



INDEX THEOREMS ON MANIFOLDS WITH STRAIGHT ENDS.

Werner Ballmann, Jochen Brüning, Gilles Carron

► To cite this version:

Werner Ballmann, Jochen Brüning, Gilles Carron. INDEX THEOREMS ON MANIFOLDS WITH STRAIGHT ENDS.. *Compositio Mathematica*, 2012, 148, pp.1897-1968. hal-00534166v2

HAL Id: hal-00534166

<https://hal.science/hal-00534166v2>

Submitted on 5 Apr 2012

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

INDEX THEOREMS ON MANIFOLDS WITH STRAIGHT ENDS.

WERNER BALLMANN, JOCHEN BRÜNING, AND GILLES CARRON

ABSTRACT. We study Fredholm properties and index formulas for Dirac operators over complete Riemannian manifolds with straight ends. An important class of examples of such manifolds are complete Riemannian manifolds with pinched negative sectional curvature and finite volume.

Date: April 5, 2012.

1991 Mathematics Subject Classification. 53C20, 58J20.

Key words and phrases. Dirac operator, Fredholm type, non-parabolic, index.

CONTENTS

1. Introduction	3
2. Preliminaries	12
2.1. Dirac Bundles	12
2.2. Decomposition of Spinors	15
2.3. Complex Hyperbolic Spaces	15
3. Dirac Systems and Distance Functions	17
3.1. Dirac Systems	17
3.2. Distance Functions	20
3.3. From Distance Functions to Dirac Systems	22
4. Boundary Values and Fredholm Properties	27
4.1. Boundary Values	28
4.2. More Function Spaces	30
4.3. Fredholm Properties of \mathcal{D}	32
5. Decomposition and Index	38
6. Manifolds with ε -Thin Ends	45
7. Cuspidal Ends	48
7.1. The Flat Connection	49
7.2. The Splitting	50
8. Explicit Index Formulas	53
9. Homogeneous Cusps	57
9.1. Asymptotic η -Invariants	62
9.2. Vanishing of η -Invariants	64
10. Heisenberg Manifolds	65
10.1. Lattices in Heisenberg groups	65
10.2. Twisted Right Regular Representation	66
10.3. Spectrum of Twisted Laplacians	68
10.4. η -Invariants for Heisenberg Manifolds	70
11. Low Energy η -Invariants	75
11.1. General Remarks and Computations	75
11.2. Contribution of Complex Hyperbolic Cusps	80
11.3. Résumé and Examples	86
References	89

1. INTRODUCTION

The celebrated Atiyah-Singer index theorem establishes a connection between analysis, geometry, and topology of closed manifolds. It contains the Chern-Gauss-Bonnet formula, Hirzebruch's signature theorem, and the Hirzebruch-Riemann-Roch formula as special cases. Later, Atiyah, Patodi, and Singer found a generalization of the index theorem for certain first order differential operators on compact manifolds with boundary [APS1]. In this article, they also discuss index theory for their class of operators on non-compact manifolds with cylindrical ends, and our work builds on that part of their work.

It is obvious that the structure of the underlying manifold and of the differential operator close to infinity plays an important role in this theory. Without restrictions on these data, not much can be expected.

Motivated by previous work of Barbasch-Moscovici [BaMo], Lott [Lo1, Lo2], and the first two authors [BB1, BB2], our main objective are Dirac operators on complete manifolds with pinched negative sectional curvature and finite volume. The structure of the ends of such manifolds has been determined by Patrick Eberlein and is related to the existence of so-called strictly invariant horospheres, see [Eb].

To set the stage, let M be a complete and connected Riemannian manifold of dimension m with Levi-Civita connection ∇ and curvature tensor R . Let $E \rightarrow M$ be a complex Dirac bundle¹ with Hermitian connection ∇^E , curvature tensor R^E , and Dirac operator D . For convenience, we assume throughout that R and R^E are uniformly bounded,

$$(1.1) \quad |R(X, Y)Z| \leq C_R |X| |Y| |Z|, \quad |R^E(X, Y)\sigma| \leq C_R^E |X| |Y| |\sigma|,$$

for all vector fields X, Y, Z on M and sections σ of E . The bound on R is equivalent to assuming a uniform bound on the modulus of the sectional curvature K_M of M .

Recall that D is an elliptic differential operator of first order. Consider D as an unbounded operator on $L^2(M, E)$ with domain $C_c^\infty(M, E)$, where $L^2(M, E)$ denotes the space of square-integrable sections of E and $C_c^\infty(M, E)$ the space of smooth sections of E with compact support, and note that D is symmetric on the latter. The closure of D has domain $H^1(M, E)$, by (1.1) and the general Bochner identity, see (2.13) and (2.14). Furthermore, $D : H^1(M, E) \rightarrow L^2(M, E)$ is self-adjoint, see [Wo] or Theorem II.5.7 in [LaMi].

We may ask under which conditions $D : H^1(M, E) \rightarrow L^2(M, E)$ is a Fredholm operator. By self-adjointness, this is the case if and only if 0 is not in the essential spectrum of D ; according to a result of Nicolae

¹in the sense of Gromov and Lawson, compare Section 2.1

Anghel, this holds if and only if there is a compact subset L in M and a constant $C = C(L)$ such that

$$(1.2) \quad \|\sigma\|_{L^2(M,E)} \leq C \|D\sigma\|_{L^2(M,E)},$$

for all smooth sections σ of E with compact support in $M \setminus L$, see [An]. If such an estimate holds, we say that D is of *Fredholm type*.

Better adapted to our investigations and more flexible is a somewhat weaker notion, introduced by the third named author in [Ca1]:

DEFINITION 1.3. We say that D is *non-parabolic at infinity* if there is a compact subset L in M such that, for any relatively compact open subset K of M , there is a constant $C = C(K, L)$ such that

$$(1.4) \quad \|\sigma\|_{L^2(K,E)} \leq C \|D\sigma\|_{L^2(M,E)},$$

for all smooth sections σ of E with compact support in $M \setminus L$.

It follows from [Ca1, Théorème 1.2] that D is non-parabolic at infinity if and only if there is a Hilbert space W of sections of E which are locally H^1 , such that $H^1(M, E)$ is a dense subspace of W , such that the inclusions

$$(1.5) \quad H^1(M, E) \subseteq W \subseteq H_{\text{loc}}^1(M, E)$$

are continuous, and such that D extends to a Fredholm operator

$$(1.6) \quad D_{\text{ext}} : W \rightarrow L^2(M, E).$$

It then follows that $\text{im } D_{\text{ext}}$ is equal to the closure of $\text{im } D$ in $L^2(M, E)$ and that $H^1(M, E) = W$ if and only if D is of Fredholm type².

If D is non-parabolic at infinity, with associated Hilbert space W , then elements of $\ker D_{\text{ext}}$ will be called *extended solutions* of D . In the case of cylindrical ends, they correpond exactly to the extended solutions in [APS1]. By the density of $C_c^\infty(M, E)$ in W , the orthogonal complement of the image of D_{ext} in $L^2(M, E)$ is equal to the space of L^2 -solutions of D . Since D_{ext} is a Fredholm operator, the spaces of extended solutions and L^2 -solutions of D are of finite dimension, and their difference, $\text{ind } D_{\text{ext}}$, is called the *extended index* of D . As a consequence of one of our main results concerning non-parabolicity, Theorem 1.14 below, we obtain the following assertion:

THEOREM 1.7. *If the sectional curvature of M is negatively pinched and the volume of M is finite, then D is non-parabolic at infinity. In particular, the space of L^2 -solutions of D is finite-dimensional.*

²Compare Section 5.

Under a more general assumption on the geometry of the ends of M , similar to Condition (1) in Theorem 1.13 below, John Lott showed that the space of square-integrable harmonic differential forms is finite dimensional, see Theorem 1 in [Lo1].

For manifolds with ends as in the case of finite volume manifolds of pinched negative sectional curvature, Lott also discusses the essential spectrum of $(d + d^*)^2$ on the space of differential forms, see Theorem 2 in [Lo1]. Under the same assumption on the geometry of the ends and for Dirac bundles as in Condition (2) of Theorem 1.13 below, he investigates the essential spectrum of the associated Dirac operator, see Theorem 5 in [Lo2]. Similar results have been obtained in [BB2]. In this article, we do not deal with the essential spectrum, but would like to mention that our investigations lead to extensions of these results.

It is clear from the definition of non-parabolicity that it only depends on the structure of D at infinity³. To state our results in that context, we need to introduce a further notion.

DEFINITION 1.8. We say that the ends of M are *straight* if M can be decomposed into a compact part M_0 and an unbounded part U_0 with common boundary N such that there is an open set $U \supseteq U_0$ and a C^2 distance function⁴ $f : U \rightarrow \mathbb{R}$ whose gradient flow establishes a C^1 diffeomorphism

$$(1.9) \quad F : (-r, \infty) \times N \rightarrow U,$$

where $r > 0$, $U_0 = f^{-1}([0, \infty))$, $N = f^{-1}(0)$, and $f(F(t, x)) = t$. If f is smooth, then we say that the ends of M are *smooth*.

If the ends of M are straight, then M is diffeomorphic to the interior of the compact manifold M_0 , and the connected components of N correspond to the different ends of M . Furthermore, the induced Riemannian metric on $\mathbb{R}_+ \times N$ is of the form

$$(1.10) \quad dt^2 + g_t,$$

where $(g_t)_{t \geq 0}$ is a family of Riemannian metrics on N . The regularity of this family is a technical problem which we address in Section 3.2 and which motivated our previous work [BBC2] on Dirac systems with Lipschitz coefficients. *Cylindrical ends* as mentioned above correspond to the case of Riemannian products, that is, f is smooth and $g_t = g_0$, for all $t \in (-r, \infty)$.

If the ends of M are straight, we fix the setup as in Definition 1.8, identify $(-r, \infty) \times N \simeq U$ via F , and call the hypersurfaces $N_t =$

³The same applies to the essential spectrum of D .

⁴Compare Section 3.2.

$f^{-1}(t)$, endowed with the Riemannian metric g_t , the *cross sections* of U . For convenience, we will always assume in this situation that the Weingarten operators $W = W_t$ of the cross sections are uniformly bounded,

$$(1.11) \quad |WX| \leq C_W |X|,$$

for all vector fields X on U .

DEFINITION 1.12. Let $\varepsilon > 0$. We say that the ends of M are ε -*thin* if they are straight and the connected components of the cross sections N_t have diameter at most ε , for all sufficiently large t . We say that the ends of M are *cuspidal* if they are straight and there are positive constants c and C such that the metrics g_t as in (1.10) satisfy $g_t \leq Ce^{c(s-t)}g_s$, for all sufficiently large $s < t$.

For example, if M has finite volume and pinched negative sectional curvature, say $-b^2 \leq K_M \leq -a^2 < 0$, then the ends of M are cuspidal with $c = 2a$ and $C = 1$. We note that, in this example, the distance function arises from Busemann functions on the universal covering space of M and that such Busemann functions are C^2 , see [HeIH] or Proposition IV.3.2 in [Ba]. Better regularity is, in general, not expected and, at least for non-positively curved manifolds, better regularity does not hold, see [BBB].

THEOREM 1.13. *There is a positive constant $\varepsilon = \varepsilon(m, C_R, C_W)$ such that D is non-parabolic at infinity if the following two conditions hold:*

- (1) *All ends of M are ε -thin, for all sufficiently large t .*
- (2) *E is a Hermitian vector bundle associated to M via a unitary representation of $O(m)$, $SO(m)$ (if M is oriented), or $Spin(m)$ (with respect to a spin structure of M), respectively.*

Extending Theorem 1.7 above, we also have:

THEOREM 1.14. *If the ends of M are cuspidal, then D is non-parabolic at infinity.*

Suppose now that the dimension m of M is even and that $E = E^+ \oplus E^-$ is a super-symmetry⁵. Then D maps sections of E^+ to sections of E^- and conversely, and thus we obtain operators

$$(1.15) \quad D^\pm : H^1(M, E^\pm) \rightarrow L^2(M, E^\mp).$$

In the case of closed manifolds, these are the operators the Atiyah-Singer index theorem is concerned with. The local index theorem associates an index form ω_{D^+} to D^+ , determined by local data of D^+ ,

⁵See Section 2.1.

whose integral is equal to the index of D^+ . If D is non-parabolic at infinity, we obtain corresponding Fredholm operators

$$(1.16) \quad D_{\text{ext}}^\pm : W^\pm \rightarrow L^2(M, E^\pm),$$

where W^\pm consists of those sections in W which take values in E^\pm . One of our general results on the extended index of D^+ is

THEOREM 1.17. *If M has at most finitely many ends, D is non-parabolic at infinity, and ω_{D^+} is integrable, then*

$$\text{ind } D_{\text{ext}}^+ = \int_M \omega_{D^+} + \sum_{\mathcal{C}} \text{Corr}(\mathcal{C}),$$

where ω_{D^+} is the index form associated to D^+ , \mathcal{C} runs over the ends of M , and $\text{Corr}(\mathcal{C})$ is a correction term determined by the end \mathcal{C} .

Theorem 1.17 is a kind of relative index theorem and, assuming the non-parabolicity of D , can also be proved along the lines of relative index formulas as in Theorem 4.18 in [GrLa] (see also Proposition 4.33), Theorem 6.2 in [Do], or Theorem 0.5 in [Ca1].

Clearly, the assumptions of Theorem 1.17 are satisfied if the ends of M are cuspidal. We assume the latter in the following discussion.

In dimension $m = 2$, the correction terms are known explicitly in terms of the type of E along the ends, see [BB1]. In higher dimensions and under strong pinching assumptions on the sectional curvature of M , they are known explicitly for the Gauss-Bonnet operator, see [BB2].

In the case where the ends of M are homogeneous cusps and, over them, E is a twisted associated bundle as in Chapter 9, we obtain explicit formulas for the extended index of D^+ . Over each cusp \mathcal{C} , D^+ can then be decomposed into two orthogonal parts, a low energy part $D_{\mathcal{C}}^{\text{le},+}$ and a high energy⁶ part $D_{\mathcal{C}}^{\text{he},+}$. Moreover, these are naturally equivalent to operators of the form

$$(1.18) \quad \frac{d}{dt} + A_{\mathcal{C}}^{\text{le},+} \quad \text{and} \quad \frac{d}{dt} + A_{\mathcal{C},t}^{\text{he},+},$$

respectively, where $A_{\mathcal{C}}^{\text{le},+}$ is a self-adjoint operator on some finite dimensional Hilbert space and where $A_{\mathcal{C},t}^{\text{he},+}$, $t \geq 0$, is a family of invertible self-adjoint operators on some (other) Hilbert space, with eigenvalues of finite multiplicity and a spectral gap that tends to infinity as $t \rightarrow \infty$.

⁶Our usage of the notion *high energy* follows the terminology introduced in [Lo1].

THEOREM 1.19. *If the ends \mathcal{C} of M are homogeneous cusps and, over them, E is a twisted associated bundle as in Chapter 9, then*

$$\operatorname{ind} D_{\text{ext}}^+ = \int_M \omega_{D^+} + \frac{1}{2} \sum_{\mathcal{C}} \left(\lim_{t \rightarrow \infty} \eta(A_{\mathcal{C},t}^{\text{he},+}) + \eta(A_{\mathcal{C}}^{\text{le},+}) + \dim \ker A_{\mathcal{C}}^{\text{le},+} \right).$$

Moreover, D^+ is of Fredholm type if and only if $\dim \ker A_{\mathcal{C}}^{\text{le},+} = 0$, for all ends \mathcal{C} of M .

Denote by h_{∞}^{\pm} the difference in dimension of the spaces of extended and L^2 -solutions of D^{\pm} . These quantities determine the difference between the extended and L^2 -indices of D^{\pm} ,

$$(1.20) \quad \operatorname{ind} D_{\text{ext}}^{\pm} = \operatorname{ind}_{L^2} D^{\pm} + h_{\infty}^{\pm},$$

where $\operatorname{ind}_{L^2} D^{\pm} := \dim \ker D^{\pm} - \dim \ker D^{\mp}$.

THEOREM 1.21. *If the ends \mathcal{C} of M are homogeneous cusps and, over them, E is a twisted associated bundle as in Chapter 9, then*

$$\operatorname{ind}_{L^2} D^+ = \int_M \omega_{D^+} + \frac{1}{2} \sum_{\mathcal{C}} \left(\lim_{t \rightarrow \infty} \eta(A_{\mathcal{C},t}^{\text{he},+}) + \eta(A_{\mathcal{C}}^{\text{le},+}) \right) - \frac{1}{2} (h_{\infty}^+ - h_{\infty}^-).$$

The term $h_{\infty}^+ - h_{\infty}^-$ is non-local, and therefore the L^2 -index does not admit an analog of Theorem 1.17, where the correction terms are determined by the different ends separately.

The most important class of examples to which Theorems 1.19 and 1.21 apply are finite volume quotients of symmetric spaces of negative sectional curvature, that is, of real, complex, or quaternionic hyperbolic spaces or of the Cayley hyperbolic plane. The work of Barbasch-Moscovici [BaMo] is a milestone in the index theory of Dirac operators of homogeneous Dirac bundles over such spaces. Their arguments rely on harmonic analysis on symmetric spaces, notably the Selberg trace formula.

Werner Müller has obtained an index formula for manifolds with cusps of \mathbb{Q} -rank one. His approach follows the one of Atiyah-Patodi-Singer: using a good approximation of the heat operator e^{-tD^2} and harmonic analysis, he computed the correction terms (see theorems 10.32 and 11.77 in [Mü2]). In particular, he determined the non-local term $h_{\infty}^+ - h_{\infty}^-$ using the scattering matrix at zero energy, see also [Mü1].

Müller's approach has been extended by Mark Stern [St1]. The derivations of index formulas by Boris Vaillant in [Va] and Leichtnam, Mazzeo, and Piazza in [LMP] are based on the pseudodifferential calculus of Mazzeo-Melrose [MaMe]. However, the fibered boundary geometry discussed in [Va] and [LMP] does not cover all finite volume locally symmetric space with \mathbb{Q} -rank one. In their recent paper [GrHu],

Daniel Grieser and Eugenie Hunsicker settle the main properties of the pseudodifferential calculus for multi-fibred boundary structures.

Our proof is different in nature and is based on our previous work: applying our results from [BBC2], we are able to discuss the contribution of each end individually. This leads to a more general setting and more transparent index formulas. Note, in particular, that our results also apply in the case where D is not of Fredholm type.

In this article, we concentrate on complex hyperbolic cusps, more precisely, cusps as they arise for quotients of complex hyperbolic space $\mathbb{C}H^n$ of dimension $m = 2n$. Here we discuss only two specific results for such quotients and refer the reader to our résumé in Section 11.3 for our general index formulas pertaining to complex hyperbolic cusps. The first result concerns the L^2 -arithmetic genus and follows from the Hirzebruch proportionality principle and Theorems 1.19 and 11.79; see also Example 1 in Section 11.3.

THEOREM 1.22. *If X is a quotient of complex hyperbolic space $\mathbb{C}H^n$ by a neat lattice⁷, then the Dolbeault operator on X is of Fredholm type and its index $\chi_{L^2}(X, \mathcal{O})$ is given by*

$$\chi_{L^2}(X, \mathcal{O}) = (-1)^n \frac{\text{vol } X}{\text{vol } \mathbb{C}P^n} + \begin{cases} 0 & \text{if } n \text{ is odd,} \\ \zeta(1-n) \sum_{\mathcal{C}} |\Gamma_{\mathcal{C}}| & \text{if } n \text{ is even.} \end{cases}$$

Here $|\Gamma_{\mathcal{C}}|$ denotes a basic invariant of the fundamental group $\Gamma_{\mathcal{C}}$ of the cusp \mathcal{C} , for each cusp \mathcal{C} of X ; compare (10.4).

The second result concerns the signature operator on X when n is even, that is, when m is a multiple of 4. It is also a consequence of the Hirzebruch proportionality principle and Theorems 1.19 and 11.79; see also Example 2 in Section 11.3.

THEOREM 1.23. *If X is a quotient of complex hyperbolic space $\mathbb{C}H^n$ by a neat lattice, where n is even, then the signature operator on X is of Fredholm type and its index $\sigma_{L^2}(X)$ is given by*

$$\sigma_{L^2}(X) = \frac{\text{vol } X}{\text{vol } \mathbb{C}P^n} + 2^n \zeta(1-n) \sum_{\mathcal{C}} |\Gamma_{\mathcal{C}}| + \nu(-1)^{n/2} \left(\binom{n-2}{n/2} - \binom{n-2}{n/2-1} \right),$$

where ν is equal to the number of ends of X .

For complex hyperbolic manifolds with finite volume, the space of L^2 -harmonic forms has a topological interpretation. In fact, in this case,

⁷Here neat means that the group generated by the eigenvalues of any non-identity element of the given lattice contains no roots of unity. Neat lattices are torsion-free.

the L^2 -signature coincides with the topological signature [Zu]. Formulas for $\sigma_{L^2}(X)$ are also stated in Theorem 7.6 of [BaMo] and Stern's article [St2] (compare Formula 6.4 there). Our correction terms consist of two terms: What we call the *high energy η -invariant* can be identified with a zeta contribution in [St2] and with the unipotent contribution in the Arthur-Selberg trace formula in [BaMo]. Our *low energy η -invariant* corresponds to the eta term in [St2] and the weighted unipotent contribution in [BaMo]. Since our correction terms are obtained by different methods, we obtain, in particular, different interpretations of the corresponding terms in [BaMo] and [St2].

The formulas in Theorems 1.22 and 1.23 show that the volume of the quotient X of $\mathbb{C}H^n$ in question is a rational multiple of the volume of $\mathbb{C}P^n$. This was already known by Harder's Gauss-Bonnet theorem which says that $(n+1) \operatorname{vol}(X) / \operatorname{vol} \mathbb{C}P^n = (-1)^n \chi(X)$, where $\chi(X) \in \mathbb{Z}$ denotes the Euler characteristic of X , see [Ha]. Theorem 1.22 implies that $\operatorname{vol}(X) / \operatorname{vol} \mathbb{C}P^n$ is integral for odd n . The question of the integrality of $\operatorname{vol}(X) / \operatorname{vol} \mathbb{C}P^n$ has been brought to our attention by Martin Olbrich: The half-integrality of $\operatorname{vol}(X) / \operatorname{vol} \mathbb{C}P^n$ implies that certain Selberg type zeta functions are meromorphic.

As another example where our methods apply, we mention the Dirac operator D on the spinor bundle, supposing that M admits a spin structure. The case $n = 1$, that is, of surfaces of finite area with cusps of constant negative curvature, has been dealt with in [Bä], see also [BB1]. In particular, D is of Fredholm type if and only if the spin structure is not periodic along (the cross sections of) any of the cusps, see Theorem 2 in [Bä] or Theorem 0.1 in [BB1]. We will discuss the correction terms for spinor bundles along complex hyperbolic cusps in more detail in Examples 10.65 and Example 11.81.3.

Our formulas for complex hyperbolic cusps apply to more examples, but we refer the reader to Theorem 11.79 for the full scope of our results.

In Chapter 2 we discuss some notions and results which are basic for our later investigations. Chapter 3 is devoted to distance functions and their relation to Dirac systems. In particular, Section 3.2 contains a detailed study of C^2 distance functions as we need it in our application to Busemann functions. In this section, we clarify and correct some of the statements from [BB2]. Some essential parts of our later analysis depend on our previous results in [BBC2]⁸. That

⁸In some cases, the work of Marius Mitrea could also be used: In Section 5 of [Mi], Mitrea investigates the regularity of the Calderón projector for Dirac operators on Lipschitz domains with $C^{1,1}$ symbol and metric tensor, using paradifferential calculus.

the applications of these results are justified is the topic of Section 3.3. In Chapter 4, we discuss boundary value problems and Fredholm properties of Dirac systems which are underlying Dirac operators over straight ends. Proposition 4.46 is one of the corner stones of our later discussion. Chapter 5 contains the first applications to index formulas and a proof of Theorem 1.17. Chapter 6 and Chapter 7 contain the proofs of Theorems 1.13 and 1.14, respectively. In Chapter 8, we prove Theorems 1.19 and 1.21. The last three chapters are devoted to a discussion of the index contributions of homogeneous cusps. Ideas from the work of Deninger-Singhof [DeSi] are basic in our computation of high energy η -invariants of Dirac operators on compact quotients of Heisenberg groups. Following the discussion of Gordon-Wilson in [GoWi], we compute in Section 10.3 the spectrum of twisted Laplacians on compact quotients of Heisenberg groups. This is needed in our computation of high energy η -invariants in Section 10.4. In Chapter 11, we discuss the low energy η -invariants of Dirac bundles over complex hyperbolic cusps. One of the main ingredients in this latter discussion is a theorem of Kostant concerning Lie algebra cohomology (Theorem 4.139 in [KnVo]).

We would like to thank Patrick Ghanaat, Jean Louis Milhorat, Henri Moscovici, and Martin Olbrich for helpful discussions. We are also very grateful to the referee whose many detailed comments and suggestions helped to improve the exposition of the article considerably. We enjoyed the hospitality of the ESI in Vienna (W.B.), the University of Nantes (W.B. and J.B.), the MPIM in Bonn (J.B. and G.C.), the University of Kyoto (J.B.), and the MSRI in Berkeley (G.C.). W.B. would like to thank the MPIM and HCM in Bonn for their continuous support. J.B. appreciates the support by SFB 647. G.C. acknowledges the support by the grant GeomEinstein 06-BLAN-0154 gratefully.

2. PRELIMINARIES

Let M be a Riemannian manifold of dimension m with Levi-Civita connection ∇ and curvature tensor R . Let $E \rightarrow M$ be a Hermitian vector bundle over M , endowed with a Hermitian connection ∇^E and associated curvature R^E . Recall that we assume that the norms of R and R^E are uniformly bounded, compare (1.1).

We denote by $C^\infty(M, E)$ and $L^2(M, E)$ the spaces of smooth and square-integrable sections of E , respectively. We let $H^1(M, E)$ be the closure of $C^\infty(M, E)$ with respect to the H^1 -norm, that is, the norm associated to the inner product

$$(2.1) \quad (\sigma, \tau)_{H^1(M, E)} := (\sigma, \tau)_{L^2(M, E)} + (\nabla^E \sigma, \nabla^E \tau)_{L^2(M, E \otimes T^*M)}.$$

We denote by $C_c^\infty(M, E)$, $L_c^2(M, E)$, and $H_c^1(M, E)$ the subspaces of corresponding sections with compact support and by $L_{\text{loc}}^2(M, E)$ and $H_{\text{loc}}^1(M, E)$ the spaces of measurable sections σ of E such that $\varphi\sigma$ belongs to $L^2(M, E)$ and $H^1(M, E)$, respectively, for any smooth function φ on M with compact support. In the case where the boundary of M is non-empty, we use a double index cc to indicate compact support in the interior of M and an index 0 to indicate vanishing along the boundary.

For better readability, we have arranged the rest of the preliminaries into sections. In Section 2.1 we introduce Dirac bundles and operators, in Section 2.2 we collect some generalities about spinors, and in Section 2.3 we introduce complex hyperbolic spaces.

2.1. Dirac Bundles. We say that E is a *Dirac bundle* over M if E is endowed with a compatible *Clifford multiplication*, that is, a field

$$(2.2) \quad TM \times E \rightarrow E, \quad (x, v) \mapsto x \cdot v,$$

of bilinear maps such that

$$(2.3) \quad XX\sigma = -|X|^2\sigma,$$

$$(2.4) \quad |X\sigma| = |X||\sigma|,$$

$$(2.5) \quad \nabla_X^E(Y\sigma) = (\nabla_X Y)\sigma + Y(\nabla_X^E \sigma),$$

for all vector fields X, Y on M and sections σ of E , where we use $X\sigma$ as a shorthand for $X \cdot \sigma$.

Suppose now that E is a Dirac bundle over M . Then the *Dirac operator* D associated to E is given by

$$(2.6) \quad D\sigma = \sum_{1 \leq i \leq m} X_i \nabla_{X_i}^E \sigma,$$

where (X_1, \dots, X_m) is a local orthonormal frame of M and σ is a section of E . For any function φ on M and section σ of E ,

$$(2.7) \quad D(\varphi\sigma) = \text{grad } \varphi \cdot \sigma + \varphi D\sigma.$$

In particular, the principal symbol of D at $\xi \in T^*M$ is given by Clifford multiplication with the dual vector $\xi^\# \in TM$, and hence D is elliptic. Note also that D is formally self-adjoint, that is, D is symmetric on $C_{cc}^\infty(M, E)$.

Suppose now that M has boundary, $N := \partial M$, let T be the inward normal field along N , and set $W := \nabla T$, the *Weingarten operator* of N with respect to T . We assume that the operator norm of W is uniformly bounded by a constant C_W . Change Clifford multiplication and connection of E along N by

$$(2.8) \quad X * \sigma := TX\sigma,$$

$$(2.9) \quad \nabla_X^T \sigma := \nabla_X^E \sigma - \frac{1}{2}(WX) * \sigma = \nabla_X^E \sigma - \frac{1}{2}(T\nabla_X T)\sigma.$$

It is well known that, with these new data, the restriction of E to N is again a Dirac bundle such that Clifford multiplication by T is ∇^T -parallel, see for example Section 3.10.1 in [Gi2]. The associated Dirac operator is given by

$$(2.10) \quad D^T \sigma = \sum_{2 \leq i \leq m} X_i * \nabla_{X_i}^T \sigma = \sum_{2 \leq i \leq m} TX_i \nabla_{X_i}^E \sigma + \frac{\kappa}{2} \sigma,$$

where (X_1, X_2, \dots, X_m) is a local orthonormal frame of M along N with $X_1 = T$, and

$$(2.11) \quad \kappa = \text{tr } W$$

is the *mean curvature* of N with respect to T . The curvature of ∇^T is

$$(2.12) \quad R^T(X, Y)\sigma = R^E(X, Y)\sigma - \frac{1}{2}(R(X, Y)T) * \sigma - \frac{1}{4}[WX, WY]\sigma.$$

The general Bochner identity [LaMi, Theorem II.8.2] implies that

$$(2.13) \quad (\nabla^E \sigma_1, \nabla^E \sigma_2)_{L^2(M, E \otimes T^*M)} + (K^E \sigma_1, \sigma_2)_{L^2(M, E)} \\ = (D\sigma_1, D\sigma_2)_{L^2(M, E)} + (D^T \sigma_1 - \frac{\kappa}{2} \sigma_1, \sigma_2)_{L^2(N, E)},$$

for all $\sigma_1, \sigma_2 \in C_c^\infty(M, E)$, where K^E is a curvature term,

$$(2.14) \quad K^E \sigma = \sum_{1 \leq i < j \leq m} X_i X_j R^E(X_i, X_j)\sigma.$$

We see that the operator norm of K^E is bounded by $m(m-1)C_R^E/2$ and conclude that on sections with compact support in the interior of M , the graph norm of D is equivalent to the H^1 -norm.

Since N has no boundary, (2.13) applied to N turns into

$$(2.15) \quad (\nabla^T \sigma_1, \nabla^T \sigma_2)_{L^2(N, E \otimes T^*N)} + (K^T \sigma_1, \sigma_2)_{L^2(N, E)} \\ = (D^T \sigma_1, D^T \sigma_2)_{L^2(N, E)},$$

where K^T denotes the curvature term built from R^T as K^E is built from R^E in (2.14). We see that the operator norm of K^T is bounded here by

$$(2.16) \quad \frac{(m-1)(m-2)}{2} (C_R^E + \frac{1}{2}C_R + \frac{1}{4}C_W^2),$$

where C_W is a uniform bound for the operator norm of W (compare (1.11)), and conclude now that, along N , the graph norm of D^T is equivalent to the H^1 -norm.

Let E be a Dirac bundle over M . A *super-symmetry* of E is an orthogonal decomposition $E = E^+ \oplus E^-$, where E^\pm are parallel Hermitian subbundles of E such that $XE^+ \subseteq E^-$ and $XE^- \subseteq E^+$, for all vector fields X of M . In particular, E^+ and E^- are of the same dimension. If $E = E^+ \oplus E^-$ is a super-symmetry, then the Dirac operator D of E maps sections of E^+ into sections of E^- and conversely and therefore can be written as

$$(2.17) \quad D = \begin{pmatrix} 0 & D^- \\ D^+ & 0 \end{pmatrix}$$

with respect to the super-symmetry. We can also think of a super-symmetry as a parallel field of unitary involutions of E which anti-commute with Clifford multiplication, where E^\pm is the subbundle of eigenspaces of the involutions for the eigenvalue ± 1 , respectively.

If M is oriented and $m = \dim M$ is even, then the *complex volume form* of M is defined to be

$$(2.18) \quad \omega_{\mathbb{C}} := i^{m/2} X_1 \cdots X_m \in \mathbb{C}l(M),$$

where (X_1, \dots, X_m) is an oriented local orthonormal frame of M . For any Dirac bundle E over M , multiplication by $\omega_{\mathbb{C}}$ is a parallel field of unitary involutions of E which anti-commutes with Clifford multiplication with vector fields, and hence it defines a super-symmetry $E = E^+ \oplus E^-$.

Suppose now M is complete and that the boundary of M is empty, and consider D as an unbounded operator in $L^2(M, E)$ with domain $C_c^\infty(M, E)$. Since D is symmetric on $C_c^\infty(M, E)$, it is closable in

$L^2(M, E)$. Since the graph norm of D is equivalent to the H^1 -norm, $H^1(M, E)$ is the domain of the closure of D . By [Wo] or Theorem II.5.7 in [LaMi], D on $H^1(M, E)$ is self-adjoint in $L^2(M, E)$.

2.2. Decomposition of Spinors. Let m be even, $m = 2n$, and consider the complex Clifford algebra $\mathbb{Cl}(2n) = \mathbb{Cl}(\mathbb{R}^{2n})$, where we denote the complex structure on $\mathbb{Cl}(2n)$ by $\sqrt{-1}$. Fix an orthonormal basis (e_1, \dots, e_{2n}) of \mathbb{R}^{2n} and set⁹

$$(2.19) \quad \omega_j := \sqrt{-1}e_{2j-1}e_{2j} \in \mathbb{Cl}(2n), \quad 1 \leq j \leq n.$$

Then

$$(2.20) \quad \omega_j^2 = 1 \quad \text{and} \quad \omega_j \omega_k = \omega_k \omega_j,$$

for all $1 \leq j, k \leq n$, and the complex volume form is given by

$$(2.21) \quad \omega_{\mathbb{C}} = \omega_1 \cdots \omega_n,$$

compare (2.18). Let $\Sigma = \Sigma_{2n}$ be the spinor representation. Then Clifford multiplication by the ω_j defines unitary involutions of Σ . By (2.20), there is an orthogonal decomposition of Σ into simultaneous eigenspaces Σ_{ε} , where ε runs over all n -tuples in $\{1, -1\}^n$ and where ω_j acts by multiplication with ε_j on Σ_{ε} , $1 \leq j \leq n$. Because Clifford multiplication with e_{2j-1} or e_{2j} anti-commutes with ω_j and commutes with ω_k , for $1 \leq k \neq j \leq n$, we have

$$(2.22) \quad e_{2j-1}\Sigma_{\varepsilon} = e_{2j}\Sigma_{\varepsilon} = \Sigma_{\delta},$$

where $\delta_k = \varepsilon_k$ for all $1 \leq k \neq j \leq n$ and $\delta_j = -\varepsilon_j$. In particular, all the subspaces Σ_{ε} have the same dimension, which is, for that reason, equal to $\dim \Sigma / 2^n = 1$. Clifford multiplication by the complex volume form acts by $\varepsilon_1 \cdots \varepsilon_n$ on Σ_{ε} , by (2.21), and hence the summands of the usual super-symmetry

$$(2.23) \quad \Sigma = \Sigma^+ \oplus \Sigma^-$$

are given by

$$(2.24) \quad \Sigma^+ = \bigoplus_{\varepsilon_1 \cdots \varepsilon_n = 1} \Sigma_{\varepsilon} \quad \text{and} \quad \Sigma^- = \bigoplus_{\varepsilon_1 \cdots \varepsilon_n = -1} \Sigma_{\varepsilon}.$$

2.3. Complex Hyperbolic Spaces. The discussion in this section will be used in Section 11.2. We represent complex hyperbolic space $\mathbb{C}H^n$ by the symmetric pair $(\mathrm{SU}(1, n), \mathrm{S}(\mathrm{U}(1) \times \mathrm{U}(n)))$ and endow the Lie algebra $\mathfrak{su}(1, n)$ of $\mathrm{SU}(1, n)$ with the non-degenerate symmetric bilinear from

$$(2.25) \quad (X, Y) := \frac{1}{2} \mathrm{Re} \, \mathrm{tr} \, XY,$$

⁹Note that the sign convention is opposite to the one in [LaMi], page 43.

a multiple of the Killing form of $\mathfrak{su}(1, n)$. We identify

$$(2.26) \quad \mathrm{S}(\mathrm{U}(1) \times \mathrm{U}(n)) = \left\{ \begin{pmatrix} \det A^{-1} & 0 \\ 0 & A \end{pmatrix} \mid A \in \mathrm{U}(n) \right\} \cong \mathrm{U}(n)$$

and, correspondingly,

$$(2.27) \quad \mathfrak{s}(\mathfrak{u}(1) \oplus \mathfrak{u}(n)) = \left\{ \begin{pmatrix} -\mathrm{tr} A & 0 \\ 0 & A \end{pmatrix} \mid A \in \mathfrak{u}(n) \right\} \cong \mathfrak{u}(n).$$

The orthogonal complement \mathfrak{p} of $\mathfrak{u}(n)$ in $\mathfrak{su}(1, n)$ is

$$(2.28) \quad \mathfrak{p} = \left\{ \begin{pmatrix} 0 & x^* \\ x & 0 \end{pmatrix} \mid x \in \mathbb{C}^n \right\} \cong \mathbb{C}^n,$$

where we note that the latter isomorphism corresponds to the standard complex structure and Riemannian metric of $\mathbb{C}H^n$. With respect to the identifications (2.26) – (2.28), we get

$$(2.29) \quad [A, B] = AB - BA, \quad [A, x] = Ax + x \cdot \mathrm{tr} A, \quad [x, y] = xy^* - yx^*$$

for the different Lie brackets and

$$(2.30) \quad \alpha(A)x := \mathrm{Ad}_A x = Ax \det A$$

for the adjoint representation $\alpha : \mathrm{U}(n) \rightarrow \mathrm{SO}(\mathfrak{p}) \cong \mathrm{SO}(2n)$. We note that α is an $n + 1$ to 1 immersion. If n is odd, then α lifts to $\hat{\alpha} : \mathrm{U}(n) \rightarrow \mathrm{Spin}(\mathfrak{p})$. If n is even, then α does not lift.

