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ERROR STUDIES OF THE COUPLING DARCY-STOKES SYSTEM

FIDA EL CHAMI[†], GIHANE MANSOUR[‡] AND TONI SAYAH[‡]

ABSTRACT. In this paper we study the Coupling Darcy-Stokes Systems. We establish a coupled variational formulation with the velocity and the pressure. The velocity is approximated with **curl** conforming finite elements and the pressure with standard continuous elements. We establish optimal a priori and a posteriori error estimates.

Keywords Stokes equations, Darcy equations, a priori and a posteriori errors.

1. INTRODUCTION.

The numerical simulation of underground flows can be treated as cracks in porous media. Indeed, the flow of a viscous incompressible fluid in a porous medium is usually modelled by Darcy equations and, when the thickness of the crack is too large to be neglected, the Stokes system must be considered in the crack and coupled with these equations. In this work, we consider the following system already studied in [2], [3] and [4]:

Let Ω and Ω_F be bounded connected open domains in \mathbb{R}^3 with Lipschitz-continuous boundaries, such that $\bar{\Omega}_F$ is contained in Ω . For simplicity, we also assume that Ω_F is simply connected and has a connected boundary. We set $\Omega_P = \Omega \setminus \bar{\Omega}_F$ and we denote by $\Gamma = \partial\Omega_F$ the interface between Ω_F and Ω_P . Let also \mathbf{n} stand for the unit outward normal vector to Ω_P on its boundary $\partial\Omega_P$.

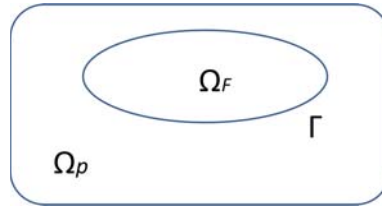


FIGURE 1. The geometry

We consider the following system of equations:

$$(P) \left\{ \begin{array}{ll} \mu \mathbf{u} + \nabla p = \mathbf{f} & \text{in } \Omega_P \\ -\nu \Delta \mathbf{u} + \nabla p = \mathbf{f} & \text{in } \Omega_F \\ \operatorname{div} \mathbf{u} = 0 & \text{in } \Omega_P \cup \Omega_F \\ \mathbf{u} \cdot \mathbf{n} = 0 & \text{on } \partial\Omega \\ (\mathbf{u}|_{\Omega_P} - \mathbf{u}|_{\Omega_F}) \cdot \mathbf{n} = 0 & \text{on } \Gamma \\ p|_{\Omega_P} - p|_{\Omega_F} = 0 & \text{on } \Gamma \\ \operatorname{curl} \mathbf{u}|_{\Omega_F} \times \mathbf{n} = 0 & \text{on } \Gamma \end{array} \right. \quad (1.1)$$

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[†] Faculté des Sciences II, Université Libanaise, Département de mathématiques, B.P. 90656, Fanar-Maten, Liban.

[‡] Faculté des Sciences, Université Saint-Joseph, B.P 11-514 Riad El Solh, Beyrouth 1107 2050, Liban.

where \mathbf{u} is the velocity, p the pressure, \mathbf{f} the density of body forces and ν and μ positive constants.

For the Vorticity-velocity-pressure formulation for the Stokes problem, we refer to [6], [7] and [12]. In [1] and [8], we treat the Stokes problem with non-standard boundary conditions and we introduce velocity-pressure weak formulation. In [2], [3] and [4], the basic idea consists in introducing the vorticity $\mathbf{w} = \mathbf{curl} \mathbf{u}$ as a new unknown on the fluid domain Ω_F . However, we treat in this work the same systems with the unknowns \mathbf{u} and p without introducing the vorticity \mathbf{w} . Since, we can discretize the pressure and the velocity independently without a discrete inf-sup condition to obtain matrix systems with an optimal dimension and optimal time of resolution.

The outline of the paper is as follows:

- In Section 2, we introduce the problem, establish a decoupled variational formulation and prove its wellposedness.
- In section 3, we introduce the finite elements and a fully discrete system using the \mathbf{curl} conforming finite elements for the velocity and the standard continuous elements for the pressure.
- We prove the a priori and a posteriori estimates in sections 4 and 5, respectively.

2. ANALYSIS OF THE MODEL

In all the paper, we suppose that $\mathbf{f} \in L^2(\Omega)^3$ and we denote by C and c generic positive constants.

In order to write the variational formulation of the previous problem, we introduce the following spaces:

$$\begin{aligned} W^{m,p}(\Omega) &= \{\mathbf{v} \in L^p(\Omega)^3, \partial^\alpha \mathbf{v} \in L^p(\Omega)^3, \forall |\alpha| \leq m\}, \\ H^m(\Omega) &= W^{m,2}(\Omega), \end{aligned}$$

equipped with the undermentioned semi-norm and norm:

$$|\mathbf{v}|_{m,p,\Omega} = \left(\sum_{|\alpha|=m} \int_{\Omega} |\partial^\alpha \mathbf{v}(x)|^p dx \right)^{1/p} \quad \text{and} \quad \|\mathbf{v}\|_{m,p,\Omega} = \left(\sum_{k \leq m} |\mathbf{v}|_{k,p,\Omega}^p \right)^{1/p}$$

As usual, we shall omit p when $p = 2$ and denote by (\cdot, \cdot) the scalar product of $L^2(\Omega)$. Also, recall the familiar notation:

$$H_0^1(\Omega) = \{\mathbf{v} \in H^1(\Omega); \mathbf{v} = 0 \text{ on } \partial\Omega\},$$

with the Poincaré inequality

$$\forall \mathbf{v} \in H_0^1(\Omega), \|\mathbf{v}\|_{0,\Omega} \leq C|\mathbf{v}|_{1,\Omega}. \quad (2.1)$$

Finally, we introduce the spaces:

$$\begin{aligned} H(\text{div}, \Omega) &= \{\mathbf{v} \in L^2(\Omega)^3, \text{div } \mathbf{v} \in L^2(\Omega)\}, \quad H_0(\text{div}, \Omega) = \{\mathbf{v} \in H(\text{div}, \Omega), \mathbf{v} \cdot \mathbf{n} = 0 \text{ on } \partial\Omega\}, \\ H(\mathbf{curl}, \Omega) &= \{\mathbf{v} \in L^2(\Omega)^3, \mathbf{curl} \mathbf{v} \in L^2(\Omega)^3\}, \quad H_0(\mathbf{curl}, \Omega) = \{\mathbf{v} \in H(\mathbf{curl}, \Omega), \mathbf{v} \times \mathbf{n} = 0 \text{ on } \partial\Omega\}, \end{aligned}$$

normed respectively by:

$$\|\mathbf{v}\|_{H(\text{div}, \Omega)} = \left(\|\mathbf{v}\|_{0,\Omega}^2 + \|\text{div } \mathbf{v}\|_{0,\Omega}^2 \right)^{1/2},$$

and

$$\|\mathbf{v}\|_{H(\mathbf{curl}, \Omega)} = \left(\|\mathbf{v}\|_{0,\Omega}^2 + \|\mathbf{curl} \mathbf{v}\|_{0,\Omega}^2 \right)^{1/2}.$$

We recall that the trace operator: $\mathbf{v} \rightarrow \mathbf{v} \cdot \mathbf{n}$ is continuous from $H(\text{div}, \Omega)$ onto $H^{-1/2}(\partial\Omega)$ and the jump $(\mathbf{v}|_{\Omega_P} - \mathbf{v}|_{\Omega_F}) \cdot \mathbf{n}$ vanishes on Γ .

We recall an important lemma which will be useful in the rest of the paper:

Lemma 2.1. *There exists a unique solution \mathbf{w} in $H^1(\Omega)/\mathbb{R}$ such that*

$$(\nabla \mathbf{w}, \nabla \mathbf{v}) = (\mathbf{f}, \nabla \mathbf{v}), \quad \forall \mathbf{v} \in H^1(\Omega)/\mathbb{R},$$

and there exists a positive constant C such that

$$\|\mathbf{w}\|_{1,\Omega} \leq C \|\mathbf{f}\|_{0,\Omega}.$$

For the following regularity theorem, we refer to Nedelec [10].

