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Lectures on Gaussian approximations with Malliavin calculus

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Overview. In a seminal paper of 2005, Nualart and Peccati [37] discovered a surprising central limit theorem (called the "Fourth Moment Theorem" in the sequel) for sequences of multiple stochastic integrals of a fixed order: in this context, convergence in distribution to the standard normal law is equivalent to convergence of just the fourth moment. Shortly afterwards, Peccati and Tudor [44] gave a multidimensional version of this characterization.

Since the publication of these two beautiful papers, many improvements and developments on this theme have been considered. Among them is the work by Nualart and Ortiz-Latorre [36], giving a new proof only based on Malliavin calculus and the use of integration by parts on Wiener space. A second step is my joint paper [25] (written in collaboration with Peccati) in which, by bringing together Stein's method with Malliavin calculus, we have been able (among other things) to associate quantitative bounds to the Fourth Moment Theorem. It turns out that Stein's method and Malliavin calculus fit together admirably well. Their interaction has led to some remarkable new results involving central and non-central limit theorems for functionals of infinite-dimensional Gaussian fields.

The current survey aims to introduce the main features of this recent theory. It originates from a series of lectures I delivered* at the Collège de France between January and March 2012, within the framework of the annual prize of the Fondation des Sciences Mathématiques de Paris. It may be seen as a teaser for the book [29], in which the interested reader will find much more than in this short survey.

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^{*}You may watch the videos of the lectures at http://www.sciencesmaths-paris.fr/index.php?page=175.

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1 Breuer-Major Theorem

The aim of this first section is to illustrate, through a guiding example, the power of the approach we will develop in this survey.

Let $\{X_k\}_{k\geqslant 1}$ be a centered stationary Gaussian family. In this context, stationary just means that there exists $\rho: \mathbb{Z} \to \mathbb{R}$ such that $E[X_k X_l] = \rho(k-l), \ k, l \geqslant 1$. Assume further that $\rho(0) = 1$, that is, each X_k is $\mathcal{N}(0,1)$ distributed.

Let $\varphi : \mathbb{R} \to \mathbb{R}$ be a measurable function satisfying

$$E[\varphi^{2}(X_{1})] = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} \varphi^{2}(x)e^{-x^{2}/2}dx < \infty.$$
 (1.1)

Let H_0, H_1, \ldots denote the sequence of Hermite polynomials. The first few Hermite polynomials are $H_0 = 1$, $H_1 = X$, $H_2 = X^2 - 1$ and $H_3 = X^3 - 3X$. More generally, the qth Hermite polynomial H_q is defined through the relation $XH_q = H_{q+1} + qH_{q-1}$. It is a well-known fact that, when it verifies (1.1), the function φ may be expanded in $L^2(\mathbb{R}, e^{-x^2/2}dx)$ (in a unique way) in terms of Hermite polynomials as follows:

$$\varphi(x) = \sum_{q=0}^{\infty} a_q H_q(x). \tag{1.2}$$

Let $d \ge 0$ be the first integer $q \ge 0$ such that $a_q \ne 0$ in (1.2). It is called the *Hermite rank* of φ ; it will play a key role in our study. Also, let us mention the following crucial property of Hermite

polynomials with respect to Gaussian elements. For any integer $p, q \ge 0$ and any jointly Gaussian random variables $U, V \sim \mathcal{N}(0, 1)$, we have

$$E[H_p(U)H_q(V)] = \begin{cases} 0 & \text{if } p \neq q \\ q! E[UV]^q & \text{if } p = q. \end{cases}$$

$$(1.3)$$

In particular (choosing p=0) we have that $E[H_q(X_1)]=0$ for all $q\geqslant 1$, meaning that $a_0=E[\varphi(X_1)]$ in (1.2). Also, combining the decomposition (1.2) with (1.3), it is straightforward to check that

$$E[\varphi^2(X_1)] = \sum_{q=0}^{\infty} q! a_q^2.$$
 (1.4)

We are now in position to state the celebrated Breuer-Major theorem.

Theorem 1.1 (Breuer, Major, 1983; see [6]) Let $\{X_k\}_{k\geqslant 1}$ and $\varphi: \mathbb{R} \to \mathbb{R}$ be as above. Assume further that $a_0 = E[\varphi(X_1)] = 0$ and that $\sum_{k\in\mathbb{Z}} |\rho(k)|^d < \infty$, where ρ is the covariance function of $\{X_k\}_{k\geqslant 1}$ and d is the Hermite rank of φ (observe that $d\geqslant 1$). Then, as $n\to\infty$,

$$V_n = \frac{1}{\sqrt{n}} \sum_{k=1}^n \varphi(X_k) \stackrel{\text{law}}{\to} \mathcal{N}(0, \sigma^2), \tag{1.5}$$

with σ^2 given by

$$\sigma^2 = \sum_{q=d}^{\infty} q! a_q^2 \sum_{k \in \mathbb{Z}} \rho(k)^q \in [0, \infty). \tag{1.6}$$

(The fact that $\sigma^2 \in [0, \infty)$ is part of the conclusion.)

The proof of Theorem 1.1 is far from being obvious. Before doing so, let us make some relevant comments.

Remark 1.2 1. First, it is worthwhile noticing that Theorem 1.1 (strictly) contains the classical central limit theorem (CLT), which is not an evident claim at first glance. Indeed, let $\{Y_k\}_{k\geqslant 1}$ be a sequence of i.i.d. centered random variables with common variance $\sigma^2 > 0$, and let F_Y denote the common cumulative distribution function. Consider the pseudo-inverse F_Y^{-1} of F_Y , defined as

$$F_Y^{-1}(u) = \inf\{y \in \mathbb{R} : u \leqslant F_Y(y)\}, \quad u \in (0,1).$$

When $U \sim \mathcal{U}_{[0,1]}$ is uniformly distributed, it is well-known that $F_Y^{-1}(U) \stackrel{\text{law}}{=} Y_1$. Observe also that $\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{X_1} e^{-t^2/2} dt$ is $\mathcal{U}_{[0,1]}$ distributed. By combining these two facts, we get that $\varphi(X_1) \stackrel{\text{law}}{=} Y_1$ with

$$\varphi(x) = F_Y^{-1} \left(\frac{1}{\sqrt{2\pi}} \int_{-\infty}^x e^{-t^2/2} dt \right), \quad x \in \mathbb{R}.$$

Assume now that $\rho(0) = 1$ and $\rho(k) = 0$ for $k \neq 0$, that is, assume that the sequence $\{X_k\}_{k \geq 1}$ is composed of i.i.d. $\mathcal{N}(0,1)$ random variables. Theorem 1.1 yields that

$$\frac{1}{\sqrt{n}} \sum_{k=1}^{n} Y_k \stackrel{\text{law}}{=} \frac{1}{\sqrt{n}} \sum_{k=1}^{n} \varphi(X_k) \stackrel{\text{law}}{\to} \mathcal{N} \left(0, \sum_{q=d}^{\infty} q! a_q^2 \right),$$

thereby concluding the proof of the CLT since $\sigma^2 = E[\varphi^2(X_1)] = \sum_{q=d}^{\infty} q! a_q^2$, see (1.4). Of course, such a proof of the CLT is like to crack a walnut with a sledgehammer. This approach has nevertheless its merits: it shows that the independence assumption in the CLT is not crucial to allow a Gaussian limit. Indeed, this is rather the summability of a series which is responsible of this fact, see also the second point of this remark.

- 2. Assume that $d \ge 2$ and that $\rho(k) \sim |k|^{-D}$ as $|k| \to \infty$ for some $D \in (0, \frac{1}{d})$. In this case, it may be shown that $n^{dD/2-1} \sum_{k=1}^{n} \varphi(X_k)$ converges in law to a <u>non-Gaussian</u> (non degenerated) random variable. This shows in particular that, in the case where $\sum_{k \in \mathbb{Z}} |\rho(k)|^d = \infty$, we usually get a non-Gaussian limit. In other words, the summability assumption in Theorem 1.1 is, roughly speaking, equivalent (when $d \ge 2$) to the asymptotic normality.
- 3. There exists a functional version of Theorem 1.1, in which the sum $\sum_{k=1}^{n}$ is replaced by $\sum_{k=1}^{[nt]}$ for $t \ge 0$. It is actually not that much harder to prove and, unsurprisingly, the limiting process is then the standard Brownian motion multiplied by σ .
- 4. The original proof of Theorem 1.1 consisted to show that all the moments of V_n converge to those of the Gaussian law $\mathcal{N}(0, \sigma^2)$. As anyone might guess, this required a high ability and a lot of combinatorics. In the proof we will offer, the complexity is the same than checking that the variance and the fourth moment of V_n converges to σ^2 and $3\sigma^4$ respectively, which is a drastic simplification with respect to the original proof.

Let us now prove Theorem 1.1. We first compute the limiting variance, which will justify the formula (1.6) we claim for σ^2 . Thanks to (1.2) and (1.3), we can write

$$E[V_n^2] = \frac{1}{n} E\left[\left(\sum_{q=d}^{\infty} a_q \sum_{k=1}^n H_q(X_k)\right)^2\right] = \frac{1}{n} \sum_{p,q=d}^{\infty} a_p a_q \sum_{k,l=1}^n E[H_p(X_k) H_q(X_l)]$$

$$= \frac{1}{n} \sum_{q=d}^{\infty} q! a_q^2 \sum_{k,l=1}^n \rho(k-l)^q = \sum_{q=d}^{\infty} q! a_q^2 \sum_{r \in \mathbb{Z}} \rho(r)^q \left(1 - \frac{|r|}{n}\right) \mathbf{1}_{\{|r| < n\}}.$$

When $q \geqslant d$ and $r \in \mathbb{Z}$ are fixed, we have that

$$q! a_q^2 \rho(r)^q \left(1 - \frac{|r|}{n}\right) \mathbf{1}_{\{|r| < n\}} \to q! a_q^2 \rho(r)^q \text{ as } n \to \infty.$$

On the other hand, using that $|\rho(k)| = |E[X_1X_{k+1}]| \leqslant \sqrt{E[X_1^2]E[X_{1+k}^2]} = 1$, we have

$$q!a_q^2|\rho(r)|^q \left(1 - \frac{|r|}{n}\right) \mathbf{1}_{\{|r| < n\}} \leqslant q!a_q^2|\rho(r)|^q \leqslant q!a_q^2|\rho(r)|^d,$$

with $\sum_{q=d}^{\infty} \sum_{r \in \mathbb{Z}} q! a_q^2 |\rho(r)|^d = E[\varphi^2(X_1)] \times \sum_{r \in \mathbb{Z}} |\rho(r)|^d < \infty$, see (1.4). By applying the dominated convergence theorem, we deduce that $E[V_n^2] \to \sigma^2$ as $n \to \infty$, with $\sigma^2 \in [0, \infty)$ given by (1.6).

Let us next concentrate on the proof of (1.5). We shall do it in three steps of increasing generality (but of decreasing complexity!):

- (i) when $\varphi = H_q$ has the form of a Hermite polynomial (for some $q \ge 1$);
- (ii) when $\varphi = P \in \mathbb{R}[X]$ is a real polynomial;

(iii) in the general case when $\varphi \in L^2(\mathbb{R}, e^{-x^2/2}dx)$.

We first show that (ii) implies (iii). That is, let us assume that Theorem 1.1 is shown for polynomial functions φ , and let us show that it holds true for any function $\varphi \in L^2(\mathbb{R}, e^{-x^2/2}dx)$. We proceed by approximation. Let $N \geqslant 1$ be a (large) integer (to be chosen later) and write

$$V_n = \frac{1}{\sqrt{n}} \sum_{q=d}^{N} a_q \sum_{k=1}^{n} H_q(X_k) + \frac{1}{\sqrt{n}} \sum_{q=N+1}^{\infty} a_q \sum_{k=1}^{n} H_q(X_k) =: V_{n,N} + R_{n,N}.$$

Similar computations as above leads to

$$\sup_{n\geqslant 1} E[R_{n,N}^2] \leqslant \sum_{q=N+1}^{\infty} q! a_q^2 \times \sum_{r\in\mathbb{Z}} |\rho(r)|^d \to 0 \text{ as } N \to \infty.$$

$$(1.7)$$

(Recall from (1.4) that $E[\varphi^2(X_1)] = \sum_{q=d}^{\infty} q! a_q^2 < \infty$.) On the other hand, using (ii) we have that, for fixed N and as $n \to \infty$,

$$V_{n,N} \stackrel{\text{law}}{\to} \mathcal{N} \left(0, \sum_{q=d}^{N} q! a_q^2 \sum_{k \in \mathbb{Z}} \rho(k)^q \right).$$
 (1.8)

It is then a routine exercise (details are left to the reader) to deduce from (1.7)-(1.8) that $V_n = V_{n,N} + R_{n,N} \xrightarrow{\text{law}} \mathcal{N}(0, \sigma^2)$ as $n \to \infty$, that is, that (iii) holds true.

Next, let us prove (i). We actually need to work with a specific realization of the sequence $\{X_k\}_{k\geqslant 1}$. The space

$$\mathcal{H} := \overline{\operatorname{span}\{X_1, X_2, \ldots\}}^{L^2(\Omega)}$$

being a real separable Hilbert space, it is isometrically isomorphic to either \mathbb{R}^N (with $N \geq 1$) or $L^2(\mathbb{R}_+)$. Let us assume that $\mathcal{H} \simeq L^2(\mathbb{R}_+)$, the case where $\mathcal{H} \simeq \mathbb{R}^N$ being easier to handle. Let $\Phi: \mathcal{H} \to L^2(\mathbb{R}_+)$ be an isometry. Set $e_k = \Phi(X_k)$ for each $k \geq 1$. We have

$$\rho(k-l) = E[X_k X_l] = \int_0^\infty e_k(x)e_l(x)dx, \quad k, l \geqslant 1$$
(1.9)

If $B = (B_t)_{t \ge 0}$ denotes a standard Brownian motion, we deduce that

$$\{X_k\}_{k\geqslant 1} \stackrel{\text{law}}{=} \left\{ \int_0^\infty e_k(t) dB_t \right\}_{k\geqslant 1},$$

these two families being indeed centered, Gaussian and have the same covariance structure (by construction of the e_k 's). On the other hand, it is a well-known result of stochastic analysis (which follows from an induction argument through the Itô's formula) that, for any function $e \in L^2(\mathbb{R}_+)$ such that $||e||_{L^2(\mathbb{R}_+)} = 1$, we have

$$H_q\left(\int_0^\infty e(t)dB_t\right) = q! \int_0^\infty dB_{t_1}e(t_1) \int_0^{t_1} dB_{t_2}e(t_2) \dots \int_0^{t_{q-1}} dB_{t_q}e(t_q). \tag{1.10}$$

(For instance, by Itô's formula we can write

$$\left(\int_0^\infty e(t)dB_t\right)^2 = 2\int_0^\infty dB_{t_1}e(t_1)\int_0^{t_1} dB_{t_2}e(t_2) + \int_0^\infty e(t)^2 dt$$
$$= 2\int_0^\infty dB_{t_1}e(t_1)\int_0^{t_1} dB_{t_2}e(t_2) + 1,$$

which is nothing but (1.10) for q = 2, since $H_2 = X^2 - 1$.) At this stage, let us adopt the two following notational conventions:

- (a) If φ (resp. ψ) is a function of r (resp. s) arguments, then the tensor product $\varphi \otimes \psi$ is the function of r+s arguments given by $\varphi \otimes \psi(x_1,\ldots,x_{r+s}) = \varphi(x_1,\ldots,x_r)\psi(x_{r+1},\ldots,x_{r+s})$. Also, if $q \geqslant 1$ is an integer and e is a function, the tensor product function $e^{\otimes q}$ is the function $e \otimes \ldots \otimes e$ where e appears q times.
- (b) If $f \in L^2(\mathbb{R}^q_+)$ is symmetric (meaning that $f(x_1,\ldots,x_q) = f(x_{\sigma(1)},\ldots,x_{\sigma(q)})$ for all permutation $\sigma \in \mathfrak{S}_q$ and almost all $x_1,\ldots,x_q \in \mathbb{R}_+$) then

$$I_q(f) = \int_{\mathbb{R}^q_{\perp}} f(t_1, \dots, t_q) dB_{t_1} \dots dB_{t_q} := q! \int_0^{\infty} dB_{t_1} \int_0^{t_1} dB_{t_2} \dots \int_0^{t_{q-1}} dB_{t_q} f(t_1, \dots, t_q).$$

With these new notations at hand, observe that we can rephrase (1.10) in a simple way as

$$H_q\left(\int_0^\infty e(t)dB_t\right) = I_q(e^{\otimes q}). \tag{1.11}$$

It is now time to introduce a very powerful tool, the so-called *Fourth Moment Theorem* of Nualart and Peccati. This wonderful result lies at the heart of the approach we shall develop in these lecture notes. We will prove it in Section 5.

Theorem 1.3 (Nualart, Peccati, 2005; see [37]) Fix an integer $q \ge 2$, and let $\{f_n\}_{n \ge 1}$ be a sequence of symmetric functions of $L^2(\mathbb{R}^q_+)$. Assume that $E[I_q(f_n)^2] = q! \|f_n\|_{L^2(\mathbb{R}^q_+)}^2 \to \sigma^2$ as $n \to \infty$ for some $\sigma > 0$. Then, the following three assertions are equivalent as $n \to \infty$:

- (1) $I_q(f_n) \stackrel{\text{law}}{\to} \mathcal{N}(0, \sigma^2);$
- (2) $E[I_q(f_n)^4] \stackrel{\text{law}}{\to} 3\sigma^4$;
- (3) $||f_n \otimes_r f_n||_{L^2(\mathbb{R}^{2q-2r}_+)} \to 0$ for each $r = 1, \dots, q-1$, where $f_n \otimes_r f_n$ is the function of $L^2(\mathbb{R}^{2q-2r}_+)$ defined by

$$f_n \otimes_r f_n(x_1, \dots, x_{2q-2r})$$

$$= \int_{\mathbb{R}^r_+} f_n(x_1, \dots, x_{q-r}, y_1, \dots, y_r) f_n(x_{q-r+1}, \dots, x_{2q-2r}, y_1, \dots, y_r) dy_1 \dots dy_r.$$

In other words, Theorem 1.3 states that the convergence in law of a normalized sequence of multiple Wiener-Itô integrals $I_q(f_n)$ towards the Gaussian law $\mathcal{N}(0, \sigma^2)$ is equivalent to convergence of just the fourth moment to $3\sigma^4$. This surprising result has been the starting point of a new line of research, and has quickly led to several applications, extensions and improvement. One of these improvements is the following quantitative bound associated to Theorem 1.3 that we shall prove in Section 5 by combining Stein's method with the Malliavin calculus.

Theorem 1.4 (Nourdin, Peccati, 2009; see [25]) If $q \ge 2$ is an integer and f is a symmetric element of $L^2(\mathbb{R}^q_+)$ satisfying $E[I_q(f)^2] = q! ||f||^2_{L^2(\mathbb{R}^q_+)} = 1$, then

$$\sup_{A\subset\mathcal{B}(\mathbb{R})}\left|P[I_q(f)\in A]-\frac{1}{\sqrt{2\pi}}\int_A e^{-x^2/2}dx\right|\leqslant 2\sqrt{\frac{q-1}{3q}}\sqrt{\left|E[I_q(f)^4]-3\right|}.$$

Let us go back to (i), that is, to the proof of (1.5) for $\varphi = H_q$. Recall that the sequence $\{e_k\}$ has be chosen for (1.9) to hold. Using (1.10) (see also (1.11)), we can write $V_n = I_q(f_n)$, with

$$f_n = \frac{1}{\sqrt{n}} \sum_{k=1}^n e_k^{\otimes q}.$$

We already showed that $E[V_n^2] \to \sigma^2$ as $n \to \infty$. So, according to Theorem 1.3, to get (i) it remains to check that $||f_n \otimes_r f_n||_{L^2(\mathbb{R}^{2q-2r}_+)} \to 0$ for any $r = 1, \ldots, q-1$. We have

$$f_n \otimes_r f_n = \frac{1}{n} \sum_{k,l=1}^n e_k^{\otimes q} \otimes_r e_l^{\otimes q} = \frac{1}{n} \sum_{k,l=1}^n \langle e_k, e_l \rangle_{L^2(\mathbb{R}_+)}^r e_k^{\otimes q-r} \otimes e_l^{\otimes q-r}$$
$$= \frac{1}{n} \sum_{k,l=1}^n \rho(k-l)^r e_k^{\otimes q-r} \otimes e_l^{\otimes q-r},$$

implying in turn

$$||f_{n} \otimes_{r} f_{n}||_{L^{2}(\mathbb{R}^{2q-2r}_{+})}^{2}$$

$$= \frac{1}{n^{2}} \sum_{i,j,k,l=1}^{n} \rho(i-j)^{r} \rho(k-l)^{r} \langle e_{i}^{\otimes q-r} \otimes e_{j}^{\otimes q-r}, e_{k}^{\otimes q-r} \otimes e_{l}^{\otimes q-r} \rangle_{L^{2}(\mathbb{R}^{2q-2r}_{+})}^{2q-2r}$$

$$= \frac{1}{n^{2}} \sum_{i,j,k,l=1}^{n} \rho(i-j)^{r} \rho(k-l)^{r} \rho(i-k)^{q-r} \rho(j-l)^{q-r}.$$

Observe that $|\rho(k-l)|^r |\rho(i-k)|^{q-r} \leq |\rho(k-l)|^q + |\rho(i-k)|^q$. This, together with other obvious manipulations, leads to the bound

$$||f_{n} \otimes_{r} f_{n}||_{L^{2}(\mathbb{R}^{2q-2r}_{+})}^{2} \leq \frac{2}{n} \sum_{k \in \mathbb{Z}} |\rho(k)|^{q} \sum_{|i| < n} |\rho(i)|^{r} \sum_{|j| < n} |\rho(j)|^{q-r}$$

$$\leq \frac{2}{n} \sum_{k \in \mathbb{Z}} |\rho(k)|^{d} \sum_{|i| < n} |\rho(i)|^{r} \sum_{|j| < n} |\rho(j)|^{q-r}$$

$$= 2 \sum_{k \in \mathbb{Z}} |\rho(k)|^{d} \times n^{-\frac{q-r}{q}} \sum_{|i| < n} |\rho(i)|^{r} \times n^{-\frac{r}{q}} \sum_{|j| < n} |\rho(j)|^{q-r}.$$

Thus, to get that $||f_n \otimes_r f_n||_{L^2(\mathbb{R}^{2q-2r}_+)} \to 0$ for any $r = 1, \ldots, q-1$, it suffices to show that

$$s_n(r) := n^{-\frac{q-r}{q}} \sum_{|i| < n} |\rho(i)|^r \to 0 \text{ for any } r = 1, \dots, q-1.$$

Let r = 1, ..., q - 1. Fix $\delta \in (0, 1)$ (to be chosen later) and let us decompose $s_n(r)$ into

$$s_n(r) = n^{-\frac{q-r}{q}} \sum_{|i| < [n\delta]} |\rho(i)|^r + n^{-\frac{q-r}{q}} \sum_{[n\delta] \le |i| < n} |\rho(i)|^r =: s_{1,n}(\delta, r) + s_{2,n}(\delta, r).$$

Using Hölder inequality, we get that

$$s_{1,n}(\delta,r) \leqslant n^{-\frac{q-r}{r}} \left(\sum_{|i| < [n\delta]} |\rho(i)|^q \right)^{r/q} \left(1 + 2[n\delta] \right)^{\frac{q-r}{q}} \leqslant \text{cst} \times \delta^{1-r/q},$$

as well as

$$s_{2,n}(\delta,r) \leqslant n^{-\frac{q-r}{r}} \left(\sum_{[n\delta] \leqslant |i| < n} |\rho(i)|^q \right)^{r/q} (2n)^{\frac{q-r}{q}} \leqslant \operatorname{cst} \times \left(\sum_{|i| \geqslant [n\delta]} |\rho(i)|^q \right)^{r/q}.$$

Since 1 - r/q > 0, it is a routine exercise (details are left to the reader) to deduce that $s_n(r) \to 0$ as $n \to \infty$. Since this is true for any $r = 1, \ldots, q - 1$, this concludes the proof of (i).

It remains to show (ii). We shall use the multivariate counterpart of Theorem 1.3, which was obtained shortly afterwards by Peccati and Tudor. Since only a weak version (where all the involved multiple Wiener-Itô integrals have different orders) is needed here, we state the result of Peccati and Tudor only in this situation. We refer to Section 6 for a proof of a more general version.

Theorem 1.5 (Peccati, Tudor, 2005; see [44]) Consider l integers $q_1, \ldots, q_l \ge 1$, with $l \ge 2$. Assume that all the q_i 's are pairwise different. For each $i = 1, \ldots, l$, let $\{f_n^i\}_{n \ge 1}$ be a sequence of symmetric functions of $L^2(\mathbb{R}^{q_i}_+)$ satisfying $E[I_{q_i}(f_n^i)^2] = q_i! ||f_n^i||_{L^2(\mathbb{R}^{q_i}_+)}^2 \to \sigma_i^2$ as $n \to \infty$ for some $\sigma_i > 0$. Then, the following two assertions are equivalent as $n \to \infty$:

(1)
$$I_{q_i}(f_n^i) \stackrel{\text{law}}{\to} \mathcal{N}(0, \sigma_i^2)$$
 for all $i = 1, \dots, l$;

(2)
$$(I_{q_1}(f_n^1), \dots, I_{q_l}(f_n^l)) \stackrel{\text{law}}{\to} \mathcal{N}(0, \text{diag}(\sigma_1^2, \dots, \sigma_l^2)).$$

In other words, Theorem 1.5 proves the surprising fact that, for a sequence of vectors of multiple Wiener-Itô integrals, componentwise convergence to Gaussian always implies joint convergence. We shall combine Theorem 1.5 with (i) to prove (ii). Let φ have the form of a real polynomial. In particular, it admits a decomposition of the type $\varphi = \sum_{q=d}^{N} a_q H_q$ for some finite integer $N \geqslant d$. Together with (i), Theorem 1.5 yields that

$$\left(\frac{1}{\sqrt{n}}\sum_{k=1}^n H_d(X_k), \dots, \frac{1}{\sqrt{n}}\sum_{k=1}^n H_N(X_k)\right) \stackrel{\text{law}}{\to} \mathcal{N}\left(0, \operatorname{diag}(\sigma_d^2, \dots, \sigma_N^2)\right),$$

where $\sigma_q^2 = q! \sum_{k \in \mathbb{Z}} \rho(k)^q$, $q = d, \dots, N$. We deduce that

$$V_n = \frac{1}{\sqrt{n}} \sum_{q=d}^N a_q \sum_{k=1}^n H_q(X_k) \xrightarrow{\text{law}} \mathcal{N} \left(0, \sum_{q=d}^N a_q^2 q! \sum_{k \in \mathbb{Z}} \rho(k)^q \right),$$

which is the desired conclusion in (ii) and conclude the proof of Theorem 1.1.

To go further. In [30], one associates quantitative bounds to Theorem 1.1 by using a similar approach.

2 Universality of Wiener chaos

Before developing the material which will be necessary for the proof of the Fourth Moment Theorem 1.3 (as well as other related results), to motivate the reader let us study yet another consequence of this beautiful result.

For any sequence X_1, X_2, \ldots of i.i.d. random variables with mean 0 and variance 1, the central limit theorem asserts that $V_n = (X_1 + \ldots + X_n)/\sqrt{n} \stackrel{\text{law}}{\to} \mathcal{N}(0,1)$ as $n \to \infty$. It is a particular

instance of what is commonly referred to as a 'universality phenomenon' in probability. Indeed, we observe that the limit of the sequence V_n does not rely on the specific law of the X_i 's, but only of the fact that its first two moments are 0 and 1 respectively.

Another example that exhibits a universality phenomenon is given by the Wigner's theorem in the random matrices theory. More precisely, let $\{X_{ij}\}_{j>i\geqslant 1}$ and $\{X_{ii}/\sqrt{2}\}_{i\geqslant 1}$ be two independent families composed of i.i.d. random variables with mean 0, variance 1, and all the moments. Set $X_{ji} = X_{ij}$ and consider the $n \times n$ random matrix $M_n = (\frac{X_{ij}}{\sqrt{n}})_{1\leqslant i,j\leqslant n}$. The matrix M_n being symmetric, its eigenvalues $\lambda_{1,n},\ldots,\lambda_{n,n}$ (possibly repeated with multiplicity) belong to \mathbb{R} . Wigner's theorem then asserts that the spectral measure of M_n , that is, the random probability measure defined as $\frac{1}{n}\sum_{k=1}^n \delta_{\lambda_{k,n}}$, always converge almost surely to the semicircular law $\frac{1}{2\pi}\sqrt{4-x^2}\mathbf{1}_{[-2,2]}(x)dx$, whatever the exact distribution of the entries of M_n are.

In this section, our aim is to prove yet another universality phenomenon, which is in the spirit of the two afore-mentioned results. To do so, we need to introduce the following two blocks of basic ingredients:

- (i) Three sequences $\mathbf{X} = (X_1, X_2, \ldots)$, $\mathbf{G} = (G_1, G_2, \ldots)$ and $\mathbf{E} = (\varepsilon_1, \varepsilon_2, \ldots)$ of i.i.d. random variables, all with mean 0, variance 1 and finite fourth moment. We are more specific with \mathbf{G} and \mathbf{E} , by assuming further that $G_1 \sim \mathcal{N}(0,1)$ and $P(\varepsilon_1 = 1) = P(\varepsilon_1 = -1) = 1/2$. (As we will see, \mathbf{E} will actually play no role in the statement of Theorem 2.1; we will however use it to build a interesting counterexample.)
- (ii) A fixed integer $d \ge 1$ as well as a sequence $g_n : \{1, \ldots, n\}^d \to \mathbb{R}, n \ge 1$ of real functions, each g_n satisfying in addition: for all $i_1, \ldots, i_d = 1, \ldots, n$,
 - (a) $g_n(i_1,\ldots,i_d)=g_n(i_{\sigma(1)},\ldots,i_{\sigma(d)})$ for all permutation $\sigma\in\mathfrak{S}_d$;
 - (b) $g_n(i_1, \ldots, i_d) = 0$ whenever $i_k = i_l$ for some $k \neq l$;
 - (c) $d! \sum_{i_1, \dots, i_d=1}^n g_n(i_1, \dots, i_d)^2 = 1.$

(Of course, conditions (a) and (b) are becoming immaterial when d=1.) If $\mathbf{x}=(x_1,x_2,\ldots)$ is a given real sequence, we also set

$$Q_d(g_n, \mathbf{x}) = \sum_{i_1, \dots, i_d=1}^n g_n(i_1, \dots, i_d) x_{i_1} \dots x_{i_d}.$$

Using (b) and (c), it is straightforward to check that, for any $n \ge 1$, we have $E[Q_d(g_n, \mathbf{X})] = 0$ and $E[Q_d(g_n, \mathbf{X})^2] = 1$.

We are now in position to state our new universality phenomenon.

Theorem 2.1 (Nourdin, Peccati, Reinert, 2010; see [31]) Assume that $d \ge 2$. Then, as $n \to \infty$, the following two assertions are equivalent:

- (α) $Q_d(g_n, \mathbf{G}) \stackrel{\text{law}}{\to} \mathcal{N}(0, 1);$
- (β) $Q_d(g_n, \mathbf{X}) \stackrel{\text{law}}{\to} \mathcal{N}(0, 1)$ for <u>any</u> sequence \mathbf{X} as given in (i).

Before proving Theorem 2.1, let us address some comments.

Remark 2.2 1. In reality, the universality phenomenon in Theorem 2.1 is a bit more subtle than in the CLT or in the Wigner's theorem. To illustrate what we have in mind, let us consider an explicit situation (in the case d = 2). Let $g_n : \{1, \ldots, n\}^2 \to \mathbb{R}$ be the function given by

$$g_n(i,j) = \frac{1}{2\sqrt{n-1}} \mathbf{1}_{\{i=1,j\geqslant 2 \text{ or } j=1,i\geqslant 2\}}.$$

It is easy to check that g_n satisfies the three assumptions (a)-(b)-(c) and also that

$$Q_2(g_n, \mathbf{x}) = x_1 \times \frac{1}{\sqrt{n-1}} \sum_{k=2}^n x_k.$$

The classical CLT then implies that $Q_2(g_n, \mathbf{G}) \stackrel{\text{law}}{\to} G_1G_2$ and $Q_2(g_n, \mathbf{E}) \stackrel{\text{law}}{\to} \varepsilon_1G_2$. Moreover, it is a classical and easy exercise to check that ε_1G_2 is $\mathcal{N}(0,1)$ distributed. Thus, what we just showed is that, although $Q_2(g_n, \mathbf{E}) \stackrel{\text{law}}{\to} \mathcal{N}(0,1)$ as $n \to \infty$, the assertion (β) in Theorem 2.1 fails when choosing $\mathbf{X} = \mathbf{G}$ (indeed, the product of two independent $\mathcal{N}(0,1)$ random variables is not gaussian). This means that, in Theorem 2.1, we cannot replace the sequence \mathbf{G} in (α) by any other sequence (at least, not by \mathbf{E} !).

2. Theorem 2.1 is completely false when d=1. For an explicit counterexample, consider for instance $g_n(i) = \mathbf{1}_{\{i=1\}}, i = 1, \ldots, n$. We then have $Q_1(g_n, \mathbf{x}) = x_1$. Consequently, the assertion (α) is trivially verified (it is even an equality in law!) but the assertion (β) is never true unless $X_1 \sim \mathcal{N}(0, 1)$.

Proof of Theorem 2.1. Of course, only the implication $(\alpha) \rightarrow (\beta)$ must be shown. Let us divide its proof into three steps.

