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Reality-Based Real-Time Cell Indentation Simulator

Hamid Ladjal, Jean-Luc Hanus, Anand Pillarisetti, Carol Keefer, Antoine Ferreira, *Member, IEEE*,
and Jaydev P. Desai, *Senior Member, IEEE*

Abstract—Training simulators that provide realistic visual and haptic feedback during cell indentation tasks are currently investigated. Complex cell geometry inherent to biological cells and intricate mechanical properties drive the need for precise mechanical and numerical modeling to assure accurate cell deformation and force calculations. Advances in alternative finite-element formulation, such as the mass–tensor approach, have reached a state, where they are applicable to model soft-cell deformation in real time. The geometrical characteristics and the mechanical properties of different cells are determined with atomic force microscopy (AFM) indentation. A real-time, haptics-enabled simulator for cell centered indentation has been developed, which utilizes the AFM data (mechanical and geometrical properties of embryonic stem cells) to accurately replicate the indentation task and predict the cell deformation during indentation in real time. This tool can be used as a mechanical marker to characterize the biological state of the cell. The operator is able to feel the change in the stiffness during cell deformation between fixed and live cells in real time. A comparative study with finite-element simulations using a commercial software and the experimental data demonstrate the effectiveness of the proposed physically based model.

Index Terms—Atomic force microscope (AFM), finite-element modeling (FEM), haptics, modeling, real-time interaction, stem cell.

I. INTRODUCTION

MECHANICAL manipulation and characterization of biological cells is currently one of the most promising research areas in the field of biorobotics applied to cellular level interactions. The mechanical properties, such as elasticity, membrane tension, cell shape, and adhesion strength, may play an

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The Promise of Stem Cell Research

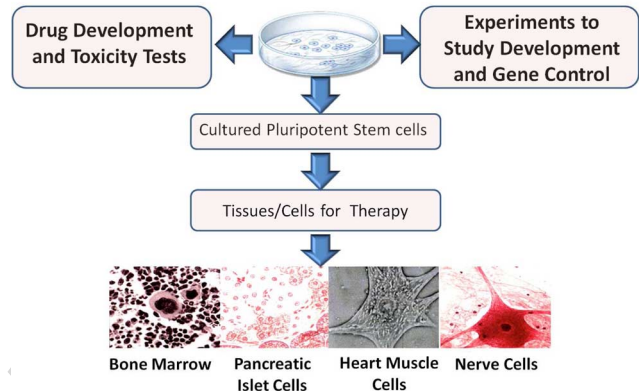


Fig. 1. Stem cells are valuable resources for toxicology, developmental biology, and regenerative medicine research.

important role in cell development and differentiation [1]–[3]. In contrast to adult stem cells, embryonic stem cells differentiate into every cell type in the body, and thus, have a huge therapeutic potential. Stem-cell-based systems offer a very promising and innovative solution for obtaining large number of cells for early efficacy and toxicity screening. Stem cell technology provides a new tool for drug development and a better insight to understand mammalian gene function (see Fig. 1). The mechanical properties of biological cells have been studied with different techniques [4], the most popular being optical tweezers [5], magnetic beads [6], and micropipette aspiration [7]. Different microrobotic systems with various end effectors have been proposed [33], [34]. However, these methods cannot compete with the precision that can be obtained with the atomic force microscopy (AFM) method [8], [32]. AFMs have been widely used in the study of micro- and nanostructures including living cells. Modern AFM techniques allow solving several problems of cell biomechanics by simultaneous evaluation of the local mechanical properties and the topography of living cells, at a high spatial resolution and force sensitivity [9]. In these experiments, an AFM cantilever serves as a microindenter to probe the cell, and further analysis of force-indentation data yields the local Young's modulus. In addition, the AFM indentation technique can be used to characterize the viscoelastic behavior of the cell cytoskeleton [10], including viscosity [11], loss and storage moduli [12], and stress relaxation times. Mechanical phenotyping prove actually to be a valuable tool in the development of improved methods of targeted cellular differentiation of embryonic, induced pluripotent (iPS), and adult stem cells for therapeutic and diagnostic purposes [30], [31]. In this study, we propose to characterize the mechanical behavior of an individual mouse embryonic stem cells (mESC) in undifferentiated state

66 using AFM indentation [13]. As we hypothesize that the cell
 67 stiffness of mESC may reflect the nuclear physical changes of
 68 the cell, it should be possible to detect modifications in stiffness
 69 by probing the cell directly. However, different limitations are
 70 inherent in AFM-based indentation techniques. First, deforma-
 71 tions induced by the AFM probe and cell positions cannot be
 72 seen by the user. Second, the force sensing of the AFM can-
 73 not resolve the force interaction between the AFM tip and the
 74 cell's surface independently in all three dimensions to avoid
 75 cytoskeleton-disrupted experiments. Therefore, the user per-
 76 forms extremely challenging manipulations requiring advanced
 77 manual skills. Modeling the physics of the interaction between
 78 tip and cell is necessary to realize complicated indentation tasks
 79 like cancer cell probing [14].

80 We developed finite-element models (FEMs) that include the
 81 topological information of adherent stem cells (shape and di-
 82 mensions), AFM tips (conical and spherical), and biological
 83 structure (cell models with cytoplasm layers, cytoskeletons, and
 84 nucleus) [15]. To validate the viscoelastic and nonlinear dynam-
 85 ics of the cell in real time using finite-element procedures, we
 86 report herein the mechanical characterization of single mESC
 87 indentation-relaxation for different colonies of fixed and live
 88 undifferentiated mESCs (termed, *undiff mESC*). The proposed
 89 indentation tests use microsphere-modified AFM probes in or-
 90 der to estimate the global elastic modulus of the cell to reflect
 91 the true global response of a mESC [14], [15]. This paper is a
 92 subsequent contribution to [27], where various analytical mod-
 93 els are tested and compared to estimate the mechanical prop-
 94 erties of the biological cells. Based on the proposed physical
 95 model, we have developed a computer-based training system
 96 using force feedback to simulate cell indentation procedures in
 97 virtual environments (VEs) for biologist training. The paper con-
 98 sists of five sections. In Section II, we present the biomechanical
 99 finite-element approach dedicated to real-time indentation. In
 100 Section III, we present the methodology of the VE system for
 101 cell indentation. In Section IV, we present the experimental
 102 setup using the force microscopy. In Section V, we make a
 103 comparative study between the physics-based model and exper-
 104 imental data. In Section VI, we present our real-time cell
 105 indentation simulator using a spherical tip. Finally, we give
 106 some concluding remarks and the directions for future work.

107 II. PHYSICS-BASED DEFORMABLE MODELS FOR CELL 108 NANOINDENTATION SIMULATION

109 In this section, we present further developments to the linear
 110 elastic mass–tensor model introduced by [16], [17] and extended
 111 by Schwartz *et al.* [18]. In our paper, the cell is modeled by a
 112 volume object discretized into a conformal tetrahedral mesh as
 113 defined by finite-element theory.

114 Using classical notations, inside each tetrahedron T^k , the
 115 displacement field, defined by a linear interpolation $[\mathbf{N}^k]$ of the
 116 nodal displacement vector $\{\mathbf{u}^k\}$ of the four vertices of tetrahe-
 117 dron, is written as follows:

$$\{\mathbf{U}(\mathbf{x})^k\} = [\mathbf{N}^k(\mathbf{x})] \{\mathbf{u}^k\}. \quad (1)$$

For a linear elastic material, the relation between the Cauchy 118
 stress tensor and the linearized strain tensor is written with 119
 Lamé's coefficient in condensed vector notation as follows: 120

$$\{\boldsymbol{\sigma}\} = \lambda (\{\boldsymbol{\varepsilon}\}_1 + \{\boldsymbol{\varepsilon}\}_2 + \{\boldsymbol{\varepsilon}\}_3) [\mathbf{I}] + 2\mu \{\boldsymbol{\varepsilon}\} \quad (2)$$

where $[\mathbf{I}]$ is the identity matrix. 121

The principle of virtual work applied to a single tetrahedron 122
 T^k leads to the elementary stiffness matrix $[\mathbf{K}^k]$ such that 123
 the elementary nodal force vector acting on a tetrahedron is as 124
 follows: 125

$$\{\mathbf{f}^k\} = [\mathbf{K}^k] \{\mathbf{u}^k\}. \quad (3)$$

This stiffness matrix is composed of a plurality of elementary 126
 submatrices, each connecting the elementary force acting on the 127
 node i to the displacement of the node j 128

$$[\mathbf{K}_{ij}^k] = \frac{1}{36 V^k} \left(\lambda \{\mathbf{m}_i\} \{\mathbf{m}_j\}^T + \mu \{\mathbf{m}_j\} \{\mathbf{m}_i\}^T + \mu \{\mathbf{m}_i\}^T \{\mathbf{m}_j\} [\mathbf{I}] \right) \quad (4)$$

where $\{\mathbf{m}\}$ are unit outward-pointing normals to triangular 129
 faces and V^k is the volume of the tetrahedron T^k . 130

Taking into account the contribution of all adjacent tetrahedra, 131
 the global internal force acting on a node l can be expressed as 132
 follows: 133

$$\{\mathbf{F}_{\text{int}}^l\} = \sum_{k \in \mathcal{V}_l} \left(\sum_{j=1}^4 [\mathbf{K}_{ij}^k] \{\mathbf{u}_j\} \right) \quad (5)$$

where \mathcal{V}_l is the neighborhood of vertex l (i.e., the tetrahedra 134
 containing node l). 135

The tensors $[\mathbf{K}_{ij}^k]$ depending on the remaining geometry and 136
 Lamé's coefficients are constant. They can be precomputed in an 137
 offline phase. This is the essential advantage of the mass–tensor 138
 approach, which makes it useful for real-time application. 139

140 A. Cell Modeling Using FEM

In our paper, the mESC is meshed with 3-D first-order tetra- 141
 hedral elements, as shown in Fig. 2. The geometrical character- 142
 istics of the cells were obtained from experimental data using 143
 the AFM system. This includes shape, diameter, and height of 144
 the cells. The mesh that has been used for the following simu- 145
 lations is composed of 835 vertices and 2456 tetrahedra, where 146
 all vertices are free to move except the ones in contact with the 147
 petri dish, which are fixed. 148

149 B. Dynamic Model

The equation of motion of a vertex l of the cell mesh can be 150
 written as follows: 151

$$M^l \{\ddot{\mathbf{u}}_l\} + \gamma^l \{\dot{\mathbf{u}}_l\} + \sum_{k \in \mathcal{V}_l} \left(\sum_{j=1}^4 [\mathbf{K}_{ij}^k] \{\mathbf{u}_j\} \right) = \{\mathbf{F}_{\text{ext}}^l\} \quad (6)$$

where M^l and γ^l are, respectively, the mass and damping coef- 152
 ficients of each vertex. 153

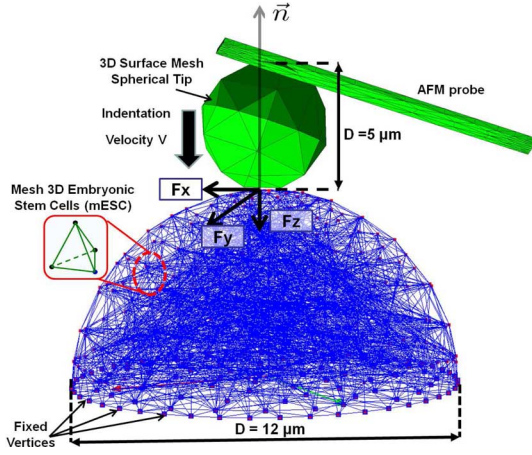


Fig. 2. Representation of the 3-D mESC mesh using P1 tetrahedral finite elements.

154 To solve the dynamic system, we tested different integration
 155 schemes (implicit and explicit) taking into account the tradeoff
 156 between real-time simulation and haptic stability requirements.
 157 We choose the explicit centered finite-difference scheme defined
 158 by

$$\begin{cases} u(t+h) = u(t) + \dot{u}(t)h + \frac{1}{2}\ddot{u}(t)h^2 + O(h^3) \\ \dot{u}(t+h) = \dot{u}(t) + \frac{1}{2}\ddot{u}(t)h + \frac{1}{2}\ddot{u}(t+h)h + O(h^3) \end{cases} \quad (7)$$

159 where h is the temporal integration step size chosen to satisfy
 160 the Courant–Friedrich–Lewy condition.

161 The discrete equation of motion of a vertex l reads

$$\begin{aligned} & \left(M^l + \frac{h}{2}\gamma^l \right) \{ \ddot{\mathbf{u}}_l \}_{(t+h)} \\ &= \{ \mathbf{F}_{\text{ext}}^l \}_{(t+h)} - \gamma^l \left(\{ \dot{\mathbf{u}}_l \}_{(t)} + \frac{h}{2} \{ \ddot{\mathbf{u}}_l \}_{(t)} \right) \\ & - \sum_{k \in \mathcal{V}_l} \left(\sum_{j=1}^4 [\mathbf{K}_{ij}^k] \left(\{ \mathbf{u}_j \}_{(t)} + h \{ \dot{\mathbf{u}}_j \}_{(t)} + \frac{h^2}{2} \{ \ddot{\mathbf{u}}_j \}_{(t)} \right) \right) \end{aligned} \quad (8)$$

162 from which we can compute the new vertices for accelerations,
 163 velocities, and positions.

164 III. VE SYSTEM FOR CELL INDENTATION

165 Fig. 3 shows the architecture of the real-time virtual-reality-
 166 based cell indentation simulator. This includes the computer
 167 generated models of stem cells, the AFM nanoindenter, the col-
 168 lision detection algorithm, physics-based models of deformable
 169 mESC, and the haptic interaction controller. The operator will
 170 be able to interact with the 3-D model of the cell using his
 171 sense of vision as well as actively manipulate using his sense
 172 of touch. Usually, real-time graphics translates to an update
 173 rate of 25 Hz, stable haptic interaction in VE's requires much
 174 higher update rate of around 1 kHz. As soft biological tissues
 175 exhibit complex viscoelastic and nonlinear properties, the real-
 176 time interaction imposes severe restrictions on the FEM model.

