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► **To cite this version:**

Colin Guillarmou, Sergiu Moroianu. Chern-Simons line bundle on Teichmüller space. 2011. hal-00564717v2

**HAL Id: hal-00564717**

**<https://hal.science/hal-00564717v2>**

Preprint submitted on 6 Mar 2011 (v2), last revised 15 Jul 2011 (v3)

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# CHERN-SIMONS LINE BUNDLE ON TEICHMÜLLER SPACE

COLIN GUILLARMOU AND SERGIU MOROIANU

ABSTRACT. Let  $(X, g)$  be a convex co-compact hyperbolic 3-manifold with boundary a (not necessarily connected) closed Riemann surface  $M$  of genus  $\mathbf{g}$ . Let  $\mathcal{T}_{\mathbf{g}}$  be the Teichmüller space of the Riemann surface  $M$ . We view the deformation space of the hyperbolic metric  $g$  as a submanifold  $\mathcal{H}$  of the tangent bundle  $T\mathcal{T}_X$  of the quotient  $\mathcal{T}_X$  of Teichmüller space  $\mathcal{T}_{\mathbf{g}}$  by a subgroup of the mapping class group, with  $\mathcal{H}$  projecting diffeomorphically on  $\mathcal{T}_X$  through  $T\mathcal{T}_X \rightarrow \mathcal{T}_X$ . We show that the renormalized volume  $\text{Vol}_R$  on  $X$  is a Kähler potential for the Weil-Petersson metric, and that  $d\text{Vol}_R$  is the Liouville 1-form restricted to  $\mathcal{H}$ , implying that  $\mathcal{H}$  is Lagrangian in  $T\mathcal{T}_X$  with respect to the canonical symplectic form. This recovers and extends previous results of Takhtajan-Zograf, Takhtajan-Teo, McMullen and Krasnov-Schlenker. We introduce a holomorphic line bundle  $\mathcal{L}_{\mathbf{g}}$  over  $\mathcal{T}_X$ , with a Hermitian connection  $\nabla^{\text{SO}(3)}$ , we show that its curvature is  $\frac{i}{8\pi}$  times the Weil-Petersson form and that its lift to  $\mathcal{H}$  is flat with respect to the modified connection  $\nabla^{\text{SO}(3)} + \frac{2}{\pi}\mu^{1,0}$  where  $\mu$  is the canonical Liouville 1-form. The properties of  $\text{Vol}_R$  arise by interpreting the real and imaginary parts of the flatness of  $\mathcal{L}_{\mathbf{g}}$  along  $\mathcal{H}$ . The line bundle  $\mathcal{L}_{\mathbf{g}}$  is constructed from Chern-Simons invariants on the compactified 3-manifold  $\overline{X}$  with boundary  $M$ , using ideas of Ramadas-Singer-Weitsman. The Chern-Simons invariant of the Levi-Civita connection is an imaginary counterpart of the renormalized volume on  $\mathcal{H}$  (or  $\mathcal{T}_X$ ) in the sense that  $e^{\frac{1}{\pi}\text{Vol}_R + 2\pi i\text{CS}}$  is flat for  $\nabla^{\text{SO}(3)} + \frac{2}{\pi}\mu^{1,0}$  on  $\mathcal{H}$  and holomorphic when pulled-back on  $\mathcal{T}_X$ . Finally, in the case of Schottky uniformisation, we use a formula of Zograf to construct an explicit isomorphism of holomorphic Hermitian line bundles between  $\mathcal{L}_{\mathbf{g}}$  and a power of the determinant line bundle.

## 1. INTRODUCTION

In [7], S.S. Chern and J. Simons defined new invariants of connections on principal bundles, arising from Chern-Weil theory. These were used for instance to give necessary conditions for a Riemannian manifold of dimension  $n$  to be conformally immersed in  $\mathbb{R}^{n+k}$ . The work of Chern and Simons has been extensively developed to what is now called Chern-Simons theory, with many applications in geometry and topology, but also in theoretical physics. In particular, Witten [36] studied a 2 + 1 dimensional quantum Yang-Mills theory, with an action consisting of the Chern-Simons term of connections on a  $\text{SU}(2)$ -principal bundle  $P$  over a 3-dimensional manifold

$$\text{CS}(\omega) = \frac{1}{8\pi^2} \int_X S^* (\text{Tr}(\omega \wedge d\omega + \frac{2}{3}\omega \wedge \omega \wedge \omega)),$$

where  $S : X \rightarrow P$  is an arbitrary section. He obtained the Jones polynomial of knots on  $S^3$  as expectation values of Wilson loop functionals. On closed 3-manifolds, the invariant  $\text{CS}(\omega)$  is only defined up to integers in the sense that the values for two different sections  $S$  differ by an integer, and one usually considers  $e^{2\pi i\text{CS}(\omega)}$ , which is independent of the choice of  $S$ , as the Chern-Simons invariant. The same invariant for  $\text{SO}(3)$ -principal bundles is also defined (with a different normalizing constant) and in particular the Chern-Simons invariant of the

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*Date:* March 6, 2011.

Levi-Civita connection on the frame bundle is an invariant of the Riemannian metric when considered with values in  $\mathbb{R}/2\pi\mathbb{Z}$ . Moreover, by the Atiyah-Patodi-Singer theorem for the signature operator, the Chern-Simons invariant of the Levi-Civita connection is related to the eta invariant by the identity  $3\eta = 2\text{CS}$  modulo  $\mathbb{Z}$  (see for instance [38]).

The theory has been extended to 3-manifolds with boundary by Ramadas-Singer-Weitsman [31], in which case  $\text{CS}(\omega)$  depends on the boundary value of the section  $S$ . When considering the moduli space of flat  $\text{SU}(2)$  connections up to gauge equivalence over Riemann surfaces, the Chern-Simons invariant  $e^{2\pi i \text{CS}(\cdot)}$  of 3-manifolds bordered by the surface can be viewed as a section of a complex line bundle (with a Hermitian structure) over the moduli space and it is shown in [31] that it is isomorphic to the determinant line bundle introduced by Quillen [30]. Some more systematic studies and extensions have then been developed by Freed [9] and Kirk-Klassen [18].

Another interesting feature of Chern-Simons theory is for hyperbolic 3-manifolds  $X = \Gamma \backslash \mathbb{H}^3$ , which possess natural flat connections  $\theta$  over principal  $\text{PSL}_2(\mathbb{C})$ -bundles. For closed manifolds, one can define their Chern-Simons invariant as above by

$$\text{CS}(\theta) = -\frac{1}{16\pi^2} \int_X S^* (\text{Tr}(\theta \wedge d\theta + \frac{2}{3}\theta \wedge \theta \wedge \theta))$$

where  $S : X \rightarrow P$  are particular sections coming from the frame bundle over  $X$ . It turns out to be a complex number with imaginary part  $-\frac{1}{2\pi^2}\text{Vol}(X)$ , and real part equal to the Chern-Simons invariant of the Levi-Civita connection on the frame bundle. As we shall now recall from Yoshida's work in [38],  $\text{CS}(\theta_\tau)$  solves some  $\bar{\partial}_\tau$  type equation on certain deformation spaces of hyperbolic structures  $\tau$ , and thus the Levi-Civita (or  $\text{SO}(3)$ ) Chern-Simons invariant is the imaginary counterpart of the volume on hyperbolic manifolds. Indeed, one can not consider the set of smooth hyperbolic 3-manifold as a regular manifold by Mostow rigidity, but Yoshida [38] considered the closed cases as discrete points within a deformation space of the hyperbolic structure (in general not complete) on a 3-manifold  $N$  with  $h$  cusps. To do that, a hyperbolic Dehn surgery needs to be performed to 'fill the cusps' with solid tori, and the closed hyperbolic manifold has a particular embedded link  $\gamma$  corresponding to the filling. The deformation space  $U$  of the hyperbolic structure on  $N$  has complex dimension  $h$  and Yoshida proved that there is a holomorphic function on  $U$  such that for all  $u \in U$  for which the hyperbolic structure can be filled to a closed hyperbolic manifold  $M_u$  (with link  $\gamma_u$ )

$$(1) \quad F(u) = \exp\left(\frac{2}{\pi}\text{Vol}(M_u) + \text{length}(\gamma_u) + i(4\pi\text{CS}(M_u) + \text{torsion}(\gamma_u))\right)$$

where  $\text{CS}(M_u)$  is the Levi-Civita (or  $\text{SO}(3)$ ) Chern-Simons invariant of  $M_u$ . Therefore, up to the contribution of the link, the function  $\exp(\frac{2}{\pi}\text{Vol}(M) + 4\pi i \text{CS}(M))$  extends to a holomorphic function on a natural deformation space containing closed hyperbolic manifolds as a discrete set.

Our setting is that of 3-dimensional infinite-volume hyperbolic 3-manifolds, more particularly *convex co-compact hyperbolic* manifolds, which are conformally compactifiable to a smooth manifold with boundary. These are particular cases of quotients of  $\mathbb{H}^3$  by Kleinian groups; typical examples are quotients of  $\mathbb{H}^3$  by quasi-Fuchsian or Schottky groups. Their conformal boundary, called *conformal infinity*, is a disjoint union of Riemann surfaces and they have a finite dimensional deformation space which is essentially the deformation space of their conformal boundary. Before defining a Chern-Simons invariant, it is natural to ask about a replacement of the volume in this case. For Einstein conformally compact manifolds, the notion of *renormalized volume*  $\text{Vol}_R(X)$  has been introduced by Henningson-Skenderis [16] in

the physics literature and by Graham [11] in the mathematical literature. In the particular setting of hyperbolic 3-manifolds, this has been studied by Krasnov [19] and extended by Takhtajan-Teo [33], in relation with earlier work of Takhtajan-Zograf [34], to show that  $\text{Vol}_R$  is a Kähler potential for the Weil-Petersson metric on Teichmüller space in Schottky and quasi-Fuchsian cases. Krasnov and Schlenker [21] gave a more geometric proof of this, using the Schläfli formula on convex co-compact hyperbolic 3-manifolds to compute the variation of  $\text{Vol}_R$  in the deformation space (notice that there were works of Anderson [2] and Albin [1] on the variation of  $\text{Vol}_R$  in even dimensional Einstein settings).

Before we introduce the Chern-Simons invariant in our setting, let us first recall the definition of  $\text{Vol}_R$ . An asymptotically hyperbolic (AH in short) manifold is a Riemannian manifold  $(X, g)$  with  $X$  being the interior of a compact smooth manifold with boundary  $\overline{X}$  and such that, if  $x$  is any smooth boundary defining function, then  $\hat{g} := x^2g$  extends smoothly to a metric on  $\overline{X}$  and  $|dx|_{\hat{g}} = 1$  on  $\partial\overline{X}$ . The curvature of  $g$  tends to  $-1$  at  $\partial\overline{X}$  under these assumptions and a choice of  $x$  induces a conformal representative  $h_0 = \hat{g}|_{T\partial\overline{X}}$  of the *conformal infinity*, defined to be  $(\partial\overline{X}, \{h_0\})$  where  $\{h_0\}$  is the conformal class of  $h_0$ . Conversely, to each conformal representative  $h_0$  in  $\{h_0\}$ , there is a boundary defining function  $x$ , uniquely characterized near  $\partial\overline{X}$  by the equation  $|dx|_{x^2g} = 1$  near  $\partial\overline{X}$  and the initial condition  $x^2g|_{T\partial\overline{X}} = h_0$ , these are called *geodesic boundary defining functions*. A convex co-compact hyperbolic manifold is a particular case of AH manifold, with constant  $-1$  sectional curvatures. The renormalized volume of  $(X, g)$  is defined to be the renormalized integral of 1 as follows

$$\text{Vol}_R(X) := \text{FP}_{\epsilon \rightarrow 0} \int_{x > \epsilon} \text{dvol}_g$$

where FP means finite-part (i.e. the coefficient of  $\epsilon^0$  in the asymptotic expansion as  $\epsilon \rightarrow 0$ ) and  $x$  is a geodesic boundary defining function. Notice that this a priori depends on the choice of  $x$ . It can be shown to be independent of the choice of  $x$  in even dimension if  $g$  is Einstein (or asymptotically Einstein), see [16, 11, 1]. In odd dimension, there is always a conformal anomaly in the sense that  $\text{Vol}_R(X)$  depends on the choice of  $x$ , however in dimension 3 and when the boundary components have genus  $\geq 2$ , one can take a geodesic boundary defining function  $x$  associated to the hyperbolic conformal representative  $h_0$  on  $\partial\overline{X}$ , which fixes a natural choice. This is done by Krasnov-Schlenker [21].

We then come to Chern-Simons invariants on 3 dimensional AH manifolds. We assume that the metric  $g$  has *totally geodesic boundary* in the sense that  $L_{\nabla x} \hat{g}|_{\partial\overline{X}} = 0$  if  $L$  is Lie derivative and  $\nabla x$  is the gradient of  $x$  with respect to  $\hat{g} = x^2g$ , this condition is independent of the choice of geodesic boundary defining function  $x$  by [12]. If  $\omega$  is the  $\text{so}(3)$ -valued connection 1-form of the Levi-Civita connection of an AH metric  $g$  in an oriented orthonormal frame  $S = (S_1, S_2, S_3)$ , we define

$$(2) \quad \text{CS}(g, S) := -\frac{1}{16\pi^2} \text{FP}_{\epsilon \rightarrow 0} \int_{x > \epsilon} \text{Tr}(\omega \wedge d\omega + \frac{2}{3}\omega \wedge \omega \wedge \omega)$$

where  $x$  is a geodesic boundary-defining function associated to the constant curvature representative of the conformal infinity. For the conformal metric  $\hat{g}$  which is smooth up to the boundary, we define similarly

$$\text{CS}(\hat{g}, \hat{S}) := -\frac{1}{16\pi^2} \int_X \text{Tr}(\hat{\omega} \wedge d\hat{\omega} + \frac{2}{3}\hat{\omega} \wedge \hat{\omega} \wedge \hat{\omega})$$

if  $\hat{S} = x^{-1}S$  is an orthonormal frame for  $\hat{g}$ , and  $\hat{\omega}$  the associated Levi-Civita connection 1-form. We also define a  $\mathrm{PSL}_2(\mathbb{C})$  Chern-Simons invariant on  $(X, g)$  by

$$(3) \quad \mathrm{CS}^{\mathrm{PSL}_2(\mathbb{C})}(g, S) := -\frac{1}{16\pi^2} \mathrm{FP}_{\epsilon \rightarrow 0} \int_{x > \epsilon} \mathrm{Tr}(\theta \wedge d\theta + \frac{2}{3}\theta \wedge \theta \wedge \theta)$$

where  $\theta = \omega + iT$  with  $\omega$  as above and  $T$  the  $\mathfrak{so}(3)$ -valued 1-form defined by  $T_{ij}(V) := g(V \times S_j, S_i)$  if  $S = (S_1, S_2, S_3)$  is an oriented orthonormal frame for  $g$  and  $\times$  denotes the vector product with respect to  $g$ . In the particular case of a convex co-compact hyperbolic manifold  $X = \Gamma \backslash \mathbb{H}^3$ , there is a natural flat connection on a  $\mathrm{PSL}_2(\mathbb{C})$  principal bundle  $F^{\mathbb{C}}(X)$  over  $X$  (which can be seen as a complexified frame bundle), with  $\mathfrak{sl}_2(\mathbb{C})$ -valued connection 1-form  $\Theta$ , and we show that the definition (3) coincides with the renormalized integral of the pull-back of the Chern-Simons form  $-\frac{1}{4\pi^2} \mathrm{Tr}(\Theta \wedge d\Theta + \frac{2}{3}\Theta^3)$  of the flat connection  $\Theta$ , see Section 3. We first show

**Proposition 1.** *On an 3-dimensional AH manifold  $(X, g)$  which has totally geodesic boundary, one has  $\mathrm{CS}(g, S) = \mathrm{CS}(\hat{g}, \hat{S})$ . If moreover  $(X, g)$  is of constant curvature  $-1$ , then*

$$(4) \quad \mathrm{CS}^{\mathrm{PSL}_2(\mathbb{C})}(g, S) = -\frac{i}{2\pi^2} \mathrm{Vol}_R(X) + \frac{i}{4\pi} \chi(\partial \bar{X}) + \mathrm{CS}(g, S)$$

where  $\chi(\partial \bar{X})$  is the Euler characteristic of the boundary  $\partial \bar{X}$ .

The relation between  $\mathrm{CS}(g, S)$  and  $\mathrm{CS}(\hat{g}, \hat{S})$  comes rather easily from the conformal change formula in the Chern-Simons form (the boundary term turns out to not contribute), while (4) is a generalization of a formula in Yoshida [38], but we give an independent easy proof. Similar identities to (4) can be found in the physics literature (see for instance [20]).

Like the function  $F(u)$  of (1), it is natural to consider the variation of  $\mathrm{CS}^{\mathrm{PSL}_2(\mathbb{C})}(g, S)$  in the set of convex co-compact hyperbolic 3-manifolds, especially since, in contrast with the finite volume case, there is a finite dimensional deformation space of smooth hyperbolic 3-manifolds, which essentially coincides with the Teichmüller space of their boundaries. One of the problem, related to the work of Ramadas-Singer-Weitsman [31] is that  $e^{2\pi i \mathrm{CS}^{\mathrm{PSL}_2(\mathbb{C})}(g, S)}$  depends on the choice of frame  $S$ , since  $X$  is not closed and has boundary. This leads us to define a complex line bundle  $\mathcal{L}_{\mathbf{g}}$  over Teichmüller space  $\mathcal{T}_{\mathbf{g}}$  of Riemann surfaces of genus  $\mathbf{g}$ , for which  $e^{2\pi i \mathrm{CS}^{\mathrm{PSL}_2(\mathbb{C})}}$  and  $e^{2\pi i \mathrm{CS}}$  are sections.

Let  $\mathcal{T}_{\mathbf{g}}$  be the Teichmüller space of a, not necessarily connected, oriented Riemann surface  $M$  with genus  $\mathbf{g} = (g_1, \dots, g_N)$ ,  $g_j \geq 2$ , defined as the space of hyperbolic metrics on  $M$  modulo the group  $\mathcal{D}_0(M)$  of diffeomorphisms isotopic to the identity. We denote by  $[h]$  the class of a hyperbolic metric  $h$ . Recall that this is a complex simply connected manifold of complex dimension  $3|\mathbf{g}| - 3$  equipped with a natural Kähler metric called the Weil-Petersson metric (see Subsection 7.1). The *mapping class group*  $\mathrm{Mod}_{\mathbf{g}}$  of isotopy classes of orientation preserving diffeomorphisms of  $M$  acts properly discontinuously on  $\mathcal{T}_{\mathbf{g}}$ .

**Theorem 2.** *There exists a holomorphic Hermitian line bundle  $\mathcal{L}_{\mathbf{g}}$  over  $\mathcal{T}_{\mathbf{g}}$  equipped with a Hermitian connection  $\nabla^{\mathrm{SO}(3)}$  such that, if  $\bar{X}_{[h]} := (\bar{X}, g_{[h]})$  is a smooth family of Riemannian 3-manifolds with totally geodesic boundary  $\partial \bar{X} = M$ , parametrized by classes of constant curvature metrics  $[h] \in \mathcal{T}_{\mathbf{g}}$  so that  $[g_{[h]}|_M] = [h]$ , then  $[h] \rightarrow e^{2\pi i \mathrm{CS}(g_{[h], \cdot})}$  is a section of  $\mathcal{L}_{\mathbf{g}}$ . The curvature of  $\nabla^{\mathrm{SO}(3)}$  is  $\frac{i}{8\pi}$  times the Weil-Petersson symplectic form  $\omega_{\mathrm{WP}}$  on  $\mathcal{T}_{\mathbf{g}}$ . Moreover, elements of  $\mathrm{Mod}_{\mathbf{g}}$  which extend as diffeomorphisms of  $\bar{X}$  act on this bundle.*

The line bundle is defined using the cocycle which appears in the Chern-Simons action under gauge transformations, this is explained in Subsection 7.3. We remark that the computation of the curvature of  $\mathcal{L}_{\mathbf{g}}$  reduces to the computation of the curvature of the vertical tangent bundle in a fibration related to the universal Teichmüller curve over  $\mathcal{T}_{\mathbf{g}}$ , and we show that the fiberwise integral of the first Pontrjagin form of this bundle is given by the Weil-Petersson form, which is similar to a result of Wolpert [37]. An analogous line bundle, but in a more general setting, has been recently studied by Bunke [6].

Of particular interest for us will be the deformation space  $\mathcal{T}_X$  of a given convex co-compact hyperbolic 3-manifold  $(X, g)$ . By Theorem 3.1 of [23], this space is identified with a quotient of the Teichmüller space  $\mathcal{T}_{\mathbf{g}}$  of the conformal boundary  $\partial\bar{X}$  by the subgroup  $\text{Mod}_X$  of  $\text{Mod}_{\mathbf{g}}$  consisting of elements which extend to diffeomorphisms on  $\bar{X}$  homotopic to the identity. This subgroup acts freely, properly discontinuously on  $\mathcal{T}_{\mathbf{g}}$  and the quotient is a complex manifold of dimension  $3|\mathbf{g}| - 3$ , by [23]. Then, the bundle  $\mathcal{L}_{\mathbf{g}}$  descends with its connection on  $\mathcal{T}_X$ .

To analyze the variation of  $e^{2\pi i \text{CS}^{\text{PSL}_2(\mathbb{C})}}$ , we need to work on the tangent space  $T\mathcal{T}_{\mathbf{g}}$  of Teichmüller space. A *hyperbolic funnel* is some collar  $(0, \epsilon)_x \times M$  equipped with a metric

$$g = \frac{dx^2 + h(x)}{x^2}, \quad h(x) \in C^\infty(M, S_+^2 T^*M), \quad h(x) = h_0 \left( (\text{Id} + \frac{x^2}{2}A) \cdot, (\text{Id} + \frac{x^2}{2}A) \cdot \right)$$

where  $M$  is a Riemann surface of genus  $\geq 2$  with a hyperbolic metric  $h_0$ ,  $A$  is an endomorphism of  $TM$  satisfying  $\text{div}_{h_0} A = 0$  and  $\text{Tr}(A) = -\frac{1}{2} \text{scal}_{h_0}$ . Such a metric  $g$  is of constant sectional curvature  $-1$ , and every end of a convex co-compact hyperbolic manifold is isometric to a unique hyperbolic funnel, see [8, 21]. A couple  $(h_0, A_0)$  can be considered as an element of  $T_{h_0} \mathcal{T}_{\mathbf{g}}$ , if  $A_0 = A - \frac{1}{2} \text{tr}(A) \text{Id}$  is the trace-free part of the divergence-free tensor  $A$ . We therefore identify the tangent bundle  $T\mathcal{T}_{\mathbf{g}}$  of  $\mathcal{T}_{\mathbf{g}}$  with the set of hyperbolic funnels modulo the action of the group  $\mathcal{D}_0(M)$ , acting trivially in the  $x$  variable.

