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## To cite this version:

Philippe Laurençot. Non-diffusive large time behaviour for a degenerate viscous HamiltonJacobi equation. Communications in Partial Differential Equations, 2009, 34 (1-3), pp.281-304. 10.1080/03605300902793808 . hal-00627281v2

HAL Id: hal-00627281
https://hal.science/hal-00627281v2
Submitted on 10 Mar 2018

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# Non-diffusive large time behaviour for a degenerate viscous Hamilton-Jacobi equation 

Philippe Laurençot*

March 10, 2018


#### Abstract

The convergence to non-diffusive self-similar solutions is investigated for non-negative solutions to the Cauchy problem $\partial_{t} u=\Delta_{p} u+|\nabla u|^{q}$ when the initial data converge to zero at infinity. Sufficient conditions on the exponents $p>2$ and $q>1$ are given that guarantee that the diffusion becomes negligible for large times and the $L^{\infty}$-norm of $u(t)$ converges to a positive value as $t \rightarrow \infty$.


## 1 Introduction

The quasilinear degenerate parabolic equation

$$
\begin{equation*}
\partial_{t} u=\Delta_{p} u+|\nabla u|^{q}, \quad(t, x) \in Q_{\infty}:=(0, \infty) \times \mathbb{R}^{N} \tag{1.1}
\end{equation*}
$$

includes two competing mechanisms acting on the space variable $x$, a degenerate diffusion $\Delta_{p} u$ involving the $p$-Laplacian operator defined by

$$
\Delta_{p} u:=\operatorname{div}\left(|\nabla u|^{p-2} \nabla u\right), \quad p>2
$$

and a source term $|\nabla u|^{q}, q>1$, depending solely on the gradient of $u$. The aim of this work is to identify a range of the parameters $p$ and $q$ for which the large time behaviour of non-negative solutions to (1.1) is dominated by the source term. More precisely, we consider the Cauchy problem and supplement (1.1) with the initial condition

$$
\begin{equation*}
u(0)=u_{0} \geq 0, \quad x \in \mathbb{R}^{N} \tag{1.2}
\end{equation*}
$$

[^0]Throughout the paper, the initial condition $u_{0}$ is assumed to fulfill

$$
\begin{equation*}
u_{0} \in \mathcal{C}_{0}\left(\mathbb{R}^{N}\right) \cap W^{1, \infty}\left(\mathbb{R}^{N}\right), \quad u_{0} \geq 0, \quad u_{0} \not \equiv 0 \tag{1.3}
\end{equation*}
$$

where

$$
\mathcal{C}_{0}\left(\mathbb{R}^{N}\right):=\left\{w \in \mathcal{B C}\left(\mathbb{R}^{N}\right): \lim _{R \rightarrow \infty} \sup _{\{|x| \geq R\}}\{|w(x)|\}=0\right\},
$$

and $\mathcal{B C}\left(\mathbb{R}^{N}\right):=\mathcal{C}\left(\mathbb{R}^{N}\right) \cap L^{\infty}\left(\mathbb{R}^{N}\right)$.
For such an initial condition, the Cauchy problem (1.1), (1.2) has a unique non-negative (viscosity) solution $u \in \mathcal{B C}\left([0, \infty) \times \mathbb{R}^{N}\right.$ ) (see Proposition 2.1 below). Moreover, $t \longmapsto$ $\|u(t)\|_{\infty}$ is a non-increasing function and has a limit $M_{\infty} \in\left[0,\left\|u_{0}\right\|_{\infty}\right]$ as $t \rightarrow \infty$. Our main result is then the following:

Theorem 1.1 Assume that $p>2$ and $q \in(1, p)$. Consider a non-negative function $u_{0}$ satisfying (1.3) and let $u$ be the corresponding (viscosity) solution to (1.1), (1.2). Assume further that

$$
\begin{equation*}
M_{\infty}:=\lim _{t \rightarrow \infty}\|u(t)\|_{\infty}>0 \tag{1.4}
\end{equation*}
$$

Then

$$
\begin{equation*}
\lim _{t \rightarrow \infty}\left\|u(t)-h_{\infty}(t)\right\|_{\infty}=0 \tag{1.5}
\end{equation*}
$$

where $h_{\infty}$ is given by

$$
\begin{equation*}
h_{\infty}(t, x):=H_{\infty}\left(\frac{x}{t^{1 / q}}\right) \quad \text { and } \quad H_{\infty}(x):=\left(M_{\infty}-\gamma_{q}|x|^{q /(q-1)}\right)_{+} \tag{1.6}
\end{equation*}
$$

for $(t, x) \in Q_{\infty}$ and $\gamma_{q}:=(q-1) q^{-q /(q-1)}$.
Here and below, $r_{+}:=\max \{r, 0\}$ denotes the positive part of the real number $r$.
The convergence (1.5) clearly indicates that the large time behaviour of non-negative solutions to (1.1), (1.2) fulfilling the condition (1.4) is governed by the gradient source term. Indeed, $h_{\infty}$ is actually a self-similar solution to the Hamilton-Jacobi equation

$$
\begin{equation*}
\partial_{t} h=|\nabla h|^{q}, \quad(t, x) \in Q_{\infty}, \tag{1.7}
\end{equation*}
$$

and an alternative formula for $h_{\infty}$ reads

$$
\begin{equation*}
h_{\infty}(t, x)=\sup _{y \in \mathbb{R}^{N}}\left\{M_{\infty} \mathbf{1}_{\{0\}}(y)-\gamma_{q} \frac{|x-y|^{q /(q-1)}}{t^{1 /(q-1)}}\right\} \tag{1.8}
\end{equation*}
$$

for $(t, x) \in[0, \infty) \times \mathbb{R}^{N}, \mathbf{1}_{\{0\}}$ denoting the indicator function of the singleton set $\{0\}$. The formula (1.8) is the well-known Hopf-Lax-Oleinik representation formula for viscosity solutions to (1.7) (see, e.g., [10, Chapter 3]) and $h_{\infty}$ turns out to be the unique viscosity solution in $\mathcal{B U C}\left(Q_{\infty}\right)$ to (1.7) with the bounded and upper semicontinuous initial condition $h_{\infty}(0, x)=1_{\{0\}}(x)$ for $x \in \mathbb{R}^{N}$ [24].

Remark 1.2 The convergence (1.5) also holds true for the viscosity solution to the HamiltonJacobi equation (1.7) with a non-negative initial condition $u_{0} \in \mathcal{C}_{0}\left(\mathbb{R}^{N}\right)$ but with $\left\|u_{0}\right\|_{\infty}$ instead of $M_{\infty}$ in the formula (1.6) giving $H_{\infty}$. For (1.1), (1.2), the constant $M_{\infty}$ takes into account that, though negligible for large times, the diffusion erodes the supremum of $u$ during the time evolution.

For $p=2$, Theorem 1.1 is also valid and is proved in [7], the proof relying on a rescaling technique: The crucial step is then to identify the possible limits of the rescaled sequence and this is done by an extensive use of the Hopf-Lax-Oleinik representation formula. The proof we perform here is of a completely different nature and relies on the relaxed halflimits method introduced in [3]. A similar approach has been used in [22] and [23] to investigate the large time behaviour of solutions to first-order Hamilton-Jacobi equations $\partial_{t} w+H(x, \nabla w)=0$ in $Q_{\infty}$. It has also been used in [20] to study the convergence to nondiffusive localized self-similar patterns for non-negative and compactly supported solutions to $\partial_{t} w-\Delta_{p} w+|\nabla w|^{q}=0$ in $Q_{\infty}$ when $p>2$ and $q \in(1, p-1)$.

In order to apply Theorem 1.1, one should check whether there are non-negative solutions to (1.1), (1.2) for which (1.4) holds true. The next result provides sufficient conditions for (1.4) to be fulfilled.

Theorem 1.3 Assume that $p>2$ and $q>1$. Consider a non-negative function $u_{0}$ satisfying (1.3) and let $u$ be the corresponding solution to (1.1), (1.2). Introducing

$$
\begin{equation*}
q_{\star}:=p-\frac{N}{N+1}, \tag{1.9}
\end{equation*}
$$

then $u$ fulfills (1.4) if
(a) either $q \in\left(1, q_{\star}\right]$,
(b) or $q \in\left(q_{\star}, p\right), u_{0} \in W^{2, \infty}\left(\mathbb{R}^{N}\right)$, and

$$
\begin{equation*}
\left\|u_{0}\right\|_{\infty}>\kappa_{0}\left|\inf _{y \in \mathbb{R}^{N}}\left\{\Delta_{p} u_{0}(y)\right\}\right|^{(p-q) / q} . \tag{1.10}
\end{equation*}
$$

for some $\kappa_{0}>0$ which depends only on $N$, $p$, and $q$.
A similar result is already available for $p=2$ and has been established in [7, 12]. It is also shown in $[7,12,18]$ that, still for $p=2,\|u(t)\|_{\infty} \rightarrow 0$ as $t \rightarrow \infty$ if either $q \geq p$ or $q \in\left(q_{\star}, p\right)$ and a suitable norm of $u_{0}$ is small enough. A similar result is likely to be true also for $p>2$, so that different large time behaviours are only to be found in the intermediate range $q \in\left(q_{\star}, p\right)$.
The proof of Theorem 1.3 for $q \in(p-1, p)$ and $p>2$ borrows some steps from the case $p=2$. However, it relies on semiconvexity estimates for solutions to (1.1), (1.2) which seem to be new for $p>2$ and $q \in(1, p)$ and are stated now.

Proposition 1.4 Assume that $p>2$ and $q \in(1, p]$. Let $u$ be the viscosity solution to (1.1), (1.2) with initial condition $u_{0} \in \mathcal{B Z C}\left(\mathbb{R}^{N}\right)$ (that is, $u_{0} \in \mathcal{B C}\left(\mathbb{R}^{N}\right)$ and is uniformly continuous in $\left.\mathbb{R}^{N}\right)$. Then $\nabla u(t)$ belongs to $L^{\infty}\left(\mathbb{R}^{N}\right)$ for each $t>0$ and there is $\kappa_{1}>0$ depending only on $N$, $p$, and $q$ such that

$$
\begin{equation*}
\Delta_{p} u(t, x) \geq-\kappa_{1}\|u(s)\|_{\infty}^{(p-q) / q}(t-s)^{-p / q}, \quad t>s \geq 0 \tag{1.11}
\end{equation*}
$$

in the sense of distributions. In addition, if $u_{0} \in W^{1, \infty}\left(\mathbb{R}^{N}\right)$, there holds

$$
\begin{equation*}
\Delta_{p} u(t, x) \geq-\frac{N(p-1)}{q(q-1)} \frac{\left\|\nabla u_{0}\right\|_{\infty}^{p-q}}{t} \tag{1.12}
\end{equation*}
$$

for $t>0$ in the sense of distributions.
The proof of Proposition 1.4 relies on the comparison principle combined with a gradient estimate established in [6].
Similar semiconvexity estimates for solutions to (1.1), (1.2) have already been obtained in [14] and [21, Lemma 5.1] for $p=q=2$, in [7, Proposition 3.2] for $p=2$ and $q \in(1,2]$, and in [9, Theorem 1] for $p=q>2$. We extend these results to the range $p>2$ and $q \in(1, p]$. As we shall see below, the estimate (1.11) plays an important role in the proof of Theorem 1.3 and is also helpful to construct a subsolution in the proof of Theorem 1.1.
Let us finally emphasize that the validity of Proposition 1.4 is not restricted to non-negative solutions and that the solutions to the Hamilton-Jacobi equation (1.7) also enjoy the semiconvexity estimates (1.11) and (1.12). These two estimates thus stem from the reaction term $|\nabla u|^{q}$ and not from the diffusion.

In the next section, we recall the well-posedness of $(1.1),(1.2)$ in $\mathcal{B U C}\left(\mathbb{R}^{N}\right)$, as well as some properties of the solutions established in [6]. We also show the finite speed of propagation of the support for non-negative compactly supported initial data. Section 3 is devoted to the proof of the semiconvexity estimates (Proposition 1.4) and Section 4 to that of Theorem 1.1. Theorem 1.3 is shown in the last section, its proof combining arguments of $[7,12,19]$ used to established analogous results when $p=2$.

Throughout the paper, $C$ and $C_{i}, i \geq 1$, denote positive constants depending only on $p, q$, and $N$. Dependence upon additional parameters will be indicated explicitly. Also, $\mathcal{M}_{N}(\mathbb{R})$ denotes the space of real-valued $N \times N$ matrices and $\delta_{i j}=1$ if $i=j$ and $\delta_{i j}=0$ if $i \neq j$, $1 \leq i, j \leq N$. Given a matrix $A=\left(a_{i j}\right) \in \mathcal{M}_{N}(\mathbb{R}), \operatorname{tr}(A)$ denotes its trace and is given by $\operatorname{tr}(A):=\sum a_{i i}$.

