



**HAL**  
open science

# Spontaneous Firings of Carnivorous Aquatic Utricularia Traps: Temporal Patterns and Mechanical Oscillations

Olivier Vincent, Ivan Roditchev, Philippe Marmottant

► **To cite this version:**

Olivier Vincent, Ivan Roditchev, Philippe Marmottant. Spontaneous Firings of Carnivorous Aquatic Utricularia Traps: Temporal Patterns and Mechanical Oscillations. PLoS ONE, 2011, 6 (5), pp.e20205. 10.1371/journal.pone.0020205 . hal-00642061

**HAL Id: hal-00642061**

**<https://hal.science/hal-00642061>**

Submitted on 17 Nov 2011

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# Spontaneous Firings of Carnivorous Aquatic *Utricularia* Traps: Temporal Patterns and Mechanical Oscillations

Olivier Vincent, Ivan Roditchev, Philippe Marmottant\*

CNRS/Université Grenoble 1, UMR 5588, Laboratoire Interdisciplinaire de Physique, Grenoble, France

## Abstract

Aquatic species of *Utricularia* are carnivorous plants living in environments poor in nutrients. Their trapping mechanism has fascinated generations of scientists and is still debated today. It was reported recently that *Utricularia* traps can fire spontaneously. We show here that these spontaneous firings follow an unexpected diversity of temporal patterns, from “metronomic” traps which fire at fixed time intervals to “random” patterns, displaying more scattered firing times. Some “bursting” traps even combine both aspects, with groups of fast regular firings separated by a variable amount of time. We propose a physical model to understand these very particular behaviors, showing that a trap of *Utricularia* accomplishes mechanical oscillations, based on continuous pumping and sudden opening of the trap door (buckling). We isolate the key parameters governing these oscillations and discuss the effect of their fluctuations.

**Citation:** Vincent O, Roditchev I, Marmottant P (2011) Spontaneous Firings of Carnivorous Aquatic *Utricularia* Traps: Temporal Patterns and Mechanical Oscillations. PLoS ONE 6(5): e20205. doi:10.1371/journal.pone.0020205

**Editor:** Tom Waigh, University of Manchester, United Kingdom

**Received:** February 11, 2011; **Accepted:** April 12, 2011; **Published:** May 27, 2011

**Copyright:** © 2011 Vincent et al. This is an open-access article distributed under the terms of the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original author and source are credited.

**Funding:** This work was funded by the French Ministry of Higher Education and Research (<http://www.enseignementsup-recherche.gouv.fr/>). The funders had no role in study design, data collection and analysis, decision to publish, or preparation of the manuscript.

**Competing Interests:** The authors have declared that no competing interests exist.

\* E-mail: philippe.marmottant@ujf-grenoble.fr

## Introduction

Aquatic species from the genus *Utricularia* are widespread carnivorous plants, catching their preys with millimeter-sized traps. Since the discovery of their carnivorous character [1,2], there has been much interest in the mechanism underlying their extremely fast motion: the entrance of a trap is closed by a door which is capable of opening and closing at the time scale of 1 ms only [3]. It is known that slow pumping of water out of the trap enables storage of elastic energy in the trap walls, which is suddenly released when the trap is triggered by a slight touch on one of its four trigger hairs [4,5]. However there is still debate on the mechanism at the origin of the door opening [6,7]. Recent work has focused on time-resolved analysis of the door dynamics at small time scales, bringing to light the mechanical role of the door as a buckling valve [3], and long time analysis, showing that a single trap is able to fire spontaneously many times without any external action [3,8]. In order to understand this surprising behavior and how it is connected to the trapping mechanism, we studied spontaneous firings of *Utricularia inflata* and *Utricularia australis*, recording the times of the firings and the temporal evolution of the trap shape. The aim of this paper is to present the original behavior of the recorded traps which proved to be much more complex than previously thought, and to show how these behaviors can be described by a simple physical model combining the deterministic mechanics of the elastic door and statistical fluctuations.

## Results

### Time repartition of spontaneous firings

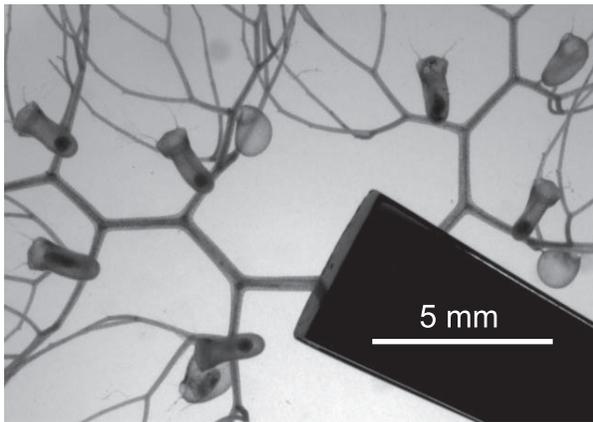
The plants were immersed in unstirred de-ionized water to avoid the presence of animals or fluid motion capable of triggering

traps (see Figure 1 and Video S1). All observed traps showed spontaneous firings, with a maximum of about 200 firings for a single trap in three weeks. Some typical examples of their temporal behavior are presented on Figure 2. On Figure 2A, a vertical bar is plotted each time  $T$  a firing occurs, for three different traps. Denoting  $n$  the firing number, we define the time interval between consecutive firings  $t_f$  as  $(t_f)_n = T_{n+1} - T_n$ . On Figure 2B,  $t_f$  is plotted as a function of  $n$ . Both panels of Figure 2B show that different traps on a same composed leaf of *Utricularia inflata* can have very different behaviors. First, firings in trap A are spaced and scattered in time. This behavior will be called “random” in the following. On the contrary, “metronomic” traps such as trap C show very regular firings occurring at well-defined time intervals. The limit between these behaviors is sometimes thin, as shown by trap B: events are not well organized in time such as in trap C, but the time interval between firings is not as widespread as in trap A. This suggests that more than two distinct behaviors, “random” and “metronomic” traps are two extreme cases of a continuous range of behaviors.

We noticed that “metronomic” traps often show a slow drift of their period  $t_f$ , which is for example doubled after about 40 firings for trap C. This fact prevents to use the standard deviation of  $t_f$  as an indicator of the behavior of a trap: much of the calculated standard deviation for trap C would indeed come from the regular drift of its period. To limit this bias, we define a “randomness index” as

$$r = \frac{\sqrt{\langle \Delta t_f^2 \rangle}}{\langle t_f \rangle} \quad (1)$$

where  $\Delta t_f$  is the variation of time intervals for successive firings:

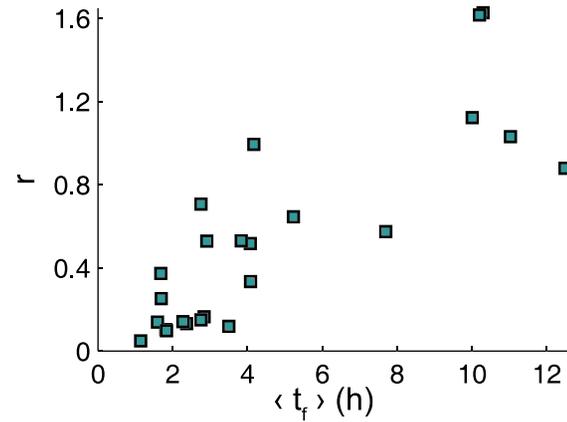


**Figure 1. General view of an excised composed leaf of *Utricularia inflata*.** The plant is held by tweezers (in black) and immersed in de-ionized water. See also Video S1. doi:10.1371/journal.pone.0020205.g001

$(\Delta t_f)_n = (t_f)_{n+1} - (t_f)_n$  and  $\langle \rangle$  represents an average over all successive firings of the considered sample. Values of  $r$  for traps A, B, C are shown on the right panel of Figure 2, showing that the visual feeling of randomness is well reproduced by the value of  $r$ , which is less than 0.1 for very “metronomic” traps and of the order of 1 for very random traps.