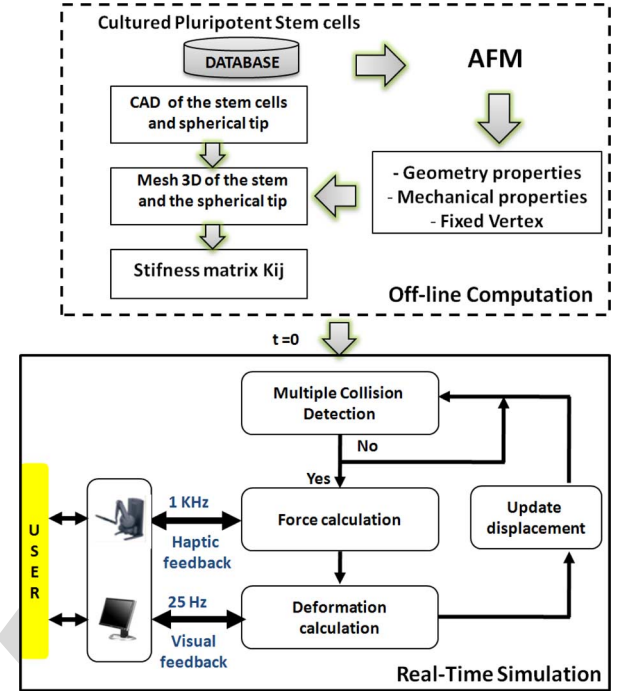


Fig. 3. Computational architecture for simulating force-reflecting deformable cell indentation in a VE. The figure shows the two simulation phases used for the real-time indentation of the cell: 1) *offline* precalculations of stiffness matrices and 2) simulation of visual and haptic interaction.

A significant difficulty of using the finite-element technique for
 177 real-time simulation is that it is computationally costly. These
 178 mesh-based schemes also require an expensive numerical in-
 179 tegration operation for the computation of the system stiffness
 180 matrices. For this reason, we adopted a computational architec-
 181 ture with two simulation stages: 1) an *offline* computation of
 182 stiffness matrices for each triangular element and 2) visual and
 183 haptic interaction for real-time indentation simulation.
 184

185 A. Offline Computation

186 The proposed computational methodology contains an offline
 187 precalculation step. The most costly and time-consuming opera-
 188 tions are realized during this step. The database contains various
 189 developed models of the different stem cells geometry (bone
 190 marrow, pancreatic, heart muscle, nerve cells, etc.). Using AFM
 191 image analysis (ImagePro software from AFM Asylum), we de-
 192 rived different geometric stem cell models based on a commer-
 193 cial computer-aided design (CAD) package (MARC-ADAMS).
 194 The meshing of the 3-D internal structure of the mESC is then
 195 carried out through a dedicated 3-D meshing software (GID
 196 software) modeled in exact dimensions. Although they were
 197 displayed as 3-D texture-mapped objects to the user, they were
 198 modeled as connected line segments to reduce the number of
 199 collision computations during real-time interactions.

200 The adherent cell was approximated as an assembly of dis-
 201 crete tetrahedral elements interconnected to each other through
 202 a fixed number of nodes (see Fig. 2). The displacements of
 203 these nodal points for applied external indentation forces were
 204 the basic unknowns of our FEM analysis. The coordinates

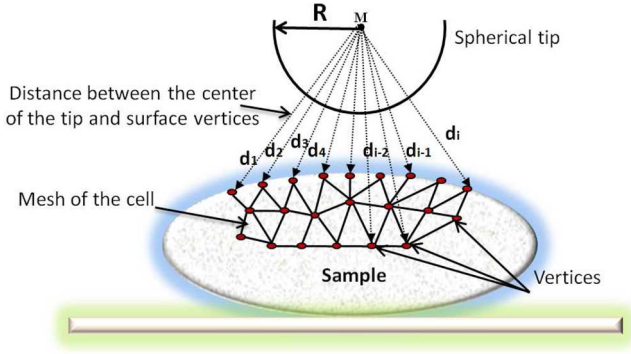


Fig. 4. Scheme illustrating the multiple collision detection algorithm using a simple sphere-points intersection problem: the distance between the center of the spherical tip and the nodes of the cell surface is calculated in real time.

of vertices, the tetrahedron indexing, and the connectivity of vertices were derived from the geometric model. Then, the mechanical (Young's modulus and stiffness) and geometrical (height and shape) properties of the cell structure are determined through AFM indentation experiments. Finally, these properties are passed to the FEM algorithm for calculating the physically based behavior of biological soft tissues. The offline calculation of all tensors K_{ij} allows precalculation stiffness values for preselected cells such as, for example, live and fixed cells.

B. Real-Time Indentation

Since haptics and graphics have different update frequencies, we implemented separate threads to update the loops. During the indentation task, the contact between the tip and the cell must occur at a special set of points (nodal points).

1) *Collision Detection*: An important part of biological cantilever tip-soft cell modeling is the fast collision detection algorithm. The collision detection is the first step to carry out realistic interactions. However, locating the contact primitive (e.g., facets) between two objects may be computationally expensive, especially if the objects are composed of a large number of polygons. In general, to determine a contact between two virtual objects, we compute the distance between them. If this distance is negative, then the objects are in contact. We may continue by determining the set of geometric objects or entities that collide. The cantilever tip in VE can be modeled as a complex 3-D object, composed of numerous surfaces, edges, and vertices, but this constitutes a challenging task. For the purpose of detecting collisions, we used a point-based representation of the cantilever tip as a simple spherical object and employed a simple sphere-points intersection algorithm with local search technique. It calculates the distance between the center of the spherical tip and the nodes of the cell surface (see Fig. 4).

2) *Penalty-Based Models*: In an interaction model, a collision between two objects may generate local and global deformations of the colliding objects, new internal forces, changes in the velocities of the involved objects, and topology modifications, such as puncture or cutting. Physical simulations suppose that these events are related to the mechanical properties of the objects. In our simulations, we used the penalty-based models.

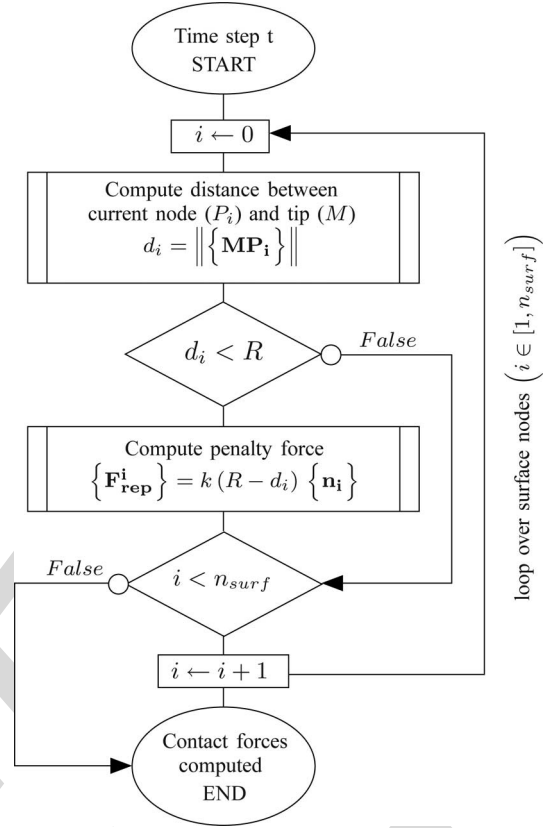


Fig. 5. Main flowchart for multi collision detection: M represents the position of the center of the spherical tip and R its radius, i represents the current external cell surface node, d_i , $\{F_{rep}^i\}$ and $\{n_i\}$ represent, respectively, the computed distance between the center of the tip and the current node, the repulsive force, and the outward-pointing contact normal at node i .

These models consider the reactive force that first reduces the relative speed of the two objects, and finally, repulses them from each other, as proportional to the deformation. Thus, the reactive force is a function of the local interpenetration depth δ . The normal repulsive force F_{rep} can be written as follows:

$$\{F_{rep}\} = \lambda \delta \{n\} \quad (9)$$

where λ is the penalty of the interpenetration and $\{n\}$ is the outward-pointing normal.

These contact forces are evaluated at each time step t (see Fig. 5). They represent the new updated external forces acting upon our stem cell.

3) *Virtual Coupling for Stability of the Haptic Rendering*: The internal operating loop of the haptic interface requires an update frequency around 1 kHz. However, the FEM update frequency is in the range of 25–30 Hz. This frequency difference threatens the coherence between both systems leading to instabilities of the user haptic rendering. We adopted the solution to use a virtual coupling model defined in [24].

This approach introduces a virtual passive link between the simulation model and the haptic interface in order to ensure the stability and the performance of the system [see Fig. 6(a)]. When we combine the impedance display implementation with an appropriate virtual coupling network, we get the admittance

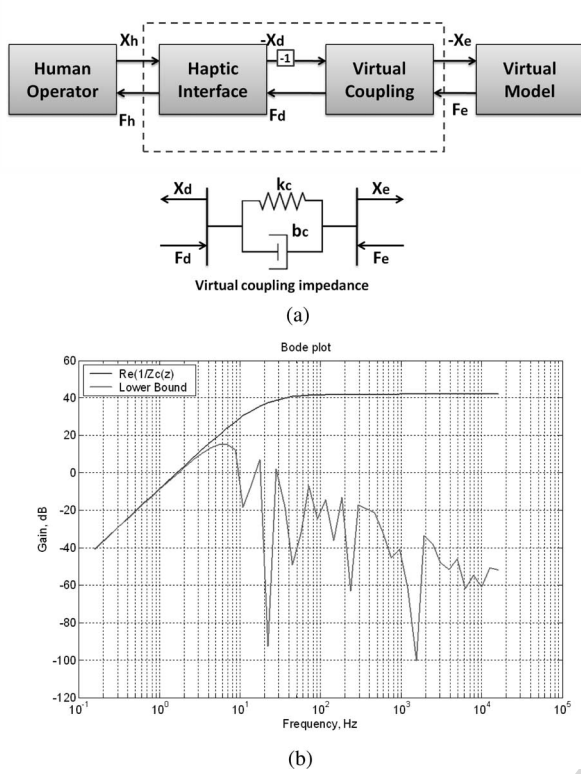


Fig. 6. Stability conditions using the virtual coupling model. (a) Model of virtual coupling. (b) Impedance display virtual coupling.

266 matrix for the combined interface. The linear two-port is said
 267 to be absolutely stable if there exists no set of passive termi-
 268 nating one-port impedances for which the system is unsta-
 269 ble. Llewellyn's stability criteria [25] provides both necessary
 270 and sufficient conditions for absolute stability of linear two-
 271 ports. We get the conditions for absolute stability of the haptic
 272 interface

$$\text{Re}(Z_{di}(z)) \geq 0, \quad \frac{1}{Z_{cvi}(z)} \geq 0 \quad (10)$$

$$\cos(\angle ZOH(z)) + \frac{2\text{Re}(Z_{di}(z))\text{Re}(1/Z_{cvi}(z))}{|ZOH(z)|} \geq 1 \quad (11)$$

274 where $Z_{cvi}(z)$ is the virtual coupling impedance (k_c , b_c),
 275 $ZOH(z)$ is a zero-order holder, and $Z_{di}(z)$ is the PHANTOM
 276 impedance. The inequality (11) can be rewritten to get an ex-
 277 plicit expression of absolute stability of the haptic interface

$$\text{Re}\left(\frac{1}{Z_{cvi}(z)}\right) \geq \frac{1 - \cos(\angle ZOH(z))}{2\text{Re}(Z_{di}(z))} |ZOH(z)|. \quad (12)$$

278 For the virtual coupling, the impedance display induces a limit
 279 on the maximum impedance which can be rendered. We use (12)
 280 to find the virtual coupling, which makes the haptic interface
 281 absolutely stable. Using the design bound of the best performing,
 282 absolutely stabilizing virtual coupling. These parameters are
 283 found as $b_c = 0.008$ N/(mm/s) and $k_c = 2.3$ N/mm (plotted on
 284 Fig. 6(a) as a thin red line). The left side of (12) with the resulting
 285 values is plotted on Fig. 6(b) as a bold blue line.

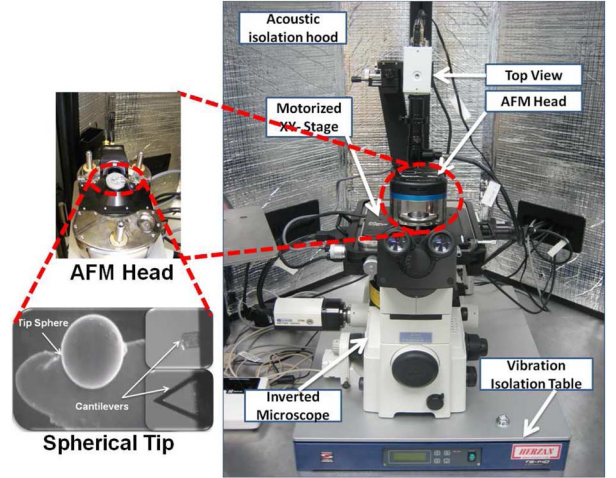


Fig. 7. Experimental setup based on AFM system for mESC indentation studies. The nanoindenter is constituted by a spherical probe ($5 \mu\text{m}$ diameter) attached to a silicon nitride cantilever from Novascan Technologies.

IV. EXPERIMENTAL SETUP

A. Description

288 The experimental tests of the cell indentations were per-
 289 formed using an AFM (Model: MFP-3D-BIO, Asylum Re-
 290 search, Santa Barbara, CA). The AFM is integrated with a top
 291 view module and mounted on an active vibration isolation ta-
 292 ble manufactured by Herzan (Laguna Hills, CA), as shown in
 293 Fig. 7. The top view module enables viewing of cells and easy
 294 alignment of the laser beam on the AFM cantilever. The XY
 295 stage (manual) allows the user to position the cell beneath the
 296 cantilever tip of the AFM. The entire AFM setup is enclosed in
 297 an acoustic isolation chamber to prevent acoustic noise from in-
 298 terfering with the AFM measurements. The x - and y -axis
 299 ranges of the scan head are $90 \mu\text{m}$. The z -axis scan range is $40 \mu\text{m}$.
 300 The AFM system is used to obtain force and cell deformation
 301 data from biological samples.

302 The cantilever is moved by the piezoelectric scanner in z -
 303 direction toward the cell. The deflection of the cantilever is
 304 detected by a photodiode when the tip comes in contact with the
 305 cell. When the tip of the cantilever is in contact with the cell,
 306 the initial cantilever deflection (d_0), and initial cantilever movement
 307 in (z) direction (z_0) are stored. As the cantilever moves further
 308 in z -direction and deforms the cell, the final cantilever deflection
 309 (d_1) and the cantilever movement (z_1) are obtained [21].

310 The stiffness depends not only on the Young's modulus but
 311 also on the geometry of the tip-surface contact. Therefore, the
 312 geometry and spring constant of the cantilever are calibrated in
 313 the same way for live and fixed cells. In the following exper-
 314 iments, we employed two types of cantilevers attached with a
 315 spherical probe ($5 \mu\text{m}$ in diameter).

- 316 1) Silicon nitride cantilever with a spring constant of 0.06
 317 N/m (Novascan Technologies, Inc., Ames, IA) for live
 318 cells.
- 319 2) Silicon cantilever with a spring constant of 1.75 N/m (No-
 320 vascan Technologies, Inc.) for fixed cells.

