



HAL
open science

Boltzmann equation without angular cutoff in the whole space: I, Global existence for soft potential

Radjesvarane Alexandre, Yoshinori Morimoto, Seiji Ukai, Chao-Jiang Xu,
Tong Yang

► **To cite this version:**

Radjesvarane Alexandre, Yoshinori Morimoto, Seiji Ukai, Chao-Jiang Xu, Tong Yang. Boltzmann equation without angular cutoff in the whole space: I, Global existence for soft potential. 2010. hal-00496950v2

HAL Id: hal-00496950

<https://hal.science/hal-00496950v2>

Preprint submitted on 27 Oct 2010

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

BOLTZMANN EQUATION WITHOUT ANGULAR CUTOFF IN THE WHOLE SPACE: I, GLOBAL EXISTENCE FOR SOFT POTENTIAL

R. ALEXANDRE, Y. MORIMOTO, S. UKAI, C.-J. XU, AND T. YANG

ABSTRACT. It is known that the singularity in the non-cutoff cross-section of the Boltzmann equation leads to the gain of regularity and gain of weight in the velocity variable. By defining and analyzing a non-isotropy norm which precisely captures the dissipation in the linearized collision operator, we first give a new and precise coercivity estimate for the non-cutoff Boltzmann equation for general physical cross sections. Then the Cauchy problem for the Boltzmann equation is considered in the framework of small perturbation of an equilibrium state. In this part, for the soft potential case in the sense that there is no positive power gain of weight in the coercivity estimate on the linearized operator, we derive some new functional estimates on the nonlinear collision operator. Together with the coercivity estimates, we prove the global existence of classical solutions for the Boltzmann equation in weighted Sobolev spaces.

CONTENTS

1. Introduction	2
2. non-isotropic norm and estimates of linearized collision operators	7
2.1. Bounds on the non-isotropic norm	8
2.2. Equivalence to the linearized operator	16
2.3. Non-isotropic norms with different kinetic factors	20
3. Estimates of nonlinear collision operator in velocity space	23
3.1. Upper bounds in general case	23
3.2. A simple proof of Theorem 1.2 for $\gamma > -3/2$	29
3.3. Proof of Theorem 1.2	34
3.4. Estimation of commutators	39
4. Functional estimates in full space	41
4.1. Estimations without weight	41
4.2. Estimation with weight	43
4.3. Estimation with modified weight	44
4.4. Weighted coercivity of the linearized operator	46
5. Local existence	48
5.1. Classical solutions	49
5.2. L^2 -solutions	51
6. Global solutions	52
6.1. L^2 -solutions	52
6.2. Classical solutions	59
7. Appendix	64
References	66

2000 *Mathematics Subject Classification.* 35A05, 35B65, 35D10, 35H20, 76P05, 84C40.

Key words and phrases. Boltzmann equation, coercivity estimate, non-cutoff cross sections, global existence, non-isotropic norm, soft potential.

1. INTRODUCTION

This is the first part of a series of papers related to the inhomogeneous Boltzmann equation without angular cut-off, in the whole space and for general physical cross-sections. This global project is a natural continuation of our previous study [7] which was specialized to Maxwellian type cross sections.

In this part, we first establish an essential coercivity estimate of the linearized collision operator, in the framework of general cross sections. As shown in [7, 8] for the special Maxwellian case, this estimate will play an important role for the related Cauchy problem.

Based on this estimation, together with Part II [9], we will prove the global existence of classical non-negative solutions to the Boltzmann equation without angular cutoff, for the soft and hard potentials respectively, so that we are able to cover a general physical setting. Finally, in the paper [10], we will study the qualitative properties of solutions, that include full regularity, non-negativity, uniqueness and convergence rates to the equilibrium. This series of works establish a satisfactory theory on the well-posedness and full regularity of classical solutions.

In our presentation, we consider the problem in the physical case with dimension 3. However, our results hold true for any dimension bigger than 2.

