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Convergence rate of the Allen-Cahn equation to generalized motion by mean curvature

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Abstract

We investigate the singular limit, as $\varepsilon \rightarrow 0$, of the Allen-Cahn equation $u_t^\varepsilon = \Delta u^\varepsilon + \varepsilon^{-2} f(u^\varepsilon)$, with f a balanced bistable nonlinearity. We consider rather general initial data u_0 that is independent of ε . It is known that this equation converges to the generalized motion by mean curvature — in the sense of viscosity solutions— defined by Evans, Spruck and Chen, Giga, Goto. However the convergence rate has not been known. We prove that the transition layers of the solutions u^ε are sandwiched between two sharp “interfaces” moving by mean curvature, provided that these “interfaces” sandwich at $t = 0$ an $\mathcal{O}(\varepsilon |\ln \varepsilon|)$ neighborhood of the initial layer. In some special cases, which allow both *extinction* and *pinches off* phenomenon, this enables to obtain an $\mathcal{O}(\varepsilon |\ln \varepsilon|)$ estimate of the location and the *thickness measured in space-time* of the transition layers. A result on the *regularity of the generalized motion by mean curvature* is also provided in the Appendix.

Key Words: Allen-Cahn equation, singular perturbation, generalized motion by mean curvature, viscosity solutions, localization and thickness of the transition layers. ¹ ²

1 Introduction

In this paper we study the behavior, as $\varepsilon \rightarrow 0$, of the solution $u^\varepsilon(x, t)$ of the Allen-Cahn type equation

$$(P^\varepsilon) \quad \begin{cases} u_t^\varepsilon = \Delta u^\varepsilon + \frac{1}{\varepsilon^2} f(u^\varepsilon) & \text{in } \mathbb{R}^N \times (0, \infty) \\ u^\varepsilon(x, 0) = u_0(x) & \text{in } \mathbb{R}^N, \end{cases}$$

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where $N \geq 2$. Here, the nonlinearity is given by $f(u) := -W'(u)$, where $W(u)$ is a double-well potential with equal well-depth, taking its global minimum value at $u = \alpha^*$, $u = \beta^*$. More precisely we assume that f is C^2 and has exactly three zeros $\alpha^* < a < \beta^*$ such that

$$f'(\alpha^*) < 0, \quad f'(a) > 0, \quad f'(\beta^*) < 0 \quad (\text{bistable nonlinearity}), \quad (1.1)$$

and that

$$\int_{\alpha^*}^{\beta^*} f(u) du = 0 \quad (\text{balanced case}). \quad (1.2)$$

The condition (1.1) implies that the potential $W(u)$ attains its local minima at $u = \alpha^*$, $u = \beta^*$, and (1.2) implies that $W(\alpha^*) = W(\beta^*)$. In other words, the two stable zeros of f , namely α^* and β^* , have “balanced” stability. A typical example is the cubic nonlinearity $f(u) = u(1 - u^2)$.

As for the initial data u_0 , we assume that it is bounded and of class C^2 on \mathbb{R}^N . Furthermore we define the “initial interface” Γ_0 by

$$\Gamma_0 := \{x \in \mathbb{R}^N : u_0(x) = a\},$$

and suppose that

$$\left\{ \begin{array}{l} \Gamma_0 \text{ is a smooth hypersurface without boundary of } \mathbb{R}^N, \\ \nabla u_0(x) \neq 0 \text{ for all } x \in \Gamma_0, \\ u_0 > a \text{ in } \Omega_0 \text{ and } u_0 < a \text{ in } (\Omega_0 \cup \Gamma_0)^c, \end{array} \right. \quad (1.3)$$

where Ω_0 denotes the region enclosed by Γ_0 . The non-zero gradient assumption in (1.3) is needed to obtain fine estimates for the development of steep transition layers at the very beginning period.

The evolution equations we consider decrease the following interfacial free energy of the Ginzburg-Landau type

$$\mathcal{E}(u) = \int_{\mathbb{R}^N} \frac{1}{2} |\nabla u|^2 + \frac{1}{\varepsilon^2} W(u),$$

where $\varepsilon > 0$ is a small parameter related to the thickness of a diffuse interfacial layer. Taking the gradient flow of \mathcal{E} with respect to the L^2 inner product leads to (P^ε) whose solutions fulfill

$$\frac{d}{dt} \mathcal{E}(u^\varepsilon) = - \int_{\mathbb{R}^N} (u_t^\varepsilon)^2 \leq 0.$$

Heuristics. As $\varepsilon \rightarrow 0$, a formal asymptotic analysis shows the following: in the very early stage, the diffusion term Δu^ε is negligible compared with the reaction term $\varepsilon^{-2} f(u^\varepsilon)$ so that, in the rescaled time scale $\tau = t/\varepsilon^2$, the equation is well approximated by the ordinary differential equation $u_\tau^\varepsilon = f(u^\varepsilon)$.

Hence, in view of the profile of f , the value of u^ε quickly becomes close to either β^* or α^* in most part of \mathbb{R}^N , creating a steep interface (transition layer) between the regions $\{u^\varepsilon \approx \alpha^*\}$ and $\{u^\varepsilon \approx \beta^*\}$ (*Generation of interface*). Once such an interface develops, the diffusion term becomes large near the interface, and comes to balance with the reaction term. As a result, the interface ceases rapid development and starts to propagate in a much slower time scale (*Motion of interface*).

Convergence to classical motion by mean curvature. The singular limit of the Allen-Cahn equation was first studied in the pioneering work of Allen and Cahn [2] and, slightly later, in Kawasaki and Ohta [19] from the point of view of physicists. They derived the interface equation by formal asymptotic analysis, thereby revealing that the interface moves by its mean curvature. More precisely, the limit solution $\tilde{u}(x, t)$ turns out to be a step function taking the value β^* on one side of the interface, and α^* on the other side. This sharp interface, which we will denote by Γ_t , obeys the following law of motion:

$$(P_{\text{classical}}^0) \quad \begin{cases} V_n = -(N-1)\kappa & \text{on } \Gamma_t \\ \Gamma_t|_{t=0} = \Gamma_0, \end{cases}$$

where V_n is the normal velocity of Γ_t in the exterior direction, κ the mean curvature at each point of Γ_t (chosen to be positive when Γ_t encloses a convex domain). If Γ_0 is smooth enough, it is well known that Problem $(P_{\text{classical}}^0)$ possesses locally in time a unique smooth solution. For more details, we refer to [12] and the references therein.

These early observations triggered a flow of mathematical studies aiming at rigorous justification of the above limiting procedure; see, for example, [20, 21], [9] and [10, 11] for results on the convergence of the partial differential equation (P^ε) to the free boundary Problem $(P_{\text{classical}}^0)$. Later, in [1], the authors prove an improved estimate for this convergence for solutions with general initial data. By performing a rigorous analysis of both the generation and the motion of interface, they show that the solution develops a steep transition layer within the time scale of $\mathcal{O}(\varepsilon^2 |\ln \varepsilon|)$, and that the layer obeys the law of motion that coincides with the formal asymptotic limit $(P_{\text{classical}}^0)$ within an error margin of $\mathcal{O}(\varepsilon)$ (previously, the best thickness estimate in the literature was of $\mathcal{O}(\varepsilon |\ln \varepsilon|)$, [10]).

Generalized motion by mean curvature. Nevertheless, it is well-known that the classical motion by mean curvature may develop singularities in finite time, even if Γ_0 is smooth. In \mathbb{R}^2 , an embedded curve evolving by its curvature can develop singularities only at the time of “shrinking to a point” [18]. In \mathbb{R}^3 , singularities may even occur before “shrinking to a point”: for instance, the boundary of a “dumbbell-shaped” region pinches off in finite time, if the neck is narrow enough. Therefore the classical framework is not

sufficient for dealing with such phenomena. Thus, one has to introduce a generalized notion of the *motion by mean curvature* (MMC in short). This enables us to define the MMC past the developments of singularities and then to study the singular limit of reaction-diffusion equations for all $t \geq 0$.

To define such a generalized MMC, the level set approach is quite convenient: one represents Γ_t as the level set of an auxiliary function which solves (in the viscosity sense) a nonlinear partial differential equation. This direct partial differential equation approach has been developed by Evans and Spruck [16], and, independently, by Chen, Giga and Goto [13]. In this framework, the involved partial differential equation is the degenerate, and even singular in the points where $Dv = 0$, parabolic problem given by

$$(P^0) \quad \begin{cases} v_t - \operatorname{tr} \left[(I - \widehat{Dv} \otimes \widehat{Dv}) D^2 v \right] = 0 & \text{in } \mathbb{R}^N \times (0, \infty) \\ v = d_0 & \text{in } \mathbb{R}^N \times \{t = 0\}, \end{cases}$$

with $\widehat{p} := \frac{p}{|p|}$, and d_0 the truncated signed distance function to Γ_0 , which is positive in the set $\{u_0 > a\}$ and negative in the set $\{u_0 < a\}$. If v is a viscosity solution of (P^0) , then each level set of v evolves according to the mean curvature in a certain generalized sense, and also in the classical sense whenever Dv does not vanish. Note that the equation can be written

$$v_t - \Delta v + D^2 v \widehat{Dv} \cdot \widehat{Dv} = 0,$$

or

$$v_t = |Dv| \operatorname{div} \left(\frac{Dv}{|Dv|} \right). \quad (1.4)$$

We refer to Section 3 for a short overview of the techniques and results of [16] and [13].

Convergence to generalized motion by mean curvature. Let us make a brief overview of known results on the convergence of the Allen-Cahn equation to generalized MMC. Evans, Soner and Souganidis [15] prove that, as $\varepsilon \rightarrow 0$, the solution of (P^ε) converges to β^* locally uniformly in $\{v > 0\}$ and to α^* locally uniformly in $\{v < 0\}$, where v is the solution of (P^0) . Since $\Gamma_t := \{x \in \mathbb{R}^N : v(x, t) = 0\}$ moves, in a weak sense, by mean curvature, this result is the natural generalization of the convergence to classical MMC mentioned above. Barles, Bronsard and Souganidis [5], Barles, Soner and Souganidis [7] generalized the result of [15] by allowing (x, t) -dependent nonlinearities and/or considering the unbalanced case instead of (1.2). Nevertheless, these early results consider only a very restricted class of initial data, namely those having a specific profile with well-developed transition layer. In other words the generation of interface from arbitrary initial data is not studied there.

Later, Soner [22, 23], Barles and Souganidis [8], Barles and Da Lio [6] study both the generation and the motion of interface; they prove the convergence of a large class of reaction-diffusion equations. By using the so-called

“open set approach”, the authors in [8] and [6] also provide a new definition for the global in time propagation of fronts; this definition turns out to be equivalent to the level set approach when there is no fattening of the interface.

From the above results, we know that the transition layers of u^ε converge to a level set of v , the solution of (P^0) , as $\varepsilon \rightarrow 0$, for all $t \geq 0$. However, none of these papers achieves a fine estimate of the convergence rate nor the thickness of the transition layers of the solutions to (P^ε) , for all $t \geq 0$. This is in contrast to the classical framework, for which $\mathcal{O}(\varepsilon |\ln \varepsilon|)$ or $\mathcal{O}(\varepsilon)$ estimates are known, as long as the limit interface remains smooth. The goal of the present paper is to obtain such fine estimates on the transition layers of solutions u^ε to Problem (P^ε) , as $\varepsilon \rightarrow 0$, for all $t \geq 0$.

Organization of the paper. In Section 2, we give our main results and make some comments, underlying the novelties. In Section 3, we recall the basic ideas of the level set approach together with some useful known properties. Section 4 is devoted to the construction of refined barriers (sub- and super-solutions) for the Allen-Cahn equation. By quoting a generation of interface result from [1] and using these barriers, we prove Theorem 2.1. In Section 5, we prove Theorem 2.5. In Appendix, we present some results on the *regularity of the generalized MMC* which, to the best of our knowledge, are not explicitly stated in the literature. They are related to our singular limit problem but are also interesting by themselves.

2 Main results and comments

2.1 Results

We consider the solution u^ε of the Allen-Cahn equation (P^ε) with initial data u_0 independent of ε . As mentioned before, u^ε quickly develops a steep transition layer. The time needed for such a generation (see Section 4) is

$$t^\varepsilon := f'(a)^{-1} \varepsilon^2 |\ln \varepsilon| \quad (\text{generation time}), \quad (2.1)$$

after which the transition layer starts to move approximately by the mean curvature. Theorem 2.1 is a first fine description of this motion of the Allen-Cahn transition layer: we show that it can be sandwiched between two sharp “interfaces” moving by mean curvature, provided that these “interfaces” sandwich at $t = 0$ an $\mathcal{O}(\varepsilon |\ln \varepsilon|)$ neighborhood of the initial layer.

