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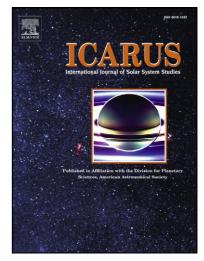
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1	Penetrometry of granular and moist planetary surface materials:
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3	
4	Karl R. Atkinson ^a , John C. Zarnecki ^a , Martin C. Towner ^{a,1} , Timothy J. Ringrose ^a , Axel
5	Hagermann ^a , Andrew J. Ball ^{a,2} , Mark R. Leese ^{a,*} , Gunter Kargl ^b , Mark D. Paton ^{a,3} , Ralph D.
6	Lorenz ^{c,a} and Simon F. Green ^a
7	
8	^a PSSRI, The Open University, Walton Hall, Milton Keynes, MK7 6AA, U.K.
9	^b Space Research Institute, Austrian Academy of Sciences, Schmiedlstraße 6, A-8042 Graz-
10	Messendorf, Austria.
11	^c Space Department, Johns Hopkins University Applied Physics Laboratory, 111100 Johns
12	Hopkins Road, Laurel, MD 20723-6099, USA.
13	
14	¹ Present address: Impacts and Astromaterials Research Centre (IARC), Earth Sciences and
15	Engineering Department, Imperial College London, South Kensington, London, SW7 2AZ, UK.
16	² Present address: European Space Research & Technology Centre, Keplerlaan 1, Postbus 299,
17	2200 AG Noordwijk, The Netherlands.
18	³ Present address: P.O.Box 64 (Gustaf Hällströmin katu 2a), FI-00014 University of Helsinki,
19	Helsinki, Finland.
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- PSSRI, The Open University, Walton Hall, Milton Keynes, MK7 6AA, U.K. Tel +44 (0)1908 652561 Fax: +44 (0)1908 858022 e-mail: <u>m.r.leese@open.ac.uk</u> -mails: arl R. Atkinson: fluid.aspect@googlemail.com

- John C. Zarnecki: j.c.zarnecki@open.ac.uk
- Martin C. Towner: m.towner@ic.ac.uk
- Timothy J. Ringrose: t.j.ringrose@open.ac.uk
- Axel Hagermann: a.hagermann@open.ac.uk
- Andrew J. Ball: Andrew.ball@esa.int
- Mark R. Leese: m.r.leese@open.ac.uk
- Gunter Kargl: guenter.kargl@oeaw.ac.at
- Mark D. Paton: mark.paton@helsinki.fi
- Ralph D. Lorenz: ralph.lorenz@jhuapl.edu
- Simon F. Green: s.f.green@open.ac.uk

46 Abstract

47 The Huygens probe landed on the then unknown surface of Titan in January 2005. A 48 small, protruding penetrometer, part of the Surface Science Package (SSP), was pushed into the 49 surface material measuring the mechanical resistance of the ground as the probe impacted the 50 landing site. We present laboratory penetrometry into room temperature surface analogue 51 materials using a replica penetrometer to investigate further the nature of Titan's surface and 52 examine the sensor's capabilities. The results are then compared to the flight instrument's 53 signature and suggest the Titan surface substrate material consists of sand-sized particles with a 54 mean grain size ~ 2 mm. A possible thin 7 mm coating with mechanical properties similar to 55 terrestrial snow may overlie this substrate, although due to the limited data we are unable to detect any further layering or grading within the near-surface material. The unusual weakening 56 57 with depth of the signature returned from Titan has, to date, only been reproduced using a damp 58 sand target that becomes progressively wetter with depth, and supports the suggestion that the surface may consist of a damp and cohesive material with interstitial liquid contained between its 59 grains. Comparison with terrestrial analogues highlights the unusual nature of the landing site 60 61 material.

⁶³ *Keywords:* Titan; Regoliths; Ices, mechanical properties

64 1 Introduction

On the 14th January 2005, the Huygens probe successfully completed the first landing on 65 66 the unexplored surface of Titan (Lebreton et al. 2005). The probe and Cassini orbiter have since 67 revealed the Titan surface to consist of a wide range of geological features including sand dunes. 68 lakes, mountain chains and dendritic channel networks (Coustenis and Hirtzig, 2009; Jaumann et 69 al., 2008). The probe's scientific instrumentation included the Surface Science Package (SSP) (Zarnecki et al., 1997; Zarnecki et al., 2002), a suite of small sensors primarily aimed at 70 71 characterising the nature of any solid or possibly liquid surface. This instrument included a 72 penetrometer, designated 'ACC-E' or ACCelerometer-External for historical reasons, as part of 73 its array of sensors. The original intention of the penetrometer was to provide a qualitative identification of a terrestrial analogue material for Titan's surface in the event of a (soft) solid 74 75 landing (Lorenz et al., 1994). The penetrometer was designed to measure forces ranging from a 76 few newtons up to a limit of 2000 N. Impact loads any higher than this would probably have caused a failure of the probe (Lorenz et al., 1994). No scientific return was expected from the 77 78 penetrometer in the case of a liquid landing due to the force on the tip being less than the trigger 79 threshold of the sensor (Zarnecki et al., 1997). Additional design details are given in Krysinski et al. (2009). 80

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83

Flight sensor and surface signature

[Insert Fig. 1 here]

The ACC-E penetrometer consists of a piezoelectric element sandwiched between a 14 mm diameter hemispherical tip and collar, mounted at the end of a short aluminium pylon (Fig. 1). This element generates a charge proportional to the stress placed on it, thereby acting as a force transducer allowing a direct measurement of the mechanical resistance of the material

encountered. Constraints were placed on the positioning and length of the flight penetrometer due

89 to the crystal being located directly behind the probe's heat shield (Lorenz *et al.*, 1994). For 90 electromagnetic compatibility, a mesh screen, through which the penetrometer passed, covered 91 the base of the SSP instrument. The electronics were designed to sample the sensor at a rate of 10 kHz such that at the expected impact speed on Titan of 5 m s⁻¹, a theoretical depth resolution of 1 92 93 mm might be obtained. This resolution was intended to allow the possibility of identifying layering in the surface or particle size if the material was granular. The probe landed slightly 94 slower than this at 4.60 \pm 0.05 m s⁻¹ (Towner *et al.*, 2006; Zarnecki *et al.*, 2005). For an impact 95 96 under parachutes at this constant vertical speed, the penetrometer force signature can be plotted 97 against depth penetrated as shown in Fig. 2.

