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Hydrological effects on gravity and correlations between gravitational variations and level of the Alzette River at the station of Walferdange, Luxembourg

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Abstract

The gravitational effects of water storage variations driven by local precipitation events are modeled for the Walferdange Underground Laboratory for Geodynamics (WULG) in the Grand Duchy of Luxembourg. A modified mass continuity model is implemented, which uses rainfall data from Walferdange as input. In the absence of soil moisture and groundwater level information, the model is empirically parameterized. Model outputs are compared to the gravity time series registered with the Observatory Superconducting Gravimeter CT040 located in the WULG. We find that the model explains 77% of the gravity residuals. In addition, a statistical analysis is carried out to determine the relationship between precipitation, gravity variations and water level changes in the nearby Alzette River. A time delay of 88 ± 34 min between the maximum variation rates of the water level and gravity signal has been calculated. The signals have an admittance of 45 ± 5 cm μGal^{-1} .

Keywords: Gravity and Superconducting gravimeter, Hydrology, Groundwater, Precipitation

1. Introduction

The effect of water storage variations on gravity has become an important issue for improving the investigation of gravity measurements. Specifically, local water storage variations within 10 km of a station alter the local mass field and thus can significantly affect gravity observations [Van Camp *et al.*, 2006]. Various authors have analyzed the effects of parameters such as local precipitation, soil moisture and groundwater storage on gravity observations [e.g. Meurers *et al.*, 2007; Van Camp *et al.*, 2006; Kroner and Jahr, 2006; Bower and Courtier, 1998; Boy and Hinderer, 2006; Kroner, 2001; Peter *et al.*, 1995; Harnish and Harnish, 2002; Crossley and Xu, 1998; Imanishi *et al.*, 2004; Llubes *et al.*, 2004; Delcourt-Honorez, 1989]. Other authors [van Dam *et al.*, 2001; Crossley *et al.*, 2005; Hinderer *et al.*, 2006; Naujoks *et al.*, 2007] have investigated the consequences of longer wavelength (several 100 km) water storage components on the gravity signal. At these wavelengths, the displacement of the Earth's surface due to the excess mass is also important.

The response of local gravity to local changes in water storage is site dependent [Hokkanen *et al.*, 2005]. Techniques and schemes for mitigating the signal in gravity data at one location are usually not entirely applicable to another location. In this study, we look at the physical relationship between water storage variations driven by local precipitation events and local gravity changes at Walferdange, in the Grand Duchy of Luxembourg. For the gravity observations, we use the data collected by the superconducting gravimeter CT040 (SG) located in the Walferdange Underground Laboratory for Geodynamics (WULG).

The present paper builds on the results of previous authors who have also investigated the correspondence between gravity and local water storage variations. A new and simple scheme is provided to remove the effects of precipitation events in the gravity observations. A synthesis of the different approaches that can be found in recent literature is proposed. The model is then extended to the Walferdange case by including a new parameter to take into account the seasonal variability of groundwater behavior.

An empirical model is developed, based on a modification of the basic mass continuity model, to estimate water storage variations due to local precipitation [Meurers *et al.*, 2007]. This model is used to correct the gravity time series residuals of the superconducting gravimeter. A statistical analysis is then undertaken to determine the correlation between gravity variations registered by the SG and water level changes in the nearby Alzette River.

In Sections 2 and 3, a synthetic description of the hydrological recharge and discharge processes related to the water cycle is given. Here the hypotheses leading to a tank model representation are described. In order to underline the dependence of the model parameters on hydrological parameters, such as porosity and hydraulic conductivity, as indicated by Van Camp *et al.* [2006], the steps leading to the fundamental model equations are detailed. This is achieved by combining the mass continuity equation and Darcy's law [Fetter, 2001; Roche, 1963], which describes the flow of a fluid through a porous medium.

In Section 4, modeled gravity variations are estimated using an admittance factor between the gravity change and the precipitation height. This is done by decomposing the area above the gravimeter in a discrete number of prismatic elements (the precipitation

height corresponding to the element thickness). The gravity effect on the SG is calculated using Newton's law of universal attraction. A further hypothesis is developed, regarding the seasonal variability of a model parameter, which represents the gravity recovery rate after a rainfall and its physical significance is discussed.

In Section 5, the model outputs for a set of empirically evaluated parameters are presented and discussed.

In Section 6, a statistical analysis is presented to determine the correlation between the gravity signal variations registered by the SG and the water level of the Alzette River. The gravity variation due to the precipitation should appear before the change in river level. The gravity observation should also contain information on the degree of soil saturation. Understanding the relationship and temporal dependence between the observed precipitation and the gravity changes might improve our ability to predict of extreme events like flooding. We estimate the mean time delay between the maximum variation of the gravity signal and the Alzette water level as a function of a set of sampled showers. Finally, the correlation between the gravity variation and water level variation is calculated and discussed.

2. Hydrological recharge and discharge processes

Variations in the amount of water stored in the ground exert an effect on the gravity signal through essentially two different mechanisms: 1) the Newtonian attraction exerted by the water mass and 2) the ground deformation resulting from the water load and the associated mass redistribution. The effect of the ground deformation on gravity is a long wavelength effect (> 50 km).

Groundwater flow is possible because Earth materials are usually not impermeable, but characterized by the presence of voids (pores) included in a solid matrix. The groundwater and soil moisture occur in the voids [Fetter, 2001]. For a control volume V_t of a given material, the porosity P is defined as the ratio between the volume V_v of the voids and the control volume V_t :

$$P = \frac{V_v}{V_t} \quad (1)$$

If the voids are completely/partially filled with water, the medium is defined as saturated/unsaturated.

The mechanical energy per unit weight possessed by a fluid is given by Bernoulli's law:

$$h = \frac{v^2}{2 \cdot g} + z + \frac{p}{\rho \cdot g} \quad (2)$$

where v is the fluid velocity, g is the gravity acceleration, z is the elevation of the fluid center of mass with respect to a reference level, p is the pressure, ρ is the fluid density and h , expressed in units of length, is defined as the hydraulic head. The kinetic term $v^2/2g$ is generally negligible in problems related to groundwater flow [Fetter, 2001].

