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▶ To cite this version:

Aurélia Fraysse. Regularity criteria of almost every function in a Sobolev space. Journal of Functional Analysis, 2010, pp.1806-1821. 10.1016/j.jfa.2009.11.017. hal-00555299v2

HAL Id: hal-00555299 https://hal.science/hal-00555299v2

Submitted on 20 Apr 2012

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Regularity criteria for almost every function in Sobolev spaces

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April 20, 2012

Abstract

In this paper we determine the multifractal nature of almost every function (in the prevalence setting) in a given Sobolev or Besov space according to different regularity exponents. These regularity criteria are based on local L^p regularity or on wavelet coefficients and give a precise information on pointwise behavior.

1 Introduction

The study of regularity, and more precisely of pointwise regularity of signals or functions raised a large amount of interest in scientific communities. This topic allows a better understanding of behavior of functions and it gives also a powerful classification tool in various domains. A recent theory, based on the study of pointwise smoothness is supplied by the multifractal analysis. The multifractal analysis was initially introduced in order to study the velocity of turbulent flows. It was then applied in several fields, such as in signal or image processing [1, 2]. But in each case, the criterium of regularity taken into account is the Hölder exponent, and this exponent is only well defined for locally bounded functions. It would be convenient to define new criteria on more general cases. This is very important for applications in several fields. For instance, the velocity of turbulent fluids is now known not

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to be bounded near vorticity filaments, see [3]. Concerning turbulent flows, in [24], Leray shows that weak solutions of Navier-Stokes equations with initial value in $L^2(\mathbb{R}^3)$ may develop singularities in finite time. This problem was then widely studied, [23, 29] and different behaviors were produced following the initial value problem involved. Furthermore, in [9], it was shown that an alternative definition of regularity can give better results in elliptic PDEs. This regularity criteria is also taken in [8] where it provides a nice tool to study viscosity solutions. It would thus be natural to take this notion, which involves local L^p norms to study irregularities of Navier-Stokes solutions when initial data are supposed in L^p .

Furthermore, Hölder regularity is not adapted for classification of natural images as it does not take into account the geometry of sets. It would however be convenient to establish these criteria in image processing, where regularity properties are widely used. A natural idea would be to determine properties of the characteristic function of sets. Nonetheless, the Hölder regularity of the characteristic function of any set only takes two values. Furthermore, most natural images, such as clouds images or medical images are discontinuous, see [3] and thus need to be studied with more general notion of regularity.

Let us recall the principle of multifractal analysis. The natural notion of regularity used in the study of pointwise behavior is provided by the Hölder exponent, defined as follows.

Definition 1. Let $\alpha \geq 0$; a function $f : \mathbb{R}^d \to \mathbb{R}$ is $C^{\alpha}(x_0)$ if for each $x \in \mathbb{R}^d$ such that $|x - x_0| \leq 1$ there exists a polynomial P of degree less than $[\alpha]$ and a constant C such that,

$$|f(x) - P(x - x_0)| \le C|x - x_0|^{\alpha}.$$
 (1)

The Hölder exponent of f at x_0 is

$$h_f(x_0) = \sup\{\alpha : f \in C^{\alpha}(x_0)\}.$$

In some cases, functions may have an Hölder regularity which changes wildly from point to point. Rather than measure the exact value of the Hölder exponent, one studies the fractal dimension of sets where it takes a given value. The *spectrum of singularities*, also called multifractal spectrum and denoted d(H), is the function which gives for each H the Hausdorff dimension of those sets. A function is then

called *multifractal* if the support of its spectrum of singularities is an interval with no empty interior.

However, the Hölder exponent has some drawbacks that prevent from using it in any situation. First, it is only defined for locally bounded functions. If a function f belongs only to L_{loc}^p this exponent is no more defined. Furthermore, as pointed by Calderòn and Zygmund in [9], it is not preserved under pseudodifferential operator of order zero, and as stated in [27] cannot thus be characterized with conditions on wavelet coefficients.

Another drawback can be emphasized with the example of Raleigh-Taylor instability. This phenomena occurs when two fluids which are not miscible are placed on top of each other. In this case, thin filaments appears giving to the interface between the two fluids a fractal structure, see [28] for a study. To study geometric properties of this interface, one would be interested on multifractal properties of its characteristic function. Nonetheless as such functions are not continuous and take only two values, their Hölder exponent is not define, and a multifractal approach can't be carry out.

For all this reasons, it would be convenient to define a new kind of multifractal analysis constructed with more general exponents. Such construction is started in [20, 21], where the authors proposed a multifractal formalism based on Calderòn-Zygmund exponents. These exponents were introduced in [9] as an extension of Hölder exponent to L_{loc}^p functions, invariant under pseudodifferential operator of order 0.

Definition 2. Let $p \in [1, \infty]$ and $u \ge -\frac{d}{p}$ be fixed. A function $f \in L^p_{loc}(\mathbb{R}^d)$ belongs to $T^p_u(x_0)$ if there exist a real R > 0 and a polynomial P, such that $deg(P) < u + \frac{d}{p}$, and c > 0 such that:

$$\forall \rho \le R: \quad \left(\frac{1}{\rho^d} \int_{\|x-x_0\| \le \rho} |f(x) - P(x)|^p dx\right)^{1/p} \le c\rho^u. \tag{2}$$

The p-exponent of f at x_0 is $u_f^p(x_0) = \sup\{u: f \in T_u^p(x_0)\}.$

With this definition, the usual Hölder condition $f \in C^s(x_0)$ corresponds to $f \in T_u^p(x_0)$ where $p = \infty$. One can also check that the p exponent is decreasing as a function of p. As it was done for the Hölder exponent one can define for each p the p-spectrum of singularities as the Hausdorff dimension of the set of points where the p-exponent take

a given value. In [20], the authors defined the weak accessibility exponent, given as a parameter of the geometry of the set. Specifically, this weak-scaling exponent deals with the local behavior of the boundary of a set. It is thus well adapted for fractal interfaces that might appear in chemical experimentation. They showed that this geometrical based exponent coincide with Calderòn-Zygmund exponents of the characteristic function of the boundary of the set.

