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CONTINUUM PERCOLATION IN HIGH DIMENSIONS

JEAN-BAPTISTE GOUÉRE AND RÉGINE MARCHAND

ABSTRACT. Consider a Boolean model Σ in \mathbb{R}^d , where the centers are given by a homogeneous Poisson point process with intensity λ and the radii of distinct balls are i.i.d. with common distribution ν . Some numerical simulations and some heuristic arguments suggest that the critical covered volume $c_d^c(\nu)$, which is the proportion of space covered by Σ at critical intensity, may be minimal when ν is a Dirac measure.

1. INTRODUCTION AND STATEMENT OF THE MAIN RESULTS

Introduction. Consider a homogeneous Poisson point process on \mathbb{R}^d . At each point of this process, we center a ball with random radius, the radii of distinct balls being i.i.d. and independent of the point process. The union Σ of these random balls is called a Boolean model. This Boolean model only depends on two parameters : the intensity λ of the point process of centers, and the common distribution ν of the radii of the balls.

Denote by $\lambda_d^c(\nu)$ the critical intensity for percolation in Σ . We then consider $c_d^c(\nu)$, the proportion of space which is covered by Σ when $\lambda = \lambda_d^c(\nu)$. This quantity is called the critical covered volume. Some numerical simulations in low dimension and some heuristic arguments in any dimension suggest that the critical covering fraction may be minimal when ν is a Dirac measure, that is when all the balls have the same radius.

We show that this is not true in high dimensions. We prove this by studying the critical covered volume $c_d^c(r^{-d}\mu(dr))$ of a given measure μ as d tends to infinity. See the definition (8) below for some remarks about this normalization. When μ is a Dirac measure, this has been studied by Penrose in [14]. Our results yield that whenever μ is not a Dirac measure, for all d large enough, $c_d^c(r^{-d}\mu(dr))$ is strictly smaller than $c_d^c(r^{-d}\delta_1(dr)) = c_d^c(\delta_1)$.

When μ is a Dirac measure, the asymptotic behavior of $c_d^c(r^{-d}\delta_1(dr)) = c_d^c(\delta_1)$ is given by the associated Galton-Watson process. This is due to the fact that geometrical dependencies vanish in high dimensions. This is not true anymore in the general case and geometry still plays a role.

The Boolean model. Let us give here a different – but equivalent – construction of the Boolean model. Let ν be a finite¹ measure on $(0, +\infty)$. We assume that the mass of ν is positive. Let $d \geq 2$ be an integer, $\lambda > 0$ be a real number and ξ be a Poisson point process on $\mathbb{R}^d \times (0, +\infty)$ whose intensity measure is the Lebesgue measure on \mathbb{R}^d times $\lambda\nu$. We define a random subset of \mathbb{R}^d as follows:

$$\Sigma(\lambda\nu) = \bigcup_{(c,r) \in \xi} B(c, r),$$

where $B(c, r)$ is the open Euclidean ball centered at $c \in \mathbb{R}^d$ and with radius $r \in (0, +\infty)$. The random subset $\Sigma(\lambda\nu)$ is a Boolean model driven by $\lambda\nu$.

1. There is no greater generality in considering finite measures instead of probability measures ; this is simply more convenient.

Critical intensity. We say that $\Sigma(\lambda\nu)$ percolates if the connected component of $\Sigma(\lambda\nu)$ that contains the origin is unbounded with positive probability. This is equivalent to the existence of an unbounded connected component of $\Sigma(\lambda\nu)$ with probability one. We refer to the book by Meester and Roy [11] for background on continuum percolation. Set:

$$\lambda_d^c(\nu) = \inf\{\lambda > 0 : \Sigma(\lambda\nu) \text{ percolates}\}.$$

One easily checks that $\lambda_d^c(\nu)$ is finite. In [6] it is proven that $\lambda_d^c(\nu)$ is positive if and only if

$$(1) \quad \int r^d \nu(dr) < +\infty.$$

We assume that this assumption is fulfilled.

Critical intensity as a function of ν , critical covered volume. By ergodicity, the Boolean model $\Sigma(\lambda\nu)$ has a deterministic natural density. This is also the probability that a given point belongs to the Boolean model and this is given by :

$$P(0 \in \Sigma(\lambda\nu)) = 1 - \exp\left(-\lambda \int v_d r^d \nu(dr)\right),$$

where v_d denotes the volume of the unit ball in \mathbb{R}^d . The critical covered volume $c_d^c(\nu)$ is the density of the Boolean model when the intensity is critical :

$$c_d^c(\nu) = 1 - \exp\left(-\lambda_d^c(\nu) \int v_d r^d \nu(dr)\right).$$

It is more convenient to study the critical covered volume through the normalised threshold defined by :

$$\tilde{\lambda}_d^c(\nu) = \lambda_d^c(\nu) \int v_d (2r)^d \nu(dr).$$

Note :

$$c_d^c(\nu) = 1 - \exp\left(-\frac{\tilde{\lambda}_d^c(\nu)}{2^d}\right).$$

The factor 2^d may seem arbitrary here. Its interest will appear in the statement of the next theorems.

We will now give two scaling relations. For all $a > 0$, define $H^a\nu$ as the image of ν under the map defined by $x \mapsto ax$. By scaling, we get:

$$(2) \quad \tilde{\lambda}_d^c(H^a\nu) = \tilde{\lambda}_d^c(\nu).$$

This is a consequence of Proposition 2.11 in [11], and it may become more obvious when considering the two following facts : a critical Boolean model remains critical when rescaling and the density is invariant by rescaling ; therefore the critical covered volume and then the normalised threshold are invariant. One also easily checks the following invariance:

$$(3) \quad \tilde{\lambda}_d^c(a\nu) = \tilde{\lambda}_d^c(\nu).$$

That is because of these two invariance relations that we prefer to study the normalized critical intensity $\tilde{\lambda}_d^c(\nu)$ instead of the critical intensity $\lambda_d^c(\nu)$

It has been conjectured by Kertész and Vicsek [9] that the critical covered volume – or equivalently the normalized critical intensity – should be independent of ν , as soon as the support of ν is bounded. Phani and Dhar [4] gave a heuristic argument suggesting that

the conjecture were false. A rigorous proof was then given by Meester, Roy and Sarkar in [12]. More precisely, they gave examples of measures ν with two atoms such that:

$$(4) \quad \tilde{\lambda}_d^c(\nu) > \tilde{\lambda}_d^c(\delta_1).$$

As a consequence of Theorem 1.1 in the paper by Menshikov, Popov and Vachkovskaia [13], we even get that $\tilde{\lambda}_d^c(\nu)$ can be arbitrarily large. More precisely, if

$$(5) \quad \nu(n, a) = \sum_{k=0}^{n-1} a^{dk} \delta_{a^{-k}}.$$

then

$$(6) \quad \tilde{\lambda}_d^c(\nu(n, a)) \rightarrow n\tilde{\lambda}_d^c(\delta_1) \text{ as } a \rightarrow \infty.$$

Actually the result of [13] is the following much stronger statement: $\lambda_d^c(\nu(+\infty, a)) \rightarrow \lambda_c(\delta_1)$ when $a \rightarrow \infty$. The convergence (6) is implicit in the work of Meester, Roy and Sarkar in [12], at least when $n = 2$. There were also heuristics for such a result in [4].

By Theorem 2.1 in [6], we get the existence of a positive constant C_d , that depends only on the dimension d , such that:

$$\tilde{\lambda}_d^c(\nu) \geq C_d.$$

To sum up, $\tilde{\lambda}_c^d(\cdot)$ is not bounded from above but is bounded from below by a positive constant. In other words, the critical covering volume $c_d^c(\cdot) \in (0, 1)$ can be arbitrarily close to 1 but is bounded from below by a positive constant. Our investigations were motivated by the following question:

$$(7) \quad \text{Is the critical covered volume } c_d^c(\nu) \text{ minimized when } \nu = \delta_1 \text{ ?}$$

In the physical literature, it is strongly believed that the answer is positive at least when $d = 2$ and $d = 3$. This conjecture is supported by numerical evidence (to the best of our knowledge, the most accurate estimations are given in a paper by Quintanilla and Ziff [15] when $d = 2$ and in a paper by Consiglio, Baker, Paul and Stanley [2] when $d = 3$). On Figure 1, we plot the critical covered volume in dimension 2 as a function of α and for different values of ρ when $\nu = (1 - \alpha)\delta_1 + \alpha\rho^{-2}\delta_\rho$. The data for finite values of ρ come from numerical estimations in [15], while the data for infinite ρ come from the study on the multi-scale Boolean model. See section 1.4 in [8] for further references. It is also supported by some heuristic arguments in any dimension (see for example Dhar [3]). See also [1]. In the above cited paper [12], it is noted that the rigorous proof of (4) suggests that the answer might be positive for any $d \geq 2$.

In this paper, we show that the answer is actually negative for all d large enough.

Critical intensity in high dimensions : the case of a deterministic radius.

Assume here that the measure ν is a Dirac mass at 1, that is that the radii of the balls are all equal to 1. Penrose proved the following result in [14] :

Theorem 1.1 (Penrose). $\lim_{d \rightarrow \infty} \tilde{\lambda}_d^c(\delta_1) = 1$.

With the invariance properties of $\tilde{\lambda}_d^c$, this limit can readily be generalized to any constant radius : for any $\alpha, a > 0$,

$$\lim_{d \rightarrow \infty} \tilde{\lambda}_d^c(\alpha\delta_a) = \lim_{d \rightarrow \infty} \tilde{\lambda}_d^c(\delta_1) = 1.$$

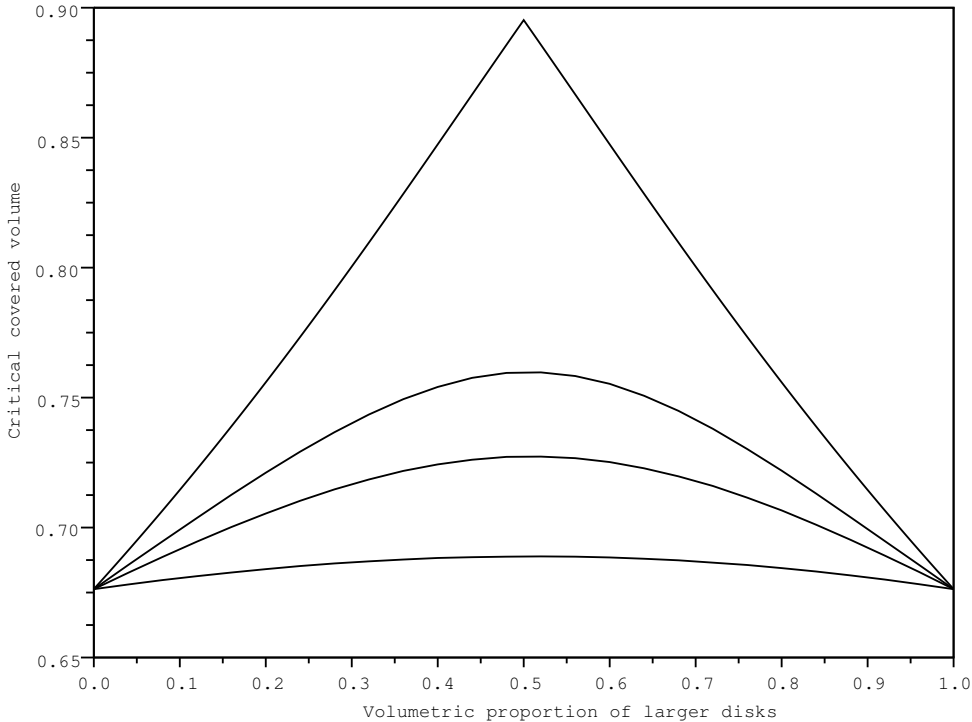


FIGURE 1. Critical covered volume as a function of α for different values of ρ . From bottom to top: $\rho = 2, \rho = 5, \rho = 10$ and the limit as $\rho \rightarrow \infty$.

