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Nizar Demni. Radon transform on spheres and generalized Bessel function associated with dihedral groups. 2010. hal-00519722

HAL Id: hal-00519722 https://hal.science/hal-00519722

Preprint submitted on 21 Sep 2010

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RADON TRANSFORM ON SPHERES AND GENERALIZED BESSEL FUNCTION ASSOCIATED WITH DIHEDRAL GROUPS

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ABSTRACT. Motivated by Dunkl operators theory, we consider a generating series involving a modified Bessel function and a Gegenbauer polynomial, that generalizes a known series already considered by L. Gegenbauer. We actually use inversion formulas for Fourier and Radon transforms to derive a closed formula for this series when the parameter of the Gegenbauer polynomial is a strictly positive integer. As a by-product, we get a relatively simple integral representation for the generalized Bessel function associated with even dihedral groups $D_2(2p), p \ge 1$ when both multiplicities sum to an integer. In particular, we recover a previous result obtained for $D_2(4)$ and we give a special interest to $D_2(6)$. The paper is closed with adapting our method to odd dihedral groups thereby exhausting the list of Weyl dihedral groups.

1. INTRODUCTION

The dihedral group $D_2(n)$ of order $n \geq 2$ is defined as the group of regular *n*-gone preserving-symmetries ([8]). It figures among reflections groups associated with root systems for which a spherical harmonics theory, generalizing the one of Harish-Chandra on semisimple Lie groups from a discrete to a continuous range of multiplicities, was introduced by C. F. Dunkl in the late eightees (see Ch.I in [3]). Since then, a huge amount of research papers on this new topic and on its stochastic side as well emerged yielding fascinating results (Ch. II, III in [3]). For instance, probabilistic considerations allowed the author to derive the so-called generalized Bessel function associated with dihedral groups ([4]). For even values $n = 2p, p \geq 1$, this function depending on two real variables, say $(x, y) \in \mathbb{R}^2$, is expressed in polar coordinates $x = \rho e^{i\phi}, y = r e^{i\theta}, \rho, r \geq 0, \phi, \theta \in [0, \pi/2p]$ as

(1)
$$D_k^W(\rho, \phi, r, \theta) = c_{p,k} \left(\frac{2}{r\rho}\right)^{\gamma} \sum_{j\geq 0} I_{2jp+\gamma}(\rho r) p_j^{l_1, l_0}(\cos(2p\phi)) p_j^{l_1, l_0}(\cos(2p\theta))$$

where

- $k = (k_0, k_1)$ is a positive-valued multiplicity function, $l_i = k_i 1/2, i \in \{1, 2\}, \gamma = p(k_0 + k_1)$. • $I_{2jp+\gamma}, p_j^{l_1, l_0}$ are the modified Bessel function of index $2jp + \gamma$ and the
- $I_{2jp+\gamma}, p_j^{\epsilon_1, \epsilon_0}$ are the modified Bessel function of index $2jp + \gamma$ and the *j*-th orthonormal Jacobi polynomial of parameters l_1, l_0 respectively (the orthogonality (Beta) measure need not to be normalized here. In fact, the normalization only alters the constant $c_{p,k}$ below).

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• The constant $c_{p,k}$ depends on p, k and is such that $D_k^W(0, y) = 1$ for all $y = (r, \theta) \in [0, \infty) \times [0, \pi/2p]$ (see [5])

$$c_{p,k} = 2^{k_0+k_1} \frac{\Gamma(p(k_1+k_0)+1)\Gamma(k_1+1/2)\Gamma(k_0+1/2)}{\Gamma(k_0+k_1+1)}.$$

In a subsequent paper ([5]), the special case p = 2 corresponding to the group of square-preserving symmetries was considered. The main ingredient used there was the famous Dijksma-Koornwinder's product formula for Jacobi polynomials ([7]) which may be written in the following way ([5]):

$$c(\alpha,\beta)p_j^{\alpha,\beta}(\cos 2\phi)p_j^{\alpha,\beta}(\cos 2\theta) = (2j+\alpha+\beta+1)\int\int C_{2j}^{\alpha+\beta+1}(z_{\phi,\theta}(u,v))\mu^{\alpha}(du)\mu^{\beta}(dv)$$

where $\alpha, \beta > -1/2$,

$$c(\alpha,\beta) = 2^{\alpha+\beta+1} \frac{\Gamma(\alpha+1)\Gamma(\beta+1)}{\Gamma(\alpha+\beta+1)},$$

