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Loop Antennas for Near-Field Multipolar-Expansion Identification: First Experimental Validations

Benjamin Vincent, Olivier Chadebec, Jean-Luc Schanen, *Senior Member, IEEE*, and Kévin Berger

Abstract—In this paper, electromagnetic compatibility is assessed using novel magnetic induction sensors based on the standard CISPR16-1 coils. The new coil shape allows the sensors to be sensitive only to specific components of the multipolar expansion, which is similar to a spatial filtering. The aim is to simplify the geometry of sensor coils introduced in the previous work by rotating the spherical-harmonic functions. This paper describes the mathematical tools required as well as the design method and concludes with the first experimental studies.

Index Terms—Electromagnetic compatibility, magnetic-field measurement, multipolar expansion, power electronics, transducers.

I. INTRODUCTION

IN ELECTROMAGNETIC compatibility studies, a compact model of the radiated magnetic field generated by power-electronic systems enables the prediction of the field surrounding the device. This model can sometimes be obtained using numerical methods, but in most cases, the systems and associated phenomena are too complex to be numerically modeled. Another approach uses measurements of the radiated magnetic field close to the device in order to identify its equivalent magnetic model and then extrapolate the field in the surrounding space. In this paper, we focus on a multipolar-expansion model of the radiated magnetic field. This model of decomposition is well known for its ability to represent magnetic sources as a combination of known standardized sources (dipole, quadrupole, octupole, . . .) [1].

It must be stressed that inaccuracies in measurement can strongly influence the determination of each term of the spherical-harmonic development. Several solutions can reduce these effects. In the case of punctual measurements, Keradec et al. [2] propose multiple measurement locations. In this work, the magnetic-field sensor is a small coil moving around the

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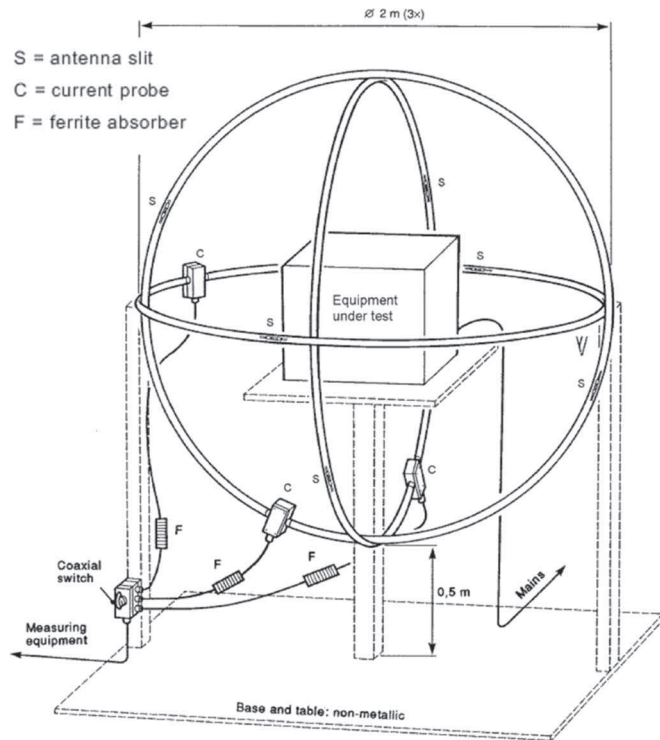


Fig. 1. From standard CISPR 16-1, the loop-antenna system, consisting of three mutually perpendicular large loop antennas.

object of study in order to measure the field at many points (up to 200). However, this measurement setup is complicated to build and can lead to inaccuracies in the positioning of the probe. An alternative is offered by Rouve et al. [3], who employ the punctual measurement but with an optimized number and positioning of fixed sensors to increase robustness of identification.

Another strategy consists in measuring the magnetic flux induced in large coils placed around the device in order to achieve a spatial integration and thus reduce positioning inaccuracies [4]. For example, to measure the first order of the decomposition (dipole) [5], the standard CISPR 16-1 [6] uses three orthogonal loops (Fig. 1), one for each dipole. The use of three identical coils with simple geometry facilitates the building of the sensor. Nevertheless, to ensure an accurate identification, the coils have to be large (at least 2 m in diameter). The resulting sensor is bulky and distant from the object of study. The magnetic field is weak, making it difficult to discriminate from electromagnetic noise. Although the sensor's measurement frequency range is limited, this standard applies for the medium frequency range: 9 kHz to 30 MHz [7].

