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HARDY INEQUALITY AND WEYL ASYMPTOTIC FOR DISCRETE LAPLACIANS

SYLVAIN GOLÉRIA

ABSTRACT. In this paper we study in detail some spectral properties of the magnetic discrete Laplacian. We identify its form-domain, characterize the absence of essential spectrum and provide the Weyl asymptotic for the eigenvalues.

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1. INTRODUCTION

The uncertainty principle is a central point in quantum physics. It can be expressed by the following Hardy inequality:

$$(1.1) \quad \left(\frac{n-2}{2}\right)^2 \int_{\mathbb{R}^n} \left| \frac{1}{|x|} f(x) \right|^2 dx \leq \int_{\mathbb{R}^n} |\nabla f|^2 dx = \langle f, -\Delta_{\mathbb{R}^n} f \rangle, \text{ where } n \geq 3,$$

and $f \in C_c^\infty(\mathbb{R}^n)$. Roughly speaking, the Laplacian controls some local singularities of a potential. In this paper, we investigate which potentials a discrete Laplacian is able to control. Obviously, we will not focus on local singularities as the value of a potential on a vertex has to be finite. However, unlike in the continuous case, we will control potentials that explodes at infinity.

We start with some definitions and fix notation for graphs. We refer to [CdV, Chu, MW] for surveys on the matter. Let \mathcal{V} be a countable set. Let $\mathcal{E} := \mathcal{V} \times \mathcal{V} \rightarrow [0, \infty)$ and assume that $\mathcal{E}(x, y) = \mathcal{E}(y, x)$, for all $x, y \in \mathcal{V}$. We say that $G := (\mathcal{E}, \mathcal{V})$ is an unoriented weighted graph with *vertices* \mathcal{V} and *weighted edges* \mathcal{E} . In the setting of electrical networks, the weights correspond to the conductances. We say that $x, y \in \mathcal{V}$ are *neighbors* if $\mathcal{E}(x, y) \neq 0$ and denote it by $x \sim y$. We say there is a *loop* in $x \in \mathcal{V}$ if $\mathcal{E}(x, x) \neq 0$. The set of *neighbors* of $x \in E$ is denoted by $\mathcal{N}_G(x) := \{y \in E \mid x \sim y\}$. The *degree* of $x \in V$ is by definition $|\mathcal{N}_G(x)|$, the number of neighbors of x . The graph is of *bounded degree*, if $\sup_{x \in \mathcal{V}} |\mathcal{N}_G(x)|$ is finite. A graph is *locally finite* if $|\mathcal{N}_G(x)|$ is finite for all $x \in V$. We also need a weight on the vertices $m : \mathcal{V} \rightarrow (0, \infty)$. Finally, as we are dealing with magnetic fields, we fix a phase $\theta : \mathcal{V} \times \mathcal{V} \rightarrow [-\pi, \pi]$, such that $\theta(x, y) = -\theta(y, x)$. We denote by $\theta_{x,y} := \theta(x, y)$. A graph is *connected*,

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if for all $x, y \in V$, there exists an x - y -*path*, i.e., there is a finite sequence $(x_1, \dots, x_{N+1}) \in \mathcal{V}^{N+1}$ such that $x_1 = x$, $x_{N+1} = y$ and $x_n \sim x_{n+1}$, for all $n \in \{1, \dots, N\}$. In the sequel, we shall always consider (magnetic) graphs $G = (\mathcal{V}, \mathcal{E}, m, \theta)$, which are unoriented, locally finite, connected and have no loop. A graph G is *simple* if \mathcal{E} has values in $\{0, 1\}$ and $\theta = 1$. Finally, we recall that a *bi-partite* graph is a graph whose vertex set can be partitioned into two subsets in such a way that no two points in the same subset are neighbors. Trees are bi-partite graphs.

We now associate a certain Hilbert space and some operators on it to a given graph $G = (\mathcal{V}, \mathcal{E}, m, \theta)$. Let $\ell^2(G, m^2) := \ell^2(\mathcal{V}, m^2; \mathbb{C})$ be the set of functions $f : \mathcal{V} \rightarrow \mathbb{C}$, such that $\|f\|^2 := \sum_{x \in \mathcal{V}} m(x)^2 |f(x)|^2$ is finite. The associated scalar product is given by $\langle f, g \rangle = \sum_{x \in \mathcal{V}} m^2(x) \overline{f(x)} g(x)$, for $f, g \in \ell^2(\mathcal{V}, m^2)$. We also denote by $\mathcal{C}_c(\mathcal{V})$ the set of functions $f : \mathcal{V} \rightarrow \mathbb{C}$, which have finite support. We define the quadratic form:

$$(1.2) \quad \mathcal{Q}(f, f) := \mathcal{Q}_{\mathcal{E}, \theta}(f, f) := \frac{1}{2} \sum_{x, y \in \mathcal{V}} \mathcal{E}(x, y) |f(x) - e^{i\theta_{x,y}} f(y)|^2 \geq 0, \text{ for } f \in \mathcal{C}_c(V).$$

It is closable and there exist a unique self-adjoint operator $\Delta_{\mathcal{E}, \theta}$, such that $\mathcal{Q}_{\mathcal{E}, \theta}(f, f) = \langle f, \Delta_{\mathcal{E}, \theta} f \rangle$, for $f \in \mathcal{C}_c(V)$ and $\mathcal{D}(\Delta_{\mathcal{E}, \theta}^{1/2}) = \mathcal{D}(\mathcal{Q}_{\mathcal{E}, \theta})$, where the latter is the completion of $\mathcal{C}_c(\mathcal{V})$ under $\|\cdot\|^2 + \mathcal{Q}_{\mathcal{E}, \theta}(\cdot, \cdot)$. This operator is the *Friedrichs extension* associated to the form $\mathcal{Q}_{\mathcal{E}, \theta}$. It acts as follows:

$$(1.3) \quad \Delta_{\mathcal{E}, \theta} f(x) := \frac{1}{m^2(x)} \sum_{y \in \mathcal{V}} \mathcal{E}(x, y) (f(x) - e^{i\theta_{x,y}} f(y)), \text{ for } f \in \mathcal{C}_c(\mathcal{V}).$$

When $m = 1$, it is essentially self-adjoint on $\mathcal{C}_c(V)$ (see Section 2.2 for further discussions). If G is simple, we shall simply write Δ_G . There exist other definitions for the discrete Laplacian, e.g., [CdV, Chu, MW], the one we study here is sometimes call the “physical Laplacian”.

In $\ell^2(G, m^2)$, we define the *weighted degree* by

$$d_G(x) := \frac{1}{m^2(x)} \sum_{y \in \mathcal{V}} \mathcal{E}(x, y).$$

Given a function $V : \mathcal{V} \rightarrow \mathbb{C}$, we denote by $V(Q)$ the operator of multiplication by V . It is elementary that $\mathcal{D}(d_G^{1/2}(Q)) \subset \mathcal{D}(\Delta_{\mathcal{E}, \theta}^{1/2})$. Indeed, one has:

$$(1.4) \quad \langle f, \Delta_{\mathcal{E}, \theta} f \rangle = \frac{1}{2} \sum_{x \in \mathcal{V}} \sum_{y \sim x} \mathcal{E}(x, y) |f(x) - f(y)|^2 \leq \sum_{x \in \mathcal{V}} \sum_{y \sim x} \mathcal{E}(x, y) (|f(x)|^2 + |f(y)|^2) = 2 \langle f, d_G(Q) f \rangle,$$

for $f \in \mathcal{C}_c(\mathcal{V})$. This inequality also gives a necessary condition for the absence of essential spectrum for $\Delta_{\mathcal{E}, \theta}$ (see Corollary 2.1). In Proposition 4.5, we also prove that the constant 2 cannot in general be improved. It is also easy to see that $\Delta_{\mathcal{E}, \theta}$ is bounded if and only if $d_G(Q)$ is (see Proposition 2.1).

In this paper we are interested in minorating the Laplacian with the help of the weighted degree. In the case of non-magnetic Laplacian, one standard approach is to use isoperimetric inequalities. The classical version gives estimates on the bottom of the spectrum. This is not adapted to our situation. We rely on a modified version. Given a subset $W \subset \mathcal{V}$, let $\partial W := \{x \in W, \text{ there is } y \sim x \text{ such that } y \notin W\}$. We define the following isoperimetric constant associated to (the weighted degree of) G by

$$\alpha(G) := \inf_{W \subset \mathcal{V}, \#W < \infty} \frac{\langle \mathbf{1}_{\partial W}, d_G(Q) \mathbf{1}_{\partial W} \rangle}{\langle \mathbf{1}_W, d_G(Q) \mathbf{1}_W \rangle},$$

where $\mathbf{1}_X$ denotes the characteristic function of X . By [KL, page 14, line -4], see also [Do, DK, Kel] and reference therein, we obtain:

$$(1.5) \quad \left(1 - \sqrt{1 - \alpha^2(G)}\right) \langle f, d_G(Q) f \rangle \leq \langle f, \Delta_{\mathcal{E}, 0} f \rangle \leq \left(1 + \sqrt{1 - \alpha^2(G)}\right) \langle f, d_G(Q) f \rangle,$$

for all $f \in \mathcal{C}_c(\mathcal{V})$. So, if $\alpha(G) > 0$, (1.4) is improved and we have that $\mathcal{D}(\Delta_{\mathcal{E}, 0}^{1/2}) = \mathcal{D}(d_G^{1/2}(Q))$. In particular, by Proposition 2.3, one sees that the essential spectrum $\sigma_{\text{ess}}(d_G(Q))$ is empty if and only if $\sigma_{\text{ess}}(\Delta_{\mathcal{E}, 0})$ is. We refer to Theorem 4.1 for the equality of the domains and to Proposition 5.2 for the stability of the essential spectrum.

We point out that a reciprocal is also true. Namely, Proposition 3.4 ensures that if there is $a > 0$ so that $a \langle f, \Delta_{\mathcal{E}, 0} f \rangle \geq \langle f, d_G(Q) f \rangle$, for all $f \in \mathcal{C}_c(\mathcal{V})$, then $\alpha(G) > 0$.

Assume that $\alpha(G) > 0$. Supposing that $\sigma_{\text{ess}}(\Delta_{\mathcal{E},0}) = \emptyset$ (or equivalently that $\lim_{|x| \rightarrow \infty} d_G(x) = +\infty$), the inequality (1.5) and the min-max principle, see Proposition 2.3, provide the bound

$$\left(1 - \sqrt{1 - \alpha^2(G)}\right) \leq \liminf_{\lambda \rightarrow \infty} \frac{\mathcal{N}_\lambda(\Delta_{\mathcal{E},\theta})}{\mathcal{N}_\lambda(d_G(Q))} \leq \limsup_{\lambda \rightarrow \infty} \frac{\mathcal{N}_\lambda(\Delta_{\mathcal{E},\theta})}{\mathcal{N}_\lambda(d_G(Q))} \leq \left(1 + \sqrt{1 - \alpha^2(G)}\right),$$

where $\mathcal{N}_\lambda(A) := \dim \text{Ran } \mathbf{1}_{(-\infty, \lambda]}(A)$, for a self-adjoint operator A . This estimate has to be refined to give the asymptotic of eigenvalues and to deal with magnetic fields. Moreover, (1.5) is not stable by small perturbation for the question of the equality of the form-domain. For instance, take a simple graph G_1 , such that $\alpha(G_1) > 0$ and the (simple) half-line graph G_2 . Note that $\alpha(G_2) = 0$. Now connect the disjoint union of G_1 and G_2 by one edge to obtain a new graph G . One sees easily that $\alpha(G) = 0$ and $\mathcal{D}(\Delta_G^{1/2}) = \mathcal{D}(d_G^{1/2}(Q))$. This is why we seek a minoration by $ad_G(Q) - b$ for some $a, b > 0$.

