



HAL
open science

First-Principle Computations of Rotational-Vibrational Transition Probabilities

Roberto Linguerri

► **To cite this version:**

Roberto Linguerri. First-Principle Computations of Rotational-Vibrational Transition Probabilities. Molecular Physics, 2008, 106 (16-18), pp.2001-2009. 10.1080/00268970802054040 . hal-00513191

HAL Id: hal-00513191

<https://hal.science/hal-00513191>

Submitted on 1 Sep 2010

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



First-Principle Computations of Rotational-Vibrational Transition Probabilities

Journal:	<i>Molecular Physics</i>
Manuscript ID:	TMPH-2008-0081
Manuscript Type:	Full Paper
Date Submitted by the Author:	12-Mar-2008
Complete List of Authors:	Linguerri, Roberto; Universite Paris-Est, Equipe de Chimie Theorique
Keywords:	radiative transition probabilities, ions, spectroscopy, molecular physics
<p>Note: The following files were submitted by the author for peer review, but cannot be converted to PDF. You must view these files (e.g. movies) online.</p>	
<p>manuscript.tex</p>	



1
2
3
4
5
6
7
8
9
10
11
12
13
14
15
16
17
18
19
20
21
22
23
24
25
26
27
28
29
30
31
32
33
34
35
36
37
38
39
40
41
42
43
44
45
46
47
48
49
50
51
52
53
54
55
56
57
58
59
60

First-Principle Computations of Rotational-Vibrational Transition Probabilities

Pavel Rosmus* and Roberto Linguerri
*Laboratoire de Chimie Théorique,
Université de Marne la Vallée,
F 77454 - Champs sur Marne, France*

Najia Komiha
*Laboratoire de Chimie Théorique,
Université Mohamed V-Agdal, Rabat,
Morocco*

(Dated: March 12, 2008)

The theoretical calculations of dipole moments provide reliable compact information required for the evaluation of radiative transition probabilities in the electronic states of small molecular ions. The comparison of the computed line intensities and dipole derivatives with available experimental data in neutral molecules demonstrates that they can be calculated *ab initio* with a precision of a few percent, hence with an accuracy sufficient for most of the experimental applications.

PACS numbers:

I. INTRODUCTION

The high quality first-principle quantum mechanical calculations of the absolute radiative transition probabilities in the electronic states of molecular ions yield more accurate results than experiments. Similar to neutral molecules such data are needed in many fields of research, for instance, for assignments of high resolution spectra, spectral line broadening, measurements of ion abundances, radiative decay rates of excited rotational-vibrational states, in atmospheric chemistry, astrophysics, or in studies of ion-molecule reactions.

In order to show how accurately the radiative transition probabilities can be calculated by a purely theoretical approach for small molecular ions, we first discuss the results for neutral triatomics O₃[1], H₂S[2], and CO₂[3] because of the numerous precise spectral data[4–6] available for comparison. In the first-principle approach the potential energy and electric dipole moments surfaces are computed from highly correlated electronic wavefunctions[7]. This compact information is then used to solve the nuclear motion problem variationally, using exact kinetic energy operator in internal coordinates[8–10], yielding rotational-vibrational energies and transition dipole matrix elements. Such a variational approach is essential for accurate results.

II. OZONE

In the analysis of the infrared spectra of the electronic ground state of ozone[11] usually Hamiltonian matrices are set up for polyads of Darling-Dennison and Coriolis

resonances, and diagonalized. Using the mixing coefficients, the necessary spectroscopic data are evaluated for the line assignments and the transition moment parameters. From the least-squares fitting of the transition moment parameters to the measured subset of line intensities one can calculate all line intensities for a given vibrational band. The experimental data were used to set up an effective electric dipole moment function[12] expanded as a linear combination in the normal modes and the corresponding dipole derivatives. At most, one derivative is evaluated from the transition moment parameters of one vibrational band. For instance, the second derivatives were obtained from the first overtones or corresponding combination levels.

We have calculated more than 4.8×10^6 electric dipole transition matrix elements for the absolute rotational-vibrational line intensities of all bound levels up to about 8500 cm^{-1} and $J=64$. All anharmonic and Coriolis resonances were implicitly taken into account in the variational calculations with the nuclear motion Hamiltonian. In Figure 1 the calculated pattern of rotational levels up to $J=64$ for the vibrational ground state is shown. Since the value of dissociation energy of ozone is about 8500 cm^{-1} , all rotational-vibrational states with lower J will play a role in the energy redistribution during the formation of ozone in vibrational ground state by three-body collisions of the O₂-O complex with other atoms and molecules. In the calculations of the line intensities at 273 K, we used the value of the partition function from the HITRAN database. In all calculations of absorption spectra a Boltzmann distribution at 273 K was assumed, although the dipole matrix elements could serve in calculations of emission or absorption spectra for any population of the initial level.

In order to compare the calculated dipole functions with previous theoretical and empirical functions, we have transformed our dipole function to the Eckart frame and used our harmonic wavefunctions to calculate the deriva-

*Electronic address: rosmus@univ-mlv.fr

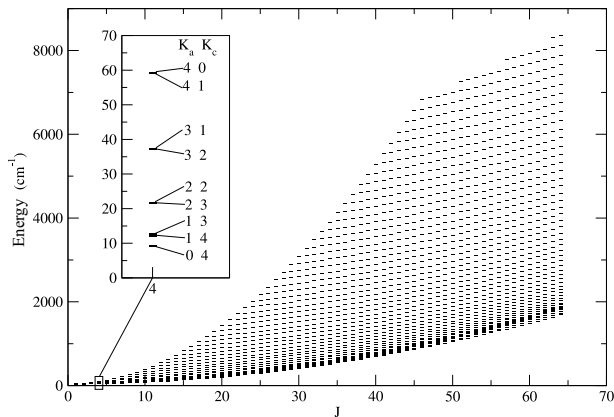


FIG. 1: Pattern of variationally calculated rotational $K_a K_c$ levels up to $J=64$ in the vibrational ground state of $^{16}\text{O}_3$. Starting from $J=46$ highest levels are not plotted

TABLE I: First and second derivatives with respect to dimensionless normal coordinates of the electric dipole moment function of ozone in Eckart frame (in 10^{-2} debye)[1].

	MRCI	MRCI+Q ^a	MRCI+Q [13]	Exp. [12]
μ_1^x	-1.839	-2.097	-1.897	-2.185
μ_2^x	-7.178	-6.812	-7.456	-6.529
μ_{11}^x	-0.409	-0.344	-0.476	-0.366
μ_{12}^x	0.828	0.786	0.777	0.894
μ_{22}^x	-0.145	-0.154	-0.097	-0.122
μ_{33}^x	0.756	1.149	0.662	0.769
μ_3^z	-26.92	-25.40	-28.85	-26.62
μ_{13}^z	4.98	4.85	4.50	4.81
μ_{23}^z	-0.258	-0.559	0.381	0.154

^aThe data in the third column have been recalculated from the dipole moment function of [13].

tives with respect to dimensionless normal coordinates. These values are given in Table I. The derivatives from experiment are in good overall agreement with our theoretical values as far as the signs and absolute values are concerned.

For atmospheric applications it is necessary to know the intensities with an accuracy better than 1 to 3 percent. In Figures 2 and 3 the first-principle rotationally resolved absorption spectra of ozone are compared with the spectra from the HITRAN database. Both sets of spectral intensities are found to be in excellent agreement. For purely rotational transitions the differences are about 0.3 percent. For the transitions to (010), (001), (110), (011) the agreement is within 1 to 3 percent. Surprisingly, also, the extremely weak transitions to the (020) state (the line intensities are ten thousands times weaker than those of the fundamental antisymmetric stretching transition)

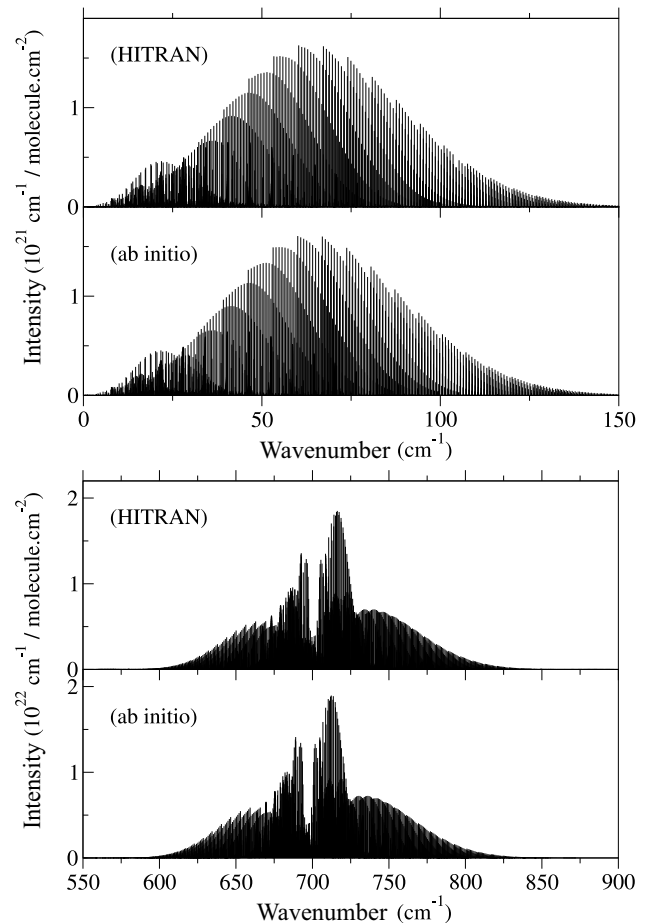


FIG. 2: Purely rotational absorption spectrum at 273 K (upper part) and bending fundamental (lower part).

agree very well with the database values. Very pleasing is the agreement between theory and experiment for the (001) band also on the large scale for transitions lying within about 10 cm^{-1} in Figure 4. The ozone spectrum exhibits congested regions where several lines lie within 1 cm^{-1} of each other. The first-principle variational calculations for a molecule with three heavy atoms are unable to yield the proper ordering of such states, although the line intensities are still correctly calculated and useful for assignments.