Step 1. Set $e_i = \mathbf{1}_{[i-1,i]}, i \ge 1$, and let $f_n \in L^2(\mathbb{R}^d_+)$ be the symmetric function defined as

$$f_n = \sum_{i_1, \dots, i_d=1}^n g_n(i_1, \dots, i_d) e_{i_1} \otimes \dots \otimes e_{i_d}.$$

By the very definition of $I_d(f_n)$, we have

$$I_d(f_n) = d! \sum_{i_1, \dots, i_{d-1}}^n g_n(i_1, \dots, i_d) \int_0^\infty dB_{t_1} e_{i_1}(t_1) \int_0^{t_1} dB_{t_2} e_{i_2}(t_2) \dots \int_0^{t_{d-1}} dB_{t_d} e_{i_d}(t_d).$$

Observe that

$$\int_0^\infty dB_{t_1} e_{i_1}(t_1) \int_0^{t_1} dB_{t_2} e_{i_2}(t_2) \dots \int_0^{t_{d-1}} dB_{t_d} e_{i_d}(t_d)$$

is not almost surely zero (if and) only if $i_d \leq i_{d-1} \leq \ldots \leq i_1$. By combining this fact with assumption (b), we deduce that

$$I_{d}(f_{n}) = d! \sum_{1 \leq i_{d} < \dots < i_{1} \leq n} g_{n}(i_{1}, \dots, i_{d}) \int_{0}^{\infty} dB_{t_{1}} e_{i_{1}}(t_{1}) \int_{0}^{t_{1}} dB_{t_{2}} e_{i_{2}}(t_{2}) \dots \int_{0}^{t_{d-1}} dB_{t_{d}} e_{i_{d}}(t_{d})$$

$$= d! \sum_{1 \leq i_{d} < \dots < i_{1} \leq n} g_{n}(i_{1}, \dots, i_{d}) (B_{i_{1}} - B_{i_{1}-1}) \dots (B_{i_{d}} - B_{i_{d}-1})$$

$$= \sum_{i_{1}, \dots, i_{d}=1}^{n} g_{n}(i_{1}, \dots, i_{d}) (B_{i_{1}} - B_{i_{1}-1}) \dots (B_{i_{d}} - B_{i_{d}-1}) \stackrel{\text{law}}{=} Q_{d}(g_{n}, \mathbf{G}).$$

That is, the sequence $Q_d(g_n, \mathbf{G})$ in (α) has actually the form of a multiple Wiener-Itô integral. On the other hand, going back to the definition of $f_n \otimes_{d-1} f_n$ and using that $\langle e_i, e_j \rangle_{L^2(\mathbb{R}_+)} = \delta_{ij}$ (Kronecker symbol), we get

$$f_n \otimes_{d-1} f_n = \sum_{i,j=1}^n \left(\sum_{k_2,\dots,k_d=1}^n g_n(i,k_2,\dots,k_d) g_n(j,k_2,\dots,k_d) \right) e_i \otimes e_j,$$

so that

$$||f_{n} \otimes_{d-1} f_{n}||_{L^{2}(\mathbb{R}^{2}_{+})}^{2}| = \sum_{i,j=1}^{n} \left(\sum_{k_{2},\dots,k_{d}=1}^{n} g_{n}(i,k_{2},\dots,k_{d}) g_{n}(j,k_{2},\dots,k_{d}) \right)^{2}$$

$$\geqslant \sum_{i=1}^{n} \left(\sum_{k_{2},\dots,k_{d}=1}^{n} g_{n}(i,k_{2},\dots,k_{d})^{2} \right)^{2} \text{ (by summing only over } i=j)$$

$$\geqslant \max_{1 \leq i \leq n} \left(\sum_{k_{2},\dots,k_{d}=1}^{n} g_{n}(i,k_{2},\dots,k_{d})^{2} \right)^{2} = \tau_{n}^{2}, \qquad (2.12)$$

where

$$\tau_n := \max_{1 \le i \le n} \sum_{k_2, \dots, k_d = 1}^n g_n(i, k_2, \dots, k_d)^2.$$
(2.13)

Now, assume that (α) holds. By Theorem 1.3 and because $Q_d(g_n, \mathbf{G}) \stackrel{\text{law}}{=} I_d(f_n)$, we have in particular that $||f_n \otimes_{d-1} f_n||_{L^2(\mathbb{R}^2_+)} \to 0$ as $n \to \infty$. Using the inequality (2.12), we deduce that $\tau_n \to 0$ as $n \to \infty$.

<u>Step 2</u>. We claim that the following result (whose proof is detailed in Step 3) allows to conclude the proof of $(\alpha) \to (\beta)$.

Theorem 2.3 (Mossel, O'Donnel, Oleszkiewicz, 2010; see [18]) Let X and G be given as in (i) and let $g_n : \{1, ..., n\}^d \to \mathbb{R}$ be a function satisfying the three conditions (a)-(b)-(c). Set $\gamma = \max\{3, E[X_1^4]\} \geqslant 1$ and let τ_n be the quantity given by (2.13). Then, for all function $\varphi : \mathbb{R} \to \mathbb{R}$ of class C^3 with $\|\varphi'''\|_{\infty} < \infty$, we have

$$|E[\varphi(Q_d(g_n, \mathbf{X}))] - E[\varphi(Q_d(g_n, \mathbf{G}))]| \leq \frac{\gamma}{3} (3 + 2\gamma)^{\frac{3}{2}(d-1)} d^{3/2} \sqrt{d!} \|\varphi'''\|_{\infty} \sqrt{\tau_n}.$$

Indeed, assume that (α) holds. By Step 1, we have that $\tau_n \to 0$ as $n \to \infty$. Next, Theorem 2.3 together with (α) , lead to (β) and therefore conclude the proof of Theorem 2.1.

Step 3: Proof of Theorem 2.3. During the proof, we will need the following auxiliary lemma, which is of independent interest.

Lemma 2.4 (Hypercontractivity) Let $n \ge d \ge 1$, and consider a multilinear polynomial $P \in \mathbb{R}[x_1, \dots, x_n]$ of degree d, that is, P is of the form

$$P(x_1,...,x_n) = \sum_{\substack{S \subset \{1,...,n\} \ |S|=d}} a_S \prod_{i \in S} x_i.$$

Let X be as in (i). Then,

$$E[P(X_1,...,X_n)^4] \le (3 + 2E[X_1^4])^{2d} E[P(X_1,...,X_n)^2]^2.$$
 (2.14)

Proof. The proof is by induction on n. The case n=1 is trivial. Indeed, in this case we have d=1 so that $P(x_1) = ax_1$; the conclusion therefore asserts that (recall that $E[X_1^2] = 1$, implying in turn that $E[X_1^4] \ge E[X_1^2]^2 = 1$)

$$a^4 E[X_1^4] \leqslant a^4 \left(3 + 2E[X_1^4]\right)^2$$

which is evident. Assume now that $n \ge 2$. We can write

$$P(x_1, \dots, x_n) = R(x_1, \dots, x_{n-1}) + x_n S(x_1, \dots, x_{n-1})$$

where $R, S \in \mathbb{R}[x_1, \dots, x_{n-1}]$ are multilinear polynomials of n-1 variables. Observe that R has degree d, while S has degree d-1. Now write $\mathbf{P} = P(X_1, \dots, X_n)$, $\mathbf{R} = R(X_1, \dots, X_{n-1})$, $\mathbf{S} = S(X_1, \dots, X_{n-1})$ and $\alpha = E[X_1^4]$. Clearly, \mathbf{R} and \mathbf{S} are independent of X_n . We have, using $E[X_n] = 0$ and $E[X_n^2] = 1$:

$$E[\mathbf{P}^2] = E[(\mathbf{R} + \mathbf{S}X_n)^2] = E[\mathbf{R}^2] + E[\mathbf{S}^2]$$

 $E[\mathbf{P}^4] = E[(\mathbf{R} + \mathbf{S}X_n)^4] = E[\mathbf{R}^4] + 6E[\mathbf{R}^2\mathbf{S}^2] + 4E[X_n^3]E[\mathbf{R}\mathbf{S}^3] + E[X_n^4]E[\mathbf{S}^4].$

Observe that $E[\mathbf{R}^2\mathbf{S}^2] \leqslant \sqrt{E[\mathbf{R}^4]}\sqrt{E[\mathbf{S}^4]}$ and

$$E[X_n^3]E[\mathbf{R}\mathbf{S}^3] \leqslant \alpha^{\frac{3}{4}} \left(E[\mathbf{R}^4] \right)^{\frac{1}{4}} \left(E[\mathbf{S}^4] \right)^{\frac{3}{4}} \leqslant \alpha \sqrt{E[\mathbf{R}^4]} \sqrt{E[\mathbf{S}^4]} + \alpha E[\mathbf{S}^4],$$

where the last inequality used both $x^{\frac{1}{4}}y^{\frac{3}{4}} \leqslant \sqrt{xy} + y$ (by considering x < y and x > y) and $\alpha^{\frac{3}{4}} \leqslant \alpha$ (because $\alpha \geqslant E[X_n^4] \geqslant E[X_n^2]^2 = 1$). Hence

$$E[\mathbf{P}^4] \leqslant E[\mathbf{R}^4] + 2(3+2\alpha)\sqrt{E[\mathbf{R}^4]}\sqrt{E[\mathbf{S}^4]} + 5\alpha E[\mathbf{S}^4]$$

$$\leqslant E[\mathbf{R}^4] + 2(3+2\alpha)\sqrt{E[\mathbf{R}^4]}\sqrt{E[\mathbf{S}^4]} + (3+2\alpha)^2 E[\mathbf{S}^4]$$

$$= \left(\sqrt{E[\mathbf{R}^4]} + (3+2\alpha)\sqrt{E[\mathbf{S}^4]}\right)^2.$$

By induction, we have $\sqrt{E[\mathbf{R}^4]} \leqslant (3+2\alpha)^d E[\mathbf{R}^2]$ and $\sqrt{E[\mathbf{S}^4]} \leqslant (3+2\alpha)^{d-1} E[\mathbf{S}^2]$. Therefore

$$E[\mathbf{P}^4] \leqslant (3+2\alpha)^{2d} \left(E[\mathbf{R}^2] + E[\mathbf{S}^2] \right)^2 = (3+2\alpha)^{2d} E[\mathbf{P}^2]^2,$$

and the proof of the lemma is concluded.

We are now in position to prove Theorem 2.3. We use the Lindeberg replacement trick. Without loss of generality, we assume that **X** and **G** are stochastically independent. For i = 0, ..., n, let $\mathbf{W}^{(i)} = (G_1, ..., G_i, X_{i+1}, ..., X_n)$. Fix a particular i = 1, ..., n and write

$$U_{i} = \sum_{\substack{1 \leq i_{1}, \dots, i_{d} \leq n \\ i_{1} \neq i, \dots, i_{d} \neq i}} g_{n}(i_{1}, \dots, i_{d}) W_{i_{1}}^{(i)} \dots W_{i_{d}}^{(i)},$$

$$V_{i} = \sum_{\substack{1 \leq i_{1}, \dots, i_{d} \leq n \\ \exists j: i_{j} = i}} g_{n}(i_{1}, \dots, i_{d}) W_{i_{1}}^{(i)} \dots \widehat{W_{i}}^{(i)} \dots W_{i_{d}}^{(i)}$$

$$= d \sum_{i_{2}, \dots, i_{d} = 1}^{n} g_{n}(i, i_{2}, \dots, i_{d}) W_{i_{2}}^{(i)} \dots W_{i_{d}}^{(i)},$$

where $\widehat{W_i^{(i)}}$ means that this particular term is dropped (observe that this notation bears no ambiguity: indeed, since g_n vanishes on diagonals, each string i_1, \ldots, i_d contributing to the definition of V_i contains the symbol i exactly once). For each i, note that U_i and V_i are independent of the variables X_i and G_i , and that

$$Q_d(g_n, \mathbf{W}^{(i-1)}) = U_i + X_i V_i$$
 and $Q_d(g_n, \mathbf{W}^{(i)}) = U_i + G_i V_i$.

By Taylor's theorem, using the independence of X_i from U_i and V_i , we have

$$\left| E\left[\varphi(U_i + X_i V_i)\right] - E\left[\varphi(U_i)\right] - E\left[\varphi'(U_i)V_i\right]E[X_i] - \frac{1}{2}E\left[\varphi''(U_i)V_i^2\right]E[X_i^2] \right| \\
\leqslant \frac{1}{6} \|\varphi'''\|_{\infty} E[|X_i|^3]E[|V_i|^3].$$

Similarly,

$$\left| E\left[\varphi(U_i + G_i V_i)\right] - E\left[\varphi(U_i)\right] - E\left[\varphi'(U_i)V_i\right]E[G_i] - \frac{1}{2}E\left[\varphi''(U_i)V_i^2\right]E[G_i^2] \right| \\
\leqslant \frac{1}{6} \|\varphi'''\|_{\infty} E[|G_i|^3]E[|V_i|^3].$$

Due to the matching moments up to second order on one hand, and using that $E[|X_i|^3] \leq \gamma$ and $E[|G_i|^3] \leq \gamma$ on the other hand, we obtain that

$$\left| E\left[\varphi(Q_d(g_n, \mathbf{W}^{(i-1)}))\right] - E\left[\varphi(Q_d(g_n, \mathbf{W}^{(i)}))\right] \right| = \left| E\left[\varphi(U_i + G_i V_i)\right] - E\left[\varphi(U_i + X_i V_i)\right] \right| \\
\leqslant \frac{\gamma}{3} \|\varphi'''\|_{\infty} E[|V_i|^3].$$

By Lemma 2.4, we have

$$E[|V_i|^3] \leqslant E[V_i^4]^{\frac{3}{4}} \leqslant (3+2\gamma)^{\frac{3}{2}(d-1)} E[V_i^2]^{\frac{3}{2}}.$$

Using the independence between **X** and **G**, the properties of g_n (which is symmetric and vanishes on diagonals) as well as $E[X_i] = E[G_i] = 0$ and $E[X_i^2] = E[G_i^2] = 1$, we get

$$E[V_i^2]^{3/2} = \left(dd! \sum_{i_2,\dots,i_d=1}^n g_n(i,i_2,\dots,i_d)^2\right)^{3/2}$$

$$\leqslant (dd!)^{3/2} \sqrt{\max_{1 \leqslant j \leqslant n} \sum_{j_2,\dots,j_d=1}^n g_n(j,j_2,\dots,j_d)^2} \times \sum_{i_2,\dots,i_d=1}^n g_n(i,i_2,\dots,i_d)^2,$$

implying in turn that

$$\sum_{i=1}^{n} E[V_i^2]^{3/2} \leqslant (dd!)^{3/2} \sqrt{\max_{1 \leqslant j \leqslant n} \sum_{j_2, \dots, j_d = 1}^{n} g_n(j, j_2, \dots, j_{d_k})^2} \times \sum_{i_1, \dots, i_d = 1}^{n} g_n(i_1, i_2, \dots, i_d)^2,$$

$$= d^{3/2} \sqrt{d!} \sqrt{\tau_n}.$$

By collecting the previous bounds, we get

$$|E[\varphi(Q_{d}(g_{n}, \mathbf{X}))] - E[\varphi(Q_{d}(g_{n}, \mathbf{G}))]|$$

$$\leq \sum_{i=1}^{n} \left| E\left[\varphi(Q_{d}(g_{n}, \mathbf{W}^{(i-1)}))\right] - E\left[\varphi(Q_{d}(g_{n}, \mathbf{W}^{(i)}))\right] \right|$$

$$\leq \frac{\gamma}{3} \|\varphi'''\|_{\infty} \sum_{i=1}^{n} E[|V_{i}|^{3}] \leq \frac{\gamma}{3} (3 + 2\gamma)^{\frac{3}{2}(d-1)} \|\varphi'''\|_{\infty} \sum_{i=1}^{n} E[V_{i}^{2}]^{\frac{3}{2}}$$

$$\leq \frac{\gamma}{3} (3 + 2\gamma)^{\frac{3}{2}(d-1)} d^{3/2} \sqrt{d!} \|\varphi'''\|_{\infty} \sqrt{\tau_{n}}.$$

As a final remark, let us observe that Theorem 2.3 contains the CLT as a special case. Indeed, fix d=1 and let $g_n:\{1,\ldots,n\}\to\mathbb{R}$ be the function given by $g_n(i)=\frac{1}{\sqrt{n}}$. We then have $\tau_n=1/n$. It is moreover clear that $Q_1(g_n,\mathbf{G})\sim\mathcal{N}(0,1)$. Then, for any function $\varphi:\mathbb{R}\to\mathbb{R}$ of class \mathcal{C}^3 with $\|\varphi'''\|_{\infty}<\infty$ and any sequence \mathbf{X} as in (i), Theorem 2.3 implies that

$$\left| E\left[\varphi\left(\frac{X_1 + \ldots + X_n}{\sqrt{n}} \right) \right] - \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} \varphi(y) e^{-y^2/2} dy \right| \leqslant \max\{E[X_1^4]/3, 1\} \|\varphi'''\|_{\infty},$$

from which it is straightforward to deduce the CLT.

To go further. In [31], Theorem 2.1 is extended to the case where the target law is the centered Gamma law. In [46], there is a version of Theorem 2.1 in which the sequence **G** is replaced by **P**, a sequence of i.i.d. Poisson random variables. Finally, let us mention that Theorems 2.1 and 2.3 have been extended to the free probability framework (see Section 11) in the reference [12].

3 Stein's method

In this section, we shall introduce some basic features of the so-called Stein's method, which is the first step toward the proof of the Fourth Moment Theorem 1.3. Actually, we will not need the full force of this method, only a basic estimate.

A random variable X is $\mathcal{N}(0,1)$ distributed if and only if $E[e^{itX}] = e^{-t^2/2}$ for all $t \in \mathbb{R}$. This simple fact leads to the idea that a random variable X has a law which is close to $\mathcal{N}(0,1)$ if and only if $E[e^{itX}]$ is approximatively $e^{-t^2/2}$ for all $t \in \mathbb{R}$. This last claim is nothing but the usual criterion for the convergence in law through the use of characteristic functions.

Stein's seminal idea is somehow similar. He noticed in [50] that X is $\mathcal{N}(0,1)$ distributed if and only if E[f'(X)-Xf(X)]=0 for all function f belonging to a sufficiently rich class of functions (for instance, the functions which are \mathcal{C}^1 and whose derivative growths at most polynomially). He then wondered whether a suitable quantitative version of this identity may have fruitful consequences. This is actually the case and, even for specialists (at least for me!), the reason why it works so well remains a bit mysterious. Surprisingly, the simple following statement happens to contain all the elements of Stein's method that are needed for our discussion. (For more details or extensions of the method, one can consult the recent books [8, 29] and the references therein.)

Lemma 3.1 (Stein, 1972; see [50]) Let $N \sim \mathcal{N}(0,1)$ be a standard Gaussian random variable. Let $h : \mathbb{R} \to [0,1]$ be any continuous function. Define $f : \mathbb{R} \to \mathbb{R}$ by

$$f(x) = e^{\frac{x^2}{2}} \int_{-\infty}^{x} (h(a) - E[h(N)]) e^{-\frac{a^2}{2}} da$$
 (3.15)

$$= -e^{\frac{x^2}{2}} \int_x^\infty \left(h(a) - E[h(N)] \right) e^{-\frac{a^2}{2}} da. \tag{3.16}$$

Then f is of class C^1 , and satisfies $|f(x)| \leq \sqrt{\pi/2}$, $|f'(x)| \leq 2$ and

$$f'(x) = xf(x) + h(x) - E[h(N)]$$
(3.17)

for all $x \in \mathbb{R}$.

Proof: The equality between (3.15) and (3.16) comes from

$$0 = E[h(N) - E[h(N)]] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} (h(a) - E[h(N)]) e^{-\frac{a^2}{2}} da.$$

Using (3.16) we have, for $x \ge 0$:

$$|xf(x)| = \left| xe^{\frac{x^2}{2}} \int_x^{+\infty} \left(h(a) - E[h(N)] \right) e^{-\frac{a^2}{2}} da \right|$$

$$\leqslant xe^{\frac{x^2}{2}} \int_x^{+\infty} e^{-\frac{a^2}{2}} da \leqslant e^{\frac{x^2}{2}} \int_x^{+\infty} ae^{-\frac{a^2}{2}} da = 1.$$

Using (3.15) we have, for $x \leq 0$:

$$\begin{split} \left| x f(x) \right| &= \left| x e^{\frac{x^2}{2}} \int_{-\infty}^{x} \left(h(a) - E[h(N)] \right) e^{-\frac{a^2}{2}} da \right| \\ &\leqslant \left| x | e^{\frac{x^2}{2}} \int_{|x|}^{+\infty} e^{-\frac{a^2}{2}} da \leqslant e^{\frac{x^2}{2}} \int_{|x|}^{+\infty} a e^{-\frac{a^2}{2}} da = 1. \end{split}$$

The identity (3.17) is readily checked. We deduce, in particular, that

$$|f'(x)| \le |xf(x)| + |h(x) - E[h(N)]| \le 2$$

for all $x \in \mathbb{R}$. On the other hand, by (3.15)-(3.16), we have, for every $x \in \mathbb{R}$,

$$|f(x)| \leqslant e^{x^2/2} \min \left(\int_{-\infty}^x e^{-y^2/2} dy, \int_x^\infty e^{-y^2/2} dy \right) = e^{x^2/2} \int_{|x|}^\infty e^{-y^2/2} dy \leqslant \sqrt{\frac{\pi}{2}},$$

where the last inequality is obtained by observing that the function $s: \mathbb{R}_+ \to \mathbb{R}$ given by $s(x) = e^{x^2/2} \int_x^{\infty} e^{-y^2/2} dy$ attains its maximum at x = 0 (indeed, we have

$$s'(x) = xe^{x^2/2} \int_x^\infty e^{-y^2/2} dy - 1 \leqslant e^{x^2/2} \int_x^\infty ye^{-y^2/2} dy - 1 = 0$$

so that s is decreasing on \mathbb{R}_+) and that $s(0) = \sqrt{\pi/2}$.

The proof of the lemma is complete.

To illustrate how Stein's method is a powerful approach, we shall use it to prove the celebrated Berry-Esseen theorem. (Our proof is based on an idea introduced by Ho and Chen in [15], see also Bolthausen [3].)

Theorem 3.2 (Berry, Esseen, 1956; see [14]) Let $\mathbf{X} = (X_1, X_2, \ldots)$ be a sequence of i.i.d. random variables with $E[X_1] = 0$, $E[X_1^2] = 1$ and $E[|X_1|^3] < \infty$, and define

$$V_n = \frac{1}{\sqrt{n}} \sum_{k=1}^n X_k, \quad n \geqslant 1,$$

to be the associated sequence of normalized partial sums. Then, for any $n \geqslant 1$, one has

$$\sup_{x \in \mathbb{R}} \left| P(V_n \leqslant x) - \frac{1}{\sqrt{2\pi}} \int_{-\infty}^x e^{-u^2/2} du \right| \leqslant \frac{33 E[|X_1|^3]}{\sqrt{n}}.$$
 (3.18)

Remark 3.3 One may actually show that (3.18) holds with the constant 0.4784 instead of 33. This has been proved by Korolev and Shevtsova [17] in 2010. (They do not use Stein's method.) On the other hand, according to Esseen [14] himself, it is impossible to expect a universal constant smaller than 0.4097.

Proof of (3.18). For each $n \ge 2$, let $C_n > 0$ be the best possible constant satisfying, for all i.i.d. random variables X_1, \ldots, X_n with $E[|X_1|^3] < \infty$, $E[X_1^2] = 1$ and $E[X_1] = 0$, that

$$\sup_{x \in \mathbb{R}} \left| P(V_n \leqslant x) - \frac{1}{\sqrt{2\pi}} \int_{-\infty}^x e^{-u^2/2} du \right| \leqslant \frac{C_n E[|X_1|^3]}{\sqrt{n}}.$$
 (3.19)

As a first (rough) estimation, we first observe that, since X_1 is centered with $E[X_1^2] = 1$, one has $E[|X_1|^3] \ge E[X_1^2]^{\frac{3}{2}} = 1$, so that $C_n \le \sqrt{n}$. This is of course not enough to conclude, since we need to show that $C_n \le 33$.

For any $x \in \mathbb{R}$ and $\varepsilon > 0$, introduce the function

$$h_{x,\varepsilon}(u) = \begin{cases} 1 & \text{if } u \leqslant x - \varepsilon \\ \text{linear} & \text{if } x - \varepsilon < u < x + \varepsilon \\ 0 & \text{if } u \geqslant x + \varepsilon \end{cases}.$$

It is immediately checked that, for all $n \ge 2$, $\varepsilon > 0$ and $x \in \mathbb{R}$, we have

$$E[h_{x-\varepsilon,\varepsilon}(V_n)] \leqslant P(V_n \leqslant x) \leqslant E[h_{x+\varepsilon,\varepsilon}(V_n)].$$

Moreover, for $N \sim \mathcal{N}(0,1)$, $\varepsilon > 0$ and $x \in \mathbb{R}$, we have, using that the density of N is bounded by $\frac{1}{\sqrt{2\pi}}$,

$$E[h_{x+\varepsilon,\varepsilon}(N)] - \frac{4\varepsilon}{\sqrt{2\pi}} \leqslant E[h_{x-\varepsilon,\varepsilon}(N)] \leqslant P(N \leqslant x)$$

$$\leqslant E[h_{x+\varepsilon,\varepsilon}(N)] \leqslant E[h_{x-\varepsilon,\varepsilon}(N)] + \frac{4\varepsilon}{\sqrt{2\pi}}.$$

Therefore, for all $n \ge 2$ and $\varepsilon > 0$, we have

$$\sup_{x \in \mathbb{R}} \left| P(V_n \leqslant x) - \frac{1}{\sqrt{2\pi}} \int_{-\infty}^x e^{-u^2/2} du \right| \leqslant \sup_{x \in \mathbb{R}} \left| E[h_{x,\varepsilon}(V_n)] - E[h_{x,\varepsilon}(N)] \right| + \frac{4\varepsilon}{\sqrt{2\pi}}.$$

Assume for the time being that, for all $\varepsilon > 0$,

$$\sup_{x \in \mathbb{R}} |E[h_{x,\varepsilon}(V_n)] - E[h_{x,\varepsilon}(N)]| \leqslant \frac{6 E[|X_1|^3]}{\sqrt{n}} + \frac{3 C_{n-1} E[|X_1|^3]^2}{\varepsilon n}.$$
 (3.20)

We deduce that, for all $\varepsilon > 0$,

$$\sup_{x \in \mathbb{R}} \left| P(V_n \leqslant x) - \frac{1}{\sqrt{2\pi}} \int_{-\infty}^x e^{-u^2/2} du \right| \leqslant \frac{6 \, E[|X_1|^3]}{\sqrt{n}} + \frac{3 \, C_{n-1} \, E[|X_1|^3]^2}{\varepsilon \, n} + \frac{4\varepsilon}{\sqrt{2\pi}}.$$

By choosing $\varepsilon = \sqrt{\frac{C_{n-1}}{n}} E[|X_1|^3]$, we get that

$$\sup_{x \in \mathbb{R}} \left| P(V_n \leqslant x) - \frac{1}{\sqrt{2\pi}} \int_{-\infty}^x e^{-u^2/2} du \right| \leqslant \frac{E[|X_1|^3]}{\sqrt{n}} \left[6 + \left(3 + \frac{4}{\sqrt{2\pi}} \right) \sqrt{C_{n-1}} \right],$$

so that $C_n \leq 6 + \left(3 + \frac{4}{\sqrt{2\pi}}\right)\sqrt{C_{n-1}}$. It follows by induction that $C_n \leq 33$ (recall that $C_n \leq \sqrt{n}$ so that $C_2 \leq 33$ in particular), which is the desired conclusion.

We shall now use Stein's Lemma 3.1 to prove that (3.20) holds. Fix $x \in \mathbb{R}$ and $\varepsilon > 0$, and let f denote the Stein's solution associated with $h = h_{x,\varepsilon}$, that is, f satisfies (3.15). Observe that h is continuous, and therefore f is \mathcal{C}^1 . Recall from Lemma 3.1 that $||f||_{\infty} \leqslant \sqrt{\frac{\pi}{2}}$ and $||f'||_{\infty} \leqslant 2$. Set also $\widetilde{f}(x) = xf(x)$, $x \in \mathbb{R}$. We then have

$$|\widetilde{f}(x) - \widetilde{f}(y)| = |f(x)(x - y) + (f(x) - f(y))y| \le \left(\sqrt{\frac{\pi}{2}} + 2|y|\right)|x - y|.$$
 (3.21)

On the other hand, set

$$V_n^i = V_n - \frac{X_i}{\sqrt{n}}, \quad i = 1, \dots, n.$$

Observe that V_n^i and X_i are independent by construction. One can thus write

$$\begin{split} &E[h(V_n)] - E[h(N)] = E[f'(V_n) - V_n f(V_n)] \\ &= \sum_{i=1}^n E\left[f'(V_n)\frac{1}{n} - f(V_n)\frac{X_i}{\sqrt{n}}\right] \\ &= \sum_{i=1}^n E\left[f'(V_n)\frac{1}{n} - \left(f(V_n) - f(V_n^i)\right)\frac{X_i}{\sqrt{n}}\right] \text{ because } E[f(V_n^i)X_i] = E[f(V_n^i)]E[X_i] = 0 \\ &= \sum_{i=1}^n E\left[f'(V_n)\frac{1}{n} - f'\left(V_n^i + \theta\frac{X_i}{\sqrt{n}}\right)\frac{X_i^2}{n}\right] \text{ with } \theta \sim \mathscr{U}_{[0,1]} \text{ independent of } X_1, \dots, X_n. \end{split}$$

We have $f'(x) = \widetilde{f}(x) + h(x) - E[h(N)]$, so that

$$E[h(V_n)] - E[h(N)] = \sum_{i=1}^{n} \left(a_i(\tilde{f}) - b_i(\tilde{f}) + a_i(h) - b_i(h) \right), \tag{3.22}$$

where

$$a_i(g) = E[g(V_n) - g(V_n^i)] \frac{1}{n}$$
 and $b_i(g) = E\left[\left(g\left(V_n^i + \theta \frac{X_i}{\sqrt{n}}\right) - g(V_n^i)\right)X_i^2\right] \frac{1}{n}$.

(Here again, we have used that V_n^i and X_i are independent.) Hence, to prove that (3.20) holds true, we must bound four terms.

1st term. One has, using (3.21) as well as $E[|X_1|] \leqslant E[X_1^2]^{\frac{1}{2}} = 1$ and $E[|V_n^i|] \leqslant E[(V_n^i)^2]^{\frac{1}{2}} \leqslant 1$,

$$\left|a_i(\widetilde{f})\right| \leqslant \frac{1}{n\sqrt{n}} \left(E[|X_1|] \sqrt{\frac{\pi}{2}} + 2E[|X_1|] E[|V_n^i|] \right) \leqslant \left(\sqrt{\frac{\pi}{2}} + 2 \right) \frac{1}{n\sqrt{n}}.$$

2nd term. Similarly and because $E[\theta] = \frac{1}{2}$, one has

$$\left| b_i(\widetilde{f}) \right| \;\; \leqslant \;\; \frac{1}{n\sqrt{n}} \left(E[\theta] E[|X_1|^3] \sqrt{\frac{\pi}{2}} + 2 E[\theta] E[|X_1|^3] E[|V_n^i|] \right) \leqslant \left(\frac{1}{2} \sqrt{\frac{\pi}{2}} + 1 \right) \frac{E[|X_1|^3]}{n\sqrt{n}}.$$

 $3rd \ term$. By definition of h, we have

$$h(v) - h(u) = (v - u) \int_0^1 h'(u + s(v - u)) ds = -\frac{v - u}{2\varepsilon} E\left[\mathbf{1}_{[x - \varepsilon, x + \varepsilon]}(u + \widehat{\theta}(v - u))\right],$$

with $\widehat{\theta} \sim \mathscr{U}_{[0,1]}$ independent of θ and X_1, \ldots, X_n , so that

$$\begin{aligned} \left| a_i(h) \right| &\leqslant \frac{1}{2\varepsilon \, n\sqrt{n}} E\left[|X_i| \mathbf{1}_{[x-\varepsilon, x+\varepsilon]} \left(V_n^i + \widehat{\theta} \frac{X_i}{\sqrt{n}} \right) \right] \\ &= \frac{1}{2\varepsilon \, n\sqrt{n}} E\left[|X_i| \, P\left(x - \frac{y}{\sqrt{n}} - \varepsilon \leqslant V_n^i \leqslant x - \frac{y}{\sqrt{n}} + \varepsilon \right) \, \Big|_{y=\widehat{\theta}X_i} \right] \\ &\leqslant \frac{1}{2\varepsilon \, n\sqrt{n}} \, \sup_{y \in \mathbb{R}} \, P\left(x - \frac{y}{\sqrt{n}} - \varepsilon \leqslant V_n^i \leqslant x - \frac{y}{\sqrt{n}} + \varepsilon \right). \end{aligned}$$

We are thus left to bound $P(a \leqslant V_n^i \leqslant b)$ for all $a,b \in \mathbb{R}$ with $a \leqslant b$. For that, set $\widetilde{V}_n^i = \frac{1}{\sqrt{n-1}} \sum_{j \neq i} X_j$, so that $V_n^i = \sqrt{1 - \frac{1}{n}} \, \widetilde{V}_n^i$. We then have, using in particular (3.19) (with n-1 instead of n) and the fact that the standard Gaussian density is bounded by $\frac{1}{\sqrt{2\pi}}$,

$$\begin{split} P(a\leqslant V_n^i\leqslant b) &= P\left(\frac{a}{\sqrt{1-\frac{1}{n}}}\leqslant \widetilde{V}_n^i\leqslant \frac{b}{\sqrt{1-\frac{1}{n}}}\right) \\ &= P\left(\frac{a}{\sqrt{1-\frac{1}{n}}}\leqslant N\leqslant \frac{b}{\sqrt{1-\frac{1}{n}}}\right) \\ &+ P\left(\frac{a}{\sqrt{1-\frac{1}{n}}}\leqslant \widetilde{V}_n^i\leqslant \frac{b}{\sqrt{1-\frac{1}{n}}}\right) - P\left(\frac{a}{\sqrt{1-\frac{1}{n}}}\leqslant N\leqslant \frac{b}{\sqrt{1-\frac{1}{n}}}\right) \\ &\leqslant \frac{b-a}{\sqrt{2\pi}\sqrt{1-\frac{1}{n}}} + \frac{2\,C_{n-1}\,E[|X_1|^3]}{\sqrt{n-1}}. \end{split}$$

We deduce that

$$\left| a_i(h) \right| \leqslant \frac{1}{\sqrt{2\pi}n\sqrt{n-1}} + \frac{C_{n-1}E[|X_1|^3]}{n\sqrt{n}\sqrt{n-1}\varepsilon}.$$

4th term. Similarly, we have

$$\begin{aligned} \left| b_i(h) \right| &= \frac{1}{2n\sqrt{n\varepsilon}} \left| E\left[X_i^3 \, \theta \, \mathbf{1}_{[x-\varepsilon,x+\varepsilon]} \left(V_n^i + \widehat{\theta} \, \theta \, \frac{X_i}{\sqrt{n}} \right) \right] \right| \\ &\leqslant \frac{E[|X_1|^3]}{4n\sqrt{n\varepsilon}} \sup_{y \in \mathbb{R}} P\left(x - \frac{y}{\sqrt{n}} - \varepsilon \leqslant V_n^i \leqslant x - \frac{y}{\sqrt{n}} + \varepsilon \right) \\ &\leqslant \frac{E[|X_1|^3]}{2\sqrt{2\pi}n\sqrt{n-1}} + \frac{C_{n-1} \, E[|X_1|^3]^2}{2n\sqrt{n}\sqrt{n-1} \, \varepsilon}. \end{aligned}$$

Plugging these four estimates into (3.22) and by using the fact that $n \ge 2$ (and therefore $n-1 \ge \frac{n}{2}$) and $E[|X_1|^3] \ge 1$, we deduce the desired conclusion.

