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# Non-diffusive large time behaviour for a degenerate viscous Hamilton-Jacobi equation

Philippe Laurençot\*

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#### Abstract

The convergence to non-diffusive self-similar solutions is investigated for non-negative solutions to the Cauchy problem  $\partial_t u = \Delta_p u + |\nabla u|^q$  when the initial data converge to zero at infinity. Sufficient conditions on the exponents p > 2 and q > 1 are given that guarantee that the diffusion becomes negligible for large times and the  $L^{\infty}$ -norm of u(t) converges to a positive value as  $t \to \infty$ .

#### 1 Introduction

The quasilinear degenerate parabolic equation

(1.1)  $\partial_t u = \Delta_p u + |\nabla u|^q , \quad (t,x) \in Q_\infty := (0,\infty) \times \mathbb{R}^N,$ 

includes two competing mechanisms acting on the space variable x, a degenerate diffusion  $\Delta_p u$  involving the *p*-Laplacian operator defined by

$$\Delta_p u := \operatorname{div} \left( |\nabla u|^{p-2} \nabla u \right), \quad p > 2,$$

and a source term  $|\nabla u|^q$ , q > 1, depending solely on the gradient of u. The aim of this work is to identify a range of the parameters p and q for which the large time behaviour of non-negative solutions to (1.1) is dominated by the source term. More precisely, we consider the Cauchy problem and supplement (1.1) with the initial condition

(1.2) 
$$u(0) = u_0 \ge 0, \quad x \in \mathbb{R}^N.$$

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Throughout the paper, the initial condition  $u_0$  is assumed to fulfill

(1.3) 
$$u_0 \in \mathcal{C}_0(\mathbb{R}^N) \cap W^{1,\infty}(\mathbb{R}^N), \quad u_0 \ge 0, \quad u_0 \ne 0,$$

where

$$\mathcal{C}_0(\mathbb{R}^N) := \left\{ w \in \mathcal{BC}(\mathbb{R}^N) : \lim_{R \to \infty} \sup_{\{|x| \ge R\}} \left\{ |w(x)| \right\} = 0 \right\}$$

and  $\mathcal{BC}(\mathbb{R}^N) := \mathcal{C}(\mathbb{R}^N) \cap L^{\infty}(\mathbb{R}^N).$ 

For such an initial condition, the Cauchy problem (1.1), (1.2) has a unique non-negative (viscosity) solution  $u \in \mathcal{BC}([0,\infty) \times \mathbb{R}^N)$  (see Proposition 2.1 below). Moreover,  $t \mapsto ||u(t)||_{\infty}$  is a non-increasing function and has a limit  $M_{\infty} \in [0, ||u_0||_{\infty}]$  as  $t \to \infty$ . Our main result is then the following:

**Theorem 1.1** Assume that p > 2 and  $q \in (1, p)$ . Consider a non-negative function  $u_0$  satisfying (1.3) and let u be the corresponding (viscosity) solution to (1.1), (1.2). Assume further that

(1.4) 
$$M_{\infty} := \lim_{t \to \infty} \|u(t)\|_{\infty} > 0.$$

(1.5) 
$$\lim_{t \to \infty} \|u(t) - h_{\infty}(t)\|_{\infty} = 0,$$

where  $h_{\infty}$  is given by

(1.6) 
$$h_{\infty}(t,x) := H_{\infty}\left(\frac{x}{t^{1/q}}\right) \quad and \quad H_{\infty}(x) := \left(M_{\infty} - \gamma_q \ |x|^{q/(q-1)}\right)_+$$

for  $(t, x) \in Q_{\infty}$  and  $\gamma_q := (q - 1) q^{-q/(q-1)}$ .

Here and below,  $r_+ := \max\{r, 0\}$  denotes the positive part of the real number r.

The convergence (1.5) clearly indicates that the large time behaviour of non-negative solutions to (1.1), (1.2) fulfilling the condition (1.4) is governed by the gradient source term. Indeed,  $h_{\infty}$  is actually a self-similar solution to the Hamilton-Jacobi equation

(1.7) 
$$\partial_t h = |\nabla h|^q , \quad (t, x) \in Q_\infty,$$

and an alternative formula for  $h_{\infty}$  reads

(1.8) 
$$h_{\infty}(t,x) = \sup_{y \in \mathbb{R}^{N}} \left\{ M_{\infty} \ \mathbf{1}_{\{0\}}(y) - \gamma_{q} \ \frac{|x-y|^{q/(q-1)}}{t^{1/(q-1)}} \right\}$$

for  $(t, x) \in [0, \infty) \times \mathbb{R}^N$ ,  $\mathbf{1}_{\{0\}}$  denoting the indicator function of the singleton set  $\{0\}$ . The formula (1.8) is the well-known Hopf-Lax-Oleinik representation formula for viscosity solutions to (1.7) (see, e.g., [10, Chapter 3]) and  $h_{\infty}$  turns out to be the unique viscosity solution in  $\mathcal{BUC}(Q_{\infty})$  to (1.7) with the bounded and upper semicontinuous initial condition  $h_{\infty}(0, x) = \mathbf{1}_{\{0\}}(x)$  for  $x \in \mathbb{R}^N$  [24]. **Remark 1.2** The convergence (1.5) also holds true for the viscosity solution to the Hamilton-Jacobi equation (1.7) with a non-negative initial condition  $u_0 \in C_0(\mathbb{R}^N)$  but with  $||u_0||_{\infty}$  instead of  $M_{\infty}$  in the formula (1.6) giving  $H_{\infty}$ . For (1.1), (1.2), the constant  $M_{\infty}$  takes into account that, though negligible for large times, the diffusion erodes the supremum of u during the time evolution.

For p = 2, Theorem 1.1 is also valid and is proved in [7], the proof relying on a rescaling technique: The crucial step is then to identify the possible limits of the rescaled sequence and this is done by an extensive use of the Hopf-Lax-Oleinik representation formula. The proof we perform here is of a completely different nature and relies on the relaxed halflimits method introduced in [3]. A similar approach has been used in [22] and [23] to investigate the large time behaviour of solutions to first-order Hamilton-Jacobi equations  $\partial_t w + H(x, \nabla w) = 0$  in  $Q_{\infty}$ . It has also been used in [20] to study the convergence to nondiffusive localized self-similar patterns for non-negative and compactly supported solutions to  $\partial_t w - \Delta_p w + |\nabla w|^q = 0$  in  $Q_{\infty}$  when p > 2 and  $q \in (1, p - 1)$ .

In order to apply Theorem 1.1, one should check whether there are non-negative solutions to (1.1), (1.2) for which (1.4) holds true. The next result provides sufficient conditions for (1.4) to be fulfilled.

**Theorem 1.3** Assume that p > 2 and q > 1. Consider a non-negative function  $u_0$  satisfying (1.3) and let u be the corresponding solution to (1.1), (1.2). Introducing

(1.9) 
$$q_{\star} := p - \frac{N}{N+1},$$

then u fulfills (1.4) if

- (a) either  $q \in (1, q_{\star}]$ ,
- (b) or  $q \in (q_{\star}, p)$ ,  $u_0 \in W^{2,\infty}(\mathbb{R}^N)$ , and

(1.10) 
$$||u_0||_{\infty} > \kappa_0 \left| \inf_{y \in \mathbb{R}^N} \left\{ \Delta_p u_0(y) \right\} \right|^{(p-q)/q}$$

for some  $\kappa_0 > 0$  which depends only on N, p, and q.

A similar result is already available for p = 2 and has been established in [7, 12]. It is also shown in [7, 12, 18] that, still for p = 2,  $||u(t)||_{\infty} \to 0$  as  $t \to \infty$  if either  $q \ge p$  or  $q \in (q_{\star}, p)$  and a suitable norm of  $u_0$  is small enough. A similar result is likely to be true also for p > 2, so that different large time behaviours are only to be found in the intermediate range  $q \in (q_{\star}, p)$ .

The proof of Theorem 1.3 for  $q \in (p-1, p)$  and p > 2 borrows some steps from the case p = 2. However, it relies on semiconvexity estimates for solutions to (1.1), (1.2) which seem to be new for p > 2 and  $q \in (1, p)$  and are stated now.

**Proposition 1.4** Assume that p > 2 and  $q \in (1, p]$ . Let u be the viscosity solution to (1.1), (1.2) with initial condition  $u_0 \in \mathcal{BUC}(\mathbb{R}^N)$  (that is,  $u_0 \in \mathcal{BC}(\mathbb{R}^N)$  and is uniformly continuous in  $\mathbb{R}^N$ ). Then  $\nabla u(t)$  belongs to  $L^{\infty}(\mathbb{R}^N)$  for each t > 0 and there is  $\kappa_1 > 0$  depending only on N, p, and q such that

(1.11) 
$$\Delta_p u(t,x) \ge -\kappa_1 \| u(s) \|_{\infty}^{(p-q)/q} (t-s)^{-p/q}, \quad t > s \ge 0,$$

in the sense of distributions. In addition, if  $u_0 \in W^{1,\infty}(\mathbb{R}^N)$ , there holds

(1.12) 
$$\Delta_p u(t,x) \ge -\frac{N(p-1)}{q(q-1)} \frac{\|\nabla u_0\|_{\infty}^{p-q}}{t}$$

for t > 0 in the sense of distributions.

The proof of Proposition 1.4 relies on the comparison principle combined with a gradient estimate established in [6].

Similar semiconvexity estimates for solutions to (1.1), (1.2) have already been obtained in [14] and [21, Lemma 5.1] for p = q = 2, in [7, Proposition 3.2] for p = 2 and  $q \in (1, 2]$ , and in [9, Theorem 1] for p = q > 2. We extend these results to the range p > 2 and  $q \in (1, p]$ . As we shall see below, the estimate (1.11) plays an important role in the proof of Theorem 1.3 and is also helpful to construct a subsolution in the proof of Theorem 1.1.

Let us finally emphasize that the validity of Proposition 1.4 is not restricted to non-negative solutions and that the solutions to the Hamilton-Jacobi equation (1.7) also enjoy the semiconvexity estimates (1.11) and (1.12). These two estimates thus stem from the reaction term  $|\nabla u|^q$  and not from the diffusion.

In the next section, we recall the well-posedness of (1.1), (1.2) in  $\mathcal{BUC}(\mathbb{R}^N)$ , as well as some properties of the solutions established in [6]. We also show the finite speed of propagation of the support for non-negative compactly supported initial data. Section 3 is devoted to the proof of the semiconvexity estimates (Proposition 1.4) and Section 4 to that of Theorem 1.1. Theorem 1.3 is shown in the last section, its proof combining arguments of [7, 12, 19] used to established analogous results when p = 2.

