



HAL
open science

A Probabilistic Scheme for Fully Non-linear Non-local Parabolic PDEs with singular Levy measures

Arash Fahim

► **To cite this version:**

Arash Fahim. A Probabilistic Scheme for Fully Non-linear Non-local Parabolic PDEs with singular Levy measures. 2011. hal-00530432v2

HAL Id: hal-00530432

<https://hal.science/hal-00530432v2>

Preprint submitted on 19 Oct 2011

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

A PROBABILISTIC SCHEME FOR FULLY NON-LINEAR NON-LOCAL PARABOLIC PDES WITH SINGULAR LÉVY MEASURES

ARASH FAHIM

ABSTRACT. We introduce a Monte Carlo scheme for the approximation of the solutions of fully non-linear parabolic non-local PDEs. The method is the generalization of the method proposed by [Fahim-Touzi-Warin, 2011] for fully non-linear parabolic PDEs. As an independent result, we also introduce a Monte Carlo Quadrature method to approximate the integral with respect to Lévy measure which may appear inside the scheme. We consider the equations whose non-linearity is of the Hamilton–Jacobi–Belman type. We avoid the difficulties of infinite Lévy measures by truncation of the Lévy integral by some $\kappa > 0$ near 0. The first result provides the convergence of the scheme for general parabolic non-linearities. The second result provides bounds on the rate of convergence for concave non-linearities from above and below. For both results, it is crucial to choose κ appropriately dependent on h .

1. INTRODUCTION

The present paper generalizes the probabilistic numerical method in [10] for approximation of the solution of fully non-linear parabolic PDEs to non-local PDEs (integro partial differential equations). Here by non-local PDE, we mean the integro-partial differential equations which sometimes are simply referred to as IPDEs. The method is originated from [8] where a similar probabilistic numerical method is suggested based on 2BSDEs, i.e. second order backward stochastic differential equations.

A probabilistic numerical (Monte Carlo) method is introduced in [10] for fully non-linear parabolic local PDEs. A convergence result is established for a class of non-linearities and bounds for the rate of convergence for concave and Lipschitz non-linearities for the Monte Carlo method is derived. Moreover, it is shown that the error due to estimation of the conditional expectations does not change the results if enough number of sample paths are used. Finally, some numerical results are provided for mean curvature flow problem and two and five dimensional continuous-time portfolio optimization problems.

Fully non-linear PDEs arise in many problems in applied mathematics and engineering including finance. For example the problem of motion by curvature, portfolio optimization under different type of constraints, option pricing under illiquidity cost, etc. Non-local fully non-linear PDEs arise from stochastic optimization problems for controlled jump-diffusion processes e.g. problem

Date: February 11, 2011.

2000 Mathematics Subject Classification. 65C05; 49L25; 34K28.

Key words and phrases. Viscosity solution, non-local PDE, Monte Carlo approximation.

I am grateful to thank Nizar Touzi and Xavier Warin for their fruitful comments and suggestions.

of portfolio optimization in Lévy markets. There are only few examples with explicit and quasi-explicit solution; for example [3] and [4]. In some applications the dimension of the PDE is so large that classical algorithms like finite element and finite difference fail to approximate the solution in a reasonable time. The advantage of Monte Carlo methods is that they are less sensitive with respect to dimension on comparison with other methods.

As in [10], the main idea is to separate the equation into a purely linear part and a fully non-linear part. Then, we use the time discretization of a suitable jump-diffusion process to approximate the derivatives and integral term in the non-linear part. The separation into linear and non-linear part is arbitrary up to the satisfaction of some assumptions. The assumptions needed for this result are degenerate ellipticity condition for the remaining non-linearity and that the diffusion coefficient is needed to dominate the partial gradient of the remaining non-linearity with respect to its Hessian component.

The other contribution of this paper is the Monte Carlo method for approximation of the integral with respect to Lévy measure which appears in the non-local PDEs. The method is referred to in this chapter as *Monte Carlo Quadrature (MCQ)*. We treat the jumps as in [7] for finite activity jump-diffusion processes. For infinite activity jump-diffusion processes, we truncate the Lévy measure near zero and then treat them as in the finite measure case. We introduce bounds for the truncation error with respect to the derivatives of integrand and truncation level.

Although MCQ is independent of the numerical scheme, we choose to approximate the Lévy integral inside the non-linearity by MCQ. In this case, we also need to choose appropriate truncation bound with respect to time step which retains the convergence and rate of convergence as in the local case in [10].

The idea of the proof is captured from [1] and [10] for the convergence result and from [5] and [10] for the rate of convergence. However, in the non-local PDEs, we need to conquer the new difficulties due to lack of Lipschitz continuity of non-linearities appearing in many interesting IPDEs e.g. HJB equations. More precisely, if Lévy measure inside the non-local integral is not finite, then the non-local non-linearity is of HJB type will not be Lipschitz. This difficulty makes it impossible to use directly the method in [10]. We show that if the truncation level; κ ; is properly dependent on time step; h ; then one can produce the approximate solution which converges to the solution of the non-local problem.

The first result concerns the convergence of the approximate solution obtained from the Monte Carlo scheme to the viscosity solution of the final value problem. As mentioned before, infinite Lévy measure in the non-linearity makes the direct use of the method in [10] impossible due to the lack of Lipschitz continuity of the non-linearity. Moreover, if we truncate the Lévy measure, the non-linearity is Lipschitz but as truncation level tends to 0, the Lipschitz constant blows up. We solve this problem through manipulation of the original final value problem to another one whose corresponding scheme is monotone. Turning the manipulation back, we obtain a bounded approximate solution. This approximation is near the approximate function created by the proposed scheme, if the truncation level depends appropriately on h .

The second result provides the rate of convergence in the case of concave non-linearity. The proof of the rate of convergence uses the results in [5] and [6] which generalizes the result of [1] to non-local case. The method is based on approximation of the solution of the equation with regular sub and super-solutions. Plugging the regular sub or super-solution into scheme and then usage of consistency provide the upper and lower bounds. Here, we also need to impose the condition that truncation level depend appropriately on the time step in order to preserve the rate of convergence after truncation. For the rate of convergence, we also need to manipulate the equation to obtain strict monotonicity for the scheme which is a crucial requirement in using the method in [5].

Finally, as mentioned in [10] for non-local case, it is worthy of noticing the relation with the generalization of [12] to non-local case introduced in [11] which provides a deterministic game theoretic interpretation for fully non-linear parabolic problems. The game consists of two players. At each time step in a predetermined time horizon, one tries to maximize her gain and the other to minimize it by imposing a penalty term to her gain. More precisely, she starts in an initial position and chooses a vector p , a matrix Γ , and a function φ . Then, he will plug an arbitrary vector w together with p , Γ and φ in a non-linear penalty term which she should pay to be allowed to change her position by taking one step with appropriate length in the direction of vector w . At the final stage, she will earn as much as a function of her final position. As time step goes to zero, her value function at any time and any position will converge to the solution of a fully nonlinear parabolic IPDE whose non-linearity relates to the penalty term. Vector p , a matrix Γ and a function φ represent the first and second derivatives and the solution function, respectively.

The paper is organized as follows: In Section 2, the problematic features of non-local fully non-linear PDE is discussed on a naïve generalization of the Monte Carlo method from local case in [10] to non-local case. In Section 3 the Monte Carlo quadrature (MCQ) is presented as a purely Monte Carlo approximation of with the error analysis. Section 4 contains the results of convergence and asymptotic properties of the scheme.

Notations For scalars $a, b \in \mathbb{R}$, we write $a \wedge b := \min\{a, b\}$, $a \vee b := \max\{a, b\}$, $a^- := \max\{-a, 0\}$, and $a^+ := \max\{a, 0\}$.

By $\mathbb{M}(n, d)$, we denote the collection of all $n \times d$ matrices with real entries. The collection of all symmetric matrices of size d is denoted \mathbb{S}_d , and its subset of nonnegative symmetric matrices is denoted by \mathbb{S}_d^+ . For a matrix $A \in \mathbb{M}(n, d)$, we denote by A^T its transpose. For $A, B \in \mathbb{M}(n, d)$, we denote $A \cdot B := \text{Tr}[A^T B]$. In particular, for $d = 1$, A and B are vectors of \mathbb{R}^n and $A \cdot B$ reduces to the Euclidean scalar product.

We denote by \mathcal{C}_d , the space of bounded continuous functions from $[0, T]$ to \mathbb{R}^d . For a suitably smooth function φ on $Q_T := (0, T] \times \mathbb{R}^d$, we define $|\varphi|_\infty := \sup_{(t,x) \in Q_T} |\varphi(t, x)|$ and $|\varphi|_1 := |\varphi|_\infty + \sup_{Q_T \times Q_T} \frac{|\varphi(t,x) - \varphi(t',x')|}{(x-x') + |t-t'|^{\frac{1}{2}}}$.

2. PRELIMINARIES AND FEATURES FOR NON-LOCAL PDES

Let μ, σ be functions from $[0, T] \times \mathbb{R}^d$ to \mathbb{R}^d and $\mathbb{M}(d, d)$ and η be a function from $[0, T] \times \mathbb{R}^d \times \mathbb{R}^d$ to \mathbb{R}^d and $\mathbb{M}(d, d)$ and $a = \sigma^T \sigma$. Suppose the following non-local Cauchy problem:

$$-\mathcal{L}^X v(t, x) - F(t, x, v(t, x), Dv(t, x), D^2v(t, x), v(t, \cdot)) = 0, \text{ on } [0, T] \times \mathbb{R}^d, \quad (2.1)$$

$$v(T, \cdot) = g, \text{ on } \mathbb{R}^d. \quad (2.2)$$

where $F : \mathbb{R}_+ \times \mathbb{R}^d \times \mathbb{R} \times \mathbb{R}^d \times \mathbb{S}_d \times \mathcal{C}_d \rightarrow \mathbb{R}$ and \mathcal{L}^X given by:

$$\begin{aligned} \mathcal{L}^X \varphi(t, x) &:= \left(\frac{\partial \varphi}{\partial t} + \mu \cdot D\varphi + \frac{1}{2} a \cdot D^2 \varphi \right)(t, x) \\ &+ \int_{\mathbb{R}_*^d} \left(\varphi(t, x + \eta(t, x, z)) - \varphi(t, x) - \mathbb{1}_{\{|z| \leq 1\}} D\varphi(t, x) \eta(t, x, z) \right) d\nu(z). \end{aligned}$$

\mathcal{L}^X is the infinitesimal generator of a jump-diffusion, X_t , satisfying SDE:

$$dX_t = \mu(t, X_t)dt + \sigma(t, X_t)dW_t + \int_{\{|z| > 1\}} \eta(t, X_{t-}, z)J(dt, dz) + \int_{\{|z| \leq 1\}} \eta(t, X_{t-}, z)\tilde{J}(dt, dz),$$

where J and \tilde{J} are respectively a Poisson jump measure and its compensation who associate to Lévy measure ν by:

$$\begin{aligned} \nu(A) &= \mathbb{E} \left[\int_A J([0, 1], dz) \right] \\ \tilde{J}(dt, dz) &= J(dt, dz) - dt \times \nu(dz). \end{aligned}$$

For more details on jump-diffusion processes, see [2] and the references therein or the classic work of [13].

The classical solution for the problem (2.1)–(2.2) does not exist in general and therefore we appeal to the notion of viscosity solutions for non-local parabolic PDEs.

Definition 2.1. • *The viscosity sub(super)-solution of (2.1)–(2.2) is a upper semi-continuous (lower semi-continuous) function $\underline{v}(\bar{v}) : [0, T] \times \mathbb{R}^d \rightarrow \mathbb{R}$ such that:*

(1) *for any $(t_0, x_0) \in [0, T] \times \mathbb{R}^d$ and any smooth function φ with:*

$$0 = \max(\min)\{\underline{v} - \varphi\} = (\bar{v} - \varphi)(t_0, x_0)$$

We have:

$$0 \geq (\leq) -\mathcal{L}^X \varphi(t_0, x_0) - F(\cdot, \varphi, D\varphi, D^2\varphi, \varphi(\cdot))(t_0, x_0).$$

(2) $g(\cdot) \geq \underline{v}(T, \cdot) (\leq \bar{v}(T, \cdot))$.

The function v which is both viscosity sub and super solution, is called viscosity solution of (2.1)–(2.2).

• *We say that (2.1) has comparison for bounded functions if for any bounded upper semi-continuous viscosity super-solution \bar{v} and any bounded lower semi-continuous sub-solution \underline{v} , satisfying*

$$\bar{v}(T, \cdot) \geq \underline{v}(T, \cdot),$$

we have $\bar{v} \geq \underline{v}$ on $[0, T] \times \mathbb{R}^d$.

2.1. Discretization of the jump–diffusion process. Our purpose is to introduce a Monte Carlo method which approximates the solution of problem (2.1)-(2.2). For this purpose, we first need to provide a discretization for the process X .

Suppose that $h = \frac{T}{n}$, $t_i = ih$, and $\kappa \geq 0$. We define the Euler discretization of jump–diffusion process X_t with truncated Lévy measure by:

$$\hat{X}_h^{t,x,\kappa} = x + \tilde{\mu}(t, x)h + \sigma(t, x)W_h + \int_{\{|z|>\kappa\}} \eta(t, x, z) \tilde{J}([0, h], dz), \quad (2.3)$$

$$\hat{X}_{t_{i+1}}^{x,\kappa} = \hat{X}_h^{t_i, \hat{X}_{t_i}^{x,\kappa}, \kappa} \quad \text{and} \quad \hat{X}_0^{x,\kappa} = x. \quad (2.4)$$

where $\tilde{\mu}(t, x) = \mu(t, x) + \int_{\{|z|>1\}} \eta(t, x, z) \nu(dz)$ and we make the choice of $\kappa = 0$ when ν is a finite measure. Let \tilde{N}_t^κ and N_t^κ be respectively the Poisson process derived from jump measure J by counting all jumps of size greater than κ which happen in time interval $[0, t]$ and its compensation, i.e.

$$N_t^\kappa = \int_{\{|z|>\kappa\}} J([0, t], dz) \quad \text{and} \quad \tilde{N}_t^\kappa = \int_{\{|z|>\kappa\}} \tilde{J}([0, T], dz). \quad (2.5)$$

One can write the jump part of $\hat{X}_h^{t,x,\kappa}$ as a compound Poisson process (see for example [9])

$$\hat{X}_h^{t,x,\kappa} = x + \mu_\kappa(t, x)h + \sigma(t, x)W_h + \sum_{i=1}^{N_h^\kappa} \eta(t, x, Z_i), \quad (2.6)$$

where $\mu_\kappa(t, x) = \mu(t, x) - \int_{\{\kappa < |z| \leq 1\}} \eta(t, x, z) \nu(dz)$, Z_i s are i.i.d. \mathbb{R}_*^d -valued random variables, independent of W and N^κ , and distributed as $\mathbf{1}_{\{|z|>\kappa\}} \frac{1}{\lambda_\kappa} \nu(dz)$.

2.2. The scheme for non–local fully non–linear parabolic PDEs. In this section, we introduce a probabilistic scheme by following directly the same idea as the scheme for the local PDEs. Then, we consider some problems which prevents us to utilize the scheme in many interesting applications. Therefore, we introduce a modified version of the scheme which works for the class of non–linearities of HJB type (Hamilton–Jacobi–Bellman).

Following the same idea as in [10], one can obtain the following immature scheme.

$$v^h(T, \cdot) = g \quad \text{and} \quad v^h(t_i, x) = \mathbf{T}_h[v^h](t_i, x), \quad (2.7)$$

where for every function $\psi : \mathbb{R}_+ \times \mathbb{R}^d \rightarrow \mathbb{R}$ with exponential growth:

$$\begin{aligned} \mathbf{T}_h[\psi](t, x) &:= \mathbb{E} \left[\psi \left(t + h, \hat{X}_h^{t,x} \right) \right] + h F_h(t, x, \mathcal{D}_h \psi, \psi(t + h, \cdot)), \\ \mathcal{D}_h \psi &:= (\mathcal{D}_h^0 \psi, \mathcal{D}_h^1 \psi, \mathcal{D}_h^2 \psi), \end{aligned} \quad (2.8)$$

where

$$\mathcal{D}_h^k \psi(t, x) := \mathbb{E} \left[\psi(t + h, \hat{X}_h^{t,x,\kappa}) H_k^h(t, x) \right], \quad k = 0, 1, 2, \quad (2.9)$$

where

$$H_0^h = 1, \quad H_1^h = (\sigma^T)^{-1} \frac{W_h}{h}, \quad H_2^h = (\sigma^T)^{-1} \frac{W_h W_h^T - h \mathbf{I}_d}{h^2} \sigma^{-1}.$$

The details of approximation of derivatives with (4.5) can be found in Lemma 2.1 in [10].