In the sequel, we take two families of (not necessarily smooth) hypersurfaces without boundaries $(\gamma_{\varepsilon,0}^-)_{\varepsilon > 0}$, $(\gamma_{\varepsilon,0}^+)_{\varepsilon > 0}$, which sandwich an $\mathcal{O}(\varepsilon |\ln \varepsilon|)$ neighborhood of Γ_0 , and such that

$$\gamma_{\varepsilon,0}^- \ll \Gamma_0 \ll \gamma_{\varepsilon,0}^+,$$

where $\Gamma_1 \ll \Gamma_2$ means that Γ_1 is enclosed by Γ_2 and $\Gamma_1 \cap \Gamma_2 = \emptyset$. More precisely, we consider two families of open sets $(\omega_{\varepsilon,0}^-)_{\varepsilon>0}$, $(\omega_{\varepsilon,0}^+)_{\varepsilon>0}$, such that

$$\{x \in \mathbb{R}^N : \text{dist}(x, \omega_{\varepsilon,0}^-) \leq C_0 \varepsilon |\ln \varepsilon|\} \subset \Omega_0, \quad (2.2)$$

and

$$\{x \in \mathbb{R}^N : \text{dist}(x, \Omega_0) \leq C_0 \varepsilon |\ln \varepsilon|\} \subset \omega_{\varepsilon,0}^+, \quad (2.3)$$

for some constant $C_0 > 0$ not depending on ε and to be specified in (4.37). Then we define

$$\gamma_{\varepsilon,0}^\pm := \partial \omega_{\varepsilon,0}^\pm,$$

$$\{\gamma_{\varepsilon,t}^\pm\}_{t \geq 0} := \text{the generalized MMC starting from } \gamma_{\varepsilon,0}^\pm,$$

(see Section 3). In the same way as we define Ω_t as the “inside at time t ” of Γ_t in (3.5), we define $\omega_{\varepsilon,t}^\pm$ as the “space-time inside” of $\gamma_{\varepsilon,t}^\pm$ by replacing Γ_0 in (3.1) by $\gamma_{\varepsilon,0}^\pm$.

Theorem 2.1 (“Sandwiching” the Allen-Cahn layers). *Let $f \in C^2(\mathbb{R})$ satisfy (1.1) and (1.2), and let $u_0 \in C_b^2(\mathbb{R}^N)$ be such that (1.3) holds. Let $(\omega_{\varepsilon,0}^-)_{\varepsilon>0}$, respectively $(\omega_{\varepsilon,0}^+)_{\varepsilon>0}$, be any family of open sets satisfying (2.2), respectively (2.3). Let $\{\omega_{\varepsilon,t}^-\}_{t \geq 0}$ and $\{\omega_{\varepsilon,t}^+\}_{t \geq 0}$ be defined as above. Fix $\zeta \in (0, \min(a - \alpha^*, \beta^* - a))$ arbitrarily. Then, for $\varepsilon > 0$ small enough,*

$$\begin{cases} \alpha^* - \zeta \leq u^\varepsilon(x, t) \leq \beta^* + \zeta & \text{for all } x \in \mathbb{R}^N \\ \beta^* - \zeta \leq u^\varepsilon(x, t) \leq \beta^* + \zeta & \text{for all } x \in \omega_{\varepsilon,t}^- \cup \gamma_{\varepsilon,t}^- \\ \alpha^* - \zeta \leq u^\varepsilon(x, t) \leq \alpha^* + \zeta & \text{for all } x \notin \omega_{\varepsilon,t}^+, \end{cases} \quad (2.4)$$

for all $t \geq t^\varepsilon$, where t^ε is as in (2.1).

In the sequel, for $\alpha \in (\alpha^*, \beta^*)$, we define the sets

$$\Omega_t^\varepsilon(\alpha) := \{x \in \mathbb{R}^N : u^\varepsilon(x, t) > \alpha\},$$

$$\tilde{\Omega}_t^\varepsilon(\alpha) := \{x \in \mathbb{R}^N : u^\varepsilon(x, t) \geq \alpha\}.$$

Roughly speaking, given $\alpha < \beta$ in (α^*, β^*) , $\Gamma_t^\varepsilon(\alpha, \beta) := \tilde{\Omega}_t^\varepsilon(\alpha) \setminus \Omega_t^\varepsilon(\beta)$ represents the transition layer of the Allen-Cahn solution u^ε , namely the “zone” $\alpha \leq u^\varepsilon \leq \beta$. As an immediate consequence of Theorem 2.1 we can localize these sets in terms of $\omega_{\varepsilon,t}^-$, $\gamma_{\varepsilon,t}^-$ and $\omega_{\varepsilon,t}^+$.

Corollary 2.2 (“Sandwiching” the Allen-Cahn layers). *Let the assumptions of Theorem 2.1 hold. Fix $\alpha < \beta$ in (α^*, β^*) arbitrarily. Then, for $\varepsilon > 0$ small enough,*

$$(\omega_{\varepsilon,t}^- \cup \gamma_{\varepsilon,t}^-) \subset \Omega_t^\varepsilon(\beta) \subset \tilde{\Omega}_t^\varepsilon(\alpha) \subset \omega_{\varepsilon,t}^+, \quad (2.5)$$

for all $t \geq t^\varepsilon$, where t^ε is as in (2.1).

The statement (2.5) gives lower and upper estimates for the Allen-Cahn layer $\Gamma_t^\varepsilon(\alpha, \beta)$, but it does not necessarily give fine estimate for the location nor the thickness of the layer. To explain this, let $\{\Gamma_t\}_{t \geq 0}$ denote the generalized MMC starting from $\Gamma_0 = \partial\Omega_0$ (see Section 3) and define, as in (3.5), Ω_t as the “inside at time t ” of Γ_t . Assume (which is natural) that (2.2) and (2.3) are sharp — in the sense that $\gamma_{\varepsilon,0}^-$ and $\gamma_{\varepsilon,0}^+$ actually lie in an $\mathcal{O}(\varepsilon \ln \varepsilon)$ neighborhood of Γ_0 . Then, for every $t \geq 0$ the property

$$\lim_{\varepsilon \rightarrow 0} (\omega_{\varepsilon,t}^- \cup \gamma_{\varepsilon,t}^-) = \Omega_t \quad \text{and} \quad \lim_{\varepsilon \rightarrow 0} \omega_{\varepsilon,t}^+ = \Omega_t \cup \Gamma_t, \quad (2.6)$$

follows as an immediate consequence of the continuity of the viscosity solution of (P^0) with respect to the initial data (see [16] or [3]). Thus (2.5) implies, in particular, that, for any $\alpha < \beta$ in (α^*, β^*) ,

$$\limsup_{\varepsilon \rightarrow 0} \Gamma_t^\varepsilon(\alpha, \beta) \subset \Gamma_t. \quad (2.7)$$

However, no precise estimate of the convergence rate in (2.6) is known in general. Therefore Theorem 2.1 does not give good convergence rate for (2.7). Note that this is the case even if all the motions around Γ_t are *regular* — see (A.1) in Appendix. Nevertheless, as we explain below, explicit fine estimates can be derived for *admissible* initial domains.

Definition 2.3 (Admissible domains). Let Ω_0 be a bounded domain in \mathbb{R}^N whose boundary $\Gamma_0 := \partial\Omega_0$ is a smooth hypersurface without boundary. We say that Ω_0 is admissible if there exist $a_1 \geq 0$, $a_2 \geq 0$ and a skew-symmetric matrix Z such that $(a_1, a_2) \neq (0, 0)$ and that, for all $x \in \Gamma_0$,

$$(-a_1 x + Zx - a_2 H(x)n(x)) \cdot n(x) < 0, \quad (2.8)$$

where, for each $x \in \Gamma_0$, $H(x)$ denotes the sum of the principal curvatures (hence $N - 1$ times the mean curvature) at x , and $n(x)$ the unit outer normal to Γ_0 at x .

Remark 2.4. Assume Ω_0 is admissible. For $t \geq 0$, define the evolution operator $\Phi_t : \Gamma_0 \mapsto \Gamma_t$, where $\{\Gamma_t\}_{t \geq 0}$ denotes the generalized MMC starting from $\Gamma_0 = \partial\Omega_0$. Then, for $\nu \geq 0$, define

$$\psi_\nu(\Gamma_0) := e^{\nu Z} [e^{-a_1 \nu} \Phi_{a_2 \nu}(\Gamma_0)],$$

obtained by letting Γ_0 evolve by its mean curvature for the time $a_2 \nu$, then dilating by factor $e^{-a_1 \nu}$, and rotating by the matrix $e^{\nu Z} \in SO_n(\mathbb{R})$. Since $e^{\nu Z}$, $e^{-a_1 \nu}$, $\Phi_{a_2 \nu}$ are commutative, one can check that the collection $(\psi_\nu(\Gamma_0))_{\nu \geq 0}$ forms a semigroup whose infinitesimal generator evaluated at $x \in \Gamma_0$ is

$$G(x) := -a_1 x + Zx - a_2 H(x)n(x).$$

In view of (2.8) and the compactness of Γ_0 , there exists $\delta > 0$ such that, for all $x \in \Gamma_0$,

$$G(x) \cdot n(x) \leq -\delta.$$

It follows that, by choosing $\nu_0 > 0$ small if necessary,

$$0 \leq \nu < \nu' \leq \nu_0 \implies \psi_{\nu'}(\Gamma_0) \ll \psi_\nu(\Gamma_0), \quad (2.9)$$

and that there exists $\widehat{C}_0 > 0$ such that, for all $0 \leq \nu \leq \nu_0$,

$$\text{dist}(\psi_\nu(\Gamma_0), \Gamma_0) \geq \widehat{C}_0 \nu. \quad (2.10)$$

Before proceeding further, let us emphasize one important difference between the classical MMC and the generalized one. In the classical framework, one first defines $[0, T^{max})$ to be the maximal time-interval on which Γ_t remains smooth, and then picks up an arbitrary closed sub-interval $0 \leq t \leq T < T^{max}$, on which the derivatives of Γ_t remain uniformly bounded. In this time range one can get a good convergence rate for (2.6) for each $0 \leq t \leq T$, because of the smoothness of Γ_t . And we even have an optimal estimate as $d_{\mathcal{H}_{\mathbb{R}^N}}(\Gamma_t^\varepsilon(\alpha, \beta), \Gamma_t) = \mathcal{O}(\varepsilon)$, where $d_{\mathcal{H}_{\mathbb{R}^N}}$ denotes the Hausdorff distance (see [1], as mentioned before). However, in the generalized framework, such fine estimates collapse whenever Γ_t develops a singularity. For example, consider two dumbbell-shaped hypersurfaces $\gamma_{\varepsilon,0}^- \ll \Gamma_0$ whose Hausdorff distance $d_{\mathcal{H}_{\mathbb{R}^N}}(\gamma_{\varepsilon,0}^-, \Gamma_0)$ is very small, say of $\mathcal{O}(\varepsilon |\ln \varepsilon|)$. Within finite time the “neck” of the smaller dumbbell $\gamma_{\varepsilon,t}^-$ pinches off, splitting the hypersurface into two parts. Shortly after this moment, before Γ_t also pinches off, the Hausdorff distance $d_{\mathcal{H}_{\mathbb{R}^N}}(\gamma_{\varepsilon,t}^-, \Gamma_t)$ is rather large compared with the distance at $t = 0$.

Therefore, in order to get fine quantitative estimates in the presence of singularities, the spatial distance at each fixed time slice is not the right measurement to use. It turns out that, by using the space-time distance, we can overcome this difficulty, at least for admissible initial domains. For this purpose, we define the “space-time insides”

$$\omega_\varepsilon^\pm := \cup_{t \geq 0} (\omega_{\varepsilon,t}^\pm \times \{t\}), \quad \Omega := \cup_{t \geq 0} (\Omega_t \times \{t\}),$$

and the “space-time interface”

$$\Gamma := \{(x, t) \in \mathbb{R}^N \times [0, \infty) : v(x, t) = 0\}, \quad (2.11)$$

with v the viscosity solution of (P^0) .

When Ω_0 is admissible, there exists a constant $C > 0$ such that, for the approximating domains $(\omega_{\varepsilon,0}^-)_{\varepsilon > 0}$, $(\omega_{\varepsilon,0}^+)_{\varepsilon > 0}$ satisfying (2.2), and (2.3) respectively, we have

$$d_{\mathcal{H}_{\mathbb{R}^{N+1}}}(\gamma_\varepsilon^\pm, \Gamma) \leq C d_{\mathcal{H}_{\mathbb{R}^N}}(\gamma_{\varepsilon,0}^\pm, \Gamma_0),$$

where $d_{\mathcal{H}_{\mathbb{R}^{N+1}}}$ denotes the Hausdorff distance in the space-time \mathbb{R}^{N+1} (see Section 5 for details). Combining this and Theorem 2.1, we can obtain the following fine description of the Allen-Cahn layers for admissible initial domains.

