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**Semiempirical calculations on the electronic properties of
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RESEARCH ARTICLE

Semiempirical calculations on the electronic properties of finite-length models of carbon nanotubes based on Clar sextet theory

Materials Studio Special Issue

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Abstract

Electronic structure calculations on carbon nanotubes (CNTs) and related materials constitute an active and challenging field of research. Computational approaches to the problem require i) the definition of consistent models of CNTs ii) calculation of the properties on such models with accurate electronic structure methods. In this work, we perform semiempirical AM1 calculations on finite-length models of CNTs based on Clar sextet theory. In particular, the use of the Accelrys Materials Studio package allows to perform both the model building and computing steps through a simple and user-friendly interface. The consistency of such an approach are demonstrated by the smooth and monotonic decrease of the highest-occupied molecular orbital (HOMO) – lowest-unoccupied molecular orbital (LUMO) energy gap with length of models for metallic nanotubes and fast convergence to a finite value for models of semiconducting CNTs. As a further example of the applicability of the method outlined in this paper, the dependence of the HOMO-LUMO gap with the diameter for a series of analogous CNTs was also analysed.

Keywords: Materials Studio; carbon nanotubes; Clar sextet theory; electronic structure; semiempirical methods

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Introduction

The past few years have witnessed an impressive boost of research disciplines related to nanotechnology. In particular, the properties of carbon nanotubes (CNTs) have been deeply investigated [1] in view of their potential use in several applications of nanotechnology [2,3] or in nanostructured devices [4]. Consequently, CNTs and related materials have been the subject of comprehensive research work, in which their peculiar properties have been analysed in detail. Beside experimental studies, computer simulations have also been spotted as successful investigation techniques in the study of nanostructured materials. In particular, electronic structure methods, such as density functional theory (DFT) or semi-empirical techniques represent a preferential tool for numerical simulations, being able to provide accurate results for systems constituted of up to a few hundreds of atoms. However, similarly to other nanostructured systems, theoretical investigations on CNTs rely on an accurate definition of models to be employed in calculations. Such models should be ideally constituted by a small number of atoms and should be also representative of the system under study.

Computational investigations on single-walled CNTs are usually performed on either periodic or finite models of the hexagonal carbon atom network. On the one hand, periodic models are generally accurate and account properly for the properties of a virtually infinite quasi-unidimensional system. On the other hand, calculations on periodic systems are computational demanding, especially when the periodic unit cell is constituted by hundreds of atoms, as in the case of chiral CNTs. An alternative is based on the use of finite-length models, as used in previous work [5-7]. However, the representativity of finite-length models in describing the properties of nanosized objects is a non-trivial problem, and this also applies to the case of CNTs [8]. In particular, the construction of finite-length models of CNTs should take into account the peculiarities of the electronic structure at the sidewall, concerning especially π -conjugation and aromaticity.

Recent work, has indicated the use of Clar sextet theory [9,10] as appropriate for the description of the electronic structure of the sidewall of CNTs [11-13], in analogy to the case of polycyclic aromatic hydrocarbons [14]. This approach allowed the construction of finite length models of CNTs which are good representatives of their infinite counterparts. In particular, models defined in terms of finite-length Clar clusters (FLCCs) [15] provide accurate electronic properties, as demonstrated for calculations performed at the DFT level [15], and exhibit several computational advantages. Moreover, methods based on semiempirical potentials and particularly the AM1 method [16] allow to describe correctly the electronic structure and the reactivity of CNTs, as for example demonstrated by Orsmy and King in the study of hydrogenation reactions at the sidewall [12].

In this paper, we make use of the Materials Studio package to build models of zig-zag, armchair and chiral CNTs based on FLCCs and evaluate their structural and electronic properties at the semiempirical AM1 level through the VAMP [17] module of Materials Studio. Our results shows as the FLCC approach can be easily integrated in the Builder module of Materials Studio and demonstrate also the consistency of the electronic properties, calculated at the AM1 level, with the results of previous DFT calculations.

Computational details

Models of semiconducting zig-zag, armchair and chiral CNTs based on FLCCs were constructed through the Builder module of the Materials Studio package. In particular, as representatives of more general cases, models of (7,7), (8,0), (6,4) and (6,5) CNTs were taken into account. Moreover, the dependence of electronic properties with the nanotube diameter for an analogous set of systems was studied for semiconducting chiral ($n,n-2$) CNTs, with n ranging from 5 to 9. Finite-length models of nanotubes were built by starting from the corresponding periodic unit cell, as provided by the Materials Studio Builder, recutting and saturating the edges of the FLCCs with hydrogen atoms as described in Ref. [15].

Geometries were optimised at the AM1 level of theory as implemented in the VAMP module[17] of Materials Studio and electronic properties were evaluated at the same level.