We intend to extend the result of [10] to the non-local case. First observe that there is an obvious extension which could be done immediately by adding the following assumptions to Assumption **F** in [10], i.e.

Assumption F (i) *The nonlinearity F is Lipschitz-continuous with respect to (x, r, p, γ, ψ) uniformly in t , and $|F(\cdot, \cdot, 0, 0, 0, 0)|_\infty < \infty$;*

(ii) *F is elliptic and dominated by the diffusion of the linear operator \mathcal{L}^X , i.e.*

$$\nabla_\gamma F \leq a \quad \text{on} \quad \mathbb{R}^d \times \mathbb{R} \times \mathbb{R}^d \times \mathbb{S}_d \times \mathbb{C}_d; \quad (2.10)$$

(iii) $F_p \in \text{Image}(F_\gamma)$ and $|F_p^T F_\gamma^- F_p|_\infty < +\infty$.

We remind that the non-local non-linearity F is called elliptic if

(1) F is non-decreasing on the second derivative component, i.e.

$$F(t, x, r, p, \gamma_1, \psi) \leq F(t, x, r, p, \gamma_2, \psi) \quad \text{for} \quad \gamma_1 \leq \gamma_2.$$

(2) F is non-decreasing on the non-local component, i.e.

$$F(t, x, r, p, \gamma, \psi_1) \leq F(t, x, r, p, \gamma, \psi_2) \quad \text{for} \quad \psi_1 \leq \psi_2.$$

Then we have the following Theorem.

Theorem 2.1. *Let Assumption **F** in hold true, and $|\mu|_1, |\sigma|_1 < \infty$ and σ is invertible. Also assume that the fully nonlinear PDE (2.1) has comparison for bounded functions. Then for every bounded Lipschitz function g , there exists a bounded function v so that*

$$v^h \longrightarrow v \quad \text{locally uniformly.}$$

In addition, v is the unique bounded viscosity solution of problem (2.1)-(2.2).

The proof is an straight forward implementation of the Subsection 3.2 of [10].

One of the major class of fully non-linear PDEs, is the class of HJB equations which come from stochastic control problems arising in many application including finance. However, The non-linearity of HJB equations do not satisfies Assumption **F** in general. Even for local PDEs of HJB type, Assumption **F** is not valid, because F is not uniformly Lipschitz with respect to x . In addition, when the Lévy measure ν is an infinite Lévy measure, there is no chance for F to be uniformly Lipschitz within respect to ψ . Therefore, we need to develop another theory for HJB equations.

The other problem which occurs in many applications is the lack of explicit form for non-linearity F . We present the following example in order to mention this problem.

Example 2.1. *Suppose we want to implement the scheme for the fully non-linear equation of the form:*

$$\begin{aligned} -v_t - F(x, Dv(t, x), D^2v(t, x), v(t, \cdot)) &= 0 \\ v(T, \cdot) &= g(\cdot), \end{aligned}$$

where

$$F(x, p, \gamma, \psi) := \sup_{\theta \in \mathbb{R}_+} \left\{ \mathcal{L}^\theta(p, \gamma) + \int_{\mathbb{R}_*} \psi(x + \theta z) \nu(dz) \right\} \quad (2.11)$$

$$\mathcal{L}^\theta(p, \gamma) := \theta b p + \frac{1}{2} \theta^2 a^2 \gamma \quad (2.12)$$

$$\mathcal{I}(x, \psi)^\theta := \int_{\mathbb{R}_*} \psi(x + \theta z) \nu(dz). \quad (2.13)$$

This fully non-linear equation solves the problem of portfolio management for one asset in the Black–Scholes model including jumps in asset price. For the sake of simplicity, for the moment we forget about infinite activity jumps. Observe that if $\nu = 0$ (the asset price do not jump) then F becomes of the form:

$$F(x, p, \gamma, \psi) := \sup_{\theta \in \mathbb{R}_+} \left\{ \theta b p + \frac{1}{2} \theta^2 a^2 \gamma \right\}.$$

which could be given in explicit form by:

$$F(x, p, \gamma, \psi) := -\frac{(bp)^2}{2a^2\gamma},$$

and the scheme could be easily implemented as in [10] as well as more complicated examples. But, when $\nu \neq 0$ (jump do exists), the explicit form for F is not known and the supremum should be approximated. This problem is in common with other numerical methods for fully non-linear PDEs e.g. finite difference. Although his problem is obviously beyond the subject of this paper, we addressed it in this paper in order to mention that why we need to approximate the integral inside the supremum.

The other problem, which appears in high dimensions, is the calculation of Lévy integral inside supremum. Some numerical methods to approximate the supremum based on the calculation of the linear operator $\mathcal{L}^\theta + \mathcal{I}^\theta$ inside the supremum for different θ s. Therefore, we proposed a Monte Carlo Quadrature method to approximate the integral in a purely probabilistic way. The MCQ could be considered independently in other applications.

From now on, we relax the assumption that ν is a finite measure. To be precise, we need to suppose that (2.13) is of the form

$$\mathcal{I}(x, \psi) := \int_{\mathbb{R}_*} (\psi(x + \theta z) - \psi(x) - \mathbf{1}_{\{|z| \leq 1\}} \theta D\psi(x) \cdot z) \nu(dz).$$

In this case, there are two ways to work with singular Lévy measure in numerical experiments; one is to truncate Lévy measure near zero (as we also did for discretization of X) and the other is to approximate infinite small jumps by a Brownian motion. In both cases, the general form for the approximate F is

$$F_\kappa(x, r, p, \gamma, \psi) := \sup_{\theta \in \mathbb{R}_+} \left\{ c_\kappa r + \theta b_\kappa p + \frac{1}{2} \theta^2 a^2 \gamma + \int_{\{|z| > \kappa\}} \psi(x + \theta z) \nu(dz) \right\}.$$

where

$$c_\kappa := \int_{\{|z|>\kappa\}} \nu(dz) \text{ and } b_\kappa := b \int_{\{1 \geq |z|>\kappa\}} z \nu(dz).$$

We will introduce the modified scheme (4.3) in Section 4 based on the approximation of non-linearity F obtained from truncation of infinite Lévy measure and MCQ and then provide asymptotic results as in [10] for non-local case.

The generalization of the result of [10] for non-local PDEs would be easy if the function F_κ were Lipschitz uniform on κ . But, for infinite Lévy measures, this is never the case. To overcome this problem, we will show that κ could be chosen dependent on h , so that the corresponding scheme satisfies the requirements of [1] for the proof of convergence.

3. MONTE CARLO QUADRATURE (MCQ)

In this section, we propose a Monte Carlo method the value of the following Lévy generator:

$$\mathcal{I}[\varphi](x) := \int_{\mathbb{R}_*^d} (\varphi(x + \eta(z)) - \varphi(x) - \mathbf{1}_{\{|z| \leq 1\}} \eta(z) \cdot D\varphi(x)) \nu(dz). \quad (3.1)$$

The method is pure Monte Carlo method to approximate (3.1) and, therefore could be used in the approximation of Lévy integral inside the scheme (4.3). Because, the result of this section is independent of the numerical scheme (4.3) introduced in this paper, we organize this Section so that one can read it independently from other Section.

Through out this Section, we drop the dependency with respect to (t, x) or other variables and for the sake of simplicity and just write $\eta(z)$. (For example in assumption **F** in Section 4 $\eta^{\alpha, \beta}(t, x, z)$ which depends on (t, x, z, α, β) will be considered as $\eta(z)$).

Notice that in order for (3.1) to be well-defined for regular functions, we impose the following assumption on η :

$$\frac{|\eta(z)|}{|z| \wedge 1} \leq C, \text{ for some constant } C. \quad (3.2)$$

We present MCQ in three cases with respect to the behavior of Lévy measure near zero:

- finite measure; $\int_{\{|z| \leq 1\}} \nu(dz) < \infty$,
- infinite measure;
 - case I: $\int_{\{|z| \leq 1\}} |\eta(z)| \nu(dz) < \infty$,
 - case II: $\int_{\{|z| \leq 1\}} |\eta(z)|^2 \nu(dz) < \infty$.

3.1. Finite Lévy Measure. When Lévy measure is finite, we choose $\kappa = 0$. In this case, we introduce Lemma 3.1 which proposes a way to approximate the Lévy integral of general form:

$$\int_{\mathbb{R}_*^d} \varphi(x + \eta(z)) \zeta(z) d\nu(z), \quad (3.3)$$

and then we use this Lemma to approximate the Lévy infinitesimal generator (3.1).

Let J be a jump Poisson measure with intensity given by Lévy measure ν , and $\{N_t\}_{t \geq 0}$ be the Poisson process given by $N_t = \int_0^t \int_{\mathbb{R}_*^d} J(ds, dz)$ whose intensity is $\lambda := \int_{\mathbb{R}_*^d} \nu(dz)$. By (2.6), we can write \hat{X}^x by

$$\hat{X}_t^x = x + \mu_0 t + \sigma W_t + \sum_{i=1}^{N_t} \eta(Z_i) \quad (3.4)$$

where Z_i s are i.i.d. random variables with law $\frac{1}{\lambda} \nu(dz)$. We also introduce a Lévy process Y_t by

$$Y_t = \sum_{i=1}^{N_t} \zeta(Z_i). \quad (3.5)$$

Next Lemma shows that (3.3) could be approximated by a Monte Carlo formula purely free of integration.

Lemma 3.1. *Let*

$$\hat{\nu}_h^{\eta, \zeta}(\varphi)(x) := \mathbb{E} \left[\int_{\mathbb{R}_*^d} \varphi(\hat{X}_h^x + \eta(z)) \zeta(z) d\nu(z) \right]. \quad (3.6)$$

Then, for every bounded function $\varphi : \mathbb{R}^d \rightarrow \mathbb{R}$:

$$\hat{\nu}_h^{\eta, \zeta}(\varphi)(x) = \frac{1}{h} \mathbb{E}[\varphi(\hat{X}_h^x) Y_h].$$

Proof. For the sake of simplicity, we just concentrate on the jump part of process \hat{X}^x and without loss of generality, we write $\hat{X}_h^x = x + \sum_{i=1}^{N_h} \eta(Y_i)$. The right hand side can be expressed as:

$$\mathbb{E} \left[\varphi(\hat{X}_h^x) Y_h \right] = e^{-\lambda h} \sum_{n=0}^{\infty} \mathbb{E} \left[\varphi(\hat{X}_h^x) Y_h | N_h = n \right] \frac{(\lambda h)^n}{n!}.$$

Then by (3.4)–(3.5),

$$\begin{aligned} \mathbb{E} \left[\varphi(\hat{X}_h^x) Y_h \right] &= e^{-\lambda h} \lambda h \sum_{n=1}^{\infty} \mathbb{E} \left[\varphi \left(x + \sum_{i=1}^n \eta(Z_i) \right) \left(\sum_{j=1}^n \zeta(Z_j) \right) \right] \frac{(\lambda h)^{n-1}}{n!} \\ &= e^{-\lambda h} \lambda h \sum_{n=1}^{\infty} \frac{(\lambda h)^{n-1}}{n!} \sum_{j=1}^n \mathbb{E} \left[\varphi \left(x + \sum_{i=1}^n \eta(Z_i) \right) \zeta(Z_j) \right]. \end{aligned}$$

Notice that in the above expression, the summation starts from $n = 1$ because $Y_h = 0$ when $N_h = 0$. Because Z_i s are i.i.d. one can conclude that,

$$\sum_{j=1}^n \mathbb{E} \left[\varphi \left(x + \sum_{i=1}^n \eta(Z_i) \right) \zeta(Z_j) \right] = n \mathbb{E} \left[\varphi \left(x + \sum_{i=1}^n \eta(Z_i) \right) \zeta(Z_1) \right]$$

Then, one can write

$$\mathbb{E} \left[\varphi \left(x + \eta(Z_1) + \sum_{i=2}^n \eta(Z_i) \right) \zeta(Z_1) \right] = \mathbb{E} \left[\varphi \left(\eta(Z) + \hat{X}_h^x \right) \zeta(Z) | N_h = n - 1 \right],$$

where Z is dependent of Z_i s but has the same law as Z_i s. Therefore, we can conclude that:

$$\mathbb{E} \left[\varphi(\hat{X}_h^x) Y_h \right] = e^{-\lambda h} \lambda h \sum_{n=1}^{\infty} \mathbb{E} \left[\varphi(\eta(Z) + \hat{X}_h^x) \zeta(Z) | N_h = n - 1 \right] \frac{(\lambda h)^{n-1}}{(n-1)!}.$$

But, we know that

$$e^{-\lambda h} \sum_{n=1}^{\infty} \mathbb{E} \left[\varphi(\eta(Z) + \hat{X}_h^x) \zeta(Z) | N_h = n - 1 \right] \frac{(\lambda h)^{n-1}}{(n-1)!} = \mathbb{E} \left[\varphi(\eta(Z) + \hat{X}_h^x) \zeta(Z) \right]$$

Therefore,

$$\mathbb{E} \left[\varphi(\hat{X}_h^x) Y_h \right] = \lambda h \mathbb{E} \left[\varphi(\eta(Z) + \hat{X}_h^x) \zeta(Z) \right].$$

Because the density of Z is $\frac{\nu(dz)}{\lambda}$,

$$\mathbb{E} \left[\varphi(\hat{X}_h^x) Y_h \right] = h \mathbb{E} \left[\int_{\mathbb{R}_*^d} \varphi(\eta(z) + \hat{X}_h^x) \zeta(z) d\nu(z) \right].$$

□

In the light of Lemma (3.1), we propose the following approximation for (3.1):

$$\mathcal{I}_h[\varphi](x) := \hat{\nu}_h^{\eta,1} - \varphi(x) \int_{\mathbb{R}_*^d} \nu(dz) - D\varphi(x) \cdot \int_{\mathbb{R}_*^d} \eta(z) \nu(dz).$$

Next Lemma provide error bound for this approximation.

Lemma 3.2. *For any Lipschitz function φ we have:*

$$|(\mathcal{I}_h - \mathcal{I})[\varphi]|_{\infty} \leq C\sqrt{h}|D\varphi|_{\infty}. \quad (3.7)$$

Proof. As a direct consequence of Lemma (3.1), $\hat{\nu}_h^{\eta,1} = \frac{1}{h} \mathbb{E}[\varphi(\hat{X}_h^x) N_h]$. Therefore, one can conclude that,

$$|(\mathcal{I} - \mathcal{I}_h)[\varphi]|_{\infty} \leq C|D\varphi|_{\infty} \mathbb{E} \left[|\hat{X}_h^x - x| \right].$$

So, because

$$\mathbb{E} \left[|\hat{X}_h^x - x| \right] \leq C \left(h \int_{\mathbb{R}_*^d} |\eta(z)| \nu(dz) + \sqrt{h} \right), \quad (3.8)$$

which provides the result. □

3.2. Infinite Lévy Measure. In the case of singular Lévy measure, we truncate Lévy measure near zero and reduce the problem to a finite measure. In other words, for any $\kappa > 0$ we have the truncation approximation of integral operator (3.1).

$$\mathcal{I}_{\kappa}[\varphi](x) := \int_{\{|z|>\kappa\}} (\varphi(x + \eta(z)) - \varphi(x) - \mathbb{1}_{\{|z|\leq 1\}} \eta(z) \cdot D\varphi(x)) \nu(dz).$$

Then, we use Lemma (3.1) to present the MCQ approximation for (3.1).

$$\mathcal{I}_{\kappa,h}[\varphi](x) := \hat{\nu}_{\kappa,h}^{\eta,1} - \varphi(x) \int_{\{|z|>\kappa\}} d\nu(z) - \int_{\{1 \geq |z|>\kappa\}} \eta(t,x,z) \cdot D\varphi(x) d\nu(z),$$

where by Lemma (3.1)

$$\hat{\nu}_{\kappa,h}^{\eta,1} := \int_{\{|z|>\kappa\}} \varphi(\hat{X}_h^{x,\kappa} + \eta(t,x,z))\nu(dz) = h^{-1}\mathbb{E} \left[\varphi(\hat{X}_h^{x,\kappa})N_h^\kappa \right]$$

Following Lemma provides the error of MCQ approximation of (3.1) in the case of infinite Lévy measure.

