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► **To cite this version:**

Muriel Boulakia, Anne-Claire Egloffé, Céline Grandmont. Stability estimates for a Robin coefficient in the two-dimensional Stokes system. *Mathematical Control and Related Fields*, 2013, 3 (1), <https://www.aims sciences.org/journals/pdfs.jsp?paperID=8296&mode=full>. hal-00582559v3

HAL Id: hal-00582559

<https://hal.science/hal-00582559v3>

Submitted on 4 Dec 2012

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Stability estimates for a Robin coefficient in the two-dimensional Stokes system ^{*}

Muriel Boulakia [†] Anne-Claire Egloffé [‡] Céline Grandmont [‡]

December 4, 2012

Abstract

In this paper, we consider the Stokes equations and we are concerned with the inverse problem of identifying a Robin coefficient on some non accessible part of the boundary from available data on the other part of the boundary. We first study the identifiability of the Robin coefficient and then we establish a stability estimate of logarithm type thanks to a Carleman inequality due to A. L. Bukhgeim [12] and under the assumption that the velocity of a given reference solution stays far from 0 on a part of the boundary where Robin conditions are prescribed.

Keywords: Inverse boundary coefficient problem, Stokes system, Robin boundary condition, Identifiability, Carleman inequality, Logarithmic stability estimate.

1 Introduction

Let us consider an open Lipschitz bounded connected domain Ω of \mathbb{R}^d , $d \geq 2$. We assume that the boundary $\partial\Omega$ is composed of two open non-empty parts Γ_0 and Γ_e such that $\Gamma_e \cup \Gamma_0 = \partial\Omega$ and $\bar{\Gamma}_e \cap \bar{\Gamma}_0 = \emptyset$ (Figure 1 gives an example of such a geometry in dimension 2).

We denote by n the exterior unit normal to Ω and let $\tau = (\tau_1, \dots, \tau_{d-1})$ be $d-1$ vectors of \mathbb{R}^d such that (n, τ) is an orthogonal basis of \mathbb{R}^d .

We introduce the following boundary problem:

$$\left\{ \begin{array}{ll} \partial_t u(t, x) - \Delta u(t, x) + \nabla p(t, x) & = 0, \quad \forall x \in \Omega, \forall t > 0, \\ \operatorname{div} u(t, x) & = 0, \quad \forall x \in \Omega, \forall t > 0, \\ \frac{\partial u}{\partial n}(t, x) - p(t, x)n(x) & = g(t, x), \quad \forall x \in \Gamma_e, \forall t > 0, \\ \frac{\partial u}{\partial n}(t, x) - p(t, x)n(x) + q(x)u(t, x) & = 0, \quad \forall x \in \Gamma_0, \forall t > 0, \\ u(0, x) & = u_0(x), \quad \forall x \in \Omega. \end{array} \right. \quad (1.1)$$

Notice that we assume that the Robin coefficient q defined on Γ_0 only depends on the space variable. Our objective is to determine the coefficient q from the values of u and p on Γ_e .

Such kinds of systems naturally appear in the modeling of biological problems like, for example, blood flow in the cardiovascular system (see [22] and [25]) or airflow in the lungs (see [4]). For an introduction on the modeling of the airflow in the lungs and on different boundary conditions which may be prescribed, we refer to [17]. The part of the boundary Γ_e represents a physical

^{*}This work was partially funded by the ANR-08-JCJC-013-01 (M3RS) project headed by C. Grandmont and the ANR-BLAN-0213-02 (CISIFS) project headed by L. Rosier.

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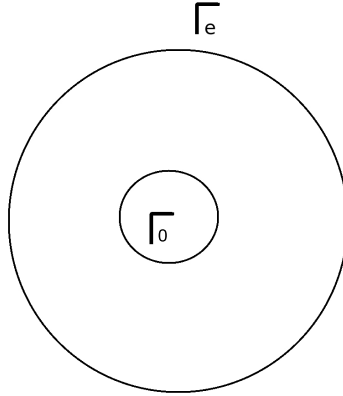


Figure 1: Example of an open set Ω such that $\Gamma_e \cup \Gamma_0 = \partial\Omega$ and $\bar{\Gamma}_e \cap \bar{\Gamma}_0 = \emptyset$ in dimension 2.

boundary on which measurements are available and Γ_0 represents an artificial boundary on which Robin boundary conditions or mixed boundary conditions involving the fluid stress tensor and its flux at the outlet are prescribed.

Similar inverse problems have been widely studied for the Laplace equation [2], [5], [13], [14], [15] and [24]. This kind of problems arises in general in corrosion detection which consists in determining a Robin coefficient on the inaccessible portion of the boundary thanks to electrostatic measurements performed on the accessible boundary. Most of these papers prove a logarithmic stability estimate ([2], [5], [13] and [15]). We mention that, in [14], S. Chaabane and M. Jaoua obtained both local and monotone global Lipschitz stability for regular Robin coefficient and under the assumption that the flux g is non negative. Under the *a priori* assumption that the Robin coefficient is piecewise constant, E. Sincich has obtained in [24] a Lipschitz stability estimate. To prove stability estimates, different approaches are developed in these papers. A first one consists in using the complex analytic function theory (see [2], [13]). A characteristic of this method is that it is only valid in dimension 2. Another classical approach is based on Carleman estimates (see [5] and [15]). In [5], the authors use a result proved by K.D. Phung in [21] to obtain a logarithmic stability estimate which is valid in any dimension for an open set Ω of class \mathcal{C}^∞ . This result has been generalized in [8] and [9] to $\mathcal{C}^{1,1}$ and Lipschitz domains. Moreover, in [5], the authors use semigroup theory to obtain a stability estimate in long time for the heat equation from the stability estimate for the Laplace equation.

In this article, we prove an identifiability result and a logarithmic stability estimate for the Stokes equations with Robin boundary conditions 1.1 under the assumption that the velocity of a given reference solution stays far from 0 on a part of the boundary where Robin conditions are prescribed. We would like to highlight why this assumption appears for the inverse problem of recovering a Robin coefficient. Let us consider (u_i, p_i) be solutions of system 1.1 associated to $q = q_i$, for $i = 1, 2$. Using the boundary conditions on Γ_0 , we obtain

$$(q_2 - q_1)u_1 = q_2(u_1 - u_2) + \left(\frac{\partial u_1}{\partial n} - \frac{\partial u_2}{\partial n} \right) - (p_1 - p_2)n.$$

When u_1 vanishes, difficulties occur to estimate the difference between the Robin coefficients $q_2 - q_1$. In the case of the scalar Laplace equation, it is possible to determine the sign of u solution of

$$\begin{cases} \Delta u & = 0, & \text{in } \Omega, \\ \frac{\partial u}{\partial n} & = g, & \text{on } \Gamma_e, \\ \frac{\partial u}{\partial n} + qu & = 0, & \text{on } \Gamma_0, \end{cases}$$

on any compact subset $K \subset \Gamma_0$ under some positivity assumption on the flux g (see [14]). Such

a result comes from properties specific to harmonic functions, like for instance the maximum principle. When the flux g has a variable sign, G. Alessandrini, L. Del Piero and L. Rondi provide in [2] a quantitative control of the vanishing rate of u that allows to estimate the difference between the Robin coefficients on $\{x \in \Gamma_0/d(x, \partial\Omega \setminus \Gamma_0) > d\}$, for any $d > 0$, by using methods of complex analytic function theory. Moreover, G. Alessandrini and E. Sincich proved in [3] that the oscillation of u on Γ_0 is bounded from below by a constant depending on the *a priori* data only. To do so, they use unique continuation estimates for the Laplace equation. Due to the methods employed, it does not seem that we can extend these results to the Stokes system. This is why we estimate the Robin coefficient on a compact subset $K \subset \Gamma_0$ on which u_1 does not vanish. This estimate and the set K depend on u_1 , and knowing whether one can control our solution, for well chosen data, on the whole set Γ_0 or on any compact subset $K \subset \Gamma_0$ remains an open problem. Note however that in a really particular case (detailed in Remark 4.9), one can obtain a logarithmic estimate on the whole set Γ_0 .

The paper is organized as follows. The second section contains preliminary results on the regularity of the solution. In the third section, we are interested in the identifiability of the Robin coefficient q . Under some regularity assumptions and using the theorem of unique continuation for the Stokes equations proved in [18], we prove that if two measurements of the velocity are equal on $(0, T) \times \Gamma$, where $\Gamma \subseteq \Gamma_e$ is a non-empty open subset of the boundary, then the two corresponding Robin coefficients are also equal on Γ_0 . Section 4 corresponds to the main part of our article. The results of this section are only valid in dimension 2. We prove a stability estimate, first for the stationary problem and then for the evolution problem. To do this, we use a global Carleman inequality due to A. L. Bukhgeim which is only valid in dimension 2 (see [12]). The stability estimate for the unsteady problem is deduced from the stability estimate for the stationary problem thanks to the semigroup theory. We end Section 4 by concluding remarks and perspectives to this work.

When we are not more specific, C is a generic constant, whose value may change and which only depends on the geometry of the open set Ω and of the boundaries Γ_e and Γ_0 . Moreover, we denote indifferently by $|| \cdot ||$ a norm on \mathbb{R}^n , for any $n \geq 1$.

We are going to start with some preliminary results which will be useful in the subsequent sections.

2 Preliminary results

In this section we study the well-posedness of the system and the regularity of the solution.

2.1 Regularity of the stationary problem

Let us first consider the stationary case:

$$\left\{ \begin{array}{ll} -\Delta u + \nabla p & = f, \quad \text{in } \Omega, \\ \operatorname{div} u & = 0, \quad \text{in } \Omega, \\ \frac{\partial u}{\partial n} - pn & = g, \quad \text{on } \Gamma_e, \\ \frac{\partial u}{\partial n} - pn + qu & = 0, \quad \text{on } \Gamma_0. \end{array} \right. \quad (2.1)$$

For $g \in H^{-\frac{1}{2}}(\Gamma_e)^d$ and $v \in H^{\frac{1}{2}}(\Gamma_e)^d$, we denote by $\langle g, v \rangle_{-\frac{1}{2}, \frac{1}{2}, \Gamma_e}$ the image of v by the linear form g .

Let us introduce some functional spaces:

$$V = \{v \in H^1(\Omega)^d / \operatorname{div} v = 0 \text{ in } \Omega\}$$

and

$$H = \overline{V}^{L^2(\Omega)^d}.$$

Proposition 2.1. *Let $\alpha > 0$, $f \in L^2(\Omega)^d$, $g \in H^{-\frac{1}{2}}(\Gamma_e)^d$ and $q \in L^\infty(\Gamma_0)$ be such that $q \geq \alpha$ on Γ_0 . System 2.1 admits a unique solution $(u, p) \in V \times L^2(\Omega)$. Moreover, there exists a constant $C(\alpha) > 0$ such that*

$$\|u\|_{H^1(\Omega)^d} \leq C(\alpha)(\|g\|_{H^{-\frac{1}{2}}(\Gamma_e)^d} + \|f\|_{L^2(\Omega)^d}). \quad (2.2)$$

Proof of Proposition 2.1. The variational formulation of the problem is: find $u \in V$ such that for every $v \in V$,

$$\int_{\Omega} \nabla u : \nabla v + \int_{\Gamma_0} qu \cdot v = \langle g, v|_{\Gamma_e} \rangle_{-\frac{1}{2}, \frac{1}{2}, \Gamma_e} + \int_{\Omega} f \cdot v.$$

For all $(u, v) \in V \times V$, we denote by

$$a_q(u, v) = \int_{\Omega} \nabla u : \nabla v + \int_{\Gamma_0} qu \cdot v, \quad (2.3)$$

and for all $v \in V$,

$$L_1(v) = \langle g, v|_{\Gamma_e} \rangle_{-\frac{1}{2}, \frac{1}{2}, \Gamma_e} + \int_{\Omega} f \cdot v.$$

We easily verify that a_q is a continuous symmetric bilinear form. Since $q \geq \alpha > 0$, according to the generalized Poincaré inequality, the bilinear form a_q is coercive on V . On the other hand, L_1 is a continuous linear form on V . Thus we prove the existence and uniqueness of $u \in V$ solution of 2.1 by using Lax-Milgram Theorem. We obtain simultaneously estimate 2.2. We prove the existence and uniqueness of $p \in L_0^2(\Omega)$ in a classical way, using De Rham Theorem. The fact that p is unique in $L^2(\Omega)$ comes from the boundary conditions. We refer to [10] for a complete proof in the case of Neumann boundary condition. \square

Next we want to derive regularity properties of the solution. Let us first recall existence and regularity results for the Stokes problem with Neumann boundary condition proved in [10].

