total energy E and momentum P satisfy

$$E \simeq E(Q_{c_1}) + E(Q_{c_2}) \quad \text{and} \quad P \simeq P(Q_{c_1}) + P(Q_{c_2}).$$

Assume next that at some later time, interaction has occurred and

$$(\eta, v) \simeq B_1(\cdot - b_1) + B_2(\cdot - b_2),$$

where still $b_2 - b_1 \gg 1$, but B_1 and B_2 need no longer be solitons. Set

$$\delta P := P(Q_{c_1}) - P(B_1).$$

Since the total momentum is an invariant of the flow, we also have

$$\delta P \simeq P(B_2) - P(Q_{c_2}),$$

so that δP may be understood as the amount of momentum transferred from Q_{c_1, a_1} towards Q_{c_2, a_2} (we let aside here the possibility of transfer to the background).

From (12), we infer that the energies $E(B_1)$ and $E(B_2)$ are necessarily larger than the energies of the corresponding solitons with momentum $P(B_1)$, respectively $P(B_2)$. Since, in particular from (11),

$$\frac{d}{dc} \left(E(Q_c) \right) = -c \frac{d}{dc} \left(P(Q_c) \right),$$

we obtain at first order in δP ,

$$E(B_1) \gtrsim E(Q_{c_1}) + c_1 \delta P \quad (13)$$

and

$$E(B_2) \gtrsim E(Q_{c_2}) - c_2 \delta P. \quad (14)$$

On the other hand, by conservation of the total energy, we have

$$E(Q_{c_1}) + E(Q_{c_2}) \simeq E(B_1) + E(B_2),$$

so that, by (14),

$$E(B_1) \lesssim E(Q_{c_1}) + E(Q_{c_2}) - E(B_2) \lesssim E(Q_{c_1}) + c_2 \delta P. \quad (15)$$

Combining (13) and (15) yields

$$(c_2 - c_1) \delta P \gtrsim 0. \quad (16)$$

The last and key observation is a monotonicity formula, which states that in the above configuration necessarily

$$\delta P \lesssim 0. \quad (17)$$

Therefore, if the speeds are ordered as in Theorem 2, namely $c_2 > c_1$, then (16) may hold only if $\delta P = 0$. Orbital stability of the chain then follows from the orbital stability of the individual solitons.

Remark 1. The monotonicity formula (17) has a simple physical interpretation. Momentum is a signed quantity which, in the present normalization, is non-negative for “waves” or “particles” travelling to the left and non-positive for those travelling to the right. Since Q_{c_1, a_1} is located at the left of Q_{c_2, a_2} (recall that $a_2 - a_1 \gg 1$ by assumption), the momentum transferred from Q_{c_1, a_1} towards Q_{c_2, a_2} (i.e. δP) needs to travel from the left to the right,⁵ and is therefore non-positive. Equivalently (and alternatively from the point of view of Q_{c_2, a_2}), the momentum transferred from Q_{c_2, a_2} towards Q_{c_1, a_1} , i.e. $-\delta P$, needs to travel from the right to the left, and is therefore non-negative.

We are now in position to present the detailed description of the elements in the proof. For later reference, we mention the explicit values

$$P(Q_c) = \arctan \left(\frac{c}{\sqrt{2 - c^2}} \right) + \frac{c\sqrt{2 - c^2}}{2} - \text{sign}(c) \frac{\pi}{2}, \quad (18)$$

and

$$E(Q_c) = \frac{(2 - c^2)^{\frac{3}{2}}}{3}. \quad (19)$$

⁵To be more precise, the transferred momentum needs to travel from the left to the right in a reference frame attached to Q_{c_1, a_1} , i.e. moving at speed c_1 . Dispersive waves have a speed larger or equal in absolute value than the speed of sound $\sqrt{2}$ given by the dispersion relation. Since solitons are subsonic, a dispersive (hence sonic or supersonic) wave going to the right in the reference frame attached to a soliton also goes to the right in the frame of the lab. This peculiarity is a major difference between the Gross-Pitaevskii equation and the nonlinear Schrödinger equations with zero density at spatial infinity, where the speeds of solitons and dispersive waves may overlap.

A Minimizing properties of solitons

We will rely on a quantitative version of the minimizing property (12). Define the quadratic form

$$H_c := E''(Q_c) + cP''(Q_c).$$

In explicit form,

$$H_c(\varepsilon) := \frac{1}{4} \int_{\mathbb{R}} \frac{(\partial_x \varepsilon_\eta)^2}{1 - \eta_c} + \frac{1}{4} \int_{\mathbb{R}} \left(2 - \frac{\partial_{xx}^2 \eta_c}{(1 - \eta_c)^2} - \frac{(\partial_x \eta_c)^2}{(1 - \eta_c)^3} \right) \varepsilon_\eta^2 + \int_{\mathbb{R}} \left((1 - \eta_c) \varepsilon_v^2 + (c - 2v_c) \varepsilon_\eta \varepsilon_v \right), \quad (20)$$

where we have written $\varepsilon = (\varepsilon_\eta, \varepsilon_v)$. We have

Proposition 1. *There exists $\Lambda_c > 0$ such that*

$$H_c(\varepsilon) \geq \Lambda_c \|\varepsilon\|_X^2, \quad (21)$$

for any $\varepsilon \in X(\mathbb{R})$ satisfying the orthogonality conditions

$$\langle \varepsilon, \partial_x Q_c \rangle_{L^2(\mathbb{R})} = P'(Q_c)(\varepsilon) = 0. \quad (22)$$

Moreover, Λ_c is uniformly bounded from below by a positive constant for c in a compact subset of $(-\sqrt{2}, \sqrt{2}) \setminus \{0\}$.

Remark 2. The first orthogonality condition in (22) is related to the invariance with respect to translations of (HGP). It may be rephrased according to the equality

$$\langle \varepsilon, \partial_x Q_c \rangle_{L^2} = \frac{1}{2} \frac{d}{dy} \left(\left\| (\eta, v) - Q_c(\cdot - y) \right\|_{L^2}^2 \right) \Big|_{y=0} = 0.$$

It is possible to choose the condition involving the momentum P as an orthogonality condition due to the property

$$\frac{d}{dc} (P(Q_c)) \neq 0,$$

which follows from the explicit value (18) of $P(Q_c)$.

In particular, if $\varepsilon := (\eta, v) - Q_c$ satisfies the orthogonality conditions (22), then

$$(E + cP)(\eta, v) \geq (E + cP)(Q_c) + \Lambda_c \|\varepsilon\|_X^2 + \mathcal{O}(\|\varepsilon\|_X^3), \quad (23)$$

as $\|\varepsilon\|_X \rightarrow 0$.

Our next goal is to obtain a coercivity inequality in the spirit of (23) for functions (η, v) which do not belong to a small neighbourhood of a single soliton but instead to a neighbourhood of a finite sum of well-separated solitons like the one described by (7) and (8). We begin with the treatment of the orthogonality conditions in that case.

B Orthogonal decomposition of a chain of solitons

Given a positive parameter L , we define the set of admissible positions as

$$\text{Pos}_N(L) := \{ \mathbf{a} := (a_1, \dots, a_N) \in \mathbb{R}^N, \text{ s.t. } a_{k+1} > a_k + L, \forall 1 \leq k \leq N-1 \},$$

and, for $\alpha > 0$ and $\mathbf{c} = (c_1, \dots, c_N) \in \text{Adm}_N$, we set

$$\mathcal{U}_c(\alpha, L) := \left\{ (\eta, v) \in X(\mathbb{R}), \text{ s.t. } \inf_{\mathbf{a} \in \text{Pos}_N(L)} \left\| (\eta, v) - R_{\mathbf{c}, \mathbf{a}} \right\|_X < \alpha \right\},$$

and

$$\mu_{\mathbf{c}} := \min \{|c_k|, 1 \leq k \leq N\}.$$

When α is small enough and L is sufficiently large, we claim that the set $\mathcal{U}_{\mathbf{c}}(\alpha, L)$ provides a suitable framework to perform satisfactory decompositions. More precisely, we have

Proposition 2. *Let $\mathbf{c}^* \in \text{Adm}_N$. There exist positive numbers α_1 , L_1 and K_1 , depending only on \mathbf{c}^* , and two continuously differentiable functions $\mathfrak{C} \in \mathcal{C}^1(\mathcal{U}_{\mathbf{c}^*}(\alpha_1, L_1), \text{Adm}_N)$ and $\mathfrak{A} \in \mathcal{C}^1(\mathcal{U}_{\mathbf{c}^*}(\alpha_1, L_1), \mathbb{R}^N)$ such that for any $(\eta, v) \in \mathcal{U}_{\mathbf{c}^*}(\alpha_1, L_1)$, the function*

$$\varepsilon := (\eta, v) - R_{\mathbf{c}, \mathbf{a}},$$

where $\mathbf{c} = (c_1, \dots, c_N) := \mathfrak{C}(\eta, v)$ and $\mathbf{a} = (a_1, \dots, a_N) := \mathfrak{A}(\eta, v)$, satisfies the orthogonality conditions

$$\langle \varepsilon, \partial_x Q_{c_k, a_k} \rangle_{L^2} = P'(Q_{c_k, a_k})(\varepsilon) = 0, \quad (24)$$

for any $1 \leq k \leq N$. Moreover, if

$$\|(\eta, v) - R_{\mathbf{c}^*, \mathbf{a}^*}\|_X \leq \alpha,$$

for some $\mathbf{a}^* = (a_1^*, \dots, a_N^*) \in \text{Pos}_N(L)$, with $L \geq L_1$ and $\alpha \leq \alpha_1$, then

$$\|\varepsilon\|_X + \sum_{k=1}^N |c_k - c_k^*| + \sum_{k=1}^N |a_k - a_k^*| \leq K_1 \alpha. \quad (25)$$

Increasing L_1 and decreasing α_1 in the statement of Proposition 2 if necessary, one may further infer from (25) the following facts which we will use throughout.

Corollary 1. *Under the assumptions of Proposition 2, whenever $(\eta, v) \in \mathcal{U}_{\mathbf{c}^*}(\alpha_1, L)$ with $L \geq L_1$, it follows that $(\eta, v) \in \mathcal{NV}(\mathbb{R})$, since*

$$1 - \eta \geq \frac{\mu_{\mathbf{c}^*}^2}{8}. \quad (26)$$

Moreover, we have

$$\mathfrak{A}(\eta, v) \in \text{Pos}_N(L - 1), \quad (27)$$

$$\nu_{\mathfrak{C}(\eta, v)} \geq \frac{\nu_{\mathbf{c}^*}}{2}, \quad (28)$$

and

$$\mu_{\mathfrak{C}(\eta, v)} \geq \frac{\mu_{\mathbf{c}^*}}{2}. \quad (29)$$

C Almost minimizing properties close to a sum of solitons

Given $\mathbf{c}^* := (c_1^*, \dots, c_N^*) \in \text{Adm}_N$ and $(\eta, v) \in \mathcal{U}_{\mathbf{c}^*}(\alpha_1, L_1)$, we write

$$(\eta, v) = R_{\mathbf{c}, \mathbf{a}} + \varepsilon,$$

according to Proposition 2, where ε satisfies the orthogonality conditions (24). Our next goal is to show that the sum $R_{\mathbf{c}, \mathbf{a}}$ possesses almost minimizing properties. Since each soliton Q_{c_k, a_k} has its own speed c_k , the function $E + cP$, whose coercivity properties were exhibited in (23) in the case of a single soliton Q_c , has to be replaced by some function which, close to each location point a_k , resembles $E + c_k P$.

More precisely, we define the functions ⁶

$$\Psi_k(x) := \begin{cases} 1 & \text{if } k = 1, \\ \frac{1}{2} \left(1 + \operatorname{th} \left(\frac{\nu_{\mathfrak{c}^*}}{16} \left(x - \frac{a_{k-1} + a_k}{2} \right) \right) \right) & \text{if } 2 \leq k \leq N, \\ 0 & \text{if } k = N + 1. \end{cases}$$

We next set

$$F(\eta, v) := E(\eta, v) + \sum_{k=1}^N c_k P_k(\eta, v),$$

where

$$P_k(\eta, v) := \int_{\mathbb{R}} p(\eta, v) (\Psi_k - \Psi_{k+1}).$$

Notice that by construction,

$$\sum_{k=1}^N (\Psi_k - \Psi_{k+1}) = \Psi_1 - \Psi_{N+1} = 1,$$

and that, provided the points a_k are sufficiently far apart one from each other, we have for each k , $\Psi_k - \Psi_{k+1} \simeq 1$ in a neighborhood of Q_{c_k, a_k} and $\Psi_k - \Psi_{k+1} \simeq 0$ in the neighborhood of Q_{c_j, a_j} , for $j \neq k$. Notice also that, even though it does not explicitly appear in the notation, Ψ_k depends on (η, v) through the points a_k . This will be of particular importance when we introduce time dependence later on.

In order to estimate the function F , we localize the function ε according to the following decomposition. Let $0 < \tau < \nu_{\mathfrak{c}^*}/16$ to be defined later (see (38)), depending only on \mathfrak{c}^* . We set

$$\Phi_k(x) := \frac{1}{2} \left(\operatorname{th} \left(\tau \left(x - a_k + \frac{L_1}{4} \right) \right) - \operatorname{th} \left(\tau \left(x - a_k - \frac{L_1}{4} \right) \right) \right),$$

for $1 \leq k \leq N$, and

$$\Phi_{k,k+1} := \begin{cases} \frac{1}{2} \left(1 - \operatorname{th} \left(\tau \left(x - a_1 + \frac{L_1}{4} \right) \right) \right) & \text{if } k = 0, \\ \frac{1}{2} \left(\operatorname{th} \left(\tau \left(x - a_k - \frac{L_1}{4} \right) \right) - \operatorname{th} \left(\tau \left(x - a_{k+1} + \frac{L_1}{4} \right) \right) \right) & \text{if } 1 < k < N, \\ \frac{1}{2} \left(1 + \operatorname{th} \left(\tau \left(x - a_N - \frac{L_1}{4} \right) \right) \right) & \text{if } k = N. \end{cases}$$

Once more, we note that by construction,

$$\sum_{k=1}^N \Phi_k + \sum_{k=0}^N \Phi_{k,k+1} = 1, \tag{30}$$

and that each term in the sum is non-negative. We finally define

$$\varepsilon_k := \Phi_k(\cdot + a_k)^{\frac{1}{2}} \varepsilon(\cdot + a_k), \tag{31}$$

for $1 \leq k \leq N$, and

$$\varepsilon_{k,k+1} := \Phi_{k,k+1}^{\frac{1}{2}} \varepsilon, \tag{32}$$

for $0 \leq k \leq N$.

⁶In the definition of Ψ_k , one could a priori replace the constant $\nu_{\mathfrak{c}^*}/16$ by an arbitrary positive constant ν^* . In order that the monotonicity formula later stated in Proposition 5 holds, the choice of ν^* is actually restricted. Up to a constant multiplicity factor, $\nu_{\mathfrak{c}^*}/16$ is presumably optimal.

Proposition 3. Assume that $(\eta, v) = R_{c,a} + \varepsilon \in \mathcal{U}_{c^*}(\alpha_1, L)$, with $L \geq L_1$. Then,

$$\begin{aligned} E(\eta, v) &= \sum_{k=1}^N E(Q_{c_k}) + \frac{1}{2} \sum_{k=1}^N E''(Q_{c_k})(\varepsilon_k, \varepsilon_k) + \frac{1}{2} \sum_{k=0}^N E''(0)(\varepsilon_{k,k+1}, \varepsilon_{k,k+1}) \\ &\quad + \mathcal{O}\left(\left(\tau + \exp\left(-\tau \frac{L_1}{2}\right)\right) \|\varepsilon\|_X^2\right) + \mathcal{O}\left(\|\varepsilon\|_X^3\right) + \mathcal{O}\left(L \exp\left(-\frac{\nu_{c^*} L}{2}\right)\right), \end{aligned} \quad (33)$$

and

$$\begin{aligned} P_k(\eta, v) &= P(Q_{c_k}) + \frac{1}{2} P''(Q_{c_k})(\varepsilon_k, \varepsilon_k) + \frac{1}{2} P_k''(0)(\varepsilon_{k,k+1}, \varepsilon_{k,k+1}) + \frac{1}{2} P_k''(0)(\varepsilon_{k-1,k}, \varepsilon_{k-1,k}) \\ &\quad + \mathcal{O}\left(L \exp\left(-\frac{\nu_{c^*} L}{16}\right)\right) + \mathcal{O}\left(\exp\left(-\frac{\tau L_1}{2}\right) \|\varepsilon\|_X^2\right). \end{aligned} \quad (34)$$

Remark 3. Here and in the sequel, we have found convenient to simplify the presentation of our estimates by introducing the notation \mathcal{O} . We adopt as a definition that one is allowed to substitute a quantity S by the notation $\mathcal{O}(T)$ if and only if there exists a positive constant K , depending only on c^* , such that

$$|S| \leq KT.$$

Later on, we will analyze the evolution in time of the evaluation of F along a solution (η, v) of (HGP). The analysis is complicated by the fact that the speeds c_k are themselves functions of (η, v) . To circumvent this difficulty, it turns out that it is sufficient to consider instead of F the function G defined on $\mathcal{U}_{c^*}(\alpha_1, L_1)$ by

$$G(\eta, v) := E(\eta, v) + \sum_{k=1}^N c_k^* P_k(\eta, v).$$

Notice here that the speeds are the reference speeds c_k^* . Concerning the functional G , we have the following consequence of Proposition 3.