We note that the coefficients of the matrix $xy^* - yx^* \in \mathfrak{u}(n)$ in (2.29) are $x_j \bar{y}_k - y_j \bar{x}_k$. In particular, for the standard unit vectors e_j and e_k in \mathbb{C}^n and complex numbers x, y , we have

$$(2.31) \quad [xe_j, ye_k] = x\bar{y}E_{jk} - y\bar{x}E_{kj} \in \mathfrak{u}(n),$$

where E_{jk} denotes the matrix with entries δ_{jk} .

Let $T = e_1 \in \mathbb{C}^n \cong \mathfrak{p}$ and set $\mathfrak{a} := \mathbb{R}T$. The orthogonal complement of \mathfrak{a} in \mathbb{C}^n consists of all $x \in \mathbb{C}^n$ with $x_1 \in \mathrm{Im} \mathbb{C}$, that is, x_1 is purely imaginary. Let $\mathfrak{z} := \mathbb{R}Z$ with

$$(2.32) \quad Z := ie_1 - iE_{11} \in \mathfrak{p} \oplus \mathfrak{u}(n).$$

We have $[\mathfrak{z}, \mathfrak{z}] = 0$ and

$$(2.33) \quad [T, Z] = 2Z.$$

Let \mathfrak{x} be the space of all

$$(2.34) \quad X_x := x + \bar{x}_2 E_{12} - x_2 E_{21} + \dots + \bar{x}_n E_{1n} - x_n E_{n1} \in \mathfrak{p} \oplus \mathfrak{u}(n),$$

where $x \in \mathbb{C}^{n-1} = \{x \in \mathbb{C}^n \mid x_1 = 0\}$. Then $[\mathfrak{z}, \mathfrak{x}] = 0$ and

$$(2.35) \quad [T, X_x] = X_x,$$

$$(2.36) \quad [X_x, X_y] = 2 \mathrm{Im}(\bar{x}y)Z.$$

Set $\mathfrak{n} := \mathfrak{z} \oplus \mathfrak{x}$ and $\mathfrak{s} := \mathfrak{a} \oplus \mathfrak{n}$. By the above, \mathfrak{n} is a two-step nilpotent subalgebra of $\mathfrak{su}(1, n)$ and \mathfrak{s} is a solvable extension of \mathfrak{n} . The subgroups A , N , and S of $SU(1, n)$ corresponding to \mathfrak{a} , \mathfrak{n} , and \mathfrak{s} satisfy $S = AN$ and $SU(1, n) = U(n)AN$ (Iwasawa decomposition of $SU(1, n)$).

Let $p \in \mathbb{C}H^n$ be the point fixed by $U(n)$. Then the orbit map

$$(2.37) \quad \Phi : S \rightarrow \mathbb{C}H^n, \quad \Phi(s) = sp,$$

is a diffeomorphism, that is, S acts simply transitively on $\mathbb{C}H^n$. Endow S with the left-invariant Riemannian metric such that the differential $d\Phi : T_e S \rightarrow T_p \mathbb{C}H^n$ is isometric. Since S acts isometrically on $\mathbb{C}H^n$, we then get that Φ is an S -equivariant isometry. That is, we can think of $\mathbb{C}H^n$ as S , endowed with the chosen left-invariant metric. With respect to this metric, we get that \mathfrak{a} , \mathfrak{z} , and \mathfrak{x} are perpendicular and that

$$(2.38) \quad |T| = 1, \quad |Z| = 1, \quad \langle X_x, X_y \rangle = \operatorname{Re} \bar{x}y.$$

Define

$$(2.39) \quad JX_x = J_Z X_x := X_{ix}.$$

Then J is skew-symmetric with $J^2 = -1$. Moreover, by (2.36) and (2.38),

$$(2.40) \quad \langle [X_x, X_y], Z \rangle = 2\langle JX_x, X_y \rangle.$$

As a preferred basis of \mathfrak{s} , we choose the $2n$ -tuple of vectors $X_1 := T$, $Y_1 = Z$,

$$(2.41) \quad X_j := e_j + E_{1j} - E_{j1} \quad \text{and} \quad Y_j := JX_j = ie_j - iE_{1j} - iE_{j1},$$

where $2 \leq j \leq n$. By (2.36) and (2.39),

$$(2.42) \quad [X_j, Y_k] = 2\delta_{jk}Z.$$

In conclusion, N is isomorphic to the standard Heisenberg group of dimension $2n - 1$. Since left-invariant vector fields on S are T -parallel, the Weingarten operator $W = \nabla T$ of N in S is given by

$$(2.43) \quad WZ = [Z, T] = -2Z \quad \text{and} \quad WX_x = [X_x, T] = -X_x,$$

by (2.33) and (2.35).

3. DIRAC SYSTEMS AND DISTANCE FUNCTIONS

3.1. Dirac Systems. The setup and the results from [BBC2] are fundamental for the discussion of this section. Let $I \subseteq \mathbb{R}$ be an interval and H be a separable complex Hilbert space. Fix an origin $t_0 \in I$.

For each $t \in I$, let $(\cdot, \cdot)_t$ be a scalar product on H which is compatible with the Hilbert space structure of H and such that $(\cdot, \cdot)_{t_0}$ coincides with the given scalar product of H . Let $\|\cdot\|_t$ be the norm associated

to $(\cdot, \cdot)_t$. Let H_t be H , but equipped with $(\cdot, \cdot)_t$, and denote by \mathcal{H} the family of Hilbert spaces H_t , $t \in I$. Assume that, for all $a < b$ in I , there is a constant $C = C(a, b)$ such that

$$(3.1) \quad |(\sigma_1, \sigma_2)_s - (\sigma_1, \sigma_2)_t| \leq C \|\sigma_1\|_s \|\sigma_2\|_s |s - t|,$$

for all $s, t \in [a, b]$ and $\sigma_1, \sigma_2 \in H$. In other words, if $G_t \in \mathcal{L}(H)$ denotes the positive definite and symmetric operator on $H = H_{t_0}$ with

$$(3.2) \quad (G_t \sigma_1, \sigma_2)_{t_0} = (\sigma_1, \sigma_2)_t,$$

for all $\sigma_1, \sigma_2 \in H$, then the map

$$G : I \rightarrow \mathcal{L}(H), \quad G(t) := G_t,$$

is in $\text{Lip}_{\text{loc}}(I, \mathcal{L}(H))$. In particular, G is weakly differentiable almost everywhere in I with weak derivative G' in $L_{\text{loc}}^\infty(I, \mathcal{L}(H))$. Moreover, G'_t is symmetric on H_{t_0} (for almost all $t \in I$) and we have

$$(3.3) \quad \Gamma := \frac{1}{2} G^{-1} G' \in L_{\text{loc}}^\infty(I, \mathcal{L}(H)).$$

We set

$$(3.4) \quad \partial := \left(\frac{d}{dt} + \Gamma \right) : \text{Lip}_{\text{loc}}(I, H) \rightarrow L_{\text{loc}}^\infty(I, H).$$

By the definition of ∂ , the function $(\sigma_1, \sigma_2) = (\sigma_1(t), \sigma_2(t))_t$ satisfies

$$(3.5) \quad (\sigma_1, \sigma_2)' = (\partial \sigma_1, \sigma_2) + (\sigma_1, \partial \sigma_2),$$

for all $\sigma_1, \sigma_2 \in \text{Lip}_{\text{loc}}(I, H)$, where the prime indicates differentiation with respect to t .

As a second data, let \mathcal{A} be a family of operators A_t , $t \in I$, on H with common dense domain H_A such that A_t is self-adjoint in H_t and such that the inclusion $H_A \hookrightarrow H$ is compact with respect to the graph norms of the A_t . Assume that, for all $a < b$ in I , there is a constant $C = C(a, b)$ such that

$$(3.6) \quad |(A_s \sigma_1, \sigma_2)_s - (A_t \sigma_1, \sigma_2)_t| \leq C(\|\sigma_1\|_s + \|A_s \sigma_1\|_s) \|\sigma_2\|_s |s - t|,$$

for all $s, t \in [a, b]$ and $\sigma_1, \sigma_2 \in H_A$.

As a final data, let

$$(3.7) \quad T \in \text{Lip}_{\text{loc}}(I, \mathcal{L}(H)) \cap L_{\text{loc}}^\infty(I, \mathcal{L}(H_A)),$$

and suppose that

$$(3.8) \quad T_t^* = T_t^{-1} = -T_t \quad \text{on } H_t, \forall t \in I,$$

$$(3.9) \quad A_t T_t = -T_t A_t \quad \text{on } H_A, \forall t \in I,$$

$$(3.10) \quad \partial T = T \partial \quad \text{on } \text{Lip}_{\text{loc}}(I, H).$$

Following [BBC2], a *Dirac system* over I consists of data \mathcal{H} , \mathcal{A} , and T as above.

Let $\mathcal{D} := (\mathcal{H}, \mathcal{A}, T)$ be a Dirac system over I . Set

$$(3.11) \quad \mathcal{L}_{\text{loc}}(\mathcal{D}) := \text{Lip}_{\text{loc}}(I, H) \cap L_{\text{loc}}^{\infty}(I, H_A),$$

and denote by $\mathcal{L}_c(\mathcal{D})$ and $\mathcal{L}_{cc}(\mathcal{D})$ the subspaces of $\mathcal{L}_{\text{loc}}(\mathcal{D})$ of maps with compact support in I and the interior of I , respectively. On $\mathcal{L}_c(\mathcal{D})$, we define the inner product

$$(3.12) \quad (\sigma_1, \sigma_2) := \int_I (\sigma_1, \sigma_2) = \int_I (\sigma_1(t), \sigma_2(t))_t dt,$$

and let $L^2(\mathcal{D})$ be the corresponding Hilbert space of square-integrable maps, also denoted by $L^2(\mathcal{H})$.

The *Dirac operator* of \mathcal{D} is the operator

$$(3.13) \quad D := T(\partial + A) : \mathcal{L}_{\text{loc}}(\mathcal{D}) \rightarrow L_{\text{loc}}^{\infty}(I, H).$$

By (3.5) and (3.8)–(3.10),

$$(3.14) \quad \int_{[a,b]} (D\sigma_1, \sigma_2) = \int_{[a,b]} (\sigma_1, D\sigma_2) - (\sigma_1, T\sigma_2)|_a^b,$$

for all $\sigma, \tau \in \mathcal{L}_{\text{loc}}(\mathcal{D})$ and $a < b$ in I .

A *super-symmetry* for a Dirac system \mathcal{D} as above is a decomposition $H = H^+ \oplus H^-$ such that, with $H_A^{\pm} := H_A \cap H^{\pm}$,

$$(3.15) \quad H^+ \perp H^- \text{ in } H_t \quad \text{and} \quad T_t H^{\pm} = H^{\mp},$$

$$(3.16) \quad H_A = H_A^+ \oplus H_A^- \quad \text{and} \quad A_t H_A^{\pm} \subseteq H^{\pm}.$$

We write H_t^{\pm} for H^{\pm} endowed with the inner product $(\cdot, \cdot)_t$. By (3.15), $H_t = H_t^+ \oplus H_t^-$ as a Hilbert space. By (3.15),

$$(3.17) \quad A_t = \begin{pmatrix} A_t^+ & 0 \\ 0 & A_t^- \end{pmatrix},$$

where A_t^{\pm} is a self-adjoint operator in H_t^{\pm} with domain H_A^{\pm} and where

$$(3.18) \quad A_t^- = T_t A_t^+ T_t = T_t (-A_t^+) T_t^{-1},$$

by (3.9). We can decompose

$$(3.19) \quad L^2(\mathcal{D}) = L^{2+}(\mathcal{D}) \oplus L^{2-}(\mathcal{D}) = L^2(\mathcal{H}^+) \oplus L^2(\mathcal{H}^-),$$

where $L^2(\mathcal{H}^{\pm})$ consists of the subspace of sections in $L^2(\mathcal{H})$ with image in H^{\pm} . Similar notation will be employed for other spaces.

By (3.16) and the definition of ∂ , see (3.4),

$$(3.20) \quad \partial = \begin{pmatrix} \partial^+ & 0 \\ 0 & \partial^- \end{pmatrix}, \quad \text{where} \quad \partial^- = T \partial^+ T^{-1},$$

by (3.10). Hence, by (3.15),

$$(3.21) \quad D = \begin{pmatrix} 0 & D^- \\ D^+ & 0 \end{pmatrix}.$$

Clearly, D^- is the formal adjoint of D^+ .

3.2. Distance Functions. Let U be an open subset in a Riemannian manifold M . We say that a function $f : U \rightarrow \mathbb{R}$ is a *distance function* if f is C^1 and $T := \text{grad } f$ is a unit vector field. There is a synthetic characterization of distance functions, compare [BGS, pp 24–25] or also Proposition IV.3.1 in [Ba]. If f is a distance function, then the solution curves of the vector field T are unit speed geodesics, called *T -geodesics*.

Busemann functions are C^2 distance functions, see Proposition 3.1 in [HeIH] or Proposition IV.3.2 in [Ba]. We assume from now on that $f : U \rightarrow \mathbb{R}$ is a C^2 distance function. Then $T := \text{grad } f$ is a C^1 unit vector field and the *cross sections* $N_t = f^{-1}(t)$ are C^2 hypersurfaces. For simplicity, we assume throughout that the cross sections N_t are compact and that the flow of T induces a C^1 diffeomorphism

$$(3.22) \quad F : I \times N \rightarrow U,$$

where I is some interval and $N = N_{t_0}$ for some $t_0 \in I$. In what follows, we often identify U with $I \times N$ by identifying $(t, x) \in I \times N$ with $F(t, x) \in U$.

Let $c = c(s)$ be a C^1 curve in U and $T(s) := T(c(s))$ be T along c , a C^1 curve of unit vectors. Then the variation field $J = J(t) := (\partial_s \gamma)(0, t)$ of the geodesic variation $\gamma_s = \gamma(s, t) := \exp(tT(s))$ satisfies $J(0) = \dot{c}(0)$. A vector field which arises in this way will be called a *T -Jacobi field*.

LEMMA 3.23. *A T -Jacobi field J satisfies the Jacobi equation*

$$J'' + R(J, T)T = 0.$$

Moreover, J and J' depend continuously on $J(0)$.

Proof. Recall that the Riemannian manifold M is smooth, hence its geodesic flow $\Phi = \Phi(t, v)$, $t \in \mathbb{R}$ and $v \in TM$, is also smooth. Since $\gamma'(s, t) = \Phi(t, T(s))$, we get that γ and γ' are C^1 . Therefore

$$(3.24) \quad J = \partial_s \gamma \quad \text{and} \quad J' = \nabla_t \partial_s \gamma = \nabla_s \partial_t \gamma = \nabla_s T$$

exist and are continuous. Moreover,

$$(3.25) \quad (\partial_s \gamma')(s, t) = \Phi_{t*}(\partial_s T(s)) = (J(s, t), J'(s, t))$$

with respect to the standard decomposition of TTM in horizontal and vertical component, see for example Proposition IV.1.13 in [Ba]. Hence

J and J' depend continuously on $\dot{c}(0)$ and J satisfies the asserted Jacobi equation. \square

REMARK 3.26. With respect to the (t, x) -coordinates, the Riemannian metric on U is of the form $g = dt^2 + g_t$, where g_t , $t \in I$, is a family of Riemannian metrics on N . In [BB2], pages 596 and 609, it is stated erroneously that g_t and $\partial_t g_t$ are C^1 on U . This is wrong in general, since it would imply that T is C^2 . Clearly, since T is C^1 , $g_t(x)$ is C^1 in (t, x) . Lemma 3.23 implies that $g_t(x)$ is two times continuously differentiable in t . This is sufficient for the discussion in [BB2] and the arguments below.

For $t \in I$, we let $S = S_t$ and $W = W_t$ be, respectively, the *second fundamental form* and the *Weingarten operator* of the C^2 submanifold N_t with respect to the normal vector field T ,

$$(3.27) \quad WX = \nabla_X T, \quad S(X, Y) = \langle \nabla_X Y, T \rangle = -\langle WX, Y \rangle,$$

where X and Y are C^1 vector fields tangent to N_t . Since T is C^1 , S and W are continuous tensor fields over U . By (3.24), Jacobi fields J as in Lemma 3.23 satisfy $J' = WJ$.

Let $E \rightarrow M$ be a smooth vector bundle with smooth connection ∇^E .

LEMMA 3.28. *Let X be a vector field and σ be a section of E over U , respectively. Assume that the restrictions of X and σ to N are C^1 and that X and σ are parallel in the T -direction. Then X and σ are C^1 . Moreover, $\nabla_T^E \nabla_X^E \sigma$ exists, is continuous, and satisfies*

$$\nabla_T^E \nabla_X^E \sigma + \nabla_{WX}^E \sigma + R^E(X, T)\sigma = 0.$$

Proof. Let $\Psi : \mathbb{R} \times (TM \oplus E) \rightarrow E$ be the smooth map which associates to $t \in \mathbb{R}$ and $(v, e) \in TM \oplus E$ (where $v \in TM$ and $e \in E$ have the same foot point) the parallel translate $\sigma(t)$ of e along the geodesic γ with $\gamma'(0) = v$. Then, with σ as in the assertion, we have $\sigma(F(t, x)) = \Psi(t - t_0, T(x), \sigma(x))$, where we recall that $N = N_{t_0}$. Hence X and σ are C^1 , where X corresponds to the special case $E = TM$.

Since $\nabla_T^E \sigma = 0$ and T is C^1 , the T derivatives of the coefficients of σ with respect to a smooth local frame of E are C^1 . Hence $\nabla_T^E \nabla_X^E \sigma$ exists, is continuous, and is given by

$$\begin{aligned} \nabla_T^E \nabla_X^E \sigma &= \nabla_X^E \nabla_T^E \sigma - \nabla_{\nabla_X T}^E \sigma + R^E(T, X)\sigma \\ &= -\nabla_{WX}^E \sigma - R^E(X, T)\sigma. \end{aligned} \quad \square$$

Among others, the case $E = TM$ is interesting. In this case, vector fields over N which are tangent to N can, in general, only be chosen to be C^1 .

COROLLARY 3.29. *The tensor field W has a continuous derivative W' in the T -direction and satisfies the Riccati equation*

$$W' + W^2 + R(\cdot, T)T = 0.$$

Proof. Choose $\sigma = T$ in Lemma 3.28 and recall that $W = \nabla T$. \square

The eigenvalues $\kappa_2, \dots, \kappa_m$ of W_t are the *principal curvatures* of the cross section N_t . We let

$$(3.30) \quad \kappa := \kappa_2 + \dots + \kappa_m = \operatorname{tr} W = \operatorname{div} T.$$

The maps

$$(3.31) \quad F_t : N = N_{t_0} \rightarrow N_t, \quad F_t(x) := F(t, x)$$

are diffeomorphisms whose Jacobian determinants will be denoted by $j = j(t, x)$. Since $\kappa = \operatorname{div} T$, the latter satisfy the differential equation

$$(3.32) \quad j' = \kappa j.$$

By Corollary 3.29, we also have

$$(3.33) \quad \kappa' = -\|W\|^2 - \operatorname{Ric}(T, T),$$

where $\|W\| = (\operatorname{tr} W^2)^{1/2}$ is the Euclidean norm of W .

Let C_R , C_R^E , and C_W be uniform upper bounds for the operator norms of the curvature R of M , the curvature R^E of E , and W , respectively. Then κ , the t -derivative κ' of κ , and $\|W\|$ are uniformly bounded, and as respective uniform upper bounds C_κ , C'_κ , and C_w we may take

$$(3.34) \quad C_\kappa = mC_W, \quad C'_\kappa = m(C_W^2 + C_R), \quad C_w = \sqrt{m}C_W,$$

where we use (3.33) for the second assertion. By (3.32), we have

$$(3.35) \quad e^{-C(t-s)}j(s, x) \leq j(t, x) \leq e^{C(t-s)}j(s, x),$$

or all $s < t$ in I and $x \in N$, where $C = C_\kappa$.

3.3. From Distance Functions to Dirac Systems. Let $E \rightarrow M$ be a smooth Dirac bundle. Denote the Hermitian product on E by $\langle \cdot, \cdot \rangle$. Our aim is to identify these data over U with a Dirac system over I as in Section 3.1.

For any $t \in I$ and any given Riemannian or Hermitian vector bundle over U with any given metric connection, we let P_t^\parallel be parallel translation along the T -geodesics from N to N_t . For a section σ of the vector bundle over N , we define a section $P^\parallel \sigma$ over U by

$$(3.36) \quad (P^\parallel \sigma)(t, x) := P_t^\parallel(\sigma(x)), \quad x \in N.$$

Thus $P^\parallel \sigma$ is the extension of σ to U which is parallel along the T -geodesics, and this point of view is convenient in arguments and formulations below. Furthermore, time dependent sections over N correspond to the space of all sections over U ,

$$(3.37) \quad (P^\parallel \sigma)(t, x) := P_t^\parallel(\sigma(t, x)), \quad t \in I, x \in N.$$

We also let $P_t^\parallel \sigma := P^\parallel \sigma|_{N_t}$.

Now let $H := L^2(N, E)$, the Hilbert space of square integrable sections of E over $N = N_{t_0}$. For $\sigma, \tau \in H$, the L^2 product of the sections $P_t^\parallel \sigma, P_t^\parallel \tau$ with respect to N_t is given by

$$(3.38) \quad (\sigma, \tau)_t := \int_N \langle \sigma(x), \tau(x) \rangle j(t, x) dx,$$

where dx denotes the volume element of N . Hence, for each $t \in I$, the correspondence $\sigma \leftrightarrow P_t^\parallel \sigma$ identifies the Hilbert space $L^2(N_t, E)$ topologically with H . The following estimate settles the requirement on the family \mathcal{H} formulated in (3.1).

LEMMA 3.39. *For all $s < t$ in I and $\sigma_1, \sigma_2 \in H$,*

$$|(\sigma_1, \sigma_2)_t - (\sigma_1, \sigma_2)_s| \leq (e^{C(t-s)} - 1) \|\sigma_1\|_s \|\sigma_2\|_s,$$

where $C = C_\kappa$.

Proof. By (3.38) and (3.35),

$$\begin{aligned} |(\sigma_1, \sigma_2)_t - (\sigma_1, \sigma_2)_s| &\leq \int_N |\langle \sigma_1(x), \sigma_2(x) \rangle| (j(t, x) - j(s, x)) dx \\ &\leq \int_N |\sigma_1(x)| |\sigma_2(x)| (e^{C(t-s)} - 1) j(s, x) dx \\ &\leq (e^{C(t-s)} - 1) \|\sigma_1\|_s \|\sigma_2\|_s. \end{aligned} \quad \square$$

LEMMA 3.40. *For all $s < t$ in I and C^1 sections σ of E over U which are parallel in the T -direction,*

$$e^{C_0(s-t)} (\|\sigma\|_s^2 + \|\nabla^E \sigma\|_s^2) \leq \|\sigma\|_t^2 + \|\nabla \sigma\|_t^2 \leq e^{C_0(t-s)} (\|\sigma\|_s^2 + \|\nabla^E \sigma\|_s^2),$$

where $C_0 = C_\kappa + mC_R^E + 2C_W$.

Proof. Using $\langle \sigma, \sigma \rangle' = 0$, we obtain

$$(\|\sigma\|_t^2 + \|\nabla^E \sigma\|_t^2)' = \int_N (\langle \nabla^E \sigma, \nabla^E \sigma \rangle' + (\langle \sigma, \sigma \rangle + \langle \nabla^E \sigma, \nabla^E \sigma \rangle) \kappa) j.$$

By Lemma 3.28,

$$\langle \nabla^E \sigma, \nabla^E \sigma \rangle' = 2 \sum_{2 \leq i \leq m} (\langle R^E(T, X_i) \sigma, \nabla_{X_i}^E \sigma \rangle - \langle \nabla_{W X_i}^E \sigma, \nabla_{X_i}^E \sigma \rangle),$$

where (T, X_1, \dots, X_n) is a local orthonormal frame of M . Hence

$$\begin{aligned} |(\|\sigma\|_t^2 + \|\nabla^E \sigma\|_t^2)'| &\leq mC_R^E \|\sigma\|_t^2 + C_R^E \|\nabla^E \sigma\|_t^2 \\ &\quad + 2C_W \|\nabla^E \sigma\|_t^2 + C_\kappa (\|\sigma\|_t^2 + \|\nabla^E \sigma\|_t^2) \\ &\leq (C_\kappa + mC_R^E + 2C_W)(\|\sigma\|_t^2 + \|\nabla^E \sigma\|_t^2). \quad \square \end{aligned}$$

Along the cross sections N_t , we change Clifford multiplication and connection of E according to (2.8) and (2.9). Denote by ∇^t the new connection and by D_t the associated Dirac operator as in (2.10). We note that Clifford multiplication with T is ∇^t -parallel. For convenience, we will not keep the $*$ notation, but will write $TX\sigma$ instead of $X * \sigma$. With this in mind, the Dirac operators D and D_t are related by

$$(3.41) \quad D = T(\nabla_T^E + \sum TX_i \nabla_{X_i}^E) = T((\nabla_T^E + \frac{\kappa}{2}) - D_t),$$

where (T, X_2, \dots, X_m) is a local orthonormal frame of M .

LEMMA 3.42. *For any C^1 section σ of E over U which is parallel in the T -direction,*

$$\|\nabla^E \sigma|_{N_t} - \nabla^t \sigma\|^2 = \frac{1}{4} \|W\|^2 |\sigma|^2 \quad \text{and} \quad |TD\sigma - D_t \sigma|^2 = \frac{1}{4} \kappa^2 |\sigma|^2.$$

Proof. The second assertion is immediate from (3.41). As for the first, let (T, X_2, \dots, X_m) be an orthonormal frame of M . Then

$$\begin{aligned} 4\|\nabla^E \sigma|_{N_t} - \nabla^t \sigma\|^2 &= 4 \sum \langle \nabla_{X_i}^E \sigma - \nabla_{X_i}^t \sigma, \nabla_{X_i}^E \sigma - \nabla_{X_i}^t \sigma \rangle \\ &= \sum \langle TWX_i \sigma, TWX_i \sigma \rangle \\ &= \sum |WX_i|^2 |\sigma|^2 = \|W\|^2 |\sigma|^2. \quad \square \end{aligned}$$

Since the cross sections N_t are C^2 submanifolds of U , the restrictions of E to them are C^2 bundles. However, because of the term involving $W = \nabla T$, the connection ∇^t is, in the generality we strive for, only continuous. If ∇^t were a C^1 connection, we would get (2.12) for its curvature, now denoted R^t . The right hand side of (2.12) makes sense in the case where W is only continuous, so that we may consider it as defining R^t . Approximating N_t by smooth submanifolds and C^1 sections by smooth sections, (2.15) implies that

$$(3.43) \quad (\nabla^t \sigma_1, \nabla^t \sigma_2)_t + (K^t \sigma_1, \sigma_2)_t = (D_t \sigma_1, D_t \sigma_2)_t$$

for all C^1 sections σ and τ of the restriction of E to N_t , where the curvature term in the Lichnerowicz formula as in (2.15) is now denoted by K^t . We recall from (2.16) that K^t is uniformly bounded.

We extend our correspondence $\sigma \leftrightarrow P^\# \sigma$ as in (3.36) and (3.37): Since T is parallel in the T -direction, Clifford multiplication by T along N satisfies

$$(3.44) \quad TP^\# \sigma = P^\# T\sigma \quad \text{and} \quad \nabla_T P^\# \sigma = P^\# \sigma',$$

for any time dependent section σ of E over N . Finally, we define A_t to be the differential operator on sections of E over N which corresponds to the operator $-D_t$,

$$(3.45) \quad P_t^\#(A_t \sigma) = -D_t P_t^\# \sigma.$$

In this notation, D corresponds to the operator

$$(3.46) \quad T(\partial + A), \quad \text{where } \partial := \frac{d}{dt} + \frac{\kappa}{2},$$

where $\kappa/2$ takes the role of Γ in (3.3) and (3.4). Thus we have associated the Dirac system

$$(3.47) \quad \mathcal{D} := (\mathcal{H}, \mathcal{A}, T)$$

to the distance function f on and the Dirac bundle E over U , where we recall from (3.32) that κ (which occurs in the definition of ∂) is defined by these data. We will now proceed with discussing the requirement for Dirac systems as formulated in Section 3.1. We already observed that Lemma 3.39 settles (3.1). Furthermore, Clifford multiplication by T satisfies the requirements (3.7)–(3.10), by (3.44) and since Clifford multiplication by T is ∇^t -parallel.

It follows from (3.43) that, on sections of the restriction of E to N_t , the graph norm of D_t is equivalent to the H^1 norm. In particular, D_t is self-adjoint with domain $H^1(N_t, E)$ in $L^2(N_t, E)$. Moreover, since the inclusion $H^1(N_t, E) \hookrightarrow L^2(N_t, E)$ is compact, the spectrum of D_t consists of eigenvalues with finite multiplicities. We also observe that, for any section σ of E over N , $P\sigma|_{N_t} \in H^1(N_t, E)$ if and only if $\sigma \in H^1(N, E)$, by Lemma 3.40. Thus the operators A_t are all self-adjoint with the same domain, $H_A := H^1(N, E)$, in $H = L^2(N, E)$, and the embedding $H_A \rightarrow H$ is compact with respect to the graph norm of any of the operators A_t . This settles the first part of the requirements for the A_t in Section 3.1.

LEMMA 3.48. *For any C^1 section σ of E over U which is parallel in the T -direction,*

$$D'_t \sigma = \sum_{2 \leq i \leq m} TX_i \{R^E(T, X_i) \sigma - \nabla_{WX_i}^E \sigma\} + \frac{\kappa'}{2} \sigma,$$

where (T, X_2, \dots, X_m) is a local orthonormal frame of M .

Proof. By Lemma 3.28,

$$\begin{aligned} D'_t \sigma &= \sum_{2 \leq i \leq m} TX_i \nabla_T^E \nabla_{X_i}^E \sigma + \frac{\kappa'}{2} \sigma \\ &= \sum_{2 \leq i \leq m} TX_i \{R^E(T, X_i) \sigma - \nabla_{W X_i}^E \sigma\} + \frac{\kappa'}{2} \sigma. \end{aligned} \quad \square$$

COROLLARY 3.49. *For any C^1 section σ of E over U , which is parallel in the T -direction,*

$$\|D'_t \sigma\|_t \leq C_1 \|\sigma\|_t + C_w \|\nabla^E \sigma\|_t \leq C_2 \|\sigma\|_t + C_w \|D_t \sigma\|_t,$$

where $C_1 = mC_R^E + C'_\kappa$ and $C_2 = mC_R^E + C'_\kappa + C_w^2 + C_w C_K^{1/2}$.

Proof. By Lemmas 3.48 and 3.42, we have, at any point p of N_t ,

$$\begin{aligned} |D'_t \sigma| &\leq (mC_R^E + \frac{1}{2}C'_\kappa)|\sigma| + \sum |\kappa_i| |\nabla_{X_i}^E \sigma| \\ &\leq (mC_R^E + \frac{1}{2}C'_\kappa)|\sigma| + \|W\| \|\nabla^E \sigma\| \\ &\leq (mC_R^E + \frac{1}{2}C'_\kappa)|\sigma| + C_w \|\nabla^E \sigma\| \\ &\leq (mC_R^E + \frac{1}{2}C'_\kappa + \frac{1}{2}C_w^2)|\sigma| + C_w \|\nabla^t \sigma\|, \end{aligned}$$

where (T, X_2, \dots, X_m) is an orthonormal frame at p such that the X_i are eigenvectors of W with corresponding eigenvalues κ_i . By (3.43),

$$\|\nabla^t \sigma\|_t^2 \leq \|D_t \sigma\|_t^2 + C_K \|\sigma\|_t^2. \quad \square$$

LEMMA 3.50. *For all $s < t$ in I and C^1 sections $\sigma_1, \sigma_2 \in H$ of E ,*

$$|(A_t \sigma_1, \sigma_2)_t - (A_s \sigma_1, \sigma_2)_s| \leq C(e^{C_0(t-s)/2} - 1)(\|\sigma_1\|_s + \|A_s \sigma_1\|_s) \|\sigma_2\|_s,$$

where $C = C(C_R, C_R^E, C_W, m)$.

Proof. Extend σ_1 and σ_2 by parallel translation along the T -geodesics. Then D_t corresponds to $-A_t$, and we get

$$\begin{aligned} |(D_t \sigma_1, \sigma_2)_t - (D_s \sigma_1, \sigma_2)_s| &\leq \left| \int_s^t \int_N (\langle D_r \sigma_1, \sigma_2 \rangle j)' \right| \\ &\leq \int_s^t \int_N |\langle D'_r \sigma_1, \sigma_2 \rangle + \langle D_r \sigma_1, \sigma_2 \rangle \kappa| j \\ (3.51) \quad &\leq \int_s^t \int_N (\|D'_r \sigma_1\| + C_\kappa \|D_r \sigma_1\|) \|\sigma_2\| j. \end{aligned}$$

By Corollary 3.49 and Lemma 3.40, the first term on the right hand side of (3.51) can be estimated by

$$\begin{aligned} \int_s^t \int_N \|D'_r \sigma_1\| \|\sigma_2\| j &\leq 2(C_1 + C_w) \int_s^t (\|\sigma_1\|_r^2 + \|\nabla^E \sigma_1\|_r^2)^{1/2} \|\sigma_2\|_r \\ &\leq 2(C_1 + C_w) \int_s^t e^{C_0(r-s)/2} (\|\sigma_1\|_s^2 + \|\nabla^E \sigma_1\|_s^2)^{1/2} \|\sigma_2\|_s \\ &= 4 \frac{C_1 + C_w}{C_0} (e^{C_0(t-s)/2} - 1) (\|\sigma_1\|_s^2 + \|\nabla^E \sigma_1\|_s^2)^{1/2} \|\sigma_2\|_s. \end{aligned}$$

Concerning the second term on the right hand side of (3.51), namely the integral of $\|D_r \sigma_1\| \|\sigma_2\| j$, we note that $\|D_r \sigma\| \leq \sqrt{m-1} \|\nabla^r \sigma\|$. Hence we can estimate this term in a similar way, using Lemma 3.42. We arrive at an estimate

$$\begin{aligned} |(D_t \sigma_1, \sigma_2)_t - (D_s \sigma_1, \sigma_2)_s| \\ \leq C' (e^{C_0(t-s)/2} - 1) (\|\sigma_1\|_s^2 + \|\nabla^E \sigma_1\|_s^2)^{1/2} \|\sigma_2\|_s, \end{aligned}$$

where $C' = C'(C_R, C_R^E, C_W, m)$. Finally, the Bochner formula (3.43) and the ensuing lines show that

$$\begin{aligned} \|\sigma_1\|_s^2 + \|\nabla^E \sigma_1\|_s^2 &\leq C(C_R^E, m) (\|\sigma_1\|_s + \|D_s \sigma_1\|_s) \\ &= C(C_R^E, m) (\|\sigma_1\|_s + \|A_s \sigma_1\|_s). \end{aligned} \quad \square$$

Since C^1 sections are dense in $H^1(N, E)$, Lemma 3.50 confirms the remaining requirements for the operators A_t in Section 3.1. Thus the system $\mathcal{D} = (\mathcal{H}, \mathcal{A}, T)$ over I from (3.47) is a Dirac system in the sense of Section 3.1 and, therefore, in the sense of Section 2.1 in [BBC2].