Let us still denote by  $\mathcal{L}_{\mathbf{g}}$  the Chern-Simons line bundle pulled-back to  $T\mathcal{T}_{\mathbf{g}}$  by the projection  $\pi_{\mathcal{T}} : T\mathcal{T}_{\mathbf{g}} \rightarrow \mathcal{T}_{\mathbf{g}}$ , and define a modified connection

$$(5) \quad \nabla^{\text{PSL}_2(\mathbb{C})} := \nabla^{\text{SO}(3)} + \frac{2}{\pi} \mu^{1,0}$$

on  $\mathcal{L}_{\mathbf{g}}$  over  $T\mathcal{T}_{\mathbf{g}}$ , where  $\mu^{1,0}$  is the  $(1,0)$  part of the Liouville 1-form  $\mu$  on  $T\mathcal{T}_{\mathbf{g}}$  (we identify  $T\mathcal{T}_{\mathbf{g}} \rightarrow T^*\mathcal{T}_{\mathbf{g}}$  via the Weil-Petersson metric). As before, the connection descends to  $T\mathcal{T}_X$ , and it is not Hermitian since  $\mu^{1,0}$  is not purely imaginary.

We compute the variation of  $e^{2\pi i \text{CS}^{\text{PSL}_2(\mathbb{C})}}$  with respect to  $\nabla^{\text{PSL}_2(\mathbb{C})}$  along any path of convex co-compact hyperbolic 3-manifolds:

**Theorem 3.** *Let  $(X, g^t)$  be a curve of convex co-compact hyperbolic 3-manifolds, and write  $g^t$  in funnel form near the boundary  $M = \partial\bar{X}$*

$$g^t = \frac{dx^2 + h_0^t((\text{Id} + \frac{x^2}{2}A^t) \cdot, (\text{Id} + \frac{x^2}{2}A^t) \cdot)}{x^2}$$

for some hyperbolic metric  $h_0^t$  on  $M$ . Then, denoting  $\dot{g} := \partial_t(h_0^t, A_0^t)|_{t=0} \in T(T\mathcal{T}_{\mathbf{g}})$  we have

$$\nabla_{\dot{g}}^{\text{PSL}_2(\mathbb{C})} e^{2i\pi \text{CS}^{\text{PSL}_2(\mathbb{C})}} = 0$$

In particular,  $\nabla^{\text{PSL}_2(\mathbb{C})}$  is flat on any submanifold  $\mathcal{H} \subset T\mathcal{T}_{\mathbf{g}}$  whose points correspond to hyperbolic convex co-compact 3-manifolds.

In the theorem, we notice that for  $g^t$  written in funnel form, the function  $x$  is taken independent of  $t$ , this can always be arranged by pulling back the metric by a diffeomorphism depending on  $t$ .

If  $\mathcal{T}_{\mathbf{g}}$  is Teichmüller space for a given oriented surface  $M$  of genus  $\mathbf{g}$  (possibly not connected) and  $h_0 \in \mathcal{T}_{\mathbf{g}}$ , we assume that we fix a compact 3-manifold  $\overline{X}$  with boundary  $M$  with hyperbolic convex-cocompact metric  $g_{h_0}$  whose conformal infinity is  $\{h_0\}$ . The deformation space of  $(X, g_{h_0})$  is well understood and is essentially given by  $\mathcal{T}_{\mathbf{g}}$  (or some quotient by a subgroup of the mapping class group): by [29] for a local result or [23, Th 3.1] for a global result, there exists a neighborhood  $\mathcal{U} \subset \mathcal{T}_{\mathbf{g}}$  of  $h_0$  and a smooth map  $g : \mathcal{U} \rightarrow \mathcal{C}^\infty(\overline{X}, S_+^2({}^0T^*\overline{X}))$  such that  $g_h$  is hyperbolic convex-cocompact with conformal infinity  $[h]$  for all  $h \in \mathcal{U}$ . This map induces a section  $\sigma$  in the tangent bundle of  $\mathcal{T}_{\mathbf{g}}$  defined by  $\sigma : h \rightarrow A_0^h$  if  $g_h$  is written in funnel form

$$g^h = \frac{dx^2 + h(\text{Id} + \frac{x^2}{2}A^h) \cdot, (\text{Id} + \frac{x^2}{2}A^h) \cdot}{x^2}.$$

The graph

$$(6) \quad \mathcal{H} := \{(h_0, \sigma(h_0)) \in T\mathcal{T}_{\mathbf{g}}; h_0 \in \mathcal{T}_{\mathbf{g}}\}$$

is then a smooth submanifold of  $T\mathcal{T}_{\mathbf{g}}$  of dimension  $\dim \mathcal{T}_{\mathbf{g}}$ , which can be identified with  $\mathcal{T}_{\mathbf{g}}$ . Actually, we can also mod out by the subgroup of the mapping class group corresponding to classes of diffeomorphisms on  $M$  which extend to  $\overline{X}$  to diffeomorphisms homotopic to the Identity, acting freely on  $\mathcal{T}_{\mathbf{g}}$ , see [23, Theorem 3.1], and all objects we construct descend to the quotient.

As a corollary of Theorem 3 we obtain

**Corollary 4.** *Let  $\mathcal{H} \subset T\mathcal{T}_{\mathbf{g}}$  be the submanifold defined in (6), then  $\mathcal{H}$  is Lagrangian with respect to the Liouville symplectic form  $d\mu$ . More precisely, viewing the renormalized volume  $\text{Vol}_R$  as a function on  $\mathcal{H}$ , we have  $d\text{Vol}_R = -\frac{1}{4}\mu$ .*

This fact was already known for quasi-Fuchsian manifolds by results of McMullen [28] and extended by Krasnov-Schlenker [22] to the convex co-compact setting.

Krasnov [19], extended by Takhtajan-Teo [33], combined with the results of Takhtajan-Zograf [34], on the Liouville functional, proved that the renormalized volume is a Kähler potential for the Weil-Petersson metric on the Schottky and quasi-Fuchsian deformation spaces. This was reproved by Krasnov-Schlenker [21] in the quasi-Fuchsian case. Another consequence of the flatness of  $\nabla^{\text{PSL}_2(\mathbb{C})}$  in Theorem 3 is a proof of this fact which is valid for hyperbolic 3-manifolds with funnels and cusps:

**Corollary 5.** *Consider the renormalized volume as a function on  $\mathcal{T}_{\mathbf{g}}$  by the composition  $\text{Vol}_R \circ \sigma$ , then we have on  $\mathcal{T}_{\mathbf{g}}$*

$$\bar{\partial}\partial(\text{Vol}_R \circ \sigma) = \frac{i}{16}\omega_{\text{WP}}$$

where  $\omega_{\text{WP}}$  is the Weil-Petersson symplectic form.

The same result holds on the quotient  $\mathcal{T}_X = \mathcal{T}_{\mathbf{g}}/\text{Mod}_X$ . Our final result relates the Chern-Simons line bundle  $\mathcal{L}_{\mathbf{g}}$  and the Quillen determinant line bundle  $\det \partial$  of  $\partial$  on functions in the particular case of Schottky hyperbolic manifolds. If  $M$  is a connected surface of genus  $\mathbf{g} \geq 2$ , one can realize any complex structure on  $M$  as a quotient of an open set  $\Omega_\Gamma \subset \mathbb{C}$  by a Schottky group  $\Gamma \subset \text{PSL}_2(\mathbb{C})$  and using a marking of  $\pi_1(M)$  and a certain normalization, there is complex manifold  $\mathfrak{S}_{\mathbf{g}}$ , called the Schottky space, of such groups. This corresponds

to  $\mathcal{T}_X$  when the boundary  $M$  is a connected Riemann surface. This is of complex dimension  $3\mathfrak{g} - 3$  and has Teichmüller space  $\mathcal{T}_{\mathfrak{g}}$  as universal cover. The Chern-Simons line bundle  $\mathcal{L}_{\mathfrak{S}_{\mathfrak{g}}}$  can be defined on  $\mathfrak{S}_{\mathfrak{g}}$  and since to any point  $\Gamma \in \mathfrak{S}_{\mathfrak{g}}$  there corresponds a hyperbolic 3-manifold  $\Gamma \backslash \mathbb{H}^3$  by viewing  $\Gamma \subset \mathrm{PSL}_2(\mathbb{C})$  as a group of isometries of  $\mathbb{H}^3$ , with conformal boundary given by  $\Gamma \backslash \Omega_{\Gamma} \simeq M$ , we can define a manifold  $\mathcal{H} \subset T\mathfrak{S}_{\mathfrak{g}}$  projecting diffeomorphically to the base  $\mathfrak{S}_{\mathfrak{g}}$ . Using this diffeomorphism, we can define the section  $e^{2\pi i \mathrm{CS}^{\mathrm{PSL}_2(\mathbb{C})}}$  of  $\mathcal{L}_{\mathfrak{S}_{\mathfrak{g}}}$  over  $\mathfrak{S}_{\mathfrak{g}}$ , there is a holomorphic structure on the bundle  $\mathcal{L}_{\mathfrak{S}_{\mathfrak{g}}}$  and the Hermitian connection  $\nabla^{\mathcal{L}_{\mathfrak{S}_{\mathfrak{g}}}} := \nabla^{\mathrm{SO}(3)}$ . The Quillen determinant bundle  $\det \partial$  is equipped with its Quillen metric and a natural holomorphic structure induced by  $\mathfrak{S}_{\mathfrak{g}}$  (see Subsection (9.2)), therefore inducing a Hermitian connection compatible with the holomorphic structure.

Using a formula of Zograf [39, 40], we show

**Theorem 6.** *There is an explicit isometric isomorphism of holomorphic Hermitian line bundles between the inverse  $\mathcal{L}_{\mathfrak{S}_{\mathfrak{g}}}^{-1}$  of the Chern-Simons line bundle and the 6-th power  $(\det \partial)^{\otimes 6}$  of the determinant line bundle  $\det \partial$ , given by*

$$(F\varphi)^{\otimes 6} \mapsto e^{-2\pi i \mathrm{CS}^{\mathrm{PSL}_2(\mathbb{C})}}.$$

where  $\varphi$  is the canonical section of  $\det \partial$  with respect to the Schottky marking,  $c_{\mathfrak{g}}$  is a constant on  $\mathfrak{S}_{\mathfrak{g}}$  and  $F$  is a holomorphic function on  $\mathfrak{S}_{\mathfrak{g}}$  which is given, on the open set where the product converges, by

$$F(\Gamma) = c_{\mathfrak{g}} \prod_{\{\gamma\}} \prod_{m=0}^{\infty} (1 - q_{\gamma}^{1+m}),$$

where  $q_{\gamma}$  is the multiplier of  $\gamma \in \Gamma$ ,  $\{\gamma\}$  runs over all distinct primitive conjugacy classes in  $\Gamma \in \mathfrak{S}_{\mathfrak{g}}$  excluding the identity, and  $c_{\mathfrak{g}}$  is a constant depending only on  $\mathfrak{g}$ .

We finally mention that there is a work by Krasnov [20] in the physics literature on Chern-Simons theory on asymptotically AdS spaces.

The paper splits in two parts: in Section 2–6, we introduce Chern-Simons invariants associated to the Levi-Civita connection and some complexification on asymptotically hyperbolic 3-manifolds with totally geodesic boundary and we study their relation with the renormalized volume in the case of convex co-compact hyperbolic metrics. In the second part, Section 7–9, we define the line bundle over  $\mathcal{T}_{\mathfrak{g}}$ , its connections, we compute the variation of our Chern-Simons invariants and deduce its links to the Weil-Petersson metric and the determinant line bundle.

**Extension to hyperbolic manifolds of finite geometry.** We expect the results of this paper to extend to geometrically finite hyperbolic 3-manifolds. While cusps of maximal rank are rather straightforward to treat, several technical difficulties appear when we extend our analysis to cusps of rank 1. In order not to complicate the notation we included the extension of our results to manifolds with funnels and cusps of rank 2 in the Appendix, leaving out general hyperbolic 3-manifolds of finite geometry from the present paper.

**Acknowledgements.** The subject of this paper arose from our joint paper [14] with J. Park, to whom we owe the idea of connecting the determinant and Chern-Simons line bundles in this context. We also thank S. Baseilhac, U. Bunke, K. Krasnov, G. Massuyeau and J.-M. Schlenker for useful discussions. C.G. is supported by grant ANR-09-JCJC-0099-01, S.M.



is partially supported by grant PN-II-ID-PCE 1188 265/2009. This work was done while S.M. was a visiting researcher for CNRS at the DMA of the Ecole Normale Supérieure and a visitor at IHES, we thank these institutions for their support.

## 2. ASYMPTOTICALLY HYPERBOLIC MANIFOLDS

Let  $(X, g)$  be an  $(n+1)$ -dimensional *asymptotically hyperbolic* manifold, i.e.,  $X$  is the interior of a compact smooth manifold with boundary  $\overline{X}$ , and there exists a smooth boundary-defining function  $x$  such that near the boundary  $\{x=0\}$  the Riemannian metric  $g$  has the form

$$g = \frac{dx^2 + h(x)}{x^2}$$

in a product decomposition  $[0, \epsilon)_x \times M \hookrightarrow \overline{X}$  near the boundary  $M = \partial\overline{X}$ , for some smooth one-parameter family  $h(x)$  of metrics on  $M$ . This condition on the metric is in fact equivalent to  $|dx|_{x^2g} = 1$  near  $\partial\overline{X}$  and such  $x$  are called *geodesic boundary defining functions*. When  $\partial_x h(x)|_{x=0} = 0$ , the boundary  $M$  is totally geodesic for the metric  $x^2g$  and we shall say that  $g$  has *totally geodesic boundary*. This condition is shown in [12] to be invariant with respect to the choice of  $x$  and is called evenness modulo  $O(x^3)$  in that paper. Examples of asymptotically hyperbolic manifold with totally geodesic boundary are the hyperbolic space  $\mathbb{H}^{n+1}$ , or more generally convex co-compact hyperbolic manifolds (cf. Eq. (8)). The *conformal infinity* of  $(X, g)$  is the compact manifold  $M = \partial\overline{X}$  equipped with the conformal class  $\{h_0\}$  of  $h_0 := h(0) = x^2g|_{TM}$ .

**2.1. Convex co-compact hyperbolic quotients.** Let  $X$  be an oriented complete hyperbolic 3-manifold, equipped with its constant curvature metric  $g$ . The universal cover  $\tilde{X}$  is isometric to the 3-dimensional hyperbolic space  $\mathbb{H}^3$ , and the deck transformation group is conjugated via this isometry to a Kleinian group  $\Gamma \subset \mathrm{PSL}_2(\mathbb{C})$  (we recall below that  $\mathrm{PSL}_2(\mathbb{C})$  can be viewed as the group of orientation-preserving isometries of  $\mathbb{H}^3$ ). In this way we get a representation of the fundamental group

$$(7) \quad \rho : \pi_1(X) \rightarrow \mathrm{PSL}_2(\mathbb{C})$$

with image  $\Gamma$ , well-defined up to conjugation.

Assume that  $X$  is convex co-compact, meaning that there exists a compact non-empty subset of  $X$  (called the convex core) which is geodesically convex. Then the manifold  $X$  has a smooth compactification into a manifold  $\overline{X}$ , with boundary  $M$  which is a disjoint union of Riemann surfaces. The boundary can be realized as the quotient  $\Gamma \backslash \Omega(\Gamma)$  where  $\Omega(\Gamma) \subset S^2$  is the domain of discontinuity of the convex co-compact subgroup  $\Gamma$ , acting as conformal transformations on  $S^2$ . Each connected component of  $M$  has a projective structure induced by the group  $\Gamma \subset \mathrm{PSL}_2(\mathbb{C})$ . It is proved in [8, 21] that the constant sectional curvature condition implies the following structure for the metric near infinity: there exists a product decomposition  $[0, \epsilon)_x \times M$  of  $\overline{X}$  near  $M$ , induced by the choice of a geodesic boundary-defining function  $x$  of  $M$ , a metric  $h_0$  on  $M$  and a symmetric endomorphism  $A$  of  $TM$  such that the metric  $g$  is of the form

$$(8) \quad g = \frac{dx^2 + h(x)}{x^2}, \quad h(x) = h_0 \left( \left(1 + \frac{x^2}{2}A\right) \cdot, \left(1 + \frac{x^2}{2}A\right) \cdot \right),$$

and moreover  $A$  satisfies

$$(9) \quad \text{Tr}(A) = -\frac{1}{2}\text{scal}_{h_0}, \quad d^{\nabla^*} A = 0.$$

**2.2. Tangent, cotangent and frame bundles.** There exists a smooth vector bundle over  $\overline{X}$  spanned over  $\mathcal{C}^\infty(\overline{X})$  by smooth vector fields vanishing on the boundary  $\partial\overline{X}$  (those are locally spanned near  $\partial\overline{X}$  by  $x\partial_x$  and  $x\partial_y$  if  $x$  is a boundary defining function and  $y$  are coordinates on the boundary), we denote by  ${}^0T\overline{X}$  this bundle. Its dual is denoted  ${}^0T^*\overline{X}$  and is locally spanned over  $\mathcal{C}^\infty(\overline{X})$  by the forms  $dx/x$  and  $dy/x$ . An asymptotically hyperbolic metric can be also defined to be a smooth section of the bundle of positive definite symmetric tensors  $S_+^2({}^0T^*\overline{X})$  such that  $|dx/x|_g = 1$  at  $\partial\overline{X}$ . The frame bundle  $F_0(X)$  for an asymptotically hyperbolic metric  $g$  is a  $\text{SO}(3)$ -principal bundle and its sections are smooth  $g$ -orthogonal vector fields in  ${}^0T\overline{X}$ . It is clearly canonically isomorphic to the frame bundle  $F(X)$  of the compactified metric  $\hat{g} := x^2g$  if  $x$  is a boundary defining function. A smooth frame  $S \in F_0(X)$  is said to be *even to first order* if, in local coordinates  $(y_1, y_2, x)$  near  $\partial\overline{X}$  induced by any geodesic defining function  $x$ ,  $S$  is of the form  $S = x(u_1\partial_{y_1} + u_2\partial_{y_2} + u_3\partial_x)$  where  $u_j$  are such that  $\partial_x u_j|_M = 0$ , or equivalently  $[\partial_x, \hat{S}]|_M = 0$  if  $\hat{S} := x^{-1}S$  is the related frame for  $\hat{g} := x^2g$ . In general, we refer the reader to [25, 24] for more details about the 0-structures and bundles.

**2.3. Orientation convention.** For an oriented asymptotically hyperbolic manifold the orientation of the boundary at infinity  $M$  is defined by the requirement that  $(\partial_x, Y_1, Y_2)$  is a positive frame on  $X$  if and only if  $(Y_1, Y_2)$  is a positive frame on  $M$ . With this convention, Stokes's formula gives

$$\int_{\overline{X}} d\alpha = - \int_M \alpha$$

for every  $\alpha \in \mathcal{C}^\infty(\overline{X}, \Lambda^2)$ .

**2.4. Renormalized integrals.** Let  $\omega \in x^{-N}C^\infty(\overline{X}, \Lambda^n\overline{X}) + C^\infty(\overline{X}, \Lambda^n\overline{X})$  for some  $N \in \mathbb{R}^+$ . The 0-integral (or renormalized integral) of  $\omega$  on  $X$  is defined by

$$\int_X^0 \omega := \text{FP}_{\epsilon \rightarrow 0} \int_{x>\epsilon} \omega$$

where FP denotes the finite part, i.e., the coefficient of  $\epsilon^0$  in the expansion of the integral at  $\epsilon = 0$ . This is independent of the choice of function  $x$  when  $N$  is not integer or  $N > -1$  but it depends a priori on the choice of  $x$  when  $N$  is a negative integer. More generally, one can define renormalized integrals of polyhomogeneous forms but this will not be used here. We refer the reader [1, 13] for detailed discussions on this topic. An example which has been introduced by Henningson-Skenderis [16] and Graham [11] for asymptotically hyperbolic Einstein manifolds is the renormalized volume (or 0-volume) defined by

$$\text{Vol}_R(X) := \int_X^0 \text{dvol}_g$$

where  $\text{dvol}_g$  is the volume form on  $(X, g)$ .

## 3. THE BUNDLE OF INFINITESIMAL KILLING VECTOR FIELDS FOR HYPERBOLIC MANIFOLDS

The hyperbolic 3-space  $\mathbb{H}^3$  can be viewed as a subset of quaternions

$$\mathbb{H}^3 \simeq \{y_1 + iy_2 + y_3j; y_3 > 0, y_1, y_2 \in \mathbb{R}\}, \quad g_{\mathbb{H}^3} = \frac{dy^2}{y_3^2}.$$

The action of  $\gamma = \begin{bmatrix} a & b \\ c & d \end{bmatrix} \in \mathrm{PSL}_2(\mathbb{C})$  on  $\zeta = y_1 + iy_2 + y_3j \in \mathbb{H}^3$  is given by

$$\gamma.\zeta = (a\zeta + b)(c\zeta + d)^{-1}.$$

This action identifies  $\mathrm{PSL}_2(\mathbb{C})$  with the group of oriented isometries of  $\mathbb{H}^3$ , which is diffeomorphic to the frame bundle  $F(\mathbb{H}^3) = \mathbb{H}^3 \times \mathrm{SO}(3)$  of  $\mathbb{H}^3$  via the map

$$\Phi : \mathrm{PSL}_2(\mathbb{C}) \rightarrow F(\mathbb{H}^3), \quad \gamma \mapsto (\gamma.j, \gamma_*(\partial_{y_1}, \partial_{y_2}, \partial_{y_3})).$$

There exists a natural embedding

$$(10) \quad \tilde{q} : F(\mathbb{H}^3) \rightarrow \mathbb{H}^3 \times \mathrm{PSL}_2(\mathbb{C}), \quad (m, V_m) \mapsto (m, \Phi^{-1}(m, V_m)).$$

which is equivariant with respect to the right action of  $\mathrm{SO}(3)$ . If  $X = \Gamma \backslash \mathbb{H}^3$  is an oriented hyperbolic quotient,  $\tilde{q}$  descends to a bundle map

$$(11) \quad q : F(X) \rightarrow \mathbb{H}^3 \times_{\Gamma} \mathrm{PSL}_2(\mathbb{C}) =: F^{\mathbb{C}}(X).$$

where  $F^{\mathbb{C}}(X)$  is a principal bundle over  $X$  with fiber  $\mathrm{PSL}_2(\mathbb{C})$ . The trivial flat connection on the product  $\mathbb{H}^3 \times \mathrm{PSL}_2(\mathbb{C})$  also descends to a flat connection on  $F^{\mathbb{C}}(X)$ , denoted  $\theta$  (i.e., a  $\mathfrak{sl}_2(\mathbb{C})$ -valued 1-form on  $F^{\mathbb{C}}(X)$ ), with holonomy representation conjugated to  $\rho$ , where  $\rho$  is defined in (7).

Let  $\nabla$  be the Levi-Civita connection on  $TX$  with respect to the hyperbolic metric  $g$ , and let

$$(12) \quad T \in \Lambda^1(X, \mathrm{End}(TX)), \quad T_V W := -V \times W,$$

where  $\times$  is the vector product with respect to the metric  $g$ .