Another regularity criterium, closely related to the previous ones is given by the following definition from [27]. With this exponent we can have a better understanding of the link between Calderòn-Zygmund exponents, Hölder exponent and the pointwise behavior of functions.

Definition 3. Let $f : \mathbb{R}^d \to \mathbb{R}$ be a function or a distribution and $x_0 \in \mathbb{R}^d$ be fixed. The weak scaling exponent of f at x_0 is the smallest real number $\beta(f, x_0)$ satisfying:

1.
$$\beta(f, x_0) \ge u_f^p(x_0) \quad \forall p \ge 1.$$

2.
$$\beta(f, x_0) = s \Leftrightarrow \beta\left(\frac{\partial f}{\partial x_j}, x_0\right) = s - 1 \quad j = 1, ..., d.$$

Similarly, we define the weak-scaling spectrum, denoted by $d_{ws}(\beta)$ as the Hausdorff dimension of sets of points where $\beta(f,x)$ takes a given value β . As we will see later, the weak scaling exponent can be characterized by conditions on wavelet coefficients. In practical applications, the spectrum of singularities cannot be computed directly, as it takes into account intricate limits. Thus, some some formula, the multifractal formalism were introduced which link the spectrum of singularities to some calculable quantities. There are indeed two formalisms based on conditions on wavelet coefficients. Historically the classical multifractal formalism stated in [12] was based directly on wavelet coefficients. The common point of view is that it is the weak scaling exponent which is involved in this formula. The second multifractal formalism, developed in [19] is based on "wavelet leaders", which can be seen as the theoretical counterpart of the "Wavelet Transform Modulus Maxima" used in [4]. This "wavelet leader" based formalism actually gives the spectrum of singularities in term of Hölder exponent and is numerically more stable. Furthermore it gives an additional information on the behavior of functions thanks to the following definition.

Definition 4. Let $f : \mathbb{R}^d \to \mathbb{R}$ be a function and $x_0 \in \mathbb{R}^d$. We say that x_0 is a cusp singularity for f if $\beta(f, x_0) = h_f(x_0)$. If $\beta(f, x_0) > h_f(x_0)$, x_0 is said to be an oscillating singularity.

An example of oscillating function at $x_0 = 0$ is given by $f(x) = |x|\sin(1/|x|)$. Here, $h_f(0) = 1$ while $\beta(f,0) = +\infty$. And we have a cusp singularity when the behavior of the function at x_0 is like $|x|^{\alpha}$ but also like $|x|^{\alpha} + |x|\sin(1/|x|)$. Indeed we talk about a cusp singularity when the function do not have oscillations at a point, or if those oscillations are hidden by the Hölder behaviour.

Many authors have studied generic values of the Hölder exponent in function spaces. In 1931 Banach, [5], proved that the pointwise regularity of quasi all, in a topological sense, continuous functions is zero. Here quasi all means that this property is true in a countable intersection of dense open sets. Since then Hunt in [13] showed that the same result is satisfied by measure theoretic almost every continuous functions. Recently, results such as those of [22] and [11] studied Hölder regularity of generic functions in Sobolev spaces in both senses. Whereas a large study of regularity properties for generic sets, there exists no result on genericity of Calderòn-Zygmund exponents or of weak scaling exponent. Our purpose here is to provide a genericity result of those exponents in given Sobolev and Besov spaces, with the measure-theoretic notion of genericity supplied by prevalence.

Prevalence is a measure theoretic notion of genericity on infinite dimensional spaces. In a finite dimensional space, the notion of genericity in a measure theoretic sense is supplied by the Lebesgue measure. The particular role played by this measure is justified by the fact that this is the only one which is σ -finite and invariant under translation. In a metric infinite dimensional space no measure enjoys this properties. The proposed alternative is to replace conditions on the measure by conditions on sets, see [6, 10, 14] and to take the following definition.

Definition 5. Let V be a complete metric vector space. A Borel set B in V is called Haar-null if there exists a probability measure μ with compact support such that

$$\mu(B+v) = 0 \quad \forall v \in V. \tag{3}$$

In this case the measure μ is said transverse to B. A subset of V is called Haar-null if it is contained in a Borel Haar-null set.

The complement of a Haar-null set is called a prevalent set.

With a slight abuse of language we will say that a property is satisfied almost everywhere when it holds on a prevalent set.

Let us recall properties of Haar-null sets, see [10, 14] and show how it generalizes notion of Lebesgue measure zero sets.

Proposition 1. 1. If S is Haar-null, then $\forall x \in V$, x + S is Haar-null.

- 2. If $\dim(V) < \infty$, S is Haar-null if and only if $\operatorname{meas}(S) = 0$ (where meas denotes the Lebesgue measure).
- 3. Prevalent sets are dense.
- 4. If S is Haar-null and $S' \subset S$ then S' is Haar null.
- 5. The union of a countable collection of Haar-null sets is Haar null.
- 6. If $\dim(V) = \infty$, compact subsets of V are Haar-null.

Remarks. Several kinds of measures can be used as transverse measures for a Borel set. Let us give two examples of transverse measure.

1. A finite dimensional space P is called a probe for a set $T \subset V$ if the Lebesgue measure on P is transverse to the complement of T.

Those measures are not compactly supported probability measures. However one immediately checks that this notion can also be defined the same way but stated with the Lebesgue measure defined on the unit ball of P. Note that in this case, the support of the measure is included in the unit ball of a finite dimensional subspace. The compactness assumption is therefore fulfilled.

2. If V is a function space, a probability measure on V can be defined by a random process X_t whose sample paths are almost surely in V. The condition $\mu(f+A)=0$ means that the event $X_t-f\in A$ has probability zero. Therefore, a way to check that a property $\mathcal P$ holds only on a Haar-null set is to exhibit a random process X_t whose sample paths are in V and is such that

 $\forall f \in V, \ a.s. \ X_t + f \ does \ not \ satisfy \ \mathcal{P}.$

1.1 Statement of main results

The purpose of this paper is stated by the two following theorems which give the multifractal properties of almost every functions with regard to exponents defined in the previous section.

Theorem 1. Let $s_0 \ge 0$ and $1 \le p_0 < \infty$ be fixed.

