

Mixed Valence Tungsten (IV,V) Compounds with Layered Structures (Part III*): Synthesis and Crystal Structure of Ag1-x[W2O2Cl6] and Investigations on the Ag+ Ion Mobility

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Mixed Valence Tungsten (IV,V) Compounds with Layered Structures (Part III*): Synthesis and Crystal Structure of Ag_{1-x}[W₂O₂Cl₆] and Investigations on the Ag⁺ Ion Mobility

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Received.....

Dedicated to Prof. Bernd Harbrecht on the Occasion of his 60th Birthday

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Abstract. The reactions of metallic silver with tungsten tetrachlorideoxide WOCl₄ or the reaction of AgCl with WOCl₃, both at 690 K, lead to Ag_{1-x}[W₂O₂Cl₆] as black lustrous crystal needles. The compound shows a substantial phase field width in the silver content depending on reaction temperature and reaction time and crystals with x = 0, 0.38 and 0.82 were isolated. The crystal structure determinations (monoclinic, C2/m) show the structure to be isotypic to those of Tl[W₂O₂Cl₆] and $K_{0.84}[W_2O_2Cl_6]$ with the presence of 1-D polymeric $[W_2O_2Cl_6]_n$ strands and Ag⁺ ions with a variety in the occupation of the respective crystallographic site. Ag_{1-x}[W₂O₂Cl₆] represents a mixed valence compound with a variable ratio of W(IV) and W(V) in the W₂ dumbbells with a short W–W separation of 2.85 Å. In Fourier maps the electron density of the Ag ion appears smeared over a large area. By structure determinations of stochiometric Ag_{1.0}[W₂O₂Cl₆] at temperatures of 123, 193, and 293 K a double minimum potential was found with the silver ion dynamically disordered over two flat basins with a thermal activation barrier of 13 meV. No long range ion mobility is present since the Ag⁺ ions are trapped within a coordinating cage consisting of ten Cl atoms of the surrounding [W₂O₂Cl₆] strands forming a distorted bicapped cube. By MAS-NMR measurements on the ¹⁰⁹Ag nuclei at different temperature the spin lattice relaxation times were determined giving a thermally activated barrier of 15 meV. The electronic conductivity of pressed powder samples of Ag[W₂O₂Cl₆] by a four-probe measurement applying direct current is 1 Ω^{-1} cm⁻¹ at room temperature and $10^{-3} \Omega^{-1} \text{ cm}^{-1}$ at 25 K. For the conduction mechanism at low temperatures a variable range hopping mechanism is suggested while at higher temperatures a standard semi-conductivity with a band gap of 60 meV is present.

Keywords: Halooxotungstates; Mixed valence compounds; Nonstoichiometry; Ion mobility; ¹⁰⁹Ag solid state NMR; Small Gap Semiconductivity

Introduction

In preceding works we reported on the structure of WOCl₃ and the structures and properties of the chlorooxotungstates $M_x[W_2O_2Cl_6]$ with $M = Tl^+$, K^+ , Pb^{2+} , Hg^{2+} and the non-stochiometric phases Ag_{0.8}[W₂O₂Br₆] [1], Cu_{1-x}[W₂O₂Cl₆] and Cu_{1-x}[W₄O₄Cl₁₀] [2]. All compounds represent a family of mixed valence compounds with structures build of negatively charged infinite $W_2O_2X_6$ double strands consisting of WO_2X_4 octahedra condensed via two edge-sharing halogen atoms to form double octahedra and corner-sharing oxygen atoms to build 1-D strands. The mono or divalent cations are located between the strands and connect them to layers. The first representatives of this structure type were the bromooxotungstates Na[W₂O₂Br₆] [3] and the non-stochiometric Ag_{0.74}[W₂O₂Br₆] [4]. The structures of these compounds can in principle be understood as intercalates with WOX₃ (X = Cl, Br) acting as the host for atoms of the corresponding metals or, in the case of $Cu_{1-x}[W_4O_4Cl_{10}]$, as intermediates between WOX₃ and WOX₂. The tungsten(IV) dihalideoxides crystallize most probably isotypic to the structure of MoOCl₂ which has been investigated by *Thiele* et al. [5,6] despite the structure of WOCl₂ is not known in detail. WOCl₂ is a product of thermal decomposition of WOCl₃.

