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Rate of decay to 0 of the solutions to a nonlinear parabolic equation

I.Ben Arbi¹

Résumé: On étudie l'ordre de convergence vers 0, quand $t \rightarrow +\infty$ de la solution de l'équation $\psi_t - \Delta\psi + |\psi|^{p-1}\psi = 0$ avec les conditions aux limites de Neumann dans un ouvert connexe borné de \mathbb{R}^n où $p > 1$. On montre que soit $\psi(t, \cdot)$ converge vers 0 exponentiellement, soit $\psi(t, \cdot)$ décroît comme $t^{-\frac{1}{(p-1)}}$.

Abstract: We study the decay rate to 0, as $t \rightarrow +\infty$ of the solution of equation $\psi_t - \Delta\psi + |\psi|^{p-1}\psi = 0$ with Neumann boundary conditions in a bounded smooth open connected domain of \mathbb{R}^n where $p > 1$. We show that either $\psi(t, \cdot)$ converges to 0 exponentially fast or $\psi(t, \cdot)$ decreases like $t^{-\frac{1}{(p-1)}}$.

Keywords: rate of decay, parabolic equation

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1 Introduction and main results.

In this paper we consider the following nonlinear parabolic equation

$$\begin{cases} \psi_t - \Delta\psi + g(\psi) = 0 & \text{in } \mathbb{R}^+ \times \Omega, \\ \frac{\partial\psi}{\partial n} = 0 & \text{on } \mathbb{R}^+ \times \partial\Omega. \end{cases} \quad (1.1)$$

where Ω is a bounded smooth open connected domain of \mathbb{R}^n and $g \in C^1(\mathbb{R})$ satisfies

$$g(0) = 0 \quad (1.2)$$

and for some $p > 1$

$$\exists c > 0, \forall s \in \mathbb{R}, \quad 0 \leq g'(s) \leq c|s|^{p-1}. \quad (1.3)$$

From (1.2)-(1.3) we deduce that $g(s)$ has the sign of s and

$$\forall s \in \mathbb{R}, \quad |g(s)| \leq \frac{c}{p}|s|^p. \quad (1.4)$$

We define the operator A by

$$D(A) = \{\psi \in H^2(\Omega), \frac{\partial\psi}{\partial n} = 0 \text{ on } \partial\Omega\}$$

and

$$\forall \psi \in D(A), \quad A\psi = -\Delta\psi$$

On the other hand the operator B defined by

$$D(B) = \{\psi \in L^2(\Omega), -\Delta\psi + g(\psi) \in L^2(\Omega) \text{ and } \frac{\partial\psi}{\partial n} = 0 \text{ on } \partial\Omega\}$$

and

$$\forall \psi \in D(B), \quad B\psi = -\Delta\psi + g(\psi)$$

is well-known to be maximal monotone in $L^2(\Omega)$. As a consequence of [1, 2] for any $\psi_0 \in L^2(\Omega)$ there exists a unique weak solution of the equation

$$\psi' + B\psi = 0 \text{ on } \mathbb{R}^+; \quad \psi(0, x) = \psi_0. \quad (1.5)$$

In addition it is well known that if $\psi_0 \in L^\infty(\Omega)$, $\psi(t, \cdot)$ remains in $L^\infty(\Omega)$ for all $t > 0$. Finally [8] contains an estimate of the solution in $C(\bar{\Omega})$ and $C^1(\bar{\Omega})$ for $t > 0$, which is valid for any sufficiently regular domain.

Concerning the behaviour for t large, in [5], A.Haraux established in the case of a pure power nonlinearity the exponential convergence to 0 of the projection on the range of A of the solution of equation (1.1). Moreover in [4], the study of the equation $u'' + u' - \Delta u + g(u) = 0$ with Neumann boundary conditions and where g satisfies $\exists C, c > 0, \forall s \in \mathbb{R}, c|s|^{p-1} \leq g'(s) \leq C|s|^{p-1}$ for some $p > 1$, showed that either $u(t)$ converges to 0 exponentially fast, or $\|u(t)\|_{H_0^1(\Omega)} \geq \gamma t^{-1/(p-1)}$ with $\gamma > 1$ for $t \geq 1$.

Several authors have treated some variants of equation (1.1), for example in [6] with $g(u) = c|u|^{p-1}u - \lambda_1 u$ and with Dirichlet boundary conditions, was studied the decay rate at the infinity of solutions to (1.1), where $\lambda_1 > 0$ is the first eigenvalue of $-\Delta$ in $H_0^1(\Omega)$. The result obtained there is optimal for positive solutions.