On the other hand we know that the Laplacian acting in $\ell^2(\mathcal{V}, m^2)$ is unitarily equivalent to a Schrödinger operator acting in another $\ell^2(\mathcal{V}, m_0^2)$ (see Proposition 3.2). This has already been noticed before, e.g., [CTT]. By extracting some positivity, we obtain our analog of the Hardy inequality:

Proposition 1.1. *Let $G = (\mathcal{V}, \mathcal{E}, m_0, \theta)$ be a locally finite graph. Given $m : \mathcal{V} \rightarrow (0, \infty)$, one has*

$$(1.6) \quad \langle f, V_m(Q)f \rangle \leq \langle f, \Delta_{\mathcal{E},\theta} f \rangle, \text{ for } f \in \mathcal{C}_c(\mathcal{V}),$$

where

$$(1.7) \quad V_m(x) := d_G(x) - W_m(x), \text{ with } W_m(x) := \frac{1}{m_0^2(x)} \sum_{y \in \mathcal{V}} \mathcal{E}(x, y) \frac{m(y)}{m(x)} \frac{m_0(x)}{m_0(y)}.$$

Moreover, if G is bi-partite, we also get:

$$\langle f, \Delta_{\mathcal{E},\theta} f \rangle \leq \langle f, (d_G(Q) + W_m(Q))f \rangle, \text{ for } f \in \mathcal{C}_c(\mathcal{V}).$$

Note that by choosing $m = m_0$, we recover that $\Delta_{\mathcal{E},\theta} \geq 0$. Moreover, V_m is independent of the magnetic field. The minoration is also different from the Kato's inequality of [DM]. We stress that the inequality (1.6) is in some cases trivial, e.g., Proposition 3.3. One has to find a favorable situation in order to exploit it. This is the case for some perturbations of weighted trees. We present our main result:

Theorem 1.1. *Let $G_\circ = (\mathcal{V}, \mathcal{E}_\circ, m, \theta_\circ)$ be a weighted tree. Assume that there is $\varepsilon_0 \in (0, 1)$, so that*

$$(1.8) \quad C_0 := \sup_{x \in \mathcal{V}} d_{G_\circ}^{\varepsilon_0 - 1}(x) \mathcal{E}_\circ(\overline{x}, x) m^{-2}(x) < \infty.$$

Let $G = (\mathcal{V}, \mathcal{E}, m, \theta)$ be a perturbed graph and $V : \mathcal{V} \rightarrow \mathbb{R}$ be a potential, satisfying:

$$(1.9) \quad |V(x)| + \Lambda(x) = o(1 + d_{G_\circ}(x)), \text{ as } |x| \rightarrow \infty, \text{ where } \Lambda(x) := \frac{1}{m^2(x)} \sum_{y \sim x} |\mathcal{E}(x, y) - \tilde{\mathcal{E}}(x, y)|.$$

Then, one has that:

- (a) *The quadratic form associated to $\Delta_{\mathcal{E},\theta} + V$ is bounded from below by some constant $-C$. We denote by H the associated Friedrichs extension.*
- (b) *For all $\varepsilon > 0$, there is $c_\varepsilon \geq 0$, so that*

$$(1.10) \quad (1 - \varepsilon) \langle f, d_{G_\circ}(Q)f \rangle - c_\varepsilon \|f\|^2 \leq \langle f, Hf \rangle \leq (1 + \varepsilon) \langle f, d_{G_\circ}(Q)f \rangle + c_\varepsilon \|f\|^2,$$

for $f \in \mathcal{C}_c(\mathcal{V})$. We have $\mathcal{D}((H + C)^{1/2}) = \mathcal{D}((d_{G_\circ}(Q))^{1/2})$.

- (c) *The essential spectrum of H is equal to that of $\Delta_{\mathcal{E}_\circ, \theta_\circ}$.*
- (d) *The essential spectrum of H is empty if and only if $\lim_{|x| \rightarrow \infty} d_{G_\circ}(x) = +\infty$. In this case we obtain:*

$$(1.11) \quad \lim_{\lambda \rightarrow \infty} \frac{\mathcal{N}_\lambda(H)}{\mathcal{N}_\lambda(d_{G_\circ}(Q))} = 1.$$

Note that we improve on the bound (1.4). There is no hypothesis on θ or on θ_\circ . We point out that Hypothesis (1.8) is fulfilled by simple trees, i.e., when $\mathcal{E} = m_0 = 1$. The inequality (1.10) is valid for a larger class of perturbations (see Proposition 5.1). We mention also that Zorn's Lemma ensures that every simple graph has a maximal subtree.

We point out that the first part of d), namely the absence of the essential spectrum, has been studied in many works, e.g., [Kel, KL, KL2, KLW]. They generalize some ideas of [DK, Fuj]. Their approach is based on some isoperimetric estimates and on the Persson's Lemma. The latter characterizes the infimum of the essential spectrum.

The asymptotic of eigenvalues is a novelty and was not considered in the literature before. We stress that one can prescribe any asymptotic of eigenvalues by choosing a proper tree G (and in fact d_G). We mention that the spectral asymptotic estimates obtained in [DM] are for some operators with non-empty essential spectrum. They study graphs which are equipped with a free action of a discrete group and establish a bound on the $\text{tr} e^{-t\Delta_{\mathcal{E},\theta}}$, where the trace is adapted to a fundamental domain.

We turn to the question of the form-domain. We stress that we do not suppose that the isoperimetric constant is non-zero. To our knowledge, this is the first time that the form-domain of the unbounded discrete Laplacian on a simple tree is identified. It is remarkable that the form-domain coincides with that of $d_G(Q)$, a multiplication operator. A useful consequence is the stability of the essential spectrum, obtained in c). This is also new. On the other hand, we stress that there are simple bi-partite graphs, such that the form-domain of the Laplacian is different from that of $d_G(Q)$ (see Proposition 4.5). In this case, (1.10) is not fulfilled.

Having the same form-domain does not necessarily ensure that the domain are also equal. In Proposition 4.3, we construct a simple tree which is such an example. However, under some further hypothesis on the graph, Theorem 4.1 ensures that the domain of the magnetic Laplacian is also the one of $d_G(Q)$. In Proposition 4.2, we give an example of simple tree T , which have 0 as associated isoperimetric constant and such that the domain of the Laplacian is the same as that of the weighted degree. Moreover, one obtains that $\sigma(\Delta_T) = \sigma_{\text{ac}}(\Delta_T) = [0, \infty)$.

Finally we present the organization of the paper. In section 2.2 we provide a new criterion of essential self-adjointness. Next, in section 2.3, we recall some well-known facts about the min-max principle, its relation to the bottom of the essential spectrum and compactness. Then, in Section 3.1 we prove the Hardy inequality and discuss its triviality. In Section 3.2, we prove Theorem 1.1 in the context of trees. Besides, in Section 4.1 we discuss the question of the domain of the Laplacian on a general graph and that of form-domain on bi-partite graphs in Section 4.2. Perturbation theory is developed in Section 5. Finally we provide two appendices, one concerning the C^1 regularity and another one concerning the Helffer-Sjöstrand's formula.

Notation: We denote by \mathbb{N} the non-negative integer. In particular, $0 \in \mathbb{N}$. We set $\langle x \rangle := (1 + x^2)^{1/2}$. Given a set X and $Y \subseteq X$ let $\mathbf{1}_Y: X \rightarrow \{0, 1\}$ be the characteristic function of Y . We denote also by Y^c the complement set of Y in X . We consider only separable complex Hilbert space. We denote by $\mathcal{B}(\mathcal{H}, \mathcal{K})$, the space of bounded operator from the Hilbert spaces \mathcal{H} to \mathcal{K} . Given a function $V: \mathcal{V} \rightarrow \mathbb{C}$, we denote by $V(Q)$ the operator of multiplication by V .

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2. GENERAL PROPERTIES

2.1. Few words about the Friedrichs extension. Given a dense subspace \mathcal{D} of a Hilbert space \mathcal{H} and a non-negative symmetric operator H on \mathcal{D} . Let \mathcal{H}_1 be the completion of \mathcal{D} under the norm given by $\mathcal{Q}(\varphi)^2 = \langle H\varphi, \varphi \rangle + \|\varphi\|^2$. The domain of the Friedrichs extension of H , is given by

$$\begin{aligned} \mathcal{D}(H_{\mathcal{F}}) &= \{f \in \mathcal{H}_1 \mid \mathcal{D} \ni g \mapsto \langle Hg, f \rangle + \langle g, f \rangle \text{ extends to a norm continuous function on } \mathcal{H}\} \\ &= \mathcal{H}_1 \cap \mathcal{D}(H^*). \end{aligned}$$

For each $f \in \mathcal{D}(H_{\mathcal{F}})$, there is a unique u_f such that $\langle Hg, f \rangle + \langle g, f \rangle = \langle g, u_f \rangle$, by Riesz' Theorem. The Friedrichs extension of H , is given by $H_{\mathcal{F}}f = u_f - f$. It is a self-adjoint extension of H , e.g., [RS, Theorem X.23].

It remains to describe the domain of the adjoint of a discrete Schrödinger operator. This is well-known, e.g., [CTT, KL2]. Let $G = (\mathcal{E}, \mathcal{V}, m, \theta)$ be a weighted graph and $V: \mathcal{V} \rightarrow \mathbb{R}$ be a potential. We set the Schrödinger operator $\mathcal{H}_G|_{\mathcal{C}_c(\mathcal{V})} := (\Delta_{\mathcal{E},\theta} + V(Q))|_{\mathcal{C}_c(\mathcal{V})}$. The domain of its adjoint is given by

$$\mathcal{D}((\mathcal{H}_G|_{\mathcal{C}_c(\mathcal{V})})^*) = \left\{ f \in \ell^2(\mathcal{V}, m^2), x \mapsto \frac{1}{m^2(x)} \left(\sum_{y \in \mathcal{V}} \mathcal{E}(x, y) f(y) \right) + V(x) f(x) \in \ell^2(\mathcal{V}, m^2) \right\}.$$

Then, given $f \in \mathcal{D}((\mathcal{H}_G|_{\mathcal{C}_c(\mathcal{V})})^*)$, one has:

$$((\mathcal{H}_G|_{\mathcal{C}_c(\mathcal{V})})^* f)(x) = \left(\frac{1}{m^2(x)} \sum_{y \sim x} \mathcal{E}(x, y) f(y) \right) + V(x) f(x), \text{ for all } x \in \mathcal{V}.$$

By definition, the operator $\mathcal{H}_G|_{\mathcal{C}_c(\mathcal{V})}$ is *essentially self-adjoint* if its closure is equal to its adjoint.

2.2. Essential self-adjointness. Before talking about essential self-adjointness, we deal with the trivial case, the boundedness of the Laplacian and refer to [KL2, KLW] for more discussions in the setting of Dirichlet forms and ℓ^p spaces:

Proposition 2.1. *Let $G = (\mathcal{V}, \mathcal{E}, m, \theta)$ be a weighted graph. One has that $\Delta_{\mathcal{E}, \theta}$ is bounded if and only if $d_G(Q)$ is bounded.*

Proof. First, (1.4) gives one direction. On the other hand, $\langle \mathbf{1}_{\{x\}}, \Delta_{\mathcal{E}, \theta} \mathbf{1}_{\{x\}} \rangle = d_G(x)$, for all $x \in \mathcal{V}$. \square

Essential self-adjointness of the discrete Laplacian was proved in many situations by Jørgensen (see [Jor] and references therein). In [Jor], he says that every discrete Laplacian is essentially self-adjoint on simple graphs. The proof of [Jor] was incomplete (see [JP]). It was fixed later in [JP]. In the meantime, [Woj] proves this fact. An alternative proof can be found in [Web] where one uses the maximum principle. Similar ideas are found in [KL2], where one generalizes this fact to some weighted graphs by studying Dirichlet forms. Then come the works of [Tor, Ma] for weighted graphs which are metrically complete, see also [CTT] for the non-metrically complete case. Some other criteria, based on commutators, are given in [GS] (see also [Gol]). Finally for the magnetic case, we mention the works [CTT2, Mil, Mil2]. We point out that in the older work of [Aom] one gives some characterization of possible self-adjoint extensions of a weighted discrete Laplacian in the limit point/circle spirit in the case of trees. More recently, in [GS], the question of the deficiency indices is discussed. We point out that, in the latter, one can consider potentials that tend to $-\infty$.

We now improve a self-adjointness criteria given in [KL2] and extend it in two directions: we allow magnetic operators and potentials that are unbounded from below.

Proposition 2.2. *Let $G = (\mathcal{V}, \mathcal{E}, m, \theta)$ be a weighted graph and $V : \mathcal{V} \rightarrow \mathbb{R}$. Take $\lambda \in \mathbb{R}$ so that*

$$(2.1) \quad \{x \in \mathcal{V}, \lambda + d_G(x) + V(x) = 0\} = \emptyset$$

Suppose that, for any $(x_n)_{n \in \mathbb{N}} \in \mathcal{V}^{\mathbb{N}}$, such that the weight $\mathcal{E}(x_n, x_{n+1}) > 0$ for all $n \in \mathbb{N}$, the property

$$(2.2) \quad \sum_{n \in \mathbb{N}} m^2(x_n) a_n^2 = \infty, \text{ where } a_n := \prod_{i=0}^{n-1} \left(1 + \frac{\lambda + V(x_i)}{d_G(x_i)}\right)$$

holds true. Then, the operator $H := \Delta_{\mathcal{E}, \theta} + V(Q)$ is essentially self-adjoint on $\mathcal{C}_c(\mathcal{V})$.

First, note it is always possible to find a λ fulfilling (2.1), as $\#\{d_G(x) + V(x), x \in \mathcal{V}\} \leq \aleph$. In [KL2], the hypothesis is stronger, i.e., they take $a_n = 1$, do not consider magnetic fields, and consider potentials that are bounded from below. Our technique relies on an improvement of [Woj, Theorem 1.3.1].