Several of the compounds investigated so far show non-stochiometric compositions with a phase field width pertaining the cations M, which are observed to underoccupy their positions. The compounds have been found to be weakly paramagnetic with temperature independent Pauli characteristics and show a remarkably high electrical conductivity of about 0.1 Ω^{-1} cm⁻¹ or higher at room temperature. The temperature dependence of the conductivity corresponds to semiconducting behaviour with a very small thermal activation barrier for the conduction process of about 0.03 eV at room temperature. The physical properties of the compounds can be understood by an interpretation as being one-dimensional metals. The real structure of the crystals, however, with a high number of discontinuances of the one-dimensional building units, limits the conductivity in close analogy to Krogmann's salts [7].

Here we report on Ag_{1-x}[W₂O₂Cl₆] as the silver containing representative of the $(M^+)_x$ [W₂O₂Cl₆] family of compounds and its remarkable properties.

Experimental

All educts and reaction products were handled in an argon filled glove box. Ampoules were made of borosilicate glass and flame-dried in vacuo prior to use. WOCl₄ was prepared from freshly precipitated WO₃, which was obtained from a NaWO₄ solution (Merck) by precipitation with concentrated hydrochloric acid. WO₃ was refluxed in SOCl₂ (Merck) and the orange solid was isolated by

filtration under argon. Before use, WOCl₄ was sublimated twice [8]. WOCl₃ was prepared by reduction of WOCl₄ with Al powder in evacuated borosilicate glass tubes at 500 K. AlCl₃ formed in the reaction was sublimated to the cold tip of the ampoule before opening in the glove box. AgCl was freshly prepared by precipitation from AgNO₃ and hydrochloric acid. It was handled in the dark and dried over P_4O_{10} .

Synthesis

 $Ag_{1-x}[W_2O_2Cl_6]$ was prepared by two ways, either by the reduction of WOCl₄ with metallic silver (Method A) or by reaction of WOCl₃ with AgCl (Method B), respectively.

Method A: For the reduction of WOCl₄ with Ag 1965 mg WOCl₄ (5.752 mmol) and 463 mg Ag powder (4.292 mmol) were filled into an ampoule of 15 cm length an 1.8 cm diameter, which was sealed, placed in a horizontal tube furnace and heated to a temperature of 690 K for a period of 7 days. The product was obtained as black needles with metallic luster. The crystals were separated manually for crystal structure analysis and solid state NMR investigations.

Method B: For the reaction of WOCl₃ with silver chloride 136 mg WOCl₃ (0.45 mmol) and 70 mg AgCl (0.49 mmol, dried in vacuo) and were filled together with a small amount (32 mg) WCl₆ (0.08 mmol) into an ampoule of 13 cm length and 6 mm diameter. This ampoule was placed in a horizontal tube furnace and heated starting from 510 K up to 700 K during a period of 40 days. Along the ampoule a temperature gradient of 20 K was present with the starting materials positioned in the hotter zone. At a temperature of 620 K crystallization of black needles started. The crystals grew on the surface of the sintered starting materials, chemical transport could not be observed. Crystals were separated manually and the residual ampoule content was grinded and sealed once more in an evacuated ampoule for further reaction at 690 K. A second crop of crystals could be separated manually.

Crystal structure determinations

Crystals from different synthetic runs were mounted in thin-walled glass capillaries, which were sealed. Data were collected using a Enraf-Nonius-Kappa-CCD diffractometer using graphite monochromatized Mo-K α radiation. Low temperature was maintained with a nitrogen gas flow cooling system (Oxford Cryosystems). Temperatures were calibrated with a micro-thermocouple mounted in the position of a crystal. Data sets were recorded at ambient temperature and additionally at

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 193 K and 123 K. Due to the high absorption coefficients and the anisotropic shape of the crystals a numerical absorption correction with the aid of the program HABITUS [10] was applied to all data sets. The optimization of the crystal shape was performed by minimization of the reliability factor of averaging symmetry equivalent reflections. Structure models were obtained by direct methods and refined against F^2 using SHELX97 [11] as implemented in the WINGX program package [12]. Tab. 1 contains the crystallographic data and details of the structure analyses, positional parameters are given in Table 2. Additionally, the crystal structures were refined against F^2 using the program JANA2000 [13] and the Gram-Charlier expansion to calculate anharmonic displacement tensors.