Noticeably, the last presented firings of trap C become more scattered as  $t_f$  increases, as was also observed on other “metronomic” traps. This suggests a link between the irregularity of a trap, characterized by  $r$ , and its period  $t_f$ . To check this hypothesis, we calculated  $r$  and the mean value of firing intervals  $\langle t_f \rangle$  for several samples from 21 different traps of *Utricularia inflata*. Results are presented on Figure 3 and confirm the tendency of higher irregularity for higher firing periods. “Metronomic” traps ( $r < 0.2 - 0.3$ ) typically fire between 1 h and 3 h on average, while “random” traps ( $r > 0.5$ ) display values of  $t_f$  usually bigger than 5 h.

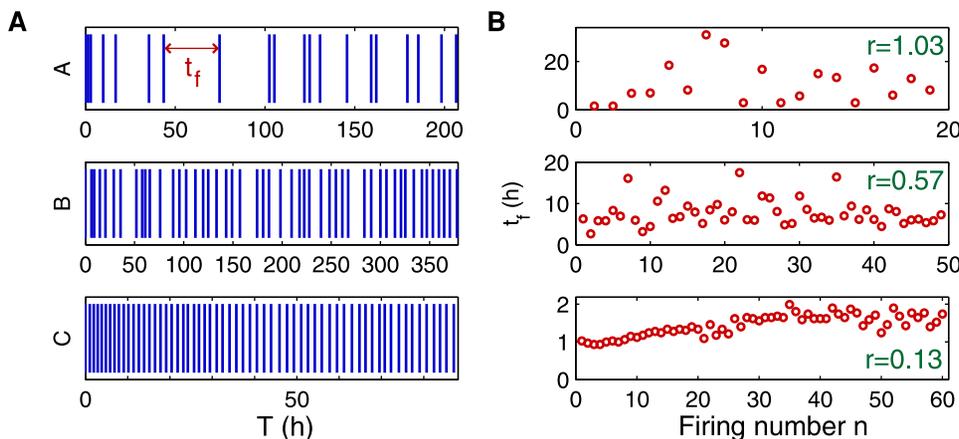
Looking closer at some apparently very fluctuating traps, we found that some of them displayed a surprising grouping effect, where firings often happen by groups of 2 up to 7 very close and



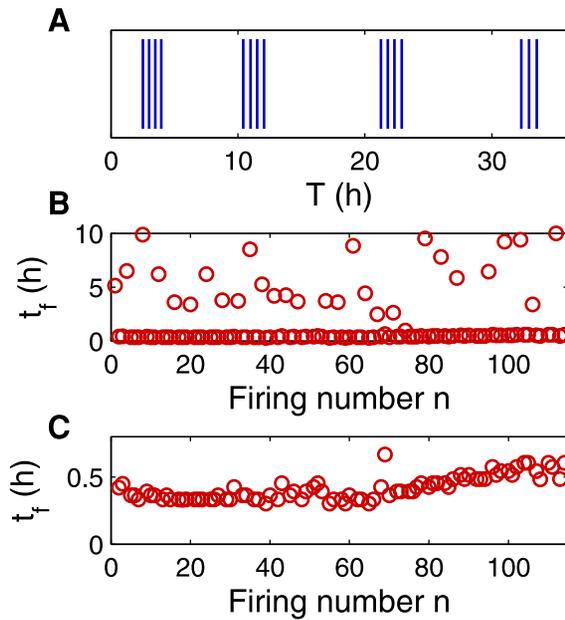
**Figure 3. Randomness index  $r$  versus mean firing time interval  $\langle t_f \rangle$ .** Values of  $r$  and  $\langle t_f \rangle$  are calculated for 24 samples containing between 10 and 107 firings from 20 different traps of *Utricularia inflata*. doi:10.1371/journal.pone.0020205.g003

regular events, separated by a variable amount of time. In our experiment with *Utricularia inflata*, five traps presented these “bursts”, but the most striking example was given by a trap of *Utricularia australis* (trap D) followed ten consecutive days (see Figure 4A and Video S2). The time intervals between firings  $t_f$  measured on this experiment are shown on Figures 4B and 4C, showing that time intervals between consecutive events inside a burst follow a very regular line as for “metronomic” traps, with a randomness index  $r$  close to 0.1. On the contrary, the time separation between bursts is very scattered as for “random” traps. Note that the number of firings per burst is roughly constant, 3 or 4 in this case, and that  $t_f$  inside a burst is very small, of order 30 min. Interestingly, even for *Utricularia inflata* traps, this latter time is small, typically between 15 min and 1 h. This suggests that bursts might be a characteristic behavior for high-frequency (low  $t_f$ ) traps.

During the three weeks of observation, traps have not shown significant changes of behavior. Slow transitions from “metronomic” to more “random” periods often happen, correlated with an increase in the time between firings  $t_f$ . Also, a few traps stopped



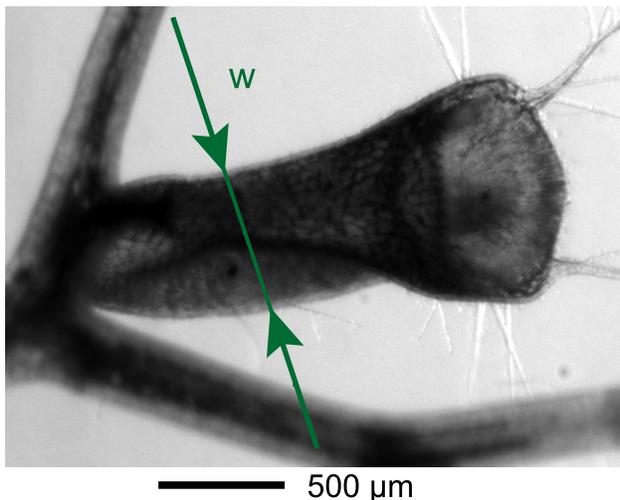
**Figure 2. Extract of the firing events of 3 different traps of *Utricularia inflata*.** Time (T) is set to 0 at the beginning of each sample. **A:** a vertical bar is drawn each time a firing occurs. **B:** corresponding time intervals  $t_f$  between successive firings. The value of the randomness index  $r$  associated to each sample is indicated. For trap C,  $r = 0.07$  for the 20 first firings. doi:10.1371/journal.pone.0020205.g002



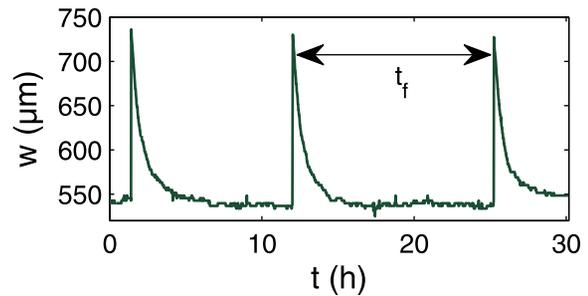
**Figure 4. “Bursting” behavior in trap D (*Utricularia australis*).** **A:** excerpt of the firing events. A vertical bar is drawn each time a firing occurs (time is set to 0 at the beginning of the sample). Firings occur by bursts of 3 or 4 events. **B:** time intervals  $t_f$  between successive firings for all the recorded firings. The scattered points correspond to the times between consecutive bursts, while the regular line at the bottom is drawn the very regular firings inside a burst. **C:** Magnification of the bottom line of panel B. The associated randomness index is  $r=0.11$ . doi:10.1371/journal.pone.0020205.g004

firing after a few days, sometimes temporarily, sometimes definitely. We will refer to these latter traps as “waiting” in the following.