 $z_{\phi,\theta}(u,v) = u\cos\theta\cos\phi + v\sin\theta\sin\phi,$

and μ^{α} is the symmetric Beta probability measure whose density is given by

$$\mu^{\alpha}(du) = \frac{\Gamma(\alpha+1)}{\sqrt{\pi}\Gamma(\alpha+1/2)} (1-u^2)^{\alpha-1/2} \mathbf{1}_{[-1,1]}(u) du, \quad \alpha > -1/2.$$

Inverting the order of integration, we were in front of the following series

(2)
$$\left(\frac{2}{r\rho}\right)^{\gamma} \sum_{j\geq 0} (2j+k_0+k_1) I_{2jp+\gamma}(\rho r) C_{2j}^{k_0+k_1}(z_{p\phi,p\theta}(u,v))$$

for $(u, v) \in]-1, 1[^2$, which specializes for p = 2 to

$$\frac{1}{2} \sum_{j \equiv 0[4]} (j+\gamma) I_{j+\gamma}(\rho r) C_{j/2}^{\gamma/2}(z_{2\phi,2\theta}(u,v)).$$

Using the identity noticed by Y. Xu ([13]):

$$C_{j}^{\nu}(\cos\zeta) = \int C_{2j}^{2\nu}\left(\sqrt{\frac{1+\cos\zeta}{2}}z\right)\mu^{\nu-1/2}(dz), \quad \nu > -1/2, \, \xi \in [0,\pi],$$

we were led to

$$\sum_{j\equiv 0[4]} (j+\gamma) I_{j+\gamma}(\rho r) C_j^{\gamma}(z_{2\phi,2\theta}(u,v))$$

which we wrote as

$$\frac{1}{4} \sum_{s=1}^{4} \sum_{j \ge 0} (j+\gamma) I_{j+\gamma}(\rho r) C_j^{\gamma}(z_{2\phi,2\theta}(u,v)) e^{isj\pi/2}$$

after the use of the elementary identity

(3)
$$\frac{1}{n}\sum_{s=1}^{m}e^{2i\pi sj/m} = \begin{cases} 1 & \text{if } j \equiv 0[m], \\ 0 & \text{otherwise,} \end{cases}$$

valid for any integer $m \ge 1$. Accordingly (Corollary 1.2 in [5])

$$D_{k}^{W}(\rho,\phi,r,\theta) = \int \int i_{(\gamma-1)/2} \left(\rho r \sqrt{\frac{1+z_{2\phi,2\theta}(u,v)}{2}} \right) \mu^{l_{1}}(du) \mu^{l_{0}}(dv)$$
²

where

$$i_{\alpha}(x) := \sum_{m=0}^{\infty} \frac{1}{(\alpha+1)_m m!} \left(\frac{x}{2}\right)^{2m}$$

is the normalized modified Bessel function ([8]) and $\gamma = 2(k_0 + k_1) \ge 2$ is even. This is a relatively simple integral representation of D_k^W since the latter function may be expressed as a bivariate hypergeometric function of Bessel-type. Recall also that it follows essentially from closed formulas due to L. Gegenbauer (equations (4), (5), p.369 in [12]):

$$\left(\frac{2}{r\rho}\right)^{\gamma} \sum_{j\geq 0} (j+\gamma) I_{j+\gamma}(\rho r) C_j^{\gamma}(\cos\zeta) (\pm 1)^j = \frac{1}{\Gamma(\gamma)} e^{\pm\rho r \cos\zeta}.$$

In this paper, we shall see that a relatively simple integral representation of D_k^W still exists for general integer $p \ge 2$ and integer $\nu := k_0 + k_1 \ge 1^2$. In fact, with regard to (2), one has to derive closed formulas for both series below

(4)
$$f_{\nu,p}^{\pm}(R, \cos\zeta) := \left(\frac{2}{R}\right)^{p\nu} \sum_{j\geq 0} (j+\nu) I_{p(j+\nu)}(R) C_j^{\nu}(\cos\zeta) (\pm 1)^j$$

with $R = \rho r$ and $\cos \zeta := \cos \zeta(u, v) = z_{p\phi,p\theta}(u, v)$. The obtained formulas reduce to Gegenbauer's results when $p = 1, \nu \ge 1$ is an integer, and do not exist up to our knowledge. Moreover, our approach is somewhat geometric since we shall interpret the sequence:

$$(\pm 1)^{j} I_{p(j+\nu)}(R), \, j \ge 0$$

for fixed R as the Gegenbauer-Fourier coefficients of $\zeta \mapsto f_{\nu,p}^{\pm}(R, \cos \zeta)$, and since spherical functions on the sphere viewed as a homogeneous space are expressed by means of Gegenbauer polynomials ([1]). Then, following [1], solving the problem when ν is a strictly positive integer amounts to appropriately use inversion formulas for Fourier and Radon transforms. Our main result is stated as

