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Benchmark 3D: CeVeFE-DDFV, a discrete duality scheme with cell/vertex/face+edge unknowns

Yves Coudiere, Florence Hubert and Gianmarco Manzini `

1 Presentation of the scheme

The method that we investigate in this contribution was proposed by Y. Coudiere and F. Hubert in [1] as a three-dimensional (3D) extension of the finite volume scheme previously studied by F. Hermeline in [4] and K. Domelevo and P. Omnès in [3]. This method belongs to the family of Discrete Duality Finite Volume (DDFV) methods, which can naturally handle anisotropic or non-linear problems on general distorted meshes.

In this benchmark paper, we present the results obtained by using the formulation in [1] and the variant for discontinuous permeabilities that is presented in the proceeding paper [2].

The DDFV method that we consider herein makes use of three polyhedral meshes for the solution approximation, denoted by $\mathcal{M}, \mathcal{N}, \mathcal{F}\mathcal{E}$, and the mesh of diamonds for the solution gradient approximation, denoted by \mathscr{D} .

We denote the control volumes of the primal mesh \mathcal{M} by K and L, and with every primal cell we associate an internal point, e.g., $x_K \in K$. Different choices are possible, which give rise to different versions of the same scheme, such as the barycenters or the arithmetic average of the position vector of cell vertices (also called "isobarycenters"). For the results shown here, we used the second choice, but apparently there is no significant difference between the two choices mentioned above as far as accuracy and convergence behavior are concerned. The vertices, the edges, and the faces of mesh $\mathcal M$ are denoted by x_A , E and F, respectively. Also, we denote the

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midpoint of E by x_{E} and the barycenter of F by x_{E} . We associate a degree of freedom (the scheme unknowns) with each one of these points; hence, the unknown scalar variable takes the form:

$$
u^{\mathscr{T}} = ((u_{K})_{K \in \mathscr{M}}, (u_{A})_{A \in \mathscr{N}}, (u_{E})_{E \in \mathscr{E}}, (u_{F})_{F \in \mathscr{F}}).
$$

We denote the collections of the boundary items (vertices, edges and faces) by $\partial \mathcal{N}$, $\partial \mathcal{F} \mathcal{E}$ and we introduce the set of boundary cells $\partial \mathcal{M}$ which is composed by the boundary faces here considered as degenerated control volumes. Dirichlet boundary conditions are easily introduced into the scheme through the set of boundary data

$$
\delta u^{\mathscr{T}} = \left((u_{K})_{x_{K} \in \partial \mathscr{M}}, (u_{A})_{x_{A} \in \partial \mathscr{N}}, (u_{E})_{x_{E} \in \partial \mathscr{F} \mathscr{E}}, (u_{F})_{x_{F} \in \partial \mathscr{F} \mathscr{E}} \right).
$$

The scalar solution field *u* is approximated by the degrees of freedom $(u^{\mathcal{T}}, \delta u^{\mathcal{T}})$.

The gradient formula is given on each diamond cell $D \in \mathcal{D}$, which is the convex hull of the points K, L, x_A, x_B, x_F, x_E , by

$$
\nabla_{\delta u}^{\mathcal{D}} u^{\mathcal{T}} = \frac{1}{3|\mathcal{D}|} ((u_{\mathcal{L}} - u_{\mathcal{K}})N_{\mathcal{K}\mathcal{L}} + (u_{\mathcal{B}} - u_{\mathcal{A}})N_{\mathcal{A}\mathcal{B}} + (u_{\mathcal{F}} - u_{\mathcal{E}})N_{\mathcal{E}\mathcal{F}})
$$
(1)

using the normal vectors $N_{KL} = \frac{1}{2}(x_B - x_A) \times (x_F - x_E)$, $N_{AB} = \frac{1}{2}(x_F - x_E) \times (x_L - x_E)$ (x_K) and $N_{\text{EF}} = \frac{1}{2}(x_L - x_K) \times (x_B - x_A)$. Gradient formula (1) allows us to define the numerical flux through each interface of the control volumes of the three meshes \mathcal{M} , $\mathcal N$ and $\mathscr{F}\mathscr{E}$. Let Q be the linear space of piecewise constant vector fields defined on the mesh of diamonds $\mathscr D$ and X be the linear space of triples of piecewise constant scalar fields defined on the three meshes \mathcal{M}, \mathcal{N} and $\mathcal{F}\mathcal{E}$. Three finite volume schemes are written by using a discrete divergence operator that maps each vector field in Q to a triple of scalar functions in X . Formally, we introduce the operator