4. BOUNDARY VALUES AND FREDHOLM PROPERTIES

Let $\mathcal{D} = (\mathcal{H}, \mathcal{A}, T)$ be a Dirac system over

$$(4.1) \quad I = \mathbb{R}_+ := [0, \infty).$$

with origin $t_0 = 0$, where we note that an analogous discussion holds true for other intervals with non-empty boundary. By (3.5), the restriction $D_{0,c}$ of the Dirac operator D to

$$(4.2) \quad \mathcal{L}_{0,c}(\mathcal{D}) := \{\sigma \in \mathcal{L}_c(\mathcal{D}) : \sigma(0) = 0\}$$

is symmetric. The adjoint operator of $D_{0,c}$ with respect to $L^2(\mathcal{D}) \supseteq \mathcal{L}_{0,c}(\mathcal{D})$ is called the *maximal extension* of D on $\mathcal{L}_c(\mathcal{D})$. We denote it by D_{\max} and let $\text{dom } D_{\max}$ be the domain of D_{\max} , endowed with the graph norm of D_{\max} . The adjoint operator D_{\min} of D_{\max} is equal to the closure of D on $\mathcal{L}_c(\mathcal{D})$. It is called the *minimal extension* of D , and its

domain is denoted by $\text{dom } D_{\min}$. We also let $H^1(\mathcal{D})$ be the completion of $\mathcal{L}_c(\mathcal{D})$ with respect to the norm

$$(4.3) \quad \|\sigma\|_{H^1(\mathcal{D})}^2 := \|\sigma\|_{L^2(\mathcal{D})}^2 + \|\partial\sigma\|_{L^2(\mathcal{D})}^2 + \|A\sigma\|_{L^2(\mathcal{D})}^2.$$

Obviously,

$$(4.4) \quad \mathcal{L}_c(\mathcal{D}) \subseteq H^1(\mathcal{D}) \subseteq \text{dom } D_{\max} \subseteq L^2(\mathcal{D}).$$

To formulate the main results on $\text{dom } D_{\max}$ from [BBC2], we need to discuss boundary values of sections at $t = 0$. As for proofs of the corresponding assertions, we refer to the discussion in Chapters 1 and 2 of [BBC2] and, in particular, to Proposition 2.30 loc.cit.

4.1. Boundary Values. Recall the convention $H = H_0$. Recall also that A_0 is self-adjoint in H with domain H_A . It will be convenient, in this section, to denote elements of H by letters x, y and to call them vectors. Fix an orthonormal basis (x_i) of H which consists of eigenvectors of A_0 , $A_0 x_i = \lambda_i x_i$.

For $s \geq 0$, let $H^s = H^s(A_0) \subseteq H = H_0$ be the domain of $|A_0|^s$. Then $H^0 = H$, $H^1 = H_A$, and $H^\infty = H^\infty(A_0) := \bigcap_{s \geq 0} H^s$ is a dense subspace of H . For $s \in \mathbb{R}$, define an inner product $\langle \cdot, \cdot \rangle_s$ on H^∞ ,

$$(4.5) \quad \langle x, y \rangle_s := ((I + A_0^2)^{s/2} x, (I + A_0^2)^{s/2} y).$$

For $s \geq 0$, the norm $\|\cdot\|_s$ associated to $\langle \cdot, \cdot \rangle_s$ is equivalent to the graph norm of $|A_0|^s$, and H^s is equivalent to the completion of H^∞ with respect to $\|\cdot\|_s$. For $s < 0$, define $H^s = H^s(A_0)$ to be the completion of H^∞ with respect to $\|\cdot\|_s$ and set $H^{-\infty} = H^{-\infty}(A_0) := \bigcup_{s \in \mathbb{R}} H^s$. In terms of the above basis (x_i) of eigenvectors, H^s consists of all linear combinations $x = \sum \xi_i x_i$ with

$$(4.6) \quad \sum (1 + \lambda_i^2)^s |\xi_i|^2 < \infty.$$

The pairing

$$(4.7) \quad B_s : H^s \times H^{-s} \rightarrow \mathbb{C}, \quad B_s(x, y) := ((I + A_0^2)^{s/2} x, (I + A_0^2)^{-s/2} y),$$

is perfect, that is, identifies H^{-s} with the dual space of H^s .

For a subset $J \subset \mathbb{R}$, let $Q_J = Q_J(A_0)$ be the corresponding spectral projection of A_0 in the spaces H^s . The image of H^s under Q_J is

$$(4.8) \quad H_J^s = H_J^s(A_0) := \{x = \sum \xi_i x_i \in H^s : \xi_i = 0 \text{ if } \lambda_i \notin J\}.$$

For $x \in H^s$, we also let $x_J := Q_J x$. For any bounded subset J of \mathbb{R} , we have $H_J^s \subseteq H^\infty$. Since $T = T_0$ anti-commutes with A_0 ,

$$(4.9) \quad TQ_J = Q_{-J}T \quad \text{and} \quad TH_J^s = H_{-J}^s$$

As shorthand, we use, for $a \in \mathbb{R}$,

$$(4.10) \quad Q_{>a} := Q_{(a,\infty)}, \quad Q_{\geq a} := Q_{[a,\infty)},$$

$$(4.11) \quad Q_{<a} := Q_{(-\infty,a)}, \quad Q_{\leq a} := Q_{(-\infty,a]},$$

and similarly for the spaces $H_J^s = Q_J(H^s)$. We also need to introduce the hybrid Sobolev space

$$(4.12) \quad \check{H} = \check{H}(A_0) := H_{\leq 0}^{1/2} \oplus H_{> 0}^{-1/2}.$$

Since $H_J \subseteq H^\infty$, for any bounded $J \subseteq \mathbb{R}$,

$$(4.13) \quad \check{H} = H_{\leq a}^{1/2} \oplus H_{> a}^{-1/2} = H_{< a}^{1/2} \oplus H_{\geq a}^{-1/2},$$

for any $a \in \mathbb{R}$. By (4.7) and (4.9),

$$(4.14) \quad \omega(x, y) := B_{1/2}(x_{\leq -a}, Ty_{\geq a}) + B_{-1/2}(x_{> -a}, Ty_{< a})$$

is well defined for $x, y \in \check{H}$ and independent of the choice of a . We note that ω is continuous, non-degenerate, and skew-Hermitian on \check{H} .

PROPOSITION 4.15. *The maximal domain $\text{dom } D_{\max}$ satisfies:*

- (1) $\mathcal{L}_c(\mathcal{D})$ is dense in $\text{dom } D_{\max}$.
- (2) Evaluation at $t = 0$ on $\mathcal{L}_c(\mathcal{D})$ induces a continuous surjection

$$\mathcal{R}_{\max} : \text{dom } D_{\max} \rightarrow \check{H}, \quad \mathcal{R}_{\max}(\sigma) =: \sigma(0).$$

- (3) $\sigma \in \text{dom } D_{\max}$ is in $H_{\text{loc}}^1(\mathcal{D})$ iff $\sigma(0) \in H^{1/2}$.
- (4) $\sigma \in \text{dom } D_{\max}$ is in $\text{dom } D_{\min}$ iff $\sigma(0) = 0$.
- (5) For all $\sigma_1, \sigma_2 \in \text{dom } D_{\max}$

$$(D_{\max}\sigma_1, \sigma_2)_{L^2(\mathcal{D})} - (\sigma_1, D_{\max}\sigma_2)_{L^2(\mathcal{D})} = \omega(\sigma_1(0), \sigma_2(0)).$$

Closed extensions of D between D_{\min} and D_{\max} correspond precisely to closed linear subspaces B of \check{H} , called *boundary conditions*. For any such boundary condition B , the domain of the corresponding extension $D_{B, \max}$ is given by

$$(4.16) \quad \text{dom } D_{B, \max} = \{\sigma \in \text{dom } D_{\max} : \sigma(0) \in B\}.$$

We are also interested in the operator D_B with domain

$$(4.17) \quad \text{dom } D_B = \text{dom } D_{B, \max} \cap H_{\text{loc}}^1(\mathcal{D}).$$

A boundary condition $B \subseteq \check{H}$ is called *regular* if $D_B = D_{B, \max}$. By Proposition 4.15, $\sigma \in \text{dom } D_{\max}$ is in $\text{dom } D_B$ if and only if $\sigma(0)$ belongs to $B \cap H^{1/2}$. In particular, B is a regular boundary condition if B is a closed subspace of \check{H} that is contained in $H^{1/2} \subseteq \check{H}$.

Let $B \subseteq \check{H}$ be a boundary condition. Since ω is non-degenerate, the adjoint operator of $D_{B, \max}$ is given by $D_{B^a, \max}$, where

$$(4.18) \quad B^a = \{x \in \check{H} : \omega(x, y) = 0 \text{ for all } y \in B\},$$

by Proposition 4.15. We say that a boundary condition B is *elliptic* if B and B^a are regular. Typical examples of elliptic boundary conditions are the Atiyah-Patodi-Singer boundary condition $B_{\text{APS}} = H_{<0}^{1/2}$ and the more general $B = H_{<a}^{1/2}$ and $B = H_{\leq a}^{1/2}$. The adjoint boundary conditions for the latter are given by $B = H_{\leq -a}^{1/2}$ and $B = H_{<-a}^{1/2}$, respectively. The maximal operators corresponding to the boundary conditions $B = H_{<a}^{1/2}$ and $B = H_{\leq a}^{1/2}$ will be denoted by $D_{<a,\max}$ and $D_{\leq a,\max}$, respectively, and similarly in other cases. By ellipticity, we actually have $D_{<a,\max} = D_{<a}$ and $D_{\leq a,\max} = D_{\leq a}$.

As for boundary conditions in the super-symmetric case,

$$(4.19) \quad H = H^+ \oplus H^-,$$

we may choose orthonormal bases x_i^\pm of H^\pm consisting of eigenvectors of A_0^\pm . By (3.18), we may actually choose $x_i^- = T_0 x_i^+ T_0^{-1}$. We get

$$(4.20) \quad H^s = H^{s+} \oplus H^{s-} \quad \text{and} \quad \check{H} = \check{H}^+ \oplus \check{H}^-,$$

where

$$(4.21) \quad H^{s+} = H^s(A_0^+), \quad H^{s-} = H^s(A_0^-), \quad \check{H}^+ = \check{H}(A_0^+),$$

and

$$(4.22) \quad \begin{aligned} \check{H}^- &= \check{H}(A_0^-) = T_0 \check{H}(-A_0^+) T_0^{-1} \\ &\simeq \hat{H}^+(A_0^+) = H_{\leq 0}^{-1/2} \oplus H_{>0}^{1/2}. \end{aligned}$$

Furthermore, \check{H}^+ and \check{H}^- are Lagrangian with respect to the non-degenerate skew-Hermitian form ω .

In the super-symmetric case, we consider super-symmetric boundary conditions $B \subseteq \check{H}$, that is,

$$(4.23) \quad B = B^+ \oplus B^-,$$

where $B^\pm = B \cap \check{H}^\pm$. Then the adjoint boundary condition is super-symmetric as well. Moreover, a super-symmetric boundary condition B is regular or elliptic if and only if B^+ and B^- are regular or elliptic in \check{H}^+ and \check{H}^- , respectively. For example, $B = H_{<a}^{1/2}$ and $B = H_{\leq a}^{1/2}$ are elliptic super-symmetric boundary conditions. The maximal operators corresponding to these will be denoted by $D_{<a,\max}^\pm$ and $D_{\leq a,\max}^\pm$, respectively, and similarly in other cases.

4.2. More Function Spaces. For convenience, we assume from now on that \mathcal{D} is the Dirac system associated to a Dirac bundle E over a straight end U of M with distance function f and C^1 diffeomorphism $F : (-r, \infty) \times N \rightarrow U$ as in Definition 1.8. We also recall the notation $U_0 = f^{-1}([0, \infty))$.

Let $H^1(U_0, E)$ be the space of sections σ in $L^2(U_0, E)$ with square integrable weak derivative, $\nabla^E \sigma \in L^2(U_0, T^*M \otimes E)$; that is, we have

$$(4.24) \quad (\nabla^E \sigma, \tau)_{L^2(U_0, E)} = (\sigma, (\nabla^E)^* \tau)_{L^2(U_0, E)},$$

for all $\tau \in C_{cc}^\infty(U_0, T^*M \otimes E)$, where $(\nabla^E)^*$ is the formal adjoint of the operator ∇^E . Recall that $H^1(U_0, E)$ is a Hilbert space with respect to the norm defined by the inner product

$$(4.25) \quad (\sigma, \tau)_{H^1(U_0, E)} = (\sigma, \tau)_{L^2(U_0, E)} + (\nabla^E \sigma, \nabla^E \tau)_{L^2(U_0, T^*M \otimes E)}.$$

There is the corresponding space $H^1(U, E)$, and $C_c^\infty(U, E)$ is dense in $H^1(U, E)$. Any section in $H^1(U_0, E)$ is the restriction of some section from $H^1(U, E)$; see Theorem 11.12 in [Ag] or Theorem 7.25 in [GiTr], noting that the problem is local and that H_{loc}^1 is invariant under C^1 diffeomorphisms. It follows that the space $C_c^\infty(U_0, E)$ of restrictions of sections in $C_c^\infty(U, E)$ to U_0 is dense in $H^1(U_0, E)$. The trace map

$$(4.26) \quad \mathcal{R} : H^1(U_0, E) \rightarrow H^{1/2}(N, E)$$

is a well defined bounded operator; see Theorem 3.10 in [Ag] or Proposition 4.4.5 in [Ta], noting again that the problem is local and that H_{loc}^1 is invariant under C^1 diffeomorphisms. The closure of $C_{cc}^1(U_0, E)$ in $H^1(U_0, E)$, and therefore also of $C_{cc}^\infty(U_0, E)$ in $H^1(U_0, E)$, is

$$(4.27) \quad H_0^1(U_0, E) := \{\sigma \in H^1(U_0, E) : \mathcal{R}\sigma = 0\}.$$

As for partial integration,

$$(4.28) \quad (\nabla^E \sigma, \tau)_{L^2(U_0, T^*M \otimes E)} = (\sigma, (\nabla^E)^* \tau)_{L^2(U_0, E)} - (\sigma, \tau(T))_{L^2(N, E)},$$

for all $\sigma \in H^1(U_0, E)$ and $\tau \in H^1(U_0, T^*M \otimes E)$. It follows that

$$(4.29) \quad (D\sigma, \tau)_{L^2(U_0, E)} = (\sigma, D\tau)_{L^2(U_0, E)} + (\sigma, T\tau)_{L^2(N, E)},$$

for all $\sigma, \tau \in H^1(U_0, E)$. In particular, any $\sigma \in H^1(U_0, E)$ belongs to the domain $\text{dom } D_{\text{max}}$ of the adjoint operator D_{max} of D , the latter considered as an unbounded operator on $L^2(U_0, E)$ with domain $C_{cc}^\infty(U_0, E)$ or, equivalently, $H_0^1(U_0, E)$.

We switch now to the associated Dirac system \mathcal{D} over $\mathbb{R}_+ = [0, \infty)$. With respect to the natural identifications,

$$(4.30) \quad C_c^\infty(U_0, E) \subseteq \mathcal{L}_c(\mathcal{D}) \subseteq H^1(\mathcal{D}) = H^1(U_0, E),$$

where we use (3.43) and (2.16) for the latter equality. The convenience we had in mind further up refers to the density of $C_c^\infty(U_0, E)$ in $H^1(\mathcal{D})$. Another convenience: We often write $\|\cdot\|_I$ for the L^2 -norm of maps defined on an interval I (if meaningful).

The Jacobian determinants j are not C^1 in general, hence the commutator $[A, \partial]$ may not be well defined. In Proposition 4.31 below and,

in particular, in its proof, we circumvent this problem by using the commutator A' of d/dt and A instead.

PROPOSITION 4.31. *For all $w \in \mathbb{R}$ and $\sigma \in H^1(\mathcal{D})$,*

$$\begin{aligned} \|D\sigma - wT\sigma\|_{\mathbb{R}_+}^2 &= \|\partial\sigma\|_{\mathbb{R}_+}^2 + \|(A - w)\sigma\|_{\mathbb{R}_+}^2 \\ &\quad - \operatorname{Re}(A'\sigma, \sigma)_{\mathbb{R}_+} - (\sigma(0), (A_0 - w)\sigma(0))_0. \end{aligned}$$

REMARK 4.32. As for the meaning of the last term on the right, we note that the trace $\sigma(0)$ of σ is in $H^{1/2}(A_0)$, hence A_0 applied to it is in $H^{-1/2}(A_0)$, and hence $(\sigma(0), (A_0 - w)\sigma(0))_0$ is well defined.

Proof of Proposition 4.31. Replacing A by $A_w := A - w$ and D by $D_w = T(\partial + A_w)$ reduces the assertion to the case $w = 0$. Furthermore, by the density of $C_c^\infty(U_0, E)$ in $H^1(U_0, E)$, we may assume that σ is smooth with compact support. We have

$$\|D\sigma\|_{\mathbb{R}_+}^2 = \|\partial\sigma\|_{\mathbb{R}_+}^2 + \|A\sigma\|_{\mathbb{R}_+}^2 + 2\operatorname{Re}(\partial\sigma, A\sigma)_{\mathbb{R}_+}$$

By (3.5) and (3.38),

$$\begin{aligned} (\partial\sigma, A\sigma)_{\mathbb{R}_+} &= \int_N \int_{\mathbb{R}_+} (\langle \sigma, A\sigma \rangle j)' dt dx - (\sigma, \partial A\sigma)_{\mathbb{R}_+} \\ &= -(\sigma(0), A_0\sigma(0))_0 - (\sigma, \partial A\sigma)_{\mathbb{R}_+}. \end{aligned}$$

Since $(A\sigma)' = A'\sigma + A\sigma'$, we conclude that

$$\begin{aligned} (\sigma, \partial A\sigma)_{\mathbb{R}_+} &= (\sigma, A'\sigma)_{\mathbb{R}_+} + (\sigma, A\sigma')_{\mathbb{R}_+} + \left(\frac{\kappa}{2}\sigma, A\sigma\right)_{\mathbb{R}_+} \\ &= (\sigma, A'\sigma)_{\mathbb{R}_+} + (A\sigma, \partial\sigma)_{\mathbb{R}_+} + i\operatorname{Im}(\kappa\sigma, A\sigma)_{\mathbb{R}_+}. \quad \square \end{aligned}$$

REMARK 4.33. It is possible to approximate a distance function f as in Definition 1.8 by smooth functions f_ε such that

$$|f - f_\varepsilon| + |\nabla f - \nabla f_\varepsilon| + |\nabla^2 f - \nabla^2 f_\varepsilon| < \varepsilon$$

uniformly on U_0 . However, the separation of variables formula, as we need it for our integration by parts formula in Proposition 4.31, would be more involved than the one derived there and would contain additional terms which would be difficult to control.

4.3. Fredholm Properties of \mathcal{D} . We say that \mathcal{D} is of *Fredholm type* if there is a constant $C > 0$ such that

$$(4.34) \quad \|\sigma\|_{\mathbb{R}_+} \leq C\|D\sigma\|_{\mathbb{R}_+}, \quad \forall \sigma \in \mathcal{L}_{0,c}(\mathcal{D}),$$

and that \mathcal{D} is *non-parabolic* if, for each $t > 0$, there is a constant $C > 0$ such that

$$(4.35) \quad \|\sigma\|_{[0,t]} \leq C\|D\sigma\|_{\mathbb{R}_+}, \quad \forall \sigma \in \mathcal{L}_{0,c}(\mathcal{D}).$$

Obviously, if \mathcal{D} is of Fredholm type, then it is non-parabolic.

In Proposition 5.6 below (and the paragraph preceding it), we will make the connection to the property *non-parabolic at infinity* of Dirac operators as considered in Definition 1.3 in the introduction.

In Lemma 2.38 of [BBC2], we showed that \mathcal{D} is non-parabolic if and only if, for each $t > 0$, there is a constant $C > 0$ such that

$$(4.36) \quad \|\sigma\|_{[0,t]} \leq C(\|\sigma(0)\|_{\check{H}}^2 + \|D\sigma\|_{\mathbb{R}_+}^2)^{1/2} =: C\|\sigma\|_W,$$

for all $\sigma \in \mathcal{L}_c(\mathcal{D})$. If \mathcal{D} is non-parabolic, we let $W \subseteq L_{\text{loc}}^2(\mathcal{D})$ be the completion of $\mathcal{L}_c(\mathcal{D})$ with respect to the norm $\|\cdot\|_W$. As we will see below, the space W is well suited for issues concerning the closedness of the image of D , compare (4.38) and the text preceding it.

Assume now that \mathcal{D} is non-parabolic. Then, since $\|\cdot\|_W$ is weaker than the graph norm of D , Proposition 4.15.1 implies that $\text{dom } D_{\text{max}} \subseteq W$. Furthermore, equality holds if and only if \mathcal{D} is of Fredholm type. Moreover, if $\varphi : \mathbb{R}_+ \rightarrow \mathbb{C}$ is Lipschitz continuous with compact support and $\sigma \in W$, then $\varphi\sigma \in \text{dom } D_{\text{max}}$. In particular, the trace \mathcal{R} is well defined and continuous on W and takes values in $\check{H} = \check{H}(A_0)$.

By the non-parabolicity of \mathcal{D} and the definition of W , D extends to a bounded operator $D_{\text{ext}} : W \rightarrow L^2(\mathcal{D})$. For a boundary condition $B \subseteq \check{H}$, we define $D_{B,\text{ext}}$ to be the operator in W with target $L^2(\mathcal{D})$ and domain

$$(4.37) \quad \text{dom } D_{B,\text{ext}} = \{\sigma \in W : \sigma(0) \in B\}.$$

Obviously, $D_{B,\text{ext}}$ is closed and extends $D_{B,\text{max}}$, and $D_{B,\text{ext}} = D_{B,\text{max}}$ if and only if \mathcal{D} is of Fredholm type.

In Theorem 2.43 of [BBC2] we showed that, for \mathcal{D} non-parabolic and B regular, $D_{B,\text{ext}}$ has finite dimensional kernel and closed image with

$$(4.38) \quad (\text{im } D_{B,\text{ext}})^\perp = \ker D_{B^a,\text{max}}.$$

Thus, if \mathcal{D} is non-parabolic and B is elliptic, then $D_{B,\text{ext}}$ is a Fredholm operator and the L^2 -index

$$(4.39) \quad \text{ind}_{L^2} D_{B,\text{max}} := \dim \ker D_{B,\text{max}} - \dim \ker D_{B^a,\text{max}}$$

of $D_{B,\text{max}}$ is well defined and finite.

PROPOSITION 4.40. *Assume that, for some $a \geq 0$,*

$$(A_t\sigma, A_t\sigma)_t \geq \text{Re}(A'_t\sigma, \sigma)_t + a\|\sigma\|_t^2,$$

for all $t \geq 0$ and $\sigma \in H_A$. Then \mathcal{D} is non-parabolic and $D_{<0,\text{ext}}$ is an isomorphism. Moreover, if $a > 0$, then \mathcal{D} is of Fredholm type.

Proof. Recall the Hardy inequality,

$$(4.41) \quad \int_{\mathbb{R}_+} |\phi'|^2 \geq \int_{\mathbb{R}_+} \frac{|\phi|^2}{4t^2},$$

for any C^1 function ϕ on \mathbb{R}_+ with $\phi(0) = 0$. By Proposition 4.31,

$$\|D\sigma\|_{\mathbb{R}_+}^2 \geq \|\partial\sigma\|_{\mathbb{R}_+}^2 + a\|\sigma\|_{\mathbb{R}_+}^2,$$

for all $\sigma \in H_0^1(U_0, E)$. Applying (3.32) and (4.41) we get

$$\begin{aligned} \|\partial\sigma\|_{\mathbb{R}_+}^2 &= \int_N \int_0^\infty \|(j^{1/2}\sigma)'\|^2 dt dx \\ &\geq \int_N \int_0^\infty \frac{1}{4t^2} \|\sigma\|^2 j dt dx = \int_0^\infty \frac{1}{4t^2} \|\sigma\|_t^2 dt. \end{aligned}$$

It follows that \mathcal{D} is non-parabolic. Clearly, if $a > 0$, then \mathcal{D} is of Fredholm type.

Using the density of $\mathcal{L}_c(\mathcal{D})$ in W , Proposition 4.31 together with the assumed inequality implies that

$$\|\partial\sigma\|_{\mathbb{R}_+}^2 - (\sigma(0), A_0\sigma(0))_0 \leq \|D\sigma\|_{\mathbb{R}_+},$$

for any $\sigma \in W$. Hence $D\sigma = 0$ and $\sigma(0) \in \check{H}_{<0}$ implies that $\partial\sigma = 0$ and $\sigma(0) = 0$. That is, σ solves

$$(4.42) \quad \sigma' = -\frac{\kappa}{2}\sigma,$$

with $\sigma(0) = 0$, hence $\sigma = 0$, and therefore $\ker D_{<0,\text{ext}}$ is trivial.

Conversely, the cokernel of $D_{<0,\text{ext}}$ is isomorphic to $\ker D_{\leq 0,\text{max}}$, by what we said further up. Now the same argument as above shows that any $\sigma \in \ker D_{\leq 0,\text{max}}$ with $\sigma(0) \in \check{H}_{\leq 0}$ satisfies $\partial\sigma = 0$. It follows that σ solves (4.42) and hence, by (3.32), that

$$\sigma(t, x) = j^{-1/2}(t, x)\sigma(0, x),$$

for all $t \in \mathbb{R}_+$ and $x \in N$. Since the L^2 -norm of σ is finite, we conclude that $\sigma = 0$. Hence $\text{coker } D_{<0,\text{ext}}$ is trivial as well. \square

By Corollary 3.49,

$$(4.43) \quad c_0 := \sup_{t \in \mathbb{R}_+, \sigma \in H_A \setminus \{0\}} \frac{\|A'_t \sigma\|_t}{\|\sigma\|_t + \|A_t \sigma\|_t} \leq C(C_R, C_R^E, C_W, n) < \infty.$$

COROLLARY 4.44. *Suppose that $\text{spec } A_t \cap (-\lambda, \lambda) = \emptyset$, for all $t \in \mathbb{R}_+$, where*

$$2\lambda \geq c_0 + \sqrt{4c_0 + c_0^2}.$$

Then \mathcal{D} is non-parabolic and $D_{<0,\text{ext}}$ is an isomorphism. Moreover, if the inequality is strict, then \mathcal{D} is of Fredholm type.

Proof. Choose $a \geq 0$ with

$$2\lambda \geq c_0 + \sqrt{4c_0 + c_0^2 + 4a}.$$

Then we have, for all $t \geq 0$ and $\sigma \in H_A$,

$$\begin{aligned} \|A_t \sigma\|_t^2 - \operatorname{Re}(A'_t \sigma, \sigma)_t &\geq \|A_t \sigma\|_t^2 - c_0(\|A_t \sigma\|_t + \|\sigma\|_t)\|\sigma\|_t \\ &\geq (\lambda^2 - c_0 \lambda - c_0)\|\sigma\|_t^2 \geq a\|\sigma\|_t^2, \end{aligned}$$

by the definition of c_0 , and hence Proposition 4.40 applies. \square

The following estimate relates boundary conditions to Fredholm properties of D , as we will see further on.

LEMMA 4.45. *For all $\sigma \in H_c^1(U_0, E)$ and $w \in \mathbb{R}$,*

$$\begin{aligned} \|\partial \sigma\|_{\mathbb{R}_+}^2 + \frac{1}{2}\|(A - w)\sigma\|_{\mathbb{R}_+}^2 \\ \leq \|(D - wT)\sigma\|_{\mathbb{R}_+}^2 + c_1\|\sigma\|_{\mathbb{R}_+}^2 + (\sigma_0, (A_0 - w)\sigma_0)_0, \end{aligned}$$

where $2c_1 = c_0(c_0 + 2 + 2|w|)$.

Proof. By Proposition 4.31 and the definition of c_0 ,

$$\begin{aligned} \|\partial \sigma\|_{\mathbb{R}_+}^2 + \|(A - w)\sigma\|_{\mathbb{R}_+}^2 - \|(D - wT)\sigma\|_{\mathbb{R}_+}^2 - (\sigma_0, (A_0 - w)\sigma_0)_0 \\ = \operatorname{Re}(A' \sigma, \sigma)_{\mathbb{R}_+} \\ \leq c_0(\|A \sigma\|_{\mathbb{R}_+} + \|\sigma\|_{\mathbb{R}_+})\|\sigma\|_{\mathbb{R}_+} \\ \leq c_0(\|(A - w)\sigma\|_{\mathbb{R}_+} + (1 + |w|)\|\sigma\|_{\mathbb{R}_+})\|\sigma\|_{\mathbb{R}_+} \\ \leq \frac{1}{2}\|(A - w)\sigma\|_{\mathbb{R}_+}^2 + c_0\left(\frac{c_0}{2} + 1 + |w|\right)\|\sigma\|_{\mathbb{R}_+}^2. \end{aligned} \quad \square$$

PROPOSITION 4.46. *Assume that there are $\Lambda > \lambda \geq 0$ such that*

$$(4.47) \quad (\Lambda - \lambda)^2 > 4c_0(c_0 + 2 + \lambda + \Lambda) \quad \text{and} \quad \operatorname{spec} A_t \cap (\lambda, \Lambda) = \emptyset,$$

for all $t \in \mathbb{R}_+$. Suppose $w \in (\lambda, \Lambda)$ satisfies

$$(4.48) \quad |\lambda - w|^2, |\Lambda - w|^2 > 2c_1 = c_0(c_0 + 2 + 2w).$$

Then we have:

- (1) *If $\sigma \in e^{-wt}L^2(\mathcal{D})$ solves $D\sigma = 0$ in the sense of distributions with $\sigma(0) \in \check{H}_{<\Lambda}$, then $\sigma = 0$.*
- (2) *If $\sigma \in e^{wt}L^2(\mathcal{D})$ solves $D\sigma = 0$ in the sense of distributions with $\sigma(0) \in \check{H}_{<-\lambda}$, then $\sigma = 0$.*
- (3) *\mathcal{D} is non-parabolic and $D_{<-\lambda, \text{ext}}$ is injective.*

The first assumption in Proposition 4.46 implies that the set of w in (λ, Λ) satisfying the required inequalities is non-empty. We also recall from (3.9) that the spectrum of A_t , $t \in \mathbb{R}$, is symmetric about 0 so that the second assumption implies that $\text{spec } A_t$ has empty intersection with $-(\lambda, \Lambda)$ as well. We get that $H_{\leq \lambda}^s = H_{< \Lambda}^s$ and that $H_{< -\lambda}^s = H_{\leq -\Lambda}^s$, for all $s \in \mathbb{R}$.

Proof of Proposition 4.46. Let $\sigma \in H_c^1(U_0, E)$, $v \in \mathbb{R}$, and set $\tau = e^{vt}\sigma$. Then

$$\begin{aligned} & \|e^{vt}D\sigma\|_{\mathbb{R}_+}^2 + (\tau_0, (A_0 - v)\tau_0)_0 - \|\partial\tau\|_{\mathbb{R}_+}^2 \\ &= \|(D - vT)\tau\|_{\mathbb{R}_+}^2 + (\tau_0, (A_0 - v)\tau_0)_0 - \|\partial\tau\|_{\mathbb{R}_+}^2 \\ &\geq \frac{1}{2}\|(A - v)\tau\|_{\mathbb{R}_+}^2 - \frac{c_0}{2}(c_0 + 2 + 2|v|)\|\tau\|_{\mathbb{R}_+}^2, \end{aligned}$$

by Lemma 4.45. Suppose now that $w \in (\lambda, \Lambda)$ satisfies the required inequalities, and choose $\varepsilon > 0$ such that

$$(4.49) \quad |\Lambda - w|^2, |\lambda - w|^2 \geq c_0(c_0 + 2 + 2w) + 2\varepsilon.$$

Then, with $v = \pm w$, we continue the above computation and get

$$(4.50) \quad \|(D - vT)\tau\|_{\mathbb{R}_+}^2 + (\tau_0, (A_0 - v)\tau_0)_0 \geq \|\partial\tau\|_{\mathbb{R}_+}^2 + \varepsilon\|\tau\|_{\mathbb{R}_+}^2.$$

By the density of $H_c^1(U_0, E)$ in $\text{dom } D_{\max}$, any $\tau \in \text{dom } D_{\max}$ satisfies

$$(4.51) \quad \|(D - vT)\tau\|_{\mathbb{R}_+}^2 + (\tau_0, (A_0 - v)\tau_0)_0 \geq \varepsilon\|\tau\|_{\mathbb{R}_+}^2,$$

where $v = \pm w$ and ε are as above. Now with σ as in the first two assertions and $v = w$ and $v = -w$, respectively, $\tau = e^{vt}\sigma$ is in $\text{dom } D_{\max}$ and satisfies $D_{\max}\tau = vT\tau$. The boundary condition for σ implies that the boundary term in (4.51) is non-positive, hence $\tau = 0$, and hence $\sigma = 0$. This shows the two first assertions. As for the last assertion, we note that

$$\begin{aligned} & \|D\sigma\|_{\mathbb{R}_+}^2 + (\sigma_0, (A_0 + w)\sigma_0)_0 \\ (4.52) \quad & \geq \|e^{-wt}D\sigma\|_{\mathbb{R}_+}^2 + (\sigma_0, (A_0 + w)\sigma_0)_0 \\ & \geq \varepsilon\|e^{-wt}\sigma\|_{\mathbb{R}_+}^2, \end{aligned}$$

for any $\sigma \in H_c^1(U_0, E)$. □

For later purposes we note that the computations in the above proof also show that

$$(4.53) \quad \|e^{wt}D\sigma\|_{\mathbb{R}_+}^2 + (\sigma_0, (A_0 - w)\sigma_0)_0 \geq \varepsilon\|e^{wt}\sigma\|_{\mathbb{R}_+}^2,$$

for any $\sigma \in H_c^1(U_0, E)$, where $w \in (\lambda, \Lambda)$ and ε is as in (4.49).

Suppose now that the assumptions of Proposition 4.46 are satisfied and that $w \in (\lambda, \Lambda)$ satisfies the corresponding inequalities. Then

(4.52) and (4.53) lead us to consider the *weighted Lebesgue spaces* $L_{\pm w}^2(\mathcal{D}) := e^{\mp wt} L^2(\mathcal{D})$, with norm associated to the inner product

$$(4.54) \quad (\sigma, \tau)_{\pm w} := (e^{\pm wt} \sigma, e^{\pm wt} \tau)_{\mathbb{R}_+},$$

and the *weighted Sobolev spaces* $H_{w, < \mu}^1(\mathcal{D})$, the completions of $H_{< \mu, c}^1(\mathcal{D})$ with respect to the norms

$$(4.55) \quad \|\sigma\|_{H_w^1(\mathcal{D})} := \|\sigma\|_w + \|D\sigma\|_w.$$

COROLLARY 4.56. *If the assumptions of Proposition 4.46 hold and $w \in (\lambda, \Lambda)$ satisfies the corresponding inequalities, then the operators*

$$\begin{aligned} D_{w, < \Lambda} : H_{w, < \Lambda}^1(\mathcal{D}) &\rightarrow L_w^2(\mathcal{D}) \quad \text{and} \\ D_{-w, < -\lambda} : H_{-w, < -\lambda}^1(\mathcal{D}) &\rightarrow L_{-w}^2(\mathcal{D}) \end{aligned}$$

are adjoints of each other and isomorphisms.

We note that D_w on $L_w^2(\mathcal{D})$ is conjugate to the operator $D - wT$ on $L^2(\mathcal{D})$, and similarly for D_{-w} . Hence the operators $D_{\pm w}$ are Dirac-Schrödinger operators in the sense of [BBC2] (compare also Remark 2.27 of loc.cit.).