To go further. Stein's method has developed considerably since its first appearance in 1972. A comprehensive and very nice reference to go further is the book [8] by Chen, Goldstein and Shao, in which several applications of Stein's method are carefully developed.

4 Malliavin calculus in a nutshell

The second ingredient for the proof of the Fourth Moment Theorem 1.3 is the Malliavin calculus (the first one being the Stein's method, as developed in the previous section). So, let us introduce the reader to the basic operators of Malliavin calculus. For the sake of simplicity and to avoid technicalities that would be useless in this survey, we will only consider the case where the underlying Gaussian process (fixed once for all throughout the sequel) is a classical Brownian motion $B = (B_t)_{t \geq 0}$ defined on some probability space (Ω, \mathcal{F}, P) ; we further assume that the σ -field \mathcal{F} is generated by B.

For a detailed exposition of Malliavin calculus (in a more general context) and for missing proofs, we refer the reader to the textbooks [29, 35].

Dimension one. In this first section, we would like to introduce the basic operators of Malliavin calculus in the simplest situation (where only *one* Gaussian random variable is involved). While easy, it is a sufficiently rich context to encapsulate all the essence of this theory. We first need to recall some useful properties of Hermite polynomials.

Proposition 4.1 The family $(H_q)_{q\in\mathbb{N}}\subset\mathbb{R}[X]$ of Hermite polynomials has the following properties.

- (a) $H'_{q} = qH_{q-1}$ and $H_{q+1} = XH_{q} qH_{q-1}$ for all $q \in \mathbb{N}$.
- (b) The family $\left(\frac{1}{\sqrt{q!}}H_q\right)_{q\in\mathbb{N}}$ is an orthonormal basis of $L^2(\mathbb{R},\frac{1}{\sqrt{2\pi}}e^{-x^2/2}dx)$.
- (c) Let (U, V) be a Gaussian vector with $U, V \sim \mathcal{N}(0, 1)$. Then, for all $k, l \in \mathbb{N}$,

$$E[H_p(U)H_q(V)] = \begin{cases} q! E[UV]^q & \text{if } p = q \\ 0 & \text{otherwise.} \end{cases}$$

Proof. See, e.g., [29, Proposition 1.4.2].

Let $\varphi : \mathbb{R} \to \mathbb{R}$ be an element of $L^2(\mathbb{R}, \frac{1}{\sqrt{2\pi}}e^{-x^2/2}dx)$. Proposition 4.1(b) implies that φ may be expanded (in a unique way) in terms of Hermite polynomials as follows:

$$\varphi = \sum_{q=0}^{\infty} a_q H_q. \tag{4.23}$$

When φ is such that $\sum qq!a_q^2 < \infty$, let us define

$$D\varphi = \sum_{q=0}^{\infty} q a_q H_{q-1}. \tag{4.24}$$

Since the Hermite polynomials satisfy $H'_q = qH_{q-1}$ (Proposition 4.1(a)), observe that

$$D\varphi = \varphi'$$

(in the sense of distributions). Let us now define the Ornstein-Uhlenbeck semigroup $(P_t)_{t\geq 0}$ by

$$P_t \varphi = \sum_{q=0}^{\infty} e^{-qt} a_q H_q. \tag{4.25}$$

Plainly, $P_0 = Id$, $P_t P_s = P_{t+s}$ $(s, t \ge 0)$ and

$$DP_t = e^{-t}P_tD. (4.26)$$

Since $(P_t)_{t\geqslant 0}$ is a semigroup, it admits a generator L defined as

$$L = \frac{d}{dt}|_{t=0}P_t.$$

Of course, for any $t \ge 0$ one has that

$$\frac{d}{dt}P_{t} = \lim_{h \to 0} \frac{P_{t+h} - P_{t}}{h} = \lim_{h \to 0} P_{t} \frac{P_{h} - Id}{h} = P_{t} \lim_{h \to 0} \frac{P_{h} - Id}{h} = P_{t} \frac{d}{dh} \Big|_{h=0} P_{h} = P_{t}L,$$

and, similarly, $\frac{d}{dt}P_t = LP_t$. Moreover, going back to the definition of $(P_t)_{t\geqslant 0}$, it is clear that the domain of L is the set of functions $\varphi \in L^2(\mathbb{R}, \frac{1}{\sqrt{2\pi}}e^{-x^2/2}dx)$ such that $\sum q^2q!a_q^2 < \infty$ and that, in this case,

$$L\varphi = -\sum_{q=0}^{\infty} q a_q H_q.$$

We have the following integration by parts formula.

Proposition 4.2 Let φ be in the domain of L and ψ be in the domain of D. Then

$$\int_{\mathbb{R}} L\varphi(x)\psi(x)\frac{e^{-x^2/2}}{\sqrt{2\pi}}dx = -\int_{\mathbb{R}} D\varphi(x)D\psi(x)\frac{e^{-x^2/2}}{\sqrt{2\pi}}dx. \tag{4.27}$$

Proof. By bilinearity and approximation, it is enough to show (4.27) for $\varphi = H_p$ and $\psi = H_q$, $p, q \in \mathbb{N}$. When $p \neq q$, we have

$$\int_{\mathbb{D}} L\varphi(x)\psi(x) \frac{e^{-x^2/2}}{\sqrt{2\pi}} dx = -p \int_{\mathbb{D}} H_p(x)H_q(x) \frac{e^{-x^2/2}}{\sqrt{2\pi}} dx = 0$$

and

$$\int_{\mathbb{R}} D\varphi(x) D\psi(x) \frac{e^{-x^2/2}}{\sqrt{2\pi}} dx = pq \int_{\mathbb{R}} H_{p-1}(x) H_{q-1}(x) \frac{e^{-x^2/2}}{\sqrt{2\pi}} dx = 0$$

by Proposition 4.1(c), so the desired conclusion holds true in this case. When p = q, we have

$$\int_{\mathbb{R}} L\varphi(x)\psi(x) \frac{e^{-x^2/2}}{\sqrt{2\pi}} dx = -p \int_{\mathbb{R}} H_p(x)^2 \frac{e^{-x^2/2}}{\sqrt{2\pi}} dx = -pp!$$

and

$$\int_{\mathbb{R}} D\varphi(x)D\psi(x)\frac{e^{-x^2/2}}{\sqrt{2\pi}}dx = p^2 \int_{\mathbb{R}} H_{p-1}(x)^2 \frac{e^{-x^2/2}}{\sqrt{2\pi}}dx = p^2(p-1)! = pp!$$

by Proposition 4.1(c), so the desired conclusion holds true in this case as well.

We shall now extend all the previous operators in a situation where, instead of dealing with a random variable of the form $F = \varphi(N)$ (that involves only *one* Gaussian random variable N),

we deal more generally with a random variable F that is measurable with respect to the Brownian motion $(B_t)_{t\geq 0}$.

Wiener integral. For any adapted[†] and square integrable stochastic process $u = (u_t)_{t \geq 0}$, let us denote by $\int_0^\infty u_t dB_t$ its Itô integral. Recall from any standard textbook of stochastic analysis that the Itô integral is a linear functional that takes its values on $L^2(\Omega)$ and has the following basic features, coming mainly from the independence property of the increments of B:

$$E\left[\int_0^\infty u_s dB_s\right] = 0 (4.28)$$

$$E\left[\int_0^\infty u_s dB_s \times \int_0^\infty v_s dB_s\right] = E\left[\int_0^\infty u_s v_s ds\right]. \tag{4.29}$$

In the particular case where $u = f \in L^2(\mathbb{R}_+)$ is deterministic, we say that $\int_0^\infty f(s)dB_s$ is the Wiener integral of f; it is then easy to show that

$$\int_0^\infty f(s)dB_s \sim \mathcal{N}\left(0, \int_0^\infty f^2(s)ds\right). \tag{4.30}$$

Multiple Wiener-Itô integrals and Wiener chaoses. Let $f \in L^2(\mathbb{R}^q)$. Let us see how one could give a 'natural' meaning to the q-fold multiple integral

$$I_q^W(f) = \int_{\mathbb{R}^q}' f(s_1, \dots, s_q) dW_{s_1} \dots dW_{s_q},$$

where the prime indicates that one does not integrate over the hyperdiagonals $t_i = t_j$, $i \neq j$. To achieve this goal, we shall use an iterated Itô integral; the following heuristic 'calculations' are thus natural within this framework:

$$\int_{\mathbb{R}^{q}}^{\prime} f(s_{1}, \dots, s_{q}) dW_{s_{1}} \dots dW_{s_{q}}$$

$$= \sum_{\sigma \in \mathfrak{S}_{q}} \int_{\mathbb{R}^{q}}^{\prime} f(s_{1}, \dots, s_{q}) \mathbf{1}_{\{s_{\sigma(1)} > \dots > s_{\sigma(q)}\}} dW_{s_{1}} \dots dW_{s_{q}}$$

$$= \sum_{\sigma \in \mathfrak{S}_{q}} \int_{-\infty}^{\infty} dW_{s_{\sigma(1)}} \int_{-\infty}^{s_{\sigma(1)}} dW_{s_{\sigma(2)}} \dots \int_{-\infty}^{s_{\sigma(q-1)}} dW_{s_{\sigma(q)}} f(s_{1}, \dots, s_{q})$$

$$= \sum_{\sigma \in \mathfrak{S}_{q}} \int_{-\infty}^{\infty} dW_{t_{1}} \int_{-\infty}^{t_{1}} dW_{t_{2}} \dots \int_{-\infty}^{t_{q-1}} dW_{t_{q}} f(t_{\sigma^{-1}(1)}, \dots, t_{\sigma^{-1}(q)})$$

$$= \sum_{\sigma \in \mathfrak{S}_{q}} \int_{-\infty}^{\infty} dW_{t_{1}} \int_{-\infty}^{t_{1}} dW_{t_{2}} \dots \int_{-\infty}^{t_{q-1}} dW_{t_{q}} f(t_{\sigma(1)}, \dots, t_{\sigma(q)}).$$

$$(4.31)$$

Now, we can use (4.31) as a natural candidate for being $I_q^W(f)$.

Definition 4.3 Let $q \ge 1$ be an integer.

1. When $f \in L^2(\mathbb{R}^q)$, we set

$$I_q^W(f) = \sum_{\sigma \in \mathfrak{S}_q} \int_{-\infty}^{\infty} dW_{t_1} \int_{-\infty}^{t_1} dW_{t_2} \dots \int_{-\infty}^{t_{q-1}} dW_{t_q} f(t_{\sigma(1)}, \dots, t_{\sigma(q)}). \tag{4.32}$$

 $^{^{\}dagger}$ Any adapted process u that is either càdlàg or càglàd admits a progressively measurable version. We will always assume that we are dealing with it.

The random variable $I_q^W(f)$ is called the qth multiple Wiener-Itô integral of f. 2. The set \mathcal{H}_q^W of random variables of the form $I_q^W(f)$, $f \in L^2(\mathbb{R}^q)$, is called the qth Wiener chaos of W. We also use the convention $\mathcal{H}_0^W = \mathbb{R}$.

The following properties are readily checked.

Proposition 4.4 Let $q \ge 1$ be an integer and let $f \in L^2(\mathbb{R}^q)$.

1. If f is symmetric, then

$$I_q^W(f) = q! \int_{-\infty}^{\infty} dW_{t_1} \int_{-\infty}^{t_1} dW_{t_2} \dots \int_{-\infty}^{t_{q-1}} dW_{t_q} f(t_1, \dots, t_q).$$
(4.33)

2. We have

$$I_q^W(f) = I_q^W(\widetilde{f}), \tag{4.34}$$

where \widetilde{f} stands for the symmetrization of f given by

$$\widetilde{f}(t_1, \dots, t_q) = \frac{1}{q!} \sum_{\sigma \in \mathfrak{S}_q} f(t_{\sigma(1)}, \dots, t_{\sigma(q)}). \tag{4.35}$$

3. For any $p,q \ge 1$, $f \in L^2(\mathbb{R}^p)$ and $g \in L^2(\mathbb{R}^q)$,

$$E[I_q^W(f)] = 0 (4.36)$$

$$E[I_p^W(f)I_q^W(g)] = p!\langle \widetilde{f}, \widetilde{g} \rangle_{L^2(\mathbb{R}^p)} \quad \text{if } p = q$$

$$\tag{4.37}$$

$$E[I_p^W(f)I_q^W(g)] = 0 \quad \text{if } p \neq q.$$
 (4.38)

The space $L^2(\Omega)$ can be decomposed into the infinite orthogonal sum of the spaces \mathcal{H}_q . (It is a statement which is analogous to the content of Proposition 4.1(b), and it is precisely here that we need to assume that the σ -field \mathcal{F} is generated by B.) It follows that any square-integrable random variable $F \in L^2(\Omega)$ admits the following chaotic expansion:

$$F = E[F] + \sum_{q=1}^{\infty} I_q(f_q), \tag{4.39}$$

where the functions $f_q \in L^2(\mathbb{R}^q_+)$ are symmetric and uniquely determined by F.

The following result contains a very useful property of multiple Wiener-Itô integrals.

Theorem 4.5 (Nelson, 1973; see [19]) Let $f \in L^2(\mathbb{R}^q_+)$ with $q \geqslant 1$. Then, for all $r \geqslant 2$,

$$E[|I_q(f)|^r] \le [(r-1)^q q!]^{r/2} ||f||_{L^2(\mathbb{R}^q_\perp)}^r < \infty.$$
(4.40)

Proof. See, e.g., [29, Corollary 2.8.14]. (The proof uses the hypercontractivity property of $(P_t)_{t\geqslant 0}$ defined as (4.48).)

Multiple Wiener-Itô integrals are linear by construction. Let us see how they behave with respect to multiplication. To this aim, we need to introduce the concept of *contractions*.

Definition 4.6 When $r \in \{1, ..., p \land q\}$, $f \in L^2(\mathbb{R}^p_+)$ and $g \in L^2(\mathbb{R}^q_+)$, we write $f \otimes_r g$ to indicate the rth contraction of f and g, defined as being the element of $L^2(\mathbb{R}^{p+q-2r}_+)$ given by

$$(f \otimes_{r} g)(t_{1}, \dots, t_{p+q-2r})$$

$$= \int_{\mathbb{R}^{r}_{+}} f(t_{1}, \dots, t_{p-r}, x_{1}, \dots, x_{r}) g(t_{p-r+1}, \dots, t_{p+q-2r}, x_{1}, \dots, x_{r}) dx_{1} \dots dx_{r}.$$

$$(4.41)$$

By convention, we set $f \otimes_0 g = f \otimes g$ as being the tensor product of f and g, that is,

$$(f \otimes_0 g)(t_1, \dots, t_{p+q}) = f(t_1, \dots, t_p)g(t_{p+1}, \dots, t_{p+q}).$$

Observe that

$$||f \otimes_r g||_{L^2(\mathbb{R}^{p+q-2r}_+)} \leqslant ||f||_{L^2(\mathbb{R}^p_+)} ||g||_{L^2(\mathbb{R}^q_+)}, \quad r = 0, \dots, p \land q$$

$$(4.42)$$

by Cauchy-Schwarz, and that $f \otimes_p g = \langle f, g \rangle_{L^2(\mathbb{R}^p_+)}$ when p = q. The next result is the fundamental product formula between multiple Wiener-Itô integrals.

Theorem 4.7 Let $p, q \ge 1$ and let $f \in L^2(\mathbb{R}^p_+)$ and $g \in L^2(\mathbb{R}^q_+)$ be two symmetric functions. Then

$$I_p(f)I_q(g) = \sum_{r=0}^{p \wedge q} r! \binom{p}{r} \binom{q}{r} I_{p+q-2r}(f \widetilde{\otimes}_r g), \tag{4.43}$$

where $f \approx_r g$ stands for the symmetrization of $f \otimes_r g$ (see (4.35)).

Malliavin derivatives. We shall extend the operator D introduced in (4.24). Let $F \in L^2(\Omega)$ and consider the chaotic expansion (4.39).

Definition 4.8 1. When $m \ge 1$ is an integer, we say that F belongs to the Sobolev-Watanabe space $\mathbb{D}^{m,2}$ if

$$\sum_{q=1}^{\infty} q^m q! \|f_q\|_{L^2(\mathbb{R}^q_+)}^2 < \infty. \tag{4.44}$$

2. When (4.44) holds with m = 1, the Malliavin derivative $DF = (D_t F)_{t \ge 0}$ of F is the element of $L^2(\Omega \times \mathbb{R}_+)$ given by

$$D_t F = \sum_{q=1}^{\infty} q I_{q-1} (f_q(\cdot, t)). \tag{4.45}$$

3. More generally, when (4.44) holds with $m \geqslant 2$ we define the mth Malliavin derivative $D^m F = (D_{t_1,...,t_m} F)_{t_1,...,t_m \geqslant 0}$ of F as the element of $L^2(\Omega \times \mathbb{R}^m_+)$ given by

$$D_{t_1,\dots,t_m}F = \sum_{q=m}^{\infty} q(q-1)\dots(q-m+1)I_{q-m}\left(f_q(\cdot,t_1,\dots,t_m)\right). \tag{4.46}$$

It is clear by construction that D is a linear operator. Using (4.37)-(4.38), it is easy to compute the L^2 -norm of DF in terms of the kernels f_q appearing in the chaotic expansion (4.39) of F:

Proposition 4.9 Let $F \in \mathbb{D}^{1,2}$. We have

$$E\left[\|DF\|_{L^{2}(\mathbb{R}_{+})}^{2}\right] = \sum_{q=1}^{\infty} qq! \|f_{q}\|_{L^{2}(\mathbb{R}_{+}^{q})}^{2}.$$

Proof. By (4.45), we can write

$$E\left[\|DF\|_{L^{2}(\mathbb{R}_{+})}^{2}\right] = \int_{\mathbb{R}_{+}} E\left[\left(\sum_{q=1}^{\infty} qI_{q-1}\left(f_{q}(\cdot,t)\right)\right)^{2}\right] dt$$
$$= \sum_{p,q=1}^{\infty} pq \int_{\mathbb{R}_{+}} E\left[I_{p-1}\left(f_{p}(\cdot,t)\right)I_{q-1}\left(f_{q}(\cdot,t)\right)\right] dt.$$

Using (4.38), we deduce that

$$E\left[\|DF\|_{L^{2}(\mathbb{R}_{+})}^{2}\right] = \sum_{q=1}^{\infty} q^{2} \int_{\mathbb{R}_{+}} E\left[I_{q-1}\left(f_{q}(\cdot,t)\right)^{2}\right] dt.$$

Finally, using (4.37), we get that

$$E\left[\|DF\|_{L^{2}(\mathbb{R}_{+})}^{2}\right] = \sum_{q=1}^{\infty} q^{2}(q-1)! \int_{\mathbb{R}_{+}} \|f_{q}(\cdot,t)\|_{L^{2}(\mathbb{R}_{+}^{q-1})}^{2} dt = \sum_{q=1}^{\infty} qq! \|f_{q}\|_{L^{2}(\mathbb{R}_{+}^{q})}^{2}.$$

Let H_q be the qth Hermite polynomial (for some $q \ge 1$) and let $e \in L^2(\mathbb{R}_+)$ have norm 1. Recall (1.10) and Proposition 4.1(a). We deduce that, for any $t \ge 0$,

$$D_t \left(H_q \left(\int_0^\infty e(s) dW_s \right) \right) = D_t (I_q(e^{\otimes q})) = q I_{q-1}(e^{\otimes q-1}) e(t)$$

$$= q H_{q-1} \left(\int_0^\infty e(s) dB_s \right) e(t) = H'_q \left(\int_0^\infty e(s) dB_s \right) D_t \left(\int_0^\infty e(s) dB_s \right).$$

More generally, the Malliavin derivative D verifies the *chain rule*:

Theorem 4.10 Let $\varphi : \mathbb{R} \to \mathbb{R}$ be both of class \mathcal{C}^1 and Lipschitz, and let $F \in \mathbb{D}^{1,2}$. Then, $\varphi(F) \in \mathbb{D}^{1,2}$ and

$$D_t \varphi(F) = \varphi'(F) D_t F, \quad t \geqslant 0. \tag{4.47}$$

Proof. See, e.g., [35, Proposition 1.2.3].

Ornstein-Uhlenbeck semigroup. We now introduce the extension of (4.25) in our infinite-dimensional setting.

Definition 4.11 The Ornstein-Uhlenbeck semigroup is the family of linear operators $(P_t)_{t\geqslant 0}$ defined on $L^2(\Omega)$ by

$$P_t F = \sum_{q=0}^{\infty} e^{-qt} I_q(f_q), \tag{4.48}$$

where the symmetric kernels f_q are given by (4.39).

A crucial property of $(P_t)_{t\geqslant 0}$ is the Mehler's formula, that gives an alternative and often useful representation formula for P_t . To be able to state it, we need to introduce a further notation. Let (B, B') be a two-dimensional Brownian motion defined on the product probability space $(\Omega, \mathcal{F}, \mathbf{P}) = (\Omega \times \Omega', \mathcal{F} \otimes \mathcal{F}', P \times P')$. Let $F \in L^2(\Omega)$. Since F is measurable with respect to the Brownian motion B, we can write $F = \Psi_F(B)$ with Ψ_F a measurable mapping determined $P \circ B^{-1}$ a.s.. As a consequence, for any $t \geqslant 0$ the random variable $\Psi_F(e^{-t}B + \sqrt{1 - e^{-2t}}B')$ is well-defined $P \times P'$ a.s. (note indeed that $e^{-t}B + \sqrt{1 - e^{-2t}}B'$ is again a Brownian motion for any t > 0). We then have the following formula.

Theorem 4.12 (Mehler's formula) For every $F = F(B) \in L^2(\Omega)$ and every $t \ge 0$, we have

$$P_t(F) = E' \left[\Psi_F(e^{-t}B + \sqrt{1 - e^{-2t}}B') \right], \tag{4.49}$$

where E' denotes the expectation with respect to P'.

Proof. By using standard arguments, one may show that the linear span of random variables F having the form $F = \exp \int_0^\infty h(s) dB_s$ with $h \in L^2(\mathbb{R}_+)$ is dense in $L^2(\Omega)$. Therefore, it suffices to consider the case where F has this particular form. On the other hand, we have the following identity, see, e.g., [29, Proposition 1.4.2(vi)]: for all $c, x \in \mathbb{R}$,

$$e^{cx-c^2/2} = \sum_{q=0}^{\infty} \frac{c^q}{q!} H_q(x),$$

with H_q the qth Hermite polynomial. By setting $c = ||h||_{L^2(\mathbb{R}_+)} = ||h||$ and $x = \int_0^\infty \frac{h(s)}{||h||} dB_s$, we deduce that

$$\exp \int_0^\infty h(s)dB_s = e^{\frac{1}{2}||h||^2} \sum_{q=0}^\infty \frac{||h||^q}{q!} H_q\left(\int_0^\infty \frac{h(s)}{||h||} dB_s\right),$$

implying in turn, using (1.10), that

$$\exp \int_0^\infty h(s)dB_s = e^{\frac{1}{2}\|h\|^2} \sum_{q=0}^\infty \frac{1}{q!} I_q(h^{\otimes q}). \tag{4.50}$$

Thus, for $F = \exp \int_0^\infty h(s) dB_s$,

$$P_t F = e^{\frac{1}{2}||h||^2} \sum_{q=0}^{\infty} \frac{e^{-qt}}{q!} I_q(h^{\otimes q}).$$

On the other hand,

$$E'\left[\Psi_{F}(e^{-t}B + \sqrt{1 - e^{-2t}}B')\right] = E'\left[\exp\int_{0}^{\infty} h(s)(e^{-t}dB_{s} + \sqrt{1 - e^{-2t}}dB'_{s})\right]$$

$$= \exp\left(e^{-t}\int_{0}^{\infty} h(s)dB_{s}\right) \exp\left(\frac{1 - e^{-2t}}{2} \|h\|^{2}\right)$$

$$= \exp\left(\frac{1 - e^{-2t}}{2} \|h\|^{2}\right) e^{\frac{e^{-2t}}{2} \|h\|^{2}} \sum_{q=0}^{\infty} \frac{e^{-qt}}{q!} I_{q}\left(h^{\otimes q}\right) \quad \text{by (4.50)}$$

$$= P_{t}F.$$

The desired conclusion follows.

Generator of the Ornstein-Uhlenbeck semigroup. Recall the definition (4.44) of the Sobolev-Watanabe spaces $\mathbb{D}^{m,2}$, $m \ge 1$, and that the symmetric kernels $f_q \in L^2(\mathbb{R}^q_+)$ are uniquely defined through (4.39).

Definition 4.13 1. The generator of the Ornstein-Uhlenbeck semigroup is the linear operator L defined on $\mathbb{D}^{2,2}$ by

$$LF = -\sum_{q=0}^{\infty} qI_q(f_q).$$

2. The pseudo-inverse of L is the linear operator L^{-1} defined on $L^2(\Omega)$ by

$$L^{-1}F = -\sum_{q=1}^{\infty} \frac{1}{q} I_q(f_q).$$

It is obvious that, for any $F \in L^2(\Omega)$, we have that $L^{-1}F \in \mathbb{D}^{2,2}$ and

$$LL^{-1}F = F - E[F]. (4.51)$$

Our terminology for L^{-1} is explained by the identity (4.51). Another crucial property of L is contained in the following result, which is the exact generalization of Proposition 4.2.

Proposition 4.14 Let $F \in \mathbb{D}^{2,2}$ and $G \in \mathbb{D}^{1,2}$. Then

$$E[LF \times G] = -E[\langle DF, DG \rangle_{L^2(\mathbb{R}_+)}]. \tag{4.52}$$

Proof. By bilinearity and approximation, it is enough to show (4.52) for $F = I_p(f)$ and $G = I_q(g)$ with $p, q \ge 1$ and $f \in L^2(\mathbb{R}^p_+)$, $g \in L^2(\mathbb{R}^q_+)$ symmetric. When $p \ne q$, we have

$$E[LF \times G] = -pE[I_p(f)I_q(g)] = 0$$

and

$$E[\langle DF, DG \rangle_{L^2(\mathbb{R}_+)}] = pq \int_0^\infty E[I_{p-1}(f(\cdot, t))I_{q-1}(g(\cdot, t))]dt = 0$$

by (4.38), so the desired conclusion holds true in this case. When p = q, we have

$$E[LF \times G] = -pE[I_p(f)I_p(g)] = 0 = -pp!\langle f, g \rangle_{L^2(\mathbb{R}^p_+)}$$

and

$$E[\langle DF, DG \rangle_{L^{2}(\mathbb{R}_{+})}] = p^{2} \int_{0}^{\infty} E[I_{p-1}(f(\cdot,t))I_{p-1}(g(\cdot,t))]dt$$
$$= p^{2}(p-1)! \int_{0}^{\infty} \langle f(\cdot,t), g(\cdot,t) \rangle_{L^{2}(\mathbb{R}_{+}^{p-1})} dt = pp! \langle f, g \rangle_{L^{2}(\mathbb{R}_{+}^{p})}$$

by (4.37), so the desired conclusion holds true also in this case.

We are now in position to state and prove an integration by parts formula which will play a crucial role in the sequel.

Theorem 4.15 Let $\varphi : \mathbb{R} \to \mathbb{R}$ be both of class C^1 and Lipschitz, and let $F \in \mathbb{D}^{1,2}$ and $G \in L^2(\Omega)$. Then

$$Cov(G, \varphi(F)) = E[\varphi'(F)\langle DF, -DL^{-1}G\rangle_{L^{2}(\mathbb{R}_{+})}]. \tag{4.53}$$

Proof. Using the assumptions made on F and φ , we can write:

$$Cov(G, \varphi(F)) = E[L(L^{-1}G) \times \varphi(F)] \quad (by (4.51))$$

$$= E[\langle D\varphi(F), -DL^{-1}G \rangle_{L^{2}(\mathbb{R}_{+})}] \quad (by (4.52))$$

$$= E[\varphi'(F)\langle D\varphi(F), -DL^{-1}G \rangle_{L^{2}(\mathbb{R}_{+})}] \quad (by (4.47)),$$

which is the announced formula.

Theorem 4.15 admits a useful extension to indicator functions. Before stating and proving it, we recall the following classical result from measure theory.

Proposition 4.16 Let C be a Borel set in \mathbb{R} , assume that $C \subset [-A, A]$ for some A > 0, and let μ be a finite measure on [-A, A]. Then, there exists a sequence (h_n) of continuous functions with support included in [-A, A] and such that $h_n(x) \in [0, 1]$ and $\mathbf{1}_C(x) = \lim_{n \to \infty} h_n(x)$ μ -a.e.

Proof. This is an immediate corollary of Lusin's theorem, see e.g. [48, page 56].

Corollary 4.17 Let C be a Borel set in \mathbb{R} , assume that $C \subset [-A, A]$ for some A > 0, and let $F \in \mathbb{D}^{1,2}$ be such that E[F] = 0. Then

$$E\left[F\int_{-\infty}^{F} \mathbf{1}_{C}(x)dx\right] = E\left[\mathbf{1}_{C}(F)\langle DF, -DL^{-1}F\rangle_{L^{2}(\mathbb{R}_{+})}\right].$$

Proof. Let λ denote the Lebesgue measure and let P_F denote the law of F. By Proposition 4.16 with $\mu = (\lambda + P_F)|_{[-A,A]}$ (that is, μ is the restriction of $\lambda + P_F$ to [-A,A]), there is a sequence (h_n) of continuous functions with support included in [-A,A] and such that $h_n(x) \in [0,1]$ and $\mathbf{1}_C(x) = \lim_{n \to \infty} h_n(x) \mu$ -a.e. In particular, $\mathbf{1}_C(x) = \lim_{n \to \infty} h_n(x) \lambda$ -a.e. and P_F -a.e. By Theorem 4.15, we have moreover that

$$E\left[F\int_{-\infty}^{F} h_n(x)dx\right] = E\left[h_n(F)\langle DF, -DL^{-1}F\rangle_{L^2(\mathbb{R}_+)}\right].$$

The dominated convergence applies and yields the desired conclusion.

As a corollary of both Theorem 4.15 and Corollary 4.17, we shall prove that the law of any multiple Wiener-Itô integral is always absolutely continuous with respect to the Lebesgue measure except, of course, when its kernel is identically zero.

Corollary 4.18 (Shigekawa; see [49]) Let $q \ge 1$ be an integer and let f be a non zero element of $L^2(\mathbb{R}^q_+)$. Then the law of $F = I_q(f)$ is absolutely continuous with respect to the Lebesgue measure.

Proof. Without loss of generality, we further assume that f is symmetric. The proof is by induction on q. When q = 1, the desired property is readily checked because $I_1(f) \sim \mathcal{N}(0, ||f||^2_{L^2(\mathbb{R}_+)})$, see (4.30). Now, let $q \geq 2$ and assume that the statement of Corollary 4.18 holds true for q - 1, that is, assume that the law of $I_{q-1}(g)$ is absolutely continuous for any symmetric element g of $L^2(\mathbb{R}_+^{q-1})$

such that $\|g\|_{L^2(\mathbb{R}^{q-1}_+)} > 0$. Let f be a symmetric element of $L^2(\mathbb{R}^q_+)$ with $\|f\|_{L^2(\mathbb{R}^q_+)} > 0$. Let $h \in L^2(\mathbb{R})$ be such that $\|\int_0^\infty f(\cdot,s)h(s)ds\|_{L^2(\mathbb{R}^{q-1}_+)} \neq 0$. (Such an h necessarily exists because, otherwise, we would have that $f(\cdot,s) = 0$ for almost all $s \geqslant 0$ which, by symmetry, would imply that $f \equiv 0$; this would be in contradiction with our assumption.) Using the induction assumption, we have that the law of

$$\langle DF, h \rangle_{L^2(\mathbb{R}_+)} = \int_0^\infty D_s F h(s) ds = q I_{q-1} \left(\int_0^\infty f(\cdot, s) h(s) ds \right)$$

is absolutely continuous with respect to the Lebesgue measure. In particular,

$$P(\langle DF, h \rangle_{L^2(\mathbb{R}_+)} = 0) = 0,$$

implying in turn, because $\{||DF||_{L^2(\mathbb{R}_+)}=0\}\subset \{\langle DF,h\rangle_{L^2(\mathbb{R}_+)}=0\}$, that

$$P(\|DF\|_{L^2(\mathbb{R}_+)} > 0) = 1. \tag{4.54}$$

Now, let C be a Borel set in \mathbb{R} . Using Corollary 4.17, we can write, for every $n \ge 1$,

$$E\left[\mathbf{1}_{C\cap[-n,n]}(F)\frac{1}{q}\|DF\|_{L^{2}(\mathbb{R}_{+})}^{2}\right] = E\left[\mathbf{1}_{C\cap[-n,n]}(F)\langle DF, -DL^{-1}F\rangle_{L^{2}(\mathbb{R}_{+})}\right]$$
$$= E\left[F\int_{-\infty}^{F}\mathbf{1}_{C\cap[-n,n]}(y)dy\right].$$

Assume that the Lebesgue measure of C is zero. The previous equality implies that

$$E\left[\mathbf{1}_{C\cap[-n,n]}(F)\frac{1}{q}\|DF\|_{L^{2}(\mathbb{R}_{+})}^{2}\right]=0, \quad n\geqslant 1.$$

But (4.54) holds as well, so $P(F \in C \cap [-n, n]) = 0$ for all $n \ge 1$. By monotone convergence, we actually get $P(F \in C) = 0$. This shows that the law of F is absolutely continuous with respect to the Lebesgue measure. The proof of Corollary 4.18 is concluded.

