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### Local Doppler effect, index of refraction through the Earth crust, PDF and the CNGS neutrino anomaly?

Armando V.D.B. Assis

Departamento de Física, Universidade Federal de Santa Catarina - UFSC, Trindade 88040-900, Florianópolis, SC, Brazil. Email: armando.assis@pgfsc.ufsc.br

In this brief paper, we show the neutrino velocity discrepancy obtained in the OPERA experiment may be due to the local Doppler effect between a local clock attached to a given detector at Gran Sasso, say  $C_G$ , and the respective instantaneous clock crossing  $C_G$ , say  $C_C$ , being this latter at rest in the instantaneous inertial frame having got the velocity of rotation of CERN about Earth's axis in relation to the fixed stars. With this effect, the index of refraction of the Earth crust may accomplish a refractive effect by which the neutrino velocity through the Earth crust turns out to be small in relation to the speed of light in the empty space, leading to an encrusted discrepancy that may have contamined the data obtained from the block of detectors at Gran Sasso, leading to a time interval excess  $\epsilon$  that did not provide an exact match between the shift of the protons PDF (probability distribution function) by TOF<sub>c</sub> and the detection data at Gran Sasso via the maximum likelihood matching.

#### 1 Definitions and Solution

Firstly, the effect investigated here is not the same one that was investigated in [2], but, throughout this paper, we will use some useful configurations defined in [2]. The relative velocity between Gran Sasso and CERN due to the Earth daily rotation may be written:

$$\vec{v}_G - \vec{v}_C = 2\omega R \sin \alpha \hat{e}_z,\tag{1}$$

where  $\hat{e}_z$  is a convenient unitary vector, the same used in [2],  $\omega$  is the norm of the Earth angular velocity vector about its daily rotation axis, being *R* given by:

$$R_E = \frac{R}{\cos\lambda},\tag{2}$$

where  $R_E$  is the radius of the Earth, its averaged value  $R_E = 6.37 \times 10^6$  m, and  $\alpha$  given by:

$$\alpha = \frac{1}{2} \left( \alpha_G - \alpha_C \right), \tag{3}$$

where  $\alpha_C$  and  $\alpha_G$  are, respectively, CERN's and Gran Sasso's longitudes ( $\leftarrow WE \rightarrow$ ). Consider the inertial (in relation to the fixed stars) reference frame  $O_C x_C y_C z_C \equiv O xyz$  in [2]. This is the lab reference frame and consider this frame with its local clocks at each spatial position as being ideally synchronized, viz., under an ideal situation of synchronicity between the clocks of  $O_C x_C y_C z_C \equiv O xyz$ . This situation is the expected ideal situation for the OPERA collaboration regarding synchronicity in the instantaneous lab (CERN) frame.

Now, consider an interaction between a single neutrino and a local detector at Gran Sasso. This event occurs at a given spacetime point  $(t_v, x_v, y_v, z_v)$  in  $O_C x_C y_C z_C \equiv Oxyz$ . The interaction instant  $t_v$  is measured by a local clock  $C_C$  at rest at  $(x_v, y_v, z_v)$  in the lab frame, viz., in the  $O_C x_C y_C z_C \equiv$ Oxyz frame. But, under gedanken, at this instant  $t_{y}$ , according to  $O_C x_C y_C z_C \equiv O x y z$ , there is a clock  $C_G$  attached to the detector at Gran Sasso that crosses the point  $(t_v, x_v, y_v, z_v)$ with velocity given by Eq. (1). Since  $C_G$  crosses  $C_C$ , the Doppler effect between the proper tic-tac rates measured at each location of  $C_C$  and  $C_G$ , viz., measured at their respective locations in their respective reference frames (the reference frame of  $C_G$  is the  $O_G x_G y_G z_G \equiv \tilde{O} \tilde{x} \tilde{y} \tilde{z}$  in [2], also inertial in relation to the fixed stars), regarding a gedanken control tic-tac rate continuosly sent by  $C_C$ , say via electromagnetic pulses from  $C_C$ , is not transverse. Since the points at which  $C_C$  and  $C_G$  are at rest in their respective reference frames will instantaneously coincide, better saying, will instantaneously intersect, at  $t_{y}$  accordingly to  $C_{C}$ , they must be previously approximating, shortening their mutual distance during the interval  $t_{\nu} - \delta t_{\nu} \rightarrow t_{\nu} \ (\delta t_{\nu} \rightarrow 0, \ \delta t_{\nu} > 0)$  along the line passing through these clocks as described in the  $C_C$  world.

