# Two results on the Dunkl maximal operator 

Luc Deleaval

## To cite this version:

Luc Deleaval. Two results on the Dunkl maximal operator. 2010. hal-00504867v1

## HAL Id: hal-00504867 https://hal.science/hal-00504867v1

Preprint submitted on 21 Jul 2010 (v1), last revised 10 Sep 2013 (v2)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire HAL, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# TWO RESULTS ON THE DUNKL MAXIMAL OPERATOR 

LUC DELEAVAL


#### Abstract

In this article, we first improve the scalar maximal theorem for the Dunkl maximal operator by giving some precisions on the behavior of the constants of this theorem for a general reflection group. Next we complete the vector-valued theorem for the Dunkl-type Fefferman-Stein operator in the case $\mathbb{Z}_{2}^{d}$ by establishing a result of exponential integrability corresponding to the case $p=+\infty$.


## 1. Introduction

Dunkl operators provide an essential tool to extend Fourier analysis on Euclidean spaces and analysis on Riemannian symmetric spaces of Euclidean type. Since their invention in 1989, these operators have largely contributed, in the setting of root systems and associated reflection groups, to the development of harmonic analysis and to the theory of multivariable hypergeometric functions.

In this paper, we focus on the Dunkl maximal operator $M_{\kappa}^{W}$ which is defined by

$$
M_{\kappa}^{W} f(x)=\sup _{r>0} \frac{1}{\mu_{\kappa}^{W}\left(B_{r}\right)}\left|\int_{\mathbb{R}^{d}} f(y) \tau_{x}^{W}\left(\chi_{B_{r}}\right)(-y) \mathrm{d} \mu_{\kappa}^{W}(y)\right|, \quad x \in \mathbb{R}^{d}
$$

where $\chi_{B_{r}}$ is the characteristic function of the Euclidean ball of radius $r$ centered at the origin, $\tau_{x}^{W}$ is the Dunkl translation and $\mu_{\kappa}^{W}$ is a weighted Lebesgue measure invariant under the action of the reflection group $W$ (see Section 2 for more details). This operator, which reduces to the well-known Hardy-Littlewood maximal operator in the case where the multiplicity function $\kappa$ is equal to 0 (see Section 2 for details), is of particular interest in view of developing harmonic analysis associated with root systems. Nevertheless, the structure of the Dunkl translation prevents us from using the tools of real analysis (covering lemma, weighted inequality, Calderón-Zygmund decomposition...) and makes difficult the study of $M_{\kappa}^{W}$.

However, Thangavelu and Xu succeeded in proving the following scalar maximal theorem in [21, where we denote by $L^{p}\left(\mu_{\kappa}^{W}\right)$ the space $L^{p}\left(\mathbb{R}^{d} ; \mu_{\kappa}^{W}\right)$ (for $1 \leqslant p \leqslant+\infty)$ and we use the shorter notation $\|\cdot\|_{W, p}$ instead of $\|\cdot\|_{L^{p}\left(\mu_{\kappa}^{W}\right)}$.
Theorem 1.1 (Scalar maximal theorem). Let $f$ be a measurable function defined on $\mathbb{R}^{d}$.
(1) If $f \in L^{1}\left(\mu_{\kappa}^{W}\right)$, then for every $\lambda>0$ we have

$$
\mu_{\kappa}^{W}\left(\left\{x \in \mathbb{R}^{d}: M_{\kappa}^{W} f(x)>\lambda\right\}\right) \leqslant \frac{C}{\lambda}\|f\|_{W, 1}
$$

where $C=C(d, \kappa)$ is a constant independent of $f$ and $\lambda$.

[^0](2) If $f \in L^{p}\left(\mu_{\kappa}^{W}\right)$, with $1<p \leqslant+\infty$, then $M_{\kappa}^{W} f \in L^{p}\left(\mu_{\kappa}^{W}\right)$ and we have
$$
\left\|M_{\kappa}^{W} f\right\|_{W, p} \leqslant C\|f\|_{W, p}
$$
where $C=C(d, \kappa, p)$ is a constant independent of $f$.
In order to prove this theorem, Thangavelu and Xu have used the following Hopf-Dunford-Schwartz ergodic theorem (see [3]).
Theorem 1.2 (Hopf-Dunford-Schwartz ergodic theorem). Let $X$ be a measurable space and let $m$ be a positive measure on $X$. Let $\left\{T_{t}\right\}_{t \geqslant 0}$ be a contraction semi-group of operators on $L^{p}(X ; m)$, that is to say a semi-group which satisfies for every $p \in[1,+\infty]$ and every $f \in L^{p}(X ; m)$
$$
\left\|T_{t} f\right\|_{L^{p}(X ; m)} \leqslant\|f\|_{L^{p}(X ; m)}
$$

Let us denote by $\mathscr{M} f$ the function defined by

$$
\mathscr{M} f(x)=\sup _{t>0}\left|\frac{1}{t} \int_{0}^{t} T_{s} f(x) \mathrm{d} s\right| .
$$

(1) If $f \in L^{1}(X ; m)$, then for every $\lambda>0$ we have

$$
m(\{x \in X: \mathscr{M} f(x)>\lambda\}) \leqslant \frac{2}{\lambda}\|f\|_{L^{1}(X ; m)}
$$

(2) If $f \in L^{p}(X ; m)$, with $1<p \leqslant+\infty$, then $\mathscr{M} f \in L^{p}(X ; m)$ and we have

$$
\|\mathscr{M} f\|_{L^{p}(X ; m)} \leqslant C\|f\|_{L^{p}(X ; m)}
$$

where $C=C(p)$ is a constant independent of $f$.
We will see that we can use sharply the previous theorem in order to refine the scalar maximal theorem. More precisely, our first result will be the following one, where we denote by $2 \gamma$ the degree of homogeneity of the measure $\mu_{\kappa}^{W}$.
Theorem 1.3. Let $f$ be a measurable function defined on $\mathbb{R}^{d}$.
(1) There exists a numerical constant $C$ such that if $f \in L^{1}\left(\mu_{\kappa}^{W}\right)$, then we have for every $\lambda>0$

$$
\mu_{\kappa}^{W}\left(\left\{x \in \mathbb{R}^{d}: M_{\kappa}^{W} f(x)>\lambda\right\}\right) \leqslant C \frac{d+2 \gamma}{\lambda}\|f\|_{W, 1}
$$

(2) There exists a numerical constant $C$ such that if $f \in L^{p}\left(\mu_{\kappa}^{W}\right)$, with $1<p \leqslant$ $+\infty$, then we have

$$
\left\|M_{\kappa}^{W} f\right\|_{W, p} \leqslant C\left(\frac{p}{p-1}\right) \sqrt{d+2 \gamma}\|f\|_{W, p}
$$

In the particular case where $\gamma=0$, the previous theorem coincides with a theorem due to Stein and Strömberg for the Hardy-Littlewood maximal operator ([18]).

Our second result deals with the vector-valued extension of the scalar maximal theorem which has been proved in [2] in the case where the reflection group is $\mathbb{Z}_{2}^{d}$. Let us recall this theorem. If we denote by $\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}}$ the Dunkl-type Fefferman-Stein operator given for a sequence $f=\left(f_{n}\right)_{n \geqslant 1}$ of measurable functions by

$$
\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}} f=\left(M_{\kappa}^{\mathbb{Z}_{2}^{d}} f_{n}\right)_{n \geqslant 1},
$$

then the theorem is the following.

Theorem 1.4 (Vector-valued maximal theorem). Let $W=\mathbb{Z}_{2}^{d}$ and let $f=$ $\left(f_{n}\right)_{n \geqslant 1}$ be a sequence of measurable functions defined on $\mathbb{R}^{d}$.
(1) Let $1<r<+\infty$. If $\|f\|_{\ell^{r}} \in L^{1}\left(\mu_{\kappa^{2}}^{\mathbb{Z}_{2}^{d}}\right)$, then for every $\lambda>0$ we have

$$
\mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\left\{x \in \mathbb{R}^{d}:\left\|\mathcal{M}_{\kappa^{\mathbb{Z}_{2}^{d}}}(x)\right\|_{\ell^{r}}>\lambda\right\}\right) \leqslant \frac{C}{\lambda}\| \| f\left\|_{\ell^{r}}\right\|_{\mathbb{Z}_{2}^{d}, 1}
$$

where $C=C(d, \kappa, r)$ is a constant independent of $\left(f_{n}\right)_{n \geqslant 1}$ and $\lambda$.
(2) Let $1<r<+\infty$ and let $1<p<+\infty$. If $\|f\|_{\ell^{r}} \in L^{p}\left(\mu_{\kappa^{\mathbb{d}}}^{\mathbb{Z}^{d}}\right)$, then we have

$$
\left\|\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}} f\right\|_{\ell^{r}}\right\|_{\mathbb{Z}_{2}^{d}, p} \leqslant C\| \| f\left\|_{\ell^{r}}\right\|_{\mathbb{Z}_{2}^{d}, p}
$$

where $C=C(d, \kappa, p, r)$ is a constant independent of $\left(f_{n}\right)_{n \geqslant 1}$.
We will see that no analogue of the second point of the previous theorem holds when $p=+\infty$. However, we will give in this case the following result of exponential integrability on every compact set which generalizes the classical one due to Fefferman and Stein (see [10] or [20, page 75]).