Consider

$$(1.1) \quad f_t + v \cdot \nabla_x f = Q(f, f), \quad f|_{t=0} = f_0.$$

Here, $f = f(t, x, v)$ is the density distribution function of particles, having position $x \in \mathbb{R}^3$ and velocity $v \in \mathbb{R}^3$ at time t . The right hand side of (1.1) is the Boltzmann bilinear collision operator, which is given in the classical σ -representation by

$$Q(g, f) = \int_{\mathbb{R}^3} \int_{\mathbb{S}^2} B(v - v_*, \sigma) \{g'_* f' - g_* f\} d\sigma dv_*,$$

where $f'_* = f(t, x, v'_*)$, $f' = f(t, x, v')$, $f_* = f(t, x, v_*)$, $f = f(t, x, v)$, and for $\sigma \in \mathbb{S}^2$,

$$v' = \frac{v + v_*}{2} + \frac{|v - v_*|}{2} \sigma, \quad v'_* = \frac{v + v_*}{2} - \frac{|v - v_*|}{2} \sigma,$$

which gives the relation between the post and pre collisional velocities that follow from the conservation of momentum and kinetic energy.

For monatomic gas, the non-negative cross section $B(z, \sigma)$ depends only on $|z|$ and the scalar product $\frac{z}{|z|} \cdot \sigma$. As in our previous works, we assume that it takes the form

$$(1.2) \quad B(v - v_*, \cos \theta) = \Phi(|v - v_*|) b(\cos \theta), \quad \cos \theta = \frac{v - v_*}{|v - v_*|} \cdot \sigma, \quad 0 \leq \theta \leq \frac{\pi}{2},$$

in which it contains a kinetic factor given by

$$\Phi(|v - v_*|) = \Phi_\gamma(|v - v_*|) = |v - v_*|^\gamma,$$

and a factor related to the collision angle containing a singularity,

$$b(\cos \theta) \approx K \theta^{-2-2s} \quad \text{when } \theta \rightarrow 0+,$$

for some constant $K > 0$.

An important example of this singular cross section is the inverse power law potential ρ^{-r} with $r > 1$, ρ being the distance between two interacting particles, in which $s = \frac{1}{r} \in]0, 1[$ and $\gamma = 1 - 4s \in]-3, 1[$, cf. [12].

In the theory on the non-cutoff Boltzmann equation, the sign of $\gamma + 2s$ plays a crucial role. Hence, from now on, the case when $\gamma + 2s \leq 0$ is referred to the *non-cutoff soft*

potential, while the case $\gamma + 2s > 0$ to the *non-cutoff hard potential*. Note that this is different from the traditional classification on the index for the inverse power law.

In our present series of works, the well-posedness theory established applies to the general cross-sections with $\gamma > -3$ and $0 < s < 1$, that includes the inverse power law as a special example. Note that $\gamma > -3$ is needed for the Boltzmann operator to be well-posed, cf. [40].

Being concerned with a close to equilibrium framework, as in [7], the setting of the problem can be formulated as follows. First of all, without loss of generality, consider the perturbation around a normalized Maxwellian distribution

$$\mu(v) = (2\pi)^{-\frac{3}{2}} e^{-\frac{|v|^2}{2}},$$

by setting $f = \mu + \sqrt{\mu}g$. Since $Q(\mu, \mu) = 0$, we have

$$Q(\mu + \sqrt{\mu}g, \mu + \sqrt{\mu}g) = Q(\mu, \sqrt{\mu}g) + Q(\sqrt{\mu}g, \mu) + Q(\sqrt{\mu}g, \sqrt{\mu}g).$$

Denote

$$\Gamma(g, h) = \mu^{-1/2} Q(\sqrt{\mu}g, \sqrt{\mu}h).$$

Then the linearized Boltzmann operator takes the form

$$\mathcal{L}g = \mathcal{L}_1 g + \mathcal{L}_2 g = -\Gamma(\sqrt{\mu}, g) - \Gamma(g, \sqrt{\mu}).$$

Now the original problem (1.1) is reduced to the Cauchy problem for the perturbation g

$$(1.3) \quad \begin{cases} g_t + v \cdot \nabla_x g + \mathcal{L}g = \Gamma(g, g), & t > 0, \\ g|_{t=0} = g_0. \end{cases}$$

This close to equilibrium framework is classical for the Boltzmann equation with angular cutoff, but much less is known for the Boltzmann equation without angular cutoff, though the spectrum of the linearized operator without angular cut-off was analyzed a long time ago by Pao in [33].

