

Detection of a southern peak in Mercury's sodium exosphere with the TNG in 2005

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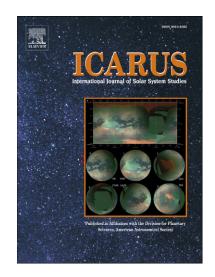
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1	Detection of a southern peak in Mercury's sodium exosphere with the TNG in 2005
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2	Abstract

A long term plan of observations of the sodium exosphere of Mercury began in 2002 by using the high resolution echelle spectrograph SARG and a devoted sodium filter at the 3.5 m Galileo National Telescope (TNG) located in La Palma, Canary Islands. This program is meant to investigate the variations of the sodium exosphere appearance under different conditions of observations, namely Mercury's position along its orbit, phase angle and different solar conditions, as reported by previous observations in August 2002 and August 2003 (Barbieri et al., 2004 and Leblanc et al., 2006).

Here we present the analysis of data taken in June 29th and 30th and in July 1st 2005, when Mercury's True Anomaly Angle (TAA) was in the range 124-130°. The spectra show particularly intense sodium lines with a distinctive peak in emission localized in the southern hemisphere at mid-latitudes. This seems to be a persistent feature related to consecutive favourable IMF conditions inducing localized enhancements of surface sodium density. The comparison with previous data taken by Potter et al. (2002) evidences a surprising consistency of the anti-sunward component, which appears to remain constant regardless of the changing illumination and space weather conditions at Mercury.

Keyword: MERCURY, EXOSPHERE; SPECTROSCOPY; SODIUM OBSERVATION

Introduction

In 1974 the existence of a thin atmosphere around Mercury was discovered by Mariner 10, which revealed UV emissions of three atomic elements: H, He and O (Broadfoot *et al.*, 1976). Three other elements (Na, K, and Ca) were later discovered with ground-based observations in the visible spectral range (Potter & Morgan 1985, Potter & Morgan 1986, and Bida *et al.* 2000, respectively). Due to the low surface number density (approximately $n = 10^5$ - 10^7 atoms/cm³, $P = 10^{-12}$ bar, on the dayside), this atmosphere can be considered collisionless. Therefore, the whole atmosphere is comparable with an exosphere having the exobase coincident with the planet's surface, i.e. a surface bounded exosphere (Stern, 1999).

More than thirty years after its discovery, the dynamics of such exosphere, the source and sink mechanisms and its relationship with solar activity, are still not clearly understood. Although NASA MESSENGER has already reached Mercury's orbit and detected gases (including Na) in the exospheric tail (McClintock et al., 2008), Earth-based observations are the only way to provide continuous data and to support exospheric simulations, and still remain an effective method to study and try to understand the Hermean environment.

Therefore, a campaign of observations started in 2002 at the Telescopio Nazionale Galileo (TNG) with the goal to provide the largest possible database of Na exosphere by covering a full Mercury's year at different true anomaly angles (TAA), phase angles and solar conditions. The observations were performed with the TNG high resolution echelle spectrograph SARG, equipped with a narrow filter of 60 Å centred on the sodium doublet at 5890-5896 Å (D lines), in order to suppress the adjacent spectral orders. The projected slit is 26.7 arcseconds long and 0.4 arcseconds wide. Table 1 gives the main parameters of the spectroscopic configuration; more detailed information on SARG can be found at the website: http://www.pd.astro.it/sarg/.

The observations performed in 2002 at TAA 171-174° and in 2003 at TAA 163-168° have been published by Barbieri et al. (2004; hereafter Paper I) and Leblanc et al. (2006; hereafter Paper II). Paper I showed the capability of SARG/TNG to detect Mercury's Na exosphere and reported column density values consistent with previous observations (Killen and Ip, 1999) as well as with model simulations by Leblanc and Johnson (2003). Paper II reported about three nights of observations, two of which with a remarkably stable Na exosphere; the third one showed a localized feature in the early morning varying on a very short time scale.

In the present paper we describe the observations performed in 2005 and the main remarkable features identified in Mercury's exosphere (Sections 1 and 2). Then we discuss (Section

1 3) and conclude (Section 4) on the possible origins of these features within the context of what is 2 known on the sources of Mercury's sodium exosphere.

I. Observations

In 2005 observations were performed during three evenings: June 29^{th} ant 30^{th} , and July 1^{st} , just before and soon after the sunset, until the moment in which the telescope safety mode blocked the observation, at Mercury's elevation of 13.5° above the horizon. At such large airmasses (2-4), large values of the seeing are expected. In the present paper seeing values are given as one σ of equation (1) in Sprague et al. (1997). These numbers should therefore be multiplied by 1.67 to obtain the FWHM seeing values used in many publications on Mercury's sodium exosphere (Potter et al. 2006; 2007; 2008). The seeing 1σ values for the three nights were in the range 1.5-2.0" in June 29^{th} , 0.9-1.3" in June 30^{th} and 1.5-2.1" in July 1^{st} .