Proof of Corollary 4.56. The operators are adjoints of each other since $\check{H}_{< -\lambda} = \check{H}_{\leq -\Lambda}$, by the assumptions of Proposition 4.46. By (4.52) and (4.53), the images of the operators are closed. The first two assertions of Proposition 4.46 say that their kernels are trivial. By integration by parts as in (5) of Proposition 4.15, we see that $\sigma \in L_w^2(\mathcal{D})$ is in the orthogonal complement of $D(H_{w, < \Lambda}^1(\mathcal{D}))$ if $\tau := e^{2wt} \sigma$ solves $D\tau = 0$ weakly with $\tau(0) \in H_{\leq -\Lambda} = H_{< -\lambda}$. Now $\tau \in e^{wt} L^2(\mathcal{D})$, hence $\tau = 0$, by the second assertion of Proposition 4.46. This shows that the first operator is an isomorphism. The claim for the second follows in a similar fashion, using the first assertion of Proposition 4.46. \square

COROLLARY 4.57. *If the assumptions of Proposition 4.46 hold, then*

$$\text{ind } D_{< 0, \text{ext}} = \dim H_{[-\lambda, 0)} - \dim \ker D_{\leq \lambda, \text{max}}.$$

In the super-symmetric case,

$$\text{ind } D_{< 0, \text{ext}}^+ = \dim H_{[-\lambda, 0)}^+ - \dim \ker D_{\leq \lambda, \text{max}}^-.$$

Proof. By Theorem 3.14 of [BBC2], we have

$$\text{ind } D_{< 0, \text{ext}} = \text{ind } D_{< -\lambda, \text{ext}} + \dim H_{[-\lambda, 0)}.$$

By Proposition 4.46, $D_{< -\lambda, \text{ext}}$ is injective. On the other hand, the orthogonal complement of $\text{im } D_{< -\lambda, \text{ext}}$ is given by the space of σ in $L^2(\mathcal{D})$ with $D\sigma = 0$ and $\sigma(0) \in H_{\leq \lambda}$. This shows the first claim, and the proof of the second is similar. \square

5. DECOMPOSITION AND INDEX

We assume from now on that we have a decomposition $M = M_0 \cup U_0$, where M_0 and U_0 are domains in M such that M_0 is compact and connected. We will need later that there is a distance function which is defined on a neighborhood of U_0 . However, up to and including Theorem 5.13, we only assume that $N := M_0 \cap U_0 \neq \emptyset$ is a level surface of a C^2 distance function f which is defined in some open neighborhood of N in M .¹⁰ We assume that $T := \text{grad } f$ points into the direction of U_0 , set $A_0 := -D_N$ as in (3.45) and get the associated Sobolev spaces $H^s = H^s(A_0)$ as in Section 4.1.

LEMMA 5.1. *There is a constant $C > 1$ such that*

$$\|\sigma\|_{H^1(M_0, E)} \leq C(\|\sigma|_N\|_{H^{1/2}} + \|D\sigma\|_{L^2(M_0, E)}),$$

for all $\sigma \in H^1(M_0, E)$.

Proof. Let $\mathcal{R}_0 : H^1(M_0, E) \rightarrow H^{1/2}$ be restriction to N , $\mathcal{R}_0\sigma := \sigma|_N$, and $\mathcal{E}_0 : H^{1/2} \rightarrow H^1(M_0, E)$ be an extension operator (for extension operators, see e.g. (1.36) and Lemma 1.37 in [BBC2]). Since \mathcal{E}_0 and \mathcal{R}_0 are continuous and $H_0^1(M_0, E)$ is the kernel of \mathcal{R}_0 ,

$$H^1(M_0, E) \rightarrow H^{1/2} \times H_0^1(M_0, E), \quad \sigma \mapsto (\mathcal{R}_0\sigma, \sigma - \mathcal{E}_0\mathcal{R}_0\sigma),$$

is a continuous bijection and therefore an isomorphism of topological vector spaces. This reduces the discussion to the case where σ belongs to $H_0^1(M_0, E)$, and then $\sigma|_N = 0$.

If a constant as asserted would not exist, there would be a sequence (σ_n) in $H_0^1(M_0, E)$ such that $\|\sigma_n\|_{L^2(M_0, E)} = 1$ and $\|D\sigma_n\|_{L^2(M_0, E)} \rightarrow 0$. Extending σ_n by 0 to U_0 , we obtain an H^1 -section τ_n of E over M with support in the compact domain M_0 such that $\|\tau_n\|_{L^2(M_0, E)} = 1$. Using partial integration as in Proposition 4.15.5 (or, more precisely, as in Proposition 5.7.3 in [BBC2]), we get that $D\tau_n = D\sigma_n$ on M_0 and that $D\tau_n = 0$ on U_0 . In particular, $\|D\tau_n\|_{L^2(M, E)} \rightarrow 0$. Therefore, up to passing to a subsequence, (τ_n) converges in $L^2(M, E)$. Furthermore, the limit τ vanishes in U_0 , is of L^2 -norm 1, and solves $D\tau = 0$ weakly. Hence τ is smooth and $D\tau = 0$, by elliptic regularity. Since the interior of U_0 is non-empty, this is in contradiction to the unique continuation property for Dirac operators, see Theorem 8.2 in [BoWo]. \square

It will be convenient to write D_{M_0} for the restriction of D to M_0 , and similarly in corresponding cases.

Consider the manifold \tilde{M} which is the disjoint union of M_0 and U_0 , endowed with the Dirac bundle $\tilde{E} \rightarrow \tilde{M}$ induced by E . We want to

¹⁰This seems to be the right setup for general applications.

apply the results from Chapter 5 of [BBC2] to the Dirac operator \tilde{D} of \tilde{E} and, therefore, need to check whether the requirements of Axiom VI there are satisfied. The only requirement in that axiom which might be non-obvious is dealt with in the following lemma.

LEMMA 5.2. *Let $\chi : \mathbb{R} \rightarrow \mathbb{R}$ be a smooth function with compact support which is equal to 1 close to 0. Then $(1 - \chi \circ f)\sigma \in \text{dom } D_{U_0, \min}$ for all $\sigma \in \text{dom } D_{U_0, \max}$, and similarly for M_0 .*

Proof. We note first that $(1 - \chi \circ f)\sigma$ is a section in $\text{dom } D_{U_0, \max}$ which vanishes in a neighborhood of the boundary N of U_0 . Hence the extension $\tilde{\sigma}$ of $(1 - \chi \circ f)\sigma$ by 0 to M_0 is in $\text{dom } D_{\max}$. Now we have $\text{dom } D_{\max} = \text{dom } D_{\min}$, by Theorem II.5.7 of [LaMi]. Hence there is a sequence of smooth sections $\sigma_k \in C_c^\infty(M, E)$ such that $\sigma_k \rightarrow \tilde{\sigma}$ in $\text{dom } D_{\min}$. It follows that $(1 - \tilde{\chi} \circ f)\sigma_k \rightarrow (1 - \chi \circ f)\sigma$ in $\text{dom } D_{U_0, \min}$, where $\tilde{\chi} : \mathbb{R} \rightarrow \mathbb{R}$ is a smooth function with compact support such that $(1 - \tilde{\chi})(1 - \chi) = (1 - \chi)$. \square

Because T is the exterior normal to M_0 , the space of boundary values of the maximal extension $D_{M_0, \max}$ of D over M_0 is the hybrid Sobolev space

$$(5.3) \quad \hat{H} = H_{<0}^{-1/2} \oplus H_{\geq 0}^{1/2} = \check{H}(-A_0).$$

LEMMA 5.4. *For any $\lambda \geq 0$, we have*

$$\text{ind } D_{M_0, \geq -\lambda} = \frac{1}{2} \dim H_{[-\lambda, \lambda]}.$$

Proof. Since $D_{M_0, \geq -\lambda}$ is the adjoint operator of $D_{M_0, > \lambda}$, we have

$$\text{ind } D_{M_0, \geq -\lambda} = -\text{ind } D_{M_0, > \lambda}.$$

On the other hand,

$$\text{ind } D_{M_0, \geq -\lambda} - \text{ind } D_{M_0, > \lambda} = \dim H_{[-\lambda, \lambda]},$$

by Theorem 5.16 in [BBC2]. \square

The same argument applies to $D_{U_0, \leq \lambda, \max}$ if D is of Fredholm type.

Specifying the data in the definition of non-parabolicity of the third named author, compare [Ca1], we say that D is *non-parabolic* with respect to some subset $L \subseteq M$ if, for any relatively compact open subset $K \subseteq M$, there exists a constant $C = C(K, L)$ such that

$$(5.5) \quad \|\sigma\|_{L^2(K, E)} \leq C \|D\sigma\|_{L^2(M, E)},$$

for any smooth section σ of E with compact support such that $\sigma|_L = 0$. Obviously, if D is of Fredholm type, then D is non-parabolic with respect to any sufficiently large compact subset, and if D is non-parabolic

with respect to some subset, then also with respect to any larger subset. Furthermore, if M is connected, then D is non-parabolic with respect to any subset whose complement is relatively compact, by Lemma 5.1. Recalling Definition 1.3, we see that D is *non-parabolic at infinity* if D is non-parabolic with respect to some compact subset of M .

PROPOSITION 5.6. *Suppose that the ends of M are straight in the sense of Definition 1.8 and let \mathcal{D} be the Dirac system over \mathbb{R}_+ associated to D over U_0 as in Section 3.3. Then D is non-parabolic with respect to M_0 if and only if \mathcal{D} is non-parabolic in the sense of Section 4.3. In particular, if \mathcal{D} satisfies the assumptions of Proposition 4.46, then D is non-parabolic with respect to M_0 . \square*

Assume from now on that D is non-parabolic with respect to M_0 . Let $W(M, E)$ and $W(U_0, E)$ be the completion of $H_c^1(M, E)$ and $H_c^1(U_0, E)$ with respect to the norms associated to the inner products

$$(5.7) \quad \begin{aligned} (\sigma, \tau)_{W(M, E)} &:= (\sigma, \tau)_{H^1(M_0, E)} + (D\sigma, D\tau)_{L^2(U_0, E)}, \\ (\sigma, \tau)_{W(U_0, E)} &:= (\sigma, \tau)_{\tilde{H}} + (D\sigma, D\tau)_{L^2(U_0, E)}, \end{aligned}$$

respectively. We have

$$(5.8) \quad \begin{aligned} W(M, E) &= \{(\sigma, \tau) \in H^1(M_0, E) \oplus W(U_0, E) : \sigma|_N = \tau|_N\} \\ &\subseteq H_{\text{loc}}^1(M, E), \end{aligned}$$

since the transmission condition $\sigma|_N = \tau|_N$ is a regular boundary condition for the manifold \tilde{M} as above, see Example 1.85 in [BBC2]. By definition, D induces continuous operators

$$(5.9) \quad \begin{aligned} D_{\text{ext}} &: W(M, E) \rightarrow L^2(M, E), \\ D_{U_0, \text{ext}} &: W(U_0, E) \rightarrow L^2(U_0, E). \end{aligned}$$

We arrive at the following version of Théorème 0.3 of [Ca2].

THEOREM 5.10. *Suppose that D is non-parabolic with respect to M_0 . Then $D_{\text{ext}} : W(M, E) \rightarrow L^2(M, E)$ is a Fredholm operator with*

$$(\text{im } D_{\text{ext}})^\perp = \ker D_{\text{max}} = \{\sigma \in L^2(M, E) : D\sigma = 0 \text{ weakly}\}.$$

Proof. Theorem 5.12 in [BBC2] implies that the image of D_{ext} is closed and that $\ker D_{\text{ext}}$ is of finite dimension. The last claim follows from the density of $H_c^1(M, E)$ in $W(M, E)$ and since D is formally self-adjoint. Finally, since $\ker D_{\text{max}} \subseteq \ker D_{\text{ext}}$ and the latter is of finite dimension, D_{ext} is a Fredholm operator. \square

In the super-symmetric case $E = E^+ \oplus E^-$, we get operators

$$(5.11) \quad D_{\text{ext}}^\pm : W(M, E^\pm) \rightarrow L^2(M, E^\mp).$$

Since D_{ext} is a Fredholm operator, the operators D_{ext}^{\pm} are Fredholm operators as well and

$$(5.12) \quad \text{ind } D_{\text{ext}}^+ = \dim \ker D_{\text{ext}}^+ - \dim \ker D_{\text{ext}}^-,$$

by Theorem 5.10 (and since D is formally self-adjoint).

The transmission condition $\sigma|_N = \tau|_N$ as above is elliptic. Therefore it can be decoupled into separate boundary conditions for M_0 and U_0 , respectively, compare Theorems 3.24 and 5.12 in [BBC2]. This leads to the following index formulas.

THEOREM 5.13. *Suppose that D is non-parabolic with respect to M_0 . Then we have, for any $\lambda \geq 0$,*

$$\text{ind } D_{\text{ext}} = \frac{1}{2} \dim H_{[-\lambda, \lambda]} + \text{ind } D_{U_0, < -\lambda, \text{ext}}.$$

In the super-symmetric case,

$$\text{ind } D_{\text{ext}}^+ = \text{ind } D_{M_0, \geq 0}^+ + \dim H_{[-\lambda, 0]}^+ + \text{ind } D_{U_0, < -\lambda, \text{ext}}^+.$$

Proof. The assertions are immediate consequences of Theorems 3.24, 4.17, and 5.12 in [BBC2] and Lemma 5.4 above. \square

Suppose now that the ends of M are straight in the sense of Definition 1.8. We may then consider weighted Lebesgue and Sobolev spaces, following the discussion just before and in Corollary 4.56. For $w \in \mathbb{R}$, let $L_w^2(M, E)$ be the space of measurable sections of E which are square integrable over M with respect to the weight which is equal to 1 over M_0 and equal to e^{2wt} over U_0 . Endow $L_w^2(M, E)$ with the corresponding inner product

$$(5.14) \quad (\sigma, \tau)_{L_w^2(M, E)} := (\sigma, \tau)_{L^2(M_0, E)} + (e^{wt}\sigma, e^{wt}\tau)_{L^2(U_0, E)}.$$

Furthermore, let $H_w^1(M, E)$ be the completion of $H_c^1(M, E)$ with respect to the norm associated to the inner product

$$(5.15) \quad (\sigma, \tau)_{H_w^1(M, E)} := (\sigma, \tau)_{L_w^2(M, E)} + (D\sigma, D\tau)_{L_w^2(M, E)}.$$

Assume from now on that the assumptions of Proposition 4.46 are satisfied and that $w \in \mathbb{R}$ satisfies the corresponding inequalities. Then, by (4.52), (4.53), and Lemma 5.1, the $H_{\pm w}^1(M, E)$ -norm is equivalent to the norm

$$(5.16) \quad \|\sigma\|_{\pm w} := \|\sigma|_N\|_{H^{1/2}} + \|D\sigma\|_{L^2(M_0, E)} + \|e^{\pm wt}D\sigma\|_{L^2(U_0, E)}.$$

Thus, by restriction to M_0 and U_0 , respectively, $H_w^1(M, E)$ is isomorphic to the space of pairs (σ, τ) in $H^1(M_0, E) \oplus H_w^1(U_0, E)$ satisfying the transmission condition $\sigma|_N = \tau|_N$.

THEOREM 5.17. *Suppose that the Dirac system \mathcal{D} over \mathbb{R}_+ associated to E over U_0 satisfies the assumptions of Proposition 4.46 and that $w > 0$ satisfies the corresponding inequalities. Then*

$$D_{-w} : H_{-w}^1(M, E) \rightarrow L_{-w}^2(M, E)$$

is a Fredholm operator with index

$$\text{ind } D_{-w} = \frac{1}{2} \dim H_{[-\lambda, \lambda]}.$$

In the super-symmetric case,

$$\text{ind } D_{-w}^+ = \text{ind } D_{M_0, \geq 0}^+ + \dim H_{[-\lambda, 0]}^+.$$

Proof. By (4.52), D_{-w} as above is a Fredholm operator. We also note that $D_{U_0, -w}$ is conjugate to the operator $D_{U_0} + w \text{grad } f$, where f is the given distance function over U_0 . Hence the results of Section 3 in [BBC2] apply (compare also Remark 2.27 of loc.cit.) and show that the Calderón projections associated to L_{-w}^2 -solutions of the equation $D\sigma = 0$ over M_0 and U_0 are elliptic. Hence, by Theorems 3.24 and 5.12 in [BBC2], D as above has index

$$\text{ind } D_{-w} = \text{ind } D_{M_0, \geq -\lambda} + \text{ind } D_{U_0, -w, < -\lambda}.$$

By Corollary 4.56, $\text{ind } D_{U_0, -w, < -\lambda} = 0$, hence the formula for $\text{ind } D_{-w}$ follows from Lemma 5.4. In the super-symmetric case,

$$\begin{aligned} \text{ind } D_{-w}^+ &= \text{ind } D_{M_0, \geq -\lambda}^+ + \text{ind } D_{U_0, -w, < -\lambda}^+ \\ &= \text{ind } D_{M_0, \geq -\lambda}^+ = \text{ind } D_{M_0, \geq 0}^+ + \dim H_{[-\lambda, 0]}^+. \end{aligned} \quad \square$$

COROLLARY 5.18. *Under the assumptions of Theorem 5.17, we have*

$$\text{ind } D_{\text{ext}}^+ = \text{ind } D_{-w} - \dim \ker D_{U_0, \leq \lambda, \max}^-.$$

Proof. By Theorems 5.13 and 5.17, we have

$$(5.19) \quad \text{ind } D_{\text{ext}}^+ = \text{ind } D_{-w} + \text{ind } D_{U_0, < -\lambda, \text{ext}}^+.$$

By Proposition 4.46.3, $D_{U_0, < -\lambda, \text{ext}}^+$ is injective. Hence, by (4.38),

$$(5.20) \quad \text{ind } D_{U_0, < -\lambda, \text{ext}}^+ = -\dim \ker D_{U_0, \leq \lambda, \max}^+ \quad \square$$

In the case where the boundary $N = N_0$ of M_0 is smooth, Theorem 3.1 in Atiyah-Patodi-Singer [APS1] applies and gives

$$(5.21) \quad \begin{aligned} \text{ind } D_{M_0, \geq 0}^+ &= \int_{M_0} \omega_{D^+} + \int_{N_0} \tau_{D^+} \\ &\quad + \frac{1}{2} (\eta(A_0^+) + \dim \ker A_0^+), \end{aligned}$$

where ω_{D^+} is the index form and τ_{D^+} the transgression form. We remark that ω_{D^+} is a universal polynomial in the curvatures of M and E and that τ_{D^+} is a universal polynomial in the curvature of M and E and the second fundamental form of N ; compare [Gi1] and Section 3.10 in [Gi2]. Now we may approximate M_0 by smooth domains such that the second fundamental forms of their boundaries approximate the second fundamental form of N . Then the integrals of ω_{D^+} and τ_{D^+} over the approximating domains and their boundaries converge to the integral of the corresponding forms over M_0 and N_0 , respectively. On the other hand, the coefficients of A_0^+ are only C^1 in general, and therefore the η -invariant of A_0^+ may not be well defined. However, since the other terms on the right hand side of (5.21) are well defined, we may define $\eta(A_0^+)$ to be the number such that (5.21) holds. In [Hi], Michel Hilsaum defined η -invariants for Lipschitz manifolds in a similar way, and he showed that they enjoy many of the properties of “smooth” η -invariants. We do not pursue this issue any further since we apply the APS-formula only in the smooth case.

Assuming now that the ends of M are smooth, we may combine the index formula for D^+ in Theorems 5.13 and 5.17 with (5.21). To that end, we continue to assume that the assumptions of Proposition 4.46 are satisfied. Then the spectrum of A_t has two parts, the part consisting of eigenvalues of modulus at most λ and the part consisting of those of modulus at least Λ . Following a corresponding convention in [Lo2], we call the first the *low energy* and the second the *high energy* part and get the corresponding spectral projections and spaces,

$$(5.22) \quad P_t := Q_{[-\lambda, \lambda]}(A_t^+), \quad H_t^{\text{le}} := P_t(H_t), \quad A_t^{\text{le}} := A_t|_{H_t^{\text{le}}},$$

$$(5.23) \quad Q_t := I - P_t, \quad H_t^{\text{he}} := Q_t(H_t), \quad A_t^{\text{he}} := A_t|_{H_t^{\text{he}}},$$

where we note that $H_t = H_t^{\text{le}} \oplus H_t^{\text{he}}$ is an orthogonal decomposition which is invariant under A_t . In the super-symmetric case we get similar decompositions and call

$$(5.24) \quad \eta(A_t^{\text{le},+}) \quad \text{and} \quad \eta(A_t^{\text{he},+}),$$

the *low* and *high energy* η -invariant of A_t^+ , respectively. We have

$$(5.25) \quad \eta(A_t^+) = \eta(A_t^{\text{le},+}) + \eta(A_t^{\text{he},+}).$$

COROLLARY 5.26. *Assume that the ends of M are smooth and straight and that the Dirac system over \mathbb{R}_+ associated to E over U_0 satisfies*

the assumptions of Proposition 4.46. Then we have, in the supersymmetric case,

$$\operatorname{ind} D_{-w}^+ = \int_{M_0} \omega_{D^+} + \int_{N_0} \tau_{D^+} + \frac{1}{2} \left(\dim H_{[-\lambda, \lambda]}^+ + \eta(A_0^{\text{he}, +}) \right). \quad \square$$

Since $\operatorname{ind} D_{-w}^+$ does not change when replacing the parameter t along the ends by $t - t_0$, for any $t_0 > 0$, it follows that $\operatorname{ind} D_{U_0, < -\lambda, \text{ext}}^+$ is an asymptotic invariant of D (for λ as in Proposition 4.46). Compare also Corollary 4.57.

The formula in Corollary 5.26 can be used to define high energy η -invariants in the case where the ends of M are not smooth. We expect that these enjoy nice properties because the family of high energy operators A_t^{he} has no spectral flow.

We conclude this chapter by explaining the

Proof of Theorem 1.17. Since M has only finitely many ends, there is a decomposition $M = M_0 \cup U_0$, where M_0 and U_0 are domains in M such that M_0 is compact, such that the common boundary $N := M_0 \cap U_0$ of M_0 and U_0 is smooth, such that each connected component of N bounds exactly one connected component of U_0 , and such that the latter are in one to one correspondence with the ends of M .

For each connected component C of N , let A_C^+ be the restriction of A_0^+ to sections of E with support on C . Then A_0^+ is the direct sum of the A_C^+ over the connected components C of N . Hence

$$\eta(A_0^+) = \sum_C \eta(A_C^+) \quad \text{and} \quad \dim \ker A_0^+ = \sum_C \dim \ker A_C^+.$$

For the connected component \mathcal{C} of U_0 with $\partial \mathcal{C} = C$ we now set

$$\begin{aligned} \operatorname{Corr}(\mathcal{C}) := \operatorname{ind} D_{\mathcal{C}, < 0, \text{ext}}^+ - \int_C \omega_{D^+} + \int_C \tau_{D^+} \\ + \frac{1}{2} \left(\eta(A_C^+) + \dim \ker A_C^+ \right). \end{aligned}$$

Then, by Theorem 5.13 and (5.21),

$$\operatorname{ind} D_{\text{ext}}^+ = \int_M \omega_{D^+} + \sum_C \operatorname{Corr}(\mathcal{C}).$$

By Theorem 3.24 of [BBC2], the terms $\operatorname{Corr}(\mathcal{C})$ only depend on the ends of M and not on the chosen decomposition of M as above. \square

6. MANIFOLDS WITH ε -THIN ENDS

Let N be a closed and connected Riemannian manifold of dimension n . We say that N is ε -flat if

$$(6.1) \quad \sqrt{K} \operatorname{diam} N \leq \varepsilon,$$

where K is some upper bound of the modulus of the sectional curvature of N . By Gromov's theorem on almost flat manifolds, there is a constant $\varepsilon(n)$ such that N is an infra-nilmanifold if N is $\varepsilon(n)$ -flat [Gr]. In what follows we need some details from the proof of Gromov's theorem from [BuKa] and from Section 4 of Ruh's improvement of Gromov's theorem in [Ru]. The estimates which we assert below hold if $\varepsilon(n)$ is chosen sufficiently small. The arguments in the proofs of these assertions are elementary albeit intricate.

For any curve $c : [a, b] \rightarrow N$, denote by $L(c)$ the length of c and by $h(c)$ parallel translation along c . For orthogonal transformations A and B between equi-dimensional Euclidean spaces V and W , we follow [Ru] and let $d(A, B)$ be the maximal angle $\angle(Av, Bv)$, where v runs over non-zero vectors in V . This is a non-smooth Finsler metric on the space of all orthogonal transformations from V to W , invariant under precomposition and postcomposition by orthogonal transformations of V and W , respectively, with injectivity radius and diameter π .

We begin with results from Chapters 2 and 3 in [BuKa]. Normalize the Riemannian metric of N so that $\operatorname{diam} N = 1$, and assume, correspondingly, that $\sqrt{K} \leq \varepsilon(n)$. As in [Ru], let

$$(6.2) \quad w = 2 \cdot 14^{\dim \operatorname{SO}(n)} \quad \text{and} \quad \rho \geq 10^4 w.$$

Let x and y be points in N . Then, if c_0 and c_1 are geodesic segments from x to y of length $< \rho$ such that $h(c_0)$ and $h(c_1)$ are 10^{-1} -close, then $h(c_0)$ and $h(c_1)$ are actually 10^{-5} -close. The relation $h(c_0) \sim h(c_1)$ iff $h(c_0)$ and $h(c_1)$ are 10^{-1} -close is an equivalence relation among the holonomies of geodesic segments from x to y of length $< \rho$. For each such equivalence class of holonomies, there is a geodesic segment from x to y of length $< 2 \cdot 10^{-4} \rho$ such that its holonomy belongs to the given equivalence class.

Let c_0 and c_1 be geodesic loops at x such that $L(c_0) + L(c_1) < \rho$. Then there is a unique geodesic loop $c_0 * c_1$ at x of length $< \rho$ homotopic to the concatenation of c_0 and c_1 , and $h(c_0 * c_1)$ is 10^{-5} -close to $h(c_1) \circ h(c_0)$. This turns the set H of equivalence classes of holonomies along geodesic loops at x of length $< \rho$ into a group, and the order of H is at most w .

Next we explain Ruh's construction of a flat metric connection on N from [Ru]. Fix an orthonormal frame $F_0 : \mathbb{R}^n \rightarrow T_x N$ to identify $T_x N$

with \mathbb{R}^n . For each equivalence class $h \in H$ of holonomies along geodesic loops at x of length $< \rho$, let $b_0(h) \in O(T_x N) \simeq O(n)$ be its barycenter. This defines an almost homomorphism $b_0 : H \rightarrow O(n)$ in the sense of [GKR] and b_0 is 10^{-4} -close to a homomorphism $b : H \rightarrow O(n)$, by Theorem 3.8 of [GKR]. It follows that b is injective, and we use b to identify H with its image in $O(n)$.

Let c_0 be a geodesic segment from x to y of length $< \rho$. For each geodesic segment c of length $< \rho$ from x to y , there is precisely one $h \in H$ such that $h(c) \circ c_0$ is 10^{-4} -close to c_0 . Enrich the equivalence class of $h(c_0)$ as above by all such $h(c) \circ c_0$.

Choose a smooth monotone function $\chi : \mathbb{R} \rightarrow \mathbb{R}$ with $\chi(r) = 1$ for $r \leq \rho/3$, $\chi(r) = 0$ for $r \geq 2\rho/3$, and $|\chi'| \leq 10/\rho$. For any enriched equivalence class $[h(c) \circ c_0]$ of holonomies as above, let $b([h(c) \circ c_0])$ be its barycenter with respect to the weights $\chi(L(c))/\nu$, where ν is the order of H times the sum of the $\chi(L(c))$, over all geodesic segments c from x to y of length $< \rho$. By the equivariance of barycenters with respect to orthogonal transformations, the set of the barycenters $b([h(c) \circ c_0])$ is invariant under right multiplication by elements from H , and hence the frames $b \circ F_0$, where b runs over the above barycenters, define a reduction of the principal bundle of orthonormal frames of N to a principal subbundle with structure group H . In other words, we get a flat metric connection $\bar{\nabla}$ on N with holonomy in H .

To estimate the norm of the difference between $\bar{\nabla}$ and the Levi-Civita connection ∇ of N , we go back one step and consider the situation before taking barycenters. Let $v \in T_y N$ and $\sigma = \sigma(s)$ be a curve through y with s -derivative $\dot{\sigma}(0) = v$. Let $c_0 : [0, 1] \rightarrow N$ be a geodesic segment from x to y with $L(c_0) < \rho$. There is a unique geodesic variation $c = c_s(t)$ of c_0 with $c_s(0) = x$ and $c_s(1) = \sigma(s)$, and then $L(c_s) < \rho$ for all (sufficiently small) s . Let $u \in T_x N$ and $X = X(s, t)$ be the vector field along c such that $X(s, 0) = u$ and such that X is parallel along the segments c_s . Note that parallel translation along σ with respect to $\bar{\nabla}$ corresponds to taking barycenters of such $X(s, 1)$ along σ , arising from geodesic segments from x to y of length $< \rho$.

We have $\nabla_t \nabla_s X = R(c', J)X$, where the s -derivative $J := \dot{c}$ of c is a Jacobi field along each of the c_s which vanishes at $t = 0$ and is equal to $\dot{\sigma}(s)$ at $t = 1$. It follows that

$$(6.3) \quad |(\nabla_t \nabla_s X)(0, t)| \leq C_0 K \rho t |v| |X|,$$

where C_0 is a universal constant. Since taking barycenters depends smoothly on points and weights, we conclude that

$$(6.4) \quad |\bar{\nabla}_v X - \nabla_v X| \leq C_1 \left(K \rho + \frac{1}{\rho} \right) |v| |X|.$$

Now, for any given $\delta > 0$, we may choose ρ so large and, accordingly, $\varepsilon = \varepsilon(n, \delta)$ so small, that the right hand side of (6.4) is $< \delta|v||X|$. Hence, reversing the normalization of the diameter, we get that

$$(6.5) \quad |\bar{\nabla} - \nabla| \leq \delta \operatorname{diam} N,$$

where we recall that scaling does not change the Levi-Civita connection. This finishes the exposition of results from [BuKa] and [Ru].

Proof of Theorem 1.13. In the above constructions, it is understood, in the literature, that the Riemannian manifold N is smooth. We want to apply it in our situation of straight ends, where the Riemannian metric of the cross sections $N_t \subseteq U \simeq [0, \infty) \times N$ is, in general, only C^1 . To overcome this technical difficulty, we note that f can be approximated, locally uniformly in the C^2 topology, by a sequence of smooth functions $f_k : U \rightarrow \mathbb{R}$. Then, for any given cross section N_t , the level sets $L_k = f_k^{-1}(t)$ approximate N_t in the sense that there is a C^1 diffeomorphism between them such that Riemannian metric, Levi-Civita connection, and Weingarten operator on N_t are approximated by the corresponding objects on L_k . In particular, diameter and modulus of the sectional curvature of the connected components of the levels L_k are bounded from above by

$$d_t + \alpha \quad \text{and} \quad K = C_R + 2C_W^2 + \alpha,$$

for any given $\alpha > 0$ and all sufficiently large k , where d_t is an upper bound for the diameter of the connected components of N_t and where we use the Gauss equation for the second estimate. Thus the above constructions apply to N_t if

$$\sqrt{K}d_t \leq \varepsilon < \varepsilon(m-1, 1),$$

where $K = C_R + 2C_W^2 + 1$ and $\varepsilon(m-1, 1) = \varepsilon(n, \delta)$ is as in the discussion of (6.5) above, and they guarantee a flat connection $\bar{\nabla}^{N_t}$ on N_t such that

$$|\bar{\nabla}^{N_t} - \nabla^{N_t}| \leq d_t,$$

where ∇^{N_t} denotes the Levi-Civita connection of N_t .

Suppose now that $E \rightarrow M$ is a Dirac bundle of the type required in Theorem 1.13. Then the restrictions of E to any given cross section N_t is of the corresponding type and hence $\bar{\nabla}^{N_t}$ induces a flat Hermitian connection $\bar{\nabla}^{E_t}$ on the restriction $E_t = E|_{N_t}$. Moreover, the holonomy of $\bar{\nabla}^{E_t}$ on each connected component of N_t is of order at most w .

For convenience, assume now that N is connected. Decompose E_t into holonomy irreducible components, and let $F \rightarrow N_t$ be any such

component. Then F has a twisted parallel orthonormal frame

$$(6.6) \quad \Phi = (\sigma_1, \dots, \sigma_k),$$

that is, the sections σ_i of F are well defined and parallel on the induced bundle with induced flat connection over the universal covering of N_t . We think of them as sections of E over N_t which transform according to the holonomy of F . Approximating the Riemannian metric on N_t by a smooth ϵ -flat Riemannian metric as above, we see that we can apply the usual estimates for the Rayleigh quotient of sections of F , that is, the estimate of Li and Yau [LiYa] in the case where F is the trivial complex line bundle and the corresponding estimate in [BBC1] in the other cases. The outcome is an estimate as follows: If σ is a section of E over N_t and σ is orthogonal to the globally $\bar{\nabla}^{E_t}$ -parallel sections of E over N_t , then

$$(6.7) \quad \|\bar{\nabla}^{E_t} \sigma\|_{N_t}^2 \geq \frac{C(C_R, C_W, m)}{\epsilon^2} \|\sigma\|_{N_t}^2.$$

Here we use, in the twisted case, that the holonomy of F is non-trivial in the sense that, for each unit vector v in F , there is a loop c in N_t (of length at most ρ) such that the angle between v and hv is at least $\pi/2$, since otherwise the holonomy orbit of v would be contained in an open spherical ball of radius $\pi/2$ and would have a fixed point. Hence, for each unit vector v in F , there is a loop c in N_t of length at most $2d_t$ such that the angle between v and $h(v)$ is at least $\pi/2w$.

Now the estimate $|\bar{\nabla}^{N_t} - \nabla^{N_t}| \leq d_t$ implies that

$$|\bar{\nabla}^{N_t} - \nabla|_{N_t} \leq d_t + C_W,$$

where ∇ denotes the Levi-Civita connections of M . Hence

$$|\bar{\nabla}^{E_t} - \nabla^E|_{N_t} \leq C(d_t + C_W),$$

where C is a constant which depends only on the type of E . It follows that the difference between the Rayleigh quotients for $\nabla^E|_{N_t}$ and $\bar{\nabla}^{E_t}$ is uniformly bounded. We conclude that the assumptions of Proposition 4.46 are satisfied. \square

7. CUSPIDAL ENDS

Assume from now that the ends of M are cuspidal. In the setup of Definition 1.8 and of Section 3.2, denote by \mathcal{D} the Dirac system associated to E over U as in Section 3.3. Clearly, for any $\epsilon > 0$, the cross sections N_t are ϵ -flat for all sufficiently large t so that Theorem 1.13 applies. On the other hand, in this chapter, we aim at more specific results. In addition, we do not need to rely on the proof of Gromov's theorem on almost flat manifolds.

7.1. The Flat Connection. Over U , define a tensor field \bar{S} of bilinear maps on $TM \oplus TM$ with values in TM by

$$(7.1) \quad \langle \bar{S}(u, v), w \rangle = - \int_s^\infty \langle R(J, T)X, Y \rangle(t, x) dt,$$

where $u, v, w \in T_{(s,x)}M$, J is the T -Jacobi field along $\gamma_{(s,x)} := F(s, x)$ with $J(s) = u$, and X, Y are the parallel vector fields along $\gamma_{(s,x)}$ with $X(s) = v$, $Y(s) = w$. The integral converges uniformly, by (1.1) and since the ends are cuspidal. Hence \bar{S} is continuous and uniformly bounded. We let C_S be an upper bound for the operator norm of \bar{S} .

In the analogous way, define a field \bar{S}^E of bilinear maps on $TM \oplus E$ with values in E ,

$$(7.2) \quad \langle \bar{S}^E(u, v), w \rangle = - \int_s^\infty \langle R^E(J, T)\sigma_1, \sigma_2 \rangle(t, x) dt,$$

where now $v, w \in E_{(s,x)}$ and σ_1, σ_2 are the parallel sections along $\gamma_{(s,x)}$ with $\sigma(s) = v$, $\tau(s) = w$. Again, the integral converges uniformly, by (1.1) and since the ends are cuspidal. Hence \bar{S}^E is also continuous and uniformly bounded. We let C_S^E be an upper bound for the operator norm of \bar{S}^E .