Theorem 2.2. *All functions $\mathbf{v} \in L^2(\Omega_F)^3$ satisfying:*

$$\operatorname{div} \mathbf{v} = 0, \quad \mathbf{curl} \mathbf{v} \in L^2(\Omega_F)^3, \quad \mathbf{v} \cdot \mathbf{n} = 0 \text{ on } \Gamma,$$

verify

$$\|\mathbf{v}\|_{0,\Omega_F} \leq C \|\mathbf{curl} \mathbf{v}\|_{0,\Omega_F}$$

If Ω_F is convex, \mathbf{v} belongs to $H^1(\Omega_F)$ and we have

$$\|\mathbf{v}\|_{1,\Omega_F} \leq C \|\mathbf{curl} \mathbf{v}\|_{0,\Omega_F}.$$

We have also from [2] (page 10 lemma 2.2) the next lemma:

Lemma 2.3. *For all functions $\mathbf{v} \in L^2(\Omega)^3$ satisfying:*

$$\mathbf{curl} \mathbf{v}|_{\Omega_F} \in L^2(\Omega_F)^3, \quad \operatorname{div} \mathbf{v} = 0 \text{ in } \Omega, \quad \mathbf{v} \cdot \mathbf{n} = 0 \text{ on } \partial\Omega$$

there exist a constant α_0 such that:

$$\|\mathbf{v}\|_{0,\Omega_F} \leq \alpha_0 \left(\|\mathbf{v}\|_{0,\Omega}^2 + \|\mathbf{curl} \mathbf{v}\|_{0,\Omega_F}^2 \right)^{1/2}.$$

In order to give the variational formulation of the problem (P), we introduce the space:

$$X = \{\mathbf{v} \in L^2(\Omega)^3, \mathbf{curl} \mathbf{v}|_{\Omega_F} \in L^2(\Omega_F)^3\},$$

provided with the norm

$$\|\mathbf{v}\|_X = \left(\|\mathbf{v}\|_{0,\Omega}^2 + \|\mathbf{curl} \mathbf{v}\|_{0,\Omega_F}^2 \right)^{1/2}.$$

We consider the following weak variational formulation, denoted by (V):

Find $\mathbf{u} \in X$ and $p \in H^1(\Omega)/\mathbb{R}$ such that

$$\mu(\mathbf{u}, \mathbf{v})_{\Omega_P} + \nu (\mathbf{curl} \mathbf{u}, \mathbf{curl} \mathbf{v})_{\Omega_F} + (\nabla p, \mathbf{v}) = (\mathbf{f}, \mathbf{v}), \quad \forall \mathbf{v} \in X \quad (2.2)$$

$$(\nabla q, \mathbf{u}) = 0, \quad \forall q \in H^1(\Omega)/\mathbb{R} \quad (2.3)$$

Theorem 2.4. *The problem (P) is equivalent to the weak variational formulation (V).*

Proof: Suppose that $(\mathbf{u}, p) \in X \times H^1(\Omega)/\mathbb{R}$ is a solution of (V).

The equation (2.3) gives the third equation of (P) in the distribution sense. Since $\mathcal{D}(\Omega_P) \subset X$ and $\mathcal{D}(\Omega_F) \subset X$, we deduce from (2.2) and the relation $-\Delta \mathbf{u} = \mathbf{curl} \mathbf{curl} \mathbf{u}$ (as we have $\operatorname{div} \mathbf{u} = 0$) the first and the second equation of (P). On the other hand, $\mathbf{u} \in H(\operatorname{div}, \Omega)$ and $(\mathbf{u}, \nabla q) = 0$ for all $q \in H^1(\Omega)$ gives us $(\mathbf{u}|_{\Omega_P} - \mathbf{u}|_{\Omega_F}) \cdot \mathbf{n} = 0$ on Γ and $\mathbf{u} \cdot \mathbf{n} = 0$ in $H^{-1/2}(\partial\Omega)$. Since $p \in H^1(\Omega)/\mathbb{R}$ we have $(p|_{\Omega_P} - p|_{\Omega_F})|_{\Gamma} = 0$. Finally, the second equation of (P) and the equation (2.2) gives

$$(\mathbf{curl} \mathbf{curl} \mathbf{u}, \mathbf{v})_{\Omega_F} = (\mathbf{curl} \mathbf{u}, \mathbf{curl} \mathbf{v})_{\Omega_F}, \quad \forall \mathbf{v} \in H(\mathbf{curl}, \Omega_F),$$

which leads to the last condition $\mathbf{curl} \mathbf{u}|_{\Omega_F} \times \mathbf{n} = 0$ on Γ .

Inversely, let $(\mathbf{u}, p) \in X \times H^1(\Omega)/\mathbb{R}$ be a solution of (P). Multiplying the first equation of (P) by $\mathbf{v} \in X$ and integrating over Ω_F and using the Green formula, and the second equation by \mathbf{v} and integrating over Ω_P gives the first equation of (V). The next three equations of (P) give the second equation of (V). \square

The variational formulation (V) can be splitted into a system for the velocity and a Poisson equation for the pressure. Let us introduce the space

$$U = \{\mathbf{v} \in X; (\nabla q, \mathbf{v}) = 0, \quad \forall q \in H^1(\Omega)/\mathbb{R}\}.$$

and remark that every $\mathbf{v} \in U$ verify the lemma 2.3.

The lemmas 2.1 and 2.3 allow us to establish the following theorem:

Theorem 2.5. *The problem (V) is equivalent to the problem*

Find $\mathbf{u} \in U$ such that

$$\mu(\mathbf{u}, \mathbf{v})_{\Omega_F} + \nu(\mathbf{curl} \mathbf{u}, \mathbf{curl} \mathbf{v})_{\Omega_F} = (\mathbf{f}, \mathbf{v}), \quad \forall \mathbf{v} \in U \quad (2.4)$$

Find $p \in H^1(\Omega)/\mathbb{R}$ such that

$$\mu(\mathbf{u}, \nabla q)_{\Omega_F} + (\nabla p, \nabla q) = (\mathbf{f}, \nabla q), \quad \forall q \in H^1(\Omega)/\mathbb{R}. \quad (2.5)$$

Furthermore, there exists a unique solution and we have the following bounds

$$\begin{aligned} \|\mathbf{u}\|_{0, \Omega_F} + \|\mathbf{curl} \mathbf{u}\|_{0, \Omega_F} &\leq C_1 \|\mathbf{f}\|_{0, \Omega}, \\ |p|_{1, \Omega} &\leq C_2 \|\mathbf{f}\|_{0, \Omega}. \end{aligned}$$

Proof: The equivalence of the two problems comes from the fact that every element $\mathbf{v} \in X$ can be written as $\mathbf{v} = \mathbf{w} + \nabla q$ where $\mathbf{w} \in U$ and $q \in H^1(\Omega)/\mathbb{R}$.

The Lax-Milgram theorem gives the existence and the uniqueness of the solution of (2.4). Having the velocity, the Lax-Milgram theorem gives the existence and the uniqueness of the solution of (2.5). We obtain the inequalities by first taking $\mathbf{v} = \mathbf{u}$ in the equation (2.4), next by taking $q = p$ in the equation (2.5). \square

We denote by (V_1) the problem defined by (2.4) and (2.5). Then, it is easy to show that (V_1) is equivalent to the following problem denoted (V_2) :

Find $\mathbf{u} \in X$, $p \in H^1(\Omega)/\mathbb{R}$ such that

$$\mu(\mathbf{u}, \mathbf{v})_{\Omega_F} + \nu(\mathbf{curl} \mathbf{u}, \mathbf{curl} \mathbf{v})_{\Omega_F} + (\nabla p, \mathbf{v}) = (\mathbf{f}, \mathbf{v}), \quad \forall \mathbf{v} \in X \quad (2.6)$$

$$-\mu(\mathbf{u}, \nabla q)_{\Omega_F} + (\nabla p, \nabla q) = (\mathbf{f}, \nabla q), \quad \forall q \in H^1(\Omega)/\mathbb{R}. \quad (2.7)$$

3. FINITE ELEMENT DISCRETIZATION

In what follows and for simplicity, we make the further assumption that both Ω and Ω_F are polyhedra. We introduce a regular family of triangulation $(\tau_h)_h$ in the sense that:

- for each h , $\bar{\Omega}$ is the union of all elements of τ_h ;
- for each h , the intersection of two different elements of τ_h , if not empty, is a corner, a whole edge or a whole face of both of them;
- the ratio of the diameter h_κ of an element κ in τ_h to the diameter of its inscribed sphere is bounded by a constant independent of κ and h ;

As usual, h denotes the maximum of the diameters of the elements of τ_h . We denote by τ_h^F (resp. τ_h^P) the set of elements κ of τ_h which are contained in Ω_F (resp. Ω_P).