In this paper, two improvements to this device are proposed: the identification of the quadrupolar terms (while maintaining identification of the dipolar terms) and the reduction of the sensor size. The measurement principle is similar to a spatial filtering: According to the coil shape, each sensor is sensitive to one specific component of the multipolar expansion. In a previous paper, we propose a set of coils that identify the first two orders ($n = 1$ and 2 , i.e., dipolar and quadrupolar terms) but are not sensitive to the following two orders (3 and 4). This results in a sensor with eight flux coils located on a measurement sphere (radius r_M). While the possibility of this new set of sensors has already been presented [8], we show here that this solution remains theoretical due to an infeasible measurement process. In this paper, we propose another approach. While only a few coils with a simple geometry are considered, the identification is performed by rotating these coils around the device. The measured fluxes in different positions are combined, and the equivalent multipolar terms can be determined due to a simpler measurement process. Here, this new solution is proposed and validated, and a pilot experiment is presented.

II. COMPLETE COIL SET FOR MULTIPOLAR-DECOMPOSITION IDENTIFICATION

A. Assumptions for the Electromagnetic Field

Power-electronic devices present low clock/commutation frequencies in comparison with standard electronic devices (about tens of kilohertz). Keeping in mind that the wavelength λ of 100 kHz in air is 3000 m, the near-field approximation, defined by $r < \lambda/2\pi$ (r being the distance between the device and the point where the field is expressed), can be used near and below 100 kHz. In the identification and extrapolation of the components in our application, we can assume that the electrical and the magnetic fields are quasi-static, i.e., there is no delay between sources and electromagnetic effects. In other words, in the selected frequency range, propagation effects can be neglected. Moreover, a radiated electromagnetic field is created by both electric and magnetic sources [9]. However, for low-frequency applications, it is commonly assumed that magnetic effects are dominant, and therefore, the electrical sources can be neglected. The radiated electromagnetic field is mainly created by current loops flowing inside the device. With a maximum frequency of 30 MHz, the quasi-static approximation is valid for up to 1.7 m away from the tested device.

B. Magnetic Multipolar Expansion

Multipolar expansion is a common tool used for electromagnetic-field representation [1]. Outside a sphere (radius r_V) that includes all radiation sources, the magnetic field can be fully described by its magnetic scalar potential ψ solution of the Laplace equation [9]. For each frequency, we have

$$\Psi(r, \theta, \varphi, f) = \frac{\mu_0}{4\pi} \sum_{n=1}^{+\infty} \sum_{m=-n}^{+n} A_{nm}(f) \cdot \frac{1}{r^{n+1}} \cdot S_{nm}(\theta, \varphi) \quad (1)$$

where S_{nm} is the real harmonic spherical function of order n and degree m , and (r, θ, φ) are the spherical coordinates whose origin is the center of the expansion. $A_{nm}(f)$ are the frequency-dependent unknowns to be measured. Applying the gradient operator to (1), we get the multipolar expansion of the magnetic field for a near-field area

$$\mathbf{B}(r, \theta, \varphi, f) = -\mathbf{grad}(\Psi(r, \theta, \varphi, f)) \quad (2)$$

$$\mathbf{B}_{nm} = -\frac{\mu_0}{4\pi} A_{nm}(f) \cdot \mathbf{grad} \left(\frac{1}{r^{n+1}} \cdot S_{nm}(\theta, \varphi) \right) \quad (3)$$

where the S_{nm} spherical-harmonic functions are defined by

$$S_{nm}(\theta, \varphi) = \Theta_{nm}(\theta) \cdot \Phi_m(\varphi) \quad (4)$$

with

$$\Theta_{nm}(\theta) = (-1)^{|m|} \sqrt{\frac{(n-|m|)!}{(n+|m|)!}} P_{n|m|}(\cos \theta) \quad (5)$$

$$\Phi_m(\varphi) = \begin{cases} \sqrt{2} \cos(m\varphi), & \text{if } m > 0 \\ 1, & \text{if } m = 0 \\ \sqrt{2} \sin(|m|\varphi), & \text{if } m < 0 \end{cases} \quad (6)$$

where P_{nm} are the associated Legendre polynomials.