Lemma 3.3. *Let function φ be Lipschitz.*

(1) *If $\int_{\{|z|\leq 1\}} |z|\nu(dz) < \infty$, then*

$$|(\mathcal{I}_{\kappa,h} - \mathcal{I})[\varphi]|_\infty \leq C|D\varphi|_\infty \left(\sqrt{h} + \int_{\{0<|z|\leq\kappa\}} |z|\nu(dz) \right). \quad (3.9)$$

(2) *If $\int_{\{|z|\leq 1\}} |z|^2\nu(dz) < \infty$, then*

$$|(\mathcal{I}_{\kappa,h} - \mathcal{I})[\varphi]|_\infty \leq C \left(|D\varphi|_\infty \left(\sqrt{h} + h \int_{\{|z|>\kappa\}} |z|\nu(dz) \right) + |D^2\varphi|_\infty \int_{\{0<|z|\leq\kappa\}} |z|^2\nu(dz) \right). \quad (3.10)$$

Proof.

(1) Notice that,

$$|(\mathcal{I} - \mathcal{I}_{\kappa,h})[\varphi]|_\infty \leq |(\mathcal{I} - \mathcal{I}_\kappa)[\varphi]|_\infty + |(\mathcal{I}_\kappa - \mathcal{I}_{\kappa,h})[\varphi]|_\infty.$$

By (3.2), the truncation error is given by:

$$|(\mathcal{I} - \mathcal{I}_\kappa)[\varphi]|_\infty \leq 2|D\varphi|_\infty \int_{\{0<|z|\leq\kappa\}} |\eta(z)|\nu(dz). \quad (3.11)$$

On the other hand, by (3.8) and (3.2), we observe that

$$\begin{aligned} |(\mathcal{I}_\kappa - \mathcal{I}_{\kappa,h})[\varphi]|_\infty &\leq C|D\varphi|_\infty \left(h \int_{\{|z|>\kappa\}} |\eta(z)|\nu(dz) + \sqrt{h} \right) \\ &\leq C|D\varphi|_\infty \left(h \int_{\{|z|>\kappa\}} |z|\nu(dz) + \sqrt{h} \right) \end{aligned}$$

which together with (3.11) provides the result.

(2) By (3.2), the truncation error is given by:

$$|(\mathcal{I} - \mathcal{I}_\kappa)[\varphi]|_\infty \leq C|D^2\varphi|_\infty \int_{\{0<|z|\leq\kappa\}} |z|^2\nu(dz), \quad (3.12)$$

for any function φ with bounded derivatives up to second order. On the other hand, (3.8) allows us to calculate the Monte Carlo error by:

$$|(\mathcal{I}_\kappa - \mathcal{I}_{\kappa,h})[\varphi]|_\infty \leq C|D\varphi|_\infty \left(h \int_{\{|z|>\kappa\}} |z|\nu(dz) + \sqrt{h} \right)$$

which completes the proof. \square

4. ASYMPTOTIC RESULTS

This section is devoted to the convergence result for the scheme (4.3). We first remind the notion of viscosity solution and provide the assumptions required for the main results together with the statement of main results. Then we provide the proof of the results in two following subsection.

we need to impose the following assumption on the non-linearity F to obtain the convergence Theorem.

Assumption IHJB1: *Function F satisfies:*

$$\frac{1}{2}a(t, x) \cdot \gamma + \mu(t, x) \cdot p + F(t, x, r, p, \gamma, \psi) := \inf_{\alpha \in \mathcal{A}} \sup_{\beta \in \mathcal{B}} \left\{ \mathcal{L}^{\alpha, \beta}(t, x, r, p, \gamma) + \mathcal{I}^{\alpha, \beta}(t, x, r, p, \gamma, \psi) \right\}$$

for given sets \mathcal{A} and \mathcal{B} where

$$\mathcal{L}^{\alpha, \beta}(t, x, r, p, \gamma) := \frac{1}{2}a^{\alpha, \beta}(t, x) \cdot \gamma + b^{\alpha, \beta}(t, x) \cdot p + c^{\alpha, \beta}(t, x)r + k^{\alpha, \beta}(t, x),$$

and

$$\mathcal{I}^{\alpha, \beta}(t, x, r, p, \psi) := \int_{\mathbb{R}_*^d} \left(\psi \left(x + \eta^{\alpha, \beta}(t, x, z) \right) - r - \mathbb{1}_{\{|z| \leq 1\}} \eta^{\alpha, \beta}(t, x, z) \cdot p \right) \nu(dz)$$

where for any $(\alpha, \beta) \in \mathcal{A} \times \mathcal{B}$, $a^{\alpha, \beta}$, $b^{\alpha, \beta}$, $c^{\alpha, \beta}$, $k^{\alpha, \beta}$ and $\eta^{\alpha, \beta}$ satisfy

$$\sup_{\alpha \in \mathcal{A}, \beta \in \mathcal{B}} \left\{ |a^{\alpha, \beta}|_1 + |b^{\alpha, \beta}|_1 + |c^{\alpha, \beta}|_1 + |k^{\alpha, \beta}|_1 + \frac{|\eta^{\alpha, \beta}(\cdot, z)|_1}{|z| \wedge 1} \right\} < \infty.$$

The non-linearity is dominated by the diffusion of the linear operator \mathcal{L}^X , i.e. for any t, x, z, α and β

$$|a^- \cdot a^{\alpha^*, \beta^*}|_1 < \infty \text{ and } 0 \leq a^{\alpha, \beta} \leq a, \quad (4.1)$$

$$\eta^{\alpha, \beta}, b^{\alpha, \beta} \in \text{Image}(a^{\alpha, \beta}) \text{ and } \sup_{\alpha \in \mathcal{A}, \beta \in \mathcal{B}} |(b^{\alpha, \beta})^\top (a^{\alpha, \beta}) - b^{\alpha, \beta}|_\infty < \infty, \quad (4.2)$$

$$\sup_{\alpha \in \mathcal{A}, \beta \in \mathcal{B}} \frac{|(\eta^{\alpha, \beta})^\top (a^{\alpha, \beta}) - b^{\alpha, \beta}|_\infty}{1 \wedge |z|} < \infty$$

$$\sup_{\alpha \in \mathcal{A}, \beta \in \mathcal{B}} \frac{|(\eta^{\alpha, \beta})^\top (a^{\alpha, \beta}) - \eta^{\alpha, \beta}|_\infty}{1 \wedge |z|^2} < \infty.$$

Remark 4.1. A function F which satisfies Assumption **IHJB1** is not well-defined for arbitrary $(t, x, r, p, \gamma, \psi) \in \mathbb{R}_+ \times \mathbb{R}^d \times \mathbb{R} \times \mathbb{R}^d \times \mathbb{S}_d \times \mathcal{C}_d$. But, for any second order differentiable function, ψ , with bounded derivatives with respect to x , $F(t, x, \psi(t, x), D\psi(t, x), D^2\psi(t, x), \psi(t, \cdot))$ is well-defined.

Now, we propose a Monte Carlo scheme for (2.1)-(2.2) based on the same idea as in [10], and also the approximation of the non-linearity.

$$v^{\kappa, h}(T, \cdot) = g \quad \text{and} \quad v^{\kappa, h}(t_i, x) = \mathbf{T}_{\kappa, h}[v^{\kappa, h}](t_i, x), \quad (4.3)$$

where for every function $\psi : \mathbb{R}_+ \times \mathbb{R}^d \rightarrow \mathbb{R}$ with exponential growth:

$$\begin{aligned} \mathbf{T}_{\kappa, h}[\psi](t, x) &:= \mathbb{E} \left[\psi \left(t + h, \hat{X}_h^{t, x, \kappa} \right) \right] + hF_{\kappa, h}(t, x, \mathcal{D}_h\psi, \psi(t + h, \cdot)), \\ \mathcal{D}_h\psi &:= (\mathcal{D}_h^0\psi, \mathcal{D}_h^1\psi, \mathcal{D}_h^2\psi), \end{aligned} \quad (4.4)$$

$$F_{\kappa,h}(t, x, r, p, \gamma, \psi) = \inf_{\alpha \in \mathcal{A}} \sup_{\beta \in \mathcal{B}} \left\{ \frac{1}{2} a^{\alpha,\beta}(t, x) \cdot \gamma + b^{\alpha,\beta}(t, x) \cdot p + c^{\alpha,\beta}(t, x) r + k^{\alpha,\beta}(t, x) \right. \\ \left. + \int_{\{|z| \geq \kappa\}} \left(\hat{\nu}_h^{\eta^{\alpha,\beta}, 1}(\psi(t, \cdot))(x) - r - \eta^{\alpha,\beta}(t, x, z) \cdot p \right) \nu(dz) \right\},$$

and

$$\mathcal{D}_h^k \psi(t, x) := \mathbb{E} \left[\psi(t+h, \hat{X}_h^{t,x,\kappa}) H_k^h(t, x) \right], \quad k = 0, 1, 2, \quad (4.5)$$

where

$$H_0^h = 1, \quad H_1^h = (\sigma^T)^{-1} \frac{W_h}{h}, \quad H_2^h = (\sigma^T)^{-1} \frac{W_h W_h^T - h \mathbf{I}_d}{h^2} \sigma^{-1}.$$

The details of approximation of derivatives with (4.5) can be found in Lemma 2.1 in [10]. In order to have the convergence result, we also need to impose the following assumption over $F_{\kappa,h}$.

Assumption Inf-Sup: For any $\kappa > 0$, $t \in [0, T]$, x and $x' \in \mathbb{R}^d$ and any Lipschitz functions ψ and φ , there exists a $(\alpha^*, \beta^*) \in \mathcal{A} \times \mathcal{B}$ such that

$$\Phi_{\kappa}^{\alpha^*, \beta^*}[\psi, \varphi](t, x, x') = \mathcal{J}_{\kappa}^{\alpha^*, \beta^*}[\psi](t, x) - \mathcal{J}_{\kappa}^{\alpha^*, \beta^*}[\varphi](t, x')$$

where

$$\Phi_{\kappa}^{\alpha, \beta}[\psi, \varphi](t, x) := \inf_{\alpha} \mathcal{J}_{\kappa}^{\alpha, \beta}[\psi](t, x) - \sup_{\beta} \mathcal{J}_{\kappa}^{\alpha, \beta}[\varphi](t, x'), \quad (4.6)$$

and

$$\mathcal{J}_{\kappa}^{\alpha, \beta}[\phi](t, x) := \frac{1}{2} a^{\alpha, \beta} \cdot D^2 \phi(t, x) + b^{\alpha, \beta} \cdot D \phi(t, x) + c^{\alpha, \beta} \phi(t, x) + k^{\alpha, \beta}(t, x) \\ + \int_{\{|z| \geq \kappa\}} \left(\hat{\nu}_h^{\eta^{\alpha, \beta}, 1}(\phi(t, \cdot))(x) - \phi(t, x) - \eta^{\alpha, \beta}(t, x, z) \cdot D \phi(t, x) \right) \nu(dz).$$

The first result concerns the convergence of the convergence of $v^{\kappa, h}$ for κ appropriately chosen with respect to h .

Theorem 4.1 (Convergence). Let η , μ and σ be bounded and Lipschitz continuous on x uniformly on t and z , σ is invertible and Assumptions **IHJB1** and **Inf-Sup** hold true, and assume that (2.1) has comparison for bounded functions. Then, if κ_h is such that:

$$\lim_{h \rightarrow 0} \kappa_h = 0 \quad \text{and} \quad \limsup_{h \rightarrow 0} \theta_{\kappa_h}^2 h = 0 \quad (4.7)$$

where

$$\theta_{\kappa} := \sup_{\alpha, \beta} |\theta_{\kappa}^{\alpha, \beta}|_{\infty}, \quad (4.8)$$

with

$$\theta_{\kappa}^{\alpha, \beta} := c^{\alpha, \beta} + \int_{\{|z| \geq \kappa\}} \nu(dz) + \frac{1}{4} \left(b^{\alpha, \beta} - \int_{\{1 > |z| \geq \kappa\}} \eta^{\alpha, \beta}(z) \nu(dz) \right)^T \\ \times (a^{\alpha, \beta}) - \left(b^{\alpha, \beta} - \int_{\{1 > |z| \geq \kappa\}} \eta^{\alpha, \beta}(z) \nu(dz) \right),$$

then $v^{\kappa_h, h}$ converges to some function v locally uniform. In addition, v is the unique viscosity solution of (2.1)-(2.2).

Specially, if Lévy measure is finite for the choice of $\kappa_h = 0$ the assertion of the Theorem hold true.

Remark 4.2. It is always possible to choose κ_h such that (4.7) is satisfied. To see this, notice that θ_κ in (4.8) is non-increasing on κ

$$\lim_{\kappa \rightarrow 0} \theta_\kappa = +\infty \text{ and } \limsup_{\kappa \rightarrow \infty} \theta_\kappa < \infty.$$

Then, we define $\kappa_h := \inf\{\kappa | \theta_\kappa \leq h^{-\frac{1}{2}}\} + h$. By the definition of κ_h , $\theta_{\kappa_h} \leq h^{-\frac{1}{2}}$. Because Observe that κ_h is non-decreasing with respect to h and $\lim_{h \rightarrow 0} \kappa_h = 0$.

If there exists a q such that, $q := \lim_{h \rightarrow 0} \kappa_h > 0$, then, for $\kappa < q$, we would have $\theta_\kappa = \infty$ which obviously contradicts the fact that for $\kappa > 0$, $\theta_\kappa < \infty$. Therefore, κ_h satisfies (4.8).

Remark 4.3. The choice of κ_h in the above Theorem seems to be crucial for the convergence. Otherwise, we only have the following convergence result.

Proposition 4.1. *Under the same assumption as Theorem 4.1, when Lévy measure ν is infinite, for every Lipschitz bounded function g , we have*

$$\lim_{\kappa \rightarrow 0} \lim_{h \rightarrow 0} v^{\kappa, h} = v$$

where v is the unique viscosity solution of (2.1)-(2.2) assuming that it exists.

Proof. . Let v^κ be the solution of the following problem:

$$-\mathcal{L}^X v^\kappa(t, x) - F_\kappa(t, x, v^\kappa(t, x), Dv^\kappa(t, x), D^2v^\kappa(t, x), v^\kappa(t, \cdot)) = 0, \text{ on } [0, T] \times \mathbb{R}^d, \quad (4.9)$$

$$v^\kappa(T, \cdot) = g(\cdot), \text{ on } \mathbb{R}^d. \quad (4.10)$$

where $F_\kappa : \mathbb{R}_+ \times \mathbb{R}^d \times \mathbb{R} \times \mathbb{R}^d \times \mathcal{S}_d \times \mathcal{C}_d \rightarrow \mathbb{R}$ is given by:

$$F_\kappa(t, x, r, p, \gamma, \psi) := \inf_{\alpha \in \mathcal{A}} \sup_{\beta \in \mathcal{B}} \left\{ \mathcal{L}^{\alpha, \beta}(t, x, r, p, \gamma) + \mathcal{I}_\kappa^{\alpha, \beta}(t, x, r, p, \gamma, \psi) \right\}$$

where

$$\mathcal{I}_\kappa^{\alpha, \beta}(t, x, r, p, \gamma, \psi) := \int_{\{|z| \geq \kappa\}} (\psi(x + \eta^{\alpha, \beta}(t, x, z)) - r - \mathbf{1}_{\{|z| \leq 1\}} \eta^{\alpha, \beta}(t, x, z) \cdot p) \nu(dz) \quad (4.11)$$

where $a^{\alpha, \beta}$, $b^{\alpha, \beta}$, $c^{\alpha, \beta}$, $k^{\alpha, \beta}$ and $\eta^{\alpha, \beta}$ are as in Assumption **IHJB1**. Let $v^{\kappa, h}$ be the approximate solution given by the scheme (4.3). Let $\kappa > 0$ be fixed. Because the truncated Lévy measure is finite, by Theorem 4.1, $v^{\kappa, h}$ converges to v^κ locally uniformly as $h \rightarrow 0$. Let v^κ be the solution of (4.9)-(4.10). By Theorem 5.1 of [6] and Assumption **F**, we have:

$$|v - v^\kappa|_\infty \leq C \sup_{\alpha, \beta} \left\{ \left(\int_{0 < |z| < \kappa} |\eta^{\alpha, \beta}(\cdot, z)|_\infty^2 \nu(dz) \right)^{\frac{1}{2}} \right\} \quad (4.12)$$

$$\leq C \left(\int_{0 < |z| < \kappa} |z|_\infty^2 \nu(dz) \right)^{\frac{1}{2}}. \quad (4.13)$$

Therefore, one can choose $\kappa > 0$ so that $|v^\kappa - v|_\infty$ be small enough. Then, when h goes to 0, $v^{\kappa,h}$ converges to v^κ . \square

The above limit proposes to implement the numerical scheme in two steps:

- First by choosing κ so that v^κ is near enough to v , we obtain a uniform approximation of v .
- Second by sending $h \rightarrow 0$, we obtain locally uniform convergence of $v^{\kappa,h}$ to v^κ .