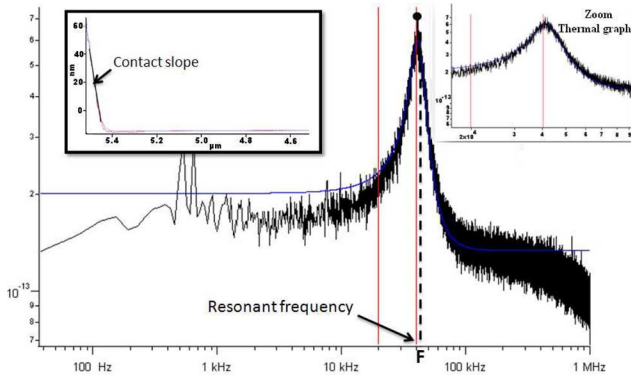


Fig. 8. Thermal graph measurement of a spherical tip cantilever ($K_c = 1.29$ N/m) used for fixed cells indentation. The left corner inset represents the contact slope from a force curve to determine the sensitivity of the cantilever and (deflection volt = 98.01 nm/v) and the right corner inset represents the zoom-in view of the resonant frequency plot ($F = 46.043$ kHz).

321 B. AFM Calibration

322 Before any experiment is done, it is necessary to perform
 323 a calibration of the cantilever to determine its resonant frequency.
 324 The resonant frequency measurement enables the user
 325 to determine if it lies within the specified range provided by
 326 the manufacturer in order to detect probe defects. The determi-
 327 nation of spring constant K_n by the thermal method involves
 328 both measurements : 1) sensitivity in nanometer per volt and 2)
 329 first-resonant frequency of the cantilever.

- 330 1) Determine the contact slope of the indentation curve on a
 331 hard surface to determine the sensitivity of the cantilever
 332 (see left corner inset of Fig. 8).
- 333 2) Measure the thermal power spectral density (PSD) to de-
 334 termine the resonant frequency of the cantilever and to
 335 confirm that the cantilever and light source are aligned.
 336 This confirms that the system is functioning properly. The
 337 results are plotted as the deflection amplitude data using
 338 Fourier transform in meter per square root hertz versus the
 339 frequency in hertz (see right corner inset of Fig. 8). The
 340 spring constant of the cantilever was determined experi-
 341 mentally for each tip used in our studies using the IGOR
 342 software interface supplied by Asylum Research.

343 Finally, the exact K_c values after the calibration procedure
 344 are $K_c = 0.08745$ N/m for live cells and $K_c = 1.29$ N/m for
 345 fixed cells.

346 V. MESC CHARACTERIZATION

347 A. Cell Culture Preparation

348 It is recommended to use specially coated dishes or slides
 349 to facilitate adhesion. Such petri dishes or slides are available
 350 commercially. Cells should be attached to some rigid substrate,
 351 usually either a slide or the bottom of a petri dish. In our study,
 352 the mESC R1 (SCRC-1011, American Type Culture Collection
 353 (ATCC), Manassas, VA) were grown on 0.1% gelatin-coated
 354 plates in the absence of feeder cells. The ES medium consisted
 355 of 1000 U/mL leukemia inhibitory factor (LIF, ESGRO, Chemi-
 356 con, Temecula, CA), 15% fetal bovine serum (FBS) (Invitro-

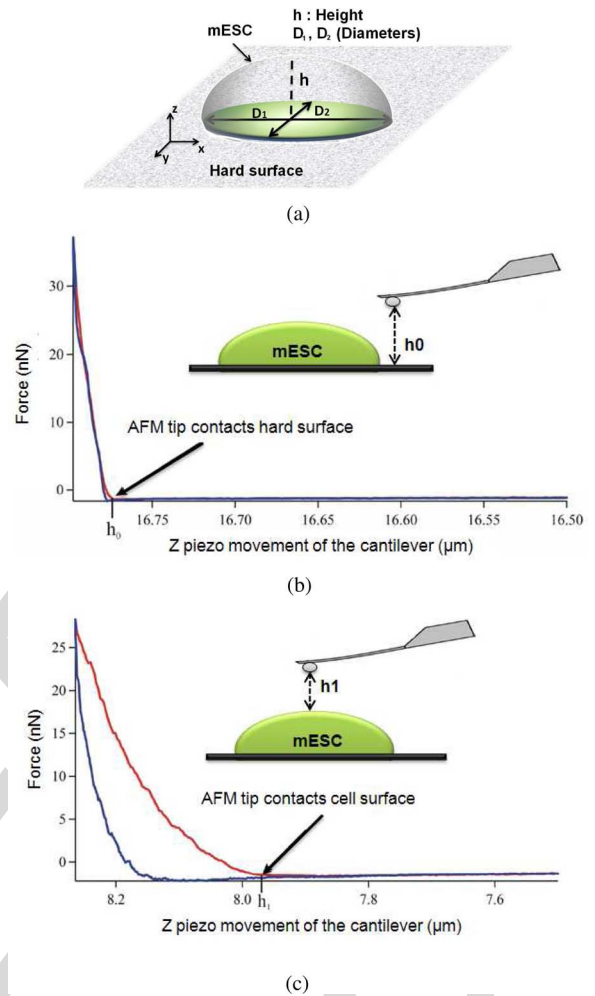


Fig. 9. How to determine the height of the stem cell through indentation curves? (a) mESC geometric parameters. (b) Force (F) versus piezo movement of the cantilever (Z) on a hard surface. The red and blue lines represent the loading and unloading curve, respectively. (c) Force (F) versus piezo movement of the cantilever (Z) on an mESC surface.

357 gen), and basic medium that included Knockout Dulbecco's
 358 modified Eagle's medium (Invitrogen), 2 mM L-glutamine, 1x
 359 nonessential amino acids, and 0.1 mM mercaptoethanol. Differ-
 360 entiation was induced by removal of LIF from the medium.
 361 Prior to the experiments, cells were dispersed using trypsin to
 362 obtain single cells and were plated on 60 mm tissue culture petri
 363 dish. Fixed mESC were obtained by treating live cells with with
 364 4% formaldehyde for 10 min, washed and stored in phosphate
 365 buffered saline.

366 B. Geometry of the Cell and Cantilever Interaction 367 With the Surface

368 The cells were determined to be 10 and 15 μm in diameter
 369 using the cantilever tip to measure the interaction force between
 370 the cell and the tip. The geometric modeling of the cell re-
 371 quires the determination of the height, adhesion surface, and
 372 cell's contour [see Fig. 9(a)]. The mESC height was calculated
 373 from the force-displacement curve obtained on hard and mESC
 374 surfaces. First, the initial height h_0 is measured by bringing

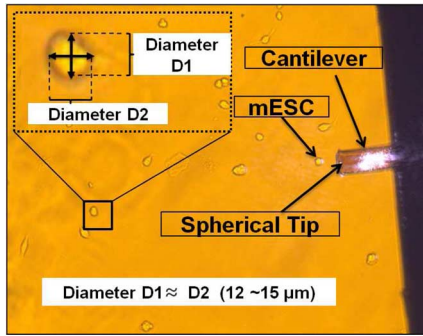


Fig. 10. How to determine the diameter of the stem cell through image analysis (ImagePro software from Asylum)?

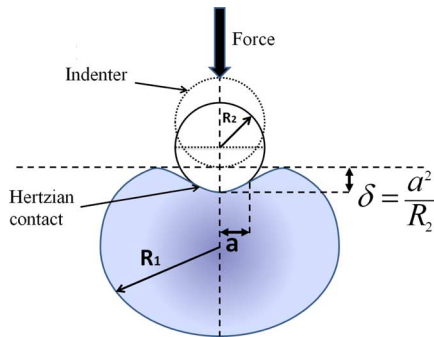


Fig. 11. Hertzian contact with a spherical tip: F is the loading force, and δ and a are the indentation and contact radius, respectively.

375 the AFM tip in contact with the hard surface near mESC [see
376 Fig. 9(b)], and then, the second height h_1 by bringing the AFM
377 tip on mESC surface [see Fig. 9(c)]. The mESC height is the
378 difference $h = h_0 - h_1$. Second, to determine the mESC shape,
379 AFM imaging experiments are carried out in order to deter-
380 mine the mESC's contour and diameter using image processing
381 techniques [19]. From Fig. 10, we selected mESC shapes with
382 circular adhesion surfaces. As illustration, the left corner inset
383 shows a typical circular mESC, where D_1 and D_2 stands for the
384 horizontal cells diameters ($D_1 \simeq D_2$ close to $12 \mu\text{m}$).

385 C. Mechanical Characterization of mESC Using Spherical Tip

386 To estimate the mechanical properties of biological cells using
387 the AFM, various analytical models can be used to identify the
388 Young's modulus of mESC in live as well as fixed cells.

389 1) *Hertz Contact*: The Hertz contact model has been used
390 extensively by the AFM community to quantify the mechanical
391 property of biological samples using AFM [8], [9]. The Hertz
392 contact model describes the simple case of elastic deformation
393 of two perfectly homogeneous smooth surfaces touching under
394 load (see Fig. 11). The forces measured are dominated by the
395 elastic properties. In our work, the geometry of the tip used for
396 all experiments is spherical. The mechanical interaction between
397 the spherical tip and mESC can be described by the Hertz contact
398 model of two elastic bodies [20]. The model assumes that:

399 1) the material properties of the tip and the cell are *isotropic*
400 *and homogeneous*;

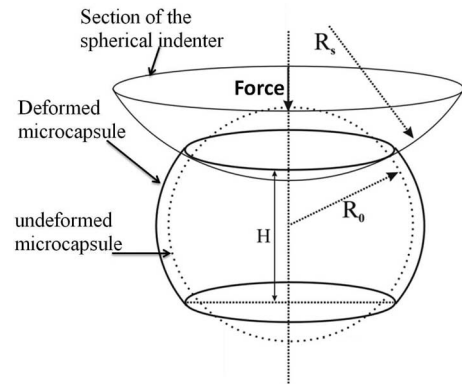


Fig. 12. Schematic of the microcapsule indented by a sphere with a loading force F .

- 2) the normal contact of the two bodies is *adhesionless and frictionless*;
- 3) the contact geometry is assumed to be *axisymmetric, smooth, and continuous*.

The relationship between the indentation δ , and the loading force F is given by

$$F = \frac{4}{3} E^* R^{\frac{1}{2}} \delta^{\frac{3}{2}} \quad (13)$$

where E^* and R are the combined modulus and the relative curvature of the tip and the cell, respectively,

$$\frac{1}{E^*} = \frac{1 - \nu_1^2}{E_1} + \frac{1 - \nu_2^2}{E_2} \quad (14)$$

$$\frac{1}{R} = \frac{1}{R_1} + \frac{1}{R_2} \quad (15)$$

where (E_1, ν_1, R_1) and (E_2, ν_2, R_2) represents the elastic modulus, Poisson's ratio, and the radius of the cell and spherical indenter, respectively. The elastic modulus of the silicon nitride cantilever and silicon cantilever were 222.22 and 168.17 GPa used for the live and fixed cells, respectively. The elastic modulus of the cells is in range of kilopascals [11]. Hence, our assumption that the tip used for probing is infinitely stiff ($E_2 \gg E_1$) compared to the cell and Hertz contact model is valid. Thus, (13) can be rewritten as follows:

$$F = \frac{4E_1}{3(1 - \nu_1^2)} R^{\frac{1}{2}} \delta^{\frac{3}{2}} \quad (16)$$

where E_1 and ν_1 represent the elastic modulus and Poisson's ratio of the cell assuming $\nu_1 = 0.5$.

2) *Capsule Contact*: The second model considers the biological cell to be composed of a cell membrane and cytoplasm (see Fig. 12). It is the capsule model. The model assumes that the cell membrane is a thin film and that the inner cytoplasm provides a uniform hydrostatic pressure on the membrane [22]. The model assumes the following.

- 1) The cell membrane is linearly elastic.
- 2) The deformation of the cell membrane is caused by stretching and bending of the cell membrane.
- 3) The cell is free of initial membrane stress or residual stress.
- 4) The cell volume is constant.

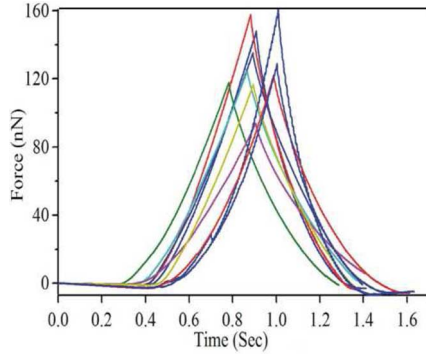


Fig. 13. Load-unload versus time for fixed undiff.

431 The force (F) and relative deformation (ε) relationship is
 432 given by [22]

$$F = \frac{2\pi h R_0 [1 + R_0 / (2R_s)]^2 E}{(1 - \nu)(1 + R_0 / R_s)^4} \varepsilon^3 + \frac{\pi}{2\sqrt{2}} E h^2 \sqrt{\varepsilon} \quad (17)$$

433 where E , ν , and h represent the elastic modulus, Poisson's rati-
 434 on, and the thickness of the capsule membrane, respectively.
 435 R_0 and R_s are the radius of the capsule and the spherical indenter,
 436 respectively. ε is the relative deformation of the capsule
 437 related to the indentation by $\varepsilon = \delta / 2R_0$. The first and the second
 438 term in (17) represent stretching and bending of the cell mem-
 439 brane, respectively, however, in our experiments, the bending
 440 deformation term can be neglected for $\varepsilon > 0.15$ [22].

441 D. Comparison Between Experimental Data 442 and Analytical Models

443 We evaluated and compared the different analytical mod-
 444 els to determine whether they appropriately predict the force-
 445 indentation relationship of mESC. Figs. 13 and 16 present the
 446 force (loading and unloading) versus time of fixed and live
 447 undiff mESC. From all these experiments, we observed adhesion
 448 force in few of the cases. For live mESC probed by a spherical
 449 tip, the adhesion force existed only for one sample. However,
 450 their amplitudes were negligible (less than 0.05 times the maxi-
 451 mum indentation) compared to the maximum indentation forces
 452 (more than 0.1 times the maximum indentation) observed by
 453 other researchers [28], [29]. In our experiments, we observed
 454 that the adhesion force exists for only one of the live undiff and
 455 does not exist for the fixed cells. This could be due to the wear
 456 of the spherical tip [26]. Since we did not observe any adhesion
 457 forces, we have not used the JKR or DMT models.

458 Figs. 14 and 15 show the experimental force versus cell indenta-
 459 tion curves for live and fixed undiff mESC and the simulated
 460 ones using the capsule and the Hertz models, respectively, for
 461 small and rather larger indentation range. Based on the ex-
 462 perimental investigation, a least-square fit on the dataset (for
 463 the whole cells) and the corresponding r^2 values are calcu-
 464 lated. These results show that the capsule model is inappropriate
 465 for describing mESC indentation. This conclusion is confirmed
 466 by the calculation of r^2 values performed on each mESC (see
 467 Table I). Finally, from our experiments, we infer that Hertz
 468 model appropriately describes the mechanical behavior of the

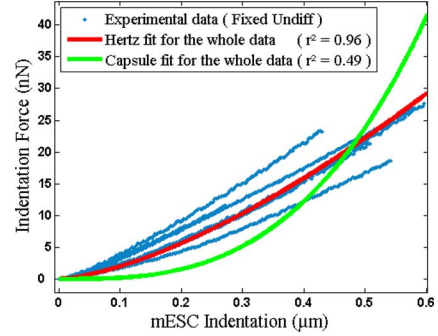


Fig. 14. Fixed undiff (small deformation).