Corollary 2. Assume that $(\eta, v) = R_{c,a} + \varepsilon \in \mathcal{U}_{c^*}(\alpha_1, L)$, with $L \geq L_1$. Then,

$$\begin{aligned} G(\eta, v) &= \sum_{k=1}^N (E(Q_{c_k^*}) + c_k^* P(Q_{c_k^*})) + \frac{1}{2} \sum_{k=1}^N H_{c_k}(\varepsilon_k) + \frac{1}{2} \sum_{k=0}^N H_0^k(\varepsilon_{k,k+1}) + \mathcal{O}\left(\sum_{k=1}^N |c_k - c_k^*|^2\right) \\ &\quad + \mathcal{O}\left(\left(\tau + \exp\left(-\frac{\tau L_1}{2}\right)\right) \|\varepsilon\|_X^2\right) + \mathcal{O}\left(\|\varepsilon\|_X^3\right) + \mathcal{O}\left(L \exp\left(-\frac{\nu_{c^*} L}{16}\right)\right), \end{aligned}$$

where we have set

$$\begin{aligned} H_0^0(\varepsilon_{0,1}) &= (E''(0) + c_1 P_1''(0))(\varepsilon_{0,1}, \varepsilon_{0,1}), \\ H_0^k(\varepsilon_{k,k+1}) &= (E''(0) + c_k P_k''(0) + c_{k+1} P_{k+1}''(0))(\varepsilon_{k,k+1}, \varepsilon_{k,k+1}), \quad \forall 1 \leq k \leq N-1, \\ H_0^N(\varepsilon_{N,N+1}) &= (E''(0) + c_N P_N''(0))(\varepsilon_{N,N+1}, \varepsilon_{N,N+1}). \end{aligned}$$

It is not possible to invoke Proposition 1 directly to obtain a lower bound on the terms $H_{c_k}(\varepsilon_k)$ since ε_k only satisfies the orthogonality conditions (22) asymptotically as $L_1 \rightarrow +\infty$. We obtain

Lemma 1. Under the assumptions of Corollary 2, there exists a positive constant Λ^* , depending only on c^* , such that

$$H_{c_k}(\varepsilon_k) \geq \Lambda^* \|\varepsilon_k\|_X^2 + \mathcal{O}\left(L_1^{\frac{1}{2}} \exp\left(-\frac{\tau L_1}{2}\right) \|\varepsilon\|_X^2\right), \quad (35)$$

for any $1 \leq k \leq N$, and such that

$$H_0^k(\varepsilon_{k,k+1}) \geq \Lambda^* \|\varepsilon_{k,k+1}\|_X^2, \quad (36)$$

for any $0 \leq k \leq N$. Moreover,

$$\sum_{k=1}^N \|\varepsilon_k\|_X^2 + \sum_{k=0}^N \|\varepsilon_{k,k+1}\|_X^2 \geq \|\varepsilon\|_X^2. \quad (37)$$

In view of Corollary 2 and Lemma 1, we now fix the value of τ , sufficiently small, and then increase the value of L_1 , if necessary, in such a way that ⁷

$$\mathcal{O}\left(\left(\tau + L_1^{\frac{1}{2}} \exp\left(-\frac{\tau L_1}{2}\right)\right) \|\varepsilon\|_X^2\right) \leq \frac{\Lambda^*}{4} \|\varepsilon\|_X^2. \quad (38)$$

With these choices we obtain

Corollary 3. *Assume that $(\eta, v) = R_{\mathbf{c}, \mathbf{a}} + \varepsilon \in \mathcal{U}_{\mathbf{c}^*}(\alpha_1, L)$, with $L \geq L_1$. Then,*

$$\begin{aligned} G(\eta, v) &\geq \sum_{k=1}^N (E(Q_{c_k^*}) + c_k^* P(Q_{c_k^*})) + \frac{\Lambda^*}{4} \|\varepsilon\|_X^2 + \mathcal{O}\left(\|\varepsilon\|_X^3\right) + \mathcal{O}\left(\sum_{k=1}^N |c_k - c_k^*|^2\right) \\ &\quad + \mathcal{O}\left(L \exp\left(-\frac{\nu_{\mathbf{c}^*} L}{16}\right)\right), \end{aligned} \quad (39)$$

as well as

$$G(\eta, v) \leq \sum_{k=1}^N (E(Q_{c_k^*}) + c_k^* P(Q_{c_k^*})) + \mathcal{O}\left(\|\varepsilon\|_X^2\right) + \mathcal{O}\left(\sum_{k=1}^N |c_k - c_k^*|^2\right) + \mathcal{O}\left(L \exp\left(-\frac{\nu_{\mathbf{c}^*} L}{16}\right)\right). \quad (40)$$

Notice that, up to this point, the entire analysis has been independent of time. In the next section, we combine the previous results with information extracted from equation (HGP).

D Dynamics of the modulation parameters

In this subsection, as well as in Subsections E and F below, $\mathbf{c}^* \in \text{Adm}_N$ is fixed and we assume that $T > 0$, and $(\eta, v) \in \mathcal{C}^0([0, T], X(\mathbb{R}))$ is a given solution to equation (HGP) such that for some $\alpha \leq \alpha_1$ and $L \geq L_1$,

$$(\eta(\cdot, t), v(\cdot, t)) \in \mathcal{U}_{\mathbf{c}^*}(\alpha, L),$$

for any $t \in [0, T]$.

For $t \in [0, T]$, we define

$$\mathbf{c}(t) = (c_1(t), \dots, c_N(t)) := \mathfrak{C}(v(\cdot, t), \eta(\cdot, t)) \quad \text{and} \quad \mathbf{a}(t) = (a_1(t), \dots, a_N(t)) := \mathfrak{A}(v(\cdot, t), \eta(\cdot, t)),$$

and

$$\varepsilon(\cdot, t) := (\eta(\cdot, t), v(\cdot, t)) - R_{\mathbf{c}(t), \mathbf{a}(t)},$$

where the functions \mathfrak{C} and \mathfrak{A} are given by Proposition 2.

The following proposition expresses the fact that the modulation parameters should follow those of the underlying solitons, as long as the solution remains close to them.

⁷By the left-hand side of (38), we actually mean all the terms which we have previously estimated in Corollary 2 and Lemma 1 as the corresponding \mathcal{O} .

Proposition 4. *There exist positive numbers $\alpha_2 \leq \alpha_1$ and $L_2 \geq L_1$, depending only on \mathbf{c}^* , such that, if $\alpha \leq \alpha_2$ and $L \geq L_2$, then $\mathbf{c} \in \mathcal{C}^1([0, T], \text{Adm}_N)$, $\mathbf{a} \in \mathcal{C}^1([0, T], \mathbb{R}^N)$, and we have the estimates*

$$\sum_{k=1}^N \left(|a'_k(t) - c_k(t)| + |c'_k(t)| \right) = \mathcal{O}\left(\|\varepsilon(\cdot, t)\|_X\right) + \mathcal{O}\left(L \exp\left(-\frac{\nu_{\mathbf{c}^*} L}{2}\right)\right), \quad (41)$$

for any $t \in [0, T]$.

From now on, we make the additional assumption that our base set of speeds \mathbf{c}^* satisfies the inequalities

$$\mathbf{c}_1^* < \cdots < \mathbf{c}_N^*, \quad (42)$$

and we set

$$\sigma^* = \frac{1}{2} \min \{c_{k+1}^* - c_k^*, 1 \leq k \leq N-1\} > 0.$$

As a consequence of (25), (41) and (42), we obtain

Corollary 4. *There exist positive numbers $\alpha_3 \leq \alpha_2$ and $L_3 \geq L_2$, depending only on \mathbf{c}^* , such that if $\alpha \leq \alpha_3$ and $L \geq L_3$, then for any $t \in [0, T]$,*

$$a_{k+1}(t) - a_k(t) > a_{k+1}(0) - a_k(0) + \sigma^* t \geq (L-1) + \sigma^* t, \quad (43)$$

for each $1 \leq k \leq N-1$, and

$$\sqrt{2 - (a'_k(t))^2} \geq \frac{\nu_{\mathbf{c}^*}}{2}, \quad (44)$$

for each $1 \leq k \leq N$.

E Evolution in time of \mathcal{G} and monotonicity formula

We define the function \mathcal{G} on $[0, T]$ by

$$\mathcal{G}(t) := G(\eta(\cdot, t), v(\cdot, t)),$$

where G was introduced in Subsection C. We rewrite \mathcal{G} according to the equality

$$\mathcal{G}(t) = E(t) + c_1^* P(t) + \sum_{k=2}^N (c_k^* - c_{k-1}^*) Q_k(t), \quad (45)$$

where

$$E(t) := E(\eta(\cdot, t), v(\cdot, t)), \quad P(t) := P(\eta(\cdot, t), v(\cdot, t)),$$

and

$$Q_k(t) := \int_{\mathbb{R}} \Psi_k(x) p(\eta(x, t), v(x, t)) dx = \sum_{j=k}^N P_j(t).$$

In case $N = 1$ (i.e. there is only one soliton), it follows from the conservation of energy and momentum that \mathcal{G} is constant in time. In the general case, we have the following monotonicity formula.

Proposition 5. *There exist positive numbers $\alpha_4 \leq \alpha_3$ and $L_4 \geq L_3$, depending only on \mathbf{c}^* , such that if $\alpha \leq \alpha_4$ and $L \geq L_4$, then*

$$\frac{d}{dt} \left(Q_k(t) \right) \leq \mathcal{O}\left(\exp\left(-\frac{\nu_{\mathbf{c}^*} (L + \sigma^* t)}{16}\right)\right),$$

for any $t \in [0, T]$.

As a consequence of the ordering condition (42), we thus obtain

Corollary 5. *If $\alpha \leq \alpha_4$ and $L \geq L_4$, then*

$$\frac{d}{dt}(\mathcal{G}(t)) \leq \mathcal{O}\left(\exp\left(-\frac{\nu_{\mathbf{c}^*}(L + \sigma^* t)}{16}\right)\right),$$

for any $t \in [0, T]$.

F Uniform control on $\|\varepsilon\|_X$

We are now in position to obtain a uniform control on $\|\varepsilon(\cdot, t)\|_X$ combining the previous results with simple algebraic identities. We divide the analysis in a number of steps, each one providing a uniform control on an intermediate quantity.

Step 1. *Control on $|Q_k(t) - Q_k(0)|$.*

Since E and P are conserved by the flow, we obtain from (45),

$$Q_j(0) - Q_j(t) = \mathcal{G}(0) - \mathcal{G}(t) + \frac{1}{c_j^* - c_{j-1}^*} \sum_{k \neq j} (c_k^* - c_{k-1}^*) (Q_k(t) - Q_k(0)), \quad (46)$$

for any $2 \leq j \leq N$. On the one hand, we have by (40) for $\mathcal{G}(0)$, and (39) for $\mathcal{G}(t)$,

$$\mathcal{G}(0) - \mathcal{G}(t) \leq \mathcal{O}\left(\|\varepsilon(\cdot, 0)\|_X^2\right) + \mathcal{O}\left(\|\varepsilon(\cdot, t)\|_X^3\right) + \mathcal{O}\left(\sum_{k=1}^N |c_k(t) - c_k^*|^2\right) + \mathcal{O}\left(L \exp\left(-\frac{\nu_{\mathbf{c}^*} L}{16}\right)\right). \quad (47)$$

On the other hand, from Proposition 5 and assumption (42) on the speeds, we have

$$\frac{1}{c_j^* - c_{j-1}^*} \sum_{k \neq j} (c_k^* - c_{k-1}^*) (Q_k(t) - Q_k(0)) \leq \mathcal{O}\left(\exp\left(-\frac{\nu_{\mathbf{c}^*} L}{16}\right)\right). \quad (48)$$

Combining (46), (47) and (48) for the upper bound part, and using once more Proposition 5 for the lower bound part, we derive

$$\sum_{k=2}^N |Q_k(t) - Q_k(0)| = \mathcal{O}\left(\|\varepsilon(\cdot, 0)\|_X^2\right) + \mathcal{O}\left(\|\varepsilon(\cdot, t)\|_X^3\right) + \mathcal{O}\left(\sum_{k=1}^N |c_k(t) - c_k^*|^2\right) + \mathcal{O}\left(L \exp\left(-\frac{\nu_{\mathbf{c}^*} L}{16}\right)\right). \quad (49)$$

Step 2. *Control on $|P(Q_{c_k(t)}) - P(Q_{c_k(0)})|$.*

It follows from Proposition 3 that

$$P(t) - Q_2(t) = P_1(t) = P(Q_{c_1(t)}) + \mathcal{O}\left(\|\varepsilon(\cdot, t)\|_X^2\right) + \mathcal{O}\left(L \exp\left(-\frac{\nu_{\mathbf{c}^*} L}{16}\right)\right), \quad (50)$$

and that for $2 \leq k \leq N$,

$$Q_k(t) - Q_{k+1}(t) = P_k(t) = P(Q_{c_k(t)}) + \mathcal{O}\left(\|\varepsilon(\cdot, t)\|_X^2\right) + \mathcal{O}\left(L \exp\left(-\frac{\nu_{\mathbf{c}^*} L}{16}\right)\right). \quad (51)$$

Combining (50) and (51) with (49) and the conservation of the momentum P , we obtain

$$\begin{aligned} \sum_{k=1}^N |P(Q_{c_k(t)}) - P(Q_{c_k(0)})| &= \mathcal{O}\left(\|\varepsilon(\cdot, 0)\|_X^2\right) + \mathcal{O}\left(\|\varepsilon(\cdot, t)\|_X^2\right) + \mathcal{O}\left(\sum_{k=1}^N |c_k(t) - c_k^*|^2\right) \\ &\quad + \mathcal{O}\left(L \exp\left(-\frac{\nu_{\mathbf{c}^*} L}{16}\right)\right). \end{aligned} \quad (52)$$

Step 3. Control on $|c_k(t) - c_k^*|$.

Combining (52) with the fact that

$$\frac{d}{dc} \left(P(Q_c) \right) = (2 - c^2)^{\frac{1}{2}} \neq 0,$$

and the inequality

$$\sum_{k=1}^N |c_k(t) - c_k^*| \leq \sum_{k=1}^N |c_k(t) - c_k(0)| + \sum_{k=1}^N |c_k(0) - c_k^*|,$$

we obtain the existence of positive numbers $\alpha_5 \leq \alpha_4$ and $L_5 \geq L_4$, depending only on \mathbf{c}^* , such that if $\alpha \leq \alpha_5$ and $L \geq L_5$, then

$$\sum_{k=1}^N |c_k(t) - c_k^*| = \mathcal{O}(\|\varepsilon(\cdot, 0)\|^2) + \mathcal{O}(\|\varepsilon(\cdot, t)\|^2) + \mathcal{O}\left(\sum_{k=1}^N |c_k(0) - c_k^*|\right) + \mathcal{O}\left(L \exp\left(-\frac{\nu_{\mathbf{c}^*} L}{16}\right)\right).$$

Finally, from Proposition 2 we may bound $\sum_{k=1}^N |c_k(0) - c_k^*|$ by $\mathcal{O}(\|\varepsilon(\cdot, 0)\|_X)$, and therefore obtain

$$\sum_{k=1}^N |c_k(t) - c_k^*| = \mathcal{O}(\|\varepsilon(\cdot, 0)\|) + \mathcal{O}(\|\varepsilon(\cdot, t)\|^2) + \mathcal{O}\left(L \exp\left(-\frac{\nu_{\mathbf{c}^*} L}{16}\right)\right). \quad (53)$$

Step 4. Control on $\|\varepsilon(\cdot, t)\|_X$.

Combining Corollary 3 for $(\eta, v) := (\eta(\cdot, t), v(\cdot, t))$, and Corollary 5, we obtain, if $\alpha \leq \alpha_4$ and $L \geq L_4$,

$$\frac{\Lambda^*}{4} \|\varepsilon(\cdot, t)\|_X^2 = \mathcal{O}(\|\varepsilon(\cdot, 0)\|_X^2) + \mathcal{O}\left(\sum_{k=1}^N |c_k(t) - c_k^*|^2\right) + \mathcal{O}\left(L \exp\left(-\frac{\nu_{\mathbf{c}^*} L}{16}\right)\right) + \mathcal{O}(\|\varepsilon(\cdot, t)\|_X^3).$$

It follows therefore from (53) that there exist positive numbers $\alpha_6 \leq \alpha_5$ and $L_6 \geq L_5$, depending only on \mathbf{c}^* , such that if $\alpha \leq \alpha_6$ and $L \geq L_6$, then

$$\|\varepsilon(\cdot, t)\|_X^2 \leq \mathcal{O}(\|\varepsilon(\cdot, 0)\|_X^2) + \mathcal{O}\left(L \exp\left(-\frac{\nu_{\mathbf{c}^*} L}{16}\right)\right). \quad (54)$$

G Proof of Theorem 2 completed

Let (η^0, v^0) be as in the statement of Theorem 2. We first impose the condition $\alpha^* < \alpha_6$ and $L^* > L_6 + 2$, so that by continuity of the flow, since $L^0 \geq L^*$ by assumption,

$$T_{\text{stop}} := \sup \{t \geq 0, \text{ s.t. } (\eta(\cdot, s), v(\cdot, s)) \in \mathcal{U}_{\mathbf{c}^*}(\alpha_6, L^0 - 2), \forall s \in [0, t]\} > 0.$$

Notice that $T_{\text{stop}} < T_{\text{max}}$ since $\mathcal{U}_{\mathbf{c}^*}(\alpha_6, L^0 - 2) \subset \mathcal{NV}(\mathbb{R})$.

By (54) we have

$$\|\varepsilon(\cdot, t)\|_X = \mathcal{O}\left(\alpha^0 + \exp\left(-\frac{\nu_{\mathbf{c}^*} L^0}{33}\right)\right). \quad (55)$$

Combining (55) with (53) and (54) yields a positive number K_6 , depending only on \mathbf{c}^* , such that

$$\|(\eta(\cdot, t), v(\cdot, t)) - R_{\mathbf{c}^*, a(t)}\|_X \leq K_6 \left(\alpha^0 + \exp\left(-\frac{\nu_{\mathbf{c}^*} L^0}{33}\right)\right), \quad (56)$$

for any $t \in [0, T_{\text{stop}})$. On the other hand, by definition of L^0 , we have $(\eta(0), v(0)) \in \mathcal{U}_{c^*}(\alpha^*, L^0)$. In particular, by Proposition 2,

$$\min \{a_{k+1}(0) - a_k(0), 1 \leq k \leq N-1\} \geq L^0 - 1,$$

so that, by Corollary 4, we finally have

$$\min \{a_{k+1}(t) - a_k(t), 1 \leq k \leq N-1\} = \min \{a_{k+1}(0) - a_k(0), 1 \leq k \leq N-1\} > L^0 - 2, \quad (57)$$

for any $t \in [0, T_{\text{stop}})$.

We therefore additionally impose the condition that α^* and L^* satisfy

$$K_6 \left(\alpha^* + \exp \left(- \frac{\nu_{c^*} L^*}{33} \right) \right) < \alpha_6,$$

so that

$$T_{\text{stop}} = +\infty,$$

and in particular, (56) and (57) hold for any $t \in [0, +\infty)$. Finally, combining Proposition 4 and (53), we obtain

$$\sum_{k=1}^N |a'_k(t) - c_k^*| \leq K_7 \left(\alpha^0 + \exp \left(- \frac{\nu_{c^*} L^0}{33} \right) \right),$$

and the proof of Theorem 2 is completed by setting $A^* = \max\{K_6, K_7\}$.

1.5 Outline of the paper

In Section 2, we present the proofs of the results stated in Subsections A, B and C of the introduction, which are related to the minimizing properties of solitons and sums of solitons. Section 3 is instead devoted to the proofs of the results related to the dynamical properties of (HGP), and which are stated in Subsections D and E of the introduction. Finally, for the sake of completeness, we provide in Appendix A a quantitative version of the implicit function theorem which we use in Section 2 in order to define the modulation parameters.

2 Minimizing properties around a soliton and a sum of solitons

2.1 Minimizing properties of solitons

In this subsection we present the proof of Proposition 1. It is reminiscent from arguments developed in [12].

