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Extinction probabilities for a distylous plant population modeled by an inhomogeneous random walk on the positive quadrant

Pauline Lafitte-Godillon*, Kilian Raschel†, Viet Chi Tran‡

January 30, 2012

Abstract

In this paper, we study a distylous flower population in which self-reproduction is not permitted. Individuals are diploid, and two alleles, A and a , can be found at the considered locus S . Pollen and ovules of flowers with the same genotype at locus S cannot mate. This prevents the pollen of a given flower to fecundate its stigmates. Only genotypes AA and Aa can be maintained in the population, so that the latter can be described by a random walk in the positive quadrant whose components are the number of individuals of each genotype. This random walk is not homogeneous and its transitions depend on the location of the process. We are interested in the computation of the extinction probabilities, where extinction happens when one of the axis is reached by the process. These extinction probabilities, which depend on the initial condition, satisfy a doubly-indexed recurrence equation that cannot be solved directly. We consider the associated generating function and show that it satisfies a partial differential equation that is solved but whose solution is explicit, though hardly tractable. Numerical results comparing stochastic and deterministic approximations of the extinction probabilities are studied.

Keywords: Inhomogeneous random walk on the positive quadrant; boundary absorption; transport equation; method of characteristics; self-incompatibility in flower populations; extinction in diploid population with sexual reproduction

AMS codes: 60G50; 60J80; 35Q92; 92D25

1 Introduction

We consider the model of distyle plants with sporophytic reproduction and absence of pollen limitation introduced in Billiard and Tran [4]. The reproduction is sexual and self-incompatible. Each plant is diploid and characterized by the two alleles that it carries at the locus S , which encode the recognition proteins present on the pollen and stigmata of the plant, and hence the possible types of partners with whom the plant may reproduce. We consider the case where only two alleles are available, A or a . The plants thus have genotypes AA , Aa or aa . The only interesting case is when A is dominant over a (see [4]), and we restrict to this case in this work. The phenotype, *i.e.* the type of proteins carried by the pollen and stigmata, of individuals with genotypes AA (resp. Aa and aa) is A (resp.

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A and a). The self-incompatibility mechanism prevents fecundation of the plant's ovule by its own pollen: only pollen and stigmata with different proteins can give viable seeds, *i.e.* pollen of a plant of phenotype A can only fecundate stigmata of a plant of phenotype a and vice-versa. It can be seen that seeds AA cannot be created, since the genotype of individuals of phenotype a is necessarily aa that combine only with individuals of phenotype A that have genotypes AA or Aa , therefore we can consider without restriction populations consisting only of individuals of genotypes Aa and aa . Each viable seed is then necessarily of genotype Aa or aa with probability $1/2$. It is assumed that ovules are produced in continuous time at rate $r > 0$ and that each ovule is fecundated to give a seed, provided there exists compatible pollen in the population. The lifetime of each individual follows an exponential distribution with mean $1/d$, where $d > 0$. In all the article, we consider

$$r > d \tag{1.1}$$

which, we will see, is the interesting case.

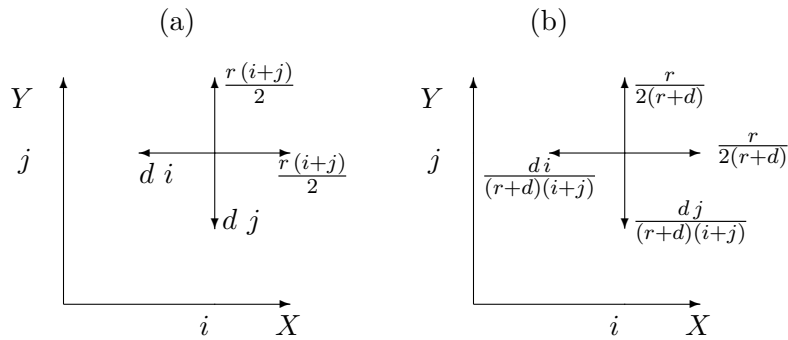


Figure 1: (a) Transition rates for the continuous-time pure-jump Markov process $(X_t, Y_t)_{t \in \mathbb{R}_+}$. (b) Transition probabilities of the embedded random walk, that we note $(X_t, Y_t)_{t \in \mathbb{N}}$ (here and throughout, we denote by \mathbb{N} the set $\{0, 1, 2, \dots\}$), with an abuse of notation.

Let us denote by X_t and Y_t the number of individuals of genotype Aa (phenotype A) and aa (phenotype a) at time $t \in \mathbb{R}_+$. The process $(X_t, Y_t)_{t \in \mathbb{R}_+}$ is a pure-jump Markov process with transitions represented in Fig. 1(a). A stochastic differential equation (SDE) representation of $(X_t, Y_t)_{t \in \mathbb{R}_+}$ is given in [4]. Here we forget the continuous-time process, and we are interested in the embedded discrete-time Markov chain, which we denote, with an abuse of notation, by $(X_t, Y_t)_{t \in \mathbb{N}}$, and with transitions represented in Fig. 1(b):

$$\begin{aligned} \mathbb{P}_{i,j}[(X_1, Y_1) = (i-1, j)] &= \frac{di}{(r+d)(i+j)}, & \mathbb{P}_{i,j}[(X_1, Y_1) = (i+1, j)] &= \frac{r}{2(r+d)}, \\ \mathbb{P}_{i,j}[(X_1, Y_1) = (i, j-1)] &= \frac{dj}{(r+d)(i+j)}, & \mathbb{P}_{i,j}[(X_1, Y_1) = (i, j+1)] &= \frac{r}{2(r+d)}, \end{aligned}$$

where $\mathbb{P}_{i,j}$ means that the process starts with the initial condition $(X_0, Y_0) = (i, j)$. The main difficulty is that this random walk is not homogeneous in space, while techniques developed in the literature for random walks on positive quadrants mostly focus on the homogeneous case (see *e.g.* Fayolle *et al.* [6], Klein Haneveld and Pittenger [8], Kurkova and Raschel [9], Walraevens, van Leeuwen and Boxma [12]). We introduce a generating function (1.4) that satisfies here a partial differential equation (PDE) of a new type that we solve. Although the particularity of the problem is exploited, these techniques and the links between probability and PDE may be extended to carry out general study of inhomogeneous random walks in cones. The introduction of PDEs through generating functions had been already used by Feller [7] for a trunking problem with an inhomogeneous random walk in

dimension 1. To our knowledge, the case of inhomogeneous random walks in the cone with absorbing boundaries has been left open.

When one of the phenotype A or a disappears, reproduction becomes impossible and the system gets extinct. We are interested in the probability of extinction of $(X_t, Y_t)_{t \in \mathbb{N}}$ (or, equivalently, in that of $(X_t, Y_t)_{t \in \mathbb{R}_+}$). Let us introduce the first time at which one of the two types gets extinct:

$$\tau_0 = \inf\{t \in \mathbb{N} : X_t = 0 \text{ or } Y_t = 0\}. \quad (1.2)$$

For $i, j \in \mathbb{N}$, let us denote by

$$p_{i,j} = \mathbb{P}_{i,j}[\tau_0 < \infty] \quad (1.3)$$

the absorption probabilities, and by

$$P(x, y) = \sum_{i, j \geq 1} p_{i,j} x^i y^j \quad (1.4)$$

their generating function. By symmetry arguments, we have, for all $i, j \in \mathbb{N}$,

$$p_{i,j} = p_{j,i}. \quad (1.5)$$

Moreover, for any $i, j \in \mathbb{N}$ such that $i = 0$ or $j = 0$, we have

$$p_{i,j} = 1. \quad (1.6)$$

In Section 2, we will see that the $p_{i,j}$'s satisfy the Dirichlet problem associated with the following doubly-indexed recurrence equation

$$q_{i,j} = \frac{di}{(r+d)(i+j)} q_{i-1,j} + \frac{dj}{(r+d)(i+j)} q_{i,j-1} + \frac{r}{2(r+d)} q_{i,j+1} + \frac{r}{2(r+d)} q_{i+1,j} \quad (1.7)$$

and with the boundary condition (1.6). This problem does not admit simple solutions. There is no uniqueness of solutions to this problem. Note that the constant sequence equal to 1 is a solution. However, we are interested in solutions that tend to 0 as i or j tends to infinity, since, [4] (see Prop. 2.2 in this paper), estimates for $p_{i,j}$ were obtained through probabilistic coupling techniques; they show that in the case (1.1) we consider, $p_{i,j}$ is strictly less than 1. In fact, the $p_{i,j}$'s correspond to the smallest positive solution of the Dirichlet problem, and are completely determined if we give the probabilities $(p_{i,1})_{i \geq 1}$. We conclude the section with more precise estimates of the absorption probabilities $p_{i,j}$ as the initial state (i, j) goes to infinity along one axis (Prop. 2.3). These new estimates rely on Prop. 2.2 and on comparisons with one-dimensional random walks. In Section 3, we consider the generating function $P(x, y)$ associated with the $p_{i,j}$'s and show that it satisfies a PDE, that has one and only one solution, that is computed (Prop. 3.5) explicitly with a dependence on the $(p_{i,1})_{i \geq 1}$, prompting us to use the name ‘‘Green’s function’’. This provides a new formulation of the solution of (1.7), that is however uneasy to work with numerically. Hence, in Section 4, we propose two different approaches leading to numerical approximations of the solution of the Dirichlet problem (1.7)–(1.6), that are based on stochastic and deterministic approaches.

2 Existence of a solution

2.1 Dirichlet problem

We first establish that the extinction probabilities $p_{i,j}$'s (1.3) solve the Dirichlet problem (1.6)–(1.7).

Proposition 2.1. (i) *The extinction probabilities $(p_{i,j})_{i,j \geq 1}$ are solutions to the Dirichlet problem (1.7) with boundary condition (1.6). Uniqueness of the solution may not hold, but the extinction probabilities $(p_{i,j})_{i,j \in \mathbb{N}}$ define the smallest positive solution to this problem.*

(ii) Assume that the probabilities $(p_{i,1})_{i \geq 1}$ are given. Then the probabilities $(p_{i,j})_{i,j \geq 1}$ are completely determined.

