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Radial symmetry and symmetry breaking for some interpolation inequalities

Jean Dolbeault · Maria J. Esteban ·
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Abstract We analyze the radial symmetry of extremals for a class of interpolation inequalities known as Caffarelli-Kohn-Nirenberg inequalities, and for a class of weighted logarithmic Hardy inequalities which appear as limiting cases of the first ones. In both classes we show that there exists a continuous surface that splits the set of admissible parameters into a region where extremals are symmetric and a region where symmetry breaking occurs. In previous results, the symmetry breaking region was identified by showing the linear instability of the radial extremals. Here we prove that symmetry can be broken even within the set of parameters where radial extremals correspond to local minima for the variational problem associated with the inequality. For interpolation inequalities, such a symmetry breaking phenomenon is entirely new.

Keywords Sobolev spaces · interpolation · Hardy-Sobolev inequality · Caffarelli-Kohn-Nirenberg inequality · logarithmic Hardy inequality · Gagliardo-Nirenberg inequality · logarithmic Sobolev inequality · extremal functions · Kelvin transformation · scale invariance · Emden-Fowler transformation · radial symmetry · symmetry breaking · linearization · existence · compactness

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1 Introduction and main results

In this paper we are interested in the symmetry properties of extremals for a family of interpolation inequalities established by Caffarelli, Kohn and Nirenberg in [1]. We also address the same issue for a class of weighted logarithmic Hardy inequalities which appear as limiting cases of the first ones, see [4,5].

More precisely, let $d \in \mathbb{N}^*$, $\theta \in (0, 1)$ and define

$$\vartheta(d, p) := d \frac{p-2}{2p}, \quad a_c := \frac{d-2}{2}, \quad \Lambda(a) := (a - a_c)^2, \quad p(a, b) := \frac{2d}{d-2+2(b-a)}.$$

Notice that

$$0 \leq \vartheta(d, p) \leq \theta < 1 \quad \iff \quad 2 \leq p < p^*(d, \theta) := \frac{2d}{d-2\theta} \leq 2^*,$$

where, as usual, $2^* = p^*(d, 1) = \frac{2d}{d-2}$ if $d \geq 3$, while we set $2^* = p^*(2, 1) = \infty$ if $d = 2$. If $d = 1$, θ is restricted to $[0, 1/2)$ and we set $2^* = p^*(1, 1/2) = \infty$. In this paper, we are concerned with the following interpolation inequalities:

Theorem 1 [1, 4, 5] *Let $d \geq 1$ and $a < a_c$.*

(i) *Let $b \in (a+1/2, a+1]$ when $d = 1$, $b \in (a, a+1]$ when $d = 2$ and $b \in [a, a+1]$ when $d \geq 3$. In addition, assume that $p = p(a, b)$. For any $\theta \in [\vartheta(d, p), 1]$, there exists a finite positive constant $C_{\text{CKN}}(\theta, p, \Lambda)$ with $\Lambda = \Lambda(a)$ such that*

$$\left(\int_{\mathbb{R}^d} \frac{|u|^p}{|x|^{bp}} dx \right)^{\frac{2}{p}} \leq C_{\text{CKN}}(\theta, p, \Lambda) \left(\int_{\mathbb{R}^d} \frac{|\nabla u|^2}{|x|^{2a}} dx \right)^{\theta} \left(\int_{\mathbb{R}^d} \frac{|u|^2}{|x|^{2(a+1)}} dx \right)^{1-\theta} \quad (1)$$

for any $u \in \mathcal{D}_a^{1,2}(\mathbb{R}^d)$. Equality in (1) is attained for any $p \in (2, 2^*)$ and $\theta \in (\vartheta(p, d), 1)$ or $\theta = \vartheta(p, d)$ and $a_c - a > 0$ not too large. It is not attained if $p = 2$, or $a < 0$, $p = 2^*$ and $d \geq 3$, or $d = 1$ and $\theta = \vartheta(p, d)$.

(ii) *Let $\gamma \geq d/4$ and $\gamma > 1/2$ if $d = 2$. There exists a positive constant $C_{\text{WLH}}(\gamma, \Lambda)$ with $\Lambda = \Lambda(a)$ such that, for any $u \in \mathcal{D}_a^{1,2}(\mathbb{R}^d)$, normalized by $\int_{\mathbb{R}^d} \frac{|u|^2}{|x|^{2(a+1)}} dx = 1$, we have:*

$$\int_{\mathbb{R}^d} \frac{|u|^2}{|x|^{2(a+1)}} \log \left(|x|^{d-2-2a} |u|^2 \right) dx \leq 2\gamma \log \left[C_{\text{WLH}}(\gamma, \Lambda) \int_{\mathbb{R}^d} \frac{|\nabla u|^2}{|x|^{2a}} dx \right] \quad (2)$$

and equality is attained if $\gamma \geq 1/4$ and $d = 1$, or $\gamma > 1/2$ if $d = 2$, or for $d \geq 3$ and either $\gamma > d/4$ or $\gamma = d/4$ and $a_c - a > 0$ not too large.

Caffarelli-Kohn-Nirenberg interpolation inequalities (1) and the weighted logarithmic Hardy inequality (2) are respectively the main results of [1] and [4]. Existence of extremals has been studied in [5]. We shall assume that all constants in the inequalities are taken with their optimal values. For brevity, we shall call *extremals* the functions which attain equality in (1) or in (2). Note that the set $\mathcal{D}_a^{1,2}(\mathbb{R}^d)$ denotes the completion with respect to the norm

$$u \mapsto \| |x|^{-a} \nabla u \|_{L^2(\mathbb{R}^d)}^2 + \| |x|^{-(a+1)} u \|_{L^2(\mathbb{R}^d)}^2$$

of the set $\mathcal{D}(\mathbb{R}^d \setminus \{0\})$ of smooth functions with compact support contained in $\mathbb{R}^d \setminus \{0\}$.

The parameters $a < a_c$ and $\Lambda = \Lambda(a) > 0$ are in one-to-one correspondence and it could look more natural to ask the constants C_{CKN} and C_{WLH} to depend on a rather than on Λ . As we shall see later, it turns out to be much more convenient to express all quantities in terms of Λ , once the problem has been reformulated using the Emden-Fowler transformation. Furthermore, we can notice that the restriction $a < a_c$ can be removed using a transformation of Kelvin type: see [7] and Section 2.1 for details.

In the sequel we will denote by $C_{\text{CKN}}^*(\theta, p, \Lambda)$ and $C_{\text{WLH}}^*(\gamma, \Lambda)$ the optimal constants in (1) and (2) respectively, when considered among radially symmetric functions. In this case the corresponding extremals are known (see [4]) and the constants can be explicitly computed:

$$C_{\text{CKN}}^*(\theta, p, \Lambda) := \left[\frac{\Lambda(p-2)^2}{2+(2\theta-1)p} \right]^{\frac{p-2}{2p}} \left[\frac{2+(2\theta-1)p}{2p\theta\Lambda} \right]^\theta \left[\frac{4}{p+2} \right]^{\frac{6-p}{2p}} \left[\frac{\Gamma\left(\frac{2}{p-2} + \frac{1}{2}\right)}{\sqrt{\pi} \Gamma\left(\frac{2}{p-2}\right)} \right]^{\frac{p-2}{p}},$$

$$C_{\text{WLH}}^*(\gamma, \Lambda) = \frac{1}{4\gamma} \frac{[\Gamma(\frac{d}{2})]^{2\gamma}}{(2\pi^{d+1}e)^{\frac{1}{4\gamma}}} \left(\frac{4\gamma-1}{\Lambda} \right)^{\frac{4\gamma-1}{4\gamma}} \quad \text{if } \gamma > \frac{1}{4}$$

and $C_{\text{WLH}}^*(\frac{1}{4}, \Lambda) = \frac{[\Gamma(\frac{d}{2})]^2}{2\pi^{d+1}e}$.

Notice that $\gamma = 1/4$ is compatible with the condition $\gamma \geq d/4$ only if $d = 1$. The constant $C_{\text{WLH}}^*(1/4, \Lambda)$ is then independent of Λ .

By definition, we know that

$$C_{\text{CKN}}^*(\theta, p, \Lambda) \leq C_{\text{CKN}}(\theta, p, \Lambda) \quad \text{and} \quad C_{\text{WLH}}^*(\gamma, \Lambda) \leq C_{\text{WLH}}(\gamma, \Lambda).$$

The main goal of this paper is to distinguish the set of parameters (θ, p, Λ) and (γ, Λ) for which equality holds in the above inequalities from the set where the inequality is strict.

To this purpose, we recall that when $\theta = 1$ and $d \geq 2$, symmetry breaking for extremals of (1) has been proved in [2, 9, 7] when

$$a < 0 \quad \text{and} \quad p > \frac{2}{a_c - a} \sqrt{\Lambda(a) + d - 1}.$$

In other words, for $\theta = 1$ and

$$a < A(p) := a_c - 2 \sqrt{\frac{d-1}{(p+2)(p-2)}} < 0,$$

we have $C_{\text{CKN}}^*(\theta, p, \Lambda) < C_{\text{CKN}}(\theta, p, \Lambda)$. This result has been extended to the case $\theta \in [\vartheta(p, d), 1]$ in [4]. Let

$$\Theta(a, p, d) := \frac{p-2}{32(d-1)p} \left[(p+2)^2 (d^2 + 4a^2 - 4a(d-2)) - 4p(p+4)(d-1) \right]$$

and

$$a_-(p) := a_c - \frac{2(d-1)}{p+2}.$$

Proposition 1 [4] *Let $d \geq 2$, $2 < p < 2^*$ and $a < a_-(p)$. Optimality for (1) is not achieved among radial functions if*

- (i) either $\vartheta(p, d) \leq \theta < \Theta(a, p, d)$ and $a \geq A(p)$,
(ii) or $\vartheta(p, d) \leq \theta \leq 1$ and $a < A(p)$.

More precisely, one sees that symmetry breaking occurs if $\theta < \Theta(a, p, d)$. We observe that, for $p \in [2, 2^*)$, we have $\vartheta(p, d) < \Theta(a, p, d)$ if and only if $a < a_-(p)$. The condition $\Theta(a, p, d) \leq 1$ is equivalent to $a \geq A(p)$.

By rewriting the condition $\theta < \Theta(a, p, d)$ in terms of a , we find that in the set $\{(\theta, p) : \vartheta(p, d) \leq \theta \leq 1, p \in (2, 2^*)\}$ the function

$$\underline{a}(\theta, p) := a_c - \frac{2\sqrt{d-1}}{p+2} \sqrt{\frac{2p\theta}{p-2} - 1} \quad (3)$$

takes values in $(-\infty, a_c)$ and is such that symmetry breaking holds for any $a < \underline{a}(\theta, p)$. Notice in particular that $a_-(p) = \underline{a}(\vartheta(p, d), p)$ and that we recover the condition $a < A(p)$ for $\theta = 1$.

Before going further, let us comment on the nature of the above symmetry breaking result. Among radially symmetric functions, extremals are uniquely defined up to a multiplication by a constant and a scaling. Denote by u^* the unique radial extremal in (1) under an appropriate normalization (see [4] for details). Conditions $a < A(p) = \underline{a}(1, p)$ for $\theta = 1$ and $a < \underline{a}(\theta, p)$ for $\theta < 1$ correspond exactly to the values of the parameters for which the linearized operator associated to the functional $\mathcal{F}_{\theta, p, \Lambda}$ (see Section 2.2) around u^* in the space orthogonal to the radial functions admits a *negative* eigenvalue, while it is *positive definite* for $a > \underline{a}(\theta, p)$. Thus, in the first case, u^* no longer corresponds to a minimizer for the variational problem associated with the inequality. Also notice that, if for a sequence of *non-radial* extremals $(u_n)_n, (a_n)_n$ converges to some a and $(u_n)_n$ converges to a *radial* extremal u^* , then $a = \underline{a}(\theta, p)$.

As in [4], it is worthwhile to observe that if $a < -1/2$, then

$$\frac{d}{4} = \frac{\partial}{\partial p} \vartheta(p, d)|_{p=2} < \frac{\partial}{\partial p} \Theta(a, p, d)|_{p=2} = \frac{1}{4} + \frac{\Lambda(a)}{d-1}.$$

This is consistent with the limiting case $\theta = \gamma(p-2)$ and $p \rightarrow 2_+$ corresponding to the the weighted logarithmic Hardy inequality (2).

Proposition 2 [4] *Let $d \geq 2$ and $a < -1/2$. Assume that $\gamma > 1/2$ if $d = 2$ and*

$$\frac{d}{4} \leq \gamma < \frac{1}{4} + \frac{\Lambda(a)}{d-1},$$

then the optimal constant $C_{\text{WLH}}(\gamma, \Lambda(a))$ in inequality (2) is not achieved by a radial function.

In other words, letting

$$\tilde{a}(\gamma) := a_c - \frac{1}{2} \sqrt{(d-1)(4\gamma-1)} \quad (4)$$

then, for any given $\gamma > d/4$, symmetry breaking occurs whenever $a \in (-\infty, \tilde{a}(\gamma))$.

A first step of our analysis is to counterbalance the above symmetry breaking results with some symmetry results. To this purpose we recall that for $\theta = 1$, radial symmetry for extremals of (1) was proved by various methods in [3, 11, 6] if $0 \leq a < a_c$. We shall extend these results to the case $\theta < 1$ using the method of [6]. Our first new result is based on Schwarz' symmetrization, and states the following:

Theorem 2 For any $d \geq 3$, $p \in (2, 2^*)$, there is a curve $\theta \mapsto \bar{a}(\theta, p)$ such that, for any $a \in [\bar{a}(\theta, p), a_c)$, $C_{\text{CKN}}(\theta, p, \Lambda(a)) = C_{\text{CKN}}^*(\theta, p, \Lambda(a))$. Moreover, $\lim_{\theta \rightarrow 1^-} \bar{a}(\theta, p) = 0$, and $\lim_{\theta \rightarrow 0^+} \bar{a}(\theta, p) = a_c$.