Proof of Proposition 1. Let us recall that the quadratic form H_c is defined in (20) as

$$H_c(\varepsilon) := \frac{1}{4} \int_{\mathbb{R}} \frac{(\partial_x \varepsilon_\eta)^2}{1 - \eta_c} + \frac{1}{4} \int_{\mathbb{R}} \left(2 - \frac{\partial_{xx}^2 \eta_c}{(1 - \eta_c)^2} - \frac{(\partial_x \eta_c)^2}{(1 - \eta_c)^3} \right) \varepsilon_\eta^2 + \int_{\mathbb{R}} \left((1 - \eta_c) \varepsilon_v^2 + (c - 2v_c) \varepsilon_\eta \varepsilon_v \right).$$

Equation (11) writes for η_c as

$$\partial_{xx}^2 \eta_c - (2 - c^2) \eta_c + 3\eta_c^2 = 0. \quad (2.1)$$

Invoking (5) and (2.1), the expression of H_c may be recast as

$$H_c(\varepsilon) = L_c(\varepsilon_\eta) + \int_{\mathbb{R}} (1 - \eta_c) \left(\varepsilon_v + \frac{c}{2(1 - \eta_c)^2} \varepsilon_\eta \right)^2, \quad (2.2)$$

where

$$L_c(\varepsilon_\eta) := \frac{1}{4} \int_{\mathbb{R}} \frac{(\partial_x \varepsilon_\eta)^2}{1 - \eta_c} + \frac{1}{4} \int_{\mathbb{R}} \frac{2 - c^2 - 6\eta_c + 3\eta_c^2}{(1 - \eta_c)^2} \varepsilon_\eta^2.$$

The Sturm-Liouville operator \mathcal{L}_c associated to the quadratic form L_c , namely

$$\mathcal{L}_c(\varepsilon_\eta) := -\partial_x \left(\frac{\partial_x \varepsilon_\eta}{4(1 - \eta_c)} \right) + \frac{2 - c^2 - 6\eta_c + 3\eta_c^2}{4(1 - \eta_c)^2} \varepsilon_\eta,$$

is a self-adjoint, unbounded operator on $L^2(\mathbb{R})$, with domain $H^2(\mathbb{R})$. Since $\eta_c(x) \rightarrow 0$ as $|x| \rightarrow +\infty$ by (5), it follows from the Weyl criterion that its essential spectrum is equal to $[(2 - c^2)/4, +\infty)$. On the other hand, one can translate the invariance with respect to translations of equation (HGP) into the property that the function $\partial_x \eta_c$ belongs to the kernel of \mathcal{L}_c . Again in view of (5), the function $\partial_x \eta_c$ has exactly one zero. As a consequence, one can infer from standard Sturm-Liouville theory (see e.g. [6]) that the operator \mathcal{L}_c has exactly one negative eigenvalue $-\mu^-$. Moreover, the associated eigenspace, as well as the kernel of \mathcal{L}_c , are of dimension one. We will denote η^- an eigenfunction of \mathcal{L}_c for the eigenvalue $-\mu^-$. Notice in particular, that there exists a constant $0 < \mu^+ \leq 2 - c^2$ such that

$$L_c(\varepsilon_\eta) \geq \mu^+ \|\varepsilon_\eta\|_{L^2}^2, \quad (2.3)$$

as soon as $\langle \varepsilon_\eta, \eta^- \rangle_{L^2} = \langle \varepsilon_\eta, \partial_x \eta_c \rangle_{L^2} = 0$.

Coming back to the quadratic form H_c , we notice that the associated operator \mathcal{H}_c , namely

$$\mathcal{H}_c(\varepsilon) := \left(\mathcal{L}_c(\varepsilon_\eta) + \frac{c^2}{4(1 - \eta_c)^3} \varepsilon_\eta + \frac{c}{2(1 - \eta_c)} \varepsilon_v, \frac{c}{2(1 - \eta_c)} \varepsilon_\eta + (1 - \eta_c) \varepsilon_v \right),$$

is a self-adjoint, unbounded operator on $L^2(\mathbb{R})^2$, with domain $H^2(\mathbb{R}) \times L^2(\mathbb{R})$. Moreover, it again follows from the Weyl criterion that its essential spectrum is equal to $[(2 - c^2)/(3 + \sqrt{1 + 4c^2}), +\infty)$. In view of (5) and (2.2), we next check that

$$H_c \left(\eta^-, -\frac{c\eta^-}{2(1 - \eta_c)^2} \right) < 0, \quad \text{and} \quad H_c(\partial_x \eta_c, \partial_x v_c) = 0,$$

so that \mathcal{H}_c has at least one negative eigenvalue, and the dimension of its kernel is at least one. Assume next that either \mathcal{H}_c owns another negative eigenvalue, or its kernel is at least of dimension two. Then, there exists a non-positive direction $\varepsilon = (\varepsilon_\eta, \varepsilon_v)$ for H_c such that $\langle \varepsilon_\eta, \eta^- \rangle_{L^2(\mathbb{R})} = \langle \varepsilon_\eta, \partial_x \eta_c \rangle_{L^2(\mathbb{R})} = 0$. In view of (2.2), this is in contradiction with (2.3). Therefore, \mathcal{H}_c has exactly one negative eigenvalue $-\lambda^-$, with eigenfunction $\chi^- := (\chi_\eta^-, \chi_v^-)$, while its kernel is spanned by $\partial_x Q_c = (\partial_x \eta_c, \partial_x v_c)$. As a consequence, there exists a positive constant $\lambda^+ \leq (2 - c^2)/(3 + \sqrt{1 + 4c^2})$ such that

$$H_c(\varepsilon) \geq \lambda^+ \|\varepsilon\|_{L^2 \times L^2}^2 := \lambda^+ (\|\varepsilon_\eta\|_{L^2}^2 + \|\varepsilon_v\|_{L^2}^2), \quad (2.4)$$

for any pair ε in the closed subspace

$$P_c^+ := \left\{ (\varepsilon_\eta, \varepsilon_v) \in H^1(\mathbb{R}) \times L^2(\mathbb{R}), \text{ s.t. } \langle \varepsilon, \chi^- \rangle_{L^2 \times L^2} = \langle \varepsilon, \partial_x Q_c \rangle_{L^2 \times L^2} = 0 \right\}. \quad (2.5)$$

The next step in the proof consists in checking that estimate (2.4) remains available (up to a further choice of λ^+), when the orthogonality conditions in (2.5) are replaced by conditions (22). The proof is reminiscent from [11]. We consider the map $S(c) := E(Q_c) + cP(Q_c)$, which is well-defined and smooth on $(-\sqrt{2}, \sqrt{2}) \setminus \{0\}$ in view of (18) and (19). Using (11), we compute

$$S'(c) = E'(Q_c)(\partial_c Q_c) + cP'(Q_c)(\partial_c Q_c) + P(Q_c) = P(Q_c),$$

so that, by (18),

$$S''(c) = P'(Q_c)(\partial_c Q_c) = (2 - c^2)^{\frac{1}{2}}.$$

On the other hand, taking the derivative with respect to c of the identity $E'(Q_c) + cP'(Q_c) = 0$, we get

$$\langle \mathcal{H}_c(\partial_c Q_c), \varepsilon \rangle_{L^2 \times L^2} + P'(Q_c)(\varepsilon) = 0, \quad (2.6)$$

for any $\varepsilon \in X(\mathbb{R})$. For $\varepsilon = \partial_c Q_c$, this gives

$$H_c(\partial_c Q_c) = -P'(Q_c)(\partial_c Q_c) = -(2 - c^2)^{\frac{1}{2}} < 0. \quad (2.7)$$

At this stage, we can decompose $\partial_c Q_c$ as $\partial_c Q_c = \alpha \chi^- + \beta \partial_x Q_c + r_c$, with $r_c \in P_c^+$. As a consequence,

$$H_c(\partial_c Q_c) = -\lambda_- \alpha^2 + H_c(r_c),$$

so that inequality (2.7) is equivalent to the fact that $\alpha \neq 0$, combined with the existence of a number $0 \leq \delta < 1$ such that

$$H_c(r_c) = \delta \lambda^- \alpha^2. \quad (2.8)$$

According to the value of δ , two situations can occur. If $\delta = 0$, then r_c is equal to 0. Since $\alpha \neq 0$, conditions (22) are actually identical to the conditions in (2.5), so that (2.4) remains available under conditions (22). If $\delta > 0$, we can decompose a pair ε which satisfies conditions (22), as $\varepsilon := a \chi^- + r$, with $r \in P_c^+$, and compute

$$H_c(\varepsilon) = -\lambda^- a^2 + \frac{1 - \delta}{2} H_c(r) + \frac{1 + \delta}{2} H_c(r).$$

In view of (2.4), the quadratic form H_c is positive on P_c^+ . Therefore, we can apply the Cauchy-Schwarz inequality to write

$$H_c(\varepsilon) \geq -\lambda^- a^2 + \frac{(1 + \delta) \langle \mathcal{H}_c(r_c), r \rangle_{L^2 \times L^2}^2}{2H_c(r_c)} + \frac{1 - \delta}{2} H_c(r).$$

At this stage, it follows from (22) and (2.6) that

$$\langle \mathcal{H}_c(r_c), r \rangle_{L^2 \times L^2} = \langle \mathcal{H}_c(\partial_c Q_c), \varepsilon \rangle_{L^2 \times L^2} + \lambda^- a \alpha = \lambda^- a \alpha.$$

Combining with (2.8), we are led to

$$H_c(\varepsilon) \geq \frac{1 - \delta}{2} (\lambda^- a^2 + H_c(r)),$$

so that by (2.4),

$$H_c(\varepsilon) \geq \frac{1 - \delta}{2} \min \{ \lambda^-, \lambda^+ \} \|\varepsilon\|_{L^2 \times L^2}^2.$$

As a conclusion, up to a possible further choice of the positive constant λ^+ , estimate (2.4) still holds when ε satisfies conditions (22).

In order to complete the proof of estimate (21), it remains to replace the L^2 -norm of ε_η in (2.4) by an H^1 -norm. Given a number $0 < \tau < 1$, we write

$$H_c(\varepsilon) = \frac{\tau}{4} \int_{\mathbb{R}} \frac{(\partial_x \varepsilon_\eta)^2}{1 - \eta_c} + \tau \left(H_c(\varepsilon) - \frac{1}{4} \int_{\mathbb{R}} \frac{(\partial_x \varepsilon_\eta)^2}{1 - \eta_c} \right) + (1 - \tau) H_c(\varepsilon).$$

In view of formulae (5) and (20), there exists a positive constant K_c , possibly depending on c , such that

$$\left| H_c(\varepsilon) - \frac{1}{4} \int_{\mathbb{R}} \frac{(\partial_x \varepsilon_\eta)^2}{1 - \eta_c} \right| \leq K_c \|\varepsilon\|_{L^2 \times L^2}^2,$$

so that, under conditions (22), we can take benefit of (2.4) to get

$$H_c(\varepsilon) \geq \frac{\tau}{4} \int_{\mathbb{R}} \frac{(\partial_x \varepsilon \eta)^2}{1 - \eta_c} + ((1 - \tau)\lambda^+ - \tau K) \|\varepsilon\|_{L^2 \times L^2}^2.$$

Since $\eta_c(0) = (2 - c^2)/2 = \max\{|\eta_c(x)|, x \in \mathbb{R}\}$ by (5), we deduce that

$$H_c(\varepsilon) \geq \frac{\tau}{2c^2} \int_{\mathbb{R}} (\partial_x \varepsilon \eta)^2 + ((1 - \tau)\lambda^+ - \tau K) \|\varepsilon\|_{L^2 \times L^2}^2.$$

Choosing τ sufficiently small yields estimate (21). Finally, since H_c depends analytically on c and the first two eigenvalues of H_c are simple, the optimal constant in (21) depends continuously on c , from which the last statement in Proposition 1 follows. \square

2.2 Orthogonal decomposition of a chain of solitons

We now turn to the proof of Proposition 2. One key property in the proof is the exponential decay of the functions η_c and v_c . As a matter of fact, one can derive from formulae (5) that, given any integer p , there exists a positive constant K , depending only on p , such that

$$\sum_{0 \leq j+k \leq p} \left(|\partial_x^j \partial_c^k \eta_c(x)| + c^{1+2j+2k} |\partial_x^j \partial_c^k v_c(x)| \right) \leq K(2 - c^2) \exp\left(- (2 - c^2)^{\frac{1}{2}} |x|\right). \quad (2.9)$$

for any $0 < |c| < \sqrt{2}$ and $x \in \mathbb{R}$. A crucial consequence of (2.9) is the property that two solitons with same speed, as well as two sums of solitons with same speeds, can be closed in $X(\mathbb{R})$ only if their center(s) of mass are closed. More precisely, we have

Lemma 2.1. *Let $\mathbf{c} = (c_1, \dots, c_N) \in \text{Adm}_N$. Given a positive number δ , there exist two positive numbers β and M such that, if*

$$\|R_{\mathbf{c},\mathbf{a}} - R_{\mathbf{c},\mathbf{b}}\|_X < \beta,$$

for some positions $\mathbf{a} \in \text{Pos}_N(M)$ and $\mathbf{b} \in \text{Pos}_N(M)$, then,

$$\sum_{k=1}^N |a_k - b_k| < \delta.$$

Proof. The proof is by induction on the integer N . When $N = 1$, we have

Step 1. *Given a positive number δ , there exists a positive number α such that, if*

$$\|Q_{c,\mathbf{a}} - Q_{c,\mathbf{b}}\|_X < \alpha,$$

then,

$$|a - b| < \delta.$$

The proof is by contradiction. Assuming that Step 1 is false, there exist a positive number δ and a sequence $(a_n)_{n \in \mathbb{N}} \in \mathbb{R}^{\mathbb{N}}$ such that

$$\|Q_c(\cdot - a_n) - Q_c\|_X \rightarrow 0, \quad (2.10)$$

as $n \rightarrow +\infty$, and simultaneously, $|a_n| \geq \delta$ for any $n \in \mathbb{N}$. If the sequence $(a_n)_{n \in \mathbb{N}}$ were unbounded, then, up to some subsequence, it would tend to $+\infty$, or $-\infty$. In any case, taking the limit $n \rightarrow +\infty$ into (2.10), we would get that $2\|Q_c\|_X = 0$, which is not possible. As a

consequence, the sequence $(a_n)_{n \in \mathbb{N}}$ is bounded. Up to a subsequence, it converges to a real number a such that $|a| \geq \delta$, and simultaneously,

$$\|Q_c(\cdot - a) - Q_c\|_X = 0.$$

This identity implies that $Q_c(\cdot - na) = Q_c$ for any $n \in \mathbb{N}$. Since $a \neq 0$, this ensures, taking the limit $n \rightarrow +\infty$ and invoking (2.9), that $Q_c = 0$, which provides the desired contradiction.

Step 2. *End of the proof.*

We now assume that the conclusions of Lemma 2.1 are available up to the integer $N - 1$, and that they are not for the integer N . Then, there exist a positive number δ , a sequence $(L_j)_{j \in \mathbb{N}}$ tending to $+\infty$ as $j \rightarrow +\infty$, and two sequences $(\mathbf{a}^{(n,j)})_{(n,j) \in \mathbb{N}^2}$ and $(\mathbf{b}^{(n,j)})_{(n,j) \in \mathbb{N}^2}$ such that

$$\mathbf{a}^{(n,j)} \in \text{Pos}_N(L_j) \quad \text{and} \quad \mathbf{b}^{(n,j)} \in \text{Pos}_N(L_j),$$

for any $n \in \mathbb{N}$,

$$\|R_{\mathbf{c}, \mathbf{a}^{(n,j)}} - R_{\mathbf{c}, \mathbf{b}^{(n,j)}}\|_{X(\mathbb{R})} \rightarrow 0, \quad (2.11)$$

as $n \rightarrow +\infty$, and simultaneously,

$$\sum_{k=1}^N |a_k^{(n,j)} - b_k^{(n,j)}| > \delta. \quad (2.12)$$

Without loss of generality, we can moreover assume that $a_N^{(n,j)} \leq b_N^{(n,j)} = 0$.

In this situation, if the sequence $(a_N^{(n,j)})_{n \in \mathbb{N}}$ were unbounded for an integer j , then, up to a subsequence, it would tend to $-\infty$ as $n \rightarrow +\infty$. At this stage, we can invoke (2.9) to certify the existence of a universal constant K such that

$$\begin{aligned} \|R_{\mathbf{c}, \mathbf{a}^{(n,j)}} - R_{\mathbf{c}, \mathbf{b}^{(n,j)}}\|_{X(\mathbb{R})} &\geq \|R_{\mathbf{c}, \mathbf{a}^{(n,j)}} - R_{\mathbf{c}, \mathbf{b}^{(n,j)}}\|_{X((-L_j/2, +\infty))} \\ &\geq \|Q_{c_N}\|_{X(\mathbb{R})} - K \sum_{k=1}^N \frac{2 - c_k^2}{c_k} \left(\int_{|x| \geq L_j/2} \exp(-2(2 - c_k^2)^{\frac{1}{2}}|x|) dx \right)^{\frac{1}{2}}, \end{aligned}$$

for n large enough. In the limit $n \rightarrow +\infty$, we deduce combining with (2.11) that

$$\|Q_{c_N}\|_{X(\mathbb{R})} \leq K \sum_{k=1}^N \frac{(2 - c_k^2)^{\frac{3}{4}}}{c_k} \exp(- (2 - c_k^2)^{\frac{1}{2}} L_j),$$

which is again not possible when $j \rightarrow +\infty$. Hence, there exists an integer J such that, given any $j \geq J$, the sequence $(a_N^{(n,j)})_{n \in \mathbb{N}}$ is bounded. Up to some subsequence, it converges to a non-positive number α^j .

Assume next that $\alpha^j \leq -L_j/2$. For n large enough, we can rely on (2.9) to write

$$\begin{aligned} &\sum_{\ell=0}^1 |\partial_x^\ell \eta_{\mathbf{c}, \mathbf{a}^{(n,j)}}(x) - \partial_x^\ell \eta_{\mathbf{c}, \mathbf{b}^{(n,j)}}(x) + \partial_x^\ell \eta_{c_N}(x)| \\ &+ |v_{\mathbf{c}, \mathbf{a}^{(n,j)}}(x) - v_{\mathbf{c}, \mathbf{b}^{(n,j)}}(x) + v_{c_N}(x)| \leq K \sum_{k=1}^N \frac{2 - c_k^2}{c_k} \exp\left(- (2 - c_k^2)^{\frac{1}{2}} \left(x - \frac{3L_j}{8}\right)\right), \end{aligned}$$

for any $x \geq -L_j/8$. As above, this gives

$$\|Q_{c_N}\|_{X(\mathbb{R})} \leq \|R_{\mathbf{c}, \mathbf{a}^{(n,j)}} - R_{\mathbf{c}, \mathbf{b}^{(n,j)}}\|_{X(\mathbb{R})} + K \sum_{k=1}^N \frac{(2 - c_k^2)^{\frac{3}{4}}}{c_k} \exp\left(- \frac{(2 - c_k^2)^{\frac{1}{2}} L_j}{2}\right),$$

which is not possible at the limit $n \rightarrow +\infty$, when L_j is large enough. Therefore, we can assume, up to a further choice of J , that $\alpha^j \geq -L_j/2$ for any $j \geq J$.

