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INFLUENCE OF THE NONLINEAR DYNAMIC BEHAVIOR OF JOURNAL BEARINGS ON GEAR-BEARING ASSEMBLIES

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ABSTRACT

Journal bearings cannot be considered as passive elements in gear-bearing assemblies, and the lubricant is recognized as playing an important role in the interactions between the shafts and the bearings. In order to take this influence into account, bearings are usually modeled by means of eight dynamic coefficients, i.e., asymmetric stiffness and damping matrices.

In this paper, a nonlinear approach is proposed enabling the behavior of a gear-shaft-bearing assembly to be analyzed. A discrete finite element model is used for the shafts, and a specific gear element is introduced which accounts for non-linear time-varying mesh stiffness as well as tooth shape deviations. The meshing forces are internal system forces whereas the effects of the bearings on the shafts are taken to be external. A combination of the Newmark time integration scheme and the Newton-Raphson algorithm is used to simultaneously solve the contact problem for the gear, and the Reynolds equations for the bearings. The resulting algorithm is applied to a single stage geared system with two shafts, four bearings, a pinion and a gear while taking mass unbalance, eccentricity and meshing excitations into account. Several examples are presented which demonstrate the influence of bearing nonlinearity and the efficiency of the proposed model and numerical procedure.

INTRODUCTION

The dynamic response of geared systems remains a major concern in terms of noise and vibration, especially in high-speed applications. In this context, journal bearings are recognized as interesting technological solutions since they provide significant stiffness and damping for reduced noise levels compared with rolling element bearings. The literature on journal bearings comprises numerous contributions, only a few of which can be cited here [1-7] and it is now accepted

that, for most practical cases, the theoretical foundations are firmly established allowing accurate predictions of bearing performance. A survey about the theory of hydrodynamics and its applications to various bearing geometries and operating conditions can be found in the classic textbooks of Cameron [8], Pinkus and Sternlicht [9], Frêne *et al.* [10].

On the other hand, the simulation of gear dynamics has yielded a vast literature [11], [12] with the majority of the gear models based upon lumped parameter representations combining rigid gears, discrete elastic and dissipative elements [13], [14], [15] or, more recently, on combinations of shaft finite elements and elastic foundations [16]. Forcing terms are often derived from transmission errors (Gregory *et al.* [15], Munro [17], Özgüven and Houser [18]) and /or the contact conditions between the active flanks [19]. Non linear features in spur gearing such as backlash and shocks on tooth flanks have been analyzed by Kahraman and Singh [20], Kahraman and Blankenship [21], Parker *et al.* [22].

Over the last decade, computing procedures have been developed which include the characteristics of shafts, gears, bearings and housings in the prediction of the global vibrational behavior of a transmission [23-26]. The majority of these codes, however, are restricted to linear analysis, i.e., static solution and modal analysis, and bearing and gear models remain global. The present work contributes to the analysis of the interactions between gears, shafts and journal bearings in geared drives. In contrast to the majority of the models in the literature, the time-varying properties and non-linearities of bearings and gears are introduced in the simulations. To this end, the equations of motion, the normal contact problems on tooth flanks and the Reynolds equation for hydrodynamic bearings are simultaneously solved by combining a time-step integration scheme, a Newton-Raphson procedure and a normal contact algorithm. Finally, the importance of gear-bearing

dynamic interactions and the interest of non-linear models for journal bearings are discussed.

NOMENCLATURE

- $(\vec{s}, \vec{t}, \vec{z})$: frame of reference (\vec{s} is along the horizontal center line, \vec{z} is in the axial direction, see figure 1)
- $(\vec{x}, \vec{y}, \vec{z})$: cartesian local frame of the bearing
- $(\vec{r}, \vec{q}, \vec{z})$: polar local frame of the bearing
- ω : rotation speed (ras/s)
- C : radial clearance
- ϵ : eccentricity ratio (varies between 0 and 1)
- ϕ : attitude angle (rad)

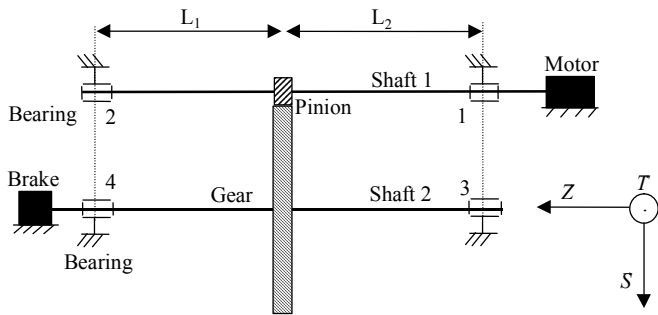


Figure 1 - Sketch of the geared system

1 - DESCRIPTION OF THE MODEL

1-1 Shaft-gear model

The dynamic model of single stage transmissions (figure 1) is based on the developments of Vexex and Maatar [19] and Ajmi and Vexex [16]. Shafts are modeled by classic two-node finite elements with 6 degrees of freedom per node for bending (Timoschenko's beam), torsion and axial displacements. A pinion - gear pair are represented by two 2-node shaft elements connected by lumped stiffnesses which account for contact, tooth and gear body deflections along the contact lines in the base plane (figure 2). An elemental mesh stiffness and an equivalent normal shape deviation are associated with every potential point of contact to simulate actual tooth flank traces and their evolutions with time. Based upon rigid-body kinematics, the lines of contact are translated and all relevant parameters (stiffness and deviations) are re-calculated at each time step of the meshing process. It is assumed that the contacts between mating flanks are line contacts within the theoretical base plane and that tooth friction forces can be neglected compared with normal forces. The normal contact condition at every potential point of contact on the tooth flanks is directly included in the mesh stiffness formulation by canceling the elemental stiffness when the corresponding deflection is negative or nil. The proposed formulation leads to time-varying

non-linear mesh stiffness matrices and forcing terms, the latter accounting for tooth shape deviations and mounting errors [16], [19], [23]. The pinion-gear mass matrix is classic since second order terms and gyroscopic components are neglected.

