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NEW STACKED CENTRAL CONFIGURATIONS FOR THE PLANAR 5–BODY PROBLEM

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Abstract. A stacked central configuration in the n–body problem is one that has a proper subset of the n–bodies forming a central configuration. In this paper we study the case where three bodies with masses m_1, m_2, m_3 (bodies 1, 2, 3) form an equilateral central configuration, and the other two with masses m_4, m_5 are symmetric with respect to the mediatrix of the segment joining 1 and 2, and they are above the triangle generated by $\{1, 2, 3\}$. We show the existence and non-existence of this kind of stacked central configurations for the planar 5–body problem.

Keywords: Planar central configurations, n–body problem, stacked central configurations.

1. Introduction

The classical *planar Newtonian n–body problem* in celestial mechanics consists in studying the motion of n pointlike masses in a fixed plane, interacting among themselves through no other forces than their mutual gravitational attraction according to Newton's gravitational law (Newton 1687).

The center of mass of the system, given by $\sum_{j=1}^{n} m_j r_j / M$, where $M = m_1 + \cdots + m_n$ is the total mass and r_j is the position vector of the mass m_j, is considered at the origin of an inertial system. Usually this inertial system is called the *inertial barycentric* system.

The simplest motion of the n bodies, called a *homographic* solution, is a motion such that the configuration of the n bodies remains the

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same (with respect to the inertial barycentric system) up to a dilation and a rotation of $\mathbb{R}^2$.

The first homographic solutions for the 3–body problem were found by Euler (1767), for which three bodies are *collinear* at any time, and by Lagrange (1873), where the three bodies are at any time at the vertices of an *equilateral triangle*.

At a given instant $t = t_0$ the configuration of the n bodies is *central* if the gravitational acceleration acting on every mass point is proportional to its position (referred to the inertial barycentric system). Central configurations and homographic solutions are linked by the *Laplace theorem* (see for instance Boccaletti and Pucacco 1996; Wintner 1941): the configuration of the n bodies in a homographic solution is central at any instant of time.

If we have a central configuration, any dilation and any rotation (centered at the center of mass) of it provides another central configuration. We say that two central configurations are *similar* if we can pass from one to another through a dilation and a rotation. So we can study the classes of central configurations defined by the above equivalence relation. Thus the 3–body problem has exactly 5 classes of central configurations for any value of the positive masses.

Central configurations of the n–body problem are important because: they allow the computation of homographic solutions; if the n bodies are heading for a simultaneous collision then the bodies tend to a central configuration (see Saari 1980); there is a relation between central configurations and the bifurcations of the hypersurfaces of constant energy and angular momentum (see Smale 1970). See also the Refs. Moeckel 1990; Moulton 1910.

The main general open problem for the planar central configurations is due to Wintner (1941) and Smale (1998): *Is the number of classes of planar central configurations finite for any choice of the (positive) masses* $m_1, \ldots, m_n$?

Hampton and Moeckel in (2006), proved this conjecture for the 4–body problem. The conjecture remains open for $n > 4$. But if one mass can be negative, Roberts (1999), proved that there exists a one–parameter not equivalent family of planar central configurations for the 5–body problem. Also considering the particles endowed with masses and charges, Alfaro and Perez-Chavela (2002) proved the existence of a continuum of central configurations in a particular 4–body problem. Other recent papers on central configurations are due to Corbera, Cors and Llibre (2010), Corbera and Llibre (2010), Gidea and Llibre (2010), Piña and Lonngi (2010), ...

In 2005 Hampton (2005) provides a new family of planar central configurations, called *stacked central configuration*, for the 5–body problem with an interesting property: the central configuration has a subset of three bodies forming a central configuration of the 3–body problem, in fact these three bodies are in an equilateral triangle, and the remaining two bodies are in the interior of the triangle and are located symmetrically with respect to a perpendicular bisector.

Recently the first and second author of this paper gave new examples of stacked central configurations of the 5–bodies which, as the ones studied by Hampton (2005), have three bodies in the vertices of an equilateral triangle, but the other two are on the perpendicular bisector (Llibre and Mello 2008).

