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Numerical treatment of rounded and sharp corners in the modeling of 2D electrostatic fields

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***Abstract*—This work deals with numerical techniques to compute electrostatic fields in devices with rounded corners in 2D situations. The approach leads to the solution of two problems: one on the device where rounded corners are replaced by sharp corners and the other on an unbounded domain representing the shape of the rounded corner after an appropriate rescaling. Details are given on several ways to solve both problems and numerical results are provided to assess the efficiency and the accuracy of the techniques.**

I. Introduction

The precise description of an object containing *rounded corners* leads to consider meshes with a large number of nodes when the finite element method (FEM) is straightforwardly applied. Dealing with such meshes makes the computation time- and resource-consuming. Moreover, these computations have to be repeated if the curvature radius of the rounded corner is modified.

In order to avoid this computational cost, the rounded corners are usually replaced by *sharp corners*. The obtained computational results are then "globally correct" but locally inaccurate in the neighborhood of the corners. We aim at remedying this drawback.

In this work, we are dealing with a 2D electrostatic problem in a domain with a rounded corner (see Fig. 1(a)). The proposed method to approximate the electric scalar potential is based on two observations, which can be checked by simple numerical experiments:

- the exact solutions *close to the corner*, computed for several values of the curvature radius ε, are quasi-similar, up to a "scaling factor" (related to ε). It is also noticed that the "shape" of the solutions *close to the corner* (their "shape" but not their amplitude) weakly depend on other elements of the studied structure, such as the distance to the boundaries: it is said that the dominant term of the solutions close to the corner are *self-similar*.
- the exact solutions *far from the corner* are weakly influenced by the change of the curvature radius ε, and they converge to the solution on the domain with the sharp corner when ε goes to zero.

It can be deduced that an accurate approximate solution of the exact solution for any curvature radius ε can be build from the solutions of two problems:

- an electrostatic problem on the real domain except that the rounded corner is replaced by a sharp corner;
- an electrostatic problem on a domain that only takes the shape of the rounded corner into account. The corresponding solution will be used with an appropriate "scaling".

This intuitive argument coincides with the mathematical results proposed partially by Timouyas [1] in 2003 and recently completed by Dauge and her collaborators [2]. In addition to the predictable behavior of the solution, some error estimates are also provided in [1], [2] for the approximate solutions we consider.

This work is addressed to the community of engineers that use numerical methods (as it was done in [3]), and the aim is to illustrate the finite element implementation of the theoretical concepts which are proved in other references [4], [2].

In Section II are given the principles of the method. Then, we provide explicitly how to solve the different problems, that enable to build the approximate solution:

- the computation of the "singularity factor" of the solution on the domain with the sharp angle, in Section III;
- the problem on the unbounded domain taking only into account the shape of the rounded corner, in Section IV.

In Sections II, III and IV, numerical results are given to assess the accuracy of the proposed techniques.

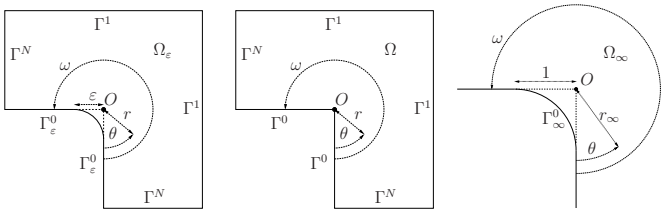
II. Principles of the method

For the sake of clarity, the method is explained on a specific structure, see Fig. 1, but it can be straightforwardly extended to other structures, in particular including more than two electric conductors.

A. Definitions of the considered problems

Denote by v_ε the harmonic function $-\triangle v_\varepsilon = 0$– in the domain Ω_ε equal to zero on Γ^0_ε, to V^1 on Γ^1 and satisfying homogeneous Neumann conditions on Γ^N; see Fig.1(a). The factor ε defines the true size of the rounded corner, starting from the profile domain, Fig. 1(c) where the characteristic size of the rounded corner is 1.

Let also v denote the solution of $\triangle v = 0$ in the domain Ω with the sharp corner, see Fig. 1(b). The function v is enforced to be equal to zero on Γ^0 and to V^1 on Γ^1, and v verifies homogeneous Neumann conditions on Γ^N. In the



(a) Domain with a rounded corner Ω_ε.

