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A NOTE ON H^p_w -BOUNDEDNESS OF RIESZ TRANSFORMS AND θ -CALDERÓN-ZYGMUND OPERATORS THROUGH MOLECULAR CHARACTERIZATION

LUONG DANG KY

ABSTRACT. Let 0 and <math>w in the Muckenhoupt class A_1 . Recently, by using the weighted atomic decomposition and molecular characterization; Lee, Lin and Yang [11] (J. Math. Anal. Appl. 301 (2005), 394–400) established that the Riesz transforms $R_j, j = 1, 2, ..., n$, are bounded on $H_w^p(\mathbb{R}^n)$. In this note we extend this to the general case of weight w in the Muckenhoupt class A_∞ through molecular characterization. One difficulty, which has not been taken care in [11], consists in passing from atoms to all functions in $H_w^p(\mathbb{R}^n)$. Furthermore, the H_w^p -boundedness of θ -Calderón-Zygmund operators are also given through molecular characterization and atomic decomposition.

1. INTRODUCTION

Calderón-Zygmund operators and their generalizations on Euclidean space \mathbb{R}^n have been extensively studied, see for example [7, 14, 18, 15]. In particular, Yabuta [18] introduced certain θ -Calderón-Zygmund operators to facilitate his study of certain classes of pseudo-differential operator.

Definition 1.1. Let θ be a nonnegative nondecreasing function on $(0, \infty)$ satisfying $\int_0^1 \frac{\theta(t)}{t} dt < \infty$. A continuous function $K : \mathbb{R}^n \times \mathbb{R}^n \setminus \{(x, x) : x \in \mathbb{R}^n\} \to \mathbb{C}$ is said to be a θ -Calderón-Zygmund singular integral kernel if there exists a constant C > 0 such that

$$|K(x,y)| \le \frac{C}{|x-y|^n}$$

for all $x \neq y$,

$$|K(x,y) - K(x',y)| + |K(y,x) - K(y,x')| \le C \frac{1}{|x-y|^n} \theta\Big(\frac{|x-x'|}{|x-y|}\Big)$$

for all $2|x - x'| \le |x - y|$.

A linear operator $T : \mathcal{S}(\mathbb{R}^n) \to \mathcal{S}'(\mathbb{R}^n)$ is said to be a θ -Calderón-Zygmund operator if T can be extended to a bounded operator on $L^2(\mathbb{R}^n)$ and there exists

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a θ -Calderon-Zygmund singular integral kernel K such that for all $f \in C_c^{\infty}(\mathbb{R}^n)$ and all $x \notin supp f$, we have

$$Tf(x) = \int_{\mathbb{R}^n} K(x, y) f(y) dy.$$

When $K_j(x,y) = \pi^{-(n+1)/2} \Gamma\left(\frac{n+1}{2}\right) \frac{x_j - y_j}{|x-y|^{n+1}}, j = 1, 2, ..., n$, then they are the classical Riesz transforms denoted by R_j .

It is well-known that the Riesz transforms $R_j, j = 1, 2, ..., n$, are bounded on unweighted Hardy spaces $H^p(\mathbb{R}^n)$. There are many different approaches to prove this classical result (see [11, 9]). Recently, by using the weighted molecular theory (see [10]) and combined with García-Cuerva's atomic decomposition [5] for weighted Hardy spaces $H^p_w(\mathbb{R}^n)$, the authors in [11] established that the Riesz transforms $R_j, j = 1, 2, ..., n$, are bounded on $H^p_w(\mathbb{R}^n)$. More precisely, they proved that $||R_j f||_{H^p_w} \leq C$ for every $w \cdot (p, \infty, ts - 1)$ -atom where $s, t \in \mathbb{N}$ satisfy $n/(n+s) and <math>((s-1)r_w+n)/(s(r_w-1))$ with r_w is the critical index of w for the reverse Hölder condition. Remark that this leaves a gap in the proof. Similar gaps exist in some litteratures, for instance in [10, 15] when the authors establish H^p_w -boundedness of Calderón-Zygmund type operators. Indeed, it is now well-known that (see [1]) the argument "the operator T is uniformly bounded in $H^p_w(\mathbb{R}^n)$ on $w(p,\infty,r)$ -atoms, and hence it extends to a bounded operator on $H^p_w(\mathbb{R}^n)$ " is wrong in general. However, Meda, Sjögren and Vallarino [13] establishes that (in the setting of unweighted Hardy spaces) this is correct if one replaces L^{∞} -atoms by L^{q} -atoms with $1 < q < \infty$. See also [19] for L^2 -atoms with a different method from [13]. Later, the authors in [2] extended these results to the weighted anisotropic Hardy spaces. More precisely, it is claimed in [2] that the operator T can be extended to a bounded operator on $H^p_w(\mathbb{R}^n)$ if it is uniformly bounded on w(p,q,r)-atoms for $q_w < q < \infty, r \ge [n(q_w/p - 1)]$ where q_w is the critical index of w.

Motivated by [11, 10, 15, 1, 2], in this paper, we extend Theorem 1 in [11] to A_{∞} weights (see Theorem 1.1); Theorem 4 in [10] (see Theorem 1.2), Theorem 3 in [15] (see Theorem 3.1) to θ -Calderón-Zygmund operators; and fill the gaps of the proofs by using the atomic decomposition and molecular characterization of $H_w^p(\mathbb{R}^n)$ as in [11].