## 2 Preliminary results

Let us first recall the well-posedness (in the framework of viscosity solutions) of (1.1), (1.2), together with some properties of the solutions established in [6].

Proposition 2.1 Consider a non-negative initial condition $u_{0} \in \mathcal{B U C}\left(\mathbb{R}^{N}\right)$. There is a unique non-negative viscosity solution $u \in \mathcal{B C}\left([0, \infty) \times \mathbb{R}^{N}\right)$ to (1.1), (1.2) such that

$$
\begin{gather*}
0 \leq u(t, x) \leq\left\|u_{0}\right\|_{\infty}, \quad(t, x) \in Q_{\infty},  \tag{2.1}\\
\|\nabla u(t)\|_{\infty} \leq \min \left\{C_{1}\|u(s)\|_{\infty}^{1 / q}(t-s)^{-1 / q},\|\nabla u(s)\|_{\infty}\right\}, \tag{2.2}
\end{gather*}
$$

and

$$
\begin{equation*}
\int_{\mathbb{R}^{N}}(u(t, x)-u(s, x)) \vartheta(x) d x+\int_{s}^{t} \int_{\mathbb{R}^{N}}\left(|\nabla u|^{p-2} \nabla u \cdot \nabla \vartheta-|\nabla u|^{q} \vartheta\right) d x d \tau=0 \tag{2.3}
\end{equation*}
$$

for $t>s \geq 0$ and $\vartheta \in \mathcal{C}_{0}^{\infty}\left(\mathbb{R}^{N}\right)$. In addition, $t \longmapsto\|u(t)\|_{\infty}$ is a non-increasing function.
Proof. We put $\tilde{u}_{0}:=\left\|u_{0}\right\|_{\infty}-u_{0}$. As $\tilde{u}_{0}$ is a non-negative function in $\mathcal{B U C}\left(\mathbb{R}^{N}\right)$, it follows from [6, Theorem 1.1] that there is a unique non-negative viscosity solution $\tilde{u}$ to

$$
\begin{equation*}
\partial_{t} \tilde{u}-\Delta_{p} \tilde{u}+|\nabla \tilde{u}|^{q}=0 \quad, \quad(t, x) \in Q_{\infty}, \tag{2.4}
\end{equation*}
$$

with initial condition $\tilde{u}(0, x)=\tilde{u}_{0}(x)$ for $x \in \mathbb{R}^{N}$. It also satisfies $0 \leq \tilde{u}(t, x) \leq\left\|u_{0}\right\|_{\infty}$ and

$$
\int_{\mathbb{R}^{N}}(\tilde{u}(t, x)-\tilde{u}(s, x)) \vartheta(x) d x+\int_{s}^{t} \int_{\mathbb{R}^{N}}\left(|\nabla \tilde{u}|^{p-2} \nabla \tilde{u} \cdot \nabla \vartheta+|\nabla \tilde{u}|^{q} \vartheta\right) d x d \tau=0
$$

for $t>s \geq 0, x \in \mathbb{R}^{N}$, and $\vartheta \in \mathcal{C}_{0}^{\infty}\left(\mathbb{R}^{N}\right)$. In addition, $\nabla \tilde{u}(t)$ belongs to $L^{\infty}\left(\mathbb{R}^{N}\right)$ for each $t>0$ and

$$
\|\nabla \tilde{u}(t)\|_{\infty} \leq C_{1}\left\|\tilde{u}_{0}\right\|_{\infty}^{1 / q} t^{-1 / q}
$$

by [6, Lemma 4.1]. Setting $u:=\left\|u_{0}\right\|_{\infty}-\tilde{u}$, we readily deduce from the properties of $\tilde{u}$ that $u$ is a non-negative viscosity solution to (1.1), (1.2) satisfying (2.1) and (2.3). Also, $\nabla u(t)$ belongs to $L^{\infty}\left(\mathbb{R}^{N}\right)$ for each $t>0$. The uniqueness and the time monotonicity of $\|u\|_{\infty}$ then both follow from the comparison principle, see [8] or [11, Theorem 2.1]. Finally, given $s \geq 0$, $(t, x) \mapsto\|u(s)\|_{\infty}-u(t+s, x)$ is the unique non-negative viscosity solution to the Cauchy problem (2.4) with initial condition $x \mapsto\|u(s)\|_{\infty}-u(s, x)$ and we infer from [6, Lemma 4.1] that

$$
\|\nabla u(t+s)\|_{\infty} \leq C_{1}\| \| u(s)\left\|_{\infty}-u(s)\right\|_{\infty}^{1 / q} t^{-1 / q} \leq C_{1}\|u(s)\|_{\infty}^{1 / q} t^{-1 / q}
$$

for $t>0$, whence (2.2).
We next turn to the propagation of the support of non-negative solutions to (1.1), (1.2) with non-negative compactly supported initial data.

Proposition 2.2 Consider a non-negative solution u to (1.1), (1.2) with an initial condition $u_{0}$ satisfying (1.3). Assume further that $u_{0}$ is compactly supported in a ball $B\left(0, R_{0}\right)$ of $\mathbb{R}^{N}$ for some $R_{0}>0$. Then $u(t)$ is compactly supported for each $t \geq 0$.

Proof. We argue by comparison with travelling wave solutions. By [13, Application 9.4], there is a travelling wave solution $w$ to the convection-diffusion equation

$$
\begin{equation*}
\partial_{t} w-\partial_{1}^{2}\left(w^{p-1}\right)+\partial_{1}\left(w^{q}\right)=0, \quad\left(t, x_{1}\right) \in(0, \infty) \times \mathbb{R}, \tag{2.5}
\end{equation*}
$$

with wave speed unity. It is given by $w\left(t, x_{1}\right)=f\left(x_{1}-t\right)$ for $\left(t, x_{1}\right) \in(0, \infty) \times \mathbb{R}$, the function $f$ being implicitly defined by

$$
(p-1) \int_{0}^{f(y)} \frac{z^{p-3}}{1-z^{q-1}} d z=(-y)_{+}, \quad y \in \mathbb{R}
$$

In particular, $f$ satisfies $f(y)=0$ if $y>0$ and $f(y) \rightarrow 1$ as $y \rightarrow-\infty$. Introducing

$$
F(y):=\int_{y}^{\infty} f(z) d z, \quad y \in \mathbb{R}
$$

the properties of $f$ ensure that $F$ is a decreasing function on $(-\infty, 0)$ with $F(y)=0$ if $y>0$, $F(y) \leq|y|$ if $y<0$, and $F(y) \rightarrow \infty$ as $y \rightarrow-\infty$. There is therefore a unique $\mu \in(-\infty, 0)$ such that $F\left(R_{0}+\mu\right)=\left\|u_{0}\right\|_{\infty}$. In addition, it readily follows from (2.5) and the invariance by translation of (1.1) that $W_{\mu}(t, x):=F\left(x_{1}+\mu-t\right)$ is a travelling wave solution to (1.1). Now, $u$ and $W_{\mu}$ are both solutions to (1.1) in $(0, \infty) \times H_{+}$, the half-space $H_{+}$being defined by $H_{+}:=\left\{x \in \mathbb{R}^{N}: x_{1}>R_{0}\right\}$. Owing to the monotonicity of $F$, the bound $0 \leq f \leq 1$, and (2.1), we have also

$$
u_{0}(x)-W_{\mu}(0, y)=0-W_{\mu}(0, y) \leq W_{\mu}(0, x)-W_{\mu}(0, y) \leq|x-y|
$$

for $x \in H_{+}$and $y \in H_{+}$,

$$
\begin{aligned}
u(t, x)-W_{\mu}(t, y) & \leq\left\|u_{0}\right\|_{\infty}-W_{\mu}(t, y) \\
& \leq F\left(R_{0}+\mu-t\right)-W_{\mu}(t, y)=W_{\mu}(t, x)-W_{\mu}(t, y) \leq|x-y|
\end{aligned}
$$

for $t>0, x \in \partial H_{+}, y \in H_{+}$, and

$$
u(t, x)-W_{\mu}(t, y) \leq\left\|u_{0}\right\|_{\infty}-F\left(R_{0}+\mu-t\right) \leq 0
$$

for $t>0, x \in H_{+}, y \in \partial H_{+}$. We are then in a position to use the comparison principle stated in [11, Theorem 2.1] to conclude that $u(t, x) \leq W_{\mu}(t, x)$ for $(t, x) \in(0, \infty) \times H_{+}$. Consequently, $u(t, x) \leq F\left(x_{1}+\mu-t\right)=0$ if $t \geq 0$ and $x_{1}>\max \left\{R_{0}, t-\mu\right\}$, and the rotational invariance of (1.1) allows us to conclude that $u(t, x)=0$ for $t \geq 0$ and $|x|>$ $\max \left\{R_{0}, t-\mu\right\}$.

We finally recall the convergence to self-similar solutions for non-negative and compactly supported solutions to the $p$-Laplacian equation [17]

$$
\begin{equation*}
\partial_{t} \varphi=\Delta_{p} \varphi, \quad(t, x) \in Q_{\infty} \tag{2.6}
\end{equation*}
$$

Proposition 2.3 Let $\varphi_{0}$ be a non-negative and compactly supported function in $L^{1}\left(\mathbb{R}^{N}\right)$ and $\varphi$ denote the unique weak solution to (2.6) with initial condition $\varphi_{0}$. Then

$$
\begin{equation*}
\lim _{t \rightarrow \infty} t^{(N(r-1)) /(r(N(p-2)+p))}\left\|\varphi(t)-\mathcal{B}_{\left\|\varphi_{0}\right\|_{1}}(t)\right\|_{r}=0 \quad \text { for } \quad r \in[1, \infty] \tag{2.7}
\end{equation*}
$$

where $\mathcal{B}_{L}$ denotes the Barenblatt solution to (2.6) given by

$$
\begin{aligned}
\mathcal{B}_{L}(t, x) & :=t^{-N /(N(p-2)+p)} b_{L}\left(x t^{-1 /(N(p-2)+p)}\right) \\
b_{L}(x) & :=\left(C_{2} L^{(p(p-2)) /((p-1)(N(p-2)+p))}-C_{3}|x|^{p /(p-1)}\right)_{+}^{(p-1) /(p-2)}
\end{aligned}
$$

for $(t, x) \in Q_{\infty}$ and $L>0$.
The convergence (2.7) is proved in [17, Theorem 2] for $r=\infty$. As $\varphi_{0}$ is compactly supported, so is $\varphi(t)$ for each $t>0$ and the support of $\varphi(t)$ is included in $B\left(0, C_{4}\left(\varphi_{0}\right) t^{1 /(N(p-2)+p)}\right)$ for $t \geq 1$ [17, Proposition 2.2]. Combining this property with [17, Theorem 2] readily provide the convergence (2.7) for all $r \in[1, \infty)$.

## 3 Semiconvexity

In this section, we prove the semiconvexity estimates (1.11) and (1.12). To this end, we would like to derive an equation for $\Delta_{p} u$ to which we could apply the comparison principle. The poor regularity of $u$ however does not allow to perform directly such a computation and an approximation procedure is needed. As a first step, we report the following result:

Lemma 3.1 Let $a$ and $b$ be two non-negative function in $\mathcal{C}^{\infty}([0, \infty))$ satisfying

$$
\begin{align*}
& a(r)>0, \quad a^{\prime}(r)>0, \quad a^{\prime}(r) b^{\prime}(r)-a(r) b^{\prime \prime}(r)>0,  \tag{3.1}\\
& c(r):=2\left(\frac{b^{\prime}}{a}\right)(r)+\frac{4 r\left(a b^{\prime \prime}-a^{\prime} b^{\prime}\right)(r)}{a^{2}(r)+2 r a(r) a^{\prime}(r)} \geq 0 . \tag{3.2}
\end{align*}
$$

Consider a classical solution $v$ to

$$
\begin{equation*}
\partial_{t} v-\operatorname{div}\left(a\left(|\nabla v|^{2}\right) \nabla v\right)=b\left(|\nabla v|^{2}\right), \quad(t, x) \in Q_{\infty} \tag{3.3}
\end{equation*}
$$

and put

$$
w:=\operatorname{div}\left(a\left(|\nabla v|^{2}\right) \nabla v\right) \quad \text { and } \quad z_{i}:=a\left(|\nabla v|^{2}\right) \partial_{i} v
$$

for $i \in\{1, \ldots, N\}$. Then

$$
\begin{equation*}
\partial_{t} w-\mathcal{L} w-\mathcal{V} \cdot \nabla w-\frac{c\left(|\nabla v|^{2}\right)}{N} w^{2} \geq 0 \quad \text { in } \quad Q_{\infty} \tag{3.4}
\end{equation*}
$$

where

$$
\begin{aligned}
& \mathcal{L} w:=\sum_{i, j} \partial_{i}\left(a\left(|\nabla v|^{2}\right) E_{i j} \partial_{j} w\right), \quad \mathcal{V}:=2 b^{\prime}\left(|\nabla v|^{2}\right) \nabla v, \\
& E_{i j}:=\delta_{i j}+2 \frac{a^{\prime}}{a}\left(|\nabla v|^{2}\right) \partial_{i} v \partial_{j} v, \quad 1 \leq i, j \leq N .
\end{aligned}
$$

The proof of Lemma 3.1 borrows some steps from the proof of [9, Theorem 1] for $p=q>2$ but requires additional arguments to handle the term coming from the fact that $q \neq p$. In particular, we recall the following elementary result which will be helpful to estimate this term.