At this point, $d = 2$ is not covered and we have no corresponding result for the weighted logarithmic Hardy inequality. Actually, numerical computations (see Fig. 1) do not indicate that our method, which is based on Schwarz' symmetrization, could eventually apply to the logarithmic Hardy inequality.

As for the case $\theta = 1$, $d \geq 2$, where symmetry is known to hold for (1) in a neighborhood of $a = 0_-$, for $b > 0$: see [12, 13, 14, 7, 6], the symmetry result of Theorem 2 is far from sharp. Indeed, for $\theta = 1$, it has recently been proved in [6] that symmetry also holds for p in a neighborhood of 2_+ , and that there is a continuous curve $p \mapsto a(p)$ such that symmetry holds for any $a \in (a(p), a_c)$, while extremals are not radially symmetric if $a \in (-\infty, a(p))$. We shall extend this result to the more general interpolation inequalities (1) and to (2).

Notice that establishing radial symmetry in the case $0 < \theta < 1$ in (1) poses a more delicate problem than when $\theta = 1$, because of the term $\| |x|^{-(a+1)} u \|_{L^2(\mathbb{R}^d)}^2$. Nonetheless, by adapting the arguments of [6], we shall prove that a continuous surface splits the set of parameters into two sets that identify respectively the symmetry and symmetry breaking regions. The case $d = 2$ is also covered, while it was not in Theorem 2.

Theorem 3 For all $d \geq 2$, there exists a continuous function a^* defined on the set $\{(\theta, p) \in (0, 1] \times (2, 2^*) : \theta > \vartheta(p, d)\}$ with values in $(-\infty, a_c)$ such that $\lim_{p \rightarrow 2^+} a^*(\theta, p) = -\infty$ and

- (i) If $(a, p) \in (a^*(\theta, p), a_c) \times (2, 2^*)$, (1) has only radially symmetric extremals.
- (ii) If $(a, p) \in (-\infty, a^*(\theta, p)) \times (2, 2^*)$, none of the extremals of (1) is radially symmetric.
- (iii) For every $p \in (2, 2^*)$, $\underline{a}(\theta, p) \leq a^*(\theta, p) \leq \bar{a}(\theta, p) < a_c$.

Surprisingly, the symmetry in the regime $a \rightarrow a_c$ appears as a consequence of the asymptotic behavior of the extremals in (1) for $\theta = 1$ as $a \rightarrow -\infty$, which has been established in [2]. Symmetry holds as $p \rightarrow 2_+$ for reasons which are similar to the ones found in [6].

Concerning the weighted logarithmic Hardy inequality (2), we observe that it can be obtained as the limiting case of inequality (1) as $p \rightarrow 2_+$, provided $\theta = \gamma(p - 2)$. Actually, in this limit, the inequality degenerates into an equality, so that (2) is obtained by differentiating both sides of the inequality with respect to p at $p = 2$. It is therefore remarkable that symmetry and symmetry breaking results can be extended to (2), which is a kind of first order correction to Hardy's inequality. Inequality (2) has been established recently and so far no symmetry results were known for its extremals. Here is our first main result:

Theorem 4 Let $d \geq 2$, there exists a continuous function $a^{**} : (d/4, \infty) \rightarrow (-\infty, a_c)$ such that for any $\gamma > d/4$ and $a \in [a^{**}(\gamma), a_c)$, there is a radially symmetric extremal for (2), while for $a < a^{**}(\gamma)$ no extremal of (2) is radially symmetric. Moreover, $a^{**}(\gamma) \geq \bar{a}(\gamma)$ for any $\gamma \in (d/4, \infty)$.

Theorems 3 and 4 do not allow to decide whether $(\theta, p) \mapsto a^*(\theta, p)$ and $\gamma \mapsto a^{**}(\gamma)$ coincide with $(\theta, p) \mapsto \underline{a}(\theta, p)$ and $\gamma \mapsto \tilde{a}(\gamma)$ given by (3) and (4) respectively. If the set of non-radial extremals bifurcates from the set of radial extremals, then $a^* = \underline{a}$ in case of (1) and $a^{**} = \tilde{a}$ in case of (2). Moreover, most of the known symmetry breaking results rely on linearization and the method developed in [6] for proving symmetry and applied in Theorems 3 and 4 also relies on linearization. It would therefore be tempting to conjecture that $a^* = \underline{a}$ and $a^{**} = \tilde{a}$. It turns out that this is not the case. We are now going to establish a *new symmetry breaking phenomenon*, outside the zone of instability of the radial extremal, i.e. when $a > \underline{a}$, for some values of $\theta < 1$ for (1) and for some $a > \tilde{a}(\gamma)$ in case of (2). These are striking results, as they clearly depart from previous methods.

Theorem 5 *Let $d \geq 2$. There exists $\eta > 0$ such that for every $p \in (2, 2 + \eta)$ there exists an $\varepsilon > 0$ with the property that for $\theta \in [\vartheta(p, d), \vartheta(p, d) + \varepsilon)$ and $a \in [\underline{a}(\theta, p), \underline{a}(\theta, p) + \varepsilon)$, no extremal for (1) corresponding to the parameters (θ, p, a) is radially symmetric.*

Notice that there is always an extremal function for (1) if $\theta > \vartheta(p, d)$, and also in some cases if $\theta = \vartheta(p, d)$. See [5] for details. The plots in Fig. 2 provide a value for η .

We have a similar statement for logarithmic Hardy inequalities, which is our third main result. Let

$$A_{\text{SB}}(\gamma, d) := \frac{1}{8} (4\gamma - 1) e^{\left(\frac{\pi^4 \gamma - d - 1}{16}\right) \frac{1}{4\gamma - 1}} \left(\frac{d}{\gamma}\right)^{\frac{4\gamma}{4\gamma - 1}} \Gamma\left(\frac{d}{2}\right)^{\frac{2}{4\gamma - 1}}. \quad (5)$$

Theorem 6 *Let $d \geq 2$ and assume that $\gamma > 1/2$ if $d = 2$. If $A(a) > A_{\text{SB}}(\gamma, d)$, then there is symmetry breaking: no extremal for (2) corresponding to the parameters (γ, a) is radially symmetric. As a consequence, there exists an $\varepsilon > 0$ such that, if $a \in [\tilde{a}(\gamma), \tilde{a}(\gamma) + \varepsilon)$ and $\gamma \in [d/4, d/4 + \varepsilon)$, with $\gamma > 1/2$ if $d = 2$, there is symmetry breaking.*

This result improves the one of Proposition 2, at least for γ in a neighborhood of $(d/4)_+$. Actually, the range of γ for which $A_{\text{SB}}(\gamma, d) < \tilde{A}(\gamma)$ can be deduced from our estimates, although explicit expressions are hard to read. See Fig. 4 and further comments at the end of Section 5.

This paper is organized as follows. Section 2 is devoted to preliminaries (Emden-Fowler transform, symmetry breaking results based on the linear instability of radial extremals) and to the proof of Theorem 2 using Schwarz' symmetrisation. Sections 3 and 4 are devoted to the proofs of Theorems 3 and 4 respectively. Theorems 5 and 6 are established in Section 5.

2 Preliminaries

2.1 The Emden-Fowler transformation

Consider the Emden-Fowler transformation

$$u(x) = |x|^{-(d-2-2a)/2} w(y) \quad \text{where} \quad y = (s, \omega) \in \mathbb{R} \times \mathbb{S}^{d-1} =: \mathcal{C}, \\ x \in \mathbb{R}^d, \quad s = -\log|x| \in \mathbb{R} \quad \text{and} \quad \omega = x/|x| \in \mathbb{S}^{d-1}.$$

The scaling invariance in \mathbb{R}^d becomes a translation invariance in the cylinder \mathcal{C} , in the s -direction, and radial symmetry for a function in \mathbb{R}^d becomes dependence on the s -variable only. Also, invariance under a certain Kelvin transformation in \mathbb{R}^d corresponds to the symmetry $s \mapsto -s$ in \mathcal{C} . More precisely, a radially symmetric function u on \mathbb{R}^d , invariant under the Kelvin transformation $u(x) \mapsto |x|^{2(a-a_c)} u(x/|x|^2)$, corresponds to a function w on \mathcal{C} which depends only on s and satisfies $w(-s) = w(s)$. We shall call such a function a s -symmetric function.

Under this transformation, (1) can be stated just as an interpolation inequality in $H^1(\mathcal{C})$. Namely, for any $w \in H^1(\mathcal{C})$,

$$\|w\|_{L^p(\mathcal{C})}^2 \leq C_{\text{CKN}}(\theta, p, \Lambda) \left(\|\nabla w\|_{L^2(\mathcal{C})}^2 + \Lambda \|w\|_{L^2(\mathcal{C})}^2 \right)^\theta \|w\|_{L^2(\mathcal{C})}^{2(1-\theta)} \quad (6)$$

with $\Lambda = \Lambda(a) = (a_c - a)^2$. With these notations, recall that

$$\Lambda = 0 \iff a = a_c \quad \text{and} \quad \Lambda > 0 \iff a < a_c .$$

At this point it becomes clear that $a < a_c$ or $a > a_c$ plays no role and only the value of $\Lambda > 0$ matters. Similarly, by the Emden-Fowler transformation, (2) becomes

$$\int_{\mathcal{C}} |w|^2 \log |w|^2 \, dy \leq 2\gamma \log \left[C_{\text{WLH}}(\gamma, \Lambda) \left(\|\nabla w\|_{L^2(\mathcal{C})}^2 + \Lambda \right) \right], \quad (7)$$

for any $w \in H^1(\mathcal{C})$ normalized by $\|w\|_{L^2(\mathcal{C})}^2 = 1$, for any $d \geq 1$, $a < a_c$, $\gamma \geq d/4$, and $\gamma > 1/2$ if $d = 2$.

2.2 Linear instability of radial extremals

Symmetry breaking for extremals of (7) has been discussed in detail for $\theta = 1$ in [2, 9], and by the same methods in [4], where symmetry breaking has been established also when $\theta \in (0, 1)$. The method goes as follows. Consider an extremal w^* for (6) among s -symmetric functions. It realizes a minimum for the functional

$$\mathcal{F}_{\theta, p, \Lambda}[w] := \frac{\left(\|\nabla w\|_{L^2(\mathcal{C})}^2 + \Lambda \|w\|_{L^2(\mathcal{C})}^2 \right) \|w\|_{L^2(\mathcal{C})}^{2 \frac{1-\theta}{\theta}}}{\|w\|_{L^p(\mathcal{C})}^{2/\theta}} \quad (8)$$

among functions depending only on s and $\mathcal{F}_{\theta, p, \Lambda}[w^*] = C_{\text{CKN}}^*(\theta, p, \Lambda)^{-1/\theta}$. Once the maximum of w^* is fixed at $s = 0$, since w^* solves an autonomous ordinary differential equation, by uniqueness, it automatically satisfies the symmetry $w^*(-s) = w^*(s)$ for any $s \in \mathbb{R}$. Next, one linearizes $\mathcal{F}_{\theta, p, \Lambda}$ around w^* . This gives rise to a linear operator, whose kernel is generated by dw^*/ds and which admits a negative eigenvalue in $H^1(\mathcal{C})$ if and only if $a < \underline{a}(\theta, p)$, that is for

$$\Lambda > \underline{\Lambda}(\theta, p) := (a_c - \underline{a}(\theta, p))^2,$$

where the function $\underline{a}(\theta, p)$ is defined in (3). Hence, if $a < \underline{a}(\theta, p)$, it is clear that $\mathcal{F}_{\theta, p, \Lambda} - \mathcal{F}_{\theta, p, \Lambda}[w^*]$ takes negative values in a neighbourhood of w^* in $H^1(\mathcal{C})$ and extremals for (6) cannot be s -symmetric, even up to translations in the s -direction. By the Emden-Fowler transformation, extremals for (1) cannot be radially symmetric.

Remark 1 Theorem 5 asserts that there are cases where $a > \underline{a}(\theta, p)$, so that the extremal s -symmetric function w^* is stable in $H^1(\mathcal{C})$, but for which symmetry is broken, in the sense that we prove $\mathbf{C}_{\text{CKN}}^*(\theta, p, \Lambda) < \mathbf{C}_{\text{CKN}}(\theta, p, \Lambda)$. This will be studied in Section 5.