According to La Salle's invariance principle, cf. [3, 7], any solution ψ of (1.1), having a precompact range on \mathbb{R}^+ with values to $L^\infty(\Omega)$, converges to a continuum of stationary solutions of equation (1.1), which reduces here to the constants of some sub-interval of $g^{-1}(0)$. By monotonicity, it is in fact known that ψ converges to some constant $a \in g^{-1}(0)$.

Our first result is valid without any additional hypothesis on g

Theorem 1.1. *Let g satisfy (1.2) and (1.3). Then any solution ψ of (1.1) satisfies the following alternative as $t \rightarrow \infty$: either*

$$\|\psi(t, \cdot)\|_\infty \leq C e^{-\lambda_2 t}, \quad (1.6)$$

or

$$\exists c' > 0, \forall t \geq 1, \quad \left| \int_\Omega \psi(t, x) dx \right| \geq c' t^{-\frac{1}{p-1}}, \quad (1.7)$$

where $\lambda_2 > 0$ is the second eigenvalue of A in $D(A)$.

Our second result provides a more accurate estimate when $g(\psi) = |\psi|^{p-1}\psi$

Theorem 1.2. *Let us consider the nonlinear parabolic problem (1.1) with $g(\psi) = |\psi|^{p-1}\psi$, then any solution ψ of (1.1) satisfies the following alternative as $t \rightarrow \infty$: either*

$$\|\psi(t, \cdot)\|_\infty \leq C e^{-\lambda_2 t}, \quad (1.8)$$

or

$$\forall t \geq 1, \quad \left\| |\psi(t, \cdot)| - ((p-1)t)^{-\frac{1}{p-1}} \right\|_\infty \leq K t^{-\frac{1}{(p-1)^2}}, \quad (1.9)$$

where $K, C > 0$, $p > 1$ and $\lambda_2 > 0$ is the second eigenvalue of A in $D(A)$.

In the following Proposition, we consider two special cases showing that both possibilities in the second result in the Theorem 1.1 can actually happen.

Proposition 1.3. *Let g satisfy (1.2) and (1.3). Then we have*

- (i) *If Ω is symmetric around 0, g is odd and $\psi(0, \cdot)$ is an odd function in Ω , then any solution of (1.1) satisfies (1.6).*
- (ii) *If $\psi(0, \cdot) \geq 0$ and ψ does not vanish a.e in Ω , then any solution of (1.1) satisfies (1.7).*

Finally, our last result shows that the second possibility is sharp for a class of functions g more general than the pure power

Proposition 1.4. *Under the additional hypothesis*

$$\exists k_1 > 0, \forall s \in \mathbb{R}, |g(s)| \geq k_1 |s|^p \quad (1.10)$$

for any solution ψ of (1.1), we have

$$\forall t \geq 1, \quad \|\psi(t, \cdot)\|_\infty \leq \left\{ \frac{1}{k_1(p-1)} \right\}^{\frac{1}{p-1}} t^{-\frac{1}{p-1}} \quad (1.11)$$

2 Proof of Proposition 1.4.

Proof. If $\psi(0, \cdot) = 0$, we have $\psi(t, \cdot) \equiv 0$ and the result is obvious. Otherwise let z be defined by

$$z(t) = \left\{ \frac{1}{\|\psi(0, \cdot)\|_\infty^{1-p} + k_1(p-1)t} \right\}^{\frac{1}{p-1}}$$

Then z is a solution of the following nonlinear ODE problem

$$\begin{cases} z' + k_1 z^p = 0, \\ z(0) = \|\psi(0, \cdot)\|_\infty. \end{cases} \quad (2.1)$$

Under the additional condition(1.10), we will show that z is a super solution of (1.1). Indeed, we have

$$\begin{aligned} z_t - \Delta z + g(z) &= -k_1(z(0)^{1-p} + k_1(p-1)t)^{-\frac{p}{p-1}} + g(z) \\ &\geq -k_1(z(0)^{1-p} + k_1(p-1)t)^{-\frac{p}{p-1}} + k_1|z|^p \\ &\geq 0. \end{aligned}$$

Since $\psi(0, \cdot) \leq z(0)$ we deduce, by standard comparison principle, that $\psi(t, \cdot) \leq z(t) \forall t \geq 1$.