Suppose  $C_C$  sends N electromagnetic pulses to  $C_G$ . During the  $C_C$  proper time interval  $(t_v - \delta t_v) - 0 = t_v - \delta t_v^*$  within which  $C_C$  emits the N electromagnetic pulses, the first emitted pulse travels the distance  $c(t_v - \delta t_v)$  and reaches the clock  $C_G$ , as described by  $C_C$ . Within this distance, there are N

<sup>\*</sup>The initial instant  $C_C$  starts to emit the electromagnetic pulses is set to zero in both the frames  $O_C x_C y_C z_C \equiv O xyz$  and  $O_G x_G y_G z_G \equiv \tilde{O} \tilde{x} \tilde{y} \tilde{z}$ ; zero also is the instant the neutrino starts the travel to Gran Sasso in  $O_C x_C y_C z_C \equiv$ O xyz; hence the instant the neutrino starts the travel to Gran Sasso and the emission of the first pulse by  $C_C$  are simultaneous events in  $O_C x_C y_C z_C \equiv$ O xyz. These events are simultaneous in  $O_G x_G y_G z_G \equiv \tilde{O} \tilde{x} \tilde{y} \tilde{z}$  too, since they have got the same spatial coordinate  $z_c = z = 0$  along the  $O_C z_C \equiv O z$ direction as defined in [2]. The relative motion between CERN and Gran Sasso is parallel to this direction. The only one difference between these events is the difference in their  $x_C = x$  coordinates, being  $x_C = 0$  for the neutrino departure and  $x_C = L = 7.3 \times 10^5$  m for  $C_C$ , being these locations perpendicularly located in relation to the relative velocity given by the Eq, (1).

equally spaced distances between consecutive pulses as described in the  $C_C$  world, say  $\lambda_C$ :

$$\mathcal{N}\lambda_C = c\left(t_v - \delta t_v\right). \tag{4}$$

Also, since the clocks  $C_C$  and  $C_G$  will intersect at  $t_v$ , as described in  $O_C x_C y_C z_C \equiv O x y z$ , during the interval  $\delta t_v$ , the clock  $C_G$  must travel the distance  $2\omega R \sin \alpha \, \delta t_v$  in the  $C_C$  world to accomplish the matching spatial intersection at the instant  $t_v$ , hence the clock  $C_G$  travels the  $2\omega R \sin \alpha \, \delta t_v$  in the  $C_C$  world, viz., as described by  $C_C$  in  $O_C x_C y_C z_C \equiv O x y z$ :

$$\mathcal{N}\lambda_C = 2\omega R \sin \alpha \, \delta t_\nu \Rightarrow \delta t_\nu = \mathcal{N} \frac{\lambda_C}{2\omega R \sin \alpha}.$$
 (5)

Solving for  $t_{\nu}$ , from the Eqs. (4) and (5), one reaches:

$$t_{\nu} = \frac{N\lambda_C}{c} \left( 1 + \frac{c}{2\omega R \sin \alpha} \right).$$
(6)

Now, from the perspective of  $C_G$ , in  $O_G x_G y_G z_G \equiv \tilde{O} \tilde{x} \tilde{y} \tilde{z}$ , there must be N electromagnetic pulses covering the distance:

$$c\left(t_{\nu}^{G}-\delta t_{\nu}^{G}\right)-2\omega R\sin\alpha\left(t_{\nu}^{G}-\delta t_{\nu}^{G}\right),\tag{7}$$

where  $t_{\nu}^{G} - \delta t_{\nu}^{G}$  is the time interval between the non-proper instants  $t^{G} = t_{\nu} = 0$ , at which the  $C_{C}$  clock sends the first pulse, and the instant  $t_{\nu}^{G} - \delta t_{\nu}^{G}$ , at which this first pulse reaches  $C_{G}$ , as described by  $C_{G}$  in its world  $O_{G}x_{G}y_{G}z_{G} \equiv \tilde{O}\tilde{x}\tilde{y}\tilde{z}$ . Within this time interval,  $t_{\nu}^{G} - \delta t_{\nu}^{G}$ ,  $C_{G}$  describes, in its  $O_{G}x_{G}y_{G}z_{G} \equiv \tilde{O}\tilde{x}\tilde{y}\tilde{z}$ world, the clock  $C_{C}$  approximating the distance:

$$2\omega R\sin\alpha \left(t_{\nu}^{G}-\delta t_{\nu}^{G}\right),\tag{8}$$

with the first pulse traveling:

$$c\left(t_{\nu}^{G}-\delta t_{\nu}^{G}\right),\tag{9}$$

giving the distance within which there must be N equally spaced pulses, say, spaced by  $\lambda_G$ , as described by  $C_G$  in its  $O_G x_G y_G z_G \equiv \tilde{O} \tilde{x} \tilde{y} \tilde{z}$  world:

$$\mathcal{N}\lambda_G = (c - 2\omega R \sin \alpha) \left( t_{\nu}^G - \delta t_{\nu}^G \right). \tag{10}$$

