



HAL
open science

Diffusion asymptotics of a kinetic model for gaseous mixtures

Laurent Boudin, Bérénice Grec, Milana Pavic, Francesco Salvarani

► **To cite this version:**

Laurent Boudin, Bérénice Grec, Milana Pavic, Francesco Salvarani. Diffusion asymptotics of a kinetic model for gaseous mixtures. 2012. hal-00704952v1

HAL Id: hal-00704952

<https://hal.science/hal-00704952v1>

Preprint submitted on 6 Jun 2012 (v1), last revised 10 Jul 2012 (v2)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

DIFFUSION ASYMPTOTICS OF A KINETIC MODEL FOR GASEOUS MIXTURES

LAURENT BOUDIN, BÉRÉNICE GREC, MILANA PAVIĆ, AND FRANCESCO SALVARANI

ABSTRACT. In this work, we investigate the asymptotic behaviour of the solutions to the non-reactive fully elastic Boltzmann equations for mixtures in the diffusive scaling. We deal with cross sections such as hard spheres or cut-off power law potentials. We use Hilbert expansions near the common thermodynamic equilibrium granted by the H-theorem. The lower-order non trivial equality obtained from the Boltzmann equations leads to a linear functional equation in the velocity variable which is solved thanks to the Fredholm alternative. Since we consider multicomponent mixtures, the classical techniques introduced by Grad cannot be applied, and we propose a new method to treat the terms involving particles with different masses. The next-order equality in the Hilbert expansion then allows to write the macroscopic continuity equations for each component of the mixture.

1. INTRODUCTION

The study of the asymptotic behaviour of the Boltzmann equation for small mean free path is known as *Hilbert's sixth problem*, after its formulation by Hilbert himself during the International Congress of Mathematicians held in Paris in 1900 [18]. Since then, it is a very active field of research, and many results have been obtained, both at a formal level and in the context of rigorous limits.

The main tools are based on asymptotic (Hilbert, Chapman-Enskog) expansions with respect to the mean free path, see for instance [9]. The translation in a rigorous mathematical language has been performed in a series of pioneering papers by Bardos, Golse and Levermore [1, 2, 3], where the authors established the ground of the subsequent results concerning the rigorous asymptotic limits. They stated a program that led to many interesting results, such as [21]. This trail culminated in [14], where Golse and Saint-Raymond established a Navier-Stokes limit for the Boltzmann equation considered over the infinite spatial domain \mathbb{R}^3 : appropriately scaled families of DiPerna-Lions renormalized solutions are shown to have fluctuations whose limit points are governed by Leray solutions of the limiting Navier-Stokes equations. We can also refer to [15], which extends the results of [14] for hard cutoff potentials in Grad's sense.

Apart from the research concerning with the classical Boltzmann equation (see [10] as a review article), which can be seen as a model describing a mono-species, monoatomic and ideal gas, one can focus on the study, at a kinetic level, of gaseous mixtures, without excluding the possibility of chemical reactions. In such a framework, the models are much more intricate. It is indeed necessary to treat systems of Boltzmann-like equations, rather than one single equation, with multi-species kernels and cross interactions between the different distribution functions describing each component of the mixture. The complexity of the models grows dramatically if exchanges of internal energy and chemical reactions are allowed [23, 25, 7, 24].

The derivation of macroscopic equations from kinetic models remains crucial for mixtures, both at a mathematical level and for deducing relevant macroscopic equations based on the modelling of microscopic binary interactions. In this spirit, in [11], the authors propose a model describing a reacting mixture of polyatomic gases and recover in the limit, via the appropriate scaling $(t, x) \rightarrow (t/\varepsilon, x/\varepsilon)$ for $\varepsilon > 0$, the reactive Euler equations.

The diffusive scaling we investigate here, i.e. $(t, x) \rightarrow (t/\varepsilon^2, x/\varepsilon)$, is more complicated, and even the formal structure of the asymptotic hierarchy is not trivial at all. After the first attempts of Chapman and Cowling [9], in [4], the authors consider a binary mixture of red and blue particles which interact via

strong short range (hard core) and weak long range pair potentials, and study the small free path limit in various situations. Ref. [12] handles the same kind of model of binary mixture with a heavy species and a lighter one.

In this paper, we consider a mixture of several different non reactive gases, which evolves in time via the classical Boltzmann system for non-reacting multicomponent mixtures, and study the formal Hilbert expansion leading to the Navier-Stokes system for mixtures. Even if this strategy dates back to Grad [16, 17], who studied the formal small free path limit for the monatomic and monospecies Boltzmann equation, the development leading to write the Chapman-Enskog hierarchy owns some particular properties, which are inherited from the multispecies feature of the mixture.

First of all, we can observe that the structure of the asymptotics is different for three or more species, in comparison to the case of a monospecies gas or a binary mixture. The most apparent feature which is peculiar to a (at least) ternary mixture is the phenomenon of uphill diffusion, independently foreseen by Maxwell [22] and Stefan [26], and experimentally recovered by Duncan and Toor [13]. This kind of diffusion is described by second-order coupled terms, which make the asymptotics quite different from the limit related to a binary mixture [6].

A second important aspect is that we cannot apply Grad's methodology when we have to deal with different species, with different masses and different cross sections. We hence propose a new approach to the problem which only works when the binary collisions involve particles with different masses, and we then use Grad's procedure [17] when considering collisions between same mass particles. Note that both approaches are needed to give a complete answer to the question of resolvability of the first-order (in ε) equations of the hierarchy, obtained by means of the Fredholm alternative. We point out that our approach holds for generic cut-off cross sections, as those described in Section 2.

Note that this type of asymptotics is also numerically investigated, see, for instance, [5, 19], using BGK approaches.

The article is structured as follows. In the next section, we describe the model, its basic properties and the main results which are proved in the rest of the paper. Then, Sections 3 and 4 are devoted to the proof the mathematical results stated in Section 2. More precisely, Fredholm's alternative is applied in Section 3 (provided the compactness of some operator is known), and the proof of the compactness is treated in Section 4.

2. MODEL

We consider an ideal gas mixture constituted with $I \geq 2$ species. Each species \mathcal{A}_i of the mixture, $1 \leq i \leq I$, is described by a microscopic density function f_i . It depends on time $t \in \mathbb{R}_+$, space position $x \in \mathbb{R}^3$ and molecular velocity $v \in \mathbb{R}^3$, and is nonnegative. More precisely, $f_i(t, x, v) dx dv$ allows to quantify the number of molecules of species \mathcal{A}_i at time t in an elementary volume of size dx , and whose velocities equal v up to dv . We can also define the macroscopic density n_i of each species \mathcal{A}_i by

$$n_i(t, x) = \int_{\mathbb{R}^3} f_i(t, x, v) dv.$$

We assume that the mixture only involves molecular elastic collisions. Let us consider two colliding molecules of species \mathcal{A}_i and \mathcal{A}_j , $1 \leq i, j \leq I$. Their masses are m_i and m_j , and their pre-collisional velocities v' and v'_* . After a collision, the particles belong to the same species (no chemical reactions), so their masses remains as they were, and their velocities become v and v_* . Since the collisions are elastic, both momentum and kinetic energy are conserved, i.e.

$$(1) \quad m_i v' + m_j v'_* = m_i v + m_j v_*, \quad \frac{1}{2} m_i v'^2 + \frac{1}{2} m_j v'^2_* = \frac{1}{2} m_i v^2 + \frac{1}{2} m_j v^2_*.$$

Consequently, v' and v'_* can be written in terms of v and v_* :

$$(2) \quad v' = \frac{m_i v + m_j v_*}{m_i + m_j} + \frac{m_j}{m_i + m_j} T_\omega(v - v_*), \quad v'_* = \frac{m_i v + m_j v_*}{m_i + m_j} - \frac{m_i}{m_i + m_j} T_\omega(v - v_*),$$

where $\omega \in S^2$ is arbitrary, and T_ω is the symmetry with respect to the plane $\{\omega\}^\perp$, i.e.

$$T_\omega w = w - 2(\omega \cdot w)\omega, \quad \forall w \in \mathbb{R}^3.$$

2.1. Collision operators. Let $1 \leq i, j \leq I$. The collision operator associated to species \mathcal{A}_i and \mathcal{A}_j is defined by

$$(3) \quad Q_{ij}(f, g)(v) = \iint_{\mathbb{R}^3 \times S^2} [f(v')g(v'_*) - f(v)g(v_*)] \check{\mathcal{B}}_{ij}(v, v_*, \omega) \, d\omega \, dv_*,$$

where v' and v'_* are defined by (2), and f and g are two functions of the velocity variable. The cross-section $\check{\mathcal{B}}_{ij}(v, v_*, \omega)$ only depends on v , v_* and ω .