In conclusion, our experiments exhibit a rich variety of behaviors, the most surprising ones being “metronomic” spontaneous firings, following precise temporal patterns, and “bursting” ones which combine regularity and randomness at different time scales. We suggested above that the mean time interval between



**Figure 5. View of trap B (*Utricularia inflata*).** The figure shows the definition of the lateral width  $w$  used for the data presented on Figure 6. doi:10.1371/journal.pone.0020205.g005



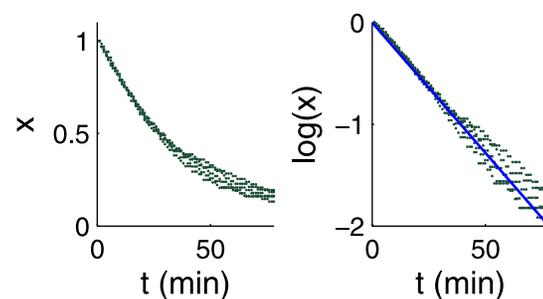
**Figure 6. Extract of the evolution in time of the lateral width  $w(t)$  of trap B.** Three successive spontaneous firings can be observed. doi:10.1371/journal.pone.0020205.g006

firings was an important parameter determining the behavior of a trap.

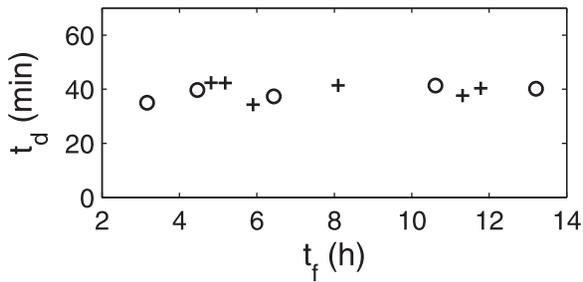
### Study of the change of width of the traps

In order to understand more in detail the origin of the behaviors described above we focused on the physical process of trap setting: due to active pumping of water, the volume of the trap decreases with time, thereby lowering the pressure inside the trap. We thus extracted when possible a measurement of the width of the traps as a function of time  $w(t)$ , obtained by image analysis on traps viewed from above (see Figure 5), so it represents a projected width, used as an indicator of the trap state. The curve obtained for trap B is shown on Figure 6. Each peak corresponds to a spontaneous firing, followed by a decrease of the trap width from its inflated state to its deflated state. As shown in a previous paper [9], this relaxation is exponential, with a characteristic time  $t_d$ . This is well verified in our experiments, as can be seen on Figure 7. We also checked that  $t_d$  did not vary much for a single trap, and that its variations were not related to those of the time intervals between firings (see Figure 8). This shows that the source of fluctuations in  $t_f$  has to be found elsewhere, and that  $t_d$  can be considered as a constant, characteristic deflation time for the considered trap.

Similar results were found on other traps of *Utricularia inflata* and *Utricularia australis*, and some values of  $t_d$  are reported in table 1. Since  $t_d$  is a natural unit of time for each trap, it is interesting to measure the firing intervals  $t_f$  in units of  $t_d$ : the ratio  $\tau = t_f/t_d$  indicates at which level of the deflation process firings occur. Our results for “metronomic” and “random” traps presented in table 1 show that  $\tau$  seems to be strongly linked to the behavior of a trap :



**Figure 7. Degree of inflation  $x(t)$  for 5 successive firings of trap B.**  $x$  is defined as  $x = \Delta w / \Delta w(t=0)$  where  $\Delta w$  is the difference between the current value of the trap lateral width  $w$  and its value at full deflation. Left: linear plot. Right: logarithmic plot. The blue line corresponds to a fit of  $x(t) = \exp(-t/t_d)$  with  $t_d = 39$  min. Time  $t$  is reset to 0 at each firing. doi:10.1371/journal.pone.0020205.g007



**Figure 8. Measured deflation time  $t_d$  for 12 different firings of trap B versus the corresponding firing time interval  $t_f$ .** Two series of 5 and 7 firings are shown (first series in circles, second series in crosses), and are separated by five days. Uncertainties on  $t_f$  are negligible compared to those on  $t_d$  ( $\Delta t_d \sim 3$  min). doi:10.1371/journal.pone.0020205.g008

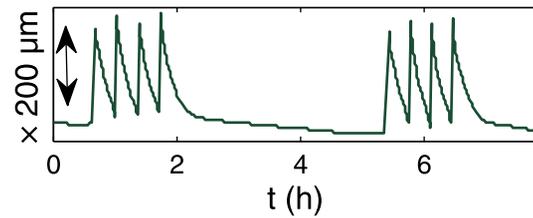
the higher the value of  $\tau$ , the higher the irregularity of the trap. Observations on other traps where  $\tau$  was not precisely measurable show this general trend: “metronomic” traps fire at an early stage of the deflation process and “random” traps usually fire close to their fully deflated state.

We now focus on “bursting” traps such as trap D. The evolution of trap thickness presented on Figure 9 displays the same characteristics: inside a burst, successive firings are fast and complete deflation is never achieved ( $\tau \sim 1$ ), visible in the fact that the slope of the curve of Figure 9 is always considerable, whereas the time interval between two bursts is long compared to  $t_d$  ( $\tau \gg 1$ ), as can be seen in the last two rows of table 1. We noticed that *Utricularia inflata* bursting traps also feature  $\tau \approx 1$ .

We thus suggest that instead of simply  $t_f$  (see previous section), the relevant parameter predicting the behavior of a trap is  $\tau = t_f/t_d$ : high values of  $\tau$  correspond to an irregular behavior, while low values are associated with regular traps. In the following discussion, we develop physical arguments supporting this hypothesis.

## Discussion

Our experiments show a very rich variety of behaviors in traps of *Utricularia inflata* and *Utricularia australis*. Environmental fluctuations such as day/night oscillations, temperature changes or light intensity variations cannot account for these observations, since all observed traps were on a same composed leaf under the same conditions. To explain our observations we have to understand how the trapping mechanism works. It has been put forward that the opening of the trap door of *Utricularia* was based on an elastic instability: buckling, which is a mechanical process where an elastic membrane, in our case the trap door, resisting a



**Figure 9. Extract of the evolution in time of the lateral width of trap D (*Utricularia australis*).** Two successive bursts comprising 4 firings each are shown. See also Video S2. doi:10.1371/journal.pone.0020205.g009

pressure difference  $\Delta p$  suddenly changes its curvature at a critical pressure difference  $\Delta p_c$  [3]. We will explain our experimental results in the light of these considerations, suggesting that the repetitive character of observed firings is a direct consequence of the spontaneous buckling of the trap door.