In this paper we find a new class of stacked central configurations in a 5–body problem which has three bodies in the vertices of an equilateral triangle and the other two are located symmetrically with respect to a perpendicular bisector in the exterior and above of the triangle, see Figure 1.

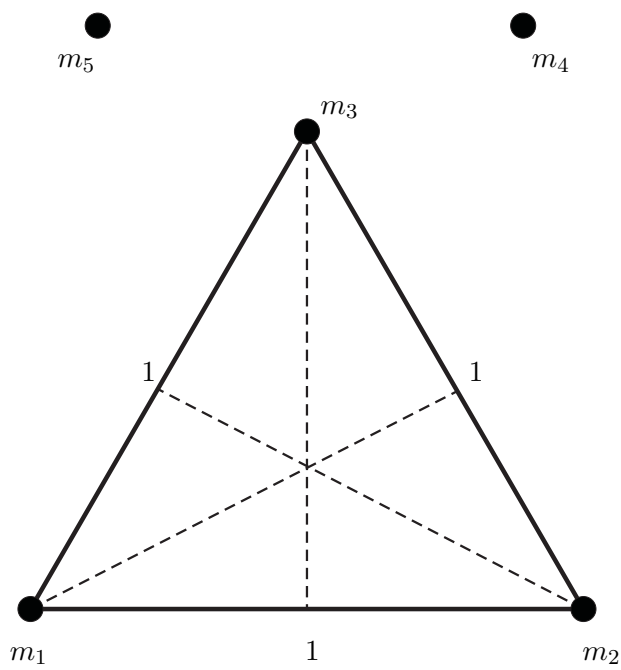

Figure 1. Three bodies at the vertices of an equilateral triangle and two bodies located symmetrically with respect to a perpendicular bisector.

The paper is organized as follows, in Section 2 we prove the existence and non-existence of this new class of stacked central configurations.

In section 3 we show some numerical evidence of the existence and non-existence of another families of this kind of central configurations.

2. Existence and non-existence of stacked central configurations

The equations of motion of the planar Newtonian n–body problem are given by

$$\ddot{r}_i = -\sum_{\substack{j=1\\ j\neq i}}^{n} m_j \frac{r_i - r_j}{r_{ij}^3}, \tag{1}$$

for $i = 1, 2, \ldots, n$. Here the gravitational constant is taken equal to 1, $r_j \in \mathbb{R}^2$ is the position vector of the punctual mass m_j in the inertial barycentric system, and as before $r_{ij} = |r_i - r_j|$ is the Euclidean distance between r_i and r_j.

For the central configurations we have $\ddot{r}_j = \lambda r_j$ with $\lambda \neq 0$ for all $j = 1, \ldots, n$. So from equation (1) we have

$$\lambda r_i = -\sum_{\substack{j=1\\ j\neq i}}^{n} m_j \frac{r_i - r_j}{r_{ij}^3}, \tag{2}$$

for $i = 1, 2, \ldots, n$.

For the planar central configurations instead of working with equation (2) we shall use the Dziobek equations (see Hagihara 1970, p. 241)

$$f_{ij} = \sum_{\substack{k=1\\ k\neq i,j}}^{n} m_k \left(R_{ik} - R_{jk}\right) \Delta_{ijk} = 0, \tag{3}$$

for $1 \leq i < j \leq n$, where $R_{ij} = 1/r_{ij}^3$ and $\Delta_{ijk} = (r_i - r_j) \wedge (r_i - r_k)$. As usual $\wedge$ denotes the cross product of two vectors.

For the 5–body problem equations (3) is a set of ten equations. Our class of configurations with five bodies as in Figure 1 without collisions must satisfy

$$r_{12} = r_{23} = r_{13} = 1, \quad r_{14} = r_{25}, \quad r_{15} = r_{24}, \quad r_{34} = r_{35}, \quad r_{14} > r_{24},$$

$$\Delta_{124} = \Delta_{125}, \quad \Delta_{143} = \Delta_{235}, \quad \Delta_{145} = \Delta_{245}, \quad \Delta_{135} = \Delta_{243}.$$

We are also assuming $\Delta_{145} \neq 0$, that is, the bodies m_1, m_4 and m_5 are not on the same straight line.