III. HYDROGEN SULFIDE

The experimental infrared spectra of H_2S molecule present interesting anomalies in the rovibrational intensity distributions[2]. They concern an unusual behaviour of band intensities versus a number of vibrational quanta $\Delta v = \Delta v_1 + \Delta v_2 + \Delta v_3$: observed intensities of H_2S fundamental bands ($\Delta v = 1$) are much weaker than those of combination bands with $\Delta v = 2$. For example, P_1 transitions of ν_1 and ν_3 fundamental bands are by one or two orders of magnitude weaker than those of $\nu_1 + \nu_2$, $\nu_2 + \nu_3$,

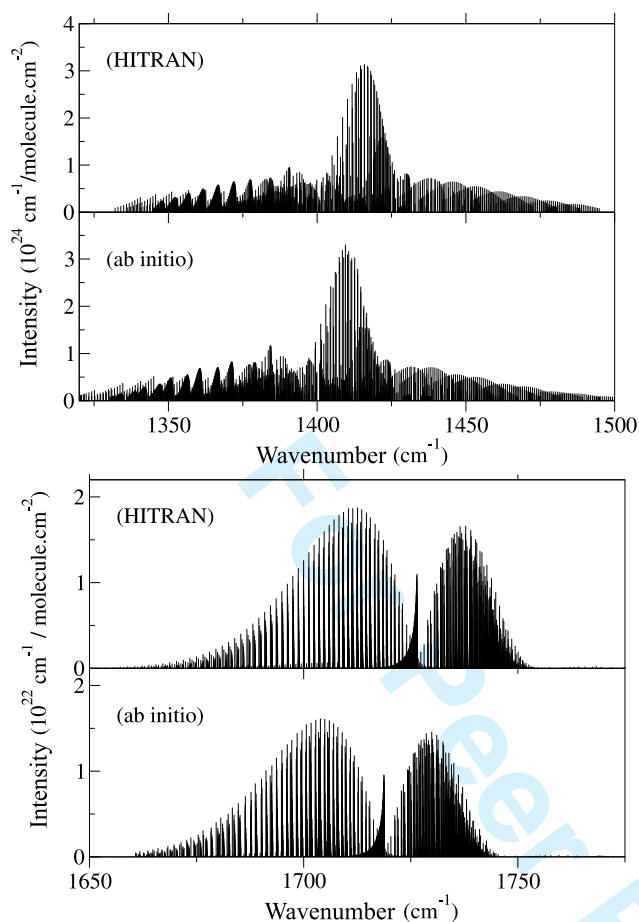


FIG. 3: Absorption spectra at 273 K for the combination (020) band (upper part) and (011) band (lower part).

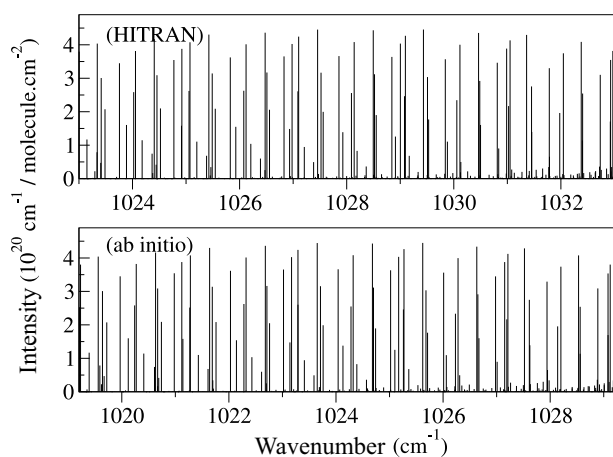


FIG. 4: Part of the absorption spectrum for the antisymmetric stretching fundamental. Upper spectrum has been taken from the HITRAN database and lower is from the first-principles computations.

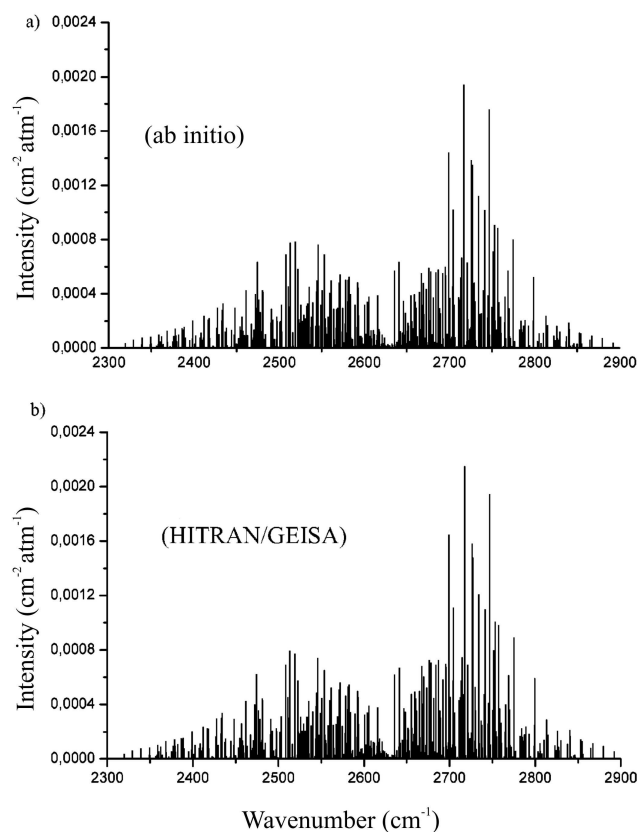


FIG. 5: Absorption spectrum at 273 K for the ν_3 band; $J_{upper} \leq 10, J_{lower} \leq 11$.

or $\nu_1 + \nu_3$ combination bands unlike the majority of bent triatomic molecules. The fundamental band $\Delta v_3 = 1$ is particularly weak, The experimentally measured fundamental band $\Delta v_3 = 1$ is weaker than the corresponding transitions of $\Delta v = 2$ bands, many $\Delta v = 3$ bands and even of some $\Delta v = 4$ bands. Fundamental bands exhibit also rotational anomalies in intensity distributions. The P-branches of the ν_1 and ν_2 bands are much less intense than the R-branches and in the ν_3 band some $\Delta K_a = 2$ transitions are more intense than corresponding $\Delta K_a = 0$ transitions. The components of the dipole moment function are very flat close to the equilibrium geometry, therefore contributions to transition dipole matrix elements of fundamental bands ν_1, ν_3 from second and higher-order derivatives with respect to normal or displacement coordinates become comparable and even more important than the first derivatives. All anomalies are due to the unusual form of this function. The coupled cluster method with perturbative triples has been employed with a large basis set augmented by diffuse functions and the components of the dipole moment have been evaluated as the energy derivatives with respect to an applied electric field. In Figure 5 we compare the rotationally resolved absorption spectrum of the very weak ν_3 band calculated *ab initio* with the spectrum available in the data bases. The overall agreement is considered as very

good. Already an error of about 10^{-2} debye/Å produces a very large deviations in the calculated line intensities. Other examples are given in Ref.[2]. The dipole moment function for the electronic ground state of hydrogen sulfide is computationally extremely challenging task due to the accuracy required for the evaluations of radiative transition probabilities particularly for higher frequency regions.

IV. DIPOLE MOMENTS OF SMALL IONS

The electric dipole moments in the vibrational ground states are of importance, for instance, in microwave spectroscopy or radioastronomy. They can be determined experimentally only in indirect way. The CO^+ dipole moment[14] was deduced from polarizabilities and precise mass comparison of CO^+ to N_2^+ . The cyclotron frequency shifts between $N = 0$, and $N = 1$, $M_N = 0, \pm 1$ states yield a value of $\mu_0 = 1.025(15) ea_0$, which compares well with our recent MRCI theoretical value of $1.035(8) ea_0$ [15]. Prior to the measurements of polarizability shifts of CO^+ , experimental determination of dipole moments in diatomic molecular ions have been based on isotopic substitution in the rotational Zeeman effect, or else extensive analysis of pure-rotation na rotation-vibration spectra of isotopic variants. The example of ArH^+ , which received most attention, shows the difficulties with previous experimental methods. The experimentally determined dipole moments of $1.4(0.4)$ debye[16] and $3.0(0.6)$ debye[17] may be compared with our recent MRCI result of 2.177 debye and the experimental determination from detailed spectral analysis $2.12(0.55)$ debye[18] ($1 ea_0 = 2.54158$ debye).

