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# On the Design Conditions of Planar Parallel Manipulators in Near-Singular Configurations 

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#### Abstract

Singularity is a major problem for parallel manipulators as it causes severe problems: uncontrollable or unstable motions, loss of stiffness, infinite forces/torques in actuated joints, and possibility of a breakdown. Many studies have been devoted to these problems. In particular, various criteria for measuring the closeness to singularities have been proposed. They are based on the limitation of the values of the input force/torque, pressure angle, natural frequency, etc. The present paper expands the previous work by incorporating a new criterion into mentioned problem, which allows the limitation of the reaction in passive joints. For this purpose, a new approach is developed, which shows that for given external force and moment applied on the platform, the reactions in passive joints depend not only on the pressure angle but also the position of the instantaneous centre of rotation of the platform. The developed approach is not only an effective way to reduce computation time but also to give a qualitative and geometrical interpretation of the problem. The suggested procedure for the determination of the reachable manipulator workspace taking into account the admissible reactions in passive joints is illustrated via a planar $3-R \underline{R} R$ parallel manipulator.

Keywords: design, effort transmission, instantaneous centre of rotation, planar parallel robots, pressure angle, singularities.


## I Introduction

It is well known that a parallel manipulator at a singular configuration can gain one or more degrees of freedom. Moreover, at such a configuration, it cannot be controlled, and it can have infinite force/torque in actuated joints, as well as infinite reactions in passive joints, leading to a manipulator breakdown. Therefore the singular configurations and the properties of parallel manipulators close to these configurations have been studied. The singularity analysis of parallel manipulators have been carried out from kinematic [1-5], kinetostatic [ $6-8]$ and dynamic points of view $[9-11]$. The task of avoiding singular configurations raises the question of

[^0]how one has to stay away from singular configurations. From this point of view, different criteria and several approaches have been proposed. The optimization under constraints has been used [12] to determine the linear complex that is closest to a singularity. This methodology was applied to describe the instantaneous behavior near singularities [13]. The optimization under constraints was also adapted in order to generalize the concept of center-of-compliance [14, 15]. The dexterity and conditioning indices were also proposed [16, 17]. These indices are based on the Jacobian matrix or its "norm", which relate the actuator velocities (efforts, resp.) to the platform twist (wrench, resp.) by the following relations:
\[

$$
\begin{equation*}
\mathbf{t}=\mathbf{J} \dot{\mathbf{q}} \text { and } \mathbf{w}=\mathbf{J}^{-T} \boldsymbol{\tau} \tag{1}
\end{equation*}
$$

\]

where $\mathbf{J}$ is the Jacobian matrix, $\mathbf{t}$ the platform twist, $\dot{\mathbf{q}}$ the input velocities, $\tau$ the actuator efforts, and $\mathbf{w}$ the wrench applied on the platform.

However, it is known that because of the non homogeneity of the terms of the Jacobian matrix, the previous indices are not well appropriated for mechanisms having both translational and rotational degrees of freedom (DOF). Moreover, as they are based on the Euclidian norm of the input vector ( $\|\dot{\mathbf{q}}\|$ or $\|\tau\|$ being considered equal to 1 ), while it is clear that each actuator may have a velocity or an effort comprised between a minimal and a maximal value $\left(\dot{q}_{i} \in[-\right.$ $\left.\dot{q}_{i}^{\text {max }}, \dot{q}_{i}^{\text {max }}\right]$ and $\tau_{i} \in\left[-\tau_{i}^{\text {max }}, \tau_{i}^{\text {max }}\right], \dot{q}_{i}$ and $\tau_{i}$ being the velocity and effort of actuator $i$ ), they do not take into account the 'technological reality' of the mechanism.

In order to overcome these drawbacks, in [8] it is proposed to characterize the force workspace of robots, which is the workspace for which, considering a given fixed wrench applied on the platform, the actuator efforts are always comprised in the interval $\left[-\tau_{i}^{\max }, \tau_{i}^{\max }\right]$. However, this workspace depends on the given direction and norm of the external wrench and will change if another wrench is considered. Moreover, for many robot applications, the wrench direction is not known, contrary to its norm. Therefore, in [18], a way to compute the maximal workspace taking into account the actuator effort limitations for a given norm of the external wrench was proposed. This approach is based on the computation of the transmission factors of matrix $\mathbf{J}^{-T}$, which are
obtained geometrically through the mapping of a cube by the Jacobian matrix.

The main drawback of the last two approaches appears in the fact that they analyze the mechanism breakdown by considering input torque limitations only. However, it can exist such poses of the manipulator close to the singularity for which the limitations of input torques can be satisfied but where the reactions in passive joints are not admissible. To have a better understanding of this fact, let us consider a simple example. For the planar $5 R$ manipulator close to the singularity (Fig. 1), a small effort $\mathbf{w}$ applied on the end-effector will create a large reaction $\mathbf{R}_{\mathbf{1}}$ in the passive joint $B$. But in this pose, the actuator torques $\tau_{1}$ will stay under acceptable values, as it depends only on the small component $\mathbf{F}_{1}$ of the reaction $\mathrm{R}_{1}$.

