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William M. Goldman, François Labourie

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GEODESICS IN MARGULIS SPACETIMES

WILLIAM M. GOLDMAN AND FRANÇOIS LABOURIE

Dedicated to the memory of Dan Rudolph

ABSTRACT. Let M^3 be a Margulis spacetime whose associated complete hyperbolic surface Σ^2 has compact convex core. Generalizing the correspondence between closed geodesics on M^3 and closed geodesics on Σ^2 , we establish an orbit equivalence between recurrent spacelike geodesics on M^3 and recurrent geodesics on Σ^2 . In contrast, no timelike geodesic recurs in neither forward nor backwards time.

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Introduction

A Margulis spacetime is a complete flat affine 3-manifold M^3 with free nonabelian fundamental group Γ . It necessarily carries a unique parallel Lorentz metric. Parallelism classes of timelike geodesics form a noncompact complete hyperbolic surface Σ^2 . This complete hyperbolic surface is naturally associated to the flat 3-manifold M^3 and we regard M^3 as an affine deformation of Σ^2 . This note relates the dynamics of the geodesic flow of the flat affine manifold M^3 to the dynamics of the geodesic flow on the hyperbolic surface Σ^2 .

We restrict to the case that Σ^2 has a compact convex core (that is, Σ^2 has finite type and no cusps). Equivalently, the Fuchsian group Γ_0 corresponding to $\pi_1(\Sigma^2)$ is convex cocompact. In particular Γ_0 is finitely generated and contains no parabolic elements. Under this assumption,

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every free homotopy class of an essential closed curve in Σ^2 contains a unique closed geodesic. Since Σ^2 and M^3 are homotopy-equivalent, free homotopy classes of essential closed curves in M correspond to free homotopy classes of essential closed curves in Σ^2 . Every essential closed curve in M^3 is likewise homotopic to a unique closed geodesic in M^3 .

In her thesis [4, 7], Charette studied the next case of dynamical behavior: geodesics spiralling around closed geodesics both in forward and backward time. She proved bispiralling geodesics in M^3 exist, and correspond to bispiralling geodesics in Σ^2 .

This paper extends the above correspondence between geodesics on Σ^2 and M^3 to recurrent geodesics.

A geodesic (either in Σ^2 or in M^3) is recurrent if and only if it (together with its velocity vector) is recurrent in both directions. These correspond to recurrent points for the corresponding geodesic flows as in Katok-Hasselblatt [15], §3.3. Under our hypotheses on Σ^2 , a geodesic on Σ^2 is recurrent if and only if the corresponding orbit of the geodesic flow is precompact.

Theorem 1. Let M^3 be a Margulis spacetime whose associated complete hyperbolic surface Σ has compact convex core.

- The recurrent part of the geodesic flow for Σ^2 is topologically orbit-equivalent to the recurrent spacelike part of the geodesic flow of M^3 .
- The set of recurrent spacelike geodesics in a Margulis spacetime is the closure of the set of periodic geodesics.
- No timelike geodesic recurs.

A semiconjugacy between these flows was observed by D. Fried [11]. This note is the sequel to [13], which characterizes properness of affine deformations by positivity of a marked Lorentzian length spectrum, the generalized Margulis invariant. A crucial step in the proof that properness implies positivity is the construction of sections of the associated flat affine bundle, called neutralized sections. A further modification of neutralized sections produces an orbit equivalence between recurrent geodesics in Σ and recurrent geodesics in M.

It follows that the set of recurrent spacelike geodesics is a Smale hyperbolic set in TM.

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1. Geodesics on Affine Manifolds

An affinely flat manifold is a smooth manifold with a distinguished atlas of local coordinate systems whose charts map to an affine space \mathbb{E} such that the coordinate changes are restrictions of affine automorphisms of \mathbb{E} . Denote the group of affine automorphisms of \mathbb{E} by Aff(\mathbb{E}). This structure is equivalent to a flat torsionfree affine connection. The affine coordinate atlas globalizes to a developing map

$$\tilde{M} \xrightarrow{\mathsf{dev}} \mathbb{E}$$

where $\tilde{M} \to M$ denotes a universal covering space of M. The coordinate changes globalize to an affine holonomy homomorphism

$$\pi_1(M) \xrightarrow{\rho} \mathsf{Aff}(\mathbb{E})$$

where $\pi_1(M)$ denotes the group of deck transformations of $\tilde{M} \to M$. The developing map is equivariant respecting ρ .