The equation $f_{45} = 0$ of (3) becomes

$$(R_{14} - R_{15})\Delta_{145}(m_1 - m_2) = 0.$$

Thus $m_1 = m_2$. Similarly the equation $f_{12} = 0$ of (3) goes over to

$$(R_{14} - R_{24})\Delta_{124}(m_4 - m_5) = 0.$$

Therefore $m_4 = m_5$. Substituting $m_1 = m_2$ and $m_4 = m_5$ into the other eight equations of (3) it follows that

$$f_{13} = 0 \Leftrightarrow f_{23} = 0, \quad f_{15} = 0 \Leftrightarrow f_{24} = 0,$$

$$f_{14} = 0 \Leftrightarrow f_{25} = 0, \quad f_{34} = 0 \Leftrightarrow f_{35} = 0,$$

which can be written respectively as

$$\begin{aligned} L = (R_{14} - R_{34})\Delta_{134} + (R_{15} - R_{34})\Delta_{135} &= 0, \quad (4)\\ (R_{12} - R_{15})\Delta_{142}\, m_1 + (R_{12} - R_{34})\Delta_{143}\, m_3 + (R_{15} - R_{45})\Delta_{145}\, m_4 &= 0, \quad (5)\\ (R_{12} - R_{14})\Delta_{142}\, m_1 + (R_{12} - R_{34})\Delta_{153}\, m_3 + (R_{14} - R_{45})\Delta_{154}\, m_4 &= 0, \quad (6)\\ [(R_{12} - R_{14})\Delta_{134} + (R_{12} - R_{15})\Delta_{153}]\, m_1 + (R_{34} - R_{45})\Delta_{345}\, m_4 &= 0. \quad (7) \end{aligned}$$

Denote by $A = [a_{ij}]$ the matrix of the coefficients of the homogeneous linear system in the variables m_1, m_3 and m_4 defined by equations (5), (6) and (7). Thus

$$a_{11} = (R_{12} - R_{15})\Delta_{142}, \quad a_{12} = (R_{12} - R_{34})\Delta_{143}, \quad a_{13} = (R_{15} - R_{45})\Delta_{145},$$

$$a_{21} = (R_{12} - R_{14})\Delta_{142}, \quad a_{22} = (R_{12} - R_{34})\Delta_{153}, \quad a_{23} = (R_{14} - R_{45})\Delta_{154},$$

$$a_{31} = (R_{12} - R_{14})\Delta_{134} + (R_{12} - R_{15})\Delta_{153}, \quad a_{32} = 0, \quad a_{33} = (R_{34} - R_{45})\Delta_{345}.$$

In order to have a solutions different from $m_1 = m_3 = m_4 = 0$ the determinant of the matrix A must be zero, where

$$\det A = (R_{12} - R_{34})\,[(R_{12} - R_{15})\Delta_{153} - (R_{12} - R_{14})\Delta_{143}]\ \Gamma, \qquad (8)$$

with Γ equal to

$$(R_{34} - R_{45})\Delta_{142}\Delta_{345} + [(R_{14} - R_{45})\Delta_{143} + (R_{15} - R_{45})\Delta_{153}]\,\Delta_{154}. \quad (9)$$

From equation (4) one has $R_{14}\Delta_{134} = R_{34}\Delta_{134} + (R_{34} - R_{15})\Delta_{135}$. Substituting this equation into equation (9) it follows that $\Gamma = 0$ and therefore, from (8), we get $\det A = 0$. So the homogeneous linear system in the variables m_1, m_3 and m_4, defined by equations (5), (6) and (7), has nontrivial solutions under the assumption $L = 0$, which can be viewed as a constraint on the geometry of the configuration.

