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1 Down-scaling of fracture energy during brittle creep 2 experiments

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3 **Abstract.** We present mode I brittle-creep fracture experiments along
4 fracture surfaces that contain strength heterogeneities. Our observations pro-
5 vide a link between smooth macroscopic time-dependent failure and inter-
6 mittent microscopic stress-dependent processes. We find the large scale re-
7 sponse of slow propagating sub-critical cracks to be well described by an Ar-
8 rhenius law that relates the fracture speed to the energy release rate. At the
9 microscopic scale, high resolution optical imaging of the transparent mate-
10 rial used (PMMA) allows detailed description of the fracture front. This re-
11 veals a local competition between subcritical and critical propagation (pseudo
12 stick-slip front advances) independently of loading rates. Moreover, we show
13 that the local geometry of the crack front is self-affine and the local crack
14 front velocity is power law distributed. We estimate the local fracture en-
15 ergy distribution by combining high resolution measurements of the crack
16 front geometry and an elastic line fracture model. We show that the aver-
17 age local fracture energy is significantly larger than the value derived from
18 a macroscopic energy balance. This suggests that homogenization of the frac-
19 ture energy is not straightforward and should be taken cautiously. Finally
20 we discuss the implications of our results in the context of fault mechanics.

1. Introduction

21 Heterogeneities in geological media have a strong impact on deformation processes in
22 Earth's crust that are difficult to assess. Indeed, deformation is often multi-scaled due to
23 numerous local rheological variations (e.g. variability of rocks, mineralogy) and geomet-
24 rical discontinuities (e.g. faults). These discontinuities are themselves often characterized
25 by complex morphologies that lead to local stress concentrations [*Okubo and Aki*, 1987;
26 *Schmittbuhl et al.*, 2006]. Micro-structures like grains, cracks and mineral assemblage also
27 contribute to the inhomogeneity of rock structures [*Scholz*, 2002; *Paterson and Wong*,
28 2005].

29 Measurements of large scale crustal deformation, usually done at the Earth surface e.g.
30 by GPS and InSAR [*Blewitt*, 2007; *Simons and Rosen*, 2007], allow limited constraints on
31 the deformation at small scales, in particular at depth. Measurements in laboratory rock
32 mechanics experiments suffer from the same limitations as sensors are located at the sur-
33 face of rock specimens. Macroscopic information averaged over the sample size might be
34 different from direct observations at the smaller scales at which the physical mechanism
35 responsible for the deformation is at play.

36 Creep of faults and rocks produces smooth and continuous deformation when recorded at
37 the system size, at large scales along faults [e.g. *Titus et al.*, 2006] and at small scales dur-
38 ing creep laboratory tests [*Heap et al.*, 2009]. However, the presence of microearthquakes
39 and transient deformation in creeping faults [*Lengliné and Marsan*, 2009] and acoustic
40 emission (AE) in laboratory experiments [*Lei et al.*, 2000; *Heap et al.*, 2009] suggest that
41 the deformation process is more intermittent and distributed over a wide range of veloci-

ties. The deformation induced by microearthquakes is generally not significant, although damage-related inelastic deformation contributes to the global deformation [Wang *et al.*, 2010].

The differences in behavior between macroscopic and microscopic deformations highlight the importance of analyzing the physical processes at local scales at which observations can unambiguously be linked to a given mechanism. Moreover these differences question homogenization procedures, i.e. the upscaling of physical parameters from microscopic to macroscopic scales. The focus of our study is the relation between macroscopic smooth creep deformation and microscopic complex intermittent dynamics during crack propagation experiments. We address the link between macroscopic and microscopic scales where local strength heterogeneities are important.

Creep deformation results from sub-critical crack growth mechanisms and plays a significant role in time-dependent failure of rocks [Atkinson, 1984]. Among the numerous models of creep, the brittle creep model results from time-dependent microfracturing at the crack tip [Scholz, 1968a]. The smooth large scale response in this model is a statistical response of a complex population of interacting micro-cracks [Scholz, 2002; Lockner, 1993; Baud and Meredith, 1997]. The time-dependence might result from chemical processes like corrosive mechanisms. In particular, rocks feature a strong action of water on Si-O bonds. For polymers or metals, lateral motion of atomic kinks along the crack front, microplasticity or diffusion induce strength variation in time [Lawn, 1975; Atkinson and Meredith, 1987a]. This time-dependence of slow kinetic crack propagation is usually well represented by a thermally activated mechanism where the strain rate obeys an Arrhenius law dependent on the energy release rate G (i.e. the energy available to drive the crack

65 per unit area) [Lawn, 1993]. In this approach, the crack growth is directly influenced
66 by environmental factors (applied stress, temperature, chemical concentrations) that af-
67 fect the free energy, and thus the energy barrier, via numerous competing mechanisms
68 like stress corrosion, diffusion, dissolution, plasticity and thermal effects [Atkinson and
69 Meredith, 1987a; Lawn, 1993]. Experimentally, several empirical relations have been of-
70 ten reported to divide the slow crack propagation in three main characteristic regimes
71 [Freiman, 1984]: 1) at very slow velocity, external variables are dominant and result in
72 an increasing speed of crack growth with increasing G ; 2) a weak stress-sensitive regime
73 follows where transport is limiting and 3) at higher stress crack growth kinetics is little
74 influenced by environment.

75 The energy release rate, G , is thus of primary importance when studying the creep mech-
76 anism and its scaling effect. G is often associated with a critical energy release rate, G_c ,
77 defined from the Griffith energy balance concept: the energy release rate equals the energy
78 available to drive the crack per unit area $G = G_c$. Such fracture energy criterion is relevant
79 to many geological processes like fault rupture or dike propagation [e.g. Rice and Simons,
80 1976; Rubin, 1993]. Estimation of G_c is difficult and relies on laboratory rock experiments
81 [Wong, 1982] or earthquake rupture modeling using recorded waveforms or geological ob-
82 servations. Most measurements of the fracture energy are associated with a macroscopic
83 quantity defined as a system-size representative value, hence ignoring all fluctuations of
84 G_c present at smaller scales. Here, we examine the relation between the macroscopic
85 quantity and the smaller scale variations in G_c . We do not investigate dynamic rupture
86 (e.g. earthquakes) but slow, creeping fractures. This slow regime is relevant to many
87 geophysical phenomena such as earthquake nucleation [Bouchon et al., 2011], slow slip

88 events and postseismic slip [e.g. *Das and Scholz*, 1981].