Proposition 1. Assume $\nu \geq 1$ is a strictly positive integer, then

$$\left(\frac{R}{2}\right)^{p\nu} f_{\nu,p}^{\pm}(R, \cos\zeta) = \frac{1}{2^{\nu}(\nu-1)!} \left[-\frac{1}{\sin\zeta} \frac{d}{d\zeta}\right]^{\nu} \frac{1}{p} \sum_{s=1}^{p} e^{\pm R\cos[(\zeta+2\pi s)/p]}.$$

A first glance at the main result may be ambiguous for the reader since the LHS depends on $\cos \zeta$ while the RHS depends on $\cos(\zeta/p), p \ge 1$. But $\cos(\zeta/p), p \ge 1$ may be expressed, though in a very complicated way (inverses of linearization formulas), as a function of $\cos \zeta$. For instance, when p = 2,

$$\cos(\zeta/2) = \sqrt{\frac{1+\cos\zeta}{2}}, \quad \zeta \in [0,\pi].$$

One then recovers Corollary 1.2. in [5] after using appropriate formulas for modified Bessel functions. When p = 3, one has to solve a special cubic equation. To proceed, we rely on results from analytic function theory and the required solution is expressed by means of Gauss hypergeometric functions ([10]) in contrast to Cardan's solution. Therefore, we get a somewhat explicit formula for the series (2), though much more complicated than the one derived for p = 2. The paper is closed with adapting our method to odd dihedral groups, in particular to $D_2(3)$ thereby

²When p = 2, this condition is equivalent to γ is even as stated in [5].

exhausting the list of dihedral groups that are Weyl groups (p = 1 corresponds to the product group \mathbb{Z}_2^2).

2. Proof of the main result

Recall the orthogonality relation for Gegenbauer polynomials ([8]):

$$\int_{0}^{\pi} C_{j}^{\nu}(\cos\zeta) C_{m}^{\nu}(\cos\zeta) (\sin\zeta)^{2\nu} d\zeta = \delta_{jm} \frac{\pi\Gamma(j+2\nu)2^{1-2\nu}}{\Gamma^{2}(\nu)(j+\nu)j!}$$
$$= \delta_{jm} \frac{\pi 2^{1-2\nu}\Gamma(2\nu)}{(j+\nu)\Gamma^{2}(\nu)} C_{j}^{\nu}(1)$$
$$= \delta_{jm} \nu \frac{\sqrt{\pi}\Gamma(\nu+1/2)}{\Gamma(\nu+1)} \frac{C_{j}^{\nu}(1)}{(j+\nu)}$$

where we used $\Gamma(\nu + 1) = \nu \Gamma(\nu)$, the Gauss duplication's formula ([8])

 $\sqrt{\pi}\Gamma(2\nu) = 2^{2\nu-1}\Gamma(\nu)\Gamma(\nu+1/2),$

and the special value ([8])

$$C_j^{\nu}(1) = \frac{(2\nu)_j}{j!}$$

Equivalently, if $\mu^{\nu}(d\cos\zeta)$ is the image of $\mu^{\nu}(d\zeta)$ under the map $\zeta \mapsto \cos\zeta$, then

$$(j+\nu)\int C_j^{\nu}(\cos\zeta)C_m^{\nu}(\cos\zeta)\mu^{\nu}(d\cos\zeta) = \nu C_j^{\nu}(1)\delta_{jm}$$

so that (4) yields

(5)
$$\nu(\pm 1)^{j} \left(\frac{2}{R}\right)^{p\nu} I_{p(j+\nu)}(R) = \int W_{j}^{\nu}(\cos\zeta) f_{\nu,p}^{\pm}(R, \cos\zeta) \mu^{\nu}(d\cos\zeta)$$

where

$$W_i^{\nu}(\cos\zeta) := C_i^{\nu}(\cos\zeta)/C_i^{\nu}(1)$$

is the *j*-th normalized Gegenbauer polynomial. Thus, the *j*-th Gegenbauer-Fourier coefficients of $\zeta \mapsto f_{\nu,p}^{\pm}(R, \cos \zeta)$ are given by

$$\nu(\pm 1)^j \left(\frac{2}{R}\right)^{p\nu} I_{p(j+\nu)}(R), \quad p \ge 2$$

Following [1] p.356, the Mehler's integral representation of W_i^{ν} ([9], p.177)