Proposition 2.2. *Let $k \in \mathbb{N}$. Assume that Ω is of class $\mathcal{C}^{k+1,1}$. We assume that:*

$$(f, h) \in H^k(\Omega)^d \times H^{k+\frac{1}{2}}(\partial\Omega)^d.$$

Then the solution (u, p) of

$$\begin{cases} -\Delta u + \nabla p &= f, & \text{in } \Omega, \\ \operatorname{div} u &= 0, & \text{in } \Omega, \\ \frac{\partial u}{\partial n} - pn &= h, & \text{on } \partial\Omega, \end{cases}$$

belongs to $H^{k+2}(\Omega)^d \times H^{k+1}(\Omega)$ and there exists a constant $C > 0$ such that:

$$\|u\|_{H^{k+2}(\Omega)^d} + \|p\|_{H^{k+1}(\Omega)} \leq C(\|h\|_{H^{k+\frac{1}{2}}(\partial\Omega)^d} + \|f\|_{H^k(\Omega)^d}).$$

In order to study the Stokes system with Robin boundary conditions, one needs to specify to which space the Robin coefficient q belongs. As stated in Proposition 2.4, we will assume that q belongs to some Sobolev space $H^s(\Gamma_0)$ where s is large enough so that $qu|_{\Gamma_0}$ belongs to $H^r(\Gamma_0)$ if $u|_{\Gamma_0}$ belongs to $H^r(\Gamma_0)$. This stability in the Sobolev spaces will allow to apply the previous proposition (Proposition 2.2). Before stating the regularity result, let us state the following lemma:

Lemma 2.3. *Let $r, s \in \mathbb{R}$, with $s > \frac{d-1}{2}$ and $0 \leq r \leq s$. Let $q \in H^s(\Gamma_0)$. The linear operator*

$$\begin{aligned} T : H^r(\Gamma_0) &\rightarrow H^r(\Gamma_0) \\ u &\mapsto qu \end{aligned}$$

is continuous. Furthermore, the following estimate holds true

$$\|qu\|_{H^r(\Gamma_0)} \leq C\|q\|_{H^s(\Gamma_0)}\|u\|_{H^r(\Gamma_0)}.$$

Proof of Lemma 2.3. Since $s > \frac{d-1}{2}$, $H^s(\Gamma_0)$ is a Banach algebra (see [1]) and thus $T \in \mathcal{L}(H^s(\Gamma_0), H^s(\Gamma_0))$ and $\|T\|_s = \sup_{u \in H^s(\Gamma_0), u \neq 0} \frac{\|Tu\|_{H^s(\Gamma_0)}}{\|u\|_{H^s(\Gamma_0)}} \leq \|q\|_{H^s(\Gamma_0)}$. Moreover, since $H^s(\Gamma_0) \hookrightarrow L^\infty(\Gamma_0)$, $T \in \mathcal{L}(L^2(\Gamma_0), L^2(\Gamma_0))$ and

$$\|T\|_0 = \sup_{u \in L^2(\Gamma_0), u \neq 0} \frac{\|Tu\|_{L^2(\Gamma_0)}}{\|u\|_{L^2(\Gamma_0)}} \leq \|q\|_{L^\infty(\Gamma_0)} \leq C\|q\|_{H^s(\Gamma_0)}.$$

Thus, the result follows by interpolation (see [7] or [19]). \square

From Proposition 2.2 and Lemma 2.3 we deduce the following result:

Proposition 2.4. *Let $k \in \mathbb{N}$ and $s \in \mathbb{R}$ with $s > \frac{d-1}{2}$ and $s \geq k + \frac{1}{2}$. Assume that Ω is of class $\mathcal{C}^{k+1,1}$. Let $\alpha > 0$, $M > 0$, $f \in H^k(\Omega)^d$, $g \in H^{k+\frac{1}{2}}(\Gamma_e)^d$ and $q \in H^s(\Gamma_0)$ such that $q \geq \alpha$ on Γ_0 . Then the solution (u, p) of system 2.1 belongs to $H^{k+2}(\Omega)^d \times H^{k+1}(\Omega)$. Moreover, there exists a constant $C(\alpha, M) > 0$ such that for every $q \in H^s(\Gamma_0)$ satisfying $\|q\|_{H^s(\Gamma_0)} \leq M$,*

$$\|u\|_{H^{k+2}(\Omega)^d} + \|p\|_{H^{k+1}(\Omega)} \leq C(\alpha, M)(\|g\|_{H^{k+\frac{1}{2}}(\Gamma_e)^d} + \|f\|_{H^k(\Omega)^d}).$$

Proof of Proposition 2.4. Let us prove the result for $k = 0$. Let $h = -qu|_{\Gamma_0} + g$. According to Proposition 2.1, u belongs to $H^1(\Omega)^d$. We obtain from Lemma 2.3 for $r = 1/2$ that $qu|_{\Gamma_0} \in H^{\frac{1}{2}}(\Gamma_0)^d$, which implies, since $g \in H^{\frac{1}{2}}(\Gamma_e)^d$ and $\bar{\Gamma}_e \cap \bar{\Gamma}_0 = \emptyset$, that $h \in H^{\frac{1}{2}}(\partial\Omega)^d$. Using Proposition 2.2 with $k = 0$ we obtain that $(u, p) \in H^2(\Omega)^d \times H^1(\Omega)$ and:

$$\|u\|_{H^2(\Omega)^d} + \|p\|_{H^1(\Omega)} \leq C(\|h\|_{H^{\frac{1}{2}}(\partial\Omega)^d} + \|f\|_{L^2(\Omega)^d}).$$

But, since by assumption, $\|q\|_{H^s(\Gamma_0)} \leq M$, we have from Lemma 2.3 with $r = 1/2$, that:

$$\|h\|_{H^{\frac{1}{2}}(\partial\Omega)^d} \leq C(M)(\|u\|_{H^{\frac{1}{2}}(\partial\Omega)^d} + \|g\|_{H^{\frac{1}{2}}(\Gamma_e)^d}).$$

We obtain:

$$\|u\|_{H^2(\Omega)^d} + \|p\|_{H^1(\Omega)} \leq C(M)(\|g\|_{H^{\frac{1}{2}}(\Gamma_e)^d} + \|u\|_{H^1(\Omega)^d} + \|f\|_{L^2(\Omega)^d}).$$

Thus we obtain the result for $k = 0$ using the inequality of Proposition 2.1. We then proceed by induction to prove the result for any $k \in \mathbb{N}$. \square

Remark 2.5. *Note that the space to which the Robin coefficient q belongs is not optimal. One could surely obtain similar regularity result for a less regular Robin coefficient. In fact, the key argument to proceed by induction in the proof of Proposition 2.4 is that $qu|_{\Gamma_0} \in H^{k+\frac{1}{2}}(\Gamma_0)^d$, for $u \in H^{k+\frac{1}{2}}(\Gamma_0)^d$ (this property allows to apply the regularity result given by Proposition 2.2).*

2.2 Regularity of the evolution problem.

Concerning the initial problem 1.1, we can prove, using the Galerkin method, the following regularity results. For the sake of completeness, the proof of Theorem 2.6 is given in the appendix.

Theorem 2.6. *Let $s \in \mathbb{R}$ be such that $s > \frac{d-1}{2}$, $T > 0$, $\alpha > 0$ and $u_0 \in V$. We assume that Ω is of class $\mathcal{C}^{1,1}$, $g \in H^1(0, T; H^{\frac{1}{2}}(\Gamma_e)^d)$ and $q \in H^s(\Gamma_0)$ is such that $q \geq \alpha$ on Γ_0 . Then problem 1.1 admits a unique solution $(u, p) \in L^2(0, T; H^2(\Omega)^d) \cap H^1(0, T; L^2(\Omega)^d) \cap L^\infty(0, T; V) \times L^2(0, T; H^1(\Omega))$.*

The following corollary will be useful when we will prove stability estimates for the evolution problem 1.1.

Corollary 2.7. *Let $s \in \mathbb{R}$ be such that $s > \frac{d-1}{2}$ and $s \geq \frac{3}{2}$, $T > 0$, $\alpha > 0$ and $u_0 \in H^3(\Omega)^d \cap H$. We assume that Ω is of class $\mathcal{C}^{2,1}$, $g \in H^2(0, T; H^{\frac{3}{2}}(\Gamma_e)^d)$ and that $q \in H^s(\Gamma_0)$ is such that $q \geq \alpha$ on Γ_0 . Then, problem 1.1 admits a unique solution $(u, p) \in L^\infty(0, T; H^3(\Omega)^d) \cap H^1(0, T; H^2(\Omega)^d) \cap H^2(0, T; L^2(\Omega)^d) \times L^\infty(0, T; H^2(\Omega)) \cap H^1(0, T; H^1(\Omega))$.*

Proof of Corollary 2.7. Let (u, p) be the solution of 1.1. Let us consider the following system:

$$\begin{cases} \partial_t v - \Delta v + \nabla \zeta &= 0, & \text{in } (0, T) \times \Omega, \\ \operatorname{div} v &= 0, & \text{in } (0, T) \times \Omega, \\ \frac{\partial v}{\partial n} - \zeta n &= \partial_t g, & \text{on } (0, T) \times \Gamma_e, \\ \frac{\partial v}{\partial n} - \zeta n + qv &= 0, & \text{on } (0, T) \times \Gamma_0, \\ v(0) &= \Delta u_0 - \nabla p_0, & \text{in } \Omega, \end{cases} \quad (2.4)$$

where $p_0 \in H^2(\Omega)$ is defined as the solution of the following elliptic boundary problem:

$$\begin{cases} \Delta p_0 &= 0, & \text{in } \Omega, \\ p_0 &= \frac{\partial u_0}{\partial n} \cdot n - g|_{t=0} \cdot n, & \text{on } \Gamma_e, \\ p_0 &= \frac{\partial u_0}{\partial n} \cdot n + qu_0 \cdot n, & \text{on } \Gamma_0. \end{cases}$$

According to Theorem 2.6, we obtain that (v, ζ) belongs to $L^2(0, T; H^2(\Omega)^d) \cap H^1(0, T; L^2(\Omega)^d) \cap L^\infty(0, T; V) \times L^2(0, T; H^1(\Omega))$. Remark that $(\partial_t u, \partial_t p)$ is solution of system 2.4 in the distribution sense on $(0, T)$. Thus, by uniqueness, $(v, \zeta) = (\partial_t u, \partial_t p)$. Then, since $q \in H^s(\Gamma_0)$ and $(\partial_t u, g) \in L^\infty(0, T; V) \times L^\infty(0, T; H^{\frac{3}{2}}(\Gamma_e)^d)$ we deduce from Proposition 2.4 that $(u, p) \in L^\infty(0, T; H^3(\Omega)^d) \times L^\infty(0, T; H^2(\Omega))$. \square

3 Identifiability

3.1 Unique continuation

We start by recalling a unique continuation result for the Stokes equations proved in [18].

Theorem 3.1. *We denote by $Q = (0, T) \times \Omega$ and let O be an open subset in Q . The horizontal component of O is*

$$C(O) = \{(t, x) \in Q / \exists x_0 \in \Omega, (t, x_0) \in O\}.$$

Let $(u, p) \in L^2(0, T; H_{loc}^1(\Omega))^d \times L_{loc}^2(Q)$ be a weak solution of

$$\begin{cases} \partial_t u - \Delta u + \nabla p &= 0, & \text{in } (0, T) \times \Omega, \\ \operatorname{div} u &= 0, & \text{in } (0, T) \times \Omega, \end{cases}$$

satisfying $u = 0$ in O then $u = 0$ and p is constant in $C(O)$.

From this theorem, we easily deduce the following result which will be useful in the next subsection.

Corollary 3.2. *Let $\delta > 0$, $x_0 \in \partial\Omega$, $t_0 \in (0, T)$ and $r > 0$ be such that $\gamma = (t_0 - \delta, t_0 + \delta) \times (\mathcal{B}(x_0, r) \cap \partial\Omega)$ is an open set in $(0, T) \times \partial\Omega$. Let $(u, p) \in L^2(0, T; H^2(\Omega)^d) \times L^2(0, T; H^1(\Omega))$ be solution of:*

$$\begin{cases} \partial_t u - \Delta u + \nabla p &= 0, & \text{in } (0, T) \times \Omega, \\ \operatorname{div} u &= 0, & \text{in } (0, T) \times \Omega, \end{cases}$$

satisfying $u = 0$ and $\frac{\partial u}{\partial n} - pn = 0$ on γ . Then $u = 0$ and $p = 0$ in $(t_0 - \delta, t_0 + \delta) \times \Omega$.

Proof of Corollary 3.2. We extend u and p by 0 on $(t_0 - \delta, t_0 + \delta) \times (\mathcal{B}(x_0, r) \cap \Omega^c)$:

$$\tilde{u} \text{ (resp } \tilde{p}) = \begin{cases} u \text{ (resp } p), & \text{in } (t_0 - \delta, t_0 + \delta) \times \Omega, \\ 0, & \text{in } (t_0 - \delta, t_0 + \delta) \times (\mathcal{B}(x_0, r) \cap \Omega^c), \end{cases}$$

and we denote $\tilde{\Omega} = \Omega \cup \mathcal{B}(x_0, r)$. Let us verify that $(\tilde{u}, \tilde{p}) \in L^2(0, T; H^1(\Omega)^d) \times L^2(0, T; L^2(\Omega))$ is still a solution of the Stokes equations in $\tilde{\Omega}$. Let $v \in \mathcal{D}(\tilde{\Omega})^d$. We check by integration by parts in space that almost everywhere in $t \in (t_0 - \delta, t_0 + \delta)$:

$$\int_{\tilde{\Omega}} \partial_t \tilde{u} \cdot v + \int_{\tilde{\Omega}} \nabla \tilde{u} : \nabla v - \int_{\tilde{\Omega}} \tilde{p} \operatorname{div} v = 0.$$

Moreover $\operatorname{div} \tilde{u} = 0$ in $(t_0 - \delta, t_0 + \delta) \times \tilde{\Omega}$. Therefore, we can apply Theorem 3.1 to (\tilde{u}, \tilde{p}) : $(\tilde{u}, \tilde{p}) = (0, 0)$ in $(t_0 - \delta, t_0 + \delta) \times \tilde{\Omega}$ which implies that $u = 0$ and p is constant in $(t_0 - \delta, t_0 + \delta) \times \Omega$. At last, the fact that $\frac{\partial u}{\partial n} - pn = 0$ on γ implies that $p = 0$ in $(t_0 - \delta, t_0 + \delta) \times \Omega$. \square

3.2 Application

Proposition 3.3. *Let $s > \frac{d-1}{2}$, $T > 0$, $\alpha > 0$, $x_e \in \Gamma_e$, $r > 0$, $g \in H^1(0, T; H^{\frac{1}{2}}(\Gamma_e)^d)$ be non identically zero, $u_0 \in V$ and $q_j \in H^s(\Gamma_0)$ be such that $q_j \geq \alpha$ on Γ_0 for $j = 1, 2$. Let (u_j, p_j) be the weak solutions of 1.1 with $q = q_j$ for $j = 1, 2$. We assume that $u_1 = u_2$ on $(0, T) \times (\mathcal{B}(x_e, r) \cap \Gamma_e)$. Then $q_1 = q_2$ on Γ_0 .*

Proof of Proposition 3.3. We are going to prove Proposition 3.3 by contradiction: we assume that q_1 is not identically equal to q_2 on Γ_0 .

