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On the Hausdorff dimension of Julia sets of some real polynomials

Genadi Levin, Michel Zinsmeister

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Résumé

We show that the supremum for c real of the Hausdorff dimension of the Julia set of the polynomial $z \mapsto z^d + c$ (d is an even natural number) is greater than 2d/(d+1).

1 Introduction and statement of the result

The Julia set of a non-linear polynomial $P: \mathbf{C} \to \mathbf{C}$ is the set of points having no neighborhood on which the family of iterates (P^n) is normal. It is a compact non-empty set which is, except for very special polynomials P, a fractal set. For $c \in \mathbf{C}$ we denote by f_c the polynomial

$$f_c(z) = z^d + c,$$

where $d \geq 2$ is an even integer number, which we fix. Denote by J_c the Julia set of f_c .

In the quadratic case d=2, the values of c for which the Hausdorff dimension (HD) of J_c is big has attracted a lot of attention. It has been known since the pioneering work of Douady-Hubbard [3] that the Hausdorff dimension of the Julia set is less than 2 for every hyperbolic polynomial. Thus $HD(J_c) < 2$ outside ∂M and the (hypothetic) non-hyperbolic components of (the interior) M. We recall that M stands for the Mandelbrot set, that is the compact subset of $c \in \mathbb{C}$ such that J_c is connected. Shishikura [9] was the first to find quadratic Julia sets with Hausdorff dimension 2. He indeed proved that this property holds on a dense \mathcal{G}_{δ} subset of ∂M or even on a dense \mathcal{G}_{δ} subset on the boundary of every hyperbolic component of M.

More recently, Buff and Chéritat [12] have found quadratic Julia sets with positive Lebesgue measure (see also http://annals.math.princeton.edu/articles/3682).

Both Shishikura's and Buff-Chéritat's results are based on the phenomenon of parabolic implosion which has been discovered and studied by Douady-Hubbard [3]. It should be pointed out that Buff-Chéritat's result is very involved and that we will make no use of it. There is no doubt that if they exist, values of $c \in \mathbf{R}$ such that the Julia set J_c has positive measure must be as hard to find as Buff-Cheritat's ones. The aim of present note is much more modest. Its starting point is the second author re-readind of Shishikura's result [10]: it states that if one implodes a polynomial with a parabolic cycle having q petals, then the dimension of its Julia set automatically gets bigger than 2q/(q+1). Shishikura's result follows from this by a Baire argument.

Very little is known about Hausdorff dimension of J_c for real c. In particular, it is not known if, for a given degree d,

$$\sup\{HD(J_c), c \in \mathbf{R}\} = \mathbf{2}.$$

Possible candidates for high dimension are of course (just look at them!) infinitely renormalizable polynomials but the analysis seems to be very delicate and at least no result concerning dimension 2 has been proven so far (for results in the opposite direction, see [1] though). It is for example unknown if the Julia set of the Feigenbaum polynomial has dimension 2 or not (see [4]- [5] for the Julia set of the Feigenbaum universal map though). The only known result about this set is a very general result of Zdunik [14]: it has dimension bigger than 1.

If one tries to use the same ideas as in [9] for real polynomials f_c , one immediately faces the problem that if f_c , for c real, has a parabolic cycle that may be imploded along the real axis then the number of petals is 1 and this does not imply more than Zdunik's general result. The only trick of this paper is to make use of a virtual doubling of petals when the critical point is mapped to a parabolic point (by Lavaurs map). It was inspired by Douady et al paper [13] and implies the following theorem, which is the main result of this work.

Theorem 1 Let $f_c(z) = z^d + c$, d even. Let N be the set of parameters $c \in \mathbf{R}$, such that f_c has a parabolic cycle of period at least 2 and multiplier 1. Then there exists an open set Y of \mathbf{R} whose closure contains N such that J_c is connected and

$$HD(J_c) > \frac{2d}{d+1},$$

for every $c \in Y$.

Comment 1 In fact, we prove a stronger statement: hyperbolic dimension [9] of J_c is bigger than 2d/(d+1).

