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## To cite this version:

Marie-Françoise Bidaut-Véron, Nguyen Anh Dao. $L^{\infty}$ estimates and uniqueness results for nonlinear parabolic equations with gradient absorption terms. 2012. hal-00669365v2

HAL Id: hal-00669365<br>https://hal.science/hal-00669365v2

Preprint submitted on 22 Mar 2013

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# $L^{\infty}$ estimates and uniqueness results for nonlinear parabolic equations with gradient absorption terms 

Marie Françoise BIDAUT-VERON* Nguyen Anh DAO ${ }^{\dagger}$


#### Abstract

We study the nonnegative solutions of the viscous Hamilton-Jacobi problem $$
\left\{\begin{array}{c} u_{t}-\nu \Delta u+|\nabla u|^{q}=0, \\ u(0)=u_{0}, \end{array}\right.
$$ in $Q_{\Omega, T}=\Omega \times(0, T)$, where $q>1, \nu \geqq 0, T \in(0, \infty]$, and $\Omega=\mathbb{R}^{N}$ or $\Omega$ is a smooth bounded domain, and $u_{0} \in L^{r}(\Omega), r \geqq 1$, or $u_{0} \in \mathcal{M}_{b}(\Omega)$. We show $L^{\infty}$ decay estimates, valid for any weak solution, without any conditions as $|x| \rightarrow \infty$, and without uniqueness assumptions. As a consequence we obtain new uniqueness results, when $u_{0} \in \mathcal{M}_{b}(\Omega)$ and $q<(N+2) /(N+1)$, or $u_{0} \in L^{r}(\Omega)$ and $q<(N+2 r) /(N+r)$. We also extend some decay properties to quasilinear equations of the model type $$
u_{t}-\Delta_{p} u+|u|^{\lambda-1} u|\nabla u|^{q}=0
$$


where $p>1, \lambda \geqq 0$, and $u$ is a signed solution.

Keywords Viscous Hamilton-Jacobi equation; quasilinear parabolic equations with gradient terms; regularity; decay estimates; regularizing effects; uniqueness results.
A.M.S. Subject Classification 35K15, 35K55, 35B33, 35B65, 35D30

## 1 Introduction

In this article we study a class of heat equations involving a nonlinear gradient absorption term. We are mainly concerned by the nonnegative solutions of the viscous parabolic Hamilton-Jacobi equation

$$
\begin{equation*}
u_{t}-\nu \Delta u+|\nabla u|^{q}=0 \tag{1.1}
\end{equation*}
$$

in $Q_{\Omega, T}=\Omega \times(0, T), T \leqq \infty$, where $q>1, \nu \geqq 0$, and $\Omega=\mathbb{R}^{N}$, or $\Omega$ is a smooth bounded domain of $\mathbb{R}^{N}$.

[^0]We study the Cauchy problem in $\mathbb{R}^{N}$ and the Cauchy-Dirichlet problem when $\Omega$ is bounded, with initial data $u(., 0)=u_{0} \geqq 0$, where $u_{0} \in L^{r}(\Omega), r \geqq 1$, or $u_{0}$ is a bounded Radon measure on $\Omega$.

We also consider the (signed) solutions of quasilinear equations of the type

$$
\begin{equation*}
u_{t}-\nu \Delta_{p} u+|u|^{\lambda-1} u|\nabla u|^{q}=0 \tag{1.2}
\end{equation*}
$$

where $p>1$ and $\Delta_{p}$ is the $p$-Laplacian, and more generally

$$
\begin{equation*}
u_{t}-\operatorname{div}(\mathrm{A}(x, t, u, \nabla u))+g(x, u, \nabla u)=0 \tag{1.3}
\end{equation*}
$$

with natural growth conditions on the function A, and nonnegativity conditions

$$
\begin{equation*}
\mathrm{A}(x, t, u, \eta) \cdot \eta \geqq \nu|\eta|^{p}, \quad g(x, u, \eta) u \geqq \gamma|u|^{\lambda+1}|\nabla u|^{q} \quad \gamma \geqq 0, \nu \geqq 0, \lambda \geqq 0, \tag{1.4}
\end{equation*}
$$

without monotonicity assumption.
In the sequel we give some decay estimates, under very few assumptions on the solutions. Then from Moser's technique, we deduce regularizing effects : $L^{\infty}$ estimates, in terms of $u_{0}$, and universal estimates when $\Omega$ is bounded. We show that two types of regularizing effect can occur: the first one is due to the gradient term $|\nabla u|^{q}$ (when $\gamma>0$ ), the second one is due to the operator itself (when $\nu>0$ ).

A part of these estimates are well known for equation (1.1) when the solutions can be approximated by smooth solutions, or satisfy growth conditions as $|x| \rightarrow \infty$ when $\Omega=\mathbb{R}^{N}$, for example semi-group solutions. Our approach is different, and our results are valid for all the solutions of the equation in a weak sense: in the sense of distributions for equation (1.1), in the renormalized sense for equation (1.3). And we make no assumption of uniqueness. In the case of equation (1.1) in $\mathbb{R}^{N}$, we require no condition as $|x| \rightarrow \infty$, all our assumptions are local.

As a consequence we deduce new uniqueness results for equation (1.1) in $\mathbb{R}^{N}$ or in a bounded domain $\Omega$.

## 2 Main results

We denote by $\mathcal{M}_{b}(\Omega)$ the set of bounded Radon measures in $\Omega$, and $\mathcal{M}_{b}^{+}(\Omega)$ the cone of nonnegative ones.
We set $Q_{\Omega, s, \tau}=\Omega \times(s, \tau)$, for any $0 \leqq s<\tau \leqq \infty$, thus $Q_{\Omega, T}=Q_{\Omega, 0, T}$.
As usual, for any $\theta \geqq 1$ we note by $\theta^{\prime}=\theta /(\theta-1)$ the conjugate of $\theta$.
In Section 3, we give some key tools for obtaining regularizing properties. The main one is an iteration property based of Moser's method, inspired by [38]:

Lemma 2.1 Let $m>1, \theta>1$ and $\lambda \in \mathbb{R}$ and $C_{0}>0$. Let $v \in C\left([0, T) ; L_{\text {loc }}^{1}(\Omega)\right)$ be nonnegative, and $v_{0}=v(x, 0) \in L^{r}(\Omega)$ for some $r \geqq 1$ such that

$$
\begin{equation*}
r>\theta^{\prime}(1-m-\lambda) . \tag{2.1}
\end{equation*}
$$

If $r>1$ we assume that for any $0 \leqq s<t<T$ and any $\alpha \geqq r-1$, there holds

$$
\begin{equation*}
\frac{1}{\alpha+1} \int_{\Omega} v^{\alpha+1}(., t) d x+\frac{C_{0}}{\beta^{m}} \int_{s}^{t}\left(\int_{\Omega} v^{\beta m \theta}(., \tau) d x\right)^{\frac{1}{\theta}} d \tau \leqq \frac{1}{\alpha+1} \int_{\Omega} v^{\alpha+1}(., s) d x \tag{2.2}
\end{equation*}
$$

where $\beta=\beta(\alpha)=1+(\alpha+\lambda) / m$, and the right-hand side can be infinite.
If $r=1$ we make one of the two following assumptions:
( $H_{1}$ ) (2.2) holds for any $\alpha \geqq 0$,
( $H_{2}$ ) $\int_{\Omega} v(., t) d x \leqq \int_{\Omega} v_{0} d x$ for any $t \in(0, T)$, and $v_{0} \in L^{\rho}(\Omega)$ for some $\rho>1$ such that $\rho \theta^{\prime}(1-m-\lambda)<1$ and (2.2) holds for any $\alpha \geqq \rho-1$.
Then there exists $C>0$, depending on $N, m, r, \lambda, C_{0}$, and possibly $\rho$, such that for any $t \in(0, T)$,

$$
\begin{equation*}
\|v(., t)\|_{L^{\infty}(\Omega)} \leqq C t^{-\sigma_{r, m, \lambda, \theta}}\left\|v_{0}\right\|_{L^{r}(\Omega)}^{\varpi_{r, m, \lambda, \theta}} \tag{2.3}
\end{equation*}
$$

where

$$
\begin{equation*}
\sigma_{r, m, \lambda, \theta}=\frac{1}{\frac{r}{\theta^{\prime}}+\lambda+m-1}=\frac{\theta^{\prime}}{r} \varpi_{r, m, \lambda, \theta} . \tag{2.4}
\end{equation*}
$$

This Lemma allows to obtain $L^{\infty}$ estimates for the solutions of equation (1.1), when $q \leqq N$, or $2 \leqq N$, and for equation (1.2) when $p \leqq N$. In the other cases the $L^{\infty}$ estimates follow from the Gagliardo-Nirenberg inequality, see Lemma 3.4. Moreover we deduce universal $L^{\infty}$ estimates when $\Omega$ is bounded, see Lemma 3.3.

In Section 4 we study the Cauchy Hamilton-Jacobi problem in $\mathbb{R}^{N}$ :

$$
\left\{\begin{array}{l}
u_{t}-\nu \Delta u+|\nabla u|^{q}=0, \quad \text { in } Q_{\mathbb{R}^{N}, T},  \tag{2.5}\\
u(x, 0)=u_{0} \geqq 0 \quad \text { in } \mathbb{R}^{N},
\end{array}\right.
$$

This equation is the objet of a huge literature, see [2], [12], [7], [15], [36], and the references therein, and also [7], [14], [28].

The first studies concern smooth initial data $u_{0} \in C_{b}^{2}\left(\mathbb{R}^{N}\right)$. From [2], (2.5) has a unique global solution $u \in C^{2,1}\left(\mathbb{R}^{N} \times[0, \infty)\right)$, and $u$ satisfies decay properties:

$$
\begin{aligned}
&\|u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq\left\|u_{0}\right\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \\
&\|\nabla u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq\left\|\nabla u_{0}\right\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} .
\end{aligned}
$$

Estimates of the gradient have been obtained for this solution, by using the Bersnstein technique, which consists in computing the equation satisfied by $|\nabla u|^{2}$ : first from [31],

$$
\|\nabla u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}^{q} \leqq t^{-1}\left\|u_{0}\right\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}
$$

then from [12], when $\nu>0$,

$$
\begin{equation*}
\left\|\nabla\left(u^{\frac{1}{q^{\prime}}}\right)(., t)\right\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq C(q, \nu) t^{-\frac{1}{2}}\left\|u_{0}\right\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}^{\frac{1}{q^{\prime}}}, \tag{2.6}
\end{equation*}
$$

$$
\begin{equation*}
\left\|\nabla\left(u^{\frac{1}{q^{T}}}\right)(., t)\right\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq \frac{(q-1)^{\frac{1}{q^{\prime}}}}{q} t^{-\frac{1}{q}}, \quad \text { that is } \quad|\nabla u(., t)|^{q} \leqq \frac{t^{-1} u(., t)}{q-1}, \quad \text { a.e.in } \mathbb{R}^{N} . \tag{2.7}
\end{equation*}
$$

If one only assumes $u_{0} \in C_{b}\left(\mathbb{R}^{N}\right)$, then (2.5) still has a unique solution $u$ such that $u \in C^{2,1}\left(Q_{\mathbb{R}^{N}, \infty}\right)$ and $u \in C\left(\mathbb{R}^{N} \times[0, \infty) \cap L^{\infty}\left(\mathbb{R}^{N} \times(0, \infty)\right)\right.$ see [29], and estimates (2.6) and (2.7) are still valid, from [7].

In case of rough initial data $u_{0} \in \mathcal{M}_{b}^{+}\left(\mathbb{R}^{N}\right)$ or $u \in L^{r}\left(\mathbb{R}^{N}\right), r \geqq 1$, assuming $\nu>0$, the solutions have been searched in an integral form

$$
\begin{equation*}
u(., t)=e^{t \Delta} u_{0}(.)-\nu \int_{0}^{t} e^{(t-s) \Delta}|\nabla u(., s)|^{q} d s \tag{2.8}
\end{equation*}
$$

involving the semi-group of heat equation $e^{t \Delta}$. Existence results hold in corresponding classes of solutions, involving integral conditions on the gradient in space and time, of global type:

- If $u_{0} \in \mathcal{M}_{b}^{+}\left(\mathbb{R}^{N}\right)$ and $1<q<(N+2) /(N+1)$, the existence of a solution $u \in C^{2,1}\left(Q_{\mathbb{R}^{N}, \infty}\right)$ is proved in [12] by approximation, and independently in [15], from the Banach fixed point theorem.
- If $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right), r \geqq 1$, existence holds for any $q \leqq 2$ from [15]. When $q>2$, it is required that $u_{0}$ is a limit of a monotone sequence of continuous functions, and existence is not known in the general case.

In those classes, decay properties and a regularizing effect follow directly from the semigroup $e^{t \Delta}$, since $u(., t) \leqq e^{t \Delta} u_{0}$. Our first main results shows that decay properties and $L^{\infty}$ estimates are valid for any weak solution, for any $\nu \geqq 0$, without any condition as $|x| \rightarrow \infty$ :

Theorem 2.2 Let $u \in L_{l o c}^{1}\left(Q_{\mathbb{R}^{N}, T}\right)$, with $|\nabla u| \in L_{l o c}^{q}\left(Q_{\mathbb{R}^{N}, T}\right)$, be any nonnegative solution of equation (1.1) in $\mathcal{D}^{\prime}\left(Q_{\mathbb{R}^{N}, T}\right)$.
(i) Let $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right), r \geqq 1$. Assume that $u \in C\left([0, T) ; L_{\text {loc }}^{r}\left(\mathbb{R}^{N}\right)\right)$ and $u(., 0)=u_{0}$. Then $u \in C\left([0, T) ; L^{r}\left(\mathbb{R}^{N}\right)\right)$; and for any $t \in(0, T), u(., t) \in L^{\infty}\left(\mathbb{R}^{N}\right)$ and

$$
\begin{align*}
& \|u(., t)\|_{L^{r}\left(\mathbb{R}^{N}\right)} \leqq\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)},  \tag{2.9}\\
& \|u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq\left\{\begin{array}{cc}
C t^{-\sigma_{r, q, N}}\left\|u_{0}\right\|_{L_{r}\left(, \mathbb{R}^{N}\right)}^{\omega_{r}, m_{r, q, N}}, & C=C(N, q, r),
\end{array} \quad \text { if } q \neq N,\right. \tag{2.10}
\end{align*}
$$

where

$$
\begin{equation*}
\sigma_{r, q, N}=\frac{1}{\frac{r q}{N}+q-1}=\frac{N}{r q} \varpi_{r, q, N} . \tag{2.11}
\end{equation*}
$$

And if $\nu>0$, then

$$
\|u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq\left\{\begin{array}{ccc}
C t^{-\frac{N}{2 r}}\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}, & C=C(N, r, \nu), & \text { if } N \neq 2,  \tag{2.12}\\
C_{\varepsilon} t^{-\frac{1+\varepsilon}{r}}\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}, & \forall \varepsilon>0, \quad C_{\varepsilon}=C(N, r, \nu, \varepsilon), & \text { if } N=2 .
\end{array}\right.
$$

(ii) Let $u_{0} \in \mathcal{M}_{b}^{+}\left(\mathbb{R}^{N}\right)$. Assume that $u(., t)$ converges weakly ${ }^{*}$ to $u_{0}$ as $t \rightarrow 0$. Then $u \in$ $C\left((0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right)$, and for any $t \in(0, T)$, the conclusions above with $r=1$ are still valid with $\left\|u_{0}\right\|_{L^{1}\left(\mathbb{R}^{N}\right)}$ replaced by $\int_{\mathbb{R}^{N}} d u_{0}$.

Note that estimates (2.9) are not valid for any weak subsolution of the heat equation. Here we prove that the result of (2.9) is essentially due to the gradient term $|\nabla u|^{q}$, which has a main regularizing effect on the equation. And then a second regularizing effect holds, due to the Laplacian, when $\nu>0$.

For any $q \leq 2$, we deduce estimates of the gradient, obtained from (2.6). As a consequence we deduce new uniqueness results, where the assumptions are only of local type:

Theorem 2.3 (i) Let $1<q<(N+2) /(N+1)$, and $u_{0} \in \mathcal{M}_{b}^{+}\left(\mathbb{R}^{N}\right)$. Then there exists a unique nonnegative function $u \in L_{l o c}^{1}\left(Q_{\mathbb{R}^{N}, T}\right)$, such that $|\nabla u| \in L_{l o c}^{q}\left(Q_{\mathbb{R}^{N}, T}\right)$, solution of equation (1.1) in $\mathcal{D}^{\prime}\left(Q_{\mathbb{R}^{N}, T}\right)$ such that

$$
\lim _{t \rightarrow 0} \int_{\mathbb{R}^{N}} u(., t) \psi d x=\int_{\mathbb{R}^{N}} \psi d u_{0}, \quad \forall \psi \in C_{c}\left(\mathbb{R}^{N}\right) .
$$

(ii) Let $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right), r \geqq 1$ and $1<q<(N+2 r) /(N+r)$. Then there exists a unique nonnegative solution $u$ as above, such that $u \in C\left([0, T) ; L_{\text {loc }}^{r}\left(\mathbb{R}^{N}\right)\right)$ and $u(., 0)=u_{0}$.

This improves the former uniqueness results of [12] and [15, Theorem 4.1], given in classes of semigroup solutions, satisfying conditions up to $t=0$ for the gradient: $|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right)$ in case (i), and $|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{r}\left(\mathbb{R}^{N}\right)\right)$ in case (ii).

We also find again in a shorter way the existence result of [15, Theorem 4.1], see Proposition 4.26. Finally we improve the estimate (2.9) when $q<(N+2 r) /(N+r)$, see Theorem 4.28.

In Section 5 we study the Cauchy-Dirichlet problem in a bounded domain $\Omega$ :

$$
\left\{\begin{array}{l}
u_{t}-\nu \Delta u+|\nabla u|^{q}=0, \quad \text { in } Q_{\Omega, T},  \tag{2.13}\\
u=0, \quad \text { on } \partial \Omega \times(0, T), \\
u(x, 0)=u_{0} \geqq 0,
\end{array}\right.
$$

Here also the problem is the object of many works, such as [22], [8], [37], [9], [33].
If $u_{0} \in C_{0}^{1}(\bar{\Omega})$, from [22], (2.13) admits a unique nonnegative solution $u \in C^{2,1}(\Omega \times(0, \infty)) \cap$ $C(\bar{\Omega} \times[0, \infty))$, such that $|\nabla u| \in C(\bar{\Omega} \times[0, \infty))$. Universal a priori estimates hold: there exist $C=C(N, q, \Omega)>0$ and a function $D \in C((0, \infty)$ such that

$$
\begin{equation*}
u(., t) \leqq C\left(1+t^{-\frac{1}{q-1}}\right) d(x, \partial \Omega), \quad|\nabla u(., t)| \leqq D(t), \tag{2.14}
\end{equation*}
$$

see [16, Remark 2.8]. The estimate on $u$ is based on the construction of supersolutions, and the estimate of the gradient is deduced from the first one by the Bernstein technique.

In case of rough initial data, a notion of mild solutions has been introduced by [8] (see definition 5.8). Such solutions satisfy $|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{1}(\Omega)\right)$.

- If $u_{0} \in \mathcal{M}_{b}^{+}(\Omega)$ and $1<q<(N+2) /(N+1)$, there is a unique nonnegative mild solution, see [8], [1]. If $u_{0} \in L^{1}(\Omega)$, and $1<q \leqq 2$, there exists at least a solution, such that $u \in C\left([0, T) ; L^{1}(\Omega)\right)$.
- If $1<q<(N+2 r) /(N+r)$ uniqueness holds in the class of mild solutions such that $u \in C\left([0, T) ; L^{r}(\Omega)\right) \cap L_{l o c}^{q}\left([0, T) ; W^{q r}(\Omega)\right)$.

Next we give decay properties and regularizing effects valid for any weak solution of the problem, in particular the universal estimate

$$
\|u(., t)\|_{L^{\infty}(\Omega)} \leqq C t^{-\frac{1}{q-1}} \quad \text { in }(0, T)
$$

where $C=C(N, q)$, see Theorem 5.12. As above we deduce uniqueness results:
Theorem 2.4 Assume that $\Omega$ is bounded.
(i) Let $1<q<(N+2) /(N+1)$, and $u_{0} \in \mathcal{M}_{b}^{+}(\Omega)$. Then there exists a unique nonnegative function $u \in C\left((0, T) ; L^{1}(\Omega)\right) \cap L_{\text {loc }}^{q}\left((0, T) ; W_{0}^{1, q}(\Omega)\right)$, solution of equation (1.1) in $\mathcal{D}^{\prime}\left(Q_{\Omega, T}\right)$, such that

$$
\lim _{t \rightarrow 0} \int_{\Omega} u(., t) \psi d x=\int_{\Omega} \psi d u_{0}, \quad \forall \psi \in C_{b}(\Omega)
$$

(ii) Let $u_{0} \in L^{r}(\Omega), r \geqq 1, u_{0} \geqq 0$, and $1<q<(N+2 r) /(N+r)$. Then there exists a unique nonnegative solution $u$ as above, such that $u \in C\left([0, T) ; L^{r}(\Omega)\right)$ and $u(., 0)=u_{0}$.

This improves the results of [8], which required assumptions up to $t=0$ for the gradient: $|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{1}(\Omega)\right)$ in case (i), $|\nabla u| \in L_{l o c}^{q}\left([0, T) ; L^{r}(\Omega)\right)$ in case (ii).

Finally we show the existence of weak solutions for any $u_{0} \in L^{r}(\Omega), r \geqq 1$, such that $u \in$ $C\left([0, T) ; L^{r}(\Omega)\right)$, see Proposition 5.17.

In Section 6 we extend some results of Section 5 to the case of the quasilinear equations (1.3), with initial data $u_{0} \in L^{r}(\Omega)$ or $u_{0} \in \mathcal{M}_{b}(\Omega)$, and $u$ may be a signed solution. In the case of equation

$$
u_{t}-\Delta_{p} u=0,
$$

several local or global $L^{\infty}$ estimates and Harnack properties have been obtained in the last decades, see for example [38], [24], [25], [30], and [23], [20] and references therein. Regularizing properties for equation (1.2) are given in [33] in a Hilbertian context in case $g=0$ or $p=2$.

Here we combine our iteration method of Section 3 with a notion of renormalized solution, developped by many authors [18], [32],[35], well adapted to rough initial data. We do not require that $u(., t) \in L^{2}(\Omega)$, but we only assume that the truncates $T_{k}(u)$ of $u$ by $\pm k(k>0)$ lie in $L^{p}\left((0, T) ; W^{1, p}(\Omega)\right)$. We prove decay and $L^{\infty}$ estimates of the following type: if $\gamma>0$, for any $r \geqq 1, p>1$ and (for simplicity) $q \neq N$, then

$$
\|u(., t)\|_{L^{\infty}(\Omega)} \leqq C t^{-\sigma}\left\|u_{0}\right\|_{L^{r}(\Omega)}^{\varpi}, \quad \sigma=\frac{1}{\frac{r q}{N}+\lambda+q-1}=\frac{N}{r q} \varpi
$$

If $\nu>0$, then for any $r \geqq 1$, and $p \neq N$ such that $p>2 N /(N+2)$,

$$
\|u(., t)\|_{L^{\infty}(\Omega)} \leqq C t^{-\tilde{\sigma}}\left\|u_{0}\right\|_{L^{r}(\Omega)}^{\tilde{\tilde{\omega}}}, \quad \tilde{\sigma}=\frac{1}{\frac{r p}{N}+p-2}=\frac{N}{r p} \tilde{\varpi} .
$$

And we deduce universal estimates as before:

$$
\|u(., t)\|_{L^{\infty}(\Omega)} \leqq C t^{-\frac{1}{q-1+\lambda}} \quad \text { if } \gamma>0 ; \quad\|u(., t)\|_{L^{\infty}(\Omega)} \leqq C t^{-\frac{1}{p-2}} \quad \text { if } \nu>0 \text { and } p>2 .
$$

Such methods can also be extended to porous media equations, and doubly nonlinear equations involving operators of the form $u \mapsto-\Delta_{p}\left(|u|^{m-1} u\right)$.