Theorem 1.5. Let $f=\left(f_{n}\right)_{n \geqslant 1}$ be a sequence of measurable functions defined on $\mathbb{R}^{d}$ and let $r$ satisfy $1<r<+\infty$. If $\|f\|_{\ell^{r}} \in L^{\infty}\left(\mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\right)$ is such that

$$
\mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\operatorname{supp}\|f\|_{\ell^{r}}^{r}\right)<+\infty
$$

then the function $\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}} f\right\|_{\ell^{r}}^{r}$ is exponentially integrable on every compact set. More precisely there exists a constant which depends only on $d$, $\kappa$ and $r$, denoted by $C_{d, \kappa, r}$ and such that for every compact set $K$ of $\mathbb{R}^{d}$ and for every $\varepsilon$ satisfying

$$
0 \leqslant \varepsilon<\frac{\log (2)}{2 C_{d, \kappa, r}\| \| f\left\|_{\ell^{r}}^{r}\right\|_{\mathbb{Z}_{2}^{d}, \infty}}
$$

we have the inequality

$$
\begin{aligned}
& \int_{K} \mathrm{e}^{\varepsilon\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}^{d}} f(x)\right\|_{\ell^{r}}^{r}} \mathrm{~d} \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}(x) \\
& \quad \leqslant \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}(K)+\frac{2 \varepsilon C_{d, \kappa, r}\| \| f\left\|_{\ell^{r}}^{r}\right\|_{\mathbb{Z}_{2}^{d}, \infty} \max \left\{2 \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}(K) ; \mu_{\kappa^{2}}^{\mathbb{Z}_{2}^{d}}\left(\operatorname{supp}\|f\|_{\ell^{r}}^{r}\right)\right\}}{\log (2)-2 \varepsilon C_{d, \kappa, r}\| \| f\left\|_{\ell^{r}}^{r}\right\|_{\mathbb{Z}_{2}^{d}, \infty}} .
\end{aligned}
$$

The paper is organized as follows. In the next section, we collect some definitions and results related to Dunkl's analysis which will be useful for the sequel. We then give in Section 3 the proof of Theorem1.3. A counterexample in the case $p=+\infty$ is given for the vector-valued maximal theorem in Section 4 and the substitute result contained in Theorem 1.5 is established.

## 2. Preliminaries

This section is devoted to the preliminaries and background. We only focus on the aspects of the Dunkl theory which will be relevant for the sequel. For a large survey about this theory, the reader may especially consult [9, 16] and the references therein.

Let $W \subset \mathcal{O}\left(\mathbb{R}^{d}\right)$ be a finite reflection group associated with a reduced root system $\mathcal{R}$ (not necessarily crystallographic) and let $\kappa: \mathcal{R} \rightarrow \mathbb{C}$ be a multiplicity function, that is a $W$-invariant function. We assume in this article that $\kappa$ takes
value in $[0,+\infty[$.
The (rational) Dunkl operators $T_{\xi}^{\mathcal{R}}$ on $\mathbb{R}^{d}$, which were introduced in 4], are the following $\kappa$-deformations of directional derivatives $\partial_{\xi}$ by reflections

$$
T_{\xi}^{\mathcal{R}} f(x)=\partial_{\xi} f(x)+\sum_{\alpha \in \mathcal{R}_{+}} \kappa(\alpha) \frac{f(x)-f\left(\sigma_{\alpha}(x)\right)}{\langle x, \alpha\rangle}\langle\xi, \alpha\rangle, \quad x \in \mathbb{R}^{d}
$$

where $\sigma_{\alpha}$ denotes the reflection with respect to the hyperplane orthogonal to $\alpha$ and $\mathcal{R}_{+}$denotes a positive subsystem of $\mathcal{R}$. The definition is of course independent of the choice of a positive subsystem since $\kappa$ is $W$-invariant. The most important property of these operators is their commutativity, that is to say $T_{\xi}^{\mathcal{R}} T_{\xi^{\prime}}^{\mathcal{R}}=T_{\xi^{\prime}}^{\mathcal{R}} T_{\xi}^{\mathcal{R}}$ ([4). Therefore, we are naturally led to consider the eigenfunction problem

$$
\begin{equation*}
T_{\xi}^{\mathcal{R}} f=\langle y, \xi\rangle f \quad \forall \xi \in \mathbb{R}^{d} \tag{2.1}
\end{equation*}
$$

with $y \in \mathbb{C}^{d}$ a fixed parameter. This problem has been completely solved by Opdam (11]).
Theorem 2.1. Let $y \in \mathbb{C}^{d}$. There exists a unique $f=E_{\kappa}^{W}(\cdot, y)$ solution of

$$
T_{\xi}^{\mathcal{R}} f=\langle y, \xi\rangle f \quad \forall \xi \in \mathbb{R}^{d}
$$

which is real-analytic on $\mathbb{R}^{d}$ and satisfies $f(0)=1$. Moreover the Dunkl kernel $E_{\kappa}^{W}$ extends to a holomorphic function on $\mathbb{C}^{d} \times \mathbb{C}^{d}$.

In fact, the existence of a solution has been already proved by Dunkl (5). Indeed, he noticed the existence of an intertwining operator $V_{\kappa}^{W}$ which satisfies

$$
V_{\kappa}^{W}\left(\mathcal{P}_{n}^{d}\right)=\mathcal{P}_{n}^{d},\left.\quad V_{\kappa}^{W}\right|_{\mathcal{P}_{0}^{d}}=\operatorname{Id}_{\left.\right|_{\mathcal{P}_{0}^{d}}}, \quad T_{\xi}^{\mathcal{R}} V_{\kappa}^{W}=V_{\kappa}^{W} \partial_{\xi} \quad \forall \xi \in \mathbb{R}^{d}
$$

where $\mathcal{P}_{n}$ denotes the space of homogeneous polynomials of degree $n$ in $d$ variables. Since the exponential function is solution of (2.1) when $\kappa=0$ (that is to say when $T_{\xi}^{\mathcal{R}}=\partial_{\xi}$ ), he naturally set $E_{\kappa}(\cdot, y)=V_{\kappa}^{W}\left(\mathrm{e}^{\langle\cdot, y\rangle}\right)$. Unfortunately, the Dunkl kernel is explicitly known only in some special cases; when the root system is of $A_{2}$-type ([7]), of $B_{2}$-type ( $[8]$ ) and when the reflection group is $\mathbb{Z}_{2}^{d}([6,22])$. Nevertheless we know that this kernel has many properties in common with the classical exponential to which it reduces when $\kappa=0$. For significant results on this kernel and the intertwining operator, the reader may especially consult [1, 5, 11, 12, 14, 17. The Dunkl kernel is of particular interest as it gives rise to an integral transform which is taken with respect to a weighted Lebesgue measure invariant under the action of $W$ and which generalizes the Euclidean Fourier transform.

More precisely, let us introduce the measure $\mathrm{d} \mu_{\kappa}^{W}(x)=h_{\kappa}^{2}(x) \mathrm{d} x$ where the weight given by

$$
h_{\kappa}^{2}(x)=\prod_{\alpha \in \mathcal{R}_{+}}|\langle x, \alpha\rangle|^{2 \kappa(\alpha)}
$$

is homogeneous of degree $2 \gamma$ with

$$
\gamma=\sum_{\alpha \in \mathcal{R}_{+}} \kappa(\alpha)
$$

Then, for every $f \in L^{1}\left(\mu_{\kappa}^{W}\right)$, the Dunkl transform of $f$, denoted by $\mathcal{F}_{\kappa}^{W}(f)$ is defined by

$$
\mathcal{F}_{\kappa}^{W}(f)(x)=c_{\kappa}^{W} \int_{\mathbb{R}^{d}} E_{\kappa}^{W}(-i x, y) f(y) \mathrm{d} \mu_{\kappa}^{W}(y), \quad x \in \mathbb{R}^{d}
$$

where $c_{\kappa}^{W}$ is the Mehta-type constant

$$
c_{\kappa}^{W}=\left(\int_{\mathbb{R}^{d}} \mathrm{e}^{-\frac{\|x\|^{2}}{2}} \mathrm{~d} \mu_{\kappa}^{W}(x)\right)^{-1}
$$

Let us point out that the Dunkl transform coincides with the Euclidean Fourier transform when $\kappa=0$ and that it is more or less a Hankel transform when $d=1$. The two main properties of the Dunkl transform are given in the following theorem ([1, 6]).
Theorem 2.2. (1) Inversion formula Let $f \in L^{1}\left(\mu_{\kappa}^{W}\right)$. If $\mathcal{F}_{\kappa}^{W}(f)$ is in $L^{1}\left(\mu_{\kappa}^{W}\right)$, then we have the following inversion formula

$$
f(x)=c_{\kappa}^{W} \int_{\mathbb{R}^{d}} E_{\kappa}^{W}(i x, y) f(y) \mathrm{d} \mu_{\kappa}^{W}(y)
$$

(2) Plancherel theorem The Dunkl transform has a unique extension to an isometric isomorphism of $L^{2}\left(\mu_{\kappa}^{W}\right)$.

The Dunkl transform shares many other properties with the Fourier transform. Therefore, it is natural to associate a generalized translation operator with this transform.

There are many ways to define the Dunkl translation but we use the definition which most underlines the analogy with the Fourier transform. It is the definition given in 21 with a different convention. Let $x \in \mathbb{R}^{d}$. The Dunkl translation $f \mapsto \tau_{x}^{W} f$ is defined on $L^{2}\left(\mu_{\kappa}^{W}\right)$ by the equation

$$
\mathcal{F}_{\kappa}^{W}\left(\tau_{x}^{W} f\right)(y)=E_{\kappa}^{W}(i x, y) \mathcal{F}_{\kappa}^{W}(f)(y), \quad y \in \mathbb{R}^{d}
$$

It is useful to have a set of functions for which the above equality holds pointwise. It can be down for the following set

$$
\mathscr{A}_{\kappa}^{W}\left(\mathbb{R}^{d}\right)=\left\{f \in L^{1}\left(\mu_{\kappa}^{W}\right): \mathcal{F}_{\kappa}^{W}(f) \in L^{1}\left(\mu_{\kappa}^{W}\right)\right\}
$$

which is a subset of $L^{2}\left(\mu_{\kappa}^{W}\right)$ (since it is contained in the intersection of $L^{1}\left(\mu_{\kappa}^{W}\right)$ and $\left.L^{\infty}\right)$. For $f \in \mathscr{A}_{\kappa}^{W}\left(\mathbb{R}^{d}\right)$, the inversion formula allows us to write

$$
\tau_{x}^{W} f(y)=c_{\kappa}^{W} \int_{\mathbb{R}^{d}} E_{\kappa}^{W}(i x, z) E_{\kappa}^{W}(i y, z) \mathcal{F}_{\kappa}^{W}(f)(z) \mathrm{d} \mu_{\kappa}^{W}(z)
$$

In Fourier analysis, the translation operator $f \mapsto f(\cdot+x)$ (to which the Dunkl translation reduces when $\kappa=0$ ) is positive and $L^{p}$-bounded. In the Dunkl setting, $\tau_{x}^{W}$ is not a positive operator $([12,21])$ and the $L^{p}\left(\mu_{\kappa}^{W}\right)$-boundedness is still a challenging problem, apart from the trivial case where $p=2$ (thanks to the Plancherel theorem and the fact that $\left.\left|E_{\kappa}^{W}(i x, y)\right| \leqslant 1\right)$. The most general result we have is given in the following theorem ([17, 21]), where we denote by $L_{\text {rad }}^{p}\left(\mu_{\kappa}^{W}\right)$ the space of radial functions of $L^{p}\left(\mu_{\kappa}^{W}\right)$.

