

Penalty Methods for the Hyperbolic System Modelling the Wall-Plasma Interaction in a Tokamak

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Objectives and motivations

Wall-plasma interactions in a tokamak [TAMAIN, PhD thesis 2007]

• A challenge for ITER project : control wall-plasma interactions. \Rightarrow development of a fast solver for numerical simulations.

A first approach [ISOARDI et al, JCP 2010]

- Show that this approach is possible with encouraging results.
- But their penalization cuts the flux term at the plasma-limiter interface which seems to be "hazardous".

A first approach [Isoardi et al, JCP 2010]

 χ : characteristic function of the limiter; η : penalization parameter.

$$\begin{split} \partial_t N + \partial_x \Gamma + \frac{\chi}{\eta} N &= (1 - \chi) S \\ \partial_t \Gamma + (1 - \chi) \partial_x \left(\frac{\Gamma^2}{N} + N \right) + \frac{\chi}{\eta} (\Gamma - M_0 N) &= 0 \\ \text{Initial conditions} : N(0, .) &= N_0 \text{ and } \Gamma(0, .) = \Gamma_0 \end{split}$$

Ę	Some numerical results
	$N(-x) = N(x) \qquad N(0.5-x) = N(0.5+x)$ $\Gamma(-x) = -\Gamma(x) \qquad \Gamma(0.5-x) = -\Gamma(0.5+x)$

Flux is cut in the limiter \Rightarrow we can think that a DIRAC measure may

The original hyperbolic system

N : plasma density ; Γ : particle flux ; $M = \frac{\Gamma}{N}$: plasma velocity

 $(t,x) \in \mathbb{R}^+ \times [-L,L]$ $\partial_t N + \partial_x \Gamma = S_N$

 $\partial_t \Gamma + \partial_x \left(\frac{\Gamma^2}{N} + N \right) = S_{\Gamma}$

Boundary conditions : M(-L) = -1 and M(L) = 1 (Bohm condition) Initial conditions : $N(0, .) = N_0$ and $\Gamma(0, .) = \Gamma_0$

Numerical tests : 2nd order finite volume scheme (VF ROE nev [GALLOUËT et al, Computers & fluids 2003] with entropy correction, MUSCL (with slope limiter) and RK2 (HEUN) time discretization.

Well-posedness issue

The waves speeds for the hyperbolic system are the eigenvalues : M-1and M+1. From the boundary conditions |M| = 1, we infer that there is no ingoing wave (one characteristic wave and one outgoing wave at $x = \pm L$). So, we can't impose any boundary condition. Thus, we change the boundary conditions to get one ingoing wave :

 $(t,x) \in \mathbb{R}^+ \times [-L,L]$

Problem of meaning for the term $(1 - \chi)\partial_x \left(\frac{\Gamma^2}{N} + N\right)$: For example, there is no solution piecewise \mathcal{C}^1 such that $\frac{\Gamma^2}{N} + N$ is continuous on x = L.

So this first approach was not completely satisfactory. This is confirmed by numerical tests (see next page).

A new penalty method

For the semi-linear case, [FORNET and GUÈS, DCDS, 2009] proposed a penalization method without any boundary layer. Though our system is quasi-linear, we apply this method to our case and we obtain interesting results. The penalized system is :

 $\partial_t N + \partial_x \Gamma = (1 - \chi) S_N$ $\langle \Gamma^2 \rangle \rangle \sim \langle \Gamma \rangle$

appear on the plasma-limiter interface. Computations stopped when $\max_{i \in \{1,...,J\}} (|M_i^n|) > 10 \text{ (normally } M \le 1) \text{ with } \eta = 10^{-3}.$ Mesh convergence study :



M versus x with $\Delta x = 1.952 \cdot 10^{-4}$ (left graph at t = 0.004107) and with $\Delta x = 4.88 \cdot 10^{-5}$ (right graph, at t = 0.0015834).

When the resolution increases, the peak is nearer and nearer to the interface. This phenomenon is observed when $M_0 = 10(0.5 - x)$ and when $M_0 = 1$ but not when $M_0 = 0$.

Asymptotic analysis

Change of unknown from both sides of the interface : $\left(\begin{array}{c} ln(N(t,x)) \\ M(t,x) - M_0 \end{array} \right)$ $\mathbf{U}_{\eta}^{\pm}(t,x) = \Big($ $\partial_t \mathbf{U}_{\eta}^{\pm} + \mathbf{f}'(\mathbf{U}_{\eta}^{\pm})\partial_x \mathbf{U}_{\eta}^{\pm} + \frac{\chi}{M_0 \eta} P \mathbf{U}_{\eta}^{\pm} = \mathbf{S}$

P =projection matrix.