98

88

[insert Fig. 2 here]

This signature can be considered in four distinct stages; however for the purposes of this 99 100 analysis only the first three stages are useful, before the addition of strong structural interactions caused by the arrival of the probe foredome at the surface that make the later data (stage 4) 101 102 unusable. The first stage is a weak but rising force corresponding to a material thickness of 7-8 103 mm. The second stage is a sudden resistance spike of less than 1 ms duration, equivalent to $\sim 2-3$ mm penetration depth at the 4.6 m s⁻¹ impact speed. This is followed by the longest period of 104 105 'clean' penetration before an abrupt rise in force seen in stage 4 which corresponds with the 106 arrival of the Electromagnetic Compatibility (EMC) screen of the Surface Science Package some 107 55 mm after penetrometer tip entry (Zarnecki *et al.*, 2005). The pre-impact signature by design 108 showed 64 force samples before the triggering threshold. Due to shorter, fixed wiring and better 109 screening, noise levels are lower than those measured with the laboratory equipment, occupying

- only the least significant bit level of the 8 bit analogue to digital converter (ADC) of the SSP
 electronics, equivalent to a force of 2 N at Titan surface temperature.
- 112

113 **3** Analogue work

114 Using a replica ACC-E penetrometer and a portable test rig with an electromagnetic 115 release, penetrometry drops were carried out at the probe impact speed both in the laboratory and 116 in the field (for a summary of drops presented in this paper see Table 1). To produce drops 117 comparable to that made by the 200.5 kg probe on Titan, it is only necessary to attach the 118 penetrometer to a small 5 kg mass that provides sufficient inertia to maintain the impact speed for 119 the duration of penetration (Lorenz et al., 1994). The required impact speed was achieved by 120 adjusting the height of the drop, equating the gravitational potential energy lost with the kinetic 121 energy on impact assuming negligible drag. The impact speed for each drop was verified using a 122 laser, photodiode sensor array, and barcode etched mirror attached to the penetrometer weight (Fig. 3). For a fuller description see Atkinson (2008). 123 124 The penetrometry in the laboratory was intended to produce impact signatures

125 representative of similar impacts into a semi-infinite planetary surface. In practice, a laboratory 126 target is necessarily bounded by the rigid sidewalls of the target container and the depth restricted 127 by the container floor. Boundaries may lead to spurious effects on the penetrometer signal caused 128 by the confinement and induced order in the target material near the container wall (Zou and Yu, 129 1995). This leads to a compromise between as large a target as possible, to minimise these 130 effects, and the practicality of producing a target that is manageable in size and not too time 131 consuming to prepare. Boundary effects have been examined extensively and various 'correction 132 factors' have been proposed to account for the effects with varying degrees of success (see e.g. 133 Huang and Hsu, 2005).

134 An analysis of boundary effects is beyond the scope of this study as the effect can be 135 dependent on multiple factors including the type of target material, its preparation, and the 136 sensitivity of the penetrometer. For the purposes of the analogue experiments however, it was not 137 necessary to examine the boundary effects *per se*, but only to establish a target container size of 138 sufficient dimension for them to be negligible. This was done by comparison of the mean force 139 detected between multiple sets of drops into cylindrical targets of varying diameter and depth. 140 The target material was chosen to be 4mm diameter 141 spherical glass beads, as these pack consistently to a narrow range of bulk density, and are easily 142 prepared by pouring. To examine edge effects, three open-bottomed cylindrical containers of 143 internal diameters 104, 152 and 235 mm were placed within a larger 320 mm diameter closed-144 bottomed cylindrical container. All four containers were filled with glass beads to a depth of 350 145 mm. Ten drops were carried out into each size container and the mean signature force recorded. 146 Using a one-way analysis of variance (ANOVA) followed by a Tukey Honestly Significant Difference (HSD) multiple comparison test (see e.g. LeBlanc, 2004), a statistically significant 147 increase in mean signature force was measured between the smallest diameter container and the 148 149 other three containers. To test the effect of the container floor, the 235 mm diameter openbottomed cylindrical container was set on increasing levels of Foamglas, a rigid but brittle 150 151 cellular glass insulating material that would crush, avoiding damage to the penetrometer should it 152 not be stopped by the beads. This produced targets of four depths: 350 mm (no Foamglas), 270, 153 190 and 110 mm respectively. A similar ANOVA and multiple comparison test analysis found 154 the mean penetration force experienced when using the small 110 mm depth container was 155 significantly higher than for the other three containers. These results implied a minimum target 156 container size of ~152 mm diameter and 190 mm depth was needed for boundary effects to be

157 considered negligible. Based on these results, for most experiments the largest container (320 mm 158 diameter, 350 mm depth) was used. 159 [Insert Table 1 here] 160 161 Signal processing path and sensor calibration 4 162 The penetrometer front-end electronics consists of a charge amplifier followed by a 163 pseudo-logarithmic amplifier used to accommodate the large range of input signals that could be 164 generated from the possible impact forces. This pseudo-logarithmic amplifier has three linear 165 gain branches selected depending on the input voltage. In the case of the flight data, only the 166 high-gain branch of this amplifier was used, as the force on the penetrometer was at the low 167 (softer impact) end of the range. An anti-aliasing filter is used to remove frequencies above the sensor's Nyquist frequency (5 kHz) before an 8-bit Analogue to Digital converter (ADC). Figure 168 169 3 summarises the signal processing paths for the flight and laboratory data. 170 [Insert Fig. 3 here] 171 The replica penetrometer was made using a piezoelectric crystal from the same batch as 172 the flight and flight spare penetrometers and constructed in accordance with the original assembly 173 instructions. The charge generated by the piezoelectric crystal leaks away rapidly after the 174 application of a force to the penetrometer tip. For this reason, it is not possible to calibrate the 175 crystal by applying static loads- a known dynamic force must be used instead (Lorenz et al., 176 1994; Zarnecki *et al.*, 1997). The original ACC-E penetrometers were calibrated using a small 177 pendulum jig to strike the sensor tip with different impactors (Lorenz 1994a; Lorenz et al., 1994). 178 For the replica penetrometer, a different method was used based on a mass moving on a 179 frictionless horizontal linear bearing colliding with the force transducer (Fujii, 2006; Fujii and 180 Fujimoto, 1999). This highly accurate calibration method used an optical interferometer to

181 determine the change in momentum of the moving mass before and after the collision. This 182 change is equal in magnitude to the time integrated impulse acting on a force transducer. Given 183 that this degree of calibration accuracy was not required for the penetrometer, the interferometer 184 was replaced with a suitable accelerometer which allowed a direct determination of the impact 185 force for a known impact mass. The limitations of impacting mass that could be supported on the 186 frictionless linear track used meant only the high gain branch was calibrated. Piezoelectric crystal 187 sensitivity varies with temperature and therefore a conversion factor of 1.83 (Lorenz *et al.*, 1994) 188 was required to account for the reduction in sensitivity of the flight signature returned from 189 Titan's -180°C surface compared to the room temperature analogue signatures.

