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Assessing velocity and impedance changes due to CO₂ saturation using interferometry on repeated seismic sources.

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The injection of CO_2 into saline aquifers produces small changes in seismic properties. The aim of time lapse monitoring is to look for these small variations of the seismic attributes through time, such as the velocity and the impedance, and to relate them to perturbations in the medium properties. However, as injection proceeds, at higher levels of saturation, relative changes in the data become very small. Assuming that a time-lapse seismic survey is a doublet, we can determine the time delay between repeat survey differences, using wave interferometry. These are then relates to velocity changes. Additionally, we can do the same for amplitude variations and determine impedance modifications. We investigate the feasibility of this technique on 1) numerical data computed from a simple layered porous medium; and 2) for synthetic data calculated in a more realistic 3D reservoir using a poro-elastic finite-difference solution. In both cases, the reconstructed values are in good agreement with the true modifications.



Introduction

The role played by the industrial emission of carbon dioxide (CO_2) in climate change has been well documented. Geological sequestration is a process to store CO_2 in the subsurface in order to reduce emissions of this gas into the atmosphere. Gas and oil reservoirs and saline aquifers are potential subsurface storage sites. For example, CO_2 has been injected into the Sleipner area (Arts et al., 2004; Carcione et al, 2006) for more than 15 years. To prevent risks of leakage, it is necessary to monitor the distribution of the CO_2 when injected and its flow after that. Seismic waves are sensitive to fluid properties through their amplitude and velocity. However, the modifications in the seismic waves can be small as shown by Davis et al. (2003). Particularly, as injection proceeds at higher levels of CO_2 , changes become very small. It is quite difficult to describe the absolute fluid properties, but it is easier to look at relative changes in the medium due to injection and/or fluid movement. We are interested in measuring the differences between data corresponding to two injection times, rather than absolute differences with the baseline data.

Coda wave interferometry has been used in seismology to look at medium changes using seismic multiplets. This method has been used to find wave velocity changes, as shown by Poupinet et al. (1984), and later by Snieder et al. (2002) and Pandolfi et al. (2006). Time-lapse active seismic data can be considered as seismic multiplets and therefore can be used to look for small changes in the medium properties. The density, through the impedance, is a parameter that is more sensitive to the fluid than the velocity. Moreover, in a reflection survey, the wave amplitude can be related to the impedance contrasts. We derive a simple relation between the amplitude change and the impedance contrast. Then, we describe the methodology used to measure both the amplitude and time change in layered media to infer changes in fluid properties following the injection of CO₂. Finally, we numerically investigate the accuracy of the results in a complex 3D medium.

Wave interferometry

Interferometry is based on the measurement of time delays, which is then converted into velocity changes δV through the simple relation (Snieder et al., 2002)

$$\delta V/V = -\delta t/t$$
, (1)

where V is the velocity of the medium, and δt is the measured travel time delay. This relation has been defined for an overall change of the medium assuming scattering waves. However, this assumption of a diffusive regime is not valid for simple layered models. We consider a model with a layer L_1 with length H embedded in a infinite homogeneous medium (L_0). Assuming a perturbation δV_1 in the layer L_1 , the delay for the reflected waves on the bottom of L_1 is:

$$\delta t = -2H_1/V_1 (\delta V_1/V_1).$$
 (2)

Equation (2) is equivalent to equation (1) assuming that $t=2H_1/V_1$. That means that t in eq. 1 is the time from the last reflected waves without time delays, i.e. the reflected waves on top of L_1 .