Proof. Let $f \in \mathcal{D}(H^*) \setminus \{0\}$ such that $H^*f + (i + \lambda)f = 0$ or $H^*f + (-i + \lambda)f = 0$. We get easily:

$$|f(x)| \leq \frac{1}{m^2(x)} \sum_{y \in \mathcal{V}} \frac{\mathcal{E}(x, y)}{|\lambda + d_G(x) + V(x)|} |f(y)|.$$

We derive:

$$(2.3) \quad |f(y)| \geq \left(1 + \frac{\lambda + V(x)}{d_G(x)}\right) |f(x)|, \text{ for all } x, y \in \mathcal{V}, \text{ so that } \mathcal{E}(x, y) \neq 0.$$

Now, since $f \neq 0$, there is $x_0 \in \mathcal{V}$ so that $f(x_0) \neq 0$. Therefore, inductively, we obtain a sequence $(x_n)_{n \in \mathbb{N}} \in \mathcal{V}^{\mathbb{N}}$ such that $\mathcal{E}(x_n, x_{n+1}) > 0$, for all $n \in \mathbb{N}$, and so that (2.3) holds for $y = x_{n+1}$ and $x = x_n$. Hence, we get

$$\sum_{n=0}^N m^2(x_n) |f(x_n)|^2 \geq \sum_{i=0}^N \prod_{j=0}^{i-1} \left(1 + \frac{\lambda + V(x_j)}{d_G(x_j)}\right)^2 |f(x_0)|^2.$$

By letting N go to infinity and remembering (2.2), we obtain a contradiction to the fact that $f \in \ell^2(\mathcal{V}, m^2)$. We conclude with the help of [RS, Theorem X.1]. \square

2.3. Min-max principle. We recall some well-known results. We refer to [RS][Chapter XIII.1] for more details and to [RS][Chapter XIII.15] for more applications. We start with the form-version of the standard variational characterization of the n -th eigenvalue.

Theorem 2.1. *Let A be a non-negative self-adjoint operator with form-domain $\mathcal{Q}(A)$. For all $n \geq 1$, we define:*

$$\mu_n(A) := \sup_{\varphi_1, \dots, \varphi_n} \inf_{\psi \in [\varphi_1, \dots, \varphi_n]^\perp} \langle \psi, A\psi \rangle,$$

where $[\varphi_1, \dots, \varphi_n]^\perp = \{\psi \in \mathcal{D}(A), \text{ so that } \|\psi\| = 1 \text{ and } \langle \psi, \varphi_i \rangle = 0, \text{ with } i = 1, \dots, n\}$. Note that φ_i are not required to be linearly independent.

If μ_n is (strictly) below the essential spectrum of A , it is the n -th eigenvalue, counted with multiplicity, and we have that:

$$\dim \text{Ran } \mathbf{1}_{[0, \mu_n(A)]}(A) = n.$$

Otherwise, μ_n is the infimum of the essential spectrum. Moreover, $\mu_j = \mu_n$, for all $j \geq n$ and there are at most $n - 1$ eigenvalues, counted with multiplicity, below the essential spectrum. In that case,

$$\dim \text{Ran } \mathbf{1}_{[0, \mu_n(A) + \varepsilon]}(A) = +\infty, \text{ for all } \varepsilon > 0.$$

Remark 2.1. One has a priori no control on the multiplicity of the (possible) eigenvalue which is at the bottom of the essential spectrum.

This ensures the useful following criteria:

Proposition 2.3. Let A, B be two self-adjoint operators, with form-domains $\mathcal{D}(A)$ and $\mathcal{D}(B)$, respectively. Suppose that

$$\mathcal{D}(A) \supset \mathcal{D}(B) \text{ and } 0 \leq \langle \psi, A\psi \rangle \leq \langle \psi, B\psi \rangle,$$

for all $\psi \in \mathcal{D}(B)$. Then one has $\inf \sigma_{\text{ess}}(A) \leq \inf \sigma_{\text{ess}}(B)$ and

$$(2.4) \quad \mathcal{N}_\lambda(A) \geq \mathcal{N}_\lambda(B), \text{ for } \lambda \in [0, \infty) \setminus \{\inf \sigma_{\text{ess}}(B)\},$$

where $\mathcal{N}_\lambda(A) := \dim \text{Ran } \mathbf{1}_{[0, \lambda]}(A)$.

In particular, if A and B have the same form-domain, then $\sigma_{\text{ess}}(A) = \emptyset$ if and only if $\sigma_{\text{ess}}(B) = \emptyset$.

Proof. It is enough to notice that $\mu_n(A) \leq \mu_n(B)$, for all $n \geq 0$. Theorem 2.1 permits us to conclude for the first part. Supposing now they have the same form-domain, by uniform boundedness principle, there is $a, b > 0$ such that:

$$\langle \psi, |A|\psi \rangle \leq a \langle \psi, |B|\psi \rangle + b \|\psi\|^2 \text{ and } \langle \psi, |B|\psi \rangle \leq a \langle \psi, |A|\psi \rangle + b \|\psi\|^2$$

for all $\psi \in \mathcal{D}(|A|) = \mathcal{D}(|B|)$. By using twice the previous statement we get the result. \square

Because of Remark 2.1, we stress that we have to remove “ $\inf \sigma_{\text{ess}}(B)$ ” in (2.4), as the min-max principle does not decide whether or not the B has an eigenvalue at this energy.

We now turn to a criteria of compactness. We recall that a compact and self-adjoint operator A is in the p -Schatten class, for $p \in [1, \infty)$, if $\text{tr}(|A|^p) < \infty$.

Proposition 2.4. Let \mathcal{H} be a Hilbert space and A, B be two bounded self-adjoint operators. Suppose that B is compact and

$$|\langle f, Af \rangle| \leq \langle f, Bf \rangle, \text{ for all } f \in \mathcal{H}.$$

Then A is also compact. Moreover, given $p \in [1, \infty)$, if B is in the p -Schatten class, so is A .

Proof. First note that $-B \leq A$ and that $-B \leq -A$ in the form sense. By Proposition 2.3, this implies that the essential spectrum of A is $\{0\}$. Finally, as A is self-adjoint, we infer that A is compact. Now assume that B is in the p -Schatten class. The min-max principle, applied to A and $-A$, guarantees that $\text{tr}(|A|^p) \leq 2^p \text{tr}(B^p)$. \square

Using (1.4), we give a straightforward consequence of Propositions 2.3 and 2.4:

Corollary 2.1. Let $G = (\mathcal{V}, \mathcal{E}, m, \theta)$ be a weighted graph. One has:

$$\inf \sigma_{\text{ess}}(\Delta_{\mathcal{E}, \theta}) \leq 2 \inf \sigma_{\text{ess}}(d_G(Q)) \text{ and } \mathcal{N}_\lambda(\Delta_{\mathcal{E}, \theta}) \geq \mathcal{N}_\lambda(2d_G(Q)),$$

for all $\lambda \in [0, \infty) \setminus \{\inf \sigma_{\text{ess}}(d_G(Q))\}$. In particular, if $0 \in \sigma_{\text{ess}}(d_G(Q))$, then $0 \in \sigma_{\text{ess}}(\Delta_{\mathcal{E}, \theta})$ and if $\Delta_{\mathcal{E}, \theta}$ has compact resolvent then, $d_G(Q)$ has also compact resolvent. In other words, one has that $\sigma_{\text{ess}}(d_G(Q)) \neq \emptyset$, then $\sigma_{\text{ess}}(\Delta_{\mathcal{E}, \theta}) \neq \emptyset$.

Moreover, if $d_G(Q)$ is compact, then $\Delta_{\mathcal{E}, \theta}$ is also compact.

We mention, that for a simple graph, $d_G(Q) \geq 1$ and it is not a compact operator.

3. SURROUNDING THE LAPLACIAN

3.1. A Hardy inequality. We start with a remark about bi-partite graphs:

Proposition 3.1. *Given a bi-partite graph $G = (\mathcal{V}, \mathcal{E}, m, \theta)$ and a function $V : \mathcal{V} \rightarrow [0, \infty)$, the following assertions are equivalent:*

$$(3.1) \quad \langle f, (d_G(Q) - V(Q))f \rangle \leq \langle f, \Delta_{\mathcal{E}, \theta} f \rangle, \quad \text{for } f \in \mathcal{C}_c(\mathcal{V}),$$

$$(3.2) \quad \langle f, \Delta_{\mathcal{E}, \theta} f \rangle \leq \langle f, (d_G(Q) + V(Q))f \rangle, \quad \text{for } f \in \mathcal{C}_c(\mathcal{V}),$$

$$(3.3) \quad |\langle f, \mathcal{A}_{\mathcal{E}, \theta} f \rangle| \leq \langle f, Vf \rangle, \quad \text{for } f \in \mathcal{C}_c(\mathcal{V}),$$

where $\mathcal{A}_{\mathcal{E}, \theta}$ is the magnetic adjacency matrix defined by

$$(3.4) \quad (\mathcal{A}_{\mathcal{E}, \theta} f)(x) := \frac{1}{m^2(x)} \sum_{y \in \mathcal{V}} \mathcal{E}(x, y) e^{i\theta_{x,y}} f(y),$$

for $f \in \mathcal{C}_c(\mathcal{V})$ and $x \in \mathcal{V}$.

Proof. Let $x_0 \in \mathcal{V}$. Given $x \in \mathcal{V}$, by definition of connectedness, there is a sequence x_1, \dots, x_n of elements of \mathcal{V} , so that $\mathcal{E}(x_i, x_{i+1}) \neq 0$, for $i = 0, \dots, n-1$ and $x_n = x$. The minimal possible n is the *unweighted length* between x_0 and x . We denote it by $\ell(x)$. Set $Uf(x) := (-1)^{\ell(x)} f(x)$. Note that $U^2 = \text{Id}$ and $U^{-1} = U^* = U$. Notice now that

$$U^{-1} \mathcal{A}_{\mathcal{E}, \theta} U = -\mathcal{A}_{\mathcal{E}, \theta} \text{ and } \Delta_{\mathcal{E}, \theta} = d_G(Q) - \mathcal{A}_{\mathcal{E}, \theta}.$$

We start with (3.1) and rewrite it as follows: $\langle f, \mathcal{A}_{\mathcal{E}, \theta} f \rangle \leq \langle f, Vf \rangle$, for $f \in \mathcal{C}_c(\mathcal{V})$. Applying this to Uf , we infer immediately (3.3). We start now from (3.3). We get:

$$\langle f, \Delta_{\mathcal{E}, \theta} f \rangle = \langle f, (d_G(Q) - \mathcal{A}_{\mathcal{E}, \theta})f \rangle \geq \langle f, (d_G(Q) - V(Q))f \rangle$$

for $f \in \mathcal{C}_c(\mathcal{V})$. In the same way, we have that (3.2) is equivalent to (3.3). \square

We now turn to the key estimate of the paper and prove Proposition 1.1. First we mention that a Laplacian in a certain $\ell^2(\mathcal{V}, m^2)$ is unitarily equivalent to a Schrödinger operator in any other $\ell^2(\mathcal{V}, m_0^2)$. This has already been noticed before, e.g., [CTT].

Proposition 3.2. *Let $G = (\mathcal{V}, \mathcal{E}, m, \theta)$ be a weighted graph and $m_0 : \mathcal{V} \rightarrow (0, \infty)$ be a weight. Then the Friedrichs extension of $\Delta_{\mathcal{E}_0, \theta}$, acting in $\ell^2(\mathcal{V}, m_0^2)$, is unitarily equivalent to that of $\Delta_{\mathcal{E}, \theta} + V(Q)$, in $\ell^2(\mathcal{V}, m^2)$, where*

$$\mathcal{E}(x, y) := \mathcal{E}_0(x, y) \frac{m(x)m(y)}{m_0(x)m_0(y)} \quad \text{and} \quad V(x) := \frac{1}{m_0^2(x)} \sum_{y \in \mathcal{V}} \mathcal{E}(x, y) \left(1 - \frac{m(y)m_0(x)}{m(x)m_0(y)} \right).$$

Proof. Consider the unitary map $U : \ell^2(G, m^2) \rightarrow \ell^2(G, m_0^2)$ given by $(Uf)(x) = (m(x)/m_0(x))f(x)$. Straightforwardly, using (1.3), one has

$$(3.5) \quad U^{-1} \Delta_{\mathcal{E}_0, \theta} Uf := (\Delta_{\mathcal{E}, \theta} + V(Q))f$$

for all $f \in \mathcal{C}_c(\mathcal{V})$. Then it holds true for the closures of $\Delta_{\mathcal{E}_0, \theta}|_{\mathcal{C}_c}$ and $(\Delta_{\mathcal{E}, \theta} + V(Q))|_{\mathcal{C}_c}$. Note now that (3.5) also holds for the adjoints of the last two operators. Therefore, the Friedrichs extensions $\Delta_{\mathcal{E}_0, \theta}$ and $\Delta_{\mathcal{E}, \theta} + V(Q)$ are unitarily equivalent. \square

The novelty in the Hardy inequality (1.6) is more in the point of view. Rather than studying $\Delta_{\mathcal{E}, \theta}$ in $\ell^2(\mathcal{V}, m^2)$ with the help of a simpler $\ell^2(\mathcal{V}, m_0^2)$ (where typically $m_0 = 1$), we study it with the help of all other weighted spaces. The applications we consider in the paper are also new.