Throughout the paper, C and  $C_i$ ,  $i \ge 1$ , denote positive constants depending only on p, q, and N. Dependence upon additional parameters will be indicated explicitly. Also,  $\mathcal{M}_N(\mathbb{R})$ denotes the space of real-valued  $N \times N$  matrices and  $\delta_{ij} = 1$  if i = j and  $\delta_{ij} = 0$  if  $i \ne j$ ,  $1 \le i, j \le N$ . Given a matrix  $A = (a_{ij}) \in \mathcal{M}_N(\mathbb{R})$ , tr(A) denotes its trace and is given by  $tr(A) := \sum a_{ii}$ .

#### 2 Preliminary results

Let us first recall the well-posedness (in the framework of viscosity solutions) of (1.1), (1.2), together with some properties of the solutions established in [6].

**Proposition 2.1** Consider a non-negative initial condition  $u_0 \in \mathcal{BUC}(\mathbb{R}^N)$ . There is a unique non-negative viscosity solution  $u \in \mathcal{BC}([0,\infty) \times \mathbb{R}^N)$  to (1.1), (1.2) such that

(2.1) 
$$0 \le u(t,x) \le ||u_0||_{\infty}, \quad (t,x) \in Q_{\infty},$$

(2.2) 
$$\|\nabla u(t)\|_{\infty} \leq \min \left\{ C_1 \|u(s)\|_{\infty}^{1/q} (t-s)^{-1/q}, \|\nabla u(s)\|_{\infty} \right\},$$

and

(2.3) 
$$\int_{\mathbb{R}^N} (u(t,x) - u(s,x)) \,\vartheta(x) \,dx + \int_s^t \int_{\mathbb{R}^N} \left( |\nabla u|^{p-2} \nabla u \cdot \nabla \vartheta - |\nabla u|^q \,\vartheta \right) \,dxd\tau = 0$$

for  $t > s \ge 0$  and  $\vartheta \in \mathcal{C}_0^{\infty}(\mathbb{R}^N)$ . In addition,  $t \longmapsto \|u(t)\|_{\infty}$  is a non-increasing function.

**Proof.** We put  $\tilde{u}_0 := ||u_0||_{\infty} - u_0$ . As  $\tilde{u}_0$  is a non-negative function in  $\mathcal{BUC}(\mathbb{R}^N)$ , it follows from [6, Theorem 1.1] that there is a unique non-negative viscosity solution  $\tilde{u}$  to

(2.4) 
$$\partial_t \tilde{u} - \Delta_p \tilde{u} + |\nabla \tilde{u}|^q = 0 \quad , \quad (t, x) \in Q_\infty \, ,$$

with initial condition  $\tilde{u}(0,x) = \tilde{u}_0(x)$  for  $x \in \mathbb{R}^N$ . It also satisfies  $0 \leq \tilde{u}(t,x) \leq ||u_0||_{\infty}$  and

$$\int_{\mathbb{R}^N} (\tilde{u}(t,x) - \tilde{u}(s,x)) \ \vartheta(x) \ dx + \int_s^t \int_{\mathbb{R}^N} \left( |\nabla \tilde{u}|^{p-2} \nabla \tilde{u} \cdot \nabla \vartheta + |\nabla \tilde{u}|^q \ \vartheta \right) \ dx d\tau = 0$$

for  $t > s \ge 0$ ,  $x \in \mathbb{R}^N$ , and  $\vartheta \in \mathcal{C}_0^{\infty}(\mathbb{R}^N)$ . In addition,  $\nabla \tilde{u}(t)$  belongs to  $L^{\infty}(\mathbb{R}^N)$  for each t > 0 and

 $\|\nabla \tilde{u}(t)\|_{\infty} \le C_1 \|\tilde{u}_0\|_{\infty}^{1/q} t^{-1/q}$ 

by [6, Lemma 4.1]. Setting  $u := ||u_0||_{\infty} - \tilde{u}$ , we readily deduce from the properties of  $\tilde{u}$  that u is a non-negative viscosity solution to (1.1), (1.2) satisfying (2.1) and (2.3). Also,  $\nabla u(t)$  belongs to  $L^{\infty}(\mathbb{R}^N)$  for each t > 0. The uniqueness and the time monotonicity of  $||u||_{\infty}$  then both follow from the comparison principle, see [8] or [11, Theorem 2.1]. Finally, given  $s \ge 0$ ,  $(t, x) \mapsto ||u(s)||_{\infty} - u(t + s, x)$  is the unique non-negative viscosity solution to the Cauchy problem (2.4) with initial condition  $x \mapsto ||u(s)||_{\infty} - u(s, x)$  and we infer from [6, Lemma 4.1] that

$$\|\nabla u(t+s)\|_{\infty} \le C_1 \ \|\|u(s)\|_{\infty} - u(s)\|_{\infty}^{1/q} \ t^{-1/q} \le C_1 \ \|u(s)\|_{\infty}^{1/q} \ t^{-1/q}$$

for t > 0, whence (2.2).

We next turn to the propagation of the support of non-negative solutions to (1.1), (1.2) with non-negative compactly supported initial data.

**Proposition 2.2** Consider a non-negative solution u to (1.1), (1.2) with an initial condition  $u_0$  satisfying (1.3). Assume further that  $u_0$  is compactly supported in a ball  $B(0, R_0)$  of  $\mathbb{R}^N$  for some  $R_0 > 0$ . Then u(t) is compactly supported for each  $t \ge 0$ .

**Proof.** We argue by comparison with travelling wave solutions. By [13, Application 9.4], there is a travelling wave solution w to the convection-diffusion equation

(2.5) 
$$\partial_t w - \partial_1^2 \left( w^{p-1} \right) + \partial_1 \left( w^q \right) = 0, \quad (t, x_1) \in (0, \infty) \times \mathbb{R},$$

with wave speed unity. It is given by  $w(t, x_1) = f(x_1 - t)$  for  $(t, x_1) \in (0, \infty) \times \mathbb{R}$ , the function f being implicitly defined by

$$(p-1) \int_0^{f(y)} \frac{z^{p-3}}{1-z^{q-1}} \, dz = (-y)_+, \quad y \in \mathbb{R}.$$

In particular, f satisfies f(y) = 0 if y > 0 and  $f(y) \to 1$  as  $y \to -\infty$ . Introducing

$$F(y) := \int_{y}^{\infty} f(z) \, dz \, , \quad y \in \mathbb{R} \, ,$$

the properties of f ensure that F is a decreasing function on  $(-\infty, 0)$  with F(y) = 0 if y > 0,  $F(y) \leq |y|$  if y < 0, and  $F(y) \to \infty$  as  $y \to -\infty$ . There is therefore a unique  $\mu \in (-\infty, 0)$ such that  $F(R_0 + \mu) = ||u_0||_{\infty}$ . In addition, it readily follows from (2.5) and the invariance by translation of (1.1) that  $W_{\mu}(t, x) := F(x_1 + \mu - t)$  is a travelling wave solution to (1.1). Now, u and  $W_{\mu}$  are both solutions to (1.1) in  $(0, \infty) \times H_+$ , the half-space  $H_+$  being defined by  $H_+ := \{x \in \mathbb{R}^N : x_1 > R_0\}$ . Owing to the monotonicity of F, the bound  $0 \leq f \leq 1$ , and (2.1), we have also

$$u_0(x) - W_\mu(0, y) = 0 - W_\mu(0, y) \le W_\mu(0, x) - W_\mu(0, y) \le |x - y|$$

for  $x \in H_+$  and  $y \in H_+$ ,

$$\begin{aligned} u(t,x) - W_{\mu}(t,y) &\leq \|u_0\|_{\infty} - W_{\mu}(t,y) \\ &\leq F(R_0 + \mu - t) - W_{\mu}(t,y) = W_{\mu}(t,x) - W_{\mu}(t,y) \leq |x-y| \end{aligned}$$

for  $t > 0, x \in \partial H_+, y \in H_+$ , and

$$u(t,x) - W_{\mu}(t,y) \le ||u_0||_{\infty} - F(R_0 + \mu - t) \le 0$$

for t > 0,  $x \in H_+$ ,  $y \in \partial H_+$ . We are then in a position to use the comparison principle stated in [11, Theorem 2.1] to conclude that  $u(t,x) \leq W_{\mu}(t,x)$  for  $(t,x) \in (0,\infty) \times H_+$ . Consequently,  $u(t,x) \leq F(x_1 + \mu - t) = 0$  if  $t \geq 0$  and  $x_1 > \max\{R_0, t - \mu\}$ , and the rotational invariance of (1.1) allows us to conclude that u(t,x) = 0 for  $t \geq 0$  and  $|x| > \max\{R_0, t - \mu\}$ .

We finally recall the convergence to self-similar solutions for non-negative and compactly supported solutions to the p-Laplacian equation [17]

(2.6) 
$$\partial_t \varphi = \Delta_p \varphi \ , \ (t, x) \in Q_\infty .$$

**Proposition 2.3** Let  $\varphi_0$  be a non-negative and compactly supported function in  $L^1(\mathbb{R}^N)$  and  $\varphi$  denote the unique weak solution to (2.6) with initial condition  $\varphi_0$ . Then

(2.7) 
$$\lim_{t \to \infty} t^{(N(r-1))/(r(N(p-2)+p))} \|\varphi(t) - \mathcal{B}_{\|\varphi_0\|_1}(t)\|_r = 0 \quad for \quad r \in [1,\infty],$$

where  $\mathcal{B}_L$  denotes the Barenblatt solution to (2.6) given by

$$\mathcal{B}_{L}(t,x) := t^{-N/(N(p-2)+p)} b_{L} \left( x t^{-1/(N(p-2)+p)} \right) ,$$
  
$$b_{L}(x) := \left( C_{2} L^{(p(p-2))/((p-1)(N(p-2)+p))} - C_{3} |x|^{p/(p-1)} \right)_{+}^{(p-1)/(p-2)}$$

for  $(t, x) \in Q_{\infty}$  and L > 0.

The convergence (2.7) is proved in [17, Theorem 2] for  $r = \infty$ . As  $\varphi_0$  is compactly supported, so is  $\varphi(t)$  for each t > 0 and the support of  $\varphi(t)$  is included in  $B\left(0, C_4(\varphi_0) t^{1/(N(p-2)+p)}\right)$  for  $t \ge 1$  [17, Proposition 2.2]. Combining this property with [17, Theorem 2] readily provide the convergence (2.7) for all  $r \in [1, \infty)$ .

