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# AN ANDREOTTI-GRAUERT THEOREM WITH $L^{r}$ ESTIMATES. 

ERIC AMAR


#### Abstract

By a theorem of Andreotti and Grauert if $\omega$ is a $(p, q)$-current, $q<n$, in a Stein manifold, $\bar{\partial}$ closed and with compact support, then there is a solution $u$ to $\bar{\partial} u=\omega$ still with compact support. The aim of this work is to show that if moreover $\omega \in L^{r}(d m)$, where $m$ is a suitable Lebesgue measure on the Stein manifold, then we have a solution $u$ with compact support and in $L^{r}(d m)$.


## 1. Introduction.

Let $\omega$ be a $\bar{\partial}$ closed $(p, q)$ form in $\mathbb{C}^{n}$ with compact support $K:=\operatorname{Supp} \omega$ and such that $\omega \in L^{r}\left(\mathbb{C}^{n}\right)$. Setting $K$ in a ball $\mathbb{B}:=B(0, R)$ with $R$ big enough, we know, by a theorem of Ovrelid [9], that we have a $(p, q-1)$ form $u \in L^{r}(\mathbb{B})$ such that $\bar{\partial} u=\omega$. On the other hand we also know, at least when $q<n$, that there is a current $v$ with compact support such that $\bar{\partial} v=\omega$, by a theorem of Andreotti-Grauert [3].
So a natural question is : may we have a solution $u$ of $\bar{\partial} u=\omega$ with compact support and in $L^{r}\left(\mathbb{C}^{n}\right)$ ?

We already answered this question by the affirmative in a join work with S. Mongodi [2] explicitly by the "method of coronas".

The aim of this work is to extend this result to Stein manifolds and for it we use a completely different approach inspired by the Serre duality [10]. Because Hahn Banach theorem is used, these results are no longer so explicit and constructive as in [2]. On the other hand the control of the support is better : the support of the solution $u$ is contained in any $r^{\prime}$-regular domain containing the support of the data $\omega$. The definition of $r^{\prime}$-regular domain will be given later on, but strictly pseudo convex domains are such domains.

I am indebted to G. Tomassini who started my interest in this subject by e-mails on precisely this kind of questions and also to S . Mongodi for a lot of discussions by mails on the subject during the preparation of our join paper [2].

## 2. Duality.

We shall study a duality between currents inspired by the Serre duality [10].
Let $\Omega \subset \mathbb{C}^{n}$ be an open set in $\mathbb{C}^{n}$ and $t$ a $(p, q)$ current with compact support in $\Omega$, noted $t \in$ $\mathcal{D}_{c,(p, q)}^{\prime}(\Omega)$.
Let also $\phi \in \mathcal{C}_{(n-p, n-q)}^{\infty}(\Omega)$ a $(n-p, n-q)$ form in $\mathcal{C}^{\infty}(\Omega)$. We have that $t \wedge \phi$ is a $(n, n)$ current with compact support in $\Omega$.

As usual we use the following notation for the pairing

$$
\langle t, \phi\rangle:=t(\phi),
$$

where $t(\phi)$ is the action, as a current, of $t$ on the smooth form $\phi$ of complementary bi-degree.
Lemma 2.1. Let $\Omega$ be an open set in $\mathbb{C}^{n}, t \in \mathcal{D}_{c,(p, q)}^{\prime}(\Omega)$ a $(p, q)$ current with compact support in $\Omega$ and $u \in \mathcal{D}_{c,(p, q-1)}^{\prime}(\Omega)$. Then we have $\bar{\partial} t=u$ iff
$(*) \quad \forall \phi \in \mathcal{C}_{(n-p, n-q)}^{\infty}(\Omega),\langle t, \phi\rangle=(-1)^{p+q-1}\langle u, \bar{\partial} \phi\rangle$.
Proof.
If $\bar{\partial} u=t$, let $\chi \in \mathcal{C}_{c}^{\infty}(\Omega)$ such that $\chi \equiv 1$ on the support of $u$, hence on the support of $t$. Then