In the case of the weighted logarithmic Hardy inequality, symmetry breaking can be investigated as in [4] by studying the linearization of the functional

$$\mathcal{G}_{\gamma, \Lambda}[w] := \frac{\|\nabla w\|_{L^2(\mathcal{C})}^2 + \Lambda \|w\|_{L^2(\mathcal{C})}^2}{\|w\|_{L^2(\mathcal{C})}^2 \exp \left\{ \frac{1}{2\gamma} \int_{\mathcal{C}} \frac{w^2}{\|w\|_{L^2(\mathcal{C})}^2} \log \left(\frac{w^2}{\|w\|_{L^2(\mathcal{C})}^2} \right) dy \right\}}. \quad (9)$$

around an s -symmetric extremal w^* . In this way one finds that extremals for inequality (7) are not s -symmetric whenever $d \geq 2$,

$$\Lambda > \tilde{\Lambda}(\gamma) := \frac{1}{4} (d-1)(4\gamma-1) = \Lambda(\tilde{a}(\gamma)).$$

2.3 Proof of Theorem 2

As in [6], we shall prove Theorem 2 by Schwarz' symmetrization after rephrasing (1) as follows. To $u \in \mathcal{D}_a^{1,2}(\mathbb{R}^d)$, we may associate the function $v \in \mathcal{D}_0^{1,2}(\mathbb{R}^d)$ by setting:

$$u(x) = |x|^a v(x) \quad \forall x \in \mathbb{R}^d.$$

Inequality (1) is then equivalent to

$$\||x|^{a-b} v\|_{L^p(\mathbb{R}^d)}^2 \leq \mathbf{C}_{\text{CKN}}(\theta, p, \Lambda) (\mathcal{A} - \lambda \mathcal{B})^\theta \mathcal{B}^{1-\theta}$$

with $\mathcal{A} := \|\nabla v\|_{L^2(\mathbb{R}^d)}^2$, $\mathcal{B} := \||x|^{-1} v\|_{L^2(\mathbb{R}^d)}^2$ and $\lambda := a(2a_c - a)$. We observe that the function $\mathcal{B} \mapsto h(\mathcal{B}) := (\mathcal{A} - \lambda \mathcal{B})^\theta \mathcal{B}^{1-\theta}$ satisfies

$$\frac{h'(\mathcal{B})}{h(\mathcal{B})} = \frac{1-\theta}{\mathcal{B}} - \frac{\lambda\theta}{\mathcal{A} - \lambda\mathcal{B}}.$$

By Hardy's inequality, we know that

$$\mathcal{A} - \lambda \mathcal{B} \geq \inf_{a>0} (\mathcal{A} - a(2a_c - a)\mathcal{B}) = \mathcal{A} - a_c^2 \mathcal{B} > 0$$

for any $v \in \mathcal{D}_0^{1,2}(\mathbb{R}^d) \setminus \{0\}$. As a consequence, $h'(\mathcal{B}) \leq 0$ if

$$(1-\theta)\mathcal{A} < \lambda \mathcal{B}. \quad (10)$$

If this is the case, Schwarz' symmetrization applied to v decreases \mathcal{A} , increases \mathcal{B} , and therefore decreases $(\mathcal{A} - \lambda \mathcal{B})^\theta \mathcal{B}^{1-\theta}$, while it increases $\||x|^{a-b} v\|_{L^p(\mathbb{R}^d)}^2$. Optimality in (1) is then reached among radial functions. Notice that $\lambda > 0$ is required by our method and hence only the case $a_c > 0$, i.e. $d \geq 3$, is covered.

Let

$$t := \frac{\mathcal{A}}{\mathcal{B}} - a_c^2.$$

Condition (10) amounts to

$$t \leq \frac{\theta a_c^2 - (a_c - a)^2}{1 - \theta}. \quad (11)$$

If u is a minimizer for (1), it has been established in [5, Lemma 3.4] that

$$(t + \Lambda)^\theta \leq \frac{(\mathcal{C}_{\text{CKN}}(1, 2^*, a_c^2))^{\vartheta(d,p)}}{\mathcal{C}_{\text{CKN}}^*(\theta, p, 1)} (a_c - a)^{2\theta - \frac{2}{d}\vartheta(p,d)} (t + a_c^2)^{\vartheta(d,p)}. \quad (12)$$

For completeness, we shall briefly sketch the proof of (12) below in Remark 5. The two conditions (11) and (12) determine two upper bounds for t , which are respectively monotone decreasing and monotone increasing in terms of a . As a consequence, they are simultaneously satisfied if and only if $a \in [a_0, a_c)$, where a_0 is determined by the equality case in (11) and (12). See Fig. 1. This completes the proof of Theorem 2. \square

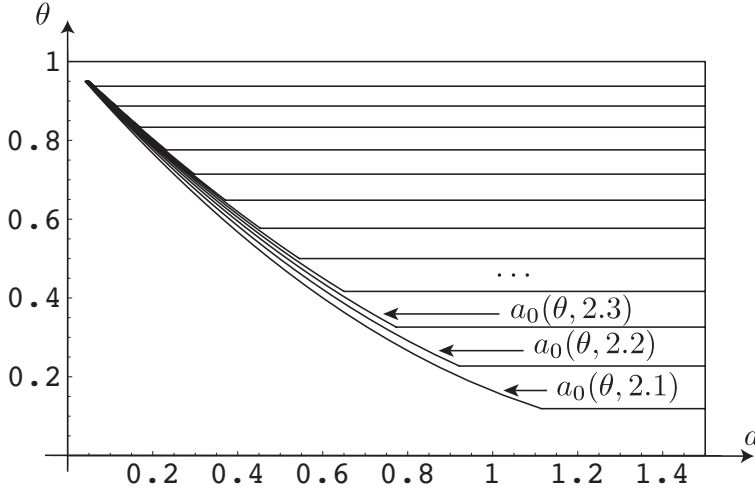


Fig. 1 According to the proof of Theorem 2, symmetry holds if $a \in [a_0(\theta, p), a_c)$, $\theta \in (\vartheta(p, d), 1)$. The curves $\theta \mapsto a_0(\theta, p)$ are parametrized by $\theta \in [\vartheta(p, d), 1)$, with $d = 5$, $a_c = 1.5$ and $p = 2.1, 2.2, \dots, 3.2$. Horizontal segments correspond to $\theta = \vartheta(p, d)$, $a_0(\theta, p) \leq a < a_c$.

Remark 2 Although this is not needed for the proof of Theorem 2, to understand why symmetry can be expected as $a \rightarrow a_c$, it is enlightening to consider the moving planes method. With the above notations, if u is an extremal for (1), then v is a solution of the Euler-Lagrange equation

$$-\frac{\theta}{\mathcal{A} - \lambda \mathcal{B}} \Delta v + \left(\frac{1 - \theta}{\mathcal{B}} - \frac{\theta \lambda}{\mathcal{A} - \lambda \mathcal{B}} \right) \frac{v}{|x|^2} = \frac{v^{p-1}}{\| |x|^{-(b-a)} v \|_{L^p(\mathbb{R}^d)}^p}.$$

If $d = 2$, then $\lambda = -a^2 < 0$ and $\frac{1-\theta}{\mathcal{B}} - \frac{\theta \lambda}{\mathcal{A} - \lambda \mathcal{B}}$ is always positive. If $d \geq 3$, $\frac{1-\theta}{\mathcal{B}} - \frac{\theta \lambda}{\mathcal{A} - \lambda \mathcal{B}}$ is negative if and only if (10) holds. Assume that this is the case. Using the Emden-Fowler transformation defined in Section 2.1, we know that the corresponding solution on the cylinder is smooth, so that v has no singularity except maybe at $x = 0$. We

can then use the moving planes technique and prove that v is radially symmetric by adapting the results of [10, 8].

Using Hardy's inequality, $(d-2)^2 \mathcal{B} \leq 4\mathcal{A}$, also notice that (10) cannot hold unless

$$\theta > \frac{(d-2-2a)^2}{(d-2)^2} = \frac{(a_c - a)^2}{a_c^2}.$$

This imposes that $a \rightarrow a_c$ as $\theta \rightarrow 0_+$. Compared to (10), a numerical investigation (see Fig. 1) shows that this last condition is qualitatively correct.

3 Radial symmetry for the Caffarelli-Kohn-Nirenberg inequalities

In this section, we shall first establish some a priori estimates which will allow us to adapt the method of [6] to the case of inequality (1).

3.1 A priori estimates

Recall that if u and w are related via the Emden-Fowler transformation, u is radially symmetric if and only if w is independent of the angular variables. The following result is taken from [2, Theorem 1.2, (i), p. 231], where $\theta = 1$. Here we are interested in the regime corresponding to $a \rightarrow -\infty$.

Lemma 1 *Let $d \geq 1$ and $p \in (2, 2^*)$. For any $t > 0$, there exists a constant $c(d, p, t)$ such that*

$$\frac{1}{c(d, p, t)} \|w\|_{L^p(C)}^2 \leq \|\nabla w\|_{L^2(C)}^2 + t\|w\|_{L^2(C)}^2 \quad \forall w \in H^1(C)$$

and

$$\lim_{t \rightarrow \infty} t^{\frac{d}{p} - a_c} c(d, p, t) = \sup_{u \in H^1(\mathbb{R}^d) \setminus \{0\}} \frac{\|u\|_{L^p(\mathbb{R}^d)}^2}{\|\nabla u\|_{L^2(\mathbb{R}^d)}^2 + \|u\|_{L^2(\mathbb{R}^d)}^2} =: S_p(\mathbb{R}^d).$$

In other words, as $t \rightarrow +\infty$, we have

$$t^{\frac{d}{p} - a_c} \|w\|_{L^p(C)}^2 \leq S_p(\mathbb{R}^d) (1 + o(1)) \left(\|\nabla w\|_{L^2(C)}^2 + t\|w\|_{L^2(C)}^2 \right)$$

for any given $p \in (2, 2^*)$.

Remark 3 $S_p(\mathbb{R}^d)$ is the best constant in the Gagliardo-Nirenberg inequality

$$\|u\|_{L^p(\mathbb{R}^d)}^2 \leq S_p(\mathbb{R}^d) \left(\|\nabla u\|_{L^2(\mathbb{R}^d)}^2 + \|u\|_{L^2(\mathbb{R}^d)}^2 \right)$$

and $t^{\frac{d}{p} - a_c}$ is the factor which appears by the scaling $u \mapsto t^{-(d-2)/4} u(\cdot/\sqrt{t})$, that is

$$t^{\frac{d}{p} - a_c} \|u\|_{L^p(\mathbb{R}^d)}^2 \leq S_p(\mathbb{R}^d) \left(\|\nabla u\|_{L^2(\mathbb{R}^d)}^2 + t\|u\|_{L^2(\mathbb{R}^d)}^2 \right)$$

for all $t > 0$. This is natural in view of the analysis done in [2]. We also observe that $\lim_{p \rightarrow 2} S_p(\mathbb{R}^d) = 1$.

From Lemma 1, we can actually deduce that the asymptotic behavior of $c(d, p, t)$ as $t \rightarrow \infty$ is uniform in the limit $p \rightarrow 2$.

Corollary 1 *Let $d \geq 1$ and $q \in (2, 2^*)$. For any $p \in [2, q]$,*

$$c(d, p, t) \leq t^{-\zeta} [c(d, q, t)]^{1-\zeta} \quad \forall t > 0$$

with $\zeta = \frac{2(q-p)}{p(q-2)}$. As a consequence,

$$\lim_{t \rightarrow \infty} \sup_{p \in [2, q]} t^{\frac{d}{p} - a_c} c(d, p, t) \leq \left[S_q(\mathbb{R}^d) \right]^{\frac{q(p-2)}{p(q-2)}}.$$

Proof Using the trivial estimate

$$\|w\|_{L^2(C)}^2 \leq \frac{1}{t} \left[\|\nabla w\|_{L^2(C)}^2 + t \|w\|_{L^2(C)}^2 \right],$$

the estimate of Lemma 1

$$\|w\|_{L^q(C)}^2 \leq c(d, q, t) \left[\|\nabla w\|_{L^2(C)}^2 + t \|w\|_{L^2(C)}^2 \right]$$

and Hölder's interpolation: $\|w\|_{L^p(C)} \leq \|w\|_{L^2(C)}^\zeta \|w\|_{L^q(C)}^{1-\zeta}$, we easily get the first estimate. Since

$$\frac{d}{p} - a_c - \zeta = (1 - \zeta) \left(\frac{d}{q} - a_c \right),$$

we find

$$t^{\frac{d}{p} - a_c} c(d, p, t) \leq \left(t^{\frac{d}{q} - a_c} c(d, q, t) \right)^{1-\zeta}.$$

and the second estimate follows. \square

Remark 4 Notice that for $d \geq 3$, the second estimate in Corollary 1 also holds with $q = 2^*$ and $\zeta = 1 - \vartheta(p, d)$. In such a case, we can actually prove that

$$c(d, p, t) \leq t^{a_c - \frac{d}{p}} (\vartheta(p, d) S_*(d))^{\vartheta(p, d)} (1 - \vartheta(p, d))^{1-\vartheta(p, d)}$$

where $S_*(d) = C_{\text{CKN}}(1, 2^*, a_c^2)$ is the optimal constant in Sobolev's inequality: for any $u \in H^1(\mathbb{R}^d)$, $\|u\|_{L^{2^*}(\mathbb{R}^d)}^2 \leq S_*(d) \|\nabla u\|_{L^2(\mathbb{R}^d)}^2$.

