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### Transport properties of MgB<sub>2</sub> single crystals doped with electrons and holes

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#### Abstract

The thermoelectric power of C, Mn, C:Li, and Al:Li substituted MgB<sub>2</sub> single crystals has been investigated in the temperature range 10-300 K. Both the in-plane ( $S_{ab}$ ) and the outof-plane ( $S_c$ ) thermopowers are positive for the non-substituted crystal and both  $S_{ab}$  and  $S_c$ change a sign for crystals doped with electrons where C is substituted for B in the amount larger than 5 % at. When Li is substituted for Mg, the  $\pi$  band rather than the  $\sigma$  band is doped with holes and the doping effects are much more subtle. The anisotropy ratio of the nonsubstituted crystal  $S_{ab}/S_c \approx 3$  and this ratio is strongly reduced by the substitution of C. Isovalent magnetic Mn ions, which substitute for Mg with a drastic reduction of  $T_c$ , do not influence the values and temperature dependences of both  $S_{ab}$  and  $S_c$ .

Substitutional chemistry can modify in the controlled way electronic structure and properties such as superconducting transition temperature,  $T_c$ , upper critical field,  $H_{c2}$ , gaps width, intra- and inter-band scattering, defect structure, etc. In order to study the influence of intra- and interband scattering and doping on the normal and superconducting properties of MgB<sub>2</sub>, investigation of substituted single crystals are particularly valuable.[1-6] Introducing C, Al, Li, Mn, and Fe into the melt we have grown MgB<sub>2</sub> crystals where Mg or B were partially substituted with iso- or heterovalent and non-magnetic or magnetic ions. Thus the scattering centers were added and/or the crystals were doped with electrons or holes. In this work we study the resistance, *R*, and thermoelectric power, *S*, of MgB<sub>2</sub> single crystals substituted with electron-doping C for B and isovalent magnetic Mn<sup>2+</sup> for Mg. The preliminary results for crystals co-substituted with Al:Li and C:Li are also presented for comparison.

Single crystals of substituted  $MgB_2$  were grown under high pressure using a cubic anvil press. The applied temperature and pressure conditions to obtain C, Al, and Mn substituted crystals were described in Ref. 3-5. The successful method to grow Li, Al:Li, and C:Li substituted crystals has been presented in Ref. 7. The C content in the crystals was estimated from changes in the *a* lattice parameter. The contents of Al, Mn and Fe were determined with EDX. The Li content was determined from the structure refinement.[7] The methods to measure *R* and *S* were described in details in Ref. 3 and 8, respectively.

Normalized *R* versus temperature is shown in Fig. 1 for the Mg(B<sub>1-x</sub>C<sub>x</sub>)<sub>2</sub> single crystals with x = 0 and 0.05. The resistance ratio RRR defined as R(300 K)/R(40 K) decreases from about 7, observed for the non-substituted (x = 0) crystal, to RRR = 2 and 1.4 obtained for crystals with x = 0.05 and 0.095 (not shown in the Figure), respectively. This nonlinear decrease of RRR with the carbon content may result in a resistance saturation effect, which can be attributed to the created by substitution microscopic defects acting as electron scattering centers. Thus, the increase of *R* with the carbon content is not a simple effect of

electron doping but also the scattering mechanism should be considered as an important one. This conclusion is supported by the S(T) results, which show a decrease of S with C content indicative of more metallic character of the substituted crystals. The S(T) results for the Mg(B<sub>1-x</sub>C<sub>x</sub>)<sub>2</sub> crystals with x = 0 and 0.05 are shown in Fig. 2. As thermoelectric power is only moderately sensitive to the microscopic defects, the main reason for the observed changes of S with the substitution is the electron doping. This doping reduces considerably the absolute value of S and also its anisotropy, defined as a ratio  $S_{ab}/S_c$ , where  $S_{ab}$  and  $S_c$  are results obtained with a temperature gradient along the ab plane and the c axis, respectively. For the non-substituted crystals,  $S_{ab}/S_c \cong 3$  at 300 K and this ratio decreases to about -1.4 for the substituted crystal with x = 0.05. Note, that for this crystal,  $S_c$  is negative suggestive of the significant change in the carrier distribution.