Lemma 3.2 Let $A$ and $B$ be two symmetric matrices in $\mathcal{M}_{N}(\mathbb{R})$ and put $M:=A B A$. Then $M$ is a symmetric matrix in $\mathcal{M}_{N}(\mathbb{R})$ and

$$
\begin{equation*}
|M X|^{2} \leq \operatorname{tr}\left(M^{2}\right)|X|^{2} \quad \text { for } \quad X \in \mathbb{R}^{N} \tag{3.5}
\end{equation*}
$$

Proof of Lemma 3.1. We first note that

$$
\begin{align*}
\partial_{j} z_{i} & =a\left(|\nabla v|^{2}\right) \sum_{k} E_{i k} \partial_{k} \partial_{j} v  \tag{3.6}\\
\partial_{t} z_{i} & =a\left(|\nabla v|^{2}\right) \sum_{k} E_{i k} \partial_{k} \partial_{t} v \tag{3.7}
\end{align*}
$$

for $1 \leq i, j \leq N$. According to the definition of $w$, we infer from (3.3), (3.6), and (3.7) that

$$
\begin{aligned}
\partial_{t} w & =\sum_{i, k} \partial_{i}\left(a\left(|\nabla v|^{2}\right) E_{i k} \partial_{k} \partial_{t} v\right) \\
& =\sum_{i, k} \partial_{i}\left(a\left(|\nabla v|^{2}\right) E_{i k} \partial_{k}\left(w+b\left(|\nabla v|^{2}\right)\right)\right) \\
& =\mathcal{L} w+2 \sum_{i, k} \partial_{i}\left(\left(a b^{\prime}\right)\left(|\nabla v|^{2}\right) E_{i k} \sum_{j} \partial_{j} v \partial_{j} \partial_{k} v\right) \\
& =\mathcal{L} w+2 \sum_{i, j} \partial_{i}\left(\left(a b^{\prime}\right)\left(|\nabla v|^{2}\right) \partial_{j} v \sum_{k} E_{i k} \partial_{k} \partial_{j} v\right) \\
& =\mathcal{L} w+2 \sum_{i, j} \partial_{i}\left(\left(\frac{b^{\prime}}{a}\right)\left(|\nabla v|^{2}\right) z_{j} \partial_{j} z_{i}\right) \\
& =\mathcal{L} w+4 \sum_{i, j}\left(\frac{b^{\prime}}{a}\right)^{\prime}\left(|\nabla v|^{2}\right) \sum_{k} \partial_{k} v \partial_{k} \partial_{i} v z_{j} \partial_{j} z_{i} \\
& +2 \sum_{i, j}\left(\frac{b^{\prime}}{a}\right)\left(|\nabla v|^{2}\right) \partial_{i} z_{j} \partial_{j} z_{i}+2 \sum_{i, j}\left(\frac{b^{\prime}}{a}\right)\left(|\nabla v|^{2}\right) z_{j} \partial_{j} \partial_{i} z_{i} .
\end{aligned}
$$

Since $w=\sum \partial_{i} z_{i}$, the last term of the right-hand side of the above inequality is equal to $\mathcal{V} \cdot \nabla w$ and

$$
\begin{align*}
\partial_{t} w & =\mathcal{L} w+\mathcal{V} \cdot \nabla w+4\left[a\left(\frac{b^{\prime}}{a}\right)^{\prime}\right]\left(|\nabla v|^{2}\right) \sum_{i, j, k} \partial_{j} v \partial_{k} v \partial_{k} \partial_{i} v \partial_{j} z_{i}  \tag{3.8}\\
& +2\left(\frac{b^{\prime}}{a}\right)\left(|\nabla v|^{2}\right) \sum_{i, j} \partial_{i} z_{j} \partial_{j} z_{i} .
\end{align*}
$$

On the one hand, introducing the matrix $\mathcal{E}:=\left(E_{i j}\right)$ and the Hessian matrix $D^{2} v=\left(\partial_{i} \partial_{j} v\right)$ of $v$, we infer from (3.6) that

$$
\begin{align*}
\sum_{i, j} \partial_{i} z_{j} \partial_{j} z_{i} & =a^{2}\left(|\nabla v|^{2}\right) \sum_{i, j, k, l} E_{i k} \partial_{k} \partial_{j} v E_{j l} \partial_{l} \partial_{i} v \\
& =a^{2}\left(|\nabla v|^{2}\right) \sum_{i, j}\left(\mathcal{E} D^{2} v\right)_{i j}\left(\mathcal{E} D^{2} v\right)_{j i} \\
\sum_{i, j} \partial_{i} z_{j} \partial_{j} z_{i} & =a^{2}\left(|\nabla v|^{2}\right) \operatorname{tr}\left(\left(\mathcal{E} D^{2} v\right)^{2}\right) . \tag{3.9}
\end{align*}
$$

On the other hand, using once more (3.6), we obtain

$$
\begin{aligned}
\sum_{i, j, k} \partial_{j} v \partial_{k} v \partial_{k} \partial_{i} v \partial_{j} z_{i} & =a\left(|\nabla v|^{2}\right) \sum_{i, j, k, l} \partial_{j} v \partial_{k} v \partial_{k} \partial_{i} v E_{i l} \partial_{l} \partial_{j} v \\
& =a\left(|\nabla v|^{2}\right) \sum_{i, l}\left(\sum_{k} \partial_{i} \partial_{k} v \partial_{k} v\right) E_{i l}\left(\sum_{j} \partial_{l} \partial_{j} v \partial_{j} v\right) \\
\sum_{i, j, k} \partial_{j} v \partial_{k} v \partial_{k} \partial_{i} v \partial_{j} z_{i} & =a\left(|\nabla v|^{2}\right)\left\langle D^{2} v \nabla v,\left(\mathcal{E} D^{2} v\right) \nabla v\right\rangle
\end{aligned}
$$

Inserting (3.9) and (3.10) in (3.8), we end up with

$$
\begin{align*}
\partial_{t} w & =\mathcal{L} w+\mathcal{V} \cdot \nabla w+2\left(a b^{\prime}\right)\left(|\nabla v|^{2}\right) \operatorname{tr}\left(\left(\mathcal{E} D^{2} v\right)^{2}\right)  \tag{3.11}\\
& +4\left(a b^{\prime \prime}-a^{\prime} b^{\prime}\right)\left(|\nabla v|^{2}\right)\left\langle D^{2} v \nabla v,\left(\mathcal{E} D^{2} v\right) \nabla v\right\rangle .
\end{align*}
$$

We next observe that

$$
\begin{equation*}
\mathcal{E} \nabla v=\left(1+2|\nabla v|^{2}\left(\frac{a^{\prime}}{a}\right)\left(|\nabla v|^{2}\right)\right) \nabla v \tag{3.12}
\end{equation*}
$$

and that, for $X \in \mathbb{R}^{N}$,

$$
\langle\mathcal{E} X, X\rangle=|X|^{2}+2\left(\frac{a^{\prime}}{a}\right)\left(|\nabla v|^{2}\right)\langle X, \nabla v\rangle^{2} \geq|X|^{2}
$$

as $a$ and $a^{\prime}$ are both positive by (3.1). Consequently, $\mathcal{E}$ is a positive definite symmetric matrix in $\mathcal{M}_{N}(\mathbb{R})$ and there exists a positive definite matrix $\mathcal{E}_{1 / 2}$ such that $\mathcal{E}_{1 / 2}^{2}=\mathcal{E}$. We then infer from the definition of $\mathcal{E}_{1 / 2}$, (3.12), and Lemma 3.2 (with $A=\mathcal{E}_{1 / 2}, B=D^{2} v$ and
$\left.X=\mathcal{E}_{1 / 2}^{-1} \nabla v\right)$ that

$$
\begin{aligned}
\left\langle D^{2} v \nabla v,\left(\mathcal{E} D^{2} v\right) \nabla v\right\rangle & =\left|\left(\mathcal{E}_{1 / 2} D^{2} v\right) \nabla v\right|^{2} \\
& =\left|\left(\left(\mathcal{E}_{1 / 2} D^{2} v \mathcal{E}_{1 / 2}\right) \mathcal{E}_{1 / 2}^{-1}\right) \nabla v\right|^{2} \\
& \leq \operatorname{tr}\left(\mathcal{E}_{1 / 2} D^{2} v \mathcal{E}_{1 / 2} \mathcal{E}_{1 / 2} D^{2} v \mathcal{E}_{1 / 2}\right)\left\langle\mathcal{E}_{1 / 2}^{-1} \nabla v, \mathcal{E}_{1 / 2}^{-1} \nabla v\right\rangle \\
& \leq \operatorname{tr}\left(\left(\mathcal{E} D^{2} v\right)^{2}\right)\left\langle\nabla v, \mathcal{E}^{-1} \nabla v\right\rangle \\
& \leq \operatorname{tr}\left(\left(\mathcal{E} D^{2} v\right)^{2}\right)|\nabla v|^{2}\left(1+2|\nabla v|^{2}\left(\frac{a^{\prime}}{a}\right)\left(|\nabla v|^{2}\right)\right)^{-1}
\end{aligned}
$$

Owing to the non-positivity (3.1) of $a b^{\prime \prime}-a^{\prime} b^{\prime}$, we deduce from (3.11) and the above inequality that

$$
\partial_{t} w \geq \mathcal{L} w+\mathcal{V} \cdot \nabla w+\left(a^{2} c\right)\left(|\nabla v|^{2}\right) \operatorname{tr}\left(\left(\mathcal{E} D^{2} v\right)^{2}\right)
$$

the function $c$ being defined in (3.2). We finally use the inequality

$$
\operatorname{tr}\left(A^{2}\right) \geq \frac{1}{N} \operatorname{tr}(A)^{2}, \quad A \in \mathcal{M}_{N}(\mathbb{R})
$$

the identity

$$
w=\sum_{i} \partial_{i} z_{i}=a\left(|\nabla v|^{2}\right) \operatorname{tr}\left(\mathcal{E} D^{2} v\right)
$$

and the non-negativity (3.2) of $c$ to conclude that

$$
\begin{aligned}
\partial_{t} w & \geq \mathcal{L} w+\mathcal{V} \cdot \nabla w+\frac{1}{N}\left(a^{2} c\right)\left(|\nabla v|^{2}\right) \operatorname{tr}\left(\mathcal{E} D^{2} v\right)^{2} \\
& \geq \mathcal{L} w+\mathcal{V} \cdot \nabla w+\frac{c\left(|\nabla v|^{2}\right)}{N} w^{2},
\end{aligned}
$$

and complete the proof.
Proof of Proposition 1.4. To be able to use Lemma 3.1, we shall first construct a suitable approximation of (1.1), (1.2). Such a construction has already been performed in [6] for similar purposes and we recall it now. Given $u_{0}$ satisfying (1.3), there is a sequence of functions $\left(u_{0, k}\right)_{k \geq 1}$ such that, for each integer $k \geq 1, u_{0, k} \in \mathcal{B C} C^{\infty}\left(\mathbb{R}^{N}\right), u_{0} \leq u_{0, k+1} \leq u_{0, k}$, and $\left(u_{0, k}, \nabla u_{0, k}\right)_{k}$ converge towards $\left(u_{0}, \nabla u_{0}\right)$ uniformly on every compact subset of $\mathbb{R}^{N}$ as $k \rightarrow \infty$. Next, for $\varepsilon \in(0,1)$ and $r \geq 0$, we set

$$
a_{\varepsilon}(r):=\left(r+\varepsilon^{2}\right)^{(p-2) / 2} \quad \text { and } \quad b_{\varepsilon}(r):=\left(r+\varepsilon^{2}\right)^{q / 2}-\varepsilon^{q} .
$$