Next, for each κ in τ_h , we introduce the spaces $\mathbb{P}_0(\kappa)$ of the restrictions to κ of constant functions on \mathbb{R}^3 , $\mathbb{P}_1(\kappa)$ of the restrictions to κ of affine functions on \mathbb{R} and the space $\mathbb{P}_K(\kappa)$ of the restrictions to κ of polynomials \mathbf{v} of the form:

$$\mathbf{v}(\mathbf{x}) = \mathbf{a} + \mathbf{b} \times \mathbf{x}, \quad \mathbf{a} \in \mathbb{R}^3, \mathbf{b} \in \mathbb{R}^3.$$

The space $\mathbb{P}_K(\kappa)$ and the corresponding finite elements are studied in [9]. Their degrees of freedom are the average flux along the edges $\int_l (\mathbf{v}, \mathbf{t}) dl$, for the six edges l of κ , \mathbf{t} is the direction vector of l .

Hence, we associate the operator r_κ where $r_\kappa(\mathbf{u})$ is the unique polynomial of \mathbb{P}_K that has the same flux along the edges as \mathbf{u} . We define also the operator I_κ where $I_\kappa(q)$ is the unique polynomial of $\mathbb{P}_1(\kappa)$ that has the same values on the vertex of κ as q . We have for all $\kappa \in \tau_h$:

$$r_\kappa(\nabla q) = \nabla I_\kappa(q) \quad \forall q \in W^{2,t}(\kappa) \quad \text{for some } t > 2.$$

Next, let us introduce the discrete spaces:

$$X_h = \{\mathbf{u}_h \in X; \mathbf{u}_h|_\kappa \in \mathbb{P}_K(\kappa), \forall \kappa \in \tau_h\}, \quad (3.1)$$

$$Q_h = \{q_h \in C^0(\bar{\Omega}); q_h|_\kappa \in \mathbb{P}_1(\kappa), \forall \kappa \in \tau_h\}, \quad (3.2)$$

With these spaces, the finite dimensional analogues of U is:

$$U_h = \{\mathbf{v}_h \in X_h; (\nabla q_h, \mathbf{v}_h) = 0, \forall q_h \in Q_h\},$$

We define the interpolation operators r_h from $H^1(\Omega)^3$ onto X_h , I_h from $H^2(\Omega)$ onto Q_h by

$$r_h u = r_\kappa(u) \text{ on } \kappa, \forall \kappa \in \tau_h \quad (\text{similarly for } I_h).$$

We have the following result:

Theorem 3.1. *Assume that the triangulation τ_h is regular. We have:*

$$\|\mathbf{u} - r_h \mathbf{u}\|_{0,\Omega} + h \|\mathbf{curl}(\mathbf{u} - r_h \mathbf{u})\|_{0,\Omega} \leq Ch \|\mathbf{u}\|_{1,t,\Omega}, \forall \mathbf{u} \in W^{1,t}(\Omega)^3, \text{ for some } t > 2.$$

Moreover, when $\mathbf{u} \in H^k(\Omega)^3$, $k = 1, 2$, we have:

$$\|\mathbf{u} - r_h \mathbf{u}\|_{0,\Omega} \leq Ch^k \|\mathbf{u}\|_{k,\Omega}$$

and, when $\mathbf{u} \in H^2(\Omega)^3$ we have:

$$\|\mathbf{curl}(\mathbf{u} - r_h \mathbf{u})\|_{0,\Omega} \leq Ch \|\mathbf{u}\|_{2,\Omega}$$

Theorem 3.2. *Let Ω be a polyhedron and Ω_F a convex polyhedron. Let τ_h be a uniformly regular family of triangulation of Ω . We have:*

$$\|\mathbf{u}_h\|_{0,\Omega_F} \leq \alpha_0 (\|\mathbf{u}_h\|_{0,\Omega_P}^2 + \|\mathbf{curl} \mathbf{u}_h\|_{0,\Omega_F}^2)^{1/2}, \quad \forall \mathbf{u}_h \in U_h \quad (3.3)$$

Proof: The inequality (3.3) is valid for every $\mathbf{v} \in U$. Let Ω_F be convex, for every function \mathbf{u}_h in U_h , we consider the Dirichlet problem:

$$(\nabla z, \nabla \mu)_{\Omega_F} = (\mathbf{u}_h, \nabla \mu)_{\Omega_F} \quad \forall \mu \in H^1(\Omega_F)/\mathbb{R}.$$

The difference $\mathbf{w} = \mathbf{u}_h - \nabla z$ belongs to the space

$$U_{\Omega_F} = \{\mathbf{v} \in H(\mathbf{curl}, \Omega_F); (\mathbf{v}, \nabla q)_{\Omega_F} = 0 \forall q \in H^1(\Omega_F)/\mathbb{R}\},$$

and $\mathbf{curl} \mathbf{w} = \mathbf{curl} \mathbf{u}_h$. It follows from theorem 2.2 that

$$\|\mathbf{w}\|_{1,\Omega_F} \leq C_1 \|\mathbf{curl} \mathbf{w}\|_{0,\Omega_F}.$$

Therefore, we can apply the interpolation operator r_h to \mathbf{w} , and \mathbf{u}_h splits into:

$$\mathbf{u}_h = r_h \mathbf{w} + \nabla z_h \quad \text{with } z_h \in Q_h.$$

Hence

$$\|\mathbf{u}_h\|_{0,\Omega_F} \leq \|r_h \mathbf{w} - \mathbf{w}\|_{0,\Omega_F} + \|\mathbf{w}\|_{0,\Omega_F} + \|\nabla z_h\|_{0,\Omega_F}.$$

Since on one hand

$$\|\mathbf{w} - r_h \mathbf{w}\|_{0,\Omega_F} \leq C_2 \|\mathbf{w}\|_{1,\Omega_F} \leq C_3 \|\mathbf{curl} \mathbf{w}\|_{0,\Omega_F},$$

and on the other hand

$$\|\mathbf{w}\|_{0,\Omega_F} \leq \|\mathbf{w}\|_{1,\Omega_F} \leq C_4 \|\mathbf{curl} \mathbf{w}\|_{0,\Omega_F}.$$

We see that it suffices to estimate $\|\nabla z_h\|_{0,\Omega_F}$. For all $\mu_h \in Q_h$ we have:

$$\begin{aligned} (\nabla z_h, \nabla \mu_h)_{\Omega_F} &= (\mathbf{u}_h - r_h \mathbf{w}, \nabla \mu_h)_{\Omega_F} \\ &= (\mathbf{u}_h, \nabla \mu_h)_{\Omega_F} + (\mathbf{w} - r_h \mathbf{w}, \nabla \mu_h)_{\Omega_F} \\ &= -(\mathbf{u}_h, \nabla \mu_h)_{\Omega_P} + (\mathbf{w} - r_h \mathbf{w}, \nabla \mu_h)_{\Omega_F}, \\ &\leq \|\mathbf{u}_h\|_{0,\Omega_P} \|\nabla \mu_h\|_{0,\Omega_P} + C_3 \|\mathbf{curl} \mathbf{w}\|_{0,\Omega_F} \|\nabla \mu_h\|_{0,\Omega_F} \end{aligned}$$

We choose $\mu_h \in Q_h$ such that $\mu_h|_{\bar{\Omega}_F} = z_h|_{\bar{\Omega}_F}$ and $\mu_h = 0$ in Ω_P .