Denote the vector space of translations $\mathbb{E} \to \mathbb{E}$ by \mathbb{V} . The action of \mathbb{V} by translations on \mathbb{E} defines a trivialization of the tangent bundle $TM \cong M \times \mathbb{V}$. In these local coordinate charts, a geodesic is a path

$$p \longmapsto p + tv$$

where $p \in \mathbb{E}$ and $\mathbf{v} \in \mathbb{V}$ is a vector. In terms of the trivialization the geodesic flow is:

$$\mathbb{E} \times \mathbb{V} \xrightarrow{\tilde{\psi}_t} \mathbb{E} \times \mathbb{V}$$
$$(p, \mathbf{v}) \longmapsto (p + t\mathbf{v}, \mathbf{v})$$

for $t \in \mathbb{R}$. Clearly this \mathbb{R} -action commutes with $\mathsf{Aff}(\mathbb{E})$.

Geodesic completeness implies that dev is a diffeomorphism. Thus the universal covering \tilde{M} is affinely isomorphic to affine space \mathbb{E} and $M \cong \mathbb{E}/\Gamma$, where $\Gamma := \rho(\pi_1(M))$ is a discrete group of affine transformations acting properly and freely on \mathbb{E} .

2. Flat Lorentz 3-manifolds

Let $\mathsf{Aff}(\mathbb{E}) \xrightarrow{\mathsf{L}} \mathsf{GL}(\mathbb{V})$ denote the homomorphism given by *linear part*, that is, $\mathsf{L}(\gamma) = A$ where

$$p \xrightarrow{\gamma} \mathsf{L}(\gamma)(p) + u(\gamma).$$

Any $L(\Gamma)$ -invariant nondegenerate inner product \langle , \rangle on \mathbb{V} defines a Γ -invariant flat pseudo-Riemannian structure on \mathbb{E} which descends to $M = \mathbb{E}/\Gamma$. In particular affine manifolds with $L(\Gamma) \subset O(n-1,1)$ are precisely the *flat Lorentzian manifolds*, and the underlying affine structures their Levi-Civita connections.

For this reason we henceforth fix the invariant Lorentzian inner product on \mathbb{V} , and hence the (parallel) flat Lorentzian structure on \mathbb{E} . The group $\mathsf{Isom}(\mathbb{E})$ of Lorentzian isometries is the semidirect product of the group \mathbb{V} of translations of \mathbb{E} with the orthogonal group $\mathsf{SO}(n-1,1)$ of linear isometries. Linear part $\mathsf{Isom}(\mathbb{E}) \xrightarrow{\mathsf{L}} \mathsf{SO}(n-1,1)$ defines the projection homomorphism for the semidirect product. For $l \in \mathbb{R}$, define

$$S_l := \{ v \in V \mid \langle v, v \rangle = l \}.$$

When l > 0, S_l is a Riemannian submanifold of constant curvature $-l^{-2}$ and when l < 0, it is a Lorentzian submanifold of constant curvature l^{-2} . In particular S_{-1} is a disjoint union of two isometrically embedded copies of hyperbolic n-1-space H^{n-1} and S_1 is de Sitter space, a model space of Lorentzian curvature +1.

The subset $\mathsf{T}_l(M)$ consisting of tangent vectors v such that $\langle v, v \rangle = l$ is invariant under the geodesic flow. Indeed, using parallel translation, these bundles trivialize over the universal covering \mathbb{E} :

$$\mathsf{T}_l(\mathbb{E}) \xrightarrow{\cong} \mathbb{E} \times \mathsf{S}_l$$

By Abels-Margulis-Soifer [2, 3], if a discrete group of Lorentz isometries acts properly on Minkowski space \mathbb{E} , and $\mathsf{L}(\Gamma)$ is Zariski dense in $\mathsf{SO}(n-1,1)$, then n=3. Consequently every complete flat Lorentz manifold is a flat Euclidean affine fibration over a complete flat Lorentz 3-manifold. Thus we henceforth restrict to n=3.