The arguments in Section 3 of [BB2] carry over word by word and show that the continuous metric connections

$$(7.3) \quad \bar{\nabla} := \nabla - \bar{S} \quad \text{and} \quad \bar{\nabla}^E := \nabla^E - \bar{S}^E$$

on TM and E over U are flat in the sense of the existence of parallel C^1 frames over simply connected domains in U . The difference to the situation in Section 6 is that we do not assume that E is geometric and that we have to pay for it by making stronger assumptions on the smallness of the Riemannian metrics g_t and by loosing control on the holonomy of $\bar{\nabla}$ and $\bar{\nabla}^E$.

It is easy to see that

$$(7.4) \quad \bar{S}^E(X, Y\sigma) = \bar{S}(X, Y)\sigma + Y\bar{S}^E(X, \sigma),$$

hence the new connections are compatible with Clifford multiplication as well, that is,

$$(7.5) \quad \bar{\nabla}_X^E(Y\sigma) = (\bar{\nabla}_X Y)\sigma + Y\bar{\nabla}_X^E \sigma$$

By definition,

$$(7.6) \quad \bar{\nabla}_T = \nabla_T, \quad \bar{\nabla}_T^E = \nabla_T^E, \quad \text{and} \quad \bar{\nabla}T = 0.$$

For each $t \in \mathbb{R}_+$, the restriction of $\bar{\nabla}$ and \bar{S} to N_t will be denoted by $\bar{\nabla}_t$ and \bar{S}_t , and similarly for $\bar{\nabla}^E$ and \bar{S}^E . We also consider $\bar{\nabla}_t^E$ as a first

order differential operator on $H^1(N_t, E)$ with values in $L^2(T^*N_t \otimes E)$. The formal adjoint of $\bar{\nabla}_t^E$ is denoted $(\bar{\nabla}_t^E)^*$.

REMARK 7.7. The above construction of a flat connection is taken from [BB2] (where it is considered for a narrower class of bundles E). In Appendix C of [BeKa], Igor Belegradek and Vitali Kapovitch remark that this connection coincides with the flat connection introduced by Brian Bowditch in [Bo] (in the case of the tangent bundle of a simply connected, complete Riemannian manifold with pinched negative sectional curvature), who uses a kind of parallel translations through infinity (which, in turn, coincides with the *horospherical translations* in Section 2 of [BrKa]): Roughly speaking, two vectors $v, w \in E$ with footpoint on a common horosphere are defined to be parallel if the distance between their parallel translates along the unit speed geodesics to the center of the horosphere converges to 0.

7.2. The Splitting. To keep the notation simple, it will be convenient to assume in this section that N is connected. It will be obvious that, mutatis mutandis, the results also apply in the case where N is not connected.

For each $t \in \mathbb{R}_+$, we let H_t^c be the space of $\bar{\nabla}^E$ -parallel sections of E over N_t , that is, H_t^c is the kernel of $\bar{\nabla}_t^E$. Here the superscript c stands for *constant*. We note that the spaces H_t^c are invariant under Clifford multiplication by T , by (7.5) and (7.6). It is also clear that parallel translation in the T -direction identifies the different spaces H_t^c , $t \in \mathbb{R}_+$. In particular, we may and will fix a family of $\bar{\nabla}^E$ -parallel sections $(\sigma_1, \dots, \sigma_k)$ of E over U which are pointwise orthonormal and whose restriction to N_t forms an orthogonal basis of H_t^c , for all $t \in \mathbb{R}_+$ simultaneously.

We let H_t^h be the orthogonal complement of H_t^c in $L^2(N_t, E)$. Thus we obtain two families $\mathcal{H}^c = (H_t^c)$ and $\mathcal{H}^h = (H_t^h)$ of Hilbert spaces, both of them invariant under Clifford multiplication by T . Note, however, that \mathcal{H}^h is not parallel in the T direction if \mathcal{H}^c is non-trivial and the volume density $j = j(t, x)$ as in Section 3.2 does not only depend on t , but also on x , compare (7.9).

As before, we use parallel translation to identify the spaces H_t^c with H_0^c , endowed with the inner products $(\cdot, \cdot)_t = (j_t \cdot, \cdot)_0$. Since T is parallel in the T direction, Clifford multiplication by T does not depend on t after this identification.

Let \bar{P}_t and $\bar{Q}_t := I - \bar{P}_t$ be the orthogonal projections in H_t onto H_t^c and H_t^h , respectively. By definition,

$$(7.8) \quad \bar{P}_t \sigma = \frac{1}{\text{vol } N_t} \sum_{1 \leq i \leq k} (\sigma_i, \sigma)_t \sigma_i.$$

For any function $\psi = \psi(t, x)$ on U we denote by $\bar{\psi} = \bar{\psi}(t)$ the function which associates to $t \in \mathbb{R}_+$ the mean of ψ over the cross section N_t . By (3.32) and (7.8), we have

$$(7.9) \quad (\nabla_T \bar{P})\sigma = \bar{P}(\kappa\sigma) - \bar{\kappa}\bar{P}\sigma.$$

Associated to the projections \bar{P} and \bar{Q} , we consider the operators

$$(7.10) \quad D^c := \bar{P}D\bar{P}, \quad D^h := \bar{Q}D\bar{Q}, \quad D^{ch} := \bar{P}D\bar{Q}, \quad D^{hc} := \bar{Q}D\bar{P}.$$

We use corresponding notations and conventions in other cases.

PROPOSITION 7.11. *The family*

$$\mathcal{D}^c := (\mathcal{H}^c, \mathcal{A}^c, T)$$

is a Dirac system in the sense of Section 3.1 with finite rank and

$$\partial^c = \frac{d}{dt} + \frac{\bar{\kappa}}{2} \quad \text{and} \quad D^c = T(\partial^c + A^c).$$

Proof. The sections $\sigma_1, \dots, \sigma_k$ as above are C^1 , so that the image H_t^c of \bar{P}_t consists of C^1 sections of E over N_t . Hence H_t^c is contained in H_A , for all $t \in \mathbb{R}_+$. Furthermore, $A_t^c = \bar{P}_t A_t \bar{P}_t$ is a bounded and symmetric operator on H_t^c . Clearly, for $\sigma_1, \sigma_2 \in H_0^c$,

$$|(\bar{P}_t A_t \bar{P}_t \sigma_1, \sigma_2)_t - (\bar{P}_s A_s \bar{P}_s \sigma_1, \sigma_2)_s| = |(A_t \sigma_1, \sigma_2)_t - (A_s \sigma_1, \sigma_2)_s|. \quad \square$$

Associated to the decomposition into constant sections and sections perpendicular to them, we get an orthogonal splitting

$$(7.12) \quad L^2(\mathcal{D}) = L^2(\mathcal{H}) = L^{2,c}(\mathcal{H}) \oplus L^{2,h}(\mathcal{H}).$$

where

$$(7.13) \quad \begin{aligned} L^{2,c}(\mathcal{H}) &:= L^2(\mathcal{H}^c) \quad \text{and} \\ L^{2,h}(\mathcal{H}) &:= \{\sigma \in L^2(\mathcal{H}) : \bar{P}\sigma = 0\}. \end{aligned}$$

We use corresponding notations for other spaces of sections.

LEMMA 7.14. *The projections \bar{P} and \bar{Q} are continuous on $H^1(\mathcal{D})$. In particular, as topological vector spaces,*

$$\begin{aligned} H^1(\mathcal{D}) &= H^{1,c}(\mathcal{D}) \oplus H^{1,h}(\mathcal{D}), \\ H_{\text{loc}}^1(\mathcal{D}) &= H_{\text{loc}}^{1,c}(\mathcal{D}) \oplus H_{\text{loc}}^{1,h}(\mathcal{D}). \end{aligned}$$

Proof. Since $\sigma_1, \dots, \sigma_k$ and $\text{vol } N_t$ are C^1 , we conclude that

$$\bar{P}(H^1(\mathcal{D})) \subseteq H^1(\mathcal{D}) \quad \text{and} \quad \bar{Q}(H^1(\mathcal{D})) \subseteq H^1(\mathcal{D}),$$

by (7.8). Hence \bar{P} and $\bar{Q} = I - \bar{P}$ are continuous with respect to the H^1 -norm, by the closed graph theorem. \square

LEMMA 7.15. *The Rayleigh quotients*

$$(1) \quad \bar{\rho}_t := \inf\{\|\bar{\nabla}_t^E \sigma\|_t^2 / \|\sigma\|_t^2 : \sigma \in H_t^h \cap H_A, \sigma \neq 0\},$$

$$(2) \quad \rho_t := \inf\{\|\nabla_t^E \sigma\|_t^2 / \|\sigma\|_t^2 : \sigma \in H_t^h \cap H_A, \sigma \neq 0\}$$

tend to infinity as t tends to infinity. Here ∇_t^E and $\bar{\nabla}_t^E$ denote the restrictions of ∇^E and $\bar{\nabla}^E$ to N_t .

Proof. We discuss the Rayleigh quotients associated to $\bar{\nabla}^E$ first. Split $H_t^h \cap H_A = U_t \oplus V_t$, where U_t consists of sections in $H^1(N_t, E)$ which are linear combinations $\sum \varphi_i \sigma_i$ of the basis $(\sigma_1, \dots, \sigma_k)$ as above and where V_t consists of sections in $H^1(N_t, E)$ which are pointwise perpendicular to $\sigma_1, \dots, \sigma_k$. Note that U_t and V_t are invariant under $\bar{\nabla}_t^E$ and perpendicular to each other, and thus it suffices to consider them separately.

Let $\sigma = \sum \varphi_i \sigma_i \in U_t$, $\sigma \neq 0$. To be perpendicular to H_t^c in $L^2(N_t, E)$ means that the coefficient functions φ_i integrate to 0. Moreover, the Rayleigh quotient of σ is given by the sum of the Rayleigh quotients corresponding to the Laplace operator on functions on N_t . Hence

$$\frac{\|\bar{\nabla}_t^E \sigma\|_t^2}{\|\sigma\|_t^2} = \frac{\sum \|\text{grad } \varphi_i\|_t^2}{\sum \|\varphi_i\|_t^2} \geq ce^{ct},$$

for some constant $c > 0$, by Theorem 7 in [LiYa].

Now we consider V_t . Perpendicular to $(\sigma_1, \dots, \sigma_k)$, the holonomy of $\bar{\nabla}$ does not have non-trivial invariant vectors. Since loops in $N = N_0$ of length at most $2 \text{diam } N$ generate the fundamental group of N , there is a constant $\alpha > 0$ such that, for each vector u in some fiber of E over N , there is a loop c in N of length at most $2 \text{diam } N$ such that the holonomy h_c of $\bar{\nabla}$ along c satisfies $|h_c u - u| \geq \alpha|u|$. For each $t \geq 0$, the $\bar{\nabla}$ -holonomy about the curve c shifted to N_t is the same. We conclude that, for each $t \in \mathbb{R}_+$ and vector u in some fiber of E over N_t , there is a loop c in N_t of length at most $2\varphi(t) \text{diam } N$ such that the holonomy h_c of $\bar{\nabla}$ along c satisfies the same inequality,

$$|h_c u - u| \geq \alpha|u|.$$

Hence Theorem 5 in [BBC1] applies and shows that the Rayleigh quotient of $\bar{\nabla}_t^E$ on V_t tends uniformly to infinity as t tends to infinity. This shows the first claim. As for the Rayleigh quotients associated to ∇_t^E , we recall that the difference $\|\bar{\nabla}_t^E - \nabla_t^E\| \leq C_S^E$. \square

THEOREM 7.16. *There are constants $\lambda_0, \Lambda_t \geq 0$ with $\lim_{t \rightarrow \infty} \Lambda_t = \infty$ such that $\text{spec } A_t \cap (\lambda_0, \Lambda_t) = \emptyset$ or, more precisely, such that*

- (1) $\|D_t \sigma\|_t \leq \lambda_0 \|\sigma\|_t$ for all $\sigma \in H_t^c$,
- (2) $\|D_t \sigma\|_t \geq \Lambda_t \|\sigma\|_t$ for all $\sigma \in H_t^h$.

In particular, for all sufficiently large t ,

- (1) \mathcal{D}_t satisfies the hypothesis of Proposition 4.46,
- (2) D is non-parabolic with respect to $M_t := M \setminus \{f > t\}$,

where f is the distance function as in Definition 1.8.

Proof. By (3.43) and (2.16),

$$|\|D_t \sigma\|_t^2 - \|\nabla^t \sigma\|_t^2| \leq C_K \|\sigma\|_t^2. \quad \square$$

Assume now that the ends of M are smooth, that is, the associated distance function f on U is smooth. Since the ends of M are cuspidal and the curvatures of M and E and the second fundamental forms of the cross sections are uniformly bounded, we have, with $M_t = M \setminus \{f > t\}$ as above,

$$(7.17) \quad \lim_{t \rightarrow \infty} \int_{M_t} \omega_{D^+} = \int_M \omega_{D^+} \quad \text{and} \quad \lim_{t \rightarrow \infty} \int_{N_t} \tau_{D^+} = 0,$$

compare (5.21). By Theorem 7.16, we may fix the starting time $t = 0$ such that the condition

$$(7.18) \quad (\Lambda_t - \lambda_0)^2 > 4c_0(c_0 + 2 + \lambda_0 + \Lambda_t)$$

of Proposition 4.46 is satisfied for all $t \in \mathbb{R}_+$, where λ and Λ there correspond to λ_0 and Λ_t here.

PROPOSITION 7.19. *If $w > 0$ satisfies $(w - \lambda_0)^2 > c_0(c_0 + 2 + 2w)$, then*

$$\text{ind } D_{-w}^+ = \int_M \omega_{D^+} + \frac{1}{2} (\dim H_{[-\lambda_0, \lambda_0]}^+(A_0^+) + \lim_{t \rightarrow \infty} \eta(A_t^{\text{he}, +})).$$

Proof. Observing that $\dim H_{[-\lambda_0, \lambda_0]}^+(A_t)$ is independent of $t \in \mathbb{R}_+$, the assertion follows immediately from Corollary 5.26 and (7.17). \square

As for the index of D_{ext}^+ , we refer to Corollary 5.18.

8. EXPLICIT INDEX FORMULAS

It is certainly desirable to get more explicit index formulas than the one in Proposition 7.19. To that end one might hope that the decomposition into low and high energy parts along the cusps should lead to a decomposition of the corresponding Dirac system \mathcal{D} into low and high energy Dirac systems. Whereas the low energy part leads to the

Dirac system \mathcal{D}^c as in Proposition 7.11, the high energy part does not seem to define a Dirac system in general. If it did, we would want to apply Proposition 4.40 to show that the high energy part would not contribute to the index. However, it is also a problem to control the mixed parts of the Dirac operator. Under additional technical assumptions on \mathcal{D} and with quite some effort, we can solve these problems. However, at this point this part of our work is rather technical and does not seem to lead to interesting applications. For the applications in this article, we can restrict to a more special situation.

Since we want to discuss the contribution of each of the ends of M to the extended index of D as in Theorem 1.17 separately, we will also consider from now on the objects introduced so far for the different ends separately. We let \mathcal{C} be an end of M and keep the notation as before, except for adding an index \mathcal{C} when appropriate. For better readability, we will not decorate the open sets U and U_t , the cross sections N_t , and the distance function f on U with an index \mathcal{C} .

We assume that \mathcal{C} is smooth and cuspidal and that, in addition,

$$(8.1) \quad \kappa_{\mathcal{C}} = \bar{\kappa}_{\mathcal{C}} \quad \text{and} \quad A_{\mathcal{C}}^{\text{le}} \sigma_i = \sum_j a_i^j \sigma_j,$$

for some fixed Hermitian matrix $(a_i^j) \in \text{Gl}(k, \mathbb{C})$. Here $(\sigma_1, \dots, \sigma_k)$ is a family of $\bar{\nabla}_{\mathcal{C}}^E$ -parallel sections of E over U which are pointwise orthonormal and whose restriction to N_t forms an orthogonal basis of $H_{\mathcal{C},t}^c$, for all $t \in \mathbb{R}_+$, as further up. The above conditions hold for homogeneous cusps as discussed in Chapters 9 – 11.

The second condition of (8.1) requires that the space $H_{\mathcal{C},t}^c$ of constant sections in $H_{\mathcal{C},t}$ is invariant under $A_{\mathcal{C},t}$. By Theorem 7.16, we get that

$$(8.2) \quad \begin{aligned} H_{\mathcal{C},t}^{\text{le}} &= H_{\mathcal{C},t}^c = H_{\mathcal{C},[-\lambda_0, \lambda_0]}(A_{\mathcal{C},t}), \\ H_{\mathcal{C},t}^{\text{he}} &= H_{\mathcal{C},t}^{\text{h}} = H_{\mathcal{C},\mathbb{R} \setminus [-\lambda_0, \lambda_0]}(A_{\mathcal{C},t}), \end{aligned}$$

compare (5.22) and (5.23). The additional assumption $\kappa_{\mathcal{C}} = \bar{\kappa}_{\mathcal{C}}$ implies that the high energy family $\mathcal{H}_{\mathcal{C}}^{\text{he}} = (H_{\mathcal{C},t}^{\text{he}})$ is invariant under parallel translation so that it defines a Dirac subsystem $\mathcal{D}_{\mathcal{C}}^{\text{he}}$ of $\mathcal{D}_{\mathcal{C}}$, as in the case of the low energy system $\mathcal{D}_{\mathcal{C}}^{\text{le}} := \mathcal{D}_{\mathcal{C}}^c$; compare (7.9) and Proposition 7.11. Along \mathcal{C} , we obtain corresponding low and high energy Dirac operators D^{le} and D^{he} , decomposing the original Dirac operator D . Since $A_{\mathcal{C},t}^{\text{le}}$ does not depend on t and the spectral gap of $A_{\mathcal{C},t}^{\text{he}}$ tends to infinity as $t \rightarrow \infty$, we may suppose that the assumptions on the spectral gap in Proposition 4.46 are satisfied, for all $t \in \mathbb{R}_+$.

LEMMA 8.3. *Under the above assumptions,*

$$D_{U_t, < \lambda, \text{ext}}^{\text{he}} = D_{U_t, \leq \lambda, \text{ext}}^{\text{he}} = D_{U_t, < \Lambda_t, \text{ext}}^{\text{he}}.$$

and $D_{U_t, < \lambda, \text{ext}}^{\text{he}, \pm}$ and $D_{U_t, < \lambda, \text{ext}}^{\text{he}, \pm}$ are isomorphisms, for all $t \geq 0$ and $-\Lambda_t < \lambda < \Lambda_t$. In particular, for all such t and λ ,

$$\text{ind } D_{U_t, < \lambda, \text{ext}}^+ = \text{ind } D_{U_t, < \lambda, \text{ext}}^{\text{le}, +}.$$

Proof. The first assertion is clear since the spectrum of $A_{\mathcal{C}, t}^{\text{he}}$ does not intersect the interval $(-\Lambda_t, \Lambda_t)$. Furthermore, $D_{U_t, < 0, \text{ext}}^{\text{he}}$ is injective, by Corollary 4.44. Now $D_{U_t, < \lambda, \text{ext}}^{\text{he}}$ and $D_{U_t, \leq -\lambda, \text{ext}}^{\text{he}}$ are adjoints of each other, hence $D_{U_t, \leq 0, \text{ext}}^{\text{he}}$ is surjective. \square

Since $\kappa_{\mathcal{C}}$ depends (at most) on t and $j_{\mathcal{C}}$ solves the initial value problem $j'_{\mathcal{C}} = \kappa_{\mathcal{C}} j_{\mathcal{C}}$ with $j_{\mathcal{C}, 0} = 1$, we conclude that $j_{\mathcal{C}} = j_{\mathcal{C}}(t, x)$ depends only on t as well. Then the linear map

$$(8.4) \quad \Phi : L^2(\mathbb{R}_+, \mathbb{R}^k) \rightarrow L^2(\mathcal{H}^{\text{le}})_{\mathcal{C}}, \quad \Phi(\varphi) = j_{\mathcal{C}}^{-1/2} \sum_i \varphi^i \sigma_i,$$

is a unitary isomorphism such that

$$(8.5) \quad \Phi^{-1} D^{\text{le}} \Phi = \bar{T} \left(\frac{d}{dt} + A_{\mathcal{C}, 0}^{\text{le}} \right),$$

where $\bar{T} = \Phi^{-1} T \Phi$. This is a finite rank Dirac system with constant coefficients. In the super-symmetric case, we get a system of the form

$$(8.6) \quad \Phi^{-1} D^{\text{le}} \Phi = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \left(\frac{d}{dt} + \begin{pmatrix} A_{\mathcal{C}, 0}^{\text{le}, +} & 0 \\ 0 & A_{\mathcal{C}, 0}^{\text{le}, -} \end{pmatrix} \right),$$

where $A_{\mathcal{C}, 0}^{\text{le}, -} = -\bar{A}_{\mathcal{C}, 0}^{\text{le}, +}$. As for high energy,

$$(8.7) \quad \Psi : L^2(\mathbb{R}_+, H_{\mathcal{C}, 0}^{\text{he}}) \rightarrow L^2(\mathcal{H}^{\text{he}})_{\mathcal{C}}, \quad \Psi(\varphi) = j_{\mathcal{C}}^{-1/2} P_{\mathcal{C}}^{\parallel} \varphi,$$

is also a unitary isomorphism, and we have

$$(8.8) \quad \Psi^{-1} D^{\text{he}} \Psi = \bar{T} \left(\frac{d}{dt} + A_{\mathcal{C}, t}^{\text{he}} \right),$$

where $\bar{T} = \Psi^{-1} T \Psi$. There is a corresponding formula in the super-symmetric case, substituting $\bar{A}_{\mathcal{C}, t}^{\text{he}, \pm}$ for $\bar{A}_{\mathcal{C}, 0}^{\pm}$ on the right in (8.6).

PROPOSITION 8.9. *Under the above assumptions, $D_{< 0, \text{ext}}^{\text{le}}$ and $D_{< 0, \text{ext}}^{\text{le}, +}$ are isomorphisms.*

Proof. The Dirac system 8.5 does not have extended or L^2 -solutions σ with $\sigma(0)$ in $H_{< 0}^{\text{le}}$ or $H_{\leq 0}^{\text{le}}$, respectively. \square

THEOREM 8.10. *If all ends of M are smooth and cuspidal and satisfy (8.1), then*

$$\text{ind } D_{\text{ext}}^+ = \int_M \omega_{D^+} + \frac{1}{2} \sum_{\mathcal{C}} \left(\lim_{t \rightarrow \infty} \eta(A_{\mathcal{C}, t}^{\text{he}, +}) + \eta(A_{\mathcal{C}, 0}^{\text{le}, +}) + \dim \ker A_{\mathcal{C}, 0}^{\text{le}, +} \right).$$

Proof. By Theorem 5.13, Lemma 8.3, and Proposition 8.9, we have

$$\operatorname{ind} D_{\text{ext}}^+ = \operatorname{ind} D_{M_t,0}^+,$$

for all $t \in \mathbb{R}_+$. Hence the assertion follows from (5.21) and (7.17), where we separate η -invariant and $\dim \ker A_t^+$ according to cusp \mathcal{C} and low and high energy, observing that $\dim \ker A_{\mathcal{C},t}^{\text{he},+} = 0$ and that $A_{\mathcal{C},t}^{\text{le},+} = A_{\mathcal{C},0}^{\text{le},+}$, for all $t \in \mathbb{R}_+$. \square

Recall the quantities $h_\infty^\pm := \dim \ker D_{\text{ext}}^\pm - \dim \ker D_{\text{max}}^\pm$ from (1.20).

THEOREM 8.11. *If all ends of M are smooth and cuspidal and satisfy (8.1), then*

$$\operatorname{ind}_{L^2} D^+ = \int_M \omega_{D^+} + \frac{1}{2} \sum_{\mathcal{C}} \left(\lim_{t \rightarrow \infty} \eta(A_{\mathcal{C},t}^{\text{he},+}) + \eta(A_{\mathcal{C},0}^{\text{le},+}) \right) - \frac{1}{2} (h_\infty^+ - h_\infty^-).$$

Proof. Since D is formally self-adjoint, the L^2 -index of D vanishes and therefore

$$\operatorname{ind} D_{\text{ext}} = \operatorname{ind} D_{\text{ext}}^+ + \operatorname{ind} D_{\text{ext}}^- = h_\infty^+ + h_\infty^-.$$

On the other hand, we have

$$\omega_{D^-} = -\omega_{D^+} \quad \text{and} \quad A_{\mathcal{C},t}^- = -A_{\mathcal{C},t}^+,$$

for all $t \in \mathbb{R}_+$ and cusps \mathcal{C} of M . Therefore, applying Theorem 8.10 to D^+ and D^- , we obtain that

$$\begin{aligned} \operatorname{ind} D_{\text{ext}} &= \operatorname{ind} D_{\text{ext}}^+ + \operatorname{ind} D_{\text{ext}}^- \\ &= \frac{1}{2} \sum_{\mathcal{C}} \left(\dim \ker A_{\mathcal{C},0}^{\text{le},+} + \dim \ker A_{\mathcal{C},0}^{\text{le},-} \right) = \dim \ker A_{\mathcal{C},0}^{\text{le},+} \end{aligned}$$

since the integral and η terms for D^+ and D^- cancel each other. We conclude that

$$(8.12) \quad h_\infty^+ + h_\infty^- = \dim \ker \bar{A}_{\mathcal{C},0}^{\text{le},+}$$

and hence that

$$(8.13) \quad \operatorname{ind} D_{\text{ext}}^+ - \operatorname{ind}_{L^2} D^+ = h_\infty^+ = \frac{1}{2} (h_\infty^+ - h_\infty^- + \dim \ker A_{\mathcal{C},0}^{\text{le},+}). \quad \square$$

REMARKS 8.14. 1) It is immediate from (8.6) that D^+ is of Fredholm type if and only if the kernel of \bar{A}^+ vanishes or, equivalently, if and only if $h_\infty^+ = h_\infty^- = 0$.

2) Theorems 8.10 and 8.11 generalize Theorems 1.19 and 1.21 from the introduction. In fact, the discussion in the first part of Chapter 9 implies that homogeneous cusps are smooth and satisfy (8.1).

9. HOMOGENEOUS CUSPS

Let N be a simply connected nilpotent Lie group with Lie algebra \mathfrak{n} . Fix a left-invariant Riemannian metric g on N , and let W be a negative definite and symmetric derivation of \mathfrak{n} . Then $(\exp(-tW))_{t \in \mathbb{R}}$ is a one-parameter group of automorphisms of \mathfrak{n} which induces a one-parameter group $(\Phi_t)_{t \in \mathbb{R}}$ of automorphisms of N . The associated semi-direct product $S := \mathbb{R} \ltimes N$, where

$$(9.1) \quad (s, x)(t, y) := (s + t, x\Phi_s(y)),$$

is a simply connected solvable Lie group containing $N \cong \{0\} \times N$ as a subgroup of codimension one. The vector field $T := \partial/\partial t$ on S is left-invariant, and the Lie algebra \mathfrak{s} of S extends \mathfrak{n} by

$$(9.2) \quad [T, X] = -WX,$$

where $X \in \mathfrak{n}$. For later use, we note that left translation, right translation, and conjugation with $(t, e) \in S$ are given by

$$(9.3) \quad \begin{aligned} L_{(t,e)}(s, x) &= (s + t, \Phi_t(x)), \\ R_{(t,e)}(s, x) &= (s + t, x), \\ (t, e)(s, x)(-t, e) &= (s, \Phi_t(x)), \end{aligned}$$

respectively. In particular, the shift by t along the T -lines is obtained by right translation with (t, e) . Moreover, for $X \in \mathfrak{n} \subseteq \mathfrak{s}$,

$$(9.4) \quad \begin{aligned} R_{(t,e)*}X_{(s,x)} &= L_{(s,x)*}L_{(t,e)*}L_{(-t,e)*}R_{(t,e)*}X \\ &= L_{(s+t,x)*}(\text{Ad}_{(t,e)}^{-1}X) = L_{(s+t,x)*}(\exp(tW)X), \end{aligned}$$

where we recall that (Φ_t) is the one-parameter group of automorphism of N associated to $-W$ (and where we identify $\mathfrak{s} \ni X = X_e \in T_e S$).

Endow S with the left-invariant Riemannian metric which agrees with g along N and such that \mathbb{R} and \mathfrak{n} are pairwise perpendicular with $|T| = 1$. Note that T is a unit normal field along the cross sections $N_t := \{t\} \times N$ and that the T -lines are unit speed geodesics. In particular,

$$(9.5) \quad f : S \rightarrow \mathbb{R}, \quad f(t, x) := t,$$

is a smooth distance function on S such that $\text{grad } f = T$ and such that the associated diffeomorphism F is the identity on $S = \mathbb{R} \times N$. By the Koszul formula and the symmetry of W ,

$$(9.6) \quad \nabla_T X = 0,$$

for any $X \in \mathfrak{s}$. For any $X \in \mathfrak{n} \subseteq \mathfrak{s}$,

$$(9.7) \quad \nabla_X T = WX,$$

by (9.2) and (9.6); that is, except for the compactness of the cross sections, we are in the situation of Section 3.2. By (9.6) and (9.7),

$$(9.8) \quad R(T, X)Y = \nabla_{[X, T]}Y = \nabla_{WX}Y,$$

for all $X \in \mathfrak{n}$ and $Y \in \mathfrak{s}$. In particular,

$$(9.9) \quad R(X, T)T = -W^2X,$$

and hence the sectional curvature of tangential 2-planes of S containing T is strictly negative.

Let $\Gamma \subseteq N$ be a discrete subgroup such that the quotient $\Gamma \backslash N$ is compact. Since $\Gamma \subseteq N$, the distance function f as in (9.5) is well defined on $\Gamma \backslash S$. We keep the notation f and $T = \text{grad } f$ on the quotient. The cross sections of f are given by $\{t\} \times \Gamma \backslash N$, and right translation by (t, e) induces the shift F_t from $\Gamma \backslash N$ to $\{t\} \times \Gamma \backslash N$, see (9.3). By (9.4), F_t has derivative $F_{t*} = \exp(tW)$. The Jacobian of F_t is given by $j(t) = \exp(\kappa t)$, where $\kappa = \text{tr } W$ as in Section 3.2. It only depends on t and not on $x \in \Gamma \backslash N$. Moreover, since W is negative definite, F_t is contracting for $t > 0$: If we order the eigenvalues of W ,

$$(9.10) \quad \kappa_2 \leq \dots \leq \kappa_m < 0,$$

then any part $[t_0, \infty) \times N$ of $\mathbb{R} \times N$ models cuspidal ends as in Definition 1.12 with $c = -2\kappa_m$ and $C = 1$. We call such ends *homogeneous cusps*.

If $X_i \in \mathfrak{n}$ is a unit eigenvector of W for the eigenvalue κ_i , then $\exp(\kappa_i t)X_i$ is a Jacobi field along each T -line and

$$(9.11) \quad \langle \nabla_X Y, Z \rangle = - \int_0^\infty \langle R(T, e^{\kappa_i t} X_i)Y, Z \rangle,$$

for all $Y, Z \in \mathfrak{s}$. It follows that the flat connection $\bar{\nabla}$ associated to the cusp as in Section 7.1 defines left-invariant vector fields on S or, rather, their image in $\Gamma \backslash S$ to be $\bar{\nabla}$ -parallel.

Let K_0 be a connected Lie subgroup of the orthogonal group $\text{SO}(\mathfrak{s})$ which contains the holonomy group of S at e . Denote the Lie algebra of K_0 by \mathfrak{k} . Consider the principal bundle $\mathcal{P}_0 := S \times K_0$ over S , with structure group K_0 , where we view $p = (s, k) \in \mathcal{P}_0$ as representing the frame $L_s \circ k : T_e S \rightarrow T_s S$ of S , where L_s denotes left-translation by s (and its derivative). This interpretation corresponds to an embedding of \mathcal{P}_0 into the principal bundle of orthonormal frames of S . The group S acts on \mathcal{P}_0 by left translation, $s(s', k) := (ss', k)$, and the orbits of this action are the left-invariant frames $F_k := \{(s, k) \mid s \in S\}$ over S .

LEMMA 9.12. *The Levi-Civita connection ∇ and flat connection $\bar{\nabla}$ of S reduce to \mathcal{P}_0 . That is, if $c : I \rightarrow S$ is a smooth curve and F is a*

parallel frame along c with respect to ∇ or $\bar{\nabla}$ such that $F(t_0) \in \mathcal{P}_0$ for some $t_0 \in I$, then $F(t) \in \mathcal{P}_0$ for all $t \in I$.

Proof. Let F be an orthonormal frame along c , and write $F(t) = L_{c(t)}f(t)$, where $f : I \rightarrow \mathcal{O}(\mathfrak{s})$. Then the covariant derivative of F along c with respect to ∇ is given by

$$(9.13) \quad F'(t) = L_{c(t)}(f'(t) + A_{c'(t)}f(t)),$$

where

$$(9.14) \quad A_X = \begin{cases} R(T, W^{-1}X) & \text{for } X \in \mathfrak{n}, \\ 0 & \text{for } X = T, \end{cases}$$

by (9.6) and (9.8). By (9.13), F is ∇ -parallel if $f' + A_c f = 0$.

Now $R(Y, Z)$ is in the Lie algebra of the holonomy group of S at e , for all $Y, Z \in \mathfrak{s}$, hence also $A_{c'(t)}$, for all $t \in I$. Since K_0 contains the holonomy group of S at e , we get that $A_{c'(t)} \in \mathfrak{k}$, for all $t \in I$. It follows that a solution of $f' + A_c f = 0$ is contained in K_0 if $f(t_0)$ is in K_0 , for some $t_0 \in I$. This proves the assertion for ∇ .

By what we said above, a frame is $\bar{\nabla}$ -parallel if and only if it is left-invariant under S . Hence the $\bar{\nabla}$ -parallel frames along c are of the form $F(t) = L_{c(t)}k$, $t \in I$, where $k \in \mathcal{O}(\mathfrak{s})$. Hence, if $F(t_0) \in \mathcal{P}_0$ for some $t_0 \in I$, then $k \in K_0$, and then $F = F_k$ is contained in \mathcal{P}_0 . \square

Let $K \rightarrow K_0$ be a covering homomorphism, where K is a connected Lie group, and let $\mathcal{P} := S \times K$ be the corresponding covering space of \mathcal{P}_0 , a principal bundle over S with structure group K . Via the projection $K \rightarrow K_0$, identify the Lie algebra of K with the Lie algebra \mathfrak{k} of K_0 . As in the case of \mathcal{P}_0 , S acts by left translations on \mathcal{P} , and we have the corresponding orbits F_k , $k \in K$. Moreover, since $\mathcal{P} \rightarrow \mathcal{P}_0$ is a covering projection, Levi-Civita and flat connection lift from \mathcal{P}_0 to \mathcal{P} .

Denote by $\hat{\alpha}_* : \mathfrak{k} \rightarrow \mathfrak{u}(\Sigma_{\mathfrak{s}})$ the composition of the differential of $\alpha : K \rightarrow K_0 \subseteq \mathrm{SO}(\mathfrak{s})$ with the differential of the spinor representation $\Sigma_{\mathfrak{s}}$ of $\mathfrak{so}(\mathfrak{s}) \simeq \mathfrak{spin}(\mathfrak{s})$. Let V be a finite dimensional Hermitian vector space and $\pi_* : \mathfrak{k} \rightarrow \mathfrak{u}(V)$ be a unitary representation. Suppose that there is a unitary representation $\beta : K \rightarrow \Sigma_{\mathfrak{s}} \otimes V$ with

$$(9.15) \quad \hat{\alpha}_* \otimes \mathrm{id} + \mathrm{id} \otimes \pi_* = \beta_*,$$

and let $E = \mathcal{P} \times_{\beta} (\Sigma_{\mathfrak{s}} \otimes V)$ be the associated Hermitian vector bundle over S . Levi-Civita and flat connection on \mathcal{P} induce Hermitian connections ∇^E and $\bar{\nabla}^E$ on E , respectively. We extend Clifford multiplication to $\Sigma_{\mathfrak{s}} \otimes V$ by

$$(9.16) \quad X \cdot (u \otimes v) := (X \cdot u) \otimes v,$$

where $X \in \mathfrak{s}$, $u \in \Sigma_{\mathfrak{s}}$, and $v \in V$. By (9.15) and since K is connected, Clifford multiplication commutes with β , that is

$$(9.17) \quad \beta(k)(Xw) = X(\beta(k)w),$$

for all $k \in K$, $X \in \mathfrak{s}$, and $w \in \Sigma_{\mathfrak{s}} \otimes V$. Hence (9.16) induces a Clifford multiplication on E which turns E into a Dirac bundle over S . The canonical action of S on E preserves the Dirac data of E ; we say that E is a *homogeneous* Dirac bundle over S .