### 3 Semiconvexity

In this section, we prove the semiconvexity estimates (1.11) and (1.12). To this end, we would like to derive an equation for  $\Delta_p u$  to which we could apply the comparison principle. The poor regularity of u however does not allow to perform directly such a computation and an approximation procedure is needed. As a first step, we report the following result:

**Lemma 3.1** Let a and b be two non-negative function in  $\mathcal{C}^{\infty}([0,\infty))$  satisfying

(3.1) 
$$a(r) > 0, \quad a'(r) > 0, \quad a'(r) \ b'(r) - a(r) \ b''(r) > 0,$$

(3.2) 
$$c(r) := 2 \left(\frac{b'}{a}\right)(r) + \frac{4r (a \ b'' - a' \ b')(r)}{a^2(r) + 2r \ a(r) \ a'(r)} \ge 0.$$

Consider a classical solution v to

(3.3) 
$$\partial_t v - div \left( a \left( |\nabla v|^2 \right) |\nabla v \right) = b \left( |\nabla v|^2 \right), \quad (t, x) \in Q_{\infty},$$

and put

$$w := div \left( a \left( |\nabla v|^2 \right) |\nabla v \right) \quad and \quad z_i := a \left( |\nabla v|^2 \right) |\partial_i v$$

for  $i \in \{1, ..., N\}$ . Then

(3.4) 
$$\partial_t w - \mathcal{L}w - \mathcal{V} \cdot \nabla w - \frac{c \left( |\nabla v|^2 \right)}{N} \ w^2 \ge 0 \quad in \quad Q_\infty \,,$$

where

$$\mathcal{L}w := \sum_{i,j} \partial_i \left( a \left( |\nabla v|^2 \right) E_{ij} \partial_j w \right), \quad \mathcal{V} := 2 \ b' \left( |\nabla v|^2 \right) \ \nabla v,$$
$$E_{ij} := \delta_{ij} + 2 \ \frac{a'}{a} \left( |\nabla v|^2 \right) \ \partial_i v \ \partial_j v, \quad 1 \le i, j \le N.$$

The proof of Lemma 3.1 borrows some steps from the proof of [9, Theorem 1] for p = q > 2 but requires additional arguments to handle the term coming from the fact that  $q \neq p$ . In particular, we recall the following elementary result which will be helpful to estimate this term.

**Lemma 3.2** Let A and B be two symmetric matrices in  $\mathcal{M}_N(\mathbb{R})$  and put M := ABA. Then M is a symmetric matrix in  $\mathcal{M}_N(\mathbb{R})$  and

(3.5) 
$$|MX|^2 \le tr\left(M^2\right) \ |X|^2 \quad for \quad X \in \mathbb{R}^N.$$

Proof of Lemma 3.1. We first note that

(3.6) 
$$\partial_j z_i = a \left( |\nabla v|^2 \right) \sum_k E_{ik} \, \partial_k \partial_j v \,,$$

(3.7) 
$$\partial_t z_i = a \left( |\nabla v|^2 \right) \sum_k E_{ik} \, \partial_k \partial_t v \,,$$

for  $1 \leq i, j \leq N$ . According to the definition of w, we infer from (3.3), (3.6), and (3.7) that

$$\begin{aligned} \partial_{t}w &= \sum_{i,k} \partial_{i} \left( a \left( |\nabla v|^{2} \right) E_{ik} \partial_{k} \partial_{t} v \right) \\ &= \sum_{i,k} \partial_{i} \left( a \left( |\nabla v|^{2} \right) E_{ik} \partial_{k} \left( w + b \left( |\nabla v|^{2} \right) \right) \right) \\ &= \mathcal{L}w + 2 \sum_{i,k} \partial_{i} \left( (ab') \left( |\nabla v|^{2} \right) E_{ik} \sum_{j} \partial_{j} v \partial_{j} \partial_{k} v \right) \\ &= \mathcal{L}w + 2 \sum_{i,j} \partial_{i} \left( (ab') \left( |\nabla v|^{2} \right) \partial_{j} v \sum_{k} E_{ik} \partial_{k} \partial_{j} v \right) \\ &= \mathcal{L}w + 2 \sum_{i,j} \partial_{i} \left( \left( \frac{b'}{a} \right) \left( |\nabla v|^{2} \right) z_{j} \partial_{j} z_{i} \right) \\ &= \mathcal{L}w + 4 \sum_{i,j} \left( \frac{b'}{a} \right)' \left( |\nabla v|^{2} \right) \sum_{k} \partial_{k} v \partial_{k} \partial_{i} v z_{j} \partial_{j} z_{i} \\ &+ 2 \sum_{i,j} \left( \frac{b'}{a} \right) \left( |\nabla v|^{2} \right) \partial_{i} z_{j} \partial_{j} z_{i} + 2 \sum_{i,j} \left( \frac{b'}{a} \right) \left( |\nabla v|^{2} \right) z_{j} \partial_{j} \partial_{i} z_{i} . \end{aligned}$$

Since  $w = \sum \partial_i z_i$ , the last term of the right-hand side of the above inequality is equal to  $\mathcal{V} \cdot \nabla w$  and

$$(3.8) \qquad \partial_t w = \mathcal{L}w + \mathcal{V} \cdot \nabla w + 4 \left[ a \left( \frac{b'}{a} \right)' \right] \left( |\nabla v|^2 \right) \sum_{i,j,k} \partial_j v \, \partial_k v \, \partial_k \partial_i v \, \partial_j z_i + 2 \left( \frac{b'}{a} \right) \left( |\nabla v|^2 \right) \sum_{i,j} \partial_i z_j \, \partial_j z_i \, .$$

On the one hand, introducing the matrix  $\mathcal{E} := (E_{ij})$  and the Hessian matrix  $D^2 v = (\partial_i \partial_j v)$ of v, we infer from (3.6) that

$$\sum_{i,j} \partial_i z_j \ \partial_j z_i = a^2 \left( |\nabla v|^2 \right) \sum_{i,j,k,l} E_{ik} \ \partial_k \partial_j v \ E_{jl} \ \partial_l \partial_i v$$
$$= a^2 \left( |\nabla v|^2 \right) \sum_{i,j} \left( \mathcal{E} \ D^2 v \right)_{ij} \ \left( \mathcal{E} \ D^2 v \right)_{ji}$$
$$(3.9) \qquad \sum_{i,j} \partial_i z_j \ \partial_j z_i = a^2 \left( |\nabla v|^2 \right) \ tr \left( \left( \mathcal{E} \ D^2 v \right)^2 \right).$$

On the other hand, using once more (3.6), we obtain

$$\sum_{i,j,k} \partial_j v \ \partial_k v \ \partial_k \partial_i v \ \partial_j z_i = a \left( |\nabla v|^2 \right) \sum_{i,j,k,l} \partial_j v \ \partial_k v \ \partial_k \partial_i v \ E_{il} \ \partial_l \partial_j v$$
$$= a \left( |\nabla v|^2 \right) \sum_{i,l} \left( \sum_k \partial_i \partial_k v \ \partial_k v \right) \ E_{il} \left( \sum_j \partial_l \partial_j v \ \partial_j v \right)$$
$$(3.10) \sum_{i,j,k} \partial_j v \ \partial_k v \ \partial_k \partial_i v \ \partial_j z_i = a \left( |\nabla v|^2 \right) \ \left\langle D^2 v \ \nabla v, (\mathcal{E} \ D^2 v) \ \nabla v \right\rangle.$$

Inserting (3.9) and (3.10) in (3.8), we end up with

(3.11) 
$$\partial_t w = \mathcal{L}w + \mathcal{V} \cdot \nabla w + 2 \ (ab') \left( |\nabla v|^2 \right) \ tr \left( \left( \mathcal{E} \ D^2 v \right)^2 \right) \\ + 4 \ (a \ b'' - a' \ b') \left( |\nabla v|^2 \right) \ \left\langle D^2 v \ \nabla v, \left( \mathcal{E} \ D^2 v \right) \ \nabla v \right\rangle .$$

We next observe that

(3.12) 
$$\mathcal{E} \nabla v = \left(1 + 2 |\nabla v|^2 \left(\frac{a'}{a}\right) (|\nabla v|^2)\right) \nabla v$$

and that, for  $X \in \mathbb{R}^N$ ,

$$\langle \mathcal{E} | X, X \rangle = |X|^2 + 2 \left( \frac{a'}{a} \right) \left( |\nabla v|^2 \right) \langle X, \nabla v \rangle^2 \ge |X|^2$$

as a and a' are both positive by (3.1). Consequently,  $\mathcal{E}$  is a positive definite symmetric matrix in  $\mathcal{M}_N(\mathbb{R})$  and there exists a positive definite matrix  $\mathcal{E}_{1/2}$  such that  $\mathcal{E}_{1/2}^2 = \mathcal{E}$ . We then infer from the definition of  $\mathcal{E}_{1/2}$ , (3.12), and Lemma 3.2 (with  $A = \mathcal{E}_{1/2}$ ,  $B = D^2 v$  and

$$\begin{split} X &= \mathcal{E}_{1/2}^{-1} \nabla v) \text{ that} \\ \left\langle D^2 v \ \nabla v, \left(\mathcal{E} \ D^2 v\right) \nabla v \right\rangle &= \left| \left(\mathcal{E}_{1/2} \ D^2 v \right) \nabla v \right|^2 \\ &= \left| \left( \left(\mathcal{E}_{1/2} \ D^2 v \ \mathcal{E}_{1/2} \right) \ \mathcal{E}_{1/2}^{-1} \right) \ \nabla v \right|^2 \\ &\leq tr \left(\mathcal{E}_{1/2} \ D^2 v \ \mathcal{E}_{1/2} \ \mathcal{E}_{1/2} \ D^2 v \ \mathcal{E}_{1/2} \right) \left\langle \mathcal{E}_{1/2}^{-1} \ \nabla v, \mathcal{E}_{1/2}^{-1} \ \nabla v \right\rangle \\ &\leq tr \left( \left(\mathcal{E} \ D^2 v \right)^2 \right) \ \left\langle \nabla v, \mathcal{E}^{-1} \ \nabla v \right\rangle \\ &\leq tr \left( \left(\mathcal{E} \ D^2 v \right)^2 \right) \ |\nabla v|^2 \ \left( 1 + 2 \ |\nabla v|^2 \ \left( \frac{a'}{a} \right) \left( |\nabla v|^2 \right) \right)^{-1} . \end{split}$$

Owing to the non-positivity (3.1) of a b'' - a' b', we deduce from (3.11) and the above inequality that

$$\partial_t w \ge \mathcal{L}w + \mathcal{V} \cdot \nabla w + (a^2 c) (|\nabla v|^2) tr ((\mathcal{E} D^2 v)^2),$$

the function c being defined in (3.2). We finally use the inequality

$$tr(A^2) \ge \frac{1}{N} tr(A)^2, \quad A \in \mathcal{M}_N(\mathbb{R}),$$

the identity

$$w = \sum_{i} \partial_i z_i = a \left( |\nabla v|^2 \right) tr \left( \mathcal{E} D^2 v \right) ,$$

and the non-negativity (3.2) of c to conclude that

$$\partial_t w \geq \mathcal{L}w + \mathcal{V} \cdot \nabla w + \frac{1}{N} (a^2 c) (|\nabla v|^2) tr (\mathcal{E} D^2 v)^2$$
  
$$\geq \mathcal{L}w + \mathcal{V} \cdot \nabla w + \frac{c (|\nabla v|^2)}{N} w^2,$$

and complete the proof.