This development is hierarchical; for elements of the n th order, the field decreases by $1/r^{n+2}$. A decomposition limited to the second order is sufficiently precise for normal power-electronic structures, beyond the distance of extrapolation. Indeed, the higher orders rapidly become negligible. Therefore, there are eight coefficients to be determined (three for the first order, i.e., the dipolar terms, and five for second order, i.e., the quadrupolar terms).

C. Initial Proposal for Component Identification

As mentioned before, the measurement principle is similar to spatial filtering: According to the coil shape, each sensor is sensitive to one specific component of the multipolar expansion. The basic principle is to measure the magnetic flux through different coils in short-circuit mode on the measurement sphere S_M (radius r_M) around the device. According to Lenz's law, the current in each coil will be directly proportional to the flux, and it is possible to calculate the flux (sum or difference) by measuring positive or negative values of the current.

In this approach, a set of coils is dedicated to only one \mathbf{B}_{nm} component and not to the others. The flux across the surface delimited by the set of coils is expressed by

$$\begin{aligned} \text{Flux}_{nm} &= \iint_{\text{Coil surface}} (\mathbf{B}_{nm}(r_M, \theta, \varphi) \cdot \mathbf{n}) \, ds \\ &= \frac{\mu_0}{4\pi} \frac{(n+1)}{r_M^{n+2}} A_{nm} \iint_{S_M \text{ portions}} S_{nm}(\theta, \varphi) \, dS \end{aligned} \quad (7)$$

where \mathbf{n} is the unit outgoing normal of S_M and Flux_{nm} corresponds to the sum of fluxes through the sensor coils.

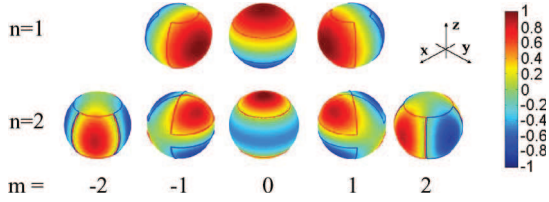


Fig. 2. Eight coil sensors for the identification of first- and second-order components. The flux is counted as positive for the red coils and negative for the blue ones. Functions S_{nm} are presented using projection on the measurement sphere.

Equation (7) shows that such a coil configuration will enable the direct determination of the A_{nm} coefficient, which is more robust than any indirect calculation. Moreover, we can see that the field distribution on a sphere S_M is directly linked to the characteristics of the functions S_{nm} . Their study is therefore sufficient for determining the shape and configuration of the set of coils.

Using the proposed method, it is not possible to identify one component while filtering an infinite number of parasitic terms of the multipolar expansion. Thus, we consider only the first four orders of this expansion. The higher orders (above the fifth order) are considered negligible. Indeed, their influence decreases by $1/r^7$, and the coils are located at two times r_V . Also, these types of power-electronic sources do not normally contain higher order components.

Using the initial system presented by Bergervoet and Van Veen [5], simple algebra shows that this kind of device dedicated to the determination of dipolar terms is influenced by octupolar components ($n = 3$) that decrease by $1/r^5$. This destabilizing component can lead to measurement inaccuracies. In fact, if a large device is tested, the flux would be created by the tracked dipole and this additional term. The solution proposed by Bergervoet and Van Veen [5] is simply to increase the size of the loop (or the sphere) up to 2 m in diameter in order to reduce the octupolar contribution. However, this set of loops becomes very sensitive to external magnetic fields since the measurement surface area is so high.

In a new approach described in a previous paper [8], we proposed a system of coils to identify the first and second orders only. The theory was based on a careful study of the mathematical properties of spherical-harmonic functions. The resulting sensor had eight flux coils on a measurement sphere. To simplify the coil-design process, the shape of the sensor coils would have constant θ or φ angles, i.e., these coils would follow only the meridians or parallels of the measurement sphere. This design constraint is suggested by (4), where the spherical-harmonic function S_{nm} depends on the θ variable and on the φ variable. These different sensor shapes are shown in Fig. 2. The proposed solution lead to a much more compact system associated with a theoretical improvement in accuracy. Moreover, with the size of the loop being reduced, the propagation effect would not disturb the measurement. The proposed system would exhibit a higher frequency range than the large coils used by Bergervoet and Van Veen [5].