(b) Domain with a sharp corner Ω.

(c) Unbounded profile domain Ω_∞.

Fig. 1. Considered domains Ω, Ω_∞ and Ω_ε.

neighborhood of the sharp corner[1], v can be expanded in series as follows [5]:

$$v(r,\theta) = \lambda r^\alpha \sin(\alpha\theta) + \sum_{k=2}^{\infty} a_k r^{k\alpha} \sin(k\alpha\theta), \ \left(\alpha = \frac{\pi}{\omega}\right), \quad (1)$$

and the first term of this expansion is said singular because it makes the amplitude of the electric field go to the infinite at the sharp corner for $\alpha < 1$. The factor λ is called the *singularity factor* in the following. Methods for the computation of λ are given in Section III. A specific notation will also be used for the first function of expansion (1):

$$S(r,\theta) = r^\alpha \sin(\alpha\theta). \quad (2)$$

Finally, let v_∞ denote the solution of $\triangle v_\infty = 0$ in the unbounded domain Ω_∞ where the characteristic size of the rounded corner is 1, see Fig. 1(c). The rounded corner has the same shape as in the real domain Ω_ε, the potential is zero on Γ^0_∞, and the behavior of the potential at infinity is given by the first term of the expansion (1) up to the scalar λ:

$$\lim_{r_\infty \to \infty} (v_\infty(r_\infty,\theta) - S(r_\infty,\theta)) = 0. \quad (3)$$

The FEM enables to compute the term v_∞ whatever rounded corner shape (see Section IV). Note that for particular shapes analytic expressions of v_∞ can be obtained, for instance using conformal maps.

It is assumed in the remainder of this section that the approximate finite element solutions for v and v_∞ have already been computed.

B. Proposed approximate solution

Solution v_ε can be computed by the FEM, using a "fine" discretization close to the rounded corner. However, this computation has to be performed for each corner "scale" ε.

Our aim is to approximate the solution v_ε by using only:

- the solution v, that will approximate v_ε far from the corner, *i.e.* when $r \gg \varepsilon$;
- the solution v_∞, after a proper scaling by ε and the singularity coefficient λ, that will approximate the solution close to the corner, *i.e.* when $r \leq \varepsilon$.

[1]It is valid on any intersection of a disc centered in O with the domain Ω which coincides with the intersection of the same disc with an infinite sector whose corner is in O and whose angle is ω.

The "scaling" used to map a point of coordinate (r,θ) in Ω_ε to a point in Ω_∞ is given by $(r_\infty,\theta) = (r/\varepsilon,\theta)$. This change of coordinates applied in (2) leads to:

$$\varepsilon^\alpha S(r_\infty,\theta) = S(r,\theta). \quad (4)$$

These considerations enable to understand the following approximate expressions, that are detailed in [2]:

$$v_\varepsilon(r,\theta) \approx \lambda\varepsilon^\alpha v_\infty(r_\infty,\theta), \ \text{for } r < k\varepsilon \,(\text{usually}, k=1), \quad (5)$$

$$v_\varepsilon(r,\theta) \approx v(r,\theta) + \lambda\varepsilon^\alpha[v_\infty(r_\infty,\theta) - S(r_\infty,\theta)], \quad (6)$$
$$\text{for } r > k\varepsilon.$$

C. Remarks

The proposed expressions remain approximate, indeed:

- expressions (5) and (6) are not equal for $r = k\varepsilon$, in particular when the geometry of the domain Ω_ε makes the solution for $r < k\varepsilon$ to be not symmetric (in this first order approximation, the effect of the domain Ω_ε close to the corner is only related to the value of the singularity factor λ. It is not sufficient: every problem with the same angle ω and the same λ will give the same approximate solution close to the corner).
- in (6), $(v_\infty(r_\infty,\theta) - S(r_\infty,\theta))$ is not exactly zero on Γ^1_ε: thus, the conditions on this boundary, which are satisfied by v, are not satisfied by approximation (6). However, it can be seen, that whenever ε goes to zero, the remaining error becomes negligible: the scaling of the domain makes to send the boundaries further from the origin in the dimensionless solution v_∞ and the behavior at infinity (3) can be used.