Proof. We begin with Point (i). Equation (1.7) is obtained by using the strong Markov property at the time of the first event. Let us denote by K the transition kernel of the discrete-time Markov chain $(X_t, Y_t)_{t \in \mathbb{N}^*}$; we have:

$$Kf(i, j) = (f(i+1, j) + f(i, j+1)) \frac{r}{2(r+d)} + f(i-1, j) \frac{di}{(r+d)(i+j)} + f(i, j-1) \frac{dj}{(r+d)(i+j)}.$$

Following classical proofs (e.g. [3, 10]), the extinction probabilities $(p_{i,j})_{i,j \in \mathbb{N}}$ satisfy the equation:

$$\forall i, j \in \mathbb{N}^*, \quad f(i, j) = Kf(i, j) \quad \text{and} \quad \forall i, j \in \mathbb{N}, \quad f(i, 0) = f(0, j) = 1. \quad (2.1)$$

The constant solution equal to 1 is a solution to (2.1). Let us prove that $(p_{i,j})_{i,j \in \mathbb{N}}$ is the smallest positive solution to (2.1). Let f be another positive solution. Let us consider $M_t = f(X_{\inf\{t, \tau_0\}}, Y_{\inf\{t, \tau_0\}})$, with τ_0 defined in (1.2). Denoting by $(\mathcal{G}_t)_{t \in \mathbb{N}}$ the filtration of $(M_t)_{t \in \mathbb{N}}$, we have:

$$\begin{aligned} \mathbb{E}[M_{t+1} | \mathcal{G}_t] &= \mathbb{E}[M_{t+1} \mathbb{1}_{\tau_0 \leq t} + M_{t+1} \mathbb{1}_{\tau_0 > t} | \mathcal{G}_t] \\ &= \mathbb{E}[M_t \mathbb{1}_{\tau_0 \leq t} + f(X_{t+1}, Y_{t+1}) \mathbb{1}_{\tau_0 > t} | \mathcal{G}_t] \\ &= M_t \mathbb{1}_{\tau_0 \leq t} + \mathbb{1}_{\tau_0 > t} \mathbb{E}[f(X_{t+1}, Y_{t+1}) | \mathcal{G}_t] \\ &= M_t \mathbb{1}_{\tau_0 \leq t} + \mathbb{1}_{\tau_0 > t} Kf(X_t, Y_t) \\ &= M_t \mathbb{1}_{\tau_0 \leq t} + \mathbb{1}_{\tau_0 > t} f(X_t, Y_t) = M_t. \end{aligned}$$

Hence $(M_t)_{t \in \mathbb{N}}$ is a martingale, which converges on $\{\tau_0 < \infty\}$ to $f(X_{\tau_0}, Y_{\tau_0}) = 1$. Thus by using the positivity of f and Fatou's lemma, we obtain that for every $i, j \in \mathbb{N}$:

$$f(i, j) = \mathbb{E}_{i,j}[M_0] = \lim_{t \rightarrow \infty} \mathbb{E}_{i,j}[M_t] \geq \mathbb{E}[\liminf_{t \rightarrow \infty} M_t \mathbb{1}_{\tau_0 < \infty}] = \mathbb{E}_{i,j}[\mathbb{1}_{\tau_0 < \infty}] = p_{i,j}.$$

This achieves the proof of Point (i).

Let us now consider Point (ii). Assume that the probabilities $(p_{i,1})_{i \geq 1}$ are given, and let us prove, by recursion, that every $p_{i,j}$ can be computed. By symmetry, we only need to prove that this is the case for $i \geq j$. Assume

(Hrec j): for $j \in \mathbb{N}^*$ all the probabilities $p_{k,\ell}$ for $\ell \leq j$ and $k \geq \ell$ can be computed from the $p_{i,1}$'s and let us prove that we can determine the probabilities $p_{i,j+1}$ for $i \geq j+1$. From (1.7) we get that:

$$p_{i,j+1} = \frac{2(r+d)}{r} p_{i,j} - \frac{2di}{r(i+j)} p_{i-1,j} - \frac{2dj}{r(i+j)} p_{i,j-1} - p_{i+1,j}. \quad (2.2)$$

All the terms in the r.h.s. of (2.2) are known by (Hrec j), and hence $p_{i,j+1}$ can be computed for any $i \geq j+1$. This achieves the recursion. \blacksquare

The following result shows that there is almost sure extinction in the case $r \leq d$. In the interesting case $r > d$, it also shows that there is a nontrivial solution to the Dirichlet problem (1.6)–(1.7).

Proposition 2.2 (Prop. 9 of [4]). *We have the following regimes given the parameters r and d :*

- (i) If $r \leq d$, we have almost sure extinction of the population.
- (ii) If $r > d (> 0)$, then there is a strictly positive survival probability. Denoting by (i, j) the initial condition, we have:

$$\left(\frac{d}{r}\right)^{i+j} \leq p_{i,j} \leq \left(\frac{d}{r}\right)^i + \left(\frac{d}{r}\right)^j - \left(\frac{d}{r}\right)^{i+j}.$$

In Point (ii), only bounds, and no explicit formula, are available for the extinction probability $p_{i,j}$. The purpose of this article is to address (1.7) by considering the Green's function $P(x, y)$ introduced in (1.4).

2.2 Asymptotic behavior of the absorption probability as the initial state goes to infinity along one axis

In this part, using the result of Prop. 2.2, we provide more precise estimates of the asymptotic behavior of the absorption probability $p_{1,j} = p_{j,1}$ when $j \rightarrow \infty$. In particular, these estimates will be very useful when we tackle the deterministic numerical simulations (see Section 4.2).

Proposition 2.3. *If $j \rightarrow \infty$, then*

$$p_{1,j} = p_{j,1} = \frac{2d}{r} \frac{1}{j} - \frac{2d(r^2 + dr + 2d^2)}{r^2(r+d)} \frac{1}{j^2} + O\left(\frac{1}{j^3}\right). \quad (2.3)$$

Proof. In addition of τ_0 , defined in (1.2), we introduce

$$S = \inf\{t \in \mathbb{N} : Y_t = 0\}, \quad T = \inf\{t \in \mathbb{N} : X_t = 0\},$$

the hitting times of the horizontal axis and vertical axis, respectively. Note that we have $\tau_0 = \inf\{S, T\}$. Let $f : \mathbb{N} \rightarrow \mathbb{N}$ be a function such that $f(j) < j$ for any $j \geq 1$. In the sequel, we will choose $f(j) = \lfloor \epsilon j \rfloor$, with $\epsilon \in (0, 1)$ and where $\lfloor \cdot \rfloor$ denotes the integer part. We obviously have the identity:

$$p_{1,j} = \mathbb{P}_{(1,j)}[\tau_0 < \infty] = \mathbb{P}_{(1,j)}[\tau_0 \leq f(j)] + \mathbb{P}_{(1,j)}[f(j) < \tau_0 < \infty]. \quad (2.4)$$

To prove Prop. 2.3, we shall give estimates for both terms present in the right-hand side of (2.4).

First step: Study of $\mathbb{P}_{(1,j)}[\tau_0 \leq f(j)]$. Since $f(j) < j$, it is impossible, starting from $(1, j)$, to reach the horizontal axis before time $f(j)$, and we have $\mathbb{P}_{(1,j)}[\tau_0 \leq f(j)] = \mathbb{P}_{(1,j)}[T \leq f(j)]$. In order to compute the latter probability, we introduce two one-dimensional random walks on \mathbb{N} , namely X^- and X^+ , which are killed at 0, and which have the jumps

$$\mathbb{P}_i[X_1^\pm = i - 1] = q_i^\pm, \quad \mathbb{P}_i[X_1^\pm = i + 1] = p_i^\pm, \quad \mathbb{P}_i[X_1^\pm = i] = r_i^\pm, \quad q_i^\pm + p_i^\pm + r_i^\pm = 1,$$

where

$$q_i^\pm = \frac{di}{(r+d)(i+j \mp f(j))}, \quad p_i^\pm = \frac{r}{2(r+d)}. \quad (2.5)$$

Both X^- and X^+ are (inhomogeneous) birth-and-death processes on \mathbb{N} . Although it is implicit, these random walks are parameterized by j . If $T^\pm = \inf\{t \in \mathbb{N} : X_t^\pm = 0\}$, then

$$\mathbb{P}_1[T^- \leq f(j)] \leq \mathbb{P}_{(1,j)}[T \leq f(j)] \leq \mathbb{P}_1[T^+ \leq f(j)]. \quad (2.6)$$

The quantities $\mathbb{P}_1[T^\pm \leq f(j)]$ are computable: we shall prove that

$$\mathbb{P}_1[T^\pm \leq f(j)] = \frac{2d}{r} \frac{1}{(j \mp f(j))} - \frac{2d(r^2 + dr + 2d^2)}{r^2(r+d)} \frac{1}{(j \mp f(j))^2} + O\left(\frac{1}{(j \mp f(j))^3}\right). \quad (2.7)$$

The main idea for proving (2.7) is that the q_i^\pm being very small as $j \rightarrow \infty$, the only paths which will significantly contribute to the probability $\mathbb{P}_1[T^\pm \leq f(j)]$ are the ones with very few jumps to the right. Let us define

$$\begin{aligned} \Lambda_p^t &= \{\text{the chain } X^\pm \text{ makes exactly } p \text{ jumps to the right between } 0 \text{ and } t\} \\ &= \{\text{there exist } 0 \leq q_1 < \dots < q_p \leq t-1 \text{ such that } X_{q_1+1}^\pm - X_{q_1}^\pm = \dots = X_{q_p+1}^\pm - X_{q_p}^\pm = 1\}. \end{aligned}$$

We are entitled to write

$$\mathbb{P}_1[T^\pm \leq f(j)] = \mathbb{P}_1[T^\pm \leq f(j), \Lambda_0^{f(j)}] + \mathbb{P}_1[T^\pm \leq f(j), \Lambda_1^{f(j)}] + \mathbb{P}_1[T^\pm \leq f(j), \cup_{p \geq 2} \Lambda_p^{f(j)}], \quad (2.8)$$

and we now separately analyze the three terms in the right-hand side of (2.8). First:

$$\begin{aligned}\mathbb{P}_1[T^\pm \leq f(j), \Lambda_0^{f(j)}] &= \sum_{k=1}^{f(j)} \mathbb{P}_1[T^\pm = k, \Lambda_0^{f(j)}] = \sum_{k=1}^{f(j)} (r_1^\pm)^{k-1} q_1^\pm \\ &= \frac{q_1^\pm}{1 - r_1^\pm} (1 - (r_1^\pm)^{f(j)}) = \frac{q_1^\pm}{1 - r_1^\pm} (1 + O((r_1^\pm)^{f(j)})).\end{aligned}\quad (2.9)$$

A Taylor expansion of $q_1^\pm/(1 - r_1^\pm)$ according to the powers of $1/(j \mp f(j))$ together with the fact that $(r_1^\pm)^{f(j)} = o(1/(j \mp f(j))^3)$ provides that:

$$\mathbb{P}_1[T^\pm \leq f(j), \Lambda_0^{f(j)}] = \frac{2d}{r(j \mp f(j))} \left(1 - \frac{1 + \frac{2d}{r}}{j \mp f(j)} + O\left(\frac{1}{(j \mp f(j))^2}\right) \right). \quad (2.10)$$