The most complete and accurate theoretical study of emission rates for all rotation-vibration transitions has been reported for HeH^+ [19]. This ion is thought to be the first molecular species to appear in the Universe. Consequently, stars formed from primordial material should contain some HeH^+ . These theoretical predictions should help to determine the best conditions to search for and detect the ion. It should also facilitate its detection in other astrophysical environments.

Recently theoretical electric dipole moment functions of the first - row diatomic hydride ions have been reported. The dipole moments in the vibrational ground states are given in Table II. The LiH^+ ion is thought to be constituent of the primordial matter like HeH^+ . The calculated dipole moment function is, however, very flat near the equilibrium distance and the dipole moment μ_0 is not very large. The protonating mode in HeH^+ has much stronger vibrational transitions and it seems to be hardly probable to detect LiH^+ in infrared spectral region of the electromagnetic radiation from outer space. The parts of the dipole moment functions near the equilibrium geometries of the ions in Figure 6 reflect also the first ionization energies of the heavy atoms relative to hydrogen. The dipole moments of ions are

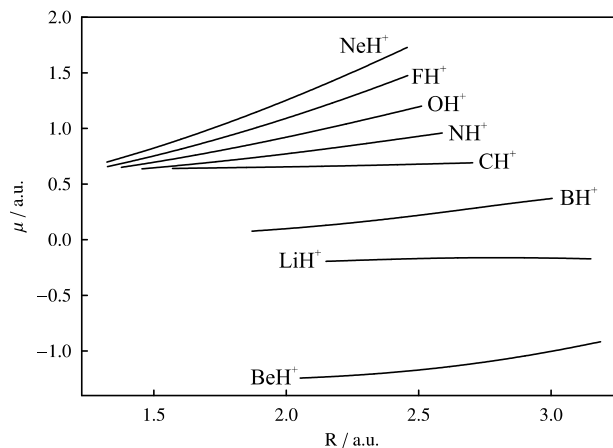


FIG. 6: Dipole moment as a function of internuclear separation: LiH^+ to NeH^+ .

calculated for each isotope variant with origin at the center of mass. Asymptotically they are zero for $r = 0$ Å and go to infinity for large internuclear distances. They can, however, possess a maxima or minima which will influence the infrared spectra in high frequency regions.

The dipole moments of the hydride ions were needed in calculations of orientation polarizabilities for very precise atomic mass measurements[15]. The bottleneck of this approach is the accuracy of the theoretical dipole moments which are expected to be $\pm 0.008 ea_0$

Several dipole moment functions of molecular ions have been used to calculate radiative relaxation of vibrational states. Perhaps the best known example is NO^+ for which the decay rates for five vibrational levels have been measured by monitoring the change in vibrational populations through internal energy dependent fast reactions. The latest experimental[20] and theoretical[21] radiative lifetimes for $v = 1$ to 5 lying in the range of 90 to 20 ms were found to be in good overall agreement. To our knowledge no other experimental information about the shapes of larger parts of the dipole moment functions in molecular ions have been reported so far.

Though recent experiments[22] proved the existence of six equivalent C_{2v} minima in the pseudorotating CH_4^+ ion no transition probabilities are known for this important species. In a series of RCCSD(T) computations with cc- V_xZ AO basis sets we have reoptimised the equilibrium geometry of the lowest C_{2v} minimum of CH_4^+ . Our best r_e geometry is for the two CH_2 angles 62.76° and 27.55° and the two CH distances 1.0796 Å and 1.1825 Å (cf. ref. for the definition of the geometry parameters). This CH_4^+ structure has an unexpectedly large dipole moment of $-0.6600 ea_0$, the CD_2H_2^+ $-0.7788 ea_0$. Regardless of complicated intramolecular dynamics playing a role in the radiative transition probabilities the large dipole advocates the possibility to detect the ion in the outer space.

TABLE II: Calculated MRCI electric dipole moments μ_0 , and internuclear separations R_0 , for the lowest vibrational state of CO^+ and the hydride cations LiH^+ to NeH^+ . A positive sign for the dipole implies it is directed towards the lighter nucleus. μ_0 and R_0 are given in atomic units (ea_0, a_0), B_0 is given in cm^{-1} (see Ref. [15] for details).

Diatomic ion	State	μ_0	R_0	R_0 Exp.	B_0 Exp.
$^{12}\text{C}^{16}\text{O}^+$	$^2\Sigma_{1/2}^+$	1.0352	2.1160	2.1125	1.9675
$^7\text{LiH}^+$	$^2\Sigma_{1/2}^+$	-0.3049	4.2729		
$^9\text{BeH}^+$	$^1\Sigma_0^+$	-1.1646	2.4838	2.4944	10.674
$^{11}\text{BH}^+$	$^2\Sigma_{1/2}^+$	0.1673	2.2814	2.2957	12.371
$^{12}\text{CH}^+$	$^1\Sigma_0^+$	0.6623	2.1530	2.1559	13.930
$^{14}\text{NH}^+$	$^2\Pi_{1/2}$	0.7897	2.0406	2.0437	15.331
$^{16}\text{OH}^+$	$^3\Sigma_{0,1}$	0.9134	1.9641	1.9663	16.422
$^{19}\text{FH}^+$	$^2\Pi_{3/2}$	1.0438	1.9157	1.9168	17.119
$^{20}\text{NeH}^+$	$^1\Sigma_0^+$	1.1823	1.9022	1.9023	17.338

V. INFRARED EMISSION IN THE $X^2\Pi_g$ STATE CO_2^+

Spacecraft observations of CO_2^+ in Mars' atmosphere have been made for more than 35 years (see Ref. [3] for a review). The densities of the principal ions (O_2^+ , CO_2^+ , O^+) in Mars' atmosphere were measured by the Retarding Potential Analyzers on the Viking 1 and 2 landers on their descent. The production rate in planetary atmospheres of the CO_2^+ ion has been determined by using ultraviolet emission bands.

The Thermal Emission Spectrometer on Mars Global Surveyor has observed numerous limb spectra of Mars' atmosphere from near the surface to above 125 km. The resulting averaged IR emission spectrum covers an energy range from $150\text{-}1650\text{ cm}^{-1}$ and shows the strong $X^1\Sigma_g^+$ CO_2 (001) \rightarrow (100)-(020) Fermi resonance pair near 960 and 1063 cm^{-1} and the $X^1\Sigma_g^+$ CO_2 fundamental bending transition at 667 cm^{-1} . The IR emission spectrum has also features in the energy range $400\text{-}500\text{ cm}^{-1}$.

Using *ab initio* calculated potential energy and electric dipole and transition moment functions, absolute radiative transition probabilities between the rovibronic Renner-Teller states of the electronic ground state $X^2\Pi_g$ of CO_2^+ have been evaluated. The rovibronic eigenstates were obtained up to 6500 cm^{-1} and $J \leq 13/2$ from variational calculations including the Renner-Teller and spin-orbit angular momentum couplings. The absolute intensities for the rotational-vibrational transitions in the neutral CO_2 electronic ground state [3], calculated by a similar approach, are found to be in excellent agreement with experimental data. For CO_2^+ the presented intensities are expected to have similar accuracy. The calculations show that overtone intensities of the antisymmetric stretching mode of CO_2^+ are of similar magnitude as the fundamental transitions, in contrast with CO_2 .

The derivatives of the dipole moment functions with respect to the dimensionless normal coordinates at the equilibrium geometry of the electronic ground states

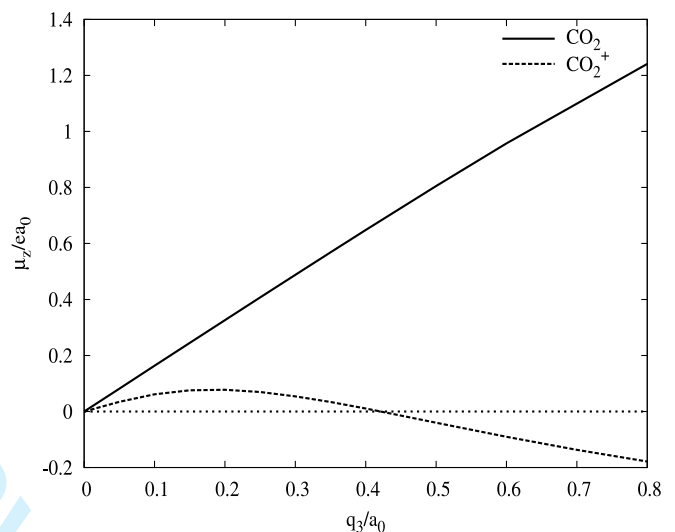


FIG. 7: μ_z along the antisymmetric coordinate q_3 using MCSCF/cc-pVQZ method. $q_3 = (r_2 - r_1)/\sqrt{2}$ and $q_1 = (r_2 + r_1)/\sqrt{2}$ is fixed at the equilibrium value.

of CO_2 and CO_2^+ are given in Tables III and IV, respectively. The *ab initio* values for CO_2 are compared with previous data extracted from experimental measurements using direct numerical diagonalization (DND) (cf. Ref. [3]). Both sets are in excellent agreement. For CO_2 , the second and third derivatives are always smaller than the first ones. The μ_3^z derivative is more than twice as large as μ_2^z reflecting larger radiative probabilities for transitions involving the antisymmetric stretching mode compared to the bending mode.