190

191 **5** Identifying substrate material grain size

192 Grain size is of key interest to the Huygens penetrometry investigation due to the strong 193 indication of a sedimentary granular environment at the landing site (Tomasko et al., 2005). One approach to this type of investigation has attempted to reconstruct the grain size distributions of 194 195 Martian soil analogues from their quasi-static penetrometry signatures by examining frequency 196 content (Kargl *et al.*, 2009). The limitations imposed by the length, the relatively low sampling 197 rate of the ACC-E penetrometer and higher speed of impact meant this technique could not be 198 applied. An alternative method was used to examine laboratory penetrometry drops collected at 199 the probe's impact speed into targets of glass and plastic spherical beads of known material 200 densities and with diameters between 1 and 8 mm. The amplitude of each peak in the resulting 201 signatures was identified after discarding peaks due to signal noise determined by examining the 202 pre-impact signal. First, to ensure the data were stationary, a moving average was subtracted from 203 each signature. The signature data points were then stepped through and all consecutive rises

204	between sample points summed until a falling value was found. If the sum of the rises was
205	greater than a 95% confidence level of the pre-impact signal (i.e. peak to trough as the
206	penetrometer is falling towards the target), the peak was considered significant and the point was
207	marked as a peak candidate. For this value to be accepted as a peak, an equivalent fall or sum of
208	consecutive falls must follow this candidate point. A similar method was used to find trough
209	candidates (Fig. 4). The mean value of peak to trough height for each target was then determined.
210	[Insert Fig. 4 here]
211	During this analysis, a regular enhancement of the penetration force in the first 8-10 mm
212	of target penetration or 'tip entry effect' was evident that considerably affected the mean peak to
213	trough height and variability. This distance corresponds approximately to the length of the ACC-
214	E hemispherical tip as it enters the target surface. In general, the effect was found to be greater
215	for larger beads. One possible explanation is given by comparison of the size of the penetrometer
216	tip to the target beads. Smaller bead targets present more of a smooth continuum surface to the tip
217	with each bead having proportionately less effect on it. However, large beads whose size is
218	comparable to the tip form an uneven target surface with variations in both the number and
219	position of the beads impacted. After the tip has fully entered the target it is completely
220	surrounded and therefore gives a more uniform response. Figure 5 shows the magnitude of the
221	mean peak to trough force plotted against the bead mass after removal of the 'tip entry effect'.
222	[Insert Fig. 5 here]
223	A statistically significant empirical relationship between bead mass and the mean peak-
224	trough amplitude, <i>h</i> , was found:
225	$h = 37.79 M^{0.33} \tag{1}$
226	

227 6 Substrate material properties with depth

228	The third stage (Fig. 2) of the flight signature lasts 8.3 ms and is unusual due to a
229	significant downward trend; resistance usually increases as the target material becomes
230	compacted in front of the advancing penetrometer, and/or overburden pressure increases the
231	penetration resistance. Several analogues were tested to try to reproduce this trend. One
232	possibility, given the fluvial nature of the surface seen in DISR images, was that the granular
233	material could be size sorted with depth, which might affect the resistance. This effect, known as
234	graded bedding, is due to changes in the flow speed of the liquid transporting the grains and can
235	result in either 'normal grading', where the grain size coarsens with depth, or 'inverse grading'
236	where the material becomes finer with depth.
237	Artificial graded targets were produced with layers of the four sizes of available glass
238	beads (8, 6, 4 and 1 mm diameter). Whilst some of the inverse graded drops had a downward
239	trend the gradients varied too much to be pursued as a suitable analogue.
240	Further attempts to reproduce a downward trend using clay targets with added voids were
241	only partially successful in modifying the typically flat clay signature to a slightly downward-
242	sloped one. The gradient of the Titan signature is significantly greater over a short length than
243	seen in a clay void target. Early material catalogue work had shown that some water-wetted sand
244	targets were able to produce a downward signature slope although the amount and distribution of
245	liquid had not been well characterized (Paton, 2005). Examining the effect of water on
246	penetration resistance in soils is often done using penetrometers with combined Time Domain
247	Reflectometry (TDR) probes that measure soil water content <i>in situ</i> as they penetrate the soil (see
248	e.g., Vaz et al., 2001). The ACC-E did not have this capability; however by using a ThetaProbe
249	(Delta-T, 1999) to measure the moisture through small sampling points in the target container
250	wall, water distribution could be determined immediately prior to penetrometer impact with

251 minimal target disturbance. Sand targets were prepared in a rigid plastic cylindrical container 252 (diameter 320 mm and sand surface at 350 mm) using a similar method to that used in an 253 investigation to measure elastic wave velocities in saturated sand (Velea et al., 2000). Two target 254 materials were used: well-sorted coarse grained Leighton Buzzard DA 16/30 sand (median grain) 255 size, $D50 = 614 \mu m$) and well-sorted fine grained RH T sand ($D50 = 230 \mu m$) both locally 256 sourced from the WBB Minerals Ltd., Double Arches quarry in Leighton Buzzard, U.K.. The 257 sand targets were poured loosely in the container giving approximate dry bulk densities of 1.51 and 1.57 g cm⁻³ for the LB DA 16/30 and RH T sand respectively. To make a reproducible wet 258 259 target, water from an external feed bottle was introduced into the sand from below through a 260 perforated hose by capillary action causing minimal disturbance to the sand. The sand was 261 allowed to saturate completely to ensure even settlement and a glass straw (open to the 262 atmosphere) in the feed bottle was set at the required water table level. The sand was then 263 allowed to dry overnight until air in the straw of the feed bottle started bubbling up. At this point, 264 the water table is level with the bottom of the straw in the feed bottle. This produces, by capillary 265 action, a moisture gradient between the water table and the surface that can be controlled by 266 adjusting the depth of water table from the surface (or equivalently its height from the fixed 267 bottom of the target container). Immediately prior to the drop into the target, the water content 268 was sampled at several points along the height of the container using the ThetaProbe that after 269 calibration gave a DC voltage related to the soil water moisture content. Figure 6 illustrates the 270 experimental arrangement used.

271

[Insert Fig. 6 here]

Drops were carried out into both sands with four water table configurations, at z=150,
200, 250 and 350 mm (saturated sand) where z is the height of the water table measured from the

274	bottom of the target container. Due to the time-consuming target preparation and subsequent
275	necessity to allow each sand to dry completely before it could be used to prepare a new target,
276	only four drops were carried out into each of the two sands. For comparison with the wet sand
277	signatures, several drops were also carried out into the two dry sands at loose and dense
278	compaction states. Representative signatures from these drops are shown in Fig. 7 and Fig. 8.
279	[Insert Fig. 7 here]
280	[Insert Fig. 8 here]
281	The coarser grained LB sand has a higher penetration resistance than the finer RH T sand
282	and, as expected, the compacted state of each sand type has greater mechanical strength than the
283	loose state. Once water is introduced into these sands even in small quantities, their penetrometry
284	signatures change considerably. Figure 9 shows the RH T sand signatures for the four
285	configurations tested. With the exception of the drop into sand with the lowest water table (d), the
286	characteristic dry sand shape is not seen; the initial tip entry peak of the sand has gone. The other
287	signatures are similar to that of cohesive clay, with a gradual rising slope followed by an
288	essentially constant resistance plateau. This plateau seems to vary in force depending on the
289	height of the water table. The saturated sand, (a), has the lowest resistance, only slightly higher
290	than that of the same sand when dry and in a loosely packed state (Fig. 7). In contrast, the sand
291	with the lowest water table at z=150 mm, (d), has a much greater resistance with a plateau phase
292	nearly twice that of the same sand in the dry dense state. The tip entry peak also returns in this
293	signature. The entry rise of each signature also varies with the moisture content of the sand. The
294	sand with the lowest water table and therefore driest surface, (d), has the sharpest rising force. As
295	the water table is raised and the moisture content of the sand near the surfaces increases, this
296	entry rise becomes increasingly gradual.