We now consider the second term in the right-hand side of (2.8). On the event $\Lambda_1^{f(j)}$, X^\pm first stays a time k_1 at 1, then jumps to 2, where it remains k_2 unit of times; it next goes to 1, and, after a time k_3 , jumps to 0. Further, since $T^\pm \leq f(j)$, we have $k_1 + 1 + k_2 + 1 + k_3 + 1 \leq f(j)$. Denoting by $\tilde{k}_1 = k_1 + k_3$ the time spent in position 1, we thus have:

$$\begin{aligned}\mathbb{P}_1[T^\pm \leq f(j), \Lambda_1^{f(j)}] &= \sum_{k_1+k_2+k_3 \leq f(j)-3} (r_1^\pm)^{k_1} p_1^\pm (r_2^\pm)^{k_2} q_2^\pm (r_1^\pm)^{k_3} q_1^\pm = p_1^\pm q_1^\pm q_2^\pm \sum_{\tilde{k}_1+k_2 \leq f(j)-3} (r_1^\pm)^{\tilde{k}_1} (r_2^\pm)^{k_2} \\ &= \frac{p_1^\pm q_1^\pm q_2^\pm}{(1 - r_1^\pm)(1 - r_2^\pm)} (1 + O((r_1^\pm \vee r_2^\pm)^{f(j)-2})).\end{aligned}\quad (2.11)$$

As for the first term, using the fact that $(r_1^\pm \vee r_2^\pm)^{f(j)-2} = o(1/(j \mp f(j))^3)$ and a Taylor expansion according to the powers of $1/(j \mp f(j))$ gives that:

$$\mathbb{P}_1[T^\pm \leq f(j), \Lambda_1^{f(j)}] = \frac{4d^2}{r(r+d)(j \mp f(j))^2} \left(1 + O\left(\frac{1}{(j \mp f(j))}\right) \right). \quad (2.12)$$

Finally, let us consider the third term $\mathbb{P}_1[T^\pm \leq f(j), \cup_{p \geq 2} \Lambda_p^{f(j)}]$. On $\cup_{p \geq 2} \Lambda_p^{f(j)}$, the two first jumps to the right are either from 1 to 2 and 2 to 3, or twice from 1 to 2. Thus, extinction means that there is at least 3 jumps from 3 to 2, 2 to 1 and 1 to 0 or two jumps from 2 to 1 and 1 to 0. Since q_i^\pm is an increasing function of i , we deduce that:

$$\mathbb{P}_1[T^\pm \leq f(j), \cup_{p \geq 2} \Lambda_p^{f(j)}] \leq (q_3^\pm)^3 = \left(\frac{d}{r+d}\right)^3 \left(\frac{3}{3+j \mp f(j)}\right)^3 = O\left(\frac{1}{(j \mp f(j))^3}\right). \quad (2.13)$$

From (2.8), (2.9), (2.11) and (2.13), we obtain (2.7).

Second step: Study of $\mathbb{P}_{(1,j)}[f(j) < \tau_0 < \infty]$.

$$\begin{aligned}\mathbb{P}_{(1,j)}[f(j) < \tau_0 < \infty] &= \sum_{k, \ell \geq 1} \mathbb{P}_{(1,j)}[f(j) < \tau_0 < \infty | (X_{f(j)}, Y_{f(j)}) = (k, \ell)] \mathbb{P}_{(1,j)}[(X_{f(j)}, Y_{f(j)}) = (k, \ell)] \\ &= \sum_{k, \ell \geq 1} p_{k, \ell} \mathbb{P}_{(1,j)}[(X_{f(j)}, Y_{f(j)}) = (k, \ell)],\end{aligned}\quad (2.14)$$

by using the strong Markov property. Introduce now a function $g : \mathbb{N} \rightarrow \mathbb{N}$ such that $f(j) + g(j) < j$ for any $j \geq 1$. We can split (2.14) into

$$\sum_{k, \ell \geq g(j)} p_{k, \ell} \mathbb{P}_{(1,j)}[(X_{f(j)}, Y_{f(j)}) = (k, \ell)] + \sum_{\substack{g(j) \geq k \geq 1 \\ \text{and/or} \\ g(j) \geq \ell \geq 1}} p_{k, \ell} \mathbb{P}_{(1,j)}[(X_{f(j)}, Y_{f(j)}) = (k, \ell)]. \quad (2.15)$$

With Prop. 2.2 we obtain the following upper bound for the first sum in (2.15):

$$\sum_{k, \ell \geq g(j)} p_{k, \ell} \mathbb{P}_{(1, j)}[(X_{f(j)}, Y_{f(j)}) = (k, \ell)] \leq 2 \left(\frac{d}{r} \right)^{g(j)}.$$

In particular, if we choose g such that as $j \rightarrow \infty$, $g(j) \rightarrow \infty$ fast enough, then clearly the term above is negligible w.r.t. (2.7). For the second sum in (2.15), we first note that

$$\sum_{\substack{g(j) \geq k \geq 1 \\ \text{and/or} \\ g(j) \geq \ell \geq 1}} p_{k, \ell} \mathbb{P}_{(1, j)}[(X_{f(j)}, Y_{f(j)}) = (k, \ell)] = \sum_{g(j) \geq k \geq 1} p_{k, \ell} \mathbb{P}_{(1, j)}[(X_{f(j)}, Y_{f(j)}) = (k, \ell)],$$

since by assumption $j - f(j) > g(j)$ so that Y cannot reach values $\ell \leq g(j)$ in $f(j)$ steps. Then we have

$$\begin{aligned} \sum_{g(j) \geq k \geq 1} p_{k, \ell} \mathbb{P}_{(1, j)}[(X_{f(j)}, Y_{f(j)}) = (k, \ell)] &\leq \sum_{g(j) \geq k \geq 1} \mathbb{P}_{(1, j)}[(X_{f(j)}, Y_{f(j)}) = (k, \ell)] \\ &\leq \mathbb{P}_{(1, j)}[0 \leq X_{f(j)} \leq g(j)]. \end{aligned} \quad (2.16)$$

To obtain an upper bound for (2.16) we are going to use, again, a one-dimensional random walk. Introduce \tilde{X} , a random walk on \mathbb{N} which is killed at 0, homogeneous on \mathbb{N}^* with jumps

$$\mathbb{P}_k[\tilde{X}_1 = k - 1] = \tilde{q}, \quad \mathbb{P}_k[\tilde{X}_1 = k + 1] = \tilde{p}, \quad \mathbb{P}_k[\tilde{X}_1 = k] = \tilde{r}, \quad \tilde{q} + \tilde{p} + \tilde{r} = 1,$$

where

$$\tilde{q} = \frac{d(1 + f(j))}{(r + d)(1 + j - 2f(j))}, \quad \tilde{p} = \frac{r}{2(r + d)}.$$

This walk is again parameterized by j . By construction of \tilde{X} , we have

$$\mathbb{P}_{(1, j)}[0 \leq X_{f(j)} \leq g(j)] \leq \mathbb{P}_1[\tilde{X}_{f(j)} \leq g(j)].$$

Denoting by \tilde{m} and $\tilde{\sigma}^2$ the mean and the variance of $(\tilde{X}_2 - \tilde{X}_1)$, respectively (they could easily be computed), we can write

$$\mathbb{P}_1[\tilde{X}_{f(j)} \leq g(j)] = \mathbb{P}_0 \left[\frac{\tilde{X}_{f(j)} - \tilde{m}f(j)}{\tilde{\sigma}\sqrt{f(j)}} \leq \frac{g(j) - 1 - \tilde{m}f(j)}{\tilde{\sigma}\sqrt{f(j)}} \right].$$

By a suitable choice of the functions f and g , for instance $f(j) = \lfloor \epsilon j \rfloor$ with $\epsilon \in (0, 1)$ and $g(j) = \lfloor j^{3/4} \rfloor$, the central limit theorem gives that the latter is negligible w.r.t. $\mathbb{P}_{(1, j)}[\tau_0 \leq f(j)]$. For this last term, using (2.7), (2.6) and letting ϵ tend to 0 provides (2.3). The proof is concluded. \blacksquare

Let us do some remarks on possible extensions of Prop. 2.3.

Remark 2.4. 1. The proof of Prop. 2.3 can easily be extended to the asymptotic of $p_{i, j}$ as $j \rightarrow \infty$, for any fixed value of i . In particular, we have the following asymptotic behavior:

$$p_{i, j} = \left(\frac{2d}{r} \right)^i \frac{i!}{j^i} + O \left(\frac{1}{j^{i+1}} \right). \quad (2.17)$$

2. It is possible to generalize (2.8) by

$$\mathbb{P}_1[T^\pm \leq f(j)] = \sum_{p=0}^{k-1} \mathbb{P}_1[T^\pm \leq f(j), \Lambda_p^{f(j)}] + \mathbb{P}_1[T^\pm \leq f(j), \cup_{p \geq k} \Lambda_p^{f(j)}]. \quad (2.18)$$

We can show as in the proof of Prop. 2.3 that $\mathbb{P}_1[T^\pm \leq f(j), \cup_{p \geq k} \Lambda_p^{f(j)}] \leq (q_{k+1}^\pm)^{k+1} = o(1/j^{k+1})$ (see (2.13)) and that the probabilities $\mathbb{P}_1[T^\pm \leq f(j), \Lambda_p^{f(j)}]$ admit Taylor expansions in powers of $1/(j \mp f(j))$ where the development for the p^{th} probability has a main term in $1/(j \mp f(j))^{p+1}$. This allows us to push the developments in (2.7) to higher orders.