$$
\forall \phi \in \mathcal{C}_{(n-p, n-q)}^{\infty}(\Omega),\langle t, \phi\rangle=\langle t, \chi \phi\rangle=\langle\bar{\partial} u, \chi \phi\rangle=(-1)^{p+q-1}\langle u, \bar{\partial}(\chi \phi)\rangle
$$

by definition of the action of derivatives on currents. Hence
$\langle t, \phi\rangle=(-1)^{p+q-1}\langle u, \bar{\partial} \chi \wedge \phi\rangle+(-1)^{p+q-1}\langle u, \chi \bar{\partial} \phi\rangle=(-1)^{p+q-1}\langle u, \chi \bar{\partial} \phi\rangle$
because $\bar{\partial} \chi=0$ on the support of $u$; so
$\langle t, \phi\rangle=(-1)^{p+q-1}\langle u, \chi \bar{\partial} \phi\rangle=(-1)^{p+q-1}\langle u, \bar{\partial} \phi\rangle$
because $\chi=1$ on the support of $u$.
Conversely if we have $(*)$ we take $\chi \in \mathcal{C}_{c}^{\infty}(\Omega)$ such that $\chi \equiv 1$ on union of the support of $u$ and the support of $t$ then

$$
\forall \phi \in \mathcal{C}_{(n-p, n-q)}^{\infty}(\Omega),\langle t, \phi\rangle=\langle t, \chi \phi\rangle
$$

because $\chi=1$ on the support of $t$, and by $(*)$
$\langle t, \chi \phi\rangle=(-1)^{p+q-1}\langle u, \bar{\partial}(\chi \phi)\rangle=\langle\bar{\partial} u, \chi \phi\rangle=\langle\bar{\partial} u, \phi\rangle$
because $\chi=1$ on the support of $u$ hence on the support of $\bar{\partial} u$.
So

$$
\forall \phi \in \mathcal{C}_{(n-p, n-q)}^{\infty}(\Omega),\langle t, \phi\rangle=\langle\bar{\partial} u, \phi\rangle,
$$

which means that $\bar{\partial} u=t$ as a current.

## 3. Solution of the $\bar{\partial}$ equation with compact support.

### 3.1. Domain $r$ regular.

We shall use the definition.
Definition 3.1. Let $X$ be a complex manifold and $\Omega$ a domain in $X$; we shall say that $\Omega$ is $r$ -regular, for $r \in[1, \infty]$ if for $p, q \in\{0, \ldots, n\}$ there is a constant $C=C_{p, q}$ such that for any $(p, q)$ form $\omega, \bar{\partial}$ closed in $\Omega$ and in $L^{r}(\Omega)$ there is a $(p, q-1)$ form $u \in L^{r}(\Omega)$ such that $\bar{\partial} u=\omega$ and $\|u\|_{L^{r}(\Omega)} \leq C\|\omega\|_{L^{r}(\Omega)}$.
The $L^{r}$ norms for forms will be defined later on.
Examples of 2 -regular domains are the bounded pseudo-convex domains by Hörmander [6].
Examples of $r$-regular domains in $\mathbb{C}^{n}$ are the strictly pseudo-convex domains with smooth boundary by Ovrelid [9]. The polydiscs in $\mathbb{C}^{n}$ by Charpentier [4], finite intersections of strictly pseudo-convex bounded domains with transverse intersections by Menini [8].

Examples of $r$ regular domains in Stein manifold are the strictly pseudo-convex domains with smooth boundary by Kerzman [7] for $(0,1)$ forms and Demailly-Laurent [5] for the general case.