Thanks to Theorem 2.6, we have $(u_j, p_j) \in L^2(0, T; H^2(\Omega)^d) \times L^2(0, T; H^1(\Omega))$ for $j = 1, 2$. We define by $u = u_1 - u_2$ and $p = p_1 - p_2$. Let us notice that (u, p) is the solution of the following problem:

$$\left\{ \begin{array}{ll} \partial_t u - \Delta u + \nabla p & = 0, \quad \text{in } (0, T) \times \Omega, \\ \operatorname{div} u & = 0, \quad \text{in } (0, T) \times \Omega, \\ \frac{\partial u}{\partial n} - pn & = 0, \quad \text{on } (0, T) \times \Gamma_e, \\ \frac{\partial u}{\partial n} - pn + q_1 u_1 - q_2 u_2 & = 0, \quad \text{on } (0, T) \times \Gamma_0. \end{array} \right.$$

By assumption, $u = 0$ and $\frac{\partial u}{\partial n} - pn = 0$ on $(0, T) \times (\mathcal{B}(x_e, r) \cap \Gamma_e)$. Thus, according to Corollary 3.2, $u_1 = u_2$ and $p_1 = p_2$ in $(0, T) \times \Omega$. Consequently, we deduce from

$$\begin{aligned} \frac{\partial u_1}{\partial n} - p_1 n + q_1 u_1 &= 0, \quad \text{on } (0, T) \times \Gamma_0, \\ \frac{\partial u_1}{\partial n} - p_1 n + q_2 u_1 &= 0, \quad \text{on } (0, T) \times \Gamma_0, \end{aligned}$$

that

$$u_1(q_1 - q_2) = 0 \quad \text{on } (0, T) \times \Gamma_0. \quad (3.1)$$

By assumption, q_1 is not identically equal to q_2 . Since $s > \frac{d-1}{2}$, q_1 and q_2 are continuous on Γ_0 . Thus, we can find an open set $\kappa \subset \Gamma_0$ with a positive measure such that:

$$(q_1 - q_2)(x) \neq 0, \quad \forall x \in \kappa.$$

Equation 3.1 implies that $u_1 \equiv 0$ on $(0, T) \times \kappa$ and then u_1 is the solution of

$$\left\{ \begin{array}{ll} \partial_t u_1 - \Delta u_1 + \nabla p_1 & = 0, \quad \text{in } (0, T) \times \Omega, \\ \operatorname{div} u_1 & = 0, \quad \text{in } (0, T) \times \Omega, \\ u_1 & = 0, \quad \text{on } (0, T) \times \kappa, \\ \frac{\partial u_1}{\partial n} - p_1 n & = 0, \quad \text{on } (0, T) \times \kappa. \end{array} \right.$$

Applying again Corollary 3.2, we obtain that $u_1 = 0$ and $p_1 = 0$ in $(0, T) \times \Omega$. This is in contradiction with the fact that g is non identically zero. \square

4 Stability estimates

In this section, we assume that $d = 2$ and that the open set $\Omega \subset \mathbb{R}^2$ is of class $\mathcal{C}^{3,1}$.

We are going to prove stability estimates for the inverse problem we are interested in by using a global Carleman inequality which is stated in Lemma 4.1.

First, in Theorem 4.3, we state a stability estimate for the stationary problem. Then we deduce from this theorem two stability estimates for the evolution problem 1.1 by using an inequality coming from the analytic semigroup theory. To be more precise, we treat separately the case where g does not depend on time (see Theorem 4.18) and the case where g depends on time (see Theorem 4.21).

4.1 Carleman inequality

Let us state a global Carleman inequality proved by A. L. Bukhgeim in [12]:

Lemma 4.1. *Let $\Psi \in \mathcal{C}^2(\overline{\Omega})$. We have:*

$$\begin{aligned} \int_{\Omega} (\Delta \Psi |u|^2 + (\Delta \Psi - 1) |\nabla u|^2) e^{\Psi} \\ \leq \int_{\Omega} |\Delta u|^2 e^{\Psi} + \int_{\partial \Omega} \frac{\partial \Psi}{\partial n} \left(|u|^2 + |\nabla u|^2 + 2 \left| \frac{\partial |\nabla u|^2}{\partial \tau} \right| \right) e^{\Psi} \end{aligned} \quad (4.1)$$

for all $u \in \mathcal{C}^2(\overline{\Omega})$.

The proof of this result, which is only valid in dimension 2, uses computational properties of function defined on \mathbb{C} (in particular, the fact that $4\partial_{\bar{z}}\partial_z = \Delta$).

Remark 4.2. *The result is still true for $u \in H^3(\Omega)$. Indeed, for all $u \in H^3(\Omega)$, there exists $(u_n)_{n \in \mathbb{N}} \in \mathcal{C}^2(\overline{\Omega})^{\mathbb{N}}$ such that*

$$u_n \rightarrow u \text{ in } H^3(\Omega). \quad (4.2)$$

We can apply Lemma 4.1 to u_n , for all $n \in \mathbb{N}$. Let us prove that:

$$\lim_{n \rightarrow \infty} \int_{\partial \Omega} \frac{\partial \Psi}{\partial n} \left| \frac{\partial |\nabla u_n|^2}{\partial \tau} \right| e^{\Psi} = \int_{\partial \Omega} \frac{\partial \Psi}{\partial n} \left| \frac{\partial |\nabla u|^2}{\partial \tau} \right| e^{\Psi}. \quad (4.3)$$

Note first that $\int_{\partial \Omega} \frac{\partial \Psi}{\partial n} \left| \frac{\partial |\nabla u|^2}{\partial \tau} \right| e^{\Psi}$ has a meaning for $u \in H^3(\Omega)$:

$$\int_{\partial \Omega} \frac{\partial \Psi}{\partial n} \left| \frac{\partial |\nabla u|^2}{\partial \tau} \right| e^{\Psi} \leq 2 \|\Psi\|_{\mathcal{C}^1(\overline{\Omega})} \|e^{\Psi}\|_{\mathcal{C}^0(\overline{\Omega})} \left(\sum_{i=1}^2 \int_{\partial \Omega} |\nabla u| \cdot |\nabla \partial_i u| \right) < \infty.$$

We have:

$$\begin{aligned} \int_{\partial \Omega} \left| \frac{\partial \Psi}{\partial n} \right| \left| \left| \frac{\partial |\nabla u_n|^2}{\partial \tau} \right| - \left| \frac{\partial |\nabla u|^2}{\partial \tau} \right| \right| e^{\Psi} \\ \leq C \|\Psi\|_{\mathcal{C}^1(\overline{\Omega})} \|e^{\Psi}\|_{\mathcal{C}^0(\overline{\Omega})} \sum_{i,j=1}^2 \left(\int_{\partial \Omega} |\partial_j u|^2 \right)^{\frac{1}{2}} \left(\int_{\partial \Omega} |\partial_{ij} u_n - \partial_{ij} u|^2 \right)^{\frac{1}{2}} \\ + C \|\Psi\|_{\mathcal{C}^1(\overline{\Omega})} \|e^{\Psi}\|_{\mathcal{C}^0(\overline{\Omega})} \sum_{i,j=1}^2 \left(\int_{\partial \Omega} |\partial_{ij} u_n|^2 \right)^{\frac{1}{2}} \left(\int_{\partial \Omega} |\partial_j u_n - \partial_j u|^2 \right)^{\frac{1}{2}}. \end{aligned}$$

According to 4.2, the sequence $(\partial_{ij} u_n)_{n \in \mathbb{N}}$ converges in $L^2(\partial \Omega)$ towards $\partial_{ij} u$ and $\|\partial_{ij} u_n\|_{L^2(\partial \Omega)}$ is bounded by a constant independent of n . Then, equality 4.3 follows from 4.2.

4.2 The stationary case

For the stationary problem:

$$\begin{cases} -\Delta u + \nabla p &= 0, & \text{in } \Omega, \\ \operatorname{div} u &= 0, & \text{in } \Omega, \\ \frac{\partial u}{\partial n} - pn &= g, & \text{on } \Gamma_e, \\ \frac{\partial u}{\partial n} - pn + qu &= 0, & \text{on } \Gamma_0, \end{cases} \quad (4.4)$$

we have the following stability estimate.

Theorem 4.3. *Let $\alpha > 0$, $M_1 > 0$, $M_2 > 0$, $(g, q_j) \in H^{\frac{5}{2}}(\Gamma_e)^2 \times H^{\frac{5}{2}}(\Gamma_0)$ for $j = 1, 2$ be such that g is not identically zero, $\|g\|_{H^{\frac{5}{2}}(\Gamma_e)} \leq M_1$, $q_j \geq \alpha$ on Γ_0 and $\|q_j\|_{H^{\frac{5}{2}}(\Gamma_0)} \leq M_2$. We denote by (u_j, p_j) the solution of system 4.4 associated to $q = q_j$ for $j = 1, 2$. Let K be a compact subset of $\{x \in \Gamma_0 / u_1(x) \neq 0\}$ and $m > 0$ be a constant such that $|u_1| \geq m$ on K .*

Then there exist positive constants $C(M_1, M_2, \alpha)$ and $C_1(M_1, M_2, \alpha)$ such that

$$\|q_1 - q_2\|_{L^2(K)} \leq \frac{1}{m} \frac{C(M_1, M_2, \alpha)}{\left(\ln \left(\frac{C_1(M_1, M_2, \alpha)}{\|u_1 - u_2\|_{L^2(\Gamma_e)}^2 + \|p_1 - p_2\|_{L^2(\Gamma_e)} + \left\| \frac{\partial p_1}{\partial n} - \frac{\partial p_2}{\partial n} \right\|_{L^2(\Gamma_e)} \right)} \right)^{\frac{1}{2}}. \quad (4.5)$$

Remark 4.4. *Since g is not identically zero, Corollary 3.2 ensures that $\{x \in \Gamma_0 / u_1(x) \neq 0\}$ is not empty. Moreover, according to Proposition 2.4, u_1 is continuous on $\bar{\Omega}$, thus we obtain the existence of a compact K and a constant m as in Theorem 4.3. We notice however that the constants involved in the estimate 4.5 and the set K depend on u_1 . Finding a uniform lower bound for any solution u of system 4.4 remains an open question. We refer to [14], [2] and [3] for the case of the scalar Laplace equation.*

Remark 4.5. *In [15], the same kind of inequality is proved for the Laplacian problem with Robin boundary conditions under the hypothesis that the measurements are small enough. Here, we free ourselves from this smallness assumption on the measurements.*

Remark 4.6. *If we compare this result with the identifiability property (Proposition 3.3), we notice that we need additional measurements on the solution. In Proposition 3.3, we only have to assume that $u_1 = u_2$ and $\frac{\partial u_1}{\partial n} - p_1 n = \frac{\partial u_2}{\partial n} - p_2 n$ on $\Gamma \subseteq \Gamma_e$, where Γ is a non-empty open part of the boundary, in order to get the identifiability of the Robin coefficient q on Γ_0 . Here, besides a measurement on $u_1 - u_2$, we need measurements on $\frac{\partial u_1}{\partial n} - \frac{\partial u_2}{\partial n}$ (or $p_1 - p_2$) and $\frac{\partial p_1}{\partial n} - \frac{\partial p_2}{\partial n}$.*

Let us begin by proving this intermediate result which gives us a logarithmic estimate of the traces of u , ∇u , p , ∇p on Γ_0 with respect to the ones on Γ_e .

Lemma 4.7. *Let $(u, p) \in H^4(\Omega)^2 \times H^3(\Omega)$ be the solution in Ω of*

$$\begin{cases} -\Delta u + \nabla p &= 0, \\ \operatorname{div} u &= 0. \end{cases}$$

Then, there exist $C > 0$, $C_1 > 0$ and $d_0 > 0$ such that for all $\tilde{d} > d_0$:

$$\begin{aligned} & \|u\|_{L^2(\Gamma_0)^2} + \|\nabla u\|_{L^2(\Gamma_0)^4} + \|p\|_{L^2(\Gamma_0)} + \|\nabla p\|_{L^2(\Gamma_0)^2} \\ & \leq \tilde{d} C \frac{\|u\|_{H^3(\Omega)^2} + \|p\|_{H^3(\Omega)}}{\left(\ln \left(C_1 \tilde{d}^2 \frac{\|u\|_{H^3(\Omega)^2} + \|p\|_{H^3(\Omega)}}{\|u\|_{L^2(\Gamma_e)^2} + \left\| \frac{\partial u}{\partial n} \right\|_{L^2(\Gamma_e)}^2 + \|p\|_{L^2(\Gamma_e)} + \left\| \frac{\partial p}{\partial n} \right\|_{L^2(\Gamma_e)} \right)} \right)^{\frac{1}{2}}. \end{aligned} \quad (4.6)$$

Proof of Lemma 4.7. The proof is based on the Carleman inequality of Lemma 4.1 for an appropriate choice of Ψ . Note that we will apply 4.1 twice: one time for the velocity u and one time for the pressure p . The weight function Ψ is chosen in order to estimate the traces on Γ_0 with respect to the ones on Γ_e .

Step 1: choice of Ψ .

We choose Ψ as in [15]. There exists $\Psi_0 \in \mathcal{C}^2(\overline{\Omega})$ non identically zero such that:

$$\Delta\Psi_0 = 0 \text{ in } \Omega, \quad \Psi_0 = 0 \text{ on } \Gamma_0, \quad \Psi_0 \geq 0 \text{ on } \Gamma_e, \quad \frac{\partial\Psi_0}{\partial n} < 0 \text{ on } \Gamma_0.$$

Indeed, let $\chi \in \mathcal{C}^2(\partial\Omega)$ such that

$$\chi = 0 \text{ on } \Gamma_0, \quad \chi \geq 0 \text{ on } \Gamma_e,$$

and χ non identically zero on Γ_e . The boundary value problem :

$$\begin{cases} \Delta\Psi_0 = 0, & \text{in } \Omega, \\ \Psi_0 = \chi, & \text{on } \partial\Omega, \end{cases} \quad (4.7)$$

has a unique solution $\Psi_0 \in \mathcal{C}^2(\overline{\Omega})$. Note that Ψ_0 is not constant because χ is non identically zero. So, from the strong maximum principle, $\Psi_0 > 0$ in Ω . According to Hopf Lemma, we have $\frac{\partial\Psi_0}{\partial n} < 0$ on Γ_0 .

Let $\lambda > 0$. We denote by $\Psi_1 \in \mathcal{C}^2(\overline{\Omega})$ the unique solution of the boundary value problem:

$$\begin{cases} \Delta\Psi_1 = \lambda, & \text{in } \Omega, \\ \Psi_1 = 0, & \text{on } \partial\Omega. \end{cases}$$

From the comparison principle and the strong maximum principle, we have $\Psi_1 < 0$ in Ω . Moreover, according to the Hopf Lemma, we have $\frac{\partial\Psi_1}{\partial n} > 0$ on $\partial\Omega$.