Then the Cauchy problem

$$
\begin{align*}
\partial_{t} u_{k, \varepsilon} & =\operatorname{div}\left(a_{\varepsilon}\left(\left|\nabla u_{k, \varepsilon}\right|^{2}\right) \nabla u_{k, \varepsilon}\right)+b_{\varepsilon}\left(\left|\nabla u_{k, \varepsilon}\right|^{2}\right), \quad(t, x) \in Q_{\infty}  \tag{3.13}\\
u_{k, \varepsilon}(0) & =u_{0, k}+\varepsilon^{\nu}, \quad x \in \mathbb{R}^{N} \tag{3.14}
\end{align*}
$$

has a unique classical solution $u_{k, \varepsilon}$, the parameter $\nu>0$ depending $p, q$, and $N$ and being appropriately chosen. Furthermore,

$$
\begin{align*}
\left\|\nabla u_{k, \varepsilon}(t)\right\|_{\infty} & \leq\left\|\nabla u_{0, k}\right\|_{\infty}, \quad t \geq 0  \tag{3.15}\\
\lim _{k \rightarrow \infty} \lim _{\varepsilon \rightarrow 0} u_{k, \varepsilon}(t, x) & =u(t, x) \tag{3.16}
\end{align*}
$$

the latter convergence being uniform on every compact subset of $[0, \infty) \times \mathbb{R}^{N}$, see $[6$, Section 3] (after performing the same change of unknown function as in the proof of Proposition 2.1).
Introducing

$$
c_{\varepsilon}(r)=2\left(\frac{b_{\varepsilon}^{\prime}}{a_{\varepsilon}}\right)(r)+\frac{4 r\left(a_{\varepsilon} b_{\varepsilon}^{\prime \prime}-a_{\varepsilon}^{\prime} b_{\varepsilon}^{\prime}\right)(r)}{a_{\varepsilon}^{2}(r)+2 r a_{\varepsilon}(r) a_{\varepsilon}^{\prime}(r)}, \quad r \geq 0
$$

let us check that $a_{\varepsilon}$ and $b_{\varepsilon}$ fulfill the conditions (3.1) and (3.2). Clearly, $a_{\varepsilon}>0$ and $a_{\varepsilon}^{\prime}>0$ as $p>2$. Next, since $1<q \leq p$,

$$
\begin{aligned}
\left(a_{\varepsilon}^{\prime} b_{\varepsilon}^{\prime}-a_{\varepsilon} b_{\varepsilon}^{\prime \prime}\right)(r) & =\frac{q(p-q)}{4}\left(r+\varepsilon^{2}\right)^{(p+q-6) / 2} \geq 0 \\
c_{\varepsilon}(r) & =q \frac{r(q-1)+\varepsilon^{2}}{r(p-1)+\varepsilon^{2}}\left(r+\varepsilon^{2}\right)^{(q-p) / 2} \geq 0
\end{aligned}
$$

We may then apply Lemma 3.1 to deduce that $w_{k, \varepsilon}:=\operatorname{div}\left(a_{\varepsilon}\left(\left|\nabla u_{k, \varepsilon}\right|^{2}\right) \nabla u_{k, \varepsilon}\right)$ satisfies

$$
\partial_{t} w_{k, \varepsilon}-\mathcal{L}_{k, \varepsilon} w_{k, \varepsilon}-\mathcal{V}_{k, \varepsilon} \cdot \nabla w_{k, \varepsilon}-\frac{c_{\varepsilon}\left(\left|\nabla u_{k, \varepsilon}\right|^{2}\right)}{N} w_{k, \varepsilon}^{2} \geq 0
$$

in $Q_{\infty}$. Observe next that the condition $1<q \leq p$ implies that $c_{\varepsilon}$ is a non-increasing function. It then follows from (3.15) that $c_{\varepsilon}\left(\left|\nabla u_{k, \varepsilon}\right|^{2}\right) \geq c_{\varepsilon}\left(\left\|\nabla u_{0, k}\right\|_{\infty}^{2}\right)$ and we end up with

$$
\begin{equation*}
\partial_{t} w_{k, \varepsilon}-\mathcal{L}_{k, \varepsilon} w_{k, \varepsilon}-\mathcal{V}_{k, \varepsilon} \cdot \nabla w_{k, \varepsilon}-\frac{c_{\varepsilon}\left(\left\|\nabla u_{0, k}\right\|_{\infty}^{2}\right)}{N} w_{k, \varepsilon}^{2} \geq 0 \tag{3.17}
\end{equation*}
$$

in $Q_{\infty}$. Clearly, $t \longmapsto-N /\left(c_{\varepsilon}\left(\left\|\nabla u_{0, k}\right\|_{\infty}^{2}\right) t\right)$ is a subsolution to (3.17) and the comparison principle warrants that

$$
\begin{equation*}
w_{k, \varepsilon}(t, x) \geq-\frac{N}{c_{\varepsilon}\left(\left\|\nabla u_{0, k}\right\|_{\infty}^{2}\right) t}, \quad(t, x) \in Q_{\infty} \tag{3.18}
\end{equation*}
$$

Letting $\varepsilon \rightarrow 0$ and $k \rightarrow \infty$ in the previous inequality with the help of (3.16) gives (1.12).
Next, since (1.1) is autonomous, we infer from (2.2) (with $s=0$ ) and (1.12) that

$$
\begin{aligned}
\Delta_{p} u(t, x) & \geq-\frac{2 N(p-1)}{q(q-1)} \frac{\|\nabla u(t / 2)\|_{\infty}^{p-q}}{t} \\
& \geq-\frac{2^{p / q} N(p-1)}{q(q-1)} C_{1}^{p-q}\left\|u_{0}\right\|_{\infty}^{(p-q) / q} t^{-p / q},
\end{aligned}
$$

whence (1.11) for $s=0$. To prove the general case $s \in(0, t)$, we use again the fact that (1.1) is autonomous.

We have a similar result when $u_{0}$ is more regular.

Corollary 3.3 Assume that $p>2$ and $q \in(1, p]$. Let $u$ be the solution to (1.1), (1.2) with an initial condition $u_{0}$ satisfying $u_{0} \in W^{2, \infty}\left(\mathbb{R}^{N}\right)$ in addition to (1.3). Then

$$
\begin{equation*}
\Delta_{p} u(t, x) \geq-\left|\inf _{y \in \mathbb{R}^{N}} \Delta_{p} u_{0}(y)\right| \tag{3.19}
\end{equation*}
$$

in the sense of distributions.
Proof. Keeping the notations introduced in the proof of Proposition 1.4, we readily infer from (3.17) and the comparison principle that

$$
\begin{equation*}
w_{k, \varepsilon}(t, x) \geq-\left|\inf _{y \in \mathbb{R}^{N}} \Delta_{p} u_{0, k}(y)\right|, \quad(t, x) \in Q_{\infty} . \tag{3.20}
\end{equation*}
$$

Owing to the regularity of $u_{0}$, it is possible to construct the sequence $\left(u_{0, k}\right)_{k}$ such that it satisfies

$$
\lim _{k \rightarrow \infty} \inf _{y \in \mathbb{R}^{N}} \Delta_{p} u_{0, k}(y)=\inf _{y \in \mathbb{R}^{N}} \Delta_{p} u_{0}(y) .
$$

We may then pass to the limit first as $\varepsilon \rightarrow 0$ and then as $k \rightarrow \infty$ in (3.20) and use (3.16) and the above convergence to complete the proof.

Another useful consequence of the semiconvexity estimates derived in Proposition 1.4 is that the solution $u$ to (1.1), (1.2) is a supersolution to a first-order Hamilton-Jacobi equation.

Corollary 3.4 Consider an initial condition $u_{0}$ satisfying (1.3). Setting $F\left(t, \xi_{0}, \xi\right):=\xi_{0}-$ $|\xi|^{q}+\kappa_{1}\left\|u_{0}\right\|_{\infty}^{(p-q) / q} t^{-p / q}$ for $t \in(0, \infty), \xi_{0} \in \mathbb{R}$, and $\xi \in \mathbb{R}^{N}$ (recall that $\kappa_{1}$ is defined in (1.11)), the solution $u$ to (1.1), (1.2) is a supersolution to $F\left(t, \partial_{t} w, \nabla w\right)=0$ in $Q_{\infty}$.

Proof. We still use the notations introduced in the proof of Proposition 1.4. As $w_{k, \varepsilon}=$ $\operatorname{div}\left(a_{\varepsilon}\left(\left|\nabla u_{k, \varepsilon}\right|^{2}\right) \nabla u_{k, \varepsilon}\right)$, we infer from (3.13) and (3.18) that

$$
\partial_{t} u_{k, \varepsilon}-b_{\varepsilon}\left(\left|\nabla u_{k, \varepsilon}\right|^{2}\right) \geq-\frac{N}{c_{\varepsilon}\left(\left\|\nabla u_{0, k}\right\|_{\infty}^{2}\right) t}
$$

in $Q_{\infty}$. We then use (3.16) and the stability of viscosity solutions $[1,2,8]$ to pass to the limit as $\varepsilon \rightarrow 0$ and $k \rightarrow \infty$ in the previous inequality and conclude that $u$ is a supersolution to

$$
\partial_{t} w-|\nabla w|^{q}+\frac{N(p-1)}{q(q-1)} \frac{\left\|\nabla u_{0}\right\|_{\infty}^{p-q}}{t}=0 \quad \text { in } \quad Q_{\infty} .
$$

Now, fix $T \geq 0$. As (1.1) is an autonomous equation, the function $(t, x) \longmapsto u(t+T, x)$ is the solution to (1.1) with initial condition $u(T)$ and the above analysis allows us to conclude that $u$ is a supersolution to

$$
\partial_{t} w-|\nabla w|^{q}+\frac{N(p-1)}{q(q-1)} \frac{\|\nabla u(T)\|_{\infty}^{p-q}}{t-T}=0 \quad \text { in } \quad(T, \infty) \times \mathbb{R}^{N} .
$$

We then use (2.2) (with $T=t / 2$ ) to complete the proof.

## 4 Convergence to self-similarity

We change the variables and the unknown function so that the convergence (1.5) is transformed to the convergence towards a steady state. More precisely, we introduce the selfsimilar (or scaling) variables

$$
\tau=\frac{1}{q} \log (1+t), \quad y=\frac{x}{(1+t)^{1 / q}}
$$

and the new unknown function $v$ defined by

$$
\begin{equation*}
u(t, x)=v\left(\frac{\log (1+t)}{q}, \frac{x}{(1+t)^{1 / q}}\right), \quad(t, x) \in[0, \infty) \times \mathbb{R}^{N} \tag{4.1}
\end{equation*}
$$

Equivalently, $v(\tau, y)=u\left(e^{q \tau}-1, y e^{\tau}\right)$ for $(\tau, y) \in[0, \infty) \times \mathbb{R}^{N}$ and it follows from (1.1), (1.2) that $v$ solves

$$
\begin{align*}
\partial_{\tau} v & =y \cdot \nabla v+q|\nabla v|^{q}+q e^{-(p-q) \tau} \Delta_{p} v, \quad(\tau, y) \in Q_{\infty},  \tag{4.2}\\
v(0) & =u_{0}, \quad y \in \mathbb{R}^{N} . \tag{4.3}
\end{align*}
$$

We also infer from (2.1) and (2.2) that there is a positive constant $C_{5}\left(u_{0}\right)$ depending only on $N, p, q$, and $u_{0}$ such that

$$
\begin{equation*}
\|v(\tau)\|_{\infty}+\|\nabla v(\tau)\|_{\infty} \leq C_{5}\left(u_{0}\right), \quad \tau \geq 0 \tag{4.4}
\end{equation*}
$$

while (1.4) reads

$$
\begin{equation*}
\lim _{\tau \rightarrow \infty}\|v(\tau)\|_{\infty}=M_{\infty}>0 \tag{4.5}
\end{equation*}
$$

Formally, since $p>q$, the diffusion term vanishes in the large time limit and we expect the large time behaviour of the solution $v$ to (4.2), (4.3) to look like that of the solutions to the first-order Hamilton-Jacobi equation

$$
\begin{equation*}
\partial_{\tau} w-y \cdot \nabla w-q|\nabla w|^{q}=0 \quad \text { in } \quad Q_{\infty} . \tag{4.6}
\end{equation*}
$$