The flow through a porous medium is described by Darcy's law:

$$q = K \cdot \left(\frac{dh}{dl} \right) \quad (3)$$

115

116 where q (length/time) is the flow per surface unit or specific flow, h is the hydraulic head.
 117 and l is the displacement in the flow direction. The proportionality coefficient (K)
 118 between the specific flow and the hydraulic gradient (dh/dl) quantifies the capacity of
 119 fluids to move through porous media. The coefficient is known as the hydraulic
 120 conductivity or permeability coefficient. It is important to underline that hydraulic
 121 conductivity is a function of the characteristics of both the fluid and the porous medium
 122 [Fetter, 2001]. The general formulation of problems related to groundwater flow is based
 123 on a combination of Darcy's law and continuity equations with reference to control
 124 volumes [Anderson, 2007].

125 For unconfined aquifers the water table represents the interface between the saturated
 126 zone (below) and the unsaturated zone (above) [Milly and Shmakin, 2002]. The water
 127 volume contained in the saturated zone represents the groundwater reservoir. When
 128 precipitation occurs, part of the fallen volume of water is retained by the vegetation
 129 canopy or human artifacts, another part is subject to direct surface runoff or other
 130 subsurface flow, and one part infiltrates through the unsaturated zone to increase the
 131 groundwater reservoir storage. This infiltration represents the recharge process, with
 132 augmentation of the water storage. After the precipitation event and the recharge process,
 133 the stored water will decrease through different mechanisms, such as evaporation and
 134 evapotranspiration from the unsaturated zone or efflux through filtration from the
 135 saturated zone, to reach draining flows. This represents the discharge process, with
 136 consequent diminution of the water storage. In conclusion, both the recharge and the

discharge must be seen as the superposition of different processes with different characteristic time scales. The number of parameters needed to describe those processes is very high, and are subject to strong time and space variability. Different models have been proposed to analyze the complex processes related to water (and energy) storage [Manabe, 1969; Milly and Shmakin, 2002, Iffly *et al.*, 2004]. In the absence of direct information about soil moisture and groundwater level, the development of an empirical local model is required to describe the water storage variations. Important simplifications of the processes involved are necessary.

3. The Tank Model

The fraction of precipitation that infiltrates the soil, percolates vertically through the unsaturated zone to reach the saturated zone below the water table. This process increases the groundwater reservoir storage. In the case of unconfined aquifers (presence of one impermeable underlying layer), the efflux from the saturated zone can essentially be described as horizontal filtration through a porous medium. The tank model [Roche, 1963] represents a strong simplification: the discharge process is reduced to the efflux of water contained in a tank (and not distributed in a solid matrix) through a porous plug. A simple illustration of the tank model is given in Figure 1.

Figure 1

The relation between the flux from the porous plug and the hydraulic head is expressed by Darcy's law (Eq. (3)) in the following form [Roche, 1963]:

$$Q = K \cdot \frac{s}{L} \cdot h \quad (4)$$

162

163 where Q represents the outgoing flux (volume/time), h is the water height (corresponding
 164 to the hydraulic head), s and L are the cross section and length of the plug, respectively
 165 (Figure 1). The parameters of the model have to be determined empirically. The outgoing
 166 flux is related to the water level through the hydraulic conductivity, which depends both
 167 on hydrological soil parameters, such as the medium's porosity, and on fluid
 168 characteristics such as viscosity. A hydrological volume balance is calculated, with
 169 reference to the elementary time interval dt :

170

$$S \cdot dh(t) = -Q(t) \cdot dt + S \cdot r(t) \cdot dt \quad (5)$$

172

173 where $r(t)$ represents the rainfall rate (length/time) and S the tank surface. Defining the
 174 level decrease rate, c (1/time), by:

175

$$c = \frac{K \cdot s}{S \cdot L} \quad (6)$$

177

178 Substituting the definition of Q (Eq. (4)) into Eq. (5), the differential equation describing
 179 the water level variation becomes:

180

$$\frac{dh(t)}{dt} = -c \cdot h(t) + r(t) \quad (7)$$

182

Equation (7) relates the water level, corresponding to the hydraulic head in this simplified representation, to the rainfall. In this equation, $r(t)$ is not a continuous function of time and the integration has to be performed numerically. In absence of precipitation, the equation can be analytically integrated leading to an exponential decay expression describing the draining of a tank through a porous plug.

4. Relation between precipitation height and gravity variation

The WULG is located underground in a derelict gypsum mine, at latitude 40.6700°N , longitude 6.1500°E . The laboratory housing the Observatory Superconducting Gravimeter *CT040* lies at the end of an 800 m long tunnel cut into the side of a ridge. The gravity sensor is 295 m above sea level and about 80 m below the local surface. The area above the gravimeter has significant topographic slope gradients (See Figure 2), and is covered by uneven vegetation. The closest human artifacts (i.e. roads and buildings) are located at about 500 m from the gravimeter. All of these factors contribute to making the WULG seismically quiet.

In order to express Eq.(7) in terms of gravity variations, we estimate the admittance between the gravity change and the precipitation height. The gravity effect of a weathered area can be calculated using digital elevation models (DEM): the weathered ground layer is discretized in prismatic elements [Van Camp *et al.*, 2006; Banerjee and Das Gupta, 1977; Talwani, 1973]. The gravity effect Δg_e of a single element is given by:

$$\Delta g_e = G \cdot \rho_{wet} \cdot \int_{x_1}^{x_2} \int_{y_1}^{y_2} \int_{z_1}^{z_2} \frac{z}{(x^2 + y^2 + z^2)^{3/2}} \cdot dx \cdot dy \cdot dz \quad (8)$$

where G is the Newtonian universal constant of gravity ($6.67428 \cdot 10^{-11} \text{ m}^3 \text{ kg}^{-1} \text{ s}^{-2}$) and ρ_{wet} is the density of the wet component of the weathered zone (mass/volume). The integration is carried out over the prism volume, delimited by the coordinates (x_1, y_1, z_1) and (x_2, y_2, z_2) in a Cartesian frame centered at the gravity station.

In absence of accurate information about ρ_{wet} in the weathered zone above the WULG, the calculation is made directly on prismatic water elements having the precipitation height as thickness. This simplification is adopted because the thickness of the weathered zone appears to be significantly smaller than the distance between the gravity sensor and the surface.