297

[Insert Fig. 9 here]

298	Coarser LB sand in the same wet drop configurations produces different results (Fig. 10).
299	A drop into saturated sand (a) shows a considerably reduced tip entry peak compared with drops
300	into the same sand when dry, followed by a constant resistance plateau similar to that seen in the
301	wet RH T sand. In subsequent drops, (b) and (c), as the water level is lowered the initial tip entry
302	peak increases in resistance but falls gradually producing a downward gradient similar to that
303	seen in the Huygens signature. Finally, in drop (d) with the saturation level 200 mm below the
304	target surface, the sand starts to behave as if in the dry state again, with a slightly more
305	pronounced tip-entry peak and a gradual increase in resistance with depth in the plateau phase.
306	These signatures again show how the addition of small quantities of water between the grains of
307	sand can modify the penetrometry signature even over a short penetration distance and, in some
308	cases, more than double the average penetration resistance compared to the same sand in the dry
309	state.

310

[Insert Fig. 10 here]

311

312 7 Comparison of laboratory analogue results to the Huygens landing site signature

313 The laboratory analogue penetrometry work can be compared with the signature returned by the ACC-E penetrometer from the Huygens landing site and some of the findings described 314 previously can be applied. Although the first two stages of the flight data are of extremely short 315 316 duration, the closest match suggests that the penetrometer's first contact with the surface may 317 have been into a thin coating on the substrate material with weaker mechanical properties than terrestrial snow (Fig. 11). At a 4.6 m s⁻¹ impact speed, this putative coating has a thickness of 318 319 ~7mm. One possibility is that this layer is a loose covering of deposited atmospheric aerosols and 320 tholins. However, recent models of Titan's atmosphere indicate a surface deposition rate of 0.01 μ m yr⁻¹ (Soderblom *et al.*, 2007), which would suggest that this coating has been deposited over 321

322	a period greater than 700,000 years. This would seem to be at odds with the current picture of a
323	dynamic Titan surface at the Huygens Landing site. Alternative origins for this layer may be
324	aeolian or fluvial in nature.
325	[Insert Fig. 11 here]
326	In previous analogue work (Zarnecki et al., 2005), the force spike seen in the second stage
327	was reproduced by impacts with small pebbles or with hard crusts; however, the images of the
328	surface taken with the DISR cameras suggest a surface shaped by fluvial processes (Tomasko et
329	al., 2005) with rounded pebbles 10-15 cm in diameter lying on top of a darker, grainy substrate
330	material, making an impact with the first of these analogues seem more likely.
331	Although the presence of water ice at the Titan landing site has yet to be confirmed
332	(Keller et al., 2008), impacts at the Huygens speed were carried out into spherical water ice
333	particles of various masses frozen in liquid nitrogen, to reproduce the peak seen in stage 2 of the
334	Titan signature. While not conclusive, due to the difficulty in controlling the exact point of
335	impact between the penetrometer and ice particle, two drops into the smallest mass particle (~2g)
336	produced a signature peak of similar magnitude to that seen in the signature. This mass
337	corresponds to a spherical particle of approximately 16 mm in diameter, far smaller than the
338	scattered pebbles imaged on the surface by the DISR camera. Given that the accelerometer on the
339	SSP registered a small precursor peak of approximately 0.6 g in magnitude prior to the
340	penetrometer triggering, this would suggest that the probe foredome may have impacted a larger
341	pebble before the first contact of the penetrometer with a granule from the substrate material.
342	The downward pointing DISR Medium Resolution Imager (MRI) also separately imaged
343	the granular substrate material. Due to the limiting resolution of this imager, the smallest size of
344	grain that can be measured from the images is 3 mm. However, applying the results of the
345	laboratory bead drops, that the grain mass affects the amplitude of the peaks, an estimate of the

346 substrate material grain size can be made assuming spherical grains that consist only of water ice which, at Titan's surface temperature, has a density of 0.93 g cm⁻³ (Collins, 2005). Measuring 347 348 the amplitudes of the peak to trough forces in stage 3, identifies 21 peaks of which 11 are above the assumed 1-bit noise level of the penetrometer electronics. These excursions have a mean 349 350 amplitude of 5.53 N which, using the empirical relationship found earlier gives a mean particle mass of 0.003g. Using the relationship between particle diameter, $d_{particle}$, and particle mass, M, 351 15 together with the water ice grain density, ρ , 352

353
$$d_{particle} = 2 \left(\frac{3M}{4\pi\rho} \right)^{\frac{1}{3}} \qquad (2)$$

354	gives a mean grain diameter of ~2 mm. This is a reasonable estimate given the very limited
355	quantity of flight data available and characterizes the grains as 'sand' according to the Udden-
356	Wentworth classification scale (Wentworth, 1922). Other work (Paton, 2005) using a different
357	technique based on peak frequency finds a mean diameter of ~5 mm for dry particles.
358	Considering next the downward trend in this third stage, Fig. 12 shows the closest
359	matching analogue to the Titan substrate material in both gradient and smaller scale undulations.
360	This is coarse wetted LB sand (D50 = 614 μ m) with a water table 150 mm below the sand surface
361	(i.e. z=200). Here the signature has been shifted horizontally, but not rescaled in any way, to
362	align with the third stage of the Huygens signature on which it is overlaid.
363	[Insert Fig. 12 here]
364	To demonstrate how unusual the downward trend seen in the Titan surface material is
365	compared with the terrestrial analogues tested, a material identification map is shown in Fig. 13.
366	Two signature parameters, the mean force of peak to trough undulations and the gradient of the
367	signature were used to classify 17 types of analogue.

368

[Insert Fig. 13 here]

369 The identification map shows that even the closest terrestrial analogue match for this 370 substrate material, a gradient wetted coarse sand, while close to matching the mean peak-trough 371 undulations in the signature caused by granularity, is still some way from matching the degree of 372 loss of strength with depth. This cannot be explained by Titan's reduced gravity compared to 373 penetrometry drops carried out on Earth, as the flight data's signature gradient is negative. 374 Gravity acts to increase overburden pressure with depth and compress underlying material, which 375 would cause a positive signature gradient (and, for the small depth penetrated the effect of 376 overburden pressure would be negligible). It must be noted that these wet analogues used water 377 as their interstitial fluid, which at room temperature has a ~5 times greater viscosity and a surface 378 tension four times that of liquid methane on Titan's surface. Liquid methane at 95K has dynamic viscosity of 2×10^{-4} Pa s and surface tension 0.018 N m⁻¹ (Ghafoor and Zarnecki, 2000; Keller *et* 379 380 al., 2008). On the other hand, the viscosity of liquid ethane at Titan surface temperatures is a little 381 higher than water (e.g. Lorenz et al., 2010). Thus, while we acknowledge that fluid properties 382 and the currently unknown contact angle between water ice and liquid hydrocarbons at Titan 383 surface temperature may be different from those in our experiments, our analogue could in fact 384 be rather similar. Nevertheless, carrying out penetration experiments into ice particles mixed 385 with liquid methane/ethane under Titan's surface temperature and pressure conditions may help 386 constrain the quantity of liquid present between the grains. However, the cost and complexity of 387 producing a sufficiently large environmental chamber within which to perform such experiments 388 would be prohibitive.