89 We present a series of quasi-static mode I fracture experiments performed on an exper-
90 imental setting introduced by *Schmittbuhl and Måløy* [1997]; *Delaplace et al.* [1999] and
91 *Grob et al.* [2009]. The system allows to track a propagating crack in a brittle hetero-
92 geneous medium. We use samples of polymethylmethacrylate (PMMA), which, contrary
93 to polycrystalline rocks, are transparent and allow for direct observation of crack defor-
94 mation at the scale of the heterogeneities. The high temporal and spatial resolution of
95 the experiment provides detailed information on the deformation process induced by the
96 crack at the microscopic scale.

97 We first present the macroscopic behavior of the fracture and an estimation of the rela-
98 tionship between fracture energy and average fracture velocity at the sample scale [*Scholz*,
99 1968a]. The brittle-creep propagation of the crack is characterized by an intermittent dy-
100 namics at the microscopic scale and is well modeled by an activation energy mechanism.
101 Then, we present the microscopic behavior of the fracture front roughness, the local frac-
102 ture velocity and the local fluctuations of the energy release rate along the crack front line.
103 The microscopic distribution of fracture energies spreads over a wide range and is directly
104 linked to the disordered morphology of the interface. The fracture energy determined
105 at the local scale is on average higher than the fracture energy computed at the global
106 scale. This highlights the importance of local defects in controlling the rupture process in
107 heterogeneous media. It also suggests that homogenization of the fracture energy is not
108 straightforward and should be taken cautiously.

2. Experimental Setup

2.1. Sample Preparation

109 We use two transparent PMMA plates of dimensions $20 \times 10 \times 1.0 \text{ cm}$ and $23 \times 2.8 \times 0.5 \text{ cm}$
110 (Figure 1). First, we sand-blast one surface of the narrow plate with glass beads of
111 diameter $\phi \in [180 - 300] \mu\text{m}$. We clean the blasted plate to remove any electrically
112 attached glass beads. Then we assemble the two plates in a stiff metallic loading cell with
113 the blasted surface facing a surface of the larger plate. Finally, we impose an homogeneous
114 normal load on the assembled plates and heat the loaded sample at 190°C for 45 minutes
115 to anneal the plates. The thermal annealing produces a weak cohesive interface along
116 which the sample breaks in mode I. The sand-blasting introduces the random roughness
117 of the plate surface that controls the local strength along the weak interface. It also
118 induces microstructures at the plate surface which make the sample opaque, but the
119 newly formed block, after annealing, recovers its transparency since the contrast of the
120 refraction index along the interface disappears.

121 We choose PMMA in our experiment because: 1) it is easy to handle and has a high
122 transparency adequate for optical imaging, 2) its low Young modulus (3.2 GPa) allows
123 fine tuning of small stress changes under imposed deformation and 3) its surface energy,
124 substantially larger than that of glass [*Brace and Walsh*, 1962; *Katsamanis and Delides*,
125 1988], limits the propagation of cracks at low applied stress. These combined properties
126 enable a better control of the experimental system. Another advantage of PMMA com-
127 pared to glass is its lower melting point which allows the welding of the PMMA plates at
128 a much lower temperature and accordingly in an oven with a better temperature moni-
129 toring. In addition to these technical motivations, PMMA exhibits viscoelastic behavior.

130 It is brittle at short time scales and semi-brittle or even plastic at longer times. Macro-
131 scopically this long time scale regime is described by a ductile rheology. PMMA shows
132 a time-temperature equivalence which is a very nice property for addressing either high
133 temperature processes or very long term evolution. This richness of the PMMA rheol-
134 ogy enables the observation of a mixed rupture regime macroscopically. It provides an
135 attractive analogy for the study of numerous time-dependent mechanisms in the Earth's
136 crust as those originating at the brittle-ductile transition. However, natural materials are
137 generally very disordered at all scales due to the presence of micro-cracks, grains, pores
138 and different minerals [Bean, 1996] that may increase the energy absorbed in the damage
139 zone. Our PMMA material does not exhibit such strong disorder, and consequently the
140 expected energy consumed around the crack tip by local plastic deformations is reduced.

2.2. Acquisition and Image Processing

141 Once the sample is ready, we clamp the widest PMMA plate to a stiff aluminum frame.
142 A stepping motor applies the loading over the top side of the narrow plate in a direction
143 normal to the plate interface (Figure 1). We measure the vertical displacement of the
144 loading point with a linear variable differential transformer (LVDT) and the loading force
145 by an STC 1205 traction/compression transducer. Displacement and force are measured
146 with a resolution of $1.3\ \mu\text{m}$ and $2.4 \cdot 10^{-3}\text{N}$ respectively. The vertical displacement
147 imposed on the narrower plate induces stable propagation of a mode I planar fracture
148 along the prescribed weak interface. We monitor the fracture front propagation using a fast
149 optical camera (CamRecord 600) with up to 1000 fps. Images have a maximum dimension
150 of 1024×1280 pixels and sample resolution is $\sim 20\ \mu\text{m}/\text{pixel}$. For some experiments, we

151 use a slow speed camera (Nikon D700) with up to 5 fps to follow the progression of the
152 average front position over long time scales (Figure 1).

153 Optical images of the interfacial mode I rupture show dark and bright regions respec-
154 tively corresponding to open crack and unbroken parts of the sample (Figure 2). Image
155 processing determines the transition between dark and bright areas that defines the frac-
156 ture front. We first compute the difference between each image and the first image of the
157 experiment. The image difference highlights the fracture front while removing permanent
158 artifacts. Then, grayscale images are transformed into black and white images accord-
159 ing to a gray level threshold separating bright and dark regions. Then, we compute the
160 gradient in the direction of front propagation to highlight the transition zone. We finally
161 extract connected pixels from the gradient images that correspond to the front position,
162 $a(x, t)$ (Figure 2). The front propagates along the y axis with the origin defined at the
163 load point and is positive in the sense of crack propagation. The x axis is perpendicular
164 to y and defines the coordinate of a point along the front (see *Måløy et al.* [2006] and
165 *Grob et al.* [2009] for details).

2.3. Loading Procedure

166 The loading rate is constant during an experiment but different among the experiments,
167 ranging from $6\mu\text{m}/\text{s}$ – $300\mu\text{m}/\text{s}$. The crack advance for each experiment is on the order of
168 1 cm. We imposed a small transient load in one of the experiments. In several experiments
169 we set the loading velocity to zero during the last stage to explore the fracture propagation
170 under a fixed load point displacement (*i.e.* relaxation test). Here the average crack
171 front velocity ranges from zero, at rupture initiation, to hundreds of $\mu\text{m}/\text{s}$. The rupture

172 velocity in all experiments is much lower than the Rayleigh wave speed of the PMMA
 173 ($V_r \sim 1.7 \cdot 10^9 \mu\text{m/s}$).