Notice that the above convergence is not uniformly on (κ, h) . However, the convergence in Theorem 4.1, is uniform on h when the choice of κ is made suitably dependent on h .

Remark 4.4. By Remark 3.7 in [10], the boundedness condition on g can be relaxed.

In order to obtain the rate of convergence result, we impose Assumptions **IHJB2** and **IHJB2+** which restrict us to concave non-linearities.

Assumption IHJB2 *The non-linearity F satisfies Assumption **IHJB1** with \mathcal{B} be a singleton set*

Remark 4.5. Therefore, when the non-linearity F satisfies **IHJB**, we can drop the super script β and write F by

$$F(t, x, r, p, \gamma, \psi) := \inf_{\alpha \in \mathcal{A}} \{ \mathcal{L}^\alpha(t, x, r, p, \gamma) + \mathcal{I}^\alpha(t, x, r, p, \gamma, \psi) \}$$

where

$$\mathcal{L}^\alpha(t, x, r, p, \gamma) := \frac{1}{2} \text{Tr} [(a^\alpha)^\top] (t, x) \gamma + b^\alpha(t, x) p + c^\alpha(t, x) r + k^\alpha(t, x),$$

and

$$\mathcal{I}^\alpha(t, x, r, p, \psi) := \int_{\mathbb{R}_*^d} (\psi(x + \eta^\alpha(t, x, z)) - r - \mathbf{1}_{\{|z| \leq 1\}} \eta^\alpha(t, x, z) \cdot p) \nu(dz).$$

In this case, the non-linearity is a concave function of (r, p, γ, ψ) .

Assumption IHJB+ *The non-linearity F satisfies **IHJB2** and for any $\delta > 0$, there exists a finite set $\{\alpha_i\}_{i=1}^{M_\delta}$ such that for any $\alpha \in \mathcal{A}$:*

$$\inf_{1 \leq i \leq M_\delta} \left\{ |\sigma^\alpha - \sigma^{\alpha_i}|_\infty + |b^\alpha - b^{\alpha_i}|_\infty + |c^\alpha - c^{\alpha_i}|_\infty + |k^\alpha - k^{\alpha_i}|_\infty + \int_{\mathbb{R}_*^d} |(\eta^\alpha - \eta^{\alpha_i})(\cdot, z)|_\infty^2 d\nu(z) \right\} \leq \delta.$$

Remark 4.6. The Assumption **IHJB+** is satisfied if \mathcal{A} is a compact separable topological space and $\sigma^\alpha(\cdot)$, $b^\alpha(\cdot)$, and $c^\alpha(\cdot)$ are continuous maps from \mathcal{A} to $C_b^{\frac{1}{2}, 1}([0, T] \times \mathbb{R}^d)$; the space of bounded maps which are Lipschitz on x and $\frac{1}{2}$ -Hölder on t and $\eta^\alpha(\cdot)$ is continuous maps from \mathcal{A} to $\left\{ \varphi : [0, T] \times \mathbb{R}^d \times \mathbb{R}_*^d \rightarrow \mathbb{R} \mid \int_{\mathbb{R}_*^d} |\varphi(\cdot, z)|_\infty^2 \nu(dz) < \infty \right\}$.

Theorem 4.2 (Rate of Convergence). *Assume that the final condition g is bounded and Lipschitz-continuous. Then, there is a constant $C > 0$ such that*

- under Assumption IHJB,
 $v - v^{\kappa,h} \leq C \left(h^{\frac{1}{4}} + h\theta_\kappa^2 + h\varepsilon^{-3} + h^{\frac{3}{4}}\theta_\kappa + h\sqrt{\theta_\kappa} + h^{-\frac{1}{4}} \int_{\{|z| \leq \kappa\}} |z|^2 \nu(dz) \right)$.
- under Assumption IHJB+,
 $-C \left(h^{1/10} + h^{\frac{7}{10}}\theta_\kappa + h\sqrt{\theta_\kappa} + h^{-\frac{3}{10}} \int_{\{|z| \leq \kappa\}} |z|^2 \nu(dz) \right) \leq v - v^{\kappa,h}$.

In addition, if it is possible to find κ_h such that

$$\lim_{h \rightarrow 0} \kappa_h = 0, \quad \limsup_{h \rightarrow 0} h^{\frac{3}{4}} \theta_{\kappa_h}^2 < \infty \quad \text{and} \quad \limsup_{h \rightarrow 0} h^{-\frac{1}{2}} \int_{0 < |z| < \kappa_h} |z|^2 \nu(dz) < \infty, \quad (4.14)$$

then, there is a constant $C > 0$ such that

- under Assumption IHJB, $v - v^{\kappa_h,h} \leq Ch^{1/4}$.
- under Assumption IHJB+, $-Ch^{1/10} \leq v - v^{\kappa_h,h}$.

Next example shows the case where the conditions of the above Theorem on the choice of κ is satisfied. If it is not the case in some situations, it does mean that the rate of convergence is less than what is proposed by Theorem 4.2.

Example 4.1. For the Lévy measure

$$\nu(dz) = \mathbf{1}_{\mathbb{R}_*^d} |z|^{-d-1} dz,$$

one can always find κ_h such that the condition of Theorem 4.2 is satisfied. In the other words, it is always enough to choose κ_h such that

$$\limsup_{h \rightarrow 0} h^{-\frac{1}{2}} \kappa_h < \infty.$$

4.1. Convergence. We suppose the all the assumptions of Theorem 4.1 holds true throughout this subsection.

We first manipulate the scheme to provide strict monotonicity by the similar idea as in Remark 3.13 and Lemma 3.19 in [10]. Let $u^{\kappa,h}$ be the solution of

$$u^{\kappa,h}(T, \cdot) = g \quad \text{and} \quad u^{\kappa,h}(t_i, x) = \bar{\mathbf{T}}_{\kappa,h}[u^{\kappa,h}](t_i, x), \quad (4.15)$$

where

$$\bar{\mathbf{T}}_{\kappa,h}[\psi](t, x) := \mathbb{E} \left[\psi \left(t + h, \hat{X}_h^{t,x,\kappa} \right) \right] + h \bar{F}_{\kappa,h}(t, x, \mathcal{D}_h \psi, \psi(t+h, \cdot)) \quad (4.16)$$

and

$$\begin{aligned} \bar{F}_{\kappa,h}(t, x, r, p, \gamma, \psi) = & \sup_{\alpha} \inf_{\beta} \left\{ \frac{1}{2} a^{\alpha,\beta} \cdot \gamma + b^{\alpha,\beta} \cdot p + (c^{\alpha,\beta} + \theta_\kappa) r + e^{\theta_\kappa(T-t)} k^{\alpha,\beta}(t, x) \right. \\ & \left. + \int_{\{|z| \geq \kappa\}} (\hat{\nu}_h^{\eta^{\alpha,\beta},1}(\psi) - r - \mathbf{1}_{\{|z| \leq 1\}} \eta^{\alpha,\beta}(z) \cdot p) \nu(dz) \right\}. \end{aligned}$$

Remark 4.7. Assumption **Inf-Sup** is also true if we replace $\mathcal{J}_\kappa^{\alpha,\beta}$ by

$$\begin{aligned} \bar{\mathcal{J}}_\kappa^{\alpha,\beta}[\psi](t, x) = & \frac{1}{2} a^{\alpha,\beta} \cdot D^2 \phi(t, x) + b^{\alpha,\beta} \cdot D \phi(t, x) + (c^{\alpha,\beta} + \theta_\kappa) \phi(t, x) + e^{\theta_\kappa(T-t)} k^{\alpha,\beta}(t, x) \\ & + \int_{\{|z| \geq \kappa\}} (\hat{\nu}_h^{\eta^{\alpha,\beta},1}(\phi(t, \cdot))(x) - \phi(t, x) - \eta^{\alpha,\beta}(t, x, z) \cdot D \phi(t, x)) \nu(dz). \end{aligned}$$

The proof is straight forward.

We have the following Lemma which shows that for proper choice of θ_κ the scheme (4.15) is strictly monotone.

Lemma 4.1. *Let θ_κ be as in (4.8) and φ and $\psi : [0, T] \times \mathbb{R}^d \rightarrow \mathbb{R}$ be two bounded functions. Then:*

$$\varphi \leq \psi \implies \bar{\mathbf{T}}_{\kappa, h}[\varphi] \leq \bar{\mathbf{T}}_{\kappa, h}[\psi].$$

Proof. Let $f := \psi - \varphi \geq 0$ where φ and ψ are as in the statement of the lemma. For simplicity, we drop the dependence on (t, x) when it is not necessary. By Assumption **IHJB1** and Lemma (3.1), we can write:

$$\begin{aligned} \bar{\mathbf{T}}_{\kappa, h}[\psi] - \bar{\mathbf{T}}_{\kappa, h}[\varphi] &= \mathbb{E}[f(t+h, \hat{X}_h)] \\ &\quad + h \left(\inf_{\alpha} \sup_{\beta} \bar{\mathcal{J}}_{\kappa}^{\alpha, \beta}[\hat{\psi}](t+h, x) - \inf_{\alpha} \sup_{\beta} \bar{\mathcal{J}}_{\kappa}^{\alpha, \beta}[\hat{\varphi}](t+h, x) \right), \end{aligned}$$

where $\hat{\phi}(t, x) := \mathbb{E}[\phi(t, \hat{X}_h^x)]$ for $\phi = \varphi$ or ψ . Therefore,

$$\bar{\mathbf{T}}_{\kappa, h}[\psi] - \bar{\mathbf{T}}_{\kappa, h}[\varphi] \geq \mathbb{E}[f(t+h, \hat{X}_h)] + h \bar{\Phi}_{\kappa}^{\alpha, \beta}[\hat{\psi}, \hat{\varphi}](t+h, x, x),$$

where $\bar{\Phi}_{\kappa}^{\alpha, \beta}$ is defined by

$$\bar{\Phi}_{\kappa}^{\alpha, \beta}[\psi, \varphi](t, x) := \inf_{\alpha} \bar{\mathcal{J}}_{\kappa}^{\alpha, \beta}[\psi](t, x) - \sup_{\beta} \bar{\mathcal{J}}_{\kappa}^{\alpha, \beta}[\varphi](t, x').$$

By Assumption **Inf-Sup**, there exists (α^*, β^*) so that

$$\bar{\mathbf{T}}_{\kappa, h}[\psi] - \bar{\mathbf{T}}_{\kappa, h}[\varphi] \geq \mathbb{E}[f(t+h, \hat{X}_h)] + h \left(\bar{\mathcal{J}}_{\kappa}^{\alpha^*}[\hat{\psi}](t+h, x) - \bar{\mathcal{J}}_{\kappa}^{\beta^*}[\hat{\varphi}](t+h, x) \right).$$

Observe that by the linearity of $\bar{\mathcal{J}}_{\kappa}^{\alpha, \beta}$, one can write:

$$\bar{\mathcal{J}}_{\kappa}^{\alpha, \beta}[\hat{\phi}](t+h, x) = \mathbb{E} \left[\bar{\mathcal{J}}_{\kappa}^{\alpha, \beta}[\phi](t+h, \hat{X}_h) \right].$$

By the definition of $\bar{\mathcal{J}}_{\kappa}^{\alpha, \beta}$ and Lemma 2.1 in [10],

$$\begin{aligned} \bar{\mathbf{T}}_{\kappa, h}[\psi] - \bar{\mathbf{T}}_{\kappa, h}[\varphi] &\geq \mathbb{E} \left[f(\hat{X}_h) \left(1 + h \left(c_{\kappa}^{\alpha^*, \beta^*} + \theta_{\kappa} + b_{\kappa}^{\alpha^*, \beta^*} \cdot (\sigma^{\mathbf{T}})^{-1} \frac{W_h}{h} \right. \right. \right. \\ &\quad \left. \left. \left. + \frac{1}{2} a^{\alpha^*, \beta^*} \cdot (\sigma^{\mathbf{T}})^{-1} \frac{W_h W_h^{\mathbf{T}} - h \mathbf{I}_d}{h^2} \sigma^{-1} \right) \right) + h \hat{\nu}_h^{\eta^{\alpha^*, \beta^*}, 1}(f) \right], \end{aligned}$$

where $b_{\kappa}^{\alpha, \beta} = b^{\alpha, \beta} - \int_{\{|1 > |z| \geq \kappa\}} \eta^{\alpha, \beta}(z) \nu(dz)$ and $c_{\kappa}^{\alpha, \beta} = c^{\alpha, \beta} - \int_{\{|z| \geq \kappa\}} \nu(dz)$.

Therefore, by the same argument as in Lemma 3.12 in [10], one can write:

$$\begin{aligned} \bar{\mathbf{T}}_{\kappa, h}[\psi] - \bar{\mathbf{T}}_{\kappa, h}[\varphi] &\geq \mathbb{E} \left[f(\hat{X}_h) \left(1 - \frac{1}{2} a^{\alpha^*, \beta^*} \cdot a^{-1} + h \left(|A_h^{\alpha^*, \beta^*}|^2 + c_{\kappa}^* + \theta_{\kappa} \right. \right. \right. \\ &\quad \left. \left. \left. - \frac{1}{4} (b_{\kappa}^{\alpha^*, \beta^*})^{\mathbf{T}} (a^{\alpha^*, \beta^*})^{-1} b_{\kappa}^{\alpha^*, \beta^*} \right) \right) + h \hat{\nu}_h^{\eta^{\alpha^*, \beta^*}, 1}(f) \right], \end{aligned}$$

where

$$A_h^{\alpha^*, \beta^*} := \frac{1}{h} (\sigma^{\alpha^*, \beta^*})^{1/2} (\sigma^{\mathbf{T}})^{-1} W_h + \frac{1}{2} ((\sigma^{\alpha^*, \beta^*})^{-})^{1/2} b_{\kappa}^{\alpha^*, \beta^*}. \quad (4.17)$$

Therefore, by positivity of f and Assumption **IHJB1**, one can deduce:

$$\overline{\mathbf{T}}_{\kappa,h}[\psi] - \overline{\mathbf{T}}_{\kappa,h}[\varphi] \geq h\mathbb{E}\left[f(\hat{X}_h)\left(c_\kappa^* + \theta_\kappa - \frac{1}{4}(b_\kappa^{\alpha^*,\beta^*})^\top(a^{\alpha^*,\beta^*}) - b_\kappa^{\alpha^*,\beta^*}\right)\right]$$

By the choice of θ_κ in (4.8), we have

$$\overline{\mathbf{T}}_{\kappa,h}[\psi] - \overline{\mathbf{T}}_{\kappa,h}[\varphi] \geq 0.$$

Then, sending ε to zero provides the result. \square

The following Corollary shows the monotonicity of scheme 4.3.

Corollary 4.1. *Let $\varphi, \psi : [0, T] \times \mathbb{R}^d \rightarrow \mathbb{R}$ be two bounded functions. Then:*

$$\varphi \leq \psi \implies \mathbf{T}_{\kappa,h}[\varphi] \leq \mathbf{T}_{\kappa,h}[\psi] - \frac{\theta_\kappa^2 h^2}{2} e^{-\theta_\kappa h} \mathbb{E}[(\psi - \varphi)(t + h, \hat{X}_h^{t,x,\kappa})].$$

In particular, if κ_h satisfies (4.7), then

$$\varphi \leq \psi \implies \mathbf{T}_{\kappa_h,h}[\varphi] \leq \mathbf{T}_{\kappa_h,h}[\psi] + Ch\mathbb{E}[(\psi - \varphi)(t + h, \hat{X}_h^{t,x,\kappa_h})]$$

for some constant C .

Proof. Let θ_κ be as in Lemma 4.1 and define $\varphi_\kappa(t, x) := e^{\theta_\kappa(T-t)}\varphi(t, x)$ and $\psi_\kappa(t, x) := e^{\theta_\kappa(T-t)}\psi(t, x)$. By Lemma 4.1,

$$\overline{\mathbf{T}}_{\kappa,h}[\varphi_\kappa] \leq \overline{\mathbf{T}}_{\kappa,h}[\psi_\kappa].$$

By multiplying both sides by $e^{-\theta_\kappa(T-t)}$, we have

$$\begin{aligned} & \left(e^{-\theta_\kappa h}(1 + \theta_\kappa h) - 1\right)\mathbb{E}[\varphi(t + h, \hat{X}_h^{t,x,\kappa})] + \mathbf{T}_{\kappa,h}[\varphi] \\ & \leq \left(e^{-\theta_\kappa h}(1 + \theta_\kappa h) - 1\right)\mathbb{E}[\psi(t + h, \hat{X}_h^{t,x,\kappa})] + \mathbf{T}_{\kappa,h}[\psi]. \end{aligned}$$

So,

$$\mathbf{T}_{\kappa,h}[\varphi] \leq \left(e^{-\theta_\kappa h}(1 + \theta_\kappa h) - 1\right)\mathbb{E}[(\psi - \varphi)(t + h, \hat{X}_h^{t,x,\kappa})] + \mathbf{T}_{\kappa,h}[\psi].$$

But, $e^{-\theta_\kappa h}(1 + \theta_\kappa h) - 1 \leq -\frac{\theta_\kappa^2 h^2}{2} e^{-\theta_\kappa h}$. So,

$$\mathbf{T}_{\kappa,h}[\varphi] \leq -\frac{\theta_\kappa^2 h^2}{2} e^{-\theta_\kappa h} \mathbb{E}[(\psi - \varphi)(t + h, \hat{X}_h^{t,x,\kappa})] + \mathbf{T}_{\kappa,h}[\psi].$$

which (4.7) provides the result. \square

In order to provide a uniform bound on $v^{\kappa,h}$, we bound $u^{\kappa,h}$ with respect to θ_κ as in the following Lemma.