To save time for observations, the read-out time was reduced by binning the CCD in the spatial direction (2x binning). As a consequence, the effective spatial pixel was $0.16\times2=0.32$ arcseconds. Due to the composition of the Earth's and Mercury's orbits inclinations plus the Mercury's axis inclination, the slit was placed at an angle of $11.2-12.4^{\circ}$ with respect of the E-W direction (with the Sun being to the W). The planet disk was 6.7-7.0 arcseconds in diameter, and more than 50% of it was illuminated by sunlight. Dawn terminator was visible (as in 2002 and 2003 observations). Mercury's main parameters are listed in Table 2.

We collected six useful spectra during the first night, thirteen during the second and five during the third night. The exposure time was 120 seconds for all the spectra. The peculiarity of this data set with respect to the previous ones of 2002 and 2003 is the much stronger sodium emission and, consequently, the high values of the ratios between exospheric sodium emission and sky background intensity. These ratios range from 2.4 to 14 for the first night, from 2.5 to 140 for the second, and from 28 to 60 for the third night. The occurrence for such remarkable enhancement is probably due to the conjunction of several reasons: clear sky conditions, a closer position of Mercury to the Sun and, most important, a larger radial heliocentric velocity of Mercury, which increases the solar flux effects and, hence, raises the excitation of the D1 and D2 resonant transitions.

II. Data Analysis

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The extraction procedure of exospheric sodium emission is the same that was used for the 2003 data set (Paper II); hence, here it will be repeated just briefly. The preliminary treatment consists of flat field, bias and sky subtraction. Spectral calibration is obtained by comparing the spectra with a Th-Ar reference lamp spectrum.

A noticeable improvement with respect to the previous data sets is a more accurate sky subtraction: special care was taken, in fact, to avoid subtracting any sodium at the edge of the slit. Such a procedure allows us to retain the faint Na emission present also in the regions out of the disk and, in particular, in the anti-sunward direction up to \sim 2 R_M from the planet surface.

The main points of the reduction phase are summarized in the following:

- G-factors: they are two values related to the photon efficiency in the process of resonant scattering producing the Na exospheric emission (Brown and Yung, 1976). They can be calculated by assuming a constant solar flux at the D wavelengths during the three nights. Values of g-factors, seeing (both conventional seeing and seeing sigma) and Doppler shifts (in mÅ) are given in Table 3.
- Continuum: its estimation is performed close to the D emission lines. In order to convert instrumental counts (ADUs) into Rayleigh (R) brightness (the Rayleigh unit is a measure of the omni-directional emission rate in a column of unit area along the line of sight, with 1 R = 10⁶ photon/cm²/s; Chamberlain and Hunten, 1987), a reflectivity model developed by Hapke (1986) and adapted to Mercury by Sprague et al. (1997) is used to derive the solar flux reflected by Mercury's surface. The comparison between this model and the measured continuum provides the conversion factors (from now on addressed as 'calibration factors') needed to calibrate the exospheric emission. This approach to calibrate Mercury's exospheric emission has been largely used in previous publications (Leblanc et al., 2006; Potter et al., 2006). In particular, it allows to get rid of the atmospheric absorption by using Mercury's surface as a calibrated candle and, therefore, to compare observations obtained at different air mass values. In addition, the comparison between the continuum and the Hapke surface reflectivity model provides an estimation of the seeing because the model depends on both the geometry and the seeing value at the time of the observations. Therefore, a good fit between the continuum and the Hapke model is also an indication of a good choice of the seeing value during our observations.
- Slit positioning above Mercury's disk: this is a delicate point. The correct position of the slit in the North-South (N-S) direction (that is, perpendicularly to the slit) is roughly derived by comparing the instrumental theoretical position with a picture of the planet reflected by the

- mask containing the slit. The correct position of the slit in the East-West (E-W) direction, instead, is derived by superposing the maximum of the measured continuum over the maximum of the calculated Hapke model. As in Paper II, the uncertainty related to the instrument pointing is estimated to be 0.5".
- Finally, the extraction of the exospheric emission: it is performed by using three different methods: (1) a Voigt profile to fit the D1 and D2 solar Fraunhofer lines (Hummer, 1962), as in the 2003 data set (Paper II); (2) a solar spectrum properly Doppler-shifted to simulate both the continuum and the Fraunhofer lines, and, (3) finally, thanks to its peculiar position on one side of the broad Fraunhofer absorption, a linear interpolation of the solar line in the interested region (as in Paper I). The comparison among the three cited methods shows no significant difference. Hence, linear interpolation was chosen as being the most straightforward solution, and is applied to all the spectra of the three nights.

The results in column density (atoms/cm²) are obtained with the assumption that D1 is optically thin and then by following the conversion described by Kameda et al. (2007).