For instance, the next term in (2.3) can be obtained by a long computation. First, we generalize (2.11) by writing that

$$\mathbb{P}_1[T^\pm \leq f(j), \Lambda_2^{f(j)}] = q_1^\pm ((p_1^\pm)^2 (q_2^\pm)^2 + p_1^\pm p_2^\pm q_3^\pm q_2^\pm) \sum_{\tilde{k}_1 + \tilde{k}_2 + \tilde{k}_3 \leq f(j) - 5} (r_1^\pm)^{\tilde{k}_1} (r_2^\pm)^{\tilde{k}_2} (r_3^\pm)^{\tilde{k}_3} \quad (2.19)$$

where \tilde{k}_1, \tilde{k}_2 and \tilde{k}_3 are the times spent by the random walk in the states 1, 2 and 3. Then, pushing further the Taylor expansion leads to:

$$p_{1,j} = \frac{2d}{r} \frac{1}{j \mp f(j)} - \frac{2d(r^2 + dr + 2d^2)}{r^2(r+d)} \frac{1}{(j \mp f(j))^2} + d \left(\frac{2}{r} \left(1 + \frac{2d}{r}\right)^2 - \frac{24d(\frac{r}{2} + d)}{r^2(r+d)} + \frac{5r^2 d^2}{2(r+d)^2(\frac{r}{2} + d)^3} \right) \frac{1}{(j \mp f(j))^3} + O\left(\frac{1}{(j \mp f(j))^4}\right).$$

3 Green's function

3.1 A functional equation for the Green's function

In this section, we consider the Green's function $P(x, y)$ defined in (1.4) associated with a solution of (1.7) in the same spirit as what can be found in Feller [7, Ch. XVII]. We show that it satisfies a non-classical linear PDE that can be solved (see Prop. 3.5).

Proposition 3.1. (i) *The function $P(x, y)$ satisfies formally:*

$$AP(x, y) = h(x, y, P), \quad (3.1)$$

where:

$$\begin{aligned} AP(x, y) &= Q(x, y) \frac{\partial P}{\partial x}(x, y) + Q(y, x) \frac{\partial P}{\partial y}(x, y) + R(x, y)P(x, y), \\ Q(x, y) &= (r+d)x - \frac{r}{2} - \frac{r}{2} \frac{x}{y} - dx^2, \\ R(x, y) &= \frac{r}{2x} + \frac{r}{2y} - dx - dy, \end{aligned} \quad (3.2)$$

and where:

$$h(x, y, P) = -\frac{r}{2} \left(x \frac{\partial^2 P}{\partial x \partial y}(x, 0) + y \frac{\partial^2 P}{\partial y \partial x}(0, y) \right) + dxy \left(\frac{1}{1-x} + \frac{1}{1-y} \right). \quad (3.3)$$

(ii) *For given $(p_{i,1})_{i \geq 1}$, we have a unique classical solution to (3.1)-(3.3) defined in $]0, 1[\times]0, 1[$.*

The function h in (3.3) only depends on a boundary condition ($\partial^2 P / \partial x \partial y$ at the boundaries $x = 0$ or $y = 0$, i.e. the $p_{i,1}$'s for $i \in \mathbb{N}^*$), which is non-classical, while the operator A is of first order and hence associated with some transport equations.

Proof of Prop. 3.1. Let us first establish (i). Using the Markov property at time $t = 1$:

$$p_{i,j} = \frac{r}{2(r+d)} (p_{i+1,j} + p_{i,j+1}) + \frac{dj}{(r+d)(i+j)} p_{i,j-1} + \frac{di}{(r+d)(i+j)} p_{i-1,j},$$

then multiplying by $x^i y^j$, and suming over $i, j \in \mathbb{N}^*$ leads to:

$$(r+d) \sum_{i,j \geq 1} (i+j) p_{i,j} x^i y^j = \frac{r}{2} \sum_{i,j \geq 1} (i+j) (p_{i+1,j} + p_{i,j+1}) x^i y^j + d \sum_{i,j \geq 1} j p_{i,j-1} x^i y^j + d \sum_{i,j \geq 1} i p_{i-1,j} x^i y^j. \quad (3.4)$$

The l.h.s. of (3.4) equals:

$$(r+d) \sum_{i,j \geq 1} (i+j) p_{i,j} x^i y^j = (r+d) \left(x \frac{\partial}{\partial x} + y \frac{\partial}{\partial y} \right) P(x, y). \quad (3.5)$$

For the r.h.s. of (3.4):

$$\begin{aligned} & \frac{r}{2} \sum_{\substack{j \geq 1 \\ i \geq 2}} p_{i,j} (i-1+j) x^{i-1} y^j + \frac{r}{2} \sum_{\substack{i \geq 1 \\ j \geq 2}} p_{i,j} (i+j-1) x^i y^{j-1} + d \sum_{\substack{i \geq 1 \\ j \geq 0}} (j+1) p_{i,j} x^i y^{j+1} + d \sum_{\substack{i \geq 0 \\ j \geq 1}} (i+1) p_{i,j} x^{i+1} y^j \\ &= \frac{r}{2} \sum_{\substack{j \geq 1 \\ i \geq 2}} p_{i,j} (i-1) x^{i-1} y^j + \frac{r}{2} \sum_{\substack{j \geq 1 \\ i \geq 2}} p_{i,j} j x^{i-1} y^j + \frac{r}{2} \sum_{\substack{i \geq 1 \\ j \geq 2}} p_{i,j} i x^i y^{j-1} + \frac{r}{2} \sum_{\substack{i \geq 1 \\ j \geq 2}} p_{i,j} (j-1) x^i y^{j-1} \\ &+ d \sum_{\substack{i \geq 1 \\ j \geq 0}} (j+1) p_{i,j} x^i y^{j+1} + d \sum_{\substack{i \geq 0 \\ j \geq 1}} (i+1) p_{i,j} x^{i+1} y^j. \end{aligned} \quad (3.6)$$

For the first term in the r.h.s. of (3.6):

$$\begin{aligned} \frac{r}{2} \sum_{\substack{j \geq 1 \\ i \geq 2}} p_{i,j} (i-1) x^{i-1} y^j &= \frac{r}{2} \sum_{\substack{j \geq 1 \\ i \geq 1}} p_{i,j} (i-1) x^{i-1} y^j = \frac{r}{2} \sum_{\substack{j \geq 1 \\ i \geq 1}} p_{i,j} i x^{i-1} y^j - \frac{r}{2} \sum_{\substack{j \geq 1 \\ i \geq 1}} p_{i,j} x^{i-1} y^j \\ &= \frac{r}{2} \frac{\partial}{\partial x} P(x, y) - \frac{r}{2} \frac{P(x, y)}{x}. \end{aligned} \quad (3.7)$$

Similar computation can be done for the fourth term of the r.h.s. of (3.6). For the second term:

$$\begin{aligned} \frac{r}{2} \sum_{\substack{j \geq 1 \\ i \geq 2}} p_{i,j} j x^{i-1} y^j &= \frac{r}{2} \frac{y}{x} \sum_{\substack{j \geq 1 \\ i \geq 2}} p_{i,j} j x^i y^{j-1} = \frac{r}{2} \frac{y}{x} \left(\sum_{\substack{j \geq 1 \\ i \geq 1}} p_{i,j} j x^i y^{j-1} - \sum_{j \geq 1} p_{1,j} j x y^{j-1} \right) \\ &= \frac{r}{2} \frac{y}{x} \frac{\partial}{\partial y} P(x, y) - \frac{r}{2} \sum_{j \geq 1} p_{1,j} j y^j. \end{aligned} \quad (3.8)$$

Similar computation holds for the third term of the r.h.s. of (3.6). Now for the fifth term:

$$d \sum_{\substack{i \geq 1 \\ j \geq 0}} p_{i,j} (j+1) x^i y^{j+1} = d \left(y^2 \frac{\partial}{\partial y} P(x, y) + y P(x, y) + \sum_{i \geq 1} x^i y \right). \quad (3.9)$$

Similar computation holds for the last term of (3.6). From (3.5), (3.6), (3.7), (3.8) and (3.9) we deduce that:

$$\begin{aligned} (r+d) \left(x \frac{\partial}{\partial x} + y \frac{\partial}{\partial y} \right) P(x, y) &= \frac{r}{2} \frac{\partial}{\partial x} P(x, y) - \frac{r}{2} \frac{P(x, y)}{x} + \frac{r}{2} \frac{\partial}{\partial y} P(x, y) - \frac{r}{2} \frac{P(x, y)}{y} \\ &+ \frac{r}{2} \frac{y}{x} \frac{\partial}{\partial y} P(x, y) - \frac{r}{2} \sum_{j \geq 1} p_{1,j} j y^j + \frac{r}{2} \frac{x}{y} \frac{\partial}{\partial x} P(x, y) - \frac{r}{2} \sum_{i \geq 1} p_{i,1} i x^i \\ &+ d \left(y^2 \frac{\partial}{\partial y} P(x, y) + y P(x, y) + \sum_{i \geq 1} x^i y \right) + d \left(x^2 \frac{\partial}{\partial x} P(x, y) + x P(x, y) + \sum_{j \geq 1} x y^j \right), \end{aligned} \quad (3.10)$$

and finally $\sum_{i \geq 1} p_{i,1} i x^i = x \frac{\partial^2 P}{\partial x \partial y}(x, 0)$.

For point (ii), local existence and uniqueness stem from classical theorems [13]. Note that we construct an explicit, albeit complicated solution (3.5) using the method of characteristics. This concludes the proof. \blacksquare

Remark 1. In the case of homogeneous random walks, the operator $AP(x, y)$ has the product form $R(x, y)P(x, y)$, see [9]. In some sense, this means that the inhomogeneity leads to partial derivatives in the functional equation.

Remark 2. This technique allowing to compute the solution of a discrete, linear problem thanks to generating series is also called the Z-transform method.

3.2 Characteristic curves

In Sections 3.2 and 3.3 we establish, by the methods of characteristic equations, an explicit formula for the solutions to (3.1), which proves the existence of the solution. Since A is a first-order differential operator, we have a transport-like PDE. We introduce the following characteristic ordinary differential equations (ODEs). Let $(x_s, y_s)_{s \in \mathbb{R}_+}$ be the solution to the system:

$$\begin{cases} \dot{x}_s &= \frac{dx}{ds}(s) = Q(x_s, y_s), \\ \dot{y}_s &= \frac{dy}{ds}(s) = Q(y_s, x_s), \end{cases} \quad (3.11)$$

where Q has been defined in (3.2). The dynamical system (3.11) and its solutions will turn out to be decisive in the sequel—*e.g.* in Prop. 3.4, where we will use these characteristic equations in order to express the solutions to the fundamental functional equation (3.1).