Consider the functional $\mathcal{F}_{\theta, p, \Lambda}$ defined by (8) on $H^1(C)$. A minimizer exists for any $p > 2$ if $d = 1$ or $d = 2$, and $p \in (2, 2^*)$ if $d \geq 3$. See [2] for details if $\theta = 1$ and [5, Theorem 1.3 (ii)] if $\theta \in (\vartheta(p, d), 1)$. The special, limiting case $\theta = \vartheta(p, 1)$ is discussed in [4] if $d = 1$ and in [5] if $d \geq 1$. From now on, we denote by $w = w_{\theta, p, \Lambda}$ an extremal for (6), whenever it exists, that is, a minimizer for $\mathcal{F}_{\theta, p, \Lambda}$. It satisfies the following Euler-Lagrange equations,

$$-\theta \Delta w + ((1 - \theta)t + \Lambda)w = (t + \Lambda)^{1-\theta} w^{p-1}$$

with $t := \|\nabla w\|_{L^2(C)}^2 / \|w\|_{L^2(C)}^2$, when we assume the normalization condition

$$\left(\|\nabla w\|_{L^2(C)}^2 + \Lambda \|w\|_{L^2(C)}^2 \right)^\theta \|w\|_{L^2(C)}^{2(1-\theta)} = \|w\|_{L^p(C)}^p.$$

Such a condition can always be achieved by homogeneity and implies

$$\|w\|_{L^p(C)}^{p-2} = \frac{1}{C_{\text{CKN}}(\theta, p, \Lambda)}. \quad (13)$$

As a consequence of the Euler-Lagrange equations, we also have

$$\|\nabla w\|_{L^2(C)}^2 + \Lambda \|w\|_{L^2(C)}^2 = (t + \Lambda)^{1-\theta} \|w\|_{L^p(C)}^p. \quad (14)$$

Remark 5 If w is a minimizer for $\mathcal{F}_{\theta, p, \Lambda}$, then we know that

$$(t + \Lambda)^\theta \|w\|_{L^2(C)}^2 = \frac{\|w\|_{L^p(C)}^2}{C_{\text{CKN}}(\theta, p, \Lambda)} \leq \frac{\|w\|_{L^p(C)}^2}{C_{\text{CKN}}^*(\theta, p, \Lambda)} = \frac{\|w\|_{L^p(C)}^2}{C_{\text{CKN}}^*(\theta, p, 1)} \Lambda^{\theta - \frac{p-2}{2p}}.$$

On the other hand, by Hölder's inequality: $\|w\|_{L^p(C)} \leq \|w\|_{L^{2^*}(C)}^{\vartheta(d,p)} \|w\|_{L^2(C)}^{1-\vartheta(d,p)}$, and by Sobolev's inequality (cf. Remark 4) written on the cylinder, we know that

$$\begin{aligned} \|w\|_{L^p(C)}^2 &\leq (\mathbf{S}_*(d))^{\vartheta(p,d)} \|w\|_{L^{2^*}(C)}^{2\vartheta(d,p)} \|w\|_{L^2(C)}^{2(1-\vartheta(d,p))} \\ &= (\mathbf{S}_*(d))^{\vartheta(p,d)} \left(\|\nabla w\|_{L^2(C)}^2 + a_c^2 \|w\|_{L^2(C)}^2 \right)^{\vartheta(p,d)} \|w\|_{L^2(C)}^{2(1-\vartheta(p,d))} \\ &= (\mathbf{S}_*(d))^{\vartheta(p,d)} \left(t + a_c^2 \right)^{\vartheta(p,d)} \|w\|_{L^2(C)}^2. \end{aligned}$$

Collecting the two estimates proves (12) for $v(x) = |x|^{-a_c} w(s, \omega)$, where $s = -\log|x|$ and $\omega = x/|x|$, for any $x \in \mathbb{R}^d$ (Emden-Fowler transformation written for $a = 0$).

As in [2], we can assume that the extremal $w = w_{\theta, p, \Lambda}$ depends only on s and on an azimuthal angle $\phi \in (0, \pi)$ of the sphere, and thus satisfies

$$-\theta (\partial_{ss} w + D_\phi (\partial_\phi w)) + ((1 - \theta)t + \Lambda) w = (t + \Lambda)^{1-\theta} w^{p-1}. \quad (15)$$

Here we denote by $\partial_s w$ and $\partial_\phi w$ the partial derivatives with respect to s and ϕ respectively, and by D_ϕ the derivative defined by: $D_\phi w := (\sin \phi)^{2-d} \partial_\phi ((\sin \phi)^{d-2} w)$. Moreover, using the translation invariance of (6) in the s -variable, the invariance of the functional $\mathcal{F}_{\theta, p, \Lambda}$ under the transformation $(s, \omega) \mapsto (-s, \omega)$ and the sliding method, we can also assume without restriction that w is such that

$$\begin{cases} w(s, \phi) = w(-s, \phi) \quad \forall (s, \phi) \in \mathbb{R} \times (0, \pi), \\ \partial_s w(s, \phi) < 0 \quad \forall (s, \phi) \in (0, +\infty) \times (0, \pi), \\ \max_C w = w(0, \phi_0), \end{cases} \quad (16)$$

for some $\phi_0 \in [0, \pi]$. In particular notice that

$$\|\nabla w\|_{L^2(C)}^2 = \omega_{d-2} \int_0^{+\infty} \int_0^\pi \left(|\partial_s w|^2 + |\partial_\phi w|^2 \right) (\sin \phi)^{d-2} d\phi ds$$

where ω_{d-2} is the area of \mathbb{S}^{d-2} . From Lemma 1, we obtain the following estimate:

Corollary 2 Assume that $d \geq 2$, $\Lambda > 0$, $p \in (2, 2^*)$ and $\theta \in (\vartheta(p, d), 1)$. Let $t = t(\theta, p, \Lambda)$ be the maximal value of $\|\nabla w\|_{L^2(C)}^2 / \|w\|_{L^2(C)}^2$ among all extremals of (6). Then $t(\theta, p, \Lambda)$ is bounded from above and moreover

$$\limsup_{p \rightarrow 2^+} t(\theta, p, \Lambda) < \infty \quad \text{and} \quad \limsup_{\Lambda \rightarrow 0^+} t(\theta, p, \Lambda) < \infty,$$

where the limits above are taken respectively for $\Lambda > 0$ and $\theta \in (0, 1)$ fixed, and for $p \in (2, 2^*)$ and $\theta \in (\vartheta(p, d), 1)$ fixed.

Proof Let $t_n := \|\nabla w_n\|_{L^2(C)}^2 / \|w_n\|_{L^2(C)}^2$, where w_n are extremals of (6) with $\Lambda = \Lambda_n \in (0, +\infty)$, $p = p_n \in (2, 2^*)$ and $\theta \in (0, 1]$. We shall be concerned with one of the following regimes:

- (i) $\Lambda_n = \Lambda$ and $p_n = p$ do not depend on $n \in \mathbb{N}$, and $\theta \in (\vartheta(p, d), 1)$,
- (ii) $\Lambda_n = \Lambda$ does not depend on $n \in \mathbb{N}$, $\theta \in (0, 1]$ and $\lim_{n \rightarrow \infty} p_n = 2$,
- (iii) $p_n = p$ does not depend on $n \in \mathbb{N}$, $\theta \in [\vartheta(p, d), 1]$ and $\lim_{n \rightarrow \infty} \Lambda_n = 0$.

Assume that $\lim_{n \rightarrow \infty} t_n = \infty$, consider (14) and apply Lemma 1 to get

$$t_n^{\frac{d}{p_n} - a_c} \frac{\min\{\theta, 1 - \theta\}}{S_{p_n}(\mathbb{R}^d)} (1 + o(1)) \leq (t_n + \Lambda_n)^{1 - \theta} \|w_n\|_{L^{p_n}(C)}^{p_n - 2} = \frac{(t_n + \Lambda_n)^{1 - \theta}}{C_{\text{CKN}}(\theta, p_n, \Lambda_n)} \quad (17)$$

where we have used the assumption that $1 - \vartheta(p_n, d) = \frac{d}{p_n} - a_c > 1 - \theta$. This gives a contradiction in case (i).

Using the fact that

$$C_{\text{CKN}}(\theta, p, \Lambda) \geq C_{\text{CKN}}^*(\theta, p, \Lambda)$$

where $C_{\text{CKN}}^*(\theta, p, \Lambda)$ is the best constant in (1) among radial functions given in Section 1, and observing that

$$C_{\text{CKN}}^*(\theta, p, \Lambda) = C_{\text{CKN}}^*(\theta, p, 1) \Lambda^{\frac{p-2}{2p} - \theta},$$

we get

$$1/C_{\text{CKN}}^*(\theta, p, \Lambda) \sim \Lambda^{\theta - \frac{p-2}{2p}} \rightarrow 0 \quad \text{as} \quad \Lambda \rightarrow 0.$$

In case (iii), if we assume that $t_n \rightarrow +\infty$, then this provides a contradiction with (17).

In case (ii), we know that

$$\lim_{p \rightarrow 2^+} C_{\text{CKN}}^*(\theta, p, \Lambda) = \Lambda^{-\theta}$$

and, by (17) and Lemma 1,

$$t_n^{\frac{d}{p_n} - a_c} \frac{\min\{\theta, 1 - \theta\}}{S_{p_n}(\mathbb{R}^d)} (1 + o(1)) \leq \frac{(t_n + \Lambda)^{1 - \theta}}{C_{\text{CKN}}^*(\theta, p_n, \Lambda)} = \Lambda^\theta t_n^{1 - \theta} (1 + o(1))$$

as $n \rightarrow \infty$. Again this provides a contradiction in case we assume $\lim_{n \rightarrow \infty} t_n = \infty$. \square

Let $k(p, \Lambda) := C_{\text{CKN}}^*(\theta = 1, p, \Lambda)$ and recall that $k(p, \Lambda) = \Lambda^{-(p+2)/(2p)} k(p, 1)$ and $\lim_{p \rightarrow 2^+} k(p, 1) = 1$. As a consequence of the symmetry result in [6], we have

Lemma 2 *There exists a positive continuous function $\bar{\varepsilon}$ on $(2, 2^*)$ with*

$$\lim_{p \rightarrow 2} \bar{\varepsilon}(p) = \infty \quad \text{and} \quad \lim_{p \rightarrow 2^*} \bar{\varepsilon}(p) = a_c^{-(p+2)/p} k(2^*, 1)$$

such that, for any $p \in (2, 2^*)$,

$$\|w\|_{L^p(\mathcal{C})}^2 \leq \varepsilon \|\nabla w\|_{L^2(\mathcal{C})}^2 + Z(\varepsilon, p) \|w\|_{L^2(\mathcal{C})}^2 \quad \forall w \in H^1(\mathcal{C})$$

holds for any $\varepsilon \in (0, \bar{\varepsilon}(p))$ with $Z(\varepsilon, p) := \varepsilon^{-\frac{p-2}{p+2}} k(p, 1)^{\frac{2p}{p+2}}$.

Proof From [6], we know that there exists a continuous function $\bar{\lambda} : (2, 2^*) \rightarrow (a_c^2, \infty)$ such that $\lim_{p \rightarrow 2} \bar{\lambda}(p) = \infty$, $\lim_{p \rightarrow 2^*} \bar{\lambda}(p) = a_c^2$ and, for any $\lambda \in (0, \bar{\lambda}(p)]$, the inequality

$$\|w\|_{L^p(\mathcal{C})}^2 \leq k(p, \lambda) \left(\|\nabla w\|_{L^2(\mathcal{C})}^2 + \lambda \|w\|_{L^2(\mathcal{C})}^2 \right) \quad \forall w \in H^1(\mathcal{C})$$

holds true. Therefore, letting

$$\bar{\varepsilon}(p) := k(p, \bar{\lambda}(p)) = \bar{\lambda}(p)^{-(p+2)/(2p)} k(p, 1),$$

our estimate holds with $\lambda = \bar{\lambda}(p)$, $\varepsilon = \lambda^{-(p+2)/(2p)} k(p, 1)$ and $Z(\varepsilon, p) = \varepsilon \lambda$. \square

Lemma 3 *Assume that $d \geq 2$, $p \in (2, 2^*)$ and $\theta \in [\vartheta(p, d), 1]$. If w is an extremal function of (6) and if w is not s -symmetric, then*

$$\theta(d-1) + (1-\theta)t + \Lambda < (t+\Lambda)^{1-\theta} (p-1) \|w\|_{L^\infty(\mathcal{C})}^{p-2}. \quad (18)$$

Proof Let w be an extremal for (6), normalized so that (16) holds. We denote by $\phi \in (0, \pi)$ the azimuthal coordinate on \mathbb{S}^{d-1} . By the Poincaré inequality in \mathbb{S}^{d-1} , we know that:

$$\int_{\mathbb{S}^{d-1}} |D_\phi(\partial_\phi w)|^2 d\omega \geq (d-1) \int_{\mathbb{S}^{d-1}} |(\partial_\phi w)|^2 d\omega$$

while, by multiplying the equation in (15) by $D_\phi(\partial_\phi w)$, after obvious integration by parts, we find:

$$\begin{aligned} & \theta \left(\int_{\mathcal{C}} \left(|\partial_s(\partial_\phi w)|^2 + |D_\phi(\partial_\phi w)|^2 \right) dy \right) + ((1-\theta)t + \Lambda) \int_{\mathcal{C}} |\partial_\phi w|^2 dy \\ & = (t+\Lambda)^{1-\theta} (p-1) \int_{\mathcal{C}} w^{p-2} |\partial_\phi w|^2 dy \leq (t+\Lambda)^{1-\theta} (p-1) \|w\|_{L^\infty(\mathcal{C})}^{p-2} \int_{\mathcal{C}} |\partial_\phi w|^2 dy. \end{aligned}$$

By combining the two above estimates, the conclusion holds if $\|\partial_\phi w\|_{L^2(\mathcal{C})} \neq 0$. \square

3.2 The critical regime: approaching $a = a_c$

Proposition 3 *Assume that $d \geq 2$, $p \in (2, 2^*)$ and $\theta \in [\vartheta(p, d), 1]$. Let $(\Lambda_n)_n$ be a sequence converging to 0_+ and let $(w_n)_n$ be a sequence of extremals for (6), satisfying the normalization condition (13). Then both $t_n := \|\nabla w_n\|_{L^2(\mathcal{C})}^2 / \|w_n\|_{L^2(\mathcal{C})}^2$ and $\|w_n\|_{L^\infty(\mathcal{C})}$ converge to 0 as $n \rightarrow +\infty$.*