The errors on the final values of column density have three different origins: (1) the instrumental noise; (2) the error caused by the extraction of the exospheric emission with a linear interpolation, and (3) the errors introduced by the conversion from ADUs to Rayleigh and from Rayleigh to atoms/cm². The instrumental noise and the error of extraction are negligible with respect to the errors due to the calibration. Since the Hapke model cannot be applied outside Mercury's visible disk, only calibration factors for the spectra placed on it can be considered. The plot in Fig. 1 (for the second night), shows that all calibration factors are close to 2.45 ± 0.70 for the spectra placed on the disk. Moreover, no significant dependence of these calibration factors with respect to airmass can be noticed. Therefore, a mean value can be assumed as the best value to be applied to all spectra of each night. We derived an average calibration factor of 1.92 for the first night, 2.45 for the second and 2.02 for the third. The first and the third values are very close (about 5% of difference), while the second one appears to be larger by about 20%, which highlights the significant better atmospheric conditions during this night of observation.

FIGURE 1

Regarding the error in the conversion from ADU to Rayleigh, which is mostly due to the calibration factor, we derive an estimate of 35%. When the conversion to atoms/cm² is applied, the contribution of the g-factor uncertainty should be considered, and the total uncertainty in column density calculation rises to 46%. For this uncertainty is really high, we decided to stop our calculation to brightness values (Rayleigh).

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III. Data Interpretation/Discussion

The Figs. 2, 3 and 4 show the final images of the measured emission for the three different nights of observation superposed to the planet disk.

During the first night we sampled the exosphere at almost regular distances between the different slit positions, while during the second night the larger number of spectra led to an almost complete coverage of the sodium emission from the exosphere of the planet. Unfortunately, during the third night the coverage of the mid-latitude regions (at both hemispheres) was very poor, even if a general trend of the spatial distribution of the exospheric emission can be deduced as well.

The column density values are in good agreement with previous measurements (1-3·10¹¹ atoms/cm² with respect of 2.8-3.8·10¹¹ by Killen et al., 1990) and with the predictions of current exospheric models (Leblanc and Johnson, 2003; Mura et al., 2007).

Our observations reveal a constant pattern during the three nights, characterized by enhancements along the equator and a peak in the mid-latitude southern hemisphere (the peak being less evident during the first night).

FIGURES 2, 3, 4

The equatorial enhancements have been already observed in the past (Potter et al., 2003; Killen et al., 2001) and seem to validate the theories of a major role in the exospheric production played by thermal desorption (TD) and photon-stimulated desorption (PSD) (McGrath et al. 1989; Hunten and Sprague, 1997). Both mechanisms act at their maximum in the sub-solar regions, whereas also atmospheric blurring would move the brightness peak toward the visible disk (Potter et al., 2006).

However, the observed maximum at mid-latitudes in the southern region cannot be explained by a uniform exosphere dominated by PSD or TD, and could be related, instead, to a local peak of release of the sodium trapped in the surface. This peak is clearly visible on the second night, when a better coverage of the planet is available. During this night (June 30th) the peak brightness is almost 15% higher than the average peak brightness of the other spectra (this is probably due to the better seeing conditions), while it is 12% for the third night, and 10% for the first. The occurrence of this peak during all the three nights, always in the same region of the disk with a regular profile both in the D2 and D1 component along the slits (i.e. the spatial dimension), is a strong indication in the direction of the real existence of the peak.

The possible explanations of this feature could be: (1) a localized sodium-rich region in the surface, and/or (2) a local peak of efficiency of one or more sodium release processes. In the first case, the sodium-rich region should be within 260-300° E longitude and 10-40° S latitude. This region is not included in any of the radar bright spots detected by Sprague et al. (1998) as having higher abundances of sodium bearing minerals, and denoted as A (15° E, 25° S), B (15° E, 55° N) and C (120° E, 15°N). However, Harmon et al. (2007) refer of a fainter radar feature located at 18° E, 27° S which could be probably related to the Bartok crater and that may be associated to the observed emission peak, if we take into account the seeing effect (seeing sigmas between 0.3 and 0.6 Hermean radii).

Localized emission peaks have been already observed in the past, at both hemispheres at mid-latitudes (see Potter et al., 2006 for a review; and Leblanc et al., 2008 for latest detection). Following these authors, the observed higher emission could be related to the interplanetary space conditions and in particular to the sign of the B_x component of the Interplanetary Magnetic Field (IMF), the x component being in the planet-Sun direction. In fact, the inter-connection between the Hermean magnetic field and the IMF may allow solar wind particles to enter the cusps down to the surface, and reconnection is virtually always present at Mercury apart from peculiar cases in which the B_Z component is dominant (e.g. see Fig. 5 of Massetti et al., 2007). A realistic IMF configuration with significantly $B_X > 0$ (< 0) causes reconnection mainly in the southern (northern) dayside magnetopause, with a consequent particle precipitation which could in turn cause an intense surface release due to ion-sputtering and, hence, a southern (northern) localized enhancement of atoms in the exosphere. The magnetic reconnection induced by IMF B_X can be amplified by a negative B_Z component, as well as by an increased solar wind density (Massetti et al., 2007).