Let us consider $\Psi = \Psi_1 + s\Psi_0$, for $s > 0$. To summarize, the function Ψ has the following properties:

$$\Delta\Psi = \lambda \text{ in } \Omega, \quad \Psi = 0 \text{ on } \Gamma_0, \quad \Psi \geq 0 \text{ on } \Gamma_e, \quad \text{and} \quad s\frac{\partial\Psi_0}{\partial n} \leq \frac{\partial\Psi}{\partial n} \leq \frac{\partial\Psi_1}{\partial n} \text{ on } \Gamma_0.$$

Step 2: We first apply Lemma 4.1 to u . Using the fact that $\Delta u = \nabla p$, we have:

$$\begin{aligned} & \int_{\Omega} (\Delta\Psi|u|^2 + (\Delta\Psi - 1)|\nabla u|^2)e^{\Psi} \\ & \leq \int_{\Omega} |\nabla p|^2 e^{\Psi} + \int_{\partial\Omega} \frac{\partial\Psi}{\partial n} \left(|u|^2 + |\nabla u|^2 + 2 \left| \frac{\partial|\nabla u|^2}{\partial\tau} \right| \right) e^{\Psi}. \end{aligned} \quad (4.8)$$

Then, we apply once again Lemma 4.1 to p :

$$\begin{aligned} & \int_{\Omega} (\Delta\Psi|p|^2 + (\Delta\Psi - 1)|\nabla p|^2)e^{\Psi} \\ & \leq \int_{\Omega} |\Delta p|^2 e^{\Psi} + \int_{\partial\Omega} \frac{\partial\Psi}{\partial n} \left(|p|^2 + |\nabla p|^2 + 2 \left| \frac{\partial|\nabla p|^2}{\partial\tau} \right| \right) e^{\Psi}. \end{aligned} \quad (4.9)$$

We have $\Delta p = \operatorname{div}(\Delta u) = 0$ hence $\int_{\Omega} |\Delta p|^2 e^{\Psi} = 0$. We now choose $\lambda \geq 2$. By summing up inequalities 4.8 and 4.9 and by eliminating the integrals on Ω in the left hand side which are positive terms, we obtain:

$$\int_{\partial\Omega} \frac{\partial\Psi}{\partial n} \left(|u|^2 + |\nabla u|^2 + 2 \left| \frac{\partial|\nabla u|^2}{\partial\tau} \right| \right) e^\Psi + \int_{\partial\Omega} \frac{\partial\Psi}{\partial n} \left(|p|^2 + |\nabla p|^2 + 2 \left| \frac{\partial|\nabla p|^2}{\partial\tau} \right| \right) e^\Psi \geq 0.$$

We now specify the dependence with respect to s . We denote by $\theta = \min_{\Gamma_0} \left| \frac{\partial\Psi_0}{\partial n} \right|$. We note that on Γ_0 , $e^\Psi = 1$. Consequently:

$$\begin{aligned} & -s\theta \int_{\Gamma_0} (|u|^2 + |\nabla u|^2 + |p|^2 + |\nabla p|^2) + \int_{\Gamma_0} \frac{\partial\Psi_1}{\partial n} (|u|^2 + |\nabla u|^2 + |p|^2 + |\nabla p|^2) \\ & + 2 \int_{\Gamma_0} \frac{\partial\Psi}{\partial n} \left(\left| \frac{\partial|\nabla p|^2}{\partial\tau} \right| + \left| \frac{\partial|\nabla u|^2}{\partial\tau} \right| \right) + 2 \int_{\Gamma_e} \frac{\partial\Psi}{\partial n} \left(\left| \frac{\partial|\nabla p|^2}{\partial\tau} \right| + \left| \frac{\partial|\nabla u|^2}{\partial\tau} \right| \right) e^\Psi \\ & + \int_{\Gamma_e} \frac{\partial\Psi}{\partial n} (|u|^2 + |\nabla u|^2 + |p|^2 + |\nabla p|^2) e^\Psi \geq 0. \end{aligned} \quad (4.10)$$

Let us study each of the terms. We have:

$$\int_{\Gamma_0} \frac{\partial\Psi_1}{\partial n} (|u|^2 + |\nabla u|^2 + |p|^2 + |\nabla p|^2) \leq C(\|u\|_{H^3(\Omega)}^2 + \|p\|_{H^3(\Omega)}^2).$$

Moreover, since $\frac{\partial\Psi}{\partial n} \leq \frac{\partial\Psi_1}{\partial n}$ on Γ_0 , we obtain:

$$2 \int_{\Gamma_0} \frac{\partial\Psi}{\partial n} \left(\left| \frac{\partial|\nabla p|^2}{\partial\tau} \right| + \left| \frac{\partial|\nabla u|^2}{\partial\tau} \right| \right) \leq C(\|u\|_{H^3(\Omega)}^2 + \|p\|_{H^3(\Omega)}^2).$$

Since, on Γ_e , $\left| \frac{\partial\Psi}{\partial n} \right| \leq sC$ for $s \geq 1$, we have:

$$\int_{\Gamma_e} \frac{\partial\Psi}{\partial n} (|u|^2 + |\nabla u|^2 + |p|^2 + |\nabla p|^2) e^\Psi \leq Cs \int_{\Gamma_e} (|u|^2 + |\nabla u|^2 + |p|^2 + |\nabla p|^2) e^\Psi.$$

Using Cauchy-Schwarz inequality, we obtain:

$$\begin{aligned} & 2 \int_{\Gamma_e} \frac{\partial\Psi}{\partial n} \left(\left| \frac{\partial|\nabla p|^2}{\partial\tau} \right| + \left| \frac{\partial|\nabla u|^2}{\partial\tau} \right| \right) e^\Psi \\ & \leq sC(\|u\|_{H^3(\Omega)}^2 + \|p\|_{H^3(\Omega)}^2) \left(\int_{\Gamma_e} (|\nabla p|^2 + |\nabla u|^2) e^{2\Psi} \right)^{\frac{1}{2}}. \end{aligned}$$

Note that e^Ψ depends on s on Γ_e . Hence, reassembling these inequalities, inequality 4.10 becomes:

$$\theta \int_{\Gamma_0} (|u|^2 + |\nabla u|^2 + |p|^2 + |\nabla p|^2) \leq C \left(K_s + \frac{1}{s} (\|u\|_{H^3(\Omega)}^2 + \|p\|_{H^3(\Omega)}^2) \right), \quad (4.11)$$

where

$$\begin{aligned} K_s &= (\|u\|_{H^3(\Omega)}^2 + \|p\|_{H^3(\Omega)}^2) \left(\int_{\Gamma_e} (|\nabla p|^2 + |\nabla u|^2) e^{2\Psi} \right)^{\frac{1}{2}} \\ & \quad + \int_{\Gamma_e} (|u|^2 + |\nabla u|^2 + |p|^2 + |\nabla p|^2) e^\Psi. \end{aligned} \quad (4.12)$$

In order to study the dependence with respect to s of K_s , we define B by:

$$B = \|u\|_{L^2(\Gamma_e)}^2 + \|p\|_{L^2(\Gamma_e)} + \left\| \frac{\partial u}{\partial n} \right\|_{L^2(\Gamma_e)}^2 + \left\| \frac{\partial p}{\partial n} \right\|_{L^2(\Gamma_e)}. \quad (4.13)$$

Let us estimate the first term in the expression of K_s . Remark that, thanks to classical interpolation inequalities (see [1]), there exists $C > 0$ such that for all $f \in H^2(\Gamma_\varepsilon)$:

$$\left\| \frac{\partial f}{\partial \tau} \right\|_{L^2(\Gamma_\varepsilon)} \leq \|f\|_{H^1(\Gamma_\varepsilon)} \leq C \|f\|_{L^2(\Gamma_\varepsilon)}^{\frac{1}{2}} \|f\|_{H^2(\Gamma_\varepsilon)}^{\frac{1}{2}}.$$

Applying the previous inequality, there exists $C > 0$ such that:

$$\int_{\Gamma_\varepsilon} \left| \frac{\partial u}{\partial \tau} \right|^2 \leq C \|u\|_{H^3(\Omega)^2} \|u\|_{L^2(\Gamma_\varepsilon)^2} \quad \text{and} \quad \int_{\Gamma_\varepsilon} \left| \frac{\partial p}{\partial \tau} \right|^2 \leq C \|p\|_{H^3(\Omega)} \|p\|_{L^2(\Gamma_\varepsilon)}. \quad (4.14)$$

We obtain, using the fact that $\nabla v = \frac{\partial v}{\partial n} n + \frac{\partial v}{\partial \tau} \tau$ on $\partial\Omega$ for all $v \in H^2(\Omega)$, and thanks to inequality 4.14, that there exists $C > 0$ such that:

$$\begin{aligned} \int_{\Gamma_\varepsilon} (|\nabla p|^2 + |\nabla u|^2) e^{2\Psi} &\leq e^{2ks} \left(\left\| \frac{\partial p}{\partial n} \right\|_{L^2(\Gamma_\varepsilon)}^2 + \left\| \frac{\partial u}{\partial n} \right\|_{L^2(\Gamma_\varepsilon)^2}^2 \right) \\ &\quad + e^{2ks} (C \|u\|_{H^3(\Omega)^2} \|u\|_{L^2(\Gamma_\varepsilon)^2} + C \|p\|_{H^3(\Omega)} \|p\|_{L^2(\Gamma_\varepsilon)}) \\ &\leq C e^{2ks} \|u\|_{H^3(\Omega)^2} \left(\left\| \frac{\partial u}{\partial n} \right\|_{L^2(\Gamma_\varepsilon)^2} + \|u\|_{L^2(\Gamma_\varepsilon)^2} \right) \\ &\quad + C e^{2ks} \|p\|_{H^3(\Omega)} \left(\left\| \frac{\partial p}{\partial n} \right\|_{L^2(\Gamma_\varepsilon)} \|p\|_{L^2(\Gamma_\varepsilon)} \right) \\ &\leq C e^{ks} (\|u\|_{H^3(\Omega)^2} + \|p\|_{H^3(\Omega)}) B, \end{aligned}$$

where $k = \max_{\Gamma} \Psi_0$. Similarly, for the second term in the expression of K_s we prove that

$$\begin{aligned} \int_{\Gamma_\varepsilon} (|u|^2 + |\nabla u|^2 + |p|^2 + |\nabla p|^2) e^\Psi &\leq C e^{ks} (\|u\|_{H^3(\Omega)^2} + \|p\|_{H^3(\Omega)}) B \\ &\leq C e^{ks} (\|u\|_{H^3(\Omega)^2} + \|p\|_{H^3(\Omega)})^{\frac{3}{2}} B^{\frac{1}{2}}. \end{aligned}$$

Thus, using the two previous inequalities, according to the definition 4.12 of K_s , we obtain,

$$K_s \leq C e^{ks} (\|u\|_{H^3(\Omega)^2} + \|p\|_{H^3(\Omega)})^{\frac{3}{2}} B^{\frac{1}{2}}.$$

Let us denote by

$$A = \|u\|_{H^3(\Omega)^2} + \|p\|_{H^3(\Omega)}.$$

Hence we get from 4.11:

$$\int_{\Gamma_0} (|u|^2 + |\nabla u|^2 + |p|^2 + |\nabla p|^2) \leq C \left(A^{\frac{3}{2}} e^{ks} B^{\frac{1}{2}} + \frac{A^2}{s} \right),$$

for all $s \geq 1$. Remark that this inequality is trivially verified for $0 < s \leq 1$ by continuity of the trace mapping. Let $\tilde{d} \geq 1$. To summarize, we have proved that:

$$\int_{\Gamma_0} (|u|^2 + |\nabla u|^2 + |p|^2 + |\nabla p|^2) \leq C A^{\frac{3}{2}} \left(e^{ks} B^{\frac{1}{2}} + \frac{\tilde{d} A^{\frac{1}{2}}}{s} \right), \quad \forall s > 0.$$

We now optimize the upper bound with respect to s . We denote by

$$f(s) = e^{ks} B^{\frac{1}{2}} + \frac{\tilde{d} A^{\frac{1}{2}}}{s}.$$

Let us study the function f in \mathbb{R}_+^* . We have:

$$\begin{cases} \lim_{s \rightarrow 0} f(s) = +\infty, \\ \lim_{s \rightarrow \infty} f(s) = +\infty. \end{cases}$$

So since f is continuous on \mathbb{R}_*^+ , f reaches its minimum at a point $s_0 > 0$. At this point,

$$f'(s_0) = 0 \Leftrightarrow B^{\frac{1}{2}} = \frac{e^{-ks_0} \tilde{d} A^{\frac{1}{2}}}{ks_0^2}, \text{ thus } f(s_0) = \frac{\tilde{d} A^{\frac{1}{2}}}{ks_0^2} + \frac{\tilde{d} A^{\frac{1}{2}}}{s_0}.$$

Hence:

$$\int_{\Gamma_0} (|u|^2 + |\nabla u|^2 + |p|^2 + |\nabla p|^2) \leq \frac{C \tilde{d} A^2}{s_0^\beta} \left(\frac{1}{k} + 1 \right),$$

where $\beta = 1$ if $s_0 \geq 1$ and $\beta = 2$ otherwise. But, we notice that

$$\frac{1}{B^{\frac{1}{2}}} = \frac{ks_0^2 e^{ks_0}}{\tilde{d} A^{\frac{1}{2}}} \leq \frac{ke^{(k+2)s_0}}{\tilde{d} A^{\frac{1}{2}}},$$

that is to say:

$$\frac{1}{s_0} \leq \frac{k+2}{\ln \left(\frac{\tilde{d} A^{\frac{1}{2}}}{kB^{\frac{1}{2}}} \right)},$$

if \tilde{d} is larger than a constant which only depends on k and on the continuity constants of the trace mapping. In the same way, when $s_0 < 1$, we obtain:

$$\frac{1}{s_0^2} \leq \frac{1}{\ln \left(\frac{\tilde{d} A^{\frac{1}{2}}}{ke^k B^{\frac{1}{2}}} \right)},$$

if \tilde{d} is larger than a constant which only depends on k and on the continuity constants of the trace mapping. Using the fact that $\ln(x^{\frac{1}{2}}) = \frac{1}{2} \ln(x)$ for all $x > 0$ and according to the definition 4.13 of B , the desired result follows. \square

Remark 4.8. Let $\Gamma \subseteq \Gamma_e$ be a non-empty open part of Γ_e . Inequality 4.6 of Lemma 4.7 still holds if we replace Γ_e by Γ in the right-hand side. To prove this, it is sufficient in the definition 4.7 of Ψ_0 to define $\chi \in C^2(\partial\Omega)$ such that $\chi = 0$ on $\Gamma_0 \cup (\Gamma_e \setminus \bar{\Gamma})$ and $\chi \geq 0$ on Γ .

Let us now prove Theorem 4.3.

