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► **To cite this version:**

Mikhail Isaev, Roman Novikov. Stability estimates for determination of potential from the impedance boundary map. 2012. hal-00652562v2

HAL Id: hal-00652562

<https://hal.science/hal-00652562v2>

Preprint submitted on 6 Feb 2012 (v2), last revised 22 Feb 2012 (v3)

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Stability estimates for determination of potential from the impedance boundary map

M.I. Isaev and R.G. Novikov

Abstract

We study the impedance boundary map (or Robin-to-Robin map) for the Schrödinger equation in open bounded domain at fixed energy in multidimensions. We give global stability estimates for determining potential from these boundary data and, as corollary, from the Cauchy data set. Our results include also, in particular, an extension of the Alessandrini identity to the case of the impedance boundary map.

1. Introduction

We consider the Schrödinger equation

$$-\Delta\psi + v(x)\psi = E\psi, \quad x \in D, \quad E \in \mathbb{R}, \quad (1.1)$$

where

$$D \text{ is an open bounded domain in } \mathbb{R}^d, \quad d \geq 2, \quad (1.2)$$
$$\text{with } \partial D \in C^2,$$

$$v \in L^\infty(D). \quad (1.3)$$

We consider the impedance boundary map $\hat{M}_\alpha = \hat{M}_{\alpha,v}(E)$ defined by

$$\hat{M}_\alpha[\psi]_\alpha = [\psi]_{\alpha-\pi/2} \quad (1.4)$$

for all sufficiently regular solutions ψ of equation (1.1) in $\bar{D} = D \cup \partial D$, where

$$[\psi]_\alpha = [\psi(x)]_\alpha = \cos \alpha \psi(x) - \sin \alpha \frac{\partial \psi}{\partial \nu} |_{\partial D}(x), \quad x \in \partial D, \quad \alpha \in \mathbb{R} \quad (1.5)$$

and ν is the outward normal to ∂D . One can show (see Lemma 3.2) that there is not more than a countable number of $\alpha \in \mathbb{R}$ such that E is an eigenvalue for the operator $-\Delta + v$ in D with the boundary condition

$$\cos \alpha \psi |_{\partial D} - \sin \alpha \frac{\partial \psi}{\partial \nu} |_{\partial D} = 0. \quad (1.6)$$

Therefore, for any energy level E we can assume that for some fixed $\alpha \in \mathbb{R}$

$$E \text{ is not an eigenvalue for the operator } -\Delta + v \text{ in } D \quad (1.7)$$
$$\text{with boundary condition (1.6)}$$

and, as a corollary, \hat{M}_α can be defined correctly.

Note that the impedance boundary map \hat{M}_α is reduced to the Dirichlet-to-Neumann (DtN) map if $\alpha = 0$ and is reduced to the Neumann-to-Dirichlet (NtD) map if $\alpha = \pi/2$. The map \hat{M}_α

can be called also as the Robin-to-Robin map. General Robin-to-Robin map was considered, in particular, in [9].

We consider the following inverse boundary value problem for equation (1.1).

Problem 1.1. Given \hat{M}_α for some fixed E and α , find v .

This problem can be considered as the Gel'fand inverse boundary value problem for the Schrödinger equation at fixed energy (see [8], [16]). At zero energy this problem can be considered also as a generalization of the Calderon problem of the electrical impedance tomography (see [5], [16]).

Problem 1.1 includes, in particular, the following questions: (a) uniqueness, (b) reconstruction, (c) stability.

Global uniqueness theorems and global reconstruction methods for Problem 1.1 with $\alpha = 0$ were given for the first time in [16] in dimension $d \geq 3$ and in [4] in dimension $d = 2$.

Global stability estimates for Problem 1.1 with $\alpha = 0$ were given for the first time in [1] in dimension $d \geq 3$ and in [23] in dimension $d = 2$. A principal improvement of the result of [1] was given recently in [21] (for the zero energy case). Due to [14] these logarithmic stability results are optimal (up to the value of the exponent). An extension of the instability estimates of [14] to the case of the non-zero energy as well as to the case of Dirichlet-to-Neumann map given on the energy intervals was given in [11].

Note also that for the Calderon problem (of the electrical impedance tomography) in its initial formulation the global uniqueness was firstly proved in [27] for $d \geq 3$ and in [15] for $d = 2$.

It should be noted that in most of previous works on inverse boundary value problems for equation (1.1) at fixed E it was assumed in one way or another that E is not a Dirichlet eigenvalue for the operator $-\Delta + v$ in D , see [1], [14], [16], [21], [23], [24], [25]. Nevertheless, the results of [4] can be considered as global uniqueness and reconstruction results for Problem 1.1 in dimension $d = 2$ with general α .

In the present work we give global stability estimates for Problem 1.1 in dimension $d \geq 2$ with general α . These results are presented in detail in Section 2.

In addition, in the present work we establish some basic properties of the impedance boundary map with general α . In particular, we extend the Alessandrini identity to this general case. These results are presented in detail in Section 3.

In a subsequent paper we plan to give also global reconstruction method for Problem 1.1 in multidimensions with general α .

2. Stability estimates

In this section we always assume that D satisfies (1.2).

We will use the fact that if v_1, v_2 are potentials satisfying (1.3), (1.7) for some fixed E and α , then

$$\hat{M}_{\alpha, v_1}(E) - \hat{M}_{\alpha, v_2}(E) \text{ is a bounded operator in } \mathbb{L}^\infty(\partial D), \quad (2.1)$$

where $\hat{M}_{\alpha, v_1}(E), \hat{M}_{\alpha, v_2}(E)$ denote the impedance boundary maps for v_1, v_2 , respectively. Actually, under our assumptions, $\hat{M}_{\alpha, v_1}(E) - \hat{M}_{\alpha, v_2}(E)$ is a compact operator in $\mathbb{L}^\infty(\partial D)$ (see, Corollary 3.1).

Let

$$\begin{aligned} \|A\| \text{ denote the norm of an operator} \\ A : \mathbb{L}^\infty(\partial D) \rightarrow \mathbb{L}^\infty(\partial D). \end{aligned} \quad (2.2)$$

Let the Cauchy data set \mathcal{C}_v for equation (1.1) be defined by:

$$\mathcal{C}_v = \left\{ \left(\psi|_{\partial D}, \frac{\partial \psi}{\partial \nu} |_{\partial D} \right) : \begin{array}{l} \text{for all sufficiently regular solutions } \psi \text{ of} \\ \text{equation (1.1) in } \bar{D} = D \cup \partial D \end{array} \right\}. \quad (2.3)$$

In addition, the Cauchy data set \mathcal{C}_v can be represented as the graph of the impedance boundary map $\hat{M}_\alpha = \hat{M}_{\alpha,v}(E)$ defined by (1.4) under assumptions (1.7).

2.1. Estimates for $d \geq 3$

In this subsection we assume for simplicity that

$$v \in W^{m,1}(\mathbb{R}^d) \text{ for some } m > d, \text{ supp } v \subset D, \quad (2.4)$$

where

$$W^{m,1}(\mathbb{R}^d) = \{v : \partial^J v \in L^1(\mathbb{R}^d), |J| \leq m\}, \quad m \in \mathbb{N} \cup 0, \quad (2.5)$$

where

$$J \in (\mathbb{N} \cup 0)^d, |J| = \sum_{i=1}^d J_i, \quad \partial^J v(x) = \frac{\partial^{|J|} v(x)}{\partial x_1^{J_1} \dots \partial x_d^{J_d}}. \quad (2.6)$$

Let

$$\|v\|_{m,1} = \max_{|J| \leq m} \|\partial^J v\|_{L^1(\mathbb{R}^d)}. \quad (2.7)$$

Note also that (2.4) \Rightarrow (1.3).

Theorem 2.1. *Let D satisfy (1.2), where $d \geq 3$. Let v_1, v_2 satisfy (2.4) and (1.7) for some fixed E and α . Let $\|v_j\|_{m,1} \leq N$, $j = 1, 2$, for some $N > 0$. Let $\hat{M}_{\alpha,v_1}(E)$ and $\hat{M}_{\alpha,v_2}(E)$ denote the impedance boundary maps for v_1 and v_2 , respectively. Then*

$$\|v_1 - v_2\|_{L^\infty(D)} \leq C_\alpha (\ln(3 + \delta_\alpha^{-1}))^{-s}, \quad 0 < s \leq (m - d)/m, \quad (2.8)$$

where $C_\alpha = C_\alpha(N, D, m, s, E)$, $\delta_\alpha = \|\hat{M}_{\alpha,v_1}(E) - \hat{M}_{\alpha,v_2}(E)\|$ is defined according to (2.2).

Remark 2.1. Estimate (2.8) with $\alpha = 0$ is a variation of the result of [1] (see also [21]).

Proof of Theorem 2.1 is given in Section 5. This proof is based on results presented in Sections 3, 4.

Theorem 2.1 implies the following corollary:

Corollary 2.1. *Let D satisfy (1.2), where $d \geq 3$. Let potentials v_1, v_2 satisfy (2.4). Then*

$$\|v_1 - v_2\|_{L^\infty(D)} \leq \min_{\alpha \in \mathbb{R}} C_\alpha (\ln(3 + \delta_\alpha^{-1}))^{-s}, \quad 0 < s \leq (m - d)/m, \quad (2.9)$$

where C_α and δ_α at fixed α are the same that in Theorem 2.1.

Actually, Corollary 2.1 can be considered as global stability estimate for determining potential v from its Cauchy data set \mathcal{C}_v for equation (1.1) at fixed energy E , where $d \geq 3$.

2.2. Estimates for $d = 2$

In this subsection we assume for simplicity that

$$v \in C^2(\bar{D}), \quad \text{supp } v \subset D. \quad (2.10)$$

Note also that (2.10) \Rightarrow (1.3).