To go further. In the literature, the most quoted reference on Malliavin calculus is the excellent book [35] by Nualart. It contains many applications of this theory (such as the study of the smoothness of probability laws or the anticipating stochastic calculus) and constitutes, as such, an inevitable reference to go further.

5 Stein meets Malliavin

We are now in a position to prove the Fourth Moment Theorem 1.3. As we will see, to do so we will combine the results of Section 3 (Stein's method) with those of Section 4 (Malliavin calculus). This explains the title of the current section!

We start by introducing the distance we shall use to measure the closeness of the laws of random variables.

Definition 5.1 The total variation distance between the laws of two real-valued random variables Y and Z is defined by

$$d_{TV}(Y,Z) = \sup_{C \in \mathcal{B}(\mathbb{R})} |P(Y \in C) - P(Z \in C)|, \tag{5.55}$$

where $\mathscr{B}(\mathbb{R})$ stands for the set of Borel sets in \mathbb{R} .

When $C \in \mathcal{B}(\mathbb{R})$, we have that $P(Y \in C \cap [-n, n]) \to P(Y \in C)$ and $P(Z \in C \cap [-n, n]) \to P(Z \in C)$ as $n \to \infty$ by the monotone convergence theorem. So, without loss we may restrict the supremum in (5.55) to be taken over *bounded* Borel sets, that is,

$$d_{TV}(Y,Z) = \sup_{\substack{C \in \mathcal{B}(\mathbb{R}) \\ C \text{ bounded}}} |P(Y \in C) - P(Z \in C)|. \tag{5.56}$$

We are now ready to derive a bound for the Gaussian approximation of any centered element F belonging to $\mathbb{D}^{1,2}$.

Theorem 5.2 (Nourdin, Peccati, 2009; see [25]) Consider $F \in \mathbb{D}^{1,2}$ with E[F] = 0. Then, with $N \sim \mathcal{N}(0,1)$,

$$d_{TV}(F,N) \leqslant 2E\left[\left|1 - \langle DF, -DL^{-1}F \rangle_{L^{2}(\mathbb{R}_{+})}\right|\right]. \tag{5.57}$$

Proof. Let C be a bounded Borel set in \mathbb{R} . Let A > 0 be such that $C \subset [-A, A]$. Let λ denote the Lebesgue measure and let P_F denote the law of F. By Proposition 4.16 with $\mu = (\lambda + P_F)|_{[-A-,A]}$ (the restriction of $\lambda + P_F$ to [-A,A]), there is a sequence (h_n) of continuous functions such that $h_n(x) \in [0,1]$ and $\mathbf{1}_C(x) = \lim_{n \to \infty} h_n(x)$ μ -a.e. By the dominated convergence theorem, $E[h_n(F)] \to P(F \in C)$ and $E[h_n(N)] \to P(N \in C)$ as $n \to \infty$. On the other hand, using Lemma 3.1 (and denoting by f_n the function associated with h_n) as well as (4.53) we can write, for each n,

$$\begin{aligned} |E[h_n(F)] - E[h_n(N)]| &= |E[f'_n(F)] - E[Ff_n(F)]| \\ &= |E[f'_n(F)(1 - \langle DF, -DL^{-1}F \rangle_{L^2(\mathbb{R}_+)}]] \\ &\leqslant 2E[|1 - \langle DF, -DL^{-1}F \rangle_{L^2(\mathbb{R}_+)}|]. \end{aligned}$$

Letting n goes to infinity yields

$$|P(F \in C) - P(N \in C)| \leq 2E[|1 - \langle DF, -DL^{-1}F \rangle_{L^{2}(\mathbb{R}_{+})}|],$$

which, together with (5.56), implies the desired conclusion.

Wiener chaos and the Fourth Moment Theorem. In this section, we apply Theorem 5.2 to chaotic random variables, that is, to random variables having the specific form of a multiple Wiener-Itô integral. We begin with a technical lemma.

Lemma 5.3 Let $q \ge 1$ be an integer and consider a symmetric function $f \in L^2(\mathbb{R}^q_+)$. Set $F = I_q(f)$ and $\sigma^2 = E[F^2] = q! ||f||_{L^2(\mathbb{R}^q_+)}^2$. The following two identities hold:

$$E\left[\left(\sigma^{2} - \frac{1}{q}\|DF\|_{L^{2}(\mathbb{R}_{+})}^{2}\right)^{2}\right] = \sum_{r=1}^{q-1} \frac{r^{2}}{q^{2}} r!^{2} {r \choose r}^{4} (2q - 2r)! \|f\widetilde{\otimes}_{r}f\|_{L^{2}(\mathbb{R}_{+}^{2q-2r})}^{2}$$

$$(5.58)$$

and

$$E[F^{4}] - 3\sigma^{4} = \frac{3}{q} \sum_{r=1}^{q-1} rr!^{2} {r \choose r}^{4} (2q - 2r)! ||f\widetilde{\otimes}_{r} f||_{L^{2}(\mathbb{R}^{2q-2r}_{+})}^{2}$$

$$(5.59)$$

$$= \sum_{r=1}^{q-1} q!^2 \binom{q}{r}^2 \left\{ \|f \otimes_r f\|_{L^2(\mathbb{R}^{2q-2r}_+)}^2 + \binom{2q-2r}{q-r} \|f \widetilde{\otimes}_r f\|_{L^2(\mathbb{R}^{2q-2r}_+)}^2 \right\}.$$
(5.60)

In particular,

$$E\left[\left(\sigma^{2} - \frac{1}{q}\|DF\|_{L^{2}(\mathbb{R}_{+})}^{2}\right)^{2}\right] \leqslant \frac{q-1}{3q}\left(E[F^{4}] - 3\sigma^{4}\right). \tag{5.61}$$

Proof. For any $t \ge 0$, we have $D_t F = q I_{q-1}(f(\cdot,t))$ so that, using (4.43),

$$\frac{1}{q} \|DF\|_{L^{2}(\mathbb{R}_{+})}^{2} = q \int_{0}^{\infty} I_{q-1}(f(\cdot,t))^{2} dt$$

$$= q \int_{0}^{\infty} \sum_{r=0}^{q-1} r! \binom{q-1}{r}^{2} I_{2q-2-2r}(f(\cdot,t) \widetilde{\otimes}_{r} f(\cdot,t)) dt$$

$$= q \int_{0}^{\infty} \sum_{r=0}^{q-1} r! \binom{q-1}{r}^{2} I_{2q-2-2r}(f(\cdot,t) \otimes_{r} f(\cdot,t)) dt$$

$$= q \sum_{r=0}^{q-1} r! \binom{q-1}{r}^{2} I_{2q-2-2r} \left(\int_{0}^{\infty} f(\cdot,t) \otimes_{r} f(\cdot,t) dt \right)$$

$$= q \sum_{r=0}^{q-1} r! \binom{q-1}{r}^{2} I_{2q-2-2r}(f \otimes_{r+1} f)$$

$$= q \sum_{r=1}^{q} (r-1)! \binom{q-1}{r-1}^{2} I_{2q-2r}(f \otimes_{r} f)$$

$$= q! \|f\|_{L^{2}(\mathbb{R}_{+}^{q})}^{2} + q \sum_{r=1}^{q-1} (r-1)! \binom{q-1}{r-1}^{2} I_{2q-2r}(f \otimes_{r} f). \tag{5.62}$$

Since $E[F^2] = q! ||f||_{L^2(\mathbb{R}^q_+)}^2 = \sigma^2$, the identity (5.58) follows now from (5.62) and the orthogonality properties of multiple Wiener-Itô integrals. Recall the hypercontractivity property (4.40) of multiple Wiener-Itô integrals, and observe the relations $-L^{-1}F = \frac{1}{q}F$ and $D(F^3) = 3F^2DF$. By combining formula (4.53) with an approximation argument (the derivative of $\varphi(x) = x^3$ being not bounded), we infer that

$$E[F^4] = E[F \times F^3] = \frac{3}{q} E[F^2 || DF ||_{L^2(\mathbb{R}_+)}^2].$$
(5.63)

Moreover, the multiplication formula (4.43) yields

$$F^{2} = I_{q}(f)^{2} = \sum_{s=0}^{q} s! \binom{q}{s}^{2} I_{2q-2s}(f \widetilde{\otimes}_{s} f).$$
 (5.64)

By combining this last identity with (5.62) and (5.63), we obtain (5.59) and finally (5.61). It remains to prove (5.60). Let σ be a permutation of $\{1,\ldots,2q\}$ (this fact is written in symbols as $\sigma \in \mathfrak{S}_{2q}$). If $r \in \{0,\ldots,q\}$ denotes the cardinality of $\{\sigma(1),\ldots,\sigma(q)\} \cap \{1,\ldots,q\}$ then it is readily checked that r is also the cardinality of $\{\sigma(q+1),\ldots,\sigma(2q)\} \cap \{q+1,\ldots,2q\}$ and that

$$\int_{\mathbb{R}^{2q}_{+}} f(t_{1}, \dots, t_{q}) f(t_{\sigma(1)}, \dots, t_{\sigma(q)}) f(t_{q+1}, \dots, t_{2q}) \\
\times f(t_{\sigma(q+1)}, \dots, t_{\sigma(2q)}) dt_{1} \dots dt_{2q} \\
= \int_{\mathbb{R}^{2q-2r}_{+}} (f \otimes_{r} f)(x_{1}, \dots, x_{2q-2r})^{2} dx_{1} \dots dx_{2q-2r} \\
= \|f \otimes_{r} f\|_{L^{2}(\mathbb{R}^{2q-2r}_{+})}^{2}.$$
(5.65)

Moreover, for any fixed $r \in \{0, \ldots, q\}$, there are $\binom{q}{r}^2(q!)^2$ permutations $\sigma \in \mathfrak{S}_{2q}$ such that $\#\{\sigma(1), \ldots, \sigma(q)\} \cap \{1, \ldots, q\} = r$. (Indeed, such a permutation is completely determined by the choice of: (a) r distinct elements y_1, \ldots, y_r of $\{1, \ldots, q\}$; (b) q - r distinct elements y_{r+1}, \ldots, y_q of $\{q+1, \ldots, 2q\}$; (c) a bijection between $\{1, \ldots, q\}$ and $\{y_1, \ldots, y_q\}$; (d) a bijection between $\{q+1, \ldots, 2q\}$ and $\{1, \ldots, 2q\} \setminus \{y_1, \ldots, y_q\}$.) Now, observe that the symmetrization of $f \otimes f$ is given by

$$f\widetilde{\otimes} f(t_1,\ldots,t_{2q}) = \frac{1}{(2q)!} \sum_{\sigma \in \mathfrak{S}_{2q}} f(t_{\sigma(1)},\ldots,t_{\sigma(q)}) f(t_{\sigma(q+1)},\ldots,t_{\sigma(2q)}).$$

Therefore,

$$\|f\widetilde{\otimes}f\|_{L^{2}(\mathbb{R}^{2q}_{+})}^{2} = \frac{1}{(2q)!^{2}} \sum_{\sigma,\sigma'\in\mathfrak{S}_{2q}} \int_{\mathbb{R}^{2q}_{+}} f(t_{\sigma(1)},\ldots,t_{\sigma(q)}) f(t_{\sigma(q+1)},\ldots,t_{\sigma(2q)}) \\ \times f(t_{\sigma'(1)},\ldots,t_{\sigma'(q)}) f(t_{\sigma'(q+1)},\ldots,t_{\sigma'(2q)}) dt_{1} \ldots dt_{2q} \\ = \frac{1}{(2q)!} \sum_{\sigma\in\mathfrak{S}_{2q}} \int_{\mathbb{R}^{2q}_{+}} f(t_{1},\ldots,t_{q}) f(t_{q+1},\ldots,t_{2q}) \\ \times f(t_{\sigma(1)},\ldots,t_{\sigma(q)}) f(t_{\sigma(q+1)},\ldots,t_{\sigma(2q)}) dt_{1} \ldots dt_{2q} \\ = \frac{1}{(2q)!} \sum_{r=0}^{q} \sum_{\substack{\sigma\in\mathfrak{S}_{2q} \\ \{\sigma(1),\ldots,\sigma(q)\}\cap\{1,\ldots,q\}=r}} \int_{\mathbb{R}^{2q}_{+}} f(t_{1},\ldots,t_{q}) f(t_{q+1},\ldots,t_{2q}) \\ \times f(t_{\sigma(1)},\ldots,t_{\sigma(q)}) f(t_{\sigma(q+1)},\ldots,t_{\sigma(2q)}) dt_{1} \ldots dt_{2q}.$$

Using (5.65), we deduce that

$$(2q)! \|f\widetilde{\otimes}f\|_{L^{2}(\mathbb{R}^{2q}_{+})}^{2} = 2(q!)^{2} \|f\|_{L^{2}(\mathbb{R}^{q}_{+})}^{4} + (q!)^{2} \sum_{r=1}^{q-1} {q \choose r}^{2} \|f\otimes_{r} f\|_{L^{2}(\mathbb{R}^{2q-2r}_{+})}^{2}.$$

$$(5.66)$$

Using the orthogonality and isometry properties of multiple Wiener-Itô integrals, the identity (5.64) yields

$$E[F^{4}] = \sum_{r=0}^{q} (r!)^{2} {q \choose r}^{4} (2q - 2r)! \|f\widetilde{\otimes}_{r}f\|_{L^{2}(\mathbb{R}^{2q-2r}_{+})}^{2}$$

$$= (2q)! \|f\widetilde{\otimes}f\|_{L^{2}(\mathbb{R}^{2q}_{+})}^{2} + (q!)^{2} \|f\|_{L^{2}(\mathbb{R}^{q}_{+})}^{4}$$

$$+ \sum_{r=1}^{q-1} (r!)^{2} {q \choose r}^{4} (2q - 2r)! \|f\widetilde{\otimes}_{r}f\|_{L^{2}(\mathbb{R}^{2q-2r}_{+})}^{2}.$$

By inserting (5.66) in the previous identity (and because $(q!)^2 ||f||_{L^2(\mathbb{R}^q_+)}^4 = E[F^2]^2 = \sigma^4$), we get (5.60).

As a consequence of Lemma 5.3, we deduce the following bound on the total variation distance for the Gaussian approximation of a normalized multiple Wiener-Itô integral. This is nothing but Theorem 1.4 but we restate it for convenience.

Theorem 5.4 (Nourdin, Peccati, 2009; see [25]) Let $q \ge 1$ be an integer and consider a symmetric function $f \in L^2(\mathbb{R}^q_+)$. Set $F = I_q(f)$, assume that $E[F^2] = 1$, and let $N \sim \mathcal{N}(0,1)$. Then

$$d_{TV}(F, N) \le 2\sqrt{\frac{q-1}{3q}|E[F^4] - 3|}.$$
 (5.67)

Proof. Since $L^{-1}F = -\frac{1}{q}F$, we have $\langle DF, -DL^{-1}F \rangle_{L^2(\mathbb{R}_+)} = \frac{1}{q}\|DF\|_{L^2(\mathbb{R}_+)}^2$. So, we only need to apply Theorem 5.2 and then formula (5.61) to conclude.

The estimate (5.67) allows to deduce the following characterization of CLTs on Wiener chaos. (This is the Fourth Moment Theorem 1.3 of Nualart and Peccati!)

Corollary 5.5 (Nualart, Peccati, 2005; see [37]) Let $q \ge 1$ be an integer and consider a sequence (f_n) of symmetric functions of $L^2(\mathbb{R}^q_+)$. Set $F_n = I_q(f_n)$ and assume that $E[F_n^2] \to \sigma^2 > 0$ as $n \to \infty$. Then, as $n \to \infty$, the following three assertions are equivalent:

- (i) $F_n \stackrel{\text{Law}}{\to} N \sim \mathcal{N}(0, \sigma^2);$
- (ii) $E[F_n^4] \to E[N^4] = 3\sigma^4$;
- (iii) $||f_n \widetilde{\otimes}_r f_n||_{L^2(\mathbb{R}^{2q-2r})} \to 0 \text{ for all } r = 1, \dots, q-1.$
- (iv) $||f_n \otimes_r f_n||_{L^2(\mathbb{R}^{2q-2r}_+)} \to 0 \text{ for all } r = 1, \dots, q-1.$

Proof. Without loss of generality, we may and do assume that $\sigma^2 = 1$ and $E[F_n^2] = 1$ for all n. The implication (ii) \rightarrow (i) is a direct application of Theorem 5.4. The implication (i) \rightarrow (ii) comes from the Continuous Mapping Theorem together with an approximation argument (observe that $\sup_{n\geqslant 1} E[F_n^4] < \infty$ by the hypercontractivity relation (4.40)). The equivalence between (ii) and (iii) is an immediate consequence of (5.59). The implication (iv) \rightarrow (iii) is obvious (as $||f_n \otimes_r f_n|| \le ||f_n \otimes_r f_n||$) whereas the implication (ii) \rightarrow (iv) follows from (5.60).

Quadratic variation of the fractional Brownian motion. In this section, we aim to illustrate Theorem 5.2 in a concrete situation. More precisely, we shall use Theorem 5.2 in order to derive an explicit bound for the second-order approximation of the quadratic variation of a fractional Brownian motion on [0,1]. Let $B^H = (B_t^H)_{t\geqslant 0}$ be a fractional Brownian motion with Hurst index $H \in (0,1)$. This means that B^H is a centered Gaussian process with covariance function given by

$$E[B_t^H B_s^H] = \frac{1}{2} (t^{2H} + s^{2H} - |t - s|^{2H}), \quad s, t \geqslant 0.$$

It is easily checked that B^H is selfsimilar of index H and has stationary incremements. Consider the suitably normalized quadratic variation of B^H , given by

$$F_n = \frac{n^{2H}}{\sigma_n} \sum_{k=0}^{n-1} \left[(B_{(k+1)/n}^H - B_{k/n}^H)^2 - n^{-2H} \right] \stackrel{\text{(law)}}{=} \frac{1}{\sigma_n} \sum_{k=0}^{n-1} \left[(B_{k+1}^H - B_k^H)^2 - 1 \right],$$

where $\sigma_n > 0$ is so that $E[F_n^2] = 1$. We have the following result.

Theorem 5.6 Let $N \sim \mathcal{N}(0,1)$ and assume that $H \leq 3/4$. Then, $\lim_{n\to\infty} \sigma_n^2/n = 2\sum_{r\in\mathbb{Z}} \rho^2(r)$ if $H \in (0,\frac{3}{4})$, with $\rho:\mathbb{Z} \to \mathbb{R}$ given by

$$\rho(r) = \frac{1}{2} (|r+1|^{2H} + |r-1|^{2H} - 2|r|^{2H}), \tag{5.68}$$

and $\lim_{n\to\infty} \sigma_n^2/(n\log n) = \frac{9}{16}$ if $H = \frac{3}{4}$. Moreover, there exists a constant $c_H > 0$ (depending only on H) such that, for every $n \ge 1$,

$$d_{TV}(F_n, N) \leqslant c_H \times \begin{cases} \frac{1}{\sqrt{n}} & \text{if } H \in (0, \frac{5}{8}) \\ \frac{(\log n)^{3/2}}{\sqrt{n}} & \text{if } H = \frac{5}{8} \\ n^{4H-3} & \text{if } H \in (\frac{5}{8}, \frac{3}{4}) \end{cases}$$

$$(5.69)$$

In order to show Theorem 5.6, we will need the following ancillary result.

Lemma 5.7 1. For any $r \in \mathbb{Z}$, let $\rho(r)$ be defined by (5.68). If $H \neq \frac{1}{2}$, one has $\rho(r) \sim H(2H-1)|r|^{2H-2}$ as $|r| \to \infty$. If $H = \frac{1}{2}$ and $|r| \ge 1$, one has $\rho(r) = 0$. Consequently, $\sum_{r \in \mathbb{Z}} \rho^2(r) < \infty$ if and only if H < 3/4.

2. For all $\alpha > -1$, we have $\sum_{r=1}^{n-1} r^{\alpha} \sim \frac{n^{\alpha+1}}{\alpha+1}$ as $n \to \infty$.

Proof. 1. The sequence ρ is symmetric, that is, one has $\rho(n) = \rho(-n)$. When $r \to \infty$,

$$\rho(r) = H(2H - 1)r^{2H-2} + o(r^{2H-2}).$$

Using the usual criterion for convergence of Riemann sums, we deduce that $\sum_{r\in\mathbb{Z}}\rho^2(r)<\infty$ if and only if 4H-4<-1 if and only if $H<\frac{3}{4}$.

2. For $\alpha > -1$, we have:

$$\frac{1}{n} \sum_{r=1}^{n} \left(\frac{r}{n}\right)^{\alpha} \to \int_{0}^{1} x^{\alpha} dx = \frac{1}{\alpha + 1} \quad \text{as } n \to \infty.$$

We deduce that $\sum_{r=1}^{n} r^{\alpha} \sim \frac{n^{\alpha+1}}{\alpha+1}$ as $n \to \infty$.

We are now in position to prove Theorem 5.6.

Proof of Theorem 5.6. Without loss of generality, we will rather use the second expression of F_n :

$$F_n = \frac{1}{\sigma_n} \sum_{k=0}^{n-1} \left[(B_{k+1}^H - B_k^H)^2 - 1 \right].$$

Consider the linear span \mathcal{H} of $(B_k^H)_{k\in\mathbb{N}}$, that is, \mathcal{H} is the closed linear subspace of $L^2(\Omega)$ generated by $(B_k^H)_{k\in\mathbb{N}}$. It is a real separable Hilbert space and, consequently, there exists an isometry $\Phi: \mathcal{H} \to L^2(\mathbb{R}_+)$. For any $k \in \mathbb{N}$, set $e_k = \Phi(B_{k+1}^H - B_k^H)$; we then have, for all $k, l \in \mathbb{N}$,

$$\int_0^\infty e_k(s)e_l(s)ds = E[(B_{k+1}^H - B_k^H)(B_{l+1}^H - B_l^H)] = \rho(k-l)$$
(5.70)

with ρ given by (5.68). Therefore,

$$\{B_{k+1}^H - B_k^H : k \in \mathbb{N}\} \stackrel{\text{law}}{=} \left\{ \int_0^\infty e_k(s) dB_s : k \in \mathbb{N} \right\} = \{I_1(e_k) : k \in \mathbb{N}\},$$

where B is a Brownian motion and $I_p(\cdot)$, $p \ge 1$, stands for the pth multiple Wiener-Itô integral associated to B. As a consequence, we can replace F_n by Consequently we can, without loss of generality, replace F_n by

$$F_n = \frac{1}{\sigma_n} \sum_{k=0}^{n-1} \left[(I_1(e_k))^2 - 1 \right].$$

Now, using the multiplication formula (4.43), we deduce that

$$F_n = I_2(f_n)$$
, with $f_n = \frac{1}{\sigma_n} \sum_{k=0}^{n-1} e_k \otimes e_k$.

Let us compute the exact value of σ_n . By the isometry formula (4.37) we can write

$$E[F_n^2] = 2\|f_n\|_{L^2(\mathbb{R}^2_+)}^2 = \frac{2}{\sigma_n^2} \sum_{k,l=0}^{n-1} \langle e_k \otimes e_k, e_l \otimes e_l \rangle_{L^2(\mathbb{R}^2_+)}^2 = \frac{2}{\sigma_n^2} \sum_{k,l=0}^{n-1} \langle e_k, e_l \rangle_{L^2(\mathbb{R}_+)}^2$$
$$= \frac{2}{\sigma_n^2} \sum_{k,l=0}^{n-1} \rho^2(k-l).$$

That is,

$$\sigma_n^2 = 2\sum_{k,l=0}^{n-1} \rho^2(k-l) = 2\sum_{|r| < n} (n-|r|)\rho^2(r).$$

Assume that $H < \frac{3}{4}$ and write

$$\frac{\sigma_n^2}{n} = 2\sum_{r \in \mathbb{Z}} \rho^2(r) \left(1 - \frac{|r|}{n}\right) \mathbf{1}_{\{|r| < n\}}.$$

Since $\sum_{r\in\mathbb{Z}}\rho^2(r)<\infty$ by Lemma 5.7, we obtain by dominated convergence that, when $H<\frac{3}{4}$,

$$\lim_{n \to \infty} \frac{\sigma_n^2}{n} = 2 \sum_{r \in \mathbb{Z}} \rho^2(r). \tag{5.71}$$

Assume now that $H = \frac{3}{4}$. We then have $\rho^2(r) \sim \frac{9}{64|r|}$ as $|r| \to \infty$, implying in turn

$$n \sum_{|r| < n} \rho^2(r) \sim \frac{9n}{64} \sum_{0 < |r| < n} \frac{1}{|r|} \sim \frac{9n \log n}{32}$$

and

$$\sum_{|r| < n} |r| \rho^2(r) \sim \frac{9}{64} \sum_{|r| < n} 1 \sim \frac{9n}{32}$$

as $n \to \infty$. Hence, when $H = \frac{3}{4}$,

$$\lim_{n \to \infty} \frac{\sigma_n^2}{n \log n} = \frac{9}{16}.\tag{5.72}$$

On the other hand, recall that the convolution of two sequences $\{u(n)\}_{n\in\mathbb{Z}}$ and $\{v(n)\}_{n\in\mathbb{Z}}$ is the sequence u*v defined as $(u*v)(j) = \sum_{n\in\mathbb{Z}} u(n)v(j-n)$, and observe that $(u*v)(l-i) = \sum_{k\in\mathbb{Z}} u(k-l)v(k-i)$ whenever u(n) = u(-n) and v(n) = v(-n) for all $n\in\mathbb{Z}$. Set

$$\rho_n(k) = |\rho(k)| \mathbf{1}_{\{|k| \leqslant n-1\}}, \quad k \in \mathbb{Z}, \, n \geqslant 1.$$

We then have (using (5.58) for the first equality, and noticing that $f_n \otimes_1 f_n = f_n \widetilde{\otimes}_1 f_n$),

$$E\left[\left(1 - \frac{1}{2}\|D[I_{2}(f_{n})]\|_{L^{2}(\mathbb{R}_{+})}^{2}\right]\right]$$

$$= 8\|f_{n} \otimes_{1} f_{n}\|_{L^{2}(\mathbb{R}_{+}^{2})}^{2} = \frac{8}{\sigma_{n}^{4}} \sum_{i,j,k,l=0}^{n-1} \rho(k-l)\rho(i-j)\rho(k-i)\rho(l-j)$$

$$\leqslant \frac{8}{\sigma_{n}^{4}} \sum_{i,l=0}^{n-1} \sum_{j,k\in\mathbb{Z}} \rho_{n}(k-l)\rho_{n}(i-j)\rho_{n}(k-i)\rho_{n}(l-j)$$

$$= \frac{8}{\sigma_{n}^{4}} \sum_{i,l=0}^{n-1} (\rho_{n} * \rho_{n})(l-i)^{2} \leqslant \frac{8n}{\sigma_{n}^{4}} \sum_{k\in\mathbb{Z}} (\rho_{n} * \rho_{n})(k)^{2} = \frac{8n}{\sigma_{n}^{4}} \|\rho_{n} * \rho_{n}\|_{\ell^{2}(\mathbb{Z})}^{2}.$$

Recall Young's inequality: if $s, p, q \ge 1$ are such that $\frac{1}{p} + \frac{1}{q} = 1 + \frac{1}{s}$, then

$$||u * v||_{\ell^{s}(\mathbb{Z})} \le ||u||_{\ell^{p}(\mathbb{Z})} ||v||_{\ell^{q}(\mathbb{Z})}.$$
 (5.73)

Let us apply (5.73) with $u = v = \rho_n$, s = 2 and $p = \frac{4}{3}$. We get $\|\rho_n * \rho_n\|_{\ell^2(\mathbb{Z})}^2 \leqslant \|\rho_n\|_{\ell^{\frac{4}{3}}(\mathbb{Z})}^4$, so that

$$E\left[\left(1 - \frac{1}{2} \|D[I_2(f_n)]\|_{L^2(\mathbb{R}_+)}^2\right)^2\right] \leqslant \frac{8n}{\sigma_n^4} \left(\sum_{|k| < n} |\rho(k)|^{\frac{4}{3}}\right)^3.$$
 (5.74)

Recall the asymptotic behavior of $\rho(k)$ as $|k| \to \infty$ from Lemma 5.7(1). Hence

$$\sum_{|k| < n} |\rho(k)|^{\frac{4}{3}} = \begin{cases} O(1) & \text{if } H \in (0, \frac{5}{8}) \\ O(\log n) & \text{if } H = \frac{5}{8} \\ O(n^{(8H-5)/3}) & \text{if } H \in (\frac{5}{8}, 1). \end{cases}$$
(5.75)

Assume first that $H < \frac{3}{4}$ and recall (5.71). This, together with (5.74) and (5.75), imply that

$$E\left[\left|1 - \frac{1}{2}\|D[I_{2}(f_{n})]\|_{L^{2}(\mathbb{R}_{+})}^{2}\right|\right] \leqslant \sqrt{E\left[\left(1 - \frac{1}{2}\|D[I_{2}(f_{n})]\|_{L^{2}(\mathbb{R}_{+})}^{2}\right)^{2}\right]}$$

$$\leqslant c_{H} \times \begin{cases} \frac{1}{\sqrt{n}} & \text{if } H \in (0, \frac{5}{8})\\ \frac{(\log n)^{3/2}}{\sqrt{n}} & \text{if } H = \frac{5}{8} \\ n^{4H-3} & \text{if } H \in (\frac{5}{8}, \frac{3}{4}) \end{cases}$$

Therefore, the desired conclusion holds for $H \in (0, \frac{3}{4})$ by applying Theorem 5.2. Assume now that $H = \frac{3}{4}$ and recall (5.72). This, together with (5.74) and (5.75), imply that

$$E\left[\left|1 - \frac{1}{2}\|D[I_2(f_n)]\|_{L^2(\mathbb{R}_+)}^2\right|\right] \leqslant \sqrt{E\left[\left(1 - \frac{1}{2}\|D[I_2(f_n)]\|_{L^2(\mathbb{R}_+)}^2\right)^2\right]} = O(1/\log n),$$

and leads to the desired conclusion for $H = \frac{3}{4}$ as well.

To go further. In [25], one may find a version of Theorem 5.2 where N is replaced by a centered Gamma law (see also [24]). In [1], one associate to Corollary 5.5 an almost sure central limit theorem. In [5], the case where H is bigger than 3/4 in Theorem 5.6 is analyzed.

6 The smart path method

The aim of this section is to prove Theorem 1.5, and even a more general version of it. Following the approach developed in the previous section for the one-dimensional case, a possible way for achieving this goal would have consisted in extending Stein's method to the multivariate setting, so to combine them with the tools of Malliavin calculus. This is indeed the approach developed in [32] and it works well. In this survey, we will actually proceed differently, by using the so-called 'smart path method' (which is a popular method in spin glasses theory, see, e.g., Talagrand [51]).

Let us first illustrate this approach in dimension one. Let $F \in \mathbb{D}^{1,2}$ with E[F] = 0, let $N \sim \mathcal{N}(0,1)$ and let $h : \mathbb{R} \to \mathbb{R}$ be a \mathcal{C}^2 function satisfying $\|\varphi''\|_{\infty} < \infty$. Imagine we want to estimate E[h(F)] - E[h(N)]. Without loss of generality, we may assume that N and F are stochastically independent. We further have:

$$E[h(F)] - E[h(N)] = \int_0^1 \frac{d}{dt} E[h(\sqrt{t}F + \sqrt{1 - t}N)] dt$$

$$= \int_0^1 \left(\frac{1}{2\sqrt{t}} E[h'(\sqrt{t}F + \sqrt{1 - t}N)F] - \frac{1}{2\sqrt{1 - t}} E[h'(\sqrt{t}F + \sqrt{1 - t}N)N] \right) dt.$$

For any $x \in \mathbb{R}$ and $t \in [0, 1]$, Theorem 4.15 implies that

$$E[h'(\sqrt{t}F + \sqrt{1-t}x)F] = \sqrt{t} E[h''(\sqrt{t}F + \sqrt{1-t}x)\langle DF, -DL^{-1}F\rangle_{L^2(\mathbb{R}_+)}]$$

whereas a classical integration by parts yields

$$E[h'(\sqrt{t}x + \sqrt{1-t}N)N] = \sqrt{1-t} E[h''(\sqrt{t}x + \sqrt{1-t}N)].$$

We deduce, since N and F are independent, that

$$E[h(F)] - E[h(N)] = \frac{1}{2} \int_0^1 E[h''(\sqrt{t}x + \sqrt{1 - t}N)(\langle DF, -DL^{-1}F \rangle_{L^2(\mathbb{R}_+)} - 1)]dt, \tag{6.76}$$

implying in turn

$$\left| E[h(F)] - E[h(N)] \right| \leqslant \frac{1}{2} \|h''\|_{\infty} E\left[\left| 1 - \langle DF, -DL^{-1}F \rangle_{L^{2}(\mathbb{R}_{+})} \right| \right], \tag{6.77}$$

compare with (5.57).

It happens that this approach extends easily to the multivariate setting. To see why, we will adopt the following short-hand notation: for every $h : \mathbb{R}^d \to \mathbb{R}$ of class \mathcal{C}^2 , we set

$$||h''||_{\infty} = \max_{i,j=1,\dots,d} \sup_{x \in \mathbb{R}^d} \left| \frac{\partial^2 h}{\partial x_i \partial x_j}(x) \right|.$$

Theorem 6.1 below is a first step towards Theorem 1.5, and is nothing but the multivariate counterpart of (6.76)-(6.77).