1. For all $p \ge 1$ such that $s_0 - \frac{d}{p_0} > -\frac{d}{p}$ the p-spectrum of singularities of almost every function in $L^{s_0,p_0}(\mathbb{R}^d)$ is given by

$$\forall u \in \left[s_0 - \frac{d}{p_0}, s_0 \right] \quad d_p(u) = p_0(u - s_0) + d. \tag{4}$$

2. For almost every function in $L^{s_0,q_0}(\mathbb{R}^d)$ the spectrum of singularities for the weak-scaling exponent is given by

$$\forall \beta \in [s_0 - \frac{d}{p_0}, s_0] \quad d_{ws}(\beta) = p_0(\beta - s_0) + d.$$
 (5)

This result in Sobolev spaces has an analogous in the Besov setting. Furthermore, Besov spaces are useful when wavelets are involved as it is the case here, those spaces having a simpler characterization.

Theorem 2. Let $s_0 \ge 0$ and $0 < q, p_0 < \infty$ be fixed.

1. For all $p \ge 1$ such that $s_0 - \frac{d}{p_0} > -\frac{d}{p}$ the p-spectrum of singularities of almost every function in $B_{p_0}^{s_0,q}(\mathbb{R}^d)$ is given by

$$\forall u \in \left[s_0 - \frac{d}{p_0}, s_0 \right] \quad d_p(u) = p_0(u - s_0) + d. \tag{6}$$

2. For almost every function in $B_{p_0}^{s_0,q}(\mathbb{R}^d)$ the spectrum of singularities for the weak-scaling exponent is given by

$$\forall \beta \in [s_0 - \frac{d}{p_0}, s_0] \quad d_{ws}(\beta) = p_0(\beta - s_0) + d. \tag{7}$$

These theorems seem a bit surprising. Let us compare them with the following proposition from [11].

Proposition 2. • If $s - d/p \le 0$, then almost every function in $L^{p,s}$ is nowhere locally bounded, and therefore its spectrum of singularities is not defined.

• If s - d/p > 0, then the Hölder exponent of almost every function f of $L^{p,s}$ takes values in [s - d/p, s] and

$$\forall H \in [s - d/p, s], \qquad d_f(H) = Hp - sp + d; \tag{8}$$

Thus the main change from [11] is given by the fact that here H can take negative values. Indeed, our present theorems give a generic regularity in Sobolev or in Besov spaces that are not imbedded in global Hölder spaces. There we have nevertheless an idea of the pointwise behavior of almost every distribution in such spaces. On the other case, when $s_0 - \frac{d}{p_0} > 0$ and the spectrum of singularities exists, it coincides with the above spectra for almost every function in Besov spaces.

In [25], it was proved that in those spaces quasi all functions have not oscillating singularities. Furthermore, presence of oscillating singularities is linked with the failure of the multifractal formalism in [30]. And in [?], it was already proven that almost every function in Besov spaces satisfy the multifractal formalism. In the present paper, we see that in a prevalent setting, we have both the multifractal formalism, and the absence of oscillating singularities.

Another remark can be made thanks to the following proposition from [21] and from [31] that give an upper bound for the p-spectrum.

Proposition 3. Let $f \in B_{p_0}^{s_0,p_0}(\mathbb{R}^d)$, where $s_0 > 0$ and let $p \ge 1$ be such that $s_0 - \frac{d}{p_0} > -\frac{d}{p}$. Then

$$\forall u \in \left[s_0 - \frac{d}{p_0}, s_0 \right] \quad d_p(u) \le p_0 u - s_0 p_0 + d. \tag{9}$$

This proposition with Theorem 2 show that the generic regularity for p criteria is as bad as possible.

In Part 2 we will prove Theorems 1 and 2. For the sake of completeness, we first have to define our main tool which is given by wavelet expansions of functions. Wavelets are naturally present in multifractal analysis, see for instance [2]. Furthermore, in our case it allows a characterization of both functional spaces and pointwise regularities.

1.2 Wavelet expansions

There exist 2^d-1 oscillating functions $(\psi^{(i)})_{i\in\{1,\dots,2^d-1\}}$ in the Schwartz class such that the functions

$$2^{dj}\psi^{(i)}(2^jx-k), \quad j \in \mathbb{Z}, \ k \in \mathbb{Z}^d$$

form an orthonormal basis of $L^2(\mathbb{R}^d)$, see [26]. Wavelets are indexed by dyadic cubes $\lambda = \left[\frac{k}{2^j}; \frac{k+1}{2^j}\right]^d$. Thus, any function $f \in L^2(\mathbb{R}^d)$ can be written:

$$f(x) = \sum_{i,k} c_{j,k}^{(i)} \psi^{(i)}(2^{j}x - k)$$

where

$$c_{j,k}^{(i)} = 2^{dj} \int f(x)\psi^{(i)}(2^j x - k)dx.$$

(Note that we use an L^{∞} normalization instead of an L^2 one, which simplifies the formulas). If p > 1 and s > 0, Sobolev space have thus the following characterization, see [26]:

$$f \in L^{s,p}(\mathbb{R}^d) \Leftrightarrow \left(\sum_{\lambda \in \Lambda} |c_{\lambda}|^2 (1 + 4^{js}) \chi_{\lambda}(x)\right)^{1/2} \in L^p(\mathbb{R}^d),$$
 (10)

where $\chi_{\lambda}(x)$ denotes the characteristic function of the cube λ and Λ is the set of all dyadics cubes. Homogeneous Besov spaces, which will also be considered, are characterized (for p, q > 0 and $s \in \mathbb{R}$) by

$$f \in B_p^{s,q}(\mathbb{R}^d) \iff \sum_j \left(\sum_{\lambda \in \Lambda_j} |c_\lambda|^p 2^{(sp-d)j} \right)^{q/p} \le C$$
 (11)

where Λ_j denotes the set of dyadics cubes at scale j, see [26].

Hölder pointwise regularity can also be expressed in term of wavelet coefficients, see [15].