Unfortunately, this measurement setup was difficult to build. The coils were not located on the plane surfaces and were too numerous to construct. Even for a simple dipole identification,

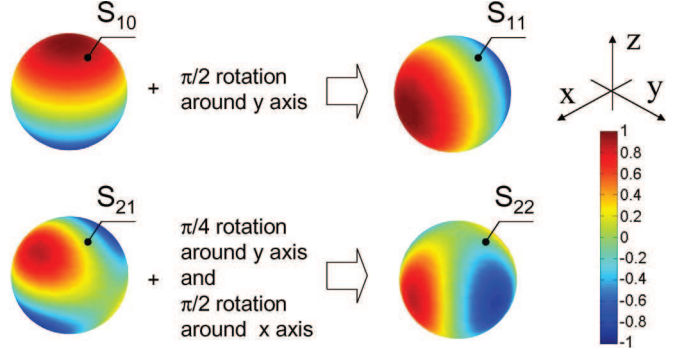


Fig. 3. Examples of rotation properties of spherical-harmonic functions.

the solution proposed for $n = 1$ remained quite difficult. A simplification of this theoretical setup was necessary.

III. NEW APPROACH BASED ON ROTATING SENSORS

A. Basic Idea

The simplification of the measurement process has been presented in [10]. The aim was to simplify the previous solution (presented in [8]) in order to build a practical sensor. The basic idea comes from the properties of spherical-harmonic functions. With the rotations shown in Fig. 3, it is possible to generate all components of the same order n with only one component. The first-order rotation dependence is obvious. For the $n = 1$ case (Fig. 2), even if the associated spherical-harmonic functions are very similar, the sensor shapes obtained present markedly different geometries, although, obviously, the dipole basis must have symmetrical behavior along the x -, y -, and z -axes. A solution with a single sensor shape associated with two rotations can lead to exactly the same identification. This solution has already been mentioned by Pasmooij [11], where a new shape could determine equivalent dipole terms. It corresponds to our S_{10} sensors (equivalent to Helmholtz coils) associated with the two $\pi/2$ rotations. This sensor presents an improvement over the system by Bergervoet and Van Veen [5] by ensuring the spatial filtering of the $n = 3$ component.

The basic idea is that it is possible to generalize this property to design all sensors. In Fig. 2, two sensors seem particularly easy to build: the coils of the A_{10} and A_{20} sensors. Each of them is composed of two circular coils. In the following paragraphs, we will focus on these simple shapes.

B. Example of A_{21} Sensor Design

The A_{21} sensor is associated with the complex coil topology shown in Fig. 2. Using the ideas described earlier, it is possible to summarize the fluxes measured by the A_{20} sensor oriented along different geometric axes to get the equivalent fluxes measured by the A_{21} sensor. The geometrical rotation properties are detailed in [12]. The orthogonal properties of the spherical-harmonic base show that, if one component sensor is rotated, only components of the same order n could create a parasitic flux through the coils. Therefore, the system of loops for detecting the A_{21} component maintains the filtering capability for the other orders. The algebra is quite complicated

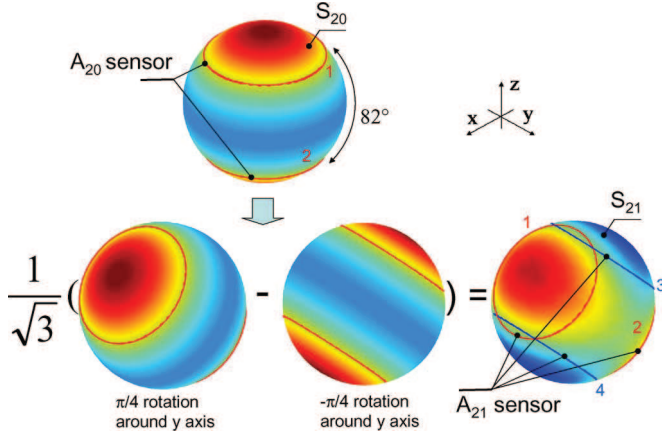


Fig. 4. Identification of A_{21} parameter using A_{20} sensor topology and associated rotation and flux combinations.