### 3.2. Weighted $L^{r}$ spaces.

Let $\mathbb{B}:=B(0, R)$, the ball of $\mathbb{C}^{n}$ centered at 0 and of radius $R$. We shall need the following notations.
We note $d m$ the Lebesgue measure on $\mathbb{C}^{n}$ and we shall define $L^{r}(\mathbb{B}, w)$ to be the set of functions $f$ defined on $\mathbb{B}$ such that
$\|f\|_{L^{r}(\mathbb{B}, w)}^{r}:=\int_{\mathbb{B}}|f(z)|^{r} w(z) d m(z)<\infty$,
with a weight $w(z)>0$. As usual we set $L^{r}(\mathbb{B})$ for $L^{r}(\mathbb{B}, 1)$.
Let $\mathcal{I}_{p}$ be the set of multi-indices of length $p$ in $(1, \ldots, n)$. We shall use the following measure defined on $\Gamma:=\mathbb{B} \times \mathcal{I}_{p} \times \mathcal{I}_{q}$ the following way :

$$
d \mu(z, k, l)=d \mu_{w, p, q}(z, k, l):=w(z) d m(z) \otimes \sum_{|I|=p,|J|=q} \delta_{I}(k) \otimes \delta_{J}(l),
$$

where $\delta_{I}(k)=1$ if the multi-index $k$ is equal to $I$ and $\delta_{I}(k)=0$ if not.
This means, if $f(z, I, J)$ is a function defined on $\Gamma$, that

$$
\int f(z, k, l) d \mu_{w, p, q}(z, k, l):=\sum_{|I|=p,|J|=q} \int_{\mathbb{B}} f(z, I, J) w(z) d m(z) .
$$

If $I$ is a multi-index of length $p$, let $I^{c}$ be the unique multi-index, ordered increasingly, such that $I \cup I^{c}=(1,2, \ldots, n)$; then $I^{c}$ is of length $n-p$.

To $t=\sum_{|I|=p,|J|=q} t_{I, J}(z) d z^{I} \wedge d \bar{z}^{J}$ a $(p, q)$ form, we associate the function on $\Gamma$ :
$T(z, I, J):=(-1)^{s(I, J)} t_{I, J}(z)$,
where
$s(I, J)=0$ if $d z^{I} \wedge d \bar{z}^{J} \wedge d z^{I^{c}} \wedge d \bar{z}^{J^{c}}=d z_{1} \wedge \cdots \wedge d z_{n} \wedge d \bar{z}_{1} \wedge \cdots \wedge d \bar{z}_{n}$ as a $(n, n)$ form and
$s(I, J)=1$ if not.
If $\phi=\sum_{|I|=p,|J|=q} \phi_{I^{c}, J c}(z) d z^{I} \wedge d \bar{z}^{J}$ is of complementary bi-degree, associate in the same manner :
$\Phi^{*}(z, I, J):=\phi_{I^{c}, J^{c}}(z)$. This is still a function on $\Gamma$.
Now we have, for $1<r<\infty$, if $T(z, I, J)$ is a function in $\mathbb{B}$ with $L^{r}(\mathbb{B})$ coefficients and with $\mu=\mu_{w, p, q}$,

$$
\|T\|_{L^{r}(d \mu)}^{r}:=\int|T(z, I, J)|^{r} d \mu_{w, p, q}(x, I, J)=\sum_{|I|=p,|J|=q}\|T(\cdot, I, J)\|_{L^{r}(\mathbb{B}, w)}^{r}
$$

The dual of $L^{r}(\mu)$ is $L^{r^{\prime}}(\mu)$ where $r^{\prime}$ is the conjugate of $r, \frac{1}{r}+\frac{1}{r^{\prime}}=1$, and the norm is defined analogously with $r^{\prime}$ replacing $r$.