Further details of the crystal structure investigations have been deposited and may be obtained from the Fachinformationszentrum Karlsruhe, D-76344 Eggenstein-Leopoldshafen, Germany (e-mail: crysdata@fiz-karlsruhe.de), on quoting the deposition numbers CSD-421449 (293 K), CSD-421448 (193 K) and CSD-421447 (123 K) for Ag[W₂O₂Cl₆].

Nuclear magnetic resonance analysis

An amount of selected crystals obtained by method A was grinded and filled into a rotor and sealed by a vespel cap. The rotor was placed inside a 4 mm MAS probe (Chemagnetics) allowing for a maximum spinning rate of 18 kHz in the temperature range from 120 to 520 K. Data were collected with a 400 MHz spectrometer (Varian Infinity+). A 9 mol L^{-1} AgNO₃ solution spiked with 0.25 mol L^{-1} Fe(NO₃)₃ for acceleration of the relaxation process was used as reference.

Conductivity Measurements

A sample of Ag[W₂O₂Cl₆], prepared via method B, was grinded and pressed to a pellet of 6 mm diameter. The bulk sample was checked by X-ray powder diffractometry and EDX analysis for purity. Only marginal reflections of AgCl were present with a maximal relative intensity of 1 %. The silver content was found to be stochiometric by EDX. Conductivity measurements were performed by four-probe technique applying direct current within the temperature range from 10 K to 300 K, in both directions with cooling and heating.

Results and Discussion

Synthesis

The silver oxochlorotungstate $Ag[W_2O_2Cl_6]$ can be obtained by two different methods. In the presence of AgCl thermal decomposition of WOCl₃ occurs according to Eq. (A).

$$AgCl + 3 WOCl_3 \rightarrow Ag[W_2O_2Cl_6] + WOCl_4$$
(A)

Metallic silver reduces WOCl₄ according to Eq. (B) under formation of AgCl.

$$3 \operatorname{Ag} + 2 \operatorname{WOCl}_4 \to \operatorname{Ag}[W_2O_2Cl_6] + 2 \operatorname{AgCl}$$
(B)

While reaction A gives the product in higher purity, reaction B gives a higher yield according to the higher reactivity of WOCl₄. The higher purity in reaction A is a consequence of the high volatility of WOCl₄, which can be sublimated off from the product whereas AgCl as the side-product of reaction B is not volatile. Additionally, the thermal decomposition of WOCl₃ to WOCl₂ and WOCl₄ at the applied temperatures is a distracting rival to the desired reaction and thus lessens the yield. It can be suppressed by adding some WCl₆ or some excess WOCl₄ into the reaction ampoule. The two different methods of synthesis show a distinct influence on the composition of the products. As indicated by EDX analyses and the results of the single crystal analyses, crystals obtained by method A show constantly a silver tungsten ratio of 1:2, indicating stochiometric composition according to the formula Ag[W₂O₂Cl₆]. Crystals obtained by reaction B showed a variable silver tungsten ratio of 1:2, 0.66:2, 0.62:2 and 0.18:2, respectively, according to the site occupancy factors obtained by crystal structure determinations on crystals from different runs. The ratio 1:2, however, was only found in crystals obtained after opening the ampoules after the first reaction cycle, grinding the contents, re-sealing and heating for a second time, as described in the Experimental Part. The occurrence of such a large variety in the silver content is remarkable. The consequence for the charge balance of the crystal is the presence of a corresponding variety in the oxidation state of the tungsten atoms, ranging between +5 in WOCl₃ = $W_2O_2Cl_6$ and +4.5 in $(Ag^+)_{1,0}[W_2O_2Cl_6]$.

