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A PDE APPROACH TO LARGE-TIME ASYMPTOTICS FOR BOUNDARY-VALUE PROBLEMS FOR NONCONVEX HAMILTON-JACOBI EQUATIONS

GUY BARLES AND HIROYOSHI MITAKE

ABSTRACT. We investigate the large-time behavior of three types of initial-boundary value problems for Hamilton-Jacobi Equations with nonconvex Hamiltonians. We consider the Neumann or oblique boundary condition, the state constraint boundary condition and Dirichlet boundary condition. We establish general convergence results for viscosity solutions to asymptotic solutions as time goes to infinity via an approach based on PDE techniques. These results are obtained not only under general conditions on the Hamiltonians but also under weak conditions on the domain and the oblique direction of reflection in the Neumann case.

1. INTRODUCTION AND MAIN RESULTS

In this paper we investigate the large time behavior of viscosity solutions of the initial-boundary value problems for Hamilton-Jacobi Equations which we write under the form

$$(IB) \quad \begin{cases} u_t + H(x, Du) = 0 & \text{in } \Omega \times (0, \infty), & (1) \\ B(x, u, Du) = 0 & \text{on } \partial\Omega \times (0, \infty), & (2) \\ u(x, 0) = u_0(x) & \text{on } \overline{\Omega}, & (3) \end{cases}$$

where Ω is a bounded domain of \mathbb{R}^N , $H = H(x, p)$ is a given real-valued continuous function on $\overline{\Omega} \times [0, \infty)$, which is *coercive*, i.e.,

$$(A1) \quad \lim_{r \rightarrow +\infty} \inf \{ H(x, p) \mid x \in \overline{\Omega}, |p| \geq r \} = +\infty.$$

Here B and u_0 are given real-valued continuous functions on $\partial\Omega \times \mathbb{R} \times \mathbb{R}^N$ and $\overline{\Omega}$, respectively. The solution u is a real-valued function on $\overline{\Omega} \times [0, \infty)$ and we respectively denote by $u_t := \partial u / \partial t$ and $Du := (\partial u / \partial x_1, \dots, \partial u / \partial x_N)$ its time derivative and gradient with respect to

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the space variable. We are dealing only with viscosity solutions of Hamilton-Jacobi equations in this paper and thus the term “viscosity” may be omitted henceforth. We also point out that the boundary conditions have to be understood in the viscosity sense : we refer the reader to the beginning of Appendix where this definition is recalled.

We consider three types of boundary conditions (BC in short)

$$\text{(Neumann or oblique BC)} \quad B(x, r, p) = \gamma(x) \cdot p - g(x), \quad (4)$$

where g is a given real-valued continuous function on $\partial\Omega$ and $\gamma : \overline{\Omega} \rightarrow \mathbb{R}^N$ is a continuous vector field which is oblique to $\partial\Omega$, i.e.,

$$n(x) \cdot \gamma(x) > 0 \quad \text{for any } x \in \partial\Omega, \quad (5)$$

where $n(x)$ is the outer unit normal vector at x to $\partial\Omega$,

$$\text{(state constraint BC)} \quad B(x, r, p) = -1 \text{ for all } (x, r, p) \in \overline{\Omega} \times \mathbb{R} \times \mathbb{R}^N, \quad (6)$$

and finally

$$\text{(Dirichlet BC)} \quad B(x, r, p) = r - g(x). \quad (7)$$

We respectively denote Problem (IB) with boundary conditions (4), (6) and (7) by (CN), (SC) and (CD).

We impose different regularity assumptions on $\partial\Omega$ depending on the boundary condition we are treating. When we consider Neumann/oblique derivative problems, we *always* assume that Ω is a domain with a C^1 -boundary. When we consider state constraint or Dirichlet problems, we *always* assume that Ω is a domain with a C^0 -boundary. Moreover when we consider Dirichlet problems, we assume a compatibility condition on the initial value u_0 and the boundary value g , i.e.,

$$u_0(x) \leq g(x) \text{ for all } x \in \partial\Omega. \quad (8)$$

Since we are going to assume γ in (4) to be only continuous (and not Lipschitz continuous as it is classically the case), we have to solve Problem (IB) with very weak conditions. This is possible as a consequence of the coercivity assumption on H since the solutions are expected to be in $W^{1,\infty}$, which denotes the set of bounded functions whose the first distributional derivatives are essentially bounded. We postpone the proof of the existence, uniqueness and regularity of the solution u of (IB) to Appendix since the proofs are a little bit technical and the main topic of this paper concerns the asymptotic behavior as $t \rightarrow +\infty$ of the solutions. The reader may first assume that $\partial\Omega$ and γ are regular enough to concentrate on this asymptotic behavior and then consider Appendix for the generalization to less regular situations.

More precisely we have the following existence result for (IB).

Theorem 1.1 (Existence of Solutions of (IB)). *Assume that (A1) holds and that $u_0 \in W^{1,\infty}(\Omega)$.*

(i) *Assume that Ω is a bounded domain with a C^1 -boundary and that γ, g are continuous functions which satisfy (5). There exists a unique solution u of (CN) which is continuous on $\bar{\Omega} \times [0, \infty)$ and in $W^{1,\infty}(\Omega \times [0, \infty))$.*

(ii) *Assume that Ω is a bounded domain with a C^0 -boundary. There exists a unique solution u of (SC) which is continuous on $\bar{\Omega} \times [0, \infty)$ and in $W^{1,\infty}(\Omega \times [0, \infty))$.*

(iii) *Assume that Ω is a bounded domain with a C^0 -boundary and that $g \in C(\partial\Omega)$ and u_0 which satisfy (8). There exists a unique solution u of (CD) which is continuous on $\bar{\Omega} \times [0, \infty)$ and in $W^{1,\infty}(\Omega \times [0, \infty))$.*

Throughout this work, we are going to use *anytime* the same assumptions as in Theorem 1.1 which clearly depend on the kind of problem we are considering. In order to simplify the statements of our results, we will say from now on the “basic assumptions” ((BA) in short) are satisfied if (A1) holds and if Ω is a bounded domain with a C^1 -boundary in the case of Neumann/oblique derivatives boundary conditions or Ω is a bounded domain with a C^0 -boundary in the case of state constraint or Dirichlet boundary conditions. When it is relevant, $u_0 \in W^{1,\infty}(\Omega)$ is also assumed to be part of (BA) (results for time-dependent problems) and so are the continuity assumptions for γ, g on $\partial\Omega$ as well as (5) and (8).

Next we recall that the standard asymptotic behavior for solutions of Hamilton-Jacobi Equations is given by an *additive eigenvalue* or *ergodic problem* : for t large enough, the solution $u(x, t)$ is expected to look like $-ct + v(x)$ where the constant c and the function v are solutions of

$$H(x, Dv(x)) = c \quad \text{in } \Omega,$$

with suitable boundary conditions. A typical result, which was first proved in \mathbb{R}^N for the periodic case by P.-L. Lions, G. Papanicolaou and S. R. S. Varadhan [27], is that there exists a unique constant c for which this problem has a solution, while the solution v may be non unique, even up to an additive constant. This non-uniqueness feature is a key difficulty in the study of the asymptotic behavior.

In the Neumann/oblique derivative or state constraint case, the additive eigenvalue or ergodic problem reads

$$(E) \quad \begin{cases} H(x, Dv(x)) = a & \text{in } \Omega, \\ B(x, v(x), Dv(x)) = 0 & \text{on } \partial\Omega. \end{cases} \quad (9) \quad (10)$$

Here one seeks for a pair (v, a) of $v \in C(\bar{\Omega})$ and $a \in \mathbb{R}$ such that v is a solution of (E). If (v, a) is such a pair, we call v an *additive eigenfunction*

or *ergodic function* and a an *additive eigenvalue* or *ergodic constant*. We denote Problem (E) with (4), (6) by (E-N) and (E-SC).

For the large-time asymptotics of solutions of (CD), as we see it later, there are different cases which are treated either by using (E-SC) or the stationary Dirichlet problems

$$(D)_a \quad \begin{cases} H(x, Dv(x)) = a & \text{in } \Omega, \\ v(x) = g(x) & \text{on } \partial\Omega \end{cases} \quad (11)$$

$$(12)$$

for $a \in \mathbb{R}$. Since the existence of solutions of (E-N), (E-SC) and $(D)_a$ are playing a central role, we first state the following result.

Theorem 1.2 (Existence of Solutions of Additive Eigenvalue Problems). *Assume that (BA) holds.*

(i) *There exists a solution $(v, c_n) \in W^{1,\infty}(\Omega) \times \mathbb{R}$ of (E-N). Moreover, the additive eigenvalue is unique and is represented by*

$$c_n = \inf\{a \in \mathbb{R} \mid \text{(E-N) has a subsolution}\}. \quad (13)$$

(ii) *There exists a solution $(v, c_{sc}) \in W^{1,\infty}(\Omega) \times \mathbb{R}$ of (E-SC). An additive eigenvalue is unique and is represented by*

$$c_{sc} = \inf\{a \in \mathbb{R} \mid \text{(9) has a subsolution}\}. \quad (14)$$

(iii) *There exists a solution $v \in W^{1,\infty}(\Omega)$ of $(D)_a$ if and only if $a \geq c_{sc}$.*

Existence results of solutions of (E) have been established by P.-L. Lions in [25] (see also [21]) for the Neumann/oblique derivative boundary condition and by I. Capuzzo-Dolcetta and P.-L. Lions in [7] (see also [28]) for the state constraint boundary condition. Additive eigenvalue problems (E) and $(D)_a$ give the “stationary states” for solutions of (IB) as our main result shows. A simple observation related to this is that, for any $(v, c) \in C(\overline{\Omega})$, the function $v(x) - ct$ is a solution of (IB) if and only if (v, c) is a solution of (E) or $(D)_c$.

As we suggested above, our main purpose is to prove that the solution u of (IB) has, for the Neumann/oblique derivative or the state constraint boundary condition, the following behavior

$$u(x, t) - (v(x) - ct) \rightarrow 0 \quad \text{uniformly on } \overline{\Omega} \text{ as } t \rightarrow \infty, \quad (15)$$

where $(v, c) \in W^{1,\infty}(\Omega) \times \mathbb{R}$ is a solution of (E-N) or (E-SC) while, for the Dirichlet boundary condition case, this behavior depends on the sign of c_{sc} . We call such a function $v(x) - ct$ an *asymptotic solution* of (IB). It is worth mentioning that though we can relatively easily get the convergence

$$\frac{u(x, t)}{t} \rightarrow -c \quad \text{uniformly on } \overline{\Omega} \text{ as } t \rightarrow \infty$$

as a result of Theorem 1.2 and the comparison principle for (IB) (see Theorem 7.1), it is generally rather difficult to show (15), since (as we already mention it above) additive eigenfunctions of (E) or solutions of $(D)_c$ may not be unique.

Before stating our main result on the asymptotic behavior of solution of (E), we recall that, in the last decade, the large time behavior of solutions of Hamilton-Jacobi equation in compact manifold (or in \mathbb{R}^N , mainly in the periodic case) has received much attention and general convergence results for solutions have been established. G. Namah and J.-M. Roquejoffre in [31] are the first to prove (15) under the following additional assumption

$$H(x, p) \geq H(x, 0) \text{ for all } (x, p) \in \mathcal{M} \times \mathbb{R}^N \text{ and } \max_{\mathcal{M}} H(x, 0) = 0, \quad (16)$$

where \mathcal{M} is a smooth compact N -dimensional manifold without boundary. Then A. Fathi in [10] proved the same type of convergence result by dynamical systems type arguments introducing the “weak KAM theory”. Contrarily to [31], the results of [10] use strict convexity (and smoothness) assumptions on $H(x, \cdot)$, i.e., $D_{pp}H(x, p) \geq \alpha I$ for all $(x, p) \in \mathcal{M} \times \mathbb{R}^N$ and $\alpha > 0$ (and also far more regularity) but do not require (16). Afterwards J.-M. Roquejoffre [32] and A. Davini and A. Siconolfi in [9] refined the approach of A. Fathi and they studied the asymptotic problem for (1) on \mathcal{M} or N -dimensional torus. We also refer to the articles [3, 20, 16, 17, 18] for the asymptotic problems in the whole domain \mathbb{R}^N without the periodic assumptions in various situations. The first author and P. E. Souganidis obtained in [4] more general results, for possibly non-convex Hamiltonians, by using an approach based on partial differential equations methods and viscosity solutions, which was not using in a crucial way the explicit formulas of representation of the solutions. Later, the first author and J.-M. Roquejoffre provided results for unbounded solutions, for convex and non-convex equations, using partially the ideas of [4] but also of [31]. In this paper, we follow the approach of [4].

There also exists results on the asymptotic behavior of solutions of convex Hamilton-Jacobi Equation with boundary conditions. The second author [28] studied the case of the state constraint boundary condition and then the Dirichlet boundary conditions [29, 30]. Roquejoffre in [32] was also dealing with solutions of the Cauchy-Dirichlet problem which satisfy the Dirichlet boundary condition pointwise (in the classical sense) : this is a key difference with the results of [29, 30] where the solutions were satisfying the Dirichlet boundary condition in a generalized (viscosity solutions) sense. For Cauchy-Neumann problems, H. Ishii in

[22] establishes very recently the asymptotic behavior. We point out that all these works use a generalized dynamical approach similar to [10, 9], which is very different from ours. Our results are more general since they also hold also for non-convex Hamiltonians. It is worthwhile to mention that our results on (15) include those in [28, 29, 22].

We also refer to the articles [32, 6] for the large time behavior of solutions to time-dependent Hamilton-Jacobi equations and [13, 12] for the convergence rate in (15). Recently the second author with Y. Giga and Q. Liu in [14, 15] has gotten the large time behavior of solutions of Hamilton-Jacobi equations with noncoercive Hamiltonian which is motivated by a model describing growing faceted crystals.

To state our main result, we need the following assumptions on the Hamiltonian. We denote by $H_a(x, p)$ the Hamiltonian $H(x, p) - a$ for $a \in \mathbb{R}$.

(A2) $_a^+$ There exists $\eta_0 > 0$ such that, for any $\eta \in (0, \eta_0]$, there exists $\psi_\eta > 0$ such that if $H_a(x, p + q) \geq \eta$ and $H_a(x, q) \leq 0$ for some $x \in \bar{\Omega}$ and $p, q \in \mathbb{R}^N$, then for any $\mu \in (0, 1]$,

$$\mu H_a(x, \frac{p}{\mu} + q) \geq H_a(x, p + q) + \psi_\eta(1 - \mu).$$

(A2) $_a^-$ There exists $\eta_0 > 0$ such that, for any $\eta \in (0, \eta_0]$, there exists $\psi_\eta > 0$ such that if $H_a(x, p + q) \leq -\eta$ and $H_a(x, q) \geq 0$ for some $x \in \bar{\Omega}$ and $p, q \in \mathbb{R}^N$, then for any $\mu \geq 1$,

$$\mu H_a(x, \frac{p}{\mu} + q) \leq H_a(x, p + q) - \frac{\psi_\eta(\mu - 1)}{\mu}.$$

Our main result is the following theorem.