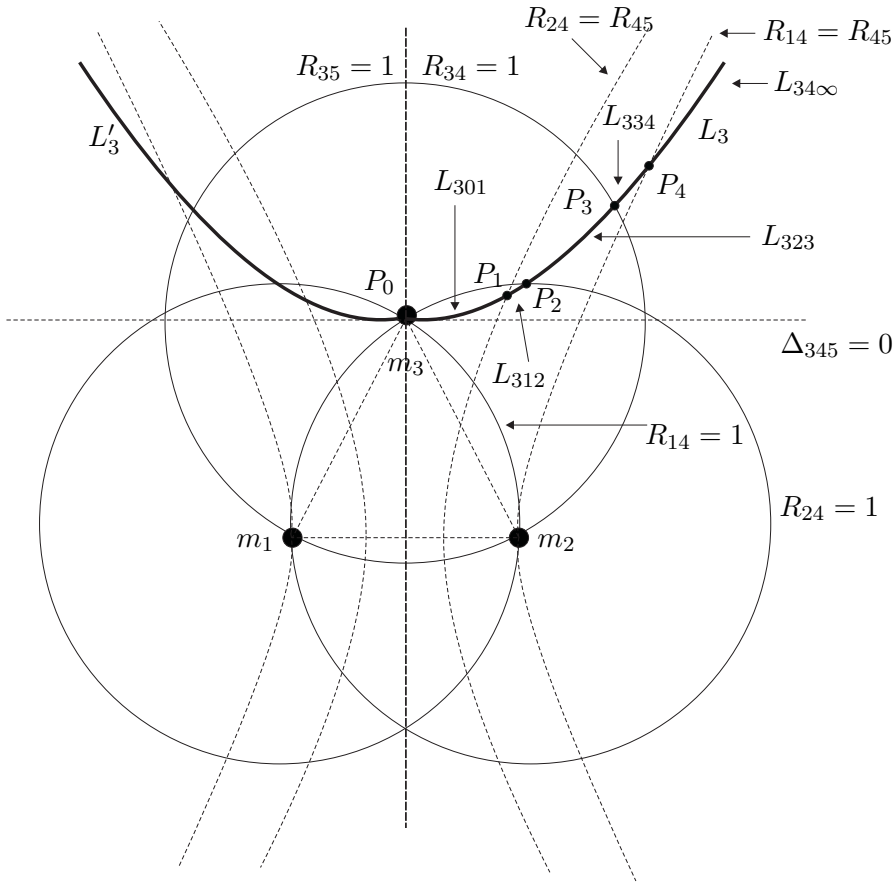

Figure 2. The component of the curve $\{L = 0\}$ in the upper semiplane and its components L_3 and L_3'.

LEMMA 1. *Consider the regions*

$$S_1 = \{R_{45} > R_{34} > R_{14} > 1\}, \quad S_2 = \{1 > R_{14} > R_{34} > R_{45}\},$$

$$S_3 = \{\Delta_{345} > 0, \Delta_{134} < 0, R_{34} > R_{24}\}$$

and S_1', S_2', S_3' *the symmetrical regions with respect to the mediatrix of the segment joining* r_1 *and* r_2. *See Figure 3. The curves*

$$L_i = \{L = 0\} \cap S_i, \quad L_i' = \{L = 0\} \cap S_i',$$

$i = 1, 2, 3$, *are well–defined and are not empty. See Figure 2.*

Proof. Without loss of generality we can take a coordinate system such that $r_1 = (-1/2, 0)$, $r_2 = (1/2, 0)$ and $r_3 = (0, \sqrt{3}/2)$ respectively.

Thus $r_4 = (x, y)$ and $r_5 = (-x, y)$, where $x > 0$ and $y \in \mathbb{R}$. With these coordinates we have

$$\begin{aligned} L(x,y) &= \left[\frac{1}{((x+1/2)^2+y^2)^{3/2}} - \frac{1}{(x^2+(y-\sqrt{2}/2)^2)^{3/2}}\right] M_1 \\ &+ \left[\frac{1}{((x-1/2)^2+y^2)^{3/2}} - \frac{1}{(x^2+(y-\sqrt{2}/2)^2)^{3/2}}\right] M_2, \end{aligned}$$

where $M_1 = \frac{y}{2} - \frac{x\sqrt{3}}{2} - \frac{\sqrt{3}}{2}$ and $M_2 = \frac{y}{2} + \frac{\sqrt{3}x}{2} - \frac{\sqrt{3}}{2}$. It is easy to see from the above equation that the curve $L(x, y) = 0$ is symmetric with respect to the mediatrix of the segment joining r_1 and r_2. Using *Mathematica* we have plotted the graph of the curve $L(x, y) = 0$, see Figure 2.