$$W_{j}^{\nu}(\cos\zeta) = 2^{\nu} \frac{\Gamma(\nu+1/2)}{\Gamma(\nu)\sqrt{\pi}} (\sin\zeta)^{1-2\nu} \int_{0}^{\zeta} [\cos(j+\nu)t] (\cos t - \cos\zeta)^{\nu-1} dt$$

valid for real $\nu > 0$, transforms (5) to

$$\begin{pmatrix} \frac{2}{R} \end{pmatrix}^{p\nu} (\pm 1)^{j} I_{p(j+\nu)}(R) = \frac{2^{\nu}}{\pi} \int_{0}^{\pi} f_{\nu,p}^{\pm}(R, \cos\zeta) \sin\zeta \int_{0}^{\zeta} [\cos(j+\nu)t] (\cos t - \cos\zeta)^{\nu-1} dt d\zeta = \frac{2^{\nu}}{\pi} \int_{0}^{\pi} [\cos(j+\nu)t] \int_{t}^{\pi} f_{\nu,p}^{\pm}(R, \cos\zeta) \sin\zeta (\cos t - \cos\zeta)^{\nu-1} d\zeta dt.$$

The second integral displayed in the RHS of the second equality is known as the Radon transform of $\zeta \mapsto f_{\nu,p}^{\pm}(R, \cos \zeta)$ and inversion formulas already exist ([1]). As a matter of fact, we firstly need to express $(\pm 1)^{j+\nu}I_{p(j+\nu)}$, when $\nu \geq 1$ is an integer, as the Fourier-cosine coefficient of order $j + \nu$ of some function. This is a

consequence of the Lemma below. Secondly, we shall use the appropriate inversion formula for the Radon transform.

Lemma. For any integer $p \ge 1$ and any $t \in [0, \pi]$:

$$2\sum_{j\geq 0} (\pm 1)^j I_{pj}(R) \cos(jt) = I_0(R) + \frac{1}{p} \sum_{s=1}^p e^{\pm R \cos[(t+2\pi s)/p]}.$$

Proof of the Lemma: we will prove the (+) part, the proof of the (-) part follows the same lines with minor modifications. Write

$$2\sum_{j\geq 0} I_{pj}(R)\cos(jt) = \sum_{j\geq 0} I_{pj}(R)[e^{ijt} + e^{-ijt}]$$
$$= I_0(R) + \sum_{j\in\mathbb{Z}} I_{pj}(R)e^{ijt}$$

where used the fact that $I_j(r) = I_{-j}(r), j \ge 0$. Using the identity (3), one obviously gets

$$\sum_{j \in \mathbb{Z}} I_{pj}(R) e^{ijt} = \frac{1}{p} \sum_{s=1}^{p} \sum_{j \in \mathbb{Z}} I_j(R) e^{ij(t+2\pi s)/p}.$$

The (+) part of the Lemma then follows from the generating series for modified Bessel functions ([12]):

$$e^{(z+1/z)R/2} = \sum_{j \in \mathbb{Z}} I_j(R) z^j, \ z \in \mathbb{C}.$$

The Lemma yields

$$I_{pj}(R) = I_0(R)\delta_{j0} + \frac{1}{\pi}\int_0^\pi \cos(jt)\frac{1}{p}\sum_{s=1}^p e^{\pm R\cos[(t+2\pi s)/p]}dt$$

for any integer $j \ge 0$. Assuming that ν is a strictly positive integer, one has

(7)
$$I_{p(j+\nu)}(R) = \frac{1}{\pi} \int_0^{\pi} \cos((j+\nu)t) \frac{1}{p} \sum_{s=1}^p e^{\pm R \cos[(t+2\pi s)/p]} dt.$$

Note that

$$t \mapsto \int_{t}^{\pi} f(R, \cos \zeta) \sin \zeta (\cos t - \cos \zeta)^{\nu - 1} d\zeta$$

as well as

$$t\mapsto \frac{1}{p}\sum_{s=1}^{p}e^{\pm R\cos[(t+2\pi s)/p]}$$

are even functions. This is true since

$$\zeta \mapsto f(R, \cos \zeta) \sin \zeta (\cos t - \cos \zeta)^{\nu - 1}$$

is an odd function so that

$$\int_{-t}^{t} f(R, \cos \zeta) \sin \zeta (\cos t - \cos \zeta)^{\nu - 1} d\zeta = 0,$$