Theorem 2.3. (1) For every $p$ satisfying $1 \leqslant p \leqslant 2$ and for every $x \in \mathbb{R}^{d}$, the Dunkl translation $\tau_{x}^{W}: L_{\mathrm{rad}}^{p}\left(\mu_{\kappa}^{W}\right) \rightarrow L^{p}\left(\mu_{\kappa}^{W}\right)$ is a bounded operator.
(2) Let $f \in L^{1}\left(\mu_{\kappa}^{W}\right)$ be a bounded, radial and positive function. Then, for every $x \in \mathbb{R}^{d}$ we have $\tau_{x}^{W} f \geqslant 0$.

The last result we mention about the Dunkl translation is the following one.

Theorem 2.4. Let $f \in L_{\mathrm{rad}}^{1}\left(\mu_{\kappa}^{W}\right)$. Then, for every $x \in \mathbb{R}^{d}$ we have

$$
\int_{\mathbb{R}^{d}} \tau_{x}^{W} f(y) \mathrm{d} \mu_{\kappa}^{W}(y)=\int_{\mathbb{R}^{d}} f(y) \mathrm{d} \mu_{\kappa}^{W}(y)
$$

Another important tool in the Dunkl analysis is the Dunkl-type heat semi-group which has been mainly studied by Rösler ([13, 15]). We are searching for solutions $u \in \mathcal{C}^{2}\left(\mathbb{R}^{d} \times\right] 0,+\infty[) \cap \mathcal{C}_{b}\left(\mathbb{R}^{d} \times[0,+\infty[)\right.$ which solve the following Cauchy problem for the generalized heat equation

$$
(C H)_{\kappa}: \begin{cases}\Delta_{\kappa}^{W} u(x, t) & \left.=\partial_{t} u(x, t) \quad \forall(x, t) \in \mathbb{R}^{d} \times\right] 0,+\infty[ \\ u(\cdot, 0) & =f\end{cases}
$$

with initial data $f$ in the Schwartz space $\mathcal{S}\left(\mathbb{R}^{d}\right)$ and where $\Delta_{\kappa}^{W}=\sum_{j=1}^{d}\left(T_{e_{j}}^{\mathcal{R}}\right)^{2}$ is the Dunkl Laplacian. It is easily noticed that a solution of $\Delta_{\kappa}^{W} u(x, t)=\partial_{t} u(x, t)$ is given on $\left.\mathbb{R}^{d} \times\right] 0,+\infty\left[\right.$ by the generalized Gaussian which is defined for every $x \in \mathbb{R}^{d}$ by

$$
q_{t}^{W}(x)=\frac{c_{\kappa}^{W}}{(2 t)^{\frac{d}{2}+\gamma}} \mathrm{e}^{-\frac{\|x\|^{2}}{4 t}}
$$

and which satisfies the two following properties.
Proposition 2.1. (1) For every $t>0$, we have

$$
\int_{\mathbb{R}^{d}} q_{t}^{W}(x) \mathrm{d} \mu_{\kappa}^{W}(x)=1
$$

(2) For every $t>0$ and for every $x \in \mathbb{R}^{d}$, we have

$$
\mathcal{F}_{\kappa}^{W}\left(q_{t}^{W}\right)(x)=c_{\kappa}^{W} \mathrm{e}^{-t\|x\|^{2}}
$$

The Dunkl-type heat kernel $Q_{\kappa}^{W}$ is defined by taking the Dunkl translation of $q_{t}^{W}$, that is, accordingly to [13]
$Q_{\kappa}^{W}(x, y, t)=\tau_{x}^{W} q_{t}^{W}(-y)=\frac{c_{\kappa}^{W}}{(2 t)^{\frac{d}{2}+\gamma}} \mathrm{e}^{-\frac{\left(\|x\|^{2}+\|y\|^{2}\right)}{4 t}} E_{\kappa}^{W}\left(\frac{x}{\sqrt{2 t}}, \frac{y}{\sqrt{2 t}}\right), x, y \in \mathbb{R}^{d}, t>0$.
This positive kernel ([13]) allows us to define a generalized heat operator (or Dunkl-type heat operator). More precisely for every $f \in L^{p}\left(\mu_{\kappa}^{W}\right)$, with $1 \leqslant p \leqslant$ $+\infty$, and for every $t \geqslant 0$, we set

$$
H_{t}^{W} f= \begin{cases}\int_{\mathbb{R}^{d}} f(y) Q_{\kappa}^{W}(\cdot, y, t) \mathrm{d} \mu_{\kappa}^{W}(y) & \text { if } t>0 \\ f & \text { if } t=0\end{cases}
$$

The fundamental result about this operator is the following one. It is due to Rösler (see 13] and [15]).
Theorem 2.5. For every $p$ satisfying $1 \leqslant p \leqslant+\infty$, the family $\left\{H_{t}^{W}\right\}_{t \geqslant 0}$ is a positive and contraction semi-group on $L^{p}\left(\mu_{\kappa}^{W}\right)$.
Moreover, for every $f \in \mathcal{S}\left(\mathbb{R}^{d}\right)$, the function $u$ given for every $(x, t) \in \mathbb{R}^{d} \times[0,+\infty[$ by

$$
u(x, t)=H_{t}^{W} f(x)
$$

belongs to $\mathcal{C}^{2}\left(\mathbb{R}^{d} \times\right] 0,+\infty[) \cap \mathcal{C}_{b}\left(\mathbb{R}^{d} \times[0,+\infty[)\right.$ and is solution of the Cauchy problem $(\mathrm{CH})_{\kappa}$.

We can easily improve the previous theorem.

Theorem 2.6. For every $p$ satisfying $1 \leqslant p \leqslant+\infty$, the family $\left\{H_{t}^{W}\right\}_{t \geqslant 0}$ is a symmetric diffusion semi-group on $L^{p}\left(\mu_{\kappa}^{W}\right)$, that is to say a semi-group which satisfies
(1) $H_{t}^{W}$ is a contraction on $L^{p}\left(\mu_{\kappa}^{W}\right)$, where $p$ satisfies $1 \leqslant p \leqslant+\infty$;
(2) $H_{t}^{W}$ is symmetric, that is to say self-adjoint on $L^{2}\left(\mu_{\kappa}^{W}\right)$;
(3) $H_{t}^{W}$ is positive;
(4) $H_{t}^{W}(1)=1$.

The reader is referred to the book of Stein 19 for a detailed study of this kind of semi-group.

## 3. Behavior of the Dunkl maximal operator in the scalar case

We give in this section the proof of Theorem 1.3 for a general reflection group. For more convenience, we prove each point of this theorem separately. Thus, we first establish the following result.
Theorem 3.1. There exists a numerical constant $C$ such that for every $f \in L^{1}\left(\mu_{\kappa}^{W}\right)$ and every $\lambda>0$ we have

$$
\mu_{\kappa}^{W}\left(\left\{x \in \mathbb{R}^{d}: M_{\kappa}^{W} f(x)>\lambda\right\}\right) \leqslant C \frac{d+2 \gamma}{\lambda}\|f\|_{W, 1}
$$

In order to prove this theorem, we need two lemmas. The first one is nothing more than basic calculus. Before stating it, we introduce some notations.
Notations. We denote by $a\left(S^{d-1}\right)$ the following constant

$$
a\left(S^{d-1}\right)=\int_{S^{d-1}} h_{\kappa}^{2}(x) \mathrm{d} \omega(x)
$$

where $\omega$ is the usual Lebesgue measure on $S^{d-1}$. We also use the following notation

$$
\left.q_{t}^{W}\right|_{S^{d-1}}=\frac{c_{\kappa}^{W}}{(2 t)^{\frac{d}{2}+\gamma}} \mathrm{e}^{-\frac{1}{4 t}}
$$

With these notations in mind, we can now formulate the lemma.
Lemma 3.1. We have the following equalities

$$
\begin{aligned}
& \mu_{\kappa}^{W}\left(B_{1}\right)=\frac{a\left(S^{d-1}\right)}{d+2 \gamma} \\
& \left(c_{\kappa}^{W}\right)^{-1}=2^{\frac{d}{2}+\gamma-1} \Gamma\left(\frac{d}{2}+\gamma\right) a\left(S^{d-1}\right) \\
& \left.\int_{0}^{+\infty} q_{t}^{W}\right|_{S^{d-1}} \\
& \mathrm{~d} t=\frac{c_{\kappa}^{W} 2^{\frac{d}{2}+\gamma}}{4} \Gamma\left(\frac{d}{2}+\gamma-1\right) .
\end{aligned}
$$

Moreover if we assume $d+2 \gamma \geqslant 8$, then we have the following inequality

$$
\left.\int_{\frac{1}{d+2 \gamma}}^{+\infty} q_{t}^{W}\right|_{S^{d-1}} \mathrm{~d} t \leqslant \frac{c_{\kappa}^{W} 2^{\frac{d}{2}+\gamma}}{4}\left(\frac{d+2 \gamma}{4}\right)^{\frac{d}{2}+\gamma-1} \mathrm{e}^{-\frac{d+2 \gamma}{4}}
$$

Proof. The equalities are easy to prove by passing to polar coordinates and by substitution. Therefore we only prove the inequality. We have by definition

$$
\left.\int_{\frac{1}{d+2 \gamma}}^{+\infty} q_{t}^{W}\right|_{S^{d-1}} \mathrm{~d} t=c_{\kappa}^{W} \int_{\frac{1}{d+2 \gamma}}^{+\infty} \frac{1}{(2 t)^{\frac{d}{2}+\gamma}} \mathrm{e}^{-\frac{1}{4 t}} \mathrm{~d} t
$$

which leads after a change of variables to

$$
\left.\int_{\frac{1}{d+2 \gamma}}^{+\infty} q_{t}^{W}\right|_{S^{d-1}} \mathrm{~d} t=\frac{c_{\kappa}^{W} 2^{\frac{d}{2}+\gamma}}{4} \int_{0}^{\frac{d+2 \gamma}{4}} t^{\frac{d}{2}+\gamma-2} \mathrm{e}^{-t} \mathrm{~d} t
$$

Since we have for $d+2 \gamma \geqslant 8$ and $t \in\left[0, \frac{d+2 \gamma}{4}\right]$

$$
t^{\frac{d}{2}+\gamma-2} \mathrm{e}^{-t} \leqslant\left(\frac{d+2 \gamma}{4}\right)^{\frac{d}{2}+\gamma-2} \mathrm{e}^{-\frac{d+2 \gamma}{4}},
$$

we then obtain

$$
\int_{0}^{\frac{d+2 \gamma}{4}} t^{\frac{d}{2}+\gamma-2} \mathrm{e}^{-t} \mathrm{~d} t \leqslant\left(\frac{d+2 \gamma}{4}\right)^{\frac{d}{2}+\gamma-1} \mathrm{e}^{-\frac{d+2 \gamma}{4}}
$$

and the inequality is proved.
The second lemma allows us to reduce the inequality of Theorem 3.1 to a more convenient one.