Theorem 6.1 Fix $d \ge 2$ and let $F = (F_1, \ldots, F_d)$ be such that $F_i \in \mathbb{D}^{1,2}$ with $E[F_i] = 0$ for any i. Let $C \in \mathcal{M}_d(\mathbb{R})$ be a symmetric and positive matrix, and let N be a centered Gaussian vector with covariance C. Then, for any $h : \mathbb{R}^d \to \mathbb{R}$ belonging to C^2 and such that $||h''||_{\infty} < \infty$, we have

$$\left| E[h(F)] - E[h(N)] \right| \leq \frac{1}{2} \|h''\|_{\infty} \sum_{i,j=1}^{d} E\left[\left| C(i,j) - \langle DF_j, -DL^{-1}F_i \rangle_{L^2(\mathbb{R}_+)} \right| \right]. \tag{6.78}$$

Proof. Without loss of generality, we assume that N is independent of the underlying Brownian motion B. Let h be as in the statement of the theorem. For any $t \in [0,1]$, set $\Psi(t) = E[h(\sqrt{1-t}F+\sqrt{t}N)]$, so that

$$E[h(N)] - E[h(F)] = \Psi(1) - \Psi(0) = \int_0^1 \Psi'(t)dt.$$

We easily see that Ψ is differentiable on (0,1) with

$$\Psi'(t) = \sum_{i=1}^{d} E\left[\frac{\partial h}{\partial x_i} \left(\sqrt{1-t}F + \sqrt{t}N\right) \left(\frac{1}{2\sqrt{t}}N_i - \frac{1}{2\sqrt{1-t}}F_i\right)\right].$$

By integrating by parts, we can write

$$E\left[\frac{\partial h}{\partial x_{i}}\left(\sqrt{1-t}F+\sqrt{t}N\right)N_{i}\right] = E\left\{E\left[\frac{\partial h}{\partial x_{i}}\left(\sqrt{1-t}x+\sqrt{t}N\right)N_{i}\right]_{|x=F}\right\}$$

$$= \sqrt{t}\sum_{j=1}^{d}C(i,j)E\left\{E\left[\frac{\partial^{2}h}{\partial x_{i}\partial x_{j}}\left(\sqrt{1-t}x+\sqrt{t}N\right)\right]_{|x=F}\right\}$$

$$= \sqrt{t}\sum_{j=1}^{d}C(i,j)E\left[\frac{\partial^{2}h}{\partial x_{i}\partial x_{j}}\left(\sqrt{1-t}F+\sqrt{t}N\right)\right].$$

By using Theorem 4.15 in order to perform the integration by parts, we can also write

$$E\left[\frac{\partial h}{\partial x_{i}}\left(\sqrt{1-t}F+\sqrt{t}N\right)F_{i}\right] = E\left\{E\left[\frac{\partial h}{\partial x_{i}}\left(\sqrt{1-t}F+\sqrt{t}x\right)F_{i}\right]_{|x=N}\right\}$$

$$= \sqrt{1-t}\sum_{j=1}^{d}E\left\{E\left[\frac{\partial^{2}h}{\partial x_{i}\partial x_{j}}\left(\sqrt{1-t}F+\sqrt{t}x\right)\left\langle DF_{j},-DL^{-1}F_{i}\right\rangle_{L^{2}(\mathbb{R}_{+})}\right]_{|x=N}\right\}$$

$$= \sqrt{1-t}\sum_{j=1}^{d}E\left[\frac{\partial^{2}h}{\partial x_{i}\partial x_{j}}\left(\sqrt{1-t}F+\sqrt{t}N\right)\left\langle DF_{j},-DL^{-1}F_{i}\right\rangle_{L^{2}(\mathbb{R}_{+})}\right].$$

Hence

$$\Psi'(t) = \frac{1}{2} \sum_{i,j=1}^{d} E\left[\frac{\partial^{2} h}{\partial x_{i} \partial x_{j}} \left(\sqrt{1-t}F + \sqrt{t}N\right) \left(C(i,j) - \langle DF_{j}, -DL^{-1}F_{j} \rangle_{L^{2}(\mathbb{R}_{+})}\right)\right],$$

and the desired conclusion follows.

We are now in position to prove Theorem 1.3. We will actually even show the following more general version.

Theorem 6.2 (Peccati, Tudor, 2005; see [44]) Let $d \ge 2$ and $q_d, \ldots, q_1 \ge 1$ be some fixed integers. Consider vectors

$$F_n = (F_{1,n}, \dots, F_{d,n}) = (I_{q_1}(f_{1,n}), \dots, I_{q_d}(f_{d,n})), \quad n \geqslant 1,$$

with $f_{i,n} \in L^2(\mathbb{R}^{q_i}_+)$ symmetric. Let $C \in \mathcal{M}_d(\mathbb{R})$ be a symmetric and positive matrix, and let N be a centered Gaussian vector with covariance C. Assume that

$$\lim_{n \to \infty} E[F_{i,n}F_{j,n}] = C(i,j), \quad 1 \leqslant i,j \leqslant d. \tag{6.79}$$

Then, as $n \to \infty$, the following two conditions are equivalent:

- (a) F_n converges in law to N;
- (b) for every $1 \leq i \leq d$, $F_{i,n}$ converges in law to $\mathcal{N}(0, C(i, i))$.

Proof. The implication $(a) \Rightarrow (b)$ being trivial, we only concentrate on $(b) \Rightarrow (a)$. So, assume (b) and let us show that (a) holds true. Thanks to (6.78), we are left to show that, for each $i, j = 1, \ldots, d$,

$$\langle DF_{j,n}, -DL^{-1}F_{i,n}\rangle_{L^2(\mathbb{R}_+)} = \frac{1}{q_i}\langle DF_{j,n}, DF_{i,n}\rangle_{L^2(\mathbb{R}_+)} \stackrel{L^2(\Omega)}{\to} C(i,j) \quad \text{as } n \to \infty.$$
 (6.80)

We consider all the possible cases for q_i and q_j .

First case: $q_i = q_j = 1$. We have $\langle DF_{j,n}, DF_{i,n} \rangle_{L^2(\mathbb{R}_+)} = \langle f_{i,n}, f_{j,n} \rangle_{L^2(\mathbb{R}_+)} = E[F_{i,n}F_{j,n}]$. But it is our assumption that $E[F_{i,n}F_{j,n}] \to C(i,j)$ so (6.80) holds true in this case.

Second case: $q_i = 1$ and $q_j \ge 2$ (a similar analysis might be done whenever $q_j = 1$ and $q_i \ge 2$). We have $\langle DF_{j,n}, DF_{i,n} \rangle_{L^2(\mathbb{R}_+)} = \langle f_{i,n}, DF_{j,n} \rangle_{L^2(\mathbb{R}_+)} = I_{q_j-1}(f_{i,n} \otimes_1 f_{j,n})$. We deduce that

$$\begin{split} E[\langle DF_{j,n}, DF_{i,n} \rangle_{L^{2}(\mathbb{R}_{+})}^{2}] &= (q_{j}-1)! \|f_{i,n} \widetilde{\otimes}_{1} f_{j,n}\|_{L^{2}(\mathbb{R}_{+}^{q_{j}-1})}^{2} \leqslant (q_{j}-1)! \|f_{i,n} \otimes_{1} f_{j,n}\|_{L^{2}(\mathbb{R}_{+}^{q_{j}-1})}^{2} \\ &= (q_{j}-1)! \langle f_{i,n} \otimes f_{i,n}, f_{j,n} \otimes_{q_{j}-1} f_{j,n} \rangle_{L^{2}(\mathbb{R}_{+}^{2})} \\ &\leqslant (q_{j}-1)! \|f_{i,n}\|_{L^{2}(\mathbb{R}_{+})}^{2} \|f_{j,n} \otimes_{q_{j}-1} f_{j,n}\|_{L^{2}(\mathbb{R}_{+}^{2})} \\ &= (q_{j}-1)! E[F_{i,n}^{2}] \|f_{j,n} \otimes_{q_{j}-1} f_{j,n}\|_{L^{2}(\mathbb{R}_{+}^{2})}. \end{split}$$

At this stage, observe the following two facts. First, because $q_i \neq q_j$, we have C(i,j) = 0 necessarily. Second, since $E[F_{j,n}^2] \to C(j,j)$ and $F_{j,n} \xrightarrow{\text{Law}} \mathcal{N}(0,C(j,j))$, we have by Theorem 5.5 that $||f_{j,n} \otimes_{q_j-1} f_{j,n}||_{L^2(\mathbb{R}^2_+)} \to 0$. Hence, (6.80) holds true in this case as well.

Third case: $q_i = q_j \ge 2$. We have, using the product formula (4.43),

$$\frac{1}{q_{i}}\langle DF_{j,n}, DF_{i,n}\rangle_{L^{2}(\mathbb{R}_{+})} = q_{i} \int_{0}^{\infty} I_{q_{i}}(f_{i,n}(\cdot,t))I_{q_{i}}(f_{j,n}(\cdot,t))dt
= q_{i} \sum_{r=0}^{q_{i}-1} r! \binom{q_{i}-1}{r}^{2} I_{2q_{i}-2-2r} \left(\int_{0}^{\infty} f_{i,n}(\cdot,t) \otimes_{r} f_{j,n}(\cdot,t)dt \right)
= q_{i} \sum_{r=0}^{q_{i}-1} r! \binom{q_{i}-1}{r}^{2} I_{2q_{i}-2-2r} (f_{i,n} \otimes_{r+1} f_{j,n})
= q_{i}! \langle f_{i,n}, f_{j,n} \rangle_{L^{2}(\mathbb{R}_{+}^{q_{i}})} + q_{i} \sum_{r=1}^{q_{i}-1} (r-1)! \binom{q_{i}-1}{r-1}^{2} I_{2q_{i}-2r}(f_{i,n} \otimes_{r} f_{j,n})
= E[F_{i,n}F_{j,n}] + q_{i} \sum_{r=1}^{q_{i}-1} (r-1)! \binom{q_{i}-1}{r-1}^{2} I_{2q_{i}-2r}(f_{i,n} \otimes_{r} f_{j,n}).$$

We deduce that

$$E\left[\left(\frac{1}{q_i}\langle DF_{j,n}, DF_{i,n}\rangle_{L^2(\mathbb{R}_+)} - C(i,j)\right)^2\right]$$

$$= \left(E[F_{i,n}F_{j,n}] - C(i,j)\right)^2 + q_i^2 \sum_{r=1}^{q_i-1} (r-1)!^2 \binom{q_i-1}{r-1}^4 (2q_i-2r)! \|f_{i,n}\widetilde{\otimes}_r f_{j,n}\|_{L^2(\mathbb{R}_+^{2q_i-2r})}^2.$$

The first term of the right-hand side tends to zero by assumption. For the second term, we can write, whenever $r \in \{1, ..., q_i - 1\}$,

$$\begin{aligned} \|f_{i,n} \widetilde{\otimes}_{r} f_{j,n}\|_{L^{2}(\mathbb{R}^{2q_{i}-2r}_{+})}^{2} & \leqslant \|f_{i,n} \otimes_{r} f_{j,n}\|_{L^{2}(\mathbb{R}^{2q_{i}-2r}_{+})}^{2} \\ & = \langle f_{i,n} \otimes_{q_{i}-r} f_{i,n}, f_{j,n} \otimes_{q_{i}-r} f_{j,n} \rangle_{L^{2}(\mathbb{R}^{2r}_{+})} \\ & \leqslant \|f_{i,n} \otimes_{q_{i}-r} f_{i,n}\|_{L^{2}(\mathbb{R}^{2r}_{+})} \|f_{j,n} \otimes_{q_{i}-r} f_{j,n}\|_{L^{2}(\mathbb{R}^{2r}_{+})}. \end{aligned}$$

Since $F_{i,n} \stackrel{\text{Law}}{\to} \mathcal{N}(0, C(i, i))$ and $F_{j,n} \stackrel{\text{Law}}{\to} \mathcal{N}(0, C(j, j))$, by Theorem 5.5 we have that $||f_{i,n} \otimes_{q_i - r} f_{i,n}||_{L^2(\mathbb{R}^{2r}_+)} ||f_{j,n} \otimes_{q_i - r} f_{j,n}||_{L^2(\mathbb{R}^{2r}_+)} \to 0$, thereby showing that (6.80) holds true in our third case.

Fourth case: $q_j > q_i \ge 2$ (a similar analysis might be done whenever $q_i > q_j \ge 2$). We have, using the product formula (4.43),

$$\frac{1}{q_{i}}\langle DF_{j,n}, DF_{i,n}\rangle_{L^{2}(\mathbb{R}_{+})} = q_{j} \int_{0}^{\infty} I_{q_{i}}(f_{i,n}(\cdot,t))I_{q_{j}}(f_{j,n}(\cdot,t))dt
= q_{j} \sum_{r=0}^{q_{i}-1} r! \binom{q_{i}-1}{r} I_{q_{i}+q_{j}-2-2r} \left(\int_{0}^{\infty} f_{i,n}(\cdot,t) \otimes_{r} f_{j,n}(\cdot,t)dt \right)
= q_{j} \sum_{r=0}^{q_{i}-1} r! \binom{q_{i}-1}{r} I_{q_{i}+q_{j}-2-2r} (f_{i,n} \otimes_{r+1} f_{j,n})
= q_{j} \sum_{r=1}^{q_{i}} (r-1)! \binom{q_{i}-1}{r-1} \binom{q_{j}-1}{r-1} I_{q_{i}+q_{j}-2r}(f_{i,n} \otimes_{r} f_{j,n}).$$

We deduce that

$$E\left[\frac{1}{q_i}\langle DF_{j,n}, DF_{i,n}\rangle_{L^2(\mathbb{R}_+)}^2\right]$$

$$= q_j^2 \sum_{r=1}^{q_i} (r-1)!^2 \binom{q_i-1}{r-1}^2 \binom{q_j-1}{r-1}^2 (q_i+q_j-2r)! \|f_{i,n}\widetilde{\otimes}_r f_{j,n}\|_{L^2(\mathbb{R}_+^{q_i+q_j-2r})}^2.$$

For any $r \in \{1, \ldots, q_i\}$, we have

$$\begin{split} \|f_{i,n} \widetilde{\otimes}_{r} f_{j,n}\|_{L^{2}(\mathbb{R}^{q_{i}+q_{j}-2r}_{+})}^{2} & \leqslant & \|f_{i,n} \otimes_{r} f_{j,n}\|_{L^{2}(\mathbb{R}^{q_{i}+q_{j}-2r}_{+})}^{2} \\ & = & \langle f_{i,n} \otimes_{q_{i}-r} f_{i,n}, f_{j,n} \otimes_{q_{j}-r} f_{j,n} \rangle_{L^{2}(\mathbb{R}^{2r}_{+})} \\ & \leqslant & \|f_{i,n} \otimes_{q_{i}-r} f_{i,n}\|_{L^{2}(\mathbb{R}^{2r}_{+})} \|f_{j,n} \otimes_{q_{j}-r} f_{j,n}\|_{L^{2}(\mathbb{R}^{2r}_{+})} \\ & \leqslant & \|f_{i,n}\|_{L^{2}(\mathbb{R}^{q_{i}}_{+})}^{2} \|f_{j,n} \otimes_{q_{j}-r} f_{j,n}\|_{L^{2}(\mathbb{R}^{2r}_{+})} \end{split}$$

Since $F_{j,n} \stackrel{\text{Law}}{\to} \mathcal{N}(0,C(j,j))$ and $q_j - r \in \{1,\ldots,q_j-1\}$, by Theorem 5.5 we have that $||f_{j,n}\otimes_{q_j-r}f_{j,n}||_{L^2(\mathbb{R}^{2r}_+)} \to 0$. We deduce that (6.80) holds true in our fourth case.

Summarizing, we have that (6.80) is true for any i and j, and the proof of the theorem is done.

When the integers q_d, \ldots, q_1 are pairwise disjoint in Theorem 6.2, notice that (6.79) is automatically verified with C(i,j) = 0 for all $i \neq j$, see indeed (4.38). As such, we recover the version of Theorem 6.2 (that is, Theorem 1.5) which was stated and used in Lecture 1 to prove Breuer-Major theorem.

To go further. In [32], Stein's method is combined with Malliavin calculus in a multivariate setting to provide bounds for the Wasserstein distance between the laws of $N \sim \mathcal{N}_d(0, C)$ and $F = (F_1, \ldots, F_d)$ where each $F_i \in \mathbb{D}^{1,2}$ verifies $E[F_i] = 0$. Compare with Theorem 6.1.

7 Cumulants on the Wiener space

In this section, our aim is to analyze the cumulants of a given element F of $\mathbb{D}^{1,2}$ and to show how the formula we shall obtain allows to give yet another proof of the Fourth Moment Theorem 1.3.

Let F be a random variable with, say, all the moments (to simplify the exposition). Let ϕ_F denote its characteristic function, that is, $\phi_F(t) = E[e^{itF}]$, $t \in \mathbb{R}$. Then, it is well-known that we may recover the moments of F from ϕ_F through the identity

$$E[F^j] = (-i)^j \frac{d^j}{dt^j}|_{t=0} \phi_F(t).$$

The cumulants of F, denoted by $\{\kappa_j(F)\}_{j\geqslant 1}$, are defined in a similar way, just by replacing ϕ_F by $\log \phi_F$ in the previous expression:

$$\kappa_j(F) = (-i)^j \frac{d^j}{dt^j}|_{t=0} \log \phi_F(t).$$

The first few cumulants are

$$\kappa_1(F) = E[F],$$

$$\kappa_2(F) = E[F^2] - E[F]^2 = Var(F),$$

$$\kappa_3(F) = E[F^3] - 3E[F^2]E[F] + 2E[F]^3.$$

It is immediate that

$$\kappa_j(F+G) = \kappa_j(F) + \kappa_j(G) \quad \text{and} \quad \kappa_j(\lambda F) = \lambda^j \kappa_j(F)$$
(7.81)

for all $j \ge 1$, when $\lambda \in \mathbb{R}$ and F and G are independent random variables (with all the moments). Also, it is easy to express moments in terms of cumulants and vice-versa. Finally, let us observe that the cumulants of $F \sim \mathcal{N}(0, \sigma^2)$ are all zero, except the second one which is σ^2 . This fact, together with the two properties (7.81), gives a quick proof of the classical CLT and illustrates that cumulants are often relevant when wanting to decide whether a given random variable is approximately normally distributed.

The following simple lemma is a useful link between moments and cumulants.

Lemma 7.1 Let F be a random variable (in a given probability space (Ω, \mathcal{F}, P)) having all the moments. Then, for all $m \in \mathbb{N}$,

$$E[F^{m+1}] = \sum_{s=0}^{m} {m \choose s} \kappa_{s+1}(F) E[F^{m-s}].$$

Proof. We can write

$$E[F^{m+1}] = (-i)^{m+1} \frac{d^{m+1}}{dt^{m+1}}|_{t=0} \phi_F(t) = (-i)^{m+1} \frac{d^m}{dt^m}|_{t=0} \left(\frac{d}{dt}(\phi_F(t)\log\phi_F(t))\right)$$

$$= (-i)^{m+1} \sum_{s=0}^m \binom{m}{s} \left(\frac{d^{s+1}}{dt^{s+1}}|_{t=0}\log\phi_F(t)\right) \left(\frac{d^{m-s}}{dt^{m-s}}|_{t=0}\phi_F(t)\right) \text{ by Leibniz}$$

$$= \sum_{s=0}^m \binom{m}{s} \kappa_{s+1}(F) E[F^{m-s}].$$

From now on, we will deal with a random variable F with all moments that is further measurable with respect to the Brownian motion $(B_t)_{t\geqslant 0}$. We let the notation of Section 4 prevail and we consider the chaotic expansion (4.39) of F. We further assume (only to avoid technical issues) that F belongs to \mathbb{D}^{∞} , meaning that $F \in \mathbb{D}^{m,2}$ for all $m \geqslant 1$ and that $E[\|D^m F\|_{L^2(\mathbb{R}^m_+)}^p] < \infty$ for all $m \geqslant 1$ and all $p \geqslant 2$. This assumption allows us to introduce recursively the following (well-defined) sequence of random variables related to F. Namely, set $\Gamma_0(F) = F$ and

$$\Gamma_{j+1}(F) = \langle DF, -DL^{-1}\Gamma_j(F)\rangle_{L^2(\mathbb{R}_+)}.$$

The following result contains a neat expression of the cumulants of F in terms of the family $\{\Gamma_s(F)\}_{s\in\mathbb{N}}$.

Theorem 7.2 (Nourdin, Peccati, 2010; see [26]) Let $F \in \mathbb{D}^{\infty}$. Then, for any $s \in \mathbb{N}$,

$$\kappa_{s+1}(F) = s! E[\Gamma_s(F)].$$

Proof. The proof is by induction. It consists to compute $\kappa_{s+1}(F)$ using the induction hypothesis, together with Lemma 7.1 and (4.53). First, the result holds true for s=0, as it only says that $\kappa_1(F)=E[\Gamma_0(F)]=E[F]$. Assume now that $m \ge 1$ is given and that $\kappa_{s+1}(F)=s!E[\Gamma_s(F)]$ for all $s \le m-1$. We can then write

$$\kappa_{m+1}(F) = E[F^{m+1}] - \sum_{s=0}^{m-1} {m \choose s} \kappa_{s+1}(F) E[F^{m-s}] \text{ by Lemma 7.1}$$

$$= E[F^{m+1}] - \sum_{s=0}^{m-1} s! {m \choose s} E[\Gamma_s(F)] E[F^{m-s}] \text{ by the induction hypothesis.}$$

On the other hand, by applying (4.53) repeatedly, we get

$$E[F^{m+1}] = E[F^m]E[\Gamma_0(F)] + \operatorname{Cov}(F^m, \Gamma_0(F)) = E[F^m]E[\Gamma_0(F)] + mE[F^{m-1}\Gamma_1(F)]$$

$$= E[F^m]E[\Gamma_0(F)] + mE[F^{m-1}]E[\Gamma_1(F)] + m\operatorname{Cov}(F^{m-1}, \Gamma_1(F))$$

$$= E[F^m]E[\Gamma_0(F)] + mE[F^{m-1}]E[\Gamma_1(F)] + m(m-1)E[F^{m-2}\Gamma_2(F)]$$

$$= \dots$$

$$= \sum_{s=0}^m s! \binom{m}{s} E[F^{m-s}]E[\Gamma_s(F)].$$

Thus

$$\kappa_{m+1}(F) = E[F^{m+1}] - \sum_{s=0}^{m-1} s! \binom{m}{s} E[\Gamma_s(F)] E[F^{m-s}] = m! E[\Gamma_m(F)],$$

and the desired conclusion follows.

Let us now focus on the computation of cumulants associated to random variables having the form of a multiple Wiener-Itô integral. The following statement provides a compact representation for the cumulants of such random variables.

Theorem 7.3 Let $q \ge 2$ and assume that $F = I_q(f)$, where $f \in L^2(\mathbb{R}^q_+)$. We have $\kappa_1(F) = 0$, $\kappa_2(F) = q! ||f||^2_{L^2(\mathbb{R}^q_+)}$ and, for every $s \ge 3$,

$$\kappa_s(F) = q!(s-1)! \sum_{r=1}^{\infty} c_q(r_1, \dots, r_{s-2}) \langle (\dots((f \widetilde{\otimes}_{r_1} f) \widetilde{\otimes}_{r_2} f) \dots \widetilde{\otimes}_{r_{s-3}} f) \widetilde{\otimes}_{r_{s-2}} f, f \rangle_{L^2(\mathbb{R}^q_\perp)}, \quad (7.82)$$

where the sum \sum runs over all collections of integers r_1, \ldots, r_{s-2} such that:

(i)
$$1 \leqslant r_1, \ldots, r_{s-2} \leqslant q$$
;

(ii)
$$r_1 + \ldots + r_{s-2} = \frac{(s-2)q}{2}$$
;

(iii)
$$r_1 < q, r_1 + r_2 < \frac{3q}{2}, \ldots, r_1 + \ldots + r_{s-3} < \frac{(s-2)q}{2};$$

(iv)
$$r_2 \leqslant 2q - 2r_1, \ldots, r_{s-2} \leqslant (s-2)q - 2r_1 - \ldots - 2r_{s-3};$$

and where the combinatorial constants $c_q(r_1,\ldots,r_{s-2})$ are recursively defined by the relations

$$c_q(r) = q(r-1)! \binom{q-1}{r-1}^2,$$

and, for $a \ge 2$,

$$c_q(r_1,\ldots,r_a) = q(r_a-1)! \binom{aq-2r_1-\ldots-2r_{a-1}-1}{r_a-1} \binom{q-1}{r_a-1} c_q(r_1,\ldots,r_{a-1}).$$

Remark 7.4 1. If sq is odd, then $\kappa_s(F) = 0$, see indeed condition (ii). This fact is easy to see in any case: use that $\kappa_s(-F) = (-1)^s \kappa_s(F)$ and observe that, when q is odd, then $F \stackrel{\text{(law)}}{=} -F$ (since $B \stackrel{\text{(law)}}{=} -B$).

2. If q = 2 and $F = I_2(f)$ with $f \in L^2(\mathbb{R}^2_+)$, then the only possible integers r_1, \ldots, r_{s-2} verifying (i) - (iv) in the previous statement are $r_1 = \ldots = r_{s-2} = 1$. On the other hand, we immediately compute that $c_2(1) = 2$, $c_2(1,1) = 4$, $c_2(1,1,1) = 8$, and so on. Therefore,

$$\kappa_s(I_2(f)) = 2^{s-1}(s-1)! \langle (...(f \otimes_1 f) \dots f) \otimes_1 f, f \rangle_{L^2(\mathbb{R}^2)}, \tag{7.83}$$

and we recover the classical expression of the cumulants of a double integral.

3. If $q \ge 2$ and $F = I_q(f)$, $f \in L^2(\mathbb{R}^q_+)$, then (7.82) for s = 4 reads

$$\kappa_{4}(I_{q}(f)) = 6q! \sum_{r=1}^{q-1} c_{q}(r, q-r) \langle (f\widetilde{\otimes}_{r}f)\widetilde{\otimes}_{q-r}f, f \rangle_{L^{2}(\mathbb{R}^{q}_{+})}
= \frac{3}{q} \sum_{r=1}^{q-1} rr!^{2} {q \choose r}^{4} (2q-2r)! \langle (f\widetilde{\otimes}_{r}f) \otimes_{q-r} f, f \rangle_{L^{2}(\mathbb{R}^{q}_{+})}
= \frac{3}{q} \sum_{r=1}^{q-1} rr!^{2} {q \choose r}^{4} (2q-2r)! \langle f\widetilde{\otimes}_{r}f, f \otimes_{r} f \rangle_{L^{2}(\mathbb{R}^{2q-2r}_{+})}
= \frac{3}{q} \sum_{r=1}^{q-1} rr!^{2} {q \choose r}^{4} (2q-2r)! \|f\widetilde{\otimes}_{r}f\|_{L^{2}(\mathbb{R}^{2q-2r}_{+})}^{2},$$
(7.84)

and we recover the expression for $\kappa_4(F)$ given in (5.59) by a different route.

Proof of Theorem 7.3. Let us first show the following formula: for any $s \ge 2$, we claim that

$$\Gamma_{s-1}(F) = \sum_{r_1=1}^{q} \dots \sum_{r_{s-1}=1}^{[(s-1)q-2r_1-\dots-2r_{s-2}]\wedge q} c_q(r_1,\dots,r_{s-1}) \mathbf{1}_{\{r_1

$$(7.85)$$$$

We shall prove (7.85) by induction. When s=2, identity (7.85) simply reads $\Gamma_1(F)=\sum_{r=1}^q c_q(r)I_{2q-2r}(f\widetilde{\otimes}_r f)$ and is nothing but (5.62). Assume now that (7.85) holds for $\Gamma_{s-1}(F)$, and let us prove that it continues to hold for $\Gamma_s(F)$. We have, using the product formula (4.43) and following the same line of reasoning as in the proof of (5.62),

$$\begin{split} \Gamma_{s}(F) &= \langle DF, -DL^{-1}\Gamma_{s-1}F\rangle_{L^{2}(\mathbb{R}_{+})} \\ &= \sum_{r_{1}=1}^{q} \dots \sum_{r_{s-1}=1}^{[(s-1)q-2r_{1}-...-2r_{s-2}]\wedge q} qc_{q}(r_{1}, \dots, r_{s-1})\mathbf{1}_{\{r_{1}< q\}} \dots \mathbf{1}_{\{r_{1}+...+r_{s-2}<\frac{(s-1)q}{2}\}} \\ &\quad \times \mathbf{1}_{\{r_{1}+...+r_{s-1}<\frac{sq}{2}\}} \langle I_{q-1}(f), I_{sq-2r_{1}-...-2r_{s-1}-1} \big((...(f\widetilde{\otimes}_{r_{1}}f)\widetilde{\otimes}_{r_{2}}f) \dots f)\widetilde{\otimes}_{r_{s-1}}f \big) \rangle_{L^{2}(\mathbb{R}_{+})} \\ &= \sum_{r_{1}=1}^{q} \dots \sum_{r_{s-1}=1}^{[(s-1)q-2r_{1}-...-2r_{s-2}]\wedge q} \underbrace{\sum_{r_{s}=1}^{[sq-2r_{1}-...-2r_{s-1}]\wedge q} c_{q}(r_{1}, \dots, r_{s-1}) \times q(r_{s}-1)! \\ &\quad \times \binom{sq-2r_{1}-...-2r_{s-1}-1}{r_{s}-1} \binom{q-1}{r_{s}-1} \mathbf{1}_{\{r_{1}< q\}} \dots \mathbf{1}_{\{r_{1}+...+r_{s-2}<\frac{(s-1)q}{2}\}} \\ &\quad \times \mathbf{1}_{\{r_{1}+...+r_{s-1}<\frac{sq}{2}\}} I_{(s+1)q-2r_{1}-...-2r_{s}} \big((...(f\widetilde{\otimes}_{r_{1}}f)\widetilde{\otimes}_{r_{2}}f) \dots f) \widetilde{\otimes}_{r_{s}}f \big), \end{split}$$

which is the desired formula for $\Gamma_s(F)$. The proof of (7.85) for all $s \ge 1$ is thus finished. Now, let us take the expectation on both sides of (7.85). We get

$$\kappa_{s}(F) = (s-1)! E[\Gamma_{s-1}(F)]
= (s-1)! \sum_{r_{1}=1}^{q} \dots \sum_{r_{s-1}=1}^{[(s-1)q-2r_{1}-\dots-2r_{s-2}]\wedge q} c_{q}(r_{1},\dots,r_{s-1}) \mathbf{1}_{\{r_{1}< q\}} \dots \mathbf{1}_{\{r_{1}+\dots+r_{s-2}<\frac{(s-1)q}{2}\}}
\times \mathbf{1}_{\{r_{1}+\dots+r_{s-1}=\frac{sq}{2}\}} \times (\dots (f \widetilde{\otimes}_{r_{1}} f) \widetilde{\otimes}_{r_{2}} f) \dots f) \widetilde{\otimes}_{r_{s-1}} f.$$

Observe that, if $2r_1 + \ldots + 2r_{s-1} = sq$ and $r_{s-1} \leq (s-1)q - 2r_1 - \ldots - 2r_{s-2}$ then $2r_{s-1} = q + (s-1)q - 2r_1 - \ldots - 2r_{s-2} \geq q + r_{s-1}$, so that $r_{s-1} \geq q$. Therefore,

$$\kappa_{s}(F) = (s-1)! \sum_{r_{1}=1}^{q} \dots \sum_{r_{s-2}=1}^{[(s-2)q-2r_{1}-...-2r_{s-3}] \wedge q} c_{q}(r_{1},...,r_{s-2},q) \mathbf{1}_{\{r_{1} < q\}} \dots \mathbf{1}_{\{r_{1}+...+r_{s-3} < \frac{(s-2)q}{2}\}} \times \mathbf{1}_{\{r_{1}+...+r_{s-2} = \frac{(s-2)q}{2}\}} \langle (...(f\widetilde{\otimes}_{r_{1}}f)\widetilde{\otimes}_{r_{2}}f) \dots f) \widetilde{\otimes}_{r_{s-2}}f, f \rangle_{L^{2}(\mathbb{R}^{q}_{+})},$$

which is the announced result, since $c_q(r_1, \ldots, r_{s-2}, q) = q!c_q(r_1, \ldots, r_{s-2})$.

We conclude this section by providing yet another proof (based on our new formula (7.82)) of the Fourth Moment Theorem 1.3. More precisely, let us show by another route that, if $q \ge 2$ is fixed and if $(F_n)_{n\ge 1}$ is a sequence of the form $F_n = I_q(f_n)$ with $f_n \in L^2(\mathbb{R}^q_+)$ such that $E[F_n^2] = q! ||f_n||^2_{L^2(\mathbb{R}^q_+)} = 1$ for all $n \ge 1$ and $E[F_n^4] \to 3$ as $n \to \infty$, then $F_n \to \mathcal{N}(0,1)$ in law as $n \to \infty$.

To this end, observe that $\kappa_1(F_n) = 0$ and $\kappa_2(F_n) = 1$. To estimate $\kappa_s(F_n)$, $s \ge 3$, we consider the expression (7.82). Let r_1, \ldots, r_{s-2} be some integers such that (i)–(iv) in Theorem 7.3 are satisfied. Using Cauchy-Schwarz and then successively

$$\|g\widetilde{\otimes}_r h\|_{L^2(\mathbb{R}^{p+q-2r})} \le \|g\otimes_r h\|_{L^2(\mathbb{R}^{p+q-2r})} \le \|g\|_{L^2(\mathbb{R}^p_+)} \|h\|_{L^2(\mathbb{R}^q_+)}$$

whenever $g \in L^2(\mathbb{R}^p_+)$, $h \in L^2(\mathbb{R}^q_+)$ and $r = 1, \ldots, p \wedge q$, we get that

$$\left| \langle (\dots(f_n \widetilde{\otimes}_{r_1} f_n) \widetilde{\otimes}_{r_2} f_n) \dots f_n) \widetilde{\otimes}_{r_{s-2}} f_n, f_n \rangle_{L^2(\mathbb{R}^q_+)} \right| \\
\leqslant \| (\dots(f_n \widetilde{\otimes}_{r_1} f_n) \widetilde{\otimes}_{r_2} f_n) \dots f_n) \widetilde{\otimes}_{r_{s-2}} f_n \|_{L^2(\mathbb{R}^q_+)} \| f_n \|_{L^2(\mathbb{R}^q_+)} \\
\leqslant \| f_n \widetilde{\otimes}_{r_1} f_n \|_{L^2(\mathbb{R}^{2q-2r_1}_+)} \| f_n \|_{L^2(\mathbb{R}^q_+)}^{s-2} \\
= (q!)^{1-\frac{s}{2}} \| f_n \widetilde{\otimes}_{r_1} f_n \|_{L^2(\mathbb{R}^{2q-2r_1}_+)}. \tag{7.86}$$

Since $E[F_n^4] - 3 = \kappa_4(F_n) \to 0$, we deduce from (7.84) that $||f_n \widetilde{\otimes}_r f_n||_{L^2(\mathbb{R}^{2q-2r}_+)} \to 0$ for all $r = 1, \ldots, q-1$. Consequently, by combining (7.82) with (7.86), we get that $\kappa_s(F_n) \to 0$ as $n \to \infty$ for all $s \geqslant 3$, implying in turn that $F_n \to \mathcal{N}(0,1)$ in law.