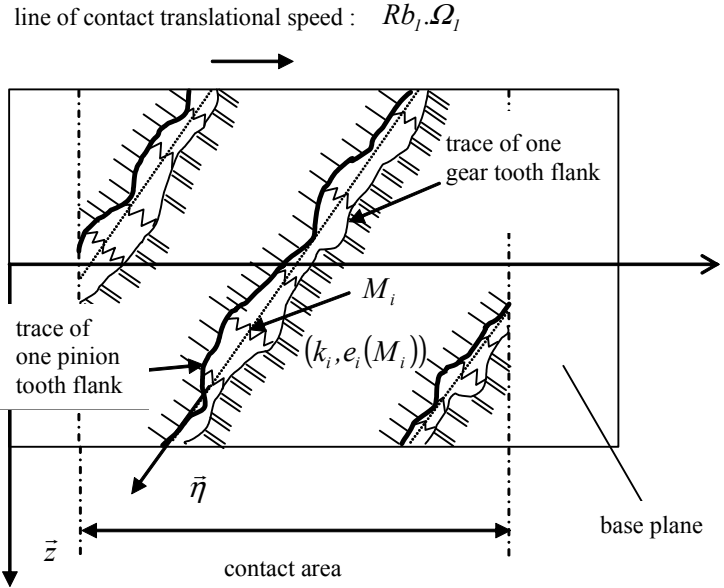


Figure 2 – Model of mesh elasticity and tooth shape deviations

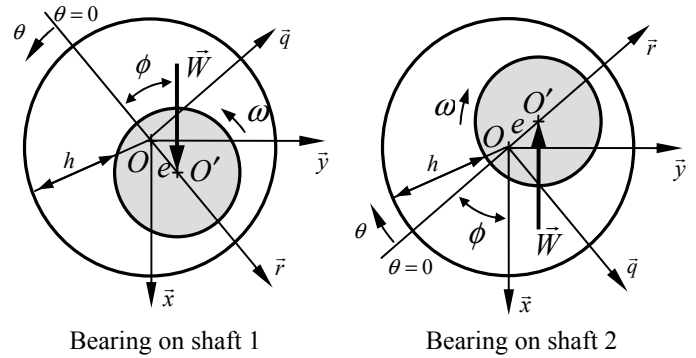


Figure 3 - Direction of external load and definition of local frames for the bearings on shaft 1 and 2

1-2 Journal bearing model

1-2-1 Reynolds equation – Short Bearing approximation

Figure 3 gives a schematic view of a journal bearing consisting of a fixed journal of radius R , length L and of a rotating shaft of radius R_s , with $C = R - R_s$ being the radial clearance. At constant speed ω and under a constant load \bar{W} , the shaft center O' keeps a fixed position in the bearing

defined by the constant eccentricity e_0 and attitude angle ϕ_0 whereas, for dynamic conditions, it follows an orbit within the journal clearance characterized by a time-dependent eccentricity e and attitude angle ϕ . The position of the shaft center with respect to the bearing center O can be described either by its Cartesian co-ordinates in the frame (\bar{x}, \bar{y})

$$\overline{OO'} = x\bar{x} + y\bar{y} = e\cos\phi\bar{x} + e\sin\phi\bar{y} \quad (1)$$

or its local (or polar) co-ordinates in the frame (\bar{r}, \bar{q})

$$\overline{OO'} = e\bar{r} \quad \phi = (\bar{x}, \bar{r}) \quad (2)$$

Introducing the angular co-ordinate θ whose origin is on the axis \bar{r} , the lubricant film thickness is given as:

$$h = C(1 + \varepsilon\cos\theta) \quad (3)$$

where the eccentricity ratio $\varepsilon = e/C$ varies between 0 and 1.

Assuming an isothermal, laminar flow of an iso-viscous incompressible fluid, the behavior of the lubricant film is governed by the Reynolds equation [2]. For the sake of simplicity, the short bearing approximation is introduced, i.e., the width-to-diameter ratio L/D is supposed to be small, which leads to the following Reynolds equation for dynamic conditions:

$$\frac{\partial}{\partial z} \left[\frac{h^3}{\mu} \frac{\partial p}{\partial z} \right] = 6 \left[(\omega - 2\dot{\phi}) \frac{dh}{d\theta} + 2C\varepsilon\cos\theta \right] \quad (4)$$

Using the boundary conditions $p(\theta, z = \pm L/2) = 0$, Equation (4) can be integrated and the hydrodynamic pressure field is expressed as:

$$p(\theta, z) = -\frac{3\mu}{C^2} \left(z^2 - \frac{L^2}{4} \right) \frac{(\omega - 2\dot{\phi})\varepsilon\sin\theta - 2\dot{\varepsilon}\cos\theta}{(1 + \varepsilon\cos\theta)^3} \quad (5)$$

The hydrodynamic force F_b is obtained by integration of $p(\theta, z)$ for $\theta \in [0, 2\pi]$ and $z \in [-L/2, L/2]$ using the Gumbel boundary conditions which discard negative pressures from the force integral [10].

1-2-1 Steady-state solution

The solution for steady-state conditions can be deduced from (5) by setting the time-varying perturbations $\dot{\varepsilon}$ and $\dot{\phi}$ to zero. The associated hydrodynamic bearing force integration is straightforward and gives

$$F_b(\varepsilon_0) = \mu LR \omega \left(\frac{L}{D} \right)^2 \left(\frac{R}{C} \right)^2 \frac{\varepsilon_0}{(1 - \varepsilon_0^2)^2} \sqrt{16\varepsilon_0^2 + \pi^2(1 - \varepsilon_0^2)} \quad (6)$$

with the attitude angle being defined as

$$\tan\phi_0 = \frac{\pi}{4} \frac{\sqrt{1 - \varepsilon_0^2}}{\varepsilon_0} \quad (7)$$

The steady-state position and bearing force depend on ε_0 . This is derived from the static equilibrium between the hydrodynamic reaction generated by the bearing and the external force W_0 imposed by the transmitted torques and transferred to the bearings by the pinion and the gear, i.e.,

$$F_b(\varepsilon_0) = W_0 \quad (8)$$

This equation is non-linear and it is solved by a dichotomy scheme. In turn, the eccentricity of the shaft in the journal modifies the pinion-gear center distance slightly and consequently, the gear mesh stiffness function and excitations.