The influence of the C substitution on thermoelectric power is summarized in Fig. 3. The main effect is a rapid reduction of both S and  $S_{ab}/S_c$ , for the MgB<sub>2</sub> crystals with a moderate amount of substituted C (x = 0.05,  $T_c = 34.5$  K), and then saturation of these quantities for higher substitutions. At this stage of our studies we are not able to interpret that effect. Figure 3 shows also  $S_{ab}$  for C:Li co-substituted crystals, where electrons doped with C, which is substituted for B, seem to be partially compensated with holes added whit Li, which is co-substituted for Mg.[7] The results fit the  $S_{ab}(T_c)$  dependence well suggestive of  $T_c$  as a parameter that characterizes the transport properties of MgB<sub>2</sub> adequately. Interestingly, the  $S_{ab}(T_c)$  dependence is not followed by results obtained for Al:Li co-substituted crystals. This can be explained by the fact that Al substituted for Mg may dope electrons into the  $\pi$  band, so this substituted crystals, so the Li substitution is negligible in this case.

Finally we studied the S(T) properties of MgB<sub>2</sub> crystals substituted with isovalent magnetic ions. Figure 4 shows  $S_{ab}(T)$  and  $S_c(T)$  for Mg<sub>1-x</sub>Mn<sub>x</sub>B<sub>2</sub> with x = 0.007 and  $T_c = 29.3$ K. Surprisingly, both the values and the temperature dependences of  $S_{ab}$  and  $S_c$  are very

similar to those obtained for the non-substituted crystal, despite of the drastic reduction of  $T_c$ . This can be explained by the fact that the Mn<sup>2+</sup> ions substituted for Mg are isovalent and that the very small amount of substituted Mn can create only a small amount of the electron scattering centers. The main effect that reduces  $T_c$  is magnetic pair-breaking, which does not influence the electronic structure and thus can not modify the S(T) properties.

In conclusion, availability of sizeable MgB<sub>2</sub> crystals doped with electrons and holes enable us to perform advanced studies of their superconducting and normal-state properties including transport characteristics. Both resistance and thermoelectric power are strongly influenced by doping when C is substituted for B. This substitution, which introduces electrons into the  $\sigma$  band, results in a large increase of *R*, most likely due to the creation of additional scattering centers. On the other hand, for the C-substituted MgB<sub>2</sub>, *S* decreases with the substitution and even changes a sign to negative indicative of more metallic and electronconducting character of substituted crystals. Another interesting observation is that the substitution of Li for Mg, which introduces holes to the  $\pi$  band, changes the superconducting and transport properties in a much subtle way. Thus, the electron- and hole-doping effects seem to be asymmetric. The substitution of isovalent magnetic Mn<sup>2+</sup> ions, which drastically influences the superconducting properties like  $T_c$  and  $H_{c2}$ , does not appear to have any effects on both the value and the temperature dependence of *S*. This suggests, not surprisingly, that the pair-breaking effect, which is responsible for the rapid reduction of  $T_c$  in the Mn substituted crystals, does not modify the Fermi surface to any substantial degree.

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Fig. 1. Temperature dependence of the in-plane resistance ratio R(T)/R(300 K) for Mg(B<sub>1</sub>.  $_xC_x)_2$  crystals, non-substituted (circles) and carbon-substituted with x = 0.05 (triangles). The superconducting transition temperature measured with low current ( $i = 2.5 \text{ A/cm}^2$ ) is  $T_c = 38.2$  and 34.5 K for the samples with x = 0 and 0.05, respectively. Adapted from Ref. 3.



Fig. 2. Thermoelectric power of Mg(B<sub>1-x</sub>C<sub>x</sub>)<sub>2</sub> crystals, non-substituted with  $T_c = 38.2$  K (circles) and carbon-substituted with x = 0.05 and  $T_c = 34.5$  K (triangles), measured along the *ab* plane (*S<sub>ab</sub>*) and the *c* axis (*S<sub>c</sub>*).



Fig. 3. Thermoelectric power along the *ab* plane and the *c* axis at 300 K (circles, star) and 100 K (squares, pentagon) versus  $T_c$  of the Mg(B<sub>1-x</sub>C<sub>x</sub>)<sub>2</sub> crystals with  $0 \le x \le 0.1$  (circles, squares), and the MgB<sub>2</sub> crystals substituted with C:Li (crossed symbols) and Al:Li (star, pentagon) to compare. For Mg(B<sub>1-x</sub>C<sub>x</sub>)<sub>2</sub> with  $0 \le x \le 0.09$ ,  $T_c$  decreases with the C content roughly linearly.



Fig. 4. Thermoelectric power of the Mg<sub>1-x</sub>Mn<sub>x</sub>B<sub>2</sub> crystals with x = 0.007 and  $T_c = 29.3$  K, measured along the *ab* plane and the *c* axis. Comparing with the non-substituted crystal, the values of both  $S_{ab}$  and  $S_c$  are only slightly enhanced despite of the drastically reduced  $T_c$ .