Theorem 2.5 (Fine estimates for admissible initial domains). *Let $f \in C^2(\mathbb{R})$ satisfy (1.1) and (1.2), and let $u_0 \in C_b^2(\mathbb{R}^N)$ be such that (1.3) holds. Assume moreover that Ω_0 is admissible. Fix $\zeta \in (0, \min(a - \alpha^*, \beta^* - a))$ arbitrarily. Then, there exists $C > 0$ such that, for $\varepsilon > 0$ small enough, for all $x \in \mathbb{R}^N$ and all $t \geq t^\varepsilon$,*

$$u^\varepsilon(x, t) \in \begin{cases} [\alpha^* - \zeta, \beta^* + \zeta] & \text{if } (x, t) \in \mathbb{R}^N \times [t^\varepsilon, \infty) \\ [\beta^* - \zeta, \beta^* + \zeta] & \text{if } (x, t) \in \Omega \setminus \mathcal{N}_{C\varepsilon|\ln\varepsilon|}(\Gamma) \\ [\alpha^* - \zeta, \alpha^* + \zeta] & \text{if } (x, t) \in (\Omega \cup \Gamma)^c \setminus \mathcal{N}_{C\varepsilon|\ln\varepsilon|}(\Gamma) \end{cases} \quad (2.12)$$

where $\mathcal{N}_r(\mathcal{A}) := \{(x, t) \in \mathbb{R}^N \times [0, \infty) : \text{dist}((x, t), \mathcal{A}) < r\}$ denotes the r -neighborhood of the set \mathcal{A} in $\mathbb{R}^N \times [0, \infty)$, and t^ε is as in (2.1).

Now, for $\alpha < \beta$ in (α^*, β^*) , we define the “zone” $\alpha \leq u^\varepsilon \leq \beta$ by

$$\Gamma^\varepsilon(\alpha, \beta) := \{(x, t) \in \mathbb{R}^N \times [t^\varepsilon, \infty) : \alpha \leq u^\varepsilon(x, t) \leq \beta\},$$

which more or less represents the transition layer of the Allen-Cahn solution u^ε in space-time. Then the following holds as a direct consequence of Theorem 2.5.

Corollary 2.6 (Location and thickness of the layers). *Let $f \in C^2(\mathbb{R})$ satisfy (1.1) and (1.2), and let $u_0 \in C_b^2(\mathbb{R}^N)$ be such that (1.3) holds. Assume moreover that Ω_0 is admissible. Fix $\alpha < \beta$ in (α^*, β^*) arbitrarily. Then there exists $C > 0$ such that, for $\varepsilon > 0$ small enough,*

$$\Gamma^\varepsilon(\alpha, \beta) \subset \mathcal{N}_{C\varepsilon|\ln\varepsilon|}(\Gamma). \quad (2.13)$$

Note that (2.13) does not only give fine estimates for the location of the Allen-Cahn layer $\Gamma^\varepsilon(\alpha, \beta)$, but it also gives fine estimates for its thickness. Indeed, under the assumption of Ω_0 being admissible, Γ is known to have no interior [7, Theorem 4.3]; in other words, the so-called *fattening phenomenon* does not occur for Γ . Therefore (2.13) provides an $\mathcal{O}(\varepsilon|\ln\varepsilon|)$ estimate of the thickness of the Allen-Cahn layer.

Incidentally, if Ω_0 is admissible, not only Γ is known to have no interior, but it is also known to be *regular from inside*, that is, $\text{Cl}_{\mathbb{R}^{N+1}}[\Omega] = \Omega \cup \Gamma$, where $\text{Cl}_{\mathbb{R}^{N+1}}[A]$ denotes the closure of the set A in \mathbb{R}^{N+1} [17, Corollary 4.5.11]. In fact, one can also prove that Γ is *regular* both from inside and from outside if the initial domain is admissible (see Appendix).

Note that Theorem 2.5 allows Ω_t to pinch off, as will be clear from Example 2.9 below. Note also that Theorem 2.5 holds even after the extinction

time t^* of the solution of the MMC starting from the initial interface Γ_0 . Therefore, in the admissible case, we can provide an answer to the question: how quickly does u^ε approach α^* after the extinction of the interface Γ_t ?

Corollary 2.7 (Behavior after the extinction time). *Let $f \in C^2(\mathbb{R})$ satisfy (1.1) and (1.2), and let $u_0 \in C_b^2(\mathbb{R}^N)$ be such that (1.3) holds. Assume moreover that Ω_0 is admissible. Fix $\zeta \in (0, \min(a - \alpha^*, \beta^* - a))$ arbitrarily. Then, there exists $C > 0$ such that, for $\varepsilon > 0$ small enough,*

$$|u^\varepsilon(x, t) - \alpha^*| \leq \zeta, \quad (2.14)$$

for all $x \in \mathbb{R}^N$ and all $t \geq t^* + C\varepsilon|\ln \varepsilon|$, with $t^* > 0$ the extinction time defined in (3.9).

2.2 Examples of admissible domains

Here are some examples of admissible domains. In what follows Ω_0 is always assumed to be a bounded domain in \mathbb{R}^N with smooth boundary Γ_0 .

Example 2.8 (Strongly star-shaped domains). A domain Ω_0 is called strongly star-shaped with respect to the origin 0, if it is star-shaped with respect to 0, and if every ray emanating from 0 intersects Γ_0 transversely. The above condition is equivalent to $x \cdot n(x) > 0$, for all $x \in \Gamma_0$. Thus, any strongly star-shaped domain is admissible with $(a_1, a_2, Z) = (1, 0, 0)$.

Example 2.9 (Dumbbell-shaped domains). For $N \geq 3$, let Ω_0 consist of a pair of disjoint domains D_1, D_2 and a narrow channel D_3 connecting D_1, D_2 . For simplicity, we assume that Ω_0 is rotationally symmetric around the x_1 -axis and given in the form $\Omega_0 := \{0 \leq r < g(x_1)\}$, $r = (x_2^2 + \dots + x_N^2)^{1/2}$, where g is a function satisfying

$$g > 0 \text{ on } (-L, L), \quad g(\pm L) = 0, \quad g'(\pm L) = \mp \infty.$$

Furthermore, for some $0 < L_1 < L_2 < L$ and $0 < a < 1$,

$$\begin{cases} g(x_1) = 1 & \text{if } |x_1| \leq L_1 \\ g(x_1) = \cosh(a(|x_1| - L_1)) & \text{if } L_1 \leq |x_1| < L_2 \\ g''(x_1) < 0 & \text{if } L_2 \leq |x_1| < L. \end{cases}$$

We then modify g slightly around $|x_1| = L_1$ and $|x_1| = L_2$ so that g is smooth for $|x_1| < L$ and that Γ_0 is a smooth hypersurface (see Figure 1). Then we can easily check that, for $a \in (0, 1)$ small enough,

$$H(x) = \frac{1}{(1 + (g'(x))^2)^{1/2}} \left(\frac{N-2}{g(x)} - \frac{g''(x)}{1 + (g'(x))^2} \right) \geq \delta,$$

for some constant $\delta > 0$. Hence Ω_0 is admissible with $(a_1, a_2, Z) = (0, 1, 0)$. It is well known (see Angenent [4]) that the generalized MMC starting from

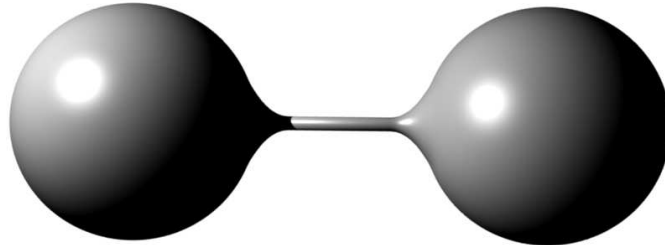


Figure 1: Dumbbell.

such Γ_0 pinches off and splits into two parts if L_1 and D_1, D_2 are large enough — thus creating a singularity .

A different type of dumbbell-shaped domain, which we call a diablo, can be constructed within the class of strongly star-shaped domains, that is $(a_1, a_2, Z) = (1, 0, 0)$, see Figure 2. The difference from the previous domain

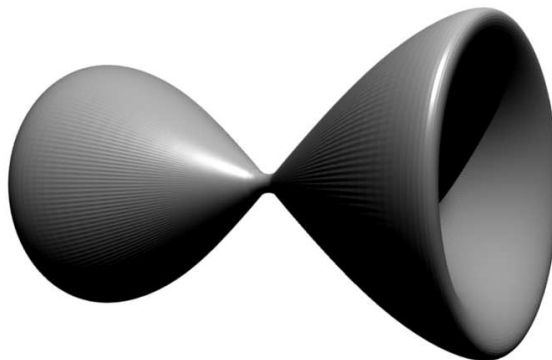


Figure 2: Diabolo.

is that the center neck, namely D_3 , cannot be too long; on the other hand, the outer end of D_1, D_2 need not to be of convex shape. This domain also

leads to pinching off if the center neck is narrow enough.

Example 2.10 (Galaxies). Here Ω_0 is a domain in \mathbb{R}^3 having the shape as in Figure 3. This is constructed by appropriately fattening the 2-dimensional

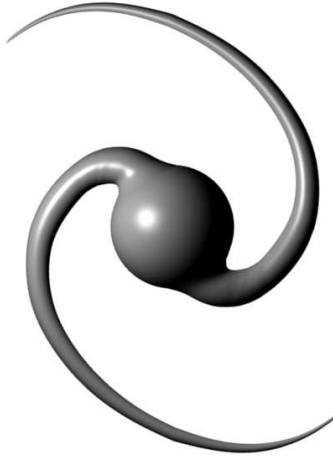


Figure 3: Galaxy.

skeleton in Figure 4, which consists of a disk at the center and two arms both of which are a portion of the logarithmic spiral of the form $r = e^{\beta(\theta-\theta_i)}$ ($i = 1, 2$), where $\beta > 0$, θ_1, θ_2 are some constants and $r = \sqrt{x^2 + y^2}$. It is easily seen that Ω_0 is admissible with

$$a_1 = 1, \quad a_2 = 0, \quad Z = \begin{bmatrix} 0 & -\beta^{-1} & 0 \\ \beta^{-1} & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}.$$

Example 2.11 (Gearwheels). Here, Ω_0 is a smooth 3-dimensional, but nearly flat, domain whose profile is as in Figure 5, with the origin 0 being the center of the inner circular hole, and with the z -axis perpendicular to this circle. The inner part of the boundary Γ_0 has positive mean curvature because of the large positive sectional curvature in the z direction, compared with the small negative sectional curvature in the rotational direction around the z -axis. The outer part is strongly star-shaped with respect to 0. With an appropriate combination of the outer shape and the size of the sectional curvature around the inner part, we see that Ω_0 is admissible with a suitable choice of $a_1 > 0$, $a_2 > 0$ and $Z = 0$.

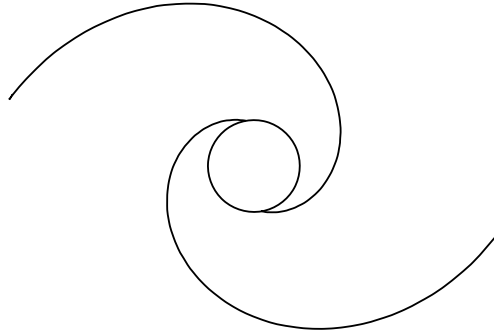


Figure 4: Skeleton of the galaxy.



Figure 5: Gearwheel.

3 Generalized motion by mean curvature

For the convenience of the reader, we briefly recall here the level set approach which enables to define uniquely a generalized MMC. We also recall some useful properties of the associated signed distance function. For more details and proofs, we refer to Evans and Spruck [16], Chen, Giga and Goto [13],

Evans, Soner and Souganidis [15] (from whom we borrow the notations) and the references therein.

Given a compact set $\Gamma_0 \subset \mathbb{R}^N$, we choose a continuous function $g : \mathbb{R}^N \rightarrow \mathbb{R}$, constant outside some ball and such that

$$\Gamma_0 = \{x \in \mathbb{R}^N : g(x) = 0\}. \quad (3.1)$$

Then, we consider the *mean curvature evolution partial differential equation*

$$v_t - \operatorname{tr} \left[(I - \widehat{D}v \otimes \widehat{D}v) D^2 v \right] = 0 \quad \text{in } \mathbb{R}^N \times (0, \infty), \quad (3.2)$$

which is nonlinear, degenerate and even undefined in the points where Dv vanishes (we recall that $\widehat{p} := \frac{p}{|p|}$). Nevertheless, Problem

$$(P^0) \quad \begin{cases} v_t - \operatorname{tr} \left[(I - \widehat{D}v \otimes \widehat{D}v) D^2 v \right] = 0 & \text{in } \mathbb{R}^N \times (0, \infty) \\ v = g & \text{in } \mathbb{R}^N \times \{t = 0\}, \end{cases}$$

admits a unique viscosity solution $v \in C(\mathbb{R}^N \times [0, \infty))$, constant outside some large enough ball, and each level set of v evolves according to the mean curvature in a generalized sense. As far as viscosity solutions are concerned, we refer the reader to the *User's guide* of Crandall, Ishii and Lions [14] and the references therein.