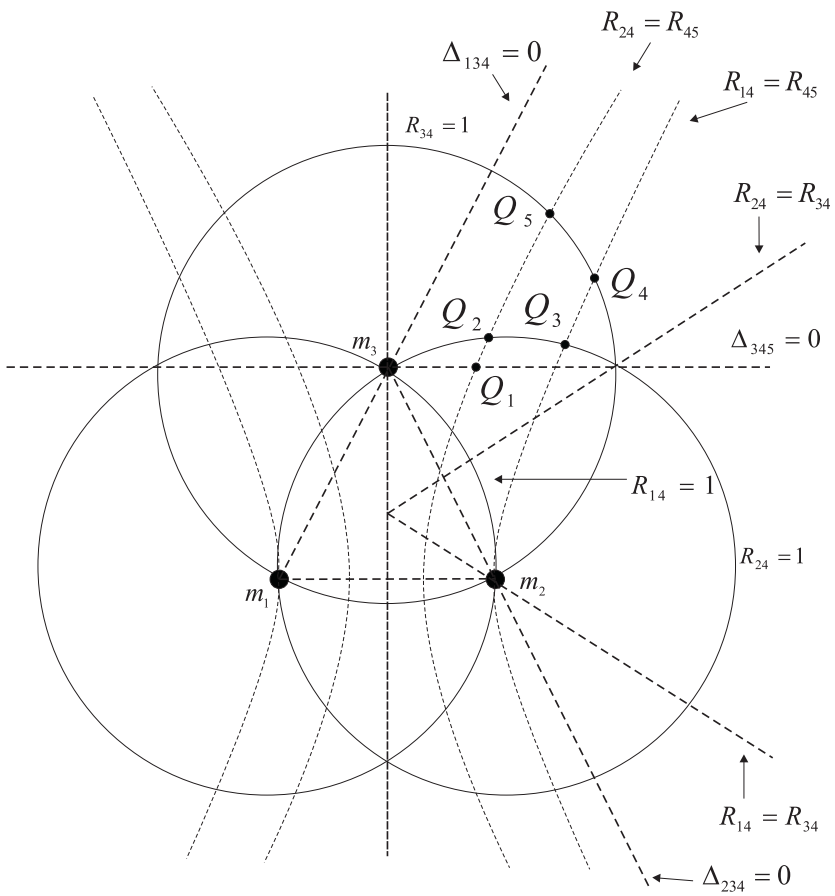

Figure 3. Some straight lines and curves for the calculation of the sign of L.

Define the points

$$\begin{aligned}
Q_1 &= \{\Delta_{345}=0\}\cap\{R_{24}=R_{45}\}\cap\overline{S_3},\\
Q_2 &= \{R_{24}=1\}\cap\{R_{24}=R_{45}\}\cap S_3,\\
Q_3 &= \{R_{24}=1\}\cap\{R_{14}=R_{45}\}\cap S_3,\\
Q_4 &= \{R_{34}=1\}\cap\{R_{14}=R_{45}\}\cap S_3,\\
Q_5 &= \{R_{34}=1\}\cap\{R_{24}=R_{45}\}\cap S_3,
\end{aligned}$$

according to Figure 3. The symmetrical points are denoted by Q_i', $i=1,2,3,4,5$. In (Hampton 2005) Hampton studied stacked central configurations when r_4 and r_5 are on the components L_1 and L_1' respectively. If $r_4\in\{\Delta_{234}=0\}$ (see Figure 3) then $r_5\in\{\Delta_{135}=0\}$. Thus $L=(R_{14}-R_{34})\Delta_{134}<0$. If $r_4\in\{R_{14}=R_{34}\}$ (see Figure 3) then $L=(R_{15}-R_{34})\Delta_{135}>0$. Therefore there is a connected component of $L=0$ between the lines $\{\Delta_{234}=0\}$ and $\{R_{14}=R_{34}\}$ denoted by L_2. See Figure 2. Its symmetrical component is denoted by L_2'.