1-2-1 Dynamic characteristics

Denoting (x_0, y_0) , the shaft center co-ordinates in steady-state conditions and $(\Delta x, \Delta y)$, the corresponding time-dependent dynamic perturbations, the actual position of the shaft center is defined as:

$$\overline{OO'} = (x_0 + \Delta x)\bar{x} + (y_0 + \Delta y)\bar{y} = (C\varepsilon_0\cos\phi_0 + \Delta x)\bar{x} + (C\varepsilon_0\sin\phi_0 + \Delta y)\bar{y} \quad (9)$$

From which, the dynamic parameters ε , ϕ , and their time-derivatives $\dot{\varepsilon}$ and $\dot{\phi}$ are deduced as:

$$\begin{aligned} \varepsilon &= \frac{\sqrt{(C\varepsilon_0\cos\phi_0 + \Delta x)^2 + (C\varepsilon_0\sin\phi_0 + \Delta y)^2}}{C} \\ \cos\phi &= \frac{C\varepsilon_0\cos\phi_0 + \Delta x}{\sqrt{(C\varepsilon_0\cos\phi_0 + \Delta x)^2 + (C\varepsilon_0\sin\phi_0 + \Delta y)^2}} \\ \dot{\varepsilon} &= \frac{\partial\varepsilon}{\partial t} = \frac{(C\varepsilon_0\cos\phi_0 + \Delta x)\Delta\dot{x} + (C\varepsilon_0\sin\phi_0 + \Delta y)\Delta\dot{y}}{C\sqrt{(C\varepsilon_0\cos\phi_0 + \Delta x)^2 + (C\varepsilon_0\sin\phi_0 + \Delta y)^2}} \\ \dot{\phi} &= \frac{\partial\phi}{\partial t} = \frac{-\Delta\dot{x}\varepsilon + (C\varepsilon_0\cos\phi_0 + \Delta x)\dot{\varepsilon}}{\varepsilon(C\varepsilon_0\sin\phi_0 + \Delta y)} \end{aligned} \quad (10)$$

2 – DYNAMIC BEHAVIOR OF THE GEAR-BEARING ASSEMBLY

2-1 Equations of motion

The assembly of all the elemental matrices and forcing term vectors leads to the following system of equations in the frame $(\bar{s}, \bar{t}, \bar{z})$ relative to the structure

$$[M]\{\ddot{X}\} + [C]\{\dot{X}\} + [K(t, X)]\{X\} = \{F(t, X)\} + \{F_b(X, \dot{X})\} \quad (11)$$

where

M and C are the global constant mass and damping matrices
 X, \dot{X}, \ddot{X} are the degree-of-freedom, velocity and acceleration
vectors

$K(t, X)$ is the global stiffness matrix of the system,
 $K(t, X) = K_{shaft} + K_{gear}(t, X)$ with

K_{shaft} : constant global stiffness matrix of the shafts

$K_{gear}(t, X)$: global stiffness matrix of the gear which is
time-dependant because of the contact length evolution on
the base plane and potentially non-linear because of the
interaction between the degrees of freedom, tooth
deviations and the instantaneous contact extent (Velex and
Maatar, [19])

$F(t, X) = F_0(t) + F_{gear}(t, X)$ with

$F_0(t)$: external load vector (input and output torques, mass
unbalance, ...)

$F_{gear}(t, X)$: embodies the inertial effects caused by unsteady
rotations (associated with no-load transmission error) and the
contributions of geometrical errors and tooth shape deviations
(pitch errors, tooth modifications ...) [19]

$F_b(X, \dot{X})$ are the hydrodynamic forces produced by the
bearings which are opposing the motions of the shaft as
described in section 2.

If the inertia center of the shaft is different from its
geometrical center, then mass imbalance forces are embedded
in the external load vector $F_0(t)$. In the frame $(\vec{s}, \vec{t}, \vec{z})$, these
forces read

$$F_{mi} = \begin{Bmatrix} F_{miX} \\ F_{miY} \\ 0 \end{Bmatrix} = \begin{Bmatrix} Me_{mi}\omega^2 \cos(\omega t) \\ Me_{mi}\omega^2 \sin(\omega t) \\ 0 \end{Bmatrix} \quad (12)$$

for each node of the shaft. In this expression, M is the
mass of the shaft plus the pinion or gear and e_{mi} represents the
distance between the inertia and geometrical centers of the
shaft.

In the proposed formulation (11), the contributions of
bearings appear as external forces to the system following the
original approach of Abdul-Wahed [27] in his study of the non-
linear behavior of rigid shafts supported by hydrodynamic
bearings. The system of equations (11) is non-linear because
the hydrodynamic forces $F_b(X, \dot{X})$ depend on the position

X and on the velocity \dot{X} , and because of the contact
conditions between the gear teeth. The resolution requires a
complex procedure combining several iterative schemes and a
time-step integration process as described in Section 3 -3,
which will be referred to as non-linear analysis. The solution

technique can be significantly simplified by assuming small
displacements in the vicinity of the static positions, i.e., using
the linear analysis which is probably no longer suited to the
case of significant dynamic loads or mass imbalances. Both
kinds of analyses, linear and non-linear, are detailed in the
following section.

2-2 Linear analysis

Using a first order expansion in the vicinity of the static
solution X_0 leads to the following dynamic bearing force
vectors:

$$F_b(X, \dot{X}) = \{F_{b0}\} - [K]_{XY} \{\Delta X\} - [C]_{XY} \{\Delta \dot{X}\} \quad (13)$$

where F_{b0} is obtained by (6), $\Delta X = X - X_0$, $[K]_{XY}$
and $[C]_{XY}$ represent the stiffness and damping matrices due to
the lubricant film.