Proof First of all notice that, under the given assumption, we can use the results in [5, Theorem 1.3 (i)] in order to ensure the existence of an extremal for (6) even for $\theta = \vartheta(p, d)$. Moreover with the notations of Corollary 2, we know that $(t_n)_n$ is bounded and, by (13) and (14),

$$\|\nabla w_n\|_{L^2(\mathcal{C})}^2 + \Lambda_n \|w_n\|_{L^2(\mathcal{C})}^2 = \frac{(t_n + \Lambda_n)^{1-\theta}}{C_{\text{CKN}}(\theta, p, \Lambda_n)^{p/(p-2)}}$$

with $C_{\text{CKN}}(\theta, p, \Lambda_n)^{-p/(p-2)} \leq C_{\text{CKN}}^*(\theta, p, \Lambda_n)^{-p/(p-2)} \sim \Lambda_n^{\frac{\theta p}{p-2} - \frac{1}{2}} \rightarrow 0$ as $\Lambda_n \rightarrow 0_+$, where we have used the fact that $\frac{\theta p}{p-2} - \frac{1}{2} > 0$ for $\theta \geq \vartheta(p, d)$. Thus, using (13) we have $\lim_{n \rightarrow \infty} \|\nabla w_n\|_{L^2(\mathcal{C})} = 0$ and $\lim_{n \rightarrow \infty} \|w_n\|_{L^p(\mathcal{C})} = 0$. Hence, $(w_n)_n$ converges to $w \equiv 0$, weakly in $H_{\text{loc}}^1(\mathcal{C})$ and also in $C_{\text{loc}}^{1,\alpha}$ for some $\alpha \in (0, 1)$. By (16), it follows

$$\lim_{n \rightarrow \infty} \|w_n\|_{L^\infty(\mathcal{C})} = 0.$$

Now, let $t_\infty := \lim_{n \rightarrow \infty} t_n$ and assume by contradiction that $t_\infty > 0$. The function $W_n = w_n / \|w_n\|_{H^1(\mathbb{R}^d)}$ solves

$$-\theta \Delta W_n + ((1 - \theta) t_n + \Lambda_n) W_n = (t_n + \Lambda_n)^{1-\theta} w_n^{p-2} W_n.$$

Multiply the above equation by W_n and integrate on \mathcal{C} , to get

$$\begin{aligned} \theta \|\nabla W_n\|_{L^2(\mathcal{C})}^2 + ((1 - \theta) t_\infty (1 + o(1)) + \Lambda_n) \|W_n\|_{L^2(\mathcal{C})}^2 \\ \leq (t_\infty (1 + o(1)) + \Lambda_n)^{1-\theta} \|w_n\|_{L^\infty(\mathcal{C})}^{p-2} \|W_n\|_{L^2(\mathcal{C})}^2. \end{aligned}$$

This is in contradiction with the fact that $\|W_n\|_{H^1(\mathbb{R}^d)} = 1$, for any $n \in \mathbb{N}$. \square

Corollary 3 *Assume that $d \geq 2$, $p \in (2, 2^*)$ and $\theta \in [\vartheta(p, d), 1]$. There exists $\varepsilon = \varepsilon(\theta, p) > 0$ such that extremals of (6) are s -symmetric for every $0 < \Lambda < \varepsilon$.*

Proof Any sequence $(w_n)_n$ as in Proposition 3 violates (18) for n large enough, unless $\partial_\phi w_n \equiv 0$. The conclusion readily follows. \square

3.3 The Hardy regime: approaching $p = 2$

We proceed similarly as in Proposition 3 and Corollary 3.

Proposition 4 *Assume that $d \geq 2$, fix $\Lambda > 0$ and $\theta \in (0, 1]$. There exists $\eta \in (0, 4\theta/(d - 2\theta))$ such that all extremals of (6) are s -symmetric if $p \in (2, 2 + \eta)$.*

Proof The case $\theta = 1$ is already established in [6]. So, for fixed $\Lambda > 0$ and $0 < \theta < 1$, let w_n be an extremal of (6) with $p = p_n \rightarrow 2_+$. By Corollary 2, we know that $(t_n)_n$ is bounded and

$$\|\nabla w_n\|_{L^2(\mathcal{C})}^2 + \Lambda \|w_n\|_{L^2(\mathcal{C})}^2 = (t_n + \Lambda)^{1-\theta} \|w_n\|_{L^{p_n}(\mathcal{C})}^{p_n-2} \|w_n\|_{L^{p_n}(\mathcal{C})}^2.$$

First we prove that t_n converges to 0 as $n \rightarrow +\infty$. Assume by contradiction that $\lim_{n \rightarrow \infty} t_n = t > 0$ after extracting a subsequence if necessary, and choose $\varepsilon \in (0, 1/\Lambda)$ so that

$$(t + \Lambda)^\theta > \Lambda^\theta (\varepsilon t + 1). \quad (19)$$

Recalling that $C_{\text{CKN}}^*(\theta, p, \Lambda) \sim \Lambda^{-\theta}$ as $p \rightarrow 2_+$, we find that

$$\|w_n\|_{L^{p_n}(\mathcal{C})}^{p_n-2} = 1/C_{\text{CKN}}(\theta, p_n, \Lambda) \leq 1/C_{\text{CKN}}^*(\theta, p_n, \Lambda) \rightarrow \Lambda^\theta.$$

Using Lemma 2 to estimate $\|w_n\|_{L^{p_n}(\mathcal{C})}^2$ by $\varepsilon \|\nabla w\|_{L^2(\mathcal{C})}^2 + Z(\varepsilon, p) \|w\|_{L^2(\mathcal{C})}^2$, for n large enough, we get

$$\begin{aligned} (t_n + \Lambda) \|w_n\|_{L^2(\mathcal{C})}^2 &= \|\nabla w_n\|_{L^2(\mathcal{C})}^2 + \Lambda \|w_n\|_{L^2(\mathcal{C})}^2 = (t_n + \Lambda)^{1-\theta} \|w_n\|_{L^{p_n}(\mathcal{C})}^{p_n-2} \|w_n\|_{L^{p_n}(\mathcal{C})}^2 \\ &\leq (t_n + \Lambda)^{1-\theta} \Lambda^\theta (1 + o(1)) (\varepsilon t_n + Z(\varepsilon, p_n, d)) \|w_n\|_{L^2(\mathcal{C})}^2. \end{aligned}$$

Hence, by passing to the limit as $n \rightarrow \infty$, and using the fact that

$$\lim_{n \rightarrow \infty} Z(\varepsilon, p_n, d) = 1,$$

we deduce that

$$(t + \Lambda) \leq (t + \Lambda)^{1-\theta} \Lambda^\theta (\varepsilon t + 1)$$

in contradiction with (19). This proves that $\lim_{n \rightarrow +\infty} t_n = 0$.

Summarizing, w_n is a solution of

$$-\theta \Delta w_n + ((1 - \theta) t_n + \Lambda) w_n = (t_n + \Lambda)^{1-\theta} w_n^{p_n-1}$$

such that $t_n = \|\nabla w_n\|_{L^2(\mathcal{C})}^2 / \|w_n\|_{L^2(\mathcal{C})}^2 \rightarrow 0$ as $n \rightarrow +\infty$. Let $c_n := \|w_n\|_{L^{p_n}(\mathcal{C})}$ and $W_n := w_n / c_n$. We know that

$$c_n^{p_n-2} = \frac{1}{C_{\text{CKN}}(\theta, p_n, \Lambda)} \leq \frac{1}{C_{\text{CKN}}^*(\theta, p_n, \Lambda)} \rightarrow \Lambda^\theta$$

and

$$\|\nabla W_n\|_{L^2(\mathcal{C})}^2 + \Lambda \|W_n\|_{L^2(\mathcal{C})}^2 = (t_n + \Lambda)^{1-\theta} c_n^{p_n-2} \|W_n\|_{L^{p_n}(\mathcal{C})}^{p_n} = (t_n + \Lambda)^{1-\theta} c_n^{p_n-2}.$$

Hence we have

$$\lim_{n \rightarrow \infty} \|\nabla W_n\|_{L^2(\mathcal{C})}^2 + \Lambda \|W_n\|_{L^2(\mathcal{C})}^2 = \lim_{n \rightarrow \infty} (t_n + \Lambda)^{1-\theta} c_n^{p_n-2} \leq \Lambda.$$

Furthermore, from $\lim_{n \rightarrow \infty} t_n = 0$, we deduce that $\lim_{n \rightarrow \infty} \|\nabla W_n\|_{L^2(\mathcal{C})}^2 = 0$ and $\limsup_{n \rightarrow \infty} \|W_n\|_{L^2(\mathcal{C})}^2 \leq 1$. This proves that $(W_n)_n$ is bounded in $H^1(\mathcal{C})$ and that, up to subsequences, its weak limit is 0. By elliptic estimates and (16), we conclude that $\limsup_{n \rightarrow \infty} \|W_n\|_{L^\infty(\mathcal{C})} = 0$. Therefore, $\limsup_{n \rightarrow \infty} \|W_n\|_{L^\infty(\mathcal{C})}^{p_n-2} \leq 1$.

We can summarize the properties we have obtained so far for an extremal w_n of (6) with $p = p_n \rightarrow 2_+$ as follows:

$$\lim_{n \rightarrow \infty} t_n = 0 \quad \text{and} \quad \limsup_{n \rightarrow \infty} \|w_n\|_{L^{p_n}(\mathcal{C})}^{p_n-2} \leq \Lambda^\theta.$$

Incidentally, by means of the maximum principle for (15), we also get that

$$\|w_n\|_{L^\infty(\mathcal{C})}^{p_n-2} \geq \frac{(1-\theta)t_n + \Lambda}{(t_n + \Lambda)^{1-\theta}} \geq \Lambda^\theta,$$

which establishes that

$$\lim_{n \rightarrow \infty} \|w_n\|_{L^\infty(\mathcal{C})}^{p_n-2} = \Lambda^\theta.$$

Inequality (18) is clearly violated for n large enough unless $\partial_\phi w_n \equiv 0$. This concludes the proof. \square

3.4 A reformulation of Theorem 3 on the cylinder. Scalings and consequences

As in [6], it is convenient to rewrite Theorem 3 using the Emden-Fowler transformation.

Theorem 7 *For all $d \geq 2$, there exists a continuous function Λ^* defined on the set $\{(\theta, p) \in (0, 1] \times (2, 2^*) : \theta \geq \vartheta(p, d)\}$ with values in $(0, +\infty)$ such that $\lim_{p \rightarrow 2_+} \Lambda^*(\theta, p) = +\infty$ and*

- (i) *If $(\Lambda, p) \in (0, \Lambda^*(\theta, p)) \times (2, 2^*)$, then (1) has only s -symmetric extremals.*
- (ii) *If $\Lambda = \Lambda^*(\theta, p)$, then $\mathbf{C}_{\text{CKN}}(\theta, p, \Lambda) = \mathbf{C}_{\text{CKN}}^*(\theta, p, \Lambda)$.*
- (iii) *If $(\Lambda, p) \in (\Lambda^*(\theta, p), +\infty) \times (2, 2^*)$, none of the extremals of (1) is s -symmetric.*
- (iv) *$0 < \Lambda^*(\theta, p) \leq \underline{\Lambda}(\theta, p)$.*

Notice that s -symmetric and non s -symmetric extremals may coexist in case (ii). In (iv), we use the notation $\underline{\Lambda}(\theta, p) = (a_c - \underline{a}(\theta, p))^2$, where the function $\underline{a}(\theta, p)$ is defined in (3).

A key step for the proof of Theorem 7 relies on scalings in the s variable of the cylinder. If $w \in H^1(\mathcal{C}) \setminus \{0\}$, let $w_\sigma(s, \omega) := w(\sigma s, \omega)$ for $\sigma > 0$. A simple calculation shows that

$$\mathcal{F}_{\theta, p, \sigma^2 \Lambda}[w_\sigma] = \sigma^{2-\frac{1}{\theta} + \frac{2}{p\theta}} \mathcal{F}_{\theta, p, \Lambda}[w] - \sigma^{2-\frac{1}{\theta} + \frac{2}{p\theta}} (\sigma^2 - 1) \frac{\|\nabla_\omega w\|_{L^2(\mathcal{C})}^2 \|w\|_{L^2(\mathcal{C})}^{2\frac{1-\theta}{\theta}}}{\|w\|_{L^p(\mathcal{C})}^{2/\theta}}. \quad (20)$$

As a consequence, we observe that

$$\begin{aligned} \mathbf{C}_{\text{CKN}}^*(\theta, p, \sigma^2 \Lambda)^{-\frac{1}{\theta}} &= \mathcal{F}_{\theta, p, \sigma^2 \Lambda}[w_{\theta, p, \sigma^2 \Lambda}^*] \\ &= \sigma^{2-\frac{1}{\theta} + \frac{2}{p\theta}} \mathbf{C}_{\text{CKN}}^*(\theta, p, \Lambda)^{-\frac{1}{\theta}} = \sigma^{2-\frac{1}{\theta} + \frac{2}{p\theta}} \mathcal{F}_{\theta, p, \Lambda}[w_{\theta, p, \Lambda}^*]. \end{aligned}$$

Lemma 4 *If $d \geq 2$, $\Lambda > 0$ and $p \in (2, 2^*)$, then the following holds:*

- (i) *If $\mathbf{C}_{\text{CKN}}(\theta, p, \Lambda) = \mathbf{C}_{\text{CKN}}^*(\theta, p, \Lambda)$, then $\mathbf{C}_{\text{CKN}}(\theta, p, \lambda) = \mathbf{C}_{\text{CKN}}^*(\theta, p, \lambda)$ and, after a proper normalization, $w_{\theta, p, \lambda} = w_{\theta, p, \lambda}^*$ for any $\lambda \in (0, \Lambda)$.*

(ii) If there is an extremal $w_{\theta,p,\Lambda}$, which is not s -symmetric, even up to translations in the s -direction, then $C_{\text{CKN}}(\theta, p, \lambda) > C_{\text{CKN}}^*(\theta, p, \lambda)$ for all $\lambda > \Lambda$.

