



HAL
open science

Kac's Program in Kinetic Theory

Stéphane Mischler, Clément Mouhot

► **To cite this version:**

| Stéphane Mischler, Clément Mouhot. Kac's Program in Kinetic Theory. 2011. hal-00608971v1

HAL Id: hal-00608971

<https://hal.science/hal-00608971v1>

Preprint submitted on 16 Jul 2011 (v1), last revised 22 Jul 2012 (v2)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

KAC'S PROGRAM IN KINETIC THEORY
(VERSION OF July 16, 2011)

S. MISCHLER AND C. MOUHOT

ABSTRACT. This paper is devoted to the study of propagation of chaos and mean-field limit for systems of indistinguishable particles undergoing collision processes, as formulated by M. Kac [42] for a simplified model and extended by H. P. McKean [53] to the Boltzmann equation. We prove quantitative and uniform in time estimates measuring the distance between the many-particle system and the limit system. These estimates imply in particular the propagation of chaos for marginals in weak measure distances but are more general: they hold for non-chaotic initial data and control the complete many-particle distribution. We also prove the propagation of *entropic chaos*, as defined in [12], answering a question of Kac about the microscopic derivation of the H -theorem. We finally prove estimates of relaxation to equilibrium (in Wasserstein distance and relative entropy) *independent of the number of particles*. Our results cover the two main Boltzmann physical collision processes with unbounded collision rates: hard spheres and *true* Maxwell molecules interactions. Starting from an inspirational paper of A. Grünbaum [36] we develop a new method which reduces the question of propagation of chaos to the one of proving a purely functional estimate on some generator operators (*consistency estimates*) together with fine differentiability estimates on the flow of the limit non-linear equation (*stability estimates*). These results provide the first answer to the question raised by Kac of relating the long-time behavior of a collisional particle system with the one of its mean-field limit, however using dissipativity at the level of the mean-field limit instead of using it at the level of the many-particle Markov process.

Keywords: Kac's program; kinetic theory; master equation; mean-field limit; quantitative; uniform in time; jump process; collision process; Boltzmann equation; Maxwell molecules; non cutoff; hard spheres.

AMS Subject Classification: 82C40 Kinetic theory of gases, 76P05 Rarefied gas flows, Boltzmann equation, 54C70 Entropy, 60J75 Jump processes.

CONTENTS

1. Introduction and main results	3
2. The abstract setting	11
3. The abstract theorem	26
4. The N -particle approximation at initial time	34
5. True Maxwell molecules	45
6. Hard spheres	64
7. H -theorem and entropic chaos	81
8. The BBGKY approach revisited	89
References	96

L. Boltzmann is best known for the equation bearing his name in kinetic theory [7, 8]. Inspired by Maxwell’s discovery [51] of (what is now called) the Boltzmann equation and its “maxwellian” (gaussian) equilibrium, Boltzmann [7] discovered the “ H -theorem” (increase of the entropy) for this equation which explains how the solutions should be driven towards the equilibrium of Maxwell. In the same work he also proposed the deep idea of “stosszahlansatz” (molecular chaos) for explaining how the irreversible Boltzmann equation could emerge from the Newton laws on the dynamics of particles. Giving a precise mathematical meaning to this notion and proving this limit remains a tremendous open problem to this date (the best and astonishing result so far [45] is only valid for very short times).

M. Kac proposed in 1956 [42] the simpler and seemingly more tractable question of deriving the *spatially homogeneous* Boltzmann equation from a many-particle jump process, and he introduced a rigorous notion of molecular chaos¹ in this context. The “chaoticity” of the many-particle equilibrium towards the maxwellian distribution, i.e. the fact that the first marginals of the uniform measure on the sphere $\mathbb{S}^{N-1}(\sqrt{N})$ converges to a gaussian as N goes to infinity, has been known for a long time, at least since Maxwell.² However in [42] Kac proposed the first proof of the propagation of chaos along time for a simplified collision process for which series expansions of the solution are available, and he showed how the many-particle limit rigorously follows from this property of propagation of chaos. This proof was later extended to a more realistic collision model, the so-called cutoff Maxwell molecules, by McKean [53].

Since in this setting both the many-particle system and the limit equation are dissipative, Kac also raised the natural question of relating their asymptotic behaviors. In his mind this program was to be achieved by understanding dissipativity at the level of the linear many-particle jump process and he insisted on the importance of estimating the rate of relaxation of the many-particle process. This has motivated beautiful works on this “Kac spectral gap problem” [40, 50, 13, 15, 11], but so far this strategy has proved unsuccessful in relating the many-particle process asymptotic behavior and that of the limit equation, see the interesting discussion in [12]. At the time of Kac the study of nonlinear partial differential equations was rather young and it was plausible that the study of a linear many-dimension Markov process would be easier. However the mathematical development somehow followed the inverse direction and the theory of existence, uniqueness and relaxation to equilibrium for the spatially homogeneous Boltzmann is now well-developed (see the many references along this paper).

This paper is an attempt to develop a quantitative theory of mean-field limit which *strongly relies on detailed knowledge of the limit nonlinear equation, rather than on detailed properties of the many-particle Markov process*. As the main outcome of this theory we prove uniform in time quantitative propagation of chaos as well as propagation of entropic chaos, and we prove relaxation rates *independent of the number of particles* (measured in Wasserstein distance and relative entropy). All this is done for the two important realistic and achetypal models of collision, namely hard spheres and true (without cutoff) Maxwell molecules. This provides a first complete answer to the question raised by Kac, however our answer is an “inverse” answer in the sense that our methodology is “top-down” from the limit equation to the many-particle system rather than “bottom-up” as was expecting Kac.

¹Kac in fact called this notion “Boltzmann’s property” in [42] as a clear tribute to the fundamental intuition of Boltzmann.

²We refer to [23] for a bibliographic discussion, see also [12] where [54] is quoted as the first paper proving this result.

A first version of this work was posted on arXiv as “Quantitative uniform in time chaos propagation for Boltzmann collision processes”. This new expanded version includes new results about entropic chaos and relaxation times.

Acknowledgments. The authors would like to thank the mathematics department of Chalmers University for the invitation in november 2008, where the abstract method was devised and the related joint work [58] with Bernt Wennberg was initiated. They also thank Ismaël Bailleul, François Bolley, Anne Boutet de Monvel, José Alfredo Cañizo, Eric Carlen, Nicolas Fournier, Arnaud Guillin, Maxime Hauray, Joël Lebowitz, Pierre-Louis Lions, Richard Nickl, Mario Pulvirenti, Judith Rousseau and Cédric Villani for fruitful comments and discussions. Finally they wish to mention the inspirative courses of Pierre-Louis Lions at Collège de France on “Mean-Field Games” in 2007-2008 and 2008-2009.

1. INTRODUCTION AND MAIN RESULTS

1.1. The Boltzmann equation. The Boltzmann equation (Cf. [19] and [20]) describes the behavior of a dilute gas when the only interactions taken into account are binary collisions. It writes

$$\frac{\partial f}{\partial t} + v \cdot \nabla_x f = Q(f, f)$$

where $Q = Q(f, f)$ is the bilinear *Boltzmann collision operator* acting only on the velocity variable.

In the case when the distribution function is assumed to be independent on the position x , we obtain the so-called *spatially homogeneous Boltzmann equation*, which reads

$$(1.1) \quad \frac{\partial f}{\partial t}(t, v) = Q(f, f)(t, v), \quad v \in \mathbb{R}^d, \quad t \geq 0,$$

where $d \geq 2$ is the dimension.

Let us now focus on the collision operator Q . It is defined by the bilinear symmetrized form

$$(1.2) \quad Q(g, f)(v) = \frac{1}{2} \int_{\mathbb{R}^d \times \mathbb{S}^{d-1}} B(|v - v_*|, \cos \theta) (g'_* f' + g' f'_* - g_* f - g f_*) dv_* d\sigma,$$

where we have used the shorthands $f = f(v)$, $f' = f(v')$, $g_* = g(v_*)$ and $g'_* = g(v'_*)$. Moreover, v' and v'_* are parametrized by

$$(1.3) \quad v' = \frac{v + v_*}{2} + \frac{|v - v_*|}{2} \sigma, \quad v'_* = \frac{v + v_*}{2} - \frac{|v - v_*|}{2} \sigma, \quad \sigma \in \mathbb{S}^{d-1}.$$

Finally, $\theta \in [0, \pi]$ is the deviation angle between $v' - v'_*$ and $v - v_*$ defined by $\cos \theta = \sigma \cdot \hat{u}$, $u = v - v_*$, $\hat{u} = u/|u|$, and B is the Boltzmann collision kernel determined by physics (related to the cross-section $\Sigma(v - v_*, \sigma)$ by the formula $B = |v - v_*| \Sigma$).

Boltzmann's collision operator has the fundamental properties of conserving mass, momentum and energy

$$\int_{\mathbb{R}^d} Q(f, f) \phi(v) dv = 0, \quad \phi(v) = 1, v, |v|^2,$$

and satisfying the so-called Boltzmann's H theorem which writes (at the formal level)

$$-\frac{d}{dt} \int_{\mathbb{R}^d} f \log f dv = - \int_{\mathbb{R}^d} Q(f, f) \log(f) dv \geq 0.$$

We shall consider collision kernels

$$B = \Gamma(|v - v_*|) b(\cos \theta)$$

with Γ, b nonnegative functions. In dimension $d = 3$, here is a short classification of possible collision kernels, together with some important examples that we shall consider in this paper.

- (1) Short (finite) range interaction are usually modeled by the **hard spheres collision kernel**

$$(1.4) \quad \text{(HS)} \quad B(|v - v_*|, \cos \theta) = \text{cst } |v - v_*|.$$

- (2) Long-range interactions are usually modeled by collision kernels deriving from interaction potentials

$$V(r) = \text{cst } r^{-s}, \quad s > 2.$$

They satisfy the formula

$$\Gamma(z) = |z|^\gamma \text{ with } \gamma = (s - 4)/s$$

and

$$b(\cos \theta) \sim_{\theta \sim 0} C_b \theta^{-2-\nu} \text{ with } \nu = 2/s$$

(b is L^1 apart from $\theta \sim 0$). Such formulas (and others) can be found in [19].

This general class of collision kernels includes in particular the **true Maxwell molecules collision kernel** when $\gamma = 0$ and $\nu = 1/2$:

$$(1.5) \quad \text{(tMM)} \quad B(|v - v_*|, \cos \theta) = B(\cos \theta) \sim_{\theta \sim 0} C_b \theta^{-5/2}.$$

It also includes the so-called **Grad's cutoff Maxwell molecules** when the singularity in the θ variable is removed. Without much restriction we shall consider as a typical such model

$$(1.6) \quad \text{(GMM)} \quad B(|v - v_*|, \cos \theta) = 1.$$

1.2. Deriving the Boltzmann equation from many-particle systems. The question of deriving the Boltzmann equation from particle systems (interacting *via* Newton's laws) is a famous problem. It is related to the so-called 6-th Hilbert problem proposed by Hilbert at the International Congress of Mathematics at Paris in 1900: axiomatize mechanics by “*developing mathematically the limiting processes [...] which lead from the atomistic view to the laws of motion of continua*”.

At least at the formal level, the correct limiting procedure has been identified by Grad [34] in the late forties and a clear mathematical formulation of the open problem was proposed in [18] in the early seventies. It is now called the *Boltzmann-Grad* or *low density* limit. However the original question of Hilbert remains largely open, in spite of a striking breakthrough due to Lanford [45], who proved the limit for short times (see also Illner and Pulvirenti [39] for a close-to-vacuum result). The tremendous difficulty underlying this limit is the *irreversibility* of the Boltzmann equation, whereas the particle system interacting *via* Newton's laws is a reversible Hamiltonian system.

In 1954-1955, Kac [42] proposed a simpler and more tractable problem: start from the Markov process corresponding to collisions only, and try to prove the limit towards the *spatially homogeneous* Boltzmann equation. *Kac's jump process* runs as follows: consider N particles with velocities $v_1, \dots, v_N \in \mathbb{R}^d$. Compute random times for each pair of particles (v_i, v_j) following an exponential law with parameter $\Gamma(|v_i - v_j|)$, take the smallest, and perform a collision $(v_i, v_j) \rightarrow (v_i^*, v_j^*)$ given by a random choice of a direction parameter whose rule is related to $b(\cos \theta)$, then recommence. This process can be considered on \mathbb{R}^{dN} ; however it leaves invariant some submanifolds of \mathbb{R}^{dN} (depending on the number of conserved quantities during collision) and can be restricted to them. For instance in the original simplified model of Kac $d = 1$ (scalar velocities) and the process can be restricted to $\mathbb{S}^{N-1}(\sqrt{\mathcal{E}N})$ the sphere with radius $\sqrt{\mathcal{E}N}$, where \mathcal{E} is the energy. In the more realistic

hard spheres of Maxwell molecules models, $d = 3$ and the process can be restricted to the sphere

$$\mathcal{S}^N := \mathbb{S}^{dN-1} \left(\sqrt{N\mathcal{E}} \right) \cap \{v_1 + \cdots + v_N = 0\}.$$

Kac formulated the notion of *propagation of chaos* that we shall now explain. Consider a sequence $(f^N)_{N \geq 1}$ of probabilities on E^N , where E is some given Polish space (think to $E = \mathbb{R}^d$ for most applications): the sequence is said *f-chaotic* if

$$f^N \sim f^{\otimes N} \quad \text{when } N \rightarrow \infty$$

for some given one-particle probability f on E . The meaning of this convergence is the following: convergence in the weak measure topology for any marginal depending on a finite number of variables. This is a *low correlation* assumption.

It was clear since Boltzmann that in the case when the joint probability density f^N of the N -particle system is tensorized *during some time interval* into N copies $f^{\otimes N}$ of a 1-particle probability density, then the latter would satisfy the limit nonlinear Boltzmann equation during this time interval. But Kac made the key remark that although in general interactions between a finite number of particles prevents any possibility of propagation of the “tensorization” property, the weaker property of chaoticity can be propagated (hopefully!) in the correct scaling limit.

The framework set by Kac is our starting point. Let us emphasize that the limit performed in this setting is different from the Boltzmann-Grad limit. The Kac’s limit is in fact a *mean-field limit*. This limiting procedure is most well-known for deriving Vlasov-like equations. In a companion paper [58] we develop systematically our new functional approach for Vlasov equations, McKean-Vlasov equations, and granular gases driven by a thermal bath.

1.3. The notion of chaos and how to measure it. Our goal in this paper is to set up a general robust method for proving the propagation of chaos with *quantitative rate* in terms of the number of particles N and of the final time of observation T .

Let us explain briefly what it means. The original formulation of Kac [42] of chaoticity is: a sequence $f^N \in P_{\text{sym}}(E^N)$ of symmetric probabilities on E^N is *f-chaotic*, for a given probability $f \in P(E)$, if for any $\ell \in \mathbb{N}^*$ and any $\varphi \in C_b(E)^{\otimes \ell}$ there holds

$$\lim_{N \rightarrow \infty} \left\langle f^N, \varphi \otimes \mathbf{1}^{N-\ell} \right\rangle = \left\langle f^{\otimes \ell}, \varphi \right\rangle$$

which amounts to the weak convergence of any marginals. This can be expressed for instance in Wasserstein distance:

$$\lim_{N \rightarrow \infty} W_1 \left(\Pi_\ell (f^N), f^{\otimes \ell} \right) = 0$$

where Π_ℓ denotes the marginal on the ℓ first variables. This is **Kac’s definition of chaos** that we shall call **finite-dimensional chaos**.

We shall deal in this paper with **quantitative chaos**, in the sense that we measure precisely the rate of convergence in the above limit. Namely, we say that f^N is *f-chaotic* with rate $\varepsilon(N)$, where $\varepsilon(N) \rightarrow 0$ when $N \rightarrow \infty$ (typically $\varepsilon(N) = N^{-r}$, $r > 0$ or $\varepsilon(N) = (\log N)^{-r}$, $r > 0$), if for any $\ell \in \mathbb{N}^*$ there exists $K_\ell \in (0, \infty)$ such that

$$W_1 \left(\Pi_\ell (f^N), f^{\otimes \ell} \right) \leq K_\ell \varepsilon(N).$$

Similar statements with other metrics can be also formulated (and shall be used in this paper): for some normed space of *smooth functions* $\mathcal{F} \subset C_b(E)$ (to be specified) and for any $\ell \in \mathbb{N}^*$ there exists $K_\ell \in (0, \infty)$ such that for any $\varphi \in \mathcal{F}^{\otimes \ell}$, $\|\varphi\|_{\mathcal{F}} \leq 1$, there holds

$$\left| \left\langle \Pi_\ell [f^N] - f^{\otimes \ell}, \varphi \right\rangle \right| \leq K_\ell \varepsilon(N).$$

The Wasserstein distance W_1 is recovered when \mathcal{F} is the space of Lipschitz functions.

Observe that in the latter statements the number of variables ℓ considered in the marginal is kept fixed as N goes to infinity. A stronger notion of **infinite-dimensional chaos** would be

$$\lim_{N \rightarrow \infty} \frac{W_1(f^N, f^{\otimes N})}{N} = 0$$

with corresponding quantitative formulations. This amounts to say that one can prove a sublinear control on K_ℓ in terms of ℓ in the previous statements. Variants for other metrics could also be considered.

Finally one can formulate an even stronger notion of **(infinite-dimensional) entropic chaos** (see [12] and definition (1.7) of the relative entropy below):

$$\frac{1}{N} H(f^N | \gamma^N) \xrightarrow{N \rightarrow \infty} H(f | \gamma)$$

with obvious quantitative versions. This notion of chaos is particularly interesting since it corresponds to the derivation of Boltzmann's entropy from the many-particle system entropies, we shall come back to this point.

Now, considering a sequence of symmetric³ N -particle densities

$$f^N \in C([0, \infty); P_{\text{sym}}(E^N))$$

and a 1-particle density of the expected mean field limit

$$f \in C([0, \infty); P(E)),$$

we say that there is *propagation of chaos* on some time interval $[0, T]$ if the f_0 -chaoticity of the initial family f_0^N implies the f_t -chaoticity of the family f_t^N for any time $t \in [0, T]$, according to one of definitions of chaoticity above.

1.4. Kac's program. As already said Kac proposed the problem of deriving the spatially homogeneous Boltzmann equation from a many-particle Markov jump process modeling the binary collisions, through its master equation (the equation on the law of the process). This amounts intuitively to consider the spatial variable as a hidden variable inducing ergodicity and markovian properties on the velocity variable. Although the latter point has not been proved so far to our knowledge, it is worth pointing out that it is at the same a very natural guess and an extremely interesting (and probably also difficult) open problem.

Interestingly enough here is in the words of Kac [42, p. 173] how his approach was inspired by Boltzmann's ideas: *"This formulation led to the well-known paradoxes which were fully discussed in the classical article of P. and T. Ehrenfest. These writers made it clear (a) that the "Stosszahlansatz" cannot be strictly derivable from purely dynamic considerations and (b) that the "Stosszahlansatz" has to be interpreted probabilistically. The recent attempts of Born and Green, Kirkwood and Bogoliubov to derive Boltzmann's equation from Liouville's equation and hence to justify the "Stosszahlansatz" dynamically are, in our opinion, incomplete, inasmuch as they do not make it clear at what point statistical assumptions are introduced. The "master equation" approach which we have chosen seems to us to follow closely the intentions of Boltzmann."* Even we partly disagree with this statement and we still believe that a fully satisfying justification of the "Stosszahlansatz" (molecular chaos) has to be derived from "purely dynamic consideration" (Newton's laws) as was pioneered by Lanford [45], the latter open problem seems tremendously difficult and it is clear that Kac raised a fascinating more tractable question: if we have to introduce stochasticity, at least *can we keep it under control all along the process of derivation of the Boltzmann equation and relate it to the dissipativity of the limit equation.* This is one of the main questions we attempt to answer in this work.

³i.e. invariant according to permutations of the particles.

Let us now discuss more in details the content of Kac's seminal paper [42]. He goes on to introduce a jump process on the $(N - 1)$ -dimensional sphere $\mathbb{S}^{N-1}(\sqrt{N})$ with radius \sqrt{N} corresponding to random rotations among pairs of variables chosen randomly, and occurring at random times following exponential laws. He then prove finite-dimensional chaos (with no rate) by a beautiful combinatorial argument, based on an infinite series "tree" representation of the solution according to the collision history of particles, and some Leibniz formula for the N -particle operator acting on tensor products.

He then raises several questions that we schematize as follows:

- (1) The first one concerns the restriction of the models as compared to realistic collision processes. The fact that the geometry of collision is simplified was relaxed by McKean [53] soon afterwards, however for cutoff Maxwell molecules for which the collision kernel does not depend on the modulus of the relative velocity. This seemingly technical issue is in fact related to deep difficulty for dealing with jump process whose jump times follow laws depending on the variables. In the words of Kac [42, p. 179]: "*The above proof suffers from the defect that it works only if the restriction on time is independent of the initial distribution. It is therefore inapplicable to the physically significant case of hard spheres because in this case our simple estimates yield a time restriction which depends on the initial distribution. A general proof that Boltzmann's property propagates in time is still lacking.*" In other words the question is: **can one prove propagation of chaos for the hard spheres collision process?**
- (2) Following closely the spirit of the previous question about going beyond the limit of Kac's original combinatorial insight in order to deal with realistic collision processes, it seems to us very natural to ask whether **one can prove propagation of chaos for the true Maxwell molecules collision process?** The difficulty lies now in the fact that the particle system can undergo infinite number of collisions in a finite time interval, and no "tree" representation of solutions is available. This is related with the physical interesting situation of *long-range interactions*, as well as with the mathematical interesting framework of *fractional derivative operator* and *Lévy walk*.
- (3) Kac then discusses the H -theorem of Boltzmann, which is not surprising as its original goal is the derivation of Boltzmann's equation and its irreversible feature. He makes the simple observation [42, p. 182] that the "*ergodic property of the Markov process under consideration*" steadily implies that it admits an infinite number of Liapunov functions, including the L^2 norm and Boltzmann's entropy. In contrast with it, the limit equation admits only (in general) Boltzmann's entropy as a Liapunov function. Kac then heuristically conjectures [42, Eq. (6.39)] that (in our notations) $H(f_t^N)/N \rightarrow H(f_t)$ along time, which would imply Boltzmann's H -theorem from the monotonicity of $H(f^N)/N$ for the Markov process. He concludes with: "*If the above steps could be made rigorous we would have a thoroughly satisfactory justification of Boltzmann's H-theorem.*" In our notation the question is **can one prove propagation of entropic chaos along time?**
- (4) He finally discusses the relaxation times, with the goal of deriving relaxation times of the limit equation from the many-particle system. This imposes to have estimates *independent of the number of particles* on this relaxation times. As a first natural step he therefore goes on to consider the L^2 spectral gap of the Markov process on the sphere and remark [42, p. 187]: "*Surprisingly enough this seems quite difficult and we have not succeeded in finding a proof. Even for the simplified model we have been considering, the question remains unsettled although we are able to give a reasonably explicit solution of the master equation.*" This question has triggered many beautiful works (see the next subsection), however it is easy

to convince oneself (see the discussion in [12] for instance) that there is no hope of passing to the limit $N \rightarrow \infty$ in this spectral gap estimate, even if the spectral gap is independent of N . The L^2 norm is catastrophic in infinite dimension. Therefore following quite closely the intention of Kac, we reframe the question in a setting which “tensorizes correctly in the limit $N \rightarrow \infty$ ”, that is in our notation: **can one prove relaxation times independent of the number of particles on the normalized Wasserstein distance $W(f^N, \gamma^N)/N$ or on the normalized relative entropy $H(f^N|\gamma^N)/N$, where γ^N denotes the N -particle invariant measure?**

This paper is concerned with solving the four questions outlined above.

Before concluding this subsection, let us briefly illustrate by some quotation [42, p. 178] that Kac was firmly believing in a “bottom-up” approach, which deduces properties on the Boltzmann equation from the Markov process: “*Since the master equation is truly descriptive of the physical situation, and since existence and uniqueness of the solution of the master equation are almost trivial, the preoccupation with existence and uniqueness theorems for the Boltzmann equation appears to be unjustified on grounds of physical interest and importance.*” Although as the time Kac was writing, almost no mathematical result (apart from [9, 10]) was available for the nonlinear spatially homogeneous Boltzmann equation, mathematical development has not followed this path, and as we have said, our approach shall be rather “top-down”, taking advantage of the bynow well-developed theory for this equation.

1.5. Review of previous results. For Boltzmann collision processes, Kac [42]-[43] has proved the point (1) in the case of his baby one-dimensional model. The key point in his analysis is a clever combinatorial use of a semi-explicit form of the solution (Wild sums). It was generalized by McKean [53] to the Boltzmann collision operator but only for “Maxwell molecules with cutoff”, i.e. roughly when the collision kernel B above is constant. In this case the combinatorial argument of Kac can be extended. Kac raised in [42] the question of proving propagation of chaos in the case of hard spheres and more generally unbounded collision kernels, although his method seemed impossible to extend (no semi-explicit combinatorial formula of the solution exists in these cases).

In the seventies, Grünbaum [36] then proposed in a very compact and abstract paper another method for dealing with hard spheres, based on the Trotter-Kato formula for semigroups and a clever functional framework (partially remindful of the tools used for mean-field limit for McKean-Vlasov equations). Unfortunately this paper was incomplete for two reasons: (1) It was based on two “unproved assumptions on the Boltzmann flow” (page 328): (a) existence and uniqueness for measure solutions and (b) a smoothness assumption. Assumption (a) was indeed recently proved in [32] using Wasserstein metrics techniques and in [29] adapting the classical DiBlasio trick [22], but concerning assumption (b), although it was inspired by cutoff Maxwell molecules (for which it is true), it fails for hard spheres (cf. the counterexample built by Lu and Wennberg in [49]) and is somehow “too rough” in this case. (2) A key part in the proof in this paper is the expansion of the “ H_f ” function, which is a clever idea of Grünbaum (and the starting point for our idea of developing an abstract differential calculus in order to control fluctuations) — however it is again too rough and is adapted for cutoff Maxwell molecules but not for hard spheres.

A completely different approach was undertaken by Sznitman in the eighties [66] (see also Tanaka [68] for partial results concerning non-cutoff Maxwell molecules). Starting from the observation that Grünbaum’s proof was incomplete, he gave a full proof of propagation of chaos for hard spheres. His work was based on: (1) a new uniqueness result for measures for the hard spheres Boltzmann equation (based on a probabilistic reasoning on an enlarged space of “trajectories”); (2) an idea already present in Grünbaum’s approach: reduce by a combinatorial argument on symmetric probabilities the question of

propagation of chaos to a law of large numbers on measures; (3) a new compactness result at the level of the empirical measures; (4) the identification of the limit by an “abstract test function” construction showing that the (infinite particle) system has trajectories included in the chaotic ones. Hence the method of Sznitman proves convergence but does not provide any rate for chaoticity.

Let us also emphasize that McKean in [52] studied fluctuations around deterministic limit for 2-speed Maxwellian gas and for the usual hard balls gas. Graham and Méléard in [35] have obtained a rate of convergence (of order $1/N$ for the ℓ -th marginal) on any bounded finite interval of the N -particle system to the deterministic Boltzmann dynamic in the case of Maxwell molecules under Grad's cut-off hypothesis. Finally Fournier and Méléard in [30, 31] have obtained the convergence of the Monte-Carlo approximation (with numerical cutoff) of the Boltzmann equation for true Maxwell molecules with a rate of convergence (depending on the numerical cutoff and on the number N of particles).

After we had finished writing our paper, we were told by I. Bailleul about the recent book [44] by V. N. Kolokoltsov. This book is focused on fluctuation estimates of central limit theorem type. It does not prove quantitative propagation of chaos but weaker estimates (and on finite time intervals), however the comparison of generators for the many-particle and the limit semigroup is reminiscent of our work.

1.6. The method. The main inspiration at first came from Grünbaum's paper [36]. Our original goal was to construct a general and robust method able to deal with mixture of jump and diffusion processes, as it occurs for granular gases (see for this point the companion paper [58]). It turns out that it leads us to develop a new theory, inspiring from more recent tools such as the course of Lions on “Mean-field games” at Collège de France, and the master courses of Méléard [55] and Villani [70] on mean-field limits. One of the byproduct of our paper is that we make fully rigorous the original intuition of Grünbaum in order to prove propagation of chaos for the Boltzmann velocities jump process associated to hard spheres contact interactions.

As Grünbaum [36] we shall use a duality argument. We introduce S_t^N the semigroup associated to the flow of the N -particle system and T_t^N its “dual” semigroup. We also introduce S_t^{NL} the (nonlinear) semigroup associated to the meanfield dynamic (the exponent “NL” recalling that the limit semigroup is nonlinear in the most physics interesting cases) as well as T_t^∞ the associated (linear) “pushforward” semigroup (see below for the definition). Then we will prove the above kind of convergence on the linear semigroups T_t^N and T_t^∞ .

The first step consists in defining a common functional framework in which the N -particle dynamic and the limit dynamic make sense so that we can compare them. Hence we work at the level of the “full” limit space $P(P(E))$ (see below). Then we shall identify the regularity required in order to prove the “consistency estimate” between the generators G^N and G^∞ of the dual semigroups T_t^N and T_t^∞ , and finally prove a corresponding “stability estimate” at the level of the limiting semigroup S_t^{NL} .

The latter crucial step shall lead us to introduce an abstract differential calculus for functions acting on measures endowed with various metrics. More precisely, we shall define functions of class $C^{1,\delta}$ on a probability space by working on affine subspaces of the probability space for which the tangent space has a Banach space structure. This notion is related but different from the notion of derivability developed in the theory of gradient flow by Ambrosio, Otto, Villani and co-authors in [2, 41, 61] as well as to the differentiability notion introduced by Lions in [46].

Another viewpoint on this method is to consider it as some kind of accurate version (in the sense that it establishes a rate of convergence) of the BBGKY hierarchy method for proving propagation of chaos and mean-field limit on statistical solutions. This viewpoint is extensively explored and made rigorous in Section 8 where we revisit the BBGKY method

for the spatially homogeneous equation as developed in [3]. The proof of uniqueness for statistical solutions to the hierarchy becomes straightforward within our framework by using differentiability of the limit semigroup as a function acting on probabilities.

This general method is the purpose of Theorem 3.1. It is, we hope, interesting per se for several reasons: (1) it is fully quantitative, (2) it is highly flexible in terms of the functional spaces used in the proof, (3) it requires a minimal amount of informations on the N -particle systems but more stability information on the limiting PDE (we intentionally presented the assumption as for the proof of the convergence of a numerical scheme, which was our “methodological model”), (4) the “differential stability” conditions that are required on the limiting PDE seem (to our knowledge) new, at least at the level of Boltzmann or more generally transport equations.

1.7. Main results. Let us give some simplified versions of the main results in this paper. All the abstract objects shall be fully introduced in the next sections.

Theorem 1.1 (Summary of the main results). *Consider some initial distribution $f_0 \in P(\mathbb{R}^d) \cap L^\infty(\mathbb{R}^d)$ with compact support or polynomial moment bounds. Consider the corresponding solution f_t to the spatially homogeneous Boltzmann equation for hard spheres or Maxwell molecules, and the solution f_t^N of the corresponding Kac’s jump process starting either from the N -fold tensorization of f_0 or the latter conditioned to \mathcal{S}^N .*

The results can be classified into three main statements:

- (1) **Quantitative uniform in time propagation of chaos, finite or infinite dimensional, in weak measure distance** (cf. Theorems 5.1-5.3-6.1-6.2):

$$\forall N \geq 1, \forall 1 \leq \ell \leq N, \sup_{t \geq 0} \frac{W_1 \left(\Pi_\ell f_t^N, \left(f_t^{\otimes \ell} \right) \right)}{\ell} \leq \alpha(N)$$

for some $\alpha(N) \rightarrow 0$ as $N \rightarrow \infty$. In the case $f_0^N = f_0^{\otimes N}$ one has moreover explicit power law (for Maxwell molecules) or logarithmic rate (for hard spheres) estimates on α ,

- (2) **Propagation of entropic chaos** (cf. Theorem 7.10-(i)): *Consider the case where the initial datum of the many-particle system is restricted to \mathcal{S}^N . If the initial datum is entropy-chaotic in the sense*

$$\frac{1}{N} H(f_0^N | \gamma^N) \xrightarrow{N \rightarrow +\infty} H(f_0 | \gamma)$$

with

$$(1.7) \quad H(f_0^N | \gamma^N) := \int_{\mathcal{S}^N} \log \frac{df_0^N}{d\gamma^N} f_0^N(dV) \quad \text{and} \quad H(f_0 | \gamma) := \int_{\mathbb{R}^d} f_0 \log \frac{f_0}{\gamma} dv$$

and where γ is the gaussian equilibrium with energy \mathcal{E} and γ^N is the uniform probability measure on \mathcal{S}^N , then the solution is also entropy chaotic for any later time:

$$\forall t \geq 0, \quad \frac{1}{N} H(f_t^N | \gamma^N) \xrightarrow{N \rightarrow +\infty} H(f_t | \gamma).$$

This proves the derivation of the H-theorem in this context.

- (3) **Quantitative estimates on relaxation times, independent of the number of particles** (cf. Theorems 5.3-6.2 and Theorem 7.10-(ii)): *Consider the case where the initial datum of the many-particle system is restricted to \mathcal{S}^N . Then we have*

$$\forall N \geq 1, \forall 1 \leq \ell \leq N, \forall t \geq 0, \quad \frac{W_1 \left(\Pi_\ell f_t^N, \Pi_\ell(\gamma^N) \right)}{\ell} \leq \beta(t)$$

for some $\beta(t) \rightarrow 0$ as $t \rightarrow \infty$. Moreover in the case of Maxwell molecules, and assuming moreover that the Fisher information of the initial datum f_0 is finite:

$$(1.8) \quad I(f_0) := \int_{\mathbb{R}^d} \frac{|\nabla_v f_0|^2}{f_0} dv < +\infty,$$

the following estimate also holds:

$$\forall N \geq 1, \quad \frac{1}{N} H(f_t^N | \gamma^N) \leq \beta(t)$$

for some function $\beta(t) \rightarrow 0$ as $t \rightarrow \infty$.

1.8. Some open questions. Here are a few questions among those raised by this work:

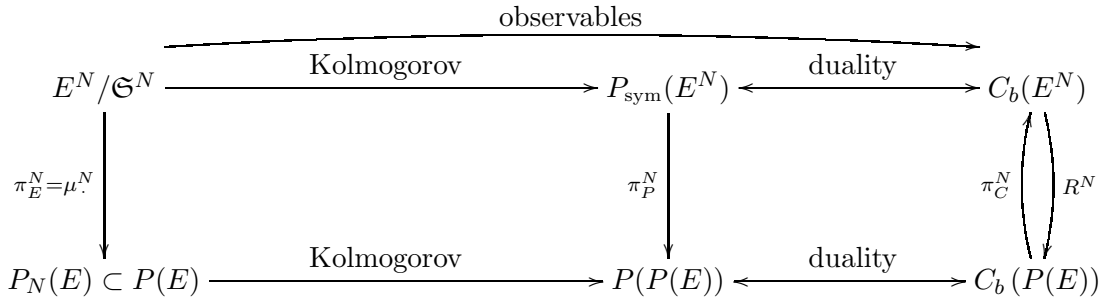
- (1) What about the optimal rate in the chaoticity estimates along time? Our method reduces this question to the chaoticity estimates at initial time, and therefore to the optimal rate in the quantitative law of large numbers for measures according to various weak measure distances.
- (2) What about the optimal rate in the relaxation times (uniformly in the number of particles)? Spectral gap studies predict exponential rates, both for the many-particle system and for the limit system, however our rates are far from it!
- (3) Can uniform in time propagation of chaos be proved for non-reversible jump processes (such as inelastic collision processes) for which the invariants measures γ^N and γ are not explicitly known (e.g. granular gases)?

1.9. Plan of the paper. In Section 2 we set the abstract functional framework together with the general assumption and in Section 3 we state and prove the abstract Theorem 3.1. In Section 4 we present some tools and results on weak measure distances, the construction of initial data restricted to submanifolds for the many-particle system, and the rate for sampling the limit distribution by empirical measures. In Section 5 we apply the method to (true) Maxwell molecules (Theorems 5.1 and 5.3). In Section 6 we apply the method to hard spheres (Theorems 6.1 and 6.2). Section 7 is devoted to the study of entropic chaos. Finally in Section 8 we revisit the BBGKY hierarchy method for the spatially homogeneous Boltzmann equation in the light of our framework.

2. THE ABSTRACT SETTING

In this section we shall state and prove the key abstract result. This will motivate the introduction of a general functional framework.

2.1. The general functional framework of the duality approach. Let us set up the framework. Here is a diagram which sums up the duality approach (norms and duality brackets shall be specified in Subsections 2.4):



In this diagram:

- E denotes a Polish space:

This is a separable completely metrizable topological space. We shall denote the distance on this space by d_E in the sequel.

- \mathfrak{S}^N denotes the N -permutation group.
- $P_{\text{sym}}(E^N)$ denotes the set of *symmetric probabilities* on E^N :

Given a permutation $\sigma \in \mathfrak{S}^N$, a vector

$$V = (v_1, \dots, v_N) \in E^N,$$

a function $\varphi \in C_b(E^N)$ and a probability $f^N \in P(E^N)$, we successively define

$$V_\sigma = (v_{\sigma(1)}, \dots, v_{\sigma(N)}) \in E^N,$$

and

$$\varphi_\sigma \in C_b(E^N) \text{ by setting } \varphi_\sigma(V) = \varphi(V_\sigma)$$

and finally

$$f_\sigma^N \in P(E^N) \text{ by setting } \langle f_\sigma^N, \varphi \rangle = \langle f^N, \varphi_\sigma \rangle.$$

We then say that a probability f^N on E^N is symmetric if it is invariant under permutations:

$$\forall \sigma \in \mathfrak{S}^N, \quad f_\sigma^N = f^N.$$

- The probability measure μ_V^N denotes the *empirical measure*:

$$\mu_V^N := \frac{1}{N} \sum_{i=1}^N \delta_{v_i}, \quad V = (v_1, \dots, v_N) \in E^N$$

where δ_{v_i} denotes the Dirac mass on E at point $v_i \in E$.

- $P_N(E)$ denotes the subset $\{\mu_V^N, V \in E^N\}$ of empirical measures of $P(E)$.
- $P(P(E))$ denotes the set of probabilities on the polish space $P(E)$.
- $C_b(P(E))$ denotes the space of continuous and bounded functions on $P(E)$:

This space shall be endowed with either the weak or strong topologies (see Subsection 2.4), and later with some metric differential structure.

- The map π_E^N from E^N/\mathfrak{S}^N to $P_N(E)$ is defined by

$$\forall V \in E^N/\mathfrak{S}^N, \quad \pi_E^N(V) := \mu_V^N.$$

- The map π_C^N from $C_b(P(E))$ to $C_b(E^N)$ is defined by

$$\forall \Phi \in C_b(P(E)), \quad \forall V \in E^N, \quad (\pi_C^N \Phi)(V) := \Phi(\mu_V^N).$$

- The map R^N from $C_b(E^N)$ to $C_b(P(E))$ is defined by:

$$\forall \varphi \in C_b(E^N), \quad \forall \rho \in P(E), \quad R_\varphi^N(\rho) := \langle \rho^{\otimes N}, \varphi \rangle.$$

- The map π_P^N from $P_{\text{sym}}(E^N)$ to $P(P(E))$ is defined by:

$$\langle \pi_P^N f^N, \Phi \rangle = \langle f^N, \pi_C^N \Phi \rangle = \int_{E^N} \Phi(\mu_V^N) f^N(dV)$$

for any $f^N \in P_{\text{sym}}(E^N)$ and any $\Phi \in C_b(P(E))$, where the first bracket means $\langle \cdot, \cdot \rangle_{P(P(E)), C_b(P(E))}$ and the second bracket means $\langle \cdot, \cdot \rangle_{P(E^N), C_b(E^N)}$.

Let us now discuss the “horizontal” arrows:

- The arrows pointing from the first column to the second one consists in writing the *Kolmogorov* equation associated with the many-particle stochastic Markov process.
- The arrows pointing from the second column to the third column consists in writing the dual evolution semigroup (note that the N -particle dynamics is linear). As we shall discuss later the dual spaces of the spaces of probabilities on the phase space can be interpreted as the spaces of observables on the original systems.

Remark 2.1. Consider a random variable V on E^N with law $f^N \in P(E^N)$. It is common notation in probability to simply write μ_V^N for the random variable on $P(E)$ with law $\pi_P^N f^N \in P(P(E))$. Our notation is slightly less compact and intuitive, but at the same time more accurate.

Remark 2.2. Our functional framework shall be applied to weighted probability spaces rather than directly in $P(E)$. More precisely, for a given weight function $m : E \rightarrow \mathbb{R}_+$ we shall use affine subsets of the weighted probability space

$$\{f \in P(E); M_m(f) := \langle f, m \rangle < \infty\}$$

as our basis functional spaces. Typical examples are

$$m(v) := \tilde{m}(d_E(v, v_0)) \text{ for some fixed } v_0 \in E \text{ with } \tilde{m}(z) = z^k \text{ or } \tilde{m}(z) = e^{az^k}, a, k > 0.$$

We shall sometimes abuse notation by writing M_k for M_m when $\tilde{m}(z) = z^k$ in the above example.

2.2. The N -particle semigroups. Let us introduce the mathematical semigroups describing the evolution of objects living in these spaces, for any $N \geq 1$.

Step 1. Consider the trajectories $\mathcal{V}_t^N \in E^N$, $t \geq 0$, of the particles (Markov process viewpoint). We make the further assumption that this flow commutes with permutations:

$$\text{For any } \sigma \in \mathfrak{S}^N, \text{ the solution at time } t \text{ starting from } (\mathcal{V}_0^N)_\sigma \text{ is } (\mathcal{V}_t^N)_\sigma.$$

This reflects mathematically the fact that particles are indistinguishable.

Step 2. This flow on E^N yields a corresponding semigroup S_t^N acting on $P_{\text{sym}}(E^N)$ for the probability density of particles in the phase space E^N (statistical viewpoint), defined through the formula

$$\forall f^N \in P_{\text{sym}}(E^N), \varphi \in C_b(E^N), \quad \langle S_t^N(f^N), \varphi \rangle = \mathbb{E}(\varphi(\mathcal{V}_t^N))$$

where the bracket obviously denotes the duality bracket between $P(E^N)$ and $C_b(E^N)$ and \mathbb{E} denotes the expectation associated to the probability space in which the process \mathcal{V}_t^N is built. Since the flow (\mathcal{V}_t^N) commutes with permutation, the semigroup S_t^N acts on $P_{\text{sym}}(E^N)$. In other words, if the law f_0^N of \mathcal{V}_0^N belongs to $P_{\text{sym}}(E^N)$, then for later times the law f_t^N of \mathcal{V}_t^N also belongs to $P_{\text{sym}}(E^N)$. And one can associate to this semigroup a linear evolution equation with generator denoted by A_N :

$$\partial_t f^N = A^N f^N, \quad f^N \in P_{\text{sym}}(E^N),$$

which can be interpreted as the forward Kolmogorov equation on the law of \mathcal{V}_t^N .

Step 3. We assume that there exists a semigroup T_t^N acting on the functions space $C_b(E^N)$ of *observables* on the evolution system (\mathcal{V}_t^N) on E^N (see the discussion in the next remark) which is *dual* to the semigroup S_t^N . More precisely, we assume

$$\forall f^N \in P(E^N), \varphi \in C_b(E^N), \quad \langle f^N, T_t^N(\varphi) \rangle = \langle S_t^N(f^N), \varphi \rangle.$$

We associate to this semigroup the following *linear* evolution equation with generator denoted by G^N :

$$\partial_t \varphi = G^N(\varphi), \quad \varphi \in C_b(E^N).$$

2.3. The mean-field limiting semigroup. We now define the evolution of the limiting mean-field equation.

Step 1. Consider a semigroup S_t^{NL} acting on $P(E)$ associated with an evolution equation and some operator Q :

For any $f_0 \in P(E)$ (assuming possibly some additional moment bounds), then $S_t^{NL}(f_0) := f_t$ where $f_t \in C(\mathbb{R}_+, P(E))$ is the solution to

$$(2.1) \quad \partial_t f_t = Q(f_t), \quad f_{|t=0} = f_0.$$

This semigroup and the operator Q are typically *nonlinear* for mean-field models, namely bilinear in case of Boltzmann's collisions interactions.

Step 2. Consider then the associated *pushforward semigroup* T_t^∞ acting on $C_b(P(E))$:

$$\forall f \in P(E), \Phi \in C_b(P(E)), \quad T_t^\infty[\Phi](f) := \Phi(S_t^{NL}(f)).$$

(Again additional moment bounds can be required on f in order to make this definition rigorous.) Note carefully that T_t^∞ is always *linear* as a function of Φ , although of course $T_t^\infty[\Phi](f)$ is not linear in general as a function of f . We shall associate (when possible) the following *linear* evolution equation on $C_b(P(E))$ with some generator denoted by G^∞ :

$$\partial_t \Phi = G^\infty(\Phi).$$

Remark 2.3. The semigroup T_t^∞ can be interpreted physically as the semigroup of the evolution of *observables* of the nonlinear equation (2.1). Let us give a short heuristic explanation. Consider a nonlinear ordinary differential equation

$$\frac{dV}{dt} = F(V) \quad \text{on } \mathbb{R}^d \quad \text{with } \nabla_v \cdot F \equiv 0 \quad \text{and } V_{|t=0} = v$$

with divergence-free vector field for simplicity. One can then define formally the *linear* Liouville transport partial differential equation

$$\partial_t f + \nabla_v \cdot (F f) = 0,$$

where $f = f_t(v)$ is a time-dependent probability density over the phase space \mathbb{R}^d , whose solution is given (at least formally) by following the characteristics backward $f_t(v) = f_0(V_{-t}(v))$.⁴ Now, instead of the Liouville viewpoint, one can adopt the viewpoint of *observables*, that is functions depending on the position of the system in the phase space (e.g. energy, momentum, etc.) For some observable function φ_0 defined on \mathbb{R}^d , the evolution of the value of this observable along the trajectory is given by $\varphi_t(v) = \varphi_0(V_t(v))$ and φ_t is solution to the following *dual* linear PDE

$$\partial_t \varphi - F \cdot \nabla_v \varphi = 0.$$

Now let us consider a nonlinear evolution system

$$\frac{df}{dt} = Q(f) \quad \text{in an abstract space } f \in \mathcal{H}.$$

⁴The formula would be slightly more complicated in the case where F is not divergence-free and would involve the jacobian of F , but this is not relevant for the present heuristic discussion.

By analogy we define two linear evolution systems on the larger functional spaces $P(\mathcal{H})$ and $C_b(\mathcal{H})$: first the abstract Liouville equation

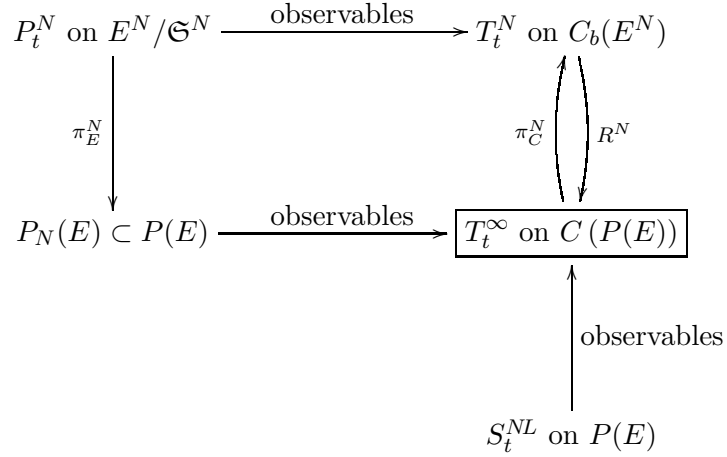
$$\partial_t \pi + \frac{\delta}{\delta f} \cdot (Q(f) \pi) = 0, \quad \pi \in P(\mathcal{H})$$

and second the abstract equation for the evolution of observables

$$\partial_t \Phi - Q(f) \cdot \frac{\delta \Phi}{\delta f}(f) = 0, \quad \Phi \in C_b(\mathcal{H}).$$

However in order to make sense of this heuristic, the scalar product have to be defined correctly as duality brackets, and, most importantly, a differential calculus on \mathcal{H} has to be defined rigorously. Taking $\mathcal{H} = P(E)$, this provides an intuition for our functional construction, as well as for the formula of the generator G^∞ below (compare the previous construction, as well as for the formula of the generator G^∞ below (compare the previous equation with formula (2.8)). Be careful that when $\mathcal{H} = P(E)$, the abstract Liouville and observable equations refers to trajectories *in the space of probabilities* $P(E)$ (i.e. solutions to the nonlinear equation (2.1)), and not trajectories of a particle in E . Note also that for a *dissipative equation* at the level of \mathcal{H} (such as the Boltzmann equation), it seems more convenient to use the observable equation rather than the Liouville equation since “forward characteristics” can be readily used in order to construct the solutions to this observable equation.