Proof of Theorem 4.3. Since $g \in H^{\frac{5}{2}}(\Gamma_e)^2$ and $q_j \in H^{\frac{5}{2}}(\Gamma_0)$ for $j = 1, 2$, thanks to Proposition 2.4 applied for $k = 2$, there exists $C(\alpha, M_1, M_2) > 0$ such that:

$$\|u_j\|_{H^4(\Omega)^2} + \|p_j\|_{H^3(\Omega)} \leq C(\alpha, M_1, M_2), \text{ for } j = 1, 2. \quad (4.15)$$

In the following, we denote by $u = u_1 - u_2$ and $p = p_1 - p_2$. We have:

$$(q_2 - q_1)u_1 = q_2 u + \frac{\partial u}{\partial n} - p n, \text{ on } \Gamma_0. \quad (4.16)$$

Consequently, since $|u_1| \geq m > 0$ on K :

$$\|q_1 - q_2\|_{L^2(K)} \leq \frac{1}{m} C(M_2) \left(\|u\|_{L^2(\Gamma_0)^2} + \left\| \frac{\partial u}{\partial n} \right\|_{L^2(\Gamma_0)^2} + \|p\|_{L^2(\Gamma_0)} \right). \quad (4.17)$$

Let us denote by

$$A = \|u\|_{H^3(\Omega)^2} + \|p\|_{H^3(\Omega)}$$

and

$$B = \|u\|_{L^2(\Gamma_e)^2} + \|p\|_{L^2(\Gamma_e)} + \left\| \frac{\partial u}{\partial n} \right\|_{L^2(\Gamma_e)^2} + \left\| \frac{\partial p}{\partial n} \right\|_{L^2(\Gamma_e)}.$$

By applying Lemma 4.7, we obtain that there exists there exists $C(M_2) > 0$, $C_1 > 0$ and $d_0 > 0$ such that for all $\tilde{d} > d_0$:

$$\|q_1 - q_2\|_{L^2(K)} \leq \frac{\tilde{d}C(M_2)}{m} \frac{A}{\left(\ln\left(C_1\tilde{d}^2\frac{A}{B}\right)\right)^{\frac{1}{2}}}. \quad (4.18)$$

We are going to conclude the proof by studying the variation of the function defined by $f_y(x) = \frac{x}{\left(\ln\left(\frac{x}{y}\right)\right)^{\frac{1}{2}}}$ on $(y, +\infty)$, for $y = \frac{B}{C_1\tilde{d}^2}$. We have

$$f'_y(x) = \frac{\ln\left(\frac{x}{y}\right) - \frac{1}{2}}{\ln\left(\frac{x}{y}\right)^{\frac{3}{2}}}.$$

Let us denote by $x_0 = ye^{\frac{1}{2}}$. The function f_y is decreasing on $(y, x_0]$ and is increasing on $[x_0, +\infty)$. For \tilde{d} large enough, we have $A \geq x_0$ by continuity of the trace mapping. Using 4.15 and since f_y is increasing on $[x_0, +\infty)$, we directly deduce that $f_y(A) \leq f_y(C(\alpha, M_1, M_2))$.

Using this result in inequality 4.18, we get that there exist constants $C(\alpha, M_1, M_2) > 0$ and $C_1(\alpha, M_1, M_2) > 0$ such that:

$$\|q_1 - q_2\|_{L^2(K)} \leq \frac{1}{m} \frac{C(\alpha, M_1, M_2)}{\left(\ln\left(\frac{C_1(\alpha, M_1, M_2)}{\|u\|_{L^2(\Gamma_e)^2} + \|p\|_{L^2(\Gamma_e)} + \left\|\frac{\partial u}{\partial n}\right\|_{L^2(\Gamma_e)^2} + \left\|\frac{\partial p}{\partial n}\right\|_{L^2(\Gamma_e)}\right)}\right)^{\frac{1}{2}}}.$$

Since $\frac{\partial u}{\partial n} = pn$ on Γ_e , we obtain the desired inequality. \square

Remark 4.9. Note that the assumption that $|u_1| \geq m > 0$ on K is essential to pass from 4.16 to 4.17. Outside the set K , an estimate of $q_1 - q_2$ may be undetermined or highly unstable. In particular, an estimate of the Robin coefficients on the whole set Γ_0 might be worst than of logarithmic type (see [6]).

Note however that for a simplified problem, it is in fact possible to obtain a logarithmic stability estimate on the whole set Γ_0 which does not depend on a given reference solution. Assume that $g = g_e n$ and $q \in \mathbb{R}$ are such that

$$(A) \quad g_e \in \mathbb{R} \text{ satisfies } \beta \leq g_e \leq M_1,$$

$$(B) \quad \alpha \leq q \leq M_2,$$

for some $\alpha > 0$, $\beta > 0$, $M_1 > 0$ and $M_2 > 0$.

We denote by $(u_{g_e, q}, p_{g_e, q})$ the solution of system 4.4 associated to q and $g = g_e n$. Thanks to the weak formulation of the problem, $\int_{\Gamma_e} u_{g_e, q} \cdot n > 0$. Moreover, one can prove by contradiction and thanks to the continuity of the solution with respect to the data that there exists $m_1 > 0$ which depends on M_1 , M_2 , α and β such that for all $(g_e, q) \in \mathbb{R}^2$ which satisfies (A) and (B),

$$\int_{\Gamma_e} u_{g_e, q} \cdot n \geq m_1.$$

For $i = 1, 2$, let $q_i \in \mathbb{R}$ satisfy the assumption (B) above. We define by $(u_i, p_i) = (u_{g_e, q_i}, p_{g_e, q_i})$ the solution of system 4.4 associated with $g = g_e n$ and $q = q_i$ for $i=1, 2$. If we multiply 4.16 by the unit normal n and we integrate on Γ_0 , we obtain:

$$(q_2 - q_1) \int_{\Gamma_0} u_1 \cdot n = q_2 \int_{\Gamma_0} (u_1 - u_2) \cdot n + \int_{\Gamma_0} \left(\frac{\partial u_1}{\partial n} - \frac{\partial u_2}{\partial n} \right) \cdot n - \int_{\Gamma_0} (p_1 - p_2).$$

Since u_1 is divergence free, $\left| \int_{\Gamma_0} u_1 \cdot n \right| = \left| \int_{\Gamma_e} u_1 \cdot n \right| \geq m_1$. Thus, we get

$$\begin{aligned} & |q_1 - q_2| \\ & \leq C(M_1, M_2, \alpha, \beta) \left(\|u_1 - u_2\|_{L^2(\Gamma_0)^2} + \left\| \frac{\partial u_1}{\partial n} - \frac{\partial u_2}{\partial n} \right\|_{L^2(\Gamma_0)^2} + \|p_1 - p_2\|_{L^2(\Gamma_0)} \right). \end{aligned}$$

We conclude as in the proof of Theorem 4.3 and obtain that positive constants $C(M_1, M_2, \alpha)$ and $C_1(M_1, M_2, \alpha, \beta)$ such that

$$|q_1 - q_2| \leq \frac{C(M_1, M_2, \alpha, \beta)}{\left(\ln \left(\frac{C_1(M_1, M_2, \alpha)}{\|u_1 - u_2\|_{L^2(\Gamma_e)^2} + \|p_1 - p_2\|_{L^2(\Gamma_e)} + \left\| \frac{\partial p_1}{\partial n} - \frac{\partial p_2}{\partial n} \right\|_{L^2(\Gamma_e)} \right)} \right)^{\frac{1}{2}}}.$$

4.3 Evolution problem

In order to use semigroup properties, we begin by introducing the Stokes operator associated with the Robin boundary conditions on Γ_0 .

4.3.1 Properties of the Stokes operator

We recall that the bilinear form a_q is defined by 2.3.

Definition 4.10. We define the set $\mathcal{D}(A_q)$ as follows:

$$\mathcal{D}(A_q) = \{u \in V / \exists C > 0, \forall v \in V, |a_q(u, v)| \leq C \|v\|_{L^2(\Omega)^2}\},$$

and the operator $A_q : \mathcal{D}(A_q) \subset H \rightarrow H$ by:

$$\forall u \in \mathcal{D}(A_q), a_q(u, v) = (A_q u, v)_{L^2(\Omega)^2}, \forall v \in V.$$

Proposition 4.11. Let $\alpha > 0$ and $q \in L^\infty(\Gamma_0)$ such that $q \geq \alpha$ almost everywhere on Γ_0 . The operator A_q has the following properties:

1. $A_q \in \mathcal{L}(\mathcal{D}(A_q), H)$ is invertible and its inverse is compact on H .
2. A_q is selfadjoint.

As a consequence, A_q admits a family of eigenvalues ϕ_q^l

$$A_q \phi_q^l = \lambda_q^l \phi_q^l \text{ with } 0 < \lambda_q^1 \leq \lambda_q^2 \leq \dots \leq \lambda_q^j \text{ and } \lim_{j \rightarrow \infty} \lambda_q^j = +\infty,$$

which is complete and orthogonal both in H and V .

Proof of Proposition 4.11. It relies on classical arguments for which we refer to [11] or [23]. \square

Remark 4.12. Let $\alpha > 0$. There exists a constants $\mu > 0$ such that for all $q \in L^\infty(\Gamma_0)$ such that $q \geq \alpha$, for $l \in \mathbb{N}$:

$$\lambda_q^l \geq \mu. \tag{4.19}$$

Indeed, $\lambda_q^l \geq \lambda_q^1 = (A_q \phi_q^1, \phi_q^1)_{L^2(\Omega)^2} = a_q(\phi_q^1, \phi_q^1) \geq a_\alpha(\phi_q^1, \phi_q^1) \geq \mu \|\phi_q^1\|_{L^2(\Omega)^2}^2 = \mu$, where μ is the coercivity constant associated with the bilinear form a_α .

Proposition 4.13. The operator $A_q^{\frac{1}{2}} : (V, a_q(\cdot, \cdot)^{\frac{1}{2}}) \rightarrow (H, \|\cdot\|_{L^2(\Omega)^2})$ is an isometry.

Proposition 4.14. *Let $\alpha > 0$ and $q \in L^\infty(\Gamma_0)$ be such that $q \geq \alpha$ almost everywhere on Γ_0 . The operator $-A_q$ generates an analytic semigroup on H . This analytic semigroup is explicitly given by:*

$$e^{-tA_q} f = \sum_{l \geq 1} e^{-t\lambda_q^l} (\phi_q^l, f)_{L^2(\Omega)^2} \phi_q^l, \quad (4.20)$$

for all $f \in H$.

Proof of Proposition 4.14. It follows from the construction of the operator A_q . We refer to [20] and [16] for details. \square

Proposition 4.15. *Let $\alpha > 0$, $M > 0$, $k \in \mathbb{N}$ and $s \in \mathbb{R}$ be such that $s > \frac{d-1}{2}$ and $s \geq \frac{1}{2} + k$. We assume that Ω is of class $C^{k+1,1}$ and that $q \in H^s(\Gamma_0)$ is such that $q \geq \alpha$ on Γ_0 .*

Then for each $f \in H \cap H^k(\Omega)^2$, there exists $u \in H^{k+2}(\Omega)^2$ solution of $A_q u = f$ if and only if there exists $p \in H^{k+1}(\Omega)$ such that (u, p) is solution of the following problem:

$$\begin{cases} -\Delta u + \nabla p & = f, & \text{in } \Omega, \\ \operatorname{div} u & = 0, & \text{in } \Omega, \\ \frac{\partial u}{\partial n} - pn & = 0, & \text{on } \Gamma_e, \\ \frac{\partial u}{\partial n} - pn + qu & = 0, & \text{on } \Gamma_0. \end{cases} \quad (4.21)$$

Moreover, there exists a constant $C(\alpha, M) > 0$ such that for every $q \in H^s(\Gamma_0)$ satisfying $\|q\|_{H^s(\Gamma_0)} \leq M$:

$$\|u\|_{H^{2+k}(\Omega)^2} \leq C(\alpha, M) \|f\|_{H^k(\Omega)^2}.$$

Proof of Proposition 4.15. This result follows from the construction of the operator A_q and from Proposition 2.4. \square

Corollary 4.16. *Let $\alpha > 0$, $k \in \mathbb{N}^*$ and $s \in \mathbb{R}$ be such that $s > \frac{d-1}{2}$ and $s \geq \frac{1}{2} + 2(k-1)$. We assume that Ω is of class $C^{2k-1,1}$ and that $q \in H^s(\Gamma_0)$ is such that $q \geq \alpha$ on Γ_0 .*

Then $\mathcal{D}(A_q^k) \hookrightarrow H^{2k}(\Omega)^2 \cap H$.

Proof of Corollary 4.16. For $k = 1$, it is clear. Take now $k = 2$. Let $u \in \mathcal{D}(A_q^2)$. We have

$$A_q^2 u = f \Leftrightarrow \begin{cases} A_q u = v \\ A_q v = f \end{cases}$$

But $v \in \mathcal{D}(A_q) \subset H^2(\Omega)^2 \cap H$ by assumption, so $u \in H^4(\Omega)^2 \cap H$ thanks to the regularity properties of the solution of the Stokes problem summarize in Proposition 2.4. We conclude by induction on k . \square

Remark 4.17. *Let us remark that, due to the prescribed boundary conditions, $\mathcal{D}(A_q)$ is not equal to $H^2(\Omega)^2 \cap H$.*

4.3.2 The flux g does not depend on t

In this paragraph, we consider the evolution problem 1.1 given in the introduction. We assume in this part that g does not depend on time. Let $\alpha > 0$, $M_1 > 0$ and $M_2 > 0$. In the following, we assume that

$$g \in H^{\frac{5}{2}}(\Gamma_e)^2 \text{ is non identically zero and } \|g\|_{H^{\frac{5}{2}}(\Gamma_e)^2} \leq M_1, \quad (4.22)$$

$$q \in H^{\frac{5}{2}}(\Gamma_0) \text{ is such that } \|q\|_{H^{\frac{5}{2}}(\Gamma_0)} \leq M_2 \text{ and } q \geq \alpha \text{ on } \Gamma_0. \quad (4.23)$$

Let us prove the following theorem:

Theorem 4.18. *Let $\alpha > 0$, $M_1 > 0$, $M_2 > 0$ and $u_0 \in H \cap H^3(\Omega)^2$. We assume that g satisfies 4.22 and that q_j satisfies 4.23 for $j = 1, 2$. We denote by (u_j, p_j) the solution of system 1.1 associated to $q = q_j$, for $j = 1, 2$. Let K be a compact subset of $\{x \in \Gamma_0 / v_1(x) \neq 0\}$, where (v_1, ζ_1) is the solution of system 4.4 with $q = q_1$ and let $m > 0$ be a constant such that $|v_1| \geq m$ on K . Then, there exist $C(\alpha, M_1, M_2) > 0$ and $C_1(M_1, M_2, \alpha) > 0$ such that*

$$\|q_1 - q_2\|_{L^2(K)} \leq \frac{1}{m} \frac{C(\alpha, M_1, M_2)}{\left(\ln \left(\frac{C_1(M_1, M_2, \alpha)}{\|u_1 - u_2\|_{L^\infty(0, +\infty; L^2(\Gamma_e)^2)} + \|p_1 - p_2\|_{L^\infty(0, +\infty; L^2(\Gamma_e))} + \left\| \frac{\partial p_1}{\partial n} - \frac{\partial p_2}{\partial n} \right\|_{L^\infty(0, +\infty; L^2(\Gamma_e))}} \right) \right)^{\frac{1}{2}}}.$$

Remark 4.19. *Due to the method which relies on semigroup theory, we need to take measurements during an infinite time.*

Proof of Theorem 4.18. For $j = 1, 2$, let (v_j, ζ_j) be the solution of the stationary problem 4.4 with $q = q_j$. According to Proposition 2.4, (v_j, ζ_j) belongs to $H^4(\Omega)^2 \times H^3(\Omega)$ and moreover, thanks to assumptions 4.22 and 4.23, there exists a constant $C(\alpha, M_1, M_2) > 0$ such that

$$\|v_j\|_{H^4(\Omega)^2} + \|\zeta_j\|_{H^3(\Omega)} \leq C(\alpha, M_1, M_2). \quad (4.24)$$

We denote $(w_j, \pi_j) = (u_j - v_j, p_j - \zeta_j)$. Thanks to Theorem 4.3, we are able to estimate $\|q_1 - q_2\|_{L^2(K)}$ with respect to an increasing function of $(v_1 - v_2)|_{\Gamma_e}$, $(\zeta_1 - \zeta_2)|_{\Gamma_e}$ and $\left(\frac{\partial \zeta_1}{\partial n} - \frac{\partial \zeta_2}{\partial n} \right)|_{\Gamma_e}$. Our objective is now to compare the asymptotic behavior of $u_1 - u_2$ and $p_1 - p_2$ to the solution of the stationary problem $v_1 - v_2$ and $\zeta_1 - \zeta_2$. More precisely, we are going to prove that:

$$\|w_j(t, \cdot)\|_{H^3(\Omega)^2} + \|\pi_j(t, \cdot)\|_{H^2(\Omega)} \leq G(t),$$

where G is a function which tends to 0 when t goes to $+\infty$. This inequality, combined with Theorem 4.3, will allow us to conclude the proof of Theorem 4.18.