Thus we see that the limitation of the passive joints' reactions is also an important consideration. It is a complementary criterion for measuring closeness to singularity. The most trivial way to compute the maximal passive joint reactions is to apply the Newton-Euler theorem to the manipulator platform and, by variation of the force vector directions and moment values, to compute the joint reactions. However, this is a complicated way which requires considerably huge computation time.

Therefore, in this paper a new method for obtaining the maximal efforts in passive joints taking into account the norm of the external wrench applied on the platform of planar parallel manipulator is proposed. It will be shown that, for a given ratio between the norm of the external wrench and of the maximal admissible effort into the passive joints, the mechanical breakdown can be avoided if the pressure angle and the minimal distance between the considered passive joint and the platform instantaneous axis of rotation are bounded. The developed approach is not only an effective way to reduce computation time but also to give a qualitative and geometrical interpretation of the problem.

## II. A new criterion for evaluation of the passive joints efforts

For any planar manipulator, the forces $\mathbf{R}_{i}$ transmitted by the legs through the passive joints linked to the platform (Fig. 2) can be related to the external wrench $\mathbf{w}^{T}=\left[\mathbf{f}^{T}\right.$, $m]^{T}$ (f is the external force, and $m$ the scalar value of the external moment applied on the platform) using the Newton-Euler theorem at any point $Q$ :

$$
\begin{equation*}
\mathbf{f}=-\sum_{i=1}^{3} \mathbf{R}_{i} \text { and } m=-\sum_{i=1}^{3} \overline{\mathbf{d}}_{Q B i}^{T} \mathbf{R}_{i} \tag{2}
\end{equation*}
$$

where $\overline{\mathbf{d}}_{Q B i}^{T}=\left[-y_{Q B i}, x_{Q B i}\right]$, with $x_{Q B i}$ and $y_{Q B i}$ are the coordinates of vector $\mathbf{Q B}_{i}$ along $\boldsymbol{x}$ and $\boldsymbol{y}$ axes. Considering that $\mathbf{R}_{i}=R_{i} \mathbf{r}_{i}$, where $\left\|\mathbf{R}_{i}\right\|=R_{i}$, and applying the Newton-Euler theorem at point $B_{1}$, it comes that:


Fig. 1. Planar $5 R$ manipulator close to the Type 2 singular pose.


Fig. 2. Determination of the pressure angle for the planar 3- $R \underline{R} R$ robot.

$$
\begin{align*}
\mathbf{w} & =-\left[\begin{array}{ccc}
\mathbf{r}_{1} & \mathbf{r}_{2} & \mathbf{r}_{3} \\
0 & \overline{\mathbf{d}}_{B_{1} B_{2}}^{T} \mathbf{r}_{2} & \overline{\mathbf{d}}_{B_{1} B_{3}}^{T} \mathbf{r}_{3}
\end{array}\right]\left[\begin{array}{l}
R_{1} \\
R_{2} \\
R_{3}
\end{array}\right]  \tag{3}\\
& =-\left[\begin{array}{lll}
\mathbf{f}_{1} & \mathbf{f}_{2} & \mathbf{f}_{3}
\end{array}\right] \mathbf{R}=-\mathbf{A}^{T} \mathbf{R}
\end{align*}
$$

$\mathbf{f}_{i}^{T}=\left[\mathbf{r}_{i}^{T}, \overline{\mathbf{d}}_{B_{1} B_{i}}^{T} \mathbf{r}_{i}\right]$ being a unit screw corresponding to the direction of the wrench applied by the actuators on the platform.

Matrix A used in equation (3) is defined in [1] as the parallel Jacobian matrix than can be found through the differentiation of the loop closure equations of the robot with respect to the platform coordinates. As a result, this matrix is always invertible if the robot is not in a Type 2 singularity.

The reactions $\mathbf{R}$ of the passive joints can be found from (3):

$$
\begin{equation*}
\mathbf{R}=-\mathbf{A}^{-T} \mathbf{w} \tag{4}
\end{equation*}
$$

where matrix $\mathbf{A}^{-1}$ can be expressed in the form:

$$
\mathbf{A}^{-1}=\left[\begin{array}{lll}
\mathbf{t}_{\mathbf{1}} & \mathbf{t}_{2} & \mathbf{t}_{3} \tag{5}
\end{array}\right] \text {, with } \mathbf{t}_{i}^{T}=\left[\mathbf{v}_{i}^{T}, \omega_{i}\right]
$$

where $\mathbf{t}_{i}$ is a screw collinear to the twist of the platform expressed at point $B_{i}$ when leg $i$ is disconnected, $\mathbf{v}_{i}$ is the translational velocity of point $B_{i}$, and $\omega_{i}$ the platform rotational velocity [7] (Fig. 2). On this picture, point $I_{1}$ corresponds to the instantaneous centre of rotation of the platform when leg 1 is disconnected.