3. Macroscopic behavior

3.1. Energy Release Rate Evolution

The macroscopic mechanical energy release rate during mode I crack propagation \bar{G} is related to the total strain energy U stored in the system by [Lawn, 1993]

$$\bar{G}(t) = -\frac{1}{b} \frac{dU(t)}{d\bar{a}(t)}, \quad (1)$$

where b is the plate width (Figure 1) and $\bar{a}(t)$ is the average position of the front, i.e. the spatial average of each profile $a(x, t)$. Owing to the geometry of our setup, \bar{G} can be estimated by beam theory. In the approximation of small, purely elastic deformation and neglecting the weight of the plate, small compared to the loading force, the force F and load point displacement u (the deflection at the free end of the beam) are related by [Lawn, 1993]

$$F = \frac{E b h^3}{4 \bar{a}^3} u, \quad (2)$$

where E is the Young's modulus and h the height of the lower PMMA plate. We observe a linear relation between the observed average front position as a function of the predicted position (i.e. $[u/F]^{1/3}$) that proves the validity of the elastic beam theory applied to our experiments (Eq. 2). The energy stored in the plate is

$$U = \frac{uF}{2}. \quad (3)$$

Combining equations (1) to (3) we obtain

$$\bar{G} = \frac{3 E h^3 u^2}{8 \bar{a}^4}. \quad (4)$$

This is the usual expression of the energy release rate for a simple cantilever beam under imposed displacement and neglecting mode II loading [Lawn, 1993]. Equation (4) can be written in terms of the observables F , u and \bar{a} as

$$\bar{G} = \frac{3 F u}{2 b \bar{a}} \quad (5)$$

174 using equation (2). We prefer expression (5) instead of (4) because it involves measured
 175 variables raised to lower powers, which reduces the uncertainties. We track the evolution
 176 of \bar{G} (Eq. 5) during crack propagation as F , u and a are continuously monitored (Figure
 177 3).

178 The crack propagates at steady speed when the macroscopic energy release rate \bar{G}
 179 reaches a plateau that defines the macroscopic fracture energy \bar{G}_c (Figure 3), consistently
 180 with a macroscopic Griffith equilibrium. \bar{G}_c is in the range $122 J \cdot m^{-2} - 190 J \cdot m^{-2}$
 181 for the set of experiments presented here. The fracture front also propagates during the
 182 relaxation regime (final stage of the experiment in Figure 3) with decreasing velocity. This
 183 propagation happens while $\bar{G} < \bar{G}_c$, which is not predicted by the Griffith theory with
 184 time-independent \bar{G}_c .

185 We also estimate the global energy release rate directly from Eq. (1). The strain energy
 186 loss due to the crack propagation is the area, ΔA , defined in the deflection-force ($u - F$)
 187 space, between the loading and unloading curves. We numerically integrate ΔA and divide
 188 it by the crack length increase, $\Delta \bar{a}$, and the plate width, b . This approach provides an
 189 independent estimate \bar{G}_{area} and validates the beam approach employed above under the
 190 assumption of a constant energy release rate during the entire crack propagation regime.
 191 We report the energy release rates and fracture energies for a subset of experiments in
 192 Table (1). Although some discrepancies exist between the two methods, we observe that a

193 simple elastic beam model is appropriate to first order. Typical uncertainties in G are on
 194 the order of 7.5% given the uncertainties in F , u , b and a (typical values are $F = 30 \pm 0.2N$,
 195 $u = 10^{-2} \pm 10^{-6}m$, $b = 2.8 \pm 0.2cm$ and $\bar{a} = 10 \pm 0.2cm$). This relative uncertainty in
 196 \bar{G} partially accounts for the small differences observed between the two methods used
 197 for its calculation. However, discrepancies might also be due to viscoelastic effects and
 198 yielding at the crack tip which contribute to energy loss during crack propagation but are
 199 neglected in our calculation. These effects are unfortunately difficult to estimate and we
 200 did not attempt to quantify their importance. It is worth noting that fracture energies in
 201 PMMA are of the same order of magnitude as those for rocks in mode I experiments at
 202 room conditions [Atkinson and Meredith, 1987b] and significantly higher than those for
 203 minerals [Scholz, 2002].

3.2. Brittle Creep Crack Propagation Model

The crack propagation we observe at $\bar{G} < \bar{G}_c$ is not accounted for by the Griffith
 criterion (Figure 3). This is particularly pronounced towards the end of experiments
 where the loading displacement is fixed, which induces crack propagation at decreasing
 speed while \bar{G} also decreases. Many mechanisms produce time-dependent subcritical crack
 advance where the velocity, v , follows an Arrhenius law [Wiederhorn and Bolz, 1970]

$$v(G) = v_o \exp \left[\frac{\beta G - E'}{RT} \right], \quad (6)$$

204 where G is the energy release rate, E' is an energy barrier, R is the universal gas constant
 205 and T is the temperature. β and v_o are material-dependent empirical constants and the
 206 energy release rate G is directly linked to the stress intensity factor K by $K = \sqrt{G E}$ in
 207 plane stress.

We consider our fracture interface as a set of discrete sites of varying energy release rate along which the crack front line is pinned [Scholz, 1968b]. We describe the energy release rate by a probability density function, $f(G, \bar{G})$. It is the probability of a point along the front line to have an energy release rate G given an average energy release rate \bar{G} along the front. The probability distribution arises from the interplay between local toughness fluctuations on the interface and elastic interactions among sites of the front line. Equation (6) is understood as the local propagation criterion. It can be viewed as proportional to the inverse time to failure for a site locally loaded by a constant energy release rate G . The average crack front velocity, $d\bar{a}/dt$, is defined by

$$\frac{d\bar{a}}{dt} = \int_{G=-\infty}^{+\infty} f(G, \bar{G})v(G)dG. \quad (7)$$

208 We assume a normal distribution for $f(G, \bar{G})$ with mean \bar{G} and standard deviation σ_G
 209 [Scholz, 1968a],

$$\frac{d\bar{a}}{dt} = \int_{-\infty}^{+\infty} \frac{v_0}{\sqrt{2\pi}\sigma_G} \exp\left[-\frac{1}{2}\left(\frac{G - \bar{G}}{\sigma_G}\right)^2 + \frac{\beta G - E'}{RT}\right] dG \quad (8)$$

which integrates to

$$\frac{d\bar{a}}{dt} = v_o \exp\left[\frac{\beta\bar{G} - E'}{RT} + \frac{\beta^2\sigma_G^2}{2R^2T^2}\right] = \bar{v}_o \exp\left[\frac{\beta}{RT}\bar{G}\right] \quad (9)$$

where $\bar{v}_o = v_o \exp\left[\frac{\beta^2\sigma_G^2}{2R^2T^2} - \frac{E'}{RT}\right]$. Equation (9) plainly shows that the evolution of the crack front velocity at the macroscopic scale is described by the same functional form as in the microscopic scale (equation 6). Equation (9) implies $\bar{v} \propto \exp(\bar{G})$ independently of the loading conditions, extending its validity to the entire duration of the experiment. Equation (9) also reproduces the average evolution of the crack speed as a function of the macroscopic energy release rate \bar{G} (Figure 4). It is also applicable when $du/dt = 0$.