$$
\mathrm{div}^{\mathscr{T}} : \xi = (\xi_{\mathrm{D}})_{\mathrm{D} \in \mathscr{D}} \in \mathbf{Q} \mapsto (\mathrm{div}^{\mathscr{M}} \xi, \mathrm{div}^{\mathscr{N}} \xi, \mathrm{div}^{\mathscr{F} \mathscr{E}} \xi) \in X
$$

whose components

$$
\text{div}^{\mathscr{M}}\,\xi=(\text{div}_{\scriptscriptstyle{K}}\,\xi)_{\scriptscriptstyle{K}},\text{div}^{\mathscr{N}}\,\xi=(\text{div}_{\scriptscriptstyle{A}}\,\xi)_{\scriptscriptstyle{A}}\text{ and }\text{div}^{\mathscr{F}\mathscr{E}}\,\xi=\{(\text{div}_{\scriptscriptstyle{E}}\,\xi)_{\scriptscriptstyle{E}},(\text{div}_{\scriptscriptstyle{F}}\,\xi)_{\scriptscriptstyle{F}}\}
$$

are given by

$$
|\mathbf{K}| \operatorname{div}_{\mathbf{K}} \xi = \sum_{\mathbf{D} \in \mathbf{D}_{\mathbf{K}}} \xi_{\mathbf{D}} \cdot N_{\mathbf{KL}}, \quad |\mathbf{A}| \operatorname{div}_{\mathbf{A}} \xi = \sum_{\mathbf{D} \in \mathbf{D}_{\mathbf{A}}} \xi_{\mathbf{D}} \cdot N_{\mathbf{AB}}, \tag{2}
$$

$$
|E| \operatorname{div}_{E} \xi = \sum_{D \in D_{E}} \xi_{D} \cdot N_{EF}, \quad |F| \operatorname{div}_{F} \xi = \sum_{D \in D_{F}} \xi_{D} \cdot (-N_{EF}). \tag{3}
$$

In the previous statements, the symbols D_K , D_A , D_E , D_F refer to the diamond cells which overlap the cells labeled by the corresponding subscripted indices K, A, E, and L.

Since each of the div_c ξ approximates $\frac{1}{|C|}\int_C \text{div}\,\xi$ (for $C = K, A, E, F$), the right hand side of the discrete problem is given by the piecewise constant projection of Benchmark 3D: The CeVeFE-DDFV scheme 3

the function *f* onto the space X , $\pi^{\mathcal{T}} f = \{(f_K)_{K \in \mathcal{M}}, (f_A)_{A \in \mathcal{N}}, (f_E, f_F)_{E \in \mathcal{E}, F \in \mathcal{F}}\}$ with $f_{\rm C} = \frac{1}{|{\rm C}|} \int_{\rm C} f(x) dx$ for any cell ${\rm C} = {\rm K} \in \mathcal{M}$ or ${\rm A} \in \mathcal{N}$ or ${\rm F}$ or ${\rm E} \in \mathcal{F} \mathcal{E}$. The CeVeFE-DDFV scheme reads:

$$
-\operatorname{div}^{\mathcal{T}}\left(\mathbf{K}_{\mathrm{D}}\nabla_{\delta u}^{\mathrm{D}}u^{\mathcal{T}}\right)=\pi^{\mathcal{T}}f,\tag{4}
$$

where $\mathbf{K}_{\text{D}} = \frac{1}{|\text{D}|} \int_{\text{D}} \mathbf{K}(x) dx$ is a piecewise constant tensor field on the mesh of the diamond cells. The scheme in (4) originates a symmetric and positive-definite linear system of equations (see [1] for a thourough discussion of the other properties). The case of the discontinuous permeability tensor of test 5 deserves a special treatment that is thouroughly discussed in [2].

