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Abstract: The Phase Doppler anemometry (PDA) has been developed to measure simultaneously the velocity and the size of droplets. When the concentration of particles is high, tightly focused beams must be used, as in the Dual Burst PDA. The latter permits an access to the refractive index of the particle, but the effect of wave front curvature of the incident beams becomes evident. In this paper, we introduce a Glare Spot Phase Doppler Anemometry which uses two large beams. The images of the particle formed by the reflected and refracted light, known as glare spots, are separated in space. When a particle passes through the probe volume, the two parts in a signal obtained by a detector in forward direction are then separated in time. If two detectors are used the phase differences and the intensity ratios between two signals, the distance between the reflected and refracted spots can be obtained. These measured values provide information about the particle diameter and its refractive index, as well as its two velocity components. This paper is devoted to the numerical study of such a configuration with two theoretical models: geometrical optics and rigorous electromagnetism solution.

Numerical study of Glare Spot Phase Doppler Anemometry

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ABSTRACT

The Phase Doppler anemometry has (PDA) been developed to measure simultaneously the velocity and the size of droplets. When the concentration of particles is high, tightly focused beams must be used, as in the Dual Burst PDA. The latter permits an access to the refractive index of the particle, but the effect of wave front curvature of the incident beams becomes evident. In this paper, we introduce a Glare Spot Phase Doppler Anemometry which uses two large beams. The images of the particle formed by the reflected and refracted light, known as glare spots, are separated in space. When a particle passes through the probe volume, the two parts in a signal obtained by a detector in forward direction are then separated in time. If two detectors are used the phase differences and the intensity ratios between two signals, the distance between the reflected and refracted spots can be obtained. These measured values provide information about the particle diameter and its refractive index, as well as its two velocity components. This paper is devoted to the numerical study of such a configuration with two theoretical models: geometrical optics and rigorous electromagnetism solution.

1. Introduction

In the research of two-phase flow, the measurement of the velocity and the size of the droplets is essential. The Phase Doppler anemometry (PDA) is one of the most important and the most widely used techniques to reach this end because it gives the correlation between the size and the velocity. But when the concentration of the particles is high, tightly focused beams must

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1 be used to ensure that there is only one particle in the probe volume. The wave front effect can
2 then be considerable. Furthermore, in certain circumstances, as in spray combustion or mixing,
3 we need also to have access to the particle temperature/refractive index or mixing fraction.
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7 With an elastic light scattering based technique such as PDA or rainbow refractometry, the
8 only way to get information about the particle temperature or its material properties is by the
9 measurement of its refractive index. The rainbow refractometry can measure precisely the
10 refractive index (resolution up to $2.0 \cdot 10^{-3}$ for nearly mono-dispersed droplets^[1] and the mean
11 refractive index for polydispersed clouds^{[2][3]}). But it does not give information about the
12 velocity of the particles.
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22 There exist also non elastic light scattering based techniques, such as the Laser Induced
23 Fluorescence (LIF)^{[4][5][6]} or the Raman spectroscopy, which allow to measure the temperature
24 and the particle composition. But these diagnostics are quite difficult to implement. LIF, for
25 example, is based on the dependence of fluorescence spectra on the temperature and the dye
26 must be added in the droplets. The properties of fluorescence depends also strongly on the
27 pressure and composition of the droplets. Furthermore, those techniques require complementary
28 methods to access to the particle size and diameter (such as PDA)^[7].
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40 Then the *in situ* estimation of the refractive index of the scattering particles with the PDA
41 technique is a challenge both to increase the PDA size measurement accuracy and to extract
42 correlated information on the particle velocity, size, temperature and/or composition.
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48 For this purpose, different techniques have been developed^[7]: extended PDA^{[8][9]}, Dual
49 Mode PDA^[10], Dual Burst PDA^{[11][12]}, coupling of PDA and rainbow refractometry^[13], coupling
50 of PDA and LIF^[14] which give velocity, size and refractive index of droplets. Among them, Dual
51 Burst technique is based on the temporal separation between the reflected and the refracted light
52 of an individual particle thanks to a small probe volume formed by two strongly focused laser
53 beams. It imposes that the diameter of the particle must be of the same order or greater than the
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beam waist diameter. But in this case, the measurement quality depends on the longitudinal location of the particles in the probe volume because of the strong curvature of the wave front.

To avoid the sensitivity to the wave front curvature of the laser beam, we introduce an alternative technique, called Glare Spot Phase Doppler Anemometry (GS-PDA). In this configuration, two large beams are used to form the probe volume and two detectors are located in the same plane as the incident beams, similar as in the planar PDA. But the detectors are in the image plane of the particle and in the forward direction at an angle where the intensities of the reflection and the first order refraction are about in the same order. So the reflected and refracted light by the particle is separated in space. When a particle passes through the probe volume, each detector obtains a signal composed of two separate parts, one due to the reflection and the other due to the refraction.

In order to comprehend out the characteristics of the GS-PDA and its sensibilities to particle size, refractive index, as well as velocity, a numerical study of this technique is presented in this paper by two complementary methods: 1). geometrical optics model (GOM) permitting to understand the principle and to treat the problem in a simple way and 2). full electromagnetic model (EIM) thanks to a simulation code of imaging for small particle^[19]. Since the main purpose of the paper is to examine the principle of the PS-PDA and its sensibility to different parameters as particle size, refractive index and velocity, the study will be undertaken step by step for three configurations: 1). one beam - one detector, 2). two beams - one detector and 3). two beams - two detectors. To simplify the analysis, the form of the detectors has not taken into account, i.e. the intensity in the image plane is considered as measured signal.

The paper is organized as follows. A brief presentation of the physical origin and the properties of the glare spots will be given in section 2. Section 3 is devoted to the geometrical optics model of GS-PDA. The full electromagnetic imaging model will be introduced and

applied in section 4 to the study of the measurement of the particle refractive index and diameter.

The conclusions and perspectives are given in the last section.

2. Glare spots

When a particle is illuminated by a beam of light, it scatters light in all space. If the particle is much larger than the wavelength of the incident beam we can observe distinctively in certain angle two bright spots on the particle surface, which are called glare spots^[15]. The formation of the glare spots can be explained in a geometrical optics point of view due to the separation of reflected and refracted light.