Throughout the whole paper, C denotes a positive geometric constant which is independent of the main parameters, but may change from line to line. In \mathbb{R}^n , we denote by B = B(x, r) an open ball with center x and radius r > 0. For any measurable set E, we denote by |E| its Lebesgue measure, and by E^c the set $\mathbb{R}^n \setminus E$.

Let us first recall some notations, definitions and well-known results.

Let $1 \leq p < \infty$. A nonnegative locally integrable function w belongs to the Muckenhoupt class A_p , say $w \in A_p$, if there exists a positive constant C so

that

$$\frac{1}{|B|} \int_{B} w(x) dx \left(\frac{1}{|B|} \int_{B} (w(x))^{-1/(p-1)} dx \right)^{p-1} \le C, \quad \text{if } 1$$

and

$$\frac{1}{|B|} \int_{B} w(x) dx \le C \operatorname{ess-inf}_{x \in B} w(x), \quad \text{if } p = 1,$$

for all balls B in \mathbb{R}^n . We say that $w \in A_\infty$ if $w \in A_p$ for some $p \in [1, \infty)$.

It is well known that $w \in A_p$, $1 \leq p < \infty$, implies $w \in A_q$ for all q > p. Also, if $w \in A_p$, $1 , then <math>w \in A_q$ for some $q \in [1, p)$. We thus write $q_w := \inf\{p \geq 1 : w \in A_p\}$ to denote the critical index of w. For a measurable set E, we note $w(E) = \int_E w(x) dx$ its weighted measure.

The following lemma gives a characterization of the class A_p , $1 \le p < \infty$. It can be found in [6].

Lemma A. The function $w \in A_p$, $1 \leq p < \infty$, if and only if, for all non-negative functions and all balls B,

$$\left(\frac{1}{|B|} \int_{B} f(x) dx\right)^{p} \leq C \frac{1}{w(B)} \int_{B} f(x)^{p} w(x) dx.$$

A close relation to A_p is the reverse Hölder condition. If there exist r > 1and a fixed constant C > 0 such that

$$\left(\frac{1}{|B|} \int_{B} w^{r}(x) dx\right)^{1/r} \leq C\left(\frac{1}{|B|} \int_{B} w(x) dx\right) \quad \text{for every ball } B \subset \mathbb{R}^{n},$$

we say that w satisfies reverse Hölder condition of order r and write $w \in RH_r$. It is known that if $w \in RH_r$, r > 1, then $w \in RH_{r+\varepsilon}$ for some $\varepsilon > 0$. We thus write $r_w := \sup\{r > 1 : w \in RH_r\}$ to denote the critical index of w for the reverse Hölder condition.

The following result provides us the comparison between the Lebesgue measure of a set E and its weighted measure w(E). It also can be found in [6].

Lemma B. Let $w \in A_p \cap RH_r$, $p \ge 1$ and r > 1. Then there exist constants $C_1, C_2 > 0$ such that

$$C_1\left(\frac{|E|}{|B|}\right)^p \le \frac{w(E)}{w(B)} \le C_2\left(\frac{|E|}{|B|}\right)^{(r-1)/r},$$

for all cubes B and measurable subsets $E \subset B$.

Given a weight function w on \mathbb{R}^n , as usual we denote by $L^q_w(\mathbb{R}^n)$ the space of all functions f satisfying $||f||_{L^q_w} := (\int_{\mathbb{R}^n} |f(x)|^q w(x) dx)^{1/q} < \infty$. When $q = \infty$, $L^\infty_w(\mathbb{R}^n)$ is $L^\infty(\mathbb{R}^n)$ and $||f||_{L^\infty_w} = ||f||_{L^\infty}$. Analogously to the classical Hardy spaces, the weighted Hardy spaces $H^p_w(\mathbb{R}^n)$, p > 0, can be defined in terms of maximal functions. Namely, let ϕ be a function in $\mathcal{S}(\mathbb{R}^n)$, the Schwartz space of rapidly decreasing smooth functions, satisfying $\int_{\mathbb{R}^n} \phi(x) dx = 1$. Define

$$\phi_t(x) = t^{-n}\phi(x/t), \quad t > 0, x \in \mathbb{R}^n,$$

and the maximal function f^* by

$$f^*(x) = \sup_{t>0} |f * \phi_t(x)|, \quad x \in \mathbb{R}^n.$$

Then $H^p_w(\mathbb{R}^n)$ consists of those tempered distributions $f \in \mathcal{S}'(\mathbb{R}^n)$ for which $f^* \in L^p_w(\mathbb{R}^n)$ with the (quasi-)norm

$$||f||_{H^p_w} = ||f^*||_{L^p_w}.$$

In order to show the H_w^p -boundedness of Riesz transforms, we characterize weighted Hardy spaces in terms of atoms and molecules in the following way.

Definition of a weighted atom. Let $0 and <math>p \ne q$ such that $w \in A_q$. Let q_w be the critical index of w. Set [·] the integer function. For $s \in \mathbb{N}$ satisfying $s \ge [n(q_w/p - 1)]$, a function $a \in L^q_w(\mathbb{R}^n)$ is called w-(p, q, s)-atom centered at x_0 if

- (i) supp $a \subset B$ for some ball B centered at x_0 ,
- (ii) $||a||_{L^q_w} \le w(B)^{1/q-1/p}$,

(iii) $\int_{\mathbb{R}^n} a(x) x^{\alpha} dx = 0$ for every multi-index α with $|\alpha| \leq s$.