Now, for each $t \geq 0$, we define the “interface at time t ” by

$$\Gamma_t := \{x \in \mathbb{R}^N : v(x, t) = 0\}, \quad (3.3)$$

which is a compact set in \mathbb{R}^N . Then the collection $\{\Gamma_t\}_{t \geq 0}$ does not depend on the choice of the function g . The family $\{\Gamma_t\}_{t \geq 0}$ is called the *generalized motion by mean curvature* starting from Γ_0 . The “space-time interface” Γ , which is defined in (2.11), can be expressed as $\Gamma = \cup_{t \geq 0} (\Gamma_t \times \{t\})$.

Assume moreover that Γ_0 is the boundary of a bounded open set $\Omega_0 \subset \mathbb{R}^N$, and choose a continuous function g such that

$$g(x) > 0 \quad \text{if } x \in \Omega_0, \quad g(x) < 0 \quad \text{if } x \in (\Omega_0 \cup \Gamma_0)^c. \quad (3.4)$$

If v denotes the solution of (P^0) , we then define, for each $t \geq 0$, the “inside at time t ” by

$$\Omega_t := \{x \in \mathbb{R}^N : v(x, t) > 0\}, \quad (3.5)$$

which is an open set in \mathbb{R}^N . We also define the “space-time inside” by

$$\Omega := \{(x, t) \in \mathbb{R}^N \times [0, \infty) : v(x, t) > 0\} = \cup_{t \geq 0} (\Omega_t \times \{t\}). \quad (3.6)$$

For the generalized MMC, the following comparison principle is known to hold (see [16] or [3, Lemma 3.2]):

$$\gamma_0 \ll \tilde{\gamma}_0 \implies (\omega_t \cup \gamma_t) \subset \tilde{\omega}_t, \forall t \geq 0. \quad (3.7)$$

Here $\{\gamma_t\}_{t \geq 0}$, respectively $\{\tilde{\gamma}_t\}_{t \geq 0}$, denotes the generalized MMC starting from γ_0 , respectively $\tilde{\gamma}_0$, and ω_t , respectively $\tilde{\omega}_t$, denotes the “inside at time t ” of γ_t , respectively $\tilde{\gamma}_t$. As a consequence,

$$\gamma_0 \ll \tilde{\gamma}_0 \implies (\omega \cup \gamma) \subset \tilde{\omega}, \quad (3.8)$$

where ω , respectively $\tilde{\omega}$, denotes the “space-time inside” associated with γ_t , respectively $\tilde{\gamma}_t$, and γ the “space-time interface”.

Next, let t^* denote the extinction time, namely

$$t^* := \inf\{t > 0 : \Gamma_t = \emptyset\}. \quad (3.9)$$

Finally, we let $d(x, t)$ be the signed distance function to Γ_t , defined by

$$d(x, t) = \begin{cases} \text{dist}(x, \Gamma_t) & \text{if } x \in \Omega_t \\ 0 & \text{if } x \in \Gamma_t \\ -\text{dist}(x, \Gamma_t) & \text{if } x \in (\Omega_t \cup \Gamma_t)^c, \end{cases} \quad (3.10)$$

for all $x \in \mathbb{R}^N$, $0 \leq t \leq t^*$. Note that d is well-defined at time $t = t^*$ (because of the continuity of v), and that d may be not continuous in time (for instance, if Γ_t is made of two pieces, one “disappearing” before the other).

As proved in [15], the signed distance function d is a viscosity super-solution, sub-solution, of the heat equation in the set $\{d > 0\}$, $\{d < 0\}$ respectively. More precisely, the following holds.

Lemma 3.1. *We have*

$$d_t - \Delta d \geq 0 \quad \text{in } \Omega \cap (\mathbb{R}^N \times (0, t^*]), \quad (3.11)$$

and

$$d_t - \Delta d \leq 0 \quad \text{in } (\Omega \cup \Gamma)^c \cap (\mathbb{R}^N \times (0, t^*]), \quad (3.12)$$

these inequalities being understood in the viscosity sense.

4 Refined Allen-Cahn barriers

The goal of this section is to show that, for *any* generalized MMC $\{\gamma_t\}_{t \geq 0}$, there exist a super-solution of (P^ε) whose transition layer lies *inside* of γ_t within distance of $\mathcal{O}(\varepsilon |\ln \varepsilon|)$, and a sub-solution of (P^ε) whose transition layer lies *outside* of γ_t within distance of $\mathcal{O}(\varepsilon |\ln \varepsilon|)$. The results are stated

in the two following propositions which are fundamental for our analysis. Combining these propositions with a generation of interface result proved in [1], we will prove Theorem 2.1 in subsection 4.3.

Note that the constant $\lambda > 0$ which appears below is completely determined by the underlying travelling wave solution (see Lemma 4.4).

Proposition 4.1 (Super-solutions). *Let (γ_0, ω_0) be an arbitrary pair with γ_0 the boundary of the bounded open set $\omega_0 \subset \mathbb{R}^N$. Denote by $\{\gamma_t\}_{t \geq 0}$, $\{\omega_t\}_{t \geq 0}$ the associated “interface at time t ”, “inside at time t ” respectively. Denote by $d(x, t)$ the signed distance function to γ_t (see Section 3). Fix $\zeta > 0$ arbitrarily small and $T > 0$ arbitrarily. Then, for all $\varepsilon > 0$ small enough, there exists a function $w_\varepsilon^+(x, t)$ such that*

(i) w_ε^+ is a viscosity super-solution of the Allen-Cahn equation on $\mathbb{R}^N \times (0, T]$

(ii) w_ε^+ has, for all $t \geq 0$, the following upper bounds:

$$\begin{cases} w_\varepsilon^+(x, t) \leq \beta^* + \zeta & \text{for all } x \in \mathbb{R}^N \\ w_\varepsilon^+(x, t) \leq \alpha^* + \zeta & \text{for all } x \notin \omega_t \end{cases} \quad (4.1)$$

(iii) $w_\varepsilon^+(\cdot, 0)$ has the following lower bounds:

$$\begin{cases} \alpha^* + \frac{\zeta}{3} \leq w_\varepsilon^+(x, 0) & \text{for all } x \in \mathbb{R}^N \\ \beta^* + \frac{\zeta}{3} \leq w_\varepsilon^+(x, 0) & \text{for all } x \text{ such that } d(x, 0) \geq \frac{8}{\lambda} \varepsilon |\ln \varepsilon|. \end{cases} \quad (4.2)$$

Proposition 4.2 (Sub-solutions). *Let the notations of Proposition 4.1 hold. Fix $\zeta > 0$ arbitrarily small and $T > 0$ arbitrarily. Then, for all $\varepsilon > 0$ small enough, there exists a function $w_\varepsilon^-(x, t)$ such that*

(i) w_ε^- is a viscosity sub-solution of the Allen-Cahn equation on $\mathbb{R}^N \times (0, T]$

(ii) w_ε^- has, for all $t \geq 0$, the following lower bounds:

$$\begin{cases} \alpha^* - \zeta \leq w_\varepsilon^-(x, t) & \text{for all } x \in \mathbb{R}^N \\ \beta^* - \zeta \leq w_\varepsilon^-(x, t) & \text{for all } x \in \omega_t \cup \gamma_t \end{cases} \quad (4.3)$$

(iii) $w_\varepsilon^-(\cdot, 0)$ has the following upper bounds:

$$\begin{cases} w_\varepsilon^-(x, 0) \leq \beta^* - \frac{\zeta}{3} & \text{for all } x \in \mathbb{R}^N \\ w_\varepsilon^-(x, 0) \leq \alpha^* - \frac{\zeta}{3} & \text{for all } x \text{ such that } d(x, 0) \leq -\frac{8}{\lambda} \varepsilon |\ln \varepsilon|. \end{cases} \quad (4.4)$$

One of the role of such a pair of sub- and super-solution shall be to control the solution u^ε to (P^ε) during the latter time range — after the generation of interface— when the motion of interface occurs. In the sequel we prove Proposition 4.1, the proof of Proposition 4.2 being similar. We begin with some preparations.

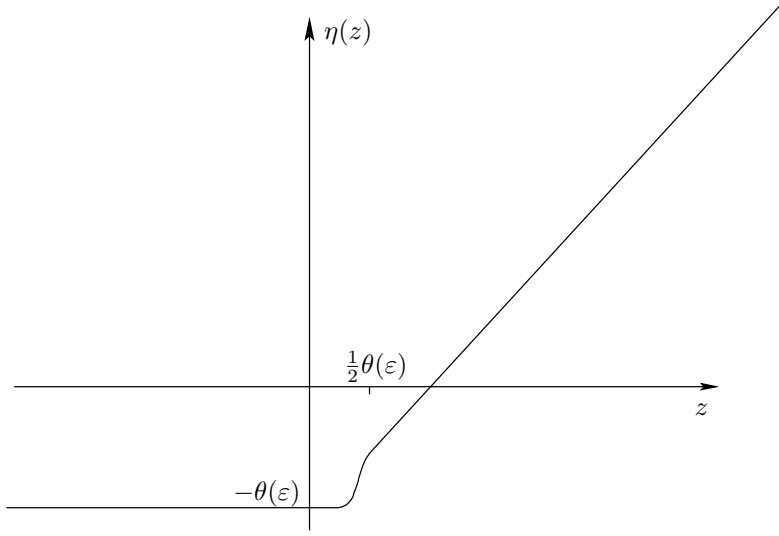


Figure 6: Graph of η .

4.1 Some preliminaries

A modified signed distance function. Let d be the signed distance function to an *arbitrary* generalized MMC. In order to construct super-solutions of (P^ε) involving the signed distance function, it is necessary to cut-off d in the set $\{d < 0\}$, where it is a sub-solution of the heat equation (Lemma 3.1). To that purpose, we slightly improve the cut-off argument used in [15].

In the following, $\theta(\varepsilon)$ is a positive function defined for $\varepsilon \in (0, \varepsilon_0)$, $\varepsilon_0 > 0$ small enough; its possible explicit forms will be indicated later. Consider a smooth auxiliary function $\eta = \eta^\varepsilon : \mathbb{R} \rightarrow \mathbb{R}$ satisfying

$$\left\{ \begin{array}{l} \eta(z) = -\theta(\varepsilon) \quad \text{for all } -\infty < z \leq \frac{1}{4}\theta(\varepsilon) \\ \eta(z) = z - \theta(\varepsilon) \quad \text{for all } z \geq \frac{1}{2}\theta(\varepsilon) \\ 0 \leq \eta' \leq C \\ |\eta''| \leq \frac{C}{\theta(\varepsilon)}, \end{array} \right. \quad (4.5)$$

where C is a constant independent of ε . Rather than d we shall use $\eta(d)$, for the construction of our super-solutions.

From [15, Lemma 3.1], there exists a constant $C > 0$ such that, for all $\varepsilon \in (0, \varepsilon_0)$,

$$\eta(d)_t - \Delta \eta(d) \geq -\frac{C}{\theta(\varepsilon)} \quad \text{in } \mathbb{R}^N \times (0, t^*], \quad (4.6)$$

and

$$\eta(d)_t - \Delta\eta(d) \geq 0 \quad \text{in } \left\{ d > \frac{1}{2}\theta(\varepsilon) \right\} \subset \mathbb{R}^N \times (0, t^*], \quad (4.7)$$

in the viscosity sense (which, in particular, contains the fact that $\eta(d)$ is lower semi continuous). For our results to hold after the extinction time, we need to extend $\eta(d(x, t))$ to all times. By abusing the notations slightly, we define

$$\eta(d(x, t)) = \begin{cases} \eta(d(x, t)) & \text{if } t \leq t^* \\ -\theta(\varepsilon) & \text{if } t > t^*. \end{cases} \quad (4.8)$$

Let us notice that, from the definition of t^* , $d(x, t^*) \leq 0$ for all $x \in \mathbb{R}^N$. But, as proved in [15], d is continuous from below (with respect to time) on $\mathbb{R}^N \times (0, t^*]$; therefore there exists a neighborhood of (x, t^*) in $\mathbb{R}^N \times (0, t^*]$ on which $\eta(d) \equiv -\theta(\varepsilon)$. As a consequence, it is obvious that (4.6)–(4.7) still hold for $t > t^*$, for the extension (4.8).

Lemma 4.3. *There exists a constant $C > 0$ such that, for all $\varepsilon \in (0, \varepsilon_0)$,*

$$\eta(d)_t - \Delta\eta(d) \geq -\frac{C}{\theta(\varepsilon)} \quad \text{in } \mathbb{R}^N \times (0, \infty), \quad (4.9)$$

and

$$\eta(d)_t - \Delta\eta(d) \geq 0 \quad \text{in } \left\{ d > \frac{1}{2}\theta(\varepsilon) \right\} \subset \mathbb{R}^N \times (0, \infty), \quad (4.10)$$

in the viscosity sense.