**Proof of Proposition 1.4.** To be able to use Lemma 3.1, we shall first construct a suitable approximation of (1.1), (1.2). Such a construction has already been performed in [6] for similar purposes and we recall it now. Given  $u_0$  satisfying (1.3), there is a sequence of functions  $(u_{0,k})_{k\geq 1}$  such that, for each integer  $k \geq 1$ ,  $u_{0,k} \in \mathcal{BC}^{\infty}(\mathbb{R}^N)$ ,  $u_0 \leq u_{0,k+1} \leq u_{0,k}$ , and  $(u_{0,k}, \nabla u_{0,k})_k$  converge towards  $(u_0, \nabla u_0)$  uniformly on every compact subset of  $\mathbb{R}^N$  as  $k \to \infty$ . Next, for  $\varepsilon \in (0, 1)$  and  $r \geq 0$ , we set

$$a_{\varepsilon}(r) := (r + \varepsilon^2)^{(p-2)/2}$$
 and  $b_{\varepsilon}(r) := (r + \varepsilon^2)^{q/2} - \varepsilon^q$ .

Then the Cauchy problem

(3.13) 
$$\partial_t u_{k,\varepsilon} = \operatorname{div} \left( a_{\varepsilon} \left( |\nabla u_{k,\varepsilon}|^2 \right) |\nabla u_{k,\varepsilon} \right) + b_{\varepsilon} \left( |\nabla u_{k,\varepsilon}|^2 \right), \quad (t,x) \in Q_{\infty},$$

(3.14) 
$$u_{k,\varepsilon}(0) = u_{0,k} + \varepsilon^{\nu}, \quad x \in \mathbb{R}^N,$$

has a unique classical solution  $u_{k,\varepsilon}$ , the parameter  $\nu > 0$  depending p, q, and N and being appropriately chosen. Furthermore,

(3.15) 
$$\|\nabla u_{k,\varepsilon}(t)\|_{\infty} \leq \|\nabla u_{0,k}\|_{\infty}, \quad t \ge 0,$$

(3.16) 
$$\lim_{k \to \infty} \lim_{\varepsilon \to 0} u_{k,\varepsilon}(t,x) = u(t,x),$$

the latter convergence being uniform on every compact subset of  $[0, \infty) \times \mathbb{R}^N$ , see [6, Section 3] (after performing the same change of unknown function as in the proof of Proposition 2.1).

Introducing

$$c_{\varepsilon}(r) = 2 \left(\frac{b_{\varepsilon}'}{a_{\varepsilon}}\right)(r) + \frac{4r \left(a_{\varepsilon} \ b_{\varepsilon}'' - a_{\varepsilon}' \ b_{\varepsilon}'\right)(r)}{a_{\varepsilon}^{2}(r) + 2r \ a_{\varepsilon}(r) \ a_{\varepsilon}'(r)}, \quad r \ge 0,$$

let us check that  $a_{\varepsilon}$  and  $b_{\varepsilon}$  fulfill the conditions (3.1) and (3.2). Clearly,  $a_{\varepsilon} > 0$  and  $a'_{\varepsilon} > 0$  as p > 2. Next, since  $1 < q \leq p$ ,

$$(a_{\varepsilon}' b_{\varepsilon}' - a_{\varepsilon} b_{\varepsilon}'')(r) = \frac{q (p-q)}{4} (r+\varepsilon^2)^{(p+q-6)/2} \ge 0,$$

$$c_{\varepsilon}(r) = q \frac{r(q-1)+\varepsilon^2}{r(p-1)+\varepsilon^2} (r+\varepsilon^2)^{(q-p)/2} \ge 0.$$

We may then apply Lemma 3.1 to deduce that  $w_{k,\varepsilon} := \operatorname{div} (a_{\varepsilon}(|\nabla u_{k,\varepsilon}|^2) |\nabla u_{k,\varepsilon})$  satisfies

$$\partial_t w_{k,\varepsilon} - \mathcal{L}_{k,\varepsilon} w_{k,\varepsilon} - \mathcal{V}_{k,\varepsilon} \cdot \nabla w_{k,\varepsilon} - \frac{c_{\varepsilon} \left( |\nabla u_{k,\varepsilon}|^2 \right)}{N} \ w_{k,\varepsilon}^2 \ge 0$$

in  $Q_{\infty}$ . Observe next that the condition  $1 < q \leq p$  implies that  $c_{\varepsilon}$  is a non-increasing function. It then follows from (3.15) that  $c_{\varepsilon}(|\nabla u_{k,\varepsilon}|^2) \geq c_{\varepsilon}(||\nabla u_{0,k}||_{\infty}^2)$  and we end up with

(3.17) 
$$\partial_t w_{k,\varepsilon} - \mathcal{L}_{k,\varepsilon} w_{k,\varepsilon} - \mathcal{V}_{k,\varepsilon} \cdot \nabla w_{k,\varepsilon} - \frac{c_\varepsilon \left( \|\nabla u_{0,k}\|_{\infty}^2 \right)}{N} w_{k,\varepsilon}^2 \ge 0$$

in  $Q_{\infty}$ . Clearly,  $t \mapsto -N/(c_{\varepsilon}(\|\nabla u_{0,k}\|_{\infty}^2) t)$  is a subsolution to (3.17) and the comparison principle warrants that

(3.18) 
$$w_{k,\varepsilon}(t,x) \ge -\frac{N}{c_{\varepsilon} \left( \|\nabla u_{0,k}\|_{\infty}^2 \right) t}, \quad (t,x) \in Q_{\infty}$$

Letting  $\varepsilon \to 0$  and  $k \to \infty$  in the previous inequality with the help of (3.16) gives (1.12).

Next, since (1.1) is autonomous, we infer from (2.2) (with s = 0) and (1.12) that

$$\begin{aligned} \Delta_p u(t,x) &\geq -\frac{2N(p-1)}{q(q-1)} \frac{\|\nabla u(t/2)\|_{\infty}^{p-q}}{t} \\ &\geq -\frac{2^{p/q}N(p-1)}{q(q-1)} C_1^{p-q} \|u_0\|_{\infty}^{(p-q)/q} t^{-p/q}, \end{aligned}$$

whence (1.11) for s = 0. To prove the general case  $s \in (0, t)$ , we use again the fact that (1.1) is autonomous.

We have a similar result when  $u_0$  is more regular.

**Corollary 3.3** Assume that p > 2 and  $q \in (1, p]$ . Let u be the solution to (1.1), (1.2) with an initial condition  $u_0$  satisfying  $u_0 \in W^{2,\infty}(\mathbb{R}^N)$  in addition to (1.3). Then

(3.19) 
$$\Delta_p u(t,x) \ge - \left| \inf_{y \in \mathbb{R}^N} \Delta_p u_0(y) \right|$$

in the sense of distributions.

**Proof.** Keeping the notations introduced in the proof of Proposition 1.4, we readily infer from (3.17) and the comparison principle that

(3.20) 
$$w_{k,\varepsilon}(t,x) \ge - \left| \inf_{y \in \mathbb{R}^N} \Delta_p u_{0,k}(y) \right|, \quad (t,x) \in Q_{\infty}$$

Owing to the regularity of  $u_0$ , it is possible to construct the sequence  $(u_{0,k})_k$  such that it satisfies

$$\lim_{k \to \infty} \inf_{y \in \mathbb{R}^N} \Delta_p u_{0,k}(y) = \inf_{y \in \mathbb{R}^N} \Delta_p u_0(y) \,.$$

We may then pass to the limit first as  $\varepsilon \to 0$  and then as  $k \to \infty$  in (3.20) and use (3.16) and the above convergence to complete the proof.

Another useful consequence of the semiconvexity estimates derived in Proposition 1.4 is that the solution u to (1.1), (1.2) is a supersolution to a first-order Hamilton-Jacobi equation.

**Corollary 3.4** Consider an initial condition  $u_0$  satisfying (1.3). Setting  $F(t,\xi_0,\xi) := \xi_0 - |\xi|^q + \kappa_1 ||u_0||_{\infty}^{(p-q)/q} t^{-p/q}$  for  $t \in (0,\infty)$ ,  $\xi_0 \in \mathbb{R}$ , and  $\xi \in \mathbb{R}^N$  (recall that  $\kappa_1$  is defined in (1.11)), the solution u to (1.1), (1.2) is a supersolution to  $F(t,\partial_t w, \nabla w) = 0$  in  $Q_{\infty}$ .

**Proof.** We still use the notations introduced in the proof of Proposition 1.4. As  $w_{k,\varepsilon} = \text{div } (a_{\varepsilon} (|\nabla u_{k,\varepsilon}|^2) |\nabla u_{k,\varepsilon})$ , we infer from (3.13) and (3.18) that

$$\partial_t u_{k,\varepsilon} - b_{\varepsilon} \left( |\nabla u_{k,\varepsilon}|^2 \right) \ge -\frac{N}{c_{\varepsilon} \left( ||\nabla u_{0,k}||_{\infty}^2 \right) t}$$

in  $Q_{\infty}$ . We then use (3.16) and the stability of viscosity solutions [1, 2, 8] to pass to the limit as  $\varepsilon \to 0$  and  $k \to \infty$  in the previous inequality and conclude that u is a supersolution to

$$\partial_t w - |\nabla w|^q + \frac{N(p-1)}{q(q-1)} \frac{\|\nabla u_0\|_{\infty}^{p-q}}{t} = 0 \text{ in } Q_{\infty}.$$

Now, fix  $T \ge 0$ . As (1.1) is an autonomous equation, the function  $(t, x) \mapsto u(t + T, x)$  is the solution to (1.1) with initial condition u(T) and the above analysis allows us to conclude that u is a supersolution to

$$\partial_t w - |\nabla w|^q + \frac{N(p-1)}{q(q-1)} \frac{\|\nabla u(T)\|_{\infty}^{p-q}}{t-T} = 0 \quad \text{in} \quad (T,\infty) \times \mathbb{R}^N.$$

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We then use (2.2) (with T = t/2) to complete the proof.