Results and discussion

As discussed in previous work [11,15], Clar sextet theory can be conveniently applied to the hexagonal network of the CNT sidewall by defining two translational basis vectors, which reflect the translational symmetry of a fully benzenoid graphene sheet and are connected to the 2-dimensional (n,m) vector basis of the nanotubes sidewall [18] by algebraic relationships [11]. The Clar vector basis allows the definition of a primitive cell of the CNTs, from which the infinite CNT can be generated by application of a screw axis operation along the nanotube principal axis [15]. Ormsby and King demonstrated [11] that the description of the sidewall in terms of the primitive Clar cell leads to a classification of the CNTs based on the parameter R , which is defined as: $R=(n-m,3)$, where n and m are the nanotube chiral vectors. This classification accounts for the main features of the electronic structure at the sidewall and indicates whether the nanotube is metallic ($R=0$) or semiconducting ($R=1$ and $R=2$). The calculations presented in this work were performed on FLCC models of CNTs ranging from 2 to 6 replica of the Clar unit cell.

Figure 1 shows the optimised geometry of FLCC models of (7,7), (8,0), (6,4) and (6,5) CNTs for the case of three Clar cells. Geometries obtained at the AM1 level are in excellent agreement with the geometries evaluated at the DFT level [15], the larger deviations being related to the position of terminal hydrogens. The root mean square deviation (RMSD) between the DFT and AM1 optimised geometries was evaluated by first rotating the structures to minimise the RMSD, according to the method of Kabsch [19], then calculating the RMSD, defined as:

$$d_{RMS} = \left[\sum_{i=1}^N \|r_{DFT} - r_{AM1}\|^2 \right]^{1/2}$$

where r_{DFT} and r_{AM1} are the positions of atoms for structures computed at the DFT and AM1 levels, respectively, and N is the total number of atoms in the structure. As shown in Table 1, RMSDs between DFT (B3LYP/3-21G) and AM1 geometries are below 0.085 Å in all cases considered, thus pointing semiempirical methods as reliable methods for obtaining accurate structures of finite-length models of CNTs.

As indicator of the convergence of the electronic properties with the number of Clar cells, the highest-occupied molecular orbital (HOMO) - lowest unoccupied molecular orbital (LUMO) energy gap was evaluated for FLCC models of CNTs by varying the number of Clar cells replica. As stated previously, finite-length models should be representative of the corresponding infinite nanotubes and, consequently, exhibit a consistent and monotonic convergence of their properties towards the asymptotic limit for an infinite number of units. Figure 2 shows the dependence of the HOMO-LUMO energy gap from the number of Clar cells for FLCC models of the CNTs considered.

For the metallic (7,7) nanotube ($R=0$), the HOMO-LUMO gap exhibits a regular and monotonic decrease with the model length. This result is consistent with the properties expected for finite-length models of a metal, where the energy gap is zero at infinite length. Moreover, the monotonic trend indicates that FLCC models are well behaved and map properly, even at short lengths, the electronic properties of metallic nanotubes, in agreement with our previous DFT calculations [15]. Inversely, the HOMO-LUMO gap of semiconducting (8,0), (6,4) ($R=2$) and (6,5) ($R=1$) nanotubes converge quickly with the number of Clar cells. This indicates as the electronic properties of infinite nanotubes can be conveniently represented by using short-sized FLCC models. The HOMO-LUMO gap of models for (6,5) and (6,4) CNTs converges within 0.05 eV with 5 Clar cells whereas the model of the (8,0) CNT exhibit a slightly slower convergence (0.15 eV for 5 Clar cells).

Although the HOMO-LUMO gap was the only parameter considered in this study to assess the reliability of our computational approach, it must be noted that frontier orbitals are usually good indicator of many

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2 electronic properties of molecular and complex systems, including reactivity [20]. Therefore, calculations
3 performed at the AM1 level on CNT models based on chemical intuition, as in the case of FLCCs, are
4 expected to provide an accurate description of the electronic structure of nanotubes and, consequently, reliable
5 information on the sidewall reactivity [12]. As a further confirmation, we also checked the convergence of the
6 dipole moment of CNTs, calculated at the AM1 level, with the length of the corresponding FLCC model. As
7 shown in Figure 4 for the examples of the (8,0) and (6,5) CNTs, total dipole moments of FLCC models
8 converge promptly to a finite value, with an accuracy of about 0.02 Debyes for 5 Clar cells.

9 As stated above, semiempirical AM1 calculations were also performed on models of $(n,n-2)$ ($n=5,6,7,8,9$)
10 CNTs to assess the effect of the nanotube diameter on the electronic properties for a set of analogous systems.
11 In this case, the HOMO-LUMO energy gap decreases monotonically (see Figure 3), as expected for rolled
12 graphene sheets of decreasing diameters. This finding is in agreement with previous DFT calculations [15]
13 and confirms the possibility of using FLCC models in the study of electronic properties of CNTs at the
14 semiempirical level.