8 Other Huygens data that support the penetrometer findings

391 That the substrate material was damp is supported by the measurement of a 40% increase 392 in the level of methane detected by the Gas Chromatograph Mass Spectrometer (GCMS) two 393 minutes after landing while the detection of nitrogen remained constant (Niemann *et al.*, 2005). 394 Modelling of the thermal environment of the inlet supports this possibility, indicating that the 395 heated inlet may have evaporated a small quantity of methane in the local material (Lorenz et al., 396 2006). This suggests that liquid methane may be mixed with the surface material and was 397 evaporated by the GCMS inlet line heater until the liquid became depleted. Detection of a 398 possible dewdrop (Karkoschka and Tomasko, 2009) falling from the descent imager baffle is also 399 consistent with methane moisture being sweated out of the ground by the 20W surface science 400 lamp. The Huygens probe overall endured a deceleration at impact of about 150 m s⁻² (i.e. 15) 401

402 Earth 'g'), the deceleration pulse lasting about 20 ms (Zarnecki et al., 2005; Lorenz et al., 2009). Modelling (Lorenz, 1994b) established the range of target parameters over which the deceleration 403 would be sensitive to the target strength, rather than the deformation of the probe. Indeed, the 404 405 Huygens landing site was soft enough that probe deformation was minimal (although transient structural 'ringing' is evident in some accelerometer data (Bettanini *et al.*, 2008)). Inspection of 406 407 the peak deceleration and the rise time of the pulse (Lorenz et al., 2009) suggest that the surface material averaged over the base of the probe ($\sim 1m^2$, as opposed to the $\sim 1cm^2$ of the penetrometer) 408 409 had some cohesion (i.e. bearing strength at zero penetration). This further supports the 'damp 410 sand' model discussed in the present paper. Dry sand has a rather longer rise time than was 411 observed.

Rock counts from the SLI imager (Keller *et al.*, 2008) identify a few gravel-sized particles
great than ~5mm in the substrate material, although the majority of the material imaged is finer

414	than this. Nevertheless, sand at the landing site may have been rather coarse – Tomasko et al.
415	(2005) note that particles 3mm across (the limiting resolution of the camera) can be identified in
416	the post-impact images.
417	
418	9 Conclusions
419	The ACC-E penetrometer made a single direct measurement of the mechanical properties
420	and texture of the surface material at the Huygens landing site. Comparative analysis of the
421	returned signature with those of terrestrial analogues taken together with measurements from
422	other instruments allows an interpretation of the surface material to be made. This work suggests
423	that the penetrometer is likely to have impacted a thin, extremely weak surface coating overlying
424	a small, hard substrate material granule before being driven into a coarse granular sand substrate
425	possibly wet with liquid methane.
426	
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433	
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Figure captions

543	Figure 1. Schematic of the SSP instrument showing the location of the penetrometer at the base
544	of the probe immediately behind the ablative heat shield. Dimensions of the penetrometer and the
545	approximate position of the Electromagnetic Compatibility (EMC) screen are shown. (Upper
546	Image credit James Garry)
547	
548	Figure 2. ACC-E penetrometry signature returned from the surface of Titan calibrated and
549	processed into force against depth, including correction for the transfer function of the
550	electronics. No smoothing has been applied. The penetration analysis stages are marked.
551	
552	Figure 3. Free fall drop rig arrangement and data acquisition equipment. For the purposes of
553	comparison with the flight data impact speed of 4.6 m s ⁻¹ , the drop height of the penetrometer
554	above the target surface was set to 1.07 m. Bottom: schematic diagram of signal and data
555	processing paths for the flight and laboratory data.
556	
557	Figure 4. Illustration of the peak and trough finding algorithm. The search starts from the
558	beginning of target penetration and identifies candidate peaks and troughs. To be accepted each
559	candidate peak and trough must have a significant rise and fall on either side. The significance is
560	based on the standard deviation of the noise on the signal as the penetrometer falls prior to
561	impact.
562	
563	Figure 5. Mean peak-to-trough force amplitude plotted against bead mass for 31 drops into plastic
564	and glass bead targets. Standard errors are marked when larger than the data point.
565	

566	Figure 6. Equipment used to produce a moisture gradient in sand with minimum mechanical
567	disturbance using a siphon principle. Moisture content was measured prior to each drop using a
568	ThetaProbe inserted into the sides of the container. Sampling points were offset to minimise sand
569	disturbance.
570	
571	Figure 7. Penetrometry signatures for dry RH T (fine grained) sand in two densities, achieved by
572	physical compaction of the target. Impact is at 4.6 m s ⁻¹ .
573	9
574	Figure 8. Penetrometer signature for dry LB (coarse grained) sand in loose and dense state
575	showing noticeably greater tip entry impact peaks than those seen in the finer RH T sand. Impact
576	speed is 4.6 m s^{-1} .
577	
578	Figure 9. Penetrometer signature for RH T (fine grained) sand with water: (a) saturated sand
579	(water table height z=350 mm) (b) water table at $z = 250$ mm (c) water table at $z = 200$ mm (d)
580	water table at $z = 150$ mm.
581	
582	Figure 10. Penetrometer signatures for wet LB (coarse grained) sand: (a) saturated sand (water
583	table height $z = 350$ mm) (b) water table at $z = 250$ mm (c) water table at $z = 200$ mm (d) water
584	table at $z = 150$ mm.
585	
586	Figure 11. Penetrometer signature of a drop into loosely packed snow at Huygens impact speed.
587	The signature from the probe (grey) has also been plotted for comparison of the initial entry force

- 588 (arrowed). The snow signature indicates the variation in density with depth of the material
- 589 probably caused by packing it into the target container.

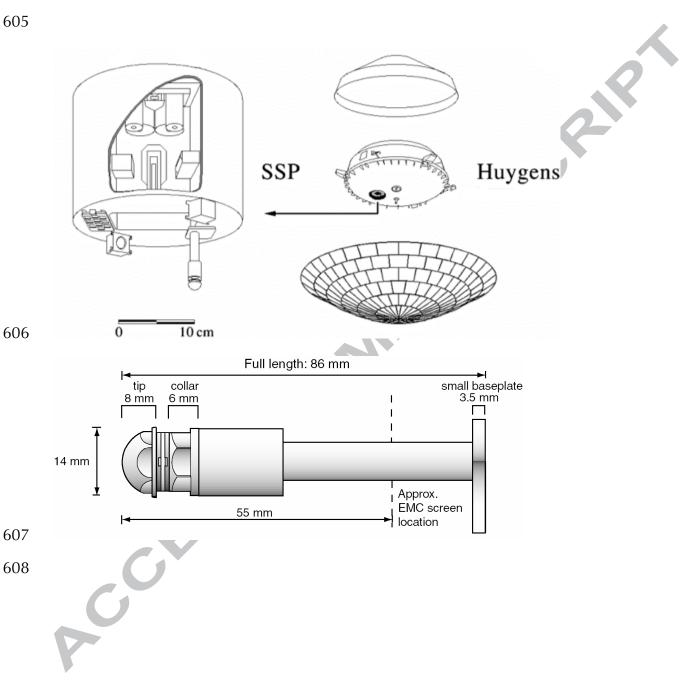
590

591 Figure 12. Penetrometry signature from Huygens overlaid on the closest analogue match for the 592 third stage (substrate) material, wet Leighton Buzzard coarse sand with a water table 150 mm 593 below the sand surface (i.e. z=200). The flight signature has been shifted horizontally (but not 594 rescaled) to align the third stage with the start of the analogue drop. Penetration speed in both 595 cases is 4.6 m s^{-1} .