Mesure on the error

To put the discrete and the exact solutions "at the same level", we use the projection $\pi^{\mathcal{T}} u_e$ of the exact solution and the associated discrete gradient reconstruction $\nabla^{\mathcal{T}} \pi^{\mathcal{T}} u_e$. Approximation errors are evaluated through the following norms:

$$
\text{erl2} = \|e^{\mathcal{F}}\|_{L^2} / \|\pi^{\mathcal{F}} u_e\|_{L^2} \text{ with } \|e^{\mathcal{F}}\|_{L^2}^2 = \frac{1}{3} \sum_{\substack{c \in \mathcal{M} \cup \mathcal{N} \cup \mathcal{F}\mathcal{E} \\ b \in \mathcal{D}}} |c| |e_c|^2
$$
\n
$$
\text{ergrad} = \|\nabla^{\mathcal{F}} e^{\mathcal{F}}\|_{L^2} / \|\nabla^{\mathcal{F}} \pi^{\mathcal{F}} u_e\|_{L^2} \text{ with } \|\nabla^{\mathcal{F}} e^{\mathcal{F}}\|^2 = \sum_{\substack{D \in \mathcal{D} \\ D \in \mathcal{D}}} |D| |\nabla^D e^{\mathcal{F}}|^2
$$
\n
$$
\text{ener} = (\mathbf{K}^{\mathcal{D}} \nabla^{\mathcal{F}} e^{\mathcal{F}}, \nabla^{\mathcal{F}} e^{\mathcal{F}})_{L^2} / (\mathbf{K}^{\mathcal{D}} \nabla^{\mathcal{F}} \pi^{\mathcal{F}} u_e, \nabla^{\mathcal{F}} \pi^{\mathcal{F}} u_e)_{L^2}
$$
\n
$$
\text{with } (\mathbf{K}^{\mathcal{D}} \nabla^{\mathcal{F}} e^{\mathcal{F}}, \nabla^{\mathcal{F}} e^{\mathcal{F}})_{L^2} = \sum_{D \in \mathcal{D}} |D| (\mathbf{K}_D \nabla^D e^{\mathcal{F}}, \nabla^D e^{\mathcal{F}})
$$

In the case of the discontinuous tensor of test 5, the diamond cell D is divided in two subdiamond cells, namely, D_K and D_L . The gradient $\nabla^D u$ is constant on D_K (respectively, D_L) with value $\nabla^D_{\mathbf{k}} u$ (respectively, $\nabla^D_{\mathbf{k}} u$). The quantities $\|\nabla^{\mathcal{F}} e^{\mathcal{F}}\|^2_{L^2}$ and $(\mathbf{K}^{\mathscr{D}}\nabla^{\mathscr{F}}e^{\mathscr{F}}, \nabla^{\mathscr{F}}e^{\mathscr{F}})_{L^2}$ become

$$
\|\nabla^{\mathscr{T}} e^{\mathscr{T}}\|_{L^2}^2 = \sum_{\mathrm{D} \in \mathscr{D}} \left(|\mathrm{D}_{\mathrm{K}}| |\nabla_{\mathrm{D}_{\mathrm{K}}} e^{\mathscr{T}}|^2 + |\mathrm{D}_{\mathrm{L}}| |\nabla_{\mathrm{D}_{\mathrm{L}}} e^{\mathscr{T}}|^2 \right)
$$

and

$$
(\mathbf{K}^{\mathscr{D}}\nabla^{\mathscr{T}}e^{\mathscr{T}}, \nabla^{\mathscr{T}}e^{\mathscr{T}})_{L^2}=\sum_{\mathrm{D}\in\mathscr{D}}\Big(|\mathrm{D}_{\mathrm{K}}|(\mathbf{K}_{\mathrm{D}_{\mathrm{K}}}\nabla_{\mathrm{D}_{\mathrm{K}}}e^{\mathscr{T}}, \nabla_{\mathrm{D}_{\mathrm{K}}}e^{\mathscr{T}})+|\mathrm{D}_{\mathrm{L}}|(\mathbf{K}_{\mathrm{D}_{\mathrm{L}}}\nabla_{\mathrm{D}_{\mathrm{L}}}e^{\mathscr{T}}, \nabla_{\mathrm{D}_{\mathrm{L}}}e^{\mathscr{T}})\Big).
$$

2 Numerical results

The following results were obtained by using a BiCG-stab solver with ILU(0) preconditioner (routine MI26 of HSL implementation).