FIGURES 5, 6

Given the impossibility to have direct information on the IMF at Mercury, we try to derive them by comparing the space weather condition at the Earth (detected by the ACE satellite, located at the L1 Sun-Earth Lagrangian point) with the solar coronal magnetic field as extrapolated from photospheric measurements (Wilcox Solar Observatory magnetograms). In this way, we find that, taking into account the time shift to the orbit of Mercury, the Sun was relatively active in the period of our observations, leading to several interplanetary perturbations (Figs. 5 and 6).

The ACE data reveals that an interplanetary discontinuity occurred between 182-183 DOY (DOY= day of the year, corresponding to July 1st and 2nd), marked by an increase of the solar wind density (up to about 50 cm⁻³) which was shortly followed by a steep increase of both temperature (up to $\sim 5\times 10^5$ K) and speed (up to 650 km/s) (Fig. 5). The event was likely caused by the transit of a fast stream coming from the low-latitude extension of a coronal hole that was facing the Earth on

June 30th (180 DOY, SoHO/EIT data) and represented by the black pattern in Fig. 6 (bottom) above 1 2 the image of the solar magnetic field. By taking into account the position of both the Earth and 3 Mercury with respect to the Sun (180° and 120° solar longitude, respectively), the solar rotation and 4 the relative time lag due to the propagation speed of the interplanetary perturbation, we can expect that the fast stream probably reached the inner planet in the middle of 176 DOY (June 25th), that is 5 6 about 4 days before our TNG ground-based observations (180-182 DOY). As ACE data show, at 7 the Earth orbit the fast stream was followed by a long period (6-7 days) characterized by a slightly 8 positive (~2 nT, on the average) IMF B_X component. By scaling back this condition to 0.42 AU we 9 could expect an IMF $B_X \sim 6-12$ nT at Mercury (the IMF strength at Mercury's orbit being about 3-6 10 times larger than at 1 AU, e.g.: Burlaga, 2001), that lasted for the whole observation period. This 11 IMF configuration is consistent with an asymmetric reconnection pattern with a large open cusp region in the southern hemisphere, where most of the solar wind plasma precipitation takes places. 12 13 peaking at mid-latitudes (Massetti et al. 2007). A fraction of that ion flux (≈10%) is expected to be 14 able to reach the dayside planetary surface, contributing to the release of sodium directly via ion-15 sputtering, or indirectly by enhanced diffusion of sodium from interior hence increasing the photon-16 stimulated desorption rate (Potter, 1995; Killen et al., 2004; Mura et al., 2008).

On the other hand, the analysis of the solar coronal magnetic field, computed from the observed photospheric field (WSO), does not completely match the IMF B_X time evolution. In fact, by using the position of the coronal hole to synchronize the solar wind and coronal field data sets (see Fig. 6), we found that after the transit of the fast stream associated with the coronal hole and embedded in a negative magnetic sector (IMF $B_X > 0$), there is a net transition to a large positive magnetic sector, leading a negative IMF B_X. This means that the IMF configuration changed significantly while the solar wind travelled till the Earth (ACE), and that somewhere B_X changed its sign. Hence, on the base of the available observations it cannot be clearly stated if the IMF B_X sign at Mercury during the period of observation was positive or negative. To try solve this ambiguity we inspected the Ecliptic-plane IMF forecast, computed by the Geophysical Institute (University of Alaska Fairbanks) on the basis of the Hakamada-Akasofu-Fry (HAF) solar wind model (Fry et al., 2001): Fig. 7 shows that the IMF likely had a 4 sector configuration, with the Earth embedded in a negative IMF B_X (away from the Sun), while Mercury was inside the adjacent positive IMF B_X sector (toward the Sun). The computed IMF geometry then confirms the sector structure observed by ACE spacecraft (Fig. 6, top panel): and we can be sufficient confident that the IMF B_X component was mainly positive during the period our observations, leading to a magnetic merging dominating in the southern hemisphere.

FIGURE 7

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Moreover, by inspecting the SOHO/LASCO CME catalogue (available on the website: http://cdaw.gsfc.nasa.gov/CME_list/) we find that during our Na observations, several Coronal Mass Ejections (CMEs) occurred close to ecliptic plane and to the east limb of the Sun, likely triggered by the active region AR0871 that was approximately facing Mercury. CMEs are violent ejections of plasma from the solar corona characterized by a steep increase in plasma density, temperature and velocity. They also drag typical magnetic field structures known as "magnetic clouds", characterised by twisted field lines. Some solar flares also occurred in the same period. They are sudden and intense brightness enhancements occurring in the solar corona, and emitting radiation (with energies of the order of 10²⁷ erg/s) as well as particles, including heavy ions. Flares occur in active regions, like the one that was facing Mercury. However, the intensity of flares (detected by GOES satellite located around the L1 point too) which occurred between June 26th and July 1st is always sensibly moderate (between C and B class).