#### 4 Convergence to self-similarity

We change the variables and the unknown function so that the convergence (1.5) is transformed to the convergence towards a steady state. More precisely, we introduce the selfsimilar (or scaling) variables

$$\tau = \frac{1}{q} \log (1+t), \qquad y = \frac{x}{(1+t)^{1/q}}$$

and the new unknown function v defined by

(4.1) 
$$u(t,x) = v\left(\frac{\log(1+t)}{q}, \frac{x}{(1+t)^{1/q}}\right), \quad (t,x) \in [0,\infty) \times \mathbb{R}^N.$$

Equivalently,  $v(\tau, y) = u(e^{q\tau} - 1, ye^{\tau})$  for  $(\tau, y) \in [0, \infty) \times \mathbb{R}^N$  and it follows from (1.1), (1.2) that v solves

(4.2) 
$$\partial_{\tau} v = y \cdot \nabla v + q |\nabla v|^{q} + q e^{-(p-q)\tau} \Delta_{p} v, \quad (\tau, y) \in Q_{\infty},$$

(4.3) 
$$v(0) = u_0, \quad y \in \mathbb{R}^N$$

We also infer from (2.1) and (2.2) that there is a positive constant  $C_5(u_0)$  depending only on N, p, q, and  $u_0$  such that

(4.4) 
$$||v(\tau)||_{\infty} + ||\nabla v(\tau)||_{\infty} \le C_5(u_0), \quad \tau \ge 0,$$

while (1.4) reads

(4.5) 
$$\lim_{\tau \to \infty} \|v(\tau)\|_{\infty} = M_{\infty} > 0.$$

Formally, since p > q, the diffusion term vanishes in the large time limit and we expect the large time behaviour of the solution v to (4.2), (4.3) to look like that of the solutions to the first-order Hamilton-Jacobi equation

(4.6) 
$$\partial_{\tau} w - y \cdot \nabla w - q \ |\nabla w|^q = 0 \text{ in } Q_{\infty}.$$

Now, to investigate the large time behaviour of first-order Hamilton-Jacobi equations, an efficient approach has been developed in [22, 23] which relies on the relaxed half-limits method introduced in [3]. More precisely, for  $(\tau, y) \in Q_{\infty}$ , we define the relaxed half-limits  $v_*$  and  $v^*$  by

(4.7) 
$$v_*(y) := \liminf_{(\sigma, z, \lambda) \to (\tau, y, \infty)} v(\sigma + \lambda, z) \text{ and } v^*(y) := \limsup_{(\sigma, z, \lambda) \to (\tau, y, \infty)} v(\sigma + \lambda, z).$$

These relaxed half-limits are well-defined thanks to (4.4) and we first note that the right-hand sides of the above definitions indeed do not depend on  $\tau > 0$ . In addition,

(4.8) 
$$0 \le v_*(x) \le v^*(x) \le M_{\infty} \quad \text{for} \quad y \in \mathbb{R}^N$$

by (4.5), while (4.4) and the Rademacher theorem ensure that  $v_*$  and  $v^*$  both belong to  $W^{1,\infty}(\mathbb{R}^N)$ . Finally, by [2, Théorème 4.1] applied to equation (4.2),  $v^*$  and  $v_*$  are viscosity subsolution and supersolution, respectively, to the Hamilton-Jacobi equation

(4.9) 
$$\mathcal{H}(y, \nabla w) := -y \cdot \nabla w - q \ |\nabla w|^q = 0 \quad \text{in} \quad \mathbb{R}^N$$

We now aim at showing that  $v^*$  and  $v_*$  coincide. However, the equation (4.9) has infinitely many solutions as  $y \mapsto (c - \gamma_q |y|^{q/(q-1)})_+$  solves (4.9) for any c > 0. The information obtained so far on  $v_*$  and  $v^*$  are thus not sufficient and are supplemented by the next two results.

**Lemma 4.1** Given  $\varepsilon \in (0,1)$ , there is  $R_{\varepsilon} > 1/\varepsilon$  such that

(4.10) 
$$v(\tau, y) \le \varepsilon \quad for \quad \tau \ge 0 \quad and \quad y \in \mathbb{R}^N \setminus B(0, R_{\varepsilon}),$$

and  $0 \le v_*(y) \le v^*(y) \le \varepsilon$  for  $y \in \mathbb{R}^N \setminus B(0, R_{\varepsilon})$ .

In other words,  $v(\tau)$  belongs to  $\mathcal{C}_0(\mathbb{R}^N)$  for each  $\tau \ge 0$  in a way which is uniform with respect to  $\tau \ge 0$ .

**Proof.** We first construct a supersolution to (4.2) in  $(0, \infty) \times (\mathbb{R}^N \setminus B(0, R))$  for R large enough. To this end, consider  $R \geq R_c := 1 + (q \ (2 \|u_0\|_{\infty})^{q-1} + 3pq \ (2 \|u_0\|_{\infty})^{p-2})^{1/q}$  and put  $\Sigma_R(y) = \|u_0\|_{\infty} R^2 \|y\|^{-2}$  for  $y \in \mathbb{R}^N \setminus B(0, R)$ . Let  $\mathcal{L}$  be the parabolic operator defined by

$$\mathcal{L}w(\tau, y) := \partial_{\tau}w(\tau, y) - y \cdot \nabla w(\tau, y) - q \ |\nabla w(\tau, y)|^{q} - q \ e^{-(p-q)\tau} \ \Delta_{p}w(\tau, y)$$

for  $(\tau, y) \in Q_{\infty}$  (so that  $\mathcal{L}v = 0$  by (4.2)). Then, if  $y \in \mathbb{R}^N \setminus B(0, R)$ , we have

$$\mathcal{L}\Sigma_{R}(y) = 2 \Sigma_{R}(y) - q \frac{2^{q}}{|y|^{q}} \Sigma_{R}(y)^{q} + q 2^{p-1} \frac{N+2-3p}{|y|^{p}} \Sigma_{R}(y)^{p-1} e^{-(p-q)\tau} \geq 2 \Sigma_{R}(y) \left\{ 1 - q (2 ||u_{0}||_{\infty})^{q-1} \frac{R^{2(q-1)}}{|y|^{3q-2}} - 3pq e^{-(p-q)\tau} (2 ||u_{0}||_{\infty})^{p-2} \frac{R^{2(p-2)}}{|y|^{3p-4}} \right\} \geq 2 \Sigma_{R}(y) \left\{ 1 - q (2 ||u_{0}||_{\infty})^{q-1} R^{-q} - 3pq e^{-(p-q)\tau} (2 ||u_{0}||_{\infty})^{p-2} R^{-p} \right\} \geq 0$$

by the choice of R. Consequently,  $\Sigma_R$  is a supersolution to (4.2) in  $(0, \infty) \times (\mathbb{R}^N \setminus B(0, R))$  for  $R \geq R_c$ .

Now, fix  $\varepsilon \in (0, 1)$ . Since  $u_0 \in \mathcal{C}_0(\mathbb{R}^N)$ , there is  $\rho_{\varepsilon} \geq \max \{R_c, \varepsilon^{-1}\}$  such that  $u_0(y) \leq \varepsilon/2$  if  $|y| \geq \rho_{\varepsilon}$ . We then infer from the monotonicity of  $\Sigma_R$  and (2.1) that

$$u_0(y) - \frac{\varepsilon}{2} - \Sigma_{\rho_{\varepsilon}}(z) \le -\Sigma_{\rho_{\varepsilon}}(z) \le 0$$

if  $|y| \ge \rho_{\varepsilon}$  and  $|z| \ge \rho_{\varepsilon}$ ,

$$v(\tau, y) - \frac{\varepsilon}{2} - \Sigma_{\rho_{\varepsilon}}(z) \le ||u_0||_{\infty} - \Sigma_{\rho_{\varepsilon}}(z) = \Sigma_{\rho_{\varepsilon}}(y) - \Sigma_{\rho_{\varepsilon}}(z) \le \frac{2 ||u_0||_{\infty}}{\rho_{\varepsilon}} |y - z|$$

if  $|y| = \rho_{\varepsilon}$ ,  $|z| \ge \rho_{\varepsilon}$ , and  $\tau \ge 0$ , and

$$v(\tau, y) - \frac{\varepsilon}{2} - \sum_{\rho_{\varepsilon}} (z) \le ||u_0||_{\infty} - ||u_0||_{\infty} \le 0$$

if  $|y| \ge \rho_{\varepsilon}$ ,  $|z| = \rho_{\varepsilon}$ , and  $\tau \ge 0$ . As  $v - \varepsilon/2$  and  $\Sigma_{\rho_{\varepsilon}}$  are subsolution and supersolution, respectively, to (4.2), the comparison principle [11, Theorem 4.1] warrants that  $v(\tau, y) - \varepsilon/2 \le \Sigma_{\rho_{\varepsilon}}(y)$  for  $\tau \ge 0$  and  $|y| \ge \rho_{\varepsilon}$ . It remains to choose  $R_{\varepsilon} \ge \rho_{\varepsilon}$  such that  $\Sigma_{\rho_{\varepsilon}}(y) \le \varepsilon/2$ for  $|y| \ge R_{\varepsilon}$  to complete the proof of (4.10). The last assertion of Lemma 4.1 is then a straightforward consequence of the definition (4.7) and (4.10).

We next use the semiconvexity estimate (1.11) (and more precisely its consequence stated in Corollary 3.4) to show that  $v_*$  lies above the profile  $H_{\infty}$  defined in (1.6).

**Lemma 4.2** For  $y \in \mathbb{R}^N$ , we have

(4.11)  $H_{\infty}(y) \le v_*(y) \le v^*(y)$ .

**Proof.** For  $\tau \geq 0$ ,  $y \in \mathbb{R}^N$ ,  $\xi_0 \in \mathbb{R}$  and  $\xi \in \mathbb{R}^N$ , we set  $\mathcal{F}(\tau, y, \xi_0, \xi) := \xi_0 - y \cdot \xi - q |\xi|^q + \kappa_2 e^{-(p-q)\tau}$  with  $\kappa_2 := q \kappa_1 ||u_0||_{\infty}^{(p-q)/q} (e^q/(e^q-1))^{p/q}$ , the constant  $\kappa_1$  being defined in (1.11). It then readily follows from Corollary 3.4 that

(4.12) v is a supersolution to  $\mathcal{F}(\tau, y, \partial_{\tau} w, \nabla w) = 0$  in  $(1, \infty) \times \mathbb{R}^N$ .