It is possible to compute complementary terms, to increase the accuracy of the approximate solution. In particular, the following term in the expansion, which will vary as $\sin(2\alpha\theta)$ (it comes from (1)), will be able to introduce an asymmetry of the approximation close to the rounded corner. Nonetheless, our numerical experiments with the first terms of the expansion show the relevance of this first order approach and illustrate the theoretical estimates for the numerical error [2].

We emphasize that the method is generic and could be applied to other particular shapes which would not be "rounded" (and even to a simple surface irregularity, on a surface plane).

D. Application to the computation of the maximum field as a function of the curvature radius

With a reduced amount of numerical computations, a mesh without the rounded corner of the whole studied device and a mesh of the rounded corner but built independently from the rest of the structure, it is possible to determine the variation of the electric field close to the rounded corner as a function of the curvature radius. Indeed, the electric field $E_\varepsilon(r,\theta)$ in any point close to the rounded corner depends only on:

- the singularity factor λ whose evaluation is introduced in Section III starting from v the solution with the sharp corner;
- the field $E_\infty(r_\infty,\theta)$ transferred to the domain Ω_ε (the solution may be analytic in some cases, see section IV);

- the scale factor ε.

The corresponding relation is given close to the corner by:

$$E_\varepsilon(r,\theta) = \lambda \varepsilon^{(\alpha-1)} E_\infty(r_\infty,\theta). \tag{7}$$

For instance, if the statistical distribution of the curvature radii is known for a series of N pieces used in a high-voltage device, it will be possible to compute the statistical distribution corresponding to the maximal electric field in these pieces.

III. Computation of the singularity factor

A. The Fourier method: weighted line integral of v

The Fourier method enables to compute an approximation of the singularity factor λ from the solution v of the problem with the sharp corner (see Fig. 1(b)) given by the FEM. A weighted integral of v is computed on an arc of circle of radius r_0 centered on O, from one border of the angle to the other:

$$\lambda = \frac{2}{\omega} r_0^{-\alpha} \int_0^\omega v(r_0,\theta)\sin(\alpha\theta)\, d\theta. \tag{8}$$

When θ varies between 0 and ω, $\alpha\theta$ varies between 0 and π; in this range of variation, the integration of $\sin(\alpha\theta)$ with $\sin(k\alpha\theta)$ is zero for all k strictly greater than 1. Relation (8) is then obtained by considering expansion (1) in the line integral.

There is no restriction in the choice of r_0 except that the arc of circle must belong to the domain Ω.

B. Dual solution method

The theoretical explanations concerning this method can be found in [1, Subsubsection 2.3.8]. This approach enables to compute the singularity factor λ from an integral on the domain Ω with the sharp angle (see Fig. 1(b)) of a function v^* to compute, weighted by a function f which is the laplacian of a "regular lifting" function w:

$$\lambda = \int_\Omega f v^*\, ds. \tag{9}$$

It is proposed here not to construct this regular lifting but the less regular lifting which is commonly (but often implicitly) used in the FEM. It is shown that this FEM lifting can be used in the same manner but first, a few elements are given to understand where this expression comes from.

Our explanations take place on a domain Ω_δ, see Fig. 2. It can be seen that when the radius δ of the circle C_δ goes to zero, Ω_δ converges to Ω, see Fig. 1(b).

For two functions v and v^* of zero laplacian in Ω_δ, it can be written:

$$0 = \int_{\partial\Omega_\delta} \left(\frac{\partial v}{\partial \mathbf{n}} v^* - v \frac{\partial v^*}{\partial \mathbf{n}} \right) dl, \tag{10}$$

where $\mathbf{n}$ is the unitary outward normal on the boundary of the domain. As in [1], integral (10) can be made explicit on each boundary of the domain for a non zero δ and then take the limit when δ goes to zero, by considering the particular properties of the functions v and v^*.

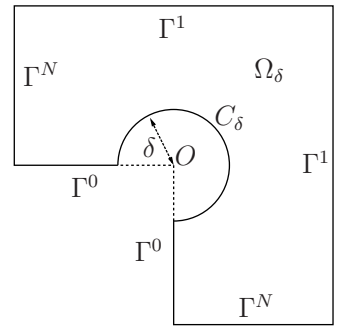


Fig. 2. Domain Ω_δ where $\delta \to 0$ is considered.