Recall that, according to [4], the extremal $w_{\theta,p,\lambda}^*$ among s -symmetric functions is uniquely defined up to translations in the s variable, multiplications by a constant and scalings with respect to s . We assume that it is normalized in such a way that it is uniquely defined. As for non s -symmetric minimizers, we have no uniqueness result. With a slightly loose notation, we shall write $w_{\theta,p,\lambda}$ for an extremal, but the reader has to keep in mind that, eventually, we pick *one* extremal among several, which are not necessarily related by one of the above transformations.

Proof To prove (i), apply (20) with $w_\sigma = w_{\theta,p,\lambda}$, $\lambda = \sigma^2 \Lambda$, $0 < \sigma < 1$ and $w(s, \omega) = w_{\theta,p,\lambda}(s/\sigma, \omega)$:

$$\begin{aligned} \frac{1}{C_{\text{CKN}}(\theta, p, \lambda)^{\frac{1}{\theta}}} &= \mathcal{F}_{\theta,p,\lambda}[w_{\theta,p,\lambda}] \\ &= \sigma^{2-\frac{1}{\theta}+\frac{2}{p\theta}} \mathcal{F}_{\theta,p,\Lambda}[w] + \sigma^{-\frac{1}{\theta}+\frac{2}{p\theta}} (1-\sigma^2) \frac{\|\nabla_\omega w\|_{L^2(\mathcal{C})}^2 \|w\|_{L^2(\mathcal{C})}^{2\frac{1-\theta}{\theta}}}{\|w\|_{L^p(\mathcal{C})}^{2/\theta}} \\ &\geq \frac{\sigma^{2-\frac{1}{\theta}+\frac{2}{p\theta}}}{C_{\text{CKN}}^*(\theta, p, \Lambda)^{\frac{1}{\theta}}} + \sigma^{-\frac{1}{\theta}+\frac{2}{p\theta}} (1-\sigma^2) \frac{\|\nabla_\omega w\|_{L^2(\mathcal{C})}^2 \|w\|_{L^2(\mathcal{C})}^{2\frac{1-\theta}{\theta}}}{\|w\|_{L^p(\mathcal{C})}^{2/\theta}} \\ &= \frac{1}{C_{\text{CKN}}^*(\theta, p, \lambda)^{\frac{1}{\theta}}} + \sigma^{-\frac{1}{\theta}+\frac{2}{p\theta}} (1-\sigma^2) \frac{\|\nabla_\omega w\|_{L^2(\mathcal{C})}^2 \|w\|_{L^2(\mathcal{C})}^{2\frac{1-\theta}{\theta}}}{\|w\|_{L^p(\mathcal{C})}^{2/\theta}}. \end{aligned}$$

By definition, $C_{\text{CKN}}(\theta, p, \lambda) \geq C_{\text{CKN}}^*(\theta, p, \lambda)$ and from the above inequality we find that necessarily $\nabla_\omega w \equiv 0$, and the first claim follows.

Assume that $w_{\theta,p,\Lambda}$ is an extremal with explicit dependence in ω and apply (20) with $w = w_{\theta,p,\Lambda}$, $w_\sigma(s, \omega) := w(\sigma s, \omega)$, $\lambda = \sigma^2 \Lambda$ and $\sigma > 1$:

$$\begin{aligned} \frac{1}{C_{\text{CKN}}(\theta, p, \lambda)^{\frac{1}{\theta}}} &\leq \mathcal{F}_{\theta,p,\sigma^2\Lambda}[w_\sigma] \\ &= \frac{\sigma^{2-\frac{1}{\theta}+\frac{2}{p\theta}}}{C_{\text{CKN}}(\theta, p, \lambda)^{\frac{1}{\theta}}} - \sigma^{-\frac{1}{\theta}+\frac{2}{p\theta}} (\sigma^2 - 1) \frac{\|\nabla_\omega w_{\theta,p,\Lambda}\|_{L^2(\mathcal{C})}^2 \|w_{\theta,p,\Lambda}\|_{L^2(\mathcal{C})}^{2\frac{1-\theta}{\theta}}}{\|w_{\theta,p,\Lambda}\|_{L^p(\mathcal{C})}^{2/\theta}} \\ &\leq \frac{\sigma^{2-\frac{1}{\theta}+\frac{2}{p\theta}}}{C_{\text{CKN}}^*(\theta, p, \Lambda)^{\frac{1}{\theta}}} - \sigma^{-\frac{1}{\theta}+\frac{2}{p\theta}} (\sigma^2 - 1) \frac{\|\nabla_\omega w_{\theta,p,\Lambda}\|_{L^2(\mathcal{C})}^2 \|w_{\theta,p,\Lambda}\|_{L^2(\mathcal{C})}^{2\frac{1-\theta}{\theta}}}{\|w_{\theta,p,\Lambda}\|_{L^p(\mathcal{C})}^{2/\theta}} \\ &< C_{\text{CKN}}^*(\theta, p, \lambda)^{-\frac{1}{\theta}}, \end{aligned}$$

since $\nabla_\omega w_{\theta,p,\Lambda} \not\equiv 0$. This proves the second claim. \square

By virtue of Corollary 3, we know that, for $p \in (2, 2^*)$ and $\vartheta(p, d) \leq \theta \leq 1$, the set $\{\Lambda > 0 : \mathcal{F}_{\theta,p,\Lambda}$ has only s -symmetric minimizers $\}$ is *not* empty, and hence we can define:

$$\Lambda^*(\theta, p) := \sup \{\Lambda > 0 : \mathcal{F}_{\theta,p,\Lambda} \text{ has only } s\text{-symmetric minimizers}\}.$$

In particular, by Proposition 1 (also see Section 2.2), Lemma 4 and Proposition 4, we have:

$$0 < \Lambda^*(\theta, p) \leq \underline{\Lambda}(\theta, p) \quad \text{and} \quad \lim_{p \rightarrow 2^+} \Lambda^*(\theta, p) = +\infty.$$

Corollary 4 *With the above definition of $\Lambda^*(\theta, p)$, we have:*

- (i) if $\lambda \in (0, \Lambda^*(\theta, p))$, then $\mathbf{C}_{\text{CKN}}(\theta, p, \lambda) = \mathbf{C}_{\text{CKN}}^*(\theta, p, \lambda)$ and, after a proper normalization, $w_{\theta, p, \lambda} = w_{\theta, p, \lambda}^*$,
- (ii) if $\lambda = \Lambda^*(\theta, p)$, then $\mathbf{C}_{\text{CKN}}(\theta, p, \lambda) = \mathbf{C}_{\text{CKN}}^*(\theta, p, \lambda)$,
- (iii) if $\lambda > \Lambda^*(\theta, p)$ and $\theta > \vartheta(p, d)$, then $\mathbf{C}_{\text{CKN}}(\theta, p, \lambda) > \mathbf{C}_{\text{CKN}}^*(\theta, p, \lambda)$.

Proof (i) is a consequence of Lemma 4 (i). It is easy to check that $\mathbf{C}_{\text{CKN}}(\theta, p, \lambda)$ is a non-increasing function of λ . By considering $\lim_{\lambda \rightarrow \Lambda^+} \mathcal{F}_{\theta, p, \Lambda}[w_{\theta, p, \lambda}^*]$, we get (ii). If $p \in (2, 2^*)$ and $\theta \in (\vartheta(p, d), 1]$, it has been shown in [5] that $\mathcal{F}_{\theta, p, \Lambda}$ always attains its minimum in $H^1(\mathcal{C}) \setminus \{0\}$, so that (iii) follows from Lemma 4 (ii). \square

3.5 The proof of Theorem 7

In case $\theta = \vartheta(p, d)$, extremals might not exist: see [5]. To complete the proof of Theorem 7, we have to prove that the property of Lemma 4 (iii) also holds if $\theta = \vartheta(p, d)$ and to establish the continuity of Λ^* .

Lemma 5 *If $\lambda > \Lambda^*(\theta, p)$ and $\theta = \vartheta(p, d)$, then $\mathbf{C}_{\text{CKN}}(\theta, p, \lambda) > \mathbf{C}_{\text{CKN}}^*(\theta, p, \lambda)$.*

Proof Consider the Gagliardo-Nirenberg inequality

$$\|u\|_{L^p(\mathbb{R}^d)}^2 \leq \mathbf{C}_{\text{GN}}(p) \|\nabla u\|_{L^2(\mathbb{R}^d)}^{2\vartheta(p, d)} \|u\|_{L^2(\mathbb{R}^d)}^{2(1-\vartheta(p, d))} \quad \forall u \in H^1(\mathbb{R}^d) \quad (21)$$

and assume that $\mathbf{C}_{\text{GN}}(p)$ is the optimal constant. According to [5] (see Lemma 6 below for more details), we know that

$$\mathbf{C}_{\text{GN}}(p) \leq \mathbf{C}_{\text{CKN}}(\vartheta(p, d), p, \lambda).$$

According to [5, Theorem 1.4 (i)] there are extremals for (6) with $\theta = \vartheta(p, d)$, $p \in (2, 2^*)$ and $\lambda > 0$, whenever the above inequality is strict. By Corollary 4 (ii) we know that $\mathbf{C}_{\text{GN}}(p) \leq \mathbf{C}_{\text{CKN}}^*(\vartheta(p, d), p, \lambda)$ if $\lambda = \Lambda^*(\vartheta(p, d), p)$.

Case 1: Assume that $\mathbf{C}_{\text{GN}}(p) = \mathbf{C}_{\text{CKN}}^*(\vartheta(p, d), p, \Lambda^*(\vartheta(p, d), p))$. Then for all $\lambda > \Lambda^*(\vartheta(p, d), p)$,

$$\mathbf{C}_{\text{CKN}}^*(\vartheta(p, d), p, \lambda) < \mathbf{C}_{\text{GN}}(p) \leq \mathbf{C}_{\text{CKN}}(\vartheta(p, d), p, \lambda)$$

because $\mathbf{C}_{\text{CKN}}^*(\theta, p, \lambda)$ is decreasing in λ , which proves the result.

Case 2: Assume that $\mathbf{C}_{\text{GN}}(p) < \mathbf{C}_{\text{CKN}}^*(\vartheta(p, d), p, \Lambda^*(\vartheta(p, d), p))$. We can always choose $\lambda > \Lambda^*(\vartheta(p, d), p)$, sufficiently close to $\Lambda^*(\vartheta(p, d), p)$, so that

$$\mathbf{C}_{\text{GN}}(p) < \mathbf{C}_{\text{CKN}}^*(\vartheta(p, d), p, \lambda) \leq \mathbf{C}_{\text{CKN}}(\vartheta(p, d), p, \lambda).$$

Then [5, Theorem 1.4 (i)] ensures the existence of an extremal $w_{\theta, p, \lambda}$ of (6) with $\theta = \vartheta(p, d)$. By the definition of $\Lambda^*(\vartheta(p, d), p)$, such an extremal is non s -symmetric. The result follows from Lemma 4 (ii). \square

To complete the proof of Theorem 7, we only need to establish the continuity of Λ^* with respect to the parameters (θ, p) with $p \in (2, 2^*)$ and $\vartheta(p, d) \leq \theta < 1$. The argument is similar to the one used in [6] for the case $\theta = 1$. First of all, by using the definition of $\Lambda^*(\theta, p)$, Lemma 4 (i) and the s -symmetric extremals, it is easy to see that, for any sequences $(\theta_n)_n$ and $(p_n)_n$ such that $\theta_n \rightarrow \theta$ and $p_n \rightarrow p \in (2, 2^*)$,

$$\limsup_{n \rightarrow +\infty} \Lambda^*(\theta_n, p_n) \leq \Lambda^*(\theta, p).$$

To see that equality actually holds, we argue by contradiction and assume that for a given sequence $\theta_n \in [\vartheta(p_n, d), 1]$ and $p_n \in (2, 2^*)$, we have:

$$\Lambda_\infty := \lim_{n \rightarrow +\infty} \Lambda^*(\theta_n, p_n) < \Lambda^*(\theta, p).$$

For n large, fix λ such that $\Lambda^*(\theta_n, p_n) < \lambda < \Lambda^*(\theta, p) \leq \underline{\Lambda}(\theta, p)$.

If $\theta > \vartheta(p, d)$, then $\theta_n > \vartheta(p_n, d)$ for n large, and we find a sequence of *non s -symmetric* extremals $w_{\theta_n, p_n, \lambda}$ that, along a subsequence, must converge to an s -symmetric extremal $w_{\theta, p, \Lambda}^*$, a contradiction with $\lambda < \underline{\Lambda}(\theta, p)$ as already noted in the introduction.