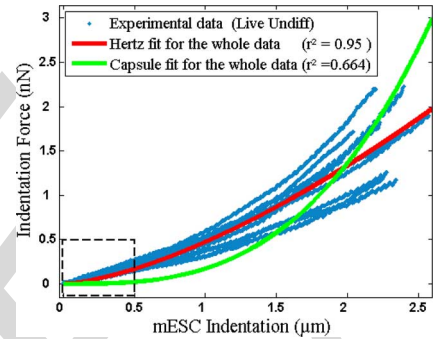


Fig. 15. Live undiff mESC.

TABLE I
 LEAST-SQUARE VALUE (COEFFICIENT CORRELATION) r^2 OBTAINED BY HERTZ
 FIT AND CAPSULE FIT FOR TEN SAMPLES OF LIVE AND FIXED UNDIFF mESC

	Live Undiff mESC		Fixed Undiff mESC	
	Hertz model	Capsule model	Hertz model	Capsule model
Correlation r^2	[0.98...0.99]	[0.68...0.90]	[0.95...0.99]	[0.46...0.73]

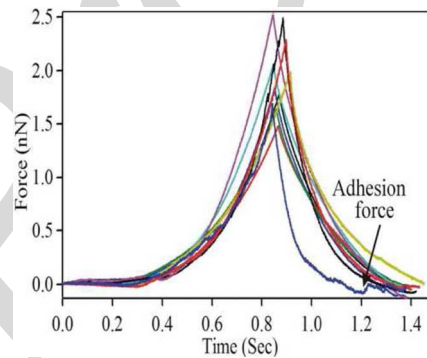


Fig. 16. Load-unload versus time for live undiff.

live and fixed undiff mESC (see Figs. 16– 18). Equation (16) was used to determine the global elastic modulus of the mESC. The average elastic modulus were 17.87 and 0.217 kPa for the fixed and live undiff, respectively, (see Fig. 19). The standard deviations for the fixed and live undiff cells were 3.37 and 0.05 kPa.

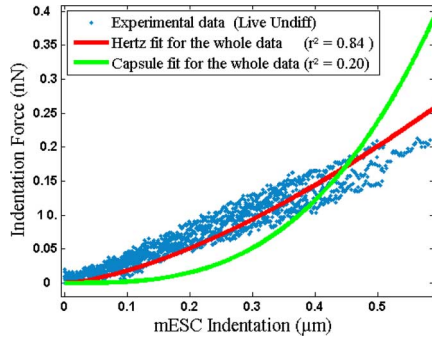


Fig. 17. Live undiff (small deformation).

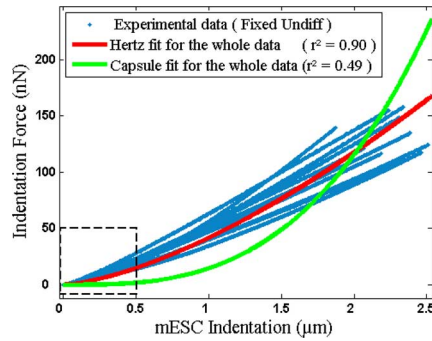
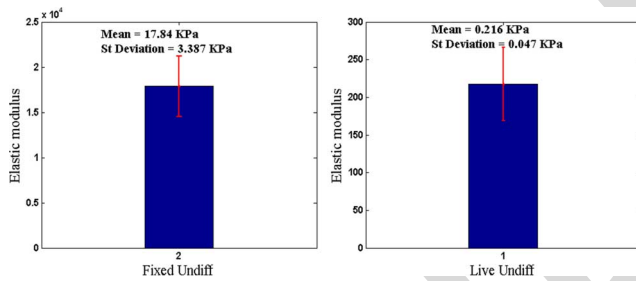


Fig. 18. Fixed undiff mESC.

Fig. 19. Elastic modulus for live and fixed undifferentiated mESC calculated from the Hertz contact model for a small deformation (range $0.5 \mu\text{m}$) with error bar.

475 It should be noticed that alterations in the mechanical properties of cells directly reflect changes in their cellular composition, 476 internal structure (cytoskeleton), and external interactions (cell– 477 cell and/or cell–surface) that occur during differentiation, aging, 478 and other changes in physiological status that are not accounted 479 in the cell indentation model. 480

481 VI. FEM SIMULATION OF CELL INDENTATION

482 Fig. 20 presents a photo of the visual and haptic user interface 483 setup. To test the accuracy and reliability of the proposed 484 user interface system with haptics-enabled simulation, a set of 485 experiments is designed using the setup given in Section II. 486 The geometrical dimensions of the cell are determined through 487 AFM image processing and force curve deflection (diameter and 488 height), and the mechanical properties through AFM indentation 489 studies due to the use of an analytical Hertz contact model.



Fig. 20. Real-time indentation simulator using the Omni haptic interface.

TABLE II
MECHANICAL AND GEOMETRICAL PROPERTIES OF SINGLE CELLS SELECTED
USED IN THE HAPTICS-ENABLED SIMULATOR

	Live mESC	Fixed mESC
Young modulus (KPa)	0.169	26.7
Poisson coefficient	0.49	0.49
Diameter (μm)	$\simeq 12$	$\simeq 12$
Height (μm)	8	8

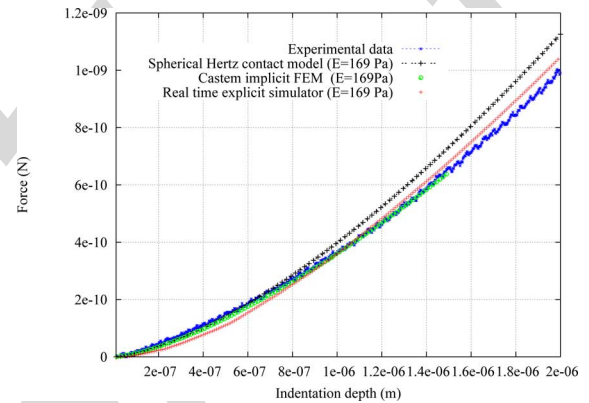


Fig. 21. Force versus indentation for live mESC: comparison between the real-time FEM, the Hertz contact, the quasi-static FEM, and experimental data.

490 These properties (see Table II) are passed to the finite-element 491 simulations.

492 In order to demonstrate the validity of our finite-element 493 indentation model in real time, we compare the simulation 494 indentation using the physics-based FEM model to the incremental 495 data provided by AFM indentation tool and to the response 496 simulated with a commercial FE software. Figs. 21 and 22 present 497 the nonlinear relationship between the penetration distance (in- 498 dentation) and the reaction force for live and fixed undiff cells. 499 The finite-element simulations using the commercial software 500 show good agreement with both Hertz contact model and the 501 experimental indentation data. The developed real-time physics- 502 based FEM simulator demonstrates its accuracy in predicting

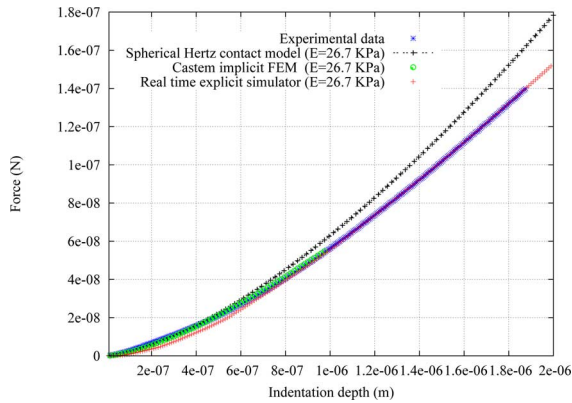


Fig. 22. Force versus indentation for fixed mESC: comparison between the real-time FEM, the Hertz contact, the quasi-static FEM, and experimental data.

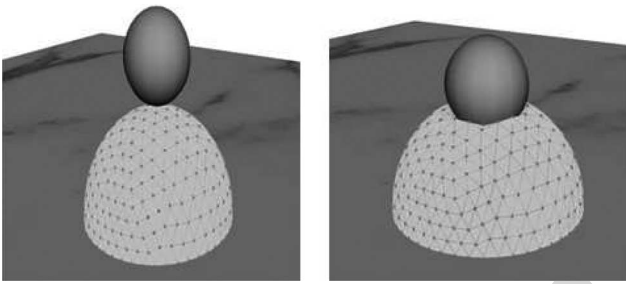


Fig. 23. Simulation of the AFM nanoindentation in VEs. The AFM tip and the stem cell was modeled using finite-element methods. A scene is shown from current training method: the 3-D cell deformation simulation before and after force indentation.

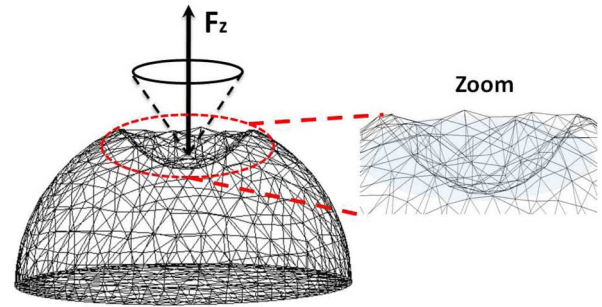


Fig. 24. Depth deformation of the cell: the left picture presents the depth indentation ($2 \mu\text{m}$) of the cell and the right picture shows the zoom of the cell deformation.

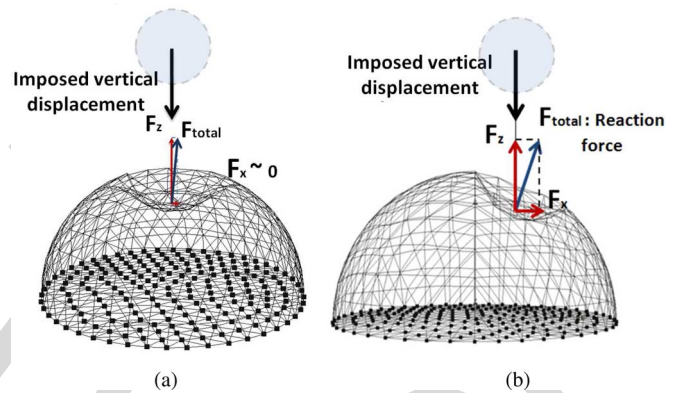


Fig. 25. Force indentation, where the probe is put (a) in the center and (b) off center of the cell.

503 indentation forces for large deformations ($>25\%$) of relative
 504 nominal strain (see Figs. 21 and 22). Fig. 23 illustrates a real-
 505 time cell indentation procedure in a VE for training. It shows a
 506 screenshot of the cell surface deformation. The interaction be-
 507 tween tip and surface can be visualized to the user in 3-D (sphere
 508 and cone represent AFM tip and forces acting on tip, respec-
 509 tively). It shows the indentation (elastic deformation) on the
 510 cell surface and repulsive contact forces while pressing on the
 511 surface. In order to improve the force feedback rendering, we
 512 added a visual metaphor on the computer screen.

513 The rendered reflection force is depicted by a cone represent-
 514 ing the force vector acting on the AFM tip. The visual metaphor
 515 has efficiently proved its effectiveness in AFM-based nanoma-
 516 nipulation simulation [23]. The vector amplitude reflects the
 517 force feedback amplitude and the vector orientation represents
 518 the vertical/horizontal force components. Fig. 24 shows in 3-D
 519 real time the probe tip, zoomed cell indentation, and micro-
 520 forces. The indentation force rendered to the operator is com-
 521 posed of a vertical F_z and horizontal F_x component. For the
 522 specific case corresponding to the center indentation, the tan-
 523 gential component is very small $F_x \approx 0$. Only vertical reflection
 524 force is provided to the user. He cannot haptically feel it dur-
 525 ing simulated indentation tests (see Figs. 25(a) and 26). In the
 526 case, where the AFM probe is pushed away from the center of
 527 the cell, a tangential component exists ($F_x \neq 0$) and the reflec-

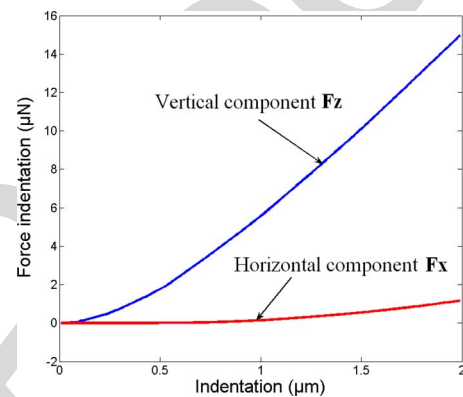


Fig. 26. Horizontal and vertical component force versus indentation: the probe is put in center of the cell.

tion force rendered to the user contains a vertical and tangential
 528 component (see Figs. 25(b) and 27). We did not reproduce this
 529 specific case in our results since the main interest of the paper is
 530 to prove experimentally the real-time explicit simulator results
 531 during centered cell indentations (since only vertical compo-
 532 nents of the force experiments are measured through the AFM
 533 cantilever).
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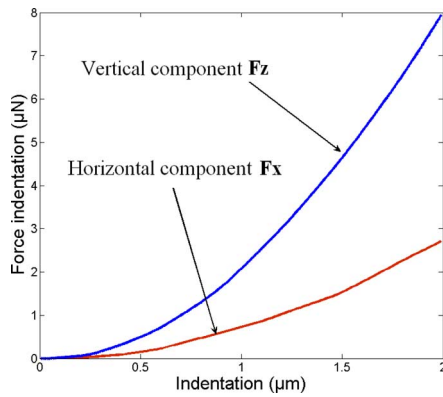


Fig. 27. Horizontal and vertical component force versus indentation: the probe is put off center of the cell.

VII. CONCLUSION

We have developed a computer-based training system to simulate real-time cell indentation procedures in VEs for training biologists. The simulator provides the user with visual and haptic feedback. The simulation of this procedure involves real-time rendering of computer generated graphical images of the AFM tip, physics-based modeling of soft biological tissue, and display of touch and force sensations to the trainee through the simulation of haptic interactions. The real-time haptics-enabled simulator is realistic, since it is mainly based on experimental AFM data (mechanical and geometrical property of mESC) to accurately replicate the indentation task and predict the cell deformation during indentation. We first investigated the challenging issues in the real-time modeling of the biomechanical properties of the cell indentation through FEMs. Compared to experimental AFM indentation results performed on mESC, we can see clearly that the proposed physically based FEM model is able to simulate the cell deformation through real-time simulation constraints. Currently, we are working on integrating nanoscale effects, such as friction, viscosity, puncture, and adhesion forces in the nonlinear FEM model. All modalities will be merged in an ergonomic tool and intelligent biological simulator for stem cell characterization.