A standing wave. We shall also need $U_0(z)$ the unique solution of the stationary problem

$$\begin{cases} U_0'' + f(U_0) = 0 \\ U_0(-\infty) = \alpha^* \quad U_0(0) = a \quad U_0(+\infty) = \beta^*. \end{cases} \quad (4.11)$$

This solution represents the first approximation of the profile of a transition layer around the interface observed in the stretched coordinates; it naturally arises when performing a formal asymptotic expansion of the solution (see [1] and the references therein). Note that the “balanced stability assumption”, namely the integral condition (1.2), guarantees the existence of a solution of (4.11). In the simple case where $f(u) = u(1 - u^2)$, we know that $U_0(z) = \tanh(z/\sqrt{2})$. In the general case, the following standard estimates hold.

Lemma 4.4. *There exist positive constants C and λ such that*

$$\begin{aligned} 0 < \beta^* - U_0(z) &\leq Ce^{-\lambda|z|} \quad \text{for } z \geq 0 \\ 0 < U_0(z) - \alpha^* &\leq Ce^{-\lambda|z|} \quad \text{for } z \leq 0. \end{aligned}$$

In addition, U_0 is a strictly increasing function and, for $j = 1, 2$,

$$|D^j U_0(z)| \leq Ce^{-\lambda|z|} \quad \text{for } z \in \mathbb{R}. \quad (4.12)$$

4.2 Construction of super-solutions

We look for super-solutions w_ε^+ for Problem (P^ε) in the form

$$w_\varepsilon^+(x, t) = U_0 \left(\frac{\eta(d(x, t)) + \varepsilon p(t)}{\varepsilon} \right) + q(t), \quad (4.13)$$

for all $(x, t) \in \mathbb{R}^N \times [0, \infty)$, where

$$\begin{aligned} p(t) &= -e^{-\beta t/\varepsilon^2} + e^{Lt} + K \\ q(t) &= \sigma \left(\beta e^{-\beta t/\varepsilon^2} + \varepsilon^2 L e^{Lt} \right), \end{aligned} \quad (4.14)$$

and where d is the signed distance function to an *arbitrary* generalized MMC $\{\gamma_t\}_{t \geq 0}$. Note that by $\eta(d(x, t))$ we understand the extension (4.8).

Let us first specify the choice of β and σ and give a useful inequality. Note that these choices are reminiscent of the ones in [1] where the convergence to a classical solution of the MMC is studied. By assumption (1.1), there exist positive constants b, m such that

$$f'(U_0(z)) \leq -m \quad \text{if } U_0(z) \in [\alpha^*, \alpha^* + b] \cup [\beta^* - b, \beta^*]. \quad (4.15)$$

On the other hand, since the region $\{z \in \mathbb{R} : U_0(z) \in [\alpha^* + b, \beta^* - b]\}$ is compact and since $U_0' > 0$ on \mathbb{R} , there exists a constant $\delta_1 > 0$ such that

$$U_0'(z) \geq \delta_1 \quad \text{if } U_0(z) \in [\alpha^* + b, \beta^* - b]. \quad (4.16)$$

We set

$$\beta := \frac{m}{4}, \quad (4.17)$$

and

$$\sigma := \frac{\zeta}{2\beta}. \quad (4.18)$$

By reducing $\zeta > 0$ if necessary, we can assume $\sigma \leq \min(\sigma_0, \sigma_1, \sigma_2)$, where

$$\sigma_0 := \frac{\delta_1}{m + F_1} \quad \sigma_1 := \frac{1}{\beta + 1} \quad \sigma_2 := \frac{4\beta}{F_2(\beta + 1)}$$

$$F_1 := \|f'\|_{L^\infty(\alpha^*, \beta^*)} \quad F_2 := \|f''\|_{L^\infty(\alpha^* - 1, \beta^* + 1)}.$$

Combining (4.15) and (4.16), and considering that $\sigma \leq \sigma_0$, we obtain

$$U_0'(z) - \sigma f'(U_0(z)) \geq \sigma m \quad \text{for } -\infty < z < \infty. \quad (4.19)$$

We now turn to the proof of Proposition 4.1 (i).

Proof. For ease of notation we here denote w_ε^+ by w . Let $K > 1$ be arbitrary. What we shall prove is that, for all $\varepsilon \in (0, \varepsilon_0)$, the inequality

$$\mathcal{L}w := w_t - \Delta w - \frac{1}{\varepsilon^2} f(w) \geq 0 \quad \text{in } \mathbb{R}^N \times (0, T] \quad (4.20)$$

holds in the viscosity sense, provided that the constants $\varepsilon_0 > 0$ and $L > 0$ are appropriately chosen. Note that the remaining freedom for the choice of $K > 1$ is crucial for the proof of Proposition 4.1 (iii).

We recall that $\alpha^* < U_0 < \beta^*$ and go on under the following assumption

$$\varepsilon_0^2 L e^{LT} \leq 1. \quad (4.21)$$

Then, given any $\varepsilon \in (0, \varepsilon_0)$, since $\sigma \leq \sigma_1$, we have $0 \leq q(t) \leq 1$, so that

$$\alpha^* \leq w(x, t) \leq \beta^* + 1. \quad (4.22)$$

In order to prove (4.20), choose $\phi \in C^\infty(\mathbb{R}^N \times (0, \infty))$ such that

$$w - \phi \quad \text{has a minimum at } (x_0, t_0) \in \mathbb{R}^N \times (0, T]. \quad (4.23)$$

Subtracting if necessary a constant from ϕ we can assume that

$$w - \phi = 0 \quad \text{at point } (x_0, t_0). \quad (4.24)$$

What we have to prove is

$$\mathcal{L}\phi = \phi_t - \Delta\phi - \frac{1}{\varepsilon^2}f(\phi) \geq 0 \quad \text{at point } (x_0, t_0), \quad (4.25)$$

for all $\varepsilon \in (0, \varepsilon_0)$, with ε_0 small enough, L large enough, both independent on ϕ . In view of (4.24), we have

$$\phi(x_0, t_0) - q(t_0) = U_0 \left(\frac{\eta(d(x_0, t_0)) + \varepsilon p(t_0)}{\varepsilon} \right) \in (\alpha^*, \beta^*),$$

and one can define a smooth function ψ in a neighborhood of (x_0, t_0) by

$$\psi(x, t) := \varepsilon U_0^{-1}(\phi(x, t) - q(t)), \quad (4.26)$$

so that (4.23), (4.24) transfer to

$$\eta(d) - (\psi - \varepsilon p) \quad \text{has a minimum at } (x_0, t_0), \quad (4.27)$$

and

$$\eta(d) - (\psi - \varepsilon p) = 0 \quad \text{at point } (x_0, t_0). \quad (4.28)$$

It follows from Lemma 4.3 applied to test functions $\psi - \varepsilon p$ that

$$\psi_t - \Delta\psi \geq \varepsilon p_t - \frac{C}{\theta(\varepsilon)} \quad \text{at point } (x_0, t_0), \quad (4.29)$$

and

$$\psi_t - \Delta\psi \geq \varepsilon p_t \quad \text{at point } (x_0, t_0) \text{ if } d(x_0, t_0) > \frac{1}{2}\theta(\varepsilon) \text{ holds.} \quad (4.30)$$

Using $\phi = U_0(\frac{\psi}{\varepsilon}) + q$, we have the expansion

$$f(\phi) = f(U_0(\frac{\psi}{\varepsilon})) + qf'(U_0(\frac{\psi}{\varepsilon})) + \frac{1}{2}q^2f''(\theta),$$

for some $U_0 < \theta < U_0 + q$. In view of the ordinary differential equation (4.11), some straightforward computations yield, at point (x_0, t_0) ,

$$\mathcal{L}\phi = E_1 + E_2 + E_3,$$

with

$$E_1 = -\frac{1}{\varepsilon^2}q \left(f'(U_0) + \frac{1}{2}qf''(\theta) \right) + U_0'p_t + q_t$$

$$E_2 = \frac{U_0''}{\varepsilon^2}(1 - |\nabla\psi|^2)$$

$$E_3 = \frac{U_0'}{\varepsilon}(\psi_t - \Delta\psi - \varepsilon p_t).$$

The term E_1 . Plugging the expressions (4.14) for p and q in E_1 , we obtain

$$E_1 = \frac{\beta}{\varepsilon^2} e^{-\beta t/\varepsilon^2} (I - \sigma\beta) + Le^{Lt}(I + \varepsilon^2\sigma L),$$

where

$$I = U_0' - \sigma f'(U_0) - \frac{\sigma^2}{2} f''(\theta)(\beta e^{-\beta t/\varepsilon^2} + \varepsilon^2 Le^{Lt}).$$

In virtue of (4.19) and (4.22), we have

$$I \geq \sigma m - \frac{\sigma^2}{2} F_2(\beta + \varepsilon^2 Le^{LT}).$$

Combining this, (4.21) and the inequality $\sigma \leq \sigma_2$, we obtain $I \geq 2\sigma\beta$. Consequently, we have

$$E_1 \geq \frac{\sigma\beta^2}{\varepsilon^2} e^{-\beta t/\varepsilon^2} + 2\sigma\beta Le^{Lt}.$$

The term E_2 . First, assume $d(x_0, t_0) > \frac{1}{2}\theta(\varepsilon)$. From the definition of η , we have $\eta(d) = d - \theta(\varepsilon)$ in a neighborhood of (x_0, t_0) . Arguing as in the proof of [15, Theorem 2.2], we see that

$$|\nabla\psi(x_0, t_0)| = 1, \tag{4.31}$$

so that $E_2 = 0$.

Now assume $d(x_0, t_0) \leq \frac{1}{2}\theta(\varepsilon)$, which implies $\eta(d(x_0, t_0)) \leq -\frac{1}{2}\theta(\varepsilon)$. In view of statement (4.27) and the definition of η , we have $|\nabla\psi| \leq C$ at point (x_0, t_0) . We deduce from Lemma 4.4 that

$$|E_2| \leq \frac{C}{\varepsilon^2} e^{-\lambda|\eta(d)+\varepsilon p|/\varepsilon} \leq \frac{C}{\varepsilon^2} e^{-\lambda(\frac{1}{2}\theta(\varepsilon)-\varepsilon p)/\varepsilon}.$$

We remark that $0 < K - 1 \leq p \leq e^{LT} + K$. Consequently, if we assume

$$e^{LT} + K \leq \frac{\theta(\varepsilon)}{4\varepsilon}, \quad (4.32)$$

then

$$|E_2| \leq \frac{C}{\varepsilon^2} e^{-\lambda \frac{\theta(\varepsilon)}{4\varepsilon}}.$$

The term E_3 . If $d(x_0, t_0) > \frac{1}{2}\theta(\varepsilon)$ it directly follows from (4.30) that $E_3 \geq 0$.

Now assume $d(x_0, t_0) \leq \frac{1}{2}\theta(\varepsilon)$, which implies $\eta(d)(x_0, t_0) \leq -\frac{1}{2}\theta(\varepsilon)$. It follows from (4.29) that

$$E_3 \geq -\frac{C}{\theta(\varepsilon)} \frac{U_0'}{\varepsilon}.$$

Using again Lemma 4.4 and arguing as above for the term E_2 , we see that

$$E_3 \geq -\frac{C}{\theta(\varepsilon)} \frac{1}{\varepsilon} e^{-\lambda \frac{\theta(\varepsilon)}{4\varepsilon}}.$$

Assumptions on $\theta(\varepsilon)$. We now specify a possible choice for $\theta(\varepsilon)$. We assume that, as $\varepsilon \rightarrow 0$,

$$\theta(\varepsilon) |\ln \varepsilon| \leq C, \quad (4.33)$$

and

$$\frac{1}{\varepsilon^2} e^{-\lambda \frac{\theta(\varepsilon)}{4\varepsilon}} \leq C, \quad (4.34)$$

for some constant $C > 0$; we remark that the latter assumption implies that $\frac{\theta(\varepsilon)}{\varepsilon} \rightarrow \infty$ and that $\frac{1}{\theta(\varepsilon)} \frac{1}{\varepsilon} e^{-\lambda \frac{\theta(\varepsilon)}{4\varepsilon}}$ is also bounded (so that E_3 is bounded from below). In the following we select

$$\theta(\varepsilon) = \frac{8}{\lambda} \varepsilon |\ln \varepsilon|, \quad (4.35)$$

so that (4.33) and (4.34) hold. As easily understood, the above possible choice is related to our improved estimate of the convergence rate of the Allen-Cahn equation to generalized MMC (see Section 2).

Completion of the proof. Collecting all these estimates gives

$$\mathcal{L}\phi \geq \frac{\sigma\beta^2}{\varepsilon^2} e^{-\beta t/\varepsilon^2} + 2\sigma\beta L e^{Lt} - C \geq 2\sigma\beta L - C.$$

Now we set

$$L := \frac{1}{T} \ln \frac{\theta(\varepsilon_0)}{8\varepsilon_0}.$$

If ε_0 is chosen small enough, the assumptions on $\theta(\varepsilon)$ combined with the above choice for L validate assumptions (4.21) and (4.32) and insure $\mathcal{L}\phi \geq 0$. The proof of Proposition 4.1 (i) is now complete. \square

To conclude this subsection, we now prove Proposition 4.1 (ii) and (iii).