If $\theta = \vartheta(p, d)$, then, by strict monotonicity of $\mathbf{C}_{\text{CKN}}^*$ with respect to λ , we find: $\mathbf{C}_{\text{GN}}(p) \leq \mathbf{C}_{\text{CKN}}^*(\vartheta(p, d), p, \Lambda^*(p, d)) < \mathbf{C}_{\text{CKN}}^*(\vartheta(p, d), p, \lambda)$ and so, for n sufficiently large: $\mathbf{C}_{\text{GN}}(p) < \mathbf{C}_{\text{CKN}}^*(\theta_n, p_n, \lambda) \leq \mathbf{C}_{\text{CKN}}(\theta_n, p_n, \lambda)$. Again by [5, Theorem 1.4 (i)], there exist non s -symmetric extremals $w_{\theta_n, p_n, \lambda}$ of (6) relative to the parameters (θ_n, p_n, λ) , that, along a subsequence, must converge to an extremal of (6) relative to the parameters (θ, p, Λ) . Since $\lambda < \Lambda^*(\theta, p)$, the limiting extremal must be s -symmetric and we obtain a contradiction as above. This completes the proof of Theorem 7. \square

Remark 6 As already noticed above, at $\Lambda = \Lambda^*(\theta, p)$, we have

$$\mathbf{C}_{\text{CKN}}(\theta, p, \Lambda^*(\theta, p)) = \mathbf{C}_{\text{CKN}}^*(\theta, p, \Lambda^*(\theta, p))$$

and, as long as there are extremal functions, either $\Lambda^*(\theta, p) = \underline{\Lambda}(\theta, p)$, or a s -symmetric extremal and a non s -symmetric one may coexist. This is precisely what occurs in the framework of Theorem 5, at least for $\theta > \vartheta(p, d)$.

4 Radial symmetry for the weighted logarithmic Hardy inequalities

As in Section 3.4, we rephrase Theorem 4 on the cylinder.

Theorem 8 *For all $d \geq 2$, there exists a continuous function Λ^{**} defined on the set $\{\gamma > d/4\}$ and with values in $(0, +\infty)$ such that for all $\Lambda \in (0, \Lambda^{**}(\gamma)]$, there is an s -symmetric extremal of (2), while for any $\Lambda > \Lambda^{**}(\gamma)$, no extremal of (7) is s -symmetric. Moreover, $\Lambda^{**}(\gamma) \leq \frac{1}{4}(4\gamma - 1)(d - 1) = \tilde{\Lambda}(\gamma)$.*

4.1 The critical regime: approaching $\Lambda = 0$

In order to prove the above theorem, we first start by showing that for $\gamma > d/4$ and Λ close to 0, the extremals for (7) are s -symmetric. From [5, Theorem 1.3 (ii)], we know that such extremals exist.

Proposition 5 *Let $\gamma > d/4$ and $d \geq 2$. Then, for $\Lambda > 0$ sufficiently small, any extremal $w_{\gamma, \Lambda}$ of (7) is s -symmetric.*

Proof Let us consider $\gamma > d/4$ and a sequence of positive numbers $(\Lambda_n)_n$ converging to 0. Let us denote by $(w_n)_n$ a sequence of extremals for (7) with parameter Λ_n . For simplicity, let us normalize the functions w_n so that $\|w_n\|_{L^2(\mathcal{C})} = 1$. Moreover, we can assume that $w_n = w_n$ depends only on s and the azimuthal angle $\phi \in \mathbb{S}^{d-1}$ and $\max_{\mathcal{C}} w_n = w_n(0, \phi_0)$ for some $\phi_0 \in [0, \pi]$. Finally, w_n is a minimum for $\mathcal{G}_{\theta, p, \Lambda}$ defined in (9), and we have $\mathcal{G}_{\theta, p, \Lambda}[w_n] = 1/C_n$ with $C_n := C_{\text{WLH}}(\gamma, \Lambda_n)$ for any $n \in \mathbb{N}$. Therefore w_n satisfies the Euler-Lagrange equation

$$-\Delta w_n - C_n^{-1} w_n (1 + \log |w_n|^2) \exp\left(\frac{1}{2\gamma} \int_{\mathcal{C}} |w_n|^2 \log |w_n|^2 dy\right) = \mu_n w_n \quad (22)$$

for some $\mu_n \in \mathbb{R}$. Multiplying this equation by w_n and integrating by parts we get

$$\|\nabla w_n\|_{L^2(\mathcal{C})}^2 - C_n^{-1} \exp\left(\frac{1}{2\gamma} \int_{\mathcal{C}} |w_n|^2 \log |w_n|^2 dy\right) \int_{\mathcal{C}} w_n^2 (1 + \log |w_n|^2) dy = \mu_n. \quad (23)$$

In addition, the condition $\mathcal{G}_{\theta, p, \Lambda}[w_n] = 1/C_n$ gives

$$\|\nabla w_n\|_{L^2(\mathcal{C})}^2 + \Lambda_n = C_n^{-1} \exp\left(\frac{1}{2\gamma} \int_{\mathcal{C}} |w_n|^2 \log |w_n|^2 dy\right).$$

As in [5], consider Hölder's inequality, $\|w\|_{L^q(\mathcal{C})} \leq \|w\|_{L^2(\mathcal{C})}^{\zeta} \|w\|_{L^p(\mathcal{C})}^{1-\zeta}$ with $\zeta = 2(p-q)/(q(p-2))$ for any q such that $2 \leq q \leq p \leq 2^*$. For $q = 2$, this inequality becomes an equality, with $\zeta = 1$, so that we can differentiate with respect to q at $q = 2$ and obtain

$$\int_{\mathcal{C}} |w|^2 \log\left(\frac{|w|^2}{\|w\|_{L^2(\mathcal{C})}^2}\right) dy \leq \frac{p}{p-2} \|w\|_{L^2(\mathcal{C})}^2 \log\left(\frac{\|w\|_{L^p(\mathcal{C})}^2}{\|w\|_{L^2(\mathcal{C})}^2}\right).$$

Let $C_{\text{GN}}(p)$ be the best constant in (21). Combining the two inequalities, we obtain the following *logarithmic Sobolev inequality* on the cylinder: for all $d \geq 1$,

$$\int_{\mathcal{C}} w^2 \log\left(\frac{w^2}{\|w\|_{L^2(\mathcal{C})}^2}\right) dy \leq \frac{d}{2} \|w\|_{L^2(\mathcal{C})}^2 \log\left(\frac{\|\nabla w\|_{L^2(\mathcal{C})}^2}{\|w\|_{L^2(\mathcal{C})}^2}\right) + K(d) \|w\|_{L^2(\mathcal{C})}^2, \quad (24)$$

where

$$K(d) := \inf_{p \in (2, 2^*)} \frac{p}{p-2} C_{\text{GN}}(p).$$

See [4, Lemma 5] for more details and a sharp version, but not in Weissler's logarithmic form as it is here, of the logarithmic Sobolev inequality on the cylinder.

Applying this inequality to w_n , we obtain

$$\|\nabla w_n\|_{L^2(\mathcal{C})}^2 + \Lambda_n \leq C_n^{-1} e^{\frac{K(d)}{2\gamma}} \left(\|\nabla w_n\|_{L^2(\mathcal{C})}^2\right)^{\frac{d}{4\gamma}}.$$

Since $\gamma > d/4$, $\Lambda_n \rightarrow 0$ and $C_n \rightarrow +\infty$ (see [4, Theorem B']), we see that $(\nabla w_n)_n$ converges to 0 as $n \rightarrow +\infty$. On the other hand, $\|w_n\|_{L^2(\mathcal{C})} = 1$, so, up to subsequences, $(w_n)_n$ converges weakly and in $C_{\text{loc}}^{2, \alpha}$ to $w \equiv 0$.

Now, like in the proofs of Corollary 3 by using (22), we see that the function $\chi_n := D_\phi w_n$ satisfies:

$$\begin{aligned} & \int_{\mathcal{C}} (|\partial_s \chi_n|^2 + |\partial_\phi \chi_n|^2) dy \\ & - C_n^{-1} \exp\left(\frac{1}{2\gamma} \int_{\mathcal{C}} |w_n|^2 \log |w_n|^2 dy\right) \int_{\mathcal{C}} |\chi_n|^2 (3 + 2 \log |w_n|^2) dy \\ & = \mu_n \|\chi_n\|_{L^2(\mathcal{C})}^2. \end{aligned} \quad (25)$$

Hence, by means of the Poincaré inequality we derive

$$\begin{aligned} & (d-1-\mu_n) \|\chi_n\|_{L^2(\mathcal{C})}^2 \\ & \leq C_n^{-1} \exp\left(\frac{1}{2\gamma} \int_{\mathcal{C}} |w_n|^2 \log |w_n|^2 dy\right) \int_{\mathcal{C}} |\chi_n|^2 (3 + 2 \log |w_n|^2) dy \\ & \leq C_n^{-1} \exp\left(\frac{1}{2\gamma} \int_{\mathcal{C}} |w_n|^2 \log |w_n|^2 dy\right) \|\chi_n\|_{L^2(\mathcal{C})}^2 (3 + 2 \log (\|w_n\|_{L^\infty(\mathcal{C})})) \leq 0 \end{aligned}$$

for n large, since $\|w_n\|_{L^\infty(\mathcal{C})}$ converges to 0 as $n \rightarrow +\infty$. Next observe that by the strong convergence of $(\nabla w_n)_n$ to 0 in $L^2(\mathbb{R}^d)$ (23) and by the logarithmic Sobolev inequality (24), we obtain $\lim_{n \rightarrow \infty} \mu_n = 0$. So, necessarily $\chi_n \equiv 0$ for n large and the proof is complete. \square

4.2 The proof of Theorem 8

Consider the functional $\mathcal{G}_{\gamma, \Lambda}$ defined in (9). If $w \in H^1(\mathcal{C}) \setminus \{0\}$, let $w_\sigma(s, \omega) := w(\sigma s, \omega)$ for any $\sigma > 0$. A simple calculation shows that for all $\sigma > 0$,

$$\mathcal{G}_{\gamma, \sigma^2 \Lambda}[w_\sigma] = \sigma^{2-\frac{1}{2\gamma}} \mathcal{G}_{\gamma, \Lambda}[w] - \frac{(\sigma^2 - 1) \sigma^{-\frac{1}{2\gamma}} \|\nabla_\omega w\|_{L^2(\mathcal{C})}^2}{\|w\|_{L^2(\mathcal{C})}^2 \exp\left\{\frac{1}{2\gamma} \int_{\mathcal{C}} \frac{w^2}{\|w\|_{L^2(\mathcal{C})}^2} \log\left(\frac{w^2}{\|w\|_{L^2(\mathcal{C})}^2}\right) dy\right\}}.$$

The above expression is the counterpart of (20) in the case of the weighted logarithmic Hardy inequality and we can even observe that $2 - \frac{1}{2\gamma} = \lim_{p \rightarrow 2^+} (2 - \frac{1}{\theta} + \frac{2}{p\theta})$ when $\theta = \gamma(p-2)$. We use it exactly as in Section 3.4 to prove that for any $d \geq 2$, $\Lambda > 0$ and $\gamma > d/4$, the following properties hold:

- (i) If $C_{\text{CKN}}(\gamma, \Lambda) = C_{\text{WLH}}^*(\gamma, \Lambda)$, then $C_{\text{CKN}}(\gamma, \lambda) = C_{\text{WLH}}^*(\gamma, \lambda)$ and, after a proper normalization, $w_{\gamma, \lambda} = w_{\gamma, \lambda}^*$, for any $\lambda \in (0, \Lambda)$.
- (ii) If there is there is an extremal $w_{\gamma, \Lambda}$, which is not s -symmetric, even up to translations in the s -direction, then $C_{\text{WLH}}(\gamma, \lambda) > C_{\text{WLH}}^*(\gamma, \lambda)$ for all $\lambda > \Lambda$.

At this point, in view of Proposition 5 and by recalling the role of the function \tilde{a} in (4), we can argue as in Section 3.5 to prove the existence of a *continuous* function Λ^{**} defined on $(d/4, \infty)$, such that

- (i) $0 < \Lambda^{**}(\gamma) < \tilde{\Lambda}(\gamma)$,
- (ii) if $\lambda \in (0, \Lambda^*(\gamma))$, then $C_{\text{WLH}}(\gamma, \lambda) = C_{\text{WLH}}^*(\gamma, \lambda)$ and, after a proper normalization, $w_{\gamma, \lambda} = w_{\gamma, \lambda}^*$,
- (iii) if $\lambda = \Lambda^{**}(\gamma)$, then $C_{\text{WLH}}(\gamma, \lambda) = C_{\text{WLH}}^*(\gamma, \lambda)$,
- (iv) if $\lambda > \Lambda^{**}(\gamma)$, then $C_{\text{WLH}}(\gamma, \lambda) > C_{\text{WLH}}^*(\gamma, \lambda)$.

This concludes the proof of Theorems 8. \square

5 New symmetry breaking results

This section is devoted to the proof of Theorems 5 and 6. We prove symmetry breaking in the range of parameters where the radial extremal is a strict, local minimum for the variational problem associated to inequalities (1) and (2). Consider the optimal constants in the limit cases given respectively by $\theta = \vartheta(p, d)$ and $\gamma = d/4$. We recall that

$$\frac{1}{C_{\text{CKN}}(\vartheta(p, d), p, \Lambda)} = \inf_{u \in D_a^{1,2}(\mathbb{R}^d) \setminus \{0\}} \frac{\| |x|^{-a} \nabla u \|_{L^2(\mathbb{R}^d)}^2 \| |x|^{-(a+1)} u \|_{L^2(\mathbb{R}^d)}^2}{\| |x|^{-b} u \|_{L^p(\mathbb{R}^d)}^2}$$

and

$$\frac{1}{C_{\text{WLH}}(d/4, \Lambda)} = \inf \| |x|^{-a} \nabla u \|_{L^2(\mathbb{R}^d)}^2 \exp \left[-\frac{2}{d} \int_{\mathbb{R}^d} \frac{|u|^2}{|x|^{2(a+1)}} \log(|x|^{2(a_c-a)} |u|^2) dx \right]$$

where the last infimum is taken on the set of the functions $u \in D_a^{1,2}(\mathbb{R}^d)$ such that $\| |x|^{-(a+1)} u \|_{L^2(\mathbb{R}^d)} = 1$. We also define the best constants in Gagliardo-Nirenberg and logarithmic Sobolev inequalities respectively by

$$\frac{1}{C_{\text{GN}}(p)} := \inf_{u \in H^1(\mathbb{R}^d) \setminus \{0\}} \frac{\| \nabla u \|_{L^2(\mathbb{R}^d)}^2 \| u \|_{L^2(\mathbb{R}^d)}^2}{\| u \|_{L^p(\mathbb{R}^d)}^2}$$

and

$$\frac{1}{C_{\text{LS}}} := \inf_{\substack{u \in H^1(\mathbb{R}^d) \\ \| u \|_{L^2(\mathbb{R}^d)} = 1}} \int_{\mathbb{R}^d} |\nabla u|^2 dx \exp \left[-\frac{2}{d} \int_{\mathbb{R}^d} |u|^2 \log |u|^2 dx \right]$$

It is well known (see for instance [15]) that $C_{\text{LS}} = \frac{2}{\pi d e}$.