Invoking once again (2.9), we next write

$$\begin{aligned} & \sum_{\ell=0}^1 |\partial_x^\ell \eta_{\mathbf{c}, \mathbf{a}^{(n,j)}}(x) - \partial_x^\ell \eta_{\mathbf{c}, \mathbf{b}^{(n,j)}}(x) - \partial_x^\ell \eta_{c_N, a_N^{(n,j)}}(x) + \partial_x^\ell \eta_{c_N}(x)| \\ & + |v_{\mathbf{c}, \mathbf{a}^{(n,j)}}(x) - v_{\mathbf{c}, \mathbf{b}^{(n,j)}}(x) - v_{c_N, a_N^{(n,j)}}(x) + v_{c_N}(x)| \leq K \sum_{k=1}^{N-1} \frac{2 - c_k^2}{c_k} \exp\left(- (2 - c_k^2)^{\frac{1}{2}}(x - L_j)\right), \end{aligned}$$

for any $x \geq -3L_j/4$, so that

$$\|Q_{c_N, a_N^{(n,j)}} - Q_{c_N}\|_{X(\mathbb{R})} \leq \|R_{\mathbf{c}, \mathbf{a}^{(n,j)}} - R_{\mathbf{c}, \mathbf{b}^{(n,j)}}\|_{X(\mathbb{R})} + K \sum_{k=1}^{N-1} \frac{(2 - c_k^2)^{\frac{3}{4}}}{c_k} \exp\left(- \frac{(2 - c_k^2)^{\frac{1}{2}} L_j}{2}\right). \quad (2.13)$$

The expression in the right-hand side of (2.13) tends to 0 when $n \rightarrow +\infty$ and $L_j \rightarrow +\infty$. Up to a further choice of J , one can invoke Step 1 to show that

$$|a_N^{(j,n)}| \leq \frac{\delta}{4}, \quad (2.14)$$

for any $j \geq J$ and any $n \geq n_j$. Moreover, one can rephrase (2.13) as

$$\begin{aligned} & \left\| \sum_{k=1}^{N-1} (Q_{c_k, a_k^{(n,j)}} - Q_{c_k, b_k^{(n,j)}}) \right\|_{X(\mathbb{R})} \\ & \leq 2 \|R_{\mathbf{c}, \mathbf{a}^{(n,j)}} - R_{\mathbf{c}, \mathbf{b}^{(n,j)}}\|_{X(\mathbb{R})} + K \sum_{k=1}^N \frac{(2 - c_k^2)^{\frac{3}{4}}}{c_k} \exp\left(- \frac{(2 - c_k^2)^{\frac{1}{2}} L_j}{2}\right). \end{aligned} \quad (2.15)$$

Once again, the expression in the right-hand side of (2.15) tends to 0 when $n \rightarrow +\infty$ and $L_j \rightarrow +\infty$. Since Lemma 2.1 is true for the integer $N - 1$, we deduce that

$$\sum_{k=1}^{N-1} |a_k^{(j,n)} - b_k^{(j,n)}| \leq \frac{\delta}{4},$$

for any j and n large enough. Combining with (2.14), we obtain a contradiction with (2.12). Lemma 2.1 follows by induction on the integer N . \square

Estimate (2.9) and Lemma 2.1 are enough to consider the proof of Proposition 2.

Proof of Proposition 2. The proof is reminiscent from [16]. It relies on the quantified version of the implicit function theorem provided by Proposition A.1. We consider the map Ξ defined by

$$\Xi(\eta, v, \sigma, \mathbf{b}) := \left(\langle \varepsilon, \partial_x Q_{\sigma_1, b_1} \rangle_{L^2}, \dots, \langle \varepsilon, \partial_x Q_{\sigma_N, b_N} \rangle_{L^2}, P'(Q_{\sigma_1, b_1})(\varepsilon), \dots, P'(Q_{\sigma_N, b_N})(\varepsilon) \right).$$

In order to simplify the notation, we have set here $\varepsilon := (\eta, v) - R_{\sigma, \mathbf{b}}$. The map Ξ is well-defined and smooth from $X(\mathbb{R}) \times \text{Adm}_N \times \mathbb{R}^N$ to \mathbb{R}^{2N} . Moreover, it fulfills the assumptions of Proposition A.1 so that we can state

Step 1. Let $0 < \tau < 1$ and set

$$\text{Adm}_N(\tau) := \{\mathbf{c} \in (-\sqrt{2}, \sqrt{2})^N, \text{ s.t. } \mu_{\mathbf{c}} > \tau \text{ and } \nu_{\mathbf{c}} > \tau\}.$$

There exist positive numbers δ , Λ and M , depending only on τ , such that, given any $(\mathbf{c}, \mathbf{a}) \in \text{Adm}_N(\tau) \times \text{Pos}_N(M)$, there exists a map $\gamma_{\mathbf{c}, \mathbf{a}} \in \mathcal{C}^1(B(R_{\mathbf{c}, \mathbf{a}}, \delta), \text{Adm}_N(\tau/2) \times \mathbb{R}^N)$ such that, given any $(\eta, v) \in B(R_{\mathbf{c}, \mathbf{a}}, \delta)$, $(\sigma, \mathbf{b}) = \gamma_{\mathbf{c}, \mathbf{a}}(\eta, v)$ is the unique solution in $B((\mathbf{c}, \mathbf{a}), \Lambda\delta)$ of the equation

$$\Xi(\eta, v, \sigma, \mathbf{b}) = 0.$$

Moreover, the map $\gamma_{\mathbf{c}, \mathbf{a}}$ is Lipschitz on $B(R_{\mathbf{c}, \mathbf{a}}, \delta)$, with Lipschitz constant at most Λ .

For $\mathbf{c} \in \text{Adm}_N(\tau)$ and $\mathbf{a} \in \mathbb{R}^N$, we check that

$$\Xi(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a}) = 0.$$

Similarly, we compute

$$\begin{cases} \partial_{\sigma_j} \Xi_k(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a}) = -\langle \partial_{\mathbf{c}} Q_{c_j, a_j}, \partial_x Q_{c_k, a_k} \rangle_{L^2}, \\ \partial_{b_j} \Xi_k(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a}) = \langle \partial_x Q_{c_j, a_j}, \partial_x Q_{c_k, a_k} \rangle_{L^2}, \end{cases}$$

as well as

$$\begin{cases} \partial_{\sigma_j} \Xi_{N+k}(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a}) = -P'(Q_{c_k, a_k})(\partial_{\mathbf{c}} Q_{c_j, a_j}), \\ \partial_{b_j} \Xi_{N+k}(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a}) = P'(Q_{c_k, a_k})(\partial_x Q_{c_j, a_j}), \end{cases}$$

for any $1 \leq j, k \leq N$. When $j = k$, we infer from formulae (5) that

$$\partial_{\sigma_k} \Xi_k(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a}) = \partial_{a_k} \Xi_{N+k}(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a}) = 0, \quad (2.16)$$

and

$$\partial_{a_k} \Xi_k(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a}) = \|\partial_x Q_{c_k}\|_{L^2}^2 = \frac{(2 - c_k^2)^{\frac{3}{2}}}{3} > 0,$$

whereas by (18),

$$\partial_{\sigma_k} \Xi_{N+k}(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a}) = -\frac{d}{dc}(P(Q_c))|_{c=c_k} = -(2 - c_k^2)^{\frac{1}{2}} < 0.$$

In particular, the diagonal matrix $A_{\mathbf{c}}$ with the same diagonal elements as $d_{\sigma, \mathbf{b}} \Xi(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a})$ is a continuous isomorphism from \mathbb{R}^{2N} to \mathbb{R}^{2N} , with operator norm bounded from below by $\tau^3/3$.

On the other hand, when $j \neq k$, it follows from (2.9) that

$$\begin{aligned} & \left| \langle \partial_{\mathbf{c}} Q_{c_j, a_j}, \partial_x Q_{c_k, a_k} \rangle_{L^2} \right| + \left| \langle \partial_x Q_{c_j, a_j}, \partial_x Q_{c_k, a_k} \rangle_{L^2} \right| \\ & \quad + \left| P'(Q_{c_k, a_k})(\partial_{\mathbf{c}} Q_{c_j, a_j}) \right| + \left| P'(Q_{c_k, a_k})(\partial_x Q_{c_j, a_j}) \right| \\ & \leq K^2 (2 - c_j^2)^{\frac{1}{2}} (2 - c_k^2)^{\frac{1}{2}} \left(1 + \frac{1}{c_j^3 c_k^3} \right) \int_{\mathbb{R}} \exp(-\nu_{j,k}(|x - a_j| + |x - a_k|)) dx \\ & = K^2 (2 - c_j^2)^{\frac{1}{2}} (2 - c_k^2)^{\frac{1}{2}} \left(1 + \frac{1}{c_j^3 c_k^3} \right) \left(|a_j - a_k| + \frac{1}{\nu_{j,k}} \right) \exp(-\nu_{j,k}|a_j - a_k|), \end{aligned}$$

where we have set $\nu_{j,k} := \min\{(2 - c_j^2)^{\frac{1}{2}}, (2 - c_k^2)^{\frac{1}{2}}\}$. Assuming that $\mathbf{a} \in \text{Pos}_N(L)$ for some positive number L , there exists a positive constant K_{τ} , depending only on τ , such that

$$\begin{aligned} & \left| \partial_{\sigma_j} \Xi_k(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a}) \right| + \left| \partial_{a_j} \Xi_k(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a}) \right| \\ & \quad + \left| \partial_{\sigma_j} \Xi_{N+k}(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a}) \right| + \left| \partial_{a_j} \Xi_{N+k}(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a}) \right| \leq K_{\tau} (L + 1) \exp(-\nu_{\mathbf{c}} L). \end{aligned}$$

Combining with (2.16), we deduce when L is large enough, that

$$d_{\sigma, \mathbf{b}} \Xi(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a}) := A_{\mathbf{c}}(Id + H(\mathbf{c}, \mathbf{a})),$$

where the operator norm of $H(\mathbf{c}, \mathbf{a})$ is less than $1/2$.

We now turn to the differential

$$d_{\eta, v} \Xi(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a})(f, g) := \left(\langle (f, g), \partial_x Q_{c_1, a_1} \rangle_{L^2}, \dots, \langle (f, g), \partial_x Q_{c_N, a_N} \rangle_{L^2}, \right. \\ \left. P'(Q_{c_1, a_1})(f, g), \dots, P'(Q_{c_N, a_N})(f, g) \right).$$

Since the operator norm of $A_{\mathbf{c}}^{-1}$ is bounded by $3/\tau^3$, we infer from (2.9) that $d_{\eta, v} \Xi(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a})$ may be written as

$$d_{\eta, v} \Xi(R_{\mathbf{c}, \mathbf{a}}, \mathbf{c}, \mathbf{a}) = A_{\mathbf{c}} T_{\mathbf{c}, \mathbf{a}},$$

where $T_{\mathbf{c}, \mathbf{a}}$ is a continuous linear mapping from $X(\mathbb{R})$ to \mathbb{R}^N with operator norm depending only on τ .

Finally, again by (2.9), the operator norm of the second order differential $d^2 \Xi(\eta, v, \sigma, \mathbf{b})$ is bounded by a constant K_{τ} , depending only on τ , when $(\eta, v, \sigma, \mathbf{b}) \in X(\mathbb{R}) \times \text{Adm}_N(\tau/2) \times \mathbb{R}^N$. It then remains to notice that Assumption (iv) of Proposition A.1 is satisfied when $U = \text{Adm}_N(\tau)$ and $V = \text{Adm}_N(\tau/2)$, and to apply Proposition A.1 to the map Ξ , in order to obtain the statements in Step 1.

Step 2. *End of the proof.*

Let $\tau = \min\{\mu_{\mathbf{c}}^*/2, \nu_{\mathbf{c}}^*/2\}$. We denote δ_1 , Λ_1 and M_1 , the constants provided by Step 1, and β_1 , the number provided by Lemma 2.1 for $\delta = \Lambda_1 \delta_1/3$. We set $\alpha_1 := \min\{\delta_1/3, \beta_1/4\}$ and $L_1 := M_1$. When $(\eta, v) \in \mathcal{U}_{\mathbf{c}^*}(\alpha_1, L_1)$, there exists $\mathbf{b} \in \text{Pos}_N(L_1)$ such that $(\eta, v) \in B(R_{\mathbf{c}^*, \mathbf{b}}, 2\alpha_1)$. Since $2\alpha_1 \leq \delta_1$ and $L_1 = M_1$, we deduce from Step 1 that the numbers \mathbf{c} and \mathbf{a} given by

$$(\mathbf{c}, \mathbf{a}) = \gamma_{\mathbf{c}^*, \mathbf{b}}(\eta, v),$$

are well-defined, so that we can set $\mathfrak{C}(\eta, v) = \mathbf{c}$ and $\mathfrak{A}(\eta, v) = \mathbf{a}$. We claim that \mathbf{c} and \mathbf{a} do not depend on the choice of $\mathbf{b} \in \text{Pos}_N(L_1)$ such that $(\eta, v) \in B(R_{\mathbf{c}^*, \mathbf{b}}, 2\alpha_1)$. As a consequence, the maps \mathfrak{C} and \mathfrak{A} are well-defined on $\mathcal{U}_{\mathbf{c}^*}(\alpha_1, L_1)$ with values in $\text{Adm}_N(\tau/2) \subset \text{Adm}_N$, respectively, \mathbb{R}^N .

Indeed, given another choice of $\mathbf{b}_2 \in \text{Pos}_N(L_1)$ such that $(\eta, v) \in B(R_{\mathbf{c}^*, \mathbf{b}_2}, 2\alpha_1)$, we have

$$\|R_{\mathbf{c}^*, \mathbf{b}} - R_{\mathbf{c}^*, \mathbf{b}_2}\|_X < 4\alpha_1 \leq \beta_1,$$

so that, by Lemma 2.1,

$$|\mathbf{b} - \mathbf{b}_2| < \frac{\Lambda_1 \delta_1}{3}. \quad (2.17)$$

On the other hand, the map $\gamma_{\mathbf{c}^*, \mathbf{b}}$ is Lipschitz on $B(R_{\mathbf{c}^*, \mathbf{b}}, 2\alpha_1)$, with Lipschitz constant at most Λ_1 . Hence,

$$|(\mathbf{c}, \mathbf{a}) - (\mathbf{c}^*, \mathbf{b})| = |\gamma_{\mathbf{c}^*, \mathbf{b}}(\eta, v) - \gamma_{\mathbf{c}^*, \mathbf{b}}(R_{\mathbf{c}^*, \mathbf{b}})| \leq \Lambda_1 \alpha_1 \leq \frac{\Lambda_1 \delta_1}{3}.$$

Combining with (2.17), we obtain

$$|(\mathbf{c}, \mathbf{a}) - (\mathbf{c}^*, \mathbf{b}_2)| < \Lambda_1 \delta_1.$$

Since $\Xi(\eta, v, \mathbf{c}, \mathbf{a}) = 0$, we deduce from Step 1 that $(\mathbf{c}, \mathbf{a}) = \gamma_{\mathbf{c}^*, \mathbf{b}_2}(\eta, v)$, so that \mathbf{c} and \mathbf{a} do not depend on the choice of $\mathbf{b} \in \text{Pos}_N(L_1)$ such that $(\eta, v) \in B(R_{\mathbf{c}^*, \mathbf{b}}, 2\alpha_1)$.

Concerning the smoothness of \mathfrak{C} and \mathfrak{A} , we consider $(\mu, w) \in \mathcal{U}_{\mathbf{c}^*}(\alpha_1, L_1)$ such that

$$\|(\mu, w) - (\eta, v)\|_X < \alpha_1.$$

Given $\mathbf{b} \in \text{Pos}_N(L_1)$ such that $(\eta, v) \in B(R_{\mathbf{c}^*, \mathbf{b}}, \alpha_1)$, we deduce that

$$\|(\mu, w) - R_{\mathbf{c}^*, \mathbf{b}}\|_X < 2\alpha_1,$$

so that, by definition,

$$(\mathfrak{C}(\eta, v), \mathfrak{A}(\eta, v)) = \gamma_{\mathbf{c}^*, \mathbf{b}}(\eta, v) \quad \text{and} \quad (\mathfrak{C}(\mu, w), \mathfrak{A}(\mu, w)) = \gamma_{\mathbf{c}^*, \mathbf{b}}(\mu, w).$$

At this stage, recall that the map $\gamma_{\mathbf{c}^*, \mathbf{b}}$ is of class \mathcal{C}^1 on $B(R_{\mathbf{c}^*, \mathbf{b}}, 2\alpha_1)$. As a consequence, the functions \mathfrak{C} and \mathfrak{A} are also of class \mathcal{C}^1 on $B((\eta, v), \alpha_1) \cap \mathcal{U}_{\mathbf{c}^*}(\alpha_1, L_1)$, therefore, on $\mathcal{U}_{\mathbf{c}^*}(\alpha_1, L_1)$.

We next consider the proof of (25). When $(\eta, v) \in B(R_{\mathbf{c}^*, \mathbf{a}^*}, \alpha)$, with $\mathbf{a}^* \in \text{Pos}_N(L)$, $\alpha \leq \alpha_1$ and $L \geq L_1$, the values of $\mathfrak{C}(\eta, v)$ and $\mathfrak{A}(\eta, v)$ are equal to

$$(\mathfrak{C}(\eta, v), \mathfrak{A}(\eta, v)) = \gamma_{\mathbf{c}^*, \mathbf{a}^*}(\eta, v).$$

In view of the Lipschitz continuity on $B(R_{\mathbf{c}^*, \mathbf{a}^*}, \alpha)$ of the map $\gamma_{\mathbf{c}^*, \mathbf{a}^*}$, we infer that

$$|\mathbf{c} - \mathbf{c}^*| + |\mathbf{a} - \mathbf{a}^*| \leq \Lambda_1 \|(\eta, v) - R_{\mathbf{c}^*, \mathbf{a}^*}\|_X \leq \Lambda_1 \alpha. \quad (2.18)$$

On the other hand, letting $\varepsilon = (\eta, v) - R_{\mathbf{c}, \mathbf{a}}$, we have

$$\|\varepsilon\|_X \leq \|(\eta, v) - R_{\mathbf{c}^*, \mathbf{a}^*}\|_X + \|R_{\mathbf{c}^*, \mathbf{a}^*} - R_{\mathbf{c}, \mathbf{a}}\|_X < \alpha + \|R_{\mathbf{c}^*, \mathbf{a}^*} - R_{\mathbf{c}, \mathbf{a}}\|_X. \quad (2.19)$$

In view of formulae (5), there exists a universal constant K such that

$$\begin{aligned} \|R_{\mathbf{c}^*, \mathbf{a}^*} - R_{\mathbf{c}, \mathbf{a}}\|_X &\leq \int_0^1 \left\| \partial_\sigma R_{\mathbf{c}+t(\mathbf{c}^*-\mathbf{c}), \mathbf{a}+t(\mathbf{a}^*-\mathbf{a})}(\mathbf{c}^* - \mathbf{c}) + \partial_{\mathbf{a}}^* R_{\mathbf{c}+t(\mathbf{c}^*-\mathbf{c}), \mathbf{a}+t(\mathbf{a}^*-\mathbf{a})}(\mathbf{a}^* - \mathbf{a}) \right\|_X dt \\ &\leq K |(\mathbf{c}^*, \mathbf{a}^*) - (\mathbf{c}, \mathbf{a})|. \end{aligned}$$

Combining with (2.18) and (2.19), we are led to

$$\|\varepsilon\|_X < (1 + K\Lambda_1)\alpha.$$

In view of (2.18), it is sufficient to set $K_1 := 1 + (K + 1)\Lambda_1$ in order to derive (25).