Theorem 2.2. *Let D satisfy (1.2), where $d = 2$. Let v_1, v_2 satisfy (2.10) and (1.7) for some fixed E and α . Let $\|v_j\|_{C^2(\bar{D})} \leq N$, $j = 1, 2$, for some $N > 0$. Let $\hat{M}_{\alpha, v_1}(E)$ and $\hat{M}_{\alpha, v_2}(E)$ denote the impedance boundary maps for v_1 and v_2 , respectively. Then*

$$\|v_1 - v_2\|_{L^\infty(D)} \leq C_\alpha (\ln(3 + \delta_\alpha^{-1}))^{-s} (\ln(3 \ln(3 + \delta_\alpha^{-1})))^2, \quad 0 < s \leq 3/4, \quad (2.11)$$

where $C_\alpha = C_\alpha(N, D, s, E)$, $\delta_\alpha = \|\hat{M}_{\alpha, v_1}(E) - \hat{M}_{\alpha, v_2}(E)\|$ is defined according to (2.2).

Remark 2.2. Theorem 2.2 for $\alpha = 0$ was given in [23] with $s = 1/2$ and in [25] with $s = 3/4$.

Proof of Theorem 2.2 is given in Section 7. This proof is based on results presented in Sections 3, 6.

Theorem 2.2 implies the following corollary:

Corollary 2.2. *Let D satisfy (1.2), where $d = 2$. Let potentials v_1, v_2 satisfy (2.10). Then*

$$\|v_1 - v_2\|_{L^\infty(D)} \leq \min_{\alpha \in \mathbb{R}} C_\alpha (\ln(3 + \delta_\alpha^{-1}))^{-s} (\ln(3 \ln(3 + \delta_\alpha^{-1})))^2, \quad 0 < s \leq 3/4, \quad (2.12)$$

where C_α and δ_α at fixed α are the same that in Theorem 2.2.

Actually, Corollary 2.2 can be considered as global stability estimate for determining potential v from its Cauchy data set \mathcal{C}_v for equation (1.1) at fixed energy E , where $d = 2$.

2.3. Concluding remarks

Theorems 2.1, 2.2 and Lemma 3.2 imply the following corollary:

Corollary 2.3. *Under assumptions (1.2), (1.3), real-valued potential v is uniquely determined by its Cauchy data \mathcal{C}_v at fixed real energy E .*

To our knowledge the result of Corollary 2.3 for $d \geq 3$ was not yet completely proved in the literature.

Let $\sigma_{\alpha, v}$ denote the spectrum of the operator $-\Delta + v$ in D with boundary condition (1.6).

Remark 2.3. In Theorems 2.1 and 2.2 we do not assume that $E \notin \sigma_{\alpha, v_1} \cup \sigma_{\alpha, v_2}$ namely for $\alpha = 0$ in contrast with [1], [21], [23], [24], [25]. In addition, in fact, in Corollaries 2.1 and 2.2 there are no special assumptions on E and α at all. Actually, the stability estimates of [1], [21], [23], [24], [25] make no sense for $E \in \sigma_{0, v_1} \cup \sigma_{0, v_2}$ and are too weak if $\text{dist}(E, \sigma_{0, v_1} \cup \sigma_{0, v_2})$ is too small.

Remark 2.4. The stability estimates of Subsections 2.1 and 2.2 admit principal improvement in the sense described in [21], [22], [26]. In particular, Theorem 2.1 with $s = m - d$ (for $d = 3$ and $E = 0$) follows from results presented in Sections 3, 4 of the present work and results presented in Section 8 of [21].

Remark 2.5. The stability estimates of Subsections 2.1 and 2.2 can be extended to the case when we do not assume that $\text{supp } v \subset D$ or, by other words, that v is zero near the boundary. In this connection see, for example, [1], [23].

In the present work we do not develop Remarks 2.4 and 2.5 in detail because of restrictions in time.

Note also that Theorems 2.1 and 2.2 remain valid with complex-valued potentials v_1, v_2 and complex E, α . Finally, we note that in Theorems 2.1, 2.2 and Corollaries 2.1, 2.2 with real α constant C_α can be considered as independent of α .

3. Some basic properties of the impedance boundary map

Lemma 3.1. *Let D satisfy (1.2). Let potential v satisfy (1.3) and (1.7) for some fixed E and α . Let $\hat{M}_\alpha = \hat{M}_{\alpha,v}(E)$ denote the impedance boundary map for v . Then*

$$\begin{aligned} (\sin \alpha \hat{M}_\alpha + \cos \alpha \hat{I}) [\psi]_\alpha &= \psi|_{\partial D}, \\ (\cos \alpha \hat{M}_\alpha - \sin \alpha \hat{I}) [\psi]_\alpha &= \frac{\partial \psi}{\partial \nu}|_{\partial D}, \end{aligned} \quad (3.1)$$

$$\int_{\partial D} [\psi^{(1)}]_\alpha \hat{M}_\alpha [\psi^{(2)}]_\alpha dx = \int_{\partial D} [\psi^{(2)}]_\alpha \hat{M}_\alpha [\psi^{(1)}]_\alpha dx \quad (3.2)$$

for all sufficiently regular solutions $\psi, \psi^{(1)}, \psi^{(2)}$ of equation (1.1) in \bar{D} , where \hat{I} denotes the identity operator on ∂D and $[\psi]_\alpha$ is defined by (1.5).

Note also that identities (3.1) imply that

$$\left(\sin(\alpha_1 - \alpha_2) \hat{M}_{\alpha_1} + \cos(\alpha_1 - \alpha_2) \hat{I} \right) \left(\sin(\alpha_2 - \alpha_1) \hat{M}_{\alpha_2} + \cos(\alpha_2 - \alpha_1) \hat{I} \right) = \hat{I}, \quad (3.3)$$

under the assumptions of Lemma 3.1 fulfilled simultaneously for $\alpha = \alpha_1$ and $\alpha = \alpha_2$.

Proof of Lemma 3.1. Identities (3.1) follow from definition (1.4) of the map \hat{M}_α .

To prove (3.2) we use, in particular, the Green formula

$$\int_{\partial D} \left(\phi^{(1)} \frac{\partial \phi^{(2)}}{\partial \nu} - \phi^{(2)} \frac{\partial \phi^{(1)}}{\partial \nu} \right) dx = \int_D \left(\phi^{(1)} \Delta \phi^{(2)} - \phi^{(2)} \Delta \phi^{(1)} \right) dx, \quad (3.4)$$

where $\phi^{(1)}$ and $\phi^{(2)}$ are arbitrary sufficiently regular functions in \bar{D} . Using (3.4) and the identities

$$\psi^{(1)} \Delta \psi^{(2)} = (v - E) \psi^{(1)} \psi^{(2)} = \psi^{(2)} \Delta \psi^{(1)} \quad \text{in } D, \quad (3.5)$$

we obtain that

$$\int_{\partial D} \left(\psi^{(1)} \frac{\partial \psi^{(2)}}{\partial \nu} - \psi^{(2)} \frac{\partial \psi^{(1)}}{\partial \nu} \right) dx = 0. \quad (3.6)$$

Using (3.6), we get that

$$\begin{aligned} & \int_{\partial D} \left(\cos \alpha \psi^{(1)} - \sin \alpha \frac{\partial \psi^{(1)}}{\partial \nu} \right) \left(\sin \alpha \psi^{(2)} + \cos \alpha \frac{\partial \psi^{(2)}}{\partial \nu} \right) dx = \\ & = \int_{\partial D} \left(\cos \alpha \psi^{(2)} - \sin \alpha \frac{\partial \psi^{(2)}}{\partial \nu} \right) \left(\sin \alpha \psi^{(1)} + \cos \alpha \frac{\partial \psi^{(1)}}{\partial \nu} \right) dx. \end{aligned} \quad (3.7)$$

Identity (3.2) follows from (3.7) and definition (1.4) of the map \hat{M}_α . ■

Theorem 3.1. *Let D satisfy (1.2). Let two potentials v_1, v_2 satisfy (1.3), (1.7) for some fixed E and α . Let $\hat{M}_{\alpha, v_1} = \hat{M}_{\alpha, v_1}(E)$, $\hat{M}_{\alpha, v_2} = \hat{M}_{\alpha, v_2}(E)$ denote the impedance boundary maps for v_1, v_2 , respectively. Then*

$$\int_D (v_1 - v_2) \psi_1 \psi_2 dx = \int_{\partial D} [\psi_1]_\alpha \left(\hat{M}_{\alpha, v_1} - \hat{M}_{\alpha, v_2} \right) [\psi_2]_\alpha dx \quad (3.8)$$

for all sufficiently regular solutions ψ_1 and ψ_2 of equation (1.1) in \bar{D} with $v = v_1$ and $v = v_2$, respectively, where $[\psi]_\alpha$ is defined by (1.5).

Proof of Theorem 3.1. As in (3.5) we have that

$$\begin{aligned} \psi_1 \Delta \psi_2 &= (v_2 - E) \psi_1 \psi_2, \\ \psi_2 \Delta \psi_1 &= (v_1 - E) \psi_1 \psi_2. \end{aligned} \quad (3.9)$$

Combining (3.9) with (3.4), (3.1) and (3.2), we obtain that

$$\begin{aligned} & \int_D (v_1(x) - v_2(x)) \psi_1(x) \psi_2(x) dx = \int_{\partial D} \left(\psi_2 \frac{\partial \psi_1}{\partial \nu} - \psi_1 \frac{\partial \psi_2}{\partial \nu} \right) dx = \\ &= \int_{\partial D} \left(\sin \alpha \hat{M}_{\alpha, v_2} + \cos \alpha \hat{I} \right) [\psi_2]_\alpha \left(\cos \alpha \hat{M}_{\alpha, v_1} - \sin \alpha \hat{I} \right) [\psi_1]_\alpha dx - \\ & - \int_{\partial D} \left(\sin \alpha \hat{M}_{\alpha, v_1} + \cos \alpha \hat{I} \right) [\psi_1]_\alpha \left(\cos \alpha \hat{M}_{\alpha, v_2} - \sin \alpha \hat{I} \right) [\psi_2]_\alpha dx = \\ &= \int_{\partial D} [\psi_1]_\alpha \left(\hat{M}_{\alpha, v_1} - \hat{M}_{\alpha, v_2} \right) [\psi_2]_\alpha dx. \end{aligned} \quad (3.10)$$

■

Remark 3.1. Identity (3.8) for $\alpha = 0$ is reduced to Alessandrini's identity (Lemma 1 of [1]).