In this work, $\check{\mathcal{B}}_{ij}$ is a function of $|v - v_*|$ and the angle θ between ω and $V := v - v_*$. Let us set

$$\mathcal{B}_{ij}(\omega, v - v_*) = \check{\mathcal{B}}_{ij}(v, v_*, \omega), \quad \forall \omega \in S^2, \quad \forall v, v_* \in \mathbb{R}^3.$$

The collisions are also supposed microreversible. That ensures that

$$(4) \quad \mathcal{B}_{ij}(\omega, v - v_*) = \mathcal{B}_{ji}(\omega, v - v_*), \quad \mathcal{B}_{ij}(\omega, v - v_*) = \mathcal{B}_{ij}(\omega, v' - v'_*), \quad \forall \omega \in S^2, \quad \forall v, v_* \in \mathbb{R}^3.$$

Moreover, we assume that \mathcal{B}_{ij} satisfies a general condition

$$(5) \quad \mathcal{B}_{ij}(\omega, V) \leq a |\sin \theta| |\cos \theta| \left(|V| + \frac{1}{|V|^{1-\delta}} \right), \quad \forall \omega \in S^2, \quad \forall V \in \mathbb{R}^3,$$

where $a > 0$, $0 < \delta < 1$. As emphasized in [17], this corresponds to intermolecular potentials with finite range and it means that \mathcal{B}_{ij} linearly approaches 0 near $\theta = 0$ and $\theta = \pi/2$, and is of restricted growth for both small and large $|V|$.

Condition (5) is, for instance, satisfied by hard spheres of diameter $\sigma_{ij} > 0$:

$$\mathcal{B}_{ij}(\omega, V) = \sigma_{ij}^2 |V| \sin \theta \cos \theta$$

and, by all cutoff power-law potentials:

$$\mathcal{B}_{ij}(\omega, V) = |V|^{\gamma_{ij}} \beta_{ij}(\theta), \quad \gamma_{ij} = \frac{s_{ij} - 5}{s_{ij} - 1},$$

where $\beta_{ij}(\theta)$ is a bounded function and linearly approaches 0 when θ goes to $\pi/2$, and $s_{ij} > 3$. Note that the previous condition does not allow to recover the monoatomic case (that would mean that $s_{ij} = 3$, i.e. $\gamma_{ij} = -1$), but includes the diatomic case (where $s_{ij} = 5$ and $\gamma_{ij} = 0$).

The collision operators can be also written under weak forms, obtained from (3) using the changes of variables $(v, v_*) \mapsto (v_*, v)$ and $(v, v_*) \mapsto (v', v'_*)$ for a fixed $\omega \in S^2$. Weak forms in the cases $i = j$ and $i \neq j$ are intrinsically different, and are explained in detail in [11, 6]. Mention that, if we choose suitable test-functions, the weak forms of (3) allow to formally write, for any i and j , and any functions f and g for which the following equations make sense:

$$(6) \quad \int_{\mathbb{R}^3} Q_{ij}(f, g)(v) \, dv = 0,$$

$$(7) \quad \int_{\mathbb{R}^3} Q_{ij}(f, g)(v) \begin{pmatrix} m_i v \\ m_i v^2/2 \end{pmatrix} \, dv + \int_{\mathbb{R}^3} Q_{ji}(g, f)(v) \begin{pmatrix} m_j v \\ m_j v^2/2 \end{pmatrix} \, dv = 0.$$

2.2. H-theorem. Let us now write down the H-theorem corresponding to the collisional operators we defined in the previous subsection, and discuss the mechanical equilibrium. More precisely, the following result holds [6] (see also [11]).

Proposition 1. *We assume that the cross sections $(\check{\mathcal{B}}_{ij})_{1 \leq i, j \leq I}$ are nonnegative almost everywhere.*

(a) For all $f_i \equiv f_i(v) \geq 0$, $1 \leq i \leq I$, such that the following quantities are defined, one has

$$\sum_{i=1}^I \sum_{j=1}^I \int_{\mathbb{R}^3} Q_{ij}(f_i, f_j)(v) \log \left(\frac{f_i(v)}{m_i^3} \right) dv \leq 0.$$

(b) Moreover, the three following properties are equivalent.

i. For any $1 \leq i, j \leq I$ and $v \in \mathbb{R}^3$

$$(8) \quad Q_{ij}(f_i, f_j)(v) = 0.$$

ii. The previous inequality is an equality:

$$\sum_{i=1}^I \sum_{j=1}^I \int_{\mathbb{R}^3} Q_{ij}(f_i, f_j)(v) \log \left(\frac{f_i(v)}{m_i^3} \right) dv = 0.$$

iii. There exist $T > 0$ and $u \in \mathbb{R}^3$ such that, for any i , there exists $n_i \geq 0$ such that

$$(9) \quad f_i(v) = n_i \left(\frac{m_i}{2\pi kT} \right)^{3/2} e^{-\frac{m_i}{2kT}|v-u|^2}.$$

2.3. Statement of the problem. In this work, we focus on the diffusion limit of the Boltzmann equations for mixtures. That limit is obtained from the framework of the classical diffusive scaling, where the scaling parameter is the mean free path. Let us choose $\varepsilon > 0$ as the mean free path. Hence, for any i , each distribution function f_i^ε must solve the following scaled Boltzmann equation, that is

$$(10) \quad \varepsilon \partial_t f_i^\varepsilon + v \cdot \nabla_x f_i^\varepsilon = \frac{1}{\varepsilon} \sum_{j=1}^I Q_{ij}(f_i^\varepsilon, f_j^\varepsilon), \quad t > 0, x \in \mathbb{R}^3, v \in \mathbb{R}^3.$$

Thanks to Proposition 1, the Maxwellian distributions (9) are zero-th order in ε solutions to (10). Therefore, each distribution function f_i^ε , $1 \leq i \leq I$, can be seen as a perturbation of the equilibrium (9). Without any loss of generality, we set $u = 0$ (diffusion limit) and $kT = 1$. Thus we can write f_i^ε as

$$(11) \quad f_i^\varepsilon(t, x, v) = M_i(v) n_i(t, x) + \varepsilon M_i(v)^{1/2} g_i(t, x, v) + O(\varepsilon^2), \quad \forall t \geq 0, \forall x, v \in \mathbb{R}^3.$$

where $M_i(v)$ is the normalized, centred Maxwell function

$$M_i(v) = \left(\frac{m_i}{2\pi} \right)^{3/2} e^{-\frac{m_i}{2} v^2}, \quad \forall v \in \mathbb{R}^3.$$

In (11), we choose to put $M_i(v)^{1/2}$ within the first-order term of f_i^ε . As a matter of fact, it allows us to work in a plain L^2 framework in the variable v for g_i .

We insert the first-order Hilbert expansion (11) into (10), then consider expressions of the same order (up to one) in ε . Taking (8) into account, we obtain the following equations, holding for any $1 \leq i \leq I$,

$$(12) \quad M_i^{-1/2} \sum_{j=1}^I \left(n_i Q_{ij}(M_i, M_j^{1/2} g_j) + n_j Q_{ij}(M_i^{1/2} g_i, M_j) \right) = M_i^{1/2} (v \cdot \nabla_x n_i),$$

$$(13) \quad M_i \partial_t n_i + M_i^{1/2} v \cdot \nabla_x g_i = \sum_{j=1}^I Q_{ij}(M_i^{1/2} g_i, M_j^{1/2} g_j).$$

Let us now focus on equation (12). We denote $g = (g_1, \dots, g_I)^T$ and write the L^2 norm of g in the variable v :

$$\|g\|_{L^2}^2 = \sum_{j=1}^I \|g_j\|_{L^2}^2 = \sum_{j=1}^I \int_{\mathbb{R}^3} g_j(t, x, v)^2 dv,$$

which in fact depends on t, x . However, in (12), the dependence on t and x is not crucial, in the sense that $t, x, (n_j)$ and $\nabla_x n_i$ can be seen as parameters, and g is the unknown function of v .

We can write the left-hand side of (12) in a more suitable form if we introduce the operator \mathcal{K} , where the i -th component of $\mathcal{K}g$ is given by

$$(14) \quad [\mathcal{K}g]_i(v) = \sum_{j=1}^I \left(\frac{m_j}{2\pi}\right)^{3/4} \iint_{\mathbb{R}^3 \times S^2} \mathcal{B}_{ij}(\omega, v - v_*) e^{-\frac{1}{4}m_i v^2} e^{-\frac{1}{2}m_j v_*^2} \left[n_i \left(\frac{m_i}{2\pi}\right)^{3/4} \left(e^{\frac{1}{4}m_j v_*^2} g_j(v'_*) - e^{\frac{1}{4}m_j v_*^2} g_j(v_*) \right) + n_j \left(\frac{m_j}{2\pi}\right)^{3/4} e^{\frac{1}{4}m_i v^2} g_i(v') \right] d\omega dv_*,$$

for any i , and the positive function $\nu = \nu(v)$, whose i -th component is

$$\nu_i(v) = \sum_{j=1}^I n_j \left(\frac{m_j}{2\pi}\right)^{3/2} \iint_{\mathbb{R}^3 \times S^2} e^{-\frac{1}{2}m_j v_*^2} \mathcal{B}_{ij}(\omega, v - v_*) d\omega dv_*.$$

Consequently, g solves the following functional equation

$$(15) \quad (\mathcal{K} - \nu \text{Id})g = \left(M_i^{1/2} (v \cdot \nabla_x n_i) \right)_{1 \leq i \leq I}.$$

Let us now state the main result of this work.

Theorem 1. *Consider $(\mathcal{B}_{ij})_{1 \leq i, j \leq I}$ nonnegative functions satisfying (5). If $(n_i(t, x)M_i)_{1 \leq i \leq I}$ are equilibria of the collisional operators $(Q_{ij})_{1 \leq i, j \leq I}$ satisfying*

$$(16) \quad \sum_{i=1}^I n_i(t, x) \text{ does not depend on } x,$$

then (10) has a solution f_i^ε close to the equilibrium $n_i M_i$. More precisely, if f_i^ε satisfies (11) for any i , then, for every t, x , there exists $g(t, x, \cdot) \in L^2(\mathbb{R}_v^3)$ satisfying (15).