In Table IV the same derivatives are given for both electronic components of CO_2^+ . The two sets of derivatives are similar. The μ_y derivatives are always much smaller than other derivatives. The μ_x derivatives of CO_2^+ are similar to the μ_x derivatives of CO_2 . The comparison of the derivatives for μ_z between CO_2 and CO_2^+ reveals unexpectedly large third derivatives with respect to the dimensionless antisymmetric stretching

TABLE III: Derivatives of the dipole moment functions (10^{-2} debye) for the $X^1\Sigma_g^+$ of CO_2 with respect to the dimensionless normal coordinates at the equilibrium geometry ($R_{CO} = 2.1970 a_0$)[3].

	μ_x			μ_z	
	CCSD(T)	Exp.(DND)		CCSD(T)	Exp.(DND)
μ_2	18.48	18.4	μ_3	-45.65	-46.1
μ_{12}	-0.78	-0.783	μ_{13}	1.86	1.82
μ_{112}	0.004	0.00591	μ_{113}	0.03	0.00518
μ_{233}	0.16		μ_{223}	0.73	0.668
μ_{222}	-0.93	-0.652	μ_{333}	0.17	0.151

TABLE IV: MRCI derivatives of the dipole moment functions (10^{-2} debye) for the $X^2\Pi_g$ of CO_2^+ with respect to the dimensionless normal coordinates at the equilibrium geometry ($R_{CO} = 2.2207 a_0$)[3].

	μ_x			μ_z	
	A'	A''		A'	A''
μ_2	25.25	21.71	μ_3	-24.40	-24.40
μ_{12}	-1.20	-0.86	μ_{13}	-1.35	-1.35
μ_{112}	-0.02	-0.24	μ_{113}	-0.12	-0.12
μ_{233}	1.48	1.62	μ_{223}	0.91	0.77
μ_{222}	-1.48	-1.06	μ_{333}	28.42	28.42

	μ_y	
	A'^a	A''^b
μ_2	-0.64	-0.58
μ_{12}	0.16	0.14
μ_{222}	0.07	0.05
μ_{112}	-0.04	-0.04
μ_{233}	-0.32	-0.29

^aderivatives with respect to the normal coordinates of A' component.

^bderivatives with respect to the normal coordinates of A'' component.

coordinate in CO_2^+ . In Figure 7, μ_z is plotted against $q_3 = (r_2 - r_1)/\sqrt{2}$ for CO_2 and CO_2^+ . The curves show that μ_x is nearly linear with a large slope for CO_2 whereas μ_z has a shallow maximum close to the linear structure. Similar to H_2S rovibronic intensity anomalies are to be expected for the transition probabilities including the antisymmetric stretching mode.

As an illustration of the results, a compact representation in form of the radiative lifetimes of the first rotational-vibrational and rovibronic states of CO_2 and CO_2^+ for $J = |P|$ are shown in Figure 8. The v_2 transitions are connected. The lifetimes decrease with the values of v_2 and v_1 (if $v_1 > 0$) in the neutral and the ionic species. In CO_2 the lifetimes for levels involving the antisymmetric stretching mode are smaller than 3 ms whereas they are about 70 ms in CO_2^+ . In Figure 8 (lower part), the outliers are due to Fermi coupling. Particularly, the radiative lifetime of the level $\Pi_{g,3/2}(1,0,0)$ is lower than the one of the level $\Pi_{g,1/2}(1,0,0)$, although its energetic position is lower.

The Mars spectra from the middle atmosphere through about 127 km show a strong broad structured peak extending from $\sim 420 - \sim 520 \text{ cm}^{-1}$. The feature is optically thin so limb tangent observations include emission from

many altitudes, including the CO_2^+ peak at $\sim 130 \text{ km}$.

For the comparison of the calculations to these observations, the spontaneous emission rates of CO_2^+ for equal initial population of the excited levels and for the values of $J 1/2 \leq J \leq 5/2$, with the maximum energy at 2000 cm^{-1} were used together with the energy positions to create a synthetic theoretical spectrum (cf. Ref.[3] for details). The envelope considering all dipole allowed transitions adapted to the resolution of the experiment is plotted in figure 9.

Both the spontaneous rates and the synthetic spectrum reproduce features between 465 and 509 cm^{-1} . The transitions are between levels involving the bending quantum number v_2 , with $\Delta v_2 = \pm 1$ and the complicated structures are the result of Fermi, spin-orbit and Renner-Teller couplings. CO_2^+ has also significant emission in the 1000 cm^{-1} region where neutral CO_2 hot bands are present. Other parts of the broad feature are less well reproduced, indicating nonthermal population of the CO_2^+ energy levels and/or additional emission sources. In comparing the peak positions of the synthetic and observed features, exact coincidence is not expected due to the spectral averaging introduced by the low-resolution sampling designed into the TES instrument. The peaks around $400\text{-}500 \text{ cm}^{-1}$ were

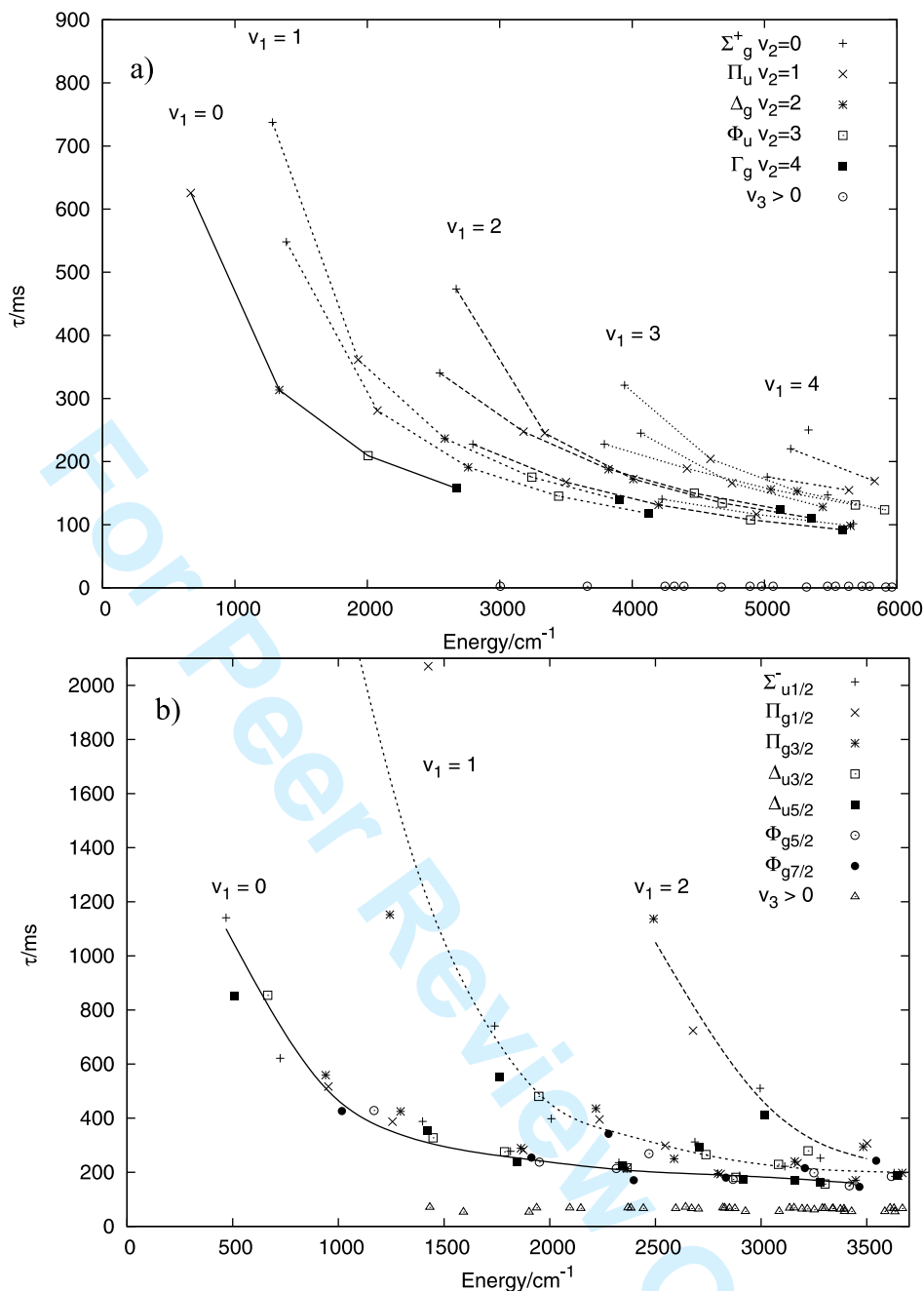


FIG. 8: a) CO_2^+ rovibrational radiative lifetimes for levels with $J = |P| = l$ and $0 \leq v_2 \leq 4$. The points associated with the same value of v_1 and the same position in the polyads are linked. The radiative lifetime τ_n of the level n is given by $\tau_n = (\sum_{n'} A_{nn'})^{-1}$ where $A_{nn'}$ is the Einstein coefficient for spontaneous emission and $E_{n'} < E_n$. b) CO_2^+ rovibrational radiative lifetimes for levels with $J = |P|$ and $0 \leq v_2 \leq 6$. Averaged curves, passing through the points associated with the same value of v_1 are shown.

assigned to the emission of the CO_2^+ bending levels. The remaining unassigned peaks are expected to be due to emission of bending levels of other neutral or ionic polyatomic species.