Now, to investigate the large time behaviour of first-order Hamilton-Jacobi equations, an efficient approach has been developed in $[22,23]$ which relies on the relaxed half-limits method introduced in [3]. More precisely, for $(\tau, y) \in Q_{\infty}$, we define the relaxed half-limits $v_{*}$ and $v^{*}$ by

$$
\begin{equation*}
v_{*}(y):=\liminf _{(\sigma, z, \lambda) \rightarrow(\tau, y, \infty)} v(\sigma+\lambda, z) \quad \text { and } \quad v^{*}(y):=\limsup _{(\sigma, z, \lambda) \rightarrow(\tau, y, \infty)} v(\sigma+\lambda, z) . \tag{4.7}
\end{equation*}
$$

These relaxed half-limits are well-defined thanks to (4.4) and we first note that the right-hand sides of the above definitions indeed do not depend on $\tau>0$. In addition,

$$
\begin{equation*}
0 \leq v_{*}(x) \leq v^{*}(x) \leq M_{\infty} \quad \text { for } \quad y \in \mathbb{R}^{N} \tag{4.8}
\end{equation*}
$$

by (4.5), while (4.4) and the Rademacher theorem ensure that $v_{*}$ and $v^{*}$ both belong to $W^{1, \infty}\left(\mathbb{R}^{N}\right)$. Finally, by [2, Théorème 4.1] applied to equation (4.2), $v^{*}$ and $v_{*}$ are viscosity subsolution and supersolution, respectively, to the Hamilton-Jacobi equation

$$
\begin{equation*}
\mathcal{H}(y, \nabla w):=-y \cdot \nabla w-q|\nabla w|^{q}=0 \quad \text { in } \quad \mathbb{R}^{N} . \tag{4.9}
\end{equation*}
$$

We now aim at showing that $v^{*}$ and $v_{*}$ coincide. However, the equation (4.9) has infinitely many solutions as $y \longmapsto\left(c-\gamma_{q}|y|^{q /(q-1)}\right)_{+}$solves (4.9) for any $c>0$. The information obtained so far on $v_{*}$ and $v^{*}$ are thus not sufficient and are supplemented by the next two results.

Lemma 4.1 Given $\varepsilon \in(0,1)$, there is $R_{\varepsilon}>1 / \varepsilon$ such that

$$
\begin{equation*}
v(\tau, y) \leq \varepsilon \quad \text { for } \quad \tau \geq 0 \quad \text { and } \quad y \in \mathbb{R}^{N} \backslash B\left(0, R_{\varepsilon}\right) \tag{4.10}
\end{equation*}
$$

and $0 \leq v_{*}(y) \leq v^{*}(y) \leq \varepsilon$ for $y \in \mathbb{R}^{N} \backslash B\left(0, R_{\varepsilon}\right)$.
In other words, $v(\tau)$ belongs to $\mathcal{C}_{0}\left(\mathbb{R}^{N}\right)$ for each $\tau \geq 0$ in a way which is uniform with respect to $\tau \geq 0$.
Proof. We first construct a supersolution to (4.2) in ( $0, \infty) \times\left(\mathbb{R}^{N} \backslash B(0, R)\right)$ for $R$ large enough. To this end, consider $R \geq R_{c}:=1+\left(q\left(2\left\|u_{0}\right\|_{\infty}\right)^{q-1}+3 p q\left(2\left\|u_{0}\right\|_{\infty}\right)^{p-2}\right)^{1 / q}$ and put $\Sigma_{R}(y)=\left\|u_{0}\right\|_{\infty} R^{2}|y|^{-2}$ for $y \in \mathbb{R}^{N} \backslash B(0, R)$. Let $\mathcal{L}$ be the parabolic operator defined by

$$
\mathcal{L} w(\tau, y):=\partial_{\tau} w(\tau, y)-y \cdot \nabla w(\tau, y)-q|\nabla w(\tau, y)|^{q}-q e^{-(p-q) \tau} \Delta_{p} w(\tau, y)
$$

for $(\tau, y) \in Q_{\infty}$ (so that $\mathcal{L} v=0$ by (4.2)). Then, if $y \in \mathbb{R}^{N} \backslash B(0, R)$, we have

$$
\begin{aligned}
\mathcal{L} \Sigma_{R}(y) & =2 \Sigma_{R}(y)-q \frac{2^{q}}{|y|^{q}} \Sigma_{R}(y)^{q}+q 2^{p-1} \frac{N+2-3 p}{|y|^{p}} \Sigma_{R}(y)^{p-1} e^{-(p-q) \tau} \\
& \geq 2 \Sigma_{R}(y)\left\{1-q\left(2\left\|u_{0}\right\|_{\infty}\right)^{q-1} \frac{R^{2(q-1)}}{|y| 3 q-2}-3 p q e^{-(p-q) \tau}\left(2\left\|u_{0}\right\|_{\infty}\right)^{p-2} \frac{R^{2(p-2)}}{|y|^{3 p-4}}\right\} \\
& \geq 2 \Sigma_{R}(y)\left\{1-q\left(2\left\|u_{0}\right\|_{\infty}\right)^{q-1} R^{-q}-3 p q e^{-(p-q) \tau}\left(2\left\|u_{0}\right\|_{\infty}\right)^{p-2} R^{-p}\right\} \\
& \geq 0
\end{aligned}
$$

by the choice of $R$. Consequently, $\Sigma_{R}$ is a supersolution to (4.2) in $(0, \infty) \times\left(\mathbb{R}^{N} \backslash B(0, R)\right)$ for $R \geq R_{c}$.
Now, fix $\varepsilon \in(0,1)$. Since $u_{0} \in \mathcal{C}_{0}\left(\mathbb{R}^{N}\right)$, there is $\rho_{\varepsilon} \geq \max \left\{R_{c}, \varepsilon^{-1}\right\}$ such that $u_{0}(y) \leq \varepsilon / 2$ if $|y| \geq \rho_{\varepsilon}$. We then infer from the monotonicity of $\Sigma_{R}$ and (2.1) that

$$
u_{0}(y)-\frac{\varepsilon}{2}-\Sigma_{\rho_{\varepsilon}}(z) \leq-\Sigma_{\rho_{\varepsilon}}(z) \leq 0
$$

if $|y| \geq \rho_{\varepsilon}$ and $|z| \geq \rho_{\varepsilon}$,

$$
v(\tau, y)-\frac{\varepsilon}{2}-\Sigma_{\rho_{\varepsilon}}(z) \leq\left\|u_{0}\right\|_{\infty}-\Sigma_{\rho_{\varepsilon}}(z)=\Sigma_{\rho_{\varepsilon}}(y)-\Sigma_{\rho_{\varepsilon}}(z) \leq \frac{2\left\|u_{0}\right\|_{\infty}}{\rho_{\varepsilon}}|y-z|
$$

if $|y|=\rho_{\varepsilon},|z| \geq \rho_{\varepsilon}$, and $\tau \geq 0$, and

$$
v(\tau, y)-\frac{\varepsilon}{2}-\Sigma_{\rho_{\varepsilon}}(z) \leq\left\|u_{0}\right\|_{\infty}-\left\|u_{0}\right\|_{\infty} \leq 0
$$

if $|y| \geq \rho_{\varepsilon},|z|=\rho_{\varepsilon}$, and $\tau \geq 0$. As $v-\varepsilon / 2$ and $\Sigma_{\rho_{\varepsilon}}$ are subsolution and supersolution, respectively, to (4.2), the comparison principle [11, Theorem 4.1] warrants that $v(\tau, y)-\varepsilon / 2 \leq$ $\Sigma_{\rho_{\varepsilon}}(y)$ for $\tau \geq 0$ and $|y| \geq \rho_{\varepsilon}$. It remains to choose $R_{\varepsilon} \geq \rho_{\varepsilon}$ such that $\Sigma_{\rho_{\varepsilon}}(y) \leq \varepsilon / 2$ for $|y| \geq R_{\varepsilon}$ to complete the proof of (4.10). The last assertion of Lemma 4.1 is then a straightforward consequence of the definition (4.7) and (4.10).

We next use the semiconvexity estimate (1.11) (and more precisely its consequence stated in Corollary 3.4) to show that $v_{*}$ lies above the profile $H_{\infty}$ defined in (1.6).

Lemma 4.2 For $y \in \mathbb{R}^{N}$, we have

$$
\begin{equation*}
H_{\infty}(y) \leq v_{*}(y) \leq v^{*}(y) . \tag{4.11}
\end{equation*}
$$

Proof. For $\tau \geq 0, y \in \mathbb{R}^{N}, \xi_{0} \in \mathbb{R}$ and $\xi \in \mathbb{R}^{N}$, we set $\mathcal{F}\left(\tau, y, \xi_{0}, \xi\right):=\xi_{0}-y \cdot \xi-q|\xi|^{q}+$ $\kappa_{2} e^{-(p-q) \tau}$ with $\kappa_{2}:=q \kappa_{1}\left\|u_{0}\right\|_{\infty}^{(p-q) / q}\left(e^{q} /\left(e^{q}-1\right)\right)^{p / q}$, the constant $\kappa_{1}$ being defined in (1.11). It then readily follows from Corollary 3.4 that

$$
\begin{equation*}
v \text { is a supersolution to } \mathcal{F}\left(\tau, y, \partial_{\tau} w, \nabla w\right)=0 \quad \text { in }(1, \infty) \times \mathbb{R}^{N} \tag{4.12}
\end{equation*}
$$

We next fix $\tau_{0}>1$ and denote by $V$ the (viscosity) solution to

$$
\begin{aligned}
\partial_{\tau} V-y \cdot \nabla V-q|\nabla V|^{q} & =0, \quad(\tau, y) \in\left(\tau_{0}, \infty\right) \times \mathbb{R}^{N}, \\
V\left(\tau_{0}\right) & =v\left(\tau_{0}\right), \quad y \in \mathbb{R}^{N} .
\end{aligned}
$$

On the one hand, a straightforward computation shows that the function $\tilde{V}$ defined by

$$
\tilde{V}(\tau, y):=V(\tau, y)-\kappa_{2} \int_{\tau_{0}}^{\tau} e^{-(p-q) s} d s, \quad(\tau, y) \in\left(\tau_{0}, \infty\right) \times \mathbb{R}^{N}
$$

is the (viscosity) solution to $\mathcal{F}\left(\tau, y, \partial_{\tau} \tilde{V}, \nabla \tilde{V}\right)=0$ in $\left(\tau_{0}, \infty\right) \times \mathbb{R}^{N}$ with initial condition $\tilde{V}\left(\tau_{0}\right)=v\left(\tau_{0}\right)$. Recalling (4.12), we infer from the comparison principle that

$$
\begin{equation*}
\tilde{V}(\tau, y) \leq v(\tau, y) \quad \text { for } \quad(\tau, y) \in\left(\tau_{0}, \infty\right) \times \mathbb{R}^{N} \tag{4.13}
\end{equation*}
$$

On the other hand, it follows from Proposition A. 1 that

$$
\lim _{\tau \rightarrow \infty} \sup _{y \in \mathbb{R}^{N}}\left|V(\tau, y)-\left(\left\|v\left(\tau_{0}\right)\right\|_{\infty}-\gamma_{q}|y|^{q /(q-1)}\right)_{+}\right|=0
$$

We may then pass to the limit as $\tau \rightarrow \infty$ in (4.13) and use the definition (4.7) to conclude that

$$
\left(\left\|v\left(\tau_{0}\right)\right\|_{\infty}-\gamma_{q}|y|^{q /(q-1)}\right)_{+}-\kappa_{2} \int_{\tau_{0}}^{\infty} e^{-(p-q) s} d s \leq v_{*}(y) \leq v^{*}(y)
$$

for $y \in \mathbb{R}^{N}$. Letting $\tau_{0} \rightarrow \infty$ in the above inequality with the help of (4.5) completes the proof of the lemma.