Condition (16) means that we consider a situation where the total number density of gaseous particles is uniform in space, see 2.4 for further discussions. Let us briefly draw the sketch of the proof of this theorem. We shall first need the following proposition.

Proposition 2. *The operator \mathcal{K} , defined by (14), is a compact operator $L^2(\mathbb{R}_v^3) \rightarrow L^2(\mathbb{R}_v^3)$.*

The proof of Proposition 2 is given in Section 4. Let us emphasize that, in Proposition 2, t and x are considered as parameters and the compactness is only related to the variable v . Since \mathcal{K} is compact, we can apply the Fredholm alternative to the operator $\mathcal{K} - \nu \text{Id}$. This is explained in the next section.

2.4. Diffusion limit. To conclude this section, let us add a few words about the diffusion limit itself. Theorem 1 ensures the existence of g_i , $1 \leq i \leq I$, in the expansion (11). Integrating (13) with respect to v and using (6)–(7), we can write

$$\partial_t n_i(t, x) + \nabla_x \cdot \int_{\mathbb{R}^3} v g_i(t, x, v) M_i(v)^{1/2} dv = 0,$$

which is the usual continuity equation

$$\partial_t n_i(t, x) + \nabla_x \cdot N_i(t, x) = 0,$$

where the flux N_i of species \mathcal{A}_i is given for any i by

$$N_i(t, x) = \int_{\mathbb{R}^3} v g_i(t, x, v) M_i(v)^{1/2} dv, \quad 1 \leq i \leq I.$$

Macroscopic relationships between the fluxes (N_i) and the (macroscopic) densities (n_i) may be obtained, with further assumptions on the cross sections, by multiplying the Boltzmann equations (10) by v and integrating them with respect to v . The reader can refer to [6] for more details. A closure relationship is also needed, and the standard equimolar diffusion assumption can be used, as it is common in closed

experimental settings [13, 20]: $\sum N_i = 0$. This implies that $\sum n_i$ equals its initial value. Hence, if we suppose that the initial value does not depend on x , assumption (16) is fulfilled. Of course, there are other situations allowing (16) to hold.

3. PROOF OF THEOREM 1

This section is devoted to the proof of the main result of our article. We here assume that the compactness of operator \mathcal{K} is known. Let us denote $\mathcal{L} = \mathcal{K} - \nu \text{Id}$, and study the null space of \mathcal{L} , as required by the Fredholm alternative.

Step 1 – Study of $\ker \mathcal{L}$. Writing down the i -th component of $\mathcal{L}g$ and performing the change of variable $(v, v_*) \mapsto (v', v'_*)$ while ω remains fixed, we obtain

$$[\mathcal{L}g]_i(v) = \sum_{j=1}^I \iint_{\mathbb{R}^3 \times S^2} M_i(v)^{1/2} M_j(v_*) \left[n_j M_i(v')^{-1/2} g_i(v') + n_i M_j(v'_*)^{-1/2} g_j(v'_*) \right. \\ \left. - n_j M_i(v)^{-1/2} g_i(v) - n_i M_j(v_*)^{-1/2} g_j(v_*) \right] \mathcal{B}_{ij}(\omega, v - v_*) d\omega dv_*.$$

Thanks to the H-theorem, $g \in \ker \mathcal{L}$ if and only if there exist $\alpha \in \mathbb{R}^I$, $\beta \in \mathbb{R}^3$, $\gamma \in \mathbb{R}$ such that, for any i ,

$$g_i(t, x, v) = n_i(t, x) M_i(v)^{1/2} \left(\alpha_i + m_i \beta \cdot v + \gamma \frac{m_i}{2} v^2 \right), \quad \forall t > 0, x, v \in \mathbb{R}^3.$$

Consequently, $\ker \mathcal{L} \neq \{0\}$, and the Fredholm alternative allows to state that (15) has a solution if and only if

$$(17) \quad \left(M_i^{1/2} (v \cdot \nabla_x n_i) \right)_{i=1, \dots, I} \in (\ker \mathcal{L}^*)^\perp, \quad \forall t > 0, x \in \mathbb{R}^3.$$

Step 2 – Computation of \mathcal{L}^* . Let us compute the adjoint operator \mathcal{L}^* by studying the inner product between $\mathcal{L}g$ and a vector $h \in L^2(\mathbb{R}_v^3)^I$. We successively write, using the change of variables $(v, v_*) \mapsto (v', v'_*)$ and $(v, v_*) \mapsto (v_*, v)$

$$\begin{aligned} & \sum_{i=1}^I \int_{\mathbb{R}^3} [\mathcal{L}g]_i(v) h_i(v) dv \\ &= \sum_{i,j=1}^I \iiint_{\mathbb{R}^3 \times \mathbb{R}^3 \times S^2} h_i(v) M_i(v)^{-1/2} \left[n_i M_i(v') M_j^{1/2}(v'_*) g_j(v'_*) - n_i M_i(v) M_j(v_*)^{1/2} \hat{g}_j(v_*) \right. \\ & \quad \left. + n_j M_j(v'_*) M_i(v')^{1/2} g_i(v') - n_j M_j(v_*) M_i(v)^{1/2} g_i(v) \right] \mathcal{B}_{ij}(\omega, v - v_*) d\omega dv_* dv \\ &= \sum_{i,j=1}^I \iiint_{\mathbb{R}^3 \times \mathbb{R}^3 \times S^2} g_i(v) n_j M_i(v)^{-1/2} \left[M_i(v') M_j^{1/2}(v'_*) h_j(v'_*) - M_i(v) M_j(v_*)^{1/2} h_j(v_*) \right. \\ & \quad \left. + M_j(v'_*) M_i(v')^{1/2} h_i(v') - M_j(v_*) M_i(v)^{1/2} h_i(v) \right] \mathcal{B}_{ij}(\omega, v - v_*) d\omega dv_* dv. \end{aligned}$$

Consequently, we have

$$[\mathcal{L}^* h]_i = M_i^{-1/2} \sum_{j=1}^I n_j \left(Q_{ij}(M_i, M_j^{1/2} h_j) + Q_{ij}(M_i^{1/2} h_i, M_j) \right).$$

Thanks to the H-theorem, $h \in \ker \mathcal{L}^*$ if and only if there exist $a \in \mathbb{R}^I$, $b \in \mathbb{R}^3$, $c \in \mathbb{R}$ such that, for any i ,

$$(18) \quad h_i(v) = M_i(v)^{1/2} \left(a_i + m_i b \cdot v + c \frac{m_i}{2} v^2 \right), \quad \forall v \in \mathbb{R}^3.$$

Step 3 – Conclusion. Now, taking (18) into account, condition (17) can be rewritten as

$$\sum_{i=1}^I \sum_{k=1}^3 \frac{\partial n_i}{\partial x_k} \int_{\mathbb{R}^3} \begin{pmatrix} v_k \\ m_i v_k v_j \\ m_i v_k v^2/2 \end{pmatrix} M_i(v) dv = 0, \quad 1 \leq j \leq 3.$$

Using parity arguments, the first and third integrals are immediately satisfied, as well as the second ones if $k \neq j$. In the case when $k = j$, the condition $\nabla_x \sum n_i = 0$, which is assumed in (16), allows to complete the proof.

4. PROOF OF PROPOSITION 2

We still have to prove that \mathcal{K} is compact. In this section, $(n_i)_{1 \leq i \leq I}$ are assumed to be nonnegative constants. Let $g \in L^2(\mathbb{R}^3)^I$. Note that we do not need g to be the function defined in (11). First, we write \mathcal{K} as the sum of four operators $\mathcal{K}_1, \dots, \mathcal{K}_4$. For any i , the i -th component of each $\mathcal{K}_\ell g$, $1 \leq \ell \leq 4$, is given by

$$\begin{aligned} [\mathcal{K}_1 g]_i(v) &= -n_i \sum_{j=1}^I \left(\frac{m_i m_j}{4\pi^2} \right)^{3/4} \iint_{\mathbb{R}^3 \times S^2} e^{-\frac{1}{4}m_i v^2} e^{-\frac{1}{4}m_j v_*^2} g_j(v_*) \mathcal{B}_{ij}(\omega, v - v_*) d\omega dv_*, \\ [\mathcal{K}_2 g]_i(v) &= n_i \sum_{j \notin \mathcal{M}_i} \left(\frac{m_i m_j}{4\pi^2} \right)^{3/4} \iint_{\mathbb{R}^3 \times S^2} e^{-\frac{1}{4}m_i v^2} e^{-\frac{1}{2}m_j v_*^2} e^{\frac{1}{4}m_j v_*'^2} g_j(v_*') \mathcal{B}_{ij}(\omega, v - v_*) d\omega dv_*, \\ [\mathcal{K}_3 g]_i(v) &= \sum_{j \in \mathcal{M}_i} \left(\frac{m_i}{2\pi} \right)^{3/2} \iint_{\mathbb{R}^3 \times S^2} e^{-\frac{1}{4}m_i v^2} e^{-\frac{1}{2}m_i v_*^2} \\ &\quad \left[n_i e^{\frac{1}{4}m_i v_*'^2} g_j(v_*') + n_j e^{\frac{1}{4}m_i v_*'^2} g_i(v_*') \right] \mathcal{B}_{ij}(\omega, v - v_*) d\omega dv_*, \\ [\mathcal{K}_4 g]_i(v) &= \sum_{j \notin \mathcal{M}_i} n_j \left(\frac{m_j}{2\pi} \right)^{3/2} \iint_{\mathbb{R}^3 \times S^2} e^{-\frac{1}{4}m_i v^2} e^{-\frac{1}{2}m_j v_*^2} e^{\frac{1}{4}m_i v_*'^2} g_i(v_*') \mathcal{B}_{ij}(\omega, v - v_*) d\omega dv_*. \end{aligned}$$

We denoted, for any i ,

$$\mathcal{M}_i := \{1 \leq j \leq I \mid m_j = m_i\},$$

which is non empty since $i \in \mathcal{M}_i$. It is crucial to dissociate cases when $m_i \neq m_j$ or $m_i = m_j$, because proofs are quite different.