Proposition 3.2. *For any initial condition $(x_0, y_0) \in \mathbb{R}^2$ such that $x_0, y_0 \neq 0$, there exists a unique solution to (3.11), defined for $s \in \mathbb{R}$ by:*

$$x_s = \frac{\lambda r e^{rs} + \mu d e^{ds}}{d(\lambda e^{rs} + \mu e^{ds} + 1)}, \quad y_s = \frac{\lambda r e^{rs} + \mu d e^{ds}}{d(\lambda e^{rs} + \mu e^{ds} - 1)}, \quad (3.12)$$

with:

$$\lambda = \frac{2d x_0 y_0 - d(x_0 + y_0)}{(x_0 - y_0)(r - d)}, \quad \mu = \frac{-2d x_0 y_0 + r(x_0 + y_0)}{(x_0 - y_0)(r - d)}. \quad (3.13)$$

Proof. First of all, since $Q(x, y)$ and $Q(y, x)$ are locally Lipschitz-continuous for $x, y \neq 0$, there is local existence and uniqueness. Let us introduce the new variable $z_s = x_s/y_s$, which is defined as long as $y_s \neq 0$. Using the expression (3.2) of Q , the system (3.11) becomes:

$$\dot{x}_s = (r + d)x_s - \frac{r}{2}(1 + z_s) - dx_s^2, \quad (3.14)$$

$$\begin{aligned} \dot{z}_s &= \frac{\dot{x}_s y_s - x_s \dot{y}_s}{y_s^2} \\ &= (r + d)z_s - \frac{r}{2} \frac{z_s}{x_s} - \frac{r}{2} \frac{z_s}{y_s} - dx_s z_s - (r + d)z_s + \frac{r}{2} \frac{z_s}{y_s} + \frac{r}{2} \frac{z_s}{x_s} + dy_s z_s \\ &= dx_s(1 - z_s). \end{aligned} \quad (3.15)$$

From (3.15), we obtain:

$$x_s = \frac{\dot{z}_s}{d(1 - z_s)}, \quad (3.16)$$

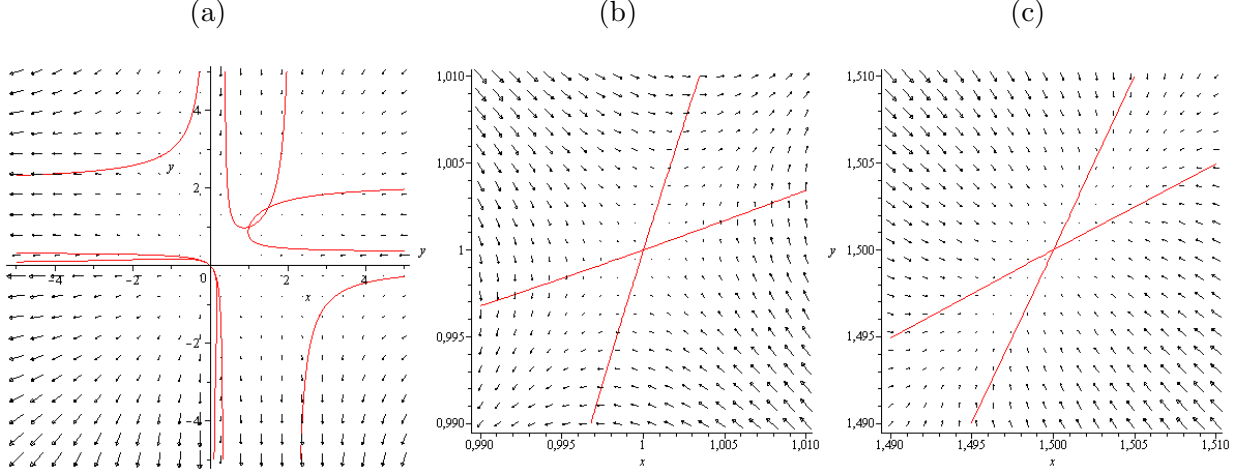


Figure 2: *Differential system (3.11): (a) Vector field $(Q(x, y), Q(y, x))$. The plain lines correspond to the sets $\{Q(x, y) = 0\}$ or $\{Q(y, x) = 0\}$. (b) Neighborhood of the saddle point $(1, 1)$. (c) Neighborhood of the attractive equilibrium $(r/d, r/d)$.*

and differentiating (3.15) with respect to the time s gives:

$$\dot{x}_s = \frac{d\dot{z}_s(1 - z_s) + d\dot{z}_s^2}{d^2(1 - z_s)^2}.$$

Plugging this expression and (3.16) into (3.14) provides:

$$\frac{d\dot{z}_s(1 - z_s) + d\dot{z}_s^2}{d^2(1 - z_s)^2} = \frac{(r + d)\dot{z}_s}{d(1 - z_s)} - \frac{r}{2}(1 + z_s) - d\frac{\dot{z}_s^2}{d^2(1 - z_s)^2},$$

and, therefore,

$$\ddot{z}_s(1 - z_s) + 2\dot{z}_s^2 - (r + d)(1 - z_s)\dot{z}_s + \frac{rd}{2}(1 + z_s)(1 - z_s)^2 = 0. \quad (3.17)$$

The system (3.16)–(3.17) can be solved explicitly. Let us define $u_s = (1 + z_s)/(1 - z_s)$, so that

$$z_s = \frac{u_s - 1}{u_s + 1}, \quad (3.18)$$

from which we obtain:

$$\dot{z}_s = \frac{2\dot{u}_s}{(1 + u_s)^2}, \quad \ddot{z}_s = \frac{2\ddot{u}_s(1 + u_s) - 4\dot{u}_s^2}{(1 + u_s)^3}. \quad (3.19)$$

Using these expressions in (3.17) provides:

$$\begin{aligned} & \frac{2\ddot{u}_s(1 + u_s) - 4\dot{u}_s^2}{(1 + u_s)^3} \frac{2}{u_s + 1} + 2\frac{4\dot{u}_s^2}{(1 + u_s)^4} - (r + d)\frac{2}{u_s + 1} \frac{2\dot{u}_s}{(1 + u_s)^2} + \frac{rd}{2} \frac{2u_s}{u_s + 1} \frac{4}{(u_s + 1)^2} = 0 \\ \Leftrightarrow & 4\ddot{u}_s(1 + u_s) - 4(r + d)\dot{u}_s(1 + u_s) + 4rd u_s(u_s + 1) = 0 \\ \Leftrightarrow & \ddot{u}_s - (r + d)\dot{u}_s + rd u_s = 0. \end{aligned} \quad (3.20)$$

Hence u satisfies a second-order ordinary differential equation, that solves in:

$$u_s = \lambda e^{rs} + \mu e^{ds}, \quad \lambda, \mu \in \mathbb{R}. \quad (3.21)$$

From (3.16), (3.18) and (3.21):

$$z_s = \frac{\lambda e^{rs} + \mu e^{ds} - 1}{\lambda e^{rs} + \mu e^{ds} + 1}, \quad x_s = \frac{\lambda r e^{rs} + \mu d e^{ds}}{d(\lambda e^{rs} + \mu e^{ds} + 1)}. \quad (3.22)$$

The integration constants λ and μ can be expressed in terms of the initial conditions x_0 and $z_0 = x_0/y_0$:

$$\lambda = \frac{2d x_0 - d(1 + z_0)}{(1 - z_0)(r - d)}, \quad \mu = \frac{r - 2d x_0 + z_0 r}{(1 - z_0)(r - d)}. \quad (3.23)$$

This yields the announced result with $y_s = x_s/z_s$. ■

Remark 3. Explicit expressions for x_s and y_s as functions of time and parameterized by r and d can be obtained. Using the expressions (3.13) of λ and μ in (3.22) and the relations between x_s , y_s and z_s provides:

$$x_s = \frac{-\frac{x_0 + y_0 - 2x_0 y_0}{x_0 + y_0 - 2(d/r)x_0 y_0} \exp(rs) + \exp(ds)}{-\frac{d}{r} \frac{x_0 + y_0 - 2x_0 y_0}{x_0 + y_0 - 2(d/r)x_0 y_0} \exp(rs) + \exp(ds) + \frac{(1 - d/r)(y_0 - x_0)}{x_0 + y_0 - 2(d/r)x_0 y_0}}, \quad (3.24)$$

$$y_s = \frac{-\frac{x_0 + y_0 - 2x_0 y_0}{x_0 + y_0 - 2(d/r)x_0 y_0} \exp(rs) + \exp(ds)}{-\frac{d}{r} \frac{x_0 + y_0 - 2x_0 y_0}{x_0 + y_0 - 2(d/r)x_0 y_0} \exp(rs) + \exp(ds) - \frac{(1 - d/r)(y_0 - x_0)}{x_0 + y_0 - 2(d/r)x_0 y_0}}. \quad (3.25)$$

For the sequel, it is useful to notice that the constant

$$-\frac{\lambda r}{\mu d} = \frac{x_0 + y_0 - 2x_0 y_0}{x_0 + y_0 - 2(d/r)x_0 y_0} \quad (3.26)$$

which appears in (3.24) and (3.25) belongs to $]0, 1[$ for all $(x_0, y_0) \in]0, 1[^2$; this is due to the fact that $d/r < 1$.

Below, we list some properties of the solutions to the dynamical system 3.11. The trajectories of the solutions to (3.11) can be decomposed into four steps. In order to describe them, let us introduce

$$s_0 = \frac{1}{r - d} \log \left(-\frac{\mu d}{\lambda r} \right) = \frac{\log \left(\frac{x_0 + y_0 - 2x_0 y_0}{x_0 + y_0 - 2(d/r)x_0 y_0} \right)}{d - r} \quad (3.27)$$

and s_{\pm} as the only positive roots of the denominators of x_s and y_s in (3.12):

$$\lambda e^{rs} + \mu e^{ds} \pm 1. \quad (3.28)$$

Proposition 3.3. *Let $(x_0, y_0) \in]0, 1[^2$.*

- (i) *The stationary solutions to (3.11) are the saddle point $(1, 1)$ and the attractive point $(r/d, r/d)$.*
- (ii) *$s_0 \in]0, \infty[$ and $s_{\pm} \in]s_0, \infty[$.*
- (iii) *$\inf\{s_+, s_-\} = s_+$ (resp. s_-) if and only if $y_0 < x_0$ (resp. $y_0 > x_0$).*
- (iv) *$\lim_{s \uparrow s_+} x_s = -\infty$ and $\lim_{s \downarrow s_-} y_s = -\infty$.*

The next points specify the behavior of the solution to (3.11).

- (v) *On $[0, s_0[$, (x_s, y_s) belongs to $]0, 1[^2$ and goes to $(x_{s_0}, y_{s_0}) = (0, 0)$ as $s \rightarrow 0$.*

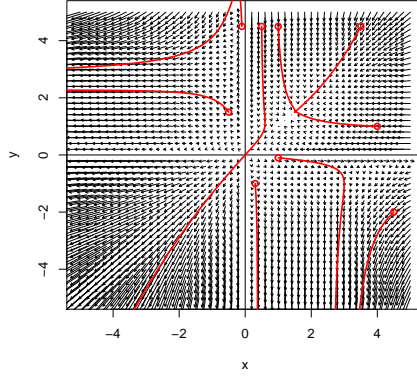


Figure 3: Solutions corresponding to several initial conditions. We see that depending on the initial condition, the solutions converge or diverge to infinity.