 $\partial_t N + \partial_x \Gamma = S_N$ $\partial_t \Gamma + \partial_x \left(\frac{\Gamma^2}{N} + N \right) = S_{\Gamma}$

Boundary conditions : $M(-L) = -1 + \epsilon$ and $M(L) = 1 - \epsilon$, $\epsilon > 0$ Initial conditions : $N(0, .) = N_0$ and $\Gamma(0, .) = \Gamma_0$

On each boundary we have only one condition to impose (this was not the case in the previous penalization). Further numerical results obtained with $\epsilon = 0.1$ and $M_0 = 1 - \epsilon = 0.9$.

$$\partial_t \Gamma + \partial_x \left(\frac{1}{N} + N \right) + \frac{\chi}{\eta} \left(\frac{1}{M_0} - N \right) = (1 - \chi) S_{\Gamma}$$

Initial conditions : $N(0, .) = N_0$ and $\Gamma(0, .) = \Gamma_0$

"Theorem" (WKB analysis). For any $k \in \mathbb{N}$, there exists an approximate solution of the penalized system

$$N_{\eta,app}(t,x) = \sum_{n=0}^{k} \eta^n \mathbf{N}^n(t,x), \qquad \Gamma_{\eta,app}(t,x) = \sum_{n=0}^{k} \eta^n \mathbf{\Gamma}^n(t,x)$$

satisfying the equations up to an error in η^k for arbitrarily large

Formal asymptotic expansion of a continuous solution of the form : $\mathbf{U}_{\eta}^{\pm}(t,x) = \sum_{n=0}^{+\infty} \eta^{n} \mathbf{U}^{n,\pm}(t,x)$ Substituting the expansion and classifying gives : • Inside the plasma : $\sum_{n=0}^{\infty} \eta^n \left(\partial_t \mathbf{U}^{n,-} + \dots \right) = \mathbf{S}$ • In the limiter set : $\frac{\eta^{-1}}{M_0} P \mathbf{U}^{0,+} + \sum_{n=0}^{\infty} \eta^n \left(\partial_t \mathbf{U}^{n,+} + \dots + \frac{1}{M_0} P \mathbf{U}^{n+1,+} \right) = \mathbf{S}$ Construction by induction with continuous connection at the interface $(\mathbf{U}^{n,-}(t,0) = \mathbf{U}^{n,+}(t,0))$: • $P\mathbf{U}^{0,+} = \mathbf{0}$ • Then, construction of $\mathbf{U}^{0,-}$, $(Id - P)\mathbf{U}^{0,+}$ and $P\mathbf{U}^{1,+}$.

• ... $\mathbf{U}^{n,-}$, $(Id - P)\mathbf{U}^{n,+}$ and $P\mathbf{U}^{n+1,+}$

If there is a boundary layer : presence of terms in $\underline{\mathbf{U}}^{n,\pm}(t,x,\frac{x}{n^{\alpha}})$. \Rightarrow Probably no boundary layer (error analysis needs to be done).





Conclusion and perspectives

- The first penalization does not work correctly. Besides the system considered is ill-posed.
- After a modification of the system, we experiment two penalty methods which give similar results.
- Penalizing both N = 0 and $M = M_0$ creates a boundary layer.

Optimal convergence rate for N and $\Gamma : \mathcal{O}(\eta)$ Non optimal rate for the x-derivatives of N in the L^2 error \Rightarrow boundary layer or artefact ?

• These results need to be extended to a more complex model for edge simulation of ITER (3-space dimensions, energy equation...).

Bibliography

P. TAMAIN, Etude des flux de matière dans le plasma de bord des tokamaks, alimentation, transport et turbulence, PhD thesis, Université de Provence, 2007.

L. Isoardi, G. Chiavassa, G. Ciraolo, P. Haldenwang, E. SERRE, Ph. GHENDRIH, Y. SARAZIN, F. SCHWANDER, P. TAMAIN, Penalization modeling of a limiter in the Tokamak edge plasma, Journal of Computational Physics, 229(6): 2220 - 2235, 2010. B. FORNET, O. GUÈS, Penalization approach of semi-linear symmetric hyperbolic problems with dissipative boundary conditions, Discrete and Continuous Dynamical Systems, 23(3): 827 - 845, 2009. T. GALLOUËT, J-M. HÉRARD, N. SEGUIN, Some approximate Godunov schemes to compute shallow-water equations with topogra*phy*, Computers & Fluids, 32(4) : 479 - 513, 2003.