We successively prove that \mathcal{K}_ℓ , $1 \leq \ell \leq 4$, is compact. In order to obtain the compactness, we prove two properties, a uniform decay at infinity

$$(19) \quad \| [\mathcal{K}_\ell g]_i \|_{L^2(B(0,R)^c)} \leq \sigma_\ell(R) \|g\|_{L^2(\mathbb{R}^3)}, \quad \forall R > 0,$$

where $B(0, R)$ denotes the open ball of \mathbb{R}_v^3 centred at 0 and of radius R , and $\sigma_\ell(R)$ goes to 0 when R goes to $+\infty$, and the equicontinuity

$$(20) \quad \| [(\tau_w - I)\mathcal{K}_\ell g]_i \|_{L^2(\mathbb{R}^3)} \leq \varrho_\ell(w) \|g\|_{L^2(\mathbb{R}^3)}, \quad \forall w \in \mathbb{R}^3,$$

where τ_w denotes the translation operator, i.e.

$$\tau_w \mathcal{K}_\ell g(v) = \mathcal{K}_\ell g(v + w), \quad \forall v, w \in \mathbb{R}^3,$$

and $\varrho_\ell(w)$ goes to 0 when w goes to 0.

The compactness of each \mathcal{K}_ℓ is an immediate consequence of (19)–(20).

4.1. **Compactness of \mathcal{K}_1 .** Let us denote, for any i, j ,

$$k_1^{ij}(v, v_*) = \int_{S^2} e^{-\frac{1}{4}m_i v^2} e^{-\frac{1}{4}m_j v_*^2} \mathcal{B}_{ij}(\omega, v - v_*) d\omega, \quad \forall v, v_* \in \mathbb{R}^3.$$

We immediately have, for any i ,

$$[\mathcal{K}_1 g]_i(v) = -n_i \sum_{j=1}^I \left(\frac{m_i m_j}{4\pi^2} \right)^{3/4} \int_{\mathbb{R}^3} g_j(v_*) k_1^{ij}(v, v_*) dv_*, \quad \forall v \in \mathbb{R}^3.$$

4.1.1. *Properties of k_1^{ij} .* First of all, note that $k_1^{ij}(v, v_*) = k_1^{ji}(v_*, v)$, for any i, j and v, v_* , thanks to (4). We need preliminary properties of k_1^{ij} to obtain (19)–(20) for \mathcal{K}_1 .

Lemma 1. *There exists $C > 0$ such that, for any i, j ,*

$$(21) \quad \int_{\mathbb{R}^3} k_1^{ij}(v, v_*) dv_* \leq C e^{-\frac{1}{4}m_i v^2} (1 + |v|), \quad \forall v \in \mathbb{R}^3.$$

Proof. Thanks to (5) and using the change of variables $v_* \mapsto V_* = v_* - v$, we can write

$$\begin{aligned} \int_{\mathbb{R}^3} k_1^{ij}(v, v_*) dv_* &\leq C \int_{\mathbb{R}^3} e^{-\frac{1}{4}m_i v^2} e^{-\frac{1}{4}m_j (V_* + v)^2} (|V_*| + |V_*|^{\delta-1}) dV_* \\ &\leq C e^{-\frac{1}{4}m_i v^2} \left[\int_{|V_*| \leq 1} (|V_*| + |V_*|^{\delta-1}) dV_* + \int_{|V_*| \geq 1} (1 + |V_*|) e^{-\frac{1}{4}m_j (V_* + v)^2} dV_* \right]. \end{aligned}$$

This inequality allows us to get (21). □

Lemma 2. *For any i, j , k_1^{ij} belongs to $L^2(\mathbb{R}^3 \times \mathbb{R}^3)$.*

Proof. The proof follows the same strategy as the previous one, using (5) and the same change of variables. We can write

$$\begin{aligned} \iint_{\mathbb{R}^3 \times \mathbb{R}^3} k_1^{ij}(v, v_*)^2 dv_* dv &\leq C \iint_{\mathbb{R}^3 \times \mathbb{R}^3} e^{-\frac{1}{2}m_i v^2} e^{-\frac{1}{2}m_j v_*^2} (|v - v_*|^2 + |v - v_*|^{2\delta-2}) dv_* \\ &\leq C \int_{\mathbb{R}^3} e^{-\frac{1}{2}m_i v^2} (1 + v^2) dv, \end{aligned}$$

which is clearly finite. □

4.1.2. *Uniform decay.* The L^2 norm of $\mathcal{K}_1 g$ decreases at infinity. More precisely, the following proposition holds.

Proposition 3. *Let $g \in L^2(\mathbb{R}^3)^I$. For any $R > 0$ and any i , we have*

$$(22) \quad \|\mathcal{K}_1 g\|_{L^2(B(0, R)^c)} \leq \frac{C n_i}{R} \|g\|_{L^2(\mathbb{R}^3)},$$

where $C > 0$ is a constant.

Proof. Let $1 \leq i \leq I$ and write

$$\int_{\mathbb{R}^3} v^2 [\mathcal{K}_1 g]_i(v)^2 dv \leq C n_i^2 \sum_{j=1}^I \int_{\mathbb{R}^3} v^2 \left[\int_{\mathbb{R}^3} k_1^{ij}(v, v_*) g_j(v_*) dv_* \right]^2 dv.$$

Thanks to the Cauchy-Schwarz inequality, the previous inequality becomes

$$\int_{\mathbb{R}^3} v^2 [\mathcal{K}_1 g]_i(v)^2 dv \leq C n_i^2 \sum_{j=1}^I \int_{\mathbb{R}^3} v^2 \left[\int_{v_* \in \mathbb{R}^3} g_j(v_*)^2 k_1^{ij}(v, v_*) dv_* \right] \left[\int_{\mathbb{R}^3} k_1^{ij}(v, v_*) dv_* \right] dv.$$

Using Lemma 1 and afterwards the Fubini theorem, we get

$$\int_{\mathbb{R}^3} v^2 [\mathcal{K}_1 g]_i(v)^2 dv \leq C n_i^2 \sum_{j=1}^I \int_{\mathbb{R}^3} g_j(v_*)^2 \left[\int_{\mathbb{R}^3} k_1^{ij}(v, v_*) \phi_i(v) dv \right] dv_*,$$

where $\phi_i(v) = v^2 (1 + |v|) e^{-\frac{1}{4}m_i v^2}$ is clearly bounded. Consequently, since $k_1^{ij}(v, v_*) = k_1^{ji}(v_*, v)$, we have

$$\int_{\mathbb{R}^3} v^2 [\mathcal{K}_1 g]_i(v)^2 dv \leq C n_i^2 \sum_{j=1}^I \int_{\mathbb{R}^3} g_j(v_*)^2 \left(\int_{\mathbb{R}^3} k_1^{ji}(v_*, v) dv \right) dv_*.$$

Eventually, using Lemma 1 again, we obtain

$$\int_{\mathbb{R}^3} v^2 [\mathcal{K}_1 g]_i(v)^2 dv \leq C n_i^2 \|g\|_{L^2}^2.$$

Besides, we can write, for any $R > 0$,

$$\int_{\mathbb{R}^3} v^2 [\mathcal{K}_1 g]_i(v)^2 dv \geq \int_{|v| \geq R} v^2 [\mathcal{K}_1 g]_i(v)^2 dv \geq R^2 \int_{|v| \geq R} [\mathcal{K}_1 g]_i(v)^2 dv.$$

It is then easy to recover (22). □

4.1.3. *Equicontinuity.* The following property of equicontinuity of \mathcal{K}_1 holds.

Proposition 4. *For any $w \in \mathbb{R}^3$, set*

$$\varrho_1(w) = C \max_{i,j} \left[n_i \iint_{\mathbb{R}^3 \times \mathbb{R}^3} \left(k_1^{ij}(v+w, v_*) - k_1^{ij}(v, v_*) \right)^2 dv_* dv \right]^{1/2},$$

where C is a suitable nonnegative constant. Then, for any i , we have

$$(23) \quad \|[(\tau_w - \text{Id})\mathcal{K}_1 g]_i\|_{L^2(\mathbb{R}_v^3)} \leq \varrho_1(w) \|g\|_{L^2(\mathbb{R}_v^3)}, \quad \forall w \in \mathbb{R}^3,$$

and $\varrho_1(w)$ goes to 0 when w goes to 0.