## 3 Regularization lemmas

We begin by a simple bootstrap property, used for example in [38]. We recall the proof for simplicity:
Lemma 3.1 Let $\omega \in(0,1)$ and $\sigma>0$, and $K, M>0$. Let $y$ be any positive function on $(0, T)$ such that $y(t) \leqq M t^{-\sigma}$, and for any $0<s<t<T$,

$$
y(t) \leqq K(t-s)^{-\sigma} y^{\omega}(s),
$$

Then $y$ satisfies an estimate independent of $M$ : for any $t \in(0, T)$,

$$
\begin{equation*}
y(t) \leqq 2^{\sigma(1-\omega)^{-2}}\left(K t^{-\sigma}\right)^{(1-\omega)^{-1}} \tag{3.1}
\end{equation*}
$$

Proof. We get by induction, for any $n \geqq 1$

$$
y^{\omega^{n-1}}\left(t / 2^{n-1}\right) \leqq K^{\omega^{n-1}} 2^{n \sigma \omega^{n-1}} t^{-\sigma \omega^{n-1}} y^{\omega^{n}}\left(t / 2^{n}\right),, \quad y^{\omega^{n}}\left(t / 2^{n}\right) \leqq 2^{n \sigma \omega^{n}} t^{-\sigma \omega^{n}} M^{\omega^{n}}
$$

Then

$$
y(t) \leqq K^{\sum_{k=0}^{n-1} \varpi^{k}} t^{-\sigma \sum_{k=0}^{n} \varpi^{k}} 2^{\left.\sigma \sum_{k=0}^{n}(k+1) \varpi^{k}\right)} M^{\omega^{n+1}},
$$

implying (3.1) as $n \rightarrow \infty$, since $\lim _{n \rightarrow \infty} M^{\omega^{n+1}}=1$.
Next we show the Moser's type property:
Proof of Lemma 2.1. (i) Let $\alpha$ be any real such that $\alpha \geqq r-1$, and $v(., s) \in L^{\alpha+1}(\Omega)$. From (2.2), $\int_{\Omega} v^{\alpha+1}(t) d x$ is decreasing for $t>s$. And $\int_{\Omega} v^{\beta m \theta}(., \xi) d x$ is finite for almost any $\xi \in(s, t)$, hence for a sequence $\left(\xi_{n}\right)$ decreasing to $s$,

$$
\int_{\Omega} v^{\alpha+1}(., t) d x+\frac{C_{0}(\alpha+1)}{\beta^{m}} \int_{\xi_{n}}^{t}\left(\int_{\Omega} v^{\beta m \theta}(., \xi) d x\right)^{\frac{1}{\theta}} d \xi \leqq \int_{\Omega} v^{\alpha+1}\left(., \xi_{n}\right) d x \leqq \int_{\Omega} v^{\alpha+1}(., s) d x .
$$

From (2.1), there holds $\beta m \theta>r$. Applying again (2.2) with $\beta m \theta-1$ instead of $\alpha$, and $\xi_{n}$ instead of $s$, we deduce that $\int_{\Omega} v^{\beta m \theta}(t) d x$ is decreasing for $t>s$, thus

$$
\int_{\Omega} v^{\alpha+1}(., t) d x+\frac{C_{0}(\alpha+1)}{\beta^{m}}\left(t-\xi_{n}\right)\left(\int_{\Omega} v^{\beta m \theta}\left(., \xi_{n}\right) d x\right)^{\frac{1}{\theta}} \leqq \int_{\Omega} v^{\alpha+1}(., s) d x
$$

As $n \rightarrow \infty, v\left(., \xi_{n}\right) \rightarrow v(., s)$ in $L_{l o c}^{1}(\Omega)$, and after extraction, a.e. in $\Omega$. Then from the Fatou lemma,

$$
\int_{\Omega} v^{\alpha+1}(., t) d x+\frac{C_{0}(\alpha+1)}{\beta^{m}}(t-s)\left(\int_{\Omega} v^{\beta m \theta}(., s) d x\right)^{\frac{1}{\theta}} \leqq \int_{\Omega} v^{\alpha+1}(., s) d x .
$$

Hence

$$
\begin{equation*}
\|v(t)\|_{L^{\beta m \theta}(\Omega)}^{\beta m \theta} \leqq\left(\frac{\beta^{m}}{C_{0}(\alpha+1)} \frac{1}{t-s}\|v(s)\|_{L^{\alpha+1}(\Omega)}^{\alpha+1}\right)^{\theta} \tag{3.2}
\end{equation*}
$$

- Case $r>1$. We start from $s=0$, we have $v_{0} \in L^{r}(\Omega)$. We take $\alpha_{0}=r-1$, thus $\int_{\Omega} v^{\alpha_{0}+1}(t) d x$ is finite, and set $\beta_{0}=1+\left(\alpha_{0}+\lambda\right) / m$. We define sequences $\left(t_{n}\right),\left(\alpha_{n}\right),\left(r_{n}\right),\left(\beta_{n}\right)$, by $t_{0}=0, r_{0}=r$ and for any $n \geqq 1$,

$$
t_{n}=t\left(1-\frac{1}{2^{n}}\right), \quad r_{n}=\alpha_{n}+1, \quad \beta_{n}=1+\frac{\alpha_{n}+\lambda}{m}, \quad r_{n+1}=\beta_{n} m \theta=\left(r_{n}+\lambda+m-1\right) \theta
$$

hence $\left(r_{n}\right),\left(\beta_{n}\right)$ are increasing, since $r_{1}>r$ from (2.1). In (3.2), we replace $s, t, \alpha, \beta m \theta$, by $t_{n}, t_{n+1} r_{n}, r_{n+1}$, and get

$$
\begin{equation*}
\left\|v\left(t_{n+1}\right)\right\|_{L^{r_{n+1}}(\Omega)} \leqq\left(\frac{1}{C_{0}(m \theta)^{m}} \frac{r_{n+1}^{m}}{r_{n}} \frac{1}{t_{n+1}-t_{n}}\right)^{\frac{\theta}{r_{n+1}}}\left\|v\left(t_{n}\right)\right\|_{L^{r_{n}}(\Omega)}^{\frac{\theta \cdot r_{n}}{r_{n}+1}} \tag{3.3}
\end{equation*}
$$

From (2.2), it follows that

$$
\begin{equation*}
\|v(t)\|_{L^{r_{n+1}}(\Omega)} \leqq\left\|v\left(t_{n+1}\right)\right\|_{L^{r_{n+1}}(\Omega)} \leqq I_{n} J_{n} L_{n}\left\|v_{0}\right\|_{L^{r}(\Omega)}^{\frac{\theta^{n+1} . r}{r_{n+1}}} \tag{3.4}
\end{equation*}
$$

where

$$
I_{n}=\prod_{k=1}^{n+1}\left(\frac{r_{k}^{m}}{r_{k-1}}\right)^{\frac{\theta^{n+2-k}}{r_{n+1}}}, \quad J_{n}=\prod_{k=1}^{n+1}\left(\frac{1}{t_{k}-t_{k-1}}\right)^{\frac{\theta^{n+2-k}}{r_{n+1}}}, \quad L_{n}=\left(C_{0}(m \theta)^{m}\right)^{-\sum_{k=1}^{n+1} \frac{\theta^{n+2-k}}{r_{n+1}}} .
$$

Since $r_{n}=\theta^{n}\left(r+(\lambda+m-1) \theta^{\prime}\left(1-\theta^{-n}\right)\right)$, we find

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \frac{\theta^{n+1} r}{r_{n+1}}=\varpi_{r, m, \lambda, \theta}, \quad \lim _{n \rightarrow \infty} \frac{1}{r_{n+1}} \sum_{k=1}^{n+1} \theta^{n+2-k}=\sigma_{r, m, \lambda, \theta}, \quad \lim _{n \rightarrow \infty} \sum_{k=1}^{n+1} k \theta^{1-k}=\theta^{\prime 2} \tag{3.5}
\end{equation*}
$$

As a consequence

$$
\begin{equation*}
\lim _{n \rightarrow \infty} J_{n}=2^{-\frac{\varpi_{r, m, \lambda, \theta}}{r} \theta^{\prime 2}} t^{-\sigma_{r, m, \lambda, \theta}}, \quad \lim _{n \rightarrow \infty} L_{n}=\left(C_{0}(m \theta)^{m}\right)^{-\sigma_{r, m, \lambda, \theta}} \tag{3.6}
\end{equation*}
$$

Otherwise

$$
\ln I_{n}=\frac{m}{r_{n+1}} \sum_{k=1}^{n+1} \theta^{n+2-k} \ln r_{k}-\frac{1}{r_{n+1}} \sum_{k=0}^{n} \theta^{n+1-k} \ln r_{k}=\frac{\theta^{n+1}}{r_{n+1}}\left(m \theta \sum_{k=1}^{n+1} \theta^{-k} \ln r_{k}-\sum_{k=0}^{n} \theta^{-k} \ln r_{k}\right)
$$

and the sum $S=\sum_{k=0}^{\infty} \theta^{-k} \ln r_{k}$ is finite, since $r_{k} \leqq \theta^{k}\left(r+|\lambda+m-1| \theta^{\prime}\right)$. Then $I_{n}$ has a finite limit $\ell=\ell(N, m, r, \lambda, \theta)=\exp \left(r^{-1} \varpi_{r, m, \lambda, \theta}((m \theta-1) S-m \theta \ln r)\right)$. Thus we can go to the limit in (3.4), and the conclusion follows.

- Case $r=1$. If $\left(\mathrm{H}_{1}\right)$ holds we can take $\alpha_{0}=r-1=0$ and the proof is done. Next assume $\left(\mathrm{H}_{2}\right)$. Then we obtain, for any $0 \leqq s<t<T$, and a constant $C$ as before,

$$
\begin{aligned}
\|v(., t)\|_{L^{\infty}(\Omega)} & \leqq C(t-s)^{-\sigma_{\rho, m, \lambda, \theta}}\|v(., s)\|_{L^{\rho}(\Omega)}^{\varpi_{\rho, m, \lambda, \theta}} \\
& \leqq C(t-s)^{-\sigma_{\rho, m, \lambda, \theta}}\|v(., s)\|_{L^{\infty}(\Omega)}^{\varpi_{\rho, m, \lambda}(\rho-1) / \rho}\|v(., s)\|_{L^{1}(\Omega)}^{\varpi_{\rho, m, \lambda, \theta} / \rho} \\
& \leqq C\left\|v_{0}\right\|_{L^{1}(\Omega)}^{\varpi_{\rho, m, \lambda, \theta} / \rho}(t-s)^{-\sigma_{\rho, m, \lambda, \theta}}\|v(., s)\|_{L^{\infty}(\Omega)}^{\varpi_{\rho, m, \lambda, \theta}(\rho-1) / \rho}
\end{aligned}
$$

Let $y(t)=\|v(., t)\|_{L^{\infty}(\Omega)}$. We can apply Lemma 3.1 to $y$, with

$$
\sigma=\sigma_{\rho, m, \lambda, \theta}, \quad \omega=\frac{\varpi_{\rho, m, \lambda, \theta}}{\rho^{\prime}}, \quad K=C\left\|v_{0}\right\|_{L^{1}(\Omega)}^{\varpi_{\rho, m, \lambda, \theta} / \rho}, \quad M=C\left\|v_{0}\right\|_{L^{\rho}(\Omega)}^{\varpi_{\rho, m, \lambda, \theta / \rho}}
$$

Indeed $\omega<1$ since $\rho \theta^{\prime}(1-m-\lambda)<1$. Then there holds

$$
\|v(., t)\|_{L^{\infty}(\Omega)} \leqq 2^{\sigma(1-\omega)^{-2}}\left(K t^{-\sigma}\right)^{(1-\omega)^{-1}}=2^{\sigma(1-\omega)^{-2}} C^{(1-\omega)^{-1}} t^{-\sigma(1-\omega)^{-1}}\left\|v_{0}\right\|_{L^{1}(\Omega, \lambda)}^{\varpi_{\rho, \lambda, \theta} / \rho((1-\omega))} .
$$

Noticing that $\sigma(1-\omega)^{-1}=\sigma_{1, m, \lambda, \theta}$ and $\varpi_{\rho, m, \lambda, \theta} / \rho((1-\omega))=\varpi_{1, m, \lambda, \theta}$, we deduce

$$
\begin{equation*}
\|v(., t)\|_{L^{\infty}(\Omega)} \leqq C t^{-\sigma_{1, m, \lambda, \theta}}\left\|v_{0}\right\|_{L^{1}(\Omega)}^{\varpi_{1, \lambda, \lambda, \theta}} \tag{3.7}
\end{equation*}
$$

with a new constant $C$, now depending on $\rho$.
Remark 3.2 This lemma can be compared with the result of [33, Theorem 2.1] obtained by the Stampacchia's method. In order to obtain decay estimates for the solutions $u$ of a parabolic equation such as (1.1) or (1.3), the Moser's method consists to take as test functions powers $|u|^{\alpha-1} u$ of $u$; the Stampacchia's method uses test functions of the form $(u-k)^{+} \operatorname{sign} u$. If one applies to sufficiently smooth solutions, both techniques leed to decay estimates of the same type. In the case of weaker solutions, the Stampacchia method supposes that the functions $(u-k)^{+}$are admissible in the equation, which leads to assume that $u(., t) \in W^{1,2}(\Omega)$, see [33]. In the sequel we combine Moser's method with regularization or truncature of $u$, in order to admit powers as test functions. So we do not need to make this assumption, thus the Moser's method appears to be more performant.

Such type of $L^{\infty}$ estimates as (2.3) may imply a universal one, that means independent of the initial data, in case $\Omega$ is bounded. This was observed for example in [38]:
Lemma 3.3 Let $\Omega$ be bounded. (i) Let $v \in C\left([0, T) ; L_{\text {loc }}^{1}(\Omega)\right)$ be nonnegative, and $v_{0}=v(x, 0) \in$ $L^{1}(\Omega)$, such that for some $C>0$, for any $0 \leqq s<t<T$,

$$
\|v(., t)\|_{L^{\infty}(\Omega)} \leqq C(t-s)^{-\sigma}\|v(., s)\|_{L^{1}(\Omega)}^{\varpi},
$$

where $\sigma>0, \varpi \in(0,1)$. Then there exists $M=M(C, \sigma, \varpi,|\Omega|)$ such that for any $t \in(0, T)$,

$$
\begin{equation*}
\|v(., t)\|_{L^{\infty}(\Omega)} \leqq M t^{-\frac{\sigma}{1-w}} . \tag{3.8}
\end{equation*}
$$

(ii) As a consequence, if $v$ satisfies (2.3), with $m-1+\lambda>0$, then

$$
\begin{equation*}
\|v(., t)\|_{L^{\infty}(\Omega)} \leqq M t^{-\frac{1}{m-1+\lambda}} . \tag{3.9}
\end{equation*}
$$

Proof. (i) For any $0<s<t<T$,

$$
\|v(., t)\|_{L^{\infty}(\Omega)} \leqq C(t-s)^{-\sigma}\|v(., s)\|_{L^{r}(\Omega)}^{\infty} \leqq C(t-s)^{-\sigma}|\Omega|^{\infty}\|v(., s)\|_{L^{\infty}(\Omega)}^{\infty}
$$

Since $\varpi<1$, (3.8) follows from Lemma 3.1: for any $t \in(0, T)$,

$$
\|v(., t)\|_{L^{\infty}(\Omega)} \leqq 2^{\sigma(1-\varpi)^{-2}}\left(C|\Omega|^{\varpi} t^{-\sigma}\right)^{(1-\varpi)^{-1}} .
$$

(ii) If $v$ satisfies (2.3), with $m-1+\lambda>0$, we take $\sigma=\sigma_{r, m, \lambda, \theta}$, and $\varpi=\varpi_{r, m, \lambda, \theta}$ defined at (2.4), then $\varpi=\left(1+(m-1+\lambda) \theta^{\prime} / r\right)^{-1}<1$ and $\sigma\left((1-\varpi)^{-1}=(m-1+\lambda)^{-1}\right.$, which proves (3.9).

In the sequel Lemma 2.1 is applied in situations where (2.2) comes from an estimate of $v$ in a Sobolev Space $W^{1, m}\left(Q_{\Omega, s, t}\right)$, when $1<m<N$, with $\theta=N /(N-m)$, or $m=N$ and $\theta>1$ is arbitrary.

In the case $m>N$, where Lemma 2.1 does not bring information, we use in the sequel a limit form of Gagliardo-Nirenberg inequality, see the proof of Theorems 4.16, 5.12 and 6.7:

Lemma 3.4 Let $m>N$, and $r \geqq 1$. Let $\Omega$ be any domain in $\mathbb{R}^{N}$. Then there exists $C=$ $C(N, m, r)>0$ such that for any $w \in L^{r}(\Omega) \cap W_{0}^{1, m}(\Omega)$,

$$
\|w\|_{L^{\infty}(\Omega)} \leqq C\|\nabla w\|_{L^{m}(\Omega)}^{k}\|w\|_{L^{r}(\Omega)}^{1-k}, \quad \frac{1}{k}=1+r\left(\frac{1}{N}-\frac{1}{m}\right)
$$

Proof. By extension by 0 outside of $\Omega$, we can assume $\Omega=\mathbb{R}^{N}$. Since $m>N$, for any $\varphi \in \mathcal{D}\left(\mathbb{R}^{N}\right)$, and any $x \in \mathbb{R}^{N}$,

$$
|\varphi(x)| \leqq C(N, m)\left(\left|\int_{B(x, 1)} \varphi d x\right|+\|\nabla \varphi\|_{L^{m}(B(x, 1))}\right) \leqq C(N, m, r)\left(\|\varphi\|_{L^{r}\left(\mathbb{R}^{N}\right)}+\|\nabla \varphi\|_{L^{m}\left(\mathbb{R}^{N}\right)}\right)
$$

by density, there holds

$$
\|w\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq C\left(\|w\|_{L^{r}\left(\mathbb{R}^{N}\right)}+\|\nabla w\|_{L^{m}\left(\mathbb{R}^{N}\right)}\right)
$$

for any $w \in L^{r}\left(\mathbb{R}^{N}\right) \cap W^{1, m}\left(\mathbb{R}^{N}\right)$. Setting $w_{t}(x)=w(t x)$ for any $t>0$, we find

$$
\|w\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}=\left\|w_{t}\right\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq C\left(t^{-\frac{N}{r}}\|w\|_{L^{r}\left(\mathbb{R}^{N}\right)}+t^{\frac{m-N}{m}}\|\nabla w\|_{L^{m}\left(\mathbb{R}^{N}\right)}\right)
$$

the result follows by taking $t=\left(\|w\|_{L^{r}\left(\mathbb{R}^{N}\right)} /\|\nabla w\|_{L^{m}\left(\mathbb{R}^{N}\right)}\right)^{1 /(1-N / m+N / r)}$.

## 4 The Hamilton-Jacobi equation in $\mathbb{R}^{N}$

### 4.1 Different notions of solution

In this section we study the Cauchy problem (2.5).
Here we consider the solutions in a weak sense, which does not use any formulation in terms of semigroups:
Definition 4.1 We say that a nonnegative function $u$ is a weak solution (resp. subsolution) of equation of (1.1) in $Q_{\mathbb{R}^{N}, T}$, if $u \in L_{l o c}^{1}\left(Q_{\mathbb{R}^{N}, T}\right)$, and $|\nabla u| \in L_{l o c}^{q}\left(Q_{\mathbb{R}^{N}, T}\right)$, and

$$
\begin{equation*}
\int_{0}^{T} \int_{\Omega}\left(-u \varphi_{t}-u \Delta \varphi+|\nabla u|^{q} \varphi\right) d x d t=0, \quad(\text { resp. } \leqq), \quad \forall \varphi \in \mathcal{D}\left(Q_{\mathbb{R}^{N}, T}\right), \varphi \geqq 0 \tag{4.1}
\end{equation*}
$$

Remark 4.2 From [16], any nonnegative weak solution satisfies

$$
\begin{equation*}
u \in L_{l o c}^{\infty}\left(Q_{\mathbb{R}^{N}, T}\right), \quad \nabla u \in L_{l o c}^{2}\left(Q_{\mathbb{R}^{N}, T}\right), \quad u \in C\left((0, T) ; L_{l o c}^{\rho}\left(\mathbb{R}^{N}\right)\right) \quad \forall \rho \geqq 1 \tag{4.2}
\end{equation*}
$$

Hence (4.1) is equivalent to:

$$
\begin{equation*}
\int_{0}^{T} \int_{\Omega}\left(-u \varphi_{t}+\nabla u . \nabla \varphi+|\nabla u|^{q} \varphi\right) d x d t=0, \quad \forall \varphi \in \mathcal{D}\left(Q_{\mathbb{R}^{N}, T}\right) \tag{4.3}
\end{equation*}
$$

and there holds, for any $s, \tau \in(0, T)$,

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} u(., \tau) \varphi(., \theta) d x-\int_{\mathbb{R}^{N}} u(., s) \varphi(., s) d x+\int_{s}^{\tau} \int_{\mathbb{R}^{N}}\left(-u \varphi_{t}+\nabla u . \nabla \varphi+|\nabla u|^{q} \varphi\right) d x d t=0 \tag{4.4}
\end{equation*}
$$

and for any $\psi \in C_{c}^{2}\left(\mathbb{R}^{N}\right)$,

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} u(., \tau) \psi d x-\int_{\mathbb{R}^{N}} u(., s) \psi d x+\int_{s}^{\tau} \int_{\mathbb{R}^{N}}\left(\nabla u . \nabla \psi+|\nabla u|^{q} \psi d x d t=0\right. \tag{4.5}
\end{equation*}
$$

Definition 4.3 Let $u_{0} \in L_{l o c}^{r}\left(\mathbb{R}^{N}\right), r \geqq 1$.
We say that $u$ is a weak $L_{\text {loc }}^{r}$ solution if $u$ is a weak solution of (1.1) and the extension of $u$ by $u_{0}$ at time 0 satisfies $u \in C\left([0, T) ; L_{\text {loc }}^{r}\left(\mathbb{R}^{N}\right)\right)$.

We say that $u$ is a weak $r$ solution of problem (2.5) if it is a weak solution of equation (1.1) such that

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{\mathbb{R}^{N}} u^{r}(., t) \psi d x=\int_{\mathbb{R}^{N}} u_{0}^{r} \psi d x, \quad \forall \psi \in C_{c}\left(\mathbb{R}^{N}\right) . \tag{4.6}
\end{equation*}
$$

Definition 4.4 Let $u_{0}$ be any nonnegative Radon measure in $\mathbb{R}^{N}$, we say that $u$ is a weak $\mathcal{M}_{\text {loc }}$ solution of problem (2.5) if it is a weak solution of (1.1) such that

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{\mathbb{R}^{N}} u(., t) \psi d x=\int_{\mathbb{R}^{N}} \psi d u_{0}, \quad \forall \psi \in C_{c}\left(\mathbb{R}^{N}\right) . \tag{4.7}
\end{equation*}
$$

Remark 4.5 Obviously, any weak $L_{\text {loc }}^{r}$ solution is a weak $r$ solution. When $r=1$, the notions of weak 1 -solution and weak $\mathcal{M}_{\text {loc }}$ solution coincide. When $r>1$, it can be easily checked that $u$ is a weak $L_{\text {loc }}^{r}$ solution if and only if it is a weak $r$ solution and

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{\mathbb{R}^{N}} u(., t) \psi d x=\int_{\mathbb{R}^{N}} u_{0} \psi d x, \quad \forall \psi \in C_{c}\left(\mathbb{R}^{N}\right) \tag{4.8}
\end{equation*}
$$

Other types of solutions using the semigroup of the heat equation have been introduced in ([15]):

Definition 4.6 Let $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right)$. A function $u$ is called mild $L^{r}$ solution of problem (2.5) if $u \in C\left([0, T) ; L^{r}\left(\mathbb{R}^{N}\right)\right)$, and $|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{r}\left(\mathbb{R}^{N}\right)\right)$ and

$$
u(., t)=e^{t \Delta} u_{0}-\int_{0}^{t} e^{(t-s) \Delta}|\nabla u(., s)|^{q} d s \quad \text { in } L^{r}\left(\mathbb{R}^{N}\right) ;
$$

here $e^{t \Delta}$ is the semi-group of the heat equation acting on $L^{r}\left(\mathbb{R}^{N}\right)$.
Definition 4.7 Let $u_{0} \in \mathcal{M}_{b}^{+}\left(\mathbb{R}^{N}\right)$. A function $u$ is called mild $\mathcal{M}$ solution of (2.5) if $u \in$ $C_{b}\left((0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right)$ and $|\nabla u|^{q} \in L_{\text {loc }}^{1}\left([0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right)$, and for any $0<t<T$,

$$
\begin{equation*}
u(., t)=e^{t \Delta} u_{0}(.)-\int_{0}^{t} e^{(t-s) \Delta}|\nabla u(., s)|^{q} d s \quad \text { in } L^{1}\left(\mathbb{R}^{N}\right), \tag{4.9}
\end{equation*}
$$

where $e^{t \Delta}$ is defined on $\mathcal{M}_{b}^{+}\left(\mathbb{R}^{N}\right)$ as the adjoint of the operator $e^{t \Delta}$ on $C_{0}\left(\mathbb{R}^{N}\right)$, the space of continuous functions on $\mathbb{R}^{N}$ which tend to 0 as $|x| \rightarrow \infty$.

Remark 4.8 Every mild $L^{r}$ solution is a weak $L_{\text {loc }}^{r}$ solution. Any mild $\mathcal{M}$ solution is a weak $\mathcal{M}_{\text {loc }}$ solution. Indeed for any $0<\epsilon<t<T$, we find

$$
u(., t)=e^{(t-\epsilon) \Delta} u(., \epsilon)-\int_{\epsilon}^{t} e^{(t-s) \Delta}|\nabla u(., s)|^{q} d s \quad \text { in } L^{1}\left(\mathbb{R}^{N}\right) ;
$$

and $u(., \epsilon) \in L^{1}\left(\mathbb{R}^{N}\right)$, thus $u$ is a weak solution on $(\epsilon, T)$, then on $(0, T)$. As $t \rightarrow 0, u(., t)-e^{t \Delta} u_{0}($. converges to 0 in $L^{1}\left(\mathbb{R}^{N}\right)$, then weakly *, and $e^{t \Delta} u_{0}(.) \rightarrow u_{0}$ weakly ${ }^{*}$, hence (4.7) holds.