Finally, conditions (24) are direct consequences of the definitions of the maps $\gamma_{\mathbf{c}^*, \mathbf{b}}$. This completes the proof of Proposition 2. \square

Corollary 1 is a direct consequence of Proposition 2.

Proof of Corollary 1. Let $(\eta, v) \in \mathcal{U}_{\mathbf{c}^*}(\alpha_1, L)$, with $L \geq L_1$, and set $(\mathbf{c}, \mathbf{a}) := (\mathfrak{C}(\eta, v), \mathfrak{A}(\eta, v))$ as in the proof of Proposition 2. In view of (25), one can decrease α_1 so that

$$|\mathbf{c} - \mathbf{c}^*| \leq \min \left\{ \frac{\mu_{\mathbf{c}^*}}{2}, \frac{\nu_{\mathbf{c}^*}^2}{8\sqrt{2}} \right\},$$

which is enough to obtain (28) and (29).

One can next increase L_1 so that $L_1 > 2$, and decrease again α_1 so that $K_1\alpha_1 \leq 1/2$. In this case, we can deduce from (25) that

$$|\mathbf{a} - \mathbf{a}^*| \leq \frac{1}{2},$$

for any $\mathbf{a}^* \in \text{Pos}_N(L)$ such that

$$\|(\eta, v) - R_{\mathbf{c}^*, \mathbf{a}^*}\|_X \leq \alpha.$$

It follows that \mathbf{a} belongs to $\text{Pos}_N(L-1)$ as mentioned in (27).

Finally, one can infer from formula (5) that

$$\min \left\{ 1 - \eta_c(x), x \in \mathbb{R} \right\} = 1 - \eta_c(0) = \frac{c^2}{2},$$

for any $c \in (-\sqrt{2}, \sqrt{2})$. Combining with the exponential decay of the functions η_c provided by (2.9), we deduce that we can enlarge again L_1 so that

$$\min \left\{ 1 - \eta_{\mathbf{c}, \mathbf{a}}(x), x \in \mathbb{R} \right\} \geq \frac{\mu_{\mathbf{c}^*}^2}{4}, \quad (2.20)$$

when \mathbf{c} satisfies (29) and $\mathbf{a} \in \text{Pos}_N(L_1-1)$. On the other hand, we can a last time decrease α_1 so that, combining (25) with the Sobolev embedding theorem, we have

$$\|\varepsilon_\eta\|_{L^\infty(\mathbb{R})} \leq \frac{\mu_{\mathbf{c}^*}^2}{8}.$$

Since $\eta = \eta_{\mathbf{c}, \mathbf{a}} + \varepsilon_\eta$, this is sufficient to guarantee that η satisfies (26), so that the pair (η, v) belongs to $\mathcal{NV}(\mathbb{R})$. This completes the proof of Corollary 1. \square

2.3 Almost minimizing properties close to a sum of solitons

The main goal of this subsection is to establish the almost minimizing properties close to a sum of solitons stated in Subsection C of the introduction. When establishing the estimates of E and P_k given by (33) and (34), we will make use of the following elementary inequality to bound various interaction terms using (2.9).

Lemma 2.2. *Let $(a, b) \in \mathbb{R}^2$, with $a < b$, $(\nu_a, \nu_b) \in (0, +\infty)^2$, and set $y^\pm := \max\{\pm y, 0\}$. Then,*

$$\left\| \exp(-\nu_a(\cdot - a)^+) \exp(-\nu_b(\cdot - b)^-) \right\|_{L^p} \leq \left(\frac{2}{p \min\{\nu_a, \nu_b\}} + b - a \right)^{\frac{1}{p}} \exp(-\min\{\nu_a, \nu_b\}(b - a)),$$

for any $1 \leq p \leq +\infty$.

We will also use the following pointwise estimates on the functions Φ_k , $\Phi_{k, k+1}$ and Ψ_k .

Lemma 2.3. *Let $1 \leq k \leq N$ and $x \in \mathbb{R}$. We have*

$$\Phi_k(x) \leq \exp\left(-2\tau\left(|x - a_k| - \frac{L_1}{4}\right)^+\right),$$

and

$$|1 - \Phi_k(x)| \leq \exp\left(-2\tau\left(x - a_k + \frac{L_1}{4}\right)^+\right) + \exp\left(-2\tau\left(x - a_k - \frac{L_1}{4}\right)^-\right).$$

Similarly, for $1 \leq k \leq N-1$,

$$\begin{aligned} \Phi_{0,1}(x) &\leq \exp\left(-2\tau\left(x - a_1 + \frac{L_1}{4}\right)^+\right), \\ \Phi_{k, k+1}(x) &\leq \exp\left(-2\tau\left(x - a_k - \frac{L_1}{4}\right)^-\right) \exp\left(-2\tau\left(x - a_{k+1} + \frac{L_1}{4}\right)^+\right), \\ \Phi_{N, N+1}(x) &\leq \exp\left(-2\tau\left(x - a_N - \frac{L_1}{4}\right)^-\right), \end{aligned}$$

while, for $2 \leq k \leq N-1$,

$$\begin{aligned}\Psi_1(x) - \Psi_2(x) &\leq \exp\left(-\frac{\nu_{\mathbf{c}^*}}{8}\left(x - \frac{a_1 + a_2}{2}\right)^+\right), \\ \Psi_k(x) - \Psi_{k+1}(x) &\leq \exp\left(-\frac{\nu_{\mathbf{c}^*}}{8}\left(x - \frac{a_{k-1} + a_k}{2}\right)^-\right) \exp\left(-\frac{\nu_{\mathbf{c}^*}}{8}\left(x - \frac{a_{k-1} + a_k}{2}\right)^+\right), \\ \Psi_N(x) - \Psi_{N+1}(x) &\leq \exp\left(-\frac{\nu_{\mathbf{c}^*}}{8}\left(x - \frac{a_{N-1} + a_N}{2}\right)^+\right),\end{aligned}$$

and

$$\begin{aligned}|1 - \Psi_1(x) + \Psi_2(x)| &\leq \exp\left(-\frac{\nu_{\mathbf{c}^*}}{8}\left(x - \frac{a_1 + a_2}{2}\right)^-\right), \\ |1 - \Psi_k(x) + \Psi_{k+1}(x)| &\leq \exp\left(-\frac{\nu_{\mathbf{c}^*}}{8}\left(x - \frac{a_{k-1} + a_k}{2}\right)^+\right) + \exp\left(-\frac{\nu_{\mathbf{c}^*}}{8}\left(x - \frac{a_{k-1} + a_k}{2}\right)^-\right), \\ |1 - \Psi_N(x) + \Psi_{N+1}(x)| &\leq \exp\left(-\frac{\nu_{\mathbf{c}^*}}{8}\left(x - \frac{a_{N-1} + a_N}{2}\right)^+\right).\end{aligned}$$

Proof. The estimates in Lemma 2.3 follow from the inequality

$$|1 - \operatorname{sign}(x)\operatorname{th}(x)| \leq 2\exp(-2|x|),$$

which holds for any $x \in \mathbb{R}$. \square

We are now in position to present the

Proof of Proposition 3. We begin with (33). We expand the energy E according to the Taylor rule by

$$E(\eta, v) = E(R_{\mathbf{c}, \mathbf{a}} + \varepsilon) = E(R_{\mathbf{c}, \mathbf{a}}) + E'(R_{\mathbf{c}, \mathbf{a}})(\varepsilon) + \frac{1}{2}E''(R_{\mathbf{c}, \mathbf{a}})(\varepsilon, \varepsilon) + \mathcal{R}(\eta, v),$$

where

$$E(R_{\mathbf{c}, \mathbf{a}}) = \int_{\mathbb{R}} \left(\frac{(\partial_x \eta_{\mathbf{c}, \mathbf{a}})^2}{8(1 - \eta_{\mathbf{c}, \mathbf{a}})} + \frac{1}{2}(1 - \eta_{\mathbf{c}, \mathbf{a}})v_{\mathbf{c}, \mathbf{a}}^2 + \frac{1}{4}\eta_{\mathbf{c}, \mathbf{a}}^2 \right), \quad (2.21)$$

$$E'(R_{\mathbf{c}, \mathbf{a}})(\varepsilon) = \int_{\mathbb{R}} \left(\frac{\partial_x \eta_{\mathbf{c}, \mathbf{a}} \partial_x \varepsilon_\eta}{4(1 - \eta_{\mathbf{c}, \mathbf{a}})} + \frac{(\partial_x \eta_{\mathbf{c}, \mathbf{a}})^2 \varepsilon_\eta}{8(1 - \eta_{\mathbf{c}, \mathbf{a}})} + (1 - \eta_{\mathbf{c}, \mathbf{a}})v_{\mathbf{c}, \mathbf{a}} \varepsilon_v + \frac{1}{2}(\eta_{\mathbf{c}, \mathbf{a}} - v_{\mathbf{c}, \mathbf{a}}^2) \varepsilon_\eta \right), \quad (2.22)$$

$$E''(R_{\mathbf{c}, \mathbf{a}})(\varepsilon, \varepsilon) = \int_{\mathbb{R}} \left(\frac{(\partial_x \varepsilon_\eta)^2}{4(1 - \eta_{\mathbf{c}, \mathbf{a}})} + \frac{\partial_x \eta_{\mathbf{c}, \mathbf{a}} \varepsilon_\eta \partial_x \varepsilon_\eta}{2(1 - \eta_{\mathbf{c}, \mathbf{a}})^2} + \frac{(\partial_x \eta_{\mathbf{c}, \mathbf{a}})^2 \varepsilon_\eta^2}{4(1 - \eta_{\mathbf{c}, \mathbf{a}})^3} - 2v_{\mathbf{c}, \mathbf{a}} \varepsilon_\eta \varepsilon_v + (1 - \eta_{\mathbf{c}, \mathbf{a}}) \varepsilon_v^2 + \frac{\varepsilon_\eta^2}{2} \right), \quad (2.23)$$

and

$$\mathcal{R}(\eta, v) := \frac{1}{8} \int_{\mathbb{R}} \left(\frac{(\partial_x \eta_{\mathbf{c}, \mathbf{a}})^2 \varepsilon_\eta^3}{(1 - \eta)(1 - \eta_{\mathbf{c}, \mathbf{a}})^3} + \frac{2\partial_x \eta_{\mathbf{c}, \mathbf{a}} \partial_x \varepsilon_\eta \varepsilon_\eta^2}{(1 - \eta)(1 - \eta_{\mathbf{c}, \mathbf{a}})^2} + \frac{(\partial_x \varepsilon_\eta)^2 \varepsilon_\eta}{(1 - \eta)(1 - \eta_{\mathbf{c}, \mathbf{a}})} - 4\varepsilon_\eta \varepsilon_v^2 \right).$$

Following Corollary 1, we may assume that (26), (27), (28) and (29), as well as estimate (2.20), hold. From (2.9) and (29), we obtain

$$|\eta_{c_k, a_k}(x)| + |\partial_x \eta_{c_k, a_k}(x)| + |v_{c_k, a_k}(x)| = \mathcal{O}\left(\exp\left(-\frac{\nu_{\mathbf{c}^*}|x - a_k|}{2}\right)\right), \quad (2.24)$$

for any $x \in \mathbb{R}$ and any $1 \leq k \leq N$. From (26), (2.20), (2.24), and the expression (2.21) for $E(R_{\mathbf{c}, \mathbf{a}})$, we compute using Lemma 2.2 with $p = 1$,

$$E(R_{\mathbf{c}, \mathbf{a}}) = \sum_{k=1}^N E(Q_{c_k}) + \mathcal{O}\left(L \exp\left(-\frac{\nu_{\mathbf{c}^*}L}{2}\right)\right). \quad (2.25)$$

Similarly, expression (2.22) for E' may be simplified using the Hölder inequality and Lemma 2.2 with $p = 2$ as

$$E'(R_{\mathbf{c},\mathbf{a}})(\varepsilon) = \sum_{k=1}^N E'(Q_{c_k, a_k})(\varepsilon) + \mathcal{O}\left(L^{\frac{1}{2}} \exp\left(-\frac{\nu_{\mathbf{c}^*} L}{2}\right) \|\varepsilon\|_X\right).$$

Since $\|\varepsilon\|_X \leq K_1 \alpha_1$ by (25) and since

$$E'(Q_{c_k, a_k})(\varepsilon) = E'(Q_{c_k, a_k})(\varepsilon) + c_k P'(Q_{c_k, a_k})(\varepsilon) = 0,$$

by (11) and (24), we deduce that

$$E'(R_{\mathbf{c},\mathbf{a}})(\varepsilon) = \mathcal{O}\left(L^{\frac{1}{2}} \exp\left(-\frac{\nu_{\mathbf{c}^*} L}{2}\right)\right). \quad (2.26)$$

From (26), (2.20) and (2.24), as well as the Sobolev embedding theorem, we also infer that

$$\mathcal{R}(\eta, v) = \mathcal{O}(\|\varepsilon\|_X^3). \quad (2.27)$$

We finally turn to the term $E''(R_{\mathbf{c},\mathbf{a}})(\varepsilon, \varepsilon)$ which requires somewhat more care. We rewrite expression (2.23) as

$$E''(R_{\mathbf{c},\mathbf{a}})(\varepsilon, \varepsilon) := \int_{\mathbb{R}} J(R_{\mathbf{c},\mathbf{a}}, \varepsilon),$$

and we decompose according to the partition of unity (30) as

$$\int_{\mathbb{R}} J(R_{\mathbf{c},\mathbf{a}}, \varepsilon) = \sum_{k=1}^N \int_{\mathbb{R}} J(R_{\mathbf{c},\mathbf{a}}, \varepsilon) \Phi_k + \sum_{k=0}^N \int_{\mathbb{R}} J(R_{\mathbf{c},\mathbf{a}}, \varepsilon) \Phi_{k,k+1}.$$

Invoking once more the decay estimates (2.24), we obtain from Lemma 2.2 with $p = +\infty$, and Lemma 2.3,

$$\int_{\mathbb{R}} J(R_{\mathbf{c},\mathbf{a}}, \varepsilon) \Phi_k = \int_{\mathbb{R}} J(Q_{c_k, a_k}, \varepsilon) \Phi_k + \mathcal{O}\left(\exp\left(-\frac{\tau L_1}{2}\right) \|\varepsilon\|_X^2\right). \quad (2.28)$$

Taking into account the formula

$$\partial_x(\Phi_k^{\frac{1}{2}} \varepsilon_\eta) = \Phi_k^{\frac{1}{2}} \partial_x \varepsilon_\eta + \frac{\partial_x \Phi_k}{2\Phi_k^{\frac{1}{2}}} \varepsilon_\eta,$$

we obtain, after a change of variable corresponding to a translation by a_k ,

$$\begin{aligned} & \int_{\mathbb{R}} J(Q_{c_k, a_k}, \varepsilon) \Phi_k \\ &= \int_{\mathbb{R}} J(Q_{c_k}, \varepsilon_k) - \int_{\mathbb{R}} \left(\frac{\partial_x \Phi_k \varepsilon_\eta^2}{32\Phi_k(1-\eta_{c_k, a_k})} + \frac{\varepsilon_\eta \partial_x \varepsilon_\eta}{8(1-\eta_{c_k, a_k})} + \frac{\partial_x \eta_{c_k, a_k} \varepsilon_\eta^2}{8(1-\eta_{c_k, a_k})^2} \right) \partial_x \Phi_k. \end{aligned} \quad (2.29)$$

Since $|\partial_x \Phi_k| = \mathcal{O}(\tau)$ and $|\partial_x \Phi_k|/|\Phi_k| = \mathcal{O}(\tau)$, we deduce from (2.29),

$$\int_{\mathbb{R}} J(Q_{c_k, a_k}, \varepsilon) \Phi_k = E''(Q_{c_k})(\varepsilon_k, \varepsilon_k) + \mathcal{O}(\tau \|\varepsilon\|_X^2). \quad (2.30)$$

In a similar manner, we write using Lemma 2.2 with $p = +\infty$, and Lemma 2.3,

$$\int_{\mathbb{R}} J(R_{\mathbf{c},\mathbf{a}}, \varepsilon) \Phi_{k,k+1} = \int_{\mathbb{R}} J(0, \varepsilon) \Phi_{k,k+1} + \mathcal{O}\left(\exp\left(-\frac{\tau L_1}{2}\right) \|\varepsilon\|_X^2\right),$$

and then

$$\int_{\mathbb{R}} J(0, \varepsilon) \Phi_{k,k+1} = \int_{\mathbb{R}} J(0, \varepsilon_{k,k+1}) + \mathcal{O}(\tau \|\varepsilon\|_X^2) = E''(0)(\varepsilon_{k,k+1}, \varepsilon_{k,k+1}) + \mathcal{O}(\tau \|\varepsilon\|_X^2).$$

Estimate (33) follows combining with (2.25), (2.26), (2.27), (2.28) and (2.30).