**Proposition 7.** *The vector bundle  $E(X)$  associated to the principal bundle  $F^{\mathbb{C}}(X)$  with respect to the adjoint representation is isomorphic, as a complex bundle, to the complexified tangent bundle  $T_{\mathbb{C}}X$ . The connection induced by  $\theta$  is  $D := \nabla + iT$ .*

*Proof.* The associated bundle with respect to the adjoint representation is given by

$$E(X) = \mathbb{H}^3 \times_{\Gamma} \mathrm{PSL}_2(\mathbb{C}) \times_{\mathrm{PSL}_2(\mathbb{C})} \mathfrak{sl}(2, \mathbb{C}) = \mathbb{H}^3 \times_{\Gamma} \mathfrak{sl}(2, \mathbb{C}) = (\mathbb{H}^3 \times \mathfrak{sl}_2(\mathbb{C})) / \sim$$

where the equivalence relation is  $[m, h] \sim [\gamma m, \gamma h \gamma^{-1}]$  for all  $\gamma \in \Gamma$ . We also have  $TX = \Gamma \backslash T\mathbb{H}^3$  where the action of  $\mathrm{PSL}_2(\mathbb{C})$  on  $T\mathbb{H}^3$  is given by  $\gamma.(m, v_m) := (\gamma m, \gamma_*(v_m))$ . For every vector field  $u$  on  $X$  define its *canonical lift*  $s_u$  to  $T_{\mathbb{C}}X$  by

$$s_u := u + \frac{i}{2} \mathrm{curl}(u)$$

where  $\mathrm{curl}(u) = (*du^{\sharp})^{\flat}$  (the map  $u \mapsto s_u$  is a first-order differential operator). Note that the sign in front of  $\mathrm{curl}$  is different from the one used in [17]. For every  $h \in \mathfrak{sl}_2(\mathbb{C})$  let  $\kappa_h$  be the Killing field on  $\mathbb{H}^3$  corresponding to the infinitesimal isometry  $h$ .

**Lemma 8.** *We have  $\kappa_{ih} = -\frac{1}{2} \mathrm{curl}(\kappa_h)$ , thus*

$$s_{\kappa_h} = \kappa_h - i\kappa_{ih}, \quad s_{\kappa_{ih}} = i s_{\kappa_h}.$$

*Proof.* Direct verification on a basis of  $\mathfrak{sl}_2(\mathbb{C})$ , using the explicit formula for  $\kappa_h$  at  $q \in \mathbb{H}^3$ :

$$\kappa_h = b + aq + qa - qcq$$

where  $h = \begin{bmatrix} a & b \\ c & d \end{bmatrix} \in \mathfrak{sl}_2(\mathbb{C})$  and  $T_q\mathbb{H}^3$  was identified with  $\mathbb{C} \oplus j\mathbb{R}$ .  $\square$

Define a vector bundle morphism

$$\Psi : \mathcal{C}^\infty(\mathbb{H}^3, E(\mathbb{H}^3)) \rightarrow \mathcal{C}^\infty(\mathbb{H}^3, T_{\mathbb{C}}\mathbb{H}^3), \quad (m, h) \mapsto s_{\kappa_h}(m) = \kappa_h(m) + \frac{i}{2}\text{curl}(\kappa_h)(m).$$

This map is injective, for a Killing field which vanishes at a point together with its curl must vanish identically. By dimensional reasons,  $\Psi$  must be a bundle isomorphism. Moreover,  $\Psi$  is  $\text{PSL}_2(\mathbb{C})$ -equivariant in the sense that for all  $\gamma \in \text{PSL}_2(\mathbb{C})$  we have  $\Psi(\gamma m, \text{ad}_\gamma h) = \gamma_* \Psi(m, h)$  (this is clear for the real part by definition, while for the imaginary part we use the fact that  $\gamma$  is an isometry to commute it across curl), hence  $\Psi$  descends to the  $\Gamma$ -quotient as an isomorphism  $\mathcal{C}^\infty(X, E(X)) \rightarrow \mathcal{C}^\infty(X, T_{\mathbb{C}}X)$ . By Lemma 8, this isomorphism is compatible with the complex structures. It remains to identify the push-forward  $D$  of the flat connection from  $E(X)$  to  $T_{\mathbb{C}}X$  under this map. It is enough to prove that on  $\mathbb{H}^3$  we have  $D = \nabla + iT$ , since both terms are  $\text{PSL}_2(\mathbb{C})$  invariant.

**Lemma 9.** *Let  $\kappa$  be a Killing vector field on an oriented 3-manifold of constant sectional curvature  $\epsilon$ . Then the field  $v := -\frac{1}{2}\text{curl}(\kappa)$  is also Killing, and satisfies for every vector  $U$*

$$\nabla_U \kappa = U \times v, \quad \nabla_U v = \epsilon U \times \kappa.$$

*Proof.* Directly from the Koszul formula, the Levi-Civita covariant derivative of a Killing vector field  $\kappa$  satisfies  $\langle \nabla_U \kappa, V \rangle = \frac{1}{2}d\kappa^\sharp(U, V)$ , which in dimension 3 implies

$$\nabla_U \kappa = -\frac{1}{2}U \times \text{curl}(\kappa) = U \times v.$$

Let now  $(U_1, U_2, U_3)$  be a radially parallel orthonormal frame near a point  $p$ , so that  $\nabla_{U_i} U_j = 0$  at  $p$ . On one hand, by assumption on the sectional curvatures, one has  $\langle R_{U_1 U_2} \kappa, U_3 \rangle = 0$  where  $R$  is the curvature tensor of the metric. On the other hand, at the point  $p$  we have

$$\begin{aligned} \langle R_{U_1 U_2} \kappa, U_3 \rangle &= U_1 \langle \nabla_{U_2} \kappa, U_3 \rangle - U_2 \langle \nabla_{U_1} \kappa, U_3 \rangle \\ &= U_1 \langle v \times U_2, U_3 \rangle - U_2 \langle v \times U_1, U_3 \rangle \\ &= \langle \nabla_{U_1} v, U_1 \rangle + \langle \nabla_{U_2} v, U_2 \rangle. \end{aligned}$$

Similarly,  $\langle \nabla_{U_2} v, U_2 \rangle + \langle \nabla_{U_3} v, U_3 \rangle = 0$  and  $\langle \nabla_{U_3} v, U_3 \rangle + \langle \nabla_{U_1} v, U_1 \rangle = 0$  so we deduce that  $\langle \nabla_{U_j} v, U_j \rangle = 0$  at  $p$ . So  $\nabla v$  is skew-symmetric at the (arbitrary) point  $p$ , or equivalently  $v$  is Killing.

Since  $\nabla_{U_i} \kappa = U_i \times v$  we see that

$$\begin{aligned} v &= \sum \langle v, U_i \rangle U_i = \langle v, U_2 \times U_3 \rangle U_1 + \langle v, U_3 \times U_1 \rangle U_2 + \langle v, U_1 \times U_2 \rangle U_3 \\ &= \langle U_3 \times v, U_2 \rangle U_1 + \langle U_1 \times v, U_3 \rangle U_2 + \langle U_2 \times v, U_1 \rangle U_3 \\ &= \langle \nabla_{U_3} \kappa, U_2 \rangle U_1 + \langle \nabla_{U_1} \kappa, U_3 \rangle U_2 + \langle \nabla_{U_2} \kappa, U_1 \rangle U_3 \end{aligned}$$

hence at the point  $p$  where  $U_j$  are parallel and commute, using  $\langle \nabla_{U_2} \kappa, U_2 \rangle = 0$ ,

$$\langle \nabla_{U_2} v, U_1 \rangle = \langle \nabla_{U_2} \nabla_{U_3} \kappa, U_2 \rangle = \langle R_{U_2 U_3} \kappa, U_2 \rangle = \epsilon \langle \kappa, U_3 \rangle = \epsilon \langle U_2 \times \kappa, U_1 \rangle.$$

Similarly,  $\langle \nabla_{U_2} v, U_3 \rangle = \epsilon \langle U_2 \times \kappa, U_3 \rangle$ . Together with  $\langle \nabla_{U_2} v, U_2 \rangle = 0$  proved above, we deduce  $\nabla_{U_2} v = U_2 \times \kappa$ . This identity clearly holds for any  $U$  in place of  $U_2$ .  $\square$

For every  $h \in \mathfrak{sl}_2(\mathbb{C})$ , the section  $S_h : m \mapsto (m, h)$  is by definition a flat section in  $E(\mathbb{H}^3)$ , so (again by definition)  $D_U \Psi(S_h) = 0$  for every vector  $U$ . Using the above lemma, we also have  $(\nabla_U + iT(U))(\kappa_h + \frac{i}{2} \text{curl}(\kappa_h)) = 0$ . Thus the connections  $\nabla + iT$  and  $D$  have the same parallel (generating) sections, hence they coincide.  $\square$

#### 4. CHERN-SIMONS FORMS AND INVARIANTS

Let  $Z$  be a manifold,  $n \in \mathbb{N}^*$  and  $\theta \in \Lambda^1(Z, M_n(\mathbb{C}))$  a matrix-valued 1-form, and set  $\Omega := d\theta + \theta \wedge \theta$ . Define

$$\text{cs}(\theta) := \text{Tr}(\theta \wedge d\theta + \frac{2}{3}\theta^3) = \text{Tr}(\theta \wedge \Omega - \frac{1}{3}\theta^3).$$

(Notation: if  $\alpha_j$  are  $M_n(\mathbb{C})$ -valued forms of degree  $d_j$  on  $Z$ ,  $j = 1, \dots, k$ , their exterior product is defined by its action on vectors  $V_1, \dots, V_N$ ,  $N := \sum_{j=1}^k d_j$  as follows

$$\begin{aligned} & (\alpha_1 \wedge \dots \wedge \alpha_k)(V_1, \dots, V_N) \\ & := \frac{1}{d_1! \dots d_k!} \sum_{\sigma \in \Sigma(N)} \epsilon(\sigma) \alpha_1(V_{\sigma(1)}, \dots, V_{\sigma(d_1)}) \dots \alpha_k(V_{\sigma(N-d_k+1)}, \dots, V_{\sigma(N)}) \end{aligned}$$

where  $\epsilon(\sigma)$  is the sign of the permutation  $\sigma$ .

##### 4.1. Properties of Chern-Simons forms.

*Relation to Chern-Weil forms.* An easy computation shows that  $d(\text{cs}(\theta)) = \text{Tr}(\Omega \wedge \Omega)$ .

*Variation.* If  $\theta^t$  is a 1-parameter family of 1-forms and  $\dot{\theta} = \partial_t \theta^t|_{t=0}$ , the variation of cs is computed using the trace identity:

$$(13) \quad \partial_t \text{cs}(\theta^t)|_{t=0} = \text{Tr}(\dot{\theta} \wedge d\theta + \theta \wedge d\dot{\theta} + 2\dot{\theta} \wedge \theta^2) = d\text{Tr}(\dot{\theta} \wedge \theta) + 2\text{Tr}(\dot{\theta} \wedge \Omega).$$

*Pull-back.* If  $\Phi : Z' \rightarrow Z$  is a smooth map, we have  $\text{cs}(\Phi^* \theta) = \Phi^* \text{cs}(\theta)$ .

*Action of representations.* If  $\theta$  takes values in a linear Lie algebra  $\mathfrak{g} \subset M_n(\mathbb{C})$  and  $\rho$  is a representation of  $\mathfrak{g}$  in  $M_m(\mathbb{C})$  such that there exists some  $\mu_\rho \in \mathbb{C}$  with  $\text{Tr}(\rho(a)\rho(b)) = \mu_\rho \text{Tr}(ab)$  for every  $a, b \in \mathfrak{g}$ , then

$$(14) \quad \text{cs}(\rho(\theta)) = \mu_\rho \text{cs}(\theta).$$

*Gauge transformation.* If  $a : Z \rightarrow M_n(\mathbb{C})$  is a smooth map and  $\gamma := a^{-1}\theta a + a^{-1}da$ , then

$$(15) \quad \text{cs}(\gamma) = \text{cs}(\theta) + d\text{Tr}(\theta \wedge daa^{-1}) - \frac{1}{3}\text{Tr}((a^{-1}da)^3).$$

The particular cases of  $\theta$  considered below are connection 1-forms either in a principal bundle or in a trivialization of a vector bundle.

**4.2. The  $\mathrm{PSL}_2(\mathbb{C})$  invariant.** Let  $\theta$  be the flat connection in the  $\mathrm{PSL}_2(\mathbb{C})$  bundle  $F^{\mathbb{C}}(X)$  of an oriented hyperbolic 3-manifold  $X$ . Let  $S : X \rightarrow F(X)$  be a section in the orthonormal frame bundle (recall that oriented 3-manifolds are parallelizable), so  $q \circ S$  is a section in  $F^{\mathbb{C}}(X)$  where  $q$  is the natural map from  $F(X)$  to  $F^{\mathbb{C}}(X)$  defined in (11). We shall denote by  $\mathrm{cs}(\theta, S)$  the complex valued 3-form  $(q \circ S)^*\mathrm{cs}(\theta)$  on  $X$ .

Recall that for  $X$  compact, the Chern-Simons invariant of  $\theta$  with respect to  $S$  is defined by

$$\mathrm{CS}^{\mathrm{PSL}_2(\mathbb{C})}(\theta, S) := -\frac{1}{4\pi^2} \int_X \mathrm{cs}(\theta, S).$$

The normalization coefficient in front of  $\mathrm{CS}^{\mathrm{PSL}_2(\mathbb{C})}$  is so chosen because

$$-\frac{1}{4\pi^2} \int_K \mathrm{Tr}(-\frac{1}{3}(\omega_{MC})^3) = 1,$$

where  $\omega_{MC}$  is the  $\mathfrak{sl}_2(\mathbb{C})$ -valued Maurer-Cartan 1-form on  $\mathrm{PSL}_2(\mathbb{C})$ , and  $K$  is the (compact) stabilizer of  $j \in \mathbb{H}^3$ . Because  $\mathrm{PSL}_2(\mathbb{C}) = K \times \mathbb{H}^3$  is homotopy equivalent to  $K$ , this identity implies that  $\mathrm{Tr}(\omega_{MC})/12\pi^2$  is an integer cohomology class on  $\mathrm{PSL}_2(\mathbb{C})$ . Thus for  $X$  closed, (15) implies that  $\mathrm{CS}^{\mathrm{PSL}_2(\mathbb{C})}(\theta, S)$  is independent of  $S$  modulo  $\mathbb{Z}$ .

**Definition 10.** Let  $X$  be a convex co-compact hyperbolic 3-manifold and  $S$  a section in the orthonormal frame bundle which is even to first order. The Chern-Simons invariant of  $\theta$  with respect to  $S$  is defined by

$$\mathrm{CS}(\theta, S) := -\frac{1}{4\pi^2} \int_X \mathrm{cs}(\theta, S) = -\frac{1}{4\pi^2} \mathrm{FP}_{\epsilon \rightarrow 0} \int_{x>\epsilon} (q \circ S)^*\mathrm{cs}(\theta).$$

We can express  $\mathrm{cs}((q \circ S)^*\theta)$  in terms of the Riemannian connection of  $g$  as follows: let

$$h_1 := \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \quad h_2 := \begin{bmatrix} 0 & i \\ -i & 0 \end{bmatrix}, \quad h_3 := \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

be a complex basis in  $\mathfrak{sl}_2(\mathbb{C})$ . The corresponding Killing vector fields on  $\mathbb{H}^3$  evaluated at  $j$  take the values

$$\kappa_{h_k} = 2\partial_{y_k}, \quad \kappa_{ih_k} = 0$$

for  $k = 1, 2, 3$ . If  $U_k$  is the section over  $X$  in the bundle  $F^{\mathbb{C}}(X) \times_{\mathrm{ad}} \mathfrak{sl}_2(\mathbb{C})$  corresponding to the vector  $h_k$  in the trivialization  $q \circ S$ , the above relations show that the complex vector field corresponding to  $U_k$  by the isomorphism from Proposition 7 is just  $2S_k$ . Thus

$$\mathrm{ad}((q \circ S)^*\theta) = \omega + iT$$

where  $\omega$  is the  $\mathfrak{so}(3)$ -valued connection 1-form of the Levi-Civita covariant derivative  $\nabla$  in the frame  $S$ , and  $T$  denotes the  $\mathfrak{so}(3)$  valued 1-form  $T$  from Equation (12) in the basis  $S$ .

$$(16) \quad \omega_{ij}(Y) := g(\nabla_Y S_j, S_i), \quad T_{ij}(Y) := g(Y \times S_j, S_i)$$

**Lemma 11.** *The Chern-Simons form of a hyperbolic metric on a 3-manifold satisfies*

$$\mathrm{cs}((q \circ S)^*\theta) = \frac{1}{4}\mathrm{cs}(\omega + iT).$$

*Proof.* We use the following identity, valid for every  $u, v \in \mathfrak{sl}_2(\mathbb{C})$ :

$$\mathrm{Tr}^{\mathrm{M}_2(\mathbb{C})}(uv) = \frac{1}{4}\mathrm{Tr}^{\mathrm{M}_3(\mathbb{C})}(\mathrm{ad}_u \mathrm{ad}_v).$$

The lemma follows from the above discussion and Equation (14).  $\square$

**Proposition 12.** *The  $\mathrm{PSL}_2(\mathbb{C})$  Chern-Simons form on a hyperbolic 3-manifold, pulled back by a section  $S \in F_0(X)$ , has the following real and imaginary parts:*

$$\mathrm{cs}(\theta, S) = 2i\mathrm{dvol}_g + \frac{i}{4}d(\mathrm{Tr}(T \wedge \omega)) + \frac{1}{4}\mathrm{Tr}(\omega \wedge d\omega + \frac{2}{3}\omega^3).$$

*Proof.* Since the connection  $D$  is flat, it follows that  $\mathrm{cs}(\omega + iT) = -\frac{1}{3}\mathrm{Tr}((\omega + iT)^3)$ . By the above lemma we can write (using the cyclicity of the trace)

$$(17) \quad \begin{aligned} -12\mathrm{cs}(\theta, S) &= \mathrm{Tr}^{\mathrm{M}_3(\mathbb{C})}(\omega + iT)^3 \\ &= \left( \mathrm{Tr}^{\mathrm{M}_3(\mathbb{C})}(\omega^3 - 3T^2 \wedge \omega) + i\mathrm{Tr}(3\omega^2 \wedge T - T^3) \right). \end{aligned}$$

The vanishing of the  $D$ -curvature implies  $d\omega + \omega^2 - T^2 = 0$  and  $dT + T \wedge \omega + \omega \wedge T = 0$ . Taking the exterior product of the first identity with  $\omega$  and taking the trace we deduce that

$$(18) \quad \mathrm{Tr}(\omega^3 - 3T^2 \wedge \omega) = -3\mathrm{Tr}(\omega \wedge d\omega + \frac{2}{3}\omega^3)$$

Similarly, since both  $T \wedge (d\omega + \omega^2 - T^2)$  and  $\omega \wedge (dT + T \wedge \omega + \omega \wedge T)$  are 0, one can take their trace and make the difference to deduce

$$0 = \mathrm{Tr}(T \wedge d\omega - dT \wedge \omega - T^3 - \omega^2 \wedge T) = -d(\mathrm{Tr}(T \wedge \omega)) - \mathrm{Tr}(T^3) - \mathrm{Tr}(\omega^2 \wedge T)$$

and then

$$(19) \quad \mathrm{Tr}(3\omega^2 \wedge T - T^3) = -4\mathrm{Tr}(T^3) - 3d(\mathrm{Tr}(T \wedge \omega)).$$

We also easily see that  $\mathrm{Tr}(T^3) = 6\mathrm{dvol}_g$  and therefore combining (19), (18) and (17), we have proved the Proposition.  $\square$

**Definition 13.** Let  $(X, g)$  be an oriented asymptotically hyperbolic 3 manifold with totally geodesic boundary. Let  $S$  be an even to first order orthonormal frame on  $X$  and  $\omega + iT$  the  $\mathfrak{so}(3) \otimes \mathbb{C}$ -valued 1-form defined in (16). The  $\mathrm{PSL}_2(\mathbb{C})$  Chern-Simons form of  $g$  in the trivialization  $S$  is  $\mathrm{cs}(\omega + iT)$ . The  $\mathrm{PSL}_2(\mathbb{C})$  Chern-Simons invariant is

$$\mathrm{CS}^{\mathrm{PSL}_2(\mathbb{C})}(g, S) := -\frac{1}{16\pi^2} \int_X^0 \mathrm{cs}(\omega + iT).$$

This invariant is our main object of study in the present paper.

**4.3. The  $\mathrm{SO}(3)$  Chern-Simons invariant.** Let  $(X, g)$  be an oriented Riemannian manifold of dimension 3. Denote by  $\mathrm{cs}(g)$  the Chern-Simons form of the Levi-Civita connection 1-form of  $g$  on the orthonormal frame bundle  $F(X)$ . Note that  $S^*\mathrm{cs}(g) = \mathrm{cs}(\omega)$  where  $\omega$  is the connection 1-form in the trivialization  $S$ .

If  $X$  is compact, for every section  $S : X \rightarrow F(X)$  in the orthonormal frame bundle we define the Chern-Simons invariant

$$\mathrm{CS}(g, S) := -\frac{1}{16\pi^2} \int_X S^* \mathrm{cs}(g).$$

When  $a : X \rightarrow \mathrm{SO}(3)$  is a compactly supported map (i.e.,  $a \cong 1$  outside a compact), we have

$$(20) \quad \frac{1}{16\pi^2} \int_X \mathrm{Tr}(-\frac{1}{3}(a^{-1}da)^3) = -\mathrm{deg}(a) \in \mathbb{Z},$$

so for  $X$  is closed, using (15) we see that the Chern-Simons invariant is independent of  $S$  modulo  $\mathbb{Z}$ .

We aim to define a  $\mathrm{SO}(3)$  Chern-Simons invariant associated to the Levi-Civita connection  $\nabla^g$  on an asymptotically hyperbolic 3-manifold  $(X, g)$  with totally geodesic boundary. First, we fix a geodesic boundary defining function  $x$  and we set

$$\hat{g} := x^2 g.$$

Let  $\hat{S} : X \rightarrow F(X, \hat{g})$  be a smooth section of the orthonormal frame bundle associated to the metric  $\hat{g}$ . We say that  $\hat{S}$  is *even to first order* if  $\mathcal{L}_{\partial_x} \hat{S}|_M = 0$  where  $\mathcal{L}$  denotes the Lie derivative. We define, starting from  $\hat{S}$ , a section  $S := x\hat{S}$  in the frame bundle  $F_0(X)$  associated to  $g$ .