For short bearings, and applying Gumbel's conditions,
the expressions of $[K]$ and $[C]$ are easier to obtain in the
local frame (\vec{r}, \vec{q}) of the bearing. According to Frene *et al.*
[10], they read

$$\begin{aligned} [K]_{rq} &= \frac{W_0}{C} \begin{bmatrix} A_{rr} & A_{rq} \\ A_{qr} & A_{qq} \end{bmatrix} & [C]_{rq} &= \frac{W_0}{C\omega} \begin{bmatrix} B_{rr} & B_{rq} \\ B_{qr} & B_{qq} \end{bmatrix} \\ A_{rr} &= \frac{8(1+\varepsilon_0^2)}{1-\varepsilon_0^2} f(\varepsilon_0) & A_{rq} &= \frac{\pi}{\varepsilon_0} \sqrt{1-\varepsilon_0^2} f(\varepsilon_0) \\ A_{qr} &= \frac{-\pi(1+2\varepsilon_0^2)}{\varepsilon_0 \sqrt{1-\varepsilon_0^2}} f(\varepsilon_0) & A_{qq} &= 4f(\varepsilon_0) \\ B_{rr} &= -2A_{qr} & B_{qq} &= 2A_{rq} & B_{rq} &= B_{qr} = -2A_{qq} \\ f(\varepsilon_0) &= \frac{1}{\sqrt{16\varepsilon_0^2 + \pi^2(1-\varepsilon_0^2)}} \end{aligned} \quad (14)$$

where W_0 is the static load generated by the meshing and
applied to the bearing. The coefficients A_{ij} and B_{ij} depend on
 ε_0 only, which characterizes the steady-state position, such
that $[K] = [K(\varepsilon_0)]$ and $[C] = [C(\varepsilon_0)]$.

The matrices $[K]_{rq}$ and $[C]_{rq}$ can be transferred
from the local frame (\vec{r}, \vec{q}) into the Cartesian frame (\vec{x}, \vec{y})
by

$$\begin{aligned} [K(\varepsilon_0)]_{xy} &= [P][K(\varepsilon_0)]_{rq}[P]^{-1} \\ [C(\varepsilon_0)]_{xy} &= [P][C(\varepsilon_0)]_{rq}[P]^{-1} \end{aligned} \quad (15)$$

where

$$[P] = \begin{bmatrix} \cos \phi_0 & -\sin \phi_0 \\ \sin \phi_0 & \cos \phi_0 \end{bmatrix} \quad (16)$$

for bearings on the input (pinion) shaft and

$$[P] = \begin{bmatrix} -\cos \phi_0 & \sin \phi_0 \\ \sin \phi_0 & \cos \phi_0 \end{bmatrix} \quad (17)$$

for bearings on the output (gear) shaft, with the attitude angle ϕ_0 obtained by (7).

Finally, the matrices $[K]_{XY}$ and $[C]_{XY}$ in Equation (13) are obtained by applying a second change of basis, from $(\vec{x}, \vec{y}, \vec{z})$ to $(\vec{s}, \vec{t}, \vec{z})$

$$\begin{aligned} [K(\varepsilon_0)]_{XY} &= [Q][K(\varepsilon_0)]_{xy}[Q]^{-1} \\ [C(\varepsilon_0)]_{XY} &= [Q][C(\varepsilon_0)]_{xy}[Q]^{-1} \end{aligned} \quad (18)$$

with

$$[Q] = \begin{bmatrix} -\sin \alpha & \cos \alpha \\ -\cos \alpha & -\sin \alpha \end{bmatrix} \quad (19)$$

for all bearings, where α is the pressure angle.

Replacing the expression (13) of $F_b(X, \dot{X})$ into the equations of motion (11) along with the expressions of ΔX and $\Delta \dot{X}$, one obtains the linear system of equations in the frame $(\vec{s}, \vec{t}, \vec{z})$

$$\begin{aligned} [M]\{\ddot{X}\} + [C_0]\{\dot{X}\} + [K_0]\{X\} = \\ \{F(t, X)\} + \{F_{b_0}\} + [K(\varepsilon_0)]_{XY}\{X_0\} \end{aligned} \quad (20)$$

where

$$\begin{aligned} [C_0] &= [C] + [C(\varepsilon_0)]_{XY} \\ [K_0] &= [K]_{shaft} + [K]_{gear} + [K(\varepsilon_0)]_{XY} \end{aligned} \quad (21)$$

The differential system (20) is solved by the implicit Newmark time-step integration scheme combined with a normal contact algorithm which verifies that all contact forces on the teeth are positive and no contact deflections occur outside the contact area [19].