Proposition 4. Let x be in \mathbb{R}^d . If f is in $C^{\alpha}(x)$ then there exists c > 0 such that for each λ :

$$|c_{\lambda}| \le c2^{-\alpha j} (1 + |2^{j}x - k|)^{\alpha}.$$
 (12)

This proposition is not a characterization. If for any $\varepsilon > 0$, a function does not belongs to $C^{\varepsilon}(\mathbb{R}^d)$ one cannot express its pointwise Hölder regularity in term of condition on wavelet coefficients. This

is an advantage of Calderòn Zygmund exponent since, as showed in [18], it can be linked to wavelet expansion without global regularity assumption.

Definition 6. Let x_0 be in \mathbb{R}^d and $j \geq 0$. We denote by $\lambda_j(x_0)$ the unique dyadic cube of width 2^{-j} which contains x_0 . And we denote

$$3\lambda_j(x_0) = \lambda_j(x_0) + \left[-\frac{1}{2^j}, \frac{1}{2^j} \right]^d.$$

Furthermore, we define the local square function by

$$S_f(j, x_0)(x) = \left(\sum_{\lambda \subset 3\lambda_j(x_0)} |c_\lambda|^2 \mathbf{1}_\lambda(x)\right)^{1/2}.$$

Proposition 5. Let $p \ge 1$ and $s \ge 0$; if $f \in T^p_{s-\frac{d}{p}}(x_0)$, then $\exists C > 0$ such that wavelet coefficients of f satisfy for all $j \ge 0$

$$||S_f(j,x_0)||_{L^p} \le c2^{-j(u+d/p)}.$$
 (13)

Conversely if (13) holds and if $s - \frac{d}{p} \notin \mathbb{N}$ then $f \in T^p_{s - \frac{d}{2}}(x_0)$.

As far as we are concerned, we don't need a characterization but a weaker condition which is given by the following proposition from [20].

Proposition 6. Let $p \ge 1$ and $s \ge 0$; if $f \in T^p_{s-\frac{d}{p}}(x_0)$, then $\exists A, C > 0$ such that wavelet coefficients of f satisfy

$$\exists C \ \forall j \quad 2^{j(sp-d)} \sum_{|k-2^j x_0| < A2^j} |c_{j,k}|^p (1+|k-2^j x_0|)^{-sp} \le Cj. \quad (14)$$

Furthermore, it is also proved in [20] that the p-exponent can be derived from wavelet coefficients.

Proposition 7. Let $p \ge 1$ and $f \in L^p_{loc}$. Define

$$\Sigma_{j}^{p}(s,A) = 2^{j(sp-d)} \sum_{|k-2^{j}x_{0}| \le A2^{j}} |c_{j,k}|^{p} (1+|k-2^{j}x_{0}|)^{-sp}, \qquad (15)$$

for A > 0 small enough. And denote

$$i_p(x_0) = \sup \left\{ s : \liminf \frac{\log \left(\sum_{j=0}^{p} (s, A)^{1/p} \right)}{-j \log 2} \ge 0 \right\}.$$
 (16)

Then the following inequality always holds

$$u_f^p(x_0) \le i_p(x_0) - \frac{d}{p}.$$
 (17)

If furthermore there exists $\delta > 0$ such that $f \in B_p^{\delta,p}$ then the p-exponent of f satisfies

$$u_f^p(x_0) = i_p(x_0) - \frac{d}{p}.$$
 (18)

As seen previously, the *p*-exponent is also related to the weak-scaling exponent. This one can also be expressed in term of wavelet coefficients, thanks to its relation with two-microlocal spaces, defined in [7].

Definition 7. Let s and s' be two real numbers. A distribution $f: \mathbb{R}^d \to \mathbb{R}$ belongs to the two-microlocal space $C^{s,s'}(x_0)$ if its wavelet coefficients satisfy that there exists c > 0 such that

$$\forall j, k \quad |c_{j,k}| \le c2^{-sj}(1 + 2^j x_0 - k|)^{-s'}.$$
 (19)

In [27] the following characterization of the weak scaling exponent is given.

Proposition 8. A tempered distribution f belongs to $\Gamma^s(x_0)$ if and only if there exists s' < 0 such that f belongs to $C^{s,s'}(x_0)$.

The weak-scaling exponent of f is

$$\beta(f, x_0) = \sup\{s : f \in \Gamma^s(x_0)\}. \tag{20}$$

But we will rather take the following alternative characterization from [20] that give a simpler condition in term of wavelet coefficients.

Proposition 9. Let f be a tempered distribution. The weak scaling exponent of f at x_0 is the supremum of s > 0 such that:

$$\forall \varepsilon > 0 \; \exists c > 0 \; \forall (j,k) \; such \; that \; |2^j x_0 - k| < 2^{\varepsilon j}, \; |c_{j,k}| \le c 2^{-(s-\varepsilon)j}$$

$$\tag{21}$$

Proofs of Theorems 1 and 2 2

2.1The p-spectrum

In this section, we will only prove the first point of Theorem 2. The will see how this proof can be adapted to Theorem 1 in a second time.

The result that we will prove in this section is more precise than the one stated. Indeed, we prove that for each $\alpha \in (1, \infty)$ and for each $p \geq 1$, the p-exponent of almost every function of $B_{p_0}^{s_0,p_0}(\mathbb{R}^d)$ is smaller than

$$s - \frac{d}{p} + \frac{d}{\alpha p}$$

on a set of Hausdorff dimension $\frac{d}{d}$.

These fractal sets are closely related to the dyadic approximation of points.

Definition 8. Let $\alpha \in (1, \infty)$ be fixed. We denote

$$F_{\alpha} = \left\{ x : \exists \ a \ sequence \ ((k_n, j_n))_{n \in \mathbb{N}} \ \left| x - \frac{k_n}{2^{j_n}} \right| \le \frac{1}{2^{\alpha j_n}} \right\}. \tag{22}$$

This set F_{α} can also be defined as

$$\limsup_{i \to \infty} \bigcup_{l \in \mathbb{N}^d} F_{\alpha}^{i,l}$$

where $F_{\alpha}^{i,l}$ denotes the cube $\frac{l}{2^i} + \left[-\frac{1}{2^{\alpha i}}; \frac{1}{2^{\alpha i}} \right]^d$. If $x \in F_{\alpha}$ it is said α -approximable by dyadics. The dyadic exponent of x is defined by $\alpha(x_0) = \sup\{\alpha : x_0 \text{ is } \alpha\text{-approximable by dyadics}\}\$

As stated in [17], the Hausdorff dimension of F_{α} is at least $\frac{d}{\alpha}$.