- (vi) On $]s_0, \inf\{s_+, s_-\}[$, (x_s, y_s) goes decreasingly to $(x_{\inf\{s_+, s_-\}}, y_{\inf\{s_+, s_-\}})$ —by “decreasingly” we mean that both coordinates decrease.
- (vii) On $] \inf\{s_+, s_-\}, \sup\{s_+, s_-\}[$, (x_s, y_s) goes decreasingly to $(x_{\sup\{s_+, s_-\}}, y_{\sup\{s_+, s_-\}})$.
- (viii) On $] \sup\{s_+, s_-\}, \infty[$, (x_s, y_s) goes decreasingly to $(r/d, r/d)$.

Proof of Item (i). To find the stationary solutions to (3.11), let us solve $\dot{x} = 0$ and $\dot{y} = 0$. With (3.11), we get $Q(x, y) = 0$ and $Q(y, x) = 0$. This directly implies that $(x, y) = (1, 1)$ or $(x, y) = (r/d, r/d)$, see Table 2. Let us now study the stability of these two equilibria.

At $(1, 1)$, the Jacobian of (3.11) is:

$$\text{Jac}(1, 1) = \frac{r}{2}J - dI,$$

where J is the 2×2 matrix full of ones and where I is the 2×2 identity matrix. The eigenvalues of $\text{Jac}(1, 1)$ are $-d < 0$ and $r - d > 0$, associated with the eigenvectors $(1, -1)$ and $(1, 1)$, respectively. By classical linearization methods (e.g. [11, Chap. 3]), we deduce that the point $(1, 1)$ is a saddle point.

At $(r/d, r/d)$, the Jacobian of (3.11) is:

$$\text{Jac}(r/d, r/d) = \frac{d}{2}J - rI.$$

The eigenvalues are $-r < 0$ and $d - r < 0$, associated with the eigenvectors $(1, -1)$ and $(1, 1)$, respectively. The point $(r/d, r/d)$ is therefore attractive. \blacksquare

Proof of Item (ii) to (iv). Now we prove the different facts dealing with s_0 , s_+ and s_- . First, (3.26) and the fact that $r > d$ immediately imply that $s_0 \in]0, \infty[$. Next, we show that (3.28) has on $[0, \infty[$ only one root, which belongs to $]s_0, \infty[$. For this, we shall start with proving that (3.28) is positive on $[0, s_0]$. Then, we shall show that (3.28) is decreasing in $]s_0, \infty[$ and goes to $-\infty$ as $s \rightarrow \infty$.

In order to prove the first point above, it is enough to show that (3.28) is positive at $s = 0$ and increasing on $[0, s_0]$. (3.28) is positive at $s = 0$ simply because

$$\lambda + \mu \pm 1 = \frac{x_0 + y_0}{y_0 - x_0} \pm 1 = \frac{2y_0}{y_0 - x_0} \text{ or } \frac{2x_0}{y_0 - x_0} = \frac{(1 - d/r)[x_0 + y_0 \pm (y_0 - x_0)]}{x_0 + y_0 - 2(d/r)x_0y_0} > 0.$$

To check that (3.28) is increasing on $[0, s_0]$, we note that the derivative of (3.28) is positive on $[0, s_0]$ —actually by construction of s_0 .

Now we prove the second point. From (3.26) and since $r > d$, we obtain that (3.28) goes to $-\infty$ as $s \rightarrow \infty$. Also, by definition of s_0 , the derivative of (3.28) is negative on $]s_0, \infty[$, (3.28) is therefore decreasing on $]s_0, \infty[$.

The fact that $\inf\{s_+, s_-\}$ equals s_+ (resp. s_-) if and only if $y_0 < x_0$ (resp. $y_0 > x_0$) follows directly from (3.28).

Finally, since the numerators of x_s and y_s are negative on $]s_0, \infty[$, hence in particular at s_{\pm} , it is immediate that $\lim_{s \uparrow s_+} x_s = -\infty$ and $\lim_{s \downarrow s_-} y_s = -\infty$. \blacksquare

Proof of Item (v) to (viii). Let us first consider *Item (v)*. By definition of s_0 , the numerators of x_s and y_s in (3.24) and (3.25) vanish for the first time at s_0 . Moreover, since $s_{\pm} > s_0$, both denominators are non-zero at s_0 and $x_{s_0} = y_{s_0} = 0$. In particular, on $[0, s_0[$, we have $(x_s, y_s) \in \mathbb{R}_+^2$. In fact, $(x_s, y_s) \in]0, 1]^2$. Indeed, on the segment $\{1\} \times]0, 1[$, $\dot{x}_s < 0$ whereas on $]0, 1[\times \{1\}$, $\dot{y}_s < 0$: it is therefore not possible to go through these segments and on $[0, s_0]$, (x_s, y_s) remains inside of the square $[0, 1]^2$.

We turn to the proof of *Item (vi)*. Thanks to (3.24) and (3.25), just after the time s_0 , (x_s, y_s) belongs to the negative quadrant \mathbb{R}_-^2 . But for any $(x, y) \in \mathbb{R}_-^2$, $Q(x, y) \leq 0$ and $Q(y, x) \leq 0$, see Table 2, in such a way that both x_s and y_s are decreasing as soon as they stay in this quarter plane, in other words for $s \in]s_0, \inf\{s_+, s_-\}[$. At time $\inf\{s_+, s_-\}$, one (or even the two if $s_+ = s_-$, *i.e.* if $x_0 = y_0$) of x_s and y_s becomes infinite. In the sequel, let us assume that $\inf\{s_+, s_-\} = s_+$; a similar reasoning would hold for the symmetrical case $\inf\{s_+, s_-\} = s_-$.

Let us show *Item (vii)*. Just after s_+ , $(x_s, y_s) \in (\mathbb{R}_+ \times \mathbb{R}_-) \cap \{(x, y) \in \mathbb{R}^2 : Q(x, y) < 0, Q(y, x) < 0\}$. The latter set is simply connected and bounded by the curve $\{(x, y) \in \mathbb{R}^2 : Q(x, y) = 0\}$, see Table 2. Using classical arguments (see *e.g.* [11]), we obtain that it is not possible to go through this limiting curve on which $\dot{x}_s = 0$; this is why for any $s \in]s_+, s_-[$, (x_s, y_s) remains inside of this set.

We conclude with the proof of *Item (viii)*. Just after the time s_- , $(x_s, y_s) \in \mathbb{R}_+^2 \cap \{(x, y) \in \mathbb{R}^2 : Q(x, y) < 0, Q(y, x) < 0\}$. For the same reasons as above, (x_s, y_s) cannot leave this set and actually converges to $(r/d, r/d)$. \blacksquare

3.3 Use of the characteristic curves to simplify the functional equation

Let us assume the existence of a solution $P(x, y)$ to (3.1), and let us define $g_s = P(x_s, y_s)$. Then:

$$\dot{g}_s = \frac{dg}{ds}(s) = \frac{\partial P}{\partial x}(x_s, y_s) \frac{dx_s}{ds} + \frac{\partial P}{\partial y}(x_s, y_s) \frac{dy_s}{ds} = \frac{\partial P}{\partial x}(x_s, y_s) Q(x_s, y_s) + \frac{\partial P}{\partial y}(x_s, y_s) Q(y_s, x_s).$$

Thus, if P is a solution to (3.1), then:

$$\dot{g}_s + R(x_s, y_s)g_s = h(x_s, y_s, P),$$

which looks like a first-order ODE for g , except that h depends on the boundary condition of P .

We first freeze the dependence on the solution in h . Using the solutions to the characteristic equations, we shall prove the following result:

Proposition 3.4. *Let $h(x, y)$ be an analytical function on $[0, 1]^2$. Let $(x_0, y_0) \in \mathbb{R}^2$. The solution to the ODE*

$$\dot{g}_s + R(x_s, y_s)g_s = h(x_s, y_s), \quad g_0 = P(x_0, y_0), \quad (3.29)$$

where $(x_s, y_s)_{s \geq 0}$ are the solutions (3.12) to the characteristic curve starting at (x_0, y_0) , is given by

$$g_s^h = P(x_0, y_0) \exp\left(-\int_0^s R(x_u, y_u) du\right) + \int_0^s h(x_u, y_u) e^{-\int_u^s R(x_\alpha, y_\alpha) d\alpha} du =: F(s, x_0, y_0, h). \quad (3.30)$$

Proof. Equation (3.29) is an inhomogeneous first-order ODE. The solution to the associated homogeneous equation is:

$$g_s = R(x_0, y_0) \exp\left(-\int_0^s R(x_u, y_u) du\right).$$

The announced result is deduced from the variation of constant method. ■

A solution P to (3.1) hence satisfies the following functional equation for all s, x_0 and y_0 :

$$P(x_s, y_s) = P(x_0, y_0) e^{-\int_0^s R(x_u, y_u) du} + \int_0^s h(x_u, y_u, P) e^{-\int_u^s R(x_\alpha, y_\alpha) d\alpha} du, \quad (3.31)$$

with the function h defined in (3.3). Plugging the definitions (1.4) and (3.3) in (3.31), we obtain:

$$\begin{aligned} \sum_{i, j \geq 1} p_{i, j} x_s^i y_s^j &= P(x_0, y_0) e^{-\int_0^s R(x_\alpha, y_\alpha) d\alpha} \\ &\quad - \frac{r}{2} \sum_{i \geq 1} p_{i, 1} \int_0^s i(x_u)^i e^{-\int_u^s R(x_\alpha, y_\alpha) d\alpha} du - \frac{r}{2} \sum_{j \geq 1} p_{1, j} \int_0^s j(y_u)^j e^{-\int_u^s R(x_\alpha, y_\alpha) d\alpha} du \\ &\quad + d \int_0^s x_u y_u \left(\frac{1}{1-x_u} + \frac{1}{1-y_u} \right) e^{-\int_u^s R(x_\alpha, y_\alpha) d\alpha} du. \end{aligned} \quad (3.32)$$

Notice that the r.h.s. of (3.32) depends only on the $p_{i,1}$'s and $p_{1,j}$'s, while the l.h.s. depends on all $p_{i,j}$'s.