We turn now to (34). We expand the functional P_k , which is exactly quadratic, as

$$P_k(\eta, v) = P_k(R_{c,a} + \varepsilon) = P_k(R_{c,a}) + P'_k(R_{c,a})(\varepsilon) + \frac{1}{2} P''_k(R_{c,a})(\varepsilon, \varepsilon),$$

where

$$P_k(R_{c,a}) = \frac{1}{2} \int_{\mathbb{R}} \eta_{c,a} v_{c,a} (\Psi_k - \Psi_{k+1}), \quad (2.31)$$

$$P'_k(R_{c,a})(\varepsilon) = \frac{1}{2} \int_{\mathbb{R}} (\eta_{c,a} \varepsilon_v + v_{c,a} \varepsilon_\eta) (\Psi_k - \Psi_{k+1}), \quad (2.32)$$

$$P''_k(R_{c,a})(\varepsilon, \varepsilon) = \int_{\mathbb{R}} \varepsilon_\eta \varepsilon_v (\Psi_k - \Psi_{k+1}). \quad (2.33)$$

Invoking here again (2.24), Lemma 2.2 with $p = 1$, and Lemma 2.3, we estimate (2.31) as

$$P_k(R_{c,a}) = \sum_{j=1}^N P_k(Q_{c_j, a_j}) + \mathcal{O}\left(L \exp\left(-\frac{\nu_{c^*} L}{2}\right)\right).$$

For $j \neq k$, we obtain from Lemma 2.2 with $p = 1$, and Lemma 2.3,

$$P_k(Q_{c_j, a_j}) = \mathcal{O}\left(L \exp\left(-\frac{\nu_{c^*} L}{16}\right)\right),$$

while

$$P_k(Q_{c_k, a_k}) = P(Q_{c_k}) + \mathcal{O}\left(L \exp\left(-\frac{\nu_{c^*} L}{16}\right)\right).$$

Concerning (2.32), we obtain similarly from Lemma 2.2 with $p = 2$, and Lemma 2.3,

$$P'_k(R_{c,a})(\varepsilon) = P'(Q_{c_k, a_k})(\varepsilon) + \mathcal{O}\left(L^{\frac{1}{2}} \exp\left(-\frac{\nu_{c^*} L}{16}\right)\right) = \mathcal{O}\left(L^{\frac{1}{2}} \exp\left(-\frac{\nu_{c^*} L}{16}\right)\right),$$

where we have used (24) for the last equality.

Finally, we decompose the integral in (2.33) as

$$\int_{\mathbb{R}} \varepsilon_\eta \varepsilon_v (\Psi_k - \Psi_{k+1}) = \sum_{j=1}^N \int_{\mathbb{R}} \varepsilon_\eta \varepsilon_v (\Psi_k - \Psi_{k+1}) \Phi_j + \sum_{j=0}^N \int_{\mathbb{R}} \varepsilon_\eta \varepsilon_v (\Psi_k - \Psi_{k+1}) \Phi_{j,j+1}.$$

When $j \neq k$, we obtain

$$\int_{\mathbb{R}} \varepsilon_\eta \varepsilon_v (\Psi_k - \Psi_{k+1}) \Phi_j = \mathcal{O}\left(\exp\left(-\frac{\tau L_1}{2}\right) \|\varepsilon\|_X^2\right),$$

while for $j = k$,

$$\begin{aligned} \int_{\mathbb{R}} \varepsilon_\eta \varepsilon_v (\Psi_k - \Psi_{k+1}) \Phi_k &= \int_{\mathbb{R}} (\varepsilon_\eta \Phi_k^{\frac{1}{2}}) (\varepsilon_v \Phi_k^{\frac{1}{2}}) + \int_{\mathbb{R}} \varepsilon_\eta \varepsilon_v \Phi_k (\Psi_k - \Psi_{k+1} - 1) \\ &= P''(Q_{c_k})(\varepsilon_k, \varepsilon_k) + \mathcal{O}\left(\exp\left(-\frac{\tau L_1}{2}\right) \|\varepsilon\|_X^2\right). \end{aligned}$$

Also, when $j \neq k$ and $j \neq k - 1$, we obtain

$$\int_{\mathbb{R}} \varepsilon_{\eta} \varepsilon_v (\Psi_k - \Psi_{k+1}) \Phi_{j,j+1} = \mathcal{O} \left(\exp \left(-\frac{\tau L_1}{2} \right) \|\varepsilon\|_X^2 \right).$$

The remaining two terms (corresponding to $j = k$ and $j = k - 1$) are simply grouped together as

$$\int_{\mathbb{R}} \varepsilon_{\eta} \varepsilon_v (\Psi_k - \Psi_{k+1}) (\Phi_{k-1,k} + \Phi_{k,k+1}) = P_k''(0)(\varepsilon_{k-1,k}, \varepsilon_{k-1,k}) + P_k''(0)(\varepsilon_{k,k+1}, \varepsilon_{k,k+1}).$$

The proof of estimate (34), and therefore of Proposition 3, is completed. \square

As a consequence of Proposition 3, we derive

Proof of Corollary 2. We write

$$E(Q_{c_k}) = E(Q_{c_k^*}) + E'(Q_{c_k^*})(Q_{c_k} - Q_{c_k^*}) + \mathcal{O}(|c_k - c_k^*|^2),$$

and

$$P(Q_{c_k}) = P(Q_{c_k^*}) + P'(Q_{c_k^*})(Q_{c_k} - Q_{c_k^*}) + \mathcal{O}(|c_k - c_k^*|^2).$$

In view of Proposition 3, the conclusion follows from the identity

$$E'(Q_{c_k^*}) + c_k^* P'(Q_{c_k^*}) = 0,$$

and from the inequalities (in order to recover the operators H_{c_k} and H_0^k)

$$|(c_k^* - c_k) P''(Q_{c_k})(\varepsilon_k, \varepsilon_k)| \leq |c_k^* - c_k| \int_{\mathbb{R}} |\varepsilon_{\eta}| |\varepsilon_v| \leq \mathcal{O}(|c_k - c_k^*|^2) + \mathcal{O}(\|\varepsilon\|_X^4),$$

and

$$|(c_k^* - c_k) P_k''(0)(\varepsilon_{j,j+1}, \varepsilon_{j,j+1})| \leq |c_k^* - c_k| \int_{\mathbb{R}} |\varepsilon_{\eta}| |\varepsilon_v| \leq \mathcal{O}(|c_k - c_k^*|^2) + \mathcal{O}(\|\varepsilon\|_X^4).$$

\square

We next prove Lemma 1.

Proof of Lemma 1. Let $1 \leq k \leq N$. Under the assumptions of Corollary 2, the pair ε satisfies the orthogonality conditions (24). In particular, we can write

$$\langle \varepsilon_k, \partial_x Q_{c_k} \rangle_{L^2} = \int_{\mathbb{R}} \langle \varepsilon, \partial_x Q_{c_k, a_k} \rangle_{\mathbb{R}^2} \Phi_k = \int_{\mathbb{R}} \langle \varepsilon, \partial_x Q_{c_k, a_k} \rangle_{\mathbb{R}^2} (\Phi_k - 1),$$

and

$$P'(Q_{c_k})(\varepsilon_k) = \frac{1}{2} \int_{\mathbb{R}} \langle \varepsilon, (v_{c_k, a_k}, \eta_{c_k, a_k}) \rangle_{\mathbb{R}^2} \Phi_k = \frac{1}{2} \int_{\mathbb{R}} \langle \varepsilon, (v_{c_k, a_k}, \eta_{c_k, a_k}) \rangle_{\mathbb{R}^2} (\Phi_k - 1).$$

Arguing as in the proof of Proposition 3, we infer that

$$\int_{\mathbb{R}} |\varepsilon| (|\partial_x Q_{c_k, a_k}| + |(v_{c_k, a_k}, \eta_{c_k, a_k})|) |\Phi_k - 1| = \mathcal{O} \left(L_1^{\frac{1}{2}} \exp \left(-\frac{\tau L_1}{2} \right) \|\varepsilon\|_{L^2} \right),$$

so that

$$\langle \varepsilon_k, \partial_x Q_{c_k} \rangle_{L^2} = \mathcal{O} \left(L_1^{\frac{1}{2}} \exp \left(-\frac{\tau L_1}{2} \right) \|\varepsilon\|_{L^2} \right),$$

and

$$P'(Q_{c_k})(\varepsilon_k) = \mathcal{O}\left(L_1^{\frac{1}{2}} \exp\left(-\frac{\tau L_1}{2}\right) \|\varepsilon\|_{L^2}\right). \quad (2.34)$$

We next write the following orthogonal decomposition of ε_k ,

$$\varepsilon_k := \frac{\langle \varepsilon_k, \partial_x Q_{c_k} \rangle_{L^2}}{\|\partial_x Q_{c_k}\|_{L^2}} u_k + \frac{\langle \varepsilon_k, (v_{c_k}, \eta_{c_k}) \rangle_{L^2}}{\|(v_{c_k}, \eta_{c_k})\|_{L^2}} v_k + r_k, \quad (2.35)$$

in which $u_k := \partial_x Q_{c_k} / \|\partial_x Q_{c_k}\|_{L^2}$, $v_k := (v_{c_k}, \eta_{c_k}) / \|(v_{c_k}, \eta_{c_k})\|_{L^2}$, and r_k satisfies the orthogonality conditions (22). In view of Proposition 1, we have

$$H_{c_k}(r_k) \geq \Lambda_{c_k} \|r_k\|_X^2.$$

On the other hand, since $\partial_x Q_{c_k}$ and hence u_k belongs to the kernel of the quadratic form H_{c_k} , we can write

$$H_{c_k}(\varepsilon_k) = H_{c_k}(r_k) + 2 \frac{P'(Q_{c_k})(\varepsilon_k)}{\|(v_{c_k}, \eta_{c_k})\|_{L^2}} \langle \mathcal{H}_{c_k}(v_k), \varepsilon_k \rangle_{L^2} + \frac{(P'(Q_{c_k})(\varepsilon_k))^2}{\|(v_{c_k}, \eta_{c_k})\|_{L^2}^2} H_{c_k}(v_k).$$

It follows from (2.34) that

$$H_{c_k}(\varepsilon_k) \geq \Lambda_{c_k} \|r_k\|_X^2 + \mathcal{O}\left(L_1^{\frac{1}{2}} \exp\left(-\frac{\tau L_1}{2}\right) \|\varepsilon\|_{L^2}^2\right). \quad (2.36)$$

Next, we similarly deduce from (2.35) that

$$\|r_k\|_X - \|\varepsilon_k\|_X \leq \left| \frac{\langle \varepsilon_k, \partial_x Q_{c_k} \rangle_{L^2}}{\|\partial_x Q_{c_k}\|_{L^2}} \right| \|u_k\|_X + \left| \frac{P'(Q_{c_k})(\varepsilon_k)}{\|(v_{c_k}, \eta_{c_k})\|_{L^2}} \right| \|v_k\|_X = \mathcal{O}\left(L_1^{\frac{1}{2}} \exp\left(-\frac{\tau L_1}{2}\right) \|\varepsilon\|_{L^2}\right),$$

and therefore from (2.36),

$$H_{c_k}(\varepsilon_k) \geq \frac{\Lambda_{c_k}}{2} \|\varepsilon_k\|_X^2 + \mathcal{O}\left(L_1^{\frac{1}{2}} \exp\left(-\frac{\tau L_1}{2}\right) \|\varepsilon\|_X^2\right). \quad (2.37)$$

We now turn to the terms $H_0^k(\varepsilon_{k,k+1})$. For $1 \leq k \leq N-1$, we write

$$\begin{aligned} & [c_k P_k''(0) + c_{k+1} P_{k+1}''(0)](\varepsilon_{k,k+1}, \varepsilon_{k,k+1}) \\ &= \int_{\mathbb{R}} c_k (\Psi_k - \Psi_{k+1}) + c_{k+1} (\Psi_{k+1} - \Psi_{k+2})(\varepsilon_{k,k+1})_{\eta} (\varepsilon_{k,k+1})_v. \end{aligned}$$

Since

$$|c_k (\Psi_k - \Psi_{k+1}) + c_{k+1} (\Psi_{k+1} - \Psi_{k+2})| \leq \max\{|c_k|, |c_{k+1}|\},$$

we obtain

$$H_0^k(\varepsilon_{k,k+1}) \geq [E''(0) - \max\{|c_k|, |c_{k+1}|\} P''(0)](\varepsilon_{k,k+1}, \varepsilon_{k,k+1}).$$

On the other hand, the quadratic form $E''(0) - \sqrt{2} P''(0)$ is non-negative, so that

$$H_0^k(\varepsilon_{k,k+1}) \geq \left(1 - \frac{\max\{|c_k|, |c_{k+1}|\}}{\sqrt{2}}\right) E''(0)(\varepsilon_{k,k+1}, \varepsilon_{k,k+1}) \geq \frac{1}{4} \left(1 - \frac{\max\{|c_k|, |c_{k+1}|\}}{\sqrt{2}}\right) \|\varepsilon_{k,k+1}\|_X^2. \quad (2.38)$$

Similarly,

$$H_0^0(\varepsilon_{0,1}) \geq \left(1 - \frac{|c_1|}{\sqrt{2}}\right) E''(0)(\varepsilon_{0,1}, \varepsilon_{0,1}) \geq \frac{1}{4} \left(1 - \frac{|c_1|}{\sqrt{2}}\right) \|\varepsilon_{0,1}\|_X^2, \quad (2.39)$$

and

$$H_0^N(\varepsilon_{N,N+1}) \geq \left(1 - \frac{|c_N|}{\sqrt{2}}\right) E''(0)(\varepsilon_{N,N+1}, \varepsilon_{N,N+1}) \geq \frac{1}{4} \left(1 - \frac{|c_N|}{\sqrt{2}}\right) \|\varepsilon_{N,N+1}\|_X^2. \quad (2.40)$$

In view of (2.37), (2.38), (2.39) and (2.40), we set

$$\Lambda^* = \min \left\{ \frac{\Lambda_{c_k}}{2}, \frac{1}{4} \left(1 - \frac{|c_k|}{\sqrt{2}}\right), 1 \leq k \leq N \right\},$$

so that (35) and (36) follow. Notice that Corollary 1 and the last statement of Proposition 1 ensure that Λ^* only depends on \mathfrak{c}^* .

In order to complete the proof of Lemma 1, we now derive (37). In view of definitions (31) and (32), we may write

$$\sum_{k=1}^N \|\varepsilon_k\|_X^2 = \sum_{k=1}^N \int_{\mathbb{R}} \left((\partial_x(\varepsilon_\eta \Phi_k^{\frac{1}{2}}))^2 + \Phi_k(\varepsilon_\eta^2 + \varepsilon_v^2) \right),$$

and

$$\sum_{k=0}^N \|\varepsilon_{k,k+1}\|_X^2 = \sum_{k=0}^N \int_{\mathbb{R}} \left((\partial_x(\varepsilon_\eta \Phi_{k,k+1}^{\frac{1}{2}}))^2 + \Phi_{k,k+1}(\varepsilon_\eta^2 + \varepsilon_v^2) \right).$$

Expanding of the derivatives $\partial_x(\varepsilon_\eta \Phi_k^{\frac{1}{2}})$ and $\partial_x(\varepsilon_\eta \Phi_{k,k+1}^{\frac{1}{2}})$ and summing, we are led to

$$\begin{aligned} \sum_{k=1}^N \|\varepsilon_k\|_X^2 + \sum_{k=0}^N \|\varepsilon_{k,k+1}\|_X^2 &= \int_{\mathbb{R}} \left(\sum_{k=1}^N \Phi_k + \sum_{k=0}^N \Phi_{k,k+1} \right) ((\partial_x \varepsilon_\eta)^2 + \varepsilon_\eta^2 + \varepsilon_v^2) \\ &+ \int_{\mathbb{R}} \partial_x \left(\sum_{k=1}^N \Phi_k + \sum_{k=0}^N \Phi_{k,k+1} \right) \varepsilon_\eta \partial_x \varepsilon_\eta + \sum_{k=1}^N \int_{\mathbb{R}} \frac{(\partial_x \Phi_k)^2}{4\Phi_k} \varepsilon_\eta^2 + \sum_{k=0}^N \int_{\mathbb{R}} \frac{(\partial_x \Phi_{k,k+1})^2}{4\Phi_{k,k+1}} \varepsilon_\eta^2. \end{aligned}$$

Since

$$\sum_{k=1}^N \Phi_k + \sum_{k=0}^N \Phi_{k,k+1} = 1,$$

and $\Phi_k > 0$ on \mathbb{R} (so that the last two sums in the previous identity are non-negative), we finally deduce that

$$\sum_{k=1}^N \|\varepsilon_k\|_X^2 + \sum_{k=0}^N \|\varepsilon_{k,k+1}\|_X^2 \geq \int_{\mathbb{R}} ((\partial_x \varepsilon_\eta)^2 + \varepsilon_\eta^2 + \varepsilon_v^2) = \|\varepsilon\|_X^2.$$

This completes the proof of Lemma 1. \square

Proof of Corollary 3. Inequality (39) follows from Corollary 2, Lemma 1 and inequality (38), while inequality (40) follows from Corollary 2 and the fact that H_{c_k} and H_0 are continuous quadratic forms, with continuous bounds depending only on \mathfrak{c}^* , in view of Corollary 1. \square

3 Dynamical properties

In this section, we focus on the evolution in time under (HGP) of various quantities introduced so far. In particular, we present the proofs of Propositions 4 and 5.

3.1 Dynamics of the modulation parameters

Proof of Proposition 4. We first derive the conclusions of Proposition 4 when the pair (η, v) is a smooth solution to (HGP), for which we are allowed to compute the differential equations satisfied by the modulation parameters $a_k(t)$ and $c_k(t)$. We next complete the proof using a density argument.

Step 1. *Regularity and estimates for smooth solutions.*

In this step, we make the further assumption that $(\eta, v) \in \mathcal{C}^0([0, T], X^3(\mathbb{R}))$, from which it follows by equation (HGP) that $(\eta, v) \in \mathcal{C}^1([0, T], X(\mathbb{R}))$, and then from Proposition 2 that $\mathbf{c} \in \mathcal{C}^1([0, T], \text{Adm}_N)$ and $\mathbf{a} \in \mathcal{C}^1([0, T], \mathbb{R}^N)$. We derive from (HGP) that $\varepsilon := (\varepsilon_\eta, \varepsilon_v)$ is solution to the equations

$$\begin{aligned} \partial_t \varepsilon_\eta &= 2\partial_x(\varepsilon_v - \varepsilon_\eta v_{\mathbf{c}, \mathbf{a}} - \varepsilon_v \eta_{\mathbf{c}, \mathbf{a}} - \varepsilon_\eta \varepsilon_v) - \sum_{k=1}^N \sum_{j \neq k} \partial_x(\eta_{c_k, a_k} v_{c_j, a_j}) \\ &\quad - \sum_{k=1}^N c'_k \partial_c \eta_{c_k, a_k} + \sum_{k=1}^N (a'_k - c_k) \partial_x \eta_{c_k, a_k}, \end{aligned} \quad (3.1)$$

and

$$\begin{aligned} \partial_t \varepsilon_v &= \partial_x(\varepsilon_\eta - 2v_{\mathbf{c}, \mathbf{a}} \varepsilon_v - (\varepsilon_v)^2) - \sum_{k=1}^N \sum_{j \neq k} \partial_x(v_{c_k, a_k} v_{c_j, a_j}) + \sum_{k=1}^N \partial_x^2 \left(\frac{\partial_x \eta_{c_k, a_k}}{2(1 - \eta_{c_k, a_k})} \right) \\ &\quad - \partial_x^2 \left(\frac{\partial_x \eta_{\mathbf{c}, \mathbf{a}} + \partial_x \varepsilon_\eta}{2(1 - \eta_{\mathbf{c}, \mathbf{a}} - \varepsilon_\eta)} \right) + \sum_{k=1}^N \partial_x \left(\frac{3(\partial_x \eta_{c_k, a_k})^2}{4(1 - \eta_{c_k, a_k})^2} \right) - \partial_x \left(\frac{3(\partial_x \eta_{\mathbf{c}, \mathbf{a}} + \partial_x \varepsilon_\eta)^2}{4(1 - \eta_{\mathbf{c}, \mathbf{a}} - \varepsilon_\eta)^2} \right) \\ &\quad - \sum_{k=1}^N c'_k \partial_c v_{c_k, a_k} + \sum_{k=1}^N (a'_k - c_k) \partial_x v_{c_k, a_k}. \end{aligned} \quad (3.2)$$

Here and in the sequel, we have dropped the explicit mention to the dependence in t when this does not lead to a confusion.