**Definition 14.** The  $\mathrm{SO}(3)$  Chern-Simons invariant of an asymptotically hyperbolic metric  $g$  with totally geodesic boundary, with respect to an even to first order trivialization  $S$  of  $F_0(X)$ , is

$$\mathrm{CS}(g, S) := -\frac{1}{16\pi^2} \int_X S^* \mathrm{cs}(g).$$

## 5. COMPARISON OF THE ASYMPTOTICALLY HYPERBOLIC AND COMPACT CHERN-SIMONS $\mathrm{SO}(3)$ INVARIANTS

For any pair of conformal metrics  $\hat{g} = x^2 g$  we can relate  $\mathrm{cs}(g, S)$  to  $\mathrm{cs}(\hat{g}, \hat{S})$  as follows. We denote by  $\omega, \hat{\omega}$  the connection 1-forms of  $g, \hat{g}$  in the trivialization  $S, \hat{S}$ , respectively  $\hat{S} = x^{-1}S$ . For every  $Y \in TX$  that means

$$\hat{\omega}_{ij}(Y) = \hat{g}(\nabla_Y^{\hat{g}} \hat{S}_j, \hat{S}_i), \quad \omega_{ij}(Y) = g(\nabla_Y^g S_j, S_i).$$

**Lemma 15.** *The connection forms of the conformal metrics  $g$  and  $\hat{g} = x^2 g$  satisfy  $\hat{\omega} = \omega + \alpha$ , where*

$$(21) \quad \alpha_{ij}(Y) := \hat{g}(Y, \hat{S}_i) \hat{S}_j(a) - \hat{g}(Y, \hat{S}_j) \hat{S}_i(a) = g(Y, S_i) S_j(a) - g(Y, S_j) S_i(a)$$

with  $a := \log(x)$ .

*Proof.* An easy computation using Koszul's formula for the Riemannian connection in a frame.  $\square$

Let  $g^t = x^{2t} g$  where  $t \in [0, 1]$ , then  $S^t := x^{-t} S$  defines a section of the frame bundle  $F^t(X)$  of  $g^t$ . Consider  $\omega^t$  the connection form of  $g^t$  in the basis  $S^t$ , and write  $\alpha^t = \omega^t - \omega$ . Notice from (21) that  $\alpha^t = t\alpha$  is linear in  $t$ , so we compute the variation of  $\mathrm{cs}(\omega^t)$  using (13):

$$\partial_t \mathrm{cs}(\omega^t)|_{t=0} = d\mathrm{Tr}(\alpha \wedge \omega) + 2\mathrm{Tr}(\alpha \wedge \Omega)$$

where  $\Omega = d\omega + \omega \wedge \omega$  is the curvature of  $\omega$ .

**Lemma 16.** *We have  $\mathrm{Tr}(\alpha \wedge \Omega) \equiv 0$ .*

*Proof.* At points where  $\nabla a = 0$  this is clear. At other points, take an orthonormal basis  $(X_1, X_2, X_3)$  of  $TX$  for  $g$  such that  $X_3$  is proportional to  $\mathrm{grad}(a)$ . Since  $\alpha \wedge \Omega$  is a tensor, we can compute its trace in the basis  $X_j$  instead of  $S_j$ :

$$\begin{aligned} \mathrm{Tr}(\alpha \wedge \Omega)(X_1, X_2, X_3) &= \sum_{i,j} \alpha_{ij}(X_1) \Omega_{ji}(X_2, X_3) - \alpha_{ij}(X_2) \Omega_{ji}(X_1, X_3) \\ &= 2(\langle R_{X_2 X_3} X_1, \mathrm{grad}(a) \rangle - \langle R_{X_1 X_3} X_2, \mathrm{grad}(a) \rangle) \end{aligned}$$



and this vanishes using the symmetry of the Riemannian curvature together with the fact that  $X_3$  and  $\text{grad}(a)$  are collinear.  $\square$

We deduce that  $\partial_t \text{cs}(\omega^t)|_{t=0} = d\text{Tr}(\alpha \wedge \omega)$  and so

$$\partial_t \text{cs}(\omega^t) = \partial_s \text{cs}(\omega^{t+s})|_{s=0} = d(\alpha^t \wedge \omega^t) = d\text{Tr}(\alpha \wedge (\omega + t\alpha)).$$

Since  $\text{Tr}(\alpha \wedge \alpha) = 0$  by cyclicity of the trace, we find

$$(22) \quad \text{cs}(\hat{g}, \hat{S}) = \text{cs}(g, S) + d\text{Tr}(\alpha \wedge \omega).$$

**Proposition 17.** *Let  $g$  be an asymptotically hyperbolic metric on  $X$  with totally geodesic boundary, let  $x$  be a smooth geodesic boundary defining function and set  $\hat{g} := x^2 g$ . Let  $\hat{S}$  be an even to first order section in  $F(X)$  with respect to  $\hat{g}$ , and let  $S = x\hat{S}$  be the corresponding section in  $F_0(X)$ . Then the  $\text{SO}(3)$  Chern-Simons invariants of  $g$  and  $\hat{g}$  with respect to  $S, \hat{S}$  coincide:*

$$\text{CS}(g, S) = \text{CS}(\hat{g}, \hat{S}).$$

*Proof.* By integration on  $X$  and using Stokes, we get

$$(23) \quad 16\pi^2 \text{CS}(\hat{g}, \hat{S}) = 16\pi^2 \text{CS}(g, S) - \text{FP}_{\epsilon \rightarrow 0} \int_{x=\epsilon} \text{Tr}(\alpha \wedge \omega).$$

The proof is finished by showing that the trace  $\text{Tr}(\alpha \wedge \omega)$  is odd in  $x$  to order  $O(x^4)$ , so

$$\text{FP}_{\epsilon \rightarrow 0} \int_{x=\epsilon} \text{Tr}(\alpha \wedge \omega) = 0.$$

For this, note that  $x\alpha$  is smooth in  $x$  and has an even expansion at  $x = 0$  in powers of  $x$  up to  $O(x^3)$  by assumption on the section  $S$ , while  $x\omega$  is smooth in  $x$  but a priori not even. Setting  $a := \log x$  we write for  $Y_1, Y_2$  vector fields on  $\partial\bar{X}$  (thus orthogonal to  $\nabla a$ ) and  $\langle \cdot, \cdot \rangle := g(\cdot, \cdot)$ :

$$\begin{aligned} & \text{Tr}(\alpha \wedge \omega)(Y_1, Y_2) \\ &= \sum_{1 \leq i, j \leq 3} \alpha_{ij}(Y_1) \omega_{ji}(Y_2) - \alpha_{ij}(Y_2) \omega_{ji}(Y_1) \\ &= \sum_{1 \leq j \leq 3} S_j(a) \langle \nabla_{Y_2} S_j, Y_1 \rangle - S_j(a) \langle \nabla_{Y_1} S_j, Y_2 \rangle - \langle S_j, Y_1 \rangle \langle \nabla_{Y_2} S_j, \nabla a \rangle + \langle S_j, Y_2 \rangle \langle \nabla_{Y_1} S_j, \nabla a \rangle \\ &= 2 \langle \nabla_{Y_2} \nabla a, Y_1 \rangle - 2 \langle \nabla_{Y_1} \nabla a, Y_2 \rangle - 2 \sum_{1 \leq j \leq 3} Y_2(S_j(a)) \langle S_j, Y_1 \rangle - Y_1(S_j(a)) \langle S_j, Y_2 \rangle \\ &= \sum_{1 \leq j \leq 3} -2Y_2(S_j(a)) \langle S_j, Y_1 \rangle + 2Y_1(S_j(a)) \langle S_j, Y_2 \rangle \end{aligned}$$

and this is odd in  $x$  to order  $O(x^2)$  since  $da = -\frac{dx}{x}$ ,  $\hat{S}_j = x^{-1}S_j$  is even to first order, and the metric has totally geodesic boundary (i.e.  $x^2g$  is even to order  $O(x^3)$ ).  $\square$

## 6. COMPARISON OF THE $\text{PSL}_2(\mathbb{C})$ AND $\text{SO}(3)$ INVARIANTS IN THE HYPERBOLIC SETTING

In this section we establish the relation between the  $\text{PSL}_2(\mathbb{C})$  and the  $\text{SO}(3)$  Chern-Simons invariants. This was known in the compact case and in the finite volume case since the work of Yoshida [38].

**Proposition 18.** *Let  $(X, g) = \Gamma \backslash \mathbb{H}^3$  be a convex co-compact hyperbolic 3-manifold with  $\Gamma \subset \mathrm{PSL}_2(\mathbb{C})$  and let  $\theta$  be the associated flat connection on the bundle  $F^{\mathbb{C}}(X) = \mathbb{H}^3 \times_{\Gamma} \mathrm{PSL}_2(\mathbb{C})$ . Let  $S : X \rightarrow F_0(X)$  be an even section of  $F_0(X)$ . Then*

$$\mathrm{CS}^{\mathrm{PSL}_2(\mathbb{C})}(\theta, S) = -\frac{i}{2\pi^2} \mathrm{Vol}_R(X) + \frac{i}{2\pi} \chi(M) + \mathrm{CS}(g, S).$$

*Proof.* Using Proposition 12 and Stokes's formula, we have

$$\begin{aligned} \mathrm{CS}^{\mathrm{PSL}_2(\mathbb{C})}(\theta, S) &= \mathrm{FP}_{\epsilon \rightarrow 0} \int_{\{x > \epsilon\}} -\frac{i}{2\pi^2} \mathrm{dvol}_{\mathbb{H}^3} + \frac{i}{16\pi^2} d(\mathrm{Tr}(\omega \wedge T)) - \frac{1}{16\pi^2} \mathrm{cs}(\omega) \\ &= -\frac{i}{2\pi^2} \mathrm{Vol}_R(X) - \mathrm{FP}_{\epsilon \rightarrow 0} \frac{i}{16\pi^2} \int_{x=\epsilon} \mathrm{Tr}(\omega \wedge T) + \mathrm{CS}(g, S). \end{aligned}$$

The conclusion follows from Lemma 19.  $\square$

**Lemma 19.** *We have*

$$(24) \quad \mathrm{FP}_{\epsilon \rightarrow 0} \int_{x=\epsilon} \mathrm{Tr}(T \wedge \omega) = 2 \int_M \mathrm{scal}_{h_0} \mathrm{dvol}_{h_0} = 8\pi \chi(M).$$

*Proof.* Let  $U_j := x^{-1} S_j$  denote the orthonormal frame for the compact metric  $\hat{g} = x^2 g$ ,  $(S^1, S^2, S^3)$  the dual basis to  $S$ ,  $\hat{\omega}_{ij}(Y) := \hat{g}(\nabla_Y^{\hat{g}} U_i, U_j)$  the Levi-Civita connection 1-form of  $\hat{g}$  in the frame  $U$ , and  $U^j = x S^j$  the dual co-frame. Let  $Y_1, Y_2$  be a local orthonormal frame on  $M$  for  $h_0$  of eigenvectors for the map  $A$  defined on  $TM$  by (8), extended on  $X$  constantly in  $x$  near  $M$ .

We split  $\omega = \hat{\omega} - \alpha$  and we first compute

$$\begin{aligned} \mathrm{Tr}(T \wedge \alpha)(Y_1, Y_2) &= \sum_{i,j} \langle Y_1 \times S_i, S_j \rangle (\langle Y_2, S_j \rangle S_i(a) - \langle Y_2, S_i \rangle S_j(a)) \\ &\quad - \langle Y_2 \times S_i, S_j \rangle (\langle Y_1, S_j \rangle S_i(a) - \langle Y_1, S_i \rangle S_j(a)) \\ &= -4 \langle Y_1 \times Y_2, \nabla a \rangle. \end{aligned}$$

Define  $\tilde{Y}_j := (1 + \frac{1}{2} x^2 A)^{-1} Y_j$ . Then  $x \tilde{Y}_1, x \tilde{Y}_2, \nabla a$  form an orthonormal frame near  $x = 0$  and  $x \tilde{Y}_1 \times x \tilde{Y}_2 = \nabla a$ . Thus,  $x^2 (1 + \frac{x^2 \lambda_1}{2})^{-1} (1 + \frac{x^2 \lambda_2}{2})^{-1} \langle Y_1 \times Y_2, \nabla a \rangle = 1$  which shows that

$$\mathrm{FP}_{\epsilon \rightarrow 0} \mathrm{Tr}(T \wedge \alpha) = 2 \mathrm{tr}(A) \mathrm{dvol}_{h_0}.$$

Let us now compute the form  $\mathrm{Tr}(\hat{\omega} \wedge T)$  on the hypersurface  $x = \epsilon$ . Notice that

$$T_{12} = S^3, \quad T_{23} = S^1, \quad T_{31} = S^2$$

so  $xT$  is smooth in  $x$ , and we easily see that

$$\frac{1}{2} \mathrm{Tr}(\hat{\omega} \wedge T) = S^1 \wedge \hat{\omega}_{23} + S^2 \wedge \hat{\omega}_{31} + S^3 \wedge \hat{\omega}_{12}.$$

Using Koszul formula and the evenness of  $g$  and  $S$ , for a vector  $Y \in TM$  independent of  $x$  the term  $\hat{\omega}_{ji}$  can be decomposed under the form

$$\begin{aligned} 2\hat{\omega}_{ji}(Y) &= U_i(\hat{g}(Y, U_j)) - U_j(\hat{g}(Y, U_i)) - \hat{g}([U_i, U_j], Y) + \text{even function of } x \\ &= dY^{\#}(U_i, U_j) + \text{even function of } x \end{aligned}$$

so the odd component is tensorial in  $U_j$ . Therefore we can compute  $\text{FP}_{\epsilon \rightarrow 0} \text{Tr}(\hat{\omega} \wedge T)$  using the orthonormal frame  $\tilde{Y}_1, \tilde{Y}_2, \tilde{Y}_3 := \partial_x$ :

$$\begin{aligned} \text{Tr}(T \wedge \hat{\omega})(Y_1, Y_2) &= x^{-1} \sum_{i,j} \langle Y_1 \times \tilde{Y}_i, \tilde{Y}_j \rangle \langle \nabla_{Y_2} \tilde{Y}_i, \tilde{Y}_j \rangle - \langle Y_2 \times \tilde{Y}_i, \tilde{Y}_j \rangle \langle \nabla_{Y_1} \tilde{Y}_i, \tilde{Y}_j \rangle \\ &= x^{-1} \sum_i \langle \nabla_{Y_2} \tilde{Y}_i, Y_1 \times \tilde{Y}_i \rangle - \langle \nabla_{Y_1} \tilde{Y}_i, Y_2 \times \tilde{Y}_i \rangle. \end{aligned}$$

(here the vector product is with respect to  $\hat{g}$ ). Since  $\tilde{Y}_j - Y_j$  is of order  $x^2$ , the finite part is unchanged if we replace  $Y_1, Y_2$  by  $\tilde{Y}_1, \tilde{Y}_2$  in the above, thus getting

$$x^{-1} \sum_i \langle \nabla_{\tilde{Y}_2} \tilde{Y}_i, \tilde{Y}_1 \times \tilde{Y}_i \rangle - \langle \nabla_{\tilde{Y}_1} \tilde{Y}_i, \tilde{Y}_2 \times \tilde{Y}_i \rangle.$$

For  $k = 1, 2$  the coefficient of  $x$  in  $\langle \nabla_{\tilde{Y}_k} \tilde{Y}_k, \partial_x \rangle = -\langle \nabla_{\tilde{Y}_k} \partial_x, \tilde{Y}_k \rangle$  is  $-\lambda_k$ . We therefore get

$$\text{FP}_{\epsilon \rightarrow 0} \text{Tr}(T \wedge \hat{\omega}) = -2\text{tr}(A) \text{dvol}_{h_0}.$$

Together with the identity  $2\text{tr}(A) = -\text{scal}_{h_0}$  and Gauss-Bonnet this ends the proof.  $\square$

## 7. THE CHERN-SIMONS LINE BUNDLE AND ITS CONNECTION

**7.1. The tangent space of Teichmüller space as the set of hyperbolic funnels.** In this subsection, we shall see that the tangent space  $T\mathcal{T}_{\mathbf{g}}$  of Teichmüller space of Riemann surfaces of genus  $\mathbf{g}$  can be identified with ends of hyperbolic 3-manifolds of funnel type. For Teichmüller space definition and conventions, we follow the book of Tromba [35].

The Teichmüller space  $\mathcal{T}_g$  is defined here as the quotient  $\mathcal{M}_{-1}(\Sigma_g)/\mathcal{D}_0(\Sigma_g)$  where  $\mathcal{M}_{-1}(\Sigma_g)$  is the set of metrics with Gaussian curvature  $-1$  on a fixed smooth surface  $\Sigma_g$  of genus  $g$ , and  $\mathcal{D}_0(\Sigma_g)$  is the group of orientation-preserving diffeomorphisms of  $\Sigma_g$  which are isotopic to the identity. Let  $M_0, \dots, M_N$  be  $N$  Riemann surfaces of genera  $g_1, \dots, g_N$ . Denoting  $\mathbf{g} = (g_1, \dots, g_N) \in \mathbb{N}^N$ , we aim to define the Chern-Simons line bundle  $\mathcal{L}_{\mathbf{g}}$  over the Teichmüller space  $\mathcal{T}_{\mathbf{g}} = \mathcal{T}_{g_1} \times \mathcal{T}_{g_N}$  of the disjoint union  $M = \sqcup_{j=1}^N M_j$ .

First, we shall identify each point of  $T\mathcal{T}_{g_i}$  with an isometry class of 3-dimensional hyperbolic ends, with conformal infinity given by the base point.

**Definition 20.** A *hyperbolic funnel* is a couple  $(M, g)$  where  $M$  is a Riemann surface (not necessarily connected) equipped with a metric  $h_0$  of Gaussian curvature  $0, 1$  or  $-1$  and  $g$  is a metric on the product  $M \times (0, \epsilon)_x$  for some small  $\epsilon > 0$ , which is of the form

$$(25) \quad g = \frac{dx^2 + h(x)}{x^2}, \quad h(x) := h_0 + x^2 h_2 + \frac{1}{4} x^4 h_2 \circ h_2$$

where  $h_2$  is a symmetric tensor satisfying

$$\text{Tr}_{h_0}(h_2) = \kappa, \quad \text{div}_{h_0}(h_2) = 0.$$

It is shown in Fefferman-Graham [8] that  $M \times (0, \epsilon)$  equipped with such a metric  $g$  is a (non-complete) hyperbolic manifold if  $\epsilon > 0$  is chosen small enough, and conversely every end of a convex co-compact hyperbolic manifold with conformal infinity  $(M, \{h_0\})$  and  $\text{genus}(M) > 1$  is isometric to a unique funnel (25) with  $h_0$  the hyperbolic metric representing the conformal

class  $\{h_0\}$ . There is an action of the group  $\mathcal{D}$  of diffeomorphisms of  $M$  on the space of funnels, simply given by

$$\psi^*(M, g) := \left( M, \frac{dx^2 + \psi^*h(x)}{x^2} \right)$$

for all  $\psi \in \mathcal{D}$ , where  $\psi^*h(x)$  is the pull-back of the metric  $h(x)$  on  $M$ . Notice also that a funnel induces a representation of  $\pi_1(M)$  into  $\mathrm{PSL}_2(\mathbb{C})$  up to conjugation.

The tangent space  $T\mathcal{M}_{-1}$  has a natural inner product, the  $L^2$ -metric, defined as follows (see [35, Sec. 2.6]): let  $h_0 \in \mathcal{M}_{-1}$ , and  $h, k \in T_{h_0}\mathcal{M}_{-1}$ . Since  $\mathcal{M}_{-1}$  is a Fréchet submanifold in the space of symmetric tensors on  $M$ , it follows that  $T\mathcal{M}_{-1} \subset S^2(TM)$ ; define

$$(26) \quad \langle h, k \rangle := \int_M \langle h, k \rangle_{h_0} \mathrm{dvol}_{h_0}.$$

This scalar product is  $\mathcal{D}(M)$ -invariant.

For any  $h_0 \in \mathcal{M}_{-1}(M)$  consider the vector space

$$V_{h_0} := \{h \in C^\infty(S^2T^*M); \mathrm{Tr}_{h_0}(h) = 0; \mathrm{div}_{h_0}(h) = 0\},$$

i.e., the set of transverse traceless symmetric tensors with respect to  $h_0$ . This is a real vector space of finite dimension which is precisely the orthogonal complement in  $T_{h_0}\mathcal{M}_{-1}$  of the orbit of  $\mathcal{D}_0$  with respect to the  $L^2(M, h_0)$  inner product. When  $h_0$  varies, these spaces form a locally trivial vector bundle  $V$  over  $\mathcal{M}_{-1}(M)$  of rank  $6g-6$  (assuming that the genera of the connected components  $M_j$  are strictly larger than 1), which we think of as the horizontal tangent bundle in the principal Riemannian fibration  $\mathcal{M}_{-1}(M) \rightarrow \mathcal{T}_{\mathbf{g}}$ . The group  $\mathcal{D}(M)$  acts isometrically on this bundle by pull-back of tensors, and the restriction of this action to the subgroup  $\mathcal{D}_0(M)$  is free. The quotient of  $V$  by  $\mathcal{D}_0$  is identified [35, Sec 2.4] with the tangent bundle  $T\mathcal{T}_{\mathbf{g}}$  of the Teichmüller space of genus  $\mathbf{g}$ . Thus, Teichmüller space inherits a Riemannian metric called the *Weil-Petersson metric*. Explicitly, on vectors in  $T_{[h_0]}\mathcal{T}$  described by trace-free, divergence free symmetric tensors  $h, k$  with respect to a representative  $h_0 \in [h_0]$ , the Weil-Petersson metric is defined by

$$(27) \quad \langle h, k \rangle_{\mathrm{WP}} := \int_M \langle h, k \rangle_{h_0} \mathrm{dvol}_{h_0}.$$

The following is a direct consequence of the discussion above:

**Lemma 21.** *There is a canonical bijection  $\Psi$  from the total space of the horizontal tangent bundle  $V \rightarrow \mathcal{M}_{-1}(M)$  to the set  $F_{\mathbf{g}}$  of hyperbolic funnels of genus  $\mathbf{g}$ , defined explicitly by*

$$(28) \quad \Psi : (h_0, h_2^0) \mapsto \left( M, \frac{dx^2 + h(x)}{x^2} \right), \quad h(x) = h_0 + x^2 h_2 + \frac{x^4}{4} h_2 \circ h_2, \quad h_2 = h_2^0 + \frac{h_0}{2}.$$

*This bijection commutes with the action of  $\mathcal{D}_0$  on both sides and hence descends to a bijection from  $T\mathcal{T}_{\mathbf{g}}$  to the space of  $\mathcal{D}_0(M)$ -equivalence classes of hyperbolic funnels.*

Any divergence free traceless tensor  $k = udx^2 - udy^2 - 2vdx dy$  with respect to a metric  $h_0$  is the real part of a quadratic holomorphic differential (QHD in short)

$$\frac{1}{2}k = \mathrm{Re}(k^{0,1}) \quad \text{with } k^{0,1} := \frac{1}{2}(u + iv)dz^2 \text{ in local complex coordinates } z = x + iy.$$

The complex structure  $J$  on Teichmüller space is then given by multiplication by  $-i$  on QHD, which on the level of transverse traceless tensors means

$$(29) \quad Jk := vdx^2 - vdy^2 + 2udx dy$$

or setting  $K$  to be the symmetric endomorphism of  $TM$  defined by  $k(\cdot, \cdot) = h_0(K\cdot, \cdot)$ ,

$$JK = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} u & -v \\ -v & -u \end{pmatrix} = \begin{pmatrix} v & u \\ u & -v \end{pmatrix}.$$

The space  $T^{0,1}\mathcal{T}_{\mathbf{g}}$  is then defined to be the subspace of complexified tangent space  $T_{\mathbb{C}}\mathcal{T}_{\mathbf{g}}$  spanned by the elements  $k + iJk$  with  $k \in T\mathcal{T}_{\mathbf{g}}$ , and  $T^{1,0}\mathcal{T}_{\mathbf{g}}$  is spanned by the  $k - iJk$ . Notice also that, with the notations used just above, one has

$$(30) \quad k^{0,1} = \frac{1}{2}(k + iJk) \in T^{0,1}\mathcal{T}_{\mathbf{g}}, \quad k^{1,0} = \overline{k^{0,1}} = \frac{1}{2}(k - iJk) \in T^{1,0}\mathcal{T}_{\mathbf{g}}.$$

The Weil-Petersson metric on  $T\mathcal{T}_{\mathbf{g}}$  induces an isomorphism  $\Phi$  between  $T\mathcal{T}_{\mathbf{g}}$  and  $T^*\mathcal{T}_{\mathbf{g}}$ . It is also a Hermitian metric for the complex structure  $J$ , in the sense that  $\langle Jh, Jk \rangle_{\text{WP}} = \langle h, k \rangle_{\text{WP}}$  for all  $h, k \in T\mathcal{T}_{\mathbf{g}}$ , the associated symplectic form is  $\omega_{\text{WP}}(\cdot, \cdot) := \langle J\cdot, \cdot \rangle_{\text{WP}}$ . By convention, the metric  $\langle \cdot, \cdot \rangle_{\text{WP}}$  on  $T\mathcal{T}_{\mathbf{g}}$  is extended to be bilinear on  $T_{\mathbb{C}}\mathcal{T}_{\mathbf{g}}$ , so that  $\langle k^{0,1}, h^{0,1} \rangle_{\text{WP}} = \langle k^{1,0}, h^{1,0} \rangle_{\text{WP}} = 0$  for all  $h, k \in T\mathcal{T}_{\mathbf{g}}$  and  $\langle k^{0,1}, k^{1,0} \rangle_{\text{WP}} \geq 0$  for all  $k \in T\mathcal{T}_{\mathbf{g}}$ .