Summing up we obtain the following picture for the semigroups:



Hence a key point of our construction is that, through the evolution of *observables*, we shall “interface” the two evolution systems (the nonlinear limiting equation and the N -particle system) via the applications π_C^N and R^N . From now on we shall denote $\pi^N = \pi_C^N$.

2.4. The metric issue. $C(P(E))$ is our fundamental space in which we shall compare (through their observables) the semigroups of the N -particle system and the limiting mean-field equation. Let us make the topological and metric structures used on $P(E)$ more precise. At the topological level there are two canonical choices (which determine two different sets $C(P(E))$):

- (1) The strong topology which is associated to the total variation norm, denoted by $\|\cdot\|_{M^1}$; the corresponding set shall be denoted by $C_b(P(E), TV)$.
- (2) The weak topology, i.e. the trace on $P(E)$ of the weak topology on $M^1(E)$ (the space of Radon measures on E with finite mass) induced by $C_b(E)$; the corresponding set shall be denoted by $C_b(P(E), w)$.

It is clear that

$$C_b(P(E), w) \subset C_b(P(E), TV).$$

The supremum norm $\|\Phi\|_{L^\infty(P(E))}$ does *not* depend on the choice of topology on $P(E)$, and induces a Banach topology on the space $C_b(P(E))$. The transformations π^N and R^N satisfy:

$$(2.2) \quad \|\pi^N \Phi\|_{L^\infty(E^N)} \leq \|\Phi\|_{L^\infty(P(E))} \quad \text{and} \quad \|R_\phi^N\|_{L^\infty(P(E))} \leq \|\phi\|_{L^\infty(E^N)}.$$

The transformation π^N is well defined from $C_b(P(E), w)$ to $C_b(E^N)$, but in general, it does not map $C_b(P(E), TV)$ into $C_b(E^N)$ since the map

$$V \in E^N \mapsto \mu_V^N \in (P(E), TV)$$

is not continuous.

In the other way round, the transformation R^N is well defined from $C_b(E^N)$ to $C_b(P(E), w)$, and therefore also from $C_b(E^N)$ to $C_b(P(E), TV)$: for any $\phi \in C_b(E^N)$ and for any sequence $f_k \rightharpoonup f$ weakly, we have $f_k^{\otimes N} \rightharpoonup f^{\otimes N}$ weakly, and then $R^N[\phi](f_k) \rightarrow R^N[\phi](f)$.

There are many different possible metric structures inducing the weak topology on $C_b(P(E), w)$. The mere notion of continuity does not require discussing these metrics, but any subspace of $C_b(P(E), w)$ with differential regularity shall strongly depend on this choice, which motivates the following definitions.

Definition 2.4. For a given *weight function* $m : E \rightarrow \mathbb{R}_+$, we define the subspaces of probabilities:

$$P_m := \{f \in P(E); \langle f, m \rangle < \infty\}.$$

We also define the corresponding bounded subsets for $a > 0$

$$\mathcal{B}P_{m,a} := \{f \in P_m; \langle f, m \rangle \leq a\}.$$

For a given *constraint function* $\mathbf{m} : E \rightarrow \mathbb{R}^D$ such that the components \mathbf{m} are controlled by m , we define the corresponding *constrained subsets*

$$P_{m,\mathbf{m},\mathbf{r}} := \{f \in P_m; \langle f, \mathbf{m} \rangle = \mathbf{r}\}, \quad \mathbf{r} \in \mathbb{R}^D.$$

We also define their corresponding bounded subsets for $a > 0$

$$\mathcal{B}P_{m,\mathbf{m},a,\mathbf{r}} := \{f \in \mathcal{B}P_{m,a}; \langle f, \mathbf{m} \rangle = \mathbf{r}\},$$

and the corresponding vectorial space of “increments”

$$\mathcal{I}P_{m,\mathbf{m}} := \{f_2 - f_1; \exists \mathbf{r} \in \mathbb{R}^D \text{ s.t. } f_1, f_2 \in P_{m,\mathbf{m},\mathbf{r}}\}$$

and

$$\mathcal{I}P_{m,\mathbf{m},a} := \{f_2 - f_1; \exists \mathbf{r} \in \mathbb{R}^D \text{ s.t. } f_1, f_2 \in \mathcal{B}P_{m,\mathbf{m},a,\mathbf{r}}\}.$$

Let us now define the notion of distances over probabilities that we shall consider.

Definition 2.5. Consider a weight function $m_{\mathcal{G}}$ and a constraint function $\mathbf{m}_{\mathcal{G}}$. We shall use for the associated spaces of the previous definition the simplified notation $P_{\mathcal{G}}$ for P_m , $P_{\mathcal{G},\mathbf{r}}$ for $P_{m,\mathbf{m},\mathbf{r}}$, $\mathcal{B}P_{\mathcal{G},a}$ for $\mathcal{B}P_{m,a}$, $\mathcal{B}P_{\mathcal{G},a,\mathbf{r}}$ for $\mathcal{B}P_{m,\mathbf{m},a,\mathbf{r}}$, $\mathcal{I}P_{\mathcal{G}}$ for $\mathcal{I}P_{m,\mathbf{m}}$ and $\mathcal{I}P_{\mathcal{G},a}$ for $\mathcal{I}P_{m,\mathbf{m},a}$.

We shall consider that a distance $d_{\mathcal{G}}$ which

- (1) either is defined the whole space $P_{\mathcal{G}}$ (i.e. whatever the values of the constraints),
- (2) or such that there is a Banach space $\mathcal{G} \supset \mathcal{I}P_{\mathcal{G}}$ endowed with a norm $\|\cdot\|_{\mathcal{G}}$ such that $d_{\mathcal{G}}$ is defined on any $P_{\mathcal{G},\mathbf{r}}$, $\mathbf{r} \in \mathbb{R}^D$, by setting

$$\forall f, g \in P_{\mathcal{G},\mathbf{r}}, \quad d_{\mathcal{G}}(f, g) := \|g - f\|_{\mathcal{G}}.$$

Let us finally define a quantitative Hölder notion of equivalence for the distances over probabilities that we shall consider.

Definition 2.6. Consider some weight and constraint functions $m_{\mathcal{G}}$, $\mathbf{m}_{\mathcal{G}}$. We say that two metrics d_0 and d_1 defined on $P_{\mathcal{G}}$ are *Hölder equivalent on bounded sets* if there exists $\kappa \in (0, \infty)$ and, for any $a \in (0, \infty)$, there exists $C_a \in (0, \infty)$ such that

$$\forall f, g \in \mathcal{B}P_{\mathcal{G},a} \quad d_0(f, g) \leq C_a [d_1(f, g)]^{\kappa}, \quad d_1(f, g) \leq C_a [d_0(f, g)]^{\kappa}$$

for some constant C_a depending on $a > 0$. If d_0 and d_1 are resulting from some normed spaces \mathcal{G}_0 and \mathcal{G}_1 we generalize this definition as

$$\forall \mathbf{r} \in \mathbb{R}^D, \forall f, g \in \mathcal{B}P_{\mathcal{G},\mathbf{r},a}, \quad d_0(f, g) \leq C_a [d_1(f, g)]^{\kappa}, \quad d_1(f, g) \leq C_a [d_0(f, g)]^{\kappa}$$

for some $\kappa \in (0, \infty)$ and some constant C_a depending on $a > 0$.

Example 2.7. The choice

$$m_e = m_{\mathcal{G}} := 1, \quad \mathbf{m}_{\mathcal{G}} := 0, \quad \|\cdot\|_{\mathcal{G}} := \|\cdot\|_{M^1}$$

recovers $P_{\mathcal{G}}(E) = P(E)$. More generally one can choose

$$m_e = 1, \quad m_{\mathcal{G}_k}(v) := d_E(v, v_0)^k, \quad \mathbf{m}_{\mathcal{G}_k} := 0, \quad \|\cdot\|_{\mathcal{G}_k} := \|\cdot\|_{M^1} \cdot d_E(v, v_0)^k.$$

For $k_3 > 0$ and $k_1 > k_2$, the spaces $P_{\mathcal{G}_{k_2}}$ and $P_{\mathcal{G}_{k_3}}$ are topologically uniformly equivalent on bounded sets of $P_{\mathcal{G}_{k_1}}$.

Example 2.8. There are many distances on $P(E)$ which induce the weak topology, see for instance [64]. In the next section, we present some of them which have a practical interest for us, and which are all topologically uniformly equivalent on bounded sets of $P(E)$ in the sense of the previous definition, with the choice of a convenient (strong enough) weight function.

2.5. Distances on probabilities. Let us discuss some well-known distances on $P(\mathbb{R}^d)$ (or defined on subsets of $P(\mathbb{R}^d)$) which shall be useful for the sequel. These distances are all topologically equivalent to the weak topology $\sigma(P(E), C_b(E))$ on the sets $\mathcal{B}P_{k,a}(E)$ for k large enough and for any $a \in (0, \infty)$, and they are all uniformly topologically equivalent (see [69, 17] and section 2.5.6). We refer to [64, 72, 17] and the references therein for more details on these distances.

2.5.1. Dual-Hölder –or Zolotarev’s– distances. Denote by d_E a distance on E and let us fix $v_0 \in E$ (e.g. $v_0 = 0$ when $E = \mathbb{R}^d$ in the sequel). Denote by $C_0^{0,s}(E)$, $s \in (0, 1)$ (resp. $\text{Lip}_0(E)$) the set of s -Hölder functions (resp. Lipschitz functions) on E vanishing at one arbitrary point $v_0 \in E$ endowed with the norm

$$[\varphi]_s := \sup_{x,y \in E} \frac{|\varphi(y) - \varphi(x)|}{d_E(x,y)^s}, \quad s \in (0, 1], \quad [\varphi]_{\text{Lip}} := [\varphi]_1.$$

We then define the dual norm: take $m_{\mathcal{G}} := 1$, $\mathbf{m}_{\mathcal{G}} := 0$ and $P_{\mathcal{G}}(E)$ endowed with

$$\forall f, g \in P_{\mathcal{G}}, \quad [g - f]_s^* := \sup_{\varphi \in C_0^{0,s}(E)} \frac{\langle g - f, \varphi \rangle}{[\varphi]_s}.$$

2.5.2. *Wasserstein distances.* Given $q \in (0, \infty)$, define W_q on

$$P_{\mathcal{G}}(E) = P_q(E) := \{f \in P(E); \langle f, d_E(\cdot, v_0)^q \rangle < \infty\}$$

by

$$\forall f, g \in P_q(E), \quad W_q(f, g) := \inf_{\Pi \in \Pi(f, g)} \int_{E \times E} d_E(x, y)^q \Pi(dx, dy),$$

where $\Pi(f, g)$ denote the set of probability measures $\Pi \in P(E \times E)$ with marginals f and g :

$$\Pi(A, E) = f(A) \quad \text{and} \quad \Pi(E, A) = g(A) \quad \text{for any Borel set } A \subset E.$$

Note that for $V_1, V_2 \in E^N$ and any $q \in [1, \infty)$, one has

$$(2.3) \quad W_q(\mu_{V_1}^N, \mu_{V_2}^N) = d_{\ell^q(E^N/\mathfrak{S}_N)}(V_1, V_2) := \min_{\sigma \in \mathfrak{S}_N} \left(\frac{1}{N} \sum_{i=1}^N d_E((V_1)_i, (V_2)_{\sigma(i)})^q \right)^{1/q},$$

and that

$$\forall f, g \in P_1(E), \quad W_1(f, g) = [f - g]_1^* = \sup_{\varphi \in \text{Lip}_0(E)} \langle f - g, \varphi \rangle.$$

2.5.3. *Fourier-based norms.* Given $E = \mathbb{R}^d$, $m_{\mathcal{G}_1} := |v|$, $\mathbf{m}_{\mathcal{G}_1} := 0$, let us define

$$\forall f \in \mathcal{I}P_{\mathcal{G}_1}, \quad \|f\|_{\mathcal{G}_1} = |f|_s := \sup_{\xi \in \mathbb{R}^d} \frac{|\hat{f}(\xi)|}{|\xi|^s}, \quad s \in (0, 1],$$

where \hat{f} denotes the Fourier transform of f defined through the expression (when say $f \in L^1$)

$$\hat{f}(\xi) := \int_{\mathbb{R}^d} f(x) e^{-ix \cdot \xi} dx.$$

Similarly, given $E = \mathbb{R}^d$, $m_{\mathcal{G}_2} := |v|^2$, $\mathbf{m}_{\mathcal{G}_2} := v$, we define

$$\forall f \in \mathcal{I}P_{\mathcal{G}_2}, \quad \|f\|_{\mathcal{G}_2} = |f|_s := \sup_{\xi \in \mathbb{R}^d} \frac{|\hat{f}(\xi)|}{|\xi|^s}, \quad s \in (1, 2].$$

Obviously high-order versions of this norm could be defined similarly by increasing the number of constraints. However we shall see in the next subsection how to extend this notion of distance without constraints.

2.5.4. *More Fourier-based norms.* More generally, given $E = \mathbb{R}^d$ and $k \in \mathbb{N}^*$, we set

$$m_{\mathcal{G}} := |v|^k, \quad \mathbf{m}_{\mathcal{G}} := (v^\alpha)_{\alpha \in \mathbb{N}^d, |\alpha| \leq k-1}$$

with $|\alpha| = \alpha_1 + \dots + \alpha_d$ and

$$v^\alpha = (v_1^{\alpha_1}, \dots, v_d^{\alpha_d}), \quad \alpha = (\alpha_1, \dots, \alpha_d) \in \mathbb{N}^d,$$

and we define

$$\forall f \in \mathcal{I}P_{\mathcal{G}}, \quad \|f\|_{\mathcal{G}} = |f|_s := \sup_{\xi \in \mathbb{R}^d} \frac{|\hat{f}(\xi)|}{|\xi|^s}, \quad s \in (0, k].$$

In fact, we may extend the above norm to $M_k^1(\mathbb{R}^d)$ in the following way. We first define for

$$f \in M_{k-1}^1(\mathbb{R}^d) \quad \text{and} \quad \alpha \in \mathbb{N}^d, |\alpha| = \alpha_1 + \dots + \alpha_d \leq k-1$$

the following moment:

$$M_\alpha[f] := \int_{\mathbb{R}^d} v^\alpha f(dv).$$

Consider a fixed (once for all) function $\chi \in C_c^\infty(\mathbb{R}^d)$ (compact support), such that $\chi \equiv 1$ on the set $\{v \in \mathbb{R}^d, |v| \leq 1\}$. This implies in particular

$$\int_{\mathbb{R}^d} \mathcal{F}^{-1}(\chi)(v) dv = \chi(0) = 1.$$

Then we define $\mathcal{M}_k[f]$ through its Fourier transform

$$\hat{\mathcal{M}}_k[f](\xi) := \chi(\xi) \left(\sum_{|\alpha| \leq k-1} M_\alpha[f] \frac{\xi^\alpha}{\alpha!} \right), \quad \alpha! := \alpha_1! \dots \alpha_d!$$

Note that this is a mollified version of the $(k-1)$ -Taylor expansion of \hat{f} at $\xi = 0$. Then we may define the semi-norms

$$|f|_k := \sup_{\xi \in \mathbb{R}^d} \frac{|\hat{f}(\xi) - \hat{\mathcal{M}}_k[f](\xi)|}{|\xi|^k}$$

and the *norms*

$$\|f\|_k := |f|_k + \sum_{\alpha \in \mathbb{N}^d, |\alpha| \leq k-1} |M_\alpha[f]|.$$

2.5.5. *Negative Sobolev norms.* Given $s \in (d/2, d/2 + 1/2)$ and

$$E = \mathbb{R}^d, \quad m_{\mathcal{G}_1} := |v|, \quad \mathbf{m}_{\mathcal{G}_1} := 0$$

we define the following negative homogeneous Sobolev norm

$$\forall f \in \mathcal{IP}_{\mathcal{G}_1}, \quad \|f\|_{\mathcal{G}_1} = \|f\|_{\dot{H}^{-s}(\mathbb{R}^d)} := \left\| \frac{\hat{f}(\xi)}{|\xi|^s} \right\|_{L^2}.$$

Observe that probabilities are included in the corresponding non homogeneous negative Sobolev space $H^{-s}(\mathbb{R}^d)$ as soon as $s > d/2$.

Similarly, given $s \in [d/2 + 1/2, d/2 + 1)$ and

$$E = \mathbb{R}^d, \quad m_{\mathcal{G}_2} := |v|^2, \quad \mathbf{m}_{\mathcal{G}_2} := v$$

we define

$$\forall f \in \mathcal{IP}_{\mathcal{G}_2}, \quad \|f\|_{\mathcal{G}_2} = \|f\|_{\dot{H}^{-s}(\mathbb{R}^d)} := \left\| \frac{\hat{f}(\xi)}{|\xi|^s} \right\|_{L^2}.$$

2.5.6. *Comparison of distances when $E = \mathbb{R}^d$.* All the previous distances are *Hölder equivalent on bounded sets* in the sense of Definition 2.6. Precise quantitative statements of these equivalences are given in Lemma 4.1 in Section 4.1.

2.6. **Differential calculus in probability spaces.** We start with a purely metric definition in the case of usual Hölder regularity.

Definition 2.9. Given some metric spaces $\tilde{\mathcal{G}}_1$ and $\tilde{\mathcal{G}}_2$, some weight function

$$\Lambda : \tilde{\mathcal{G}}_1 \mapsto [1, +\infty),$$

we denote by

$$UC_\Lambda(\tilde{\mathcal{G}}_1, \tilde{\mathcal{G}}_2)$$

the weighted space of uniformly continuous functions from $\tilde{\mathcal{G}}_1$ to $\tilde{\mathcal{G}}_2$, that is the set of functions $\mathcal{S} : \tilde{\mathcal{G}}_1 \rightarrow \tilde{\mathcal{G}}_2$ such that there exists a modulus of continuity ω so that

$$(2.4) \quad \forall f_1, f_2 \in \tilde{\mathcal{G}}_1 \quad d_{\mathcal{G}_2}(\mathcal{S}(f_1), \mathcal{S}(f_2)) \leq \Lambda(f_1, f_2) \omega(d_{\mathcal{G}_1}(f_1, f_2)),$$

with

$$\Lambda(f_1, f_2) := \max\{\Lambda(f_1), \Lambda(f_2)\}$$

and where $d_{\mathcal{G}_k}$ denotes the metric of $\tilde{\mathcal{G}}_k$. Note that the tilde sign in the notation of the distance has been removed in order to present unified notation with the next definition.

For any $\eta \in (0, 1]$, we denote by

$$C_{\Lambda}^{0,\eta}(\tilde{\mathcal{G}}_1, \tilde{\mathcal{G}}_2)$$

the weighted space of functions from $\tilde{\mathcal{G}}_1$ to $\tilde{\mathcal{G}}_2$ with η -Hölder regularity, that are the uniformly continuous functions for which the modulus of continuity satisfies $\omega(s) \leq C s^\eta$ for some constant $C > 0$. We then define the *semi-norm*

$$[S]_{C_{\Lambda}^{0,\eta}(\tilde{\mathcal{G}}_1, \tilde{\mathcal{G}}_2)} \quad \text{for } S \in C_{\Lambda}^{0,\eta}(\tilde{\mathcal{G}}_1, \tilde{\mathcal{G}}_2)$$

as the infimum of the constants $C > 0$ such that (2.4) holds with $\omega(s) = C s^\eta$.

We now define a first order differential calculus, for which we require a norm structure on the functional spaces.

Definition 2.10. Given some *Banach* spaces $\mathcal{G}_1, \mathcal{G}_2$ and some *metric* sets $\tilde{\mathcal{G}}_1, \tilde{\mathcal{G}}_2$ such that

$$\tilde{\mathcal{G}}_i - \tilde{\mathcal{G}}_i \subset \mathcal{G}_i, \quad i = 1, 2,$$

some weight function

$$\Lambda : \tilde{\mathcal{G}}_1 \mapsto [1, \infty),$$

we define

$$UC_{\Lambda}^1(\tilde{\mathcal{G}}_1, \mathcal{G}_1; \tilde{\mathcal{G}}_2, \mathcal{G}_2)$$

(later simply denoted by $UC_{\Lambda}^1(\tilde{\mathcal{G}}_1; \tilde{\mathcal{G}}_2)$), the space of continuously differentiable functions from $\tilde{\mathcal{G}}_1$ to $\tilde{\mathcal{G}}_2$, whose derivative satisfies some weighted uniform continuity.

More explicitly, this is the set of uniformly continuous functions

$$\mathcal{S} : \tilde{\mathcal{G}}_1 \rightarrow \tilde{\mathcal{G}}_2$$

such that there exists a uniformly continuous function

$$D\mathcal{S} : \tilde{\mathcal{G}}_1 \rightarrow \mathcal{L}(\mathcal{G}_1, \mathcal{G}_2)$$

(where $\mathcal{L}(\mathcal{G}_1, \mathcal{G}_2)$ denotes the space of bounded linear applications from \mathcal{G}_1 to \mathcal{G}_2 endowed with the usual operator norm), some constants $C_i > 0$, $i = 1, 2$, and some modulus of differentiability Ω (that is a function $\Omega : \mathbb{R}_+ \rightarrow \mathbb{R}_+$ such that $\Omega(s)/s \rightarrow 0$ when $s \rightarrow 0$) so that for any $f_1, f_2 \in \tilde{\mathcal{G}}_1$:

$$(2.5) \quad \|\mathcal{S}(f_2) - \mathcal{S}(f_1)\|_{\mathcal{G}_2} \leq C_1 \Lambda(f_1, f_2) \|f_2 - f_1\|_{\mathcal{G}_1}$$

$$(2.6) \quad \|\langle D\mathcal{S}[f_1], f_2 - f_1 \rangle\|_{\mathcal{G}_2} \leq C_2 \Lambda(f_1, f_2) \|f_2 - f_1\|_{\mathcal{G}_1}$$

$$(2.7) \quad \|\mathcal{S}(f_2) - \mathcal{S}(f_1) - \langle D\mathcal{S}[f_1], f_2 - f_1 \rangle\|_{\mathcal{G}_2} \leq C_3 \Lambda(f_1, f_2) \Omega(\|f_2 - f_1\|_{\mathcal{G}_1}).$$

For any $\eta \in (0, 1]$, we also denote by

$$C_{\Lambda}^{1,\eta}(\tilde{\mathcal{G}}_1, \mathcal{G}_1; \tilde{\mathcal{G}}_2, \mathcal{G}_2)$$

(later simply denoted by $C_{\Lambda}^{1,\eta}(\tilde{\mathcal{G}}_1; \tilde{\mathcal{G}}_2)$), the space of continuously differentiable functions from $\tilde{\mathcal{G}}_1$ to $\tilde{\mathcal{G}}_2$, whose derivative satisfies some weighted η -Hölder regularity, which is the set of continuously differentiable functions such that the modulus of differentiability satisfies $\Omega(s) \leq C s^{1+\eta}$ for some constant $C > 0$. We define respectively $C_1^{\mathcal{S}}, C_2^{\mathcal{S}}, C_3^{\mathcal{S}}$, as the infimum of the constant $C_1, C_2, C_3 > 0$ such that respectively (2.5), (2.6), (2.7) with $\Omega(s) = C_3 s^{1+\eta}$ hold. We then define the *semi-norms*

$$[\mathcal{S}]_{C_{\Lambda}^{0,1}} := C_1^{\mathcal{S}}, \quad [\mathcal{S}]_{C_{\Lambda}^{1,0}} := C_2^{\mathcal{S}}, \quad [\mathcal{S}]_{C_{\Lambda}^{1,\eta}} := C_3^{\mathcal{S}}, \quad \|\mathcal{S}\|_{C_{\Lambda}^{1,\eta}} := C_1^{\mathcal{S}} + C_2^{\mathcal{S}} + C_3^{\mathcal{S}}.$$

In the sequel we omit the subscript Λ or we replace it by the subscript b in the case when $\Lambda \equiv 1$.

Remark 2.11. In the sequel, we shall apply this abstract differential calculus with some suitable subspaces $\tilde{\mathcal{G}}_i \subset P(E)$. This choice of subspaces is crucial in order to make rigorous the intuition of Grünbaum [36] (see the — unjustified — expansion of H_f in [36]). It is worth emphasizing that our differential calculus is based on the idea of considering $P(E)$ (or subsets of $P(E)$) as “plunged sub-manifolds” of some larger normed spaces \mathcal{G}_i . Our approach thus differs from the approach of P.-L. Lions recently developed in his course at Collège de France [46] or the one developed by L. Ambrosio et al in order to deal with gradient flows in probability measures spaces, see for instance [2]. We develop a differential calculus in probability measures spaces into a simple and robust framework, well suited to deal with the different objects we have to manipulate (1-particle semigroup, polynomial, generators...). And one of the main innovations of our work is the use of this differential calculus to state some “differential” stability conditions on the limiting semigroup. Roughly speaking the latter estimates measure how this limiting semigroup handles fluctuations departing from chaoticity. They are the corner stone of our analysis.

This differential calculus behaves well for composition in the sense that for any given

$$\mathcal{U} \in C_{\Lambda_{\mathcal{U}}}^{1,\eta}(\tilde{\mathcal{G}}_1, \tilde{\mathcal{G}}_2) \quad \text{and} \quad \mathcal{V} \in C_{\Lambda_{\mathcal{V}}}^{1,\eta}(\tilde{\mathcal{G}}_2, \tilde{\mathcal{G}}_3)$$

there holds

$$\mathcal{S} := \mathcal{V} \circ \mathcal{U} \in C_{\Lambda_{\mathcal{S}}}^{1,\eta}(\tilde{\mathcal{G}}_1, \tilde{\mathcal{G}}_3)$$

for some appropriate weight function $\Lambda_{\mathcal{S}}$. We conclude the section by stating a precise result well adapted to our applications.

Lemma 2.12. *For any given*

$$\mathcal{U} \in C_{\Lambda}^{1,\eta}(\tilde{\mathcal{G}}_1, \tilde{\mathcal{G}}_2) \quad \text{and} \quad \mathcal{V} \in C^{1,\eta}(\tilde{\mathcal{G}}_2, \tilde{\mathcal{G}}_3)$$

there holds

$$\mathcal{S} := \mathcal{V} \circ \mathcal{U} \in C_{\Lambda^{1+\eta}}^{1,\eta}(\tilde{\mathcal{G}}_1, \tilde{\mathcal{G}}_3) \quad \text{and} \quad DS[f] = DV[\mathcal{U}(f)] \circ DU[f].$$

More precisely, there holds

$$[\mathcal{S}]_{C_{\Lambda}^{0,1}} \leq [\mathcal{V}]_{C^{0,1}} [\mathcal{U}]_{C_{\Lambda}^{0,1}}, \quad [\mathcal{S}]_{C_{\Lambda}^{1,0}} \leq [\mathcal{V}]_{C^{1,0}} [\mathcal{U}]_{C_{\Lambda}^{1,0}}$$

and

$$[\mathcal{S}]_{C_{\Lambda^{1+\eta}}^{1,\eta}} \leq [\mathcal{V}]_{C^{1,0}} [\mathcal{U}]_{C_{\Lambda}^{1,\eta}} + [\mathcal{V}]_{C^{1,\eta}} [\mathcal{U}]_{C_{\Lambda}^{0,1}}^{1+\eta}.$$

When further $\mathcal{V} \in C^{1,1}(\tilde{\mathcal{G}}_2, \tilde{\mathcal{G}}_3)$, we also have

$$\mathcal{S} := \mathcal{V} \circ \mathcal{U} \in C_{\Lambda^2}^{1,\eta}(\tilde{\mathcal{G}}_1, \tilde{\mathcal{G}}_3)$$

with

$$[\mathcal{S}]_{C_{\Lambda^2}^{1,\eta}} \leq [\mathcal{V}]_{C^{1,0}} [\mathcal{U}]_{C_{\Lambda}^{1,\eta}} + [\mathcal{V}]_{C^{1,1}} [\mathcal{U}]_{C_{\Lambda}^{0,(1+\eta)/2}}^2.$$

Proof of Lemma 2.12. For $f_1, f_2 \in \tilde{\mathcal{G}}_1$ we have

$$\mathcal{U}(f_2) = \mathcal{U}(f_1) + \langle DU[f_1], (f_2 - f_1) \rangle + \mathcal{R}_{\mathcal{U}}(f_1, f_2)$$

with

$$\|\langle DU[\rho_1], (f_2 - f_1) \rangle\|_{\mathcal{G}_2} \leq [\mathcal{U}]_{C_{\Lambda}^{1,0}} \Lambda(f_1, f_2) \|f_2 - f_1\|_{\mathcal{G}_1}$$

and

$$\|\mathcal{R}_{\mathcal{U}}(f_1, f_2)\|_{\mathcal{G}_2} \leq [\mathcal{U}]_{C_{\Lambda}^{1,\eta}} \Lambda(f_1, f_2) \|f_2 - f_1\|_{\mathcal{G}_1}^{1+\theta},$$

and a similar Taylor expansion holds for \mathcal{V} : for $g_1, g_2 \in \tilde{\mathcal{G}}_2$,

$$\mathcal{V}(g_2) = \mathcal{V}(g_1) + \langle DV[g_1], (g_2 - g_1) \rangle + \mathcal{R}_{\mathcal{V}}(g_1, g_2)$$

with

$$\|\langle DV[g_1], (g_2 - g_1) \rangle\|_{\mathcal{G}_3} \leq [\mathcal{V}]_{C^{1,0}} \|g_2 - g_1\|_{\mathcal{G}_2}$$

and

$$\|\mathcal{R}_\mathcal{V}(g_1, g_2)\|_{\mathcal{G}_3} \leq [\mathcal{V}]_{C^{1,\eta}} \|g_2 - g_1\|_{\mathcal{G}_2}^{1+\eta}.$$

We then write

$$\|\mathcal{S}(f_2) - \mathcal{S}(f_1)\|_{\mathcal{G}_3} \leq [\mathcal{V}]_{C^{0,1}} \|\mathcal{U}(f_2) - \mathcal{U}(f_1)\|_{\mathcal{G}_2} \leq [\mathcal{V}]_{C^{0,1}} [\mathcal{U}]_{C_\Lambda^{0,1}} \Lambda(f_1, f_2) \|f_2 - f_1\|_{\mathcal{G}_1}$$

which implies

$$[\mathcal{S}]_{C_\Lambda^{0,1}} \leq [\mathcal{V}]_{C^{0,1}} [\mathcal{U}]_{C_\Lambda^{0,1}}$$

and

$$\begin{aligned} \mathcal{S}(f_2) &= (\mathcal{V} \circ \mathcal{U})(f_2) = \mathcal{V}\left(\mathcal{U}(f_1) + \langle D\mathcal{U}[f_1], (f_2 - f_1) \rangle + \mathcal{R}_\mathcal{U}(f_1, f_2)\right) \\ &= \mathcal{V}(\mathcal{U}(f_1)) + \mathcal{R}_\mathcal{V}(\mathcal{U}(f_2), \mathcal{U}(f_1)) \\ &\quad + \left\langle D\mathcal{V}[\mathcal{U}(f_1)], \left(\langle D\mathcal{U}[f_1], (f_2 - f_1) \rangle + \mathcal{R}_\mathcal{U}(f_1, f_2)\right) \right\rangle. \end{aligned}$$

which implies

$$\langle D\mathcal{S}[f_1], (f_2 - f_1) \rangle = \left\langle D\mathcal{V}[\mathcal{U}(f_1)], \left(\langle D\mathcal{U}[f_1], (f_2 - f_1) \rangle\right) \right\rangle,$$

and

$$[\mathcal{V} \circ \mathcal{U}]_{C_\Lambda^{1,0}} \leq [\mathcal{V}]_{C^{1,0}} [\mathcal{U}]_{C_\Lambda^{1,\eta}}.$$

Finally we estimate the remaining term:

$$\begin{aligned} &|\mathcal{S}(f_2) - \mathcal{S}(f_1) - \langle D\mathcal{S}[f_1], (f_2 - f_1) \rangle| \\ &= |\mathcal{R}_\mathcal{V}(\mathcal{U}(f_2), \mathcal{U}(f_1)) + \langle D\mathcal{V}[\mathcal{U}(f_1)], \mathcal{R}_\mathcal{U}(f_1, f_2) \rangle| \\ &\leq \left(\Lambda(f_1, f_2)^{1+\eta} [\mathcal{V}]_{C^{1,\eta}} [\mathcal{U}]_{C_\Lambda^{0,1}}^{1+\eta} + \Lambda(f_1, f_2) [\mathcal{V}]_{C^{1,0}} [\mathcal{U}]_{C_\Lambda^{1,\eta}} \right) \|f_2 - f_1\|_{\mathcal{G}_1}^{1+\eta}. \end{aligned}$$

We hence conclude that

$$[\mathcal{S}]_{C_{\Lambda^{1+\eta}}^{1,\eta}} \leq [\mathcal{V}]_{C^{1,0}} [\mathcal{U}]_{C_\Lambda^{1,\eta}} + [\mathcal{V}]_{C^{1,\eta}} [\mathcal{U}]_{C_\Lambda^{0,1}}^{1+\eta}.$$

The variant

$$[\mathcal{S}]_{C_{\Lambda^2}^{1,\eta}} \leq [\mathcal{V}]_{C^{1,0}} [\mathcal{U}]_{C_\Lambda^{1,\eta}} + [\mathcal{V}]_{C^{1,1}} [\mathcal{U}]_{C_\Lambda^{0,(1+\eta)/2}}^2.$$

is easily obtained by estimating instead

$$\begin{aligned} |\mathcal{R}_\mathcal{V}(\mathcal{U}(f_2), \mathcal{U}(f_1))| &\leq [\mathcal{V}]_{C^{1,1}} \|\mathcal{U}(f_2) - \mathcal{U}(f_1)\|_{\mathcal{G}_2}^2 \\ &\leq [\mathcal{V}]_{C^{1,1}} \Lambda(f_1, f_2)^2 [\mathcal{U}]_{C_\Lambda^{0,(1+\eta)/2}} \|f_2 - f_1\|_{\mathcal{G}_1}^{1+\eta}. \end{aligned}$$

□

2.7. The pushforward generator. As a first application of this differential calculus, let us compute the generator of the pushforward limiting semigroup.

Lemma 2.13. *Given some Banach space \mathcal{G} and some probability space $P_{\mathcal{G}}(E)$ (see Definitions 2.4-2.5) associated to a weight function m and constraint function \mathbf{m} , and endowed with the metric induced from \mathcal{G} , then for some $\delta \in (0, 1]$ and some $\bar{a} \in (0, \infty)$ we assume that for any $a \in (\bar{a}, \infty)$:*

(i) *The equation (2.1) generates a semigroup*

$$S_t^{NL} : \mathcal{B}P_{\mathcal{G},a} \rightarrow \mathcal{B}P_{\mathcal{G},a}$$

which is δ -Hölder continuous locally uniformly in time, in the sense that for any $\tau \in (0, \infty)$ there exists $C_\tau \in (0, \infty)$ such that

$$\forall f, g \in \mathcal{B}P_{\mathcal{G},a}, \quad \sup_{t \in [0, \tau]} \|S_t^{NL} f - S_t^{NL} g\|_{\mathcal{G}_1} \leq C_\tau \|f - g\|_{\mathcal{G}_1}^\delta.$$

(ii) The application Q is bounded and δ -Hölder continuous from $\mathcal{BP}_{\mathcal{G},a}$ into \mathcal{G} .
Then for any $a \in (\bar{a}, \infty)$ the pushforward semigroup T_t^∞ defined by

$$\forall f \in \mathcal{BP}_{\mathcal{G},a}(E), \quad \Phi \in UC_b(\mathcal{BP}_{\mathcal{G},a}(E)), \quad T_t^\infty[\Phi](f) := \Phi(S_t^{NL}(f))$$

is a C_0 -semigroup of contractions on the Banach space $UC_b(\mathcal{BP}_{\mathcal{G},a}(E))$.

Its generator G^∞ is an unbounded linear operator on $UC_b(\mathcal{BP}_{\mathcal{G},a}(E))$ with domain $Dom(G^\infty)$ containing $UC_b^1(\mathcal{BP}_{\mathcal{G},a}(E))$. On the latter space, it is defined by the formula

$$(2.8) \quad \forall \Phi \in UC_b^1(\mathcal{BP}_{\mathcal{G},a}(E)), \quad \forall f \in \mathcal{BP}_{\mathcal{G},a}(E), \quad (G^\infty \Phi)(f) := \langle D\Phi[f], Q(f) \rangle.$$

Remark 2.14. Note that the restriction to uniformly continuous functions Φ on probability spaces will be harmless in the sequel for two reasons: first in most cases our choice of weight, constraints and distance yields a compact space $\mathcal{BP}_{\mathcal{G},a}(E)$, and second and most importantly we shall only manipulate this pushforward semigroup for functions Φ having at least Hölder regularity.

Proof of Lemma 2.13. The proof is split in several steps.

Step 1. We claim that for any $f_0 \in \mathcal{BP}_{\mathcal{G},a}(E)$ and $\tau > 0$ the application

$$\mathcal{S}(f_0) : [0, \tau] \rightarrow P_{\mathcal{G}}, \quad t \mapsto \mathcal{S}_t^{NL}(f_0)$$

is right differentiable in $t = 0$ with

$$\mathcal{S}(f_0)'(0^+) = Q(f_0).$$

Denote $f_t := S_t^{NL}f_0$. First, since $f_t \in \mathcal{BP}_{\mathcal{G},a}$ for any $t \in [0, \tau]$ and Q is bounded on $\mathcal{BP}_{\mathcal{G},a}(E)$ (assumption (ii)), we deduce that uniformly on $f_0 \in \mathcal{BP}_{\mathcal{G},a}(E)$ we have

$$(2.9) \quad \|f_t - f_0\|_{\mathcal{G}} = \left\| \int_0^t Q(f_s) ds \right\|_{\mathcal{G}} \leq K t.$$

We then use the previous inequality together with the fact that Q is δ -Hölder continuous (assumption (ii) again), to get

$$\begin{aligned} \|f_t - f_0 - tQ(f_0)\|_{\mathcal{G}} &= \left\| \int_0^t (Q(f_s) - Q(f_0)) ds \right\|_{\mathcal{G}} \\ &= L \int_0^t \|f_s - f_0\|_{\mathcal{G}}^\delta ds \\ &\leq L \int_0^t (Ks)^\delta ds = L K^\delta \frac{t^{1+\delta}}{1+\delta}, \end{aligned}$$

which implies the claim.

Then the semigroup property of (S_t^{NL}) implies that $t \mapsto f_t$ is continuous from \mathbb{R}_+ into $P_{\mathcal{G}}(E)$ and right differentiable at any point.

Step 2. We claim that (T_t^∞) is a C_0 -semigroup of contractions on $UC_b(\mathcal{BP}_{\mathcal{G},a}(E))$.

First for any $\Phi \in UC_b(\mathcal{BP}_{\mathcal{G},a}(E))$, we denote by ω_Φ the modulus of continuity of Φ . We have thanks to the assumption (i):

$$\begin{aligned} \forall t \in [0, \tau], \quad |(T_t^\infty \Phi)(g) - (T_t^\infty \Phi)(f)| &= |\Phi(S_t^{NL}(g)) - \Phi(S_t^{NL}(f))| \\ &\leq \omega_\Phi \left(\|S_t^{NL}(g) - S_t^{NL}(f)\|_{\mathcal{G}_1} \right) \\ &\leq \omega_\Phi \left(C_\tau \|g - f\|_{\mathcal{G}_1}^\delta \right), \end{aligned}$$

so that $T_t^\infty \Phi \in UC_b(\mathcal{BP}_{\mathcal{G},a}(E))$ for any $t \in [0, \tau]$, and then, by iteration, for any $t \geq 0$. Next, we have

$$\|T_t^\infty\| = \sup_{\|\Phi\| \leq 1} \|T_t^\infty \Phi\| = \sup_{\|\Phi\| \leq 1} \sup_{f \in \mathcal{BP}_{\mathcal{G},a}(E)} |\Phi(S_t^{NL}(f))| \leq 1$$

since

$$\|\Phi\| = \sup_{f \in \mathcal{BP}_{\mathcal{G},a}(E)} |\Phi(f)|.$$

Finally, from (2.9), for any $\Phi \in UC_b(\mathcal{BP}_{\mathcal{G},a}(E))$, we have

$$\|T_t^\infty \Phi - \Phi\| = \sup_{f \in \mathcal{BP}_{\mathcal{G},a}(E)} |\Phi(S_t^{NL}(f)) - \Phi(f)| \leq \omega_\Phi(Kt) \rightarrow 0.$$

As a consequence (T_t^∞) has a closed generator G^∞ with dense domain

$$\text{Dom}(G^\infty) \subset UC_b(\mathcal{BP}_{\mathcal{G},a}(E))$$

(see for instance [62, Chapter 1, Corollary 2.5]).

Step 3. We shall now identify this generator, at least on a part of its domain. Let us construct a natural candidate provided by the heuristic of Remark 2.3. Let us define $\tilde{G}^\infty \Phi$ by

$$\forall \Phi \in UC_b^1(\mathcal{BP}_{\mathcal{G},a}(E)), \forall f \in \mathcal{BP}_{\mathcal{G},a}(E), \quad (\tilde{G}^\infty \Phi)(f) := \langle D\Phi[f], Q(f) \rangle.$$

The right hand side is well defined since

$$D\Phi[f] \in \mathcal{L}(\mathcal{G}, \mathbb{R}) = \mathcal{G}' \quad \text{and} \quad Q(f) \in \mathcal{G}.$$

Moreover, since both applications

$$f \mapsto D\Phi[f] \quad \text{and} \quad f \mapsto Q(f)$$

are uniformly continuous on $\mathcal{BP}_{\mathcal{G},a}(E)$, so is the application

$$f \mapsto (\tilde{G}^\infty \Phi)(f).$$

Hence $\tilde{G}^\infty \Phi \in UC_b(\mathcal{BP}_{\mathcal{G},a}(E))$.

Step 4. Finally, by composition, for any fixed $\Phi \in UC_b^1(\mathcal{BP}_{\mathcal{G},a}(E))$ and $f \in \mathcal{BP}_{\mathcal{G},a}(E)$, the map

$$t \mapsto T_t^\infty \Phi(\rho) = \Phi \circ S_t^{NL}(\rho)$$

is right differentiable in $t = 0$ and

$$\begin{aligned} \frac{d}{dt}(T_t^\infty \Phi)(\rho)|_{t=0} &:= \frac{d}{dt}(\Phi \circ \mathcal{S}(\rho)(t))|_{t=0} \\ &= \left\langle D\Phi[\mathcal{S}(\rho)(0)], \frac{d}{dt}\mathcal{S}(\rho)(0) \right\rangle \\ &= \langle D\Phi[\rho], Q(\rho) \rangle = (\tilde{G}^\infty \Phi)(\rho), \end{aligned}$$

which precisely means that $\Phi \in \text{Dom}(G^\infty)$ and that (2.8) holds. \square

2.8. Duality inequalities. Our transformations π^N and R^N behave nicely for the supremum norm on $C_b(P(E), TV)$, see (2.2). More generally we shall consider “*duality pairs*” of metric spaces as follows:

Definition 2.15. We say that a pair (\mathcal{F}, P_G) of a normed vectorial space $\mathcal{F} \subset C_b(E)$ endowed with the norm $\|\cdot\|_{\mathcal{F}}$ and a probability space $P_G \subset P(E)$ endowed with a metric d_G satisfy a *duality inequality* if

$$(2.10) \quad \forall f, g \in P_G, \forall \varphi \in \mathcal{F}, \quad |\langle g - f, \varphi \rangle| \leq C d_G(f, g) \|\varphi\|_{\mathcal{F}},$$

where here $\langle \cdot, \cdot \rangle$ stands for the usual duality brackets between probabilities and continuous functions. In the case where the distance d_G is associated with a normed vector space \mathcal{G} , this amounts to the usual duality inequality $|\langle h, \varphi \rangle| \leq \|h\|_{\mathcal{G}} \|\varphi\|_{\mathcal{F}}$.

The “compatibility” of the transformation R^N for any such pair follows from the multilinearity: if \mathcal{F} and \mathcal{G} are in duality, $\mathcal{F} \subset C_b(E)$ and P_G is endowed with the metric associated to $\|\cdot\|_{\mathcal{G}}$, then for any

$$\varphi = \varphi_1 \times \cdots \times \varphi_N \in \mathcal{F}^{\otimes N},$$

the polynomial function R_{φ}^N in $C_b(P(E))$ is $C^{1,1}(P_G, \mathbb{R})$. Indeed, given $f_1, f_2 \in P_{G_1}$, we define

$$\mathcal{G} \rightarrow \mathbb{R}, \quad h \mapsto DR_{\varphi}^{\ell}[f_1](h) := \sum_{i=1}^N \left(\prod_{j \neq i} \langle f_1, \varphi_j \rangle \right) \langle h, \varphi_i \rangle,$$

and we have

$$R_{\varphi}^N(f_2) - R_{\varphi}^N(f_1) = \sum_{i=1}^N \left(\prod_{1 \leq k < i} \langle f_2, \varphi_k \rangle \right) \langle f_2 - f_1, \varphi_i \rangle \left(\prod_{i < k \leq \ell} \langle f_1, \varphi_k \rangle \right),$$

and

$$\begin{aligned} & R_{\varphi}^N(f_2) - R_{\varphi}^N(f_1) - DR_{\varphi}^N[f_1](f_2 - f_1) = \\ &= \sum_{1 \leq j < i \leq N} \left(\prod_{1 \leq k < j} \langle f_2, \varphi_k \rangle \right) \langle f_2 - f_1, \varphi_j \rangle \left(\prod_{j < k < i} \langle f_1, \varphi_k \rangle \right) \langle f_2 - f_1, \varphi_i \rangle \left(\prod_{i < k \leq \ell} \langle f_1, \varphi_k \rangle \right). \end{aligned}$$

Hence for instance $R_{\varphi}^N \in C^{1,1}(P_G; \mathbb{R})$ with

$$\begin{aligned} |R_{\varphi}^N(f_2) - R_{\varphi}^N(f_1)| &\leq N \|\varphi\|_{\mathcal{F} \otimes (L^{\infty})^{N-1}} \|f_2 - f_1\|_{\mathcal{G}}, \\ |DR_{\varphi}^N[f_1](h)| &\leq N \|\varphi\|_{\mathcal{F} \otimes (L^{\infty})^{N-1}} \|h\|_{\mathcal{G}}, \end{aligned}$$

and

$$(2.11) \quad |R_{\varphi}^N(f_2) - R_{\varphi}^N(f_1) - DR_{\varphi}^N[f_1](f_2 - f_1)| \leq \frac{N(N-1)}{2} \|\varphi\|_{\mathcal{F}^2 \otimes (L^{\infty})^{N-2}} \|f_2 - f_1\|_{\mathcal{G}}^2,$$

where we have defined

$$\|\varphi\|_{\mathcal{F}^k \otimes (L^{\infty})^{N-k}} := \max_{i_1, \dots, i_k \text{ distincts in } \llbracket 1, N \rrbracket} \left(\|\varphi_{i_1}\|_{\mathcal{F}} \cdots \|\varphi_{i_k}\|_{\mathcal{F}} \prod_{j \neq (i_1, \dots, i_k)} \|\varphi_j\|_{L^{\infty}(E)} \right).$$

Remarks 2.16. (1) It is easily seen in this computation that the assumption that φ is tensor product is not necessary. In fact it is likely that this assumption could be relaxed all along our proof. We do not pursue this line of research.