We also know that, for $p, q$ fixed,

$$
\begin{equation*}
\|T\|_{L^{r}(\mu)}=\sup _{\Phi \in L^{r^{\prime}}(\mu)} \frac{\left|\int T \Phi d \mu\right|}{\|\Phi\|_{L^{r^{\prime}}(\mu)}} \tag{3.1}
\end{equation*}
$$

For a $(p, q)$ form $t=\sum_{|J|=p,|K|=q} t_{J, K} d z^{J} \wedge d \bar{z}^{K}$, and a weight $w>0$ we define its norm by :

$$
\begin{equation*}
\|t\|_{L^{r}(\mathbb{B}, w)}^{r}:=\sum_{|J|=p,|K|=q}\left\|t_{J, K}\right\|_{L^{r}(\mathbb{B}, w)}^{r}=\|T\|_{L^{r}(\mu)}^{r} \tag{3.2}
\end{equation*}
$$

Now we can state
Lemma 3.2. Let $w>0$ be a weight. If $u$ is a $(p, q)$ current defined on $(n-p, n-q)$ forms in $L^{r^{\prime}}(\mathbb{B}, w)$ and such that

$$
\forall \alpha \in L_{(n-p, n-q)}^{r^{\prime}}(\mathbb{B}, w),|\langle u, \alpha\rangle| \leq C\|\alpha\|_{L^{r^{\prime}}(\mathbb{B}, w)}
$$

then $\|u\|_{L^{r}\left(\mathbb{B}, w^{1-r}\right)} \leq C$.
Proof.
Let us take the measure $\mu=\mu_{w, p, q}$ as above. Let $\Phi^{*}$ be the function on $\Gamma$ associated to $\alpha$ and $T$ the one associated to $u$. We have, by definition of the measure $\mu$ applied to the function

$$
\begin{aligned}
& f(z, I, J):=T(z, I, J) w^{-1} \Phi^{*}(z, I, J) \\
& \int T w^{-1} \Phi^{*} d \mu=\int f(z, k, l) d \mu(z, k, l):=\sum_{|I|=p,|J|=q} \int_{\mathbb{B}} f(z, I, J) w(z) d m(z)=
\end{aligned}
$$

$$
=\sum_{|I|=p,|J|=q} \int_{\mathbb{B}} T(z, I, J) w^{-1}(z) \Phi^{*}(z, I, J) w(z) d m(z)=\langle u, \alpha\rangle,
$$

by definition of $T$ and $\Phi^{*}$.
Hence we have, by (3.1)

$$
\left\|T w^{-1}\right\|_{L^{r}(\mu)}=\sup _{\Psi \in L^{r^{\prime}}(\mu)} \frac{|\langle u, \alpha\rangle|}{\|\Psi\|_{L^{r^{\prime}}(\mu)}}
$$

But $\left\|T w^{-1}\right\|_{L^{r}(\mu)}=\left\|u w^{-1}\right\|_{L^{r}(\mathbb{B}, w)}$ by (3.2), and

$$
\left\|f w^{-1}\right\|_{L^{r}(\mathbb{B}, w)}^{r}=\int_{\mathbb{B}}\left|f \eta^{-1}\right|^{r} w d m=\int_{\mathbb{B}}|f|^{r} w^{1-r} d m=\|f\|_{L^{r}\left(\mathbb{B}, w^{1-r}\right)},
$$

so we get

$$
\|u\|_{L^{r}\left(\mathbb{B}, w^{1-r}\right)}=\sup _{\Psi \in L^{r^{\prime}}(\mu)} \frac{|\langle u, \alpha\rangle|}{\|\Psi\|_{L^{r^{\prime}}(\mu)}}
$$