Let M^3 be a complete affinely flat 3-manifold. By Fried-Goldman [12], either Γ is solvable or $\mathsf{L} \circ h$ embeds Γ as a discrete subgroup in (a conjugate of) the orthogonal group

$$SO(2,1) \subset GL(3,\mathbb{R}).$$

The cases when Γ is solvable are easily classified (see [12]) and we assume we are in the latter case. In that case, M^3 is a complete flat Lorentz 3-manifold.

In the early 1980's Margulis, answering a question of Milnor [20]), constructed the first examples [17, 18], which are now called *Margulis spacetimes*. Explicit geometric constructions of these manifolds have been given by Drumm [8, 9] and his coauthors [5, 6, 10].

Since the hyperbolic plane H^2 is the symmetric space of SO(2,1), Γ acts properly and discretely on H^2 . Since M^3 is aspherical, its fundamental group $\pi_1(M^3) \cong \Gamma$ is torsionfree, so Γ acts freely as well. Therefore the quotient $H^2/L(\Gamma)$ is a complete hyperbolic surface Σ^2 . Furthermore Σ is noncompact (Mess [19], see also Goldman-Margulis [14] and Labourie [16] for alternate proofs), and every noncompact complete hyperbolic surface occurs for a Margulis spacetime (Drumm [8])

The points of Σ^2 correspond to parallelism classes of timelike geodesics on M^3 as follows. A timelike geodesic is given by an ψ -orbit in

$$\mathsf{T}_{-1}M \cong (\mathbb{E} \times \mathsf{H}^2)/\Gamma$$

Since $\Gamma \xrightarrow{\mathsf{L}} \mathsf{SO}(2,1)$ is a discrete embedding [12], $\mathsf{SO}(2,1)$ acting properly on H^2 implies that Γ acts properly on H^2 . Cartesian projection $\mathbb{E} \times \mathsf{H}^2 \to \mathsf{H}^2$ defines a projection $\mathsf{T}^{-1}M \longrightarrow \mathsf{H}^2/\mathsf{L}(\Gamma) = \Sigma$ which is invariant under the restriction of the geodesic flow ψ to $\mathsf{T}^{-1}M$; and defines an \mathbb{E} -bundle. Its fiber over a fixed future-pointing unit-timelike vector v is the union of geodesics parallel to v . In particular properness of the $\mathsf{L}(\Gamma)$ -action on H^2 implies nonrecurrence of timelike geodesics, the last statement in Theorem 1.

More generally, any $L(\Gamma)$ -invariant subset $\Omega \subset \mathbb{V}$ defines a subset $T_{\Omega}(M) \subset TM$ invariant under the geodesic flow. If Ω is an open set upon which $L(\Gamma)$ acts properly, then the geodesic flow defines a proper \mathbb{R} -action on $T_{\Omega}(M)$. In particular every geodesic whose velocity lies in Ω is properly immersed and is neither positively nor negatively recurrent.

3. From geodesics in Σ^2 to geodesics in M^3

While timelike directions correspond to points of Σ^2 , spacelike directions correspond to geodesics in H^2 . The recurrent geodesics in Σ intimately relate to the recurrent spacelike geodesics on M^3 .

Denote the set of oriented spacelike geodesics in \mathbb{E} by \mathscr{S} . It identifies with the orbit space of the geodesic flow $\tilde{\psi}$ on $T_{+1}\mathbb{E} \cong \mathbb{E} \times S_{+1}$. The natural map $\mathscr{S} \xrightarrow{\Upsilon} S_{+1}$ which associates to a spacelike vector its direction is equivariant respecting $\mathsf{Isom}(\mathbb{E}) \xrightarrow{\mathsf{L}} \mathsf{SO}(2,1)$.