Lemma 3.2. If there exists $t_{0}>0$ such that

$$
\frac{1}{\mu_{\kappa}^{W}\left(B_{1}\right)} \leqslant\left.\frac{C(d, \kappa)}{t_{0}} \int_{0}^{t_{0}} q_{t}^{W}\right|_{S^{d-1}} \mathrm{~d} t
$$

then we have for every $f \in L^{1}\left(\mu_{\kappa}^{W}\right)$ and every $\lambda>0$ the following inequality

$$
\mu_{\kappa}^{W}\left(\left\{x \in \mathbb{R}^{d}: M_{\kappa}^{W} f(x)>\lambda\right\}\right) \leqslant 2 \frac{C(d, \kappa)}{\lambda}\|f\|_{W, 1}
$$

where $C(d, \kappa)$ is the same positive constant in both hypothesis and conclusion of the lemma.

Proof. We can assume that $f$ is nonnegative. If there exists $t_{0}>0$ such that

$$
\begin{equation*}
\frac{1}{\mu_{\kappa}^{W}\left(B_{1}\right)} \leqslant\left.\frac{C(d, \kappa)}{t_{0}} \int_{0}^{t_{0}} q_{t}^{W}\right|_{S^{d-1}} \mathrm{~d} t \tag{3.1}
\end{equation*}
$$

then we claim that we have for every $y \in \mathbb{R}^{d}$

$$
\begin{equation*}
\frac{1}{\mu_{\kappa}^{W}\left(B_{1}\right)} \chi_{B_{1}}(y) \leqslant \frac{C(d, \kappa)}{t_{0}} \int_{0}^{t_{0}} q_{t}^{W}(y) \mathrm{d} t \tag{3.2}
\end{equation*}
$$

Indeed if $\|y\|>1$, there is nothing to do since $\chi_{B_{1}}(y)=0$. If $\|y\| \leqslant 1$, it is enough to use (3.1) and the fact that $\left.q_{t}^{W}\right|_{S^{d-1}} \leqslant q_{t}^{W}(y)$.
As a result, we can write for every $r>0$ and every $y \in \mathbb{R}^{d}$

$$
\frac{1}{\mu_{\kappa}^{W}\left(B_{r}\right)} \chi_{B_{r}}(y)=\frac{1}{r^{d+2 \gamma} \mu_{\kappa}^{W}\left(B_{1}\right)} \chi_{B_{1}}\left(\frac{y}{r}\right) \leqslant \frac{C(d, \kappa)}{r^{d+2 \gamma} t_{0}} \int_{0}^{t_{0}} q_{t}^{W}\left(\frac{y}{r}\right) \mathrm{d} t
$$

which leads after a change of variables to

$$
\frac{1}{\mu_{\kappa}^{W}\left(B_{r}\right)} \chi_{B_{r}}(y) \leqslant \frac{C(d, \kappa)}{r^{2} t_{0}} \int_{0}^{r^{2} t_{0}} q_{t}^{W}(y) \mathrm{d} t
$$

Let $x \in \mathbb{R}^{d}$. Due to the second point of Theorem 2.3 we can assert that

$$
\frac{1}{\mu_{\kappa}^{W}\left(B_{r}\right)} \tau_{x}^{W}\left(\chi_{B_{r}}\right)(-y) \leqslant \frac{C(d, \kappa)}{r^{2} t_{0}} \tau_{x}^{W}\left(\int_{0}^{r^{2} t_{0}} q_{t}^{W}(\cdot) \mathrm{d} t\right)(-y)
$$

Let us temporarily assume that

$$
\begin{equation*}
\tau_{x}^{W}\left(\int_{0}^{r^{2} t_{0}} q_{t}^{W}(\cdot) \mathrm{d} t\right)(-y)=\int_{0}^{r^{2} t_{0}} \tau_{x}^{W}\left(q_{t}^{W}\right)(-y) \mathrm{d} t \tag{3.3}
\end{equation*}
$$

It implies the following inequality

$$
\frac{1}{\mu_{\kappa}^{W}\left(B_{r}\right)} \tau_{x}^{W}\left(\chi_{B_{r}}\right)(-y) \leqslant \frac{C(d, \kappa)}{r^{2} t_{0}} \int_{0}^{r^{2} t_{0}} \tau_{x}^{W}\left(q_{t}^{W}\right)(-y) \mathrm{d} t
$$

We obtain by multiplying both sides by $f(y)$ and by integrating over $\mathbb{R}^{d}$

$$
\frac{1}{\mu_{\kappa}^{W}\left(B_{r}\right)} \int_{\mathbb{R}^{d}} f(y) \tau_{x}^{W}\left(\chi_{B_{r}}\right)(-y) \mathrm{d} \mu_{\kappa}^{W}(y) \leqslant \frac{C(d, \kappa)}{r^{2} t_{0}} \int_{0}^{r^{2} t_{0}} H_{t}^{W} f(x) \mathrm{d} t
$$

It easily yields to

$$
\frac{1}{\mu_{\kappa}^{W}\left(B_{r}\right)} \int_{\mathbb{R}^{d}} f(y) \tau_{x}^{W}\left(\chi_{B_{r}}\right)(-y) \mathrm{d} \mu_{\kappa}^{W}(y) \leqslant C(d, \kappa) \sup _{s>0}\left(\frac{1}{s} \int_{0}^{s} H_{t}^{W} f(x) \mathrm{d} t\right)
$$

from which we can deduce that

$$
M_{\kappa}^{W} f(x) \leqslant C(d, \kappa) \sup _{s>0}\left(\frac{1}{s} \int_{0}^{s} H_{t}^{W} f(x) \mathrm{d} t\right)
$$

Since $\left\{H_{t}^{W}\right\}_{t \geqslant 0}$ is a semi-group which satisfies the contraction property on $L^{p}\left(\mu_{\kappa}^{W}\right)$ (Theorem 2.5), the first point of the Hopf-Dunford-Schwartz ergodic theorem implies the desired conclusion and we are left with the task of establishing (3.3). Let $n$ be an integer. We first prove that we have for every $y \in \mathbb{R}^{d}$

$$
\begin{equation*}
\tau_{x}^{W}\left(\int_{\frac{1}{n}}^{r^{2} t_{0}} q_{t}^{W}(\cdot) \mathrm{d} t\right)(y)=\int_{\frac{1}{n}}^{r^{2} t_{0}} \tau_{x}^{W}\left(q_{t}^{W}\right)(y) \mathrm{d} t \tag{3.4}
\end{equation*}
$$

Due to Proposition 2.1 it is easily seen that the radial function $\int_{\frac{1}{n}}^{r^{2} t_{0}} q_{t}^{W}(\cdot) \mathrm{d} t$ is in $L^{1}\left(\mu_{\kappa}^{W}\right)$ so using Theorem 2.3 leads us to

$$
\tau_{x}^{W}\left(\int_{\frac{1}{n}}^{r^{2} t_{0}} q_{t}^{W}(\cdot) \mathrm{d} t\right) \in L^{1}\left(\mu_{\kappa}^{W}\right)
$$

Since we have on one hand

$$
\mathcal{F}_{\kappa}^{W}\left(\int_{\frac{1}{n}}^{r^{2} t_{0}} q_{t}^{W}(\cdot) \mathrm{d} t\right)=\int_{\frac{1}{n}}^{r^{2} t_{0}} c_{\kappa}^{W} \mathrm{e}^{-t\|\cdot\|^{2}} \mathrm{~d} t \in L^{1}\left(\mu_{\kappa}^{W}\right)
$$

and on the other hand

$$
\mathcal{F}_{\kappa}^{W}\left(\tau_{x}^{W}\left(\int_{\frac{1}{n}}^{r^{2} t_{0}} q_{t}^{W}(\cdot) \mathrm{d} t\right)\right)=E_{\kappa}^{W}(i x, \cdot) \mathcal{F}_{\kappa}^{W}\left(\int_{\frac{1}{n}}^{r^{2} t_{0}} q_{t}^{W}(\cdot) \mathrm{d} t\right)
$$

we claim that

$$
\mathcal{F}_{\kappa}^{W}\left(\tau_{x}^{W}\left(\int_{\frac{1}{n}}^{r^{2} t_{0}} q_{t}^{W}(\cdot) \mathrm{d} t\right)\right) \in L^{1}\left(\mu_{\kappa}^{W}\right)
$$

Consequently, $\tau_{x}^{W}\left(\int_{\frac{1}{n}}^{r^{2} t_{0}} q_{t}^{W}(\cdot) \mathrm{d} t\right) \in \mathscr{A}_{\kappa}^{W}\left(\mathbb{R}^{d}\right)$ and we can use the inversion formula to obtain

$$
\tau_{x}^{W}\left(\int_{\frac{1}{n}}^{r^{2} t_{0}} q_{t}^{W}(\cdot) \mathrm{d} t\right)(y)=c_{\kappa}^{W} \int_{\mathbb{R}^{d}} E_{\kappa}^{W}(i x, z) E_{\kappa}^{W}(i y, z) \int_{\frac{1}{n}}^{r^{2} t_{0}} c_{\kappa}^{W} \mathrm{e}^{-t\|z\|^{2}} \mathrm{~d} t \mathrm{~d} \mu_{\kappa}^{W}(z)
$$

that is to say

$$
\begin{aligned}
\tau_{x}^{W}\left(\int_{\frac{1}{n}}^{r^{2} t_{0}} q_{t}^{W}(\cdot) \mathrm{d} t\right. & )(y) \\
& =\int_{\frac{1}{n}}^{r^{2} t_{0}}\left(c_{\kappa}^{W} \int_{\mathbb{R}^{d}} E_{\kappa}^{W}(i x, z) E_{\kappa}^{W}(i y, z) c_{\kappa}^{W} \mathrm{e}^{-t\|z\|^{2}} \mathrm{~d} \mu_{\kappa}^{W}(z)\right) \mathrm{d} t
\end{aligned}
$$