and since

$$\cos[(-t + 2s\pi)/p] = \cos[(t + 2(p - s)\pi)/p]$$

so that one performs the index change $s \to p - s$ and notes that the terms corresponding to s = 0 and s = p are equal. Similar arguments yield the 2π -periodicity

of these functions, therefore, the Fourier-cosine transforms of their restrictions on $(-\pi, \pi)$ coincide with their Fourier transforms on that interval. By injectivity of the Fourier transform and 2π -periodicity,

$$\left(\frac{R}{2}\right)^{p\nu} \int_{t}^{\pi} f_{\nu,p}(R, \cos\zeta) \sin\zeta (\cos t - \cos\zeta)^{\nu-1} d\zeta = \frac{1}{2^{\nu} p} \sum_{s=1}^{p} e^{\pm R \cos[(t+2\pi s)/p]}$$

for all t since both functions are continuous. Finally, the Proposition follows from Theorem 3.1. p.363 in [1]. $\hfill\blacksquare$

Remark. When $\nu = (d-1)/2$ for some integer $d \ge 1$, the Gegenbauer-Fourier transform is interpreted as the Fourier Transform on the sphere S^{d+1} considered as a homogenous space SO(d+1)/SO(d). More precisely, the spherical functions of this space are given by ([1] p.356):

$$W_i^{\nu}\left(\langle z, N \rangle\right), z \in S^{d+1}$$

where $N = (0, \dots, 0, 1) \in S^{d+1}$ is the north pole and $\langle \cdot, \cdot \rangle$ denotes the Euclidian inner product on \mathbb{R}^{d+1} .

Corollary 1. For any integer $\nu \geq 1$

$$\sum_{j\geq 0} (2j+\nu) I_{p(2j+\nu)}(R) C_{2j}^{\nu}(\cos\zeta) = \frac{1}{2^{\nu} \Gamma(\nu)} \left[-\frac{1}{\sin\zeta} \frac{d}{d\zeta} \right]^{\nu} \frac{1}{p} \sum_{s=1}^{p} \cosh\left(R\cos\left[(\zeta+2\pi s)/p\right]\right) + \frac{1}{2^{\nu} \Gamma(\nu)} \left[-\frac{1}{\sin\zeta} \frac{d}{d\zeta} \right]^{\nu} \frac{1}{p} \sum_{s=1}^{p} \left[-\frac{1}{2^{\nu} \Gamma(\nu)} \left[-\frac{1}{2^{\nu$$

3. Weyl group settings p = 2, 3: explicit formulas

3.1. **p=2.** Letting p = 2 and using the fact that $u \mapsto \cosh u$ is an even function, our main result yields

$$\left(\frac{4}{R^2}\right)^{\nu} \sum_{j\geq 0} (2j+\nu) I_{2(2j+\nu)}(R) C_{2j}^{\nu}(\cos\zeta) = \frac{1}{2^{\nu} \Gamma(\nu)} \left[-\frac{4}{R^2 \sin\zeta} \frac{d}{d\zeta}\right]^{\nu} \cosh\left(R\cos(\cdot/2)\right)(\zeta).$$

Noting that

$$-\frac{4}{R^2 \sin \zeta} \frac{d}{d\zeta} \cosh\left(R \cos(\cdot/2)\right)(\zeta) = \frac{1}{R \cos t/2} \frac{d}{dt} \left(u \mapsto \cosh u\right)_{|u|=R \cos(\zeta/2)},$$

after the use of the identity $\sin \zeta = 2 \sin \zeta / 2 \cos \zeta / 2$, it follows that

$$\begin{bmatrix} -\frac{4}{R^2 \sin \zeta} \frac{d}{d\zeta} \end{bmatrix}^{\nu} \cosh\left(R \cos(\cdot/2)\right)(\zeta) = \begin{bmatrix} \frac{1}{u} \frac{d}{du} \end{bmatrix}^{\nu} (u \mapsto \cosh u)_{|u=R \cos(\zeta/2)}$$
$$= \begin{bmatrix} \frac{1}{u} \frac{d}{du} \end{bmatrix}^{\nu-1} (u \mapsto \frac{\sinh u}{u})_{|u=R \cos(\zeta/2)}$$
$$= \sqrt{\frac{\pi}{2}} \begin{bmatrix} \frac{d}{du} \end{bmatrix}^{\nu-1} \left(u \mapsto \frac{I_{1/2}(u)}{\sqrt{u}}\right)_{|u=R \cos(\zeta/2)}$$
$$= \sqrt{\frac{\pi}{2}} \frac{1}{u^{\nu-1/2}} I_{\nu-1/2}(u)_{|u=R \cos(\zeta/2)}$$
$$= \frac{\sqrt{\pi}}{2^{\nu} \Gamma(\nu+1/2)} i_{\nu-1/2}(R \cos(\zeta/2))$$