In this computation, v is the FEM solution of the problem on Ω detailed in Subsection II-A. The function v^* is the same as in [1]. It is the sum of the "dual singular function" S^*:

$$S^*(r,\theta) = r^{-\alpha}\sin(\alpha\theta), \tag{11}$$

and a more regular function v_R^* computed on the domain Ω which is the solution of the following problem:

$$\begin{cases} \triangle v_R^* = 0, \text{ on } \Omega, \\ v_R^* = -S^*, \text{ on } \Gamma^0 \cup \Gamma^1, \\ \dfrac{\partial v_R^*}{\partial \mathbf{n}} = -\dfrac{\partial S^*}{\partial \mathbf{n}}, \text{ on } \Gamma^N. \end{cases} \tag{12}$$

Problem (12) can be solved by the FEM. Thus it is obtained:

$$v^* = v_R^* + S^*, \tag{13}$$

which is a singular non zero function. Moreover v^* has the following properties:

$$\begin{cases} \triangle v^* = 0, \text{ on } \Omega, \\ v^* = 0, \text{ on } \Gamma^0 \cup \Gamma^1, \\ \dfrac{\partial v^*}{\partial \mathbf{n}} = 0, \text{ on } \Gamma^N. \end{cases} \tag{14}$$

Considering again the path integral (10) when δ goes to zero, the integrations are restricted to Γ^1 and C_δ:

$$\begin{aligned} 0 &= \lim_{\delta\to 0} \int_{\partial\Omega_\delta} \left(\frac{\partial v}{\partial \mathbf{n}} v^* - v \frac{\partial v^*}{\partial \mathbf{n}} \right) dl, \\ &= \lim_{\delta\to 0} \int_{\Gamma^1} -v \frac{\partial v^*}{\partial \mathbf{n}}\, dl + \lim_{\delta\to 0} \int_{C_\delta} \left(\frac{\partial v}{\partial \mathbf{n}} v^* - v \frac{\partial v^*}{\partial \mathbf{n}} \right) dl. \end{aligned} \tag{15}$$

It can be shown that the path integral on C_δ converges to $\pi\lambda$ when δ goes to zero (λ being the sought singular factor):

$$\lim_{\delta\to 0} \int_{C_\delta} \left(\frac{\partial v}{\partial \mathbf{n}} v^* - v \frac{\partial v^*}{\partial \mathbf{n}} \right) dl = \pi\lambda. \tag{16}$$

Moreover, the integral on Γ^1 in (15) can be written:

$$\int_{\Gamma^1} -v \frac{\partial v^*}{\partial \mathbf{n}}\, dl = -V^1 \int_{\Gamma^1} \frac{\partial v^*}{\partial \mathbf{n}}\, dl. \tag{17}$$

It is an integral only on the boundaries with non zero Dirichlet conditions (that are far from the corner), and it corresponds to the flux on each of these boundaries of the dual solution v^*.

This is a "classical" computation in the FEM, that can be done by a *surface* integral, on a transition layer $\mathrm{C}\Gamma^1$ (often limited to one layer of elements in direct contact with the boundary), of the scalar product of the gradients of the functions v^* and $V^1 w$, where w is equal to 1 on Γ^1 and goes to zero through the layer:

$$-V^1 \int_{\Gamma^1} \frac{\partial v^*}{\partial \mathbf{n}} dl \approx V^1 \int_{\mathrm{C}\Gamma^1} \nabla w \cdot \nabla v^* \, ds. \tag{18}$$

In fact, $V^1 w$ is the FEM lifting and it can be seen as an application of (9), by taking the laplacian of $V^1 w$ as a distribution. Eventually, the singularity factor is given by:

$$\lambda = -\frac{1}{\pi} V^1 \int_{\mathrm{C}\Gamma^1} \nabla w \cdot \nabla v^* \, ds. \tag{19}$$

C. Numerical comparison of both methods

The singularity factor λ is computed on a geometry similar to this on Fig. 1(b), where the length of each straight part of Γ^0 is equal to 50 mm, the length of each straight part of Γ^N is equal to 50 mm and finally the length of each straight part of Γ^1 is equal to 100 mm. The boundary conditions are also the same as these described in Subsection II-A. The potential V^1 is chosen equal to 1.

It is shown on Fig. 3 that for the method proposed in Subsection III-A, the computed result for λ is dependent on the choice of r_0. As r_0 increases, the computed value for λ quickly coincides for both methods.