Let $G_\alpha(x, y, E)$ be the Green function for the operator $\Delta - v + E$ in D with the impedance boundary condition (1.6) under assumptions (1.2), (1.3) and (1.7). Note that

$$G_\alpha(x, y, E) = G_\alpha(y, x, E), \quad x, y \in \bar{D}. \quad (3.11)$$

The symmetry (3.11) is proved in Section 9.

Theorem 3.2. *Let D satisfy (1.2). Let potential v satisfy (1.3) and (1.7) for some fixed E and α such that $\sin \alpha \neq 0$. Let $G_\alpha(x, y, E)$ be the Green function for the operator $\Delta - v + E$ in D with the impedance boundary condition (1.6). Then for $x, y \in \partial D$*

$$M_\alpha(x, y, E) = \frac{1}{\sin^2 \alpha} G_\alpha(x, y, E) - \operatorname{ctg} \alpha \delta_{\partial D}(x - y), \quad (3.12)$$

where $M_\alpha(x, y, E)$ and $\delta_{\partial D}(x - y)$ denote the Schwartz kernels of the impedance boundary map $\hat{M}_\alpha = \hat{M}_{\alpha, v}(E)$ and the identity operator \hat{I} on ∂D , respectively, where \hat{M}_α and \hat{I} are considered as linear integral operators.

Proof of Theorem 3.2. Note that

$$[\phi]_{\alpha-\pi/2} = \frac{1}{\sin^2 \alpha} \sin \alpha \phi|_{\partial D} - \operatorname{ctg} \alpha [\phi]_{\alpha}. \quad (3.13)$$

for all sufficiently regular functions ϕ in some neighbourhood of ∂D in D . Since G_{α} is the Green function for equation (1.1) we have that

$$\psi(y) = \int_{\partial D} \left(\psi(x) \frac{\partial G_{\alpha}}{\partial \nu_x}(x, y, E) - G_{\alpha}(x, y, E) \frac{\partial \psi}{\partial \nu}(x) \right) dx, \quad y \in D, \quad (3.14)$$

for all sufficiently regular solutions ψ of equation (1.1). Using (3.14) and impedance boundary condition (1.6) for G_{α} , we get that

$$\begin{aligned} \sin \alpha \psi(y) &= \sin \alpha \int_{\partial D} \left(\psi(x) \frac{\partial G_{\alpha}}{\partial \nu_x}(x, y, E) - G_{\alpha}(x, y, E) \frac{\partial \psi}{\partial \nu}(x) \right) dx = \\ &= \int_{\partial D} [\psi(x)]_{\alpha} G_{\alpha}(x, y, E) dx, \quad y \in D. \end{aligned} \quad (3.15)$$

Due to (3.2) we have that

$$M_{\alpha}(x, y, E) = M_{\alpha}(y, x, E), \quad x, y \in \partial D. \quad (3.16)$$

Combining (1.4), (3.13), (3.15) and (3.16), we obtain (3.12). \blacksquare

Corollary 3.1. *Let assumptions of Theorem 3.1 hold. Then*

$$\hat{M}_{\alpha, v_1}(E) - \hat{M}_{\alpha, v_2}(E) \text{ is a compact operator in } \mathbb{L}^{\infty}(\partial D). \quad (3.17)$$

Scheme of the proof of Corollary 3.1. Let $G_{\alpha, v_1}(x, y, E)$ and $G_{\alpha, v_2}(x, y, E)$ be the Green functions for the operator $\Delta - v + E$ in D with the impedance boundary condition (1.6) for $v = v_1$ and $v = v_2$, respectively. Using (3.11), we find that

$$\begin{aligned} G_{\alpha, v_1}(x, y, E) &= \int_D G_{\alpha, v_1}(x, \xi, E) (\Delta_{\xi} - v_2(\xi) + E) G_{\alpha, v_2}(\xi, y, E) d\xi, \\ G_{\alpha, v_2}(x, y, E) &= \int_D (\Delta_{\xi} - v_1(\xi) + E) G_{\alpha, v_1}(x, \xi, E) G_{\alpha, v_2}(\xi, y, E) d\xi, \\ \int_{\partial D} \left(G_{\alpha, v_1}(x, \xi, E) \frac{\partial G_{\alpha, v_2}}{\partial \nu_{\xi}}(\xi, y, E) - G_{\alpha, v_2}(\xi, y, E) \frac{\partial G_{\alpha, v_1}}{\partial \nu_{\xi}}(x, \xi, E) \right) d\xi &= 0, \end{aligned} \quad (3.18)$$

$x, y \in D.$

Combining (3.18) with (3.4), we get that

$$\begin{aligned} G_{\alpha, v_1}(x, y, E) - G_{\alpha, v_2}(x, y, E) &= \int_D (v_1(\xi) - v_2(\xi)) G_{\alpha, v_1}(x, \xi, E) G_{\alpha, v_2}(\xi, y, E) d\xi, \\ & \quad x, y \in D. \end{aligned} \quad (3.19)$$

The proof of (3.17) for the case of $\sin \alpha \neq 0$ can be completed proceeding from (3.3), (3.12), (3.19) and estimates of [12] and [3] on $G_\alpha(x, y, E)$ for $v \equiv 0$.

Corollary 3.1 for the Dirichlet-to-Neumann case ($\sin \alpha = 0$) was given in [16]. \blacksquare

Lemma 3.2. *Let D satisfy (1.2). Let v be a real-valued potential satisfying (1.3). Then for any fixed $E \in \mathbb{R}$ there is not more than countable number of $\alpha \in \mathbb{R}$ such that E is an eigenvalue for the operator $-\Delta + v$ in D with boundary condition (1.6).*

Proof of Lemma 3.2. Let $\psi^{(1)}, \psi^{(2)}$ be eigenfunctions for the operator $-\Delta + v$ in D with boundary condition (1.6) for $\alpha = \alpha^{(1)}$ and $\alpha = \alpha^{(2)}$, respectively. Then

$$\sin(\alpha^{(1)} - \alpha^{(2)}) \int_{\partial D} \psi^{(1)} \psi^{(2)} dx = \sin \alpha^{(1)} \sin \alpha^{(2)} \int_{\partial D} \left(\psi^{(1)} \frac{\partial \psi^{(2)}}{\partial \nu} - \psi^{(2)} \frac{\partial \psi^{(1)}}{\partial \nu} \right) dx = 0. \quad (3.20)$$

Since in the separable space $\mathbb{L}^2(\partial D)$ there is not more than countable orthogonal system of functions, we obtain the assertion of Lemma 3.2. \blacksquare

Remark 3.1 The assertion of Lemma 3.2 remains valid for the case of $\alpha \in \mathbb{C}$.

4. Faddeev functions

We consider the Faddeev functions G, ψ, h (see [6], [7], [10], [16]):

$$\psi(x, k) = e^{ikx} + \int_{\mathbb{R}^d} G(x - y, k) v(y) \psi(y, k) dy, \quad (4.1)$$

$$G(x, k) = e^{ikx} g(x, k), \quad g(x, k) = -(2\pi)^{-d} \int_{\mathbb{R}^d} \frac{e^{i\xi x} d\xi}{\xi^2 + 2k\xi}, \quad (4.2)$$

where $x \in \mathbb{R}^d, k \in \mathbb{C}^d, \operatorname{Im} k \neq 0, d \geq 3$,

$$h(k, l) = (2\pi)^{-d} \int_{\mathbb{R}^d} e^{-ilx} v(x) \psi(x, k) dx, \quad (4.3)$$

where

$$k, l \in \mathbb{C}^d, k^2 = l^2, \operatorname{Im} k = \operatorname{Im} l \neq 0. \quad (4.4)$$

One can consider (4.1), (4.3) assuming that

$$v \text{ is a sufficiently regular function on } \mathbb{R}^d \text{ with sufficient decay at infinity.} \quad (4.5)$$

For example, in connection with Problem 1.1, one can consider (4.1), (4.3) assuming that

$$v \in \mathbb{L}^\infty(D), \quad v \equiv 0 \text{ on } \mathbb{R} \setminus D. \quad (4.6)$$

We recall that (see [6], [7], [10], [16]):

- The function G satisfies the equation

$$(\Delta + k^2)G(x, k) = \delta(x), \quad x \in \mathbb{R}^d, \quad k \in \mathbb{C}^d \setminus \mathbb{R}^d; \quad (4.7)$$

- Formula (4.1) at fixed k is considered as an equation for

$$\psi = e^{ikx} \mu(x, k), \quad (4.8)$$

where μ is sought in $\mathbb{L}^\infty(\mathbb{R}^d)$;

- As a corollary of (4.1), (4.2), (4.7), ψ satisfies (1.1) for $E = k^2$;
- The Faddeev functions G, ψ, h are (non-analytic) continuation to the complex domain of functions of the classical scattering theory for the Schrödinger equation (in particular, h is a generalized „scattering“ amplitude).

In addition, G, ψ, h in their zero energy restriction, that is for $E = 0$, were considered for the first time in [2]. The Faddeev functions G, ψ, h were, actually, rediscovered in [2].