Proof. First, thanks to Lemma 2, it is clear that ϱ_1 is a continuous function of w , and goes to 0 when w goes to zero. Let us now focus on (23). For any i , using the Cauchy-Schwarz inequality, we have

$$\|[(\tau_w - \text{Id})\mathcal{K}_1 g]_i\|_{L^2(\mathbb{R}^3)}^2 \leq C n_i^2 \sum_{j=1}^I \|g_j\|_{L^2}^2 \iint_{\mathbb{R}^3 \times \mathbb{R}^3} \left(k_1^{ij}(v+w, v_*) - k_1^{ij}(v, v_*) \right)^2 dv_* dv.$$

Estimate (23) is an immediate consequence of the previous inequality. □

4.2. **Compactness of \mathcal{K}_2 .** As in section 4.1, we first write \mathcal{K}_2 in a more convenient form. Indeed, thanks to (1), we have

$$-\frac{1}{4}m_i v^2 - \frac{1}{2}m_j v_*^2 + \frac{1}{4}m_j v_*'^2 = -\frac{1}{4}m_j v_*^2 - \frac{1}{4}m_i v'^2.$$

Hence, $[\mathcal{K}_2]_i$ becomes

$$(24) \quad [\mathcal{K}_2 g]_i(v) = \sum_{j \notin \mathcal{M}_i} n_i \left(\frac{m_i m_j}{4\pi^2} \right)^{3/4} \iint_{\mathbb{R}^3 \times S^2} e^{-\frac{1}{4}m_j v_*^2} e^{-\frac{1}{4}m_i v'^2} g_j(v_*') \mathcal{B}_{ij}(\omega, v - v_*) d\omega dv_*.$$

The main idea is to recover a kernel form of (24) to be able to apply the same strategy as for \mathcal{K}_1 . In order to get this particular form, we need the following lemma.

Lemma 3. *There exists $b > 0$ such that, for any i, j satisfying $m_i \neq m_j$,*

$$(25) \quad m_i v'^2 + m_j v_*^2 \geq b (m_i v^2 + m_j v_*'^2)$$

for any $v, v_* \in \mathbb{R}^3$ and v', v_*' given by (2).

Remark. The assumption on the masses is here crucial, as we shall see in the proof. Indeed, (25) somehow gives a property of norm equivalence in $\mathbb{R}^3 \times \mathbb{R}^3$, linking (v, v'_*) and (v', v_*) . Such a property does not hold when we deal with molecules of the same mass.

Proof. Choose $j \notin \mathcal{M}_i$. Equation (2) can be rewritten as

$$(26) \quad v' = \left(\mathbf{I}_3 - 2 \frac{m_j}{m_i + m_j} \omega \omega^T \right) v + 2 \frac{m_j}{m_i + m_j} \omega \omega^T v_*,$$

$$(27) \quad v'_* = \left(\mathbf{I}_3 - 2 \frac{m_i}{m_i + m_j} \omega \omega^T \right) v_* + 2 \frac{m_i}{m_i + m_j} \omega \omega^T v,$$

where \mathbf{I}_3 is the identity matrix of \mathbb{R}^3 . Then, from (27), we get

$$\left(\mathbf{I}_3 - 2 \frac{m_i}{m_i + m_j} \omega \omega^T \right) v_* = v'_* - 2 \frac{m_i}{m_i + m_j} \omega \omega^T v.$$

Let us now set

$$\mathbf{A} = \mathbf{I}_3 - 2 \frac{m_i}{m_i + m_j} \omega \omega^T.$$

This matrix \mathbf{A} is invertible, since $\det \mathbf{A} = (m_j - m_i) / (m_i + m_j)$ and $j \notin \mathcal{M}_i$. Consequently, we can write

$$(28) \quad v_* = (\mathbf{I}_3 - \mathbf{A}^{-1}) v + \mathbf{A}^{-1} v'_*,$$

where we used the equality

$$-2 \frac{m_i}{m_i + m_j} \mathbf{A}^{-1} \omega \omega^T = \mathbf{I}_3 - \mathbf{A}^{-1}.$$

Then we put (28) in (26) to obtain

$$v' = \left(\frac{m_i + m_j}{m_i} \mathbf{I}_3 - \frac{m_j}{m_i} \mathbf{A}^{-1} \right) v - \frac{m_j}{m_i} (\mathbf{I}_3 - \mathbf{A}^{-1}) v'_*.$$

Consider now the following block matrix

$$\mathbb{A} = \begin{bmatrix} \frac{m_i + m_j}{m_i} \mathbf{I}_3 - \frac{m_j}{m_i} \mathbf{A}^{-1} & -\sqrt{\frac{m_j}{m_i}} (\mathbf{I}_3 - \mathbf{A}^{-1}) \\ \sqrt{\frac{m_j}{m_i}} (\mathbf{I}_3 - \mathbf{A}^{-1}) & \mathbf{A}^{-1} \end{bmatrix},$$

which is invertible: $\det \mathbb{A} = -1$ and $\mathbb{A}^{-1} = \mathbb{A}$. The following vector equality holds:

$$\begin{bmatrix} \sqrt{m_i} v' \\ \sqrt{m_j} v_* \end{bmatrix} = \mathbb{A} \begin{bmatrix} \sqrt{m_i} v \\ \sqrt{m_j} v'_* \end{bmatrix}.$$

In fact, (25) is obtained by finding a lower bound of

$$\frac{\left| \mathbb{A} \begin{bmatrix} \sqrt{m_i} v & \sqrt{m_j} v'_* \end{bmatrix}^T \right|^2}{\left| \begin{bmatrix} \sqrt{m_i} v & \sqrt{m_j} v'_* \end{bmatrix}^T \right|^2},$$

which is $\|\mathbb{A}^{-1}\|_2^{-2} = \|\mathbb{A}\|_2^{-2}$. Since \mathbb{A} is clearly a continuous function of ω , we may conclude that the matrix norm $\|\mathbb{A}\|_2$ is a positive continuous function of ω on S^2 , which is compact. Therefore, it reaches its minimum, which of course remains positive. Choosing

$$b = \min_{\omega \in S^2} \|\mathbb{A}\|_2^{-2} > 0$$

leads to the required estimate (25). □

Using (5) and Lemma 3, we obtain the upper bound

$$[\mathcal{K}_2 g]_i(v) \leq C n_i \sum_{j \notin \mathcal{M}_i} e^{-\frac{1}{4} m_j v^2} \int_{\mathbb{R}^3} e^{-\frac{1}{4} m_i v_*'^2} g_j(v_*') \left(|v - v_*| + |v - v_*|^{\delta-1} \right) dv_*.$$

Let us then perform the change of variable $v_* \mapsto v_*'$, whose Jacobian is $1/\det A$. Since

$$v - v_* = A^{-1} (v - v_*') \quad \text{and} \quad \|A\|_2^{-1} \leq \frac{|A^{-1}(v - v_*')|}{|v - v_*'|} \leq \|A^{-1}\|_2,$$

we can write

$$|v - v_*| + |v - v_*|^{\delta-1} \leq \|A^{-1}\|_2 |v - v_*'| + \|A\|_2^{1-\delta} |v - v_*'|^{\delta-1}.$$

Eventually, we obtain

$$[\mathcal{K}_2 g]_i(v) \leq C n_i \sum_{j \notin \mathcal{M}_i} \iint_{\mathbb{R}^3 \times S^2} e^{-\frac{1}{4} m_i v^2} e^{-\frac{1}{4} m_j v_*'^2} g_j(v_*') \left(|v - v_*'| + |v - v_*'|^{\delta-1} \right) d\omega dv_*'.$$

The upper bound in the previous equality has exactly a kernel form, which allows us to conclude on the compactness of \mathcal{K}_2 in the same way as in section 4.1.

4.3. Compactness of \mathcal{K}_3 . This operator deals with a situation where the molecular masses are equal. Note that it does not mean that species \mathcal{A}_i and \mathcal{A}_j are the same, since \mathcal{B}_{ii} , \mathcal{B}_{ij} and \mathcal{B}_{jj} can be different. We only have to adapt ideas from [17] and [8] used in the monospecies case.

4.3.1. *Obtaining a kernel form.* Note that if $m_i = m_j$, (2) becomes

$$(29) \quad v' = v - (\omega \cdot (v - v_*)) \omega, \quad v_*' = v_* + (\omega \cdot (v - v_*)) \omega.$$

Symmetry properties allow us to write $[\mathcal{K}_3 g]_i$ in terms of v , v_* and v' , and not v_*' anymore. More precisely, we have the following lemma.

Lemma 4. *For any i , there exist nonnegative functions $(\tilde{\mathcal{B}}_{ij})_{1 \leq j \leq I}$ satisfying (5), such that*

$$(30) \quad [\mathcal{K}_3 g]_i(v) = \sum_{j \in \mathcal{M}_i} \iint_{\mathbb{R}^3 \times S^2} e^{-\frac{1}{4} m_i v^2 - \frac{1}{2} m_j v_*^2 + \frac{1}{4} m_i v'^2} g_j(v') \tilde{\mathcal{B}}_{ij}(\omega, v - v_*) d\omega dv_*, \quad \forall v \in \mathbb{R}^3.$$

Proof. The key idea of the proof lies in (29). Indeed, if we consider the relative velocity $V = v - v_*$, we can choose one unit vector $\omega^\perp \in \text{Span}(V, \omega)$ orthogonal to ω (the choice of either ω^\perp or $-\omega^\perp$ is relevant, but must be performed in a continuous way with respect to ω , and not randomly). Consequently, we can write

$$V = \omega(\omega \cdot V) + \omega^\perp(\omega^\perp \cdot V),$$

from which we immediately get

$$(31) \quad v - (\omega \cdot V)\omega = v_* + (\omega^\perp \cdot V)\omega^\perp, \quad v_* + (\omega \cdot V)\omega = v - (\omega^\perp \cdot V)\omega^\perp.$$

We can see that if we look for the post-collision relative velocity for the same pre-collisional V , but with respect to ω^\perp , we just exchange velocities v' and v_*' : for instance, the new v_*' , depending on ω^\perp , will be the old v' , depending on ω . Hence, it is clear that $\omega \mapsto \omega^\perp$ implies $v' \mapsto v_*'$ and $v_*' \mapsto v'$.