The 2000 m x 2000 m zone above the gravimeter is divided into 10000 prismatic elements of size 20 m x 20 m. The map projection of the discretized layer defines a 100 x 100 grid. The square elements of the grid are identified by the indices i and j corresponding to the coordinates in meters $x_i=20i$, $y_j=20j$ and elevation z_{ij} in a local Cartesian frame (the coordinates are referred to the element central points). The gravimeter is located at (x_g, y_g, z_g) . The gravity effect Δg_{ij} of each prism (m s^{-2}) at the gravimeter location is calculated with:

$$\Delta g_{ij} = G \cdot \Delta S \cdot \rho \cdot l_e \cdot \frac{z_{ij}}{((x_i - x_g)^2 + (y_j - y_g)^2 + (z_{ij} - z_g)^2)^{3/2}} \quad (9)$$

where ΔS is the surface area of the element (400 m^2), ρ represents the water density (1000 kg m^{-3}) and l_e is the element thickness (m). The total gravity effect Δg of the layer (m s^{-2}) on the gravity sensor is given by:

228

$$\Delta g = \sum_{i=1}^{100} \sum_{j=1}^{100} G \cdot \Delta S \cdot \rho \cdot l_e \cdot \frac{z_{ij}}{((x_i - x_g)^2 + (y_j - y_g)^2 + (z_{ij} - z_g)^2)^{3/2}} \quad (10)$$

230

231 The admittance, α , between a gravity change (μGal) and a precipitation height (m) is
 232 directly calculated with Eq.(10), by defining $l_e=1$ m and expressing the gravity in μGal .
 233 The result is $\alpha=36 \mu\text{Gal m}^{-1}$. For reference, the value given by the Bouguer infinite plate
 234 model is $42 \mu\text{Gal m}^{-1}$. This value is generally adopted for almost flat surfaces. The
 235 significant height differences of the area above the WULG justify the calculation made
 236 on the actual topography.

237

238 **Figure 2**

239

240

241 With the calculation of the admittance between the hydrological component of the
 242 gravity variation and the precipitation height, the equation describing the gravity changes
 243 can be obtained by writing the continuity equation (7) in terms of gravity. This is done by
 244 multiplying both sides of Eq.(7) by α , setting $\Delta g=\alpha \cdot h$, $c=\gamma$, and changing the sign of the
 245 term $\alpha \cdot r$. This last step is required by the fact that the gravimeter is located in an
 246 underground laboratory: an increase of the stored water level above the gravimeter will
 247 produce a reduction of the gravity signal. We thus obtain:

248

$$249 \quad \frac{d(\Delta g)}{dt} = -\gamma(t) \cdot \Delta g - \alpha \cdot r(t) \quad (11)$$

250

where Δg represents the hydrological part of the gravity variation and γ (corresponding to c in Eq.(7)) is here defined as the gravity recovery rate after rainfall (1/time). The coefficient γ is related to hydrological parameters like hydraulic conductivity and porosity, in addition to the geometrical parameters depending on the local configuration. The inverse of γ , defined as τ , has the dimensions of a time and it characterizes the duration of the discharge phase. Equation (11) has been previously implemented by *Imanishi et al.* [2004] to calculate the gravity effects of underground water in Matsushiro (Japan). The basic tank model is modified using the hypothesis of a seasonal variability for the parameter γ . The seasonality is included to account for the annual changes in the meteorological and hydro geological conditions (i.e. evapotranspiration and hydraulic conductivity). This variability is expressed in the form:

$$\gamma(t) = \gamma_0 \cdot (1 + A \cdot \cos(\frac{2 \cdot \pi}{T} \cdot t + \varphi)) \quad (12)$$

where γ_0 is the mean value of γ in the period T (one year). The values of A (amplitude) and φ (phase) are evaluated empirically (see Section 5.2.). Finally, the equation for the gravity variation becomes:

$$\frac{d(\Delta g)}{dt} = -\gamma_0 \cdot (1 + A \cdot \cos(\frac{2 \cdot \pi}{T} \cdot t + \varphi)) \cdot \Delta g - \alpha \cdot r(t) \quad (13)$$

As already pointed out, the rainfall rate $r(t)$ is not defined by a continuous function of time but by a discrete set of values. To allow us to solve the equation using finite

difference algorithms, Eq. (13) has to be rewritten as a discrete expression. We define the terms of a progression:

$$\Delta g(i+1) = (1 - \gamma(i)) \cdot \Delta g(i) - \alpha \cdot r(i) \quad (14)$$

where i represents the non dimensional time indices, $\gamma(i)$ is the non dimensional recovery rate and $r(i)$ are the precipitation heights. The time variable t in Eq. (13) is given by $t = \Delta t \cdot i$, where Δt is the reference time interval (1 hour in the hourly implementation), and the recovery rate function $\gamma(i)$ is formally identical to $\gamma(t)$.

5. Data and Results

5. 1. Input Data

The available data for the model implementation are the local rainfall time series from the 1 January 2003 to the 1 January 2009. Two different data sources were available: the pluviometer of the Walferdange Meteorological Station and the WULG pluviometer (Figure 2). Both the pluviometers provide the rainfall heights (mm) at one-minute sampling intervals. A comparison between the two series indicates a qualitative similarity. A significant quantitative difference exists as well. The cumulative precipitation provided by the WULG is approximately 10% greater than the cumulative precipitation provided by the Walferdange Meteorological Station. This difference will be used to assess the uncertainties of the rainfall data. The time series provided by the WULG is displayed in Figure (3).

Figure 3

5. 2. Model parameterization for the hourly implementation

The value of the admittance, $\alpha=36 \mu\text{Gal m}^{-1}$ derived in Section 4, has been used for the model parameterization in Eq.(13). The parameters γ_0 , A and φ in the same equation are obtained empirically by a least square adjustment. The gravity variations observed with the SG are compared with the values predicted by Eq.(13) using the rainfall data, for the period 19 December 2003 to 1 December 2004. The parameters are determined from the data in order to minimize the residuals. The values of $A=0.6$, $\varphi= 2\pi/5$, and $\gamma_0=0.0011 \text{ hours}^{-1}$ provide the best results. For reference, *Meurers* (2007) adopted the value $\tau=720 \text{ hours}$ for the discharge time parameter at the gravity station in Vienna, which corresponds to $\gamma_0=0.0015 \text{ hours}^{-1}$ (see Section 4).