We next fix  $\tau_0 > 1$  and denote by V the (viscosity) solution to

$$\partial_{\tau} V - y \cdot \nabla V - q \ |\nabla V|^{q} = 0, \qquad (\tau, y) \in (\tau_{0}, \infty) \times \mathbb{R}^{N}, V(\tau_{0}) = v(\tau_{0}), \qquad y \in \mathbb{R}^{N}.$$

On the one hand, a straightforward computation shows that the function  $\tilde{V}$  defined by

$$\tilde{V}(\tau, y) := V(\tau, y) - \kappa_2 \int_{\tau_0}^{\tau} e^{-(p-q)s} ds, \qquad (\tau, y) \in (\tau_0, \infty) \times \mathbb{R}^N,$$

is the (viscosity) solution to  $\mathcal{F}(\tau, y, \partial_{\tau} \tilde{V}, \nabla \tilde{V}) = 0$  in  $(\tau_0, \infty) \times \mathbb{R}^N$  with initial condition  $\tilde{V}(\tau_0) = v(\tau_0)$ . Recalling (4.12), we infer from the comparison principle that

(4.13) 
$$\tilde{V}(\tau, y) \le v(\tau, y) \text{ for } (\tau, y) \in (\tau_0, \infty) \times \mathbb{R}^N.$$

On the other hand, it follows from Proposition A.1 that

$$\lim_{\tau \to \infty} \sup_{y \in \mathbb{R}^N} \left| V(\tau, y) - \left( \| v(\tau_0) \|_{\infty} - \gamma_q \| y \|^{q/(q-1)} \right)_+ \right| = 0.$$

We may then pass to the limit as  $\tau \to \infty$  in (4.13) and use the definition (4.7) to conclude that

$$\left(\|v(\tau_0)\|_{\infty} - \gamma_q \, |y|^{q/(q-1)}\right)_+ - \kappa_2 \, \int_{\tau_0}^{\infty} e^{-(p-q)s} \, ds \le v_*(y) \le v^*(y)$$

for  $y \in \mathbb{R}^N$ . Letting  $\tau_0 \to \infty$  in the above inequality with the help of (4.5) completes the proof of the lemma.

We are now in a position to complete the proof of Theorem 1.1. To this end, fix  $\varepsilon \in (0, 1)$ . Lemma 4.1 ensures that  $v^*(y) \leq \varepsilon$  for  $|y| \geq R_{\varepsilon} \geq 1/\varepsilon$  while the continuity of  $H_{\infty}$  implies that there is  $r_{\varepsilon} \in (0, \varepsilon)$  such that  $H_{\infty}(y) \geq M_{\infty} - \varepsilon$  for  $|y| \leq r_{\varepsilon}$ . Recalling (4.8), we realize that

(4.14) 
$$\begin{cases} v^*(y) - \varepsilon \le 0 \le H_{\infty}(y) & \text{if } |y| = R_{\varepsilon} \\ v^*(y) - \varepsilon \le M_{\infty} - \varepsilon \le H_{\infty}(y) & \text{if } |y| = r_{\varepsilon}. \end{cases}$$

Moreover, introducing  $\psi(y) = -\gamma_q |y|^{q/(q-1)}/2$ , we have

(4.15) 
$$\mathcal{H}(y,\nabla\psi(y)) = \frac{q\gamma_q}{2(q-1)} |y|^{q/(q-1)} \left(1 - \frac{1}{2^{q-1}}\right) > 0 \quad \text{if} \quad r_{\varepsilon} < |y| < R_{\varepsilon},$$

the Hamiltonian  $\mathcal{H}$  being defined in (4.9). Summarizing, we have shown that  $H_{\infty}$  and  $v^* - \varepsilon$ are supersolution and subsolution, respectively, to (4.9) in  $\Omega_{\varepsilon} := \{y \in \mathbb{R}^N : r_{\varepsilon} < |y| < R_{\varepsilon}\}$ with  $v^* - \varepsilon \leq H_{\infty}$  on  $\partial \Omega_{\varepsilon}$  by (4.14). Owing to (4.15) and the concavity of  $\mathcal{H}$  with respect to its second variable, we may apply [15, Theorem 1] to conclude that  $v^* - \varepsilon \leq H_{\infty}$  in  $\Omega_{\varepsilon}$ . This property being valid for each  $\varepsilon \in (0, 1)$ , we actually have  $v^* \leq H_{\infty}$  in  $\mathbb{R}^N$  by passing to the limit as  $\varepsilon \to 0$  thanks to the properties of  $r_{\varepsilon}$  and  $R_{\varepsilon}$ . Recalling (4.11), we have thus established that  $v^* = v_* = H_{\infty}$  in  $\mathbb{R}^N$ . In particular, the property  $v^* = v_*$  and the definition (4.7) provide the uniform convergence of  $\{v(\tau)\}_{\tau\geq 0}$  towards  $v^* = H_{\infty}$  on every compact subset of  $\mathbb{R}^N$  as  $\tau \to \infty$ , see [2, Lemme 4.1] or [1, Lemma V.1.9]. Combining this local convergence with Lemma 4.1 actually gives

(4.16) 
$$\lim_{\tau \to \infty} \|v(\tau) - H_{\infty}\|_{\infty} = 0$$

Theorem 1.1 then readily follows after writing the convergence (4.16) in the original variables (t, x) for the function u and noticing that  $||h_{\infty}(1+t) - h_{\infty}(t)||_{\infty} \longrightarrow 0$  as  $t \to \infty$ .

### 5 Limit value of $||u(t)||_{\infty}$

This section is devoted to the proof of Proposition 1.4, for which three cases are to be distinguished and handled differently:  $q \in (1, p - 1], q \in (p - 1, q_*]$ , and  $q \in (q_*, p)$ .

**Proof of Proposition 1.4:**  $q \in (1, p - 1]$ . We proceed as in [19, Proposition 1] (where a similar result is proved for p = 2 and q = 1). For  $\alpha > N/2$ ,  $\delta > 0$ , and  $x \in \mathbb{R}^N$ , we set  $\varrho_{\delta}(x) := (1 + \delta |x|^2)^{-\alpha}$ . Clearly,  $\varrho_{\delta} \in L^1(\mathbb{R}^N)$  and it follows from (2.3) that

$$\frac{d}{dt} \int_{\mathbb{R}^N} \varrho_{\delta}(x) \ u(t,x) \ dx = \int_{\mathbb{R}^N} \left\{ \varrho_{\delta}(x) \ |\nabla u(t,x)|^q - |\nabla u(t,x)|^{p-2} \ \nabla u(t,x) \cdot \nabla \varrho_{\delta}(x) \right\} \ dx$$
$$\geq \int_{\mathbb{R}^N} \varrho_{\delta}(x) \ |\nabla u(t,x)|^q \ \left( 1 - |\nabla u(t,x)|^{p-1-q} \ \frac{|\nabla \varrho_{\delta}(x)|}{\varrho_{\delta}(x)} \right) \ dx \,.$$

Recalling that  $\|\nabla u(t)\|_{\infty} \leq \|\nabla u_0\|_{\infty}$  by (1.3) and (2.2) and noticing that  $|\nabla \varrho_{\delta}| \leq \alpha \ \delta^{1/2} \ \varrho_{\delta}$ , we further obtain

$$\frac{d}{dt} \int_{\mathbb{R}^N} \varrho_{\delta}(x) \ u(t,x) \ dx \ge \int_{\mathbb{R}^N} \varrho_{\delta}(x) \ |\nabla u(t,x)|^q \ \left(1 - \alpha \ \delta^{1/2} \ \|\nabla u_0\|_{\infty}^{p-1-q}\right) \ dx$$

Choosing  $\delta = \|\nabla u_0\|_{\infty}^{2(q+1-p)}/\alpha^2$  and integrating with respect to time give

$$\|u(t)\|_{\infty} \|\varrho_{\delta}\|_{1} \ge \int_{\mathbb{R}^{N}} \varrho_{\delta}(x) \ u(t,x) \ dx \ge \int_{\mathbb{R}^{N}} \varrho_{\delta}(x) \ u_{0}(x) \ dx > 0.$$

We then pass to the limit as  $t \to \infty$  to conclude that  $M_{\infty} > 0$ .

We next turn to the case  $q \in (p-1, q_*]$  which turns out to be more complicated and requires two preparatory results.

**Lemma 5.1** Assume that  $q \in (1, q_*]$  and let u be a non-negative solution to (1.1), (1.2) with a compactly supported initial condition  $u_0$  satisfying (1.3). Then  $u(t) \in L^1(\mathbb{R}^N)$  for each  $t \ge 0$ , the function  $t \mapsto ||u(t)||_1$  is non-decreasing and

(5.1) 
$$\lim_{t \to \infty} \|u(t)\|_1 = \infty.$$

**Proof.** For every  $t \ge 0$ , u(t) is bounded and compactly supported by (2.1) and Proposition 2.2, and is thus in  $L^1(\mathbb{R}^N)$ . The time monotonicity of the  $L^1$ -norm of u then readily follows from (2.3) with  $\vartheta = 1$ , a valid choice in this particular case as u(t) is compactly supported. It further follows from (2.3) with  $\vartheta = 1$  that

(5.2) 
$$||u(t)||_1 \ge ||u(T)||_1 + \int_T^t ||\nabla u(s)||_q^q \, ds \quad \text{for } t > T \ge 0.$$

Consider next T > 0 and t > T. Recalling the Gagliardo-Nirenberg inequality

(5.3) 
$$\|w\|_q \le C_6 \|\nabla w\|_q^{N(q-1)/(N(q-1)+q)} \|w\|_1^{q/(N(q-1)+q)}, \quad w \in W^{1,q}(\mathbb{R}^N) \cap L^1(\mathbb{R}^N),$$

we infer from (5.2), (5.3), and the time monotonicity of the  $L^1$ -norm of u that

$$\begin{aligned} \|u(t)\|_{1}^{1+(q^{2}/N(q-1))} &\geq \|u(t)\|_{1}^{q^{2}/N(q-1)} \left( \|u(T)\|_{1} + \int_{T}^{t} \|\nabla u(s)\|_{q}^{q} ds \right) \\ &\geq \int_{T}^{t} \|u(s)\|_{1}^{q^{2}/N(q-1)} \|\nabla u(s)\|_{q}^{q} ds \\ &\geq C_{7} \int_{T}^{t} \left( \|u(s)\|_{q}^{q} \right)^{(N(q-1)+q)/N(q-1)} ds \,. \end{aligned}$$

If  $\varphi$  denotes the solution to the *p*-Laplacian equation  $\partial_t \varphi - \Delta_p \varphi = 0$  in  $Q_{\infty}$  with initial condition  $\varphi(0) = u_0$ , the comparison principle readily implies that

(5.4) 
$$\varphi(t,x) \le u(t,x), \quad (t,x) \in Q_{\infty}.$$

Inserting this estimate in the previous lower bound for  $||u(t)||_1$ , we end up with

(5.5) 
$$\|u(t)\|_{1}^{1+(q^{2}/N(q-1))} \geq C_{7} \int_{T}^{t} \left(\|\varphi(s)\|_{q}^{q}\right)^{(N(q-1)+q)/N(q-1)} ds .$$

Now, by Proposition 2.3 we have

$$\lim_{s \to \infty} s^{N(q-1)/(N(p-2)+p)} \|\varphi(s) - \mathcal{B}_{\|u_0\|_1}(s)\|_q^q = 0$$

and

$$\|\mathcal{B}_{\|u_0\|_1}(s)\|_q^q = C_8 s^{-N(q-1)/(N(p-2)+p)},$$

so that

$$\begin{aligned} \|\varphi(s)\|_{q}^{q} &\geq \left(\left\|\mathcal{B}_{\|u_{0}\|_{1}}(s)\right\|_{q} - \left\|\varphi(s) - \mathcal{B}_{\|u_{0}\|_{1}}(s)\right\|_{q}\right)^{q} \\ &\geq s^{-N(q-1)/(N(p-2)+p)} \left(C_{8} - \left\|\varphi(s) - \mathcal{B}_{\|u_{0}\|_{1}}(s)\right\|_{q}\right)^{q} \\ &\geq \left(\frac{C_{8}}{2}\right)^{q} s^{-N(q-1)/(N(p-2)+p)} \end{aligned}$$

for  $s \geq T$ , provided T is chosen sufficiently large. Inserting this estimate in (5.5) gives

$$\begin{aligned} \|u(t)\|_{1}^{1+(q^{2}/N(q-1))} &\geq C_{9} \int_{T}^{t} s^{-(N(q-1)+q)/(N(p-2)+p)} ds \\ &\geq C_{10} \begin{cases} t^{(N+1)(q_{\star}-q)/(N(p-2)+p)} - T^{(N+1)(q_{\star}-q)/(N(p-2)+p)} & \text{if } q \in (1,q_{\star}), \\ &\log(t/T) & \text{if } q = q_{\star}. \end{cases} \end{aligned}$$

We then let  $t \to \infty$  to obtain the claimed result.