### Buckling cycles: “metronomic” and “waiting” traps

Our results confirm previous observations [3,8,10] that deflation starts immediately after firing. Since deflation originates from active pumping of water out of the trap [5], this indicates that pumping is a continuous process. Starting from a fully inflated trap, this continuous pumping entails a progressive deflation, represented by  $w(t)$ . As a consequence, the pressure inside the trap lowers to a pressure  $p_{\text{int}} < p_{\text{ext}}$ , where  $p_{\text{ext}}$  is the pressure in the surrounding liquid. This entails an increase of  $\Delta p = p_{\text{ext}} - p_{\text{int}}$ , representing the net pressure exerted on the trap door. If there is a well-defined pressure difference  $\Delta p_c$  at which the trap door buckles, then the door spontaneously opens when  $\Delta p$  reaches that value. The door being open, the trap inflates and the pressure difference is reset to zero. As the door closes, the same cycle of deflation - buckling can start again, *ad infinitum*. This picture shows how the observed “metronomic” oscillations of the trap width may arise from the combination between continuous pumping and spontaneous buckling of the trap door. Note that if  $\Delta p$  is never high enough to reach  $\Delta p_c$ , then the trap never fires spontaneously and stays in a “waiting” phase (see Figure 10). Experimentally, the hypothesis of door buckling is supported in our experiments by the fact that the level of deflation achieved when a firing happens does not vary much for successive firings of a single trap (see Figure 9 for example), meaning that firings probably occur at comparable  $\Delta p$ . We now derive a simple model to extract the physical parameters governing these oscillations.

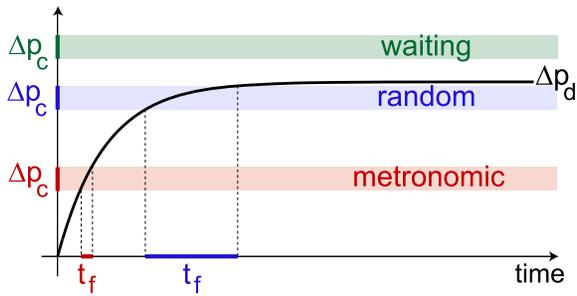
First, we have to understand the evolution of the pressure difference  $\Delta p$  in time. Our experiments only access to the trap width  $w$ , but due to the elasticity of the trap wall,  $w$  is directly

**Table 1. Deflation times and firing intervals.**

trap	A	B	C'	D	
type	random ( $r \approx 1$ )	random ( $r \approx 0.6$ )	metronomic ( $r \approx 0.1$ )	burst (intra)	burst (inter)
$t_d$ (h)	$0.87 \pm 0.10$	$0.65 \pm 0.05$	$0.47 \pm 0.07$	$0.42 \pm 0.08$	
$\langle t_f \rangle$ (h)	10.9	7.7	1.6	0.41	6.5
$\tau = \langle t_f \rangle / t_d$	$12.5 \pm 1.4$	$11.8 \pm 0.9$	$3.4 \pm 0.5$	$0.98 \pm 0.19$	$15.5 \pm 3.1$

Typical values of deflation times  $t_d$  and the mean value of firing intervals  $\langle t_f \rangle$ , calculated over successive firings of the considered trap. Due to the angle of observation,  $t_d$  was not precisely measurable on trap C so another “metronomic” trap was considered, denoted C'. For trap D, we distinguish the firing time inside a burst (intra) and between bursts (inter).

doi:10.1371/journal.pone.0020205.t001



**Figure 10. Model explaining the trap behaviors.** The black curve is the evolution of the pressure difference  $\Delta p$  due to the deflation process, saturating at a value  $\Delta p_d$ . Firing of the trap door occurs at a time  $t_f$  when  $\Delta p$  reaches the buckling pressure  $\Delta p_c$ . Fluctuations in  $\Delta p_c$  entail fluctuations in  $t_f$  which are bigger when  $\Delta p_c$  is close to  $\Delta p_d$ , explaining the scattered values of  $t_f$  for “random” traps. If  $\Delta p_c$  is bigger than  $\Delta p_d$ , buckling is impossible and the trap is in a “waiting” state. doi:10.1371/journal.pone.0020205.g010

linked to  $\Delta p$  with a relationship that we assume linear (see Methods). Consequently, since  $w$  decays exponentially with a time constant  $t_d$ , one should also have

$$\Delta p(t) = \Delta p_d (1 - \exp(-t/t_d)) \quad (2)$$

$\Delta p_d$  being the maximum pressure difference attainable by the trap, corresponding to a fully deflated state. The value of  $\Delta p_d$  has been estimated to 16 kPa [10].

Note that expression (2) can be predicted theoretically using simple hypotheses (see Methods). The model presented in the Methods section also justifies our experimental observation that  $t_d$  does not vary much in time for a single trap, since it does not depend on the pumping rate but mainly on the permeability and elasticity of the trap body, which can be considered as constant.

The spontaneous firing of the trap occurs at a time  $t_f$  where  $\Delta p(t_f) = \Delta p_c$ , which is possible only if  $\Delta p_d > \Delta p_c$ , i.e. when pumping is strong enough to make the door buckle. If this is the case, then from equation (2) we have

$$t_f = t_d \ln \left( \frac{1}{1 - \Delta p_c / \Delta p_d} \right). \quad (3)$$

As we demonstrated in our experiments,  $t_d$  is constant for a trap, so that  $t_f$  only depends on the ratio  $\Delta p_c / \Delta p_d$ . If  $\Delta p_c$  and  $\Delta p_d$  are constant,  $t_f$  is a constant too and firings occur at very regular intervals: this corresponds to the “metronomic” behavior. Notice that this metronomic characteristic doesn’t depend on the mathematical expression of  $\Delta p(t)$ , and is always true if the evolution of  $\Delta p$  in time is the same after each firing, which is justified by our observations that  $t_d$  is a trap constant, and if there exists a time  $t_f$  where  $\Delta p$  reaches the critical buckling pressure  $\Delta p_c$ . If this latter condition is not verified, the trap is not able to fire spontaneously and is in a “waiting” state.

### Fluctuations and “random” traps

In order to explain random firings with this model it is necessary to introduce fluctuations. At these scales, buckling is insensitive to thermal noise [9] and for an incompressible spherical shell with a thickness  $h$ , a radius  $R$  and a Young’s modulus  $E$ , the critical buckling pressure is given by [11]

$$\Delta p_c \sim E \left( \frac{h}{R} \right)^2 \quad (4)$$

so any change in elasticity, affecting  $E$ , or shape, affecting  $R$  but also probably the exact prefactor in equation (4), is able to impact the value of  $\Delta p_c$ .

Changes in shape can occur at each firing since the trap door doesn’t necessarily come back exactly at the same position when it closes. Changes in elasticity are also possible if there are variations of turgor pressure inside the door wall. It is clear from equation (3) that fluctuations in  $\Delta p_c$  directly impact the time interval between firings  $t_f$ . Figure 10 shows how fluctuations in  $\Delta p_c$  affect the distribution of  $t_f$ . In particular, fluctuations around a small value of  $\Delta p_c$  have a much weaker effect than the same fluctuations around a value of  $\Delta p_c$  close to  $\Delta p_d$ , due to the exponential evolution of  $\Delta p(t)$ . Thus “metronomic” traps should have a low value of  $\Delta p_c / \Delta p_d$  or equivalently a low value of  $\tau = t_f / t_d$  (see equation (3)), while “random” traps would have  $\Delta p_c / \Delta p_d$  closer to 1, leading to a higher value of  $\tau$ , which is well supported by our experimental results.