where the fourth equality is a consequence of the differentiation formula (6) p.79 in [12]. With the help of Gauss duplication's formula, one easily gets:

$$\left(\frac{4}{R^2}\right)^{\nu} \sum_{j \ge 0} (2j+\nu) I_{2(2j+\nu)}(R) C_{2j}^{\nu}(\cos\zeta) = \frac{1}{2\Gamma(2\nu)} i_{\nu-1/2}(R\cos(\zeta/2))$$

and finally recovers Corollary 1.2 in [5] since $c_{2,k}/c(k_1-1/2, k_0-1/2) = \Gamma(2\nu+1)/\nu$.

3.2. **p=3.** The corresponding dihedral group $D_2(6)$ is isomorphic to the Weyl group of type $G_2([2])$. Let $\zeta \in]0, \pi[$ and start with the linearization formula:

$$4\cos^3(\zeta/3) = \cos\zeta + 3\cos(\zeta/3).$$

Thus, we are led to find a root lying in [-1, 1] of the cubic equation

$$Z^{3} - (3/4)Z - (\cos\zeta)/4 = 0$$

for |Z|<1. Set $Z=(\sqrt{-1}/2)T, |T|<2$, the above cubic equation transforms to $T^3+3T-2\sqrt{-1}\cos\zeta=0.$

The obtained cubic equation already showed up in analytic function theory in relation to the local inversion Theorem ([10] p.265-266). Amazingly (compared to Cardan's formulas), its real and both complex roots are expressed through the Gauss Hypergeometric function $_2F_1$. Since we are looking for real $Z = (\sqrt{-1}/2)T$, we shall only consider the complex roots (see the bottom of p. 266 in [10]):

$$T^{\pm} = \pm \sqrt{-1} \left[\sqrt{3} \,_{2}F_{1} \left(-\frac{1}{6}, \frac{1}{6}, \frac{1}{2}; \cos^{2} \zeta \right) - \frac{1}{3} \cos \zeta \,_{2}F_{1} \left(\frac{1}{3}, \frac{2}{3}, \frac{3}{2}; \cos^{2} \zeta \right) \right]$$

so that

$$Z^{\pm} = \pm \left[\frac{\sqrt{3}}{2} {}_{2}F_{1} \left(-\frac{1}{6}, \frac{1}{6}, \frac{1}{2}; \cos^{2} \zeta \right) - \frac{1}{6} \cos \zeta {}_{2}F_{1} \left(\frac{1}{3}, \frac{2}{3}, \frac{3}{2}; \cos^{2} \zeta \right) \right].$$

Since for $\zeta = \pi/2$, $\cos \zeta/3 = \cos \pi/6 = \sqrt{3}/2$, it follows that

$$\cos(\zeta/3) = \left[\frac{\sqrt{3}}{2} {}_{2}F_{1}\left(-\frac{1}{6}, \frac{1}{6}, \frac{1}{2}; \cos^{2}\zeta\right) - \frac{1}{6}\cos\zeta {}_{2}F_{1}\left(\frac{1}{3}, \frac{2}{3}, \frac{3}{2}; \cos^{2}\zeta\right)\right]$$

for all $\zeta \in (0, \pi)$. Now, write $Z = Z(\cos \zeta)$ so that

$$\cos[(\zeta + 2s\pi)/3] = \cos(2s\pi/3)\cos(\zeta/3) - \sin(2s\pi/3)\sqrt{1 - \cos^2(\zeta/3)}$$
$$= \cos(2s\pi/3)Z(\cos\zeta) - \sin(2s\pi/3)\sqrt{1 - Z^2(\cos\zeta)}$$

for any $1 \leq s \leq 3$. It follows that

$$f_{\nu,3}(R,\cos\zeta) = \frac{1}{3\Gamma(\nu)} \left[-\frac{4}{R^3 \sin\zeta} \frac{d}{d\zeta} \right]^{\nu} \sum_{s=1}^3 g_s(RZ(\cos\zeta))$$

where

$$g_s(u) = \cosh\left[\left(\cos(2s\pi/3)u - \sin(2s\pi/3)\sqrt{R^2 - u^2}\right)\right], u \in (-1, 1).$$