We obtain : $\|\nabla \mu_h\|_{0,\Omega_P} \leq C_1 \|z_h\|_{1/2,\Gamma} \leq C_2 \|\nabla z_h\|_{0,\Omega_F}$ and we deduce

$$\|\nabla z_h\|_{0,\Omega_F} \leq C_3 (\|\mathbf{u}_h\|_{0,\Omega_P} + \|\mathbf{curl} \mathbf{w}\|_{0,\Omega_F})$$

and finally the result. \square

We discretize (V) by:

Find $\mathbf{u}_h \in U_h$ and $p_h \in Q_h/\mathbb{R}$ such that

$$\mu(\mathbf{u}_h, \mathbf{v}_h)_{\Omega_P} + \nu(\mathbf{curl} \mathbf{u}_h, \mathbf{curl} \mathbf{v}_h)_{\Omega_F} + (\nabla p_h, \mathbf{v}_h) = (\mathbf{f}, \mathbf{v}_h) \quad \forall \mathbf{v}_h \in X_h. \quad (3.4)$$

As in the continuous way, the last problem can be splitted to

Find $\mathbf{u}_h \in U_h$ such that

$$\mu(\mathbf{u}_h, \mathbf{v}_h)_{\Omega_P} + \nu(\mathbf{curl} \mathbf{u}_h, \mathbf{curl} \mathbf{v}_h)_{\Omega_F} = (\mathbf{f}, \mathbf{v}_h), \quad \forall \mathbf{v}_h \in U_h, \quad (3.5)$$

Find $p_h \in Q_h/\mathbb{R}$ such that

$$\mu(\mathbf{u}_h, \nabla q_h)_{\Omega_P} + (\nabla p_h, \nabla q_h) = (\mathbf{f}, \nabla q_h), \quad \forall q_h \in Q_h/\mathbb{R}. \quad (3.6)$$

Let Ω_F be convex, it is easy to show, using the theorem 3.2, that these two last discrete problems have a unique solution and we have:

$$\|\mathbf{u}_h\|_{0, \Omega_P} + \|\mathbf{curl} \mathbf{u}_h\|_{0, \Omega_F} \leq C_3 \|\mathbf{f}\|_{0, \Omega}$$

and

$$|p_h|_{1, \Omega} \leq C_4 \|\mathbf{f}\|_{0, \Omega}.$$

It is obvious that the last problem is equivalent to:

Find $\mathbf{u}_h \in U_h$ and $p_h \in Q_h/\mathbb{R}$ such that

$$\mu(\mathbf{u}_h, \mathbf{v}_h)_{\Omega_P} + \nu(\mathbf{curl} \mathbf{u}_h, \mathbf{curl} \mathbf{v}_h)_{\Omega_F} + (\nabla p_h, \mathbf{v}_h) = (\mathbf{f}, \mathbf{v}_h), \quad \forall \mathbf{v}_h \in X_h, \quad (3.7)$$

Find $p_h \in Q_h/\mathbb{R}$ such that

$$-\mu(\mathbf{u}_h, \nabla q_h)_{\Omega_P} + (\nabla p_h, \nabla q_h) = (\mathbf{f}, \nabla q_h), \quad \forall q_h \in Q_h/\mathbb{R}. \quad (3.8)$$

4. A PRIORI ERROR ANALYSIS

In this section, we will establish the error estimates for the pressure and the velocity. First of all, we consider the quantity $\mathbf{u}_h - r_h \mathbf{u}$ and we consider the finite dimensional problem:

Find $\lambda_h \in Q_h/\mathbb{R}$ such that

$$\forall q_h \in Q_h/\mathbb{R}, \quad \int_{\Omega} \nabla \lambda_h \nabla q_h = \int_{\Omega} (\mathbf{u}_h - r_h \mathbf{u}) \nabla q_h$$

which admits a unique solution λ_h such that $\mathbf{w}_h = (\mathbf{u}_h - r_h \mathbf{u}) - \nabla \lambda_h$ is in the space U_h with $\mathbf{curl}(\mathbf{u}_h - r_h \mathbf{u}) = \mathbf{curl} \mathbf{w}_h$.

Furthermore we consider, for all $q_h \in Q_h/\mathbb{R}$, the relation

$$\int_{\Omega} \nabla \lambda_h \nabla q_h = \int_{\Omega} (\mathbf{u}_h - r_h \mathbf{u}) \nabla q_h = - \int_{\Omega} (r_h \mathbf{u} - \mathbf{u}) \nabla q_h$$

which gives by taking $q_h = \lambda_h$ and supposing that $\mathbf{u} \in H^1(\Omega)^3$:

$$|\lambda_h|_{1, \Omega} \leq C h \|\mathbf{u}\|_{1, \Omega}$$

To obtain the a priori error estimate for the velocity, it suffices to show an error estimate of \mathbf{w}_h and we conclude an error estimate of $\mathbf{u} \in H^2(\Omega)^3$ by using the theorem 3.2. Let Ω_F be convex. For all $\mathbf{u}_h \in U_h$,

we have:

$$\begin{aligned}
 \|\mathbf{u} - \mathbf{u}_h\|_X^2 &\leq \|\mathbf{u} - r_h \mathbf{u}\|_{0,\Omega}^2 + \|r_h \mathbf{u} - \mathbf{u}_h\|_{0,\Omega}^2 + \|\mathbf{curl}(\mathbf{u} - r_h \mathbf{u})\|_{0,\Omega_F}^2 + \|\mathbf{curl}(r_h \mathbf{u} - \mathbf{u}_h)\|_{0,\Omega_F}^2 \\
 &\leq C_1(u, \Omega) \left(h^2 + \|r_h \mathbf{u} - \mathbf{u}_h\|_{0,\Omega}^2 + \|\mathbf{curl}(r_h \mathbf{u} - \mathbf{u}_h)\|_{0,\Omega_F}^2 \right) \\
 &\leq C_2(u, \Omega) \left(h^2 + \|\nabla \lambda_h\|_{0,\Omega}^2 + \|\mathbf{w}_h\|_{0,\Omega}^2 + \|\mathbf{curl} \mathbf{w}_h\|_{0,\Omega_F}^2 \right) \\
 &\leq C_3(u, \Omega) \left(h^2 + \|\nabla \lambda_h\|_{0,\Omega}^2 + \|\mathbf{w}_h\|_{0,\Omega_P}^2 + \|\mathbf{curl} \mathbf{w}_h\|_{0,\Omega_F}^2 \right)
 \end{aligned} \tag{4.1}$$

Next, to obtain the error estimate for \mathbf{w} , we consider the difference of the equation (2.2) with $\mathbf{v} = \mathbf{v}_h \in X_h$ and the equation (3.4):

$$\mu \int_{\Omega_P} (\mathbf{u} - \mathbf{u}_h) \mathbf{v}_h + \nu \int_{\Omega_F} \mathbf{curl}(\mathbf{u} - \mathbf{u}_h) \mathbf{curl} \mathbf{v}_h + \int_{\Omega} \nabla(p - p_h) \mathbf{v}_h = 0$$

We insert $\pm r_h \mathbf{u}$ in the first and the second terms, $\pm I_h p$ in the third term and we obtain:

$$\begin{aligned}
 \mu \int_{\Omega_P} (r_h \mathbf{u} - \mathbf{u}_h, \mathbf{v}_h) + \nu \int_{\Omega_F} \mathbf{curl}(r_h \mathbf{u} - \mathbf{u}_h) \mathbf{curl} \mathbf{v}_h = \\
 \mu (r_h \mathbf{u} - \mathbf{u}, \mathbf{v}_h)_{\Omega_P} + \nu (\mathbf{curl}(r_h \mathbf{u} - \mathbf{u}_h), \mathbf{curl} \mathbf{v}_h)_{\Omega_F} - (\nabla(p - I_h p), \mathbf{v}_h) - (\nabla(I_h p - p_h), \mathbf{v}_h)
 \end{aligned}$$

We replace $\mathbf{u}_h - r_h \mathbf{u} = \nabla \lambda_h + \mathbf{w}_h$ and we choose $\mathbf{v}_h = \mathbf{w}_h$ to obtain:

$$\begin{aligned}
 \mu \|\mathbf{w}_h\|_{0,\Omega_P}^2 + \nu \|\mathbf{curl} \mathbf{w}_h\|_{0,\Omega_F}^2 = \\
 -\mu \int_{\Omega_P} \nabla \lambda_h \mathbf{w}_h + \mu (r_h \mathbf{u} - \mathbf{u}, \mathbf{w}_h)_{\Omega_P} + \nu (\mathbf{curl}(r_h \mathbf{u} - \mathbf{u}_h), \mathbf{curl} \mathbf{w}_h)_{\Omega_F} - (\nabla(p - I_h p), \mathbf{w}_h)
 \end{aligned}$$