We consider a particle of diameter d and refractive index m illuminated by a laser beam. The scattered light is collected by a lens of diameter ϕ and focal distance f located in the direction θ relative to the incident beam (Fig. 1). Two bright points are then formed on the image plan: one (A' in Fig. 1) is created by the refracted light and the other (R') by the reflected light.

When a particle is illuminated by two beams (similar to the Planar Phase Doppler Anemometry, but here the detector is in the image plane of the particle, see Fig. 2), the light due to the reflection and the refraction from the two beams interfere. When a particle passes through the probe volume, the signal obtained by a detector is composed of two Doppler-like bursts, one due to the reflection and the other to the refraction. The detailed analysis of such a configuration will be examined with the geometrical optics model (GOM) and the full electromagnetic imagine model (EIM) in the following sections.

3. Geometrical optics model

The geometrical optics is the most simple and instructive approach to understand the behaviour of GS-PDA. At the first stage, we do not take into consideration the lens.

When a particle of refractive index m and diameter d is illuminated by a plane wave of wavelength λ , the glare spots' positions can be determined by the refracted and reflected rays (Fig. 3). The positions of the glare spots depend not only on the refractive index and the diameter of the particle, but also on the detection angle θ .

The detector being far from the particle, we suppose here for the simplicity that the reflected and refracted rays are parallel.

For convenience, we define the reference ray as the line coming from the source at infinity, passing by the centre of the particle and going out in the direction of emerging rays. The distances between the reference ray and the reflected ray is noted P_r , and respectively P_a for the refracted ray (the subscript r and a represent respectively reflection and refraction). P_r and P_a are then given by:

$$P_r = \frac{d}{2} \cos \frac{\theta}{2} \quad (1)$$

and

$$P_a = \frac{d}{2} \cos \tau \quad (2)$$

Since $2\tau' - 2\tau = \theta$, according to the Snell-Descartes law :

$$m \cos \tau' = \cos \tau \quad (3)$$

P_a can then be written as :

$$P_a = \frac{dm \sin\left(\frac{\theta}{2}\right)}{2\sqrt{1 + m^2 - 2m \cos\left(\frac{\theta}{2}\right)}} \quad (4)$$

It is clear that the distance P_r for the reflected ray does not depends on the refractive index but the distance P_a does ^{[16][17]}. The distance between the two spots $D_s = P_r + P_a$ is proportional to the diameter and depends on the refractive index and the collection angle.

The intensities of the reflected and refracted light can be also calculated according to the geometrical optics. The reflected intensity is given by:

$$I_r = I_0 r_i^2 \frac{d^2}{l^2} \quad (5)$$

and the refracted intensity, by taking into account the divergence factor^[18], is given by:

$$I_a = I_0 (1 - r_i^2)^2 \frac{d^2}{l^2} \frac{m(1 - m \cos(\theta/2))^2}{(m^2 - 1) \cos(\theta/2)(1 + m^2 - 2m \cos(\theta/2))} \quad (6)$$

where I_0 is the intensity of incident light, r_i is the Fresnel reflection factor for parallel ($i = 1$) or perpendicular ($i = 2$) polarization, l the distance between the particle and the observation point.

When the particle is illuminated by two beams with an angle 2α between them, the reflected and refracted light from the two beams interferes. The optical path difference for the reflected and refracted rays are respectively given by:

$$\delta_r = d \left[\sin\left(\frac{\theta - \alpha}{2}\right) - \sin\left(\frac{\theta + \alpha}{2}\right) \right] \quad (7)$$

for reflection and

$$\delta_a = d \left[\sqrt{1 + m^2 - 2m \cos\left(\frac{\theta + \alpha}{2}\right)} - \sqrt{1 + m^2 - 2m \cos\left(\frac{\theta - \alpha}{2}\right)} \right] \quad (8)$$

for refraction.

The corresponding phase differences for reflected and refracted light are then $\phi_r = \frac{2\pi\delta_r}{\lambda}$ and $\phi_a = \frac{2\pi\delta_a}{\lambda}$ respectively. They are both proportional to the particle diameter so they can be rewritten as $\phi_r = A_r d$ and $\phi_a = A_a d$. Note that the factor A_r does not depends on the refractive index but A_a does. The glare spot intensities on the image plane vary according to the particle size and are modulated by the phase differences.

Again to simplify the analysis, we suppose that the reflected intensities from the two beams are equal on the image plane, so does the refracted intensities. Then we can write:

$$I_r = T_r d^2 \cos^2\left(\frac{A_r d}{2}\right) \quad I_a = T_a d^2 \cos^2\left(\frac{A_a d}{2}\right) \quad (9)$$

where T_r and T_a are two coefficients depending on the reflection and refraction Fresnel coefficients and the optical configuration.

The figure 4 shows the behaviour of the reflected and refracted light intensities from a water droplet ($m = 1.33$) illuminated by two planes waves with a half-angle $\alpha = 5^\circ$, the collection angle is $\theta = 66.5^\circ$. The two factors A_r and A_a are respectively equal to $A_r = 0.299\pi$ and $A_a = 0.3533\pi$. The intensity of the spots varies periodically with particle diameter. The corresponding periods of d are then $6.689 \mu\text{m}$ and $5.661 \mu\text{m}$ for reflection and refraction respectively. It is worth to point out that if the intensities due to the two beams are not equal, the periods remain the same but the contrast will be worse.

This simple model of geometrical optics predicts the skeleton of behaviour of glare spots in the GS-PDA. The significant parameters as the lens dimension, the focal length, and the high order refraction effects have not be taken into account. These questions will be examined in the next section with a full electromagnetic model.

4. Electromagnetic imaging model

In this section an electromagnetic imaging model based on full solution of Maxell's equations and Huygens-Fresnel diffraction principle will be used. The model will be presented in the first subsection and it will then be applied to one or two beam(s) configurations of a GS-PDA system.

The electromagnetic imaging model (EIM) is based on the work of Ren et al.^{[19][20][21]}. We consider a particle of diameter d and refractive index m located arbitrarily in a laser beam of wavelength λ . The calculation of the intensity in the image plane consists of three steps. The scattered field on the surface of the lens is firstly calculated by the Lorenz-Mie theory or generalised Lorenz-Mie theory^[22]. The role of the lens is then taken into account by the thin lens transformation^[23]. The intensity on the detector is finally calculated by Huygens-Fresnel integration^{[24][25][26]} of the electromagnetic field on exit surface the lens.