Lemma 4.2. *Let φ and $\psi : [0, T] \times \mathbb{R}^d \rightarrow \mathbb{R}$ be two L^∞ -bounded functions. Then*

$$|\overline{\mathbf{T}}_{\kappa,h}[\varphi] - \overline{\mathbf{T}}_{\kappa,h}[\psi]|_\infty \leq |\varphi - \psi|_\infty (1 + (C + \theta_\kappa)h)$$

where $C = \sup_{\alpha, \beta} |c^{\alpha, \beta}|_{\infty}$. In particular, if g is L^{∞} -bounded, for a fixed κ the family $(u^{\kappa, h}(t, \cdot))_h$ defined in (4.3) is L^{∞} -bounded, uniformly in h by

$$(\overline{C} + |g|_{\infty})e^{(C+\theta_{\kappa})(T-t_i)}.$$

Proof. Let $f := \varphi - \psi$. Then, by Assumption **Inf-Sup** and the same argument as in the proof of Lemma 4.1,

$$\begin{aligned} \overline{\mathbf{T}}_{\kappa, h}[\varphi] - \overline{\mathbf{T}}_{\kappa, h}[\psi] &\leq \mathbb{E} \left[f(\hat{X}_h) \left(1 - a^{-1} \cdot a^{\alpha^*, \beta^*} + h \left(|A_h^{\alpha^*, \beta^*}|^2 + c^{\alpha^*, \beta^*} + \theta_{\kappa} \right. \right. \right. \\ &\quad \left. \left. - \int_{\{|z| \geq \kappa\}} \nu(dz) - \frac{1}{4} \left(b^{\alpha^*, \beta^*} - \int_{\{1 > |z| \geq \kappa\}} \eta^{\alpha^*, \beta^*}(z) \nu(dz) \right)^{\mathbf{T}} (a^{\alpha^*, \beta^*})^{-1} \right. \right. \\ &\quad \left. \left. \times \left(b^{\alpha^*, \beta^*} - \int_{\{1 > |z| \geq \kappa\}} \eta^{\alpha^*, \beta^*}(z) \nu(dz) \right) \right) \right) + h \hat{\nu}_h^{\eta^{\alpha^*, \beta^*}, 1}(f) \right], \end{aligned}$$

where $A_h^{\alpha^*, \beta^*}$ is given by (4.17). On the other hand,

$$\left| \hat{\nu}_h^{\eta^{\alpha^*, \beta^*}, 1}(f) \right| \leq |f|_{\infty} \int_{\{|z| \geq \kappa\}} \nu(dz)$$

Therefore ,

$$\begin{aligned} \overline{\mathbf{T}}_{\kappa, h}[\varphi] - \overline{\mathbf{T}}_{\kappa, h}[\psi] &\leq |f|_{\infty} \mathbb{E} \left[\left| 1 - a^{-1} \cdot a^{\alpha^*, \beta^*} + h \left(|A_h^{\alpha^*, \beta^*}|^2 + c^{\alpha^*, \beta^*} + \theta_{\kappa} \right. \right. \right. \\ &\quad \left. \left. - \frac{1}{4} \left(b^{\alpha^*, \beta^*} - \int_{\{1 > |z| \geq \kappa\}} \eta^{\alpha^*, \beta^*}(z) \nu(dz) \right)^{\mathbf{T}} (a^{\alpha^*, \beta^*})^{-1} \left(b^{\alpha^*, \beta^*} - \int_{\{1 > |z| \geq \kappa\}} \eta^{\alpha^*, \beta^*}(z) \nu(dz) \right) \right) \right] \right|. \end{aligned}$$

By Assumption **IHJB1** and (4.8), $1 - a^{-1} \cdot a^{\alpha^*, \beta^*}$ and

$$c^{\alpha^*, \beta^*} + \theta_{\kappa} - \frac{1}{4} \left(b^{\alpha^*, \beta^*} - \int_{\{1 > |z| \geq \kappa\}} \eta^{\alpha^*, \beta^*}(z) \nu(dz) \right)^{\mathbf{T}} (a^{\alpha^*, \beta^*})^{-1} \left(b^{\alpha^*, \beta^*} - \int_{\{1 > |z| \geq \kappa\}} \eta^{\alpha^*, \beta^*}(z) \nu(dz) \right)$$

are positive. Therefore, one can write

$$\begin{aligned} \overline{\mathbf{T}}_{\kappa, h}[\varphi] - \overline{\mathbf{T}}_{\kappa, h}[\psi] &\leq |f|_{\infty} \left(1 - a^{-1} \cdot a^{\alpha^*, \beta^*} + h \left(\mathbb{E}[|A_h^{\alpha^*, \beta^*}|^2] + c^{\alpha^*, \beta^*} + \theta_{\kappa} \right. \right. \\ &\quad \left. \left. - \frac{1}{4} \left(b^{\alpha^*, \beta^*} - \int_{\{1 > |z| \geq \kappa\}} \eta^{\alpha^*, \beta^*}(z) \nu(dz) \right)^{\mathbf{T}} (a^{\alpha^*, \beta^*})^{-1} \left(b^{\alpha^*, \beta^*} - \int_{\{1 > |z| \geq \kappa\}} \eta^{\alpha^*, \beta^*}(z) \nu(dz) \right) \right) \right). \end{aligned} \quad (4.18)$$

But, Notice that

$$\begin{aligned} \mathbb{E}[|A_h^{\alpha^*, \beta^*}|^2] &= h^{-1} a^{-1} \cdot a^{\alpha^*, \beta^*} \\ &\quad + \frac{1}{4} \left(b^{\alpha^*, \beta^*} - \int_{\{1 > |z| \geq \kappa\}} \eta^{\alpha^*, \beta^*}(z) \nu(dz) \right)^{\mathbf{T}} a^{\alpha^*, \beta^* - 1} \left(b^{\alpha^*, \beta^*} - \int_{\{1 > |z| \geq \kappa\}} \eta^{\alpha^*, \beta^*}(z) \nu(dz) \right). \end{aligned}$$

By replacing $\mathbb{E}[|A_h^{\alpha^*, \beta^*}|^2]$ into (4.18), one obtains

$$\begin{aligned} \overline{\mathbf{T}}_{\kappa, h}[\varphi] - \overline{\mathbf{T}}_{\kappa, h}[\psi] &\leq |f|_{\infty} (1 + h(c^{\alpha^*, \beta^*} + \theta_{\kappa})) \\ &\leq |f|_{\infty} (1 + (C + \theta_{\kappa})h), \end{aligned}$$

with $C = \sup_{\alpha, \beta} |c^{\alpha, \beta}|_\infty$. By changing the role of φ and ψ and implementing the same argument, one obtains

$$|\overline{\mathbf{T}}_{\kappa, h}[\varphi] - \overline{\mathbf{T}}_{\kappa, h}[\psi]|_\infty \leq |f|_\infty (1 + (C + \theta_\kappa)h).$$

To prove that the family $(u^{\kappa, h})_h$ is bounded, we proceed by backward induction as in Lemma 3.14 in [10]. By choosing in the first part of the proof $\varphi \equiv \bar{u}^{\kappa, h}(t_{i+1}, \cdot)$ and $\psi \equiv 0$, we see that

$$|u^{\kappa, h}(t_i, \cdot)|_\infty \leq h\overline{C}e^{\theta_\kappa(T-t_i)} + |u^{\kappa, h}(t_{i+1}, \cdot)|_\infty (1 + (C + \theta_\kappa)h),$$

where $\overline{C} := \sup_{\alpha, \beta} |k^{\alpha, \beta}|_\infty$. It follows from the discrete Gronwall inequality that

$$|u^{\kappa, h}(t_i, \cdot)|_\infty \leq (\overline{C}(T - t_i) + |g|_\infty)e^{(C + \theta_\kappa)(T-t_i)}.$$

□

Define

$$\bar{v}^{\kappa, h} := e^{-\theta_\kappa(T-t)} u^{\kappa, h}. \quad (4.19)$$

Next Corollary provides a bound for $v^{\kappa, h}$ uniformly on κ and h .

Corollary 4.2. *$\bar{v}^{\kappa, h}$ is bounded uniformly on h and κ , and*

$$|v^{\kappa, h} - \bar{v}^{\kappa, h}|_\infty \leq K\theta_\kappa^2 h \text{ for some constant } K.$$

If also, κ_h satisfies (4.7), then

$$\lim_{h \rightarrow 0} |v^{\kappa_h, h} - \bar{v}^{\kappa_h, h}|_\infty = 0.$$

Proof. By Lemma 4.2 for fixed κ , we have:

$$|u^{\kappa, h}(t, \cdot)|_\infty \leq (\overline{C} + |g|_\infty)e^{(C + \theta_\kappa)(T-t)}.$$

Therefore,

$$|\bar{v}^{\kappa, h}(t, \cdot)|_\infty \leq (\overline{C} + |g|_\infty)e^{C(T-t)}.$$

For the next part, define $\bar{u}^{\kappa, h}(t, x) = e^{\theta_\kappa(T-t)} v^{\kappa, h}(t, x)$. Direct calculations shows that

$$\bar{u}^{\kappa, h} = e^{\theta_\kappa h} (1 - \theta_\kappa h) \mathbb{E} \left[\bar{u}^{\kappa, h} \left(t + h, \hat{X}_h^{t, x, \kappa} \right) \right] + h \overline{F}_{\kappa, h} \left(t, x, \mathcal{D}_h \bar{u}^{\kappa, h}, \bar{u}^{\kappa, h}(t + h, \cdot) \right).$$

By an argument similar to Lemma 3.19 in [10], we have

$$\begin{aligned} |(u^{\kappa, h} - \bar{u}^{\kappa, h})(t, \cdot)|_\infty &\leq \frac{1}{2} \theta_\kappa^2 h^2 |\bar{u}^{\kappa, h}(t + h, \cdot)|_\infty \\ &\quad + (1 + (C + \theta_\kappa)h) |(u^{\kappa, h} - \bar{u}^{\kappa, h})(t + h, \cdot)|_\infty, \end{aligned} \quad (4.20)$$

where C is as in Lemma 4.2. By repeating the proof of Lemma 4.2 for $\bar{u}^{\kappa, h}$, one can conclude,

$$|\bar{u}^{\kappa, h}(t, \cdot)|_\infty \leq (\overline{C} + |g|_\infty)e^{(C + \theta_\kappa)(T-t)} \left(1 + \frac{\theta_\kappa h}{2} \right).$$

So, by multiplying 4.20 by $e^{\theta_\kappa(T-t)}$, we have

$$\begin{aligned} |(\bar{v}^{\kappa,h} - v^{\kappa,h})(t, \cdot)|_\infty &\leq \frac{1}{2} \tilde{C} \theta_\kappa^2 h^2 e^{C(T-t)} \left(1 + \frac{\theta_\kappa h}{2}\right) e^{-\theta_\kappa h} \\ &\quad + e^{-\theta_\kappa h} (1 + (C + \theta_\kappa)h) |(\bar{v}^{\kappa,h} - v^{\kappa,h})(t+h, \cdot)|_\infty, \end{aligned}$$

for some constant \tilde{C} . Because $e^{-\theta_\kappa h} (1 + (C + \theta_\kappa)h) \leq e^{Ch}$, one can deduce from discrete Gronwall inequality that

$$|(\bar{v}^{\kappa,h} - v^{\kappa,h})(t, \cdot)|_\infty \leq K \theta_\kappa^2 h,$$

for some constant K independent of κ which provides the second part of the theorem. \square

We continue with the following consistency Lemma.

Lemma 4.3. *Let φ be a smooth function with the bounded derivatives. Then for all $(t, x) \in [0, T] \times \mathbb{R}^d$:*

$$\lim_{\substack{(t', x') \rightarrow (t, x) \\ (h, c) \rightarrow (0, 0) \\ t' + h \leq T}} \frac{\varphi(t', x') - \mathbf{T}_{\kappa, h}[c + \varphi](t', x')}{h} = -(\mathcal{L}^X \varphi + F(\cdot, \varphi, D\varphi, D^2\varphi, \varphi(t, \cdot))) (t, x).$$

Proof. The proof is straightforward by Lebesgue dominated convergence Theorem. \square

To complete the convergence argument, we need to prove the approximate solution $v^{\kappa_h, h}$ converge to the final condition as

Lemma 4.4. *Let κ_h satisfy (4.7), then $\bar{v}^{\kappa_h, h}$ is uniformly Lipschitz with respect to x .*

Proof. We report the following calculation in the one-dimensional case $d = 1$ in order to simplify the presentation.

For fixed $t \in [0, T - h]$, we argue as in the proof of Lemma 4.2 to see that for $x, x' \in \mathbb{R}$ with $x > x'$:

$$\begin{aligned} u^{\kappa, h}(t, x) - u^{\kappa, h}(t, x') &= \mathbb{E} \left[\left(u^{\kappa, h}(t+h, \hat{X}^{t, x}) - u^{\kappa, h}(t+h, \hat{X}^{t, x'}) \right) \right] \\ &\quad + h \left(\inf_\alpha \sup_\beta \bar{\mathcal{J}}_\kappa^{\alpha, \beta}[\widehat{u^{\kappa, h}}](t+h, x) - \inf_\alpha \sup_\beta \bar{\mathcal{J}}_\kappa^{\alpha, \beta}[\widehat{u^{\kappa, h}}](t+h, x') \right) \\ &\leq \mathbb{E} \left[\left(u^{\kappa, h}(t+h, \hat{X}^{t, x}) - u^{\kappa, h}(t+h, \hat{X}^{t, x'}) \right) \right] \\ &\quad + h \left(\sup_\beta \bar{\mathcal{J}}_\kappa^{\alpha, \beta}[\widehat{u^{\kappa, h}}](t+h, x) - \inf_\alpha \bar{\mathcal{J}}_\kappa^{\alpha, \beta}[\widehat{u^{\kappa, h}}](t+h, x') \right). \end{aligned}$$

Observe that by (4.6), one can write

$$\begin{aligned} u^{\kappa, h}(t, x) - u^{\kappa, h}(t, x') &\leq \mathbb{E} \left[\left(u^{\kappa, h}(t+h, \hat{X}^{t, x}) - u^{\kappa, h}(t+h, \hat{X}^{t, x'}) \right) \right] \\ &\quad + h \left(\bar{\Phi}^{\alpha, \beta}[\widehat{u^{\kappa, h}}, \widehat{u^{\kappa, h}}](t+h, x, x') \right), \end{aligned}$$

where $\bar{\Phi}$ is defined in the proof of Lemma 4.2. By Assumption **Inf-Sup**, there exists (α^*, β^*) such that

$$\bar{\Phi}^{\alpha^*, \beta^*}[\widehat{u^{\kappa, h}}, \widehat{u^{\kappa, h}}](t+h, x, x') = \bar{\mathcal{J}}_{\kappa}^{\alpha^*, \beta^*}[\widehat{u^{\kappa, h}}](t+h, x) - \bar{\mathcal{J}}_{\kappa}^{\alpha^*, \beta^*}[\widehat{u^{\kappa, h}}](t+h, x').$$

Therefore,

$$\begin{aligned} u^{\kappa, h}(t, x) - u^{\kappa, h}(t, x') &\leq \mathbb{E}\left[\left(u^{\kappa, h}(t+h, \hat{X}^{t, x}) - u^{\kappa, h}(t+h, \hat{X}^{t, x'})\right)\right] \\ &\quad + h\left(\bar{\mathcal{J}}_{\kappa}^{\alpha^*, \beta^*}[\widehat{u^{\kappa, h}}](t+h, x) - \bar{\mathcal{J}}_{\kappa}^{\alpha^*, \beta^*}[\widehat{u^{\kappa, h}}](t+h, x')\right). \end{aligned}$$

For the other in equality we do the same except that when we

$$u^{\kappa, h}(t, x) - u^{\kappa, h}(t, x') \leq A + hB + hC,$$

where

$$\begin{aligned} A &:= \mathbb{E}\left[\left(u^{\kappa, h}(t+h, \hat{X}^{t, x}) - u^{\kappa, h}(t+h, \hat{X}^{t, x'})\right)\right] \\ &\quad + h\left(\bar{\mathcal{J}}_{\kappa}^{\alpha^*, \beta^*}[\widehat{u^{\kappa, h}}](t+h, x) - \bar{\mathcal{J}}_{\kappa}^{\alpha^*, \beta^*}[\widehat{u^{\kappa, h}}](t+h, x')\right), \end{aligned}$$

with $\tilde{u}^{\kappa, h}(y) = u^{\kappa, h}(y + x' - x)$,

$$B := \bar{\mathcal{J}}_{\kappa}^{\alpha^*, \beta^*}[\widehat{\tilde{u}^{\kappa, h}}](t+h, x) - \bar{\mathcal{J}}_{\kappa}^{\alpha^*, \beta^*}[\widehat{u^{\kappa, h}}](t+h, x'),$$

and

$$C := \hat{\nu}_h^{\alpha^*, \beta^*, 1}(u^{\kappa, h}(t+h, \cdot))(x) - \hat{\nu}_h^{\alpha^*, \beta^*, 1}(u^{\kappa, h}(t+h, \cdot))(x').$$

We continue the proof in the following steps.