This model also predicts that if the mean firing period of a single trap increases, fluctuations of the firing times should also increase. Trap C provides a good illustration of that point on Figure 2B, bottom. It could also be an explanation of our observation that as time passes, “metronomic” traps often slow down their firings, leading them to become more “random”, temporarily or permanently.

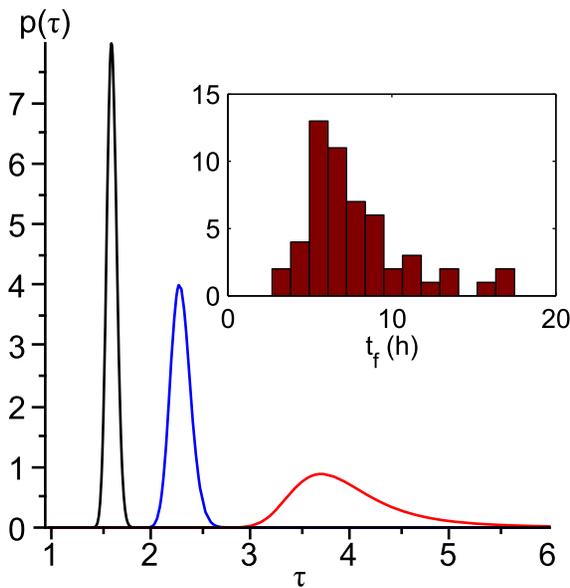
Notice that  $\Delta p_d$  could also fluctuate on the same order of magnitude than  $\Delta p_c$ , due to changes in the pumping rate for example. However, one cannot actually separate the effect of  $\Delta p_c$  and  $\Delta p_d$  as can be seen from equation (3), and the relevant parameter is in fact  $\Pi = \Delta p_c / \Delta p_d$ . Using the other natural parameter  $\tau = t_f / t_d$ , equation (3) reduces as

$$\tau = \ln \left( \frac{1}{1 - \Pi} \right) \quad (5)$$

and we now derive some properties of such a dependence between  $\tau$  and  $\Pi$ , keeping in mind that the detailed results depend on the exact expression of  $\Delta p(t)$ . However this simple expression helps us to illustrate the previous arguments. Moreover, the ideas developed below remain valid for any  $\Delta p(t)$  provided that  $d^2 \Delta p / dt^2 < 0$ . For example one can calculate how fluctuations propagate from  $\Pi$  to  $\tau$  (see Methods), and one can readily show, assuming  $\Pi < 1$ , that the mean values  $\langle \tau \rangle$  and  $\langle \Pi \rangle$  follow equation (5), and that the standard deviations  $\sigma_\tau$  and  $\sigma_\Pi$  of respectively  $\tau$  and  $\Pi$ , are related by

$$\sigma_\tau = \left( \frac{1}{1 - \langle \Pi \rangle} \right) \sigma_\Pi. \quad (6)$$

Due to the denominator  $1 - \langle \Pi \rangle$ , fluctuations get largely amplified as the mean value of the pressure ratio  $\langle \Pi \rangle$  gets closer to 1. Within the framework of this model, the wide range of time intervals between firings of some traps (the “random” ones) just reflects the amplification of pressure fluctuations which become very important when the buckling and pumping pressure  $\Delta p_c$  and  $\Delta p_d$  are comparable. We also show that in addition to the amplification of fluctuations, the shape of the probability distribution is modified (see Figure 11). In particular, symmetrical distributions on  $\Pi$  give distributions on  $\tau$  that expand towards large values of  $\tau$ , explaining the non-symmetrical aspect of the



**Figure 11. Probability distributions on  $\tau$ .** They are calculated assuming gaussian fluctuations on  $\Pi$  with standard deviation 0.01 and centered on  $\langle \Pi \rangle = 0.8$  (black),  $\langle \Pi \rangle = 0.9$  (blue),  $\langle \Pi \rangle = 0.98$  (red). The distribution on  $\tau$  becomes broader as  $\langle \Pi \rangle$  increases, but also less symmetrical, as observed for “random” traps: a histogram of the firing times  $t_f$  of trap B is presented in the inset of the figure for qualitative comparison.

doi:10.1371/journal.pone.0020205.g011

repartition of firing times  $t_f$  in random traps (see inset of Figure 11).

Fluctuations around  $\Pi = 1$  can also occur, making the trap oscillate between “waiting” and “random”, giving much more scattered events, which is probably the case for trap A which displays long waiting periods and a large value of  $\tau$ .

Interestingly, even “metronomic” traps have  $\tau$  bigger than 1, meaning that  $\Delta p_c / \Delta p_d = \Pi = 1 - \exp(-\tau)$  is never far from 1. Thus, all traps seem to have  $\Delta p_c$  comparable to  $\Delta p_d$ . The reason could be that a too low  $\Pi$  would have the trap firing very often but not achieving significant deflation: only a small amount of liquid would be sucked at each firing making the trap not efficient to catch preys. On the other hand, a high value of  $\Pi$  would mean that the deflation pressure  $\Delta p_d$  is small compared to the buckling pressure  $\Delta p_c$ , making the door wall very stable and the trap difficult to trigger. This completes the discussion presented in [3] showing how the elasticities and shapes of the trap wall and the door are optimized for efficient prey capture.

## Bursts

At first sight, bursts could also be interpreted as fluctuations of time intervals. However, the facts that time intervals in the bursts are very well defined and that the number of peaks in a burst is roughly constant rules this idea out. In a burst, the behavior of a trap strongly looks like a “metronomic” one. Looking at a larger time scale, bursts groups seem on the contrary to be randomly distributed. To account for this unexpected behavior, we suggest that after a trap is triggered, there should be a process relaxing with a time  $t_{ex}$ . For example,  $t_{ex}$  could be a typical relaxation time of the door rigidity *via* its turgor pressure. Thus, any firing of the trap would be associated with a reduction of the buckling threshold  $\Delta p_c$  for a time  $t_{ex}$ . When triggered or spontaneously fired, such a trap would go from “waiting” or “random” to a

“metronomic” state for a time  $t_{ex}$  due to the lowering of  $\Delta p_c$ , thus  $\Pi$ , then go back to its initial state, giving the observed bursts. This hypothesis is supported by the experimental fact that the number  $n_{burst}$  of firings per burst is usually constant over long periods of time: this number would be approximately given by  $n_{burst} \sim t_{ex} / t_f$ .

To check this, we analyzed 53 bursts on six different traps, and recorded  $n_{burst}$  and the mean value of firing intervals  $\langle t_f \rangle$  inside the considered burst. We then plotted  $n_{burst}$  versus  $1 / \langle t_f \rangle$  (see Figure 12), the slope of which should be approximately  $t_{ex}$ . We find a good agreement with this prediction for a wide range of  $n_{burst}$ , comprised between 2 and 7, and we deduce that  $t_{ex}$  should be of order 1–2 h. The scattering of the points on the plot probably comes from the variation of  $t_{ex}$  from trap to trap, but also for a single trap in time.