Proof of Proposition 1.1. On $\ell^2(G, m^2)$, we consider the quadratic form:

$$\mathcal{Q}_{\mathcal{E}, \theta}(f, f) := \frac{1}{2} \sum_{x, y \in \mathcal{V}} \tilde{\mathcal{E}}(x, y) |f(x) - e^{i\theta_{x,y}} f(y)|^2 \geq 0, \text{ for } f \in \mathcal{C}_c(\mathcal{V}),$$

where

$$\tilde{\mathcal{E}}(x, y) := \mathcal{E}(x, y) \frac{m(x)m(y)}{m_0(x)m_0(y)}.$$

Note that $\mathcal{Q}(f, f) = \langle f, H_{m^2} f \rangle_{\ell^2(\mathcal{V}, m^2)}$, where

$$(H_{m^2} f)(x) = \frac{1}{m^2(x)} \sum_{y \sim x} \tilde{\mathcal{E}}(x, y) (f(x) - e^{i\theta_{x,y}} f(y)), \text{ for } f \in \mathcal{C}_c(\mathcal{V}).$$

Consider now the unitary map $U : \ell^2(G, m^2) \rightarrow \ell^2(G, m_0^2)$ given by $(Uf)(x) = (m(x)/m_0(x))f(x)$. Set $H := UH_{m^2}U^{-1}$. We have:

$$\begin{aligned} (Hf)(x) &= \frac{1}{m^2(x)} \sum_{y \in \mathcal{V}} \tilde{\mathcal{E}}(x, y) f(x) - \frac{1}{m_0^2(x)} \sum_{y \in \mathcal{V}} \mathcal{E}(x, y) e^{i\theta_{x,y}} f(y) \\ &= \frac{1}{m_0^2(x)} \sum_{y \in \mathcal{V}} \mathcal{E}(x, y) (f(x) - e^{i\theta_{x,y}} f(y)) - V_m(x) f(x). \end{aligned}$$

Since $H \geq 0$, we obtain (1.6). The rest of the statement is given by Proposition 3.1. \square

In the paper, we are interested in minorating the Laplacian with the help of the weighted degree. Before proving the main result for simple trees in the next section, we start by giving some negative answers. First, it is obvious that one cannot find a non-negative V_m in (1.6) if the graph is finite. This continues to hold true for some infinite graphs.

Proposition 3.3. *Let $G = (\mathcal{V}, \mathcal{E}, m, \theta)$ be a weighted graph and take $V : \mathcal{V} \rightarrow \mathbb{R}$ so that $0 \leq V \leq \Delta_{\mathcal{E}, \theta}$, in the form sense on $\mathcal{C}_c(\mathcal{V})$. If the constant function 1 is in $\mathcal{D}((\Delta_{\mathcal{E}, \theta})^{1/2})$ then $V = 0$.*

If one supposes that $1 \in \ell^2(\mathcal{V}, m^2)$ and $\Delta_{\mathcal{E}, \theta}$ is essentially self-adjoint on $\mathcal{C}_c(\mathcal{V})$, then 1 is in $\mathcal{D}((\Delta_{\mathcal{E}, \theta})^{1/2})$, as $1 \in \mathcal{D}((\Delta_{\mathcal{E}, \theta}|_{\mathcal{C}_c})^*)$.

Proof. By construction of the Friedrichs extension, there is $f_n \in \mathcal{C}_c(\mathcal{V})$, so that f_n tends to 1 in the graph norm of $\Delta_{\mathcal{E}, \theta}^{1/2}$. Moreover,

$$0 \leq \sum_{x \in \mathcal{V}} m^2(x) V(x) |f_n(x)|^2 = \langle f_n, V f_n \rangle \leq \langle f_n, \Delta_{\mathcal{E}, \theta} f_n \rangle \rightarrow 0,$$

as n goes to ∞ . Now since $V^{1/2}(Q)$ is closed, the closed graph theorem gives that the l.h.s. tends to $\sum_{x \in \mathcal{V}} m^2(x) V(x)$. Therefore $V = 0$. \square

We now treat the impossibility of minorating by using the weighted degree. We say that $\{K_n\}_{n \in \mathbb{N}}$ is a *filtration* of a graph $(\mathcal{V}, \mathcal{E}, m, \theta)$ if $K_n \subset \mathcal{V}$ is finite, $K_n \subset K_{n+1}$, for all $n \in \mathbb{N}$, and such that $\bigcup_{n \in \mathbb{N}} K_n = \mathcal{V}$. Recall that given a subset $K \in \mathcal{V}$, $\partial K := \{x \in K, \text{ there is } y \sim x \text{ such that } y \notin K\}$.

Proposition 3.4. *Let $G = (\mathcal{V}, \mathcal{E}, m, 0)$ be a weighted graph, such that $\mathcal{E} \neq \emptyset$.*

(a) *Assume that $\alpha(G) = 0$. Then, there is no $a > 0$ such that*

$$\langle f, d_G(Q)f \rangle \leq a \langle f, \Delta_{\mathcal{E}, 0} f \rangle, \text{ for } f \in \mathcal{C}_c(\mathcal{V}).$$

(b) *Suppose that there is a filtration $\{K_n\}_{n \in \mathbb{N}}$ associated with G , such that:*

$$(3.6) \quad \limsup_{n \rightarrow +\infty} \frac{\sum_{x \in K_n \setminus \partial K_n} m^2(x) d_G(x)}{\sum_{x \in K_n} m^2(x)} = +\infty,$$

then there is no $a \in [0, 1]$ and $b \geq 0$ such that

$$(3.7) \quad \langle f, d_G(Q)f \rangle \leq a \langle f, \Delta_{\mathcal{E}, 0} f \rangle + b \|f\|^2, \text{ for all } f \in \mathcal{C}_c(\mathcal{V}).$$

In particular, (3.6) is fulfilled, if $1 \in \ell^1(\mathcal{V}, m^2)$ and $\sup_{x \in \mathcal{V}} m^2(x) d_G(x) = +\infty$.

Proof. Take $K \subset \mathcal{V}$ finite. Using (1.2), note that

$$(3.8) \quad \langle \mathbf{1}_K, d_G(Q) \mathbf{1}_K \rangle = \sum_{x \in K} m^2(x) d_G(x) \text{ and } \langle \mathbf{1}_K, \Delta_{\mathcal{E}, 0} \mathbf{1}_K \rangle \leq \sum_{x \in \partial K} m^2(x) d_G(x).$$

The first part follows from the definition of $\alpha(G)$.

For the second assertion, consider $K = K_n$ and note that (3.7) implies that

$$\sum_{x \in K_n \setminus \partial K_n} m^2(x) d_G(x) \leq b \sum_{x \in K_n} m^2(x).$$

By letting n go to infinity, we obtain a contradiction with (3.6). \square

3.2. The case of trees. We now turn to a minoration of the magnetic Laplacian and present it in the context of weighted trees. Perturbation theory will be considered in Section 5. We fix some notation.

We first recall that a *tree* is a connected graph $G = (\mathcal{E}, \mathcal{V}, m, \theta)$ such that for each edge $e \in \mathcal{V} \times \mathcal{V}$ with $\mathcal{E}(e) \neq 0$ the graph $(\tilde{\mathcal{E}}, \mathcal{V}, m, \theta|_{\tilde{\mathcal{E}}})$, with $\tilde{\mathcal{E}} := \mathcal{E} \times 1_{\{e\}^c}$, i.e., with the edge e removed, is disconnected. It is convenient to choose a root in the tree. Due to its structure, one can take any point of \mathcal{V} . We denote it by ϵ . We define inductively the *spheres* S_n by $S_{-1} = \emptyset$, $S_0 := \{\epsilon\}$, and $S_{n+1} := \mathcal{N}_G(S_n) \setminus S_{n-1}$. Given $n \in \mathbb{N}$, $x \in S_n$, and $y \in \mathcal{N}_G(x)$, one sees that $y \in S_{n-1} \cup S_{n+1}$. We write $x \rightsquigarrow y$ and say that x is a *son* of y , if $y \in S_{n-1}$, while we write $x \prec y$ and say that x is a *father* of y , if $y \in S_{n+1}$. Notice that ϵ has no father. Given $x \neq \epsilon$, note that there is a unique $y \in V$ with $x \rightsquigarrow y$, i.e., everyone apart from ϵ has one and only one father. We denote the father of x by \overleftarrow{x} .

Proof of Theorem 1.1. We focus on the left hand side of (1.10) for G_\circ , i.e., for $H = \Delta_{\mathcal{E}_\circ, \theta_\circ}$. Set $K := \{x \in \mathcal{V}, d_{G_\circ}(x) \leq 1\}$ and $\tilde{\mathcal{V}} := \mathcal{V} \setminus K$. Note that induced graph \tilde{G}_\circ of G_\circ w.r.t. $\tilde{\mathcal{V}} = \cup_{n \in \mathbb{N}} \tilde{\mathcal{V}}_n$ is a forest of disjoint and maximal subtrees $\tilde{G}_n := (\mathcal{E}_n, \tilde{\mathcal{V}}_n, m|_{\tilde{\mathcal{V}}_n}, \theta_\circ|_{\tilde{\mathcal{E}}_n})$, where $\tilde{\mathcal{E}}_n := \mathcal{E}_\circ|_{\tilde{\mathcal{V}}_n \times \tilde{\mathcal{V}}_n}$. For each \tilde{G}_n of \tilde{G}_\circ , we define:

$$\tilde{m}(\epsilon_{\tilde{\mathcal{V}}_n}) := 1 \text{ and } \tilde{m}(x) := \eta \tilde{m}(\overleftarrow{x}) \frac{m(x)}{m(\overleftarrow{x})} d_{G_\circ}^{-\varepsilon_0/2}(x), \text{ for all } x \in \tilde{\mathcal{V}}_n.$$

With this definition and $V_{\tilde{m}}$ as in (1.7), we obtain:

$$\begin{aligned} \frac{V_{\tilde{m}}(x)}{d_{\tilde{G}_\circ}(x)} &= 1 - \frac{1}{d_G(x)m^2(x)} \left(\mathcal{E}(\overleftarrow{x}, x) \frac{\tilde{m}(\overleftarrow{x})}{\tilde{m}(x)} \frac{m(x)}{m(\overleftarrow{x})} + \sum_{y \rightsquigarrow x} \mathcal{E}(y, x) \frac{\tilde{m}(y)}{\tilde{m}(x)} \frac{m(x)}{m(y)} \right) \\ &= 1 - \frac{1}{\eta d_G^{(1-\varepsilon_0/2)}(x)m^2(x)} \mathcal{E}(\overleftarrow{x}, x) - \frac{\eta}{d_G(x)m^2(x)} \sum_{y \rightsquigarrow x} \mathcal{E}(y, x) d_G^{-\varepsilon_0/2}(y) \\ (3.9) \quad &\geq 1 - \eta - \frac{1}{\eta} d_G^{-\varepsilon_0/2}(x) C_0, \end{aligned}$$

for all $x \in \tilde{\mathcal{V}}_n$. Consider $f \in \mathcal{C}_c(\mathcal{V})$. We apply Proposition 1.1 to \tilde{G} . Since $\mathcal{E} \geq 0$ we derive that there is a constant $C(\eta) > 0$, so that:

$$\begin{aligned} \langle f, \Delta_{\mathcal{E}_\circ, \theta_\circ} f \rangle &\geq \langle f|_{\tilde{\mathcal{V}}}, \Delta_{\mathcal{E}_\circ, \theta_\circ} (f|_{\tilde{\mathcal{V}}}) \rangle \geq \langle f|_{\tilde{\mathcal{V}}}, V_{\tilde{m}}(Q) (f|_{\tilde{\mathcal{V}}}) \rangle \\ &\geq \langle d_G^{1/2}(Q) f|_{\tilde{\mathcal{V}}}, (1 - \eta - C_0/d_G^{\varepsilon_0/2}(Q)\eta) d_G^{1/2}(Q) f|_{\tilde{\mathcal{V}}} \rangle \\ &\geq (1 - 2\eta) \langle f|_{\tilde{\mathcal{V}}}, d_{\tilde{G}}(Q) f|_{\tilde{\mathcal{V}}} \rangle - C(\eta) \|f|_{\tilde{\mathcal{V}}}\|^2 \\ (3.10) \quad &\geq (1 - 2\eta) \langle f, d_{G_\circ}(Q) f \rangle - (C_\eta + 1) \|f\|^2. \end{aligned}$$

This gives (1.10). The reverse inequality is obtained by applying Proposition 3.1.