which implies the lemma because, still by (3.1), we can take $\Psi=\Phi^{*}$ and $\|\Psi\|_{L^{r^{\prime}}(\mu)}=\|\alpha\|_{L^{r^{\prime}}(\mathbb{B}, w)}$.
Let us denote $\mathcal{H}_{p}(\Omega)$ the set of all $(p, 0) \bar{\partial}$ closed forms in $\Omega$. If $p=0, \mathcal{H}_{0}(\Omega)=\mathcal{H}(\Omega)$ is the set of holomorphic functions in $\Omega$. If $p>0$, we have $\phi \in \mathcal{H}_{p}(\Omega) \Longrightarrow \phi(z)=\sum_{|J|=p} a_{J}(z) d z^{J}$, where $d z^{J}:=d z_{j_{1}} \wedge \cdots \wedge d z_{j_{p}}$ and the functions $a_{J}(z)$ are holomorphic in $\Omega$. Hence in $\mathbb{C}^{n}$, a $(p, 0) \bar{\partial}$ closed form is a vector of global holomorphic functions in $\Omega$.
Theorem 3.3. Let $t$ be a $(p, q)$ form in $L_{c}^{r}\left(\mathbb{C}^{n}\right), r>1, K:=\operatorname{Supp} t$. Let $\mathbb{B}$ be a ball containing $K$ and let $\Omega \subset \mathbb{B}$ be a $r^{\prime}$-regular domain containing $K$. Suppose that $t$ is such that $\bar{\partial} t=0$ if $1 \leq q<n$ and $\forall h \in \mathcal{H}_{p}(\Omega),\langle t, h\rangle=0$ if $q=n$.
Then there is a $(p, q-1)$ form $u$ in $L_{c}^{r}(\mathbb{B}), 0$ a.e. on $\mathbb{B} \backslash \Omega$ and such that $\bar{\partial} u=t$ as distributions.

## Proof.

Let us consider the weight $\eta=\eta_{\epsilon}:=\mathbb{1}_{\Omega}(z)+\epsilon_{\mathbb{B} \backslash \Omega}(z)$ for a fixed $\epsilon>0$.
We shall consider the linear form $\mathcal{L}$ defined on $(n-p, n-q+1)$ form $\alpha \in L^{r^{\prime}}(\mathbb{B}, \eta), \bar{\partial}$ closed in $\Omega$ as follows :
$\mathcal{L}(\alpha):=(-1)^{p+q-1}\langle t, \phi\rangle$, where $\phi \in L^{r^{\prime}}(\mathbb{B})$ is such that $\bar{\partial} \phi=\alpha$.
Because $\epsilon>0$ we have $\alpha \in L^{r^{\prime}}(\mathbb{B}, \eta) \Longrightarrow \alpha \in L^{r^{\prime}}(\mathbb{B})$ and such a $\phi$ exists because $\mathbb{B}$ being $r^{\prime}$-regular there is $\phi \in L^{r^{\prime}}(\mathbb{B})$ with $\bar{\partial} \phi=\alpha$.

Let us see that $\mathcal{L}$ is well defined :
suppose first that $q<n$. In order for $\mathcal{L}$ to be well defined we need
$\forall \phi, \psi \in L_{(n-p, n-q)}^{r^{\prime}}(\mathbb{B}), \bar{\partial} \phi=\bar{\partial} \psi \Longrightarrow\langle t, \phi\rangle=\langle t, \psi\rangle$.
This is meaningful because $t \in L_{c}^{r}\left(\mathbb{C}^{n}\right), r>1$, Supp $t \subset \mathbb{B}$.
Then we have $\bar{\partial}(\phi-\psi)=0$ hence, because $\mathbb{B}$ is $r^{\prime}$-regular, we can solve $\bar{\partial}$ in $L^{r^{\prime}}(\mathbb{B})$ :
$\exists \gamma \in L_{(n-p, n-q-1)}^{r^{\prime}}(\mathbb{B}):: \bar{\partial} \gamma=(\phi-\psi)$.
So $\langle t, \phi-\psi\rangle=\langle t, \bar{\partial} \gamma\rangle=(-1)^{p+q-1}\langle\bar{\partial} t, \gamma\rangle=0$. Hence $\mathcal{L}$ is well defined in that case.
Suppose now that $q=n$, then of course $\bar{\partial} t=0$ and we have that $\phi, \psi$ are ( $p, 0$ ) forms hence $\bar{\partial}(\phi-\psi)=0$ means that $h:=\phi-\psi$ is a $\bar{\partial}$ closed $(p, 0)$ form hence $h \in \mathcal{H}_{p}(\mathbb{B}) \subset \mathcal{H}_{p}(\Omega)$. But by assumption we have $\langle t, h\rangle=0$ hence $\mathcal{L}$ is also well defined in that case.