VI. CONCLUSIONS

The theoretical radiative transition probabilities within the electronic states of molecular ions can be calculated with sufficient precision for most applications in molecular physics. Progress may be expected in the development of new experimental techniques for more reliable electric dipole moment functions of molecular ions.

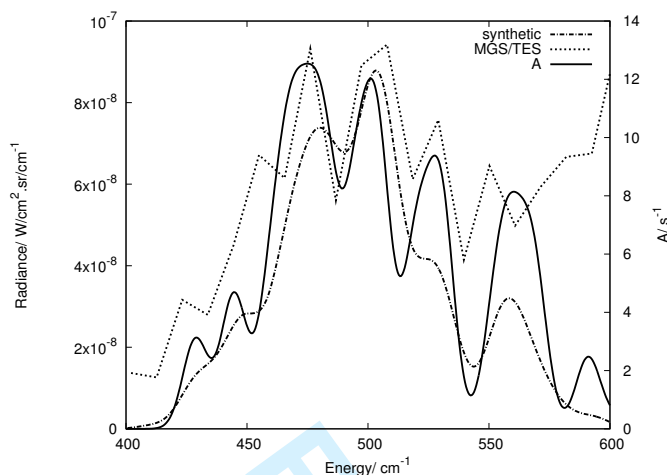
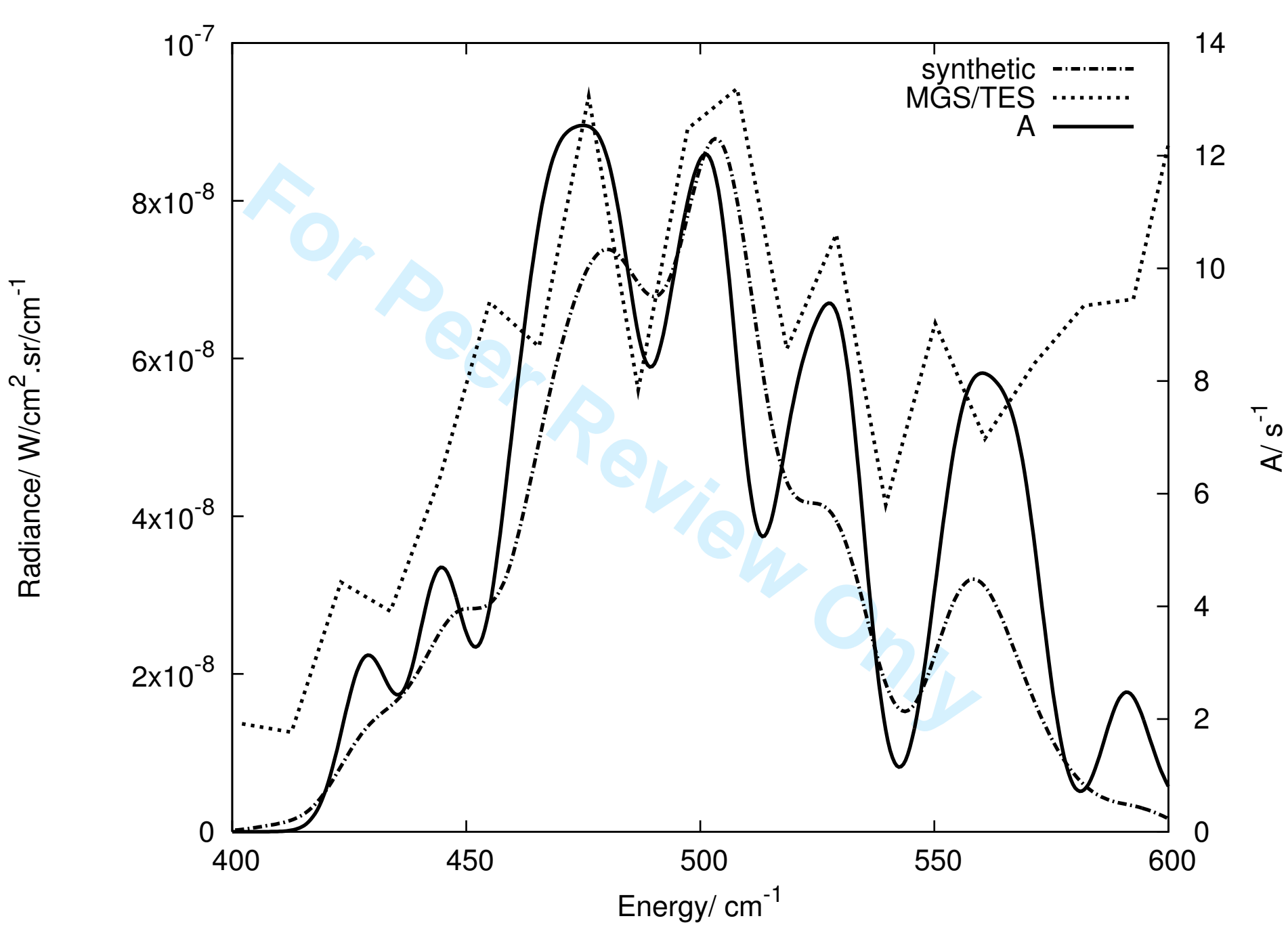
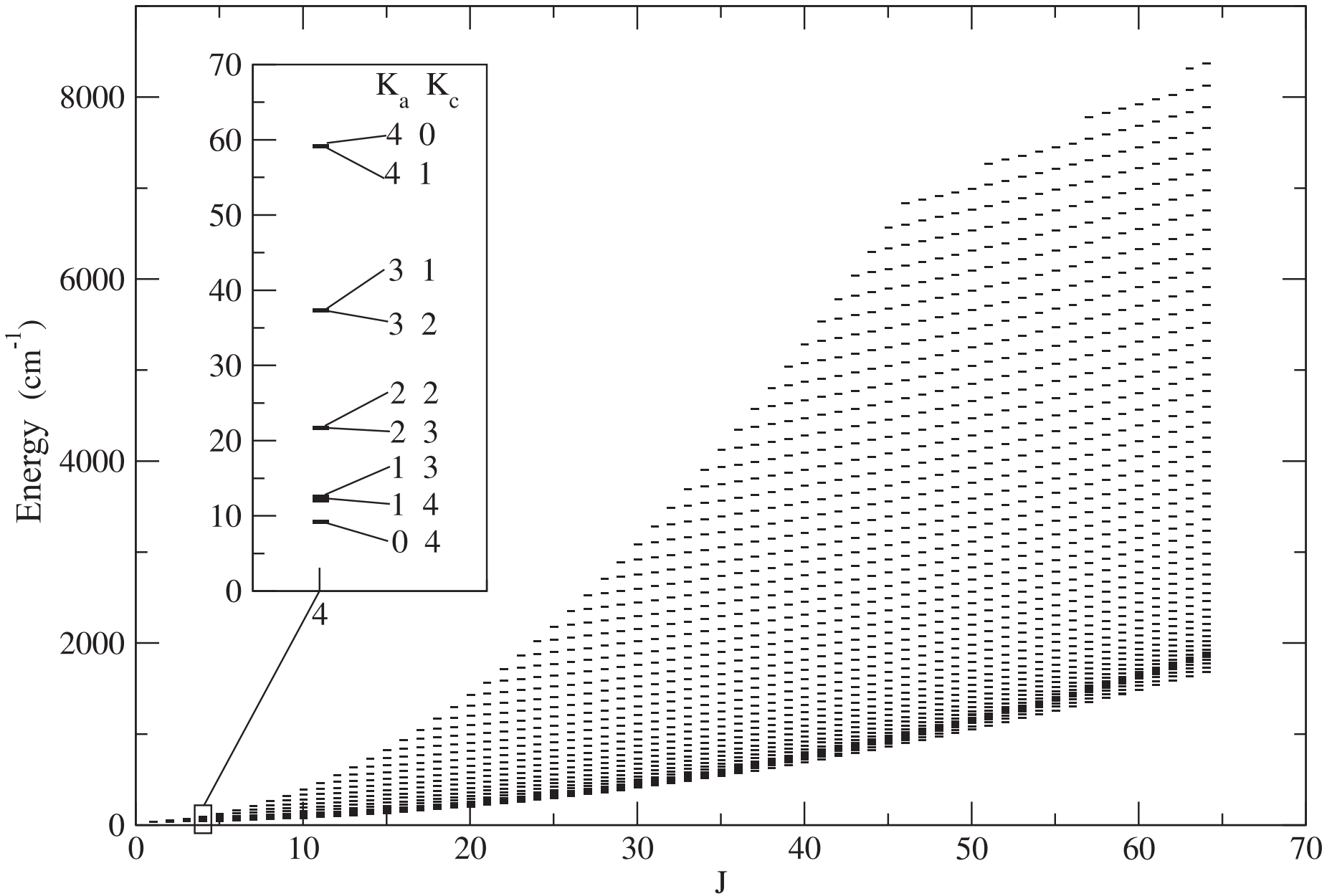