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Reality-Based Real-Time Cell Indentation Simulator

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and Jaydev P. Desai, *Senior Member, IEEE*

Abstract—Training simulators that provide realistic visual and haptic feedback during cell indentation tasks are currently investigated. Complex cell geometry inherent to biological cells and intricate mechanical properties drive the need for precise mechanical and numerical modeling to assure accurate cell deformation and force calculations. Advances in alternative finite-element formulation, such as the mass–tensor approach, have reached a state, where they are applicable to model soft-cell deformation in real time. The geometrical characteristics and the mechanical properties of different cells are determined with atomic force microscopy (AFM) indentation. A real-time, haptics-enabled simulator for cell centered indentation has been developed, which utilizes the AFM data (mechanical and geometrical properties of embryonic stem cells) to accurately replicate the indentation task and predict the cell deformation during indentation in real time. This tool can be used as a mechanical marker to characterize the biological state of the cell. The operator is able to feel the change in the stiffness during cell deformation between fixed and live cells in real time. A comparative study with finite-element simulations using a commercial software and the experimental data demonstrate the effectiveness of the proposed physically based model.

Index Terms—Atomic force microscope (AFM), finite-element modeling (FEM), haptics, modeling, real-time interaction, stem cell.

I. INTRODUCTION

MECHANICAL manipulation and characterization of biological cells is currently one of the most promising research areas in the field of biorobotics applied to cellular level interactions. The mechanical properties, such as elasticity, membrane tension, cell shape, and adhesion strength, may play an

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The Promise of Stem Cell Research

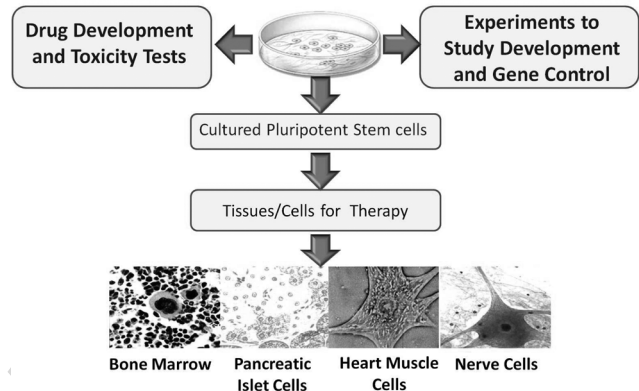


Fig. 1. Stem cells are valuable resources for toxicology, developmental biology, and regenerative medicine research.

important role in cell development and differentiation [1]–[3]. In contrast to adult stem cells, embryonic stem cells differentiate into every cell type in the body, and thus, have a huge therapeutic potential. Stem-cell-based systems offer a very promising and innovative solution for obtaining large number of cells for early efficacy and toxicity screening. Stem cell technology provides a new tool for drug development and a better insight to understand mammalian gene function (see Fig. 1). The mechanical properties of biological cells have been studied with different techniques [4], the most popular being optical tweezers [5], magnetic beads [6], and micropipette aspiration [7]. Different microrobotic systems with various end effectors have been proposed [33], [34]. However, these methods cannot compete with the precision that can be obtained with the atomic force microscopy (AFM) method [8], [32]. AFMs have been widely used in the study of micro- and nanostructures including living cells. Modern AFM techniques allow solving several problems of cell biomechanics by simultaneous evaluation of the local mechanical properties and the topography of living cells, at a high spatial resolution and force sensitivity [9]. In these experiments, an AFM cantilever serves as a microindenter to probe the cell, and further analysis of force-indentation data yields the local Young's modulus. In addition, the AFM indentation technique can be used to characterize the viscoelastic behavior of the cell cytoskeleton [10], including viscosity [11], loss and storage moduli [12], and stress relaxation times. Mechanical phenotyping prove actually to be a valuable tool in the development of improved methods of targeted cellular differentiation of embryonic, induced pluripotent (iPS), and adult stem cells for therapeutic and diagnostic purposes [30], [31]. In this study, we propose to characterize the mechanical behavior of an individual mouse embryonic stem cells (mESC) in undifferentiated state

66 using AFM indentation [13]. As we hypothesize that the cell
 67 stiffness of mESC may reflect the nuclear physical changes of
 68 the cell, it should be possible to detect modifications in stiffness
 69 by probing the cell directly. However, different limitations are
 70 inherent in AFM-based indentation techniques. First, deforma-
 71 tions induced by the AFM probe and cell positions cannot be
 72 seen by the user. Second, the force sensing of the AFM can-
 73 not resolve the force interaction between the AFM tip and the
 74 cell's surface independently in all three dimensions to avoid
 75 cytoskeleton-disrupted experiments. Therefore, the user per-
 76 forms extremely challenging manipulations requiring advanced
 77 manual skills. Modeling the physics of the interaction between
 78 tip and cell is necessary to realize complicated indentation tasks
 79 like cancer cell probing [14].

80 We developed finite-element models (FEMs) that include the
 81 topological information of adherent stem cells (shape and di-
 82 mensions), AFM tips (conical and spherical), and biological
 83 structure (cell models with cytoplasm layers, cytoskeletons, and
 84 nucleus) [15]. To validate the viscoelastic and nonlinear dynam-
 85 ics of the cell in real time using finite-element procedures, we
 86 report herein the mechanical characterization of single mESC
 87 indentation-relaxation for different colonies of fixed and live
 88 undifferentiated mESCs (termed, *undiff mESC*). The proposed
 89 indentation tests use microsphere-modified AFM probes in or-
 90 der to estimate the global elastic modulus of the cell to reflect
 91 the true global response of a mESC [14], [15]. This paper is a
 92 subsequent contribution to [27], where various analytical mod-
 93 els are tested and compared to estimate the mechanical prop-
 94 erties of the biological cells. Based on the proposed physical
 95 model, we have developed a computer-based training system
 96 using force feedback to simulate cell indentation procedures in
 97 virtual environments (VEs) for biologist training. The paper con-
 98 sists of five sections. In Section II, we present the biomechanical
 99 finite-element approach dedicated to real-time indentation. In
 100 Section III, we present the methodology of the VE system for
 101 cell indentation. In Section IV, we present the experimental
 102 setup using the force microscopy. In Section V, we make a
 103 comparative study between the physics-based model and exper-
 104 imental data. In Section VI, we present our real-time cell
 105 indentation simulator using a spherical tip. Finally, we give
 106 some concluding remarks and the directions for future work.

107 II. PHYSICS-BASED DEFORMABLE MODELS FOR CELL 108 NANOINDENTATION SIMULATION

109 In this section, we present further developments to the linear
 110 elastic mass–tensor model introduced by [16], [17] and extended
 111 by Schwartz *et al.* [18]. In our paper, the cell is modeled by a
 112 volume object discretized into a conformal tetrahedral mesh as
 113 defined by finite-element theory.

114 Using classical notations, inside each tetrahedron T^k , the
 115 displacement field, defined by a linear interpolation $[\mathbf{N}^k]$ of the
 116 nodal displacement vector $\{\mathbf{u}^k\}$ of the four vertices of tetrahe-
 117 dron, is written as follows:

$$\{\mathbf{U}(\mathbf{x})^k\} = [\mathbf{N}^k(\mathbf{x})] \{\mathbf{u}^k\}. \quad (1)$$

For a linear elastic material, the relation between the Cauchy
 stress tensor and the linearized strain tensor is written with
 Lamé's coefficient in condensed vector notation as follows:

$$\{\boldsymbol{\sigma}\} = \lambda (\{\boldsymbol{\varepsilon}\}_1 + \{\boldsymbol{\varepsilon}\}_2 + \{\boldsymbol{\varepsilon}\}_3) [\mathbf{I}] + 2\mu \{\boldsymbol{\varepsilon}\} \quad (2)$$

where $[\mathbf{I}]$ is the identity matrix.

The principle of virtual work applied to a single tetrahedron
 T^k leads to the elementary stiffness matrix $[\mathbf{K}^k]$ such that
 the elementary nodal force vector acting on a tetrahedron is as
 follows:

$$\{\mathbf{f}^k\} = [\mathbf{K}^k] \{\mathbf{u}^k\}. \quad (3)$$

This stiffness matrix is composed of a plurality of elementary
 submatrices, each connecting the elementary force acting on the
 node i to the displacement of the node j

$$[\mathbf{K}_{ij}^k] = \frac{1}{36 V^k} \left(\lambda \{\mathbf{m}_i\} \{\mathbf{m}_j\}^T + \mu \{\mathbf{m}_j\} \{\mathbf{m}_i\}^T + \mu \{\mathbf{m}_i\}^T \{\mathbf{m}_j\} [\mathbf{I}] \right) \quad (4)$$

where $\{\mathbf{m}\}$ are unit outward-pointing normals to triangular
 faces and V^k is the volume of the tetrahedron T^k .

Taking into account the contribution of all adjacent tetrahedra,
 the global internal force acting on a node l can be expressed as
 follows:

$$\{\mathbf{F}_{\text{int}}^l\} = \sum_{k \in \mathcal{V}_l} \left(\sum_{j=1}^4 [\mathbf{K}_{ij}^k] \{\mathbf{u}_j\} \right) \quad (5)$$

where \mathcal{V}_l is the neighborhood of vertex l (i.e., the tetrahedra
 containing node l).

The tensors $[\mathbf{K}_{ij}^k]$ depending on the remaining geometry and
 Lamé's coefficients are constant. They can be precomputed in an
 offline phase. This is the essential advantage of the mass–tensor
 approach, which makes it useful for real-time application.

A. Cell Modeling Using FEM

In our paper, the mESC is meshed with 3-D first-order tetra-
 hedral elements, as shown in Fig. 2. The geometrical character-
 istics of the cells were obtained from experimental data using
 the AFM system. This includes shape, diameter, and height of
 the cells. The mesh that has been used for the following simu-
 lations is composed of 835 vertices and 2456 tetrahedra, where
 all vertices are free to move except the ones in contact with the
 petri dish, which are fixed.

B. Dynamic Model

The equation of motion of a vertex l of the cell mesh can be
 written as follows:

$$M^l \{\ddot{\mathbf{u}}_l\} + \gamma^l \{\dot{\mathbf{u}}_l\} + \sum_{k \in \mathcal{V}_l} \left(\sum_{j=1}^4 [\mathbf{K}_{ij}^k] \{\mathbf{u}_j\} \right) = \{\mathbf{F}_{\text{ext}}^l\} \quad (6)$$

where M^l and γ^l are, respectively, the mass and damping coef-
 ficients of each vertex.

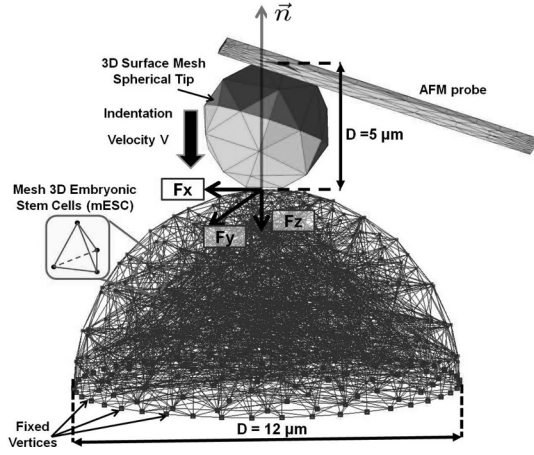


Fig. 2. Representation of the 3-D mESC mesh using P1 tetrahedral finite elements.

154 To solve the dynamic system, we tested different integration
 155 schemes (implicit and explicit) taking into account the tradeoff
 156 between real-time simulation and haptic stability requirements.
 157 We choose the explicit centered finite-difference scheme defined
 158 by

$$\begin{cases} u(t+h) = u(t) + \dot{u}(t)h + \frac{1}{2}\ddot{u}(t)h^2 + O(h^3) \\ \dot{u}(t+h) = \dot{u}(t) + \frac{1}{2}\ddot{u}(t)h + \frac{1}{2}\ddot{u}(t+h)h + O(h^3) \end{cases} \quad (7)$$

159 where h is the temporal integration step size chosen to satisfy
 160 the Courant–Friedrich–Lewy condition.

161 The discrete equation of motion of a vertex l reads

$$\begin{aligned} & \left(M^l + \frac{h}{2}\gamma^l \right) \{ \ddot{\mathbf{u}}_l \}_{(t+h)} \\ &= \{ \mathbf{F}_{\text{ext}}^l \}_{(t+h)} - \gamma^l \left(\{ \dot{\mathbf{u}}_l \}_{(t)} + \frac{h}{2} \{ \ddot{\mathbf{u}}_l \}_{(t)} \right) \\ & - \sum_{k \in \mathcal{V}_l} \left(\sum_{j=1}^4 [\mathbf{K}_{ij}^k] \left(\{ \mathbf{u}_j \}_{(t)} + h \{ \dot{\mathbf{u}}_j \}_{(t)} + \frac{h^2}{2} \{ \ddot{\mathbf{u}}_j \}_{(t)} \right) \right) \end{aligned} \quad (8)$$

162 from which we can compute the new vertices for accelerations,
 163 velocities, and positions.

164 III. VE SYSTEM FOR CELL INDENTATION

165 Fig. 3 shows the architecture of the real-time virtual-reality-
 166 based cell indentation simulator. This includes the computer
 167 generated models of stem cells, the AFM nanoindenter, the col-
 168 lision detection algorithm, physics-based models of deformable
 169 mESC, and the haptic interaction controller. The operator will
 170 be able to interact with the 3-D model of the cell using his
 171 sense of vision as well as actively manipulate using his sense
 172 of touch. Usually, real-time graphics translates to an update
 173 rate of 25 Hz, stable haptic interaction in VE's requires much
 174 higher update rate of around 1 kHz. As soft biological tissues
 175 exhibit complex viscoelastic and nonlinear properties, the real-
 176 time interaction imposes severe restrictions on the FEM model.

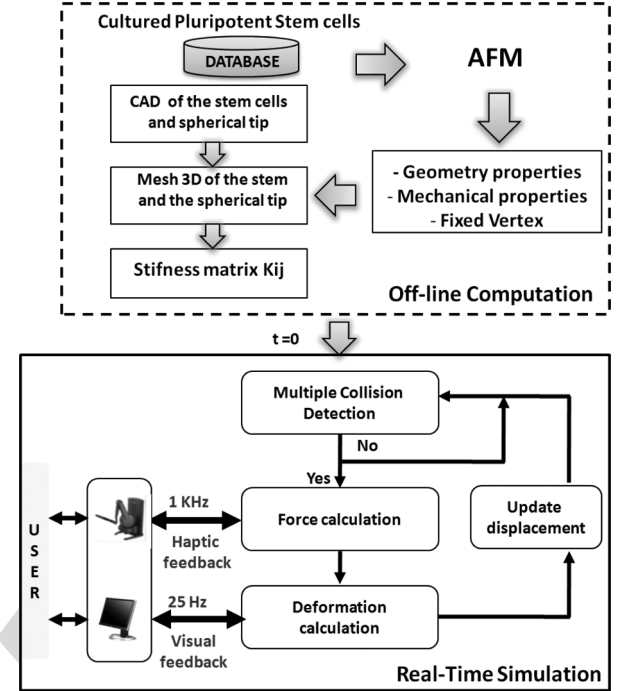


Fig. 3. Computational architecture for simulating force-reflecting deformable cell indentation in a VE. The figure shows the two simulation phases used for the real-time indentation of the cell: 1) *offline* precalculations of stiffness matrices and 2) simulation of visual and haptic interaction.