The use of the inversion formula leads us to

$$
\tau_{x}^{W}\left(\int_{\frac{1}{n}}^{r^{2} t_{0}} q_{t}^{W}(\cdot) \mathrm{d} t\right)(y)=\int_{\frac{1}{n}}^{r^{2} t_{0}} \tau_{x}^{W}\left(q_{t}^{W}\right)(y) \mathrm{d} t
$$

and (3.4) is therefore established. Of course, we will take limit in (3.4) in order to prove (3.3). On one hand, $\int_{\frac{1}{n}}^{r^{2} t_{0}} q_{t}^{W}(\cdot) \mathrm{d} t \rightarrow \int_{0}^{r^{2} t_{0}} q_{t}^{W}(\cdot) \mathrm{d} t$ in $L^{1}\left(\mu_{\kappa}^{W}\right)$ as $n$ goes to infinity. Therefore, we obtain by using the first point of Theorem 2.3

$$
\tau_{x}^{W}\left(\int_{\frac{1}{n}}^{r^{2} t_{0}} q_{t}^{W}(\cdot) \mathrm{d} t\right) \rightarrow \tau_{x}^{W}\left(\int_{0}^{r^{2} t_{0}} q_{t}^{W}(\cdot) \mathrm{d} t\right)
$$

in $L^{1}\left(\mu_{\kappa}^{W}\right)$ as $n$ goes to infinity.
On the other hand

$$
0 \leqslant \int_{0}^{r^{2} t_{0}} \tau_{x}^{W}\left(q_{t}^{W}\right)(y) \mathrm{d} t-\int_{\frac{1}{n}}^{r^{2} t_{0}} \tau_{x}^{W}\left(q_{t}^{W}\right)(y) \mathrm{d} t=\int_{0}^{\frac{1}{n}} \tau_{x}^{W}\left(q_{t}^{W}\right)(y) \mathrm{d} t
$$

But we claim that

$$
\int_{\mathbb{R}^{d}} \int_{0}^{\frac{1}{n}} \tau_{x}^{W}\left(q_{t}^{W}\right)(y) \mathrm{d} t \mathrm{~d} \mu_{\kappa}^{W}(y)=\int_{0}^{\frac{1}{n}} \int_{\mathbb{R}^{d}} \tau_{x}^{W}\left(q_{t}^{W}\right)(y) \mathrm{d} \mu_{\kappa}^{W}(y) \mathrm{d} t=\frac{1}{n}
$$

since we have thanks to the second point of Theorem 2.4

$$
\int_{\mathbb{R}^{d}} \tau_{x}^{W}\left(q_{t}^{W}\right)(y) \mathrm{d} \mu_{\kappa}^{W}(y)=\int_{\mathbb{R}^{d}} q_{t}^{W}(y) \mathrm{d} \mu_{\kappa}^{W}(y)=1
$$

Consequently

$$
\int_{\frac{1}{n}}^{r^{2} t_{0}} \tau_{x}^{W}\left(q_{t}^{W}\right)(y) \mathrm{d} t \rightarrow \int_{0}^{r^{2} t_{0}} \tau_{x}^{W}\left(q_{t}^{W}\right)(y) \mathrm{d} t
$$

in $L^{1}\left(\mu_{\kappa}^{W}\right)$ as $n$ goes to infinity.
The equality (3.3) is true, and the lemma is proved.
We are now in a position to prove Theorem 3.1.
Proof. We can assume without loss of generality that $f$ is nonnegative. Due to Lemma 3.2, it is enough to find $t_{0}>0$ such that

$$
\begin{equation*}
\frac{1}{\mu_{\kappa}^{W}\left(B_{1}\right)} \leqslant\left. C \frac{d+2 \gamma}{t_{0}} \int_{0}^{t_{0}} q_{t}^{W}\right|_{S^{d-1}} \mathrm{~d} t \tag{3.5}
\end{equation*}
$$

where $C$ is a numerical constant. Set $t_{0}=\frac{1}{d+2 \gamma}$. The inequality of Lemma 3.1 asserts that we have for $d+2 \gamma \geqslant 8$

$$
\left.\int_{\frac{1}{d+2 \gamma}}^{+\infty} q_{t}^{W}\right|_{S^{d-1}} \mathrm{~d} t \leqslant \frac{c_{\kappa}^{W} 2^{\frac{d}{2}+\gamma}}{4}\left(\frac{d+2 \gamma}{4}\right)^{\frac{d}{2}+\gamma-1} \mathrm{e}^{-\frac{d+2 \gamma}{4}}
$$

On one hand, Stirling's formula implies that

$$
\left(\frac{n}{4}\right)^{\frac{n}{2}-1} \mathrm{e}^{-\frac{n}{4}}=\underset{n \rightarrow+\infty}{\mathrm{o}}\left(\Gamma\left(\frac{n}{2}-1\right)\right)
$$

and, on the other hand, the third equality of Lemma 3.1 yields to

$$
\left.\int_{0}^{+\infty} q_{t}^{W}\right|_{S^{d-1}} \mathrm{~d} t=\frac{c_{\kappa}^{W} 2^{\frac{d}{2}+\gamma}}{4} \Gamma\left(\frac{d}{2}+\gamma-1\right)
$$

Therefore, we can conclude that there exists a numerical constant $C$ such that

$$
\begin{equation*}
c_{\kappa}^{W} 2^{\frac{d}{2}+\gamma} \Gamma\left(\frac{d}{2}+\gamma-1\right) \leqslant\left. C \int_{0}^{\frac{1}{d+2 \gamma}} q_{t}^{W}\right|_{S^{d-1}} \mathrm{~d} t \tag{3.6}
\end{equation*}
$$

We can write by using the first two equalities of Lemma 3.1

$$
c_{\kappa}^{W} 2^{\frac{d}{2}+\gamma} \Gamma\left(\frac{d}{2}+\gamma-1\right)=\frac{2 \Gamma\left(\frac{d}{2}+\gamma-1\right)}{\mu_{\kappa}^{W}\left(B_{1}\right)(d+2 \gamma) \Gamma\left(\frac{d}{2}+\gamma\right)},
$$

and, by inserting the previous equality in (3.6) we are led to

$$
\frac{1}{\mu_{\kappa}^{W}\left(B_{1}\right)} \leqslant\left. C \frac{(d+2 \gamma) \Gamma\left(\frac{d}{2}+\gamma\right)}{\Gamma\left(\frac{d}{2}+\gamma-1\right)} \int_{0}^{\frac{1}{d+2 \gamma}} q_{t}^{W}\right|_{S^{d-1}} \mathrm{~d} t
$$

which finally implies that

$$
\frac{1}{\mu_{\kappa}^{W}\left(B_{1}\right)} \leqslant\left. C(d+2 \gamma)^{2} \int_{0}^{\frac{1}{d+2 \gamma}} q_{t}^{W}\right|_{S^{d-1}} \mathrm{~d} t
$$

This proves (3.5) and the theorem is established.
We now turn to the second point of Theorem 1.3 that we recall below.
Theorem 3.2. There exists a numerical constant $C$ such that for every $f \in$ $L^{p}\left(\mu_{\kappa}^{W}\right)$, with $1<p \leqslant+\infty$, we have

$$
\left\|M_{\kappa}^{W} f\right\|_{W, p} \leqslant C\left(\frac{p}{p-1}\right) \sqrt{d+2 \gamma}\|f\|_{W, p}
$$

Remark. The result of the previous theorem is better that one would obtain by using Theorem 3.1, the $L^{\infty}$ case and the Marcinkiewicz interpolation theorem.

In order to prove Theorem 3.2 we need the following lemma which reduces the inequality of this theorem to a more convenient one.

Lemma 3.3. If there exists $t_{0}>0$ such that

$$
\frac{1}{\mu_{\kappa}^{W}\left(B_{1}\right)} \leqslant\left. C(d, \kappa) q_{t_{0}}^{W}\right|_{S^{d-1}},
$$

then there exists a numerical constant $C$ such that for every $p$ satisfying $1<p \leqslant$ $+\infty$ and every $f \in L^{p}\left(\mu_{\kappa}^{W}\right)$

$$
\left\|M_{\kappa}^{W} f\right\|_{W, p} \leqslant C\left(\frac{p}{p-1}\right) C(d, \kappa)\|f\|_{W, p}
$$

where $C(d, \kappa)$ is the same positive constant in both hypothesis and conclusion of the lemma.

Proof. The proof is quite similar to the one of Lemma3.2. Without any restriction, one can assume that $f$ is nonnegative. If there exists $t_{0}>0$ such that

$$
\begin{equation*}
\frac{1}{\mu_{\kappa}^{W}\left(B_{1}\right)} \leqslant\left. C(d, \kappa) q_{t_{0}}^{W}\right|_{S^{d-1}}, \tag{3.7}
\end{equation*}
$$

then we can deduce that for every $y \in \mathbb{R}^{d}$

$$
\begin{equation*}
\frac{1}{\mu_{\kappa}^{W}\left(B_{1}\right)} \chi_{B_{1}}(y) \leqslant C(d, \kappa) q_{t_{0}}^{W}(y) \tag{3.8}
\end{equation*}
$$

Indeed, if $\|y\|>1$, it is obvious since $\chi_{B_{1}}(y)=0$. If $\|y\| \leqslant 1$, it is enough to use (3.7) and the fact that $q_{\left.t_{0}\right|_{S_{d-1}}} \leqslant q_{t_{0}}^{W}(y)$.