Theorem 1.1 is the continuum analogue of a result of Kesten [10] for Bernoulli bond percolation on the nearest-neighbor integer lattice \mathbb{Z}^d , which says that the critical percolation parameter is asymptotically equivalent to $1/(2d)$.

Let us say a word about the ideas of the proof of Theorem 1.1.

The inequality $\tilde{\lambda}_d^c(\delta_1) > 1$ holds for any $d \geq 2$. The proof is simple, and here is the idea. We consider the following natural genealogy. The deterministic ball $B(0, 1)$ is said to be the ball of generation 0. The random balls of $\Sigma(\lambda\delta_1)$ that touches $B(0, 1)$ are then the balls of generation 1. The random balls that touches one ball of generation 1 without being one of them are then the balls of generation 2 and so on. Let us denote by N_d the number of all balls that are descendants of $B(0, 1)$. There is no percolation if and only if N_d is almost surely finite.

Now denote by m the Poisson distribution with mean $\lambda v_d 2^d$: this is the law of the number of balls of $\Sigma(\lambda\delta_1)$ that touch a given ball of radius 1. Therefore, if there were no interference between children of different balls, N_d would be equal to Z , the total population in a Galton-Watson process with offspring distribution m . Because of the interferences due to the fact that the Boolean model lives in \mathbb{R}^d , this is not true: in fact, N_d is only stochastically dominated by Z . Therefore, if $\lambda v_d 2^d < 1$, then Z is finite almost surely, then N_d is finite almost surely and therefore there is no percolation. This implies

$$\tilde{\lambda}_d^c(\delta_1) = v_d 2^d \lambda_d^c(\delta_1) > 1.$$

The difficult part of Theorem 1.1 is to prove that if d is large, then the interferences are small, then N_d is close to Z and therefore there is percolation as soon as $\tilde{\lambda}_d^c(\delta_1) = v_d 2^d \lambda_d^c(\delta) > 1$.

To sum up, at first order, the asymptotic behavior of the critical intensity of the Boolean model with constant radius is given by the threshold of the associated Galton-Watson process, as in the case of Bernoulli percolation on \mathbb{Z}^d : roughly speaking, as the dimension increases, the geometrical constraints of the finite dimension space decrease and at the limit, we recover the non-geometrical case of the corresponding Galton-Watson process.

Critical intensity in high dimensions : the case of true random radii. If μ is a finite measure on $(0, +\infty)$ and if $d \geq 2$ is an integer, we define a measure μ_d on $(0, +\infty)$ by setting :

$$(8) \quad \mu_d(dr) = r^{-d} \mu(dr).$$

Note that, for any d , the assumption (1) is fulfilled by μ_d , and that $(\delta_1)_d = \delta_1$. Note also that μ_d is not necessarily a finite measure. However the definitions made above still make sense in this case and we still have $\lambda_c^d(\mu_d) \in (0, +\infty)$ thanks to Theorem 1.1 in [7].

We will study the behavior of $\tilde{\lambda}_d^c(\mu_d)$ as d tends to infinity. Let us motivate the definition of μ_d with the following two related properties :

- (1) Consider the Boolean model $\Sigma(\lambda\mu_d)$ on \mathbb{R}^d driven by $\lambda\mu_d$ where $\lambda > 0$. For any $0 < s < t < \infty$, the number of balls of $\Sigma(\lambda\mu_d)$ with radius in $[s, t]$ that contains a given point is a Poisson random variable with intensity:

$$\int_{[s,t]} v_d r^d \lambda \mu_d(dr) = v_d \lambda \mu([s, t]).$$

Loosely speaking, this means that contrary to what happens in the Boolean model driven by $\lambda\mu$, the relative importance of radii of different sizes does not depend on the dimension d in the Boolean model driven by $\lambda\mu_d$.

- (2) A closely related property is the following one. Consider for example the case $\mu = \alpha\delta_a + \beta\delta_b$. Then, a way to build $\Sigma(\lambda\mu_d)$ is to proceed as follows. Consider two independent Boolean model: Σ^A , driven by $\lambda\alpha\delta_1$, and Σ^B , driven by $\lambda\beta\delta_1$. Then set $\Sigma(\lambda\mu_d) = a\Sigma^A + b\Sigma^B$.

We prove the following result :

Theorem 1.2. *Let μ be a finite measure on $(0, +\infty)$. We assume that the mass of μ is positive and that μ is not concentrated on a singleton. Then :*

$$\limsup_{d \rightarrow +\infty} \frac{1}{d} \ln \left(\tilde{\lambda}_d^c(\mu_d) \right) < 0.$$

Note that $(\delta_1)_d = \delta_1$. A straightforward consequence of Theorem 1.2 and Theorem 1.1 – or, actually, of the much weaker and easier convergence of $\ln(\tilde{\lambda}_d^c(\delta_1))$ to 0 – is the following result:

Corollary 1.3. *Let μ be a finite measure on $(0, +\infty)$. We assume that the mass of μ is positive and that μ is not concentrated on a singleton. Then, for any d large enough, we have:*

$$\tilde{\lambda}_d^c(\mu_d) < \tilde{\lambda}_d^c(\delta_1), \text{ or equivalently } c_d^c(\mu_d) < c_d^c(\delta_1).$$

In fact, Theorem 1.2 follows from the particular case of radii taking only two different values, for which we have more precise results.

Critical intensity in high dimensions : the case when the radii can take two values. To state the result, we need some further notations. Fix $\rho > 1$. Fix $k \geq 1$. Set $r_1 = r_{k+1} = 1 + \rho$, and for $i \in \{2, \dots, k\}$, $r_i = 2$. For $(a_i)_{2 \leq i \leq k+1} \in [0, 1]^k$, we build an increasing sequence of distances $(d_i)_{1 \leq i \leq k+1}$ by setting $d_1 = 1 + \rho$ and, for every $i \in \{2, \dots, k+1\}$:

$$d_i^2 = d_{i-1}^2 + 2r_i a_i d_{i-1} + r_i^2.$$

Note that the sequence $(d_i)_{1 \leq i \leq k+1}$ depends on ρ , k , and the a_i 's. We denote by $\mathcal{D}(a_2, \dots, a_{k+1}) = d_{k+1}$. Now set, for every $k \geq 1$,

$$(9) \quad \kappa_\rho^c(k) = \inf_{0 \leq a_2, \dots, a_{k+1} < 1} \max \left(\left(\frac{4\rho}{(1+\rho)^2 \sqrt{\prod_{2 \leq i \leq k+1} (1-a_i^2)}} \right)^{\frac{1}{k+1}}, \frac{2\rho}{\mathcal{D}(a_2, \dots, a_{k+1})} \right).$$

Finally, let:

$$(10) \quad \kappa_\rho^c = \inf_{k \geq 1} \kappa_\rho^c(k).$$

We give some intuition on κ_ρ^c in Section 2.1, make a few further comments and simulations for $\kappa_\rho^c(1)$, $\kappa_\rho^c(2)$ and $\kappa_\rho^c(3)$ in Section 2.2. We will see that $\kappa_\rho^c \in (0, 1)$ for all $\rho > 1$. This κ_ρ^c gives the asymptotic behavior of $\tilde{\lambda}_d^c(\mu_d)$ when μ charges two distinct points :

Theorem 1.4. *Let $b > a > 0$, $\alpha > 0$ and $\beta > 0$. Set $\mu = \alpha\delta_a + \beta\delta_b$ and $\rho = b/a > 1$. Then*

$$(11) \quad \lim_{d \rightarrow +\infty} \frac{1}{d} \ln \left(\tilde{\lambda}_d^c(\mu_d) \right) = \ln(\kappa_\rho^c) < 0.$$

Remember that in the case of a deterministic radius (assumptions of Theorem 1.1), the first order of the asymptotic behavior of the critical intensity in high dimensions is given by the threshold of the associated Galton-Watson process.

In the case of two distinct radii (assumptions of Theorem 1.4), the associated Galton-Watson process is two-type, one for each radius. At first order and on a logarithmic scale, the asymptotic behavior of the critical intensity in high dimensions is given by κ_ρ^c , and we could expect that, as in the deterministic case, κ_ρ^c would be given by the threshold of this two-type Galton-Watson process. We will see in Lemma 2.1 that it is actually the case, but only when ρ is not too far from 1. When ρ is large, this is no longer the case and geometry still plays a role. See the discussion in Section 2.2.

The proof of Theorem 1.4 is given in Section 2. The main ideas of the proofs are give in Section 2.1. Theorem 1.2 is an easy consequence of Theorem 1.4. The proof is given in Section 3. Corollary 1.3 is a straightforward consequence of Theorems 1.1 and 1.2. We give some intuition and results about κ_ρ^c is Section 2.1 and 2.2.

2. THE CASE WHEN THE RADII TAKE TWO VALUES

2.1. Notations and ideas of the proof. In the whole proof, we fix $\rho > 1$ and $\kappa > 0$.

Once the dimension $d \geq 1$ is given, we consider two independent stationary Poisson point processes on \mathbb{R}^d : χ_1 and χ_ρ , with respective intensities

$$\lambda_1 = \frac{\kappa^d}{v_d 2^d} \quad \text{and} \quad \lambda_\rho = \frac{\kappa^d}{v_d 2^d \rho^d}.$$

With χ_1 and χ_ρ , we respectively associate the two Boolean models

$$\Sigma_1 = \bigcup_{x \in \chi_1} B(x, 1) \quad \text{and} \quad \Sigma_\rho = \bigcup_{x \in \chi_\rho} B(x, \rho).$$

Note that Σ_ρ is an independent copy of $\rho\Sigma_1$. Note also that the expected number of balls of Σ_1 that touches a given ball of radius 1 is κ^d . Thus the expected number of balls of Σ_ρ that touches a given ball of radius ρ is also κ^d .

We focus on the percolation properties of the following two-type Boolean model

$$\Sigma = \Sigma_1 \cup \Sigma_\rho.$$

We begin by studying the existence of infinite k -alternating paths. For $k \geq 1$, an infinite k -alternating path is an infinite path made of balls such that the radius of the first ball is ρ , the radius of the next k balls is 1, the radius of the next ball is ρ and so on. For a fixed $k \geq 1$, we wonder whether infinite k -alternating paths exist and seek the critical threshold $\kappa_\rho^c(k)$ for their existence. A natural first step is to study the following quantities :

$$(12) \quad N_0 = \#\{x_1 \in \chi_\rho : \|x_1\| < 2\rho\},$$

$$\text{and for } k \geq 1, \quad N_k = \#\left\{ \begin{array}{l} x_{k+1} \in \chi_\rho : \exists (x_i)_{1 \leq i \leq k} \in \chi_1 \text{ distinct such that} \\ \|x_1\| < 1 + \rho, \forall i \in \{1, \dots, k-1\} \|x_{i+1} - x_i\| < 2, \\ \|x_{k+1} - x_k\| < 1 + \rho \end{array} \right\}.$$

Fix $k \geq 1$. Remember that $\kappa_\rho^c(k)$ is defined in (9).