5. 3. Results

The rainfall time series provided by the WULG, integrated to hourly values and referred to the mid-hour points, are used as model input. The rainfall heights are expressed in meters. An uncertainty of 10% on the rainfall data, as determined by the quantitative difference in the rainfall data from two closely space pluviometers (see Section 5.1) is assigned.

The modeled hourly time series of the gravity (μGal) is compared with the observed hourly time series obtained from the SG for the period 19 December 2003 to 1 January 2009. The instrumental drift of the SG is modeled using an exponential function [*Van*

Camp and Francis, 2006]. The modeled gravity (green), the observed gravity (blue) and their difference (red), expressed in μGal , are displayed in Figure 4.a.

Figure 4

The RMS of the uncorrected gravity time series is $1.89 \mu\text{Gal}$. After removing the modeled hydrological effects from the observations, an RMS value of $0.89 \mu\text{Gal}$ is obtained, which corresponds to a 77% signal reduction.

The gravity changes after a high intensity precipitation event (Figure 4.b) show that the gravity recovery process has a significantly longer duration than the gravity decrease process. The tank model describes only the recovery process. The gravity decrease process is considered as instantaneous. When a precipitation occurs, the rainfall height is simply added to the current water level, producing an instantaneous gravity change. A higher time sampling of the data is therefore not necessary.

5. 4. Discussion

The model provides a reasonable description of the medium and long term effects of rainfall on the gravity observations. However, it does not describe short time scale effects like air pressure changes or air mass redistribution which may themselves be related to the precipitation events [*Meurers, 2007; Meurers et al., 2007*]. The hypothesis of the seasonal variability of the gravity recovery rate after rainfall provides a better fit to the actual gravity changes. This variability could be explained from four different factors.

First, it could be due to variability of the hydraulic conductivity K , which is related to the parameter γ [Van Camp *et al.*, 2006]. K can be expressed in the form $K=k \cdot \rho \cdot g \cdot \mu^{-1}$, where k (length²) is the permeability of the solid matrix, which depends only on the characteristics of the porous medium. The term $\rho \cdot g \cdot \mu^{-1}$ depends only on the fluid characteristics, where ρ and μ represent the density and the viscosity of the fluid, respectively. The viscosity depends on groundwater temperature. At mid-latitude locations such as Walferdange, groundwater temperature has a definite seasonal variability, even if it is less sensitive in the saturated zone. Bartolino (2003) investigated the annual groundwater temperature fluctuations as a function of depth below the surface beneath the Rio Grande in New Mexico. He estimated an annual groundwater temperature range varying from 22°C at 1 m depth below the surface to 4°C at 15 m depth below the surface..

A second seasonal effect which could positively affect our comparison, could be due to the infiltration variability. The runoff coefficients, and consequently the fraction of precipitation infiltrating the soil, can have a significant seasonal variability, related to the degree of saturation of the root zone.

Thirdly, the improvement could be due to evapotranspiration variability [Hupet and Vanclooster, 2005]. The process of evapotranspiration has a strong dependence on air and soil surface temperatures that are subjected to significant seasonal oscillations.

Finally, it could be due to regional effects. The model is based on local precipitation. However, the gravity recovery rate variability hypothesis may absorb longer wavelength seasonal effects as well [van Dam *et al.*, 2001; Crossley *et al.*, 2005].

A confirmation of the variability of parameter γ (ranging from 0.4 γ_0 to 1.6 γ_0 as results from Eq. (12) with $A=0.6$) can be obtained by a comparison with the Antecedent Precipitations Index (API) calculation model [Musy and Higy, 2003; Rosenthal *et al.* 1982]. The API relates the soil moisture level to the precipitation via the expression:

$$API_i = API_{i-1} \cdot K_p + P_{i-1} \quad (15)$$

where API_i is the index (mm) at time i , API_{i-1} the index at time $i-1$, P_{i-1} the precipitation (mm) at time $i-1$ and K_p is a coefficient <1 , that appears strongly correlated to parameter γ in Eq. (11). Rosenthal *et al.* [1982] have estimated that in Washita Basin (USA), for a superficial layer of 15 cm thickness, the value of K_p ranges from 0.84 in July to 0.99 in December. The correlation between γ and K_p indicates that the local value of γ may range from 0.1 γ_0 to 1.6 γ_0 , where γ_0 represents the yearly average value. Implementing the tank model with a constant value of γ , the seasonal variations are not absorbed and the maximum percentage of signal scatter reduction in Walferdange is only 25%.

6. Correlation between gravity variation and level of the Alzette River

In this Section, we determine the statistical correlation between the gravity variations registered by the SG and the water level changes of the Alzette River. The motivation for this analysis is that hydrological gravity may show a better correlation with river water levels than precipitation. This is because both gravity and water level changes are dependent on soil hydrological parameters, whereas rainfall is not.

The water level data is available from the Walferdange Meteorological Station 1.63 km away and downhill from the WULG (see Figure 2). The direct effect of the water level change on observed gravity variations can be neglected due to the distance between the SG and the river.

A shower [Musy and Higy, 2003] is defined as an ensemble of precipitation events related to the same meteorological perturbation. Two showers are considered distinct if, during a defined time interval Δt_s the precipitation height is inferior to a defined level Δh_s [Musy and Higy, 2003]. The water level increase duration Δt_i is defined as the time delay between the start of river level increase and the maximum river level increase due to a particular shower. The mean water level increase duration $\langle \Delta t_i \rangle = 185 \pm 35$ min is calculated for 45 showers sampled between January 2004 and April 2007. In general the same value can be adopted for the parameter Δt_s . The value of Δh_s is the minimum continuous precipitation for which a water level increase can be observed. The values of 3 hours and 0.5 mm hour^{-1} are adopted for Δt_s and Δh_m , respectively.

In addition to Δt_s and Δh_m , defined as continuity parameters, the quantitative parameters characterizing a shower are the duration (time), the precipitation height (mm), the mean intensity, ratio between the precipitation height and the shower duration, and the maximal intensity [Musy and Higy, 2003].