Consequently, for every $r>0$, we can write

$$
\frac{1}{\mu_{\kappa}^{W}\left(B_{r}\right)} \chi_{B_{r}}(y)=\frac{1}{r^{d+2 \gamma} \mu_{\kappa}^{W}\left(B_{1}\right)} \chi_{B_{1}}\left(\frac{y}{r}\right) \leqslant \frac{C(d, \kappa)}{r^{d+2 \gamma}} q_{t_{0}}^{W}\left(\frac{y}{r}\right)=C(d, \kappa) q_{r^{2} t_{0}}^{W}(y) .
$$

By using the same arguments given in the proof of Lemma 3.2, we can easily deduce that for every $x \in \mathbb{R}^{d}$

$$
\begin{equation*}
M_{\kappa}^{W} f(x) \leqslant C(d, \kappa) \sup _{t>0} H_{t}^{W} f(x) \tag{3.9}
\end{equation*}
$$

But $\left\{H_{t}^{W}\right\}_{t \geqslant 0}$ is a symmetric diffusion semi-group on $L^{p}\left(\mu_{\kappa}^{W}\right)$ (Theorem 2.6). Therefore, we claim thanks to a result due to Stein (see [19, chapter 4]) that we have for every $p$ satisfying $1<p \leqslant+\infty$

$$
\left\|\sup _{t>0} H_{t}^{W} f\right\|_{W, p} \leqslant C\left(\frac{p}{p-1}\right)\|f\|_{W, p}
$$

where $C$ is a numerical constant. We obtain the desired result by using the previous inequality in (3.9).

We can now turn to the proof of Theorem 3.2.
Proof. Without loss of generality, we can assume that $f$ is nonnegative. Thanks to the previous lemma, it is enough to find $t_{0}>0$ such that

$$
\frac{1}{\mu_{\kappa}^{W}\left(B_{1}\right)} \leqslant\left. C \sqrt{d+2 \gamma} q_{t_{0}}^{W}\right|_{S^{d-1}}
$$

or, equivalently, such that

$$
\frac{1}{\mu_{\kappa}^{W}\left(B_{1}\right)} \leqslant C c_{\kappa}^{W} \sqrt{d+2 \gamma}\left(\frac{1}{2 t_{0}}\right)^{\frac{d}{2}+\gamma} \mathrm{e}^{-\frac{1}{4 t_{0}}}
$$

On one hand, the first two equalities of Lemma 3.1 allow us to write

$$
\frac{1}{\mu_{\kappa}^{W}\left(B_{1}\right)}=c_{\kappa}^{W}(d+2 \gamma) 2^{\frac{d}{2}+\gamma-1} \Gamma\left(\frac{d}{2}+\gamma\right)
$$

and on the other hand, Stirling's formula gives us

$$
2^{\frac{n}{2}} \Gamma\left(\frac{n}{2}\right)=\underset{n \rightarrow+\infty}{\mathrm{O}}\left(n^{\frac{n-1}{2}} \mathrm{e}^{-\frac{n}{2}}\right)
$$

We finally obtain the desired result by choosing $t_{0}=\frac{1}{2 d+4 \gamma}$.

## 4. Exponential integrability in the vector-valued case

This section is devoted to the proof of Theorem 1.5 in the case where the reflection group is $\mathbb{Z}_{2}^{d}$ (that is to say the root system is $\mathcal{R}=\left\{ \pm e_{j}: 1 \leqslant j \leqslant d\right\}$ ). Let us recall some facts related to Dunkl's analysis associated with this particular reflection group.

In this case, an explicit formula of the intertwining operator $V_{\kappa}^{\mathbb{Z}_{2}^{d}}$ is known (see [22]) and there is an explicit formula for the Dunkl translation. In the onedimensional case, the following formula has been proved by Rösler in the setting of signed hypergroups (see [12])

$$
\begin{align*}
\tau_{x}^{\mathbb{Z}_{2}} f(y)= & \frac{1}{2} \int_{-1}^{1} f\left(\sqrt{x^{2}+y^{2}+2 x y t}\right)\left(1+\frac{x+y}{\sqrt{x^{2}+y^{2}+2 x y t}}\right) \psi_{\kappa}(t) \mathrm{d} t  \tag{4.1}\\
& +\frac{1}{2} \int_{-1}^{1} f\left(-\sqrt{x^{2}+y^{2}+2 x y t}\right)\left(1-\frac{x+y}{\sqrt{x^{2}+y^{2}+2 x y t}}\right) \psi_{\kappa}(t) \mathrm{d} t
\end{align*}
$$

where $\psi$ is given by $\psi_{\kappa}(t)=\left(B\left(\kappa, \frac{1}{2}\right)\right)^{-1}(1+t)\left(1-t^{2}\right)^{\kappa-1}$ (with $B$ the beta function) and where $\kappa$ is the only value taken by the multiplicity function $\kappa$. This formula implies an explicit one in the case $\mathbb{Z}_{2}^{d}$ which gives us the boundedness of the Dunkl translation. In order to give an equivalent formula to (4.1), we need to introduce some notations.

Notations. (1) For $x, y, z \in \mathbb{R}$, we put

$$
\sigma_{x, y, z}= \begin{cases}\frac{1}{2 x y}\left(x^{2}+y^{2}-z^{2}\right) & \text { if } x, y \neq 0 \\ 0 & \text { if } x=0 \text { or } y=0\end{cases}
$$

as well as

$$
\varrho(x, y, z)=\frac{1}{2}\left(1-\sigma_{x, y, z}+\sigma_{z, x, y}+\sigma_{z, y, x}\right)
$$

(2) For $x, y, z>0$, we put

$$
K_{\kappa}(x, y, z)=2^{2 \kappa-2}\left(B\left(\kappa, \frac{1}{2}\right)\right)^{-1} \frac{\Delta(x, y, z)^{2 \kappa-2}}{(x y z)^{2 \kappa-1}} \chi_{[|x-y|, x+y]}(z)
$$

where $\Delta(x, y, z)$ denotes the area of the triangle (perhaps degenerated) with sides $x, y, z$.

With these notations in mind, (4.1) can be reformulated (using a change of variables) as follows

$$
\begin{equation*}
\tau_{x}^{\mathbb{Z}_{2}} f(y)=\int_{\mathbb{R}} f(z) \mathrm{d} \nu_{x, y}^{\mathbb{Z}_{2}}(z) \tag{4.2}
\end{equation*}
$$

where the measure $\nu_{x, y}^{\mathbb{Z}_{2}}$ is given by

$$
\mathrm{d} \nu_{x, y}^{\mathbb{Z}_{2}}(z)= \begin{cases}\mathcal{K}_{\kappa}(x, y, z) \mathrm{d} \mu_{\kappa}^{\mathbb{Z}_{2}}(z) & \text { if } x, y \neq 0 \\ \mathrm{~d} \delta_{x}(z) & \text { if } y=0 \\ \mathrm{~d} \delta_{y}(z) & \text { if } x=0\end{cases}
$$

with

$$
\mathcal{K}_{\kappa}(x, y, z)=K_{\kappa}(|x|,|y|,|z|) \varrho(x, y, z) .
$$

Thanks to Rösler, we have the following one-dimensional result.

Theorem 4.1. The measure $\nu_{x, y}^{\mathbb{Z}_{2}}$ satisfies
(1) $\operatorname{supp} \nu_{x, y}^{\mathbb{Z}_{2}}=[-|x|-|y|,-||x|-|y||] \bigcup[| | x|-|y||,|x|+|y|]$ for $x, y \neq 0$.
(2) $\nu_{x, y}^{\mathbb{Z}_{2}}(\mathbb{R})=1$ and $\left\|\nu_{x, y}^{\mathbb{Z}_{2}}\right\|=\int_{\mathbb{R}}\left|\mathrm{d} \nu_{x, y}^{\mathbb{Z}_{2}}\right| \leqslant 4$, for $x, y \in \mathbb{R}$.

We now prove that no analogue of the second point of Theorem 1.4 holds when $p=+\infty$.

Proposition 4.1. The conclusion of the second point of Theorem 1.4 does not hold when $p=+\infty$.

Proof. Let $d=1$ and let $f=\left(f_{n}\right)_{n \geqslant 1}$ be the sequence of functions where for every $n \geqslant 1, f_{n}=\chi_{\left[2^{n-1}, 2^{n}[ \right.}$. We easily see that

$$
\|f\|_{\ell^{r}}=\chi_{[1,+\infty[ } \in L^{\infty}\left(\mu_{\kappa}^{\mathbb{Z}_{2}}\right)
$$

while we will prove that

$$
\begin{equation*}
\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}} f\right\|_{\ell^{r}} \notin L^{\infty}\left(\mu_{\kappa}^{\mathbb{Z}_{2}}\right) \tag{4.3}
\end{equation*}
$$

For every $x \in \mathbb{R}$ and for every $n \geqslant 1$ we have, by definition of $M_{\kappa}^{\mathbb{Z}_{2}} \chi_{\left[2^{n-1}, 2^{n}[ \right.}$, the following inequality

$$
\begin{aligned}
& M_{\kappa}^{\mathbb{Z}_{2}} \chi_{\left[2^{n-1}, 2^{n}\right.}(x) \\
& \quad \geqslant \frac{1}{\mu_{\kappa}^{\mathbb{Z}_{2}}(]-|x|-2^{n},|x|+2^{n}[)} \int_{2^{n-1}}^{2^{n}} \tau_{x}^{\mathbb{Z}_{2}}\left(\chi_{]-|x|-2^{n},|x|+2^{n}[ }\right)(-y) \mathrm{d} \mu_{\kappa}^{\mathbb{Z}_{2}}(y) .
\end{aligned}
$$

But we claim that for every $y \in\left[2^{n-1}, 2^{n}[\right.$

$$
\begin{equation*}
\tau_{x}^{\mathbb{Z}_{2}}\left(\chi_{]-|x|-2^{n},|x|+2^{n} \mid}\right)(-y)=1 \tag{4.4}
\end{equation*}
$$