4.1. Glare spots of a particle illuminated by one beam

When a particle is illuminated by one beam, two glare spots are formed on the image plane (see Fig. 1). Their intensities and the distance between them depend on the properties of droplets (size and refractive index) and the optical configuration.

We consider a water droplet illuminated by a plane wave of wavelength $0.488 \mu\text{m}$. The detector is located at $\theta = 66.875^\circ$, in the image plane, 400 mm from the lens. The diameter and the focal length of the collection lens are respectively 15 mm and 200 mm. The magnification of this configuration is then equal to 1. The collection angle of the lens is then $\Delta\theta = 1.074^\circ$.

Firstly, we compare the evolutions of the maximum intensities in the glare spots predicted by the GOM and EIM. Fig. 5 presents the evolution of the intensities as function of the particle diameter for refractive index $m=1.33$. The incident beam is polarized perpendicularly to the scattering plane. We note that the mean intensity of reflection of EIM agrees well with that of GOM but the refraction intensity of EIM is more important than that of GOM. The predictions of EIM present strong fluctuations for both reflection and refraction because in the simulation of EIM all the scattered light have been taken into account while in the GOM only the reflection or the first order refraction is considered. At the chosen angle, the amplitudes of reflected and refracted light dominate and higher orders are relatively small so they contribute mainly to the oscillation. It is worth to note that the relative ratio between reflected and refracted intensities in Figs. 4 and 5 are different because in one beam configuration (Fig. 5), the scattering angle is 66.8° whereas for the two beam configuration (Fig. 4) the collection angles relative to the two incident beams are respectively $66.5+5^\circ$ and $66.5-5^\circ$.

According to the geometrical optics, the reflected and refracted intensity ratio has a simple relation with the refractive index, potentially permitting the measurement of the refractive index. Schaller and al ^{[27][28]} have examined this relation with geometrical optics and Lorenz-Mie theory and by experiments. They found strong fluctuation of the intensity ratio as function of the

1 refractive index and the particle size. We have also examined this relation by the EIM with full
2 GLMT (Generalized Lorenz-Mie Theory). Fig. 6 presents the intensity ratio of reflected and
3 refracted intensities versus the refractive index for a particle of diameter 50 μm or 100 μm ,
4 illuminated by a plane wave ($\lambda = 0.488 \mu\text{m}$). We find that even the tendency of the variation
5 predicted by EIM and GOM is the similar : the intensity ratio increases as function of refractive
6 index, the discrepancy is significant and the result of EIM is somewhat “chaotic”, especially for
7 a small particle.
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10 Then, the distance between the two spots predicted by EIM and by GOM are compared.
11 This distance depends both on the particle diameter and the refractive index. Fig. 7 shows the
12 distance between them for a particle of refractive index 1.33. We find that when the diameter is
13 greater than a critic value, here of 20 μm , the distance is proportional to the particle diameter and
14 the agreement between EIM and GOM is good, the deviation is less than 10%. The simulations
15 have also been done for a magnification of 2 and a collection angle $\Delta\theta = 1.432^\circ$. We found that
16 the under limit of the particle diameter is practically independent of the magnification but
17 depends on the collection angle. When the collection angle $\Delta\theta$ is 1.432° the under limit is
18 pushed to 15 μm for both magnifications of 1 and 2. This effect is due mainly to the diffraction
19 of the lens. When the lens is sufficiently large this effect will be negligible. So the distance
20 between the spots can be used to determine the diameter of the particle.
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23 When a particle passes through the probe volume created by two beams, two Doppler-like
24 bursts are obtained by each detector since in one hand, the scattered light from the beams
25 interfere to create the “fringes”, on the other hand, each detector located on the image plane of
26 the particle “sees” subsequently the reflected and refracted glair spots . Hess *et al* have used the
27 pulse displacement (or time delay) to measure the diameter of the particle and they supposed that
28 the velocity of the particle is perpendicular to the fringes^[28]. We will show in the section 4.3 that
29 two beam – one detector configuration permits to measure simultaneously x component of the
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velocity V_x and the diameter of the particle D . If two detectors are used, the V_x and D can be measured respectively by the frequencies and the phase differences and y component of the velocity can be determined by the time delay between the two maxima of the pedestal which correspond to the two glare spots of reflection and refraction.

4.2. Glare spots of a particle illuminated by two beams

When the particle is illuminated by two beams, the total intensity on the detector is the result of the sum of the electric fields scattered by the particle from the two beams. The two scattered fields then interfere on the detector. We examine in this section the characteristics of the glare spots obtained when a particle is illuminated by two large beams. The configuration considered is shown in Fig. 2. The lens axis makes an angle of $\theta = 66.5^\circ$ relative to the bisect of the two beams of a semi-angle $\alpha = 5^\circ$ between them. The scattering angles are then 61.5° and 71.5° . The particle moves in the direction perpendicularly to the bisect of the two beams.

To show the interference of the light scattered from the two beams, we examine now the intensity in the image plane when the particle is illuminated by only the beam 1 or beam 2, or by the two beams. Fig. 8 shows the intensities on image plane when a water droplet ($m=1.33$) of diameter $50 \mu\text{m}$ is illuminated by one beam or by two beams. The glare spots due to the reflection (right) and refraction(left) are well separated and their positions correspond to the prediction of geometrical optics. The glare spot due to the reflection is at $20.5 \mu\text{m}$ for one beam and at $22.0 \mu\text{m}$ for the other because the two beams are not symmetrical about the detector. The glare spots due to the refraction is at $-24.5 \mu\text{m}$.

We note that the intensities of reflected and refracted glare spots in one beam case (at 71.5° or 61.5°) are very different. The intensity ratios of reflected and refracted glare spots are 0.7 at for 71.5° and 2.48 for 61.5° . The total intensity of the two beams in the glare spot due to the reflection is smaller than that of one beam illumination, the electromagnetic fields from the two

beams are destructive. The total intensity of the glare spot due to the refraction is greater than that one beam illumination, the fields scattered from the two beams are constructive.

In fact, these interferences (construction or destruction) depend also on the particle diameter. Fig. 9 shows the total intensity versus the diameter. The profiles of the intensity evolution of the glare spots are periodical in function of the particle diameter as predicted by geometrical optics (Fig. 4) and the period corresponds well to that of geometrical optics. But it is worth noting that minimum intensity is not zero and increases with the particle diameter. There exists also a high frequency fluctuation. These can be explained by the fact that the size of the lens is finite (15 mm for this simulation) and all orders of scattered light is taken into account in EIM.