If $\alpha$ is a $(n-p, n-q+1)$ form in $L^{r^{\prime}}(\mathbb{B}, \eta), \bar{\partial}$ closed in $\mathbb{B}$, then, $\alpha \in L^{r^{\prime}}(\Omega)$ and is still $\bar{\partial}$ closed in $\Omega$, hence there is a $\psi \in L^{r^{\prime}}(\Omega):: \bar{\partial} \psi=\alpha$ because $\Omega$ is $r^{\prime}$-regular.
For $q<n$, we have $\bar{\partial}(\phi-\psi)=\alpha-\alpha=0$ on $\Omega$ and, because Supp $t \subset \Omega \subset \mathbb{B}, \bar{\partial} t=0$ we get $\langle t, \phi-\psi\rangle=0$ hence

$$
\mathcal{L}(\alpha)=\langle t, \phi\rangle=\langle t, \psi\rangle
$$

If $q=n$, we still have $\bar{\partial}(\phi-\psi)=\alpha-\alpha=0$ on $\Omega$, and $\phi-\psi \in \mathcal{H}_{p}(\Omega)$ hence again by the hypothesis we still get

$$
\mathcal{L}(\alpha)=\langle t, \phi\rangle=\langle t, \psi\rangle .
$$

In any cases, by Hölder inequalities

$$
|\mathcal{L}(\alpha)| \leq\|t\|_{L^{r}(\Omega)}\|\psi\|_{L^{r^{r}}(\Omega)}
$$

But by the $r^{\prime}$-regularity of $\Omega$ there is a constant $C$ such that

$$
\|\psi\|_{L^{r^{\prime}}(\Omega)} \leq C\|\alpha\|_{L^{r^{\prime}}(\Omega)}
$$

Of course we have

$$
\|\alpha\|_{L^{r^{\prime}}(\Omega)} \leq\|\alpha\|_{L^{r^{\prime}}(\mathbb{B}, \eta)}
$$

because $\eta=1$ on $\Omega$, hence

$$
|\mathcal{L}(\alpha)| \leq C\|t\|_{L^{r}(\Omega)}\|\alpha\|_{L^{r^{\prime}}(\mathbb{B}, \eta)}
$$

So we have that the norm of $\mathcal{L}$ is bounded on the subspace of $\bar{\partial}$ closed forms in $L^{r^{\prime}}(\mathbb{B}, \eta)$ by $C\|t\|_{L^{r}(\Omega)}$ which is independent of $\epsilon$.

We apply the Hahn-Banach theorem to extend $\mathcal{L}$ with the same norm to all $(n-p, n-q+1)$ forms in $L^{r^{\prime}}(\mathbb{B}, \eta)$. As in Serre duality theorem ( $[10]$, p. 20) this is one of the main ingredient in the proof.

This means, by the definition of current, that there is a ( $p, q-1$ ) current $u$ which represents the extended form $\mathcal{L}$ and such that

$$
\sup _{\alpha \in L^{r^{\prime}}(\mathbb{B}, \eta),\|\alpha\|=1}|\langle u, \alpha\rangle| \leq C\|t\|_{L^{r}(\Omega)}
$$

and by lemma 3.2 with the weight $\eta$, this implies

$$
\|u\|_{L^{r}\left(\mathbb{B}, \eta^{1-r}\right)} \leq C\|t\|_{L^{r}(\Omega)}
$$