However, since the late 1990s, the regularizing effect on the solution, produced by the singularity of the cross-section, has become reachable by rigorous analysis. Let us mention the systematic work on the entropy dissipation method initiated by Alexandre [1] and developed firstly by Lions [26], and then by many others, cf [3, 39, 40] and references therein. Since then, various works have been done on deriving the coercivity estimates in different settings and in different norms for different purposes. In particular, this kind of coercivity estimates has displayed some non-isotropic property in the very loose sense that, on one hand one gets a gain of the regularity in Sobolev norm of fractional order; and on the other hand, one also get a gain the moment to some fractional power in the velocity variable, cf. [2, 3, 5, 6, 7, 16, 21, 22, 24, 31, 32, 38, 39, 40] and references therein. Furthermore, these coercivity estimates have been proven to be very useful in getting the global existence and gain of full regularity in all variables for the Boltzmann equation without angular cutoff, as shown in our previous work [7]. For details about the recent progress in some of the directions mentioned previously, readers are referred to the survey paper by Alexandre, [2].

Since the coercivity estimate plays an important role in the study on the angular non-cutoff Boltzmann equation, such estimate in terms of the indices γ and s , has been pursued by many people. One of the purposes of this paper is to present a precise estimate that gives the essential properties of this singular behavior, that will be stated in the next theorem. Let us note that this result is proved in a general setting and it improves on previous results, such as those obtained in [5, 6, 7, 31, 32]. And this estimate will be used herein and in our

papers [9, 10] on the global existence in the hard potential case, and qualitative study of solutions.

To derive the desired coercivity estimate, we generalize the non-isotropic norm introduced in [7] as

$$\begin{aligned} \|g\|^2 &= \iiint \Phi(|v - v_*|) b(\cos \theta) \mu_* (g' - g)^2 \\ &\quad + \iiint \Phi(|v - v_*|) b(\cos \theta) g_*^2 (\sqrt{\mu'} - \sqrt{\mu})^2, \end{aligned}$$

where the integration is over $\mathbb{R}_v^3 \times \mathbb{R}_{v_*}^3 \times \mathbb{S}_\theta^2$. Note that it is a norm with respect to the velocity variable $v \in \mathbb{R}^3$ only. We can compare this non-isotropic norm with classical weighted Sobolev norms, see precisely Proposition 2.1.

The introduction of this norm was motivated by the study on the Landau equation which can be viewed as the grazing limit of the Boltzmann equation. It is known that for the Landau equation, see for example [19], that the essential norm in order to capture the dissipation of the linearized Landau operator can be defined just as the Dirichlet form of the linearized operator. By doing so, a norm can be well defined without loss of any dissipative information in the operator and this can be done directly for the Landau equation mainly because the corresponding Landau operator is a differential operator. However, for the Boltzmann equation without angular cutoff, the collision operator is a singular integral operator so that a direct analog is not obvious or feasible. Therefore, in the first part of this paper, we analyze the properties of the non-isotropic norm and obtain the precise coercivity estimate of the linearized collision operator. At this point, let us mention the different approach undertaken by Gressman-Strain [21, 22].

We shall use the following standard weighted Sobolev space defined, for $k, \ell \in \mathbb{R}$, as

$$H_\ell^k = H_\ell^k(\mathbb{R}_v^3) = \{f \in \mathcal{S}'(\mathbb{R}_v^3); W_\ell f \in H^k(\mathbb{R}_v^3)\}$$

and

$$H_\ell^k(\mathbb{R}_{x,v}^6) = \{f \in \mathcal{S}'(\mathbb{R}_{x,v}^6); W_\ell f \in H^k(\mathbb{R}_{x,v}^6)\}$$

where $W_\ell(v) = \langle v \rangle^\ell = (1 + |v|^2)^{\ell/2}$ is always the weight for v variables. Herein, $(\cdot, \cdot)_{L^2} = (\cdot, \cdot)_{L^2(\mathbb{R}_v^3)}$ denotes the usual scalar product in $L^2 = L^2(\mathbb{R}^3)$ for v variables. Recall that $L_\ell^2 = H_\ell^0$.