Theorem 1.3 (Large-Time Asymptotics). *Assume that (BA) holds.*

(i) **(Neumann/oblique derivative problem)** *If (A2) $_{c_n}^+$ or (A2) $_{c_n}^-$ holds, then there exists a solution $v \in W^{1,\infty}(\Omega)$ of (E-N) with $c = c_n$ such that (15) holds with $c = c_n$.*

(ii) **(State constraint problem)** *If (A2) $_{c_{sc}}^+$ or (A2) $_{c_{sc}}^-$ holds, then there exists a solution $v \in W^{1,\infty}(\Omega)$ of (E-SC) with $c = c_{sc}$ such that (15) holds with $c = c_{sc}$.*

(iii) **(Dirichlet problem)** *Assume that (A2) $_{c_{sc}}^+$ or (A2) $_{c_{sc}}^-$ holds.*

(iii-a) *If $c_{sc} > 0$, then there exists a solution $v \in W^{1,\infty}(\Omega)$ of (E-SC) with $c = c_{sc}$ such that (15) holds with $c = c_{sc}$.*

(iii-b) *If $c_{sc} \leq 0$, then there exists a solution $v \in W^{1,\infty}(\Omega)$ of (D) $_0$ such that (15) holds with $c = 0$.*

Assumption (A2) $_a^+$ is introduced in [4] to replace the convexity assumption : it mainly concerns the set $\{H_a \geq 0\}$ and the behavior of

H_a in this set. Assumption $(A2)_a^-$ is a modification of $(A2)_a^+$ which, on the contrary, concerns the set $\{H_a \leq 0\}$. We can generalize them as in [4] (see Remark 1 (ii)) but to simplify our arguments we only use the simplified version in Theorem 1.3. It is worthwhile to mention that we can deal with $u_0 \in C(\overline{\Omega})$ and prove that uniform continuous solutions of (IB) (instead of Lipschitz continuous solutions) have the same large-time asymptotics given by Theorem 1.3. See Remark 1 (i). We also mention that assumptions $(A2)_c^+$ and $(A2)_c^-$ are equivalent as $(A7)_+$ and $(A7)_-$ in [18] under the assumption that H is coercive and convex with respect to p -variable, where c is the associated constant as in Theorem 1.3 (see [16, Appencix C]).

This paper is organized as follows: in Section 2 we recall the proof of Theorem 1.2. In Section 3 we present *asymptotically monotone property*, which is a key tool to prove Theorem 1.3. Section 4 is devoted to the proof of a key lemma to prove asymptotically monotone properties. In Section 5, we give the proof of Theorem 1.3 and we give some remarks on the large-time behavior of solutions of convex Hamilton-Jacobi equations in Section 6. In Appendix we present the definition, existence, uniqueness and regularity results for (IB).

Notations. Let \mathbb{R}^N denote the N -dimensional Euclidean space for some $k \in \mathbb{N}$. We denote by $|\cdot|$ the usual Euclidean norm. We write $B(x, r) = \{y \in \mathbb{R}^N \mid |x - y| < r\}$ for $x \in \mathbb{R}^N$, $r > 0$. For $A \subset \mathbb{R}^N$, we denote by $C(A)$, $BUC(A)$, $LSC(A)$, $USC(A)$ and $C^k(A)$ the space of real-valued continuous, bounded uniformly continuous, lower semicontinuous, upper semicontinuous and k -th continuous differentiable functions on A for $k \in \mathbb{N}$, respectively. We write $a \wedge b = \min\{a, b\}$ and $a \vee b = \max\{a, b\}$ for $a, b \in \mathbb{R}$. We call a function $m : [0, \infty) \rightarrow [0, \infty)$ a modulus if it is continuous and nondecreasing on $[0, \infty)$ and vanishes at the origin. For any $k \in \mathbb{N} \cup \{0\}$ we call Ω a domain with a C^k boundary if Ω has a boundary $\partial\Omega$ locally represented as the graph of a C^k function, i.e., for each $z \in \partial\Omega$ there exist $r > 0$ and a function $b \in C^k(\mathbb{R}^{N-1})$ such that—upon relabelling and re-orienting the coordinates axes if necessary—we have

$$\Omega \cap B(z, r) = \{(x', x_N) \in \mathbb{R}^{N-1} \times \mathbb{R} \mid x \in B(z, r), x_N > b(x')\}.$$

2. ADDITIVE EIGENVALUE PROBLEM

Proof of Theorem 1.2. We just sketch the proof since it is an easy adaptation of classical arguments.

We start with (E-N) and for any $\varepsilon \in (0, 1)$, we consider

$$\begin{cases} \varepsilon u^\varepsilon + H(x, Du^\varepsilon) = 0 & \text{in } \Omega, \\ B(x, u^\varepsilon, Du^\varepsilon) = 0 & \text{on } \partial\Omega. \end{cases} \quad (17)$$

In order to apply Perron's method, we first build a subsolution and a supersolution.

We claim that $u_{1,\varepsilon}^- = -M_1/\varepsilon$ and $u_{1,\varepsilon}^+ = M_1/\varepsilon$ are respectively sub and supersolution of (17). This is obvious in Ω for a suitable large $M_1 > 0$ since $H(x, 0)$ is continuous, hence bounded, on $\bar{\Omega}$. If $x \in \partial\Omega$, we recall (see [25, 8] for instance) that the super-differential of $u_{1,\varepsilon}^-$ at x consists in elements of the form $\lambda n(x)$ with $\lambda \leq 0$ and we have to show that

$$\min\{-M_1 + H(x, \lambda n(x)), \lambda n(x) \cdot \gamma(x) - g(x)\} \leq 0.$$

By (5) it is clear enough that there exists $\bar{\lambda} < 0$ such that, if $\lambda \leq \bar{\lambda}$, then $\lambda n(x) \cdot \gamma(x) - g(x) \leq 0$. Then choosing $M_1 \geq \max\{H(x, \lambda n(x)) \mid x \in \partial\Omega, \bar{\lambda} \leq \lambda \leq 0\}$, the above inequality holds. A similar argument shows that $u_{1,\varepsilon}^+$ is a supersolution of (17) for M_1 large enough.

We remark that, because of (A1) and the regularity of the boundary of Ω , the subsolutions w of (17) such that $u_{1,\varepsilon}^- \leq w \leq u_{1,\varepsilon}^+$ on $\bar{\Omega}$ satisfy $|Dw| \leq M_2$ in Ω and therefore they are equi-Lipschitz continuous on $\bar{\Omega}$. With these informations, Perron's method provides us with a solution $u_\varepsilon \in W^{1,\infty}(\Omega)$ of (17). Moreover, by construction, we have

$$|\varepsilon u_\varepsilon| \leq M_1 \text{ on } \bar{\Omega} \quad \text{and} \quad |Du_\varepsilon| \leq M_2 \text{ in } \Omega. \quad (18)$$

Next we set $v_\varepsilon(x) := u_\varepsilon(x) - u_\varepsilon(x_0)$ for a fixed $x_0 \in \bar{\Omega}$. Because of (18) and the regularity of the boundary $\partial\Omega$, $\{v_\varepsilon\}_{\varepsilon \in (0,1)}$ is a sequence of equi-Lipschitz continuous and uniformly bounded functions on $\bar{\Omega}$. By Ascoli-Arzelà's Theorem, there exist subsequences $\{v_{\varepsilon_j}\}_j$ and $\{u_{\varepsilon_j}\}_j$ such that

$$v_{\varepsilon_j} \rightarrow v, \quad \varepsilon_j u_{\varepsilon_j} \rightarrow -c_n \text{ uniformly on } \bar{\Omega}$$

as $j \rightarrow \infty$ for some $v \in W^{1,\infty}(\Omega)$ and $c_n \in \mathbb{R}$. By a standard stability result of viscosity solutions we see that (v, c_n) is a solution of (E-N).

We prove the uniqueness of additive eigenvalues. Let $(v_a, a) \in W^{1,\infty}(\Omega) \times \mathbb{R}$ be a subsolution of (E-N). Adding a positive constant to v if necessary we may assume that $v_a \leq v$ on $\bar{\Omega}$. Note that $v - c_n t$ and $v_a - at$ are, respectively, a solution and a subsolution of (1) and (2). By the comparison principle for (IB) (see Theorem 7.1), we have

$$v_a - at \leq v - c_n t \text{ on } \bar{\Omega} \times [0, \infty).$$

Therefore we have

$$v_a - v \leq (a - c_n)t \text{ on } \bar{\Omega} \times [0, \infty). \quad (19)$$

If $a < c_n$, then (19) implies a contradiction for a large $t > 0$ since $(a - c_n)t \rightarrow -\infty$ as $t \rightarrow +\infty$. Therefore we obtain $c_n \leq a$ and thus we see that c can be represented by

$$c_n = \inf\{a \in \mathbb{R} \mid (\text{E-N}) \text{ has a subsolution}\}.$$

The uniqueness of additive eigenvalues follows immediately.

Next we consider (E-SC). The proof follows along the same lines, except for the following point.

(i) For a subsolution, it is enough to choose the constant $-m/\varepsilon$, where $m = \max_{x \in \overline{\Omega}} H(x, 0)$ since only the inequality in Ω is needed. For a supersolution, we choose M/ε , where $M := -\min_{\overline{\Omega} \times \mathbb{R}^N} H(x, p)$.

(ii) Since the boundary has only C^0 -regularity, the solutions are in $W^{1,\infty}(\Omega)$ but may not be Lipschitz continuous up to the boundary. But they are equi-continuous (see [23, Proposition 8.1]) and this is enough to complete the argument.

(iii) We have

$$c_{sc} = \inf\{a \in \mathbb{R} \mid (9) \text{ has a subsolution}\},$$

by the same argument as above since, in the state-constraint case, there are no boundary conditions for the subsolutions.

Finally we consider $(D)_a$. In the case where $c_{sc} > a$, it is clear that $(D)_a$ has no solutions due to the definition of c_{sc} . Thus we need only to show that $(D)_a$ has a solution in the case where $c_{sc} \leq a$.

Due to the coercivity of H , there exists $p_0 \in \mathbb{R}^N$ such that $H(x, p_0) \geq a$ for all $x \in \overline{\Omega}$. We set $\psi_1(x) = p_0 \cdot x + C$, where $C > 0$ is a constant chosen so that $\psi_1(x) \geq g(x)$ for all $x \in \partial\Omega$. Since $a \geq c_{sc}$, any solution $\psi_2 \in C(\overline{\Omega})$ of (E-SC) is a subsolution of (11) and subtracting a sufficiently large constant from ψ_2 if necessary, we may assume that $\psi_2(x) \leq g(x)$ on $\partial\Omega$ and $\psi_2 \leq \psi_1$ on $\overline{\Omega}$. As a consequence of Perron's method we see that there exists a solution v of $(D)_a$ for any $a \geq c_{sc}$. \square

The following proposition shows that, taking into account the ergodic effect, we obtain bounded solutions of (IB). This result is a straightforward consequence of Theorems 1.2 and 7.1.

Proposition 2.1 (Boundedness of Solutions of (IB)). *Assume that (BA) holds. Let c_n and c_{sc} be the constants given by (13) and (14), respectively.*

- (i) *Let u be the solution of (CN). Then $u + c_n t$ is bounded on $\overline{\Omega} \times [0, \infty)$.*
- (ii) *Let u be the solution of (SC). Then $u + c_{sc} t$ is bounded on $\overline{\Omega} \times [0, \infty)$.*
- (iii) *Let u be the solution of (CD).*
- (iii-a) *If $c_{sc} > 0$, then $u + c_{sc} t$ is bounded on $\overline{\Omega} \times [0, \infty)$.*
- (iii-b) *If $c_{sc} \leq 0$, then u is bounded on $\overline{\Omega} \times [0, \infty)$.*

3. ASYMPTOTICALLY MONOTONE PROPERTY

Following Proposition 2.1, it turns out that, in order to go further in the study of the asymptotic behavior of the solutions u of (IB), we have to consider $u_c := u + ct$, where c is the associated additive eigenvalue: $c := c_n$ for (CN), $c := c_{sc}$ for (SC) and $c := c_{sc}$ if $c_{sc} > 0$ and $c := 0$ if $c_{sc} \leq 0$ for (CD). A first step to prove the convergence of the u_c 's is the following result on their asymptotic monotonicity in time which can be stated exactly in the same way for the three cases, namely (CN), (SC) and (CD).

Theorem 3.1 (Asymptotically Monotone Property). *Assume that (BA) holds.*

(i) **(Asymptotically Increasing Property)**

Assume that $(A2)_c^+$ holds. For any $\eta \in (0, \eta_0]$, there exists $\delta_\eta : [0, \infty) \rightarrow [0, 1]$ such that

$$\begin{aligned} \delta_\eta(s) &\rightarrow 0 \quad \text{as } s \rightarrow \infty \quad \text{and} \\ u_c(x, s) - u_c(x, t) + \eta(s - t) &\leq \delta_\eta(s) \end{aligned}$$

for all $x \in \overline{\Omega}$, $s, t \in [0, \infty)$ with $t \geq s$.

(ii) **(Asymptotically Decreasing Property)**

Assume that $(A2)_c^-$ holds. For any $\eta \in (0, \eta_0]$, there exists $\delta_\eta : [0, \infty) \rightarrow [0, \infty)$ such that

$$\begin{aligned} \delta_\eta(s) &\rightarrow 0 \quad \text{as } s \rightarrow \infty \quad \text{and} \\ u_c(x, t) - u_c(x, s) - \eta(t - s) &\leq \delta_\eta(s) \end{aligned}$$

for all $x \in \overline{\Omega}$, $s, t \in [0, \infty)$ with $t \geq s$.

In the same way as for u_c , we denote by v_c a subsolution of (E) which means, of course, either (E-N) or (E-SC). We keep the subscript “ c ” to remind that it is associated to either c_n or c_{sc} .