We next argue as in [12, Lemma 14] (for p = 2) to show that, if  $q \in (p-1, p)$  and  $M_{\infty} = 0$ , then the  $L^{\infty}$ -norm of u(t) decays faster than an explicit rate.

**Lemma 5.2** Assume that  $q \in (p-1, p)$  and let u be a non-negative solution to (1.1), (1.2) with an initial condition  $u_0$  satisfying (1.3). If  $M_{\infty} = 0$  in (1.4), then

(5.6) 
$$||u(t)||_{\infty} \le C_{11} t^{-(p-q)/(2q-p)} \text{ for } t > 0.$$

Observe that the assumptions p > 2 and  $q \in (p-1, p)$  imply that 2q > p and (p-q)/(2q-p) > 0.

**Proof.** Consider a non-negative function  $\eta \in \mathcal{C}^{\infty}(\mathbb{R}^N)$  with compact support in B(0,1) and  $\|\eta\|_1 = 1$ . We then define a sequence of mollifiers  $(\eta_{\delta})_{\delta}$  by  $\eta_{\delta}(x) := \eta(x/\delta)/\delta^N$  for  $x \in \mathbb{R}^N$ 

and  $\delta \in (0,1)$ . For  $(t, x_0) \in Q_{\infty}$  and T > t, we take  $\vartheta(x) = \eta_{\delta}(x - x_0)$  in (2.3) and infer from (1.11) (with s = t/2) that

$$\begin{aligned} \|u(T)\|_{\infty} &\geq \int_{\mathbb{R}^{N}} u(T,x) \ \eta_{\delta}(x-x_{0}) \ dx \\ &\geq \int_{\mathbb{R}^{N}} u(t,x) \ \eta_{\delta}(x-x_{0}) \ dx - \int_{t}^{T} \int_{\mathbb{R}^{N}} |\nabla u(s,x)|^{p-2} \ \nabla u(s,x) \cdot \nabla \eta_{\delta}(t,x-x_{0}) \ dx ds \\ &\geq \int_{\mathbb{R}^{N}} u(t,x) \ \eta_{\delta}(x-x_{0}) \ dx - 2^{p/q} \ \kappa_{1} \ \left\| u\left(\frac{t}{2}\right) \right\|_{\infty}^{(p-q)/q} \ \int_{t}^{T} (2s-t)^{-p/q} \ ds \\ &\geq \int_{\mathbb{R}^{N}} u(t,x) \ \eta_{\delta}(x-x_{0}) \ dx - C_{12} \ \left\| u\left(\frac{t}{2}\right) \right\|_{\infty}^{(p-q)/q} \ \left(t^{(q-p)/q} - T^{(q-p)/q}\right). \end{aligned}$$

Owing to the continuity of u, we may pass to the limit as  $\delta \to 0$  in the above inequality and deduce that

$$||u(T)||_{\infty} \ge u(t, x_0) - C_{12} ||u\left(\frac{t}{2}\right)||_{\infty}^{(p-q)/q} (t^{(q-p)/q} - T^{(q-p)/q}).$$

But the above inequality is valid for all  $x_0 \in \mathbb{R}^N$  and we thus end up with

$$||u(T)||_{\infty} \ge ||u(t)||_{\infty} - C_{12} \left\| u\left(\frac{t}{2}\right) \right\|_{\infty}^{(p-q)/q} \left(t^{(q-p)/q} - T^{(q-p)/q}\right).$$

Finally, as q < p, we may let  $T \to \infty$  in the previous inequality and use the assumption  $M_{\infty} = 0$  to conclude that

$$||u(t)||_{\infty} \le C_{12} ||u\left(\frac{t}{2}\right)||_{\infty}^{(p-q)/q} t^{(q-p)/q},$$

or, equivalently, as 2q > p,

$$t^{(p-q)/(2q-p)} \|u(t)\|_{\infty} \le C_{13} \left\{ \left(\frac{t}{2}\right)^{(p-q)/(2q-p)} \|u\left(\frac{t}{2}\right)\|_{\infty} \right\}^{(p-q)/q}$$

for  $t \geq 0$ . Introducing

$$A(t) := \sup_{s \in (0,t)} \left\{ s^{(p-q)/(2q-p)} \| u(s) \|_{\infty} \right\} \in [0,\infty), \quad t \ge 0,$$

we deduce from the previous inequality that  $A(t) \leq C_{13} A(t)^{(p-q)/q}$ , whence  $A(t) \leq C_{13}^{q/(2q-p)}$  for  $t \geq 0$ . This bound being valid for each t > 0, the proof of (5.6) is complete.

**Proof of Proposition 1.4:**  $q \in (p-1, q_*]$ .

Step 1: We first consider a compactly supported initial condition  $u_0$  satisfying (1.3) and assume for contradiction that  $M_{\infty} = 0$ . On the one hand, according to Lemma 5.2 and the assumption  $q \leq q_{\star}$ , there holds

(5.7) 
$$\limsup_{t \to \infty} t^{N/(N(p-2)+p)} \|u(t)\|_{\infty} \le C_{11} t^{(p(N+1)(q-q_*))/((2q-p)(N(p-2)+p))} \le C_{11}.$$

On the other hand, fix  $t_0 > 0$  and let  $\varphi$  be the solution to the *p*-Laplacian equation  $\partial_t \varphi - \Delta_p \varphi = 0$  in  $Q_\infty$  with initial condition  $\varphi(0) = u(t_0)$ . As  $u_0$  is compactly supported, so is  $u(t_0)$  by Proposition 2.2 and  $u(t_0)$  thus belongs to  $L^1(\mathbb{R}^N)$ . Moreover, the comparison principle warrants that  $u(t,x) \ge \varphi(t-t_0,x)$  for  $(t,x) \in [t_0,\infty) \times \mathbb{R}^N$ . We then infer from the above properties and Proposition 2.3 that, for  $t > t_0$ ,

$$t^{N/(N(p-2)+p)} \| u(t) \|_{\infty} \geq (t-t_0)^{N/(N(p-2)+p)} \| u(t) \|_{\infty} \\ \geq (t-t_0)^{N/(N(p-2)+p)} \| \varphi(t-t_0) \|_{\infty} \\ \geq (t-t_0)^{N/(N(p-2)+p)} \| \mathcal{B}_{\|u(t_0)\|_1}(t-t_0) \|_{\infty} \\ - (t-t_0)^{N/(N(p-2)+p)} \| \mathcal{B}_{\|u(t_0)\|_1}(t-t_0) - \varphi(t-t_0) \|_{\infty} \\ \geq C_{14} \| u(t_0) \|_1^{p/(N(p-2)+p)} \\ - (t-t_0)^{N/(N(p-2)+p)} \| \mathcal{B}_{\|u(t_0)\|_1}(t-t_0) - \varphi(t-t_0) \|_{\infty} .$$

Using once more Proposition 2.3, we may pass to the limit as  $t \to \infty$  in the previous inequality to obtain

(5.8) 
$$\liminf_{t \to \infty} t^{N/(N(p-2)+p)} \|u(t)\|_{\infty} \ge C_{14} \|u(t_0)\|_1^{p/(N(p-2)+p)}$$

Combining (5.7) and (5.8) yields  $||u(t_0)||_1 \leq C_{15}$  for all  $t_0 > 0$  which contradicts Lemma 5.1. Therefore,  $M_{\infty} > 0$ .

Step 2: Now, if  $u_0$  is an arbitrary initial condition satisfying (1.3), there clearly exists a compactly supported initial condition  $\tilde{u}_0$  satisfying (1.3) and such that  $u_0 \geq \tilde{u}_0$  in  $\mathbb{R}^N$ . Introducing the solution  $\tilde{u}$  to (1.1) with initial condition  $\tilde{u}_0$ , the comparison principle entails that  $u \geq \tilde{u}$  in  $Q_{\infty}$ , hence

$$M_{\infty} \ge \lim_{t \to \infty} \|\tilde{u}(t)\|_{\infty}$$

The first step of the proof ensures that the right-hand side of the above inequality is positive which completes the proof.  $\hfill \Box$ 

It remains to investigate the case  $q \in (q_*, p)$ , for which we adapt the proof of [7, Theorem 2.4(b)].