A significant difficulty of using the finite-element technique for real-time simulation is that it is computationally costly. These mesh-based schemes also require an expensive numerical integration operation for the computation of the system stiffness matrices. For this reason, we adopted a computational architecture with two simulation stages: 1) an *offline* computation of stiffness matrices for each triangular element and 2) visual and haptic interaction for real-time indentation simulation.

185 A. Offline Computation

186 The proposed computational methodology contains an offline
 187 precalculation step. The most costly and time-consuming opera-
 188 tions are realized during this step. The database contains various
 189 developed models of the different stem cells geometry (bone
 190 marrow, pancreatic, heart muscle, nerve cells, etc.). Using AFM
 191 image analysis (ImagePro software from AFM Asylum), we de-
 192 rived different geometric stem cell models based on a commer-
 193 cial computer-aided design (CAD) package (MARC-ADAMS).
 194 The meshing of the 3-D internal structure of the mESC is then
 195 carried out through a dedicated 3-D meshing software (GID
 196 software) modeled in exact dimensions. Although they were
 197 displayed as 3-D texture-mapped objects to the user, they were
 198 modeled as connected line segments to reduce the number of
 199 collision computations during real-time interactions.

200 The adherent cell was approximated as an assembly of discrete
 201 tetrahedral elements interconnected to each other through a
 202 fixed number of nodes (see Fig. 2). The displacements of
 203 these nodal points for applied external indentation forces were
 204 the basic unknowns of our FEM analysis. The coordinates

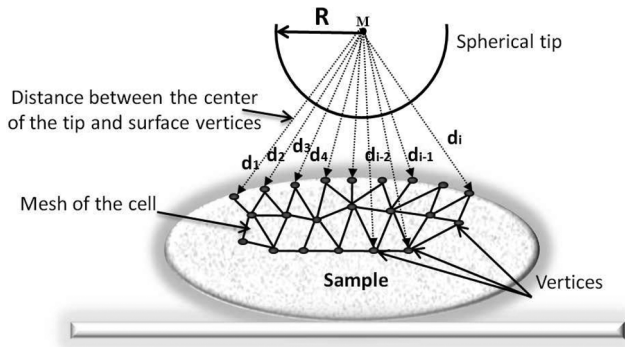


Fig. 4. Scheme illustrating the multiple collision detection algorithm using a simple sphere-points intersection problem: the distance between the center of the spherical tip and the nodes of the cell surface is calculated in real time.

205 of vertices, the tetrahedron indexing, and the connectivity of
 206 vertices were derived from the geometric model. Then, the
 207 mechanical (Young's modulus and stiffness) and geometrical
 208 (height and shape) properties of the cell structure are determined
 209 through AFM indentation experiments. Finally, these properties
 210 are passed to the FEM algorithm for calculating the physically
 211 based behavior of biological soft tissues. The offline calcula-
 212 tion of all tensors K_{ij} allows precalculation stiffness values for
 213 preselected cells such as, for example, live and fixed cells.

214 B. Real-Time Indentation

215 Since haptics and graphics have different update frequencies,
 216 we implemented separate threads to update the loops. During
 217 the indentation task, the contact between the tip and the cell
 218 must occur at a special set of points (nodal points).

219 1) *Collision Detection*: An important part of biological can-
 220 tilever tip-soft cell modeling is the fast collision detection algo-
 221 rithm. The collision detection is the first step to carry out
 222 realistic interactions. However, locating the contact primitive
 223 (e.g., facets) between two objects may be computationally ex-
 224 pensive, especially if the objects are composed of a large number
 225 of polygons. In general, to determine a contact between two vir-
 226 tual objects, we compute the distance between them. If this
 227 distance is negative, then the objects are in contact. We may
 228 continue by determining the set of geometric objects or entities
 229 that collide. The cantilever tip in VE can be modeled as a com-
 230 plex 3-D object, composed of numerous surfaces, edges, and
 231 vertices, but this constitutes a challenging task. For the purpose
 232 of detecting collisions, we used a point-based representation of
 233 the cantilever tip as a simple spherical object and employed a
 234 simple sphere-points intersection algorithm with local search
 235 technique. It calculates the distance between the center of the
 236 spherical tip and the nodes of the cell surface (see Fig. 4).

237 2) *Penalty-Based Models*: In an interaction model, a colli-
 238 sion between two objects may generate local and global defor-
 239 mations of the colliding objects, new internal forces, changes
 240 in the velocities of the involved objects, and topology modifica-
 241 tions, such as puncture or cutting. Physical simulations suppose
 242 that these events are related to the mechanical properties of the
 243 objects. In our simulations, we used the penalty-based models.

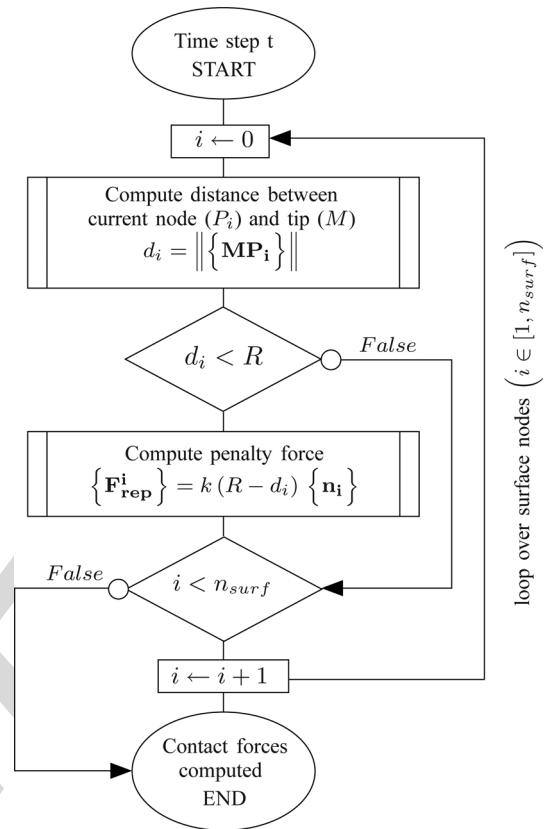


Fig. 5. Main flowchart for multi collision detection: M represents the position of the center of the spherical tip and R its radius, i represents the current external cell surface node, d_i , $\{F_{rep}^i\}$ and $\{n_i\}$ represent, respectively, the computed distance between the center of the tip and the current node, the repulsive force, and the outward-pointing contact normal at node i .

244 These models consider the reactive force that first reduces the
 245 relative speed of the two objects, and finally, repulses them from
 246 each other, as proportional to the deformation. Thus, the reactive
 247 force is a function of the local interpenetration depth δ . The
 248 normal repulsive force F_{rep} can be written as follows:

$$\{F_{rep}\} = \lambda \delta \{n\} \quad (9)$$

249 where λ is the penalty of the interpenetration and $\{n\}$ is the
 250 outward-pointing normal.

251 These contact forces are evaluated at each time step t (see
 252 Fig. 5). They represent the new updated external forces acting
 253 upon our stem cell.

254 3) *Virtual Coupling for Stability of the Haptic Rendering*:
 255 The internal operating loop of the haptic interface requires an
 256 update frequency around 1 kHz. However, the FEM update fre-
 257 quency is in the range of 25–30 Hz. This frequency difference
 258 threatens the coherence between both systems leading to insta-
 259 bilities of the user haptic rendering. We adopted the solution to
 260 use a virtual coupling model defined in [24].

261 This approach introduces a virtual passive link between the
 262 simulation model and the haptic interface in order to ensure
 263 the stability and the performance of the system [see Fig. 6(a)].
 264 When we combine the impedance display implementation with
 265 an appropriate virtual coupling network, we get the admittance

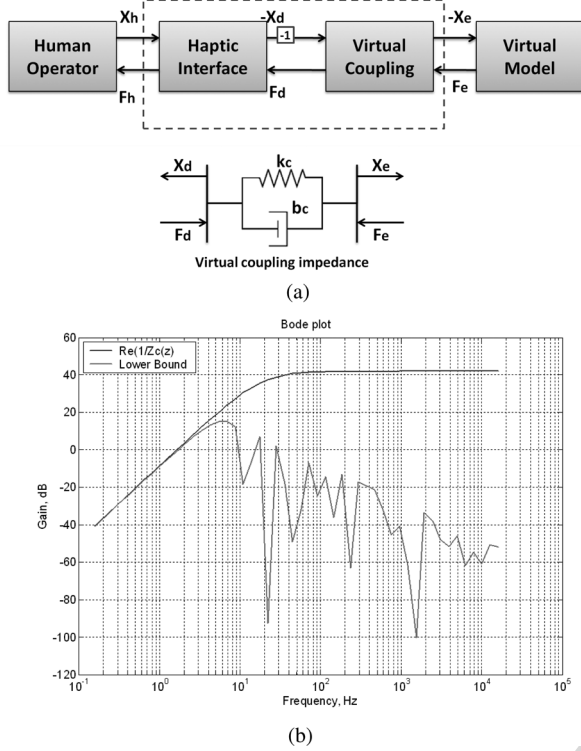


Fig. 6. Stability conditions using the virtual coupling model. (a) Model of virtual coupling. (b) Impedance display virtual coupling.

matrix for the combined interface. The linear two-port is said to be absolutely stable if there exists no set of passive terminating one-port impedances for which the system is unstable. Llewellyn's stability criteria [25] provides both necessary and sufficient conditions for absolute stability of linear two-ports. We get the conditions for absolute stability of the haptic interface

$$\text{Re}(Z_{di}(z)) \geq 0, \quad \frac{1}{Z_{cvi}(z)} \geq 0 \quad (10)$$

$$\cos(\angle \text{ZOH}(z)) + \frac{2\text{Re}(Z_{di}(z))\text{Re}(1/Z_{cvi}(z))}{|\text{ZOH}(z)|} \geq 1 \quad (11)$$

where $Z_{cvi}(z)$ is the virtual coupling impedance (k_c , b_c), $\text{ZOH}(z)$ is a zero-order holder, and $Z_{di}(z)$ is the PHANToM impedance. The inequality (11) can be rewritten to get an explicit expression of absolute stability of the haptic interface

$$\text{Re}\left(\frac{1}{Z_{cvi}(z)}\right) \geq \frac{1 - \cos(\angle \text{ZOH}(z))}{2\text{Re}(Z_{di}(z))} |\text{ZOH}(z)|. \quad (12)$$

For the virtual coupling, the impedance display induces a limit on the maximum impedance which can be rendered. We use (12) to find the virtual coupling, which makes the haptic interface absolutely stable. Using the design bound of the best performing, absolutely stabilizing virtual coupling. These parameters are found as $b_c = 0.008$ N/(mm/s) and $k_c = 2.3$ N/mm (plotted on Fig. 6(a) as a thin red line). The left side of (12) with the resulting values is plotted on Fig. 6(b) as a bold blue line.

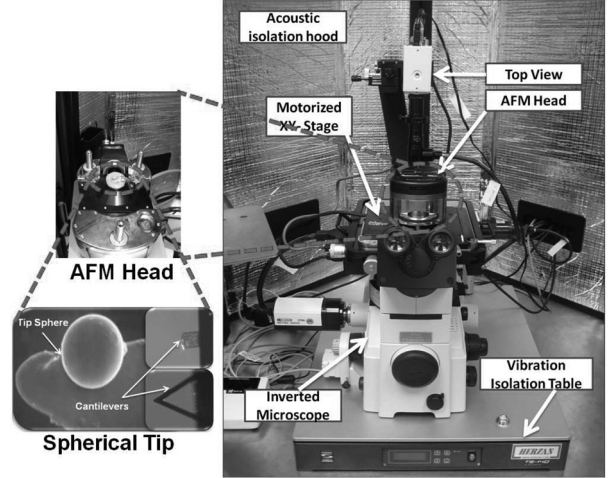


Fig. 7. Experimental setup based on AFM system for mESC indentation studies. The nanoindenter is constituted by a spherical probe ($5 \mu\text{m}$ diameter) attached to a silicon nitride cantilever from Novascan Technologies.

IV. EXPERIMENTAL SETUP

A. Description

The experimental tests of the cell indentations were performed using an AFM (Model: MFP-3D-BIO, Asylum Research, Santa Barbara, CA). The AFM is integrated with a top view module and mounted on an active vibration isolation table manufactured by Herzan (Laguna Hills, CA), as shown in Fig. 7. The top view module enables viewing of cells and easy alignment of the laser beam on the AFM cantilever. The XY stage (manual) allows the user to position the cell beneath the cantilever tip of the AFM. The entire AFM setup is enclosed in an acoustic isolation chamber to prevent acoustic noise from interfering with the AFM measurements. The x - and y -axis ranges of the scan head are $90 \mu\text{m}$. The z -axis scan range is $40 \mu\text{m}$. The AFM system is used to obtain force and cell deformation data from biological samples.

The cantilever is moved by the piezoelectric scanner in z -direction toward the cell. The deflection of the cantilever is detected by a photodiode when the tip comes in contact with the cell. When the tip of the cantilever is in contact with the cell, the initial cantilever deflection (d_0), and initial cantilever movement in (z) direction (z_0) are stored. As the cantilever moves further in z -direction and deforms the cell, the final cantilever deflection (d_1) and the cantilever movement (z_1) are obtained [21].

The stiffness depends not only on the Young's modulus but also on the geometry of the tip-surface contact. Therefore, the geometry and spring constant of the cantilever are calibrated in the same way for live and fixed cells. In the following experiments, we employed two types of cantilevers attached with a spherical probe ($5 \mu\text{m}$ in diameter).

- 1) Silicon nitride cantilever with a spring constant of 0.06 N/m (Novascan Technologies, Inc., Ames, IA) for live cells.
- 2) Silicon cantilever with a spring constant of 1.75 N/m (Novascan Technologies, Inc.) for fixed cells.