Let $H^{p,q,s}_{w}(\mathbb{R}^{n})$ denote the space consisting of tempered distributions admitting a decomposition $f = \sum_{j=1}^{\infty} \lambda_{j} a_{j}$ in $\mathcal{S}'(\mathbb{R}^{n})$, where a_{j} 's are w-(p,q,s)-atoms and $\sum_{j=1}^{\infty} |\lambda_{j}|^{p} < \infty$. And for every $f \in H^{p,q,s}_{w}(\mathbb{R}^{n})$, we consider the (quasi-)norm

$$\|f\|_{H^{p,q,s}_{w}} = \inf \left\{ \left(\sum_{j=1}^{\infty} |a_{j}|^{p}\right)^{1/p} : f \stackrel{\mathcal{S}'}{=} \sum_{j=1}^{\infty} \lambda_{j} a_{j}, \ \{a_{j}\}_{j=1}^{\infty} \text{ are } w \cdot (p,q,s) \text{ -atoms} \right\}.$$

Denote by $H^{p,q,s}_{w,\text{fin}}(\mathbb{R}^n)$ the vector space of all finite linear combinations of w-(p,q,s)-atoms, and the (quasi-)norm of f in $H^{p,q,s}_{w,\text{fin}}(\mathbb{R}^n)$ is defined by

$$\|f\|_{H^{p,q,s}_{w,\mathrm{fin}}} := \inf \left\{ \left(\sum_{j=1}^{k} |\lambda_j|^p \right)^{1/p} : f = \sum_{j=1}^{k} \lambda_j a_j, k \in \mathbb{N}, \{a_j\}_{j=1}^k \text{ are } w\text{-}(p,q,s)\text{-}\mathrm{atoms} \right\}.$$

We have the following atomic decomposition for $H^p_w(\mathbb{R}^n)$. It can be found in [5] (see also [2, 8]).

Theorem A. If the triplet (p, q, s) satisfies the conditions of w-(p, q, s)-atoms, then $H^p_w(\mathbb{R}^n) = H^{p,q,s}_w(\mathbb{R}^n)$ with equivalent norms.

The molecules corresponding to the atoms mentioned above can be defined as follows.

Definition of a weighted molecule. For $0 and <math>p \ne q$, let $w \in A_q$ with critical index q_w and critical index r_w for the reverse Hölder

condition. Set $s \ge [n(q_w/p-1)], \varepsilon > \max\{sr_w(r_w-1)^{-1}n^{-1} + (r_w-1)^{-1}, 1/p - 1\}$ 1}, $a = 1 - 1/p + \varepsilon$, and $b = 1 - 1/q + \varepsilon$. A w- (p, q, s, ε) -molecule centered at x_0 is a function $M \in L^q_w(\mathbb{R}^n)$ satisfying

(i) $M.w(B(x_0, \cdot -x_0))^b \in L^q_w(\mathbb{R}^n),$ (ii) $\|M\|_{L^q_w}^{a/b} \|M.w(B(x_0, \cdot -x_0))^b\|_{L^q_w}^{1-a/b} \equiv \mathfrak{N}_w(M) < \infty,$ (iii) $\int_{\mathbb{R}^n} M(x) x^{\alpha} dx = 0$ for every multi-index α with $|\alpha| \leq s.$

The above quantity $\mathfrak{N}_w(M)$ is called the *w*-molecular norm of *M*.

In [10], Lee and Lin proved that every weighted molecule belongs to the weighted Hardy space $H^p_w(\mathbb{R}^n)$, and the embedding is continuous.

Theorem B. Let $0 and <math>p \ne q$, $w \in A_q$, and (p, q, s, ε) be the quadruple in the definition of molecule. Then, every w- (p, q, s, ε) -molecule M centered at any point in \mathbb{R}^n is in $H^p_w(\mathbb{R}^n)$, and $\|M\|_{H^p_w} \leq C\mathfrak{N}_w(M)$ where the constant C is independent of the molecule.

Although, in general, one cannot conclude that an operator T is bounded on $H^p_w(\mathbb{R}^n)$ by checking that their norms have uniform bound on all of the corresponding $w(p, \infty, s)$ -atoms (cf. [1]). However, this is correct when dealing with w(p,q,s)-atoms with $q_w < q < \infty$. Indeed, we have the following result (see [2, Theorem 7.2]).

Theorem C. Let $0 , <math>w \in A_{\infty}$, $q \in (q_w, \infty)$ and $s \in \mathbb{Z}$ satisfying $s \geq [n(q_w/p - 1)]$. Suppose that $T : H^{p,q,s}_{w,\text{fin}}(\mathbb{R}^n) \to H^p_w(\mathbb{R}^n)$ is a linear operator satisfying

 $\sup\{\|Ta\|_{H^p_w}: a \text{ is any } w-(p,q,s)-\text{atom}\} < \infty.$

Then T can be extended to a bounded linear operator on $H^p_w(\mathbb{R}^n)$.

Our first main result, which generalizes Theorem 1 in [11], is as follows:

Theorem 1.1. Let $0 and <math>w \in A_{\infty}$. Then, the Riesz transforms are bounded on $H^p_w(\mathbb{R}^n)$.

For the next result, we need the notion $T^*1 = 0$.