Using the left-invariant orbit F_e in \mathcal{P} , we view sections of E as smooth maps $\sigma : S \rightarrow \Sigma_m \otimes V$. In this interpretation, covariant derivatives and Dirac operator are given by

$$(9.18) \quad \nabla_X^E \sigma = X(\sigma) + \beta_*(A_X)\sigma, \quad \bar{\nabla}_X^E \sigma = X(\sigma),$$

and

$$(9.19) \quad D\sigma = \sum_j X_j \cdot (X_j(\sigma) + \beta_*(A_{X_j})\sigma),$$

where X is a vector field on S , (X_1, \dots, X_m) is an orthonormal frame of S , and A_X is as in (9.14). In particular, σ is $\bar{\nabla}^E$ -parallel if and only if σ is constant.

Let τ be a unitary representation of Γ on V , the *twist*, and assume that τ and π_* commute, that is,

$$(9.20) \quad \tau(\gamma)\pi_*(Y) = \pi_*(Y)\tau(\gamma),$$

for all $\gamma \in \Gamma$ and $Y \in \mathfrak{k}$.

LEMMA 9.21. *Extend τ by the trivial representation on $\Sigma_{\mathfrak{s}}$ to $\Sigma_{\mathfrak{s}} \otimes V$. Then τ commutes with β and Clifford multiplication,*

$$\tau(\gamma)(\beta(k)w) = \beta(k)(\tau(\gamma)w),$$

$$\tau(\gamma)(Xw) = X(\tau(\gamma)w),$$

for all $\gamma \in \Gamma$, $k \in K$, $X \in \mathfrak{s}$, and $w \in \Sigma_{\mathfrak{s}} \otimes V$.

Proof. Since K is connected, the first assertion follows from the corresponding infinitesimal properties in (9.15) and (9.20). As for the second assertion, we note that τ acts trivially on the first and Clifford multiplication trivially on the second factor of $\Sigma_{\mathfrak{s}} \otimes V$. \square

By Lemma 9.21, τ induces a Hermitian bundle E_{τ} over $\Gamma \backslash S$ such that sections of E_{τ} correspond to maps $\sigma : S \rightarrow \Sigma_{\mathfrak{s}} \otimes V$ which satisfy

$$(9.22) \quad \sigma(\gamma s) = \tau(\gamma)\sigma(s),$$

for all $s \in S$. The connections ∇^E and $\bar{\nabla}^E$ on E descend to Hermitian connections on E_{τ} , also denoted by ∇^E and $\bar{\nabla}^E$, respectively. Moreover, E_{τ} inherits Clifford multiplication from E and thus turns into a Dirac bundle over $\Gamma \backslash S$.

EXAMPLES 9.23. 1) (Spinor bundles) Since S is contractible, spin structures over $\Gamma \backslash S$ are determined by homomorphisms $\tau : \Gamma \rightarrow \{+1, -1\}$. In our setup, the corresponding spinor bundles over $\Gamma \backslash S$ can be given by the data: $K_0 = \mathrm{SO}(\mathfrak{s})$ and $K = \mathrm{Spin}(\mathfrak{s})$, $\alpha : \mathrm{Spin}(\mathfrak{s}) \rightarrow \mathrm{SO}(\mathfrak{s})$ the canonical covering map, $V = \mathbb{C}$, $\pi_* = 0$, β the spinor representation, extended trivially to the factor \mathbb{C} of $\Sigma_{\mathfrak{s}} \otimes \mathbb{C}$, and finally the twist defined by τ , where γ acts by multiplication with $\tau(\gamma) = \pm 1$ on \mathbb{C} .

2) (Clifford bundle) If m is even, then $\mathrm{Cl}(\mathfrak{s}) = \Sigma_{\mathfrak{s}} \otimes \Sigma_{\mathfrak{s}}$. Thus, to obtain the Clifford bundle over $\Gamma \backslash S$, we may take $K_0 = K = \mathrm{SO}(\mathfrak{s})$, $\alpha = \mathrm{id}$, $V = \Sigma_{\mathfrak{s}}$, β_* the differential of the spinor representation, and τ the trivial representation of Γ on $\Sigma_{\mathfrak{s}}$.

If the dimension m of S is even, then the ± 1 -eigenspaces $\Sigma_{\mathfrak{s}}^{\pm} \otimes V$ of multiplication by the complex volume form (compare Section 2.2) are invariant under β , by (9.17). By Lemma 9.21, they are also invariant under τ . Thus the complex volume form yields the super-symmetry $E = E^+ \oplus E^-$ with

$$(9.24) \quad E^{\pm} = \mathcal{P} \times_{\beta} (\Sigma_{\mathfrak{s}}^{\pm} \otimes V).$$

In the case of the Clifford bundle, there is another natural super-symmetry, namely the even-odd decomposition. Our methods also allow for a discussion of the latter, but here and below we concentrate on the decomposition given by the complex volume form.

We now pass to the Dirac system associated to the distance function f and the Dirac bundle E_{τ} over $\Gamma \backslash S$. We identify sections of E_{τ} over $\{t\} \times \Gamma \backslash N$ with maps $\sigma : N \rightarrow \Sigma_{\mathfrak{s}} \otimes V$ satisfying (9.22). Under this identification, parallel translation along the T -lines is the identity, and the Hilbert space $L^2(\{t\} \times \Gamma \backslash N, E_{\tau})$ corresponds to the Hilbert space of measurable maps $N \rightarrow \Sigma_{\mathfrak{s}} \otimes V$ satisfying (9.22) which are square integrable over a fundamental domain of Γ . In the notation of (3.46),

$$(9.25) \quad A_t \sigma = - \sum_{2 \leq j \leq m} e^{-\kappa_j t} T X_j \cdot X_j(\sigma) - \sum_{2 \leq j \leq m} T X_j \cdot \beta_*(A_{X_j}) \sigma - \frac{\kappa}{2} \sigma,$$

where (X_2, \dots, X_m) is an orthonormal basis of \mathfrak{n} consisting of eigenvectors of W , $W X_i = \kappa_i X_i$.

We may also have a different view on E_{τ} over $\{t\} \times \Gamma \backslash N$: $L_{(t,e)}$ is an isometry of S which maps N to $\{t\} \times N$ and which leaves the normal field T to the cross sections $\{t\} \times N$ invariant. Suppressing the coordinate t in $\{t\} \times N$, $L_{(t,e)}$ corresponds to Φ_t , by (9.3). That is, E_{τ} over $\{t\} \times \Gamma \backslash N$ corresponds to $E_{\Phi_t \tau \Phi_t^{-1}}$ over $\Phi_t(\Gamma) \backslash N$, where N is endowed with the fixed left-invariant metric g . Under this correspondence, the exponential factors in the expression for A_t in (9.25) disappear. More

precisely, $-A_t$ corresponds to the Dirac operator

$$(9.26) \quad D_t \sigma = \sum_{2 \leq j \leq m} TX_j \cdot X_j(\sigma) + \sum_{2 \leq j \leq m} TX_j \cdot \beta_*(A_{X_j})\sigma + \frac{\kappa}{2}\sigma,$$

where σ satisfies the twist data with respect to $\Phi_t \tau \Phi_t^{-1}$. In particular, the local data for the different operators D_t coincide under the correspondence.

9.1. Asymptotic η -Invariants. Let $L^{2,\pm}(t)$ be the Hilbert space of measurable maps $N \rightarrow \Sigma_{\mathfrak{s}}^{\pm} \otimes V$ satisfying (9.22) with respect to $\Phi_t \tau \Phi_t^{-1}$ which are square integrable over a fundamental domain of $\Phi_t(\Gamma)$. Then $D_t^{\pm} = -A_t^{\pm}$ is an unbounded self-adjoint operator on $L^{2,\pm}(t)$.

For the computation of the asymptotic high energy η -invariant of D_t^+ , it will be useful to consider the *flat Dirac operator* \bar{D}_t^+ , defined by

$$(9.27) \quad \bar{D}_t^+ \sigma = \sum_{2 \leq j \leq m} TX_j \cdot X_j(\sigma).$$

We note that \bar{D}_t^+ is a formally self-adjoint operator and that $D_t^+ - \bar{D}_t^+$ is left-invariant of order zero. In particular, the principal symbols of D_t^+ and \bar{D}_t^+ are the same. We have

$$(9.28) \quad \begin{aligned} (\bar{D}_t^+)^2 \sigma &= \sum_{2 \leq j, k \leq m} TX_j \cdot X_j(TX_k \cdot X_k(\sigma)) \\ &= - \sum_{2 \leq j \leq m} X_j(X_j(\sigma)) + \sum_{2 \leq j \neq k \leq m} X_j X_k \cdot X_j(X_k(\sigma)), \\ &= \Delta \sigma + \sum_{2 \leq j < k \leq m} X_j X_k \cdot [X_j, X_k](\sigma), \end{aligned}$$

where Δ denotes the standard Laplace operator of N , here acting on maps with values in the vector space $\Sigma_{\mathfrak{s}}^{\pm} \otimes V$. If \mathfrak{n} is two-step nilpotent, then the Lie brackets $[X_j, X_k]$ in the second term on the right of (9.28) are in the center of \mathfrak{n} , and then the operator defined by the second term commutes with Δ .

The idea to consider \bar{D}_t^+ is taken from [DeSi]. The proof of our main result in this direction, Theorem 9.29 below, is a variation of arguments in §5 of [DeSi]. This line of reasoning was also used by Cheeger and Gromov in order to show that their ρ -invariant is the limit of the (signature) η -invariant under a collapse of the corresponding manifold with bounded covering geometry [ChGr].

THEOREM 9.29. *For D_t^+ and \bar{D}_t^+ as above, we have*

$$\lim_{t \rightarrow \infty} \eta(D_t^{\text{he},+}) = \lim_{t \rightarrow \infty} \eta(\bar{D}_t^+),$$

Proof. Left-invariant sections in $L^{2,+}(t)$ are in the kernel of \bar{D}_t^+ , for all t . By what we said above, sections in $L^{2,+}(t)$ are $\bar{\nabla}$ -parallel if and only if they are left-invariant. Since the difference $D_t^+ - \bar{D}_t^+$ is left-invariant, hence uniformly bounded, Theorem 7.16 implies that the kernel of the operator \bar{D}_t^+ consists precisely of the left-invariant sections in $L^{2,+}(t)$, for all sufficiently large t .

Let $P_t : L^{2,+}(t) \rightarrow L^{2,+}(t)$ be the orthogonal projection onto the space of left-invariant sections in $L^{2,+}(t)$. Then P_t commutes with \bar{D}_t^+ and $D_{t,c}^+$, where we write $D_{t,c}^+ = D_t^+ - \bar{D}_t^+$. For fixed t , consider the family of operators

$$D_{t,u}^+ := \bar{D}_t^+ + u(I - P_t)D_{t,c}^+(I - P_t) + P_t, \quad 0 \leq u \leq 1.$$

By definition,

$$\begin{aligned} \eta(D_{t,1}^+) &= \eta(D_t^{\text{he},+}) + \dim \operatorname{im} P_t, \\ \eta(D_{t,0}^+) &= \eta(\bar{D}_t^+) + \dim \operatorname{im} P_t. \end{aligned}$$

The non-zero eigenvalues of \bar{D}_t^+ tend to infinity as t tends to ∞ , whereas $D_{t,c}^+$ is uniformly bounded independently of t . It follows that $D_{t,u}^+$ is invertible, for all sufficiently large t . Now by Proposition 2.12 in [APS3] and the invertibility of $D_{t,u}^+$,

$$\frac{d}{du} \eta(D_{t,u}^+)$$

is a local invariant¹¹, given by an explicit integral formula constructed out of the complete symbols of $D_{t,u}^+$ and $(I - P_t)D_{t,c}^+(I - P_t)$. On the other hand, P_t is (infinitely) smoothing, and hence the complete symbol of $D_{t,u}^+$ and $(I - P_t)D_{t,c}^+(I - P_t)$ are the same as those of

$$L_{t,u} := \bar{D}_t^+ + uD_{t,c}^+ \quad \text{and} \quad D_{t,c}^+.$$

Now the symbols of $L_{t,u}$ and $D_{t,c}^+$ do not depend on t , by (9.26) and (9.27). It follows that the local invariant for $d\eta(D_{t,u}^+)/du$ is bounded in modulus by a continuous function $b = b(u)$ which does not depend on t . Therefore we have

$$\begin{aligned} |\eta(D_t^{\text{he},+}) - \eta(\bar{D}_t^+)| &= |\eta(D_{t,1}^+) - \eta(D_{t,0}^+)| \\ &\leq \operatorname{const} \cdot \operatorname{vol}(\Phi_t(\Gamma) \backslash N) \rightarrow 0. \end{aligned} \quad \square$$

The fact that the high energy η -invariant has no spectral flow is perhaps an indication that its limit deserves to be investigated along the lines of the discussion of the ρ -invariant in [ChGr].

¹¹In [APS3] this assertion is only stated for the η -invariant modulo \mathbb{Z} . However, as is clear from the remarks preceding Proposition 2.12 in [APS3], this is only because of the possibility of eigenvalues crossing 0, which is excluded by invertibility.

9.2. Vanishing of η -Invariants. Let Z belong to the center of \mathfrak{n} . Then Z commutes with Γ and pushes forward to a vector field on the quotient $\Gamma \backslash N$, also denoted by Z .

LEMMA 9.30. *Clifford multiplication with Z commutes with $(\bar{D}_t^+)^2$.*

Proof. We can assume that Z has norm one. Choosing $X_2 = Z$, then, in the second sum on the right in (9.28) above, the terms with $i = 2$ vanish since Z commutes with all the X_i , $i > 2$. \square

THEOREM 9.31. *If the center of N has dimension at least two, then the spectrum of \bar{D}_t^+ , including multiplicities, is symmetric about zero. In other words, the eta function of \bar{D}_t^+ vanishes identically.*

Proof. Choose orthonormal vector fields Z and Z' in the center of \mathfrak{n} and let W_\pm be the eigenspaces of the involution iZ in $\Sigma_\mathfrak{s}^+ \otimes V$ for the eigenvalues ± 1 . Since $(\bar{D}_t^+)^2$ commutes with iZ , see Lemma 9.30, it leaves the spaces of sections with values in W_+ and W_- invariant. In particular, if $\lambda > 0$ is an eigenvalue of $(\bar{D}_t^+)^2$ and $\mathcal{S}(\lambda)$ denotes the corresponding eigenspace of sections, then

$$\mathcal{S}(\lambda) = \mathcal{S}_+(\lambda) \oplus \mathcal{S}_-(\lambda),$$

where $\mathcal{S}_+(\lambda)$ and $\mathcal{S}_-(\lambda)$ consist of eigensections in $\mathcal{S}(\lambda)$ with values in W_+ and W_- , respectively.

We note that $\mathcal{S}(\lambda)$ is invariant under \bar{D}_t^+ and that \bar{D}_t^+ has eigenvalues $\pm\sqrt{\lambda}$ on $\mathcal{S}(\lambda)$. Furthermore, the multiplicities of $\sqrt{\lambda}$ and $-\sqrt{\lambda}$ as eigenvalue of \bar{D}_t^+ coincide if and only if the trace of \bar{D}_t^+ on $\mathcal{S}(\lambda)$ vanishes.

We let $X_2 = Z$. Then $X_i W_+ = W_-$ and $X_i W_- = W_+$ for $3 \leq i \leq m$, and hence the corresponding terms of \bar{D}_t^+ do not contribute to the trace of \bar{D}_t^+ on $\mathcal{S}(\lambda)$. Now the remaining term $X_2 \cdot X_2(\sigma) = Z \cdot Z(\sigma)$ of $\bar{D}_t^+ \sigma$ leaves $\mathcal{S}(\lambda)$ invariant, and its trace on $\mathcal{S}(\lambda)$ is equal to the trace of \bar{D}_t^+ on $\mathcal{S}(\lambda)$, by what we just said.

Clifford multiplication with Z' leaves $\mathcal{S}(\lambda)$ invariant, by Lemma 9.30. On the other hand,

$$Z \cdot Z(Z' \cdot \sigma) = Z \cdot (Z' \cdot Z(\sigma)) = -Z' \cdot (Z \cdot Z(\sigma)),$$

that is, the involution iZ' anticommutes with the operator which sends σ to $Z \cdot Z(\sigma)$. It follows that the trace of \bar{D}_t^+ on $\mathcal{S}(\lambda)$ vanishes. \square

COROLLARY 9.32. *If the center of N has dimension at least two, then the asymptotic high energy η -invariant $\lim_{t \rightarrow \infty} \eta(A_t^{\text{he},+}) = 0$.*

Proof. Recall that $A_t^+ = -D_t^+$ and apply Theorems 9.29 and 9.31. \square

10. HEISENBERG MANIFOLDS

The only simply connected two-step nilpotent Lie groups not covered by Theorem 9.31 are the standard Heisenberg groups $N = G_n$, where here $m - 1 = \dim N = 2n + 1$. We represent G_n as $\mathbb{R}^n \times \mathbb{R}^n \times \mathbb{R}$ with typical element (x, y, z) and multiplication given by

$$(10.1) \quad (x, y, z)(x', y', z') = (x + x', y + y', z + z' + xy').$$

The left-invariant vector fields

$$(10.2) \quad X_j := \frac{\partial}{\partial x_j}, \quad Y_j := \frac{\partial}{\partial y_j} + x_j \frac{\partial}{\partial z}, \quad \text{and} \quad Z := \frac{\partial}{\partial z}$$

form a basis of the Lie algebra of G_n . They commute pairwise, except for the n Lie brackets $[X_j, Y_j] = Z$.

10.1. Lattices in Heisenberg groups. Lattices in G_n are classified in [GoWi, Section 2]: Let D_n be the set of n -tuples $d = (d_1, \dots, d_n)$ of natural numbers such that d_i divides d_{i+1} , $1 \leq i < n$. Then, for any $d \in D_n$,

$$(10.3) \quad \Gamma_d := \{(x, y, z) \mid x, y \in \mathbb{Z}^n, z \in \mathbb{Z}, d_i \text{ divides } x_i\}$$

is a lattice in G_n . The isomorphism type of Γ_d is determined by d and, up automorphism of G_n , any lattice Γ in G_n is equal to some Γ_d , $d \in D_n$. Following the notation in [GoWi], we then set

$$(10.4) \quad |\Gamma| := d_1 \cdots d_n.$$

Fix $d \in D_n$. The $2n + 1$ elements

$$(10.5) \quad \phi_j := (d_j e_j, 0, 0), \quad \psi_j := (0, e_j, 0), \quad \zeta := (0, 0, 1)$$

generate Γ_d . They commute pairwise, except for the n relations

$$(10.6) \quad \phi_j \psi_j \phi_j^{-1} \psi_j^{-1} = \zeta^{d_j} = (0, 0, d_j).$$

Let τ be an irreducible unitary representation of Γ_d on a finite dimensional Hermitian vector space V . Since τ is irreducible and ζ is central, there is a number $c \in [0, 1)$ with

$$(10.7) \quad \tau(\zeta) = e^{2\pi i c} I.$$

Let $A_j := \tau(\phi_j)$ and $B_j := \tau(\psi_j)$, for $1 \leq j \leq n$. Then, if λ is an eigenvalue of B_j , for some j and some eigenvector $v \in V$, then

$$(10.8) \quad B_j(A_j v) = e^{-2\pi i c d_j} (A_j B_j A_j^{-1})(A_j v) = e^{-2\pi i c d_j} \lambda A_j v,$$

and hence $e^{-2\pi i c d_j} \lambda$ is an eigenvalue of B_j as well. It follows that c is rational, by the finite dimensionality of V .

Let m_j be the denominator of cd_j . Consider the sublattice $\Gamma_{md} \subseteq \Gamma_d$, where $md := (m_1d_1, \dots, m_nd_n)$. The order of Γ_{md} in Γ_d is

$$(10.9) \quad |\Gamma_d/\Gamma_{md}| = m_1 \cdots m_n,$$

and τ restricts to an Abelian representation on Γ_{md} . By irreducibility, τ is induced from a one-dimensional representation of Γ_{md} . That is, there are real numbers $\alpha_1, \beta_1, \dots, \alpha_n, \beta_n$ such that $\phi_j^{m_j}$ and ψ_j act on \mathbb{C} by multiplication with $e^{2\pi i \alpha_j}$ and $e^{2\pi i \beta_j}$, respectively, and τ is induced from this representation of Γ_{md} . In particular,

$$(10.10) \quad \dim V = m_1 \cdots m_n.$$

For any n -tuple

$$(10.11) \quad b = (b_1, \dots, b_n) = (\beta_1 + l_1cd_1, \dots, \beta_n + l_ncd_n) \in \mathbb{R}^n/\mathbb{Z}^n,$$

where $(l_1, \dots, l_n) \in \mathbb{Z}^n$, we let V_b be the subspace of V on which ψ_j acts by $e^{2\pi i b_j}$. We note that these subspaces V_b are one-dimensional and pairwise orthogonal and that they span V .

10.2. Twisted Right Regular Representation. The set

$$(10.12) \quad F := \{(x, y, z) \in G_n \mid x \in P, (y, z) \in Q\},$$

where

$$(10.13) \quad \begin{aligned} P &:= \{x \in \mathbb{R}^n \mid 0 \leq x_j \leq d_j\}, \\ Q &:= \{(y, z) \in \mathbb{R}^n \times \mathbb{R} \mid 0 \leq y_j, z \leq 1\}, \end{aligned}$$

is a fundamental domain of the action of Γ_d on G_n by left translations. Observe that, by (10.2), the standard Lebesgue measure with respect to the (x, y, z) -coordinates is left-invariant, hence bi-invariant, on G_n .

Fix an irreducible unitary representation τ of Γ_d on a finite dimensional Hermitian vector space V as above. and consider the Hilbert space $L^2(\tau)$ of maps $\sigma : G_n \rightarrow V$ such that

$$(10.14) \quad \sigma(\gamma g) = \tau(\gamma)\sigma(g)$$

for all $\gamma \in \Gamma_d$ and $g \in G_n$ which are square integrable over F . The *right regular representation* ρ of G_n acts unitarily on $L^2(\tau)$ by

$$(10.15) \quad (\rho(g)\sigma)(x, y, z) = \sigma((x, y, z)g),$$

and our next aim is to determine the multiplicities of the irreducible unitary representations of G_n in $L^2(\tau)$. Here we recall that irreducible unitary representations of the Heisenberg group G_n correspond to coadjoint orbits of G_n , by the classical theorem of Stone and von Neumann (or by the more general Kirillov theory, respectively). This correspondence will show up in the following discussion.

Let $\sigma \in L^2(\tau)$. Then

$$(10.16) \quad e^{2\pi ic}\sigma(x, y, z) = \tau(\zeta)\sigma(x, y, z) = \sigma(x, y, z + 1).$$

Let σ_b be the component of σ in V_b . Then

$$(10.17) \quad e^{2\pi ib_j}\sigma_b(x, y, z) = B_j\sigma_b(x, y, z) = \sigma_b(x, y + e_j, z).$$

The transformation rule with respect to A_j is more complicated,

$$(10.18) \quad A_j\sigma_b(x, y, z) = \sigma_{b-cd_je_j}(x + d_je_j, y, z + d_jy_j).$$

By (10.16) and (10.17), we can develop σ_b in a Fourier series,

$$(10.19) \quad \sigma_b(x, y, z) = \sum_{\substack{v \equiv b \\ w \equiv c}} e^{2\pi i(vy+wz)} \sigma_{v,w}(x)$$

where \equiv indicates congruence modulo \mathbb{Z}^n . Fix a w congruent to c and consider the space $L^2(\tau, w)$ of $\sigma \in L^2(\tau)$ with

$$(10.20) \quad \sigma(x, y, z + t) = e^{2\pi iwt}\sigma(x, y, z),$$

that is, in the above Fourier development of the components σ_b of σ , only the terms with the given w occur. We obtain an orthogonal decomposition

$$(10.21) \quad L^2(\tau) = \oplus_{w \equiv c} L^2(\tau, w).$$

Now the spaces $L^2(\tau, w)$ are ρ -invariant and, therefore, it remains to investigate ρ on them. For $\sigma \in L^2(\tau, w)$, we have

$$(10.22) \quad \begin{aligned} \sum_{u \equiv b + cd_je_j} e^{2\pi i(uy+wz)} A_j\sigma_{u,w}(x) &= A_j\sigma_{b+cd_je_j}(x, y, z) \\ &= \sigma_b(x + d_je_j, y, z + d_jy_j) \\ &= e^{2\pi iwd_jy_j}\sigma_b(x + d_je_j, y, z) \\ &= \sum_{v \equiv b} e^{2\pi i((v+wd_je_j)y+wz)} \sigma_{v,w}(x + d_je_j). \end{aligned}$$

We conclude that, for any $v \equiv b$ and $x \in \mathbb{R}^n$,

$$(10.23) \quad \sigma_{v+wd_je_j,w}(x) = A_j^{-1}\sigma_{v,w}(x + d_je_j).$$

There are two cases, $w = 0$ and $w \neq 0$, respectively.

If $w = 0$, then $w = c = 0$ and $\dim V = 1$. By (10.23), the Fourier coefficients $\sigma_{v,0}$ are d_j -periodic in x_j up to the twists by the complex numbers A_j of norm one.

Suppose now that $w \neq 0$. Then, by (10.23), the Fourier coefficients $\sigma_{u,w}$ with

$$(10.24) \quad u = b + k_1e_1 + \cdots + k_ne_n, \quad 0 \leq k_j < |w|d_j,$$

determine all the Fourier coefficients of σ . We also get

$$(10.25) \quad \|\sigma\|_{L^2(\tau, w)}^2 = \sum \|\sigma_{u, w}\|_{L^2(\mathbb{R}^n, V_b)}^2,$$

where the sum is over all u as in (10.24). Here we recall that, on the left hand side, the L^2 -norms are given by the corresponding integrals over the fundamental domain F of Γ_d as in (10.12), whereas the integrals on the right hand side are over Euclidean x -space. We obtain

$$(10.26) \quad L^2(\tau, w) \cong \oplus L^2(\mathbb{R}^n, V_b),$$

where we have $m_1 d_1 \cdots m_n d_n |w|^n$ summands $L^2(\mathbb{R}^n, V_b)$ on the right hand side, namely one for each u as in (10.24).

To identify ρ on $\oplus L^2(\mathbb{R}^n, V_b) \cong L^2(\tau, w)$, let $g = (x', y', z') \in G_n$ and recall (10.1) and (10.15). We compute

$$(10.27) \quad \begin{aligned} e^{2\pi i(u(y+y') + w(z+z'+xy'))} \sigma_{u, w}(x + x') \\ = e^{2\pi i(uy + wz)} e^{2\pi i(uy' + w(z'+xy'))} \sigma_{u, w}(x + x'), \end{aligned}$$

hence g acts on $\sigma_{u, w} \in L^2(\mathbb{R}^n, V_b)$ by

$$(10.28) \quad (\rho(g)\sigma_{u, w})(x) = e^{2\pi i(uy' + w(z'+xy'))} \sigma_{u, w}(x + x').$$

Via a unitary identification $V_b \cong \mathbb{C}$ and the substitution $x + u/c$ for x , we see that ρ on $L^2(\mathbb{R}^n, V_b)$ is unitarily equivalent to the irreducible unitary representation ρ_w of G_n on $L^2(\mathbb{R}^n, \mathbb{C})$ with

$$(10.29) \quad (\rho_w(g)f)(x) = e^{2\pi i w(z'+xy')} f(x + x').$$

This is the standard representation of G_n associated to the coadjoint orbit of linear functionals on the Lie algebra of G_n which send Z to w . Hence $L^2(\tau, w)$ is a corresponding isotypical component of ρ_w in $L^2(\tau)$. By (10.24) and (10.25), the multiplicity of ρ_w in $L^2(\tau)$ and $L^2(\tau, w)$ is

$$(10.30) \quad m_1 |w| d_1 \cdots m_n |w| d_n = |\Gamma| \dim V |w|^n.$$

10.3. Spectrum of Twisted Laplacians. Let $w \neq 0$. To determine the spectrum of the Laplacian Δ_w of a given left-invariant Riemannian metric on $L^2(\mathbb{R}^n, \mathbb{C})$ with respect to the representation ρ_w as in (10.29), we follow the discussion in the proof of Theorem 3.3 of [GoWi]: With respect to the given metric, there is an orthonormal basis

$$(10.31) \quad X'_1, \dots, X'_n, Y'_1, \dots, Y'_n, Z'$$

of the Lie algebra of G_n with $Z' = rZ$, $r = 1/|Z| > 0$, such that

$$(10.32) \quad [X'_j, X'_k] = [X'_j, Y'_k] = [Y'_j, Y'_k] = 0$$

for all $j \neq k$ and such that there are numbers $r_j > 0$ with

$$(10.33) \quad [X'_j, Y'_j] = r_j^2 Z.$$

The pull back of the metric under the automorphism Φ of G_n with

$$(10.34) \quad \Phi_*(r_j X_j) = X'_j, \quad \Phi_*(r_j Y_j) = Y'_j, \quad \Phi_*(Z) = Z$$

is the left-invariant Riemannian metric on G_n for which the fields

$$(10.35) \quad r_1 X_1, \dots, r_n X_n, r_1 Y_1, \dots, r_n Y_n, rZ$$

are orthonormal. Since $\Phi_*(Z) = Z$, $\rho_w \circ \Phi$ is still an irreducible unitary representation of G_n associated to the coadjoint orbit of linear functions on the Lie algebra of G_n which send Z to w , hence $\rho_w \circ \Phi$ is unitarily equivalent to ρ_w . In other words, we can assume without loss of generality that the given left-invariant Riemannian metric on G_n has an orthonormal basis as in (10.35). As for the Laplacian on $L^2(\mathbb{R}^n, \mathbb{C})$ with respect to ρ_w , we obtain

$$(10.36) \quad \Delta_w = - \sum_{1 \leq j \leq n} r_j^2 \frac{\partial^2}{\partial x_j^2} + 4\pi^2 w^2 (r^2 + \sum_{1 \leq j \leq n} x_j^2 r_j^2),$$

by (10.29). Now the Hermite functions

$$(10.37) \quad h_p(x) = \exp(x^2/2) \frac{\partial^{p_1+\dots+p_n}}{\partial x_1^{p_1} \dots \partial x_n^{p_n}} \exp(-x^2),$$

where $p = (p_1, \dots, p_n)$ runs over all n -tuples of non-negative integers, form an orthogonal basis of $L^2(\mathbb{R}^n, \mathbb{C})$ and satisfy

$$(10.38) \quad x_j^2 h_p - \frac{\partial^2 h_p}{\partial x_j^2} = (2p_j + 1) h_p.$$

It follows that the functions $f_p(x) = h_p(\sqrt{2\pi|w|}x)$ are an orthogonal basis of $L^2(\mathbb{R}^n, \mathbb{C})$ and that they satisfy

$$(10.39) \quad \Delta_w f_p = \lambda(w, p) f_p,$$

where

$$(10.40) \quad \lambda(w, p) := 4\pi^2 w^2 r^2 + 2\pi|w| \sum_{1 \leq j \leq n} (2p_j + 1) r_j^2.$$

Thus, by (10.30), the multiplicity of $\lambda(w, p)$ in $L^2(\tau, w)$ is equal to

$$(10.41) \quad d_1 \cdots d_n m_1 \cdots m_n |w|^n = |\Gamma| \dim V |w|^n,$$

when counted according to the n -tuples p .

In our application of the above in the proof of Theorem 10.47, we will vary the parameter $r = 1/|Z|$ of the metric, keeping

$$r_1 X_1, \dots, r_n X_n, r_1 Y_1, \dots, r_n Y_n$$

orthonormal and perpendicular to Z . Then the above functions f_p remain eigenfunctions of Δ_w in $L^2(\mathbb{R}^n, \mathbb{C})$ and the corresponding eigenvalues vary according to (10.40). Hence the eigensections in $L^2(\tau, w)$ corresponding to the above eigenfunctions f_p remain the same during this variation of the metric and the corresponding eigenvalues vary according to (10.40) as well.

10.4. η -Invariants for Heisenberg Manifolds. In this section, we study the η -invariant of the operator \bar{D}_t^+ as in (9.27). The solvable extension S of $N = G_n$ as in Chapter 9 and the connection ∇^E do not enter in this discussion. We recall though that $\Sigma_s^+ \simeq \Sigma_{\mathfrak{n}}$, where \mathfrak{n} denotes the Lie algebra of G_n and where Clifford multiplication with X in $\Sigma_{\mathfrak{n}}$ corresponds to Clifford multiplication with TX in Σ_s^+ , for all $X \in \mathfrak{n}$. This should be kept in mind, see e.g. (10.44).

Let Γ be a lattice in G_n of type d . It is clear from (10.3) that there is a smallest $s > 0$ such that $\zeta := \exp(s^2 Z)$ is contained in Γ and that ζ is a generator of the center of Γ . The automorphism $\Phi(x, y, z) = (sx, sy, s^2 z)$ of G_n maps $\exp Z$ to ζ , and, therefore, we may assume that

$$(10.42) \quad \zeta = \exp Z$$

generates the center of Γ . For any left-invariant Riemannian metric on G_n , $N = \Gamma \backslash G_n$ is a Riemannian submersion over a flat torus with closed geodesics as fibers, given as orbits of the one-parameter group generated by Z . By our normalization (10.42), the length of the fibers is given by $|Z|$.

Let τ be an irreducible unitary representation of Γ_d on a finite dimensional Hermitian vector space V as in Section 10.1 and extend τ by the trivial representation on $\Sigma_{\mathfrak{n}}$ to $\Sigma_{\mathfrak{n}} \otimes V$ as in Chapter 9. Recall from Section 10.1 that ζ acts by multiplication with $\exp(2\pi i c)$ for some $c = c(\tau) \in [0, 1) \cap \mathbb{Q}$ and that

$$(10.43) \quad \dim V = \delta(c, d) := m_1 \cdots m_n,$$

where $d = d(\Gamma)$ and m_j is the denominator of cd_j . In the notation of this chapter, and in terms of an orthonormal frame (E_j) of G_n , we study the unbounded operator

$$(10.44) \quad \bar{D}\sigma = \sum E_j \cdot E_j(\sigma),$$

in the Hilbert space $L^2(\tau)$ of measurable maps $G_n \rightarrow \Sigma_{\mathfrak{n}} \otimes V$ satisfying (9.22) which are square integrable over a fundamental domain of $\Gamma = \Gamma_d$ in the Heisenberg group G_n .

Before stating the next result, we recall the definition of the *Hurwitz zeta function*, for $c > 0$ and $\operatorname{Re} s > 1$ given by the infinite sum

$$(10.45) \quad \zeta_c(s) = \zeta(s, c) := \sum_{k \geq 0} (k + c)^{-s}.$$

We have $\zeta_1 = \zeta$, the *Riemann zeta function*. We also set $\zeta_0 := \zeta$. For each $c \geq 0$, ζ_c can be extended to a meromorphic function on the complex plane, defined for all $s \neq 1$, and with a simple pole at $s = 1$, where the residue is equal to 1.

It is maybe interesting to note that, for $0 < c < 1$,

$$(10.46) \quad \zeta_c(s) - \zeta_{1-c}(s) \quad \text{and} \quad \zeta_c(2s) + \zeta_{1-c}(2s)$$

are the eta and zeta function of the operator id/dt and $-d^2/dt^2$, respectively, on the Hermitian line bundle over $\mathbb{R}/2\pi\mathbb{Z}$ with twist $e^{-2\pi ic}$.

THEOREM 10.47. *Endow G_n with a left-invariant Riemannian metric, let Γ be a lattice in G_n such that $\zeta = \exp Z$ generates the center of Γ , and set $r := 1/|Z|$. Consider a Clifford module $\Sigma_{\mathfrak{n}} \otimes V$ as above and let $c = c(\tau)$. Then we have, for all $s \in \mathbb{C}$ with sufficiently large real part,*

$$(1) \quad \eta(\bar{D}, s) = |\Gamma| \dim V (2\pi r)^{-s} (\zeta_c(s - n) - \zeta_{1-c}(s - n))$$

if n is even

$$(2) \quad \eta(\bar{D}, s) = -|\Gamma| \dim V (2\pi r)^{-s} (\zeta_c(s - n) + \zeta_{1-c}(s - n))$$

if n is odd.