In the same way, we want to relate a small change in the amplitude A of the reflected waves to a impedance modification. At normal incidence, $A=(Z_1-Z_0)/(Z_1+Z_0)$, where Z_1 and Z_0 denote the impedance in layer L_1 and L_0 . A small change δZ_1 and/or δZ_0 in the impedances will result in an amplitude $A'=(Z_1+\delta Z_1-Z_0-\delta Z_0)/(Z_1+\delta Z_1+Z_0+\delta Z_0)$. Assuming that $\delta Z_1+\delta Z_0<<< Z_1+Z_0$, we have

$$\delta A/A = (A'-A)/A = (\delta Z_1 - \delta Z_0)/(Z_1 - Z_0) = \delta \Delta Z/\Delta Z$$
. (3)



This relation means that a small relative change in amplitude is equivalent to a relative change in impedance contrast. It can be generalised for the reflected waves from the bottom of the layer, assuming that the impedance contrast is small. Equation (3) is valid under the assumptions of 1) a normal incidence of the reflected waves, 2) a small change in the impedance contrast and 3) a small impedance contrast. However, both relations (1) and (3) are derived only for a change in a simple one-layer model. It is thus necessary to check the behaviour of these relations in the case of more complex geometry.

Plane layered model

To check the accuracy of the method, we applied it to a simple layered model. We consider a layered porous reservoir between 775 m and 1125 m depth with shale and sand layers. The reservoir is embedded in an infinite medium. CO_2 is injected into the first layer (775-825 m), and we want to quantify a change of a CO_2 saturation from 20% to 40%. Properties of CO_2 are computed using the relations of Batzle and Wang (1992), Carcione et al. (2006). Data are computed at normal incidence for an explosive source (Ricker wavelet, 30 Hz frequency peak) using a reflectivity method (Kennett, 1981) followed by a wavenumber integration method (Bouchon, 1981) in a poroelastic medium (see Pride, 2005 and De Barros and Dietrich, 2008).

True velocity and impedance changes are given in figure 1. We measure the time delay and the amplitude changes using a time windowed cross correlation of the two traces. For each window and measured point, we find the time lag corresponding to the maximum of the correlation coefficient. To be more precise, we use a spline interpolation around it to allow a precision on the delays smaller than the data time step. We then align the two signals and compute the amplitude ratio. The amplitude A is defined as the square root of the sum of the squared amplitude of the samples contained in the window. To avoid instabilities due to null signals when computing the amplitude ratio, we apply a water level method $(10^{-6} \text{ on A in eq. 3})$. For every measurement point, we repeat the measures assuming different window lengths (here, 8 measures from half period to 5 dominant periods). If the correlation coefficient is lower than a threshold (here 0.9), values are not kept. t is computed from the time of the reflected waves at the top of the reservoir. Velocity and amplitude variation are then obtained from eq. (1) and (3).

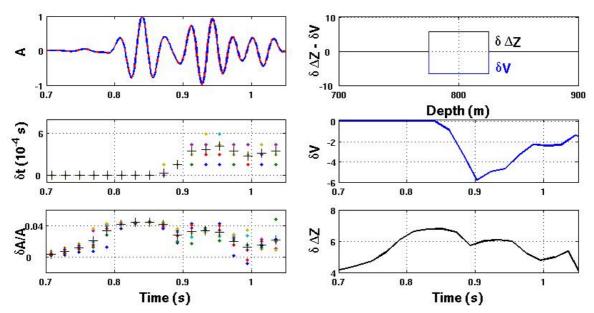


Figure 1 Application of the method to a layered model. Left) Top: Data computed with 20% (red) and 40% (blue) CO_2 saturation; Middle: measured time delay δt ; Bottom: measured amplitude contrast $\delta A/A$. Coloured dots correspond to different sliding window lengths, + signs are the mean values. Right) Top: True model showing velocity and impedance changes; Middle: Estimated velocity



change; Bottom: Estimated impedance change. Note that the true values are given according to depth, while the measured changes are given according to time. The impedance change has been divided by 10^4 for figure clarity.

In figure 1, we see that the maximum value obtained for the change (δV =-5.8 m/s and $\delta \Delta Z$ =67000 kg/m²/s) are very similar to the true values (δV =-6.5 m/s and $\delta \Delta Z$ =61000 kg/m²/s). For this simple case, this method seems accurate. As expected, the velocity change is seen in the reflected waves from the bottom of the perturbed layer, while the change in amplitude is seen by the two reflections, but is more accurate for the reflected waves at the top of the perturbed layer. Finally, the reconstructed changes are smooth, as the entire Ricker wavelet carries information about the perturbation. We computed δV and $\delta \Delta Z$ in this case as we know V and ΔZ , perfectly. However, in a real case, we can directly estimate a variation of the fluid properties from $\delta V/V$ and $\delta \Delta Z/\Delta Z$ using the Gassmann fluid substitution relations.