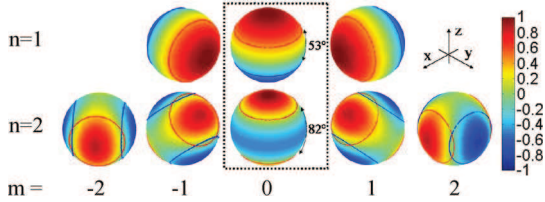


Fig. 5. Four simple circular loops associated with rotation to identify all the $n = 1$ and $n = 2$ components.

and will not be developed in this paper, but an illustration of this example is shown in Fig. 4. The S_{21} function is generated from the rotation and composition of S_{20}

$$S_{21} = \frac{1}{\sqrt{3}} \left(S_{20}^{\pi/4} - S_{20}^{-\pi/4} \right) \quad (8)$$

where $S_{20}^{\pi/4}$ represents the S_{20} function rotated by an angle of $\pi/4$ around the y -axis. The A_{21} sensor is made by using the A_{20} sensors associated with two rotations and a composition, as shown by (8).

C. Final Topology

Finally, we present here the simplification of the shapes obtained (Fig. 5). There are only four circular loops (two for the dipole identification and two for the quadrupolar one). This new set of loops provides an accurate identification of A_{nm} , like the first sensor shown in Fig. 2. The relationship between measured fluxes and A_{nm} components is given by

$$\begin{cases} A_{1m} = \frac{10^8 r_M}{32\pi} \text{Flux}_{1m} \\ A_{20} = \frac{6125 \cdot 10^4 r_M^2}{3\pi\sqrt{21}} \text{Flux}_{20} \\ A_{2m} = \frac{6125 \cdot 10^4 \cdot r_M^2}{3\pi\sqrt{21}} \left(\frac{1}{\sqrt{3}} \right) \text{Flux}_{2m}, \quad \text{for } m \neq 0 \end{cases} \quad (9)$$

where Flux_{nm} corresponds to the total flux through the sensor coils. For example, for the A_{21} sensor, according to Fig. 4, we have

$$\text{Flux}_{21} = \text{flux}_1 + \text{flux}_2 - \text{flux}_3 - \text{flux}_4 \quad (10)$$

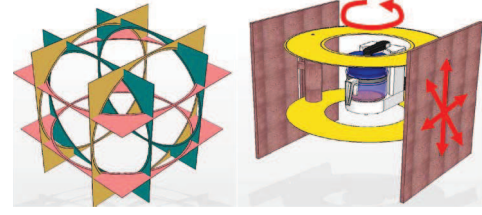


Fig. 6. Two suggestions for construction of the sensor. On the left is a static solution enabling the determination of the three dipole components. Construction is possible using three different radii of the measurement sphere. The construction on the right enables the identification of dipolar and quadrupolar terms (coils have three different axes, and device has two around the z -axis). The two A_{10} and A_{20} sensors have been located on the same plane to simplify the geometry, which changes the measurement radius.

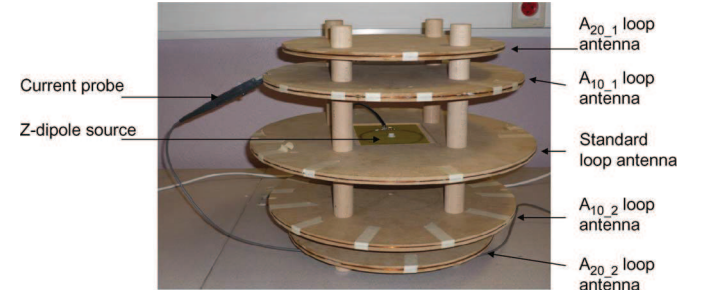


Fig. 7. Prototype photograph presenting the whole system of loop antennas (“ A_{10} ,” “ A_{20} ,” and “standard”).

where flux_i corresponds to the flux measured across the coil i in Fig. 4.

In Fig. 5, the complexity of the system is significantly reduced for the coils, but it must be pointed out that the mechanical management of rotation can lead to new difficulties. Some ideas for their construction are shown in Fig. 6.

IV. PILOT EXPERIMENTS

In this pilot experiment, we built a prototype that includes only the two A_{10} and A_{20} sensors. They correspond to the z -components of the dipole and the quadrupole (black frame in Fig. 5). This paper will illustrate the improvement of the filtering capabilities and robustness over the standard system without addressing the construction problems due to the three axes. The prototype is shown in Fig. 7. All coils are designed around a measurement sphere of radius $r_M = 0.225$ m.