Consequently, if we replace ω by ω^\perp in the integral

$$\int_{\mathbb{R}^3 \times S^2} e^{-\frac{1}{2} m_i v_*^2} e^{\frac{1}{4} m_i (v_* + (\omega \cdot V)\omega)^2} g_j(v_* + (\omega \cdot V)\omega) \mathcal{B}_{ij}(\omega, V) d\omega dv_*,$$

it becomes

$$\iint_{\mathbb{R}^3 \times S^2} e^{-\frac{1}{2} m_i v_*^2} e^{\frac{1}{4} m_i (v_* + (\omega^\perp \cdot V)\omega^\perp)^2} g_j(v_* + (\omega^\perp \cdot V)\omega^\perp) \mathcal{B}_{ij}(\omega^\perp, V) d\omega^\perp dv_*.$$

Changing the variable ω into ω^\perp is obtained thanks to a rotation, so $d\omega^\perp = d\omega$. Hence, using (31), the previous integral becomes

$$\iint_{\mathbb{R}^3 \times S^2} e^{-\frac{1}{2}m_i v_*^2} e^{\frac{1}{4}m_i v'^2} g_j(v') \mathcal{B}_{ij}(\omega^\perp, V) d\omega^\perp dv_* = \iint_{\mathbb{R}^3 \times S^2} e^{-\frac{1}{2}m_i v_*^2} e^{\frac{1}{4}m_i v'^2} g_j(v') \mathcal{B}_{ij}(\omega^\perp, V) d\omega dv_*.$$

Let us set

$$\tilde{\mathcal{B}}_{ij}(\omega, V) = \left(\frac{m_i}{2\pi}\right)^{3/2} \begin{cases} n_i \mathcal{B}_{ij}(\omega^\perp, V) & \text{if } i \neq j, \\ \sum_{k \in \mathcal{M}_i} n_k \mathcal{B}_{ik}(\omega, V) + n_i \mathcal{B}_{ii}(\omega^\perp, V) & \text{if } i = j. \end{cases}$$

Assumption (5) on both $\mathcal{B}_{ij}(\omega, V)$ and $\mathcal{B}_{ij}(\omega^\perp, V)$ ensures that, for any i, j

$$(32) \quad \tilde{\mathcal{B}}_{ij}(\omega, V) \leq 2a \left(\frac{m_i}{2\pi}\right)^{3/2} \left(\max_{k \in \mathcal{M}_i} n_k\right) |\sin \theta| |\cos \theta| \left(|V| + |V|^{\delta-1}\right),$$

as well as (30). □

Lemma 4 allows to obtain the kernel form of \mathcal{K}_3 . More precisely, we have

Proposition 5. *Denote, for any i, j ,*

$$(33) \quad k_3^{ij}(\eta, v) = e^{-\frac{1}{8}m_i(\eta-v)^2 - \frac{1}{8}m_i \frac{(\eta^2 - v^2)^2}{(\eta-v)^2}} |\eta - v|^{-1} \varphi_3^{ij}(\eta - v), \quad \forall \eta, v \in \mathbb{R}^3,$$

where

$$(34) \quad \varphi_3^{ij}(p) = \frac{2}{|p|} \int_{\{p\}^\perp} e^{-\frac{1}{2}m_i(q+z_2)^2} \tilde{\mathcal{B}}_{ij}(p, q) |\sin(p, p+q)|^{-1} dq, \quad \forall p \in \mathbb{R}^3.$$

Then we have

$$[\mathcal{K}_3 g]_i(v) = \sum_{j \in \mathcal{M}_i} \int_{\mathbb{R}^3} g_j(\eta) k_3^{ij}(\eta, v) d\eta, \quad \forall v \in \mathbb{R}^3.$$

Proof. We perform the change of variable $v_* \mapsto V_* = v_* - v$ in (30), whose Jacobian equals 1, and get

$$(35) \quad [\mathcal{K}_3 g]_i(v) = \sum_{j \in \mathcal{M}_i} \int_{\mathbb{R}^3} \int_{S^2} e^{-\frac{1}{4}m_i v^2} e^{-\frac{1}{2}m_i(V_*+v)^2} e^{\frac{1}{4}m_i v'^2} g_j(v') \tilde{\mathcal{B}}_{ij}(\omega, V_*) d\omega dV_*.$$

Next, we consider the components of V_* respectively parallel and orthogonal to ω , i.e. we write $V_* = p + q$, where $p = \omega(\omega \cdot V_*)$, $q = V_* - \omega(\omega \cdot V_*)$. The component q which is orthogonal to ω belongs to the plane $\Pi = \{\omega\}^\perp = \{p\}^\perp$.

We want to perform the change of variables (see Figure 1)

$$(36) \quad (V_*, \omega) \mapsto (p, q), \quad \mathbb{R}^3 \times S^2 \rightarrow \mathbb{R}^3 \times \Pi.$$

Let us compute its Jacobian. For ω fixed, the replacement of V_* by p and q has unit Jacobian. Note that V_* and ω are independent of each other, which is not the case with p and q . Therefore, we need to pay attention on the integration order. Hence we first integrate with respect to q since $\Pi = \{p\}^\perp$. Then we combine the one-dimensional integration in the direction ω with the integral of ω over the unit sphere to give a three-dimensional integration over the three rectangular components of $|p|\omega$. We have to introduce the factor 2, since $p = \pm|p|\omega$. The Jacobian from p to $(|p|, \omega)$ (Cartesian to spherical coordinates) is $p^2 \sin(p, p+q)$. Consequently, we can write

$$dV_* d\omega = \frac{2}{p^2 \sin(p, p+q)} dp dq.$$

Eventually, it is clear that

$$v' = v - \omega(\omega \cdot (v - v_*)) = v + \omega(\omega \cdot V_*) = v + p.$$

4.3.2. *Properties of k_3^{ij} .* Let us first prove the following lemma, and then investigate some properties of k_3^{ij} .

Lemma 5. *The function $\varphi_3^{ij} : \mathbb{R}^3 \rightarrow \mathbb{R}$, defined by (34) for any i, j , belongs to $L^\infty(\mathbb{R}^3)$.*

Proof. Let $1 \leq i, j \leq I$, and choose $p \in \mathbb{R}^3$ and $q \in \{p\}^\perp$. From (32), we obtain

$$0 \leq \frac{\tilde{\mathcal{B}}_{ij}(p, q)}{|\sin(p, p+q)|} \leq 2a \left(\frac{m_i}{2\pi}\right)^{3/2} \left(\max_{k \in \mathcal{M}_i} n_k\right) |\cos(p, p+q)| \left(|p+q| + |p+q|^{\delta-1}\right).$$

Since $|\tan(p, p+q)| = |q|/|p|$, we can write

$$0 \leq \frac{\tilde{\mathcal{B}}_{ij}(p, q)}{|\sin(p, p+q)|} \leq C_i \left(\max_{k \in \mathcal{M}_i} n_k\right) \left(1 + \frac{q^2}{p^2}\right)^{-\frac{1}{2}} \left[(p^2 + q^2)^{\frac{1}{2}} + (p^2 + q^2)^{\frac{\delta-1}{2}}\right],$$

where $C_i = 2a(m_i/2\pi)^{3/2} > 0$. In what follows, C_i will denote any nonnegative constant only depending on m_i . This implies

$$0 \leq \frac{\tilde{\mathcal{B}}_{ij}(p, q)}{|p||\sin(p, p+q)|} \leq C_i \left(\max_{k \in \mathcal{M}_i} n_k\right) \left[1 + (p^2 + q^2)^{\frac{\delta-1}{2}}\right] \leq C_i \left(\max_{k \in \mathcal{M}_i} n_k\right) \left[1 + |q|^{\delta-2}\right],$$

using the fact that $\delta < 1$. Now, we split the range of integration in (34) into $|q| \leq 1$ and $|q| \geq 1$, and finally get

$$0 \leq \varphi_3^{ij}(p) \leq C_i \max_{k \in \mathcal{M}_i} n_k \left(\int_{|q| \leq 1} (1 + |q|^{\delta-2}) dq + \int_{|q| \geq 1} e^{-\frac{1}{2}(q+z_2)^2} dq \right) \leq C_i \max_{k \in \mathcal{M}_i} n_k.$$

This ends the proof of Lemma 5. □

Let us now investigate two properties of k_3^{ij} , which are related to Lemmas 1 and 2 for k_1^{ij} .