Finally,

$$f_{\nu,3}(R,\cos\zeta) = \frac{1}{3\Gamma(\nu)} \left[\frac{4}{R^3} \frac{d}{du}\right]^{\nu} \sum_{s=1}^3 h_s(u)|_{u=\cos\zeta}$$

where $h_s(u) := g_s(RZ(u)), 1 \le s \le 3$. For instance, let $\nu = 1$, then it is not difficult to see that

$$\frac{d}{du}h_s(u)|_{u=\cos\zeta} = \frac{R}{\sin\zeta/3}\frac{dZ}{du}|_{u=\cos\zeta}\sin\left(\frac{\xi+2\pi s}{3}\right)\sinh\left[\sin\left(\frac{\xi+2\pi s}{3}\right)\right]$$

for any $s \in \{1, 2, 3\}$ and the derivative of $u \mapsto Z(u)$ is computed using the differentiation formula for ${}_2F_1$:

$$\frac{d}{du}{}_{2}F_{1}(a,b,c;u) = \frac{ab}{c}{}_{2}F_{1}(a+1,b+1,c+1;u), |u| < 1, c \neq 0$$

As the reader may conclude, formulas are cumbersome compared to the ones derived for p = 2.

4. Odd Dihedral groups

Let $n \geq 3$ be an odd integer. For odd dihedral groups $D_2(n)$, the generalized Bessel function reads ([4] p.157):

$$D_k^W(\rho,\phi,r,\theta) = c_{n,k} \left(\frac{2}{r\rho}\right)^{nk} \sum_{j\ge 0} I_{n(2j+k)}(\rho r) p_j^{-1/2,l_0}(\cos(2n\phi)) p_j^{-1/2,l_0}(\cos(2n\theta))$$

where $k \ge 0, \rho, r \ge 0, \theta, \phi \in [0, \pi/n]$, and

$$c_{n,k} = 2^k \Gamma(nk+1) \frac{\sqrt{\pi} \Gamma(k+1/2)}{\Gamma(k+1)}.$$

In order to adapt our method to those groups, we need to write down the product formula for orthonormal Jacobi polynomials in the limiting case $\alpha = -1/2$ or equivalently $k_1 = 0$. But note that, from an analytic point of view, this generalized Bessel function is obtained from the one associated with even dihedral groups via the substitutions $k_1 = 0, p = n$. Hence one expects the product formula for orthonormal Jacobi polynomials still holds in the limiting case. Indeed, the required limiting formula was derived in [7] p.194 using implicitely the fact that the Beta distribution μ^{α} converges weakly to the dirac mass δ_1 . In order to fit it into our normalizations, we proceed as follows: use the well-known quadratic transformation ([8]):

$$P_j^{-1/2,k-1/2}(1-2\sin^2(n\theta)) = (-1)^j P_j^{k-1/2,-1/2}(2\sin^2(n\theta)-1) = (-1)^j \frac{(1/2)_j}{(k)_j} C_{2j}^k(\sin(n\theta))$$

where $P_j^{\alpha,\beta}$ is the (non orthonormal) *j*-th Jacobi polynomial, together with $\cos(2n\theta) = 1 - 2\sin^2(n\theta)$ to obtain

$$P_j^{-1/2,k-1/2}(\cos(2n\theta))P_j^{-1/2,k-1/2}(\cos(2n\phi)) = \left[\frac{(1/2)_j}{(k)_j}\right]^2 C_{2j}^k(\sin(n\theta))C_{2j}^k(\sin(n\phi)).$$

Now, let k > 0 and recall that the squared L^2 -norm of $P_i^{-1/2,k-1/2}$ is given by ([8])

$$\frac{2^k}{2j+k} \frac{\Gamma(j+1/2)\Gamma(j+k+1/2)}{j!\Gamma(j+k)} = \frac{2^k \sqrt{\pi} \Gamma(k+1/2)}{\Gamma(k)} \frac{(1/2)_j}{(k)_j} \frac{(k+1/2)_j}{(2j+k)j!}$$

Recall also the special value

$$C_{2j}^{k}(1) = \frac{(2k)_{2j}}{(2j)!} = \frac{2^{2k}}{\Gamma(2k)} \frac{\Gamma(k+j)\Gamma(k+j+1/2)}{\Gamma(j+1/2)j!} = 2\frac{(k)_{j}(k+1/2)_{j}}{(1/2)_{j}j!}$$

where we use Gauss duplication formula twice to derive both the second and the third equalities. It follows that