(2) The assumption $\mathcal{F} \subset C_b(E)$ could also be relaxed. For instance, when

$$\mathcal{F} := \text{Lip}_0(E)$$

is the space of Lipschitz function which vanishes in some fixed point $x_0 \in E$, \mathcal{G} is its dual space, and

$$P_{\mathcal{G}} := \{f \in P_1(E); \langle f, \text{dist}_E(\cdot, x_0) \rangle \leq a\}$$

for some fixed $a > 0$, we have $R_{\varphi}^N \in C^{1,1}(P_1(E); \mathbb{R})$ with

$$[R_{\varphi}^N]_{C^{0,1}} \leq N a^{N-2} \|\varphi\|_{\mathcal{F} \otimes N}, \quad [R_{\varphi}^N]_{C^{1,1}} \leq \frac{N(N-1)}{2} a^{N-1} \|\varphi\|_{\mathcal{F} \otimes N},$$

or equivalently $R_{\varphi}^N \in C_{\Lambda}^{1,1}(P_1(E); \mathbb{R})$ with $\Lambda(f) := \|f\|_{M_1^1}^{N-2}$.

In the other way round, for the projection π^N it is clear that if the empirical measure map

$$V \in E^N \mapsto \mu_V^N \in P(E)$$

belongs to $C^{k,\eta}(E^N, P_{\mathcal{G}})$ for some norm structure \mathcal{G} , then by composition one has

$$(2.12) \quad \|\pi^N(\Phi)\|_{C^{k,\eta}(E^N; \mathbb{R})} \leq C_{\pi} \|\Phi\|_{C^{k,\eta}(P_{\mathcal{G}})}.$$

However the regularity of the empirical measure of course heavily depends on the choice of the metric \mathcal{G} .

Example 2.17. In the case $\mathcal{F} = (C_b(E), L^{\infty})$ and $\mathcal{G} = (M^1(E), TV)$, (2.12) trivially holds with $C^{k,\eta}$ replaced by C_b .

Example 2.18. When $\mathcal{F} = \text{Lip}_0(E)$ (Lipschitz function vanishing at some given point v_0) endowed with the norm $\|\phi\|_{Lip}$ and $P_{\mathcal{G}}(E)$ (constructed in Subsubsection 2.5.2) is endowed with the Wasserstein distance W_1 with linear cost, one has (2.12) with $k = 0$, $\eta = 1$:

$$|\Phi(\mu_X^N) - \Phi(\mu_Y^N)| \leq \|\Phi\|_{C^{0,1}(P_{\mathcal{G}})} W_1(\mu_X^N, \hat{\mu}_Y^N) \leq \|\Phi\|_{C^{0,1}(P_{\mathcal{G}})} \|X - Y\|_{\ell^1},$$

where we use (2.3), which proves that

$$\|\pi^N(\Phi)\|_{C^{0,1}(E^N)} \leq \|\Phi\|_{C^{0,1}(P_{\mathcal{G}})},$$

when E^N is endowed with the ℓ^1 distance defined in (2.3).

3. THE ABSTRACT THEOREM

3.1. Assumptions of the abstract theorem. Let us list the assumptions that we need for our main abstract theorem.

(A1) Assumptions on the N -particle system.

G^N and T_t^N are well defined on $C_b(E^N)$ and invariant under permutation, and they satisfy the following moment conditions:

- (i) *Propagation of integral moment bound:* There exists a weight function $m_{\mathcal{G}_1}$, a time $T \in (0, \infty]$ and a constant $C_{1,T} \in (0, \infty)$, possibly depending on T and $m_{\mathcal{G}_1}$, but not on the number of particles N , such that

$$\forall N \geq 1, \quad \sup_{0 \leq t < T} \langle f_t^N, M_{m_{\mathcal{G}_1}}^N \rangle \leq C_{m_{\mathcal{G}_1}, T}.$$

- (ii) *Support moment bound at initial time:* There exists a weight function $m_{\mathcal{G}_3}$ and a constant $C_3^N \in (0, +\infty)$, possibly depending on the number of particles N , such that

$$\text{Supp } f_0^N \subset \left\{ V \in E^N; M_{m_{\mathcal{G}_3}}^N(V) \leq C_{m_{\mathcal{G}_3}, 0}^N \right\}.$$

Note that the name of the weights $m_{\mathcal{G}_1}$ and $m_{\mathcal{G}_3}$ in **(A1)** above are chosen for coherence with the functional spaces in the other assumptions.

(A2) Assumptions for the existence of the pushforward semigroup.

Given some Banach space \mathcal{G}_1 and some probability space $P_{\mathcal{G}_1}(E)$ (see Definitions 2.4-2.5) associated to a weight function $m_{\mathcal{G}_1}$ (as in **(A1)**-(i)) and a constraint function $\mathbf{m}_{\mathcal{G}_1}$, and endowed with the metric induced from \mathcal{G}_1 , then for some $\delta \in (0, 1]$ and some $\bar{a} \in (0, \infty)$ we assume that for any $a \in (\bar{a}, \infty)$:

(i) The equation (2.1) generates a semigroup

$$S_t^{NL} : \mathcal{B}P_{\mathcal{G}_1, a} \rightarrow \mathcal{B}P_{\mathcal{G}_1, a}$$

which is δ -Hölder continuous locally uniformly in time, in the sense that for any $\tau \in (0, \infty)$ there exists $C_\tau \in (0, \infty)$ such that

$$\forall f, g \in \mathcal{B}P_{\mathcal{G}_1, a}, \quad \sup_{t \in [0, \tau]} \|S_t^{NL} f - S_t^{NL} g\|_{\mathcal{G}_1} \leq C_\tau \|f - g\|_{\mathcal{G}_1}^\delta.$$

(ii) The application Q is bounded and δ -Hölder continuous from $\mathcal{B}P_{\mathcal{G}_1, a}$ into \mathcal{G}_1 .

So in particular the semigroups S_t^N , T_t^N , S_t^{NL} and T_t^∞ are well defined as well as the generators G^N and G^∞ .

We then need the key following *consistency* assumption. It intuitively states that the N -particle *approximation* of the limiting mean-field equation is consistent. More rigorously this means a convergence of the generators of the N -particle approximation towards the generator of the limiting pushforward semigroup within the abstract functional framework we have introduced.

(A3) Convergence of the generators.

In the probability metrized set $P_{\mathcal{G}_1}$ introduced in **(A2)** (associated to the weight function $m_{\mathcal{G}_1}$ and constraint function $\mathbf{m}_{\mathcal{G}_1}$) we define

$$\mathbf{R}_{\mathcal{G}_1} := \{ \mathbf{r} \in \mathbb{R}^D; \quad \exists f \in P_{\mathcal{G}_1} \text{ s.t. } \mathbf{m}_{\mathcal{G}_1}(f) = \mathbf{r} \}.$$

We also define a weight function

$$m'_{\mathcal{G}_1} \leq C m_{\mathcal{G}_1}$$

possibly weaker than $m_{\mathcal{G}_1}$ and we define the associated weight on the distribution:

$$\Lambda_1(f) := \langle f, m'_{\mathcal{G}_1} \rangle.$$

Then we assume that for some function

$$\varepsilon(N) \rightarrow 0 \quad \text{as } N \rightarrow \infty,$$

the generators G^N and G^∞ satisfy

$$\begin{aligned} \forall \Phi \in \bigcap_{\mathbf{r} \in \mathbf{R}_{\mathcal{G}_1}} C_{\Lambda_1}^{1, \eta}(P_{\mathcal{G}_1, \mathbf{r}}; \mathbb{R}) \\ \left\| \left(M_{m_{\mathcal{G}_1}}^N \right)^{-1} (G^N \pi_N - \pi_N G^\infty) \Phi \right\|_{L^\infty(E^N)} \leq \varepsilon(N) \sup_{\mathbf{r} \in \mathbf{R}_{\mathcal{G}_1}} [\Phi]_{C_{\Lambda_1}^{1, \eta}(P_{\mathcal{G}_1, \mathbf{r}})}, \end{aligned}$$

where $M_{m_{\mathcal{G}_1}}^N$ is defined by

$$M_{m_{\mathcal{G}_1}}^N := \frac{1}{N} \sum_{i=1}^N m_{\mathcal{G}_1}(v_i).$$

Note the following aspect, which shall be a crucial source of difficulty in the application of the following abstract theorem: **the loss of weight in the consistency estimate (A3) has to be matched by the moment bounds propagated on the N -particle system in the assumption (A1)-(i).** (In fact the loss of weight in the consistency estimate $m_{\mathcal{G}_1}$ can be even slightly higher than the one from the stability estimate Λ_1 , see the above relation.)

More specifically, the best we are able to prove uniformly in N on the N -particle system are *polynomial* moment bounds. This thus constraints the kind of loss of weight we can afford in the following stability estimate.

We now state the second key *stability* assumption. Intuitively this corresponds to the abstract regularity that needs to be transported along the flow of the limiting mean-field equation so that the fluctuations around chaoticity can be controlled. More rigorously this means some differential regularity on the pushforward limiting semigroup, which corresponds to some differential regularity on the limiting nonlinear semigroup *according to the initial data and in probability spaces.*

(A4) Differential stability of the limiting semigroup.

We consider some Banach space $\mathcal{G}_2 \supset \mathcal{G}_1$ (where \mathcal{G}_1 was defined in (A2)) and the corresponding probability space $P_{\mathcal{G}_2}(E)$ (see Definitions 2.4-2.5) with the weight function $m_{\mathcal{G}_2}$ and the constraint function $\mathbf{m}_{\mathcal{G}_2}$, and endowed with the metric induced from \mathcal{G}_2 .

We assume that the flow S_t^{NL} is $C_{\Lambda_2}^{1,\eta}(P_{\mathcal{G}_1,\mathbf{r}}, P_{\mathcal{G}_2})$ for any $\mathbf{r} \in \mathbf{R}_{\mathcal{G}_1}$ in the sense that there exists $C_T^\infty > 0$ such that

$$\sup_{\mathbf{r} \in \mathbf{R}_{\mathcal{G}_1}} \int_0^T \left([S_t^{NL}]_{C_{\Lambda_2}^{1,\eta}(P_{\mathcal{G}_1,\mathbf{r}}, P_{\mathcal{G}_2})} + [S_t^{NL}]_{C_{\Lambda_2}^{0,\eta''}(P_{\mathcal{G}_1,\mathbf{r}}, P_{\mathcal{G}_2})}^{1+\eta'} \right) dt \leq C_T^\infty,$$

where $\eta \in (0, 1)$ is the same as in (A3), where

$$(\eta', \eta'') = (\eta, 1) \quad \text{or} \quad (\eta', \eta'') = \left(1, \frac{1+\eta}{2} \right),$$

and with the weight (recall that Λ_1 was defined in (A3))

$$\Lambda_2 = \Lambda_1^{\frac{1}{1+\eta'}}.$$

We finally state a weaker stability assumption on the limiting semigroup. It shall be used intuitively for proving that initial error in the law of large number when approximating a probability by empirical measures is propagated by the limiting semigroup. The reason for dissociating this assumption from the previous one is because we need some flexibility in the independent choice of distances for these two assumptions.

(A5) Weak stability of the limiting semigroup.

We assume that, for some probabilistic space $P_{\mathcal{G}_3}(E)$ associated to the weight function $m_{\mathcal{G}_3}$ (as in (A1)-(ii)), a constraint function $\mathbf{m}_{\mathcal{G}_3}$ and some

metric structure $d_{\mathcal{G}_3}$, for any $a, T > 0$ there exists a concave modulus of continuity $\Theta_{a,T}$ such that we have

$$\begin{aligned} & \forall f_1, f_2 \in \mathcal{BP}_{\mathcal{G}_3, a}(E) \\ & \sup_{[0, T]} d_{\mathcal{G}_3} (S_t^{NL}(f_1), S_t^{NL}(f_2)) \leq \Theta_{a, T} (d_{\mathcal{G}_3} (f_1, f_2)). \end{aligned}$$

Observe that in the latter assumption, we require that $f_1, f_2 \in \mathcal{BP}_{\mathcal{G}_3, a}(E)$ which implies in particular that we require the bounds

$$M_{m_{\mathcal{G}_3}}(f_1) = \int_{\mathbb{R}^d} f_1 m_{\mathcal{G}_3} dv \leq a, \quad M_{m_{\mathcal{G}_3}}(f_2) = \int_{\mathbb{R}^d} f_2 m_{\mathcal{G}_3} dv \leq a.$$

When applying the assumption to some empirical measure for one of the argument f_1 or f_2 , this shall induces the requirement of controlling pointwise terms like $M_{m_{\mathcal{G}_3}}^N$. This is the reason for the assumption **(A1)-(ii)**.

3.2. Statement of the result. We are now in position to state the main abstract result. This result can be considered intuitively as a *convergence* in approximation theory, in the sense of proving that approximation errors between the N -particle system and the limiting mean-field system are propagated along time without instability amplification mechanism. More specifically the approximation error means in the present context some kind of distance between the discrete N -particle system and the limiting mean-field system, within our abstract functional framework. This result implies in particular the propagation of chaos.

Theorem 3.1. *Consider a family of N -particle initial conditions*

$$f_0^N \in P_{\text{sym}}(E^N), \quad N \geq 1$$

and the associated solutions

$$f_t^N = S_t^N (f_0^N).$$

Consider a 1-particle initial condition $f_0 \in P(E)$ and the associated solution

$$f_t = S_t^{NL} (f_0)$$

of the limiting mean-field equation.

*Assume that the assumptions **(A1)-(A2)-(A3)-(A4)-(A5)** hold for some spaces $P_{\mathcal{G}_k}$, \mathcal{G}_k and \mathcal{F}_k , $k = 1, 2, 3$ with $\mathcal{F}_k \subset C_b(E)$, and where \mathcal{F}_k and \mathcal{G}_k are in duality (that is (2.10) holds).*

Assume also that the 1-particle distribution satisfies the moment bound

$$M_{m_{\mathcal{G}_3}}(f_0) = \langle f_0, m_{\mathcal{G}_3} \rangle < +\infty.$$

Then there is an explicit absolute constant $C \in (0, \infty)$ such that for any $N, \ell \in \mathbb{N}^$, with $N \geq 2\ell$, and for any*

$$\varphi = \varphi_1 \otimes \varphi_2 \otimes \cdots \otimes \varphi_\ell \in (\mathcal{F}_1 \cap \mathcal{F}_2 \cap \mathcal{F}_3)^{\otimes \ell}$$

we have

$$\begin{aligned} (3.1) \quad & \sup_{[0, T]} \left| \left\langle \left(S_t^N (f_0^N) - (S_t^{NL} (f_0))^{\otimes N} \right), \varphi \right\rangle \right| \\ & \leq C \left[\ell^2 \frac{\|\varphi\|_\infty}{N} + C_{m_{\mathcal{G}_1}, T} C_T^\infty \varepsilon(N) \ell^2 \|\varphi\|_{\mathcal{F}_2^2 \otimes (L^\infty)^{\ell-2}} \right. \\ & \quad \left. + \ell \|\varphi\|_{\mathcal{F}_3 \otimes (L^\infty)^{\ell-1}} \Theta_{a, T} \left(\mathcal{W}_{d_{\mathcal{G}_3}}(\pi_P^N f_0^N, \delta_{f_0}) \right) \right], \end{aligned}$$

where $a^N > 0$ depends on $C_{\mathcal{G}_3,0}^N$ and $M_{\mathcal{G}_3}(f_0)$, and where $\mathcal{W}_{d_{\mathcal{G}_3}}$ stands for an abstract Monge-Kantorovich distance in $P(P_{\mathcal{G}_3}(E))$ (see the third point in the next remark)

$$(3.2) \quad \mathcal{W}_{d_{\mathcal{G}_3}}(\pi_P^N f_0^N, \delta_{f_0}) = \int_{E^N} d_{\mathcal{G}_3}(\mu_V^N, f_0) f_0^N(dV).$$

Remarks 3.2. (1) The spirit of this method is to first treat the N -particle system as a perturbation (in a very degenerated sense) of the limiting problem, and second to minimize assumptions on the many-particle systems in order to avoid complications of many dimensions dynamics.

(2) In the applications the worst decay rate in the right-hand side of (3.1) is always the last one. This last term controls two kind of errors: (1) the chaoticity of the initial data, that is how well $f_0^N \sim f_0^{\otimes N}$, (2) the rate of convergence in the law of large numbers for measures in the distance $d_{\mathcal{G}_3}$.

(3) Let us discuss more the meaning of this last term and the related issue of sampling by empirical measures in statistics (see also Section 4).

Following the abstract definition of the optimal transport Wasserstein distance we define

$$\forall \mu_1, \mu_2 \in P(P_{\mathcal{G}_3}), \quad \mathcal{W}_{d_{\mathcal{G}_3}}(\mu_1, \mu_2) = \inf_{\pi \in \Pi(\mu_1, \mu_2)} \int_{P_{\mathcal{G}_3} \times P_{\mathcal{G}_3}} d_{\mathcal{G}_3}(\rho_1, \rho_2) \pi(d\rho_1, d\rho_2),$$

where $\Pi(\mu_1, \mu_2)$ denotes the probability measures on the product space $P_{\mathcal{G}_3} \times P_{\mathcal{G}_3}$ with first marginal μ_1 and second marginal μ_2 . In the case when $\mu_2 = \delta_{f_0}$ then

$$\Pi(\mu_1, \delta_{f_0}) = \{\mu_1 \otimes \delta_{f_0}\}$$

has only one element, and therefore

$$\begin{aligned} \mathcal{W}_{d_{\mathcal{G}_3}}(\pi_P^N f_0^N, \delta_{f_0}) &= \inf_{\pi \in \Pi(\pi_P^N f_0^N, \delta_{f_0})} \int_{P_{\mathcal{G}_3} \times P_{\mathcal{G}_3}} d_{\mathcal{G}_3}(\rho_1, \rho_2) \pi(d\rho_1, d\rho_2) \\ &= \int_{E^N} d_{\mathcal{G}_3}(\mu_V^N, f_0) f_0^N(dV). \end{aligned}$$

which explains the notation (3.2). We simply denote in the tensorized case:

$$\mathcal{W}_{d_{\mathcal{G}_3}}^N(f_0) := \mathcal{W}_{d_{\mathcal{G}_3}}(\pi_P^N f_0^{\otimes N}, \delta_{f_0})$$

Comparisons of the \mathcal{W}_d^N functionals and estimates on the rate

$$\mathcal{W}_d^N(f) \rightarrow 0 \quad \text{as } N \rightarrow \infty$$

depending on the choice of the distance d are discussed in Subsection 4.2.

3.3. Proof of the abstract theorem. For a given function

$$\varphi \in (\mathcal{F}_1 \cap \mathcal{F}_2 \cap \mathcal{F}_3)^{\otimes \ell},$$

we break up the term to be estimated into three parts:

$$\begin{aligned} & \left| \left\langle \left(S_t^N(f_0^N) - (S_t^{NL}(f_0))^{\otimes N} \right), \varphi \otimes 1^{\otimes N-\ell} \right\rangle \right| \leq \\ & \leq \left| \left\langle S_t^N(f_0^N), \varphi \otimes 1^{\otimes N-\ell} \right\rangle - \left\langle S_t^N(f_0^N), R_\varphi^\ell \circ \mu_V^N \right\rangle \right| \quad (=:\mathcal{T}_1) \\ & + \left| \left\langle f_0^N, T_t^N(R_\varphi^\ell \circ \mu_V^N) \right\rangle - \left\langle f_0^N, (T_t^\infty R_\varphi^\ell) \circ \mu_V^N \right\rangle \right| \quad (=:\mathcal{T}_2) \\ & + \left| \left\langle f_0^N, (T_t^\infty R_\varphi^\ell) \circ \mu_V^N \right\rangle - \left\langle (S_t^{NL}(f_0))^{\otimes \ell}, \varphi \right\rangle \right| \quad (=:\mathcal{T}_3). \end{aligned}$$

We deal separately with each part step by step:

- \mathcal{T}_1 is controlled by a purely combinatorial arguments introduced in [36]. Roughly speaking it is the combinatorial price we have to pay when we use the injection π_E^N based on empirical measures.
- \mathcal{T}_2 is controlled thanks to the consistency estimate **(A3)** on the generators, the differential stability assumption **(A4)** on the limiting semigroup and the propagation of integral moment bounds **(A1)-(i)**.
- \mathcal{T}_3 is controlled in terms of the chaoticity of the initial data thanks to the weak stability assumption **(A5)** on the limiting semigroup and the support moment bounds at initial time **(A1)-(ii)**.

Step 1: Estimate of the first term \mathcal{T}_1 . Let us prove that for any $t \geq 0$ and any $N \geq 2\ell$ there holds

$$(3.3) \quad \mathcal{T}_1 := \left| \left\langle S_t^N(f_0^N), \varphi \otimes 1^{\otimes N-\ell} \right\rangle - \left\langle S_t^N(f_0^N), R_\varphi^\ell \circ \mu_V^N \right\rangle \right| \leq \frac{2\ell^2 \|\varphi\|_{L^\infty(E^\ell)}}{N}.$$

Since $S_t^N(f_0^N)$ is a symmetric probability measure, estimate (3.3) is a direct consequence of the following lemma.

Lemma 3.3. *For any $\varphi \in C_b(E^\ell)$ we have*

$$(3.4) \quad \forall N \geq 2\ell, \quad \left| \left(\varphi \otimes 1^{\otimes N-\ell} \right)_{sym} - \pi_N R_\varphi^\ell \right| \leq \frac{2\ell^2 \|\varphi\|_{L^\infty(E^\ell)}}{N}$$

where for a function $\phi \in C_b(E^N)$, we define its symmetrized version ϕ_{sym} as:

$$\phi_{sym} = \frac{1}{|\mathfrak{S}^N|} \sum_{\sigma \in \mathfrak{S}^N} \phi_\sigma$$

where we recall that \mathfrak{S}^N is the set of N -permutations.

As a consequence for any symmetric measure we have

$$(3.5) \quad \forall f^N \in P_{sym}(E^N), \quad \left| \langle f^N, R_\varphi^\ell(\mu_V^N) \rangle - \langle f^N, \varphi \rangle \right| \leq \frac{2\ell^2 \|\varphi\|_{L^\infty(E^\ell)}}{N}.$$

Proof of Lemma 3.3. This lemma is a simple and classical combinatorial computation. We briefly sketch the proof for the sake of completeness.

For a given $\ell \leq N/2$ we introduce

$$A_{N,\ell} := \left\{ (i_1, \dots, i_\ell) \in [1, N]^\ell : \forall k \neq k' \in [1, \ell], i_k \neq i_{k'} \right\}$$

and

$$B_{N,\ell} := A_{N,\ell}^c = \left\{ (i_1, \dots, i_\ell) \in [1, N]^\ell \right\} \setminus A_{N,\ell}.$$

Since there are $N!/(N-\ell)!$ ways of choosing ℓ distinct indices among $[1, N]$ we get

$$\begin{aligned} \frac{|B_{N,\ell}|}{N^\ell} &= 1 - \frac{N!}{(N-\ell)! N^\ell} \\ &= 1 - \left(1 - \frac{1}{N}\right) \cdots \left(1 - \frac{\ell-1}{N}\right) = 1 - \exp\left(\sum_{i=0}^{\ell-1} \log\left(1 - \frac{i}{N}\right)\right) \\ &\leq 1 - \exp\left(-2 \sum_{i=0}^{\ell-1} \frac{i}{N}\right) \leq 2 \sum_{i=0}^{\ell-1} \frac{i}{N} \leq \frac{\ell^2}{N}, \end{aligned}$$

where we have used

$$\forall x \in [0, 1/2], \quad \log(1-x) \geq -2x \quad \text{and} \quad \forall x \in \mathbb{R}, \quad e^{-x} \geq 1-x.$$

Then we compute

$$\begin{aligned}
R_\varphi^\ell(\mu_V^N) &= \frac{1}{N^\ell} \sum_{i_1, \dots, i_\ell=1}^N \varphi(v_{i_1}, \dots, v_{i_\ell}) \\
&= \frac{1}{N^\ell} \sum_{(i_1, \dots, i_\ell) \in A_{N, \ell}} \varphi(v_{i_1}, \dots, v_{i_\ell}) + \frac{1}{N^\ell} \sum_{(i_1, \dots, i_\ell) \in B_{N, \ell}} \varphi(v_{i_1}, \dots, v_{i_\ell}) \\
&= \frac{1}{N^\ell} \frac{1}{(N-\ell)!} \sum_{\sigma \in \mathfrak{S}_N} \varphi(v_{\sigma(1)}, \dots, v_{\sigma(\ell)}) + \mathcal{O}\left(\frac{\ell^2}{N} \|\varphi\|_{L^\infty}\right)
\end{aligned}$$

We now use the same estimate

$$1 - \frac{N!}{(N-\ell)! N^\ell} \leq \frac{\ell^2}{N}$$

as above to get

$$\begin{aligned}
R_\varphi^\ell(\mu_V^N) &= \frac{1}{N!} \sum_{\sigma \in \mathfrak{S}_N} \varphi(v_{\sigma(1)}, \dots, v_{\sigma(\ell)}) + \mathcal{O}\left(\frac{2\ell^2}{N} \|\varphi\|_{L^\infty}\right) \\
&= \left(\varphi \otimes \mathbf{1}^{\otimes N-\ell}\right)_{\text{sym}} + \mathcal{O}\left(\frac{2\ell^2}{N} \|\varphi\|_{L^\infty}\right)
\end{aligned}$$

and the proof of (3.4) is complete.

Next for any $f^N \in P_{\text{sym}}(E^N)$ we have

$$\langle f^N, \varphi \rangle = \left\langle f^N, \left(\varphi \otimes \mathbf{1}^{\otimes N-\ell}\right)_{\text{sym}} \right\rangle,$$

and (3.5) trivially follows from (3.4). \square

Step 2: Estimate of the second term \mathcal{T}_2 . Let us prove that for any $t \in [0, T]$ and any $N \geq 2\ell$ there holds

$$\begin{aligned}
(3.6) \quad \mathcal{T}_2 &:= \left| \left\langle f_0^N, T_t^N \left(R_\varphi^\ell \circ \mu_V^N\right) \right\rangle - \left\langle f_0^N, \left(T_t^\infty R_\varphi^\ell\right) \circ \mu_V^N \right\rangle \right| \\
&\leq C_{m_{\mathcal{G}_1}, T} C_T^\infty \varepsilon(N) \ell^2 \|\varphi\|_{\mathcal{F}_2^2 \otimes (L^\infty)^{\ell-2}}.
\end{aligned}$$

We start with the following algebraic identity

$$T_t^N \pi_N - \pi_N T_t^\infty = - \int_0^t \frac{d}{ds} (T_{t-s}^N \pi_N T_s^\infty) ds = \int_0^t T_{t-s}^N [G^N \pi_N - \pi_N G^\infty] T_s^\infty ds.$$

We then use assumptions **(A1)-(i)** and **(A3)** and we get for any $t \in [0, T]$

$$\begin{aligned}
(3.7) \quad &\left| \left\langle f_0^N, T_t^N \left(R_\varphi^\ell \circ \mu_V^N\right) \right\rangle - \left\langle f_0^N, \left(T_t^\infty R_\varphi^\ell\right) \circ \mu_V^N \right\rangle \right| \\
&\leq \int_0^T \left| \left\langle M_{m_{\mathcal{G}_1}}^N S_{t-s}^N (f_0^N), \left(M_{m_{\mathcal{G}_1}}^N\right)^{-1} [G^N \pi_N - \pi_N G^\infty] \left(T_s^\infty R_\varphi^\ell\right) \right\rangle \right| ds \\
&\leq \left(\sup_{0 \leq t < T} \left\langle f_t^N, M_{m_{\mathcal{G}_1}}^N \right\rangle \right) \times \\
&\quad \left(\int_0^T \left\| \left(M_{m_{\mathcal{G}_1}}^N\right)^{-1} [G^N \pi_N - \pi_N G^\infty] \left(T_s^\infty R_\varphi^\ell\right) \right\|_{L^\infty(\mathbb{E}_N)} ds \right) \\
&\leq C_{m_{\mathcal{G}_1}, T} \varepsilon(N) \sup_{\mathbf{r} \in \mathbf{R}_{\mathcal{G}_1}} \int_0^T \left[T_s^\infty R_\varphi^\ell \right]_{C_{\Lambda_1}^{1, \eta}(P_{\mathcal{G}_1}, \mathbf{r})} ds.
\end{aligned}$$

Now, let us fix $\mathbf{r} \in \mathbf{R}_{\mathcal{G}_1}$. Since

$$T_t^\infty(R_\varphi^\ell) = R_\varphi^\ell \circ S_t^{NL} \quad \text{with } S_t^{NL} \in C_{\Lambda_2}^{1,\eta}(P_{\mathcal{G}_1,\mathbf{r}}; P_{\mathcal{G}_2})$$

and $R_\varphi^\ell \in C^{1,1}(P_{\mathcal{G}_2})$ because $\varphi \in \mathcal{F}_2^{\otimes \ell}$ (see subsection 2.8), we can apply Lemma 2.12 and use assumption **(A4)** to obtain

$$T_t^\infty(R_\varphi^\ell) \in C_{\Lambda_1}^{1,\eta}(P_{\mathcal{G}_1,\mathbf{r}})$$

with

$$\left[T_s^\infty(R_\varphi^\ell) \right]_{C_{\Lambda_1}^{1,\eta}(P_{\mathcal{G}_1,\mathbf{r}})} \leq \left([S_t^{NL}]_{C_{\Lambda_2}^{1,\eta}(P_{\mathcal{G}_1,\mathbf{r}}, P_{\mathcal{G}_2})} + [S_t^{NL}]_{C_{\Lambda_2}^{0,\eta''}(P_{\mathcal{G}_1,\mathbf{r}}, P_{\mathcal{G}_2})}^{1+\eta'} \right) \|R_\varphi^\ell\|_{C^{1,\eta}(P_{\mathcal{G}_2})}$$

and

$$\Lambda_2 = \Lambda_1^{1/(1+\eta')}.$$

Note that the two different sets of indices

$$(\eta', \eta'') = (\eta, 1) \quad \text{or} \quad (\eta', \eta'') = (1, (1+\eta)/2)$$

in Lemma 2.12 correspond to the two cases discussed in Lemma 2.12.

We then deduce thanks to (2.11) and assumption **(A4)**:

$$(3.8) \quad \int_0^T [T_s^\infty(R_\varphi^\ell)]_{C_{\Lambda_1}^{1,\eta}(P_{\mathcal{G}_1,\mathbf{r}})} ds \leq C_T^\infty \ell^2 \left(\|\varphi\|_{\mathcal{F}_2 \otimes (L^\infty)^{\ell-2}} + \|\varphi\|_{\mathcal{F}_2 \otimes (L^\infty)^{\ell-1}} \right).$$

Then we go back to the computation (3.7), and plugging (3.8) we deduce (3.6).

Step 3: Estimate of the third term \mathcal{T}_3 . Let us prove that for any $t \geq 0$ and $N \geq \ell$ we have

$$\begin{aligned} \mathcal{T}_3 &:= \left| \left\langle f_0^N, \left(T_t^\infty R_\varphi^\ell \right) \circ \mu_V^N \right\rangle - \left\langle \left(S_t^{NL}(f_0) \right)^{\otimes \ell}, \varphi \right\rangle \right| \leq \\ &\leq [R_\varphi]_{C^{0,1}} \Theta_{a^N, T} \left(\mathcal{W}_{1, P_{\mathcal{G}_3}} \left(\pi_P^N f_0^N, \delta_{f_0} \right) \right) \end{aligned}$$

where $\Theta_{a,T}$ was introduced in assumption **(A5)**, and $a = a^N$ is defined by

$$a^N := \max \left\{ M_{\mathcal{G}_3}(f_0), C_{m_{\mathcal{G}_3,0}}^N \right\}$$

where $C_{m_{\mathcal{G}_3,0}}^N$ was introduced in assumption **(A1)-(ii)**.

Assumption **(A1)-(ii)** indeed implies that

$$\text{Supp } f_0^N \subset \mathcal{K} := \left\{ V \in \mathbb{R}^{dN} \text{ s. t. } M_{m_{\mathcal{G}_3}}^N(\mu_V^N) = \frac{1}{N} \sum_{i=1}^N m_{\mathcal{G}_3}(v_i) \leq C_{m_{\mathcal{G}_3,0}}^N \right\}.$$

Hence we are in position to apply **(A5)** for the functions f_0 and μ_V^N on the support of f_0^N since $M_{m_{\mathcal{G}_3}}(f_0)$ is bounded by assumption, and $M_{\mathcal{G}_3}(\mu_V^N)$ is bounded by $C_{m_{\mathcal{G}_3,0}}^N$ when restricting to $V \in \mathcal{K}$ thanks to the previous equation.

Let us also recall that $R_\varphi^\ell \in C^{0,1}(P_{\mathcal{G}_3}, \mathbb{R})$ because $\varphi \in \mathcal{F}_3^{\otimes \ell}$.

We then write

$$\begin{aligned} \mathcal{T}_3 &= \left| \left\langle f_0^N, R_\varphi^\ell(S_t^{NL}(\mu_V^N)) \right\rangle - \left\langle f_0^N, R_\varphi^\ell(S_t^{NL}(f_0)) \right\rangle \right| \\ &= \left| \left\langle f_0^N, R_\varphi^\ell(S_t^{NL}(\mu_V^N)) - R_\varphi^\ell(S_t^{NL}(f_0)) \right\rangle \right| \\ &\leq [R_\varphi]_{C^{0,1}(P_{\mathcal{G}_3})} \langle f_0^N, d_{\mathcal{G}_3}(S_t^{NL}(f_0), S_t^{NL}(\mu_V^N)) \rangle. \end{aligned}$$

We now apply **(A5)** to get

$$\forall t \in [0, T], \quad d_{\mathcal{G}_3}(S_t^{NL}(f_0), S_t^{NL}(\mu_V^N)) \leq \Theta_{a^N, T}(d_{\mathcal{G}_3}(f_0, \mu_V^N))$$

on the support of f_0^N , and therefore

$$\mathcal{T}_3 \leq [R_\varphi]_{C^{0,1}(P_{\mathcal{G}_3})} \langle f_0^N, \Theta_{a^N, T}(\mathrm{d}_{\mathcal{G}_3}(f_0, \mu_V^N)) \rangle.$$

We can further rewrite the last equation in the following way by using the concavity of the $\Theta_{a^N, T}$ function:

$$\mathcal{T}_3 \leq [R_\varphi]_{C^{0,1}(P_{\mathcal{G}_3})} \Theta_{a^N, T}(\langle f_0^N, \mathrm{d}_{\mathcal{G}_3}(f_0, \mu_V^N) \rangle)$$

which concludes the proof of this step.

The proof of Theorem 3.1 is complete by piling the previous steps.

4. THE N -PARTICLE APPROXIMATION AT INITIAL TIME

4.1. Comparison of distances on probabilities. In the following lemma we compare the different metrics and norms defined Subsection 2.5. Let us denote

$$M_k(f, g) := \max \left\{ \langle f, \langle v \rangle^k \rangle; \langle g, \langle v \rangle^k \rangle \right\}.$$

Lemma 4.1. *Let $f, g \in P(\mathbb{R}^d)$ and $k \in (0, \infty)$, then the following estimates hold:*

(i) *For any $q \in (1, +\infty)$ and any $k \in [q - 1, \infty)$:*

$$(4.1) \quad W_1(f, g) \leq W_q(f, g) \leq 2^{\frac{k+1}{q}} M_{k+1}(f, g)^{\frac{q-1}{qk}} W_1(f, g)^{\frac{1}{q}(1-\frac{q-1}{k})}.$$

(ii) *For any $s \in (0, 1]$,*

$$(4.2) \quad |f - g|_s \leq 2^{(1-s)} W_s(f, g) \leq 2^{(1-s)} W_1(f, g)^s.$$

(iii) *For any $s \in (d/2, d/2 + 1)$,*

$$(4.3) \quad \|f - g\|_{\dot{H}^{-s}}^2 \leq \frac{8 |\mathbb{S}^{d-1}|}{(2s - d)} \left(\frac{(2s - d)}{4(d + 2 - 2s)} \right)^{s - \frac{d}{2}} |f - g|_1^{2s-d}.$$

(iv) *For any $s > 0$ and $k > 0$ we have*

$$(4.4) \quad [f - g]_1^* \leq C(d, s, k) M_{k+1}(f, g)^{\frac{d}{d+k(d+s)}} |f - g|_s^{\frac{k}{d+k(d+s)}}$$

for some constant $C(d, s, k) > 0$ depending on d, s and k .

(v) *For any*

$$s \in \left(\max \left\{ \frac{d}{2}; 1 \right\}, \frac{d}{2} + 1 \right)$$

and $k > 0$ we have

$$(4.5) \quad [f - g]_1^* \leq C(d, s, k) M_{k+1}(f, g)^{\frac{d}{d+2ks}} \|f - g\|_{\dot{H}^{-s}}^{\frac{2k}{d+2ks}}$$

for some constant $C(d, s, k) > 0$ depending on d, s and k .

Proof of Lemma 4.1. Let us consider each inequality one by one.

Point (i). The inequality (4.1) is well-known in optimal transport theory, we refer for instance to [69, 17].

Point (ii). Let us prove inequality (4.2). Let $\pi \in \Pi(f, g)$. We write

$$\begin{aligned} |\hat{f}(\xi) - \hat{g}(\xi)| &= \left| \int_{\mathbb{R}^d \times \mathbb{R}^d} \left(e^{-i v \cdot \xi} - e^{-i w \cdot \xi} \right) \pi(dv, dw) \right| \\ &\leq \int_{\mathbb{R}^d \times \mathbb{R}^d} \left| e^{-i v \cdot \xi} - e^{-i w \cdot \xi} \right| \pi(dv, dw) \\ &\leq 2^{(1-s)} \int_{\mathbb{R}^d \times \mathbb{R}^d} |v - w|^s |\xi|^s \pi(dv, dw), \end{aligned}$$

which yields the first inequality in (4.2) by taking the supremum in $\xi \in \mathbb{R}^d$ and the infimum in $\pi \in \Pi(f, g)$. The second inequality is then immediate from the concavity estimate

$$W_s(f, g) \leq (W_1(f, g))^s.$$

Point (iii). Let us prove the inequality (4.3). Consider $R > 0$ and the ball

$$B_R := \left\{ x \in \mathbb{R}^d ; |x| \leq R \right\},$$

and write

$$\begin{aligned} \|f - g\|_{\dot{H}^{-s}}^2 &= \int_{B_R} \frac{|\hat{f}(\xi) - \hat{g}(\xi)|^2}{|\xi|^{2s}} d\xi + \int_{B_R^c} \frac{|\hat{f}(\xi) - \hat{g}(\xi)|^2}{|\xi|^{2s}} d\xi \\ &\leq |f - g|_1^2 \int_{B_R} \frac{d\xi}{|\xi|^{2(s-1)}} + 4 \int_{B_R^c} \frac{d\xi}{|\xi|^{2s}} \\ &\leq |\mathbb{S}^{d-1}| R^{d-2s} \left(\frac{R^2}{(d+2-2s)} |f - g|_1^2 + \frac{4}{(2s-d)} \right). \end{aligned}$$

We optimize this estimate by choosing

$$R = \left(\frac{4(d+2-2s)}{(2s-d)} \right)^{\frac{1}{2}} |f - g|_1^{-1}$$

which yields

$$\|f - g\|_{\dot{H}^{-s}}^2 \leq \frac{8 |\mathbb{S}^{d-1}|}{(2s-d)} \left(\frac{(2s-d)}{4(d+2-2s)} \right)^{s-\frac{d}{2}} |f - g|_1^{2s-d}$$

which concludes the proof of (4.3).

Point (iv). Let us now prove inequality (4.4). We introduce a truncation function

$$\chi_R(x) = \chi\left(\frac{x}{R}\right), \quad R > 0$$

where

$$\chi \in C_c^\infty(\mathbb{R}^d), \quad [\chi]_1 \leq 1, \quad \text{and} \quad 0 \leq \chi \leq 1, \quad \chi \equiv 1 \text{ on } B(0, 1)$$

and a mollifier

$$\gamma_\varepsilon(x) = \varepsilon^{-d} \gamma\left(\frac{x}{\varepsilon}\right), \quad \varepsilon > 0 \quad \text{with} \quad \gamma(x) = \frac{e^{-\frac{|x|^2}{2}}}{(2\pi)^{d/2}}.$$

In particular we have an explicit formula for the Fourier transform of this mollifier:

$$\hat{\gamma}_\varepsilon(\xi) = \hat{\gamma}(\varepsilon\xi) = e^{-\varepsilon^2 \frac{|\xi|^2}{2}}.$$

Fix $\varphi \in W^{1,\infty}(\mathbb{R}^d)$ such that $[\varphi]_1 \leq 1$, $\varphi(0) = 0$ and define

$$\varphi_R := \varphi \chi_R, \quad \varphi_{R,\varepsilon} = \varphi_R * \gamma_\varepsilon$$

and write

$$\int \varphi(df - dg) = \int \varphi_{R,\varepsilon}(df - dg) + \int (\varphi_R - \varphi_{R,\varepsilon})(df - dg) + \int (\varphi - \varphi_R)(df - dg).$$

For the last term, we have

$$\begin{aligned} (4.6) \quad \left| \int (\varphi_R - \varphi)(df - dg) \right| &\leq \int (1 - \chi_R) |\varphi| (df + dg) \\ &\leq \int_{B_R^c} [\varphi]_1 \frac{|x|^{k+1}}{R^k} (df + dg) \leq \frac{M_{k+1}(f, g)}{R^k}. \end{aligned}$$

Concerning the second term, we observe that

$$|\nabla \varphi_R(x)| \leq \chi\left(\frac{x}{R}\right) + |\varphi(x)| |\nabla(\chi_R)(x)| \leq \chi\left(\frac{x}{R}\right) + \frac{|x|}{R} \left| \nabla \chi\left(\frac{x}{R}\right) \right|,$$

so that

$$\forall q \in [1, \infty], \quad \|\nabla \varphi_R\|_{L^q} \leq C(q, d, \chi) R^{\frac{d}{q}},$$

for some constant depending on q , d and χ . Next, using that

$$\begin{aligned} \|\varphi_R - \varphi_{R,\varepsilon}\|_\infty &\leq \|\nabla \varphi_R\|_\infty \left(\int_{\mathbb{R}^d} \gamma_\varepsilon(x) |x| dx \right) \\ &= \varepsilon \|\nabla \varphi_R\|_\infty \left(\int_{\mathbb{R}^d} \gamma(x) |x| dx \right) \leq C(q, d, \chi) \varepsilon, \end{aligned}$$

we find

$$(4.7) \quad \left| \int (\varphi_R - \varphi_{R,\varepsilon}) (df - dg) \right| \leq C(q, d, \chi) \varepsilon.$$

Concerning the first term, using Parseval's identity, we have (the "hat" denotes as usual the Fourier transform)

$$\begin{aligned} \left| \int \varphi_{R,\varepsilon} (f - g) \right| &= \frac{1}{(2\pi)^d} \left| \int \hat{\varphi}_R \hat{\gamma}_\varepsilon \overline{(\hat{f} - \hat{g})} d\xi \right| \\ &\leq \frac{1}{(2\pi)^d} \|\nabla \varphi_R\|_{L^1} \left\| \frac{\hat{f} - \hat{g}}{|\xi|^s} \right\|_{L^\infty} \left(\int_{\mathbb{R}^d} |\xi|^{s-1} e^{-\varepsilon^2 \frac{|\xi|^2}{2}} d\xi \right) \\ &\leq C(d, \chi) R^d |f - g|_s \varepsilon^{-(d+s-1)} \\ (4.8) \quad &\leq C R^d \varepsilon^{-(d+s-1)} |f - g|_s. \end{aligned}$$

Gathering (4.6), (4.7) and (4.8), we get

$$[f - g]_1^* \leq C(q, d, \chi) \left(\frac{M_{k+1}(f, g)}{R^k} + \varepsilon + R^d \varepsilon^{-(d+s-1)} |f - g|_s \right).$$

This yields (4.4) by optimizing the parameters ε and R with

$$R = M_{k+1}(f, g)^{\frac{1}{d+k}} |f - g|_s^{-\frac{1}{d+k}} \varepsilon^{\frac{d+s-1}{d+k}}$$

and then

$$\varepsilon = M_{k+1}(f, g)^{\frac{d}{d+k(d+s)}} |f - g|_s^{\frac{k}{d+k(d+s)}}.$$

Point (v). Let us now prove inequality (4.5).

Let us consider some smooth φ such that $[\varphi]_1 \leq 1$, $\varphi(0) = 0$ and let us perform the same decomposition as before:

$$\int \varphi (df - dg) = \int \varphi_{R,\varepsilon} (df - dg) + \int (\varphi_R - \varphi_{R,\varepsilon}) (df - dg) + \int (\varphi - \varphi_R) (df - dg).$$

The first term is controlled by

$$\left| \int \varphi_{R,\varepsilon} (df - dg) \right| = \left| \int \hat{\varphi}_{R,\varepsilon} |\xi|^s \frac{(\hat{f} - \hat{g})}{|\xi|^s} \right| \leq \|\varphi_{R,\varepsilon}\|_{\dot{H}^s} \|f - g\|_{\dot{H}^{-s}}$$

with

$$\begin{aligned} \|\varphi_{R,\varepsilon}\|_{\dot{H}^s} &= \left(\int |\xi|^2 |\widehat{\varphi_R}|^2 |\xi|^{2(s-1)} |\hat{\gamma}_\varepsilon|^2 d\xi \right)^{1/2} \\ &\leq \|\nabla(\varphi_R)\|_{L^2} \left\| |\xi|^{s-1} \hat{\gamma}_\varepsilon(\xi) \right\|_{L^\infty} \\ &\leq C(s) \|\nabla(\varphi_R)\|_{L^2} \varepsilon^{1-s} \leq C(d, s) R^{\frac{d}{2}} \varepsilon^{-(s-1)} \end{aligned}$$

with the same arguments as above.

The second term and the last term are controlled exactly as in (4.6) and (4.7), which yields

$$[f - g]_1^* \leq C \left(\varepsilon + \frac{M_{k+1}(f, g)}{R^k} + R^{\frac{d}{2}} \varepsilon^{-(s-1)} \|f - g\|_{\dot{H}^{-s}} \right).$$

This yields (4.5) by optimizing the parameters ε and R with

$$R = M_{k+1}(f, g)^{\frac{2}{d+2k}} \|f - g\|_{\dot{H}^{-s}}^{-\frac{2}{d+2k}} \varepsilon^{\frac{2(s-1)}{d+2k}}$$

and then

$$\varepsilon = M_{k+1}(f, g)^{\frac{d}{d+2ks}} \|f - g\|_{\dot{H}^{-s}}^{\frac{2k}{d+2ks}}.$$

□

4.2. Quantitative law of large numbers for measures. Let us recall and extend the definition of the functional $\mathcal{W}_d^N(f)$ which was introduced in (3.2). For any function

$$D : P_k(E) \times P_k(E) \rightarrow \mathbb{R}_+, \quad (f, g) \mapsto D(f, g)$$

(where $k \geq 0$ is the index of a polynomial weight possibly required for the correct definition of D) such that

$$D(f, g) = 0 \quad \text{if and only if} \quad f = g$$

it is legitimate to define

$$\forall f \in P_k(\mathbb{R}^d), \quad \mathcal{W}_D^N(f) := \langle f^{\otimes N}, D(\pi_E^N, f) \rangle = \int_{E^N} D(\mu_V^N, f) f^{\otimes N}(dV)$$

(D should be thought as some function of a proper distance).