Let

$$\begin{aligned} \Sigma_E &= \{k \in \mathbb{C}^d : k^2 = k_1^2 + \dots + k_d^2 = E\}, \\ \Theta_E &= \{k \in \Sigma_E, l \in \Sigma_E : \operatorname{Im} k = \operatorname{Im} l\}. \end{aligned} \quad (4.9)$$

Under the assumptions of Theorem 2.1, we have that:

$$\mu(x, k) \rightarrow 1 \quad \text{as} \quad |\operatorname{Im} k| \rightarrow \infty \quad (4.10)$$

and, for any $\sigma > 1$,

$$|\mu(x, k)| + |\nabla \mu(x, k)| \leq \sigma \quad \text{for} \quad |\operatorname{Im} k| \geq r_1(N, D, E, m, \sigma), \quad (4.11)$$

where $x \in \mathbb{R}^d, k \in \Sigma_E$;

$$\hat{v}(p) = \lim_{\substack{(k, l) \in \Theta_E, k-l=p \\ |\operatorname{Im} k| = |\operatorname{Im} l| \rightarrow \infty}} h(k, l) \quad \text{for any } p \in \mathbb{R}^d, \quad (4.12)$$

$$\begin{aligned} |\hat{v}(p) - h(k, l)| &\leq \frac{c_1(D, E, m)N^2}{\rho} \quad \text{for } (k, l) \in \Theta_E, \quad p = k - l, \\ |\operatorname{Im} k| = |\operatorname{Im} l| &= \rho \geq r_2(N, D, E, m), \\ p^2 &\leq 4(E + \rho^2), \end{aligned} \quad (4.13)$$

where

$$\hat{v}(p) = (2\pi)^{-d} \int_{\mathbb{R}^d} e^{ipx} v(x) dx, \quad p \in \mathbb{R}^d. \quad (4.14)$$

Results of the type (4.10) go back to [2]. Results of the type (4.12), (4.13) (with less precise right-hand side in (4.13)) go back to [10]. In the present work estimate (4.11) is given according to [18], [20]. Estimate (4.13) follows, for example, from the estimate

$$\begin{aligned} \|\Lambda^{-s} g(k) \Lambda^{-s}\|_{\mathbb{L}^2(\mathbb{R}^d) \rightarrow \mathbb{L}^2(\mathbb{R}^d)} &= O(|k|^{-1}) \quad \text{as } |k| \rightarrow \infty, \\ k &\in \mathbb{C}^d \setminus \mathbb{R}^d, \quad |k| = (|\operatorname{Re} k|^2 + |\operatorname{Im} k|^2)^{1/2}, \end{aligned} \quad (4.15)$$

for $s > 1/2$, where $g(k)$ denotes the integral operator with the Schwartz kernel $g(x - y, k)$ and Λ denotes the multiplication operator by the function $(1 + |x|^2)^{1/2}$. Estimate (4.15) was formulated, first, in [13] for $d \geq 3$. Concerning proof of (4.15), see [29].

In addition, we have that:

$$h_2(k, l) - h_1(k, l) = (2\pi)^{-d} \int_{\mathbb{R}^d} \psi_1(x, -l)(v_2(x) - v_1(x))\psi_2(x, k)dx$$

$$\text{for } (k, l) \in \Theta_E, \quad |\text{Im } k| = |\text{Im } l| \neq 0,$$

$$\text{and } v_1, v_2 \text{ satisfying (4.5),}$$
(4.16)

$$h_2(k, l) - h_1(k, l) = (2\pi)^{-d} \int_{\partial D} [\psi_1(\cdot, -l)]_\alpha \left(\hat{M}_{\alpha, v_2} - \hat{M}_{\alpha, v_1} \right) [\psi_2(\cdot, k)]_\alpha dx$$

$$\text{for } (k, l) \in \Theta_E, \quad |\text{Im } k| = |\text{Im } l| \neq 0,$$

$$\text{and } v_1, v_2 \text{ satisfying (1.7), (4.6),}$$
(4.17)

where h_j, ψ_j denote h and ψ of (4.3) and (4.1) for $v = v_j$, and \hat{M}_{α, v_j} denotes the impedance boundary map of (1.4) for $v = v_j$, where $j = 1, 2$.

Formula (4.16) was given in [17]. Formula (4.17) follows from Theorem 3.1 and (4.16). Formula (4.17) for $\alpha = 0$ was given in [19].

5. Proof of Theorem 2.1

Let

$$\mathbb{L}_\mu^\infty(\mathbb{R}^d) = \{u \in \mathbb{L}^\infty(\mathbb{R}^d) : \|u\|_\mu < +\infty\},$$

$$\|u\|_\mu = \text{ess sup}_{p \in \mathbb{R}^d} (1 + |p|)^\mu |u(p)|, \quad \mu > 0.$$
(5.1)

Note that

$$w \in \mathbb{W}^{m,1}(\mathbb{R}^d) \implies \hat{w} \in \mathbb{L}_\mu^\infty(\mathbb{R}^d) \cap \mathcal{C}(\mathbb{R}^d),$$

$$\|\hat{w}\|_\mu \leq c_2(m, d) \|w\|_{m,1} \quad \text{for } \mu = m,$$
(5.2)

where $\mathbb{W}^{m,1}, \mathbb{L}_\mu^\infty$ are the spaces of (2.5), (5.1),

$$\hat{w}(p) = (2\pi)^{-d} \int_{\mathbb{R}^d} e^{ipx} w(x) dx, \quad p \in \mathbb{R}^d.$$
(5.3)

Using the inverse Fourier transform formula

$$w(x) = \int_{\mathbb{R}^d} e^{-ipx} \hat{w}(p) dp, \quad x \in \mathbb{R}^d,$$
(5.4)

we have that

$$\|v_1 - v_2\|_{\mathbb{L}^\infty(D)} \leq \sup_{x \in \bar{D}} \left| \int_{\mathbb{R}^d} e^{-ipx} (\hat{v}_2(p) - \hat{v}_1(p)) dp \right| \leq$$

$$\leq I_1(r) + I_2(r) \quad \text{for any } r > 0,$$
(5.5)

where

$$\begin{aligned} I_1(r) &= \int_{|p| \leq r} |\hat{v}_2(p) - \hat{v}_1(p)| dp, \\ I_2(r) &= \int_{|p| \geq r} |\hat{v}_2(p) - \hat{v}_1(p)| dp. \end{aligned} \quad (5.6)$$

Using (5.2), we obtain that

$$|\hat{v}_2(p) - \hat{v}_1(p)| \leq 2c_2(m, d)N(1 + |p|)^{-m}, \quad p \in \mathbb{R}^d. \quad (5.7)$$

Due to (4.13), we have that

$$\begin{aligned} |\hat{v}_2(p) - \hat{v}_1(p)| &\leq |h_2(k, l) - h_1(k, l)| + \frac{2c_1(D, E, m)N^2}{\rho}, \\ p &\in \mathbb{R}^d, \quad p = k - l, \quad (k, l) \in \Theta_E, \\ |\operatorname{Im} k| = |\operatorname{Im} l| &= \rho \geq r_2(N, D, E, m), \\ p^2 &\leq 4(E + \rho^2). \end{aligned} \quad (5.8)$$

Let

$$\begin{aligned} c_3 &= (2\pi)^{-d} \int_{\partial D} dx, \quad L = \max_{x \in \partial D} |x|, \\ \delta_\alpha &= \|\hat{M}_{\alpha, v_2}(E) - \hat{M}_{\alpha, v_1}(E)\|, \end{aligned} \quad (5.9)$$

where $\|\hat{M}_{\alpha, v_2}(E) - \hat{M}_{\alpha, v_1}(E)\|$ is defined according to (2.2).

Due to (4.16), (4.17), we have that

$$\begin{aligned} |h_2(k, l) - h_1(k, l)| &\leq c_3 \|\psi_1(\cdot, -l)\|_{\mathbb{L}^\infty(\partial D)} \delta_\alpha \|\psi_2(\cdot, k)\|_{\mathbb{L}^\infty(\partial D)}, \\ (k, l) &\in \Theta_E, \quad |\operatorname{Im} k| = |\operatorname{Im} l| \neq 0. \end{aligned} \quad (5.10)$$

Using (1.5), (4.11), we find that

$$\begin{aligned} \|\psi(\cdot, k)\|_{\mathbb{L}^\infty(\partial D)} &\leq c_4(E) \sigma \exp\left(|\operatorname{Im} k|(L + 1)\right), \\ k &\in \Sigma_E, \quad |\operatorname{Im} k| \geq r_1(N, D, E, m, \sigma). \end{aligned} \quad (5.11)$$

Here and bellow in this section the constant σ is the same that in (4.11).

Combining (5.10) and (5.11), we obtain that

$$\begin{aligned} |h_2(k, l) - h_1(k, l)| &\leq c_3 (c_4(E)\sigma)^2 \exp\left(2\rho(L + 1)\right) \delta_\alpha, \\ (k, l) &\in \Theta_E, \quad \rho = |\operatorname{Im} k| = |\operatorname{Im} l| \geq r_1(N, D, E, m, \sigma). \end{aligned} \quad (5.12)$$

Using (5.8), (5.12), we get that

$$\begin{aligned} |\hat{v}_2(p) - \hat{v}_1(p)| &\leq c_3 (c_4(E)\sigma)^2 \exp\left(2\rho(L + 1)\right) \delta_\alpha + \frac{2c_1(D, E, m)N^2}{\rho}, \\ p &\in \mathbb{R}^d, \quad p^2 \leq 4(E + \rho^2), \quad \rho \geq r_3(N, D, E, m, \sigma), \end{aligned} \quad (5.13)$$

where $r_3(N, D, E, m, \sigma)$ is such that

$$\rho \geq r_3(N, D, E, m, \sigma) \implies \begin{cases} \rho \geq r_1(N, D, E, m, \sigma), \\ \rho \geq r_2(N, D, E, m), \\ \rho^{2/m} \leq 4(E + \rho^2). \end{cases} \quad (5.14)$$

Let

$$c_5 = \int_{p \in \mathbb{R}^d, |p| \leq 1} dp, \quad c_6 = \int_{p \in \mathbb{R}^d, |p|=1} dp. \quad (5.15)$$