Comment 2 By [11] ([6], [8] for d = 2), the set of real c such that f_c is hyperbolic is dense in \mathbf{R} . In particular, hyperbolic parameters c are dense in Y.

Acknowledgment. This work was done during the first author's one month's visit at the university of Orléans in 2008.

2 Proof of the theorem

We fix an even integer $d \geq 2$ and consider the family $f_c(x) = x^d + c$, for real c. Then the Julia set J_c is connected if and only if $c \in [a, b]$, where a < 0 is such that $f_a^2(0)$ is a fixed point, and b > 0 is such that f_b has a fixed point with multiplier 1. It is sufficient to prove that, given $c_0 \in (a, b)$ such that f_{c_0} has a neutral cycle of period k > 1 with multiplier 1, that is, with one petal, there is an open set Z accumulating at c_0 for which $HD(J_c) > 2d/(d+1)$ for $c \in \mathbf{Z}$.

We begin with three pictures. The first two ones illustrate how to choose the parameter (for d=2) and how the corresponding Julia set looks like. The third one has been kindly drawn for us by the referee and shows the case d=4.

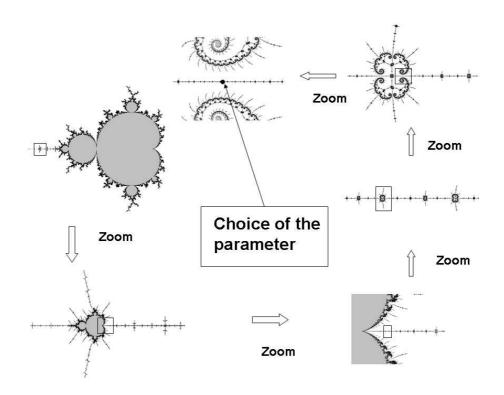


FIGURE 1 – Choice of the parameter

Since $c_0 \in (a, b)$, the set J_c is connected for c in a small neighborhood U of c_0 . It is known, that f_c has an attracting periodic orbit of period k for $c \in U$ on the left side of c_0 . Since k > 1, the corresponding filled-in Julia set K_{c_0} is such that

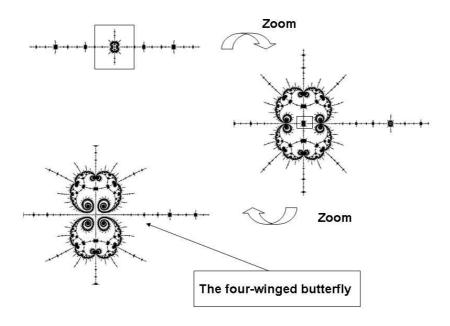


FIGURE 2 – The corresponding Julia set

its interior has a component Δ containing 0, and infinitely many preimages of Δ . The boundary $\partial \Delta$ contains a parabolic point α of period k. In particular, there is a sequence of preimages of Δ , which intersect the real line and accumulate at α .

For definicity, one can assume that $\alpha > 0$. Then $F = f_{c_0}^k$ has the following local form :

$$F(z) = z + a(z - \alpha)^{2} + b(z - \alpha)^{3} + ...,$$
(1)

where a > 0. This implies that a parabolic implosion phenomenon occurs as $\epsilon \to 0$, $\epsilon > 0$, for the maps $f_{c_0+\epsilon}$.

At this point we digress somewhat and describe briefly the theory of parabolic implosion.