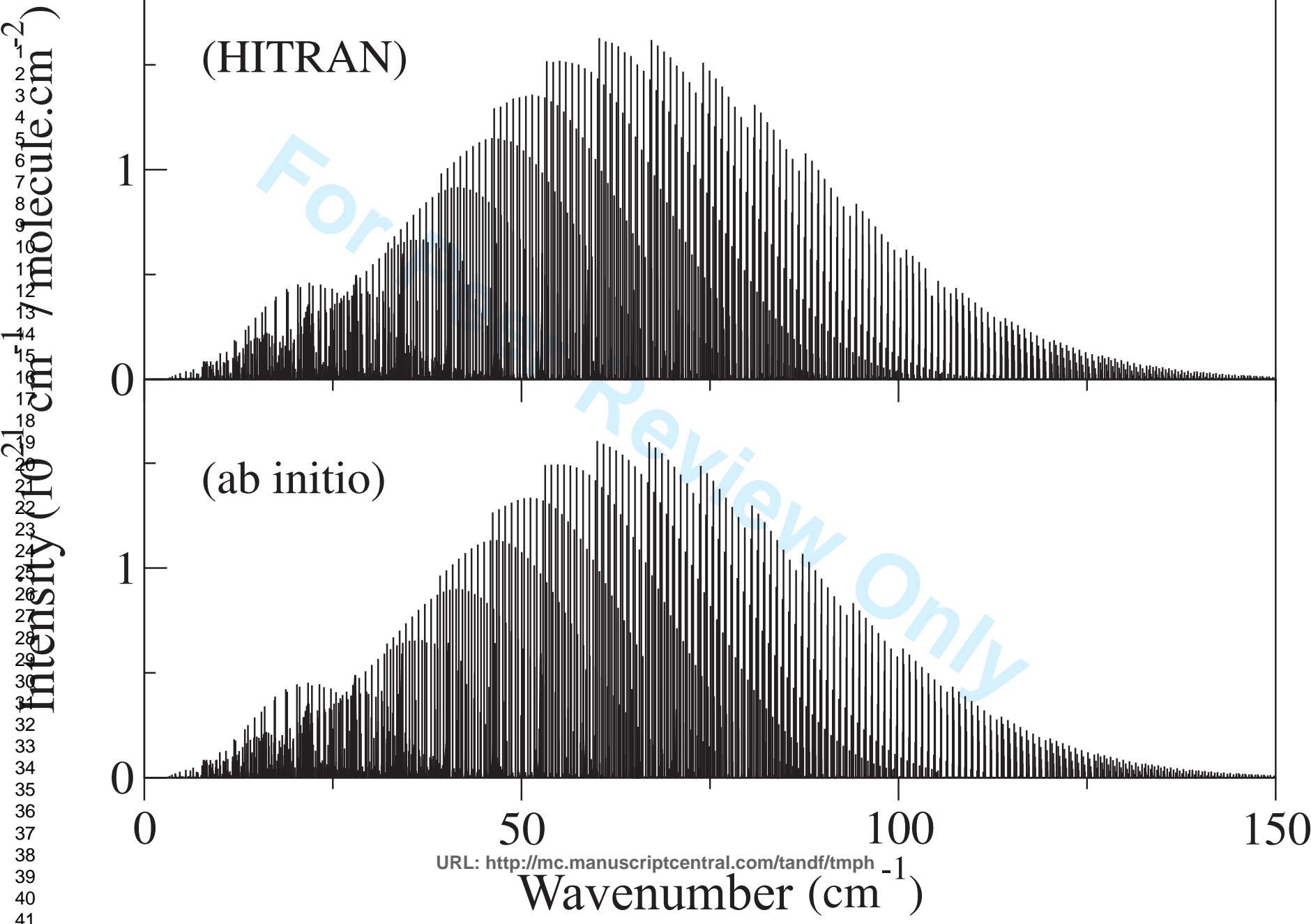
FIG. 9: Mars Global Surveyor thermal emission spectra. The dotted line curve is a night spectrum with random error about 10^{-9} $\text{W}/\text{cm}^2 \cdot \text{sr}/\text{cm}^{-1}$. It is an average over 30°S to 30°N latitude, 42-84 km. The synthetic spectrum (broken curve) is computed for a tangent altitude of 50 km and assumes Boltzmann populations. The solid line curve represents the ab initio spectrum obtained with the spontaneous emission rates (cf. text and Ref. [3]).

In this challenging task, the interplay between theory and experiment will remain essential.

- [1] M. Diehr, P. Rosmus, S. Carter and P.J. Knowles *Mol. Phys.*, **102**, 2181 (2004) and references therein.
- [2] T. Cours, P. Rosmus and Vl. G. Tyuterev *J. Chem. Phys.*, **117**, 5192 (2002) and references therein.
- [3] C. Léonard, M. Diehr, P. Rosmus and W.C. Maguire *Journal of Quantitative Spectroscopy nad Radiative transfer*, **109** 535 (2007) and references therein.
- [4] C. P. Rinsland, J. -M. Flaud, A. Perrin, M. Birk, G. Wagner, A. Goldman, A. Barbe, M. - R. De Backer-Barilly, S. N. Mikhailenko, Vl. G. Tyuterev, M. A. H. Smith, V. Malathy Devi, D. Chrisz Benner, F. Schreier, K. V. Chance, J. Orphal, T. M. Stephen, 2003, *J. Quant. Spectrosc. and Radiat. Transfer*, **82**, 207, and references therein.
- [5] L. S. Rothman, A. Barbe, D. C. Benner, L. R. Brown, C. Camy-Peyret, M. R. Carleer, K. Chance, C. Clerbaux, V. Dana, V. M. Devi, A. Fayt, J. - M. Flaud, R. R. Gamache, A. Goldman, D. Jacquemart, K. W. Jucks, W. J. Lafferty, J. Y. Mandin, S. T. Massie, V. Nemtchinov, D. A. Newnham, A. Perrin, C. P. Rinsland, J. Schroeder, K. M. Smith, M. A. H. Smith, K. Tang, R. A. Toth, J. Vander Auwera, P. Varanasi, and K. Yoshino, *J. Quant. Spectrosc. Radiat. Transfer*, **82**, 5 (2003) and references therein.
- [6] S. Mikhailenko, Yu. Babikov, Vl. G. Tyuterev, A. Barbe, *The Database of Ozone Spectroscopy on WEB (S&MPO)*, Computational Technologies, 2002, v.7, pp.64-70; S&MPO Information system v1.1 available at <http://ozone.iao.ru>, and references therein.
- [7] All electronic structure calculations were performed with MOLPRO package of ab initio programs of H.-J. Werner, and P.-J. Knowles. Further information can be obtained from <http://www.molpro.net/>
- [8] S. Carter, N.C. Handy, and B.T. Sutcliffe, *Mol. Phys.*, **49**, 745 (1983).
- [9] S. Carter, J. Senekowitsch, N. C. Handy, and P. Rosmus, *Mol. Phys.*, **65**, 143 (1988).
- [10] W. Gabriel, E. - A. Reinsch, P. Rosmus, S. Carter, and N. C. Handy, *J. Chem. Phys.*, **99**, 897 (1993).
- [11] J. - M. Flaud, and R. Bacis, *Spectrochimica Acta*, **A 54**, 3 (1998) and references therein.
- [12] O. N. Sulakshina, Yu. G. Borkov, and Vl. G. Tyuterev, *Vl. G J. Chem. Phys.* **113**, 10572 (2000) and references therein.
- [13] D. Xie, H. Guo, and K. A. Peterson, *J. Chem. Phys.* **112**, 8378 (2000) and references therein.
- [14] J. K. Thompson, S. Rainville, and D. E. Pritchard, *Nature (london)*, **430**, 58 (2004).
- [15] M. Cheng, J.M. Brown, P. Rosmus, R. Linguerri, N. Komihia and E.G. Myers, *Phys. Rev. A*, **75**, 012502 (2007) and references therein.
- [16] K. B. Laughlin, G. A. Blake, R. C. Cohen, D. C. Hovde, and R. J. Saykally, *Phys. Rev. Lett.* **58**, 996 (1987).
- [17] K. B. Laughlin, G. A. Blake, R. C. Cohen, and R. J. Saykally, *J. Chem. Phys.* **90**, 1358 (1989).
- [18] M. Molski, *Chem. Phys. Lett.* **342**, 293 (2001).
- [19] E. A. Engel, N. Doss, G. J. Harris, and J. Tennyson, *Mon. Not. R. Astron. Soc.* **357**, 471 (2005),
- [20] T. Wytttenbach, C. G. Beggs, and M. T. Bowers, *Chem. Phys. Lett.* **177**, 239 (1991).
- [21] G. Chambaud, and P. Rosmus, *Chem. Phys. Lett.* **165**, 429 (1990).
- [22] H. J. Wörner, R. van der Veen, and F. Merkt, *Phys. Rev. Lett.* **97**, 173003 (2006).