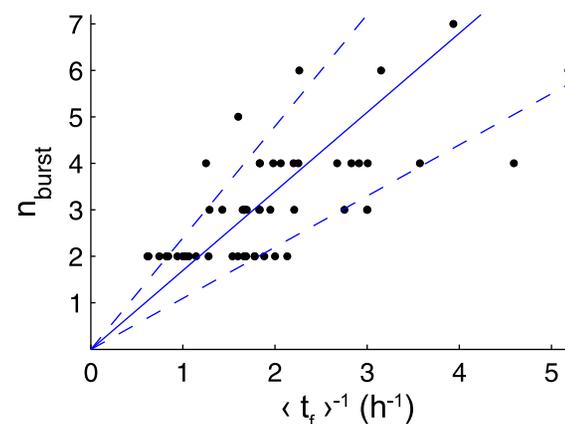
Note that if the process related to  $t_{ex}$  was triggered at each firing, the started burst would never end: it is here implicitly assumed that the excitation process cannot be reactivated until it has fully relaxed.

When  $t_{ex}$  is smaller than  $t_f$  or of the same order, there is only one firing per burst, which means that no burst is observed. This should be the case for usual “random” or “metronomic” traps which present no bursts. Indeed, as seen on Figure 3, the fastest non-bursting traps have  $t_f$  of order 1–2 h which is also the value of  $t_{ex}$  calculated above.

Bursts in *Utricularia* seem to be an evidence of a sensitive process occurring during firings of its traps, suggesting that in addition to  $t_d$  there is another characteristic time of a trap  $t_{ex}$  to take into account. We hypothesize that bursts are likely to happen when the parameter  $\xi = t_{ex} / t_d$  is greater than 1, since spontaneous firings happen with a time interval of order  $t_d$  or more. However, there are still open questions to explore, either on the precise chemo-physical mechanism explaining the origin of  $t_{ex}$  or on the benefit bursts could bring to the plant: is it a way to increase the capture rate of a trap when animals are close, or is it just an unavoidable effect of the global trapping mechanism of *Utricularia*?

## Conclusion

Early investigations on spontaneous firings of *Utricularia* traps suggested that they were randomly distributed in time. We proved



**Figure 12. Number of firings inside a burst versus firing frequency.** The firing frequency is defined as the inverse of the mean time interval between firings inside a burst  $\langle t_f \rangle$ . It is calculated for 53 bursts of 5 different traps of *Utricularia inflata* and one trap of *Utricularia australis*, along with the number of firings per burst  $n_{burst}$ . Blue lines correspond to  $n_{burst} \times \langle t_f \rangle = t_{ex}$  with values of  $t_{ex}$  of 1.7 h (central line), 1.1 h (lowest dashed line) and 2.4 h (upper dashed line). doi:10.1371/journal.pone.0020205.g012

here that these apparently random distributions were just one aspect of a larger set of behaviors, which can be very regular and organized in time. All these behaviors can be found on different traps of a composed leaf, suggesting complementary roles: in addition to catching occasional preys as “waiting” and very “random” traps do, the very regular firings of “metronomic” traps could be a way to diversify the plant’s alimentation by continuously catching smaller organisms not capable of triggering the trap by themselves, such as phytoplankton or bacteria. This underlines the importance of these organisms for the plant’s nutrients supply, as recently suggested [12,13].

We also proposed a physical model, showing how the short and long time behavior of the traps were connected: fast opening of the door and spontaneous firings are just two consequences of a single aspect which is the buckling of the door wall. Thus, to achieve its regular firings without any active signal or feedback, the plant simply uses mechanical oscillations, which only ingredients are continuous pumping and buckling of the trap door. Based on this idea, the different trap types can be explained by introducing fluctuations in the mechanical parameters, which occur naturally due to biological or geometrical changes. The key parameters to predict the behavior of a trap are  $\Delta p_c$ , the critical pressure at which its door buckles, and  $\Delta p_d$ , the maximum pressure difference it can generate by active pumping. We suggested that  $\Delta p_c$  and  $\Delta p_d$  were always of the same order, optimizing the trap efficiency, and that their relative values condition the temporal aspect of firings. It has to be noted that even if our model strongly supports the idea of buckling as the mechanism for firing *Utricularia* traps, it doesn’t exclude any sensitive effect of the trigger hairs, which could act chemically or mechanically to facilitate buckling. Note that the presence of a sensitive process is also suggested by the bursting behavior of some traps.

This ingenious way to create a periodic signal, recalling some aspects of Tantalus vase, could provide biomimetical inspiration for autonomous elastic structures, and represents in itself an original illustration of mechanical oscillators for an undergraduate Physics course.

Hopefully this work will stimulate further collaboration between biologists and physicists to clarify completely the mechanical and biological processes at the root of the unique trapping mechanism of *Utricularia*. One big challenge is a direct, non destructive measurement of the pressure inside the trap, which is for now only accessible by looking at the trap width. Future work could also be directed towards the characterization of the bio-chemical response resulting of action on the trigger hairs, or of temporal behaviors for other *Utricularia* species.

## Materials and Methods

### Preparation of excised leaves

Composed leaves from aquarium-cultivated *Utricularia inflata* and *Utricularia australis* obtained from “Nature et Paysages”, France, were excised keeping bladders, and carefully washed with deionized (DI) water before immersion in a Petri dish filled with DI water. Special care was taken to avoid accidentally triggering the traps when transferred, usually leading to the aspiration of an air bubble. The leaf was held at the bottom of the petri dish by inoxidable tweezers. Volumes of DI water used in the experiments were small (< 50 mL), so we cannot exclude the presence of solutes such as minerals in an unknown concentration, brought by the plant itself for example. As a matter of fact, authors of previous studies of *Utricularia* cited in this article (see for example [8,10]) usually add a small quantity of ions in water to reproduce natural living conditions. However, since the excised leaves continued to

live and grow for more than three weeks and most of the traps presented regular deflation - firing cycles, our liquid medium was probably adapted, even if not optimized.

### Observation

The Petri dish rested on a LED Backlight device (from LEICA, France) consisting of 20 6-watts white LED at color temperature 5000K, distributed on a 55 mm disk under a light diffuser. Such constant illumination was used to avoid any effect due to ambient light. Images were recorded with a time-lapse camera, allowing observation of *Utricularia* traps for long times, of the order of several weeks. Petri dishes were not covered to avoid condensation, so they had to be regularly refilled with care, typically each week. The room temperature during observation was  $23 \pm 4$  degrees Celsius and no effect of temperature variations on the trap behaviors was observed.

Two composed leaves of *Utricularia inflata* respectively containing 10 and 12 traps were followed continuously during 3 weeks. All of the 22 observed traps showed spontaneous firings, even if 6 of them stopped firing after 1 to 3 days. Among these 6 latter traps, 2 traps fired again a few days later, showing that they were still working. One trap also oscillated between periods of firings and periods of apparent inactivity, each during about 3 days. All other traps had constant firing activity.

Two other experiments were conducted with single traps: one with *Utricularia inflata* (trap B) and one with *Utricularia australis* (trap D).

On these 24 traps, a total amount of 1549 spontaneous firings were recorded. The bursting behavior was observed on 6 different traps.

### Image analysis and data processing

Image and data were processed using ImageJ freeware and Matlab (Mathworks), to extract the times at which observed firings occur, the time intervals between firings and their distribution.