4.3. Glare spots Phase Doppler Anemometry

If the detector is small enough relative to the size of the spots, when a particle passes through the probe volume, the two spots' images pass successively in the detector and the detector "sees" the reflection and refraction glare spots one after the other, we then obtain a signal composed of two Doppler-like bursts. A typical GS-PDA signal simulated by EIM is presented in Fig. 10. A particle of $80\text{ }\mu\text{m}$ passes along x axis through the probe volume formed by two plane waves of wavelength $0.488\text{ }\mu\text{m}$ with a half angle of 5° between them (Fig. 2). This signal is similar to the signals obtained by a dual burst PDA in which two tightly focused beams are used^{[11][12]}.

The fringe space in this configuration is equal to $\delta=2.8\text{ }\mu\text{m}$. The velocity of the particle is assumed to be 1.0 m/s and the period of the signal is then $2.8\mu\text{s}$. The two parts due to the reflected and refracted light are well distinct in the signal. By the Fourier transformation we obtained the main frequency of the signal $f=0.357\text{ MHz}$. We can also obtained the pedestals of the signal by filtering the high frequency and extract the distance between the two spots.

It is worth noting that when a particle is illuminated by two beams and one detector is used, there exist specific positions for which the reflection and/or refraction spots disappear as shown in figures 4, 8 and 9. In this case no distance between the two spots can be extracted. But when the particle moves in the probe volume we can always obtain a Doppler-like signal to extract the pedestals even for some positions the intensity is destructives, for example the point A for refraction and B for reflection in Fig. 10.

We investigate in the following the distance between the two maxima of pedestal as function of the particle size. For the simulation presented in Fig. 10, the diameter is 80μm and the refractive index $m=1.33$. The detector is located at $\theta = 67.574^\circ$ so that the reflected and refracted light is of the same order of intensity.

Firstly, we note that the x component of velocity can be determined by the frequency of the signals:

$$V_x = f \delta = 0.357 \times 10^6 \times 2.8 \times 10^{-6} = 1.0 \text{ m/s} \quad (10)$$

The time delay between the two maxima in the pedestal (dashed line in Fig. 10) T_s is related then to the displacement of the particle D_x and the distance between the spots D_s by:

$$D_s = g D_x \cos \theta = g V_x T_s \cos \theta \quad (11)$$

where g is the magnification of the used optics. Since the particle diameter d is related to $D_s = P_a + P_r$ (see Eqs. (1) and (4)), the particle diameter can then be obtained from the time delay between the two maxima in the pedestal T_s .

For the case shown in Fig. 10, the time delay between the two maxima in the pedestal is $T_s = 192 \text{ μs}$ which corresponds to a displacement of $D_x = 192 \text{ μm}$. Considering the angular position of the detector, the distance of between the two spots $D_s = 72.6 \text{ μm}$ can then be deduced and it is in agreement with that predicted in Fig. 7.

Then the simulation has been done for a water droplet of diameter from 30 μm to 120 μm and the result is presented in Fig. 11. We found that the relation D_s-d is linear and the agreement

between EIM and GOM is very good. Therefore the distance D_s permits to determine the diameter of the particle.

On the other hand, since both the position of the refraction spot and the ratio of the two maximum intensities depend also on the refractive index, we would expect to obtain information about refractive index from these values if the diameter is known. But both GOM and EIM show that the distances D_s depends little on the refractive index. Furthermore, the intensity ratio obtained from EIM fluctuates as function of the refractive index and the fluctuation is too strong to be used to determine the refractive index of an individual particle. We will show in the following section that the two beams – two detectors configuration permits to get better information of the refractive index.

Now we add a second lens and a detector same as the first but with a different collection angle. The lens and the detectors are all in the same plane as the incident beams. In this configuration, when a particle passes through the probe volume, two signals will be obtained as in the dual burst PDA^{[11][12]}, each of them is composed of two parts: one due to reflection and the other due to refraction. Besides the distances D_s and intensity ratios as described above, we can also extract the two phase differences between the two signals separately for reflection and refraction parts. Fig. 12 shows two signals simulated by EIM. The two detectors are located in 65.426° and 67.574° relative to the bisect of the two beams. Firstly, we remark that the distances between the two intensity maxima are not the same because the two detectors are not symmetric about the particle trajectory. Note that, in principle, the relation between the phase difference and the particle diameter, and that between the frequency and the transverse velocity component remain the same as in Planar PDA if the particle is located in the object plane.

The phase differences extracted from the signals are shown in Fig. 13. We find that the phase differences are relatively stable around the maximum intensity (it should be noted that the phase difference for signals of very weak intensity has no sense since the noise can be more

important than the signals). The average phase differences for reflection and refraction are respectively 43° and 24°. The same simulation has been done for different particle size and refractive index. The results are presented in Figs.14 and 15. The evolution tendency of the phase differences predicted by EIM agrees with those predicted by geometrical optics. Larger is the diameter, better is the agreement. But the evolution of phase differences with the diameter and the refractive index predicted by EIM are not linear(chaotic in certain sense) as predicted by GOM. Nevertheless we find that when the particle is very large (larger than 100 μm), the phase differences can be used to determined the diameter with the precision about 30 % (still too important, see comments in the last section).

Nevertheless, the *mean* value of the ratio of the two phase differences obtained by EIM is relatively stable. Left part of the Fig. 16 shows the ratio predicted by EIM and GOM for a particle of diameter from 25 μm to 150 μm and refractive index 1.33. The mean value of the ratio predicted by EIM is in agreement with that of GOM. Therefore the *mean value of the* refractive index can be determined according to the relation of OGM (right part of the figure).

Once the particle diameter is determined by the two phase differences, the two time delays Δ_i between the two maxima of the pedestals obtained from the signals of the two detectors can be used to determine y component of the velocity in the plane of the beams according to:

$$V_z = \frac{D_s}{\Delta_i \sin \theta_i} - \frac{V_x}{\tan \theta_i} \quad (12)$$

where $i=1, 2$ corresponds to the detector 1 or 2, $D_s = P_r + P_a$ is the distance between the two glare spots predicted by GOM. V_x is the x component of the velocity perpendicular to the fringes determined by the frequency of the signal.