A lower bound for $\kappa_\rho^c(k)$. In Subsection 2.3, we obtain lower bounds for $\kappa_\rho^c(k)$ by looking for upper bounds for $E(N_k)$. On one side, a natural genealogy is associated to the definition of N_k (see also the comments below Theorem 1.1 and in Section 2.2). We start with an ancestor x_0 located at the origin. We then seek his children in $\chi_1 \cap B(x_0, 1 + \rho)$: they constitute the first generation. If x_1 is one of those children, we then seek the children of x_1 in $\chi_1 \cap B(x_1, 2)$ to build the second generation and so on. On the other side, the process lives in \mathbb{R}^d and the geometry induces dependences: if x_1 and x'_1 are two individuals of the first generation, their children are a priori dependent. If we forget geometry and only consider genealogy, we get the following upper bound:

$$E(N_k) \leq \lambda_1 |B(\cdot, 1 + \rho)| \left(\prod_{i=2}^k \lambda_1 |B(\cdot, 2)| \right) \lambda_\rho |B(\cdot, 1 + \rho)|.$$

But the points of the last generation are in $B(0, 2\rho + 2k)$. So if we forget genealogy and only consider geometry we get the following upper bound:

$$E(N_k) \leq \lambda_\rho |B(0, 2\rho + 2k)|.$$

Expliciting the two previous bounds and combining them together, we get:

$$E(N_k) \leq \max \left(\frac{\kappa^{k+1}(1 + \rho)^2}{4\rho}, \frac{\kappa(\rho + k)}{\rho} \right)^d.$$

In this upper bound, the first argument of the maximum is due to genealogy while the second one is due to geometry. To get the geometrical term, we considered the worst case: the one in which, at each generation i , x_i is as far from the origin as possible. This gives a very poor bound. To get a better bound, we proceed as follows. Fix $a_2, \dots, a_{k+1} \in [0, 1)$. As before, we set $r_1 = r_{k+1} = 1 + \rho$, and for $i \in \{2, \dots, k\}$, $r_i = 2$ and we build the increasing sequence of distances $(d_i)_{1 \leq i \leq k+1}$ by setting $d_1 = 1 + \rho$ and, for every $i \in \{2, \dots, k+1\}$:

$$d_i^2 = d_{i-1}^2 + 2r_i a_i d_{i-1} + r_i^2.$$

See Figure 2 for a better understanding of these distances d_i .

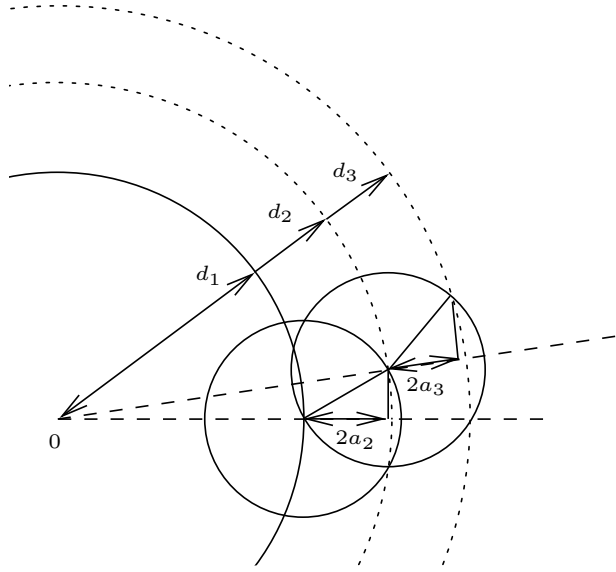


FIGURE 2. Definition of the distances d_i . Circles in plain line are of radius $1 + \rho$ and 2.

Denote by $\tilde{N}_k(a_2, \dots, a_{k+1})$ the number of points $x_{k+1} \in \chi_\rho$ for which there exists a path x_1, \dots, x_k fulfilling the same requirement as for N_k and such that $\|x_i\| \approx d_i$ for all i . Proceeding as before, we obtain the following upper bound:

$$E(\tilde{N}_k(a_2, \dots, a_{k+1})) \lesssim \min \left(\kappa^{k+1} \frac{(1 + \rho)^2}{4\rho} \sqrt{\prod_{2 \leq i \leq k+1} (1 - a_i)^2}, \frac{\kappa \mathcal{D}(a_2, \dots, a_{k+1})}{2\rho} \right)^d.$$

Here again, the first argument of the maximum is due to genealogy while the second one is due to geometry. Optimizing then on the a_i 's, we get:

$$E(N_k) \lesssim \sup_{a_2, \dots, a_{k+1}} \min \left(\kappa^{k+1} \frac{(1 + \rho)^2}{4\rho} \sqrt{\prod_{2 \leq i \leq k+1} (1 - a_i)^2}, \frac{\kappa \mathcal{D}(a_2, \dots, a_{k+1})}{2\rho} \right)^d.$$

A precise statement is given in Lemma 2.5. The precise value of the threshold $\kappa_\rho^c(k)$ given in (9) is then the value such that the above upper bound converges to 0 when $\kappa < \kappa_\rho^c(k)$. This heuristic will be precised in Subsection 2.3: we will prove there that when $\kappa < \kappa_\rho^c(k)$, $E(N_k)$ converges to 0 as d tends to infinity, and this will imply that there exists no infinite k -alternating path.

An upper bound for $\kappa_\rho^c(k)$. If, on the contrary, $\kappa > \kappa_\rho^c(k)$, then we will prove that $E(N_k)$ does not converge to 0. Actually, to prove that when $\kappa > \kappa_\rho^c(k)$, there exist infinite k -alternating paths, we will show, in Subsection 2.4, the following stronger property : With a probability that converges to 1 as d tends to infinity, we can find a path which fulfills the requirements of the definition of N_k – or more precisely of $\tilde{N}_k(a_2, \dots, a_{k+1})$ for some a_2, \dots, a_{k+1} nearly optimal – and which fulfills some extra conditions on the positions of the balls. This is Proposition 2.9 and this is the main technical part of this paper. Those extra conditions provide independence properties between the existence of different paths of the same kind. We can then show the existence of many such paths and concatenate some of them to build an infinite k -alternating path. Technically, the last step is achieved by comparing our model with a supercritical oriented percolation process on \mathbb{Z}^2 . In this comparison, an open bond in the oriented percolation process corresponds to one of the above paths in our model. This comparison with oriented percolation was already the last step in the paper of Penrose [14].

From infinite k -alternating paths to infinite paths. Recall $\kappa_\rho^c = \inf_{k \geq 1} \kappa_\rho^c(k)$. With the previous results, it is rather easy to show that there is no percolation for d large enough as soon as $\kappa < \kappa_\rho^c$. When $\kappa > \kappa_\rho^c$ then $\kappa > \kappa_\rho^c(k)$ for a $k \geq 1$. Therefore there is k -alternating percolation and therefore there is percolation.

2.2. The threshold κ_ρ^c and the associated Galton-Watson process. We begin with a couple of results about $\kappa_\rho^c(1)$ and κ_ρ^c .

Lemma 2.1. *Let $\rho > 1$. Then*

- $0 < \kappa_\rho^c(1) < 1$. More precisely :

$$\text{If } 1 < \rho \leq 2 \text{ then } \kappa_\rho^c(1) = \frac{2\sqrt{\rho}}{1+\rho}, \quad \text{while if } \rho \geq 2 \text{ then } \kappa_\rho^c(1) = \frac{\sqrt{4+\rho^2}}{1+\rho}.$$

- $0 < \kappa_\rho^c < 1$.
- There exists $\rho_0 > 2$ such that if $\rho \leq \rho_0$, then $\kappa_\rho^c = \kappa_\rho^c(1)$. This implies

$$\text{If } 1 < \rho \leq 2 \text{ then } \kappa_\rho^c = \frac{2\sqrt{\rho}}{1+\rho}, \quad \text{while if } 2 \leq \rho \leq \rho_0 \text{ then } \kappa_\rho^c = \frac{\sqrt{4+\rho^2}}{1+\rho}.$$

On Figure 3, we plot $\kappa_\rho^c(i)$, for $i \in \{1, 2, 3\}$. The data come from numerical estimations.

Proof. • By definition, $\kappa_\rho^c(1) = \inf_{0 \leq a < 1} \max(\phi_1(a), \phi_2(a))$, where $\phi_1, \phi_2 : [0, 1) \rightarrow \mathbb{R}$ are defined by:

$$\phi_1(a) = \frac{2\sqrt{\rho}}{(1+\rho)(1-a^2)^{1/4}} \quad \text{and} \quad \phi_2(a) = \frac{\rho\sqrt{2}}{(1+\rho)\sqrt{1+a}}.$$

If $\rho \leq 2$ then $\phi_1(0) \geq \phi_2(0)$. As ϕ_1 is increasing and ϕ_2 is decreasing, we get:

$$\kappa_\rho^c(1) = \inf_{0 \leq a < 1} \phi_1(a) = \phi_1(0) = \frac{\sqrt{2\rho}}{1+\rho}.$$

Assume, on the contrary, $\rho \geq 2$. Set

$$a = \frac{\rho^2 - 4}{\rho^2 + 4} \in [0, 1).$$

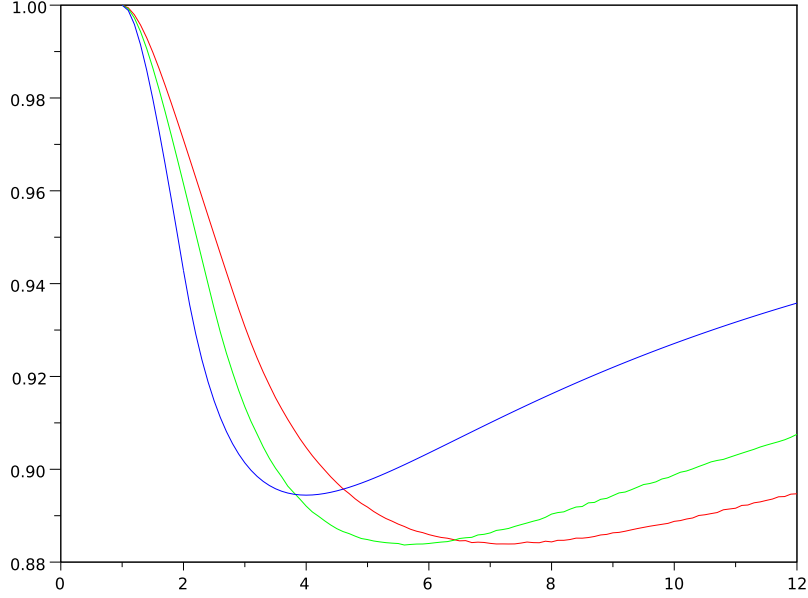


FIGURE 3. Simulations for $\kappa_\rho^c(i)$, $i \in \{1, 2, 3\}$, from left to right.

Then $\phi_1(a) = \phi_2(a)$. As ϕ_1 is increasing and ϕ_2 is decreasing, we get:

$$\kappa_\rho^c(1) = \phi_1(a) = \phi_2(a) = \frac{\sqrt{4 + \rho^2}}{1 + \rho}.$$

- Clearly we have, for every $k \geq 1$:

$$\kappa_\rho^c(k) \geq \inf_{0 \leq a_2, \dots, a_{k+1} < 1} \left(\frac{4\rho}{(1 + \rho)^2 \sqrt{\prod_{2 \leq i \leq k+1} (1 - a_i^2)}} \right)^{1/(k+1)} = \left(\frac{4\rho}{(1 + \rho)^2} \right)^{1/(k+1)}.$$

Therefore, as $\kappa_\rho^c = \inf_{k \geq 1} \kappa_\rho^c(k)$,

$$0 < \left(\frac{4\rho}{(1 + \rho)^2} \right)^{1/2} \leq \kappa_\rho^c \leq \kappa_\rho^c(1) < 1.$$

- The last inequalities imply that $\kappa_\rho^c = \kappa_\rho^c(1)$ if $1 < \rho \leq 2$. In fact, as soon as

$$(13) \quad \kappa_\rho^c(1) \leq \left(\frac{4\rho}{(1 + \rho)^2} \right)^{1/3},$$

it is true that $\kappa_\rho^c(1) \leq \kappa_\rho^c(k)$ for all $k \geq 2$ and therefore that $\kappa_\rho^c = \kappa_\rho^c(1)$. As the inequality in (13) is strict for $\rho = 2$, we obtain by continuity the existence of $\rho_0 > 2$ such that for every $\rho \in (1, \rho_0)$, $\kappa_\rho^c = \kappa_\rho^c(1)$. \square

The associated Galton-Watson process. The associated Galton-Watson process is a two type Galton-Watson process, one type for each radius. Consider for example the offspring distribution of type ρ of an individual of type 1. We define it to be the number of balls of a Boolean model directed by $\lambda_\rho \delta_\rho$ that intersects a given ball of radius 1. Therefore, this is a Poisson random variable with mean $\lambda_\rho v_d (1 + \rho)^d$. The other offspring distribution

are defined similarly. We moreover assume the the offspring of type 1 and ρ of a given individual are independent.