On  $T^*\mathcal{T}_{\mathbf{g}}$ , there is a natural symplectic form, obtained by taking the exterior derivative  $d\mu$  of the Liouville 1-form  $\mu$  defined for  $h_0 \in \mathcal{T}_{\mathbf{g}}$ ,  $k^* \in T_{h_0}^*\mathcal{T}_{\mathbf{g}}$  by

$$\mu_{(h_0, k^*)} := k^*.d\pi$$

if  $\pi : T^*\mathcal{T}_{\mathbf{g}} \rightarrow \mathcal{T}_{\mathbf{g}}$  is the natural projection. Since  $T^*\mathcal{T}_{\mathbf{g}}$  has a complex structure induced naturally by that of  $\mathcal{T}_{\mathbf{g}}$  we can also define the  $(0, 1)$  component  $\mu^{1,0}$  of the Liouville measure. The Liouville form  $\mu$  and  $\mu^{1,0}$  pull-back to natural form on  $T\mathcal{T}_{\mathbf{g}}$  through  $\Phi$ , satisfying

$$\Phi^*\mu_{(h_0, k)}(\dot{h}_0, \dot{k}) = \langle k, \dot{h}_0 \rangle_{\text{WP}}, \quad \Phi^*\mu_{(h_0, k)}^{1,0}(\dot{h}_0, \dot{k}) = \langle k^{0,1}, \dot{h}_0^{1,0} \rangle_{\text{WP}}$$

for  $(h_0, k) \in T\mathcal{T}_{\mathbf{g}}$ , and  $(\dot{h}_0, \dot{k}) \in TT_{h_0}\mathcal{T}_{\mathbf{g}} = T_{h_0}\mathcal{T}_{\mathbf{g}} \oplus T_{h_0}\mathcal{T}_{\mathbf{g}}$ . Notice that  $d\mu^{1,0}$  is a  $(1, 1)$  type form on  $T^*\mathcal{T}_{\mathbf{g}}$ .

**7.2. The cocycle.** In order to define the Chern-Simons line bundle  $\mathcal{L}_{\mathbf{g}}$  over  $\mathcal{T}_{\mathbf{g}}$  in a way similar to Freed [9] and Ramadas-Singer-Weitsman [31], we need to define a certain cocycle. The natural bundle turns out to be the  $\text{SO}(3)$  Chern-Simons line bundle associated to a 3-manifold bounding a given surface.

For each  $h_0 \in \mathcal{M}_{-1}$ , let  $(\overline{X}, \hat{g})$  be a compact Riemannian manifold with totally geodesic boundary  $(M, h_0)$ . Consider the map  $c^X : \mathcal{C}^\infty(M, F(X)) \times \mathcal{C}_{\text{ext}}^\infty(M, \text{SO}(3)) \rightarrow \mathbb{C}$  defined by

$$c^X(\hat{S}, a) = \exp \left( 2\pi i \int_M \frac{1}{16\pi^2} \text{Tr}(\hat{\omega} \wedge da a^{-1}) + 2\pi i \int_X \frac{1}{48\pi^2} \text{Tr}((\tilde{a}^{-1} d\tilde{a})^3) \right)$$

where  $\hat{\omega}$  is the connection form of the Levi-Civita connection of  $\hat{g}$  in any extension in  $\mathcal{C}^\infty(\overline{X}, F(X))$  of  $\hat{S}$  to  $\overline{X}$ , and  $\tilde{a}$  is any smooth extension of  $a$  on  $\overline{X}$  ( $\mathcal{C}_{\text{ext}}^\infty(M)$  means sections which are extendible to  $\overline{X}$ ). Note that any  $a \in \mathcal{C}^\infty(M, \text{SO}(3))$  can be extended to some  $\tilde{a}$  on a handlebody with boundary  $M$ . This definition is consistent by the following Lemma.

**Lemma 22.** *The value  $c^X(\hat{S}, a)$  depends only on  $h_0$ ,  $\hat{S}$  and  $a$ .*

*Proof.* In other words, we must prove that  $c^X(\hat{S}, a)$  does not depend on the choice of the even metric  $\hat{g}$ , on the choice of  $X$  bounding  $M$  and on the extensions of  $(\hat{S}, a)$  from  $M$  to  $X$ . The independence of  $\hat{\omega}|_{TM}$  with respect to  $\hat{g}$  and the extension of  $S$  is a consequence of Koszul

formula and the evenness to first order of  $\hat{g}$ . The independence with respect to the choice of  $\overline{X}$  and the extension  $\tilde{a}$  is a consequence of the fact that

$$\exp\left(-2\pi i \int_Z \frac{1}{48\pi^2} \text{Tr}((\tilde{a}^{-1} d\tilde{a})^3)\right) = 1$$

if  $Z$  is a compact manifold without boundary and  $\tilde{a} \in \mathcal{C}^\infty(Z, \text{SO}(3))$  (see (20)).  $\square$

We therefore get a map  $c^X : \mathcal{C}^\infty(M, F(X)) \times \mathcal{C}_{\text{ext}}^\infty(M, \text{SO}(3)) \rightarrow \mathbb{C}$  by associating to  $(\hat{S}|_M, a|_M)$  the quantity  $c^X(\hat{S}, a)$ , where the subscript ext denotes objects on  $M$  extendible to  $X$ . As we shall see below, when acting on  $\hat{S} \in \mathcal{C}_{\text{ext}}^\infty$ , this map satisfies the cocycle condition

$$(31) \quad c^X(\hat{S}, ab) = c^X(\hat{S}, a)c^X(\hat{S}a, b).$$

**7.3. The Chern-Simons line bundle  $\mathcal{L}_{\mathbf{g}}$ .** We follow the presentation given in the book of Baseilhac [4], but adapted to our setting.

**Definition 23.** The complex line  $L_{h_0}^X$  over  $h_0 \in \mathcal{M}_{-1}$  is defined for a choice of extension  $X$  by

$$L_{h_0}^X := \{f : \mathcal{C}_{\text{ext}}^\infty(M, F(X)) \rightarrow \mathbb{C}; \forall a \in \mathcal{C}_{\text{ext}}^\infty(M, \text{SO}(3)), f(\hat{S}a) = c(\hat{S}, a)f(\hat{S})\}.$$

We define the Chern-Simons line bundle (as a set) over  $\mathcal{M}_{-1}$  by

$$\mathcal{L}_{\mathbf{g}}^X := \bigsqcup_{h_0 \in \mathcal{M}_{-1}} L_{h_0}^X.$$

Using the Gauge transformation law (15), we deduce

**Lemma 24.** *For any metric  $\hat{g}$  on  $\overline{X}$  with  $\hat{g}|_{\partial\overline{X}} = h_0$ , the map  $\hat{S} \mapsto e^{2\pi i \text{CS}(\hat{g}, \hat{S})}$  is an element of the fiber over  $h_0$ .*

This fact directly implies the cocycle condition (31).

An element in  $L_{h_0}^X$  is determined by its value on any frame extendible to  $\overline{X}$  by the condition  $f(\hat{S}a) = c(\hat{S}, a)f(\hat{S})$ , therefore the dimension of the  $\mathbb{C}$ -vector space  $L_{h_0}^X$  is 1. If  $X_1$  and  $X_2$  are two fillings of  $M$ , let  $Z = X_1 \cup X_2$  be the oriented closed manifold obtained by gluing  $X_1$  and  $X_2$  along  $M$ , then any frame on  $Z$  restricted to  $M$  is extendible both to  $X_1$  and  $X_2$ . For such a frame  $\hat{S}$ , we define an isomorphism between  $\mathcal{L}_{\mathbf{g}}^{X_1}$  and  $\mathcal{L}_{\mathbf{g}}^{X_2}$  by setting

$$f^{X_1} \mapsto f^{X_2}, \quad f^{X_2}(\hat{S}) := f^{X_1}(\hat{S}).$$

By Lemma 22, this isomorphism is independent of the choice of  $\hat{S}$  extendible to  $X_1$  and  $X_2$ . Therefore we have a well defined bundle  $\mathcal{L}_{\mathbf{g}}$  over  $\mathcal{T}_{\mathbf{g}}$  independent of the filling  $X$ .

We define the smooth structure on  $\mathcal{L}_{\mathbf{g}}^X$  through global trivializations as follows: let  $\hat{S}$  be a smooth positively oriented frame (not a priori orthonormal) on  $\overline{X}$  and let  $\hat{S}_{h_0}$  be the orthonormal frame obtained from  $\hat{S}$  by Gram-Schmidt process with respect to the metric  $dx^2 + h_0$  near the boundary  $\partial\overline{X}$  and define a global trivialization by

$$\mathcal{L}_{\mathbf{g}}^X \rightarrow \mathcal{M}_{-1} \times \mathbb{C} \quad (h_0, f) \mapsto (h_0, f(\hat{S}_{h_0})).$$

Changes of trivializations corresponding to different choices of  $\hat{S}$  are smooth on  $\mathcal{M}_{-1}$ , thus we get a structure of smooth line bundle on  $\mathcal{L}_{\mathbf{g}}^X$  over  $\mathcal{M}_{-1}$ .

**7.4. Action by the mapping class group.** The mapping class group  $\text{Mod}_{\mathbf{g}}$  is the set of isotopy classes of orientation preserving diffeomorphisms of  $M = \partial\overline{X}$ , it acts on  $\mathcal{T}_{\mathbf{g}}$  properly discontinuously. By following Marden [23, Theorem 3.1], the subgroup  $\text{Mod}_X$  of  $\text{Mod}_{\mathbf{g}}$  arising from elements which extend to diffeomorphisms of  $\overline{X}$  homotopic to the identity on  $\overline{X}$  acts freely on  $\mathcal{T}_{\mathbf{g}}$  and the quotient  $\mathcal{T}_X := \mathcal{T}_{\mathbf{g}}/\text{Mod}_X$  is a complex manifold of dimension  $3|\mathbf{g}| - 3$ . Moreover the Weil-Petersson metric descends to  $\mathcal{T}_X$ .

Every diffeomorphism  $\psi : \overline{X} \rightarrow \overline{X}$  induces an isomorphism  $L_{h_0}^X \rightarrow L_{\psi^*h_0}^X$  by  $f \mapsto f_\psi := (\hat{S} \mapsto f(\psi_*\hat{S}))$ . In particular since any  $\psi \in \mathcal{D}_0$  can be extended on  $\overline{X}$  as a diffeomorphism and the map  $L_{h_0}^X \rightarrow L_{\psi^*h_0}^X$  does not depend on the extension, the bundle  $\mathcal{L}_{\mathbf{g}}^X$  descends to  $\mathcal{T}_{\mathbf{g}}$  as a complex line bundle.

We define the pull-back bundle  $\pi^*\mathcal{L}_{\mathbf{g}}$  on  $T\mathcal{T}_{\mathbf{g}}$  if  $\pi : T\mathcal{T}_{\mathbf{g}} \rightarrow \mathcal{T}_{\mathbf{g}}$  is the projection on the base and we shall use the notation  $\mathcal{L}_{\mathbf{g}}$  instead of  $\pi^*\mathcal{L}_{\mathbf{g}}$ .

In what follows, we shall work with Teichmüller space but all constructions are  $\text{Mod}_X$  invariants and descend to  $\mathcal{T}_X$ .

**7.5. Hermitian metric on  $\mathcal{L}_{\mathbf{g}}$ .** Since the cocycle is of absolute value 1, there exists on  $\mathcal{L}_{\mathbf{g}}$  a canonical Hermitian metric, denoted  $\langle \cdot, \cdot \rangle_{\text{CS}}$ , given simply by

$$(32) \quad \langle f_1, f_2 \rangle_{\text{CS}} := f_1(\hat{S})\overline{f_2(\hat{S})}$$

if  $f_1, f_2$  are two sections of  $\mathcal{L}_{\mathbf{g}}$  and  $\hat{S} \in C^\infty(M, F(X))$ .

**7.6. The connections on  $\mathcal{L}_{\mathbf{g}}$ .** We define 2 different connections on  $\mathcal{L}_{\mathbf{g}}$ . We start with a Hermitian connection coming from the base  $\mathcal{T}_{\mathbf{g}}$ .

**Definition 25.** Let  $h_0^t \in \mathcal{M}_{-1}$  for  $t \in \mathbb{R}$  be a curve of hyperbolic metrics on  $M$  extended evenly to first order to a metric  $\hat{g}^t$  on  $X$ , with  $h_0^0 =: h_0$ . For any  $\hat{g}$ -orthonormal frame  $\hat{S}$  on  $X$  we define  $\hat{S}^t$  to be the parallel transport of  $\hat{S}$  in the  $t$  direction with respect to the metric  $\hat{G} := dt^2 + \hat{g}^t$  ( $\hat{S}^t$  is then a  $\hat{g}^t$ -orthonormal frame). For any section  $f$  of  $\mathcal{L}_{\mathbf{g}}$ , we define for  $\dot{h}_0 = \partial_t h_0^t|_{t=0} \in T_{h_0}\mathcal{M}_{-1}$

$$(\nabla_{\dot{h}_0}^{\text{SO}(3)} f)(\hat{S}) := \partial_t f(h_0^t, \hat{S}^t)|_{t=0} - 2\pi i f(\hat{S}) \int_M \frac{1}{16\pi^2} \text{Tr}(\dot{\omega} \wedge \hat{\omega})$$

where  $\dot{\omega} = \partial_t \hat{\omega}^t|_{t=0}$  and  $\hat{\omega}^t$  is the Levi-Civita connection 1-form ( $\text{so}(3)$ -valued) of the metric  $\hat{g}^t$  in the frame  $\hat{S}^t$ .

This connection is  $\mathcal{D}(\overline{X})$  invariant (recall that  $\mathcal{D}(\overline{X})$  acts on  $\mathcal{L}_{\mathbf{g}}^X$  over  $\mathcal{M}_{-1}$ ), thus we get a connection in the Chern-Simons bundle over  $\mathcal{T}_{\mathbf{g}}$  and any of its quotients by a subgroup of the mapping class group acting freely on  $\mathcal{T}_{\mathbf{g}}$  whose elements can be realized as diffeomorphism of  $\overline{X}$  for some given  $\overline{X}$  bounding  $M$ .

A straightforward application of Koszul formula shows the

**Lemma 26.** *Let  $S \in C^\infty(\overline{X}, F_0(X))$  be an even to first order orthonormal frame on  $X$  with respect to an even to first order AH metric  $g$ , and let  $g^t$  be a curve of even to first order AH metrics with  $g^0 = g$ . Write the metric  $g^t$  near the boundary under the funnel form (25)*

$$g^t = \frac{dx^2 + h^t(x)}{x^2}.$$

Then the parallel transported frame  $S^t$  of  $S$  in the  $t$  direction with respect to the metric  $G = dt^2 + g^t$  is equal to  $x\hat{S}^t$  where  $x$  is a geodesic boundary defining function for  $g$  and  $\hat{S}^t$  is the parallel transported frame of  $\hat{S} := x^{-1}S$  for  $\hat{G}^t = dt^2 + x^2g^t$  in the  $t$ -direction.

**7.7. The curvature of  $\nabla^{\text{SO}(3)}$ .** Consider the trivial fibration  $\mathcal{M}_{-1} \times M \rightarrow \mathcal{M}_{-1}$  with fiber type  $M$ , with metric  $h$  along the fiber above  $h \in \mathcal{M}_{-1}$ . The action of the group  $\mathcal{D}(M)$  on  $\mathcal{M}_{-1}$  extends isometrically to the fibers, thus by quotienting through the free action of  $\mathcal{D}_0$  we obtain the so-called universal curve  $\mathcal{F} \rightarrow \mathcal{T}_{\mathbf{g}}$  with fiber type  $M$ , which is a Riemannian submersion over  $\mathcal{T}_{\mathbf{g}}$ . In the proof below we shall consider the restriction of the fibration  $\mathcal{M}_{-1} \times M \rightarrow \mathcal{M}_{-1}$  above the image of a local section in  $\mathcal{M}_{-1} \rightarrow \mathcal{T}_{\mathbf{g}}$ . The resulting trivial fibration is canonically diffeomorphic to an open set in  $\mathcal{F}$  but not isometric, although the identification is an isometry along the fibers.

**Proposition 27.** *The curvature of  $\nabla^{\text{SO}(3)}$  equals  $\frac{i}{8\pi}\omega_{\text{WP}}$ , where  $\omega_{\text{WP}}$  denotes the Weil-Petersson symplectic form on  $\mathcal{T}_{\mathbf{g}}$ ,  $\omega_{\text{WP}}(U, V) = \langle JU, V \rangle_{\text{WP}}$ .*

*Proof.* Let  $f$  be a local section in  $\mathcal{L}_{\mathbf{g}} \rightarrow \mathcal{U} \subset \mathcal{T}_{\mathbf{g}}$  constructed as follows: first, choose a local section  $s : \mathcal{U} \subset \mathcal{T}_{\mathbf{g}} \rightarrow \mathcal{M}_{-1}$  in the principal fibration  $\mathcal{M}_{-1} \rightarrow \mathcal{T}_{\mathbf{g}}$ , i.e., a smooth family of hyperbolic metrics  $\mathcal{U} \ni [h] \mapsto h$  which by projection give a local parametrization of  $\mathcal{T}_{\mathbf{g}}$ . By restricting the metric of  $\mathcal{M}_{-1} \times M$  to  $s(\mathcal{U}) \times M =: \mathcal{M}^{\mathcal{U}}$ , we obtain a metric on  $\mathcal{M}^{\mathcal{U}}$  with respect to which  $s(\mathcal{U})$  and  $M$  are orthogonal, the projection on  $\mathcal{U}$  is a Riemannian submersion on the Weil-Petersson metric (26) on  $s(\mathcal{U})$ , and the metric on the fiber  $\{h\} \times M$  is  $h = s([h])$ . Next, extend each metric  $h \in s(\mathcal{U})$  to a metric  $g_{[h]}$  on a fixed compact manifold  $X$  with boundary  $M$ , so that for each  $[h] \in \mathcal{U}$ ,  $g_{[h]}$  restricts to  $h$  on  $M$ , has totally geodesic boundary, and depends smoothly on  $[h]$ . We get in this way a metric  $G$  on  $\mathcal{X}^{\mathcal{U}} := s(\mathcal{U}) \times X$  with respect to which  $s(\mathcal{U})$  and  $X$  are orthogonal, the projection on  $\mathcal{U}$  is a Riemannian submersion, and the metric on the fiber  $\{h\} \times X$  is  $g_{[h]}$ . Define

$$f : \mathcal{U} \rightarrow \mathcal{L}_{\mathbf{g}}, \quad f([h]) := e^{2\pi i \text{CS}(g_{[h]}, \cdot)}.$$

Let  $\mathbb{R} \ni t \mapsto h^t$  be a smooth curve in  $s(\mathcal{U})$  parametrized by arc-length and  $\dot{h}$  its tangent vector at  $t = 0$ . By the variation formula (13), the covariant derivative of the section  $f$  in the direction  $[\dot{h}]$  is

$$(\nabla_{[\dot{h}]}^{\text{SO}(3)} f)(S) = -\frac{2\pi i}{16\pi^2} f(S) \int_X 2\text{Tr}(\dot{\omega} \wedge \Omega),$$

where  $\Omega$  is the curvature tensor of  $g_{[h]}$  on  $X$ , and  $\omega$  is the connection 1-form, in any orthonormal frame  $S$  for  $h^0$ , parallel transported in the direction of  $\partial_t$  with respect to the metric  $G = dt^2 + g^t$ , where  $g^t := g([h^t])$ . Therefore the connection 1-form  $\alpha \in \Lambda^1(\mathcal{U})$  of  $\nabla^{\text{SO}(3)}$  in the trivialization  $f$  is given by

$$\alpha([\dot{h}]) = \frac{1}{4\pi i} \int_X \sum_{i,j=1}^3 \dot{\omega}_{ij} \wedge \Omega_{ji}$$

(we note that this does not depend on  $S$  anymore). Let  $R^G$  be the curvature tensor of  $G$ , and  $R^V$  the curvature of the vertical connection  $\nabla^V := \Pi_{TX} \circ \nabla^G$ . As a side note, we remark that this vertical connection is independent on the choice of metric on the horizontal distribution, so we could have chosen in the definition of  $G$  any other metric, for instance the one induced



from  $\mathcal{T}_{\mathbf{g}}$  via the projection. We compute

$$\partial_t \omega_{ij}(Y) = \langle R_{\partial_t, Y}^G S_j, S_i \rangle = \langle R_{\partial_t, Y}^V S_j, S_i \rangle, \quad \Omega_{ji}(Y_2, Y_3) = \langle R_{Y_2, Y_3}^G S_i, S_j \rangle = \langle R_{Y_2, Y_3}^V S_i, S_j \rangle$$

where the scalar products are with respect to  $G$ . This implies

$$\alpha([\dot{h}]) = \frac{1}{8\pi i} \int_X \partial_t \lrcorner \text{Tr}((R^V)^2).$$

The Chern-Simons form of the connection 1-form  $\omega^V$  of  $\nabla^V$  in a vertical frame  $S$  is a transgression for the Chern-Weil form  $\text{Tr}((R^V)^2)$ :

$$d\text{cs}(\omega^V) = \text{Tr}((R^V)^2).$$

Writing  $d = d^X + d^{\mathbb{R}}$  and using Stokes, we get

$$\alpha([\dot{h}]) = \frac{1}{8\pi i} \left( \int_M \partial_t \lrcorner \text{cs}(\omega^V) \right) + \frac{1}{8\pi i} \partial_t \left( \int_X \text{cs}(\omega^V) \right) \Big|_{t=0}.$$

Thus the connection 1-form of  $\nabla^{\text{SO}(3)}$  over  $s(\mathcal{U})$  satisfies

$$\alpha = \frac{1}{8\pi i} \int_{\mathcal{M}^{\mathcal{U}}/s(\mathcal{U})} \text{cs}(\omega^V) + \frac{1}{8\pi i} d \int_{\mathcal{X}^{\mathcal{U}}/s(\mathcal{U})} \text{cs}(\omega^V).$$

The second contribution is an exact form, the curvature of  $\nabla^{\text{SO}(3)}$  is therefore the horizontal exterior differential

$$R^{\text{SO}(3)} = d\alpha = \frac{1}{8\pi i} \int_{\mathcal{M}^{\mathcal{U}}/s(\mathcal{U})} d^H \text{cs}(\omega^V).$$

By Stokes, we can add inside the integral the vertical exterior differential, thus

$$(33) \quad R^{\text{SO}(3)} = \frac{1}{8\pi i} \int_{\mathcal{M}^{\mathcal{U}}/s(\mathcal{U})} d\text{cs}(\omega^V) = \frac{1}{8\pi i} \int_{\mathcal{M}^{\mathcal{U}}/s(\mathcal{U})} \text{Tr}((R^V)^2) = \frac{1}{8\pi i} \int_{\mathcal{M}^{\mathcal{U}}/s(\mathcal{U})} \text{Tr}(R^2).$$

Here  $R$  is the curvature of the vertical tangent bundle of  $\mathcal{M}^{\mathcal{U}} \rightarrow s(\mathcal{U})$  with respect to the natural connection induced by the vertical metric and the horizontal distribution. Notice that the vertical tangent bundle of the fibration  $\mathcal{X}^{\mathcal{U}} \rightarrow s(\mathcal{U})$  splits orthogonally along  $\mathcal{M}^{\mathcal{U}} \rightarrow s(\mathcal{U})$  into a flat real line bundle corresponding to the normal bundle to  $M \subset X$ , and the tangent bundle to the fibers of  $\mathcal{M}^{\mathcal{U}}$ . Thus in the above Chern-Weil integral we can eliminate the normal bundle to  $M$  in  $X$ , which justifies the last equality in (33).