We have that (w_j, π_j) is the solution of the following problem: for $t > 0$,

$$\begin{cases} \partial_t w - \Delta w + \nabla \pi &= 0, & \text{in } \Omega, \\ \operatorname{div} w &= 0, & \text{in } \Omega, \\ \frac{\partial w}{\partial n} - \pi n &= 0, & \text{on } \Gamma_e, \\ \frac{\partial w}{\partial n} - \pi n + q_j w &= 0, & \text{on } \Gamma_0, \end{cases}$$

completed with the initial condition $w(0) = u_0 - v_j$. Let $t > 0$. We have from the theory of analytic semigroup that:

$$w_j(t, \cdot) = e^{-tA_{q_j}} w_j(0, \cdot). \quad (4.25)$$

Let $\eta > 0$. There exists a constant $C > 0$ independent of q_j such that:

$$\|A_{q_j}^\eta e^{-tA_{q_j}}\| \leq C \frac{e^{-\mu t}}{t^\eta}, \quad t > 0, \eta > 0, \quad (4.26)$$

where μ is given by 4.19 and where $\| \cdot \|$ is the norm operator. Using regularity result for the stationary problem given in Proposition 2.4, we have that:

$$\|w_j(t, \cdot)\|_{H^3(\Omega)^2} + \|\pi_j(t, \cdot)\|_{H^2(\Omega)} \leq C(\alpha, M_2) \|\partial_t w_j(t, \cdot)\|_{H^1(\Omega)^2}.$$

Note that, thanks to Proposition 4.15 we have:

$$\|\partial_t w_j(t, \cdot)\|_{H^1(\Omega)^2} = \|A_{q_j} w_j(t, \cdot)\|_{H^1(\Omega)^2}.$$

Then, since $w_j(t, \cdot)$ is given by 4.25, and using Proposition 4.13 and estimates 4.24 and 4.26 with $\eta = \frac{3}{2}$, it follows:

$$\begin{aligned} \|w_j(t, \cdot)\|_{H^3(\Omega)^2} + \|\pi_j(t, \cdot)\|_{H^2(\Omega)} &\leq C(\alpha, M_2) \|A_q^{\frac{3}{2}} e^{-tA_{q_j}} w_j(0, \cdot)\|_{L^2(\Omega)^2} \\ &\leq C(\alpha, M_2) \frac{e^{-\mu t}}{t^{\frac{3}{2}}} (\|u_0\|_{L^2(\Omega)^2} + \|v_j\|_{L^2(\Omega)^2}) \\ &\leq C(\alpha, u_0, M_1, M_2) \frac{e^{-\mu t}}{t^{\frac{3}{2}}}. \end{aligned} \quad (4.27)$$

We have from 4.27:

$$\|v_1 - v_2\|_{L^2(\Gamma_e)^2} \leq C(\alpha, u_0, M_1, M_2) \frac{e^{-\mu t}}{t^{\frac{3}{2}}} + \|u_1 - u_2\|_{L^\infty(0, +\infty; L^2(\Gamma_e)^2)}.$$

Then, passing to the limit when t goes to infinity, we get:

$$\|v_1 - v_2\|_{L^2(\Gamma_e)^2} \leq \|u_1 - u_2\|_{L^\infty(0, +\infty; L^2(\Gamma_e)^2)}.$$

We prove similarly:

$$\|\zeta_1 - \zeta_2\|_{L^2(\Gamma_e)} \leq \|p_1 - p_2\|_{L^\infty(0, +\infty; L^2(\Gamma_e))},$$

and

$$\left\| \frac{\partial \zeta_1}{\partial n} - \frac{\partial \zeta_2}{\partial n} \right\|_{L^2(\Gamma_e)} \leq \left\| \frac{\partial p_1}{\partial n} - \frac{\partial p_2}{\partial n} \right\|_{L^\infty(0, +\infty; L^2(\Gamma_e))}.$$

To summarize, we have obtained:

$$\begin{aligned} &\|v_1 - v_2\|_{L^2(\Gamma_e)^2} + \|\zeta_1 - \zeta_2\|_{L^2(\Gamma_e)} + \left\| \frac{\partial \zeta_1}{\partial n} - \frac{\partial \zeta_2}{\partial n} \right\|_{L^2(\Gamma_e)} \\ &\leq \|u_1 - u_2\|_{L^\infty(0, +\infty; L^2(\Gamma_e)^2)} + \|p_1 - p_2\|_{L^\infty(0, +\infty; L^2(\Gamma_e))} + \left\| \frac{\partial p_1}{\partial n} - \frac{\partial p_2}{\partial n} \right\|_{L^\infty(0, +\infty; L^2(\Gamma_e))}. \end{aligned}$$

Applying Theorem 4.3 to (v_j, ζ_j) for $j = 1, 2$, we obtain the existence of positive constants $C(M_1, M_2, \alpha)$ and $C_1(M_1, M_2, \alpha)$ such that

$$\|q_1 - q_2\|_{L^2(K)} \leq \frac{1}{m} \frac{C(M_1, M_2, \alpha)}{\left(\ln \left(\frac{C_1(M_1, M_2, \alpha)}{\|v_1 - v_2\|_{L^2(\Gamma_e)^2} + \|\zeta_1 - \zeta_2\|_{L^2(\Gamma_e)} + \left\| \frac{\partial \zeta_1}{\partial n} - \frac{\partial \zeta_2}{\partial n} \right\|_{L^2(\Gamma_e)} \right)} \right)^{\frac{1}{2}}}.$$

We conclude by using the fact that the function $x \rightarrow \frac{1}{\ln(\frac{1}{x})}$ increases on \mathbb{R}_+^* . \square

Remark 4.20. *Remark that*

$$(u_j, p_j) \in L^\infty(0, +\infty; H^3(\Omega)^2) \times L^\infty(0, +\infty; H^2(\Omega)). \quad (4.28)$$

Let us prove 4.28. Let $\nu > 0$. In fact, thanks to equation 4.27, we obtain that

$$(w_j, \pi_j) \in L^\infty(\nu, +\infty; H^3(\Omega)^2) \times L^\infty(\nu, +\infty; H^2(\Omega)),$$

and since $u_j = w_j + v_j$ and $p_j = \pi_j + \zeta_j$, we deduce that

$$(u_j, p_j) \in L^\infty(\nu, +\infty; H^3(\Omega)^2) \times L^\infty(\nu, +\infty; H^2(\Omega)).$$

Moreover, thanks to Corollary 2.7, we have

$$(u_j, p_j) \in L^\infty(0, \nu; H^3(\Omega)^2) \times L^\infty(0, \nu; H^2(\Omega)).$$

Thus, 4.28 follows.

4.3.3 The flux g depends on t

We restrict our study to the case where g is colinear to the exterior unit normal n : $g = \kappa n$. Let $\alpha > 0$, $M_1 > 0$ and $M_2 > 0$. We assume that:

$$\kappa \in H_{loc}^2(0, +\infty; H^{\frac{3}{2}}(\Gamma_e)), \quad (4.29)$$

and

$$q \in H^{\frac{5}{2}}(\Gamma_0) \text{ is such that } \|q\|_{H^{\frac{5}{2}}(\Gamma_0)} \leq M_2 \text{ and } q \geq \alpha \text{ on } \Gamma_0. \quad (4.30)$$

Let us introduce h such that:

$$h \in H^{\frac{5}{2}}(\Gamma_e) \text{ is non identically zero and } \|h\|_{H^{\frac{5}{2}}(\Gamma_e)} \leq M_1. \quad (4.31)$$

We assume that:

$$\begin{aligned} & \lim_{t \rightarrow \infty} (\|\kappa(t, \cdot) - h\|_{H^{\frac{3}{2}}(\Gamma_e)} + \|\partial_t \kappa(t, \cdot)\|_{H^{\frac{3}{2}}(\Gamma_e)} \\ & + \left(\int_0^t e^{-\mu(t-s)} \|\partial_t \kappa(s, \cdot)\|_{H^{\frac{3}{2}}(\Gamma_e)}^2 ds \right)^{\frac{1}{2}}) = 0, \end{aligned} \quad (4.32)$$

where μ is given by equation 4.19.

Theorem 4.21. *Let $\alpha > 0$, $M_1 > 0$, $M_2 > 0$ and $u_0 \in H^3(\Omega)^2 \cap H$. We assume that h and κ satisfy respectively 4.31 and 4.29 and for $j = 1, 2$, q_j satisfies 4.30. We denote by (u_j, p_j) the solution of system 1.1 associated to $q = q_j$. Let K be a compact subset of $\{x \in \Gamma_0 / v_1(x) \neq 0\}$, where (v_1, ζ_1) is the solution of*

$$\begin{cases} -\Delta v + \nabla \zeta & = 0, & \text{in } \Omega, \\ \operatorname{div} v & = 0, & \text{in } \Omega, \\ \frac{\partial v}{\partial n} - \zeta n & = hn, & \text{on } \Gamma_e, \\ \frac{\partial v}{\partial n} - \zeta n + q_1 v & = 0, & \text{on } \Gamma_0, \end{cases}$$

and let $m > 0$ be a constant such that $|v_1| > m$ on K . We assume that 4.32 is verified. Then there exist $C(\alpha, M_1, M_2) > 0$ and $C_1(\alpha, M_1, M_2) > 0$ such that

$$\|q_1 - q_2\|_{L^2(K)} \leq \frac{1}{m} \frac{C(\alpha, M_1, M_2)}{\left(\ln \left(\frac{C_1(\alpha, M_1, M_2)}{\|u_1 - u_2\|_{L^\infty(0, +\infty; L^2(\Gamma_e)^2)} + \|p_1 - p_2\|_{L^\infty(0, +\infty; L^2(\Gamma_e))} + \left\| \frac{\partial p_1}{\partial n} - \frac{\partial p_2}{\partial n} \right\|_{L^\infty(0, +\infty; L^2(\Gamma_e))}} \right) \right)^{\frac{1}{2}}}.$$

Remark 4.22. *Let $l \in H_{loc}^2(0, +\infty; H^{\frac{3}{2}}(\Gamma_e))$ and $h \in H^{\frac{3}{2}}(\Gamma_e)$. Assume that there exists $\theta > 0$ such that:*

$$\sup_{t \geq 0} e^{t\theta} \left(\|l(t, \cdot)\|_{H^{\frac{3}{2}}(\Gamma_e)} + \|\partial_t l(t, \cdot)\|_{H^{\frac{3}{2}}(\Gamma_e)} \right) < +\infty,$$

Then $\kappa = h + l$ satisfies 4.32. We note that a particular case of function satisfying 4.32 is given by $l(t, x) = \omega(t)\rho(x)$ where $\omega \in H_{loc}^2(0, +\infty)$, $\rho \in H^{\frac{3}{2}}(\Gamma_e)$ and $\lim_{t \rightarrow \infty} e^{t\theta}\omega(t) = \lim_{t \rightarrow \infty} e^{t\theta}\omega'(t) = 0$.

Proof of Theorem 4.21. For $j = 1, 2$, we decompose u_j into $u_j = v_j + w_j$ where $(v_j, \zeta_j) \in H^4(\Omega)^2 \times H^3(\Omega)$ is the solution of the stationary problem:

$$\begin{cases} -\Delta v + \nabla \zeta & = 0, & \text{in } \Omega, \\ \operatorname{div} v & = 0, & \text{in } \Omega, \\ \frac{\partial v}{\partial n} - \zeta n & = hn, & \text{on } \Gamma_e, \\ \frac{\partial v}{\partial n} - \zeta n + q_j v & = 0, & \text{on } \Gamma_0, \end{cases}$$

and (w_j, π_j) is solution of the following problem:

$$\begin{cases} \partial_t w - \Delta w + \nabla \pi &= 0, & \text{in } (0, +\infty) \times \Omega, \\ \operatorname{div} w &= 0, & \text{in } (0, +\infty) \times \Omega, \\ \frac{\partial w}{\partial n} - \pi n &= (\kappa - h)n, & \text{on } (0, +\infty) \times \Gamma_e, \\ \frac{\partial w}{\partial n} - \pi n + q_j w &= 0, & \text{on } (0, +\infty) \times \Gamma_0, \\ w(0, x) &= u_0(x) - v_j(x), & \text{in } \Omega. \end{cases}$$

We would like to perform the same reasoning as in Theorem 4.18. More precisely, we are going to prove that:

$$\|w_j(t, \cdot)\|_{H^3(\Omega)^2} + \|\pi_j(t, \cdot)\|_{H^2(\Omega)} \leq G(t),$$

where G is a function which tends to 0 when t goes to $+\infty$. Since the function κ depends on t , there will be one more step than in Theorem 4.18 and that is why we assume 4.32.