If $r_4\in\{\Delta_{134}=0\}$ (see Figure 3) then $L=(R_{24}-R_{34})\Delta_{135}<0$. If $r_4\in\{R_{24}=R_{34}\}$ (see Figure 3) then $L=(R_{14}-R_{34})\Delta_{134}>0$. By the same argument $L>0$ on the set $\{\Delta_{345}=0\}$. Therefore there is a connected component of $L=0$ between the lines $\{\Delta_{134}=0\}$ and $\{R_{24}=R_{34}\}\cup\{\Delta_{345}=0\}$ denoted by L_3. See Figure 2. Its symmetrical component is denoted by L_3'.

In the coordinate system that we have introduced previously one has

$$Q_1=\left(\frac{\sqrt{13}-1}{6},\frac{\sqrt{3}}{2}\right), Q_2=\left(\frac{1}{2},1\right), Q_3=\left(\frac{1+\sqrt{5}}{4},\sqrt{1-\left(\frac{\sqrt{5}-1}{4}\right)^2}\right),$$

$$Q_4=(x_4,y_4),\quad Q_5=(x_5,y_5),$$

where

$$\begin{aligned}
x_4 &= \frac{1}{4}\left(\sqrt{3}\cos(\pi/18)+2\cos(\pi/9)+\sin(\pi/18)\right),\\
y_4 &= \frac{1}{2}\left(\cos(\pi/18)-\sqrt{3}(\sin(\pi/18)-1)\right),\\
x_5 &= \frac{1}{2}\sqrt{2+\cos(\pi/9)-\sqrt{3}\sin(\pi/9)},\\
y_5 &= \frac{1}{2}\left(\cos(\pi/18)+\sqrt{3}(1+\sin(\pi/18))\right).
\end{aligned}$$

From elementary calculations one has the following things for L: $L(Q_1)>0$, $L(Q_2)<0$, $L(Q_3)>0$, $L(Q_4)>0$ and $L(Q_5)<0$.

■

Since Q_2, Q_3 are on the circle $R_{24} = 1$ and $L(Q_2) < 0, L(Q_3) > 0$, by continuity there exists at least one point in the intersection of the curve L_3 with the circle $R_{24} = 1$. There is numerical evidence that this point is unique, we denote it as P_2. Let be P_2' be the symmetric point respect to the mediatrix of the segment joining the particle with mass m_1 with the particle with mass m_2.

The main result of this paper is the following.

THEOREM 2. *Assume that the points r_1, r_2 and r_3 (with positive masses m_1, m_2 and m_3) are at the vertices of an equilateral triangle, whose sides have length 1, and the points r_4 and r_5 (with masses m_4 and m_5) are located symmetrically with respect to a perpendicular bisector, as in Figure 1. Then there exist stacked central configurations for which the particles with masses m_4 and m_5 are in a small arc of L_3 near the points P_2 and P_2' respectively.*

Proof. When the particles 4 and 5 are at P_2 and P_2' respectively, then since $R_{15} = R_{24} = 1$ we obtain $a_{31} = (1 - R_{14})\Delta_{134} < 0$. Therefore the signs of the coefficients of the matrix A are given by

$$A\Big|_{P_2 \cup P_2'} = \begin{pmatrix} 0 & - & + \\ - & + & + \\ - & 0 & + \end{pmatrix}.$$

By the first row of the above matrix the masses m_3 and m_4 have the same sign and by the last row of the matrix the masses m_1 and m_4 also have the same sign. So the masses m_1, m_3 and m_4 have the same sign. Hence when the particles 4 and 5 are at P_2 and P_2' respectively, there exists masses m_1, m_3 and m_4 which form a stacked central configuration.

Now, since a_{31} depends continuously on the mutual distances, we assure that $a_{31} < 0$ when particles 4 and 5 are in a small neighborhood of P_2 and P_2' respectively, the same is true for the other coefficients of a_{ij}, which preserve its sign in a small neighborhood of P_2 and P_2'. Using the constraint $L = 0$ we obtain the small arc cited in the statement of the theorem.

■

LEMMA 3. *The function $a_{31} = (R_{12} - R_{14})\Delta_{134} + (R_{12} - R_{15})\Delta_{153}$ is negative on the set $\{L = 0\} \bigcap \{R_{24} < 1\}$.*

Proof. For our class of central configurations one has $R_{14} < R_{24}$. Substituting this inequality into the expression of a_{31} we have $a_{31} < (\Delta_{143} + \Delta_{135})(R_{24} - 1) < 0$ if $r_{24} > 1$. Thus the lemma follows for the components of L_3 in the exterior of the circle $\{R_{24} = 1\}$.