Let $p \geq 1$ be given such that $s_0 - \frac{d}{p_0} > -\frac{d}{p}$. For $\alpha \geq 1$ fixed we denote $s(\alpha) = s_0 - \frac{d}{p_0} + \frac{d}{\alpha p_0} + \frac{d}{p}$. For $\varepsilon > 0$ fixed, let $\beta = s(\alpha) + \varepsilon$. We first check that the set of functions in $B_{p_0}^{s_0, p_0}$ satisfying (14) with exponent β at a point in F_{α} is a Haar-null Borel set. This set can be included in a countable union over A > 0 and c > 0 of sets M(A, c)which are sets of functions in $B_{p_0}^{s_0,p_0}(\mathbb{R}^d)$ satisfying

$$\exists x \in F_{\alpha} \ \forall j \ 2^{j(\beta p - d)} \sum_{|k - 2^{j}x| \le A2^{j}} |c_{\lambda}|^{p} (1 + |k - 2^{j}x|)^{-\beta p} \le c.$$

And for each $i \in \mathbb{N}$ these sets can be included in the countable union over $l \in \{0, ..., 2^i - 1\}^d$ of $M_{i,l}(A, c)$, defined by the set of f such that

$$\exists x \in F_{\alpha}^{i,l} \ \forall j \ 2^{j(\beta p - d)} \sum_{|k - 2^{j}x| \le A2^{j}} |c_{\lambda}|^{p} (1 + |k - 2^{j}x|)^{-\beta p} \le c.$$

Each $M_{i,l}(A,c)$ is a closed set. Indeed, suppose that a sequence (f_n) of elements of $M_{i,l}(A,c)$ converges to f in $B_{p_0}^{s_0,p_0}(\mathbb{R}^d)$. Denote $c_{j,k}^n$ the wavelet coefficients of f_n , for each $n \in \mathbb{N}$ and $c_{j,k}$ those of f. For each n there exists $x_n \in F_{\alpha}^{i,l}$ such that f_n satisfies (14) at x_n . Thus

$$\forall j \ 2^{j(\beta p - d)} \sum_{|k_n - 2^j x_n| \le A2^j} |c_{\lambda}^n|^p (1 + |k_n - 2^j x_n|)^{-\beta p} \le c.$$
 (23)

As $F_{\alpha}^{i,l}$ is a compact set, there exists an accumulation point $x \in F_{\alpha}^{i,l}$ of x_n . The mapping giving the wavelet coefficients of a function f in a Besov space is continuous. Furthermore, if k_n is such that $|k_n-2^jx_n| \leq A2^j$ for a subsequence $x_{\phi(n)}$ such that $\lim x_{\phi(n)} = x$, the corresponding $k_{\phi(n)}$ converges to k and $|k-2^jx| \leq A2^j$. Thus up to a subsequence, when n tends to infinity, (23) becomes

$$\forall j \ 2^{j(\beta p - d)} \sum_{|k - 2^j x| \le A2^j} |c_{\lambda}|^p (1 + |k - 2^j x|)^{-\beta p} \le c.$$

Thus f belongs to $M_{i,l}(A,c)$ and M(A,c) is a Borel set.

To prove that it is also a Haar-null set, we construct a probe based on a slight modification of the "saturating function" introduced in [17].

Let $i \in \mathbb{N}$ and $l \in \{0, ..., 2^i - 1\}^d$ be fixed. Let $n \in \mathbb{N}$ be fixed large enough such that $N = 2^{dn} > \frac{d}{p\alpha\varepsilon} + 1$. Each dyadic cube λ is split into M subcubes of size $2^{-d(j+n)}$. For each index $i \in \{1, ..., N\}$, we choose a subcube $i(\lambda)$ and the wavelet coefficient of g_i is given by:

$$d_{\lambda}^{i} = \begin{cases} \frac{1}{j^{a}} 2^{(\frac{d}{p_{0}} - s_{0})j} 2^{-\frac{d}{p_{0}}J} & \text{if } i = i(\lambda) \\ 0 & \text{else.} \end{cases}$$
 (24)

where $a = \frac{2}{p_0}$ and $J \leq j$ and $K \in \{0, ... 2^J - 1\}^d$ are such that

$$\frac{k}{2^j} = \frac{K}{2^J}$$

is an irreducible form. It is proven in [17] that these functions belong to $B_{p_0}^{s_0,p_0}$.

Furthermore, if a point $x \in (0,1)^d$ is α -approximable by dyadics, there exists a subsequence (j_n, k_n) where $j_n = [J_n \alpha]$, J_n and K_n being defined in (22) and k_n is such that $\frac{k_n}{2^{j_n}} = \frac{K_n}{2^{j_n}}$. The corresponding wavelet coefficients of all functions g_i satisfy that there exists a constant c > 0 such that if (j, k) satisfy $|x_0 - \frac{k}{2^j}| < A$:

$$d_{j,k}^{i} > c(A) \frac{2^{(\frac{d}{p_0} - s_0)j} 2^{-\frac{d}{\alpha p_0}j}}{j^a}.$$
 (25)

Let $f = \sum c_{j,k} \psi_{j,k}$ be an arbitrary function in $B_{p_0}^{s_0,p_0}(\mathbb{R}^d)$. Suppose that there exists two points $\gamma_1 \in \mathbb{R}^N$ and $\gamma_2 \in \mathbb{R}^N$ such that for $a = 1, 2, f + \sum_i \gamma_a^i g^i$ belong to $M_{i,l}(A,c)$. By definition there also exist two points x_1 and x_2 in $F_{\alpha}^{i,l}$ such that, for a = 1, 2,

$$\forall j \ 2^{j(\beta p - d)} \sum_{|k - 2^j x_a| \le A2^j} |c_\lambda + \sum_{i=1}^N \gamma_a^i d_\lambda^i|^p (1 + |k - 2^j x|)^{-\beta p} \le c.$$