A. Prototype Presentation

The first issue addressed when making the prototype was the choice of the field-measurement technique. There are two solutions: placing the coils in short-circuit mode and measuring the induced current or opening the coil and measuring the voltage. We choose the first alternative, allowing a nearly flat response in the frequency range (9 kHz–30 MHz) as explained by [13]. However, this solution leads to a new measurement difficulty. The mutual inductance between the coil sensors cannot be neglected, so the measurement must be achieved using one coil at a time, while the others remain in an open-circuit position. The current is measured using the usual current sensors (P6022, Tektronix) connected to a numerical oscilloscope.

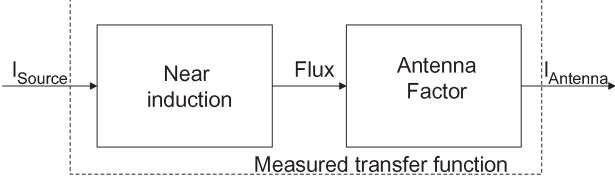


Fig. 8. Calibration process of each coil.

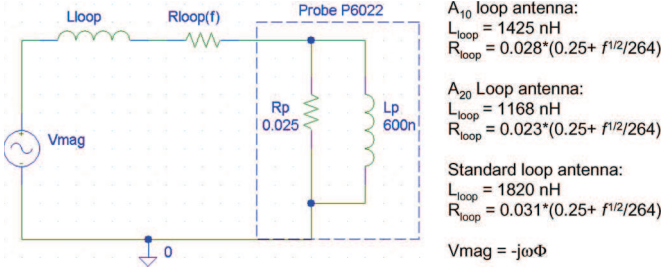


Fig. 9. Loop-antenna equivalent circuit.

In addition to the four-coil system, a simple loop antenna located in the middle plan has been added to demonstrate the robustness of this new approach. It corresponds to the initial z -dipole sensor proposed in [5].

With this configuration, the identified coefficients A_{nm} are given using (9). For the standard sensor case with z -dipole source, we have, from (7)

$$\begin{aligned}
 \text{Flux}_{\text{Standard}} &= \iint_{\text{Half sphere}} (\mathbf{B}_{10}(r_M, \theta, \varphi) \cdot \mathbf{n}) ds \\
 &= \frac{\mu_0}{4\pi} \frac{(n+1)}{r_M^{n+2}} A_{10} \int_{\theta=0}^{\pi/2} \int_{\varphi=0}^{2\pi} S_{10}(\theta, \varphi) r_M^2 \sin(\theta) d\theta d\varphi \quad (11)
 \end{aligned}$$

where \mathbf{n} is the unit of outgoing normal of the measurement sphere and $\text{Flux}_{\text{Standard}}$ corresponds to the measured flux through the standard sensor coil. Therefore, after calculation, we get

$$A_{10} = 3.58 \cdot 10^5 \text{Flux}_{\text{Standard}}. \quad (12)$$

B. Prototype Calibration

The relationship between the flux and the measured current is needed to obtain the A_{nm} coefficients (equivalent to the antenna factor). The calibration process is shown in Fig. 8.

To determine the antenna factor, two approaches are proposed. The first one consists in modeling the antenna factor using an RL equivalent circuit, with the component being calculated analytically, shown in Fig. 9. The R_{loop} equation takes into account the skin effect in the conductor. As mentioned before in Section II-A, this simple magnetic modeling of the loop is sufficient to characterize its behavior.

The second approach is to introduce a well-characterized magnetic source into our device. We have chosen a circular current loop (the value of the dipole is given by the current

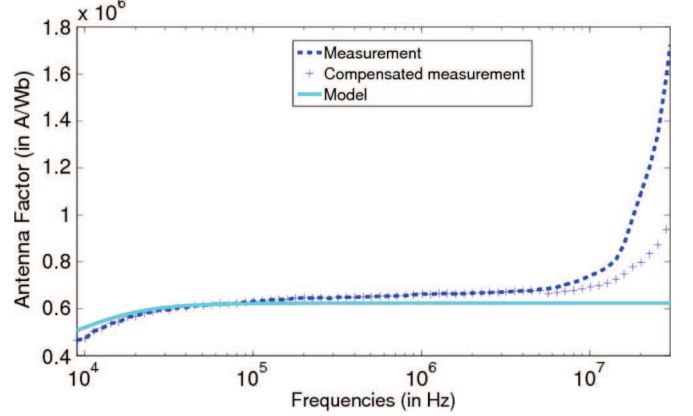
Fig. 10. A_{10_1} loop-antenna factor from modeling circuit and calibration measurement.