By supposing that $p \in H^2(\Omega)$ and $\mathbf{u} \in H^2(\Omega)^3$, we deduce using the properties of r_h and I_h , the formula $a.b \leq \frac{1}{2\varepsilon} a^2 + \frac{1}{2} \varepsilon b^2$ with a suitable choice of ε and the previous upper bound of λ :

$$\mu \|\mathbf{w}_h\|_{0,\Omega_P}^2 + \nu \|\mathbf{curl} \mathbf{w}_h\|_{0,\Omega_F}^2 \leq C(\Omega) h^2 (\|\mathbf{u}\|_{2,\Omega}^2 + \|p\|_{2,\Omega}^2)$$

Now, we will show an estimate for the pressure. We subtract the equation (2.5) with $q = q_h \in Q_h/\mathbb{R}$ from the equation (3.6) to get

$$\mu \int_{\Omega_P} (\mathbf{u} - \mathbf{u}_h) \nabla q_h + \int_{\Omega} \nabla(p - p_h) \nabla q_h = 0$$

We insert $\pm I_h p$ in the second term and we choose $q_h = p_h - I_h p$ to obtain

$$|p_h - I_h p|_{1,\Omega} \leq \mu \|(\mathbf{u} - \mathbf{u}_h)\|_{0,\Omega_P} + |p - I_h p|_{1,\Omega}$$

We deduce the error estimate:

$$\begin{aligned}
 |p - p_h|_{1,\Omega} &\leq |p - I_h p|_{1,\Omega} + |I_h p - p_h|_{1,\Omega} \\
 &\leq C(\Omega) (\|p\|_{2,\Omega} + \|\mathbf{u}\|_{2,\Omega}) h
 \end{aligned} \tag{4.2}$$

Theorem 4.1. *If $\mathbf{u} \in H^2(\Omega)^3$ and $p \in H^2(\Omega)$, the theoretical solution (\mathbf{u}, p) of the problem (2.4)-(2.5) and the numerical solution (\mathbf{u}_h, p_h) of the problem (3.5)-(3.6) verify the error estimate:*

$$\|\mathbf{u} - \mathbf{u}_h\|_X + |p - p_h|_{1,\Omega} \leq C(\Omega, \mathbf{u}, p) h \tag{4.3}$$

5. A POSTERIORI ERROR ANALYSIS

We now intend to prove a posteriori error estimates between the exact solution (\mathbf{u}, p) of the problem (2.4)-(2.5) and the numerical solution (\mathbf{u}_h, p_h) of the problem (3.7)-(3.8). In all the rest of the paper, we suppose that $\mathbf{f} \in H(\text{div}, \Omega)$.

We first introduce the space

$$Z_h = \{\mathbf{g}_h \in L^2(\Omega)^3; \forall \kappa \in \tau_h, \mathbf{g}_h|_{\kappa} \in \mathbb{P}_0(\kappa)\}$$

and we fix an approximation \mathbf{f}_h of the data \mathbf{f} in Z_h .

Next, we denote by ε_h the set of all faces of the elements. For every element κ in τ_h , we denote by ε_κ the set of faces of κ that are not contained in Γ , Δ_κ (resp. Δ_κ^F or Δ_κ^P) the set of union of elements of τ_h that intersect κ (resp. contained in Ω^F or contained in Ω^P), Δ_e (resp. Δ_e^F or Δ_e^P) the union of elements of τ_h that intersect the face e (resp. contained in Ω^F or contained in Ω^P), h_κ the diameter of κ and h_e the diameter of the face e . Also, \mathbf{n}_κ stands for the unit outward normal vector to κ on $\partial\kappa$ and $[\cdot]_e$ the jump through the face e of κ .

For the demonstration of the next theorems, we introduce for an element κ of τ_h , the bull function ψ_κ (resp. ψ_e of the face e) which is equal to the product of the $d + 1$ barycentric coordinates associated with the vertices of κ (resp. of e) and \mathcal{L}_e (resp. \mathcal{L}_e^F or \mathcal{L}_e^P) the lifting operator from polynomials defined on e into polynomials defined on the elements κ and κ' containing e (resp. elements contained in Ω^F or contained in Ω^P), which is constructed by affine transformations from a fixed operator on the reference element.

Property 5.1. Denoting by $P_r(e)$ the polynomial of degrees r on e , we have

$$\forall v \text{ polynom of } P_r(e) \quad c \|v\|_{L^2(e)} \leq \|v\psi_e^{1/2}\|_{L^2(e)} \leq c' \|v\|_{L^2(e)}$$

and for any polynom v of $P_r(e)$ which vanishes on ∂e , we have

$$\|\mathcal{L}_e v\|_{L^2(\kappa)} + h_e \|\mathcal{L}_e v\|_{H^1(\kappa)} \leq ch_e^{1/2} \|v\|_{L^2(e)}.$$

We denote by R_h the Clément operator [5]. For any function $q \in H_0^1(\Omega)$, $R_h q \in Q_h$ verifies

$$\begin{aligned} \|q - R_h q\|_{L^2(\kappa)} &\leq ch_\kappa \|q\|_{H^1(\Delta_\kappa)}, \\ \|q - R_h q\|_{L^2(e)} &\leq ch_e^{1/2} \|q\|_{H^1(\Delta_e)}. \end{aligned} \tag{5.1}$$

We also denote by \mathcal{R}_h the Raviart-Thomas operator : for any smooth enough vectorial function \mathbf{v} which is divergence-free on Ω , $\mathcal{R}_h \mathbf{v}$ belongs to X_h and satisfies

$$\forall e \in \varepsilon_h, \quad \int_e (\mathbf{v} - \mathcal{R}_h \mathbf{v}) \cdot \mathbf{n} d\tau = 0.$$

Moreover, this operator satisfies, see [11] : $\forall \mathbf{v}$ in $H^1(\Omega)^3$ and $\forall \kappa$ in τ_h ,

$$\begin{aligned} \|\mathbf{v} - \mathcal{R}_h \mathbf{v}\|_{L^2(\kappa)^3} &\leq ch_\kappa \|\mathbf{v}\|_{H^1(\kappa)^3} \\ \|\mathbf{v} - \mathcal{R}_h \mathbf{v}\|_{L^2(e)^3} &\leq ch_e^{1/2} \|\mathbf{v}\|_{H^1(\Delta_e)^3} \end{aligned} \tag{5.2}$$

To prove the a posteriori estimates, we begin by decomposing $\mathbf{u} - \mathbf{u}_h = \nabla \lambda + \mathbf{w}$ where $\lambda \in H^1(\Omega)$ and $\mathbf{w} \in U$. Then, we establish a posteriori estimate for λ and \mathbf{w} to deduce using the lemma (2.3)

$$\begin{aligned} \|\mathbf{u} - \mathbf{u}_h\|_X^2 &= \|\nabla \lambda + \mathbf{w}\|_{0,\Omega}^2 + \|\mathbf{curl} \mathbf{w}\|_{0,\Omega_F}^2 \\ &\leq C(\|\lambda\|_{1,\Omega}^2 + \|\mathbf{w}\|_{0,\Omega_F}^2 + \|\mathbf{curl} \mathbf{w}\|_{0,\Omega_F}^2) \end{aligned}$$

and we finish with the a posteriori estimate for the pressure.