Since $\bar{a}(t)$ is small in the creeping regime, we use a first order expansion around the initial position \bar{a}_0 :

$$\bar{G}(\bar{a}) = \bar{G}_0 + \frac{d\bar{G}(\bar{a}_0)}{d\bar{a}}(\bar{a} - \bar{a}_0) \quad (10)$$

where \bar{a}_0 is the average position of the front at the onset of the creeping regime and $\bar{G}_0 = \bar{G}(\bar{a}_0)$. Substituting into (9), we get

$$\frac{d\bar{a}}{dt} = \bar{v}'_o \exp[-C_0\bar{a}] \quad (11)$$

where $\bar{v}'_o = \bar{v}_o \cdot \exp(5\beta\bar{G}_0/RT)$ and $C_0 = 4\beta\bar{G}_0/RT\bar{a}_0$. After integration we get

$$\bar{a}(t) - \bar{a}_0 = \frac{1}{C_0} \ln [C_0\bar{v}'_o(t - t_0) + \exp(C_0)] \quad (12)$$

210 where t_0 is the time at the onset of the creeping regime.

211

We also investigate the prediction of the model in the case $du/dt \neq 0$. When the front reaches a steady-state regime, the crack propagates with $\bar{G} = \text{constant}$. This condition is similar to the Griffith energy criterion. A propagation with a constant energy release rate is described from equation (4) by

$$\bar{G} = \frac{u^2 3Eh^3}{\bar{a}^4 8} \quad \text{and} \quad a(u) = \sqrt{u} \left(\frac{3Eh^3}{8\bar{G}} \right)^{1/4}. \quad (13)$$

212 Eq. (12) provides a good description of the macroscopic evolution of the crack front in
 213 the relaxation regime (Figure 5, bottom) and Eq. (13) a good description in the forced
 214 regime (Figure 5, top). The macroscopic evolution of the front advance is well reproduced
 215 by a subcritical crack growth mechanism (Figures 4 and 5) that explains the propagation
 216 of the crack even for an energy release rate (or similarly the stress intensity factor) lower
 217 than the interface fracture energy (or similarly the toughness). It is also valid even when

218 the front is propagating at a constant speed. In order to supplement our estimation of
219 the fracture energy, next we present a procedure to estimate G at the local scale.

4. Microscopic Behavior

4.1. Scaling of the fracture front

220 So far, we only considered straight crack fronts. These are expected if the toughness is
221 uniform along the interface. However, departures from the straight front geometry suggest
222 lateral variation of material resistance due to the sand-blasting procedure applied during
223 the sample preparation (Figure 6). The morphology of propagating interfacial cracks has
224 been intensely studied [*Schmittbuhl and Måløy, 1997; Delaplace et al., 1999*] and repro-
225 duced numerically [*Schmittbuhl et al., 2003a*]. The in-plane morphology of crack fronts
226 is self-affine and can be characterized by a roughness exponent. We show that a similar
227 scaling of the fracture front morphology exists in our experiments. The Fourier transform
228 of each front line, for a given experiment, is computed after detrending and tapering the
229 original profile $a(x, t)$. We then obtain the average power spectrum representative of all
230 fronts in this experiment (Figure 6). We observe a power law decay of the power spectrum
231 in agreement with a self-affine morphology and the roughness exponent is in agreement
232 with previously reported values ($H = 0.6$) [*Santucci et al., 2010*]. Variable geometry of
233 the crack front line suggests that heterogeneities over the interface should be taken into
234 account when evaluating the fracture energy.

4.2. Distribution of local velocities

235 The waiting time spent by the crack front at each local site is power law distributed
236 [*Måløy et al., 2006*]. The intermittent dynamics at the microscopic scale is also indirectly

237 evidenced in rocks during creep tests [e.g. *Heap et al.*, 2009]. The recording of acoustic
 238 emissions during a slow macroscopic deformation suggests that at the local scale the defor-
 239 mation is distributed over a wide range of velocities. Figure 7 shows, for two experiments,
 240 the probability density function (pdf) of v obtained by the procedure described by *Måløy*
 241 *et al.* [2006]. The broad distribution of local speeds, spanning more than two orders of
 242 magnitude, highlights the irregular nature of the fracturing process at the small scale.

4.3. Local Energy Release Rate

243 4.3.1. Non-straight Front Geometry

We follow a first-order analysis based on a perturbation method for nearly straight crack fronts, valid for small relative variations of the fracture energy [*Gao and Rice*, 1989]. This approach is valid in our experiments since front lines are straight at first order. Furthermore, the model adequately reproduces the morphology of experimentally propagating cracks, although a discrepancy in the value of the roughness exponent remains [*Schmittbuhl et al.*, 2003a]. However, this approach ignores crack coalescence, which can play a role in redistributing stresses ahead of the crack front [*Hansen and Schmittbuhl*, 2003; *Schmittbuhl et al.*, 2003b]. Here we assume that the elastic front line model represents the dominant process. The first order approximation of the energy release rate of an almost straight crack is [*Gao and Rice*, 1989]

$$G(x_i, t_k) = \bar{G}(t_k) (1 + \gamma(x_i, t_k)) \quad (14)$$

for a given front at time t_k . $\bar{G}(t_k)$ is the energy release rate computed from the average front position at time t_k and $\gamma(x_i, t_k)$ is the contribution from elastic interactions due to

local fluctuations of the front position given by:

$$\gamma(x, t) = \frac{1}{\pi} \mathcal{P}\mathcal{V} \int_{-\infty}^{+\infty} \frac{a(x', t) - a(x, t)}{(x - x')^2} dx' \quad (15)$$

where $\mathcal{P}\mathcal{V}$ denotes the principal value. This expression is the Hilbert transform of the local slope of the front [Ampuero and Rubin, 2008], which explains why small scale variations in the $\gamma(x)$ profile are larger than in the $a(x)$ profile. The discretized version of Eq. (15) is

$$\gamma(x_i, t_k) = \frac{1}{\pi} \sum_{j=-l/2, j \neq i}^{l/2} \frac{a(x_j, t_k) - a(x_i, t_k)}{(x_j - x_i)^2} . dl \quad (16)$$