15 It is worth noting that approaches to the modelling of CNTs based on the conventional translational unit cell
16 do not guarantee an adequate description of the electronic properties of the sidewall, concerning especially
17 features related to π -conjugation and aromaticity, even if computed with a sophisticated theoretical method
18 [9]. Inversely, the use of models which fulfil the requirements of the electron conjugation pattern on the
19 sidewall even for short-sized clusters is “additive” by construction and allows to obtain consistent and
20 accurate results even with the application of relatively simple (and inexpensive) simulation techniques.

21 Conclusions

22 In summary, we demonstrated the use of the Materials Studio package as an integrated tool for the
23 construction of finite-length CNT models based on Clar sextet theory and for the evaluation of their properties
24 at the semiempirical AM1 level. Optimised geometries obtained at the AM1 level agree remarkably well with
25 the results of higher-level DFT (B3LYP/3-21G) calculations. Moreover, the electronic properties of CNT
26 models based on FLCCs calculated at the AM1 level are good representatives of their infinite counterparts. In
27 particular, the AM1 HOMO-LUMO energy gap decreases monotonically to zero gap or converges
28 monotonically to a finite value for models of metallic and semiconducting CNTs, respectively, in agreement
29 with previous studies [15]. Dipole moment calculations also confirm the reliability of the FLCC approach.

30 Furthermore, semiempirical calculations performed on FLCC models allow a correct description of the
31 change of the electronic properties with the diameter for a set of analogous CNTs.

32 Therefore, the application of user-friendly packages for the definition of computational models of CNTs based
33 on chemical considerations, such as the Material Studio builder, coupled to semiempirical methods, represents
34 a valuable method for obtaining quantitative information on the electronic properties of CNTs at a relative
35 cheap computational cost. The results and the methodology presented in this work can be extended to studies
36 where the electronic properties of CNTs play a crucial role, as in the case of investigations on the reactivity at
37 the sidewall.

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Figure Captions

- 34 Figure 1. Optimised geometries for CNT models constituted of 4 Clar cells: a) (7,7); b) (8,0); c) (6,4); d) (6,5).
35 Figure 2. HOMO-LUMO energy gap of CNT models with respect to the number of Clar cells: a) (7,7); b) (8,0); c) (6,4); d) (6,5).
36 Figure 3. Variation of the HOMO-LUMO energy gap with the diameter for models of (*n,n-2*) CNTs (*R=2*).
37 [Figure 4. Dipole moment of models of \(8,0\) and \(6,5\) CNTs with respect to the number of Clar cells.](#)

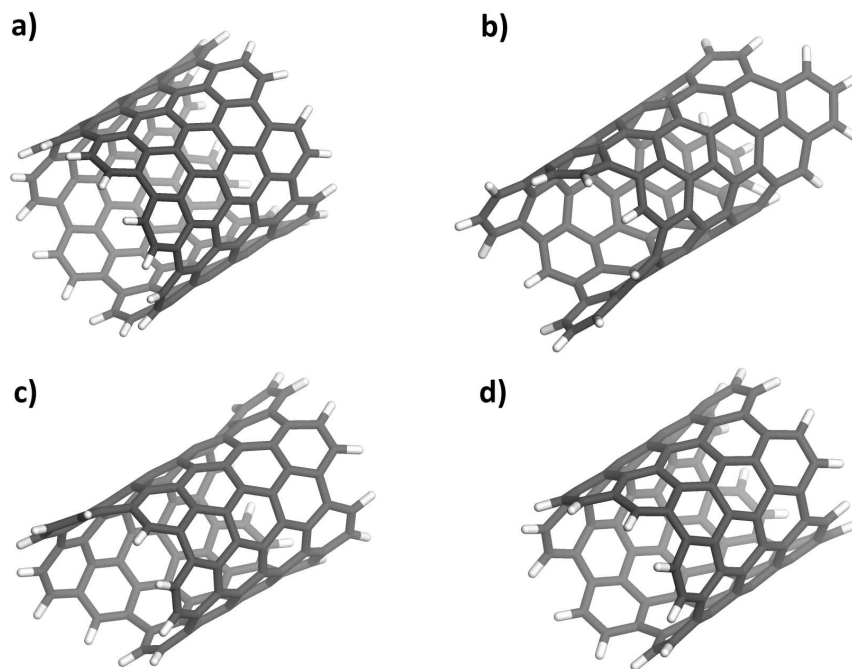
Table Captions

Table 1. RMSD, in Angstroms, between optimized geometries of CNT models calculated at the DFT and AM1 levels, respectively.