Lemma 6. *There exists $C > 0$ such that, for any i, j ,*

$$\int_{\mathbb{R}^3} k_3^{ij}(\eta, v) d\eta \leq \frac{C}{|v|}, \quad \forall v \in \mathbb{R}^3 \setminus \{0\}, \quad \int_{\mathbb{R}^3} k_3^{ij}(\eta, v) d\eta \leq C, \quad \forall v \in \mathbb{R}^3.$$

Proof. Let $1 \leq i, j \leq I$. We integrate (33) with respect to η and perform the change of variable $\eta \mapsto p = \eta - v$. Using Lemma 5, we get

$$\int_{\mathbb{R}^3} k_3^{ij}(\eta, v) d\eta \leq C \int_{\mathbb{R}^3} e^{-\frac{1}{8}m_i p^2 - \frac{1}{8}m_i \frac{(p^2 + 2p \cdot v)^2}{p^2}} \frac{1}{|p|} dp, \quad \forall v \in \mathbb{R}^3.$$

We split the right-hand integral into $I_1 + I_2$, where I_1 refers to $|p| \geq |v|$ and I_2 to $|p| \leq |v|$. On the one hand, we have

$$I_1 \leq \int_{|p| \geq |v|} e^{-\frac{1}{8}m_i p^2} \frac{1}{|p|} dp \leq C e^{-\frac{1}{8}m_i v^2}.$$

On the other hand, for the second integral, with spherical coordinates, we have

$$\begin{aligned} I_2 &= \int_0^{|v|} \int_0^{2\pi} \int_0^\pi e^{-\frac{1}{8}m_i r^2 - \frac{1}{8}m_i (r+2|v| \cos \psi_1)^2} r \sin \psi_1 d\psi_1 d\psi_2 dr \\ &= 2\pi \int_0^{|v|} r e^{-\frac{1}{8}m_i r^2} \int_0^\pi e^{-\frac{1}{8}m_i (r+2|v| \cos \psi_1)^2} \sin \psi_1 d\psi_1 dr, \end{aligned}$$

where ψ_1 corresponds to the angle between p and v . We then have to consider two situations for I_2 : $|v| \geq 1$ and $|v| \leq 1$. Simple changes of variables give, if $v \neq 0$,

$$\int_0^\pi e^{-\frac{1}{8}m_i (r+2|v| \cos \psi_1)^2} \sin \psi_1 d\psi_1 = \frac{1}{2|v|} \int_{r-2|v|}^{r+2|v|} e^{-\frac{1}{8}m_i s^2} ds.$$

If $|v| \geq 1$,

$$I_2 \leq \frac{C}{|v|} \int_0^{+\infty} r e^{-\frac{1}{8}m_i r^2} dr \int_{-\infty}^{+\infty} e^{-\frac{1}{8}m_i s^2} ds = \frac{C_i}{|v|} \leq C_i.$$

On the contrary, if $|v| \leq 1$,

$$I_2 \leq \frac{C_i}{|v|} \int_{r-2|v|}^{r+2|v|} ds = C_i, \quad \text{if } v \neq 0.$$

Consequently, we get the required estimates. \square

Using the same strategy as above, the following lemma can also be proved.

Lemma 7. *For any i, j , k_3^{ij} belongs to $L^2_{\text{loc}}(\mathbb{R}^3_v; L^2(\mathbb{R}^3_\eta))$.*

4.3.3. *Uniform decay.* Let us now prove the uniform decay property at infinity.

Proposition 6. *Let $g \in L^2(\mathbb{R}^3)^I$. For any $R > 0$ and any i , we have*

$$\|[\mathcal{K}_3 g]_i\|_{L^2(B(0,R)^c)} \leq \frac{C}{\sqrt{R}} \|g\|_{L^2(\mathbb{R}^3)},$$

where $C > 0$ is a constant.

Proof. Using the Cauchy-Schwarz inequality and Lemma 6, we can write

$$\begin{aligned} \|[\mathcal{K}_3 g]_i\|_{L^2(B(0,R)^c)}^2 &\leq C \int_{|v| \geq R} \sum_{j \in \mathcal{M}_i} \left[\int_{\mathbb{R}^3} g_j(\eta)^2 k_3^{ij}(\eta, v) d\eta \right] \left[\int_{\mathbb{R}^3} k_3^{ij}(\eta, v) d\eta \right] dv \\ &\leq \frac{C}{R} \sum_{j \in \mathcal{M}_i} \int_{\mathbb{R}^3} \left[\int_{\mathbb{R}^3} k_3^{ij}(v, \eta) dv \right] g_j(\eta)^2 d\eta = \frac{C}{R} \|g\|_{L^2(\mathbb{R}^3)}^2, \end{aligned}$$

where we also used the fact that $k_3^{ij}(\eta, v) = k_3^{ij}(v, \eta)$. This ends the proof. \square

4.3.4. *Equicontinuity.* This property is described in the following proposition.

Proposition 7. *For all $\varepsilon > 0$, there exists $\alpha > 0$ (not depending on g or i) such that*

$$(38) \quad \|[(\tau_w - \text{Id})\mathcal{K}_3 g]_i\|_{L^2(\mathbb{R}^3)} \leq \varepsilon \|g\|_{L^2(\mathbb{R}^3)}, \quad \forall w \in B(0, \alpha).$$

Remark. The previous result is not exactly like (20), but is enough to ensure the equicontinuity property.

Proof. Let $R > 0$. We obviously have, for any $w \in B(0, R)$,

$$(39) \quad \|[(\tau_w - \text{Id})\mathcal{K}_3 g]_i\|_{L^2(\mathbb{R}^3)}^2 \leq \int_{B(0,2R)} ([\mathcal{K}_3 g]_i(v+w) - [\mathcal{K}_3 g]_i(v))^2 dv + \int_{B(0,R)^c} [\mathcal{K}_3 g]_i(v)^2 dv.$$

Proposition 6 ensures that the second integral is upper-bounded by $\|g\|_{L^2}^2/R$. The first one can also be upper-estimated by

$$\sum_{j \in \mathcal{M}_i} \|g_j\|_{L^2}^2 \int_{B(0,2R)} \int_{\mathbb{R}^3} \left(k_3^{ij}(\eta, v+w) - k_3^{ij}(\eta, v) \right)^2 d\eta dv.$$

We can now choose $R \geq 2/\varepsilon^2 > 0$. Thanks to Lemma 7, there exists $\alpha > 0$ such that

$$|w| < \alpha \quad \Rightarrow \quad \int_{B(0,2R)} \int_{\mathbb{R}^3} \left(k_3^{ij}(\eta, v+w) - k_3^{ij}(\eta, v) \right)^2 d\eta dv \leq \frac{\varepsilon^2}{2}.$$

It is then immediate to recover (38) from (39). \square

4.4. **Compactness of \mathcal{K}_4 .** The proof of the compactness of \mathcal{K}_4 is very similar to the final part of the proof for \mathcal{K}_3 . The main difficulty is to obtain a kernel form of \mathcal{K}_4 . Once it is done, (19)–(20) can easily be proven as in section 4.3. Using the same change of variables as to obtain (37), we can write

$$[\mathcal{K}_4g]_i(v) = 2 \sum_{j \notin \mathcal{M}_i} n_j \left(\frac{m_j}{2\pi} \right)^{3/2} \int_{\mathbb{R}^3} \int_{\Pi} e^{-\frac{1}{4}m_i v^2 - \frac{1}{2}m_j(p+q+v)^2 + \frac{1}{4}m_i \left(v + 2\frac{m_j}{m_i+m_j}p \right)^2} g_i \left(v + 2\frac{m_j}{m_i+m_j}p \right) \mathcal{B}_{ij}(p, q) |p|^{-2} |\sin(p, p+q)|^{-1} dq dp.$$

The exponential term can be modified thanks to the following relation

$$-\frac{1}{4}m_i v^2 - \frac{1}{2}m_j(p+q+v)^2 + \frac{1}{4}m_i \left(v + 2\frac{m_j}{m_i+m_j}p \right)^2 = -\frac{m_j m_i^2}{2(m_i+m_j)^2} p^2 - \frac{m_j}{2} \left(q + v + \frac{m_j}{m_i+m_j}p \right)^2.$$

If we denote

$$z = v + \frac{m_j}{m_i+m_j}p,$$

and decompose it into the component z_1 parallel to ω and the component z_2 orthogonal to ω ($z_2 \in \Pi$), we obtain the new form of $[\mathcal{K}_4g]_i$:

$$[\mathcal{K}_4g]_i(v) = 2 \sum_{j \notin \mathcal{M}_i} n_j \left(\frac{m_j}{2\pi} \right)^{3/2} \int_{\mathbb{R}^3} e^{-\frac{m_j m_i^2}{2(m_i+m_j)^2} p^2 - \frac{1}{2}m_j z_1^2} g_i \left(v + 2\frac{m_j}{m_i+m_j}p \right) |p|^{-2} \int_{\Pi} e^{-\frac{1}{2}m_j(q+z_2)^2} \mathcal{B}_{ij}(p, q) |\sin(p, p+q)|^{-1} dq dp.$$

Next, we perform the change of variables

$$p \mapsto \eta = v + 2\frac{m_j}{m_i+m_j}p, \text{ whose Jacobian equals } \left(2\frac{m_j}{m_i+m_j} \right)^3,$$

and write z_1^2 in the following form

$$z_1^2 = \left(z \cdot \frac{\eta - v}{|\eta - v|} \right)^2 = \frac{1}{4} \frac{(\eta^2 - v^2)^2}{(\eta - v)^2}.$$