In particular $\|u\|_{L^{r}(\mathbb{B})} \leq C\|t\|_{L^{r}(\Omega)}$ because with $\epsilon<1, r>1$, we have $\eta^{1-r} \geq 1$.
So applied to a $\bar{\partial}$ closed $(n-p, n-q+1)$ current $\alpha$ we get

$$
\langle u, \alpha\rangle=(-1)^{p+q-1}\langle t, \phi\rangle \text {, with } \bar{\partial} \phi=\alpha \text {, i.e. } \forall \phi:: \bar{\partial} \phi \in L^{r^{\prime}}(\mathbb{B}, \eta) \text {, we have }\langle u, \bar{\partial} \phi\rangle=
$$ $(-1)^{p+q-1}\langle t, \phi\rangle$ and this means precisely, by lemma 2.1, that $\bar{\partial} u=t$.

Now for $\epsilon>0$ with $\eta_{\epsilon}(z):=\mathbb{1}_{\Omega}(z)+\epsilon \mathbb{\|}_{\mathbb{B} \backslash \Omega}(z)$, let $u_{\epsilon} \in L^{r}\left(\mathbb{B}, \eta_{\epsilon}^{1-r}\right)$ be the previous solution, then

$$
\left\|u_{\epsilon}\right\|_{L^{r}(\mathbb{B})}^{r} \leq \int_{\mathbb{B}}\left|u_{\epsilon}\right|^{r} \eta^{1-r} d m \leq C^{r}\|t\|_{L^{r}(\Omega)}^{r}
$$

Replacing $\eta$ by its value we get

$$
\int_{\Omega}\left|u_{\epsilon}\right|^{r} d m+\int_{\mathbb{B} \backslash \Omega}\left|u_{\epsilon}\right|^{r} \epsilon^{1-r} d m \leq C^{r}\|t\|_{L^{r}(\Omega)}^{r} \Longrightarrow \int_{\mathbb{B} \backslash \Omega}\left|u_{\epsilon}\right|^{r} \epsilon^{1-r} d m \leq C^{r}\|t\|_{L^{r}(\Omega)}^{r}
$$

hence

$$
\int_{\mathbb{B} \backslash \Omega}\left|u_{\epsilon}\right|^{r} d m \leq C^{r} \epsilon^{r-1}\|t\|_{L^{r}(\Omega)}^{r}
$$

Because $C$ and the norm of $t$ are independent of $\epsilon$, we have $\left\|u_{\epsilon}\right\|_{L^{r}(\mathbb{B})}$ is uniformly bounded and $r>1 \Longrightarrow L^{r}(\mathbb{B})$ is a dual, hence there is a sub-sequence $\left\{u_{\epsilon_{k}}\right\}_{k \in \mathbb{N}}$ of $\left\{u_{\epsilon}\right\}$ which converges weakly, when $\epsilon_{k} \longrightarrow 0$, to a $(p, q-1)$ form $u$ in $L^{r}(\mathbb{B})$.
To see that this form $u$ is 0 a.e. on $\mathbb{B} \backslash \Omega$ let us take a component $u_{I, J}$ of it ; it is the weak limit of the sequence of functions $\left\{u_{\epsilon_{k}, I, J}\right\}$ which means, with the notations $v:=u_{I, J}, v_{k}:=u_{\epsilon_{k}, I, J}$

$$
\forall f \in L^{r^{\prime}}(\mathbb{B}), \int_{\mathbb{B}_{\overline{1}}} v f d m=\lim _{k \rightarrow \infty} \int_{\mathbb{B}} v_{k} f d m .
$$

As usual take $f:=\frac{\bar{v}}{|v|} \mathbb{1}_{E}$ where $E:=\{|v|>0\} \cap(\mathbb{B} \backslash \Omega)$ then we get

$$
\int_{\mathbb{B}} v f d m=\int_{E}|v| d m=\lim _{k \longrightarrow \infty} \int_{\mathbb{B}} v_{k} f d m=\lim _{k \longrightarrow \infty} \int_{E} \frac{v_{k} \bar{v}}{|v|} d m .
$$