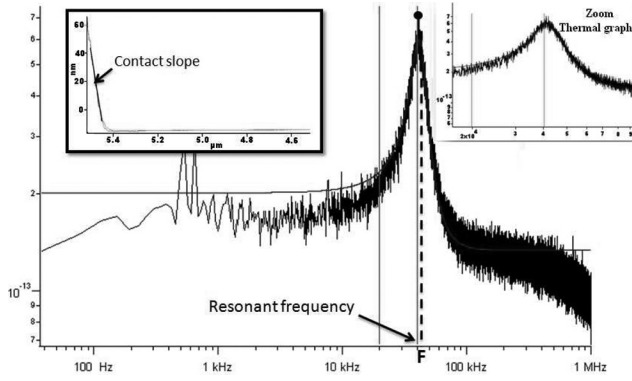


Fig. 8. Thermal graph measurement of a spherical tip cantilever ($K_c = 1.29$ N/m) used for fixed cells indentation. The left corner inset represents the contact slope from a force curve to determine the sensitivity of the cantilever and (deflection volt = 98.01 nm/v) and the right corner inset represents the zoom-in view of the resonant frequency plot ($F = 46.043$ kHz).

321 B. AFM Calibration

322 Before any experiment is done, it is necessary to perform
323 a calibration of the cantilever to determine its resonant frequency.
324 The resonant frequency measurement enables the user
325 to determine if it lies within the specified range provided by
326 the manufacturer in order to detect probe defects. The determination
327 of spring constant K_n by the thermal method involves
328 both measurements : 1) sensitivity in nanometer per volt and 2)
329 first-resonant frequency of the cantilever.

- 330 1) Determine the contact slope of the indentation curve on a
331 hard surface to determine the sensitivity of the cantilever
332 (see left corner inset of Fig. 8).
- 333 2) Measure the thermal power spectral density (PSD) to determine
334 the resonant frequency of the cantilever and to confirm that the
335 cantilever and light source are aligned. This confirms that the system
336 is functioning properly. The results are plotted as the deflection
337 amplitude data using Fourier transform in meter per square root hertz
338 versus the frequency in hertz (see right corner inset of Fig. 8).
339 The spring constant of the cantilever was determined experimentally
340 for each tip used in our studies using the IGOR software interface
341 supplied by Asylum Research.

342 Finally, the exact K_c values after the calibration procedure
343 are $K_c = 0.08745$ N/m for live cells and $K_c = 1.29$ N/m for
344 fixed cells.
345

346 V. MESC CHARACTERIZATION

347 A. Cell Culture Preparation

348 It is recommended to use specially coated dishes or slides
349 to facilitate adhesion. Such petri dishes or slides are available
350 commercially. Cells should be attached to some rigid substrate,
351 usually either a slide or the bottom of a petri dish. In our study,
352 the mESC R1 (SCRC-1011, American Type Culture Collection
353 (ATCC), Manassas, VA) were grown on 0.1% gelatin-coated
354 plates in the absence of feeder cells. The ES medium consisted
355 of 1000 U/mL leukemia inhibitory factor (LIF, ESGRO, Chemicon,
356 Temecula, CA), 15% fetal bovine serum (FBS) (Invitro-

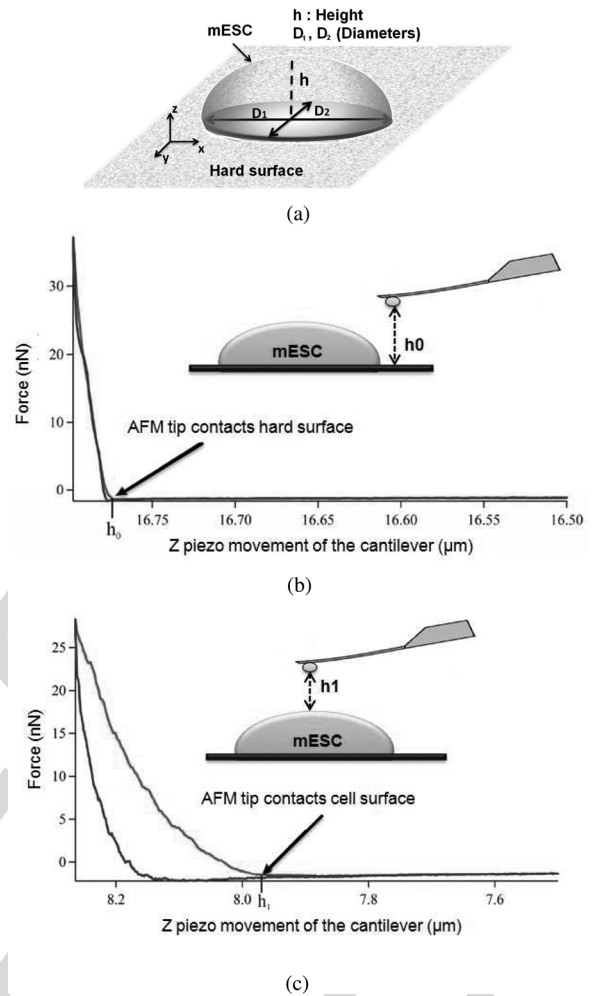


Fig. 9. How to determine the height of the stem cell through indentation curves? (a) mESC geometric parameters. (b) Force (F) versus piezo movement of the cantilever (Z) on a hard surface. The red and blue lines represent the loading and unloading curve, respectively. (c) Force (F) versus piezo movement of the cantilever (Z) on an mESC surface.

357 gen), and basic medium that included Knockout Dulbecco's
358 modified Eagle's medium (Invitrogen), 2 mM L-glutamine, 1x
359 nonessential amino acids, and 0.1 mM mercaptoethanol. Differentiation
360 was induced by removal of LIF from the medium.
361 Prior to the experiments, cells were dispersed using trypsin to
362 obtain single cells and were plated on 60 mm tissue culture petri
363 dish. Fixed mESC were obtained by treating live cells with with
364 4% formaldehyde for 10 min, washed and stored in phosphate
365 buffered saline.
366

367 B. Geometry of the Cell and Cantilever Interaction 368 With the Surface

369 The cells were determined to be 10 and 15 μm in diameter
370 using the cantilever tip to measure the interaction force between
371 the cell and the tip. The geometric modeling of the cell requires
372 the determination of the height, adhesion surface, and cell's contour
373 [see Fig. 9(a)]. The mESC height was calculated from the force-
374 displacement curve obtained on hard and mESC surfaces. First, the
375 initial height h_0 is measured by bringing

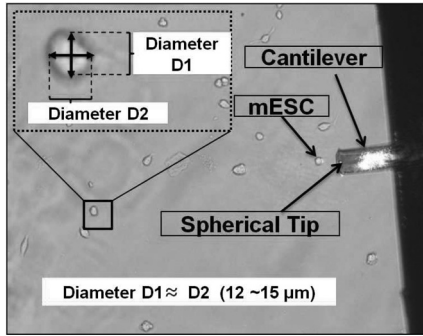


Fig. 10. How to determine the diameter of the stem cell through image analysis (ImagePro software from Asylum)?

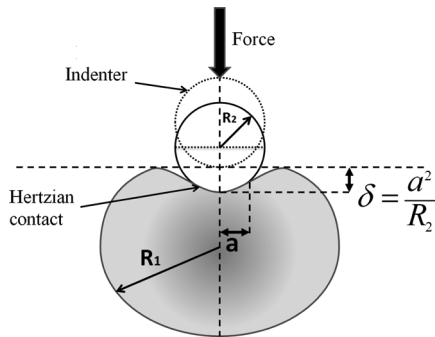


Fig. 11. Hertzian contact with a spherical tip: F is the loading force, and δ and a are the indentation and contact radius, respectively.

375 the AFM tip in contact with the hard surface near mESC [see
 376 Fig. 9(b)], and then, the second height h_1 by bringing the AFM
 377 tip on mESC surface [see Fig. 9(c)]. The mESC height is the
 378 difference $h = h_0 - h_1$. Second, to determine the mESC shape,
 379 AFM imaging experiments are carried out in order to deter-
 380 mine the mESC's contour and diameter using image processing
 381 techniques [19]. From Fig. 10, we selected mESC shapes with
 382 circular adhesion surfaces. As illustration, the left corner inset
 383 shows a typical circular mESC, where D_1 and D_2 stands for the
 384 horizontal cells diameters ($D_1 \simeq D_2$ close to $12 \mu\text{m}$).

385 C. Mechanical Characterization of mESC Using Spherical Tip

386 To estimate the mechanical properties of biological cells using
 387 the AFM, various analytical models can be used to identify the
 388 Young's modulus of mESC in live as well as fixed cells.

389 1) *Hertz Contact*: The Hertz contact model has been used
 390 extensively by the AFM community to quantify the mechanical
 391 property of biological samples using AFM [8], [9]. The Hertz
 392 contact model describes the simple case of elastic deformation
 393 of two perfectly homogeneous smooth surfaces touching under
 394 load (see Fig. 11). The forces measured are dominated by the
 395 elastic properties. In our work, the geometry of the tip used for
 396 all experiments is spherical. The mechanical interaction between
 397 the spherical tip and mESC can be described by the Hertz contact
 398 model of two elastic bodies [20]. The model assumes that:

399 1) the material properties of the tip and the cell are *isotropic*
 400 *and homogeneous*;

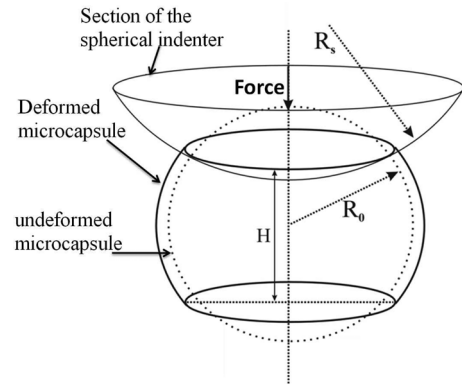


Fig. 12. Schematic of the microcapsule indented by a sphere with a loading force F .

- 2) the normal contact of the two bodies is *adhesionless and frictionless*;
- 3) the contact geometry is assumed to be *axisymmetric, smooth, and continuous*.

The relationship between the indentation δ , and the loading force F is given by

$$F = \frac{4}{3} E^* R^{\frac{1}{2}} \delta^{\frac{3}{2}} \quad (13)$$

where E^* and R are the combined modulus and the relative curvature of the tip and the cell, respectively,

$$\frac{1}{E^*} = \frac{1 - \nu_1^2}{E_1} + \frac{1 - \nu_2^2}{E_2} \quad (14)$$

$$\frac{1}{R} = \frac{1}{R_1} + \frac{1}{R_2} \quad (15)$$

where (E_1, ν_1, R_1) and (E_2, ν_2, R_2) represents the elastic modulus, Poisson's ratio, and the radius of the cell and spherical indenter, respectively. The elastic modulus of the silicon nitride cantilever and silicon cantilever were 222.22 and 168.17 GPa used for the live and fixed cells, respectively. The elastic modulus of the cells is in range of kilopascals [11]. Hence, our assumption that the tip used for probing is infinitely stiff ($E_2 \gg E_1$) compared to the cell and Hertz contact model is valid. Thus, (13) can be rewritten as follows:

$$F = \frac{4E_1}{3(1 - \nu_1^2)} R^{\frac{1}{2}} \delta^{\frac{3}{2}} \quad (16)$$

where E_1 and ν_1 represent the elastic modulus and Poisson's ratio of the cell assuming $\nu_1 = 0.5$.

2) *Capsule Contact*: The second model considers the biological cell to be composed of a cell membrane and cytoplasm (see Fig. 12). It is the capsule model. The model assumes that the cell membrane is a thin film and that the inner cytoplasm provides a uniform hydrostatic pressure on the membrane [22]. The model assumes the following.

- 1) The cell membrane is linearly elastic.
- 2) The deformation of the cell membrane is caused by stretching and bending of the cell membrane.
- 3) The cell is free of initial membrane stress or residual stress.
- 4) The cell volume is constant.

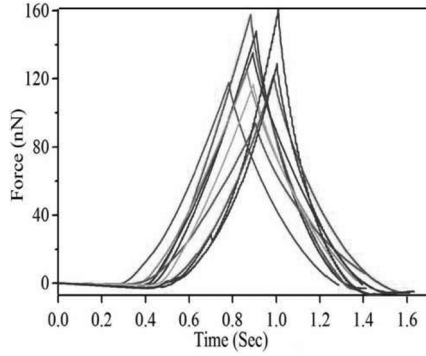


Fig. 13. Load-unload versus time for fixed undiff.

431 The force (F) and relative deformation (ε) relationship is
 432 given by [22]

$$F = \frac{2\pi h R_0 [1 + R_0 / (2R_s)]^2 E}{(1 - \nu)(1 + R_0 / R_s)^4} \varepsilon^3 + \frac{\pi}{2\sqrt{2}} E h^2 \sqrt{\varepsilon} \quad (17)$$

433 where E , ν , and h represent the elastic modulus, Poisson's rati-
 434 on, and the thickness of the capsule membrane, respectively.
 435 R_0 and R_s are the radius of the capsule and the spherical indenter,
 436 respectively. ε is the relative deformation of the capsule
 437 related to the indentation by $\varepsilon = \delta / 2R_0$. The first and the second
 438 term in (17) represent stretching and bending of the cell mem-
 439 brane, respectively, however, in our experiments, the bending
 440 deformation term can be neglected for $\varepsilon > 0.15$ [22].

441 D. Comparison Between Experimental Data 442 and Analytical Models

443 We evaluated and compared the different analytical mod-
 444 els to determine whether they appropriately predict the force-
 445 indentation relationship of mESC. Figs. 13 and 16 present the
 446 force (loading and unloading) versus time of fixed and live
 447 undiff mESC. From all these experiments, we observed adhesion
 448 force in few of the cases. For live mESC probed by a spherical
 449 tip, the adhesion force existed only for one sample. However,
 450 their amplitudes were negligible (less than 0.05 times the maxi-
 451 mum indentation) compared to the maximum indentation forces
 452 (more than 0.1 times the maximum indentation) observed by
 453 other researchers [28], [29]. In our experiments, we observed
 454 that the adhesion force exists for only one of the live undiff and
 455 does not exist for the fixed cells. This could be due to the wear
 456 of the spherical tip [26]. Since we did not observe any adhesion
 457 forces, we have not used the JKR or DMT models.

458 Figs. 14 and 15 show the experimental force versus cell indenta-
 459 tion curves for live and fixed undiff mESC and the simulated
 460 ones using the capsule and the Hertz models, respectively, for
 461 small and rather larger indentation range. Based on the ex-
 462 perimental investigation, a least-square fit on the dataset (for
 463 the whole cells) and the corresponding r^2 values are calcu-
 464 lated. These results show that the capsule model is inappropriate
 465 for describing mESC indentation. This conclusion is confirmed
 466 by the calculation of r^2 values performed on each mESC (see
 467 Table I). Finally, from our experiments, we infer that Hertz
 468 model appropriately describes the mechanical behavior of the

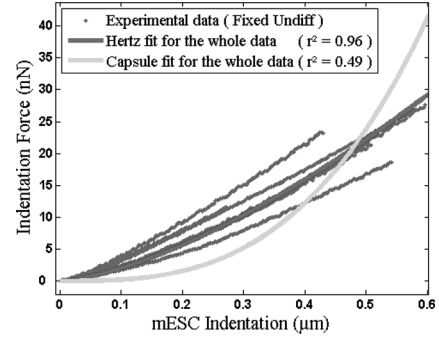


Fig. 14. Fixed undiff (small deformation).