Using (5.6), (5.13), we get that

$$I_1(r) \leq c_5 r^d \left(c_3 (c_4(E)\sigma)^2 \exp\left(2\rho(L+1)\right) \delta_\alpha + \frac{2c_1(D, E, m)N^2}{\rho} \right), \quad (5.16)$$

$$r > 0, \quad r^2 \leq 4(\rho^2 + E), \quad \rho \geq r_3(N, D, E, m, \sigma).$$

Using (5.6), (5.7), we find that for any $r > 0$

$$I_2(r) \leq 2c_2(m, d)Nc_6 \int_r^{+\infty} \frac{dt}{t^{m-d+1}} \leq \frac{2c_2(m, D)Nc_6}{m-d} \frac{1}{r^{m-d}}. \quad (5.17)$$

Combining (5.5), (5.16), (5.17) for $r = \rho^{1/m}$ and (5.14), we get that

$$\|v_1 - v_2\|_{\mathbb{L}^\infty(D)} \leq c_7(D, \sigma) \rho^{d/m} e^{2\rho(L+1)} \delta_\alpha + c_8(N, D, E, m) \rho^{-\frac{m-d}{m}}, \quad (5.18)$$

$$\rho \geq r_3(N, D, E, m, \sigma).$$

We fix some $\tau \in (0, 1)$ and let

$$\beta = \frac{1-\tau}{2(L+1)}, \quad \rho = \beta \ln(3 + \delta_\alpha^{-1}), \quad (5.19)$$

where δ_α is so small that $\rho \geq r_3(N, D, E, m, \sigma)$. Then due to (5.18), we have that

$$\begin{aligned} \|v_1 - v_2\|_{\mathbb{L}^\infty(D)} &\leq c_7(D, \sigma) (\beta \ln(3 + \delta_\alpha^{-1}))^{d/m} (3 + \delta_\alpha^{-1})^{2\beta(L+1)} \delta_\alpha + \\ &\quad + c_8(N, D, E, m) (\beta \ln(3 + \delta_\alpha^{-1}))^{-\frac{m-d}{m}} = \\ &= c_7(D, \sigma) \beta^{d/m} (1 + 3\delta_\alpha)^{1-\tau} \delta_\alpha^\tau (\ln(3 + \delta_\alpha^{-1}))^{d/m} + \\ &\quad + c_8(N, D, E, m) \beta^{-\frac{m-d}{m}} (\ln(3 + \delta_\alpha^{-1}))^{-\frac{m-d}{m}}, \end{aligned} \quad (5.20)$$

where τ, β and δ_α are the same as in (5.19).

Using (5.20), we obtain that

$$\|v_1 - v_2\|_{\mathbb{L}^\infty(D)} \leq c_9(N, D, E, m, \sigma) (\ln(3 + \delta_\alpha^{-1}))^{-\frac{m-d}{m}} \quad (5.21)$$

for $\delta_\alpha = \|\hat{M}_{\alpha, v_2} - \hat{M}_{\alpha, v_1}\| \leq \delta^{(0)}(N, D, E, m, \sigma)$, where $\delta^{(0)}$ is a sufficiently small positive constant. Estimate (5.21) in the general case (with modified c_9) follows from (5.21) for $\delta_\alpha \leq \delta^{(0)}(N, D, E, m, \sigma)$ and the property that $\|v_j\|_{\mathbb{L}^\infty(D)} \leq c_{10}(D, m)N$.

Thus, Theorem 2.1 is proved for $s = \frac{m-d}{m}$ and, since $\ln(3 + \delta_\alpha^{-1}) > 1$, for any $0 < s \leq \frac{m-d}{m}$.

6. Buckhgeim-type functions

In dimension $d = 2$, we consider the functions $G_{z_0}, \psi_{z_0}, \tilde{\psi}_{z_0}, \delta h_{z_0}$ of [23] and going back to Buckhgeim's paper [4]:

$$\begin{aligned}\psi_{z_0}(z, \lambda) &= e^{\lambda(z-z_0)^2} + \int_D G_{z_0}(z, \zeta, \lambda) v(\zeta) \psi_{z_0}(\zeta, \lambda) d\operatorname{Re}\zeta d\operatorname{Im}\zeta, \\ \tilde{\psi}_{z_0}(z, \lambda) &= e^{\bar{\lambda}(\bar{z}-\bar{z}_0)^2} + \int_D \overline{G_{z_0}(z, \zeta, \lambda)} v(\zeta) \tilde{\psi}_{z_0}(\zeta, \lambda) d\operatorname{Re}\zeta d\operatorname{Im}\zeta,\end{aligned}\tag{6.1}$$

$$G_{z_0}(z, \zeta, \lambda) = \frac{1}{4\pi^2} \int_D \frac{e^{-\lambda(\eta-z_0)^2 + \bar{\lambda}(\bar{\eta}-\bar{z}_0)^2} d\operatorname{Re}\eta d\operatorname{Im}\eta}{(z-\eta)(\bar{\eta}-\bar{\zeta})} e^{\lambda(z-z_0)^2 - \bar{\lambda}(\bar{z}-\bar{z}_0)^2},\tag{6.2}$$

$z = x_1 + ix_2, z_0 \in D, \lambda \in \mathbb{C},$

where \mathbb{R}^2 is identified with \mathbb{C} and v, D satisfy (1.2), (1.3) for $d = 2$;

$$\delta h_{z_0}(\lambda) = \int_D \tilde{\psi}_{z_0,1}(z, -\lambda) (v_2(z) - v_1(z)) \psi_{z_0,2}(z, \lambda) d\operatorname{Re}z d\operatorname{Im}z, \quad \lambda \in \mathbb{C},\tag{6.3}$$

where v_1, v_2 satisfy (1.3) for $d = 2$ and $\tilde{\psi}_{z_0,1}, \psi_{z_0,2}$ denote $\tilde{\psi}_{z_0}, \psi_{z_0}$ of (6.1) for $v = v_1$ and $v = v_2$, respectively.

We recall that (see [23], [24]):

$$\begin{aligned}4 \frac{\partial^2}{\partial z \partial \bar{z}} G_{z_0}(z, \zeta, \lambda) &= \delta(z - \zeta), \\ 4 \frac{\partial^2}{\partial \zeta \partial \bar{\zeta}} G_{z_0}(z, \zeta, \lambda) &= \delta(z - \zeta),\end{aligned}\tag{6.4}$$

where $z, z_0, \zeta \in D, \lambda \in \mathbb{C}$ and δ is the Dirac delta function; formulas (6.1) at fixed z_0 and λ are considered as equations for $\psi_{z_0}, \tilde{\psi}_{z_0}$ in $L^\infty(D)$; as a corollary of (6.1), (6.2), (6.4), the functions $\psi_{z_0}, \tilde{\psi}_{z_0}$ satisfy (1.1) for $E = 0$ and $d = 2$; δh_{z_0} is similar to the right side of (4.16).

Let potentials $v, v_1, v_2 \in C^2(\bar{D})$ and

$$\begin{aligned}\|v\|_{C^2(\bar{D})} &\leq N, \quad \|v_j\|_{C^2(\bar{D})} \leq N, \quad j = 1, 2, \\ (v_1 - v_2)|_{\partial D} &= 0, \quad \frac{\partial}{\partial \nu}(v_1 - v_2)|_{\partial D} = 0,\end{aligned}\tag{6.5}$$

then we have that:

$$\psi_{z_0}(z, \lambda) = e^{\lambda(z-z_0)^2} \mu_{z_0}(z, \lambda), \quad \tilde{\psi}_{z_0}(z, \lambda) = e^{\bar{\lambda}(\bar{z}-\bar{z}_0)^2} \tilde{\mu}_{z_0}(z, \lambda),\tag{6.6}$$

$$\mu_{z_0}(z, \lambda) \rightarrow 1, \quad \tilde{\mu}_{z_0}(z, \lambda) \rightarrow 1 \quad \text{as } |\lambda| \rightarrow \infty\tag{6.7}$$

and, for any $\sigma > 1$,

$$|\mu_{z_0}(z, \lambda)| + |\nabla \mu_{z_0}(z, \lambda)| \leq \sigma,\tag{6.8a}$$

$$|\tilde{\mu}_{z_0}(z, \lambda)| + |\nabla \tilde{\mu}_{z_0}(z, \lambda)| \leq \sigma,\tag{6.8b}$$

where $\nabla = (\partial/\partial x_1, \partial/\partial x_2)$, $z = x_1 + ix_2$, $z_0 \in D$, $\lambda \in \mathbb{C}$, $|\lambda| \geq \rho_1(N, D, \sigma)$;

$$v_2(z_0) - v_1(z_0) = \lim_{\lambda \rightarrow \infty} \frac{2}{\pi} |\lambda| \delta h_{z_0}(\lambda) \quad (6.9)$$

for any $z_0 \in D$,

$$\left| v_2(z_0) - v_1(z_0) - \frac{2}{\pi} |\lambda| \delta h_{z_0}(\lambda) \right| \leq \frac{c_{11}(N, D) (\ln(3|\lambda|))^2}{|\lambda|^{3/4}} \quad (6.10)$$

for $z_0 \in D$, $|\lambda| \geq \rho_2(N, D)$.

Formulas (6.6) can be considered as definitions of μ_{z_0} , $\tilde{\mu}_{z_0}$. Formulas (6.7), (6.9) were given in [23], [24] and go back to [4]. Estimate (6.10) was obtained in [23], [25]. Estimates (6.8) are proved in Section 8.

7. Proof of Theorem 2.2

We suppose that $\tilde{\psi}_{z_0,1}(\cdot, -\lambda)$, $\psi_{z_0,2}(\cdot, \lambda)$, $\delta h_{z_0}(\lambda)$ are defined as in Section 6 but with $v_j - E$ in place of v_j , $j = 1, 2$. We use the identity

$$\hat{M}_{\alpha,v}(E) = \hat{M}_{\alpha,v-E}(0). \quad (7.1)$$

We also use the notation $N_E = N + E$. Then, using (6.10), we have that

$$\left| v_2(z_0) - v_1(z_0) - \frac{2}{\pi} |\lambda| \delta h_{z_0}(\lambda) \right| \leq \frac{c_{11}(N_E, D) (\ln(3|\lambda|))^2}{|\lambda|^{3/4}} \quad (7.2)$$

for $z_0 \in D$, $|\lambda| \geq \rho_2(N_E, D)$.