Indeed, we can write by using (4.2)

$$
\begin{aligned}
& 2^{2-2 \kappa} B\left(\kappa, \frac{1}{2}\right) \tau_{x}^{\mathbb{Z}_{2}}\left(\chi_{]-|x|-2^{n},|x|+2^{n}[ }\right)(-y) \\
& =\int_{\mathbb{R}} \chi_{]-|x|-2^{n},|x|+2^{n}[ }(z) \chi_{]| | x|-|y||,|x|+|y| \mid}(|z|) \frac{\Delta(|x|,|y|,|z|)^{2 \kappa-2}}{(|x y z|)^{2 \kappa-1}} \varrho(x,-y, z) \mathrm{d} \mu_{\kappa}^{\mathbb{Z}_{2}}(z),
\end{aligned}
$$

and, since $|x|+|y|<|x|+2^{n}$, we can deduce that

$$
\begin{aligned}
2^{2-2 \kappa} B\left(\kappa, \frac{1}{2}\right) \tau_{x}^{\mathbb{Z}_{2}} & \left(\chi_{]-|x|-2^{n},|x|+2^{n} \mid}\right)(-y) \\
& =\int_{\mathbb{R}} \chi_{]| | x|-|y||,|x|+|y|\lceil }(|z|) \frac{\Delta(|x|,|y|,|z|)^{2 \kappa-2}}{(|x y z|)^{2 \kappa-1}} \varrho(x,-y, z) \mathrm{d} \mu_{\kappa}^{\mathbb{Z}_{2}}(z),
\end{aligned}
$$

that is to say

$$
\tau_{x}^{\mathbb{Z}_{2}}\left(\chi_{]-|x|-2^{n},|x|+2^{n}}\right)(-y)=\nu_{x,-y}^{\mathbb{Z}_{2}}(\mathbb{R}) .
$$

We obtain (4.4) by using Theorem 4.1. As a result we have the inequality

$$
M_{\kappa}^{\mathbb{Z}_{2}} \chi_{\left[2^{n-1,2^{n}}[ \right.}(x) \geqslant \frac{1}{\mu_{\kappa}^{\mathbb{Z}_{2}}(]-|x|-2^{n},|x|+2^{n}[)} \int_{2^{n-1}}^{2^{n}} \mathrm{~d} \mu_{\kappa}^{\mathbb{Z}_{2}}(y),
$$

and since $\mathrm{d} \mu_{\kappa}^{\mathbb{Z}_{2}}(y)=|y|^{2 \kappa} \mathrm{~d} y$, we are led to

$$
M_{\kappa}^{\mathbb{Z}_{2}} \chi_{\left[2^{n-1}, 2^{n}[ \right.}(x) \geqslant \frac{2^{n(2 \kappa+1)}-2^{(n-1)(2 \kappa+1)}}{2\left(|x|+2^{n}\right)^{2 \kappa+1}}
$$

For every $x$ satisfying $|x| \leqslant 2^{n}$ we then obtain

$$
M_{\kappa}^{\mathbb{Z}_{2}} \chi_{\left[2^{n-1,2^{n}}[ \right.}(x) \geqslant \frac{2^{n(2 \kappa+1)}-2^{(n-1)(2 \kappa+1)}}{2^{(n+1)(2 \kappa+1)+1}}
$$

and we get after simplifications

$$
M_{\kappa}^{\mathbb{Z}_{2}} \chi_{\left[2^{n-1}, 2^{n}[ \right.}(x) \geqslant \frac{1-2^{-(2 \kappa+1)}}{2^{2 \kappa+2}}
$$

Thus, we can write for every $x \in \mathbb{R}$

$$
\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}} f\right\|_{\ell^{r}}^{r} \geqslant \sum_{\left\{n: 2^{n} \geqslant|x|\right\}}\left(\frac{1-2^{-(2 \kappa+1)}}{2^{2 \kappa+2}}\right)^{r}=+\infty
$$

and the proof of (4.3) is finished.
The substitute result in the case where $p=+\infty$ is the one given in Theorem 1.5 that we now prove. In order to do that, we need three lemmas. The first one is just a trivial functional equality.

Lemma 4.1. Let $X$ be a measure space and let $m$ be a positive measure on $X$. Let $\varphi$ be an increasing continuously differentiable function on $[0,+\infty[$ which satisfies $\varphi(0)=0$. Then for every measurable function $f$ on $X$ we have the following equality

$$
\int_{X} \varphi(|f(x)|) \mathrm{d} m(x)=\int_{0}^{+\infty} \varphi^{\prime}(\lambda) m(\{x \in X:|f(x)|>\lambda\}) \mathrm{d} \lambda
$$

Proof. Since $\varphi^{\prime}$ is nonnegative, we can write using the Fubini theorem

$$
\int_{0}^{+\infty} \varphi^{\prime}(\lambda) m(\{x \in X:|f(x)|>\lambda\}) \mathrm{d} \lambda=\int_{X}\left(\int_{0}^{|f(x)|} \varphi^{\prime}(\lambda) \mathrm{d} \lambda\right) \mathrm{d} m(x)
$$

We then obtain the desired result by integrating and using the fact that $\varphi(0)=$ 0 .

The second lemma gives us the behavior of the constant of the second point of Theorem 1.4 when $p$ grows and $r$ is fixed. The result of exponential integrability is closely related to this behavior.

Lemma 4.2. The constant $C(d, \kappa, p, r)$ of the vector-valued maximal theorem is such that

$$
C(d, \kappa, p, r)=\underset{p \rightarrow+\infty}{O}\left(p^{\frac{1}{r}}\right)
$$

Proof. Since the parameter $r$ is fixed, it is enough to consider the proof of Theorem 1.4 when $p>r$. As explained in [2], once we have constructed the operator $M_{\kappa}^{\mathbb{Z}_{2}^{d}, R}$ (see [2] for the definition) and shown that it satisfies a scalar maximal theorem and a weighted inequality, we can follow almost verbatim the proof of the vectorvalued maximal theorem for the Hardy-Littlewood maximal operator (see [10] or [20]). Thus, it is easily seen that the dependence in $p$ is given by the constant (with exponent $\frac{1}{r}$ ) of the maximal theorem for $M_{\kappa}^{\mathbb{Z}_{2}^{d}, R}$ and for the space $L^{\frac{p}{p-r}}\left(\mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\right)$. Since this constant is obtained by interpolation, we can then write that

$$
C(d, \kappa, p, r)=C(d, \kappa, r)\left(C(d, \kappa) \frac{\frac{p}{p-r}}{\frac{p}{p-r}-1}\right)^{\frac{p-r}{p r}}
$$

But it is obvious that

$$
\left(C(d, \kappa) \frac{\frac{p}{p-r}}{\frac{p}{p-r}-1}\right)^{\frac{p-r}{p}}=\underset{p \rightarrow+\infty}{O}(p)
$$

and the lemma is therefore proved.
The last lemma gives us (under the hypothesis of Theorem 1.5) a sharp estimate of the measure of the set $\left\{x \in K:\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}} f(x)\right\|_{\ell^{r}}^{r}>\lambda\right\}$ for every $\lambda>0$ and where $K$ denotes a compact subset of $\mathbb{R}^{d}$. More precisely we have the following inequality.

Lemma 4.3. Let $f=\left(f_{n}\right)_{n \geqslant 1}$ be a sequence of measurable functions defined on $\mathbb{R}^{d}$ and let $r$ satisfy $1<r<+\infty$. If $\|f\|_{\ell^{r}} \in L^{\infty}\left(\mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\right)$ is such that

$$
\mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\operatorname{supp}\|f\|_{\ell^{r}}^{r}\right)<+\infty
$$

then for every compact set $K$ of $\mathbb{R}^{d}$ and every $\lambda>0$ we have

$$
\begin{aligned}
& \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\left\{x \in K:\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}} f(x)\right\|_{\ell^{r}}^{r}>\lambda\right\}\right) \\
& \left.\quad \leqslant \max \left\{2 \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}(K) ; \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\operatorname{supp}\|f\|_{\ell^{r}}^{r}\right)\right\} \mathrm{e}^{-\left(\frac{\log (2)}{2 C\| \| f\left\|_{\ell^{r} r}\right\|_{2}^{d}, \infty}\right)}\right)
\end{aligned}
$$

where $C=C(d, \kappa, r)$ is independent of $\left(f_{n}\right)_{n \geqslant 1}, K$ and $\lambda$.
Proof. For every $p$ satisfying $1 \leqslant p<+\infty$ we can write thanks to the Chebyshev inequality

$$
\begin{aligned}
& \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\left\{x \in K:\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}} f(x)\right\|_{\ell^{r}}^{r}>\lambda\right\}\right) \\
& \leqslant \frac{1}{\lambda^{p}} \int_{\left\{x \in K:\left\|\mathcal{M}_{\kappa^{2}}^{Z_{d}^{d}} f(x)\right\|_{\ell^{r}}^{r}>\lambda\right\}}\left\|\mathcal{M}_{\kappa^{\mathbb{Z}_{2}^{d}}} f(x)\right\|_{\ell^{r}}^{p r} \mathrm{~d} \mu_{\kappa_{2}^{\mathbb{Z}_{2}^{d}}(x)}
\end{aligned}
$$

which implies by enlarging the domain of integration

$$
\mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\left\{x \in K:\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}} f(x)\right\|_{\ell^{r}}^{r}>\lambda\right\}\right) \leqslant \frac{1}{\lambda^{p}} \int_{\mathbb{R}^{d}}\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}} f(x)\right\|_{\ell^{r}}^{p r} \mathrm{~d} \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}(x)
$$

By applying the vector-valued maximal theorem for $\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}}$ we get

$$
\mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\left\{x \in K:\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}} f(x)\right\|_{\ell^{r}}^{r}>\lambda\right\}\right) \leqslant \frac{(C(d, \kappa, p r, r))^{p r}}{\lambda^{p}} \int_{\mathbb{R}^{d}}\|f(x)\|_{\ell^{r}}^{p r} \mathrm{~d} \mu_{\kappa^{\mathbb{Z}}}^{\mathbb{Z}_{2}^{d}}(x),
$$

where $C=C(d, \kappa, p r, r)$ is the constant of Theorem 1.4 (and which is independent of $\left(f_{n}\right)_{n \geqslant 1}, K$ and $\left.\lambda\right)$. Thanks to Lemma 4.2 we are lead to

$$
\mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\left\{x \in K:\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}} f(x)\right\|_{\ell^{r}}^{r}>\lambda\right\}\right) \leqslant \frac{C p^{p}}{\lambda^{p}} \int_{\mathbb{R}^{d}}\|f(x)\|_{\ell^{r}}^{p r} \mathrm{~d} \mu_{\kappa^{Z}}^{\mathbb{Z}_{2}^{d}}(x)
$$

where $C=C(d, \kappa, r)$ is independent of $\left(f_{n}\right)_{n \geqslant 1}, K, \lambda$ and $p$. The hypothesis of the lemma allow us to write

$$
\begin{align*}
\mu_{\kappa^{d}}^{\mathbb{Z}_{2}^{d}}\left(\left\{x \in K:\left\|\mathcal{M}_{\kappa^{\mathbb{Z}}}^{\mathbb{Z}_{2}^{d}} f(x)\right\|_{\ell^{r}}^{r}>\lambda\right\}\right) &  \tag{4.5}\\
& \leqslant\left(\frac{C p\| \| f\left\|_{\ell^{r}}^{r}\right\|_{\mathbb{Z}_{2}^{d}, \infty}}{\lambda}\right)^{p} \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\operatorname{supp}\|f\|_{\ell^{r}}^{r}\right) .
\end{align*}
$$