Crystal structure of $Ag_{1-x}[W_2O_2Cl_6]$

The crystal structures have been determined for several selected crystals and, for a crystal of the composition $Ag_{1,0}[W_2O_2Cl_6]$, at three different temperatures. All of these crystallize in the mono-

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clinic space group *C*2/*m* with structures isotypic to those of Tl[W₂O₂Cl₆] or K_{0.84}[W₂O₂Cl₆] [1]. The basic building blocks are $_{\infty}^{-1}$ [W₂O₂Cl₆] strands and Ag⁺ ions coordinated by terminal Cl atoms of the negatively charged W₂O_{4/2}Cl₆ groups (Fig. 1). The refined occupation factors of the Ag atoms reveal the large phase field width of the Ag content. The investigated crystals were found to have the empiric formulae Ag_{0.18}[W₂O₂Cl₆], Ag_{0.62}[W₂O₂Cl₆] and Ag_{1.0}[W₂O₂Cl₆], respectively. As a consequence, the mean oxidation states of the W atoms span almost the complete region between +4.5 and +5, which has to be interpreted as a dynamically intervalent state since all W atoms occupy the same crystallographic site 4*i*. Neighboured strands are bound together by Ag⁺ ions to layers in the *b-c* plane (Fig. 2).

Refinement of anisotropic displacement parameters for Ag^+ ions on $(0, \frac{1}{2}, \frac{1}{2})$ resulted in an extremely elongated displacement ellipsoid and a high residual electron density. Splitting the position of the Ag^+ ion over two closely neighboured sites (0.048, $\frac{1}{2}$, 0.501) did not solve the problem: though the residual electron density decreased clearly, no change in the elongation of the displacement ellipsoids of the Ag⁺ ions could be obtained. An electron density map in the (010) plane shows a broad region of electron density smeared along [100] with two closely neighboured maxima close to the special position 2d (2/m) at $(0, \frac{1}{2}, \frac{1}{2})$. An alternative model for the refinement was set up, implying four closely neighboured (symmetry generated seven) Ag sites using isotropic displacement parameters. The actual electron density distribution was described better by this structure model, even though strong correlations between parameters of the split Ag atoms occurred. To gain further insight into the phenomenon data sets at low temperature were recorded. At 123 K the site occupation factor of the 'central' Ag site in $(0, \frac{1}{2}, \frac{1}{2})$ became less than its error indicating a decreasing electron density in the central point. So this site was set to 'not occupied' on refinements based on the data collected at this temperature and the model contained six Ag sites (Fig. 1). The occupation factors were allowed for free refinement and finally reached a sum corresponding to the formula Ag_{1.02(6)}[W₂O₂Cl₆] (Tab. 2). In a third model, the electron density of the Ag ion was refined using an anharmonic displacement model with the Gram-Charlier expansion of the harmonic displacement tensor (program JANA2000 [13]). The Ag site was split into two symmetric equivalent positions close to $(0, \frac{1}{2}, \frac{1}{2})$, each of them refined with a displacement tensor of 3^{rd} order. The residual electron density decreased to an acceptable level but some correlations between parameters persisted. Nonetheless it seems to be a method inherent phenomenon that parameters of the displacement tensor are correlated [14]. Fig. 2 displays electron density charts of the Ag ion and a section along the maxima of the electron density at 295 K, 193 K, and 123 K, and the shape of the anhar-

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monic displacement parameters. At ambient temperature, the electron density of the Ag ion is blurred over a distance of nearly 4 Å with two flat maxima, representing preferred positions for the Ag ion. On decreasing temperature the two maxima of electron density persist but these maxima become sharper and the distance between the maxima increases from 1.24(2) Å at 295 K to 1.473(8) Å at 123 K. The Ag ions are actually located in a double-minimum potential with a saddle

point of low energy. According to Boltzmann's distribution theorem $\frac{N_{E+\Delta E}}{N_E} = e^{\frac{-\Delta E}{kT}}$, the electron

density diagram can be transferred into a potential diagram assuming that the electron density represents the population of the respective position and the observed maximal electron density as the population of the lowest energy state N_E . The result is shown at the bottom of Fig. 2 from which the activation energy for the thermally activated hopping of the Ag⁺ ion between the two most preferential positions at 123 K is estimated to 10 meV or 1 kJ mol⁻¹. This fits with the thermal energy kT, which corresponds to 11 meV at 123 K.