As $\beta > 0$, this condition implies :

$$\forall j \ 2^{j(\beta p - d)} \sum_{|k - 2^j x_a| \le A2^j} |c_\lambda + \sum_{i=1}^N \gamma_a^i d_\lambda^i|^p (1 + A2^j)^{-\beta p} \le c.$$

But x_1 and x_2 belong to same dyadic cubes of size j > i. Thus the same k satisfy $|k - 2^j x_a| \le A2^j$ for a = 1, 2 and wavelet coefficients of $f_1 - f_2$ are such that for all $j > \alpha i$

$$2^{j(\beta p - d)} \sum_{|k - 2^j x_a| \le A2^j} |\sum_{m = 1}^N (\gamma_1^m - \gamma_2^m) d_\lambda^m|^p (1 + A2^j)^{-\beta p} \le 2c.$$

It is obvious that

$$2^{j(\beta p - d)} \sum_{|\frac{k}{2^{j}} - x_{1}| \leq A} |\sum_{m} (\gamma_{1}^{m} - \gamma_{2}^{m}) d_{\lambda}^{m}|^{p} 2^{-\beta p j} (2^{-j} + |\frac{k}{2^{j}} - x_{1}|)^{-\beta p}$$

$$\geq 2^{j(\tilde{s}p-d)} \sup_{|\frac{k}{2j}-x_1| \leq A} |\sum_m (\gamma_1^m - \gamma_2^m) d_{\lambda}^m|^p 2^{-\tilde{s}pj} (2^{-j} + |\frac{k}{2^j} - x_1|)^{-\tilde{s}p}$$

Using definition of function g_i , if for each j we define j' = j + n, at scale j' there is only one function g_i with non zero coefficient. And with (25) one finally obtains that there exists a subsequence j such that

$$2^{n(\beta p-d)}2^{j(\beta p-d)}\sup_{|\frac{k}{2^j}-x_1|\leq A}|\sum_m(\gamma_1^m-\gamma_2^m)d_\lambda^m|^p2^{-\beta pj}(2^{-j}+|\frac{k}{2^j}-x_1|)^{-\tilde{s}p}\geq |\gamma_1^i-\gamma_2^i|^p\tilde{c}^p\frac{1}{j^{pa}}2^{p\varepsilon j},$$

where \tilde{c} depends only of n and A. Those two inequalities imply that

$$\|\gamma_1 - \gamma_2\|_{\infty}^p \le 2cc(N)i^{1/p_0}2^{-\varepsilon\alpha pi}.$$
 (26)

Therefore the set of γ such that $f + \sum_i \gamma^i g^i$ belongs to $M_{i,l}(A,c)$ is included in a ball of radius less than $(2cc(N))^N i^{N/p_0} 2^{-\varepsilon \alpha pNi}$. Taking the countable union over l, we obtain that for each i_0 fixed, the set of γ satisfying

$$\exists x \in F_{\alpha}^{i_0}$$
 such that $f + \sum_{m} \gamma^m g^m$ satisfy (23) at x

is of Lebesgue measure bounded by

$$\sum_{i=i_0}^{\infty} (2cc(N))^N i^{N/p_0} 2^{di-\varepsilon\alpha pNi}.$$

As N is large enough, this measure tends to zero when i_0 tends to infinity. And M(A, c) is then a Haar-null set.

As this result does not depend on c or on A, we can take the union over countable $c_n > 0$ and $A_n > 0$. Then the set of functions in $B_{p_0}^{s_0,p_0}(\mathbb{R}^d)$ belonging to $T_{\beta}^p(x)$ at a point $x \in F_{\alpha}$ is a Haar-null set.

Thus,

$$\forall p \geq 1, \ \forall \alpha \geq 1 \ \forall \beta > s(\alpha) \ \text{a.s. in} \ B_{p_0}^{s_0,p_0} \ \forall x \in F_\alpha \ u_f^p(x) \leq \beta.$$

Taking $\varepsilon \to 0$ it follows by countable intersection that

$$\forall p \geq 1, \ \forall \alpha \geq 1 \ \text{ a.s. in } B^{s_0,p_0}_{p_0} \ \forall x \in F_\alpha \ u^p_f(x) \leq s(\alpha).$$

Therefore, if α_n is a dense sequence in $(1, \infty)$, using the same argument, one obtains that

$$\forall p \ge 1$$
, a.s. in $B_{p_0}^{s_0, p_0} \, \forall n \in \mathbb{N} \, \forall x \in F_{\alpha_n} \, u_f^p(x) \le s(\alpha_n)$. (27)

Let f be a function satisfying (27) and $\alpha \geq 1$ be fixed. Let $\alpha_{\phi(n)}$ a nondecreasing subsequence of α_n converging to α . Then the intersection E_{α} of F_{α_n} contains F_{α} and for all $x \in E_{\alpha}$, and thus for all $x \in F_{\alpha}$, $u_f^p(x) \leq s(\alpha)$. Furthermore, see [16], there exists a measure m_{α} positive on F_{α} but such that every set of dimension less than $\frac{d}{\alpha}$ is of measure zero. Let us denote G_H the set of points where $u_p(x) < H$. According to Proposition 3, this set can be written as a countable union of sets of m_{α} measure zero. Thus, we obtain

$$m_{\alpha}(\{x: u_{p}(x) = H\}) = m_{\alpha}(F_{\alpha}\backslash G_{H}) > 0.$$

Which gives us the p spectrum of singularities

$$\forall u \in \left[s_0 - \frac{d}{p_0}, s_0 \right] \quad d_p(u) = p_0 u + d - s_0 p_0.$$

This proof does not depends on the choice of q. It can then be extended the same way for any Besov space $B_{p_0}^{s_0,q}$ for $0 \le q < \infty$.

The proof for the Sobolev case is similar. The functions g_i defined in (24) also belong to $B_{p_0}^{s_0,1}$. Since $B_{p_0}^{s_0,1} \hookrightarrow L^{s_0,p_0}$, the g_i belong to L^{s_0,p_0} and the remaining of the proof is unchanged.