The identity component of SO(2,1) acts simply transitively on the unit tangent bundle UH^2 , and therefore we identify $SO(2,1)^0$ with UH^2 by choosing a basepoint u_0 in UH^2 . Unit-spacelike vectors in S_{+1} correspond to oriented geodesics in H^2 . Explicitly, if $v \in S_{+1}$, then there is a one-parameter subgroup $a(t) \in SO(2,1)$, having v as a fixed vector, and such that $det(v, v^-, v^+) > 0$, where v^+ is an expanding eigenvector of a(t) (for t > 0) and v^- is the contracting eigenvector. Choose a basepoint $v_0 \in S_{+1}$ corresponding to the orbit of u_0 under the geodesic flow on $U\Sigma$. Geodesics in H^2 relate to spacelike directions by an equivariant mapping

$$\begin{array}{ccc}
\mathsf{UH}^2 & \xrightarrow{\tilde{\nu}} \mathsf{S}_{+1} \\
g(u_0) & \longmapsto g(v_0)
\end{array}$$

The unit tangent bundle $U\Sigma$ of Σ identifies with the quotient

$$L(\Gamma)\backslash UH^2 \cong L(\Gamma)\backslash SO(2,1)^0$$

where the geodesic flow ψ corresponds the right-action of a(-t) (see, for example, [13],§1.2).

Observe that a geodesic in Σ^2 is recurrent if and only if the endpoints of any of its lifts to $\tilde{\Sigma} \approx \mathsf{H}^2$ lie in the limit set Λ of $\mathsf{L}(\Gamma)$. If the convex core of Σ^2 is compact, then the union $\mathsf{U}_{\mathsf{rec}}\Sigma$ of recurrent ϕ -orbits is compact.

Lemma 2. There exists an orbit-preserving map

$$\mathsf{U}_{\mathsf{rec}}\Sigma \xrightarrow{\widehat{\mathbf{N}}} \mathsf{T}_{+1}(M)$$

mapping ϕ -orbits injectively to recurrent ψ -orbits.

Proof. The associated flat affine bundle E over $U\Sigma$ associated to the affine deformation Γ is defined as follows. The affine representation of Γ defines a diagonal action of Γ on $\widetilde{U\Sigma} \times \mathbb{E}$

Its total space is the quotient of the product $U\Sigma \times \mathbb{E}$ by the diagonal action of $\pi_1(\Sigma)$

$$\pi_1(\mathsf{U}\Sigma) \longrightarrow \pi_1(\Sigma) \longrightarrow \mathsf{Isom}(\mathbb{E}).$$

Similarly the flat vector bundle V over $U\Sigma$ is the quotient of $\widetilde{U\Sigma}\times \mathbb{V}$ by the diagonal action

$$\pi_1(\mathsf{U}\Sigma) \longrightarrow \pi_1(\Sigma) \longrightarrow \mathsf{Isom}(\mathbb{E}) \xrightarrow{\mathsf{L}} \mathsf{SO}(2,1).$$

According to [13], the neutral section of V is a SO(2, 1)-invariant section which is parallel with respect to the geodesic flow on U Σ . and arises from the graph of the SO(2, 1)-equivariant map

$$\mathsf{U}\tilde{\Sigma}\cong\mathsf{U}\mathsf{H}^2\longrightarrow\mathbb{V}$$

with image S_{+1} , the space of unit-spacelike vectors in \mathbb{V} .

Here is the main construction of [13]. To every section σ of E continuously differentiable along ϕ , associate the function

$$F_{\sigma} := \langle \nabla_{\phi} \sigma, \nu \rangle$$

on $U\Sigma$. (Here the covariant derivative of a section of E along a vector field ϕ in the base is a section of the associated vector bundle V.) Different choices of section σ yield cohomologous functions F_{σ} . (Recall that two functions f_1, f_2 are cohomologous, written $f_1 \sim f_2$, if

$$f_1 - f_2 = \phi g$$

for a function g which is differentiable with respect to the vector field ϕ ([15],§2.2).