Another definition of solution with initial data measure was given in ([12]):
Definition 4.9 Let $u_{0} \in \mathcal{M}_{b}^{+}\left(\mathbb{R}^{N}\right)$. A function $u$ is called weak semi-group solution if $u \in$ $C\left((0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right)$ and $|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right)$ and for any $0<\epsilon<t<T$,

$$
\begin{gather*}
u(., t)=e^{(t-\epsilon) \Delta} u(., \epsilon)-\int_{\epsilon}^{t} e^{(t-s) \Delta}|\nabla u(., s)|^{q} d s \quad \text { in } L^{1}\left(\mathbb{R}^{N}\right)  \tag{4.10}\\
\lim _{t \rightarrow 0} \int_{\mathbb{R}^{N}} u(., t) \varphi d x=\int_{\mathbb{R}^{N}} \varphi d u_{0}, \quad \forall \varphi \in C_{b}\left(\mathbb{R}^{N}\right) \tag{4.11a}
\end{gather*}
$$

In fact the two definitions coincide, see the proof in the Appendix:
Lemma 4.10 Let $u_{0} \in \mathcal{M}_{b}^{+}\left(\mathbb{R}^{N}\right)$. Then
$u$ is a mild $\mathcal{M}$ solution of (2.5) $\Longleftrightarrow u$ is a weak semi-group solution of (2.5).
Remark 4.11 All these definitions of semi-group solutions assume a global in space condition: $|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right)$ or $|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{r}\left(\mathbb{R}^{N}\right)\right)$. Observe also that (4.11a) is assumed for any $\varphi \in C_{b}\left(\mathbb{R}^{N}\right)$. On the contrary, our definitions of weak solutions are local in space and time, they do not require such global properties.

Finally we mention another weaker form of semi-group solutions, given in ([15]), which will be used in the sequel:

Definition 4.12 Let $u_{0} \in \mathcal{M}_{b}^{+}\left(\mathbb{R}^{N}\right)$. Then $u$ is a pointwise mild solution of (2.5) if $u \in L_{\text {loc }}^{1}\left(Q_{\mathbb{R}^{N}, T}\right)$, and $|\nabla u|^{q} \in L_{l o c}^{1}\left(Q_{\mathbb{R}^{N}, T}\right)$, and

$$
u(x, t)=\left(e^{t \Delta} u_{0}\right)(x)-\int_{0}^{t} \int_{\mathbb{R}^{N}} g(x-y, t-s)|\nabla u(y, s)|^{q} d y d s \quad \text { for a.e. }(x, t) \in Q_{\mathbb{R}^{N}, T}
$$

where $g$ is the heat kernel.
Remark 4.13 For $r \geqq 1$, it is clear that every mild $L^{r}$ solution is a pointwise mild solution. If $u_{0} \in L^{1}\left(\mathbb{R}^{N}\right)$ every pointwise mild solution is a mild $L^{1}$ solution; if $u_{0} \in \mathcal{M}_{b}^{+}\left(\mathbb{R}^{N}\right)$, every pointwise mild solution, is a mild $\mathcal{M}$ solution. see [15, Proposition 1.1 and Remark 1.2].

### 4.2 Decay of the norms

Next we show a decay result for the solutions of Hamilton Jacobi equations, which is valid for any $q>1$, and for all the weak solutions, with no condition of boundedness at infinity.

When $q \leqq 2$, any weak solution $u$ of equation (1.1) is smooth: $u \in C^{2,1}\left(Q_{\mathbb{R}^{N}, T}\right)$, from [16, Theorem 2.9]. Since it may be false for $q>2$, we regularize $u$ by convolution, setting

$$
u_{\varepsilon}=u * \varrho_{\varepsilon},
$$

where $\left(\varrho_{\varepsilon}\right)_{\varepsilon>0}$ is a sequence of mollifiers. We recall that for given $0<s<\tau<T$, and $\varepsilon$ small enough, $u_{\varepsilon}$ is a subsolution of equation (1.1), see [16]:

$$
\begin{equation*}
\left(u_{\varepsilon}\right)_{t}-\nu \Delta u_{\varepsilon}+\left|\nabla u_{\varepsilon}\right|^{q} \leqq 0, \quad \text { in } Q_{\mathbb{R}^{N}, s, \tau} \tag{4.12}
\end{equation*}
$$

Theorem 4.14 Assume $q>1$. Let $r \geqq$. Let $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right)$ be nonnegative. Let $u$ be any nonnegative weak $r$ solution of problem (2.5).
(i) Then $u(., t) \in L^{r}\left(\mathbb{R}^{N}\right)$ for any $t \in(0, T)$, and

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} u^{r}(., t) d x \leqq \int_{\mathbb{R}^{N}} u_{0}^{r} d x \tag{4.13}
\end{equation*}
$$

(ii) Moreover $u^{r-1}|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right)$; and $u^{r-2}|\nabla u|^{2} \in L_{l o c}^{1}\left([0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right)$ if $r>1$ and $\nu>0$. For any $t \in(0, T)$, we have the equalities

$$
\begin{gather*}
\int_{\mathbb{R}^{N}} u^{r}(., t) d x+r \int_{0}^{t} \int_{\mathbb{R}^{N}} u^{r-1}|\nabla u|^{q} d x d t+r(r-1) \nu \int_{0}^{t} \int_{\mathbb{R}^{N}} u^{r-2}|\nabla u|^{2} d x d t=\int_{\mathbb{R}^{N}} u_{0}^{r} d x,  \tag{4.15}\\
\int_{\mathbb{R}^{N}} u(., t) d x+\int_{0}^{t} \int_{\mathbb{R}^{N}}|\nabla u|^{q} d x d t=\int_{\mathbb{R}^{N}} u_{0} d x, \quad \text { if } r=1  \tag{4.14}\\
\lim _{t \rightarrow 0} \int_{\mathbb{R}^{N}} u^{r}(., t) d x=\int_{\mathbb{R}^{N}} u_{0}^{r} d x \tag{4.16}
\end{gather*}
$$

(iii) $u^{q-1+r} \in L_{l o c}^{1}\left(\left([0, T) ; W^{1,1}\left(\mathbb{R}^{N}\right)\right)\right.$; and if $\nu>0$, then $u^{r / 2} \in L_{l o c}^{2}\left([0, T) ; W^{1,2}\left(\mathbb{R}^{N}\right)\right)$.
(iv) If $u$ is a weak $L_{l o c}^{r}$ solution, then $u \in C\left([0, T) ; L^{r}\left(\mathbb{R}^{N}\right)\right)$.

Proof. (i) First step: case $q^{\prime}>N / r$. That means $r \geqq N$ or $q$ is small enough: $1<$ $q<N /(N-r)$. Let $0<s<\tau<T$. Take $\varepsilon>0$ small enough such that (4.12) holds. Let $\delta>0$, and $u_{\varepsilon, \delta}=u_{\varepsilon}+\delta$, so that $u_{\varepsilon, \delta}^{r-2}$ is well defined for $r<2$. For any $R>0$, we consider $\xi(x)=\xi_{R}(x)=\psi(x / R)$, where $\psi(x) \in[0,1], \psi(x)=1$ for $|x| \leqq 1, \psi(x)=0$ for $|x| \geqq 2$. Multiplying (4.12) by $u_{\varepsilon, \delta}^{r-1} \xi^{\lambda}$ where $\lambda>0$, we get for any $t \in[s, \tau]$,

$$
\begin{aligned}
& \frac{d}{d t}\left(\frac{1}{r} \int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r}(., t) \xi^{\lambda} d x\right)+(r-1) \nu \int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r-2}\left|\nabla u_{\varepsilon, \delta}\right|^{2}(., t) \xi^{\lambda} d x+\int_{\mathbb{R}^{N}}\left|\nabla u_{\varepsilon, \delta}\right|^{q} u_{\varepsilon, \delta}^{r-1} \xi^{\lambda-1} \xi^{\lambda} d x \\
& \leqq-\lambda \int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r-1} \xi^{\lambda-1} \xi^{\lambda-1} \nabla u_{\varepsilon, \delta} \nabla \xi d x
\end{aligned}
$$

and from the Hölder inequality, with $C=C(q, \lambda)$

$$
\begin{aligned}
\lambda \int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r-1}\left|\nabla u_{\varepsilon, \delta}\right|(., t) \xi^{\lambda-1}|\nabla \xi| d x & \leqq \frac{1}{2} \int_{\mathbb{R}^{N}}\left|\nabla u_{\varepsilon, \delta}(., t)\right|^{q} u_{\varepsilon, \delta}^{r-1} \xi^{\lambda} d x+C \int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r-1}(., t) \xi^{\lambda-q^{\prime}}|\nabla \xi|^{q^{\prime}} d x \\
\int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r-1}(., t) \xi^{\lambda-q^{\prime}}|\nabla \xi|^{q^{\prime}} d x & \leqq\left(\int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r}(., t) \xi^{\lambda} d x\right)^{\frac{1}{r^{\prime}}}\left(\int_{\mathbb{R}^{N}} \xi^{\lambda-r q^{\prime}}|\nabla \xi|^{r q^{\prime}} d x\right)^{\frac{1}{r}}
\end{aligned}
$$

Choosing $\lambda=r q^{\prime}$ we deduce

$$
\frac{d}{d t}\left(\left(\int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r}(., t) \xi^{\lambda} d x\right)^{\frac{1}{r}}\right) \leqq C R^{\frac{N}{r}-q^{\prime}}
$$

where $C=C(N, q, r, \psi)$. By integration,

$$
\left(\int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r}(., t) \xi^{\lambda} d x\right)^{\frac{1}{r}} \leqq\left(\int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r}(., s) \xi^{\lambda} d x\right)^{\frac{1}{r}}+C \tau R^{\frac{N}{r}-q^{\prime}}
$$

with a new constant $C$ as above. Let $R_{0}>0$ be fixed and take $R>R_{0}$, thus

$$
\left(\int_{B_{R_{0}}} u_{\varepsilon, \delta}^{r}(., t) d x\right)^{\frac{1}{r}} \leqq\left(\int_{B_{2 R}} u_{\varepsilon, \delta}^{r}(., s) \xi^{\lambda} d x\right)^{1 / r}+C \tau R^{\frac{N}{r}-q^{\prime}}
$$

As $\delta \rightarrow 0$, and then as $\varepsilon \rightarrow 0$, we deduce that

$$
\begin{equation*}
\left(\int_{B_{R_{0}}} u(., t)^{r} d x\right)^{\frac{1}{r}} \leqq\left(\int_{\mathbb{R}^{N}} u(., s)^{r} \xi^{\lambda} d x\right)^{\frac{1}{r}}+C \tau R^{\frac{N}{r}-q^{\prime}} \tag{4.17}
\end{equation*}
$$

for any $0<s<t<T$; from (4.6) we obtain, as $s \rightarrow 0$,

$$
\left(\int_{B_{R_{0}}} u(., t)^{r} d x\right)^{\frac{1}{r}} \leqq\left(\int_{\mathbb{R}^{N}} u_{0}^{r} \xi^{\lambda} d x\right)^{\frac{1}{r}}+C \tau R^{\frac{N}{r}-q^{\prime}} \leqq\left(\int_{\mathbb{R}^{N}} u_{0}^{r} d x\right)^{\frac{1}{r}}+C \tau R^{\frac{N}{r}-q^{\prime}}
$$

Finally (4.13) follows as $R \rightarrow \infty$ and then as $R_{0} \rightarrow \infty$.
Second step: case $q^{\prime} \leqq N / r$. Then $r<N$ and $q \geqq N /(N-r)>1$. Let us fix some $k \in(1, N /(N-r))$. For any $\eta \in(0,1)$, we have $\eta|\nabla u|^{k} \leqq \eta+|\nabla u|^{q}$, hence the function

$$
w_{\eta}=\eta^{1 /(k-1)}(u-\eta t)
$$

satisfies

$$
\left(w_{\eta}\right)_{t}-\nu \Delta w_{\eta}+\left|\nabla w_{\eta}\right|^{k} \leqq 0
$$

in the weak sense. Thanks to Kato's inequality, see [21], [5], we deduce that

$$
\begin{equation*}
\left(w_{\eta}^{+}\right)_{t}-\nu \Delta w_{\eta}^{+}+\left|\nabla w_{\eta}^{+}\right|^{k} \leqq 0 \tag{4.18}
\end{equation*}
$$

in $\mathcal{D}^{\prime}\left(Q_{\mathbb{R}^{N}, T}\right)$. And $w_{\eta}^{+}$has the same regularity as $u$. Moreover it satisfies an analogous property to (4.6):

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{\mathbb{R}^{N}}\left(w_{\eta}^{+}\right)^{r}(., t) \psi d x=\int_{\mathbb{R}^{N}}\left(\eta^{1 /(k-1)} u_{0}\right)^{r} \psi d x, \quad \forall \psi \in C_{c}\left(\mathbb{R}^{N}\right) \tag{4.19}
\end{equation*}
$$

Indeed

$$
\begin{aligned}
\left.\mid \int_{\mathbb{R}^{N}}\left((u-\eta t)^{+}\right)^{r}-u^{r}(., t)\right) \psi d x \mid & \left.\leqq \int_{\{u \geqq \eta t\}}\left|(u(., t)-\eta t)^{r}-u^{r}(., t)\right| \psi d x+\int_{\{u \leqq \eta t\}} u^{r}(., t)\right) \psi d x \\
& \leqq r \eta t \int_{\mathbb{R}^{N}} u^{r-1}(., t) \psi d x+C(\psi) t^{r} \\
& \leqq r \eta t\left(\int_{\mathbb{R}^{N}} u^{r}(., t) \psi d x\right)^{1 / r^{\prime}}\left(\int_{\mathbb{R}^{N}} \psi d x\right)^{1 / r}+C(\psi) t^{r}
\end{aligned}
$$

then

$$
\left.\lim _{t \rightarrow 0} \int_{\mathbb{R}^{N}}\left((u-\eta t)^{+}\right)^{r}-u^{r}(., t)\right) \psi d x=0
$$

and (4.19) follows from (4.6) applied to $\eta^{1 /(k-1)} u$. Since $k^{\prime}>N / r$, we can apply the first step to $w_{\eta}^{+}$; we deduce that $w_{\eta}^{+}(t) \in L^{r}\left(\mathbb{R}^{N}\right)$ and

$$
\int_{\mathbb{R}^{N}}\left(w_{\eta}^{+}\right)^{r}(., t) d x \leqq \eta^{\frac{r}{k-1}} \int_{\mathbb{R}^{N}} u_{0}^{r} d x
$$

Then $\left\|(u-\eta t)^{+}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)} \leqq\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}$. Since $u \leqq \eta t+(u-\eta t)^{+}$, we find, for any $R>0$,

$$
\|u(., t)\|_{L^{r}\left(B_{R}\right)} \leqq\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}+\eta t\left|B_{R}\right|^{\frac{1}{r}}
$$

As $\eta \rightarrow 0$ we get $\|u(., t)\|_{L^{r}\left(B_{R}\right)} \leqq\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}$, then as $R \rightarrow \infty$ we deduce that $u(., t) \in L^{r}\left(\mathbb{R}^{N}\right)$, and (4.13) holds.
(ii) Consider again $0<s<\tau<T$ and $u_{\varepsilon, \delta}$ as above. Setting $F_{\varepsilon}=|\nabla u|^{q} * \varrho_{\varepsilon}$, there holds

$$
\left(u_{\varepsilon, \delta}\right)_{t}-\nu \Delta u_{\varepsilon, \delta}+F_{\varepsilon}=0
$$

Then for any $t \in[s, \tau]$,

$$
\begin{aligned}
& \frac{d}{d t}\left(\int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r}(., t) \xi^{\lambda} d x\right)+r(r-1) \nu \int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r-2}\left|\nabla u_{\varepsilon, \delta}\right|^{2}(., t) \xi^{\lambda} d x+r \int_{\mathbb{R}^{N}} F_{\varepsilon} u_{\varepsilon, \delta}^{r-1}(., t) \xi^{\lambda} d x \\
& =-r \nu \int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r-1} \nabla u_{\varepsilon, \delta}(., t) \cdot \nabla\left(\xi^{\lambda}\right) d x=\nu \int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r}(., t) \Delta\left(\xi^{\lambda}\right) d x
\end{aligned}
$$

thus

$$
\begin{array}{r}
\int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r}(., t) \xi^{\lambda} d x+r \int_{s}^{t} \int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r-1} F_{\varepsilon} \xi^{\lambda} d x d t \\
+r(r-1) \nu \int_{s}^{t} \int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r-2}\left|\nabla u_{\varepsilon, \delta}\right|^{2} \xi^{\lambda} d x=\int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r}(., s) \xi^{\lambda} d x+\nu \int_{s}^{t} \int_{\mathbb{R}^{N}} u_{\varepsilon, \delta}^{r} \Delta\left(\xi^{\lambda}\right) d x
\end{array}
$$

First we go to the limit as $\varepsilon \rightarrow 0$, because $u \in L_{l o c}^{\infty}\left(Q_{\mathbb{R}^{N}, T}\right)$, and $|\nabla u|^{2} \in L_{l o c}^{1}\left(Q_{\mathbb{R}^{N}, T}\right)$, and $F_{\varepsilon}$ converges to $|\nabla u|^{q}$ in $L_{l o c}^{1}\left(Q_{\mathbb{R}^{N}, T}\right)$. Setting $v_{\delta}=u+\delta$, we obtain for almost any $s, t$, and by continuity for any $0<s<t \leqq \tau$,

$$
\begin{array}{r}
\int_{\mathbb{R}^{N}} v_{\delta}^{r}(., t) \xi^{\lambda} d x+r \int_{s}^{t} \int_{\mathbb{R}^{N}} v_{\delta}^{r-1}|\nabla u|^{q} \psi d x d t \\
+r(r-1) \nu \int_{s}^{t} \int_{\mathbb{R}^{N}} v_{\delta}^{r-2}|\nabla u|^{2} \xi^{\lambda} d x=\int_{\mathbb{R}^{N}} v_{\delta}^{r}(., s) \xi^{\lambda} d x+\nu \int_{s}^{t} \int_{\mathbb{R}^{N}} v_{\delta}^{r} \Delta\left(\xi^{\lambda}\right) d x
\end{array}
$$

Next we go to the limit as $\delta \rightarrow 0$ : from the Fatou Lemma, $\int_{s}^{t} \int_{\mathbb{R}^{N}} u^{r-1}|\nabla u|^{q} \psi d x d t$ and $(r-$ 1) $\nu \int_{s}^{t} \int_{\mathbb{R}^{N}} u^{r-2}|\nabla u|^{2} \xi^{\lambda} d x$ are finite, and then from the dominated convergence theorem,

$$
\begin{array}{r}
\int_{\mathbb{R}^{N}} u^{r}(., t) \xi^{\lambda} d x+r \int_{s}^{t} \int_{\mathbb{R}^{N}} u^{r-1}|\nabla u|^{q} \xi^{\lambda} d x d t \\
+r(r-1) \nu \int_{s}^{t} \int_{\mathbb{R}^{N}} u^{r-2}|\nabla u|^{2} \xi^{\lambda} d x=\int_{\mathbb{R}^{N}} u^{r}(., \sigma) \xi^{\lambda} d x+\nu \int_{s}^{t} \int_{\mathbb{R}^{N}} u^{r} \Delta\left(\xi^{\lambda}\right) d x .
\end{array}
$$

As $s \rightarrow 0$, from (4.6), we deduce that

$$
\begin{array}{r}
\int_{\mathbb{R}^{N}} u^{r}(., t) \xi^{\lambda} d x+r \int_{0}^{t} \int_{\mathbb{R}^{N}} u^{r-1}|\nabla u|^{q} \xi^{\lambda} d x d t \\
+r(r-1) \int_{0}^{t} \int_{\mathbb{R}^{N}} u^{r-2}|\nabla u|^{2} \xi^{\lambda} d x=\int_{\mathbb{R}^{N}} u_{0}^{r}(., \sigma) \xi^{\lambda} d x+\nu \int_{\sigma}^{t} \int_{\mathbb{R}^{N}} u^{r} \Delta\left(\xi^{\lambda}\right) d x
\end{array}
$$

Now $u(., t) \in L^{r}\left(\mathbb{R}^{N}\right)$ for any $t \in[s, \tau]$, and

$$
\int_{\sigma}^{t} \int_{\mathbb{R}^{N}} u^{r} \Delta\left(\xi^{\lambda}\right) d x \leqq \frac{C}{R^{2}} \tau \int_{\mathbb{R}^{N}} u^{r}(\sigma) d x
$$

thus we can make $R \rightarrow \infty$. Then $\int_{0}^{\tau} \int_{\mathbb{R}^{N}} u^{r-1}|\nabla u|^{q} d x d t$ and $(r-1) \nu \int_{0}^{\tau} \int_{\mathbb{R}^{N}} u^{r-2}|\nabla u|^{2} d x d t$ are finite and, from the dominated convergence theorem,

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} u^{r}(., t) d x+r \int_{0}^{t} \int_{\mathbb{R}^{N}} u^{r-1}|\nabla u|^{q} d x d t+r(r-1) \nu \int_{0}^{t} \int_{\mathbb{R}^{N}} u^{r-2}|\nabla u|^{2} d x=\int_{\mathbb{R}^{N}} u_{0}^{r} d x \tag{4.22}
\end{equation*}
$$

Hence (4.14) and (4.15) follow, implying (4.16).
(iii) Setting $v=u^{b}$ with $b=(q-1+r) / q \leqq r$, there holds $|\nabla v|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right)$, and $v \in L^{\infty}\left((0, T) ; L^{r / b}\left(\mathbb{R}^{N}\right)\right)$. From the Gagliardo-Nirenberg inequality,

$$
\begin{equation*}
\|v(., t)\|_{L^{q}\left(\mathbb{R}^{N}\right)} \leqq C(N, q, r)\|v(., t)\|_{L^{\frac{r}{b}\left(\mathbb{R}^{N}\right)}}^{1-k}\|\nabla v(., t)\|_{L^{q}\left(\mathbb{R}^{N}\right)}^{k}, \quad \frac{1}{k}=1+\frac{r q^{\prime}}{N} \tag{4.23}
\end{equation*}
$$

By integration, for any $0<\tau<T$, we get, from Hölder inequality, with $C=C((\tau, N, q, r)$,

$$
\int_{0}^{\tau} \int_{\mathbb{R}^{N}} v^{q}(., t) d x d t=\int_{0}^{\tau} \int_{\mathbb{R}^{N}} u^{q-1+r}(., t) d x d t \leqq C\|v\|_{L^{\infty}\left((0, \tau) ; L^{\frac{r}{m}}\left(\mathbb{R}^{N}\right)\right.}^{(1-k) q}\left(\int_{0}^{\tau} \int_{\mathbb{R}^{N}}|\nabla v|^{q} d x d t\right)^{k}
$$

Then $u \in L^{q-1+r}\left(Q_{\mathbb{R}^{N}, \tau}\right)$, and $v^{q}=u^{q-1+r} \in L^{1}\left((0, \tau) ; W^{1,1}\left(\mathbb{R}^{N}\right)\right), v \in L^{q}\left((0, \tau) ; W^{1, q}\left(\mathbb{R}^{N}\right)\right)$. If $\nu>0$, we also have $u^{r-2}|\nabla u|^{2}=\left|\nabla\left(u^{r / 2}\right)\right|^{2} \in L^{1}\left(Q_{\mathbb{R}^{N}, \tau}\right)$, and $u^{r / 2} \in L^{2}\left(Q_{\mathbb{R}^{N}, \tau}\right)$, then $u^{r / 2} \in$ $L^{2}\left((0, \tau) ; W^{1,2}\left(\mathbb{R}^{N}\right)\right)$.
(iv) Here $u \in C\left([0, T) ; L_{l o c}^{r}\left(\mathbb{R}^{N}\right)\right)$. We only need to prove that $\lim _{t \rightarrow 0}\left\|u(., t)-u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}=0$. From a diagonal procedure, there exists $t_{n} \rightarrow 0$ such that $\left(u\left(., t_{n}\right)\right)$ converges to $u_{0}$ a.e. in $\mathbb{R}^{N}$. First assume $r>1$; since the convergence holds weakly in $L^{r}\left(\mathbb{R}^{N}\right)$, and $\lim _{n \rightarrow \infty}\left\|u\left(., t_{n}\right)\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}=$ $\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}$ from (4.16). Then it holds from any sequence, and $u \in C\left([0, T) ; L^{r}\left(\mathbb{R}^{N}\right)\right)$. Next assume $r=1$. We have for any $p>0$,

$$
\begin{aligned}
\int_{\mathbb{R}^{N}}\left|u\left(t_{n}\right)-u_{0}\right| d x & \leqq \int_{B_{p}}\left|u\left(t_{n}\right)-u_{0}\right| d x+\int_{\mathbb{R}^{N} \backslash B_{p}}\left|u\left(t_{n}\right)-u_{0}\right| d x \\
& \leqq \int_{B_{p}}\left|u\left(t_{n}\right)-u_{0}\right| d x+\int_{\mathbb{R}^{N} \backslash B_{p}} u\left(t_{n}\right) d x+\int_{\mathbb{R}^{N} \backslash B_{p}} u_{0} d x \\
& =\int_{B_{p}}\left|u\left(t_{n}\right)-u_{0}\right| d x+\int_{\mathbb{R}^{N}} u\left(t_{n}\right) d x-\int_{B_{p}} u_{0} d x \\
& -\int_{B_{p}}\left(u\left(t_{n}\right)-u_{0}\right) d x+\int_{\mathbb{R}^{N} \backslash B_{p}} u_{0} d x \\
& \leqq 2 \int_{B_{p}}\left|u\left(t_{n}\right)-u_{0}\right| d x+\int_{\mathbb{R}^{N}} u\left(t_{n}\right) d x-\int_{\mathbb{R}^{N}} u_{0} d x+2 \int_{\mathbb{R}^{N} \backslash B_{p}} u_{0} d x
\end{aligned}
$$

The result follows from (4.16), because $u_{0} \in L^{1}\left(\mathbb{R}^{N}\right)$.
The decay result is also available for initial data measures, where we do not assume that $q<(N+2) /(N+1):$

Theorem 4.15 Assume $q>1$. Let $u_{0} \in \mathcal{M}_{b}^{+}\left(\mathbb{R}^{N}\right)$ and $u$ be any non-negative weak $\mathcal{M}_{\text {loc }}$ solution of equation (2.5) in $Q_{\mathbb{R}^{N}, T}$. Then $u(., t) \in L^{1}\left(\mathbb{R}^{N}\right)$ for any $t>0$, and

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} u(., t) d x \leqq \int_{\mathbb{R}^{N}} d u_{0} . \tag{4.24}
\end{equation*}
$$

Moreover $u \in C\left((0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right),|\nabla u|^{q} \in L_{\text {loc }}^{1}\left([0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right)$ and

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} u(., t) d x+\int_{0}^{t} \int_{\mathbb{R}^{N}}|\nabla u|^{q} d x d t=\int_{\mathbb{R}^{N}} d u_{0} \tag{4.25}
\end{equation*}
$$

and

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{\mathbb{R}^{N}} u(., t) \varphi d x=\int_{\mathbb{R}^{N}} \varphi d u_{0}, \quad \forall \varphi \in C_{b}\left(\mathbb{R}^{N}\right) \tag{4.26}
\end{equation*}
$$

Proof. If $q^{\prime}<N$, we obtain in the same way (4.17) with $r=1$, and we go to the limit as $s \rightarrow 0$ from (4.7), then

$$
\int_{B_{R_{0}}} u(., t) d x \leqq \int_{\mathbb{R}^{N}} \xi^{\lambda} d u_{0}+C \tau R^{N-q^{\prime}} \leqq \int_{\mathbb{R}^{N}} d u_{0}+C \tau R^{N-q^{\prime}} .
$$

Going to the limit as $R \rightarrow \infty$, and then as $R_{0} \rightarrow \infty$, we deduce (4.24). If $q^{\prime} \geqq N$, we proceed as in the second step of Theorem 4.14, and get again (4.24). Then (4.25) follows. And $u \in$ $C\left((0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right)$, from the dominated convergence theorem, because $u \in C\left((0, T) ; L_{l o c}^{1}\left(\mathbb{R}^{N}\right)\right)$, and $u \in L^{\infty}\left((0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right)$.