In this section, the time delay between the gravity and water level variations due to the showers and the correlation between the gravity and water level changes are estimated.

6.1. Time delay between gravity and water level variations

In order to estimate the time delay of gravity and water level variations, it is more precise to work with the time derivative of the quantities. The sampling interval of the water level time series $L(t)$ is 15 min. The time derivative of $L(t)$ is defined as:

$$\dot{L}(t) = \frac{L(t + \Delta t) - L(t)}{\Delta t} \quad (16)$$

The time derivative $\Delta \dot{g}(t)$ of the gravity time series $\Delta g(t)$ is computed in a similar way. Two additional parameters are defined: t_l is the time delay between the maximum values of the precipitation intensity and the gravity time derivative related to a shower, t_s is the time delay between the maximum values of water level and gravity time derivatives related to a shower (Figure 5).

The statistical analysis is made for 30 showers sampled between January 2005 and February 2007. In general, it is difficult to estimate the time delay between the signals as the relationship can be clearly identified only for some type of showers: high intensity showers and medium intensity showers with their maximum intensity concentrated in the initial phase.

Figure 5 displays the rainfall height (mm), the gravity time derivative ($\mu\text{Gal hour}^{-1}$) and the water level time derivative (cm hour^{-1}) related to the shower, which occurred on the 25 June 2006.

Figure 5

The average values and the standard deviation of t_s calculated for the entire sample, for the showers that occurred in the time span April to September (warm months), and for the showers that occurred in the time span October to March (cold months) are presented in Table 1. No substantial differences appear in the values calculated for the three periods, but because of the sample limitation these results must be interpreted with caution.

Table 1

The average values and the standard deviation of t_l are also calculated for 3 subsets: the entire sample, for the time span April to September (warm months) and the time span October to March (cold months). The results are given in Table 2.

Table 2

As the uncertainties of the t_l values are even greater than the values themselves, the SG observations do not provide any additional or complementary information already provided by the pluviometer.

6.2. Correlation between gravity and water level changes as a function of the shower height

A statistical analysis of the correlation between the shower heights, the gravity variations and the water level changes is carried out for 45 showers sampled between January 2004 and April 2007.

The parameters r_s , Δg_s and L_s are defined as the integrated water content of a shower, the gravity change, and the water level change associated with a shower (Figure 6) respectively. The L_s values as a function of the Δg_s values are displayed in Figure 7.

Figure 6

The correlation coefficient between L_s and Δg_s is 0.59. The admittance between L_s and Δg_s was calculated to $45 \pm 5 \text{ cm } \mu\text{Gal}^{-1}$

Figure 7

The L_s values as a function of the r_s values are displayed in Figure 8. The correlation coefficient between L_s and r_s is 0.68, higher than the correlation coefficient between L_s and Δg_s . The admittance of $2.2 \pm 0.4 \text{ cm mm}^{-1}$ between L_s and r_s is calculated. Again, the slightly better correlation between L_s and r_s shows that the SG observations (as compared to predictions based solely on the pluviometer) fail to provide a better insight into river levels.

Figure 8

7. Conclusions

A simple tank model, based on a combination of a mass continuity equation and Darcy's law, is implemented in order to evaluate the effects of water storage variations on the gravity observations in Walferdange. The hypothesis of the seasonal variability

of the model parameter describing the gravity recovery rate after a rainfall is introduced. No a priori information about soil moisture and groundwater storage is used to estimate the hydrological effects on gravity. The model is based solely on the observed local rainfall and is empirically parameterized.

The model reduces the scatter of the SG data by 77%. However, in the absence of the extra seasonal parameter, the maximum percentage of signal reduction is 25%. This result indicates that seasonal effects are important even when looking at short-scale spatial relationships.

The time delay between the maximum gravity and water level time derivatives and maximum precipitation intensity and gravity time derivative is 88 ± 34 minutes and 4 ± 9 minutes, respectively. The correlation between the gravity and water level variation and the precipitation amount and water level variation is 0.59 and 0.68, respectively. The admittance between water level and gravity variation and water level variation and precipitation amount is $45 \pm 5 \text{ cm } \mu\text{Gal}^{-1}$ and $2.2 \pm 0.4 \text{ cm mm}^{-1}$, respectively.

In conclusion, the gravity signal does not allow for a better prediction of the water level of the Alzette as compared to the prediction based solely on the pluviometer.

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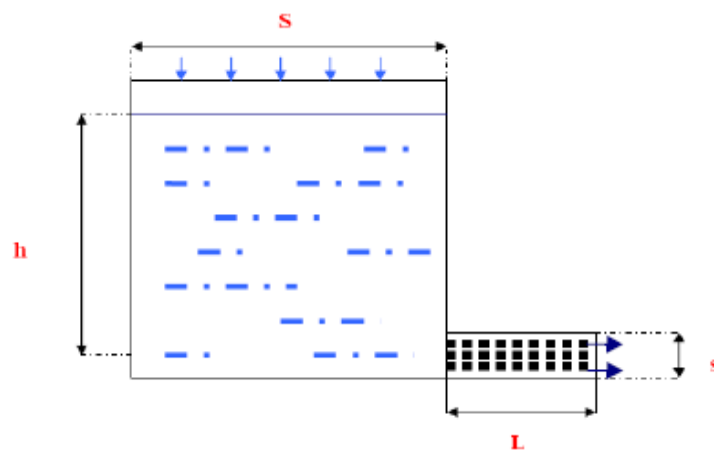


Fig. 1: The tank model: S is the tank surface, h the water level, s and L the cross section and length of the plug..