Definition 1.2. Let T be a θ -Calderón-Zygmund operator. We say that $T^*1 =$ 0 if $\int_{\mathbb{R}^n} Tf(x) dx = 0$ for all $f \in L^q(\mathbb{R}^n), 1 < q \leq \infty$, with compact support and $\int_{\mathbb{R}^n} f(x) dx = 0.$

We now can give the H_w^p -boundedness of θ -Calderón-Zygmund type operators, which generalizes Theorem 4 in [10] by taking q = 1 and $\theta(t) = t^{\delta}$, as follows:

Theorem 1.2. Given $\delta \in (0,1]$, $n/(n+\delta) , and <math>w \in A_q \cap RH_r$ with $1 \leq q < p(n+\delta)/n$, $(n+\delta)/(n+\delta-nq) < r$. Let θ be a nonnegative nondecreasing function on $(0,\infty)$ with $\int_0^1 \frac{\theta(t)}{t^{1+\delta}} dt < \infty$, and T be a θ -Calderón-Zygmund operator satisfying $T^*1 = 0$. Then T is bounded on $H^p_w(\mathbb{R}^n)$.

2. Proof of Theorem 1.1

In order to prove the main theorems, we need the following lemma (see [6, page 412]).

Lemma C. Let $w \in A_r, r > 1$. Then there exists a constant C > 0 such that

$$\int_{B^c} \frac{1}{|x - x_0|^{nr}} w(x) dx \le C \frac{1}{\sigma^{nr}} w(B)$$

for all balls $B = B(x_0, \sigma)$ in \mathbb{R}^n .

Proof of Theorem 1.1. For $q = 2(q_w + 1) \in (q_w, \infty)$, then $s := [n(q/p - 1)] \ge [n(q_w/p - 1)]$. We now choose (and fix) a positive number ε satisfying

(2.1)
$$\max\{sr_w(r_w-1)^{-1}n^{-1}+(r_w-1)^{-1}, q/p-1\} < \varepsilon < t(s+1)(nq)^{-1}+q^{-1}-1,$$

for some $t \in \mathbb{N}, t \ge 1$ and $\max\{sr_w(r_w - 1)^{-1}n^{-1} + (r_w - 1)^{-1}, q/p - 1\} < t(s+1)(nq)^{-1} + q^{-1} - 1.$

Clearly, $\ell := t(s+1) - 1 \ge s \ge [n(q_w/p-1)]$. Hence, by Theorem B and Theorem C, it is sufficient to show that for every w- (p,q,ℓ) -atom f centered at x_0 and supported in ball $B = B(x_0, \sigma)$, the Riesz transforms $R_j f = K_j * f$, j = 1, 2, ..., n, are w- (p, q, s, ε) -molecules with the norm $\mathfrak{N}_w(R_j f) \le C$.

Indeed, as $w \in A_q$ by $q = 2(q_w + 1) \in (q_w, \infty)$. It follows from L_w^q -boundedness of Riesz transforms that

(2.2)
$$\|R_j f\|_{L^q_w} \le \|R_j\|_{L^q_w \to L^q_w} \|f\|_{L^q_w} \le Cw(B)^{1/q-1/p}$$

To estimate $||R_j f.w(B(x_0, |\cdot -x_0|))^b||_{L^q_w}$ where $b = 1 - 1/q + \varepsilon$, we write

$$\begin{aligned} \|R_{j}f.w(B(x_{0},\cdot-x_{0}))^{b}\|_{L^{q}_{w}}^{q} &= \int_{|x-x_{0}| \leq 2\sqrt{n}\sigma} |R_{j}f(x)|^{q}w(B(x_{0},|x-x_{0}|))^{bq}w(x)dx + \\ &+ \int_{|x-x_{0}| > 2\sqrt{n}\sigma} |R_{j}f(x)|^{q}w(B(x_{0},|x-x_{0}|))^{bq}w(x)dx \\ &= I + II. \end{aligned}$$

By Lemma B, we have the following estimate,

$$I = \int_{|x-x_0| \le 2\sqrt{n\sigma}} |R_j f(x)|^q w(B(x_0, |x-x_0|))^{bq} w(x) dx$$

$$\le w(B(x_0, 2\sqrt{n\sigma}))^{bq} \int_{|x-x_0| \le 2\sqrt{n\sigma}} |R_j f(x)|^q w(x) dx$$

$$\le Cw(B)^{bq} \|R_j\|_{L^q_w \to L^q_w}^q \|f\|_{L^q_w}^q \le Cw(B)^{(b+1/q-1/p)q}.$$

To estimate II, as f is w- (p, q, ℓ) -atom, by the Taylor's fomular and Lemma A, we get

$$|K_{j} * f(x)| = \left| \int_{|y-x_{0}| \leq \sigma} \left(K_{j}(x-y) - \sum_{|\alpha| \leq \ell} \frac{1}{\alpha!} D^{\alpha} K_{j}(x-x_{0})(x_{0}-y)^{\alpha} \right) f(y) dy \right|$$

$$\leq C \int_{|y-x_{0}| \leq \sigma} \frac{\sigma^{\ell+1}}{|x-x_{0}|^{n+\ell+1}} |f(y)| dy$$

$$\leq C \frac{\sigma^{n+\ell+1}}{|x-x_{0}|^{n+\ell+1}} w(B)^{-1/q} ||f||_{L_{w}^{q}},$$

for all $x \in (B(x_0, 2\sqrt{n\sigma}))^c$. As $b = 1 - 1/q + \varepsilon$, it follows from (2.1) that $(n + \ell + 1)q - q^2nb > nq$. Therefore, by combining the above inequality, Lemma B and Lemma C, we obtain