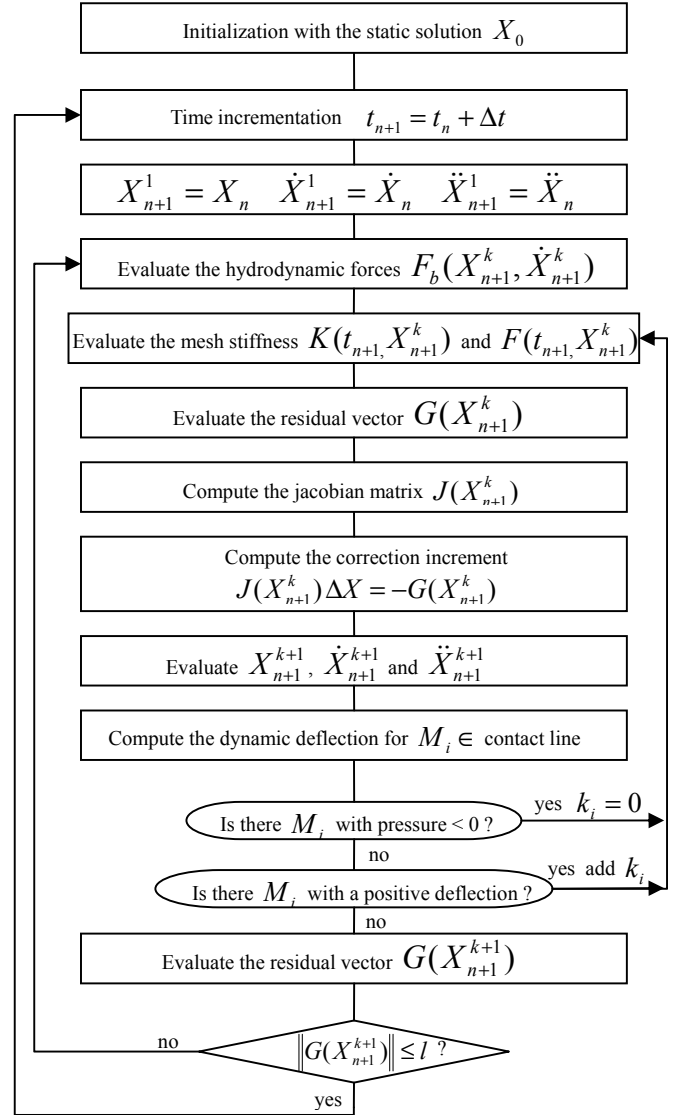


Table 1 – Iterative procedure

2-3 Non-linear analysis

Because of the assumption of small perturbations, the linear theory cannot account for large dynamic loads. Solving the complete non-linear problem requires the Reynolds equation in dynamic regime, the contact conditions on the teeth and the equations of motion to be dealt with simultaneously. In this case, the hydrodynamic forces $F_b(X, \dot{X})$ are obtained by numerical integration of the pressure field given by (5) as

$$F_b = \begin{cases} F_{br} \\ F_{bq} \end{cases} = \begin{cases} \iint p(\theta, z) \cos \theta R d\theta dz \\ \iint p(\theta, z) \sin \theta R d\theta dz \end{cases} \quad (22)$$

with $z \in [-L/2, L/2]$ and $\theta \in [0, 2\pi]$.

As for the linear approach, Gmbel's conditions are obtained by considering only the positive values of $p(\theta, z)$. The components of $F_b(X, \dot{X})$ in frames $(\vec{x}, \vec{y}, \vec{z})$ and $(\vec{s}, \vec{t}, \vec{z})$ are respectively given by

$$\begin{cases} F_{bx} \\ F_{by} \end{cases} = [P] \begin{cases} F_{br} \\ F_{bq} \end{cases} \quad (23)$$

$$\begin{cases} F_{bX} \\ F_{bY} \end{cases} = [Q] \begin{cases} F_{bx} \\ F_{by} \end{cases}$$

where $[P]$ and $[Q]$ are the matrices in (16), (17) and (19)

The non-linear equations (11) are solved by combining the Newmark time integration scheme with the Newton-Raphson incremental-iterative algorithm, and with an iterative process aimed at updating the dynamic characteristics of the meshing process (Table 1).

3 – NUMERICAL RESULTS

The gear, shaft and bearing data used in the numerical simulations are given in Tables 2 to 4. A unique modal damping factor of 0.03 has been used for all the numerical simulations and the time increment Δt was set to 1/32th of the mesh period. All simulations were performed over 256 mesh periods, i.e. 8192 time increments, which correspond approximately to 10 pinion revolutions.

3 – 1 Gear-bearing interactions

A first simulation was performed with an input (pinion) speed $\omega=700$ rad/s, an input torque $C=200$ N.m and the mass imbalance forces in (12) were computed with $e_{mi}=0.5$ mm. The corresponding transient motion of the shaft center in bearing 1 (input shaft) is plotted in Figure 4a. The motion, composed of 10 loops corresponding to the 10 pinion revolutions, becomes stable after the transient oscillations on the first orbits have disappeared. The last calculated orbit is shown in Figure 4b which shows the 26 oscillations associated with the mesh period and demonstrates the interaction between the gears and the bearings. The norm of the hydrodynamic forces over the last orbit, i.e. the last 256 mesh periods, of the same bearing is plotted in Figure 5. The low-frequency excitation seems prevalent but, as in the previous example, the mesh frequency components appear in the time signal superimposed on the basic once-per-revolution period related to mass imbalance.

	Pinion	Gear
Number of teeth	26	157
Face width (mm)	50	40
Helix angle (°)	0	
Module (mm)	4	
Pressure angle α (°)	20	
Dedendum coefficient	1.	1.
Addendum coefficient	1.4	1.4
Profile shift coefficient	0.16	-0.16
Center distance (mm)	366	
Mass (kg)	2	74

Table 2 - Gear data

	Pinion shaft	Gear shaft
Outer diameter (mm)	70	90
Inner diameter (mm)	30	30
Shaft length (mm)	$L_1 = L_2 = 160$	
Young modulus (N/m ²)	$2.1 \cdot 10^{11}$	
Density (kg/m ³)	7800	

Table 3 - Shaft data

	Bearings 1 and 2 Pinion shaft	Bearings 3 and 4 (gear shaft)
Diameter (mm)	70	90
Length (mm)	14	18
Radial clearance (μ m)	150	110

Table 4 - Bearing data

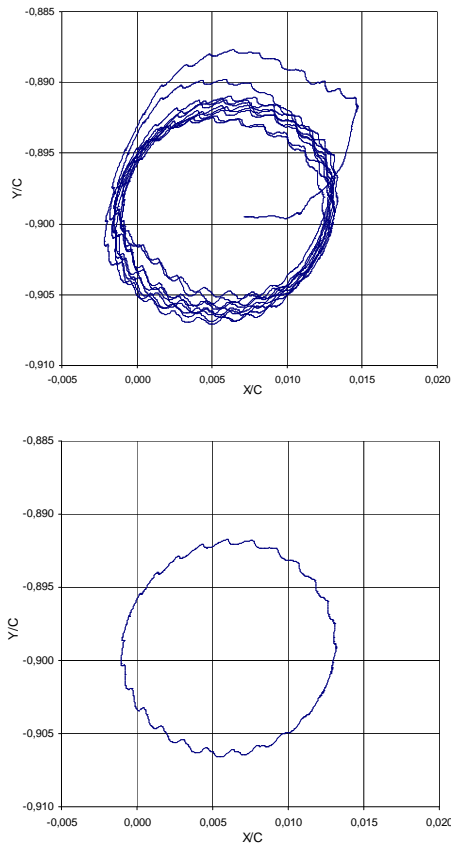


Figure 4 – a) Transient motion of the shaft center of bearing 1 and b) last calculated orbit