3D synthetic reservoir

We use a more complex 3D model with a 80 m thick reservoir, with 5 different facies where the geometry is based on a statistical ensemble of shallow petroleum reservoirs (Manzocchi et al., 2008). Waveforms are computed in a poroelastic heterogeneous medium using the staggered finite difference solver (4^{th} order in space, 2^{nd} order in time) of O'Brien (submitted). Synthetic data are computed for 1) a medium fully saturated by water (baseline data), 2) a CO_2 saturation of 20% and 3) a 40% saturation in an area with a maximum radius of 200 m where the radius is related to the permeability of the facies. Figure 2 shows the data differences between the baseline and the data with 20% and 40% CO_2 saturation.

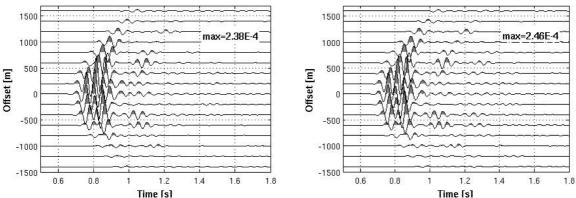


Figure 2 *Differential seismograms between reference data and data computed for 20% (left) and 40% (right) of CO₂. Centre of CO₂ patch and seismic source are located at the zero offset position.*

We apply the interferometry method in order to measure the difference between the data corresponding to 20% and 40% CO_2 saturation. To get rid of the energetic waves that do not carry information about the fluid substitution we use the differential data of figure 2. We use the traces closest to the sources, i.e. with an angle of incidence smaller than 15° .



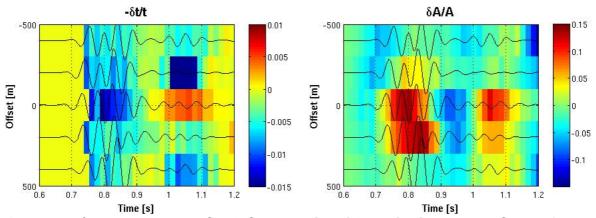


Figure 3 *Left) Time variation* $-\delta t/t$ ($=\delta V/V$) *and Right) amplitude variation* $\delta A/A$ ($=\delta \Delta Z/\Delta Z$) *measured when comparing differential seismograms of figure 2.*

The theoretical mean values for the change are $\delta V/V=-0.8\%$ and $\delta \Delta Z/\Delta Z=17.8\%$. Figure 3 shows the results. We see that both the time delay and amplitude change gave values in good agreement with the true theoretical mean changes. The order of magnitude of these values confirm that the sensitivity to a fluid change is higher for the impedance than for the velocity. A time change is present in all the seismograms and does not allow us to localise the CO_2 change. On the other hand, the amplitude change seems to be located above the medium perturbations.

Conclusions

Injection of CO_2 (or any other fluid injection or extraction) changes both the velocity and the amplitude of reflected waves. However, the change can be very small and as such, high resolution seismic techniques have to be employed to quantify these changes. One such technique is seismic wave interferometry. It is easy to apply and allows the measurement of very small variations in the velocity and amplitude of the seismic wavefield. We use this method on synthetic active seismic time-lapse data to measure both the velocity and the impedance changes. The time-lapse difference seismograms are treated as seismic multiplets. The preliminary tests on a plane layered and complex 3D models show good agreement between the measured and true values. However, it gives a smooth picture of the perturbations. The relationships between the amplitude contrast and impedance have still to be investigated in more detail, both with analytical and numerical approaches. A direct interpretation of the variations of the seismic attributes into fluid perturbations will also be performed.

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