By inspecting the emission brightness values during the three nights, June 30th appears to be globally stronger in emission with respect to the other two. The total enhancement is estimated to be 19% higher than in the third night. Unfortunately, the conversion error from ADU to Rayleigh accounts for 35% uncertainty on the measured emission brightness. Hence, we cannot state whether the observed enhancement is a real feature in the present observation, or not. However, should it resemble a real global enhancement, the present feature could be caused by a higher efficiency of all the release processes (or of the dominant ones) or by a higher efficiency of the resonant scattering process producing the exospheric D lines. Both these explanations would be reasonable if a higher solar flux reached Mercury between the first and the second night of observation. In fact, higher plasma density, velocity and He/H ratio are expected to cause a higher rate of efficiency of the release processes, leading to a global enhancement in the exospheric emission that could be consistent with the observations.

Finally, as already mentioned in Section 2, the accurate extraction of sodium emission in all the pixels of the spectra shows evidence of a sodium component out of the disk in the anti-sunward direction. The Mercury sodium tail was discovered by Potter et al. (2002) and observed by other authors (Baumgardner et al., 2008; Kameda et al., 2008; Potter and Killen, 2008). This feature is clearly visible in the spectra that lie completely outside the disk, but similar information can be extracted from the spectra crossing the disk as well, although the shadow of the planet itself decreases the efficiency of resonant scattering and, hence, the 'visibility' of the exospheric sodium component. All our spectra sample the sodium tail up to \sim 2 $R_{\rm M}$ far from the planet surface. Potter et al. (2002) plotted the sodium tail up to the surface along the axis of the tail, i.e. in the anti-sunward direction, by averaging the 10-pixel values along the N-S direction during two nights: June 5th 2000

(TAA = 114°) and May 26th 2001 (TAA = 125°). We compare our data with Potter et al. one by 1 averaging the values of the spectra at same x-value (Figs. 2, 3 and 4) and by considering the 2 3 different sizes of the planet disk at the different dates. To account for the different size of Mercury 4 at the time of observations, we compare our data E' to a reference observation E, by multiply E' by a factor D^{2}/D^{2} (where D and D' are the disk diameters during E and E' measurements, 5 respectively). Moreover, we have to convert the measured distances into heliocentric distances 6 7 (Potter et al., 2002), due to the different phase angles of the Earth with respect of Mercury and the 8 Sun.

FIGURE 8

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Figure 8 shows the results of the comparison of our three nights (solid lines) with the Potter et al. observations (filled triangles). In all cases the tail emissions have similar shape and brightness. This is consistent with the similar orbital position of Mercury during all the compared five nights (TAAs are in the range 114°-130°). In fact, it is expected that the sodium tail varies along the orbit of Mercury due to the solar radiation pressure variation, which is a function of both the heliocentric distance and the radial velocity (Potter and Morgan, 1987; Potter and Killen, 2008). Figure 8 also shows similar shape and brightness of the disk emissions during the first and third nights as compared to June 5th 2000 data, even if the greater seeing of July 1st causes broadening of the emission profile for this night. May 26th 2001 data, instead, appear to be higher (by a factor 1.7 in intensity of the maximum) when compared with June 30th, our most intense night. Even if considering all the discrepancies, our comparison with Potter et al. (2002) provides a rather and surprisingly good agreement of the 'tail' emission. Actually, the comparison of the sodium tail measured on May 26th 2001 with our measurement on June 29th 2005 displays a very good agreement. This is the best agreement and actually it is the case where the two profiles have been obtained for very close TAA. This result suggests, therefore, that the shape and brightness of the tail depend essentially on the position of Mercury along its orbit and does not depend significantly on the mechanism leading to the ejection of the sodium atoms. If confirmed by further observations, it would imply that the main population of Mercury's tail is produced by a mechanism not strongly dependent on the solar wind conditions (the most variable mechanism of ejection, in particular from solar minimum to maximum conditions).

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Conclusions

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In this paper we report the observation of Mercury sodium exosphere performed at TNG during three nights (June 29^{th} – July 1^{st} , 2005) at sunset. We observe column density values above

the disk in the range $1-3\cdot10^{11}$ atoms/cm², with a peak of emission in the southern hemisphere at mid-latitudes.