For well chosen function D , the goal of the next lemma is to quantify the rate of convergence

$$\mathcal{W}_D^N(f) \xrightarrow{N \rightarrow +\infty} 0 \quad \text{in the case } E = \mathbb{R}^d.$$

Lemma 4.2. *We have the following rates for the \mathcal{W} function:*

(i) *Let us consider*

$$\forall f, g \in P_2(\mathbb{R}^d), \quad D_1(f, g) := \|f - g\|_{\dot{H}^{-s}}^2.$$

Then for any $s \in (d/2, d/2 + 1)$ and $N \geq 1$ there holds

$$(4.9) \quad \forall f \in P_2(\mathbb{R}^d), \quad \mathcal{W}_{D_1}^N(f) = \int_{\mathbb{R}^{dN}} \|\mu_V^N - f\|_{\dot{H}^{-s}}^2 df^{\otimes N}(V) \leq \frac{C}{N}$$

for some constant C depending on the second moment of f .

(ii) *Let us consider*

$$\forall f, g \in P_2(\mathbb{R}^d), \quad D_2(f, g) := \|f - g\|_{H^{-s}}^2.$$

Then for any $s > d/2$ and $N \geq 1$ there holds

$$(4.10) \quad \forall f \in P_2(\mathbb{R}^d), \quad \mathcal{W}_{D_2}^N(f) = \int_{\mathbb{R}^{dN}} \|\mu_V^N - f\|_{H^{-s}}^2 df^{\otimes N}(V) \leq \frac{C}{N}$$

for some constant C depending on the second moment of f .

(iii) *Let us consider*

$$\forall f, g \in P_1(\mathbb{R}^d), \quad D_3(f, g) := W_1(f, g).$$

Then for any $\eta > 0$ there exists $k \geq 1$ such that for any $N \geq 1$

$$(4.11) \quad \forall f \in P_k(\mathbb{R}^d), \quad \mathcal{W}_{D_3}^N(f) = \int_{\mathbb{R}^{dN}} W_1(\mu_V^N, f) df^{\otimes N}(V) \leq \frac{C}{N^{\max\{d, 2\} + \eta}}$$

for some constant C depending on η and the k -th moment of f .

(iv) *Let us consider*

$$\forall f, g \in P_2(\mathbb{R}^d), \quad D_4(f, g) := (W_2(f, g))^2.$$

Then for any $\eta > 0$ there exists $k \geq 2$ such that for any $N \geq 1$

$$(4.12) \quad \forall f \in P_k(\mathbb{R}^d), \quad \mathcal{W}_{D_4}^N(f) = \int_{\mathbb{R}^{dN}} (W_2(\mu_V^N, f))^2 df^{\otimes N}(V) \leq \frac{C}{N^{\frac{1}{\max\{d, 2\} + \eta}}}$$

for some constant C depending on η and the k -th moment of f .

Remarks 4.3. (1) Estimate (4.12) has to be compared with the following classical estimate (see e.g. [64]): for any and any $N \geq 1$ there holds

$$(4.13) \quad \forall f \in P_{d+5}(\mathbb{R}^d), \quad \mathcal{W}_{W_2}^N(f) \leq \frac{C}{N^{\frac{1}{d+4}}}$$

where the constant $C > 0$ depends on the $(d+5)$ -th moment of f . It is worth mentioning that our estimate (4.12) improves on (4.13) when $d \leq 3$ and k is large enough.

Similarly, if we try to deduce from (4.13) some estimate on the rate for the W_1 distance, by using a Hölder inequality we get that for any $N \geq 1$

$$\forall f \in P_{d+5}(\mathbb{R}^d), \quad \mathcal{W}_{W_1}^N(f) \leq \frac{C}{N^{\frac{1}{d+4}}}$$

where the constant $C > 0$ depends on the $(d+5)$ -th moment of f . For any d this is worst than (4.11) as soon as when the probability $f \in P_k(\mathbb{R}^d)$ with k large enough.

(2) When $f, g \in P(\mathbb{R}^d)$ are compactly supported, observe that the estimate (4.5) improves into

$$\forall s \geq 1 \quad [f - g]_1^* \leq C \|f - g\|_{H^{-s}}^{1/s},$$

for a constant C depending on s and on a common bound R of the support of f and g .

If furthermore $d = 1$, we can take $s = 1$ in order to apply (4.9) in the proof of (4.11) below and we obtain the “*optimal rate*” of convergence in the functional law of large numbers in Wasserstein distance W_1 :

$$\forall N \geq 1 \quad \mathcal{W}_{W_1}^N(f) \leq \frac{C}{\sqrt{N}}.$$

In higher dimension $d \geq 2$, the restriction $s > d/2$ means that we do not produce a better estimate than (4.11) by this line of argument.

(3) As was kindly pointed out by M. Hauray, estimate (4.12) should also be compared with some estimates in [25] where the related quantity

$$\mathcal{Z}^N(f) := \int_{\mathbb{R}^{2dN}} W_1(\mu_{V_1}^N, \mu_{V_2}^N) df^{\otimes N}(V_1) df^{\otimes N}(V_2)$$

is considered. When $f \in P(\mathbb{R}^d)$ has compact support and $d \geq 3$, they prove that

$$\mathcal{Z}^N(f) \leq \frac{C}{N^{1/d}}$$

where the constant depends on the support of f .

Since for any $f, g \in P_1(\mathbb{R}^d)$ and for any $\varphi \in \text{Lip}_1(\mathbb{R}^d)$ we have

$$\begin{aligned} \int_{\mathbb{R}^{dN}} W_1(g, \mu_V^N) df^{\otimes N}(V) &\geq \int_{\mathbb{R}^{d(N+1)}} \varphi(v) (dg - d\mu_V^N)(v) df^{\otimes N}(V) \\ &= \int_{\mathbb{R}^d} \varphi(v) dg(v) - \frac{1}{N} \sum_{i=1}^N \int_{\mathbb{R}^{dN}} \varphi(v_i) df^{\otimes N}(V) \\ &= \int_{\mathbb{R}^d} \varphi(v) (dg - df)(v), \end{aligned}$$

we deduce by minimizing in φ that

$$W_1(f, g) \leq \int_{\mathbb{R}^{dN}} W_1(g, \mu_V^N) df^{\otimes N}(V),$$

and therefore

$$\mathcal{W}_{W_1}^N(f) \leq \mathcal{Z}^N(f).$$

As a consequence, when $f \in P(\mathbb{R}^d)$ has compact support and $d \geq 3$ we obtain from this line of argument the stronger estimate

$$\mathcal{W}_{W_1}^N(f) \leq \frac{C}{N^{1/d}}.$$

It is likely that one could obtain similar estimates to (4.11) by tracking the formula for the constants in the results of [25] and combining them with moment bounds and some interpolation.

On the other hand, observe that our estimate (4.11) is *almost optimal* in the sense that we can not expect a better convergence rate than (4.11) with $\eta = 0$, as it is stressed in [63, Appendix].

Proof of Lemma 4.2. We split the proof into two steps.

Step 1. Let us prove (4.9) (note that (4.10) is then readily implied by (4.9)).

Let us fix $f \in P_2(\mathbb{R}^d)$. We write in Fourier transform

$$\left(\hat{\mu}_V^N - \hat{f} \right) (\xi) = \frac{1}{N} \sum_{j=1}^N \left(e^{-i v_j \cdot \xi} - \hat{f}(\xi) \right),$$

which implies

$$\begin{aligned} \mathcal{W}_{\|\cdot\|_{\dot{H}^{-s}}^2}^N(f) &= \int_{\mathbb{R}^{Nd}} \left(\int_{\mathbb{R}^d} \frac{|\hat{\mu}_V^N - \hat{f}|^2}{|\xi|^{2s}} d\xi \right) df^{\otimes N}(V) \\ &= \frac{1}{N^2} \sum_{j_1, j_2=1}^N \int_{\mathbb{R}^{(N+1)d}} \frac{\left(e^{-i v_{j_1} \cdot \xi} - \hat{f}(\xi) \right) \overline{\left(e^{-i v_{j_2} \cdot \xi} - \hat{f}(\xi) \right)}}{|\xi|^{2s}} d\xi df^{\otimes N}(V). \end{aligned}$$

Observe then that

$$\int_{\mathbb{R}^d} \left(e^{-i v_j \cdot \xi} - \hat{f}(\xi) \right) df(v_j) = 0, \quad j = 1, \dots, d,$$

which implies that

$$\int_{\mathbb{R}^{(N+1)d}} \frac{\left(e^{-i v_{j_1} \cdot \xi} - \hat{f}(\xi) \right) \overline{\left(e^{-i v_{j_2} \cdot \xi} - \hat{f}(\xi) \right)}}{|\xi|^{2s}} d\xi df^{\otimes N}(V) = 0$$

as soon as $j_1 \neq j_2$, and

$$\begin{aligned} \int_{\mathbb{R}^d} \left| e^{-i v \cdot \xi} - \hat{f}(\xi) \right|^2 df(v) &= \int_{\mathbb{R}^d} \left[1 - e^{-i v \cdot \xi} \overline{\hat{f}(\xi)} - e^{i v \cdot \xi} \hat{f}(\xi) + |\hat{f}(\xi)|^2 \right] df(v) \\ &= 1 - |\hat{f}(\xi)|^2. \end{aligned}$$

We deduce that

$$\begin{aligned} \mathcal{W}_{\|\cdot\|_{\dot{H}^{-s}}^2}^N(f) &= \frac{1}{N^2} \sum_{j=1}^N \int_{\mathbb{R}^{(N+1)d}} \frac{\left| e^{-i v_j \cdot \xi} - \hat{f}(\xi) \right|^2}{|\xi|^{2s}} d\xi f^{\otimes N}(dV) \\ &= \frac{1}{N} \int_{\mathbb{R}^{2d}} \frac{\left| e^{-i v \cdot \xi} - \hat{f}(\xi) \right|^2}{|\xi|^{2s}} d\xi f(dv) \\ &= \frac{1}{N} \int_{\mathbb{R}^d} \frac{(1 - |\hat{f}(\xi)|^2)}{|\xi|^{2s}} d\xi. \end{aligned}$$

Finally, denoting

$$M_2 := \int_{\mathbb{R}^d} \langle v \rangle^2 df(v)$$

we have

$$\hat{f}(\xi) = 1 + i \langle f, v \rangle \cdot \xi + \mathcal{O}(M_2 |\xi|^2),$$

and therefore

$$\begin{aligned} |\hat{f}(\xi)|^2 &= (1 + i \langle f, v \rangle \cdot \xi + \mathcal{O}(M_2 |\xi|^2)) \left(1 - i \langle f, v \rangle \cdot \xi + \overline{\mathcal{O}(M_2 |\xi|^2)} \right) \\ &= 1 + \mathcal{O}(M_2 |\xi|^2), \end{aligned}$$

which implies

$$\begin{aligned} \mathcal{W}_{\|\cdot\|_{\dot{H}^{-s}}^2}^N(f) &= \frac{1}{N} \left(\int_{|\xi| \leq 1} \frac{(1 - |\hat{f}(\xi)|^2)}{|\xi|^{2s}} d\xi + \int_{|\xi| \geq 1} \frac{(1 - |\hat{f}(\xi)|^2)}{|\xi|^{2s}} d\xi \right) \\ &= \frac{1}{N} \left(\int_{|\xi| \leq 1} \frac{M_2}{|\xi|^{2(s-1)}} d\xi + \int_{|\xi| \geq 1} \frac{1}{|\xi|^{2s}} d\xi \right) \leq \frac{C}{N} \end{aligned}$$

from which (4.9) follows.

Step 2. Let us now prove (4.11).

We use first (4.5) in order to get

$$\begin{aligned} \mathcal{W}_{W_1}^N(f) &= \int_{R^{dN}} [\mu_V^N - f]_1^* df^{\otimes N}(V) \\ &\leq C \int_{R^{dN}} (M_{k+1}(f) + M_{k+1}(\mu_V^N))^{\frac{d}{d+2ks}} \left(\|\mu_V^N - f\|_{\dot{H}^{-s}}^2 \right)^{\frac{k}{d+2ks}} df^{\otimes N}(V). \end{aligned}$$

We then perform a Hölder inequality with exponents

$$p = \frac{d+2ks}{k}, \quad p' = \frac{d+2ks}{d+k(2s-1)}$$

and get

$$\begin{aligned} \mathcal{W}_{W_1}^N(f) &\leq C \left(\int_{R^{dN}} (M_{k+1}(f) + M_{k+1}(\mu_V^N))^{\frac{d}{d+k(2s-1)}} df^{\otimes N}(V) \right)^{\frac{d+k(2s-1)}{d+2ks}} \times \\ &\quad \left(\int_{R^{dN}} \|\mu_V^N - f\|_{\dot{H}^{-s}}^2 df^{\otimes N}(V) \right)^{\frac{k}{d+2ks}}. \end{aligned}$$

Since

$$\begin{aligned}
& \int_{\mathbb{R}^{dN}} (M_{k+1}(f) + M_{k+1}(\mu_V^N))^{\frac{d}{d+k(2s-1)}} df^{\otimes N}(V) \\
& \leq \int_{\mathbb{R}^{dN}} (M_{k+1}(f) + M_{k+1}(\mu_V^N)) df^{\otimes N}(V) \\
& \leq M_{k+1}(f) + \int_{\mathbb{R}^{dN}} M_{k+1}(\mu_V^N) df^{\otimes N}(V) \\
& \leq M_{k+1}(f) + \frac{1}{N} \sum_{i=1}^N \int_{\mathbb{R}^{dN}} \langle v_i \rangle^{k+1} df^{\otimes N}(V) \leq 2 M_{k+1}(f)
\end{aligned}$$

we deduce by using (4.9) that

$$\mathcal{W}_{W_1}^N(f) \leq \frac{C(f, k)}{N^{\frac{k}{d+2ks}}}$$

where the constant $C(f, k)$ depends on the $(k+1)$ -th moment of f .

This allows to conclude the proof of (4.11) since

- if $d = 1$ we can take $s = 1$ in (4.5) and then k large enough so that $k/(d+2ks) = 2 + \eta$ with some $\eta > 0$ as small as wanted,
- if $d \geq 2$ we take s close to $d/2$ and then k large enough so that $k/(d+2ks) = 1/d + \eta$ with some $\eta > 0$ as small as wanted.

Then the estimate (4.12) follows from (4.11) with the help of (4.1) and a Hölder inequality. □

4.3. Chaotic initial data with prescribed energy and momentum. In many aspects, the simplest N -particle initial data is the sequence of tensorized initial data $f^{\otimes N}$, $N \geq 1$, where f is a 1-particle distribution. This means perfect chaoticity. On the other hand it has a drawback: since in all applications we shall use pointwise bounds on the energy of the N -particle system (and also sometimes pointwise higher moment bound as in **(A1)-(ii)**), this implies for this kind of initial data that f has to be compactly supported. There is another “natural” choice of initial data, by restricting to one of the subspaces left invariant by the dynamics:

$$\mathcal{S}^N = \left\{ V \in \mathbb{R}^{dN} \text{ s. t. } \frac{1}{N} \sum_{i=1}^N |v_i|^2 = \mathcal{E}, \quad \frac{1}{N} \sum_{i=1}^N v_i = \mathcal{M} \right\}.$$

Without loss of generality we shall often set $\mathcal{M} = 0$ and $\mathcal{E} = 1$ in this formula in the sequel.

The drawback is now that an initial data on \mathcal{S}^N *cannot* be perfectly tensorized, and some additional chaoticity error is paid at initial time. However an advantage of this viewpoint is that it is simpler to study the asymptotic behavior of both the N -particle and the limiting mean-field equation in this setting. Moreover it has historical value since this approach was introduced by Kac (see the discussion in [42, Section 5 “Distributions having Boltzmann’s property”]), although in his case there was only one conservation law, namely the energy, and therefore \mathcal{S}^N was replaced by $\mathbb{S}^{N-1}(\sqrt{N})$. We shall present some results on the construction of chaotic initial data on \mathcal{S}^N , whose proofs are mostly extensions of the precise statements and estimates recently established in [12] on this issue in the setting of Kac on \mathbb{S}^{N-1} . We refer to the work in progress [16] where an extensive study and precise computations of rates shall be performed.

Lemma 4.4. *Consider an initial data*

$$f_0 \in P_4(\mathbb{R}^d) \cap L^\infty(\mathbb{R}^d)$$

which fulfills some moment conditions

$$M_{m_{\mathcal{G}_1}}(f) = \langle f, m_{\mathcal{G}_1} \rangle < +\infty, \quad M_{m_{\mathcal{G}_3}}(f) = \langle f, m_{\mathcal{G}_3} \rangle < +\infty$$

for some positive radially symmetric increasing weight functions $m_{\mathcal{G}_1}$ and $m_{\mathcal{G}_3}$, and let us denote

$$\int_{\mathbb{R}^d} |v|^2 f_0(dv) = \mathcal{E} \in (0, \infty).$$

Let us define a non-decreasing sequence $(\alpha_N)_{N \geq 1}$ as follows:

- If f_0 has compact support

$$(4.14) \quad \text{Supp} f_0 \subset \{v \in \mathbb{R}^d, |v| \leq A\}$$

for some $A > 0$, then

$$\forall N \geq 1, \quad \alpha_N := m_{\mathcal{G}_3}(A).$$

- If f_0 has not compact support, then $(\alpha_N)_{N \geq 1}$ can be chosen as any non-decreasing sequence such that

$$\lim_{N \rightarrow \infty} \alpha_N = +\infty$$

(note in particular that this sequence can grow as slow as wanted).

Then there exists

$$f_0^N \in P(\mathbb{R}^{dN}), \quad N \geq 1,$$

such that

- The sequence $(f_0^N)_{N \geq 1}$ is f_0 -chaotic.
- Its support satisfies

$$\text{Supp} f_0^N \subset \mathcal{S}^N.$$

- It satisfies the following integral moment bound based on m_1 :

$$\forall N \geq 0, \quad \left\langle f_0^N, M_{m_{\mathcal{G}_1}}^N \right\rangle \leq C \langle f_0, m_{\mathcal{G}_1} \rangle$$

where the constant $C > 0$ depends on M_{0, m_1}^{NL} .

- It satisfies the following “support moment bound”:

$$\text{Supp} f_0^N \subset \left\{ V \in \mathbb{R}^{dN}; M_{m_{\mathcal{G}_3}}^N(V) \leq \alpha_N \right\}.$$

- It satisfies a uniform relative entropy bound

$$\frac{1}{N} H(f_0^N | \gamma^N) \leq C,$$

for some constant $C > 0$ (see (1.7) for notations).

- If furthermore the Fisher information associated to f_0 is bounded, that is $I(f_0) < \infty$ (see (1.8) for notations), then f_0^N can be built in such a way that it satisfies a uniform relative Fisher information bound

$$I(f_0^N | \gamma^N) := \frac{1}{N} \int_{\mathcal{S}^N} \left| \nabla \log \frac{df_0^N}{d\gamma^N} \right|^2 f_0^N \leq C,$$

for some constant $C > 0$, where the gradient in this formula stands for the Riemannian gradient on the manifold \mathcal{S}^N .

Proof of Lemma 4.4. For the sake of simplicity, we assume with no loss of generality that the energy $\mathcal{E} = 1$ and that f_0 is centered.

We aim at defining our initial data f_0^N by conditioning the tensorized initial data $f_0^{\otimes N}$ to \mathcal{S}^N :

$$f_0^N(V) = [f_0^{\otimes N}]_{\mathcal{S}^N} := \left(\frac{\prod_{j=1}^N f_0(v_j)}{F_N(\sqrt{N})} \right) \Big|_{\mathcal{S}^N}$$

with

$$F^N(r) := \int_{\mathbb{S}^{dN-1}(r) \cap (\sum_{i=1}^N v_i = 0)} \prod_{j=1}^N f_0(v_j) d\omega.$$

Such a construction obviously satisfies (ii).

It is proved similarly as in [12] (see for instance Theorem 9 in this reference) that this conditioned measure is well-defined, and that it is f_0 -chaotic, which proves (i).

Remark 4.5. Among many interesting intermediate steps and other results, it is also proved in [12] the following estimate: assume for simplicity that $d = 1$ and that f_0 has energy 1, then the function

$$\bar{F}^N(r) := \frac{F^N(r)}{\gamma^N(r)}$$

is asymptotically divergent except for $r = \sqrt{N}$, for which

$$\bar{F}^N(\sqrt{N}) \sim_{N \rightarrow +\infty} \frac{\sqrt{2}}{\Sigma}$$

with

$$\Sigma = \sqrt{\int_{\mathbb{R}} (v^2 - 1)^2 df(v)}.$$

(In fact this result was sketched by Kac [42] but the proof is made more precise in [12].)

This shows in particular that the sequence of chaotic initial data $f_0^{\otimes N}$, $N \geq 1$, as considered many times in the sequel, asymptotically concentrates on the energy sphere with the energy of f_0 . This manifestation of the central limit theorem explains why the construction of Kac (to condition to a given energy sphere) is very natural. It also enlightens why it is possible to expect the kind of uniform in time propagation of chaos results that we shall prove in the next sections for such chaotic initial data.

Point (iii) is just a consequence of the chaoticity with the test function $M_{m_{\mathcal{G}_1}}^N$ (actually an easy truncation and passage to the limit procedure is needed in full rigor).

Concerning point (iv), first if f_0 is compactly supported (4.14) we deduce that

$$\text{Supp } f_0^N \subset \left\{ V \in \mathbb{R}^{dN}, M_{m_{\mathcal{G}_3}}^N(V) \leq m_{\mathcal{G}_3}(A) \right\}$$

and (iv) holds.

In the non compactly supported case, for any increasing sequence $(A_k)_{k \geq 1}$ of positive reals (with A_0 big enough for the following to be well-defined) we define

$$f_{0,k} := \frac{f_0 \mathbf{1}_{|v| \leq A_k}}{f_0(\{|v| \leq A_k\})}.$$

Using the previous we know that

$$f_{0,k}^N := [f_{0,k}^{\otimes N}]$$

(conditioning to the sphere) is $f_{0,k}$ -chaotic. Conditions (ii) and (iii) will therefore be immediately satisfied.

We now want to choose a sequence $k_N \rightarrow \infty$ such that (iv) is satisfied and at the same proving chaoticity *towards* f_0 . It is clear that

$$\text{Supp } f_{0,k}^N \subset \left\{ V \in \mathbb{R}^{dN}; M_{m_3}^N(V) \leq m_3(A_k) \right\}.$$

For any given sequence (α_N) which tends to infinity, we define k_N in such a way that $m_3(A_{k_N}) \leq \alpha_N$ so that $k_N \rightarrow \infty$ when $N \rightarrow \infty$. The chaos property is equivalent to the weak convergence of the 2-marginal, which can be expressed in Wasserstein distance for instance:

$$W_1 \left((f_{0,k}^N)_2, f_0^{\otimes 2} \right) \leq W_1 \left((f_{0,k}^N)_2, f_{0,k}^{\otimes 2} \right) + W_1 \left(f_{0,k}^{\otimes 2}, f_0^{\otimes 2} \right).$$

The last term of the RHS converges to zero only depending on $k \rightarrow \infty$, while the first term in the RHS converges to zero for fixed k as $N \rightarrow \infty$ from the previous part of the proof. Therefore, maybe at the price of a slower increasing sequence k_N we can have both the support moment condition (iv) and

$$W_1 \left((f_{0,k_N}^N)_2, f_0^{\otimes 2} \right) \xrightarrow{N \rightarrow \infty} 0$$

which shows the chaoticity and concludes the proof.

For the proof of (v) and (vi) we refer to [16]. \square

Remarks 4.6. (1) We note that if one only wants to get rid of the compact support requirement in f_0 (used for deriving the support bounds on f_0^N on the energy and $m_{\mathcal{G}_3}$), and not necessarily to prescribe a given energy, another strategy could have been to simply perform the cutoff in the end of the previous proof. In principle it could allow to get better information on the rate of convergence. However a drawback of this approach is that, in the absence of conditioning to an energy sphere, the bound on the support of the energy of f_0^N shall grow with N . In our applications it induces a growth in N of the moment bounds that we prove *along time* on the N -particle system. This growth should be matched by the decay of the scheme and a precise optimized balance could be searched for. We do not pursue this line of research.

(2) Observe that the process of conditioning on the energy sphere is obviously compatible with the equilibrium states in the following sense: if one denotes by γ a centered gaussian equilibrium of the limiting equation with energy normalized to 1, then

$$\gamma^N(V) := [\gamma^{\otimes N}]_{\mathcal{S}^N}$$

is the uniform measure on \mathcal{S}^N , i.e. an equilibrium of the N -particle system.

Let us also state a refinement of the previous lemma which is needed for the applications.

Lemma 4.7. *We use the same setting and assumptions as in Lemma 4.4. We consider some function $\Theta_a(x)$ such that*

$$\forall a > 0, \quad \Theta_a(x) \xrightarrow{N \rightarrow +\infty} 0.$$

Then the sequence (f_0^N) , $N \geq 1$ of Lemma 4.4 satisfies the more precise chaoticity estimate:

$$(4.15) \quad \mathcal{W}_{W_1}(\pi_P^N(f_0^N), f_0) = \int_{\mathbb{R}^{dN}} W_1(\mu_V^N, f_0) df_0^N(V) \xrightarrow{N \rightarrow +\infty} 0.$$

as well as

$$(4.16) \quad \Theta_{a^N}(\mathcal{W}_{W_1}(\pi_P^N(f_0^N), f_0)) \xrightarrow{N \rightarrow +\infty} 0.$$

with

$$a^N = \max \{ \alpha_N; \langle f_0, m_{\mathcal{G}_3} \rangle \}.$$

Remarks 4.8. (1) Using Lemma 4.1 it would be immediate to extend the previous statement to the other weak measure distances we have discussed so far.

(2) Some polynomial rates could be obtained by refining the techniques from [12], see [16].

Proof of Lemma 4.7. Let us prove that (4.15) holds. First, thanks to [66, Proposition 2.2] and the fact that the sequence f_0^N constructed in Lemma 4.4 is f_0 -chaotic, we deduce that

$$\pi_P^N f_0^N \rightharpoonup \delta_{f_0} \text{ in } P\left(P\left(\mathbb{R}^d\right)\right)$$

(which means convergence when testing against functions in $C(P(\mathbb{R}^d))$).

Next, thanks to [72, Theorem 7.12], (4.15) boils down to prove the tightness estimate

$$(4.17) \quad \lim_{R \rightarrow \infty} \sup_{N \in \mathbb{N}^*} \int_{W_1(\rho, f_0) \geq R} W_1(\rho, f_0) d\pi_P^N f_0^N(\rho) = 0.$$

Let us prove that it follows easily from the following bound

$$(4.18) \quad \sup_{N \in \mathbb{N}^*} \int_{E^N} (W_1(\mu_V^N, f_0))^2 df_0^N(V) < \infty.$$

Indeed (4.18) implies that uniformly in $N \geq 1$

$$\begin{aligned} \int_{W_1(\rho, f_0) \geq R} W_1(\rho, f_0) d\pi_P^N f_0^N(\rho) &= \int_{V \in E^N \text{ s.t. } W_1(\mu_V^N, f_0) \geq R} W_1(\mu_V^N, f_0) df_0^N(V) \\ &\leq \frac{1}{R} \int_{E^N} (W_1(\mu_V^N, f_0))^2 df_0^N(V) \leq \frac{C}{R} \xrightarrow{R \rightarrow \infty} 0 \end{aligned}$$

which concludes the proof of (4.17).

In order to show (4.18), we infer that from [72, Theorem 7.10]

$$\begin{aligned} (W_1(\mu_V^N, f_0))^2 &\leq \|\mu_V^N - f_0\|_{M_1^1}^2 \leq 2\|\mu_V^N\|_{M_1^1}^2 + 2\|f_0\|_{M_1^1}^2 \\ &\leq 2(M_1^N(V))^2 + 2\|f_0\|_{M_1^1}^2 \leq 2M_2^N(V) + 2\|f_0\|_{M_1^1}^2, \end{aligned}$$

which implies

$$\int_{E^N} (W_1(\mu_V^N, f_0))^2 df_0^N(V) \leq 2\|f_0\|_{M_1^1}^2 + 2\langle f_0^N, M_2^N \rangle,$$

which, together with (ii) in Lemma 4.4, ends the proof of (4.18) and then of (4.15).

Then it is an easy diagonal process exercise to build a sequence (α^N) such that (4.16) holds by using the assumption on the function Θ_a together with (4.15). \square

5. TRUE MAXWELL MOLECULES

5.1. The model. Let us consider $E = \mathbb{R}^d$, $d \geq 2$, and a N -particle system undergoing space homogeneous random Boltzmann type collisions according to a collision kernel

$$B = \Gamma(z) b(\cos \theta)$$

(see Subsection 1.1). More precisely, given a pre-collisional system of velocity variables

$$V = (v_1, \dots, v_N) \in E^N = (\mathbb{R}^d)^N,$$

the stochastic process is:

- (i) for any $i' \neq j'$, draw a random time $T_{\Gamma(|v_{i'} - v_{j'}|)}$ of collision accordingly to an exponential law of parameter $\Gamma(|v_{i'} - v_{j'}|)$, and then choose the collision time T_1 and the colliding pair (v_i, v_j) (which is a.s. well-defined) in such a way that

$$T_1 = T_{\Gamma(|v_i - v_j|)} := \min_{1 \leq i' \neq j' \leq N} T_{\Gamma(|v_{i'} - v_{j'}|)};$$

(ii) then draw $\sigma \in S^{d-1}$ according to the law $b(\cos \theta_{ij})$, where

$$\cos \theta_{ij} = \sigma \cdot (v_j - v_i) / |v_j - v_i|;$$

(iii) the new state after collision at time T_1 becomes

$$(5.1) \quad V_{ij}^* = (v_1, \dots, v_i^*, \dots, v_j^*, \dots, v_N),$$

where only velocities labelled i and j have changed, according to the rotation

$$(5.2) \quad v_i^* = \frac{v_i + v_j}{2} + \frac{|v_i - v_j| \sigma}{2}, \quad v_j^* = \frac{v_i + v_j}{2} - \frac{|v_i - v_j| \sigma}{2}.$$

The associated Markov process

$$(\mathcal{V}_t)_{t \geq 0} \text{ on } (\mathbb{R}^d)^N$$

is then built by iterating the above construction.

After rescaling time $t \rightarrow t/N$ in order that the number of interactions is of order $\mathcal{O}(1)$ on finite time interval (see [65]) we denote by f_t^N the law of \mathcal{V}_t and S_t^N the associated semigroup. We recall the notation G^N and T_t^N respectively for the dual generator and dual semigroup, as in the previous abstract construction.

The so-called *Master equation* on the law f_t^N is given in dual form by

$$(5.3) \quad \partial_t \langle f_t^N, \varphi \rangle = \langle f_t^N, G^N \varphi \rangle$$

with

$$(5.4) \quad (G^N \varphi)(V) = \frac{1}{N} \sum_{1 \leq i < j \leq N} \Gamma(|v_i - v_j|) \int_{S^{d-1}} b(\cos \theta_{ij}) [\varphi_{ij}^* - \varphi] d\sigma$$

$$\text{where } \varphi_{ij}^* = \varphi(V_{ij}^*) \text{ and } \varphi = \varphi(V) \in C_b(\mathbb{R}^{Nd}).$$

This collision process is invariant under velocities permutations and satisfies the microscopic conservations of momentum and energy at any collision time

$$\sum_{j=1}^N v_j^* = \sum_{j=1}^N v_j \quad \text{and} \quad |V^*|^2 = \sum_{j=1}^N |v_j^*|^2 = \sum_{j=1}^N |v_j|^2 = |V|^2.$$

As a consequence, for any symmetric initial law $f_0^N \in P_{\text{sym}}(\mathbb{R}^{Nd})$ the law f_t^N at later times is also a symmetric probability, and it conserves momentum and energy:

$$\forall \alpha = 1, \dots, d, \quad \int_{\mathbb{R}^{dN}} \left(\sum_{j=1}^N v_{j,\alpha} \right) f_t^N(dV) = \int_{\mathbb{R}^{dN}} \left(\sum_{j=1}^N v_{j,\alpha} \right) f_0^N(dV),$$

where $(v_{j,\alpha})_{1 \leq \alpha \leq d}$ denote the components of $v_j \in \mathbb{R}^d$, and

$$(5.5) \quad \forall \phi : \mathbb{R}_+ \rightarrow \mathbb{R}_+, \quad \int_{\mathbb{R}^{dN}} \phi(|V|^2) f_t^N(dV) = \int_{\mathbb{R}^{dN}} \phi(|V|^2) f_0^N(dV)$$

(equality between possibly infinite non-negative quantities).

The (expected) limiting nonlinear homogeneous Boltzmann equation is defined by (1.1), (1.2), (1.3). The equation generates a nonlinear semigroup

$$S_t^{NL}(f_0) := f_t \text{ for any } f_0 \in P_2(\mathbb{R}^d)$$

where $P_2(\mathbb{R}^d)$ denotes the space of probabilities with bounded second moment.

Concerning the Cauchy theory for the limiting Boltzmann equation:

- In the case **(GMM)** (Maxwell molecules with angular cutoff), see equation (1.6) in Subsection 1.1, we refer to [67];

- In the case **(tMM)** (true Maxwell molecules without angular cutoff), see equation (1.5) in Subsection 1.1, we refer to [69];
- In the case **(HS)** of hard spheres, see equation (1.4) in Subsection 1.1, we refer to [59] (L^1 theory) and [29, 32, 48] ($P_2(\mathbb{R}^d)$ theory).

For these solutions, one has the conservation of momentum and energy

$$\forall t \geq 0, \quad \int_{\mathbb{R}^d} v f_t(dv) = \int_{\mathbb{R}^d} v f_0(dv), \quad \int_{\mathbb{R}^d} |v|^2 f_t(dv) = \int_{\mathbb{R}^d} |v|^2 f_0(dv).$$

Observe that the change of variable

$$\sigma \in \mathbb{S}^{d-1} \mapsto -\sigma \in \mathbb{S}^{d-1}$$

maps the domain

$$\theta \in [-\pi, \pi/2] \cap [\pi/2, \pi] \quad \text{in} \quad \theta \in [-\pi/2, \pi/2].$$

Therefore without restriction we can consider, for the limiting equation as well as the N -particle system, kernel function b such that $\text{Supp } b \subset [0, 1]$. We still denote by b the symmetrized version of b by a slight abuse of notation.

In this section we aim at considering the case of the *Maxwell molecules kernel*. We shall indeed make the following general assumption:

$$(5.6) \quad \begin{cases} \Gamma \equiv 1, & b \in L_{\text{loc}}^\infty([0, 1]) \\ \forall \alpha > 0, & C_\alpha(b) := \int_{\mathbb{S}^{d-1}} b(\cos \theta) (1 - \cos \theta)^{\frac{1}{4} + \alpha} d\sigma < \infty. \end{cases}$$

Let us show that Maxwell molecules model (1.5) enters this general framework. Indeed for any positive real function ψ and any given vector $u \in \mathbb{R}^d$ we have

$$\int_{\mathbb{S}^{d-1}} \psi(\hat{u} \cdot \sigma) d\sigma = |\mathbb{S}^{d-2}| \int_0^\pi \psi(\cos \theta) \sin^{d-2} \theta d\theta.$$

Therefore the model (1.5) satisfies (in dimension $d = 3$)

$$b(z) \sim K (1 - z)^{-5/4} \quad \text{as} \quad z \rightarrow 1,$$

which hence fulfills (5.6). This assumption also trivially includes the Grad's cutoff Maxwell molecules model (1.6).

5.2. Statement of the results. Our main propagation of chaos estimate result for this model then states as follows:

Theorem 5.1 (Maxwell molecules detailed chaos estimates). *Let us consider a 1-particle initial distribution $f_0 \in P(\mathbb{R}^d)$. and a hierarchy of N -particle distributions*

$$f_t^N = S_t^N (f_0^{\otimes N})$$

as well the 1-particle of the limiting semigroup

$$f_t = S_t^{NL} (f_0)$$

where we assume that the collision satisfies (5.6).

Then for any $\eta > 0$, for any $\ell \in \mathbb{N}^*$ and for any

$$\varphi = \varphi_1 \otimes \varphi_2 \otimes \cdots \otimes \varphi_\ell \in \mathcal{F}^{\otimes \ell}, \quad \varphi_i \in \mathcal{F}, \quad i = 1, \dots, \ell,$$

where \mathcal{F} shall be specified below, we have:

- (i) Cases **(GMM)** and **(tMM)**: Consider a tensorized initial datum $f_0^N = f_0^{\otimes N}$ for the N -particle system and assume that f_0 has compact support, and take

$$\mathcal{F} := \left\{ \varphi : \mathbb{R}^d \rightarrow \mathbb{R}; \|\varphi\|_{\mathcal{F}} := \int_{\mathbb{R}^d} (1 + |\xi|^4) |\hat{\varphi}(\xi)| d\xi < \infty \right\}.$$

Then we have

$$(5.7) \quad \forall N \geq 2\ell, \quad \sup_{t \geq 0} \left| \left\langle \left(S_t^N(f_0^N) - (S_t^{NL}(f_0))^{\otimes N} \right), \varphi \right\rangle \right| \\ \leq C_\eta \left[\ell^2 \frac{\|\varphi\|_\infty}{N} + \frac{\ell^2}{N^{1-\eta}} \|\varphi\|_{\mathcal{F}^2 \otimes (L^\infty)^{\ell-2}} + \ell \|\varphi\|_{W^{1,\infty} \otimes (L^\infty)^{\ell-1}} \mathcal{W}_{W_2}^N(f_0) \right]$$

for some constant $C_\eta > 0$ (possibly blowing up as $\eta \rightarrow 0$) depending only η , on the collision kernel, and on the size of the support and some moments on f_0 .

We deduce the following rate of convergence as N goes infinity by using (4.12)-(4.13):

$$\sup_{t \geq 0} \left| \left\langle \left(S_t^N(f_0^N) - (S_t^{NL}(f_0))^{\otimes N} \right), \varphi \right\rangle \right| \leq \frac{C'_\eta \ell^2}{N^{\kappa(d,\eta)}} \|\varphi\|_{\mathcal{F}^{\otimes \ell}}$$

with

$$\kappa(d,\eta) := \begin{cases} \frac{1}{4} - \eta & \text{if } d \leq 2, \\ \frac{1}{6} - \eta & \text{if } d = 3, \\ \frac{1}{d+4} & \text{if } d \geq 4. \end{cases}$$

The constant $C'_\eta > 0$ may blow up when $\eta \rightarrow 0$, and depend on b and on f_0 through the size of its support and its moments.

- (ii) Case (GMM) with optimal rate for finite time: On a finite time interval $[0, T]$, the following variant is available: consider tensorized initial data $f_0^N = f_0^{\otimes N}$ for the N -particle system and assume that f_0 has compact support, and take $\mathcal{F} = H^s$ with $s > d/2$ high enough. Then we have

$$\sup_{0 \leq t \leq T} \left| \left\langle \left(S_t^N(f_0^N) - (S_t^{NL}(f_0))^{\otimes N} \right), \varphi \right\rangle \right| \\ \leq C \left[\ell^2 \frac{\|\varphi\|_\infty}{N} + C_{T,4}^N \frac{C_{\eta,\infty}^\infty}{N^{1-\eta}} \ell^2 \|\varphi\|_{\mathcal{F}^2 \otimes (L^\infty)^{\ell-2}} + \ell \|\varphi\|_{H^s \otimes (L^\infty)^{\ell-1}} \mathcal{W}_{H^{-s}}^N(f_0) \right].$$

By using (4.9) this proves

$$\sup_{t \in [0, T]} \left| \left\langle \left(S_t^N(f_0^N) - (S_t^{NL}(f_0))^{\otimes N} \right), \varphi \right\rangle \right| \leq \ell^2 \frac{C_{\eta,T}}{N^{1/2}} \|\varphi\|_{\mathcal{F}^{\otimes \ell}}$$

with $\mathcal{F} = H^s$, with the optimal rate of the law of large number.

- (iii) Cases (GMM) and (tMM) conditioned to the sphere: Finally consider \mathcal{F} as in (i), some initial data

$$f_0 \in P_4(\mathbb{R}^d) \cap L^\infty(\mathbb{R}^d)$$

and the N -particle initial data the sequence $(f_0^N)_{N \geq 1}$ constructed in Lemma 4.4 and 4.7 by conditioning to the energy sphere. Then we have

$$\sup_{t \geq 0} \left| \left\langle \left(S_t^N(f_0^N) - (S_t^{NL}(f_0))^{\otimes N} \right), \varphi \right\rangle \right| \\ C_\eta \left[\ell^2 \frac{\|\varphi\|_\infty}{N} + \frac{\ell^2}{N^{1-\eta}} \|\varphi\|_{\mathcal{F}^2 \otimes (L^\infty)^{\ell-2}} + \ell \|\varphi\|_{W^{1,\infty} \otimes (L^\infty)^{\ell-1}} \mathcal{W}_{W_2}^N(\pi_P^N f_0^N, \delta_{f_0}) \right]$$

which goes to zero as N goes infinity thanks to Lemma 4.7, and hence proves the propagation of chaos, uniformly in time.

Remark 5.2. Observe that by using the point (iii) below and taking $t \rightarrow \infty$ we deduce that

$$\gamma_\ell^N \xrightarrow{N \rightarrow \infty} \gamma^{\otimes \ell}$$

for any fixed $\ell \geq 1$, where γ_ℓ^N is the ℓ -marginal of the uniform measure on the energy sphere γ^N , and γ is the gaussian equilibrium with energy \mathcal{E} . However in our proof we need to build a sequence of initial data f_0^N which is supported on \mathcal{S}^N and chaotic. And the only proof we know makes use of the fact that γ^N is γ -chaotic. Nevertheless, this line of argument is not really based on an explicit computation nor on variational or entropy optimization characterization of the equilibrium. It only relies on the *action of the N -particle and limiting semigroups*. As a consequence, it can be applied to other situations where neither the steady states for the N -particle system nor the steady state for the mean-field equation are known explicitly and where a direct computation cannot be made (which is typical of open systems). We refer to the work in progress [57] for a application to some dissipative Boltzmann equation related to granular gases.

We now state the key Wasserstein version of the propagation of chaos estimate, which is valid *for any number of marginals*, although with a possibly worse (but still constructive) rate. Combined with previous results on the relaxation of the N -particle system we also deduce some estimate of relaxation to equilibrium *independent of N* and, again, for any number of marginals.

Theorem 5.3 (Maxwell molecules Wasserstein chaos). *Under the same setting as in Theorem 5.1, either*

- (a) *with f_0 compactly supported and $f_0^N = f_0^{\otimes N}$ or*
- (b) *with $f_0 \in P_4 \cap L^\infty$ and f_0^N constructed by Lemma 4.4,*

we have

$$(5.8) \quad \forall N \geq 1, \forall 1 \leq \ell \leq N, \quad \sup_{t \geq 0} \frac{W_1 \left(\Pi_\ell f_t^N, \left(f_t^{\otimes \ell} \right) \right)}{\ell} \leq \alpha(N)$$

for some $\alpha(N) \rightarrow 0$ as $N \rightarrow \infty$ (in the case (a) one has moreover explicit power law rate estimates on α).

Then in the case (b) we have

$$(5.9) \quad \forall N \geq 1, \forall 1 \leq \ell \leq N, \forall t \geq 0, \quad \frac{W_1 \left(\Pi_\ell f_t^N, \Pi_\ell (\gamma^N) \right)}{\ell} \leq \beta(t)$$

for some $\beta(t) \rightarrow 0$ as $t \rightarrow \infty$, where γ is the gaussian equilibrium with energy \mathcal{E} and γ^N is the uniform probability measure on $\mathcal{S}^N(\mathcal{E})$.

In order to prove Theorem 5.1, we have to establish the assumptions **(A1)**-**(A2)**-**(A3)**-**(A4)**-**(A5)** of Theorem 3.1 with $T = \infty$. The application of the latter theorem then exactly yields Theorem 5.1 by following carefully each constant computed below. Then the proof of Theorem 5.3 will be done in Subsection 5.9: it is deduced from Theorem 5.1 by using Lemma 4.1 together a result from [37].

5.3. Proof of (A1). When the collision kernel B is bounded the operator G^N is a linear bounded operator on $C(B_R)$ with $B_R := \{V \in \mathbb{R}^{dN}; |V| \leq R\}$ for any $R \in (0, \infty)$ with an operator norm independent of R . As a consequence, G^N is also well defined and bounded on

$$C_{-k,0}^0(\mathbb{R}^{dN}) := \left\{ \varphi \in C(\mathbb{R}^{dN}) \quad \text{s. t.} \quad \frac{\varphi(V)}{|V|^k} \rightarrow 0 \quad \text{as} \quad |V| \rightarrow \infty \right\}$$

endowed with the norm

$$\|\varphi\|_{L_{-k}^\infty} := \sup_V |\varphi(V)| \langle V \rangle^{-k}$$

for any $k \in \mathbb{R}$. It is also easy (and classical) to verify that G^N is dissipative in the sense that

$$\forall \varphi \in C_{-k,0}^0(\mathbb{R}^{dN}), \forall \lambda > 0 \quad \|(\lambda - G^N) \phi\|_{L_{-k}^\infty} \leq \lambda \|\phi\|_{L_{-k}^\infty}.$$

From the Hille-Yosida theory we deduce that G^N generates a Markov type semigroup T_t^N on $C_{-k,0}^0(\mathbb{R}^{dN})$ and we may also define S_t^N by duality as a semigroup on $P_k(\mathbb{R}^{dN})$. The nonlinear semigroup S_t^{NL} is also well defined on $P_k(\mathbb{R}^d)$, see for instance [69, 32, 29, 48].

For the true Maxwell molecules model, the operator G^N is not bounded anymore and some additional explanations are needed. The fastest way to argue is just to say this in that case B can be approximated by a sequence of bounded collision kernels

$$B_\varepsilon := b_\varepsilon(\cos \theta) \quad \text{with } b_\varepsilon \in L^\infty \quad \text{and } b_\varepsilon \nearrow b.$$

We may then define the associated generator $G^{N,\varepsilon}$, the associated semigroups $T_t^{N,\varepsilon}$ on continuous functions and $S_t^{N,\varepsilon}$ on probabilities and the nonlinear semigroup $S_t^{NL,\varepsilon}$ on probabilities. We first write estimate (5.7) for any fixed $\varepsilon > 0$. Then since (1) the right-hand side term in (5.7) does not depend on $\varepsilon > 0$ (as a consequence of the estimates established in the proof below) and (2) $S^{NL,\varepsilon}(f_0) \rightharpoonup S^{NL}(f_0)$ weakly in $P(\mathbb{R}^d)$ and (3) $S^{N,\varepsilon}(f_0^N) \rightharpoonup f_t^N$ weakly in $P(\mathbb{R}^{Nd})$, we can conclude that (5.7) holds for the true Maxwell molecules model by letting ε go to 0.

Possible other direct arguments (without using approximations) could be (1) to establish and use Wasserstein W_1 stability of the many-particle equation, or (2) use the following core $\mathcal{C} := W_{k+2}^{1,\infty}$ and prove that $\varphi \in W_{k+2}^{1,\infty}$ implies $G^N \varphi \in C_{k,0}$ (this follows from an easy decomposition between singular and non-singular angles in the formula for G^N).