Now we have by Hölder

$$
\left|\int_{E} \frac{v_{k} \bar{v}}{|v|} d m\right| \leq\left\|v_{k}\right\|_{L^{r}(E)}\left\|\mathbb{1}_{E}\right\|_{L^{r^{\prime}}(E)}
$$

But

$$
\left\|v_{k}\right\|_{L^{r}(E)}^{r} \leq \int_{\mathbb{B} \backslash \Omega}\left|u_{\epsilon_{k}}\right|^{r} d m \leq\left(\epsilon_{k}\right)^{r-1}\|t\|_{L^{r}(\Omega)} \longrightarrow 0, k \longrightarrow \infty .
$$

Hence

$$
\left|\int_{E}\right| v|d m| \leq C^{r}\left\|1_{E}\right\|_{L^{r^{\prime}}(E)}\left(\epsilon_{k}\right)^{r-1}\|t\|_{L^{r}(\Omega)}^{r} \longrightarrow 0
$$

which implies $m(E)=0$ because on $E,|v|>0$.
This being true for all components of $u$, we get that the form $u$ is 0 a.e. on $\mathbb{B} \backslash \Omega$.
So we get

$$
\forall \phi \in \mathcal{D}_{n-p, n-q}(\mathbb{B}),(-1)^{p+q-1}\langle t, \phi\rangle=\left\langle u_{\epsilon}, \bar{\partial} \phi\right\rangle \longrightarrow\langle u, \bar{\partial} \phi\rangle \Longrightarrow\langle u, \bar{\partial} \phi\rangle=(-1)^{p+q-1}\langle t, \phi\rangle
$$

hence $\bar{\partial} u=t$ as distributions.
Remark 3.4. The condition of orthogonality to the holomorphic functions in case $q=n$ was already seen for extension of $C R$ functions see [1] and the references therein.

## 4. Case of Stein manifold.

First we define the "Lebesgue measure" on a complex manifold $X$ as in Hörmander's book [6] section 5.2 , with a hermitian metric locally equivalent to the usual one on any analytic coordinates system. Associate to this metric there is a volume form $d V$ and we take it for the Lebesgue measure on $X$.

Let us denote $\mathcal{H}_{p}(\Omega)$ the set of all $(p, 0) \bar{\partial}$ closed forms in $\Omega$. Still we have that $\mathcal{H}_{0}(\Omega)=\mathcal{H}(\Omega)$, the set of holomorphic functions in $\Omega$, but if $p>1, \phi \in \mathcal{H}_{p}(\Omega)$ is a priori no longer a vector of global holomorphic functions in $\Omega$.
Theorem 4.1. Let $X$ be a Stein manifold and $\Omega$ a bounded strictly pseudo-convex domain with smooth boundary of $X$. Let $t$ be a $(p, q)$ current in $L_{c}^{r}(\Omega), r>1$, such that $\bar{\partial} t=0$ if $1 \leq q<n$ and $\forall h \in \mathcal{H}_{p}(\Omega),\langle t, h\rangle=0$ if $q=n$.
Let $B$ be another bounded strictly pseudo-convex domain with smooth boundary in $X$ such that $\bar{\Omega} \subset B$.

Then there is a $(p, q-1)$ current $u$ in $L^{r}(B)$ such that $\bar{\partial} u=t$ and which is 0 a.e. in $B \backslash \Omega$.
Proof.
We can solve the $\bar{\partial}$ equation in $B$ and $\Omega$ with $L^{r}$ estimates for $(0,1)$ forms by use of N . Kerzman kernels [7] and for all ( $p, q$ ) forms by J-P. Demailly and C. Laurent ones( [5], Remarque 4, page 596).

The remainder of the proof is identical to the proof of theorem 3.3.

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