Consider the matrix of means of offspring distributions:

$$M_d = \begin{pmatrix} \lambda_1 v_d (1+1)^d & \lambda_\rho v_d (1+\rho)^d \\ \lambda_1 v_d (1+\rho)^d & \lambda_\rho v_d (\rho+\rho)^d \end{pmatrix}.$$

For example, the top-right term is the mean offspring of type ρ of an individual of type 1. We have:

$$M_d = \kappa^d \begin{pmatrix} 1 & \left(\frac{1+\rho}{2}\right)^d \\ \left(\frac{1+\rho}{2}\right)^d & 1 \end{pmatrix}.$$

Let r_d denotes the largest eigenvalue of M_d . The extinction probability of the Galton-Watson process is 1 if and only if $r_d \leq 1$. We have:

$$r_d \sim \left(\frac{\kappa(1+\rho)}{2\sqrt{\rho}} \right)^d.$$

If we fix κ and let d tends to infinity, the critical κ for the Galton-Watson process is thus :

$$\kappa = \frac{2\sqrt{\rho}}{1+\rho}.$$

With Theorem 1.4 and Lemma 2.1 we thus see that the comparison with the Galton-Watson is asymptotically sharp on a logarithmic scale only when $1 < \rho \leq 2$.

2.3. Subcritical phase. Let $\rho > 1$ be fixed. We consider, in \mathbb{R}^d , the two-type Boolean model Σ introduced in Subsection 2.1, with radii 1 and ρ and respective intensities

$$\lambda_1 = \frac{\kappa^d}{v_d 2^d} \text{ and } \lambda_\rho = \frac{\kappa^d}{v_d 2^d \rho^d}$$

depending on some $\kappa \in (0, 1)$. The aim of this subsection is to prove the following proposition:

Proposition 2.2. *Let $\rho > 1$ be fixed. If $\kappa < \kappa_\rho^c$, then, as soon as the dimension d is large enough, percolation does not occur in the two-type Boolean model Σ .*

In the following of this subsection, we fix $\rho > 1$ and $0 < \kappa < \kappa_\rho^c$.

We start with an elementary upper bound, in which we do not take into account the geometrical constraints. We recall that the N'_k 's have been introduced in (12).

Lemma 2.3. *$E(N_0) = \kappa^d$ and, for $k \geq 1$, $E(N_k) \leq \left(\frac{\kappa^{k+1}(1+\rho)^2}{4\rho} \right)^d$.*

Proof. The result for N_0 follows directly from the equality $E(N_0) = \lambda_\rho |B(0, 2\rho)|$.

Take now $k \geq 1$. We have :

$$(14) \quad E(N_k) \leq \lambda_1 |B(\cdot, 1+\rho)| \left(\prod_{i=2}^k \lambda_1 |B(\cdot, 2)| \right) \lambda_\rho |B(\cdot, 1+\rho)|$$

where $B(\cdot, r)$ stands for a ball with radius r and center unspecified. This can for instance be seen as follows :

$$\begin{aligned} E(N_k) &\leq E \left(\sum_{\substack{x_1, \dots, x_k \in \chi_1 \\ \text{distinct}, x_{k+1} \in \chi_\rho}} \mathbf{1}_{x_1 \in B(0, 1+\rho)} \cdots \mathbf{1}_{x_{k+1} \in B(x_k, 1+\rho)} \right) \\ &= \lambda_1^k \lambda_\rho \int_{\mathbb{R}^{d(k+1)}} dx_1 \cdots dx_{k+1} \mathbf{1}_{x_1 \in B(0, 1+\rho)} \cdots \mathbf{1}_{x_{k+1} \in B(x_k, 1+\rho)}, \end{aligned}$$

which gives (14). The lemma follows. \square

To give a more accurate upper bound for the N_k 's, we are going to cut the balls into slices and to estimate which slices give the main contribution. For $x \in \mathbb{R}^d \setminus \{0\}$, $0 \leq a < b \leq 1$ and $r > 0$, we now define :

$$\begin{aligned} \text{If } a > 0 : \quad B(x, r, a, b) &= \left\{ y \in \mathbb{R}^d : \|y - x\| \leq r \text{ and } ar < \left\langle y - x, \frac{x}{\|x\|} \right\rangle \leq br \right\}, \\ \text{If } a = 0 : \quad B(x, r, a, b) &= \left\{ y \in \mathbb{R}^d : \|y - x\| \leq r \text{ and } \left\langle y - x, \frac{x}{\|x\|} \right\rangle \leq br \right\}. \end{aligned}$$

The next lemma gives asymptotics for the volume of these sets:

Lemma 2.4. For $x \in \mathbb{R}^d \setminus \{0\}$, $0 \leq a < b \leq 1$ and $r > 0$,

$$\lim_{d \rightarrow +\infty} \frac{1}{d} \ln \left(\frac{|B(x, r, a, b)|}{v_d} \right) = \ln \left(r \sqrt{1 - a^2} \right).$$

Actually we will only use :

$$\limsup_{d \rightarrow +\infty} \frac{1}{d} \ln \left(\frac{|B(x, r, a, b)|}{v_d} \right) \leq \ln \left(r \sqrt{1 - a^2} \right).$$

Proof of Lemma 2.4. Note that it is sufficient to prove the lemma for $x = e_1$, first vector of the canonical basis, and $r = 1$.

First, if $a = 0$, the result follows directly from the inequality $v_d/2 \leq |B(e_1, 1, 0, b)| \leq v_d$. Assume next that $a > 0$. On the one hand, $B(e_1, 1, a, b)$ is included in the cylinder

$$\{x = (x_1, \dots, x_d) : x_1 \in [a, 1] \text{ and } \|(0, x_2, \dots, x_d)\| \leq \sqrt{1 - a^2}\},$$

which implies

$$(15) \quad |B(e_1, 1, a, b)| \leq v_{d-1} \sqrt{1 - a^2}^{d-1} (1 - a).$$

On the other end, by convexity, $B(e_1, 1, a, b)$ contains the following difference between two homothetical cones:

$$\{x = (x_1, \dots, x_d) : x_1 \in [a, b] \text{ and } \|(0, x_2, \dots, x_d)\| \leq \sqrt{1 - a^2} \frac{1 - x_1}{1 - a}\},$$

which implies

$$(16) \quad \frac{v_{d-1} \sqrt{1 - a^2}^{d-1}}{d} ((1 - a) - (1 - b)) \leq |B(e_1, 1, a, b)|.$$

The lemma follows from (15) and (16). \square

We can now improve the control given in Lemma 2.3:

Lemma 2.5. For every $k \geq 1$,

$$\begin{aligned} & \limsup_{d \rightarrow \infty} \frac{1}{d} \ln(E(N_k)) \\ & \leq \ln \left(\sup_{0 \leq a_2, \dots, a_{k+1} < 1} \min \left(\kappa^{k+1} \frac{(1+\rho)^2}{4\rho} \sqrt{\prod_{2 \leq i \leq k+1} (1-a_i^2)}, \frac{\kappa \mathcal{D}(a_2, \dots, a_{k+1})}{2\rho} \right) \right) \end{aligned}$$

Proof. • Fix $N \geq 1$. Note that the ball $B(x, r)$ is the disjoint union of the slices $B(x, r, n/N, (n+1)/N)$ for $n \in \{0, \dots, N-1\}$. For any $n_2, \dots, n_{k+1} \in \{0, \dots, N-1\}$, we set

$$a_i = \frac{n_i}{N} \text{ and } a_i^+ = \frac{n_i + 1}{N}.$$

We focus on the contribution from a specific product of slices :

$$N_k(n_2, \dots, n_{k+1}) = \# \left\{ \begin{array}{l} x_{k+1} \in \chi_\rho : \exists (x_i)_{1 \leq i \leq k} \in \chi_1 \text{ distinct with} \\ \|x_1\| < 1 + \rho, \forall i \in \{1, \dots, k-1\} x_{i+1} \in B(x_i, 2, a_{i+1}, a_{i+1}^+), \\ x_{k+1} \in B(x_k, 1 + \rho, a_{k+1}, a_{k+1}^+) \end{array} \right\}.$$

Then we have :

$$(17) \quad N_k \leq \sum N_k(n_2, \dots, n_{k+1}),$$

where the sum is over $(n_2, \dots, n_{k+1}) \in \{0, \dots, N-1\}^k$.

• As we can check that the points contributing to $N_k(n_2, \dots, n_{k+1})$ are in $B(0, \mathcal{D}(a_2^+, \dots, a_{k+1}^+))$, we get :

$$E(N_k(n_2, \dots, n_{k+1})) \leq \lambda_\rho v_d \mathcal{D}(a_2^+, \dots, a_{k+1}^+)^d,$$

this leads to :

$$(18) \quad \limsup_{d \rightarrow +\infty} \frac{1}{d} \ln(E(N_k(n_2, \dots, n_{k+1}))) \leq \ln \left(\frac{\kappa \mathcal{D}(a_2^+, \dots, a_{k+1}^+)}{2\rho} \right).$$

• Besides, proceeding as in the proof of Lemma 2.3, we obtain :

$$E(N_k(n_2, \dots, n_{k+1})) \leq \lambda_1 |B(0, 1 + \rho)| \left(\prod_{i=2}^k \lambda_1 |B(\cdot, 2, a_i, a_i^+)| \right) \lambda_\rho |B(\cdot, 1 + \rho, a_{k+1}, a_{k+1}^+)|.$$

With Lemma 2.4, we deduce :

$$(19) \quad \limsup_{d \rightarrow \infty} \frac{1}{d} \ln E(N_k(n_2, \dots, n_{k+1})) \leq \ln \left(\kappa^{k+1} \frac{(1+\rho)^2}{4\rho} \sqrt{\prod_{2 \leq i \leq k+1} (1-a_i^2)} \right).$$

• From (17), (18) and (19) we finally get :

$$\begin{aligned} & \limsup_{d \rightarrow +\infty} \frac{\ln(E(N_k))}{d} \\ & \leq \ln \left(\max_{a_2, \dots, a_{k+1} \in \{0, \dots, \frac{N-1}{N}\}} \min \left(\kappa^{k+1} \frac{(1+\rho)^2}{4\rho} \sqrt{\prod_{2 \leq i \leq k+1} (1-a_i^2)}, \frac{\kappa \mathcal{D}(a_2^+, \dots, a_{k+1}^+)}{2\rho} \right) \right). \end{aligned}$$

As \mathcal{D} is uniformly continuous on $[0, 1]^k$, we end the proof by taking the limit when N goes to $+\infty$. \square

The next step consists in taking into account all $k \geq 0$ simultaneously; we thus introduce