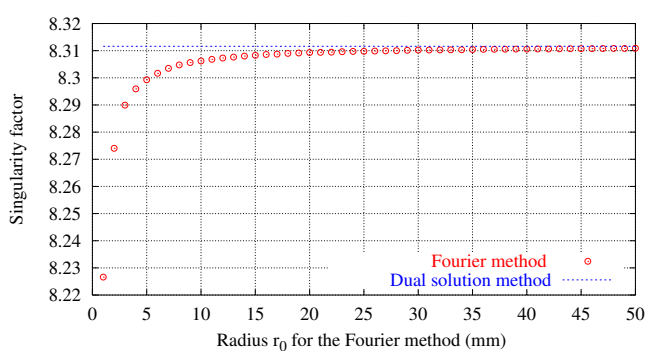


Fig. 3. Value of λ computed by the method described in Subsection III-A as a function of r_0. The dotted line indicates the value computed by the method described in Subsection III-B.

IV. Solution of the problem in Ω_∞

It was introduced in Subsection II-A a potential v_∞ solution of the Laplace equation on the unbounded domain Ω_∞.

In Subsection IV-A, the problem for v_∞ is formulated as a classical Laplace problem with non homogeneous boundary conditions, straightforwardly solved by the FEM. In Subsection IV-B, it is shown that conformal maps may give a solution for certain specific shapes of the rounded corner. In Subsection IV-C, the practical method used to obtain the approximation (6) by combining the solution v_∞ with the solution v on the domain with the sharp corner, is explained.

A. The equivalent problem with shifted singularity

Let us define a new function S_∞, with the correct behavior at infinity (3) and whose laplacian is zero in Ω_∞. Note that S itself is not suitable, because its laplacian is not defined in O (r equal to 0). We propose for instance:

$$S_\infty(r, \theta) = r_s^\alpha \sin(\alpha\theta), \tag{20}$$

where r_s is the distance between the observation point and any point inside the infinite conductor, typically the center of curvature of the rounded corner (Fig. 4).

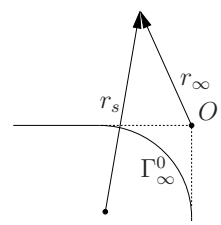


Fig. 4. Definition of r_s

The behavior at infinite is obviously correct, but this function is not zero on the surface of the conductor: v_∞ will be written as the sum of this function S_∞ and a corrective potential $v_{R\infty s}$:

$$v_\infty = S_\infty + v_{R\infty s} \tag{21}$$

which is the solution of:

$$\begin{cases} \triangle v_{R\infty s} = 0, \text{ on } \Omega_\infty, \\ v_{R\infty s} = -S_\infty, \text{ on } \Gamma^0_\infty, \\ \lim\limits_{r_\infty \to \infty} v_{R\infty s}(r_\infty, \theta) = 0. \end{cases} \tag{22}$$

This regular potential $v_{R\infty s}$ can be found with the FEM and a classical "shell transformation" for taking into account the unbounded domain [6]. Then, v_∞ is obtained by (21).

B. Method of conformal maps

The method of conformal maps [7] allows to build analytically potentials with the properties requested for v_∞, *i.e.* equal to zero on the conductor and with a prescribed behavior at infinity, on domains like Ω_∞. However, only some particular shapes are reachable by these transformations.

1) Example of transformation for a rounded corner: Let us consider the two complex planes:

$$z = x + \mathrm{i}y = r\mathrm{e}^{\mathrm{i}\theta} \text{ corresponding to the real geometry,} \tag{23}$$

$$\mathrm{w} = \mathrm{u} + \mathrm{i}\mathrm{v} = R\mathrm{e}^{\mathrm{i}\phi} \text{ where v is the electric potential.} \tag{24}$$

and the following conformal map [8, p. 320]:

$$z = k[(\mathrm{w} + a)^{1/\alpha} + (\mathrm{w} - a)^{1/\alpha}], \tag{25}$$

with $a = 2^{\alpha-1}$ and $k = 1/2$.