We next differentiate the orthogonality conditions (22) to obtain

$$\frac{d}{dt} \langle \varepsilon(\cdot, t), \partial_x Q_{c_k(t), a_k(t)} \rangle_{L^2} = \frac{d}{dt} P'(Q_{c_k(t), a_k(t)})(\varepsilon(\cdot, t)) = 0,$$

for any $1 \leq k \leq N$. Expanding the previous identities according to the chain rule and combining the resulting equations with (3.1) and (3.2) yields the system

$$M(t) \begin{pmatrix} \mathbf{c}'(t) \\ \mathbf{a}'(t) - \mathbf{c}(t) \end{pmatrix} = \Phi(t), \quad (3.3)$$

where the coefficients $(M_{k,j}(t))_{1 \leq j, k \leq 2N}$ of the $2N$ -dimensional matrix $M(t)$ are given by

$$\begin{aligned} M_{k,j}(t) &:= P'(Q_{c_k, a_k})(\partial_c Q_{c_j, a_j}) - \delta_{j,k} P'(\partial_c Q_{c_k, a_k})(\varepsilon), \\ M_{k,j+N}(t) &:= -P'(Q_{c_k, a_k})(\partial_x Q_{c_j, a_j}) + \delta_{j,k} P'(\partial_x Q_{c_k, a_k})(\varepsilon), \\ M_{k+N,j}(t) &:= \langle \partial_c Q_{c_j, a_j}, \partial_x Q_{c_k, a_k} \rangle_{L^2} - \delta_{j,k} \langle \varepsilon, \partial_c \partial_x Q_{c_k, a_k} \rangle_{L^2}, \\ M_{k+N,j+N}(t) &:= -\langle \partial_x Q_{c_j, a_j}, \partial_x Q_{c_k, a_k} \rangle_{L^2} + \delta_{j,k} \langle \varepsilon, \partial_x^2 Q_{c_k, a_k} \rangle_{L^2}, \end{aligned}$$

and the components $(\Phi_k(t))_{1 \leq k \leq 2N}$ of the $2N$ -vector $\Phi(t)$ are given by

$$\begin{aligned}
\Phi_k(t) := & -2\langle \varepsilon_v - \varepsilon_\eta v_{c,a} - \varepsilon_v \eta_{c,a} - \varepsilon_\eta \varepsilon_v, \partial_x v_{c_k, a_k} \rangle_{L^2} - \langle \varepsilon_\eta - 2v_{c,a} \varepsilon_v - (\varepsilon_v)^2, \partial_x \eta_{c_k, a_k} \rangle_{L^2} \\
& - c_k P'(\partial_x Q_{c_k, a_k})(\varepsilon) - \left\langle \frac{\partial_x \eta_{c,a} + \partial_x \varepsilon_\eta}{2(1 - \eta_{c,a} - \varepsilon_\eta)} - \sum_{j=1}^N \frac{\partial_x \eta_{c_j, a_j}}{2(1 - \eta_{c_j, a_j})}, \partial_x^2 \eta_{c_k, a_k} \right\rangle_{L^2} \\
& + \left\langle \frac{3(\partial_x \eta_{c,a} + \partial_x \varepsilon_\eta)^2}{4(1 - \eta_{c,a} - \varepsilon_\eta)^2} - \sum_{j=1}^N \frac{3(\partial_x \eta_{c_j, a_j})^2}{4(1 - \eta_{c_j, a_j})^2}, \partial_x \eta_{c_k, a_k} \right\rangle_{L^2} \\
& + \sum_{j=1}^N \sum_{\substack{l=1 \\ l \neq j}}^N \left(\langle \eta_{c_j, a_j} v_{c_l, a_l}, \partial_x v_{c_k, a_k} \rangle_{L^2} + \langle v_{c_j, a_j} v_{c_l, a_l}, \partial_x \eta_{c_k, a_k} \rangle_{L^2} \right),
\end{aligned} \tag{3.4}$$

$$\begin{aligned}
\Phi_{k+N}(t) := & 2\langle \varepsilon_\eta v_{c,a} + \varepsilon_v \eta_{c,a} + \varepsilon_\eta \varepsilon_v - \varepsilon_v, \partial_x^2 \eta_{c_k, a_k} \rangle_{L^2} - \langle \varepsilon_\eta - 2v_{c,a} \varepsilon_v - (\varepsilon_v)^2, \partial_x^2 v_{c_k, a_k} \rangle_{L^2} \\
& - c_k \langle \varepsilon, \partial_x^2 Q_{c_k, a_k} \rangle_{L^2} - \left\langle \frac{\partial_x \eta_{c,a} + \partial_x \varepsilon_\eta}{2(1 - \eta_{c,a} - \varepsilon_\eta)} - \sum_{j=1}^N \frac{\partial_x \eta_{c_j, a_j}}{2(1 - \eta_{c_j, a_j})}, \partial_x^3 v_{c_k, a_k} \right\rangle_{L^2} \\
& + \left\langle \frac{3(\partial_x \eta_{c,a} + \partial_x \varepsilon_\eta)^2}{4(1 - \eta_{c,a} - \varepsilon_\eta)^2} - \sum_{j=1}^N \frac{3(\partial_x \eta_{c_j, a_j})^2}{4(1 - \eta_{c_j, a_j})^2}, \partial_x^2 v_{c_k, a_k} \right\rangle_{L^2} \\
& + \sum_{j=1}^N \sum_{\substack{l=1 \\ l \neq j}}^N \left(\langle \eta_{c_j, a_j} v_{c_l, a_l}, \partial_x^2 \eta_{c_k, a_k} \rangle_{L^2} + \langle v_{c_j, a_j} v_{c_l, a_l}, \partial_x^2 v_{c_k, a_k} \rangle_{L^2} \right).
\end{aligned} \tag{3.5}$$

We next write the matrix $M(t)$ under the form $M(t) := D(t) + H(t)$, where $D(t)$ refers to the diagonal matrix with diagonal coefficients

$$\begin{aligned}
D_{k,k}(t) & := P'(Q_{c_k})(\partial_c Q_{c_k}) = (2 - c_k^2)^{\frac{1}{2}} \neq 0, \\
D_{k+N, k+N}(t) & := \|\partial_x Q_{c_k}\|_{L^2}^2 = \frac{(2 - c_k^2)^{\frac{3}{2}}}{3} \neq 0,
\end{aligned}$$

where $1 \leq k \leq N$. In particular, $D(t)$ is invertible, and its inverse is uniformly bounded with respect to t in view of (28). In order to estimate the matrix $H(t)$, we first invoke the identities

$$P'(Q_{c_k, a_k})(\partial_x Q_{c_k, a_k}) = \langle \partial_c Q_{c_k, a_k}, \partial_x Q_{c_k, a_k} \rangle_{L^2} = 0.$$

Since the remaining terms involve either ε or distinct solitons, we can rely on computations similar to the ones in the proof of Proposition 3 to get

$$|H_{k,j}(t)| = \mathcal{O}(\|\varepsilon\|_{L^2}) + \mathcal{O}\left(L \exp\left(-\frac{\nu_{c^*} L}{2}\right)\right).$$

Hence, there exists $\alpha_2 \leq \alpha_1$ and $L_2 \geq L_1$ such that for $\alpha \leq \alpha_2$ and $L \geq L_2$, the operator norm of the matrix $D(t)^{-1}H(t)$ is less than $1/2$, so that $M(t)$ is invertible. Moreover, the operator norms of the matrices $M(t)^{-1}$ are uniformly bounded with respect to t .

At this stage, we have proved that

$$\sum_{k=1}^N (|a'_k(t) - c_k(t)| + |c'_k(t)|) = \mathcal{O}(\|\Phi(t)\|_{\mathbb{R}^{2N}}),$$

for any $t \in [0, T]$. Inspection of $\Phi(t)$ shows that, similar to $H_{k,j}$ above, it only contains terms involving either ε or distinct solitons. Therefore we obtain

$$\sum_{k=1}^N (|a'_k(t) - c_k(t)| + |c'_k(t)|) = \mathcal{O}(\|\varepsilon\|_{L^2}) + \mathcal{O}\left(L \exp\left(-\frac{\nu_{c^*} L}{2}\right)\right),$$

which corresponds to (41).

Step 2. *Density argument.*

We consider a sequence (η_0^n, v_0^n) in $\mathcal{NV}^3(\mathbb{R})$ such that

$$(\eta_0^n, v_0^n) \rightarrow (\eta_0, v_0) \quad \text{in } X(\mathbb{R}),$$

as $n \rightarrow +\infty$, and we denote (η^n, v^n) the corresponding solutions to (HGP). It follows from Theorem 1 that the solutions (η^n, v^n) are well-defined on $[0, T]$ for n sufficiently large and that

$$(\eta^n, v^n) \rightarrow (\eta, v) \quad \text{in } \mathcal{C}^0([0, T], X(\mathbb{R})), \quad (3.6)$$

as $n \rightarrow +\infty$. In particular, for n large enough, the compactness of the segment $[0, T]$ ensures that $(\eta^n(\cdot, t), v^n(\cdot, t))$ belongs to $\mathcal{U}_{c^*}(\alpha, L)$ for any $t \in [0, T]$. As a consequence, one can apply Step 1 to (η^n, v^n) . This yields maps \mathbf{c}^n and \mathbf{a}^n of class \mathcal{C}^1 on $[0, T]$ such that the orthogonality conditions (22) are satisfied by the pair $\varepsilon^n := (\varepsilon_\eta^n, \varepsilon_v^n) = (\eta^n, v^n) - R_{\mathbf{c}^n, \mathbf{a}^n}$. Moreover, combining the Lipschitz continuity of the maps \mathfrak{C} and \mathfrak{A} on some tubular neighborhood of the compact set $\{(\eta(t), v(t)), t \in [0, T]\}$ with (3.6), we have

$$(\mathbf{c}^n(t), \mathbf{a}^n(t)) \rightarrow (\mathbf{c}(t), \mathbf{a}(t)) \quad \text{in } \mathcal{C}^0([0, T], \mathbb{R}^{2N}), \quad (3.7)$$

as $n \rightarrow +\infty$. Recall that, by formulae (5), the maps $(c, a) \rightarrow \partial_c^\alpha \partial_x^\beta Q_c(\cdot - a)$ are continuous from $((-\sqrt{2}, \sqrt{2}) \setminus \{0\}) \times \mathbb{R}$ into $L^2(\mathbb{R})^2$ for any $(\alpha, \beta) \in \mathbb{N}^2$. Combined with (3.7), this shows that the matrices $M^n(t)$ converge towards the matrices $M(t)$, uniformly with respect to t , as $n \rightarrow +\infty$. Since the map $A \mapsto A^{-1}$ is continuous from $\mathcal{GL}_N(\mathbb{R})$ to $\mathcal{M}_N(\mathbb{R})$, the inverses $(M^n(t))^{-1}$ also converge towards the inverses $M(t)^{-1}$ uniformly with respect to $t \in [0, T]$.

We next turn to the vectors $\Phi^n(t)$. In view of (3.6) and (3.7), we can write

$$\varepsilon^n = (\eta^n, v^n) - R_{\mathbf{c}^n, \mathbf{a}^n} \rightarrow (\eta, v) - R_{\mathbf{c}, \mathbf{a}} = \varepsilon \quad \text{in } \mathcal{C}^0([0, T], X(\mathbb{R})), \quad (3.8)$$

as $n \rightarrow +\infty$. On the other hand, in view of formulae (5), the derivatives $\partial_c^\alpha \partial_x^\beta Q_c$ have a vanishing limit at infinity for any $(\alpha, \beta) \in \mathbb{N}^2$. Therefore, the maps $(c, a) \rightarrow \partial_c^\alpha \partial_x^\beta Q_c(\cdot - a)$ are not only continuous for the L^2 -norm, but also for the uniform one. Applying (3.7) and (3.8) to expressions (3.4) and (3.5), and invoking the Sobolev embedding theorem of $H^1(\mathbb{R})$ into $\mathcal{C}^0(\mathbb{R})$ when necessary, it follows that

$$\Phi^n \rightarrow \Phi \quad \text{in } \mathcal{C}^0([0, T], \mathbb{R}^{2N}),$$

as $n \rightarrow +\infty$. Coming back to (3.3), we are finally led to

$$\begin{pmatrix} (\mathbf{c}^n)'(t) \\ (\mathbf{a}^n)'(t) - \mathbf{c}^n(t) \end{pmatrix} \rightarrow M(t)^{-1} \Phi(t),$$

as $n \rightarrow +\infty$, the convergence being uniform with respect to $t \in [0, T]$. In view of (3.7), this shows that the maps \mathbf{c} and \mathbf{a} are actually of class \mathcal{C}^1 on $[0, T]$. Moreover, they are given by

$$\begin{pmatrix} \mathbf{c}'(t) \\ \mathbf{a}'(t) - \mathbf{c}(t) \end{pmatrix} = M(t)^{-1} \Phi(t),$$

for any $t \in [0, T]$, so that (41) follows as in Step 1. \square

We are now in position to provide the

Proof of Corollary 4. When $(\eta(\cdot, t), v(\cdot, t)) \in \mathcal{U}_{c^*}(\alpha, L)$ for any $t \in [0, T]$, we can combine estimates (25) and (41) to obtain

$$\sum_{k=1}^N |a'_k(t) - c_k^*| \leq \sum_{k=1}^N |a'_k(t) - c_k(t)| + \sum_{k=1}^N |c_k(t) - c_k^*| = \mathcal{O}(\alpha) + \mathcal{O}\left(\left(L \exp\left(-\frac{\nu_{c^*} L}{2}\right)\right)\right). \quad (3.9)$$

As a consequence, we can fix α_3 small enough and L_3 sufficiently large, so that

$$a'_{k+1}(t) - a'_k(t) > \frac{1}{2}(c_{k+1}^* - c_k^*) \geq \sigma^*,$$

for any $1 \leq k \leq N-1$, when $\alpha \leq \alpha_3$ and $L \geq L_3$. The first inequality in (43) follows by writing

$$a_{k+1}(t) - a_k(t) - a_{k+1}(0) + a_k(0) = \int_0^t (a'_{k+1}(s) - a'_k(s)) ds > \sigma^* t.$$

Applying (27) to the pair $(\eta(\cdot, 0), v(\cdot, 0))$, we deduce the second inequality in (43).

Concerning (44), it follows from (3.9) that

$$\left| \sqrt{2 - (a'_k(t))^2} - \sqrt{2 - (c_k^*)^2} \right| = \mathcal{O}(\alpha) + \mathcal{O}\left(\left(L \exp\left(-\frac{\nu_{c^*} L}{2}\right)\right)\right).$$

Therefore, one can decrease α_3 and increase L_3 , if necessary, so that inequality (44) holds for any $1 \leq k \leq N$, when $\alpha \leq \alpha_3$ and $L \geq L_3$. This completes the proof of Corollary 4. \square

3.2 Monotonicity of localized scalar momentum

This subsection is devoted to the almost monotonicity properties of suitably localized versions of the scalar momentum

$$P(\eta, v) := \frac{1}{2} \int_{\mathbb{R}} \eta v.$$

We derive them using the conservative form of the equation governing the integrand ηv of P .

Lemma 3.1. *Let $(\eta, v) \in \mathcal{C}^0([0, T], \mathcal{NV}^3(\mathbb{R}))$ be a solution to (HGP). Then,*

$$\partial_t(\eta v) = \partial_x \left((1 - 2\eta)v^2 + \frac{\eta^2}{2} + \frac{(3 - 2\eta)(\partial_x \eta)^2}{4(1 - \eta)^2} \right) + \frac{1}{2} \partial_x^3 (\eta + \ln(1 - \eta)) \quad \text{on } [0, T] \times \mathbb{R}. \quad (3.10)$$

Proof. Under such regularity, it follows from (HGP) that $\partial_t \eta$ and $\partial_t v$ are continuous functions on $[0, T] \times \mathbb{R}$ and that

$$\partial_t(\eta v) = \partial_x \left((1 - 2\eta)v^2 + \frac{\eta^2}{2} \right) - \eta \partial_x \left(\frac{\partial_x^2 \eta}{2(1 - \eta)} + \frac{(\partial_x \eta)^2}{4(1 - \eta)^2} \right).$$

The conservation law given by (3.10) then follows from the identity

$$\eta \partial_x \left(\frac{\partial_x^2 \eta}{2(1 - \eta)} + \frac{(\partial_x \eta)^2}{4(1 - \eta)^2} \right) = -\frac{1}{2} \partial_x^3 (\eta + \ln(1 - \eta)) - \partial_x \left(\frac{(3 - 2\eta)(\partial_x \eta)^2}{4(1 - \eta)^2} \right).$$

\square

One can drop the smoothness assumptions in Lemma 3.1 by deriving an integral version of (3.10).

Corollary 3.1. *Let $(\eta, v) \in \mathcal{C}^0([0, T], \mathcal{NV}(\mathbb{R}))$ be a solution to (HGP), and let $\psi \in \mathcal{C}^0([0, T], \mathcal{C}_b^3(\mathbb{R})) \cap \mathcal{C}^1([0, T], \mathcal{C}_b^0(\mathbb{R}))$. Then, on $[0, T]$*

$$\partial_t \left(\int_{\mathbb{R}} \psi \eta v \right) = \int_{\mathbb{R}} \partial_t \psi \eta v - \int_{\mathbb{R}} \partial_x \psi \left((1-2\eta)v^2 + \frac{\eta^2}{2} + \frac{(3-2\eta)(\partial_x \eta)^2}{4(1-\eta)^2} \right) - \frac{1}{2} \int_{\mathbb{R}} \partial_x^3 \psi (\eta + \ln(1-\eta)). \quad (3.11)$$

Proof. It follows by approximation of the initial datum using Theorem 1. \square

In particular, taking $\psi = 1$ in Corollary 3.1 yields the conservation of P .