$$II = \int_{|x-x_0|>2\sqrt{n\sigma}} |R_j f(x)|^q w(B(x_0, |x-x_0|))^{bq} w(x) dx$$

$$\leq C\sigma^{(n+\ell+1)q} w(B)^{-1} ||f||^q_{L^q_w} \int_{|x-x_0|>2\sqrt{n\sigma}} \frac{1}{|x-x_0|^{(n+\ell+1)q}} w(B(x_0, |x-x_0|))^{bq} w(x) dx$$

$$\leq C\sigma^{(n+\ell+1)q-q^2nb} w(B)^{(b-1/p)q} \int_{|x-x_0|>2\sqrt{n\sigma}} \frac{1}{|x-x_0|^{(n+\ell+1)q-q^2nb}} w(x) dx$$

$$\leq Cw(B)^{(b+1/q-1/p)q}.$$

Thus,

(2.3)
$$||R_j f.w(B(x_0, |\cdot -x_0|))^b||_{L^q_w} = (I + II)^{1/q} \le Cw(B)^{b+1/q-1/p}.$$

Remark that $a = 1 - 1/p + \varepsilon$. Combining (2.2) and (2.3), we obtain

$$\mathfrak{N}_w(R_j f) \le Cw(B)^{(1/q-1/p)a/b} w(B)^{(b+1/q-1/p)(1-a/b)} \le C.$$

The proof will be concluded if we establish the vanishing moment conditions of $R_j f$. One first consider the following lemma.

Lemma. For every classical atom $(p, 2, \ell)$ -atom g centered at x_0 , we have

$$\int_{\mathbb{R}^n} R_j g(x) x^{\alpha} dx = 0 \quad \text{for } 0 \le |\alpha| \le s, 1 \le j \le n.$$

Proof of the Lemma. Since $b = 1 - 1/q + \varepsilon < (\ell + 1)(nq)^{-1} < (\ell + 1)n^{-1}$, we obtain $2(n + \ell + 1) - 2nb > n$. It is similar to the previous argument, we also obtain that R_jg and $R_jg.|\cdot -x_0|^{nb}$ belong to $L^2(\mathbb{R}^n)$. Now, we establish that $R_jg.(\cdot - x_0)^{\alpha} \in L^1(\mathbb{R}^n)$ for every multi-index α with $|\alpha| \leq s$. Indeed, since $\varepsilon > q/p - 1$ by (2.1), implies that 2(s - nb) < (s - nb)q' < -n by $q = 2(q_w + 1) > 2$, where 1/q + 1/q' = 1. We use Schwartz inequality to get

$$\int_{B(x_0,1)^c} |R_j g(x)(x-x_0)^{\alpha}| dx \leq \int_{B(x_0,1)^c} |R_j g(x)| |x-x_0|^s dx \\
\leq \left(\int_{B(x_0,1)^c} |R_j g(x)|^2 |x-x_0|^{2nb} dx \right)^{1/2} \left(\int_{B(x_0,1)^c} |x-x_0|^{2(s-nb)} dx \right)^{1/2} \\
\leq C \|R_j g_{\cdot}| \cdot -x_0|^{nb} \|_{L^2} < \infty,$$

and

$$\int_{B(x_0,1)} |R_j g(x)(x-x_0)^{\alpha}| dx \le |B(x_0,1)|^{1/2} \Big(\int_{B(x_0,1)} |R_j g(x)|^2 dx\Big)^{1/2} < \infty$$

Thus, $R_j g.(\cdot - x_0)^{\alpha} \in L^1(\mathbb{R}^n)$ for any $|\alpha| \leq s$. Deduce that $R_j g(x) x^{\alpha} \in L^1(\mathbb{R}^n)$ for any $|\alpha| \leq s$. Therefore,

$$(R_j g(x) x^{\alpha})(\xi) = C_{\alpha} . D^{\alpha} (R_j g)(\xi)$$

is continuous, with $|C_{\alpha}| \leq C_s$ (C_s depends only on s) for any $|\alpha| \leq s$, where \hat{h} is used to denote the fourier transform of h. Consequently,

$$\int_{\mathbb{R}^n} R_j g(x) x^{\alpha} dx = C_{\alpha} . D^{\alpha} \widehat{(R_j g)}(0) = C_{\alpha} . D^{\alpha} (m_j \hat{g})(0),$$

where $m_j(x) = -ix_j/|x|$. Moreover, since g is a classical $(p, 2, \ell)$ -atom, it follows from [17, Lemma 9.1] that \hat{g} is ℓ th order differentiable and $\hat{g}(\xi) = O(|\xi|^{\ell+1})$ as $\xi \to 0$. We write e_j to be the *j*th standard basis vector of \mathbb{R}^n , $\alpha = (\alpha_1, ..., \alpha_n)$ a multi-index of nonnegative integers α_j , $\Delta_{he_j}\phi(x) = \phi(x) - \phi(x - he_j)$, $\Delta_{he_j}^{\alpha_j}\phi(x) = \Delta_{he_j}^{\alpha_j-1}\phi(x) - \Delta_{he_j}^{\alpha_j-1}\phi(x - he_j)$ for $\alpha_j \ge 2$, $\Delta_{he_j}^0\phi(x) = \phi(x)$, and $\Delta_h^{\alpha} = \Delta_{he_1}^{\alpha_1}...\Delta_{he_n}^{\alpha_n}$. Then, the boundedness of m_j , and $|C_{\alpha}| \le C_s$ for $|\alpha| \le s$, implies

$$\begin{aligned} \left| \int_{\mathbb{R}^n} R_j g(x) x^{\alpha} dx \right| &= |C_{\alpha}| \left| \lim_{h \to 0} |h|^{-|\alpha|} \Delta_h^{\alpha}(m_j \hat{g})(0) \right| \\ &\leq C \lim_{h \to 0} |h|^{\ell+1-|\alpha|} = 0, \end{aligned}$$

for $|\alpha| \leq s$ by $s \leq \ell$. Thus, for any j = 1, 2, ..., n, and $|\alpha| \leq s$,

$$\int_{\mathbb{R}^n} R_j g(x) x^\alpha dx = 0.$$

This complete the proof of the lemma.