We shall use also in the following two different Sobolev spaces, one with x -derivatives only, another one with x, v derivatives and weight in the velocity variable v . For $k \in \mathbb{N}$, $\ell \in \mathbb{R}$, let

$$\begin{aligned} \mathcal{H}_\ell^k(\mathbb{R}^6) &= \left\{ f \in \mathcal{S}'(\mathbb{R}_{x,v}^6); \|f\|_{\mathcal{H}_\ell^k(\mathbb{R}^6)}^2 = \sum_{|\alpha|+|\beta| \leq N} \|W_{\ell-|\beta|} \partial_\beta^\alpha f\|_{L^2(\mathbb{R}^6)}^2 < +\infty \right\}, \\ \tilde{\mathcal{H}}_\ell^k(\mathbb{R}^6) &= \left\{ f \in \mathcal{S}'(\mathbb{R}_{x,v}^6); \|f\|_{\tilde{\mathcal{H}}_\ell^k(\mathbb{R}^6)}^2 = \sum_{|\alpha|+|\beta| \leq N} \|\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha f\|_{L^2(\mathbb{R}^6)}^2 < +\infty \right\}, \end{aligned}$$

where $\tilde{W}_\ell = (1 + |v|^2)^{|\alpha|+|\beta|/2}$.

We recall that the linearized operator \mathcal{L} has the following null space, which is spanned by the set of collision invariants:

$$\mathcal{N} = \text{Span} \left\{ \sqrt{\mu}, v_1 \sqrt{\mu}, v_2 \sqrt{\mu}, v_3 \sqrt{\mu}, |v|^2 \sqrt{\mu} \right\},$$

that is, $(\mathcal{L}g, g)_{L^2(\mathbb{R}_v^3)} = 0$ if and only if $g \in \mathcal{N}$.

Theorem 1.1. *Assume that the cross-section satisfies (1.2) with $0 < s < 1$ and $\gamma > -3$. Then there exist two generic constants $C_1, C_2 > 0$ such that for any suitable function g*

$$C_1 \|(\mathbf{I} - \mathbf{P})g\|^2 \leq (\mathcal{L}g, g)_{L^2} \leq C_2 \|g\|^2,$$

where \mathbf{P} is the L^2 -orthogonal projection onto the null space \mathcal{N} .

This coercivity estimate of the linearized collisional operator will prove to be crucial for the global existence of classical solutions to the Boltzmann equation. For this purpose, the analysis on the nonlinear operator is necessary, and we prove the following upper bound estimate.

Theorem 1.2. *For all $0 < s < 1$, assume that $\gamma > \max\{-3, -\frac{3}{2} - 2s\}$. Then, one has,*

$$\begin{aligned} |(\Gamma(f, g), h)_{L^2}| &\lesssim \left\{ \|f\|_{L^2_{s+\gamma/2}} \|g\| + \|g\|_{L^2_{s+\gamma/2}} \|f\| \right. \\ &\quad \left. + \min \{ \|f\|_{L^2} \|g\|_{L^2_{s+\gamma/2}}, \|f\|_{L^2_{s+\gamma/2}} \|g\|_{L^2} \} \|h\| \right\}, \end{aligned}$$

for suitable functions f, g, h .

We will then concentrate on the global existence of solutions, both weak and strong, for the *non-cutoff soft potential case* in the framework of small perturbation of an equilibrium state. Even though some estimates hold for the general case and will be used in the forthcoming papers, the condition $\gamma + 2s \leq 0$ will be imposed in the main existence results. In the Part II [9], we will then present the global existence theory for the hard potential case, that is, the condition $\gamma + 2s > 0$ imposed. Furthermore, the qualitative behavior of the solutions, such as the uniqueness, non-negativity, regularity and convergence rate to the equilibrium will be investigated in [10]. Note that both the global existence and the qualitative study on the solution behavior were firstly investigated in [7] for the Maxwellian molecule case where a generalized uncertainty principle obtained in [5] was used.

We begin with a local existence of classical solutions that holds true in general case.

Theorem 1.3. *Assume that the cross-section satisfies (1.2) with $\gamma + 2s \leq 0$, $0 < s < 1$ and $\gamma > -3$. Let $N \geq 6$ and $\ell \geq N$. For a small $\varepsilon > 0$, if $\|g_0\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \leq \varepsilon$, then there exists $T > 0$ such that the Cauchy problem (1.3) admits a solution*

$$g \in L^\infty([0, T]; \mathcal{H}_\ell^N(\mathbb{R}^6)).$$

Since we are interested in getting global existence results, the next statement deals with this issue asking only for control of x derivatives.