Let us show (4.26): let $\varphi \in C_{b}\left(\mathbb{R}^{N}\right)$ be nonnegative, we can assume that $\varphi$ takes its values in $[0,1]$. Let $t_{n} \rightarrow 0$. We know that $\lim _{n \rightarrow \infty} \int_{\mathbb{R}^{N}} u\left(., t_{n}\right) d x=\int_{\mathbb{R}^{N}} d u_{0}$. Let $\psi_{p} \in \mathcal{D}\left(\mathbb{R}^{N}\right)$ with values in $[0,1], \psi_{p}(x)=1$ if $|x| \leqq p, 0$ if $|x| \geqq 2 p$. Then $\lim _{p \rightarrow \infty} \int_{\mathbb{R}^{N}}\left(1-\psi_{p}\right) d u_{0}=0$, from the dominated convergence Theorem. Thus for any $\eta>0$, one can choose $p_{\eta}$ such that $\int_{\mathbb{R}^{N}}\left(1-\psi_{p_{\eta}}\right) d u_{0} \leqq \eta$; and

$$
\begin{aligned}
& \varlimsup\left|\int_{\mathbb{R}^{N}} u\left(., t_{n}\right) \varphi d x-\int_{\mathbb{R}^{N}} \varphi d u_{0}\right| \\
& \leqq \lim _{n \rightarrow \infty}\left|\int_{\mathbb{R}^{N}} u\left(., t_{n}\right) \varphi \psi_{p_{\eta}} d x-\int_{\mathbb{R}^{N}} \varphi \psi_{p_{\eta}} d u_{0}\right|+\int_{\mathbb{R}^{N}} \varphi\left(1-\psi_{p_{\eta}}\right) d u_{0}+\lim _{n \rightarrow \infty} \int_{\mathbb{R}^{N}} u\left(., t_{n}\right) \varphi\left(1-\psi_{p_{\eta}}\right) d x \leqq \eta,
\end{aligned}
$$

hence the conclusion follows.

### 4.3 Regularizing effects

Here we deduce of the decay estimates a regularizing effect without any condition at $\infty$, ending the proof of Theorem 2.2.

Theorem 4.16 Let $q>1$. Let $r \geqq 1$ and $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right)$. Let $u$ be any non-negative weak $L_{\text {loc }}^{r}$ solution of problem (2.5) in $Q_{\mathbb{R}^{N}, T}$ (4.6).

Then $u(., t) \in L^{\infty}\left(\mathbb{R}^{N}\right)$ for any $t \in(0, T)$ and $u$ satisfies the estimates (2.10), where $\sigma_{r, q, N}, \varpi_{r, q, N}$ are given by (2.11).

Moreover if $\nu>0$, then $u$ satifies the estimates (2.12). If $u_{0} \in \mathcal{M}_{b}^{+}\left(\mathbb{R}^{N}\right)$, the same results hold, where $\left\|u_{0}\right\|_{L^{1}\left(\mathbb{R}^{N}\right)}$ is replaced by $\int_{\mathbb{R}^{N}} d u_{0}$.

Proof. Since $u$ is a weak $L_{l o c}^{r}$ solution, then $u \in C\left([0, T) ; L^{r}\left(\mathbb{R}^{N}\right)\right)$, from Theorem 4.14. Thus for any $0 \leqq s<T, u$ is a weak $r$ solution in $Q_{\mathbb{R}^{N}, s, T}$; and $\int_{\mathbb{R}^{N}} u^{r}(s) d x<\infty$ with $r \geqq 1$. For any $0<s \leqq t<T$, and any $\alpha \geqq r-1$ such that $\int_{\mathbb{R}^{N}} u^{\alpha+1}(s) d x<\infty$, we can apply Theorem 4.14 to $u$ starting at point $s$, because of (4.2). Denoting $\beta=1+\alpha / q$, we have

$$
\begin{equation*}
\frac{1}{\alpha+1} \int_{\mathbb{R}^{N}} u^{\alpha+1}(., t) d x+\frac{1}{\beta^{q}} \int_{s}^{t} \int_{\mathbb{R}^{N}}\left|\nabla\left(u^{\beta}\right)\right|^{q} d x d t \leqq \frac{1}{\alpha+1} \int_{\mathbb{R}^{N}} u^{\alpha+1}(., s) d x \tag{4.27}
\end{equation*}
$$

and $u^{\beta}(., t) \in L^{q}\left(\mathbb{R}^{N}\right)$ for almost any $t \in(0, T)$.
(i) Proof of (2.10).

First assume $q<N$. From the Sobolev injection of $W^{1, q}\left(\mathbb{R}^{N}\right)$ into $L^{N q /(N-q)}\left(\mathbb{R}^{N}\right)$, there holds

$$
\left.\frac{1}{\alpha+1} \int_{\mathbb{R}^{N}} u^{\alpha+1}(., t) d x+\frac{C(N, q)}{\beta^{q}} \int_{s}^{t}\left(\int_{\mathbb{R}^{N}} u^{\beta \frac{N q}{N-q}}(., t) d x\right)^{\frac{N-q}{N}}\right) d t \leqq \frac{1}{\alpha+1} \int_{\mathbb{R}^{N}} u^{\alpha+1}(., s) d x
$$

thus Lemma 2.1 applies with $m=q$ and $\theta=N /(N-q)$. We obtain

$$
\|u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq C(t-s)^{-\sigma_{r, q, N}}\|u(., s)\|_{L_{r, q, N}^{r}\left(\mathbb{R}^{N}\right)}^{\varpi_{r}}, \quad C=C(N, q, r)
$$

and deduce (2.10) as $s$ goes to 0 .
If $q=N$, we deduce (2.10) from Lemma 2.1 with $\theta>1$ arbitrary, since $W^{1, N}\left(\mathbb{R}^{N}\right) \subset L^{N \theta}\left(\mathbb{R}^{N}\right)$.
Next assume $q>N$. We straight away obtain, for any $t \in(0, T)$,

$$
\begin{equation*}
\int_{0}^{t} \int_{\mathbb{R}^{N}}\left|\nabla\left(u^{\beta}\right)\right|^{q} d x d t \leqq \frac{1}{r} \int_{\mathbb{R}^{N}} u_{0}^{r} d x \tag{4.28}
\end{equation*}
$$

with $\beta=1+(r-1) / q$. From the Sobolev injection $W^{1, q}\left(\mathbb{R}^{N}\right) \subset L^{\infty}\left(\mathbb{R}^{N}\right), u(., s) \in L^{\infty}\left(\mathbb{R}^{N}\right)$ for almost any $s \in(0, T)$, hence $u(., s) \in L^{\rho}\left(\mathbb{R}^{N}\right)$ ) for any $\rho \geqq r$, and $u \in C\left([s, t), L^{\rho}\left(\mathbb{R}^{N}\right)\right)$ from (4.2). In turn $u(., t) \in L^{\infty}\left(\mathbb{R}^{N}\right)$ for any $t \in(0, T)$ and $t \mapsto\|u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}$ is nonincreasing, thus

$$
r t\|u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}^{q+r-1} \leqq C(N, q) \int_{\mathbb{R}^{N}} u_{0}^{r} d x
$$

This does not give the optimal estimate (2.11). However from Lemma 3.4, v=u $u^{\beta}$ satisfies the Gagliardo-Nirenberg inequality, for almost any $t \in(0, T)$,

$$
\|v(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq C\|v(., t)\|_{L^{\frac{r}{\beta}}\left(\mathbb{R}^{N}\right)}^{1-k}\|\nabla v(., t)\|_{L^{q}\left(\mathbb{R}^{N}\right)}^{k}
$$

where $1 / k=1+(1 / N-1 / q) r / \beta$ and $C=C(N, q, r)$. Then

$$
\|u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}^{\frac{\beta q}{k}} \leqq C\|u(., t)\|_{L^{r}\left(\mathbb{R}^{N}\right)}^{\frac{\beta q(1-k)}{k}} \int_{\mathbb{R}^{N}}\left|\nabla\left(u^{\beta}\right)\right|^{q} d x d t \leqq C\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}^{\frac{\beta q(1-k)}{k}} \int_{\mathbb{R}^{N}}\left|\nabla\left(u^{\beta}\right)\right|^{q} d x d t
$$

By integration, using (4.28), we find

$$
r t\|u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}^{\frac{\beta q}{k}} \leqq C\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}^{\frac{\beta q(1-k)}{k}+r}
$$

which gives precisely (2.10), since $k / \beta q=\sigma_{r, q, N}$ and $(1-k)+k r / \beta q=\varpi_{r, q, N}$.
(ii) Proof of (2.12).

First assume $N>2$. For any $\alpha \geqq r-1$ such that $\int_{\mathbb{R}^{N}} u^{\alpha+1}(s) d x<\infty$,

$$
\frac{1}{\alpha+1} \int_{\mathbb{R}^{N}} u^{\alpha+1}(t) d x+\frac{\alpha}{\tilde{\beta}^{2}} \nu \int_{s}^{t} \int_{\mathbb{R}^{N}}\left|\nabla\left(u^{\tilde{\beta}}\right)\right|^{2} d x d t \leqq \frac{1}{\alpha+1} \int_{\mathbb{R}^{N}} u^{\alpha+1}(s) d x
$$

where $\tilde{\beta}=(\alpha+1) / 2$; and $u^{\tilde{\beta}} \in L_{l o c}^{2}\left((0, \tau) ; W^{1,2}\left(\mathbb{R}^{N}\right)\right)$. From the Sobolev injection of $W^{1,2}\left(\mathbb{R}^{N}\right)$ into $L^{2 N /(N-2)}\left(\mathbb{R}^{N}\right)$, we get

$$
\frac{1}{\alpha+1} \int_{\mathbb{R}^{N}} u^{\alpha+1}(t) d x+\frac{\alpha C(N)}{\tilde{\beta}^{2}} \nu \int_{s}^{t}\left(\int_{\mathbb{R}^{N}} u^{\tilde{\beta} \frac{2 N}{N-2}}\right)^{\frac{N-2}{N}} d x \leqq \frac{1}{\alpha+1} \int_{\mathbb{R}^{N}} u^{\alpha+1}(s) d x
$$

In case $r>1$, Lemma 2.1 applies with $C_{0}=(r-1) C(N) \nu, q=2, \theta=N /(N-2)$ and $\lambda=-1$, $\tilde{\beta}=1+(\alpha-1) / 2$, since $r>N(1-2+1) / 2$; and (2.12) follows.
In case $r=1$, then $u \in C\left([0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right) \cap L_{\text {loc }}^{\infty}\left((0, T) ; L^{\infty}\left(\mathbb{R}^{N}\right)\right)$ because of estimate (2.10). Hence $C\left([0, T) ; L^{\rho}\left(\mathbb{R}^{N}\right)\right)$ for any $\rho>1$, for example with $\rho=2$, and $\|u(., t)\|_{L^{1}\left(\mathbb{R}^{N}\right)}$ is nonincreasing, from Theorem 4.14. Therefore Lemma 2.1 applies on $(\epsilon, t)$ for $0<\epsilon<t<T$ :

$$
\|u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq C(t-\epsilon)^{-\frac{N}{2}}\|u(., \epsilon)\|_{L^{1}\left(\mathbb{R}^{N}\right)} \leqq C(t-\epsilon)^{-\frac{N}{2}}\left\|u_{0}\right\|_{L^{1}\left(\mathbb{R}^{N}\right)}
$$

with $C=C(N, q, r, \nu)$, hence (2.12) follows as $\epsilon \rightarrow 0$.
If $N=2$, we proceed as above to conclude. Next assume $N=1$. In case $r>1$, there holds, for any $t \in(0, T)$,

$$
4(r-1) \nu \int_{0}^{t} \int_{\mathbb{R}^{N}}\left|\nabla\left(u^{\frac{r}{2}}\right)\right|^{2} d x d t \leqq \int_{\mathbb{R}^{N}} u_{0}^{r} d x
$$

and, from Lemma 3.4, applied to $v=u^{r / 2}$, with $m=2=1 / k$,

$$
\|u(., t)\|_{L^{\infty}(\mathbb{R})}^{2 r} \leqq C\|u(., t)\|_{L^{r}(\mathbb{R})}^{r} \int_{\mathbb{R}}\left|\nabla\left(u^{\frac{r}{2}}\right)\right|^{2} d x d t \leqq C\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}^{r} \int_{\mathbb{R}}\left|\nabla\left(u^{\frac{r}{2}}\right)\right|^{2} d x d t
$$

by integration, we get, with a new constant $C=C(r, \nu)$,

$$
t\|u(., t)\|_{L^{\infty}(\mathbb{R})}^{2 r} \leqq C\left\|u_{0}\right\|_{L^{r}(\mathbb{R})}^{2 r}
$$

which proves (2.12). In case $r=1$, taking $\rho=2$ as above, we obtain, for any $0<\epsilon<t<T$,

$$
\|u(., t)\|_{L^{\infty}(\mathbb{R})} \leqq C(\nu)(t-\epsilon)^{-\frac{1}{4}}\|u(., \epsilon)\|_{L^{2}(\mathbb{R})} \leqq C(\nu)(t-\epsilon)^{-\frac{1}{4}}\|u(., \epsilon)\|_{L^{\infty}(\mathbb{R})}^{\frac{1}{2}}\|u(., \epsilon)\|_{L^{1}(\mathbb{R})}^{\frac{1}{2}}
$$

From Lemma 3.1, we deduce

$$
\|u(., t)\|_{L^{\infty}(\mathbb{R})} \leqq C(\nu)(t-\epsilon)^{-\frac{1}{2}}\left\|u_{0}\right\|_{L^{1}(\mathbb{R})}
$$

and we conclude as $\epsilon \rightarrow 0$.
If $u_{0} \in \mathcal{M}_{b}^{+}\left(\mathbb{R}^{N}\right)$, we apply the estimates on $(\epsilon, T)$ and go to the limit as $\epsilon \rightarrow 0$.

Remark 4.17 As a consequence, for any $k \geqq 1$, and for example $q \neq N, N \neq 2$,

$$
\begin{gather*}
\|u(., t)\|_{L^{k r}\left(\mathbb{R}^{N}\right)} \leqq C t^{-\frac{\sigma_{r, q, N}}{k^{\prime}}}\left\|u_{0}\right\|_{\left.L^{r( } \mathbb{R}^{N}\right)}^{\frac{\sigma_{r, q, N}}{k}}+\frac{1}{k}  \tag{4.29}\\
\|u(., t)\|_{L^{k r}\left(\mathbb{R}^{N}\right)} \leqq C t^{-\frac{N}{2 r k^{r}}}\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}, \quad \text { if } \nu>0 . \tag{4.30}
\end{gather*}
$$

Indeed it follows from (4.13) and (2.10), (2.12) by interpolation:

$$
\|u(., t)\|_{L^{k r}\left(\mathbb{R}^{N}\right)} \leqq\|u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}^{1 / k^{\prime}}\|u(., t)\|_{L^{r}\left(\mathbb{R}^{N}\right)}^{1 / k}
$$

Remark 4.18 If $q \leqq 2$, then $u \in C^{2,1}\left(Q_{\mathbb{R}^{N}, T}\right)$, thus we do not need to introduce the regularization by $u_{\varepsilon}$; we only need to introduce $u+\delta$, when $r>1$ and make $\delta \rightarrow 0$.

Remark 4.19 Up to now, the decay estimate (4.13) and the $L^{\infty}$ estimate of $u$ were proved for $u_{0} \in C_{b}\left(\mathbb{R}^{N}\right) \cap L^{r}\left(\mathbb{R}^{N}\right)$, and for the unique bounded solution $u$ of problem (2.5), and based on the estimate (2.7) given in [15, Theorem 5.6]; indeed from the classical inequality

$$
\|u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq C(N, r)\|\nabla u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}^{\frac{N}{N+r}}\|u(., t)\|_{L^{r}\left(\mathbb{R}^{N}\right)}^{\frac{r}{N+r}}
$$

and (2.7), there holds, with $C=C(N, q, r)$,

$$
\begin{gather*}
\|u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq C t^{-\frac{N}{q(N+r)}}\|u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}^{\frac{N}{q(N+r)}}\|u(., t)\|_{L^{r}\left(\mathbb{R}^{N}\right)}^{\frac{r}{N+r}}, \\
\|u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq C t^{-\sigma_{r, q, N}}\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}^{\sigma_{r, q, N}} . \tag{4.31}
\end{gather*}
$$

### 4.4 Further estimates and convergence results for $q \leqq 2$.

Here we consider the case $1<q \leqq 2$. From the $L^{\infty}$ estimates above, and the interior regularity of $u$, we deduce new local estimates and convergence results:

Corollary 4.20 Assume $1<q \leqq 2$.
(i) Any nonnegative weak $L_{\text {loc }}^{r}$ solution (resp. $\mathcal{M}_{\text {loc }}$ solution) $u$ of problem (2.5) with initial data $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right), r \geqq 1$ (resp. $u_{0} \in \mathcal{M}_{b}^{+}\left(\mathbb{R}^{N}\right)$ ) satisfies $u \in C^{2,1}\left(Q_{\mathbb{R}^{N}, T}\right) \cap L_{\text {loc }}^{\infty}\left((0, T) ; C_{b}\left(\mathbb{R}^{N}\right)\right)$.
(ii) Let $\left(u_{0, n}\right)$ be any bounded sequence in $L^{r}\left(\mathbb{R}^{N}\right)$, $r \geqq 1$ (resp. in $\mathcal{M}_{b}^{+}\left(\mathbb{R}^{N}\right)$ ). For any $n \in \mathbb{N}$, let $u_{n}$ be any nonnegative weak $L_{\text {loc }}^{r}$ solution (resp. $\mathcal{M}_{\text {loc }}$ solution) of problem (2.5) with initial data $u_{0, n}$. Then one can extract a subsequence converging in $C_{\text {loc }}^{2,1}\left(Q_{\mathbb{R}^{N}, T}\right)$ to a weak solution $u$ of (1.1) in $Q_{\mathbb{R}^{N}, T}$.

Proof. From [16, Theorem 2.9] there there exists $\gamma \in(0,1)$ such that for any nonnegative weak solution of equation (1.1) $u$ in $Q_{\mathbb{R}^{N}, T}$ and any ball $B_{R} \subset \mathbb{R}^{N}$, and $0<s<\tau<T$,

$$
\|u\|_{C^{2+\gamma, 1+\frac{\gamma}{2}}\left(Q_{B_{R}, s, \tau}\right)} \leqq C \Phi\left(\|u\|_{L^{\infty}\left(Q_{B_{2 R}, \frac{s}{2}, \tau}\right)}\right) .
$$

where $C=C(N, q, R, s, \tau)$ and $\Phi$ is a continuous increasing function. From estimates (2.10), we deduce that $u \in L_{l o c}^{\infty}\left((0, T) ; C_{b}\left(\mathbb{R}^{N}\right)\right)$ and

$$
\begin{equation*}
\|u\|_{C^{2+\gamma, 1+\frac{\gamma}{2}}\left(Q_{B_{R}, s, \tau}\right)} \leqq C \Phi\left(\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}\right), \quad\left(\text { resp. }\|u\|_{C^{2+\gamma, 1+\frac{\gamma}{2}}\left(Q_{B_{R}, s, \tau}\right)} \leqq C \Phi\left(\int_{\mathbb{R}^{N}} d u_{0}\right)\right. \tag{4.32}
\end{equation*}
$$

and the conclusions follow.
We also deduce global gradient estimates in $\mathbb{R}^{N}$ :
Corollary 4.21 Assume $\nu>0,1<q \leqq 2$. (i) Let $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right), r \geqq 1$. Then any weak $L_{\text {loc }}^{r}$ solution $u$ of problem (2.5) satisfies for $q \neq N$

$$
\begin{gather*}
\|\nabla u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq C t^{-\vartheta_{r, q, N}}\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}^{\varkappa_{r, N}}  \tag{4.33}\\
\vartheta_{r, q, N}=\frac{N+r}{r q+N(q-1)}, \quad \varkappa_{r, q, N}=\frac{r}{r q+N(q-1)}
\end{gather*}
$$

and $|\nabla u|^{q} \in L_{l o c}^{\infty}\left((0, T) ; L^{r}\left(\mathbb{R}^{N}\right)\right)$, and

$$
\begin{equation*}
\int_{\mathbb{R}^{N}}|\nabla u(., t)|^{q r} d x \leqq C t^{-r\left(\frac{q}{2}+\sigma_{r, q, N}(q-1)\right)}\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}^{\left(1+\varpi_{r, q, N}(q-1)\right) r} \tag{4.34}
\end{equation*}
$$

where $\sigma_{r, q, N}, \varpi_{r, q, N}$ are defined at (2.11), and $C=C(N, q, r, \nu)$. For $N \neq 2$, then

$$
\begin{gather*}
\|\nabla u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq C t^{-\frac{1}{q}\left(\frac{N}{2 r}+1\right)}\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}^{\frac{1}{q}}  \tag{4.35}\\
\int_{\mathbb{R}^{N}}|\nabla u(., t)|^{q r} d x \leqq C t^{-r\left(\frac{q}{2}+\frac{N}{2 r}(q-1)\right)}\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}^{q r} \tag{4.36}
\end{gather*}
$$

If $N=2$, estimates hold up to an $\varepsilon>0$. Moreover if $q<2$, u is a pointwise mild solution.
(ii) Let $u_{0} \in \mathcal{M}_{b}^{+}\left(\mathbb{R}^{N}\right)$. Then any weak $\mathcal{M}_{\text {loc }}$ solution of (2.5) satisfies the same estimates as in case $r=1$, with $\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}$ replaced by $\int_{\mathbb{R}^{N}} d u_{0}$.

Proof. (i) Let $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right), r \geqq 1$. Then for any $\epsilon>0, u(., \epsilon) \in C_{b}\left(\mathbb{R}^{N}\right)$, from Corollary 4.20. From [29], $u$ is the unique solution $v$ such that $v \in C^{2,1}\left(\mathbb{R}^{N} \times(\epsilon, T)\right) \cap C_{b}\left(\mathbb{R}^{N} \times[\epsilon, T)\right)$, and $v(., \epsilon)=u(., \epsilon)$; since $v \in C_{b}^{2}\left(\mathbb{R}^{N} \times(\epsilon, T)\right)$, we deduce that $u \in C_{b}^{2}\left(\mathbb{R}^{N} \times(0, T)\right)$; and for any $\epsilon \leqq t<T$,

$$
\|u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq\|u(., \epsilon)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}, \quad\|\nabla u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq\|\nabla u(., \epsilon)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}
$$

and from (2.7),

$$
\begin{equation*}
|\nabla u(., t)|^{q} \leqq C(q)(t-\epsilon)^{-1} u(., t), \quad \text { a.e. in } \mathbb{R}^{N} \tag{4.37}
\end{equation*}
$$

From the decay estimates, we also have $\|u(., \epsilon)\|_{L^{r}\left(\mathbb{R}^{N}\right)} \leqq\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}$. And $u(., \epsilon) \in L^{\tilde{r}}\left(\mathbb{R}^{N}\right)$ for any $\tilde{r} \in[r, \infty]$, and $u \in C\left([\epsilon, T) ; L^{\tilde{r}}\left(\mathbb{R}^{N}\right)\right.$ ). Going to the limit in (4.37) as $\epsilon \rightarrow 0$, we deduce (4.33) from (2.10), and (4.35) from (2.12), if $q \neq N$ or $N \neq 2$. Moreover $|\nabla u|^{q} \in L_{l o c}^{\infty}\left((0, T) ; L^{r}\left(\mathbb{R}^{N}\right)\right)$, since

$$
\|\nabla u(., t)\|_{L^{q r}\left(\mathbb{R}^{N}\right)} \leqq C(q) t^{-\frac{1}{q}}\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}^{\frac{1}{q}}
$$

More precisely we get from estimate (2.6),

$$
\| \nabla\left(u^{\frac{1}{q^{\prime}}}(., t)\left\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq C(t-\epsilon)^{-\frac{1}{2}}\right\| u(., \epsilon) \|_{L^{\infty}\left(\mathbb{R}^{N}\right)}^{\frac{1}{q^{\prime}}}\right.
$$

with $C=C(q, \nu)$; then from estimate (2.12), for any $t \in(0, T)$, with other constants $C=C(q, \nu)$,

$$
\begin{gathered}
\| \nabla\left(u^{\frac{1}{q^{\prime}}}(., t)\left\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq C t^{-\frac{1}{2}}\right\| u\left(., \frac{t}{2}\right) \|_{L^{\infty}\left(\mathbb{R}^{N}\right)}^{\frac{1}{q^{\prime}}}\right. \\
|\nabla u(., t)|^{q} \leqq C t^{-\frac{q}{2}}\left\|u\left(., \frac{t}{2}\right)\right\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}^{q-1} u(., t),
\end{gathered}
$$

then from estimate (2.10) we get

$$
\int_{\mathbb{R}^{N}}|\nabla u(., t)|^{q r} d x \leqq C\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}^{\varpi_{r, q, N}(q-1) r} t^{-r\left(\frac{q}{2}+\sigma_{r, q, N}(q-1)\right)} \int_{\mathbb{R}^{N}} u(., t)^{r} d x ;
$$

then (4.34) follows. And (4.36) follows from (2.12). If $N=2$, in particular if $q=N$, the same estimates hold up to an $\varepsilon>0$, from (2.10) and (2.12).