Taking into account that $\mathbf{A}^{T} \mathbf{A}^{-T}=\mathbf{I}$, we obtain: $\mathbf{f}_{\mathbf{1}}^{T} \mathbf{t}_{\mathbf{1}}=\mathbf{r}_{\mathbf{1}}^{T} \mathbf{v}_{\mathbf{1}}=1$.


Fig. 3. Instantaneous system equivalent to the planar 3- $R \underline{R} R$ robot platform.

It was shown in [7] that the pressure angle of the leg 1, depicted as $\alpha_{1}$ (see Fig. 2), may be expressed as the acute angle between the directions of vectors $\mathbf{r}_{1}$ and $\mathbf{v}_{1}$. Therefore:

$$
\begin{equation*}
\left|\mathbf{r}_{1}^{T} \mathbf{v}_{\mathbf{1}}\right|=\left\|\mathbf { r } _ { \mathbf { 1 } } \left|\left\|\mid \mathbf{v}_{\mathbf{1}}\right\| \cos \alpha_{1} \text {, i.e. } \cos \alpha_{1}=\frac{\left|\mathbf{r}_{\mathbf{1}}^{T} \mathbf{v}_{\mathbf{1}}\right|}{\left\|\mathbf{r}_{\mathbf{1}}\right\|\left\|\mathbf{v}_{\mathbf{1}}\right\|}\right.\right. \tag{6}
\end{equation*}
$$

By definition, $\mathbf{r}_{1}$ is a unit vector. As a result, $\left\|\mathbf{v}_{\mathbf{1}}\right\|=\left(\cos \alpha_{1}\right)^{-1}$.

For a rigid body in a planar displacement, it is known that the norm of the translational velocity $\mathbf{v}_{\mathbf{Q}}$ of any point $Q$ is equal to the product of the angular velocity absolute value by the distance between $Q$ and the instantaneous centre of rotation of the body (denoted as $I$ ). Thus, we have $\left\|\mathbf{v}_{1}\right\|=d_{1}\left|\omega_{1}\right|$, where $d_{1}$ is the distance from the platform instantaneous centre of rotation $I_{1}$ to point $B_{1}$ (Fig. 2).

Developing (4), it comes that:

$$
R_{1}=-\left(\mathbf{v}_{\mathbf{1}}^{T} \mathbf{f}+\omega_{1}\left(m+\overline{\mathbf{d}}_{B 1 P}^{T} \mathbf{f}\right)\right)=-\left(\left(\mathbf{v}_{\mathbf{1}}^{T}+\omega_{1} \overline{\mathbf{d}}_{B 1 P}^{T}\right) \mathbf{f}+\omega_{1} m\right) .
$$

For a given norm $f$ of the external force $\mathbf{f}$ and a given value $m$ of the external moment, and for any direction of vector $\mathbf{f}$, the maximal value $R_{1 \text { max }}$ of $R_{1}$ appears when

$$
\begin{align*}
R_{1 \max } & =\max _{\mathbf{f}, m}\left(R_{1}\right) \\
& =f \sqrt{\left\|\mathbf{v}_{\mathbf{1}}\right\|^{2}+\left(\omega_{1} b_{1}\right)^{2}-2\left\|\mathbf{v}_{\mathbf{1}}\right\|\left|\omega_{1} b_{1}\right| \cos \beta_{1}}+\left|\omega_{1} m\right| \tag{7}
\end{align*}
$$

where $b_{1}$ is the distance between the application point of the external wrench, denoted as $P$, and point $B_{1} . \beta_{1}$ is the angle between vectors $\mathbf{v}_{1}$ and $-\omega_{1} \overline{\mathbf{d}}_{B 1 P}$. For a numerical computation of the joint reactions, using (7) could be sufficient, as all of these parameters can be defined via matrix $\mathbf{A}$ and the mechanism geometry. But, it will be shown below that we can go deeper in the analysis using the pressure angle $\alpha_{1}$ and the distance $d_{1}$. These indices will give valuable insights for the effort transmission inside of the mechanism.