Let us come back to the proof of Theorem 1.1. As $q/2 = q_w + 1 > q_w$, by Lemma A,

$$\left(\frac{1}{|B|} \int_{B} |f(x)|^2 dx\right)^{q/2} \le C \frac{1}{w(B)} \int_{B} |f(x)|^q w(x) dx.$$

Therefore, $g := C^{-1/q} |B|^{-1/p} w(B)^{1/p} f$ is a classical $(p, 2, \ell)$ -atom since f is w- (p, q, ℓ) -atom associated with ball B. Consequently, by the above lemma,

$$\int_{\mathbb{R}^n} R_j f(x) x^{\alpha} dx = C^{1/q} |B|^{1/p} w(B)^{-1/p} \int_{\mathbb{R}^n} R_j g(x) x^{\alpha} dx = 0$$

for all j = 1, 2, ..., n and $|\alpha| \leq s$. Thus, the theorem is proved.

Following a similar but easier argument, we also have the following H_w^p -boundedness of Hilbert transform. We leave details to readers.

Theorem 2.1. Let $0 and <math>w \in A_{\infty}$. Then, the Hilbert transform is bounded on $H^p_w(\mathbb{R})$.

3. Proof of theorem 1.2

We first consider the following lemma

Lemma 3.1. Let $p \in (0,1], w \in A_q, 1 < q < \infty$, and T be a θ -Calderón-Zygmund operator satisfying $T^*1 = 0$. Then, $\int_{\mathbb{R}^n} Tf(x)dx = 0$ for all w-(p,q,0)-atoms f.

Proof of Lemma 3.1. Let f be an arbitrary w-(p, q, 0)-atom associated with ball B. It is well-known that there exists 1 < r < q such that $w \in A_r$. Therefore, it follows from Lemma A that

$$\int_{B} |f(x)|^{q/r} dx \le C |B| w(B)^{1/r} ||f||_{L^{q}_{w}}^{q/r} < \infty.$$

We deduce that f is a multiple of classical (p, q/r, 0)-atom, and thus the condition $T^*1 = 0$ implies $\int_{\mathbb{R}^n} Tf(x) dx = 0$.

Proof of Theorem 1.2. Because of the hypothesis, without loss of generality we can assume q > 1. Futhermore, it is clear that $[n(q_w/p - 1)] = 0$, and there exists a positive constant ε such that

(3.1)
$$\max\left\{\frac{1}{r_w-1}, \frac{1}{p}-1\right\} < \varepsilon < \frac{n+\delta}{nq}-1.$$

Similarly to the arguments in Theorem 1.1, it is sufficient to show that, for every w-(p,q,0)-atom f centered at x_0 and supported in ball $B = B(x_0,\sigma)$,

Tf is a w- $(p,q,0,\varepsilon)$ -molecule with the norm $\mathfrak{N}_w(Tf) \leq C$. One first observe that $\int_{\mathbb{R}^n} Tf(x) dx = 0$ by Lemma 3.1, and

$$\sum_{k=0}^{\infty} \theta(2^{-k}) 2^{knbq} < \infty,$$

where $b = 1 - 1/q + \varepsilon$, by $\int_0^1 \frac{\theta(t)}{t^{1+\delta}} dt < \infty$ and (3.1). We deduce that

(3.2)
$$\sum_{k=0}^{\infty} \left(\theta(2^{-k})2^{knbq}\right)^q < \infty.$$

As $w \subset A_q$, $1 < q < \infty$, it follows from [18, Theorem 2.4] that

(3.3)
$$||Tf||_{L^q_w} \le C ||f||_{L^q_w} \le Cw(B)^{1/q-1/p}$$

To estimate $||Tf.w(B(x_0, |\cdot -x_0|))^b||_{L^q_w}$, we write

$$\begin{aligned} \|Tf.w(B(x_0, \cdot - x_0))^b\|_{L^q_w}^q &= \int_{|x-x_0| \le 2\sigma} |Tf(x)|^q w(B(x_0, |x-x_0|))^{bq} w(x) dx + \\ &+ \int_{|x-x_0| > 2\sigma} |Tf(x)|^q w(B(x_0, |x-x_0|))^{bq} w(x) dx = I + II. \end{aligned}$$