(20)

$$N = \# \left(\left\{ \begin{array}{l} y \in \chi_\rho : \exists k \geq 1, \exists (x_i)_{1 \leq i \leq k} \in \chi_1 \text{ distinct with} \\ \|x_1\| < 1 + \rho, \forall i \in \{1, \dots, k-1\} \|x_{i+1} - x_i\| < 2, \\ \|y - x_k\| < 1 + \rho \end{array} \right\} \cup \{y \in \chi_\rho : \|y\| < 2\rho\} \right).$$

Lemma 2.6. *If $\kappa < \kappa_\rho^c$, then $\limsup_{d \rightarrow +\infty} \frac{1}{d} \ln(E(N)) < 0$.*

Proof. We have :

$$(21) \quad E(N) \leq \sum_{k \geq 0} E(N_k).$$

As $\kappa < \kappa_\rho^c$, Lemma 2.5 ensures that for every $k \geq 1$:

$$(22) \quad \limsup_{d \rightarrow +\infty} \frac{1}{d} \ln(E(N_k)) < 0.$$

Moreover, the assumption $\kappa < \kappa_\rho^c$ also implies, thanks to Lemma 2.1, that $\kappa < 1$. We can then choose k_0 large enough to have :

$$\frac{\kappa^{k_0+1}(1+\rho)^2}{4\rho} \leq \exp(-1).$$

With Lemma 2.3, we thus get :

$$E(N_0) + \sum_{k \geq k_0} E(N_k) \leq \kappa^d + \exp(-d) \sum_{k \geq 0} \kappa^{kd} = \kappa^d + \exp(-d) \frac{1}{1 - \kappa^d}.$$

With (21) and (22), this ends the proof. \square

The next lemma is elementary

Lemma 2.7. *Assume $\kappa < 1$.*

Then the connected components of $\bigcup_{x \in \chi_1} B(x, 1)$ are bounded with probability 1.

Proof. For any integer $k \geq 0$, denote by $M(k)$ the number of balls with radius 1 linked to $B(0, 1)$ by a chain of k distinct balls with radius 1. Proceeding as in the proof of Lemma 2.3, we get :

$$E(M_k) \leq \kappa^{d(k+1)}.$$

Now denote by M the number of balls with radius 1 linked to $B(0, 1)$ by a chain of (perhaps no) balls with radius 1. Then :

$$E(M) \leq E \left(\sum_{k \geq 0} M_k \right) = \frac{\kappa^d}{1 - \kappa^d} < +\infty.$$

Therefore, M is finite with probability 1. So the connected components that touch $B(0, 1)$ are bounded with probability 1. So with probability 1, every connected component is bounded. \square

Proof of Proposition 2.2. Remember that we proved in Lemma 2.1 that $\kappa_\rho^c < 1$. Take κ such that $0 < \kappa < \kappa_\rho^c < 1$.

Let ξ_1 be the set of random balls with radius ρ that can be connected to $B(0, \rho)$ through a chain of random balls with radius 1 (we consider the condition as fulfilled if the ball touches $B(0, \rho)$ directly). Let ξ_2 be the set of random balls with radius ρ that are not in

ξ_1 , but that can be connected to $B(0, \rho)$ through a path of random balls in which there is only one ball with radius ρ . We define similarly ξ_3, ξ_4 and so on and denote by ξ the disjoint union of all these sets.

We have $\#\xi_1 = N$. (Remember that N has been defined in (20).) By Lemma 2.6, we have :

$$\limsup_{d \rightarrow +\infty} \frac{1}{d} \ln(E(\xi_1)) < 0.$$

Take some $\mu > 0$ and assume from now on that d is large enough to have

$$\frac{1}{d} \ln(E(\#\xi_1)) \leq -\mu.$$

Pour tout $k \geq 1$, on a alors : $E(\#\xi_k) \leq (E(\#\xi_1))^k \leq \exp(-dk\mu)$.

As $\xi = \cup_{k \geq 1} \xi_k$, we deduce from the previous inequalities that ξ is finite with probability 1. So any unbounded connected component of $\Sigma_1 \cup \Sigma_\rho$ that touches $B(0, \rho)$ is included in Σ_1 . As $\kappa < 1$, Lemma 2.7 rules out the possibility of an unbounded connected component in Σ_1 . So with probability 1, the connected components of $\Sigma_1 \cup \Sigma_\rho$ that touch $B(0, \rho)$ are bounded, which ends the proof. \square

2.4. Supercritical phase. We fix here $\rho > 1$. We consider once again the two-type Boolean model Σ introduced in Subsection 2.1 and we fix an integer $k \geq 1$.

For every $n \geq 0$, we set $R_n = \rho$ if $k + 1$ divides n and $R_n = 1$ otherwise. We say that percolation by k -alternation occurs if there exists an infinite sequence of distinct points $(x_n)_{n \in \mathbb{N}}$ in \mathbb{R}^d such that, for every $n \geq 0$:

- $x_n \in \chi_{R_n}$.
- $B(x_n, R_n) \cap B(x_{n+1}, R_{n+1}) \neq \emptyset$.

In other words, percolation by k -alternation occurs if there exists an infinite path along which k balls of radius 1 alternate with one ball of radius ρ , *ie* if there exists an infinite k -alternating path. The aim of this subsection is to prove the following proposition:

Proposition 2.8. *Let $\rho > 1$ and $k \geq 1$ be fixed. Assume that $\kappa \in (\kappa_\rho^c(k), 1)$. If the dimension d is large enough, then percolation by k -alternation occurs with probability one.*

As announced in subsection 2.1, percolation by k -alternation of the two-type Boolean model in the supercritical case will be proved by embedding in the model a supercritical 2-dimensional oriented percolation process.

We thus specify the two first coordinates, and introduce the following notations. When $d \geq 3$, for any $x \in \mathbb{R}^d$, we write

$$x = (x', x'') \in \mathbb{R}^2 \times \mathbb{R}^{d-2}.$$

We write $B'(c, r)$ for the open Euclidean balls of \mathbb{R}^2 with center $c \in \mathbb{R}^2$ and radius $r > 0$. In the same way we denote by $B''(c, r)$ the open Euclidean balls of \mathbb{R}^{d-2} with center $c \in \mathbb{R}^{d-2}$ and radius $r > 0$.

2.4.1. One step in the 2-dimensional oriented percolation model. The point here is to define the event that will govern the opening of the edges in the 2-dimensional oriented percolation process : it is naturally linked to the existence of a finite path composed of k balls of radius 1 and a ball of radius ρ , whose positions of centers are specified.

We define, for a given dimension d , the two following subsets of \mathbb{R}^d :

$$\begin{aligned} W &= d^{-1/2} ((-1, 1) \times (-1, 0) \times \mathbb{R}^{d-2}), \\ W^+ &= d^{-1/2} ((0, 0) \times (0, 1) \times \mathbb{R}^{d-2}). \end{aligned}$$

For $x \in W$ we set :

$$(23) \quad \mathcal{G}^+(x) = \left\{ \begin{array}{l} \text{There exist distinct } x_1, \dots, x_k \in \chi_1 \cap W^+ \text{ and } x_{k+1} \in \chi_\rho \cap W^+ \\ \text{such that } x_0, x_1, \dots, x_{k+1} \text{ is a path} \end{array} \right\}.$$

Our goal here is to prove that the probability of occurrence of this event is asymptotically large :

Proposition 2.9. *Let $\rho > 1$ and $k \geq 1$ be a fixed. Assume that $\kappa \in (\kappa_\rho^c(k), 1)$ and choose $p \in (0, 1)$. If the dimension d is large enough, then for every $x \in W$,*

$$P(\mathcal{G}^+(x)) \geq p.$$

Note already that by translation invariance, $P(\mathcal{G}^+(x))$ does not depend on x'' , so we can assume without loss of generality that $x'' = 0$. In the sequel of this subsection, $\rho > 1$ and $k \geq 1$ are fixed.

We first recall the definitions of the $(d_i)_{1 \leq i \leq k+1}$ and of $\kappa_\rho^c(k)$ we give in the introduction. We set $r_1 = r_{k+1} = 1 + \rho$ and for $i \in \{2, \dots, k\}$, $r_i = 2$. Then, for a given sequence $(a_i)_{2 \leq i \leq k+1} \in [0, 1]^k$, we build an increasing sequence $(d_i)_{1 \leq i \leq k+1}$ by setting $d_1 = 1 + \rho$ and, for every $i \in \{2, \dots, k+1\}$:

$$d_i^2 = d_{i-1}^2 + 2r_i a_i d_{i-1} + r_i^2.$$

Finally, we note $\mathcal{D}(a_2, \dots, a_{k+1}) = d_{k+1}$ and we set

$$\kappa_\rho^c(k) = \inf_{0 \leq a_2, \dots, a_{k+1} < 1} \max \left(\left(\frac{4\rho}{(1+\rho)^2 \sqrt{\prod_{2 \leq i \leq k+1} (1-a_i^2)}} \right)^{1/(k+1)}, \frac{2\rho}{\mathcal{D}(a_2, \dots, a_{k+1})} \right).$$

The first step consists in choosing a nearly optimal sequence $(a_i)_{2 \leq i \leq k+1} \in [0, 1]^k$ satisfying some extra inequalities :

Lemma 2.10. *We can choose $(a_i)_{2 \leq i \leq k+1} \in [0, 1]^k$ such that :*

$$(24) \quad 1 < \kappa^{k+1} \frac{(1+\rho)^2}{4\rho} \sqrt{\prod_{2 \leq j \leq k+1} (1-a_j^2)} < \kappa \frac{d_{k+1}}{2\rho}.$$

Proof. As $\kappa_\rho^c(k) < \kappa$, we can choose $(a_i^0)_{2 \leq i \leq k+1} \in (0, 1)^k$ such that the two following conditions

$$(25) \quad \kappa > \left(\frac{4\rho}{(1+\rho)^2 \sqrt{1-a_2^2} \dots \sqrt{1-a_{k+1}^2}} \right)^{\frac{1}{k+1}},$$

$$(26) \quad \kappa > \frac{2\rho}{d_{k+1}}$$

are fulfilled for $(a_i)_i = (a_i^0)_i$. We fix $(a_3, \dots, a_{k+1}) = (a_3^0, \dots, a_{k+1}^0)$. Note that

$$\begin{aligned} f : a_2 \mapsto \kappa \frac{d_{k+1}}{2\rho} & \text{ is continuous and increasing,} \\ g : a_2 \mapsto \kappa^{k+1} \frac{(1+\rho)^2}{4\rho} \sqrt{\prod_{2 \leq j \leq k+1} (1-a_j^2)} & \text{ is continuous and decreasing,} \end{aligned}$$

and that $\lim_{a_2 \rightarrow 1} g(a_2) = 0$. Moreover, Conditions (26) and (25) ensure that $f(a_2^0) > 1$ and $g(a_2^0) > 1$. Thus we only require that $f(a_2^0) > g(a_2^0)$ to prove the lemma.

If $f(a_2^0) \leq g(a_2^0)$, we can take $a_2 > a_2^0$ such that $1 < g(a_2) < f(a_2^0)$: then $f(a_2) \geq f(a_2^0) > g(a_2) > 1$ and the lemma is proved. \square

Note that (24) implies (25) and (26).

As explained in subsection 2.3, the main contribution to the number N_k of centers x_{k+1} of balls of radius ρ that are linked to a ball of radius ρ centered at the origin by a chain $(x_i)_{1 \leq i \leq k}$ of k balls of radius 1 – see the precise definition (12) – is obtained for $\|x_i\| \sim d_i$, where the d_i 's are build from a (nearly) optimal sequence $(a_i)_{2 \leq i \leq k+1} \in [0, 1]^k$. So we fix a nearly optimal family $(a_i)_{2 \leq i \leq k+1} \in (0, 1)^k$ satisfying (24), we build the associated family of distances $(d_i)_{1 \leq i \leq k+1} \in (0, 1)^k$ and we are going to look for a good sequence of centers $(x_i)_{1 \leq i \leq k+1} \in (0, 1)^k$ with $\|x_i\| \sim d_i$.