Next, the equality of the domains of the forms follow immediately and so the essential spectrum of $\Delta_{\mathcal{E}_\circ, \theta_\circ}$ is empty if and only if $\lim_{|x| \rightarrow \infty} d_G(x) = +\infty$. Finally, we use twice Proposition 2.3 with the double inequality (1.10). This yields:

$$1 - \varepsilon \leq \liminf_{\lambda \rightarrow \infty} \frac{\mathcal{N}_\lambda(\Delta_{\mathcal{E}, \theta})}{\mathcal{N}_\lambda(d_G(Q))} \leq \limsup_{\lambda \rightarrow \infty} \frac{\mathcal{N}_\lambda(\Delta_{\mathcal{E}, \theta})}{\mathcal{N}_\lambda(d_G(Q))} \leq 1 + \varepsilon.$$

By letting ε go to zero we obtain the Weyl asymptotic (1.11) for $\Delta_{\mathcal{E}_\circ, \theta_\circ}$. It remains to use Proposition 5.1 to obtain the result for $\Delta_{\mathcal{E}, \theta}$. \square

Remark 3.1. *It is important to note that ε has to be positive in (1.10). For instance, on a simple tree, if ε was equal to 0, this would imply that the magnetic adjacency matrix was bounded (see Proposition 3.1). Finally note that, by considering delta functions as test functions, the magnetic adjacency matrix is bounded if and only if the weighted degree is bounded.*

Remark 3.2. *We could not provide an example of Schrödinger operator for which we can compute the asymptotic of eigenvalues and which is not essentially self-adjoint on $\mathcal{C}_c(\mathcal{V})$.*

4. COMPARISON OF DOMAINS

4.1. From form-domain to domain. Once one knows that the form-domains are equal, the next question is to guarantee that the domains are equal. We present an approach by commutators. Some subtleties arise since we have to deal with the square root of the Laplacian.

We first prove the invariance of the domain under the C_0 -group $\{e^{itd_G(Q)}\}_{t \in \mathbb{R}}$ and the \mathcal{C}^1 regularity (see Appendix A for a definition).

Proposition 4.1. *Let $G = (\mathcal{V}, \mathcal{E}, m, \theta)$ be a weighted graph, such that $\Delta_{\mathcal{E}, \theta}$ is essentially self-adjoint on $\mathcal{C}_c(\mathcal{V})$, $\mathcal{D}(\Delta_{\mathcal{E}, \theta}^{1/2}) = \mathcal{D}(d_G^{1/2}(Q))$, and*

$$(4.1) \quad \sup_{x \in \mathcal{V}} \left(\frac{1}{m^2(x)} \sum_{y \in \mathcal{V}} \frac{\mathcal{E}(x, y)}{\min(\langle d_G(x) \rangle^{1/2}, \langle d_G(y) \rangle^{1/2})} |d_G(x) - d_G(y)| \right) < \infty.$$

Then, one has that the invariance $e^{itd_G(Q)}\mathcal{D}(\Delta_{\mathcal{E}, \theta}) \subset \mathcal{D}(\Delta_{\mathcal{E}, \theta})$, for all $t \in \mathbb{R}$. Moreover, we have that $\Delta_{\mathcal{E}, \theta} \in \mathcal{C}^1(d_G(Q))$ and that $[\Delta_{\mathcal{E}, \theta}, d_G(Q)]|_{\mathcal{C}_c(\mathcal{V})}$ extends to $[\Delta_{\mathcal{E}, \theta}, d_G(Q)]_o \in \mathcal{B}(\mathcal{D}(\Delta_{\mathcal{E}, \theta}^{1/2}), \ell^2(\mathcal{V}, m^2))$.

Proof. We denote by $[\Delta_{\mathcal{E}, \theta}, e^{itd_G(Q)}]_o$ the closure of the commutator $[\Delta_{\mathcal{E}, \theta}, e^{itd_G(Q)}]|_{\mathcal{C}_c(\mathcal{V})}$ and by c_0 the constant appearing in (4.1). We have that:

$$(4.2) \quad \begin{aligned} |\langle f, [\Delta_{\mathcal{E}, \theta}, e^{itd_G(Q)}]_o g \rangle| &= \sum_{x \in \mathcal{V}} \sum_{y \in \mathcal{V}} \overline{f(x)} \mathcal{E}(x, y) \left(e^{itd_G(x)} - e^{itd_G(y)} \right) g(y) \\ &\leq \frac{|t|}{2} \sum_{x \in \mathcal{V}} \sum_{y \in \mathcal{V}} \frac{\mathcal{E}(x, y)}{\langle d_G(y) \rangle^{1/2}} |d_G(x) - d_G(y)| \left(|f(x)|^2 + |\langle d_G(Q) \rangle^{1/2} g(y)|^2 \right) \\ &\leq c_0 |t| \left(\|f\|^2 + \left\| \langle d_G(Q) \rangle^{1/2} g \right\|^2 \right), \end{aligned}$$

for all $f, g \in \mathcal{C}_c(\mathcal{V})$. Therefore, $[\Delta_{\mathcal{E}, \theta}, e^{itd_G(Q)}]_o$ is bounded from $\mathcal{D}(\Delta_{\mathcal{E}, \theta}^{1/2})$ to $\ell^2(\mathcal{V}, m^2)$. On the other hand, we get:

$$\begin{aligned} \|\Delta_{\mathcal{E}, \theta} e^{itd_G(Q)} f\| &\leq \|e^{itd_G(Q)} \Delta_{\mathcal{E}, \theta} f\| + \|[\Delta_{\mathcal{E}, \theta}, e^{itd_G(Q)}] \langle d_G(Q) \rangle^{-1/2} \langle d_G(Q) \rangle^{1/2} f\| \\ &\leq C (\|f\| + \|\Delta_{\mathcal{E}, \theta} f\|), \end{aligned}$$

for all $f \in \mathcal{C}_c(\mathcal{V})$. By density we obtain the invariance of the domain.

Next, thanks to (4.2), we obtain that

$$\liminf_{t \rightarrow 0^+} \left\| [\Delta_{\mathcal{E}, \theta}, e^{itd_G(Q)}] / t \right\|_{\mathcal{B}(\mathcal{D}(\Delta_{\mathcal{E}, \theta}^{1/2}), \ell^2(\mathcal{V}, m^2))} < \infty.$$

Therefore, Theorem A.2 yields that $\Delta_{\mathcal{E}, \theta} \in \mathcal{C}^1(d_G(Q))$. Finally, by estimating as in (4.2), we obtain that $[\Delta_{\mathcal{E}, \theta}, d_G(Q)]_o$, belongs $\mathcal{B}(\mathcal{D}(\Delta_{\mathcal{E}, \theta}^{1/2}), \ell^2(\mathcal{V}, m^2))$. \square

We turn to the central result of this section.

Theorem 4.1. *Let $G = (\mathcal{V}, \mathcal{E}, m, \theta)$ be a weighted graph, such that $\mathcal{D}(\Delta_{\mathcal{E}, \theta}^{1/2}) = \mathcal{D}(d_G^{1/2}(Q))$, $\Delta_{\mathcal{E}, \theta}$ is essentially self-adjoint on $\mathcal{C}_c(\mathcal{V})$, and there is $\varepsilon > 0$ such that*

$$(4.3) \quad \sup_{x \in \mathcal{V}} \left(\frac{1}{m^2(x)} \sum_{y \in \mathcal{V}} \frac{\mathcal{E}(x, y)}{\min(\langle d_G(x) \rangle^{1/2-\varepsilon}, \langle d_G(y) \rangle^{1/2-\varepsilon})} |d_G(x) - d_G(y)| \right) < \infty.$$

Take also a potential $V : \mathcal{V} \rightarrow \mathbb{R}$, such that:

$$(4.4) \quad \lim_{|x| \rightarrow +\infty} \frac{V(x)}{d_G(x) + 1} = 0.$$

Then, $\mathcal{D}(\Delta_{\mathcal{E}, \theta} + V(Q)) = \mathcal{D}(d_G(Q))$ and $\sigma_{\text{ess}}(\Delta_{\mathcal{E}, \theta} + V(Q)) = \sigma_{\text{ess}}(\Delta_{\mathcal{E}, \theta})$.

Remark 4.1. *We point out that the hypothesis on the essential self-adjointness is also necessary, as $d_G(Q)$ is essentially self-adjoint on $\mathcal{C}_c(\mathcal{V})$.*

Remark 4.2. *By taking $\varepsilon = 0$ in (4.3), i.e., under the same hypothesis as Proposition 4.1, the proof ensures only that $\mathcal{D}(d_G(Q)) \subset \mathcal{D}(\Delta_{\mathcal{E}, \theta})$.*

Proof. The stability of the essential spectrum is ensured by Proposition 5.2. We focus on the domain. By Kato-Rellich's Theorem, e.g. [RS, Theorem X.12], it is enough to consider $V = 0$. We assume for the moment that

$$M := (\Delta_{\mathcal{E}, \theta} + 1)^{1/2} [d_G(Q), (\Delta_{\mathcal{E}, \theta} + 1)^{1/2}]_o + [\Delta_{\mathcal{E}, \theta}, (d_G(Q) + 1)^{1/2}]_o (d_G(Q) + 1)^{1/2}.$$

is a bounded operator from $\mathcal{D}(d_G^{1-\varepsilon}(Q))$ to $\mathcal{D}(d_G^{1-\varepsilon}(Q))^*$. The 1 is here to make the square root smooth over the spectrum. Since the form-domain are equal of $\Delta_{\mathcal{E}, \theta}$ and $d_G(Q)$, we have that there is $a, b > 0$ so that (3.7) holds true. By using twice (3.7) and working in the form sense on $\mathcal{C}_c(\mathcal{V})$, we infer:

$$\begin{aligned}
(\Delta_{\varepsilon,\theta} + 1)^2 &\leq a(\Delta_{\varepsilon,\theta} + 1)^{1/2}(d_G(Q) + 1)(\Delta_{\varepsilon,\theta} + 1)^{1/2} + (b + 1)(\Delta_{\varepsilon,\theta} + 1) \\
&= a(d_G(Q) + 1)^{1/2}\Delta_{\varepsilon,\theta}(d_G(Q) + 1)^{1/2} + aM + (b + 1)(\Delta_{\varepsilon,\theta} + 1) \\
(4.5) \qquad \qquad &\leq a'd_G^2(Q) + b',
\end{aligned}$$

for some $a', b' > 0$. By using the fact that $\varepsilon > 0$, the reverse inequality holds true for the same reasons for some $a' \in (0, 1)$ and $b' > 0$. Therefore, the domains are equal. It remains to prove the boundedness of M . We start with the r.h.s. term. Let c_1 be the constant in (4.3). We estimate as above:

$$\begin{aligned}
&|\langle f, [\Delta_{\varepsilon,\theta}, (d_G(Q) + 1)^{1/2}]_o \langle d_G(Q) \rangle^{1/2} g \rangle| \\
&= \sum_{x \in \mathcal{V}} \sum_{y \in \mathcal{V}} \overline{f(x)} \mathcal{E}(x, y) ((d_G(x) + 1)^{1/2} - (d_G(y) + 1)^{1/2}) \langle d_G(y) \rangle^{1/2} g(y) \\
&\leq \frac{1}{2} \sum_{x \in \mathcal{V}} \sum_{y \in \mathcal{V}} \frac{\mathcal{E}(x, y)}{\langle d_G(y) \rangle^{1/2 - \varepsilon}} |d_G(x) - d_G(y)| (|f(x)|^2 + \langle d_G(Q) \rangle^{1 - \varepsilon} |g(y)|^2) \\
&\leq c_1 \left(\|f\|^2 + \|\langle d_G(Q) \rangle^{1 - \varepsilon} g\|^2 \right),
\end{aligned}$$

for all $f, g \in \mathcal{C}_c(\mathcal{V})$. In particular, $[\Delta_{\varepsilon,\theta}, (d_G(Q) + 1)^{1/2}]_o \langle d_G(Q) \rangle^{1/2}$ is bounded from $\mathcal{D}(d_G^{1 - \varepsilon}(Q))$ to $\ell^2(\mathcal{V}, m^2)$.