Step 1.

$$C = h^{-1} \mathbb{E}\left[\left(u^{\kappa, h}(t+h, \hat{X}^{*, x}) - u^{\kappa, h}(t+h, \hat{X}^{*, x'})\right) N_h^{\kappa}\right],$$

where $\hat{X}^{*, x} := x + \sum_{i=1}^{N_h^{\kappa}} \eta^{\alpha^*, \beta^*}(x, Z_i)$ with Z_i s are i.i.d. random variables distributed as $\frac{\nu(dz)}{\lambda_{\kappa}}$.

Step 2. By the definition of $\bar{\mathcal{J}}_{\kappa}^{\alpha, \beta}$,

$$\begin{aligned} B &= \frac{1}{2}(a^{\alpha^*, \beta^*}(x) - a^{\alpha^*, \beta^*}(x')) \mathcal{D}_h^2 u^{\kappa, h}(t+h, x') + (b_{\kappa}^{\alpha^*, \beta^*}(x) - b_{\kappa}^{\alpha^*, \beta^*}(x')) \mathcal{D}_h^1 u^{\kappa, h}(t+h, x') \\ &\quad + (c^{\alpha^*, \beta^*}(x) - c^{\alpha^*, \beta^*}(x')) \mathcal{D}_h^0 u^{\kappa, h}(t+h, x') + k^{\alpha^*, \beta^*}(x) - k^{\alpha^*, \beta^*}(x'), \end{aligned}$$

where $b_{\kappa}^{\alpha, \beta}(x) := b^{\alpha, \beta}(x) - \int_{\{1 > |z| \geq \kappa\}} \eta^{\alpha, \beta}(x, z) \nu(dz)$. On the other hand,

$$\mathcal{D}_h^k = \mathbb{E}\left[Du^{\kappa, h}(t+h, \hat{X}_h^{x'}) \left(\frac{W_h}{h} \sigma^{-1}(x')\right)^{k-1}\right], \text{ for } k = 1, 2.$$

So,

$$\begin{aligned}
B \leq & \mathbb{E} \left[\frac{1}{2} (a^{\alpha^*, \beta^*}(x) - a^{\alpha^*, \beta^*}(x')) Du^{\kappa, h}(t+h, \hat{X}_h^{x'}) \frac{W_h}{h} \sigma^{-1}(x') \right. \\
& + (b_{\kappa}^{\alpha^*, \beta^*}(x) - b_{\kappa}^{\alpha^*, \beta^*}(x')) Du^{\kappa, h}(t+h, \hat{X}_h^{x'}) + (c^{\alpha^*, \beta^*}(x) - c^{\alpha^*, \beta^*}(x')) u^{\kappa, h}(t+h, \hat{X}_h^{x'}) \left. \right] \\
& + f^{\alpha^*, \beta^*}(x) - f^{\alpha^*, \beta^*}(x').
\end{aligned}$$

Step 3. By the definition of $\bar{\mathcal{J}}_{\kappa}^{\alpha, \beta}$, one can observe that

$$\begin{aligned}
\bar{\mathcal{J}}_{\kappa}^{\alpha^*, \beta^*} [u^{\kappa, h}](t+h, x) - \bar{\mathcal{J}}_{\kappa}^{\alpha^*, \beta^*} [\tilde{u}^{\kappa, h}](t+h, x) \\
= \frac{1}{2} a^{\alpha^*, \beta^*}(x) \delta^{(2)} + b_{\kappa}^*(x) \delta(1) + c_{\kappa}^*(x) \delta(0)
\end{aligned}$$

where c_{κ}^* and b_{κ}^* are defined in the proof of Lemma 4.1, and

$$\delta^{(k)} = \mathbb{E} \left[D^k u^{\kappa, h}(t+h, \hat{X}_h^x) - D^k u^{\kappa, h}(t+h, \hat{X}_h^{x'}) \right] \text{ for } k = 0, 1, 2.$$

By Lemma 2.1 in [10], for $k = 1$ and 2

$$\begin{aligned}
\delta^{(k)} &= \mathbb{E} \left[\left(u^{\kappa, h}(t+h, \hat{X}_h^x) - u^{\kappa, h}(t+h, \hat{X}_h^{x'}) \right) H_h^k(t, x) \right. \\
&\quad \left. + u^{\kappa, h}(t+h, \hat{X}_h^{x'}) H_h^k(t, x) \left(1 - \frac{\sigma^k(x)}{\sigma^k(x')} \right) \right] \\
&= \mathbb{E} \left[\left(u^{\kappa, h}(t+h, \hat{X}_h^x) - u^{\kappa, h}(t+h, \hat{X}_h^{x'}) \right) H_h^k(t, x) \right. \\
&\quad \left. + Du^{\kappa, h}(t+h, \hat{X}_h^{x'}) \left(\frac{W_h}{h} \right)^{k-1} \sigma(x') \left(\sigma^{-k}(x) - \sigma^{-k}(x') \right) \right].
\end{aligned}$$

Therefore, one can write

$$\begin{aligned}
A \leq & \mathbb{E} \left[\left(u^{\kappa, h}(t+h, \hat{X}_h^x) - u^{\kappa, h}(t+h, \hat{X}_h^{x'}) \right) \right. \\
&\quad \times \left(1 - \bar{a}^* + \bar{a}^* N^2 + h c_{\kappa}^* + b_{\kappa}^* N \sqrt{h} \right) (x) \\
&\quad + h b_{\kappa}^*(x') Du^{\kappa, h}(t+h, \hat{X}_h^{x'}) \sigma(x') \left(\sigma^{-1}(x) - \sigma^{-1}(x') \right) \\
&\quad \left. + a^*(x') Du^{\kappa, h}(t+h, \hat{X}_h^{x'}) \sqrt{h} N \sigma(x') \left(\sigma^{-2}(x) - \sigma^{-2}(x') \right) \right],
\end{aligned}$$

where $a^* := \frac{1}{2} a^{\alpha^*, \beta^*}$, $\bar{a}^* := \frac{1}{2} a^{-1} a^{\alpha^*, \beta^*}$, $c^* := c^{\alpha^*, \beta^*}$, $c_{\kappa}^* := c^* + \theta_{\kappa}$, and $b_{\kappa}^* := b_{\kappa}^{\alpha^*, \beta^*}$.

Step 4. By dividing both sides by $x - x'$ and taking the limit we have:

$$\begin{aligned}
Du^{\kappa, h}(t, x) &\leq \mathbb{E} \left[Du^{\kappa, h}(t+h, \hat{X}_h^x) \left(\left(1 + h \bar{\mu}'_{\kappa} + \sqrt{h} \sigma' N + \tilde{\mathcal{J}}_{\kappa, h} \right) \right. \right. \\
&\quad \times \left(1 - \bar{a}^* + \bar{a}^* N^2 + h c_{\kappa}^* + b_{\kappa}^* N \sqrt{h} \right) \\
&\quad \left. + h \left((b_{\kappa}^*)' - b_{\kappa}^* \frac{\sigma'}{\sigma} \right) + \left(\frac{1}{2} (a^{\alpha^*, \beta^*})' \sigma^{-1} - a^{\alpha^*, \beta^*} \frac{\sigma'}{\sigma^2} \right) \sqrt{h} N \right) \\
&\quad \left. + Du^{\kappa, h}(t+h, \hat{X}_h^{*, x}) \left(1 + \mu^* h + \tilde{\mathcal{J}}_{\kappa, h}^* \right) N_h^{\kappa} \right] + C e^{\theta_{\kappa}(T-t)} h,
\end{aligned}$$

where $\tilde{J}_{\kappa,h} := \int_{\{|z|>\kappa\}} \eta(z) \tilde{J}([0,h], dz)$, $\tilde{J}'_{\kappa,h} := \int_{\{|z|>\kappa\}} \eta'(z) \tilde{J}([0,h], dz)$, and N_h^κ is a Poisson process with intensity $\lambda_\kappa := \int_{\{|z|>\kappa\}} \nu(dz)$.

Let $L_t := |Du^{\kappa,h}(t, \cdot)|_\infty$. Then

$$\mathbb{E} \left[Du^{\kappa,h}(t+h, \hat{X}_h^{*,x}) \left(1 + \mu^* h + \tilde{J}'_{\kappa,h} \right) N_h^\kappa \right] \leq L_{t+h} Ch (\lambda_\kappa + \lambda_\kappa^*),$$

where $\lambda_\kappa^* := \int_{\{|z|>\kappa\}} \eta^*(z) \nu(dz)$. Let $G := N + \frac{b_\kappa^* \sigma}{2} \sqrt{h}$. By the change of measure

$$\frac{d\mathbb{Q}}{d\mathbb{P}} := \exp \left(-\frac{(b_\kappa^* \sigma)^2}{4} h + \frac{b_\kappa^* \sigma}{2} \sqrt{h} N \right),$$

we have $G \sim \mathcal{N}(0, 1)$ under \mathbb{Q} and one can write

$$\begin{aligned} Du^{\kappa,h}(t, x) &\leq \mathbb{E}^{\mathbb{Q}} \left[\frac{d\mathbb{P}}{d\mathbb{Q}} Du^{\kappa,h}(t+h, \hat{X}_h^x) \left(\left(1 + h(\tilde{\mu}'_\kappa - \frac{b_\kappa^* \sigma}{2}) + \sqrt{h} \sigma' G + \tilde{J}_{\kappa,h} \right) \right. \right. \\ &\quad \times \left(1 - \bar{a}^* + \bar{a}^* G^2 + h(c_\kappa^* - \frac{(b_\kappa^* \sigma)^2}{2}) \right) \\ &\quad \left. \left. + h \left((b_\kappa^*)' - b_\kappa^* \frac{\sigma'}{\sigma} - \frac{b_\kappa^* \sigma}{2} \right) + \left(\frac{1}{2} (a^{\alpha^*, \beta^*})' \sigma^{-1} - a^{\alpha^*, \beta^*} \frac{\sigma'}{\sigma^2} \right) \sqrt{h} G \right) \right] \\ &\quad + L_{t+h} Ch (\lambda_\kappa + \lambda_\kappa^*) + C e^{\theta_\kappa(T-t)} h, \end{aligned}$$

Step 4. Notice that $1 - \bar{a}^* + \bar{a}^* G^2 + h(c_\kappa^* - \frac{(b_\kappa^* \sigma)^2}{2})$ is positive and therefore, one can take $\frac{Z}{\mathbb{E}^{\mathbb{Q}}[Z]}$ as a density for the new measure \mathbb{Q}^Z . So,

$$\begin{aligned} Du^{\kappa,h}(t, x) &\leq \mathbb{E}^{\mathbb{Q}^Z} \left[\frac{d\mathbb{P}}{d\mathbb{Q}} Du^{\kappa,h}(t+h, \hat{X}_h^x) \left(\left(1 + h(\tilde{\mu}'_\kappa - \frac{b_\kappa^* \sigma}{2}) + \sqrt{h} \sigma' G + \tilde{J}_{\kappa,h} \right) \right. \right. \\ &\quad \left. \left. + Z^{-1} \left(h \left((b_\kappa^*)' - b_\kappa^* \frac{\sigma'}{\sigma} - \frac{b_\kappa^* \sigma}{2} \right) + \left(\frac{1}{2} (a^{\alpha^*, \beta^*})' \sigma^{-1} - a^{\alpha^*, \beta^*} \frac{\sigma'}{\sigma^2} \right) \sqrt{h} G \right) \right) \right] \\ &\quad + L_{t+h} Ch (\lambda_\kappa + \lambda_\kappa^*) + C e^{\theta_\kappa(T-t)} h. \end{aligned}$$

So,

$$\begin{aligned} Du^{\kappa,h}(t, x) &\leq \mathbb{E}^{\mathbb{Q}^Z} \left[\left(\frac{d\mathbb{P}}{d\mathbb{Q}} \right)^2 (Du^{\kappa,h}(t+h, \hat{X}_h^x))^2 \right]^{\frac{1}{2}} \mathbb{E}^{\mathbb{Q}^Z} \left[\left(\left(1 + h(\tilde{\mu}'_\kappa - \frac{b_\kappa^* \sigma}{2}) + \sqrt{h} \sigma' G + \tilde{J}_{\kappa,h} \right) \right. \right. \\ &\quad \left. \left. + Z^{-1} \left(h \left((b_\kappa^*)' - b_\kappa^* \frac{\sigma'}{\sigma} - \frac{b_\kappa^* \sigma}{2} \right) + \left(\frac{1}{2} (a^{\alpha^*, \beta^*})' \sigma^{-1} - a^{\alpha^*, \beta^*} \frac{\sigma'}{\sigma^2} \right) \sqrt{h} G \right) \right)^2 \right]^{\frac{1}{2}} \\ &\quad + L_{t+h} Ch (\lambda_\kappa + \lambda_\kappa^*) + C e^{\theta_\kappa(T-t)} h. \end{aligned}$$

Notice that

$$\mathbb{E}^{\mathbb{Q}^Z} \left[\left(\frac{d\mathbb{Q}}{d\mathbb{P}} \right)^2 (Du^{\kappa,h}(t+h, \hat{X}_h^x))^2 \right] \leq L_{t+h}^2 \exp \left(\frac{1}{4} (b_\kappa^* \sigma)^2 h \right).$$

On the other hand,

$$\begin{aligned} & \mathbb{E}^{\mathbb{Q}^Z} \left[\frac{d\mathbb{Q}}{d\mathbb{P}} \left(\left(1 + h(\tilde{\mu}'_{\kappa} - \frac{b_{\kappa}^* \sigma}{2}) + \sqrt{h} \sigma' G + \tilde{J}_{\kappa, h} \right) + Z^{-1} \left(h \left((b_{\kappa}^*)' - b_{\kappa}^* \frac{\sigma'}{\sigma} - \frac{b_{\kappa}^* \sigma}{2} \right) \right. \right. \right. \\ & \quad \left. \left. \left. + \left(\frac{1}{2} (a^{\alpha^*, \beta^*})' \sigma^{-1} - a^{\alpha^*, \beta^*} \frac{\sigma'}{\sigma^2} \right) \sqrt{h} G \right) \right)^2 \right] \\ &= \mathbb{E} \left[Z \left(\left(1 + h(\tilde{\mu}'_{\kappa} - \frac{b_{\kappa}^* \sigma}{2}) + \sqrt{h} \sigma' G + \tilde{J}_{\kappa, h} \right) + Z^{-1} \left(h \left((b_{\kappa}^*)' - b_{\kappa}^* \frac{\sigma'}{\sigma} - \frac{b_{\kappa}^* \sigma}{2} \right) \right. \right. \right. \\ & \quad \left. \left. \left. + \left(\frac{1}{2} (a^{\alpha^*, \beta^*})' \sigma^{-1} - a^{\alpha^*, \beta^*} \frac{\sigma'}{\sigma^2} \right) \sqrt{h} G \right) \right)^2 \right]. \end{aligned}$$