Proposition 3.5. *Let $s_0 > 0$ be defined in (3.27). We have:*

$$\begin{aligned} P(x_0, y_0) &= \frac{r}{2} \sum_{i \geq 1} p_{i, 1} \int_0^{s_0} i(x_u)^i e^{\int_0^u R(x_\alpha, y_\alpha) d\alpha} du + \frac{r}{2} \sum_{j \geq 1} p_{1, j} \int_0^{s_0} j(y_u)^j e^{\int_0^u R(x_\alpha, y_\alpha) d\alpha} du \\ &\quad - d \int_0^{s_0} x_u y_u \left(\frac{1}{1-x_u} + \frac{1}{1-y_u} \right) e^{\int_0^u R(x_\alpha, y_\alpha) d\alpha} du. \end{aligned} \quad (3.33)$$

Before proving Prop. 3.5, let us show that the different quantities that appear in its statement are well defined—indeed, this is *a priori* not clear: as $\alpha \rightarrow s_0$, $x_\alpha \rightarrow 0$ and $y_\alpha \rightarrow 0$, in such a way that $R(x_\alpha, y_\alpha) \rightarrow \infty$, see (3.2).

Lemma 3.6. *Let $i, j \in \mathbb{N}$. Then $\lim_{u \rightarrow s_0} (x_u)^i (y_u)^j e^{\int_0^u R(x_\alpha, y_\alpha) d\alpha}$ is finite if and only if $i + j \geq 1$ —and equals zero if and only if $i + j \geq 2$.*

Proof. First, since the only zero of any function of the form $\alpha \exp(as) + \beta \exp(bs)$ with $\alpha\beta < 0$ and $a \neq b$ has order one, the following function has a simple zero at s_0 :

$$-\frac{x_0 + y_0 - 2x_0y_0}{x_0 + y_0 - 2(d/r)x_0y_0} \exp(ru) + \exp(du). \quad (3.34)$$

Thanks to this fact and since $s_+, s_- > s_0$, both x_s and y_s have at s_0 a zero of order one.

Moreover, with λ and μ defined in (3.13), we obtain that:

$$\begin{aligned} \exp\left(\int_0^u R(x_\alpha, y_\alpha) d\alpha\right) &= \frac{\lambda/\mu + 1 - 1/\mu}{(\lambda/\mu) \exp(ru) + \exp(du) - 1/\mu} \times \\ &\quad \times \frac{\lambda/\mu + 1 + 1/\mu}{(\lambda/\mu) \exp(ru) + \exp(du) + 1/\mu} \frac{r/d\lambda/\mu + 1}{(r/d)(\lambda/\mu) \exp(ru) + \exp(du)} \exp((r+d)u). \end{aligned} \quad (3.35)$$

It is indeed easy to check that the derivative of the logarithm of (3.35) is actually equal to $R(x_u, y_u)$, for which we have an explicit expression, see (3.2), (3.24) and (3.25).

From (3.35), we see that $e^{\int_0^u R(x_\alpha, y_\alpha) d\alpha}$ has three poles, namely at s_0, s_+, s_- . The zero at s_0 has order one by using again the considerations on the zeros of (3.34). In particular, Lem. 3.6 follows immediately. ■

Proof of Prop. 3.5. Start by multiplying Equation (3.32) by $e^{\int_0^s R(x_\alpha, y_\alpha) d\alpha}$ and then let $s \rightarrow s_0$. Since $P(x, y) = xy \sum_{i, j \geq 1} p_{i, j} x^{i-1} y^{j-1}$, see (1.4), and since $\lim_{s \rightarrow s_0} x_s y_s e^{\int_0^s R(x_\alpha, y_\alpha) d\alpha} = 0$, see Lem. 3.6, we obtain that

$$\lim_{s \rightarrow s_0} P(x_s, y_s) e^{\int_0^s R(x_\alpha, y_\alpha) d\alpha} = 0,$$

which concludes the proof of Prop. 3.5. ■

Remark 3.7. When $(x_0, y_0) \in (0, 1)^2$, we also have $(x_u, y_u) \in (0, 1)^2$ for all $u \in (0, s_0)$. Thus it is possible to plug approximations of the $p_{1, j}$'s and $p_{i, 1}$'s into (3.36) thanks to the terms $(x_u)^i$ and $(y_u)^j$. Using that $p_{1, i} = 2d/(ri) + o(1/i)$ when $i \rightarrow +\infty$, it is possible to find I_0 sufficiently large so that:

$$\begin{aligned} & \frac{r}{2} \sum_{i > I_0} p_{i, 1} \int_0^{s_0} i((x_u)^i + (y_u)^i) e^{\int_0^u R(x_\alpha, y_\alpha) d\alpha} du \\ & \sim d \int_0^{s_0} \sum_{i > I_0} ((x_u)^i + (y_u)^i) e^{\int_0^u R(x_\alpha, y_\alpha) d\alpha} du = d \int_0^{s_0} \left(\frac{x_u^{I_0+1}}{1-x_u} + \frac{y_u^{I_0+1}}{1-y_u} \right) e^{\int_0^u R(x_\alpha, y_\alpha) d\alpha} du \end{aligned}$$

Thus:

$$\begin{aligned} P(x_0, y_0) &= \frac{r}{2} \sum_{i=1}^{I_0} p_{i, 1} \int_0^{s_0} i((x_u)^i + (y_u)^i) e^{\int_0^u R(x_\alpha, y_\alpha) d\alpha} du \\ &+ d \int_0^{s_0} \left(\frac{x_u(x_u^{I_0} - y_u)}{1-x_u} + \frac{y_u(y_u^{I_0} - x_u)}{1-y_u} \right) e^{\int_0^u R(x_\alpha, y_\alpha) d\alpha} du + o(x_0^{I_0+1} + y_0^{I_0+1}). \end{aligned} \quad (3.36)$$

The latter expression shows that numerically, one can restrict to the computation of a finite number of probabilities $p_{i, 1}$, for $i \leq I_0$.

4 Numerical results

In this section, we present two different ways of approximating the extinction probabilities $p_{i, j}$.

4.1 Probabilistic algorithm

A first possibility, if we are interested in a given initial condition (i, j) , is to approximate $p_{i, j}$ by Monte-Carlo simulations. For $T > 0$ large, we simulate M paths $(X_t^\ell, Y_t^\ell)_{t \in \{1, \dots, T\}}$ started at (i, j) , for $\ell \in \{1, \dots, M\}$, independent and distributed as the process $(X_t, Y_t)_{t \in \{1, \dots, T\}}$. The extinction probability is estimated by:

$$\hat{p}_{M, T} = \frac{1}{M} \sum_{\ell=1}^M \mathbb{1}_{\{\exists t \leq T, X_t^\ell Y_t^\ell = 0\}}.$$

The estimator $\hat{p}_{M, T}$ is the proportion of paths that have gone extinct before time T .

Proposition 4.1. *Let (i, j) be the initial condition. The estimator $\hat{p}_{M, T}$ has the following properties:*

- (i) *It is a convergent and unbiased estimator of $\mathbb{P}_{i, j}[\tau_0 \leq T]$.*
- (ii) *Its variance is $\mathbb{P}_{i, j}[\tau_0 \leq T](1 - \mathbb{P}_{i, j}[\tau_0 \leq T])/M$, and hence we have the following asymptotic 95% confidence interval for $p_{i, j}$:*

$$\left[\hat{p}_{M, T} - 1.96 \sqrt{\frac{\hat{p}_{M, T}(1 - \hat{p}_{M, T})}{M}}; \hat{p}_{M, T} + 1.96 \sqrt{\frac{\hat{p}_{M, T}(1 - \hat{p}_{M, T})}{M}} \right]. \quad (4.1)$$

Proof. These results are straightforward consequences of the law of large numbers and central limit theorem, given that $\mathbb{1}_{\{\exists t \leq T, X_t^\ell Y_t^\ell = 0\}}$ are independent Bernoulli random variables with parameter $\mathbb{P}_{i,j}[\tau_0 \leq T]$. \blacksquare

Computing the extinction probabilities by Monte-Carlo methods yields good results if we are interested in a given initial condition (i, j) . We then have a complexity of order $M \times T$. However, Biologists may be interested in investigating the extinction probabilities when the initial condition (i, j) varies, and the method become computationally expensive.

4.2 Deterministic algorithm

For numerical approximations, we restrict ourselves to the computation of $(p_{i,j})_{i,j \in \{1, \dots, N\}}$ for a positive (large) integer N . In this case, (1.7) can be approximated by the solution to a linear system.

Let us define $\mathbf{p}_N = (p_{1,1}, \dots, p_{1,N}, p_{2,1}, \dots, p_{2,N}, \dots, p_{N,1}, \dots, p_{N,N})^T$ and T_N is a $N^2 \times N^2$ -matrix with five non-zero diagonals:

$$T_N = \begin{pmatrix} A_1 & D & 0 & \dots & 0 \\ B_{2,1} & A_2 & D & \ddots & \vdots \\ 0 & B_{3,2} & \ddots & \ddots & 0 \\ \vdots & \ddots & \ddots & \ddots & D \\ 0 & \dots & 0 & B_{N,N-1} & A_N \end{pmatrix}$$

where $D = \frac{r}{2(r+d)} Id_N$, A_i ($i \in \{1, \dots, N\}$) and $B_{i,i-1}$ ($i \in \{2, \dots, N\}$) are the $N \times N$ -matrices

$$A_i = \begin{pmatrix} -1 & \frac{r}{2(r+d)} & 0 & \dots & 0 \\ \frac{d}{r+d} \frac{2}{i+2} & -1 & \ddots & \ddots & \vdots \\ 0 & \frac{d}{r+d} \frac{3}{i+3} & \ddots & \ddots & 0 \\ \vdots & \ddots & \ddots & \ddots & \frac{r}{2(r+d)} \\ 0 & \dots & 0 & \frac{d}{r+d} \frac{N}{i+N} & -1 \end{pmatrix},$$

$$B_{i,i-1} = \begin{pmatrix} \frac{d}{r+d} \frac{i}{i+1} & 0 & \dots & 0 \\ 0 & \frac{d}{r+d} \frac{i}{i+2} & \ddots & \vdots \\ \vdots & \ddots & \ddots & 0 \\ 0 & \dots & 0 & \frac{d}{r+d} \frac{i}{i+N} \end{pmatrix}.$$