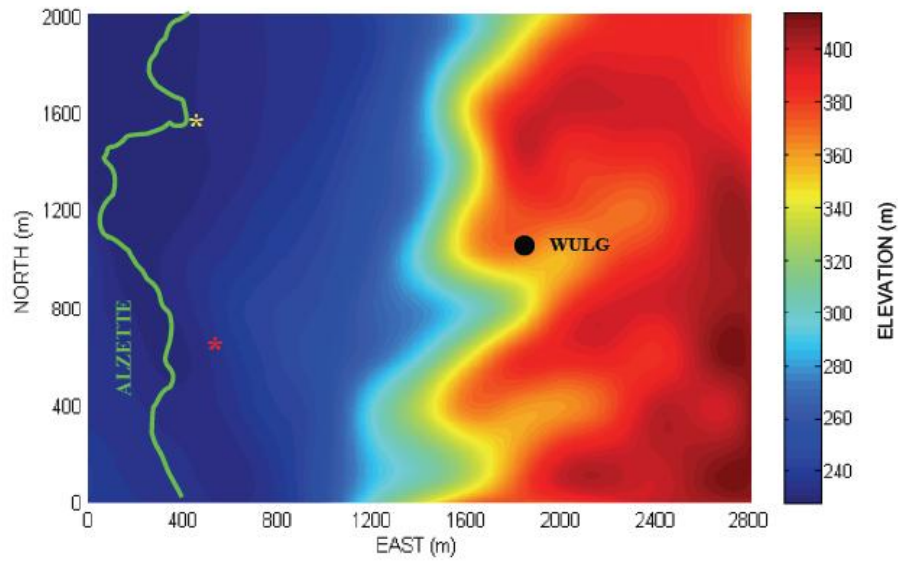


Fig. 2: Topography of the area above the WULG (black dot) and Alzette River (green line). The red asterisk represents the WULG pluviometer. The yellow asterisk represents the Walferdange Meteorological Station.

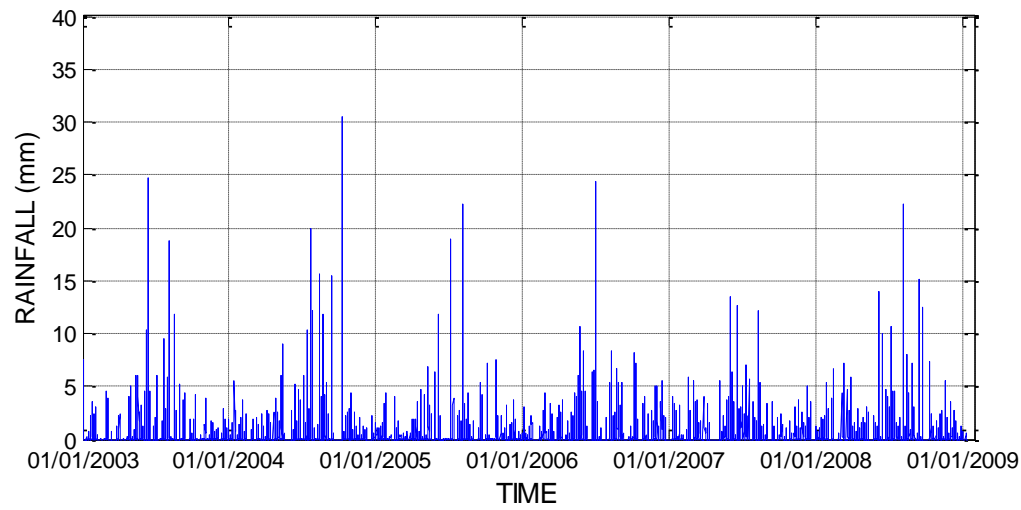


Fig. 3: Input data: Rainfall time series (mm/hour) provided by the WULG for the period 1 January 2003 to 1 January 2009.

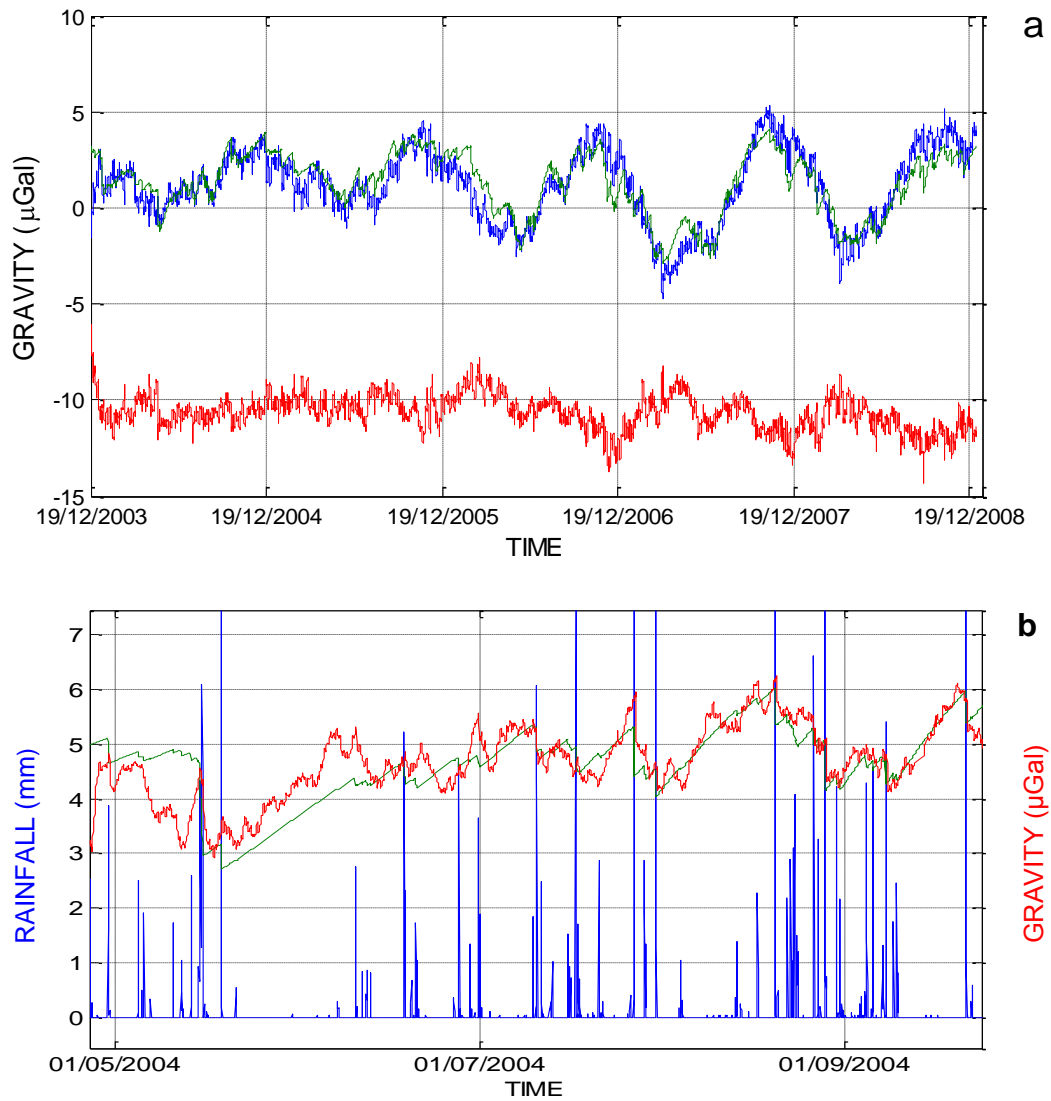


Fig. 4: Model outputs:

- a. Observed gravity (blue, μGal), modeled gravity (green, μGal) and their difference (red, μGal) at Walferdange. The comparison period is 19 December 2003 to 1 January 2009.
- b. Zoom on a: rainfall (blue, mm/hour), observed gravity (red, μGal), modeled gravity (green, μGal) for the period May to September 2004. The same scale is adopted for gravity and rainfall.

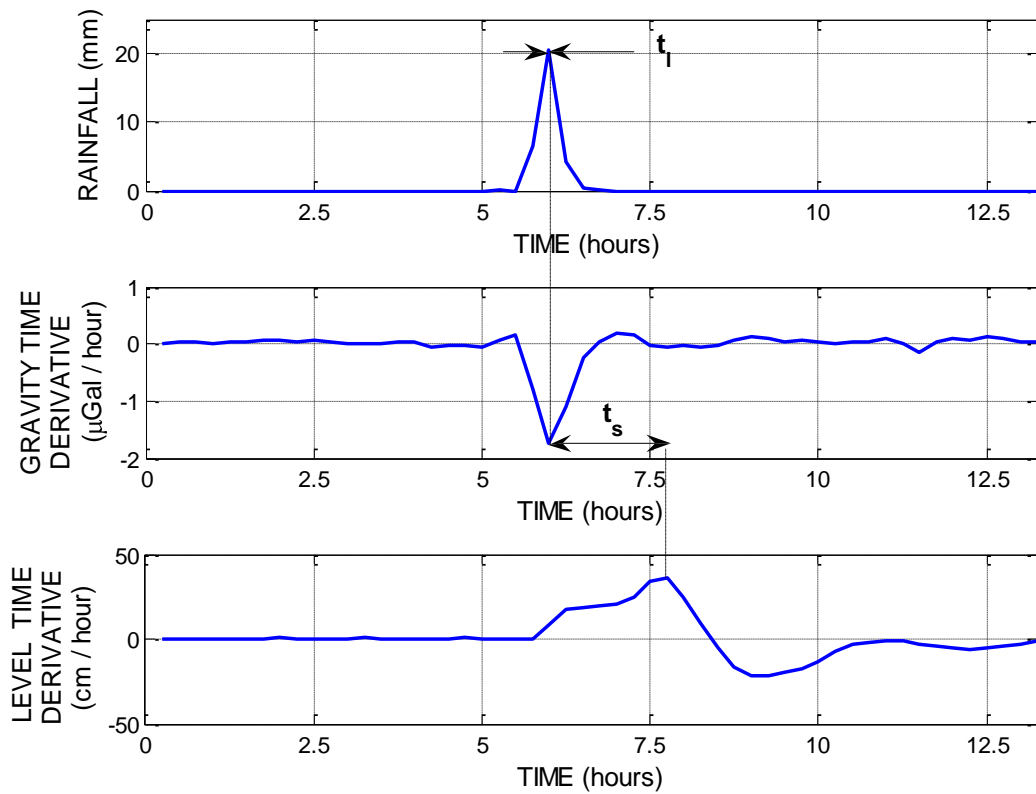


Fig. 5: Rainfall height (mm), gravity time derivative ($\mu\text{Gal}/\text{hour}$) and water level time derivative (cm/hour) related to the shower that occurred on the 25 June 2006.

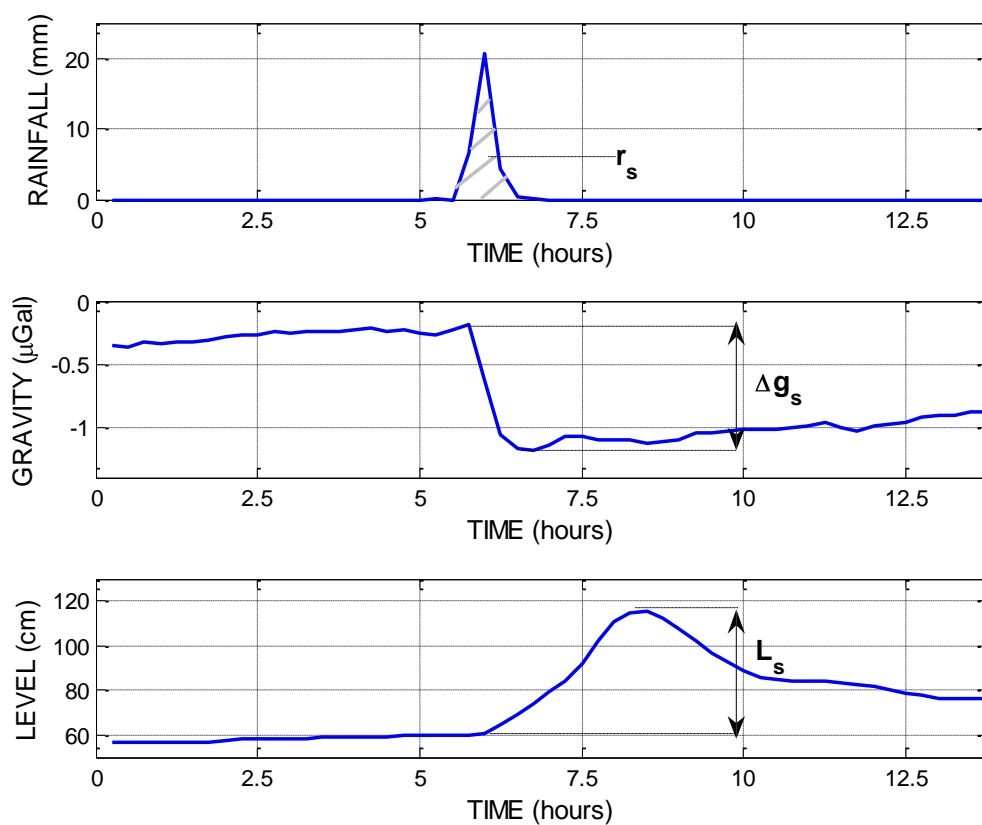


Fig. 6: Time evolution of rainfall (mm), gravity (μ Gal) and water level (cm) related to the shower that occurred on the 25 June 2006.