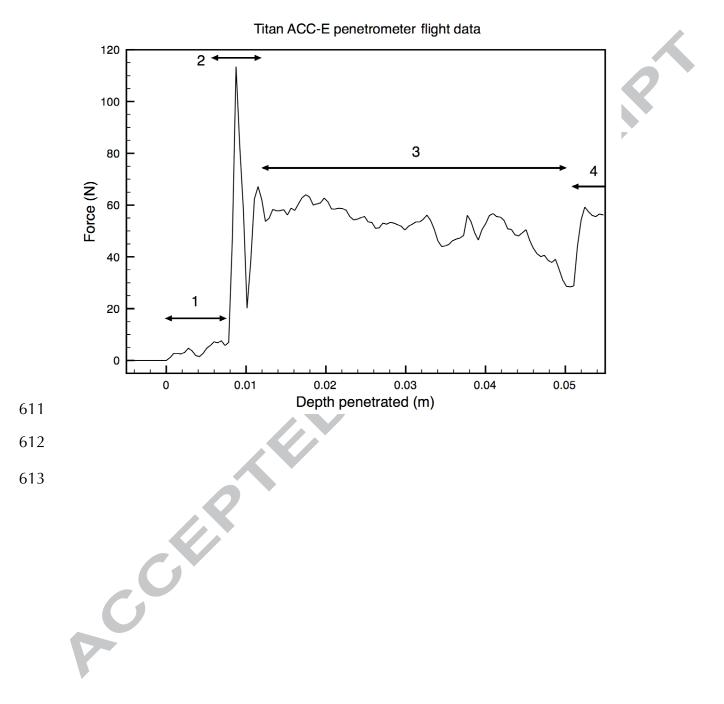
596

597 Figure 13. Surface material identification using two signature parameters. Parameters exclude the first 1 cm 'tip entry' of signature for comparative purposes with the Huygens third stage of 598 599 signature. Where multiple drops of an analogue are available averages have been taken and 600 standard errors are shown. There is a marked difference between the flight data's signature 601 gradient compared to the terrestrial analogues.