■

Let P_0 be the vertex of the triangle at r_3. If the particles 4 and 5 are on the curve $L = 0$ and close enough to P_0, there are no stacked central configurations. In other words we have the following property

PROPOSITION 4. *Assume that we have the masses m_1, m_2 and m_3 at the vertices of an equilateral triangle, whose sides have length 1, and the masses m_4 and m_5 are at $L_3 = 0$ and L_3' respectively, located symmetrically with respect to a perpendicular bisector. If r_4 and r_5 are sufficiently close to P_0, then there are no positive masses $m_1 = m_2$, $m_4 = m_5$ and m_3 such that the five bodies with these masses form stacked central configurations.*

Proof. With the hypothesis of the proposition, we have $\Delta_{142} < 0$, by the other hand the point r_4 is inside the circle $R_{24} = 1$ which corresponds to the region $R_{24} > 1$, that is $R_{24} = R_{15} < 1$, therefore $a_{11} = (R_{12} - R_{15})\Delta_{142} < 0$ in this region and the signs of the coefficients of the matrix A are given by

$$A = \begin{pmatrix} - & - & - \\ - & + & + \\ a_{31} & 0 & + \end{pmatrix}.$$

The first row in the above matrix implies that $m_1 = m_2$, $m_4 = m_5$ and m_3 cannot have the same sign.

■

3. Numerical evidence of new stacked central configurations

There is numerical evidence that L_3 intersects the hyperbola $R_{14} = R_{45}$ at a unique point denoted by P_4. Also based on numerical evidences we define $L_{34\infty}$ to be the unbounded connected component of L_3 with one endpoint at P_4. As usual, $L'_{34\infty}$ is the unbounded connected component of L_3' with one endpoint at P_4'.

LEMMA 5. *There are no stacked central configurations when r_4 and r_5 are on $L_{34\infty}$ and $L'_{34\infty}$, respectively.*

Proof. When r_4 and r_5 are at P_4 and P_4', respectively, it is easy verify that $a_{21} < 0$, $a_{22} < 0$ and $a_{23} = 0$ and this implies that the masses m_1 and m_3 have opposite signs.

On the other hand, if r_4 and r_5 are on the open arcs $L_{34\infty}$ and $L'_{34\infty}$, respectively, we have that $a_{21} < 0$, $a_{22} < 0$ and $a_{23} < 0$ and this implies that two of the masses m_1, m_3 and m_4 have opposite signs. This proves (5).

■

There is numerical evidence that L_3 intersects the hyperbola $R_{24} = R_{45}$ in a unique point denoted as P_1, and intersects the circle $R_{34} = 1$ in a unique point denoted as P_3. We now define: L_{312}, L_{323} and L_{334} be the open arcs L_{3ij} of L_3, whose endpoints are P_i and P_j, for $i, j = 1, 3, 4$. More precisely,

$$L_{312} = L_3 \cap \{\Delta_{345} > 0, R_{24} > 1, R_{24} > R_{45}, R_{14} < R_{45}\},$$

$$L_{323} = L_3 \cap \{R_{24} < 1, R_{24} > R_{45}, R_{14} < R_{45}, R_{34} > 1\},$$

$$L_{334} = L_3 \cap \{R_{34} < 1, R_{24} > R_{45}, R_{14} < R_{45}\}.$$

The same definitions are valid for L'_3, see Figure 2.

In order to give numerical evidence for the existence of new stacked central configurations, we have calculated the coordinates of the points $P_i, i = 1, 2, 3, 4$

$$P_1 = (0.45179079888471224, 0.9022946174921869),$$

$$P_2 = (0.63217335448833913, 0.9912266160486711),$$

$$P_3 = (0.93395902567645331, 1.2234055072562044),$$

$$P_4 = (0.95907733939463801, 1.2451548857175381).$$

LEMMA 6. *There are no stacked central configurations when r_4 and r_5 are on the components L_{334} and L'_{334}, respectively.*