We are now in a position to complete the proof of Theorem 1.1. To this end, fix $\varepsilon \in(0,1)$. Lemma 4.1 ensures that $v^{*}(y) \leq \varepsilon$ for $|y| \geq R_{\varepsilon} \geq 1 / \varepsilon$ while the continuity of $H_{\infty}$ implies that there is $r_{\varepsilon} \in(0, \varepsilon)$ such that $H_{\infty}(y) \geq M_{\infty}-\varepsilon$ for $|y| \leq r_{\varepsilon}$. Recalling (4.8), we realize that

$$
\left\{\begin{array}{ll}
v^{*}(y)-\varepsilon \leq 0 \leq H_{\infty}(y) & \text { if }  \tag{4.14}\\
v^{*}(y)-\varepsilon \leq M_{\infty}-\varepsilon \leq H_{\infty}(y) & \text { if }
\end{array}|y|=r_{\varepsilon}\right.
$$

Moreover, introducing $\psi(y)=-\gamma_{q}|y|^{q /(q-1)} / 2$, we have

$$
\begin{equation*}
\mathcal{H}(y, \nabla \psi(y))=\frac{q \gamma_{q}}{2(q-1)}|y|^{q /(q-1)}\left(1-\frac{1}{2^{q-1}}\right)>0 \quad \text { if } \quad r_{\varepsilon}<|y|<R_{\varepsilon} \tag{4.15}
\end{equation*}
$$

the Hamiltonian $\mathcal{H}$ being defined in (4.9). Summarizing, we have shown that $H_{\infty}$ and $v^{*}-\varepsilon$ are supersolution and subsolution, respectively, to (4.9) in $\Omega_{\varepsilon}:=\left\{y \in \mathbb{R}^{N}: r_{\varepsilon}<|y|<R_{\varepsilon}\right\}$ with $v^{*}-\varepsilon \leq H_{\infty}$ on $\partial \Omega_{\varepsilon}$ by (4.14). Owing to (4.15) and the concavity of $\mathcal{H}$ with respect to its second variable, we may apply [15, Theorem 1] to conclude that $v^{*}-\varepsilon \leq H_{\infty}$ in $\Omega_{\varepsilon}$. This property being valid for each $\varepsilon \in(0,1)$, we actually have $v^{*} \leq H_{\infty}$ in $\mathbb{R}^{N}$ by passing to the limit as $\varepsilon \rightarrow 0$ thanks to the properties of $r_{\varepsilon}$ and $R_{\varepsilon}$. Recalling (4.11), we have thus established that $v^{*}=v_{*}=H_{\infty}$ in $\mathbb{R}^{N}$. In particular, the property $v^{*}=v_{*}$ and the definition (4.7) provide the uniform convergence of $\{v(\tau)\}_{\tau \geq 0}$ towards $v^{*}=H_{\infty}$ on every compact subset of $\mathbb{R}^{N}$ as $\tau \rightarrow \infty$, see [2, Lemme 4.1] or [1, Lemma V.1.9]. Combining this local convergence with Lemma 4.1 actually gives

$$
\begin{equation*}
\lim _{\tau \rightarrow \infty}\left\|v(\tau)-H_{\infty}\right\|_{\infty}=0 . \tag{4.16}
\end{equation*}
$$

Theorem 1.1 then readily follows after writing the convergence (4.16) in the original variables $(t, x)$ for the function $u$ and noticing that $\left\|h_{\infty}(1+t)-h_{\infty}(t)\right\|_{\infty} \longrightarrow 0$ as $t \rightarrow \infty$.

## 5 Limit value of $\|u(t)\|_{\infty}$

This section is devoted to the proof of Proposition 1.4, for which three cases are to be distinguished and handled differently: $q \in(1, p-1], q \in\left(p-1, q_{\star}\right]$, and $q \in\left(q_{\star}, p\right)$.

Proof of Proposition 1.4: $q \in(1, p-1]$. We proceed as in [19, Proposition 1] (where a similar result is proved for $p=2$ and $q=1$ ). For $\alpha>N / 2, \delta>0$, and $x \in \mathbb{R}^{N}$, we set $\varrho_{\delta}(x):=\left(1+\delta|x|^{2}\right)^{-\alpha}$. Clearly, $\varrho_{\delta} \in L^{1}\left(\mathbb{R}^{N}\right)$ and it follows from (2.3) that

$$
\begin{aligned}
\frac{d}{d t} \int_{\mathbb{R}^{N}} \varrho_{\delta}(x) u(t, x) d x & =\int_{\mathbb{R}^{N}}\left\{\varrho_{\delta}(x)|\nabla u(t, x)|^{q}-|\nabla u(t, x)|^{p-2} \nabla u(t, x) \cdot \nabla \varrho_{\delta}(x)\right\} d x \\
& \geq \int_{\mathbb{R}^{N}} \varrho_{\delta}(x)|\nabla u(t, x)|^{q}\left(1-|\nabla u(t, x)|^{p-1-q} \frac{\left|\nabla \varrho_{\delta}(x)\right|}{\varrho_{\delta}(x)}\right) d x
\end{aligned}
$$

Recalling that $\|\nabla u(t)\|_{\infty} \leq\left\|\nabla u_{0}\right\|_{\infty}$ by (1.3) and (2.2) and noticing that $\left|\nabla \varrho_{\delta}\right| \leq \alpha \delta^{1 / 2} \varrho_{\delta}$, we further obtain

$$
\frac{d}{d t} \int_{\mathbb{R}^{N}} \varrho_{\delta}(x) u(t, x) d x \geq \int_{\mathbb{R}^{N}} \varrho_{\delta}(x)|\nabla u(t, x)|^{q}\left(1-\alpha \delta^{1 / 2}\left\|\nabla u_{0}\right\|_{\infty}^{p-1-q}\right) d x
$$

Choosing $\delta=\left\|\nabla u_{0}\right\|_{\infty}^{2(q+1-p)} / \alpha^{2}$ and integrating with respect to time give

$$
\|u(t)\|_{\infty}\left\|\varrho_{\delta}\right\|_{1} \geq \int_{\mathbb{R}^{N}} \varrho_{\delta}(x) u(t, x) d x \geq \int_{\mathbb{R}^{N}} \varrho_{\delta}(x) u_{0}(x) d x>0 .
$$

We then pass to the limit as $t \rightarrow \infty$ to conclude that $M_{\infty}>0$.
We next turn to the case $q \in\left(p-1, q_{\star}\right]$ which turns out to be more complicated and requires two preparatory results.

Lemma 5.1 Assume that $q \in\left(1, q_{\star}\right]$ and let $u$ be a non-negative solution to (1.1), (1.2) with a compactly supported initial condition $u_{0}$ satisfying (1.3). Then $u(t) \in L^{1}\left(\mathbb{R}^{N}\right)$ for each $t \geq 0$, the function $t \longmapsto\|u(t)\|_{1}$ is non-decreasing and

$$
\begin{equation*}
\lim _{t \rightarrow \infty}\|u(t)\|_{1}=\infty \tag{5.1}
\end{equation*}
$$

Proof. For every $t \geq 0, u(t)$ is bounded and compactly supported by (2.1) and Proposition 2.2 , and is thus in $L^{1}\left(\mathbb{R}^{N}\right)$. The time monotonicity of the $L^{1}$-norm of $u$ then readily follows from (2.3) with $\vartheta=1$, a valid choice in this particular case as $u(t)$ is compactly supported. It further follows from (2.3) with $\vartheta=1$ that

$$
\begin{equation*}
\|u(t)\|_{1} \geq\|u(T)\|_{1}+\int_{T}^{t}\|\nabla u(s)\|_{q}^{q} d s \quad \text { for } t>T \geq 0 \tag{5.2}
\end{equation*}
$$

Consider next $T>0$ and $t>T$. Recalling the Gagliardo-Nirenberg inequality

$$
\begin{equation*}
\|w\|_{q} \leq C_{6}\|\nabla w\|_{q}^{N(q-1) /(N(q-1)+q)}\|w\|_{1}^{q /(N(q-1)+q)}, \quad w \in W^{1, q}\left(\mathbb{R}^{N}\right) \cap L^{1}\left(\mathbb{R}^{N}\right) \tag{5.3}
\end{equation*}
$$

we infer from (5.2), (5.3), and the time monotonicity of the $L^{1}$-norm of $u$ that

$$
\begin{aligned}
\|u(t)\|_{1}^{1+\left(q^{2} / N(q-1)\right)} & \geq\|u(t)\|_{1}^{q^{2} / N(q-1)}\left(\|u(T)\|_{1}+\int_{T}^{t}\|\nabla u(s)\|_{q}^{q} d s\right) \\
& \geq \int_{T}^{t}\|u(s)\|_{1}^{q^{2} / N(q-1)}\|\nabla u(s)\|_{q}^{q} d s \\
& \geq C_{7} \int_{T}^{t}\left(\|u(s)\|_{q}^{q}\right)^{(N(q-1)+q) / N(q-1)} d s .
\end{aligned}
$$

If $\varphi$ denotes the solution to the $p$-Laplacian equation $\partial_{t} \varphi-\Delta_{p} \varphi=0$ in $Q_{\infty}$ with initial condition $\varphi(0)=u_{0}$, the comparison principle readily implies that

$$
\begin{equation*}
\varphi(t, x) \leq u(t, x), \quad(t, x) \in Q_{\infty} \tag{5.4}
\end{equation*}
$$

Inserting this estimate in the previous lower bound for $\|u(t)\|_{1}$, we end up with

$$
\begin{equation*}
\|u(t)\|_{1}^{1+\left(q^{2} / N(q-1)\right)} \geq C_{7} \int_{T}^{t}\left(\|\varphi(s)\|_{q}^{q}\right)^{(N(q-1)+q) / N(q-1)} d s \tag{5.5}
\end{equation*}
$$

Now, by Proposition 2.3 we have

$$
\lim _{s \rightarrow \infty} s^{N(q-1) /(N(p-2)+p)}\left\|\varphi(s)-\mathcal{B}_{\left\|u_{0}\right\|_{1}}(s)\right\|_{q}^{q}=0
$$

and

$$
\left\|\mathcal{B}_{\left\|u_{0}\right\|_{1}}(s)\right\|_{q}^{q}=C_{8} s^{-N(q-1) /(N(p-2)+p)},
$$

so that

$$
\begin{aligned}
\|\varphi(s)\|_{q}^{q} & \geq\left(\left\|\mathcal{B}_{\left\|u_{0}\right\|_{1}}(s)\right\|_{q}-\left\|\varphi(s)-\mathcal{B}_{\left\|u_{0}\right\|_{1}}(s)\right\|_{q}\right)^{q} \\
& \geq s^{-N(q-1) /(N(p-2)+p)}\left(C_{8}-\left\|\varphi(s)-\mathcal{B}_{\left\|u_{0}\right\|_{1}}(s)\right\|_{q}\right)^{q} \\
& \geq\left(\frac{C_{8}}{2}\right)^{q} s^{-N(q-1) /(N(p-2)+p)}
\end{aligned}
$$

for $s \geq T$, provided $T$ is chosen sufficiently large. Inserting this estimate in (5.5) gives

$$
\begin{aligned}
\|u(t)\|_{1}^{1+\left(q^{2} / N(q-1)\right)} & \geq C_{9} \int_{T}^{t} s^{-(N(q-1)+q) /(N(p-2)+p)} d s \\
& \geq C_{10}\left\{\begin{array}{lll}
t^{(N+1)\left(q_{\star}-q\right) /(N(p-2)+p)}-T^{(N+1)\left(q_{\star}-q\right) /(N(p-2)+p)} & \text { if } q \in\left(1, q_{\star}\right) \\
\log (t / T) & \text { if } q=q_{\star}
\end{array}\right.
\end{aligned}
$$

We then let $t \rightarrow \infty$ to obtain the claimed result.
We next argue as in [12, Lemma 14] (for $p=2$ ) to show that, if $q \in(p-1, p)$ and $M_{\infty}=0$, then the $L^{\infty}$-norm of $u(t)$ decays faster than an explicit rate.