We thus introduce the following subsets of \mathbb{R}^2 :

$$\begin{aligned} D'_0 &= (-d^{-1/2}, d^{-1/2}) \times (-d^{-1/2}, 0), \\ \forall i \in \{1, \dots, k+1\} \quad D'_i &= (0, d^{-1/2})^2, \end{aligned}$$

and the following sets in \mathbb{R}^{d-2}

$$\begin{aligned} C''_0 &= \{0\}, \\ \forall i \in \{1, \dots, k+1\} \quad C''_i &= B''(0, d_i - 2d^{-1}) \setminus B''(0, d_i - 3d^{-1}). \end{aligned}$$

Finally, for $i \in \{0, \dots, k+1\}$, we set $C_i = D'_i \times C''_i$. Note that for d large enough, these sets are disjoint. The next lemma controls the asymptotics in the dimension d of the volume of these sets

Lemma 2.11. *For every $i \in \{1, \dots, k+1\}$:*

$$\lim_{d \rightarrow +\infty} \frac{1}{d} \ln \frac{|C''_i|}{v_{d-2}} = \lim_{d \rightarrow +\infty} \frac{1}{d} \ln \frac{|C_i|}{v_d} = \ln d_i.$$

Proof. This can be proven by elementary computations. \square

Each x_i will be taken in C_i , but we also have to ensure that the $(x_i)_{1 \leq i \leq k+1}$ form a path. Note that for $i \in \{2, \dots, k+1\}$, we have $0 < d_{i-1} + a_i r_i < d_i$, which legitimates the following definition. See also Figure 2. For $i \in \{2, \dots, k+1\}$ and d large enough, we denote by θ_i the unique real number in $(0, \pi/2)$ such that

$$\cos \theta_i = \frac{d_{i-1} + a_i r_i}{d_i} + d^{-1/2}.$$

We introduce next, for $y \in C_{i-1}$, the following subset of \mathbb{R}^{d-2} :

$$D''_i(y'') = \{z'' \in C''_i : \langle z'', y'' \rangle \geq \|y''\| \cdot \|z''\| \cdot \cos \theta_i\}$$

We also set $D''_0 = C''_0$ and $D''_1(y'') = C''_1$ for every $y \in C_0$. Finally, we define for every $i \in \{1, \dots, k+1\}$ and $y \in C_{i-1}$:

$$D_i(y) = D'_i \times D''_i(y'') \subset C_i,$$

and $D_0 = D'_0 \times D''_0$.

Lemma 2.12. • *If the dimension d is large enough, for every $i \in \{1, \dots, k+1\}$ and $y \in C_{i-1}$,*

$$D_i(y) \subset B(y, r_i) \cap C_i.$$

• *Let $x_0 \in D_0$. If there exist $X_1, \dots, X_k \in \chi_1$ and $X_{k+1} \in \chi_\rho$ such that $X_1 \in D_1(x_0), X_2 \in D_2(X_1), \dots, X_{k+1} \in D_{k+1}(X_k)$, then the event $\mathcal{G}_+(x_0)$ occurs.*

Proof. • The inclusion $D_i(y) \subset C_i$ is clear for every $i \in \{1, \dots, k+1\}$. Let $i \in \{2, \dots, k+1\}$, $y \in C_{i-1}$ and $z \in D_i(y)$. Then, as soon as d is large enough,

$$\begin{aligned} \|z - y\|^2 &= \|z' - y'\|^2 + \|z'' - y''\|^2 \\ &\leq \frac{2}{d} + \|y''\|^2 + \|z''\|^2 - 2 \langle y'', z'' \rangle \\ &\leq \frac{2}{d} + (d_{i-1} - 2d^{-1})^2 + (d_i - 2d^{-1})^2 - 2(d_{i-1} - 3d^{-1})(d_i - 3d^{-1}) \cos \theta_i \\ &\leq d_i^2 + d_{i-1}^2 - 2d_{i-1}(d_{i-1} + a_i r_i) - 2d^{-1/2} d_i d_{i-1} + O_i(d^{-1}) \\ &\leq r_i^2 - 2d^{-1/2} d_i d_{i-1} + O_i(d^{-1}) \leq r_i^2. \end{aligned}$$

Let now $y \in C_0$ and $z \in D_1(y)$. As $d_1 = 1 + \rho = r_1 > 2$, we obtain, for d large enough :

$$\|z - y\|^2 = \|z' - y'\|^2 + \|z'' - y''\|^2 \leq \frac{8}{d} + (d_1 - 2d^{-1})^2 \leq r_1^2.$$

• The second point is a simple consequence of the first point, of the fact that the sets $D_i(x_{i-1})$, as the sets C_i , are disjoint and of the definition of the event \mathcal{G}_+ . \square

Note that for $i \in \{1, \dots, k+1\}$, $|D_i(y)|$ and $|D_i''(y'')|$ do not depend on the choice of $y \in C_{i-1}$. We thus denote by $|D_i|$ and $|D_i''|$ these values. We now give asymptotic estimates for these values :

Lemma 2.13. *For every $i \in \{2, \dots, k+1\}$,*

$$\lim_{d \rightarrow +\infty} \frac{1}{d} \ln \frac{|D_i''|}{v_{d-2}} = \lim_{d \rightarrow +\infty} \frac{1}{d} \ln \frac{|D_i|}{v_d} = \ln(r_i \sqrt{1 - a_i^2}).$$

Proof. We have, by homogeneity and isotropy:

$$(27) \quad |D_i''| = ((d_i - 2d^{-1})^{d-2} - (d_i - 3d^{-1})^{d-2}) |S|$$

where $S = \{x = (x_1, \dots, x_{d-2}) \in B''(0, 1) : x_1 \geq \|x\| \cos(\theta_i)\}$.

But S is included in the cylinder

$$\{(x_i)_{1 \leq i \leq d-2} \in \mathbb{R}^{d-2} : x_1 \in [0, 1], \|(x_2, \dots, x_{d-2})\| \leq \sin(\theta_i)\}$$

and S contains the cone

$$\{(x_i)_{1 \leq i \leq d-2} \in \mathbb{R}^{d-2} : x_1 \in [0, \cos(\theta_i)], \|(x_2, \dots, x_{d-2})\| \leq x_1 \sin(\theta_i) \cos(\theta_i)^{-1}\}.$$

Therefore :

$$(28) \quad v_{d-3} \cos(\theta_i) \sin(\theta_i)^{d-3} (d-2)^{-1} \leq |S| \leq v_{d-3} \sin(\theta_i)^{d-3}.$$

From (27), (28), and the limits $\cos(\theta_i) \rightarrow (d_{i-1} + a_i r_i) d_i^{-1} \neq 0$ and $d_i \sin(\theta_i) \rightarrow r_i \sqrt{1 - a_i^2}$, we get

$$\lim \frac{1}{d} \ln \left(\frac{|D_i''|}{v_{d-2}} \right) = \ln(r_i \sqrt{1 - a_i^2}).$$

The lemma follows. Note that a direct calculus with spherical coordinates can also give the announced estimates. \square

Everything is now in place to prove Proposition 2.9.

Proof of Proposition 2.9. Choose $p < 1$ and $x \in W$ such that $x'' = 0$.

• We start with a single individual, encoded by its position $\zeta_0 = \{x\} \subset C_0$, and we build, generation by generation, its descendance : we set, for $1 \leq i \leq k$,

$$\zeta_i = \chi_1 \cap \bigcup_{y \in \zeta_{i-1}} D_i(y) \subset C_i,$$

and for the $(k+1)$ -th generation, we finally set

$$\zeta_{k+1} = \chi_\rho \cap \bigcup_{y \in \zeta_k} D_{k+1}(y) \subset C_{k+1}.$$

By Lemma 2.12, if $\zeta_{k+1} \neq \emptyset$ then the event $\mathcal{G}^+(x)$ occurs. To bound from below the probability that $\zeta_{k+1} \neq \emptyset$, we now build a simpler process ξ , stochastically dominated by η .

• We set $\alpha_i = \lambda_1 |D_i|$ for $i \in \{1, \dots, k\}$ and $\alpha_{k+1} = \lambda_\rho |D_{k+1}|$: thus, α_i is the mean number of children of a point of the $(i-1)$ -th generation. Let $X_0 = x$ be the position of the first individual. Let N be a Poisson random variable with parameter α_1 : This random variable gives the number of children of X_0 .

Then we complete X_0 in a family $X = (X_0, X_1, \dots, X_{k+1})$ of points in \mathbb{R}^d , independent of N , in the following manner : X_1 is taken uniformly in $D_1(X_0)$, then X_2 is taken uniformly in $D_2(X_1)$, and so on : The family X represents the positions of the possible successive descendents of x in a single branch. Let then $(X^j)_{j \geq 1}$ be independent copies of X . We will use the N first copies, one for each child of X_0 .

Let $Y = (Y_i^j)_{2 \leq i \leq k+1, j \geq 1}$ be an independent family of independent and identically distributed random variables, such that Y_i^j follows the Bernoulli law with parameter $1 - \exp(-\alpha_i)$, which is the probability that a Poisson random variable with parameter α_i is different from 0. The random variable Y_i^j encodes the fact that in the j -th branch, the member of the $(i-1)$ -th generation has children or not in a process in which we forget dependencies due to geometry : $Y_i^j = 1$ if and only if the member of the $(i-1)$ -th generation in the j -th branch has at least one child. In other words, with the Y_i^j , we only keep, for each individual from the first generation X_1 , at most one child, and its position is given by X_i^j . We set $J_1 = \{1, \dots, N\}$ and, for every $i \in \{2, \dots, k+1\}$:

$$J_i = \{1 \leq j \leq N : Y_2^j = \dots = Y_i^j = 1\}.$$

Thus the random set J_i gives the exponents of the branches that are, among the N initial branches, still alive at the i -th generation in a process with no dependencies due to geometry.

Until now, we did not take into account the geometrical constraints between individuals. For every $i \in \{2, \dots, k+1\}$ and every $j \geq 1$, we set

$$Z_i^j = 1 \text{ if } X_i^j \notin \bigcup_{j' \in J_{i-1} \setminus \{j\}} D_i(X_{i-1}^{j'}) \quad \text{and } Z_i^j = 0 \text{ otherwise.}$$

When building generation i from generation $i-1$, we explore the Poisson point processes in the area $\bigcup_{j \in J_{i-1}} D_i(X_{i-1}^j)$. By construction of the C_i , these areas are distinct for different generations. When $Z_i^j = 1$, we are sure that the member of the i -th generation in the j -th branch does not live in an already explored zone. Thus, one can check that, for every $i \in \{2, \dots, k+1\}$, the set

$$\xi_i = \{X_i^j : j \in J_i \text{ and } Z_2^j = \dots = Z_i^j = 1\}$$

is stochastically dominated by ζ_i . Thus to prove Proposition 2.9, we now need to bound from below the probability that ξ_{k+1} is not empty.