**Proof of Proposition 1.4:**  $q \in (q_{\star}, p)$ . We put

$$m_0 := \left| \inf_{y \in \mathbb{R}^N} \Delta_p u_0(y) \right|$$

As in the proof of Lemma 5.1, let  $\eta \in \mathcal{C}^{\infty}(\mathbb{R}^N)$  be a non-negative function with compact support in B(0,1) and  $\|\eta\|_1 = 1$ , and define a sequence of mollifiers  $(\eta_{\delta})_{\delta}$  by  $\eta_{\delta}(x) := \eta(x/\delta)/\delta^N$  for  $x \in \mathbb{R}^N$  and  $\delta \in (0,1)$ . For  $(t, x_0) \in Q_{\infty}$  and  $T \in (0, t)$ , we take  $\vartheta(x) = \eta_{\delta}(x - x_0)$  in (2.3) and infer from (1.11) (with s = 0) and Corollary 3.3 that

$$\begin{aligned} \|u(t)\|_{\infty} &\geq \int_{\mathbb{R}^{N}} u(t,x) \ \eta_{\delta}(x-x_{0}) \ dx \\ &\geq \int_{\mathbb{R}^{N}} u_{0}(x) \ \eta_{\delta}(x-x_{0}) \ dx - \int_{0}^{t} \int_{\mathbb{R}^{N}} |\nabla u(s,x)|^{p-2} \ \nabla u(s,x) \cdot \nabla \eta_{\delta}(t,x-x_{0}) \ dx ds \\ &\geq \int_{\mathbb{R}^{N}} u_{0}(x) \ \eta_{\delta}(x-x_{0}) \ dx - \int_{0}^{T} m_{0} \ ds - \kappa_{1} \ \|u_{0}\|_{\infty}^{(p-q)/q} \ \int_{T}^{t} s^{-p/q} \ ds \\ &\geq \int_{\mathbb{R}^{N}} u_{0}(x) \ \eta_{\delta}(x-x_{0}) \ dx - T \ m_{0} - C_{16} \ \|u_{0}\|_{\infty}^{(p-q)/q} \ \left(T^{(q-p)/q} - t^{(q-p)/q}\right) . \end{aligned}$$

Owing to the continuity of  $u_0$ , we may pass to the limit as  $\delta \to 0$  in the above inequality and deduce that

$$||u(t)||_{\infty} \ge u_0(x_0) - T m_0 - C_{16} ||u_0||_{\infty}^{(p-q)/q} (T^{(q-p)/q} - t^{(q-p)/q}).$$

Since q < p, we may let  $t \to \infty$  in the above inequality and take the supremum with respect to  $x_0$  to conclude that

$$M_{\infty} \ge \|u_0\|_{\infty} - T \ m_0 - C_{16} \ \|u_0\|_{\infty}^{(p-q)/q} \ T^{(q-p)/q}$$

Next, for  $\beta \in (0, 1)$ , the choice  $T = \|u_0\|_{\infty}^{(p-q)/p} (\beta + m_0)^{-q/p}$  in the previous inequality yields

$$M_{\infty} \ge \|u_0\|_{\infty}^{(p-q)/p} \left(\|u_0\|_{\infty}^{q/p} - (1+C_{16}) (\beta+m_0)^{(p-q)/p}\right).$$

This inequality being valid for every  $\beta \in (0, 1)$ , we conclude that

$$M_{\infty} \ge \|u_0\|_{\infty}^{(p-q)/p} \left(\|u_0\|_{\infty}^{q/p} - (1+C_{16}) \ m_0^{(p-q)/p}\right) > 0$$

as soon as (1.10) is fulfilled with  $\kappa_0 = (1 + C_{16})^{p/q}$ .

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### A Convergence for the Hamilton-Jacobi equation (4.6)

In this section, we study the large behaviour of non-negative solutions to the Hamilton-Jacobi equation (4.6) with initial data in  $\mathcal{C}_0(\mathbb{R}^N)$  and show their convergence to a steady state uniquely determined by the  $L^{\infty}$ -norm of the initial data. Though the large time behaviour of solutions to first-order Hamilton-Jacobi equations has received considerable attention in recent years (see [4, 5, 16, 22, 23] and the references therein), the particular case of (4.6) does not seem to have been investigated in the literature. We thus provide a simple proof relying on the Hopf-Lax-Oleinik formula.

**Proposition A.1** Let q > 1 and consider a non-negative function  $h_0 \in C_0(\mathbb{R}^N)$ . Let h be the unique viscosity solution to the Cauchy problem

(A.1) 
$$\partial_{\tau}h - y \cdot \nabla h - q |\nabla h|^q = 0, \quad (\tau, y) \in (0, \infty) \times \mathbb{R}^N,$$

(A.2) 
$$h(0) = h_0, \qquad y \in \mathbb{R}^N.$$

Then

(A.3) 
$$\lim_{\tau \to \infty} \|h(\tau) - h_s\|_{\infty} = 0$$

with

$$h_s(y) := \left( \|h_0\|_{\infty} - \gamma_q \ |y|^{q/(q-1)} \right)_+, \qquad y \in \mathbb{R}^N,$$

the constant  $\gamma_q = (q-1) q^{-q/(q-1)}$  being defined in Theorem 1.1.

Thanks to the concavity of the Hamiltonian  $\mathcal{H}(y,\xi) = -y \cdot \xi - q |\xi|^q$ ,  $(y,\xi) \in \mathbb{R}^N \times \mathbb{R}^N$ , with respect to its second variable, the Hopf-Lax-Oleinik formula provides a representation formula for the solution h to (A.1), (A.2) which can be used to prove (A.3).

**Proof.** We first recall that h is given by the Hopf-Lax-Oleinik formula

$$h(\tau, y) = \sup_{z \in \mathbb{R}^N} \left\{ h_0(z) - \gamma_q \left| y - z \, e^{-\tau} \right|^{q/(q-1)} \left( 1 - e^{-q\tau} \right)^{-1/(q-1)} \right\}$$

for  $(\tau, y) \in [0, \infty) \times \mathbb{R}^N$ , see, e.g., [10, Chapter 3]. Since  $h(\tau, y) \ge h_0(ye^{\tau}) \ge 0$ , we have in fact

$$h(\tau, y) = \sup_{z \in \mathbb{R}^N} \left\{ \left( h_0(z) - \gamma_q | y - z e^{-\tau} |^{q/(q-1)} (1 - e^{-q\tau})^{-1/(q-1)} \right)_+ \right\}$$

for  $(\tau, y) \in [0, \infty) \times \mathbb{R}^N$ .

Consider now  $\beta \in (0,1)$ . As  $h_0 \in \mathcal{C}_0(\mathbb{R}^N)$ , there is  $R_\beta > (\|h_0\|_{\infty}/\gamma_q)^{(q-1)/q}$  such that

(A.4) 
$$h_0(z) \le \beta \quad \text{for} \quad |z| \ge R_\beta$$

On the one hand, if  $(\tau, y) \in [\log R_{\beta}, \infty) \times \mathbb{R}^N$  and  $z \in \mathbb{R}^N$ , we have either  $|z| \geq R_{\beta}$  and

$$\left| \begin{pmatrix} h_0(z) - \gamma_q \ |y - z \ e^{-\tau} |^{q/(q-1)} \ (1 - e^{-q\tau})^{-1/(q-1)} \end{pmatrix}_+ - \begin{pmatrix} h_0(z) - \gamma_q \ |y|^{q/(q-1)} \end{pmatrix}_+ \right|$$

$$\leq \left( h_0(z) - \gamma_q \ |y - z \ e^{-\tau} |^{q/(q-1)} \ (1 - e^{-q\tau})^{-1/(q-1)} \right)_+ + \left( h_0(z) - \gamma_q \ |y|^{q/(q-1)} \right)_+$$

$$\leq 2 \beta$$

by (A.4) or  $z \in B(0, R_{\beta})$  and

$$\left| \begin{pmatrix} h_0(z) - \gamma_q \ |y - z \ e^{-\tau}|^{q/(q-1)} \ (1 - e^{-q\tau})^{-1/(q-1)} \end{pmatrix}_+ - \begin{pmatrix} h_0(z) - \gamma_q \ |y|^{q/(q-1)} \end{pmatrix}_+ \right|$$

$$\leq \gamma_q \ |y - z \ e^{-\tau}|^{q/(q-1)} \ \left\{ (1 - e^{-q\tau})^{-1/(q-1)} - 1 \right\} + \gamma_q \ \left| |y - z \ e^{-\tau}|^{q/(q-1)} - |y|^{q/(q-1)} \right|$$

$$\leq \gamma_q \ (|y| + R_\beta \ e^{-\tau})^{q/(q-1)} \ \left\{ (1 - e^{-q\tau})^{-1/(q-1)} - 1 \right\} + \frac{q \ \gamma_q}{q-1} \ (|y| + |z| \ e^{-\tau})^{1/(q-1)} \ |z| \ e^{-\tau}$$

$$\leq \gamma_q \ (|y| + 1)^{1/(q-1)} \ \left\{ \frac{q}{q-1} + |y| + 1 \right\} \ \left\{ (1 - e^{-q\tau})^{-1/(q-1)} - 1 + R_\beta \ e^{-\tau} \right\}$$

as  $\tau \geq \log R_{\beta}$ . Combining the above two estimates gives

$$\left| h(\tau, y) - \sup_{z \in \mathbb{R}^N} \left\{ \left( h_0(z) - \gamma_q |y|^{q/(q-1)} \right)_+ \right\} \right|$$
  
  $\leq C(q) (|y|+1)^{q/(q-1)} \left\{ \left( 1 - e^{-q\tau} \right)^{-1/(q-1)} - 1 + R_\beta e^{-\tau} \right\} + 2\beta$ 

whence

(A.5) 
$$|h(\tau, y) - h_s(y)| \le C(q) (|y|+1)^{q/(q-1)} \left\{ \left(1 - e^{-q\tau}\right)^{-1/(q-1)} - 1 + R_\beta e^{-\tau} \right\} + 2\beta$$

for  $(\tau, y) \in [\log R_{\beta}, \infty) \times \mathbb{R}^N$ . On the other hand, if  $\tau \geq \log(R_{\beta})$ ,  $|y| \geq Y := 1 + (\|h_0\|_{\infty}/\gamma_q)^{(q-1)/q}$  and  $z \in \mathbb{R}^N$ , we have either  $|y - z e^{-\tau}| \geq Y - 1$  and

$$\begin{aligned} h_0(z) &- \gamma_q |y - z e^{-\tau}|^{q/(q-1)} (1 - e^{-q\tau})^{-1/(q-1)} \\ &\leq (1 - e^{-q\tau})^{-1/(q-1)} \left\{ \|h_0\|_{\infty} (1 - e^{-q\tau})^{1/(q-1)} - \gamma_q |y - z e^{-\tau}|^{q/(q-1)} \right\} \\ &\leq (1 - e^{-q\tau})^{-1/(q-1)} \left\{ \|h_0\|_{\infty} - \gamma_q (Y - 1)^{q/(q-1)} \right\} \\ &\leq 0, \end{aligned}$$

or  $|y - z \ e^{-\tau}| < Y - 1$  and

$$|z| \ge |y \ e^{\tau}| - |z - y \ e^{\tau}| \ge Y \ e^{\tau} - (Y - 1) \ e^{\tau} = e^{\tau} \ge R_{\beta},$$

so that

$$h_0(z) - \gamma_q |y - z e^{-\tau}|^{q/(q-1)} (1 - e^{-q\tau})^{-1/(q-1)} \le \beta$$

by (A.4). Therefore,

(A.6) 
$$h(\tau, y) \leq \beta \quad \text{for} \quad (\tau, y) \in [\log R_{\beta}, \infty) \times (\mathbb{R}^N \setminus B(0, Y)).$$

The claim (A.3) then easily follows from (A.5) and (A.6).

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