Hence the semigroups T_t^N and

$$S_t^N = (T_t^N)^* = T_t^N$$

are well defined on $C_{-k,0}^0(\mathbb{R}^{dN})$. Moreover since for $\varphi \in L^2(\mathcal{S}^N)$ we have

$$\langle G^N \varphi, \varphi \rangle_{L^2(\mathcal{S}^N)} = -\frac{1}{N} \sum_{i,j=1}^N \int_{\mathcal{S}^N} \int_{\mathbb{S}^{d-1}} b(\cos \theta_{ij}) [\varphi_{ij}^* - \varphi]^2 d\sigma \gamma^N(dV) \leq 0$$

it is easily seen by arguing similarly as above that they are C_0 -semigroups of contractions on this space $L^2(\mathcal{S}^N)$.

Then it remains to prove bounds on the polynomial moments of the N -particle system. We shall prove the following more general lemma:

Lemma 5.4. *Consider the collision kernel*

$$B = |v - v_*|^\gamma b(\cos \theta) \quad \text{with } \gamma = 0 \text{ or } 1$$

and $b \geq 0$ such that

$$\int_0^1 b(z) (1-z)^2 dz < +\infty.$$

This covers the three cases **(HS)**, **(tMM)** and **(GMM)**.

Assume that the initial datum of the N -particle system satisfies:

$$\text{Supp } f_0^N \subset \left\{ V \in \mathbb{R}^{Nd}; M_2^N(V) \leq \mathcal{E}_0 \right\} \quad \text{where } M_2^N = \frac{1}{N} \sum_{j=1}^N |v_j|^2$$

and

$$\langle f_0^N, M_k^N \rangle \leq C_{0,k} < \infty \quad \text{where } M_k^N = \frac{1}{N} \sum_{j=1}^N |v_j|^k, \quad k \geq 2.$$

Then we have

$$\sup_{t \geq 0} \langle f_t^N, M_k^N \rangle \leq \max \{C_{0,k}; \bar{a}_k\}$$

where $\bar{a}_k \in (0, \infty)$ depends on k and \mathcal{E}_0 .

Proof of Lemma 5.4. By using (5.5) with the function of the energy

$$\phi(z) := \mathbf{1}_{z > N \mathcal{E}_0}$$

and the assumption on f_0^N we deduce

$$(5.10) \quad \forall t \geq 0, \quad \text{Supp } f_t^N \subset \left\{ V \in \mathbb{R}^{Nd}; M_2^N(V) \leq \mathcal{E}_0 \right\}.$$

Next, we write the differential equality on the k -th moment:

$$\frac{d}{dt} \left\langle f_t^N, \frac{1}{N} \sum_{j=1}^N |v_j|^k \right\rangle = \frac{1}{N^2} \sum_{j_1 \neq j_2}^N \langle f_t^N, |v_{j_1} - v_{j_2}|^\gamma \mathcal{K}(v_{j_1}, v_{j_2}) \rangle,$$

with

$$\mathcal{K}(v_{j_1}, v_{j_2}) = \frac{1}{2} \int_{\mathbb{S}^{d-1}} b(\theta_{j_1 j_2}) \left[|v_{j_1}^*|^k + |v_{j_2}^*|^k - |v_{j_1}|^k - |v_{j_2}|^k \right] d\sigma.$$

We then apply the so-called *Povner's Lemma* proved in [59, Lemma 2.2] (valid for singular collision kernel as in our case) which implies

$$\mathcal{K}(v_{j_1}, v_{j_2}) \leq C_1 \left(|v_{j_1}|^{k-1} |v_{j_2}| + |v_{j_1}| |v_{j_2}|^{k-1} \right) - C_2 \left(|v_{j_1}|^k + |v_{j_2}|^k \right)$$

for some constants $C_1, C_2 \in (0, \infty)$ depending only on k and b .

By using the inequalities $|v_{j_1} - v_{j_2}| \geq |v_{j_1}| - |v_{j_2}|$ and $|v_{j_1} - v_{j_2}| \geq |v_{j_2}| - |v_{j_1}|$ in order to estimate the last term when $\gamma = 1$, we then deduce

$$\begin{aligned} |v_{j_1} - v_{j_2}|^\gamma \mathcal{K}(v_{j_1}, v_{j_2}) &\leq C_3 \left[(1 + |v_{j_1}|^{k+\gamma-1}) (1 + |v_{j_2}|^2) \right. \\ &\quad \left. + (1 + |v_{j_1}|^2) (1 + |v_{j_2}|^{k+\gamma-1}) \right] - C_2 (|v_{j_1}|^{k+\gamma} + |v_{j_2}|^{k+\gamma}), \end{aligned}$$

for a constant C_3 depending on C_1 and C_2 .

Using (symmetry hypothesis) that

$$\forall k \geq 0, \quad \left\langle f_t^N, |v_1|^k \right\rangle = \left\langle f_t^N, M_k^N \right\rangle,$$

and (5.10) we get

$$\begin{aligned} \frac{d}{dt} \left\langle f_t^N, |v_1|^k \right\rangle &\leq 2C_3 \left\langle f_t^N, (1 + M_{k+\gamma-1}^N) (1 + M_2^N) \right\rangle - 2C_2 \left\langle f_t^N, M_{k+\gamma}^N \right\rangle \\ &\leq 2C_3 (1 + \mathcal{E}) \left(1 + \left\langle f_t^N, |v_1|^{k+\gamma-1} \right\rangle \right) - 2C_2 \left\langle f_t^N, |v_1|^{k+\gamma} \right\rangle. \end{aligned}$$

Using finally Hölder's inequality

$$\left\langle f_1^N, |v|^{k-\gamma+1} \right\rangle \leq \left\langle f_1^N, |v|^{k+\gamma} \right\rangle^{(k-\gamma+1)/(k+\gamma)}$$

we conclude that $y(t) = \langle f_t^N, |v_1|^k \rangle$ satisfies a differential inequality of the following kind

$$y' \leq -K_1 y^{\theta_1} + K_2 y^{\theta_2} + K_3$$

with $\theta_1 \geq 1$ and $\theta_2 < \theta_1$, and for some constants $K_1, K_2, K_3 > 0$, which concludes the proof of the lemma. \square

Lemma 5.4 proves **(A1)-(i)** with

$$m_{\mathcal{G}_1}(v) := \langle v \rangle^6.$$

Moreover we do not need **(A1)-(ii)** in the present case since we may take $m_{\mathcal{G}_3} \equiv 0$.

5.4. **Proof of (A2).** Let us define

$$P_{\mathcal{G}_1} := \left\{ f \in P(\mathbb{R}^d); \langle f, m_e \rangle \leq \mathcal{E} \text{ and } \langle f, m_{\mathcal{G}_1} \rangle < +\infty \right\}$$

endowed with the distance induced by $|\cdot|_2$, as well as

$$\mathcal{B}P_{\mathcal{G}_1, a} := \{f \in P_{\mathcal{G}_1}; \langle f, m_{\mathcal{G}_1} \rangle \leq a\}.$$

Let us recall the following result proved in [67, 33, 14, 69]. We briefly outline its proof for the sake of completeness and most importantly because we shall need to modify it in order to adapt it to our purpose in the next sections.

Lemma 5.5. *For any $f_0, g_0 \in P_2(\mathbb{R}^d)$, the associated solutions f_t and g_t to the Boltzmann equation for Maxwell molecules satisfy*

$$(5.11) \quad \sup_{t \geq 0} |f_t - g_t|_2 \leq |f_0 - g_0|_2.$$

Moreover, there exists $\bar{a} \in (0, \infty)$ such that

$$\forall a \in [\bar{a}, \infty), \quad S_t^{NL} : \mathcal{B}P_{\mathcal{G}_1, a} \rightarrow \mathcal{B}P_{\mathcal{G}_1, a}.$$

Proof of Lemma 5.5. We only prove (5.11) and we refer to the above quoted references for the moment estimates in the statement of the Lemma (which is nothing but a variation on the arguments presented in the proof of Lemma 5.4).

We recall Bobylev's identity for maxwellian collision kernel (cf. [6])

$$\mathcal{F}(Q^+(f, g))(\xi) = \hat{Q}^+(F, G)(\xi) =: \frac{1}{2} \int_{S^{d-1}} b(\sigma \cdot \hat{\xi}) [F^+ G^- + F^- G^+] d\sigma,$$

with

$$F = \hat{f}, \quad G = \hat{g}, \quad F^\pm = F(\xi^\pm), \quad G^\pm = G(\xi^\pm), \quad \hat{\xi} = \frac{\xi}{|\xi|}$$

and

$$\xi^+ = \frac{1}{2}(\xi + |\xi| \sigma), \quad \xi^- = \frac{1}{2}(\xi - |\xi| \sigma).$$

With the shorthand notation $D = \hat{g} - \hat{f}$, $S = \hat{g} + \hat{f}$, the following equation holds

$$(5.12) \quad \partial_t D = \hat{Q}(S, D) = \int_{S^{d-1}} b(\sigma \cdot \hat{\xi}) \left[\frac{D^+ S^-}{2} + \frac{D^- S^+}{2} - D \right] d\sigma.$$

We perform the following cutoff decomposition on the angular collision kernel:

$$b = b_K + b_K^c \quad \text{with} \quad \int_{S^{d-1}} b_K(\sigma \cdot \hat{\xi}) d\sigma = K, \quad b_K = b \mathbf{1}_{|\theta| \geq \delta(K)}$$

for some well-chosen $\delta(K)$. As in [69] observe that

$$R_K(\xi) = \int_{S^{d-1}} b_K^c(\sigma \cdot \hat{\xi}) \left[\frac{D^+ S^-}{2} + \frac{D^- S^+}{2} - D \right] d\sigma$$

satisfies

$$\forall \xi \in \mathbb{R}^d, \quad |R_K(\xi)| \leq r_K |\xi|^2 \quad \text{where} \quad r_K \xrightarrow{K \rightarrow \infty} 0$$

and r_K depends on moments of order 2 on d and s (hence bounded by the energy).

Using that $\|S\|_\infty \leq 2$, we deduce in distributional sense

$$\frac{d}{dt} \frac{|D|}{|\xi|^2} + K \frac{|D|}{|\xi|^2} \leq \left(\sup_{\xi \in \mathbb{R}^d} \frac{|D|}{|\xi|^2} \right) \left(\sup_{\xi \in \mathbb{R}^d} \int_{S^{d-1}} b_K(\sigma \cdot \hat{\xi}) \left(|\hat{\xi}^+|^2 + |\hat{\xi}^-|^2 \right) d\sigma \right) + r_K$$

with

$$|\hat{\xi}^+| = \frac{1}{\sqrt{2}} \left(1 + \sigma \cdot \hat{\xi} \right)^{1/2}, \quad |\hat{\xi}^-| = \frac{1}{\sqrt{2}} \left(1 - \sigma \cdot \hat{\xi} \right)^{1/2}.$$

By using

$$|\hat{\xi}^+|^2 + |\hat{\xi}^-|^2 = 1,$$

we deduce

$$\frac{d}{dt} \frac{|D|}{|\xi|^2} + K \frac{|D|}{|\xi|^2} \leq K \left(\sup_{\xi \in \mathbb{R}^d} \frac{|D|}{|\xi|^2} \right) + r_K$$

which implies

$$\sup_{\xi \in \mathbb{R}^d} \frac{|D_t(\xi)|}{|\xi|^2} \leq \sup_{\xi \in \mathbb{R}^d} \frac{|D_0(\xi)|}{|\xi|^2} + C \frac{r_K}{K}$$

for any value of the cutoff parameter K . Therefore by relaxing $K \rightarrow \infty$, we deduce (5.11). \square

Hence we deduce that $S_t^{NL} \in C^{0,1}(P_{\mathcal{G}_1}, P_{\mathcal{G}_1})$ and **(A2)-(i)** is proved.

Lemma 5.6. *For any $f, g \in P(\mathbb{R}^d)$ with finite second moment*

$$\int_{\mathbb{R}^d} f |v|^2 dv < \infty, \quad \int_{\mathbb{R}^d} g |v|^2 dv < \infty$$

and same momentum

$$\int_{\mathbb{R}^d} f v_i dv = \int_{\mathbb{R}^d} g v_i dv, \quad i = 1, \dots, d,$$

we have

$$(5.13) \quad |Q(f, f)|_2 \leq C \left(\int_{\mathbb{R}^d} (1 + |v|^2) df(v) \right)^2$$

and

$$(5.14) \quad |Q(f + g, f - g)|_2 \leq C \left(\int_{\mathbb{R}^d} (1 + |v|) (df(v) + dg(v)) \right) (|f - g|_2 + |(f - g)v|_1).$$

Proof of Lemma 5.5. We prove the second inequality (5.14). The first inequalities (5.13) then follows immediately by writing

$$Q(f, f) = Q(f, f) - Q(M, M) = Q(f - M, f + M)$$

where M is the maxwellian distribution with same momentum and energy as f , and then applying (5.14) with $f - M$ and $f + M$.

We write in Fourier:

$$\begin{aligned} \mathcal{F}(Q(f + g, f - g)) &= \hat{Q}(D, S) \\ &= \frac{1}{2} \int_{\mathbb{S}^{d-1}} b(\sigma \cdot \hat{\xi}) (S(\xi^+) D(\xi^-) + S(\xi^-) D(\xi^+) - 2D(\xi)) \end{aligned}$$

where \hat{Q} is the Fourier form the symmetrized collision operator Q , which we can rewrite

$$\frac{|\hat{Q}(D, S)|}{|\xi|^2} \leq \mathcal{T}_1 + \mathcal{T}_2 + \mathcal{T}_3.$$

We have

$$\mathcal{T}_1 \leq \int_{\mathbb{S}^{d-1}} b(\sigma \cdot \hat{\xi}) |S(\xi^+)| \frac{|D(\xi^-)|}{|\xi^-|^2} \frac{|\xi^-|^2}{|\xi|^2} d\sigma \leq C |D|_2$$

for some constant $C > 0$, where we have used

$$\frac{|\xi^-|^2}{|\xi|^2} = (1 - \cos \theta)^2$$

which permits to control the angular singularity of b .

Similarly we compute

$$\mathcal{T}_2 \leq \int_{\mathbb{S}^{d-1}} b(\sigma \cdot \hat{\xi}) \frac{|D(\xi^+)|}{|\xi^+|} \frac{|S(\xi^-) - 2|}{|\xi^-|} \frac{|\xi^-|}{|\xi|} d\sigma \leq C |D|_1 \left(\int_{\mathbb{R}^d} (1 + |v|) (df(v) + dg(v)) \right)$$

for some constant $C > 0$, and

$$\begin{aligned} \mathcal{T}_3 &\leq 2 \int_{\mathbb{S}^{d-1}} b(\sigma \cdot \hat{\xi}) \frac{|D(\xi^+) - D(\xi)|}{|\xi|} d\sigma \\ &\leq \int_{\mathbb{S}^{d-1}} b(\sigma \cdot \hat{\xi}) \frac{|\xi^-|}{|\xi|} \int_0^1 \frac{|\nabla D(\theta\xi + (1-\theta)\xi^+)|}{|\theta\xi + (1-\theta)\xi^+|} d\theta d\sigma \leq C |(f-g)v|_1 \end{aligned}$$

for some constant $C > 0$. This concludes the proof of (5.14) by piling these estimates. \square

The proof of **(A2)-(ii)** is a consequence of (5.14). Indeed, from (4.2) we have

$$\forall f, g \in \mathcal{BP}_{\hat{G}, a}, \quad |(f-g)v|_1 \leq C [(f-g)v]_1^*.$$

Moreover for $f, g \in \mathcal{BP}_{\mathcal{G}_1, a}$,

$$\begin{aligned} [(f-g)v]_1^* &\leq \inf_{R>0} \left\{ [(f-g)v \chi_R]_1^* + [(f-g)v(1-\chi_R)]_1^* \right\} \\ &\leq \inf_{R>0} \left\{ C R [f-g]_1^* + C' \frac{a}{R^2} \right\} \\ &\leq C a^{1/3} ([f-g]_1^*)^{2/3} \end{aligned}$$

where we have used the bound on the fourth moment.

Finally, from (4.4), we conclude that

$$\forall f, g \in \mathcal{BP}_{\mathcal{G}_1, a}, \quad [f-g]_1^* \leq C |f-g|_2^\nu,$$

for some constants $C > 0$ and $\nu \in (0, 1)$ depending on d and a .

Gathering all these estimates as well as (5.14) we deduce that

$$\begin{aligned} |Q(f+g, f-g)|_2 &\leq C \left(\int_{\mathbb{R}^d} (1 + |v|) (df(v) + dg(v)) \right) (|f-g|_2 + |f-g|_2^\nu). \end{aligned}$$

5.5. Proof of (A3). Let us define $m'_{\mathcal{G}_1}$ and

$$\Lambda_1(f) := \langle f, m'_{\mathcal{G}_1} \rangle = \langle f, \langle v \rangle^4 \rangle$$

for any $f \in P_{\mathcal{G}_1}$.

Let us prove that for any

$$\mathbf{r} \in \mathbf{R}_{\mathcal{G}_1} \quad \text{and} \quad \Phi \in C_{\Lambda_1}^{1, \eta}(P_{\mathcal{G}_1, \mathbf{r}})$$

we have

$$\left\| \left(M_{m_{\mathcal{G}_1}}^N(V) \right)^{-1} (G^N \pi^N - \pi^N G^\infty) \Phi \right\|_{L^\infty(E_N)} \leq \frac{C_1 \mathcal{E}}{N} [\Phi]_{C_{\Lambda_1}^{1, \eta}(P_{\mathcal{G}_1, \mathbf{r}})},$$

for some constant $C_1 > 0$.

First, consider velocities $v, v_*, w, w_* \in \mathbb{R}^d$ such that

$$w = \frac{v + v_*}{2} + \frac{|v - v_*|}{2} \sigma, \quad w_* = \frac{v + v_*}{2} - \frac{|v - v_*|}{2} \sigma, \quad \sigma \in \mathbb{S}^{d-1}.$$

Then $\delta_v + \delta_{v_*} - \delta_w - \delta_{w_*} \in \mathcal{I}P_{\mathcal{G}_1}$. Performing Taylor expansions, we get

$$\begin{aligned} & e^{iv \cdot \xi} + e^{iv_* \cdot \xi} - e^{iw \cdot \xi} - e^{iw_* \cdot \xi} \\ &= i(w - v) \xi e^{iv \cdot \xi} + \mathcal{O}(|w - v|^2 |\xi|^2) + i(w_* - v_*) \xi e^{iv_* \cdot \xi} + \mathcal{O}(|w_* - v_*|^2 |\xi|^2) \\ &= i(w - v) \xi e^{iv \cdot \xi} + \mathcal{O}(|w - v|^2 |\xi|^2) \\ &\quad + i(w_* - v_*) \xi (e^{iv \cdot \xi} + \mathcal{O}(|v - v_*| |\xi|)) + \mathcal{O}(|w_* - v_*|^2 |\xi|^2) \\ &= \mathcal{O}(|v - v_*|^2 |\xi|^2 \sin \theta / 2) \end{aligned}$$

thanks to the impulsion conservation and the fact that

$$|w - v| = |w_* - v_*| = |v - v_*| \sin \frac{\theta}{2}.$$

We hence deduce

$$|\delta_v + \delta_{v_*} - \delta_w - \delta_{w_*}|_2 = \sup_{\xi \in \mathbb{R}^d} \frac{|e^{iv \cdot \xi} + e^{iv_* \cdot \xi} - e^{iw \cdot \xi} - e^{iw_* \cdot \xi}|}{|\xi|^2} \leq C |v - v_*|^2 (1 - \cos \theta).$$

As an immediate consequence, for any $V \in E^N$ and V_{ij}^* defined by (5.2), we have

$$\left| \mu_{V_{ij}^*}^N - \mu_V^N \right|_2 \leq \frac{C}{N} |v_i - v_j|^2 (1 - \cos \theta_{ij}).$$

Consider $V \in \mathbb{E}_N$ and define

$$\mathbf{r}_V := (\langle \mu_V^N, z_1 \rangle, \dots, \langle \mu_V^N, z_d \rangle, \langle \mu_V^N, |z|^2 \rangle) \in \mathbf{R}_{\mathcal{E}}.$$

Then for a given $\Phi \in C_{\Lambda_1}^{1,\eta}(P_{\mathcal{G}_1, \mathbf{r}_V})$, we set

$$\phi := D\Phi [\mu_V^N] \quad \text{and} \quad u_{ij} = (v_i - v_j)$$

and we compute:

$$\begin{aligned} G^N(\Phi \circ \mu_V^N) &= \frac{1}{2N} \sum_{i,j=1}^N \int_{\mathbb{S}^{d-1}} \left[\Phi(\mu_{V_{ij}^*}^N) - \Phi(\mu_V^N) \right] b(\cos \theta_{ij}) \, d\sigma \\ &= \frac{1}{2N} \sum_{i,j=1}^N \int_{\mathbb{S}^{d-1}} \langle \mu_{V_{ij}^*}^N - \mu_V^N, \phi \rangle b(\cos \theta_{ij}) \, d\sigma \\ &\quad + \frac{[\Phi]_{C_{\Lambda_1}^{1,\eta}(P_{\mathcal{G}_1, \mathbf{r}_V})}}{2N} \sum_{i,j=1}^N \int_{\mathbb{S}^{d-1}} \left[M_{m'_{\mathcal{G}_1}}(\mu_{V_{ij}^*}^N) + M_{m'_{\mathcal{G}_1}}(\mu_V^N) \right] \mathcal{O}\left(\left| \mu_{V_{ij}^*}^N - \mu_V^N \right|_2^{1+\eta} \right) \, d\sigma \\ &=: I_1(V) + I_2(V). \end{aligned}$$

Concerning the first term $I_1(V)$, thanks to Lemma 2.13, we have

$$\begin{aligned} I_1(V) &= \frac{1}{2N^2} \sum_{i,j=1}^N \int_{\mathbb{S}^{d-1}} b(\cos \theta_{ij}) [\phi(v_i^*) + \phi(v_j^*) - \phi(v_i) - \phi(v_j)] \, d\sigma \\ &= \frac{1}{2} \int_v \int_w \int_{\mathbb{S}^{d-1}} b(\cos \theta) [\phi(v^*) + \phi(w^*) - \phi(v) - \phi(w)] \mu_V^N(dv) \mu_V^N(dw) \, d\sigma \\ &= \langle Q(\mu_V^N, \mu_V^N), \phi \rangle = (G^\infty \Phi)(\mu_V^N). \end{aligned}$$

For the second term $I_2(V)$, using that

$$\begin{aligned}
M_{m'_{\mathcal{G}_1}} \left(\mu_{V_{ij}^*}^N \right) &:= \frac{1}{N} \sum_{\ell=1}^N m'_{\mathcal{G}_1} \left((V_{ij}^*)_{\ell} \right) \\
&\leq C \left(1 + \frac{1}{N} \left(\left(\sum_{\ell \neq i,j} |v_{\ell}|^4 \right) + |v_i^*|^4 + |v_j^*|^4 \right) \right) \\
&\leq C \left(1 + \frac{1}{N} \left(\left(\sum_{\ell \neq i,j} |v_{\ell}|^4 \right) + 2 (|v_i|^2 + |v_j|^2)^2 \right) \right) \\
&\leq C \left(1 + \frac{2^3}{N} \left(\sum_{\ell=1}^N |v_{\ell}|^4 \right) \right) \leq C M_{m'_{\mathcal{G}_1}} \left(\mu_V^N \right) = C M_{m'_{\mathcal{G}_1}}^N (V),
\end{aligned}$$

we deduce

$$\begin{aligned}
|I_2(V)| &\leq \frac{C}{N^{2+\eta}} M_{m'_{\mathcal{G}_1}}^N (V) [\Phi]_{C_{\Lambda}^{1,\eta}(P_{\mathcal{G}_1, r_V})} \\
&\quad \times \sum_{i,j=1}^N \mathcal{O} \left(\int_{\mathbb{S}^{d-1}} b(\cos \theta_{ij}) (1 + |v_i|^2 + |v_j|^2) (1 - \sigma \cdot \hat{u}_{ij}) d\sigma \right).
\end{aligned}$$

We finally use that

$$\begin{aligned}
\frac{1}{N^2} M_{m'_{\mathcal{G}_1}}^N (V) \mathcal{O} \left(\int_{\mathbb{S}^{d-1}} b(\cos \theta_{ij}) (1 + |v_i|^2 + |v_j|^2) (1 - \sigma \cdot \hat{u}_{ij}) d\sigma \right) \\
\leq C M_{m'_{\mathcal{G}_1}}^N (V) \left(\frac{1}{N} \sum_{i=1}^N |v_i|^2 \right) \leq C M_{m_{\mathcal{G}_1}}^N (V)
\end{aligned}$$

(recall for the last line that $m_{\mathcal{G}_1}(v) = \langle v \rangle^6$ and $m'_{\mathcal{G}_1}(v) = \langle v \rangle^4$), which implies

$$|I_2(V)| \leq \frac{C}{N^{\eta}} M_{m_{\mathcal{G}_1}}^N (V) [\Phi]_{C_{\Lambda}^{1,\eta}(P_{\mathcal{G}_1, r_V})}$$

and concludes the proof.

5.6. Proof of (A4) uniformly in time. Let us consider some 1-particle initial data $f_0, g_0 \in P_4(\mathbb{R}^d)$ (probability with bounded fourth moment). And let us define the associated solutions f_t and g_t to the nonlinear Boltzmann equation (1.1) under the assumption (5.6). We then define

$$h_t := \mathcal{L}_t^{NL} [f_0] (g_0 - f_0)$$

the solution to the linearized Boltzmann equation around f_t , given by

$$\begin{cases} \partial_t f_t = Q(f_t, f_t), & f_{|t=0} = f_0 \\ \partial_t g_t = Q(g_t, g_t), & g_{|t=0} = g_0 \\ \partial_t h_t = 2Q(h_t, f_t), & h_{|t=0} = h_0 := g_0 - f_0. \end{cases}$$

We shall prove assumption **(A4)** with the choice of indices

$$(\eta', \eta'') = (\eta, 1)$$

which means that the weight which has to be used is

We shall now *expand the limiting nonlinear semigroup* in terms of the initial data, around f_0 .

Lemma 5.7. *There exists $\lambda \in (0, \infty)$ and, for any $\eta \in (0, 1)$, there exists $C_\eta > 0$ such that for any*

$$\mathbf{r} \in \mathbf{R}_{\mathcal{G}_1} \quad \text{and} \quad f_0, g_0 \in P_{\mathcal{G}_1, \mathbf{r}}$$

we have

$$(5.15) \quad |f_t - g_t|_2 \leq C_\eta e^{-(1-\eta)\lambda t} (\max\{M_4(f_0), M_4(g_0)\})^{\frac{1}{1+\eta}} |f_0 - g_0|_2^\eta,$$

and

$$(5.16) \quad |h_t|_2 \leq C_\eta e^{-(1-\eta)\lambda t} (\max\{M_4(f_0), M_4(g_0)\})^{\frac{1}{1+\eta}} |f_0 - g_0|_2^\eta$$

where we recall that

$$\forall f \in P(\mathbb{R}^d), \quad M_4(f) := \langle f, \langle v \rangle^4 \rangle.$$

As a consequence, the operator $\mathcal{L}_t^{NL}[f_0](\cdot)$ is bounded from $\mathcal{G}_{1, \mathbf{r}}$ to itself.

Remark 5.8. Observe that the loss of weight exactly matches the need of assumption **(A4)** since

$$\Lambda_2 = \Lambda_1^{\frac{1}{1+\eta}}$$

(recall that we have made the choice the indices $(\eta', \eta'') = (\eta, 1)$ in **(A4)** in the present application).

Proof of Lemma 5.7. We shall proceed in several steps.

Step 1. Estimate in $|\cdot|_4$. We closely follow ideas in [67, 14]. We shall use the notation

$$\mathcal{M} = \mathcal{M}_4, \quad \hat{\mathcal{M}} = \hat{\mathcal{M}}_4,$$

introduced in Example 2.5.4, as well as

$$d := f - g, \quad s := f + g$$

and

$$\tilde{d} := d - \mathcal{M}[d], \quad D := \mathcal{F}(d), \quad S := \mathcal{F}(s) \quad \text{and} \quad \tilde{D} := \mathcal{F}(\tilde{d}) = D - \hat{\mathcal{M}}[d].$$

The equation satisfied by \tilde{D} is

$$(5.17) \quad \begin{aligned} \partial_t \tilde{D} &= \hat{Q}(D, S) - \partial_t \hat{\mathcal{M}}[d] \\ &= \hat{Q}(\tilde{D}, S) + \left(\hat{Q}(\hat{\mathcal{M}}[d], S) - \hat{\mathcal{M}}[Q(d, s)] \right). \end{aligned}$$

We infer from [67] that for any $\alpha \in \mathbb{N}^d$, there exists some absolute coefficients $(a_{\alpha, \beta})$, $\beta \leq \alpha$ (which means $\beta_i \leq \alpha_i$ for any $1 \leq i \leq d$), depending on the collision kernel b through

$$(5.18) \quad \int_{\mathbb{S}^{d-1}} b(\cos \theta) [(v^\alpha)' + (v^\alpha)'_* - (v^\alpha) - (v^\alpha)_*] d\sigma = \sum_{\beta, \beta \leq \alpha} a_{\alpha, \beta} (v^\beta) (v^{\alpha-\beta})_*$$

where $\alpha, \beta \in \mathbb{N}^d$ are *coordinates* indices and

$$v^\alpha := v_1^{\alpha_1} v_2^{\alpha_2} \dots v_d^{\alpha_d}.$$

These multi-indices are compared through the usual lexicographical order, and we use the standart notation

$$|\alpha| := \sum_{k=1}^d \alpha_k.$$

We deduce that

$$\forall |\alpha| \leq 3, \quad \nabla_\xi^\alpha \hat{\mathcal{M}}[Q(d, s)]|_{\xi=0} = M_\alpha[Q(d, s)] = \sum_{\beta, \beta \leq \alpha} a_{\alpha, \beta} M_\beta[d] M_{\alpha-\beta}[s]$$

together with

$$\begin{aligned} \forall |\alpha| \leq 3, \quad \nabla_{\hat{\xi}}^{\alpha} \hat{Q}(\hat{\mathcal{M}}[d], S) \Big|_{\xi=0} &= M_{\alpha}[Q(\mathcal{M}[d], s)] \\ &= \sum_{\beta, \beta \leq \alpha} a_{\alpha, \beta} M_{\beta}[\mathcal{M}[d]] M_{\alpha-\beta}[s] = \sum_{\beta, \beta \leq \alpha} a_{\alpha, \beta} M_{\beta}[d] M_{\alpha-\beta}[s] \end{aligned}$$

since

$$M_{\alpha}[\mathcal{M}[d]] = M_{\alpha}[d]$$

for any $|\alpha| \leq 3$ by construction. As a consequence, we get

$$(5.19) \quad \forall \xi \in \mathbb{R}^d, \quad \left| \hat{\mathcal{M}}[Q(d, s)] - \hat{Q}(\hat{\mathcal{M}}[d], S) \right| \leq C |\xi|^4 \left(\sum_{|\alpha| \leq 3} |M_{\alpha}[f_t - g_t]| \right).$$

On the other hand, from [67, Theorem 8.1] and its corollary, we know that there exists some constants $C, \lambda \in (0, \infty)$ such that

$$(5.20) \quad \forall t \geq 0, \quad \left(\sum_{|\alpha| \leq 3} |M_{\alpha}[f_t - g_t]| \right) \leq C e^{-\lambda t} \left(\sum_{|\alpha| \leq 3} |M_{\alpha}[f_0 - g_0]| \right).$$

We perform the same decomposition on the angular collision kernel

$$b = b_K + b_K^c \quad \text{with} \quad \int_{\mathbb{S}^{d-1}} b_K(\sigma \cdot \hat{\xi}) \, d\sigma = K, \quad b_K = b \mathbf{1}_{|\theta| \geq \delta(K)}$$

as in the proof of Lemma 5.5 and use the straightforward estimate

$$R_K(\xi) := \hat{Q}_{b_K^c}(\tilde{D}, S)(\xi)$$

satisfies

$$\forall \xi \in \mathbb{R}^d, \quad |R_K(\xi)| \leq r_K |\xi|^4 \quad \text{where} \quad r_K \xrightarrow{K \rightarrow \infty} 0$$

where $Q_{b_K^c}$ denotes the collision operator associated with the part b_K^c of the decomposition of the angular collision kernel, and where r_K depends on moments of order 4 on d and s .

Then we gather (5.17), (5.19) and (5.20) and we have

$$\begin{aligned} \frac{d}{dt} \frac{|\tilde{D}(\xi)|}{|\xi|^4} + K \frac{|\tilde{D}(\xi)|}{|\xi|^4} &\leq \left(\sup_{\xi \in \mathbb{R}^d} \frac{|\tilde{D}(\xi)|}{|\xi|^4} \right) \left(\sup_{\xi \in \mathbb{R}^d} \int_{\mathbb{S}^{d-1}} b_K(\sigma \cdot \hat{\xi}) \left(|\hat{\xi}^+|^4 + |\hat{\xi}^-|^4 \right) d\sigma \right) \\ &\quad + C e^{-\lambda t} \left(\sum_{|\alpha| \leq 3} |M_{\alpha}[f_0 - g_0]| \right) + r_K. \end{aligned}$$

Let us compute (the supremum has been dropped thanks to the spherical invariance)

$$\lambda_K := \int_{\mathbb{S}^{d-1}} b_K(\sigma \cdot \hat{\xi}) \left(|\hat{\xi}^+|^4 + |\hat{\xi}^-|^4 \right) d\sigma = \int_{\mathbb{S}^{d-1}} b_K(\sigma \cdot \hat{\xi}) \left(\frac{1 + (\sigma \cdot \hat{\xi})^2}{2} \right) d\sigma,$$

so that

$$\begin{aligned} \lambda_K - K &= - \left(\int_{\mathbb{S}^{d-1}} b_K(\sigma \cdot \hat{\xi}) \left(\frac{1 - (\sigma \cdot \hat{\xi})^2}{2} \right) d\sigma \right) \\ &\xrightarrow{K \rightarrow \infty} - \int_{\mathbb{S}^{d-1}} b(\sigma \cdot \hat{\xi}) \left(\frac{1 - (\sigma \cdot \hat{\xi})^2}{2} \right) d\sigma := -\bar{\lambda} \in (-\infty, 0) \end{aligned}$$

where in the last step we used the factor

$$\left(\frac{1 - (\sigma \cdot \hat{\xi})^2}{2} \right) \sim C \theta^2 \text{ as } \theta \sim 0$$

in order to control the singularity of b .

Then, thanks to Gronwall lemma, we get

$$\begin{aligned} \left(\sup_{\xi \in \mathbb{R}^d} \frac{|\tilde{D}_t(\xi)|}{|\xi|^4} \right) &\leq e^{(\lambda_K - K)t} \left(\sup_{\xi \in \mathbb{R}^d} \frac{|\tilde{D}_0(\xi)|}{|\xi|^4} \right) \\ &+ C_3 \left(\sum_{|\alpha| \leq 3} |M_\alpha[f_0 - g_0]| \right) \left(\frac{e^{-\lambda t}}{K - \lambda_K - \lambda} - \frac{e^{(\lambda_K - K)t}}{K - \lambda_K - \lambda} \right) + C \frac{r_K}{K(K - \lambda_K)}. \end{aligned}$$

Therefore, passing to the limit $K \rightarrow \infty$ and choosing (without restriction) $\lambda \in (0, \bar{\lambda})$, we obtain

$$\sup_{\xi \in \mathbb{R}^d} \frac{|\tilde{D}_t(\xi)|}{|\xi|^4} \leq C e^{-\lambda t} \left(\sup_{\xi \in \mathbb{R}^d} \frac{|\tilde{D}_0(\xi)|}{|\xi|^4} + \sum_{|\alpha| \leq 3} |M_\alpha[f_0 - g_0]| \right)$$

or equivalently (and with the notations of Example 2.5.4),

$$\|d_t\|_4 \leq C e^{-\lambda t} \|d_0\|_4.$$

Step 2. From $|\cdot|_4$ to $|\cdot|_2$ on the difference. From the preceding step and a straightforward interpolation argument, we have

$$\begin{aligned} |f - g|_2 &\leq |f - g - \mathcal{M}[f - g]|_2 + C \left(\sum_{|\alpha| \leq 3} |M_\alpha[f - g]| \right) \\ &\leq \|f - g - \mathcal{M}[f - g]\|_{M_4}^{1/2} \|f - g - \mathcal{M}[f - g]\|_4^{1/2} + C \left(\sum_{|\alpha| \leq 3} |M_\alpha[f - g]| \right) \\ &\leq C (1 + M_4(f_0) + M_4(g_0)) e^{-(\lambda/2)t}. \end{aligned}$$

Then by writing

$$|f - g|_2 \leq |f - g|_2^\eta |f - g|_2^{1-\eta},$$

using Lemma 5.11 for the first term of the right hand side and the previous decay estimates for the second term, we obtain

$$|f_t - g_t|_2 \leq C_\eta e^{-(1-\eta)\lambda t} (\max\{M_4(f_0), M_4(g_0)\})^{(1-\eta)} |f_0 - g_0|_2^\eta.$$

This concludes the proof of (5.15) by using $(1 - \eta) \leq 1/(1 + \eta)$.

Step 3. From the difference to the linearized semigroup. The same computations imply exactly the same estimate on h_t as on the difference $(f_t - g_t)$, that is inequality (5.16) in Lemma 5.7. \square

We can now consider the *second-order* term in the expansion of the semigroup. Let us recall that the crucial point here is to prove that this second-order term is controlled in terms of some power *strictly greater than 1* of the initial difference.

Lemma 5.9. *There exists $\lambda \in (0, \infty)$ and, for any $\eta \in (0, 1)$, there exists C_η such that for any*

$$\mathbf{r} \in \mathbf{R}_{\mathcal{G}_1} \quad \text{and} \quad f_0, g_0 \in P_{\mathcal{G}_1, \mathbf{r}},$$

we have

$$|\omega_t|_4 \leq C e^{-(1-\eta)\lambda t} |g_0 - f_0|_2^{1+\eta}$$

where

$$\omega_t := g_t - f_t - h_t = S_t^{NL}(g_0) - S_t^{NL}(f_0) - \mathcal{L}_t^{NL}[f_0](g_0 - f_0).$$

Remark 5.10. As proved below ω_t always has vanishing moments up to order 3, which implies that the norm $|\omega_t|_4$ is well-defined. Moreover observe that in this estimate there is *no loss of weight*.

Proof of Lemma 5.9. We consider the angular cutoff decomposition as in Lemma 5.5. Consider the error term

$$\omega := g - f - h, \quad \Omega := \hat{\omega}.$$

which satisfies the evolution equation

$$\partial_t \omega_t = Q(\omega_t, f + g) - Q^+(h, f - g), \quad \omega_0 = 0$$

and (in the Fourier side)

$$\partial_t \Omega = \hat{Q}(\Omega, S) - \hat{Q}^+(H, D).$$

Let us prove that

$$\forall |\alpha| \leq 3, \quad \forall t \geq 0, \quad M_\alpha[\omega_t] := \int_{\mathbb{R}^d} v^\alpha d\omega_t(v) = 0.$$

We shall use again the fact that, for maxwell molecules, the α -th moment of $Q(f_1, f_2)$ is a sum of terms given by product of moments of f_1 and f_2 whose orders sum to $|\alpha|$, see equation (5.18).

We obtain

$$\forall |\alpha| \leq 3, \quad \frac{d}{dt} M_\alpha[\omega_t] = \sum_{\beta \leq \alpha} a_{\alpha, \beta} M_\beta[\omega_t] M_{\alpha-\beta}[f_t + g_t] + \sum_{\beta \leq \alpha} a_{\alpha, \beta} M_\beta[h_t] M_{\alpha-\beta}[f_t - g_t]$$

and since

$$\forall |\alpha| \leq 1, \quad M_\alpha[h_t] = M_\alpha[f_t - g_t] = 0,$$

we deduce

$$\forall |\alpha| \leq 3, \quad \frac{d}{dt} M_\alpha[\omega_t] = \sum_{\beta \leq \alpha} a_{\alpha, \beta} M_\beta[\omega_t] M_{\alpha-\beta}[f_t + g_t].$$

This concludes the proof of the claim about the moments of ω_t since $\omega_0 = 0$.

We now consider the equation in Fourier form

$$\partial_t \Omega = \hat{Q}(\Omega, S) - \hat{Q}^+(H, D)$$

and we deduce in distributional sense

$$\left(\frac{d}{dt} \frac{|\Omega(\xi)|}{|\xi|^4} + K \frac{|\Omega(\xi)|}{|\xi|^4} \right) \leq \mathcal{T}_1 + \mathcal{T}_2 + r_K, \quad r_K \xrightarrow{K \rightarrow \infty} 0$$

(depending on some moments of order 1 of ω , h , d), and

$$\begin{aligned}
\mathcal{T}_1 &:= \sup_{\xi \in \mathbb{R}^3} \int_{\mathbb{S}^{d-1}} \frac{b(\sigma \cdot \hat{\xi})}{|\xi|^4} \left(\left| \frac{\Omega(\xi^+) S(\xi^-)}{2} \right| + \left| \frac{\Omega(\xi^-) S(\xi^+)}{2} \right| \right) d\sigma \\
&\leq \sup_{\xi \in \mathbb{R}^3} \int_{\mathbb{S}^{d-1}} b(\sigma \cdot \hat{\xi}) \left(\frac{|\Omega(\xi^+)|}{|\xi^+|^4} \frac{|\xi^+|^4}{|\xi|^4} + \frac{|\Omega(\xi^-)|}{|\xi^-|^4} \frac{|\xi^-|^4}{|\xi|^4} \right) d\sigma \\
&\leq \left(\sup_{\xi \in \mathbb{R}^3} \frac{|\Omega(\xi)|^2}{|\xi|^2} \right) \left(\sup_{\xi \in \mathbb{R}^3} \int_{\mathbb{S}^{d-1}} b(\sigma \cdot \hat{\xi}) \left(|\hat{\xi}^+|^4 + |\hat{\xi}^-|^4 \right) d\sigma \right) \\
&\leq \lambda_K \left(\sup_{\xi \in \mathbb{R}^3} \frac{|\Omega(\xi)|}{|\xi|^4} \right),
\end{aligned}$$

where λ_K was defined in Lemma 5.5, and

$$\begin{aligned}
\mathcal{T}_2 &:= \frac{1}{2} \sup_{\xi \in \mathbb{R}^3} \int_{\mathbb{S}^{d-1}} \frac{b(\sigma \cdot \hat{\xi})}{|\xi|^4} |H(\xi^+) D(\xi^-) + H(\xi^-) D(\xi^+)| d\sigma \\
&\leq \frac{1}{2} \sup_{\xi \in \mathbb{R}^3} \int_{\mathbb{S}^{d-1}} b(\sigma \cdot \hat{\xi}) \left(\frac{|H(\xi^+)|}{|\xi^+|^2} \frac{|D(\xi^-)|}{|\xi^-|^2} \frac{|\xi^-|^2}{|\xi|^2} + \frac{|D(\xi^+)|^2}{|\xi^+|^2} \frac{|H(\xi^-)|^2}{|\xi^-|^2} \frac{|\xi^-|^2}{|\xi|^2} \right) d\sigma \\
&\leq |h_t|_2 |d_t|_2 \int_{\mathbb{S}^{d-1}} b(\sigma \cdot \hat{\xi}_0) (1 - \sigma \cdot \hat{\xi}_0) d\sigma \\
&= C e^{-(1-\eta)\lambda t} |h_0|_2 |d_0|_2^\eta \leq C e^{-(1-\eta)\lambda t} |d_0|_2^{1+\eta}
\end{aligned}$$

by using the estimates of Lemma 5.5.

Hence we obtain

$$\left(\frac{d}{dt} \frac{|\Omega(\xi)|}{|\xi|^4} + K \frac{|\Omega(\xi)|}{|\xi|^4} \right) \leq \lambda_K \left(\sup_{\xi \in \mathbb{R}^3} \frac{|\Omega(\xi)|}{|\xi|^4} \right) + C e^{-(1-\eta)\lambda t} |d_0|_2^{1+\eta} + r_K.$$

We then deduce from the Gronwall inequality, relaxing the cutoff parameter K as in Lemma 5.7 and assuming without restriction $(1-\eta)\lambda \leq \bar{\lambda}$, that

$$\left(\sup_{\xi \in \mathbb{R}^3} \frac{|\Omega_t(\xi)|}{|\xi|^4} \right) \leq C e^{-(1-\eta)\lambda t} |g_0 - f_0|_2^{1+\eta}.$$

This concludes the proof. \square

5.7. Proof of (A5) uniformly in time in Wasserstein distance. We know from [67] that

$$\sup_{t \geq 0} W_2(S_t^{NL} f_0, S_t^{NL} g_0) \leq W_2(f_0, g_0).$$

As a consequence, by using

$$[\cdot]_1^* = W_1 \leq W_2,$$

we deduce that **(A5)** holds with

$$\Theta(x) = x, \quad \mathcal{F}_3 = \text{Lip}(\mathbb{R}^d) \quad \text{and} \quad P_{\mathcal{G}_3} = P_2(\mathbb{R}^d)$$

endowed with the distance $d_{\mathcal{G}_3} = W_2$.

By using Theorem 3.1 whose assumptions have been proved above, this proves point (i) in Theorem 5.1 and the rate follows from the estimate on $\mathcal{W}_{W_2^N}^N(f)$ from Lemma 4.2.

By using Lemma 4.1 in order to relate $\mathcal{W}_{W_2^N}(\pi_P^N(f_0^N), f_0)$ with $\mathcal{W}_{W_1}(\pi_P^N(f_0^N), f_0)$ and then Lemma 4.7 in order to estimate

$$\mathcal{W}_{W_1}(\pi_P^N(f_0^N), f_0) \xrightarrow{N \rightarrow \infty} 0$$

for the sequence of initial data conditioned on the energy sphere constructed in Lemma 4.4, we then deduce point (iii) in Theorem 5.1.

5.8. Proof of (A5) with time growing bounds in negative Sobolev norms. It is also possible (and in fact easier) to prove, in the cutoff case, that the weak stability holds in Sobolev space *on finite time*:

Lemma 5.11. *For any $T \geq 0$ and $s > d/2$ there exists $C_{T,s}$ such that for any f_t, g_t solutions of the Boltzmann equation for Maxwell molecules (5.6) and initial data f_0 and g_0 , there holds*

$$\sup_{[0,T]} \|f_t - g_t\|_{H^{-s}} \leq C_{T,s} \|f_0 - g_0\|_{H^{-s}}.$$

Sketch the proof of Lemma 5.11. We integrate (5.12) against $D/(1 + |\xi|^2)^s$:

$$\frac{d}{dt} \|D\|_{H^{-k}}^2 = \frac{1}{2} \int_{\xi} \int_{\mathbb{S}^{d-1}} b(\sigma \cdot \hat{\xi}) \frac{[D^- S^+ D + D^+ S^- D - 2|D|^2]}{(1 + |\xi|^2)^s} d\sigma d\xi$$

and we use Young's inequality together with the bounds

$$\|S^+\|_{\infty}, \quad \|S^-\|_{\infty} \leq \|f + g\|_{M^1} \leq 2$$

to conclude. □

This proves (A5) with the alternate choice

$$\Theta(x) = x, \quad \mathcal{F}_3 = H^s(\mathbb{R}^d) \quad \text{and} \quad P_{\mathcal{G}_3} = P_2(\mathbb{R}^d)$$

endowed with the distance of the normed space $\mathcal{G}_3 = H^s(\mathbb{R}^d)$. Then point (ii) in Theorem 5.1 follows from the abstract theorem 3.1 where the rate is provided by the estimate on $\mathcal{W}_{(H^{-s})^2}^N(f)$ from Lemma 4.2.

5.9. Proof of infinite-dimensional Wasserstein chaos. We shall prove Theorem 5.3 in this subsection. Let us proceed in several steps. Let us emphasize that we do not search for optimality on the rate functions given by our argument.