A similar calculation shows that $\psi(t, \cdot) \geq -z(t), \quad \forall t \geq 1$ which concludes the proof. \blacksquare

3 A general result on the range component.

Defining the orthogonal projection $P : H \longrightarrow N$, where

$$H = L^2(\Omega), N = \ker(A) \text{ and } P\psi(t, \cdot) = \frac{1}{|\Omega|} \int_{\Omega} \psi(t, x) dx,$$

as already mentioned in the introduction, [5] showed that for $g(\psi) = |\psi|^{p-1}\psi$, the following estimate holds

$$\|\psi(t) - P\psi(t)\|_{L^2(\Omega)} \leq K e^{-\lambda_2 t},$$

for some constant $K > 0$. In this section, we will show that we have the same result for any function g satisfying (1.3).

Proposition 3.1. *Let $\psi \in C(\mathbb{R}^+, L^\infty)$ be any solution of (1.1). Assume that g is a locally Lipschitz non-decreasing function. Then we have ,*

$$\|\psi(t) - P\psi(t)\|_2 \leq \|\psi(0) - P\psi(0)\|_2 e^{-\lambda_2 t}, \quad (3.1)$$

where $\|\cdot\|_2$ denotes the norm in $L^2(\Omega)$ and $\lambda_2 > 0$ is the second eigenvalue of $-\Delta$ in $L^2(\Omega)$ with Neumann boundary conditions .

Proof. We denote by (u, v) the inner product of two functions u, v of $L^2(\Omega)$. Since g is a nondecreasing function, for all $\psi \in L^\infty(\Omega)$, we have a.e. in $x \in \Omega$

$$(g(\psi) - g(P\psi))(\psi - P\psi) \geq 0$$

and then by integrating over Ω

$$(g(\psi), \psi - P\psi) - (g(P\psi), \psi - P\psi) \geq 0. \quad (3.2)$$

Since $g(P\psi)$ is a constant and $(\psi - P\psi) \in N^\perp$, we deduce that $(g(P\psi), \psi - P\psi) = 0$. Hence from (3.2),

$$(g(\psi), \psi - P\psi) \geq 0.$$

Setting

$$w = \psi - P\psi,$$

we have since $\Delta P\psi = P\Delta\psi = 0$

$$w' - \Delta w = \psi' - \Delta\psi - P\psi' + \Delta P\psi = (I - P)(\psi' - \Delta\psi) = -(I - P)g(\psi).$$

Thus, since $(w, (I - P)g(\psi)) = ((I - P)w, g(\psi)) = (w, g(\psi)) = (g(\psi), \psi - P\psi)$ we find

$$\frac{1}{2} \frac{d}{dt} \|w(t)\|_2^2 = (w, \Delta w) - (g(\psi), \psi - P\psi) \leq -\lambda_2 |w|^2.$$

By integrating we obtain (3.1). ■

4 Proof of Theorem 1.1.

We set $\psi = u + w$, where $u = P\psi$ and $w = (I - P)\psi$. By projecting (1.1) on N we obtain

$$u' + P(g(\psi)) = 0, \tag{4.1}$$

where we have used that $P(A\psi) = 0$, since $R(A) \subset N^\perp$. Noticing that

$$u' + P(g(u)) + P(g(\psi) - g(u)) = u' + g(u) + P(g(\psi) - g(u)),$$

we can rewrite the equation (4.1) as

$$u' + g(u) = -P(g(\psi) - g(u)).$$

By the assumption (1.3), we deduce that

$$|P(g(\psi) - g(u))| \leq \frac{1}{|\Omega|} \|g(\psi) - g(u)\|_1 \leq \frac{c}{|\Omega|} (\|\psi\|_{2p-2}^{p-1} + \|u\|_{2p-2}^{p-1}) \|w\|_2.$$

But ψ and u are uniformly bounded and from Proposition 3.1 we have the estimate $\|w(t)\|_2 \leq K e^{-\lambda_2 t}$. Therefore

$$|P(g(\psi) - g(u))| \leq K' e^{-\lambda_2 t},$$

with $K' > 0$. Which leads us to study the equation:

$$u' + g(u) = f(t) \quad \text{in} \quad \mathbb{R}^+, \tag{4.2}$$

where

$$f(t) = P(g(\psi) - g(u))$$

and

$$|f(t)| \leq K' e^{-\lambda_2 t}.$$

Using the same method as in [5], we show the following result:

Lemma 4.1. *Let $c > 0$, $\gamma > 0$, $p > 1$ and g satisfying (1.2) and (1.3) Let $M > 0$ such that*

$$M \leq \left(\frac{\gamma}{2c}\right)^{\frac{1}{p-1}}, \quad (4.3)$$

$c_1 > 0$ with

$$c_1 \leq \frac{\gamma}{2}M,$$

then for every function f satisfying

$$|f(t)| \leq c_1 e^{-\gamma t},$$

there exists a unique function $v \in C^1(\mathbb{R}^+)$ satisfying

$$\forall t \geq 0, \quad v' + g(v) = f(t) \quad (4.4)$$

and

$$\sup_{t \in (0, +\infty)} \{e^{\gamma t} |v(t)|\} \leq M. \quad (4.5)$$

Proof. Since any solution of (4.4).(4.5) satisfies the integral equation

$$v(t) = - \int_t^{+\infty} (f(s) - g(v(s))) ds. \quad (4.6)$$

We look for a solution of (4.6). It is then natural to introduce the following function space :

$$X = \{v \in C(0, +\infty); \sup_{t \in (0, +\infty)} e^{\gamma t} |v(t)| \leq M\},$$

equipped with the distance associated to the norm

$$\|v\|_\gamma = \sup_{t \in (0, +\infty)} e^{\gamma t} |v(t)|. \quad (4.7)$$

We consider the operator $\mathcal{T} : X \rightarrow C(0, +\infty)$ defined by

$$\mathcal{T}v(t) = - \int_t^{+\infty} (f(s) - g(v(s))) ds.$$

From (1.4), we have the estimate

$$\forall s \in \mathbb{R}^+; \quad |g(v(s))| \leq \frac{c}{p} |v(s)|^p.$$

First we will show that $\mathcal{T}(X) \subset X$. Let $v \in X$, then for all $t \geq 0$,

$$\begin{aligned}
|\mathcal{T}v(t)| &\leq \int_t^{+\infty} |f(s)|ds + \int_t^{+\infty} |g(v(s))|ds \\
&\leq \int_t^{+\infty} |f(s)|ds + \frac{c}{p} \int_t^{+\infty} |v(s)|^p ds \\
&\leq \frac{c_1}{\gamma} e^{-\gamma t} + \frac{c}{p} M^p \int_t^{+\infty} e^{-p\gamma s} ds \\
&\leq \left(\frac{c_1}{\gamma} + \frac{cM^p}{p^2\gamma} \right) e^{-\gamma t} \\
&\leq \left(\frac{M}{2} + Mc \frac{M^{p-1}}{p^2\gamma} \right) e^{-\gamma t} \\
&\leq \left(\frac{M}{2} + \frac{M}{2p^2} \right) e^{-\gamma t}.
\end{aligned}$$

Since $p > 1$, it follows that

$$|\mathcal{T}v(t)| \leq M e^{-\gamma t}.$$

Hence by (4.5), we obtain that $\mathcal{T}v \in X$, with

$$\|\mathcal{T}v(t)\|_\gamma \leq M.$$

Secondly, we will prove that \mathcal{T} is a contraction on X . In fact, for $x, \bar{x} \in X$ and for all $t \geq 0$

$$\begin{aligned}
|\mathcal{T}x(t) - \mathcal{T}\bar{x}(t)| &\leq \int_t^{+\infty} |g(x(s)) - g(\bar{x}(s))|ds \\
&\leq cM^{p-1} \int_t^{+\infty} e^{-p\gamma s} e^{\gamma s} |x(s) - \bar{x}(s)| ds \\
&\leq \frac{cM^{p-1}}{p\gamma} \|x - \bar{x}\|_\gamma e^{-\gamma t}.
\end{aligned}$$

Then we have

$$|\mathcal{T}x(t) - \mathcal{T}\bar{x}(t)| e^{\gamma t} \leq \frac{cM^{p-1}}{p\gamma} \|x - \bar{x}\|_\gamma.$$

Therefore, since M^{p-1} satisfies (4.3), we conclude that $\forall x, \bar{x} \in X$

$$\|\mathcal{T}x - \mathcal{T}\bar{x}\|_\gamma \leq \frac{1}{2} \|x - \bar{x}\|_\gamma.$$