Let $\delta > 0$ be very small and D_{\pm} be the disks centered at $\alpha \pm \delta$ with radius δ . The map F sends D_{-} into itself while $D_{+} \subset F(D_{+})$. This defines, after identification of z with F(z) at the boundary, two cylinders $U_{-} = D_{-} \backslash F(D_{-})$ and $U_{+} = F(D_{+}) \backslash D_{+}$. The fact that U_{\pm} are actual cylinders is best seen in the approximate Fatou coordinate $I: z \mapsto -1/(a(z-\alpha))$ which sends α to ∞ and

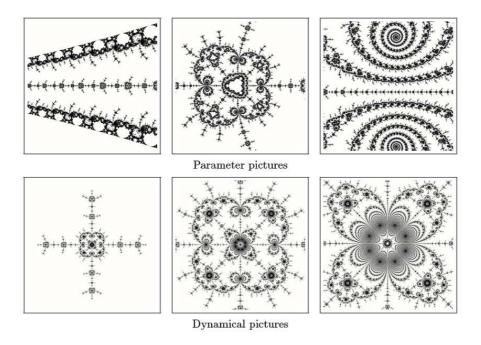


FIGURE 3 – The case d=4

conjugates F to a map which is asymptotically the translation by 1 at ∞ :

$$F_{\infty}(w) = I \circ F \circ I^{-1}(w) = w + 1 + \frac{A}{w} + O(\frac{1}{|w|^2}),$$
 (2)

where $A = 1 - b/a^2$. The real number A is a conformal invariant. In the case of real polynomial F which has a parabolic fixed point α with multiplier 1 and with a single critical point in its immediate basin Δ , it is known [9] that A > 0.

By Riemann mapping theorem these two cylinders may be uniformized by "straight" cylinders. In other words there exists φ_{\pm} mapping the cylinders U_{\pm} to vertical strips V_{\pm} of width 1 conjugating F to the translation by 1. For further use we notice that, by symmetry, c_0 being real, we may assume that $\varphi_{\pm}(\overline{z}) = \overline{\varphi_{\pm}(z)}$. We also notice that these maps are unique up to post-composition by a real translation. These two maps are called respectively repelling (+) and attracting (-) Fatou coordinates. We normalize them as follows. For every $\kappa > 0$ and R > 0, consider two sectors $\Sigma_{-}(\kappa, R) = \{w : Re(w) > R - \kappa |Im(w)|\}$, $\Sigma_{+}(\kappa, R) = \{w : Re(w) < -R + \kappa |Im(w)|\}$. Then for any $\kappa > 0$ there is a big enough $R(\kappa)$, such

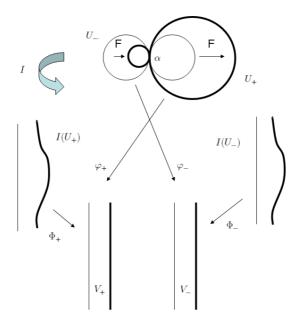


FIGURE 4 – Fatou coordinates

that, if we introduce two sectors $\Sigma_{\pm}(\kappa) = \Sigma_{\pm}(\kappa, R(\kappa))$, then $\varphi_{\pm} = \Phi_{\pm} \circ I$, where

$$\Phi_{\pm}(w) = w - A \log_{\pm}(w) + C_{\pm} + o(1) \tag{3}$$

as $w \to \infty$ within $\Sigma_{\pm}(\kappa)$ respectively. We specify the constants C_{\pm} and the logbranches in such a way, that Φ_{\pm} are real for real w. Namely, we choose $C_{-}=0$ and \log_{-} to be the standard log-branch in the slit plane $\mathbb{C} \setminus \{\mathbf{x} \le \mathbf{0}\}$. In turn, let $C_{+} = iA\pi$ and \log_{+} to be a branch of \log in $\mathbb{C} \setminus \{\mathbf{x} \ge \mathbf{0}\}$ so that $\log_{+}(w) =$ $\log |w| + i\pi$ for w < 0.

Now, we extend φ_{\pm} in the following way. Since

$$\varphi_{-}(F(z)) = T_1(\varphi_{-}(z)) \tag{4}$$

where T_{σ} denotes the translation by σ , and since every orbit converging to α passes through U_{-} exactly once, φ_{-} extends uniquely to Δ to an holomorphic function still satisfying (4).