According to Theorem 3.1 and (6.3), we get that

$$\delta h_{z_0}(\lambda) = \frac{1}{4\pi^2} \int_{\partial D} [\tilde{\psi}_{z_0,1}(\cdot, -\lambda)]_{\alpha} \left(\hat{M}_{\alpha,v_2}(E) - \hat{M}_{\alpha,v_1}(E) \right) [\psi_{z_0,2}(\cdot, \lambda)]_{\alpha} |dz|, \quad (7.3)$$

$\lambda \in \mathbb{C}$.

Let

$$c_{12} = \frac{1}{4\pi^2} \int_{\partial D} |dz|, \quad L = \max_{z \in \partial D} |z|, \quad (7.4)$$

$$\delta_{\alpha} = \|\hat{M}_{\alpha,v_2}(E) - \hat{M}_{\alpha,v_1}(E)\|,$$

where $\|\hat{M}_{\alpha,v_2}(E) - \hat{M}_{\alpha,v_1}(E)\|$ is defined according to (2.2).

Using (7.3), we get that

$$|\delta h_{z_0}(\lambda)| \leq c_{12} \|\tilde{\psi}_{z_0,1}(\cdot, -\lambda)\|_{\alpha} \| \delta_{\alpha} \| \|\psi_{z_0,2}(\cdot, \lambda)\|_{\alpha} \quad (7.5)$$

Using (1.5), (6.8), we find that:

$$\|\tilde{\psi}_{z_0,1}(\cdot, -\lambda)\|_{\alpha} \leq \sigma \exp \left(|\lambda| (4L^2 + 4L) \right),$$

$$\|\psi_{z_0,2}(\cdot, \lambda)\|_{\alpha} \leq \sigma \exp \left(|\lambda| (4L^2 + 4L) \right), \quad (7.6)$$

$\lambda \in \mathbb{C}$, $|\lambda| \geq \rho_1(N_E, D, \sigma)$.

Here and bellow in this section the constant σ is the same that in (6.8).

Combining (7.5), (7.6), we obtain that

$$\begin{aligned} |\delta h_{z_0}(\lambda)| &\leq c_{12}\sigma^2 \exp\left(|\lambda|(8L^2 + 8L)\right)\delta_\alpha, \\ \lambda &\in \mathbb{C}, \quad |\lambda| \geq \rho_1(N_E, D, \sigma). \end{aligned} \quad (7.7)$$

Using (7.2) and (7.7), we get that

$$\begin{aligned} |v_2(z_0) - v_1(z_0)| &\leq c_{12}\sigma^2 \exp\left(|\lambda|(8L^2 + 8L)\right)\delta_\alpha + \frac{c_{11}(N_E, D) (\ln(3|\lambda|))^2}{|\lambda|^{3/4}}, \\ z_0 &\in D, \quad \lambda \in \mathbb{C}, \quad |\lambda| \geq \rho_3(N_E, D, \sigma) = \max\{\rho_1, \rho_2\}. \end{aligned} \quad (7.8)$$

We fix some $\tau \in (0, 1)$ and let

$$\beta = \frac{1 - \tau}{8L^2 + 8L}, \quad \lambda = \beta \ln(3 + \delta_\alpha^{-1}), \quad (7.9)$$

where δ_α is so small that $|\lambda| \geq \rho_3(N_E, D, \sigma)$. Then due to (7.8), we have that

$$\begin{aligned} \|v_1 - v_2\|_{\mathbb{L}^\infty(D)} &\leq c_{12}\sigma^2 (3 + \delta_\alpha^{-1})^{\beta(8L^2+8L)} \delta_\alpha + \\ &+ c_{11}(N_E, D) \frac{(\ln(3\beta \ln(3 + \delta_\alpha^{-1})))^2}{(\beta \ln(3 + \delta_\alpha^{-1}))^{3/4}} = \\ &= c_{12}\sigma^2 (1 + 3\delta_\alpha)^{1-\tau} \delta_\alpha^\tau + \\ &+ c_{11}(N_E, D) \beta^{-3/4} \frac{(\ln(3\beta \ln(3 + \delta_\alpha^{-1})))^2}{(\ln(3 + \delta_\alpha^{-1}))^{3/4}}, \end{aligned} \quad (7.10)$$

where τ, β and δ_α are the same as in (7.9).

Using (7.10), we obtain that

$$\|v_1 - v_2\|_{\mathbb{L}^\infty(D)} \leq c_{13}(N_E, D, \sigma) (\ln(3 + \delta_\alpha^{-1}))^{-3/4} (\ln(3 \ln(3 + \delta_\alpha^{-1})))^2 \quad (7.11)$$

for $\delta_\alpha = \|\hat{M}_{\alpha, v_2}(E) - \hat{M}_{\alpha, v_1}(E)\| \leq \delta^{(0)}(N_E, D, \sigma)$, where $\delta^{(0)}$ is a sufficiently small positive constant. Estimate (5.21) in the general case (with modified c_{13}) follows from (7.11) for $\delta_\alpha \leq \delta^{(0)}(N_E, D, \sigma)$ and the property that $\|v_j\|_{\mathbb{L}^\infty(D)} \leq c_{14}(D)N$.

Thus, Theorem 2.2 is proved for $s = \frac{3}{4}$ and, since $\ln(3 + \delta_\alpha^{-1}) > 1$, for any $0 < s \leq \frac{3}{4}$.

8. Proof of estimates (6.8)

In this section we prove estimate (6.8a). Estimate (6.8b) can be proved a completely similar way. Let

$$\begin{aligned} C_{\bar{z}}^1(\bar{D}) &= \left\{ u : u, \frac{\partial u}{\partial \bar{z}} \in C(\bar{D}) \right\}, \\ \|u\|_{C_{\bar{z}}^1(\bar{D})} &= \max \left(\|u\|_{C(\bar{D})}, \left\| \frac{\partial u}{\partial \bar{z}} \right\|_{C(\bar{D})} \right). \end{aligned} \quad (8.1)$$

Due to estimates of Section 3 of [23], we have that, for any $\varepsilon_1 > 0$,

$$\mu_{z_0}(\cdot, \lambda) \in C_{\bar{z}}^1(\bar{D}), \quad \|\mu_{z_0}(\cdot, \lambda)\|_{C_{\bar{z}}^1(\bar{D})} \leq 1 + \varepsilon_1 \quad \text{for } |\lambda| \geq \rho_4(N, D, \varepsilon_1). \quad (8.2)$$

In view of (8.2), to prove (6.8a) it remains to prove that, for any $\varepsilon_2 > 0$,

$$\partial_z \mu_{z_0}(\cdot, \lambda) \in C(\bar{D}), \quad \|\partial_z \mu_{z_0}(\cdot, \lambda)\|_{C(\bar{D})} \leq \varepsilon_2 \quad \text{for } |\lambda| \geq \rho_5(N, D, \varepsilon_2), \quad (8.3)$$

where $\partial_z \mu_{z_0}(\cdot, \lambda)$ is considered as a function of $z \in \bar{D}$ and $\partial_z = \partial/\partial z$.

We have that (see Sections 2 and 5 of [23]):

$$\partial_z \mu_{z_0} = \frac{1}{4} \Pi \bar{\mathbb{T}}_{z_0, \lambda} v \mu_{z_0}, \quad (8.4)$$

$$\Pi u(z) = -\frac{1}{\pi} \int_D \frac{u(\zeta)}{(\zeta - z)^2} d\text{Re}\zeta d\text{Im}\zeta, \quad (8.5)$$

$$\bar{\mathbb{T}}_{z_0, \lambda} u(z) = -\frac{e^{-\lambda(z-z_0)^2 + \bar{\lambda}(\bar{z}-\bar{z}_0)^2}}{\pi} \int_D \frac{e^{\lambda(\zeta-z_0)^2 - \bar{\lambda}(\bar{\zeta}-\bar{z}_0)^2}}{\bar{\zeta} - \bar{z}} u(\zeta) d\text{Re}\zeta d\text{Im}\zeta, \quad (8.6)$$

where u is a test function, $z \in \bar{D}$.

In view of (8.2), (8.4) and Theorem 1.33 of [28], to prove (8.3) it is sufficient to show that

$$\|\bar{\mathbb{T}}_{z_0, \lambda} u\|_{C_s(\bar{D})} \leq \frac{A(D, s)}{|\lambda|^{\delta(s)}} \|u\|_{C_{\bar{z}}^1(\bar{D})}, \quad |\lambda| \geq 1, \quad z_0 \in \bar{D}, \quad (8.7)$$

for some fixed $s \in (0, \frac{1}{2})$ and $\delta(s) > 0$, where $C_s(\bar{D})$ is the Hölder space,

$$\begin{aligned} C_s(\bar{D}) &= \{u \in C(\bar{D}) : \|u\|_{C_s(\bar{D})} < +\infty\}, \\ \|u\|_{C_s(\bar{D})} &= \max \left\{ \|u\|_{C(\bar{D})}, \|u\|'_{C_s(\bar{D})} \right\}, \\ \|u\|'_{C_s(\bar{D})} &= \sup_{z_1, z_2 \in \bar{D}, 0 < |z_1 - z_2| < 1} \frac{|u(z_1) - u(z_2)|}{|z_1 - z_2|^s}. \end{aligned} \quad (8.8)$$

Due to estimate (5.6) of [23], we have that

$$\|\bar{\mathbb{T}}_{z_0, \lambda} u\|_{C(\bar{D})} \leq \frac{A_0(D)}{|\lambda|^{1/2}} \|u\|_{C_{\bar{z}}^1(\bar{D})}, \quad |\lambda| \geq 1, \quad z_0 \in \bar{D}. \quad (8.9)$$

Therefore, to prove (8.7) it remains to prove that

$$\|\bar{\mathbb{T}}_{z_0, \lambda} u\|'_{C_s(\bar{D})} \leq \frac{A_1(D, s)}{|\lambda|^{\delta(s)}} \|u\|_{C_{\bar{z}}^1(\bar{D})}, \quad |\lambda| \geq 1, \quad z_0 \in \bar{D}, \quad (8.10)$$

for some fixed $s \in (0, \frac{1}{2})$ and $\delta(s) > 0$.