We turn to the second part of M and use the Appendix B. Let φ be in $\mathcal{S}^{1/2}$ such that $\varphi(x) = \sqrt{x}$, for all $x \geq 1$. Since $\Delta_{\varepsilon,\theta}$ is non-negative, $\varphi(\Delta_{\varepsilon,\theta} + 1) = (\Delta_{\varepsilon,\theta} + 1)^{1/2}$. We cannot use the (B.5) directly with φ as the integral does not seem to exist. We proceed as in [GJ]. Take $\chi_1 \in C_c^\infty(\mathbb{R}; \mathbb{R})$ with values in $[0, 1]$ and being 1 on $[-1, 1]$. Set $\chi_R := \chi(\cdot/R)$. As R goes to infinity, χ_R converges pointwise to 1. Moreover, $\{\chi_R\}_{R \in [1, \infty)}$ is bounded in \mathcal{S}^0 . We infer $\varphi_R := \varphi \chi_R$ tends pointwise to φ and that $\{\varphi_R\}_{R \in [1, \infty)}$ is bounded in $\mathcal{S}^{1/2}$. Now, recalling $\Delta_{\varepsilon,\theta} \in \mathcal{C}^1(d_G(Q))$ and (A.3), we obtain

$$\begin{aligned}
(4.6) \qquad &[\varphi_R(\Delta_{\varepsilon,\theta}), d_G(Q)]_o \langle \Delta_{\varepsilon,\theta} \rangle^{-1/2} = \\
&= \frac{i}{2\pi} \int_{\mathbb{C}} \frac{\partial \varphi_R^{\mathbb{C}}}{\partial \bar{z}} (z - \Delta_{\varepsilon,\theta})^{-1} [d_G(Q), \Delta_{\varepsilon,\theta}]_o (z - \Delta_{\varepsilon,\theta})^{-1} \langle \Delta_{\varepsilon,\theta} \rangle^{-1/2} dz \wedge d\bar{z}. \\
&= \frac{i}{2\pi} \int_{\mathbb{C}} \frac{\partial \varphi_R^{\mathbb{C}}}{\partial \bar{z}} (z - \Delta_{\varepsilon,\theta})^{-1} [d_G(Q), \Delta_{\varepsilon,\theta}]_o \langle \Delta_{\varepsilon,\theta} \rangle^{-1/2} (z - \Delta_{\varepsilon,\theta})^{-1} dz \wedge d\bar{z}.
\end{aligned}$$

By Proposition 4.1, $[d_G(Q), \Delta_{\varepsilon,\theta}]_o \langle \Delta_{\varepsilon,\theta} \rangle^{-1/2}$ is bounded. Moreover, using (B.2) with $l = 2$, we bound the integrand, uniformly in R , by $C \langle x \rangle^{-1+1/2-2} |y|^2 |y|^{-1} |y|^{-1}$, for some constant C . It is integrable on the domain given by (B.4). By Lebesgue domination, the r.h.s. of (4.6) has a limit in norm. Note now that the l.h.s., as form on $\mathcal{C}_c(\mathcal{V})$, tends to the operator $[(\Delta_{\varepsilon,\theta} + 1)^{1/2}, d_G(Q)] \langle \Delta_{\varepsilon,\theta} \rangle^{-1/2}$. This gives that

$$\langle \Delta_{\varepsilon,\theta} \rangle^{1/2} [(\Delta_{\varepsilon,\theta} + 1)^{1/2}, d_G(Q)]_o \leq C(d_G(Q) + 1)$$

in the form sense on $\mathcal{C}_c(\mathcal{V})$ for some constant C . This ensures the announced boundedness of M . \square

Remark 4.3. Assuming also that for all $\varepsilon > 0$, there is $c_\varepsilon \geq 0$ so that (1.10) holds true, one observes that a' can be arbitrary close to 1 in (4.5). Therefore, using again the Kato-Rellich's Theorem, one can weaken (4.4) in Theorem 4.1 and replace it by

$$(4.7) \qquad \limsup_{|x| \rightarrow +\infty} \frac{|V(x)|}{d_G(x) + 1} < 1.$$

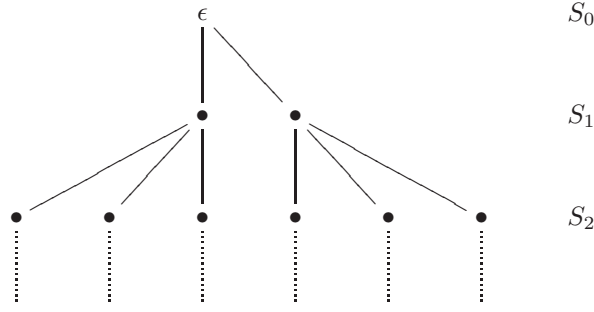
to obtain $\mathcal{D}(\Delta_{\varepsilon,\theta} + V(Q)) = \mathcal{D}(d_G(Q))$.

Finally, we examine the question of the equality of the domains in the context of simple trees. We provide with a positive result.

Proposition 4.2. *There is a simple tree T , such that $\mathcal{D}(\Delta_T) = \mathcal{D}(d_T(Q))$. Moreover the spectrum of the associated Laplacian is given by $\sigma(\Delta_T) = \sigma_{ac}(\Delta_T) = [0, \infty)$. In particular, $\alpha(T) = 0$.*

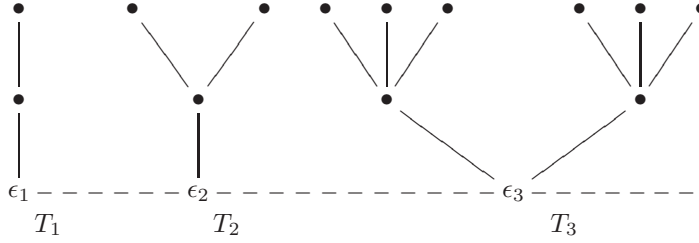
Here ac stands for absolutely continuous.

Proof. We start by fixing some notation. Given offspring sequence $(b_n)_{n \in \mathbb{N}}$, with $b_n \in \mathbb{N}^*$, we associate a simple tree with root ε such that $b_n = \#\{x, \overleftarrow{x} \in S_n\}$.



Example of a tree with $b_0 = 2$ and $b_1 = 3$.

We turn to our example and construct some trees $T_n = (\mathcal{E}_n, \mathcal{V}_n)$. For $n = 1$, we take $\mathcal{V}_1 := \mathbb{N}$, with $\varepsilon_1 := 0$ and with $\mathcal{E}_1(p, q) = 1$ if and only if $|p - q| = 1$. For $n \geq 2$, we take trees that are n -ari after the first generation. For each $n \in \mathbb{N} \setminus \{0, 1\}$, let ε_n be the root and set that the offspring $b(n)_k := n$, for all $k \in \mathbb{N} \setminus \{0\}$ and $b(n)_0 := n - 1$. Now take $T := (\mathcal{E}, \mathcal{V})$, where $\mathcal{V} := \cup_{n \in \mathbb{N} \setminus \{0, 1\}} \mathcal{V}_n$ and $\mathcal{E}(x, y) := \mathcal{E}_n(x, y)$, if $x, y \in \mathcal{V}_n$, $\mathcal{E}(\varepsilon_n, \varepsilon_{n+1}) := 1$, and $\mathcal{E}(x, y) := 0$ otherwise.



Graph of T

Note that $x \mapsto \sum_{y \sim x} |d_T(x) - d_T(y)|$ is with support on $\cup_n \{\varepsilon_n\}$ and is with values in $\{0, 1\}$. Hence, (4.3) is fulfilled. Since the graph is simple, Δ_T is essentially self-adjoint on $\mathcal{C}_c(\mathcal{V})$ (see Section 2.2). Therefore, we derive from Theorem 4.1 that the $\mathcal{D}(\Delta_T) = \mathcal{D}(d_T(Q))$.

We turn to the spectrum. First, $\sigma(\Delta_{G_1}) = \sigma_{ac}(\Delta_{G_1}) = [0, 4]$, where $G_1 := T_1$. This is easy to see by discrete Fourier transformation, e.g. [AF]. For each $n \geq 2$, T contains a subtree G_n which is n -ari and which is connected to the rest of T by only one edge. It is well-known that $\sigma_{ac}(\Delta_{G_n}) = [n - \sqrt{n}, n + \sqrt{n}]$, e.g. [AF]. Now, we denote by $\tilde{G}_i := T \setminus G_i$, the graph obtained by removing the only edge that is connecting G_i to the rest of the graph. Note that $\Delta_{\tilde{G}_i} - \Delta_{G_i}$ is a rank one operator, for all $i \geq 1$. Therefore, $\sigma_{ac}(\Delta_{G_i}) \subset \sigma_{ac}(\Delta_T)$ for all $i \geq 1$. Hence, $\sigma(\Delta_T) = \sigma_{ac}(\Delta_T) = [0, \infty)$. Finally note that if $\alpha(T) > 0$, (1.5) ensures that $\inf \sigma(\Delta_T) > 0$. This is a contradiction. \square

Remark 4.4. This construction also provide an example of graph, on which the adjacency matrix $\mathcal{A}_{\mathcal{E}, 0}$, see (3.4), is essentially self-adjoint on $\mathcal{C}_c(\mathcal{V})$ (see [GS, Lemma 2.1]) and has absolutely continuous spectrum equal to \mathbb{R} .

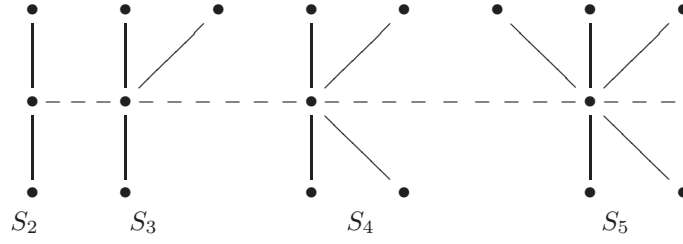
Finally, we give a negative example.

Proposition 4.3. There is a simple tree T , such that $\mathcal{D}(\Delta_T) \neq \mathcal{D}(d_T(Q))$ and such that the form-domains $\mathcal{D}(\Delta_T^{1/2}) = \mathcal{D}(d_T^{1/2}(Q))$.

Proof. The second point follows from Theorem 1.1. We start by constructing the star graph S_n . Let $S_n := (\mathcal{E}_n, \mathcal{V}_n)$ be defined as follows: $\mathcal{V}_n := \{1, \dots, n+1\}$ and so that $\mathcal{E}_n((1, j)) := 1, \forall j \in \{2, \dots, n+1\}$ and $\mathcal{E}_n((j, k)) := 0, \forall j, k \in \{2, \dots, n+1\}$. Consider now $f_n(x) := 1$ on \mathcal{V}_n . One has:

$$(4.8) \quad \|\Delta_{S_n} f_n\|^2 = 0, \quad \|f_n\|^2 = n+1, \quad \text{and} \quad \|d_{S_n}(Q) f_n\|^2 = n(n+1).$$

Now take $T := (\mathcal{E}, \mathcal{V})$, where $\mathcal{V} := \cup_{n \in \mathbb{N} \setminus \{0\}} \mathcal{V}_n$ and $\mathcal{E}(x, y) := \mathcal{E}_n(x, y)$, if $x, y \in \mathcal{V}_n$, $\mathcal{E}(x, y) := 1$, if $x = 1 \in \mathcal{V}_n$ and $y = 1 \in \mathcal{V}_{n+1}$, for all $n \geq 1$, and $\mathcal{E}(x, y) := 0$ otherwise.



Graph of T

Now if $\mathcal{D}(d_T(Q)) \subset \mathcal{D}(\Delta_T)$, by the uniform boundedness principle, there are constants $a, b > 0$, so that

$$\|d_T(Q)f\|^2 \leq a\|\Delta_T f\|^2 + b\|f\|^2, \quad \text{for all } f \in \mathcal{C}_c(T).$$

This leads to a contradiction with (4.8), as $\|\Delta_T f_n\|^2 = 4$ and $\|d_T(Q)f_n\|^2 = (n+2)^2 + n$ for $n \geq 3$. \square

4.2. The form-domain for bi-partite graphs. As (1.10) holds true for trees, it is natural to ask the question for bi-partite graphs. The answer is no. We start by relating the form-domain of the magnetic Laplacian with the inequality (4.9).

Proposition 4.4. *Let $G = (\mathcal{V}, \mathcal{E}, m, \theta)$ be a weighted bi-partite graph. Then there is $a \in (0, 1]$ and $C_a > 0$ so that:*

$$(4.9) \quad (1-a)\langle f, d_G(Q)f \rangle - C_a\|f\|^2 \leq \langle f, \Delta_{\mathcal{E}, \theta} f \rangle \leq (1+a)\langle f, d_G(Q)f \rangle + C_a\|f\|^2,$$

for all $f \in \mathcal{C}_c(\mathcal{V})$. Moreover, one can take some $a < 1$ in (4.9) if and only if $\mathcal{D}((\Delta_{\mathcal{E}, \theta})^{1/2}) = \mathcal{D}(d_G^{1/2}(Q))$.

Suppose also that $\Delta_{\mathcal{E}, \theta}$ has compact resolvent. Then, $d_G(Q)$ has also compact resolvent and, with the same a as in (4.9), one has:

$$(4.10) \quad 1-a \leq \liminf_{\lambda \rightarrow \infty} \frac{\mathcal{N}_\lambda(\Delta_{\mathcal{E}, \theta})}{\mathcal{N}_\lambda(d_G(Q))} \leq \limsup_{\lambda \rightarrow \infty} \frac{\mathcal{N}_\lambda(\Delta_{\mathcal{E}, \theta})}{\mathcal{N}_\lambda(d_G(Q))} \leq 1+a.$$

We recall that (1.4) ensures that $\mathcal{D}(d_G^{1/2}(Q)) \subset \mathcal{D}((\Delta_{\mathcal{E}, \theta})^{1/2})$.

Proof. First, note that (4.9) follows from (1.4) and Proposition 3.1. Suppose that (4.9) holds true for some $a < 1$, this gives immediately that $\mathcal{D}((\Delta_{\mathcal{E}, \theta})^{1/2}) = \mathcal{D}(d_G^{1/2}(Q))$. Reciprocally, suppose now that $\mathcal{D}((\Delta_{\mathcal{E}, \theta})^{1/2}) = \mathcal{D}(d_G^{1/2}(Q))$. By the uniform boundedness principle, there is $a_0, b_0 > 0$ so that

$$a_0\langle f, d_G(Q)f \rangle - b_0\|f\|^2 \leq \langle f, \Delta_{\mathcal{E}, \theta} f \rangle,$$

for all $f \in \mathcal{C}_c(\mathcal{V})$. Using again Proposition 3.1, (4.9) holds true with $a = 1 - a_0 < 1$.