With similar reasoning that led to the Eq. (5), now in the  $O_G x_G y_G z_G \equiv \tilde{O} \tilde{x} \tilde{y} \tilde{z} C_G$  world, prior to the spatial matching intersection between  $C_C$  and  $C_G$ , the  $C_C$  clock must travel the distance  $N\lambda_G$  during the time interval  $\delta t_v^G$ , with the  $C_C$  approximation velocity  $2\omega R \sin \alpha$ :

$$\mathcal{N}\lambda_G = 2\omega R \sin \alpha \,\delta t_{\nu}^G \Rightarrow \delta t_{\nu}^G = \mathcal{N} \frac{\lambda_G}{2\omega R \sin \alpha}.$$
 (11)

From Eqs. (10) and (11), we solve for  $t_{\nu}^{G}$ :

$$t_{\nu}^{G} = \mathcal{N} \frac{\lambda_{G}}{2\omega R \sin \alpha} \frac{1}{\left[1 - (2\omega R \sin \alpha)/c\right]}.$$
 (12)

From the Eqs. (6) and (12), we have got the relation between the neutrino arrival instant  $t_v$  as measured by the CERN reference frame,  $O_C x_C y_C z_C \equiv O xyz$ , and the neutrino arrival instant  $t_v^G$  as measured by the Gran Sasso reference frame,  $O_G x_G y_G z_G \equiv \tilde{O} \tilde{x} \tilde{y} \tilde{z}$ , at the exact location of the interaction at an interation location within the Gran Sasso block of detectors, provided the effect of the Earth daily rotation under the assumptions we are taking in relation to the intantaneous movements of these locations in relation to the fixed stars as previously discussed:

$$\frac{t_{\nu}^G}{t_{\nu}} = \frac{\lambda_G}{\lambda_C} \left[ 1 - (2\omega R \sin \alpha)^2 / c^2 \right]^{-1} = \gamma^2 \frac{\lambda_G}{\lambda_C}, \quad (13)$$

where  $\gamma \ge 1$  is the usual relativity factor as defined above.

Now,  $\lambda_G/\lambda_C$  is simply the ratio between the spatial displacement between our consecutive gedanken control pulses, being these displacements defined through our previous paragraphs, leading to the Eqs. (4) and (10). Of course, this ratio is simply given by the relativistic Doppler effect under an approximation case in which  $C_C$  is the source and  $C_G$  the detector. The ratio between the Eqs. (10) and (4) gives:

$$\frac{\lambda_G}{\lambda_C} = \left[1 - \left(2\omega R\sin\alpha\right)/c\right] \frac{\left(t_v^G - \delta t_v^G\right)}{\left(t_v - \delta t_v\right)}.$$
 (14)

But the time interval  $(t_v - \delta t_v)$  is a proper time interval measured by the source clock  $C_C$ , as previously discussed. It accounts for the time interval between the first pulse sent and the last pulse sent as locally described by  $C_C$  is its  $O_C x_C y_C z_C \equiv$ Oxyz world. These two events accur at different spatial locations in the  $C_G$  detector clock world  $O_G x_G y_G z_G \equiv \tilde{O} \tilde{x} \tilde{y} \tilde{z}$ , since  $C_C$  is approximating to  $C_G$  is this latter world. Hence,  $t_v - \delta t_v$  is the Lorentz time contraction of  $t_v^G - \delta t_v^G$ , viz.:

$$t_{\nu} - \delta t_{\nu} = \gamma^{-1} \left( t_{\nu}^{G} - \delta t_{\nu}^{G} \right) \quad \therefore$$
$$\frac{\left( t_{\nu}^{G} - \delta t_{\nu}^{G} \right)}{t_{\nu} - \delta t_{\nu}} = \gamma = \left[ 1 - \left( 2\omega R \sin \alpha \right)^{2} / c^{2} \right]^{-1/2}. \tag{15}$$

With the Eqs. (14) and (15), one reaches the usual relativistic Doppler effect expression for the approximation case:

$$\frac{\lambda_G}{\lambda_C} = \sqrt{\frac{1 - (2\omega R \sin \alpha) / c}{1 + (2\omega R \sin \alpha) / c}}.$$
 (16)

