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FAST MODEL OF SPACE-VARIANT BLURRING AND ITS APPLICATION TO DECONVOLUTION IN ASTRONOMY

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ABSTRACT

Image deblurring is essential to high resolution imaging and is therefore widely used in astronomy, microscopy or computational photography. While shift-invariant blur is modeled by convolution and leads to fast FFT-based algorithms, shift-variant blurring requires models both accurate and fast. When the point spread function (PSF) varies smoothly across the field, these two opposite objectives can be reached by interpolating from a grid of PSF samples.

Several models for smoothly varying PSF co-exist in the literature. We advocate that one of them is both physically-grounded and fast. Moreover, we show that the approximation can be largely improved by tuning the PSF samples and interpolation weights with respect to a given continuous model. This improvement comes without increasing the computational cost of the blurring operator.

We illustrate the developed blurring model on a deconvolution application in astronomy. Regularized reconstruction with our model leads to large improvements over existing results.

Index Terms— deconvolution, shift-variant PSF

1. INTRODUCTION

Image deconvolution is widely used to enhance the resolution, signal-to-noise ratio and contrast of blurred images. In many cases, blur is space-variant and thus can no longer be modeled by a convolution. Accurate modeling of the point spread function (PSF) is essential for restoration. Depending on whether PSF variations across the field are smooth or discontinuous, PSF are either interpolated, or the field is segmented into regions inside which PSF are invariant. In the former case, effort is put in finding a good tradeoff between approximation quality and speed. The most challenging aspect of the latter case is the segmentation step.

In astronomy, blurring due to atmospheric turbulence varies in the field of view. Even after adaptive optics correction, the PSF is shift-variant since the correction quality decreases away from the guiding star[1]. Optical aberrations and obstructions also lead to space-variant degradations[2]. In 3D microscopy, the 3-D PSF varies with the depth from the coverslip[3]. In all these applications, PSF vary smoothly.

Except for very compact PSF, the storage and application of a different PSF for each pixel[4] is not computationally tractable for iterative deblurring methods. It has then been proposed[5] to decompose the image into patches where blur can be considered approximatively invariant (isoplanatic regions). After deconvolution of each patch, the reconstructed regions must be tied together which leads to border artifacts[6, 3, 7]. It is therefore preferable to model smooth blur variations and then, to deblur the whole image.

Smooth PSF variations can be decomposed on a subspace of PSF[8, 9]. The cost of this modeling increases linearly with the number of basis PSF used. It has been noticed independently by several authors that PSF variations could be modeled by interpolation and yet lead to a fast blurring model[10, 11, 12]. Two formulations have been proposed, equivalent in terms of computational complexity but leading to different PSF models. We discuss in section 2.1 these two approximations and their implications in terms of PSF modeling. We show that one of them is superior both in approximation quality and modeling properties. In section 2.2, we further improve PSF approximation error by tuning the PSF samples and interpolation weights for a target PSF model. The space-variant blurring model is then applied to deconvolution of Hubble Space Telescope (HST) simulated data. We show improved reconstructions compared to the results reported on the same dataset in [10].

2. APPROXIMATION OF SMOOTHLY VARYING PSF

During the formation of an image g, the original (crisp) distribution f undergoes distortions due to atmosphere turbulence, object/camera relative motion, the instrument (limited aperture, optical aberrations). These degradations are typically modeled by a linear transform:

$$g(\mathbf{r}) = \int h(\mathbf{r}, \mathbf{s}) f(\mathbf{s}) d\mathbf{s}, \tag{1}$$

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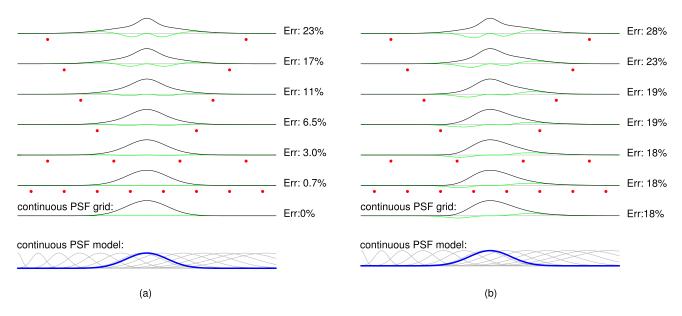


Fig. 1. Comparison of (a) \mathcal{H} and (b) \mathcal{H}^{\dagger} for PSF approximation (blue: target PSF, red: PSF sample locations, green: approximation error).

where h denotes the point spread function (PSF). When the PSF does only depend on the difference r - s, equation (1) is a convolution and the system is said isoplanatic.