$$\begin{aligned} c(k)p_{j}^{-1/2,k-1/2}(\cos(2n\theta))p_{j}^{-1/2,k-1/2}(\cos(2n\phi)) &= \frac{(1/2)_{j}}{(k)_{j}} \frac{(2j+k)j!}{(k+1/2)_{j}} C_{2j}^{k}(\sin(n\theta)) C_{2j}^{k}(\sin(n\phi)) \\ &= \frac{(2j+k)}{C_{2j}^{k}(1)} C_{2j}^{k}(\sin(n\theta)) C_{2j}^{k}(\sin(n\phi)) \\ &= (2j+k) \int C_{2j}^{k} \left(z_{n\phi,n\theta}(u,1) \right) \mu^{k}(du), \end{aligned}$$

according to [7] p.194, where

$$c(k) := \frac{2^{k+1}\sqrt{\pi}\Gamma(k+1/2)}{\Gamma(k)}.$$

As a matter of fact, we are led again to series of the form

$$\left(\frac{2}{R}\right)^{nk} \sum_{j\geq 0} (2j+k)I_{n(2j+k)}(R)C_{2j}^k(\cos\zeta) = \frac{1}{2}[f_{k,n}^+ + f_{k,n}^-](R,\cos\zeta).$$

5. Two Remarks

The first remark is concerned with $D_2(4)$ which coincides with the B_2 -type Weyl group ([8]). Recall from ([6]) that D_k^W may be expressed through a bivariate hypergeometric function as

$$D_k^W(x,y) = {}_1F_0^{(1/k_1)}\left(\frac{\gamma+1}{2},\frac{x^2}{2},\frac{y^2}{2}\right),$$

where we set $x^2 := (x_1^2, x_2^2) = (\rho^2 \cos^2 \phi, \rho^2 \sin^2 \phi)$ and similarly for y^2 . This series is defined via Jack polynomials:

$${}_{1}F_{0}^{(1/r)}(a,x,y) = \sum_{\tau} (a)_{\tau} \frac{J_{\tau}^{1/r}(x)J_{\tau}^{1/r}(y)}{J_{\tau}^{1/r}(\mathbf{1})|\tau|!}$$

where $\mathbf{1} = (1, 1)$, $\tau = (\tau_1, \tau_2)$ is a partition of length 2, $|\tau| = \tau_1 + \tau_2$ is its weight and $(a)_{\tau}$ is the generalized Pochhammer symbol (see [6] for definitions). But those polynomials, known also as Jack polynomials of type A_1 , may be expressed through Gegenbauer polynomials, a result due to M. Lassalle (see for instance formula 4.10 in [11]):

$$J_{\tau}^{1/r}(x^2) = \frac{(\tau_1 - \tau_2)!}{2^{|\tau|}(r)_{\tau_1 - \tau_2}} \sin^{|\tau|}(2\phi) C_{\tau_1 - \tau_2}^r \left(\frac{1}{\sin(2\phi)}\right)$$

where $(r)_{\tau_1-\tau_2}$ is the (usual) Pochammer symbol. As a matter fact, one wonders if it is possible to come from the hypergeometric series to Corollary 1.2 in [5] and vice-versa.

The second remark comes in the same spirit of the first one. Consider the odd dihedral system $I_2(3) = \{\pm e^{-i\pi/2}e^{i\pi l/3}, 1 \leq l \leq 3\}$ ([8]). It is isomorphic to the A_2 -type root system defined by

$$\{\pm(1,-1,0),\pm(1,0,-1),\pm(0,1,-1)\} \subset \mathbb{R}^3$$

which spans the hyperplane $(1, 1, 1)^{\perp}$. The isomorphism is given by

$$(z_1, z_2, z_3) \mapsto \frac{1}{\sqrt{2}} \left(\sqrt{\frac{3}{2}} z_2, \frac{z_3 - z_1}{\sqrt{2}} \right)$$

subject to $z_1 + z_2 + z_3 = 0$ and for the A_2 -type root system, the generalized Bessel function is given by the trivariate hypergeometric series ${}_0F_0^{(1/k)}$ (see [6] for the definition). Is it possible to relate this function to

$$\frac{c_{3,k}}{c(k)}\int [f_{k,3}^+ + f_{k,3}^-](\rho r, z_{3\phi,3\theta}(u,1))\mu^k(du) = \frac{3\Gamma(3k)}{4}\int [f_{k,3}^+ + f_{k,3}^-](\rho r, z_{3\phi,3\theta}(u,1))\mu^k(du)$$

in the same way the ${}_{0}F_{1}^{1/k_{1}}$ is related to the integral representation derived for p = 2?

Acknowledgment: the author is grateful to Professor C.F. Dunkl who made him aware of the hypergeometric formulas for the roots of the cubic equation.

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