Step 1: Finite-dimensional Wasserstein chaos. It is immediate that Theorem 5.1 implies that, under one of the two possible assumptions on the initial data, for any given $\ell \geq 1$, one has

$$\sup_{t \geq 0} \left\| \Pi_{\ell} [f_t^N] - f_t^{\otimes \ell} \right\|_{H^{-s}} \leq \alpha_0(\ell, N)$$

for some power law rate function $\alpha_0(\ell, N) \rightarrow 0$ as $N \rightarrow \infty$.

Then by using Lemma 4.1 we deduce that

$$\sup_{t \geq 0} W_1 \left(\Pi_{\ell} [f_t^N], f_t^{\otimes \ell} \right) \leq \alpha(\ell, N)$$

for some power law rate function $\alpha(\ell, N) \rightarrow 0$ as $N \rightarrow \infty$.

Note carefully that at this point our rate function still depends on ℓ and in fact a quick look at Theorem 5.1 shows that they scale like ℓ^2 , therefore making impossible to choose $\ell \sim N$.

Step 2: Infinite-dimensional Wasserstein chaos. We shall use here the following result obtained in [37], see also [56, Théorème 2.1]: for any $f \in P(\mathbb{R}^d)$ and sequence $f^N \in P_{\text{sym}}(\mathbb{R}^d)$ we have

$$\forall 1 \leq \ell \leq N, \quad \frac{W_1(\Pi_{\ell} [f^N], f^{\otimes \ell})}{\ell} \leq C \left(W_1(\Pi_2 [f^N], f^{\otimes 2})^{\alpha_1} + \frac{1}{N^{\alpha_2}} \right)$$

for some constructive constant $C, \alpha_1, \alpha_2 > 0$.

By combining this estimate with the previous step we immediately obtain

$$\sup_{1 \leq \ell \leq N} \sup_{t \geq 0} W_1 \left(\Pi_\ell [f_t^N], f_t^{\otimes \ell} \right) \leq \alpha(N)$$

for some power law rate function $\alpha(N) \rightarrow 0$ as $N \rightarrow \infty$. This concludes the proof of (5.8).

Step 3: Relaxation in Wasserstein distance. We shall prove (5.9) and we shall consider here initial data f_0^N constructed by conditioning $f_0^{\otimes N}$ to the energy sphere. We first write

$$\frac{W_1(f_t^N, \gamma^N)}{N} \leq \frac{W_1(f_t^N, f_t^{\otimes N})}{N} + \frac{W_1(f_t^{\otimes N}, \gamma^{\otimes N})}{N} + \frac{W_1(\gamma^{\otimes N}, \gamma^N)}{N}.$$

Since $f_t^N \rightarrow \gamma^N$ in L^2 and $f_t \rightarrow \gamma$ in L^1 as $t \rightarrow +\infty$, one can pass to the limit in the Wasserstein distance and get from the previous step

$$\frac{W_1(\gamma^{\otimes \ell}, \Pi_\ell[\gamma^N])}{N} \leq \alpha(N).$$

Moreover it is immediate that

$$\frac{W_1(f_t^{\otimes N}, \gamma^{\otimes N})}{N} = W_1(f_t, \gamma).$$

Finally it was proved in [60] that under our assumptions on f_0 one has

$$\|(f_t - \gamma) \langle v \rangle\|_{L^1} \leq C e^{-\lambda_1 t}$$

for some constants $C > 0$ and $\lambda_1 > 0$ which implies

$$W_1(f_t, \gamma) \leq \|(f_t - \gamma) \langle v \rangle\|_{L^1} \leq C e^{-\lambda_1 t}.$$

Hence, gathering these three estimates, we deduce that

$$(5.21) \quad \frac{W_1(f_t^N, \gamma^N)}{N} \leq 2\alpha(N) + C e^{-\lambda_1 t}.$$

It was proved in [15] that there exists $\lambda_2 > 0$ such that

$$\forall N \geq 1, \forall t \geq 0 \quad \|h^N - 1\|_{L^2(\mathcal{S}^N, \gamma^N)} \leq e^{-\lambda_2 t} \|h_0^N - 1\|_{L^2(\mathcal{S}^N, \gamma^N)},$$

where $h^N = df^N/d\gamma^N$ is the Radon-Nikodym derivative of f^N with respect to the measure γ^N so that $f^N = h^N \gamma^N$. When $f_0^N = [f_0^{\otimes N}]_{\mathcal{S}^{dN-1}(\sqrt{N})}$ with $f_0 \in P_4(\mathbb{R}^d)$ we easily upper bound the right hand side term by

$$\|h_0^N - 1\|_{L^2(\mathcal{S}^N, \gamma^N)} \leq A^N,$$

where $A = A(f_0) > 1$. Thanks to Cauchy-Schwartz inequality and the control of the Wasserstein distance in terms of a weighted total variation norm (see [72, Proposition 7.10]) we also have

$$\begin{aligned} \|h^N - 1\|_{L^2(\mathcal{S}^{dN-1}(\sqrt{N}), \gamma^N)} &\geq \|h^N - 1\|_{L^1(\mathcal{S}^{dN-1}(\sqrt{N}), \gamma^N)} \\ &\geq \int_{\mathbb{R}^{dN}} \frac{1}{N} \sum_{i=1}^N |v_i| |df^N - d\gamma^N|(V) \\ &\geq \frac{1}{N} W_1(f^N, \gamma^N) \end{aligned}$$

and we deduce

$$(5.22) \quad \forall N \geq 1, \forall t \geq 0, \quad W_1(f_t^N, \gamma^N) \leq A^N e^{-\lambda_2 t}.$$

Finally by combining (5.21) and (5.22) and taking the better of these two estimates depending on N and t , we easily obtain

$$\forall N \geq 1, \forall t \geq 0, \quad W_1(f_t^N, \gamma^N) \leq \beta(t)$$

for some $\beta(t) \rightarrow 0$ as $t \rightarrow +\infty$, which concludes the proof of (5.9).

6. HARD SPHERES

6.1. The model. The limiting equation was introduced in Subsection 1.1 and the stochastic model has been already discussed Subsection 5.1.

We consider here the case of the Master equation (5.3), (5.4) and the limit nonlinear homogeneous Boltzmann equation (1.1), (1.2), (1.3) with

$$(6.1) \quad B(z, \cos \theta) = \Gamma(z) b(\cos \theta) = \Gamma(z) = |z|.$$

6.2. Statement of the result. Our fluctuations estimate result for this model then states as follows:

Theorem 6.1 (Hard spheres detailed chaos estimates). *Let us consider a 1-particle initial distribution $f_0 \in P(\mathbb{R}^d)$ with energy less than \mathcal{E}*

$$M_2(f_0) \leq \mathcal{E},$$

and a hierarchy of N -particle distributions

$$f_t^N = S_t^N(f_0^{\otimes N})$$

issued from the initial data f_0^N , as well the 1-particle of the limiting semigroup

$$f_t = S_t^{NL}(f_0)$$

where we assume that the collision satisfies (6.1).

Let us finally fix some $\delta \in (0, 1)$.

Then

(i) Suppose that f_0 has compact support

$$\text{Supp } f_0 \subset \left\{ v \in \mathbb{R}^d, |v| \leq A \right\}$$

and that the N -particle initial data is tensorized

$$\forall N \geq 1, \quad f_0^N = f_0^{\otimes N}.$$

Then there are

- some constants $k_1 \geq 2$ depending on δ and \mathcal{E} ;
 - some constant $C_\delta > 0$ depending on δ and \mathcal{E} , and blowing up as $\delta \rightarrow 1$;
 - some constant $C_b > 0$ depending on the collision kernel,
- such that for any $\ell \in \mathbb{N}^*$, and for any

$$\varphi = \varphi_1 \otimes \varphi_2 \otimes \cdots \otimes \varphi_\ell \in W^{1,\infty}(\mathbb{R}^d)^{\otimes \ell},$$

we have

$$(6.2) \quad \forall N \geq 2\ell, \quad \sup_{[0,T]} \left| \left\langle \left(S_t^N(f_0^N) - (S_t^{NL}(f_0))^{\otimes N} \right), \varphi \right\rangle \right| \\ \leq \|\varphi\|_{W^{1,\infty}(\mathbb{R}^d)^{\otimes \ell}} \left[\frac{2\ell^2}{N} + \frac{C_\delta \ell^2 \|f_0\|_{M_{k_1}^1}}{N^{1-\delta}} + \ell e^{C_b A} \theta(N) \right].$$

The last term of the right hand side (which is also the dominant error term as N goes to infinity in our estimate) is given by

$$\theta(N) = \frac{C}{(1 + |\log N|)^\alpha}$$

for some constants $C, \alpha > 0$.

(ii) Under the same setting but assuming instead for the limiting initial data

$$f_0 \in L^\infty(\mathbb{R}^d) \quad \text{s. t.} \quad \int_{\mathbb{R}^d} e^{z|v|} df_0(v) < +\infty$$

for some $z > 0$, and taking for the N -particle initial data the sequence $(f_0^N)_{N \geq 1}$ constructed in Lemma 4.4 and 4.7, we have the same estimate (6.2) however with no information on the rate

$$\theta(N) \xrightarrow{N \rightarrow 0} 0.$$

This nevertheless proves the propagation of chaos, uniformly in time.

We now state the key Wasserstein version of the propagation of chaos estimate, which is valid for any number of marginals, although with a possibly worse (but still constructive) rate. Combined with previous results on the relaxation of the N -particle system we also deduce some estimate of relaxation to equilibrium independent of N and, again, for any number of marginals.

Theorem 6.2 (Hard spheres Wasserstein chaos). *Under the same setting as in Theorem 6.1, either*

- (a) with f_0 compactly supported and $f_0^N = f_0^{\otimes N}$ or
- (b) with

$$f_0 \in L^\infty(\mathbb{R}^d) \quad \text{s. t.} \quad \int_{\mathbb{R}^d} e^{z|v|} df_0(v) < +\infty$$

and f_0^N constructed by Lemma 4.4,

we have

$$(6.3) \quad \forall N \geq 1, \forall 1 \leq \ell \leq N, \quad \sup_{t \geq 0} \frac{W_1(\Pi_\ell f_t^N, (f_t^{\otimes \ell}))}{\ell} \leq \alpha(N)$$

for some $\alpha(N) \rightarrow 0$ as $N \rightarrow \infty$ (in the case (a) one has moreover explicit estimates on the rate α like a power of a logarithm).

Then in case (b) we have

$$(6.4) \quad \forall N \geq 1, \forall 1 \leq \ell \leq N, \forall t \geq 0, \quad \frac{W_1(\Pi_\ell f_t^N, \Pi_\ell(\gamma^N))}{\ell} \leq \beta(t)$$

for some $\beta(t) \rightarrow 0$ as $t \rightarrow \infty$, where γ is the gaussian equilibrium with energy \mathcal{E} and γ^N is the uniform probability measure on $\mathcal{S}^N(\mathcal{E})$.

In order to prove Theorem 6.1, we shall prove assumptions **(A1)**-**(A2)**-**(A3)**-**(A4)**-**(A5)** of Theorem 3.1 with $T = \infty$. The application of the latter theorem then exactly yields Theorem 6.1 by following carefully each constant computed below. We fix

$$\mathcal{F}_1 = \mathcal{F}_2 = C_b(\mathbb{R}^d) \quad \text{and} \quad \mathcal{F}_3 = \text{Lip}(\mathbb{R}^d).$$

Then the proof of Theorem 6.2 is deduced from Theorem 6.1 in a similar way as Theorem 5.3 was deduced from Theorem 5.1, see Subsection 6.9.

6.3. Proof of (A1). From the discussion made in section 5.3 we easily see that for the Hard sphere model the operator G^N is bounded from $C_{-k+1,0}(\mathbb{R}^{dN})$ onto $C_{-k,0}(\mathbb{R}^{dN})$ for any $k \in \mathbb{R}$. Since G^N is close, dissipative and $C_{-k+1,0}(\mathbb{R}^{dN})$ is dense in $C_{-k,0}(\mathbb{R}^{dN})$, the Hille-Yosida theory implies that G^N generates a Markov type semigroup T_t^N on $C_{-k,0}(\mathbb{R}^{dN})$ and we may also define S_t^N by duality as a semigroup on $P_k(\mathbb{R}^{dN})$. The nonlinear semigroup S_t^{NL} is also well defined on $P_k(\mathbb{R}^d)$, $k \geq 2$, see for instance [32, 29, 48].

Lemma 5.4 was proved both for Maxwell molecules and hard spheres. It first shows that

$$\forall t \geq 0, \quad \text{Supp } f_t^N \subset \mathbb{E}_N := \left\{ V \in \mathbb{R}^{dN}; M_2^N(V) \leq \mathcal{E} \right\}.$$

It also proves that for any $k \geq 2$,

$$\sup_{t \geq 0} \langle f_t^N, M_k^N \rangle \leq C_k^N$$

where C_k^N depends on k , \mathcal{E} , on the collision kernel and on the initial value

$$\langle f_0^N, M_k^N \rangle$$

which is uniformly bounded in N in terms of k and A . This shows that **(A1)-(i)** holds with $m_1(v) := |v|^{k_1}$ for any $k_1 \geq 2$. The precise value of k_1 shall be chosen in Section 6.7.

As for **(A1)-(ii)**, we remark that for a given N -particle velocity

$$V = (v_1, \dots, v_N) \in \mathbb{R}^{dN},$$

we have

$$V \in \text{Supp } f_0^{\otimes N} \iff \forall i = 1, \dots, N, v_i \in \text{Supp } f_0$$

which implies

$$\forall i = 1, \dots, N, \quad m_{\mathcal{G}_3}(v_i) \leq m_{\mathcal{G}_3}(A) \quad \text{with} \quad m_{\mathcal{G}_3}(v) := e^{a|v|}$$

for any constant $a > 0$, which shall be chosen later on.

We conclude that

$$\text{Supp } f_0^{\otimes N} \subset \left\{ V \in \mathbb{R}^{dN}; M_{m_{\mathcal{G}_3}}^N(V) \leq m_{\mathcal{G}_3}(A) \right\},$$

and **(A1)-(ii)** holds for the exponential growing weight $m_{\mathcal{G}_3}$.

6.4. Proof of (A2). We define

$$P_{\mathcal{G}_1} := \left\{ f \in P(\mathbb{R}^d); \langle f, \langle v \rangle^{k_1} \rangle < +\infty \right\}$$

that we endow with the total variation norm $\|\cdot\|_{\mathcal{G}_1} := \|\cdot\|_{M^1}$. Note that since $k_1 \geq 2$, elements of $P_{\mathcal{G}_1}$ have in particular finite energy.

This assertion **(A2)-(i)** reads:

For some constant $a_{k_1} > 0$ and for any $a \geq a_{k_1}$ and for any $t > 0$, the application

$$f_0 \mapsto S_t^{NL} f_0$$

maps

$$\mathcal{BP}_{\mathcal{G}_1, a} := \left\{ f \in P_{\mathcal{G}_1}, M_{k_1}(f) = \langle f, \langle v \rangle^{k_1} \rangle \leq a \right\}$$

continuously into itself.

It is postponed to section 6.6, where we prove in (6.6) a Hölder continuity of the flow in $\mathcal{BP}_{\mathcal{G}_1, a}$.

Let us consider assertion **(A2)-(ii)**, that is the fact that the operator Q is bounded and Hölder continuous from $P_{\mathcal{G}_1}$ to \mathcal{G}_1 .

For any $f, g \in P_{\mathcal{G}_1}$ we have

$$\begin{aligned} \|Q(g, g) - Q(f, f)\|_{M^1} &= \|Q(g - f, g + f)\|_{M^1} \\ &\leq 2 \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \int_{\mathbb{S}^{d-1}} b(\theta) |v - v_*| |f - g| |f_* + g_*| d\sigma dv_* dv \\ (1 + \mathcal{E}) \|b\|_{L^1} \|(f - g) \langle v \rangle\|_{M^1}. \end{aligned}$$

We deduce that

$$\|Q(g, g) - Q(f, f)\|_{M^1} \leq 2(1 + \mathcal{E})^{3/2} \|b\|_{L^1} \|f - g\|_{M^1}^{1/2}$$

which yields

$$Q \in C^{0,1/2}(P_{\mathcal{G}_1}; \mathcal{G}_1)$$

and also implies that Q is bounded on $P_{\mathcal{G}_1}$ since we can choose g to be a maxwellian distribution.

6.5. Proof of (A3). For any $k_1 \geq 2$, any energy

$$r_{d+1} \geq 0$$

and any mean velocity

$$(r_1, \dots, r_d) \in B_{\mathbb{R}^d}(0, \sqrt{r_{d+1}})$$

we define

$$P_{\mathcal{G}_1, \mathbf{r}} := \{f \in P_{\mathcal{G}_1} ; \langle f, v_j \rangle = r_j, j = 1, \dots, d, \langle f, |v|^2 \rangle = r_{d+1}\}$$

with the notation

$$\mathbf{r} := (r_1, \dots, r_{d+1}).$$

This corresponds to the vector of constraints

$$\mathbf{m} = (v_1, \dots, v_d, |v|^2)$$

in Definition 2.4.

We also denote by \mathbf{R} the set of all admissible constraint vectors $\mathbf{r} \in \mathbb{R}^{d+1}$:

$$\mathbf{R} = \left\{ (r_1, \dots, r_{d+1}) \in \mathbb{R}^{d+1} \text{ s. t. } r_{d+1} \geq 0 \text{ and } \sum_{i=1}^d r_i^2 \leq r_{d+1} \right\}.$$

Let us finally define the weight

$$\Lambda_1(f) := M_{k_1-2}(f) = \langle f, \langle v \rangle^{k_1-2} \rangle$$

(this means that we choose $m'_{\mathcal{G}_1}(v) = \langle v \rangle^{k_1-2}$ in assumption (A3)).

We claim that there exists a constant $C_{k_1} > 0$ (depending on k_1) such that for any $\eta \in (0, 1)$ and any function

$$\Phi \in \bigcap_{\mathbf{r} \in \mathbf{R}_\varepsilon} C_{\Lambda_1}^{1, \eta}(P_{\mathcal{G}_1, \mathbf{r}}; \mathbb{R}),$$

we have

$$(6.5) \quad \forall V \in \mathbb{R}^{Nd}, \quad |G^N(\Phi \circ \mu_V^N) - (G^\infty \Phi)(\mu_V^N)| \leq C_{k_1} \left(\sup_{\mathbf{r} \in \mathbf{R}_\varepsilon} [\Phi]_{C_{\Lambda_1}^{1, \eta}(P_{\mathcal{G}_1, \mathbf{r}})} \right) \frac{M_{k_1}^N(V)}{N^\eta}$$

where we recall that

$$M_{k_1}^N(V) := \frac{1}{N} \sum_{i=1}^N \langle v_i \rangle^{k_1}.$$

This would prove assumption (A3) with the rate

$$\varepsilon(N) = \frac{C_{k_1}}{N^\eta}, \quad \eta := 1 - \delta.$$

Consider $V \in \mathbb{R}^{dN}$ and define

$$\mathbf{r}_V := \left(\langle \mu_V^N, z_1 \rangle, \dots, \langle \mu_V^N, z_d \rangle, \langle \mu_V^N, |z|^2 \rangle \right) \in \mathbf{R}_{\mathcal{E}}.$$

Then let us consider a given

$$\Phi \in C_{\Lambda_1}^{1,\eta}(P_{\mathcal{G}_1, \mathbf{r}_V}; \mathbb{R}),$$

and let us set

$$\phi := D\Phi[\mu_V^N].$$

We compute

$$\begin{aligned} G^N(\Phi \circ \mu_V^N) &= \frac{1}{2N} \sum_{i,j=1}^N |v_i - v_j| \int_{\mathbb{S}^{d-1}} \left[\Phi(\mu_{V_{ij}^*}^N) - \Phi(\mu_V^N) \right] b(\theta_{ij}) d\sigma \\ &= \frac{1}{2N} \sum_{i,j=1}^N |v_i - v_j| \int_{\mathbb{S}^{d-1}} \langle \mu_{V_{ij}^*}^N - \mu_V^N, \phi \rangle b(\theta_{ij}) d\sigma \\ &\quad + \frac{[\Phi]_{C_{\Lambda_1}^{1,\eta}(P_{\mathcal{G}_1, \mathbf{r}_V})}}{2N} \sum_{i,j=1}^N |v_i - v_j| \times \\ &\quad \int_{\mathbb{S}^{d-1}} \max \left\{ M_{m'_{\mathcal{G}_1}}(\mu_{V_{ij}^*}^N); M_{m'_{\mathcal{G}_1}}(\mu_V^N) \right\} \mathcal{O} \left(\left\| \mu_{V_{ij}^*}^N - \mu_V^N \right\|_{M^1}^{1+\eta} \right) d\sigma \\ &=: I_1(V) + I_2(V). \end{aligned}$$

As for the first term $I_1(V)$, we have

$$\begin{aligned} I_1(V) &= \frac{1}{2N^2} \sum_{i,j=1}^N |v_i - v_j| \int_{\mathbb{S}^{d-1}} b(\theta_{ij}) [\phi(v_i^*) + \phi(v_j^*) - \phi(v_i) - \phi(v_j)] d\sigma \\ &= \frac{1}{2N^2} \int_v \int_w |v - w| \int_{\mathbb{S}^{d-1}} b(\theta) [\phi(v^*) + \phi(w^*) - \phi(v) - \phi(w)] \mu_V^N(dv) \mu_V^N(dw) d\sigma \\ &= \langle Q(\mu_V^N, \mu_V^N), \phi \rangle \end{aligned}$$

which is precisely the term we are searching for thanks to assumption **(A2)** and Lemma 2.13:

$$I_1(V) = \langle Q(\mu_V^N, \mu_V^N), \phi \rangle = (G^\infty \Phi)(\mu_V^N).$$

As for the second term $I_2(V)$, using that

$$\begin{aligned} M_{m'_{\mathcal{G}_1}}(\mu_{V_{ij}^*}^N) &:= M_{k_1-2}^N(V_{ij}^*) = \frac{1}{N} \left(\left(\sum_{\ell \neq i,j} \langle v_\ell \rangle^{k_1-2} \right) + \langle v_i^* \rangle^{k_1-2} + \langle v_j^* \rangle^{k_1-2} \right) \\ &\leq \frac{1}{N} \left(\left(\sum_{\ell \neq i,j} \langle v_\ell \rangle^{k_1-2} \right) + 2 (\langle v_i \rangle^2 + \langle v_j \rangle^2)^{\frac{k_1}{2}-1} \right) \\ &\leq \frac{2^{k_1/2}}{N} \left(\sum_{\ell=1}^N \langle v_\ell \rangle^{k_1-2} \right) = C M_{k_1-2}^N(V) = C M_{m'_{\mathcal{G}_1}}(\mu_V^N), \end{aligned}$$

we deduce

$$\begin{aligned}
|I_2(V)| &\leq C M_{k_1-2}^N(V) [\Phi]_{C_\Lambda^{1,\eta}(P_{\mathcal{G}_1, r_V})} \left(\frac{1}{2N} \sum_{i,j=1}^N |v_i - v_j| \left(\frac{4}{N} \right)^{1+\eta} \right) \\
&\leq C M_{k_1-2}^N(V) [\Phi]_{C_\Lambda^{1,\eta}(P_{\mathcal{G}_1, r_V})} \left(\frac{1}{N^\eta} \frac{1}{N^2} \sum_{i,j=1}^N (\langle v_i \rangle + \langle v_j \rangle) \right) \\
&\leq \frac{C}{N^\eta} M_{k_1-2}^N(V) M_2^N(V) [\Phi]_{C_\Lambda^{1,\eta}(P_{\mathcal{G}_1, r_V})}.
\end{aligned}$$

We then use the elementary inequality

$$M_{k_1-2}^N(V) M_2^N(V) \leq C_{k_1} M_{k_1}^N(V)$$

for a constant C_{k_1} depending on k_1 but *not* on N , which yields

$$|I_2(V)| \leq \frac{C}{N^\eta} M_{k_1}^N(V) [\Phi]_{C_\Lambda^{1,\eta}(P_{\mathcal{G}_1, r_V})}.$$

We conclude that (6.5) holds by piling the two last estimates for I_1 and I_2 .

6.6. Proof of (A4) with time growing bounds. Let us consider some 1-particle initial data

$$f_0, g_0 \in P_{\mathcal{G}_1}.$$

Similar to as in the previous section, we then define (under the assumption (6.1) on the collision kernel) the associated solutions f_t and g_t to the nonlinear Boltzmann equation (1.1), as well as

$$h_t := \mathcal{L}_t^{NL} [f_0] (g_0 - f_0)$$

the solution to the linearized Boltzmann equation around f_t , given by

$$\begin{cases} \partial_t f_t = Q(f_t, f_t), & f_{|t=0} = f_0 \\ \partial_t g_t = Q(g_t, g_t), & g_{|t=0} = g_0 \\ \partial_t h_t = 2Q(f_t, h_t), & h_{|t=0} = h_0 := g_0 - f_0. \end{cases}$$

We also define as before

$$\omega_t := g_t - f_t - h_t.$$

We shall now again *expand the limiting nonlinear semigroup* in terms of the initial data, around f_0 . The goal is to prove assumption **(A4)** with the choice of indices

$$(\eta', \eta'') = (1, (1 + \eta)/2).$$

This imposes the choice of weight

$$\Lambda_2(f) = \Lambda_1(f)^{\frac{1}{2}} = \sqrt{M_{k_1-2}(f)}.$$

Lemma 6.3. *For any given energy $\mathcal{E} > 0$ and any $\eta > 0$ there exists*

- *some constant $k_1 \geq 2$ (depending on \mathcal{E} and δ),*
- *some constant C (depending on \mathcal{E}),*

such that for any

$$\mathbf{r} \in \mathbf{R}_{\mathcal{E}} \quad \text{and} \quad f_0, g_0 \in P_{\mathcal{G}_1, \mathbf{r}}$$

we have for any $t \geq 0$:

$$(6.6) \quad \|g_t - f_t\|_{M_2^1} \leq e^{C(1+t)} \sqrt{\max\{M_{k_1-2}(f_0), M_{k_1-2}(g_0)\}} \|f_0 - g_0\|_{M_2^1},$$

$$(6.7) \quad \|h_t\|_{M_2^1} \leq e^{C(1+t)} \sqrt{M_{k_1-2}(f_0)} \|f_0 - g_0\|_{M_2^1},$$

$$(6.8) \quad \|\omega_t\|_{M_2^1} \leq e^{C(1+t)} \sqrt{\max\{M_{k_1-2}(f_0), M_{k_1-2}(g_0)\}} \|f_0 - g_0\|_{M_2^1}^{2-\eta}.$$

Proof of Lemma 6.3. We proceed in several steps and number the constant for clarity.

Let us define

$$\forall f \in M^1(\mathbb{R}^d), \quad \|f\|_{M_k^1} := \int_{\mathbb{R}^d} \langle v \rangle^k |f|(dv), \quad \|f\|_{M_{k,\ell}^1} := \int_{\mathbb{R}^d} \langle v \rangle^k (1 + \log \langle v \rangle)^\ell |f|(dv).$$

Step 1. The strategy. Existence and uniqueness for f_t , g_t and h_t is a consequence of the following important stability argument that we use several times. This estimate is due to DiBlasio [22] in a L^1 framework, and it has been recently extended to a measure framework in [29, Lemma 3.2] (see also [32] and [48] for other argument of uniqueness for measure solutions of the spatially homogeneous Boltzmann equation).

Let us sketch the argument for h . We first write

$$(6.9) \quad \begin{aligned} \frac{d}{dt} \int \langle v \rangle^2 |h_t|(dv) &\leq \iiint |h_t|(dv) f_t(dv_*) |u| b(\theta) \left[\langle v' \rangle^2 + \langle v'_* \rangle^2 - \langle v \rangle^2 - \langle v_* \rangle^2 \right] d\sigma \\ &\quad + 2 \iiint |h_t|(dv) f_t(dv_*) |u| b(\theta) \langle v_* \rangle^2 d\sigma \end{aligned}$$

(this formal computation can be justified by a regularization procedure, we refer to [29] for instance). Since the first term vanishes, we deduce that

$$(6.10) \quad \frac{d}{dt} \|h_t\|_{M_2^1} \leq C_1 \|f\|_{M_3^1} \|h_t\|_{M_2^1}$$

for some constant $C_1 > 0$ only depending on b .

Then in the case when

$$(6.11) \quad \|f_s\|_{M_3^1} \in L^1(0, t) \text{ on some time interval } s \in [0, t]$$

we may integrate this differential inequality and we deduce that h is unique.

More precisely, we have established

$$(6.12) \quad \sup_{s \in [0, t]} \|h_s\|_{M_2^1} \leq \|g_0 - f_0\|_{M_2^1} \exp \left(C_1 \int_0^t \|f_s\|_{M_3^1} ds \right),$$

and similar arguments imply

$$(6.13) \quad \sup_{s \in [0, t]} \|f_s - g_s\|_{M_2^1} \leq \|g_0 - f_0\|_{M_2^1} \exp \left(C_1 \int_0^t \|f_s + g_s\|_{M_3^1} ds \right).$$

It is worth mentioning that one cannot prove (6.11) under the sole assumption

$$\|f_0\|_{M_2^1} < \infty$$

on the initial data since it would contradict the non-uniqueness result of [49]. However, as we prove in (6.15) below, one may show (thanks to the Povzner inequality, as developed in [59, 47]) that (6.11) holds as soon as

$$\|f_0\|_{M_{2,1}^1} < \infty.$$

This will be a key step for establishing (6.6) and (6.7).

Now, our goal is to estimate the M_2^1 norm of

$$\omega_t := g_t - f_t - h_t$$

in terms of $\|g_0 - f_0\|_{M_2^1}$. The measure ω_t satisfies the evolution equation:

$$\partial_t \omega_t = Q(g_t, g_t) - Q(f_t, f_t) - Q(h_t, f_t) - Q(f_t, h_t), \quad \omega_0 = 0$$

which can be rewritten as

$$\partial_t \omega_t = Q(\omega_t, f_t + g_t) + Q(h_t, g_t - f_t).$$

The same arguments as in (6.9)-(6.10) yield the following differential inequality

$$\frac{d}{dt} \|\omega_t\|_{M_2^1} \leq C_2 \|\omega_t\|_{M_2^1} \|f_t + g_t\|_{M_3^1} + \|Q(h_t, g_t - f_t)\|_{M_2^1}, \quad \|\omega_0\|_{M_2^1} = 0$$

for some constant $C_2 > 0$ depending on b .

We deduce

$$\sup_{s \in [0, t]} \|\omega_s\|_{M_2^1} \leq \left(\int_0^t \|Q(h_s, f_s - g_s)\|_{M_2^1} ds \right) \exp \left(C_2 \int_0^t \|f_s + g_s\|_{M_3^1} ds \right).$$

Since

$$\begin{aligned} \int_0^t \|Q(h_s, f_s - g_s)\|_{M_2^1} ds &\leq C_2 \left(\sup_{s \in [0, t]} \|h_s\|_{M_2^1} \right) \left(\int_0^t \|g_s - f_s\|_{M_3^1} ds \right) \\ &\quad + C_2 \left(\sup_{s \in [0, t]} \|g_s - f_s\|_{M_2^1} \right) \left(\int_0^t \|h_s\|_{M_3^1} ds \right), \end{aligned}$$

we deduce from (6.12) and (6.13)

$$(6.14) \quad \begin{aligned} \sup_{s \in [0, t]} \|\omega_s\|_{M_2^1} &\leq C_2 \|g_0 - f_0\|_{M_2^1} \exp \left(C_2 \int_0^t (\|f_s\|_{M_3^1} + \|g_s\|_{M_3^1}) ds \right) \times \\ &\quad \times \left[\left(\int_0^t \|g_s - f_s\|_{M_3^1} \right) \exp \left(C_1 \int_0^t \|f_s\|_{M_3^1} \right) \right. \\ &\quad \left. + \left(\int_0^t \|h_s\|_{M_3^1} \right) \exp \left(C_1 \int_0^t (\|f_s\|_{M_3^1} + \|g_s\|_{M_3^1}) ds \right) \right]. \end{aligned}$$

Hence the problem now reduces to the obtaining of time integral controls over

$$\|f_s\|_{M_3^1}, \quad \|g_s\|_{M_3^1}, \quad \|f_s - g_s\|_{M_3^1} \quad \text{and} \quad \|h_s\|_{M_3^1}.$$

Step 2. Time integral control of f and g in M_3^1 . In this step we prove

$$(6.15) \quad \int_0^t \|f_s\|_{M_{3, \ell-1}^1} dt \leq C_3(\mathcal{E}) t + C_4 \|f_0\|_{M_{2, \ell}^1} \quad \ell = 1, 2,$$

for the solution f_t , where $C_3(\mathcal{E}) > 0$ is a constant depending on the energy, and $C_4 > 0$ is a numerical constant. The same estimate obviously holds for the solution g_t .

The estimates (6.15) are a consequence of the accurate version of the Povzner inequality which has been proved in [59, 47]. Indeed it was shown in [59, Lemma 2.2] that for any function

$$\Psi : \mathbb{R}^d \rightarrow \mathbb{R}, \quad \Psi(v) = \psi(|v|^2) \quad \text{with } \psi \text{ convex,}$$

the solution f_t to the hard spheres Boltzmann equation satisfies

$$\frac{d}{dt} \int_{\mathbb{R}^d} \Psi(v) f_t(dv) = \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} f_t(dv) f_t(dv_*) |v - v_*| K_\Psi(v, v_*)$$

with $K_\Psi = G_\Psi - H_\Psi$, where the term G_Ψ “behaves mildly” (see below) and the term H_Ψ is given by (see [59, formula (2.7)])

$$H_\Psi(v, v_*) = 2\pi \int_0^{\pi/2} \left[\psi (|v|^2 \cos^2 \theta + |v_*|^2 \sin^2 \theta) - \cos^2 \theta \psi (|v|^2) - \sin^2 \theta \psi (|v_*|^2) \right] d\theta.$$

Note that $H_\Psi \geq 0$ since its integrand is nonnegative because of the convexity of ψ .

More precisely, in the cases that we are interested with, namely

$$\Psi(v) = \psi_{2,\zeta}(|v|^2) \quad \text{with} \quad \psi_{k,\zeta}(r) = r^{k/2} (\log r)^\zeta \quad \text{and} \quad \zeta = 1, 2,$$

it is established in [59] that (with obvious notation)

$$\forall v, v_* \in \mathbb{R}^d, \quad |G_{\psi_{2,\zeta}}(v, v_*)| \leq C_\zeta \langle v \rangle (\log \langle v \rangle)^\zeta \langle v_* \rangle (\log \langle v_* \rangle)^\zeta$$

for some constant $C_5(\zeta) > 0$ depending on ζ .

On the other hand, in the case $\zeta = 1$ we compute, with the help of the the notation $x := \cos^2 \theta$ and $u = |v_*|/|v|$,

$$\forall x \in [1/4, 3/4], \quad \forall u \in [0, 1/2],$$

$$\begin{aligned} & \psi_{2,1} \left(|v|^2 \cos^2 \theta + |v_*|^2 \sin^2 \theta \right) - \cos^2 \theta \psi_{2,1} (|v|^2) - \sin^2 \theta \psi_{2,1} (|v_*|^2) = \\ & = |v|^2 \left[(1-x) \psi_{2,1} (u^2) + x \psi_{2,1} (1) - \psi_{2,1} ((1-x)u^2 + x) \right] \geq C_6 |v|^2, \end{aligned}$$

for some numerical constant $C_6 > 0$, which only depends on the strict convexity of the real function $\psi_{2,1}$. We deduce that there exists a constant $C_7 > 0$ such that

$$H_{\psi_{2,1}}(v, v_*) \geq C_7 |v|^2 \mathbf{1}_{|v| \geq 2|v_*|}.$$

Similarly, in the case $\zeta = 2$, we have

$$\forall x \in [1/4, 3/4], \quad \forall u \in [0, 1/2],$$

$$\begin{aligned} & \psi_{2,2} \left(|v|^2 \cos^2 \theta + |v_*|^2 \sin^2 \theta \right) - \cos^2 \theta \psi_{2,2} (|v|^2) - \sin^2 \theta \psi_{2,2} (|v_*|^2) = \\ & = 2 |v|^2 \log |v|^2 \left\{ (1-x) \psi_{2,1} (u^2) + x \psi_{2,1} (1) - \psi_{2,1} ((1-x)u^2 + x) \right\} \\ & + |v|^2 \left[(1-x) \psi_{2,2} (u^2) + x \psi_{2,2} (1) - \psi_{2,2} ((1-x)u^2 + x) \right] \geq C_8 |v|^2 \log |v|^2, \end{aligned}$$

for some constant $C_8 > 0$ depending on the strict convexity of $\psi_{2,1}$ and $\psi_{2,2}$. Hence we obtain for some constant $C_9 > 0$

$$H_{\psi_{2,2}}(v, v_*) \geq C_9 |v|^2 \log |v|^2 \mathbf{1}_{|v| \geq 2|v_*|}.$$

Putting together the estimates obtained on $G_{2,\zeta}$ and $H_{2,\zeta}$ we deduce

$$|v - v_*| K_{2,\zeta} \leq C_{10} \langle v \rangle^2 \langle v_* \rangle^2 \log \langle v \rangle \log \langle v_* \rangle - C_{11} |v - v_*| |v|^2 (\log |v|)^{\zeta-1} \mathbf{1}_{|v| \geq 2|v_*|}$$

for some constants $C_{10}, C_{11} > 0$. Since

$$\begin{aligned} |v - v_*| |v|^2 (\log |v|)^{\zeta-1} \mathbf{1}_{|v| \geq 2|v_*|} & \geq \text{Cst.} \langle v \rangle^3 (\log \langle v \rangle)^{\zeta-1} \mathbf{1}_{|v| \geq 2|v_*|} - \text{Cst.} \\ & \geq \text{Cst.} \langle v \rangle^3 (\log \langle v \rangle)^{\zeta-1} - \text{Cst.} \langle v \rangle^2 \langle v_* \rangle^2 \log \langle v \rangle \log \langle v_* \rangle \end{aligned}$$

we deduce

$$(6.16) \quad |v - v_*| K_{2,\zeta} \leq C_{12} (\langle v \rangle^2 + \langle v_* \rangle^2) - C_{13} \langle v \rangle^3 (\log \langle v \rangle)^{\zeta-1}$$

for some constants $C_{12}, C_{13} > 0$, and we finally obtain the differential inequality

$$\frac{d}{dt} \|f_t\|_{M_{2,\zeta}^1} \leq 2C_{12} (1 + \mathcal{E}) - C_{13} M_{3,\zeta-1},$$

from which (6.15) follows.

Step 3. Exponential time integral control of f and g in M_3^1 .

This step yields a proof of (6.6) and (6.7).

Let us first prove that

$$(6.17) \quad \forall t \geq 0, \quad e^{(2C_1+C_2) \int_0^t (\|f_s\|_{M_3^1} + \|g_s\|_{M_3^1}) ds} \leq e^{C_{14}(1+\mathcal{E})t} (\max\{M_k(f_0), M_k(g_0)\})^{\frac{1}{6}},$$

for some constant $C_{14} > 0$, for any $k \geq k_{\mathcal{E}}$, with $k_{\mathcal{E}}$ big enough depending on the energy \mathcal{E} .

We shall use the previous step and an interpolation argument. For any given probability measure

$$f \in P_k(\mathbb{R}^d) \quad \text{with} \quad M_2(f) \leq \mathcal{E},$$

we have for any $a > 2$

$$\begin{aligned} \|f\|_{M_{2,1}^1} &= \int_{\mathbb{R}^d} \langle v \rangle^2 \left(1 + \frac{\log(\langle v \rangle^2)}{2}\right) (\mathbf{1}_{\langle v \rangle^2 \leq a} + \mathbf{1}_{\langle v \rangle^2 \geq a}) f(dv) \\ &\leq (1 + \mathcal{E}) \left(1 + \frac{\log a}{2}\right) + \frac{1}{a} \int_{\mathbb{R}^d} \langle v \rangle^4 (1 + \log \langle v \rangle) f(dv) \\ &\leq (1 + \mathcal{E}) \left(1 + \frac{\log a}{2}\right) + \frac{1}{a} \|f\|_{M_5^1} \end{aligned}$$

where we have used inequality $\log x \leq x - 1$ on $x \geq 1$ in the last step.

By choosing

$$a := \|f\|_{M_5^1}^2,$$

we get

$$(6.18) \quad \|f\|_{M_{2,1}^1} \leq 2(1 + \mathcal{E}) \left(1 + \log \|f\|_{M_5^1}\right).$$

Remark 6.4. Observe here that it was absolutely crucial to be able to control the right hand side of (6.15) in terms of the $M_{2,1}^1$ moment, that is only a *logarithmic loss* of moment as compared to M_2^1 . This is what allows us to control this right hand side in terms of the *logarithm* of a higher moment of f , so that the exponential in (6.17) can be controlled in terms of some *polynomial* moment of f , hence fulfilling the requirement on the loss of weight of the stability on the semigroup for the abstract method. Recall indeed that the moment associated with the weight Λ_1 has to be *propagated* along the flow of the N -particle system. And we have been unable to control an exponential moment for such a high-dimension evolution.

On the other hand, the following elementary Hölder inequality holds

$$(6.19) \quad \forall k, k' \in \mathbb{N}, \quad k' \leq k, \quad \forall f \in M_k^1, \quad \|f\|_{M_{k'}^1} \leq \|f\|_{M^1}^{1-k'/k} \|f\|_{M_k^1}^{k'/k} \leq \|f\|_{M_k^1}^{k'/k}.$$

Then estimate (6.17) follows from (6.15), (6.18) and (6.19) with $k' = 5$ and $k \geq 5$ large enough in such a way that

$$(C_1 + C_2) C_4 2(1 + \mathcal{E}) \frac{5}{k} \leq \frac{1}{6}.$$

We then deduce (6.6) from (6.13), and (similarly) (6.7) from (6.12).

Step 4. Time integral control on d and h . Let us denote as before

$$d_t := f_t - g_t.$$

Let us prove

$$(6.20) \quad \left(\int_0^t \|d_s\|_{M_3^1} ds \right) \quad \text{and} \quad \left(\int_0^t \|h_s\|_{M_3^1} ds \right) \\ \leq C_{15} \|d_0\|_{M_2^1} e^{C_1 \int_0^t (\|f_s\|_{M_3^1} + \|g_s\|_{M_3^1}) ds} \left(C_3(\mathcal{E}) t + C_4 \|f_0\|_{M_{2,2}^1} \right) + \|d_0\|_{M_{2,1}^1}.$$

for some energy dependent constant $C_{\mathcal{E}}$ and some numerical constant C' . Performing similar computations to those leading to (6.9), we obtain

$$\frac{d}{dt} \|h_t\|_{M_{2,1}^1} \leq \iint |h_t|(dv) f_t(dv_*) |v - v_*| K_{2,1}(v, v_*) \\ + C_{16} \iiint |h_t|(dv) f_t(dv_*) |v - v_*| \langle v_* \rangle^2 (1 + \log \langle v_* \rangle)$$

for some constant $C_{16} > 0$ depending on b . Thanks to the Povzner inequality (6.16) (with $\zeta = 1$), we deduce for some constants $C_{17}, C_{18} > 0$

$$\frac{d}{dt} \|h_t\|_{M_{2,1}^1} \leq C_{17} \|h_t\|_{M_2^1} \|f_t\|_{M_{3,1}^1} - C_{18} \|h_t\|_{M_3^1}.$$

Integrating this differential inequality yields

$$\|h_t\|_{M_{2,1}^1} + C_{18} \int_0^t \|h_s\|_{M_3^1} ds \leq C_{17} \left(\sup_{s \in [0, t]} \|h_s\|_{M_2^1} \right) \left(\int_0^t \|f_s\|_{M_{3,1}^1} ds \right) + \|h_0\|_{M_{2,1}^1}.$$

Using the previous pointwise control on $\|h_t\|_{M_2^1}$ and (6.15) (with $\zeta = 2$) we get

$$\int_0^t \|h_s\|_{M_3^1} ds \leq \frac{C_{16}}{C_{17}} \|d_0\|_{M_2^1} e^{C_1 \int_0^t \|f_s\|_{M_3^1} ds} \left(C_3(\mathcal{E}) t + C_4 \|f_0\|_{M_{2,2}^1} \right) + \|d_0\|_{M_{2,1}^1}.$$

Arguing similarly for d_t , we deduce (6.20).

Step 5. Conclusion. We first rewrite (6.14) as

$$\sup_{s \in [0, t]} \|\omega_s\|_{M_2^1} \\ \leq C_2 \|d_0\|_{M_2^1} e^{(C_1 + C_2) \int_0^t (\|f_s\|_{M_3^1} + \|g_s\|_{M_3^1}) ds} \left(\int_0^t (\|d_s\|_{M_3^1} + \|h_s\|_{M_3^1}) ds \right).$$

Then we use the estimate (6.20) for the last term:

$$\sup_{s \in [0, t]} \|\omega_s\|_{M_2^1} \\ \leq C_2 C_{15} \|d_0\|_{M_2^1} \|d_0\|_{M_{2,1}^1} e^{(2C_1 + C_2) \int_0^t (\|f_s\|_{M_3^1} + \|g_s\|_{M_3^1}) ds} \left(C_3(\mathcal{E}) t + C_4 \|f_0\|_{M_{2,2}^1} \right).$$

Finally we use estimate (6.17) for the exponential term with $k = k_1$ and we obtain

$$\sup_{s \in [0, t]} \|\omega_s\|_{M_2^1} \leq C_2 C_{15} \|d_0\|_{M_2^1} \|d_0\|_{M_{2,1}^1} \times \\ \times e^{C_{14} (1 + \mathcal{E}) t} (\max \{M_k(f_0), M_k(g_0)\})^{\frac{1}{4}} \left(C_3(\mathcal{E}) t + C_4 \|f_0\|_{M_{2,2}^1} \right).$$

Then arguing as in the end of Step 3, for any $\eta \in (0, 1)$, using (6.19) with k_1 large enough, we have

$$\|d_0\|_{M_{2,1}^1} \leq (\max \{M_k(f_0), M_k(g_0)\})^{\frac{1}{6}} \|d_0\|_{M_2^1}^{1-\eta}$$

and

$$\|f_0\|_{M_{2,2}^1} \leq (\max \{M_k(f_0), M_k(g_0)\})^{\frac{1}{6}}.$$

We conclude with

$$\sup_{s \in [0, t]} \|\omega_s\|_{M_2^1} \leq e^{C(1+t)} \sqrt{\max\{M_k(f_0), M_k(g_0)\}} \|f_0 - g_0\|_{M^1}^{2-\eta},$$

from which estimate (6.8) follows. \square

6.7. Proof of (A4) uniformly in time. Let us start from an auxiliary result from [60]. Let us define the linearized Boltzmann collision operator at γ

$$\mathcal{L}_\gamma(f) = 2Q(\gamma, f)$$

where

$$\gamma = \frac{e^{-\frac{|v-u|^2}{2T}}}{(2\pi T)^d}$$

is a maxwellian distribution with momentum $u \in \mathbb{R}^d$ and temperature $T > 0$.

Theorem 6.5 (Theorem 1.2 in [60]). *First the linearized Boltzmann semigroup $e^{\mathcal{L}_\gamma t}$ for hard spheres satisfies*

$$(6.21) \quad \|e^{\mathcal{L}_\gamma t}\|_{L^1(m_z^{-1})} \leq C_z e^{-\lambda t}$$

where

$$m_z(v) := e^{z|v|}, \quad z > 0,$$

and $\lambda = \lambda(T)$ is the optimal rate, given by the first non-zero eigenvalue of the linearized operator \mathcal{L}_γ in the smaller space $L^2(\gamma^{-1})$.