We will use that

$$\|u_1 u_2\|'_{C_s(\bar{D})} \leq \|u_1\|'_{C_s(\bar{D})} \|u_2\|_{C(\bar{D})} + \|u_1\|_{C(\bar{D})} \|u_2\|'_{C_s(\bar{D})}, \quad 0 < s < 1. \quad (8.11)$$

One can see that

$$\bar{\mathbb{T}}_{z_0, \lambda} = F_{z_0, -\lambda} \bar{\mathbb{T}} F_{z_0, \lambda}, \quad (8.12)$$

where $\bar{T} = \bar{T}_{z_0,0}$ and $F_{z_0,\lambda}$ is the multiplication operator by the function

$$F(z, z_0, \lambda) = e^{\lambda(z-z_0)^2 - \bar{\lambda}(\bar{z}-\bar{z}_0)^2}. \quad (8.13)$$

One can see also that

$$\begin{aligned} \|F(\cdot, z_0, -\lambda)\|_{C(\bar{D})} &= 1, \\ \|F(\cdot, z_0, -\lambda)\|'_{C_s(\bar{D})} &\leq A_2(D, s)|\lambda|^s, \quad |\lambda| \geq 1, \quad z_0 \in \bar{D}. \end{aligned} \quad (8.14)$$

In view of (8.9), (8.11) - (8.14), to prove (8.10) it remains to prove that

$$\|\bar{T}F_{z_0,\lambda}u\|'_{C_s(\bar{D})} \leq \frac{A_3(D, s)}{|\lambda|^{\delta_1(s)}} \|u\|_{C_{\bar{z}}^1(\bar{D})}, \quad |\lambda| \geq 1, \quad z_0 \in \bar{D}, \quad (8.15)$$

for some fixed $s \in (0, \frac{1}{2})$ and $\delta_1(s) > 0$.

We have that

$$\begin{aligned} \pi \bar{T}F_{z_0,\lambda}u(z_1) - \pi \bar{T}F_{z_0,\lambda}u(z_2) &= \int_D \frac{F(\zeta, z_0, \lambda)u(\zeta)(\bar{z}_2 - \bar{z}_1)}{(\bar{\zeta} - \bar{z}_1)(\bar{\zeta} - \bar{z}_2)} d\text{Re}\zeta d\text{Im}\zeta = \\ &= I_{z_0,\lambda,\varepsilon}(z_1, z_2) + J_{z_0,\lambda,\varepsilon}(z_1, z_2), \end{aligned} \quad (8.16)$$

where

$$I_{z_0,\lambda,\varepsilon}(z_1, z_2) = \int_{D \setminus D_{z_0,z_1,z_2,\varepsilon}} \frac{F(\zeta, z_0, \lambda)u(\zeta)(\bar{z}_2 - \bar{z}_1)}{(\bar{\zeta} - \bar{z}_1)(\bar{\zeta} - \bar{z}_2)} d\text{Re}\zeta d\text{Im}\zeta, \quad (8.17)$$

$$J_{z_0,\lambda,\varepsilon}(z_1, z_2) = \int_{D_{z_0,z_1,z_2,\varepsilon}} \frac{F(\zeta, z_0, \lambda)u(\zeta)(\bar{z}_2 - \bar{z}_1)}{(\bar{\zeta} - \bar{z}_1)(\bar{\zeta} - \bar{z}_2)} d\text{Re}\zeta d\text{Im}\zeta, \quad (8.18)$$

where $B_{z,\varepsilon} = \{\zeta \in \mathbb{C} : |\zeta - z| < \varepsilon\}$, $D_{z_0,z_1,z_2,\varepsilon} = D \setminus \left(\bigcup_{j=0}^2 B_{z_j,\varepsilon} \right)$.

We will use the following inequalities:

$$\left| \frac{z_2 - z_1}{(\zeta - z_1)(\zeta - z_2)} \right| \leq n_1 |z_2 - z_1|^s \sum_{j=1}^2 \frac{1}{|\zeta - z_j|^{1+s}}, \quad (8.19)$$

$$\left| \frac{z_2 - z_1}{(\zeta - z_1)(\zeta - z_2)(\zeta - z_0)} \right| \leq n_2 |z_2 - z_1|^s \sum_{j=0}^2 \frac{1}{|\zeta - z_j|^{2+s}}, \quad (8.20)$$

$$\left| \frac{\partial}{\partial \zeta} \left(\frac{z_2 - z_1}{(\zeta - z_1)(\zeta - z_2)(\zeta - z_0)} \right) \right| \leq n_3 |z_2 - z_1|^s \sum_{j=0}^2 \frac{1}{|\zeta - z_j|^{3+s}}, \quad (8.21)$$

where $s \in (0, 1)$, $n_1, n_2, n_3 > 0$, $z_0, z_1, z_2, \zeta \in \mathbb{C}$ and $\zeta \neq z_i$ for $j = 0, 1, 2$.

Using (8.17), (8.19), we obtain that

$$I_{z_0,\lambda,\varepsilon}(z_1, z_2) \leq n_4(s) \varepsilon^{1-s} |z_2 - z_1|^s, \quad (8.22)$$

where $n_4(s) > 0$, $z_0, z_1, z_2, \zeta \in \mathbb{C}$ and $\varepsilon \in (0, 1)$. Further, we have that

$$\begin{aligned} J_{z_0,\lambda,\varepsilon}(z_1, z_2) &= -\frac{1}{2\bar{\lambda}} \int_{D_{z_0,z_1,z_2,\varepsilon}} \frac{\partial F(\zeta, z_0, \lambda)}{\partial \bar{\zeta}} \frac{u(\zeta)(\bar{z}_2 - \bar{z}_1)}{(\bar{\zeta} - \bar{z}_1)(\bar{\zeta} - \bar{z}_2)(\bar{\zeta} - \bar{z}_0)} d\text{Re}\zeta d\text{Im}\zeta = \\ &= J_{z_0,\lambda,\varepsilon}^1(z_1, z_2) + J_{z_0,\lambda,\varepsilon}^2(z_1, z_2), \end{aligned} \quad (8.23)$$

where

$$\begin{aligned}
J_{z_0, \lambda, \varepsilon}^1(z_1, z_2) &= -\frac{1}{4i\lambda} \int_{\partial D_{z_0, z_1, z_2, \varepsilon}} \frac{F(\zeta, z_0, \lambda) u(\zeta) (\bar{z}_2 - \bar{z}_1)}{(\bar{\zeta} - \bar{z}_1)(\bar{\zeta} - \bar{z}_2)(\bar{\zeta} - \bar{z}_0)} d\zeta, \\
J_{z_0, \lambda, \varepsilon}^2(z_1, z_2) &= \frac{1}{2\lambda} \int_{D_{z_0, z_1, z_2, \varepsilon}} F(\zeta, z_0, \lambda) \frac{\partial}{\partial \bar{\zeta}} \left(\frac{u(\zeta) (\bar{z}_2 - \bar{z}_1)}{(\bar{\zeta} - \bar{z}_1)(\bar{\zeta} - \bar{z}_2)(\bar{\zeta} - \bar{z}_0)} \right) d\operatorname{Re}\zeta d\operatorname{Im}\zeta,
\end{aligned} \tag{8.24}$$

Using (8.20), (8.21), (8.24), we obtain that

$$\begin{aligned}
J_{z_0, \lambda, \varepsilon}^1(z_1, z_2) &\leq |\lambda|^{-1} n_5(D, s) \varepsilon^{-1-s} |z_2 - z_1|^s \|u\|_{C(\bar{D})}, \\
J_{z_0, \lambda, \varepsilon}^2(z_1, z_2) &\leq |\lambda|^{-1} n_6(D, s) \varepsilon^{-1-s} |z_2 - z_1|^s \|u\|_{C(\bar{D})} + \\
&\quad + |\lambda|^{-1} n_7(D, s) \varepsilon^{-s} |z_2 - z_1|^s \left\| \frac{\partial u}{\partial \bar{z}} \right\|_{C(\bar{D})},
\end{aligned} \tag{8.25}$$

where $z_0, z_1, z_2, \lambda \in \mathbb{C}$, $|\lambda| \geq 1$, $\varepsilon \in (0, 1)$.

Using (8.16), (8.22), (8.23), (8.25) and putting $\varepsilon = |\lambda|^{-1/2}$ into (8.22), (8.25) we obtain (8.15) with $\delta_1(s) = (1-s)/2$.

9. Proof of symmetry (3.11)

Let D' be an open bounded domain in \mathbb{R}^d such that

- $D \subset D'$,
- D' satisfies (1.2),
- E is not a Dirichlet eigenvalue for the operator $-\Delta + v$ in D' .