Transforming expression (7) taking exp. (6) into account, and generalizing the approach to the other legs, we obtain the maximal value $R_{i \max }$ of the platform joint reaction at $B_{i}(i=1$ to 3$)$

$$
\begin{equation*}
R_{i \max }=\frac{\gamma_{i} f+|m| / d_{i}}{\cos \alpha_{i}} \tag{8}
\end{equation*}
$$

where $\quad \gamma_{i}=\sqrt{1+\left(b_{i} / d_{i}\right)^{2}-2 \cos \beta_{i} b_{i} / d_{i}}, \quad d_{i} \quad$ is the distance between point $B_{i}$ and the instantaneous centre of rotation of the platform when leg $i$ is disconnected, $\alpha_{i}$ is the pressure angle of the leg $i, b_{i}$ is the distance between $P$ and $B_{i}$, and $\beta_{i}$ is the angle between vectors $\mathbf{v}_{i}$ and $-\omega_{i} \overline{\mathbf{d}}_{B i P}$ (and as a result between vectors $\mathbf{B}_{i} \mathbf{P}$ and $\mathbf{B}_{i} \mathbf{I}_{i}-$ Fig. 3), $\overline{\mathbf{d}}_{B i P}, \mathbf{v}_{i}$ and $\omega_{i}$ being defined at (2) and (5). It should be mentioned that the distance between point $P$ and $I_{i}$ is equal to $\gamma_{i} d_{i}$.

This equation is the basic result of the present paper, which shows that, for a given set of external force and moment applied on the platform, the reactions in passive joints depend not only on the pressure angle but also the position of the instantaneous centre of rotation of the platform. This property has not been rigorously formulated and demonstrated before. It allows not only giving a qualitative evaluation of the kinetostatic properties of the parallel manipulators and reducing the computational time but also disclosing the geometrical interpretation of the problem: from the equation (8), it can be shown that the mechanical system under study can be instantaneously replaced by a virtual cantilever attached to the ground at $I_{i}$ (Fig. 3) lying on a virtual point contact at $B_{i}$, of which direction is parallel to the vector $\mathbf{R}_{i}$. The external force $\mathbf{f}$ is applied on the point $P$ of the cantilever.

As $\alpha_{i}$ and $d_{i}$ are criterion used for the kinetostatic design of robots [6, 7], it is of interest to find their boundaries with respect to some design requirements. In order to avoid the breakdown of the studied joint, technological requirements imply that the admissible value of $R_{i}$ should not be superior to a given value $R_{a d m}$, i.e. $R_{i \max } \leq R_{\text {adm }}$. Thus, we would like to find $\alpha_{i}$ and $d_{i}$, i.e. the parameters of the equivalent cantilever system, for which this inequality is respected. Introducing (8) into the last inequality leads to:

$$
\begin{equation*}
f \sqrt{1+\left(b_{i} / d_{i}\right)^{2}-2 \cos \beta_{i} b_{i} / d_{i}} \leq R_{a d m} \cos \alpha_{i}-|m| / d_{i} \tag{9}
\end{equation*}
$$

Please note that, as by definition, $f, \cos \alpha_{i}, d_{i}$ and $R_{\text {adm }}$ have positive values, a necessary condition for the existence of (9) is:

$$
\begin{equation*}
0 \leq R_{a d m} \cos \alpha_{i}-|m| / d_{i} \Rightarrow \frac{|m|}{R_{a d m} \cos \alpha_{i}} \leq d_{i} \tag{10}
\end{equation*}
$$

If not, the left term of (9) will always be superior to $R_{a d m}$. Let us now square the left and right sides of equation (9) and multiply them by $d_{i}^{2}$ :

$$
\begin{align*}
f^{2}\left(d_{1}^{2}+b_{1}^{2}-2 \cos \beta_{1} b_{1} d_{1}\right) \leq & d_{1}^{2}\left(R_{a d m} \cos \alpha_{1}\right)^{2}+\ldots  \tag{11}\\
& m^{2}-2 d_{1}|m| R_{a d m} \cos \alpha_{1}
\end{align*}
$$

The obtained expression can be rewritten as:

$$
\begin{equation*}
0 \leq p_{i}\left(d_{i}\right), \tag{12}
\end{equation*}
$$

where, $\quad p_{i}\left(d_{i}\right)=u_{i} d_{i}^{2}+v_{i} d_{i}+w_{i}, u_{i}=R_{a d m}^{2} \cos ^{2} \alpha_{i}-f^{2}$, $v_{i}=-2\left(|m| R_{a d m} \cos \alpha_{i}-\cos \beta_{i} b_{i} f^{2}\right), w_{i}=m^{2}-f^{2} b_{i}^{2}$.

The inequality (12) has different solutions, depending on the vanishing and signs of terms $u_{i}, v_{i}$ and $w_{i}$. There are three main cases $u_{i}>0, u_{i}<0$ and $u_{i}=0$.

Let us consider these cases.