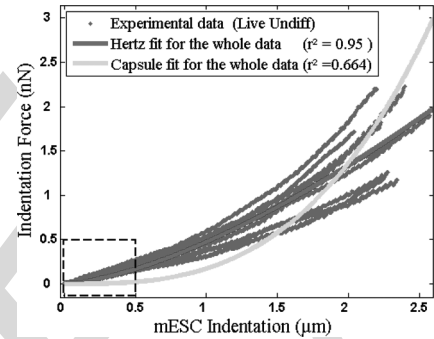


Fig. 15. Live undiff mESC.

TABLE I
 LEAST-SQUARE VALUE (COEFFICIENT CORRELATION) r^2 OBTAINED BY HERTZ
 FIT AND CAPSULE FIT FOR TEN SAMPLES OF LIVE AND FIXED UNDIFF mESC

	Live Undiff mESC		Fixed Undiff mESC	
	Hertz model	Capsule model	Hertz model	Capsule model
Correlation r^2	[0.98...0.99]	[0.68...0.90]	[0.95...0.99]	[0.46...0.73]

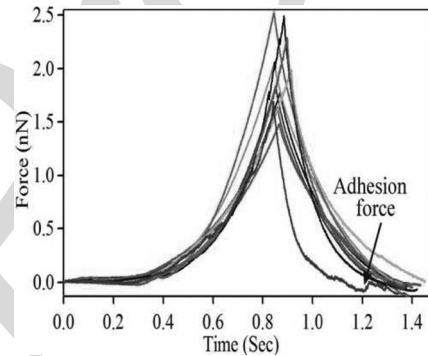


Fig. 16. Load-unload versus time for live undiff.

live and fixed undiff mESC (see Figs. 16– 18). Equation (16) was used to determine the global elastic modulus of the mESC. The average elastic modulus were 17.87 and 0.217 kPa for the fixed and live undiff, respectively, (see Fig. 19). The standard deviations for the fixed and live undiff cells were 3.37 and 0.05 kPa.

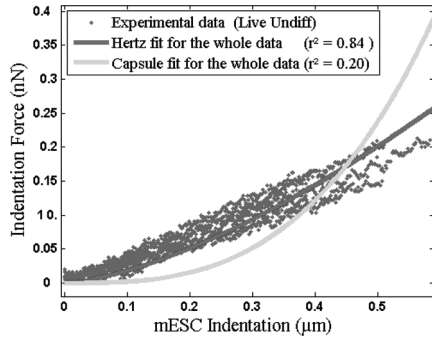


Fig. 17. Live undiff (small deformation).

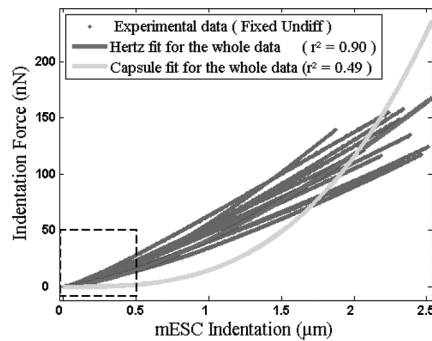
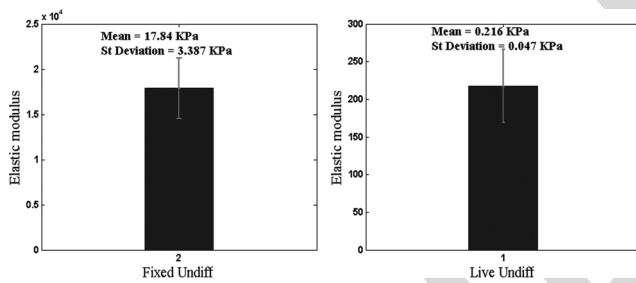


Fig. 18. Fixed undiff mESC.

Fig. 19. Elastic modulus for live and fixed undifferentiated mESC calculated from the Hertz contact model for a small deformation (range $0.5 \mu\text{m}$) with error bar.

475 It should be noticed that alterations in the mechanical properties of cells directly reflect changes in their cellular composition, 476 internal structure (cytoskeleton), and external interactions (cell– 477 cell and/or cell–surface) that occur during differentiation, aging, 478 and other changes in physiological status that are not accounted 479 in the cell indentation model. 480

481 VI. FEM SIMULATION OF CELL INDENTATION

482 Fig. 20 presents a photo of the visual and haptic user interface 483 setup. To test the accuracy and reliability of the proposed 484 user interface system with haptics-enabled simulation, a set of 485 experiments is designed using the setup given in Section II. 486 The geometrical dimensions of the cell are determined through 487 AFM image processing and force curve deflection (diameter and 488 height), and the mechanical properties through AFM indentation 489 studies due to the use of an analytical Hertz contact model.



Fig. 20. Real-time indentation simulator using the Omni haptic interface.

TABLE II
MECHANICAL AND GEOMETRICAL PROPERTIES OF SINGLE CELLS SELECTED
USED IN THE HAPTICS-ENABLED SIMULATOR

	Live mESC	Fixed mESC
Young modulus (KPa)	0.169	26.7
Poisson coefficient	0.49	0.49
Diameter (μm)	≈ 12	≈ 12
Height (μm)	8	8

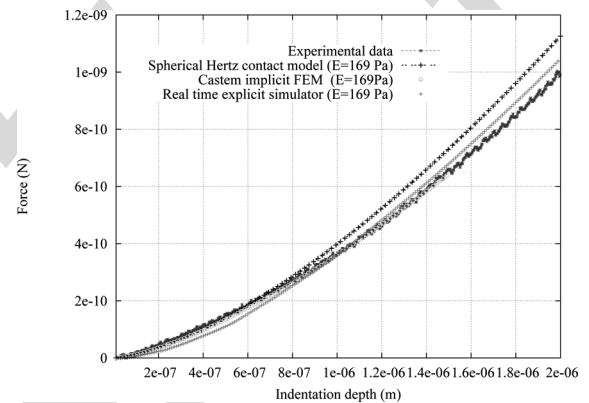


Fig. 21. Force versus indentation for live mESC: comparison between the real-time FEM, the Hertz contact, the quasi-static FEM, and experimental data.

490 These properties (see Table II) are passed to the finite-element 491 simulations.

492 In order to demonstrate the validity of our finite-element 493 indentation model in real time, we compare the simulation 494 indentation using the physics-based FEM model to the incremental 495 data provided by AFM indentation tool and to the response 496 simulated with a commercial FE software. Figs. 21 and 22 present 497 the nonlinear relationship between the penetration distance (in- 498 dentation) and the reaction force for live and fixed undiff cells. 499 The finite-element simulations using the commercial software 500 show good agreement with both Hertz contact model and the 501 experimental indentation data. The developed real-time physics- 502 based FEM simulator demonstrates its accuracy in predicting

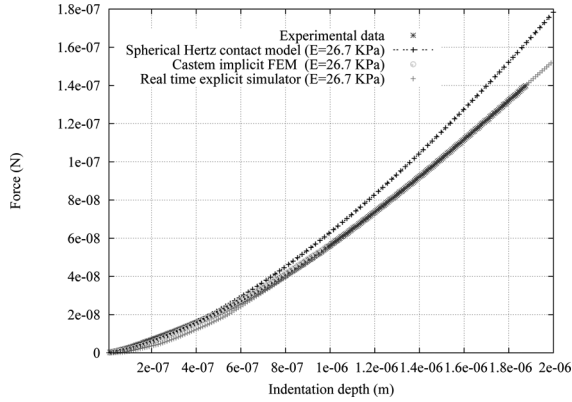


Fig. 22. Force versus indentation for fixed mESC: comparison between the real-time FEM, the Hertz contact, the quasi-static FEM, and experimental data.

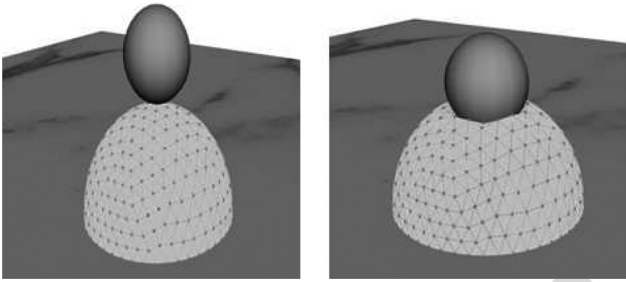


Fig. 23. Simulation of the AFM nanoindentation in VEs. The AFM tip and the stem cell was modeled using finite-element methods. A scene is shown from current training method: the 3-D cell deformation simulation before and after force indentation.

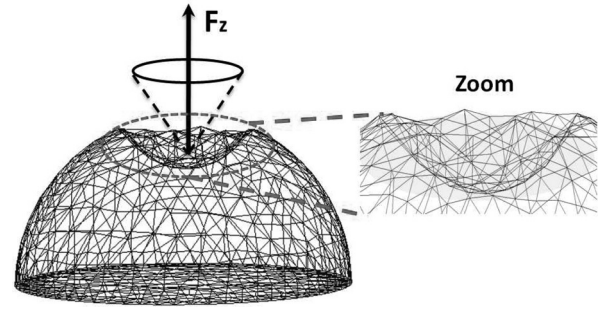


Fig. 24. Depth deformation of the cell: the left picture presents the depth indentation ($2 \mu\text{m}$) of the cell and the right picture shows the zoom of the cell deformation.

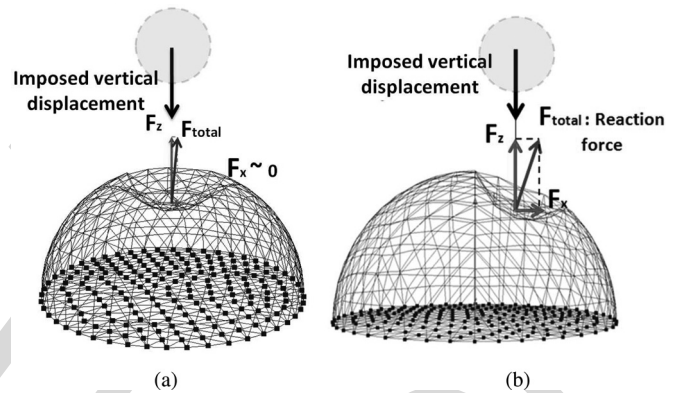


Fig. 25. Force indentation, where the probe is put (a) in the center and (b) off center of the cell.

503 indentation forces for large deformations ($>25\%$) of relative
 504 nominal strain (see Figs. 21 and 22). Fig. 23 illustrates a real-time
 505 cell indentation procedure in a VE for training. It shows a
 506 screenshot of the cell surface deformation. The interaction between
 507 tip and surface can be visualized to the user in 3-D (sphere
 508 and cone represent AFM tip and forces acting on tip, respectively).
 509 It shows the indentation (elastic deformation) on the
 510 cell surface and repulsive contact forces while pressing on the
 511 surface. In order to improve the force feedback rendering, we
 512 added a visual metaphor on the computer screen.

513 The rendered reflection force is depicted by a cone representing
 514 the force vector acting on the AFM tip. The visual metaphor
 515 has efficiently proved its effectiveness in AFM-based nanomanipulation
 516 simulation [23]. The vector amplitude reflects the force feedback
 517 amplitude and the vector orientation represents the vertical/horizontal
 518 force components. Fig. 24 shows in 3-D real time the probe tip,
 519 zoomed cell indentation, and micro-forces. The indentation force
 520 rendered to the operator is composed of a vertical F_z and horizontal
 521 F_x component. For the specific case corresponding to the center
 522 indentation, the tangential component is very small $F_x \approx 0$. Only
 523 vertical reflection force is provided to the user. He cannot haptically
 524 feel it during simulated indentation tests (see Figs. 25(a) and 26).
 525 In the case, where the AFM probe is pushed away from the center
 526 of the cell, a tangential component exists ($F_x \neq 0$) and the
 527 reflection

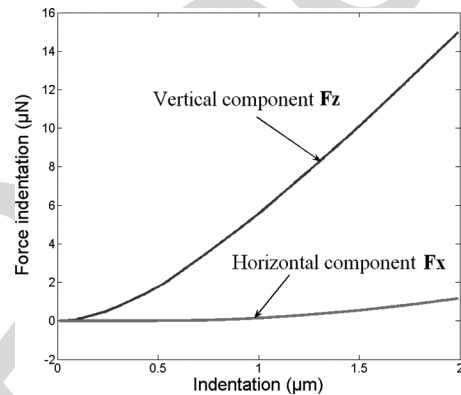


Fig. 26. Horizontal and vertical component force versus indentation: the probe is put in center of the cell.

528 tion force rendered to the user contains a vertical and tangential
 529 component (see Figs. 25(b) and 27). We did not reproduce this
 530 specific case in our results since the main interest of the paper is
 531 to prove experimentally the real-time explicit simulator results
 532 during centered cell indentations (since only vertical components
 533 of the force experiments are measured through the AFM
 534 cantilever).

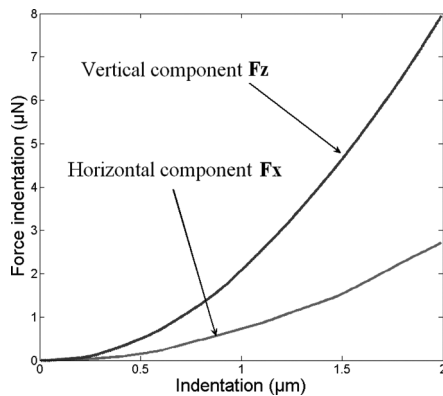


Fig. 27. Horizontal and vertical component force versus indentation: the probe is put off center of the cell.

VII. CONCLUSION

We have developed a computer-based training system to simulate real-time cell indentation procedures in VEs for training biologists. The simulator provides the user with visual and haptic feedback. The simulation of this procedure involves real-time rendering of computer generated graphical images of the AFM tip, physics-based modeling of soft biological tissue, and display of touch and force sensations to the trainee through the simulation of haptic interactions. The real-time haptics-enabled simulator is realistic, since it is mainly based on experimental AFM data (mechanical and geometrical property of mESC) to accurately replicate the indentation task and predict the cell deformation during indentation. We first investigated the challenging issues in the real-time modeling of the biomechanical properties of the cell indentation through FEMs. Compared to experimental AFM indentation results performed on mESC, we can see clearly that the proposed physically based FEM model is able to simulate the cell deformation through real-time simulation constraints. Currently, we are working on integrating nanoscale effects, such as friction, viscosity, puncture, and adhesion forces in the nonlinear FEM model. All modalities will be merged in an ergonomic tool and intelligent biological simulator for stem cell characterization.

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