We turn now to localized versions of P . We deduce from Corollary 3.1

Proposition 3.1. *Let $(\eta, v) \in \mathcal{C}^0([0, T], \mathcal{NV}(\mathbb{R}))$ be a solution to (HGP). Let $\Psi \in \mathcal{C}_b^3(\mathbb{R})$ be a non-decreasing function such that $|\partial_x^3 \Psi| \leq C_0 \partial_x \Psi$ on \mathbb{R} for some positive number C_0 . Let $X \in \mathcal{C}^1([0, T], \mathbb{R})$, $t_0 \in [0, T]$ and $\tau_0 > 0$ be such that $2C_0\tau_0^2 < 1$, and*

$$\eta(t_0, x) \leq \frac{1}{2} - \frac{X'(t_0)^2}{4(1-2C_0\tau_0^2)}, \quad (3.12)$$

for any $x \in [X(t_0) - R_0, X(t_0) + R_0]$ and some positive number R_0 . Then,

$$\partial_t \left(\int_{\mathbb{R}} \Psi(\tau_0(x - X(t))) \eta(x, t) v(x, t) dx \right) \Big|_{t=t_0} \leq \tau_0 \left(6 + \frac{6}{1-\eta_{\text{sup}}} \right) E(\eta, v) \sup_{|x| \geq \tau_0 R_0} \partial_x \Psi(x), \quad (3.13)$$

where

$$\eta_{\text{sup}} := \sup \{ \eta(t, x), (t, x) \in [0, T] \times \mathbb{R} \} < 1.$$

Proof. Set $I_0 := [X(t_0) - R_0, X(t_0) + R_0]$. It follows from (3.12) and the condition $2C_0\tau_0^2 < 1$ that $\eta \leq 1/2$ on I_0 , and real analysis then implies that

$$|\eta + \ln(1-\eta)| \leq \eta^2 \quad \text{on } I_0. \quad (3.14)$$

Applying (3.11), we may rewrite the left-hand side of (3.13) as

$$- \int_{\mathbb{R}} \tau_0 \partial_x \Psi \left((1-2\eta)v^2 + \frac{\eta^2}{2} + X'(t_0)\eta v + \frac{(3-2\eta)(\partial_x \eta)^2}{4(1-\eta)^2} \right) - \frac{1}{2} \int_{\mathbb{R}} \tau_0^3 \partial_x^3 \Psi (\eta + \ln(1-\eta)),$$

where $\partial_x \Psi$ and $\partial_x^3 \Psi$ are evaluated at the point $\tau_0(x - X(t_0))$.

For $x \in I_0$, we deduce from (3.14) and the bound $|\partial_x^3 \Psi| \leq C_0 \partial_x \Psi$ that

$$\begin{aligned} \tau_0 \partial_x \Psi \left((1-2\eta)v^2 + \frac{\eta^2}{2} + X'(t_0)\eta v \right) + \frac{1}{2} \tau_0^3 \partial_x^3 \Psi (\eta + \ln(1-\eta)) \\ \geq \tau_0 \partial_x \Psi \left((1-2\eta)v^2 + \left(\frac{1}{2} - C_0\tau_0^2 \right) \eta^2 - |X'(t_0)|\eta v \right) \geq 0, \end{aligned}$$

where the non-negativity of the quadratic form, which yields the last inequality, follows from assumption (3.12). Since $\eta \leq 1/2$ on I_0 , we also have

$$\tau_0 \partial_x \Psi \frac{(3-2\eta)(\partial_x \eta)^2}{4(1-\eta)^2} \geq 0,$$

for $x \in I_0$, and therefore,

$$\tau_0 \partial_x \Psi \left((1-2\eta)v^2 + \frac{\eta^2}{2} + X'(t_0)\eta v + \frac{(3-2\eta)(\partial_x \eta)^2}{4(1-\eta)^2} \right) + \frac{\tau_0^3}{2} \partial_x^3 \Psi (\eta + \ln(1-\eta)) \leq 0 \quad (3.15)$$

on I_0 .

It remains to consider the case $x \notin I_0$. In view of the definition of η_{sup} , we have

$$\begin{aligned} -\frac{2}{1-\eta_{\text{sup}}}\frac{(1-\eta)v^2}{2} &\leq (1-2\eta)v^2 \leq 4\frac{(1-\eta)v^2}{2}, \\ 0 &\leq \frac{(3-2\eta)(\partial_x\eta)^2}{4(1-\eta)^2} \leq \max\left\{6, \frac{6}{1-\eta_{\text{sup}}}\right\} \frac{(\partial_x\eta)^2}{8(1-\eta)}, \end{aligned}$$

as well as

$$|\eta + \ln(1-\eta)| \leq \frac{2}{1-\eta_{\text{sup}}}\frac{\eta^2}{4}, \quad \text{and} \quad |X'(t_0)\eta v| \leq 2\sqrt{2}\frac{\eta^2}{4} + \frac{\sqrt{2}}{1-\eta_{\text{sup}}}\frac{(1-\eta)v^2}{2},$$

where we have used the inequality $|X'(t_0)| \leq \sqrt{2}$. Therefore,

$$\begin{aligned} \left| \tau_0 \partial_x \Psi \left((1-2\eta)v^2 + \frac{\eta^2}{2} + X'(t_0)\eta v \right) + \frac{1}{2} \tau_0^3 \partial_x^3 \Psi (\eta + \ln(1-\eta)) \right| \\ \leq \tau_0 \left(6 + \frac{6}{1-\eta_{\text{sup}}} \right) e(\eta, v) \sup_{|x| \geq \tau_0 R_0} \partial_x \Psi(x) \end{aligned} \quad (3.16)$$

on $\mathbb{R} \setminus I_0$, and the conclusion follows from (3.15) and (3.16) by integration. \square

We are finally in position to provide the

Proof of Proposition 5. Let $t_0 \in [0, T]$. We apply Proposition 3.1 with the choices $t_0 := t_0$, $\Psi(x) := (1 + \text{th}(x))/2$, $\tau_0 := \nu_{\mathbf{c}^*}/16$ and $X(t) := (a_k(t) + a_{k+1}(t))/2$. A quick computation first shows that for this choice of Ψ , the inequality $|\partial_x^3 \Psi| \leq C_0 \partial_x \Psi$ holds for $C_0 := 4$. The condition $2C_0\tau_0^2 < 1$ is satisfied since $\nu_{\mathbf{c}^*}/16 < \sqrt{2}/16$. We turn now to condition (3.12). It follows from (44), namely

$$\sqrt{2 - (a'_k(t_0))^2} \geq \frac{\nu_{\mathbf{c}^*}}{2} \quad \text{and} \quad \sqrt{2 - (a'_{k+1}(t_0))^2} \geq \frac{\nu_{\mathbf{c}^*}}{2},$$

that

$$X'(t_0)^2 \leq 2 - \frac{\nu_{\mathbf{c}^*}^2}{4}. \quad (3.17)$$

Using the inequality $1/(1-s) \leq 1+2s$ for $s \in [0, 1/2]$, we infer from (3.17) and the inequality $\nu_{\mathbf{c}^*}/16 < 1/4$ that

$$\frac{1}{2} - \frac{X'(t_0)^2}{4(1-8(\nu^*)^2)} \geq \frac{1}{2} - \frac{2 - \frac{\nu_{\mathbf{c}^*}^2}{4}}{4(1-8(\nu^*)^2)} \geq \frac{1}{2} - \left(\frac{1}{2} - \frac{\nu_{\mathbf{c}^*}^2}{16} \right) (1 + 16(\nu^*)^2) \geq \frac{\nu_{\mathbf{c}^*}^2}{32}.$$

At this stage, we come back to the decomposition

$$(\eta(\cdot, t), v(\cdot, t)) = R_{\mathbf{c}(t), \mathbf{a}(t)}(\cdot) + \varepsilon(\cdot, t).$$

In view of (25), we have $\|\varepsilon(\cdot, t)\|_X \leq K_1\alpha$. Hence, by the Sobolev embedding theorem, we may choose $\alpha_4 \leq \alpha_3$ so that, if $\alpha \leq \alpha_4$, then

$$\|\varepsilon(\cdot, t)\|_{L^\infty} \leq \frac{\nu_{\mathbf{c}^*}^2}{64}. \quad (3.18)$$

It follows from (3.18) that (3.12) holds everywhere on the set

$$S := \left\{ x \in \mathbb{R}, \text{ s.t. } |\eta_{\mathbf{c}(t_0), \mathbf{a}(t_0)}(x)| \leq \frac{\nu_{\mathbf{c}^*}^2}{64} \right\}.$$

Combining the definition of $\eta_{c,a}$, the explicit form (5) of the solitons, condition (43) on $\mathbf{a}(t_0)$ and the inequality $\nu_c \geq \nu_{c^*}/2$ stated in (28), it follows that the set S contains the interval $[X(t_0) - R_0, X(t_0) + R_0]$ where

$$R_0 := \frac{(L-1) + \sigma^* t_0}{2} - \frac{2}{\nu_{c^*}} \ln \left(\frac{256N}{\nu_{c^*}^2} \right).$$

Finally, since $0 < \partial_x \Psi(x) \leq 2 \exp(-2|x|)$ for any $x \in \mathbb{R}$, we obtain from (3.13) that

$$\frac{d}{dt}(Q_k(t))|_{t=t_0} \leq \mathcal{O} \left(\exp \left(-\frac{\nu_{c^*}(L + \sigma^* t_0)}{16} \right) \right),$$

and the proof is completed. \square

We finally give the

Proof of Corollary 5. Since the energy E and the momentum P are conserved quantities for equation (HGP), we infer from definition (45) that

$$\frac{d}{dt}(G(t)) = \sum_{k=2}^N (c_k^* - c_{k+1}^*) \frac{d}{dt}(Q_k(t)).$$

Corollary 5 follows by combining Proposition 5 and the ordering condition (42). \square

A A quantitative version of the implicit function theorem

In this appendix, we give the proof of the following version of the implicit function theorem which we have invoked in order to establish Proposition 4.

Proposition A.1. *Let E and F be two Banach spaces, V , an open subset of F , and ϕ , a function of class \mathcal{C}^2 from $E \times V$ to F . We assume that*

(i) *there exist an open subset U of V and a map \bar{x} of class \mathcal{C}^1 from U to E such that*

$$\phi(\bar{x}(y), y) = 0,$$

for any $y \in U$.

(ii) *there exist a positive constant K , as well as a continuous isomorphism A_y from F to F , such that*

$$d_y \phi(\bar{x}(y), y) = A_y (Id + H_y), \tag{A.1}$$

and

$$d_x \phi(\bar{x}(y), y) = A_y T_y, \tag{A.2}$$

for any $y \in U$. In (A.1), the operator norm of the continuous linear mapping H_y is less than $1/2$, whereas in (A.2), the operator norm of T_y is less than K .

(iii) *the operator norm of $d^2 \phi_{(x,y)}$ is uniformly bounded with respect to $x \in E$ and $y \in V$.*

(iv) *there exists a positive constant r such that*

$$B(y, r) \subset V,$$

for any $y \in U$.

Then, there exist two positive numbers δ and $\Lambda := 4K + 4$, such that, given any $y \in U$, there exists a map $\pi_y \in \mathcal{C}^1(B(\bar{x}(y), \delta), V)$ such that, given any $t \in B(\bar{x}(y), \delta)$, $z = \pi_y(t)$ is the unique solution in $B(y, \Lambda\delta)$ of the equation

$$\phi(t, z) = 0.$$

Moreover, the map π_y is Lipschitz on $B(\bar{x}(y), \delta)$, with Lipschitz constant at most Λ .

Remark A.1. When y is fixed in U , assumptions (i) and (ii) are enough to apply the implicit function theorem in order to construct a map $\pi_y \in \mathcal{C}^1(B(\bar{x}(y), \delta), V)$ as in Proposition A.1. The difficulty comes from the fact that, in this case, the constants δ and K depend a priori on y .

Proof. The proof is similar to the usual proof of the implicit function theorem. We fix $y \in U$ and we set

$$\Phi(x, z) := (\Phi_1(x, z), \Phi_2(x, z)) := (x, A_y^{-1}\phi(x, z) - T_y(x)).$$

The function Φ is well-defined and of class \mathcal{C}^2 from $E \times V$ to $E \times F$. Moreover, its differential at the point $(\bar{x}(y), y)$ may be written as

$$d\Phi_{(\bar{x}(y), y)} = \begin{pmatrix} Id & 0 \\ 0 & Id + H(y) \end{pmatrix}.$$

In view of assumption (ii), the operator norm of the maps $d\Phi_{(\bar{x}(y), y)} - Id$ is less than $1/2$ for any $y \in U$. Combining with assumption (iii), we infer the existence of a positive number $\rho < r$ such that the operator norm $d\Phi_{(x, z)} - Id$ is less than $3/4$ for any $(x, z) \in B(\bar{x}(y), \rho) \times B(y, \rho)$.

Given a point $t \in E$, we next set

$$\Psi(x, z) := (t, -T_y(t)) + (x, z) - \Phi(x, z).$$

The map Ψ is well-defined on $E \times V$, and Lipschitz on $B(\bar{x}(y), \rho) \times B(y, \rho)$, with Lipschitz constant at most $3/4$. Moreover, one can check that

$$\Psi(\bar{x}(y), y) = (t, y - T_y(t - \bar{x}(y))),$$

so that, letting $\Lambda := 4 + 4K$,

$$\|\Psi(\bar{x}(y), y) - (\bar{x}(y), y)\|_{E \times F} \leq \frac{\rho}{4},$$

when $t \in B(\bar{x}(y), \rho/\Lambda)$. In this case, the map Ψ is a contraction from $B(\bar{x}(y), \rho) \times B(y, \rho)$ to $B(\bar{x}(y), \rho) \times B(y, \rho)$. By the fixed point theorem, there exists a unique point (x_t, z_t) in $B(\bar{x}(y), \rho) \times B(y, \rho)$ such that $\Psi(x_t, z_t) = (x_t, z_t)$, i.e.

$$x_t = t \quad \text{and} \quad \phi(t, z_t) = 0.$$

At this stage, we set $\delta := \rho/\Lambda < \rho < r$, and

$$\pi_y(t) := z_t \in V.$$

In view of the definition of z_t , given two points t_1 and t_2 in $B(\bar{x}(y), \delta)$, we have

$$\pi_y(t_1) - \pi_y(t_2) = -T_y(t_1 - t_2) + (\pi_y(t_1) - \Phi_2(t_1, \pi_y(t_1)) - \pi_y(t_2) - \Phi_2(t_2, \pi_y(t_2))),$$

so that

$$\|\pi_y(t_1) - \pi_y(t_2)\|_F \leq K\|t_1 - t_2\|_E + \frac{3}{4}(\|t_1 - t_2\|_E + \|\pi_y(t_1) - \pi_y(t_2)\|_F).$$

As a consequence, the map π_y is Lipschitz on $B(\bar{x}(y), \delta)$, with Lipschitz constant at most $4K + 3 \leq \Lambda$. The fact that π_y is of class \mathcal{C}^1 on $B(\bar{x}(y), \delta)$ follows from the implicit function theorem. This completes the proof of Proposition A.1. \square

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References

- [1] F. Béthuel, P. Gravejat, and J.-C. Saut. Existence and properties of travelling waves for the Gross-Pitaevskii equation. In A. Farina and J.-C. Saut, editors, *Stationary and time dependent Gross-Pitaevskii equations*, volume 473 of *Contemp. Math.*, pages 55–104. Amer. Math. Soc., Providence, RI, 2008.
- [2] F. Béthuel, P. Gravejat, J.-C. Saut, and D. Smets. Orbital stability of the black soliton for the Gross-Pitaevskii equation. *Indiana Univ. Math. J.*, 57(6):2611–2642, 2008.
- [3] F. Béthuel, P. Gravejat, J.-C. Saut, and D. Smets. On the Korteweg-de Vries long-wave approximation of the Gross-Pitaevskii equation I. *Int. Math. Res. Not.*, 2009(14):2700–2748, 2009.
- [4] D. Chiron. Travelling waves for the nonlinear Schrödinger equation with general nonlinearity in dimension one. *Preprint*, 2011.
- [5] D. Chiron and F. Rousset. The KdV/KP-I limit of the nonlinear Schrödinger equation. *SIAM J. Math. Anal.*, 42(1):64–96, 2010.
- [6] N. Dunford and J.T. Schwartz. *Linear operators. Part II. Spectral theory. Self-adjoint operators in Hilbert space*, volume 7 of *Pure and Applied Mathematics*. Interscience Publishers, John Wiley and Sons, New York-London-Sydney, 1963. With the assistance of W.G. Bade and R.G. Bartle.
- [7] L.D. Faddeev and L.A. Takhtajan. *Hamiltonian methods in the theory of solitons*. Classics in Mathematics. Springer-Verlag, Berlin-Heidelberg-New York, 2007. Translated by A.G. Reyman.
- [8] C. Gallo. Schrödinger group on Zhidkov spaces. *Adv. Differential Equations*, 9(5-6):509–538, 2004.
- [9] P. Gérard. The Cauchy problem for the Gross-Pitaevskii equation. *Ann. Inst. Henri Poincaré, Analyse Non Linéaire*, 23(5):765–779, 2006.
- [10] P. Gérard and Z. Zhang. Orbital stability of traveling waves for the one-dimensional Gross-Pitaevskii equation. *J. Math. Pures Appl.*, 91(2):178–210, 2009.
- [11] M. Grillakis, J. Shatah, and W.A. Strauss. Stability theory of solitary waves in the presence of symmetry I. *J. Funct. Anal.*, 74(1):160–197, 1987.
- [12] Z. Lin. Stability and instability of traveling solitonic bubbles. *Adv. Differential Equations*, 7(8):897–918, 2002.
- [13] Y. Martel and F. Merle. Stability of two soliton collision for nonintegrable gKdV equations. *Commun. Math. Phys.*, 286(1):39–79, 2009.
- [14] Y. Martel and F. Merle. Inelastic interaction of nearly equal solitons for the quartic gKdV equation. *Invent. Math.*, 183(3):563–648, 2011.
- [15] Y. Martel, F. Merle, and T.-P. Tsai. Stability and asymptotic stability in the energy space of the sum of N solitons for subcritical gKdV equations. *Commun. Math. Phys.*, 231(2):347–373, 2002.
- [16] Y. Martel, F. Merle, and T.-P. Tsai. Stability in H^1 of the sum of K solitary waves for some nonlinear Schrödinger equations. *Duke Math. J.*, 133(3):405–466, 2006.

- [17] R.M. Miura. The Korteweg- de Vries equation: a survey of results. *SIAM Rev.*, 18(3):412–459, 1976.
- [18] H.M. Tartousi. PhD thesis. In preparation.
- [19] A.H. Vartanian. Long-time asymptotics of solutions to the Cauchy problem for the defocusing nonlinear Schrödinger equation with finite-density initial data. II. Dark solitons on continua. *Math. Phys. Anal. Geom.*, 5(4):319–413, 2002.
- [20] V.E. Zakharov and A.B. Shabat. Interaction between solitons in a stable medium. *Sov. Phys. JETP*, 37:823–828, 1973.
- [21] P.E. Zhidkov. *Korteweg-De Vries and nonlinear Schrödinger equations : qualitative theory*, volume 1756 of *Lecture Notes in Mathematics*. Springer-Verlag, Berlin, 2001.