Next, we compute explicitly this integral along the fibers of the universal curve in terms of the Weil-Petersson form on  $\mathcal{T}$ . Take a 2-parameters family  $h^{t,s}$  in  $\mathcal{M}_{-1}$  and let  $\dot{H}^t, \dot{H}^s \in \text{End}(TM)$  be defined by

$$\partial_t h^{t,s} \Big|_{t=s=0} = h(\dot{H}^t \cdot, \cdot), \quad \partial_s h^{t,s} \Big|_{t=s=0} = h(\dot{H}^s \cdot, \cdot).$$

where  $h := h^{t,s} \Big|_{t=s=0}$ .

Let  $X_1, X_2$  be a local frame on  $M$ , orthogonal at some point  $p \in M$  with respect to the metric  $h$ , and  $R$  the curvature of the connection on  $TM$  over  $\mathbb{R}^2 \times M$ .

**Lemma 28.** *At the point  $p \in M$  where the frame  $X_j$  is orthonormal we have*

$$R_{\partial_s \partial_t} X_j = -\frac{1}{4} [\dot{H}^s, \dot{H}^t] X_j, \quad \langle R_{X_1, X_2} X_2, X_1 \rangle = -1$$

and if we choose the family  $h$  such that  $\dot{H}^t \in V_h$ , the space of transverse traceless symmetric 2-tensors, then

$$R_{\partial_t X_j} = 0.$$

*Proof.* We first compute from Koszul's formula

$$\langle \nabla_{\partial_t} X_i, X_j \rangle = \frac{1}{2} \partial_t \langle X_i, X_j \rangle = \frac{1}{2} \langle \dot{H}^t(X_i), X_j \rangle,$$

so  $\nabla_{\partial_t} X_i = \frac{1}{2} \dot{H}^t(X_i)$  and similarly  $\nabla_{\partial_s} X_i = \frac{1}{2} \dot{H}^s(X_i)$ . Next we compute

$$\begin{aligned} \langle R_{\partial_s \partial_t} X_i, X_j \rangle &= \langle \nabla_{\partial_s} \nabla_{\partial_t} X_i, X_j \rangle - \langle \nabla_{\partial_t} \nabla_{\partial_s} X_i, X_j \rangle \\ &= \partial_s \langle \nabla_{\partial_t} X_i, X_j \rangle - \langle \nabla_{\partial_t} X_i, \nabla_{\partial_s} X_j \rangle - \partial_t \langle \nabla_{\partial_s} X_i, X_j \rangle + \langle \nabla_{\partial_s} X_i, \nabla_{\partial_t} X_j \rangle \\ &= \frac{1}{2} \partial_s \partial_t \langle X_i, X_j \rangle - \frac{1}{4} \langle \dot{H}^t(X_i), \dot{H}^s(X_j) \rangle - \frac{1}{2} \partial_t \partial_s \langle X_i, X_j \rangle + \frac{1}{4} \langle \dot{H}^s(X_i), \dot{H}^t(X_j) \rangle \end{aligned}$$

which proves the first identity of the lemma. The second identity is simply the fact that metric along the fibers has curvature  $-1$ . For the third, assume that  $\dot{H}^t$  is transverse traceless. At a fixed point  $p \in M$  choose a holomorphic coordinate  $z = x_1 + ix_2$  for  $h$  such that  $h = |dz^2| + O(|z|^2)$ , and choose  $X_1 = \partial_{x_1}, X_2 = \partial_{x_2}$ . Using that  $\nabla_{X_i} X_j = 0$  at  $p$ , we compute at that point

$$\begin{aligned} \langle \nabla_{\partial_t} \nabla_{X_1} X_1, X_2 \rangle &= \partial_t \langle \nabla_{X_1} X_1, X_2 \rangle = \partial_t (X_1 \langle X_1, X_2 \rangle - \frac{1}{2} X_2 \langle X_1, X_1 \rangle) \\ &= \partial_{x_1} \dot{H}_{12}^t - \frac{1}{2} \partial_{x_2} \dot{H}_{11}^t, \\ \langle \nabla_{X_1} \nabla_{\partial_t} X_1, X_2 \rangle &= X_1 \langle \nabla_{\partial_t} X_1, X_2 \rangle = \frac{1}{2} \partial_{x_1} \dot{H}_{12}^t \end{aligned}$$

which implies at  $p$

$$\langle R_{\partial_t, X_1} X_1, X_2 \rangle = \frac{1}{2} (\partial_{x_1} \dot{H}_{12}^t - \partial_{x_2} \dot{H}_{11}^t).$$

This last quantity vanishes by the Cauchy-Riemann equations when we expand  $\dot{H}_{ij}^t$  using  $\dot{H}^t = \Re(f(z)dz^2)$  for some holomorphic function  $f$ .  $\square$

Lemma (28) implies for the trace of the curvature at  $p \in M$

$$\begin{aligned} \text{Tr}(R^2)(\partial_s, \partial_t, X_1, X_2) &= 2\text{Tr}(R_{\partial_s, \partial_t} R_{X_1, X_2}) = 4\langle R_{\partial_s, \partial_t} X_1, X_2 \rangle \langle R_{X_1, X_2} X_2, X_1 \rangle \\ &= \langle [\dot{H}^s, \dot{H}^t] X_1, X_2 \rangle = -\langle \dot{H}^s \dot{H}^t J X_2, X_2 \rangle - \langle \dot{H}^t \dot{H}^s X_1, J X_1 \rangle \\ &= -\text{Tr}(J \dot{H}^s \dot{H}^t). \end{aligned}$$

Since  $\dot{H}^t$  is transverse traceless, the Weil-Petersson inner product of the vectors  $\partial_t, J\partial_s \in T_h \mathcal{T}$  is just the  $L^2$  product  $\int_M \text{Tr}(\dot{H}^t J \dot{H}^s) \text{dvol}_h$ . The proof is finished by applying (33).  $\square$

The identity (33) expressing the curvature of the Chern-Simons bundle as the fiberwise integral of the Pontrjagin form  $\text{Tr}(R^2)$  was proved for arbitrary surface fibrations by U. Bunke [6] in the context of smooth cohomology.

Since the curvature of the connection  $\nabla^{\text{SO}(3)}$  is a  $(1, 1)$  form, we get the

**Corollary 29.** *The complex line bundle  $\mathcal{L}_{\mathbf{g}}$  on  $\mathcal{T}_{\mathbf{g}}$  has a holomorphic structure induced by the connection  $\nabla^{\text{SO}(3)}$ , such that the  $\bar{\partial}$  operator is the  $(0, 1)$  component of  $\nabla^{\text{SO}(3)}$ .*

8. VARIATION OF THE CHERN-SIMONS INVARIANT AND CURVATURE OF  $\mathcal{L}_{\mathbf{g}}$ 

In this section we study the covariant derivative of the Chern-Simons  $\text{CS}(\theta)$  invariant viewed as a section in the pull-back of Chern-Simons bundle to  $T\mathcal{T}_{\mathbf{g}}$ .

By Proposition 18, Proposition 17 and Lemma 24, if  $g^t$  is a curve of convex co-compact hyperbolic 3-manifolds, then the invariant  $e^{2\pi i \text{CS}^{\text{PSL}_2(\mathbb{C})}(g^t, \cdot)}$  can be seen as a section of the line bundle  $\mathcal{L}_{\mathbf{g}}$  over a curve  $h_0^t \in \mathcal{M}_{-1}$  induced by the conformal infinities of  $g^t$ .

**Theorem 30.** *Let  $(X, g^t)$ ,  $t \in (-\epsilon, \epsilon)$ , be a smooth curve of convex co-compact hyperbolic 3-manifolds with conformal infinity a disjoint union of Riemann surface  $M = M_0 \sqcup \dots \sqcup M_N$  with genus  $\mathbf{g}$ , and such that  $g$  is isometric near  $M$  to the funnel  $(0, \epsilon)_x \times M$*

$$(34) \quad g^t = \frac{dx^2 + h^t(x)}{x^2}, \quad h(x) = h_0^t + x^2 h_2^t + \frac{1}{4} x^4 h_2^t \circ h_2^t,$$

with  $(h_0^t, h_2^t - \frac{1}{2}h_0^t) \in T\mathcal{T}_{\mathbf{g}}$ . Let  $S \in C^\infty(\overline{X}, F_0(\overline{X}))$  be an orthonormal frame for  $g^0$  and let  $S^t$  be the parallel transport of  $S$  in the  $t$  direction with respect to the metric  $G = dt^2 + g^t$  on  $X \times (-\epsilon, \epsilon)$ . Then, setting  $\dot{h}_0 := \partial_t h_0^t|_{t=0}$  and  $h_2 := (h_2^t - \frac{1}{2}h_0^t)|_{t=0}$  so that  $\dot{h}_0, h_2 \in T_{h_0} \mathcal{T}_{\mathbf{g}}$ , one has

$$\partial_t \text{CS}^{\text{PSL}_2(\mathbb{C})}(g^t, S^t)|_{t=0} = \frac{1}{16\pi^2} \int_M \text{Tr}(\dot{\omega} \wedge \dot{\omega}) + \frac{i}{8\pi^2} \langle (\text{Id} - iJ)\dot{h}_0, h_2 \rangle_{\text{WP}}.$$

Notice, by Lemma 26, that  $\hat{S}^t := x^{-1}S^t$  is parallel for  $\hat{G}^t = dt^2 + x^2 g^t$  and thus Theorem 30 is sufficient to compute the covariant derivative of  $e^{2\pi i \text{CS}^{\text{PSL}_2(\mathbb{C})}}$  with respect to  $\nabla^{\text{SO}(3)}$  in the direction of conformal infinities of hyperbolic metrics on  $X$ .

Before giving the proof, let us give as an application the variation formula for the renormalized volume.

**Corollary 31.** *Let  $X^t := (X, g^t)$  be a smooth curve of convex co-compact hyperbolic 3-manifolds like in Theorem 30. Then*

$$\partial_t (\text{Vol}_R(X^t))|_{t=0} = -\frac{1}{4} \langle \dot{h}_0, h_2 \rangle_{\text{WP}}.$$

*Proof.* It suffices to combine Theorem 30 with Proposition 18 and consider the imaginary part in the variation formula of  $\text{CS}^{\text{PSL}_2(\mathbb{C})}$ .  $\square$

This formula was proved by Krasnov and Schlenker [21], using the Schläfli formula, in order to show that the renormalized volume is a Kähler potential for the Weil-Petersson metric on Teichmüller space. The Chern-Simons approach thus provides another proof.

*Proof of Theorem 30.* Let  $T \in \Lambda^1(X, \text{End}(TX))$  be defined by  $T_U(V) := -U \times V$ , where  $\times$  denotes the vector product with respect to the Riemannian metric. Clearly  $T$  is anti-symmetric. We consider a 1-parameter family of metrics on  $X$  hyperbolic outside a compact set,  $g^t = \frac{dx^2 + h^t(x)}{x^2}$ , and we define a Riemannian metric on  $\mathbb{R} \times X$  by

$$G = dt^2 + g^t.$$

Recall that for every fixed  $t$ , the metrics  $h^t(x)$  and  $h_0^t := h^t(0)$  on  $M$  are related by (8).

For a given section  $S$  in the orthonormal frame bundle for  $g_0$ , we define  $S^t$  as the parallel transport in the  $t$  direction of  $S$ , more precisely  $\nabla_{\partial_t} S_j^t = 0$  for  $j = 1, 2, 3$ . Here and in

what follows,  $\nabla$  denotes the Riemannian connection for  $G$ . Since the integral curves of  $\partial_t$  are geodesics, it follows that  $S^t$  is an orthonormal frame for  $g^t$ .

Consider the connections  $D^t = \nabla^{g^t} + iT^t$  on  $T_{\mathbb{C}}X$  corresponding to the metric  $g^t$ . In the trivialization given by the section  $S^t$ , the connection form is  $\theta^t = \omega^t + iT^t$ . It is a  $\mathfrak{so}(3) \otimes \mathbb{C}$ -valued 1-form with real and imaginary parts

$$\omega_{ij}^t(Y) = \langle \nabla_Y^{g^t} S_j^t, S_i^t \rangle_{g^t}, \quad T_{ij}^t(Y) = \langle Y \times^t S_i^t, S_j^t \rangle_{g^t}.$$

We first compute the variation with respect to  $t$  of the Chern-Simons form of  $\theta^t$  on  $X$ . In what follows, we will drop the upperscript  $t$  when we evaluate at  $t = 0$  and we shall use a dot to denote the  $t$ -derivative at  $t = 0$ . Substituting in (13) for  $\theta = \omega + iT = \hat{\omega} - \alpha + iT$  like in (21), with  $\hat{\omega}$  the connection form of the Levi-Civita connection of the conformally compactified metric  $\hat{g} = x^2g$ , we get

$$\begin{aligned} \partial_t \text{cs}(\theta^t)|_{t=0} &= d(\text{Tr}(\dot{\omega} \wedge \omega) - \text{Tr}(\dot{T} \wedge T) + i(\text{Tr}(\dot{\omega} \wedge T) + \text{Tr}(\dot{T} \wedge \omega))) + 2\text{Tr}(\dot{\theta} \wedge \Omega^\theta) \\ (35) \quad &= d\text{Tr}(\dot{\omega} \wedge \hat{\omega}) + d[\text{Tr}(\dot{\alpha} \wedge \hat{\omega}) + \text{Tr}(\dot{\omega} \wedge \alpha)] + id[\text{Tr}(\dot{\omega} \wedge T) + \text{Tr}(\dot{T} \wedge \omega)] \\ &\quad + d[\text{Tr}(\dot{\alpha} \wedge \alpha) - \text{Tr}(\dot{T} \wedge T)] + 2\text{Tr}(\dot{\theta} \wedge \Omega^\theta). \end{aligned}$$

Observe that if  $g^t$  is a variation through hyperbolic metrics on  $X$  then  $\Omega^\theta$  vanishes. We claim that in the variation formula for the Chern-Simons invariant of  $\theta^t$ , the finite parts corresponding to the terms  $\text{Tr}(\dot{\alpha} \wedge \alpha)$  and  $\text{Tr}(\dot{T} \wedge T)$  vanish. We start with the term  $\text{Tr}(T \wedge \dot{T})$ :

**Lemma 32.** *We have  $\text{FP}_{\epsilon=0} \int_{x=\epsilon} \text{Tr}(T \wedge \dot{T}) = 0$ .*

*Proof.* Let  $Y_1, Y_2$  be vector fields on  $M$ , independent of  $t$ , then using that  $\nabla_{\partial_t}^G S_j = 0$ , we have  $\dot{T}_{ij}(Y_k) = \langle \nabla_{\partial_t}^G Y_k \times S_i, S_j \rangle$  so

$$\begin{aligned} \text{Tr}(T \wedge \dot{T})(Y_1, Y_2) &= \sum_{i,j} \langle Y_1 \times S_i, S_j \rangle \langle \nabla_{\partial_t}^G Y_2 \times S_i, S_j \rangle - \langle Y_2 \times S_i, S_j \rangle \langle \nabla_{\partial_t}^G Y_1 \times S_i, S_j \rangle \\ &= \langle Y_1, \nabla_{\partial_t} Y_2 \rangle - \langle Y_2, \nabla_{\partial_t} Y_1 \rangle \end{aligned}$$

which is zero because by Koszul,  $\langle Y_1, \nabla_{\partial_t} Y_2 \rangle = \frac{1}{2}(L_{\partial_t} G)(Y_1, Y_2)$  is symmetric in  $Y_1, Y_2$ .  $\square$

**Lemma 33.** *For  $\epsilon > 0$  sufficiently small we have*

$$\text{Tr}(\dot{\alpha} \wedge \alpha)|_{x=\epsilon} = 0$$

*Proof.* Let  $Y_1, Y_2$  be tangent vector fields to  $M$ , independent of  $t$ . Notice that for  $S_i$  parallel with respect to  $\nabla_{\partial_t}^G$  then  $\hat{S}_i = x^{-1}S_i$  is parallel with respect to  $\nabla_{\partial_t}^{\hat{G}}$  where  $\hat{G} = dt^2 + \hat{g}^t$ . Then since  $\partial_x$  is also killed by  $\nabla_{\partial_t}^{\hat{G}}$

$$\begin{aligned} x\text{Tr}(\dot{\alpha} \wedge \alpha)(Y_1, Y_2) &= \partial_t[\hat{g}^t(Y_1, \hat{S}_i^t)\hat{S}_j^t(x) - \hat{g}^t(Y_1, \hat{S}_j^t)\hat{S}_i^t(x)]|_{t=0}(\hat{g}(Y_2, \hat{S}_j)\hat{S}_i(x) - \hat{g}(Y_2, \hat{S}_i)\hat{S}_j(x)) \\ &\quad - \text{Sym}(Y_1 \rightarrow Y_2) \\ &= -2\hat{G}(\nabla_{\partial_t}^{\hat{G}} Y_1, Y_2) + 2\hat{G}(\nabla_{\partial_t}^{\hat{G}} Y_2, Y_1) = -2\hat{G}(\nabla_{Y_1}^{\hat{G}} \partial_t, Y_2) + 2\hat{G}(\nabla_{Y_2}^{\hat{G}} \partial_t, Y_1) \end{aligned}$$

$$x\text{Tr}(\dot{\alpha} \wedge \alpha)(Y_1, Y_2) = 2\hat{G}(\nabla_{Y_1}^{\hat{G}} Y_2, \partial_t) - 2\hat{G}(\nabla_{Y_2}^{\hat{G}} Y_1, \partial_t) = 2\hat{G}([Y_1, Y_2], \partial_t) = 0$$

and this finishes the proof.  $\square$

We now consider the term  $\text{Tr}(\dot{\omega} \wedge \alpha) + \text{Tr}(\dot{\alpha} \wedge \hat{\omega})$ .

**Lemma 34.** *Let  $\dot{H}_0$  and  $A$  be the symmetric endomorphism on  $TM$  defined by  $\dot{h}_0(\cdot, \cdot) = h_0(\dot{H}_0 \cdot, \cdot)$  and  $h_2(\cdot, \cdot) = h_0(A \cdot, \cdot)$ . We have the following identity*

$$\text{FP}_{\epsilon \rightarrow 0} \left( \text{Tr}(\dot{\omega} \wedge \alpha) + \text{Tr}(\dot{\alpha} \wedge \dot{\omega}) \right) |_{x=\epsilon} = 2 \int_M \text{Tr}(J\dot{H}_0 A) \text{dvol}_{h_0}$$

where  $J$  is the complex structure on  $TM$ .