By calculation of the right hand side of the above equality, one can observe that all the terms of order \sqrt{h} vanish and we have:

$$\begin{aligned} & \mathbb{E}^{\mathbb{Q}^Z} \left[\frac{d\mathbb{Q}}{d\mathbb{P}} \left(\left(1 + h(\tilde{\mu}'_{\kappa} - \frac{b_{\kappa}^* \sigma}{2}) + \sqrt{h} \sigma' G + \tilde{J}_{\kappa, h} \right) + Z^{-1} \left(h \left((b_{\kappa}^*)' - b_{\kappa}^* \frac{\sigma'}{\sigma} - \frac{b_{\kappa}^* \sigma}{2} \right) \right. \right. \right. \\ & \quad \left. \left. \left. + \left(\frac{1}{2} (a^{\alpha^*, \beta^*})' \sigma^{-1} - a^{\alpha^*, \beta^*} \frac{\sigma'}{\sigma^2} \right) \sqrt{h} G \right) \right)^2 \right]^{\frac{1}{2}} \\ & \leq \left(1 + h \left(c^* + \theta_{\kappa} - \frac{(b_{\kappa}^*)^2}{4a^*} - b_{\kappa}^* \sigma \sigma' + (b_{\kappa}^*)' - \frac{b_{\kappa}^* \sigma'}{\sigma} - \frac{b_{\kappa}^* \sigma}{2} + O(h\theta_{\kappa}^2) \right) \right)^{\frac{1}{2}}. \end{aligned}$$

Therefore, by the choice of κ_h , for h small enough we have

$$\begin{aligned} L_t & \leq L_{t+h} \exp \left(\frac{1}{2} h \left(C + \theta_{\kappa_h} - b_{\kappa_h}^* \sigma \sigma' + (b_{\kappa_h}^*)' - \frac{b_{\kappa_h}^* \sigma'}{\sigma} - \frac{b_{\kappa_h}^* \sigma}{2} + 2\lambda_{\kappa_h} + 2\lambda^{*\kappa_h} \right) \right) + C e^{\theta_{\kappa_h} (T-t)} h \\ & \leq L_{t+h} \exp \left(h(C + \theta_{\kappa_h}) \right) + C e^{\theta_{\kappa_h} (T-t)} h \end{aligned}$$

By discrete Gronwall inequality,

$$L_t \leq (|Dg|_{\infty} + C(T-t)) e^{(\theta_{\kappa_h} + C)(T-t)}.$$

Therefore by definition of $\bar{v}^{\kappa, h}$, we have

$$|D\bar{v}^{\kappa, h}|_1 \leq e^{C(T-t)} (|Dg|_{\infty} + C(T-t)).$$

□

Lemma 4.5. *Let κ_h satisfies (4.7), then*

$$\lim_{t \rightarrow T} \bar{v}^{\kappa, h}(t, x) = g(x).$$

Proof. We follow the same notations as in the proof of the previous Lemma and write

$$\begin{aligned} u^{\kappa, h}(t, x) &= \mathbb{E} \left[u^{\kappa, h}(t+h, \hat{X}^{t, x}) \right] + h \inf_{\alpha} \sup_{\beta} \bar{\mathcal{J}}_{\kappa}^{\alpha, \beta} [\widehat{u^{\kappa, h}}](t+h, x) \\ &\leq \mathbb{E} \left[u^{\kappa, h}(t+h, \hat{X}^{t, x}) \right] + h \sup_{\beta} \bar{\mathcal{J}}_{\kappa}^{\alpha, \beta} [\widehat{u^{\kappa, h}}](t+h, x). \end{aligned}$$

Observe that by (4.6), one can write

$$u^{\kappa, h}(t, x) \leq \mathbb{E} \left[u^{\kappa, h}(t+h, \hat{X}^{t, x}) \right] + h \left(\bar{\Phi}^{\alpha, \beta} [\widehat{u^{\kappa, h}}, 0](t+h, x, x') \right) + h \sup_{\alpha, \beta} |f^{\alpha, \beta}|_{\infty},$$

By Assumption **Inf-Sup**, there exists (α^*, β^*) so that

$$u^{\kappa, h}(t, x) \leq \mathbb{E} \left[u^{\kappa, h}(t+h, \hat{X}^{t, x}) \right] + h \bar{\mathcal{J}}_{\kappa}^{\alpha^*, \beta^*} [\widehat{u^{\kappa, h}}](t+h, x) + h \bar{C},$$

where $\bar{C} := \sup_{\alpha, \beta} |f^{\alpha, \beta}|_{\infty}$. Therefore, for any $j = i, \dots, n-1$ one can write

$$u^{\kappa, h}(t_j, \hat{X}_{t_j}^{t_i, x}) \leq \mathbb{E}_{t_j}^{\mathbb{Q}} \left[u^{\kappa, h}(t_{j+1}, \hat{X}_{t_{j+1}}^{t_i, x}) \left(1 - a_j^* + a_j^* G_j^2 + h C_j^* \right) \right] + h \bar{C}.$$

where $a_j^* := \bar{a}^*(t_j, \hat{X}_{t_j}^{t_i, x})$, $C_j^* := (c_{\kappa}^* - \frac{(b_{\kappa}^* \sigma)^2}{2})(t_j, \hat{X}_{t_j}^{t_i, x})$ and G_j s are independent standard Gaussian random variables under the new equivalent measure \mathbb{Q} . By the consecutive use of the above inequality and the fact that $1 - a_j^* + a_j^* G_j^2 + h C_j^*$ is positive, one can write

$$u^{\kappa, h}(t_i, x) \leq \mathbb{E}^{\mathbb{Q}} \left[g(\hat{X}_T^{t_i, x}) \prod_{j=i}^{n-1} \left(1 - a_j^* + a_j^* G_j^2 + h C_j^* \right) \right] + \bar{C} h \sum_{j=i}^{n-1} e^{\theta_{\kappa} t_j}.$$

Notice that in the above inequality we used the fact that

$$\mathbb{E}_{t_j}^{\mathbb{Q}} \left[1 - a_j^* + a_j^* G_j^2 + h C_j^* \right] = 1 + h \mathbb{E}_{t_j}^{\mathbb{Q}} [C_j^*] \leq 1 + \theta_{\kappa} h.$$

On the other hand, $Z := \prod_{j=i}^{n-1} \left(1 - a_j^* + a_j^* G_j^2 + h C_j^* \right)$ is positive there for $\text{frac} Z \mathbb{E}^{\mathbb{Q}}[Z]$ could be considered as a density of a new measure \mathbb{Q}^Z with respect to \mathbb{P} . Therefore,

$$u^{\kappa, h}(t_i, x) \leq \mathbb{E}^{\mathbb{Q}}[Z] \mathbb{E}^{\mathbb{Q}^Z} \left[g(\hat{X}_T^{t_i, x}) \right] + \bar{C} h \sum_{j=i}^{n-1} e^{\theta_{\kappa} t_j}.$$

By the definition of $\bar{v}^{\kappa, h}$, one can write

$$\bar{v}^{\kappa, h}(t_i, x) \leq e^{-\theta_{\kappa}(T-t_i)} \mathbb{E}^{\mathbb{Q}}[Z] \mathbb{E}^{\mathbb{Q}^Z} \left[g(\hat{X}_T^{t_i, x}) \right] + e^{-\theta_{\kappa}(T-t_i)} \bar{C} h \sum_{j=i}^{n-1} e^{\theta_{\kappa} t_j}.$$

Therefore,

$$\bar{v}^{\kappa, h}(t_i, x) - g(x) \leq e^{-\theta_{\kappa}(T-t_i)} \mathbb{E}^{\mathbb{Q}}[Z] \mathbb{E}^{\mathbb{Q}^Z} \left[|g(\hat{X}_T^{t_i, x}) - g(x)| \right] + C |g(x)| (T-t_i) + e^{-\theta_{\kappa}(T-t_i)} \bar{C} (T-t_i).$$

Notice that $g(\hat{X}_T^{t_i, x}) - g(x)$ converges to zero \mathbb{P} -a.s. and therefore \mathbb{Q}^Z a.s. as $(t_i, h) \rightarrow (T, 0)$. So, by Lebesgue dominated convergence Theorem,

$$\limsup_{(t_i, h) \rightarrow (T, 0)} \bar{v}^{\kappa, h}(t_i, x) - g(x) \leq 0.$$

By the similar argument one can prove that:

$$\liminf_{(t_i, h) \rightarrow (T, 0)} \bar{v}^{\kappa, h}(t_i, x) - g(x) \geq 0,$$

which completes the proof. \square

Remark 4.8. By extending the above proof as in the Lemma 3.17 and Corollary 3.18 of [10], one can proof that

$$|\bar{v}^{\kappa, h}(t, x) - g(x)| \leq C(T-t)^{\frac{1}{2}}.$$

Also, observe that by the similar argument as in [10], $v^{\kappa, h}$ is $\frac{1}{2}$ -Hölder on t uniformly on h and x .

So, the approximate solution $\bar{v}^{\kappa_h, h}$ both satisfies the requirement of the convergence established in [1] and converges to a function v locally uniformly. Moreover, v is the unique viscosity solution of (2.1)–(2.2). So, by Corollary 4.2, the same assertion is true for $v^{\kappa, h}$.

4.2. Rate of Convergence. For local PDEs, the rate of convergence of probabilistic numerical scheme relies on the approximation of the solution of PDE by regular sub and super-solutions, the consistency for the scheme and the comparison principle derived from strict monotonicity. One can approximate the solution of the local PDE by a regular sub-solution and a almost regular super-solution from up and down, respectively. These approximations are provided by a switching system and Krylov method of shaking coefficients. Next, we use the consistency Lemma 3.22 in [10] to produce inequalities for the regular approximations plugged into the scheme. Then by comparison principle; Proposition 3.20 in [10]; we obtain the bounds for the difference of approximate solution derived from scheme and regular approximate solution obtained from Krylov method and switching system.

We continue this Subsection with establishing the same line of argument as in [10] for non-local case. The generalization of the method we used in [10] for the rate of convergence, is developed in [5] where the scheme needs to be consistent and satisfies comparison principle. Before, providing consistency and comparison principle result for the scheme (4.15), we show that truncation error could be handled by the Theorem of continuous dependence for (2.1)–(2.2). More precisely, if v and v^κ are solutions of (2.1)–(2.2) and (4.9)–(4.10), respectively; then by Theorem 5.1 in [6]

$$|v - v^\kappa|_\infty \leq C \left(\int_{0 < |z| < \kappa} |z|^2 \nu(dz) \right)^{\frac{1}{2}}.$$

Therefore, By choosing κ_h so that $\int_{0 < |z| < \kappa_h} |z|^2 \nu(dz) \leq Ch^{\frac{1}{2}}$, one can just concentrate on the rate of convergence of $v^{\kappa, h}$ to v^κ .

We shift to $\bar{v}^{\kappa_h, h}$ which is derived from the strictly monotone scheme (4.15) and find the rate of convergence for $\bar{v}^{\kappa_h, h}$. The following Corollary shows that this shift do not effect the rate of convergence.

Corollary 4.3. *For F which satisfies **IHJB**, $F(t, x, 0, 0, 0, 0) = 0$. Then,*

$$|\bar{v}^{\kappa_h, h} - v^{\kappa_h, h}| \leq Ch\theta_{\kappa_h}^2.$$

In addition, if κ_h is such that

$$\limsup_{h \rightarrow 0} h^{\frac{3}{4}} \theta_{\kappa_h}^2 < \infty,$$

then

$$|\bar{v}^{\kappa_h, h} - v^{\kappa_h, h}| \leq Ch^{\frac{1}{4}}$$

Proof. The proof is straightforward by the proof of Lemma 4.2. □

From now on, we concentrate on the approximate solution $\bar{v}^{\kappa,h}$ which is obtained from strictly monotone scheme 4.15 through (4.19). In order to provide the result, we need to use the consistency of the scheme for the regular approximate solutions. Then, the comparison principle for the scheme provides bounds over the difference between $u^{\kappa,h}$ and regular approximate solutions. Let

$$\begin{aligned} \mathcal{R}_{\kappa,h}[\psi](t,x) &:= \frac{\psi(t,x) - \overline{\mathbf{T}}_{\kappa,h}[\psi](t,x)}{h} + \mathcal{L}^X \psi(t,x) \\ &\quad + \overline{F}_{\kappa}(\cdot, \psi, D\psi, D^2\psi, \psi(t, \cdot))(t,x). \end{aligned}$$

Lemma 4.6. *For a family $\{\varphi_\varepsilon\}_{0 < \varepsilon < 1}$ of smooth functions satisfying*

$$\left| \partial_t^{\beta_0} D^\beta \varphi_\varepsilon \right| \leq C \varepsilon^{1-2\beta_0-|\beta|_1} \quad \text{for any } (\beta_0, \beta) \in \mathbb{N} \times \mathbb{N}^d \setminus \{0\}, \quad (4.21)$$

where $|\beta|_1 := \sum_{i=1}^d \beta_i$, and $C > 0$ is some constant, we have:

$$|\mathcal{R}_{\kappa,h}[\varphi_\varepsilon]|_\infty \leq R(h, \varepsilon) := C \left(h\varepsilon^{-3} + h\theta_\kappa \varepsilon^{-1} + h\sqrt{\theta_\kappa} + \varepsilon^{-1} \int_{\{|z| \leq \kappa\}} |z|^2 \nu(dz) \right),$$

for some constant $C > 0$ independent of κ . If in addition

$$\limsup_{h \rightarrow 0} h\theta_{\kappa_h}^2 < \infty \quad \text{and} \quad \limsup_{h \rightarrow 0} \sqrt{h} \int_{\{|z| \leq \kappa\}} |z|^2 \nu(dz) < \infty,$$

we have:

$$|\mathcal{R}_{\kappa_h,h}[\varphi_\varepsilon]|_\infty \leq R(h, \varepsilon) := C (h\varepsilon^{-3} + \sqrt{h}\varepsilon^{-1}).$$

Proof. $\mathcal{R}_{\kappa,h}[\varphi_\varepsilon]$ is bounded by

$$\begin{aligned} \sup_\alpha \left\{ \left| \mathbb{E} \left[\frac{1}{h} (\varphi_\varepsilon(t+h, X_h^{t,x,\kappa}) - \varphi_\varepsilon(t,x)) + \frac{1}{2} \text{Tr} \left[a^\alpha (D^2 \varphi_\varepsilon(t+h, X_h^{t,x,\kappa}) - D^2 \varphi_\varepsilon(t,x)) \right. \right. \right. \right. \\ \left. \left. \left. + b^\alpha (D \varphi_\varepsilon(t+h, X_h^{t,x,\kappa}) - D \varphi_\varepsilon(t,x)) + (\theta_\kappa + c^\alpha) (\varphi_\varepsilon(t+h, X_h^{t,x,\kappa}) - \varphi_\varepsilon(t,x)) \right. \right. \right. \\ \left. \left. \left. + \mathcal{I}^\alpha[\varphi_\varepsilon](t,x) - \mathcal{I}_{\kappa,h}^\alpha[\varphi_\varepsilon](t+h,x) \right] \right| \right\} \end{aligned}$$

For the Lévy integral term by Lemma 3.3, we have:

$$\begin{aligned} |\mathcal{I}^\alpha[\varphi_\varepsilon](t,x) - \mathcal{I}_{\kappa,h}^\alpha[\varphi_\varepsilon](t+h,x)| &\leq C \left(|D\varphi_\varepsilon|_\infty (\sqrt{h} + h \int_{\{|z| > \kappa\}} |z| \nu(dz)) \right. \\ &\quad \left. + h |\partial_t D^2 \varphi_\varepsilon|_\infty + |D^2 \varphi_\varepsilon|_\infty \int_{\{|z| \leq \kappa\}} |z|^2 \nu(dz) \right) \\ &\leq C \left(h\varepsilon^{-3} + h\sqrt{\theta_\kappa} + \varepsilon^{-1} \int_{\{|z| \leq \kappa\}} |z|^2 \nu(dz) \right). \end{aligned}$$

By the same argument as Lemma 3.22 in [10] all the other terms are bounded by $h\varepsilon^{-3}$ except

$$\theta_h \left(\varphi_\varepsilon(t+h, X_h^{t,x,\kappa}) - \varphi_\varepsilon(t,x) \right)$$

which is bounded by $\theta_h h \varepsilon^{-1}$. The second assertion of the Lemma is straightforward. \square

Next we need to have maximum principle for scheme 4.15. Note that Lemma 3.21 in [10] holds true for scheme 4.15 with $\lambda = \theta_\kappa + C$ and $\beta > \theta_\kappa + C$ where $C = \sup_\alpha |c^\alpha|$. Therefore, Proposition 3.20 in [10] holds true for non-local case. More precisely, we have the following Proposition.