The image of the straight line $\mathrm{v} = 0$ in the w-plane is composed of (see Fig. 5):

- the half-line $y = 0$, $x > 1$, for $\mathrm{u} > a$ ($\mathrm{u} - a > 0$ and $\mathrm{u} + a > 0$);

- the half-line $re^{i\pi/\alpha}$, $r > 1$, for $\mathrm{u} < -a$ ($\mathrm{u} - a < 0$ and $\mathrm{u} + a < 0$);
- a rounded line from $P_l(z_l = e^{i\pi/\alpha})$ to $P_r(z_r = 1)$, for $\mathrm{u} \in [-a, a]$, ($\mathrm{u} - a < 0$ and $\mathrm{u} + a > 0$).

Moreover, for $|w| \gg |a|$, $z \approx \mathrm{w}^{1/\alpha}$ and $\mathrm{v} = \Im(\mathrm{w}) = R\sin(\phi) \approx r^\alpha \sin(\alpha\theta)$, what is the expected behavior at infinity. Then, this transformation gives the solution v_∞ for the particular shape given by (25).

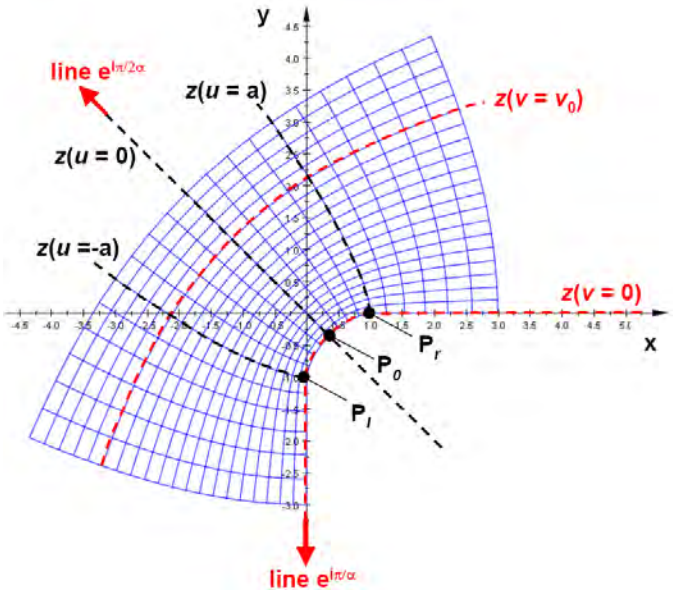


Fig. 5. $z(\mathrm{u} = \text{constant})$ and $z(\mathrm{v} = \text{constant})$ plotted in the z-plane for the conformal map (25)

2) Reference values of the electric field: The electric field on the conductor, for any point on the right of the rounded corner ($v = 0$, $\mathrm{u} > a$) is:

$$-\frac{d\mathrm{v}}{dy} = \frac{-2\alpha}{(\mathrm{u}+a)^{(1-\alpha)/\alpha} + (\mathrm{u}-a)^{(1-\alpha)/\alpha}}. \tag{26}$$

For example, the modulus values of E_∞ are:

$$|E_\infty(P_l)| = |E_\infty(P_r)| = \alpha 2^\alpha, \tag{27}$$

at the ends of the rounded corner, and:

$$|E_\infty(P_0)| = \frac{\alpha 2^{(1-\alpha)^2/\alpha}}{\sin\left(\frac{\pi}{2\alpha}\right)}, \tag{28}$$

on the middle-point of the rounded corner.

For the right angle ($\alpha = 2/3$), the field is constant on the rounded corner:

$$|E_\infty(\alpha = 2/3, \varepsilon = 1)| = 2^{5/3}/3 = 1,058267... \tag{29}$$

even though it varies for other values of the conducting corner.

The approximate field on the real structure at the same points is deduced from E_∞ using (7): for this computation, only the scale ε and the singularity factor λ are required.

We propose to use the reference field (28), or its particular value 1.06 for a right rounded angle, when the shape of the rounding is not exactly known, and when an approximation of the maximum field as function of the mean curvature of the corner is the requested information.

3) Remarks: Such conformal maps have the required behavior (at infinite and on the straight parts of the conductor) for favorably replacing the shifted singular function S_∞, (20) used in Subsection IV-A: the condition is that the line $v = 0$ remains inside the real conductor (at the reference scale $\varepsilon = 1$). By the way, the norm of the correction $v_{R\infty s}$ (22) will probably be reduced.