To go further. The multivariate version of Theorem 7.2 may be found in [21].

8 A new density formula

In this section, we shall explain how the quantity $\langle DF, -DL^{-1}F \rangle_{L^2(\mathbb{R}_+)}$ is related to the density of $F \in \mathbb{D}^{1,2}$ (provided it exists). More specifically, when $F \in \mathbb{D}^{1,2}$ is such that E[F] = 0, let us introduce the function $g_F : \mathbb{R} \to \mathbb{R}$, defined by means of the following identity:

$$g_F(F) = E[\langle DF, -DL^{-1}F \rangle_{L^2(\mathbb{R}_+)} | F].$$
 (8.87)

A key property of the random variable $g_F(F)$ is as follows.

Proposition 8.1 If $F \in \mathbb{D}^{1,2}$ satisfies E[F] = 0, then $P(g_F(F) \ge 0) = 1$.

Proof. Let C be a Borel set of \mathbb{R} and set $\phi_n(x) = \int_0^x \mathbf{1}_{C \cap [-n,n]}(t) dt$, $n \ge 1$ (with the usual convention $\int_0^x = -\int_x^0$ for x < 0). Since ϕ_n is increasing and vanishing at zero, we have $x\phi_n(x) \ge 0$ for all $x \in \mathbb{R}$. In particular,

$$0 \leqslant E[F\phi_n(F)] = E\left[F\int_0^F \mathbf{1}_{C\cap[-n,n]}(t)dt\right] = E\left[F\int_{-\infty}^F \mathbf{1}_{C\cap[-n,n]}(t)dt\right].$$

Therefore, we deduce from Corollary 4.17 that $E\left[g_F(F)\mathbf{1}_{C\cap[-n,n]}(F)\right]\geqslant 0$. By dominated convergence, this yields $E\left[g_F(F)\mathbf{1}_C(F)\right]\geqslant 0$, implying in turn that $P(g_F(F)\geqslant 0)=1$.

The following theorem gives a new density formula for F in terms of the function g_F . We will then study some of its consequences.

Theorem 8.2 (Nourdin, Viens, 2009; see [34]) Let $F \in \mathbb{D}^{1,2}$ with E[F] = 0. Then, the law of F admits a density with respect to Lebesgue measure (say, $\rho : \mathbb{R} \to \mathbb{R}$) if and only if $P(g_F(F) > 0) = 1$. In this case, the support of ρ , denoted by supp ρ , is a closed interval of \mathbb{R} containing zero and we have, for (almost) all $x \in \text{supp } \rho$:

$$\rho(x) = \frac{E[|F|]}{2g_F(x)} \exp\left(-\int_0^x \frac{y \, dy}{g_F(y)}\right). \tag{8.88}$$

Proof. Assume that $P(g_F(F) > 0) = 1$ and let C be a Borel set. Let $n \ge 1$. Corollary 4.17 yields

$$E\left[F\int_{-\infty}^{F} \mathbf{1}_{C\cap[-n,n]}(t)dt\right] = E\left[\mathbf{1}_{C\cap[-n,n]}(F)g_F(F)\right]. \tag{8.89}$$

Suppose that the Lebesgue measure of C is zero. Then $\int_{-\infty}^{F} \mathbf{1}_{C\cap[-n,n]}(t)dt = 0$, so that $E\left[\mathbf{1}_{C\cap[-n,n]}(F)g_F(F)\right] = 0$ by (8.89). But, since $P(g_F(F) > 0) = 1$, we get that $P(F \in C \cap [-n,n]) = 0$ and, by letting $n \to \infty$, that $P(F \in C) = 0$. Therefore, the Radon-Nikodym criterion is verified, hence implying that the law of F has a density.

Conversely, assume that the law of F has a density, say ρ . Let $\phi : \mathbb{R} \to \mathbb{R}$ be a continuous function with compact support, and let Φ denote any antiderivative of ϕ . Note that Φ is necessarily bounded. We can write:

$$E[\phi(F)g_F(F)] = E[\Phi(F)F] \text{ by } (4.53)$$

$$= \int_{\mathbb{R}} \Phi(x) x \, \rho(x) dx = \int_{\mathbb{R}} \phi(x) \left(\int_{x}^{\infty} y \rho(y) dy \right) dx = E\left[\phi(F) \frac{\int_{F}^{\infty} y \rho(y) dy}{\rho(F)}\right].$$

Equation (*) was obtained by integrating by parts, after observing that

$$\int_{x}^{\infty} y \rho(y) dy \to 0 \text{ as } |x| \to \infty$$

(for $x \to +\infty$, this is because $F \in L^1(\Omega)$; for $x \to -\infty$, this is because F has mean zero). Therefore, we have shown that, P-a.s.,

$$g_F(F) = \frac{\int_F^\infty y \rho(y) dy}{\rho(F)}.$$
(8.90)

(Notice that $P(\rho(F) > 0) = \int_{\mathbb{R}} \mathbf{1}_{\{\rho(x) > 0\}} \rho(x) dx = \int_{\mathbb{R}} \rho(x) dx = 1$, so that identity (8.90) always makes sense.) Since $F \in \mathbb{D}^{1,2}$, one has (see, e.g., [35, Proposition 2.1.7]) that supp $\rho = [\alpha, \beta]$ with $-\infty \leqslant \alpha < \beta \leqslant +\infty$. Since F has zero mean, note that $\alpha < 0$ and $\beta > 0$ necessarily. For every $x \in (\alpha, \beta)$, define

$$\varphi(x) = \int_{x}^{\infty} y \rho(y) \, dy. \tag{8.91}$$

The function φ is differentiable almost everywhere on (α, β) , and its derivative is $-x\rho(x)$. In particular, since $\varphi(\alpha) = \varphi(\beta) = 0$ and φ is strictly increasing before 0 and strictly decreasing afterwards, we have $\varphi(x) > 0$ for all $x \in (\alpha, \beta)$. Hence, (8.90) implies that $P(g_F(F) > 0) = 1$.

Finally, let us prove (8.88). Let φ still be defined by (8.91). On the one hand, we have $\varphi'(x) = -x\rho(x)$ for almost all $x \in \text{supp } \rho$. On the other hand, by (8.90), we have, for almost all $x \in \text{supp } \rho$,

$$\varphi(x) = \rho(x)g_F(x). \tag{8.92}$$

By putting these two facts together, we get the following ordinary differential equation satisfied by φ :

$$\frac{\varphi'(x)}{\varphi(x)} = -\frac{x}{g_F(x)} \quad \text{for almost all } x \in \text{supp } \rho.$$

Integrating this relation over the interval [0, x] yields

$$\log \varphi(x) = \log \varphi(0) - \int_0^x \frac{y \, dy}{g_F(y)}.$$

Taking the exponential and using $0 = E(F) = E(F_+) - E(F_-)$ so that $E|F| = E(F_+) + E(F_-) = 2E(F_+) = 2\varphi(0)$, we get

$$\varphi(x) = \frac{1}{2} E[|F|] \exp\left(-\int_0^x \frac{y \, dy}{g_F(y)}\right).$$

Finally, the desired conclusion comes from (8.92).

A consequence of Theorem 8.2 is the following statement, yielding sufficient conditions in order for the law of F to have a support equal to the real line.

Corollary 8.3 Let $F \in \mathbb{D}^{1,2}$ with E[F] = 0. Assume that there exists $\sigma_{\min} > 0$ such that

$$g_F(F) \geqslant \sigma_{\min}^2, \quad P\text{-}a.s.$$
 (8.93)

Then the law of F, which has a density ρ by Theorem 8.2, has \mathbb{R} for support and (8.88) holds almost everywhere in \mathbb{R} .

Proof. It is an immediate consequence of Theorem 8.2, except the fact that supp $\rho = \mathbb{R}$. For the moment, we just know that supp $\rho = [\alpha, \beta]$ with $-\infty \leq \alpha < 0 < \beta \leq +\infty$. Identity (8.90) yields

$$\int_{x}^{\infty} y \rho(y) \, dy \geqslant \sigma_{\min}^{2} \rho(x) \quad \text{for almost all } x \in (\alpha, \beta).$$
 (8.94)

Let φ be defined by (8.91), and recall that $\varphi(x) > 0$ for all $x \in (\alpha, \beta)$. When multiplied by $x \in [0, \beta)$, the inequality (8.94) gives $\frac{\varphi'(x)}{\varphi(x)} \geqslant -\frac{x}{\sigma_{\min}^2}$. Integrating this relation over the interval [0, x] yields $\log \varphi(x) - \log \varphi(0) \geqslant -\frac{x^2}{2 \sigma_{\min}^2}$, i.e., since $\varphi(0) = \frac{1}{2} E|F|$,

$$\varphi(x) = \int_{x}^{\infty} y \rho(y) \, dy \geqslant \frac{1}{2} E|F| e^{-\frac{x^2}{2\sigma_{\min}^2}}.$$
(8.95)

Similarly, when multiplied by $x \in (\alpha, 0]$, inequality (8.94) gives $\frac{\varphi'(x)}{\varphi(x)} \leqslant -\frac{x}{\sigma_{\min}^2}$. Integrating this relation over the interval [x, 0] yields $\log \varphi(0) - \log \varphi(x) \leqslant \frac{x^2}{2\sigma_{\min}^2}$, i.e. (8.95) still holds for $x \in (\alpha, 0]$. Now, let us prove that $\beta = +\infty$. If this were not the case, by definition, we would have $\varphi(\beta) = 0$; on the other hand, by letting x tend to β in the above inequality, because φ is continuous, we would have $\varphi(\beta) \geqslant \frac{1}{2}E|F|e^{-\frac{\beta^2}{2\sigma_{\min}^2}} > 0$, which contradicts $\beta < +\infty$. The proof of $\alpha = -\infty$ is similar. In conclusion, we have shown that supp $\rho = \mathbb{R}$.

Using Corollary 8.3, we deduce a neat criterion for normality.

Corollary 8.4 Let $F \in \mathbb{D}^{1,2}$ with E[F] = 0 and assume that F is not identically zero. Then F is Gaussian if and only if $Var(g_F(F)) = 0$.

Proof: By (4.53) (choose $\varphi(x) = x$, G = F and recall that E[F] = 0), we have

$$E[\langle DF, -DL^{-1}F \rangle_{5}] = E[F^{2}] = \text{Var}F. \tag{8.96}$$

Therefore, the condition $\operatorname{Var}(g_F(F)) = 0$ is equivalent to $P(g_F(F) = \operatorname{Var} F) = 1$. Let $F \sim \mathcal{N}(0, \sigma^2)$ with $\sigma > 0$. Using (8.90), we immediately check that $g_F(F) = \sigma^2$, P-a.s. Conversely, if $g_F(F) = \sigma^2 > 0$ P-a.s., then Corollary 8.3 implies that the law of F has a density ρ , given by $\rho(x) = \frac{E|F|}{2\sigma^2}e^{-\frac{x^2}{2\sigma^2}}$ for almost all $x \in \mathbb{R}$, from which we immediately deduce that $F \sim \mathcal{N}(0, \sigma^2)$.

Observe that if $F \sim \mathcal{N}(0, \sigma^2)$ with $\sigma > 0$, then $E|F| = \sqrt{2/\pi} \sigma$, so that the formula (8.88) for ρ agrees, of course, with the usual one in this case.

As a 'concrete' application of (8.88), let us consider the following situation. Let $K:[0,1]^2 \to \mathbb{R}$ be a square-integrable kernel such that K(t,s)=0 for s>t, and consider the centered Gaussian process $X=(X_t)_{t\in[0,1]}$ defined as

$$X_t = \int_0^1 K(t, s) dB_s = \int_0^t K(t, s) dB_s, \quad t \in [0, 1].$$
(8.97)

Fractional Brownian motion is an instance of such a process, see, e.g., [23, Section 2.3]. Consider the maximum

$$Z = \sup_{t \in [0,1]} X_t. \tag{8.98}$$

Assume further that the kernel K satisfies

$$\exists c, \alpha > 0, \quad \forall s, t \in [0, 1]^2, \ s \neq t, \quad 0 < \int_0^1 (K(t, u) - K(s, u))^2 du \leqslant c|t - s|^{\alpha}. \tag{8.99}$$

It ensures (see, e.g., [10]) that: (i) $Z \in \mathbb{D}^{1,2}$; (ii) the law of Z has a density with respect to Lebesgue measure; (iii) there exists a (a.s.) unique random point $\tau \in [0,1]$ where the supremum is attained, that is, such that $Z = X_{\tau} = \int_0^1 K(\tau,s)dB_s$; and (iv) $D_tZ = K(\tau,t)$, $t \in [0,1]$. We claim the following formula.

Proposition 8.5 Let Z be given by (8.98), X be defined as (8.97) and $K \in L^2([0,1]^2)$ be satisfying (8.99). Then, the law of Z has a density ρ whose support is \mathbb{R}_+ , given by

$$\rho(x) = \frac{E|Z - E[Z]|}{2h_Z(x)} \exp\left(-\int_{E[Z]}^x \frac{(y - E[Z])dy}{h_Z(y)}\right), \quad x \geqslant 0.$$

Here,

$$h_Z(x) = \int_0^\infty e^{-u} \mathbf{E} \left[R(\tau_0, \tau_u) | Z = x \right] du,$$

where $R(s,t) = E[X_sX_t]$, $s,t \in [0,1]$, and τ_u is the (almost surely) unique random point where

$$X_t^{(u)} = \int_0^1 K(t,s)(e^{-u}dB_s + \sqrt{1 - e^{-2u}}dB_s')$$

attains its maximum on [0,1], with (B,B') a two-dimensional Brownian motion defined on the product probability space $(\Omega, \mathcal{F}, \mathbf{P}) = (\Omega \times \Omega', \mathscr{F} \otimes \mathscr{F}', P \times P')$.

Proof. Set F = Z - E[Z]. We have $-D_t L^{-1}F = \sum_{q=1}^{\infty} I_{q-1}(f_q(\cdot,t))$ and $D_t F = \sum_{q=1}^{\infty} q I_{q-1}(f_q(\cdot,t))$. Thus

$$\int_0^\infty e^{-u} P_u(D_t F) du = \sum_{q=1}^\infty I_{q-1}(f_q(\cdot, t)) \int_0^\infty e^{-u} q e^{-(q-1)u} du = \sum_{q=1}^\infty I_{q-1}(f_q(\cdot, t)).$$

Consequently,

$$-D_t L^{-1} F = \int_0^\infty e^{-u} P_u(D_t F) du, \quad t \in [0, 1].$$

By Mehler's formula (4.49), and since $DF = DZ = K(\tau, \cdot)$ with $\tau = \operatorname{argmax}_{t \in [0,1]} \int_0^1 K(t, s) dB_s$, we deduce that

$$-D_t L^{-1} F = \int_0^\infty e^{-u} E'[K(\tau_u, t)] du,$$

implying in turn

$$g_{F}(F) = E[\langle DF, -DL^{-1}F \rangle_{L^{2}([0,1])}|F] = \int_{0}^{1} dt \int_{0}^{\infty} du \, e^{-u}K(\tau_{0}, t)E[E'[K(\tau_{u}, t)|F]]$$

$$= \int_{0}^{\infty} e^{-u}E\left[E'\left[\int_{0}^{1} K(\tau_{0}, t)K(\tau_{u}, t)dt|F\right]\right] du = \int_{0}^{\infty} e^{-u}E\left[E'\left[R(\tau_{0}, \tau_{u})|F\right]\right] du$$

$$= \int_{0}^{\infty} e^{-u}\mathbf{E}\left[R(\tau_{0}, \tau_{u})|F\right] du.$$

The desired conclusion follows now from Theorem 8.2 and the fact that F = Z - E[Z].

To go further. The reference [34] also contains interesting concentration inequalities for centered random variables $F \in \mathbb{D}^{1,2}$ satisfying $g_F(F) \leq \alpha F + \beta$. The paper [38] shows how Theorem 8.2 can lead to optimal Gaussian density estimates for a class of stochastic equations with additive noise.

9 Exact rates of convergence

Let $\{F_n\}_{n\geqslant 1}$ be a sequence of random variables in $\mathbb{D}^{1,2}$ such that $E[F_n]=0$, $Var(F_n)=1$ and $F_n \stackrel{\text{law}}{\to} N \sim \mathcal{N}(0,1)$ as $n\to\infty$. In this section, our aim is to develop tools for computing the exact asymptotic expression of the (suitably normalized) sequence

$$P(F_n \leqslant x) - P(N \leqslant x), \quad n \geqslant 1,$$

when $x \in \mathbb{R}$ is fixed. This will complement the content of Theorem 5.2.

A technical computation. For every fixed x, we denote by $f_x : \mathbb{R} \to \mathbb{R}$ the function

$$f_{x}(u) = e^{u^{2}/2} \int_{-\infty}^{u} (\mathbf{1}_{(-\infty,x]}(a) - \Phi(x)) e^{-a^{2}/2} da$$

$$= \sqrt{2\pi} e^{u^{2}/2} \times \begin{cases} \Phi(u)(1 - \Phi(x)) & \text{if } u \leq x \\ \Phi(x)(1 - \Phi(u)) & \text{if } u \geqslant x \end{cases}, \tag{9.100}$$

where $\Phi(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{x} e^{-a^2/2} da$. We have the following result.

Proposition 9.1 Let $N \sim \mathcal{N}(0,1)$. We have, for every $x \in \mathbb{R}$,

$$E[f_x'(N)N] = \frac{1}{3}(x^2 - 1)\frac{e^{-x^2/2}}{\sqrt{2\pi}}. (9.101)$$

Proof. Integrating by parts (the bracket term is easily shown to vanish), we first obtain that

$$E[f'_x(N)N] = \int_{-\infty}^{+\infty} f'_x(u)u \frac{e^{-u^2/2}}{\sqrt{2\pi}} du = \int_{-\infty}^{+\infty} f_x(u)(u^2 - 1) \frac{e^{-u^2/2}}{\sqrt{2\pi}} du$$
$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} (u^2 - 1) \left(\int_{-\infty}^{u} \left[\mathbf{1}_{(-\infty, x]}(a) - \Phi(x) \right] e^{-a^2/2} da \right) du.$$

Integrating by parts once again, this time using the relation $u^2 - 1 = \frac{1}{3}(u^3 - 3u)'$, we deduce that

$$\int_{-\infty}^{+\infty} (u^2 - 1) \left(\int_{-\infty}^{u} \left[\mathbf{1}_{(-\infty, x]}(a) - \Phi(x) \right] e^{-a^2/2} da \right) du$$

$$= -\frac{1}{3} \int_{-\infty}^{+\infty} (u^3 - 3u) \left[\mathbf{1}_{(-\infty, x]}(u) - \Phi(x) \right] e^{-u^2/2} du$$

$$= -\frac{1}{3} \left(\int_{-\infty}^{x} (u^3 - 3u) e^{-u^2/2} du - \Phi(x) \int_{-\infty}^{+\infty} (u^3 - 3u) e^{-u^2/2} du \right)$$

$$= \frac{1}{3} (x^2 - 1) e^{-x^2/2}, \quad \text{since } [(u^2 - 1) e^{-u^2/2}]' = -(u^3 - 3u) e^{-u^2/2}.$$

A general result. Assume that $\{F_n\}_{n\geq 1}$ is a sequence of (sufficiently regular) centered random variables with unitary variance such that the sequence

$$\varphi(n) := \sqrt{E[(1 - \langle DF_n, -DL^{-1}F_n \rangle_{L^2(\mathbb{R}_+)})^2]}, \quad n \geqslant 1,$$
(9.102)

converges to zero as $n \to \infty$. According to Theorem 5.2 one has that, for any $x \in \mathbb{R}$ and as $n \to \infty$,

$$P(F_n \leqslant x) - P(N \leqslant x) \leqslant d_{TV}(F_n, N) \leqslant 2\varphi(n) \to 0, \tag{9.103}$$

where $N \sim \mathcal{N}(0,1)$. The forthcoming result provides a useful criterion in order to compute an exact asymptotic expression (as $n \to \infty$) for the quantity

$$\frac{P(F_n \leqslant x) - P(N \leqslant x)}{\varphi(n)}, \quad n \geqslant 1.$$

Theorem 9.2 (Nourdin, Peccati, 2010; see [27]) Let $\{F_n\}_{n\geq 1}$ be a sequence of random variables belonging to $\mathbb{D}^{1,2}$, and such that $E[F_n] = 0$, $Var[F_n] = 1$. Suppose moreover that the following three conditions hold:

- (i) we have $0 < \varphi(n) < \infty$ for every n and $\varphi(n) \to 0$ as $n \to \infty$;
- (ii) the law of F_n has a density with respect to Lebesgue measure for every n;
- (iii) as $n \to \infty$, the two-dimensional vector $\left(F_n, \frac{\langle DF_n, -DL^{-1}F_n \rangle_{L^2(\mathbb{R}_+)} 1}{\varphi(n)}\right)$ converges in distribution to a centered two-dimensional Gaussian vector (N_1, N_2) , such that $E[N_1^2] = E[N_2^2] = 1$ and $E[N_1N_2] = \rho$.

Then, as $n \to \infty$, one has for every $x \in \mathbb{R}$,

$$\frac{P(F_n \leqslant x) - P(N \leqslant x)}{\varphi(n)} \to \frac{\rho}{3} (1 - x^2) \frac{e^{-x^2/2}}{\sqrt{2\pi}}.$$
 (9.104)

Proof. For any integer n and any C^1 -function f with a bounded derivative, we know by Theorem 4.15 that

$$E[F_n f(F_n)] = E[f'(F_n)\langle DF_n, -DL^{-1}F_n\rangle_{L^2(\mathbb{R}_+)}].$$

Fix $x \in \mathbb{R}$ and observe that the function f_x defined by (9.100) is not \mathcal{C}^1 due to the singularity in x. However, by using a regularization argument given assumption (ii), one can show that the identity

$$E[F_n f_x(F_n)] = E[f'_x(F_n)\langle DF_n, -DL^{-1}F_n\rangle_{L^2(\mathbb{R}_+)}]$$

is true for any n. Therefore, since $P(F_n \leq x) - P(N \leq x) = E[f'_x(F_n)] - E[F_n f_x(F_n)]$, we get

$$\frac{P(F_n \leqslant x) - P(N \leqslant x)}{\varphi(n)} = E\left[f'_x(F_n) \times \frac{1 - \langle DF_n, -DL^{-1}F_n \rangle_{L^2(\mathbb{R}_+)}}{\varphi(n)}\right].$$

Reasoning as in Lemma 3.1, one may show that f_x is Lipschitz with constant 2. Since $\varphi(n)^{-1}(1-\langle DF_n, -DL^{-1}F_n\rangle_{L^2(\mathbb{R}_+)})$ has variance 1 by definition of $\varphi(n)$, we deduce that the sequence

$$f'_x(F_n) \times \frac{1 - \langle DF_n, -DL^{-1}F_n \rangle_{L^2(\mathbb{R}_+)}}{\varphi(n)}, \quad n \geqslant 1,$$

is uniformly integrable. Definition (9.100) shows that $u \to f'_x(u)$ is continuous at every $u \neq x$. This yields that, as $n \to \infty$ and due to assumption (iii),

$$E\left[f'_x(F_n) \times \frac{1 - \langle DF_n, -DL^{-1}F_n \rangle_{L^2(\mathbb{R}_+)}}{\varphi(n)}\right] \to -E[f'_x(N_1)N_2] = -\rho E[f'_x(N_1)N_1].$$

Consequently, relation (9.104) now follows from formula (9.101).

The double integrals case and a concrete application. When applying Theorem 9.2 in concrete situations, the main issue is often to check that condition (ii) therein holds true. In the particular case of sequences belonging to the second Wiener chaos, we can go further in the analysis, leading to the following result.

Proposition 9.3 Let $N \sim \mathcal{N}(0,1)$ and let $F_n = I_2(f_n)$ be such that $f_n \in L^2(\mathbb{R}^2_+)$ is symmetric for all $n \geq 1$. Write $\kappa_p(F_n)$, $p \geq 1$, to indicate the sequence of the cumulants of F_n . Assume that $\kappa_2(F_n) = E[F_n^2] = 1$ for all $n \geq 1$ and that $\kappa_4(F_n) = E[F_n^4] - 3 \rightarrow 0$ as $n \rightarrow \infty$. If we have in addition that

$$\frac{\kappa_3(F_n)}{\sqrt{\kappa_4(F_n)}} \to \alpha \quad and \quad \frac{\kappa_8(F_n)}{(\kappa_4(F_n))^2} \to 0, \tag{9.105}$$

then, for all $x \in \mathbb{R}$,

$$\frac{P(F_n \leqslant x) - P(N \leqslant x)}{\sqrt{\kappa_4(F_n)}} \to \frac{\alpha}{6\sqrt{2\pi}} \left(1 - x^2\right) e^{-\frac{x^2}{2}} \quad as \ n \to \infty. \tag{9.106}$$

Remark 9.4 Due to (9.105), we see that (9.106) is equivalent to

$$\frac{P(F_n \leqslant x) - P(N \leqslant x)}{\kappa_3(F_n)} \to \frac{1}{6\sqrt{2\pi}} \left(1 - x^2\right) e^{-\frac{x^2}{2}} \quad \text{as } n \to \infty.$$

Since each F_n is centered, one also has that $\kappa_3(F_n) = E[F_n^3]$.

Proof. We shall apply Theorem 9.2. Thanks to (5.60), we get that

$$\frac{\kappa_4(F_n)}{6} = \frac{E[F_n^4] - 3}{6} = 8 \|f_n \otimes_1 f_n\|_{L^2(\mathbb{R}^2_+)}^2.$$

By combining this identity with (5.58) (here, it is worth observing that $f_n \otimes_1 f_n$ is symmetric, so that the symmetrization $f_n \otimes_1 f_n$ is immaterial), we see that the quantity $\varphi(n)$ appearing in (9.102) is given by $\sqrt{\kappa_4(F_n)/6}$. In particular, condition (i) in Theorem 9.2 is met (here, let us stress that one may show that $\kappa_4(F_n) > 0$ for all n by means of (5.60)). On the other hand, since F_n is a non-zero double integral, its law has a density with respect to Lebesgue measure, according to Theorem 4.18. This means that condition (ii) in Theorem 9.2 is also in order. Hence, it remains to check condition (iii). Assume that (9.105) holds. Using (7.83) in the cases p = 3 and p = 8, we deduce that

$$\frac{\kappa_3(F_n)}{\sqrt{\kappa_4(F_n)}} = \frac{8 \langle f_n, f_n \otimes_1 f_n \rangle_{L^2(\mathbb{R}^2_+)}}{\sqrt{6} \varphi(n)}$$

and

$$\frac{\kappa_8(F_n)}{(\kappa_4(F_n))^2} = \frac{17920 \|(f_n \otimes_1 f_n) \otimes_1 (f_n \otimes_1 f_n)\|_{L^2(\mathbb{R}^2_+)}^2}{\varphi(n)^4}.$$

On the other hand, set

$$Y_n = \frac{\frac{1}{2} \|DF_n\|_{L^2(\mathbb{R}_+)}^2 - 1}{\varphi(n)}.$$

By (5.62), we have $\frac{1}{2} ||DY_n||_{L^2(\mathbb{R}_+)}^2 - 1 = 2 I_2(f_n \otimes_1 f_n)$. Therefore, by (5.58), we get that

$$E\left[\left(\frac{1}{2}\|DY_n\|_{L^2(\mathbb{R}_+)}^2 - 1\right)^2\right] = \frac{128}{\varphi(n)^4} \|(f_n \otimes_1 f_n) \otimes_1 (f_n \otimes_1 f_n)\|_{L^2(\mathbb{R}_+^2)}$$
$$= \frac{\kappa_8(F_n)}{140 (\kappa_4(F_n))^2} \to 0 \text{ as } n \to \infty.$$

Hence, by Theorem 5.5, we deduce that $Y_n \stackrel{\text{Law}}{\to} \mathcal{N}(0,1)$. We also have

$$E[Y_n F_n] = \frac{4}{\varphi(n)} \langle f_n \otimes_1 f_n, f_n \rangle_{L^2(\mathbb{R}^2_+)} = \frac{\sqrt{6} \kappa_3(F_n)}{2\sqrt{\kappa_4(F_n)}} \to \frac{\alpha\sqrt{6}}{2} =: \rho \quad \text{as } n \to \infty.$$

Therefore, to conclude that condition (iii) in Theorem 9.2 holds true, it suffices to apply Theorem 6.2.

To give a concrete application of Proposition 9.3, let us go back to the quadratic variation of fractional Brownian motion. Let $B^H = (B_t^H)_{t\geqslant 0}$ be a fractional Brownian motion with Hurst index $H \in (0, \frac{1}{2})$ and let

$$F_n = \frac{1}{\sigma_n} \sum_{k=0}^{n-1} \left[(B_{k+1}^H - B_k^H)^2 - 1 \right],$$

where $\sigma_n > 0$ is so that $E[F_n^2] = 1$. Recall from Theorem 5.6 that $\lim_{n \to \infty} \sigma_n^2/n = 2 \sum_{r \in \mathbb{Z}} \rho^2(r) < \infty$, with $\rho : \mathbb{Z} \to \mathbb{R}_+$ given by (5.68); moreover, there exists a constant $c_H > 0$ (depending only on H) such that, with $N \sim \mathcal{N}(0,1)$,

$$d_{TV}(F_n, N) \leqslant \frac{c_H}{\sqrt{n}}, \quad n \geqslant 1. \tag{9.107}$$

The next two results aim to show that one can associate a *lower* bound to (9.107). We start by the following auxiliary result.

Proposition 9.5 Fix an integer $s \ge 2$, let F_n be as above and let ρ be given by (5.68). Recall that $H < \frac{1}{2}$, so that $\rho \in \ell^1(\mathbb{Z})$. Then, the sth cumulant of F_n behaves asymptotically as

$$\kappa_s(F_n) \sim n^{1-s/2} \, 2^{s/2-1} (s-1)! \, \frac{\langle \rho^{*(s-1)}, \rho \rangle_{\ell^2(\mathbb{Z})}}{\|\rho\|_{\ell^2(\mathbb{Z})}^s} \quad as \ n \to \infty.$$
(9.108)

Proof. As in the proof of Theorem 5.6, we have that $F_n \stackrel{\text{law}}{=} I_2(f_n)$ with $f_n = \frac{1}{\sigma_n} \sum_{k=0}^{n-1} e_k^{\otimes 2}$. Now, let us proceed with the proof. It is divided into several steps.

First step. Using the formula (7.83) giving the cumulants of $F_n = I_2(f_n)$ as well as the very definition of the contraction \otimes_1 , we immediately check that

$$\kappa_s(F_n) = \frac{2^{s-1}(s-1)!}{\sigma_n^s} \sum_{k_1,\dots,k_s=0}^{n-1} \rho(k_s - k_{s-1}) \dots \rho(k_2 - k_1) \rho(k_1 - k_s).$$

Second step. Since $H < \frac{1}{2}$, we have that $\rho \in \ell^1(\mathbb{Z})$. Therefore, by applying Young inequality repeatedly, we have

$$\| |\rho|^{*(s-1)} \|_{\ell^{\infty}(\mathbb{Z})} \leqslant \| \rho \|_{\ell^{1}(\mathbb{Z})} \| |\rho|^{*(s-2)} \|_{\ell^{\infty}(\mathbb{Z})} \leqslant \ldots \leqslant \| \rho \|_{\ell^{1}(\mathbb{Z})}^{s-1} < \infty.$$

In particular, we have that $\langle |\rho|^{*(s-1)}, |\rho| \rangle_{\ell^2(\mathbb{Z})} \leq ||\rho||_{\ell^1(\mathbb{Z})}^s < \infty$.

Third step. Thanks to the result shown in the previous step, observe first that

$$\sum_{k_2,\dots,k_s\in\mathbb{Z}} |\rho(k_2)\rho(k_2-k_3)\rho(k_3-k_4)\dots\rho(k_{s-1}-k_s)\rho(k_s)| = \langle |\rho|^{*(s-1)}, |\rho| \rangle_{\ell^2(\mathbb{Z})} < \infty.$$

Hence, one can apply dominated convergence to get, as $n \to \infty$, that

$$\frac{\sigma_{n}^{s}}{2^{s-1}(s-1)! n} \kappa_{s}(F_{n})$$

$$= \frac{1}{n} \sum_{k_{1}=0}^{n-1} \sum_{k_{2},\dots,k_{s}=-k_{1}}^{n-1-k_{1}} \rho(k_{2})\rho(k_{2}-k_{3})\rho(k_{3}-k_{4})\dots\rho(k_{s-1}-k_{s})\rho(k_{s})$$

$$= \sum_{k_{2},\dots,k_{s}\in\mathbb{Z}} \rho(k_{2})\rho(k_{2}-k_{3})\rho(k_{3}-k_{4})\dots\rho(k_{s-1}-k_{s})\rho(k_{s})$$

$$\times \left[1 \wedge \left(1 - \frac{\max\{k_{2},\dots,k_{s}\}}{n}\right) - 0 \vee \left(\frac{\min\{k_{2},\dots,k_{s}\}}{n}\right)\right] \mathbf{1}_{\{|k_{2}|< n,\dots,|k_{s}|< n\}}$$

$$\rightarrow \sum_{k_{2},\dots,k_{s}\in\mathbb{Z}} \rho(k_{2})\rho(k_{2}-k_{3})\rho(k_{3}-k_{4})\dots\rho(k_{s-1}-k_{s})\rho(k_{s}) = \langle \rho^{*(s-1)},\rho \rangle_{\ell^{2}(\mathbb{Z})}.$$
(9.109)

Since $\sigma_n \sim \sqrt{2n} \|\rho\|_{\ell^2(\mathbb{Z})}$ as $n \to \infty$, the desired conclusion follows.