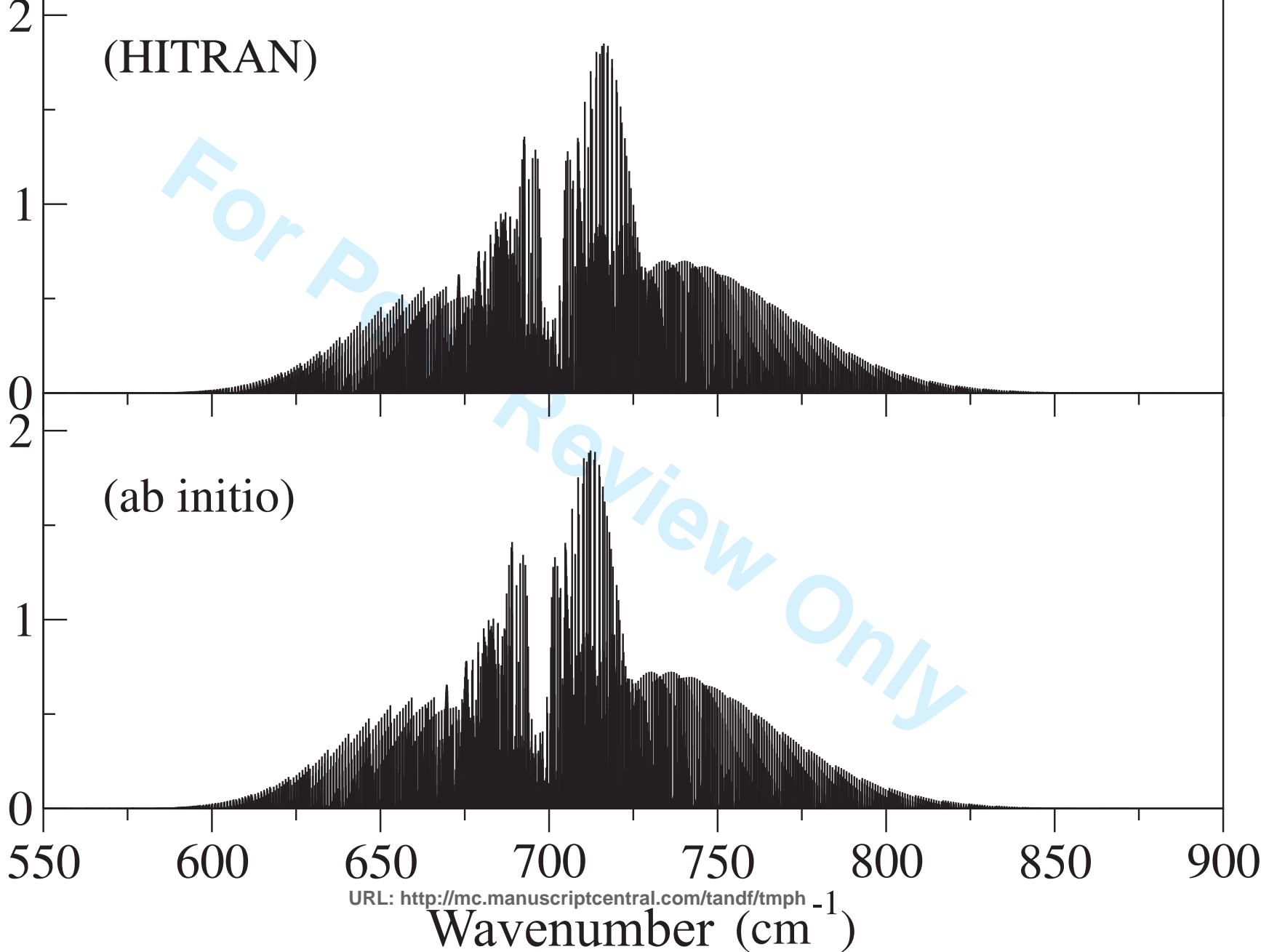


1
2
3
4
5
6
7
8
9
10
11
12
13
14
15
16
17
18
19
20
21
22
23
24
25
26
27
28
29
30
31
32
33
34
35
36
37
38
39
40
41

2
3
4
5
6
7
8
9
10
11
12
13
14
15
16
17
18
19
20
21
22
23
24
25
26
27
28
29
30
31
32
33
34
35
36
37
38
39
40
41

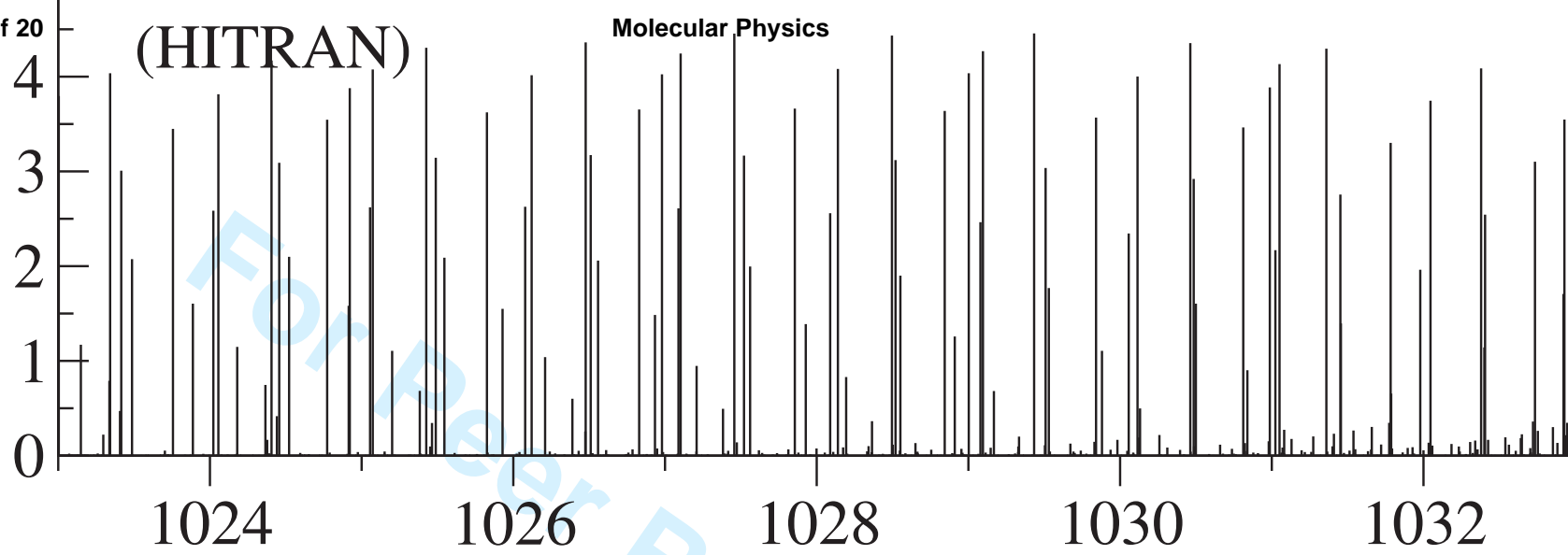
(HITRAN)

(ab initio)

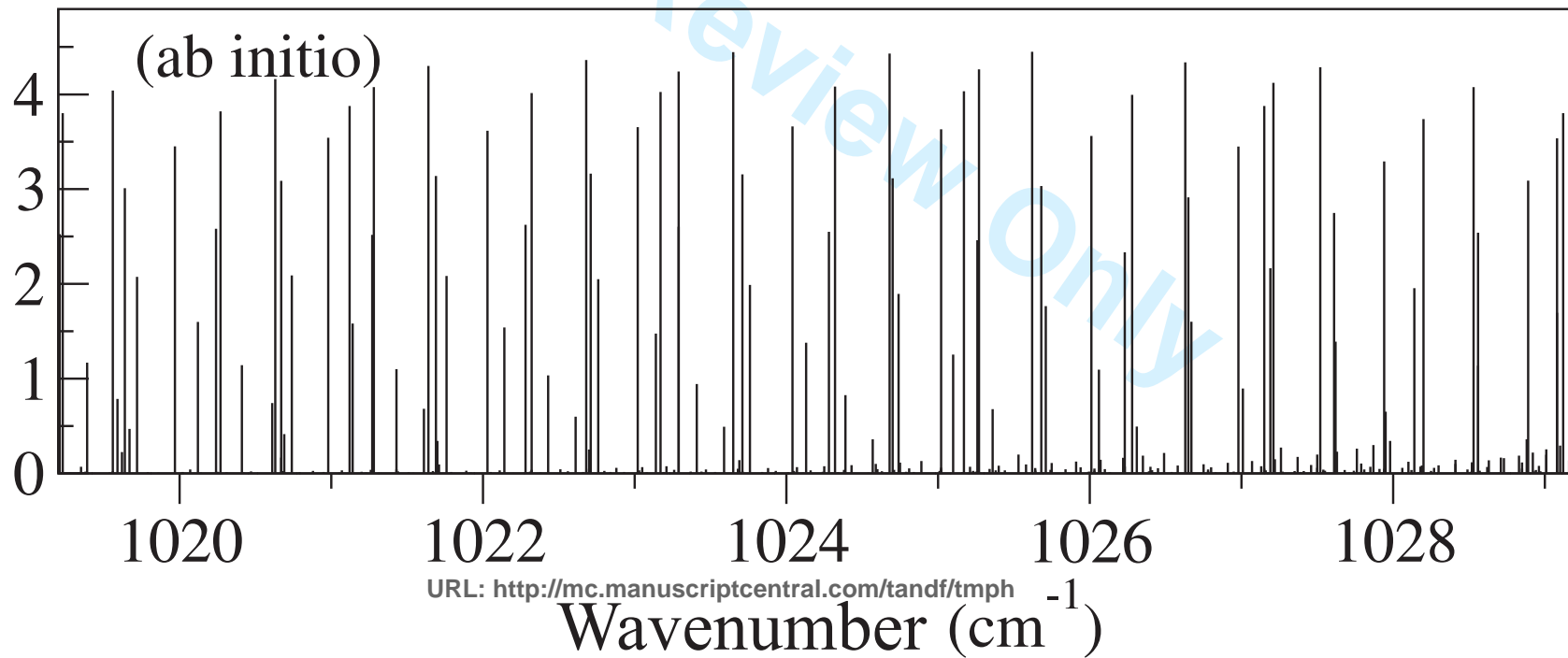


1
2
3
4
5
6
7
8
9
10
11
12
13
14
15
16
17
18
19
20
21
22
23
24
25
26
27
28
29
30
31
32
33
34
35
36
37
38
39
40