The error function $\mathbf{u} - \mathbf{u}_h$ belongs to X , there exists a unique solution $\lambda_1 \in H^1(\Omega_F)/\mathbb{R}$ of the problem:

$$\int_{\Omega_F} \nabla \lambda_1 \nabla q = \int_{\Omega_F} (\mathbf{u} - \mathbf{u}_h) \nabla q \quad \forall q \in H^1(\Omega_F)/\mathbb{R},$$

and the function $\mathbf{w}_1 = (\mathbf{u} - \mathbf{u}_h) - \nabla \lambda_1$ belongs to $U_F = \{\mathbf{v} \in H(\mathbf{rot}, \Omega_F) / (\mathbf{v}, \nabla q)_{\Omega_F} = 0 ; \forall q \in H^1(\Omega_F)\}$. We define the function $\tilde{\mathbf{w}}$, equal to \mathbf{w}_1 in Ω_F and $\mathbf{0}$ in Ω_P , which belongs to U and verifies $\mathbf{curl} \tilde{\mathbf{w}} = \mathbf{curl} \mathbf{w}_1$ in Ω_F . Furthermore, there exists a unique solution $\lambda \in H^1(\Omega)/\mathbb{R}$ of the problem:

$$\int_{\Omega} \nabla \lambda \nabla q = \int_{\Omega} (\mathbf{u} - \mathbf{u}_h) \nabla q \quad \forall q \in H^1(\Omega)/\mathbb{R},$$

and the function $\mathbf{w} = (\mathbf{u} - \mathbf{u}_h) - \nabla\lambda$ belongs to U and we have $\mathbf{curl}\mathbf{w} = \mathbf{curl}\mathbf{w}_1 = \mathbf{curl}(\mathbf{u} - \mathbf{u}_h)$ in Ω_F . We have: $\forall q \in H^1(\Omega)$,

$$\begin{aligned} \int_{\Omega} \nabla\lambda \nabla q &= \int_{\Omega} (\mathbf{w} + \nabla\lambda) \nabla q = \int_{\Omega} (\mathbf{u} - \mathbf{u}_h) \nabla q = - \int_{\Omega} \mathbf{u}_h \nabla q \\ &= - \int_{\Omega} \mathbf{u}_h \nabla (q - q_h) = - \frac{1}{2} \sum_{\kappa \in \tau_h} \left(\sum_{e \in \varepsilon_{\kappa}} \int_e [\mathbf{u}_h \cdot \mathbf{n}] (q - q_h) \right) \quad \forall q_h \in Q_h. \end{aligned} \quad (5.3)$$

We introduce the indicators

$$\xi_{\kappa} = \sum_{e \in \varepsilon_{\kappa}} h_e^{1/2} \| [\mathbf{u}_h \cdot \mathbf{n}] \|_{0,e} \quad (5.4)$$

Theorem 5.2. *The following bounds hold*

$$\|\nabla\lambda\|_{0,\Omega} \leq C \left(\sum_{\kappa \in \tau_h} \xi_{\kappa}^2 \right)^{1/2}$$

and

$$\xi_{\kappa} \leq c \|\nabla\lambda\|_{0,\Delta_{\kappa}} \quad (5.5)$$

Proof: First we take, in the equation (5.3), $q = \lambda$ and $q_h = R_h q$, the image of q by the Clément type regularization operator, and we obtain the upper bound. In order to find the lower bound, we take in the equation (5.3) $q_h = 0$ and $q = \mathcal{L}_e([\mathbf{u}_h \cdot \mathbf{n}] \psi_e)$, and we obtain

$$\| [\mathbf{u}_h \cdot \mathbf{n}] \|_{0,e} \leq C \left(h_e^{-1/2} |\lambda|_{1,\kappa \cup \kappa'} \right)$$

which leads to

$$\xi_{\kappa} \leq c \left(|\lambda|_{1,\Delta_{\kappa}} \right). \quad (5.6)$$

□

To find a posteriori estimates for \mathbf{w} , we begin to establish upper and lower bounds for $\mathbf{curl}\mathbf{w}$ in Ω_F . We introduce the indicators

$$\gamma_{\kappa,F} = h_{\kappa} \| \mathbf{f}_h - \nabla p_h \|_{0,\kappa} + \frac{\nu}{2} \sum_{e \in \varepsilon_{\kappa}} h_{\kappa}^{1/2} \| [\mathbf{curl}\mathbf{u}_h \times \mathbf{n}] \|_{0,e}, \quad \text{if } \kappa \in \Omega_F \quad (5.7)$$

where $[\cdot]_e$ is the jump through the face e of κ . If the face e is on Γ , $[\cdot]_e$ will be the trace on e from the domain Ω_F containing κ .

Theorem 5.3. *Let Ω_F be convex. The following bounds hold:*

$$\nu \| \mathbf{curl}\mathbf{w} \|_{0,\Omega_F} \leq C \left(\sum_{\kappa \in \tau_h^F} h_{\kappa}^2 \| \mathbf{f} - \mathbf{f}_h \|_{0,\kappa}^2 + \sum_{\kappa \in \tau_h^F} \gamma_{\kappa,F}^2 \right)^{1/2} \quad (5.8)$$

and

$$\gamma_{\kappa,F} \leq c \sum_{e \in \varepsilon_{\kappa}} \left(\| \mathbf{curl}\mathbf{w} \|_{0,\Delta_e^F}^2 + h_e^2 |p - p_h|_{1,\Delta_e^F}^2 + h_e^2 \| \mathbf{f} - \mathbf{f}_h \|_{0,\Delta_e^F}^2 \right)^{1/2} \quad (5.9)$$

Proof : The error function $\mathbf{u} - \mathbf{u}_h$ verifies, by using the equations (2.2) and (3.4): $\forall \mathbf{v} \in X$ and $\forall \mathbf{v}_h \in X_h$

$$\begin{aligned}
& \mu \int_{\Omega_P} (\mathbf{u} - \mathbf{u}_h) \mathbf{v} + \nu \int_{\Omega_F} \mathbf{curl}(\mathbf{u} - \mathbf{u}_h) \mathbf{curl} \mathbf{v} + \int_{\Omega} \nabla(p - p_h) \mathbf{v} \\
&= \mu \int_{\Omega_P} (\mathbf{u} - \mathbf{u}_h) (\mathbf{v} - \mathbf{v}_h) + \nu \int_{\Omega_F} \mathbf{curl}(\mathbf{u} - \mathbf{u}_h) \mathbf{curl}(\mathbf{v} - \mathbf{v}_h) + \int_{\Omega} \nabla(p - p_h) (\mathbf{v} - \mathbf{v}_h) \\
&= (\mathbf{f}, \mathbf{v} - \mathbf{v}_h) - \mu \int_{\Omega_P} \mathbf{u}_h (\mathbf{v} - \mathbf{v}_h) - \nu \int_{\Omega_F} \mathbf{curl} \mathbf{u}_h \mathbf{curl}(\mathbf{v} - \mathbf{v}_h) - \int_{\Omega} \nabla p_h (\mathbf{v} - \mathbf{v}_h) \quad (5.10) \\
&= ((\mathbf{f} - \mathbf{f}_h), \mathbf{v} - \mathbf{v}_h) + \int_{\Omega_F} (\mathbf{f}_h - \nabla p_h) (\mathbf{v} - \mathbf{v}_h) + \int_{\Omega_P} (\mathbf{f}_h - \nabla p_h - \mu \mathbf{u}_h) (\mathbf{v} - \mathbf{v}_h) \\
&\quad - \nu \int_{\Omega_F} \mathbf{curl} \mathbf{u}_h \mathbf{curl}(\mathbf{v} - \mathbf{v}_h)
\end{aligned}$$

We replace $\mathbf{u} - \mathbf{u}_h$ by $\mathbf{w} + \nabla \lambda$, take $\mathbf{v} = \tilde{\mathbf{w}}$ and $\mathbf{v}_h = \mathcal{R}_h \tilde{\mathbf{w}}$, remark that $\mathbf{curl} \mathbf{w} = \mathbf{curl} \tilde{\mathbf{w}}$ in Ω_F and use the integration by part formula to obtain:

$$\begin{aligned}
\nu \|\mathbf{curl} \mathbf{w}\|_{0, \Omega_F}^2 &= \int_{\Omega_F} (\mathbf{f} - \mathbf{f}_h) (\tilde{\mathbf{w}} - \mathcal{R}_h \tilde{\mathbf{w}}) - \int_{\Omega_F} (\mathbf{f}_h - \nabla p_h) (\tilde{\mathbf{w}} - \mathcal{R}_h \tilde{\mathbf{w}}) \\
&\quad - \frac{\nu}{2} \sum_{\kappa \in \tau_h^F} \sum_{e \in \varepsilon_\kappa} \int_e [\mathbf{curl} \mathbf{u}_h \times \mathbf{n}] (\tilde{\mathbf{w}} - \mathcal{R}_h \tilde{\mathbf{w}}) \quad (5.11)
\end{aligned}$$