By Proposition 2.1 u_c is bounded on $\overline{\Omega} \times [0, \infty)$. Note that $v_c - M$ are still subsolutions of (E) or (D) $_c$ for any $M > 0$. Therefore subtracting a positive constant to v_c if necessary, we may assume that

$$1 \leq u_c(x, t) - v_c(x) \leq C \quad \text{for all } (x, t) \in \overline{\Omega} \times [0, \infty) \text{ and some } C > 0 \quad (20)$$

and we fix such a constant C . We define the functions $\mu_\eta^+, \mu_\eta^- : [0, \infty) \rightarrow \mathbb{R}$ by

$$\begin{aligned} \mu_\eta^+(s) &:= \min_{x \in \overline{\Omega}, t \geq s} \left(\frac{u_c(x, t) - v_c(x) + \eta(t - s)}{u_c(x, s) - v_c(x)} \right), \\ \mu_\eta^-(s) &:= \max_{x \in \overline{\Omega}, t \geq s} \left(\frac{u_c(x, t) - v_c(x) - \eta(t - s)}{u_c(x, s) - v_c(x)} \right) \end{aligned} \quad (21)$$

for $\eta \in (0, \eta_0]$. By the uniform continuity of u_c and v_c , we have $\mu_\eta^\pm \in C([0, \infty))$. It is easily seen that $0 \leq \mu_\eta^+(s) \leq 1$ and $\mu_\eta^-(s) \geq 1$ for all $s \in [0, \infty)$ and $\eta \in (0, \eta_0]$.

Proposition 3.2. *Assume that (BA) holds.*

- (i) *Assume that $(A2)_c^+$ holds. We have $\mu_\eta^+(s) \rightarrow 1$ as $s \rightarrow \infty$ for any $\eta \in (0, \eta_0]$.*
- (ii) *Assume that $(A2)_c^-$ holds. We have $\mu_\eta^-(s) \rightarrow 1$ as $s \rightarrow \infty$ for any $\eta \in (0, \eta_0]$.*

Proposition 3.2 is a consequence of the following lemma.

Lemma 3.3 (Key Lemma). *Assume that (BA) holds. Let C be the constant given by (20).*

- (i) *Assume that $(A2)_c^+$ holds. The function μ_η^+ is a supersolution of*

$$\max\{w(s) - 1, w'(s) + \frac{\psi_\eta}{C}(w(s) - 1)\} = 0 \text{ in } (0, \infty) \quad (22)$$

for any $\eta \in (0, \eta_0]$.

- (ii) *Assume that $(A2)_c^-$ holds. The function μ_η^- is a subsolution of*

$$\max\{w(s) - 1, w'(s) + \frac{\psi_\eta}{C} \cdot \frac{w(s) - 1}{w(s)}\} = 0 \text{ in } (0, \infty) \quad (23)$$

for any $\eta \in (0, \eta_0]$.

We first prove Proposition 3.2 and Theorem 3.1 by using Lemma 3.3 and we postpone to prove Lemma 3.3 in the next section.

Proof of Proposition 3.2. Fix $\eta \in (0, \eta_0]$. We first notice that, by definition,

$$\mu_\eta^+(s) \leq 1 \leq \mu_\eta^-(s),$$

for any $s \geq 0$. On the other hand, one checks easily that the functions

$$1 + (\mu_\eta^+(0) - 1) \exp\left(-\frac{\psi_\eta}{C}t\right),$$

and

$$1 + (\mu_\eta^-(0) - 1) \exp\left(-\frac{\psi_\eta}{C\mu_\eta^-(0)}t\right),$$

are, respectively, sub and supersolution of (22) or (23) and they take the same value as μ_η^+ and μ_η^- for $s = 0$, respectively. Therefore, by comparison, we have for any $s > 0$

$$1 + (\mu_\eta^+(0) - 1) \exp\left(-\frac{\psi_\eta}{C}t\right) \leq \mu_\eta^+(s),$$

and

$$\mu_\eta^-(s) \leq 1 + (\mu_\eta^-(0) - 1) \exp\left(-\frac{\psi_\eta}{C\mu_\eta^-(0)}t\right),$$

which imply the conclusion. \square

Proof of Theorem 3.1. We only prove (i), since we can prove (ii) similarly. For any $x \in \overline{\Omega}$ and $t, s \in [0, \infty)$ with $t \geq s$, we have

$$\mu_\eta^+(s)(u_c(x, s) - v_c(x)) \leq u_c(x, t) - v_c(x) + \eta(t - s),$$

which implies

$$\begin{aligned} u_c(x, s) - u_c(x, t) + \eta(s - t) &\leq (1 - \mu_\eta^+(s))(u_c(x, s) - v_c(x)) \\ &\leq C(1 - \mu_\eta^+(s)) =: \delta_\eta(s). \end{aligned}$$

Therefore we have the conclusion. \square

4. PROOF OF LEMMA 3.3

In this section we divide into three subsections and we prove Lemma 3.3 for three types boundary problems. We only prove (i), since we can prove (ii) similarly in any case.

4.1. Neumann Problem. In this subsection we normalize the additive eigenvalue c_n to be 0 by replacing H by $H - c_n$, where c_n is the constant given by (13).

Proof of Lemma 3.3. Fix $\mu \in (0, \eta_0]$ and let μ_η^+ be the function given by (21). By abuse of notation we write μ for μ_η^+ . We recall that $\mu(s) \leq 1$ for any $s \geq 0$.

Let $\phi \in C^1((0, \infty))$ and $\sigma > 0$ be a strict local minimum of $\mu - \phi$, i.e., $(\mu - \phi)(s) > (\mu - \phi)(\sigma)$ for all $s \in [\sigma - \delta, \sigma + \delta] \setminus \{\sigma\}$ and some $\delta > 0$. Since there is nothing to check $\mu(\sigma) = 1$, we assume that $\mu(\sigma) < 1$. We choose $\xi \in \overline{\Omega}$ and $\tau \geq \sigma$ such that

$$\mu(\sigma) = \frac{u(\xi, \tau) - v(\xi) + \eta(\tau - \sigma)}{u(\xi, \sigma) - v(\xi)}.$$

We only consider the case where $\xi \in \partial\Omega$ in this proof.

Choose a positive definite symmetric matrix A which satisfies $A\gamma(\xi) = n(\xi)$ (see [26, Lemma 4.1]) and define the quadratic function $f : \overline{\Omega} \rightarrow \mathbb{R}$ by

$$f(x) := \frac{1}{2}Ax \cdot x. \tag{24}$$

For $\varepsilon, \alpha \in (0, 1)$, we set $K := \overline{\Omega}^3 \times \{(t, s) \mid t \geq s, s \in [\sigma - \delta, \sigma + \delta]\}$ and define the function $\Psi : K \rightarrow \mathbb{R}$ by

$$\begin{aligned} & \Psi(x, y, z, t, s) \\ & := \frac{u(x, t) - v(z) + \eta(t - s) + g(\xi)(d(x) - d(z))}{u(y, s) - v(z) + g(\xi)(d(y) - d(z))} - \phi(s) \\ & \quad + \frac{1}{\varepsilon^2}(f(x - y) + f(x - z)) + |x - \xi|^2 - \alpha(d(x) + d(z) - 2d(y)), \end{aligned}$$

where

$$d(x) := -\frac{\gamma(\xi) \cdot x}{|\gamma(\xi)|^2}. \quad (25)$$

Let Ψ achieve its minimum over K at some $(\bar{x}, \bar{y}, \bar{z}, \bar{t}, \bar{s})$. Set

$$\begin{aligned} \bar{\mu}_1 & := u(\bar{y}, \bar{s}) - v(\bar{z}) + g(\xi)(d(\bar{y}) - d(\bar{z})), \\ \bar{\mu}_2 & := u(\bar{x}, \bar{t}) - v(\bar{z}) + \eta(\bar{t} - \bar{s}) + g(\xi)(d(\bar{x}) - d(\bar{z})), \\ \bar{\mu} & := \frac{\bar{\mu}_2}{\bar{\mu}_1}. \end{aligned}$$

Since $\overline{\Omega}$ is compact, by taking a subsequence if necessary, we may assume that

$$\bar{x}, \bar{y}, \bar{z} \rightarrow \xi \text{ and } \bar{s} \rightarrow \sigma \text{ as } \varepsilon \rightarrow 0.$$

Moreover, we have

$$\bar{\mu} \rightarrow \mu(\sigma) \text{ as } \varepsilon \rightarrow 0, \quad (26)$$

$$\frac{|\bar{x} - \bar{y}|}{\varepsilon^2} + \frac{|\bar{x} - \bar{z}|}{\varepsilon^2} \leq C_1 \text{ for some } C_1 > 0. \quad (27)$$

The second property is due to the Lipschitz continuity of u and v . From (26), we may assume that $\bar{\mu} < 1$ for small $\varepsilon > 0$.

Derivating Ψ with respect to each variable x, t, y, s, z at $(\bar{x}, \bar{y}, \bar{z}, \bar{t}, \bar{s})$ formally, we get

$$\begin{aligned} D_x u(\bar{x}, \bar{t}) & = \bar{\mu}_1 \left(\frac{1}{\varepsilon^2} A(\bar{y} - \bar{x}) + \frac{1}{\varepsilon^2} A(\bar{z} - \bar{x}) + 2(\xi - \bar{x}) + \alpha Dd(\bar{x}) \right) \\ & \quad - g(\xi) Dd(\bar{x}), \\ u_t(\bar{x}, \bar{t}) & = -\eta, \\ D_y u(\bar{y}, \bar{s}) & = \frac{\bar{\mu}_1}{\bar{\mu}} \left(\frac{1}{\varepsilon^2} A(\bar{y} - \bar{x}) - \alpha Dd(\bar{y}) \right) - g(\xi) Dd(\bar{y}) \\ u_s(\bar{y}, \bar{s}) & = -\frac{1}{\bar{\mu}} (\eta + \bar{\mu}_1 \phi'(\bar{s})), \\ D_z v(\bar{z}) & = \frac{\bar{\mu}_1}{1 - \bar{\mu}} \left(\frac{1}{\varepsilon^2} A(\bar{z} - \bar{x}) - \alpha Dd(\bar{z}) \right) - g(\xi) Dd(\bar{z}). \end{aligned}$$

We remark that we should interpret $D_x u, u_t, D_y u, u_s$ and $D_z v$ as the viscosity solution sense here.

We consider the case where $\bar{x} \in \partial\Omega$ and then we obtain

$$D_x u(\bar{x}, \bar{t}) \cdot \gamma(\bar{x}) \leq g(\bar{x}) + m(\varepsilon) - \alpha n(\xi) \cdot \gamma(\xi)$$

for some modulus of continuity m by similar arguments in the proof of Theorem 7.1. Therefore, if $\varepsilon > 0$ is suitable small compared to $\alpha > 0$, we have

$$D_x u(\bar{x}, \bar{t}) \cdot \gamma(\bar{x}) < g(\bar{x}).$$

Similarly, if $\bar{y} \in \partial\Omega$ and $\bar{z} \in \partial\Omega$, then we have

$$D_y u(\bar{y}, \bar{s}) \cdot \gamma(\bar{y}) > g(\bar{y}) \quad \text{and}$$

$$D_z v(\bar{z}) \cdot \gamma(\bar{z}) > g(\bar{z}),$$

respectively, for $\varepsilon > 0$ which is small enough compared to $\alpha > 0$.

Therefore, by the definition of viscosity solutions of (CN), we have

$$-\eta + H(\bar{x}, D_x u(\bar{x}, \bar{t})) \geq 0, \quad (28)$$

$$-\frac{1}{\bar{\mu}}(\eta + \bar{\mu}_1 \phi'(\bar{s})) + H(\bar{y}, D_y u(\bar{y}, \bar{s})) \leq 0, \quad (29)$$

$$H(\bar{z}, D_z v(\bar{z})) \leq 0. \quad (30)$$

In view of (27) we may assume that

$$\frac{1}{\varepsilon^2} A(\bar{y} - \bar{x}) \rightarrow p_y, \quad \frac{1}{\varepsilon^2} A(\bar{z} - \bar{x}) \rightarrow p_z$$

as $\varepsilon \rightarrow 0$ for some $p_y, p_z \in \mathbb{R}^N$ by taking a subsequence if necessary. Set

$$\mu_1 := u(\xi, \sigma) - v(\xi),$$

$$P := \frac{\mu_1}{\mu(\sigma)} p_y,$$

$$Q := \frac{\mu_1}{1 - \mu(\sigma)} p_z,$$

$$\tilde{P} := \mu(\sigma)(P - Q) \quad \text{and}$$

$$\tilde{Q} := Q - g(\xi) Dd(\xi).$$

Sending $\varepsilon \rightarrow 0$ and then $\alpha \rightarrow 0$ in (28) and (30), we obtain

$$H(\xi, \tilde{P} + \tilde{Q}) \geq \eta \quad \text{and} \quad H(\xi, \tilde{Q}) \leq 0. \quad (31)$$

Therefore, since we have $\tilde{P}/\mu(\sigma) + \tilde{Q} = P - g(\xi) Dd(\xi)$, by using (A2) $_0^+$, we obtain

$$H(\xi, \tilde{P} + \tilde{Q}) \leq \mu(\sigma) H(\xi, P - g(\xi) Dd(\xi)) - \psi_\eta(1 - \mu(\sigma)) \quad (32)$$

for some $\psi_\eta > 0$.

Sending $\varepsilon \rightarrow 0$ and then $\alpha \rightarrow 0$ in (29), we have

$$-\frac{1}{\mu(\sigma)}(\eta + \mu_1\phi'(\sigma)) + H(\xi, P - g(\xi)Dd(\xi)) \leq 0. \quad (33)$$

Therefore by (31), (32) and (33) we obtain

$$\eta \leq H(\xi, \tilde{P} + \tilde{Q}) \leq \eta + \mu_1\phi'(\sigma) + \psi_\eta(\mu(\sigma) - 1),$$

which implies the conclusion. \square

4.2. State Constraint Problem. In this subsection we normalize the additive eigenvalue c_{sc} is 0 by replacing H by $H - c_{sc}$, where c_{sc} is the constant given by (14).

Proof of Lemma 3.3. Fix $\eta \in (0, \eta_0]$ and let μ_η^+ be the function given by (21). By abuse of notation we write μ for μ_η^+ . Let $\mu - \phi$ take a strict local minimum at $\sigma > 0$ and for some $\phi \in C^1((0, \infty))$, i.e., $(\mu - \phi)(t) > (\mu - \phi)(\sigma)$ for all $t \in [\sigma - \delta, \sigma + \delta] \setminus \{\sigma\}$ and some $\delta > 0$. Since there is nothing to check in the case where $\mu(\sigma) = 1$, we assume that $\mu(\sigma) < 1$. We choose $\xi \in \bar{\Omega}$ and $\tau \geq \sigma$ such that

$$\mu(\sigma) = \frac{u(\xi, \tau) - v(\xi) + \eta(\tau - \sigma)}{u(\xi, \sigma) - v(\xi)}.$$

We only consider the case where $\xi \in \partial\Omega$.