Second the nonlinear Boltzmann semigroup S_t^{NL} satisfies

$$(6.22) \quad \|S_t^{NL}(f_0) - \gamma\|_{L^1(m_z^{-1})} \leq C_z e^{-\lambda t} \|f_0 - \gamma\|_{L^1(m_z^{-1})}$$

where $\gamma = \gamma_{f_0}$ is the maxwellian equilibrium associated with f_0 :

$$\gamma(v) = \frac{e^{-\frac{|v-u|^2}{2T}}}{(2\pi T)^d}$$

with

$$\forall i = 1, \dots, d, \quad u_i = \int_{\mathbb{R}^d} f_0 v_i dv \quad \text{and} \quad T = \frac{1}{d} \int_{\mathbb{R}^d} f_0 |v|^2 dv$$

and C_z is some constant possibly depending on z and the energy \mathcal{E} of the solution considered, and $\lambda = \lambda(T)$ is the same rate function as before.

Let us now prove *uniform in time* estimate for the expansion of the limiting semigroup in terms of the initial data.

Lemma 6.6. *For any given energy $\mathcal{E} > 0$ and for any $\eta \in (0, 1)$, there exists*

- some constant $k_{\mathcal{E}, \eta} \geq 2$ (depending on \mathcal{E} and η),
- some constant C (depending on \mathcal{E}),

such that for any

$$f_0, g_0 \in P_{G_1}(\mathbb{R}^d)$$

satisfying

$$\forall i = 1, \dots, d, \quad \langle f_0, v_i \rangle = \langle g_0, v_i \rangle \quad \text{and} \quad \langle f_0, |v|^2 \rangle = \langle g_0, |v|^2 \rangle \leq \mathcal{E},$$

and for any $k \geq k_{\mathcal{E}, \eta}$ we have

$$(6.23) \quad \|g_t - f_t\|_{M_2^1} \leq C e^{-\frac{\lambda}{2}t} \sqrt{\max\{M_k(f_0), M_k(g_0)\}} \|g_0 - f_0\|_{M_2^1}^{1-\eta},$$

$$(6.24) \quad \|h_t\|_{M_2^1} \leq C e^{-\frac{\lambda}{2}t} \sqrt{M_k(f_0)} \|g_0 - f_0\|_{M_2^1}^{1-\eta},$$

$$(6.25) \quad \|\omega_t\|_{M_2^1} \leq C e^{-\frac{\lambda}{2}t} \sqrt{\max\{M_k(f_0), M_k(g_0)\}} \|g_0 - f_0\|_{M_2^1}^{2-\eta}.$$

Note that these estimates imply **(A4)** with $T = \infty$, $P_{\mathcal{G}_2} = P_{\mathcal{G}_1}$.

Remark 6.7. Observe that in the following proof we shall use *moment production bounds* on the limiting equation. Indeed once stability estimates for small times have been secured (as in Lemma 6.3), one can use, for $t \geq T_0 > 0$, moments production whose bounds only depend on the energy of the solution. This, together with the linearized theory in L^1 setting with exponential moment bounds of Theorem 6.5, will be the key to the following proof.

Proof of Lemma 6.6. From [48, Theorem 1.2-(b)] (see also [1] for a simpler different proof), there exists some constants z, Z (only depending on the collision kernel and the energy of the solutions) such that

$$(6.26) \quad \sup_{t \geq 1} \|f_t + g_t + h_t\|_{L_{m_{2z}}^1} \leq Z, \quad m_{2z}(v) := e^{2z|v|}$$

(note that the proof in [48] applies to the solutions f_t and g_t , however it is straightforward to apply exactly the same proof to the linearized solution h_t around f_t , once exponential moment is known on f_t).

We also know from (6.22) that (maybe by choosing a larger Z)

$$(6.27) \quad \forall t \geq 1 \quad \|f_t - \gamma\|_{L_{m_{2z}}^1} + \|g_t - \gamma\|_{L_{m_{2z}}^1} \leq 2Z e^{-\lambda t},$$

where

$$\gamma := \gamma_{f_0} = \gamma_{g_0} = \frac{e^{-\frac{|v-u|^2}{2T}}}{(2\pi T)^d}$$

with

$$\forall i = 1, \dots, d, \quad u_i = \langle f, v_i \rangle = \langle g, v_i \rangle \quad \text{and} \quad T = \frac{1}{d} \langle f, |v|^2 \rangle = \frac{1}{d} \langle g, |v|^2 \rangle$$

stands for the normalized maxwellian associated to f_0 and g_0 .

We write

$$\begin{aligned} \partial_t(f_t - g_t) &= Q(f_t - g_t, f_t + g_t) \\ &= \mathcal{L}_\gamma(f_t - g_t) + Q(f_t - g_t, f_t - \gamma) + Q(f_t - g_t, g_t - \gamma) \end{aligned}$$

and, using also (6.21) on the linearized semigroup, we deduce for

$$u(t) := \|f_t - g_t\|_{M_{m_z}^1}$$

the following differential inequality for $t \geq T_0 \geq 1$

$$\begin{aligned} u(t) &\leq C e^{-\lambda(t-T_0)} u(T_0) \\ &\quad + C \int_{T_0}^t e^{-\lambda(t-s)} \left(\|Q(f_s - g_s, f_s - \gamma)\|_{M^1(m_z)} + \|Q(f_s - g_s, g_s - \gamma)\|_{M^1(m_z)} \right) ds \end{aligned}$$

(this formal inequality and next ones can easily be justified rigorously by a regularizing procedure and using a uniqueness result for measure solutions such as [32, 29, 48]).

Therefore we obtain

$$u(t) \leq C e^{-\lambda(t-T_0)} u(T_0) + C \int_{T_0}^t e^{-\lambda(t-s)} \left(\|f_s - \gamma\|_{M^1(\langle v \rangle m_z)} + \|g_s - \gamma\|_{M^1(\langle v \rangle m_z)} \right) \|f_s - g_s\|_{M^1(\langle v \rangle m_z)} ds.$$

We then use the control of $M^1(\langle v \rangle m_z)$ by $M^1(m_{2z})$ together with the controls (6.26)-(6.27), the decay control (6.21) and the estimate

$$e^{-\lambda s - \lambda(t-s)} \leq e^{-\frac{\lambda}{2} t - \frac{\lambda}{2} s}.$$

We get

$$u(t) \leq C e^{-\lambda(t-T_0)} u(T_0) + C e^{-\frac{\lambda}{2} t} \int_{T_0}^t e^{-\frac{\lambda}{2} s} \|f_s - g_s\|_{M^1(\langle v \rangle m_z)} ds.$$

We then use the following control for any $a > 0$:

$$\begin{aligned} \|f - g\|_{M^1_{\langle v \rangle m_z}} &= \int |f - g| \langle v \rangle e^{z|v|} \\ &\leq a \int_{|v| \leq a} |f - g| e^{z|v|} + e^{-za} \int_{|v| \geq a} (f + g) e^{2z|v|} \\ &\leq a u(t) + e^{-za} Z. \end{aligned}$$

Hence we get

$$\|f - g\|_{M^1_{\langle v \rangle m_z}} \leq \begin{cases} u + e^{-z} Z \leq (1 + Z) u & \text{when } u \geq 1, \quad (\text{choosing } a := 1) \\ \frac{1}{z} |\log u| u + u Z & \text{when } u \leq 1 \quad (\text{choosing } -za := \log u) \end{cases}$$

and we deduce

$$\|f - g\|_{M^1_{\langle v \rangle m_z}} \leq K u (1 + (\log u)_-), \quad K := 1 + \frac{1}{z} + Z.$$

Then for any $\delta \in (0, 1)$, we have, by choosing T_0 large enough,

$$\forall t \geq T_0, \quad e^{-\frac{\lambda}{2} t} \leq \delta e^{-\frac{\lambda}{4} t}$$

and we conclude with the following integral inequality

$$(6.28) \quad u(t) \leq C e^{-\lambda(t-T_0)} u(T_0) + \delta e^{-\frac{\lambda}{4} t} \int_{T_0}^t e^{-\frac{\lambda}{2} s} u_s (1 + (\log u_s)_-) ds.$$

Let us prove that this integral inequality implies

$$(6.29) \quad \forall t \geq T_0, \quad u(t) \leq C e^{-\frac{\lambda}{4} t} u(T_0)^{1-\delta}.$$

Consider the case of equality in (6.28). Then we have

$$u(t) \geq e^{-\lambda(t-T_0)} u(T_0)$$

and therefore

$$(1 + (\log u_t)_-) \leq (1 + (\log u(T_0))_- + \lambda(t - T_0)).$$

We then have

$$\begin{aligned} U(t) &:= \int_{T_0}^t e^{-\frac{\lambda}{2} s} u_s (1 + (\log u_s)_-) ds \\ &\leq \int_{T_0}^t e^{-\frac{\lambda}{2} s} u_s (1 + (\log u(T_0))_- + \lambda(s - T_0)) ds \\ &\leq (1 + (\log u(T_0))_-) \int_{T_0}^t e^{-\frac{\lambda}{4} s} u_s ds. \end{aligned}$$

By a Gronwall-like argument we can therefore obtain

$$u(t) \leq C e^{-\lambda(t-T_0)} u(T_0) + C \delta e^{-\frac{\lambda}{4}t} (1 + (\log u(T_0))_-) u(T_0).$$

Then thanks to the inequality

$$\forall x \in (0, 1], \quad -(\log x) x \leq \frac{x^{1-\delta}}{\delta}$$

we can prove (6.29) when $u(T_0) \leq 1$, and in the case when $u(T_0) \geq 1$, we can use (6.26) again to get

$$u(T_0) \leq (2Z)^\delta u(T_0)^{1-\delta}.$$

This concludes the proof of the claimed inequality (6.29).

Then estimate (6.23) follows by choosing δ small enough (in relation to η) and then connecting the last estimate (6.29) from time T_0 on together with the previous finite time estimate (6.6) from time 0 until time T_0 .

Then the estimate (6.24) is proved exactly in the same way by using the equation

$$\partial_t h_t = \mathcal{L}_\gamma(f_t - g_t) + Q(h_t, f_t - \gamma)$$

(which is even simpler than the equation for $f_t - g_t$).

Concerning the estimate (6.25) we start from the equation

$$\partial_t \omega_t = 2 \mathcal{L}_\gamma(h_t) + Q(\omega_t, f_t - \gamma) + Q(\omega_t, g_t - \gamma) + Q(h_t, d_t).$$

Then we establish on

$$y(t) := \|\omega_t\|_{M_{m_z}^1}$$

the following differential inequality

$$y(t) \leq C e^{-\lambda(t-T_0)} y(T_0) + C \delta e^{-\frac{\lambda}{4}t} (1 + (\log y(T_0))_-) y(T_0) + C e^{-\frac{\lambda}{2}t} \|d_{T_0}\|_{M_{m_z}^1}^{1-\delta} \|h_{T_0}\|_{M_{m_z}^1}^{1-\delta}$$

which implies

$$y(t) \leq C e^{-\frac{\lambda}{4}t} \left(y(T_0)^{1-\delta} + \|d_{T_0}\|_{M_{m_z}^1}^{1-\delta} \|h_{T_0}\|_{M_{m_z}^1}^{1-\delta} \right).$$

Then estimate (6.25) follows by choosing δ small enough (in relation to η) and then connecting the last estimate from time T_0 on together with the previous finite time estimate (6.8) from time 0 until time T_0 . \square

6.8. Proof of (A5) uniformly in time. Let us prove that for any $\bar{z}, \mathcal{M}_{\bar{z}} \in (0, \infty)$ there exists some continuous function

$$\Theta : \mathbb{R}_+ \rightarrow \mathbb{R}_+, \quad \Theta(0) = 0,$$

such that for any

$$f_0, g_0 \in P_{m_{\bar{z}}}(\mathbb{R}^d), \quad m_{\bar{z}}(v) := e^{\bar{z}|v|}$$

such that

$$\|f_0\|_{M_{m_{\bar{z}}}^1} \leq \mathcal{M}_{\bar{z}}, \quad \|g_0\|_{M_{m_{\bar{z}}}^1} \leq \mathcal{M}_{\bar{z}},$$

there holds

$$\sup_{t \geq 0} W_1(S_t^{NL}(f_0), S_t^{NL}(g_0)) \leq \Theta(W_1(f_0, g_0)).$$

where W_t stands for the Kantorovich-Rubinstein distance. Let us denote

$$W_t := W_1(S_t^{NL}(f_0), S_t^{NL}(g_0)).$$

As we shall see, we may choose

$$(6.30) \quad \Theta(w) := \min \left\{ \bar{\Theta}, \bar{\Theta} e^{1-(1+|\log w|)^{1/2}}, \frac{C_1}{(1+|\log w|)^{\frac{\lambda}{2K}}} + C_2 w^\alpha \right\}, \quad \Theta(0) = 0,$$

for some constants $\bar{\Theta}$, C_1 , $C_2 > 0$ (only depending on \bar{z} and $\mathcal{M}_{\bar{z}}$).

We start off with the inequality

$$\forall t \geq 0 \quad W_t \leq \|(f_t - g_t)|v\|_{M^1} \leq \frac{1}{2} \|(f_t + g_t)\langle v \rangle^2\|_{M^1} = 1 + \mathcal{E} =: \bar{\Theta}.$$

Let us now improve this inequality for small value of W_0 . Therefore we assume without restriction that

$$W_0 \leq \frac{1}{2}$$

in the sequel.

On the one hand, it has been proved in [32, Theorem 2.2 and Corollary 2.3] that

$$(6.31) \quad W_t \leq W_0 + K \int_0^t W_s (1 + (\log W_s)_-) ds,$$

for some constant K .

One can then check that the function

$$\bar{W}(t) := e^{1-e^{-Kt}} (W_0)^{e^{-Kt}}$$

satisfies

$$\bar{W}'(t) = K (1 - \log \bar{W}(t)) \bar{W}(t), \quad \bar{W}(0) = W_0.$$

Therefore it is a super-solution of the differential inequality (6.31) as long as $W_t \leq 1$. It is an easy computation that this super-solution satisfies

$$\bar{W}(t) \leq 1 \quad \text{as long as} \quad t \leq t_0 := \frac{\log(1 + |\log W_0|)}{K}.$$

Observe also that $\bar{W}(t)$ is increasing on $t \in [0, t_0]$.

We then define

$$(6.32) \quad t_1 := \frac{t_0}{2} = \frac{\log(1 + |\log W_0|)}{2K}$$

and we deduce the following bound on the solution of (6.31):

$$(6.33) \quad \forall t \in [0, t_1], \quad W_t \leq \bar{W}_t \leq \bar{W}_{t_1} = e^{1-(1+|\log W_0|)^{1/2}}.$$

On the other hand, from (6.22), there are constants $\lambda, Z > 0$, $z \in (0, \bar{z})$ such that

$$(6.34) \quad \forall t \geq 0 \quad \|f_t - \gamma_{f_0}\|_{L_{m_z}^1} + \|g_t - \gamma_{g_0}\|_{L_{m_z}^1} \leq Z e^{-\lambda t},$$

where γ_{f_0} and γ_{g_0} stand again for the normalized maxwellian associated to f_0 and g_0 .

We denote by u_{f_0} and u_{g_0} the momentum of f_0 and g_0

$$u_{f_0} = \langle f_0, v \rangle, \quad u_{g_0} = \langle g_0, v \rangle,$$

by T_{f_0} and T_{g_0} the temperature of f_0 and g_0

$$T_{f_0} = \frac{1}{d} \langle f_0, |v - u_{f_0}|^2 \rangle, \quad T_{g_0} = \frac{1}{d} \langle g_0, |v - u_{g_0}|^2 \rangle,$$

and by \mathcal{E}_{f_0} and \mathcal{E}_{g_0} the energy of f_0 and g_0

$$\mathcal{E}_{f_0} = \langle f_0, |v|^2 \rangle, \quad \mathcal{E}_{g_0} = \langle g_0, |v|^2 \rangle.$$

Then we can use (4.5) for some

$$s > 0 \text{ large enough} \quad \text{and} \quad \alpha \in \left(0, \frac{1}{2s}\right)$$

and we compute

$$\begin{aligned}
W_1(\gamma_{f_0}, \gamma_{g_0}) &\leq C \left(\|\gamma_{f_0} - \gamma_{g_0}\|_{H^{-s}}^2 \right)^\alpha \\
&\leq C \left(\int_{\mathbb{R}^d} \frac{\left| e^{-\theta_{f_0} \frac{|\xi|^2}{2} - i u_{f_0} \sqrt{\theta_{f_0}} \xi} - e^{-\theta_{g_0} \frac{|\xi|^2}{2} - i u_{g_0} \sqrt{\theta_{g_0}} \xi} \right|^2}{\langle \xi \rangle^{2s}} d\xi \right)^\alpha \\
&\leq C \left(\int_{\mathbb{R}^d} \frac{|T_{f_0} - T_{g_0}|^2 |\xi|^4 + |u_{f_0} \sqrt{T_{f_0}} - u_{g_0} \sqrt{T_{g_0}}|^2 |\xi|^2}{\langle \xi \rangle^{2s}} d\xi \right)^\alpha \\
&\leq C \left(|T_{f_0} - T_{g_0}|^{2\alpha} + |u_{f_0} - u_{g_0}|^{2\alpha} \right) \\
&\leq C \left(|\mathcal{E}_{f_0} - \mathcal{E}_{g_0}|^{2\alpha} + |u_{f_0} - u_{g_0}|^{2\alpha} \right) \\
&\leq C \left(W_2(f_0, g_0)^{2\alpha} + W_1(f_0, g_0)^{2\alpha} \right) \\
(6.35) \quad &\leq C \left(W_1(f_0, g_0)^\alpha + W_1(f_0, g_0)^{2\alpha} \right)
\end{aligned}$$

where the final constant $C > 0$ depends on $s > 0$ as well as on upper and lower bounds on the temperatures of f_0 and g_0 .

Remark 6.8. We also refer to [21] for more general estimates of the Wasserstein distance between two gaussians. Such refined estimates are however not needed in our study.

Then we combine (6.34) and (6.35) by triangular inequality to get

$$(6.36) \quad \forall t \geq 0 \quad W_t \leq C_1 e^{-\lambda t} + C_2 W_0^\alpha$$

for two constants $C_1, C_2 > 0$.

We then consider times $t \geq t_1$ and we deduce from (6.32) and (6.36) the following bound from above

$$(6.37) \quad \forall t \geq t_1, \quad W_t \leq C_1 e^{-\lambda t_1} + C_2 W_0^\alpha = \frac{C_1}{(1 + |\log W_0|)^{\frac{\lambda}{2K}}} + C_2 W_0^\alpha.$$

It is then straightforward to conclude the proof of **(A5)** uniformly in time for the function (6.30) by combining (6.33) and (6.37).

By using Theorem 3.1 whose assumptions have been proved above, this concludes the proof of point (i) in Theorem 6.1 together with the estimate on $\mathcal{W}_{W_1}^N(f)$ from Lemma 4.2.

One can also conclude the proof of point (ii) in Theorem 6.1 by using

- Lemma 4.4 for the construction of the sequence initial data f_0^N which satisfies the required integral and support moment bounds,
- The previous in order to apply Theorem 3.1,
- Lemma 4.7 in order to estimate

$$\mathcal{W}_{W_1}(\pi_P^N(f_0^N), f_0) \xrightarrow{N \rightarrow \infty} 0.$$

6.9. Proof of infinite-dimensional Wasserstein chaos. Let us now prove Theorem 6.2. Its proof is similar to Theorem 5.3.

First the proof of (6.3) as a consequence of point (i) in Theorem 6.1 and [37, Theorem 1.1] is strictly similar to the one of (5.8) as a consequence of point (i) in Theorem 5.1 and [37, Theorem 1.1].

Then the proof of (6.4) is also similar to the one of (5.9), the only difference being that one needs the following result of lower bound (independent of N) on the spectral gap of the N -particle system.

Theorem 6.9 ([11]). *Consider the operator L_{HS} for the hard spheres N -particle model with collision kernel $B(v-w) = |v-w|$. Then there is a constant $\lambda > 0$ such that for any probability f^N on \mathcal{S}^N one has*

$$\langle L_{HS}f^N, f^N \rangle_{L^2(\mathcal{S}^N)} \leq -\lambda \|f^N\|_{L^2(\mathcal{S}^N)}.$$

where \mathcal{S}^N is the hypersurface

$$\mathcal{S}^N := \left\{ V \in E^N; \frac{1}{N} \sum_{k=1}^N v_k = 0 \quad \text{and} \quad \frac{1}{N} \sum_{k=1}^N |v_k|^2 = 1 \right\}.$$

and for some constant $\lambda > 0$ independent of N .

Then using Theorem 6.9 we deduce that

$$\forall N \geq 1, \forall t \geq 0 \quad \|h^N - 1\|_{L^2(\mathcal{S}^N, \gamma^N)} \leq e^{-\lambda t} \|h_0^N - 1\|_{L^2(\mathcal{S}^N, \gamma^N)},$$

where $h^N = df^N/d\gamma^N$ is the Radon-Nikodym derivative of f^N with respect to the measure γ^N and the end of proof of (6.4) is then exactly similar to the one of (5.9).

7. H -THEOREM AND ENTROPIC CHAOS

This section is concerned with the H -theorem. We answer a question raised by Kac [42] about the derivation of the H -theorem.

7.1. Statement of the results. Our main results of this section state as follows:

Theorem 7.10. *Consider the Boltzmann collision process for Maxwell molecules (with or without cutoff) or hard spheres, and some initial data*

$$f_0 \in L^\infty(\mathbb{R}^d) \quad \text{s. t.} \quad \int_{\mathbb{R}^d} e^{z|v|} df_0(v) < +\infty$$

for some $z > 0$, and the sequence of N -particle initial data $(f_0^N)_{N \geq 1}$ constructed in Lemma 4.4 and 4.7.

Then we have:

- (i) *In the case of Maxwell molecules with cut-off and hard spheres, if the initial data is entropy-chaotic in the sense*

$$\frac{1}{N} H(f_0^N | \gamma^N) \xrightarrow{N \rightarrow +\infty} H(f_0 | \gamma),$$

with

$$H(f_0^N) := \int_{\mathcal{S}^N} h_0^N \log h_0^N \gamma^N(dV), \quad h_0^N := \frac{df_0^N}{d\gamma^N},$$

then the solution is also entropy chaotic for any later time:

$$\forall t \geq 0, \quad \frac{1}{N} H(f_t^N | \gamma^N) \xrightarrow{N \rightarrow +\infty} H(f_t | \gamma).$$

This proves the derivation of the H -theorem in this context.

- (ii) *In the case of Maxwell molecules, and assuming moreover that the Fisher information of the initial data f_0 is finite:*

$$\int_{\mathbb{R}^d} \frac{|\nabla_v f_0|^2}{f_0} dv < +\infty,$$

the following estimate on the relaxation induced by the H -theorem uniformly in the number of particles also holds:

$$\forall N \geq 1, \quad \frac{1}{N} H(f_t^N | \gamma^N) \leq \beta(t)$$

for some function $\beta(t) \rightarrow 0$ as $t \rightarrow \infty$.

Remarks 7.11. (1) The assumptions on the initial data could be relaxed to just $P_4 \cap L^\infty$ as in point (iii) of Theorem 5.1. However this stronger allows to apply the previous theorems for hard spheres and Maxwell molecules. We do not search for optimal statement here, but rather emphasize the ideas of our method.

(2) A stronger notion of entropy chaoticity could be

$$\frac{1}{N} H(f^N | [f^{\otimes N}]_{\mathcal{S}^N}) \xrightarrow{N \rightarrow +\infty} 0.$$

The propagation of such a stronger property is an interesting open question.

(3) The point (ii) holds for the hard spheres conditionally to a bound on the Fisher information uniformly in time and in the number of particle. However at now, proving such a bound for the many-particle hard spheres jump process is an open problem.

(4) In point (ii) one could expect to have

$$\forall N \in \mathbb{N}^*, \forall t \geq 0 \quad \frac{1}{N} H(f_t^N | \gamma^N) \leq C e^{-\lambda t}.$$

7.2. Propagation of entropic chaos and derivation of the H -theorem. In this subsection we shall prove the point (i) of Theorem 7.10. Its proof relies on a convexity argument.

Let us define $h^N := df^N/d\gamma^N$ and then compute

$$\begin{aligned} \frac{d}{dt} \frac{1}{N} H(f_t^N | \gamma^N) &= -D^N(f_t^N) \\ &:= -\frac{1}{2N^2} \int_{\mathcal{S}^N} \sum_{i \neq j} \int_{\mathbb{S}^{d-1}} (h_t^N(V_{ij}^*) - h_t^N(V)) \log \frac{h_t^N(V_{ij}^*)}{h_t^N(V)} B(v_i - v_j, \sigma) d\sigma \gamma^N(dV) \end{aligned}$$

where we recall that V_{ij}^* was defined in (5.1), which implies

$$(7.38) \quad \forall t \geq 0, \quad \frac{1}{N} H(f_t^N | \gamma^N) + \int_0^t D^N(f_s^N) ds = \frac{1}{N} H(f_0^N | \gamma^N).$$

We also note that the same kind of equality is true at the limit (see e.g. [47])

$$\forall t \geq 0, \quad H(f_t | \gamma) + \int_0^t D^\infty(f_s) ds = H(f_0 | \gamma)$$

with

$$D^\infty(f) := \frac{1}{2} \int_{\mathbb{R}^d \times \mathbb{R}^d \times \mathbb{S}^{d-1}} (f' f'_* - f f_*) \log \frac{f' f'_*}{f f_*} B(v - v_*, \sigma) dv dv_* d\sigma$$

(be careful to the factor 1/2 in our definition of the collision operator (1.2) when computing the entropy production functional).

We then have the following lower semi-continuity property on these functionals, as a consequence of their convexity property.

Lemma 7.12. *The many-particle relative entropy and entropy production functionals defined above are lower semi-continuous: if the sequence $(f^N)_{N \geq 1}$ is f -chaotic then*

$$\liminf_{N \rightarrow \infty} \frac{1}{N} H(f^N | \gamma^N) \geq H(f | \gamma)$$

and

$$\liminf_{N \rightarrow \infty} \frac{1}{N} D^N(f^N) \geq D^\infty(f).$$

Let us first explain how to conclude the proof of point (i) of Theorem 7.10 with this lemma at hand. We first deduce from (7.38) and the entropic chaoticity of the initial data that

$$\begin{aligned} \forall t \geq 0, \quad \frac{1}{N} H(f_t^N | \gamma^N) + \int_0^t D^N(f_s^N) ds &\xrightarrow{N \rightarrow \infty} H(f_0 | \gamma) \\ &= H(f_t | \gamma) + \int_0^t D^\infty(f_s) ds \end{aligned}$$

Second we use Lemma 7.12 on the LHS to deduce that

$$\forall t \geq 0, \quad \liminf_{N \rightarrow \infty} \left(\frac{1}{N} H(f_t^N | \gamma^N) + \int_0^t D^N(f_s^N) ds \right) \geq H(f_t | \gamma) + \int_0^t D^\infty(f_s) ds$$

where each of the limit of the two non-negative terms on the LHS is greater than the corresponding non-negative term in the RHS. We deduce from the two last equations that necessarily

$$\forall t \geq 0, \quad \frac{1}{N} H(f_t^N | \gamma^N) \xrightarrow{N \rightarrow \infty} H(f_t | \gamma)$$

and

$$\forall t \geq 0, \quad \int_0^t D^N(f_s^N) ds \xrightarrow{N \rightarrow \infty} \int_0^t D^\infty(f_s) ds$$

which concludes the proof of point (i) of Theorem 7.10.

Proof of Lemma 7.12. These inequalities are consequences of convexity properties. The lower continuity property on the relative entropy on the spheres was proved in [12, Theorem 12] (actually the proof in this reference is performed on the sphere \mathbb{S}^{N-1} , but extending it to the invariant subspaces of our jump processes \mathcal{S}^N is straightforward). We refer to [16] for a detailed proof of the latter.

Let us now prove the inequality for the entropy production functional D^N . Denoting $Z = h^N(V_{12}^*)/h^N$, we first rewrite thanks to the symmetry of f^N as

$$D^N(f^N) = \frac{N(N-1)}{2N^2} \int_{\mathcal{S}^N} \int_{\mathbb{S}^{d-1}} J(Z) B(v_1 - v_2, \sigma) f_2^N(v_1, v_2) d\sigma \frac{f^N(dV)}{f_2^N(v_1, v_2)},$$

where $J(z) := (z-1) \log z$ and f_2^N denotes the 2-marginal. Since the function $z \mapsto J(z)$ is convex, we can apply a Jensen inequality according to the variables v_3, \dots, v_N with reference probability measure f^N/f_2^N , which yields

$$D^N(f^N) \geq \frac{N(N-1)}{2N^2} \int_{v_1, v_2 \in \mathbb{R}^d} \int_{\mathbb{S}^{d-1}} J(\bar{Z}) B(v_1 - v_2, \sigma) f_2^N(v_1, v_2) d\sigma dv_1 dv_2$$

with

$$\bar{Z}(v_1, v_2) := \int_{v_3, \dots, v_N \in \mathcal{S}^N(v_1, v_2)} Z \frac{f^N(V)}{f_2^N(v_1, v_2)} = \frac{f_2^N((V_{12}^*)_1, (V_{12}^*)_2)}{f_2^N(v_1, v_2)},$$

where $\mathcal{S}^N(v_1, v_2) := \{v_3, \dots, v_N \in E^{N-2}, (v_1, \dots, v_N) \in \mathcal{S}^N\}$. We therefore deduce a control from below of the N -particle entropy production functional in terms of the 2-particle entropy production functional, denoting $(f_2^N)^* = f_2^N((V_{12}^*)_1, (V_{12}^*)_2)$:

$$D^N(f^N) \geq \frac{N(N-1)}{2N^2} \int_{v_1, v_2 \in \mathbb{R}^d} \int_{\mathbb{S}^{d-1}} ((f_2^N)^* - f_2^N) \log \frac{(f_2^N)^*}{f_2^N} B(v_1 - v_2, \sigma) d\sigma dv_1 dv_2$$

Finally we take advantage of the convexity of the functional

$$h(x, y) = (x - y) \log \frac{x}{y}$$

which implies that the function

$$(f_2, g_2) \rightarrow \int_{v_1, v_2 \in \mathbb{R}^d} \int_{\sigma \in \mathbb{S}^{d-1}} (f_2 - g_2) \log \frac{f_2}{g_2} B(v_1 - v_2, \sigma)$$

is lower semi-continuous for the weak convergence of the 2-particle distributions f_2 and g_2 as proved in [24, Step 2 of the proof]. Hence we obtain thanks to the chaoticity of the second marginal

$$\begin{aligned} & \liminf_{N \rightarrow \infty} D^N(f^N) \\ & \geq \frac{1}{2} \int_{v, v_* \in \mathbb{R}^d} \int_{\mathbb{S}^{d-1}} (f(v')f(v'_*) - f(v)f(v_*)) \log \frac{f(v')f(v'_*)}{f(v)f(v_*)} B(v - v_*, \sigma) d\sigma dv dv_* \\ & = D^\infty(f) \end{aligned}$$

which concludes the proof. \square

7.3. Many-particle relaxation rate in the H -theorem. In this subsection we shall prove point (ii) in Theorem 7.10. Its proof goes in two steps. First we shall prove that it follows from an estimate on the Fisher information thanks to the so-called ‘‘HWI’’ interpolation inequality [72]. Second we shall prove such a uniform bound on the Fisher information in the case of Maxwell molecules. Let us take the opportunity to thank Maxime Hauray who kindly communicated to us a proof for the latter step.

Let us define the Fisher informations for the N -particle distribution:

$$I(f^N) := \int_{\mathbb{R}^{dN}} \frac{|\nabla f^N|^2}{f^N} dV$$

and

$$I(f^N | \gamma^N) := \int_{\mathcal{S}^N} \frac{|\nabla_{\mathcal{S}^N} h^N|^2}{h^N} \gamma^N(dV), \quad h^N := \frac{df^N}{d\gamma^N}$$

for a probability f^N having a density with respect to the Lebesgue measure in \mathbb{R}^{dN} and with respect to the measure γ^N respectively. The gradient in this formula has to be understood as the usual Riemannian geometry gradient in the manifold \mathcal{S}^N . The tangent space $T\mathcal{S}_V^N$ (of dimension $Nd - 2$) at some given point $V \in \mathcal{S}^N$ is given by

$$T\mathcal{S}_V^N = \left\{ W \in \mathbb{R}^{dN} \text{ s. t. } \sum_{i=1}^N w_i = 0 \text{ and } W \perp V \right\}.$$

For more informations and other results on the Fisher informations on \mathcal{S}^{N-1} we refer to [5]. We shall prove the following lemma whose proof is inspired from [71]).

Lemma 7.13. *Consider the N -particle jump process for Maxwell molecules as defined in Subsection 5.1 for a give $N \geq 1$, and some initial data f_0^N whose Fisher information is finite $I(f_0^N | \gamma^N) < +\infty$ on \mathcal{S}^N . Then one has the following uniform in time bound on the Fisher information of the solution*

$$\forall t \geq 0, \quad I(f_t^N | \gamma^N) \leq I(f_0^N | \gamma^N).$$

Proof of Lemma 7.13. We shall first consider the case of cutoff Maxwell molecules whose collision kernel b is integrable, and for a positive and smooth solution f^N on \mathcal{S}^N . These assumptions can be relaxed by a mollification argument.

It is possible to study directly the estimate to be proved on the manifold \mathcal{S}^N , however it means that one has to consider some local coordinates and local basis for the tangent space. Another simpler method is to take advantage of the fact that the dynamics leaves the energy unchanged.

Starting from an initial data f_0^N on \mathcal{S}^N we consider the flatened initial data

$$\tilde{f}_0^N := \alpha(\mathcal{E}(V)) \beta(m(V)) f_0^N \left(\frac{V - m(V)}{\sqrt{\mathcal{E}(V)}} \right)$$

$$\text{with } \mathcal{E}(V) = \frac{\sum_{i=1}^N |v_i|^2}{N} \text{ and } m(V) = \frac{\sum_{i=1}^N v_i}{N}.$$

Observe that from the conservation of energy and momentum and the uniqueness of the solutions to the linear master N -particle equation

$$\forall t \geq 0, \quad \tilde{f}_t^N := \alpha(\mathcal{E}(V)) \beta(m(V)) f_t^N \left(\frac{V - m(V)}{\sqrt{\mathcal{E}(V)}} \right)$$

where \tilde{f}_t^N denotes the solution in \mathbb{R}^{dN} starting from \tilde{f}_0^N . If the functions α and β are regular and compactly supported, as well as f_0^N , this produces a smooth solution on \mathbb{R}^{dN} .

Assume that the result on the Fisher information is true in \mathbb{R}^{dN} :

$$I(\tilde{f}_t^N) := \int_{\mathbb{R}^{dN}} \frac{|\nabla \tilde{f}_t^N|^2}{\tilde{f}_t^N} dV \leq \int_{\mathbb{R}^{dN}} \frac{|\nabla \tilde{f}_0^N|^2}{\tilde{f}_0^N} dV = I(\tilde{f}_0^N).$$

Then we have the orthogonal decomposition of the gradient locally in terms of radial and ortho-radial directions

$$\nabla_{\mathbb{R}^{dN}} \tilde{f}_t^N = \nabla_{\mathcal{E}} \tilde{f}_t^N + \nabla_m \tilde{f}_t^N + \nabla_{\mathcal{S}^N} \tilde{f}_t^N = (\nabla_{\mathcal{E}} \alpha + \nabla_m \beta) \tilde{f}_t^N + \nabla_{\mathcal{S}^N} \tilde{f}_t^N$$

that we can plug into the Fisher information inequality:

$$\begin{aligned} I(\tilde{f}_t^N) &:= (\nabla_{\mathcal{E}} \alpha + \nabla_m \beta)^2 + \left(\int_{\mathcal{E}, m} \alpha(\mathcal{E}) \beta(m) \right) \int_{\mathcal{S}^N} \frac{|\nabla_{\mathcal{S}^N} h_t^N|^2}{h_t^N} \gamma^N(dV) \\ &\leq (\nabla_{\mathcal{E}} \alpha + \nabla_m \beta)^2 + \left(\int_{\mathcal{E}, m} \alpha(\mathcal{E}) \beta(m) \right) \int_{\mathcal{S}^N} \frac{|\nabla_{\mathcal{S}^N} h_0^N|^2}{h_0^N} \gamma^N(dV) = I(\tilde{f}_0^N). \end{aligned}$$

Dropping the terms which do not depend on time we obtain the desired inequality on \mathcal{S}^N .

Let us now prove the inequality on \mathbb{R}^{dN} . Let us first fix some notation: the N -particle solution f_t^N satisfies

$$\begin{aligned} \partial_t f^N &= \frac{1}{N} \sum_{i,j=1, i \neq j}^N \int_{\mathbb{S}^{d-1}} (f^N(r_{ij,\sigma}(V)) b(\cos \theta_{ij}) d\sigma - f^N(V)) d\sigma \\ &=: NB(Q^{+,N}(f^N) - f^N) \end{aligned}$$

where we use the following notations. We define

$$Q^{+,N}(f^N) := \frac{1}{N^2} \sum_{i,j=1, i \neq j}^N Q_{ij}^{+,N}(f^N), \quad Q_{ij}^{+,N}(f^N) := \int_{\mathbb{S}^{d-1}} f_{ij}^N b(\cos \theta_{ij}) d\sigma,$$

$$\cos \theta_{ij} := \sigma \cdot k_{ij} \quad \text{with } k_{ij} = (v_i - v_j)/|v_i - v_j|$$

and where we assume that b is even and that

$$\int_{\mathbb{S}^{d-1}} b(\sigma \cdot k) d\sigma = C_B \quad \text{for any } k, |k| = 1.$$

For any function g^N on \mathbb{R}^{dN} shall use the shorthand notation g_{ij}^N to denote the function $V \mapsto g(r_{ij,\sigma}(V))$, which depends also implicitly on σ . We shall make use of the measure

preserving involution

$$\Theta_{ij} : \begin{cases} \mathbb{R}^n \times \mathbb{R}^n \times \mathbb{S}^{d-1} & \rightarrow \mathbb{R}^n \times \mathbb{R}^n \times \mathbb{S}^{d-1} \\ (v_i, v_j, \sigma) & \mapsto (v'_i, v'_j, \sigma') \end{cases}$$

where $\sigma' = (v_i - v_j)/|v_i - v_j| = k_{ij}$.

Finally as in [71], we shall use the following endomorphism of \mathbb{R}^d

$$\begin{aligned} M_{\sigma k}(x) &= (k \cdot \sigma)x - (k \cdot x)\sigma \\ P_{\sigma k}(x) &= (\sigma \cdot x)k + M_{\sigma k} \end{aligned}$$

and we recall that $\|P_{\sigma k}(x)\| \leq \|x\|$ with equality only if x, σ, k are coplanar.

We claim that it is enough to prove that

$$(7.39) \quad I(Q^{+,N}(f^N)) = I\left(\frac{1}{N^2} \sum_{i,j=1}^N \int_{\mathbb{S}^{d-1}} f(r_{ij,\sigma}(V)) b(\cos \theta_{ij}) d\sigma\right) \leq C_B I(f).$$

Indeed with this result at hand, we can write for $\varepsilon > 0$:

$$f_{t+\varepsilon}^N = e^{-NC_B\varepsilon} f_t^N + N C_B \int_0^\varepsilon e^{NC_B(s-\varepsilon)} Q^{+,N}(f_{t+s}^N) ds$$

and therefore from the convexity of I

$$I(f_{t+\varepsilon}^N) \leq e^{-NC_B\varepsilon} I(f_t^N) + (1 - e^{-NC_B\varepsilon}) I\left(\int_0^\varepsilon Q^{+,N}(f_{t+s}^N) \frac{N C_B e^{NC_B(s-\varepsilon)}}{(1 - e^{-NC_B\varepsilon})} ds\right).$$

Observe that

$$\int_0^\varepsilon \frac{N C_B e^{NC_B(s-\varepsilon)}}{(1 - e^{-NC_B\varepsilon})} ds = 1$$

and then we can use the convexity of I again to get

$$I(f_{t+\varepsilon}^N) \leq e^{-NC_B\varepsilon} I(f_t^N) + \int_0^\varepsilon I(Q^{+,N}(f_{t+s}^N)) N C_B e^{NC_B(s-\varepsilon)} ds.$$

Finally using the claimed result (7.39) we obtain

$$\frac{I(f_{t+\varepsilon}^N) - I(f_t^N)}{\varepsilon} \leq -\frac{(1 - e^{-NC_B\varepsilon})}{\varepsilon} I(f_t^N) + \frac{1}{\varepsilon} \int_0^\varepsilon I(f_{t+s}^N) N C_B^2 e^{NC_B(s-\varepsilon)} ds.$$

Then taking $\varepsilon \rightarrow 0$ and using Lebesgue's theorem we deduce

$$\frac{d}{dt} I(f_t^N) \leq -N C_B I(f_t^N) + N C_B I(f_t^N) \leq 0$$

which concludes the proof.

Let us now focus on the proof of the claim (7.39). Taking advantage of the convexity of I , it is enough to prove

$$\forall i \neq j \in [1, N], \quad I(Q_{ij}^{+,N}(f^N)) \leq C_B I(f^N).$$

Let us compute each partial derivative of $Q_{ij}^{+,N}(f^N)$. If $\ell \notin \{i, j\}$ then the derivative does not act on the kernel b and we obtain:

$$\begin{aligned} \nabla_{v_\ell} (Q_{ij}^{+,N}(f^N)) &= \int_{\mathbb{S}^{d-1}} \nabla_{v_\ell} (f_{ij}^N) b(\cos \theta_{ij}) d\sigma = \int_{\mathbb{S}^{d-1}} (\nabla_{v_\ell} f^N)_{ij} b(\cos \theta_{ij}) d\sigma \\ &= 2 \int_{\mathbb{S}^{d-1}} (\sqrt{f^N})_{ij} (\nabla_{v_\ell} \sqrt{f^N})_{ij} b(\cos \theta_{ij}) d\sigma. \end{aligned}$$

If $\ell \in \{i, j\}$, then it is slightly more complicated. Without restriction we perform calculations in the case $\ell = i$. Let us first prove the formula

$$(7.40) \quad \begin{aligned} & \nabla_{v_i} \left(Q_{ij}^{+,N} (f^N) \right) \\ &= \int_{\mathbb{S}^{d-1}} \left[(\nabla_{v_i} f^N)_{ij} + (\nabla_{v_j} f^N)_{ij} + P_{\sigma k} \left((\nabla_{v_i} f^N)_{ij} - (\nabla_{v_j} f^N)_{ij} \right) \right] b(\cos \theta_{ij}) \, d\sigma \end{aligned}$$

(the same equality obviously holds where i is replaced by j).

Simple computations (see for instance [71]) yield

$$\begin{aligned} \nabla_{v_i} (f_{ij}^N) &= \frac{1}{2} \left((\nabla_{v_i} f^N)_{ij} + (\nabla_{v_j} f^N)_{ij} \right) + \frac{1}{2} \left[\left((\nabla_{v_i} f^N)_{ij} - (\nabla_{v_j} f^N)_{ij} \right) \cdot \sigma \right] k_{ij}, \\ \nabla_{\sigma} (f_{ij}^N) &= \frac{|v_i - v_j|}{2} \left((\nabla_{v_i} f^N)_{ij} - (\nabla_{v_j} f^N)_{ij} \right), \\ \nabla_{v_i} [b(\cos \theta_{ij})] &= \frac{1}{|v_i - v_j|} b'(\sigma \cdot k_{ij}) \Pi_{k^\perp} \sigma, \end{aligned}$$

where Π_{k^\perp} is the projection on the hyperplane k^\perp . Using the first and third equality above, we get

$$(7.41) \quad \begin{aligned} & \nabla_{v_i} \left(Q_{ij}^{+,N} (f^N) \right) \\ &= \frac{1}{2} \int_{\mathbb{S}^{d-1}} b(\cos \theta_{ij}) \left((\nabla_{v_i} f)_{ij} + (\nabla_{v_j} f)_{ij} + \left[\left((\nabla_{v_i} f)_{ij} - (\nabla_{v_j} f)_{ij} \right) \cdot \sigma \right] k \right) d\sigma \\ & \quad + \left(\int_{\mathbb{S}^{d-1}} b'(\cos \theta_{ij}) \frac{f_{ij}}{|v_i - v_j|} \Pi_{k^\perp} \sigma \, d\sigma \right) \end{aligned}$$

and we use the following formula of integration by part on the sphere \mathbb{S}^{d-1} (see [71, Lemma 2])

$$\int_{\mathbb{S}^{d-1}} b'(\cos \theta_{ij}) F(\sigma) \Pi_{k^\perp} \sigma \, d\sigma = \int_{\mathbb{S}^{d-1}} b(\cos \theta_{ij}) M_{\sigma k} (\nabla_{\sigma} F(\sigma)) \, d\sigma$$

and the second equality above to rewrite the term involving b' in (7.41) into

$$\frac{1}{2} \int_{\mathbb{S}^{d-1}} b(\cos \theta_{ij}) M_{\sigma k} \left((\nabla_{v_i} f)_{ij} - (\nabla_{v_j} f)_{ij} \right) \, d\sigma$$

Putting all together, we get formula (7.40).

We deduce that for $\ell \neq i, j$ we have by Cauchy-Schwarz

$$\begin{aligned} & \left| \nabla_{v_\ell} \left(Q_{ij}^{+,N} (f^N) \right) \right|^2 \\ & \leq 4 \left(\int_{\mathbb{S}^{d-1}} f_{ij}^N b(\cos \theta_{ji}) \, d\sigma \right) \left(\int_{\mathbb{S}^{d-1}} \left| \left(\nabla_{v_\ell} \sqrt{f^N} \right)_{ij} \right|^2 b(\cos \theta_{ij}) \, d\sigma \right) \end{aligned}$$

and therefore

$$\frac{\left| \nabla_{v_\ell} \left(Q_{ij}^{+,N} (f^N) \right) \right|^2}{Q_{ij}^{+,N} (f^N)} \leq 4 \int_{\mathbb{S}^{d-1}} \left| \left(\nabla_{v_\ell} \sqrt{f^N} \right)_{ij} \right|^2 b(\cos \theta_{ij}) \, d\sigma.$$

Now integrating in V we obtain

$$\begin{aligned} I_\ell \left(Q_{ij}^{+,N} (f^N) \right) &\leq 4 \int_{\mathbb{R}^{dN}} \int_{\mathbb{S}^{d-1}} \left| \left(\nabla_{v_\ell} \sqrt{f^N} \right)_{ij} \right|^2 b(\cos \theta_{ij}) \, d\sigma \, dV \\ &\leq 4 \int_{\mathbb{R}^{dN}} \int_{\mathbb{S}^{d-1}} \left| \nabla_{v_\ell} \sqrt{f^N} \right|^2 b(\cos \theta_{ij}) \, d\sigma \, dV \\ &\leq 4 C_B \int_{\mathbb{R}^{dN}} \left| \nabla_{v_\ell} \sqrt{f^N} \right|^2 \, dV =: I_\ell (f^N) \end{aligned}$$

where we have used Cauchy-Schwartz's inequality and the change of variable Θ_{ij} , and where I_ℓ is defined from the last line (Fisher information restricted to the ℓ -th derivative).