• Test 1 Mild anisotropy, $u(x, y, z) = 1 + \sin(\pi x) \sin(\pi (y + \frac{1}{2})) \sin(\pi (z + \frac{1}{3}))$ $min = 0, max = 2, Tetrahedral$ meshes

		$\mathbf{1}$	nu		nmat			umin			uemin							umax uemax normg		
			7777								100569 6.09E-03 1.05E-02 1.988 1.980							1.790		
		2	15495								208527 7.48E-03 9.35E-03 1.995 1.994							1.793		
		3	31139								431667 3.19E-03 5.93E-03 1.993 1.993							1.795		
		4	62419								885735 1.48E-03 2.98E-03 1.996 1.996 1.796									
			125993 1823199 1.56E-03 2.28E-03 2.000 1.999															1.797		
			6 254657 3746829 1.93E-03 2.70E-03 1.999 1.998															1.798		
\mathbf{i}		nu		erl ₂			ratiol2		ergrad			ratiograd				ener		ratioener		
1		7777	$0.228E-02$			\sim			$0.562E-01$							0.528E-01				
2			15495 0.147E-02 1.904 0.441E-01										1.051			$0.415E - 01$			1.054	
$\overline{\mathcal{E}}$			31139 0.916E-03 2.036 0.349E-01										1.011			0.327E-01			1.021	
41			62419 0.573E-03 2.025 0.276E-01											1.006		0.258E-01			1.022	
	5 125993 0.374E-03 1.819 0.219E-01											0.994				$0.206E - 01$			0.969	
	6 254657 0.231E-03 2.067 0.174E-01											0.983				$0.163E-01$			0.990	

• Test 1 Mild anisotropy, $u(x, y, z) = 1 + \sin(\pi x) \sin(\pi (y + \frac{1}{2})) \sin(\pi (z + \frac{1}{3}))$ $min = 0, max = 2, **Voronoi meshes**$

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• Test 1 Mild anisotropy, $u(x, y, z) = 1 + \sin(\pi x) \sin(\pi (y + \frac{1}{2})) \sin(\pi (z + \frac{1}{3}))$ $min = 0, max = 2$, Kershaw meshes

nu	nmat	umın	uemin umax uemax normg		
3375		49071 5.67E-02 3.43E-02 1.940 1.974 1.767			
29791		455895 9.19E-03 7.33E-03 1.988 1.991 1.782			
	3 250047 3916359 2.42E-03 1.59E-03 1.999 1.998 1.793				
	42048383 32446751 6.52E-04 6.17E-04 2.000 1.999				- 1.790

	nu		erl2 ratiol2 ergrad ratiograd ener ratioener		
			$\begin{array}{cccc} 1 & 3375 & 0.287E - 01 & - & 0.481E + 00 & - & 0.589E + 00 \end{array}$		
			2 29791 0.113E-01 1.289 0.218E+00 1.088 0.233E+00 1.277		
			$\begin{bmatrix} 3 & 250047 & 0.330E-02 & 1.730 & 0.904E-01 & 1.243 & 0.953E-01 & 1.260 \end{bmatrix}$		
			$ 4 2048383 0.859E-03 1.922 0.395E-01 1.180 0.422E-01 1.161$		

• Test 1 Mild anisotropy, $u(x, y, z) = 1 + \sin(\pi x) \sin(\pi (y + \frac{1}{2})) \sin(\pi (z + \frac{1}{3}))$ $min = 0, max = 2,$ Checkerboard meshes

• Test 2 Heterogeneous anisotropy,

 $u(x, y, z) = x^3y^2z + x\sin(2\pi xz)\sin(2\pi xy)\sin(2\pi z)$, min = −0.862, max = 1.0487, Prism meshes

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	i nu erl2 ratiol2 ergrad ratiograd ener ratioener			
	$\boxed{1}$ 12179 0.392E-01 - 0.811E-01 - 0.803E-01 -			
	$ 2 96759$ 0.109E-01 1.854 0.392E-01 1.054 0.397E-01 1.019			
	$ 3 325739$ 0.502E-02 1.917 0.256E-01 1.051 0.261E-01 1.040			
	$ 4 771119$ 0.287E-02 1.942 0.190E-01 1.039 0.194E-01 1.034			

• Test 3 Flow on random meshes, $u(x, y, z) = \sin(\pi x) \sin(\pi y) \sin(\pi z)$, $\min = 0, \max = 1,$ Random meshes

	1	nu	nmat	umin			uemin umax uemax normg		
	$\mathbf{1}$	343					4447 -4.25E+01 -9.78E-01 49.169 0.931 38.139		
	2.	3375					49855 -2.22E+01 -9.94E-01 21.970 0.982 21.514		
	31	29791					466111 -6.96E+00 -9.95E-01 7.051 0.993 12.536		
			4250047 4019647 -2.67E+00 -9.98E-01 2.725 0.998 7.541						
\mathbf{i}		nu	erl2				ratiol2 ergrad ratiograd ener	ratioener	
$\frac{1}{2}$		343	$0.147E+03 - 0.238E+02 -$				$0.162E + 01$		
2			3375 0.956E+01 3.589 0.121E+02			0.892	$0.888E + 00$	0.787	
3			29791 0.681E+00 3.640 0.632E+01				0.891 $0.459E+00$	0.909	
			4250047 0.447E-01 3.840 0.314E+01			0.988	$0.229E + 00$	0.979	