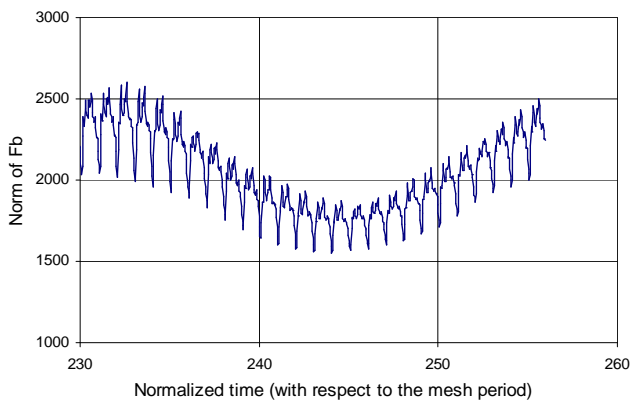


Figure 5 – Norm of the hydrodynamic forces F_b in bearing 1 for the last orbit

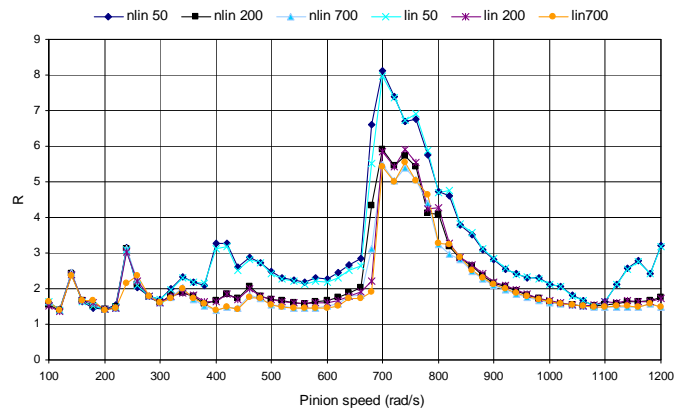


Figure 6 – Maximum dynamic-to-maximum static mesh load ratio for different input torques

3 – 2 Comparison of the linear and the non-linear theories

Dynamic tooth loading is characterized by considering the dynamic tooth load factor R defined as the ratio of the maximum dynamic mesh force to the maximum static one. Figure 6 represents the evolutions of R versus the pinion speed calculated for various loads and using different solution techniques. The solutions obtained by the linear and non-linear methods are very close and are dominated by amplitude jumps and shocks at the major tooth critical frequency. It can also be observed that dynamic amplifications are reduced when the transmitted torque is increased.

Considering bearing forces, the results on the pinion shaft (Figure 7) reveal that, here again, the linear and non-linear theories lead to very similar response curves with slightly larger differences as the transmitted load increases. The situation is different on the gear shaft, as illustrated in Figure 8, for which the non-linear and linear results deviate significantly as either speed or load increases. This discrepancy is apparently caused by the larger imbalance on the gear. This observation is confirmed by the bearing force plots in Figure 9, obtained for three imbalance amplitudes ($e_{mi}=0\text{mm}$, $e_{mi}=0.5\text{mm}$ and $e_{mi}=1\text{mm}$), which indicate that larger dynamic forces make non-linear modeling necessary.

On the other hand, dynamic tooth loads appear as moderately sensitive to imbalance, as illustrated in Figure 10, and, in contrast to bearing forces, the linear and non-linear results remain very close whatever the imbalance amplitude. The final set of results in Figures 11 and 12, shows, for two imbalance amplitudes ($e_{mi}=0.5\text{mm}$ and $e_{mi}=1\text{mm}$), the trajectories of the pinion shaft in its bearings for a pinion speed ranging from 100 to 1200 rad/s. The previous conclusions about the role of dynamic force amplitudes on shaft-bearing interactions are confirmed and suggest that linear theories might not be suited for accurate vibration transfer analysis.