We divide (w_j, π_j) into two terms: $w_j = u_j^0 + \tilde{w}_j$ and $\pi_j = p_j^0 + \tilde{\pi}_j$, where (u_j^0, p_j^0) is solution of

$$\begin{cases} \partial_t u^0 - \Delta u^0 + \nabla p^0 &= 0, & \text{in } (0, +\infty) \times \Omega, \\ \operatorname{div} u^0 &= 0, & \text{in } (0, +\infty) \times \Omega, \\ \frac{\partial u^0}{\partial n} - p^0 n &= (\kappa - h)n, & \text{on } (0, +\infty) \times \Gamma_e, \\ \frac{\partial u^0}{\partial n} - p^0 n + q_j u^0 &= 0, & \text{on } (0, +\infty) \times \Gamma_0, \\ u^0(0, x) &= 0, & \text{in } \Omega, \end{cases}$$

and $(\tilde{w}_j, \tilde{\pi}_j)$ is solution of

$$\begin{cases} \partial_t \tilde{w} - \Delta \tilde{w} + \nabla \tilde{\pi} &= 0, & \text{in } (0, +\infty) \times \Omega, \\ \operatorname{div} \tilde{w} &= 0, & \text{in } (0, +\infty) \times \Omega, \\ \frac{\partial \tilde{w}}{\partial n} - \tilde{\pi} n &= 0, & \text{on } (0, +\infty) \times \Gamma_e, \\ \frac{\partial \tilde{w}}{\partial n} - \tilde{\pi} n + q_j \tilde{w} &= 0, & \text{on } (0, +\infty) \times \Gamma_0, \\ \tilde{w}(0, x) &= u_0(x) - v_j(x), & \text{in } \Omega. \end{cases}$$

Let $t > 0$. Using the same arguments as in the previous subsection, we prove that there exists $C(\alpha, u_0, M_1, M_2) > 0$ such that:

$$\|\tilde{w}_j(t, \cdot)\|_{H^3(\Omega)^2} + \|\tilde{\pi}_j(t, \cdot)\|_{H^2(\Omega)} \leq C(\alpha, u_0, M_1, M_2) \frac{e^{-\mu t}}{t^{\frac{3}{2}}}. \quad (4.33)$$

It remains for us to bound $\|u_j^0(t, \cdot)\|_{H^3(\Omega)^2}$ and $\|p_j^0(t, \cdot)\|_{H^2(\Omega)}$. We are going to prove that there exists a constant $C(\alpha, M_2) > 0$ such that:

$$\begin{aligned} & \|u_j^0(t, \cdot)\|_{H^3(\Omega)^2} + \|p_j^0(t, \cdot)\|_{H^2(\Omega)} \\ & \leq C(\alpha, M_2) \left(\int_0^t e^{-\mu(t-s)} \|\partial_t \kappa(s, \cdot)\|_{H^{\frac{3}{2}}(\Gamma_e)}^2 ds \right)^{\frac{1}{2}} + C(\alpha, M_2) \|\partial_t \kappa(t, \cdot)\|_{H^{\frac{3}{2}}(\Gamma_e)} \\ & \quad + C(\alpha, M_2) \left(\|\kappa(t, \cdot) - h\|_{H^{\frac{3}{2}}(\Gamma_e)} + e^{-\mu t} \|\kappa(0, \cdot) - h\|_{H^{\frac{3}{2}}(\Gamma_e)} \right). \end{aligned} \quad (4.34)$$

If inequality 4.34 is satisfied, we can end the proof of Theorem 4.21:

$$\begin{aligned} \|w_1(t, \cdot) - w_2(t, \cdot)\|_{H^3(\Omega)^2} & \leq \|u_1^0(t, \cdot) - u_2^0(t, \cdot)\|_{H^3(\Omega)^2} + \|\tilde{w}_1(t, \cdot) - \tilde{w}_2(t, \cdot)\|_{H^3(\Omega)^2}, \\ \|\pi_1(t, \cdot) - \pi_2(t, \cdot)\|_{H^2(\Omega)} & \leq \|p_1^0(t, \cdot) - p_2^0(t, \cdot)\|_{H^2(\Omega)} + \|\tilde{\pi}_1(t, \cdot) - \tilde{\pi}_2(t, \cdot)\|_{H^2(\Omega)}, \end{aligned}$$

and in the following two estimates, the right hand side tends to 0 when t goes to infinity thanks to inequalities 4.33 and assumption 4.32.

We introduce (y_j, ρ_j) the solution of

$$\begin{cases} -\Delta y + \nabla \rho & = 0, & \text{in } \Omega, \\ \operatorname{div} y & = 0, & \text{in } \Omega, \\ \frac{\partial y}{\partial n} - \rho n & = (\kappa - h)n, & \text{on } \Gamma_e, \\ \frac{\partial y}{\partial n} - \rho n + q_j y & = 0, & \text{on } \Gamma_0, \end{cases}$$

for all $t > 0$. We know that $(y_j(t, \cdot), \rho_j(t, \cdot)) \in H^3(\Omega)^2 \times H^2(\Omega)$ and satisfies, thanks to Proposition 2.4:

$$\|y_j(t, \cdot)\|_{H^3(\Omega)^2} + \|\rho_j(t, \cdot)\|_{H^2(\Omega)} \leq C(\alpha, M_2) \|\kappa(t, \cdot) - h\|_{H^{\frac{3}{2}}(\Gamma_e)}. \quad (4.35)$$

Remark that $y_j(t, \cdot)$ belongs to $\mathcal{D}(A_{q_j}^{\frac{3}{2}})$. Indeed, there exists a unique $\tilde{p}(t, \cdot) \in H^3(\Omega)$ solution of

$$\begin{cases} \Delta \tilde{p} & = 0, & \text{in } \Omega, \\ \tilde{p} & = \kappa - h, & \text{on } \Gamma_e, \\ \tilde{p} & = 0, & \text{on } \Gamma_0, \end{cases} \quad (4.36)$$

for all $t > 0$ and there exists a constant $C > 0$ such that

$$\|\tilde{p}(t, \cdot)\|_{H^3(\Omega)} \leq C \|\kappa(t, \cdot) - h\|_{H^{\frac{3}{2}}(\Gamma_e)}. \quad (4.37)$$

Then $(y_j, \rho_j + \tilde{p})$ satisfies

$$\begin{cases} -\Delta y_j + \nabla(\rho_j + \tilde{p}) & = \nabla \tilde{p}, & \text{in } \Omega, \\ \operatorname{div} y_j & = 0, & \text{in } \Omega, \\ \frac{\partial y_j}{\partial n} - (\rho_j + \tilde{p})n & = 0, & \text{on } \Gamma_e, \\ \frac{\partial y_j}{\partial n} - (\rho_j + \tilde{p})n + q_j y_j & = 0, & \text{on } \Gamma_0, \end{cases}$$

for all $t > 0$. Remark that, since $\nabla \tilde{p} \in L^2(\Omega)$, we have that $y_j(t) \in \mathcal{D}(A_{q_j})$ by definition of $\mathcal{D}(A_{q_j})$. Notice that the fact that g is colinear to n is important here to do the *change of variable* in the pressure. We deduce from $A_{q_j} y_j(t) = \nabla \tilde{p}(t) \in V = \mathcal{D}(A_{q_j}^{\frac{1}{2}})$ that $y_j(t) \in \mathcal{D}(A_{q_j}^{\frac{3}{2}})$. Moreover, using Proposition 4.13 and inequality 4.37, there exists a constant $C(M_2) > 0$ such that:

$$\begin{aligned} \|A_{q_j}^{\frac{3}{2}} y_j(t, \cdot)\|_{L^2(\Omega)^2} &\leq C(M_2) \|A_{q_j} y_j(t, \cdot)\|_{H^1(\Omega)^2} = C(M_2) \|\nabla \tilde{p}(t)\|_{H^1(\Omega)^2} \\ &\leq C(M_2) \|\kappa(t, \cdot) - h\|_{H^{\frac{3}{2}}(\Gamma_e)}, \end{aligned} \quad (4.38)$$

that is to say:

$$\|y_j(t, \cdot)\|_{\mathcal{D}(A_{q_j}^{\frac{3}{2}})} \leq C(\alpha, M_2) \|\kappa(t, \cdot) - h\|_{H^{\frac{3}{2}}(\Gamma_e)}. \quad (4.39)$$

We can use the same argument, replacing $\kappa - h$ by $\partial_t \kappa$, to prove that $\partial_t y_j(t, \cdot) \in \mathcal{D}(A_{q_j}^{\frac{3}{2}})$ together with the estimate

$$\|\partial_t y_j(t, \cdot)\|_{\mathcal{D}(A_{q_j}^{\frac{3}{2}})} \leq C(\alpha, M_2) \|\partial_t \kappa(t, \cdot)\|_{H^{\frac{3}{2}}(\Gamma_e)}. \quad (4.40)$$

Let us consider $\bar{w}_j = u_j^0 - y_j$ and $\bar{p}_j = p_j^0 - \rho_j$. The couple (\bar{w}_j, \bar{p}_j) is solution of

$$\begin{cases} \partial_t \bar{w} - \Delta \bar{w} + \nabla \bar{p} & = -\partial_t y_j, & \text{in } (0, +\infty) \times \Omega, \\ \operatorname{div} \bar{w} & = 0, & \text{in } (0, +\infty) \times \Omega, \\ \frac{\partial \bar{w}}{\partial n} - \bar{p} n & = 0, & \text{on } (0, +\infty) \times \Gamma_e, \\ \frac{\partial \bar{w}}{\partial n} - \bar{p} n + q_j \bar{w} & = 0, & \text{on } (0, +\infty) \times \Gamma_0, \\ \bar{w}(0, x) & = -y_j(0, x), & \text{in } \Omega. \end{cases} \quad (4.41)$$

We know that \bar{w}_j is given by:

$$\bar{w}_j(t, \cdot) = -e^{-tA_{q_j}} y_j(0, \cdot) - \int_0^t e^{-(t-s)A_{q_j}} \partial_t y_j(s, \cdot) ds.$$

Using the family $(\phi_{q_j}^l)_{l \geq 1}$ defined by Proposition 4.11, we have: $\bar{w}_j(t, \cdot) = \sum_{l \geq 1} C_l(t) \phi_{q_j}^l$, with

$$C_l(t) = -e^{-t\lambda_{q_j}^l} (y_j(0, \cdot), \phi_{q_j}^l)_{L^2(\Omega)^2} - \int_0^t e^{-(t-s)\lambda_{q_j}^l} (\partial_t y_j(s, \cdot), \phi_{q_j}^l)_{L^2(\Omega)^2} ds.$$

Thus, recalling that $(\lambda_{q_j}^l)_{l \geq 1}$ satisfies 4.19 and using Cauchy-Schwarz inequality, there exists $C > 0$ such that:

$$C_l(t)^2 \leq 2e^{-2t\mu} (y_j(0, \cdot), \phi_{q_j}^l)_{L^2(\Omega)^2}^2 + C \int_0^t e^{-(t-s)\mu} (\partial_t y_j(s, \cdot), \phi_{q_j}^l)_{L^2(\Omega)^2}^2 ds.$$

We obtain from estimates 4.39 and 4.40:

$$\begin{aligned} \|\bar{w}_j(t, \cdot)\|_{\mathcal{D}(A_{q_j}^{\frac{3}{2}})} &\leq C(\alpha, M_2) e^{-\mu t} \|\kappa(0, \cdot) - h\|_{H^{\frac{3}{2}}(\Gamma_e)} \\ &\quad + C(\alpha, M_2) \left(\int_0^t e^{-\mu(t-s)} \|\partial_t \kappa(s, \cdot)\|_{H^{\frac{3}{2}}(\Gamma_e)}^2 ds \right)^{\frac{1}{2}}. \end{aligned} \quad (4.42)$$

Remark that, thanks to Proposition 4.15 and Proposition 4.13, we have:

$$\begin{aligned} \|\bar{w}_j(t, \cdot)\|_{H^3(\Omega)^2} &\leq C(\alpha, M_2) \|A \bar{w}_j(t, \cdot)\|_{H^1(\Omega)^2} \\ &\leq C(\alpha, M_2) \|A^{\frac{3}{2}} \bar{w}_j(t, \cdot)\|_{L^2(\Omega)^2} \leq C(\alpha, M_2) \|\bar{w}_j(t, \cdot)\|_{\mathcal{D}(A_{q_j}^{\frac{3}{2}})}. \end{aligned} \quad (4.43)$$

To summarize, using 4.43 and 4.42, we obtain the estimate:

$$\begin{aligned} \|\bar{w}_j(t, \cdot)\|_{H^3(\Omega)^2} &\leq C(\alpha, M_2) e^{-\mu t} \|\kappa(0, \cdot) - h\|_{H^{\frac{3}{2}}(\Gamma_e)} \\ &\quad + C(\alpha, M_2) \left(\int_0^t e^{-\mu(t-s)} \|\partial_t \kappa(s, \cdot)\|_{H^{\frac{3}{2}}(\Gamma_e)}^2 ds \right)^{\frac{1}{2}}. \end{aligned} \quad (4.44)$$

Using now the regularity result for the stationary problem given in Proposition 2.4, we have:

$$\|\bar{p}_j(t, \cdot)\|_{H^2(\Omega)} \leq C(\alpha, M_2) (\|\partial_t y_j(t, \cdot)\|_{H^1(\Omega)^2} + \|\partial_t \bar{w}_j(t, \cdot)\|_{H^1(\Omega)^2}).$$

Since $A_{q_j} \bar{w}_j = -\partial_t y_j - \partial_t \bar{w}_j$, we obtain:

$$\|\bar{p}_j(t, \cdot)\|_{H^2(\Omega)} \leq C(\alpha, M_2) (\|\partial_t y_j(t, \cdot)\|_{H^1(\Omega)^2} + \|A_{q_j} \bar{w}_j(t, \cdot)\|_{H^1(\Omega)^2}).$$

Thanks to Proposition 4.13, we know that

$$\|A_{q_j} \bar{w}_j(t, \cdot)\|_{H^1(\Omega)^2} \leq C(\alpha) \|A_{q_j}^{\frac{3}{2}} \bar{w}_j(t, \cdot)\|_{L^2(\Omega)^2}.$$

Therefore, using 4.40 and 4.42, we obtain:

$$\begin{aligned} \|\bar{p}_j(t, \cdot)\|_{H^2(\Omega)} &\leq C(\alpha, M_2) \left(e^{-\mu t} \|\kappa(0, \cdot) - h\|_{H^{\frac{3}{2}}(\Gamma_e)} + \|\partial_t \kappa(t, \cdot)\|_{H^{\frac{3}{2}}(\Gamma_e)} \right) \\ &\quad + C(\alpha, M_2) \left(\int_0^t e^{-\mu(t-s)} \|\partial_t \kappa(s, \cdot)\|_{H^{\frac{3}{2}}(\Gamma_e)}^2 ds \right)^{\frac{1}{2}}. \end{aligned} \quad (4.45)$$

The estimate 4.34 follows from $u_j^0 = \bar{w}_j + y_j$, $p_j^0 = \bar{w}_j + \rho_j$ and inequalities 4.35, 4.44 and 4.45. \square

4.4 Conclusion

To conclude, we have proved, under some regularity assumptions on the open set Ω and on the solution (u, p) of system 1.1, logarithmic stability estimates for the Stokes system with mixed Neumann and Robin boundary conditions. Due to the method which relies on a global Carleman inequality proved in [12], these estimates are valid in dimension 2.