It is seen from (4), that the map $\varphi_-: \Delta \to \mathbf{C}$ is a branched covering, with the critical points at 0 and all its preimages in Δ by F^n , n > 0, and the critical values at the real numbers $\varphi_-(0) - n$, $n \ge 0$. In particular, there exists a simply-connected domain $\Omega_- \subset \Delta \cap \mathbf{H}^+$, where \mathbf{H}^+ is the upper half-plane, such that $\varphi_-: \Omega_- \to \mathbf{H}^+$ is a holomorphic homeomorphism. Moreover, the intersection of the boundary of Ω_- with \mathbf{R} is the interval $(0, \alpha)$.

Concerning the repelling Fatou coordinate it is best to consider $\psi_+ = \varphi_-^{-1}$: $V_+ \to U_+$. The functionnal relation is now

$$\psi_{+}(T_{1}(z)) = F(\psi_{+}(z)) \tag{5}$$

and we can extend ψ_+ to an entire function by putting, for $n \in \mathbf{Z}$, $\psi_+(T_n(z)) = F^n(\psi_+(z))$ for $z \in V_+$. There exists a simply connected domain $\Omega_+ \subset \mathbf{H}^+$, such that $\psi_+ : \Omega_+ \to \mathbf{H}^+$ is a homeomorphism.

Let now σ be a real number. We define the Lavaurs map g_{σ} on the component Δ of the interior of the filled-in Julia set of f_{c_0} by

$$g_{\sigma} = \psi_{+} \circ T_{\sigma} \circ \varphi_{-}.$$

The "raison d'être" of this definition is the following theorem due to Douady and Lavaurs ([2]):

Theorem 2: There exists a sequence of positive ϵ_n converging to 0 and a sequence of positive integers N_n such that

$$g_{\sigma}(z) = \lim_{n \to \infty} f_{c_0 + \epsilon_n}^{kN_n}(z) \tag{6}$$

uniformly on compact sets of Δ .

Using (3) with the constants C_{\pm} and the log-branches specified as above, it is easy to get, that, for every κ , if w tends to ∞ in $\Sigma(\kappa) := \Sigma_{-}(\kappa) \cap \Sigma_{+}(\kappa) \cap \mathbf{H}^{+}$, then $g_{\infty}(w) := I \circ g_{\sigma} \circ I^{-1}(w) = w + (\sigma - iA\pi) + O(\frac{1}{|w|})$, where A is real and positive. Therefore, for every real σ and every $\kappa > \kappa(\sigma)$, the inverse map g_{∞}^{-1} leaves the sector $\Sigma(\kappa)$ invariant and $g_{\infty}^{-n} \to \infty$ as $n \to \infty$. Coming back to the z-plane, we conclude that the branch $G = I^{-1} \circ g_{\infty}^{-1} \circ I$ of g_{σ}^{-1} leaves the set $S(\kappa) = I^{-1}(\Sigma(\kappa))$ invariant, and $G^{n}(z) \to \alpha$ as $n \to \infty$, for $z \in S(\kappa)$ and every $\kappa > \kappa(\sigma)$. We have, for $w \in \Sigma(\kappa)$:

$$G_{\infty}(w) := I \circ G \circ I^{-1}(w) = w + (-\sigma + iA\pi) + O(\frac{1}{|w|}),$$
 (7)

Now, from the definition of $g_{\sigma}(z)$ and the global properties of the maps φ_{-} and ψ_{+} , it follows the existence of a simply-connected domain $\Omega_{0} \subset \Omega_{-}$, which is mapped by g_{σ} homeomorphically onto \mathbf{H}^{+} and such that $\alpha \in \bar{\Omega}_{0}$. Moreover, from the above description, $S(\kappa) \subset \Omega_{0}$, for every $\kappa > \kappa(\sigma)$. Therefore, the branch G of g^{-1} which is defined above, extends to a global univalent branch $G: \mathbf{H}^{+} \to \Omega_{0}$ of g_{σ}^{-1} . Since $\Omega_{0} \subset \mathbf{H}^{+}$, the iterates $G^{n}: \mathbf{H}^{+} \to \Omega_{0}$, n > 0, converge uniformly on compact sets in \mathbf{H}^{+} to a unique fixed point in $\bar{\Omega}_{0}$, which must be α .