If possible, the evolution of the trap thickness in time was also recorded, by extraction of the lateral dimension of a thresholded image of the trap. The characteristic pumping time  $t_d$  was then calculated by exponential fitting on these curves: if we define the deflated state  $w = w_d$  as the value of the exponential plateau (since that due to spontaneous firings,  $w_d$  is never attained, its value has to be manually adjusted) and if  $w_i$  is the inflated width, we define a degree of inflation

$$x(t) = \frac{w(t) - w_d}{w_i - w_d} \quad (7)$$

which value is  $x=0$  for a fully deflated state, and  $x=1$  for an inflated state, just after any firing.  $t$  is reset to zero for each firing so that one has

$$x(t) = \exp(-t/t_d). \quad (8)$$

Uncertainties on  $t_d$  represent the standard deviation of the fitting parameter estimated in the regression process. Figure 7 presents on a same graph five successive firings of trap B showing that the deflation process is identical after all firings.

Some values of  $t_d$  were also determined graphically with the methods of tangents, uncertainty is then an estimate of the error made on the slope of the curve at its origin. For the purpose of this article, precise determination of values and their uncertainties is not essential and the order of magnitudes extracted are enough to discuss the results.

For purposes involving more precise measurement of the trap width, the image analysis technique could also be used. Its precision is relatively poor for small magnifications (Figure 1 for example), due to the important pixel size: in this case, the precision of the measurement of  $w$  is of the order of 10 to 20  $\mu\text{m}$ , but it can be greatly improved using higher magnifications. The measurement of  $w$  as shown on Figures 5 and 6 has a precision of 2 to 3  $\mu\text{m}$ . The drawback of using high magnifications is the loss of the ability to follow several traps at the same time.

Compared to the linear position sensor used in [8], the image analysis technique has the advantage to avoid any direct contact with the plant, but it only accesses a projected width of the trap, making it sensitive to any natural rotation of the trap. The combined use of these two techniques should thus be advantageous.

### Trap pumping and pressure evolution

The observed saturation of deflation to the fully deflated state shows that there is a process balancing pumping for high deflation. Two hypotheses can be formulated: either the pumping rate depends on the pressure difference  $\Delta p$  between the inside and the outside of the trap, so that pumping could be significantly lowered in the deflated state, or there is an incoming water flow balancing pumping due to porous fluid transport. Since lateral walls of the trap are thin, the latter is probable. We show that these processes can explain an exponential decrease of pressure inside the trap, using simple hypotheses: assuming that water is expelled from a trap with a constant flow rate  $Q_0$ , the volume  $V$  of water inside the trap should decrease linearly with time as  $\frac{dV}{dt} = -Q_0$ . However, if the wall of the trap is porous, there will be an incoming flow rate  $Q$ . Due to the slowness of the pumping process and the small lengthscales involved, Darcy's law should be verified, and  $Q$  should directly be proportional to the pressure gap  $\Delta p$ :

$$Q = \alpha \Delta p \quad (9)$$

with

$$\alpha = \frac{kS}{h\eta} \quad (10)$$

where  $S$  is the surface of the trap,  $h$  and  $k$  are respectively the thickness and permeability of the wall and  $\eta$  the viscosity of water. The volume conservation equation thus implies

$$\frac{dV}{dt} = -Q_0 + \alpha \Delta p. \quad (11)$$

We now assume that due to the elasticity of the trap wall, there is a linear relationship between pressure and volume such that  $\Delta p = K(V_0 - V)/V_0$  where  $V_0$  is the initial inflated volume of the trap, and  $K$  is a positive effective elasticity modulus (in Pa). This hypothesis was justified by numerical simulations with realistic *Utricularia* shapes [3], showing that  $K$  was constant except for very small deflations (volume change inferior to 5%). These simulations also showed that the trap volume  $V$  and width  $w$  are proportional, which is due to the fact that the trap deforms mainly in the lateral direction (indicated by the arrows on Figure 5). As a result, the assumption of linearity between  $w$  and  $\Delta p$  used in the Discussion section is reasonable.

Equation (11) then rewrites as

$$\frac{d\Delta p}{dt} = -\frac{1}{t_d}(\Delta p - \Delta p_d) \quad (12)$$

with

$$\begin{cases} \Delta p_d = \frac{Q_0}{\alpha} \\ t_d = \frac{V_0}{K\alpha} \end{cases} \quad (13)$$

We recognize here a first order differential equation which admits (2) as a unique solution for the initial condition  $\Delta p_{t=0} = 0$ .

From equation (13) we can estimate the trap permeability  $k$  using equation (10) rewritten as  $k = \alpha \frac{h\eta}{S}$  and using typical values for an *Utricularia inflata* trap:  $S \simeq 1 \text{ mm}^2$ ,  $h \simeq 100 \mu\text{m}$ ,  $Q_0 \simeq 0.02 \text{ mm}^3/\text{min}$ ,  $\Delta p_d \simeq 15 \text{ kPa}$  and the viscosity of water  $\eta \simeq 10^{-3} \text{ Pa}\cdot\text{s}$ , one finds  $k \simeq 2 \cdot 10^{-18} \text{ m}^2$  (see [9] for a similar estimation). If trap permeability is not the only phenomenon causing leaks, the flow rate  $Q$  due to porous leaks should be lower, so that this estimate is a maximum value for the trap permeability. Notice also that due to the inhomogeneous character of the trap wall, the obtained value is an equivalent permeability averaged over all its surface and thickness.

This model is consistent with the experimental value of  $t_d$ : using  $V_0/K = \Delta V_d/\Delta p_d$ ,  $\Delta V_d \simeq 1 \text{ mm}^3$  being the difference of volume between the inflated and deflated state, and equation (13), one finds  $t_d \simeq 50 \text{ min}$ .

At last, note that the exponential evolution of pressure is also compatible with a model (not detailed here) using the hypotheses of zero permeability of the trap wall and of a pumping rate  $Q$  depending linearly on  $\Delta p$ .

### Permeability and equivalent radius

In the above paragraph, we showed that a trap could not go beyond a maximum pumping pressure  $\Delta p_d$  due to porous leaks. However, leaks could also come from a single hole of radius  $r$  in a perfectly impermeable trap. Then if the Reynolds number (see below) is sufficiently low,  $Q$  is also proportional to  $\Delta p$  creating a Poiseuille flow with hydraulic resistance

$$R_h = \frac{\Delta p}{Q} = \frac{8\eta h}{\pi r^4}. \quad (14)$$

Since  $\Delta p/Q$  was equal to  $\alpha^{-1}$  in the permeability model above, we find using equations (10) and (14) that

$$r = \left(\frac{8}{\pi} S k\right)^{1/4}. \quad (15)$$

Using the previous value of  $1 \text{ mm}^2$  for  $S$ , one finds  $r \simeq 1 \mu\text{m}$ .

The Reynolds number is expressed by  $Re \simeq r\rho v/\eta$  where  $v$  is the fluid velocity in the hole and  $\rho$  the fluid density. If there is a flow rate  $Q_0$  through the hole of radius  $r$  we should have  $v \sim Q_0/r^2$  so that we have  $Re \sim \rho Q_0/r\eta \simeq 0.2$ . The approximation of Poiseuille flow is thus justified.

Notice that the values of the permeability  $k$  is very low, meaning that the water fluxes in and out of the trap are very small. As can be seen on the equivalent hydraulic radius of  $1 \mu\text{m}$ , the trap door has to be perfectly closed to avoid any opening of this order of magnitude. This also underlines the difficulty of intrusive measurements of the inside pressure of the trap such as those of

[10] which are likely to give biased results, since any hole of order 1 mm entails water fluxes comparable to the maximum ones that naturally occur. Such a provoked leak would considerably lower the observed value of  $\Delta p_d$ .