2.1. PSF interpolation

Smoothly varying PSF can be approximated by interpolating PSF at locations s_i with interpolation kernel φ_{int} :

$$h(\mathbf{r}, \mathbf{s}) \approx \sum_{i} h_i(\mathbf{r} - \mathbf{s}) \varphi_{\text{int}}(\mathbf{s} - \mathbf{s}_i),$$
 (2)

with $h_i(u) \equiv h(s_i + u, s_i)$ the PSF for a source located at s_i . The image formation model is then approximated by:

$$g(\mathbf{r}) \approx \sum_{i} \int h_{i}(\mathbf{r} - \mathbf{s}) \, \varphi_{i}(\mathbf{s}) \, f(\mathbf{s}) \, d\mathbf{s} \equiv [\mathcal{H} \circ f](\mathbf{r}), \quad (3)$$

where $\varphi_i(s) = \varphi_{\text{int}}(s - s_i)$ and \mathcal{H} denotes the linear operator of our approximate model. This operator and its adjoint (which is needed for image reconstruction) expand as:

$$\mathcal{H} = \sum_{i} \mathcal{H}_{i} \circ \mathcal{W}_{i} \quad \text{and} \quad \mathcal{H}^{\sharp} = \sum_{i} \mathcal{W}_{i} \circ \mathcal{H}_{i}^{\sharp} , \qquad (4)$$

with W_i a scaling operator corresponding to the pointwise multiplication by φ_i , \mathcal{H}_i and \mathcal{H}_i^{\sharp} a convolution and a correlation by the *i*th sampled PSF. In words, Eq. (3) expresses the degraded image as the sum of convolutions between PSF samples and weighted versions of the crisp image.

An alternative model has been proposed by [10] who interpolate the result of convolving the original image by the PSF samples. The corresponding operator writes:

$$\mathcal{H}^{\dagger} = \sum_{i} \mathcal{W}_{i} \circ \mathcal{H}_{i} \,. \tag{5}$$

[12] uses formulation \mathcal{H} and [11] give both. Note that \mathcal{H}^{\dagger} is similar to the adjoint \mathcal{H}^{\sharp} of our operator (using convolutions instead of correlations); hence the computational burden are the same for both. Though closely related, the operators \mathcal{H} and \mathcal{H}^{\dagger} are however different in their ability to correctly approximate the shift-variant PSF. Figure 1 gives a monodimensional illustration of the difference between the two. We consider a system with Gaussian PSF whose standard deviation increases linearly with the position from left to right. Several of these Gaussian PSF are drawn in light gray on last row of the figure. The target PSF at the center of the field is drawn in thick blue stroke. This PSF is approximated using PSF samples at locations depicted by red dots: coarse sampling at the top row, and refined sampling at the next rows until continuous sampling for the last but one. A linear interpolation kernel is used. The formulation of \mathcal{H} is used for sub-figure 1(a), while \mathcal{H}^{\dagger} is used for sub-figure 1(b). Approximation errors are drawn in green, and the relative norm of the error is given at the right of each curve. Several properties of the approximations are noticeable: (i) PSF interpolation (\mathcal{H}) preserves PSF symmetry, while interpolation of convolution results (\mathcal{H}^{\dagger}) does not; (ii) when the PSF supports are small compared to the distance between two PSF sample locations, the two approximations are comparable; (iii) \mathcal{H} achieves exact interpolation (i.e., PSF approximation error tends to zero when the PSF grid is ever more refined) while \mathcal{H}^{\dagger} does not; (iv) PSF positivity and normalization are preserved by linear interpolation (which is not the case with \mathcal{H}^{\dagger} , even when a linear interpolation kernel is used). Based on the difference between the two models, we recommend the use of formulation \mathcal{H} rather than \mathcal{H}^{\dagger} for smoothly varying blurs.

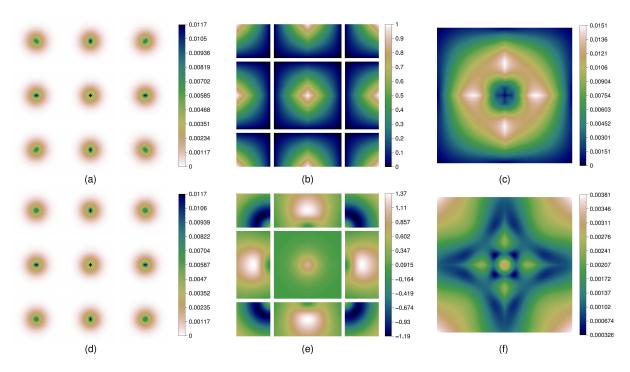


Fig. 2. PSF approximation: (a)-(c) using bilinear interpolation; (d)-(f) generalized interpolation with optimal PSF and weights; (a), (d): PSF grid; (b), (e): interpolation weights on each of the 9 patches; (c), (f): rms error with respect to target PSF.