Since Ω_F is convex, the theorem 2.2 and the lemma 2.3 get:

$$\begin{aligned}
\nu \|\mathbf{curl} \mathbf{w}\|_{0, \Omega_F} &\leq C \left(\sum_{\kappa \in \tau_h^F} (h_\kappa^2 \|\mathbf{f} - \mathbf{f}_h\|_{0, \kappa}^2 + h_\kappa^2 \|\mathbf{f}_h - \nabla p_h\|_{0, \kappa}^2)^{1/2} + \frac{\nu}{2} \sum_{\kappa \in \tau_h^F} \sum_{e \in \varepsilon_\kappa} h_\kappa \|[\mathbf{curl} \mathbf{u}_h \times \mathbf{n}]\|_{0, e}^2 \right)^{1/2} \\
&\leq C \left(\sum_{\kappa \in \tau_h^F} h_\kappa^2 \|\mathbf{f} - \mathbf{f}_h\|_{0, \kappa}^2 + \sum_{\kappa \in \tau_h^F} \gamma_{\kappa, F}^2 \right)^{1/2}, \quad (5.12)
\end{aligned}$$

and we obtain the upper bound. For the lower bound, we choose in the equation (5.10), $\mathbf{v}_h = 0$ and we take for an element $\kappa \in \Omega_F$, $\mathbf{v} = (\mathbf{f}_h - \nabla p_h) \psi_\kappa$ and remark that

$$\int_{\kappa} \mathbf{curl} \mathbf{u}_h \mathbf{curl} \mathbf{v} = \int_{\partial \kappa} (\mathbf{curl} \mathbf{u}_h \times \mathbf{n}) \mathbf{v} = 0$$

to obtain using the inverse inequality $\|\mathbf{curl} \mathbf{v}\|_{0, \kappa} \leq h_\kappa^{-1} \|\mathbf{v}\|_{0, \kappa}$:

$$\|\mathbf{f}_h - \nabla p_h\|_{0, \kappa}^2 \leq \|\mathbf{f} - \mathbf{f}_h\|_{0, \kappa}^2 + |p - p_h|_{1, \kappa}^2 + h_\kappa^{-2} \|\mathbf{curl} \mathbf{w}\|_{0, \kappa}^2$$

Next, we take $\mathbf{v} = \mathcal{L}_e([\mathbf{curl} \mathbf{u}_h \times \mathbf{n}] \psi_e)$ and integrate by part the last term of the equation (5.10) to obtain

$$\nu \|[\mathbf{curl} \mathbf{u}_h \times \mathbf{n}] \psi_e^{1/2}\|_{0, e} \leq c (h_e^{-1/2} \|\mathbf{curl} \mathbf{w}\|_{0, \Delta_e^F} + h_e^{1/2} |p - p_h|_{1, \Delta_e^F} + h_e^{1/2} \|\mathbf{f}_h - \nabla p_h\|_{0, \Delta_e^F} + h_e^{1/2} \|\mathbf{f} - \mathbf{f}_h\|_{0, \Delta_e^F})$$

and we deduce the lower bound. \square

To show an upper and a lower bound of the pressure, we define the indicators:

$$\eta_\kappa = \begin{cases} \eta_\kappa^F = \sum_{e \in \varepsilon_\kappa} h_e^{1/2} \|[(\mathbf{f}_h - \nabla p_h) \cdot \mathbf{n}]\|_{L^2(e)} & \text{if } \kappa \in \Omega_F \\ \eta_\kappa^P = \sum_{e \in \varepsilon_\kappa} h_e^{1/2} \|[(\mathbf{f}_h - \nabla p_h - \mu \mathbf{u}_h) \cdot \mathbf{n}]\|_{L^2(e)} & \text{if } \kappa \in \Omega_P \end{cases} \quad (5.13)$$

where $[\cdot]_e$ is the jump through the face e of κ . If the face e is on Γ or $\partial \Omega$, $[\cdot]_e$ will be the trace on e from the domain (Ω_P or Ω_F) containing κ .

Theorem 5.4. *The following bounds hold*

$$\begin{aligned} |p - p_h|_{1,\Omega} &\leq C \left\{ \mu \|\mathbf{curl} \mathbf{w}\|_{0,\Omega_F}^2 + \mu \|\nabla \lambda\|_{0,\Omega_F}^2 \right. \\ &\quad \left. + \sum_{\kappa \in \tau_h} \left(\eta_\kappa^2 + h_\kappa^2 \|\operatorname{div} \mathbf{f}\|_{L^2(\kappa)}^2 + \left(\sum_{e \in \varepsilon_\kappa} h_e^{1/2} \|[(\mathbf{f} - \mathbf{f}_h) \cdot \mathbf{n}]\|_{L^2(e)} \right)^2 \right) \right\}^{1/2} \end{aligned} \quad (5.14)$$

and

$$\begin{cases} \eta_\kappa^F \leq C \left(\|\mathbf{w}\|_{0,\Delta_\kappa^F} + |p - p_h|_{H^1(\Delta_\kappa^F)} + h_e \|\operatorname{div} \mathbf{f}\|_{L^2(\Delta_\kappa^F)} + \sum_{e \in \varepsilon_\kappa} h_e^{1/2} \|[(\mathbf{f} - \mathbf{f}_h) \cdot \mathbf{n}]\|_{L^2(e)} \right) \\ \eta_\kappa^P \leq C \left(\|\nabla \lambda\|_{0,\Delta_\kappa^P} + |p - p_h|_{H^1(\Delta_\kappa^P)} + h_e \|\operatorname{div} \mathbf{f}\|_{L^2(\Delta_\kappa^P)} + \sum_{e \in \varepsilon_\kappa} h_e^{1/2} \|[(\mathbf{f} - \mathbf{f}_h) \cdot \mathbf{n}]\|_{L^2(e)} \right) \end{cases} \quad (5.15)$$

Proof : The error function $p - p_h$ belongs to $H^1(\Omega)$ and satisfies, using (3.6), for all $q \in H^1(\Omega)$ and $q_h \in Q_h$:

$$\begin{aligned} \mu(\mathbf{u} - \mathbf{u}_h, \nabla q)_{\Omega_P} + (\nabla(p - p_h), \nabla q) &= \mu(\mathbf{u} - \mathbf{u}_h, \nabla(q - q_h))_{\Omega_P} + (\nabla(p - p_h), \nabla(q - q_h)) \\ &= \int_\Omega \mathbf{f} \nabla(q - q_h) - \mu \int_{\Omega_P} \mathbf{u}_h \nabla(q - q_h) - \int_\Omega \nabla p_h \nabla(q - q_h) \end{aligned} \quad (5.16)$$

We replace $\mathbf{u} - \mathbf{u}_h$ by $\nabla \lambda + \mathbf{w}$ and we obtain for all $q \in H^1(\Omega)$ and $q_h \in Q_h$:

$$\begin{aligned} (\nabla(p - p_h), \nabla q) &= \int_\Omega (\mathbf{f} - \nabla p_h) \nabla(q - q_h) - \mu \int_{\Omega_P} \mathbf{u}_h \nabla(q - q_h) - \mu \int_{\Omega_P} \nabla \lambda \nabla q - \mu \int_{\Omega_P} \mathbf{w} \nabla q \\ &= \int_\Omega (\mathbf{f} - \mathbf{f}_h) \nabla(q - q_h) - \mu \int_{\Omega_P} \nabla \lambda \nabla q + \mu \int_{\Omega_F} \mathbf{w} \nabla q \\ &\quad + \int_{\Omega_P} (\mathbf{f}_h - \nabla p_h + \mu \mathbf{u}_h) \nabla(q - q_h) + \int_{\Omega_F} (\mathbf{f}_h - \nabla p_h) \nabla(q - q_h) \end{aligned} \quad (5.17)$$

Integrating by part and remarking that $\operatorname{div} \mathbf{f}_h = \operatorname{div} \mathbf{u}_h = \operatorname{div} \nabla p_h = 0$ in every element $\kappa \in \tau_h$, we obtain:

$$\begin{aligned} (\nabla(p - p_h), \nabla q) &= -\mu \int_{\Omega_P} \nabla \lambda \nabla q + \mu \int_{\Omega_F} \mathbf{w} \nabla q + \sum_{\kappa \in \tau_h} \left\{ - \int_\kappa \operatorname{div} \mathbf{f} (q - q_h) + \frac{1}{2} \sum_{e \in \varepsilon_\kappa} \int_e [(\mathbf{f} - \mathbf{f}_h) \cdot \mathbf{n}] (q - q_h) \right\} \\ &\quad - \frac{1}{2} \sum_{\kappa \in \tau_h^F} \left\{ \sum_{e \in \varepsilon_\kappa} \int_e [(\mathbf{f}_h - \nabla p_h) \cdot \mathbf{n}] (q - q_h) \right\} - \frac{1}{2} \sum_{\kappa \in \tau_h^P} \left\{ \sum_{e \in \varepsilon_\kappa} \int_e [(\mathbf{f}_h - \nabla p_h + \mu \mathbf{u}_h) \cdot \mathbf{n}] (q - q_h) \right\} \end{aligned} \quad (5.18)$$

We take $q = p - p_h$ and $q_h = I_h q$, we use the relation $a \cdot b \leq \frac{1}{2\varepsilon} a^2 + \frac{\varepsilon}{2} b^2$ with a suitable choice of ε , the properties of I_h and the theorem (2.2) to obtain the inequality (5.14).