### Fluctuation propagation between $\tau$ and $\pi$

We assume here that probability distributions are smooth enough so that mean values and standard deviations are defined.

Expanding  $\tau(\Pi)$  around  $\langle \Pi \rangle$  at the second order in  $\Pi - \langle \Pi \rangle$  and taking the average one finds

$$\langle \tau \rangle = \tau(\langle \Pi \rangle) + \frac{1}{2} \sigma_{\Pi}^2 \left( \frac{d^2 \tau}{d\Pi^2} \right)_{\Pi = \langle \Pi \rangle}. \quad (16)$$

Notice that the first order term is cancelled in the averaging process. Using the expression of  $\tau(\Pi)$  given in equation (5), one finds that the second order term is negligible when

$$\left( \frac{\sigma_{\Pi}}{1 - \langle \Pi \rangle} \right)^2 \ll 2 \langle \tau \rangle \quad (17)$$

so that the mean value of  $\tau$  is simply given by

$$\langle \tau \rangle = \tau(\langle \Pi \rangle). \quad (18)$$

This allows calculating the mean value of  $\tau$  knowing the mean value of  $\Pi$ , or by reversing the equation deducing the mean value of  $\Pi$  by measuring the mean value of  $\tau$  experimentally.

Standard deviation can also be calculated with the same Taylor expansion technique. The result for  $\sigma_{\tau}^2$  brings into play a sum of terms proportional to  $\langle (\Pi - \langle \Pi \rangle)^n \rangle$ , starting at  $n=2$ . The  $n=3$  term equals 0 for symmetrical distributions of  $\Pi$ , which we will assume in the following. This is also the case for all odd terms. Even terms are of order  $(\sigma_{\Pi}/(1 - \langle \Pi \rangle))^2$  with respect to the previous even term so that one can keep only the term  $n=2$  at the condition

$$\left( \frac{\sigma_{\Pi}}{1 - \langle \Pi \rangle} \right)^2 \ll 1 \quad (19)$$

giving the simple result

$$\sigma_{\tau} = \frac{\sigma_{\Pi}}{1 - \langle \Pi \rangle}. \quad (20)$$

Notice that conditions (17) and (19) can be rewritten respectively as  $\sigma_{\tau}^2 \ll 2 \langle \tau \rangle$  and  $\sigma_{\tau}^2 \ll 1$ . In our experiments  $\tau$  is always bigger than 0.5 so that the second condition is the most restrictive.

Comparison with experimental values is difficult for several reasons: first, the previous inequalities are not verified for “random” traps. “Metronomic” traps have smaller fluctuations but these are difficult to measure, since it is not easy to separate actual fluctuations

from the natural drift of the firing period. Second, results strongly depend on the exact mathematical relation between  $\tau$  and  $\Pi$ , which is not accessible experimentally for now, especially for long times.

### Probability distributions

Assuming that the variable  $\Pi$  has gaussian fluctuations with standard deviation  $\sigma_{\Pi}$ , the associated probability distribution is

$$p_{\Pi}(\Pi) = \frac{\exp\left(-\frac{(\Pi - \langle \Pi \rangle)^2}{2\sigma_{\Pi}^2}\right)}{\sigma_{\Pi}\sqrt{2\pi}}$$

and one has the relation  $p_{\tau} = \frac{d\Pi}{d\tau} p_{\Pi}$  which gives, using equation (5):

$$p_{\tau}(\tau) = \frac{\exp\left(-\frac{(1 - \langle \Pi \rangle) - \exp(-\tau))^2}{2\sigma_{\Pi}^2} - \tau\right)}{\sigma_{\Pi}\sqrt{2\pi}}$$

which is a function of  $\tau$  and of the parameters of the initial  $\Pi$  distribution: its mean  $\langle \Pi \rangle$  and its standard deviation  $\sigma_{\Pi}$ . On Figure 11,  $\langle \Pi \rangle$  and  $\sigma_{\Pi}$  are chosen to illustrate the basic properties of such a distribution, namely the amplification of its standard deviation as  $\langle \Pi \rangle$  increases, and its asymmetry.

### Supporting Information

#### Video S1 Spontaneous firings of *Utricularia inflata* traps.

This is the animated version of Figure 1. Ten traps of a same branch of *Utricularia inflata* were immersed in de-ionized water and their spontaneous firings were recorded with a time-lapse camera. The field is about 16.5 × 11.5 mm and the video is accelerated 1680 times (real duration: 11 hours and 12 minutes).

(AVI)

#### Video S2 Bursts in trap D (*Utricularia australis*).

This video corresponds to Figure 9. This trap of *Utricularia australis* was recorded with a time-lapse camera and presented regular bursts of 3 or 4 spontaneous firings. Two bursts of 4 firings are present in the video. The trap is approximately 1 mm long and the video is accelerated 1680 times (real duration: 7 hours and 47 minutes).

(AVI)

### Acknowledgments

We would like to thank Lubomír Adamec for fruitful comments on the present works. We are grateful to Bénédicte Hingant, Marc Joyeux and Simon Poppinga for comments on the first versions of the article.

### Author Contributions

Conceived and designed the experiments: OV IR PM. Performed the experiments: IR OV. Analyzed the data: OV PM IR. Wrote the paper: OV. Developed the physical model: OV.

### References

- Darwin C (1875) Insectivorous Plants. London: John Murray, Albemarle Street.
- Treat M (1876) Is the valve of *Utricularia* sensitive? Harper's New Monthly Magazine 52: 382–387.
- Vincent O, Weißkopf C, Poppinga S, Masselter T, Speck T, et al. (2011) Ultrafast underwater suction traps. Proc Roy Soc B;doi:10.1098/rspb.2010.2292.
- Lloyd FE (1942) The carnivorous plants. Chronica Botanica Company.
- Juniper BE, Robins RJ, Joel DM (1989) The carnivorous plants Academic Press.
- Sydenham P, Findlay G (1975) Transport of solutes and water by resetting bladders of *Utricularia*. Functional Plant Biol 2: 335–351.
- Heide-Jørgensen HS (1981) Parasitisme og carnivori. Kompendium, Institut for Planteanatomi og Cytologi Københavns Universitet.
- Adamec L (2011) The comparison of mechanically stimulated and spontaneous firings in traps of aquatic carnivorous *Utricularia* species. Aquatic Botany 94: 44–49.

9. Joyeux M, Vincent O, Marmottant P (2011) Mechanical model of the ultra-fast underwater trap of *Utricularia*. *Physical Review E* 83: 021911.
10. Sasago A, Sibaoka T (1985) Water extrusion in the trap bladders of *Utricularia vulgaris* L. a possible pathway of water across the bladder wall. *Bot Mat Tokyo* 98: 55–66.
11. Landau LD, Lifshitz EM (1986) *Theory of Elasticity* Elsevier, 3d edition.
12. Gordon E, Pacheco S (2007) Prey composition in the carnivorous plants *Utricularia inflata* and *U. gibba* (Lentibulariaceae) from Paria Peninsula, Venezuela. *Rev Biol Trop* 55: 795–803.
13. Sirová D, Borovec J, Černá B, Rejmánková E, Adamec L, et al. (2009) Microbial community development in the traps of aquatic *Utricularia* species. *Aquatic Botany* 90: 129–136.