Higher emissions have been already correlated to solar activity (Killen et al. 2001), while asymmetries have been explained by models simulating the interaction between solar plasma and planetary magnetic field considering the different reconnection geometries driven by the IMF B_X component (Sarantos, 2001; Kallio and Janhunen, 2003). Hence, we try to correlate the observed features with solar activity by inspecting the SWO solar data, ACE solar wind data and SOHO/LASCO CMEs data at L1, and GOES flares data from a geostationary orbit. SWO evidences an active region and ACE data show that its effects remain for a long period (6-7 days) with a positive IMF B_X component (scaled at Mercury to ~6-12 nT). Moreover, many CMEs and flares occurred in the Sun in the direction of Mercury in the days before our observations, leading to the conclusion that some SEP events should have occurred and that a period of medium or high level space weathering affected the Hermean environment just before and during our observation run. Moreover, the modelled IMF geometry and the ACE spacecraft data, supports the hypothesis that the IMF B_X component was consistent with reconnection in the southern hemisphere.

The comparison of our data with previous one at similar TAA (June 5th 2000, and May 26th 2001; see Potter et al., 2002) reveals consistency both in profiles and intensities. In particular, the substantial invariability of the shape and brightness of the sodium emission in the anti-sunward direction for similar conditions of observation suggests that the tail surprisingly displays rather constant characteristics and is therefore only slightly dependent on other conditions than Mercury's TAA. This conclusion need further investigations and long-period observations, but it clearly illustrates that key information on Mercury's exosphere might be obtained by observing Mercury's sodium tail even close to its disk.

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3	

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1	Table and Figure Captions
2	
3	
4	Table 1 – SARG main parameters.
5	
6	Table 2 – Mercury's parameters at observation time.
7	
8	Table 3 – Seeing values, radial velocity shifts on spectra and g-factors for the three nights.
9	
10	Figure 1 – Plot of the calibration factor versus airmass values (Nb. time and slit numbers
11	increase to the right) for the night of June 30 th , 2005. Only the slits that cross the disk (hence, for
12	46, 47, 50, 53, 54, 57, 58, 59 and 60) show an almost constant value around 2.20. An average value
13	equal to 2.45 was calculated.
14	
15	Figure 2 – Global view of the observations performed during the first night (June 29 th July
16	1st) for the D1 lines in emission brightness (kR). Colour bars on the right give the corresponding
17	emission brightness (in kiloRayleigh), while seeing sigma are expressed on top (in Mercury radii
18	units). Numbers of the slit are also put on the figures to give the time sequence of the observation
19	during each night (increasing number corresponds to increasing time with ~3 minutes of delay
20	between two consecutive slit positions). Sun is on the right side (hence, this is also the illuminated
21	side of the disk); North Celestial Pole and Mercury North Pole are on top of the figure. Six slits
22	were available, showing higher brightness in the illuminated side of the disk, with a slightly higher
23	value (more extended to higher latitudes) in the Southern hemisphere.
24	
25	Figure 3 – Global view of the observations performed during the second night (June 30 th) for
26	the D1 lines in emission brightness (kR). Same descriptions as previous figure. Thirteen slits were
27	available, almost all on the disk, with an almost total coverage of it that evidences a more evident
28	and localized enhancement in the Southern hemisphere.
29	
30	Figure 4 – Global view of the observations performed during the third night (July 1 st), for
31	D1 lines in emission brightness (kR). Same descriptions as previous figure. Five slits were available
32	with a poor coverage of the interest region in the South, but still an enhancement is clearly visible.
33	

Figure 5 – ACE data of IMF B_X and B_Z , He/H ratio, proton temperature (T_P) , density (n_P)
and velocity (v _P) between doy 170-191 of 2005. A fast stream arrival is recognizable by the
increase of the solar wind velocity at the end of DOY 182 and is evidenced by a vertical line. A
peak in all these parameters is clearly present. Note that during the fast stream passage, the $B_{\rm X}$
component is >0.

Figure 6 – The IMF B_X component measured by ACE (second plot from the top in Figure 5) is reported, with time and y-axis reversed. Bottom left: large scale solar magnetic field as extrapolated at 20 R_{\odot} (coronal field) from SOHO data. On June 28th 2005 Earth is at 180° (central vertical line), solar equator is the horizontal line at mid altitude, while Mercury position is at about 120°. Light gray region is for IMF B_X <0, dark gray for IMF B_X >0. The solid black pattern in the middle of the figure shows the position of the coronal hole from which the fast stream originated (visible at the Earth from DOY 182). The effect of solar rotation leads to observe back in the past as we move to the right and to the future as we move to the left. Hence, we can compare the evolution of the B_X component at the Earth (upper panel) to the correspondent IMF B_X at the Sun by checking the main features, i.e. the sign inversion (solid arrows). Note that in the solar magnetic field the IMF B_X during the fast stream is >0, but there is a sign inversion after about four days, which is not present in the B_X ACE data (first arrow to the left).