Proposition 4.2. *Let Assumption **F** holds true, and consider two arbitrary bounded functions φ and ψ satisfying:*

$$h^{-1} (\varphi - \overline{\mathbf{T}}_h[\varphi]) \leq g_1 \quad \text{and} \quad h^{-1} (\psi - \overline{\mathbf{T}}_h[\psi]) \geq g_2$$

for some bounded functions g_1 and g_2 . Then, for every $i = 0, \dots, n$:

$$(\varphi - \psi)(t_i, x) \leq e^{(\theta_\kappa + C_1)} |(\varphi - \psi)^+(T, \cdot)|_\infty + (T - h) e^{(\theta_\kappa + C_1)(T - t_i)} |(g_1 - g_2)^+|_\infty$$

where $C_1 > \sup_\alpha |c^\alpha|$.

The approximation of the solution of non-local PDE by the Krylov method of shaking coefficients and switching system is developed in [5]. [5] provides the result of rate of convergence of general monotone schemes for the non-local PDEs satisfying Assumption IHJB. However, they referred the regularity of the approximate solutions to the result of [6] where the approximate solution obtained from switching system could only be locally $\frac{1}{2}$ -Hölder continuous on t . But, in the case of scheme (4.3), we need the solution of (2.1)–(2.2) be uniformly $\frac{1}{2}$ -Hölder continuous on t . It is because we need the regular approximate solutions obtained from Krylov method and switching solution to satisfy (4.21). Therefore, in the present work we need to rebuild Lemma 5.3 in [6] under the Assumption IHJB to obtain global $\frac{1}{2}$ -Hölder continuous on t for the solution of the switching system.

Therefore, we continue this subsection by introducing the switching system of non-local PDEs with the regularity result needed for the solution of this system.

Let k be a non-negative constant. Suppose the following system of PDEs:

$$\begin{aligned} \max \{ -\mathcal{L}^X v_i(t, x) - F_i(t, x, v_i(t, x), Dv_i(t, x), D^2 v_i(t, x), v_i(t, \cdot)), v_i - \mathcal{M}^i v \} = 0 \\ v_i(T, \cdot) = g_i(\cdot), \end{aligned} \quad (4.22)$$

where $i = 1, \dots, M$ and

$$\begin{aligned} F_i(t, x, r, p, \gamma, \psi) &:= \inf_{\alpha \in \mathcal{A}_i} \{ \mathcal{L}^\alpha(t, x, r, p, \gamma, \psi) + \mathcal{I}^\alpha(t, x, r, p, \gamma, \psi) \} \\ \mathcal{L}^\alpha(t, x, r, p, \gamma, \gamma) &:= \frac{1}{2} \text{Tr} [a^\alpha(t, x) \gamma] + b^\alpha(t, x) \cdot p + c^\alpha(t, x) r + k^\alpha(t, x) \\ \mathcal{I}^\alpha(t, x, r, p, \gamma, \psi) &:= \int_{\mathbb{R}_*^d} (\psi(t, x + \eta^\alpha(t, x, z)) - r - \mathbf{1}_{\{|z| \leq 1\}} \eta^\alpha(t, x, z) \cdot p) d\nu(z) \\ \mathcal{M}^i r &:= \min_{j \neq i} r_j + k. \end{aligned}$$

We would like to emphasize that g_i s need to satisfy $g_i - \mathcal{M}^i \bar{g} \leq 0$ where $\bar{g} = (g_1, \dots, g_M)$. If each $g_i = g$ then we obviously have $g_i - \mathcal{M}^i \bar{g} \leq 0$.

Existence and comparison principle result for the above switching system is provided in Proposition 6.1 [5]. Also, it is known from Theorem 6.3 in [5], that if (v^1, \dots, v^M) and v be respectively the solutions of (4.22) and (2.1)-(2.2) with $\mathcal{A} = \cup_{i=1}^M \mathcal{A}_i$ and \mathcal{A}_i s are disjoint sets, then

$$0 \leq v^i - v \leq Ck^{\frac{1}{3}} \text{ for } i = 1, \dots, M. \quad (4.23)$$

The regularity result for (4.22) is provided in [6]. There, it is proved that (v^i) is Lipschitz with respect to x and locally 1/2-Hölder continuous with respect to t . For the proof of Theorem 4.2, (v^i) should be uniformly 1/2-Hölder continuous with respect to t . The following Lemma provide the uniform 1/2-Hölder continuity for (v^i) .

Lemma 4.7. *Assume **HJB** holds for each i and let (v^i) be the viscosity solution of (4.22). Then there exist a constant C such that for any $i = 1, \dots, M$:*

$$|v^i|_1 \leq C.$$

Proof. Lipschitz continuity with respect to x is done in Lemma 5.2 in [6]. To obtain uniform 1/2-Hölder continuity with respect to t , the proof of Lemma 5.3 in [6] should be modified by using assumption **HJB**.

Fix $y \in \mathbb{R}^d$, t and t' where $t \leq t'$. For each $i = 1, \dots, M$, define:

$$\psi_i(t, x) := \lambda \frac{L}{2} \left[e^{A(t'-t)} |x - y|^2 + B(t' - t) \right] + K(t' - t) + \lambda^{-1} \frac{L}{2} + v^i(t', y)$$

Where $L = \frac{1}{2}|v|_1$ and λ , a and γ will be defined later. Then:

$$\begin{aligned} \partial_t \psi_i(t, x) &= -\lambda \frac{L}{2} \left(A e^{A(t'-t)} |x - y|^2 + B \right) - K \\ D \psi_i(t, x) &= \lambda L e^{A(t'-t)} (x - y) \\ D^2 \psi(t, x) &= \lambda L e^{A(t'-t)} I_{d \times d}. \end{aligned}$$

So,

$$\begin{aligned} & -\partial_t \psi_i + \mathcal{L}^\alpha(t, x, \psi_i, D \psi_i, D^2 \psi_i) + \mathcal{I}^\alpha(t, x, \psi_i, D \psi_i) \\ &= -\lambda L \left(A e^{A(t'-t)} |x - y|^2 + B \right) - K \\ & \quad + \lambda L e^{A(t'-t)} \text{Tr} [a^\alpha(t, x)] + 2\lambda L e^{A(t'-t)} b^\alpha(t, x) \cdot (x - y) + c^\alpha(t, x) \psi_i + k^\alpha(t, x) \\ & \quad + \lambda \frac{L}{2} e^{A(t'-t)} \int_{\mathbb{R}_*^d} (|x + \eta^\alpha(t, x, z) - y|^2 - |x - y|^2 - 2\mathbb{1}_{\{|z| \leq 1\}} \eta^\alpha(t, x, z) \cdot (x - y)) d\nu(z). \end{aligned}$$

By **HJB**, we choose K and λ so that,

$$\begin{aligned} |a^\alpha|_\infty \leq K, |b^\alpha|_\infty \leq K, |c^\alpha|_\infty \leq K, |k^\alpha|_\infty \leq K, K^{-1} \leq \lambda \leq K \\ |v|_\infty \leq K, |\eta^\alpha(t, x, z)| \leq K(1 \wedge |z|). \end{aligned}$$

Without loss of generality and with the similar argument as in Remark 3.7, we can suppose that for any α , $c^\alpha \leq 0$. So by choosing positive large D , there exists positive constants C_1 , C_2 and C_3

such that:

$$\begin{aligned} & -\partial_t \psi_i + \mathcal{L}^\alpha(t, x, \psi_i, D\psi_i, D^2\psi_i) + \mathcal{I}^\alpha(t, x, \psi_i, D\psi_i) \\ & \leq -\lambda L e^{A(t'-t)} K \left(\left(\frac{A}{K} - \frac{1}{2} \right) |x-y|^2 - \frac{3}{2} |x-y| + C_1 B - C_2 \right) + C_3. \end{aligned}$$

Therefore, choice of large B and D makes the right hand side negative.

$$-\partial_t \psi_i + \mathcal{L}^\alpha(t, x, \psi_i, D\psi_i, D^2\psi_i) + \mathcal{I}^\alpha(t, x, \psi_i, D\psi_i) \leq 0.$$

On the other hand,

$$\psi(t', x) = \frac{L}{2} (\lambda |x-y|^2 + \lambda^{-1}) + v^i(t', y).$$

Minimizing with respect to λ ,

$$\psi(t', x) \geq L|x-y| + v^i(t', y) \geq v^i(t', x).$$

We can conclude that ψ_i is a super solution of (4.22). So, by comparison Theorem in [5],

$$\psi_i(t, y) \geq v^i(t, y).$$

So,

$$\frac{L}{2} (\lambda B(t'-t) + \lambda^{-1}) + v^i(t', y) \geq v^i(t, y).$$

Therefore,

$$v^i(t, y) - v^i(t', y) \leq C\sqrt{t'-t}.$$

The other inequality can be done similarly by choosing:

$$\psi_i(t, x) := -\lambda \frac{L}{2} \left[e^{A(t'-t)} |x-y|^2 - B(t'-t) \right] - K(t'-t) - \lambda^{-1} \frac{L}{2} + v^i(t', y).$$

□

Remark 4.9. Notice that all the result of switching system is correct for (2.1)-(2.2) satisfying \mathbf{F} by simply setting $M = 1$ and $k = 0$.

Therefore, by [5] there are regular functions $\underline{w}_\varepsilon^\kappa$ and $\overline{w}_\varepsilon^\kappa$ which are respectively the regular sub- and super-solution of

$$\begin{aligned} & -\mathcal{L}^X u^\kappa(t, x) - \overline{F}_\kappa(t, x, u^\kappa(t, x), Du^\kappa(t, x), D^2u^\kappa(t, x), u^\kappa(t, \cdot)) = 0, \text{ on } [0, T] \times \mathbb{R}^d, \\ & u^\kappa(T, \cdot) = g, \text{ on } \mathbb{R}^d. \end{aligned}$$

where

$$\overline{F}_\kappa(t, x, r, p, \gamma, \psi) := \inf_{\alpha \in \mathcal{A}} \left\{ \mathcal{L}^{\alpha, \beta}(t, x, r, p, \gamma) + \mathcal{I}_\kappa^{\alpha, \beta}(t, x, r, p, \gamma, \psi) \right\}$$

(one can replace sup inf by inf sup) where

$$\mathcal{L}^{\alpha, \beta}(t, x, r, p, \gamma) := \frac{1}{2} \text{Tr} \left[\sigma^{\alpha, \beta} \sigma^{\alpha, \beta \Gamma}(t, x) \gamma \right] + b^{\alpha, \beta}(t, x) p + (c^{\alpha, \beta}(t, x) + \theta_\kappa) r,$$

and

$$\mathcal{I}_\kappa^{\alpha,\beta}(t, x, r, p, \gamma, \psi) := \int_{\{|z|>\kappa\}} (\psi(x + \eta^{\alpha,\beta}(t, x, z)) - r - \mathbf{1}_{\{|z|\leq 1\}} \eta^{\alpha,\beta}(t, x, z) \cdot p) \nu(dz).$$

Therefore, by Proposition 6.2 and Theorem 6.3 of [5], Lemma 4.6 and Proposition 4.2,

$$\begin{aligned} (u^\kappa - u^{\kappa,h})(t, x) &\leq (u^\kappa - \underline{w}_\varepsilon^\kappa + \underline{w}_\varepsilon^\kappa - u^{\kappa,h})(t, x) \\ &\leq C e^{(\theta_\kappa + C_1)(T-t)} \left(\varepsilon + h\varepsilon^{-3} + h\theta_\kappa \varepsilon^{-1} + h\sqrt{\theta_\kappa} + \varepsilon^{-1} \int_{\{|z|\leq \kappa\}} |z|^2 \nu(dz) \right) \end{aligned}$$

and

$$\begin{aligned} (u^{\kappa,h} - u^\kappa)(t, x) &\leq (u^{\kappa,h} - \bar{w}_\varepsilon^\kappa + \bar{w}_\varepsilon^\kappa - u^\kappa)(t, x) \\ &\leq C e^{(\theta_\kappa + C_1)(T-t)} \left(\varepsilon^{\frac{1}{3}} + h\varepsilon^{-3} + h\theta_\kappa \varepsilon^{-1} + h\sqrt{\theta_\kappa} + \varepsilon^{-1} \int_{\{|z|\leq \kappa\}} |z|^2 \nu(dz) \right). \end{aligned}$$

Notice that $v^\kappa(t, x) = e^{-\theta_\kappa(T-t)} u^\kappa(t, x)$. So,

$$v^\kappa - \bar{v}^{\kappa,h} \leq C \left(\varepsilon + h\varepsilon^{-3} + h\theta_\kappa \varepsilon^{-1} + h\sqrt{\theta_\kappa} + \varepsilon^{-1} \int_{\{|z|\leq \kappa\}} |z|^2 \nu(dz) \right)$$

and

$$\bar{v}^{\kappa,h} - v^\kappa \leq C \left(\varepsilon^{\frac{1}{3}} + h\varepsilon^{-3} + h\theta_\kappa \varepsilon^{-1} + h\sqrt{\theta_\kappa} + \varepsilon^{-1} \int_{\{|z|\leq \kappa\}} |z|^2 \nu(dz) \right).$$

On the other hand, because of (4.14) and by Lemma (4.2), the second part of Theorem 4.2 is provided after choice of optimal ε .

5. CONCLUSION

The algorithm is the first probabilistic numerical method for fully non-linear nonlocal parabolic problems. As in local case ([10]), it converges to the viscosity solution of the problem. A rate of convergence is known for the convex (concave) non-linearities. Also with the same argument as in Section 4 in [10], Monte Carlo approximations of expectations inside the scheme do not affect the asymptotic results if enough number of samples would be used. The error analysis for MCQ shows that the appropriate approximation of jump-diffusion process with compound Poisson process could be applied in discretization procedure. The theoretical result is followed by some numerical examples which confirms the convergence of the scheme.

On the other hand there are some features where the scheme is not implementable in non-local case, e.g. when the non-linearity is of HJB type. This could be the challenge of future works. The other open issue is to relax some assumptions. For example, relaxing the assumption of uniform ellipticity may be a future work.

REFERENCES

- [1] G. Barles and P. E. Souganidis. Convergence of approximation schemes for fully nonlinear second order equations. *Asymptotic Anal.*, 4(3):271–283, 1991.
- [2] Richard F. Bass. Stochastic differential equations with jumps. *Probab. Surv.*, 1:1–19 (electronic), 2004.
- [3] Fred Espen Benth, Kenneth Hvistendahl Karlsen, and Kristin Reikvam. Optimal portfolio management rules in a non-Gaussian market with durability and intertemporal substitution. *Finance Stoch.*, 5(4):447–467, 2001.
- [4] Fred Espen Benth, Kenneth Hvistendahl Karlsen, and Kristin Reikvam. Optimal portfolio selection with consumption and nonlinear integro-differential equations with gradient constraint: a viscosity solution approach. *Finance Stoch.*, 5(3):275–303, 2001.
- [5] Imran H. Biswas, Espen R. Jakobsen, and Kenneth H. Karlsen. Error estimates for a class of finite difference-quadrature schemes for fully nonlinear degenerate parabolic integro-PDEs. *J. Hyperbolic Differ. Equ.*, 5(1):187–219, 2008.
- [6] Imran H. Biswas, Espen R. Jakobsen, and Kenneth H. Karlsen. Difference-quadrature schemes for nonlinear degenerate parabolic integro-PDE. *SIAM J. Numer. Anal.*, 48(3):1110–1135, 2010.
- [7] Bruno Bouchard and Romuald Elie. Discrete-time approximation of decoupled forward-backward SDE with jumps. *Stochastic Process. Appl.*, 118(1):53–75, 2008.
- [8] Patrick Cheridito, H. Mete Soner, Nizar Touzi, and Nicolas Victoir. Second-order backward stochastic differential equations and fully nonlinear parabolic PDEs. *Comm. Pure Appl. Math.*, 60(7):1081–1110, 2007.
- [9] Rama Cont and Peter Tankov. *Financial modelling with jump processes*. Chapman & Hall/CRC Financial Mathematics Series. Chapman & Hall/CRC, Boca Raton, FL, 2004.
- [10] Arash Fahim, Nizar Touzi, and Xavier Warin. A probabilistic numerical method for fully non-linear parabolic pdes. *Annals of Applied Probability*, 21(4):1322–1364, 2011.
- [11] Cyril Imbert and Sylvia Serfaty. Repeated games for non-linear parabolic integro-differential equations and integral curvature flows. *Discrete Contin. Dyn. Syst.*, 29(4):1517–1552, 2011.
- [12] Robert V. Kohn and Sylvia Serfaty. A deterministic-control-based approach to motion by curvature. *Comm. Pure Appl. Math.*, 59(3):344–407, 2006.
- [13] Daniel Stroock. Diffusion processes associated with lévy generators. *Z. Wahrscheinlichkeitstheorie und Verw. Gebiete*, 32(3):209–244, 1975.

(Arash Fahim) CMAP,, ECOLE POLYTECHNIQUE, PARIS

E-mail address: fahimara@umich.edu