Thus \mathcal{T} is a $\frac{1}{2}$ -Lipschitz functional on the complete metric space X and the result follows from the Banach fixed point theorem. From (4.6) it follows easily that v satisfies (4.4). Then the uniqueness of v follows from the uniqueness of the solution of (4.6) (\mathcal{T} is a contraction) and the fact that any solution of (4.4) satisfies (4.6). The existence comes from the fact that conversely any solution of (4.6) satisfies (4.4). \blacksquare

Proof of Theorem 1.1 continued. Consequently, we have a solution v that satisfies the equation (4.4) and for all $t \geq T_0$

$$|v(t)| \leq Me^{-\lambda_2 t}, \quad (4.8)$$

where $M = M'e^{\gamma T_0}$ and $M' > 0$. If we subtract (4.4) from (4.2) we obtain

$$(u - v)' + g(u) - g(v) = 0.$$

Setting $z = u - v$, we complete the proof analyzing two cases.

1st case: If $z(T_0) = 0$, then for all $t \geq T_0$, $z(t) = 0$. Hence $u \equiv v$ and from (4.8) it follows that

$$|u(t)| \leq Me^{-\lambda_2 t}.$$

Then, using (3.1), we obtain

$$\|\psi(t)\|_2 \leq M'e^{-\lambda_2 t}.$$

Finally by reasoning as in [5], [6] we obtain (1.6).

2nd case: If $z(T_0) \neq 0$ then $\forall t \geq T_0$, $z(t) \neq 0$ and we have

$$z'(t) + \frac{g(u(t)) - g(v(t))}{u(t) - v(t)} z(t) = 0.$$

Since g is a monotonic function,

$$\alpha(t) := \frac{g(u(t)) - g(v(t))}{u(t) - v(t)}$$

is a strictly positive function. Moreover, $\exists \theta \in]0, 1[$,

$$\alpha(t) = g'(\theta u(t) + (1 - \theta)v(t)) \leq c|\theta u(t) + (1 - \theta)v(t)|^{p-1} \quad (4.9)$$

We distinguish 2 cases:

– $p > 2$ then by convexity of the p th power we have

$$|\theta u(t) + (1 - \theta)v(t)|^{p-1} \leq \theta|u|^{p-1} + (1 - \theta)|v|^{p-1} \leq |u|^{p-1} + |v|^{p-1}.$$

– $1 < p < 2$ we study the function $(x+y)^a - x^a$ for $0 < a < 1$ and $x, y > 0$ we prove that $X \rightarrow (1 + X)^a - X^a$ is a decreasing function on $(0, +\infty)$, we deduce $0 < (1 + X)^a - X^a < 1$ and it follows by homogeneity that $(x + y)^a < x^a + y^a$ by letting $X = \frac{x}{y}$. Then we conclude that $|\theta u(t) + (1 - \theta)v(t)|^{p-1} \leq |u|^{p-1} + |v|^{p-1}$. Consequently we obtain

$$\forall p > 1, \quad \alpha(t) \leq c(|u|^{p-1} + |v|^{p-1}). \quad (4.10)$$

Setting $y = |z|$, we have that

$$y' + \alpha(t)y \geq 0.$$

Then the estimate (4.10) implies that

$$y' \geq -c(|u|^{p-1} + |v|^{p-1})y \geq -c(|z + v|^{p-1} + |v|^{p-1})y.$$

Hence there exists some constants $c_2, c_3 > 0$ such that

$$y' \geq -c_2(|z|^{p-1})y - c_3|v|^{p-1}y.$$

Since $y = |z|$,

$$y' \geq -c_2y^p - c_3|v|^{p-1}y.$$

Putting $a(t) = c_3|v|^{p-1}$, we deduce that

$$y' + a(t)y \geq -c_2y^p. \quad (4.11)$$

We set

$$A(t) = -c_3 \int_t^{+\infty} |v|^{p-1} ds, \quad \omega(t) = e^{A(t)}y$$

and by replacing ω in (4.11), we obtain

$$\omega(t) \geq \left\{ \frac{1}{\omega(0)^{1-p} + (p-1)c_4 t} \right\}^{\frac{1}{p-1}}$$

with $c_4 = \left(\int_0^{+\infty} a(s) ds \right)^{-(p-1)}$. Then for t large enough we have

$$\omega(t) \geq Kt^{-\frac{1}{p-1}}.$$

Since $t \rightarrow e^{A(t)}$ is a bounded function, we conclude (1.7) by observing that $u = z + v$ and v tends to 0 exponentially at infinity. ■