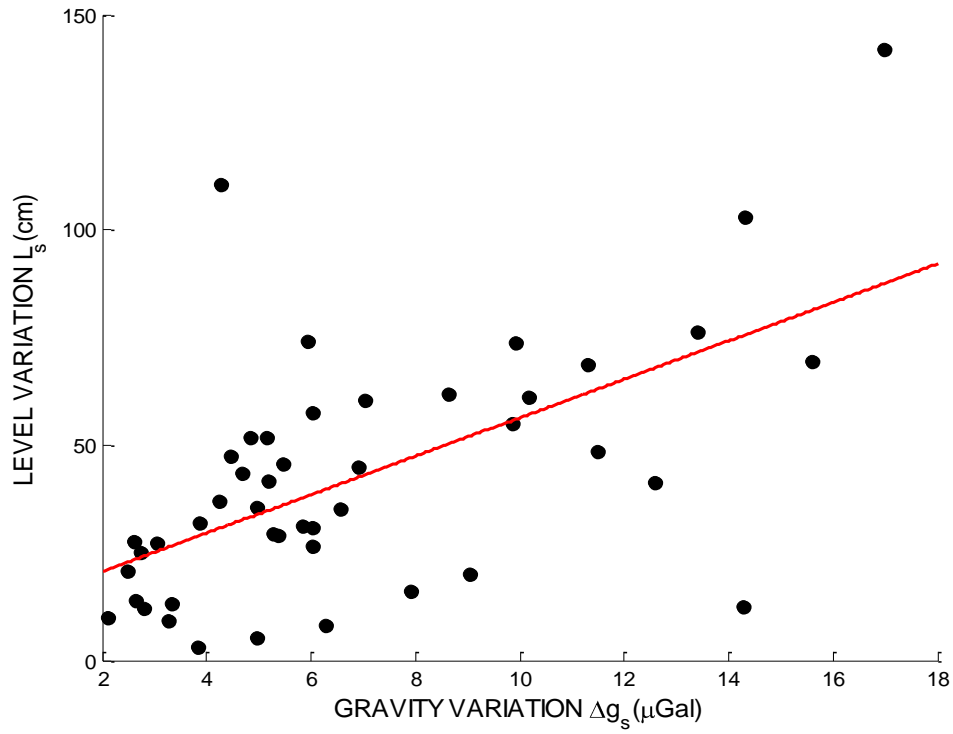
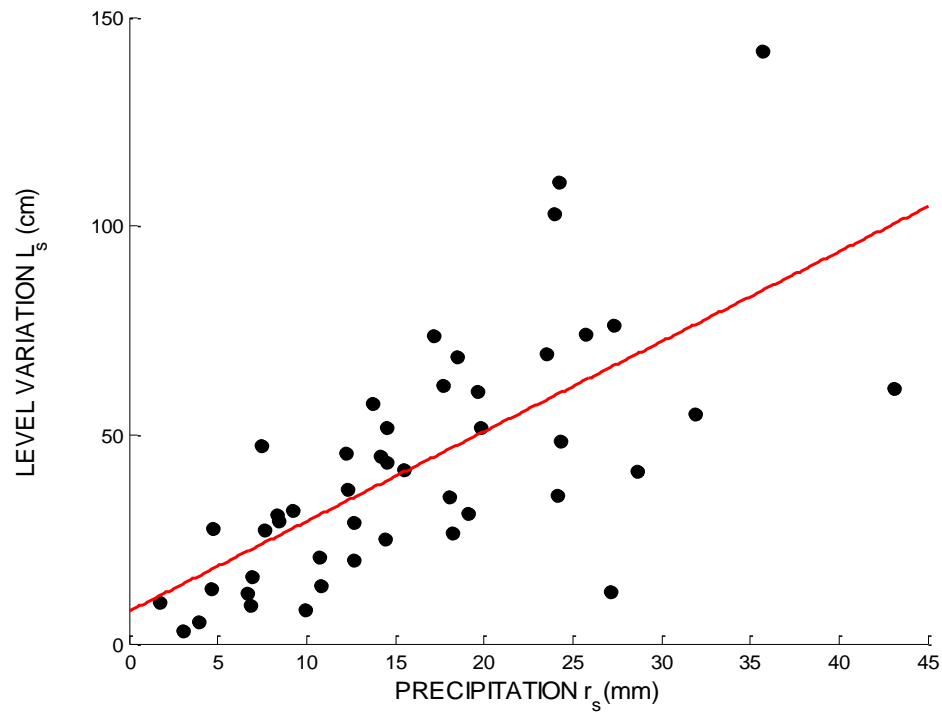


Fig. 7: Correlation between gravity and water level changes for the showers between January 2004 and April 2007.

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92 **Fig. 8:** Correlation between precipitation heights and water level changes for the
 93 showers between January 2004 and April 2007.

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	January 2005-February 2007	April-September 2005-2007	October-March 2005-2007
t_s	88 ± 34 min	85 ± 34 min	93 ± 34 min

Table 1: Time delay t_s between the maximum values of gravity and water level time derivatives, for the showers sampled between January 2005 and February 2007.

	January 2005-February 2007	April-September 2005-2007	October-March 2005-2007
t_l	4 ± 9 min	5 ± 9 min	1 ± 9 min

Table 2: Time delay t_l between the maximum values of shower intensity and gravity time derivative, for the showers sampled between January 2005 and February 2007.