Let us consider the continuous map $\sigma \mapsto g_{\sigma}(0) = \psi_{+}(\varphi_{-}(0) + \sigma)$: if σ runs in the interval $I = \{\varphi_{+}(x) - \varphi_{-}(0) : x \in U_{+} \cap \mathbf{R}\}$, then $g_{\sigma}(0)$ runs over $U_{+} \cap \mathbf{R}$. It is thus clear, and this is the key point in the proof, that we can choose σ in such a way that $g_{\sigma}(0)$ is a preimage of α : there is $j \geq 1$, such that $f_{c_0}^j \circ g_{\sigma}(0) = \alpha$. Since 0 is a critical point for g_{σ} , taking the inverse image by $f_{c_0}^j \circ g_{\sigma}$ has the same effect as multiplying the number of petals by d and we may state:

Lemma 2.1 There exists an infinite iterated function system defined on a small compact neighborhood B_0 of zero and generated by some holomorphic branches of $f_{c_0}^{-1}$ and g_{σ}^{-1} such that its limit set has Hausdorff dimension bigger than 2d/(d+1).

Comment 3 In fact, the limit set is a subset of a so-called Julia-Lavaurs set denoted by $J_{c_0,\sigma}$. It is defined as follows. The map g_{σ} can be extended in a natural way from Δ to the interior of the filled-in Fatou set of f_{c_0} : if $f_{c_0}^k(z) \in \Delta$, set $g_{\sigma}(z) = g_{\sigma} \circ f_{c_0}^k(z)$. Then $J_{c_0,\sigma}$ is simply the closure of the set of points z for which there exists $m \in \mathbb{N}$ such that $g_{\sigma}^m(z)$ is defined and belongs to $J(f_0)$.

This lemma together with the above discussion implies the theorem. Indeed, by a general property of iterated function systems [7], there exists its finite subsystem with the Hausdorff dimension of its limit set bigger than 2d/(d+1). On the other hand, the finite iterated function system persists for $f_{c_0+\epsilon}$ by (6). To be more precise, if $\{I_j: B_0 \to X_j, 1 \le j \le j_0\}$ is this finite iterated function system, then each I_j can be extended to a univelent map to a fixed neighborhood Y of B_0 as $I_j: Y \to Y_j$. Consider the inverse univalent map $I_j^{-1}: Y_j \to Y$. Since the convergence in (6) is uniform on compacts in Δ , for every ϵ_n small enough there is some integer $N_j > 0$ and a compact set $X_{j,n}$, so that $f_{c_0+\epsilon_n}^{N_j}: X_{j,n} \to B_0$ is univalent, too. Now it is clear, that, for every ϵ_n small enough, the non-escaping set K_n of the dynamical system which consists of a finitely many maps $f_{c_0+\epsilon_n}^{N_j}: X_{j,n} \to B_0$, $1 \le j \le j_0$, has the Hausdorff dimension which is bigger than 2d/(d+1). On the other hand, K_n must lie in the Julia set of $f_{c_0+\epsilon_n}$ because some iterate of the map $f_{c_0+\epsilon_n}^{N_1N_2...N_{j_0}}$ leaves the set K_n invariant and is expanding on it.