In the crystal structure, the Ag⁺ ions exhibit a distorted bicapped cuboidal coordination environment. The coordination polyhedron consists of eight Cl(2) and Cl(3) atoms with cuboidal shape, which is capped by two additional Cl(3) atoms (Fig. 1). This is caused by the lattice centering, which shifts the planar W_2Cl_6 groups alternately by half of the lattice constant in the b direction and brings two Cl(3) atoms of two neighbouring [W₂O₂Cl₆] strand in positions over two opposite faces of the cuboidal void in the crystallographic a direction. Thus the migration paths of the Ag⁺ ions are blocked and no long range ion mobility is possible since the structure of the anionic lattice does not offer non-intermittent migration paths for ionic conduction and only a local mobility is observed. The Ag-Cl distances amount 2.582(6) (Ag(4)-Cl(2)^{VII}), 2.653(6) (Ag(4)-Cl(3)), 2.976(6) Å (Ag(4)–Cl(3)^{VI}) for the different Ag sites and are comparable for example with the Ag–Cl distances in Ag₃YCl₆ (2.570 to 2.733 Å) [15]. Ag₃YCl₆ also contains disordered Ag ions with an elongated electron density inside a pair of face sharing octahedra. A comparable situation is observed for one of the Cu ions in the structure of tetrahedrite, $Cu_{12}Sb_4S_{13}$ [16]. The Cu ion is located in a trigonal bipyramidal coordination environment with three S atoms in the equatorial plane and two antimony atoms at the apexes. A dynamic disorder perpendicular to the triangular plane was observed by temperature dependent structure determinations.

From the structure analyses of the silver deficient crystals the tendency is present that on decreasing silver content an increasing W–W distance follows. With fully occupied silver sites in

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Ag_{1-x}[W₂O₂Cl₆], i.e. x = 0, a W–W distance of 2.8572(2) Å is found whereas for x = 0.38 this distance amounts 2.8631(5) Å and for x = 0.82 it enlarges to 2.8949(4) Å. The decreasing Ag⁺ content lowers the content of W(IV) with respect to W(V). The W–W distance thus enlarges to reach finally the value found for WOCl₃, containing only W(V) [1].

Nuclear magnetic resonance analysis

To gain further insight into the ionic motion process a nuclear resonance analysis was performed. MAS-NMR measurements on ¹⁰⁹Ag nuclei were performed to obtain the isotropic chemical shift of the silver ions. The MAS-NMR spectrum shows one signal at 232.0 ppm referenced to AgNO₃ (Fig. 3) This result is in accordance with the results of the X-ray single crystal structure analysis, which showed the equivalence of the two minima of the double minimum potential. Further measurements aimed on the longitudinal relaxation time T_1 . A sequence of pulses was used to make the magnetization fade out. After a pulse delay (*pd*) a rectangular pulse was set to excite ¹⁰⁹Ag nuclei. The free induction decay *FID* was plotted against *pd*. This procedure was performed at different temperatures of 265, 300, 330 and 355 K. By using the *Bloch* equation [17] the spin lattice relaxation time was calculated. The obtained relaxation times were plotted semi-logarithmic against the reciprocal values of the corresponding temperatures (Fig. 3). By fitting a linear slope to the data points, the Arrhenius plot showed a linear alignment following log $T_1 = 0.6908+175.96$ K \cdot T^{-1} . According to $\log T_1 = \log k_0 - \frac{E_A}{R}KT^{-1}$ the activation energy is calculated to 15.1 meV.

Electrical Conductivity

Like the others so far examined members of the $M_x[W_2O_2Cl_6]$ series, Ag[W_2O_2Cl_6] is a good electrical conductor. At room temperature, a specific conductivity of about 1 Ω^{-1} cm⁻¹ was found on a pressed powder sample. The temperature characteristics is semi-conducting, since the conductivity decreases with decreasing temperature and reaches about 10⁻³ Ω^{-1} cm⁻¹ at 25 K (Fig. 4).