244 where dl is the width of a pixel and l is the width of the picture. A small error may arise
 245 in the determination of γ for points close to the edge of the images of the fracture front.
 246 Since the decay of the elastic kernel with distance is fast ($\propto 1/x^2$), the loss of information
 247 at the edge only affects a small region near the edge of the picture. The computation of γ ,
 248 as presented in Eq. (16), assumes a half plane crack in an infinite body. If the finite height
 249 of the plate is taken into account, the elastic kernel transitions from the $1/x^2$ behavior
 250 of Eq. (15) to a faster (exponential) decay at distances x larger than the plate height.
 251 We calculate the local value of γ at the middle of a front line in order to quantify the
 252 impact of the truncation and the finite height. We progressively extend the bounds of the
 253 summation up to $l/2$. Two different front lines originating from two different experiments
 254 are processed in this way to represent the typical convergence of γ as a function of the
 255 integration range. We observe that $\sim 90\%$ of the final estimate of γ is recovered in the
 256 first 100 pixels which corresponds to a dimension of 2 mm (figure 8). This shows that γ
 257 is controlled by the closest neighboring sites. The infinite medium approach is valid since

258 the length-scale influencing the estimate of γ is smaller than all other dimensions of the
 259 plate. An example of a γ profile for a given front is shown in Figure 6.

260 4.3.2. 1D Time Evolution

261 The matrix $\gamma(x, t)$ represents normalized fluctuations of the energy release rate at any
 262 time and position along the front. The mean of the γ distribution, along a given front line,
 263 for a given image, is zero by definition. We track the evolution of $\gamma(x_0, t)$ for a particular
 264 position, x_0 , along the front (Figure 9). The local movement of the front $a(x_0, t)$ reveals
 265 potential correlations between the evolution of γ and local variations of a (Figure 9).

266 The local motion of the front is not continuous but rather exhibits a stick-slip pattern
 267 with long resting periods followed by jumps of the front position (Figure 7 and 9). In
 268 contrast, the large scale evolution of the average crack position is continuous.

269 4.3.3. 2D Maps

The energy release rate increases at pinned positions along the front line. This increase
 has two contributions : 1) a large-scale contribution from the external loading applied
 to the system (increase of the load point deflection) and 2) local contributions resulting
 from local elastic interaction due to the differential movement of neighboring points. As
 G increases, it rises the probability of the local site to fail in the next time interval as
 defined by Eq. (6). This probability increases until G reaches G_c , at which point the local
 site breaks. Accordingly, we define γ_c as the last value of γ just before a local variation
 of the front position and we estimate the fracture energy at this site as

$$G_c(x_i, a(x_i, t_k^*)) = \bar{G}(t_k^*) (1 + \gamma(x_i, a(x_i, t_k^*))) \quad (17)$$

270 where t_k^* is the time at which γ_c was reached. It is important to note that $\bar{G}(t_k^*)$ is ob-
 271 tained from the measurement of the average position of the front $\bar{a}(t_k^*)$. Any variation of

272 γ taking place during the time interval between estimating γ_c and the local front advance
 273 is neglected. These microscopic estimates of the fracture energy are lower bounds, since
 274 there is a finite probability that the local site actually breaks with $G < G_c$ according to
 275 the Arrhenius model described earlier. In order to get an insight on collective processes
 276 that take place at the local scale during the crack propagation, we show two examples of
 277 a spatio-temporal map $\gamma(x, t)$ in Figure 10. These maps are subregions of the effectively
 278 recorded areas (both in time and space) taken from two different experiments represen-
 279 tative of the two propagation regimes, at fixed loading point (relaxation regime) and
 280 constant loading rate respectively .

281 Notable features emerging from Figure 10 are summarized below:

282 1. At a given time (horizontal line), only a few points are at $G = G_c$, and most points
 283 have $G < G_c$ for both forced ($\bar{G} = \bar{G}_c$) and relaxation regimes. Observation of the
 284 intermittent failure of local sites during slow crack propagation is a direct measurement
 285 of a brittle/creep behavior. As most of the sites along the front line are below G_c , the
 286 macroscopic estimate of fracture energy \bar{G}_c is an average among local values of G that
 287 are mainly below the critical state ($\bar{G} \neq \bar{G}_c$). Therefore, the macroscopic fracture energy
 288 is a lower estimate compared to the microscopic values.

289 2. Large collective failure events (numerous black points on the same horizontal line)
 290 corresponding to a burst of activity in a relatively short time interval occur at discrete
 291 periods. These fast local movements arise from a connected set of unstable positions along
 292 the front. A burst of unstable positions appears when the local G is equal or higher than
 293 the fracture energy of all sites of the cluster. This scenario typically occurs after breaking

294 a strong local asperity (high G_c value). The front then travels in a region with lower G_c
295 up to a new arrest position that corresponds to the next accessible equilibrium position.

296 3. Despite very different average velocities and loading regimes, strong similarities exist
297 between the patterns of the two maps. The main difference between the two loading
298 regimes arises from the waiting time between each successive move of the front at a
299 particular local site. Time intervals between two successive moves are much larger in the
300 relaxation regime (frame rates in Figure 10 are 125 fps and 1000 fps for the relaxation
301 and the forced regime respectively). The shorter time intervals observed in the forced
302 regime are related to a higher loading rate at each local site.