By Lemma B, we have the following estimate,

$$I = \int_{|x-x_0| \le 2\sigma} |Tf(x)|^q w(B(x_0, |x-x_0|))^{bq} w(x) dx$$

$$\leq w(B(x_0, 2\sigma))^{bq} \int_{|x-x_0| \le 2\sigma} |Tf(x)|^q w(x) dx$$

$$\leq Cw(B)^{bq} ||f||_{L^q_w}^q \le Cw(B)^{(b+1/q-1/p)q}.$$

To estimate II, since f is of mean zero, by Lemma A, we have

$$\begin{aligned} |Tf(x)| &= \left| \int_{|y-x_0| \le \sigma} \left(K(x,y) - K(x,x_0) \right) f(y) dy \right| \\ &\le C \int_{|y-x_0| \le \sigma} \frac{1}{|x-x_0|^n} \theta \Big(\frac{|y-x_0|}{|x-x_0|} \Big) |f(y)| dy \\ &\le C \frac{\sigma^n}{|x-x_0|^n} \theta \Big(\frac{\sigma}{|x-x_0|} \Big) w(B)^{-1/q} ||f||_{L^q_w}, \end{aligned}$$

for all $x \in (B(x_0, 2\sigma))^c$. Therefore, by combining the above inequality, Lemma B and (3.2), we obtain

$$II = \int_{|x-x_0|>2\sigma} |Tf(x)|^q w(B(x_0, |x-x_0|))^{bq} w(x) dx$$

$$\leq Cw(B)^{-1} ||f||_{L^q_w}^q \int_{|x-x_0|>2\sigma} \frac{\sigma^{nq}}{|x-x_0|^{nq}} \left(\theta\left(\frac{\sigma}{|x-x_0|}\right)\right)^q w(B(x_0, |x-x_0|))^{bq} w(x) dx$$

$$\leq Cw(B)^{-q/p} \sum_{k=1}^{\infty} \int_{2^k \sigma < |x-x_0| \le 2^{k+1}\sigma} \frac{\sigma^{nq}}{|x-x_0|^{nq}} \left(\theta\left(\frac{\sigma}{|x-x_0|}\right)\right)^q w(B(x_0, |x-x_0|))^{bq} w(x) dx$$

$$\leq Cw(B)^{(b+1/q-1/p)q} \sum_{k=0}^{\infty} \left(\theta(2^{-k})2^{knbq}\right)^q \le Cw(B)^{(b+1/q-1/p)q}.$$

Thus,

(3.4)
$$\|Tf.w(B(x_0, |\cdot -x_0|))^b\|_{L^q_w} = (I+II)^{1/q} \le Cw(B)^{b+1/q-1/p}.$$

Remark that $a = 1 - 1/p + \varepsilon$. Combining (3.3) and (3.4), we obtain

$$\mathfrak{N}_w(Tf) \le Cw(B)^{(1/q-1/p)a/b} w(B)^{(b+1/q-1/p)(1-a/b)} \le C.$$

This finishes the proof.

It is well-known that the molecular theory of (unweighted) Hardy spaces of Taibleson and Weiss [17] is one of useful tools to establish boundedness of operators in Hardy spaces (cf. [17, 12]). In the setting of Muckenhoupt weight, this theory has been considered by the authors in [10], since then, they have been well used to establish boundedness of operators in weighted Hardy spaces (cf. [10, 11, 3]). However in some cases, the weighted molecular characterization, which obtained in [10], does not give the best possible results. For Calderón-Zygmund type operators in Theorem 1.2, for instance, it involves assumption on the critical index of w for the reverse Hölder condition as the following theorem does not.

Theorem 3.1. Given $\delta \in (0,1]$, $n/(n+\delta) , and <math>w \in A_q$ with $1 \leq q < p(n+\delta)/n$. Let θ be a nonnegative nondecreasing function on $(0,\infty)$ with $\int_0^1 \frac{\theta(t)}{t^{1+\delta}} dt < \infty$, and T be a θ -Calderón-Zygmund operator satisfying $T^*1 = 0$. Then T is bounded on $H^p_w(\mathbb{R}^n)$.

The following corollary give the boundedness of the classical Calderón-Zygmund type operators on weighted Hardy spaces (see [15, Theorem 3]).

Corollary 3.1. Let $0 < \delta \leq 1$ and T be the classical δ -Calderón-Zygmund operator, i.e. $\theta(t) = t^{\delta}$, satisfying $T^*1 = 0$. If $n/(n+\delta) and <math>w \in A_q$ with $1 \leq q < p(n+\delta)/n$, then T is bounded on $H^p_w(\mathbb{R}^n)$.

Proof of Corollary 3.1. By taking $\delta' \in (0, \delta)$ which is close enough δ . Then, we apply Theorem 3.1 with δ' instead of δ .

Proof of Theorem 3.1. Without loss of generality we can assume $1 < q < p(n+\delta)/n$. Fix $\phi \in \mathcal{S}(\mathbb{R}^n)$ with $\int_{\mathbb{R}^n} \phi(x) dx \neq 0$. By Theorem C, it is sufficient to show that for every w-(p, q, 0)-atom f centered at x_0 and supported in ball $B = B(x_0, \sigma), ||(Tf)^*||_{L^p_w} \leq C$. In order to do this, one write

$$\begin{aligned} \|(Tf)^*\|_{L^p_w}^p &= \int_{|x-x_0| \le 4\sigma} \left((Tf)^*(x) \right)^p w(x) dx + \int_{|x-x_0| > 4\sigma} \left((Tf)^*(x) \right)^p w(x) dx \\ &= L_1 + L_2. \end{aligned}$$

By Hölder inequality, $L^q_w\mbox{-}\mathrm{boundedness}$ of the maximal function and Lemma B, we get

$$L_{1} \leq \left(\int_{|x-x_{0}| \leq 4\sigma} \left((Tf)^{*}(x) \right)^{q} w(x) dx \right)^{p/q} \left(\int_{|x-x_{0}| \leq 4\sigma} w(x) dx \right)^{1-p/q} \\ \leq C \|f\|_{L_{w}^{q}}^{p} w(B(x_{0}, 4\sigma))^{1-p/q} \leq C.$$