We now exploit (4.5) by choosing $p$ in terms of $\lambda$.
If $\lambda \geqslant 2 C\| \| f\left\|_{\ell^{r}}^{r}\right\|_{\mathbb{Z}_{2}^{d}, \infty}$, we then put

$$
p=\frac{\lambda}{2 C\| \| f\left\|_{\ell^{r}}^{r}\right\|_{\mathbb{Z}_{2}^{d}, \infty}} \geqslant 1
$$

Therefore, (4.5) can be reformulated as follows

$$
\mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\left\{x \in K:\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}} f(x)\right\|_{\ell^{r}}^{r}>\lambda\right\}\right) \leqslant\left(\frac{1}{2}\right)^{p} \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\operatorname{supp}\|f\|_{\ell^{r}}^{r}\right)
$$

that is to say

$$
\begin{equation*}
\mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\left\{x \in K:\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}} f(x)\right\|_{\ell^{r}}^{r}>\lambda\right\}\right) \leqslant \mathrm{e}^{-\left(\frac{\log (2)}{2 C\| \| f\left\|_{\ell^{\prime} r}\right\|_{\mathbb{Z}_{2}^{d}, \infty}}\right) \lambda} \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\operatorname{supp}\|f\|_{\ell^{r}}^{r}\right) \tag{4.6}
\end{equation*}
$$

If $\lambda \leqslant 2 C\| \| f\left\|_{\ell^{r}}^{r}\right\|_{\mathbb{Z}_{2}^{d}, \infty}$, we immediately write in this case

$$
\mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\left\{x \in K:\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}} f(x)\right\|_{\ell^{r}}^{r}>\lambda\right\}\right) \leqslant \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}(K)
$$

Since

$$
\mathrm{e}^{-\left(\frac{\log (2)}{2 C\| \| f\left\|_{\ell}^{( } r\right\|_{\mathbb{Z}_{2}^{d}, \infty}}\right) \lambda} \geqslant \frac{1}{2}
$$

we claim that

$$
\begin{equation*}
\mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\left\{x \in K:\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}} f(x)\right\|_{\ell^{r}}^{r}>\lambda\right\}\right) \leqslant 2 \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}(K) \mathrm{e}^{-\left(\frac{\log (2)}{2 C\| \| f\left\|_{\ell^{r}}^{\prime}\right\|_{\mathbb{Z}_{2}^{d}, \infty}}\right) \lambda} \tag{4.7}
\end{equation*}
$$

The desired result is then a trivial consequence of (4.6) and (4.7).
We are now in a position to prove Theorem 1.5
Proof. Let $K$ be a compact set of $\mathbb{R}^{d}$ and let $\varepsilon$ be a real number which satisfies

$$
0 \leqslant \varepsilon<\frac{\log (2)}{2 C_{d, \kappa, r}\| \| f\left\|_{\ell^{r}}^{r}\right\|_{\mathbb{Z}_{2}^{d}, \infty}}
$$

where $C_{d, \kappa, r}$ is the constant of Lemma 4.3. We first write

$$
\int_{K} \mathrm{e}^{\varepsilon\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}} f(x)\right\|_{\ell^{r}}^{r}} \mathrm{~d} \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}(x)=\mu_{\kappa}^{\mathbb{Z}_{2}^{d}}(K)+\int_{K}\left(\mathrm{e}^{\varepsilon\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}} f(x)\right\|_{\ell^{r}}^{r}}-1\right) \mathrm{d} \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}(x)
$$

We then apply the equality of Lemma 4.1 to the function $\varphi: t \mapsto \mathrm{e}^{\varepsilon t}-1$ to obtain

$$
\begin{aligned}
& \int_{K} \mathrm{e}^{\varepsilon\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}^{d}} f(x)\right\|_{\ell^{r}}^{r}} \mathrm{~d} \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}(x) \\
&=\mu_{\kappa}^{\mathbb{Z}_{2}^{d}}(K)+\varepsilon \int_{0}^{+\infty} \mathrm{e}^{\varepsilon \lambda} \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\left\{x \in K:\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}_{2}^{d}} f(x)\right\|_{\ell^{r}}^{r}>\lambda\right\}\right) \mathrm{d} \lambda .
\end{aligned}
$$

Thanks to Lemma 4.3, we are led to

$$
\begin{aligned}
& \int_{K} \mathrm{e}^{\varepsilon\left\|\mathcal{M}_{\kappa}^{\mathbb{Z}^{d}} f(x)\right\|_{\ell^{r}}^{r}} \mathrm{~d} \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}(x) \\
& \left.\leqslant \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}(K)+\varepsilon \max \left\{2 \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}(K) ; \mu_{\kappa}^{\mathbb{Z}_{2}^{d}}\left(\operatorname{supp}\|f\|_{\ell^{r}}^{r}\right)\right\} \int_{0}^{+\infty} \mathrm{e}^{\lambda\left(\varepsilon-\frac{\log (2)}{2 C\| \| f \|_{\ell^{r} r} \mathbb{Z}_{2}^{d}, \infty}\right.}\right) \\
& \mathrm{d} \lambda
\end{aligned}
$$

The condition on $\varepsilon$ and an integration allow us to conclude.

## References

[1] Marcel de Jeu. The Dunkl transform. Invent. Math., 113(1):147-162, 1993.
[2] Luc Deleaval. Fefferman-Stein inequalities for the $\mathbb{Z}_{2}^{d}$ Dunkl maximal operator. J. Math. Anal. Appl., 360(2):711-726, 2009.
[3] Nelson Dunford and Jacob T. Schwartz. Linear Operators. I. General Theory. With the assistance of W. G. Bade and R. G. Bartle. Pure and Applied Mathematics, Vol. 7. Interscience Publishers, Inc., New York, 1958.
[4] Charles F. Dunkl. Differential-difference operators associated to reflection groups. Trans. Amer. Math. Soc., 311(1):167-183, 1989.
[5] Charles F. Dunkl. Integral kernels with reflection group invariance. Canad. J. Math., 43(6):1213-1227, 1991.
[6] Charles F. Dunkl. Hankel transforms associated to finite reflection groups. In Hypergeometric functions on domains of positivity, Jack polynomials, and applications (Tampa, FL, 1991), volume 138 of Contemp. Math., pages 123-138. Amer. Math. Soc., Providence, RI, 1992.
[7] Charles F. Dunkl. Intertwining operators associated to the group $S_{3}$. Trans. Amer. Math. Soc., 347(9):3347-3374, 1995.
[8] Charles F. Dunkl. An intertwining operator for the group B2. Glasg. Math. J., 49(2):291-319, 2007.
[9] Charles F. Dunkl and Yuan Xu. Orthogonal polynomials of several variables, volume 81 of Encyclopedia of Mathematics and its Applications. Cambridge University Press, Cambridge, 2001.
[10] C. Fefferman and E. M. Stein. Some maximal inequalities. Amer. J. Math., 93:107-115, 1971.
[11] E. M. Opdam. Dunkl operators, Bessel functions and the discriminant of a finite Coxeter group. Compositio Math., 85(3):333-373, 1993.
[12] Margit Rösler. Bessel-type signed hypergroups on R. In Probability measures on groups and related structures, XI (Oberwolfach, 1994), pages 292-304. World Sci. Publ., River Edge, NJ, 1995.
[13] Margit Rösler. Generalized Hermite polynomials and the heat equation for Dunkl operators. Comm. Math. Phys., 192(3):519-542, 1998.
[14] Margit Rösler. Positivity of Dunkl's intertwining operator. Duke Math. J., 98(3):445-463, 1999.
[15] Margit Rösler. One-parameter semigroups related to abstract quantum models of Calogero type. In Infinite dimensional harmonic analysis (Kyoto, 1999), pages 290-305. Gräbner, Altendorf, 2000.
[16] Margit Rösler. Dunkl operators: theory and applications. In Orthogonal polynomials and special functions (Leuven, 2002), volume 1817 of Lecture Notes in Math., pages 93-135. Springer, Berlin, 2003.
[17] Margit Rösler. A positive radial product formula for the Dunkl kernel. Trans. Amer. Math. Soc., 355(6):2413-2438 (electronic), 2003.
[18] E. M. Stein and J.-O. Strömberg. Behavior of maximal functions in $\mathbf{R}^{n}$ for large $n$. Ark. Mat., 21(2):259-269, 1983.
[19] Elias M. Stein. Topics in harmonic analysis related to the Littlewood-Paley theory. Annals of Mathematics Studies, No. 63. Princeton University Press, Princeton, N.J., 1970.
[20] Elias M. Stein. Harmonic analysis: real-variable methods, orthogonality, and oscillatory integrals, volume 43 of Princeton Mathematical Series. Princeton University Press, Princeton, NJ, 1993.
[21] Sundaram Thangavelu and Yuan Xu. Convolution operator and maximal function for the Dunkl transform. J. Anal. Math., 97:25-55, 2005.
[22] Yuan Xu. Orthogonal polynomials for a family of product weight functions on the spheres. Canad. J. Math., 49(1):175-192, 1997.

Institut de Mathématiques de Jussieu, Université Pierre et Marie Curie, 175, Rue du Chevaleret, 75013 Paris, France

E-mail address: deleaval@math.jussieu.fr


[^0]:    2000 Mathematics Subject Classification. 42B10, 42B25.
    Key words and phrases. Dunkl operators, Maximal operator, Fefferman-Stein-type operator, Real analysis.

    The author is pleased to express his thanks to Banafsheh for her advises in english.