Here and bellow in this section we assume that $v \equiv 0$ on $D' \setminus D$. Let $R(x, y, E)$ denote the Green function for the operator $-\Delta + v - E$ in D' with the Dirichlet boundary condition. We recall that

$$R(x, y, E) = R(y, x, E), \quad x, y \in D'. \tag{9.1}$$

Using (3.4), (9.1), we find that for $x, y \in D$

$$\begin{aligned}
&\int_{\partial D} \left(R(x, \xi, E) \frac{\partial R}{\partial \nu_\xi}(y, \xi, E) - R(y, \xi, E) \frac{\partial R}{\partial \nu_\xi}(x, \xi, E) \right) d\xi = \\
&= \int_D \left(R(x, \xi, E) (\Delta_\xi - v + E) R(y, \xi, E) - R(y, \xi, E) (\Delta_\xi - v + E) R(x, \xi, E) \right) d\xi = \\
&= -R(x, y, E) + R(y, x, E) = 0.
\end{aligned} \tag{9.2}$$

Note that $W = G_\alpha + R(E)$ is the solution of the equation

$$(-\Delta_x + v - E)W(x, y) = 0, \quad x, y \in D \tag{9.3}$$

with the boundary condition

$$\begin{aligned} & \left(\cos \alpha W(x, y) - \sin \alpha \frac{\partial W}{\partial \nu_x}(x, y) \right) \Big|_{x \in \partial D} = \\ & = \left(\cos \alpha R(x, y, E) - \sin \alpha \frac{\partial R}{\partial \nu_x}(x, y, E) \right) \Big|_{x \in \partial D}, \quad y \in D. \end{aligned} \quad (9.4)$$

Using (3.4) and (9.3), we find that for $x, y \in D$

$$\begin{aligned} & \int_{\partial D} \left(W(\xi, x) \frac{\partial W}{\partial \nu_\xi}(\xi, y) - W(\xi, y) \frac{\partial W}{\partial \nu_\xi}(\xi, x) \right) d\xi = \\ & = \int_D \left(W(\xi, x) (\Delta_\xi - v + E) W(\xi, y) - W(\xi, y) (\Delta_\xi - v + E) W(\xi, x) \right) d\xi = 0 \end{aligned} \quad (9.5)$$

Note that

$$W(x, y) = - \int_D W(\xi, y) (\Delta_\xi - v + E) R(\xi, x, E) d\xi, \quad x, y \in D. \quad (9.6)$$

Combining (3.4), (9.3) and (9.6), we obtain that

$$\begin{aligned} W(x, y) = - \int_{\partial D} \left(W(\xi, y) \frac{\partial R}{\partial \nu_\xi}(\xi, x, E) - R(\xi, x, E) \frac{\partial W}{\partial \nu_\xi}(\xi, y) \right) d\xi, \\ x, y \in D. \end{aligned} \quad (9.7)$$

Using (9.4) and (9.7), we get that

$$\begin{aligned} & \sin \alpha W(x, y) = \\ & = \int_{\partial D} W(\xi, y) \left(\cos \alpha W(\xi, x) - \sin \alpha \frac{\partial W}{\partial \nu_\xi}(\xi, x) - \cos \alpha R(\xi, x, E) \right) d\xi - \\ & \quad - \int_{\partial D} R(\xi, x, E) \left(\cos \alpha R(\xi, y, E) - \sin \alpha \frac{\partial R}{\partial \nu_\xi}(\xi, x, E) - \cos \alpha W(\xi, y) \right) d\xi, \\ & \quad x, y \in D. \end{aligned} \quad (9.8)$$

Combining similar to (9.8) formula for $\sin \alpha W(y, x)$, (9.2) and (9.5), we obtain that

$$\sin \alpha W(x, y) - \sin \alpha W(y, x) = 0, \quad x, y \in D. \quad (9.9)$$

In the case of $\sin \alpha = 0$, combining (9.4) and (9.7), we get that

$$\begin{aligned} W(x, y) = \int_{\partial D} \left(-R(\xi, y, E) \frac{\partial R}{\partial \nu_\xi}(\xi, x, E) + W(\xi, x) \frac{\partial W}{\partial \nu_\xi}(\xi, y) \right) d\xi, \\ x, y \in D. \end{aligned} \quad (9.10)$$

Hence, one can get that for any α

$$W(x, y) = W(y, x), \quad x, y \in D. \quad (9.11)$$

Combining (9.1) and (9.11), we obtain (3.11).

Acknowledgements

The second author was partially supported by the Russian Federation Government grant No. 2010-220-01-077.

References

- [1] G.Alessandrini, *Stable determination of conductivity by boundary measurements*, Appl.Anal. 27, 1988, 153-172.
- [2] R. Beals and R. Coifman, *Multidimensional inverse scattering and nonlinear partial differential equations*, Proc. Symp. Pure Math., 43, 1985, 45-70.
- [3] H. Begehr and T. Vaitekhovich, *Some harmonic Robin functions in the complex plane*, Adv. Pure Appl. Math. 1, 2010, 19–34.
- [4] A. L. Bukhgeim, *Recovering a potential from Cauchy data in the two-dimensional case*, J. Inverse Ill-Posed Probl. 16, 2008, no. 1, 19–33.
- [5] Calderón, A.P., *On an inverse boundary problem*, Seminar on Numerical Analysis and its Applications to Continuum Physics, Soc. Brasileira de Matematica, Rio de Janeiro, 1980, 61–73.
- [6] L.D. Faddeev, *Growing solutions of the Schrödinger equation*, Dokl. Akad. Nauk SSSR, 165, N.3, 1965, 514-517 (in Russian); English Transl.: Sov. Phys. Dokl. 10, 1966, 1033-1035.
- [7] L.D. Faddeev, *The inverse problem in the quantum theory of scattering. II*, Current problems in mathematics, Vol. 3, 1974, pp. 93–180, 259. Akad. Nauk SSSR Vsesojuz. Inst. Naucn. i Tehn. Informacii, Moscow(in Russian); English Transl.: J.Sov. Math. 5, 1976, 334-396.
- [8] I.M. Gelfand, *Some problems of functional analysis and algebra*, Proceedings of the International Congress of Mathematicians, Amsterdam, 1954, pp.253-276.
- [9] F. Gesztesy and M. Mitrea, *Robin-to-Robin Maps and Krein-Type Resolvent Formulas for Schrodinger Operators on Bounded Lipschitz Domains*, Modern Analysis and Applications Operator Theory: Advances and Applications, Volume 191, 2009, Part 1, 81-113.
- [10] G.M. Henkin and R.G. Novikov, *The $\bar{\partial}$ -equation in the multidimensional inverse scattering problem*, Uspekhi Mat. Nauk 42(3), 1987, 93-152 (in Russian); English Transl.: Russ. Math. Surv. 42(3), 1987, 109-180.
- [11] M.I. Isaev, *Exponential instability in the Gel'fand inverse problem on the energy intervals*, J. Inverse Ill-Posed Probl., Vol. 19(3), 2011, 453-473,; e-print arXiv: 1012.2193.
- [12] L. Lanzani and Z. Shen, *On the Robin boundary condition for Laplace's equation in Lipschitz domains*, Comm. Partial Differential Equations, 29, 2004, 91–109.
- [13] R.B. Lavine and A.I. Nachman, *On the inverse scattering transform of the n -dimensional Schrödinger operator* Topics in Soliton Theory and Exactly Solvable Nonlinear Equations ed M Ablowitz, B Fuchssteiner and M Kruskal (Singapore: World Scientific), 1987, pp 33–44
- [14] N. Mandache, *Exponential instability in an inverse problem for the Schrödinger equation* Inverse Problems. 17, 2001, 1435–1444.
- [15] A. Nachman, *Global uniqueness for a two-dimensional inverse boundary value problem*, Ann. Math. 143, 1996, 71–96.
- [16] R.G. Novikov, *Multidimensional inverse spectral problem for the equation $-\Delta\psi + (v(x) - Eu(x))\psi = 0$* Funkt. Anal. Prilozhen. 22(4), 1988, 11–22 (in Russian); Engl. Transl. Funct. Anal. Appl. 22, 1988, 263–72.

- [17] R.G. Novikov, *$\bar{\partial}$ -method with nonzero background potential. Application to inverse scattering for the two-dimensional acoustic equation*, Comm. Partial Differential Equations 21, 1996, no. 3-4, 597-618.
- [18] R.G. Novikov, *Approximate solution of the inverse problem of quantum scattering theory with fixed energy in dimension 2*, Proceedings of the Steklov Mathematical Institute 225, 1999, Solitony Geom. Topol. na Perekrest., 301-318 (in Russian); Engl. Transl. in Proc. Steklov Inst. Math. 225, 1999, no. 2, 285-302.
- [19] R.G. Novikov, *Formulae and equations for finding scattering data from the Dirichlet-to-Neumann map with nonzero background potential*, Inverse Problems 21, 2005, 257-270.
- [20] R.G. Novikov, *On non-overdetermined inverse scattering at zero energy in three dimensions*, Ann. Scuola Norm. Sup. Pisa Cl. Sci. 5, 2006, 279-328
- [21] R.G. Novikov, *New global stability estimates for the Gel'fand-Calderon inverse problem*, Inverse Problems 27, 2011, 015001(21pp); e-print arXiv:1002.0153.
- [22] R.G. Novikov and N.N. Novikova, *On stable determination of potential by boundary measurements*, ESAIM: Proceedings 26, 2009, 94-99.
- [23] R. Novikov and M. Santacesaria, *A global stability estimate for the Gel'fand-Calderon inverse problem in two dimensions*, J. Inverse Ill-Posed Probl., Volume 18, Issue 7, 2010, Pages 765-785; e-print arXiv: 1008.4888.
- [24] R. Novikov and M. Santacesaria, *Global uniqueness and reconstruction for the multi-channel Gel'fand-Calderon inverse problem in two dimensions*, Bulletin des Sciences Mathematiques 135, 5, 2011, 421-434.
- [25] M. Santacesaria, *Global stability for the multi-channel Gel'fand-Calderon inverse problem in two dimensions*, Bulletin des Sciences Mathematiques (to appear), e-print: hal-00569366.
- [26] M. Santacesaria, *New global stability estimates for the Calderon inverse problem in two dimensions*, e-print: hal-00628403.
- [27] J. Sylvester and G. Uhlmann, *A global uniqueness theorem for an inverse boundary value problem*, Ann. of Math. 125, 1987, 153-169
- [28] I.N. Vekua, *Generalized Analytic Functions*, Pergamon Press Ltd., 1962.
- [29] R. Weder, *Generalized limiting absorption method and multidimensional inverse scattering theory*, Mathematical Methods in the Applied Sciences, 14, 1991, 509-524.

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