*Proof.* First, from the proof of Proposition 17, we know that  $\text{FP}_{\epsilon \rightarrow 0} \text{Tr}(\alpha \wedge \dot{\omega}) |_{x=\epsilon} = 0$ , and therefore

$$\text{FP}_{\epsilon \rightarrow 0} \left( \text{Tr}(\dot{\omega} \wedge \alpha) + \text{Tr}(\dot{\alpha} \wedge \dot{\omega}) \right) |_{x=\epsilon} = 2 \text{FP}_{\epsilon \rightarrow 0} \text{Tr}(\dot{\omega} \wedge \alpha).$$

Now, for  $Y_1, Y_2$  tangent to  $M$  and independent of  $t$ , we can use that  $\nabla_{\partial_t}^{\hat{G}} \hat{S}_i = 0$  and  $\hat{\omega}_{ij}(Y) = \hat{g}(\nabla_Y^{\hat{G}} \hat{S}_j, \hat{S}_i) = \hat{G}(\nabla_Y^{\hat{G}} \hat{S}_j, \hat{S}_i)$  to deduce

$$\hat{\omega}_{ij}(Y) = \partial_t \langle \nabla_Y \hat{S}_j, \hat{S}_i \rangle = \langle \nabla_{\partial_t} \nabla_Y \hat{S}_j, \hat{S}_i \rangle = \langle R_{\partial_t Y} \hat{S}_j, \hat{S}_i \rangle$$

where  $\hat{R}$  is the curvature tensor of  $\hat{G}$ , therefore

$$\begin{aligned} \text{Tr}(\dot{\omega} \wedge \alpha)(Y_1, Y_2) &= \sum_{i,j} \langle R_{\partial_t Y_1} \hat{S}_j, \hat{S}_i \rangle (\langle Y_2, S_j \rangle S_i(a) - \langle Y_2, S_i \rangle S_j(a)) \\ &\quad - \langle \hat{R}_{\partial_t Y_2} \hat{S}_j, \hat{S}_i \rangle (\langle Y_1, S_j \rangle S_i(a) - \langle Y_1, S_i \rangle S_j(a)) \\ &= 2(\langle \hat{R}_{\partial_t Y_2} Y_1, x^{-1} \partial_x \rangle - \langle \hat{R}_{\partial_t Y_1} Y_2, x^{-1} \partial_x \rangle) \\ &= 2x^{-1} \langle \hat{R}_{\partial_t, \partial_x} Y_1, Y_2 \rangle \end{aligned}$$

by Bianchi. Since we are interested in the finite part, we can modify  $Y_1, Y_2$  by a term of order  $x^2$  without changing the result, and we will take  $\tilde{Y}_i^t = (1 - \frac{1}{2}x^2 A^t)Y_i$  where the endomorphism  $A^t$  of  $TM$  is defined by  $h_2^t(\cdot, \cdot) = h_0^t(A^t \cdot, \cdot)$ . Then

$$\begin{aligned} \hat{G}(\hat{R}_{\partial_t, \partial_x} Y_1, Y_2) &= -\partial_x \left( \hat{g}(\nabla_{\partial_t}^{\hat{G}} \tilde{Y}_1^t, \tilde{Y}_2^t) \right) |_{t=0} + O(x^2) \\ &= -\frac{1}{2} \partial_x \left( \partial_t(\hat{g}^t(\tilde{Y}_1^t, \tilde{Y}_2^t)) + \hat{g}^t([\partial_t, \tilde{Y}_1^t], \tilde{Y}_2^t) |_{t=0} - \hat{g}^t([\partial_t, \tilde{Y}_2^t], \tilde{Y}_1^t) |_{t=0} \right) + O(x^2). \end{aligned}$$

The term  $\partial_t(\hat{g}^t(\tilde{Y}_1^t, \tilde{Y}_2^t)) |_{t=0}$  is easily seen to be a  $\dot{h}_0(Y_1, Y_2) + O(x^3)$  by using that  $\hat{g}^t = dx^2 + h_0^t + x^2 h_0^t(A^t \cdot, \cdot) + O(x^4)$ , while the other two terms are

$$\begin{aligned} \hat{g}^t([\partial_t, \tilde{Y}_1^t], \tilde{Y}_2^t) |_{t=0} - \hat{g}^t([\partial_t, \tilde{Y}_2^t], \tilde{Y}_1^t) |_{t=0} &= \frac{1}{2}x^2 h_0(\dot{A}Y_1, Y_2) - \frac{1}{2}x^2 h_0(\dot{A}Y_2, Y_1) + O(x^4) \\ &= \frac{1}{2}x^2 h_0((\dot{A} - \dot{A}^T)Y_1, Y_2) + O(x^4). \end{aligned}$$

but since  $A^t$  is symmetric with respect to  $h_0^t$ , we deduce by differentiating at  $t = 0$  that  $\dot{A} - \dot{A}^T = (\dot{H}_0 A)^T - \dot{H}_0 A$  and therefore

$$\begin{aligned} \hat{g}^t([\partial_t, \tilde{Y}_1^t], \tilde{Y}_2^t) |_{t=0} - \hat{g}^t([\partial_t, \tilde{Y}_2^t], \tilde{Y}_1^t) |_{t=0} &= \frac{1}{2}x^2 h_0(\dot{H}_0 A Y_1, J Y_1) + \frac{1}{2}x^2 h_0(\dot{H}_0 A Y_2, J Y_2) + O(x^4) \\ &= -\frac{1}{2}x^2 \text{Tr}(J\dot{H}_0 A) + O(x^4). \end{aligned}$$

We conclude that the limit of  $\frac{2}{x} \hat{G}(\hat{R}_{\partial_t, \partial_x} Y_1, Y_2)$  as  $x \rightarrow 0$  is given by  $\text{Tr}(J\dot{H}_0 A)$ .  $\square$

Next, we reduce the sum  $\text{Tr}(\dot{T} \wedge \omega) + \text{Tr}(\dot{\omega} \wedge T)$  as follows:

**Lemma 35.** *We have the following identity*

$$\text{FP}_{\epsilon=0} \int_{x=\epsilon} \text{Tr}(\dot{T} \wedge \omega) + \text{Tr}(\dot{\omega} \wedge T) = 2\text{FP}_{\epsilon=0} \int_{x=\epsilon} \text{Tr}(\dot{\omega} \wedge T).$$

*Proof.* It suffices use (24) to deduce that  $\partial_t \text{FP}_{\epsilon=0} \int_{x=\epsilon} \text{Tr}(\omega \wedge T) = 8\pi \partial_t(\chi(M)) = 0$ .  $\square$

**Proposition 36.** *Let  $\dot{H}_0$  be the endomorphism on  $TM$  defined by  $\dot{h}_0(\cdot, \cdot) = h_0(\dot{H}_0 \cdot, \cdot)$ . Then near  $x = 0$  we have*

$$\text{Tr}(\dot{\omega} \wedge T) = [-x^{-2} \text{Tr}(\dot{H}_0) + \text{Tr}(\dot{A}) - \frac{1}{2} \text{Tr}(A) \text{Tr}(\dot{H}_0) + \text{Tr}(\dot{H}_0 A)] \text{dvol}_{h_0} + O(x^2).$$

*Proof.* Notice that for every  $Y$  tangent to  $X$  we have  $\omega_{ij}^t(Y) = \omega_{ij}(Y)$ , as a simple consequence of the Koszul formula. For a vector field  $Y$  on  $X$  extended on  $\mathbb{R} \times X$  to be constant with respect to the flow of  $\partial_t$  we compute

$$(\partial_t \omega_{ij})(Y)|_{t=0} = \partial_t \langle \nabla_Y S_j^t, S_i^t \rangle|_{t=0} = \langle \nabla_{\partial_t} \nabla_Y S_j^t, S_i^t \rangle|_{t=0} = \langle R_{\partial_t, Y} S_j, S_i \rangle$$

where  $\langle \cdot, \cdot \rangle$  denotes the metric  $G$ . In the last equality we have used the fact that  $S_i^t$  is parallel in the direction of  $\partial_t$  and the vanishing of the bracket  $[\partial_t, Y]$ . By the symmetry of the Riemannian curvature tensor, we rewrite the last term as  $-\langle R_{S_i, S_j} \partial_t, Y \rangle$ . It follows that

$$\begin{aligned} \text{Tr}(\dot{\omega} \wedge T)(Y_1, Y_2) &= \sum_{i,j=1}^3 -\langle R_{S_i, S_j} \partial_t, Y_1 \rangle \langle Y_2 \times S_j, S_i \rangle + \langle R_{S_i, S_j} \partial_t, Y_2 \rangle \langle Y_1 \times S_j, S_i \rangle \\ (36) \quad &= \sum_{j=1}^3 -\langle R_{Y_2 \times S_j, S_j} \partial_t, Y_1 \rangle + \langle R_{Y_1 \times S_j, S_j} \partial_t, Y_2 \rangle \\ &= E(Y_1, Y_2) - E(Y_2, Y_1) \end{aligned}$$

where we have defined

$$E(Y, Z) := \sum_{j=1}^3 \langle R_{Y \times S_j, S_j}^G \partial_t, Z \rangle.$$

For every vector field  $Y$  on  $M$ , define a vector field  $\tilde{Y}^t$  on a neighborhood of  $M$  in  $X$  by

$$\tilde{Y}^t = (1 + \frac{x^2}{2} A^t)^{-1} Y$$

where  $h_2^t = h_0^t(A^t \cdot, \cdot)$ . From (8) we see that for any orthonormal frame  $Y_1, Y_2$  on  $M$  for  $h_0$ , the frame  $\tilde{Y}_1^t, \tilde{Y}_2^t$  at  $t = 0$  is also orthonormal on  $X$ . The complex structure  $J$  on  $\{t\} \times \{x\} \times M$  satisfies  $J\tilde{Y}^t = x \partial_x \times^t \tilde{Y}^t$ , so in particular  $J\tilde{Y} = \widetilde{JY}$  at  $t = 0$ .

**Lemma 37.** *Let  $Y, Z$  be vector fields on  $M$ . Then near  $x = 0$  we have the expansion*

$$(37) \quad E(J\tilde{Y}, \tilde{Z}) = x^{-2} \dot{h}_0(Y, Z) - \frac{1}{2} (h_0(\dot{A}Y, Z) + h_0(Y, \dot{A}Z)) - \dot{h}_0(AY, Z) + O(x^2).$$

*Proof.* The expression defining  $E$  is independent of the orthonormal frame  $S_j$  for  $g$ , thus we can compute it using the frame  $x\tilde{Y}, xJ\tilde{Y}, x\partial_x$  (all these are at  $t = 0$ ):

$$E(J\tilde{Y}, \tilde{Z}) = 2 \langle R_{x\tilde{Y}, x\partial_x} \partial_t, \tilde{Z} \rangle.$$

Note the following identities:

$$(38) \quad \tilde{Y}^t = Y - \frac{x^2}{2} A^t Y + O(x^4), \quad [x\partial_x, x\tilde{Y}^t] = x\tilde{Y}^t - x^3 A^t Y + O(x^5), \quad \nabla_{x\partial_x}^G x\tilde{Z}^t = O(x^4).$$

Also, note that  $\nabla_{\partial_x}^G \partial_t = 0$ . Using these facts, we get

$$\begin{aligned}
\langle \nabla_{x\tilde{Y}}^G \partial_t, x\tilde{Z} \rangle &= \frac{1}{2} (L_{\partial_t} G)(x\tilde{Y}^t, x\tilde{Z}^t)|_{t=0} \\
(39) \quad &= \frac{1}{2} ((\partial_t h_0^t((1 + \frac{x^2}{2} A^t) \cdot, (1 + \frac{x^2}{2} A^t) \cdot))(\tilde{Y}, \tilde{Z}))|_{t=0} + O(x^4) \\
&= \frac{1}{2} \dot{h}_0(Y, Z) + \frac{x^2}{4} (h_0(\dot{A}Y, Z) + h_0(Y, \dot{A}Z)) + O(x^4),
\end{aligned}$$

therefore by using (38)

$$\begin{aligned}
\langle R_{x\partial_x, x\tilde{Y}} \partial_t, x\tilde{Z} \rangle &= x\partial_x \langle \nabla_{x\tilde{Y}}^G \partial_t, x\tilde{Z} \rangle - \langle \nabla_{[x\partial_x, x\tilde{Y}]}^G \partial_t, x\tilde{Z} \rangle \\
&= (x\partial_x - 1) \langle \nabla_{x\tilde{Y}}^G \partial_t, x\tilde{Z} \rangle + x^2 \langle \nabla_{x\tilde{AY}}^G \partial_t, x\tilde{Z} \rangle \\
&= -\frac{1}{2} \dot{h}_0(Y, Z) + \frac{x^2}{4} (h_0(\dot{A}Y, Z) + h_0(Y, \dot{A}Z)) + \frac{x^2}{2} \dot{h}_0(AY, Z) + O(x^4)
\end{aligned}$$

(in the last step we have used (39) for  $AY$  in the place of  $Y$ ). Using the tensoriality of the curvature to get out the factors of  $x$ , we proved the lemma.  $\square$

Let us now write (all what follows is at  $t = 0$ )

$$Y = \tilde{Y} + \frac{x^2}{2} AY + O(x^4) = \tilde{Y} + \frac{x^2}{2} \widetilde{AY} + O(x^4).$$

By linearity we get

$$(40) \quad E(Y, Z) = E(\tilde{Y}, \tilde{Z}) + \frac{x^2}{2} (E(\widetilde{AY}, Z) + E(Y, \widetilde{AZ})) + O(x^2).$$

Assume now that  $Y_j$  have been chosen at a given point on  $M$  as (orthonormal) eigenvectors of  $A$  for  $h_0$  of eigenvalue  $\lambda_j$ , with  $JY_1 = Y_2$ . Then from (40) we get

$$\begin{aligned}
E(Y_2, Y_1) &= (1 + \frac{x^2}{2}(\lambda_1 + \lambda_2)) E(\tilde{Y}_2, \tilde{Y}_1), \\
E(Y_1, Y_2) &= (1 + \frac{x^2}{2}(\lambda_1 + \lambda_2)) E(\tilde{Y}_1, \tilde{Y}_2),
\end{aligned}$$

therefore from (36) and Lemma 37

$$\begin{aligned}
\text{Tr}(\dot{\omega} \wedge T)(Y_1, Y_2) &= E(Y_1, Y_2) - E(Y_2, Y_1) \\
&= (1 + \frac{x^2}{2} \text{Tr}(A))(E(\tilde{Y}_1, \tilde{Y}_2) - E(\tilde{Y}_2, \tilde{Y}_1)) \\
&= - (1 + \frac{x^2}{2} \text{Tr}(A))(E(J\tilde{Y}_1, \tilde{Y}_1) + E(J\tilde{Y}_2, \tilde{Y}_2)) \\
&= (1 + \frac{x^2}{2} \text{Tr}(A))(-x^{-2} \text{Tr}(\dot{H}^0) + \text{Tr}(\dot{A}) + \text{Tr}(\dot{H}_0 A) + O(x^2)) \\
\text{Tr}(\dot{\omega} \wedge T)(Y_1, Y_2) &= -x^{-2} \text{Tr}(\dot{H}_0) + \text{Tr}(\dot{A}) - \frac{1}{2} \text{Tr}(A) \text{Tr}(\dot{H}_0) + \text{Tr}(\dot{H}_0 A) + O(x^2).
\end{aligned}$$

which is the claim of Proposition 36.  $\square$

We are now in position to finish the proof of Theorem 30. Since we consider a family of hyperbolic metrics  $g^t$ , we have  $\text{Tr}(A^t) = -\frac{1}{2} \text{scal}_{h_0^t}$  by (9) so by Gauss-Bonnet the following integral is constant in  $t$ :

$$(41) \quad \int_M \text{Tr}(A^t) \text{dvol}_{h_0^t} = 2\pi \chi(M).$$

Using  $\partial_t \text{dvol}_{h_t}|_{t=0} = \frac{1}{2} \text{Tr}(\dot{h}_0) \text{dvol}_{h_0}$  we deduce by differentiating (41) that

$$\int_M (\text{Tr}(\dot{A}) + \frac{1}{2} \text{Tr}(A) \text{Tr}(\dot{H}_0)) \text{dvol}_{h_0} = 0$$

so

$$\int_{x=\epsilon} \text{Tr}(\dot{\omega} \wedge T) = -\epsilon^{-2} \partial_t \text{Vol}(M, h_0) + \int_M (2\text{Tr}(\dot{A}) + \text{Tr}(\dot{H}_0 A)) \text{dvol}_{h_0} + O(\epsilon^2).$$

This achieves the proof of Theorem 30.  $\square$

## 9. CHERN-SIMONS LINE BUNDLE AND DETERMINANT LINE BUNDLE

In Ramadas-Singer-Weitsman [31], the Chern-Simons line bundle was introduced on the moduli space  $\mathcal{A}_F^s/\mathcal{G}$  of irreducible flat  $\text{SU}(2)$  connections up to gauge, they showed that it has a natural connection whose curvature is (up to a factor of  $i$ ) the standard symplectic form, and a natural Hermitian structure. Quillen [30] defined the determinant line bundle over the space  $\{\bar{\partial}_A; A \in \mathcal{A}_F^s\}$  of d-bar operators for a given complex structure on the surface  $M$ : he showed that it descends to  $\mathcal{A}_F^s/\mathcal{G}$  as a Hermitian line bundle with a natural connection and with curvature the standard symplectic form (up to a factor of  $i$ ). Ramadas-Singer-Weitsman proved that these bundles are isomorphic as Hermitian line bundle with connection over  $\mathcal{A}_F^s/\mathcal{G}$ . Moreover there curvature form is of  $(1, 1)$  type with respect to the natural complex structure on  $\mathcal{A}_F^s/\mathcal{G}$  and therefore the line bundle admits a holomorphic structure.

**9.1. The submanifold  $\mathcal{H}$  of hyperbolic 3-manifolds.** In what follows, we shall construct, in particular cases, a similar isomorphism using our Chern-Simons invariant and the determinant of the Laplacian. Let us be more precise and first make the following assumption: if  $\mathcal{T}_{\mathbf{g}}$  is Teichmüller space for a given oriented surface  $M$  of genus  $\mathbf{g}$  (possibly not connected) and  $h_0 \in \mathcal{T}_{\mathbf{g}}$ , we assume that we fix a compact 3-manifold  $\bar{X}$  with boundary  $M$  with hyperbolic convex-cocompact metric  $F(h_0)$  whose conformal infinity is  $\{h_0\}$ .

**Proposition 38.** *There exists a neighborhood  $\mathcal{U} \subset \mathcal{T}_{\mathbf{g}}$  of  $h_0$  and a smooth map  $F : \mathcal{U} \rightarrow \mathcal{C}^\infty(\bar{X}, S_+^2({}^0T^*\bar{X}))$  such that  $F(h)$  is hyperbolic convex-cocompact with conformal infinity  $[h]$  for all  $h \in \mathcal{U}$ .*

*Proof.* The proof is written for instance in [29]. A quasiconformal approach can be found for instance in Marden [23].  $\square$

This map induces by Lemma 21 a local section in the tangent bundle of  $\mathcal{T}_{\mathbf{g}}$ . By Mostow rigidity [23, Theorem 2.12] and Marden [23, Theorem 3.1], this section is unique and extends to a global smooth section  $\sigma : \mathcal{T}_{\mathbf{g}} \rightarrow T\mathcal{T}_{\mathbf{g}}$ . The graph

$$(42) \quad \mathcal{H} := \{(h_0, \sigma(h_0)) \in T\mathcal{T}_{\mathbf{g}}; h_0 \in \mathcal{T}_{\mathbf{g}}\}$$

is then a smooth submanifold of  $T\mathcal{T}_{\mathbf{g}}$  of dimension  $\dim \mathcal{T}_{\mathbf{g}}$ . By uniqueness, the subgroup of modular transformations of  $M$  consisting of classes of diffeomorphisms which extend to  $\bar{X}$  leaves this section invariant, therefore  $\sigma$  descends to any quotient of  $\mathcal{T}_{\mathbf{g}}$  by such a subgroup. For instance, this applies to the deformation space of a given convex co-compact hyperbolic 3 manifold  $X = \Gamma \backslash \mathbb{H}^3$ , which is a quotient of  $\mathcal{T}_{\mathbf{g}}$  by the subgroup of the mapping class group corresponding to diffeomorphisms of the boundary  $\partial \bar{X}$  which extend to diffeomorphisms on  $\bar{X}$  homotopic to the Identity; this subgroup acts freely and the quotient space is a complex manifold by Theorem 3.1 in Marden [23].



The examples we have in mind are the quasi-Fuchsian and Schottky cases. For the quasi-Fuchsian setting, we take  $M = M_- \sqcup M_+$  a disjoint union of 2 surfaces of the same genus ( $> 1$ ). By a result of Bers [5], for any hyperbolic metric  $h_0 := (h_0^-, h_0^+) \in \mathcal{T}_{\mathbf{g}}$  on  $M$ , there exist a unique (modulo isometries) convex co-compact hyperbolic 3-manifold with conformal infinity  $h_0 := (h_0^-, h_0^+)$  on  $M = M_- \sqcup M_+$  and which is a topological cylinder. It is isometric near  $M_{\pm}$  to a funnel

$$\left( (0, \epsilon)_x \times M_{\pm}, \frac{dx^2 + h^{\pm}(x)}{x^2} \right), \quad h^{\pm}(x) = h_0^{\pm} + x^2 h_2^{\pm} + \frac{x^4}{4} h_2^{\pm} \circ h_2^{\pm}$$

and this defines the tensor  $h_2^0 = (h_2^+, h_2^-) - \frac{1}{2}h_0 \in T_{h_0}\mathcal{T}_{\mathbf{g}}$ . The map  $\sigma : h_0 \rightarrow h_2^0$  is then a smooth section of  $T\mathcal{T}_{\mathbf{g}}$ . This corresponds to the set of points in  $T\mathcal{T}_{\mathbf{g}}$  for which the associated 2 funnels are the 2 ends of a (topologically cylindrical) hyperbolic 3-manifold. For Schottky uniformisation, this is essentially similar: any hyperbolic metric  $h_0$  on a connected Riemann surface of genus  $\mathbf{g}$  can be realized as a quotient  $\Gamma \backslash \mathbb{C}^2$  by a Schottky group in  $\mathrm{PSL}_2(\mathbb{C})$ , and it is therefore the conformal infinity of the 3-dimensional hyperbolic manifold  $X = \Gamma \backslash \mathbb{H}^3$ , where  $\Gamma$  is considered as a group of isometries on  $\mathbb{H}^3$  by Poincaré extension. This manifold  $X$  is topologically a solid torus. Actually, to realize this in global fashion on a deformation space, one needs to pass to a quotient of  $\mathcal{T}_{\mathbf{g}}$ , called the Schottky space  $\mathfrak{S}_{\mathbf{g}}$ , for which one has a covering map  $\mathcal{T}_{\mathbf{g}} \rightarrow \mathfrak{S}_{\mathbf{g}}$ , see for instance Takhtajan-Zograf [34] or Hejhal [15] for details. Like in the quasi-Fuchsian case, there is a map  $\sigma : h_0 \rightarrow h_2^0 = h_2 - \frac{1}{2}h_0$  which is a smooth section of  $T\mathfrak{S}_{\mathbf{g}}$ . In general, in the context of Proposition 38, one defines similarly a section  $\sigma$  of  $T\mathfrak{S}_{\mathbf{g}}$  and a submanifold

$$\mathcal{H} := \{(h_0, \sigma(h_0)) \in T\mathfrak{S}_{\mathbf{g}}; h_0 \in \mathfrak{S}_{\mathbf{g}}\}$$

of  $T\mathfrak{S}_{\mathbf{g}}$ .

Let us introduce a new connection on the pull-back of  $\mathcal{L}_{\mathbf{g}}$  to  $T\mathcal{T}_{\mathbf{g}}$ , for which the  $\mathrm{PSL}_2(\mathbb{C})$  Chern-Simons section is flat along the deformation space of hyperbolic metrics on  $X$ .

**Definition 39.** We define the connection  $\nabla^{\mathrm{PSL}_2(\mathbb{C})}$  on  $\mathcal{L}_{\mathbf{g}}$  by

$$\nabla^{\mathrm{PSL}_2(\mathbb{C})} = \nabla^{\mathrm{SO}(3)} + \frac{1}{2\pi} \Phi^* \mu^{1,0}$$

where  $\mu$  is the Liouville 1-form on  $T^*\mathcal{T}_{\mathbf{g}}$  and  $\Phi : T\mathcal{T}_{\mathbf{g}} \rightarrow T^*\mathcal{T}_{\mathbf{g}}$  is the isomorphism induced by the Weil-Petersson metric. In what follows we will omit the identification  $\Phi$ .

This connection is *not* Hermitian with respect to  $\langle \cdot, \cdot \rangle_{\mathrm{CS}}$  since the form  $\mu^{1,0}$  is not purely imaginary. The Chern-Simons line bundle  $\mathcal{L}_{\mathbf{g}}$  equipped with the connection  $\nabla^{\mathrm{PSL}_2(\mathbb{C})}$  has curvature

$$(43) \quad \Omega_{\mathcal{L}_{\mathbf{g}}} = \frac{i}{8\pi} \Omega_{\mathrm{WP}} + \frac{1}{2\pi} \bar{\partial} \mu^{1,0}$$

with real part  $\mathrm{Re}(\Omega_{\mathcal{L}_{\mathbf{g}}}) = \frac{1}{4\pi} d\mu$ , where here and below,  $\omega_{\mathrm{WP}}$  is understood as  $\pi^* \omega_{\mathrm{WP}}$  if  $\pi = T\mathcal{T}_{\mathbf{g}} \rightarrow \mathcal{T}_{\mathbf{g}}$  is the projection on the basis.