5 Proof of Theorem 1.2.

Considering $g(\psi) = |\psi|^{p-1}\psi$, g satisfies (1.2) and (1.3) with $c = p$. Hence Lemma 4.1 is applicable with $c = p$, therefore we assume

$$M \leq \left(\frac{\lambda_2}{2p}\right)^{\frac{1}{p-1}}.$$

In (4.2), we replace $g(\psi) = |\psi|^{p-1}\psi$, we can subtract (4.4) from (4.2), we deduce:

$$(u - v)' + |u|^{p-1}u - |v|^{p-1}v = 0. \quad (5.1)$$

We will study two cases.

1st case: If $z(T_0) = 0$ then for all $t \geq T_0$, $z(t) = 0$. Hence $u \equiv v$ and from (4.8) it follows that

$$|u(t)| \leq Me^{-\lambda_2 t}.$$

Moreover, using (3.1), we obtain (1.8).

2nd case: if $z(T_0) \neq 0$ then for all $t \geq T_0$ $z(t) \neq 0$ and we have

$$z'(t) + |u(t)|^{p-1}u(t) - |v(t)|^{p-1}v(t) = z'(t) + \alpha(t)z(t) = 0.$$

with

$$\begin{aligned} \alpha(t) &= \frac{|u(t)|^{p-1}u(t) - |v(t)|^{p-1}v(t)}{u(t) - v(t)} \\ &= \frac{|z(t) + v(t)|^{p-1}(z(t) + v(t)) - |v(t)|^{p-1}v(t)}{z(t)} \\ &= \frac{|z(t)|^{p-1}\left|1 + \frac{v(t)}{z(t)}\right|^{p-1}(z(t) + v(t)) - |v(t)|^{p-1}v(t)}{z(t)}. \end{aligned}$$

In that case $\alpha(t) > 0$, indeed $t \mapsto |u(t)|^{p-1}u(t)$ is non decreasing function. By applying Taylor's formula, we obtain

$$\alpha(t) = |z(t)|^{p-1} + \beta(t), \quad (5.2)$$

with $|\beta(t)| \leq Be^{-\eta t}$, where $\eta > 0$ is any positive number smaller than $(p-1)\lambda_2$ and $B > 0$. Replacing α by its expression in (5.2), equation (5.1) becomes

$$z' + |z|^{p-1}z + \beta(t)z = 0. \quad (5.3)$$

Let $y = |z|$, we obtain

$$y' + y^p + \beta(t)y = 0. \quad (5.4)$$

Setting $\xi(t) = e^{A(t)}y(t)$, with $A(t) = -\int_t^{+\infty} \beta(s)ds$ we find

$$|A(t)| \leq \int_t^{+\infty} |\beta(s)|ds \leq \frac{B}{\eta}e^{-\eta t}. \quad (5.5)$$

By Taylor's formula we have

$$\forall h \in [-1, 1], \quad |e^h - 1| \leq 2|h|, \quad (5.6)$$

we can give the following estimate

$$|\xi(t) - y(t)| \leq y(t)|e^{A(t)} - 1| \leq 2|A(t)|y(t) \leq \frac{2B}{\eta}\|y\|_{\infty}e^{-\eta t},$$

we conclude

$$|\xi(t) - y(t)| \leq ke^{-\delta t}, \quad (5.7)$$

where $\delta = \eta$ and $k = \frac{2B}{\eta}\|y\|_{\infty}$.

Replacing $\xi(t)$ in (5.3), we have

$$\begin{aligned} -e^{-A(t)}\xi'(t) &= e^{-pA(t)}\xi^p(t) \\ \Leftrightarrow -\frac{\xi'(t)}{\xi^p(t)} &= e^{(1-p)A(t)} \\ \Leftrightarrow \left(\frac{1}{\xi^{p-1}(t)}\right)' &= (p-1)e^{(1-p)A(t)}. \end{aligned}$$

$t \mapsto e^{-(p-1)A(t)}$ is bounded and tends to 1 at infinity, ξ is given by

$$\xi(t)^{p-1} = \frac{1}{\xi(t_0)^{1-p} + (p-1)\int_{t_0}^t e^{-(p-1)A(s)}ds}.$$

we set

$$D(t) = \xi(t_0)^{1-p} + (p-1)\int_{t_0}^t e^{-(p-1)A(s)}ds \quad \text{and} \quad h(t) = -(p-1)A(t),$$

then we show that $D(t) - (p-1)t$ is bounded.