Theorem 1.4. *Assume that the cross-section satisfies (1.2) with $\gamma + 2s \leq 0$, $0 < s < 1$ and $\gamma > \max\{-3, -\frac{3}{2} - 2s\}$. Let $N \geq 3$. For a small $\varepsilon > 0$, if $\|g_0\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \leq \varepsilon$, then the Cauchy problem (1.3) admits a global solution*

$$g \in L^\infty([0, +\infty[; H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))).$$

The above global existence result is in a non-weighted function space without v derivatives in the framework of weak solutions. On the other hand, we will prove the following global existence result on classical solutions for which the proof is more involved. Note that for the qualitative study on the solution behavior, such as the regularity as will be shown in [10], solutions in a function space with x and v derivatives together with weight in v are needed. Hence, the next theorem also serves for this purpose.

Theorem 1.5. *Assume that the cross-section satisfies (1.2) with $\gamma + 2s \leq 0$, $0 < s < 1$ and $\gamma > \max\{-3, -\frac{3}{2} - 2s\}$. Let $N \geq 6$, $\ell \geq N$. For a small $\varepsilon > 0$, if $\|g_0\|_{\tilde{\mathcal{H}}_\ell^N(\mathbb{R}^6)} \leq \varepsilon$, then the Cauchy problem (1.3) admits a global solution*

$$g \in L^\infty([0, +\infty[; \tilde{\mathcal{H}}_\ell^N(\mathbb{R}^6)).$$

Let us now review some related works on this topic. First of all, the well-posedness theory for the Boltzmann equation has now been well established under the Grad's angular cutoff assumption. Under this assumption, there exist basically three frameworks of existence of global solutions. The first one was initiated by Grad [18] and firstly completed by Ukai [35, 36, 38] in the framework of weighted L_v^∞ function space for small perturbation of an equilibrium, where the spectrum analysis was used through a bootstrap argument. An important progress on the existence theory is the introduction of the renormalized solutions for large perturbation in the framework of L_v^1 function space by DiPerna-Lions [17, 25], where the velocity averaging lemma plays a key role. Recently, solutions in L_v^2 framework were established by macro-micro decompositions and energy method for small perturbation of an equilibrium, cf [19, 20, 27, 28].

However, without Grad's angular cutoff assumption, the established mathematical theories are far less. In this direction, the spectral analysis of the linearized collisional operator was studied by Pao [33]. In 1990's, some simplified models, such as Kac's model and the Boltzmann equation in lower dimensions with symmetry, were successfully studied, [13, 14, 15]. In 2000's, the mathematical theory for the spatially homogeneous Boltzmann equation was satisfactorily solved, [3, 4, 16, 24, 29, 30]. For the original Boltzmann equation in physical space, in the framework of renormalized solutions, the only existing result can be found in [11] where the basic existence result is still lacking. There are some local existence results, [1, 37], see also the reviews [2, 40].

Since 2006, we have been working on the original Boltzmann equation without angular cutoff, cf. [5, 6, 7]. Based on a new generalized uncertainty principle proved in [5], we developed a new approach for the regularity study. In the framework of small perturbation of an equilibrium in the whole space, the first complete global well-posedness theory and regularity were established for the Maxwellian molecule case [7]. As a continuation of these works, we successfully solve, in this series of papers, the fundamental problems, that is, existence, uniqueness, regularity, non-negativity and convergence rates of solutions, so that a complete and satisfactory mathematical theory is now established under some minimal regularity requirement on the initial data. Through this analysis, mathematical tools and techniques from harmonic analysis are used and some new ones are introduced, such as the generalized uncertainty principle and the above non-isotropic norm. Here we would like to mention that recently by using a different method, an existence result on solutions in the torus case was obtained in [21, 22, 23].

Finally, we present the main strategy of analysis in this paper. Based on the essential coercivity estimate on the linearized collisional operator and the non-isotropic norm proven in a first step, what is needed in this paper is then the detailed analysis on the collisional operator in both unweighted and weighted spaces, where its upper bounds, commutators with differentiation in v and commutators with weights in v are given. Through this analysis, we can see the role played by the parameters γ and s in the cross section. With these estimates, the energy method can be applied through the macro-micro decomposition analysis introduced in [19, 20]. Basically, the microscopic component of the solution is controlled by the essential coercivity estimate on the linearized collisional operator in the non-isotropic norm, while the dissipation on the macroscopic component is recovered by the system on the fluid functions through the macro-micro decomposition. Then the nonlinear terms are

essentially of higher order in the non-isotropic norm so that the energy estimate can be closed in the framework of small perturbation.