• Test 4 Flow around a well, Well meshes, $min = 0, max = 5.415$

3 Comments

This finite volume method assigns one degree of freedom to any mesh item (cells, faces, edges, and vertices). For this reason, the scheme has a large number of degrees of freedom if compared to other finite volume methods or similar discretization techniques (such as mimetic finite differences). Nonetheless, the method was proved very effective both for two and three dimensional problems with strong anisotropic coefficients and using meshes with strongly distorted cells. Among the other advantages offered by the method, we mention the coercivity of the method that eases the convergence analysis and the fact that this finite volume method generally shows second order of accuracy in all numerical experiments where the exact solution is sufficiently regular. The results shown in the tables of the previous section confirm this general behavior.

All linear systems were solved efficiently by standard preconditioned Krylov methods as BiCG-stab or GMRES. Direct solvers for general asymmetric systems (UMFPACK) can also be used, but they normally require a huge memory storage, in particular for the biggest problems. In Table 1-2, we see an example of the performance of the different solvers offered by the benchmark site when solving Test 1 on the checkerboard meshes $8 \times 8 \times 8$ and $16 \times 16 \times 16$. The comparison reveals that PETSc implementation of the CG solver is the fastest one, in particular, when combined with the diagonal preconditioner (Jacobi). A good performance in terms of CPU costs is also provided by the ISTL-BiCGstab implementation using Jacobi or ILU(0) preconditioners. CPU times are usually smaller than those obtained by using the direct solver UMFPACK, which is also available in the benchmark site.

For example, in the case of $8 \times 8 \times 8$ -size mesh we note that UMFPACK requires a CPU time of 3.180 seconds.

solver	precond	CPU time	# iters	Rel. resid.
PETSc-CG	Jacobi	0.209	202	6.368e-11
PETSc-CG	none	0.243	242	1.715e-10
ISTL-BiCGstab	ILU(0)	0.404	53	$6.832e-11$
ISTL-BiCGstab	none	0.563	167	3.656e-11
ISTL-BiCGstab	Jacobi	0.680	120	$4.415e-11$
ISTL-GMRES	ILU(0)	0.683	152	$4.241e-11$

Table 1 CeVeFe-DDFV method, test 1 using checkerboard mesh, grid resolution $8 \times 8 \times 8$; CPU times are measured in seconds.

solver	precond	CPU time	# iters	Rel. resid.
PETSc-CG	Jacobi	3.946	369	4.248e-11
PETSc-CG	none	4.989	471	8.038e-11
ISTL-CG	ILU(0)	5.540	166	$4.041e-11$
ISTL-CG	none	7.319	471	8.038e-11
ISTL-BiCGstab	ILU(0)	8.681	107	3.873e-11
ISTL-CG	Jacobi	10.989	368	4.281e-11

Table 2 CeVeFe-DDFV method, test 1 using checkerboard mesh, grid resolution $16 \times 16 \times 16$; CPU times are measured in seconds.

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References

- 1. Y. Coudiere and F. Hubert. A 3D discrete duality finite volume method for nonlinear elliptic ` equation. Preprint HAL, URL: http://hal.archives-ouvertes.fr/hal-00456837/fr, 2010.
- 2. Y. Coudiere, F. Hubert, and G. Manzini. A cevefe-ddfv scheme for discontinuous permeability ` tensors. In *Finite Volume For Complex Applications, Problems And Perspectives. 6th International Conference (this volume)*, 2011.
- 3. K. Domelevo and P. Omnes. A finite volume method for the laplace equation on almost arbitrary ` two-dimensional grids. *M2AN, Math. Model. Numer. Anal.*, 39(6):1203–1249, 2005.
- 4. F. Hermeline. Approximation of diffusion operators with discontinuous tensor coefficients on distorted meshes. *Comput. Methods Appl. Mech. Engrg.*, 192(16):1939–1959, 2003.

The paper is in final form and no similar paper has been or is being submitted elsewhere.