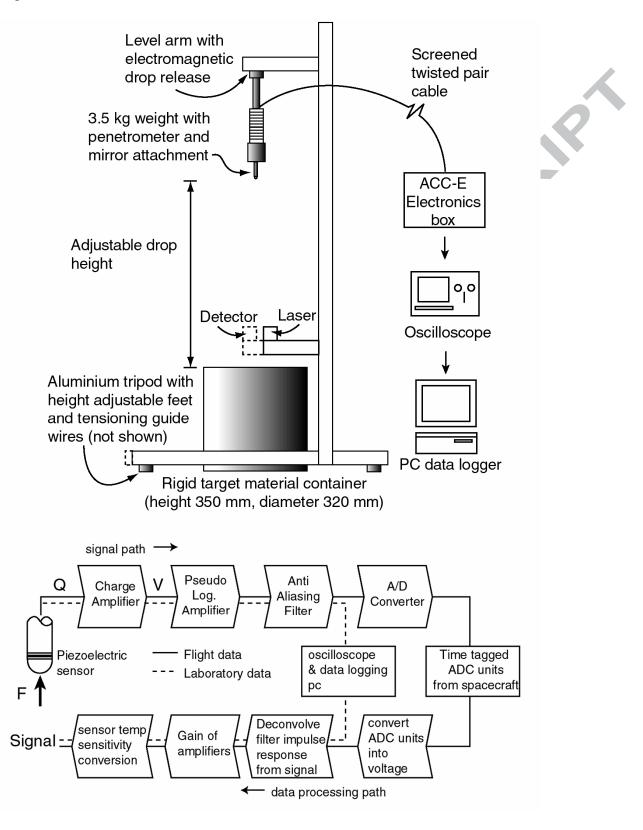
Figure 1



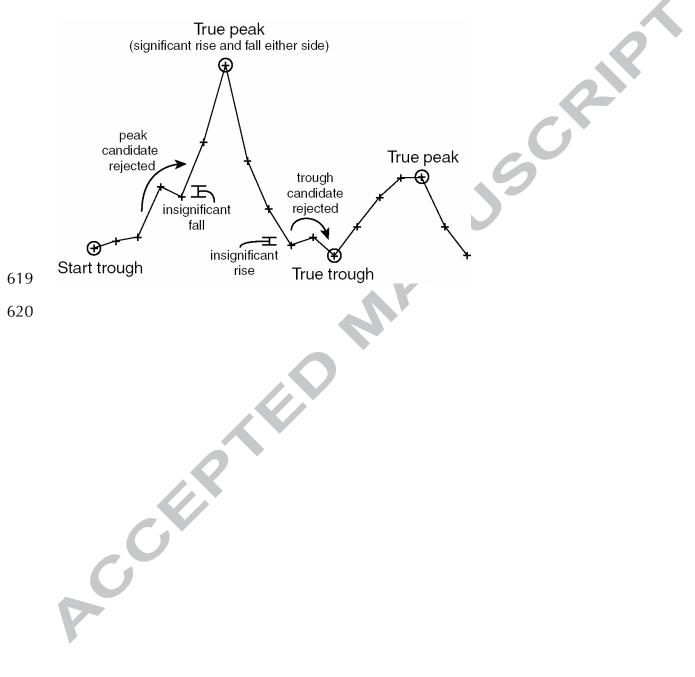
609 Figure 2

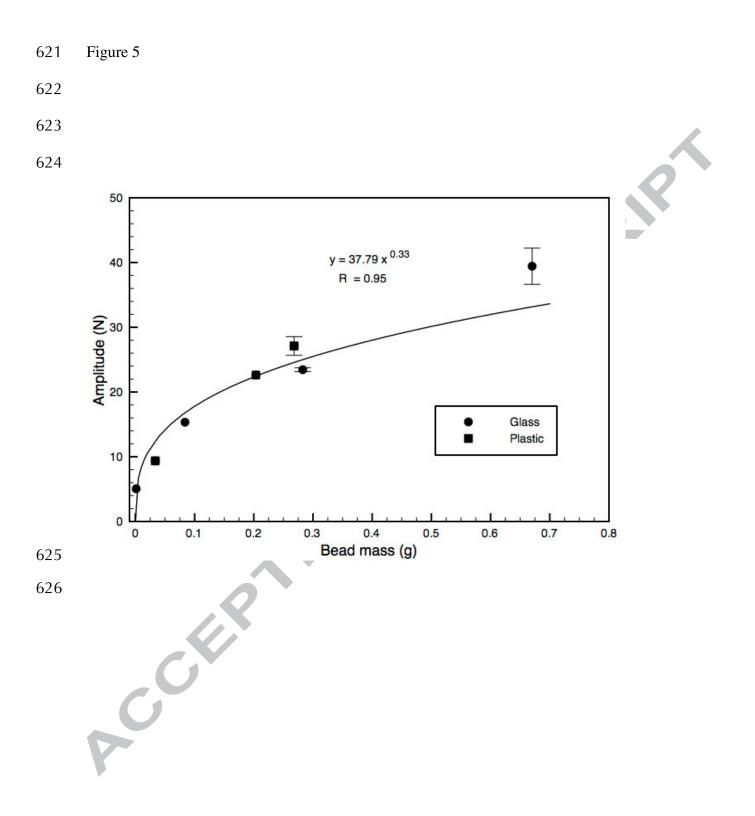


614 Figure 3



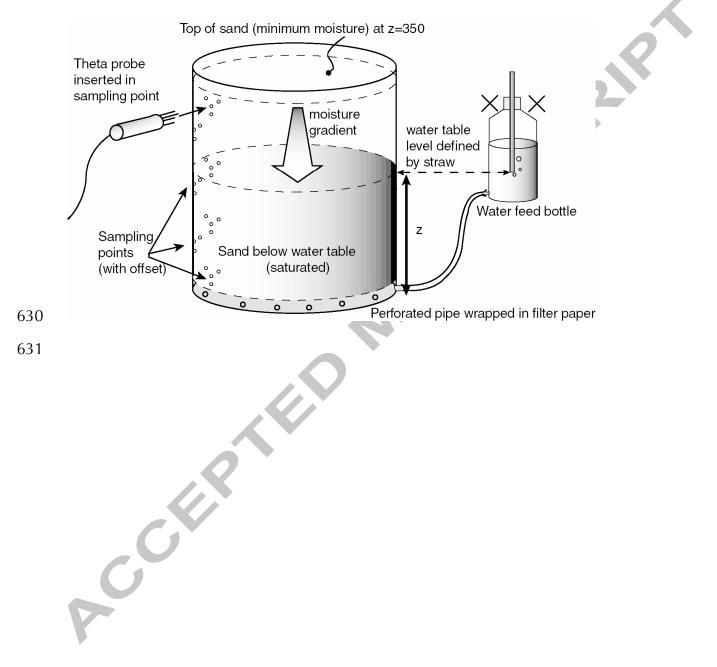


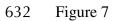




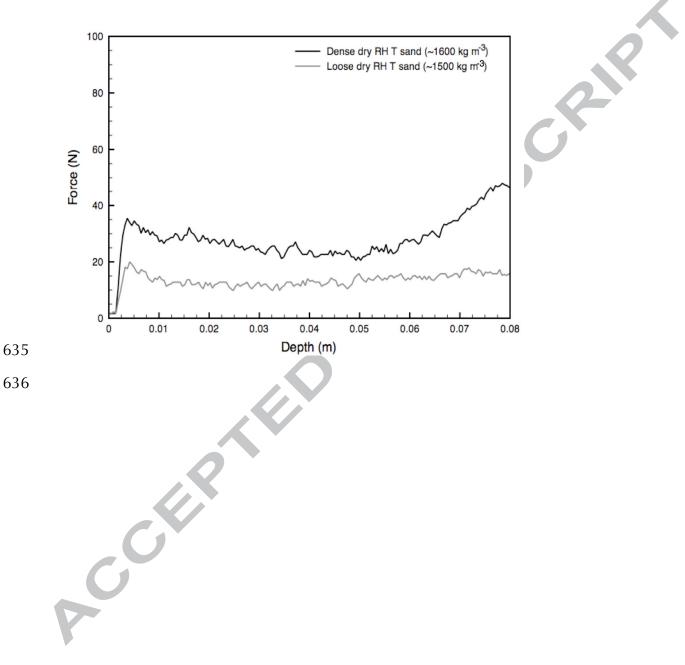
627 Figure 6

628



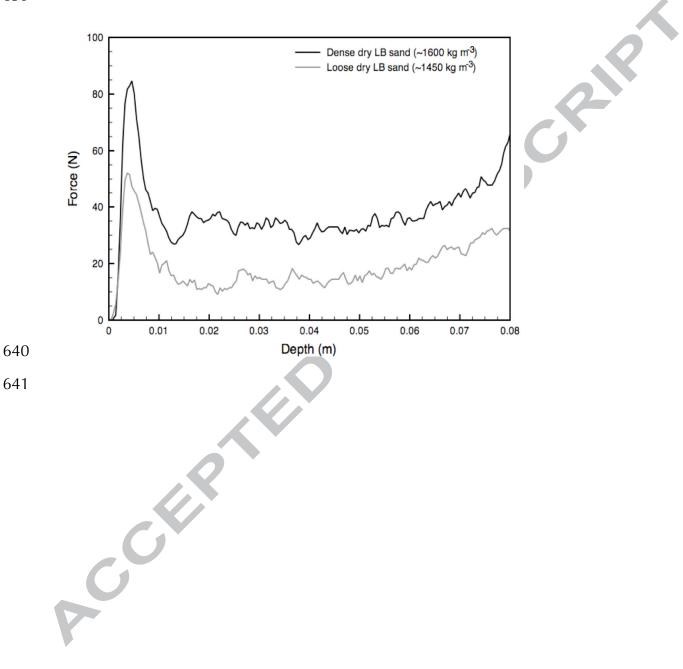


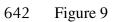














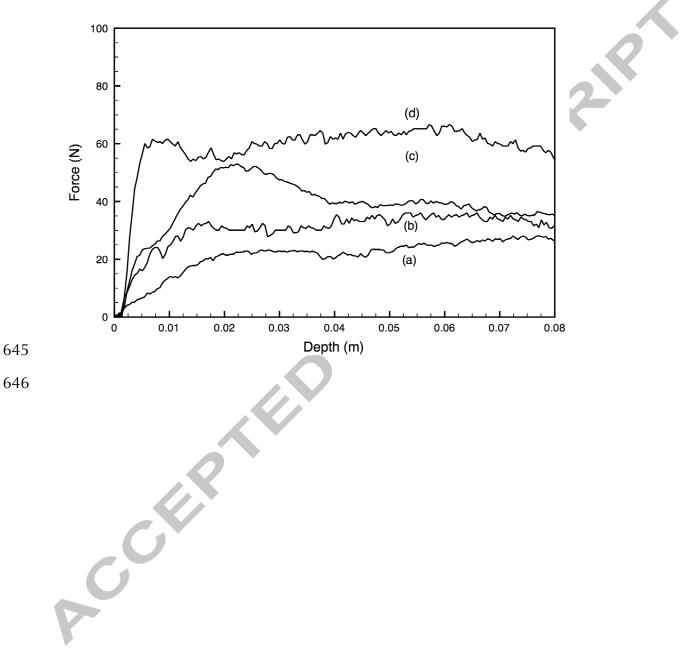
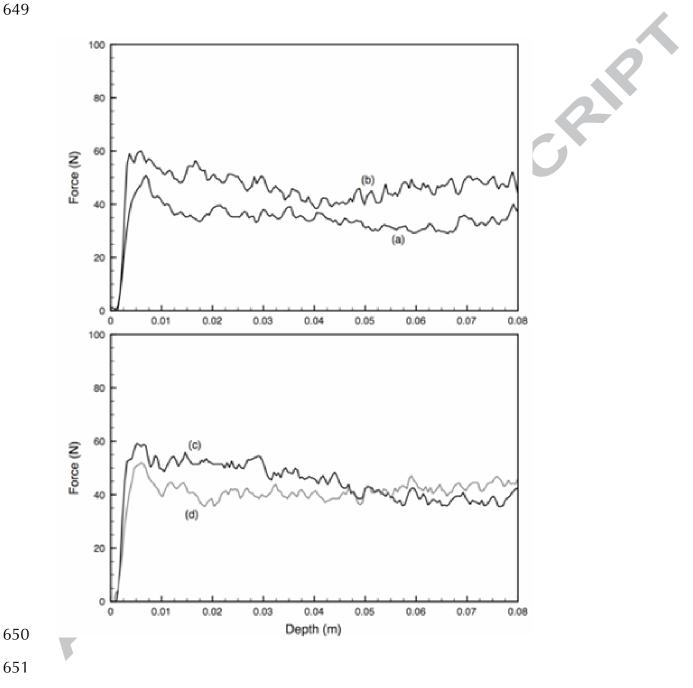
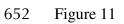
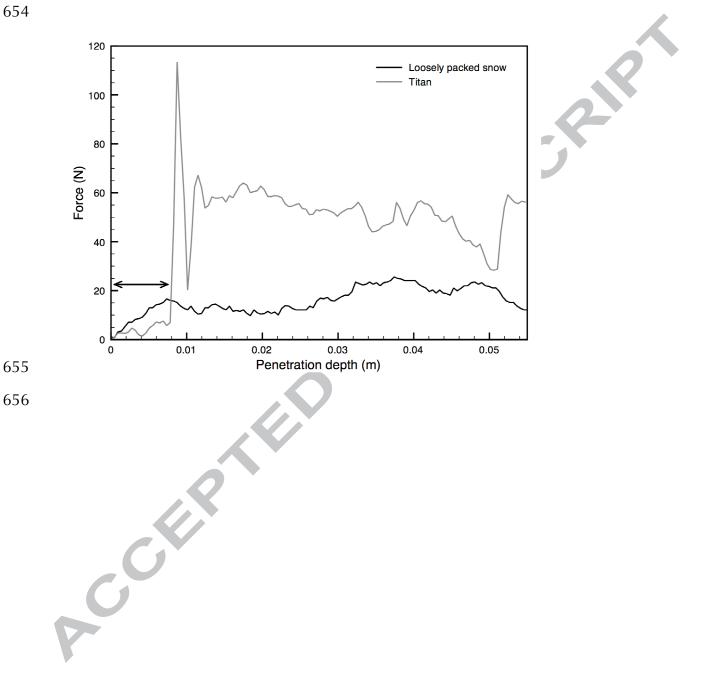