Let us also define the vector $b_N = (b_{1N}, \dots, b_{NN})^T \in \mathbb{R}^{N \times N}$ such that:

$$b_{1N} = - \begin{pmatrix} -\frac{d}{r+d} \\ \frac{d}{r+d} \frac{1}{1+2} \\ \vdots \\ \frac{d}{r+d} \frac{1}{1+(N-1)} \\ \frac{d}{r+d} \frac{1}{1+N} + \frac{r}{2(r+d)} \tilde{p}_{1,N+1} \end{pmatrix}, \quad \text{for } i \in \{2, \dots, N-1\}, \quad b_{iN} = - \begin{pmatrix} \frac{d}{r+d} \frac{1}{1+i} \\ 0 \\ \vdots \\ 0 \\ \frac{r}{2(r+d)} \tilde{p}_{i,N+1} \end{pmatrix},$$

$$b_{NN} = - \begin{pmatrix} \frac{d}{r+d} \frac{1}{1+N} + \frac{r}{2(r+d)} \tilde{p}_{N+1,1} \\ \frac{r}{2(r+d)} \tilde{p}_{N+1,2} \\ \vdots \\ \frac{r}{2(r+d)} \tilde{p}_{N+1,N-1} \\ \frac{r}{2(r+d)} (\tilde{p}_{N+1,N} + \tilde{p}_{N,N+1}) \end{pmatrix}$$

where $\tilde{p}_{i,N+1}$ and $\tilde{p}_{N+1,N}$ are approximations of $p_{i,N+1}$ given by Prop. 2.3. With these notations, (1.7) rewrites as

$$T_N \mathbf{p}_N = b_N.$$

4.3 Results

We start with $r = 3$ and $d = 2$. For the Monte-Carlo simulation, we use $M = 200$ and $T = 5000$. For the deterministic method, we use $N = 50$, so that $(i, j) \in \{1, \dots, 50\}^2$. Estimators of the extinction probabilities $\hat{p}_{i,j}^{(1)}$ and $\hat{p}_{i,j}^{(2)}$ obtained respectively from the methods of Sections 4.1 and 4.2 are plotted in Figure 4. The results given by both methods are very similar, as shown by the statistics of Table 1. In Table 1, we compute the square difference between the two predictions $(\hat{p}_{i,j}^{(1)} - \hat{p}_{i,j}^{(2)})^2$, the absolute difference $|\hat{p}_{i,j}^{(1)} - \hat{p}_{i,j}^{(2)}|$ and the relative difference $|\hat{p}_{i,j}^{(1)} - \hat{p}_{i,j}^{(2)}|/\hat{p}_{i,j}^{(2)}$. For the latter, we consider only the couples (i, j) where $\hat{p}_{i,j}^{(1)}$ and $\hat{p}_{i,j}^{(2)}$ do not vanish (else, the fraction is either not defined or either 1 whatever the value of $\hat{p}_{i,j}^{(1)}$).

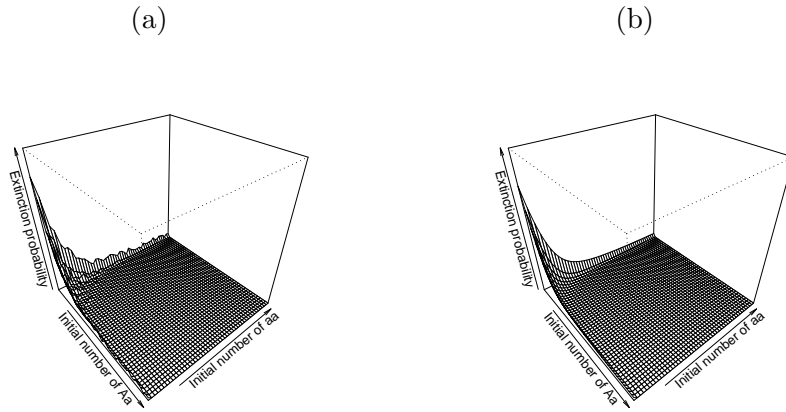


Figure 4: Estimation of the extinction probabilities $p_{i,j}$'s when $r = 3$ and $d = 2$: (a) with the Monte-Carlo method of Section 4.1. (b) with the deterministic method of Section 4.2.

	Mean	St.dev	Min	Max
Square error	$3.24 \cdot 10^{-5}$	$2.47 \cdot 10^{-4}$	$1.68 \cdot 10^{-36}$	$4.63 \cdot 10^{-3}$
Absolute error	$1.07 \cdot 10^{-8}$	$5.53 \cdot 10^{-3}$	$1.30 \cdot 10^{-18}$	$6.81 \cdot 10^{-2}$
Relative error	$4.49 \cdot 10^{-2}$	$1.71 \cdot 10^{-1}$	$9.80 \cdot 10^{-3}$	$9.71 \cdot 10^{-1}$

Table 1: Square, absolute and relative differences between the predictions of the stochastic method when $r = 3$ and $d = 2$ (Section 4.1) and of the deterministic methods (Section 4.2). Recall that with $M = 200$, the width of the confidence interval (4.1) is $6.92 \cdot 10^{-2}$.

To carry further the comparison of the stochastic and deterministic method, and to observe the influence of N on the quality of the approximation, we compute the relative quadratic error

$$\frac{\sqrt{\sum_{1 \leq i, j \leq 10} (\hat{p}_{i,j}^{(1)} - \hat{p}_{i,j}^{(3)})^2}}{\sqrt{\sum_{1 \leq i, j \leq 10} (\hat{p}_{i,j}^{(1)})^2}} \text{ or } \frac{\sqrt{\sum_{1 \leq i, j \leq 10} (\hat{p}_{i,j}^{(1)} - \hat{p}_{i,j}^{(3)})^2}}{\sqrt{\sum_{1 \leq i, j \leq 10} (\hat{p}_{i,j}^{(3)})^2}} \quad (4.2)$$

when $\hat{p}_{i,j}^{(1)}$ is the deterministic approximation for $N \in \{10, \dots, 50\}$ and $\hat{p}_{i,j}^{(3)}$ is either given by the deterministic approximation with $N = 50$, or by the stochastic approximation $\hat{p}_{i,j}^{(2)}$ with $M = 200$. In the first case when $\hat{p}_{i,j}^{(3)}$ is the deterministic approximation with $N = 50$, the relative quadratic errors decrease exponentially fast in $\exp(-0.6842 N)$ ($R^2 = 99.92\%$). In the second case, when $\hat{p}_{i,j}^{(3)} = \hat{p}_{i,j}^{(2)}$, the decrease in the quadratic errors stops around $N = 18$ around 0.0891. This corresponds roughly to the stochastic error of the law of large numbers (4.1) which depends only on M .

In a second experiment, we choose $r = 2.002$ and $d = 2$. This case is more interesting in population ecology, since small populations are of interest when they are fragile and endangered species. For the Monte-Carlo simulation, we use $M = 200$ and $T = 5000$. For the deterministic method, we use $N = 100$, so that $(i, j) \in \{1, \dots, 100\}^2$. The estimated extinction probabilities $\hat{p}_{i,j}^{(1)}$ and $\hat{p}_{i,j}^{(2)}$ are plotted in Figure 5, and statistics are computed in Table 2. Again, results from both methods are similar. This is confirmed by computing the relative quadratic errors, with the $\hat{p}_{i,j}^{(1)}$'s obtained from the deterministic method and the $\hat{p}_{i,j}^{(3)} = \hat{p}_{i,j}^{(2)}$'s from the Monte-Carlo method. The decrease of this error is exponential with N in $\exp(-0.0619 N)$ ($R^2 = 98.98\%$) as shown in Figure 5(c). It can be noticed that in this case, the performances of the Monte-Carlo method match better the one of the deterministic algorithm. This is due to the fact that Monte-Carlo methods fail to produce good estimates of small probabilities (see [5] and references therein).

When the probabilities $\hat{p}_{i,j}^{(3)}$'s are given by the deterministic method with $N = 50$ in (4.2), we have as in the previous case ($r = 3$) an exponential decrease of the relative quadratic error in $\exp(-0.1092 N)$ ($R^2 = 95.07\%$).

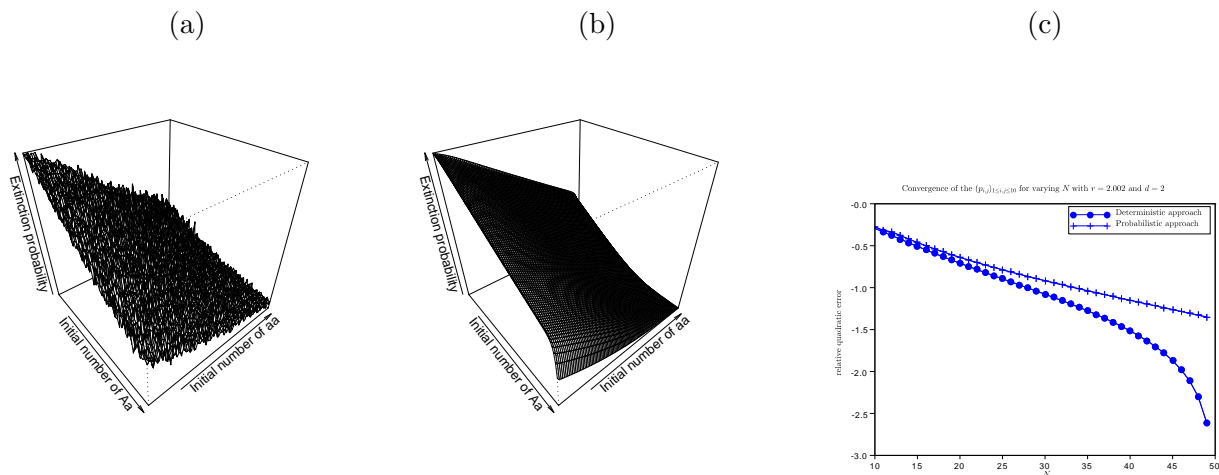


Figure 5: Estimation of the extinction probabilities $p_{i,j}$'s when $r = 2.002$ and $d = 2$: (a) with the Monte-Carlo method of Section 4.1. (b) with the deterministic method of Section 4.2. (c) Decrease with N of the relative quadratic errors (4.2) between the deterministic and Monte-Carlo methods ($M = 200$).

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	Mean	St.dev	Min	Max
Square error	$6.27 \cdot 10^{-3}$	$6.25 \cdot 10^{-3}$	$8.42 \cdot 10^{-10}$	$5.31 \cdot 10^{-2}$
Absolute error	$6.89 \cdot 10^{-2}$	$3.90 \cdot 10^{-2}$	$2.90 \cdot 10^{-5}$	$2.31 \cdot 10^{-1}$
Relative error	$2.22 \cdot 10^{-1}$	$2.17 \cdot 10^{-1}$	$1.02 \cdot 10^{-4}$	1

Table 2: Square, absolute and relative differences between the predictions of the stochastic method when $r = 2.002$ and $d = 2$ (Section 4.1) and of the deterministic methods (Section 4.2). As in Table 1, since $M = 200$, the width of the confidence interval (4.1) is $6.92 \cdot 10^{-2}$.

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