Restrict the affine bundle E to $U_{rec}\Sigma$. Lemma 8.4 of [13] guarantees the existence of a neutralized section, that is, a section N of $E|_{U_{rec}\Sigma}$ satisfying

$$\nabla_X N = f \nu$$
,

for some function f.

Although the following lemma is well known, we could not find a proof in the literature. For completeness, we supply a proof in the appendix.

Lemma 3. Let X be a compact space equipped with a flow ϕ . Let $f \in C(X)$, such that for all ϕ -invariant measures μ on X,

$$\int f \, \mathrm{d}\mu \, > \, 0.$$

Then f is cohomologous to a positive function.

Since Γ acts properly, Proposition 8.1 of [13] implies that $\int F_{\sigma}d\mu \neq 0$ for all ϕ -invariant probability measures μ on $\mathsf{U}_{\mathsf{rec}}\Sigma$. Since the set of invariant measures is connected, $\int F_{\sigma}d\mu$ is either positive for all ϕ -invariant probability measures μ on $\mathsf{U}_{\mathsf{rec}}\Sigma$ or negative for all ϕ -invariant probability measures μ on $\mathsf{U}_{\mathsf{rec}}\Sigma$. Conjugating by -I if necessary we may assume that $\int F_{\sigma}d\mu > 0$. Lemma 3 implies $F_{\sigma} + L_{\phi}g > 0$ for some function g. Write

$$\widehat{N} = N + q\nu.$$

 \widehat{N} remains neutralized, and $\nabla_X \widehat{N}$ vanishes nowhere.

Let $\widetilde{\mathsf{U}_{\mathsf{rec}}\Sigma}$ be the preimage of $\mathsf{U}_{\mathsf{rec}}\Sigma$ in UH^2 . Then \widehat{N} determines a Γ -equivariant map

$$\widetilde{\mathsf{U}_{\mathsf{rec}}\Sigma} \overset{\widetilde{N}}{\longrightarrow} \mathbb{E}.$$

Each $\tilde{\phi}$ -orbit injectively maps to a spacelike geodesic. The map

$$\mathbf{U}_{\mathsf{rec}} \Sigma \xrightarrow{\hat{\mathbf{N}}} \left(\mathbb{E} \times \mathsf{S}_{+1} \right) / \Gamma = \mathsf{T}_{+1}(M)
\widehat{\mathbf{N}}(x) := \left[(\widehat{N}(x), \nu(x)) \right].$$

is the desired orbit equivalence $\mathsf{U}_{\mathsf{rec}}\Sigma \longrightarrow \mathsf{T}_{+1}(M)$.

Lemma 4. Any spacelike recurrent geodesic parallel to a geodesic γ in the image of $\hat{\mathbf{N}}$ coincides with γ .

Proof. Let $t \stackrel{g}{\longmapsto} \phi_t(v)$ be an orbit in $\mathsf{U}_{\mathsf{rec}}\Sigma$. A geodesic ξ parallel to $\widehat{\mathbf{N}}(g)$ determines a parallel section u of \mathbb{V} along g. Since g recurs, the resulting parallel section is a bounded invariant parallel section along the closure of g. By the Anosov property, such a section is along ν , and therefore, up to reparametrization, $\gamma = \widehat{\mathbf{N}}(g)$.

Proposition 5. $\widehat{\mathbf{N}}$ is injective and its image is the set of recurrent spacelike geodesics.

Proof. An orbit of the geodesic flow ϕ recurs if and only if the corresponding Γ -orbit in the space $\mathscr S$ of spacelike geodesics in $\mathbb E$ recurs. Similarly a ϕ -orbit in $\mathsf{T}_{+1}(M)$ recurs if and only if the corresponding $\mathsf{L}(\Gamma)$ -orbit in S_{+1} recurs. The map $\mathscr S \xrightarrow{\Upsilon} \mathsf{S}_{+1}$ recording the direction of a spacelike geodesic is L -equivariant. If the Γ -orbit of $g \in \mathscr S$ corresponds to a recurrent spacelike geodesic in M, then the $\mathsf{L}(\Gamma)$ -orbit of $\Upsilon(g)$ corresponds to a recurrent ϕ -orbit in $\mathsf{U}\Sigma$.