When $\ell = i, j$, we use (7.40) to get

$$\begin{aligned} \left| \nabla_{v_\ell} \left(Q_{ij}^{+,N} (f^N) \right) \right|^2 &\leq \left(\int_{\mathbb{S}^{d-1}} f^N b(\cos \theta_{ji}) \, d\sigma \right) \\ &\quad \times \left(\int_{\mathbb{S}^{d-1}} \left| \left(\nabla_{v_i} \sqrt{f^N} \right) + \left(\nabla_{v_j} \sqrt{f^N} \right) \right. \right. \\ &\quad \left. \left. + P_{\sigma k} \left(\left(\nabla_{v_j} \sqrt{f^N} \right) - \left(\nabla_{v_i} \sqrt{f^N} \right) \right) \right|^2 b(\cos \theta_{ij}) \, d\sigma \right) \end{aligned}$$

where we have used Cauchy-Schwartz's inequality and the change of variable Θ_{ij}

Since for fixed V , $P_{\sigma k}$ is odd in σ and $b(\cos \theta_{ij})$ is even in σ , we have

$$\int_{\mathbb{S}^{d-1}} A \cdot P_{\sigma k}(B) \, d\sigma = 0$$

for any functions A, B independent of σ . Using finally that $P_{\sigma k}$ has norm less than 1 (for the subordinated norm to the euclidean norm) we get

$$\begin{aligned} \left| \nabla_{v_\ell} \left(Q_{ij}^{+,N} (f^N) \right) \right|^2 &\leq 2 \left(\int_{\mathbb{S}^{d-1}} f^N b(\cos \theta_{ji}) \, d\sigma \right) \\ &\quad \times \left(\int_{\mathbb{S}^{d-1}} \left| \nabla_{v_i} \sqrt{f^N} \right|^2 + \left| \nabla_{v_j} \sqrt{f^N} \right|^2 b(\cos \theta_{ij}) \, d\sigma \right) \end{aligned}$$

and therefore

$$I_\ell \left(Q_{ij}^{+,N} (f^N) \right) \leq C_B \frac{I_i (f^N) + I_j (f^N)}{2}.$$

Finally we end up with

$$I \left(Q_{ij}^{+,N} (f^N) \right) = C_B \sum_{\ell=1}^N I_\ell \left(Q_{ij}^{+,N} (f^N) \right) \leq C_B \sum_{\ell=1}^N I_\ell (f^N) = C_B I (f^N)$$

which concludes the proof. \square

Let us now conclude the proof of point (ii) in Theorem 7.10. We make use of the so-called ‘‘HWT’’ interpolation inequality on the manifold \mathcal{S}^N . Observe that \mathcal{S}^N has positive Ricci curvature since it has positive curvature. Then [73, Theorem 30.21] implies that

$$\frac{1}{N} H (f^N | \gamma^N) \leq \frac{W_2 (f^N, \gamma^N)}{\sqrt{N}} \sqrt{\frac{I (f^N | \gamma^N)}{N}}.$$

We can then use the uniform bound on the Fisher information provided by Lemma 7.13 (vi) and the bound on the initial data to get:

$$\frac{I (f^N | \gamma^N)}{N} \leq \frac{I (f_0^N | \gamma^N)}{N} \leq C$$

for some constant $C > 0$ independent of N . Moreover Lemma 4.1 and the propagation of moments on the N -particle system in Lemma 5.4 imply that

$$\frac{W_2(f^N, \gamma^N)}{\sqrt{N}} \leq C \left(\frac{W_1(f^N, \gamma^N)}{N} \right)^\alpha$$

for some constant $C > 0$ and exponent $\alpha > 0$ independent of N . Then using Theorem 5.3 (case (b)) we deduce that

$$\frac{W_2(f^N, \gamma^N)}{\sqrt{N}} \leq C \left(\frac{W_1(f^N, \gamma^N)}{N} \right)^\alpha \leq \beta(t)$$

with $\beta(t) \rightarrow 0$ as $t \rightarrow +\infty$, which implies that

$$\frac{1}{N} H(f^N | \gamma^N) \leq C \beta(t)$$

and concludes the proof.

8. THE BBGKY APPROACH REVISITED

The so-called BBGKY (Bogoliubov, Born, Green, Kirkwood and Yvon) and statistical solutions approach is very popular in physics and mathematics for studying many-particle systems: see for instance [3] where this approach is used for Kac's master equation for hard spheres, or see, among many other works, the recent impressive series of papers [4, 26, 27, 28] where this approach is used for the derivation of nonlinear Schrödinger equations in mean-field theory in quantum physics. The basic ideas underlying the BBGKY approach to mean-field could be summarized as:

- Write a BBGKY hierarchy on marginals of the N -particle system and prove that the N -particle system solutions converge to the solutions of an "infinite hierarchy" when N goes to infinity.
- Then prove that solutions to this infinite hierarchy are unique, which is the hardest part of this program.
- Then deduce the propagation of chaos by exhibiting, for any chaotic initial data to the infinite hierarchy, a natural solution obtained by the infinite tensorization of the 1-particle solution to the limiting nonlinear mean-field equation.

This section revisits the so-called BBGKY hierarchy and statistical solutions approach in the case of collision processes, by showing (1) how these notions are included in our functional framework (cf. the abstract semigroup T_t^∞ defined in Section 2), and (2) how to give a proof of uniqueness and propagation of chaos based on them by using the functional tools we have introduced.

8.1. The BBGKY hierarchy. Let us recall the master of the N -particle system undergoing a Boltzmann collision process (the notion of BBGKY hierarchy has wider application range, but we shall stick to this concrete case for clarity), see (5.3)-(5.4):

$$(8.1) \quad \partial_t \langle f_t^N, \varphi \rangle = \langle f_t^N, G^N \varphi \rangle$$

with

$$(G^N \varphi)(V) = \frac{1}{N} \sum_{1 \leq i < j \leq N} \Gamma(|v_i - v_j|) \int_{\mathbb{S}^{d-1}} b(\cos \theta_{ij}) [\varphi_{ij}^* - \varphi] d\sigma$$

$$\text{where } \varphi_{ij}^* = \varphi(V_{ij}^*) \text{ and } \varphi = \varphi(V) \in C_b(\mathbb{R}^{Nd}).$$

Then the BBGKY hierarchy writes as follows. Let us recall the notation

$$f_\ell^N = \Pi_\ell[f^N] = \int_{v_{\ell+1}, \dots, v_N} df^N(v_{\ell+1}, \dots, v_N)$$

for the marginals. Then integrating the master equation (8.1) against some test function $\varphi = \varphi(v_1, \dots, v_\ell)$ depending only on the first ℓ variables leads to

$$\frac{d}{dt} \langle f_\ell^N, \varphi \rangle = \langle f_{\ell+1}^N, G_{\ell+1}^N(\varphi) \rangle$$

where

$$G_{\ell+1}^N(\varphi) := \frac{1}{N} \sum_{1 \leq i \leq \ell, 1 \leq j \leq N, i \neq j} \Gamma(|v_i - v_j|) \int_{\mathbb{S}^{d-1}} b(\cos \theta_{ij}) [\varphi_{ij}^* - \varphi] d\sigma$$

Then denoting by

$$\mathcal{Z}_{ij}^N := \left\langle f^N, \Gamma(|v_i - v_j|) \int_{\mathbb{S}^{d-1}} b(\cos \theta_{ij}) [\varphi_{ij}^* - \varphi] d\sigma \right\rangle$$

we can further decompose this sum as

$$\frac{d}{dt} \langle f_\ell^N, \varphi \rangle = \frac{1}{N} \sum_{i,j \leq \ell} \mathcal{Z}_{ij}^N + \frac{1}{N} \sum_{i \leq \ell < j} \mathcal{Z}_{ij}^N = \frac{1}{N} \sum_{i,j \leq \ell} \mathcal{Z}_{ij}^N + \mathcal{O}\left(\frac{\ell^2}{N}\right)$$

by observing that $\mathcal{Z}_{ij}^N = 0$ for $i, j > \ell$, and using the symmetry of f^N we finally deduce

$$(8.2) \quad \frac{d}{dt} \langle f_\ell^N, \varphi \rangle = \frac{(N-\ell)}{N} \left(\sum_{i=1}^{\ell} \mathcal{Z}_{i(\ell+1)}^N \right) + \mathcal{O}\left(\frac{\ell^2}{N}\right).$$

We thus end with a series of N coupled equations on the marginals f_ℓ^N , where the ℓ -equation ($\ell \leq N-1$) depends on the $f_{\ell+1}^N$ marginal.

8.2. The infinite hierarchy and statistical solutions. Assume now that

$$\forall \ell \geq 1, \quad f_\ell^N \rightharpoonup \pi_\ell \text{ in } P(\mathbb{R}^{d\ell}).$$

Starting from (8.2) we obtain, for $\varphi = \varphi(v_1, \dots, v_\ell)$ depending only on the first ℓ variables,

$$\frac{d}{dt} \langle \pi_\ell, \varphi \rangle = \langle \pi_{\ell+1}, G_{\ell+1}^\infty(\varphi) \rangle$$

where $G_{\ell+1}^\infty(\varphi) \in C_b(\mathbb{R}^{d(\ell+1)})$ is defined by

$$G_{\ell+1}^\infty(\varphi) := \Gamma(|v_i - v_j|) \int_{\mathbb{S}^{d-1}} b(\cos \theta_{ij}) [\varphi_{ij}^* - \varphi] d\sigma.$$

In a more compact form, we have the following set of linear coupled evolution equations

$$(8.3) \quad \forall \ell \geq 1, \quad \partial_t \pi_\ell = A_{\ell+1}^\infty(\pi_{\ell+1}) \text{ with } A_{\ell+1}^\infty := (G_{\ell+1}^\infty)^*.$$

Since the family of ℓ -particle probabilities π_ℓ is symmetric and compatible in the sense that

$$\forall \ell \geq 1, \quad \Pi_\ell[\pi_{\ell+1}] = \pi_\ell$$

(this follows from the construction), we can associate by Hewitt-Savage's Theorem [38] a unique $\pi \in P(P(\mathbb{R}^d))$ and this set of evolution equation translates into an evolution equation

$$\partial_t \pi = A^\infty(\pi) \text{ on } P(P(\mathbb{R}^d))$$

of *statistical solutions* and the corresponding dual evolution

$$\partial_t \Phi = \bar{G}^\infty \Phi \text{ on } C_b(P(\mathbb{R}^d)).$$

In order to make this heuristic rigorous at an abstract level, one needs at least some tightness on the sequence $(f_\ell^N)_{N \geq \ell}$ for any ℓ , and some convergence

$$G_{\ell+1}^N(\varphi) \rightarrow G_{\ell+1}^\infty(\varphi)$$

on compact subset of $\mathbb{R}^{d(\ell+1)}$. Both are satisfied for Boltzmann collision processes considered in this paper (tightness follows from the moment estimates in Lemma 5.4 for instance).

8.3. Uniqueness of statistical solutions and chaos. We now want, under appropriate abstract assumptions, to identify the limit evolution in $P(P(\mathbb{R}^d))$ or $C_b(P(\mathbb{R}^d))$ obtained from the hierarchy, and show that it coincides with the the pushforward evolution semigroup T_t^∞ introduced in Subsection 2.3. Meanwhile we shall prove that the statistical solutions to the infinite hierarchy are unique, and hence prove the propagation of chaos, without any rate, but also under weaker assumptions than previously.

We shall make the following assumptions

(A1') Assumptions on the N -particle system.

G^N and T_t^N are well defined on $C_b(E^N)$ and invariant under permutation, and the associated solutions f_t^N satisfies:

$$\forall \ell \geq 1, \quad \text{the sequence } (f_\ell^N)_{N \geq \ell} \text{ is tight in } P_{\mathcal{G}_1}(E)^{\otimes \ell}$$

where \mathcal{G}_1 is a Banach space and $P_{\mathcal{G}_1}(E)$ is defined in Definitions 2.4-2.5), and is associated to a weight function $m_{\mathcal{G}_1}$ and a constraint function $\mathbf{m}_{\mathcal{G}_1}$, and endowed with the metric induced from \mathcal{G}_1 .

(A2') Assumptions for the existence of the statistical and pushforward semigroups.

For some $\delta \in (0, 1]$ and some $\bar{a} \in (0, \infty)$ we assume that for any $a \in (\bar{a}, \infty)$:

(i) The equation (2.1) generates a semigroup

$$S_t^{NL} : \mathcal{B}P_{\mathcal{G}_1, a} \rightarrow \mathcal{B}P_{\mathcal{G}_1, a}$$

which is δ -Hölder continuous locally uniformly in time, in the sense that for any $\tau \in (0, \infty)$ there exists $C_\tau \in (0, \infty)$ such that

$$\forall f, g \in \mathcal{B}P_{\mathcal{G}_1, a}, \quad \sup_{t \in [0, \tau]} \|S_t^{NL} f - S_t^{NL} g\|_{\mathcal{G}_1} \leq C_\tau \|f - g\|_{\mathcal{G}_1}^\delta.$$

(ii) The application Q is bounded and δ -Hölder continuous from $\mathcal{B}P_{\mathcal{G}_1, a}$ into \mathcal{G}_1 .

(iii) E is a locally compact Polish space and there is \mathcal{F}_1 in duality with \mathcal{G}_1 such that \mathcal{F}_1 is dense in $C_b(E)$ in the sense of uniform convergence on any compact set.

(A3') Convergence of the generators.

For any fixed $\ell \in \mathbb{N}^*$ and any $\varphi \in C_b(E^\ell)$, the sequence

$$G_{\ell+1}^N(\varphi) \in C_b(E^{\ell+1}) \quad \text{satisfies} \quad G_{\ell+1}^N \varphi \xrightarrow{N \rightarrow \infty} G_{\ell+1}^\infty \varphi$$

uniformly on compact sets, where $G_{\ell+1}^\infty \varphi$ satisfies the following *compatibility binary derivation structure*:

for any $\varphi = \varphi_1 \otimes \dots \otimes \varphi_\ell \in C_b(E)^{\otimes \ell}$ and any $V = (v_1, \dots, v_{\ell+1}) \in E^{\ell+1}$

$$(8.4) \quad G_{\ell+1}^\infty(\varphi)(V) = \sum_{i=1}^{\ell} \left(\prod_{j \neq i} \varphi_j(v_j) \right) Q^*(\varphi_i)(v_i, v_{\ell+1})$$

where Q^* is defined from Q through the duality relation

$$\forall f \in P_{\mathcal{G}_1}, \quad \forall \psi \in C_b(E), \quad \langle Q(f, f), \psi \rangle = \langle f \otimes f, Q^*(\psi) \rangle.$$

Remark 8.1. The identity (8.4) is called *compatibility binary derivation structure* for the following reasons: *compatibility* since it is a natural condition in order that any solution f_t to the nonlinear Boltzmann provides a tensorized solution to the BBGKY hierarchy (8.3). Indeed, considering such a solution

$$f_t \in C(\mathbb{R}_+; P_{\mathcal{G}_1}(E)) \quad \text{and} \quad \varphi = \varphi_1 \otimes \dots \otimes \varphi_\ell \in C_b(E)^{\otimes \ell}$$

we compute

$$\begin{aligned} \frac{d}{dt} \langle f_t^{\otimes \ell}, \varphi \rangle &= \sum_{i=1}^{\ell} \left(\prod_{j \neq i} \langle f_t, \varphi_j \rangle \right) \frac{d}{dt} \langle f_t, \varphi_i \rangle = \sum_{i=1}^{\ell} \left(\prod_{j \neq i} \langle f_t, \varphi_j \rangle \right) \langle Q(f_t, f_t), \varphi_i \rangle \\ &= \sum_{i=1}^{\ell} \left(\prod_{j \neq i} \langle f_t, \varphi_j \rangle \right) \langle f_t \otimes f_t, Q^*(\varphi_i) \rangle = \langle f_t^{\otimes \ell+1}, G_{\ell+1}^\infty \varphi \rangle. \end{aligned}$$

The word *binary* refers to the fact that $G_{\ell+1}^\infty$ decomposes in function acting on *one* variable and adding *one* variable, which corresponds to the binary nature of the collisions. Finally the word *derivation* refers to the fact that the following distributivity property holds

$$G_{\ell+1}^\infty(\varphi \otimes \psi) = G_{\ell+1}^\infty(\varphi) \otimes \psi + \varphi \otimes G_{\ell+1}^\infty(\psi).$$

Let us mention that this distributivity property is at the basis of the original combinatorial proof of Kac [42] of propagation of chaos for the simplified Boltzmann-Kac equation.

(A4') Differential stability of the limiting semigroup.

We consider some Banach space $\mathcal{G}_2 \supset \mathcal{G}_1$ (where \mathcal{G}_1 was defined in (A2)) and the corresponding probability space $P_{\mathcal{G}_2}(E)$ (see Definitions 2.4-2.5) with the weight function $m_{\mathcal{G}_2}$ and the constraint function $\mathbf{m}_{\mathcal{G}_2}$, and endowed with the metric induced from \mathcal{G}_2 .

We assume that the flow S_t^{NL} is $UC_{\Lambda_2}^1(P_{\mathcal{G}_1, r}, P_{\mathcal{G}_2})$ for any $r \in \mathbf{R}_{\mathcal{G}_1}$ in the sense that there exists $C_T^\infty > 0$ such that

$$\sup_{r \in \mathbf{R}_{\mathcal{G}_1}} \int_0^T \left([S_t^{NL}]_{C_{\Lambda_2}^{1,0}(P_{\mathcal{G}_1, r}, P_{\mathcal{G}_2})} + [S_t^{NL}]_{C_{\Lambda_2}^{0,1}(P_{\mathcal{G}_1, r}, P_{\mathcal{G}_2})} \right) dt \leq C_T^\infty.$$

Thanks to (A2'), we know from Lemma 2.13 that for any $\Phi \in UC_b(P_{\mathcal{G}_1}, \mathbb{R})$ we may define the C_0 -semigroup $T_t^\infty[\Phi] \in UC_b(P_{\mathcal{G}_1})$ by

$$T_t^\infty[\Phi](f) = \Phi(S_t^{NL} f),$$

and so that $\Phi_t = T_t^\infty[\Phi]$ satisfies the equation

$$\partial_t \Phi = G^\infty[\Phi]$$

with a generator G^∞ which is a closed operator on $UC_b(P_{\mathcal{G}_1})$ and has domain $\text{Dom}(G^\infty)$ which contains $C^1(P_{\mathcal{G}_1})$, and is defined by

$$G^\infty[\Phi](f) = \langle D\Phi(f), Q(f, f) \rangle_{\mathcal{G}'_1, \mathcal{G}_1} = \langle Q(f, f), D\Phi(f) \rangle_{P_{\mathcal{G}_1}, C_b(P_{\mathcal{G}_1})}.$$

The evolution corresponds to the following dual evolution equation

$$(8.5) \quad \frac{d}{dt} \langle \pi_t, \Phi \rangle = \langle \pi_t, G^\infty[\Phi] \rangle.$$

Our goal is to prove (1) that the evolution equations (8.3) and (8.5) (or in other words that $\bar{G}^\infty = G^\infty$, and (2) most importantly that the solution to these equation is given by the characteristics equation

$$(8.6) \quad \forall \Phi \in UC_b(P_{\mathcal{G}_1}; \mathbb{R}) \quad \langle \pi_t, \Phi \rangle = \langle \pi_0, T_t^\infty \Phi \rangle.$$

(or in other words that the heuristic generator \bar{G}^∞ introduced for the hierarchy is well-defined and equal to G^∞).

Let us explain why the relation (8.6) indeed defines uniquely a probability evolution $\bar{\pi}_t \in P(P_{\mathcal{G}_1})$. For any $\ell \in \mathbb{N}^*$ we define

$$\varphi \in \mathcal{F}^{\otimes \ell} \mapsto \langle \pi_t^\ell, \varphi \rangle := \langle \pi_0, T_t^\infty R_\varphi^\ell \rangle.$$

That is a positive linear form on $\mathcal{F}^{\otimes \ell}$. Thanks to **(A2')**-**(iii)**, the Stone-Weierstrass density theorem and the Markov-Riesz representation theorem, we conclude that π_t^ℓ is well defined as a element of $P_{\mathcal{G}_1}(E)^{\otimes \ell}$. Since now the sequence (π_t^ℓ) is symmetric and compatible, the Hewitt-Savage representation theorem implies that there exists a unique probability measure $\bar{\pi}_t \in P(P_{\mathcal{G}_1})$ such that for any $\varphi \in \mathcal{F}_1^{\otimes \ell}$

$$(8.7) \quad \langle \bar{\pi}_t, R_\varphi^\ell \rangle := \langle \pi_0, T_t^\infty R_\varphi^\ell \rangle.$$

Theorem 8.2. *Under the assumptions **(A1')**-**(A2')**-**(A3')**-**(A4')**, for any initial datum $\pi_0 \in P(P_{\mathcal{G}_1})$, the flow $\bar{\pi}_t$ defined from (8.7) is the unique solution in $C([0, \infty); P(P_{\mathcal{G}_1}))$ to the infinite hierarchy evolution (8.3) starting from π_0 .*

Moreover, if π_0 is f_0 -chaotic (that is if $\pi_0 = \delta_{f_0}$ with $f_0 \in P(E)$), then π_t is $S_t^{NL} f_0$ -chaotic for any $t \geq 0$. As a consequence we deduce that if f_0^N is f_0 -chaotic, then f_t^N is $S_t^{NL} f_0$ -chaotic. More generally, if f_0^N converges to π_0 then f_t^N converge to $\bar{\pi}_t$ the associated statistical solution.

Proof of Theorem 8.2. We shall proceed in several steps.

Step 1: Propagation of chaos. Let us recall that Hewitt-Savage's theorem [38] implies that for any $\pi \in P(P(E))$ there exists a unique sequence $(\pi^\ell) \in P(E^\ell)$ such that

$$\forall \varphi \in (C_b(E))^{\otimes \ell}, \quad \langle \pi^\ell, \varphi \rangle = \langle \pi, R_\varphi^\ell \rangle.$$

As a consequence, if π_0 is f_0 -chaotic and $\bar{\pi}$ satisfies (8.7), then

$$\begin{aligned} \langle \bar{\pi}_{t,\ell}, \varphi \rangle &= \langle \bar{\pi}_t, R_\varphi^\ell \rangle = \langle \pi_0, T_t^\infty R_\varphi^\ell \rangle = T_t^\infty [R_\varphi^\ell] (f_0) \\ &= R_\varphi^\ell (S_t^{NL} f_0) = \langle S_t^{NL} f_0, \varphi_1 \rangle \dots \langle S_t^{NL} f_0, \varphi_\ell \rangle, \end{aligned}$$

which means that $\bar{\pi}_{t,\ell} = f_t^{\otimes \ell}$, or equivalently $\bar{\pi}_t = \delta_{f_t}$, and the solution $\bar{\pi}_t$ defined by (8.7) is f_t -chaotic.

Step 2: Equivalence between (8.3) and (8.5).

First let us assume (8.5) and prove (8.3). Consider $f \in P_{\mathcal{G}_1}(E)$ and $\varphi \in \mathcal{F}^{\otimes \ell}$. Then we have $R_\varphi \in C^{1,1}(P_{\mathcal{G}_1}(E))$ and we deduce from (8.4) that

$$\langle f^{\otimes \ell+1}, G_{\ell+1}^\infty \varphi \rangle = \langle Q(f, f), DR_\varphi^\ell(f) \rangle = G^\infty [R_\varphi^\ell] (f)$$

which means

$$R_{G_{\ell+1}^\infty \varphi}^{\ell+1} = G^\infty [R_\varphi^\ell].$$

Then, using Hewitt-Savage's Theorem again, (8.5) implies that

$$(8.8) \quad \frac{d}{dt} \langle \pi_{\ell,t}, \varphi \rangle = \frac{d}{dt} \langle \pi_t, R_\varphi^\ell \rangle = \langle \pi_t, G^\infty [R_\varphi^\ell] \rangle = \langle \pi_t, R_{G_{\ell+1}^\infty [\varphi]}^{\ell+1} \rangle = \langle \pi_{\ell+1,t}, G_{\ell+1}^\infty [\varphi] \rangle$$

which means that π_t satisfies (8.5).

Assume conversely that π_t satisfies (8.3) and let us prove (8.5). One needs to prove that one can recover any $\Phi \in UC^1(P_{\mathcal{G}_1}(E))$ from the previous equation (8.8).

Therefore consider $\Phi \in UC^1(P_{\mathcal{G}_1}(E))$ and let us define

$$\varphi = \left(\pi_C^\ell \Phi \right) (V) = \Phi \left(\mu_V^\ell \right), \quad V = (v_1, \dots, v_\ell)$$

and let us write (8.8) for this choice of φ :

$$\frac{d}{dt} \langle \pi_t, R_{\pi_C^\ell}^\ell \Phi \rangle = \langle \pi_t, G^\infty [R_{\pi_C^\ell}^\ell \Phi] \rangle = \langle \pi_t, R_{G_{\ell+1}^\infty[\pi_C^\ell \Phi]}^{\ell+1} \rangle.$$

Then, on the one hand, for any $f \in P_{\mathcal{G}_1}(E)$

$$R_{\pi_C^\ell}^\ell(f) = \int_{E^\ell} \Phi(\mu_V^\ell) df^{\otimes \ell}(V) \xrightarrow{\ell \rightarrow \infty} \Phi(f)$$

by the law of large numbers.

On the other hand, for any $f \in P_{\mathcal{G}_1}(E)$, we have

$$\begin{aligned} R_{G_{\ell+1}^\infty[\pi_C^\ell \Phi]}(f) &= \langle f^{\otimes \ell+1}, G_{\ell+1}^\infty(\pi_C^\ell \Phi) \rangle \\ &= \langle DR_{\pi_C^\ell}^\ell(f), Q(f, f) \rangle \\ &= \sum_{i=1}^{\ell} \int_{E^\ell} \Phi(\mu_V^\ell) dQ(f, f)(v_i) \prod_{j \neq i} df(v_j). \end{aligned}$$

For any given $i = 1, \dots, \ell$, we define

$$\phi_{V_i}^{\ell-1} = D\Phi(\mu_{V_i}^{\ell-1}), \quad V_i := (v_1, \dots, v_{i-1}, v_{i+1}, \dots, v_\ell)$$

and we write

$$\Phi(\mu_V^\ell) = \Phi(\mu_{V_i}^{\ell-1}) + \langle \phi_{V_i}^{\ell-1}, \mu_V^\ell - \mu_{V_i}^{\ell-1} \rangle + \mathcal{O}\left(\Omega\left(\|\mu_{V_i}^{\ell-1} - \mu_V^\ell\|\right)\right).$$

Observing that

$$\mu_V^\ell - \mu_{V_i}^{\ell-1} = \frac{1}{N} \delta_{v_i} - \sum_{j \neq i} \frac{1}{\ell(\ell-1)} \delta_{v_j}$$

and that $\langle Q(f, f), 1 \rangle = 0$ from assumption **(A2')**-**(ii)**, we find

$$\begin{aligned} R_{G_{\ell+1}^\infty(\pi_C^\ell \Phi)}(f) &= \sum_{i=1}^{\ell} \int_{E^\ell} \left(\frac{1}{\ell} \phi_{V_i}^{\ell-1}(v_i) + \mathcal{O}(\Omega(\ell^{-1})) \right) dQ(f, f)(v_i) \prod_{j \neq i} df(v_j) \\ &= \sum_{i=1}^{\ell-1} \int_{E^{\ell-1}} \frac{1}{\ell-1} \langle Q(f, f), \phi_{V_i}^{\ell-1} \rangle df^{\otimes(\ell-1)}(V) + \mathcal{O}(\ell \Omega(\ell^{-1})) \\ &= \int_{E^{\ell-1}} \langle Q(f, f), D\Phi(\mu_V^{\ell-1}) \rangle df^{\otimes(\ell-1)}(V) + \mathcal{O}(\ell \Omega(\ell^{-1})) \\ &\xrightarrow{\ell \rightarrow \infty} \langle D\Phi(f), Q(f, f) \rangle \end{aligned}$$

by the law of large number again. This implies (8.5). We then conclude that (8.5) holds for any $\Phi \in C^1(P_{\mathcal{G}_1}(E))$ by density of $UC^1(P_{\mathcal{G}_1}(E))$ in this space and the fact that the domain of G^∞ contains $C^1(P_{\mathcal{G}_1}(E))$.

Step 3: Uniqueness. Let us prove that any solution of (8.3)-(8.5) satisfies the characteristics equation (8.7), or in other words that $\pi_t = \bar{\pi}_t$. This shall imply uniqueness since we have already seen that the $\bar{\pi}_t$ satisfying (8.7) is unique.

The fundamental point here is that for any $\Phi \in UC^1(P_{\mathcal{G}_1}(E))$ if we define $\Phi_t := T_t^\infty \Phi$, thanks to Lemma 2.13 we have

$$\forall t \geq 0, \quad \Phi_t \in UC^1(P_{\mathcal{G}_1}(E)) \subset \text{Dom}(G^\infty).$$

Then since

$$\tau \in [0, t] \mapsto \langle \pi_\tau, \Phi_{t-\tau} \rangle$$

is C^1 from the fact that $\Phi_{t-\tau} \in C^1(P_{\mathcal{G}_1}(E))$ belongs to the domain of G^∞ for any τ , we compute

$$\frac{d}{d\tau} \langle \pi_\tau, \Phi_{t-\tau} \rangle = \langle \pi_\tau, G^\infty [\Phi_{t-\tau}] \rangle - \langle \pi_\tau, G^\infty [\Phi_{t-\tau}] \rangle = 0$$

and we deduce that

$$\langle \pi_t, \Phi_0 \rangle = \langle \pi_0, \Phi_t \rangle$$

which proves that $\pi_t = \bar{\pi}_t$ satisfies (8.7), and concludes the proof. \square

8.4. A remark on stationary statistical solutions. As we have seen:

- The chaoticity of a sequence of symmetric N -particle distributions $f^N \in P(E^N)$, $N \geq 1$ is equivalent to the fact that the associated $\pi \in P(P(E))$ is a Dirac at some $f_0 \in P(E)$: $\pi = \delta_{f_0}$. Hence, in view of Hewitt-Savage's theorem, non-chaoticity can be reframed as saying that π is a superposition of *several*, instead of one, chaotic states.
- We have recalled the result in [42, 12] stating that a chaotic (tensorized) sequence is asymptotically concentrated on the energy sphere, which is an effect of the Central Limit Theorem.
- Finally let us make the simple observation that the N -particle dynamics leaves the energy spheres invariant and relaxes on each energy spheres to the uniform measure. This is a consequence of the energy conservation laws: at the level of the particle system, the dynamics is layered according to the value of this conservation law.

One deduces from these considerations that there is room for *non-chaotic stationary states* of the N -particle system, namely superposition of *several* stationary states on different energy spheres. Let us make this more precise.

Lemma 8.3. *There exists a non-chaotic stationary solutions to the statistical Boltzmann equation. In other words, there exists $\pi \in P(P(\mathbb{R}^d))$ such that $\pi \neq \delta_p$ for some $p \in P(\mathbb{R}^d)$ and $A_{\ell+1}^\infty(\pi_{\ell+1}) = 0$ for any $\ell \in \mathbb{N}$.*

Proof of Lemma 8.3. It is clear that any function on the form

$$V \in \mathbb{R}^{d(\ell+1)} \mapsto \pi_{\ell+1}(V) = \phi(|V|^2)$$

is a stationary solution for the equation (8.3), that is $A_{\ell+1}(\pi_{\ell+1}) = 0$ for any $\ell \geq 1$. Now we define, with $d = 1$ for the sake of simplicity, the sequence

$$V \in \mathbb{R}^\ell \mapsto \pi_\ell(V) = \frac{c_\ell}{(1 + |V|^2)^{m+\ell/2}} \quad \forall \ell \geq 1,$$

where the sequence of positive constants c_ℓ is inductively constructed in the following way.

- First c_1 is chosen in a unique way so that π_1 is a probability measure.
- Then, once c_1, \dots, c_ℓ are constructed, $c_{\ell+1}$ is constructed so that $\Pi_\ell[\pi_{\ell+1}] = \pi_\ell$, which means

$$\forall V \in \mathbb{R}^\ell, \quad \int_{v_* \in \mathbb{R}} \frac{c_{\ell+1}}{(1 + |V|^2 + |v_*|^2)^{m+\ell/2+1/2}} dv_* = \frac{c_\ell}{(1 + |V|^2)^{m+\ell/2}}.$$

This is always possible since

$$\begin{aligned} & \int_{v_* \in \mathbb{R}} \frac{c_{\ell+1}}{(1 + |V|^2 + |v_*|^2)^{m+\ell/2+1/2}} dv_* \\ &= \frac{c_{\ell+1}}{(1 + |V|^2)^{m+\ell/2+1/2}} \int_{v_* \in \mathbb{R}} \frac{1}{\left(1 + \frac{|v_*|^2}{(1+|V|^2)}\right)^{m+\ell/2+1/2}} dv_* \\ &= \frac{c_{\ell+1}}{(1 + |V|^2)^{m+\ell/2}} \int_{v_* \in \mathbb{R}} \frac{1}{(1 + |v_*|^2)^{m+\ell/2+1/2}} dv_* \end{aligned}$$

which concludes the induction.

We then deduce that the sequence π_ℓ , $\ell \geq 1$, satisfies (8.3) since every terms only depends on the energy, and also satisfies the compatibility condition $\Pi_\ell[\pi_{\ell+1}] = \pi_\ell$. This concludes the proof. \square

REFERENCES

- [1] ALONSO, R., CAÑIZO, J. A., GAMBA, I., AND MOUHOT, C. Creation of exponential moments in the Boltzmann equation. Work in progress.
- [2] AMBROSIO, L., GIGLIAND, N., AND SAVARE, G. *Gradient flows in metric spaces and in the space of Probability measures*, vol. 2005 of *Lectures in Mathematics ETH Zurich*. Birkhauser, 2005.
- [3] ARKERYD, L., CAPRINO, S., AND IANIRO, N. The homogeneous Boltzmann hierarchy and statistical solutions to the homogeneous Boltzmann equation. *J. Statist. Phys.* 63, 1-2 (1991), 345–361.
- [4] BARDOS, C., GOLSE, F., AND MAUSER, N. J. Weak coupling limit of the N -particle Schrödinger equation. *Methods Appl. Anal.* 7, 2 (2000), 275–293. Cathleen Morawetz: a great mathematician.
- [5] BARTHE, F., CORDERO-ERAUSQUIN, D., AND MAUREY, B. Entropy of spherical marginals and related inequalities. *J. Math. Pures Appl. (9)* 86, 2 (2006), 89–99.
- [6] BOBYLÉV, A. V. The theory of the nonlinear spatially uniform Boltzmann equation for Maxwell molecules. In *Mathematical physics reviews, Vol. 7*, vol. 7 of *Soviet Sci. Rev. Sect. C Math. Phys. Rev.* Harwood Academic Publ., Chur, 1988, pp. 111–233.
- [7] BOLTZMANN, L. Weitere studien über das wärme gleichgenicht unfer gasmolökuler. *Sitzungsberichte der Akademie der Wissenschaften* 66 (1872), 275–370. Translation: Further studies on the thermal equilibrium of gas molecules, in *Kinetic Theory* 2, 88174, Ed. S.G. Brush, Pergamon, Oxford (1966).
- [8] BOLTZMANN, L. *Lectures on gas theory*. University of California Press, Berkeley, 1964. Translated by Stephen G. Brush. Reprint of the 1896/1898 Edition. Reprinted by Dover Publications, 1995.
- [9] CARLEMAN, T. Sur la théorie de l'équation intégrodifférentielle de Boltzmann. *Acta Math.* 60, 1 (1933), 91–146.
- [10] CARLEMAN, T. *Problèmes mathématiques dans la théorie cinétique des gaz*. Publ. Sci. Inst. Mittag-Leffler. 2. Almqvist & Wiksells Boktryckeri Ab, Uppsala, 1957.
- [11] CARLEN, E., CARVALHO, M. C., AND LOSS, M. Spectral gap for the kac model with hard collisions. Work in progress (personnal communication).
- [12] CARLEN, E. A., CARVALHO, M. C., LE ROUX, J., LOSS, M., AND VILLANI, C. Entropy and chaos in the Kac model. *Kinet. Relat. Models* 3, 1 (2010), 85–122.
- [13] CARLEN, E. A., CARVALHO, M. C., AND LOSS, M. Determination of the spectral gap for Kac's master equation and related stochastic evolution. *Acta Math.* 191, 1 (2003), 1–54.
- [14] CARLEN, E. A., GABETTA, E., AND TOSCANI, G. Propagation of smoothness and the rate of exponential convergence to equilibrium for a spatially homogeneous Maxwellian gas. *Comm. Math. Phys.* 199, 3 (1999), 521–546.
- [15] CARLEN, E. A., GERONIMO, J. S., AND LOSS, M. Determination of the spectral gap in the Kac model for physical momentum and energy-conserving collisions. *SIAM J. Math. Anal.* 40, 1 (2008), 327–364.
- [16] CARRAPATOSO, K. Some remarks on chaotic distributions on Boltzmann's sphere. Work in progress.
- [17] CARRILLO, J. A., AND TOSCANI, G. Contractive probability metrics and asymptotic behavior of dissipative kinetic equations. *Riv. Mat. Univ. Parma (7)* 6 (2007), 75–198.
- [18] CERCIGNANI, C. On the Boltzmann equation for rigid spheres. *Transport Theory Statist. Phys.* 2, 3 (1972), 211–225.
- [19] CERCIGNANI, C. *The Boltzmann equation and its applications*, vol. 67 of *Applied Mathematical Sciences*. Springer-Verlag, New York, 1988.
- [20] CERCIGNANI, C., ILLNER, R., AND PULVIRENTI, M. *The mathematical theory of dilute gases*, vol. 106 of *Applied Mathematical Sciences*. Springer-Verlag, New York, 1994.
- [21] CHAFAÏ, D., AND MALRIEU, F. On fine properties of mixtures with respect to concentration of measure and Sobolev type inequalities. *Ann. Inst. Henri Poincaré Probab. Stat.* 46, 1 (2010), 72–96.
- [22] DI BLASIO, G. Differentiability of spatially homogeneous solutions of the Boltzmann equation in the non Maxwellian case. *Comm. Math. Phys.* 38 (1974), 331–340.
- [23] DIACONIS, P., AND FREEDMAN, D. A dozen de Finetti-style results in search of a theory. *Ann. Inst. H. Poincaré Probab. Statist.* 23, 2, suppl. (1987), 397–423.
- [24] DI PERNA, R. J., AND LIONS, P.-L. Global solutions of Boltzmann's equation and the entropy inequality. *Arch. Rational Mech. Anal.* 114, 1 (1991), 47–55.
- [25] DOBRIĆ, V., AND YUKICH, J. E. Asymptotics for transportation cost in high dimensions. *J. Theoret. Probab.* 8, 1 (1995), 97–118.

- [26] ERDŐS, L., SCHLEIN, B., AND YAU, H.-T. Derivation of the Gross-Pitaevskii hierarchy for the dynamics of Bose-Einstein condensate. *Comm. Pure Appl. Math.* 59, 12 (2006), 1659–1741.
- [27] ERDŐS, L., SCHLEIN, B., AND YAU, H.-T. Derivation of the cubic non-linear Schrödinger equation from quantum dynamics of many-body systems. *Invent. Math.* 167, 3 (2007), 515–614.
- [28] ERDŐS, L., SCHLEIN, B., AND YAU, H.-T. Derivation of the Gross-Pitaevskii equation for the dynamics of Bose-Einstein condensate. *Ann. of Math. (2)* 172, 1 (2010), 291–370.
- [29] ESCOBEDO, M., AND MISCHLER, S. Scalings for a ballistic aggregation equation. *J. Stat. Phys.* 141, 3 (2010), 422–458.
- [30] FOURNIER, N., AND MÉLÉARD, S. Monte Carlo approximations and fluctuations for 2d Boltzmann equations without cutoff. *Markov Process. Related Fields* 7 (2001), 159–191.
- [31] FOURNIER, N., AND MÉLÉARD, S. A stochastic particle numerical method for 3d Boltzmann equation without cutoff. *Math. Comp.* 71 (2002), 583–604.
- [32] FOURNIER, N., AND MOUHOT, C. On the well-posedness of the spatially homogeneous Boltzmann equation with a moderate angular singularity. *Comm. Math. Phys.* 283, 3 (2009), 803–824.
- [33] GABETTA, G., TOSCANI, G., AND WENNERBERG, B. Metrics for probability distributions and the trend to equilibrium for solutions of the Boltzmann equation. *J. Statist. Phys.* 81, 5-6 (1995), 901–934.
- [34] GRAD, H. On the kinetic theory of rarefied gases. *Comm. Pure Appl. Math.* 2 (1949), 331–407.
- [35] GRAHAM, C., AND MÉLÉARD, S. Stochastic particle approximations for generalized Boltzmann models and convergence estimates. *The Annals of Probability* 25 (1997), 115–132.
- [36] GRÜNBAUM, F. A. Propagation of chaos for the Boltzmann equation. *Arch. Rational Mech. Anal.* 42 (1971), 323–345.
- [37] HAURAY, M., AND MISCHLER, S. On Kac's chaos and entropic chaos for many-particle systems. Work in progress.
- [38] HEWITT, E., AND SAVAGE, L. J. Symmetric measures on Cartesian products. *Trans. Amer. Math. Soc.* 80 (1955), 470–501.
- [39] ILLNER, R., AND PULVIRENTI, M. Global validity of the Boltzmann equation for a two-dimensional rare gas in vacuum. *Comm. Math. Phys.* 105, 2 (1986), 189–203.
- [40] JANVRESSE, E. Spectral gap for Kac's model of Boltzmann equation. *Ann. Probab.* 29, 1 (2001), 288–304.
- [41] JORDAN, R., KINDERLEHRER, D., AND OTTO, F. The variational formulation of the Fokker-Planck equation. *SIAM J. Math. Anal.* 29, 1 (1998), 1–17.
- [42] KAC, M. Foundations of kinetic theory. In *Proceedings of the Third Berkeley Symposium on Mathematical Statistics and Probability, 1954–1955, vol. III* (Berkeley and Los Angeles, 1956), University of California Press, pp. 171–197.
- [43] KAC, M. *Probability and related topics in physical sciences*, vol. 1957 of *With special lectures by G. E. Uhlenbeck, A. R. Hibbs, and B. van der Pol. Lectures in Applied Mathematics. Proceedings of the Summer Seminar, Boulder, Colo.* Interscience Publishers, London-New York, 1959.
- [44] KOLOKOLTSOV, V. N. *Nonlinear Markov processes and kinetic equations*, vol. 182 of *Cambridge Tracts in Mathematics*. Cambridge University Press, Cambridge, 2010.
- [45] LANFORD, III, O. E. Time evolution of large classical systems. In *Dynamical systems, theory and applications (Recontres, Battelle Res. Inst., Seattle, Wash., 1974)*. Springer, Berlin, 1975, pp. 1–111. *Lecture Notes in Phys.*, Vol. 38.
- [46] LIONS, P.-L. Théorie des jeux de champ moyen et applications (mean field games). In *Cours du Collège de France*. http://www.college-de-france.fr/default/EN/all/equ_der/audio_video.jsp, 2007–2009.
- [47] LU, X. Conservation of energy, entropy identity, and local stability for the spatially homogeneous Boltzmann equation. *J. Statist. Phys.* 96, 3-4 (1999), 765–796.
- [48] LU, X., AND MOUHOT, C. On measure solutions of the Boltzmann equation part I: Moment production and stability estimates. Preprint 2011 arXiv:1102.0373.
- [49] LU, X., AND WENNERBERG, B. Solutions with increasing energy for the spatially homogeneous Boltzmann equation. *Nonlinear Anal. Real World Appl.* 3, 2 (2002), 243–258.
- [50] MASLEN, D. K. The eigenvalues of Kac's master equation. *Math. Z.* 243, 2 (2003), 291–331.
- [51] MAXWELL, J. C. On the dynamical theory of gases. *Philos. Trans. Roy. Soc. London Ser. A* 157 (1867), 49–88.
- [52] MCKEAN, H. P. Fluctuations in the kinetic theory of gases. *Comm. Pure Appl. Math.* 28, 4 (1975), 435–455.
- [53] MCKEAN, JR., H. P. An exponential formula for solving Boltzmann's equation for a Maxwellian gas. *J. Combinatorial Theory* 2 (1967), 358–382.
- [54] MEHLER, F. G. Ueber die Entwicklung einer function von beliebig vielen variablen nach Laplace'schen functionen höherer ordnung. *Crelle's Journal* (1866), 161–176.

- [55] MÉLÉARD, S. Asymptotic behaviour of some interacting particle systems; McKean-Vlasov and Boltzmann models. In *Probabilistic models for nonlinear partial differential equations (Montecatini Terme, 1995)*, vol. 1627 of *Lecture Notes in Math.* Springer, Berlin, 1996, pp. 42–95.
- [56] MISCHLER, S. Sur le programme de Kac (concernant les limites de champ moyen). Séminaire EDP-X, Décembre 2010, Preprint arXiv.
- [57] MISCHLER, S., AND MOUHOT, C. On chaos properties of a Boltzmann-Kac system for inelastic collisions. Work in progress.
- [58] MISCHLER, S., MOUHOT, C., AND WENNERBERG, M. A new approach to quantitative propagation of chaos for drift, diffusion and jump processes. Preprint 2011, arXiv:1101.4727.
- [59] MISCHLER, S., AND WENNERBERG, B. On the spatially homogeneous Boltzmann equation. *Ann. Inst. H. Poincaré Anal. Non Linéaire* 16, 4 (1999), 467–501.
- [60] MOUHOT, C. Rate of convergence to equilibrium for the spatially homogeneous Boltzmann equation with hard potentials. *Comm. Math. Phys.* 261, 3 (2006), 629–672.
- [61] OTTO, F., AND VILLANI, C. Generalization of an inequality by Talagrand and links with the logarithmic Sobolev inequality. *J. Funct. Anal.* 173, 2 (2000), 361–400.
- [62] PAZY, A. *Semigroups of linear operators and applications to partial differential equations*, vol. 44 of *Applied Mathematical Sciences.* Springer-Verlag, New York, 1983.
- [63] PEYRE, R. Some ideas about quantitative convergence of collision models to their mean field limit. *J. Stat. Phys.* 136, 6 (2009), 1105–1130.
- [64] RACHEV, S. T., AND RÜSCHENDORF, L. *Mass transportation problems. Vol. II. Probability and its Applications* (New York). Springer-Verlag, New York, 1998. Applications.
- [65] SPOHN, H. *Large scale dynamics of interacting particles.* Texts and Monograph in physics. Springer-Verlag, 1991.
- [66] SZNITMAN, A.-S. Topics in propagation of chaos. In *École d’Été de Probabilités de Saint-Flour XIX—1989*, vol. 1464 of *Lecture Notes in Math.* Springer, Berlin, 1991, pp. 165–251.
- [67] TANAKA, H. Probabilistic treatment of the Boltzmann equation of Maxwellian molecules. *Z. Wahrsch. Verw. Gebiete* 46, 1 (1978/79), 67–105.
- [68] TANAKA, H. Some probabilistic problems in the spatially homogeneous Boltzmann equation. In *Theory and application of random fields (Bangalore, 1982)*, vol. 49 of *Lecture Notes in Control and Inform. Sci.* Springer, Berlin, 1983, pp. 258–267.
- [69] TOSCANI, G., AND VILLANI, C. Probability metrics and uniqueness of the solution to the Boltzmann equation for a Maxwell gas. *J. Statist. Phys.* 94, 3-4 (1999), 619–637.
- [70] VILLANI, C. Limite de champ moyen. Cours de DEA, 2001-2002, ÉNS Lyon.
- [71] VILLANI, C. Fisher information estimates for Boltzmann’s collision operator. *J. Math. Pures Appl.* (9) 77, 8 (1998), 821–837.
- [72] VILLANI, C. *Topics in Optimal Transportation*, vol. 58 of *Graduate Studies in Mathematics series.* American Mathematical Society, 2003.
- [73] VILLANI, C. *Optimal transport*, vol. 338 of *Grundlehren der Mathematischen Wissenschaften [Fundamental Principles of Mathematical Sciences]*. Springer-Verlag, Berlin, 2009. Old and new.

STÉPHANE MISCHLER

UNIVERSITÉ PARIS-DAUPHINE
 CEREMADE, UMR CNRS 7534
 PLACE DU MARÉCHAL DE LATTRE DE TASSIGNY 75775 PARIS CEDEX 16
 FRANCE

E-MAIL: mischler@ceremade.dauphine.fr

CLÉMENT MOUHOT

UNIVERSITY OF CAMBRIDGE
 DPMMS, CENTRE FOR MATHEMATICAL SCIENCES
 WILBERFORCE ROAD, CAMBRIDGE CB3 0WA, UK

E-MAIL: C.Mouhot@dpmms.cam.ac.uk