5. Conclusions and perspectives

A new technique, Glare Spots Phase Doppler Anemometry, is proposed and studied by using the geometrical optics model and the full electromagnetic imaging model. The simulations

by EIM agree in general with those predicted by GOM and shows that the GS-PDA provides
rich information: two distances and two intensity ratios between the reflection and refraction
Doppler-like bursts, two phase differences as well as the frequency of the signals. The particle
diameter and two components of velocity can be determined by the time delays between the two
bursts, the phase differences and the frequency of the two signals. The refractive index can be
evaluated from the mean value of the phase difference ratios of the two signals.

Nevertheless, it should be pointed out that the relations of the refractive index with the
phase difference and the intensity ratio obtained by the simulation for individual particles are
somewhat “chaotic” (even the mean value has sense). Further study must be carried out to
determine if it is the intrinsic property of the glare spots or due to the diffraction effect since the
diameter of the lens is only 15 mm in the simulation because of numerical difficulty. Two
possible ways could be involved : check by experiment with realistic lens and detectors or
simulation by improved code permitting to take into account the effect of the size of the lens, the
form of the detectors and others parameters.

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France.

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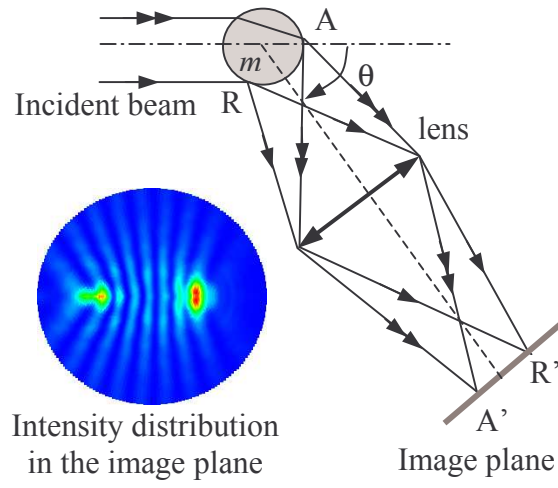


Fig. 1. Off-axis imaging. A' and R' are respectively the glair spots of refraction and reflection (colour on line).

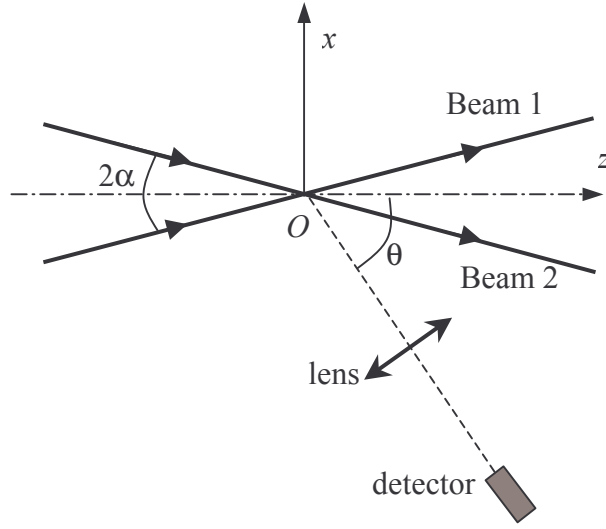


Fig. 2. Schema of a planar GS- PDA. The two incident beams, the lens and the detector are all in the same plane. The detector is located in the image plane of the particle at an angle where the reflection and first order refraction light intensities are in the same order.

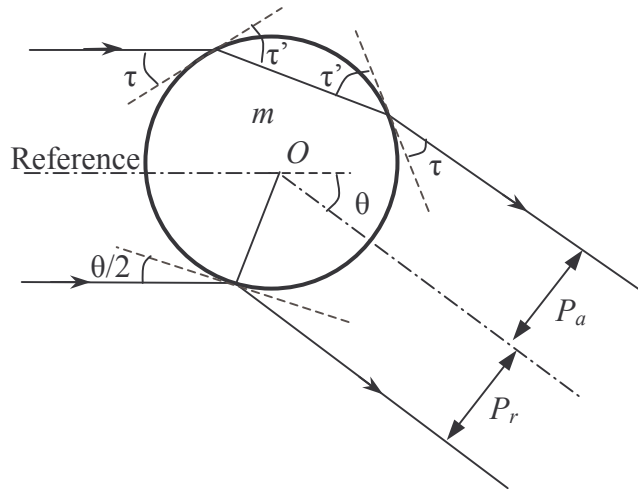


Fig. 3. Optical paths of a reflection ray and a refraction ray.

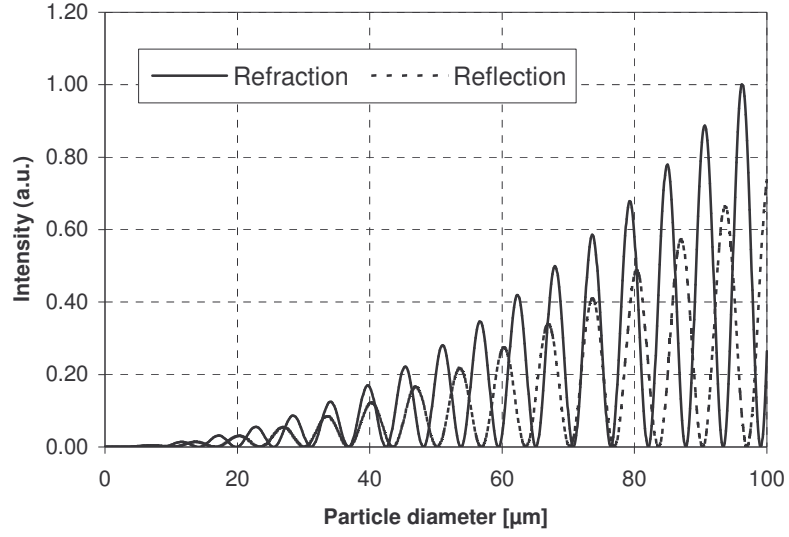


Fig. 4. Glare spot intensities as function of the particle diameter predicted by geometrical optics: a) for reflected light, b) for refracted light (the ratio of the reflected and the refracted intensities has been justified according to the Debye series calculation).