4.4. Influence of Fluctuations in Toughness

303 4.4.1. Large scale fluctuations

The front geometry has a long wavelength curvature due to a large scale variation of the fracture energy at the interface. The annealing condition in the sample during its preparation is not perfectly homogeneous. Transient heat diffusion from the sides of the plates to the center of the sample is possibly responsible for such curvature. As a consequence, edges of the plate have higher strength than the center. We used this large scale curvature of the front line to infer the large scale variation of the fracture energy along the x direction. The largest Fourier mode of the front is a single cosine function of the form $a(x) = A_0 + 2A_1 \cos(\pi x/b)$ where x is the coordinate along the front direction and b is the width of the plate [Gao and Rice, 1989]. Griffith criterion holds everywhere along the crack front (G very close to G_c), with G_c being the local fracture energy at each

site. It follows from *Gao and Rice* [1989]

$$\sqrt{\tilde{G}} = \sqrt{\tilde{G}_0} \left[\frac{1}{A_0} - \frac{\pi}{L} \right] A_1, \quad (18)$$

where \tilde{G} and \tilde{G}_0 are the Fourier modes of order 1 and 0 respectively. The variation of fracture energy at long wavelength between the center of the plate and its border is

$$\frac{G_c^{center}}{G_c^{border}} \approx \left(\frac{b - \pi A_1}{b + \pi A_1} \right)^2 \quad (19)$$

304 for $A_0 \gg b/\pi$ which is the case for our experiments. As A_1 is on the order of several
 305 hundred of micrometers and b is $\sim 2cm$, the long wavelength shape of the front is explained
 306 by a variation of fracture energy along the plate on the order of 10%.

307 4.4.2. 2D Toughness fluctuations

308 We estimate spatial maps of local fracture energies $G_c(x, a(x))$, i.e. everywhere along the
 309 interface, from Eq. 17 (Figure 11). The local fracture energy is broadly distributed over
 310 the range 18 to 395 $J.m^{-2}$, clearly expressing the heterogeneity of local fracture energies
 311 (Figure 11). Interface roughness, δz , is obtained from a broken sample, previously peeled
 312 off (Figure 11). The analyzed interface does not correspond to the same area where the G_c
 313 map was computed but the statistical properties of the interface morphology are assumed
 314 to be representative of the whole sample interface. The morphology of the interface was
 315 measured by a laser profilometer with spatial and vertical resolution of 10 μm and $\sim 1\mu m$
 316 respectively (Figure 11). Height fluctuations δz 's have a standard deviation of several of
 317 μm .

Local toughness variations arise from the squeezing of local asperities along the in-
 terface. We extract all profiles of G_c and δz along the front propagation direction and
 perpendicular to it. The power spectra, $S(k)$, is averaged over the two in-plane directions

for both the distributions of G_c and δ_z . The autocorrelation function, $C(\Delta r)$, is then obtained as the inverse Fourier transform of the power spectrum,

$$C(\Delta r) = \int_{-\infty}^{\infty} S(k) e^{j2\pi k\tau} dk, \quad (20)$$

where k is the wavenumber. Figure 12 shows correlation functions of the fracture energy and the surface morphology along with their power spectra $S(k)$. The auto-correlation functions of G_c and δz decrease with distance (Figure 12). The decrease is faster for G_c than for δz . The power spectra shows an exponential form although we cannot exclude a model with power-law decay associated with a cut-off length. An exponential of the form $S(k) \propto \exp(-k/k_c)$ is an adequate fit. The length-scale $1/k_c$ is the correlation length, the distance beyond which the distribution becomes uncorrelated. For the front morphology we find $1/k_c \sim 188 \mu m$. The distribution of G_c is characterized by a smaller correlation length $\sim 84 \mu m \sim 4$ pixels. These cutoff length-scales can also be appraised from the correlation functions as we observe that $C(\Delta r)$ gets around 0 above these distances. Although the cut-off length is not the same, the exponential decay is observed for both G_c and δ_z . The existence of a finite correlation length for G_c and δz does not necessarily imply that an approach using a representative element volume (REV) could be valid at scales greater than this length scale. Long range interactions, owing to elastic forces applied over long distances results in an interplay between the toughness fluctuations and these elastic forces. This is well demonstrated by the self-affine nature of the crack front line observed previously (section 4.1).

5. Scaling of the Fracture Energy

335 The observation that most points along the crack front have energy release rate $G < G_c$
 336 indicates that the macroscopic estimate of the fracture energy, \bar{G}_c , is different from the
 337 fracture energy at the local scale. We test this proposition by comparing G_c at the local
 338 and the global scales. In particular, we address whether the macroscopic toughness is a
 339 simple arithmetic average of the distribution of local toughness. In such a case, an ho-
 340 mogenization procedure is quite straightforward and $\bar{G}_c = \langle G_c \rangle$, i.e. the local variations
 341 of material resistance are simply scattered around the global value and the average of the
 342 distribution is equal to the global value. We calculate normalized distributions of G_c as
 343 well as global estimations of \bar{G}_c , obtained with Eq. (5) for each experiment (Figure 13).
 344 We only consider for the estimates of G_c those locations where the front was present
 345 during more than one time interval. The estimate involves only values at equilibrium
 346 positions and thus related to the interface property [Roux *et al.*, 2003]. We also calculate
 347 the normalized distribution of G computed from Eq. (14) (Figure 13). The distribution
 348 of local energy release rate is wide and is centered on the macroscopic estimate of the
 349 fracture energy. This is expected from Eq. (14) and Eq. (16), because γ has zero mean.
 350 Accordingly, the distribution of G should be approximately centered on the global value
 351 \bar{G} , as observed. A more interesting feature emerges when analyzing the distribution of
 352 local fracture energy, G_c (Figure 13). We clearly observe that the average of local fracture
 353 energy, $\langle G_c \rangle$, is higher than the macroscopic estimate of the fracture energy \bar{G}_c . We
 354 recall that the macroscopic estimate of the fracture energy \bar{G}_c is defined as the maximum
 355 value of the \bar{G} reached during the crack propagation. Such higher values of G_c compared
 356 to \bar{G}_c are observed for all experiments and the distribution of G_c is systematically shifted

357 to higher values compared to the global \bar{G}_c . This indicates that G_c samples the subset
 358 with the highest values of G . For a pinned position of the front, γ increases with time
 359 leading to higher values at the end of the time interval, just before the local advance of
 360 the front. Therefore, γ_c are on average higher than γ and the distribution of G_c is shifted
 361 towards higher values compared to G . It is worth noting that the shift to higher values,
 362 going from macro- to micro-scale, does not depend of any particular choice of calculating
 363 \bar{G} . For example, Eq. (4) and (5) produce a similar shift of $\langle G_c \rangle$ compared to \bar{G}_c because
 364 G_c values are computed using \bar{G} from Eq. (14). The observed toughening effect at the
 365 microscopic scale is thus independent on a particular calculation of \bar{G} . The same argu-
 366 ment is also valid when considering the variation of \bar{G} owing to the uncertainties of the
 367 parameters defining \bar{G} . Despite these uncertainties, local values of G_c will still be, on
 368 average, higher than \bar{G}_c . We note however that the amplitude of this toughening depends
 369 on the macroscopic value of the energy release rate.