2.2 Generic values of the weak scaling spectrum

We now prove of the second point of Theorem 1 and 2. As in the previous case, we prove Theorem 2, the same argument as in the previous part giving the Sobolev case.

Proposition 10. Let $s_0 > 0$ and $0 \le p_0, q < \infty$ be fixed. For almost every function in $B_{p_0}^{s_0,q}$ the spectrum of singularities for the weak-scaling exponent is given by

$$\forall \beta \in [s_0 - \frac{d}{p_0}, s_0] \quad d_{ws}(\beta) = p_0(\beta - s_0) + d. \tag{28}$$

Proof. Let $\alpha \geq 1$ be fixed and denote by F_{α} the set of Definition 8. Let $\varepsilon > 0$ be fixed and define $\beta = s_0 - \frac{d}{p_0} + \frac{d}{p_0 \alpha} + \varepsilon$. According to Proposition 9, we first have to show that for a given

c > 0 the set :

$$M_{\alpha,c} = \{ f = \sum_{k} c_{\lambda} \psi_{\lambda} \in B_{p_0}^{s_0,q} : \exists x \in F_{\alpha} \ \forall \varepsilon' > 0 \ \forall (j,k) \ |2^j x - k| \le 2^{\varepsilon' j} \ |c_{\lambda}| \le c2^{-(\beta - \varepsilon')j} \}$$

$$(29)$$

is a Borel Haar-null set.

Let us remark that for all $i \in \mathbb{N}$, this set is included in the countable union of:

$$M_{\alpha,c}(i,l) = \{ f \in B_{p_0}^{s_0,q} : \exists x \in F_{\alpha}^{i,l} \ \forall \varepsilon' > 0 \ \forall (j,k) \ |2^j x - k| \le 2^{\varepsilon' j} \ |c_{\lambda}| \le c 2^{-(\beta - \varepsilon') j} \}.$$

$$(30)$$

One easily checks that $M_{\alpha,c}(i,l)$ is closed and therefore that $M_{\alpha,c}$ is a Borel set.

To prove that $M_{\alpha,c}$ is also Haar-null, we will use a different transverse measure than in the previous case. As $M_{\alpha,c}$ depends only on the dyadic properties of points, we can restrict the proof to $[0,1]^d$. Consider the following stochastic process on $[0,1]^d$:

$$X_{x} = \sum_{j=0}^{\infty} \sum_{\lambda \in [0,1]^{d}} \varepsilon_{j,k} \frac{2^{-(s_{0} - \frac{d}{p_{0}})j} 2^{-\frac{d}{p_{0}}J}}{j^{a}} \psi(2^{j}x - k)$$
(31)

where J and a are defined as in (24) and $\{\varepsilon_{j,k}\}_{j,k}$ is a Rademacher sequence. That is the $\varepsilon_{j,k}$ are i.i.d. random variables such that

$$\mathbb{P}(\varepsilon_{j,k}=1) = \mathbb{P}(\varepsilon_{j,k}=-1) = \frac{1}{2}.$$

This process belongs to $B_{p_0}^{s_0,q}$. Furthermore, the measure defined by this stochastic process is supported by the continuous image of a compact set. Thus, $(X_x)_{x \in [0,1]^d}$ defines a compactly supported probability measure on $B_{p_0}^{s_0,q}$.

Let f be an arbitrary function in $B_{p_0}^{s_0,q}(\mathbb{R}^d)$. Thanks to Fubini's theorem, it is sufficient to prove that for all $x \in F_{\alpha}$, almost surely, condition (21) is not satisfied by f + X.

Let $x_0 \in F_\alpha$ be fixed and suppose that f+X satisfies condition (21) at x_0 . Then for all $\varepsilon' > 0$ and for all (j,k) such that $|k-2^jx_0| \leq 2^{\varepsilon'j}$,

$$|c_{j,k} + \varepsilon_{j,k} \frac{2^{-(s_0 - \frac{d}{p_0})j} 2^{-\frac{d}{p_0}J}}{j^a}| \le c 2^{-(\beta - \varepsilon')j}.$$

Taking (J_n, K_n) the sequence of Definition 8, $j = [\alpha J_n]$ and $k = \frac{K_n 2^j}{2^{J_n}}$ one obtains that there exists a sequence (j, k) such that $|2^j x_0 - k| \le 1$ and the following property holds:

$$\varepsilon_{j,k} = c_{j,k} j^a 2^{\left(s_0 - \frac{d}{p_0} + \frac{d}{p_0 \alpha}\right)j} + o(2^{-(\varepsilon - \varepsilon')j}).$$

Taking $\varepsilon' = \frac{\varepsilon}{2}$, one obtains that $\varepsilon_{j,k} \sim c_{j,k} j^a 2^{(s_0 - \frac{d}{p_0} + \frac{d}{p_0 \alpha})j_n}$ when $j_n \to \infty$. Since the $c_{j,k}$ are deterministic, this result implies that there exists an infinite sequence of independent stochastic variables which are deterministic. This event is of probability zero and $M_{\alpha,c}$ is a Haarnull set.

Therefore, taking countable unions over c > 0 and $\varepsilon \to 0$, it follows that for all $\alpha \ge 1$, the set of functions in $B_{p_0}^{s_0,q}$ with a weak scaling exponent greater than $s_0 - \frac{d}{p_0} + \frac{d}{p_0 \alpha}$ at some point of F^{α} is a Haar-null set.

Let $(\alpha_n)_{n\in\mathbb{N}}$ be a dense sequence in $(1,\infty)$ and take a countable union over α_n . We finally obtain

a. s. in
$$B_{p_0}^{s_0,q}(\mathbb{R}^d) \ \forall n \in \mathbb{N} \ \forall x \in F_{\alpha_n} \ \beta(f,x) \le s_0 - \frac{d}{p_0} + \frac{d}{p_0 \alpha_n}$$
.