• Let T be the smallest integer j such that $Y_2^j = \dots = Y_{k+1}^j = 1$: in other words, T is the smallest exponent of a branch that lives till generation $k + 1$. To ensure that $\xi_{k+1} \neq \emptyset$, it is sufficient that $T \leq N$ and that $Z_2^T = \dots = Z_{k+1}^T = 1$. So :

$$\begin{aligned} 1 - P(\mathcal{G}^+(x)) &\leq P(\xi_{k+1} = \emptyset) \\ &\leq P(\#J_{k+1} = 0) + P\left(\{T \leq N\} \cap \bigcup_{2 \leq i \leq k+1} \{Z_i^T = 0\}\right) \\ &\leq P(\#J_{k+1} = 0) + \sum_{2 \leq i \leq k+1} P(T \leq N \text{ and } Z_i^T = 0). \end{aligned}$$

For every $2 \leq i \leq k + 1$, we have by construction :

$$\begin{aligned} P(T \leq N \text{ and } Z_i^T = 0) &= P(T \leq N, \exists j \in J_{i-1} \setminus \{T\} \text{ such that } X_i^T \in D_i(X_{i-1}^j)) \\ &\leq \sum_{j \geq 1} P(T \leq N \text{ and } j \in J_{i-1} \setminus \{T\} \text{ and } X_i^T \in D_i(X_{i-1}^j)) \\ &= \sum_{j \geq 1} E(1_{T \leq N} 1_{j \in J_{i-1} \setminus \{T\}} P(X_i^T \in D_i(X_{i-1}^j) | Y, N)) \\ &= \sum_{j \geq 1} E(1_{T \leq N} 1_{j \in J_{i-1} \setminus \{T\}}) P(X_i^1 \in D_i(X_{i-1}^2)) \\ &\leq E(\#J_{i-1}) P(X_i^1 \in D_i(X_{i-1}^2)). \end{aligned}$$

Besides, as $(X_i^1)''$ is uniformly distributed on C_i'' and is independent of $(X_{i-1}^2)''$,

$$(29) \quad P(X_i^1 \in D_i(X_{i-1}^2)) = P((X_i^1)'' \in D_i''((X_{i-1}^2)'')) = \frac{|D_i''|}{|C_i''|}.$$

This leads to

$$(30) \quad 1 - P(\mathcal{G}^+(x)) \leq P(\#J_{k+1} = 0) + \sum_{i=2}^{k+1} E(\#J_{i-1}) \frac{|D_i''|}{|C_i''|}.$$

• For $1 \leq i \leq k + 1$, the cardinality of J_i follows a Poisson law with parameter

$$\eta_i = \alpha_1 \prod_{i'=2}^i (1 - \exp(-\alpha_{i'})).$$

Remember that $\alpha_i = \lambda_1 |D_i|$ for $i \in \{1, \dots, k\}$ and $\alpha_{k+1} = \lambda_\rho |D_{k+1}|$. By lemma 2.13, we have the following limits:

$$\begin{aligned} \lim_{d \rightarrow +\infty} \frac{1}{d} \ln \alpha_1 &= \ln \frac{\kappa(1+\rho)}{2} < 0, \\ \lim_{d \rightarrow +\infty} \frac{1}{d} \ln \alpha_i &= \ln(\kappa \sqrt{1 - a_i^2}) < 0 \text{ for } 2 \leq i \leq k, \\ \lim_{d \rightarrow +\infty} \frac{1}{d} \ln \alpha_{k+1} &= \ln(\kappa \sqrt{1 - a_{k+1}^2} \frac{1+\rho}{2\rho}) < 0. \end{aligned}$$

To see the signs of the limits, note that $\kappa < 1$, that $\frac{1+\rho}{2\rho} < 1$ and that (24) implies that

$$\kappa > \kappa^{k+1} > \frac{4\rho}{(1+\rho)^2 \sqrt{1 - a_2^2} \dots \sqrt{1 - a_{k+1}^2}} > \frac{2}{1+\rho}.$$

Consequently, we first see that

$$\begin{aligned} \lim_{d \rightarrow +\infty} \frac{1}{d} \ln(\eta_{k+1}) &= \lim_{d \rightarrow +\infty} \frac{1}{d} \ln(\alpha_1 \dots \alpha_{k+1}) \\ &= \ln \left(\kappa^{k+1} \frac{(1+\rho)^2}{4\rho} \sqrt{\prod_{2 \leq j \leq k+1} (1-a_j^2)} \right) > 0 \quad \text{with (24);} \end{aligned}$$

$$(31) \quad \text{therefore,} \quad \lim_{d \rightarrow +\infty} P(\#J_{k+1} = 0) = 0.$$

Similarly, for $2 \leq i \leq k+1$, we have

$$\lim_{d \rightarrow +\infty} \frac{1}{d} \ln(\eta_{i-1}) = \ln \left(\kappa^{i-1} \frac{1+\rho}{2} \sqrt{\prod_{2 \leq i' \leq i-1} (1-a_{i'}^2)} \right).$$

Lemmas 2.11 and 2.13 ensure that :

$$\lim_{d \rightarrow +\infty} \frac{1}{d} \ln \left(\frac{|D_i''|}{|C_i''|} \right) = \ln \left(\frac{r_i \sqrt{1-a_i^2}}{d_i} \right).$$

Thus, for $2 \leq i \leq k+1$, we have :

$$\limsup_{d \rightarrow +\infty} \frac{1}{d} \ln \left(E(\#J_{i-1}) \frac{|D_i''|}{|C_i''|} \right) \leq \ln \left(\frac{r_i (1+\rho) \kappa^{i-1}}{2d_i} \sqrt{\prod_{2 \leq i' \leq i} (1-a_{i'}^2)} \right).$$

Now,

$$(32) \quad \text{for } 2 \leq i \leq k, \quad \limsup_{d \rightarrow +\infty} \frac{1}{d} \ln \left(E(\#J_{i-1}) \frac{|D_i''|}{|C_i''|} \right) \leq \ln \left(\frac{1+\rho}{d_i} \right) < 0,$$

$$(33) \quad \limsup_{d \rightarrow +\infty} \frac{1}{d} \ln \left(E(\#J_k) \frac{|D_{k+1}''|}{|C_{k+1}''|} \right) \leq \ln \left(\frac{(1+\rho)^2 \kappa^k}{2d_{k+1}} \sqrt{\prod_{2 \leq i' \leq k+1} (1-a_{i'}^2)} \right) < 0$$

with (24). To end the proof, we put estimates (31), (32) and (33) in (30). \square

2.4.2. Several steps in the 2-dimensional oriented percolation model. We prove here Proposition 2.8 by building the supercritical 2-dimensional oriented percolation process embedded in the two-type Boolean Model.

Proof of Proposition 2.8. We first define an oriented graph in the following manner: the set of sites is

$$S = \{(a, n) \in \mathbb{Z} \times \mathbb{N} : |a| \leq n, a+n \text{ is even}\};$$

from any point $(a, n) \in S$, we put an oriented edge from (a, n) to $(a+1, n+1)$, and an oriented edge from (a, n) to $(a-1, n+1)$. We denote by $\vec{p}_c(2) \in (0, 1)$ the critical parameter for Bernoulli percolation on this oriented graph – see Durrett 84 [5].

For any $(a, n) \in S$, we define the following subsets of \mathbb{R}^d

$$\begin{aligned} W_{a,n} &= d^{-1/2} (]a-1, a+1[\times]n-1, n[\times \mathbb{R}^{d-2}), \\ W_{a,n}^- &= d^{-1/2} (]a-1, a[\times]n, n+1[\times \mathbb{R}^{d-2}), \\ W_{a,n}^+ &= d^{-1/2} (]a, a+1[\times]n, n+1[\times \mathbb{R}^{d-2}). \end{aligned}$$

Note that the $(W_{a,n})_{(a,n) \in S}$ are disjoint and that $W_{a,n}^+ \cup W_{a+2,n}^- \subset W_{a+1,n+1}$.

We now fix $k \geq 1$ and $\kappa \in (\kappa_\rho^c(k), 1)$, and for $x_0 \in W_{a,n}$, we introduce the events :

$$\begin{aligned} \mathcal{G}_{a,n}^+(x_0) &= \left\{ \begin{array}{l} \text{There exist distinct } x_1, \dots, x_k \in \chi_1 \cap W_{a,n}^+ \text{ and } x_{k+1} \in \chi_\rho \cap W_{a,n}^+ \\ \text{such that } x_0, x_1, \dots, x_{k+1} \text{ is a path} \end{array} \right\}, \\ \mathcal{G}_{a,n}^-(x_0) &= \left\{ \begin{array}{l} \text{There exist distinct } x_1, \dots, x_k \in \chi_1 \cap W_{a,n}^- \text{ and } x_{k+1} \in \chi_\rho \cap W_{a,n}^- \\ \text{such that } x_0, x_1, \dots, x_{k+1} \text{ is a path} \end{array} \right\}. \end{aligned}$$

Note that $\mathcal{G}_{0,0}^+(x)$ is exactly the event $\mathcal{G}^+(x)$ introduced in (23), and that the other events are obtained from this one by symmetry and/or translation.

Next we choose $p \in (\vec{p}_c(2), 1)$. With Proposition 2.9, and by translation and symmetry invariance, we know that for every large enough dimension d , for every $(a, n) \in S$, for every $x \in W_{a,n}$:

$$(34) \quad P(\mathcal{G}_{a,n}^\pm(x)) \geq p.$$

We fix then a dimension d large enough to satisfy (34). We can now construct the random states, open or closed, of the edges of our oriented graph. We denote by ∞ a virtual site.

Definition of the site on level 0. Almost surely, $\chi_\rho \cap W_{0,0} \neq \emptyset$. We take then some $x(0, 0) \in \chi_\rho \cap W_{0,0}$.

Definition of the edges between levels n and $n + 1$. Fix $n \geq 0$ and assume we built a site $x(a, n) \in W_{a,n} \cup \{\infty\}$ for every $(a, n) \in S$. Consider $(a, n) \in S$:

- If $x(a, n) = \infty$: we decide that each of the two edges starting from (a, n) is open with probability p and closed with probability $1 - p$, independently of everything else; we set $z^-(a, n) = z^+(a, n) = \infty$.
- Otherwise :
 - Edge to the left-hand side :
 - if the event $\mathcal{G}_{a,n}^-(x(a, n))$ occurs : we take for $z^-(a, n)$ some point $x_{k+1} \in W_{a,n}^- \subset W_{a-1, n+1}$ given by the occurrence of the event, and we open the edge from (a, n) to $(a - 1, n + 1)$;
 - otherwise : we set $z^-(a, n) = \infty$ and we close the edge from (a, n) to $(a - 1, n + 1)$.
 - Edge to the right-hand side :
 - if the event $\mathcal{G}_{a,n}^+(x(a, n))$ occurs : we take for $z^+(a, n)$ some point $x_{k+1} \in W_{a,n}^+ \subset W_{a+1, n+1}$ given by the occurrence of the event, and we open the edge from (a, n) to $(a + 1, n + 1)$;
 - otherwise : we set $z^+(a, n) = \infty$ and we close the edge from (a, n) to $(a + 1, n + 1)$.

For (a, n) outside S , we set $z^\pm(a, n) = \infty$.

Definition of the sites at level $n + 1$. Fix $n \geq 0$ and assume we determined the state of every edge between levels n and $n + 1$. Consider $(a, n + 1) \in S$:

- If $z^+(a - 1, n) \neq \infty$: set $x(a, n + 1) = z^+(a - 1, n) \in W_{a, n+1}$.
- Otherwise :
 - if $z^-(a + 1, n) \neq \infty$: set $x(a, n + 1) = z^-(a + 1, n) \in W_{a, n+1}$,
 - otherwise : set $x(a, n + 1) = \infty$.

Assume that there exists an open path of length n starting from the origin in this oriented percolation : we can check that the leftmost open path of length n starting from the origin gives a path in the two-type Boolean model with n alternating sequences of k balls with radius 1 and one ball with radius ρ . Thus, percolation in this oriented

percolation model implies percolation by k -alternation in the two-type Boolean model. Let us check that percolation occurs indeed with positive probability.