We conclude that, under the assumptions of the above theorem, the eta function of \bar{D} is holomorphic if n is even and is meromorphic with a simple pole at $s = n + 1$ if n is odd. We also see that the η -invariant $\eta(\bar{D}) = \eta(\bar{D}, 0)$ of \bar{D} only depends on n , the type of Γ , and c .

Proof of Theorem 10.47. The main argument in the proof is modeled along the lines of the proof of Proposition 4.1 of [DeSi]. For $w \equiv c$ modulo integers, we let

$$(10.48) \quad L^2(\tau, w) := \{\sigma \in L^2(\tau) : \sigma(x, y, z + t) = e^{2\pi i w t} \sigma(x, y, z)\}$$

and get an orthogonal decomposition

$$(10.49) \quad L^2(\tau) = \oplus_{w \equiv c} L^2(\tau, w),$$

where $L^2(\tau)$ is the Hilbert space of measurable maps $G_n \rightarrow \Sigma_{\mathfrak{n}} \otimes V$ satisfying (9.22) which are square integrable over a fundamental domain of $\Gamma = \Gamma_d$ in G_n as above. Since the spaces $L^2(\tau, w)$ are invariant under \bar{D} , the eta function of \bar{D} is the sum of the eta functions of the

restrictions of \bar{D} to the different $L^2(\tau, w)$. Thus we can consider the latter separately.

There are two cases, $w = 0$ and $w \neq 0$. As for $w = 0$, we note that $Z(\sigma) = 0$ for any $\sigma \in L^2(\tau, 0)$. Hence the unitary involution ω_0 of $L^2(\tau, 0)$ given by Clifford multiplication with irZ anti-commutes with \bar{D} . Hence the spectrum of \bar{D} is symmetric about 0, and, therefore, the eta function of \bar{D} on $L^2(\tau, 0)$ vanishes identically.

Suppose now that $w \neq 0$. We want to apply the results from the beginning of this chapter and note, to that end, that the spaces $L^2(\tau)$ and $L^2(\tau, w)$ here are isomorphic to the corresponding spaces in Section 10.2, tensored with Σ_n .

It follows from the discussion in Section 10.3 that, except for the determination of multiplicities, the particular lattice does not enter into the discussion. By what we explain in Section 10.3, we can assume that

$$(10.50) \quad r_1 X_1, r_1 Y_1, \dots, r_n X_n, r_n Y_n, rZ$$

is an orthonormal basis of the given left-invariant metric on G_n . Then (10.44) turns into

$$(10.51) \quad \bar{D}\sigma = \sum_{1 \leq j \leq n} r_j^2 (X_j \cdot X_j(\sigma) + Y_j \cdot Y_j(\sigma)) + r^2 Z \cdot Z(\sigma),$$

and (9.28) turns into

$$(10.52) \quad \bar{D}^2 \sigma = \Delta \sigma + \sum_{1 \leq j \leq n} r_j^4 X_j Y_j \cdot Z(\sigma),$$

where $\sigma \in L^2(\tau, w)$ is smooth.

We let ω_j , $1 \leq j \leq n$, be the unitary involutions on $\Sigma_n \otimes V$ and $L^2(\tau, w)$ given by Clifford multiplication with $ir_j^2 X_j Y_j$, respectively. Then

$$(10.53) \quad \Sigma_n = \oplus_{\varepsilon \in \{1, -1\}^n} \Sigma_\varepsilon,$$

where

$$(10.54) \quad \Sigma_\varepsilon = \{\sigma \in \Sigma_n : \omega_j \sigma = \varepsilon_j \sigma \text{ for all } 1 \leq j \leq n\}.$$

Now the unitary involutions ω_j commute with Δ . Thus on

$$(10.55) \quad L^2(\tau, w, \varepsilon) := \{\sigma \in L(\tau, w) : \sigma \text{ has values in } \Sigma_\varepsilon \otimes V\},$$

\bar{D}^2 has eigenvalues

$$(10.56) \quad \begin{aligned} \lambda(w, p, \varepsilon) &= \lambda(w, p) + 2\pi w(r_1^2 \varepsilon_1 + \dots + r_n^2 \varepsilon_n) \\ &= 4\pi^2 w^2 r^2 + 2\pi |w| \sum_{1 \leq j \leq n} (2p_j + 1 + \varepsilon_j \operatorname{sign} w) r_j^2 \end{aligned}$$

with multiplicity $2^n m_1 d_1 \cdots m_n d_n |w|^n$, where p runs over all n -tuples of non-negative integers, by (10.40) and (10.41). For all p , we have

$$(10.57) \quad \lambda(w, p, \varepsilon) \geq 4\pi^2 w^2 r^2 > 0.$$

Let W be an eigenspace of \bar{D}^2 in $L^2(\tau, w)$ for the eigenvalue λ , and recall from Section 10.3 that W is independent of the parameter r of the metric. Since \bar{D}^2 commutes with the involutions ω_j , W has an orthonormal basis consisting of eigensections of \bar{D}^2 such that each of them belongs to some $L^2(\tau, w, \varepsilon)$, where p and ε satisfy

$$(10.58) \quad S := 2\pi|w| \sum_{1 \leq j \leq n} (2p_j + 1 + \varepsilon_j \operatorname{sign} w) r_j^2 = \lambda - 4\pi^2 w^2 r^2,$$

by (10.56). Now Clifford multiplication by the unit vector rZ commutes with \bar{D}^2 and leaves the subspaces $L^2(\tau, w, \varepsilon)$ invariant, whereas Clifford multiplication by the unit vectors $r_j X_j$ and $r_j Y_j$ maps $L^2(\tau, w, \varepsilon)$ to $L^2(\tau, w, \delta)$ for $\delta \neq \varepsilon$. Hence using an orthonormal basis of eigensections of W as above, we see that the trace of \bar{D} on W is equal to an integral multiple $k2\pi wr$ of $2\pi wr$. On the other hand, the trace of \bar{D} on W is also equal to $l\sqrt{\lambda}$ for some integer l . Now 0 is not an eigenvalue of \bar{D}^2 on W , independently of $r > 0$. Hence k and l do not depend on r , and we get an equality of functions of $r \in (0, \infty)$,

$$(10.59) \quad k^2 4\pi^2 w^2 r^2 = l^2 (4\pi^2 w^2 r^2 + S)^2.$$

If $l = 0$, then the eigenvalues $\pm\sqrt{\lambda}$ of \bar{D} occur with equal multiplicity in W and, therefore, their contributions to the eta function of \bar{D} on $L^2(\tau, w)$ cancel. If $l \neq 0$, then $S = 0$, since S does not depend on r . But then, since $w \neq 0$, $p_j \geq 0$, and $\varepsilon_j = \pm 1$ for all j , we conclude that $\lambda(w, p, \varepsilon) = 4\pi^2 w^2 r^2$ and that

$$(10.60) \quad p_1 = \cdots = p_n = 0 \quad \text{and} \quad \varepsilon_1 = \cdots = \varepsilon_n = -\operatorname{sign} w$$

for $1 \leq j \leq n$. This will be denoted by $p = 0$ and $\varepsilon = -\operatorname{sign} w$.

To determine the contribution of the corresponding eigenspaces, we note that, by our identification $\Sigma_n = \Sigma_s^+$, Clifford multiplication by $irZ\omega_1 \cdots \omega_n$ is equal to the identity on Σ_n . Since Clifford multiplication with irZ commutes with Clifford multiplication with the ω_j , it leaves the subspaces Σ_ε invariant and acts by multiplication with $\varepsilon_1 \cdots \varepsilon_n$ on them. Now $Z(\sigma) = 2\pi i w \sigma$ for any σ in $L^2(\tau, w)$. Hence the eigenspace for \bar{D}^2 in $L^2(\tau, w)$ with eigenvalue $\lambda(w, 0, -\operatorname{sign} w) = 4\pi^2 w^2 r^2$ is an eigenspace of \bar{D} with eigenvalue

$$(10.61) \quad \begin{aligned} & 2\pi wr && \text{if } n \text{ is even,} \\ & -2\pi|w|r && \text{if } n \text{ is odd,} \end{aligned}$$

and dimension $m_1 \dots m_n d_1 \dots d_n |w|^n = |\Gamma| \dim V$. Thus, for all $s \in \mathbb{C}$ with sufficiently large real part,

$$(10.62) \quad \eta(\bar{D}, s) = |\Gamma| \dim V (2\pi r)^{-s} \sum_{w \equiv c, w \neq 0} \text{sign}(w) |w|^{n-s}$$

if n is even and

$$(10.63) \quad \eta(\bar{D}, s) = -|\Gamma| \dim V (2\pi r)^{-s} \sum_{w \equiv c, w \neq 0} |w|^{n-s}$$

if n is odd. □

We apply the results of this chapter to Dirac operators on homogeneous vector bundles over complex hyperbolic cusps of complex dimension n . Such cusps are homogeneous in the sense of Chapter 9, where the nilpotent Lie group is given by the Heisenberg group $N = G_{n-1}$ of dimension $2n - 1$ and $\Gamma \subseteq G_{n-1}$ is a lattice. In our formulas above we therefore need to substitute n by $n - 1$.

COROLLARY 10.64. *In the sense of Chapter 9, suppose that a complex hyperbolic cusp is determined by a lattice $\Gamma \subseteq G_{n-1}$ and that the homogeneous Dirac bundle over the cusp is given by unitary representations π_* of $\mathfrak{u}(n)$ and τ of Γ on a Hermitian vector space V . Assume that V is irreducible as a joint $\mathfrak{u}(n)$ and Γ module. Then the twist parameter c of τ is well defined and*

$$\lim_{t \rightarrow \infty} \eta(A_t^{\text{he},+}) = (-1)^n |\Gamma| \dim V (\zeta_c(1-n) + (-1)^n \zeta_{1-c}(1-n)).$$

Proof. We recall that $A_t = -D_t$, see (3.45). By Theorem 9.29, we have $\lim_{t \rightarrow \infty} \eta(D_t^{\text{he},+}) = \lim_{t \rightarrow \infty} \eta(\bar{D}_t^+)$. Now the operator \bar{D}_t^+ corresponds to the operator \bar{D} considered above, where the left-invariant metric on G_n comes from the cross section $\{t\} \times N$ in S . Since V is irreducible as a joint $\mathfrak{u}(n)$ and Γ module, it is a direct sum of isotypical irreducible representations of Γ as used in Theorem 10.47 so that the number c is the same for each summand. Hence Theorem 10.47 applies and shows that $\eta(\bar{D}_t^+)$ does not depend on t and that it is given by the formula in Corollary 10.64 □

EXAMPLE 10.65. Spinor bundles as in Example 9.23 are given by the trivial representation of $\mathfrak{u}(n)$ and classified by twists $\tau : \Gamma_d \rightarrow \{+1, -1\}$. Since $\tau(\zeta) = \pm 1$, we have $c = 0$ or $c = 1/2$. Hence the asymptotic high energy η -invariant of A_t^+ vanishes identically if n is odd. If n is even and $c = 0$, then

$$(10.66) \quad \lim_{t \rightarrow \infty} \eta(A_t^{\text{he},+}) = 2|\Gamma| \zeta(1-n),$$

which agrees with Proposition 4.1 in [DeSi] in the case $\Gamma = \Gamma_{(1, \dots, 1)}$ considered there (with a different choice of orientation). If n is even and $c = 1/2$, then

$$(10.67) \quad \lim_{t \rightarrow \infty} \eta(A_t^{\text{he},+}) = 2(2^{1-n} - 1)|\Gamma|\zeta(1-n),$$

where we use that $\zeta_{1/2}(s) = (2^s - 1)\zeta(s)$. Recall also that

$$(10.68) \quad \zeta(1-n) = -B_n/n,$$

where B_n denotes the n -th Bernoulli number.

11. LOW ENERGY η -INVARIANTS

11.1. General Remarks and Computations. We return to the situation and notation considered in Chapter 9 and let $E = \mathcal{P} \times_{\beta} (\Sigma_{\mathfrak{s}} \otimes V)$ be a homogeneous Dirac bundle over S . As in Chapter 10, we view sections of E^+ as smooth maps $\sigma : S \rightarrow \Sigma_{\mathfrak{n}} \otimes V$.

The vector field T is left-invariant and a global unit normal field along the hypersurfaces $N_t := \{t\} \times N$ of S . In accordance with this, we choose frames (X_1, \dots, X_m) of S to be left-invariant and orthonormal with $X_1 = T$. Then X_2, \dots, X_m are tangent to the hypersurfaces N_t .

Let $\Gamma \subseteq N$ be a lattice, $\tau : \Gamma \rightarrow V$ be a unitary representation, and E_{τ} be the induced Dirac bundle over $\Gamma \backslash S = \mathbb{R} \times (\Gamma \backslash N)$. Then we have, for any $t \in \mathbb{R}$, the orthogonal decomposition

$$(11.1) \quad L^2(N_t, E_{\tau}^+) = H^{\text{le},+}(A_t) \oplus H^{\text{he},+}(A_t),$$

where $H^{\text{le},+}(A_t)$ is the space of constant maps $N_t \rightarrow \Sigma_{\mathfrak{n}} \otimes V$, compare Chapter 7 and, in particular, (8.2).

PROPOSITION 11.2. *If V is irreducible as a joint \mathfrak{k} and Γ module and τ is non-trivial, then the low energy spaces $H_t^{\text{le},+}(A_t)$ are trivial and, therefore, $\eta(A_t^{\text{he},+}) = 0$, for all $t \in \mathbb{R}$. \square*

Thus the low energy η -invariant can only be non-trivial when τ is trivial. We refer to this as the *untwisted case* and assume for the rest of this section that we are in this case, whether V is irreducible as a \mathfrak{k} module or not. Then the space $H^{\text{le},+}(A_t)$ is isomorphic to $\Sigma_{\mathfrak{n}} \otimes V$, by identifying constant maps with their respective values.

For $\sigma \in H^{\text{le},+}(A_t)$ and with A_X as in (9.14), we have

$$(11.3) \quad D_t^+ \sigma = \sum_{2 \leq j \leq m} T X_j \cdot \beta_*(A_{X_j}) \sigma + \frac{\kappa}{2} \sigma,$$

by (9.26), where we recall our convention $X_1 = T$. Our objective in this chapter is the η -invariant of D_t^+ on $H^{\text{le}}(A_t^+)$. We view elements of $H^{\text{le},+}(A_t)$ as constant maps on S . Then $H^{\text{le},+}(A_t)$ becomes independent

of Γ and t . By (11.3), D_t^+ does not depend on t either. As a shorthand, we will write

$$(11.4) \quad H_N^{\text{le}} \text{ for } H^{\text{le},+}(A_t) \quad \text{and} \quad D_N^{\text{le}} \text{ for } D_t^+ \text{ on } H_N^{\text{le}}.$$

Recall that $\beta_* = \hat{\alpha}_* \otimes \text{id} + \text{id} \otimes \pi_*$, by (9.15), and that

$$(11.5) \quad \hat{\alpha}_*(A_X) = \frac{1}{2} \sum_{1 \leq j < k \leq m} \langle \nabla_X X_j, X_k \rangle X_j X_k.$$

where $X_j X_k$ stands for Clifford multiplication by $X_j X_k$. With our convention $X_1 = T$, (11.5) turns into

$$(11.6) \quad \hat{\alpha}_*(A_X) = \frac{1}{2} T \nabla_X T + \frac{1}{2} \sum_{2 \leq j < k \leq m} \langle \nabla_X X_j, X_k \rangle X_j X_k.$$

It follows that (2.8) and (2.9) define the Dirac structure on E associated to the Riemannian metric of N . As for the first term on the right hand side of (11.6), we note that

$$(11.7) \quad \sum_{j \geq 2} T \nabla_{X_j} T = -\kappa.$$

Choose an orthonormal frame (X_2, \dots, X_m) of \mathfrak{n} such that $[X_j, X_k]$ is contained in the linear hull of the X_l with $l < \min\{j, k\}$. On H_N^{le} , we then have

$$(11.8) \quad D_N^{\text{le}} - \sum_{j \geq 2} T X_j \otimes \pi_*(A_{X_j}) = \frac{1}{2} \sum_{j \geq 2 \leq k < l} T X_j \langle \nabla_{X_j} X_k, X_l \rangle X_k X_l,$$

by (11.3), (11.6), and (11.7). By the Koszul formula and since $[X_j, X_k]$ is perpendicular to X_l for $k < l$, the right hand side is equal to

$$(11.9) \quad -\frac{1}{4} \sum_{j \geq 2 \leq k < l} T X_j (\langle X_j, [X_k, X_l] \rangle + \langle X_k, [X_j, X_l] \rangle) X_k X_l.$$

Now $[X_j, X_l]$ is a linear combination of the X_k for $k < l$, hence the terms in (11.9) can be rewritten as

$$(11.10) \quad -\frac{1}{4} \sum_{2 \leq k < l} T [X_k, X_l] X_k X_l - \frac{1}{4} \sum_{2 \leq j, l} T X_j [X_j, X_l] X_l.$$

Since X_2 is central, these terms are equal to

$$(11.11) \quad \frac{1}{4} \sum_{2 < j < k} T [X_j, X_k] X_j X_k = \frac{1}{8} \sum_{j, k > 2} T [X_j, X_k] X_j X_k.$$

In conclusion,

$$(11.12) \quad D_N^{\text{le}} = \frac{1}{8} \sum_{j, k > 2} T [X_j, X_k] X_j X_k + \sum_{j \geq 2} T X_j \otimes \pi_*(A_{X_j}),$$

our main formula regarding D_N^{le} .

For general nilpotent Lie groups N , it will probably be hard to get an explicit formula for the low energy η -invariant. We expect that a simple explicit formula does exist for groups of Heisenberg type as in [Ka]. So far, we have been able to achieve this for the standard Heisenberg groups as in Chapter 10; their low energy η -invariant will be the topic in Section 11.2 below. Recall that cusps of complex hyperbolic manifolds correspond to the case where N is the standard Heisenberg group and that cusps of quaternionic and Cayley hyperbolic manifolds also give rise to groups of Heisenberg type.

Suppose now that N is of Heisenberg type. That is, we are given an orthogonal decomposition

$$(11.13) \quad \mathfrak{n} = \mathfrak{z} + \mathfrak{x},$$

where \mathfrak{z} is contained in the center of \mathfrak{n} , and a linear map J from \mathfrak{z} into the space of skew-symmetric endomorphisms of \mathfrak{x} such that the Clifford relations hold,

$$(11.14) \quad J_{Z_1} J_{Z_2} + J_{Z_2} J_{Z_1} + 2\langle Z_1, Z_2 \rangle = 0,$$

for all $Z_1, Z_2 \in \mathfrak{z}$. Moreover, the Lie brackets of vectors in \mathfrak{x} are contained in \mathfrak{z} and satisfy, by definition,

$$(11.15) \quad \langle [X_1, X_2], Z \rangle = 2c \langle J_Z X_1, X_2 \rangle,$$

for all $X_1, X_2 \in \mathfrak{x}$ and $Z \in \mathfrak{z}$, where $c > 0$ is a fixed constant. Compare with (2.39)–(2.40), where we have the case $\dim \mathfrak{z} = 1$ and $c = 1$.

The constant $c > 0$ in (11.15) is arbitrary, and the derivation W is defined to have \mathfrak{x} and \mathfrak{z} as eigenspaces with $-c$ and $-2c$ as respective eigenvalues. This normalization has the following amazing formula as a consequence.

LEMMA 11.16. *For all $Z \in \mathfrak{z}$ and $X \in \mathfrak{x}$, we have*

$$R(Z, X) = R(J_Z X, T).$$

REMARKS 11.17. 1) In Lemma 11.16, we consider $R(Z, X)$ as an endomorphism of \mathfrak{s} , the tangent space of S at the neutral element, but note that $R(Z, X)$ acts also on associated bundles.

2) If N is the standard Heisenberg group of dimension $2n + 1$, then S is isometric to the complex hyperbolic space $\mathbb{C}H^{n+1}$ of dimension $2n + 2$ with sectional curvature in $[-4c^2, -c^2]$ and complex structure J with $JT = Z$ and such that J coincides with J_Z on N . In this case, the equation in Lemma 11.16 is a special case of the more general $R(JU, V) = -R(U, JV)$ which says that the curvature tensor of $\mathbb{C}H^{n+1}$ is a differential form of type $(1, 1)$.

Proof of Lemma 11.16. By straightforward computations, using (9.7), (9.8), (11.14), and (11.15), \square

Let $Z \in \mathfrak{z}$ with $|Z| = 1$. Then J_Z is an orthogonal complex structure on \mathfrak{x} . In particular, the dimension of \mathfrak{x} is even, and we denote it by $2n$. Moreover, there is an orthonormal basis (X_1, \dots, X_{2n}) of \mathfrak{x} such that $J_Z X_{2j-1} = X_{2j}$, for $1 \leq j \leq n$. Given any such basis, set

$$(11.18) \quad D_Z := \frac{c}{2} \sum_{1 \leq j \leq n} TZX_{2j-1}X_{2j} + TZ \otimes \pi_*(A_Z).$$

Observe that, for any orthonormal basis (Y_1, \dots, Y_{2n}) of \mathfrak{x} ,

$$(11.19) \quad \begin{aligned} D_Z &= \frac{1}{8} \sum_{j,k} \langle [Y_j, Y_k], Z \rangle TZY_j Y_k + TZ \otimes \pi_*(A_Z) \\ &= \frac{c}{4} \sum_{j,k} \langle J_Z Y_j, Y_k \rangle TZY_j Y_k + TZ \otimes \pi_*(A_Z) \\ &= \frac{c}{2} \sum_{j < k} \langle J_Z Y_j, Y_k \rangle TZY_j Y_k + TZ \otimes \pi_*(A_Z) \end{aligned}$$

In what follows, let $\{A, B\} := AB + BA$.

LEMMA 11.20. *For any $X \in \mathfrak{x}$, we have*

$$\{D_Z, TX \otimes \pi_*(A_X) + TJ_Z X \otimes \pi_*(A_{J_Z X})\} = 0.$$

Proof. We have $cA_X = R(X, T)$, by (9.8), and hence

$$(11.21) \quad c\{TZXJ_Z X, TX \otimes \pi_*(A_X)\} = 2ZJ_Z X \otimes \pi_*(R(X, T)).$$

By substituting $J_Z X$ for X in (11.21), we obtain

$$(11.22) \quad c\{TZXJ_Z X, TJ_Z X \otimes \pi_*(A_{J_Z X})\} = -2ZX \otimes \pi_*(R(J_Z X, T)).$$

We also have $[Z, X] = 0$, hence $[A_Z, A_X] = R(Z, X)$. Furthermore, $R(Z, X) = R(J_Z X, T)$, by Lemma 11.16, hence

$$(11.23) \quad \{TZ \otimes \pi_*(A_Z), TX \otimes \pi_*(A_X)\} = ZX \otimes \pi_*(R(J_Z X, T)).$$

By substituting $J_Z X$ for X in (11.23), we obtain

$$(11.24) \quad \{TZ \otimes \pi_*(A_Z), TJ_Z X \otimes \pi_*(A_{J_Z X})\} = -ZJ_Z X \otimes \pi_*(R(X, T)).$$

Moreover, we have

$$(11.25) \quad \begin{aligned} \{TZYJ_Z Y, TX \otimes \pi_*(A_X)\} \\ = \{TZYJ_Z Y, TJ_Z X \otimes \pi_*(A_{J_Z X})\} = 0, \end{aligned}$$

for all $Y \in \mathfrak{x}$ perpendicular to X and $J_Z X$. Now we may assume that X is of norm 1. Then there is an orthonormal basis (X_1, \dots, X_{2n}) of \mathfrak{x} such that $J_Z X_{2j-1} = X_{2j}$, for $1 \leq j \leq n$, and such that $X = X_1$. By (11.25), the terms of D_Z involving $TZX_{2j-1}X_{2j}$, $j \geq 2$, do not contribute to

the anti-commutator $\{D_Z, TX \otimes \pi_*(A_X) + TJ_ZX \otimes \pi_*(A_{J_ZX})\}$. The four remaining terms cancel pairwise, by (11.21)–(11.24). \square

For an orthonormal basis (X_1, \dots, X_{2n}) of \mathfrak{x} , set

$$(11.26) \quad D_{\mathfrak{x}} := TX_1 \otimes \pi_*(A_{X_1}) + \dots + TX_{2n} \otimes \pi_*(A_{X_{2n}}),$$

and note that $D_{\mathfrak{x}}$ does not depend on the choice of (X_1, \dots, X_{2n}) .

REMARK 11.27. If $\mathfrak{z} = 0$, then \mathfrak{n} is Abelian and we are in the case of real hyperbolic spaces or cusps, respectively, and we get $D_N^{\text{le}} = D_{\mathfrak{x}}$ on H_N^{le} . The contribution of cusps in the case $\dim N = 1$ follows easily from the more general discussion in [BB1]. If $\dim N \geq 2$, then the arguments in the proof of Theorem 9.31 apply and show that the low energy η -invariant vanishes.

LEMMA 11.28. *For any unit vectors $Z \in \mathfrak{z}$,*

$$\{D_Z, D_{\mathfrak{x}}\} = 0.$$

Proof. Apply Lemma 11.20, using an orthonormal basis (X_1, \dots, X_{2n}) of \mathfrak{x} with $J_Z X_{2j-1} = X_{2j}$, for $1 \leq j \leq n$. \square

Assume from now on that $\mathfrak{z} \neq 0$, compare Remark 11.27. For an orthonormal basis (Z_1, \dots, Z_{ℓ}) of \mathfrak{z} , set

$$(11.29) \quad D_{\mathfrak{z}} := D_{Z_1} + \dots + D_{Z_{\ell}},$$

and note that $D_{\mathfrak{z}}$ does not depend on the choice of (Z_1, \dots, Z_{ℓ}) .

COROLLARY 11.30. *On H_N^{le} , we have*

$$D_N^{\text{le}} = D_{\mathfrak{z}} + D_{\mathfrak{x}} \quad \text{and} \quad \{D_{\mathfrak{z}}, D_{\mathfrak{x}}\} = 0. \quad \square$$

PROPOSITION 11.31. *On H_N^{le} , we have*

- (1) $\ker(D_N^{\text{le}}) = \ker D_{\mathfrak{z}} \cap \ker D_{\mathfrak{x}},$
- (2) $\eta(D_N^{\text{le}}) = \eta(D_{\mathfrak{z}}) = \eta(D_{\mathfrak{z}}|_{\ker D_{\mathfrak{x}}}).$

Proof. By Corollary 11.30, (1) is clear and

$$\begin{aligned} \eta(D_N^{\text{le}}) &= \eta(D_{\mathfrak{z}}|_{\ker D_{\mathfrak{x}}}) + \eta(D_{\mathfrak{x}}|_{\ker D_{\mathfrak{z}}}), \\ \eta(D_{\mathfrak{z}}|_{\ker D_{\mathfrak{x}}}) &= \eta(D_{\mathfrak{z}}), \\ \eta(D_{\mathfrak{x}}|_{\ker D_{\mathfrak{z}}}) &= \eta(D_{\mathfrak{x}}). \end{aligned}$$

Now $D_{\mathfrak{x}}$ anticommutes with the involution TZ_1 of $\Sigma^+ \otimes V$, hence $\eta(D_{\mathfrak{x}}) = 0$, hence (2). \square

11.2. Contribution of Complex Hyperbolic Cusps. As in Section 2.3, we represent complex hyperbolic space $\mathbb{C}H^n$ by the symmetric pair (G, K) , where $G = \mathrm{SU}(1, n)$ and $K = \mathrm{S}(\mathrm{U}(1) \times \mathrm{U}(n)) \cong \mathrm{U}(n)$, and recall the decomposition $\mathfrak{su}(1, n) = \mathfrak{k} + \mathfrak{p}$, where $\mathfrak{k} \cong \mathfrak{u}(n)$. We decompose matrices $X \in \mathfrak{su}(1, n)$ correspondingly and write

$$(11.32) \quad X = X^{\mathfrak{k}} + X^{\mathfrak{p}} \quad \text{with } X^{\mathfrak{k}} \in \mathfrak{k} \text{ and } X^{\mathfrak{p}} \in \mathfrak{p}.$$

We recall that, after identifying \mathfrak{p} with the tangent space of $\mathbb{C}H^n$ at the point fixed by $\mathrm{U}(n)$ as usual, we have

$$(11.33) \quad R(X, Y)Z = -[[X, Y], Z],$$

for all $X, Y, Z \in \mathfrak{p}$, where the Lie brackets on the right hand side are taken in $\mathfrak{su}(1, n)$.

Let $\mathfrak{n} \subset \mathfrak{s} \subset \mathfrak{su}(1, n)$ be as in Section 2.3. Recall that we identify $\mathbb{C}H^n$ in (2.37) also with the solvable Lie subgroup S of $\mathrm{SU}(1, n)$ corresponding to \mathfrak{s} , endowed with the left-invariant Riemannian metric such that the isomorphism $\mathfrak{s} \rightarrow \mathfrak{p}$, $X \mapsto X^{\mathfrak{p}}$, preserves scalar products. Furthermore, the nilpotent Lie subgroup N of S corresponding to \mathfrak{n} is isomorphic to the standard Heisenberg group. Thus we are in the situation of Section 11.1, where N is the standard Heisenberg group and where $c = 1$ in (11.15).

Let $X \in \mathfrak{n}$. By (2.33) and (2.35), we have $[T, X] = -WX$ and hence

$$(11.34) \quad [T, X^{\mathfrak{p}}] = -(WX)^{\mathfrak{k}},$$

where we use the notation from (11.32). Using (9.8) and (11.33), we obtain therefore that

$$(11.35) \quad A_{WX}Y = R(T, X^{\mathfrak{p}})Y^{\mathfrak{p}} = -[[T, X^{\mathfrak{p}}], Y^{\mathfrak{p}}] = [(WX)^{\mathfrak{k}}, Y^{\mathfrak{p}}].$$

Since W is invertible, we conclude that, for any $X \in \mathfrak{n}$,

$$(11.36) \quad A_X = X^{\mathfrak{k}}.$$

With α as in (2.30), we let $\hat{\alpha}_* : \mathfrak{u}(n) \rightarrow \mathfrak{u}(\Sigma)$ be the composition of the differential α_* of α with the differential of the spinor representation of $\mathfrak{so}(\mathfrak{p}) \simeq \mathfrak{spin}(\mathfrak{p})$ on $\Sigma := \Sigma_{\mathfrak{p}}$. Following Chapter 9, we choose $K = \mathrm{U}(n)$ and let π_* be a unitary representation of $\mathfrak{u}(n)$ on a Hermitian vector space V . We assume that there exists a unitary representation β of $K = \mathrm{U}(n)$ on $\Sigma \otimes V$ satisfying (9.15) and get the associated Dirac bundle E over $\mathbb{C}H^n$. Clifford multiplication by the complex volume form $\omega_{\mathbb{C}}$ determines a super-symmetry $E = E^+ \oplus E^-$, and this super-symmetry is induced by the corresponding decomposition $\Sigma = \Sigma^+ \oplus \Sigma^-$.

To distinguish it from multiplication with i in $\mathbb{C}^n \simeq \mathfrak{p}$, we denote the complex structure in $\mathrm{Cl}(\mathfrak{p})$ by $\sqrt{-1}$. With the corresponding changes

in notation, we follow Section 2.2 and set

$$(11.37) \quad \omega_j := \sqrt{-1} X_j^p Y_j^p, \quad 1 \leq j \leq n,$$

where $X_1 = T, Y_1 = Z, X_2, Y_2, \dots, X_n, Y_n$ are as in (2.41). By the discussion in Section 2.2, we have

$$(11.38) \quad \Sigma^+ = \oplus_{\epsilon \in \{-1, 1\}^{n-1}} \Sigma_\epsilon^+,$$

where

$$(11.39) \quad \Sigma_\epsilon^+ := \{\sigma \in \Sigma^+ : \omega_j \sigma = \epsilon_j \sigma \text{ for } 2 \leq j \leq n\}.$$

Since X_j commutes with ω_k for $k \neq j$ and anti-commutes with ω_j , all the subspaces Σ_ϵ are isomorphic. In particular, for all $\epsilon \in \{-1, 1\}^{n-1}$,

$$(11.40) \quad \dim \Sigma_\epsilon^+ = \dim \Sigma^+ / \text{card}\{-1, 1\}^{n-1} = 1.$$

For any $\epsilon \in \{-1, 1\}^{n-1}$, let $\nu(\epsilon) \in \{0, \dots, n-1\}$ be the number of j with $\epsilon_j = -1$, for $2 \leq j \leq n$. Then

$$(11.41) \quad \Sigma^+ = \oplus_k \Sigma_k^+, \quad \text{where} \quad \Sigma_k^+ := \oplus_{\nu(\epsilon)=k} \Sigma_\epsilon^+.$$

By definition, $\omega_{\mathbb{C}}$ acts as identity on Σ^+ , hence $\omega_1 = \omega_2 \cdots \omega_n$ on Σ^+ . Therefore

$$(11.42) \quad \Sigma_{\text{even}}^+ = \oplus_{k \text{ even}} \Sigma_k^+ \quad \text{and} \quad \Sigma_{\text{odd}}^+ = \oplus_{k \text{ odd}} \Sigma_k^+$$

are the eigenspaces of ω_1 for the eigenvalues 1 and -1 , respectively. In passing we note that the left side of (11.41) gives the decomposition of Σ^+ into irreducible representations of the stabilizer of T in $U(n)$, by work of Camporesi and Pedon, see [CaPe, Lemma 3.1].

We recall that the complexification of $\mathfrak{u}(n)$ is $\mathfrak{gl}(n, \mathbb{C})$, where the complex structure of $\mathfrak{gl}(n, \mathbb{C})$ is given by multiplication of matrix coefficients with i . The space $\mathfrak{h} \subseteq \mathfrak{gl}(n, \mathbb{C})$ of diagonal matrices is a Cartan subalgebra of $\mathfrak{gl}(n, \mathbb{C})$, and the roots

$$(11.43) \quad \rho_j(\text{diag}(h_1, \dots, h_n)) := h_j$$

constitute a basis of \mathfrak{h}^* . The associated Weyl group \mathcal{W} of automorphisms of \mathfrak{h} leaves the set $\{\rho_1, \dots, \rho_n\}$ invariant and acts on it as the (complete) group of permutations. As usual, we choose

$$(11.44) \quad \{h = \text{diag}(h_1, \dots, h_n) : h_j \in \mathbb{R}, h_1 > h_2 > \cdots > h_n\}$$

as positive Weyl chamber. The corresponding set of positive roots of $\mathfrak{gl}(n, \mathbb{C})$ is given by

$$(11.45) \quad \Delta^+ = \{\rho_j - \rho_k : 1 \leq j < k \leq n\}.$$

Irreducible complex representations of $\mathfrak{u}(n)$ are classified by their *highest weight* $\lambda = \sum \lambda_j \rho_j$, where λ is *dominant*, that is, $\lambda_1 \geq \lambda_2 \geq \cdots \geq$

λ_n , and algebraically integral, that is, $\lambda_i - \lambda_j \in \mathbb{Z}$ for all i, j . The dimension of the corresponding representation space V_λ is

$$(11.46) \quad \dim V_\lambda = \prod_{j < k} \frac{k - j - \lambda_k + \lambda_j}{k - j},$$

by the Weyl dimension formula. The irreducible representation with highest weight λ is induced by a representation of $U(n)$ if all the λ_j are integral. The representation α as above is the irreducible representation of $U(n)$ with highest weight $(2, 1, \dots, 1)$ (and complex dimension n).

For the discussion of $\hat{\alpha}_*$, we identify $\mathfrak{p} = \mathbb{R}^{2n}$ and $\Sigma = \Sigma_{2n}$. We let (e_1, \dots, e_{2n}) be the standard basis of $\mathbb{R}^{2n} \simeq \mathbb{C}^n$ with $e_{2j} = ie_{2j-1}$, $1 \leq j \leq n$, and denote the complex structure of Σ_{2n} by $\sqrt{-1}$ as above. For $h = (it_1, \dots, it_n)$ in $\mathfrak{h} \cap \mathfrak{u}(n)$, we get

$$(11.47) \quad \begin{aligned} \hat{\alpha}_*(h) &= \frac{1}{4} \sum_{1 \leq j \leq 2n} e_j (it_j e_j + (it_1 + \dots + it_n) e_j) \\ &= -\frac{1}{2} \sqrt{-1} \sum_{1 \leq j \leq n} (t_1 + \dots + 2t_j + \dots + t_n) \omega_j. \end{aligned}$$

by the Parthasarathy formula [Pa, Lemma 2.1], where e_j and ω_j stand for Clifford multiplication by e_j and ω_j , respectively. Hence the subspaces Σ_ε of Σ_{2n} as in Section 2.2 are weight spaces. For $0 \leq l \leq n$, we let V_l be the sum over all Σ_ε such that l is the number of j with $\varepsilon_j = -1$, that is, $\varepsilon_1 + \dots + \varepsilon_n = n - 2l$. Then V_l is the irreducible representation of $\mathfrak{u}(n)$ with highest weight

$$(11.48) \quad \lambda_1 = \dots = \lambda_l = l - \frac{n-1}{2} > \lambda_{l+1} = \dots = \lambda_n = l - \frac{n+1}{2},$$

and is of dimension $\binom{n}{l}$, in agreement with Weyl's dimension formula.