370 There are some known geophysical scaling problems where a simple arithmetic mean (av-
 371 erage) is not relevant for an upscaling of the problem. For example, bulk permeability
 372 of a heterogeneous material is bounded from above by the arithmetic mean and below
 373 by the harmonic average [*Zimmermann and Main, 2004*]. Contrary to the case of the
 374 permeability, for which the Darcy's law allows to make the transition link between scales,
 375 an equivalent expression does not exist in our problem. This prevents us to define the
 376 particular averaging procedure relevant in our system. Nevertheless, we test several forms
 377 of mean of our values of G_c . We compute the harmonic, the geometric and the arith-
 378 metic means of the microscopic distribution of fracture energy for each experiment. We
 379 observe for all experiments that all types of average are systematically overestimating the

380 macroscopic value. This suggests that scaling effects are non trivial and related to the
381 observation scale and the mechanism of brittle creep as proposed.

6. Implications for Fault Mechanics

6.1. Link Between Crack Modes

382 Several analogies can be drawn between our mode I fracture experiment and deformation
383 processes in the crust. Tensile failure of natural materials occurs during dyke propagation
384 and during secondary micro-cracking surrounding shear fractures. The mechanical process
385 in our experiments is not limited to tensile fracture. The role of heterogeneities along a
386 fault plane has similar consequences on the local values of the stress intensity factor (or
387 energy release rate). An expression similar to Eq. (14) exists for mode II and mode
388 III cracks when ignoring second order coupling among shear modes [*Gao et al.*, 1991].
389 The only difference with the mode I expression is a constant coefficient dependent on the
390 Poisson ratio. One could refer to *Schmittbuhl et al.* [2003a] for an extended discussion on
391 the relation between rupture modes.

6.2. Importance of the Process Zone

392 Fracture in rock samples or faults exhibit a process zone that encompasses a region with
393 complex microcracking because of the high stress field around the crack tip. Significant
394 energy is dissipated in the formation of the process zone and in the frictional work on
395 these microcracks. This energy loss is much higher than the surface energy associated with
396 the separation of the two blocks in contact. As the size of the process zone is observed
397 to increase with fault length, a scaling argument suggests that the energy release rate
398 dissipated in the process zone also increases with the fault/system size [*Scholz*, 2002].

399 This results in an apparent scaling of the fracture energy with the system size, where
400 the fracture energy includes all sources of energy dissipations [e.g. *Abercrombie and Rice,*
401 2005]. Our system does not involve a pronounced process zone because of the small bulk
402 disorder. As a consequence, estimates of the fracture energy in our experiments are close
403 to the surface energy, in contrast to earthquake data where the growth of the process zone
404 and friction greatly influence the dissipation process. Our observations do not directly
405 address the issue of the scaling of G_c with system size. The size of our system is fixed by
406 the sample size (plate dimensions) and does not vary. The difference of fracture energy
407 that we observe results from the change of the observation scale and not the sample scale.

6.3. Slow Slip Events

408 Our observations provide a guide for understanding the simultaneous presence of acous-
409 tic emissions and smooth deformation during creep experiments in rock samples [e.g. *Heap*
410 *et al.*, 2009]. Such brittle-creep behavior is possible due to the heterogeneous nature of the
411 interface which produces, at the local scale, a complex pattern of ruptures (see Figure 10)
412 because of elastic interactions. It is a competition between sites with sub-critical or crit-
413 ical propagation (pseudo stick-slip crack advances) resulting in an average (macroscopic)
414 steady deformation. This mode of rupture implies a variation of the fracture energy with
415 observation scale. The fracture energy computed at the sample scale actually integrates
416 all points along the front line and thus is an average among sites that are at $G = G_c$ (only
417 a few points) and a majority of points that are pinned along the interface under unloading
418 conditions with $G < G_c$. A higher proportion of sites along the front line are at $G = G_c$
419 when the loading speed is increased, and subsequently the crack speed. The difference
420 between \bar{G}_c and $\langle G_c \rangle$ vanishes when the speed of the crack approaches the Rayleigh wave

421 speed. The deformation rate is an important factor that affects the large scale estimate
422 \bar{G}_c while faintly modifying estimates at local scales. In the case of a dynamic crack, in our
423 model, at a given time step, all points along the front line are moving. Our definition of
424 the local fracture energy implies that at all time steps and for every position along front,
425 the energy release rate is equal to the fracture energy. Averaging over all local estimates
426 produces the exact same value as the macroscopic one, $\bar{G}_c = \langle G_c \rangle$. Our results are more
427 closely related to slow deformation processes occurring in the Earth crust rather than dy-
428 namic events, like creeping of fault segments, nucleation phase of earthquakes, postseismic
429 slip or slow slip processes. The fracture energy for slow ruptures in the earth are lower
430 (by few orders of magnitude) compared to values reported for dynamic fractures of the
431 same size. This is the case for a creep event along the San Andreas fault reported by [*Rice*
432 *and Simons*, 1976]. This is also suggested from the low slip to size ratio of slow events
433 [e.g. *Ide et al.*, 2007], which combined with Eq. (67) from *Rice and Simons* [1976] yields
434 low values of fracture energy compared to regular earthquake. These studies however do
435 not resolve the spatial distribution of fracture energy on the fault plane during the slow
436 events which would make possible the comparison with macroscopic estimates as in our
437 study.

7. Conclusion

438 We provide a direct description of brittle creep crack propagation at different observa-
439 tional scales. At the local scale, the fracturing process is intermittent and is characterized
440 by a complex fracture front morphology and a wide range of local crack velocities. At
441 a given time, a mixture of slow and fast ruptures can be observed simultaneously. This
442 complex pattern results from the presence of heterogeneities along the crack interface to-

443 gether with significant elastic interactions. When observed at the macroscopic scale, the
444 variability of the local scales disappears and the crack propagation is smooth. This sug-
445 gests that macroscopic observation of slow deformation processes (for example for rocks
446 samples or faults) should be analyzed in the light of the heterogeneous nature of the mate-
447 rial. Therefore, the fracture energy depends on the observation scale and its macroscopic
448 estimate is lower than the average over local values.

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Figure 1. Side view (left) and top view (right) of the experimental setup. A stiff aluminum frame is attached to the upper PMMA plate. The bottom plate is separated from the upper one using a loading force F applied by a rod connected to a stepping motor. The load causes a deflection u of the bottom plate and the propagation of a mode I crack. The crack front is located at distance \bar{a} from the free end. The front advance is monitored by a high speed camera set in vertical position, perpendicular to the crack plane. Lower plate height and width are noted h and b respectively.