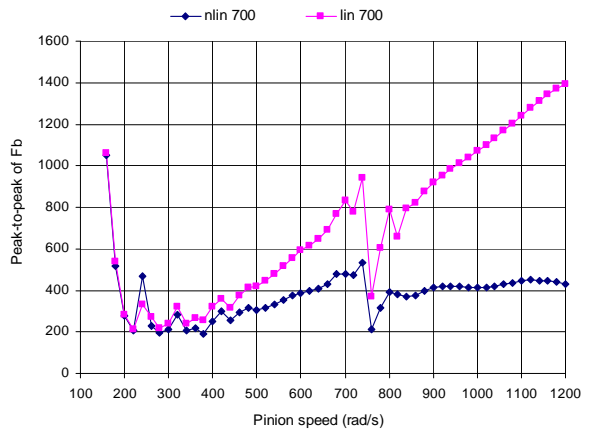
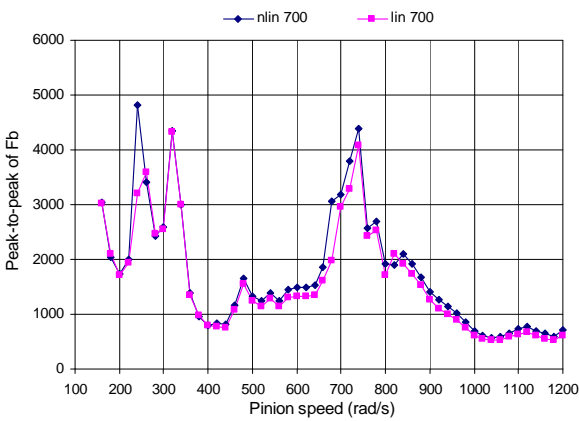
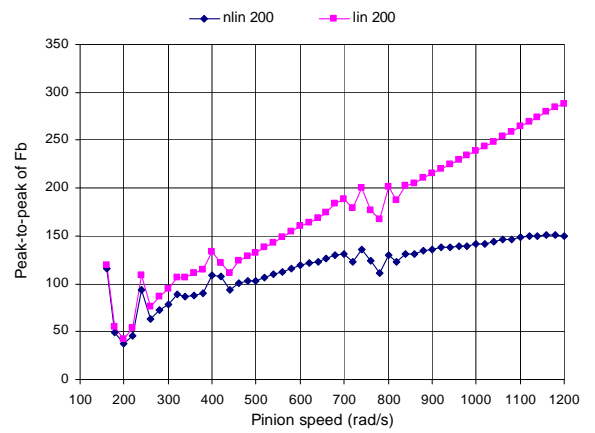
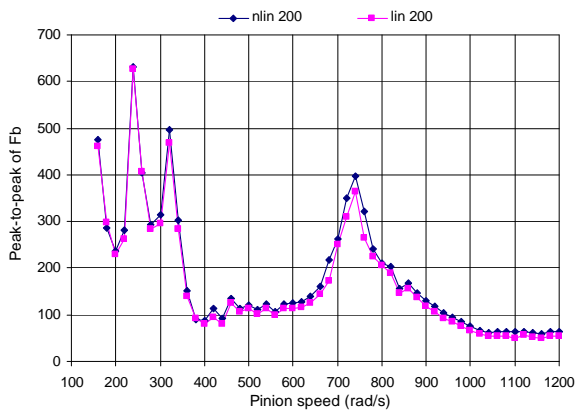
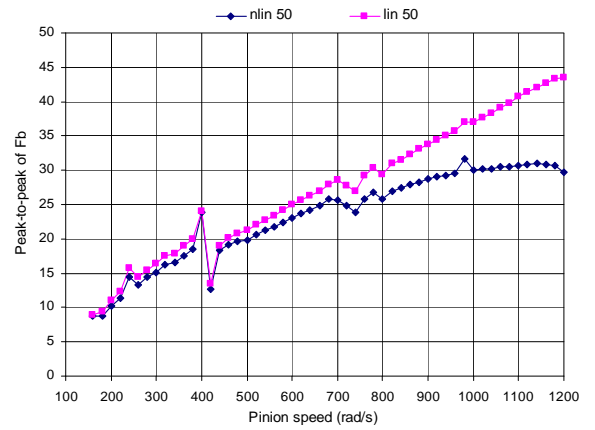
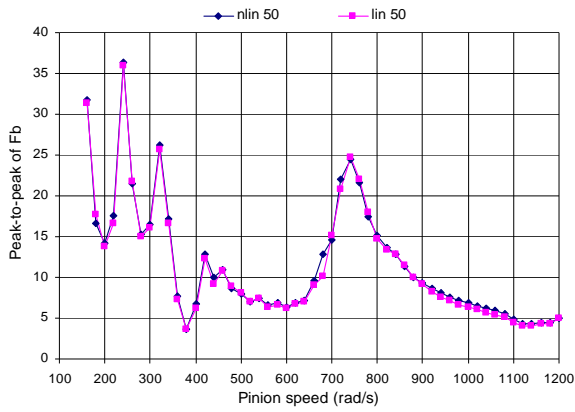


Figure 7 – Peak-to-peak of the hydrodynamic load F_b in bearing 1 for different input torques

Figure 8 – Peak-to-peak of the hydrodynamic load F_b in bearing 3 for different input torques

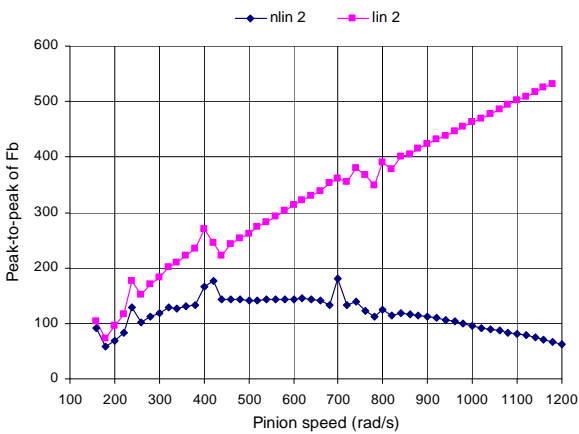
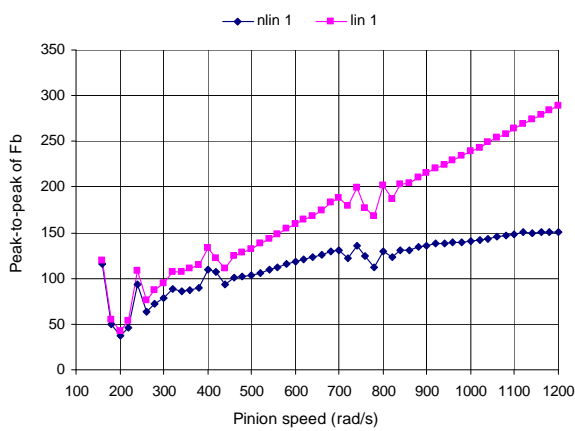
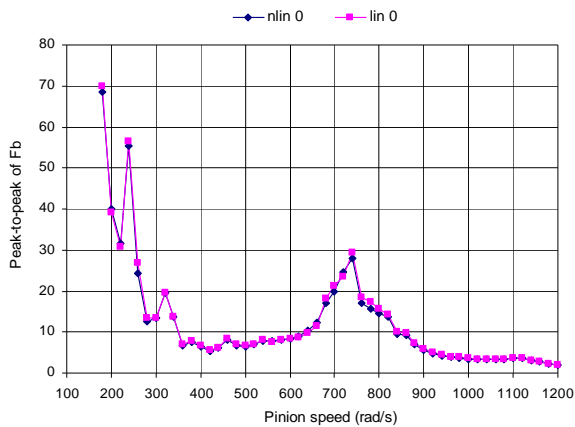