Since Ω is a domain with a C^0 -boundary, after relabelling and re-orienting the coordinates axes if necessary, we may assume that $B(\xi, r) \cap \Omega = B(\xi, r) \cap \{(x', x_N) \in \mathbb{R}^{N-1} \times \mathbb{R} \mid x_N > b(x')\}$ for some constant $r > 0$ and some continuous function b on \mathbb{R}^{N-1} . Replacing r by a smaller positive number if necessary, we may assume that there is an $a_0 > 0$ such that if $0 < a \leq a_0$, then we have $B(\xi, r) \cap \bar{\Omega} \subset -ae_N + \Omega$, where $e_N := (0, \dots, 0, 1) \in \mathbb{R}^N$. Henceforth we assume that $0 < a \leq a_0$. We may choose a bounded open neighborhood W_a of $\bar{B}(\xi, r) \cap \bar{\Omega}$ such that $\bar{W}_a \subset -ae_N + \Omega$.

Set

$$u^a(x, t) := u(x + ae_N, t) \text{ and } v^a(x) := v(x + ae_N)$$

for $x \in -ae_N + \Omega$ and $t \in [0, \infty)$. It is clear to see that

$$\begin{aligned} (u^a)_t + H(x + ae_N, Du^a(x, t)) &= 0 \text{ in } W_a \times (\sigma - \delta, \sigma + \delta), \\ H(x + ae_N, Dv^a(x, t)) &= 0 \text{ in } W_a \end{aligned}$$

in the viscosity sense. We set

$$\mu_a(s) := \min_{x \in \bar{\Omega}, t \geq s} \left(\frac{u(x, t) - v^a(x) + \eta(t - s)}{u^a(x, s) - v^a(x)} \right)$$

and let $\mu_a - \phi$ take a minimum at some $\sigma_a \in [\sigma - \delta, \sigma + \delta]$ and then we have $\sigma_a \rightarrow \sigma$ and $\mu_a(\sigma_a) \rightarrow \mu(\sigma)$ as $a \rightarrow 0$. We choose $\xi_a \in \overline{\Omega}$ and $\tau_a \geq \sigma_a$ such that

$$\mu_a(\sigma) = \frac{u(\xi_a, \tau_a) - v^a(\xi_a) + \eta(\tau_a - \sigma_a)}{u^a(\xi_a, \sigma_a) - v^a(\xi_a)}.$$

We set $K = (\overline{B}(\xi, r) \cap \overline{\Omega}) \times \overline{W}_a^2 \times \{(t, s) \mid t \geq s, s \in [\sigma - \delta, \sigma + \delta]\}$ and define the function $\Psi : K \rightarrow \mathbb{R}$ by

$$\begin{aligned} \Psi(x, y, z, t, s) := & \frac{u(x, t) - v^a(z) + \eta(t - s)}{u^a(y, s) - v^a(z)} \\ & + \frac{1}{2\varepsilon^2}(|x - y|^2 + |x - z|^2) + |x - \xi_a|^2 + |s - \sigma_a|^2 - \phi(s). \end{aligned}$$

Let Ψ achieve its minimum on K at some $(\overline{x}, \overline{y}, \overline{z}, \overline{t}, \overline{s})$ and set

$$\begin{aligned} \overline{\mu}_{1a} &:= u^a(\overline{y}, \overline{s}) - v^a(\overline{z}), \\ \overline{\mu}_{2a} &:= u(\overline{x}, \overline{t}) - v^a(\overline{z}) + \eta(\overline{t} - \overline{s}), \\ \overline{\mu}_a &:= \frac{\overline{\mu}_{2a}}{\overline{\mu}_{1a}}. \end{aligned}$$

Since we may assume by taking a subsequence if necessary that

$$\overline{x}, \overline{y}, \overline{z} \rightarrow \xi_a, \overline{s} \rightarrow \sigma_a, \text{ and } \overline{\mu}_a \rightarrow \mu_a(\sigma_a) \text{ as } \varepsilon \rightarrow 0,$$

we have $\overline{y}, \overline{z} \in W_a$ if ε is small enough.

Therefore, by the definition of viscosity solutions we have

$$-\eta + H(\overline{x}, D_x u(\overline{x}, \overline{t})) \geq 0, \quad (34)$$

$$-\frac{1}{\overline{\mu}_a}(\eta + \overline{\mu}_{1a}(\phi'(\overline{s}) - 2(\overline{s} - \sigma_a)) + H(\overline{y} + ae_N, D_y u^a(\overline{y}, \overline{s}))) \leq 0, \quad (35)$$

$$H(\overline{z} + ae_N, D_z v^a(\overline{z})) \leq 0, \quad (36)$$

where

$$\begin{aligned} D_x u(\overline{x}, \overline{t}) &= \overline{\mu}_{1a} \left(\frac{\overline{y} - \overline{x}}{\varepsilon^2} + \frac{\overline{z} - \overline{x}}{\varepsilon^2} + 2(\xi_a - \overline{x}) \right), \\ D_y u^a(\overline{y}, \overline{s}) &= \frac{\overline{\mu}_{1a}}{\overline{\mu}_a} \cdot \frac{\overline{y} - \overline{x}}{\varepsilon^2}, \\ D_z v^a(\overline{z}) &= \frac{\overline{\mu}_{1a}}{1 - \overline{\mu}_a} \cdot \frac{\overline{z} - \overline{x}}{\varepsilon^2}. \end{aligned}$$

In view of the Lipschitz continuity of u^a on \overline{W}_a , we see that $\{\frac{\overline{y}-\overline{x}}{\varepsilon^2}\}_\varepsilon$, $\{\frac{\overline{z}-\overline{x}}{\varepsilon^2}\}_\varepsilon$ are bounded uniformly in $\varepsilon > 0$. Thus we may assume that

$$\begin{aligned}\frac{\overline{y}-\overline{x}}{\varepsilon^2} &\rightarrow p_y^a, & \frac{\overline{z}-\overline{x}}{\varepsilon^2} &\rightarrow p_z^a, \\ \overline{\mu}_{1a} &\rightarrow \mu_{1a} := u^a(\xi_a, \sigma_a) - v^a(\xi_a), \\ \overline{\mu}_a &\rightarrow \mu_a(\sigma_a)\end{aligned}$$

as $\varepsilon \rightarrow 0$ for some $p_y^a, p_z^a \in \mathbb{R}^N$ by taking a subsequence if necessary. Set

$$\begin{aligned}P_a &:= \frac{\mu_{1a}}{\mu_a(\sigma_a)} p_y^a, \\ Q_a &:= \frac{\mu_{1a}}{1 - \mu_a(\sigma_a)} p_z^a \quad \text{and} \\ \tilde{P}_a &:= \mu_a(\sigma_a)(P_a - Q_a).\end{aligned}$$

Sending $\varepsilon \rightarrow 0$ in inequalities (34), (35) and (36) yields

$$\begin{aligned}-\eta + H(\xi_a, \tilde{P}_a + Q_a) &\geq 0, \\ -\frac{1}{\mu_a(\sigma_a)}(\eta + \mu_{1a}\phi'(\sigma_a)) + H(\xi_a + ae_N, P_a) &\leq 0, \\ H(\xi_a + ae_N, Q_a) &\leq 0.\end{aligned}\tag{37}$$

Note that $|P_a| + |Q_a| \leq R$ for some $R > 0$ which is independent of a . There exists a modulus ω_R such that

$$H(\xi_a + ae_N, \tilde{P}_a + Q_a) \geq \eta - \omega_R(a).\tag{38}$$

For small $a > 0$ we have

$$H(\xi_a + ae_N, \tilde{P}_a + Q_a) \geq \frac{\eta}{2}.$$

Since we have $\tilde{P}_a/\mu_a(\sigma_a) + \tilde{Q}_a = P_a$, by using (A2) $_0^+$, we obtain

$$\mu_a(\sigma_a)H(\xi_a + ae_N, P_a) \geq H(\xi_a + ae_N, \tilde{P}_a + Q_a) + \psi_\eta(1 - \mu_a(\sigma_a))\tag{39}$$

for a constant $\psi_\eta > 0$.

By (38), (39) and (37) we get

$$\begin{aligned}\eta - \omega_R(a) &\leq H(\xi_a + ae_N, \tilde{P}_a + Q_a) \\ &\leq \mu_a(\sigma_a)H(\xi_a + ae_N, P_a) - \psi_\eta(1 - \mu_a(\sigma_a)) \\ &\leq \eta + \mu_{1a}\phi'(\sigma_a) - \psi_\eta(1 - \mu_a(\sigma_a)).\end{aligned}$$

We divide by $\mu_{1a} > 0$ and then we obtain

$$\begin{aligned}0 &\leq \phi'(\sigma_a) + \frac{\psi_\eta}{\mu_{1a}}(\mu_a(\sigma_a) - 1) + \frac{\omega_R(a)}{\mu_{1a}} \\ &\leq \phi'(\sigma_a) + \frac{\psi_\eta}{C}(\mu_a(\sigma_a) - 1) + \omega_R(a).\end{aligned}$$

Sending $a \rightarrow 0$ yields

$$\phi'(\sigma) + \frac{\psi_\eta}{C}(\mu(\sigma) - 1) \geq 0,$$

which is the conclusion. \square

4.3. Dirichlet Problem. Let u be the solution of (CD) and c_{sc} be the constant defined by (14).

We first treat the case where $c_{sc} > 0$. Set $u_{c_{sc}}(x, t) = u(x, t) + c_{sc}t$ for $(x, t) \in \bar{\Omega} \times [0, \infty)$ and $g_{c_{sc}}(x, t) = g(x) + c_{sc}t$ for $(x, t) \in \partial\Omega \times (0, \infty)$. Since $u_{c_{sc}}$ is bounded on $\bar{\Omega} \times [0, \infty)$ by Proposition 2.1 (iii-a) and $g_{c_{sc}}(x, t) \rightarrow \infty$ uniformly for $x \in \partial\Omega$ as $t \rightarrow \infty$, there exists a constant $\bar{t} > 0$ such that $g_{c_{sc}}(x, t) > u_{c_{sc}}(x, t)$ for all $(x, t) \in (\bar{t}, \infty)$. As $u_{c_{sc}}$ satisfies

$$\begin{cases} (u_{c_{sc}})_t + H(x, Du_{c_{sc}}) = c_{sc} & \text{in } \Omega \times (0, \infty), \\ u_{c_{sc}}(x, t) = g_{c_{sc}}(x, t) & \text{on } \partial\Omega \times (0, \infty) \end{cases}$$

in the viscosity sense, we see easily that $u_{c_{sc}}$ is a solution of

$$\begin{cases} (u_{c_{sc}})_t + H(x, Du_{c_{sc}}) \leq c_{sc} & \text{in } \Omega \times (\bar{t}, \infty), \\ (u_{c_{sc}})_t + H(x, Du_{c_{sc}}) \geq c_{sc} & \text{on } \bar{\Omega} \times (\bar{t}, \infty). \end{cases}$$

Therefore the large-time asymptotic behavior of $u_{c_{sc}}$ is same as the large-time asymptotic behavior of solutions of (SC).

Next we consider the case where $c_{sc} \leq 0$. Let v be a subsolution of (D)₀. Since u is bounded on $\bar{\Omega} \times [0, \infty)$ by Proposition 2.1 (iii-b) and $v - M$ are still subsolutions of (D)₀ for any $M > 0$, by subtracting a positive constant to v if necessary, we may assume that $u - v$ satisfies (20).

Let μ_η^+ be the function defined by (21) for $\eta \in (0, \eta_0]$. By abuse of notation we write μ for μ_η^+ . We prove that μ satisfies (22). We first notice that, in view of the coercivity of H , we have

$$\begin{aligned} u(x, t) &\leq g(x) \text{ and} \\ v(x) &\leq g(x) \quad \text{for all } (x, t) \in \partial\Omega \times (0, \infty) \end{aligned} \tag{40}$$

by Proposition 7.2.

Let $\varepsilon > 0$, $\xi \in \bar{\Omega}$, $\sigma, \tau \in (0, \infty)$ and $(\bar{x}, \bar{t}) \in \bar{\Omega} \times [0, \infty)$ be the same as those in the proof of Lemma 3.3 in the case of (SC). We only consider the case where $\mu(\tau) < 1$ and $\xi \in \partial\Omega$. Since we have

$$0 \leq \mu(\tau) = \frac{u(\xi, \tau) - v(\xi) + \eta(\tau - \sigma)}{u(\xi, \sigma) - v(\xi)} < 1$$

and $u(\xi, \sigma) - v(\xi) > 0$, we obtain

$$u(\xi, \tau) - v(\xi) + \eta(\tau - \sigma) < u(\xi, \sigma) - v(\xi) \leq g(\xi) - v(\xi),$$

which implies $u(\xi, \tau) < g(\xi)$. Therefore we have $u(\bar{x}, \bar{t}) - g(\bar{x}) < 0$ for small $\varepsilon > 0$. The rest of the proof follows by the same method as in subsection 4.2.

5. CONVERGENCE

In this section, we prove Theorem 1.3 by using Theorem 3.1.

Proof of Theorem 1.3. Let c be the associated additive eigenvalue. When we consider (CD), let us set $c := c_{sc}$ if $c_{sc} > 0$ and $c := 0$ if $c_{sc} \leq 0$.

Since $\{u_c(\cdot, t)\}_{t \geq 0}$ is compact in $W^{1,\infty}(\Omega)$, there exists a sequence $\{u_c(\cdot, T_n)\}_{n \in \mathbb{N}}$ which converges uniformly on $\bar{\Omega}$ as $n \rightarrow \infty$. The maximum principle implies that we have

$$\|u_c(\cdot, T_n + \cdot) - u_c(\cdot, T_m + \cdot)\|_{L^\infty(\Omega \times (0, \infty))} \leq \|u_c(\cdot, T_n) - u_c(\cdot, T_m)\|_{L^\infty(\Omega)}$$

for any $n, m \in \mathbb{N}$. Therefore, $\{u_c(\cdot, T_n + \cdot)\}_{n \in \mathbb{N}}$ is a Cauchy sequence in $\text{BUC}(\bar{\Omega} \times [0, \infty))$ and it converges to a function denoted by $u_c^\infty \in \text{BUC}(\bar{\Omega} \times [0, \infty))$.

Fix any $x \in \bar{\Omega}$ and $s, t \in [0, \infty)$ with $t \geq s$. By Theorem 3.1 we have

$$u_c(x, s + T_n) - u_c(x, t + T_n) + \eta(s - t) \leq \delta_\eta(s + T_n)$$

or

$$u_c(x, t + T_n) - u_c(x, s + T_n) - \eta(t - s) \leq \delta_\eta(s + T_n)$$

for any $n \in \mathbb{N}$ and $\eta > 0$. Sending $n \rightarrow \infty$ and then $\eta \rightarrow 0$, we get, for any $t \geq s$

$$u_c^\infty(x, s) \leq u_c^\infty(x, t).$$

or

$$u_c^\infty(x, t) \leq u_c^\infty(x, s).$$

Therefore, we see that the functions $x \mapsto u_c^\infty(x, t)$ are uniformly bounded and equi-continuous which are also monotone in t . This implies that $u_c^\infty(x, t) \rightarrow w(x)$ uniformly on $\bar{\Omega}$ as $t \rightarrow \infty$ for some $w \in W^{1,\infty}(\Omega)$. Moreover, by a standard stability property of viscosity solutions, w is a solution of either (E) or (D)₀.