C. Scaling by ε: projection method

Whatever the method used to build v_∞ in the domain Ω_∞ (the FEM, conformal maps or another method), this potential is not directly constructed in the mesh where v is computed on the domain Ω with the sharp angle.

Moreover, the computation of the sum (5) and (6) for a particular ε, on a domain Ω_ε, see Fig. 1(a), requires a projection of v_∞ on the mesh of the final domain, with a scaling factor different for each value of ε. For this purpose, the projection is performed by a continuous least-square approach [9].

When only the results in some particular points is of interest, such as the amplitude of the electric field in the middle of the rounded corner, this step requires no particular computation.

D. Numerical results

It is considered the same problem as in Subsection III-C.

It is shown on Fig. 6, the good agreement between results obtained by formula (7) and results obtained by straightforward FEM computations. In order to plot these results, the solution v_∞ is computed by considering the method proposed in Subsection IV-A and λ is computed by the method proposed in Subsection III-B. To be precise, the values retained for $\max(|E_\infty|)$ and for λ are respectively $1.1467\,\mathrm{V/m}$ and 8.31235 (the value for $\max(|E_\infty|)$ can be compared to the value $1.06\,\mathrm{V/m}$ that is given by the conformal map described in Subsection IV-B).

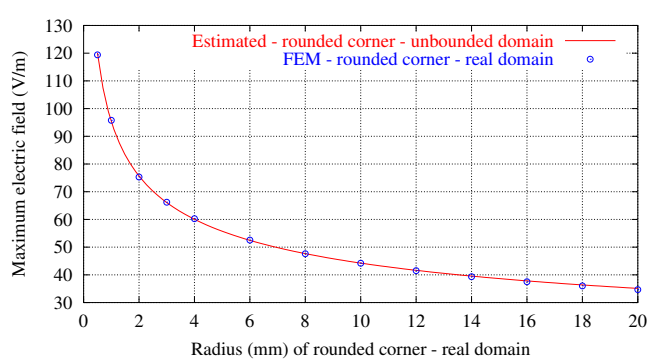


Fig. 6. Comparison of the maximum electric field on the rounded corner between results obtained by straightforward FEM computations and results obtained by relation (7).

It is also shown on Fig. 7 that the knowledge of λ and v_∞ provides an accurate approximation of the normal electric field on the conductor close to the corner (relation (5)) and that v and relation (6) provide accurate approximations of the normal electric field on the conductor far from the corner.

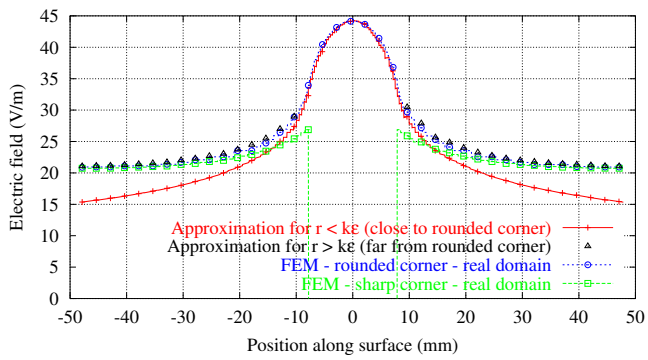


Fig. 7. Comparison of the normal electric field on the conductor for $\varepsilon = 10\,\mathrm{mm}$.

V. Conclusion and perspectives

In this work, a set of practical techniques was proposed to enable the implementation of new methods to estimate the electric field in the neighborhood of conductors containing rounded angles, in the 2D case.

The computation on the whole studied structure can be done by replacing rounded corners by sharp corners. A generic auto-similar solution enables to estimate the electric field close to the corner by post-processing for several values of the curvature radius. Some numerical tests have assessed the relevance of the proposed approach.

It has been also shown that forgotten techniques, as conformal maps, enable, in some cases, to build exactly the generic solution and in every case, to approximate it. They can be coupled with more computationally intensive numerical methods as the FEM, by providing speed and accuracy.

We are currently working to the extension of these techniques to:

- cases with several angles (such as some MEMs applications),
- 3D cases which is our real aim,
- singularity treatment (with sharp corners) for eddy-current problems, that lead also to self-similar solutions, with the penetration depth δ playing the same role as the scaling factor ε in the present study.

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