From the Arrhenius equation for the activation energy of current transport in semiconductors one gets $\sigma(T) = \sigma_0 \cdot e^{-\frac{\Delta E}{2kT}}$. Accordingly, $-\ln \sigma = f(T^{-1})$ is expected as a linear function. As Fig. 4 shows,

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a linear function is actually present in the temperature range between 70 K and 300 K. From the slope the energy of the thermal activation is calculated to 0.0298(1) eV. This corresponds to a band gap of $\Delta E = 0.0596(2)$ eV. In the range below 70 K the activation energy decreases with decreasing temperature and a linear function is no longer present. This behaviour may be interpreted with a *variable range hopping* mechanism (*VRH*), describing the low temperature behaviour of the conductivity in disordered solids. The temperature dependence of the conductivity follows the expres-

sion $\sigma(T) = \sigma_0 \cdot e^{-\left(\frac{T_0}{T}\right)^{\alpha}}$ [18]. Initially developed for amorphous solids, the model is nowadays used also for crystalline semiconductors [19]. The exponent α is critical towards the dimensionality of the process and is assumed to be 0.5 for 1D systems [20]. As can be seen from Fig. 4, a satisfying fit to the conductivity data can be obtained with the exponent $\alpha = 0.53(2)$.

Conclusion

Ag_{1-x}[W₂O₂Cl₆] represents a further member of the A_{1-x} [W₂O₂X₆] family of compounds. A large phase field width has been achieved for the Ag containing compound, ranging from full occupation of the Ag⁺ site down to an occupation of merely 18 %. The Ag⁺ ions are mobile within their cuboidal coordination sphere in a double minimum potential and are dynamically disordered with an activation energy in the order of *k*T, shown by temperature dependent structure analyses and spin lattice relaxation times based on ¹⁰⁹Ag solid state NMR spectra. The crystal structure, however, does not allow for a long range mobility of the Ag⁺ ions, since there are no continuous migration paths for the Ag⁺ ions. Like all other members of the A_{1-x} [W₂O₂X₆] family of compounds, Ag[W₂O₂Cl₆] exhibits a remarkable high electrical conductivity with semiconducting characteristics and a small band gap of approximately 60 meV.

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Table 1 Crystallographic data and details of the crystal structure determination of Ag_{1.0}[W₂O₂Cl₆].
The refinement is based on the model for the Ag electron density by 4 (295 K und 193 K) or 3 (123 K) closely neighboured sites. Ag atoms are refined with isotropic displacement parameters. The standard deviations in brackets refer to the last quoted digit.

Formula Ag[W ₂ O ₂ Cl ₆]				
Crystal system, space group	monoclinic, $C 2/m$			
Temperature of measurement / K	295(2)	193(3)	123(3)	
Lattice parameters / Å	a = 12.9211(3) b = 3.7631(1) c = 9.8663(3) $\beta = 109.064(1)$	a = 12.8737(3) b = 3.7629(1) c = 9.8525(2) $\beta = 109.161(1)$	a = 12.8499(6) b = 3.7626(2) c = 9.8477(5) $\beta = 109.227(2)$	
Cell volume / Å ³	453.42(2)	450.84(2)	449.57(4)	
Number of formula units		2	I	
Calculated density /g cm ⁻³]	5.276	5.306	5.321	
Crystal size / mm ³	$0.019 \times 0.112 \times 0.018$			
Radiation	Mo-K $\overline{\alpha}$ ($\lambda = 0.71073$ Å)			
Data collection range / °	$6.7 \le 2\theta \le 55.0$	$6.7 \le 2\theta \le 55.0$	$9.1 \le 2\theta \le 54.9$	
min. / max. h k l	$-16 \le h \le 16$ $-4 \le k \le 4$ $-12 \le l \le 12$	$-16 \le h \le 16$ $-4 \le k \le 4$ $-12 \le l \le 12$	$-16 \le h \le 16$ $-4 \le k \le 4$ $-12 \le l \le 12$	
Number of measured reflections	5736	5303	2395	
Number of independent reflections	598	594	584	
Data quality R _{int} /R _{sigma}	0.051 / 0.019	0.058 / 0.025	0.088 / 0.056	
Number of refined parameters	46	46	44	
Relation reflections / parameters	13.0	12.9	13.3	
Absorption coefficient μ / cm ⁻¹	291.5	293.2	294.0	
Min. / max. transmission	0.2599 / 0.7014	0.2561 / 0.7672	0.2857 / 0.7640	
Quality of refinement R(F) $R(F) (n(F_o)>4\sigma(F_o))$ $wR(F^2)$	0.015 0.013 (566) 0.029	0.015 0.014 (577) 0.032	0.025 0.024 (562) 0.054	
max./min. residual electron density $[eÅ^{-3}]$	+0.92 / -0.87	+0.89 / -0.84	+1.30 / -1.49	