## A. $u_{i}>0$

$u_{i}>0$ implies that $f<R_{\text {adm }} \cos \alpha_{i}$. In this case, the polynomial $p_{i}$ has got two roots but only one corresponds to the real mechanism, i.e. to a solution of (9). The other root is solution of:
$-f^{2} d_{1}^{2} \gamma_{1}^{2} \leq d_{1}^{2}\left(R_{a d m} \cos \alpha_{1}\right)^{2}+m^{2}-2 d_{1}|m| R_{a d m} \cos \alpha_{1}$
In order to obtain a root of (12) with physical values, it is necessary that the condition $v_{i}^{2}-4 u_{i} w_{i} \geq 0$ should be respected, i.e

$$
\begin{align*}
& \left(R_{a d m}^{2} f^{2} b_{i}^{2}\right) \cos ^{2} \alpha_{i}-2\left(|m| R_{a d m} \cos \beta_{i} b_{i} f^{2}\right) \cos \alpha_{i}  \tag{14}\\
& \ldots-\sin ^{2} \beta_{i} b_{i}^{2} f^{4}+f^{2} m^{2} \geq 0
\end{align*}
$$

Developing and simplifying, it can be shown that this polynomial in $\cos \alpha_{i}$ has roots with real values if and only if:

$$
\begin{equation*}
R_{a d m}^{2} \sin ^{2} \beta_{i} b_{i}^{2} f^{4} w_{i} \leq 0 \tag{15}
\end{equation*}
$$

For the analysis of this inequality, two following cases must be considered: $w_{i} \leq 0$ and $w_{i}>0$.

$$
\text { A1. } w_{i}>0
$$

The condition $w_{i}>0$ implies that $|m|>f b_{i}$. Here, (14) has no real roots, i.e. (14) is true for any value of $\alpha_{i}$. Thus, the condition for (9) to be true is that

$$
\begin{equation*}
d_{i} \geq \max \left(\left(d_{i}\right)_{1},|m| /\left(R_{a d m} \cos \alpha_{i}\right)\right) \tag{16}
\end{equation*}
$$

where $\left(d_{i}\right)_{1}=\left(-v_{i}+\sqrt{v_{i}^{2}-4 u_{i} w_{i}}\right) / 2 u_{i}$ is the root of (12) solution of (9).

$$
A 2 . w_{i} \leq 0
$$

The condition $w_{i} \leq 0$ implies that $|m| \leq f b_{i}$. (14) is true if its roots are bounded by

$$
\begin{equation*}
\cos \alpha_{i} \geq \frac{|m| \cos \beta_{i}+\left|\sin \beta_{i}\right| \sqrt{f^{2} b_{i}^{2}-m^{2}}}{R_{a d m} b_{i}} \tag{17a}
\end{equation*}
$$

or

$$
\begin{equation*}
\cos \alpha_{i} \leq \frac{|m| \cos \beta_{i}-\left|\sin \beta_{i}\right| \sqrt{f^{2} b_{i}^{2}-m^{2}}}{R_{a d m} b_{i}} \tag{17b}
\end{equation*}
$$

After mathematical simplifications, it can be proven that if (17b) is true, then,

$$
\begin{equation*}
\cos \alpha_{i} \leq \frac{|m|}{R_{a d m} b_{i}} \leq \frac{f}{R_{a d m}} \tag{18}
\end{equation*}
$$

which is in contradiction with $u_{i}>0$. Therefore, the only condition for (14) to be true is that
$\cos \alpha_{i} \geq \max \left(\frac{|m| \cos \beta_{i}+\left|\sin \beta_{i}\right| \sqrt{f^{2} b_{i}^{2}-m^{2}}}{R_{a d m} b_{i}}, \frac{f}{R_{a d m}}\right)$
However, it can be demonstrated that, for $u_{i}>0$ and $w_{i}$ $\leq 0$

$$
\begin{equation*}
\frac{|m| \cos \beta_{i}+\left|\sin \beta_{i}\right| \sqrt{f^{2} b_{i}^{2}-m^{2}}}{R_{a d m} b_{i}} \leq \frac{f}{R_{a d m}} . \tag{20}
\end{equation*}
$$

Thus (19) implies that $u_{i}>0$, which is true. Then, (14) is always true is this section. As a result, the only condition for $(9)$ to be true is (16).

## B. $u_{i}<0$

$u_{i}<0$ implies that $f>R_{\text {adm }} \cos \alpha_{i}$. Introducing this into (9) leads to $0 \leq \gamma_{i} \leq 1$, i.e. the cantilever allows decreasing the applied force $\mathbf{f}$. Here also, two cases should be studied: $w_{i} \leq 0$ and $w_{i}>0$.

$$
\text { B1. } w_{i}>0
$$

If $w_{i}>0$, i.e. $|m|>f b_{i}$, from (10) it can be shown that:

$$
\begin{equation*}
\frac{f b_{i}}{R_{a d m} \cos \alpha_{i}}<\frac{|m|}{R_{a d m} \cos \alpha_{i}} \leq d_{i} \tag{21}
\end{equation*}
$$