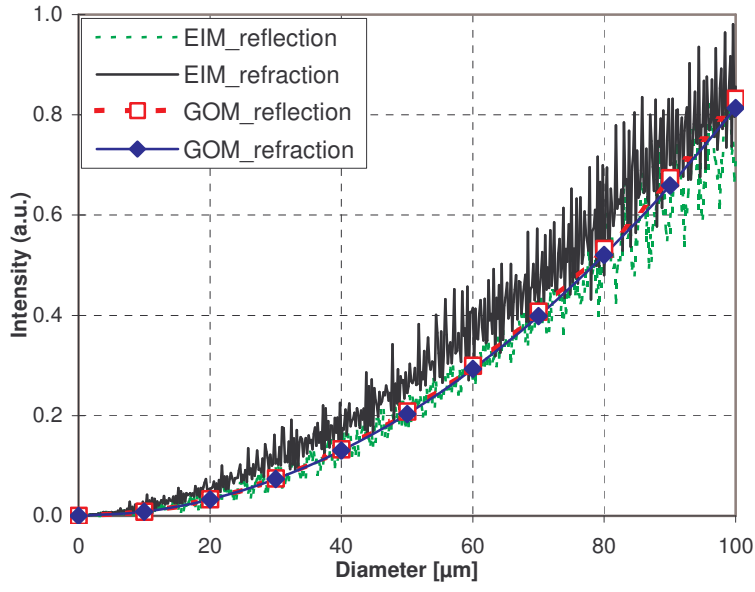


Fig. 5. Evolution of intensities of the glare spots versus the diameter of a water droplet ($m=1.33$) illuminated by a plane wave of wavelength $\lambda=0.488 \mu\text{m}$ oriented at 66.785° relative to the lens axis (colour on line).

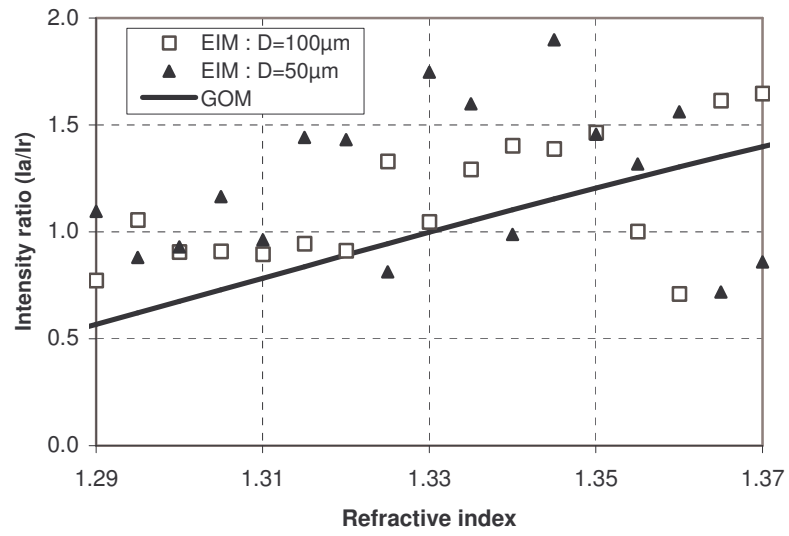


Fig. 6. Intensity ratios of the spots as function of the refractive index.

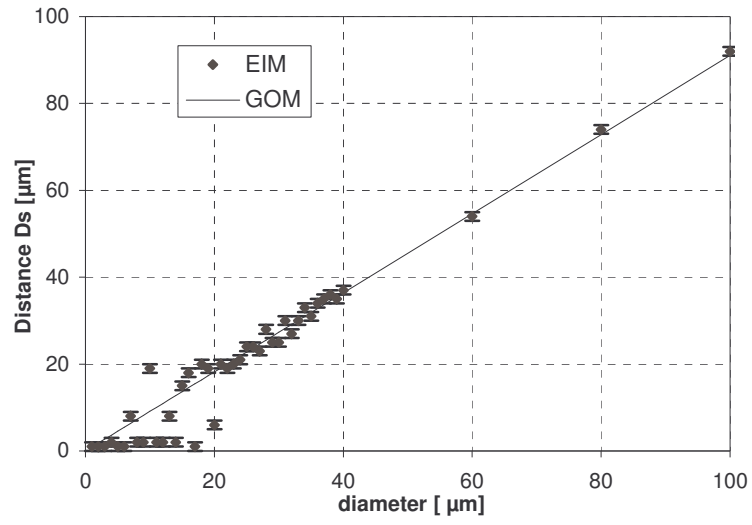


Fig. 7. Evolution of the spots' distance as function of the particle diameter. The refractive index is 1.33.

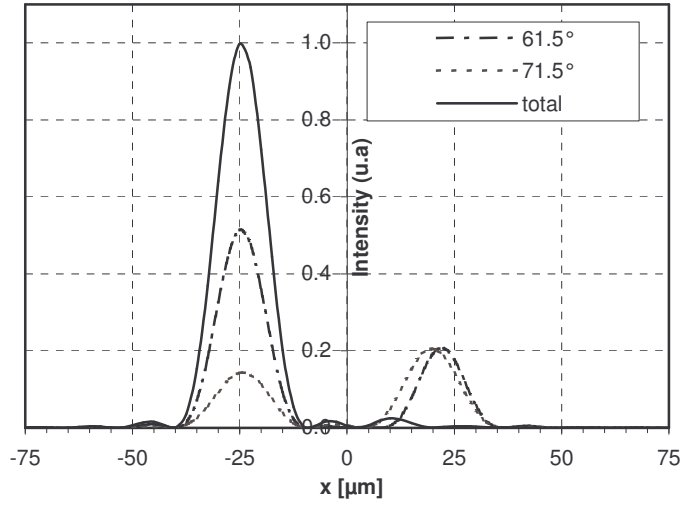


Fig. 8. Intensities on the image plane when a water droplet ($m=1.33$, $d=50\text{ }\mu\text{m}$) is illuminated by one beam at $\theta = 61.5^\circ$ (dashed line) or $\theta = 71.5^\circ$ (long dashed line) or by two beams (solid line).