Proof. For ease of notation we here denote w_ε^+ by w .

In view of $\sigma\beta = \zeta/2$ and (4.21), we have, for $\varepsilon > 0$ small enough, $q(t) \leq \zeta$ for all $t \geq 0$. Hence $w(x, t) \leq \beta^* + \zeta$ holds true for all $x \in \mathbb{R}^N$. Next, choose $x \in \mathbb{R}^N$ such that $x \notin \omega_t$, that is $d(x, t) \leq 0$. In view of the graph of η we then have $\eta(d(x, t)) = -\theta(\varepsilon) = -\frac{\delta}{\lambda}\varepsilon|\ln \varepsilon|$. Therefore, for $t \geq 0$, we have

$$\begin{aligned} w(x, t) &= U_0\left(-\frac{\delta}{\lambda}|\ln \varepsilon| + p(t)\right) + q(t) \\ &\leq U_0\left(-\frac{\delta}{\lambda}|\ln \varepsilon| + e^{LT} + K\right) + \sigma(\beta + \varepsilon^2 L e^{LT}). \end{aligned}$$

Then it follows from $\sigma\beta = \zeta/2$ and from $U_0(-\infty) = \alpha^*$ that, for $\varepsilon > 0$ small enough (not depending on $x \notin \omega_t$), the inequality $w(x, t) \leq \alpha^* + \zeta$ holds true. The proof of Proposition 4.1 (ii) is now complete.

We now prove (iii). Since

$$w_\varepsilon^+(x, 0) = U_0\left(\frac{\eta(d(x, 0))}{\varepsilon} + K\right) + \frac{\zeta}{2} + \sigma\varepsilon^2 L, \quad (4.36)$$

is immediate that $w_\varepsilon^+(x, 0) \geq \alpha^* + \zeta/3$ for all $x \in \mathbb{R}^N$. Last, choose $K > 1$ large enough so that $U_0(K) \geq \beta^* - \frac{\zeta}{6}$. If x is such that $d(x, 0) \geq \frac{\delta}{\lambda}\varepsilon|\ln \varepsilon| = \theta(\varepsilon)$, the graph of η shows that $\eta(d(x, 0)) \geq 0$ so that $w_\varepsilon^+(x, 0) \geq U_0(K) + \frac{\zeta}{2} \geq \beta^* + \frac{\zeta}{3}$. Proposition 4.1 (iii) is proved. \square

Remark 4.5. As a matter of fact, our super-solutions w_ε^+ prove a bit more than (4.1) since $w_\varepsilon^+(x, t) \leq \alpha^* + \zeta$ is valid not only for $d(x, t) \leq 0$, but also for $d(x, t) \leq \frac{\delta}{\lambda}\varepsilon|\ln \varepsilon| - C\varepsilon$, with $C > 0$ large enough (because, in this case, $\eta(d(x, t)) \leq -C\varepsilon$); see Figure 8. \square

4.3 Proof of Theorem 2.1

Let $(\omega_{\varepsilon, 0}^+)_{\varepsilon > 0}$ be any family of open sets satisfying (2.3), with

$$C_0 > \frac{8}{\lambda}, \quad (4.37)$$

where $\lambda > 0$ is the constant that appears in Lemma 4.4. Fix $\zeta \in (0, \min(a - \alpha^*, \beta^* - a))$ arbitrarily. The strategy is the following. By quoting a generation of interface result from [1] and then using the super-solutions w_ε^+ associated with the pair $(\gamma_{\varepsilon, 0}^+, \omega_{\varepsilon, 0}^+) := (\partial\omega_{\varepsilon, 0}^+, \omega_{\varepsilon, 0}^+)$ (see Proposition 4.1), we will show that

$$\begin{cases} u^\varepsilon(x, t) \leq \beta^* + \zeta & \text{for all } x \in \mathbb{R}^N \\ u^\varepsilon(x, t) \leq \alpha^* + \zeta & \text{for all } x \notin \omega_{\varepsilon, t}^+, \end{cases} \quad (4.38)$$

for all $t \geq t^\varepsilon$, with t^ε the generation time that appears in (2.1). Since sub-solutions w_ε^- can be used in an analogous way, this will be enough to prove the theorem.

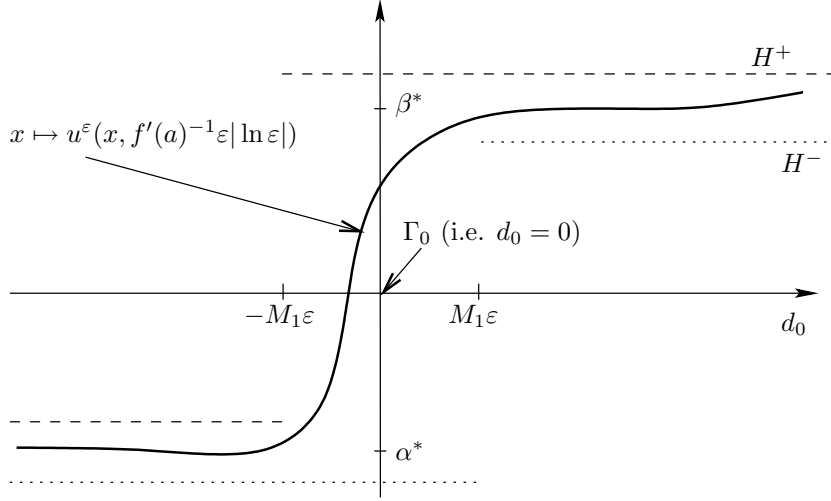


Figure 7: Prepared initial condition.

The rapid formation of internal layers that takes place in a neighborhood of $\Gamma_0 = \{x \in \mathbb{R}^N : u_0(x) = a\}$ is studied in [1]: from an arbitrary initial data $u_0 \in C_b^2(\mathbb{R}^N)$ satisfying (1.3), an interface is fully developed at time $t^\varepsilon := f'(a)^{-1}\varepsilon^2|\ln \varepsilon|$. In particular, there exists a positive constant M_1 such that, for $\varepsilon > 0$ small enough,

$$H^-(x) \leq u^\varepsilon(x, t^\varepsilon) \leq H^+(x), \quad (4.39)$$

for all $x \in \mathbb{R}^N$, where the functions $H^+(x), H^-(x)$ are given by

$$H^+(x) = \begin{cases} \beta^* + \frac{\zeta}{3} & \text{if } d_0(x) > -M_1\varepsilon \\ \alpha^* + \frac{\zeta}{3} & \text{if } d_0(x) \leq -M_1\varepsilon \end{cases}$$

$$H^-(x) = \begin{cases} \beta^* - \frac{\zeta}{3} & \text{if } d_0(x) \geq M_1\varepsilon \\ \alpha^* - \frac{\zeta}{3} & \text{if } d_0(x) < M_1\varepsilon, \end{cases}$$

with $d_0(x) := d(x, 0)$ the signed distance function associated with the hypersurface Γ_0 (see Figure 7).

For $T > 0$, we denote by w_ε^+ the super-solution associated with the pair $(\gamma_{\varepsilon,0}^+, \omega_{\varepsilon,0}^+)$ in the sense of Proposition 4.1. We denote by d^ε the signed distance function associated with $\{\gamma_{\varepsilon,t}^+\}_{t \geq 0}$, the generalized MMC starting from $\gamma_{\varepsilon,0}^+ := \partial\omega_{\varepsilon,0}^+$ (see Figure 8).

We claim that, for $\varepsilon > 0$ small enough,

$$H^+(x) \leq w_\varepsilon^+(x, 0), \quad (4.40)$$

for all $x \in \mathbb{R}^N$. In the range where $d_0(x) \leq -M_1\varepsilon$ this follows from (4.2). Now assume $d_0(x) > -M_1\varepsilon$. Since the constant C_0 which appears in (2.3)

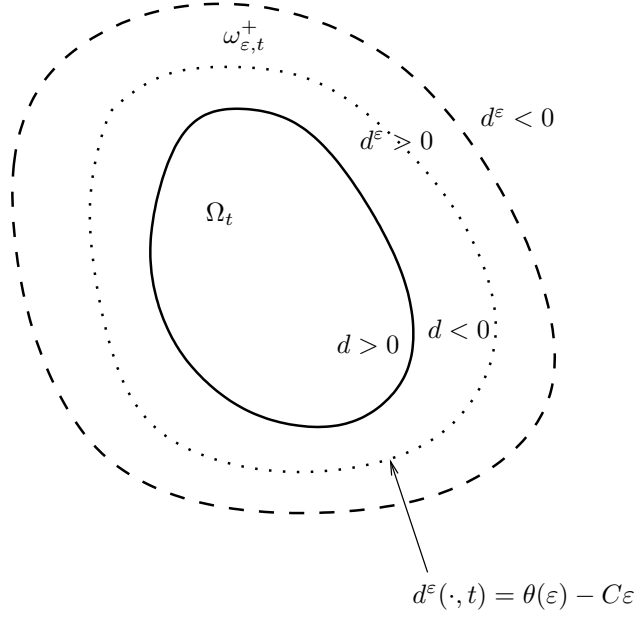


Figure 8: Ω_t , $\omega_{\varepsilon,t}^+$ and related signed distances.

is such that $C_0 > \frac{8}{\lambda}$, we see that $d^\varepsilon(x, 0) \geq \theta(\varepsilon) = \frac{8}{\lambda}\varepsilon|\ln \varepsilon|$, for $\varepsilon > 0$ small enough. Therefore, (4.2) implies (4.40).

From (4.40) and the comparison principle, we have

$$u^\varepsilon(x, t + t^\varepsilon) \leq w_\varepsilon^+(x, t) \quad \text{for } 0 \leq t \leq T - t^\varepsilon. \quad (4.41)$$

From this and (4.1) (with $\omega_{\varepsilon,t}^+$ playing the role of ω_t) we immediately infer that, for all $T > 0$, (4.38) is true on the time interval $[t^\varepsilon, T]$. If we choose $T > 0$ large enough so that $\omega_{\varepsilon,t}^+ = \emptyset$ for all $t \geq T$ (that is the generalized MMC starting from $\gamma_{\varepsilon,0}^+$ have become extinct), we see that $u^\varepsilon(x, T) \leq \alpha^* + \zeta$ for all $x \in \mathbb{R}^N$; the comparison principle then shows that this inequality persists for all $t \geq T$ and thus that (4.38) remains true on the time interval $[t^\varepsilon, \infty)$. \square

5 Proof of Theorem 2.5

In this section, we prove Theorem 2.5. Assume Ω_0 is admissible in the sense of Definition 2.3. Fix $\zeta > 0$ arbitrarily small. After making two key observations in subsection 5.1, we split the proof into the lower bounds and the upper bounds appearing in Theorem 2.5.

5.1 Two key observations

A first observation is that the mean curvature evolution partial differential equation is invariant under time-shifts, dilations, rotations. More precisely if $v(x, t)$ solves

$$v_t = |Dv| \operatorname{div} \left(\frac{Dv}{|Dv|} \right), \quad (5.1)$$

so do $v(x, t + s)$ with $s \geq 0$, $v(\lambda x, \lambda^2 t)$ with $\lambda > 0$ and $v(Rx, t)$ with $R \in SO_n(\mathbb{R})$.

Next, for $\nu \geq 0$, define the invertible map $\Pi_\nu : \mathbb{R}^N \times \mathbb{R} \rightarrow \mathbb{R}^N \times \mathbb{R}$ by

$$\Pi_\nu((x, t)) := (e^{\nu Z} e^{-a_1 \nu} x, e^{-2a_1 \nu} (t - a_2 \nu)). \quad (5.2)$$

For a given compact set $K \subset \mathbb{R}^N \times \mathbb{R}$ it is obvious that there exists $C_K > 0$ such that, for all $\nu \geq 0$,

$$\sup_{(x,t) \in K} \|(\Pi_\nu - Id)(x, t)\| \leq C_K \nu, \quad (5.3)$$

$$\sup_{(x,t) \in K} \|(\Pi_\nu^{-1} - Id)(x, t)\| \leq C_K \nu. \quad (5.4)$$

5.2 The lower bounds

We first prove the lower bounds that appear in Theorem 2.5.