We turn to the second part and work under the hypothesis that $\Delta_{\mathcal{E}, \theta}$ has compact resolvent. Corollary 2.1 ensures that $d_G(Q)$ has also compact resolvent. We conclude by using Proposition 2.3 twice. \square

Remark 4.5. *For a bi-partite graph, the constant 2 in (1.4) can be improved, in the sense of (4.9), if and only if $\mathcal{D}((\Delta_{\mathcal{E}, \theta})^{1/2}) = \mathcal{D}(d_G^{1/2}(Q))$.*

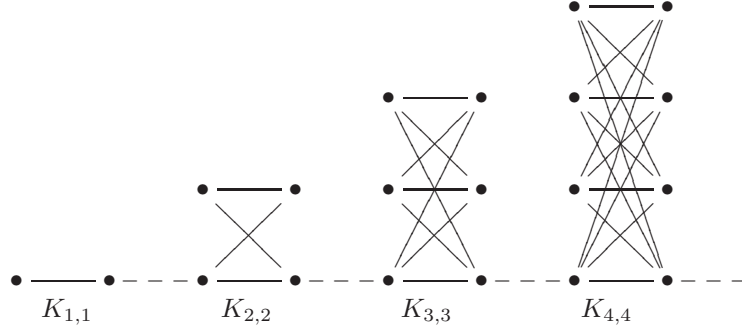
We finally provide an example.

Proposition 4.5. *There is a simple bi-partite graph K such that $\mathcal{D}(d_K^{1/2}) \subsetneq \mathcal{D}(\Delta_K^{1/2})$. In particular, the constant 2 in (1.4) is optimal in the sense of (4.9).*

Proof. We start by constructing a complete bi-partite graph. Let $K_{n,n} := (\mathcal{E}_n, \mathcal{V}_n)$ be defined as follows: $\mathcal{V}_n = \{1, \dots, n\} \times \{1, \dots, n\}$ and such that $\mathcal{E}_n((k, i), (j, i)) = 0, \forall j, k \in \{1, \dots, n\}$ and $i = 1, 2$ and $\mathcal{E}_n((k, 1), (j, 2)) = 0, \forall j, k \in \{1, \dots, n\}$. Consider now $f_n(x) = 1$ on $K_{n,n}$. One has:

$$(4.11) \quad \langle f_n, \Delta_K f_n \rangle = 0, \quad \|f_n\|^2 = 2n, \quad \text{and} \quad \langle f_n, d_K(Q)f_n \rangle = 2n^2.$$

Now take $K := (\mathcal{E}, \mathcal{V})$, where $\mathcal{V} = \cup_{n \in \mathbb{N} \setminus \{0\}} \mathcal{V}_n$ and $\mathcal{E}(x, y) := \mathcal{E}_n(x, y)$, if $x, y \in \mathcal{V}_n$, $\mathcal{E}(x, y) = 1$, if $x = (1, 2) \in \mathcal{V}_n$ and $y = (1, 1) \in \mathcal{V}_{n+1}$, for all $n \geq 1$, and $\mathcal{E}(x, y) = 0$ otherwise.



Graph of K

On the other hand, if $\mathcal{D}(d_K^{1/2}(Q)) = \mathcal{D}(\Delta_K^{1/2})$, the uniform boundedness principle ensures that there are constants $a, b > 0$, so that

$$\langle f, d_K(Q)f \rangle \leq a \langle f, \Delta_K f \rangle + b \|f\|^2, \quad \text{for all } f \in \mathcal{C}_c(K).$$

This leads to a contradiction with (4.11), as $\langle f_n, \Delta_K f_n \rangle = 2$, for $n \geq 2$. Optimality follows by Proposition and 4.4 and Remark 4.5. \square

5. PERTURBATION THEORY

We finally go into perturbation theory in order to obtain the stability of the essential spectrum, of the inequality (5.1), and of the asymptotic of eigenvalues.

Proposition 5.1. *Let $G = (\mathcal{V}, \mathcal{E}, m, \theta)$ and $G_o = (\mathcal{V}, \mathcal{E}_o, m, \theta_o)$ be weighted graphs and $V : \mathcal{V} \rightarrow \mathbb{R}$. Suppose that for all $\varepsilon > 0$ there is $c_\varepsilon > 0$ so that (1.10) holds true for $\Delta_{\mathcal{E}_o, \theta_o}$. Suppose that there is $\eta \in (0, 1)$ and $\kappa_\eta \geq 0$, so that*

$$(5.1) \quad |\langle f, V(Q)f \rangle| + 2 \langle f, \Lambda(Q)f \rangle \leq \eta \langle f, d_{G_o}(Q)f \rangle + \kappa_\eta \|f\|^2,$$

for all $f \in \mathcal{C}_c(\mathcal{V})$, where

$$(5.2) \quad \Lambda(x) := \frac{1}{m^2(x)} \sum_{y \sim x} |\mathcal{E}(x, y) - \tilde{\mathcal{E}}(x, y)|.$$

Then, one has that:

- The operator $\Delta_{\mathcal{E}, \theta} + V(Q)$ is bounded from below by some negative constant $-C$, in the form sense, and also that $\mathcal{D}((\Delta_{\mathcal{E}, \theta} + V(Q) + C)^{1/2}) = \mathcal{D}((\Delta_{\mathcal{E}_o, \theta_o})^{1/2})$.
- The essential spectrum of $\Delta_{\mathcal{E}, \theta} + V(Q)$ is empty if and only if the essential spectrum of $\Delta_{\mathcal{E}_o, \theta_o}$ is empty and if and only if $\lim_{|x| \rightarrow \infty} d_{G_o}(x) = +\infty$.
- Supposing that the essential spectrum of $\Delta_{\mathcal{E}, \theta} + V(Q)$ is empty and that for all $\eta \in (0, 1)$ there is $\kappa_\eta \geq 0$, so that (5.1) holds true, then:

$$\lim_{\lambda \rightarrow \infty} \frac{\mathcal{N}_\lambda(\Delta_{\mathcal{E}, \theta} + V(Q))}{\mathcal{N}_\lambda(d_{G_o}(Q))} = 1.$$

Proof. We start by noticing that, as in (1.4), (5.1) ensures that:

$$(5.3) \quad |\langle f, (\Delta_{\mathcal{E}, \theta} + V(Q) - \Delta_{\mathcal{E}_o, \theta_o})f \rangle| \leq |\langle f, V(Q)f \rangle| + 2 \langle f, \Lambda(G)f \rangle \leq \eta \langle f, d_{G_o}(Q)f \rangle + \kappa_\eta \|f\|^2$$

for all $f \in \mathcal{C}_c(\mathcal{V})$. Therefore, for all $\varepsilon > 0$ and all $\eta \in (0, 1)$, satisfying (5.1), there is $c_{\varepsilon, \eta}$ so that:

$$(1 - \eta - \varepsilon) \langle f, d_{G_o}(Q)f \rangle - c_{\varepsilon, \eta} \|f\|^2 \leq \langle f, (\Delta_{\mathcal{E}, \theta} + V(Q))f \rangle \leq (1 + \eta + \varepsilon) \langle f, d_{G_o}(Q)f \rangle + c_{\varepsilon, \eta} \|f\|^2,$$

for all $f \in \mathcal{C}_c(\mathcal{V})$. This gives, directly, the first point. Moreover, as above, the second and third points follow by Proposition 2.3. \square

We now turn to the stability of the essential spectrum.

Proposition 5.2. *Let $G = (\mathcal{V}, \mathcal{E}, m, \theta)$ and $G_\circ = (\mathcal{V}, \mathcal{E}_\circ, m, \theta_\circ)$ be weighted graphs and $V : \mathcal{V} \rightarrow \mathbb{R}$. Suppose that $\mathcal{D}(\Delta_{\mathcal{E}_\circ, \theta_\circ}^{1/2}) = \mathcal{D}(d_G^{1/2}(Q))$ and that*

$$(5.4) \quad |V(x)| + \Lambda(x) = o(1 + d_{G_\circ}(x)), \text{ as } |x| \rightarrow \infty,$$

where Λ is defined in (5.2). Then we obtain that $\sigma_{\text{ess}}(\Delta_{\mathcal{E}_\circ, \theta_\circ}) = \sigma_{\text{ess}}(\Delta_{\mathcal{E}, \theta} + V(Q))$.

Proof. The uniform boundedness Theorem and (5.4) implies that $\Delta_{\mathcal{E}, \theta} + V(Q)$ is bounded from below by some negative constant $-C$, in the form sense. We consider its Friedrichs extension. Next, KLMN's Theorem, e.g., [RS, Theorem X.17] ensures that $\mathcal{D}((\Delta_{\mathcal{E}, \theta} + V(Q) + C)^{1/2}) = \mathcal{D}(d_G^{1/2}(Q))$.

By Weyl's Theorem, e.g., [RS, Theorem XII.1], it is enough to show that the difference of the resolvents is compact. As we have to work with forms, one should be careful with the resolvent equation. We give a complete proof and refer to [GGo] for more discussions of this matter. To lighten notation, we set $H_0 := \Delta_{\mathcal{E}_\circ, \theta_\circ}$ and $H := \Delta_{\mathcal{E}, \theta} + V$. To start off, we give a rigorous meaning to

$$(5.5) \quad (H + i)^{-1} - (H_0 + i)^{-1} = (H + i)^{-1}(H_0 - H)(H_0 + i)^{-1}$$

Since $\mathcal{G} := \mathcal{D}((H + C)^{1/2}) = \mathcal{D}((H_0)^{1/2})$, both operators extend to an element of $\mathcal{B}(\mathcal{G}, \mathcal{G}^*)$. Here we use the Riesz lemma to identify \mathcal{H} with its anti-dual \mathcal{H}^* . We denote these extensions with a tilde.

We have $(H_0 + i)^{-1*} \mathcal{H} \subset \mathcal{G}$. This allows one to deduce that $(H_0 + i)^{-1}$ extends to a unique continuous operator $\mathcal{G}^* \rightarrow \mathcal{H}$. We denote it for the moment by R . From $R(H_0 + i)u = u$ for $u \in \mathcal{D}(H_0)$ we get, by density of $\mathcal{D}(H_0)$ in \mathcal{G} and continuity, $R(\widetilde{H}_0 + i)u = u$ for $u \in \mathcal{G}$. In particular

$$(H + i)^{-1} = R(\widetilde{H}_0 + i)(H + i)^{-1}.$$

Clearly,

$$(H_0 + i)^{-1} = (H_0 + i)^{-1}(H + i)(H + i)^{-1} = R(\widetilde{H} + i)(H + i)^{-1}.$$

We subtract the last two relations to obtain that

$$(H_0 + i)^{-1} - (H + i)^{-1} = R(\widetilde{H} - \widetilde{H}_0)(H + i)^{-1}$$

Since R is uniquely determined as the extension of $(H_0 + i)^{-1}$ to a continuous map $\mathcal{G}^* \rightarrow \mathcal{H}$, one may keep the notation $(H_0 + i)^{-1}$ for it. With this convention, the rigorous version of (5.5) that we use is:

$$(5.6) \quad (H_0 + i)^{-1} - (H + i)^{-1} = (H_0 + i)^{-1}(\widetilde{H} - \widetilde{H}_0)(H + i)^{-1}.$$

Therefore, to prove c), it is enough to show that $\widetilde{H} - \widetilde{H}_0$ is a compact operator from \mathcal{G} to \mathcal{G}^* . By (5.3), one get:

$$\begin{aligned} | \langle (1 + d_{G_\circ})^{-1/2}(Q)f, (\Delta_{\mathcal{E}, \theta} + V(Q) - \Delta_{\mathcal{E}_\circ, \theta_\circ})(1 + d_{G_\circ})^{-1/2}(Q)f \rangle | \leq \\ \langle (1 + d_{G_\circ})^{-1/2}(Q)f, (V + 2\Lambda)(Q)(1 + d_{G_\circ})^{-1/2}(Q)f \rangle, \end{aligned}$$

for all $f \in \mathcal{C}_c(\mathcal{V})$. Now, by hypothesis (5.4), $(V + 2\Lambda)(1 + d_{G_\circ})^{-1}(Q)$ is compact in $\ell^2(\mathcal{V}, m^2)$. We conclude by using Proposition 2.4. \square

Remark 5.1. *Note that (5.1) and (5.4) allow us to consider some potential $V(Q)$ that are unbounded from below, whereas $\Delta_{\mathcal{E}, \theta} + V(Q)$ is bounded from below. This is due to the fact that we know the form-domain explicitly.*

APPENDIX A. THE \mathcal{C}^1 -REGULARITY

We start with some generalities. Given a bounded operator B and a self-adjoint operator A acting in a Hilbert space \mathcal{H} , one says that $B \in \mathcal{C}^k(A)$ if $t \mapsto e^{-itA} B e^{itA}$ is strongly \mathcal{C}^k . Given a closed and densely defined operator B , one says that $B \in \mathcal{C}^k(A)$ if for some (hence any) $z \notin \sigma(B)$, $t \mapsto e^{-itA}(B - z)^{-1} e^{itA}$ is strongly \mathcal{C}^k . The two definitions coincide in the case of a bounded self-adjoint operator. We recall a result following from Lemma 6.2.9 and Theorem 6.2.10 of [ABG].