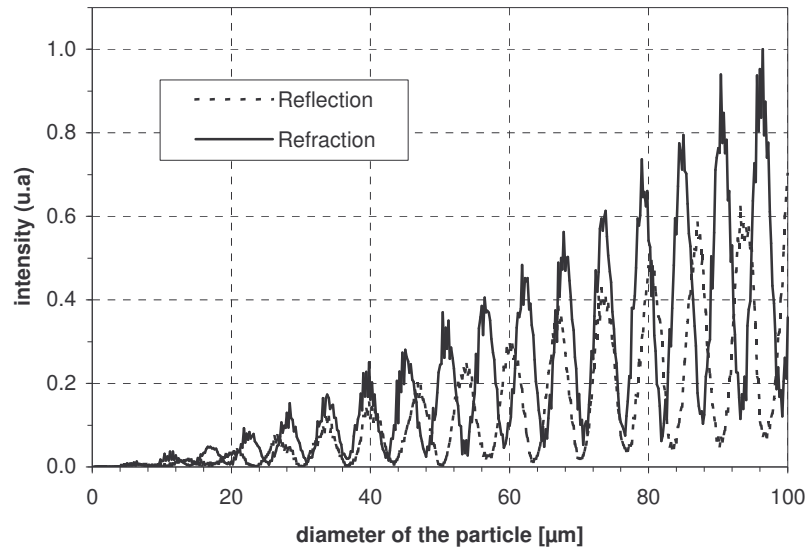


Fig. 9. Evolution of the glare spot intensities as function of the diameter of a water droplet illuminated by *two plane waves*.

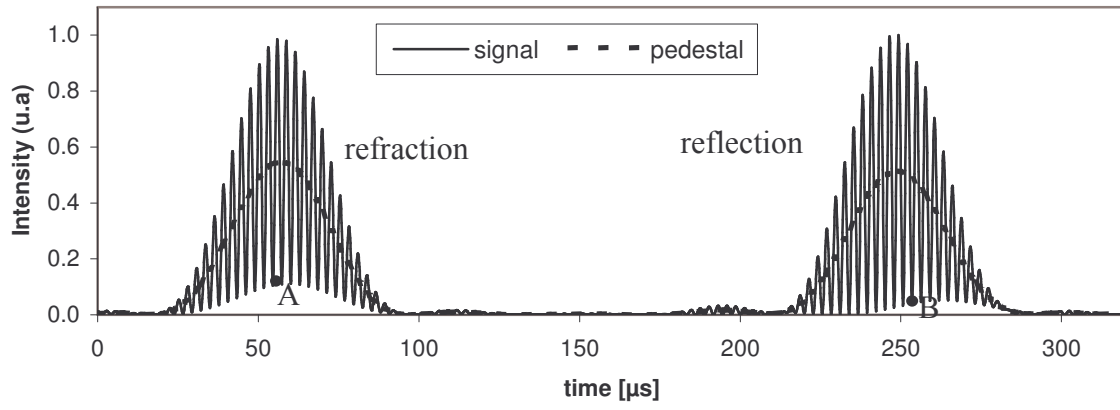


Fig. 10. Typical signal of GS-PDA.

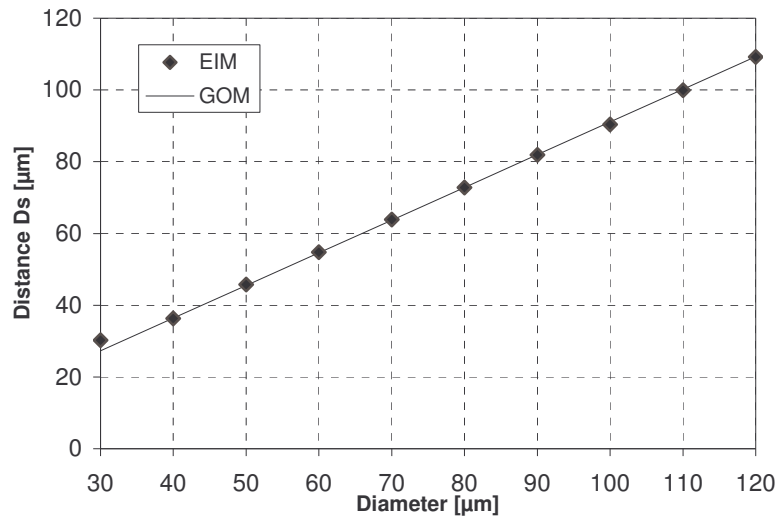


Fig. 11. Distance between the spots as function of diameter of a water droplet deduced from a GS-PDA signal.

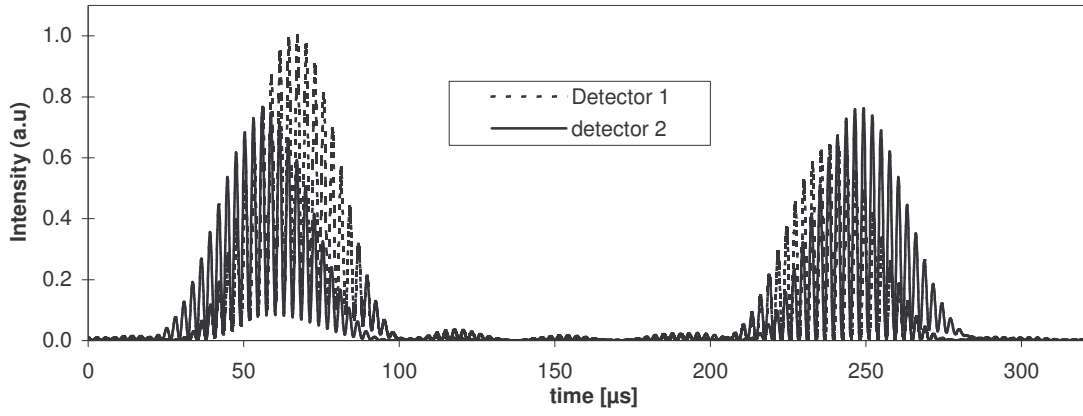


Fig. 12. Signals of GS-PDA obtained by two detectors located respectively at $\theta = 65.426^\circ$ and 67.574° with a half angle of 5° . A water droplet ($m=1.33$, $d=80 \mu\text{m}$) is illuminated by two plane wave.