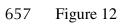
Figure 10



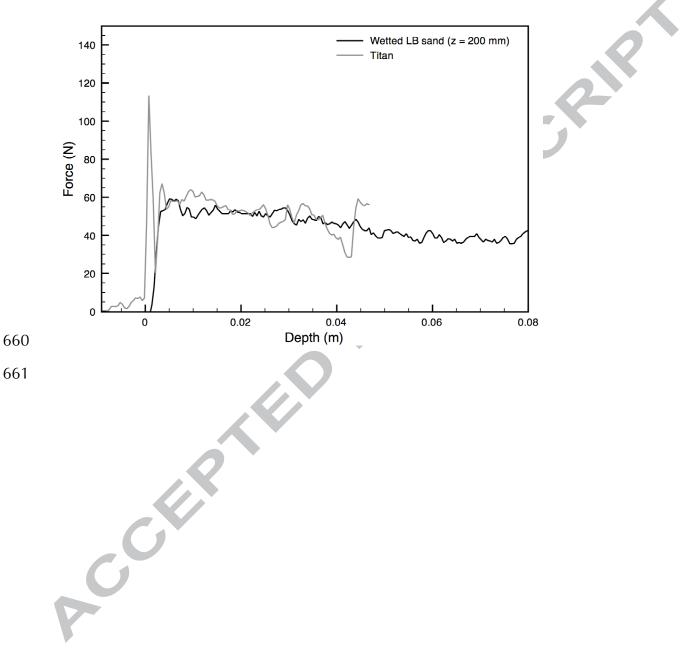












662 Figure 13

663

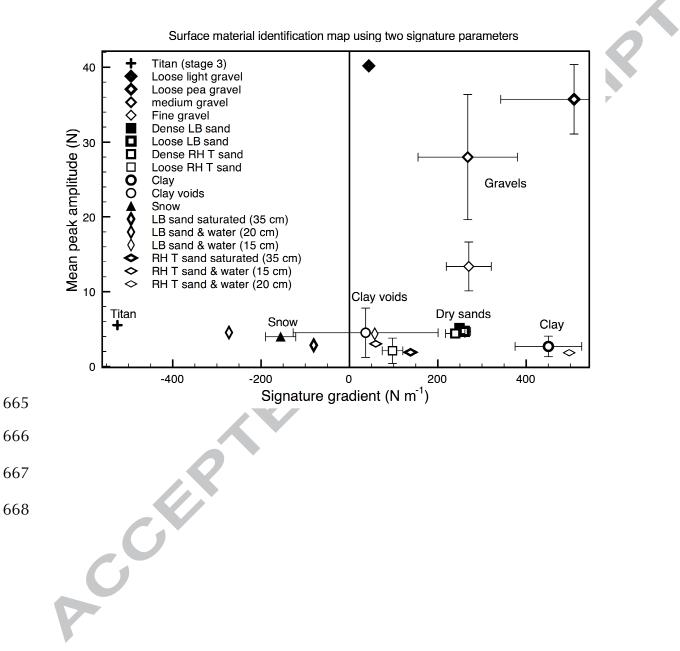


Table 1

Summary of penetration drops carried out. Target impact speeds are all 4.6 m s⁻¹ unless stated otherwise.

Target material	Number of drops	Target preparation notes (all dimensions in mm)	Remarks
Glass beads, (1,4,6,8 mm diameters)	Total 80 drops (10 into each configuration)	Edge and boundary effects: 4 mm glass beads loosely poured into rigid container of diameters 104, 152, 235, 320, and depths 110, 190, 270 and 350. Lower drop speed 4 m s ⁻¹ to avoid damaging penetrometer.	219
	16	Peaks mass relationships: 130 dia. 150 depth with loose surrounding sand to container size 320 dia. 350 depth (For comparative purposes, see plastic beads below).	Figure 5
Plastic beads (4,6 and 8 mm diameters)	15	130 dia. 150 depth target with loose surrounding sand to outer container 320 dia. 350 depth due to limitations on available target material. For further details see Atkinson (2008).	Figure 5
Sands: Dry RH T sand, Dry LB sand both in loose and dense compactions	14	Target size 320 dia. 350 depth. Loosely poured or tamped to compacted density.	Figure 7 Figure 8
Gravels: pea (4 mm diameter), fine (4.5 mm), light (5.5 mm) and medium (7 mm)	12	Target size 320 diameter, 350 depth. Poured loosely into target container.	Figure 13
Clay and clay with voids	6	Compacted clay blocks with approximate dimensions $270 \times 150 \times 130$ on sand. 6 mm voids created by pushing out clay cylinders using a metal tube.	Figure 13
Wet LB sand and Wet RH T sand	8	Target size 320 diameter 350 depth. Targets prepared by loosely pouring sand and filling with siphon technique and settling overnight-see text.	Figure 9 Figure 10 Figure 12 Figure 13
Snow (lightly compacted)	2	Target prepared by layering collected blocks of snow with minimal compaction. Approximate target size 320 dia. 350 depth.	Figure 11 Figure 13