To estimate L_2 , we first estimate $(Tf)^*(x)$ for $|x - x_0| > 4\sigma$. For any t > 0, since $\int_{\mathbb{R}^n} Tf(x) dx = 0$ by Lemma 3.1, we get

$$\begin{aligned} |Tf * \phi_t(x)| &= \left| \int_{\mathbb{R}^n} Tf(y) \frac{1}{t^n} \left(\phi\left(\frac{x-y}{t}\right) - \phi\left(\frac{x-x_0}{t}\right) \right) dy \right| \\ &\leq \left| \frac{1}{t^n} \int_{|y-x_0| < 2\sigma} |Tf(y)| \left| \phi\left(\frac{x-y}{t}\right) - \phi\left(\frac{x-x_0}{t}\right) \right| dy \\ &+ \frac{1}{t^n} \int_{2\sigma \le |y-x_0| < \frac{|x-x_0|}{2}} \cdots + \frac{1}{t^n} \int_{|y-x_0| \ge \frac{|x-x_0|}{2}} \cdots \\ &= E_1(t) + E_2(t) + E_3(t). \end{aligned}$$

As $|x - x_0| > 4\sigma$, by the mean value theorem, Lemma A and Lemma B, we get

$$\begin{split} E_{1}(t) &= \left. \frac{1}{t^{n}} \int_{|y-x_{0}|<2\sigma} |Tf(y)| \left| \phi\left(\frac{x-y}{t}\right) - \phi\left(\frac{x-x_{0}}{t}\right) \right| dy \\ &\leq \left. \frac{1}{t^{n}} \int_{|y-x_{0}|<2\sigma} |Tf(y)| \frac{|y-x_{0}|}{t} \sup_{\lambda \in (0,1)} \sup_{\lambda \in (0,1)} \left| \nabla \phi\left(\frac{x-x_{0}+\lambda(y-x_{0})}{t}\right) \right| dy \\ &\leq \left. C \frac{\sigma}{|x-x_{0}|^{n+1}} \int_{|y-x_{0}|<2\sigma} |Tf(y)| dy \\ &\leq \left. C \frac{\sigma}{|x-x_{0}|^{n+1}} |B(x_{0},2\sigma)| w(B(x_{0},2\sigma))^{-1/q} ||Tf||_{L^{q}_{w}} \\ &\leq \left. C \frac{\sigma^{n+1}}{|x-x_{0}|^{n+1}} w(B)^{-1/q} ||f||_{L^{q}_{w}} \leq C \frac{\sigma^{n+1}}{|x-x_{0}|^{n+1}} w(B)^{-1/p}. \end{split}$$

Similarly, we also get

Next, let us look at L_3 . Similarly, we also have

$$\begin{split} E_{3}(t) &\leq \left| \frac{1}{t^{n}} \int_{|y-x_{0}| \geq \frac{|x-x_{0}|}{2}} \left| \int_{\mathbb{R}^{n}} f(z) \Big(K(y,z) - K(y,x_{0}) \Big) dz \right| \left(\left| \phi \Big(\frac{y-x_{0}}{t} \Big) \right| + 2 \Big| \phi \Big(\frac{x-x_{0}}{t} \Big) \Big| \Big) dy \\ &\leq \left| C \frac{1}{|x-x_{0}|^{n}} \int_{|y-x_{0}| \geq \frac{|x-x_{0}|}{2}} \int_{|z-x_{0}| < \sigma} |f(z)| \frac{1}{|y-x_{0}|^{n}} \theta \Big(\frac{|z-x_{0}|}{|y-x_{0}|} \Big) dz dy \\ &\leq \left| C \Big(\frac{\sigma}{|x-x_{0}|} \Big)^{n} \int_{0}^{2\sigma/|x-x_{0}|} \frac{\theta(t)}{t} dt w(B)^{-1/p} \\ &\leq \left| C \Big(\frac{\sigma}{|x-x_{0}|} \Big)^{n} \int_{0}^{2\sigma/|x-x_{0}|} \frac{\theta(t)}{t^{1+\delta}} dt \Big(\frac{2\sigma}{|x-x_{0}|} \Big)^{\delta} w(B)^{-1/p} \\ &\leq \left| C \Big(\frac{\sigma}{|x-x_{0}|} \Big)^{n+\delta} w(B)^{-1/p}. \end{split}$$

Therefore, for all $|x - x_0| > 4\sigma$,

$$(Tf)^*(x) = \sup_{t>0} (E_1(t) + E_2(t) + E_3(t)) \le C \left(\frac{\sigma}{|x-x_0|}\right)^{n+\delta} w(B)^{-1/p}.$$

Combining this, Lemma C and Lemma B, we obtain that

$$L_{2} = \int_{|x-x_{0}| > 4\sigma} \left((Tf)^{*}(x) \right)^{p} w(x) dx \leq C \int_{|x-x_{0}| > 4\sigma} \frac{\sigma^{(n+\delta)p}}{|x-x_{0}|^{(n+\delta)p}} w(B)^{-1} w(x) dx$$
$$\leq Cw(B)^{-1} w(B(x_{0}, 4\sigma)) \leq C,$$

since $(n + \delta)p > nq$. This finishes the proof.

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