**Theorem 40.** *The Chern-Simons invariant  $e^{2\pi i \mathrm{CS}^{\mathrm{PSL}_2(\mathbb{C})}}$  restricted to the submanifold  $\mathcal{H}$  in (42) is a parallel section of  $\mathcal{L}_{\mathbf{g}}|_{\mathcal{H}}$  for the connection  $\nabla^{\mathrm{PSL}_2(\mathbb{C})}$ . As a consequence,  $\mathcal{H}$  is a Lagrangian submanifold of  $T\mathcal{T}_{\mathbf{g}}$  for the standard symplectic Liouville form  $d\mu$  on  $T\mathcal{T}_{\mathbf{g}}$  obtained from pull-back by the duality isomorphism  $T\mathcal{T}_{\mathbf{g}} \rightarrow T^*\mathcal{T}_{\mathbf{g}}$  induced by  $\langle \cdot, \cdot \rangle_{\mathrm{WP}}$ .*

*Proof.* This is a direct consequence of the variation formula in Theorem 30 and the definition of the connection  $\nabla^{\mathrm{PSL}_2(\mathbb{C})}$ .  $\square$

It is proved by Krasnov [19] for Schottky cases and more generally by Takhtajan-Teo [33] for Kleinian groups of class A (see also Krasnov-Schlenker [21] for quasi-Fuchsian cases), that

$$\bar{\partial}\partial(\text{Vol}_R) = \frac{i}{16}\omega_{\text{WP}}$$

using previous work of Takhtajan-Zograf [34] on Liouville functional. Here we use our convention for Weil-Petersson metric. The Theorem above generalizes these results providing a unified treatment:

**Corollary 41.** *For  $h_0$  in an open set  $U \subset \mathcal{T}_{\mathbf{g}}$ , let  $X_{h_0} = (X, g_{h_0})$  be a smooth family of convex co-compact hyperbolic 3-manifolds with conformal infinity parametrized smoothly by  $h_0 \in U$ , then*

$$\bar{\partial}\partial\text{Vol}_R(X_{h_0}) = \frac{i}{16}\omega_{\text{WP}}.$$

*Proof.* Let  $\sigma : U \subset \mathcal{T}_{\mathbf{g}} \rightarrow T\mathcal{T}_{\mathbf{g}}$  be the smooth map  $h_0 \rightarrow (h_0, h_2^0)$  parametrizing the submanifold  $\mathcal{H}$ . We consider  $\text{Vol}_R(X_{h_0})$  as a function on  $U$ . By Corollary 31, we have for  $\dot{h}_0 \in T_{h_0}\mathcal{T}_{\mathbf{g}}$

$$\partial\text{Vol}_R(X_{h_0}) \cdot \dot{h}_0 = -\frac{1}{4}\sigma^*\mu^{1,0}(\dot{h}_0).$$

From the vanishing of the curvature  $\Omega_{\mathcal{L}_{\mathbf{g}}}$  on  $\mathcal{H}$  and the formula (43), we obtain for any  $\dot{h}_0, \dot{\ell}_0 \in T_{h_0}\mathcal{T}_{\mathbf{g}}$

$$d\mu_{\sigma(h_0)}^{1,0}(d\sigma \cdot \dot{h}_0, d\sigma \cdot \dot{\ell}_0) = -\frac{i}{4}\omega_{\text{WP}}(\dot{h}_0, \dot{\ell}_0)$$

and since  $\sigma^*d\mu_{h_0}^{1,0}(\dot{h}_0, \dot{\ell}_0) = d(\sigma^*\mu^{1,0})(\dot{h}_0, \dot{\ell}_0) = \bar{\partial}\partial\text{Vol}_R(X_{h_0})(\dot{h}_0, \dot{\ell}_0)$ , the proof is finished.  $\square$

**9.2. An isomorphism with the determinant line bundle.** Finally, we construct an explicit isomorphism of Hermitian line bundles between  $\mathcal{L}_{\mathcal{H}}$  and the determinant line bundle in the particular cases of quasi-Fuchsian and Schottky manifolds.

Let  $M$  be a marked Riemann surface of genus  $\mathbf{g}$ , i.e., a surface with a distinguished set of generators  $\alpha_1, \dots, \alpha_{\mathbf{g}}, \beta_1, \dots, \beta_{\mathbf{g}}$  of  $\pi_1(M, x_0)$  for some  $x_0 \in M$ . With respect to this marking, if a complex structure is given on  $M$ , there is a basis  $\varphi_1, \dots, \varphi_{\mathbf{g}}$  of holomorphic 1-forms such that  $\int_{\alpha_j} \varphi_i = \delta_{ij}$  and this defines the period matrix  $(\tau_{ij}) = (\int_{\beta_j} \varphi_i)$  whose imaginary part is positive definite since  $2\text{Im} \tau_{ij} = \langle \varphi_i, \varphi_j \rangle$ . Schottky groups are free groups generated by  $L_1, \dots, L_{\mathbf{g}} \in \text{PSL}_2(\mathbb{C})$  which map circles  $C_1, \dots, C_{\mathbf{g}} \subset \hat{\mathbb{C}} = \mathbb{C} \cup \{\infty\}$  to other circles  $C_{-1}, \dots, C_{-\mathbf{g}} \in \hat{\mathbb{C}}$  (with orientation reversed). Each element  $\gamma \in \Gamma$  is conjugated in  $\text{PSL}_2(\mathbb{C})$  to  $z \rightarrow q_{\gamma}z$  for some  $q_{\gamma} \in \mathbb{C}$  with  $|q_{\gamma}| < 1$ , called the multiplier of  $\gamma$ . The quotient of the discontinuity set  $\Omega_{\Gamma}$  of  $\Gamma$  by  $\Gamma$  is a closed Riemann surface and every closed Riemann surface of genus  $\mathbf{g}$  can be represented in this manner by a result of Koebe, see [10]. The Schottky group is marked if each  $C_k$  is homotopic to  $\alpha_k$  in the quotient  $\Gamma \backslash \Omega_{\Gamma}$ . The marked group is unique up to a global conjugation in  $\text{PSL}_2(\mathbb{C})$  and a normalization condition (by assigning the 2 fixed points of  $L_1$  and one of  $L_2$ ) can be set to fix it. One then obtains the Schottky space  $\mathfrak{S}_{\mathbf{g}}$  which covers the moduli space (i.e., the set of isomorphism classes of compact Riemann surface of genus  $\mathbf{g}$ ) but is covered by Teichmüller space  $\mathcal{T}_{\mathbf{g}}$  whose points are isomorphism classes of marked compact Riemann surfaces.

Since any Schottky group  $\Gamma \subset \text{PSL}_2(\mathbb{C})$  acts as isometries on  $\mathbb{H}^3$  as a convex co-compact group, there is a canonical hyperbolic 3-manifold  $\Gamma \backslash \mathbb{H}^3$  with conformal infinity given by  $\Gamma \backslash \Omega_{\Gamma}$ . This manifold denoted  $X$  is a handlebody with conformal boundary  $M$ . Let  $\mathcal{D}_X(M)$  the group of diffeomorphisms of  $M$  which extend to  $\bar{X}$  factored by the group  $\mathcal{D}_0$  of diffeomorphisms of  $M$  homotopic to Id.

The Chern-Simons line bundle defined on  $\mathcal{T}_{\mathbf{g}}$  above is acted upon by  $\mathcal{D}_X(M)$ , thus it descends to the Schottky space  $\mathfrak{S}_{\mathbf{g}}$  which is a quotient of  $\mathcal{T}_{\mathbf{g}}$  by a subgroup of  $\mathcal{D}_X(M)$ , we denote it  $\mathcal{L}_{\mathfrak{S}_{\mathbf{g}}}$ . The connection on  $\mathcal{L}_{\mathbf{g}}$  over  $\mathcal{T}_{\mathbf{g}}$  defined in Subsection 7.6 is  $\mathcal{D}_X(M)$  invariant, hence it descends to  $\mathfrak{S}_{\mathbf{g}}$ . The Liouville form on  $T\mathcal{T}_{\mathbf{g}}$  is  $\mathcal{D}(M)$  invariant and thus also descends to  $T\mathfrak{S}_{\mathbf{g}}$ , then the connection  $\nabla^{\text{SO}(3)}$  descends to  $T\mathfrak{S}_{\mathbf{g}}$ , we denote it  $\nabla^{\mathfrak{S}_{\mathbf{g}}}$ . Again, we can define the Lagrangian submanifold  $\mathcal{H} \subset T\mathfrak{S}_{\mathbf{g}}$  consisting of those funnels which extend to Schottky 3-manifolds. The operator  $\partial_{\Gamma} : \mathcal{C}^{\infty}(M) \rightarrow \mathcal{C}^{\infty}(M, \Lambda^{1,0}M)$  for a given complex structure induced by  $\Gamma$  on  $M$  is Fredholm on Sobolev spaces and, considered as a family of operators parametrized by points  $\Gamma \in \mathfrak{S}_{\mathbf{g}}$ , one can define its determinant line bundle  $\det(\partial)$  of  $\partial$ , as in Quillen [30], to be at  $\Gamma$  the line<sup>1</sup>

$$\det(\partial_{\Gamma}) := \Lambda^{\mathbf{g}}(\text{coker } \partial_{\Gamma})$$

when  $\mathbf{g} \in \mathbb{N}$  and  $\text{coker } \partial_{\Gamma} = \ker(\bar{\partial}_{\Gamma} : \mathcal{C}^{\infty}(M, \Lambda^{1,0}) \rightarrow \mathcal{C}^{\infty}(M, \Lambda^2(M))) =: H^{0,1}(\Gamma \backslash \Omega_{\Gamma})$  is the vector space of holomorphic 1-forms on  $M \simeq \Gamma \backslash \Omega_{\Gamma}$ . The line bundle  $\det(\partial)$  over  $\mathfrak{S}_{\mathbf{g}}$  is a holomorphic line bundle with a holomorphic canonical section

$$(44) \quad \varphi := \varphi_1 \wedge \cdots \wedge \varphi_{\mathbf{g}}$$

and is equipped with a Hermitian norm, called Quillen metric, defined as follows: for each Riemann surface  $\Gamma \backslash \Omega_{\Gamma}$  with  $\Gamma \in \mathfrak{S}_{\mathbf{g}}$ , let  $h_0$  be the associated hyperbolic metric obtained by uniformisation and define  $\det' \Delta_{h_0}$  the determinant of its Laplacian, as defined in Ray-Singer [32], then the Hermitian metric on  $\det(\partial)$  is given at  $\Gamma \in \mathfrak{S}_{\mathbf{g}}$  by

$$(45) \quad \|\varphi\|_Q^2 := \frac{\|\varphi\|_{h_0}^2}{\det' \Delta_{h_0}} = \frac{\det \text{Im } \tau}{\det' \Delta_{h_0}}$$

where  $\|\cdot\|_{h_0}$  is the Hermitian product on  $\Lambda^{\mathbf{g}}(\text{coker } \partial_{\Gamma})$  induced by the metric  $h_0$  on differential forms on  $M$ . We denote by  $\nabla^{\det}$  the unique Hermitian connection associated to the holomorphic structure on  $\det \partial$  and the Hermitian norm  $\|\varphi\|_Q$ .

To state the isomorphism between powers of Chern-Simons line bundle and a power of the determinant line bundle, we will use a formula proved by Zograf [39, 40] and generalized by McIntyre-Takhtajan [26]

**Theorem 42. [Zograf]** *There exists a holomorphic function  $F(\Gamma) : \mathfrak{S}_{\mathbf{g}} \rightarrow \mathbb{C}$  such that*

$$(46) \quad \frac{\det' \Delta_{h_0}}{\det \text{Im } \tau} = c_{\mathbf{g}} \exp\left(\frac{\text{Vol}_R(X)}{3\pi}\right) |F(\Gamma)|^2$$

where  $c_{\mathbf{g}}$  is a constant depending only on  $\mathbf{g}$  where  $X = \Gamma \backslash \mathbb{H}^3$  when we see  $\Gamma \subset \text{PSL}_2(\mathbb{C})$  as a group of isometries of  $\mathbb{H}^3$ , and  $h_0$  is the hyperbolic metric on  $\Gamma \backslash \Omega_{\Gamma} \simeq \partial \bar{X}$ . For points in  $\mathfrak{S}_{\mathbf{g}}$  corresponding to Schottky groups  $\Gamma$  with dimension of limit set  $\delta_{\Gamma} < 1$ , the function  $F(\Gamma)$  is given by the following absolutely convergent product:

$$(47) \quad F(\Gamma) = \prod_{\{\gamma\}} \prod_{m=0}^{\infty} (1 - q_{\gamma}^{1+m})$$

where  $q_{\gamma}$  is the multiplier of  $\gamma \in \Gamma$ , and  $\{\gamma\}$  runs over all distinct primitive conjugacy classes in  $\Gamma$  excluding the identity.

<sup>1</sup>We have ignored the kernel of  $\partial$  since it is only made of constants with norm given essentially by the Euler characteristic of  $M$  by Gauss-Bonnet, therefore not depending at all on the complex structure on  $M$ .

*Remark 43.* The formula (46) was in fact given in terms of Liouville action  $S$  instead of renormalized volume, but it has been shown that  $S = -4 \text{Vol}_R(X) + c_{\mathbf{g}}$  for some constant  $c_{\mathbf{g}}$  depending only on  $\mathbf{g}$ , by Krasnov [19] for Schottky manifolds and by Takhtajan-Teo [33] for quasi-Fuchsian manifolds.

We therefore deduce from this last theorem and our construction the following

**Theorem 44.** *On the Schottky space  $\mathfrak{S}_{\mathbf{g}}$ , the bundle  $\mathcal{L}_{\mathfrak{S}_{\mathbf{g}}}^{-1}$  is isomorphic to  $(\det \partial)^{\otimes 6}$  when equipped with their connections and Hermitian products induced by those of  $(\mathcal{L}_{\mathfrak{S}_{\mathbf{g}}}, \nabla^{\mathfrak{S}_{\mathbf{g}}}, \|\cdot\|_{\mathcal{L}})$  and  $(\det \partial, \nabla^{\det}, \|\cdot\|_Q)$ . There is an explicit isometric isomorphism of holomorphic Hermitian line bundles given by*

$$(\sqrt{c_{\mathbf{g}}}F\varphi)^{\otimes 6} \mapsto e^{-2\pi i \text{CS}^{\text{PSL}_2(\mathbb{C})}}.$$

where  $F$  and  $c_{\mathbf{g}}$  are the holomorphic functions and constants of Theorem 42,  $\varphi$  is the canonical section of  $\det \partial$  defined in (44).

*Proof.* The section  $e^{2\pi i \text{CS}^{\text{PSL}_2(\mathbb{C})}} \otimes \sqrt{c_{\mathbf{g}}}F\varphi)^{\otimes 6}$  is holomorphic and has constant norm in the Hermitian line bundle  $\mathcal{L}_{\mathfrak{S}_{\mathbf{g}}} \otimes (\det \partial)^{\otimes 6}$  which is flat with respect to the Hermitian connection, then it is parallel and provides an isomorphism with the trivial line bundle.  $\square$

*Remark 45.* Notice that the function  $F$  defined by the product (47) when  $\delta(\Gamma) < 1$  is known to extend analytically in  $\mathfrak{S}_{\mathbf{g}}$  by results of Zograf [39, 40], and our theorem provides another proof, assuming formula (46) only in the subset  $\{\Gamma \in \mathfrak{S}_{\mathbf{g}}; \delta(\Gamma) < 1\}$ .

*Remark 46.* In our previous work [14], we proved that

$$F(\Gamma) = |F(\Gamma)| \exp\left(-\frac{\pi i}{2} \eta(A)\right)$$

when  $\delta(\Gamma) < 1$ , where  $\eta(A)$  is the eta invariant of the signature operator  $A = *d + d*$  on odd dimensional forms on the Schottky 3-manifold  $\Gamma \backslash \mathbb{H}^3$ .

*Remark 47.* Using the result of McIntyre-Takhajan and McIntyre-Teo [26, 27], a similar result with different powers of the bundles is easily obtained in the Schottky and quasi-Fuchsian cases if one replaces the bundle  $\det \partial$  by the determinant line bundle  $\det \Lambda_n$  of the vector space of holomorphic  $n$ -differentials on  $M$ .

## APPENDIX A. CHERN-SIMONS INVARIANTS OF 3-MANIFOLDS WITH FUNNELS AND CUSPS OF RANK 2

In this appendix we show how to extend the results of this paper to include 3-manifolds of finite geometry with funnels as well as rank 2-cusps. We will concentrate on the cusps since funnels have already been treated.

By definition, a *cusps of maximal rank* is a half-complete warped product  $(a, \infty) \times M$  with metric  $dt^2 + e^{-2t}h$ , where  $h$  is a flat metric on  $M$ . Here  $M$  will be of dimension 2. After a linear change of variables in  $t$ , we can thus assume that  $M$  is isometric to a flat torus with a closed simple geodesic of length 1.

A manifold with asymptotic funnels and cusps is a complete Riemannian manifold which outside a compact is isometric to a finite disjoint union of hyperbolic funnels (Definition 20) and cusps in the above sense. The main examples of such manifolds are geometrically finite hyperbolic 3-manifolds without cusps of rank 1.

By changing variables  $x := e^{-t} \in (0, e^{-a})$ , the cusp metric becomes

$$\frac{dx^2}{x^2} + x^2 h = x^2 \left( \frac{dx^2}{x^4} + h \right).$$

Thus a cusp is conformal to a half-infinite cylinder  $dy^2 + h$  where  $y := x^{-1} = e^t \in [e^a, \infty)$ , the conformal factor being  $x = y^{-1}$ . The function  $x$  can be used to glue to the cusps a copy of  $M$  at  $x = 0$ , thus compactifying  $X$ . Thus if we choose  $\rho : X \rightarrow (0, \infty)$  to be a function which agrees with  $x$  on funnels and with  $y$  on cusps, it follows that a manifold with asymptotic funnels and cusps is conformal to a manifold with boundary (corresponding to the funnels) and flat half-infinite cylindrical ends (corresponding to each cusp):

$$g = \rho^{-2} \hat{g}, \quad \hat{g} = d\rho^2 + h(\rho)$$

where  $h(\rho) = h$  is flat and independent of  $\rho$  on the cusps.

Let  $\hat{S}$  be a orthonormal frame for  $\hat{g}$  which is parallel in the  $y$  direction in the cusp. Then both the connection 1-form  $\hat{\omega}$  and the curvature form  $\hat{\Omega}$  vanish when contracted with  $\partial_y$ . It follows that the Chern-Simons form  $\text{cs}(\hat{g}, \hat{S})$  vanishes identically on the cusp, thus the  $\text{SO}(3)$  Chern-Simons invariant for  $\hat{g}$  is well-defined and moreover it coincides with the invariant of the compact manifold with boundary obtained by chopping off the cylindrical ends.

The line bundle  $\mathcal{L}$  is constructed now over the set of constant-curvature metrics on  $M$ , namely hyperbolic on the funnel ends and flat on the cusp ends. In the definition of the cocycle  $c^X(\hat{S}, a)$  notice that the second term vanishes identically on the cusp, since we work with frames  $S$  parallel in the direction of  $y$ , which implies that  $\partial_y \tilde{a} = 0$ , or in other words  $\tilde{a}$  is independent of  $y$ . The definition of the  $\text{SO}(3)$  connection is unchanged if we include now in  $M$  the flat components corresponding to the cusps. Its curvature is computed in terms of a fiberwise integral of the Pontrjagin form by following verbatim the proof of Proposition 27. However in Lemma 28 the curvature of the tori fibers vanishes, thus the cusps do not contribute to the curvature and so the curvature of  $\nabla^{\text{SO}(3)}$  is  $\frac{i}{8\pi}$  times the Weil-Petersson symplectic form of the Teichmüller space corresponding to the funnels, i.e., it does not “see” the cusps.

We define now the  $\text{SO}(3)$  invariant of the hyperbolic metric  $g$ . Using (22) with the roles of  $g, \hat{g}$  reversed and (21) we see that in the cusp, the Chern-Simons form  $\text{cs}(g, S)$  of  $g$  equals  $d\text{Tr}(\hat{\alpha} \wedge \hat{\omega})$ , where

$$\hat{\alpha}_{ij}(Y) = y^{-1} [\hat{g}(Y, \hat{S}_j) S_i(y) - \hat{g}(Y, S_i) S_j(y)].$$

Now  $\hat{\omega}_{ij}$  is constant in  $y$  in the sense that  $\mathcal{L}_{\partial_y} \hat{\omega}_{ij} = 0$ , while  $\hat{\alpha}$  is of homogeneity  $-1$ . It follows that  $\text{cs}(g, S)$  decreases like  $y^{-2}$  as  $y \rightarrow \infty$ , thus it is integrable without regularization. Moreover the form  $\text{Tr}(\hat{\alpha} \wedge \hat{\omega})$  from (22) is homogeneous in  $y$  of degree  $-1$ , hence Proposition 17 continues to hold in the setting of this appendix.

To define the  $\text{PSL}_2(\mathbb{C})$  invariant we use Proposition 12. We note that the volume of the cusps is finite, the  $\text{SO}(3)$  Chern-Simons form was proved above to be integrable in the cusp, and we claim that the remaining term  $\text{Tr}(T \wedge \omega)$  decreases in the cusp like  $y^{-1}$ . Indeed, we have seen above that  $\omega = \hat{\omega} + \hat{\alpha}$  is of homogeneity  $0$  and  $-1$ , while  $T = y^{-1} \hat{T}$  is of homogeneity  $-1$ . Therefore  $\text{CS}^{\text{PSL}_2(\mathbb{C})}$  does not involve regularization in the cusps, while Proposition 18 continues to hold. Note that the Euler characteristic of a torus is  $0$ , so it is irrelevant whether the tori closing the cusps are included or not in the formula from Proposition 18 when we allow cusps.

The variation formula for  $\text{CS}^{\text{PSL}_2(\mathbb{C})}$  (Theorem 30) continues to hold as in the case without cusps. because in (35) the cusp terms involved (other than the first one which is the connection 1-form) do not have contributions of degree 0 in  $y$ . This is obvious if one takes into account that  $\alpha$  and  $T$  are of homogeneity  $-1$ , while  $\hat{\omega}$  is of homogeneity 0. Hence the variation of the regularized volume of a hyperbolic manifold with funnels and cusps is given by Corollary 31 (and only depends on local data on the funnels).

Finally, the correspondence between hyperbolic metrics on  $X$  and conformal infinity in the funnels continues to hold in the presence of cusps [23].

These hyperbolic metrics with cusps and funnels form therefore a Lagrangian submanifold in  $T\mathcal{T}_{\mathbf{g}}$ , and their renormalized volume is a Kähler potential for the Teichmüller space corresponding to the funnels (see Corollary 41).

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