From (5.5), we know that $t \mapsto h(t)$ is an integrable function, then by (5.6), $(e^{-(p-1)A(t)} - 1)$ is also integrable. In order to show (1.9), we proceed

as follows

$$\begin{aligned} |D(t) - (p-1)t| &= |\xi(t_0)^{1-p} + (p-1) \int_{t_0}^t e^{-(p-1)A(s)} ds - (p-1)t| \\ &= |\xi(t_0)^{1-p} + (p-1) \int_{t_0}^t (e^{-(p-1)A(s)} - 1) ds - (p-1)t_0| \end{aligned}$$

Using (5.6) we obtain

$$\begin{aligned} |D(t) - (p-1)t| &\leq |K| + 2(p-1)^2 \int_{t_0}^t |A(s)| ds \\ &\leq M_1 = |K| + 2(p-1)^2 \frac{B}{\eta^2} \end{aligned}$$

with $K = \xi(t_0)^{1-p} - (p-1)t_0$. Setting $d(t) = D(t) - (p-1)t$, we obtain

$$\begin{aligned} \left| \xi(t) - \left(\frac{1}{(p-1)t} \right)^{\frac{1}{p-1}} \right| &= \left| \left(\frac{1}{D(t)} \right)^{\frac{1}{p-1}} - \left(\frac{1}{(p-1)t} \right)^{\frac{1}{p-1}} \right| \\ &= \left| \left(\frac{1}{D(t) - (p-1)t + (p-1)t} \right)^{\frac{1}{p-1}} - \left(\frac{1}{(p-1)t} \right)^{\frac{1}{p-1}} \right| \\ &= \left| \left(\frac{1}{d(t) + (p-1)t} \right)^{\frac{1}{p-1}} - \left(\frac{1}{(p-1)t} \right)^{\frac{1}{p-1}} \right| \\ &= \left| \left(\frac{1}{(p-1)t} \right)^{\frac{1}{p-1}} \left(\frac{1}{\frac{d(t)}{(p-1)t} + 1} \right)^{\frac{1}{p-1}} - \left(\frac{1}{(p-1)t} \right)^{\frac{1}{p-1}} \right| \\ &= \left(\frac{1}{(p-1)t} \right)^{\frac{1}{p-1}} \left| \left(\frac{d(t)}{(p-1)t} + 1 \right)^{-\frac{1}{p-1}} - 1 \right| \\ &= \left(\frac{1}{(p-1)t} \right)^{\frac{1}{p-1}} \left(1 - \left(\frac{d(t)}{(p-1)t} + 1 \right)^{-\frac{1}{p-1}} \right). \end{aligned}$$

Since $\left| \frac{d(t)}{(p-1)t} \right| < 1$ for t large enough. Let $\eta(t) = \left(1 + \frac{d(t)}{(p-1)t} \right)^{-\frac{1}{p-1}}$, by the mean value Theorem and if we suppose that $\left| \frac{d(t)}{(p-1)t} \right| \leq \frac{1}{2}$, we obtain

$$|\eta'(t)| \leq \frac{1}{p-1} \times 2^{1+\frac{1}{p-1}}.$$

Therefore

$$\left| 1 - \left(1 + \frac{d(t)}{(p-1)t} \right)^{-\frac{1}{(p-1)}} \right| \leq \frac{1}{p-1} \times 2^{1+\frac{1}{p-1}} \left| \frac{d(t)}{(p-1)t} \right|.$$

As we have seen above, $d(t)$ is bounded by M_1 then we conclude

$$\left| \xi(t) - \left(\frac{1}{(p-1)t} \right)^{\frac{1}{(p-1)}} \right| \leq Ct^{-1-\frac{1}{p-1}}.$$

With $C = \left(\frac{1}{p-1}\right)^{\frac{p}{p-1}} \times 2^{1+\frac{1}{p-1}} M_1$.

We recall that z has a constant sign on $[T_0, +\infty[$ and z and u have the same sign. As in Section 4, we set $u = v + z$ and $\psi = u + w$, we distinguish two cases

- If $z > 0$, this implies that for t large enough $u > 0$ and $|\psi| = \psi$.