Lemma 6 *Let $d \geq 3$ and $p \in (2, 2^*)$. For all $a < a_c$, we have*

$$C_{\text{GN}}(p) \leq C_{\text{CKN}}(\vartheta(p, d), p, \Lambda) \quad \text{and} \quad C_{\text{LS}} \leq C_{\text{WLH}}(d/4, \Lambda).$$

If $d = 2$, the first inequality still holds while the second one is replaced by $C_{\text{LS}} \leq \limsup_{\gamma \rightarrow (1/2)_+} C_{\text{WLH}}(\gamma, \Lambda)$.

Proof Consider an extremal u for either the Gagliardo-Nirenberg or the logarithmic Sobolev inequality. It is known that such a solution exists, is unique up to multiplication by constants, translations and scalings (in case of the logarithmic Sobolev inequalities, take for instance $u(x) = (2\pi)^{-d/4} \exp(-|x|^2/4)$ for any $x \in \mathbb{R}^d$). Let $\mathbf{e} \in \mathbb{S}^{d-1}$ and use $u_n(x) := u(x + n\mathbf{e})$, $n \in \mathbb{N}$, as a sequence of test functions for the quotients defining $C_{\text{CKN}}(\vartheta(p, d), p, \Lambda)$ and $C_{\text{WLH}}(d/4, \Lambda)$ respectively. We first use the reformulation of (1) used in Section 2.3 in terms of $v(x) = |x|^{-a} u(x)$, and observe that, for $\theta = \vartheta(p, d) =$

$1 - (b - a)$, $d \geq 2$, we have

$$\begin{aligned}
& \frac{1}{\mathbf{C}_{\text{CKN}}(\theta, p, A)} \\
&= \inf_{v \in H^1(\mathbb{R}^d) \setminus \{0\}} \frac{\left(\|\nabla v\|_{L^2(\mathbb{R}^d)}^2 + a(a - 2a_c) \| |x|^{-1} v \|_{L^2(\mathbb{R}^d)}^2 \right)^\theta \| |x|^{-1} v \|_{L^2(\mathbb{R}^d)}^{2(1-\theta)}}{\| |x|^{\theta-1} v \|_{L^p(\mathbb{R}^d)}^2} \\
&\leq \frac{\left(\|\nabla u_n\|_{L^2(\mathbb{R}^d)}^2 + a(a - 2a_c) \| |x|^{-1} u_n \|_{L^2(\mathbb{R}^d)}^2 \right)^\theta \| |x|^{-1} u_n \|_{L^2(\mathbb{R}^d)}^{2(1-\theta)}}{\| |x|^{\theta-1} u_n \|_{L^p(\mathbb{R}^d)}^2} \\
&= \frac{\left(\|\nabla u\|_{L^2(\mathbb{R}^d)}^2 + \frac{a(a-2a_c)}{n^2} \| |\frac{x}{n} - e|^{-1} u \|_{L^2(\mathbb{R}^d)}^2 \right)^\theta \| |\frac{x}{n} - e|^{-1} u \|_{L^2(\mathbb{R}^d)}^{2(1-\theta)}}{\| |\frac{x}{n} - e|^{1-\theta} u \|_{L^p(\mathbb{R}^d)}^2} \\
&\xrightarrow{n \rightarrow +\infty} \frac{\left(\|\nabla u\|_{L^2(\mathbb{R}^d)}^2 \right)^\theta \| u \|_{L^2(\mathbb{R}^d)}^{2(1-\theta)}}{\| u \|_{L^p(\mathbb{R}^d)}^2} = \frac{1}{\mathbf{C}_{\text{GN}}(p)}.
\end{aligned}$$

The inequality $\mathbf{C}_{\text{LS}} \leq \mathbf{C}_{\text{WLH}}(d/4, A)$ follows from a similar computation if $d \geq 3$. If $d = 2$, it is enough to repeat the computation for a well chosen sequence $(\gamma_n)_n$ such that $\gamma_n > 1/2$ for any $n \in \mathbb{N}$ and $\lim_{n \rightarrow \infty} \gamma_n = 1/2$. \square

Proof of Theorem 5. Let $\mathbf{g}(x) := (2\pi)^{-d/4} \exp(-|x|^2/4)$ for any $x \in \mathbb{R}^d$ and consider the function

$$h(p, d) := \frac{\|\nabla \mathbf{g}\|_{L^2(\mathbb{R}^d)}^{2\vartheta(p,d)} \|\mathbf{g}\|_{L^2(\mathbb{R}^d)}^{2(1-\vartheta(p,d))}}{\|\mathbf{g}\|_{L^p(\mathbb{R}^d)}^2}.$$

A tedious but elementary computation provides an explicit value for $h(p, d)$ in terms of Γ functions, that can be used to get the estimate

$$\frac{1}{\mathbf{C}_{\text{CKN}}(\vartheta(p, d), p, A(a_-(p)))} \leq \frac{1}{\mathbf{C}_{\text{GN}}(p)} \leq h(p, d)$$

where $a_-(p) = \underline{a}(\vartheta(p, d), p)$. Consider the function

$$\mathbf{L}(p, d) := h(p, d) \mathbf{C}_{\text{CKN}}^*(\vartheta(p, d), p, A(a_-(p))).$$

Explicit computations show that $\lim_{p \rightarrow 2^+} \mathbf{L}(p, d) = 1$ and $\ell(d) := \lim_{p \rightarrow 2^+} \frac{\partial \mathbf{L}}{\partial p}(p, d)$ is an increasing function of d such that $\lim_{d \rightarrow \infty} \ell(d) = -\frac{1}{4} \log 2 < 0$. Hence, for any given $d \geq 2$, there exists an $\eta > 0$ such that $\mathbf{L}(p, d) < 1$ for any $p \in (2, 2 + \eta)$. See Fig. 2. As a consequence, we have

$$h(p, d) < \frac{1}{\mathbf{C}_{\text{CKN}}^*(\vartheta(p, d), p, A(a_-(p)))}$$

provided $0 < p - 2 < \eta$, with η small enough, thus proving that $\mathbf{C}_{\text{CKN}}^*(\theta, p, A) < \mathbf{C}_{\text{CKN}}(\theta, p, A)$ if $\theta = \vartheta(p, d)$ and $a = a_-(p)$. By continuity and according to Theorem 7 (ii), the strict inequality also holds for θ close to $\vartheta(p, d)$ and a close to $a_-(p)$, as claimed. \square

Proof of Theorem 6. For the weighted logarithmic Hardy inequality (2), the same method applies. From the explicit estimates of \mathbf{C}_{LS} and $\mathbf{C}_{\text{WLH}}^*(\gamma, A)$, it is a tedious

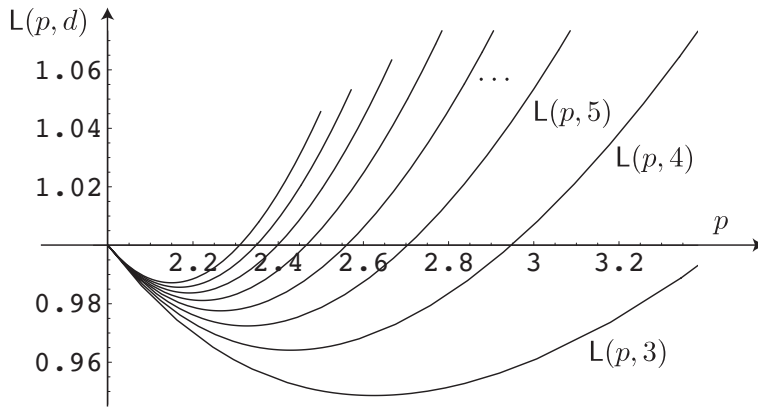


Fig. 2 Plots of $L(p, d)$ as a function of p for $d = 3, \dots, 10$.

but straightforward computation to check that $C_{\text{WLH}}^*(\gamma, \Lambda) < C_{\text{LS}}$ if and only if $\Lambda(a) > \Lambda_{\text{SB}}(\gamma, d)$, where Λ_{SB} has been defined in (5). As a special case, notice that $C_{\text{WLH}}^*(d/4, \Lambda(-1/2)) < C_{\text{LS}}$ if $d \geq 3$, while, for $d = 2$, we have:

$$\lim_{\gamma \rightarrow (1/2)_+} C_{\text{WLH}}^*(\gamma, \Lambda(-1/2)) < C_{\text{LS}}.$$

See Fig. 3. By continuity, the inequality $C_{\text{WLH}}^*(\gamma, \tilde{\Lambda}(\gamma)) < C_{\text{LS}}$ remains valid for $\gamma > d/4$, provided $\gamma - d/4 > 0$ is small enough. This completes the proof. \square

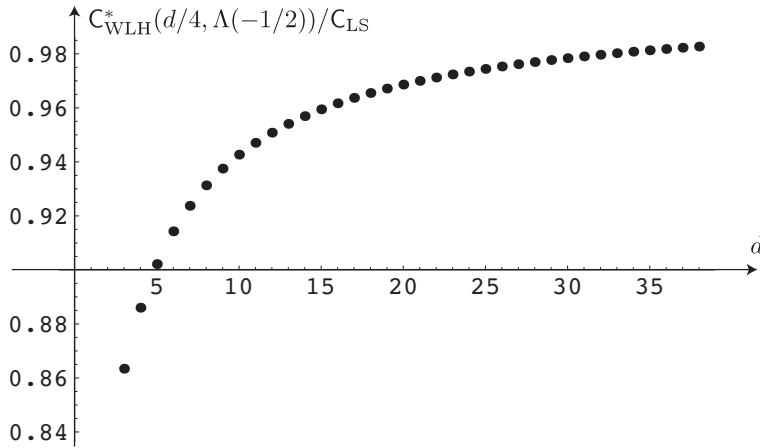


Fig. 3 Plot of $C_{\text{WLH}}^*(d/4, \Lambda(-1/2))/C_{\text{LS}}$ in terms of $d \in \mathbb{N}$, $d \geq 3$.

Remark 7 The condition $C_{\text{LS}} < C_{\text{WLH}}^*(d/4, \Lambda)$ amounts to $a \in (a_*, a_c)$ for some explicit a_* and from [5, Theorem 1.4] we know that this is a sufficient condition for the existence of an extremal function for (2). The symmetry breaking results of Theorem 6

hold for any $a \in (-\infty, a_*)$. In that case, the existence of an extremal for (2) is not known if $\gamma = d/4$, $d \geq 3$, but it is granted by [5, Theorem 1.3] for any $\gamma > d/4$, $d \geq 2$.

Compared with the result in Proposition 2, we see by numerical calculations that $\Lambda(a) > \tilde{\Lambda}(\gamma)$ is more restrictive than $\Lambda(a) > \Lambda_{\text{SB}}(\gamma, d)$ except if $d = 2$ and $\gamma \in [0.621414 \dots, 6.69625 \dots]$, $d = 3$ and $\gamma \in [0.937725 \dots, 4.14851 \dots]$, or $d = 4$ and $\gamma \in [1.31303 \dots, 2.98835 \dots]$. For $d \geq 5$, we observe that $\Lambda_{\text{SB}}(\gamma, d) < \tilde{\Lambda}(\gamma)$. See Fig. 4.

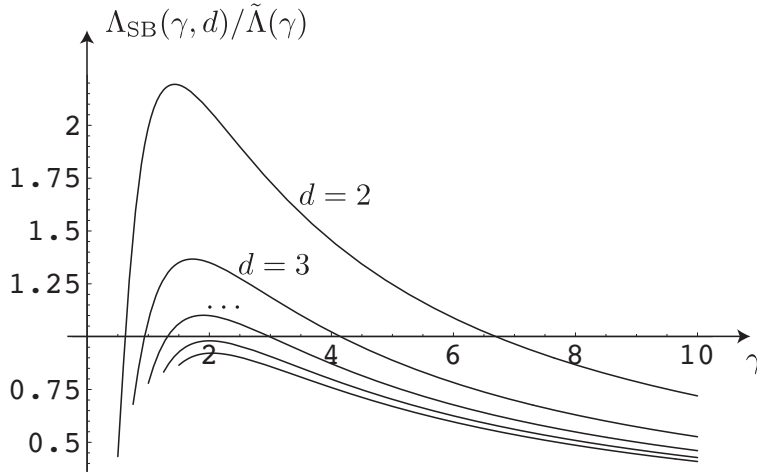


Fig. 4 Plot of $\Lambda_{\text{SB}}(\gamma, d)/\tilde{\Lambda}(\gamma)$ as a function of γ , for $d = 2, 3, \dots, 6$.

As a concluding remark for the weighted logarithmic Hardy inequality, we emphasize the fact that, in many cases, the comparison with the logarithmic Sobolev inequality gives better informations about the symmetry breaking properties of the extremals than methods based on a linearization approach.

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