	2	3	4	5	6
(7,7)	0.048	0.039	0.036	0.034	0.034
(8,0)	0.084	0.066	0.058	0.053	0.049
(6,5)	0.064	0.050	0.046	0.049	0.045
(6,4)	0.049	0.039	0.053	0.045	0.041

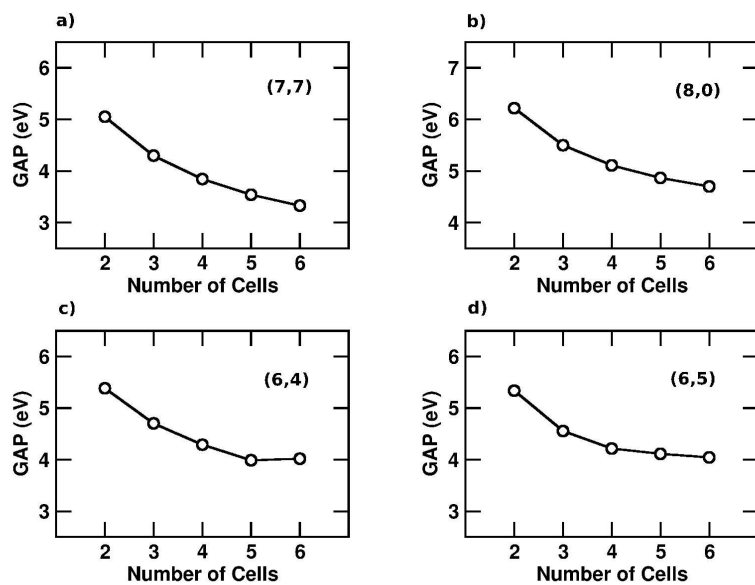
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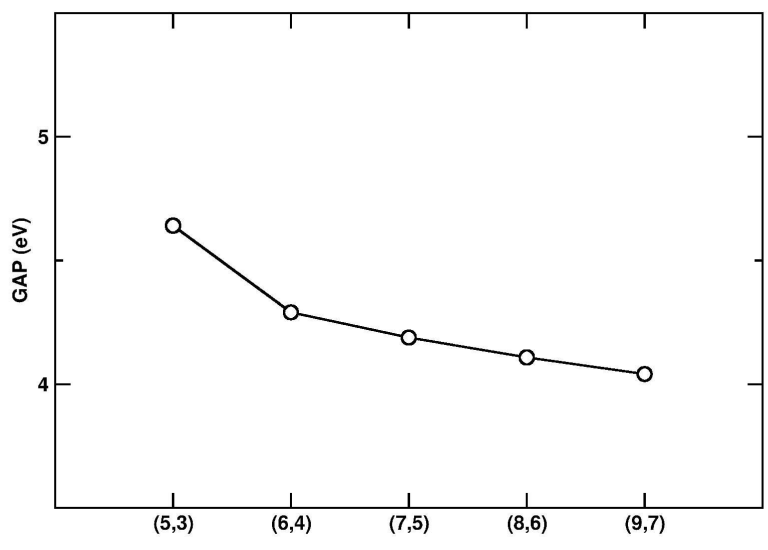
Optimised geometries for CNT models constituted of 4 Clar cells: a) (7,7); b) (8,0); c) (6,4); d) (6,5).
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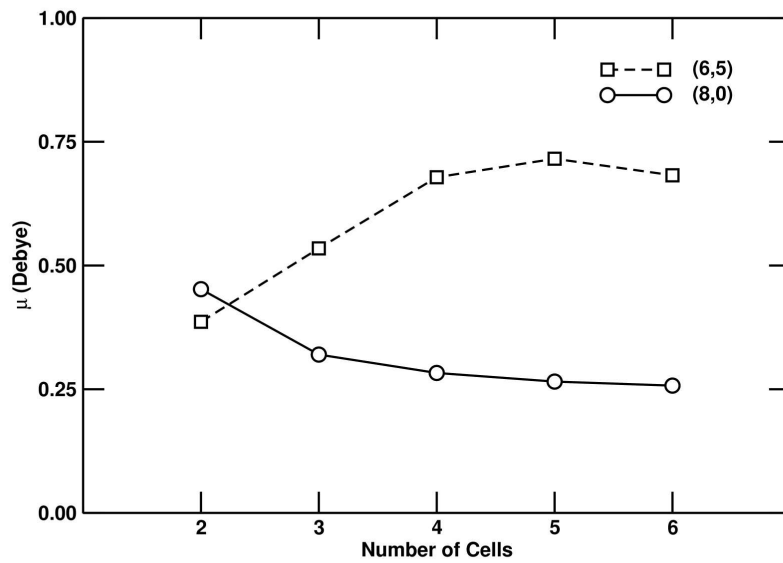
HOMO-LUMO energy gap of CNT models with respect to the number of Clar cells: a) (7,7); b) (8,0); c) (6,4); d) (6,5).
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Variation of the HOMO-LUMO energy gap with the diameter for models of (n,n-2) CNTs (R=2).
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