Next we prove that $u$ is a pointwise mild solution as $q<2$. From [29, Theorem 6], $u(., t) \in$ $C_{b}^{2}\left(\mathbb{R}^{N}\right)$ for any $t \in(\epsilon, T)$, in particular $u(., 2 \epsilon) \in C_{b}^{2}\left(\mathbb{R}^{N}\right)$, then for any $t \geqq \epsilon$, and any $x \in \mathbb{R}^{N}$,

$$
\begin{equation*}
u(x, t)=e^{(t-2 \epsilon) \Delta} u(x, 2 \epsilon)-\int_{2 \epsilon}^{t} \int_{\mathbb{R}^{N}} g(x-y, t-s)|\nabla u(y, s)|^{q} d y d s \tag{4.38}
\end{equation*}
$$

see for example [7, Proposition 4.2]. But $u(x, 2 \epsilon)$ converges to $u_{0}$ in $L^{r}\left(\mathbb{R}^{N}\right)$, and then $e^{(t-2 \epsilon) \Delta} u(., \epsilon)$ converges to $e^{t \Delta} u_{0}$ in $L^{r}\left(\mathbb{R}^{N}\right)$. Then we can go to the limit as $\epsilon \rightarrow 0$ in (4.38), for a.e. $x \in \mathbb{R}^{N}$ : the integral is convergent, then the conclusion follows.
(ii) For Theorem 4.15, we have $u(., t) \in L^{1}\left(\mathbb{R}^{N}\right)$ for $t \geqq \epsilon>0$, which gives from (i)

$$
\|u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq C(t-\varepsilon)^{-\sigma_{1, q, N}}\|u(., \epsilon)\|_{L^{1}\left(\mathbb{R}^{N}\right)}^{\varpi_{1,, N}} \leqq C(t-\varepsilon)^{-\sigma_{1, q, N}}\left(\int_{\mathbb{R}^{N}} d u_{0}\right)^{\varpi_{1, q}} .
$$

As $\epsilon \rightarrow 0$, we obtain (4.33), (4.35), (4.34) and (4.36) hold with $r=1$ and $\left\|u_{0}\right\|_{L^{1}\left(\mathbb{R}^{N}\right)}$ replaced by $\int_{\mathbb{R}^{N}} d u_{0}$. And

$$
\|\nabla u(., t)\|_{L^{q}\left(\mathbb{R}^{N}\right)} \leqq C t^{-\frac{1}{q}}\left(\int_{\mathbb{R}^{N}} d u_{0}\right)^{\frac{1}{q}},
$$

thus $|\nabla u|^{q} \in L_{\text {loc }}^{\infty}\left((0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right)$.
Remark 4.22 As a consequence, under the assumptions of Corollary 4.21, there holds $u(., t) \in$ $C_{b}^{1}\left(\mathbb{R}^{N}\right)$, for any $t \in(0, T)$, then $u$ can be extended to a global solution of problem (2.5) on $Q_{\mathbb{R}^{N}, \infty}$, see for example [36].

### 4.5 Existence and uniqueness results for $q \leq 2$

Let $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right), r \geqq 1$. We first consider the "subcritical" case

$$
\begin{equation*}
1<q<\frac{N+2 r}{N+r}, \quad \text { equivalently } \quad q<2 \text { and } r>\frac{N(q-1)}{2-q} . \tag{4.39}
\end{equation*}
$$

Theorem 4.23 Let $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right), r \geqq 1$. Suppose (4.39), and $\nu>0$. Then any weak $L_{\text {loc }}^{r}$ solution u of problem (2.5) satisfies

$$
\begin{equation*}
|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{r}\left(\mathbb{R}^{N}\right)\right) . \tag{4.40}
\end{equation*}
$$

And

$$
u \text { is a weak } L_{\text {loc }}^{r} \text { solution } \Longleftrightarrow u \text { is a mild } L^{r} \text { solution. }
$$

Proof. Let $u$ be any weak $L_{\text {loc }}^{r}$ solution. Then from (4.34),

$$
\int_{0}^{\tau}\|\nabla u(., t)\|_{L^{q r}\left(\mathbb{R}^{N}\right)}^{q} d t=\int_{0}^{\tau}\left(\int_{\mathbb{R}^{N}}|\nabla u(., t)|^{q r} d x\right)^{\frac{1}{r}} d t \leqq C \int_{0}^{\tau} t^{-\left(\frac{q}{2}+\sigma_{r, q, N}(q-1)\right)} d t
$$

with $C=C_{q}\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}, q\right)}^{\left(1+\omega_{r, N}(q-1)\right) r}$, and (4.39) is equivalent to $q / 2+\sigma_{r, q, N}(q-1)<1$. Since $\nu>0$, the estimate (4.36) leads to the same conclusion, because (4.39) is also equivalent to $q / 2+(q-1) N / 2 r<$ 1. Then (4.40) holds. Moreover from Corollary 4.21, $u$ is a mild pointwise solution:

$$
\begin{equation*}
u(., t)=e^{t \Delta} u_{0}(.)-\nu \int_{0}^{t} \int_{\mathbb{R}^{N}} g(x-y, t-s)|\nabla u(y, s)|^{q} d y d s \tag{4.41}
\end{equation*}
$$

Otherwise $u \in C\left([0, T) ; L^{r}\left(\mathbb{R}^{N}\right)\right)$ from Theorem 4.14, and $f=|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{r}\left(\mathbb{R}^{N}\right)\right)$, thus the relation (4.41) holds in $L^{r}\left(\mathbb{R}^{N}\right)$,

$$
\begin{equation*}
u(., t)=\left(e^{t \Delta} u_{0}\right)-\nu \int_{0}^{t} e^{(t-s) \Delta}|\nabla u(., s)|^{q}(s) d s \quad \text { in } L^{r}\left(\mathbb{R}^{N}\right) \tag{4.42}
\end{equation*}
$$

that means $u$ is a mild $L^{r}$ solution. The converse is clear.
Next we deduce the uniqueness results of Theorem 2.3.
Theorem 4.24 Let $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right)$. Assume (4.39) or $q=2$, and $\nu>0$. Then there exists a unique weak $L_{\text {loc }}^{r}$ solution $u$ of problem (2.5). In the first case, $u \in C\left((0, T) ; W^{1, q r}\left(\mathbb{R}^{N}\right)\right)$.

Proof. (i) Case $1<q<(N+2 r) /(N+r)$. From [15, Theorem 2.1], there exists a mild $L^{r}$ solution $u$, and it is unique in the class of mild $L^{r}$ solutions such that $u \in L_{l o c}^{\infty}\left((0, T) ; W^{1, q r}\left(\mathbb{R}^{N}\right)\right)$, see [15, Lemma 2.2 and Remark 2.5]. Then u is a $L_{\text {loc }}^{r}$ solution. Let $v$ be any weak $L_{l o c}^{r}$ solution, thus $u$ is a mild $L^{r}$ solution, from Theorem 4.23. From Theorem 4.14, Corollary 4.20, and Theorem 4.23, $v \in L^{\infty}\left((0, T) ; L^{r}\left(\mathbb{R}^{N}\right)\right) \cap L_{\text {loc }}^{\infty}\left((0, T) ; C_{b}\left(\mathbb{R}^{N}\right)\right)$, and $|\nabla v| \in L_{\text {loc }}^{\infty}\left((0, T) ; L^{q r}\left(\mathbb{R}^{N}\right)\right)$ from Theorem 4.23. Then $v \in L_{l o c}^{\infty}\left((0, T) ; W^{1, q r}\left(\mathbb{R}^{N}\right)\right)$, then $v=u$, and we reach the conclusion. Moreover $u \in C\left((0, T) ; W^{1, q r}\left(\mathbb{R}^{N}\right)\right)$, from [15, Theorem 2.1].
(ii) Case $q=2$. From [15, Theorem 4.2] there exists a unique solution $u$ such that $u \in$ $C\left([0, T) ; L^{r}\left(\mathbb{R}^{N}\right)\right) \cap u \in C^{2,1}\left(\left(Q_{\mathbb{R}^{N}, \infty}\right)\right.$ solution of (1.1) at each point. Then it is a weak $L_{\text {loc }}^{r}$ solution. Reciprocally any weak $L_{l o c}^{r}$ solution $u$ satisfies the conditions above, from Theorem 4.14 and [16].

Theorem 4.25 Assume $1<q<(N+2) /(N+1)$, $\nu>0$. Let $u_{0} \in \mathcal{M}_{b}^{+}\left(\mathbb{R}^{N}\right)$. Then there exists a unique weak $\mathcal{M}_{l o c}$ solution of problem (2.5).

Proof. The existence of a weak semi-group solution was obtained in [12] by approximation. The existence of a mild $\mathcal{M}$ solution was proved in [15, Theorem 2.2], and the two notions are equivalent from Lemma 4.10. In any case the solution is a weak $\mathcal{M}_{\text {loc }}$ solution. Next consider any solution $\mathcal{M}_{l o c}$ solution $u$. Then $u(., t) \in L^{\infty}\left(\mathbb{R}^{N}\right)$ for any $t \in(\epsilon, T)$ by applying Theorem 4.16 on $(\epsilon / 2, T)$. Then again we deduce $u(., \epsilon) \in C_{b}\left(\mathbb{R}^{N}\right)$, and then (4.37) holds. From Theorem 4.14 we still obtain that $u \in L_{l o c}^{\infty}\left((0, T) ; W^{1, q}\left(\mathbb{R}^{N}\right)\right)$. And from the uniquenes on $(\epsilon, T)$, we have $u \in C\left((\epsilon, T) ; W^{1, q}\left(\mathbb{R}^{N}\right)\right)$ from Theorem 4.24. Then $u \in C\left((0, T) ; W^{1, q}\left(\mathbb{R}^{N}\right)\right)$. And $u$ satisfies (4.26), from Theorem 4.15. Then $u$ is a weak semi-group solution, thus a mild $\mathcal{M}$ solution from Lemma 4.10. Therefore $u$ belongs to the class of uniqueness of [15, Theorem 2.2]. We can also prove the uniqueness directly: if $u_{1}, u_{2}$ are two solutions, they are mild $\mathcal{M}$ solutions, thus

$$
\left(u_{1}-u_{2}\right)(., t)=\nu \int_{0}^{t} e^{(t-s) \Delta}\left(\left|\nabla u_{1}(., s)\right|^{q}-\left|\nabla u_{2}(., s)\right|^{q}\right) d s
$$

and we know that $\left|\nabla u_{j}\right|^{q} \in C\left((0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right)$, hence

$$
\begin{aligned}
\left\|\nabla\left(u_{1}-u_{2}\right)(., t)\right\|_{L^{q}\left(\mathbb{R}^{N}\right)} & \leqq \nu \int_{0}^{t}\left\|\nabla\left(e^{(t-s) \Delta}\right)\right\|_{L^{1}\left(\mathbb{R}^{N}\right)}\left\|\left|\nabla u_{1}(., s)\right|^{q}-\left|\nabla u_{2}(., s)\right|^{q}\right\|_{L^{q r}\left(\mathbb{R}^{N}\right)} d s \\
& \leqq C \int_{0}^{t}(t-s)^{-\frac{1}{2}} \max _{j=1,2}\left\|\nabla u_{j}(., s)\right\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}^{q-1}\left\|\nabla\left(u_{1}-u_{2}\right)(., s)\right\|_{L^{q}\left(\mathbb{R}^{N}\right)} d s \\
& \leqq C \int_{0}^{t}(t-s)^{-\frac{1}{2}} s^{-(q-1) \vartheta_{1, q, N}}\left\|\nabla\left(u_{1}-u_{2}\right)(., s)\right\|_{L^{q r}\left(\mathbb{R}^{N}\right)} d s
\end{aligned}
$$

thus we can apply the singular Gronwall Lemma when $(q-1) \vartheta_{1, q, N}<1 / 2$, which means precisely $q<(N+2) /(N+1)$. Then $\nabla\left(u_{1}-u_{2}\right)(., t)=0$ in $L^{q}\left(\mathbb{R}^{N}\right)$, hence $u_{1}=u_{2}$.

Finally we give a short proof of the existence result of [15, Theorem 4.1].
Proposition 4.26 Let $\nu>0,1<q<2$. For any nonnegative $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right), r \geq 1$, there exists a mild pointwise solution $u$ of problem (2.5), and $u \in C\left([0, T) ; L^{r}\left(\mathbb{R}^{N}\right)\right)$.

Proof. Let $u_{0, n}=\min \left(u_{0}, n\right)$. Then $u_{0, n} \in L^{\rho}\left(\mathbb{R}^{N}\right)$ for any $\rho \geq r$. We choose $\rho>N(q-1) /(2-$ $q)$, that means $q<(N+2 \rho) /(N+\rho)$. From [15, Theorem 2.1], there exists a mild $L^{\rho}$ solution $u_{n}$ with initial data $u_{0, n}$, and $u_{n} \in C\left((0, T) ; C_{b}^{2}\left(\mathbb{R}^{N}\right)\right) \cap C^{2,1}\left(Q_{\mathbb{R}^{N}, T}\right)$. The sequence $\left(u_{n}\right)$ is nondecreasing from the comparison principle, and $u_{n}(., t) \leq e^{t \Delta} u_{0} \leq C t^{-N / 2 r}\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}$. From Corollary 4.20, $\left(u_{n}\right)$ converges in $C_{l o c}^{2,1}\left(Q_{\mathbb{R}^{N}, T}\right)$ to a weak solution $u$ of $(1.1)$ in $Q_{\mathbb{R}^{N}, T}$, and $u(., t) \leq e^{t \Delta} u_{0}$. Moreover $\left(\left|\nabla u_{n}\right|^{q}\right)$ is bounded in $L_{l o c}^{1}\left([0, T) ; L_{l o c}^{1}\left(\mathbb{R}^{N}\right)\right)$ : indeed for any $\xi \in \mathcal{D}\left(\mathbb{R}^{N}\right)$, with values in $[0,1]$, and any $0<s<t<T$,

$$
\begin{aligned}
& \int_{\mathbb{R}^{N}} u_{n}(t, .) \xi^{q^{\prime}} d x+\nu \int_{s}^{t} \int_{\mathbb{R}^{N}}\left|\nabla u_{n}\right|^{q} \xi^{q^{\prime}} d x \leqq-q^{\prime} \nu \int_{s}^{t} \int_{\mathbb{R}^{N}} \xi^{\frac{1}{q-1}} \nabla u_{n} . \nabla \xi d x+\int_{\mathbb{R}^{N}} u_{n}(s, .) \xi^{q^{\prime}} d x \\
& \leqq \frac{\nu}{2} \int_{s}^{t} \int_{\mathbb{R}^{N}}\left|\nabla u_{n}\right|^{q} \xi^{q^{\prime}} d x+C t \int_{\mathbb{R}^{N}}|\nabla \xi|^{q^{\prime}} d x+\int_{\mathbb{R}^{N}} u_{n}(s, .) \xi^{q^{\prime}} d x
\end{aligned}
$$

and $u_{n} \in C\left([0, T) ; L^{\rho}\left(\mathbb{R}^{N}\right)\right)$; thus we can go to the limit as $s \rightarrow 0$ :

$$
\int_{\mathbb{R}^{N}} u_{n}(t, .) \xi^{q^{\prime}} d x+\frac{1}{2} \int_{s}^{t} \int_{\mathbb{R}^{N}}\left|\nabla u_{n}\right|^{q} \xi^{q^{\prime}} d x \leqq C t \int_{\mathbb{R}^{N}}|\nabla \xi|^{q^{\prime}} d x+\int_{\mathbb{R}^{N}} u_{0} \xi^{q^{\prime}} d x
$$

Thus $|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L_{l o c}^{1}\left(\mathbb{R}^{N}\right)\right)$, hence, from [16, Proposition 2.15], $u$ admits a trace as $t \rightarrow 0$ : there exists a Radon measure $\mu_{0}$ in $\mathbb{R}^{N}$, such that $u(., t)$ converges weakly* to $\mu_{0}$. Otherwise $e^{t \Delta} u_{0}$ converges to $u_{0}$ in $L^{r}\left(\mathbb{R}^{N}\right)$, thus $\mu_{0} \in L_{l o c}^{1}\left(\mathbb{R}^{N}\right)$ and $0 \leq \mu_{0} \leq u_{0}$; and $u_{n} \leq u$, thus $u_{0, n} \leq \mu_{0}$, hence $\mu_{0}=u_{0}$. Moreover there exists a function $g \in L^{r}\left(\mathbb{R}^{N}\right)$ such that $u(., t) \leqq g$ for small $t$. Then the nonnegative function $e^{t \Delta} u_{0}-u(., t)$ converges weakly* to 0 , and then in $L_{l o c}^{1}\left(\mathbb{R}^{N}\right)$. Hence $u(., t)$ converges to $u_{0}$ in $L_{l o c}^{1}\left(\mathbb{R}^{N}\right)$, then in $L^{r}(\Omega)$ from the dominated convergence theorem. Thus $u \in C\left([0, T) ; L^{r}\left(\mathbb{R}^{N}\right)\right)$. In particular $u$ is a weak $L_{l o c}^{r}$ solution, then a pointwise mild solution, from Corollary 4.21.

Remark 4.27 The uniqueness of the solution is still an open problem when $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right)$ and $q \geqq(N+2 r) /(N+r)$.

### 4.6 More decay estimates for $q<(N+2 r) /(N+r)$

Here, we exploit theorem 4.14 to obtain a better decay estimate of the $L^{r}$ norm when $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right)$ in the subcritical case (4.39), which appears to be new for $r>1$. In case $r=1$ we find again the result of [3], proved under the assumption that the energy relation (4.25) holds.

Theorem 4.28 Let $r \geqq 1$ and assume (4.39), $\nu>0$. Let $u$ be any non-negative weak $r$ solution of problem (2.5) in $Q_{\mathbb{R}^{N}, \infty}$, with initial data $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right)$. Then there exists $C=C(N, q, r)$ such that, for any $t>0$,

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} u^{r}(., t) d x \leq C\left(\int_{\{|x|>\sqrt{t}\}} u_{0}^{r}(x) d x+t^{-\frac{a r-N}{2}}\right), \quad a=\frac{2-q}{q-1} \tag{4.43}
\end{equation*}
$$

As a consequence, $\lim _{t \rightarrow \infty}\|u(t)\|_{L^{r}\left(\mathbb{R}^{N}\right)}=0$ and

$$
r \int_{0}^{\infty} \int_{\mathbb{R}^{N}} u^{r-1}|\nabla u|^{q} d x d t+r(r-1) \nu \int_{0}^{\infty} \int_{\mathbb{R}^{N}} u^{r-2}|\nabla u|^{2} d x d t=\int_{\mathbb{R}^{N}} u_{0}^{r} d x
$$

Proof. We still consider $v=u^{b}$ with $b=(q-1+r) / q<r$, and set $E(s)=\int_{\mathbb{R}^{N}} u^{r}(., s) d x$. Then $E \in W^{1,1}((0, T))$, from the energy relation (4.14), and for almost any $s \in(0, T)$,

$$
E^{\prime}(s)=-r(r-1) \nu \int_{\mathbb{R}^{N}}|\nabla u|^{2} u^{r-2}(., s) d x-\int_{\mathbb{R}^{N}}|\nabla u|^{q} u^{r-1}(., s) d x \leq 0
$$

Next, we set $E=E_{1}+E_{2}$ with

$$
E_{1}(s)=\int_{\{|x|<2 R\}} u^{r}(x, s) d x, \quad E_{2}(s)=\int_{\{|x| \geq 2 R\}} u^{r}(x, s) d x
$$

From the Gagliardo-Nirenberg inequality (4.23), we obtain successively, with $C=C(N, q, r)$,

$$
\begin{aligned}
E_{1}(s) & =\int_{\{|x|<2 R\}} v^{\frac{r}{b}}(x, s) d x \leq\left(\int_{\{|x|<2 R\}} v^{q}(x, s) d x\right)^{\frac{r}{b q}}(2 R)^{1-\frac{r}{b q}} \\
& \leqq C\|\nabla v(s)\|_{L^{q}\left(\mathbb{R}^{N}\right)}^{\frac{k r}{b}}\|v(s)\|_{L^{\frac{(1-k) r}{r / b}}\left(\mathbb{R}^{N}\right)} R^{N\left(1-\frac{r}{b q}\right)} \\
& \leqq \frac{1}{2}\|v(s)\|_{L^{r / b}\left(\mathbb{R}^{N}\right)}^{\frac{r}{b}}+C\|\nabla v(s)\|_{L^{q}\left(\mathbb{R}^{N}\right)}^{\frac{k r}{b}} R^{\frac{N}{k}\left(1-\frac{r}{b q}\right)}
\end{aligned}
$$

thus

$$
\begin{equation*}
E(s) \leqq C\left(\|\nabla v(s)\|_{L^{q}\left(\mathbb{R}^{N}\right)}^{\frac{r}{b}} R^{\frac{N}{k}\left(1-\frac{r}{b q}\right)}+2 E_{2}(s)\right) . \tag{4.44}
\end{equation*}
$$

Let $\eta \in \mathcal{D}\left(\mathbb{R}^{N}\right)$ with values in $[0,1]$, such that $\varphi=1$ in $B_{1}$, with support in $\overline{B_{2}}$, and set $\eta=1-\varphi$, and $\varphi_{l}(x)=\varphi\left(\frac{x}{l}\right), \eta_{R}(x)=\eta\left(\frac{x}{R}\right)$. Observe that our assumption on $q$ implies $q^{\prime}>N / r$. As in the first step of theorem 4.14, we obtain for any $0<\sigma<s<t<T$, and $l>2 R$,

$$
\begin{equation*}
\left(\int_{\mathbb{R}^{N}} u^{r}(., s) \varphi_{l}^{\lambda} \eta_{R}^{\lambda} d x\right)^{\frac{1}{r}} \leq\left(\int_{\mathbb{R}^{N}} u^{r}(., \sigma) \varphi_{l}^{\lambda} \eta_{R}^{\lambda} d x\right)^{\frac{1}{r}}+C(s-\sigma)\left(R^{\frac{N}{r}-q^{\prime}}+l^{\frac{N}{r}-q^{\prime}}\right), \tag{4.45}
\end{equation*}
$$

with $\lambda=r q^{\prime}$, and $C=C(N, q, r, \eta)$. As $\sigma \rightarrow 0$ and $l \rightarrow \infty$. we deduce

$$
\left(\int_{\mathbb{R}^{N}} u^{r}(x, s) \eta_{R} d x\right)^{\frac{1}{r}} \leq\left(\int_{\mathbb{R}^{N}} u_{0}^{r}(x) \eta_{R} d x\right)^{\frac{1}{r}}+C s R^{\frac{N}{r}-q^{\prime}}
$$

Taking $R=\sqrt{t}$, and setting

$$
\rho=r+\frac{N-r q^{\prime}}{2}=\frac{(N+2 r)-q(N+r)}{2(q-1)}=\frac{a r-N}{2},
$$

we find, with a constant $C$ as above,

$$
E_{2}(s) \leq A(t)=C\left(\int_{\{|x|>\sqrt{t}\}} u_{0}^{r}(x) d x+t^{-\rho}\right),
$$

Next, we consider $F(s)=E(s)-2 A(t)$. If there exists $t_{0} \in(0, t)$ such that $F\left(t_{0}\right) \leq 0$, then $F(s) \leq 0, \forall s \in\left(t_{0}, t\right)$; thus $E(t) \leq 2 A(t)$, by continuity, hence (4.43) holds. Next assume that $F(s)>0$, for any $s \in(0, t)$. Since

$$
\begin{equation*}
-F^{\prime}(s) \geq \nu \int_{\mathbb{R}^{N}}|\nabla u|^{q} u^{r-1}(x, s) d x=\nu \int_{\mathbb{R}^{N}}|\nabla v(x, s)|^{q} d x \tag{4.46}
\end{equation*}
$$

we find $F(s) \leqq C\left(-F^{\prime}(s)\right)^{r / b q} t^{(1-r / b q) N / 2 k}$ from (4.44). By integration we get

$$
C(t-s) t^{-\frac{N}{2 k}\left(1-\frac{r}{b q}\right)} \leqq F(t)^{-\frac{q-1}{r}}-F(s)^{\frac{q-1}{r}} .
$$

As $s \longrightarrow 0$ we deduce that $F(t) \leqq C t^{-\rho}$, since $\rho=r /(q-1)-N / 2 k$, and (4.43) still holds.