Proof of Lemma 2.1. As the first step, let us fix a small enough closed ball B_0 around zero, so that it does not contain points of the postcritical set of f_{c_0} . There exists its preimage B' by F^{-1} in $\Delta \cap \mathbf{H}^+$. Then we can apply to B' the maps G^n , n > 0. By the above, $B'_n = G^n(B')$ are pairwise disjoint, compactly contained in Δ , and $B'_n \to \alpha$ as $n \to \infty$. Now, for every $n \geq n_0$, so that B'_n lies in a small enough neighborhood U of α , we make "clones" of B'_n in $U \cap \Delta$ applying to it F^r , $r \in \mathbf{Z}$, where F^r for r < 0 is a well-defined in $U \cap \Delta$ branch which fixes α . We obtain the sets $B'_{n,r} = F^r(B'_n)$. On the second step, we consider the map $f_{c_0}^{j} \circ g_{\sigma}$ from a neighborhood V of 0 onto U, This map is a remified cover with the only remification point at 0 of order d. Let $U^* = U \setminus \{x \geq \alpha\}$, and $V^* = V \cap \{z : Arg(z) \in (0, 2\pi/d)\}$. Denote by h a branch of $(f_{c_0}^j \circ g_\sigma)^{-1}$ from U^* onto V^* . Let $B_{n,r} = h(B'_{n,r})$. We obtain a system of holomorphic maps $\Psi = \{\psi_{n,r} : B_0 \to B_{n,r}\}, \text{ where } \psi_{n,r} = h \circ F^r \circ G^n \circ F^{-1}, r \in \mathbf{Z}, n \geq n_0. \text{ If the } \mathbf{Z} \in \mathbf{Z}, n \geq n_0. \text{ If the } \mathbf{Z} \in \mathbf{Z}, n \geq n_0. \text{ If the } \mathbf{Z} \in \mathbf{Z}, n \geq n_0. \text{ If the } \mathbf{Z} \in \mathbf{Z}, n \geq n_0. \text{ If the } \mathbf{Z} \in \mathbf{Z}, n \geq n_0. \text{ If the } \mathbf{Z} \in \mathbf{Z}, n \geq n_0. \text{ If the } \mathbf{Z} \in \mathbf{Z}, n \geq n_0. \text{ If the } \mathbf{Z} \in \mathbf{Z}, n \geq n_0. \text{ If the } \mathbf{Z} \in \mathbf{Z}, n \geq n_0. \text{ If the } \mathbf{Z} \in \mathbf{Z}, n \geq n_0. \text{ If the } \mathbf{Z} \in \mathbf{Z}, n \geq n_0. \text{ If the } \mathbf{Z} \in \mathbf{Z}, n \geq n_0. \text{ If } \mathbf{Z$ neighborhood U is chosen small enough, the maps $\psi_{n,r}$ extend to univalent maps in a fixed neighborhood B of B_0 into itself. In particular, the compact sets $B_{n,r}$ are pairwise disjoint and compactly contained in B_0 . Now, it is quite standard to check that Ψ form a conformal infinite iterated function system in the sense of [7] (strictly speaking, in the hyperbolic metric of B, which is equivalent to the Eucledian one

on B_0 though). Let us calculate the parameter $\theta = \inf\{t : \psi(t) < \infty\}$ of Ψ , where $\psi(t) = \sum_{(n,r)} \max_{z \in B_0} |\psi'_{n,r}(z)|^t$. The map $\psi_{n,r} = h \circ I^{-1} \circ F_\infty^r \circ G_\infty^n \circ I \circ F^{-1}$. Here F^{-1} is a univalent map of a neighborhood \tilde{B} of B_0 into Δ . Now, routine and well-know calculations based on (2)-(7) show (see e.g. [10]), that, for all $r \in \mathbf{Z}, n \geq n_0$ and some C, which depends only on a compact set in $\mathbf{H}^+ \cap \Delta$, from which w is taken, $C^{-1}|r + (-\sigma + i\pi A)n| \leq |F_\infty^r \circ G_\infty^n(w)| \leq C|r + (-\sigma + i\pi A)n|$, and $C^{-1} \leq |(F_\infty^r \circ G_\infty^n)'(w)| \leq C$. On the other hand, the map h is a composition of a univalent map with an inverse branch of $z^{1/d}$. This gives us : $C_1^{-1}|r + (-\sigma + i\pi A)n|^{-1-1/d} \leq |\psi'_{n,r}(z)| \leq C_1|r + (-\sigma + i\pi A)n|^{-1-1/d}$, for some C_1 and every $z \in B_0$. It follows that the series for $\psi(t)$ converges if and only if $t > \theta = 2d/(d+1)$, and $\psi(\theta) = \infty$. Hence [7], the Hausdorff dimension of the limit set of Ψ is strictly bigger than 2d/(d+1).

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