 $\hat{\mathbf{N}}$ is injective along orbits of the geodesic flow. Thus it suffices to prove that the restriction of Υ to the subset of Γ -recurrent geodesics in $\mathscr S$ is injective. Since the fibers of Υ are parallelism classes of spacelike geodesics, Lemma 4 implies injectivity of $\hat{\mathbf{N}}$.

Finally let g be a ψ -recurrent point in $\mathsf{T}_{+1}(M)$, corresponding to a spacelike recurrent geodesic γ in M. It corresponds to a recurrent Γ -orbit Γg in \mathscr{S} . Then $\Upsilon(\Gamma g)$ is a recurrent $\mathsf{L}(\Gamma)$ -orbit in S_{+1} , and corresponds to a recurrent ϕ -orbit in $\mathsf{U}\Sigma$. The image of this ϕ -orbit under $\widehat{\mathbf{N}}$ is a spacelike recurrent geodesic in $\mathsf{T}_{+1}(M)$ parallel to γ . Now apply Lemma 4 again to conclude that g lies in the image of $\mathsf{h}N$. \square

The proof of Theorem 1 is complete.

4. Appendix: Cohomology and positive functions

Let X be a smooth manifold equipped with a smooth flow ϕ . A function $g \in C(X)$ is continuously differentiable along ϕ if for each $x \in X$, the function

$$t \mapsto g(\phi_t(x))$$

is a continuously differentiable map $\mathbb{R} \longrightarrow X$. Denote the subspace of C(X) consisting of functions continuously differentiable along ϕ by $C_{\phi}(X)$. For $g \in C_{\phi}(X)$, denote its directional derivative by:

$$\phi(g) := \frac{d}{dt}\Big|_{t=0} g \circ \phi_t.$$

The proof of Lemma 3 will be based on two lemmas.

Lemma 6. Let $f \in C_{\phi}(X)$. For any T > 0, define

$$f_T(x) := \frac{1}{T} \int_0^T f(\phi_s(x)) ds.$$

Then $f \sim f_T$.

Proof. We must show that there exists a function $g \in C_{\phi}(X)$ such that:

$$f_T - f = \phi g.$$

By the fundamental theorem of calculus,

$$f \circ \phi_t = f + \int_0^t (\phi f \circ \phi_s) \, \mathrm{d}s.$$

Writing

$$g = \frac{1}{T} \int_0^T \int_0^t (f \circ \phi_s) \, \mathrm{d}s \, \mathrm{d}t,$$

then

$$f_T - f = \frac{1}{T} \int_0^T (f \circ \phi_t - f) dt$$
$$= \frac{1}{T} \int_0^T \int_0^t \phi(f \circ \phi_s) ds dt$$
$$= \phi g.$$

as desired.

Lemma 7. Assume that for all ϕ -invariant measures μ ,

$$\int f \mathrm{d}\mu > 0.$$

Then $f_T > 0$ for some T > 0.

Proof. Otherwise sequences $\{T_m\}_{m\in\widehat{\mathbb{N}}}$ of positive real numbers and sequences $\{x_m\}_{m\in\widehat{\mathbb{N}}}$ of points in M exist such that

$$f_{T_n}(x_n) \le 0.$$

Using the flow ϕ_t , push forward the (normalized) Lebesgue measure

$$\frac{1}{T_m}\mu_{[0,T_m]}$$

on the interval $[0, T_m]$ to X, to obtain a sequence of probability measures μ_n on X such that

$$\int f \, \mathrm{d}\mu_n \, \leq \, 0.$$

As in [13], §7, a subsequence weakly converges to an ϕ -invariant measure μ for which

$$\int f \, \mathrm{d}\mu \le 0,$$

contradicting our hypotheses.

Proof of Lemma 3. By Lemma 6, $f \sim f_T$ for any T > 0, and Lemma 7 implies that $f_T > 0$ for some T.