Remark 4.29 The case $r=1$ has been the object of many works, assuming that $u_{0} \in L^{1}\left(\mathbb{R}^{N}\right) \cap$ $W^{1, \infty}\left(\mathbb{R}^{N}\right)$. There holds

$$
\lim _{t \rightarrow \infty}\|u(t)\|_{L^{1}\left(\mathbb{R}^{N}\right)}=0 \Longleftrightarrow q \leq(N+2) /(N+1)
$$

see [2], [12], [4], [28]. When $q<(N+2) /(N+1)$, the absorption plays a role in the asymptotics. From [10], if $\lim _{|x| \rightarrow \infty}|x|^{a} u_{0}(x)=0$, where $a=(2-q) /(q-1)$, then $u(., t)$ converges as $t \rightarrow \infty$ to the very singular solution constructed in [34], [13]; then $\int_{\mathbb{R}^{N}} u(., t) d x$ behaves like $t^{-(a-N) / 2}$ for large $t$, and estimate (4.43) is sharp. When $q>(N+2) /(N+1)$, and $u_{0} \in L^{1}\left(\mathbb{R}^{N}\right)$, then $u(., t)$ behaves as the fundamental solution of heat equation, see [10].

Our result is new when $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right), r>1$. When $q>(N+2) /(N+1)$, and $u_{0}$ is bounded and behaves like $|x|^{-b}$ as $|x| \rightarrow \infty$ with $b \in(a, N)$, it has been shown that $u(., t)$ behaves as the selfsimilar solution of the heat equation with initial data $|x|^{-b}$, see [17]. In that case $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right)$ for any $r>N / b$ and $\int_{\mathbb{R}^{N}} u^{r}(., t) d x$ behaves like $t^{-(b r-N) / 2}$. Thus (4.43) is sharp as $b \rightarrow a$.

## 5 The Dirichlet problem in $Q_{\Omega, T}$

Here we study equation (1.1) in case of a regular bounded domain $\Omega$, with Dirichlet conditions on $\partial \Omega \times(0, T)$, with $\nu>0$; by homothety we can assume $\nu=1$ :

$$
\left(D_{\Omega, T}\right)\left\{\begin{array}{l}
u_{t}-\Delta u+|\nabla u|^{q}=0, \quad \text { in } Q_{\Omega, T},  \tag{5.1}\\
u=0 \quad \text { on } \partial \Omega \times(0, T),
\end{array}\right.
$$

As in section 4, we study the problem with rough initial data, and introduce different notions of solutions.

### 5.1 Solutions of the heat equation with $L^{1}$ data

The regularization method used at Section 4 does not provide estimates up to the boundary. In this section we use another argument: the notion of entropy solution, introduced in [35], for the problem

$$
\left\{\begin{array}{l}
u_{t}-\Delta u=f, \quad \text { in } Q_{\Omega, s, \tau}  \tag{5.2}\\
u=0 \quad \text { on } \partial \Omega \times(s, \tau) \\
u(., s)=u_{s} \geqq 0
\end{array}\right.
$$

when $f$ and $u_{s}$ are integrable, that we recall now. For any $k>0$ and $\theta \in \mathbb{R}$, we define as usual the truncation function $T_{k}$ and a primitive $\Theta_{k}$ by

$$
\begin{equation*}
T_{k}(\theta)=\max (-k, \min (k, \theta)), \quad \Theta_{k}(s)=\int_{0}^{r} T_{k}(\theta) d \theta \tag{5.3}
\end{equation*}
$$

Definition 5.1 Let $s, \tau \in \mathbb{R}$ with $s<\tau$, and $f \in L^{1}\left(Q_{\Omega, s, \tau}\right)$ and $u_{s} \in L^{1}(\Omega)$. A function $u \in$ $C\left([s, \tau] ; L^{1}(\Omega)\right)$ is an entropy solution of the problem (5.2) if $T_{k}(u) \in L^{2}\left((s, \tau) ; W_{0}^{1,2}(\Omega)\right)$ for any $k>0$, and

$$
\begin{gather*}
\int_{\Omega} \Theta_{k}(u-\varphi)(., \tau) d x-\int_{\Omega} \Theta_{k}\left(u_{s}-\varphi(., s) d x+\int_{s}^{\tau}\left\langle\varphi_{t}, T_{k}(u-\varphi)\right\rangle d t\right.  \tag{5.4}\\
+\int_{s}^{\tau} \int_{\Omega}\left(\nabla u . \nabla T_{k}(u-\varphi)-f T_{k}(u-\varphi) d x d t \leq 0\right.
\end{gather*}
$$

for any $\varphi \in L^{2}\left((s, \tau) ; W^{1,2}(\Omega)\right) \cap L^{\infty}\left(Q_{\Omega, \tau}\right)$ such that $\varphi_{t} \in L^{2}\left((s, \tau) ; W^{-1,2}(\Omega)\right)$.

Other notions of solutions have been used for this problem, see [8], recalled below. In fact they are equivalent: here $e^{t \Delta}$ denotes the semi-group of the heat equation with Dirichlet conditions acting on $L^{1}(\Omega)$,

Lemma 5.2 Let $-\infty<s<\tau<\infty, f \in L^{1}\left(Q_{\Omega, s, \tau}\right)$, $u_{s} \in L^{1}(\Omega)$ and $u \in C\left([s, \tau] ; L^{1}(\Omega)\right)$, $u(., s)=u_{s}$. Then the three properties are equivalent:
(i) $u \in L^{1}\left((s, \tau) ; W_{0}^{1,1}(\Omega)\right)$, such that

$$
\begin{equation*}
u_{t}-\Delta u=f, \quad \text { in } \mathcal{D}^{\prime}\left(Q_{\Omega, s, \tau}\right) \tag{5.5}
\end{equation*}
$$

(ii) $u$ is a mild solution of (5.2), that means, for any $t \in[s, \tau]$,

$$
\begin{equation*}
u(., t)=e^{(t-s) \Delta} u_{s}+\int_{s}^{t} e^{(t-\sigma) \Delta} f(\sigma) d \sigma \quad \text { in } L^{1}(\Omega) \tag{5.6}
\end{equation*}
$$

(iii) $u$ is an entropy solution of (5.2).

Such a solution exists, is unique, and will be called weak solution of (5.2).
Proof. It follows from the existence and uniqueness of the solutions of (i) from [5, Lemma 3.4], as noticed in [8], and of the entropy solutions, see [18].

As a consequence, when $u$ is bounded, we can admit test functions of the form $u^{\alpha}$ :
Lemma 5.3 Let $s, \tau \in \mathbb{R}$ with $s<\tau$, and $f \in L^{1}\left(Q_{\Omega, s, \tau}\right)$ and $u$ be any nonnegative bounded weak solution in $Q_{\Omega, s, \tau}$ of (5.2).

Then, for any $\alpha>0$, there holds $u^{\alpha-1}|\nabla u|^{2} \in L^{1}\left(Q_{\Omega, s, \tau}\right)$ and

$$
\begin{equation*}
\left.\left.\frac{1}{\alpha+1} \int_{\Omega} u^{\alpha+1}(., \tau)\right) d x+\alpha \iint_{Q_{\Omega, s, \tau}} u^{\alpha-1}|\nabla u|^{2} d x d t=\frac{1}{\alpha+1} \int_{\Omega} u^{\alpha+1}(., s)\right) d x+\int_{s}^{\tau} \int_{\Omega} f u^{\alpha} d x d t \tag{5.7}
\end{equation*}
$$

Proof. We have $u \in L^{2}\left((s, \tau) ; W_{0}^{1,2}(\Omega)\right) \cap L^{\infty}\left(Q_{\Omega, s, \tau}\right)$, and $u_{t} \in L^{2}\left((s, \tau) ; W^{-1,2}(\Omega)\right)+$ $L^{1}\left(Q_{\Omega, s, \tau}\right)$. Then any function $\varphi \in L^{2}\left((s, \tau) ; W_{0}^{1,2}(\Omega)\right) \cap L^{\infty}\left(Q_{\Omega, s, \tau}\right)$ is admissible in equation (5.5). In particular for any $\alpha>0$, we can take $\varphi=M_{\alpha, \delta}(u)=(u+\delta)^{\alpha}-\delta^{\alpha}$, with $\delta>0$. Integrating on $[s, \tau]$ we deduce that

$$
\int_{s}^{\tau}<u_{t}, \varphi>+\alpha \iint_{Q_{\Omega, s, \tau}}(u+\delta)^{\alpha-1}|\nabla u|^{2} d x d t=\int_{s}^{\tau} \int_{\Omega} f M_{\alpha, \delta}(u) d x d t
$$

Let $k>0$ such that $\sup _{Q_{\Omega, s, \tau}} u \leqq k$, thus $u=T_{k}(u)$. The function $\theta \mapsto M(\theta)=\left(T_{k}(\theta)+\delta\right)^{\alpha}-\delta^{\alpha}$ is continuous on $\mathbb{R}^{+}$and piecewise $C^{1}$ such that $M(0)=0$ and $M^{\prime}$ has a compact support. Denoting $\mathcal{M}_{\alpha, \delta}(r)=(u+\delta)^{\alpha+1} /(\alpha+1)-\delta^{\alpha} u$, we can integrate by parts from [27, Lemma 7.1], and deduce that
$\left.\left.\int_{\Omega} \mathcal{M}_{\alpha, \delta}(u)(., \tau)\right) d x-\int_{\Omega} \mathcal{M}_{\alpha, \delta}(u)(., s)\right) d x+\alpha \iint_{Q_{\Omega, s, \tau}}(u+\delta)^{\alpha-1}|\nabla u|^{2} d x d t=\int_{s}^{\tau} \int_{\Omega} f M_{\alpha, \delta}(u) d x d t ;$
We can go to the limit as $\delta \rightarrow 0$ from the Fatou Lemma, and then from the dominated convergence theorem. Thus (5.7) holds for $\alpha>0$.

Remark 5.4 From [27], the notion of entropy solution of (5.2) is also equivalent to the notion of renormalized solution, that we develop in Section 6. Lemma 5.3 is a special case of a much more general property of the truncates when $u$ is not necessarily bounded, see Lemma 6.3.

### 5.2 Different notions of solutions of problem ( $D_{\Omega, T}$ )

Definition 5.5 We say that $u$ is a weak solution of the problem $\left(D_{\Omega, T}\right)$ if $u \in C\left((0, T) ; L^{1}(\Omega)\right) \cap$ $L_{l o c}^{1}\left((0, T) ; W_{0}^{1,1}(\Omega)\right)$, such that $|\nabla u|^{q} \in L_{l o c}^{1}\left((0, T) ; L^{1}(\Omega)\right)$ and u satisfies

$$
\begin{equation*}
u_{t}-\Delta u+|\nabla u|^{q}=0, \quad \text { in } \mathcal{D}^{\prime}\left(Q_{\Omega, T}\right) \tag{5.8}
\end{equation*}
$$

Next we study the Cauchy-Dirichlet problem

$$
\left\{\begin{array}{l}
u_{t}-\Delta u+|\nabla u|^{q}=0, \quad \text { in } Q_{\Omega, T}  \tag{5.9}\\
u=0 \quad \text { on } \partial \Omega \times(0, T) \\
u(x, 0)=u_{0} \geqq 0
\end{array}\right.
$$

with $u_{0} \in L^{r}(\Omega), r \geqq 1$, or only $u_{0} \in \mathcal{M}_{b}^{+}(\Omega)$. Here in any case $u_{0} \in \mathcal{M}_{b}^{+}(\Omega)$.
Definition 5.6 If $u_{0} \in L^{r}(\Omega), r \geqq 1$, we say that $u$ is a weak $L^{r}$ solution of problem (5.9) if it is a weak solution of $\left(D_{\Omega, T}\right)$, such that the extension of $u$ by $u_{0}$ at $t=0$ satisfies $u \in C\left([0, T) ; L^{r}(\Omega)\right)$.

Definition 5.7 For any $u_{0} \in \mathcal{M}_{b}^{+}(\Omega)$, we say that $u$ is a weak $\mathcal{M}$ solution of problem (5.9) if it is a weak solution of $\left(D_{\Omega, T}\right)$, such that

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{\Omega} u(., t) \psi d x=\int_{\Omega} \psi d u_{0}, \quad \forall \psi \in C_{b}(\Omega) \tag{5.10}
\end{equation*}
$$

Semi-group type solutions have been introduced in [8], see also [1]. For $u_{0} \in \mathcal{M}_{b}^{+}(\Omega)$, we set $e^{t \Delta} u_{0}=\int_{\Omega} g_{\Omega}(., y, t) d u_{0}(y)$, where $g_{\Omega}$ is the heat kernel with Dirichlet conditions on $\partial \Omega$.

Definition 5.8 For any $u_{0} \in \mathcal{M}_{b}^{+}(\Omega)$, a function $u$ is a mild solution of problem (5.9) if $u \in C\left((0, T) ; L^{1}(\Omega)\right)$, and $|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{1}(\Omega)\right)$ and

$$
\begin{equation*}
u(., t)=e^{t \Delta} u_{0}(.)-\int_{0}^{t} e^{(t-s) \Delta}|\nabla u(., s)|^{q} d s \quad \text { in } L^{1}(\Omega) \tag{5.11}
\end{equation*}
$$

Remark 5.9 As it was shown in [8, p.1420], from Lemma 5.2,
$u$ is a mild solution $\Longleftrightarrow u$ is a weak $\mathcal{M}$ solution such that $|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{1}(\Omega)\right)$;
and then $u \in L_{l o c}^{1}\left([0, T) ; W_{0}^{1,1}(\Omega)\right)$.
Remark 5.10 As in Remark 4.11, the definition of mild solution requires an integrability property of the gradient up to $t=0$, namely $|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{1}(\Omega)\right)$. The definition of weak solution only assumes that $|\nabla u|^{q} \in L_{l o c}^{1}\left((0, T) ; L^{1}(\Omega)\right)$.

### 5.3 Decay and regularizing effect

Here $\Omega$ is bounded, then the situation is simpler than in $\mathbb{R}^{N}$ : indeed we take benefit of the regularizing effect of the semi-group $e^{t \Delta}$ associated with the first eigenvalue $\lambda_{1}$ of the Laplacian, and also of the inclusion $L^{r}(\Omega) \subset L^{1}(\Omega)$.

Lemma 5.11 Let $q>1$, and $u_{0} \in L^{r}(\Omega), r \geqq 1$. 1) Let $u$ be any non-negative weak $L^{r}$ solution of problem (5.9).
(i) Then $u(., t) \in L^{\infty}(\Omega)$ for any $t \in(0, T)$, and

$$
\begin{equation*}
\|u(., t)\|_{L^{r}(\Omega)} \leqq C e^{-\lambda_{1} t}\left\|u_{0}\right\|_{L^{r}(\Omega)}, \quad\|u(., t)\|_{L^{\infty}(\Omega)} \leqq C t^{-\frac{N}{2 r}} e^{-\lambda_{1} t}\left\|u_{0}\right\|_{L^{r}(\Omega)} \tag{5.12}
\end{equation*}
$$

(ii) Moreover $|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{1}(\Omega)\right)$, and

$$
\begin{equation*}
\int_{\Omega} u(., t) d x+\int_{0}^{t} \int_{\Omega}|\nabla u|^{q} d x d t \leqq \int_{\Omega} u_{0} d x \tag{5.13}
\end{equation*}
$$

If $r>1$, then $u^{r-1}|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{1}(\Omega)\right)$ and $u^{r-2}|\nabla u|^{2} \in L_{l o c}^{1}\left([0, T) ; L^{1}(\Omega)\right)$, and

$$
\begin{equation*}
\frac{1}{r} \int_{\Omega} u^{r}(., t) d x+\int_{0}^{t} \int_{\Omega} u^{r-1}|\nabla u|^{q} d x d t+(r-1) \int_{0}^{t} \int_{\Omega} u^{r-2}|\nabla u|^{2} d x d t=\frac{1}{r} \int_{\Omega} u_{0}^{r} d x, \tag{5.14}
\end{equation*}
$$

As a consequence, $u^{q-1+r} \in L_{\text {loc }}^{1}\left(\left([0, T) ; W_{0}^{1,1}(\Omega)\right)\right.$.
2) Let $u_{0} \in \mathcal{M}_{b}^{+}(\Omega)$ and $u$ be any non-negative weak $\mathcal{M}$ solution of problem (5.9). Then (5.12) and (5.13) still hold as in case $u_{0} \in L^{1}(\Omega)$, where the norm $\left\|u_{0}\right\|_{L^{1}(\Omega)}$ is replaced by $\int_{\Omega} d u_{0}$. In particular $u$ is a mild solution.

Proof. 1) (i) Let $0<\epsilon<\tau<T$. Since $u$ is a weak solution of ( $D_{\Omega, T}$ ), we can apply Lemma 5.2 with $f=-|\nabla u|^{q}$ in $Q_{\Omega, \epsilon, \tau}$. Thus $u$ is a mild solution of the problem in $Q_{\Omega, \epsilon, \tau}$ : for any $t \in[\epsilon, \tau]$,

$$
u(., t)=e^{(t-\epsilon) \Delta} u(., \epsilon)-\int_{\epsilon}^{t} e^{(t-\sigma) \Delta}|\nabla u|^{q} d \sigma \quad \text { in } L^{1}(\Omega)
$$

therefore $u(., t) \leqq e^{(t-\epsilon) \Delta} u(., \epsilon)$. From our assumptions $u \in C\left([0, T) ; L^{r}(\Omega)\right)$, we deduce that $u(., t) \leqq e^{t \Delta} u_{0}$ as $\epsilon \rightarrow 0$. Then (5.12) follows from the properties of the semi-group $e^{t \Delta}$.
(ii) The function $u$ is bounded in $Q_{\Omega, s, \tau}$, thus from Lemma 5.3, for any $\rho>1$,

$$
\begin{equation*}
\frac{1}{\rho} \int_{\Omega} u^{\rho}(., t) d x+\int_{\epsilon}^{t} \int_{\Omega} u^{\rho-1}|\nabla u|^{q} d x d t+(\rho-1) \int_{\epsilon}^{t} \int_{\Omega} u^{\rho-2}|\nabla u|^{2} d x d t=\frac{1}{\rho} \int_{\Omega} u^{\rho}(., \epsilon) d x \tag{5.15}
\end{equation*}
$$

As $\rho \rightarrow 1$, we deduce that $|\nabla u|^{q} \in L^{1}\left(Q_{\Omega, \epsilon, \tau}\right)$ from the Fatou Lemma, and

$$
\int_{\Omega} u(., t) d x+\int_{\epsilon}^{t} \int_{\Omega}|\nabla u|^{q} d x d t \leqq \int_{\Omega} u(., \epsilon) d x
$$

As $\epsilon \rightarrow 0$ we deduce that $|\nabla u|^{q} \in L^{1}\left(Q_{\Omega, \tau}\right)$ and (5.13) holds. If $r>1$, we can take $\rho=r$ in (5.15) and obtain (5.14) as $\epsilon \rightarrow 0$. Then $u^{q-1+r} \in L_{l o c}^{1}\left(\left([0, T) ; W_{0}^{1,1}(\Omega)\right)\right.$ as in the case of $\mathbb{R}^{N}$.
2) The same estimates hold because $\lim _{\epsilon \rightarrow 0}\|u(., \epsilon)\|_{L^{1}(\Omega)}=\int_{\Omega} d u_{0}$.

Theorem 5.12 Let $q>1$ and $u_{0} \in L^{r}(\Omega), r \geqq 1$. 1) Let $u$ be any non-negative weak $L^{r}$ solution of problem (5.9). Then

$$
\|u(., t)\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leqq\left\{\begin{array}{ccc}
C t^{-\sigma_{r, q, N}}\left\|u_{0}\right\|_{L_{r, q, N}}^{\omega_{r}(\Omega)}, & C=C(N, q, r), & \text { if } q \neq N,  \tag{5.16}\\
C_{\varepsilon} t^{-(1+\varepsilon) \sigma_{r, N, N}}\left\|u_{0}\right\|_{L^{r}(\Omega)}^{1(\Omega)},
\end{array}, \quad \forall \varepsilon>0, \quad C_{\varepsilon}=C(N, q, r, \varepsilon), \quad \text { if } q=N,\right.
$$

where $\sigma_{r, q, N}, \varpi_{r, q, N}$ are given at (2.11).
2) Any non-negative weak solution u of $\left(D_{\Omega, T}\right)$ satisfies the universal estimate, where $C=C(N, q,|\Omega|)$,

$$
\begin{equation*}
\|u(., t)\|_{L^{\infty}(\Omega)} \leqq C t^{-\frac{1}{q-1}} \tag{5.17}
\end{equation*}
$$

Proof. 1) First assume $q<N$. For any $\alpha>0$, setting $\rho=1+\alpha$, and $0<\epsilon \leqq s<t<T$, setting $\beta=1+\alpha / q$, we obtain, from (5.15),

$$
\frac{1}{\alpha+1} \int_{\Omega} u^{\alpha+1}(., t) d x+\frac{1}{\beta^{q}} \int_{s}^{t} \int_{\Omega}\left|\nabla\left(u^{\beta}\right)\right|^{q} d x d t \leqq \frac{1}{\alpha+1} \int_{\Omega} u^{\alpha+1}(., s) d x
$$

Then $u^{\beta}(., t) \in W_{0}^{1, q}(\Omega)$, since $\left.u(., t) \in L^{\infty}\left(Q_{\Omega, s, \tau}\right) \cap W_{0}^{1,1}(\Omega)\right)$. From the Sobolev injection of $W_{0}^{1, q}(\Omega)$ into $L^{N q /(N-q)}(\Omega)$,

$$
\frac{1}{\alpha+1} \int_{\Omega} u^{\alpha+1}(., t) d x+\frac{C(N, q)}{\beta^{q}} \int_{s}^{t}\left(\int_{\Omega} u^{\beta \frac{N q}{N-q}}(., \sigma) d x\right)^{\frac{N-q}{N}} d t \leqq \frac{1}{\alpha+1} \int_{\Omega} u^{\alpha+1}(., s) d x
$$

From Lemma 2.1 on $[\epsilon, T)$ with $m=q$ and $\theta=N /(N-q)$, we obtain estimates for $\epsilon<t<T$ :

$$
\|u(., t)\|_{L^{\infty}(\Omega)} \leqq C(t-\epsilon)^{-\sigma_{r, q}, N}\|u(., \epsilon)\|_{L^{r}(\Omega)}^{\varpi_{r, q, N}}, \quad\|u(., t)\|_{L^{\infty}(\Omega)} \leqq C(t-\epsilon)^{-\frac{1}{q-1}}
$$

and we deduce (5.16) and (5.17) as $\epsilon \rightarrow 0$. In the case $q=N$ the same conclusion follows from Lemma 2.1 with any $\theta>1$. If $q>N$ we proceed as in Theorem 4.16 by applying Lemma 3.4.
2) Let $u$ be any weak solution of $\left(D_{\Omega, T}\right)$. Since $u \in C\left([\epsilon, T) ; L^{1}(\Omega)\right)$ for $\epsilon>0$, we find, for any $t \in[\epsilon, T)$,

$$
\|u(., t)\|_{L^{\infty}(\Omega)} \leqq C(t-\epsilon)^{-\frac{1}{q-1}}
$$

with $C=C(N, q)$, and deduce (5.17) for any $t \in(0, T)$ as $\epsilon \rightarrow 0$.
Remark 5.13 In particular we find again estimate (5.17) obtained in [33] in case $q<2$, for solutions $u$ such that $u \in C\left((0, T) ; L^{2}(\Omega)\right) \cap L^{2}\left((0, T) ; W_{0}^{1,2}(\Omega)\right)$, and $(u-k)^{+}$is admissible as a test function in the equation; those conditions imply integrability properties of $u|\nabla u|^{q}$. Our result is valid without any of these conditions.

### 5.4 Existence and uniqueness results for $q \leqq 2$

From estimate (5.17), we deduce new convergence results when $q \leqq 2$ :
Corollary 5.14 Assume $1<q \leqq 2$. Then
(i) any weak solution $u$ of problem $\left(D_{\Omega, T}\right)$ satisfies $u \in C^{2,1}\left(Q_{\Omega, T}\right) \cap C^{1,0}(\bar{\Omega} \times(0, T))$;
(ii) for any sequence of weak solutions $\left(u_{n}\right)$ of $\left(D_{\Omega, T}\right)$, one can extract a subsequence converging in $C^{2,1}\left(Q_{\Omega, T}\right) \cap C^{1,0}(\bar{\Omega} \times(0, T))$ to a weak solution $u$ of $\left(D_{\Omega, T}\right)$.

Proof. (i) From [16, Theorem 2.9], any weak solution $u$ of $\left(D_{\Omega, T}\right)$ such that $u \in L_{l o c}^{\infty}\left((0, T) ; L^{\infty}(\Omega)\right)$ satisfies $u \in C^{2,1}\left(Q_{\Omega, T}\right) \cap C^{1,0}(\bar{\Omega} \times(0, T))$. And we obtain precisely $u \in L_{\text {loc }}^{\infty}\left((0, T) ; L^{\infty}(\Omega)\right)$, at Theorem 5.12,3.
(ii) Moreover $\left(u_{n}\right)$ is uniformly bounded in $\left.L_{l o c}^{\infty}(0, T) ; L^{\infty}(\Omega)\right)$. From [16], there exists $v \in(0,1)$ such that, for any $0<s<\tau<T$,

$$
\begin{equation*}
\left\|u_{n}\right\|_{C(\bar{\Omega} \times[s, \tau])}+\left\|\nabla u_{n}\right\|_{C^{v, v / 2}(\bar{\Omega} \times[s, \tau])} \leqq C \Phi\left(\left\|u_{n}\right\|_{L^{\infty}\left(Q_{\Omega, s / 2, \tau}\right)}\right) \tag{5.18}
\end{equation*}
$$

where $C=C((N, q, \Omega, s, \tau, v)$, and $\Phi$ is an increasing function. The conclusion follows.
Theorem 5.15 Suppose $1<q<(N+2) /(N+1)$. For any $u_{0} \in \mathcal{M}_{b}^{+}(\Omega)$, problem (5.9) admits a unique weak $\mathcal{M}$ solution.