Since $u_c(\cdot, T_n + \cdot) \rightarrow u_c^\infty$ uniformly in $\bar{\Omega} \times [0, \infty)$ as $n \rightarrow \infty$, we have

$$-o_n(1) + u_c^\infty(x, t) \leq u_c(x, T_n + t) \leq u_c^\infty(x, t) + o_n(1),$$

where $o_n(1) \rightarrow 0$ as $n \rightarrow \infty$, uniformly in x and t . Taking the half-relaxed semi-limits as $t \rightarrow +\infty$, we get

$$-o_n(1) + v(x) \leq \liminf_{t \rightarrow \infty}^* [u_c](x, t) \leq \limsup_{t \rightarrow \infty}^* [u_c](x, t) \leq v(x) + o_n(1).$$

Sending $n \rightarrow \infty$ yields

$$v(x) = \liminf_{t \rightarrow \infty}^* [u_c](x, t) = \limsup_{t \rightarrow \infty}^* [u_c](x, t)$$

for all $x \in \overline{\Omega}$. □

Remark 1.

(i) The Lipschitz regularity assumption on u_0 is convenient to avoid technicalities but it is not necessary. We can remove it as follows. We may choose a sequence $\{u_0^k\}_{k \in \mathbb{N}} \subset W^{1,\infty}(\Omega) \cap C(\overline{\Omega})$ so that $\|u_0^k - u_0\|_{L^\infty(\Omega)} \leq 1/k$ for all $n \in \mathbb{N}$. By the maximum principle, we have

$$\|u - u_k\|_{L^\infty(\Omega \times (0, \infty))} \leq \|u_0^k - u_0\|_{L^\infty(\Omega)} \leq 1/k$$

and therefore

$$u_k(x, t) - 1/k \leq u(x, t) \leq u_k(x, t) + 1/k \text{ for all } (x, t) \in \overline{\Omega} \times [0, \infty),$$

where u is the solution of (IB) and u_k are the solutions of (IB) with $u_0 = u_0^k$. Therefore we have

$$u_\infty^k(x) - 1/k \leq \liminf_{t \rightarrow \infty} u(x, t) \leq \limsup_{t \rightarrow \infty} u(x, t) \leq u_\infty^k(x) + 1/k$$

for all $x \in \overline{\Omega}$, where $u_\infty^k(x) = \lim_{t \rightarrow \infty} u_k(x, t)$. Thus, $|\liminf_{t \rightarrow \infty} u(x, t) - \limsup_{t \rightarrow \infty} u(x, t)| \leq 2/k$ for all $k \in \mathbb{N}$ and $x \in \overline{\Omega}$, which implies that

$$\liminf_{t \rightarrow \infty} u(x, t) = \limsup_{t \rightarrow \infty} u(x, t)$$

for all $x \in \overline{\Omega}$. We note that, by the same argument, we can obtain the asymptotic monotone property, Theorem 3.1, of solutions of (IB) without the Lipschitz continuity of solutions.

(ii) We remark that modifying $(A2)_a^\pm$ as in [4], we can generalize Theorem 1.3. We use the following assumptions instead of $(A2)_a^\pm$. We denote $H - a$ by H_a .

$(A3)_a^+$ There exists a closed set $K \subset \overline{\Omega}$ (K is possibly empty) having the properties

$$(i) \min_{p \in \mathbb{R}^N} H_a(x, p) = 0 \text{ for all } x \in K,$$

(ii) for each $\varepsilon > 0$ there exists a modulus $\psi_\varepsilon(r) > 0$ for all $r > 0$ and $\eta_0^\varepsilon > 0$ such that for all $\eta \in (0, \eta_0^\varepsilon]$ if $\text{dist}(x, K) \geq \varepsilon$, $H_a(x, p + q) \geq \eta$ and $H_a(x, q) \leq 0$ for some $x \in \overline{\Omega}$ and $p, q \in \mathbb{R}$, then for any $\mu \in (0, 1]$,

$$\mu H_a(x, \frac{p}{\mu} + q) \geq H_a(x, p + q) + \psi_\varepsilon(\eta)(1 - \mu).$$

$(A3)_a^-$ There exists a closed set $K \subset \overline{\Omega}$ (K is possibly empty) having the properties

- (i) $\min_{p \in \mathbb{R}^N} H_a(x, p) = 0$ for all $x \in K$,
- (ii) for each $\varepsilon > 0$ there exists a modulus $\psi_\varepsilon(r) > 0$ for all $r > 0$ and $\eta_0^\varepsilon > 0$ such that for all $\eta \in (0, \eta_0^\varepsilon]$ if $\text{dist}(x, K) \geq \varepsilon$, $H_a(x, p + q) \leq -\eta$ and $H_a(x, q) \geq 0$ for some $x \in \overline{\Omega}$ and $p, q \in \mathbb{R}$, then for any $\mu \in (0, 1]$,

$$\mu H_a(x, \frac{p}{\mu} + q) \leq H_a(x, p + q) - \frac{\psi_\varepsilon(\eta)(\mu - 1)}{\mu}.$$

Theorem 5.1. *The results of Theorem 1.3 still hold if we replace assumptions (A2)_c⁺ or (A2)_c⁺ by (A3)_c⁺ or (A3)_c⁺ where c is defined as in Theorem 1.3.*

6. REMARKS ON CONVEX HAMILTON-JACOBI EQUATIONS

In this section we deal with convex Hamilton-Jacobi equations, i.e.,
 (A4) $p \mapsto H(x, p)$ is convex for any $x \in \overline{\Omega}$.

6.1. The Namah-Roquejoffre Case. We use the following assumptions in this subsection. We consider the Hamiltonian $H(x, p) = F(x, p) - f(x)$, where F and f are assumed to satisfy

- (A5) $F(x, p) \geq F(x, 0) = 0$ for all $(x, p) \in \overline{\Omega} \times \mathbb{R}^N$,
 (A6) $f(x) \geq 0$ for all $x \in \overline{\Omega}$ and $\mathcal{A}_f := \{x \in \overline{\Omega} \mid f(x) = 0\} \neq \emptyset$.
 (A7) $g(x) \geq 0$ for all $x \in \partial\Omega$.

A typical example of this Hamiltonian is $H(x, p) = |p| - f(x)$ with $f \geq 0$ on $\overline{\Omega}$ and in this case, it is clearly seen that H does not satisfy (A2)₀⁺ and (A2)₀⁻. Meanwhile, Hamiltonian has a simple structure and therefore in a relatively easy way we can obtain

Theorem 6.1. *Assume that (BA) and (A4)-(A6) hold.*

- (i) **(Neumann/oblique derivative problem)** *Assume that (A7) holds. For the solution u of (CN), (15) holds with a solution $(v, 0)$ of (E-N).*
- (ii) **(State constraint problem or Dirichlet problem)** *For the solution u of (SC) or (CD), (15) holds with a solution $(v, 0)$ of (E-S) or (D)₀.*

Sketch of Proof. It is easy to see that

$$c_n = 0 \text{ and } c_{sc} = 0.$$

Indeed, on one hand, in view of (A5)-(A7), any constant is a subsolution of (E-N) and (E-SC) with $a = 0$, which implies that $c_n, c_{sc} \leq 0$. On the other hand, if we assume $c_n < 0$ or $c_{sc} < 0$, then (A5), (A6) yields a contradiction since $H(x, p) = F(x, p) - f(x) \geq 0$ for all $x \in \mathcal{A}_f$.

We have $u_t = -F(x, Du) + f(x) \leq 0$ in $\mathcal{A}_f \times (0, \infty)$ and therefore we can get the monotonicity of the function $t \mapsto u(x, t)$ at least formally (see [31, Lemma 2.4] and [14, Lemma 4.4] for a rigorous proof). Thus, taking into account the uniform Lipschitz continuity of u , we have $\liminf_{*t \rightarrow \infty} u(x, t) = \limsup_{*t \rightarrow \infty} u(x, t)$ for any $x \in \mathcal{A}_f$. By [21, Theorem 6.6] for Neumann problems, [23, Theorem 7.3] for state constraint problems and [29, Theorem 5.3], we obtain $\liminf_{*t \rightarrow \infty} u(x, t) = \limsup_{*t \rightarrow \infty} u(x, t)$ for any $x \in \overline{\Omega}$, since, outside \mathcal{A}_f , one has a strict subsolution, which is a key tool to obtain the comparison of the half-relaxed limits. \square

Remark 2. We remark that the result of Theorem 6.1 is included in Theorem 5.1, since $(A3)_0^+$ holds with $K = \{x \in \mathbb{R}^N \mid f(x) = 0\}$.

6.2. Asymptotic Profile. In this subsection we give representation formulas for the asymptotic solutions in each case. We define the functions $\phi^-, \phi^\infty \in C(\overline{\Omega})$ by

$$\begin{aligned} \phi^-(x) &:= \sup\{v(x) \mid v \in C(\overline{\Omega}) \text{ is a subsolution of (E), } v \leq u_0 \text{ on } \overline{\Omega}\}, \\ \phi^\infty(x) &:= \inf\{v(x) \mid v \in C(\overline{\Omega}) \text{ is a solution of (E), } v \geq \phi^- \text{ on } \overline{\Omega}\}. \end{aligned}$$

We denote ϕ^- associated with (E-N) and (E-SC) by ϕ_n^- and ϕ_s^- , respectively and ϕ^∞ associated with (E-N) and (E-SC) by ϕ_n^∞ and ϕ_s^∞ , respectively. In view of (A4) we see that ϕ_n^∞ and ϕ_s^∞ are a solution of (E-N), (E-SC), respectively. We refer to the articles [5, 23, 21] for a stability result under infimum operation. When $c_{sc} \leq 0$, we define the functions $\phi_d^-, \phi_d^\infty \in C(\overline{\Omega})$ by

$$\begin{aligned} \phi_d^-(x) &:= \sup\{v(x) \in C(\overline{\Omega}) \mid v \in C(\overline{\Omega}) \text{ is a subsolution of (D)}_0\}, \\ \phi_d^\infty(x) &:= \inf\{v(x) \mid v \in C(\overline{\Omega}) \text{ is a solution of (D)}_0, v \geq \phi_s^- \wedge \phi_d^- \text{ on } \overline{\Omega}\}. \end{aligned}$$

Theorem 6.2 (Asymptotic Profile). *Assume that (BA) and (A4) hold.*

(i) **(Neumann/oblique derivative problem)** *Let u be the solution of (CN) and then we have*

$$\lim_{t \rightarrow \infty} (u(x, t) + c_n t) = \phi_n^\infty(x) \quad \text{uniformly for all } x \in \overline{\Omega}. \quad (41)$$

(ii) **(State constraint problem)** *Let u be the solution of (SC) and then we have*

$$\lim_{t \rightarrow \infty} (u(x, t) + c_{sc} t) = \phi_s^\infty(x) \quad \text{uniformly for all } x \in \overline{\Omega}. \quad (42)$$

(iii) **(Dirichlet problem)** *Let u be the solution of (CD).*

(iii-a) *If $c_{sc} > 0$, then*

$$\lim_{t \rightarrow \infty} (u(x, t) + c_{sc} t) = \phi_s^\infty(x) \quad \text{uniformly for all } x \in \overline{\Omega}. \quad (43)$$

(iii-b) If $c_{sc} = 0$, then

$$\lim_{t \rightarrow \infty} u(x, t) = \phi_d^\infty(x) \quad \text{uniformly for all } x \in \overline{\Omega}. \quad (44)$$

(iii-c) If $c_{sc} < 0$, then

$$\lim_{t \rightarrow \infty} u(x, t) = \phi_d^-(x) \quad \text{uniformly for all } x \in \overline{\Omega}.$$

Proof. We prove (i), (ii), (iii-a) and (iii-b) at the same time in order to avoid duplication of explanations. We denote ϕ_n^-, ϕ_s^- and $\phi_s^- \wedge \phi_d^-$ by ϕ^- , and $\phi_s^\infty, \phi_n^\infty$ and ϕ_d^∞ by ϕ^∞ in any case. We write $u_\infty(x)$ for the left hand side of (41), (42), (43) and (44) in any case. Let c denote c_n and c_{sc} in any case and set $u_c := u + ct$ on $\overline{\Omega} \times [0, \infty)$.

Since ϕ^- is a subsolution of (E-N), (E-SC) or (D) $_c$, respectively, ϕ^- is a subsolution of (IB) too. By the comparison principle for (IB) (cf. Theorem 7.1) we have

$$\phi^-(x) \leq u_c(x, t) \quad \text{for all } (x, s) \in \overline{\Omega} \times [0, \infty). \quad (45)$$

Therefore, we get

$$\phi^-(x) \leq u_\infty(x) \quad \text{for all } x \in \overline{\Omega}.$$

Note that u_∞ is a solution of (E-N), (E-SC) or (D) $_c$. By the definition of ϕ^∞ we get $\phi^\infty \leq u_\infty$ on $\overline{\Omega}$.

We define the functions $v_c \in \text{BUC}(\overline{\Omega} \times [0, \infty))$ by $v_c(x, t) := \inf_{s \geq t} u_c(x, s)$. By (45) we have particularly $\phi^- \leq v_c(\cdot, 0)$ on $\overline{\Omega}$. Note that $v_c(\cdot, 0)$ is a subsolution of (E-N), (E-SC) or (D) $_c$ and $v_c(\cdot, 0) \leq u_0$ on $\overline{\Omega}$. Indeed, v_c satisfies $(v_c)_t + H(x, Dv_c(x, t)) = 0$ in $\Omega \times (0, \infty)$ (see [5, 23, 21]) and it is clear that $v_c(x, \cdot)$ is non-decreasing, from which we have $(v_c)_t(x, t) \geq 0$ in $\Omega \times (0, \infty)$ in the viscosity sense. Therefore $H(x, Dv_c(x, t)) \leq 0$ in Ω for any $t \geq 0$. Since $v_c(x, t) \rightarrow v_c(x, 0)$ uniformly on $\overline{\Omega}$ as $t \rightarrow 0$, we see that $H(x, Dv_c(x, 0)) \leq 0$ in Ω in the viscosity sense. It is easily seen that $v_c(\cdot, 0) \leq u_0$ on $\overline{\Omega}$. Thus by the definition of ϕ^- , we have $v_c(\cdot, 0) \leq \phi^-$ on $\overline{\Omega}$. Therefore we obtain $v_c(\cdot, 0) = \phi^-$ on $\overline{\Omega}$.