Figure 9 – Peak-to-peak of the hydrodynamic load F_b in bearing 3 for different values of the mass imbalance

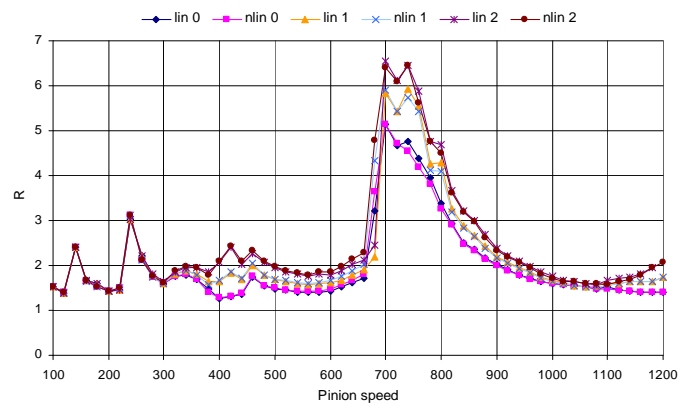


Figure 10 – Maximum dynamic-to-maximum static mesh load ratio for different amplitudes of the mass imbalance

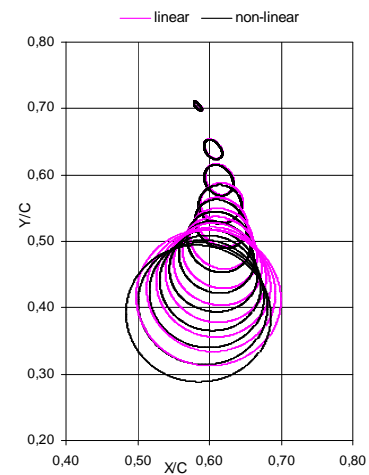
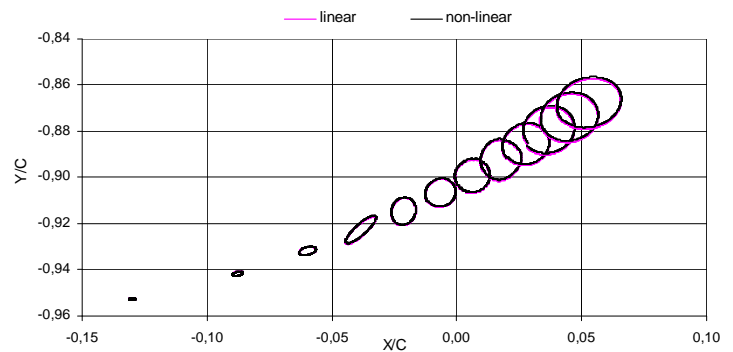


Figure 11 – Orbits of bearings 1 and 3 for values of the input speed from 100 rad/s (smallest orbit) up to 1200 rad/s (largest orbit) obtained by the linear and non-linear theory - Mass imbalance corresponding to case 1

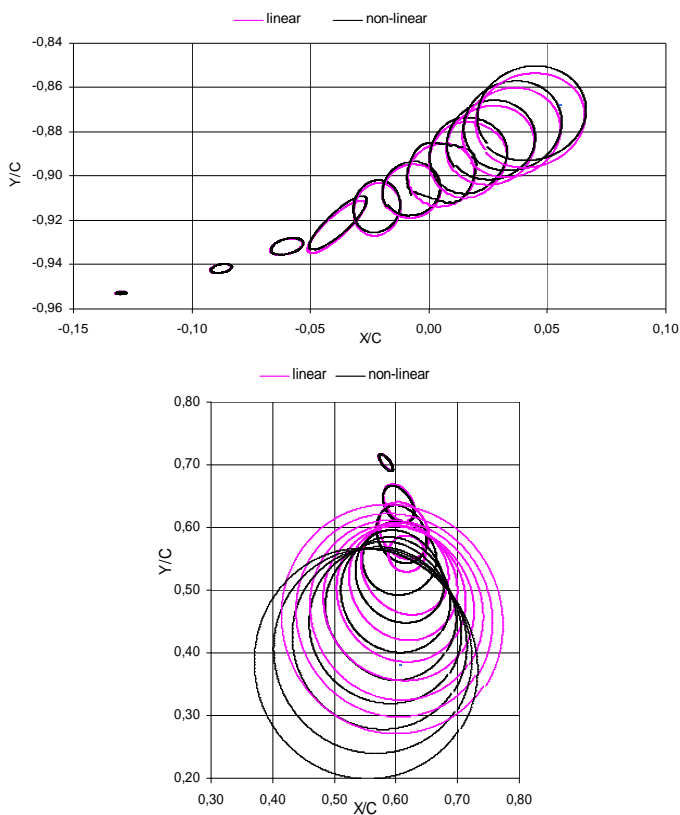


Figure 12 – Orbits of bearings 1 and 3 for values of the input speed from 100 rad/s (smallest orbit) up to 1200 rad/s (largest orbit) obtained by the linear and non-linear theory - Mass imbalance corresponding to case 2

CONCLUSION

This paper presents a model of gear-shaft-bearing systems which accounts for various contact non-linearities and couplings. The state equations are solved by an original technique which combines a Newmark scheme, a Newton-Raphson procedure for bearing non-linearities, and a unilateral contact algorithm which deals with the contact problem on gear teeth. The proposed theory is applied to a simplified F.E. model of a single-stage spur gear system with 2 shafts and 4 bearings.

From the numerical examples, the following observations can be made:

- i) for accurate predictions of vibration transfer through the bearings, bearing non-linear behavior has to be taken into account in the case of large dynamic forces,
- ii) mesh excitations appear in the dynamic bearing response although bearing force amplitudes are usually dominated by imbalance,
- iii) dynamic tooth load sensitivity to once-per-revolution excitations seems very limited.

The present study is part of a research program on dynamic couplings in geared systems and is clearly approximate in nature since it is based on the short-bearing theory which limits the practical range of applications. Further work is under way to account for the influence of helical gears, different bearing technologies and flexible mounts.

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