As in section II.B $u_{i}<0$, which is equivalent to $R_{a d m}$ $\cos \alpha_{i}<f$, from (21) it comes

$$
\begin{equation*}
b_{i}<\frac{f b_{i}}{R_{a d m} \cos \alpha_{i}}, \text { thus } b_{i}<d_{i} \tag{22}
\end{equation*}
$$

If (22) is true, the expression of $\gamma_{i}$ at (8) is bounded by

$$
\begin{equation*}
\gamma_{i} \geq \frac{\left|b_{i}-d_{i}\right|}{d_{i}}=\frac{d_{i}-b_{i}}{d_{i}}>0 \tag{23}
\end{equation*}
$$

Introducing (23) into (9), and as $|m|>f b_{i}$, it comes that:

$$
\begin{equation*}
\frac{f\left(d_{i}-b_{i}\right)+f b_{i}}{d_{i}}=f<\frac{f\left(d_{i}-b_{i}\right)+|m|}{d_{i}} \leq R_{a d m} \cos \alpha_{i} . \tag{24}
\end{equation*}
$$

Thus,

$$
\begin{equation*}
f<R_{a d m} \cos \alpha_{i}, \text { or equivalently } u_{i}>0 \tag{25}
\end{equation*}
$$

which is impossible in section 2.2.

$$
B 2 . w_{i} \leq 0
$$

If $w_{i} \leq 0$, i.e. $|m| \leq f b_{i}$, it could be shown after several mathematical simplifications that (14) is true if and only if
$\frac{f}{R_{a d m}}>\cos \alpha_{i} \geq \frac{|m| \cos \beta_{i}+\left|\sin \beta_{i}\right| \sqrt{f^{2} b_{i}^{2}-m^{2}}}{R_{a d m} b_{i}}$.

|  | $\|m\|>f b_{i}$ | $\|m\|<f b_{i}$ |  |  | $\|m\|=f b_{i}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\frac{f}{R_{a d n}}<\cos \alpha_{i}$ | $d_{i} \geq \max \left(\left(d_{i}\right)_{1}, \frac{\|m\|}{R_{a d m} \cos \alpha_{i}}\right)$ |  |  |  |  |  |  |
| $\frac{f}{R_{a d n}}>\cos \alpha_{i}$ | N/A | $\cos \alpha_{i} \geq \frac{\|m\| \cos \beta_{i}+\left\|\sin \beta_{i}\right\| \sqrt{f^{2} b_{i}^{2}-m^{2}}}{R_{a d m} b_{i}}, \max \left(\frac{\|m\|}{R_{a d m} \cos \alpha_{i}}, \frac{b_{i}}{2 \cos \beta_{i}}\right) \leq d_{i} \leq\left(d_{i}\right)_{2}$ |  |  |  |  |  |
| $\frac{f}{R_{a d m}}=\cos \alpha_{i}$ | N/A | $v_{i}>0$ | $v_{i}<0$ | $v_{i}=0$ | $v_{i}>0$ | $v_{i}<0$ | $v_{i}=0$ |
|  |  | $d_{i} \geq \max \left(\frac{\|m\|}{R_{\text {adm }} \cos \alpha_{i}}, \frac{-w_{i}}{v_{i}}\right)$ | N/A | N/A | $d_{i} \geq \frac{\|m\|}{R_{\text {add }} \cos \alpha_{i}}$ | N/A | $\begin{gathered} d_{i} \geq b_{i}, \text { for } \\ \beta_{i}=0 \end{gathered}$ |

TABLE 1. Conditions for limiting the maximal values of the revolute joint linked to the platform.

Once this condition is achieved, the condition for (9) to be true is that

$$
\begin{equation*}
\max \left(\frac{|m|}{R_{a d m} \cos \alpha_{i}}, \frac{b_{i}}{2 \cos \beta_{i}}\right) \leq d_{i} \leq\left(d_{i}\right)_{2} \tag{27}
\end{equation*}
$$

where $\left(d_{i}\right)_{2}=\left(-v_{i}-\sqrt{v_{i}^{2}-4 u_{i} w_{i}}\right) / 2 u_{i}$ is the root of (12) solution of (9).
C. $u_{i}=0$

We have to analyze one last case: $u_{i}=0$, i.e. $\cos \alpha_{i}=f / R_{a d m}$. Here also, this condition leads to $0 \leq \gamma_{i}$ $\leq 1$, i.e. the cantilever allows decreasing the applied force f. In such a case, the solution of (12) is solution of $0 \leq v_{i} d_{i}+w_{i} \quad$ with $\quad v_{i}=2 f\left(\cos \beta_{i} b_{i} f-|m|\right) \quad$ and $w_{i}=m^{2}-f^{2} b_{i}^{2}$.