Note that ϕ^∞ is a solution of (1), (2) and satisfies that $\phi^\infty(x) \geq \phi^-(x) = v_c(x, 0)$ on $\overline{\Omega}$. By the comparison principle for (IB) (cf. Theorem 7.1, again) we have for all $x \in \overline{\Omega}$,

$$u_\infty(x) = \lim_{t \rightarrow \infty} u_c(x, t) = \liminf_{t \rightarrow \infty} u_c(x, t) = \lim_{t \rightarrow \infty} v_c(x, t) \leq \phi^\infty(x),$$

which establishes formulas.

We finally consider (iii-c). In this case Problem (D) $_0$ has the unique solution. Since u_∞ and ϕ_d are solutions of (D) $_0$, we see that $u_\infty = \phi_d$ on $\overline{\Omega}$. \square

We finally give another formulas for $\phi_n^-, \phi_s^-, \phi_d^-, \phi_n^\infty, \phi_s^\infty$ and ϕ_d^∞ .

Proposition 6.3. *Assume that (BA) and (A4) hold. We have*

$$\phi_n^-(x) = \min\{d_n(x, y) + u_0(y) \mid y \in \overline{\Omega}\}, \quad (46)$$

$$\phi_n^\infty(x) = \min\{d_n(x, y) + \phi_n^-(y) \mid y \in \mathcal{A}_n\}, \quad (47)$$

$$\phi_s^-(x) = \min\{d_s(x, y) + u_0(y) \mid y \in \overline{\Omega}\}, \quad (48)$$

$$\phi_s^\infty(x) = \min\{d_s(x, y) + \phi_s^-(y) \mid y \in \mathcal{A}_s\}, \quad (49)$$

$$\phi_d^-(x) = \min\{d_0(x, y) + g(y) \mid y \in \partial\Omega\}, \quad (50)$$

$$\phi_d^\infty(x) = \min\{d_0(x, y) + \phi_s^-(y) \wedge \phi_d^-(y) \mid y \in \mathcal{A}_s\} \quad (51)$$

for any $x \in \overline{\Omega}$, where

$$d_n(x, y) := \sup\{v(x) - v(y) \mid v \text{ is a subsolution of (9), (10) with } a = c_n\},$$

$$d_s(x, y) := \sup\{v(x) - v(y) \mid v \text{ is a subsolution of (9) with } a = c_{sc}\},$$

$$d_0(x, y) := \sup\{v(x) - v(y) \mid v \text{ is a subsolution of (11) with } a = 0\},$$

$$\mathcal{A}_n := \{y \in \overline{\Omega} \mid d_n(\cdot, y) \text{ is a solution of (E-N)}\},$$

$$\mathcal{A}_s := \{y \in \overline{\Omega} \mid d_s(\cdot, y) \text{ is a solution of (E-SC)}\}.$$

Proof. We first see that (50) holds. We denote by w_d the right hand side of (50). By a standard stability result of viscosity solution ϕ_d^- is a subsolution of $(D)_0$. By Proposition 7.2 we have $\phi_d^- \leq g$ on $\partial\Omega$. Therefore we have $\phi_d^-(x) \leq d_0(x, y) + \phi_d^-(y) \leq d_0(x, y) + g(y)$ for any $x, y \in \overline{\Omega}$, which implies that $\phi_d^- \leq w_d$ on $\overline{\Omega}$. It is easily seen that w_d is a subsolution of $(D)_0$ and therefore $w_d \leq \phi_d^-$ on $\overline{\Omega}$.

We next prove (46) and (48). We denote ϕ_n^- , ϕ_s^- and $\phi_s^- \wedge \phi_d^-$ by ϕ^- and denote by w_- the right hand side of (46) and (48) in any case. Let d denote d_n or d_s in any case. By the definition of ϕ^- we see that ϕ^- is a subsolution of (E-N), (E-SC) or $(D)_0$ and $\phi^- \leq u_0$ on $\overline{\Omega}$. By the definition of d we have $\phi^-(x) \leq d(x, y) + \phi^-(y) \leq d(x, y) + u_0(y)$ for all $x, y \in \overline{\Omega}$, which implies $\phi^- \leq w_-$ on $\overline{\Omega}$. Note that w_- is a subsolution of (E-N), (E-SC) or $(D)_0$ and $w_-(x) \leq d(x, x) + u_0(x) = u_0(x)$ for all $x \in \overline{\Omega}$. By the definition of ϕ^- we obtain $\phi^- = w_-$ on $\overline{\Omega}$.

We finally prove (47), (49) and (51). We denote ϕ_n^∞ , ϕ_s^∞ and ϕ_d^∞ by ϕ^∞ and denote by w_∞ the right hand side of (47), (49) and (51) in any case. Let \mathcal{A} denote \mathcal{A}_n or \mathcal{A}_s in any case.

It is easy to see that $w_\infty(x) = \phi^-(x)$ for all $x \in \mathcal{A}$. By [21, Theorem 6.6], [23, Theorem 7.3] and [29, Theorem 5.3] we get $w_\infty \geq \phi^-$ on $\overline{\Omega}$. By the definition of ϕ^∞ we obtain $w_\infty \geq \phi^\infty$ on $\overline{\Omega}$. Note that $\phi^- \leq \phi^\infty$ on $\overline{\Omega}$. Then, we have

$$w_\infty(x) = \phi^-(x) \leq \phi^\infty(x) \quad \text{for all } x \in \mathcal{A}.$$

Therefore we get $w_\infty \leq \phi^\infty$ on $\overline{\Omega}$ by [21, Theorem 6.6], [23, Theorem 7.3] and [29, Theorem 5.3]. \square

6.3. On the compatibility condition (8). In this subsection we consider (CD) under assumptions (A1), (A4) and $(A2)_c^+$ or $(A2)_c^-$ and we remove the compatibility condition (8), where $c := c_{sc}$ if $c_{sc} > 0$ and $c := 0$ if $c_{sc} \leq 0$. Let u be a solution of (CD). We notice that, if we do not assume (8), then u may be discontinuous and therefore we interpret solutions as discontinuous viscosity solutions introduced in [19].

We consider the following approximate problems of (CD)

$$(CD)_k^1 \quad \begin{cases} u_t + H(x, Du) = 0 & \text{in } \Omega \times (0, \infty), \\ u(x, t) = g(x) & \text{on } \partial\Omega \times (0, \infty), \\ u(x, 0) = u_0^k(x) & \text{on } \overline{\Omega}, \end{cases}$$

and

$$(CD)_k^2 \quad \begin{cases} u_t + H(x, Du) = 0 & \text{in } \Omega \times (0, \infty), \\ u(x, t) = g_k(x, t) & \text{on } \partial\Omega \times (0, \infty), \\ u(x, 0) = u_0(x) & \text{on } \overline{\Omega}, \end{cases}$$

where

$$\begin{aligned} u_0^k(x) &:= \min_{y \in \overline{\Omega}} \{u_0^g(y) + k|x - y|^2\}, \\ u_0^g(x) &:= \begin{cases} u_0(x) & \text{for } x \in \Omega, \\ u_0(x) \wedge g(x) & \text{for } x \in \partial\Omega, \end{cases} \\ g_k(x, t) &:= \max\{g(x), u_0(x) - kt\} \end{aligned}$$

for all $x \in \overline{\Omega}$, $g \geq 0$ and $k \in \mathbb{N}$. Note that

$$u_0^k(x) \leq g(x) \text{ and } u_0(x) \leq g_k(x, t)$$

for all $(x, t) \in \partial\Omega \times [0, \infty)$ and $k \in \mathbb{N}$.

Let u_k^1 and u_k^2 be the solution of $(CD)_k^1$ and $(CD)_k^2$, respectively. By Theorem 6.2 we have for any $k \in \mathbb{N}$

$$\begin{cases} \text{if } c_{sc} > 0, & u_k^1(\cdot, t) + c_{sc}t \rightarrow \min\{d_s(\cdot, y) + \phi_k^-(y) \mid y \in \mathcal{A}_s\} \\ \text{if } c_{sc} = 0, & u_k^1(\cdot, t) \rightarrow \min\{d_s(\cdot, y) + \phi_k^-(y) \wedge \phi_d^-(y) \mid y \in \mathcal{A}_s\}, \\ \text{if } c_{sc} < 0, & u_k^1(\cdot, t) \rightarrow \phi_d^- \end{cases}$$

uniformly on $\overline{\Omega}$ as $t \rightarrow \infty$, where

$$\phi_k^-(x) := \min\{d_s(x, y) + u_0^k(y) \mid y \in \overline{\Omega}\} \text{ for all } x \in \overline{\Omega}.$$

By [30, Theorem 6.1] we have for any $k \in \mathbb{N}$

$$\begin{cases} \text{if } c_{sc} \geq 0, & u_k^2(\cdot, t) + c_{sc}t \rightarrow \min\{d_s(\cdot, y) + \phi_s^-(y) \wedge \phi_{g_k}^-(y) \mid y \in \mathcal{A}_s\} \\ \text{if } c_{sc} < 0, & u_k^2(\cdot, t) \rightarrow \phi_d^- \end{cases}$$

uniformly on $\overline{\Omega}$ as $t \rightarrow \infty$, where

$$\begin{aligned}\phi_{g_k}^- (x) &:= \inf\{d_s(x, y) + \underline{g}_k(y) \mid y \in \partial\Omega\} \text{ for all } x \in \overline{\Omega}, \\ \underline{g}_k(x) &:= \inf\{g_k(x, s) + c_{sc}s \mid s \geq 0\} \text{ for all } x \in \partial\Omega.\end{aligned}$$

Proposition 6.4. *We have $\phi_k^- \rightarrow \phi_s^-$ uniformly on $\overline{\Omega}$ as $k \rightarrow \infty$.*

Proof. It is easy to see that $\phi_k^-(x)$ is nondecreasing as $k \rightarrow \infty$ for any $x \in \overline{\Omega}$. We prove that $\phi_k^-(x)$ converges to $\phi_s^-(x) \wedge \phi_d^-(x)$ as $k \rightarrow \infty$ for any $x \in \overline{\Omega}$. Fix $x \in \overline{\Omega}$. First, we can easily see that $\limsup_{k \rightarrow \infty} \phi_k^-(x) \leq \phi_s^-(x) \wedge \phi_d^-(x)$, since $\phi_k^- \leq u_0^k \leq u_0^g$ on $\overline{\Omega}$.

Choose $y_k, z_k \in \overline{\Omega}$ such that

$$\begin{aligned}\phi_k^-(x) &= d_s(x, y_k) + u_0^k(y_k) \text{ and} \\ u_0^k(y_k) &= u_0^g(z_k) + k|y_k - z_k|^2.\end{aligned}$$

Since $\overline{\Omega}$ is compact, we may assume that $y_k, z_k \rightarrow y_0 \in \overline{\Omega}$ as $k \rightarrow \infty$ if taking a subsequence if necessary. We have

$$\begin{aligned}\liminf_{k \rightarrow \infty} \phi_k^-(x) &= \liminf_{k \rightarrow \infty} (d_s(x, y_k) + u_0^k(y_k)) \\ &\geq \liminf_{k \rightarrow \infty} (d_s(x, y_k) + u_0^g(z_k)) \\ &\geq d_s(x, y_0) + u_0^g(y_0) \\ &\geq \phi_s^-(x) \wedge \phi_d^-(x).\end{aligned}$$

Therefore we obtain $\limsup_{k \rightarrow \infty} \phi_k^-(x) = \liminf_{k \rightarrow \infty} \phi_k^-(x) = \phi_s^-(x) \wedge \phi_d^-(x)$. Since $\phi_s^- \wedge \phi_d^- \in C(\overline{\Omega})$, in view of Dini's theorem, we see that the convergence is uniform for all $x \in \overline{\Omega}$. \square

Proposition 6.5. *We have $\phi_{g_k}^- \rightarrow \phi_d^-$ uniformly on $\overline{\Omega}$ as $k \rightarrow \infty$.*

Proof. We only need to consider $c_{sc} \geq 0$ and prove that $\underline{g}_k \rightarrow g$ uniformly on $\partial\Omega$ as $k \rightarrow \infty$. Fix $x \in \partial\Omega$. Since $g_k(x, s) \geq g(x)$ for all $x \in \partial\Omega$ and $s \geq 0$, we have $\underline{g}_k(x) = \inf_{s \geq 0} \{g_k(x, s) + c_{sc}s\} \geq g(x)$. Thus we have $\liminf_{k \rightarrow \infty} \underline{g}_k(x) \geq g(x)$ for all $x \in \partial\Omega$.

Fix any $s \in (0, \infty)$. Then there exists $k_0 \in \mathbb{N}$ such that for all $k \geq k_0$, $g_k(x, s) = g(x)$. We get $\underline{g}_k(x) \leq g_k(x, s) + c_{sc}s = g(x) + c_{sc}s$. Therefore we have $\limsup_{k \rightarrow \infty} \underline{g}_k(x) \leq g(x) + c_{sc}s$. Sending $s \rightarrow 0$, we obtain $\limsup_{k \rightarrow \infty} \underline{g}_k(x) \leq g(x)$. Noting that $g \in C(\partial\Omega)$ and \underline{g}_k is nonincreasing as $k \rightarrow \infty$, in view of Dini's theorem we get a conclusion. \square

By the comparison principle for (IB) we have

$$u_k^1 \leq u \leq u_k^2 \text{ on } \overline{\Omega} \times [0, \infty) \text{ for all } k \in \mathbb{N}.$$

By Propositions 6.4, 6.5 we obtain

$$\limsup_{t \rightarrow \infty}^* u = \liminf_{t \rightarrow \infty}^* u = \lim_{k \rightarrow \infty} u_k^1 = \lim_{k \rightarrow \infty} u_k^2 \text{ on } \bar{\Omega}.$$

7. APPENDIX : EXISTENCE, UNIQUENESS AND REGULARITY RESULTS FOR (IB)

All the results presented in this appendix may appear at first glance as being well-known and covered by the standard results of the theory of viscosity solutions (see for instance [25, 21, 1, 7]). But some of them are not completely standard because we are dealing with an oblique vector field γ which is only continuous (and not Lipschitz continuous as in the standard cases) and with a domain Ω which is not very regular. Of course, we can extend the classical results to this more general framework because the solutions of (IB) are expected to be in $W^{1,\infty}$ since the Hamiltonian is coercive. But we have also to prove directly the existence of such $W^{1,\infty}$ -solutions, by using only comparison results which hold for $W^{1,\infty}$ sub and supersolutions, which is a little bit unusual in the theory of viscosity solutions.

Also, we remark that the calculations in the proof of comparison principles are used in the proofs of a key ingredient, Lemma 3.3, in order to prove the asymptotic monotone property, Theorem 1.3. Therefore, reminding the proofs of the comparison principles helps us to understand them.