(HITRAN)



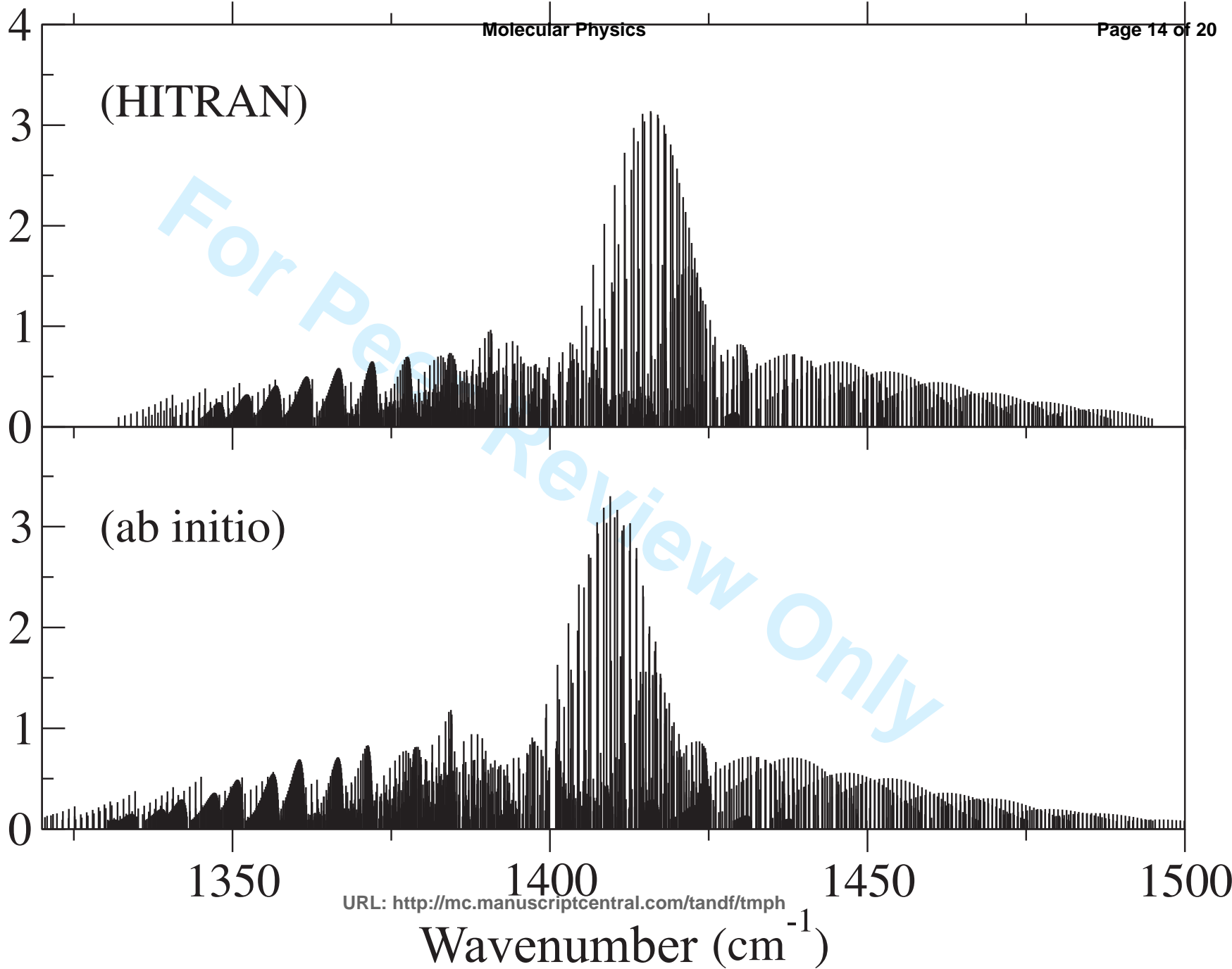
(ab initio)

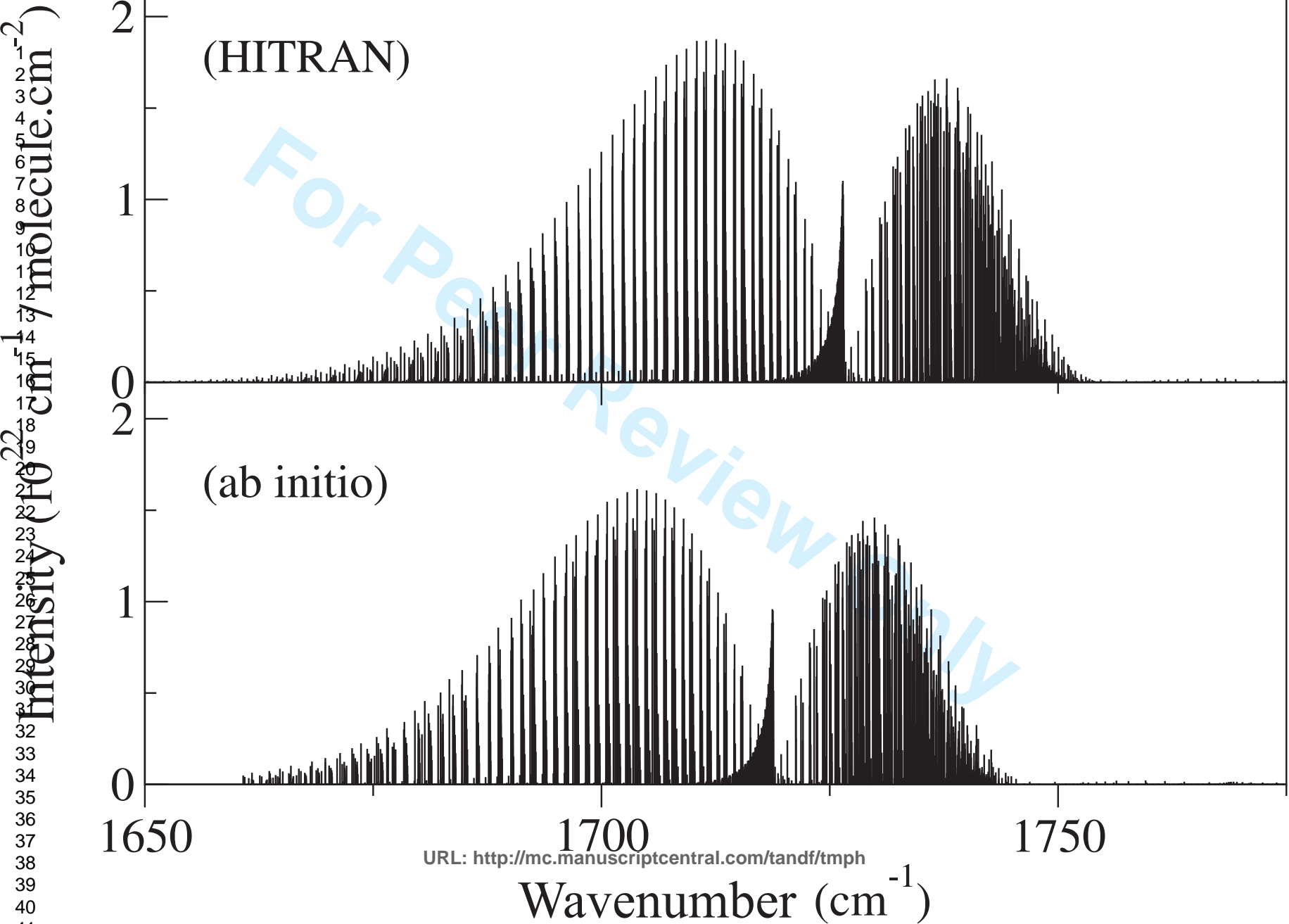


2
3
4
5
6
7
8
9
10
11
12
13
14
15
16
17
18
19
20
21
22
23
24
25
26
27
28
29
30
31
32
33
34
35
36
37
38
39
40
41

(HITRAN)

(ab initio)





2
3
4
5
6
7
8
9
10
11
12
13
14
15
16
17
18
19
20
21
22
23
24
25
26
27
28
29
30
31
32
33
34
35
36
37
38
39
40
...

1
2
3
4
5
6
7
8
9
10
11
12
13
14
15
16
17
18
19
20
21
22
23
24
25
26
27
28
29
30
31
32
33
34
35
36
37
38