Then

$$\left| |\psi| - \left(\frac{1}{(p-1)t} \right)^{\frac{1}{(p-1)}} \right| \leq \left| u - \left(\frac{1}{(p-1)t} \right)^{\frac{1}{(p-1)}} \right| + |w| \leq Kt^{-1-\frac{1}{p-1}}.$$

Indeed,

$$\begin{aligned} \left| u - \left(\frac{1}{(p-1)t} \right)^{\frac{1}{(p-1)}} \right| &\leq |u - z| + |z - \xi| + \left| \xi - \left(\frac{1}{(p-1)t} \right)^{\frac{1}{(p-1)}} \right| \\ &\leq Me^{-\lambda_2 t} + ke^{-\delta t} + Ct^{-1-\frac{1}{p-1}}. \end{aligned}$$

Since we have (5.7), we obtain (1.9).

- If we suppose that $z < 0$, then implies that $u < 0$. By similar calculations we obtain the same result.

Indeed, $|\psi| = -\psi$, and we have

$$\begin{aligned} \left| |\psi| - \left(\frac{1}{(p-1)t} \right)^{\frac{1}{(p-1)}} \right| &\leq \left| -u - \left(\frac{1}{(p-1)t} \right)^{\frac{1}{(p-1)}} \right| + |w| \\ &\leq Kt^{-1-\frac{1}{p-1}} \end{aligned}$$

since

$$\begin{aligned} \left| -u - \left(\frac{1}{(p-1)t} \right)^{\frac{1}{(p-1)}} \right| &\leq |u - z| + |-z - \xi| + \left| \xi - \left(\frac{1}{(p-1)t} \right)^{\frac{1}{(p-1)}} \right| \\ &\leq Me^{-\lambda_2 t} + ke^{-\delta t} + Ct^{-1-\frac{1}{p-1}}. \end{aligned}$$

Also since we have (5.7), finally we obtain (1.9).

6 Proof of Proposition 1.3

(i) If ψ is an odd function, then $\psi(0, -x) = -\psi(0, x)$. It implies for all $t > 0$, $\psi(t, -x) = -\psi(t, x)$.

In that case

$$u(t) = P\psi(t, x) = \frac{1}{|\Omega|} \int_{\Omega} \psi(t, x) dx = 0.$$

Moreover, we know that $\psi = u + w$, where $w(t) \in N^\perp$ for all $t \geq 0$ and (cf. Proposition 3.1) we have

$$\|w(t)\|_{L^\infty(\Omega)} \leq Ke^{-\lambda_2 t},$$

hence the solution ψ satisfies (1.8).

(ii) If $\psi(0, x) \geq 0$ and ψ does not vanish a.e in Ω , then for all $t \geq 0$ $\psi(t, x) > 0$, it implies that

$$\int_{\Omega} \psi(t, x) dx > 0. \quad (6.1)$$

We suppose that we have $\|\psi(t, \cdot)\|_\infty \leq Ce^{-\lambda_2 t}$ and we consider the problem (1.1) then we integrate on Ω , we obtain

$$\int_{\Omega} \psi_t(t, x) dx = - \int_{\Omega} g(\psi(t, x)) dx. \quad (6.2)$$

An elementary calculation shows that we have

$$\begin{aligned} \int_{\Omega} g(\psi(t, x)) dx &\leq \frac{c}{p} \int_{\Omega} |\psi(t, x)|^p dx \\ &\leq \frac{c}{p} \int_{\Omega} |\psi(t, x)|^{p-1} \psi(t, x) dx \\ &\leq \frac{c}{p} \int_{\Omega} \|\psi(t, x)\|_\infty^{p-1} \psi(t, x) dx \\ &\leq \frac{c}{p} C^{p-1} e^{-(p-1)\lambda_2 t} \int_{\Omega} \psi(t, x) dx \end{aligned}$$

Now we set $y(t) = \int_{\Omega} \psi(t, x) dx$. From (6.2), we deduce

$$y'(t) \geq -Me^{-\delta t} y(t) \quad (6.3)$$

with $M = \frac{c}{p} C^{p-1}$ and $\delta = (p-1)\lambda_2$.

Since $y(t) > 0$ by (6.1), we can integrate in the interval $[0, t]$, we obtain

$$y(t) \geq y(0) \exp \left\{ -M \int_0^t e^{-\delta s} ds \right\} \geq y(0) \exp \left\{ -\frac{M}{\delta} \right\} > 0. \quad (6.4)$$

Hence y does not tend to 0 for t large, this contradicts our hypothesis and we conclude that y satisfies (1.7).

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