Corollary 9.6 Let F_n be as above (with $H < \frac{1}{2}$), let $N \sim \mathcal{N}(0,1)$, and let ρ be given by (5.68). Then, for all $x \in \mathbb{R}$, we have

$$\sqrt{n} \left(P(F_n \leqslant x) - P(N \leqslant x) \right) \to \frac{\langle \rho^{*2}, \rho \rangle_{\ell^2(\mathbb{Z})}}{3 \|\rho\|_{\ell^2(\mathbb{Z})}^2} (1 - x^2) e^{-\frac{x^2}{2}} \quad as \ n \to \infty.$$

In particular, we deduce that there exists $d_H > 0$ such that

$$\frac{d_H}{\sqrt{n}} \leqslant \left| P(F_n \leqslant 0) - P(N \leqslant 0) \right| \leqslant d_{TV}(F_n, N), \quad n \geqslant 1. \tag{9.110}$$

Proof. The desired conclusion follows immediately by combining Propositions 9.3 and 9.5.

By paying closer attention to the used estimates, one may actually show that (9.110) holds true for any $H < \frac{5}{8}$ (not only $H < \frac{1}{2}$). See [29, Theorem 9.5.1] for the details.

To go further. The paper [27] contains several other examples of application of Theorem 9.2 and Proposition 9.3. In the reference [4], one shows that the deterministic sequence

$$\max\{E[F_n^3], E[F_n^4] - 3\}, \quad n \geqslant 1,$$

completely characterizes the rate of convergence (with respect to smooth distances) in CLTs involving chaotic random variables.

10 An extension to the Poisson space (following the invited talk by Giovanni Peccati)

Let $B = (B_t)_{t \ge 0}$ be a Brownian motion, let F be any centered element of $\mathbb{D}^{1,2}$ and let $N \sim \mathcal{N}(0,1)$. We know from Theorem 5.2 that

$$d_{TV}(F, N) \le 2 E[|1 - \langle DF, -DL^{-1}F \rangle_{L^2(\mathbb{R}_+)}|].$$
 (10.111)

The aim of this section is to explain how to deduce inequalities of the type (10.111), when F is a regular functional of a *Poisson measure* η and when the target law N is either Gaussian or Poisson. We first need to introduce the basic concepts in this framework.

Poisson measure. In what follows, we shall use the symbol $Po(\lambda)$ to indicate the Poisson distribution of parameter $\lambda > 0$ (that is, $\mathcal{P}_{\lambda} \sim Po(\lambda)$ if and only if $P(\mathcal{P}_{\lambda} = k) = e^{-\lambda} \frac{\lambda^k}{k!}$ for all $k \in \mathbb{N}$), with the convention that $Po(0) = \delta_0$ (Dirac mass at 0). Set $A = \mathbb{R}^d$ with $d \ge 1$, let \mathcal{A} be the Borel σ -field on A, and let μ be a positive, σ -finite and atomless measure over (A, \mathcal{A}) . We set $\mathcal{A}_{\mu} = \{B \in \mathcal{A} : \mu(B) < \infty\}$.

Definition 10.1 A Poisson measure η with control μ is an object of the form $\{\eta(B)\}_{B\in\mathcal{A}_{\mu}}$ with the following features:

- (1) for all $B \in \mathcal{A}_{\mu}$, we have $\eta(B) \sim Po(\mu(B))$.
- (2) for all $B, C \in \mathcal{A}_{\mu}$ with $B \cap C \neq \emptyset$, the random variables $\eta(B)$ and $\eta(C)$ are independent. Also, we note $\widehat{\eta}(B) = \eta(B) - \mu(B)$.
- **Remark 10.2** 1. As a simple example, note that for d=1 and $\mu=\lambda \times \text{Leb}$ (with 'Leb' the Lebesgue measure) the process $\{\eta([0,t])\}_{t\geqslant 0}$ is nothing but a Poisson process with intensity λ .
 - 2. Let μ be a σ -finite atomless measure over (A, \mathcal{A}) , and observe that this implies that there exists a sequence of disjoint sets $\{A_j: j \geq 1\} \subset \mathcal{A}_{\mu}$ such that $\cup_j A_j = A$. For every j = 1, 2, ... belonging to the set J_0 of those indices such that $\mu(A_j) > 0$ consider the following objects: $X^{(j)} = \{X_k^{(j)}: k \geq 1\}$ is a sequence of i.i.d. random variables with values in A_j and with common distribution $\frac{\mu_{|A_j|}}{\mu(A_j)}$; P_j is a Poisson random variable with parameter $\mu(A_j)$. Assume moreover that : (i) $X^{(j)}$ is independent of $X^{(k)}$ for every $k \neq j$, (ii) P_j is independent of P_k for every $k \neq j$, and (iii) the classes $\{X^{(j)}\}$ and $\{P_j\}$ are independent. Then, it is a straightforward computation to verify that the random point measure

$$\eta(\cdot) = \sum_{j \in J_0} \sum_{k=1}^{P_j} \delta_{X_k^{(j)}}(\cdot),$$

where δ_x indicates the Dirac mass at x and $\sum_{k=1}^{0} = 0$ by convention, is a Poisson random measure with control μ . See e.g. [47, Section 1.7].

Multiple integrals and chaotic expansion. As a preliminary remark, let us observe that $E[\widehat{\eta}(B)] = 0$ and $E[\widehat{\eta}(B)^2] = \mu(B)$ for all $B \in \mathcal{A}_{\mu}$. For any $q \ge 1$, set $L^2(\mu^q) = L^2(A^q, \mathcal{A}^q, \mu^q)$. We want to appropriately define

$$I_q(f) = \int_{A^q} f(x_1, \dots, x_q) \widehat{\eta}(dx_1) \dots \widehat{\eta}(dx_q)$$

when $f \in L^2(\mu^q)$. To reach our goal, we proceed in a classical way. We first consider the subset $\mathcal{E}(\mu^q)$ of simple functions, which is defined as

$$\mathcal{E}(\mu^q) = \operatorname{span} \left\{ \mathbf{1}_{B_1} \otimes \ldots \otimes \mathbf{1}_{B_q}, \text{ with } B_1, \ldots, B_q \in \mathcal{A}_{\mu} \text{ such that } B_i \cap B_j = \emptyset \text{ for all } i \neq j \right\}.$$

When $f = \mathbf{1}_{B_1} \otimes \ldots \otimes \mathbf{1}_{B_q}$ with $B_1, \ldots, B_q \in \mathcal{A}_{\mu}$ such that $B_i \cap B_j = \emptyset$ for all $i \neq j$, we naturally set

$$I_q(f) := \widehat{\eta}(B_1) \dots \widehat{\eta}(B_q) = \int_{A^q} f(x_1, \dots, x_q) \widehat{\eta}(dx_1) \dots \widehat{\eta}(dx_q).$$

(For such a simple function f, note that the right-hand side in the previous formula makes perfectly sense by considering $\widehat{\eta}$ as a signed measure.) We can extend by linearity the definition of $I_q(f)$ to any $f \in \mathcal{E}(\mu^q)$. It is then a simple computation to check that

$$E[I_p(f)I_q(g)] = p!\delta_{p,q}\langle \widetilde{f}, \widetilde{g}\rangle_{L^2(\mu^p)}$$

for all $f \in \mathcal{E}(\mu^p)$ and $g \in \mathcal{E}(\mu^q)$, with \widetilde{f} (resp. \widetilde{g}) the symmetrization of f (resp. g) and $\delta_{p,q}$ the Kronecker symbol. Since $\mathcal{E}(\mu^q)$ is dense in $L^2(\mu^q)$ (it is precisely here that the fact that μ has no atom is crucial!), we can define $I_q(f)$ by isometry to any $f \in L^2(\mu^q)$. Relevant properties of $I_q(f)$ include $E[I_q(f)] = 0$, $I_q(f) = I_q(\widetilde{f})$ and (importantly!) the fact that $I_q(f)$ is a true multiple integral when $f \in \mathcal{E}(\mu^q)$.

Definition 10.3 Fix $q \ge 1$. The set of random variables of the form $I_q(f)$ is called the qth Poisson-Wiener chaos.

In this framework, we have an analogue of the chaotic decomposition (4.39) – see e.g. [43, Corollary 10.0.5] for a proof.

Theorem 10.4 For all $F \in L^2(\sigma\{\eta\})$ (that is, for all random variable F which is square integrable and measurable with respect to η), we have

$$F = E[F] + \sum_{q=1}^{\infty} I_q(f_q), \tag{10.112}$$

where the kernels f_q are $(\mu^q$ -a.e.) symmetric elements of $L^2(\mu^q)$ and are uniquely determined by F.

Multiplication formula and contractions. When $f \in \mathcal{E}(\mu^p)$ and $g \in \mathcal{E}(\mu^q)$ are symmetric, we define, for all $r = 0, \dots, p \land q$ and $l = 0, \dots, r$:

$$f \star_{r}^{l} g(x_{1}, \dots, x_{p+q-r-l})$$

$$= \int_{A^{l}} f(y_{1}, \dots, y_{l}, x_{1}, \dots, x_{r-l}, x_{r-l+1}, \dots, x_{p-l}) g(y_{1}, \dots, y_{l}, x_{1}, \dots, x_{r-l}, x_{p-l+1}, \dots, x_{p+q-r-l})$$

$$\times dy_{1} \dots dy_{l}.$$

We then have the following product formula, compare with (4.43).

Theorem 10.5 (Product formula) Let $p, q \ge 1$ and let $f \in \mathcal{E}(\mu^p)$ and $g \in \mathcal{E}(\mu^q)$ be symmetric. Then

$$I_p(f)I_q(g) = \sum_{r=0}^{p \wedge q} r! \binom{p}{r} \binom{q}{r} \sum_{l=0}^r \binom{r}{l} I_{p+q-r-l} \left(\widetilde{f \star_r^l g} \right).$$

Proof. Recall that, when dealing with functions in $\mathcal{E}(\mu^p)$, $I_p(f)$ is a true multiple integral (by seeing $\hat{\eta}$ as a signed measure). We deduce

$$I_p(f)I_q(g) = \int_{A^{p+q}} f(x_1, \dots, x_p)g(y_1, \dots, y_q)\widehat{\eta}(dx_1) \dots \widehat{\eta}(dx_p)\widehat{\eta}(dy_1) \dots \widehat{\eta}(dy_q).$$

By definition of f (the same applies for g), we have that $f(x_1, ..., x_p) = 0$ when $x_i = x_j$ for some $i \neq j$. Consider $r = 0, ..., p \land q$, as well as pairwise disjoint indices $i_1, ..., i_r \in \{1, ..., p\}$ and pairwise disjoint indices $j_1, ..., j_r \in \{1, ..., q\}$. Set $\{k_1, ..., k_{p-r}\} = \{1, ..., p\} \setminus \{i_1, ..., i_r\}$ and $\{l_1, ..., l_{q-r}\} = \{1, ..., q\} \setminus \{j_1, ..., j_r\}$. We have, since μ is atomless and using $\widehat{\eta}(dx) = \eta(dx) - \mu(dx)$,

$$\int_{A^{p+q}} f(x_1, \dots, x_p) g(y_1, \dots, y_q) \mathbf{1}_{\{x_{i_1} = y_{j_1}, \dots, x_{i_r} = y_{j_r}\}} \widehat{\eta}(dx_1) \dots \widehat{\eta}(dx_p) \widehat{\eta}(dy_1) \dots \widehat{\eta}(dy_q)
= \int_{A^{p+q-2r}} f(x_{k_1}, \dots, x_{k_{p-r}}, x_{i_1}, \dots, x_{i_r}) g(y_{l_1}, \dots, y_{l_{q-r}}, x_{i_1}, \dots, x_{i_r})
\times \widehat{\eta}(dx_{k_1}) \dots \widehat{\eta}(dx_{k_{p-r}}) \widehat{\eta}(dy_{l_1}) \dots \widehat{\eta}(dy_{l_{q-r}}) \eta(dx_{i_1}) \dots \eta(dx_{i_r})
= \int_{A^{p+q-2r}} f(x_1, \dots, x_{p-r}, a_1, \dots, a_r) g(y_1, \dots, y_{q-r}, a_1, \dots, a_r)
\times \widehat{\eta}(dx_1) \dots \widehat{\eta}(dx_{p-r}) \widehat{\eta}(dy_1) \dots \widehat{\eta}(dy_{q-r}) \eta(da_1) \dots \eta(da_r).$$

By decomposing over the hyperdiagonals $\{x_i = y_j\}$, we deduce that

$$I_{p}(f)I_{q}(g) = \sum_{r=0}^{p \wedge q} r! \binom{p}{r} \binom{q}{r} \int_{A^{p+q-2r}} f(x_{1}, \dots, x_{p-r}, a_{1}, \dots, a_{r}) g(y_{1}, \dots, y_{q-r}, a_{1}, \dots, a_{r}) \times \widehat{\eta}(dx_{1}) \dots \widehat{\eta}(dx_{p-r}) \widehat{\eta}(dy_{1}) \dots \widehat{\eta}(dy_{q-r}) \eta(da_{1}) \dots \eta(da_{r}),$$

and we get the desired conclusion by using the relationship

$$\eta(da_1)\dots\eta(da_r) = (\widehat{\eta}(da_1) + \mu(da_1))\dots(\widehat{\eta}(da_r) + \mu(da_r)).$$

Malliavin operators. Each time we deal with a random element F of $L^2(\{\sigma(\eta)\})$, in what follows we always consider its chaotic expansion (10.112).

Definition 10.6 1. Set $Dom D = \{F \in L^2(\sigma\{\eta\}) : \sum qq! \|f_q\|_{L^2(\mu^q)}^2 < \infty\}$. If $F \in Dom D$, we set

$$D_t F = \sum_{q=1}^{\infty} q I_{q-1}(f_q(\cdot, t)), \quad t \in A.$$

The operator D is called the Malliavin derivative.

2. Set $Dom L = \{F \in L^2(\sigma\{\eta\}) : \sum q^2 q! \|f_q\|_{L^2(\mu^q)}^2 < \infty\}$. If $F \in Dom L$, we set

$$LF = -\sum_{q=1}^{\infty} qI_q(f_q).$$

The operator L is called the generator of the Ornstein-Uhlenbeck semigroup. 3. If $F \in L^2(\sigma\{\eta\})$, we set

$$L^{-1}F = -\sum_{q=1}^{\infty} \frac{1}{q} I_q(f_q).$$

The operator L^{-1} is called the pseudo-inverse of L.

It is readily checked that $LL^{-1}F = F - E[F]$ for $F \in L^2(\sigma\{\eta\})$. Moreover, proceeding mutatis mutandis as in the proof of Theorem 4.15, we get the following result.

Proposition 10.7 Let $F \in L^2(\sigma\{\eta\})$ and let $G \in Dom D$. Then

$$Cov(F,G) = E[\langle DG, -DL^{-1}F \rangle_{L^{2}(\mu)}].$$
(10.113)

The operator D does not satisfy the chain rule. Instead, it admits an 'add-one cost' representation which plays an identical role.

Theorem 10.8 (Nualart, Vives, 1990; see [39]) Let $F \in Dom D$. Since F is measurable with respect to η , we can view it as $F = F(\eta)$ with a slight abuse of notation. Then

$$D_t F = F(\eta + \delta_t) - F(\eta), \quad t \in A, \tag{10.114}$$

where δ_t stands for the Dirac mass at t.

Proof. By linearity and approximation, it suffices to prove the claim for $F = I_q(f)$, with $q \ge 1$ and $f \in \mathcal{E}(\mu^q)$ symmetric. In this case, we have

$$F(\eta + \delta_t) = \int_{A^q} f(x_1, \dots, x_q) (\widehat{\eta}(dx_1) + \delta_t(dx_1)) \dots (\widehat{\eta}(dx_q) + \delta_t(dx_q)).$$

Let us expand the integrator. Each member of such an expansion such that there is strictly more than one Dirac mass in the resulting expression gives a contribution equal to zero, since f vanishes on diagonals. We therefore deduce that

$$F(\eta + \delta_t) = F(\eta) + \sum_{l=1}^q \int_{A^q} f(x_1, \dots, x_{l-1}, t, x_{l+1}, \dots, x_q) \widehat{\eta}(dx_1) \dots \widehat{\eta}(dx_{l-1}) \widehat{\eta}(dx_{l+1}) \dots \widehat{\eta}(dx_q)$$

$$= F(\eta) + qI_{q-1}(f(t, \cdot)) \text{ by symmetry of } f$$

$$= F(\eta) + D_t F.$$

As an immediate corollary of the previous theorem, we get the formula

$$D_t(F^2) = (F + D_t F)^2 - F^2 = 2F D_t F + (D_t F)^2, \quad t \in A,$$

which shows how D is far from satisfying the chain rule (4.47).

Gaussian approximation. It happens that it is the following distance which is appropriate in our framework.

Definition 10.9 The Wasserstein distance between the laws of two real-valued random variables Y and Z is defined by

$$d_W(Y,Z) = \sup_{h \in \text{Lip}(1)} |E[h(Y)] - E[h(Z)]|, \tag{10.115}$$

where Lip(1) stands for the set of Lipschitz functions $h: \mathbb{R} \to \mathbb{R}$ with constant 1.

Since we are here dealing with Lipschitz functions h, we need a suitable version of the Stein's lemma. Compare with Lemma 3.1.

Lemma 10.10 (Stein, 1972; see [50]) Suppose $h : \mathbb{R} \to \mathbb{R}$ is a Lipschitz constant with constant 1. Let $N \sim \mathcal{N}(0,1)$. Then, there exists a solution to the equation

$$f'(x) - xf(x) = h(x) - E[h(N)], \quad x \in \mathbb{R}$$

that satisfies $||f'||_{\infty} \leqslant \sqrt{\frac{2}{\pi}}$ and $||f''||_{\infty} \leqslant 2$.

Proof. Let us recall that, according to Rademacher's theorem, a function which is Lipschitz continuous on \mathbb{R} is almost everywhere differentiable. Let $f: \mathbb{R} \to \mathbb{R}$ be the (well-defined!) function given by

$$f(x) = -\int_0^\infty \frac{e^{-t}}{\sqrt{1 - e^{-2t}}} E[h(e^{-t}x + \sqrt{1 - e^{-2t}}N)N]dt.$$
 (10.116)

By dominated convergence we have that $f_h \in \mathcal{C}^1$ with

$$f'(x) = -\int_0^\infty \frac{e^{-2t}}{\sqrt{1 - e^{-2t}}} E[h'(e^{-t}x + \sqrt{1 - e^{-2t}}N)N]dt.$$

We deduce, for any $x \in \mathbb{R}$,

$$|f'(x)| \le E|N| \int_0^\infty \frac{e^{-2t}}{\sqrt{1 - e^{-2t}}} dt = \sqrt{\frac{2}{\pi}}.$$
 (10.117)

Now, let $F: \mathbb{R} \to \mathbb{R}$ be the function given by

$$F(x) = \int_0^\infty E[h(N) - h(e^{-t}x + \sqrt{1 - e^{-2t}}N)]dt, \quad x \in \mathbb{R}.$$

Observe that F is well-defined since $h(N) - h(e^{-t}x + \sqrt{1 - e^{-2t}}N)$ is integrable due to

$$|h(N) - h(e^{-t}x + \sqrt{1 - e^{-2t}}N)| \leq e^{-t}|x| + (1 - \sqrt{1 - e^{-2t}})|N|$$

$$\leq e^{-t}|x| + e^{-2t}|N|,$$

where the last inequality follows from $1 - \sqrt{1 - u} = u/(\sqrt{1 - u} + 1) \le u$ if $u \in [0, 1]$. By dominated convergence, we immediately see that F is differentiable with

$$F'(x) = -\int_0^\infty e^{-t} E[h'(e^{-t}x + \sqrt{1 - e^{-2t}}N)]dt.$$

By integrating by parts, we see that F'(x) = f(x). Moreover, by using the notation introduced in Section 4, we can write

$$f'(x) - xf(x)$$

$$= LF(x), \text{ by decomposing in Hermite polynomials, since } LH_q = -qH_q = H_q'' - XH_q''$$

$$= -\int_0^\infty LP_t h(x)dt, \text{ since } F(x) = \int_0^\infty \left(E[h(N)] - P_t h(x) \right) dt$$

$$= -\int_0^\infty \frac{d}{dt} P_t h(x) dt$$

$$= P_0 h(x) - P_\infty h(x) = h(x) - E[h(N)].$$

This proves the claim for $||f'||_{\infty}$. The claim for $||f''||_{\infty}$ is a bit more difficult to achieve; we refer to Stein [50, pp. 25-28] to keep the length of this survey within bounds.

We can now derive a bound for the Gaussian approximation of any centered element F belonging to Dom D, compare with (10.111).

Theorem 10.11 (Peccati, Solé, Taqqu, Utzet, 2010; see [42]) Consider $F \in \text{Dom}D$ with E[F] = 0. Then, with $N \sim \mathcal{N}(0, 1)$,

$$d_W(F,N) \leqslant \sqrt{\frac{2}{\pi}} E\left[\left| 1 - \langle DF, -DL^{-1}F \rangle_{L^2(\mu)} \right| \right] + E\left[\int_A (D_t F)^2 |D_t L^{-1}F| \mu(dt) \right].$$

Proof. Let $h \in \text{Lip}(1)$ and let f be the function of Lemma 10.10. Using (10.114) and a Taylor's formula, we can write

$$D_t f(F) = f(F + D_t F) - f(F) = f'(F)D_t F + R(t),$$

with $|R(t)| \leq \frac{1}{2} ||f''||_{\infty} (D_t F)^2 \leq (D_t F)^2$. We deduce, using (10.113) as well,

$$E[h(F)] - E[h(N)] = E[f'(F)] - E[Ff(F)] = E[f'(F)] - E[\langle Df(F), -DL^{-1}F \rangle_{L^{2}(\mu)}]$$
$$= E[f'(F)(1 - \langle DF, -DL^{-1}F \rangle_{L^{2}(\mu)})] + \int_{\Lambda} (-D_{t}L^{-1}F)R(t)\mu(dt).$$

Consequently, since $||f'||_{\infty} \leqslant \sqrt{\frac{2}{\pi}}$,

$$\begin{split} d_W(F,N) &= \sup_{h \in \text{Lip}(1)} |E[h(F)] - E[h(N)]| \\ &\leqslant \sqrt{\frac{2}{\pi}} E\left[\left| 1 - \langle DF, -DL^{-1}F \rangle_{L^2(\mu)} \right| \right] + E\left[\int_A (D_t F)^2 |D_t L^{-1}F| \mu(dt) \right]. \end{split}$$

Poisson approximation. To conclude this section, we will prove a very interesting results, which may be seen as a Poisson counterpart of Theorem 10.11.

Theorem 10.12 (Peccati, 2012; see [40]) Let $F \in DomD$ with $E[F] = \lambda > 0$ and F taking its values in \mathbb{N} . Let $\mathcal{P}_{\lambda} \sim Po(\lambda)$. Then,

$$\sup_{C \subset \mathbb{N}} \left| P(F \in C) - P(\mathcal{P}_{\lambda} \in C) \right|$$

$$\leq \frac{1 - e^{-\lambda}}{\lambda} E|\lambda - \langle DF, -DL^{-1}F \rangle_{L^{2}(\mu)}| + \frac{1 - e^{-\lambda}}{\lambda^{2}} E \int |D_{t}F(D_{t}F - 1)D_{t}L^{-1}F|\mu(dt).$$
(10.118)

Just as a mere illustration, consider the case where $F = \eta(B) = I_1(\mathbf{1}_B)$ with $B \in \mathcal{A}_{\mu}$. We then have $DF = -DL^{-1}F = \mathbf{1}_B$, so that $\langle DF, -DL^{-1}F \rangle_{L^2(\mu)} = \int \mathbf{1}_B d\mu = \mu(B)$ and DF(DF - 1) = 0 a.e. The right-hand side of (10.118) is therefore zero, as it was expected since $F \sim Po(\lambda)$.

During the proof of Theorem 10.12, we shall use an analogue of Lemma 3.1 in the Poisson context, which reads as follows.

Lemma 10.13 (Chen, 1975; see [7]) Let $C \subset \mathbb{N}$, let $\lambda > 0$ and let $\mathcal{P}_{\lambda} \sim Po(\lambda)$. The equation with unknown $f : \mathbb{N} \to \mathbb{R}$,

$$\lambda f(k+1) - kf(k) = \mathbf{1}_C(k) - P(\mathcal{P}_{\lambda} \in C), \quad k \in \mathbb{N}, \tag{10.119}$$

admits a unique solution such that f(0) = 0, denoted by f_C . Moreover, by setting $\Delta f(k) = f(k+1) - f(k)$, we have $\|\Delta f_C\|_{\infty} \leqslant \frac{1-e^{-\lambda}}{\lambda}$ and $\|\Delta^2 f_C\|_{\infty} \leqslant \frac{2}{\lambda} \|\Delta f_C\|_{\infty}$.

Proof. We only provide a proof for the bound on Δf_C ; the estimate on $\Delta^2 f_C$ is proved e.g. by Daly in [9]. Multiplying both sides of (10.119) by $\lambda^k/k!$ and summing up yields that, for every $k \ge 1$,

$$f_C(k) = \frac{(k-1)!}{\lambda^k} \sum_{r=0}^{k-1} \frac{\lambda^r}{r!} [\mathbf{1}_C(r) - P(\mathcal{P}_{\lambda} \in C)]$$
 (10.120)

$$= \sum_{j \in C} f_{\{j\}}(k) \tag{10.121}$$

$$= -f_{C^c}(k)$$
 (10.122)

$$= -\frac{(k-1)!}{\lambda^k} \sum_{r=k}^{\infty} \frac{\lambda^r}{r!} [\mathbf{1}_C(r) - P(\mathcal{P}_{\lambda} \in C)], \qquad (10.123)$$

where C^c denotes the complement of C in \mathbb{N} . (Identity (10.121) comes from the additivity property of $C \mapsto f_C$, identity (10.122) is because $f_{\mathbb{N}} \equiv 0$ and identity (10.122) is due to

$$\sum_{r=0}^{\infty} \frac{\lambda^r}{r!} [\mathbf{1}_C(r) - P(\mathcal{P}_{\lambda} \in C)] = E[\mathbf{1}_C(\mathcal{P}_{\lambda}) - E[\mathbf{1}_C(\mathcal{P}_{\lambda})]] = 0.$$

Since $f_C(k) - f_C(k+1) = f_{C^c}(k+1) - f_{C^c}(k)$ (due to (10.122)), it is sufficient to prove that, for every $k \ge 1$ and every $C \subset \mathbb{N}$, $f_C(k+1) - f_C(k) \le (1 - e^{-\lambda})/\lambda$. One has the following fact: for every $j \ge 1$ the mapping $k \mapsto f_{\{j\}}(k)$ is negative and decreasing for k = 1, ..., j and positive and decreasing for $k \ge j + 1$. Indeed, we use (10.120) to deduce that, if $1 \le k \le j$,

$$f_{\{j\}}(k) = -e^{-\lambda} \frac{\lambda^j}{j!} \sum_{r=1}^k \lambda^{-r} \frac{(k-1)!}{(k-r)!} \quad \text{(which is negative and decreasing in } k),$$

whereas (10.123) implies that, if $k \ge j + 1$,

$$f_{\{j\}}(k) = e^{-\lambda} \frac{\lambda^j}{j!} \sum_{r=0}^{\infty} \lambda^r \frac{(k-1)!}{(k+r)!}$$
 (which is positive and decreasing in k).

Using (10.121), one therefore infers that $f_C(k+1) - f_C(k) \le f_{\{k\}}(k+1) - f_{\{k\}}(k)$, for every $k \ge 0$. Since

$$f_{\{k\}}(k+1) - f_{\{k\}}(k) = e^{-\lambda} \left[\sum_{r=0}^{k-1} \frac{\lambda^r}{r!k} + \sum_{r=k+1}^{\infty} \frac{\lambda^{r-1}}{r!} \right] = \frac{e^{-\lambda}}{\lambda} \left[\sum_{r=1}^k \frac{\lambda^r}{r!} \times \frac{r}{k} + \sum_{r=k+1}^{\infty} \frac{\lambda^r}{r!} \right]$$

$$\leqslant \frac{1 - e^{-\lambda}}{\lambda},$$

the proof is concluded.

We are now in a position to prove Theorem 10.12.

Proof of Theorem 10.12. The main ingredient is the following simple inequality, which is a kind of Taylor's formula: for all $k, a \in \mathbb{N}$,

$$|f(k) - f(a) - \Delta f(a)(k - a)| \le \frac{1}{2} ||\Delta^2 f||_{\infty} |(k - a)(k - a - 1)|.$$
 (10.124)

Assume for the time being that (10.124) holds true and fix $C \subset \mathbb{N}$. We have, using Lemma 10.13 and then (10.113)

$$\begin{aligned} \left| P(F \in C) - P(\mathcal{P}_{\lambda} \in C) \right| &= \left| E[\lambda f_C(F+1)] - E[Ff_C(F)] \right| \\ &= \left| \lambda E[\Delta f_C(F)] - E[(F-\lambda)f_C(F)] \right| \\ &= \left| \lambda E[\Delta f_C(F)] - E[\langle Df_C(F), -DL^{-1}F \rangle_{L^2(\mu)}] \right|. \end{aligned}$$

Now, combining (10.114) with (10.124), we can write

$$D_t f_C(F) = \Delta f_C(F) D_t F + S(t),$$

with $S(t) \leqslant \frac{1}{2} ||\Delta^2 f_C||_{\infty} |D_t F(D_t F - 1)| \leqslant \frac{1 - e^{-\lambda}}{\lambda^2} |D_t F(D_t F - 1)|$, see indeed Lemma 10.13 for the last inequality. Putting all these bounds together and since $||\Delta f_C||_{\infty} \leqslant \frac{1 - e^{-\lambda}}{\lambda}$ by Lemma 10.13, we get the desired conclusion.

So, to conclude the proof, it remains to show that (10.124) holds true. Let us first assume that $k \ge a + 2$. We then have

$$f(k) = f(a) + \sum_{j=a}^{k-1} \Delta f(j) = f(a) + \Delta f(a)(k-a) + \sum_{j=a}^{k-1} (\Delta f(j) - \Delta f(a))$$

$$= f(a) + \Delta f(a)(k-a) + \sum_{j=a}^{k-1} \sum_{l=a}^{j-1} \Delta^2 f(l) = f(a) + \Delta f(a)(k-a) + \sum_{l=a}^{k-2} \Delta^2 f(l)(k-l-1),$$

so that

$$|f(k) - f(a) - \Delta f(a)(k - a)| \le ||\Delta^2 f||_{\infty} \sum_{l=a}^{k-2} (k - l - 1) = \frac{1}{2} ||\Delta^2 f||_{\infty} (k - a)(k - a - 1)$$

that is, (10.124) holds true in this case. When k = a or k = a + 1, (10.124) is obviously true. Finally, consider the case $k \leq a - 1$. We have

$$f(k) = f(a) - \sum_{j=k}^{a-1} \Delta f(j) = f(a) + \Delta f(a)(k-a) + \sum_{j=k}^{a-1} (\Delta f(a) - \Delta f(j))$$

$$= f(a) + \Delta f(a)(k-a) + \sum_{j=k}^{a-1} \sum_{l=j}^{a-1} \Delta^2 f(l) = f(a) + \Delta f(a)(k-a) + \sum_{l=k}^{a-1} \Delta^2 f(l)(l-k+1),$$

so that

$$|f(k) - f(a) - \Delta f(a)(k - a)| \le ||\Delta^2 f||_{\infty} \sum_{l=k}^{a-1} (l - k + 1) = \frac{1}{2} ||\Delta^2 f||_{\infty} (a - k)(a - k + 1),$$

that is, (10.124) holds true in this case as well. The proof of Theorem 10.12 is done.

To go further. A multivariate extension of Theorem 10.11 can be found in [45]. The reference [41] contains several explicit applications of the tools developed in this section.

11 Fourth Moment Theorem and free probability

To conclude this survey, we shall explain how the Fourth Moment Theorem 1.3 extends in the theory of free probability, which provides a convenient framework for investigating limits of random matrices.

We start with a short introduction to free probability; we refer to [20] for a systematic presentation.

Random matrices. Let $(B_{ij})_{i,j\geqslant 1}$ be a sequence of independent standard Brownian motions, all defined on the same probability space (Ω, \mathcal{F}, P) . Set

$$A_n(t) = \left(\frac{B_{ij}(t)}{\sqrt{n}}\right)_{1 \leqslant i,j \leqslant n} \quad \text{and} \quad M_n(t) = \frac{1}{\sqrt{2}} \left(A_n(t) + A_n(t)^T\right), \tag{11.125}$$

where \cdot^T stands for the usual transpose operator. For each t, $A_n(t)$ and $M_n(t)$ both belong to \mathscr{A}_n , the set of random matrices with entries in $L^{\infty-}(\Omega) = \bigcap_{p>1} L^p(\Omega)$, that is, with all moments. On \mathscr{A}_n , let us consider the linear form $\varphi_n : \mathscr{A}_n \to \mathbb{R}$ defined by

$$\varphi_n(M) = E\left[\frac{1}{n}\text{Tr}(M)\right],\tag{11.126}$$

where E denotes the mathematical expectation associated to P, whereas $\text{Tr}(\cdot)$ stands for the usual trace operator. The space $(\mathcal{A}_n, \varphi_n)$ is the prototype of a *free tracial probability space* (also called non-commutative probability space).

Let $t_k > ... > t_1 > t_0 = 0$. A celebrated theorem by Voiculescu [53] asserts that the increments $M_n(t_1), M_n(t_2) - M_n(t_1), ..., M_n(t_k) - M_n(t_{k-1})$ are asymptotically free, meaning that

$$\varphi_n \left(Q_1(M_n(t_{i_1}) - M_n(t_{i_1-1})) \dots Q_m(M_n(t_{i_m}) - M_n(t_{i_m-1})) \right) \to 0 \quad \text{as } n \to \infty,$$

for all $m \ge 2$, all $i_1, \ldots, i_m \in \{1, \ldots, k\}$ with $i_1 \ne i_2, i_2 \ne i_3, \ldots, i_{m-1} \ne i_m$, and all real-valued polynomials Q_1, \ldots, Q_m such that $\varphi_n(Q_i(M_n(t_i) - M_n(t_{i-1}))) \to 0$ as $n \to \infty$ for each $i = 1, \ldots, m$.