Before providing these results and their proofs, we recall, for the reader's convenience, the definition of viscosity solutions for (IB), and in particular of boundary conditions in the viscosity sense (see [8] for instance).

Definition 1. *An upper-semicontinuous function u (resp., a lower semicontinuous function u) is a subsolution (resp., supersolution) of (IB) if the following conditions hold:*

- (i) u is a (viscosity) subsolution (resp., (viscosity) supersolution) of (1),
- (ii) $u(x, 0) \leq u_0(x)$ (resp., $u(x, 0) \geq u_0(x)$) for all $x \in \bar{\Omega}$, and
- (iii) for any $\phi \in C^1(\bar{\Omega} \times [0, \infty))$ and any $(x_0, t_0) \in \partial\Omega \times (0, \infty)$ such that $u - \phi$ takes a local maximum (resp., minimum) at (x_0, t_0) ,

$$\min\{\phi_t(x_0, t_0) + H(x_0, D\phi(x_0, t_0)), B(x_0, u(x_0, t_0), D\phi(x_0, t_0))\} \leq 0$$

(resp.,

$$\max\{\phi_t(x_0, t_0) + H(x_0, D\phi(x_0, t_0)), B(x_0, u(x_0, t_0), D\phi(x_0, t_0))\} \geq 0).$$

We call u a solution of (IB) if it is a subsolution and a supersolution of (IB).

7.1. Comparison Results for (IB).

Theorem 7.1. *Assume that (BA) holds.*

(i) **(Neumann/oblique derivative problem)** *Let $u \in C(\overline{\Omega} \times [0, \infty))$, $v \in \text{LSC}(\overline{\Omega} \times [0, \infty))$ be a subsolution and a supersolution of (CN), respectively. If $u(\cdot, 0) \leq v(\cdot, 0)$ on $\overline{\Omega}$, then $u \leq v$ on $\overline{\Omega} \times [0, \infty)$.*

(ii) **(State constraint or Dirichlet problem)** *Let $u \in C(\overline{\Omega} \times [0, \infty))$, $v \in \text{LSC}(\overline{\Omega} \times [0, \infty))$ be a subsolution and a supersolution of (SC) or (CD), respectively. If $u(\cdot, 0) \leq v(\cdot, 0)$ on $\overline{\Omega}$, then $u \leq v$ on $\overline{\Omega} \times [0, \infty)$.*

Remark 3. *It is worth mentioning that it is well known that there exists a discontinuous solution of (SC) and (CD), which implies that in the comparison principle for (SC) and (CD), we can not replace requirements of continuity of u by that of semicontinuity.*

Proof of Theorem 7.1 (i). We argue by contradiction assuming that there would exist $T > 0$ such that $\max_{\overline{Q_T}}(u - v)(x, t) > 0$, where $Q_T := \Omega \times (0, T)$.

Let u^δ denote the function

$$u^\delta(x, t) := \max_{s \in [0, T+2]} \{u(x, s) - (1/\delta)(t - s)^2\},$$

for any $\delta > 0$. This sup-convolution procedure is standard in the theory of viscosity solutions (although, here, it acts only on the time-variable) and it is known that, for δ small enough, u^δ is a subsolution of (CN) in $\Omega \times (a_\delta, T + 1)$, where $a_\delta := (2\delta \max_{Q_{T+2}} |u(x, t)|)^{1/2}$ (see [1, 8] for instance).

Moreover, it is easy to check that $|u_t^\delta| \leq C_\delta$ in $\Omega \times (a_\delta, T + 1)$ and therefore by the coercivity of H and the C^1 -regularity of $\partial\Omega$ we have, for all $x, y \in \overline{\Omega}$, $t, s \in [a_\delta, T + 1]$

$$|u^\delta(x, t) - u^\delta(y, s)| \leq C_\delta(|x - y| + |t - s|) \quad (52)$$

for some $C_\delta > 0$. Finally, as $\delta \rightarrow 0$, $\max_{\overline{Q_T}}(u^\delta - v)(x, t) \rightarrow \max_{\overline{Q_T}}(u - v)(x, t) > 0$.

Therefore it is enough to consider $\max_{\overline{Q_T}}(u^\delta - v)(x, t)$ for $\delta > 0$ small enough, and we follow the classical proof by introducing

$$\max_{\overline{Q_T}} \{(u^\delta - v)(x, t) - \eta t\},$$

for $0 < \eta \ll 1$. This maximum is achieved at $(\xi, \tau) \in \overline{\Omega} \times [0, T]$, namely $(u^\delta - v)(\xi, \tau) - \eta\tau = \max_{\overline{Q_T}} \{(u^\delta - v)(x, t) - \eta t\}$. Clearly τ depends on η but we can assume that it remains bounded away from 0, otherwise we easily get a contradiction.

We only consider the case where $\xi \in \partial\Omega$. We define the function $\Psi : \overline{\Omega}^2 \times [0, T] \rightarrow \mathbb{R}$ by

$$\begin{aligned} \Psi(x, y, t) := & u^\delta(x, t) - v(y, t) - \eta t - \frac{1}{\varepsilon^2} f(x - y) \\ & + g(\xi)(d(x) - d(y)) + \alpha(d(x) + d(y)) - |x - \xi|^2 - (t - \tau)^2, \end{aligned}$$

where f and d are given by (24) and (25). Let Ψ achieve its maximum at $(\bar{x}, \bar{y}, \bar{t}) \in \overline{\Omega}^2 \times [0, T]$. By a standard arguments, we have

$$\bar{x}, \bar{y} \rightarrow \xi \text{ and } \bar{t} \rightarrow \tau \text{ as } \varepsilon \rightarrow 0 \quad (53)$$

by taking a subsequence if necessary and, because of the Lipschitz continuity (52) of u^δ , we have

$$\frac{|\bar{x} - \bar{y}|}{\varepsilon^2} \leq C_\delta \quad (54)$$

Derivating (formally) Ψ with respect to each variable x, y at $(\bar{x}, \bar{y}, \bar{t})$, we have

$$\begin{aligned} D_x u^\delta(\bar{x}, \bar{t}) &= \frac{1}{\varepsilon^2} A(\bar{x} - \bar{y}) + 2(\bar{x} - \xi) - \alpha Dd(\bar{x}) - g(\xi) Dd(\bar{x}), \\ D_y v(\bar{y}, \bar{t}) &= \frac{1}{\varepsilon^2} A(\bar{x} - \bar{y}) + \alpha Dd(\bar{y}) - g(\xi) Dd(\bar{y}). \end{aligned}$$

We remark that we should interpret $D_x u^\delta$ and $D_y v$ in the viscosity solution sense here. We also point out that the viscosity inequalities we are going to write down below, hold up to time T , in the spirit of [1], Lemma 2.8, p. 41.

Since Ω is a domain with a C^1 -boundary, we first observe that for any $x \in \partial\Omega$ and $y \in \overline{\Omega}$

$$(x - y) \cdot n(x) \geq o(|x - y|),$$

where $o : [0, \infty) \rightarrow \mathbb{R}$ is a continuous function such that $|o(r)|/r \rightarrow 0$ as $r \rightarrow 0$.

Moreover, we have

$$\begin{aligned} & A(x - y) \cdot \gamma(x) \\ & \geq (x - y) \cdot A\gamma(\xi) - |y - x| m_\gamma(|x - \xi|) \\ & = (x - y) \cdot n(\xi) - |y - x| m_\gamma(|x - \xi|) \\ & \geq (x - y) \cdot n(x) - |y - x| (m_\gamma(|x - \xi|) + m_n(|x - \xi|)) \\ & \geq o(|x - y|) - |y - x| (m_\gamma(|x - y|) + m_n(|x - \xi|)), \end{aligned}$$

where m_γ and m_n are modulus of continuity of γ, n on $\overline{\Omega}$, respectively.

Setting $m_d := |o(r)|/r$, by (54), we obtain

$$\frac{1}{\varepsilon^2} A(\bar{x} - \bar{y}) \cdot \gamma(\bar{x}) \geq -C_\delta (m_d(C_\delta \varepsilon^2) + m_\gamma(C_\delta \varepsilon^2) + m_n(|\bar{x} - \xi|)). \quad (55)$$

Moreover, we have

$$\alpha Dd(\bar{x}) \cdot \gamma(\bar{x}) \leq -\alpha + \alpha m_\gamma(|\bar{x} - \xi|)$$

and

$$g(\xi) Dd(\bar{x}) \cdot \gamma(\bar{x}) \leq -g(\bar{x}) + m_g(|\bar{x} - \xi|),$$

where m_g is a modulus of continuity of g on $\partial\Omega$. Therefore we have

$$D_x u(\bar{x}, \bar{t}) \cdot \gamma(\bar{x}) \geq g(\bar{x}) - m(\varepsilon) + \alpha,$$

for a modulus m . Therefore, if $\varepsilon > 0$ is suitable small compared to $\alpha > 0$, we have

$$D_x u(\bar{x}, \bar{t}) \cdot \gamma(\bar{x}) > g(\bar{x}).$$

Similarly, if $\bar{y} \in \partial\Omega$, then we have

$$D_y u(\bar{y}, \bar{t}) \cdot \gamma(\bar{y}) < g(\bar{y})$$

for $\varepsilon > 0$ which is suitable small compared to $\alpha > 0$.

Therefore, by the definition of viscosity solutions of (CN), using the arguments of the User's guide to viscosity solutions [8], there exists $a_1, a_2 \in \mathbb{R}$ such that

$$\begin{aligned} a_1 + H(\bar{x}, \frac{1}{\varepsilon^2} A(\bar{x} - \bar{y}) + 2(\bar{x} - \xi) - \alpha Dd(\bar{x}) - g(\xi) Dd(\bar{x})) &\leq 0, \\ a_2 + H(\bar{y}, \frac{1}{\varepsilon^2} A(\bar{x} - \bar{y}) + \alpha Dd(\bar{x}) - g(\xi) Dd(\bar{y})) &\geq 0 \end{aligned}$$

with $a_1 - a_2 = \eta + 2(\bar{t} - \tau)$. By (54) we may assume that

$$\frac{1}{\varepsilon^2} A(\bar{x} - \bar{y}) \rightarrow p$$

as $\varepsilon \rightarrow 0$ for some $p \in \mathbb{R}^N$ by taking a subsequence if necessary. Sending $\varepsilon \rightarrow 0$ and then $\alpha \rightarrow 0$ in the above inequalities, we have a contradiction since $a - b \rightarrow \eta > 0$ while the H -terms converge to the same limit. Therefore τ cannot be assumed to remain bounded away from 0 and the conclusion follows. \square

Proof of Theorem 7.1 (ii). We only prove the comparison principle for (CD), since we can regard (SC) as a problem of (CD) with the extreme form “ $g(x) \equiv +\infty$ ”.

To justify this choice, we first prove the

Proposition 7.2. (Classical Dirichlet Boundary Conditions) *Assume that (BA) holds. If $u \in \text{USC}(\bar{\Omega} \times [0, \infty))$ is a subsolution of (CD), then $u(x, t) \leq g(x)$ for all $(x, t) \in \partial\Omega \times (0, \infty)$.*

In the same way we can prove that, if $u \in \text{USC}(\bar{\Omega})$ is a subsolution of (D)_a for any $a \geq c_{sc}$, then $u(x) \leq g(x)$ for all $x \in \partial\Omega$.

Proof. Fix $(x_0, t_0) \in \partial\Omega \times (0, \infty)$. Fix $r \in (0, t_0)$ and set $K := (\overline{B}(x_0, r) \cap \overline{\Omega}) \times [t_0 - r, t_0 + r]$. Choose a sequence $\{x_k\}_{k \in \mathbb{N}} \subset \mathbb{R}^N \setminus \Omega$ such that $|x_0 - x_k| = 1/k^2$. We define the function $\phi : K \rightarrow \mathbb{R}$ by

$$\phi(x, t) = u(x, t) - k|x - x_k| - \alpha_k(t - t_0)^2,$$

where $\{\alpha_k\}_{k \in \mathbb{N}} \subset (0, \infty)$ is a divergent sequence which will be fixed later. Let $(\xi_k, \tau_k) \in K$ be a maximum point of ϕ on K . Noting that $\phi(\xi_k, \tau_k) \geq \phi(x_0, t_0)$, we have

$$k|\xi_k - x_k| + \alpha_k(\tau_k - t_0)^2 \leq u(\xi_k, \tau_k) - u(x_0, t_0) + k|x_0 - x_k| \leq C, \quad (56)$$

where $C > 0$ is a constant independent of k . From the above, we see that $\xi_k \rightarrow x_0, \tau_k \rightarrow t_0$ as $k \rightarrow \infty$.

By the viscosity property of u , we have

$$q_k + H(\xi_k, p_k) \leq 0 \quad \text{or} \quad u(\xi_k, \tau_k) \leq g(\xi_k),$$

if $k \in \mathbb{N}$ is sufficiently large, where $p_k = k(\xi_k - x_k)/|\xi_k - x_k|$ and $q_k = 2\alpha_k(\tau_k - t_0)$. Set $f(k) := \min_{x \in B(x_0, r) \cap \overline{\Omega}} H(x, p_k)$. Noting that $|p_k| = k$, by the coercivity of H , $f(k) \rightarrow \infty$ as $k \rightarrow \infty$. Choose $\{\alpha_k\}_{k \in \mathbb{N}} \subset (0, \infty)$ such that

$\alpha_k \rightarrow \infty$ as $k \rightarrow \infty$ and $2\sqrt{C\alpha_k} + 1 \leq f(k)$ for sufficiently large $k \in \mathbb{N}$.

Since we have $\alpha_k|\tau_k - t_0| \leq \sqrt{C\alpha_k}$ for all $k \in \mathbb{N}$ by (56),

$$q_k + H(x_k, p_k) \geq -2\alpha_k|\tau_k - t_0| + f(k) \geq -2\sqrt{C\alpha_k} + f(k) \geq 1 > 0$$

for sufficiently large $k \in \mathbb{N}$, we must have $u(\xi_k, \tau_k) \leq g(\xi_k)$. Sending $k \rightarrow \infty$, we obtain $u(x_0, t_0) \leq g(x_0)$. \square

Now we return to the proof of the comparison result for (CD). We argue by contradiction, exactly in the same way as for (CN), introducing Q_T and the function u^δ which is a subsolution of (CD) in $\Omega \times (a_\delta, T+1)$. We still denote by $(\xi, \tau) \in \overline{\Omega} \times [0, T]$, a maximum point of $(u^\delta - v)(x, t) - \eta t$, namely

$$(u^\delta - v)(\xi, \tau) - \eta\tau = \max_{Q_T} \{(u^\delta - v)(x, t) - \eta t\}.$$

Again we may assume that τ remains bounded away from 0, otherwise the result would follow.

