

Comment on The plasmasheath transition in low temperature plasmas: on the existence of a collisionally modified Bohm criterion

R N Franklin

► To cite this version:

R N Franklin. Comment on The plasma sheath transition in low temperature plasmas: on the existence of a collisionally modified Bohm criterion. Journal of Physics D: Applied Physics, 2011, 44 (15), pp.158001. 10.1088/0022-3727/44/15/158001. hal-00613265

HAL Id: hal-00613265 https://hal.science/hal-00613265

Submitted on 4 Aug 2011

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés. Comment on Brinkmann (J. Phys D : Appl. Phys. 44 (2011) 042002)

R. N. Franklin, Oxford Research Unit, Open University, Oxford, OX1 5HR U.K.

In a recent Rapid Communication Brinkmann [1] claims to have discovered a collisionally-modified Bohm Criterion. He introduces a new local variable length λ defined by $\lambda = n_i dx/dn_i$, where n_i is the ion density, this has the effect of locally compressing the natural length scale where n_i is rapidly varying and has maximum effect near the point of inflexion in $n_i(x)$. This point he calls the Bohm point and the ion speed there is given the name the collisionally-modified Bohm Criterion – let us call it c_s^* , with c_s being the ion sound speed. It can readily be shown that $d^2n_i/dx^2 = (1 - d\lambda/dx).n_i/\lambda^2$. There is no doubt that the spatial distributions of both ions and electrons have a point of inflexion somewhere between the centre of the plasma and the wall in a real active plasma, but it is difficult to see how such a point can be given the significance suggested when its location is a function of the plasma size L, the ion mean free path λ_i , and the Debye length λ_D based on the plasma density at the centre n_0 . c_s^* is likewise a function of these three variables.

For specific gas and a finite plasma one needs to include the ionization rate by electron impact Z, and this varies with the nature of the gas, and not just the ion mass. Thus for a given situation it is possible to write the plasma balance equation as $-(Z(Z+\nu_i)^{1/2}L/c_s = K(\nu_i/Z, \lambda_D/L))$ where K is a number of the order of 1, and depends on the geometry, albeit plane or cylindrical, ν_i is the ion collision frequency for momentum transfer. While the ion sound speed enters here unmodified, one could devise a modified expression for $c_s^* = c_s N(\nu_i/Z, \lambda_D/L)$, N being a number between 1 and 0 with limiting values N = 1 for $\nu_i/Z \rightarrow 0$, and N $\rightarrow 0$ for $\nu_i/Z \gg 1$.. I cannot see how such a two-dimensional plot, even when for $\lambda_D/L \rightarrow 0$, N has known values given in my book [2], gives additional physical insight.

All of this was fully explored in my Topical Review [3], which was not acknowledged by Brinkmann.

References

Brinkmann R.P. 2011 J. Phys D : Appl. Phys. 44 042002.
Franklin R. N 1976. Plasma Phenomena in Gas Discharges (Oxford : Oxford University Press).
Franklin 2003 J. Phys. D : Appl. Phys. 36 R309-R320.