As noted in [10, 11, 12], operators \mathcal{H} and \mathcal{H}^{\dagger} can be efficiently computed using fast Fourier transforms (FFT). Each term of the sum in Eq. (4) is a convolution of a PSF (size H^2) with a patch whose size B^2 corresponds to the support of the interpolation kernel (for a grid of K^2 PSF samples, an N^2 image, and linear interpolation, B = 2N/K). Each discrete convolution requires $\mathcal{O}[(B+H)^2\log(B+H)]$ operations, thus the whole operator requires $\mathcal{O}[K^2(B+H)^2\log(B+H)]$ operations. If the PSF support H is much smaller than the support of the interpolation kernel, the computational cost is $\mathcal{O}[4N^2 \log N]$ operations, i.e., 4 times the cost of a single convolution of the whole image. Table 1 reports the computation time of our parallel implementation of operator \mathcal{H} (based on FFTW and Linux pthread libraries, running on an Intel Xeon 3.3GHz with 6 cores). When the PSF support is large compared to the support of the interpolation kernel, most of the computational effort is spent computing the boundaries of each patch and the complexity raises as illustrated by the last rows of table 1.

2.2. Improvement of the approximation

There are several options to improve the approximation error: (i) refine the interpolation grid (costly when the PSF support is large); (ii) increase the interpolation order (huge cost since the support of the interpolation kernel is proportional to the interpolation order); (iii) use generalized interpolation, i.e., refine interpolation weights φ_i (with constant interpolation support) and PSF samples h_i to fit a target PSF.

nb of threads		1	2	4	6
PSF grid time in ms (relative to a convolution)					olution)
	5×5	71 (3.7)	37 (1.9)	19(1)	14 (0.7)
(a)	10×10	46 (2.4)	22 (1.2)	11 (0.6)	8 (0.4)
	20×20	60 (3.2)	30 (1.6)	15 (0.8)	11 (0.6)
	5×5	182 (4.2)	97 (2.3)	57 (1.3)	47 (1.1)
(b)	10×10	636 (15)	325 (7.6)	170 (4)	117 (2.7)
	20×20	1680 (39)	850 (20)	435 (10)	293 (6.8)

Table 1. Average time to compute operator \mathcal{H} : (a) 512^2 pixels image, with 31×31 pixels PSF; (b) 1000^2 pixels image, with 101×101 pixels PSF (Intel Xeon Processor with 6 cores).

Option (iii) is the most interesting since it improves the approximation without increasing the computational cost of \mathcal{H} . We achieve this by minimizing the quadratic difference:

$$\epsilon^2 = \frac{1}{N^2} \iint \left[h(\boldsymbol{r}, \boldsymbol{s}) - \sum_i h_i(\boldsymbol{r} - \boldsymbol{s}) \, \varphi_i(\boldsymbol{s}) \right]^2 \, \mathrm{d}\boldsymbol{r} \, \mathrm{d}\boldsymbol{s} \ \, (6)$$

with respect to the PSF $\{h_i\}_{i=1}^{K^2}$ and the weights $\{\varphi_i\}_{i=1}^{K^2}$. Starting with the PSF samples, we alternately solve for the best weights, and then for the best PSF samples. Using this algorithm, the approximation can be largely improved (the rms error ϵ is divided by 4, which is better than when using a 6×6 grid with linear interpolation), as illustrated by figure 2.

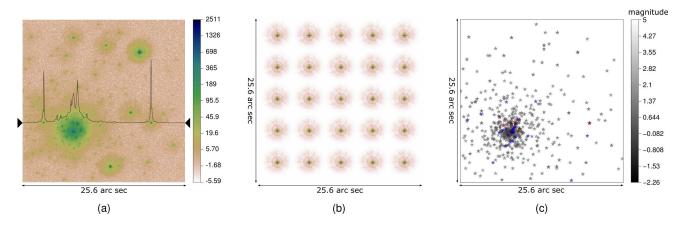


Fig. 3. Deconvolution with shift-variant PSF: (a) simulation of an observation of a star field with Hubble's Wide Field Camera before corrective optics; (b) provided 5×5 grid of PSF; (c) recovered stars after deconvolution with sparsity prior, gray: good detection, blue: non-detected star (4.9%), red: false-detection (7.2%)

2.3. Application to deconvolution in astronomy

We used the same dataset as in [10] to illustrate the application of our model on a realistic astronomical case. Contrary to [10] who used truncated preconditioned conjugate gradients, we performed regularized inversion with an ℓ_1 sparsity-inducing prior and a positivity constraint (minimization done with FISTA algorithm [13]). Figure 3 presents the obtained results: rms error is 3.6% when using 5×5 PSF grid, 6.2% when using the mean PSF, to be compared with the 16% error reported in [10]. The use of a regularization largely improves the reconstruction. The spatially-variant PSF model further improves the results.

3. CONCLUSION

We have shown that PSF interpolation is physically more justified and leads to a better approximation of space-variant PSF. The computational cost remains modest $(4 \times \text{that of a single convolution})$. We proposed to further improve the model accuracy while keeping the same computational cost. We applied our model to a realistic astronomical case. Our next step will be to extend our approach to blind or myopic deconvolution, in wide-field adaptive optics and microscopic imaging.

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