As an example, we discuss differential forms. Since α is the irreducible representation with maximal weight $(2, 1, \dots, 1)$, the bundles of differential forms of type $(p, 0)$ and $(0, q)$ are associated to the irreducible representations β of $U(n)$ with maximal weights

$$(11.49) \quad \lambda_1 = \dots = \lambda_{n-p} = -p > \lambda_{n-p+1} = \dots = \lambda_n = -(p+1)$$

and

$$(11.50) \quad \lambda_1 = \dots = \lambda_q = q+1 > \lambda_{q+1} = \dots = \lambda_n = q,$$

respectively. We see that the sum of the bundles of differential forms of type $(0, q)$, $0 \leq q \leq n$, is given by $\Sigma \otimes V_n$, where V_n is as above. That is, π_* is the one-dimensional irreducible representation of $\mathfrak{u}(n)$ with highest weight $\lambda_j = (n+1)/2$, $1 \leq j \leq n$.

REMARK 11.51. From (11.48), we see that $\hat{\alpha}_*$ comes from a representation of $U(n)$ if n is odd, and then the spinor bundle of $\mathbb{C}H^n$ descends to quotients of $\mathbb{C}H^n$ by discrete subgroups of $SU(1, n)$. On the other hand, if $\widetilde{SU}(1, n)$ denotes the non-trivial twofold cover of $SU(1, n)$, then $\hat{\alpha}_*$ comes from a representation of the corresponding twofold cover $\tilde{U}(n)$ of $U(n)$, for all n . Hence, if the discrete subgroup of $SU(n)$ under consideration admits a lift into $\widetilde{SU}(1, n)$, then the spinor bundle also descends to the corresponding quotient of $\mathbb{C}H^n$. A similar remark applies to β_* .

We note that $D_{\mathfrak{r}}$ is an odd operator with respect to the grading

$$(11.52) \quad \Sigma^+ \otimes V_\pi = (\Sigma_{\text{even}}^+ \otimes V_\pi) \oplus (\Sigma_{\text{odd}}^+ \otimes V_\pi),$$

whereas $D_{\mathfrak{z}} = D_Z$ is an even operator.

THEOREM 11.53. *With $H^k := \ker D_{\mathfrak{r}} \cap (\Sigma_k^+ \otimes V_\pi)$ and $b_k := \dim H^k(\pi)$, for $0 \leq k \leq n-1$, we have*

$$(1) \quad \ker D_{\mathfrak{r}} = \oplus H^k,$$

$$(2) \quad b_k = (n-1)! \dim V_\pi \prod_{\substack{1 \leq j \leq n \\ j \neq k+1}} |\lambda_j - \lambda_{k+1} + k + 1 - j|^{-1},$$

$$(3) \quad D_Z|_{H^k} = (-1)^k (2k - 2\lambda_{k+1} - n + 1)/2.$$

Proof. Our proof relies on Kostant's theorem, see [Ko] or Theorem 4.139 in [KnVo]. We start by describing an explicit model of Σ , compare Chapter 5 of [Wu]. For $1 \leq j \leq n$, let

$$(11.54) \quad F_j := \frac{1}{2}(X_j^{\mathfrak{p}} - \sqrt{-1}Y_j^{\mathfrak{p}}) \quad \text{and} \quad \bar{F}_j := \frac{1}{2}(X_j^{\mathfrak{p}} + \sqrt{-1}Y_j^{\mathfrak{p}}).$$

As elements of $\mathbb{C}l(\mathfrak{p})$, they satisfy

$$(11.55) \quad F_j F_j = \bar{F}_j \bar{F}_j = 0, \quad \bar{F}_j F_j = -F_j \bar{F}_j - 1,$$

for $1 \leq j \leq n$, and

$$(11.56) \quad F_j F_k = -F_k F_j, \quad \bar{F}_j \bar{F}_k = -\bar{F}_k \bar{F}_j, \quad F_j \bar{F}_k = -\bar{F}_k F_j,$$

for $1 \leq j \neq k \leq n$. We identify Σ with the left ideal in the Clifford algebra generated by $\bar{F} = \bar{F}_1 \cdots \bar{F}_n$. Then the monomials $F_I \bar{F}$ over all $0 \leq k \leq n$ and multi-indices $I = \{i_1, \dots, i_k\}$ with $i_1 < \dots < i_k$ constitute a basis of Σ . The relations (11.55) and (11.56) determine an isomorphism $\Sigma \simeq \Lambda(\mathbb{C}^n)$, where \mathbb{C}^n is spanned by F_1, \dots, F_n . We have

$$(11.57) \quad \omega_j \cdot F_I \bar{F} = \begin{cases} F_I \bar{F} & \text{if } j \notin I, \\ -F_I \bar{F} & \text{if } j \in I. \end{cases}$$

so that, under the identification $\Sigma \simeq \Lambda(\mathbb{C}^n)$,

$$(11.58) \quad \Sigma_k^+ \simeq \begin{cases} \Lambda^k(\mathbb{C}^{n-1}) & \text{if } k \text{ is even,} \\ F_1 \wedge \Lambda^k(\mathbb{C}^{n-1}) \simeq \Lambda^k(\mathbb{C}^{n-1}) & \text{if } k \text{ is odd,} \end{cases}$$

where \mathbb{C}^{n-1} is spanned by F_2, \dots, F_n .

Recall that, by complexification, π_* induces a representation of $\mathfrak{gl}(n, \mathbb{C})$. Following the exposition in [LaMi, §IV.8], we set

$$(11.59) \quad \begin{aligned} \mathcal{D}_{\mathfrak{r}} &:= \frac{1}{2} \sum_{2 \leq j \leq n} T(X_j^{\mathfrak{p}} - \sqrt{-1}Y_j^{\mathfrak{p}}) \otimes \pi_*(X_j^{\mathfrak{k}} + iY_j^{\mathfrak{k}}) \\ &= 2 \sum_{2 \leq j \leq n} TF_j \otimes \pi_*(E_{1j}), \end{aligned}$$

$$(11.60) \quad \begin{aligned} \bar{\mathcal{D}}_{\mathfrak{r}} &:= \frac{1}{2} \sum_{2 \leq j \leq n} T(X_j^{\mathfrak{p}} + \sqrt{-1}Y_j^{\mathfrak{p}}) \otimes \pi_*(X_j^{\mathfrak{k}} - iY_j^{\mathfrak{k}}) \\ &= -2 \sum_{2 \leq j \leq n} T\bar{F}_j \otimes \pi_*(E_{j1}), \end{aligned}$$

where we note that factors $\sqrt{-1}$ on the left and i on the right of \otimes multiply to -1 in the tensor product. Using (11.36), we have

$$(11.61) \quad D_{\mathfrak{r}} = \mathcal{D}_{\mathfrak{r}} + \bar{\mathcal{D}}_{\mathfrak{r}}, \quad \bar{\mathcal{D}}_{\mathfrak{r}} = \mathcal{D}_{\mathfrak{r}}^*, \quad \text{and} \quad \mathcal{D}_{\mathfrak{r}}\mathcal{D}_{\mathfrak{r}} = \bar{\mathcal{D}}_{\mathfrak{r}}\bar{\mathcal{D}}_{\mathfrak{r}} = 0.$$

Moreover,

$$(11.62) \quad \mathcal{D}_{\mathfrak{r}}(\Sigma_k^+ \otimes V_{\pi}) \subset \Sigma_{k+1}^+ \otimes V_{\pi} \quad \text{and} \quad \bar{\mathcal{D}}_{\mathfrak{r}}(\Sigma_k^+ \otimes V_{\pi}) \subset \Sigma_{k-1}^+ \otimes V_{\pi}.$$

Hence $\ker D_{\mathfrak{r}}$ is equal to the space of $\mathcal{D}_{\mathfrak{r}}$ -harmonic cocycles of the cochain complex

$$(11.63) \quad \dots \xrightarrow{\mathcal{D}_{\mathfrak{r}}} \Sigma_{k-1}^+ \otimes V_{\pi} \xrightarrow{\mathcal{D}_{\mathfrak{r}}} \Sigma_k^+ \otimes V_{\pi} \xrightarrow{\mathcal{D}_{\mathfrak{r}}} \Sigma_{k+1}^+ \otimes V_{\pi} \xrightarrow{\mathcal{D}_{\mathfrak{r}}} \dots$$

This shows the first assertion of the theorem and that $\ker D_{\mathfrak{r}}$ is isomorphic to the cohomology of the complex. Moreover, under the above identification $\Sigma^+ = \Lambda(\mathbb{C}^{n-1})$, we have

$$(11.64) \quad \mathcal{D}_{\mathfrak{r}}(\omega \otimes v) = 2 \sum_{2 \leq j \leq n} (F_j \wedge \omega) \otimes \pi_*(E_{1j})v.$$

Following the notation in [KnVo], we consider the subalgebras \mathfrak{u} and \mathfrak{l} of $\mathfrak{gl}(n, \mathbb{C})$, where \mathfrak{u} is spanned by the E_{1j} , $2 \leq j \leq n$, and

$$(11.65) \quad \mathfrak{l} := \left\{ \begin{pmatrix} x & 0 \\ 0 & B \end{pmatrix} : x \in \mathbb{C} \text{ and } B \in \mathfrak{gl}(n-1, \mathbb{C}) \right\}.$$

Then $\mathfrak{q} = \mathfrak{l} \oplus \mathfrak{u}$ is a parabolic subalgebra of $\mathfrak{gl}(n, \mathbb{C})$. By (11.64), the kernel of the restriction of $D_{\mathfrak{r}}$ is isomorphic to $H^k(\mathfrak{u}, \pi)$, the Lie algebra

cohomology of \mathfrak{u} with respect to π . Now Kostant's theorem determines the latter as an \mathfrak{l} -module, where $l \in \mathfrak{l}$ acts on $\Lambda^k(\mathfrak{u}) \otimes V_\pi$ by

$$(11.66) \quad -\text{ad}(l)^* \otimes \text{id} + \text{id} \otimes \pi_*(l),$$

see (4.138b) in [KnVo]. To apply Kostant's theorem, we introduce

$$(11.67) \quad \Delta^+(\mathfrak{u}) = \{\rho_1 - \rho_j : 2 \leq j \leq n\},$$

$$(11.68) \quad \Delta^+(\mathfrak{l}) = \{\rho_i - \rho_j : 2 \leq i < j \leq n\}.$$

For $w \in \mathcal{W}$, we also introduce

$$(11.69) \quad \Delta^+(w) := \{\lambda \in \Delta^+ : w^{-1}\lambda < 0\}, \quad \ell(w) := |\Delta^+(w)|,$$

and

$$(11.70) \quad \mathcal{W}^1 := \{w \in \mathcal{W} : \Delta^+(w) \subseteq \Delta^+(\mathfrak{u})\}.$$

Then $\mathcal{W}^1 = \{w_0, \dots, w_{n-1}\}$, where $w_0 = \text{id}$ and

$$(11.71) \quad w_k^{-1} = \begin{pmatrix} 1 & 2 & \cdots & k+1 \\ k+1 & 1 & \cdots & k \end{pmatrix},$$

for $1 \leq k \leq n-1$. We note that $\ell(w_k) = k$, for $0 \leq k \leq n-1$.

Kostant's theorem implies that, as an \mathfrak{l} -module, $H^k(\mathfrak{u}, \pi)$ is the irreducible representation of \mathfrak{l} with highest weight

$$(11.72) \quad w_k(\lambda + \delta) - \delta = (\lambda_{k+1} - k)\rho_1 + \sum_{2 \leq j \leq k+1} (\lambda_{j-1} + 1)\rho_j + \sum_{j > k+1} \lambda_j \rho_j,$$

where δ is the half sum of the positive roots of $\mathfrak{gl}(n, \mathbb{C})$,

$$(11.73) \quad \delta := \frac{1}{2} \sum_{j=1}^n (n+1-2j)\rho_j.$$

Moreover, the action of the \mathfrak{k} -component $Z^\mathfrak{k} \simeq -iE_{1,1}$ of Z on $H^k(\mathfrak{u}, \pi)$ is given by multiplication with

$$(11.74) \quad ik - i\lambda_{k+1} = ik + \text{id} \otimes \pi(Z^\mathfrak{k}).$$

In particular, $\text{id} \otimes \pi(Z^\mathfrak{k}) = -i\lambda_{k+1}$. It follows that, on $H^k(\mathfrak{u}, \pi)$,

$$(11.75) \quad D_{\mathfrak{z}} = D_Z = (-1)^k \frac{1}{2} (2k - 2\lambda_{k+1} - n + 1),$$

which is the third assertion of the theorem. We have

$$(11.76) \quad b_k = \dim H^k(\mathfrak{u}, \pi) = \prod_{\alpha \in \Delta^+(\mathfrak{l})} \frac{(\alpha, w_k(\lambda + \delta) - \delta + \delta_{\mathfrak{l}})}{(\alpha, \delta_{\mathfrak{l}})},$$

by Weyl's dimension formula, where $\delta_{\mathfrak{l}}$ is the half sum of the positive roots of \mathfrak{l} ,

$$(11.77) \quad \delta_{\mathfrak{l}} = \frac{1}{2} \sum_{j=2}^n (n+2-2j) \rho_j.$$

The second assertion of the theorem is an easy consequence. \square

11.3. Résumé and Examples. For the convenience of the reader, we give a résumé of the correction terms $\text{Corr}(\mathcal{C})$ to the extended index of D^+ for complex hyperbolic cusps \mathcal{C} of complex dimension n . We assume that the cross section of \mathcal{C} is given as $\Gamma \backslash G_{n-1}$ as in Corollary 10.64, where $N := G_{n-1}$ is the Heisenberg group of dimension $2n-1$ and where Γ is a uniform lattice in G_{n-1} . We denote by d the type of Γ and set $|\Gamma| = d_1 \cdots d_{n-1}$ as in Section 10.1.

Assume that, over \mathcal{C} , the Dirac bundle E is up to twist of the form $E = \mathcal{P} \times_{\beta} (\Sigma_{\mathfrak{s}} \otimes V)$ with $\mathfrak{s} \subset \mathfrak{su}(1, n)$ as in Section 2.3 and β as in (9.15), that is, $\beta_* = \hat{\alpha}_* \otimes \text{id} + \text{id} \otimes \pi_*$, where π_* is a unitary representation of $\mathfrak{k} = \mathfrak{u}(n)$ on V . Then sections of E^+ over \mathcal{C} correspond to maps $S \rightarrow \Sigma_{\mathfrak{n}} \otimes V$ satisfying the twist condition (9.22). Clearly, $\text{Corr}(\mathcal{C})$ is the sum of the parts which are determined by the subbundles of E^+ over \mathcal{C} of the form $\mathcal{P} \times_{\beta} (\Sigma_{\mathfrak{n}} \otimes W)$, where W is irreducible as a joint $\mathfrak{u}(n)$ and Γ module. Therefore, we may restrict ourselves to the case where V is irreducible as a joint $\mathfrak{u}(n)$ and Γ module.

Recall that we use an index \mathcal{C} to denote objects connected to the cusp \mathcal{C} and that

$$(11.78) \quad \text{Corr}(\mathcal{C}) = \frac{1}{2} \left(\lim_{t \rightarrow \infty} \eta(A_{\mathcal{C},t}^{\text{he},+}) + \eta(A_{\mathcal{C},0}^{\text{le},+}) + \dim \ker A_{\mathcal{C},0}^{\text{le},+} \right),$$

see Theorem 8.10.

THEOREM 11.79. *Suppose that V is irreducible as a joint $\mathfrak{u}(n)$ and Γ module. Then there are two cases:*

(1) *If τ is non-trivial,*

$$\begin{aligned} \lim_{t \rightarrow \infty} \eta(A_{\mathcal{C},t}^{\text{he},+}) &= (-1)^n |\Gamma| \dim V (\zeta_c(1-n) + (-1)^n \zeta_{1-c}(1-n)), \\ \eta(A_{\mathcal{C},0}^{\text{le},+}) &= \dim \ker(A_{\mathcal{C},0}^{\text{le},+}) = 0, \end{aligned}$$

where c denotes the twist parameter of τ as in (10.7).

(2) If τ is trivial, then π_* is irreducible and

$$\begin{aligned}\lim_{t \rightarrow \infty} \eta(A_{\mathcal{C},t}^{\text{he},+}) &= (-1)^n |\Gamma| \dim V (\zeta(1-n) + (-1)^n \zeta(1-n)), \\ \eta(A_{\mathcal{C},0}^{\text{le},+}) &= \sum (-1)^k b_k \text{sign}(n-1-2k+2\lambda_{k+1}), \\ \dim \ker A_{\mathcal{C},0}^{\text{le},+} &= \sum b_k \cdot \delta_{2\lambda_{k+1}, 2k+1-n},\end{aligned}$$

where $\lambda_1 \geq \dots \geq \lambda_n$ denotes the highest weight of π_* and

$$b_k = (n-1)! \dim V \prod_{\substack{1 \leq j \leq n \\ j \neq k+1}} |\lambda_j - \lambda_{k+1} + k + 1 - j|^{-1}.$$

REMARKS 11.80. 1) Recall that D^+ is of Fredholm type if and only if $\dim \ker A_{\mathcal{C},0}^{\text{le},+} = 0$, for all cusps \mathcal{C} of M .

2) The dimension of V in Assertion 2 is determined in (11.46).

Proof of Theorem 11.79. The formulas for $\lim_{t \rightarrow \infty} \eta(A_{\mathcal{C},t}^{\text{he},+})$ are taken from Corollary 10.64. The formulas for $\eta(A_{\mathcal{C},0}^{\text{le},+})$ and $\dim \ker A_{\mathcal{C},0}^{\text{le},+}$ follow from Proposition 11.31 and Theorem 11.53, where we observe that D_N^{le} in Proposition 11.31 corresponds to $-A_{\mathcal{C},0}^{\text{le},+}$. \square

EXAMPLES 11.81. We discuss the low energy invariants in some cases, where the twist τ is trivial. For convenience, we assume that all ends of M are complex hyperbolic cusps. Furthermore, in our examples, the Dirac bundle is of the same kind along all the cusps of M , that is, π_* is the same for all the cusps. We discuss the more precise quantities from Theorem 11.53. The formulas for the low energy invariants as in Theorem 11.79 will follow easily.

1) DOLBEAULT OPERATOR on forms of bi-degree $(0, q)$: In this case, π is the irreducible representation with highest weight $\lambda_j = (n+1)/2$, $1 \leq j \leq n$. We compute

$$(11.82) \quad b_k = \dim H^k(\pi) = \binom{n-1}{k}$$

and

$$(11.83) \quad D_Z|_{H^k(\pi)} = (-1)^k (k-n),$$

by Theorem 11.53.3. Hence $\eta(A_{\mathcal{C},0}^{\text{le},+}) = 0$ and $\dim \ker A_{\mathcal{C},0}^{\text{le},+} = 0$, for all cusps \mathcal{C} of M . In particular, D^+ is of Fredholm type.

2) SIGNATURE OPERATOR: In this case, π is the spin representation $\Sigma = \Sigma_{\mathfrak{p}}$, which is the sum of the irreducible representations V_l with highest weight as in (11.48), where $0 \leq l \leq n$. As for the dimension

$b_{k,l}$ of $H^k(\mathbf{u}, V_l)$, there are two cases:

$$(11.84) \quad b_{k,l} = \begin{cases} \binom{n}{l} \binom{n}{k} \frac{l-k}{n} & \text{if } k < l, \\ \binom{n}{l} \binom{n}{k+1} \frac{k+1-l}{n} & \text{if } k \geq l. \end{cases}$$

Furthermore, we have

$$(11.85) \quad D_Z|_{H^k(V_l)} = \begin{cases} (-1)^k(k-l) & \text{if } k < l, \\ (-1)^k(k+1-l) & \text{if } k \geq l, \end{cases}$$

by Theorem 11.53.3. Hence $\dim \ker A_{\mathcal{C},0}^{\text{le},+} = 0$, for all cusps \mathcal{C} of M , and therefore D^+ is a Fredholm operator. We also get

$$(11.86) \quad \eta(A_{\mathcal{C},0}^{\text{le},+}) = \sum_{k < l} (-1)^k \binom{n}{l} \binom{n}{k} \frac{l-k}{n} + \sum_{k \geq l} (-1)^{k+1} \binom{n}{l} \binom{n}{k+1} \frac{k+1-l}{n}.$$

If we change l in $n-l$ and k in $n-1-k$ in the second sum, we obtain

$$(11.87) \quad \eta(A_{\mathcal{C},0}^{\text{le},+}) = 0 \quad \text{if } n \text{ is odd.}$$

For even n , we get

$$(11.88) \quad \eta(A_{\mathcal{C},0}^{\text{le},+}) = 2 \sum_{k < l} (-1)^k \binom{n}{l} \binom{n}{k} \frac{l-k}{n}.$$

For $1 \leq l \leq n$, we have

$$\sum_{0 \leq k < l} (-1)^k \binom{n}{k} = (-1)^{l-1} \binom{n-1}{l-1}$$

and

$$\sum_{0 \leq k < l} (-1)^k \binom{n}{k} \frac{k}{n} = \sum_{0 \leq k < l-1} (-1)^{k+1} \binom{n-1}{k} = (-1)^{l-1} \binom{n-2}{l-2}.$$

Hence

$$\begin{aligned} \eta(A_{\mathcal{C},0}^{\text{le},+}) &= 2 \sum_{1 \leq l \leq n} (-1)^{l-1} \binom{n}{l} \left\{ \binom{n-1}{l-1} \frac{l}{n} - \binom{n-2}{l-2} \right\} \\ &= 2 \sum_{1 \leq l \leq n} (-1)^{l-1} \binom{n-1}{l-1}^2 + 2 \sum_{1 \leq l \leq n} (-1)^l \binom{n}{l} \binom{n-2}{l-2} \end{aligned}$$

The first sum is zero since n is even. The second sum is the coefficient of x^n in $(1-x)^n(1+x)^{n-2} = (1-x)^2(1-x^2)^{n-2}$ and hence

$$(11.89) \quad \eta(A_{\mathcal{C},0}^{\text{le},+}) = 2(-1)^{n/2} \left(\binom{n-2}{n/2} - \binom{n-2}{n/2-1} \right).$$

3) DIRAC OPERATOR on spinors: In this case, π_* is the irreducible representation with highest weight $\lambda = 0$ (where the spin structure along the cusps is periodic). If n is even, then $\ker A_{\mathcal{C},0}^{\text{le},+} = 0$, for all

cusps \mathcal{C} . In particular, D is a Fredholm operator. Furthermore, each cusp \mathcal{C} contributes a low energy η -invariant,

$$\begin{aligned}
 \eta(A_{\mathcal{C},0}^{\text{le},+}) &= \sum_{0 \leq k \leq n-1} (-1)^k \binom{n-1}{k} \text{sign}(n-1-2k) \\
 &= 2 \sum_{0 \leq 2k \leq n-2} (-1)^k \binom{n-1}{k} \\
 (11.90) \quad &= 2 \sum_{0 \leq 2k \leq n-2} (-1)^k \left(\binom{n-2}{k} + \binom{n-2}{k-1} \right) \\
 &= 2(-1)^{\frac{n-2}{2}} \binom{n-2}{\frac{n-2}{2}}.
 \end{aligned}$$

If n is odd, we have $\eta(A_{\mathcal{C},0}^{\text{le},+}) = 0$, for each cusp \mathcal{C} . Furthermore, each cusp \mathcal{C} contributes to the kernel,

$$(11.91) \quad \dim \ker A_{\mathcal{C},0}^{\text{le},+} = \binom{n-1}{\frac{n-1}{2}}.$$

In particular, D^+ is not a Fredholm operator for odd n .

REMARK 11.92. For n odd, the bilinear form β considered in Exercise 20.38 of [FuHa] defines a real $\text{Spin}(2n)$ -equivariant isomorphism between Σ_{2n}^+ and Σ_{2n}^- . By equivariance, this isomorphism induces a parallel isomorphism between the positive and negative spinor bundles Σ^\pm of a spin manifold M of dimension $2n$ and therefore also an isomorphism between the spaces of harmonic sections of bundles of the form $E^\pm = \Sigma^\pm \otimes_{\pi_*} V$. We conclude that the L^2 -index of the corresponding Dirac operator D^+ vanishes and that the extended indices of D^+ and D^- coincide. In particular, we get

$$(11.93) \quad \text{ind } D_{\text{ext}}^+ = h_\infty^+ = h_\infty^-,$$

compare Theorem 1.21. By a similar reasoning, we also obtain that the index densities ω_{D^\pm} vanish. In the case of the Dirac operator as in Example 11.81.3 above, the term in (11.91) times the number ν of ends of M gives the extended index of D which is equal to $h_\infty^+ + h_\infty^-$, by (8.12). Hence $\text{ind } D_{\text{ext}}^+$ is in this case equal to $\nu/2$ times $\binom{n-1}{\frac{n-1}{2}}$.

REFERENCES

- [Ag] S. Agmon, *Lectures on elliptic boundary value problems*. Van Nostrand Company, Princeton-Toronto-London, 1965.
- [An] N. Anghel, An abstract index theorem on non-compact Riemannian manifolds. *Houston J. Math.* **19** (2) (1993), 223–237.
- [APS1] M.F. Atiyah, V.K. Patodi, & I.M. Singer, Spectral asymmetry and Riemannian geometry. I. *Math. Proc. Cambridge Phil. Soc.* **77** (1975), 43–69.

- [APS3] M.F. Atiyah, V.K. Patodi, & I.M. Singer, Spectral asymmetry and Riemannian geometry. III. *Math. Proc. Cambridge Phil. Soc.* **79** (1976), 71–99.
- [Bä] C. Bär, The Dirac operator on hyperbolic manifolds of finite volume. *J. Differential Geom.* **54** (2000), no. 3, 439–488.
- [Ba] W. Ballmann, *Lectures on spaces of nonpositive curvature*. With an appendix by Misha Brin. DMV Seminar, 25. Birkhäuser Verlag, Basel, 1995. viii+112 pp.
- [BBB] W. Ballmann, M. Brin, & K. Burns, On the differentiability of horocycles and horocycle foliations. *J. Differential Geom.* **26** (1987), no. 2, 337–347.
- [BB1] W. Ballmann & J. Brüning, On the spectral theory of surfaces with cusps. In: *Geometric analysis and partial differential equations*, 13–37, Springer, Berlin, 2003.
- [BB2] W. Ballmann & J. Brüning, On the spectral theory of manifolds with cusps. *J. Math. Pures Appl.* **80** (2001), 593–625.
- [BBC1] W. Ballmann, J. Brüning, & G. Carron, Eigenvalues and holonomy. *Int. Math. Res. Notes* **2003**, 657–665.
- [BBC2] W. Ballmann, J. Brüning, & G. Carron, Regularity and index theory for Dirac-Schrödinger systems with Lipschitz coefficients. *J. Math. Pures Appl.* (9) **89** (2008), no. 5, 429–476.
- [BGS] W. Ballmann, M. Gromov, & V. Schroeder, *Manifolds of nonpositive curvature*. Progress in Mathematics, 61. Birkhäuser Boston, Inc., Boston, MA, 1985. vi+263 pp.
- [BaMo] D. Barbasch & H. Moscovici, L^2 -index and the Selberg trace formula. *J. Funct. Anal.* **53** (1983), no. 2, 151–201.
- [BeKa] I. Belegradek & V. Kapovitch, Classification of negatively pinched manifolds with amenable fundamental groups. *Acta Math.* **196** (2006), no. 2, 229–260.
- [BoWo] B. Booß-Bavnbek & K. Wojciechowski: *Elliptic boundary problems for Dirac operators*. Birkhäuser, Basel etc. 1993.
- [Bo] B.H. Bowditch, Discrete parabolic groups. *J. Differential Geom.* **38** (1993), no. 3, 559–583.
- [BrKa] M. Brin & H. Karcher, Frame flows on manifolds with pinched negative curvature. *Compositio Math.* **52** (1984), no. 3, 275–297.
- [BuKa] P. Buser & H. Karcher, *Gromov’s almost flat manifolds*. Astérisque, 81. Société Math. de France, Paris, 1981. 148 pp.
- [CaPe] R. Camporesi & E. Pedon, The continuous spectrum of the Dirac operator on noncompact Riemannian symmetric spaces of rank one. *Proc. Amer. Math. Soc.* **130** (2002), no. 2, 507–516.
- [Ca1] G. Carron, Un théorème de l’indice relatif. *Pacific J. Math.* **198** (2001), 81–107.
- [Ca2] G. Carron, Théorèmes de l’indice sur les variétés non-compactes. *J. Reine Angew. Math.* **541** (2001), 81–115.
- [ChGr] J. Cheeger & M. Gromov, Bounds on the von Neumann dimension of L^2 -cohomology and the Gauss-Bonnet theorem for open manifolds. *J. Differential Geom.* **21** (1985), no. 1, 1–34.

- [DeSi] C. Deninger & W. Singhof, The e -invariant and the spectrum of the Laplacian for compact nilmanifolds covered by Heisenberg groups. *Invent. Math.* **78** (1984), no. 1, 101–112.
- [Do] H. Donnelly, Essential spectrum and heat kernel. *J. Funct. Anal.* **75** (1987), no. 2, 362–381.
- [Eb] P. Eberlein, Lattices in spaces of nonpositive curvature. *Ann. of Math.* (2) **111** (1980), no. 3, 435–476.
- [FuHa] W. Fulton & J. Harris, *Representation theory. A first course*. Graduate Texts in Mathematics, 129. Readings in Mathematics. Springer-Verlag, New York, 1991. xvi+551 pp.
- [GiTr] D. Gilbarg & N.S. Trudinger, *Elliptic partial differential equations of second order*. Second edition. Grundlehren der Mathematischen Wissenschaften 224. Springer-Verlag, Berlin, 1983. xiii+513 pp.
- [Gi1] P.B. Gilkey, On the index of geometrical operators for Riemannian manifolds with boundary. *Adv. Math.* **102** (1993), no. 2, 129–183.
- [Gi2] P.B. Gilkey, *Invariance theory, the heat equation, and the Atiyah-Singer index theorem*. Second edition. Studies in Advanced Mathematics. CRC Press, Boca Raton, FL, 1995. x+516 pp.
- [GoWi] C.S. Gordon & E.N. Wilson, The spectrum of the Laplacian on Riemannian Heisenberg manifolds. *Michigan Math. J.* **33** (1986), no. 2, 253–271.
- [GrHu] D. Grieser & E. Hunsicker, Pseudodifferential operator calculus for generalized \mathbb{Q} -rank 1 locally symmetric spaces I. *J. Funct. Anal.* **257** (2009), 3748–3801.
- [Gr] M. Gromov, Almost flat manifolds. *J. Differential Geom.* **13** (1978), no. 2, 231–241.
- [GrLa] M. Gromov & H.B. Lawson, Jr., Positive scalar curvature and the Dirac operator on complete Riemannian manifolds. *Inst. Hautes Études Sci. Publ. Math.* **58** (1983), 83–196.
- [GKR] K. Grove, H. Karcher, & E.A. Ruh, Jacobi fields and Finsler metrics on compact Lie groups with an application to differentiable pinching problems. *Math. Ann.* **211** (1974), 7–21.
- [Ha] G. Harder, A Gauss-Bonnet formula for discrete arithmetically defined groups. *Ann. Sci. École Norm. Sup.* (4) **4** (1971), 409–455.
- [HeIH] E. Heintze & H. C. Im Hof, Geometry of horospheres. *J. Differential Geom.* **12** (1977), 481–491.
- [Hi] M. Hilsum, L’invariant η pour les variétés lipschitziennes. *J. Differential Geom.* **55** (2000), no. 1, 1–41.
- [Ka] A. Kaplan, Fundamental solutions for a class of hypoelliptic PDE generated by composition of quadratic forms. *Trans. Amer. Math. Soc.* **258** (1980), no. 1, 147–153.
- [KnVo] A. Knapp & D. Vogan, *Cohomological induction and unitary representations*. Princeton Mathematical Series, 45. Princeton University Press, Princeton, NJ, 1995.
- [Ko] B. Kostant, Lie algebra cohomology and the generalized Borel-Weil theorem. *Ann. of Math.* **74** (1961) 329–387
- [LaMi] H.B. Lawson Jr. & M.-L. Michelsohn, *Spin geometry*. Princeton Mathematical Series, 38. Princeton University Press, Princeton, NJ, 1989. xii+427 pp.

- [LMP] E. Leichtnam, R. Mazzeo, & P. Piazza, The index of Dirac operators on manifolds with fibered boundaries. *Bull. Belg. Math. Soc. Simon Stevin* **13** (2006), 845–855.
- [LiYa] P. Li & S.T. Yau, Eigenvalues of a compact Riemannian manifold. *Proc. Symp. Pure Math.* **36** (1980), 205–239.
- [Lo1] J. Lott, On the spectrum of a finite-volume negatively-curved manifold. *Amer. J. of Math.* **123** (2001) 185–205.
- [Lo2] J. Lott, Collapsing and Dirac-type operators. *Geom. Ded.* **91**, (2002), 175–196.
- [MaMe] R. Mazzeo & R.B. Melrose, Pseudodifferential operators on manifolds with fibered boundaries. *Asian J. Math.* **2** (1998), 833–866.
- [Mi] M. Mitrea, Generalized Dirac operators on nonsmooth manifolds and Maxwell’s equations. *J. Fourier Anal. Appl.* **7** (2001), no. 3, 207–256.
- [Mü1] W. Müller, L^2 -index and resonances. In *Geometry and Analysis on manifolds*, Katata/Kyoto (1987), Lectures Notes in Math. 1339, 203–211.
- [Mü2] W. Müller, *Manifolds with cusps of rank one, spectral theory and L^2 -index theorem*, Lecture Notes in Math. 1244, Springer-Verlag, Berlin [e.a.], 1987.
- [Pa] R. Parthasarathy, Dirac operator and the discrete series, *Ann. of Math.* **96** (1972), 1–30.
- [Ru] E. Ruh, Almost flat manifolds. *J. Differential Geom.* **17** (1982), 1–14.
- [St1] M. Stern, L^2 -index theorems on locally symmetric spaces. *Invent. Math.* **96** (1989), no. 2, 231–282.
- [St2] M. Stern, Eta invariants and Hermitian locally symmetric spaces. *J. Differential Geom.* **31** (1990), no. 3, 771–789.
- [Ta] M. Taylor, *Partial differential equations I. Basic theory*. Appl. Math. Sciences 115, Springer, Berlin etc. 1996.
- [Va] B. Vaillant, *Index- and spectral theory for manifolds with generalized fibred cusps*. Bonner Mathematische Schriften 344, Universität Bonn, Mathematisches Institut, Bonn, 2001. ii+124 pp.
- [Wo] J.A. Wolf, Essential self-adjointness for the Dirac operator and its square. *Indiana Univ. Math. J.* **22** (1972/73), 611–640.
- [Wu] H.H. Wu, *The Bochner technique in differential geometry*. Math. Rep. 3 (1988), no. 2, i–xii and 289–538.
- [Zu] S. Zucker, L_2 cohomology of warped products and arithmetic groups. *Invent. Math.* **70** (1982/83), no. 2, 169–218.

MATHEMATISCHES INSTITUT, UNIVERSITÄT BONN, ENDENICHER ALLEE 60,
53115 BONN AND MAX-PLANCK-INSTITUT FÜR MATHEMATIK, VIVATSGASSE 7,
53111 BONN, DEUTSCHLAND,

E-mail address: `hwblmnn@math.uni-bonn.de`

INSTITUT FÜR MATHEMATIK, HUMBOLDT-UNIVERSITÄT, RUDOWER CHAUSSEE
5, 12489 BERLIN, GERMANY,

E-mail address: `bruening@mathematik.hu-berlin.de`

DÉPARTEMENT DE MATHÉMATIQUES, UNIVERSITÉ DE NANTES, 2 RUE DE LA
HOUSINIÈRE, BP 92208, 44322 NANTES CEDEX 03, FRANCE,

E-mail address: `Gilles.Carron@math.univ-nantes.fr`