Figure 7 – IMF geometry on the ecliptic plane as computed by the Geophysical Institute (University of Alaska Fairbanks). Red regions are for IMF away from the Sun (IMF $B_X < 0$), blue ones for IMF toward the Sun (IMF $B_X > 0$). During the days of interest Mercury (black dot with denoted with 'M') lies permanently inside IMF B_X positive region.

Figure 8 – Comparisons of each night of TNG observation (average of the central slits, crossing the equator) with the tail profile of Potter et al. (2002) for two nights: June 5th 2000, and May 26th 2001. Sun is on the right side and double x-axis shows both the heliocentric distance and the distance expressed in Hermean radii (R_M).Good agreement is displayed for all cases in the tail region (see semilog plot on the right side); good agreement is also shown in the first and third night with the data of June 5th 2000 on the disk region (central part, see normal plot on the left side). June 30th instead shows higher values compared with this night, but not enough with respect to May 26th 2001. Larger profile is shown on July 1st due to higher seeing values.

1 Figure 1

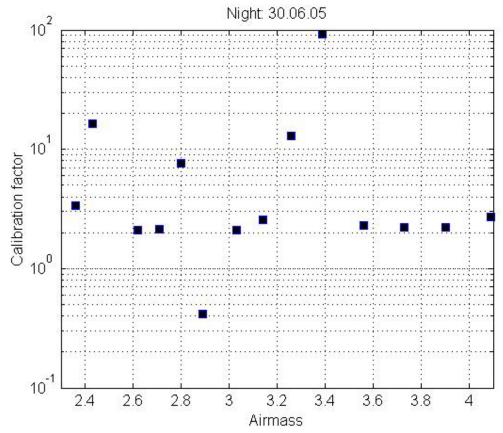
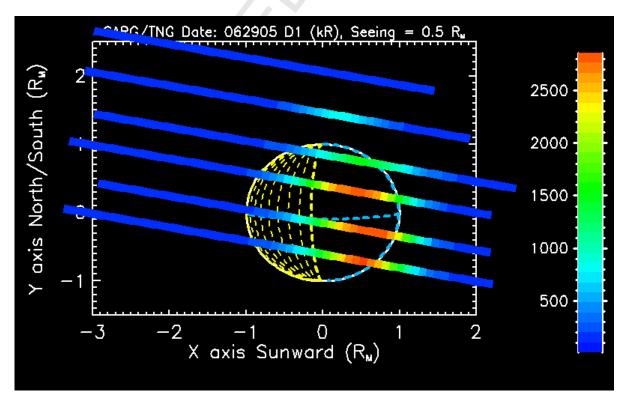
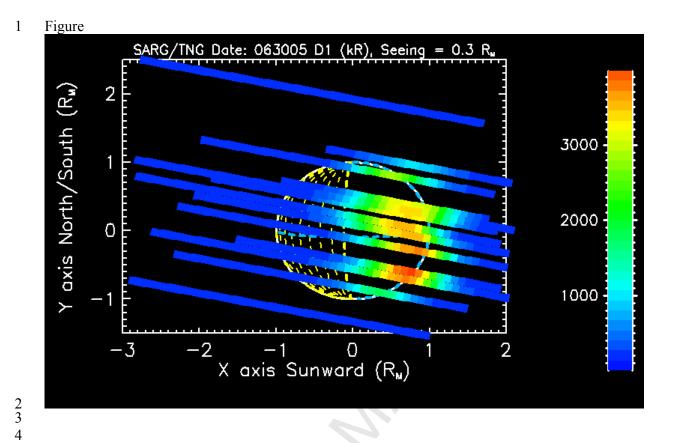
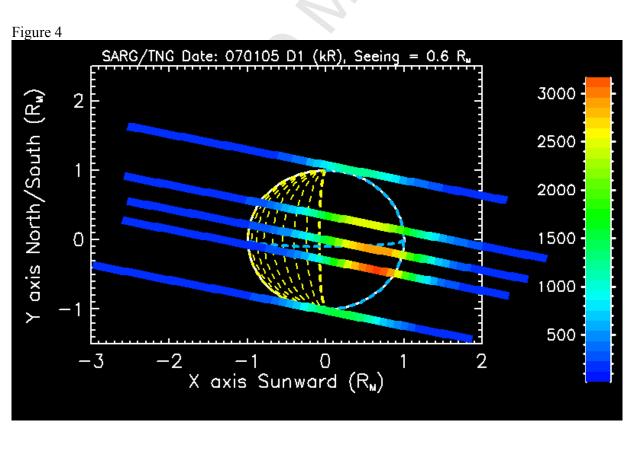