For $\varepsilon \geq 0$ small enough, we construct “inner approximations of Γ_0 ” by

$$\gamma_{\varepsilon,0}^- := \psi_{\varepsilon|\ln \varepsilon|}(\Gamma_0) = e^{\varepsilon|\ln \varepsilon|Z} \left[e^{-a_1 \varepsilon|\ln \varepsilon|} \Phi_{a_2 \varepsilon|\ln \varepsilon|}(\Gamma_0) \right], \quad (5.5)$$

see Remark 2.4. Here we remark that

$$\gamma_{\varepsilon,0}^- = P(\Pi_{\varepsilon|\ln \varepsilon|}(\Gamma_0)), \quad (5.6)$$

where

$$P : (x, t) \mapsto x$$

is the projection from \mathbb{R}^{N+1} onto \mathbb{R}^N . We then define $\omega_{\varepsilon,0}^-$ as the domain enclosed by $\gamma_{\varepsilon,0}^-$. From Remark 2.4, we deduce that

$$0 \leq \varepsilon < \varepsilon' \implies \gamma_{\varepsilon',0}^- \ll \gamma_{\varepsilon,0}^-, \quad (5.7)$$

and

$$\operatorname{dist}(\gamma_{\varepsilon,0}^-, \Gamma_0) \geq \widehat{C}_0 \varepsilon |\ln \varepsilon|, \quad (5.8)$$

for some constant $\widehat{C}_0 > 0$. By replacing if necessary $\psi_{\varepsilon|\ln \varepsilon|}(\Gamma_0)$ in (5.5) by $\psi_{C\varepsilon|\ln \varepsilon|}$ with $C \gg 1$, we can assume $\widehat{C}_0 > \frac{8}{\lambda}$, so that (2.2) is satisfied with $C_0 = \widehat{C}_0$. Therefore the lower bounds in Theorem 2.1 hold.

Next, it follows from (5.7) and the comparison principle (3.8) that $(\omega_\varepsilon^- \cup \gamma_\varepsilon^-) \subset \Omega$. Since (see proof below)

$$d_{\mathcal{H}_{\mathbb{R}^{N+1}}}(\gamma_\varepsilon^-, \Gamma) \leq C\varepsilon |\ln \varepsilon|, \quad (5.9)$$

for some $C > 0$, it follows that

$$\Omega \setminus \mathcal{N}_{C\varepsilon |\ln \varepsilon|}(\Gamma) \subset (\omega_\varepsilon^- \cup \gamma_\varepsilon^-), \quad (5.10)$$

so that the lower bounds in Theorem 2.5 follow from the ones in Theorem 2.1. More precisely, if $(x, t) \in \Omega \setminus \mathcal{N}_{C\varepsilon |\ln \varepsilon|}(\Gamma)$, we deduce from (5.10) and the lower bounds in Theorem 2.1 that $u^\varepsilon(x, t) \geq \beta^* - \zeta$.

It remains to prove (5.9). We use the observations made in subsection 5.1. The function

$$v_\varepsilon^-(x, t) := v(\Pi_{\varepsilon |\ln \varepsilon|}^{-1}(x, t)) \quad (5.11)$$

solves (5.1), and $v_\varepsilon^-(x, 0) = 0$ if and only if $x \in \gamma_{\varepsilon, 0}^-$. Hence v_ε^- “describes” the generalized MMC $\{\gamma_{\varepsilon, t}^-\}_{t \geq 0}$ starting from $\gamma_{\varepsilon, 0}^-$. Therefore the “space-time interface” γ_ε^- is given by

$$\gamma_\varepsilon^- = \{(x, t) \in \mathbb{R}^N \times [0, \infty) : v_\varepsilon^-(x, t) = 0\}.$$

In view of (5.11) this yields (see Figure 9)

$$\gamma_\varepsilon^- = \Pi_{\varepsilon |\ln \varepsilon|}(\Gamma) \cap (\mathbb{R}^N \times [0, \infty)). \quad (5.12)$$

Therefore (5.9) follows from (5.12) and (5.3).

5.3 The upper bounds

We now prove the upper bounds that appear in Theorem 2.5. Since the evolution operator $\Phi_t : \Gamma_0 \mapsto \Gamma_t$ is not invertible, the argument is more involved.

First, choose $\varepsilon > 0$ small enough so that the MMC starting from Γ_0 remains smooth on the time interval $[0, \tau^\varepsilon]$, where

$$\tau^\varepsilon := a_2 \varepsilon |\ln \varepsilon|.$$

From the classical framework analysis [1, Theorem 1.3], there exists $M > 0$ such that, for all $t^\varepsilon \leq t \leq \tau^\varepsilon$,

$$u^\varepsilon(x, t) \in \begin{cases} [\alpha^* - \frac{\zeta}{3}, \alpha^* + \frac{\zeta}{3}] & \text{if } d(x, t) \leq -M\varepsilon \\ [\alpha^* - \frac{\zeta}{3}, \beta^* + \frac{\zeta}{3}] & \text{if } -M\varepsilon < d(x, t) < M\varepsilon \\ [\beta^* - \frac{\zeta}{3}, \beta^* + \frac{\zeta}{3}] & \text{if } d(x, t) \geq M\varepsilon, \end{cases} \quad (5.13)$$

with $d(x, t)$ the signed distance function to Γ_t defined in (3.10). We recall that $t^\varepsilon := f'(a)^{-1} \varepsilon^2 |\ln \varepsilon|$ denotes the generation time.

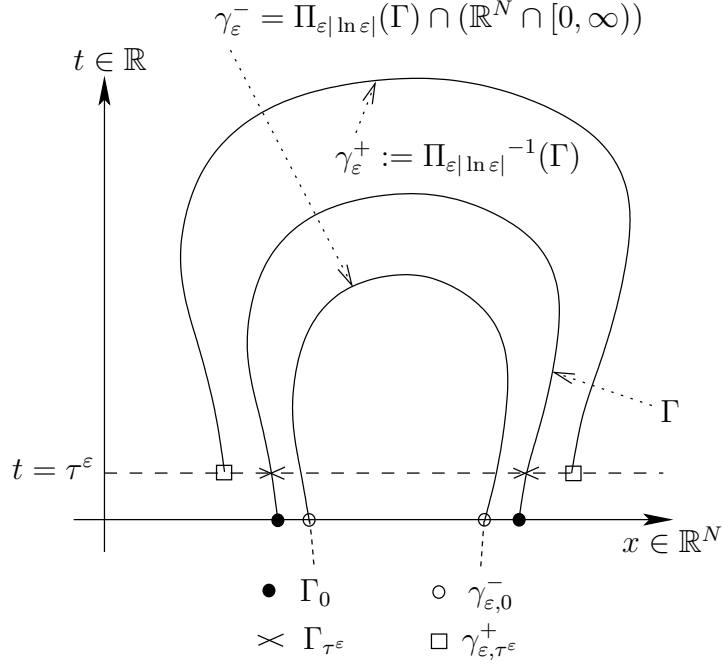


Figure 9: Sets used in the proof of Theorem 2.5.

Next, using the map $\Pi_{\varepsilon|\ln\varepsilon|}^{-1}$, we define the space-time sets (see Figure 9)

$$\gamma_\varepsilon^+ := \Pi_{\varepsilon|\ln\varepsilon|}^{-1}(\Gamma), \quad \text{and} \quad \omega_\varepsilon^+ := \Pi_{\varepsilon|\ln\varepsilon|}^{-1}(\Omega).$$

From (5.4) we deduce that there exists $C > 0$ such that

$$d_{\mathcal{H}_{\mathbb{R}^{N+1}}}(\gamma_\varepsilon^+, \Gamma) \leq C\varepsilon|\ln\varepsilon|. \quad (5.14)$$

Since the function

$$v_\varepsilon^+(x, t) := v(\Pi_{\varepsilon|\ln\varepsilon|}(x, t)) \quad (5.15)$$

solves (5.1) for $t \geq \tau^\varepsilon$, and $v_\varepsilon^+(x, \tau^\varepsilon) = 0$ if and only if $x \in e^{-\varepsilon|\ln\varepsilon|Z} e^{a_1\varepsilon|\ln\varepsilon|}\Gamma_0$, the set γ_ε^+ is actually the “space-time interface” associated with the generalized MMC $\{\gamma_{\varepsilon,t}^+\}_{t \geq \tau^\varepsilon}$ starting from

$$\gamma_{\varepsilon,\tau^\varepsilon}^+ := e^{-\varepsilon|\ln\varepsilon|Z} e^{a_1\varepsilon|\ln\varepsilon|}\Gamma_0. \quad (5.16)$$

Hence Proposition 4.1 provides an Allen-Cahn super-solution w_ε^+ on $(\tau^\varepsilon, T]$ such that, for all $t \geq \tau^\varepsilon$,

$$w_\varepsilon^+(x, t) \leq \alpha^* + \zeta \quad \text{for all } x \notin \omega_{\varepsilon,t}^+, \quad (5.17)$$

with $\omega_{\varepsilon,t}^+$ the “inside at time t ” associated with $\gamma_{\varepsilon,t}^+$, and

$$\begin{cases} \alpha^* + \frac{\zeta}{3} \leq w_\varepsilon^+(x, \tau^\varepsilon) & \text{for all } x \in \mathbb{R}^N \\ \beta^* + \frac{\zeta}{3} \leq w_\varepsilon^+(x, \tau^\varepsilon) & \text{for all } x \text{ such that } d^\varepsilon(x, \tau^\varepsilon) \geq \frac{8}{\lambda}\varepsilon|\ln\varepsilon|, \end{cases} \quad (5.18)$$

with $d^\varepsilon(x, t)$ the signed distance function to $\gamma_{\varepsilon, t}^+$.

From (5.5) and (5.16) we have $e^{\varepsilon|\ln \varepsilon|Z} e^{-a_1 \varepsilon|\ln \varepsilon|} \Gamma_{\tau^\varepsilon} = \gamma_{\varepsilon, 0}^-$ and $\gamma_{\varepsilon, \tau^\varepsilon}^+ = [e^{\varepsilon|\ln \varepsilon|Z} e^{-a_1 \varepsilon|\ln \varepsilon|}]^{-1} \Gamma_0$, and thus

$$\text{dist}(\Gamma_{\tau^\varepsilon}, \gamma_{\varepsilon, \tau^\varepsilon}^+) \geq c_\varepsilon \text{dist}(\gamma_{\varepsilon, 0}^-, \Gamma_0),$$

where $c_\varepsilon \rightarrow 1$, as $\varepsilon \rightarrow 0$. In view of (5.8) it follows that, for $\varepsilon > 0$ small enough, $\text{dist}(\Gamma_{\tau^\varepsilon}, \gamma_{\varepsilon, \tau^\varepsilon}^+) \geq \tilde{C}_0 \varepsilon |\ln \varepsilon|$, with $\tilde{C}_0 > \frac{8}{\lambda}$. It follows that, for $\varepsilon > 0$ small enough,

$$d(x, \tau^\varepsilon) \geq -M\varepsilon \implies d^\varepsilon(x, \tau^\varepsilon) \geq \frac{8}{\lambda} \varepsilon |\ln \varepsilon|.$$

Combining this with (5.13) and (5.18), we infer that

$$u^\varepsilon(x, \tau^\varepsilon) \leq w_\varepsilon^+(x, \tau^\varepsilon),$$

for all $x \in \mathbb{R}^N$. The comparison principle now implies

$$u^\varepsilon(x, \tau^\varepsilon + t) \leq w_\varepsilon^+(x, \tau^\varepsilon + t), \quad (5.19)$$

for all $x \in \mathbb{R}^N$, all $t \in [0, T - \tau^\varepsilon]$.

Finally, for $C > 0$ as in (5.14), we take

$$(x, t) \in (\Omega \cup \Gamma)^c \setminus \mathcal{N}_{C\varepsilon|\ln \varepsilon|}(\Gamma), \quad (5.20)$$

with $t \geq t^\varepsilon$, and prove that $u^\varepsilon(x, t) \leq \alpha^* + \zeta$. If $t \geq \tau^\varepsilon$, we deduce from (5.20) and (5.14) that $x \notin \omega_{\varepsilon, t}^+$ so that conclusion follows from (5.17). If $t^\varepsilon \leq t \leq \tau^\varepsilon$, (5.20) shows that $d(x, t) \leq -C\varepsilon|\ln \varepsilon| \leq -M\varepsilon$ and the conclusion follows from (5.13).

A On the generic regularity of generalized MMC

Let Ω_0 be a bounded domain in \mathbb{R}^N , whose boundary $\Gamma_0 := \partial\Omega_0$ is a smooth hypersurface without boundary. Let $\{\Gamma_t\}_{t \geq 0}$ be the generalized MMC (motion by mean curvature) starting from Γ_0 , and let Ω_t denote the ‘‘inside at time t ’’ as defined in (3.5). We denote by Γ the ‘‘space-time interface’’ and by Ω the ‘‘space-time inside’’. Let us recall some classical definitions in the ‘‘generalized MMC literature’’ (see [7] for instance). We say that the motion is *regular from inside* if $\text{Cl}_{\mathbb{R}^{N+1}}[\Omega] = \Omega \cup \Gamma$, and *regular* if

$$\text{Cl}_{\mathbb{R}^{N+1}}[\Omega] = \Omega \cup \Gamma \quad \text{and} \quad \text{Cl}_{\mathbb{R}^{N+1}}[(\Omega \cup \Gamma)^c] = (\Omega \cup \Gamma)^c \cup \Gamma. \quad (\text{A.1})$$

It is clear that regularity (or regularity from inside) implies non fattening.