Thus $[\mathcal{K}_4g]_i$ becomes

$$[\mathcal{K}_4g]_i(v) = \frac{1}{4} \sum_{j \notin \mathcal{M}_i} n_j \left(\frac{m_j}{2\pi} \right)^{3/2} \left(\frac{m_i+m_j}{m_j} \right) \int_{\eta \in \mathbb{R}^3} e^{-\frac{1}{8}\frac{m_i^2}{m_j}(\eta-v)^2 - \frac{1}{8}m_j \frac{(\eta^2-v^2)^2}{(\eta-v)^2}} g_i(\eta) |\eta - v|^{-2} \int_{\Pi} e^{-\frac{1}{2}m_j(q+z_2)^2} \mathcal{B}_{ij} \left(\frac{m_i+m_j}{2m_j}(\eta-v), q \right) \left| \sin \left(\frac{m_i+m_j}{2m_j}(\eta-v), \frac{m_i+m_j}{2m_j}(\eta-v) + q \right) \right|^{-1} dq d\eta.$$

To write $[\mathcal{K}_4g]_i$ into the convenient kernel form, we introduce the function

$$\varphi_4^{ij}(p) = \left(\frac{m_j}{2\pi} \right)^{3/2} \frac{m_i+m_j}{2m_j|p|} \int_{q \in \Pi} e^{-m_j \frac{1}{2}(q+z_2)^2} \mathcal{B}_{ij}(p, q) |\sin(p, p+q)|^{-1} dq.$$

It is easy to prove, in the same way as in Lemma 5, that there exists $C > 0$ such that $\|\varphi_4^{ij}\|_{L^\infty(\mathbb{R}^3)} \leq C$ for any i, j . The i -th component of \mathcal{K}_4g can be written in the kernel form

$$[\mathcal{K}_4g]_i(v) = \sum_{j \notin \mathcal{M}_i} n_j \int_{\mathbb{R}^3} k_4^{ij}(\eta, v) g_i(\eta) d\eta,$$

where $k_4^{ij}(\eta, v)$ is given by

$$k_4^{ij}(\eta, v) = \frac{1}{4} e^{-\frac{1}{8} \frac{m_i^2}{m_j} (\eta-v)^2 - \frac{1}{8} m_j \frac{(\eta^2-v^2)^2}{(\eta-v)^2}} |\eta - v|^{-1} \varphi_4^{ij} \left(\frac{m_i + m_j}{2m_j} (\eta - v) \right).$$

The form of each k_4^{ij} is exactly the same as in (33). Consequently, k_4^{ij} inherits the same properties as k_3^{ij} (see Lemmas 5 and 6), which allows to obtain (19)–(20) as in Propositions 6 and 7, and the compactness of \mathcal{K}_4 .

Acknowledgment. The authors would like to thank Laurent Desvillettes for stimulating and helpful discussions.

REFERENCES

- [1] C. Bardos, F. Golse, and C. Levermore. Sur les limites asymptotiques de la théorie cinétique conduisant à la dynamique des fluides incompressibles. *C. R. Acad. Sci. Paris Sér. I Math.*, 309(11):727–732, 1989.
- [2] C. Bardos, F. Golse, and C. Levermore. Fluid dynamic limits of kinetic equations. I. Formal derivations. *J. Statist. Phys.*, 63(1-2):323–344, 1991.
- [3] C. Bardos, F. Golse, and C. Levermore. Fluid dynamic limits of kinetic equations. II. Convergence proofs for the Boltzmann equation. *Comm. Pure Appl. Math.*, 46(5):667–753, 1993.
- [4] S. Bastea, R. Esposito, J. L. Lebowitz, and R. Marra. Binary fluids with long range segregating interaction. I. Derivation of kinetic and hydrodynamic equations. *J. Statist. Phys.*, 101(5-6):1087–1136, 2000.
- [5] M. Bennoune, M. Lemou, and L. Mieussens. An asymptotic preserving scheme for the Kac model of the Boltzmann equation in the diffusion limit. *Contin. Mech. Thermodyn.*, 21(5):401–421, 2009.
- [6] L. Boudin, B. Grec, and F. Salvarani. The Maxwell-Stefan diffusion limit for a kinetic model of mixtures. Submitted.
- [7] J.-F. Bourgat, L. Desvillettes, P. Le Tallec, and B. Perthame. Microreversible collisions for polyatomic gases and Boltzmann’s theorem. *European J. Mech. B Fluids*, 13(2):237–254, 1994.
- [8] C. Cercignani. *The Boltzmann equation and its applications*, volume 67 of *Applied Mathematical Sciences*. Springer-Verlag, New York, 1988.
- [9] S. Chapman and T. G. Cowling. *The mathematical theory of nonuniform gases*. Cambridge Mathematical Library. Cambridge University Press, Cambridge, third edition, 1990. An account of the kinetic theory of viscosity, thermal conduction and diffusion in gases, In co-operation with D. Burnett, With a foreword by Carlo Cercignani.
- [10] P. Degond. Macroscopic limits of the Boltzmann equation: a review. In *Modeling and computational methods for kinetic equations*, Model. Simul. Sci. Eng. Technol., pages 3–57. Birkhäuser Boston, Boston, MA, 2004.
- [11] L. Desvillettes, R. Monaco, and F. Salvarani. A kinetic model allowing to obtain the energy law of polytropic gases in the presence of chemical reactions. *Eur. J. Mech. B Fluids*, 24(2):219–236, 2005.
- [12] C. Dogbe. Fluid dynamic limits for gas mixture. I. Formal derivations. *Math. Models Methods Appl. Sci.*, 18(9):1633–1672, 2008.
- [13] J. B. Duncan and H. L. Toor. An experimental study of three component gas diffusion. *AIChE Journal*, 8(1):38–41, 1962.
- [14] F. Golse and L. Saint-Raymond. The Navier-Stokes limit of the Boltzmann equation for bounded collision kernels. *Invent. Math.*, 155(1):81–161, 2004.
- [15] F. Golse and L. Saint-Raymond. The incompressible Navier-Stokes limit of the Boltzmann equation for hard cutoff potentials. *J. Math. Pures Appl. (9)*, 91(5):508–552, 2009.
- [16] H. Grad. Asymptotic theory of the Boltzmann equation. *Phys. Fluids*, 6:147–181, 1963.
- [17] H. Grad. Asymptotic theory of the Boltzmann equation. II. In *Rarefied Gas Dynamics (Proc. 3rd Internat. Sympos., Palais de l’UNESCO, Paris, 1962)*, Vol. I, pages 26–59. Academic Press, New York, 1963.
- [18] D. Hilbert. Mathematical problems. *Bull. Amer. Math. Soc.*, 8(10):437–479, 1902.
- [19] S. Jin and Q. Li. A BGK-penalization asymptotic-preserving scheme for the multispecies Boltzmann equation. Technical report, 2012. Submitted.
- [20] R. Krishna and J. A. Wesselingh. The Maxwell-Stefan approach to mass transfer. *Chem. Eng. Sci.*, 52(6):861–911, 1997.
- [21] P.-L. Lions and N. Masmoudi. From the Boltzmann equations to the equations of incompressible fluid mechanics. I, II. *Arch. Ration. Mech. Anal.*, 158(3):173–193, 195–211, 2001.
- [22] J. C. Maxwell. On the dynamical theory of gases. *Phil. Trans. R. Soc.*, 157:49–88, 1866.
- [23] J. Ross and P. Mazur. Some deductions from a formal statistical mechanical theory of chemical kinetics. *J. Chem. Phys.*, 35(1):19–28, 1961.
- [24] A. Rossani and G. Spiga. A note on the kinetic theory of chemically reacting gases. *Phys. A*, 272(3-4):563–573, 1999.
- [25] L. Sirovich. Kinetic modeling of gas mixtures. *Phys. Fluids*, 5:908–918, 1962.

[26] J. Stefan. Ueber das Gleichgewicht und die Bewegung insbesondere die Diffusion von Gasgemengen. *Akad. Wiss. Wien*, 63:63–124, 1871.

L.B.: UPMC UNIV PARIS 06, UMR 7598 LJLL, PARIS, F-75005, FRANCE & INRIA PARIS-ROCQUENCOURT, REO PROJECT TEAM, BP 105, F-78153 LE CHESNAY CEDEX, FRANCE
E-mail address: `laurent.boudin@upmc.fr`

B.G.: MAP5, CNRS UMR 8145, UNIVERSITÉ PARIS DESCARTES, SORBONNE PARIS CITÉ, 45 RUE DES SAINTS PÈRES, F-75006 PARIS, FRANCE & INRIA PARIS-ROCQUENCOURT, REO PROJECT TEAM, BP 105, F-78153 LE CHESNAY CEDEX, FRANCE
E-mail address: `berenice.grec@parisdescartes.fr`

M.P.: CMLA, ENS CACHAN, PRES UNIVERSUD PARIS, 61 AVENUE DU PRÉSIDENT WILSON, F-94235 CACHAN CEDEX, FRANCE & DEPARTMENT OF MATHEMATICS AND INFORMATICS, FACULTY OF SCIENCES, UNIVERSITY OF NOVI SAD, TRG DOSITEJA OBRADOVIĆA 4, 21000 NOVI SAD, SERBIA
E-mail address: `milana.pavic@cmla.ens-cachan.fr`

F.S.: DIPARTIMENTO DI MATEMATICA, UNIVERSITÀ DEGLI STUDI DI PAVIA, VIA FERRATA 1, I-27100 PAVIA, ITALY
E-mail address: `francesco.salvarani@unipv.it`