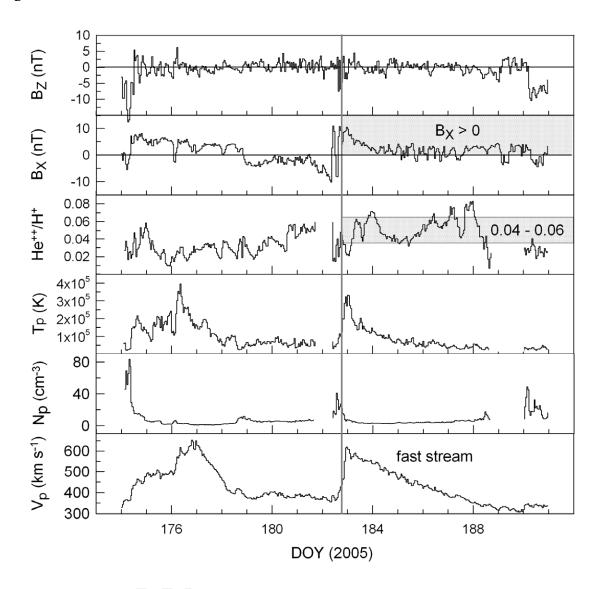
Figure 2



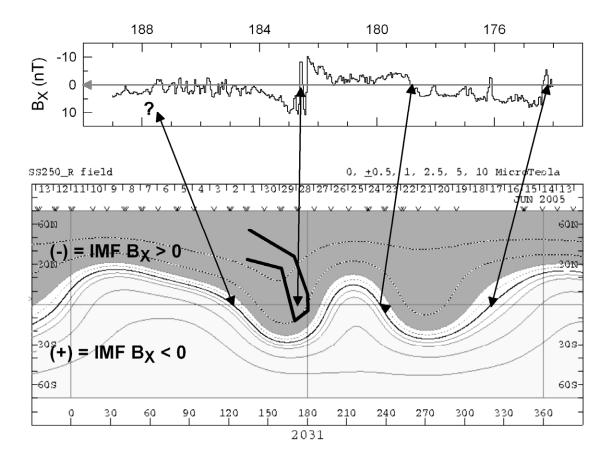




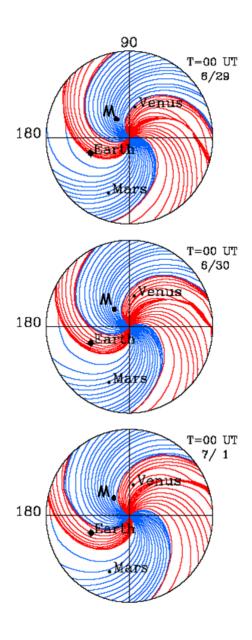
1 Figure 5



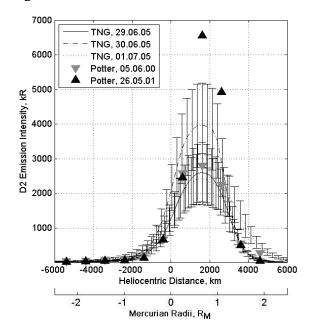
1 Figure 6

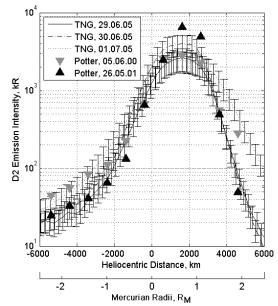


1 Figure 7



1 Figure 8





1 2 Table 1

Spectrograph resolution	115000
Slit length and width	26.7 x 0.40 arcsec
Pixel dimension and scale	0.022 Å, 0.16 arcsec
CCD dimension	2K x 4K pixels

Table 2

7

Parameters	June 29 th	June 30 th	July 1 st
Sun-Mercury Distance (AU)	0.422	0.426	0.430
Illuminated Fraction	57%	56%	54%
Angular Diameter (arcsec)	6.79	6.91	7.03
Mercury-Earth Radial Velocity	-28.72 km/s	-28.58 km/s	-28.42 km/s
Mercury-Sun Radial Velocity	8.12 km/s	7.77 km/s	7.39 km/s
Sun-Earth-Mercury Phase Angle	81.37°	83.34°	85.28°
True Anomaly Angle (TAA)	124°	127°	130°
Sub-Earth Point (West Long. and Lat.)	246.9°, 5.5°	242.0°, 5.7°	237.3°, 5.8°
Sub-Solar Point (West Long. and Lat.)	328.4°,0.0°	325.5°,0.0°	322.6°,0.0°

12 Table 3

Date	Seeing Sigma (")	Seeing FWHM (")	Sun-Mercury Radial velocity (mÅ)	Earth-Mercury Radial velocity (mÅ)	g-factor D2	g-factor D1
29/06/2005	1.5-2.0	2.5-3.3	159	-574	21.2	13.7
30/06/2005	0.9-1.3	1.5-2.2	152	-562	19.2	12.6
01/07/2005	1.5-2.1	2.5-3.5	145	-558	17.3	11.7