We take the opportunity of this Appendix to state a result on the regularity, Proposition A.2, which does not seem to exist in the literature. We start with the following simple lemma. It is well-known in general topology, but for the convenience of the reader we give a proof.

Lemma A.1. *Assume $S \subset (0, \infty)$ is an uncountable set. Then*

$$\exists a > 0, \forall b > a, S \cap [a, b] \text{ is uncountable.} \quad (\text{A.2})$$

Proof. Since $(0, \infty) = \cup_{p \geq 1} [p^{-1}, p)$, there exists an integer $p \geq 2$ such that

$$S \cap [p^{-1}, p) \text{ is uncountable.} \quad (\text{A.3})$$

Assume that (A.2) does not hold. Then

$$\forall a \in [p^{-1}, p), \exists b_a \in (a, p), S \cap [a, b_a) \text{ is countable.} \quad (\text{A.4})$$

Define

$$P := \{c \in (p^{-1}, p) : S \cap [p^{-1}, c) \text{ is at most countable}\}.$$

Since $b_{p^{-1}} \in P$, this set is non empty. We can therefore define $c^* := \sup P \leq p$. Take a sequence (c_n) of elements of P such that $c_n \nearrow c^*$. Since $[p^{-1}, c^*) = \cup_{n \geq 1} [p^{-1}, c_n)$, and since $S \cap [p^{-1}, c_n)$ is at most countable for $n \geq 1$, we have $c^* \in P$. Now if $c^* < p$, then

$$[p^{-1}, b_{c^*}) = [p^{-1}, c^*) \cup [c^*, b_{c^*})$$

contains at most countably many elements of S , which implies $b_{c^*} \in P$. Since $b_{c^*} > c^* = \sup P$, this is impossible. Therefore $c^* = p$; hence $[p^{-1}, p)$ contains at most countably many elements of S . This contradicts (A.3), therefore (A.4) is false. The lemma is proved. \square

We now state our result on the regularity.

Proposition A.2 (Generic regularity). *Let $(\omega_{\nu,0})_{\nu>0}$ be a family of bounded open sets such that $\gamma_{\nu,0} := \partial\omega_{\nu,0}$ is a hypersurface without boundary. Assume that*

$$0 < \nu < \nu' \implies \text{Cl}_{\mathbb{R}^N} [\omega_{\nu',0}] \subset \omega_{\nu,0}. \quad (\text{A.5})$$

For $\nu > 0$, denote by $\{\gamma_{\nu,t}\}_{t \geq 0}$ the generalized MMC starting from $\gamma_{\nu,0}$. We denote by γ_ν and ω_ν the associated “space-time interface” and “space-time inside”. Then the sets

$$\mathcal{J}^- := \{\nu > 0 : \text{Cl}_{\mathbb{R}^{N+1}} [\omega_\nu] \neq \omega_\nu \cup \gamma_\nu\}$$

$$\mathcal{J}^+ := \{\nu > 0 : \text{Cl}_{\mathbb{R}^{N+1}} [(\omega_\nu \cup \gamma_\nu)^c] \neq (\omega_\nu \cup \gamma_\nu)^c \cup \gamma_\nu\}$$

are at most countable.

Proof. We only prove the assertion for \mathcal{J}^- since that for \mathcal{J}^+ can be shown in the same manner. First note that it follows from assumption (A.5) and the comparison principle (3.8) that

$$0 < \nu < \nu' \implies (\omega_{\nu'} \cup \gamma_{\nu'}) \subset \omega_\nu. \quad (\text{A.6})$$

For $\nu > 0$, define

$$\delta_\nu := \sup_{(y, \tau) \in \gamma_\nu} \text{dist}((y, \tau), \text{Cl}_{\mathbb{R}^{N+1}}[\omega_\nu]) .$$

Since $\text{Cl}_{\mathbb{R}^{N+1}}[\omega_\nu] \subset (\omega_\nu \cup \gamma_\nu)$ always holds, we have $\mathcal{J}^- = \{\nu > 0 : \delta_\nu > 0\}$. Assume by contradiction that \mathcal{J}^- is uncountable. Then there exists an integer n_0 such that the set

$$\mathcal{J}_0^- := \{\nu > 0 : \delta_\nu \geq \frac{1}{n_0}\} \text{ is uncountable.}$$

From Lemma A.1, there exists $\nu^* > 0$ such that, for all $\nu > \nu^*$, the set $\mathcal{J}_0^- \cap [\nu^*, \nu)$ is uncountable. Therefore we can construct a decreasing sequence (ν_n) of elements of \mathcal{J}_0^- which tends to ν^* . From the definition of \mathcal{J}_0^- , we deduce the existence of $(x_n, t_n) \in \gamma_{\nu_n}$ such that

$$\text{dist}((x_n, t_n), \text{Cl}_{\mathbb{R}^{N+1}}[\omega_{\nu_n}]) \geq \frac{1}{n_0} . \quad (\text{A.7})$$

On the one hand, if $j < k$, the decreasing of the sequence (ν_n) and (A.6) imply that $(x_j, t_j) \in \omega_{\nu_k}$. Therefore (A.7) yields

$$d((x_k, t_k), (x_j, t_j)) \geq \frac{1}{n_0} .$$

But, on the other hand, (A.6) implies that $(x_n, t_n) \in \omega_{\nu^*}$ so that we can extract a convergent subsequence of (x_n, t_n) . This is a contradiction and the proposition is proved. \square

As stated in Section 2, it is known that, for an admissible initial domain Ω_0 , the evolution $\{\Gamma_t\}_{t \geq 0}$ starting from $\Gamma_0 = \partial\Omega_0$ is regular from inside [17, Corollary 4.5.11]. Proposition A.2 provides a simple proof of the regularity both from inside and from outside. More precisely, the following holds.

Corollary A.3. *Let Ω_0 be an admissible domain in the sense of Definition 2.3. Then the generalized MMC starting from $\Gamma_0 := \partial\Omega_0$ is regular.*

Proof. For $0 \leq \nu \leq \nu_0$, define, as in Remark 2.4,

$$\gamma_{\nu,0} := \psi_\nu(\Gamma_0) = e^{\nu Z} [e^{-a_1 \nu} \Phi_{a_2 \nu}(\Gamma_0)] ,$$

and denote by $\omega_{\nu,0}$ the domain enclosed by $\gamma_{\nu,0}$. It is clear from (2.9) that the family of bounded open sets $(\omega_{\nu,0})_{0 < \nu < \nu_0}$ satisfies the assumption (A.5) of Proposition A.3. Therefore, for almost all $\nu \in (0, \nu_0)$, the generalized MMC $\{\gamma_{\nu,t}\}_{t \geq 0}$ starting from $\gamma_{\nu,0}$ is regular, that is,

$$\text{Cl}_{\mathbb{R}^{N+1}}[\omega_\nu] = \omega_\nu \cup \gamma_\nu , \quad (\text{A.8})$$

$$\text{Cl}_{\mathbb{R}^{N+1}} [(\omega_\nu \cup \gamma_\nu)^c] = (\omega_\nu \cup \gamma_\nu)^c \cup \gamma_\nu. \quad (\text{A.9})$$

Since

$$\omega_\nu = \Pi_\nu(\Omega) \cap (\mathbb{R}^N \times [0, \infty)) = \Pi_\nu(\Omega \cap (\mathbb{R}^N \times [a_2\nu, \infty))),$$

with Π_ν as in (5.2), we have

$$\Omega \cap (\mathbb{R}^N \times [a_2\nu, \infty)) = \Pi_\nu^{-1}(\omega_\nu).$$

Since Π_ν is a homeomorphism on \mathbb{R}^{N+1} , we see from (A.8) that, for almost all $\nu \in (0, \nu_0)$,

$$\text{Cl}_{\mathbb{R}^{N+1}} [\Omega \cap (\mathbb{R}^N \times [a_2\nu, \infty))] = (\Omega \cup \Gamma) \cap (\mathbb{R}^N \times [a_2\nu, \infty)).$$

Letting $\nu \rightarrow 0$, we obtain

$$(\Omega \cup \Gamma) \cap (\mathbb{R}^N \times (0, \infty)) \subset \text{Cl}_{\mathbb{R}^{N+1}} [\Omega \cap (\mathbb{R}^N \times (0, \infty))].$$

Combining this and $\text{Cl}_{\mathbb{R}^N} [\Omega_0] = \Omega_0 \cup \Gamma_0$, we obtain $\text{Cl}_{\mathbb{R}^{N+1}} [\Omega] = \Omega \cup \Gamma$. Similarly, from (A.9), we obtain $\text{Cl}_{\mathbb{R}^{N+1}} [(\Omega \cup \Gamma)^c] = (\Omega \cup \Gamma)^c \cup \Gamma$. The proposition is proved. \square

References

- [1] M. Alfaro, D. Hilhorst and H. Matano, *The singular limit of the Allen-Cahn equation and the FitzHugh-Nagumo system*, J. Differential Equations **245** (2008), no. 2, 505–565.
- [2] S. Allen and J. Cahn, *A microscopic theory for antiphase boundary motion and its application to antiphase domain coarsening*, Acta Metallica **27** (1979), 1084–1095.
- [3] S. Altschuler, S. Angenent and Y. Giga, *Mean curvature flow through singularities for surfaces of rotation*, J. Geom. Anal. **5** (1993), no. 3, 293–358.
- [4] S. Angenent, *Shrinking doughnuts*, Nonlinear diffusion equations and their equilibrium states, 3 (Gregynog, 1989), 21–38, Progr. Nonlinear Differential Equations Appl., 7, Birkhuser Boston, Boston, MA, 1992.
- [5] G. Barles, L. Bronsard and P. E. Souganidis, *Front propagation for reaction-diffusion equations of bistable type*, Ann. Inst. Henri Poincaré, **9** (1992), 479–496.
- [6] G. Barles and F. Da Lio, *A geometrical approach to front propagation problems in bounded domains with Neumann-type boundary conditions*, Interfaces Free Bound. **5** (2003), 239–274.

- [7] G. Barles, H. M. Soner and P. E. Souganidis, *Front propagation and phase field theory*, SIAM J. Control Optim. **31** (1993), 439–469.
- [8] G. Barles and P. E. Souganidis, *A new approach to front propagation problems : theory and applications*, Arch. Rational Mech. Anal. **141** (1998), 237–296.
- [9] L. Bronsard and R. V. Kohn, *Motion by mean curvature as the singular limit of Ginzburg–Landau dynamics*, J. Differential Equations **90** (1991), 211–237.
- [10] X. Chen, *Generation and propagation of interfaces for reaction-diffusion equations*, J. Differential Equations **96** (1992), 116–141.
- [11] X. Chen, *Generation and propagation of interfaces for reaction-diffusion systems*, Trans. Amer. Math. Soc. **334** (1992), 877–913.
- [12] X. Chen and F. Reitich, *Local existence and uniqueness of solutions of the Stefan problem with surface tension and kinetic undercooling*, J. Math. Anal. Appl. **164** (1992), 350–362.
- [13] Y. G. Chen, Y. Giga and S. Goto, *Uniqueness and existence of viscosity solutions of generalized mean curvature flow equations*, J. Diff. Geometry **33** (1991), 749–786.
- [14] M. G. Crandall, H. Ishii and P. -L. Lions, *User’s guide to viscosity solutions of second order partial differential equations*, Bull. AMS **27** (1992), 1–67.
- [15] L. C. Evans, H. M. Soner and P. E. Souganidis, *Phase transitions and generalized motion by mean curvature*, Comm. Pure Appl. Math. **45** (1992), 1097–1123.
- [16] L. C. Evans and J. Spruck, *Motion of level sets by mean curvature I*, J. Differential Geometry **33** (1991), 635–681.
- [17] Y. Giga, *Surface evolution equations*, Monographs in Mathematics 99, Birkhäuser Verlag, Basel, Boston, Berlin, 2006.
- [18] M. A. Grayson, *The heat equation shrinks embedded plane curves to round points*, J. Differential Geometry **26** (1987), 285–314.
- [19] K. Kawasaki and T. Ohta, *Kinetic drumhead model of interface I*, Progress of Theoretical Physics **67** (1982), 147–163.
- [20] P. de Mottoni and M. Schatzman, *Development of interfaces in \mathbb{R}^n* , Proc. Roy. Soc. Edinburgh **116A** (1990), 207–220.

- [21] P. de Mottoni and M. Schatzman, *Geometrical evolution of developed interfaces*, Trans. Amer. Math. Soc. **347** (1995), 1533–1589.
- [22] H. M. Soner, *Ginzburg-Landau equation and motion by mean curvature, I: convergence*, J. Geom. Anal. **7** (1997), 437–475.
- [23] H. M. Soner, *Ginzburg-Landau equation and motion by mean curvature, II: development of the initial interface*, J. Geom. Anal. **7** (1997), 477–491.