With a similar argument as in part 2.1, one can prove that:

a. s. in
$$B_{p_0}^{s_0,q}(\mathbb{R}^d) \ \forall \alpha \ge 1 \ \forall x \in F_{\alpha} \ \beta(f,x) \le s_0 - \frac{d}{p_0} + \frac{d}{p_0\alpha}$$
. (32)

Furthermore, we saw in Part 2.1 that there exists a measure m_{α} which is positive on F_{α} and such that

$$m_{\alpha}(\lbrace x; u_p(x) = s_0 - \frac{d}{p_0} + \frac{d}{p_0 \alpha} \rbrace) > 0.$$

And by definition, $\forall p \geq 1, \beta(f, x) \geq u_p(x)$, thus

$$m_{\alpha}(\{x; \beta(f, x) = s_0 - \frac{d}{p_0} + \frac{d}{p_0 \alpha}\}) > 0.$$

Which states that the spectrum of singularities for the weak scaling exponent of almost every function in $B_{p_0}^{s_0,q}(\mathbb{R}^d)$ is given by

$$\forall \beta \in [s_0 - \frac{d}{p_0}, s_0] \quad d_{ws}(\beta) = p_0 \beta + d - s_0 p_0.$$

References

[1] P. Abry, Ondelettes et turbulences. Multirésolutions, algorithmes de décomposition, invariance d'échelle et signaux de pression, Nouveaux Essais. Paris: Diderot, 1997.

- [2] A. Arneodo, Ondelettes, multifractales et turbulences: De l'ADN aux croissances cristallines, inconnu, 1980.
- [3] A. Arneodo, B. Audit, N. Decoster, J-F. Muzy, and C. Vaillant, Wavelet-based multifractal formalism: applications to DNA sequences, satellite images of the cloud structure and stock market data, The science of Disasters (2002), 27–102.
- [4] A. Arnéodo, E. Bacry, and J.F. Muzy, *The thermodynamics of fractals revisited with wavelets*, Physica A. **213** (1995), 232–275.
- [5] S. Banach, Über die Baire'sche Kategorie gewisser Funktionenmengen, Studia Math. 3 (1931), 174-179.
- [6] Y. Benyamini and J. Lindenstrauss, Geometric nonlinear functional analysis. Volume 1, Colloquium Publications. American Mathematical Society (AMS), 2000.
- [7] J-M. Bony, Second microlocalization and propagation of singularities for semilinear hyperbolic equations., Proc. Taniguchi Int. Symp., 1986, pp. 11–49.
- [8] L. Caffarelli, M. G. Crandall, M. Kocan, and A. Swich, On viscosity solutions of fully nonlinear equations with measurable ingredients, Comm. Pure Appl. Math. 49 (1996), no. 4, 365–397.
- [9] A. Calderón and A. Zygmund, Local properties of solutions of elliptic partial differential equations, Studia Math. **20** (1961), 171–227.
- [10] J.P.R. Christensen, On sets of Haar measure zero in Abelian Polish groups, Israel J. Math. 13 (1972), 255–260.

- [11] A. Fraysse and S. Jaffard, How smooth is almost every function in a Sobolev space?, Rev. Mat. Iber. Amer. 22 (2006), no. 2, 663–682.
- [12] U. Frisch and G. Parisi, On the singularity structure of fully developed turbulence, Turbulence and Predictability in Geophysical Fluid Dynamics and Climate Dynamics (R. Benzi M. Ghil and G. Parisi (North-Holland), eds.), 1985, pp. 84–88.
- [13] B. Hunt, The prevalence of continuous nowhere differentiable function, Proceed. A.M.S 122 (1994), no. 3, 711–717.
- [14] B. Hunt, T. Sauer, and J. Yorke, Prevalence: A translation invariant "almost every" on infinite dimensional spaces, Bull. A.M.S 27 (1992), no. 2, 217–238.
- [15] S. Jaffard, Multifractal formalism for functions, SIAM J. Math. Anal 28 (1997), 944–970.
- [16] _____, Old friends revisited: The multifractal nature of some classical functions, J. Four. Anal. App 3 (1997), no. 1, 1–22.
- [17] _____, On the Frisch-Parisi conjecture, J. Math. Pures Appl **79** (2000), 525–552.
- [18] _____, Pointwise regularity criteria, C. R. Acad. Sci. Paris, Ser. I 336 (2003).
- [19] S. Jaffard, B. Lashermes, and P. Abry, Wavelet leaders n multifractal analysis, Wavelet Analysis and Applications (2006), 201– 246.
- [20] S. Jaffard and C. Melot, Wavelet analysis of fractal boundaries. Part 1: Local exponents, Comm. Math. Phys. **258** (2005), no. 3, 513–539.
- [21] _____, Wavelet analysis of fractal boundaries. Part 2: Multifractal formalism, Comm. Math. Phys. 258 (2005), no. 3, 541–565.
- [22] S. Jaffard and Y. Meyer, On the pointwise regularity in critical Besov spaces, J. Funct. Anal 175 (2000), 415–434.
- [23] P. G. Lemarié-Rieusset, Recent developments in the Navier-Stokes problem, Chapman & Hall/CRC Research Notes in Mathematics, vol. 431, Chapman & Hall/CRC, Boca Raton, FL, 2002.
- [24] J. Leray, Sur le mouvement d'un liquide visqueux emplissant l'espace, Acta Math. 63 (1934), no. 1, 193-248.
- [25] C. Melot, Oscillating singularities in Besov spaces, J. Math. Pures Appl 83 (2004), 367–416.

- [26] Y. Meyer, Ondelettes et opérateurs, Hermann, 1990.
- [27] _____, Wavelets, vibrations and scalings, CRM series AMS 9 (1998).
- [28] S. Mimouni, Analyse fractale d'interfaces pour les instabilités de Raleigh-Taylor, Ph.D. thesis, Ecole Polytechinique, 1995.
- [29] Rodolfo Salvi (ed.), Navier-Stokes equations: theory and numerical methods, Pitman Research Notes in Mathematics Series, vol. 388, Harlow, Longman, 1998.
- [30] S. Seuret, Detecting and creating oscillations using multifractal methods, Math. Nachr. 279 (2006), no. 11, 1195–1211.
- [31] W. Ziemer, Weakly differentiable functions. Sobolev spaces and functions of bounded variation, Graduate Texts in Mathematics, 120. Berlin etc.: Springer-Verlag., 1989.