Proof. From [8, Theorem 3.2], [1], for any (possibly signed) $u_{0} \in \mathcal{M}_{b}(\Omega)$, problem (5.9) has a unique mild $\mathcal{M}$ solution, and it is nonnegative when $u_{0} \in \mathcal{M}_{b}^{+}(\Omega)$. From Lemma 5.11, any weak $\mathcal{M}$ solution is a mild $\mathcal{M}$ solution, thus uniqueness holds in this class.

Next assume that $u_{0} \in L^{r}(\Omega)$ and consider the subcritical case (4.39). In [8, Theorem 3.3], it is proved that there exists a weak $L^{r}$ solution such that $u \in L_{l o c}^{q}\left([0, T) ; W_{0}^{1, q r}(\Omega)\right)$, and it is unique in this space. The local existence and uniqueness in an interval $\left(0, T_{1}\right)$ is obtained by the Banach fixed point theorem in a ball of radius $K_{1}$ of the space

$$
X_{K_{1}}\left(T_{1}\right)=\left\{u \in C\left(\left(0, T_{1}\right], W_{0}^{1, q r}(\Omega)\right): \sup _{\left(0, t_{1}\right]} t^{\theta}\left(\|u(., t)\|_{L^{q r}(\Omega)}+t^{\frac{1}{2}}\|\nabla u(., t)\|_{L^{q r}(\Omega)}\right)<\infty\right\}
$$

where $\theta=N / 2 r q^{\prime}$, under the condition

$$
\begin{equation*}
\left\|u_{0}\right\|_{L^{r}(\Omega)}+K_{1}^{q} T_{1}^{\gamma} \leqq C K_{1}, \quad \text { where } \gamma=1-q(\theta+1 / 2) \quad \text { and } C=C(N, q, r, \Omega) \tag{5.19}
\end{equation*}
$$

We prove the uniqueness with no condition of integrability:
Theorem 5.16 Assume that $u_{0} \in L^{r}(\Omega)$ and $1<q<(N+2 r) /(N+r)$. Then problem (5.9) admits a unique weak $L^{r}$ solution.

Proof. Let $\epsilon>0$. From Theorem 5.12, $u$ is bounded on $(\epsilon, T)$ for any $\epsilon \in(0, T)$. Then $u \in C^{2,1}\left(Q_{\Omega, T}\right) \cap C^{1,0}(\bar{\Omega} \times(0, T))$ because $q<2$, from [16, Theorem 2.10]. From (2.14), there exists a function $D \in C((0, \infty)$ such that for any $\epsilon>0$ and for $t \geqq \epsilon$

$$
\|\nabla u(., t)\|_{L^{\infty}(\Omega)} \leqq D(t-\epsilon)
$$

Then $|\nabla u|$ is bounded in $Q_{\epsilon, T, \Omega}$ for any $\epsilon>0$. Thus $u \in C\left((0, T), W_{0}^{1, q r}(\Omega)\right)$. The problem with initial data $u(., \epsilon)$ at time 0 has a unique solution $v_{\epsilon}$ such that $v_{\epsilon} \in C\left((0, T-\epsilon), W_{0}^{1, q r}(\Omega)\right)$, then $v_{\epsilon}(., t)=u(., t+\epsilon)$. Let $K_{1}$ and $T_{1}$ such that (5.19) holds. Since $\|u(., \epsilon)\|_{L^{r}(\Omega)} \leqq\left\|u_{0}\right\|_{L^{r}(\Omega)}$, we also have $\left\|v_{\epsilon}(0)\right\|_{L^{r}(\Omega)}+K_{1}^{q} T_{1}^{\gamma} \leqq C K_{1}$, thus for any $t \in\left(0, T_{1}\right)$

$$
t^{\theta}\left(\left\|v_{\epsilon}(., t)\right\|_{L^{q r}(\Omega)}+t^{\frac{1}{2}}\left\|\nabla v_{\epsilon}(., t)\right\|_{L^{q r}(\Omega)}\right) \leqq K_{1}
$$

Going to the limit as $\epsilon \rightarrow 0$ from the Fatou Lemma, we obtain

$$
t^{\theta}\left(\|u(., t)\|_{L^{q r}(\Omega)}+t^{\frac{1}{2}}\|\nabla u(., t)\|_{L^{q r}(\Omega)}\right) \leqq K_{1} .
$$

Uniqueness follows in $\left(0, T_{1}\right)$, and by induction on $(0, T)$.
Finally we give existence results for any $u_{0} \in L^{r}(\Omega), r \geq 1$, extending the results of $[8$, Theorem 3.4] for $u_{0} \in L^{1}(\Omega)$, see also [32] for more general operators. We proceed as in Proposition 4.26.

Proposition 5.17 Let $1<q \leq 2$. For any nonnegative $u_{0} \in L^{r}(\Omega), r \geq 1$, there exists a weak $L^{r}$ solution of problem (5.9). And it is unique if $q=2$.

Proof. (i) Case $q<2$. Let $u_{0, n}=\min \left(u_{0}, n\right)$. Then for $\rho>N(q-1) /(2-q)$, from [8, Theorem 3.3], there exists a mild solution $u_{n}$ with initial data $u_{0, n}$, and $u_{n} \in C\left([0, T) ; L^{\rho}(\Omega)\right) \cap$ $L^{q}\left((0, T) ; W_{0}^{1, q \rho}(\Omega) \cap C^{2,1}\left(Q_{\Omega, T}\right)\right.$. Then $u_{n}(., t) \leq e^{t \Delta} u_{0}$, and $\left(u_{n}\right)$ is nondecreasing and $\left|\nabla u_{n}\right|^{q}$ is bounded in $L_{l o c}^{1}\left([0, T) ; L^{1}(\Omega)\right)$ from (5.13). From Corollary 4.20, $\left(u_{n}\right)$ converges in $C_{l o c}^{2,1}\left(Q_{\Omega, T}\right)$ to a weak solution $u$ of (1.1) in $Q_{\Omega, T}$. As a consequence, $u(., t) \leq e^{t \Delta} u_{0}$ and $|\nabla u|^{q} \in L_{l o c}^{1}\left([0, T) ; L^{1}(\Omega)\right)$. From [16, Proposition 2.11], $u(., t)$ converges weakly* to some Radon measure $\mu_{0}$ on $\Omega$. And $e^{t \Delta} u_{0}$ converges to $u_{0}$ in $L^{r}(\Omega)$, thus $\mu_{0} \in L_{l o c}^{1}(\Omega)$ and $0 \leq \mu_{0} \leq u_{0}$. Since $u_{n} \leq u$, there holds $u_{0, n} \leq \mu_{0}$, hence $\mu_{0}=u_{0} \in L^{r}(\Omega)$. Also there exists a function $g \in L^{r}(\Omega)$ such that $u(., t) \leqq g$ for small $t$. Then $e^{t \Delta} u_{0}-u(., t)$ converges weakly* to 0 , and then in $L_{l o c}^{1}(\Omega)$. Hence $u(., t)$ converges to $u_{0}$ in $L_{l o c}^{1}(\Omega)$, then in $L^{r}(\Omega)$ from the dominated convergence theorem. Thus $u \in C\left([0, T) ; L^{r}(\Omega)\right)$.
(ii) Case $q=2$. As in [15, Theorem 4.2], using the classical transformation $v=1-e^{-u}$, it can be shown that there exists a unique solution $u$ such that $u \in C\left([0, T) ; L^{r}(\Omega)\right) \cap C^{2,1}\left(Q_{\Omega, T}\right) \cap$ $C^{1}(\bar{\Omega} \times(0, T))$. Then it is a weak $L^{r}$ solution. Reciprocally any weak $L^{r}$ solution $u$ satisfies the conditions above, from Corollary 5.14 and [16, Theorem 2.17].

## 6 Regularizing effects for quasilinear Dirichlet problems

Here we extend some results of section 5 to a general quasilinear problem, where $u$ may be a signed solution. In this section, we suppose $\Omega$ is a smooth bounded domain in $\mathbb{R}^{N}$.

Let $p>1$ and A be a Caratheodory function on $Q_{\Omega, \infty} \times \mathbb{R} \times \mathbb{R}^{N}$ such that for any $(u, \eta) \in \mathbb{R} \times \mathbb{R}^{N}$, and a.e. $(x, t) \in Q_{\Omega, \infty}$,

$$
\begin{equation*}
|\mathrm{A}(x, t, u, \eta)| \leqq C\left(|\eta|^{p-1}+b(x, t)\right), \quad C>0, \quad b \in L^{p^{\prime}}\left(Q_{\Omega, \infty}\right) \tag{6.1}
\end{equation*}
$$

and A is nonnegative operator:

$$
\begin{equation*}
\mathrm{A}(x, t, u, \eta) \cdot \eta \geqq \nu|\eta|^{p} \quad \nu \geqq 0 \tag{6.2}
\end{equation*}
$$

with no monotonicity assumption.
Let $q>1$ and $g$ be a Caratheodory function on $Q_{\Omega, \infty} \times \mathbb{R}^{+} \times \mathbb{R}^{N}$, such that

$$
\begin{equation*}
g(x, t, u, \eta) u \geqq \gamma|u|^{\lambda+1}|\eta|^{q}, \quad \lambda \geqq 0, \quad \gamma \geqq 0 \tag{6.3}
\end{equation*}
$$

We say that A is coercive if (6.2) holds with $\nu>0$, and $g$ is coercive if (6.3) holds with $\gamma>0$.

We consider the solutions of the Dirichlet problem

$$
\left(P_{\Omega, T}\right)\left\{\begin{array}{l}
u_{t}-\operatorname{div}(\mathrm{A}(x, t, u, \nabla u))+g(x, t, u, \nabla u)=0, \quad \text { in } Q_{\Omega, T},  \tag{6.4}\\
u=0, \quad \text { on } \partial \Omega \times(0, T), \\
u(x, 0)=u_{0},
\end{array}\right.
$$

where $u_{0} \in L^{r}(\Omega), r \geqq 1$ or only $u_{0} \in \mathcal{M}_{b}(\Omega)$.

### 6.1 Solutions of quasilinear heat equation with $L^{1}$ data

First consider the problem in $Q_{\Omega, s, \tau}$

$$
\left\{\begin{array}{l}
u_{t}-\operatorname{div}(\mathrm{A}(x, t, u, \nabla u))=f, \quad \text { in } Q_{\Omega, s, \tau},  \tag{6.5}\\
u=0, \quad \text { on } \partial \Omega \times(s, \tau), \\
u(x, s)=u_{s}
\end{array}\right.
$$

Let us recall the notion of renormalized solution introduced in [18] for this problem with $L^{1}$ data, where the truncations $T_{k}$ are defined by (5.3):

Definition 6.1 Let $s, \tau \in \mathbb{R}$ with $s<\tau$, and $f \in L^{1}\left(Q_{\Omega, s, \tau}\right)$ and $u_{s} \in L^{1}(\Omega)$. A function $u \in$ $L^{\infty}\left((s, \tau) ; L^{1}(\Omega)\right)$ is a renormalized solution in $Q_{\Omega, s, \tau}$ of $(6.5)$ if $T_{k}(u) \in L^{p}\left((s, \tau) ; W_{0}^{1, p}(\Omega)\right)$ for any $k \geqq 0$, and for any $S \in W^{2, \infty}(\mathbb{R})$ such that $S^{\prime}$ has a compact support,

$$
\begin{equation*}
(S(u))_{t}-\operatorname{div}\left(\mathrm{A}(x, t, u, \nabla u) S^{\prime}(u)\right)+S^{\prime \prime}(u)\left(\mathrm{A}(x, t, u, \nabla u) \cdot \nabla u-S^{\prime}(u) f=0 \quad \text { in } \mathcal{D}^{\prime}\left(Q_{\Omega, s, \tau}\right),\right. \tag{6.6}
\end{equation*}
$$

and $u(s)=u_{s}$, and

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \iint_{Q_{\Omega, s, \tau} \cap\{n \leqq u \leqq n+1\}}|\nabla u|^{p} d x d t=0 \tag{6.7}
\end{equation*}
$$

Remark 6.2 The initial condition takes sense from [18], because $S(u)$ lies in the set

$$
\begin{equation*}
E=\left\{\varphi \in L^{p}\left((0, T) ; W_{0}^{1, p}(\Omega)\right): \varphi_{t} \in L^{p^{\prime}}\left((0, T) ; W^{-1, p^{\prime}}(\Omega)\right)+L^{1}\left(Q_{\Omega, T}\right)\right\} \tag{6.8}
\end{equation*}
$$

and $E \subset C\left([0, T] ; L^{1}(\Omega)\right)$. Any function $\varphi \in L^{p}\left((0, T) ; W_{0}^{1, p}(\Omega)\right) \cap L^{\infty}\left(Q_{\Omega, T}\right)$ can be chosen as a test function in equation (6.6). Moreover, from [27, Lemma 7.1], $v=S(u)$ satisfies for any $\psi \in C^{\infty}([s, \tau] \times \bar{\Omega})$ the integration formula

$$
\begin{equation*}
\int_{s}^{\tau}\left\langle v_{t}, M(v) \psi\right\rangle=\int_{\Omega} \mathcal{M}(v(., \tau)) \psi(., \tau) d x-\int_{\Omega} \mathcal{M}(v(., s)) \psi(., s) d x-\int_{s}^{\tau} \int_{\Omega} \psi_{t} \mathcal{M}(v) d x d t \tag{6.9}
\end{equation*}
$$

for any function $M$ continuous and piecewise $C^{1}$ such that $M(0)=0$ and $M^{\prime}$ has a compact support, where $\mathcal{M}(r)=\int_{0}^{r} M(\theta) d \theta$.

A main point in the sequel is the choice of test functions: here we approximate $|u|^{\alpha-1} u$ for $\alpha>0$ by truncation. In the following lemma, we solve some technical difficulties arising because the truncates are not smooth enough to apply the integration formula, and moreover we do not assume $\alpha \geqq 1$.

Lemma 6.3 Let $s, \tau \in \mathbb{R}$ with $s<\tau$, and $f \in L^{1}\left(Q_{\Omega, s, \tau}\right)$. Let $u \in C\left([s, \tau] ; L^{1}(\Omega)\right)$ be any nonnegative renormalized solution in $Q_{\Omega, s, \tau}$ of (6.5), with $u_{s}=u(., s)$. For any $\alpha>0$ and $k>0$, we set

$$
\mathcal{T}_{k, \alpha}(r)=\int_{0}^{r}\left|T_{k}(\theta)\right|^{\alpha-1} T_{k}(\theta) d \theta
$$

Then $\left|T_{k}(u)\right|^{\alpha-1} \mathrm{~A}(x, t, u, \nabla u) . \nabla\left(T_{k}(u)\right) \in L^{1}\left(Q_{\Omega, s, \tau}\right)$ and

$$
\begin{align*}
& \left.\int_{\Omega} \mathcal{T}_{k, \alpha}(u)(., \tau)\right) d x+\alpha \iint_{Q_{\Omega, s, \tau}}\left|T_{k}(u)\right|^{\alpha-1} \mathrm{~A}(x, t, u, \nabla u) . \nabla\left(T_{k}(u)\right) d x d t \\
& \left.=\int_{\Omega} \mathcal{T}_{k, \alpha}(u)(., s)\right) d x+\int_{s}^{\tau} \int_{\Omega} f\left|T_{k}(u)\right|^{\alpha-1} T_{k}(u) d x d t \tag{6.10}
\end{align*}
$$

Proof. Let $\alpha>0, k>0$ be fixed, and for any $n \geqq 2$, and $\theta \in \mathbb{R}$,

$$
S_{n}(\theta)=\int_{0}^{\theta}\left(1-\mid T_{1}\left(s-T_{n}(s) \mid\right) d s, \quad n \geqq 2\right.
$$

This function, introduced in [18], is still a truncation, smoother than $T_{n+1}$, such that $0 \leqq S_{n}(\theta) \theta \leqq$ $T_{n+1}(\theta) \theta, \operatorname{supp} S_{n}^{\prime} \subset[-(n+1), n+1], S_{n}^{\prime \prime}=\chi_{(-n,-n-1) \cup(n, n+1)}$, and $S_{n}\left(T_{k}(\theta)\right)=T_{k}(\theta)$ for any $n>k$. Let $\delta \in(0, \min (1, k))$, and $n>k$. We set

$$
\left.T_{\delta, k, \alpha}(\theta)=\left(\left(T_{k}(|\theta|)+\delta\right)\right)^{\alpha}-\delta^{\alpha}\right) \operatorname{sign} \theta, \quad \mathcal{T}_{\delta, k, \alpha}(r)=\int_{0}^{r} T_{\delta, k, \alpha}(\theta) d \theta
$$

We can take in (6.6) $S=S_{n}$ and $\varphi=T_{\delta, k, \alpha}(u)=T_{\delta, k, \alpha}\left(S_{n}(u)\right)$ as a test function. We obtain

$$
\begin{aligned}
\int_{s}^{t} & <\left(S_{n}(u)\right)_{t}, \varphi>+\int_{s}^{t} \int_{\Omega} S_{n}^{\prime}(u) \mathrm{A}(x, t, u, \nabla u) . \nabla \varphi d x d t \\
& =\int_{s}^{t} \int_{\Omega} S_{n}^{\prime}(u) f \varphi d x d t-\int_{s}^{t} \int_{\Omega} S_{n}^{\prime \prime}(u)(\mathrm{A}(x, t, u, \nabla u) . \nabla u) \varphi d x d t
\end{aligned}
$$

then from (6.9), we deduce

$$
\begin{aligned}
& \alpha \iint_{Q_{\Omega, s, \tau}}\left(T_{k}(|u|)+\delta\right)^{\alpha-1} \mathrm{~A}(x, t, u, \nabla u) \cdot \nabla\left(T_{k}(u)\right) d x d t \\
& =\int_{\Omega} \mathcal{T}_{\delta, k, \alpha}\left(S_{n}(u)(., s)\right) d x-\int_{\Omega} \mathcal{T}_{\delta, k, \alpha}\left(S_{n}(u)(., \tau)\right) d x \\
& +\int_{s}^{\tau} \int_{\Omega} S_{n}^{\prime}(u) f \varphi d x d t-\int_{s}^{t} \int_{\Omega} S_{n}^{\prime \prime}(u)(\mathrm{A}(x, t, u, \nabla u) . \nabla u) \varphi d x d t
\end{aligned}
$$

First we make $\delta \rightarrow 0$. Notice that $\left|\mathcal{T}_{\delta, k, \alpha}(\theta)\right| \leqq(k+1)^{\alpha}|\theta|$ for any $\theta \in \mathbb{R}$, and $S_{n}(u) \in C\left([0, T] ; L^{1}(\Omega)\right)$, and $S_{n}^{\prime}$ is bounded. Thus we can go to the limit in the right hand side. In the left hand side, from the positivity of $A$, and the Fatou Lemma we deduce that

$$
T_{k}(|u|)^{\alpha-1} \mathrm{~A}(x, u, \nabla u) . \nabla T_{k}(u) \in L^{1}\left(Q_{\Omega, s, \tau}\right)
$$

Moreover we can apply dominated convergence theorem. Indeed $\mathrm{A}(x, u, \nabla u) . \nabla T_{k}(u) \in L^{1}\left(Q_{\Omega, s, \tau}\right)$ from (6.1), since $T_{k}(u) \in L^{p}\left((s, \tau) ; W_{0}^{1, p}(\Omega)\right)$, and

$$
\left(T_{k}(|u|)+\delta\right)^{\alpha-1} \mathrm{~A}(x, t, u, \nabla u) . \nabla\left(T_{k}(u)\right) \leqq \max \left(T_{k}^{\alpha-1}(|u|),(k+1)^{\alpha-1}\right) \mathrm{A}(x, u, \nabla u) . \nabla\left(T_{k}(u)\right)
$$

Hence the same relation holds with $\delta=0$, with $T_{0, k, \alpha}(r)=T_{k}^{\alpha-1}(|u|) T_{k}(u)$ :

$$
\begin{aligned}
& \int_{\Omega} \mathcal{T}_{k, \alpha}\left(S_{n}(u)(., \tau)\right) d x-\int_{\Omega} \mathcal{T}_{k, \alpha}\left(S_{n}(u)(., s)\right) d x+\alpha \int_{s}^{t} \int_{\Omega} T_{k}^{\alpha-1}(|u|) \mathrm{A}(x, u, \nabla u) \cdot \nabla\left(T_{k}(u)\right) d x d t \\
& =\int_{s}^{\tau} \int_{\Omega} S_{n}^{\prime}(u) f T_{0, k, \alpha}(u) d x d t-\int_{s}^{t} \int_{\Omega} S_{n}^{\prime \prime}(u)(\mathrm{A}(x, t, u, \nabla u) \cdot \nabla u) T_{0, k,, \alpha}(u) d x d t .
\end{aligned}
$$

Then we make $n \rightarrow \infty$. Since $u \in C\left([0, T] ; L^{1}(\Omega)\right)$, for any $t \in[s, \tau]$, we find

$$
\lim _{n \rightarrow \infty} \int_{\Omega} \mathcal{T}_{k, \alpha}\left(S_{n}(u)(., t)\right) d x=\int_{\Omega} \mathcal{T}_{k, \alpha}(u(., t)) d x
$$

moreover

$$
\lim _{n \rightarrow \infty} \int_{s}^{t} \int_{\Omega} S_{n}^{\prime \prime}(u)(\mathrm{A}(x, t, u, \nabla u) \cdot \nabla u) T_{0, k, \alpha}(u) d x d t=0
$$

from (6.7), (6.1). Finally

$$
\lim _{n \rightarrow \infty} \int_{s}^{\tau} \int_{\Omega} S_{n}^{\prime}(u) f T_{0, k, \alpha}(u) d x d t=\int_{s}^{\tau} \int_{\Omega} f T_{0, k,, \alpha}(u) d x d t
$$

since $S_{n}^{\prime}(u) \rightarrow 1$ a.e. and is uniformly bounded. Then (6.10) follows.

### 6.2 Notion of solutions of problem $\left(P_{\Omega, T}\right)$

Definition 6.4 We say that $u$ is a renormalized solution of problem $\left(P_{\Omega, T}\right)$ if:
(i) $u \in C\left((0, T) ; L^{1}(\Omega)\right), T_{k}(u) \in L_{l o c}^{p}\left((0, T) ; W_{0}^{1, p}(\Omega)\right)$ for any $k \geqq 0$, and $g(x, u, \nabla u) \in L_{l o c}^{1}\left((0, T) ; L^{1}(\Omega)\right)$;
(ii) for any $0<s<\tau<T$, u is a renormalized solution of problem

$$
\left\{\begin{array}{l}
u_{t}-\operatorname{div}(\mathrm{A}(x, t, u, \nabla u))+g(x, t, u, \nabla u)=0, \quad \text { in } Q_{\Omega, s, \tau}, \\
u=0, \quad \text { on } \partial \Omega \times(0, T),
\end{array}\right.
$$

with initial data $u(., s)$;
(iii) for $u_{0} \in L^{r}(\Omega)$, the extension of $u$ by $u_{0}$ at time 0 belongs to $C\left([0, T) ; L^{r}(\Omega)\right)$; for $u_{0} \in \mathcal{M}_{b}(\Omega)$, there holds

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{\Omega} u(., t) \psi d x=\int_{\Omega} \psi d u_{0}, \quad \forall \psi \in C_{b}(\Omega) \tag{6.11}
\end{equation*}
$$

Remark 6.5 Recall that $\nabla u$ is defined by $\nabla u=\nabla\left(T_{k}(u)\right)$ on the set $|u| \leq k$. The assumption on $g$ means that, for any $0<s<\tau<T$,

$$
\int_{Q_{\Omega, s, \tau}}|g(., u, \nabla u)| d x d t=\sum_{k=1}^{\infty} \int_{Q_{\Omega, s, \tau} \cap\{k-1 \leqq|u| \leqq k\}} \mid g\left(., u, \nabla\left(T_{k}(u)\right) \mid d x d t<\infty .\right.
$$

We first prove decay properties of the solutions.

Theorem 6.6 Let $p, q>1$, and $A$ and $g$ satisfying (6.1) (6.2) and (6.3).

1) Let $u_{0} \in L^{r}(\Omega), r \geqq 1$ and $u$ be any renormalized solution of $\left(P_{\Omega, T}\right)$. Then for any $t \in[0, T)$,

$$
\begin{equation*}
\int_{\Omega}|u|^{r}(., t) d x \leqq \int_{\Omega}\left|u_{0}\right|^{r} d x \tag{6.12}
\end{equation*}
$$

Moreover if $r>1$, or if $g$ is coercive, then $\gamma|u|^{\lambda+r-1}|\nabla u|^{q}+\nu|u|^{r-2}|\nabla u|^{p} \in L_{l o c}^{1}\left([0, T) ; L^{1}(\Omega)\right)$, and

$$
\begin{equation*}
\int_{\Omega}|u|^{r}(., t) d x+r \gamma \int_{0}^{t} \int_{\Omega}|u|^{\lambda+r-1}|\nabla u|^{q} d x d t+r(r-1) \nu \int_{0}^{t} \int_{\Omega}|u|^{r-2}|\nabla u|^{p} d x d t \leqq \int_{\Omega}\left|u_{0}\right|^{r} d x \tag{6.13}
\end{equation*}
$$

2) Let $u_{0} \in \mathcal{M}_{b}^{+}(\Omega)$ and $u$ be any nonnegative renormalized solution of ( $P_{\Omega, T}$ ) of problem (5.9). Then the same conclusions hold as in case $u_{0} \in L^{1}(\Omega)$, where the norm $\left\|u_{0}\right\|_{L^{1}(\Omega)}$ is replaced by $\int_{\Omega} d u_{0}$.