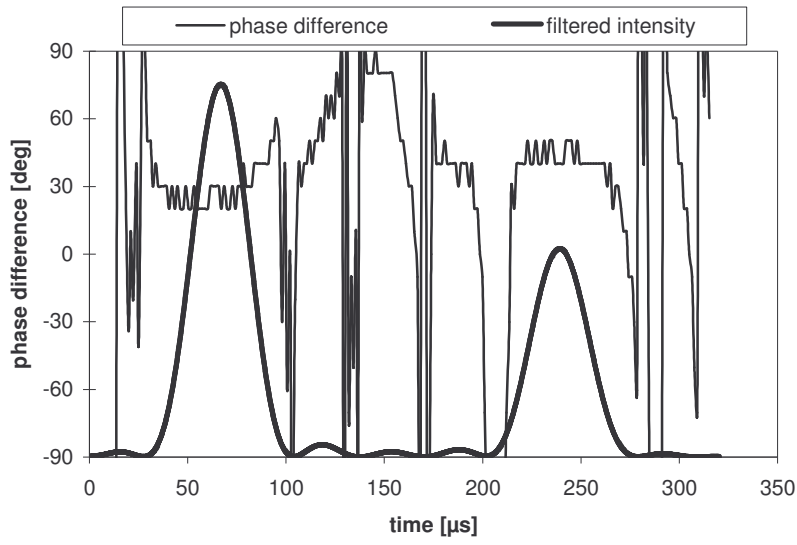


Fig. 13. Phase difference between the two signals and their pedestal obtained from the signals in Fig. 12.

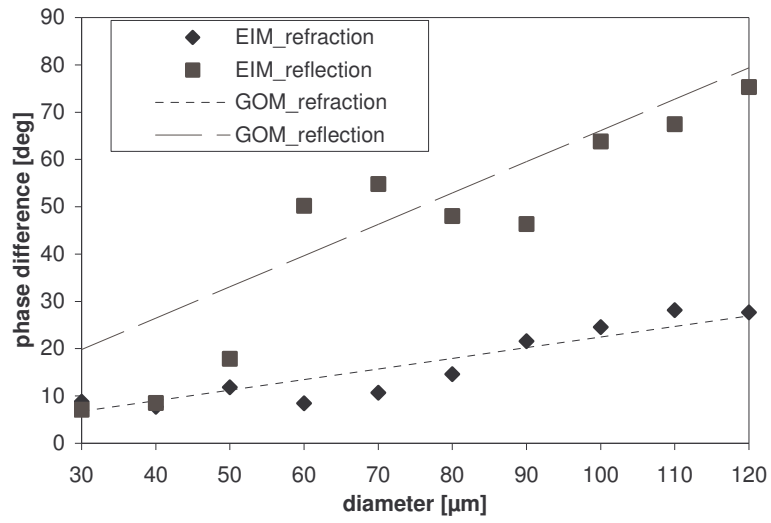


Fig. 14. Phase differences versus diameter of water droplet.

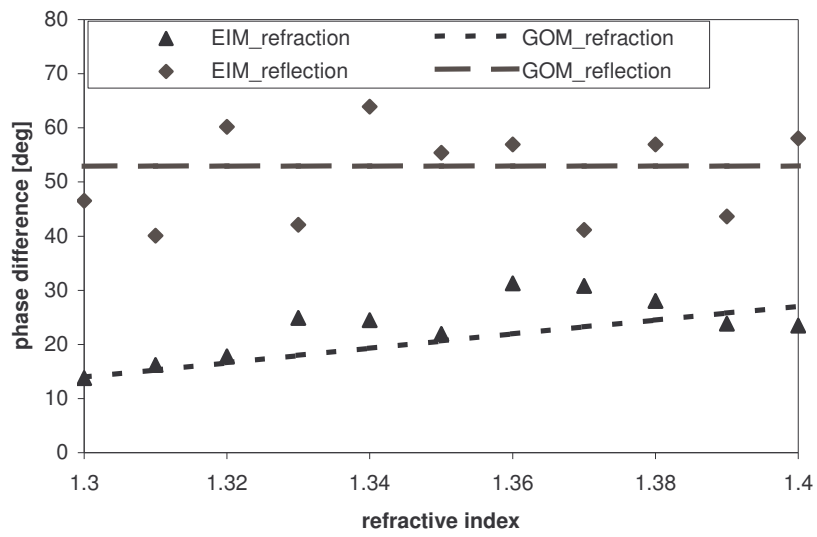


Fig. 15. Phase difference versus refractive index for a particle of diameter 80 μm .

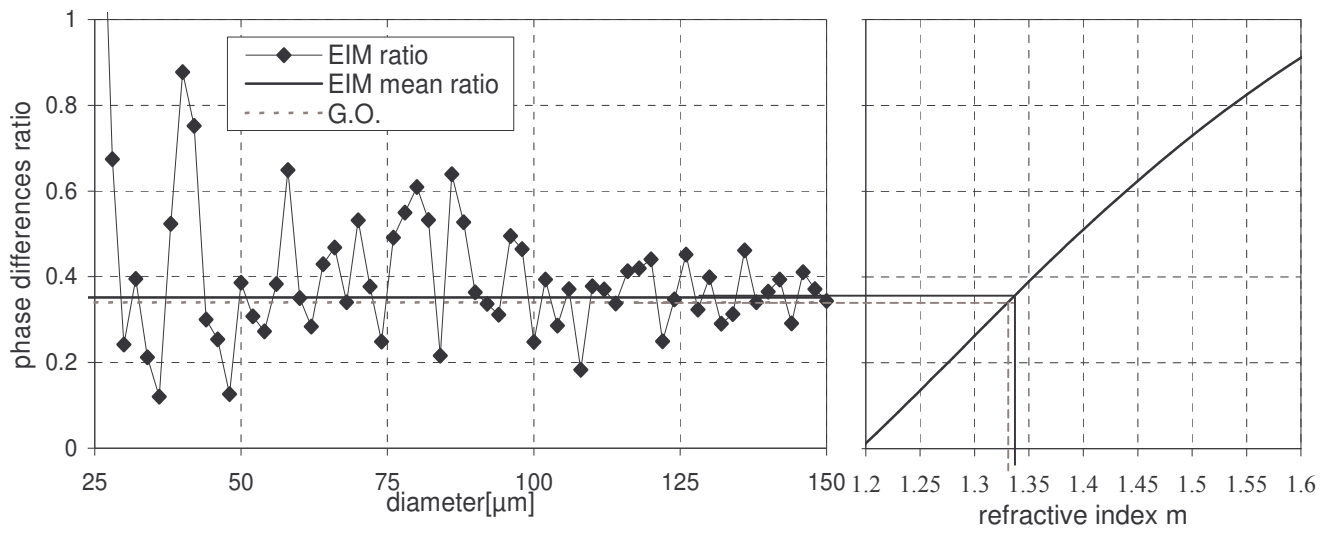


Fig. 16. Ratio of the two phase differences versus refractive index for a particle refractive index 1.33.