Lemma 5.2 Assume that $q \in(p-1, p)$ and let $u$ be a non-negative solution to (1.1), (1.2) with an initial condition $u_{0}$ satisfying (1.3). If $M_{\infty}=0$ in (1.4), then

$$
\begin{equation*}
\|u(t)\|_{\infty} \leq C_{11} t^{-(p-q) /(2 q-p)} \quad \text { for } \quad t>0 \tag{5.6}
\end{equation*}
$$

Observe that the assumptions $p>2$ and $q \in(p-1, p)$ imply that $2 q>p$ and $(p-q) /(2 q-$ $p)>0$.
Proof. Consider a non-negative function $\eta \in \mathcal{C}^{\infty}\left(\mathbb{R}^{N}\right)$ with compact support in $B(0,1)$ and $\|\eta\|_{1}=1$. We then define a sequence of mollifiers $\left(\eta_{\delta}\right)_{\delta}$ by $\eta_{\delta}(x):=\eta(x / \delta) / \delta^{N}$ for $x \in \mathbb{R}^{N}$
and $\delta \in(0,1)$. For $\left(t, x_{0}\right) \in Q_{\infty}$ and $T>t$, we take $\vartheta(x)=\eta_{\delta}\left(x-x_{0}\right)$ in (2.3) and infer from (1.11) (with $s=t / 2$ ) that

$$
\begin{aligned}
\|u(T)\|_{\infty} & \geq \int_{\mathbb{R}^{N}} u(T, x) \eta_{\delta}\left(x-x_{0}\right) d x \\
& \geq \int_{\mathbb{R}^{N}} u(t, x) \eta_{\delta}\left(x-x_{0}\right) d x-\int_{t}^{T} \int_{\mathbb{R}^{N}}|\nabla u(s, x)|^{p-2} \nabla u(s, x) \cdot \nabla \eta_{\delta}\left(t, x-x_{0}\right) d x d s \\
& \geq \int_{\mathbb{R}^{N}} u(t, x) \eta_{\delta}\left(x-x_{0}\right) d x-2^{p / q} \kappa_{1}\left\|u\left(\frac{t}{2}\right)\right\|_{\infty}^{(p-q) / q} \int_{t}^{T}(2 s-t)^{-p / q} d s \\
& \geq \int_{\mathbb{R}^{N}} u(t, x) \eta_{\delta}\left(x-x_{0}\right) d x-C_{12}\left\|u\left(\frac{t}{2}\right)\right\|_{\infty}^{(p-q) / q}\left(t^{(q-p) / q}-T^{(q-p) / q}\right) .
\end{aligned}
$$

Owing to the continuity of $u$, we may pass to the limit as $\delta \rightarrow 0$ in the above inequality and deduce that

$$
\|u(T)\|_{\infty} \geq u\left(t, x_{0}\right)-C_{12}\left\|u\left(\frac{t}{2}\right)\right\|_{\infty}^{(p-q) / q}\left(t^{(q-p) / q}-T^{(q-p) / q}\right)
$$

But the above inequality is valid for all $x_{0} \in \mathbb{R}^{N}$ and we thus end up with

$$
\|u(T)\|_{\infty} \geq\|u(t)\|_{\infty}-C_{12}\left\|u\left(\frac{t}{2}\right)\right\|_{\infty}^{(p-q) / q}\left(t^{(q-p) / q}-T^{(q-p) / q}\right) .
$$

Finally, as $q<p$, we may let $T \rightarrow \infty$ in the previous inequality and use the assumption $M_{\infty}=0$ to conclude that

$$
\|u(t)\|_{\infty} \leq C_{12}\left\|u\left(\frac{t}{2}\right)\right\|_{\infty}^{(p-q) / q} t^{(q-p) / q}
$$

or, equivalently, as $2 q>p$,

$$
t^{(p-q) /(2 q-p)}\|u(t)\|_{\infty} \leq C_{13}\left\{\left(\frac{t}{2}\right)^{(p-q) /(2 q-p)}\left\|u\left(\frac{t}{2}\right)\right\|_{\infty}\right\}^{(p-q) / q}
$$

for $t \geq 0$. Introducing

$$
A(t):=\sup _{s \in(0, t)}\left\{s^{(p-q) /(2 q-p)}\|u(s)\|_{\infty}\right\} \in[0, \infty), \quad t \geq 0
$$

we deduce from the previous inequality that $A(t) \leq C_{13} A(t)^{(p-q) / q}$, whence $A(t) \leq C_{13}^{q /(2 q-p)}$ for $t \geq 0$. This bound being valid for each $t>0$, the proof of (5.6) is complete.

Proof of Proposition 1.4: $q \in\left(p-1, q_{\star}\right]$.

Step 1: We first consider a compactly supported initial condition $u_{0}$ satisfying (1.3) and assume for contradiction that $M_{\infty}=0$. On the one hand, according to Lemma 5.2 and the assumption $q \leq q_{\star}$, there holds

$$
\begin{equation*}
\limsup _{t \rightarrow \infty} t^{N /(N(p-2)+p)}\|u(t)\|_{\infty} \leq C_{11} t^{\left(p(N+1)\left(q-q_{\star}\right)\right) /((2 q-p)(N(p-2)+p))} \leq C_{11} \tag{5.7}
\end{equation*}
$$

On the other hand, fix $t_{0}>0$ and let $\varphi$ be the solution to the $p$-Laplacian equation $\partial_{t} \varphi-$ $\Delta_{p} \varphi=0$ in $Q_{\infty}$ with initial condition $\varphi(0)=u\left(t_{0}\right)$. As $u_{0}$ is compactly supported, so is $u\left(t_{0}\right)$ by Proposition 2.2 and $u\left(t_{0}\right)$ thus belongs to $L^{1}\left(\mathbb{R}^{N}\right)$. Moreover, the comparison principle warrants that $u(t, x) \geq \varphi\left(t-t_{0}, x\right)$ for $(t, x) \in\left[t_{0}, \infty\right) \times \mathbb{R}^{N}$. We then infer from the above properties and Proposition 2.3 that, for $t>t_{0}$,

$$
\begin{aligned}
t^{N /(N(p-2)+p)}\|u(t)\|_{\infty} \geq & \left(t-t_{0}\right)^{N /(N(p-2)+p)}\|u(t)\|_{\infty} \\
\geq & \left(t-t_{0}\right)^{N /(N(p-2)+p)}\left\|\varphi\left(t-t_{0}\right)\right\|_{\infty} \\
\geq & \left(t-t_{0}\right)^{N /(N(p-2)+p)}\left\|\mathcal{B}_{\left\|u\left(t_{0}\right)\right\|_{1}}\left(t-t_{0}\right)\right\|_{\infty} \\
& -\left(t-t_{0}\right)^{N /(N(p-2)+p)}\left\|\mathcal{B}_{\left\|u\left(t_{0}\right)\right\|_{1}}\left(t-t_{0}\right)-\varphi\left(t-t_{0}\right)\right\|_{\infty} \\
\geq & C_{14}\left\|u\left(t_{0}\right)\right\|_{1}^{p /(N(p-2)+p)} \\
& -\left(t-t_{0}\right)^{N /(N(p-2)+p)}\left\|\mathcal{B}_{\left\|u\left(t_{0}\right)\right\|_{1}}\left(t-t_{0}\right)-\varphi\left(t-t_{0}\right)\right\|_{\infty}
\end{aligned}
$$

Using once more Proposition 2.3, we may pass to the limit as $t \rightarrow \infty$ in the previous inequality to obtain

$$
\begin{equation*}
\liminf _{t \rightarrow \infty} t^{N /(N(p-2)+p)}\|u(t)\|_{\infty} \geq C_{14}\left\|u\left(t_{0}\right)\right\|_{1}^{p /(N(p-2)+p)} \tag{5.8}
\end{equation*}
$$

Combining (5.7) and (5.8) yields $\left\|u\left(t_{0}\right)\right\|_{1} \leq C_{15}$ for all $t_{0}>0$ which contradicts Lemma 5.1. Therefore, $M_{\infty}>0$.

Step 2: Now, if $u_{0}$ is an arbitrary initial condition satisfying (1.3), there clearly exists a compactly supported initial condition $\tilde{u}_{0}$ satisfying (1.3) and such that $u_{0} \geq \tilde{u}_{0}$ in $\mathbb{R}^{N}$. Introducing the solution $\tilde{u}$ to (1.1) with initial condition $\tilde{u}_{0}$, the comparison principle entails that $u \geq \tilde{u}$ in $Q_{\infty}$, hence

$$
M_{\infty} \geq \lim _{t \rightarrow \infty}\|\tilde{u}(t)\|_{\infty}
$$

The first step of the proof ensures that the right-hand side of the above inequality is positive which completes the proof.

It remains to investigate the case $q \in\left(q_{\star}, p\right)$, for which we adapt the proof of $[7$, Theorem 2.4(b)].
Proof of Proposition 1.4: $q \in\left(q_{\star}, p\right)$. We put

$$
m_{0}:=\left|\inf _{y \in \mathbb{R}^{N}} \Delta_{p} u_{0}(y)\right|
$$

As in the proof of Lemma 5.1, let $\eta \in \mathcal{C}^{\infty}\left(\mathbb{R}^{N}\right)$ be a non-negative function with compact support in $B(0,1)$ and $\|\eta\|_{1}=1$, and define a sequence of mollifiers $\left(\eta_{\delta}\right)_{\delta}$ by $\eta_{\delta}(x):=$ $\eta(x / \delta) / \delta^{N}$ for $x \in \mathbb{R}^{N}$ and $\delta \in(0,1)$. For $\left(t, x_{0}\right) \in Q_{\infty}$ and $T \in(0, t)$, we take $\vartheta(x)=$ $\eta_{\delta}\left(x-x_{0}\right)$ in (2.3) and infer from (1.11) (with $s=0$ ) and Corollary 3.3 that

$$
\begin{aligned}
\|u(t)\|_{\infty} & \geq \int_{\mathbb{R}^{N}} u(t, x) \eta_{\delta}\left(x-x_{0}\right) d x \\
& \geq \int_{\mathbb{R}^{N}} u_{0}(x) \eta_{\delta}\left(x-x_{0}\right) d x-\int_{0}^{t} \int_{\mathbb{R}^{N}}|\nabla u(s, x)|^{p-2} \nabla u(s, x) \cdot \nabla \eta_{\delta}\left(t, x-x_{0}\right) d x d s \\
& \geq \int_{\mathbb{R}^{N}} u_{0}(x) \eta_{\delta}\left(x-x_{0}\right) d x-\int_{0}^{T} m_{0} d s-\kappa_{1}\left\|u_{0}\right\|_{\infty}^{(p-q) / q} \int_{T}^{t} s^{-p / q} d s \\
& \geq \int_{\mathbb{R}^{N}} u_{0}(x) \eta_{\delta}\left(x-x_{0}\right) d x-T m_{0}-C_{16}\left\|u_{0}\right\|_{\infty}^{(p-q) / q}\left(T^{(q-p) / q}-t^{(q-p) / q}\right)
\end{aligned}
$$

Owing to the continuity of $u_{0}$, we may pass to the limit as $\delta \rightarrow 0$ in the above inequality and deduce that

$$
\|u(t)\|_{\infty} \geq u_{0}\left(x_{0}\right)-T m_{0}-C_{16}\left\|u_{0}\right\|_{\infty}^{(p-q) / q}\left(T^{(q-p) / q}-t^{(q-p) / q}\right)
$$

Since $q<p$, we may let $t \rightarrow \infty$ in the above inequality and take the supremum with respect to $x_{0}$ to conclude that

$$
M_{\infty} \geq\left\|u_{0}\right\|_{\infty}-T m_{0}-C_{16}\left\|u_{0}\right\|_{\infty}^{(p-q) / q} T^{(q-p) / q} .
$$

Next, for $\beta \in(0,1)$, the choice $T=\left\|u_{0}\right\|_{\infty}^{(p-q) / p}\left(\beta+m_{0}\right)^{-q / p}$ in the previous inequality yields

$$
M_{\infty} \geq\left\|u_{0}\right\|_{\infty}^{(p-q) / p}\left(\left\|u_{0}\right\|_{\infty}^{q / p}-\left(1+C_{16}\right)\left(\beta+m_{0}\right)^{(p-q) / p}\right) .
$$

This inequality being valid for every $\beta \in(0,1)$, we conclude that

$$
M_{\infty} \geq\left\|u_{0}\right\|_{\infty}^{(p-q) / p}\left(\left\|u_{0}\right\|_{\infty}^{q / p}-\left(1+C_{16}\right) m_{0}^{(p-q) / p}\right)>0
$$

as soon as (1.10) is fulfilled with $\kappa_{0}=\left(1+C_{16}\right)^{p / q}$.
Acknowledgments. I thank the referee for carefully reading the manuscript.

## A Convergence for the Hamilton-Jacobi equation (4.6)

In this section, we study the large behaviour of non-negative solutions to the Hamilton-Jacobi equation (4.6) with initial data in $\mathcal{C}_{0}\left(\mathbb{R}^{N}\right)$ and show their convergence to a steady state uniquely determined by the $L^{\infty}$-norm of the initial data. Though the large time behaviour of solutions to first-order Hamilton-Jacobi equations has received considerable attention in recent years (see $[4,5,16,22,23]$ and the references therein), the particular case of (4.6) does not seem to have been investigated in the literature. We thus provide a simple proof relying on the Hopf-Lax-Oleinik formula.