We only consider the case where $\xi \in \partial\Omega$. Since Ω is a domain with a C^0 -boundary, there exist $r > 0, b \in C(\mathbb{R}^{N-1}, \mathbb{R})$ such that, after re-labelling and re-orienting the coordinates axes if necessary, we may assume $B(\xi, r) \cap \Omega = B(\xi, r) \cap \{(x', x_N) \mid x_N > b(x')\}$. There exists $a_0 > 0$ such that for any $a \in (0, a_0)$, there exists a bounded open neighborhood W_a of $\overline{B}(\xi, r) \cap \overline{\Omega}$ such that $W_a \subset -ae_N + \Omega := \{-ae_N + x \mid x \in \Omega\}$, where $e_N := (0, \dots, 0, 1) \in \mathbb{R}^N$.

Since $(u^\delta - g)(\xi, \tau) \leq 0$ in view of Proposition 7.2 and $u^\delta(\xi, \tau) - v(\xi, \tau) > 0$, we have $v(\xi, \tau) - g(\xi) < 0$.

For any $a \in (0, a_0)$ and $(x, t) \in \overline{W}_a \times [0, T]$, set $u_a^\delta(x, t) := u^\delta(x + ae_N, t)$. Then, it is easily seen that u_a^δ satisfies

$$(u_a^\delta)_t(x, t) + H(x + ae_N, Du_a^\delta(x, t)) \leq 0 \quad \text{in } W_a \times (a_\delta, T + 1)$$

in the viscosity sense.

Let $\varepsilon > 0$ and define the function $\Psi : \overline{W}_a \times (\overline{B}(\xi, r) \cap \overline{\Omega}) \times [0, T] \rightarrow \mathbb{R}$ by

$$\begin{aligned} \Psi(x, y, t, s) &:= u_a^\delta(x, t) - v(y, t) - \eta t \\ &\quad - \frac{1}{2\varepsilon^2} |x - y|^2 - |y - \xi|^2 - (t - \tau)^2. \end{aligned}$$

Set $V := \overline{W}_a \times (\overline{B}(\xi, r) \cap \overline{\Omega}) \times [0, T]$ and let $(\bar{x}, \bar{y}, \bar{t}) \in V$ be a maximum point of Ψ on V . By the compactness of V we may assume that $\bar{x}, \bar{y} \rightarrow x_a$ and $\bar{t} \rightarrow t_a$ as $\varepsilon \rightarrow 0$ by taking a subsequence if necessary. Taking the limit as $\varepsilon \rightarrow 0$ in the inequality $\Psi(\bar{x}, \bar{y}, \bar{t}) \geq \Psi(\xi, \xi, \tau)$ yields

$$\begin{aligned} |x_a - \xi|^2 + |t_a - \tau|^2 &\leq (u_a^\delta - v)(x_a, t_a) - \eta t_a - \{(u_a^\delta - v)(\xi, \tau) - \eta \tau\} \\ &\leq (u^\delta - v)(x_a, t_a) - \eta t_a - \{(u^\delta - v)(\xi, \tau) - \eta \tau\} \\ &\quad + 2m_u(a) \\ &\leq 2m_u(a), \end{aligned}$$

where m_u be a modulus of continuity of u^δ on \overline{Q}_T . We notice that we use the continuity of u here.

Fix a small $a \in (0, a_0)$ so that $v(x_a, t_a) - g(x_a) < 0$ and moreover we may assume that $v(\bar{y}, \bar{t}) - g(\bar{y}) < 0$ for $\bar{y} \in \partial\Omega$ if ε is sufficiently small.

By the viscosity property of u_a^δ and v and using classical arguments, there exist $b_1, b_2 \in \mathbb{R}$ such that

$$b_1 + H(\bar{x} + ae_N, \frac{\bar{x} - \bar{y}}{\varepsilon^2}) \leq 0, \quad (57)$$

$$b_2 + H(\bar{y}, \frac{\bar{x} - \bar{y}}{\varepsilon^2} - 2(\bar{y} - \xi)) \geq 0 \quad (58)$$

with $b_1 - b_2 = \eta + 2(\bar{t} - \tau)$. In view of the Lipschitz continuity of u_a^δ on $\overline{W}_a \times [0, T]$, we have $(\bar{x} - \bar{y})/\varepsilon^2$ is bounded uniformly $\varepsilon > 0$ and therefore, by sending $\varepsilon \rightarrow 0$ we may assume $(\bar{x} - \bar{y})/\varepsilon^2 \rightarrow p_{\delta, a} \in B(0, C_{\delta, a})$ for some $C_{\delta, a} > 0$. Noting that H is uniformly continuous on $\overline{\Omega} \times B(0, C_{\delta, a} + 1)$, we see that there exists a modulus m_δ such that

$$H(x + ae_N, p) \geq H(x, p) - m_\delta(a) \quad \text{for any } (x, p) \in \overline{\Omega} \times B(0, C_\delta + 1). \quad (59)$$

Combining (57) and (58) with (59), if $a > 0$ is small enough, we get

$$m_\delta(a) \geq \eta + 2(\bar{t} - \tau) + H(\bar{x}, \frac{\bar{x} - \bar{y}}{\varepsilon^2}) - H(\bar{y}, \frac{\bar{x} - \bar{y}}{\varepsilon^2} - 2(\bar{y} - \xi)).$$

After taking the limit as $\varepsilon \rightarrow 0$, send $a \rightarrow 0$ and then we get $\eta \leq 0$, which is a contradiction. We have thus completed the proof. \square

7.2. Existence of Solutions of (IB).

Proof of Theorem 1.1. We present the proof of the existence of Lipschitz continuous solutions of (IB).

Following similar arguments as in the proof of Theorem 1.2, it is easy to prove that, for $C > 0$ large enough $-Ct + u_0(x)$ and $Ct + u_0(x)$ are, respectively, viscosity subsolution and a supersolution of (CN) or (SC) or (CD) with, of course, different constant C in each case.

By Perron's method (see [19]) and Theorem 7.1 we obtain continuous solutions of (IB) for (CN), (SC) and (CD) that we denote by u_1, u_2 and u_3 , respectively. As a consequence of Perron's method, we have

$$-Ct + u_0(x) \leq u_i(x, t) \leq Ct + u_0(x) \text{ on } \bar{\Omega} \times [0, \infty),$$

for $i = 1, 2, 3$.

To conclude, we use a standard argument : comparing the solutions $u_i(x, t)$ and $u_i(x, t + h)$ for some $h > 0$ and using the above property on the u_i , we have

$$\|u_i(\cdot, \cdot + h) - u_i(\cdot, \cdot)\|_\infty \leq \|u_i(\cdot, h) - u_i(\cdot, 0)\|_\infty \leq Ch.$$

As a consequence we have $\|(u_i)_t\|_\infty \leq C$ and, by using the equation together with (A1), we obtain that Du is also bounded. This completes the proof. \square

We finally remark that we can deduce the uniform continuity of solutions of (IB) under the assumption $u_0 \in C(\bar{\Omega})$. We may choose a sequence $\{u_0^k\}_{k \in \mathbb{N}} \subset W^{1,\infty}(\Omega) \cap C(\bar{\Omega})$ so that $\|u_0^k - u_0\|_{L^\infty(\Omega)} \leq 1/k$ for all $n \in \mathbb{N}$. Let u_k be a solution of (IB) with $u_0 = u_0^k$ and by the above argument we see $u_k \in UC(\bar{\Omega} \times [0, \infty))$ for all $k \in \mathbb{N}$. The maximum principle for (IB) implies that u_k uniformly converges to u on $\bar{\Omega} \times [0, \infty)$. Thus we obtain $u \in UC(\bar{\Omega} \times [0, \infty))$.

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Physique Théorique, Université de Tours and Mathematics Department, University of California, Berkeley. He is grateful for their hospitality.

REFERENCES

1. G. Barles, *Solutions de viscosité des équations de Hamilton-Jacobi*, Mathématiques & Applications (Berlin), 17, Springer-Verlag, Paris, 1994.
2. G. Barles, H. Ishii and H. Mitake, in preparation.
3. G. Barles and J.-M. Roquejoffre, *Ergodic type problems and large time behaviour of unbounded solutions of Hamilton-Jacobi equations*, Comm. Partial Differential Equations **31** (2006), no. 7-9, 1209–1225.
4. G. Barles and P. E. Souganidis, *On the large time behavior of solutions of Hamilton-Jacobi equations*, SIAM J. Math. Anal. **31** (2000), no. 4, 925–939.
5. E. N. Barron and R. Jensen, *Semicontinuous viscosity solutions for Hamilton-Jacobi equations with convex Hamiltonians*, Comm. Partial Differential Equations **15** (1990), no. 12, 1713–1742.
6. P. Bernard, J.-M. Roquejoffre, *Convergence to time-periodic solutions in time-periodic Hamilton-Jacobi equations on the circle*, Comm. Partial Differential Equations **29** (2004), no. 3-4, 457–469.
7. I. Capuzzo-Dolcetta and P.-L. Lions, *Hamilton-Jacobi equations with state constraints*, Trans. Amer. Math. Soc. **318** (1990), no. 2, 643–683.
8. M. G. Crandall, H. Ishii and P.-L. Lions, *User's guide to viscosity solutions of second order partial differential equations*, Bull. Amer. Math. Soc. (N.S.) **27** (1992), no. 1, 1–67.
9. A. Davini and A. Siconolfi, *A generalized dynamical approach to the large time behavior of solutions of Hamilton-Jacobi equations*, SIAM J. Math. Anal. **38** (2006), no. 2, 478–502.
10. A. Fathi, *Sur la convergence du semi-groupe de Lax-Oleinik*, C. R. Acad. Sci. Paris Sér. I Math. **327** (1998), no. 3, 267–270.
11. A. Fathi and A. Siconolfi, *PDE aspects of Aubry-Mather theory for quasiconvex Hamiltonians*, Calc. Var. Partial Differential Equations **22** (2005), no. 2, 185–228.
12. Y. Fujita and P. Loreti, *Long-time behavior of solutions to Hamilton-Jacobi equations with quadratic gradient term*, NoDEA Nonlinear Differential Equations Appl. **16** (2009), no. 6, 771–791.
13. Y. Fujita and K. Uchiyama, *Asymptotic solutions with slow convergence rate of Hamilton-Jacobi equations in Euclidean n space*, Differential Integral Equations **20** (2007), no. 10, 1185–1200.
14. Y. Giga, Q. Liu and H. Mitake, *Singular Neumann problems and large-time behavior of solutions of non-coercive Hamiltonian-Jacobi equations*, submitted.
15. Y. Giga, Q. Liu and H. Mitake, *Large-time behavior of one-dimensional Dirichlet problems of Hamilton-Jacobi equations with non-coercive Hamiltonians*, preprint.
16. N. Ichihara and H. Ishii, *Asymptotic solutions of Hamilton-Jacobi equations with semi-periodic Hamiltonians*, Comm. Partial Differential Equations **33** (2008), no. 4-6, 784–807.
17. N. Ichihara and H. Ishii, *The large-time behavior of solutions of Hamilton-Jacobi equations on the real line*, Methods Appl. Anal. **15** (2008), no. 2, 223–242.

18. N. Ichihara and H. Ishii, *Long-time behavior of solutions of Hamilton-Jacobi equations with convex and coercive Hamiltonians*, Arch. Ration. Mech. Anal. **194** (2009), no. 2, 383–419.
19. H. Ishii, *Perron's method for Hamilton-Jacobi equations*, Duke Math. J. **55** (1987), no. 2, 369–384.
20. H. Ishii, *Asymptotic solutions for large time of Hamilton-Jacobi equations in Euclidean n space*, Ann. Inst. H. Poincaré Anal. Non Linéaire, **25** (2008), no 2, 231–266.
21. H. Ishii, *Weak KAM aspects of convex Hamilton-Jacobi equations with Neumann type boundary conditions*, submitted. Available at <http://www.edu.waseda.ac.jp/~ishii/>.
22. H. Ishii, *Long-time asymptotic solutions of convex Hamilton-Jacobi equations with Neumann type boundary conditions*, to appear in Calc. Var. Partial Differential Equations.
23. H. Ishii and H. Mitake, *Representation formulas for solutions of Hamilton-Jacobi equations with convex Hamiltonians*, Indiana Univ. Math. J., **56** (2007), no. 5, 2159–2184.
24. S. N. Kružkov, *Generalized solutions of nonlinear equations of the first order with several independent variables. II*, (Russian) Mat. Sb. (N.S.) **72** (114) 1967 108–134.
25. P.-L. Lions, *Neumann type boundary conditions for Hamilton-Jacobi equations*, Duke Math. J. **52** (1985), no. 4, 793–820.
26. P.-L. Lions, A.-S. Sznitman, *Stochastic differential equations with reflecting boundary conditions*, Comm. Pure Appl. Math. **37** (1984), no. 4, 511–537.
27. P.-L. Lions, G. Papanicolaou, S. R. S. Varadhan, *Homogenization of Hamilton-Jacobi Equations*, unpublished work.
28. H. Mitake, *Asymptotic solutions of Hamilton-Jacobi equations with state constraints*, Appl. Math. Optim. **58** (2008), no. 3, 393–410.
29. H. Mitake, *The large-time behavior of solutions of the Cauchy-Dirichlet problem for Hamilton-Jacobi equations*, NoDEA Nonlinear Differential Equations App. **15** (2008), no. 3, 347–362.
30. H. Mitake, *Large time behavior of solutions of Hamilton-Jacobi equations with periodic boundary data*, Nonlinear Anal. **71** (2009), no. 11, 5392–5405.
31. G. Namah and J.-M. Roquejoffre, *Remarks on the long time behaviour of the solutions of Hamilton-Jacobi equations*, Comm. Partial Differential Equations **24** (1999), no. 5-6, 883–893.
32. J.-M. Roquejoffre, *Convergence to steady states or periodic solutions in a class of Hamilton-Jacobi equations*, J. Math. Pures Appl. (9) **80** (2001), no. 1, 85–104.

(G. Barles) LABORATOIRE DE MATHÉMATIQUES ET PHYSIQUE THÉORIQUE (UMR CNRS 6083), FÉDÉRATION DENIS POISSON, UNIVERSITÉ DE TOURS, PLACE DE GRANDMONT, 37200 TOURS, FRANCE

E-mail address: barles@lmpt.univ-tours.fr

URL: <http://www.lmpt.univ-tours.fr/~barles>

(H. Mitake) DEPARTMENT OF APPLIED MATHEMATICS, GRADUATE SCHOOL OF ENGINEERING HIROSHIMA UNIVERSITY HIGASHI-HIROSHIMA 739-8527, JAPAN

E-mail address: mitake@amath.hiroshima-u.ac.jp