To prove the inequality (5.15), we take $q_h = 0$, $q = \mathcal{L}_e^F([(\mathbf{f}_h - \nabla p_h) \cdot \mathbf{n}] \psi_e)$ (resp. $q = \mathcal{L}_e^P([(\mathbf{f}_h - \nabla p_h) \cdot \mathbf{n}] \psi_e)$) in the equation (5.18) and we use the property (5.1) to obtain the first inequality of (5.15) (resp. the second inequality). \square

To complete the upper and lower bounds of \mathbf{w} , we show bound on \mathbf{w} in Ω_P . We introduce the indicators

$$\gamma_{\kappa,P} = \|\mathbf{f}_h - \nabla p_h + \mu \mathbf{u}_h\|_{0,\kappa} \quad \text{if } \kappa \in \Omega_P \quad (5.19)$$

Theorem 5.5. *The following bounds hold:*

$$\mu \|\mathbf{w}\|_{0,\Omega_P} \leq \left(\sum_{\kappa \in \tau_h^P} \|\mathbf{f} - \mathbf{f}_h\|_{0,\kappa}^2 \right)^{1/2} + \left(\sum_{\kappa \in \tau_h^F} \gamma_{\kappa,P}^2 \right)^{1/2} + \mu \|\nabla \lambda\|_{L^2(\Omega_P)} + \|\nabla(p - p_h)\|_{L^2(\Omega_P)} \quad (5.20)$$

and

$$\gamma_{\kappa,P} \leq C \left(\|\mathbf{f} - \mathbf{f}_h\|_{0,\kappa}^2 + \|\mathbf{w}\|_{0,\kappa}^2 + |p - p_h|_{1,\kappa}^2 + |\lambda|_{1,\kappa}^2 \right)^{1/2} \quad (5.21)$$

Proof : In the equation (5.10), we replace $\mathbf{u} - \mathbf{u}_h = \mathbf{w} + \nabla\lambda$ and we take $\mathbf{v} = \mathbf{w}$ in Ω_P and $\mathbf{0}$ in Ω_F . We choose $\mathbf{v}_h = 0$ and we obtain:

$$\mu \|\mathbf{w}\|_{0,\Omega_P}^2 = \int_{\Omega_P} (\mathbf{f} - \mathbf{f}_h)\mathbf{w} + \int_{\Omega_P} (\mathbf{f}_h - \nabla p_h + \mu\mathbf{u}_h)\mathbf{w} - \mu \int_{\Omega_P} \mathbf{w}\nabla\lambda - \int_{\Omega_P} \nabla(p - p_h)\mathbf{w} \quad (5.22)$$

Then we get the upper bound. For the lower bound, we choose $\mathbf{v}_h = 0$ in the equation (5.10) and we take for an element $\kappa \in \Omega_P$, $\mathbf{v} = (\mathbf{f}_h - \nabla p_h + \mu\mathbf{u}_h)\psi_\kappa$ to obtain the lower bound. \square

To simplify the notations, we define the indicators:

$$\gamma_\kappa = \begin{cases} \gamma_{\kappa,F} & \text{if } \kappa \in \Omega_F \\ \gamma_{\kappa,P} & \text{if } \kappa \in \Omega_P \end{cases}$$

Corollary 5.6. *Let Ω_F be convex. The optimal a posteriori estimate holds*

$$\begin{aligned} \|\mathbf{u} - \mathbf{u}_h\|_X + |p - p_h|_{1,\Omega} \leq & \left\{ \sum_{\kappa \in \mathcal{T}_h} (\gamma_\kappa^2 + \xi_\kappa^2 + \eta_\kappa^2 + \|\mathbf{f} - \mathbf{f}_h\|_{0,\kappa}^2 + h_\kappa^2 \|\operatorname{div} \mathbf{f}\|_{L^2(\kappa)}^2) \right. \\ & \left. + \left(\sum_{e \in \mathcal{E}_\kappa} h_e^{1/2} \|[(\mathbf{f} - \mathbf{f}_h) \cdot \mathbf{n}]\|_{L^2(e)} \right)^2 \right\}^{1/2} \end{aligned} \quad (5.23)$$

where η_κ , ξ_κ and γ_κ are given by the formulas (5.5), (5.9), (5.15) and (5.21).

Conclusion: We observe that estimate (5.23) is optimal: up to the terms involving the data, the full error is bounded by a constant times the sum of all indicators. Estimates (5.5), (5.9), (5.15) and (5.21) are local, i.e., only involve the error in a neighborhood of K or e . The indicators η_κ , ξ_κ and γ_κ can be viewed as a measure for the error of the space discretization and can be used to adapt the mesh-size in space.

REFERENCES

- [1] H. ABBOUD, F. EL CHAMI & T. SAYAH, A priori and a posteriori estimates for three dimensional Stokes equations with non standard boundary conditions, *submitted*.
- [2] C. BERNARDI, F. HECHT, O. PIRONEAU Coupling Darcy and Stokes Equations for the porous media with cracks, *ESAIM: M2AN*, Vol. 39, No 1, 2005, pp. 7-35.
- [3] C. BERNARDI, F. HECHT & Z. MGHAZLI Mortar finite element discretization for the flow in a non homogeneous porous medium, *Comput. Methods Appl. Mech. Engrg.*, Vol. 196, 2007, pp. 1554-1573.
- [4] C. BERNARDI, F. HECHT & FATMA ZOHRA NOURI A new finite element discretization of the Stokes problem coupled with Darcy equations, *IMA Journal of Numerical Analysis*, Vol. 30, 2010, pp. 61-93.
- [5] P. CLÉMENT Approximation by finite element functions using local regularisation., *R.A.I.R.O. Anal. Numer.* 9, pp, 77-84 (1975)
- [6] F. DUBOIS. Vorticity-velocity-pressure formulation for the Stokes problem, *Math. Meth. in the Applied Sciences*, **25** (2002), 1091-1119.
- [7] F. DUBOIS, M. SALAÜN & S. SALMON. Vorticity-velocity-pressure and stream function-vorticity formulations for the Stokes problem, *J. Math. Pures Appl.*, **82** (2003), 1395-1451.
- [8] F. EL CHAMI & T. SAYAH, A posteriori error estimators for the fully discrete time dependent Stokes problem with some different boundary conditions, *CALCOLO*, 47 (3) 2010, 169-192.
- [9] J. C. NEDELEC Mixed finite element in \mathbb{R}^3 . *Numer. Math.* 35, pp. 315-341 (1980).
- [10] J. C. NEDELEC Eléments finis mixtes incompressibles pour l'équation de Stokes dans \mathbb{R}^3 . *Numer. Math.* 39, pp. 97-112 (1982).
- [11] P.-A. RAVIART & J.-M. THOMAS A mixed finite element method for second order elliptic problems, *Mathematical Aspects of Finite Element Methods*. Springer, Berlin. Lect. Notes Math., 606 pp 292-315 (1977).
- [12] S. SALMON Développement numérique de la formulation tourbillon-vitesse-pression pour le problème de Stokes, *Thesis, Université Pierre et Marie Curie, Paris* (1999).