Our result which, as far as we know, is the first result of this type for Stokes system, could be improved in different ways. A first concern could be to prove a logarithmic stability estimate which is valid in any dimension. This will be the subject of a forthcoming work. Next, as mentioned in Remark 4.9, Robin coefficients are estimated on a compact subset $K \subset \Gamma_0$ which is not a fixed inner portion of Γ_0 but is unknown and depends on a given reference solution. To obtain an estimate of Robin coefficients on the whole set Γ_0 or on any compact subset $K \subset \Gamma_0$ is still an open question. Finally, in our stability estimates, we need measurements on Γ_e of u , p and $\frac{\partial p}{\partial n}$, while the identifiability result given by Proposition 3.3 only requires information on u and $\frac{\partial u}{\partial n} - pn$ on Γ , where $\Gamma \subseteq \Gamma_e$ is a non-empty open subset of the boundary. Therefore, it might be interesting to know whether it is possible to obtain a stability inequality with less measurement terms.

A Existence and uniqueness for the unsteady problem

We study the regularity of the solution of the unsteady problem:

$$\begin{cases} \partial_t u - \Delta u + \nabla p &= 0, & \text{in } (0, T) \times \Omega, \\ \operatorname{div} u &= 0, & \text{in } (0, T) \times \Omega, \\ \frac{\partial u}{\partial n} - pn &= g, & \text{on } (0, T) \times \Gamma_e, \\ \frac{\partial u}{\partial n} - pn + qu &= 0, & \text{on } (0, T) \times \Gamma_0, \\ u(0, \cdot) &= u_0, & \text{in } \Omega. \end{cases}$$

where q only depends on the space variable. We are going to prove Theorem 2.6. First of all, as a preliminary result, we prove the following existence result:

Proposition A.1. *Let $T > 0$, $\alpha > 0$ and $u_0 \in H$. We assume that $g \in L^2(0, T; L^2(\Gamma_e)^d)$ and that $q \in L^\infty(\Gamma_0)$ is such that $q \geq \alpha$ on Γ_0 . There exists $u \in L^2(0, T; V)$ such that for all $v \in V$, we have in the distribution sense on $(0, T)$:*

$$\frac{d}{dt} \int_{\Omega} u \cdot v + \int_{\Omega} \nabla u : \nabla v + \int_{\Gamma_0} qu \cdot v = \int_{\Gamma_e} g \cdot v, \quad (\text{A.1})$$

and for all $v \in V$,

$$\int_{\Omega} u(0) \cdot v = \int_{\Omega} u_0 \cdot v. \quad (\text{A.2})$$

Proof of Proposition A.1. We begin by proving, using a Galerkin method, that there exists $u \in L^2(0, T; V)$ such that

$$\begin{aligned} & \forall v \in V, \forall \psi \in C^1(0, T) \text{ such that } \psi(T) = 0 \\ & - \int_0^T \int_{\Omega} u(t, x) \cdot v(x) \psi'(t) dx dt + \int_0^T \int_{\Omega} \nabla u(t, x) : \nabla v(x) \psi(t) dx dt \\ & + \int_0^T \int_{\Gamma_0} q(x) u(t, x) \cdot v(x) \psi(t) dx dt - \psi(0) \int_{\Omega} u_0(x) \cdot v(x) dx \\ & = \int_0^T \int_{\Gamma_e} g(t, x) \cdot v(x) \psi(t) dx dt. \end{aligned} \quad (\text{A.3})$$

Let $(w_i)_{i \in \mathbb{N}^*}$ be a Hilbert basis of V which is also an orthogonal basis of H . For each $n \in \mathbb{N}^*$, we define an approximate solution as follows: we search $u_n \in V_n = \operatorname{Span} \left\{ (w_i)_{1 \leq i \leq n} \right\}$ which satisfies

$$\begin{cases} \int_{\Omega} u_{n,t} \cdot w_j + \int_{\Omega} \nabla u_n : \nabla w_j + \int_{\Gamma_0} q u_n \cdot w_j = \int_{\Gamma_e} g \cdot w_j, \forall j \in \{1, \dots, n\}, \\ u_n(0) = \sum_{k=1}^n (u_0, w_k)_{L^2(\Omega)^d} w_k, \end{cases} \quad (\text{A.4})$$

where $u_{n,t}$ denotes $\partial_t u_n$.

Let $t \in (0, T)$. We decompose $u_n(t, \cdot)$ in the Hilbert basis:

$$u_n(t, \cdot) = \sum_{i=1}^n \xi_i(t) w_i.$$

We denote by

$$A = \left(\int_{\Omega} w_i(x) \cdot w_j(x) dx \right)_{1 \leq i, j \leq n}$$

$$B = \left(\int_{\Omega} \nabla w_i(x) : \nabla w_j(x) + \int_{\Gamma_0} q(x) w_i(x) \cdot w_j(x) dx \right)_{1 \leq i, j \leq n}$$

$$\xi(t) = (\xi_i(t))_{1 \leq i \leq n}$$

and

$$L(t) = \left(\int_{\Gamma_e} g(t, x) \cdot w_i(x) dx \right)_{1 \leq i \leq n}.$$

We can rewrite system A.4 in the form:

$$\begin{cases} A\xi'(t) + B\xi(t) = L(t), \\ \xi(0) = ((u_0, w_i)_{L^2(\Omega)^d})_{1 \leq i \leq n}. \end{cases}$$

Since the matrix A is invertible, the system has a unique global solution $\xi \in H^1(0, T)^d$. We are now going to prove that there exists a constant $C > 0$ independent of $n \in \mathbb{N}^*$ such that:

$$\sup_{t \in (0, T)} \int_{\Omega} |u_n|^2 + \int_0^T \int_{\Omega} |\nabla u_n|^2 + \int_0^T \int_{\Omega} |u_n|^2 \leq C. \quad (\text{A.5})$$

Multiplying the first equation of A.4 by ξ_j , summing over j for $j = 1, \dots, n$ and then integrating on $(0, t)$, we obtain:

$$\int_0^t \int_{\Omega} u_{n,t} \cdot u_n + \int_0^t \int_{\Omega} |\nabla u_n|^2 + \int_0^t \int_{\Gamma_0} q |u_n|^2 = \int_0^t \int_{\Gamma_e} g \cdot u_n \quad (\text{A.6})$$

Let $\epsilon > 0$. We have, thanks to Cauchy-Schwartz and Young inequalities:

$$\int_0^t \int_{\Gamma_e} g \cdot u_n \leq C \int_0^T \int_{\Gamma_e} |g|^2 + \epsilon \int_0^t \int_{\Gamma_e} |u_n|^2 \leq C \int_0^T \int_{\Gamma_e} |g|^2 + \epsilon \int_0^t \|u_n\|_{H^1(\Omega)^d}^2.$$

Choosing ϵ small enough and using the fact that $q \geq \alpha$ on Γ_0 , we obtain:

$$\sup_{t \in (0, T)} \int_{\Omega} |u_n|^2 + \int_0^T \int_{\Omega} |\nabla u_n|^2 + \int_0^T \int_{\Omega} |u_n|^2 \leq C \left(\int_0^T \int_{\Gamma_e} |g|^2 + \int_{\Omega} |u_0|^2 \right). \quad (\text{A.7})$$

This gives A.5. According to inequality A.5, there exists $u \in L^2(0, T; V)$ such that, up to a subsequence,

$$u_n \rightharpoonup u \text{ in } L^2(0, T; V).$$

Let $j \in \mathbb{N}^*$. Multiplying the first equation of A.4 by $\psi \in \mathcal{C}^1(0, T)$ such that $\psi(T) = 0$ then integrating on $(0, T)$, we get, $\forall n \geq j$:

$$\begin{aligned} & \int_0^T \int_{\Omega} u_{n,t}(t, x) \cdot w_j(x) \psi(t) dx dt + \int_0^T \int_{\Gamma_0} q(x) u_n(t, x) \cdot w_j(x) \psi(t) dx dt \\ & + \int_0^T \int_{\Omega} \nabla u_n(t, x) : \nabla w_j(x) \psi(t) dx dt = \int_0^T \int_{\Gamma_e} g(t, x) \cdot w_j(x) \psi(t) dx dt. \end{aligned} \quad (\text{A.8})$$

Taking into account that:

$$\begin{aligned} & \int_0^T \int_{\Omega} u_{n,t}(t, x) \cdot w_j(x) \psi(t) dx dt \\ & = - \int_0^T \int_{\Omega} u_n(t, x) \cdot w_j(x) \psi'(t) dx dt - \int_{\Omega} u_n(0, x) \cdot w_j(x) \psi(0) dx, \end{aligned}$$

we easily pass to the limit when n goes to infinity in A.8. Remark that this inequality is still valid if we replace w_j by any $v \in V$ by continuity. This ends the proof of the existence of $u \in L^2(0, T; V)$ which satisfies A.1 in the distribution sense on $(0, T)$.

Let us finish the proof of Proposition A.1 by proving that the initial condition A.2 is satisfied. Let $v \in V$. We deduce from equality A.3 that $\frac{d}{dt}(u, v)_{L^2(\Omega)^d} \in L^2(0, T)$. Consequently, the function $t \rightarrow (u(t), v)_{L^2(\Omega)^d}$ is continuous. This gives a sense to $(u(0), v)_{L^2(\Omega)^d}$. Let $\psi \in \mathcal{C}^1(0, T)$ such that $\psi(T) = 0$. Multiplying A.1 by ψ and then integrating on $(0, T)$, we obtain:

$$\begin{aligned} - \int_0^T (u, v)_{L^2(\Omega)^d} \psi'(t) dt + \int_0^T a_q(u, v) \psi(t) dt \\ = (u(0, \cdot), v)_{L^2(\Omega)^d} \psi(0) + \int_0^T l(v) \psi(t) dt, \end{aligned} \quad (\text{A.9})$$

where we recall that a_q is defined by 2.3 and with $l(v) = \int_{\Gamma_e} g \cdot v$, for $v \in V$. Comparing equality A.9 with equality A.3, we obtain $\psi(0)(u(0, \cdot) - u_0, v)_{L^2(\Omega)^d} = 0$, for all $v \in V$. By choosing ψ such that $\psi(0) \neq 0$, equality A.2 follows. \square

We are now able to prove Theorem 2.6.

Proof of Theorem 2.6. We will begin by proving that $\partial_t u \in L^2(0, T; H)$, then we will conclude by using the regularity result for the stationary problem from Proposition 2.4.

Let $t \in (0, T)$. Multiplying the first equation of A.4 by ξ_j' , summing over j for $j = 1, \dots, n$ and then integrating on $(0, t)$, we obtain:

$$\int_0^t \int_{\Omega} |u_{n,t}|^2 + \int_0^t \int_{\Gamma_0} q u_n \cdot u_{n,t} + \int_0^t \int_{\Omega} \nabla u_n : \nabla u_{n,t} = \int_0^t \int_{\Gamma_e} g \cdot u_{n,t}.$$

We have:

$$\int_0^t \int_{\Gamma_e} g \cdot u_{n,t} = - \int_0^t \int_{\Gamma_e} \partial_t g \cdot u_n - \int_{\Gamma_e} g(0) \cdot u_n(0) + \int_{\Gamma_e} g(t) \cdot u_n(t).$$

Let $\epsilon > 0$. Then, thanks to Cauchy-Schwarz and Young inequalities, there exists $C > 0$:

$$\begin{aligned} \left| \int_0^t \int_{\Gamma_e} g \cdot u_{n,t} \right| \leq \int_0^T \int_{\Gamma_e} |\partial_t g|^2 + \epsilon \int_0^T \|u_n\|_{H^1(\Omega)^d}^2 + 2 \sup_{t \in (0, T)} \int_{\Gamma_e} |g|^2 \\ + \|u_0\|_{H^1(\Omega)^d}^2 + \epsilon \sup_{t \in (0, T)} \|u_n\|_{H^1(\Omega)^d}^2. \end{aligned}$$

If we choose ϵ small enough, we finally obtain, using estimate A.7:

$$\begin{aligned} \sup_{t \in (0, T)} \|u_n\|_{H^1(\Omega)^d}^2 + \int_0^T \|u_n\|_{H^1(\Omega)^d}^2 + \int_0^T \int_{\Omega} |u_{n,t}|^2 \\ \leq C \left(\|u_0\|_{H^1(\Omega)^d}^2 + \int_0^T \int_{\Gamma_e} |\partial_t g|^2 + \int_0^T \int_{\Gamma_e} |g|^2 + \sup_{t \in (0, T)} \int_{\Gamma_e} |g|^2 \right). \end{aligned} \quad (\text{A.10})$$

We deduce that $(u_n)_{n \in \mathbb{N}^*}$ is bounded in $H^1(0, T; H) \cap L^\infty(0, T; V)$ and therefore $u \in H^1(0, T; H) \cap L^\infty(0, T; V)$.

To get regularity in space, we use the regularity result stated in Proposition 2.4 for the stationary problem. To do so, we notice that, for all $t \in (0, T)$, $(u(t), p(t))$ is solution of system 2.1 with f and g replaced by $\partial_t u(t)$ and $g(t)$. So, by Proposition 2.4 applied with $k = 0$, since $(\partial_t u, g)$ belongs to $L^2(0, T; L^2(\Omega)^d) \times L^2(0, T; H^{\frac{1}{2}}(\Gamma_e)^d)$, we deduce that $(u, p) \in L^2(0, T; H^2(\Omega)^d) \times L^2(0, T; H^1(\Omega))$.

Let us now prove the uniqueness of solution. Assume that u_1 and u_2 are two solutions and let $w = u_1 - u_2$. Then we have for all $v \in V$:

$$\int_{\Omega} \partial_t w(t) \cdot v + \int_{\Omega} \nabla w(t) : \nabla v + \int_{\Gamma_0} q w(t) \cdot v = 0, w(0) = 0. \quad (\text{A.11})$$

Taking $v = w(t)$ in A.11, we find:

$$\frac{1}{2} \frac{d}{dt} \int_{\Omega} |w(t)|^2 + \int_{\Omega} |\nabla w(t)|^2 + \int_{\Gamma_0} q |w(t)|^2 = 0,$$

that is to say

$$\int_{\Omega} |w(t)|^2 \leq \int_{\Omega} |w(0)|^2 = 0, \text{ for all } t \in (0, T).$$

So $u_1 = u_2$ on $(0, T) \times \Omega$. To conclude, thanks to system 1.1, we obtain $p_1 = p_2$. □

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