For every n , denote by \mathcal{F}_n the σ -field generated by the restrictions of the Poisson point processes χ_1 and χ_ρ to the set

$$d^{-1/2} (\mathbb{R} \times (-\infty, n) \times \mathbb{R}^{d-2}).$$

By definition of the events \mathcal{G} – remember that the $(W_{a,n})_{(a,n) \in S}$ are disjoint – and by (34), the states of the different edges between levels n and $n+1$ are independent conditionally to \mathcal{F}_n . Moreover, conditionally to \mathcal{F}_n , each edge between levels n and $n+1$ has a probability at least p to be open. Therefore, the oriented percolation model we built stochastically dominates Bernoulli oriented percolation with parameter p . As $p > \vec{p}_c(\mathbb{Z}^2)$, with positive probability, there exists an infinite open path in the oriented percolation model we built; this ends the proof of Proposition 2.8. \square

2.5. Proof of Theorem 1.4. We first prove how Propositions 2.2 and 2.8 give (11) when $a = 1$, $b > 1$ and $\alpha = \beta = 1$, and then we see how we can deduce the general case by scaling and coupling.

When $a = 1$, $b > 1$ and $\alpha = \beta = 1$. Set $\rho = b$. In this case, $\nu = \delta_1 + \delta_\rho$, so $\nu_d = \delta_1 + \frac{1}{\rho^d} \delta_\rho$.

Note then that the two-type Boolean model Σ introduced in Subsection 2.1 and whose intensities depend on $\kappa \in (0, 1)$ coincides with the Boolean model directed by the measure

$$\frac{\kappa^d}{\nu_d 2^d} \nu_d$$

as defined in the introduction.

If $\kappa < \kappa_\rho^c$ then, by Proposition 2.2, there is no percolation for d large enough. Therefore, for any such κ and for any large enough d we have:

$$\lambda_d^c(\nu_d) \geq \frac{\kappa^d}{\nu_d 2^d} \quad \text{and then} \quad \tilde{\lambda}_d^c(\nu_d) = \lambda_d^c(\nu_d) \nu_d 2^d \int r^d \nu_d(dr) \geq 2\kappa^d.$$

Letting d goes to $+\infty$ and then κ goes to κ_ρ^c , we then obtain

$$(35) \quad \liminf_{d \rightarrow +\infty} \frac{1}{d} \ln(\lambda_d^c(\nu_d)) \geq \ln(\kappa_\rho^c).$$

As $\kappa_\rho^c < 1$ by Lemma 2.1, choose now κ such that $\kappa_\rho^c < \kappa < 1$. Then, there exists $k \geq 1$ such that $\kappa_\rho^c(k) < \kappa$. Therefore, by Proposition 2.8, there is percolation for d large enough in Σ ; by coupling, this remains true for any $\kappa > \kappa_\rho^c$. Therefore, for any $\kappa > \kappa_\rho^c$ and for any large enough d we have, as before:

$$\lambda_d^c(\nu_d) \leq \frac{\kappa^d}{\nu_d 2^d} \quad \text{and then} \quad \tilde{\lambda}_d^c(\nu_d) \leq 2\kappa^d.$$

Letting d goes to $+\infty$ and then κ goes to κ_ρ^c , we then obtain

$$(36) \quad \limsup_{d \rightarrow +\infty} \frac{1}{d} \ln(\lambda_d^c(\nu_d)) \leq \ln(\kappa_\rho^c).$$

Bringing (35) and (36) together, we get (11) when $a = 1$, $b = \rho > 1$ and $\alpha = \beta = 1$.

When $b > a > 0$ and $\alpha = \beta = 1$. Set $\rho = b/a$. Here, $\nu = \delta_a + \delta_b$; set $\mu = \delta_1 + \delta_\rho$. With the notation of the introduction,

$$\nu_d = \frac{1}{a^d}(\delta_a + \frac{1}{\rho^d}\delta_b) = \frac{1}{a^d}H^{1/a} \cdot (\delta_1 + \frac{1}{\rho^d}\delta_\rho) = \frac{1}{a^d}H^{1/a}\mu_d.$$

By the scaling relations (2) and (3), we obtain

$$\tilde{\lambda}_d^c(\nu_d) = \tilde{\lambda}_d^c(\mu_d).$$

The result when $b > a > 0$ and $\alpha = \beta = 1$ follows then from the previous case.

When $b > a > 0$ and $\alpha, \beta > 0$. Here $\nu = \alpha\delta_a + \beta\delta_b$. Set $\mu = \delta_a + \delta_b$, $m = \min(\alpha, \beta)$ and $M = \max(\alpha, \beta)$. Then $m\mu_d \leq \nu_d \leq M\mu_d$ and so

$$\begin{aligned} m \int r^d d\mu_d(r) &\leq \int r^d d\nu_d(r) \leq M \int r^d d\mu_d(r), \\ \frac{1}{M} \lambda_d^c(\mu_d) = \lambda_d^c(M\mu_d) &\leq \lambda_d^c(\nu_d) \leq \lambda_d^c(m\mu_d) = \frac{1}{m} \lambda_d^c(\mu_d). \end{aligned}$$

The two previous inequalities give:

$$\frac{m}{M} \tilde{\lambda}_d^c(\mu_d) \leq \tilde{\lambda}_d^c(\nu_d) \leq \frac{M}{m} \tilde{\lambda}_d^c(\mu_d),$$

and the theorem follows from the previous case. \square

3. PROOF OF THEOREM 1.2

Theorem 1.2 follows from Theorem 1.4 by coupling and scaling. By assumption, μ is a measure on $(0, +\infty)$ whose support is not a singleton. We can therefore choose $b' > a' > 0$ in the support, set $\rho = b'/a'$ and then take a small enough $\varepsilon > 0$ such that

$$\begin{aligned} a'(1 + \varepsilon) &< b'(1 - \varepsilon), & \mu([a'(1 - \varepsilon), a'(1 + \varepsilon)]) &> 0, \\ \mu([b'(1 - \varepsilon), b'(1 + \varepsilon)]) &> 0, & (1 + \varepsilon)(1 - \varepsilon)^{-1} \kappa_\rho^c &< 1. \end{aligned}$$

Set $a = a'(1 - \varepsilon)$, $b = b'(1 - \varepsilon)$ and $\tau = (1 + \varepsilon)(1 - \varepsilon)^{-1} > 1$. We have

$$a\tau < b, \mu([a, a\tau]) > 0, \mu([b, b\tau]) > 0 \text{ and } \tau \kappa_\rho^c < 1.$$

Set $\nu = \mu([a, a\tau])\delta_a + \mu([b, b\tau])\delta_b$ and $S = [a, a\tau] \cup [b, b\tau]$. For all $d \geq 1$ we have

$$\tau^{-d} \nu_d(\{a\}) = \mu([a, a\tau])(a\tau)^{-d} \leq \int_{[a, a\tau]} r^{-d} \mu(dr) = \mu_d([a, a\tau])$$

and, similarly, $\tau^{-d} \nu_d(\{b\}) \leq \mu_d([b, b\tau])$. By coupling, this implies that $\lambda_d^c(1_S \mu_d) \leq \lambda_d^c(\tau^{-d} \nu_d)$, and then that

$$\lambda_d^c(\mu_d) \leq \lambda_d^c(1_S \mu_d) \leq \lambda_d^c(\tau^{-d} \nu_d) = \tau^d \lambda_d^c(\nu_d).$$

But $\tilde{\lambda}_d^c(\mu_d) = \lambda_d^c(\mu_d) 2^d \nu_d \mu((0, +\infty))$ and, similarly, $\tilde{\lambda}_d^c(\nu_d) = \lambda_d^c(\nu) 2^d \nu_d \nu((0, +\infty))$, which leads to

$$\tilde{\lambda}_d^c(\mu_d) \leq \tau^d \frac{\mu((0, +\infty))}{\nu((0, +\infty))} \tilde{\lambda}_d^c(\nu_d).$$

But by Theorem 1.4 we have

$$\lim_{d \rightarrow +\infty} \frac{1}{d} \ln \left(\tilde{\lambda}_d^c(\nu_d) \right) = \ln(\kappa_\rho^c), \text{ and then } \limsup_{d \rightarrow +\infty} \frac{1}{d} \ln \left(\tilde{\lambda}_d^c(\mu_d) \right) \leq \ln(\tau \kappa_\rho^c) < 0,$$

which ends the proof. \square

Note that as a byproduct of the proof, we obtain the following upper bound :

$$\limsup_{d \rightarrow +\infty} \frac{1}{d} \ln \left(\tilde{\lambda}_d^c(\mu_d) \right) \leq \inf_{0 < a < b < +\infty, a, b \in \text{Supp}(\mu)} \ln(\kappa_{b/a}^c).$$

An interesting but certainly tough result would be to obtain a real limit, and there is no reason why this limit should coincide with this upper bound.

REFERENCES

- [1] Ajit Balram and Deepak Dhar. Scaling relation for determining the critical threshold for continuum percolation of overlapping discs of two sizes. *Pramana*, 74:109–114, 2010. 10.1007/s12043-010-0012-0.
- [2] R. Consiglio, D. R. Baker, G. Paul, and H. E. Stanley. Continuum percolation thresholds for mixtures of spheres of different sizes. *Physica A: Statistical Mechanics and its Applications*, 319:49 – 55, 2003.
- [3] Deepak Dhar. On the critical density for continuum percolation of spheres of variable radii. *Physica A: Statistical and Theoretical Physics*, 242(3-4):341 – 346, 1997.
- [4] Deepak Dhar and Mohan K. Phani. Continuum percolation with discs having a distribution of radii. *J. Phys. A*, 17:L645–L649, 1984.
- [5] Richard Durrett. Oriented percolation in two dimensions. *Ann. Probab.*, 12(4):999–1040, 1984.
- [6] Jean-Baptiste Gou  r  . Subcritical regimes in the Poisson Boolean model of continuum percolation. *Ann. Probab.*, 36(4):1209–1220, 2008.
- [7] Jean-Baptiste Gou  r  . Subcritical regimes in some models of continuum percolation. *Ann. Appl. Probab.*, 19(4):1292–1318, 2009.
- [8] Jean-Baptiste Gou  r  . Percolation in a multiscale boolean model. Preprint, 2010.
- [9] J  nos Kert  sz and Tam  s Vicsek. Monte carlo renormalization group study of the percolation problem of discs with a distribution of radii. *Zeitschrift f  r Physik B Condensed Matter*, 45:345–350, 1982. 10.1007/BF01321871.
- [10] Harry Kesten. Asymptotics in high dimensions for percolation. In *Disorder in physical systems*, Oxford Sci. Publ., pages 219–240. Oxford Univ. Press, New York, 1990.
- [11] Ronald Meester and Rahul Roy. *Continuum percolation*, volume 119 of *Cambridge Tracts in Mathematics*. Cambridge University Press, Cambridge, 1996.
- [12] Ronald Meester, Rahul Roy, and Anish Sarkar. Nonuniversality and continuity of the critical covered volume fraction in continuum percolation. *J. Statist. Phys.*, 75(1-2):123–134, 1994.
- [13] M. V. Menshikov, S. Yu. Popov, and M. Vachkovskaia. On the connectivity properties of the complementary set in fractal percolation models. *Probab. Theory Related Fields*, 119(2):176–186, 2001.
- [14] Mathew D. Penrose. Continuum percolation and Euclidean minimal spanning trees in high dimensions. *Ann. Appl. Probab.*, 6(2):528–544, 1996.
- [15] John A. Quintanilla and Robert M. Ziff. Asymmetry in the percolation thresholds of fully penetrable disks with two different radii. *Phys. Rev. E*, 76(5):051115, Nov 2007.

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