TABLE I
IDENTIFICATION RESULTS WITH TWO DIFFERENT
SOURCES PLACED AT THE CENTER

	z-dipole source (mA.m ²)		z-quadrupole source (mA.m ³)	
	Calculated	Measured	Calculated	Measured
A_{10} sensor	7.854	7.955	/	-0.025
Standard sensor	7.854	7.854	/	0.030
A_{20} sensor	/	$9.2 \cdot 10^{-4}$	0.732	0.753

multiplied by its surface area). The calibration process must be performed for each coil. The comparison between both approaches is shown in Fig. 10. They are highly consistent in the range from 10 kHz to 10 MHz. The difference in high frequencies comes from interferences between the external strong magnetic field and the current measurement through the probe shield. In this case, a good solution would be to make a differential measurement, according to the instruction manual of the probe P6022. In Fig. 10, the antenna factor is obtained using a clamped probe around the conductor of the loop antenna. The compensated antenna factor is obtained by subtracting the signal measured with a probe placed as close as possible to the first probe and without the conductor inside. Even if the interferences are fairly reduced, a difference in modeling behavior persists. The following section describes an experiment performed at 1 MHz and with the measured antenna factor.

C. Measurement Results

1) *Filtering Capabilities*: The validation of the z -dipole and z -quadrupole source identifications is summarized in Table I. These sources are centered, and their effects on the three sensors are measured one after the other. In all cases, the consistent results validate the A_{10} and A_{20} theoretical values. Thus, when the source is composed of a z -dipole, the result from the A_{20} sensor measurement is very small. As with the standard sensor, the proposed solution exhibits good filtering of the second order. This filtering capability is less significant for the z -quadrupole case. Indeed, current intensities are close to the sensitivity limits of the oscilloscope, so the accuracy is reduced.

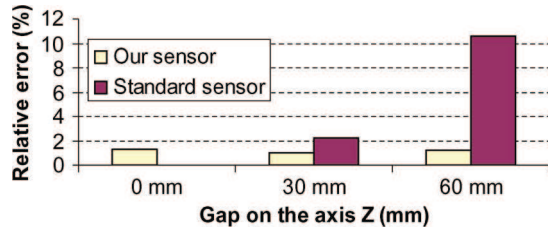


Fig. 11. Comparison of the relative error in identifications made by the standard sensor and our sensor, with respect to the theoretical z -dipole for different off-center placements of the source.

It should be noted in Table I that the result of the calibration using the z -dipole is fairly good, with only a 1.3% error.

The accuracy of the proposed sensor can be evaluated on the quadrupole identification results: less than 3% error, which is quite a good result.

However, the ideal case presented previously is far from likely source conditions, where the sources could be off-center and/or of a higher order than a pure dipole. This will be studied in the following section.

2) *Robustness*: The robustness of the proposed sensor is compared with the standard sensor by shifting the source loop away from the center of the sensors. Only the dipole identification is observed in this experiment. The measurement results in Fig. 11 show that an error occurs in the dipole determination with the standard sensor. However, despite the shifting of the source, the identification error remains the same with our sensor. With the standard sensor, the error increases up to 11%. A decentered dipole implies a more complex field (i.e., higher orders in the decomposition). Since the filtering capability of the standard sensor is lower than our proposed sensor, the contribution of the higher orders to the measured signal corrupts the identification. The error is mainly attributed to the contribution of the third-order multipolar induction term in the standard sensor. The filtering capabilities of our sensors minimize this problem. The error is reasonable and lower than using the standard antenna since the latter is sensitive to third-order terms.

V. CONCLUSION

This paper has presented a new measurement system allowing the identification of an equivalent source for any radiating device. It provides the coefficients of the multipolar expansion, up to the quadrupole order. In comparison with previous solutions, the advantages are as follows.

- 1) Higher order determination, more robust behavior, and more compact system than the standard system using three simple loops.
- 2) Simple shapes used for construction, by rotating the simple sensors A_{10} and A_{20} to provide the other components. The complex shapes proposed in the previous work were efficient to filter the other orders but were too complicated to be built.

The new system has been detailed and validated using measurements on a single z -axis. The comparison with the standard sensor has shown increased robustness against shifts of a dipolar source, as expected.

Further work will focus on building the triple-axis system, improving the current measurement, and applying it to a complex radiating source, like a power-electronic converter.

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