References

- [1] Abels, H., Properly discontinuous groups of affine transformations, A survey, Geom. Ded. 87 (2001) 309–333.
- [2] Abels, H., Margulis, G., and Soifer, G., Properly discontinuous groups of affine transformations with orthogonal linear part., C. R. Acad. Sci. Paris Sér. I Math. **324** (1997), no. 3, 253?258.
- [3] Abels, H., Margulis, G., and Soifer, G., The linear part of an affine group acting properly discontinuously and leaving a quadratic form invariant, Geom. Ded. (to appear).
- [4] Charette, V., Proper Actions of Discrete Groups on 2+1 Spacetime, Doctoral Dissertation, University of Maryland (2000).
- [5] Charette, V., Drumm, T., and Goldman, W., Affine deformations of a three-holed sphere. Geom. Topol. 14 (2010), no. 3, 1355–1382,
- [6] Charette, V., and Goldman, W., Affine Schottky groups and crooked tilings. Crystallographic groups and their generalizations (Kortrijk, 1999), 69–97, Contemp. Math., 262, Amer. Math. Soc., Providence, RI, 2000.
- [7] Charette, V., Goldman, W., and Jones, C., Recurrent Geodesics in Flat Lorentz 3-Manifolds, Canad. Math. Bull. 47 (3), (2004), 332–342.
- [8] Drumm, T., *Linear holonomy of Margulis space-times*. J. Differential Geom. 38 (1993), no. 3, 676–690.
- [9] Drumm, T., Fundamental polyhedra for Margulis space-times. Doctoral dissertation, University of Maryland (1990); Topology 31 (1992), no. 4, 677–683.
- [10] Drumm, T. and Goldman, W., Complete flat Lorentz 3-manifolds with free fundamental group. Internat. J. Math. 1 (1990), no. 2, 149–161.
- [11] Fried, D., personal communication to W. Goldman (2002).
- [12] Fried, D. and Goldman, W., Three-dimensional affine crystallographic groups, Adv. Math. 47 (1983), 1–49.
- [13] Goldman, W., Labourie, F., and Margulis, G., Proper affine actions and geodesic flows of hyperbolic surfaces, Ann. Math. 170 No. 3, (2009), 1051–1083.
- [14] Goldman, W. and Margulis, G., Flat Lorentz 3-manifolds and cocompact Fuchsian groups, in "Crystallographic Groups and their Generalizations," Contemp. Math. **262** (2000), 135—146, Amer. Math. Soc.
- [15] Katok, A. and Hasselblatt, B., "Introduction to the modern theory of dynamical systems. With a supplementary chapter by Katok and Leonardo Mendoza." Encyclopedia of Mathematics and its Applications, 54. Cambridge University Press, Cambridge (1995).
- [16] Labourie, F., Fuchsian affine actions of surface groups, J. Diff. Geo. **59** (1), (2001), 15 31.
- [17] Margulis, G., Free properly discontinuous groups of affine transformations, Dokl. Akad. Nauk SSSR **272** (1983), 937–940.
- [18] Margulis, G., Complete affine locally flat manifolds with a free fundamental group, J. Soviet Math. 134 (1987), 129–134.
- [19] Mess, G., Lorentz spacetimes of constant curvature, Geom. Ded. 126, no. 1 (2007), 3-45, in "New techniques in Lorentz manifolds: Proceedings of the BIRS 2004 workshop," (V. Charette and W. Goldman, eds.)

[20] Milnor, J., On fundamental groups of complete affinely flat manifolds, Adv. Math. **25** (1977), 178–187.

Department of Mathematics, University of Maryland, College Park, MD 20742 USA

 $E\text{-}mail\ address{:}\ \mathtt{wmg@math.umd.edu}$

Univ. Paris-Sud, Laboratoire de Mathématiques, Orsay F-91405 Cedex; CNRS, Orsay cedex, F-91405

 $E\text{-}mail\ address: \texttt{francois.labourie@math.u-psud.fr}$