Proposition A. 1 Let $q>1$ and consider a non-negative function $h_{0} \in \mathcal{C}_{0}\left(\mathbb{R}^{N}\right)$. Let $h$ be the unique viscosity solution to the Cauchy problem

$$
\begin{align*}
\partial_{\tau} h-y \cdot \nabla h-q|\nabla h|^{q} & =0, \quad(\tau, y) \in(0, \infty) \times \mathbb{R}^{N},  \tag{A.1}\\
h(0) & =h_{0}, \quad y \in \mathbb{R}^{N} . \tag{A.2}
\end{align*}
$$

Then

$$
\begin{equation*}
\lim _{\tau \rightarrow \infty}\left\|h(\tau)-h_{s}\right\|_{\infty}=0 \tag{A.3}
\end{equation*}
$$

with

$$
h_{s}(y):=\left(\left\|h_{0}\right\|_{\infty}-\gamma_{q}|y|^{q /(q-1)}\right)_{+}, \quad y \in \mathbb{R}^{N}
$$

the constant $\gamma_{q}=(q-1) q^{-q /(q-1)}$ being defined in Theorem 1.1.
Thanks to the concavity of the Hamiltonian $\mathcal{H}(y, \xi)=-y \cdot \xi-q|\xi|^{q},(y, \xi) \in \mathbb{R}^{N} \times \mathbb{R}^{N}$, with respect to its second variable, the Hopf-Lax-Oleinik formula provides a representation formula for the solution $h$ to (A.1), (A.2) which can be used to prove (A.3).
Proof. We first recall that $h$ is given by the Hopf-Lax-Oleinik formula

$$
h(\tau, y)=\sup _{z \in \mathbb{R}^{N}}\left\{h_{0}(z)-\gamma_{q}\left|y-z e^{-\tau}\right|^{q /(q-1)}\left(1-e^{-q \tau}\right)^{-1 /(q-1)}\right\}
$$

for $(\tau, y) \in[0, \infty) \times \mathbb{R}^{N}$, see, e.g., [10, Chapter 3]. Since $h(\tau, y) \geq h_{0}\left(y e^{\tau}\right) \geq 0$, we have in fact

$$
h(\tau, y)=\sup _{z \in \mathbb{R}^{N}}\left\{\left(h_{0}(z)-\gamma_{q}\left|y-z e^{-\tau}\right|^{q /(q-1)}\left(1-e^{-q \tau}\right)^{-1 /(q-1)}\right)_{+}\right\}
$$

for $(\tau, y) \in[0, \infty) \times \mathbb{R}^{N}$.
Consider now $\beta \in(0,1)$. As $h_{0} \in \mathcal{C}_{0}\left(\mathbb{R}^{N}\right)$, there is $R_{\beta}>\left(\left\|h_{0}\right\|_{\infty} / \gamma_{q}\right)^{(q-1) / q}$ such that

$$
\begin{equation*}
h_{0}(z) \leq \beta \quad \text { for } \quad|z| \geq R_{\beta} . \tag{A.4}
\end{equation*}
$$

On the one hand, if $(\tau, y) \in\left[\log R_{\beta}, \infty\right) \times \mathbb{R}^{N}$ and $z \in \mathbb{R}^{N}$, we have either $|z| \geq R_{\beta}$ and

$$
\begin{aligned}
& \left|\left(h_{0}(z)-\gamma_{q}\left|y-z e^{-\tau}\right|^{q /(q-1)}\left(1-e^{-q \tau}\right)^{-1 /(q-1)}\right)_{+}-\left(h_{0}(z)-\gamma_{q}|y|^{q /(q-1)}\right)_{+}\right| \\
\leq & \left(h_{0}(z)-\gamma_{q}\left|y-z e^{-\tau}\right|^{q /(q-1)}\left(1-e^{-q \tau}\right)^{-1 /(q-1)}\right)_{+}+\left(h_{0}(z)-\gamma_{q}|y|^{q /(q-1)}\right)_{+} \\
\leq & 2 \beta
\end{aligned}
$$

by (A.4) or $z \in B\left(0, R_{\beta}\right)$ and

$$
\begin{aligned}
& \left|\left(h_{0}(z)-\gamma_{q}\left|y-z e^{-\tau}\right|^{q /(q-1)}\left(1-e^{-q \tau}\right)^{-1 /(q-1)}\right)_{+}-\left(h_{0}(z)-\gamma_{q}|y|^{q /(q-1)}\right)_{+}\right| \\
\leq & \gamma_{q}\left|y-z e^{-\tau}\right|^{q /(q-1)}\left\{\left(1-e^{-q \tau}\right)^{-1 /(q-1)}-1\right\}+\gamma_{q}| | y-\left.z e^{-\tau}\right|^{q /(q-1)}-|y|^{q /(q-1)} \mid \\
\leq & \gamma_{q}\left(|y|+R_{\beta} e^{-\tau}\right)^{q /(q-1)}\left\{\left(1-e^{-q \tau}\right)^{-1 /(q-1)}-1\right\}+\frac{q \gamma_{q}}{q-1}\left(|y|+|z| e^{-\tau}\right)^{1 /(q-1)}|z| e^{-\tau} \\
\leq & \gamma_{q}(|y|+1)^{1 /(q-1)}\left\{\frac{q}{q-1}+|y|+1\right\}\left\{\left(1-e^{-q \tau}\right)^{-1 /(q-1)}-1+R_{\beta} e^{-\tau}\right\}
\end{aligned}
$$

as $\tau \geq \log R_{\beta}$. Combining the above two estimates gives

$$
\begin{aligned}
& \left|h(\tau, y)-\sup _{z \in \mathbb{R}^{N}}\left\{\left(h_{0}(z)-\gamma_{q}|y|^{q /(q-1)}\right)_{+}\right\}\right| \\
\leq & C(q)(|y|+1)^{q /(q-1)}\left\{\left(1-e^{-q \tau}\right)^{-1 /(q-1)}-1+R_{\beta} e^{-\tau}\right\}+2 \beta
\end{aligned}
$$

whence

$$
\begin{equation*}
\left|h(\tau, y)-h_{s}(y)\right| \leq C(q)(|y|+1)^{q /(q-1)}\left\{\left(1-e^{-q \tau}\right)^{-1 /(q-1)}-1+R_{\beta} e^{-\tau}\right\}+2 \beta \tag{A.5}
\end{equation*}
$$

for $(\tau, y) \in\left[\log R_{\beta}, \infty\right) \times \mathbb{R}^{N}$. On the other hand, if $\tau \geq \log \left(R_{\beta}\right),|y| \geq Y:=1+$ $\left(\left\|h_{0}\right\|_{\infty} / \gamma_{q}\right)^{(q-1) / q}$ and $z \in \mathbb{R}^{N}$, we have either $\left|y-z e^{-\tau}\right| \geq Y-1$ and

$$
\begin{aligned}
& h_{0}(z)-\gamma_{q}\left|y-z e^{-\tau}\right|^{q /(q-1)}\left(1-e^{-q \tau}\right)^{-1 /(q-1)} \\
\leq & \left(1-e^{-q \tau}\right)^{-1 /(q-1)}\left\{\left\|h_{0}\right\|_{\infty}\left(1-e^{-q \tau}\right)^{1 /(q-1)}-\gamma_{q}\left|y-z e^{-\tau}\right|^{q /(q-1)}\right\} \\
\leq & \left(1-e^{-q \tau}\right)^{-1 /(q-1)}\left\{\left\|h_{0}\right\|_{\infty}-\gamma_{q}(Y-1)^{q /(q-1)}\right\} \\
\leq & 0,
\end{aligned}
$$

or $\left|y-z e^{-\tau}\right|<Y-1$ and

$$
|z| \geq\left|y e^{\tau}\right|-\left|z-y e^{\tau}\right| \geq Y e^{\tau}-(Y-1) e^{\tau}=e^{\tau} \geq R_{\beta},
$$

so that

$$
h_{0}(z)-\gamma_{q}\left|y-z e^{-\tau}\right|^{q /(q-1)}\left(1-e^{-q \tau}\right)^{-1 /(q-1)} \leq \beta
$$

by (A.4). Therefore,

$$
\begin{equation*}
h(\tau, y) \leq \beta \quad \text { for } \quad(\tau, y) \in\left[\log R_{\beta}, \infty\right) \times\left(\mathbb{R}^{N} \backslash B(0, Y)\right) \tag{A.6}
\end{equation*}
$$

The claim (A.3) then easily follows from (A.5) and (A.6).

## References

[1] Martino Bardi and Italo Capuzzo-Dolcetta, Optimal Control and Viscosity Solutions of Hamilton-Jacobi-Bellman Equations, Systems Control Found. Appl., Birkhäuser, Boston, 1997.
[2] Guy Barles, Solutions de Viscosité des Equations d'Hamilton-Jacobi, Mathématiques \& Applications 17, Springer-Verlag, Berlin, 1994.
[3] Guy Barles and Benoît Perthame, Exit time problems in optimal control and vanishing viscosity method, SIAM J. Control Optim. 26 (1988), 1133-1148.
[4] Guy Barles and Jean-Michel Roquejoffre, Ergodic type problems and large time behaviour of unbounded solutions of Hamilton-Jacobi equations, Comm. Partial Differential Equations 31 (2006), 1209-1225.
[5] Guy Barles and Panagiotis E. Souganidis, On the large time behavior of solutions of Hamilton-Jacobi equations, SIAM J. Math. Anal. 31 (2000), 925-939.
[6] Jean-Philippe Bartier and Philippe Laurençot, Gradient estimates for a degenerate parabolic equation with gradient absorption and applications, J. Funct. Anal. 254 (2008), 851-878.
[7] Saïd Benachour, Grzegorz Karch, and Philippe Laurençot, Asymptotic profiles of solutions to viscous Hamilton-Jacobi equations, J. Math. Pures Appl. (9) 83 (2004), 12751308.
[8] Michael G. Crandall, Hitoshi Ishii, and Pierre-Louis Lions, User's guide to viscosity solutions of second order partial differential equations, Bull. Amer. Math. Soc. (N.S.) 27 (1992), 1-67.
[9] Juan R. Esteban and Pierangelo Marcati, Approximate solutions to first and second order quasilinear evolution equations via nonlinear viscosity, Trans. Amer. Math. Soc. 342 (1994), 501-521.
[10] Lawrence C. Evans, Partial Differential Equations, Grad. Stud. Math. 19, Amer. Math. Soc., Providence, RI, 1998.
[11] Yoshikazu Giga, Shun'ichi Goto, Hitoshi Ishii, and Moto-Hiko Sato, Comparison principle and convexity preserving properties for singular degenerate parabolic equations on unbounded domains, Indiana Univ. Math. J. 40 (1991), 443-470.
[12] Brian H. Gilding, The Cauchy problem for $u_{t}=\Delta u+|\nabla u|^{q}$, large-time behaviour, J. Math. Pures Appl. (9) 84 (2005), 753-785.
[13] Brian H. Gilding and Robert Kersner, Travelling Waves in Nonlinear DiffusionConvection Reaction, Progr. Nonlinear Differential Equations Appl. 60, Birkhäuser, Basel, 2004.
[14] Richard S. Hamilton, A matrix Harnack estimate for the heat equation, Comm. Anal. Geom. 1 (1993), 113-126.
[15] Hitoshi Ishii, A simple, direct proof of uniqueness for solutions of the Hamilton-Jacobi equations of eikonal type, Proc. Amer. Math. Soc. 100 (1987), 247-251.
[16] Hitoshi Ishii, Asymptotic solutions for large time of Hamilton-Jacobi equations in Euclidean $n$ space, Ann. Inst. H. Poincaré Anal. Non Linéaire 25 (2008), 231-266.
[17] Shoshana Kamin and Juan Luis Vázquez, Fundamental solutions and asymptotic behaviour for the p-Laplacian equation, Rev. Mat. Iberoamericana 4 (1988), 339-354.
[18] Philippe Laurençot and Philippe Souplet, On the growth of mass for a viscous HamiltonJacobi equation, J. Anal. Math. 89 (2003), 367-383.
[19] Philippe Laurençot and Philippe Souplet, Optimal growth rates for a viscous HamiltonJacobi equation, J. Evolution Equations 5 (2005), 123-135.
[20] Philippe Laurençot and Juan Luis Vázquez, Localized non-diffusive asymptotic patterns for nonlinear parabolic equations with gradient absorption, J. Dynamics Differential Equations 19 (2007), 985-1005.
[21] Chi-Tien Lin and Eitan Tadmor, $L^{1}$-stability and error estimates for approximate Hamilton-Jacobi solutions, Numer. Math. 87 (2001), 701-735.
[22] Gawtum Namah and Jean-Michel Roquejoffre, Remarks on the long time behaviour of the solutions of Hamilton-Jacobi equations, Comm. Partial Differential Equations 24 (1999), 883-893.
[23] Jean-Michel Roquejoffre, Convergence to steady states or periodic solutions in a class of Hamilton-Jacobi equations, J. Math. Pures Appl. (9) 80 (2001), 85-104.
[24] Thomas Strömberg, The Hopf-Lax formula gives the unique viscosity solution, Differential Integral Equations 15 (2002), 47-52.


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