Three cases will appear: $v_{i}>0, v_{i}<0$ and $v_{i}=0$.
C1. $v_{i}>0$
In this case, $R_{i \max } \leq R_{\text {adm }}$ can be satisfied if and only if

$$
\begin{equation*}
\frac{-w_{i}}{v_{i}}=\frac{f^{2} b_{i}^{2}-m^{2}}{2 f\left(\cos \beta_{i} b_{i} f-|m|\right)} \leq d_{i} \tag{28}
\end{equation*}
$$

C2. $v_{i}<0$
In this case, $R_{i \max } \leq R_{\text {adm }}$ can be satisfied if and only if

$$
\begin{equation*}
0 \leq d_{i} \leq \frac{-w_{i}}{v_{i}}=\frac{f^{2} b_{i}^{2}-m^{2}}{2 f\left(\cos \beta_{i} b_{i} f-|m|\right)} \tag{29}
\end{equation*}
$$

As in section 2.2.1, the condition $w_{i}>0$ is not compatible with the fact that $u_{i}=0$. For $w_{i} \leq 0$ :

$$
\begin{equation*}
\frac{f^{2} b_{i}^{2}-m^{2}}{2 f\left(\cos \beta_{i} b_{i} f-|m|\right)} \leq 0 \tag{30}
\end{equation*}
$$

Therefore, if $v_{i}<0$, the only condition is that (10) should be respected.

C3. $v_{i}=0$
Condition (12) can be satisfied if and only if $w_{i} \geq 0$. But, as previously, $w_{i}>0$ is not compatible with the fact that $u_{i}=0$. Therefore, it exist only one possible case, $w_{i}=$ 0 , i.e. $f b_{i}=m$.

In table 1 are summarized all conditions for obtaining $R_{i \text { max }} \leq R_{\text {adm }}$. It should be mentioned that for planar parallel robots, the reactions in the other joints that are not linked to the platform can be found using linear relationships with respect to $R_{i \max }$ (not detailed here).

Let us now consider an illustrative example.

## III. Illustrative example

In this part, it is studied the effort transmission in a 3$R \underline{R} R$ parallel manipulator (Fig. 2), with actuated joints located at $C_{i}$ (represented in grey) having the following characteristics:

- points $A_{1}, A_{2}$ and $A_{3}$ and points $B_{1}, B_{2}$ and $B_{3}$ form equilateral triangles
- circumcircle of $A_{1} A_{2} A_{3}$ of radius $r_{b}=0.35 \mathrm{~m}$
- circumcircle of $B_{1} B_{2} B_{3}$ of radius $r_{p}=0.1 \mathrm{~m}$
- links length $l_{A i C i}=l_{C i B i}=0.4 \mathrm{~m}$

For this robot, it is known that the directions of the forces $\mathbf{R}_{i}$ in the passive joints are directed along the vectors $\mathbf{A}_{i} \mathbf{B}_{i}$ (Fig. 2) [5]. Moreover, its Type 2 singular configurations appears when [5]:

- the orientation of the platform is equal to 0 or 180 deg, for any position of the platform centre
- when the platform centre, for a given orientation $\phi$, is located on a circle $C$ centred in $O$ of which radius is equal to $\sqrt{r_{b}^{2}+r_{p}^{2}-2 r_{b} r_{p} \cos \phi}$.
As a result, we will study the effort transmission in the passive joints of the robot in workspaces delimited by circles $C$..

For $3-R \underline{R} R$ robots, it can be shown that the matrix $\mathbf{A}^{T}$ expressed at point $B_{i}$ is equal to [5]:

$$
\mathbf{A}^{T}=\left[\begin{array}{lll}
\mathbf{R}_{1} & \mathbf{R}_{\mathbf{2}} & \mathbf{R}_{3}
\end{array}\right]=\left[\begin{array}{ccc}
\mathbf{f}_{1} & \mathbf{f}_{2} & \mathbf{f}_{\mathbf{3}}  \tag{31}\\
m_{1} & m_{2} & m_{3}
\end{array}\right]
$$

with

$$
\mathbf{f}_{i}^{T}=\left[\begin{array}{ll}
-\sin q_{i} & \cos q_{i}
\end{array}\right]
$$

and, if $B_{i}=B_{1}$ then

$$
\begin{align*}
& m_{1}=0 \\
& m_{2}=\sqrt{3} r_{p}\left[\begin{array}{c}
-\sin \phi \\
\cos \phi
\end{array}\right]^{T} \mathbf{f}_{2}  \tag{32a}\\
& m_{3}=\sqrt{3} r_{p}\left[\begin{array}{c}
-\sin (\pi / 3+\phi) \\
\cos (\pi / 3+\phi)
\end{array}\right]^{T} \mathbf{f}_{3}
\end{align*}
$$

if $B_{i}=B_{2}$ then

$$
\begin{align*}
& m_{1}=\sqrt{3} R_{p}\left[\begin{array}{c}
-\sin (\phi+\pi) \\
\cos (\phi+\pi)
\end{array}\right]^{T} \mathbf{f}_{1} \\
& m_{2}=0  \tag{32b}\\
& m_{3}=\sqrt{3} R_{p}\left[\begin{array}{c}
-\sin (2 \pi / 3+\phi) \\
\cos (2 \pi / 3+\phi)
\end{array}\right]^{T} \mathbf{f}_{3}
\end{align*}
$$