Table 2 Positional parameters and equivalent isotropic displacement factors *U* of the atoms in the structure of Ag[W₂O₂Cl₆] at 123 K. Values in brackets are standard deviations referring to the last quoted digit. All atoms are located on Wyckoff site 4*i* with site symmetry *m*. SOF is the respective site occupation factor for the split positions of the Ag atom. The sum of the occupation factors amounts $\sum SOF_{2Ag(2)+2Ag(3)+2Ag(4)} = 1.02(6)$

Atom	SOF	x	У	Ζ	$U_{ m eq}/{ m \AA}^2$
W		0.01132(3)	0	0.14890(3)	0.00698(18)
0		0.0096(5)	1/2	0.1466(5)	0.0085(13)
Cl(1)		0.15815(18)	0	0.0503(2)	0.0101(4)
Cl(2)		0.15555(18)	0	0.3716(2)	0.0141(5)
Cl(3)		-0.10774(19)	0	0.2928(2)	0.0151(4)
Ag(2)	0.13(2)	0.023(3)	1/2	0.4962(11)	0.016(2)
Ag(3)	0.265(17)	0.058(2)	1/2	0.5025(8)	0.0131(15)
Ag(4)	0.114(13)	0.099(2)	1/2	0.5240(15)	0.016(3)

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Legends for Figures

Fig. 1 Section of the chain of edge sharing pairs of $[WO_{2/2}Cl_2Cl_{2/2}]$ octahedra in the structure of $Ag_{1.0}[W_2O_2Cl_6]$ at 123 K (top), the unit cell of $Ag_{1.0}[W_2O_2Cl_6]$ in a view along the crystallographic *b* axis (middle, the Ag ions are represented by their harmonic displacement components U_{ij}), and the coordination of the dynamically disordered Ag ion by surrounding Cl atoms (bottom, the Ag ions are represented by their isotropic displacement components *U*). The displacement spheres and ellipsoids comply with a probability of localization of 90 %. Bond lengths /Å and angles / ° at 123 K: W–W^I 2.8501(4), W–O, W–O^{II} 1.8815(1), W–Cl(1) 2.390(2), W–Cl(1)^I 2.403(2), W–Cl(2) 2.404(2), W–Cl(3) 2.360(2), O–W–O^{II} 178.5(4), Cl(1)–W–Cl(1)^I 107.03(6), Cl(2)–W–Cl(3) 84.83(7), W–Cl(1)–W^I 72.97(6). Symmetry generated atoms are marked by superscript indices : I = -x, y, -z; II = x, -1+y, z; III = -x, y, 1-z; IV = x, y, 1-z; V = x, 1+y, z; VI = $\frac{1}{2}$ –x, $\frac{1}{2}$ +y, 1–z; VII = -x, 1+y, 1–z.

Fig. 2 Representation of the electron density of the Ag ion in the structure of $Ag_{1.0}[W_2O_2Cl_6]$ (left, based on Fourier coefficients F_0 , axis labellings refer to crystal coordinates, contour lines are graded by 2 e/Å³), sections of the electron density function through the maxima parallel to the crystallographic *c* axis (middle) and representations of the anharmonic displacements of the Ag atoms, drawn to include a probability density of 70 % (right). On bottom the respective double-minimum potential function for the Ag ion at 140 K is shown.

Fig. 3 109 Ag MAS NMR spectrum of Ag_{1.0}[W₂O₂Cl₆], recorded with a rotation frequency of 10 kHz (top) and the function of the spin lattice relaxation time at different temperatures versus the reciprocal temperature including the corresponding linear regression function (bottom).

Fig. 4 Electrical conductivity of a compressed sample of $Ag_{1.0}[W_2O_2Cl_6]$ as a function of temperature. The conductivity is shown on a logarithmic scale (top) and the respective Arrhenius function $-\ln \sigma = f(T^{-1})$ (bottom). Two different functions have been fitted to the data, a linear fit for the high temperature region (a) and a fit according the *VRH* model (b).

CI 1

0





Fig. 1











Fig. 4