Proof. 1) Let $0<s<t<T$. Then for any $\alpha>0$, any $k>0$, from Lemma 6.3,

$$
\begin{aligned}
& \left.\int_{\Omega} \mathcal{T}_{k, \alpha}(u)(., \tau)\right) d x+\alpha \int_{s}^{t} \int_{\Omega}\left|T_{k}(u)\right|^{\alpha-1} \mathrm{~A}(x, t, u, \nabla u) . \nabla\left(T_{k}(u)\right) d x d t \\
& \left.=\int_{\Omega} \mathcal{T}_{k, \alpha}(u)(., s)\right) d x-\int_{s}^{\tau} \int_{\Omega}\left|T_{k}(u)\right|^{\alpha-1} T_{k}(u) g(., u, \nabla u) d x d t
\end{aligned}
$$

And $\left|T_{k}(u)\right|^{\alpha-1} T_{k}(u) g(., u, \nabla u) \geq \gamma\left|T_{k}(u)\right|^{\alpha+\lambda}\left|\nabla T_{k}(u)\right|^{q}$ from (6.3). Therefore $\left.\int_{\Omega} \mathcal{T}_{k, \alpha}(u)(., t)\right)$ is decreasing for any $k, \alpha>0$, and

$$
\begin{align*}
& \left.\int_{\Omega} \mathcal{T}_{k, \alpha}(u)(., \tau)\right) d x+\gamma \int_{s}^{t} \int_{\Omega}\left|T_{k}(u)\right|^{\alpha+\lambda}\left|\nabla T_{k}(u)\right|^{q} d x d t+\alpha \nu \int_{s}^{t} \int_{\Omega}\left|T_{k}(u)\right|^{\alpha-1}\left|\nabla T_{k}(u)\right|^{p} d x d t \\
& \left.\leqq \int_{\Omega} \mathcal{T}_{k, \alpha}(u)(., s)\right) d x \tag{6.14}
\end{align*}
$$

- If $r>1$, we can take $\alpha=r-1>0$ in (6.14) and get

$$
\begin{align*}
& \left.\int_{\Omega} \mathcal{T}_{k, r-1}(u)(., t)\right) d x+\gamma \int_{s}^{t} \int_{\Omega}\left|T_{k}(u)\right|^{r-1+\lambda}\left|\nabla T_{k}(u)\right|^{q} d x d t+\alpha \nu \int_{s}^{t} \int_{\Omega}\left|T_{k}(u)\right|^{r-2}\left|\nabla T_{k}(u)\right|^{p} d x d t \\
& \left.\leqq \int_{\Omega} \mathcal{T}_{k, r-1}(u)(., s)\right) d x \leqq \frac{1}{r} \int_{\Omega}|u|^{r}(., s) d x \tag{6.15}
\end{align*}
$$

Since $u \in C\left([0, T) ; L^{r}(\Omega)\right)$ we can go to the limit as $k \rightarrow \infty$, and $s \rightarrow 0$; we obtain that $\gamma|u|^{r-1+\lambda}|\nabla u|^{q}$ and $\alpha \nu|u|^{r-2}|\nabla u|^{p}$ belong to $L_{l o c}^{1}\left([0, T) ; L^{1}(\Omega)\right)$; and for any $t \in(0, T)$,

$$
\int_{\Omega}|u|^{r}(., t) d x+r \gamma \int_{0}^{t} \int_{\Omega}|u|^{r-1+\lambda}|\nabla u|^{q} d x d t+r(r-1) \nu \int_{0}^{t} \int_{\Omega}|u|^{r-2}|\nabla u|^{p} d x d t \leqq \int_{\Omega}\left|u_{0}\right|^{r} d x
$$

- If $r=1$, we take any $\alpha>0$ in (6.14); notice that

$$
\begin{equation*}
\frac{\left|T_{k}(\theta)\right|^{\alpha+1}}{\alpha+1} \leqq \mathcal{T}_{k, \alpha}(\theta) \leqq k^{\alpha}|\theta| \tag{6.16}
\end{equation*}
$$

for any $\theta>0$. Then

$$
\left.\int_{\Omega}\left|T_{k}(u)\right|^{\alpha+1}(., t)\right) d x+(\alpha+1) \gamma \int_{s}^{t} \int_{\Omega}\left|T_{k}(u)\right|^{\alpha+\lambda}\left|\nabla T_{k}(u)\right|^{q} d x d t \leqq(\alpha+1) k^{\alpha} \int_{\Omega}|u|(., s) d x
$$

Going to the limit as $\alpha \rightarrow 0$, we deduce

$$
\begin{equation*}
\left.\int_{\Omega}\left|T_{k}(u)\right|(., t)\right) d x+\gamma \int_{s}^{t} \int_{\Omega}\left|T_{k}(u)\right|^{\lambda}\left|\nabla T_{k}(u)\right|^{q} d x d t \leqq \int_{\Omega}|u|(., s) d x \tag{6.17}
\end{equation*}
$$

and then as $s \rightarrow 0$ we find

$$
\begin{equation*}
\left.\int_{\Omega}\left|T_{k}(u)\right|(., t)\right) d x+\gamma \int_{s}^{t} \int_{\Omega}\left|T_{k}(u)\right|^{\lambda}\left|\nabla T_{k}(u)\right|^{q} d x d t \leqq \int_{\Omega}\left|u_{0}\right| d x \tag{6.18}
\end{equation*}
$$

and finally as $k \rightarrow \infty$, we obtain that $\int_{\Omega}|u|(., t) d x \leqq \int_{\Omega}\left|u_{0}\right| d x$. Moreover if $\gamma>0$, we find

$$
\int_{\Omega}|u|(., t) d x+\gamma \int_{0}^{t} \int_{\Omega}|u|^{\lambda}|\nabla u|^{q} d x d t \leqq \int_{\Omega}\left|u_{0}\right| d x
$$

thus (6.13) still holds with $r=1$.
2) We still find (6.17). And $\lim _{s \rightarrow 0} \int_{\Omega} u(., s) d x=\int_{\Omega} d u_{0}$ from (6.11), hence the conclusion.

Next we deduce $L^{\infty}$ estimates, in particular a universal one.
Theorem 6.7 Let $p, q>1$, and $A$ and $g$ satisfying (6.1) (6.2) and (6.3). Let $u_{0} \in L^{r}(\Omega), r \geqq 1$, and $u$ be any renormalized solution of $\left(P_{\Omega, T}\right)$.
(i) If $g$ is coercive, then

$$
\|u(., t)\|_{L^{\infty}(\Omega)} \leqq\left\{\begin{array}{cc}
C t^{-\sigma_{r, q, \lambda}}\left\|u_{0}\right\|_{L^{r},(, N, \lambda}^{\varpi_{r, q}}, & C=C(N, q, r, \lambda, \gamma),  \tag{6.19}\\
C_{\varepsilon} t^{-(1+\varepsilon) \sigma_{r, n, \lambda}}\left\|u_{0}\right\|_{L_{r, q, N, \lambda}(\Omega)}, & C_{\varepsilon}=C(N, q, r, \lambda, \gamma, \varepsilon), \\
\text { if } q=N
\end{array}\right.
$$

where

$$
\sigma_{r, q, N, \lambda}=\frac{1}{\frac{r q}{N}+\lambda+q-1}=\frac{N}{r q} \varpi_{r, q, N, \lambda} .
$$

Moreover

$$
\begin{equation*}
\|u(., t)\|_{L^{\infty}(\Omega)} \leqq C t^{-\frac{1}{q-1+\lambda}}, \quad C=C(N, q, \lambda,|\Omega|) \tag{6.20}
\end{equation*}
$$

(ii) If A is coercive and $r>(2-p) N / p$, in particular if $p>2 N /(N+1)$, then

$$
\|u(., t)\|_{L^{\infty}(\Omega)} \leqq\left\{\begin{array}{cc}
C t^{-\sigma_{r, p, N,-1}}\left\|u_{0}\right\|_{L^{r}(\Omega), N ?-1}^{\varpi_{r, p}(\Omega)}, & C=C(N, p, r, \nu, \Omega),  \tag{6.21}\\
C_{\varepsilon} t^{-(1+\varepsilon) \sigma_{r, N, N,-1}}\left\|u_{0}\right\|_{L_{r, p, N,-1}^{r}(\Omega)}^{\varpi_{r}}, & C_{\varepsilon}=C(N, p, r, \nu, \Omega, \varepsilon), \\
\text { if } p=N
\end{array}\right.
$$

where

$$
\sigma_{r, p, N,-1}=\frac{1}{\frac{r p}{N}+p-2}=\frac{N}{r p} \varpi_{r, p, N,-1}
$$

Moreover if $p>2$, then

$$
\begin{equation*}
\|u(., t)\|_{L^{\infty}(\Omega)} \leqq C t^{-\frac{1}{p-2}}, \quad C=C(N, p,|\Omega|) \tag{6.22}
\end{equation*}
$$

(iii) The same conclusions hold if $u$ is nonnegative and $u_{0} \in \mathcal{M}_{b}^{+}(\Omega)$, as in case $u_{0} \in L^{1}(\Omega)$, where the norm $\left\|u_{0}\right\|_{L^{1}(\Omega)}$ is replaced by $\int_{\Omega} d u_{0}$. In particular (6.22) holds for $p>2$.

Proof. (i) Let $0<s<t<T$. Since $g$ is coercive, from Theorem 6.6, for any $\alpha \geqq 0$ such that $|u|^{\alpha+1}(., s) \in L^{1}(\Omega)$, there holds

$$
\int_{\Omega}|u|^{\alpha+1}(., t) d x+(\alpha+1) \gamma \int_{s}^{t} \int_{\Omega}|u|^{\lambda+\alpha}|\nabla u|^{q} d x d t \leqq \int_{\Omega}|u|^{\alpha+1}(., s) d x,
$$

from (6.13); in particular

$$
\int_{\Omega}\left|T_{k}(u)\right|^{\alpha+1}(., t) d x+(\alpha+1) \gamma \int_{s}^{t} \int_{\Omega}\left(T_{k}(u)\right)^{\lambda+\alpha}\left|\nabla T_{k}(u)\right|^{q} d x d t \leqq \int_{\Omega}|u|^{\alpha+1}(., s) d x
$$

And $|u|^{\lambda+\alpha}|\nabla u|^{q}=\left|\nabla\left(|u|^{\beta-1} u\right)\right|^{q}$ with $\beta=1+(\alpha+\lambda) / q \geqq 1$. Then $\left|\nabla\left(\left(|u|^{\beta-1} u\right)(., t)\right)\right|$, and also $\left|\nabla\left(\left(\left|T_{k}(u)\right|^{\beta-1} T_{k}(u)\right)(., t)\right)\right|$ belong to $L^{q}(\Omega)$ for almost any $t \in(0, T)$. Since $\left|T_{k}(u)\right|^{\beta-1} T_{k}(u)(., t) \in$ $L^{\infty}(\Omega)$, it follows that $\left|T_{k}(u)\right|^{\beta-1} T_{k}(u)(., t) \in W^{1, q}(\Omega)$. Moreover $\left.T_{k}(u)(., t) \in W_{0}^{1, p}(\Omega)\right)$, hence $\left|T_{k}(u)\right|^{\beta-1} T_{k}(u)(., t) \in W_{0}^{1, q}(\Omega)$. If $q<N$, we deduce

$$
\int_{\Omega}\left|T_{k}(u)\right|^{\alpha+1}(., t) d x+\gamma \frac{(\alpha+1) C(N, q)}{\beta^{q}} \int_{s}^{t}\left(\int_{\Omega}\left|T_{k}(u)\right|^{\beta \frac{N q}{N-q}}(., \sigma) d x\right)^{\frac{N-q}{N}} d \sigma \leqq \int_{\Omega}|u|^{\alpha+1}(., s) d x
$$

Going to the limit as $k \rightarrow \infty$, we find

$$
\int_{\Omega}|u|^{\alpha+1}(., t) d x+\gamma \frac{(\alpha+1) C(N, q)}{\beta^{q}} \int_{s}^{t}\left(\int_{\Omega}|u|^{\beta \frac{N q}{N-q}}(., \sigma) d x\right)^{\frac{N-q}{N}} d \sigma \leqq \frac{1}{\alpha+1} \int_{\Omega}|u|^{\alpha+1}(., s) d x .
$$

Then we can apply Lemma 2.1 on $[\epsilon, T)$, with $m=q$ and $\theta=N /(N-q)$; indeed (2.1) is satisfied, since $\lambda \geqq 0$; we deduce the estimate for $[\epsilon, T)$,

$$
\|u(., t)\|_{L^{\infty}(\Omega)} \leqq C(t-\epsilon)^{-\sigma_{r, q, N, \lambda}}\|u(., \epsilon)\|_{L^{r}(\Omega)}^{\varpi_{r, q, N, \lambda}},
$$

with $C=C(N, q, r, \lambda, \gamma, \Omega)$. Going to the limit as $\epsilon \rightarrow 0$, we get (6.20), and (6.19) for $u_{0} \in L^{r}(\Omega)$, and the analogous when $u_{0} \in \mathcal{M}_{b}^{+}(\Omega)$. In case $q \geqq N$ we proceed as in Theorem 5.12.
(ii) Assume that A is coercive. Then for any $\alpha>0$,

$$
\left.\left.\int_{\Omega} \mathcal{T}_{k, \alpha}(u)(., t)\right) d x+\alpha \nu \int_{s}^{t} \int_{\Omega}\left|T_{k}(u)\right|^{\alpha-1}\left|\nabla T_{k}(u)\right|^{p} d x d t \leqq \int_{\Omega} \mathcal{T}_{k, \alpha}(u)(., s)\right) d x
$$

from (6.14). First assume $p<N$. From the Sobolev injection of $W_{0}^{1, p}(\Omega)$ into $L^{N p /(N-p)}(\Omega)$, we deduce

$$
\frac{1}{\alpha+1} \int_{\Omega}|u|^{\alpha+1}(., t) d x+\alpha \nu \frac{C(N, p)}{k^{p}} \int_{s}^{t}\left(\int_{\Omega}^{k \frac{N p}{N-p}}|u|(., \sigma) d x\right)^{\frac{N-p}{N}} d t \leqq \frac{1}{\alpha+1} \int_{\Omega}|u|^{\alpha+1}(., s) d x
$$

with $k=1+(\alpha-1) / p$.
First suppose $r>1$; then we start from $\alpha_{0}=r-1>0$, and we can apply Lemma 2.1 with $C_{0}=(r-1) \nu C(N, p), m=p, \theta=N /(N-p)$ and $\lambda=-1$; indeed (2.1) is satisfied, since $r>N(2-p) / p$.

Next suppose $r=1$. Then $1>(2-p) N / p$, thus $p-1+p / N>1$. For any $\alpha>0$,

$$
\left.\int_{\Omega}\left|T_{k}(u)\right|^{\alpha+1}(., t)\right) d x+\alpha(\alpha+1) \nu \int_{s}^{t} \int_{\Omega}\left|T_{k}(u)\right|^{\alpha-1}\left|\nabla T_{k}(u)\right|^{p} d x d t \leqq(\alpha+1) k^{\alpha} \int_{\Omega}|u|(., s) d x
$$

Taking $\alpha=1$, we get from (6.12),

$$
\nu \int_{s}^{t} \int_{\Omega}\left|\nabla T_{k}(u)\right|^{p} d x d t \leqq k \int_{\Omega}|u|(., s) d x \leqq k \int_{\Omega}\left|u_{0}\right| d x .
$$

And from (6.12), $u \in L^{\infty}\left((s, T) ; L^{1}(\Omega)\right)$. From standard estimates, there holds $u \in L^{\rho}\left(Q_{\Omega, s, t}\right)$ for any $\rho \in(1, p-1+p / N)$, see [19]. Then $|u|^{\rho}(., t) \in L^{1}(\Omega)$ for almost any $t \in(0, T)$. Hence we can apply Lemma 2.1 on $[\epsilon, T)$ for $\epsilon>0$, with the same parameters, after fixing such a $\rho=\rho_{p, N}$ such that $\rho N(2-p) / p<1$. We obtain that

$$
\|u(., t)\|_{L^{\infty}(\Omega)} \leqq C(t-\epsilon)^{-\sigma_{1, p,-1}}\|u(., \epsilon)\|_{L^{1}(\Omega)}^{\varpi_{1, p,-1}},
$$

where $C=C\left(N, p \rho_{p, N}\right)=C(N, p)$; finally we go to the limit as $\epsilon \rightarrow 0$ because $u \in C\left([0, T] ; L^{1}(\Omega)\right)$. Estimate (6.22) follows, since $-1+p-1>0$.

If $p=N$, we proceed as above, applying Lemma 2.1 with $m=N, \lambda=-1$ and $\theta>1$ arbitrary. Next assume $p>N$. In case $r>1$, there holds, for any $t \in(0, T)$,

$$
r(r-1) \frac{\nu}{\kappa^{p}} \int_{0}^{t} \int_{\mathbb{R}^{N}}\left|\nabla\left(|u|^{\kappa}\right)\right|^{p} d x d t \leqq \int_{\mathbb{R}^{N}}\left|u_{0}\right|^{r} d x
$$

where $\kappa=1+(r-2) / p>0$. From Lemma 3.4, applied to $v=|u|^{\kappa}$, with $m=p, 1 / k=1+r(p-$ $N) / N p \kappa$, we obtain

$$
\|u(., t)\|_{L^{\infty}(\mathbb{R})}^{\frac{\kappa p}{k}} \leqq C\|u(., t)\|_{L^{r}(\mathbb{R})}^{\frac{\kappa(1-k) p}{k}} \int_{\mathbb{R}}\left|\nabla\left(|u|^{\kappa}\right)\right|^{p} d x d t
$$

and by integration, with a new constant $C=C(N, p, r, \nu)$,

$$
t\|u(., t)\|_{L^{\infty}(\mathbb{R})}^{\frac{k p}{k}} \leqq C\left\|u_{0}\right\|_{L^{r}\left(\mathbb{R}^{N}\right)}^{r+\frac{\kappa(1-k) p}{k}},
$$

which is precisely (6.21). In case $r=1$, we choose $\rho=p \in(1, p-1+p / N)$, and obtain from above, for any $0<\epsilon<s<t<T$,

where $C=C(N, p, \nu)$. From Lemma 3.1, we deduce precisely

$$
\|u(., t)\|_{L^{\infty}(\mathbb{R})} \leqq C(t-\epsilon)^{-\sigma_{1, p, p-1}}\left\|u_{0}\right\|_{L^{1}(\mathbb{R})}^{w_{1, N,-1}}
$$

and we conclude as $\epsilon \rightarrow 0$.
(iii) We obtain the estimates on $(\epsilon, T)$ as above and go to the limit as $\epsilon \rightarrow 0$.

Remark 6.8 Our results apply in particular to the problem

$$
\left\{\begin{array}{l}
u_{t}-\operatorname{div}(\mathrm{A}(x, t, u, \nabla u))=0, \quad \text { in } Q_{\Omega, T}, \\
u=0, \quad \text { on } \partial \Omega \times(0, T), \\
u(x, 0)=u_{0}
\end{array}\right.
$$

Thus we find again and improve the estimates of [33, Theorem 5.3], with less regularity on the solutions: those estimates were proved for solutions $u \in C\left([0, T) ; L^{r}(\Omega)\right)$ such that $u \in L^{p}\left((0, T) ; W_{0}^{1, p}(\Omega)\right) \cap$ $C\left([0, T) ; L^{2}(\Omega)\right)$. The notion of renormalized solutions, equivalent to the notion of entropy solutions of [35] (see [27]), is weaker. Moreover our results in case $p>N$ are optimal.

Remark 6.9 The extension of results of section 4 to the case of equation of type (1.2) in the case $\Omega=\mathbb{R}^{N}$ will be treated a further article.

## 7 Appendix

Proof of Lemma 4.10. (i) Let $u$ be a mild $\mathcal{M}$ solution. Then clearly (4.10) holds. Moreover for any $\psi \in C_{0}\left(\mathbb{R}^{N}\right)$, from the assumption on the gradient,

$$
<e^{t \Delta} u_{0}, \psi>=<u_{0}, e^{t \Delta} \psi>=\int_{\mathbb{R}^{n}} e^{t \Delta} \psi d u_{0}=\int_{\mathbb{R}^{N}}\left(u(., t)+\int_{0}^{t} e^{(t-s) \Delta}|\nabla u(., s)|^{q} d s\right) \psi d x
$$

The relation extends to any $\varphi \in C_{b}\left(\mathbb{R}^{N}\right)$ : we can assume that $\varphi \geqq 0$; from the Beppo-Levi theorem,

$$
\begin{aligned}
\int_{\mathbb{R}^{N}} e^{t \Delta} \varphi d u_{0} & =\int_{\mathbb{R}^{N}} u(., t) \varphi d x+\int_{\mathbb{R}^{N}}\left(\int_{0}^{t} e^{(t-s) \Delta}|\nabla u(., s)|^{q} d s\right) \varphi d x \\
& =\int_{\mathbb{R}^{N}} u(., t) \varphi d x+\int_{0}^{t} \int_{\mathbb{R}^{N}}|\nabla u|^{q} \varphi d x d s,
\end{aligned}
$$

since the measure is bounded. From the integrability of the gradient and the dominated convergence theorem in $L^{1}\left(\mathbb{R}^{N}, d u_{0}\right)$, we deduce

$$
\lim _{t \rightarrow 0} \int_{0}^{t} \int_{\mathbb{R}^{N}}|\nabla u|^{q} \varphi d x d s=0, \quad \lim _{t \rightarrow 0} \int_{\mathbb{R}^{N}} e^{t \Delta} \varphi d \mu_{0}=\int_{\mathbb{R}^{N}} \varphi d \mu_{0}
$$

since $\left\|e^{t \Delta} \varphi\right\|_{L^{\infty}\left(\mathbb{R}^{N}\right)} \leq\|\varphi\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}$ and $e^{t \Delta} \varphi$ converges to $\varphi$ everywhere as $t \rightarrow 0$; thus (4.11a) holds.
(ii) Let $u$ be a weak semi-group solution. Then obviously $u \in C_{b}\left((0, T) ; L^{1}\left(\mathbb{R}^{N}\right)\right)$. As $\epsilon \rightarrow 0$, we have

$$
\lim _{\epsilon \rightarrow 0} \int_{\epsilon}^{t} e^{(t-s) \Delta}|\nabla u(., s)|^{q} d s=\int_{0}^{t} e^{(t-s) \Delta}|\nabla u(., s)|^{q} d s \quad \text { in } L^{1}\left(\mathbb{R}^{N}\right) .
$$

Then

$$
\lim _{\epsilon \rightarrow 0} e^{(t-\epsilon) \Delta} u(., \epsilon)=u(., t)+\int_{0}^{t} e^{(t-s) \Delta}|\nabla u(., s)|^{q} d s \quad \text { in } L^{1}\left(\mathbb{R}^{N}\right)
$$

Moreover (4.11a) entails that that $u(., \epsilon) \rightarrow u_{0}$ in $\mathcal{S}^{\prime}\left(\mathbb{R}^{N}\right)$ and

$$
\begin{equation*}
\lim _{\epsilon \rightarrow 0} e^{(t-\epsilon) \Delta} u(., \epsilon)=e^{t \Delta} u_{0} \quad \text { in } \mathcal{S}^{\prime}\left(\mathbb{R}^{N}\right) ; \tag{7.1}
\end{equation*}
$$

indeed for any $\varphi \in \mathcal{S}\left(\mathbb{R}^{N}\right)$,

$$
\begin{aligned}
\left|<e^{(t-\epsilon) \Delta} u(., \epsilon)-e^{t \Delta} u_{0}, \varphi\right\rangle \mid & \leq \mid<e^{t \Delta}\left(u(., \epsilon)-u_{0}(.), \varphi>|+| \int_{\mathbb{R}^{N}}\left(u(x, \epsilon)\left(\left(e^{(t-\epsilon) \Delta}-e^{t \Delta}\right) \varphi\right)(x) d x \mid\right.\right. \\
& \leq \mid<e^{t \Delta}\left(u(., \epsilon)-u_{0}(.), \varphi>\mid\right. \\
& \left.+\|u(., \epsilon)\|_{L^{1}\left(\mathbb{R}^{N}\right)}\left\|\left(e^{(t-\epsilon) \Delta}-e^{t \Delta}\right) \varphi\right\|_{L^{\infty}\left(\mathbb{R}^{N}\right)}\right)
\end{aligned}
$$

and $e^{t \Delta}$ is continuous on $\mathcal{S}\left(\mathbb{R}^{N}\right)$. Hence, for any $\varphi \in \mathcal{S}\left(\mathbb{R}^{N}\right)$, we get

$$
<e^{t \Delta} u_{0}, \varphi>=\int_{\mathbb{R}^{n}} u(., t) \varphi d x+\int_{\mathbb{R}^{n}}\left(\int_{0}^{t} e^{(t-s) \Delta}|\nabla u(., s)|^{q} d s\right) \varphi d x
$$

which extends to any $\varphi \in C_{0}\left(\mathbb{R}^{N}\right)$ by density. Thus (4.9) follows.
Acknowledgement 7.1 We thank Professor F. Weissler for helpfull discussions during the preparation of this article.

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