Let t > 0. The no less famous Wigner theorem [54] can be formulated as follows: as $n \to \infty$, $M_n(t)$ converges in law to the semicircular law of variance t, that is, for any real-valued polynomial Q,

$$\varphi_n(Q(M_n(t))) \to \frac{1}{2\pi t} \int_{-2\sqrt{t}}^{2\sqrt{t}} Q(x) \sqrt{4t - x^2} dx.$$

In the same way as calculus provides a nice setting for studying limits of sums and classical Brownian motion provides a nice setting for studying limits of random walks, free probability provides a convenient framework for investigating limits of random matrices. In the limit, free independence replaces independence, increments have a semicircular marginal law and thus, free Brownian motion which has these properties, can be used to study random matrices $M_n(t)$ for large

Free tracial probability space. A free tracial probability space is a von Neumann algebra \mathscr{A} (that is, an algebra of operators on a real separable Hilbert space, closed under adjoint and convergence in the weak operator topology) equipped with a $trace \varphi$, that is, a unital linear functional (meaning preserving the identity) which is weakly continuous, positive (meaning $\varphi(X) \geq 0$ whenever X is a non-negative element of \mathscr{A} ; i.e. whenever $X = YY^*$ for some $Y \in \mathscr{A}$), faithful (meaning

that if $\varphi(YY^*) = 0$ then Y = 0, and tracial (meaning that $\varphi(XY) = \varphi(YX)$ for all $X, Y \in \mathcal{A}$, even though in general $XY \neq YX$).

Random variables. In a free tracial probability space, we refer to the self-adjoint elements of the algebra as $random\ variables$. Any random variable X has a law: this is the unique probability measure μ on \mathbb{R} with the same moments as X; in other words, μ is such that

$$\int_{\mathbb{R}} Q(x)d\mu(x) = \varphi(Q(X)),\tag{11.127}$$

for any real polynomial Q. (The existence and uniqueness of μ follow from the positivity of φ , see [20, Proposition 3.13].) Thus φ acts as an expectation. The φ_n in (11.126), for example, play the role of φ .

Convergence in law. We say that a sequence $(X_{1,n},\ldots,X_{k,n}), n \ge 1$, of random vectors converges in law to a random vector $(X_{1,\infty},\ldots,X_{k,\infty})$, and we write

$$(X_{1,n},\ldots,X_{k,n}) \stackrel{\text{law}}{\to} (X_{1,\infty},\ldots,X_{k,\infty}),$$

to indicate the convergence in the sense of (joint) moments, that is,

$$\lim_{n \to \infty} \varphi\left(Q(X_{1,n}, \dots, X_{k,n})\right) = \varphi\left(Q(X_{1,\infty}, \dots, X_{k,\infty})\right),\tag{11.128}$$

for any real-valued polynomial Q in k non-commuting variables.

We say that a sequence (F_n) of non-commutative stochastic processes (that is, each F_n is a one-parameter family of self-adjoint operators $F_n(t)$ in the free tracial probability space (\mathscr{A}, φ)) converges in the sense of finite-dimensional distributions to a non-commutative stochastic process F_{∞} , and we write

$$F_n \stackrel{\text{f.d.d.}}{\to} F_{\infty},$$

to indicate that, for any $k \ge 1$ and any $t_1, \ldots, t_k \ge 0$,

$$(F_n(t_1),\ldots,F_n(t_k)) \stackrel{\text{law}}{\to} (F_\infty(t_1),\ldots,F_\infty(t_k)).$$

Free independence. In the free probability setting, the notion of independence (introduced by Voiculescu in [52]) goes as follows. Let $\mathscr{A}_1, \ldots, \mathscr{A}_p$ be unital subalgebras of \mathscr{A} . Let X_1, \ldots, X_m be elements chosen from among the \mathscr{A}_i 's such that, for $1 \leq j < m$, two consecutive elements X_j and X_{j+1} do not come from the same \mathscr{A}_i and are such that $\varphi(X_j) = 0$ for each j. The subalgebras $\mathscr{A}_1, \ldots, \mathscr{A}_p$ are said to be free or freely independent if, in this circumstance,

$$\varphi(X_1 X_2 \cdots X_m) = 0. \tag{11.129}$$

Random variables are called freely independent if the unital algebras they generate are freely independent. Freeness is in general much more complicated than classical independence. For example, if X, Y are free and $m, n \ge 1$, then by (11.129),

$$\varphi((X^m - \varphi(X^m)1)(Y^n - \varphi(Y^n)1)) = 0.$$

By expanding (and using the linear property of φ), we get

$$\varphi(X^m Y^n) = \varphi(X^m)\varphi(Y^n),\tag{11.130}$$

which is what we would expect under classical independence. But, by setting $X_1 = X_3 = X - \varphi(X)1$ and $X_2 = X_4 = Y - \varphi(Y)$ in (11.129), we also have

$$\varphi((X - \varphi(X)1)(Y - \varphi(Y)1)(X - \varphi(X)1)(Y - \varphi(Y)1)) = 0.$$

By expanding, using (11.130) and the tracial property of φ (for instance $\varphi(XYX) = \varphi(X^2Y)$) we get

$$\varphi(XYXY) = \varphi(Y)^2 \varphi(X^2) + \varphi(X)^2 \varphi(Y^2) - \varphi(X)^2 \varphi(Y)^2,$$

which is different from $\varphi(X^2)\varphi(Y^2)$, which is what one would have obtained if X and Y were classical independent random variables. Nevertheless, if X,Y are freely independent, then their joint moments are determined by the moments of X and Y separately, exactly as in the classical case.

Semicircular distribution. The semicircular distribution $S(m, \sigma^2)$ with mean $m \in \mathbb{R}$ and variance $\sigma^2 > 0$ is the probability distribution

$$S(m, \sigma^2)(dx) = \frac{1}{2\pi\sigma^2} \sqrt{4\sigma^2 - (x - m)^2} \mathbf{1}_{\{|x - m| \le 2\sigma\}} dx.$$
 (11.131)

If m = 0, this distribution is symmetric around 0, and therefore its odd moments are all 0. A simple calculation shows that the even centered moments are given by (scaled) Catalan numbers: for non-negative integers k,

$$\int_{m-2\sigma}^{m+2\sigma} (x-m)^{2k} \mathcal{S}(m,\sigma^2)(dx) = C_k \sigma^{2k},$$

where $C_k = \frac{1}{k+1} {2k \choose k}$ (see, e.g., [20, Lecture 2]). In particular, the variance is σ^2 while the centered fourth moment is $2\sigma^4$. The semicircular distribution plays here the role of the Gaussian distribution. It has the following similar properties:

- 1. If $S \sim \mathcal{S}(m, \sigma^2)$ and $a, b \in \mathbb{R}$, then $aS + b \sim \mathcal{S}(am + b, a^2\sigma^2)$.
- 2. If $S_1 \sim \mathcal{S}(m_1, \sigma_1^2)$ and $S_2 \sim \mathcal{S}(m_2, \sigma_2^2)$ are freely independent, then $S_1 + S_2 \sim \mathcal{S}(m_1 + m_2, \sigma_1^2 + \sigma_2^2)$.

Free Brownian Motion. A free Brownian motion $S = \{S(t)\}_{t\geq 0}$ is a non-commutative stochastic process with the following defining characteristics:

- (1) S(0) = 0.
- (2) For $t_2 > t_1 \ge 0$, the law of $S(t_2) S(t_1)$ is the semicircular distribution of mean 0 and variance $t_2 t_1$.
- (3) For all n and $t_n > \cdots > t_2 > t_1 > 0$, the increments $S(t_1)$, $S(t_2) S(t_1)$, ..., $S(t_n) S(t_{n-1})$ are freely independent.

If we rephrase Wigner and Voiculescu theorems in our new language, we get that $M_n \stackrel{\text{f.d.d.}}{\to} S$ as $n \to \infty$, where M_n is given by (11.125) and S is a free Brownian motion. So, we may think of free Brownian motion as 'infinite-dimensional matrix-valued Brownian motion'.

Wigner integral. Let $S = \{S(t)\}_{t\geq 0}$ be a free Brownian motion. Let us quickly sketch out the construction of the Wigner integral of f with respect to S. For an indicator function $f = \mathbf{1}_{[u,v]}$, the Wigner integral of f is defined by

$$\int_0^\infty \mathbf{1}_{[u,v]}(x)dS(x) = S(v) - S(u).$$

We then extend this definition by linearity to simple functions of the form $f = \sum_{i=1}^k \alpha_i \mathbf{1}_{[u_i,v_i]}$, where $[u_i,v_i]$ are disjoint intervals of \mathbb{R}_+ . Simple computations show that

$$\varphi\left(\int_0^\infty f(x)dS(x)\right) = 0 \tag{11.132}$$

$$\varphi\left(\int_0^\infty f(x)dS(x) \times \int_0^\infty g(x)dS(x)\right) = \langle f, g \rangle_{L^2(\mathbb{R}_+)}. \tag{11.133}$$

By isometry, the definition of $\int_0^\infty f(x)dS(x)$ is extended to all $f \in L^2(\mathbb{R}_+)$, and (11.132)-(11.133) continue to hold in this more general setting.

Multiple Wigner integral. Let $S = \{S(t)\}_{t \geq 0}$ be a free Brownian motion, and let $q \geq 1$ be an integer. When $f \in L^2(\mathbb{R}^q_+)$ is real-valued, we write f^* to indicate the function of $L^2(\mathbb{R}^q_+)$ given by $f^*(t_1, \ldots, t_q) = f(t_q, \ldots, t_1)$.

Following [2], let us quickly sketch out the construction of the multiple Wigner integral of f with respect to S. Let $D^q \subset \mathbb{R}^q_+$ be the collection of all diagonals, i.e.

$$D^{q} = \{(t_{1}, \dots, t_{q}) \in \mathbb{R}^{q}_{+} : t_{i} = t_{j} \text{ for some } i \neq j\}.$$
(11.134)

For an indicator function $f = \mathbf{1}_A$, where $A \subset \mathbb{R}^q_+$ has the form $A = [u_1, v_1] \times \ldots \times [u_q, v_q]$ with $A \cap D^q = \emptyset$, the qth multiple Wigner integral of f is defined by

$$I_a(f) = (S(v_1) - S(u_1)) \dots (S(v_a) - S(u_a)).$$

We then extend this definition by linearity to simple functions of the form $f = \sum_{i=1}^k \alpha_i \mathbf{1}_{A_i}$, where $A_i = [u_1^i, v_1^i] \times \ldots \times [u_q^i, v_q^i]$ are disjoint q-dimensional rectangles as above which do not meet the diagonals. Simple computations show that

$$\varphi(I_q(f)) = 0 \tag{11.135}$$

$$\varphi(I_q(f)I_q(g)) = \langle f, g^* \rangle_{L^2(\mathbb{R}^q_+)}. \tag{11.136}$$

By isometry, the definition of $I_q(f)$ is extended to all $f \in L^2(\mathbb{R}^q_+)$, and (11.135)-(11.136) continue to hold in this more general setting. If one wants $I_q(f)$ to be a random variable, it is necessary for f to be mirror symmetric, that is, $f = f^*$ (see [16]). Observe that $I_1(f) = \int_0^\infty f(x)dS(x)$ when q = 1. We have moreover

$$\varphi(I_p(f)I_q(g)) = 0 \text{ when } p \neq q, f \in L^2(\mathbb{R}^p_+) \text{ and } g \in L^2(\mathbb{R}^q_+).$$
 (11.137)

When $r \in \{1, \ldots, p \land q\}$, $f \in L^2(\mathbb{R}^p_+)$ and $g \in L^2(\mathbb{R}^q_+)$, let us write $f \stackrel{r}{\frown} g$ to indicate the rth contraction of f and g, defined as being the element of $L^2(\mathbb{R}^{p+q-2r}_+)$ given by

$$f \stackrel{r}{\frown} g(t_1, \dots, t_{p+q-2r})$$

$$= \int_{\mathbb{R}^r_+} f(t_1, \dots, t_{p-r}, x_1, \dots, x_r) g(x_r, \dots, x_1, t_{p-r+1}, \dots, t_{p+q-2r}) dx_1 \dots dx_r.$$
(11.138)

By convention, set $f \stackrel{0}{\frown} g = f \otimes g$ as being the tensor product of f and g. Since f and g are not necessarily symmetric functions, the position of the identified variables x_1, \ldots, x_r in (11.138) is important, in contrast to what happens in classical probability. Observe moreover that

$$||f \stackrel{r}{\frown} g||_{L^{2}(\mathbb{R}^{p+q-2r})} \leq ||f||_{L^{2}(\mathbb{R}^{p}_{+})} ||g||_{L^{2}(\mathbb{R}^{q}_{+})}$$
(11.139)

by Cauchy-Schwarz, and also that $f \stackrel{q}{\frown} g = \langle f, g^* \rangle_{L^2(\mathbb{R}^q_+)}$ when p = q.

We have the following product formula (see [2, Proposition 5.3.3]), valid for any $f \in L^2(\mathbb{R}^p_+)$ and $g \in L^2(\mathbb{R}^q_+)$:

$$I_p(f)I_q(g) = \sum_{r=0}^{p \wedge q} I_{p+q-2r}(f \stackrel{r}{\frown} g).$$
 (11.140)

We deduce (by a straightforward induction) that, for any $e \in L^2(\mathbb{R}_+)$ and any $q \ge 1$,

$$U_q\left(\int_0^\infty e(x)dS_x\right) = I_q(e^{\otimes q}),\tag{11.141}$$

where $U_0=1$, $U_1=X$, $U_2=X^2-1$, $U_3=X^3-2X$, ..., is the sequence of Tchebycheff polynomials of second kind (determined by the recursion $XU_k=U_{k+1}+U_{k-1}$), $\int_0^\infty e(x)dS(x)$ is understood as a Wigner integral, and $e^{\otimes q}$ is the qth tensor product of e. This is the exact analogue of (1.10) in our context.

We are now in a position to offer a free analogue of the Fourth Moment Theorem 5.5, which reads as follows.

Theorem 11.1 (Kemp, Nourdin, Peccati, Speicher, 2011; see [16]) Fix an integer $q \ge 2$ and let $\{S_t\}_{t\ge 0}$ be a free Brownian motion. Whenever $f \in L^2(\mathbb{R}^q_+)$, set $I_q(f)$ to denote the qth multiple Wigner integrals of f with respect to S. Let $\{F_n\}_{n\ge 1}$ be a sequence of Wigner multiple integrals of the form

$$F_n = I_q(f_n),$$

where each $f_n \in L^2(\mathbb{R}_+)$ is mirror-symmetric, that is, is such that $f_n = f_n^*$. Suppose moreover that $\varphi(F_n^2) \to 1$ as $n \to \infty$. Then, as $n \to \infty$, the following two assertions are equivalent:

(i)
$$F_n \stackrel{\text{Law}}{\to} S_1 \sim \mathcal{S}(0,1);$$

(ii)
$$\varphi(F_n^4) \to 2 = \varphi(S_1^4)$$
.

Proof (following [22]). Without loss of generality and for sake of simplicity, we suppose that $\varphi(F_n^2) = 1$ for all n (instead of $\varphi(F_n^2) \to 1$ as $n \to \infty$). The proof of the implication $(i) \Rightarrow (ii)$ being trivial by the very definition of the convergence in law in a free tracial probability space, we only concentrate on the proof of $(ii) \Rightarrow (i)$.

Fix $k \ge 3$. By iterating the product formula (11.140), we can write

$$F_n^k = I_q(f_n)^k = \sum_{(r_1, \dots, r_{k-1}) \in A_{k,q}} I_{kq-2r_1-\dots-2r_{k-1}} \left(\left(\dots \left(\left(f_n \stackrel{r_1}{\frown} f_n \right) \stackrel{r_2}{\frown} f_n \right) \dots \right) \stackrel{r_{k-1}}{\frown} f_n \right),$$

where

$$A_{k,q} = \{(r_1, \dots, r_{k-1}) \in \{0, 1, \dots, q\}^{k-1} : r_2 \leqslant 2q - 2r_1, r_3 \leqslant 3q - 2r_1 - 2r_2, \dots, r_{k-1} \leqslant (k-1)q - 2r_1 - \dots - 2r_{k-2}\}.$$

By taking the φ -trace in the previous expression and taking into account that (11.135) holds, we deduce that

$$\varphi(F_n^k) = \varphi(I_q(f_n)^k) = \sum_{\substack{(r_1, \dots, r_{k-1}) \in B_{k,q}}} \left(\dots \left(\left(f_n \stackrel{r_1}{\frown} f_n \right) \stackrel{r_2}{\frown} f_n \right) \dots \right) \stackrel{r_{k-1}}{\frown} f_n, \tag{11.142}$$

with

$$B_{k,q} = \{(r_1, \dots, r_{k-1}) \in A_{k,q} : 2r_1 + \dots + 2r_{k-1} = kq\}.$$

Let us decompose $B_{k,q}$ into $C_{k,q} \cup E_{k,q}$, with $C_{k,q} = B_{k,q} \cap \{0,q\}^{k-1}$ and $E_{k,q} = B_{k,q} \setminus C_{k,q}$. We then have

$$\varphi(F_n^k) = \sum_{\substack{(r_1, \dots, r_{k-1}) \in C_{k,q}}} \left(\left(\dots \left(\left(f_n \stackrel{r_1}{\frown} f_n \right) \stackrel{r_2}{\frown} f_n \right) \dots \right) \stackrel{r_{k-1}}{\frown} f_n \right) + \sum_{\substack{(r_1, \dots, r_{k-1}) \in E_{k,q}}} \left(\left(\dots \left(\left(f_n \stackrel{r_1}{\frown} f_n \right) \stackrel{r_2}{\frown} f_n \right) \dots \right) \stackrel{r_{k-1}}{\frown} f_n \right).$$

Using the two relationships $f_n \stackrel{0}{\frown} f_n = f_n \otimes f_n$ and

$$f_n \stackrel{q}{\frown} f_n = \int_{\mathbb{R}^q_+} f_n(t_1, \dots, t_q) f_n(t_q, \dots, t_1) dt_1 \dots dt_q = ||f_n||_{L^2(\mathbb{R}^q_+)}^2 = 1,$$

it is evident that $(\dots((f_n \stackrel{r_1}{\frown} f_n) \stackrel{r_2}{\frown} f_n)\dots) \stackrel{r_{k-1}}{\frown} f_n = 1$ for all $(r_1,\dots,r_{k-1}) \in C_{k,q}$. We deduce that

$$\varphi(F_n^k) = \#C_{k,q} + \sum_{(r_1,\ldots,r_{k-1})\in E_{k,q}} \left(\left(\ldots \left(\left(f_n \stackrel{r_1}{\frown} f_n \right) \stackrel{r_2}{\frown} f_n \right) \ldots \right) \stackrel{r_{k-1}}{\frown} f_n \right).$$

On the other hand, by applying (11.142) with q = 1, we get that

$$\varphi(S_1^k) = \varphi(I_1(\mathbf{1}_{[0,1]})^k) = \sum_{(r_1,\dots,r_{k-1})\in B_{k,1}} (\dots ((\mathbf{1}_{[0,1]} \stackrel{r_1}{\frown} \mathbf{1}_{[0,1]}) \stackrel{r_2}{\frown} \mathbf{1}_{[0,1]}) \dots) \stackrel{r_{k-1}}{\frown} \mathbf{1}_{[0,1]}$$

$$= \sum_{(r_1,\dots,r_{k-1})\in B_{k,1}} 1 = \#B_{k,1}.$$

But it is clear that $C_{k,q}$ is in bijection with $B_{k,1}$ (by dividing all the r_i 's in $C_{k,q}$ by q). Consequently,

$$\varphi(F_n^k) = \varphi(S_1^k) + \sum_{(r_1, \dots, r_{k-1}) \in E_{k,q}} \left((\dots ((f_n \stackrel{r_1}{\frown} f_n) \stackrel{r_2}{\frown} f_n) \dots) \stackrel{r_{k-1}}{\frown} f_n \right).$$
 (11.143)

Now, assume that $\varphi(F_n^4) \to \varphi(S_1^4) = 2$ and let us show that $\varphi(F_n^k) \to \varphi(S_1^k)$ for all $k \geqslant 3$. Using that $f_n = f_n^*$, observe that

$$f_{n} \stackrel{r}{\frown} f_{n}(t_{1}, \dots, t_{2q-2r})$$

$$= \int_{\mathbb{R}^{r}_{+}} f_{n}(t_{1}, \dots, t_{q-r}, s_{1}, \dots, s_{r}) f_{n}(s_{r}, \dots, s_{1}, t_{q-r+1}, \dots, t_{2q-2r}) ds_{1} \dots ds_{r}$$

$$= \int_{\mathbb{R}^{r}_{+}} f_{n}(s_{r}, \dots, s_{1}, t_{q-r}, \dots, t_{1}) f_{n}(t_{2q-2r}, \dots, t_{q-r+1}, s_{1}, \dots, s_{r}) ds_{1} \dots ds_{r}$$

$$= f_{n} \stackrel{r}{\frown} f_{n}(t_{2q-2r}, \dots, t_{1}) = (f_{n} \stackrel{r}{\frown} f_{n})^{*}(t_{1}, \dots, t_{2q-2r}),$$

that is, $f_n \stackrel{r}{\frown} f_n = (f_n \stackrel{r}{\frown} f_n)^*$. On the other hand, the product formula (11.140) leads to $F_n^2 = \sum_{r=0}^q I_{2q-2r}(f_n \stackrel{r}{\frown} f_n)$. Since two multiple integrals of different orders are orthogonal (see (11.137)), we deduce that

$$\varphi(F_n^4) = \|f_n \otimes f_n\|_{L^2(\mathbb{R}^{2q}_+)}^2 + (\|f_n\|_{L^2(\mathbb{R}^q_+)}^2)^2 + \sum_{r=1}^{q-1} \langle f_n \stackrel{r}{\frown} f_n, (f_n \stackrel{r}{\frown} f_n)^* \rangle_{L^2(\mathbb{R}^{2q-2r}_+)}^{} \\
= 2\|f_n\|_{L^2([0,1]^q)}^4 + \sum_{r=1}^{q-1} \|f_n \stackrel{r}{\frown} f_n\|_{L^2(\mathbb{R}^{2q-2r}_+)}^2 = 2 + \sum_{r=1}^{q-1} \|f_n \stackrel{r}{\frown} f_n\|_{L^2(\mathbb{R}^{2q-2r}_+)}^2. \quad (11.144)$$

Using that $\varphi(F_n^4) \to 2$, we deduce that

$$||f_n \stackrel{r}{\frown} f_n||_{L^2(\mathbb{R}^{2q-2r}_+)}^2 \to 0 \quad \text{for all } r = 1, \dots, q-1.$$
 (11.145)

Fix $(r_1, \ldots, r_{k-1}) \in E_{k,q}$ and let $j \in \{1, \ldots, k-1\}$ be the smallest integer such that $r_j \in \{1, \ldots, q-1\}$. Then:

Therefore, we deduce from (11.143) that $\varphi(F_n^k) \to \varphi(S_1^k)$, which is the desired conclusion and concludes the proof of the theorem.

During the proof of Theorem 11.1, we actually showed (see indeed (11.144)) that the two assertions (i)-(ii) are both equivalent to a third one, namely

(iii):
$$||f_n \stackrel{r}{\frown} f_n||_{L^2(\mathbb{R}^{2q-2r}_+)}^2 \to 0 \text{ for all } r = 1, \dots, q-1.$$

Combining (iii) with Corollary 5.5, we immediately deduce an interesting transfer principle for translating results between the classical and free chaoses.

Corollary 11.2 Fix an integer $q \ge 2$, let $\{B_t\}_{t\ge 0}$ be a standard Brownian motion and let $\{S_t\}_{t\ge 0}$ be a free Brownian motion. Whenever $f \in L^2(\mathbb{R}^q_+)$, we write $I_q^B(f)$ (resp. $I_q^S(f)$) to indicate the qth multiple Wiener integrals of f with respect to B (resp. S). Let $\{f_n\}_{n\ge 1} \subset L^2(\mathbb{R}^q_+)$ be a sequence of symmetric functions and let $\sigma > 0$ be a finite constant. Then, as $n \to \infty$, the following two assertions hold true.

- (i) $E[I_q^B(f_n)] \to q!\sigma^2$ if and only if $\varphi(I_q^S(f_n)^2) \to \sigma^2$.
- (ii) If the asymptotic relations in (i) are verified, then $I_q^B(f_n) \stackrel{\text{law}}{\to} \mathcal{N}(0, q!\sigma^2)$ if and only if $I_q^S(f_n) \stackrel{\text{law}}{\to} \mathcal{S}(0, \sigma^2)$.

To go further. A multivariate version of Theorem 11.1 (free counterpart of Theorem 6.2) can be found in [33]. In [28] (resp. [11]), one exhibits a version of Theorem 11.1 in which the semicircular law in the limit is replaced by the free Poisson law (resp. the so-called tetilla law). An extension of Theorem 11.1 in the context of the q-Brownian motion (which is an interpolation between the standard Brownian motion corresponding to q = 1 and the free Brownian motion corresponding to q = 0) is given in [13].

References

- [1] B. Bercu, I. Nourdin and M.S. Taqqu (2010). Almost sure central limit theorems on the Wiener space. Stoch. Proc. Appl. 120, no. 9, 1607-1628.
- [2] P. Biane and R. Speicher (1998). Stochastic analysis with respect to free Brownian motion and analysis on Wigner space. *Probab. Theory Rel. Fields* **112**, 373-409.
- [3] E. Bolthausen (1984). An estimate of the remainder in a combinatorial central limit theorem. Z. Wahrscheinlichkeitstheorie verw. Gebiete 66, 379-386.
- [4] H. Biermé, A. Bonami, I. Nourdin and G. Peccati (2011). Optimal Berry-Esseen rates on the Wiener space: the barrier of third and fourth cumulants. Preprint.
- [5] J.-C. Breton and I. Nourdin (2008). Error bounds on the non-normal approximation of Hermite power variations of fractional Brownian motion. *Electron. Comm. Probab.* **13**, 482-493 (electronic).
- [6] P. Breuer and P. Major (1983). Central limit theorems for non-linear functionals of Gaussian fields. J. Mult. Anal. 13, 425-441.
- [7] L.H.Y. Chen (1975). Poisson approximation for dependent trials. Ann. Probab. 3, no. 3, 534-545
- [8] L.H.Y. Chen, L. Goldstein and Q.-M. Shao (2010). *Normal Approximation by Stein's Method*. Probability and Its Applications, Springer.
- [9] F. Daly (2008). Upper bounds for Stein-type operators. *Electronic Journal of Probability* **13**(20), 566-587 (Electronic).
- [10] L. Decreusefond, D. Nualart (2008). Hitting times for Gaussian processes. *Ann. Probab.* **36** (1), 319-330.
- [11] A. Deya and I. Nourdin (2012). Convergence of Wigner integrals to the tetilla law. ALEA 9, 101-127.
- [12] A. Deya and I. Nourdin (2011). Invariance principles for homogeneous sums of free random variables. Preprint.
- [13] A. Deya, S. Noreddine and I. Nourdin (2012). Fourth Moment Theorem and q-Brownian Chaos. Preprint.
- [14] C.G. Esseen (1956). A moment inequality with an application to the central limi theorem. Skand. Aktuarietidskr. 39, 160-170.
- [15] S.-T. Ho and L.H.Y. Chen (1978). An L_p bound for the remainder in a combinatorial Central Limit Theorem. Ann. Probab. $\mathbf{6}(2)$, 231-249.

- [16] T. Kemp, I. Nourdin, G. Peccati and R. Speicher (2011). Wigner chaos and the fourth moment. *Ann. Probab.*, in press.
- [17] V. Yu. Korolev and I.G. Shevtsova (2010). On the upper bound for the absolute constant in the Berry-Esseen inequality. *Theory of Probability and its Applications* **54**, no. 4, 638-658.
- [18] E. Mossel, R. O'Donnell and K. Oleszkiewicz (2010). Noise stability of functions with low influences: invariance and optimality. *Ann. Math.* **171**, no. 1, 295-341.
- [19] E. Nelson (1973). The free Markoff field. J. Funct. Analysis, 12, 211-227.
- [20] A. Nica and R. Speicher (2006). Lectures on the Combinatorics of Free Probability. Cambridge University Press.
- [21] S. Noreddine and I. Nourdin (2011). On the Gaussian approximation of vector-valued multiple integrals. J. Multiv. Anal. 102, no. 6, 1008-1017.
- [22] I. Nourdin (2011). Yet another proof of the Nualart-Peccati criterion. *Electron. Comm. Probab.* **16**, 467-481.
- [23] I. Nourdin (2012). Selected aspects of fractional Brownian motion, and related topics. Springer Verlag.
- [24] I. Nourdin and G. Peccati (2007). Non-central convergence of multiple integrals. *Ann. Probab.*, **37**(4), 1412-1426.
- [25] I. Nourdin and G. Peccati (2009). Stein's method on Wiener chaos. *Probab. Theory Rel. Fields* 145(1), 75-118.
- [26] I. Nourdin and G. Peccati (2010). Cumulants on the Wiener space. J. Funct. Anal. 258, 3775-3791.
- [27] I. Nourdin and G. Peccati (2010). Stein's method and exact Berry-Esséen asymptotics for functionals of Gaussian fields. *Ann. Probab.* **37**(6), 2231-2261.
- [28] I. Nourdin and G. Peccati (2011). Poisson approximations on the free Wigner chaos. Preprint.
- [29] I. Nourdin and G. Peccati (2012). Normal Approximations Using Malliavin Calculus: from Stein's Method to Universality. Cambridge Tracts in Mathematics. Cambridge University Press.
- [30] I. Nourdin, G. Peccati and M. Podolskij (2011). Quantitative Breuer-Major Theorems. Stoch. Proc. App. 121, no. 4, 793-812.
- [31] I. Nourdin, G. Peccati and G. Reinert (2010). Invariance principles for homogeneous sums: universality of Gaussian Wiener chaos. *Ann. Probab.* **38**(5), 1947-1985.
- [32] I. Nourdin, G. Peccati and A. Réveillac (2010). Multivariate normal approximation using Stein's method and Malliavin calculus. *Ann. Inst. H. Poincaré* (B) Probab. Statist. **46**(1), 45-58.
- [33] I. Nourdin, G. Peccati and R. Speicher (2011). Multidimensional semicircular limits on the free Wigner chaos. To appear in *Ascona 2011 Proceedings*.
- [34] I. Nourdin and F.G. Viens (2009). Density estimates and concentration inequalities with Malliavin calculus. *Electron. J. Probab.* **14**, 2287-2309 (electronic).

- [35] D. Nualart (2006): The Malliavin calculus and related topics. Springer-Verlag, Berlin, second edition.
- [36] D. Nualart and S. Ortiz-Latorre (2008). Central limit theorems for multiple stochastic integrals and Malliavin calculus. *Stoch. Proc. Appl.* 118 (4), 614-628.
- [37] D. Nualart and G. Peccati (2005). Central limit theorems for sequences of multiple stochastic integrals. Ann. Probab. 33 (1), 177-193.
- [38] D. Nualart and L. Quer-Sardanyons (2011). Optimal Gaussian density estimates for a class of stochastic equations with additive noise. *Infinite Dimensional Analysis*, *Quantum Probability and Related Topics* 14, no.1, 25-34.
- [39] D. Nualart and J. Vives (1990). Anticipative calculus for the Poisson process based on the Fock space. Sem. de Proba. XXIV, LNM 1426, pp. 154-165. Springer-Verlag
- [40] G. Peccati (2012). The Chen-Stein method for Poisson functionals. Preprint.
- [41] R. Lachiï£; ze-Rey and G. Peccati (2011). Fine Gaussian fluctuations on the Poisson space, I: contractions, cumulants and random geometric graphs. Preprint.
- [42] G. Peccati, J.-L. Solé, M.S. Taqqu and F. Utzet (2010). Stein's method and normal approximation of Poisson functionals. *Ann. Probab.* **38**(2), 443-478.
- [43] G. Peccati and M.S. Taqqu (2010). Wiener chaos: moments, cumulants and diagrams. Springer-Verlag.
- [44] G. Peccati and C.A. Tudor (2005). Gaussian limits for vector-valued multiple stochastic integrals. Séminaire de Probabilités XXXVIII, LNM 1857. Springer-Verlag, pp. 247-262.
- [45] G. Peccati and C. Zheng (2010). Multidimensional Gaussian fluctuations on the Poisson space. *Electron. J. Probab.* **15**, paper 48, 1487-1527 (electronic)
- [46] G. Peccati and C. Zheng (2011). Universal Gaussian fluctuations on the discrete Poisson chaos. Preprint.
- [47] M. Penrose (2003). Random Geometric Graphs. Oxford University Press.
- [48] W. Rudin (1987): Real and Complex Analysis. Third edition. McGraw-Hill Book Co., New York.
- [49] I. Shigekawa (1980). Derivatives of Wiener functionals and absolute continuity of induced measures. J. Math. Kyoto Univ. **20**(2), 263-289.
- [50] Ch. Stein (1972). A bound for the error in the normal approximation to the distribution of a sum of dependent random variables. In: *Proceedings of the Sixth Berkeley Symposium on Mathematical Statistics and Probability, Vol. II: Probability theory*, 583-602. Univ. California Press, Berkeley, Calif..
- [51] M. Talagrand (2003). Spin Glasses, a Challenge for Mathematicians. Springer Verlag.
- [52] D.V. Voiculescu (1985). Symmetries of some reduced free product C*-algebras. Operator algebras and their connection with topology and ergodic theory, Springer Lecture Notes in Mathematics 1132, 556–588.

- [53] D.V. Voiculescu (1991). Limit laws for random matrices and free product. *Invent. Math.* **104**, 201-220.
- [54] E.P. Wigner (1958). On the distribution of the roots of certain symmetric matrices. *Ann. of Math.* (2) **67**, 325-327.