Proof. For $r_4 \in L_{334}$ and $r_5 \in L'_{334}$ the signs of the coefficients of the matrix A are given by

$$A\Big|_{L_{334} \cup L'_{334}} = \begin{pmatrix} - & + & + \\ - & - & + \\ - & 0 & + \end{pmatrix}.$$

Using the last row of the above matrix we conclude that m_1 and m_4 have the same sign. Unfortunately we cannot obtain more information about the sign of the mass m_3 from the above matrix. Nevertheless, equations (5) and (7) define two planes through the origin in the space (m_1, m_3, m_4). The normal vectors of these planes are, respectively, $n_1 =$

(a_{11}, a_{12}, a_{13}) and $n_3 = (a_{31}, 0, a_{33})$. Let $T = (T_1, T_2, T_3) = n_1 \wedge n_3$ be the vector parallel to the straight line defined by the intersection of these two planes. Thus there exist positive masses m_1, m_3 and m_4 as solutions of equations (5) and (7) if and only if the components of the vector T have the same sign. It is easy verify that

$$T_1 = a_{12}a_{33} > 0, \quad \text{and} \quad T_3 = -a_{12}a_{31} > 0.$$

In order to see that $T_2 = a_{13}a_{31} - a_{11}a_{33} < 0$, and since we already know the coordinates of the points P_3 and P_4, we can verify numerically that the values of T_2 on the curve $L = 0$ between the points P_3 and P_4 are in the interval $(-0.376045, -0.368863)$. As not all the above components have the same sign, Lemma 6 has been proved.

■

LEMMA 7. *When r_4 and r_5 are at P_3 and P_3', respectively, there are no stacked central configurations.*

Proof. When r_4 and r_5 are at P_3 and P_3', the matrix A takes the form

$$A = \begin{pmatrix} -0.664093 & 0 & 0.694115 \\ -1.04073 & 0 & 0.00940803 \\ -1.0721 & 0 & 0.56513 \end{pmatrix}.$$

Then, since columns 1 and 3 in the above matrix are linearly independent the masses m_1 and m_4 must vanish. Therefore there are no stacked central configurations in this case.

■

The next proposition brings numerical evidence of the existence of stacked central configurations for $r_4 \in L_{323}$.

PROPOSITION 8. *Assume that r_1, r_2 and r_3 (with positive masses m_1, m_2 and m_3) are at the vertices of an equilateral triangle whose sides have length 1, r_4 and r_5 (with positive masses m_4 and m_5) are located symmetrically with respect to a perpendicular bisector such that $r_4 \in L_{323}$ and $r_5 \in L_{323}'$, as in Figure 2. Then there are positive masses $m_1 = m_2$, $m_4 = m_5$ and m_3 such that the five bodies with these masses form stacked central configurations.*

Proof. For $r_4 \in L_{323}$ and $r_5 \in L_{323}'$ the signs of the coefficients of the matrix A are given by

$$A\Big|_{L_{323} \cup L_{323}'} = \begin{pmatrix} - & - & + \\ - & + & + \\ - & 0 & + \end{pmatrix}.$$

Using the last row of the above matrix we conclude that m_1 and m_4 have the same sign. Unfortunately we cannot obtain more information about the sign of the mass m_3 from the above matrix. Nevertheless, equations (5) and (6) define two planes through the origin in the space (m_1, m_3, m_4). The normal vectors of these planes are, respectively, $n_1 = (a_{11}, a_{12}, a_{13})$ and $n_2 = (a_{21}, a_{22}, a_{23})$. Let $T = (T_1, T_2, T_3) = n_1 \wedge n_2$ be the vector parallel to the straight line defined by the intersection of these two planes. Thus there exist positive masses m_1, m_3 and m_4 as solutions of equations (5) and (6) if and only if the components of the vector T have the same sign. An easy computation shows that

$$T_1 = a_{12}a_{23} - a_{13}a_{22} < 0, \quad T_3 = a_{11}a_{22} - a_{12}a_{21} < 0.$$

In order to see that $T_2 = a_{13}a_{21} - a_{11}a_{23} < 0$, we have verified numerically that the values of T_2 on the curve $L = 0$ between the points P_2 and P_3 are in the interval $(-0.716137, -0.44343)$. The Proposition 8 has been proved.

■

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