Figure 2. Zoom on the crack front displaying the difference between an image at the onset of the experiment and the current image of the front. During the experiment the crack front propagates from the top to bottom. The bright area marks the zone which is already cracked while in the dark area the two plates are still in contact (i.e. uncracked). The front line, $a(x, t)$, is superimposed on the image as a red line. Scales are provided by two black segments which are 1 mm and 0.5 mm long respectively in the x and y directions.

Figure 3. Evolution of \bar{G} during a mode I fracture experiment as a function of the crack front position. During this experiment, the sample is loaded at a constant speed. The front starts to move slowly and \bar{G} increases (light gray area) up to a nearly constant value when the crack front velocity reaches a steady-state value (gray area). The loading point position is finally set to a constant value and \bar{G} decreases (dark gray area) before the loading point moves back to its initial position at the end of the experiment (in a part of the curve not displayed here).

Table 1. Macroscopic energy release rate determined for all the seven experiments performed with the fast speed camera. The values of the fracture energies, \bar{G}_c , are obtained by taking the maximum of \bar{G} computed from Eq. (5) for each experiment. We also report the mean energy release rates during the crack propagation obtained by the beam theory, $\langle \bar{G} \rangle$ and by integration under the deflection-force space, \bar{G}_{area} . All values are in $J.m^{-2}$.

Experiment n°	\bar{G}_c	$\langle \bar{G} \rangle$	\bar{G}_{area}
1	174	148	144
2	122	118	126
3	153	140	123
4	147	142	113
5	182	173	167
6	130	122	115
7	190	182	159

Figure 4. Variation of the average crack velocity $d\bar{a}/dt$ as a function of the average energy release rate, \bar{G} for an experiment conducted with the high speed camera. The gray scale refers to the time since the start of the recording. During the entire crack propagation, from the initiation phase up to the relaxation phase, the best fit using Eq. (6), displayed as a dashed line, provides a good description of the data. It suggests that the crack propagation is a process driven by time-dependent brittle creep.

Figure 5. Top: Variation of the average crack front position, \bar{a} , as a function of the loading point displacement, u . The gray points refer to data recorded from different samples. For each sample we carried out several experiments. The best fit using Eq. (13) is displayed as a black line for each sample. Crack front positions are shifted vertically for each sample in order to enhance the visibility. Bottom: Evolution of the average front position, \bar{a} during the relaxation regime (gray dots). The x -axis is the time after we stopped the loading. For each experiment, the best fit using Eq. (12) is displayed as a black line. A good match of the fitted line to the data is observed after several seconds which marks a transition period from the previous regime. Each of the three bottom figures represent a different sample. For each experiment, the vertical axis is shifted in order to enhance visibility.

Figure 6. Top : example of front position, $a(x)$, extracted from a picture captured with the high speed camera. Distances are in pixels and the size of a pixel is $\sim 21 \mu m$. The difference between the less advanced and the farthest advanced point on the front is slightly less than $1 mm$ whereas the front length is $\sim 1.1 cm$. Middle : γ values corresponding to the front line displayed above and computed using Eq. (16). We observe sharp fluctuations of γ which are correlated with local variations of the slope of the front $a(x)$. Bottom : averaged power spectrum (PS) of the front morphology over all fronts extracted during an experiment. We observe a clear power-law decay of the PS which shows the presence of fluctuations of $a(x, t)$ at all spatial scales and is representative of the self-affine nature of the crack front morphology. The dashed line indicates a power law decay function with an exponent of -2.2 , consistent with a roughness exponent $H=0.6$.

Figure 7. Probability density function of the local speed computed for two different experiments (light and dark gray circles). All velocities are normalized by the average crack propagation velocity in order to compare experiments. We clearly observe a power-law decay of the pdf above $v/\langle v \rangle = 1$. The dashed line indicate a power law decay with exponent -2.55 as found by *Måløy et al.* [2006].

Figure 8. Convergence of γ to its final estimate as a function of the half size of the window used to compute it. This convergence is tested on two different front lines taken from two different experiments. Convergence is fast: $\sim 90\%$ of the final estimate of γ is recovered in the first 100 pixels ($2mm$).

Figure 9. Evolution of γ (black line) and front position (gray line) for a given position along front and during ~ 8 s. Stars represent the last values of γ at a given front position and are defined as γ_c .

Figure 10. Zooms on a sub-space of γ matrix obtained using Eq. (16) for two experiments recorded with the fast video camera at 125 fps (top) and 1000 fps (bottom). Time axis is vertical and is oriented from bottom to top. Note that the time interval between two pictures is not the same for the two maps (8 s and 0.6 s). The figure on the top is obtained while the front is relaxing (i.e. $u = \text{constant}$) while the bottom figure was obtained while the front was propagating at constant velocity. Black dots correspond to last positions of the front before a local move and thus define G_c as expressed in Eq. (17). White dots correspond to estimation of G_c obtained from only one measure, thus reflecting the non-equilibrium position of the front at these locations.

Figure 11. Top : Zoom on a map of local fracture energy, G_c , computed using Eq. (17). During this experiment the crack propagates from bottom to top. We clearly observe an heterogeneous distribution of G_c which appears as a random field although a correlation among sites can be identified. The few holes (3 % of total number of pixels) that exist at some places in the map, because of a too low frame rate, have been filled by linear interpolation. The inset graph represents the distribution of fracture energy for all sites of the main picture. Bottom : map of a sample interface morphology, δz of a similar size to the one represented for the fracture energy. An heterogeneous distribution of δz is clearly visible with an amplitude of variation of several tens of micrometers. The inset graph represents the distribution of the interface height δz .

Figure 12. Averaged correlation functions, $C(\Delta r)$, as a function of the distance Δr between two points along the interface. Correlations are averaged from stripes computed along the front direction and perpendicular to it. The gray line denotes the correlation function computed for the interface morphology whereas the black line represents the correlation functions of the fracture energy. The inset figure shows the two power spectra, $S(k)$ from which correlation functions are computed. The best exponential fit for each function $S(k)$ is plotted as a dashed line. The correlation lengths inferred from these fits are $r_c = 188\mu m$ and $84\mu m$ for the height and fracture energy respectively. The values are in agreement with the lengths where the correlations functions $C(\Delta r)$ get uncorrelated.

Figure 13. Distributions of G_c (black line), G (dash line) for all the seven experiments carried with the fast video camera. The vertical black lines is the value of the macroscopic fracture energy \bar{G}_c . We systematically observe that the microscopic distribution of G_c is centered on higher value than their respective macroscopic estimates.



