Theorem A.1. *Let A and B be two self-adjoint operators in the Hilbert space \mathcal{H} . For $z \notin \sigma(A)$, set $R(z) := (B - z)^{-1}$. The following points are equivalent to $B \in \mathcal{C}^1(A)$:*

(a) *For one (then for all) $z \notin \sigma(B)$, there is a finite c such that*

$$(A.1) \quad |\langle Af, R(z)f \rangle - \langle R(\bar{z})f, Af \rangle| \leq c \|f\|^2, \text{ for all } f \in \mathcal{D}(A).$$

(b) a. *There is a finite c such that for all $f \in \mathcal{D}(A) \cap \mathcal{D}(B)$:*

$$(A.2) \quad |\langle Af, Bf \rangle - \langle Bf, Af \rangle| \leq c(\|Bf\|^2 + \|f\|^2).$$

- b. For some (then for all) $z \notin \sigma(B)$, the set $\{f \in \mathcal{D}(A) \mid R(z)f \in \mathcal{D}(A) \text{ and } R(\bar{z})f \in \mathcal{D}(A)\}$ is a core for A .

Note that the condition (2.b) could be uneasy to check (see [GGé]). We mention [GM][Lemma A.2] and [GL][Lemma 3.2.2] to overcome this subtlety.

Note that (A.1) yields the commutator $[A, R(z)]$ extends to a bounded operator, in the form sense. We shall denote the extension by $[A, R(z)]_\circ$. In the same way, since $\mathcal{D}(B) \cap \mathcal{D}(A)$ is dense in $\mathcal{D}(B)$, (A.2) ensures that the commutator $[B, A]$ extends to a unique element of $\mathcal{B}(\mathcal{D}(B), \mathcal{D}(B)^*)$ denoted by $[B, A]_\circ$. Moreover, when $B \in \mathcal{C}^1(A)$, one has:

$$(A.3) \quad [A, (B - z)^{-1}]_\circ = \underbrace{(B - z)^{-1}}_{\mathcal{H} \leftarrow \mathcal{D}(B)^*} \underbrace{[B, A]_\circ}_{\mathcal{D}(B)^* \leftarrow \mathcal{D}(B)} \underbrace{(B - z)^{-1}}_{\mathcal{D}(B) \leftarrow \mathcal{H}}.$$

Here we use the Riesz lemma to identify \mathcal{H} with its anti-dual \mathcal{H}^* .

It turns out that an easier characterization is available if the domain of B is conserved under the action of the C_0 -group generated by A .

Theorem A.2. ([ABG, p. 258]) *Let A and B be two self-adjoint operators in the Hilbert space \mathcal{H} such that $e^{itA}\mathcal{D}(B) \subset \mathcal{D}(B)$, for all $t \in \mathbb{R}$. Then, $B \in \mathcal{C}^1(A)$ if and only if*

$$\liminf_{t \rightarrow 0^+} \|[B, e^{itB}/t]\|_{\mathcal{B}(\mathcal{D}(B), \mathcal{D}(B)^*)} < \infty.$$

Note that $e^{itA}\mathcal{D}(B)^* \subset \mathcal{D}(B)^*$ by duality.

APPENDIX B. THE HELFFER-SJÖSTRAND'S FORMULA

We present briefly the Helffer-Sjöstrand's formula. We refer to [GJ][Appendix B] and [BG][Appendix A] (see also [DG, HSS]) for commutator expansion. We first recall some well-known facts about almost analytic extensions. For $\rho \in \mathbb{R}$, let \mathcal{S}^ρ be the class of function $\varphi \in \mathcal{C}^\infty(\mathbb{R}; \mathbb{C})$ such that

$$(B.1) \quad \forall k \in \mathbb{N}, \quad C_k(\varphi) := \sup_{t \in \mathbb{R}} \langle t \rangle^{-\rho+k} |\varphi^{(k)}(t)| < \infty.$$

Equipped with the semi-norms defined by (B.1), \mathcal{S}^ρ is a Fréchet space. Leibniz' formula implies the continuous embedding: $\mathcal{S}^\rho \cdot \mathcal{S}^{\rho'} \subset \mathcal{S}^{\rho+\rho'}$. We shall use the following result, e.g., [DG, Appendix C.2].

Lemma B.1. *Let $\varphi \in \mathcal{S}^\rho$ with $\rho \in \mathbb{R}$. For all $l \in \mathbb{N}$, there is a smooth function $\varphi^{\mathbb{C}} : \mathbb{C} \rightarrow \mathbb{C}$, such that:*

$$(B.2) \quad \varphi^{\mathbb{C}}|_{\mathbb{R}} = \varphi, \quad \left| \frac{\partial \varphi^{\mathbb{C}}}{\partial \bar{z}}(z) \right| \leq c_1 \langle \Re(z) \rangle^{\rho-1-l} |\Im(z)|^l$$

$$(B.3) \quad \text{supp } \varphi^{\mathbb{C}} \subset \{x + iy \mid |y| \leq c_2 \langle x \rangle\},$$

$$(B.4) \quad \varphi^{\mathbb{C}}(x + iy) = 0, \text{ if } x \notin \text{supp } \varphi.$$

for some constants c_1, c_2 depending on the semi-norms (B.1) of φ in \mathcal{S}^ρ and not on φ .

One calls $\varphi^{\mathbb{C}}$ an *almost analytic extension* of φ . Let A be a self-adjoint operator, $\rho < 0$ and $\varphi \in \mathcal{S}^\rho$. By functional calculus, one has $\varphi(A)$ bounded. The Helffer-Sjöstrand's formula, e.g., [HS] and [DG], gives that for all almost analytic extension of $\varphi \in \mathcal{S}^\rho$, with $\rho < 0$, one has:

$$(B.5) \quad \varphi(A) = \frac{i}{2\pi} \int_{\mathbb{C}} \frac{\partial \varphi^{\mathbb{C}}}{\partial \bar{z}}(z - A)^{-1} dz \wedge d\bar{z}.$$

Note the integral exists in the norm topology, by (B.2) with $l = 1$ and by taking in account the domain of integration given in (B.3).

REFERENCES

- [AF] C. Allard and R. Froese: *A Mourre estimate for a Schrödinger operator on a binary tree*, Rev. Math. Phys. 12 (2000), no. 12, 1655–1667.
- [ABG] W. O. Amrein, A. Boutet de Monvel, and V. Georgescu: *C_0 -groups, commutator methods and spectral theory of N -body Hamiltonians*, Progress in Mathematics, vol. 135, Birkhäuser Verlag, Basel, 1996.
- [Aom] K. Aomoto: *Selfadjointness and limit pointness for adjacency operators on a tree*, J. Analyse Math. 53 (1989), 219–232.
- [BG] N. Boussaïd and S. Golénia: *Limiting absorption principle for some long range perturbations of Dirac systems at threshold energies*, Comm. Math. Phys. 299 (2010), no. 3, 677–708.
- [Chu] F.R.K. Chung: *Spectral graph theory* Regional Conference Series in Mathematics. 92. Providence, RI: American Mathematical Society (AMS). xi, 207 p.

- [CdV] Y. Colin de Verdière: *Spectres de graphes*, Cours Spécialisés, 4. Société Mathématique de France, Paris, 1998.
- [CTT] Y. Colin De Verdière, N. Toriki-Hamza and F. Truc: *Essential self-adjointness for combinatorial Schrödinger operators II- Metrically non complete graphs*, Mathematical Physics Analysis and Geometry 14, 1 (2011) 21–38.
- [CTT2] Y. Colin De Verdière, N. Toriki-Hamza and F. Truc: *Essential self-adjointness for combinatorial Schrödinger operators III- Magnetic fields*, preprint arXiv:1011.6492v1.
- [CDS] D. Cvetković, M. Doob, and H. Sachs: *Spectra of graphs. Theory and application*, Second edition. VEB Deutscher Verlag der Wissenschaften, Berlin, 1982. 368 pp.
- [DSV] G. Davidoff, P. Sarnak, and A. Valette: *Elementary number theory, group theory, and Ramanujan graphs*, London Mathematical Society Student Texts, 55. Cambridge University Press, Cambridge, 2003. x+144 pp.
- [DG] J. Dereziński and C. Gérard: *Scattering theory of classical and quantum n-particle systems.*, Texts and Monographs in Physics. Berlin: Springer. xii,444 p., 1997.
- [Do] J. Dodziuk: *isoperimetric inequality and transience of certain random walks*, Trans. Amer. Math. Soc. 284 (1984), no. 2, 787–794.
- [DK] J. Dodziuk and W.S. Kendall: *Combinatorial Laplacians and isoperimetric inequality*, from local times to global geometry, control and physics (Coventry, 1984/85), 68–74, Pitman Res. Notes Math. Ser., 150, Longman Sci. Tech., Harlow, 1986.
- [DM] J. Dodziuk and V. Matthai: *Kato’s inequality and asymptotic spectral properties for discrete magnetic Laplacians*, The ubiquitous heat kernel, 69–81, Contemp. Math., 398, Amer. Math. Soc., Providence, RI, 2006.
- [Fuj] K. Fujiwara: *Laplacians on rapidly branching trees*, Duke Math. Jour., 83 (1996), No. 1, 191–202.
- [GGé] V. Georgescu and C. Gérard: *On the virial theorem in quantum mechanics*, Comm. Math. Phys. 208 (1999) p 275–281.
- [GGo] V. Georgescu and S. Golénia: *Decay Preserving Operators and stability of the essential spectrum.*, J. Operator Theory 59 (2008), no. 1, 115–155.
- [GL] C. Gérard and I. Lapa: *Multiparticle quantum scattering in constant magnetic fields*, Mathematical Surveys and Monographs 90. Providence, RI: AMS, American Mathematical Society. xiii.
- [Gol] S. Golénia: *Unboundedness of adjacency matrices of locally finite graphs*, to appear in Letters in mathematical physics.
- [GS] S. Golénia and C. Schumacher: *The problem of deficiency indices for discrete Schrödinger operators on locally finite graphs*, to appear in Journal of Mathematical Physics.
- [GJ] S. Golénia and T. Jecko: *A new look at Mourre’s commutator theory.*, Complex Anal. Oper. Theory 1 (2007), no. 3, 399–422.
- [GM] S. Golénia and S. Moroianu: *Spectral analysis of magnetic Laplacians on conformally cusp manifolds*, Ann. Henri Poincaré 9 (2008), no. 1, 131–179.
- [HS] B. Helffer and J. Sjöstrand: *Opérateurs de schrödinger avec champs magnétiques faibles et constants*, Sémin. Équations Dériv. Partielles 1988–1989, Exp. No. 12, 11 p. (1989).
- [HSS] W. Hunziker, I.M. Sigal, and A. Soffer: *Minimal escape velocities*, Communications in Partial Differential Equations, 24:11, 2279 –2295 (1999).
- [Jor] P.E.T. Jørgensen: *Essential self-adjointness of the graph-Laplacian* J. Math. Phys. 49, No. 7, 073510, 33 p. (2008).
- [JP] P.E.T. Jørgensen and E.P.J. Pearse: *Spectral reciprocity and matrix representations of unbounded operators*, to appear in Journal of Functional Analysis.
- [Kel] M. Keller: *The essential spectrum of the Laplacian on rapidly branching tessellations*, Math. Ann. 346, Issue 1 (2010), 51–66.
- [KL] M. Keller and D. Lenz: *Unbounded Laplacians on graphs: Basic spectral properties and the heat equation*, Math. Model. Nat. Phenom. (2009) Vol. 5, No. 2.
- [KL2] M. Keller and D. Lenz: *Dirichlet forms and stochastic completeness of graphs and subgraphs*, to appear in Journal fuer die reine und angewandte Mathematik.
- [KLW] M. Keller, D. Lenz and R. Wojciechowski: *Volume Growth, Spectrum and Stochastic Completeness of Infinite Graphs*, preprint arXiv:1105.0395v1.
- [Ma] J. Masamune: *A Liouville property and its application to the Laplacian of an infinite graph*, Spectral analysis in geometry and number theory, 103–115, Contemp. Math., 484, Amer. Math. Soc.
- [Mil] O. Milatovic: *Essential self-adjointness of magnetic Schrödinger operators on locally finite graphs*, to appear in Integral Equations and Operator Theory.
- [Mil2] O. Milatovic: *Essential self-adjointness of discrete magnetic Schrödinger operators*, preprint arXiv:1105.3129v1.
- [MW] B. Mohar and W. Woess: *A survey on spectra of infinite graphs*, J. Bull. Lond. Math. Soc. 21, No.3, 209–234 (1989).
- [RS] M. Reed and B. Simon: *Methods of Modern Mathematical Physics, Tome I–IV: Analysis of operators* Academic Press.
- [Tor] N. Toriki-Hamza: *Laplaciens de graphes infinis I Graphes métriquement complets*, Confluentes Math. 2 (3) (2010) 333–350.
- [Web] A. Weber: *Analysis of the physical Laplacian and the heat flow on a locally finite graph*, J. Math. Anal. Appl. 370 (1) (2010) 146–158.
- [Woj] R. Wojciechowski: *Stochastic completeness of graphs*, Ph.D. Thesis, 2007, arXiv:0712.1570v2[math.SP].

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