With the Eq. (16), the Eq. (13) reads:

$$\frac{t_{\gamma}^{G}}{t_{\gamma}} = \left[1 - (2\omega R \sin \alpha)^{2} / c^{2}\right]^{-1/2} \left[1 + (2\omega R \sin \alpha) / c\right]^{-1} = \frac{\gamma}{1 + (2\omega R \sin \alpha) / c}.$$
(17)

Since  $(2\omega R \sin \alpha)/c \ll 1$ , we may apply an approximation for the Eq. (17), viz.:

$$\gamma \approx 1 + \frac{1}{2} \frac{(2\omega R \sin \alpha)^2}{c^2},\tag{18}$$

and:

$$[1 + (2\omega R \sin \alpha) / c]^{-1} \approx 1 - (2\omega R \sin \alpha) / c, \qquad (19)$$

from which, neglecting the higher order terms, the Eq. (17) reads:

$$\frac{t_{\nu}^{G}}{t_{\nu}} \approx 1 - \frac{2\omega R \sin \alpha}{c} \quad \therefore \tag{20}$$

$$t_{\nu}^{G} - t_{\nu} = -\frac{2\omega R \sin \alpha}{c} t_{\nu}.$$
 (21)

From this result, the clock that tag the arrival interaction instant  $t_v^G$  in Gran Sasso turns out to measure an arrival time that is shorter than the correct one, this latter given by  $t_v$ . With the discrepancy,  $\epsilon$ , given by the value measured by the OPERA Collaboration [1], since  $t_v$  is simply given by  $L/v_v$ , where L is the baseline distance between the CERN and Gran Sasso,  $v_v$  the speed of neutrino through the Earth crust, one obtains a value for  $v_v$ . We rewrite the Eq. (21):

$$\epsilon = t_{\nu}^{G} - t_{\nu} = -\frac{2\omega R \sin \alpha}{c} \frac{L}{v_{\nu}}.$$
 (22)

With the values<sup>\*</sup>  $\omega = 7.3 \times 10^{-5} \text{ s}^{-1}$ ,  $R = R_E \cos \lambda \approx 6.4 \times 10^6 \text{ m} \times \cos (\pi/4) = 4.5 \times 10^6 \text{ m}$ ,  $\sin \alpha \approx \sin (7\pi/180) = 1.2 \times 10^{-1}$ ,  $c = 3.0 \times 10^8 \text{ ms}^{-1}$  and  $L = 7.3 \times 10^5 \text{ m}$ , also with the discrepancy  $\epsilon$ , given by the Eq. (22), being, say,  $\epsilon = -62 \times 10^{-9} \text{ s}$ , the neutrino velocity through the Earth crust reads:

$$v_{\nu} \approx 3.1 \times 10^6 \,\mathrm{ms}^{-1},$$
 (23)

being the refraction index of the Earth crust for neutrino given by:

$$n_{c|\nu} = \frac{c}{v_{\nu}} \approx 97.$$
(24)

In reference to the matching PDF (probability distribution function) in the OPERA experiment, one would have a discrepancy between the maximum likelihood distribution obtained from the block of detectors at Gran Sasso and the translation of the PDF due to the protons distribution by  $TOF_c$  given by, in virtue of the Eq. (22):

$$TOF_{\nu} = TOF_{c} + \epsilon = TOF_{c} - \frac{2\omega R \sin \alpha}{c} \frac{L}{v_{\nu}} \therefore$$
$$TOF_{\nu} - TOF_{c} \approx -62 \text{ ns}, \qquad (25)$$

under the reasoning and simplifications throughout this paper. One should notice the resoning here holds if the discrepancy turns out to be encrusted within the time translation of the PDF data, but such effect would not arise if the time interval TOF<sub> $\nu$ </sub> were directly measured, since, in this latter situation, such interval would only read  $L/v_{\nu}$ .

#### Conclusion

It is interesting to observe that even with a velocity having got two orders of magnitude lesser than *c* a neutrino may be interpreted as having got a velocity greater than *c*, depending on the method used to measure neutrino's time of flight, with the Earth crust presenting an index of refraction  $n_{c|v} > 1$ , due, also, to the local Doppler effect between the clocks attached to Gran Sasso and the respective intersecting ones in the CERN reference frame, as discussed throughout this paper, in virtue of the Earth daily rotation.

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<sup>\*</sup>See the Eqs. (2) and (3). The latitudes of CERN and Gran Sasso are, respectively:  $46^{deg}14^{min}3^{sec}(N)$  and  $42^{deg}28^{min}12^{sec}(N)$ . The longitudes of CERN and Gran Sasso are, respectively:  $6^{deg}3^{min}19^{sec}(E)$  and  $13^{deg}33^{min}0^{sec}(E)$ .