if $B_{i}=B_{3}$ then

$$
\begin{align*}
& m_{1}=\sqrt{3} R_{p}\left[\begin{array}{c}
-\sin (\phi+4 \pi / 3) \\
\cos (\phi+4 \pi / 3)
\end{array}\right]^{T} \mathbf{f}_{1} \\
& m_{2}=\sqrt{3} R_{p}\left[\begin{array}{c}
-\sin (-\pi / 3+\phi) \\
\cos (-\pi / 3+\phi)
\end{array}\right]^{T} \mathbf{f}_{2}  \tag{32c}\\
& m_{3}=0
\end{align*}
$$

The pressure angles and the distances to the instantaneous centre of rotation can be computed using the geometric representation shown in Fig. 2. In Fig. 4 are presented the variations of the joint reactions at point $B_{1}$ within the workspace for several platform orientations $\phi$, for $f=100 \mathrm{~N}$ and $m=5 \mathrm{Nm}$. Indeed, for this kind of robot, it can be shown that the reactions inside of the whole joints of leg $i$ are equal to $R_{i}$. Therefore studying the value of $R_{i \max }$ at $B_{i}$ allows having good information about the force transmission inside of leg $i$.

It can be observed that, the closer to Type 2 singularities, the higher the joint reaction. Moreover, the lowest values of the joint reactions appear in zones close to the workspace centre.

Let us now compute the force workspaces of the manipulator. The method consists in discretising the workspace using polar coordinates $(r, \theta)$. For one given angle $\theta$, the algorithm tests for all rising values of $r$ that the manipulator can support the applied wrench (see table 1). In the case where the manipulator cannot support the


Fig. 4. Variations of the joint reactions (in Newtons) at $B_{1}$ within the workspace for several platform orientations $\phi$, for $f=100 \mathrm{~N}$ and $m=$ 5 Nm .

(a) $\phi=15 \mathrm{deg}$

(b) $\phi=45 \mathrm{deg}$
$y(\mathrm{~m})$

(c) $\phi=\cos ^{-1}\left(r_{p} / r_{b}\right) \approx 73 \mathrm{deg}$

Fig. 5. Constant orientation workspace for several platform orientations $\phi$, for $f=100 \mathrm{~N}$ and $m=5 \mathrm{Nm}$ and $R_{a d m}=200 \mathrm{~N}$.

(a) $\phi=15 \mathrm{deg}$
$y$ (m)

(b) $\phi=45 \mathrm{deg}$
$y$ (m)

(c) $\phi=\cos ^{-1}\left(r_{p} / r_{b}\right) \approx 73 \mathrm{deg}$

Fig. 6. Constant orientation workspace for several platform orientations $\phi$, for $f=100 \mathrm{~N}$ and $m=5 \mathrm{Nm}$ and $R_{a d m}=130 \mathrm{~N}$.
applied wrench, the boundary of the force workspace is defined by the previous computational point.

In the following of this example, it is considered that the maximal admissible value $R_{a d m}$ of the reactions of the revolute joints should be equal to 200 N in one case, and to 130 N in a second case. The same wrench as previously is applied on the platform, i.e. a force of norm $f=100 \mathrm{~N}$, and a moment of norm $m=5 \mathrm{Nm}$. The shape of the force workspace, for one given assembly mode and for several values of platform orientation $\phi$ and of $R_{\text {adm }}$ is shown in Figs. 5 and 6.

It can be observed that, the smaller $\phi$, i.e. the closer from the orientation for which the robot is in singular configuration for any position of the platform centre, the smaller the workspace.

## V. Conclusions

In this paper has been presented a method that allows computing the maximal efforts in the joints of planar parallel manipulators. Trivial solution of this task could be the calculation of joint reactions for several values of the external wrench. However such a solution is very time consuming and do not give any quantitative or geometrical evaluation of the effort transmission. In the present paper a new approach has been developed, which has shown that for given external force and moment applied on the platform, the reactions in passive joints depend not only on the pressure angle but also the position of the instantaneous centre of rotation of the platform. This is the first time that such a kinetostatic property is rigorously formulated and clearly demonstrated. The obtained results allow expanding the knowledge about the criteria for measuring closeness to singularity and to avoid the mechanical breakdown by exceeding the admissible values of reactions in passive joints of planar parallel manipulators. In this aim, the optimal boundaries for the values of the pressure angle and of the distance to the instantaneous centre of rotation are computed. The $3-R \underline{R} R$ robot has been studied as an illustrative example. Its effort transmission has been studied as well as its maximal reachable workspace.

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