The rest of the paper is arranged as follows. In the next section, we extend the definition of the non-isotropic norm introduced in [7] and then state the main estimates in this paper. The proof of the upper and lower bound estimates of the non-isotropic norm will be given in Section 3. In Section 4, we will prove the equivalence of the Dirichlet form of the linearized collision operator and the square of the non-isotropic norm. The equivalence of the non-isotropic norms with different kinetic factors and different weights will be shown in Section 5. An upper bound estimate on the nonlinear collision operator which is useful for the well-posedness theory for the Boltzmann equation will be given next. However, because of unnecessary restrictions on the values of the parameter γ , we shall amplify such estimations and obtain some new functional estimates in v variable only in another Section. The functional estimates in both v and x variables are given in Section 4. With these estimates, the local and global existence of both weak and classical solutions are given in the last two sections respectively.

Notations: Herein, letters f, g, \dots stand for various suitable functions, while C, c, \dots stand for various numerical constants, independent from functions f, g, \dots and which may vary from line to line. Notation $A \lesssim B$ means that there exists a constant C such that $A \leq CB$, and similarly for $A \gtrsim B$. While $A \sim B$ means that there exist two generic constants $C_1, C_2 > 0$ such that

$$C_1 A \leq B \leq C_2 A.$$

2. NON-ISOTROPIC NORM AND ESTIMATES OF LINEARIZED COLLISION OPERATORS

For later use, we will need to compare the original cross-section with the situation when its kinetic part is mollified. That is, for the function $\Phi(z)$ appearing in the cross-section, we denote by $\tilde{\Phi}(z) = (1 + |z|^2)^{\frac{\gamma}{2}}$ its smoothed version. To show the dependence of the estimates on the mollified or non-mollified kinetic factor in the cross-section, we will use the notations $Q^{\tilde{\Phi}}$ and Q^{Φ} to denote the Boltzmann collisional operator when the kinetic part is $\tilde{\Phi}$ and Φ respectively. In particular, $Q = Q^{\Phi}$. This upper-script will be also used for other operators as well.

First of all, let us recall that

$$(\mathcal{L}g, g)_{L^2} = -(\Gamma(\sqrt{\mu}, g) + \Gamma(g, \sqrt{\mu}), g)_{L^2} \geq 0,$$

and the definition of the non-isotropic norm

$$(2.1) \quad \begin{aligned} \|g\|^2 &= \iiint \Phi(|v - v_*|) b(\cos \theta) \mu_* (g' - g)^2 \\ &\quad + \iiint \Phi(|v - v_*|) b(\cos \theta) g_*^2 (\sqrt{\mu'} - \sqrt{\mu})^2 \\ &= J_1 + J_2, \end{aligned}$$

where the integration is over $\mathbb{R}_v^3 \times \mathbb{R}_{v_*}^3 \times \mathbb{S}_\sigma^2$.

The following proposition gives a precise version of Theorem 1.1.

Proposition 2.1. *Assume that the cross-section satisfies (1.2) with $0 < s < 1$ and $\gamma > -3$. Then there exist two generic constants $C_1, C_2 > 0$ such that*

$$C_1 \|(\mathbf{I} - \mathbf{P})g\|^2 \leq (\mathcal{L}g, g)_{L^2} \leq 2(\mathcal{L}_1 g, g)_{L^2} \leq C_2 \|g\|^2$$

for any suitable function g .

Concerning the lower and upper bounds of the non-isotropic norm we have

Proposition 2.2. *Assume that the cross-section satisfies (1.2) with $0 < s < 1$ and $\gamma > -3$. Then there exist two generic constants $C_1, C_2 > 0$ such that*

$$C_1 \left\{ \|g\|_{H_{\gamma/2}^s}^2 + \|g\|_{L_{s+\gamma/2}^2}^2 \right\} \leq \|g\|^2 \leq C_2 \|g\|_{H_{s+\gamma/2}^s}^2$$

for any suitable function g .

From this estimate and Theorem 1.1, we can get the following estimate in classical weighted Sobolev spaces

$$(2.2) \quad C_1 \left\{ \|(\mathbf{I} - \mathbf{P})g\|_{H_{\gamma/2}^s}^2 + \|(\mathbf{I} - \mathbf{P})g\|_{L_{s+\gamma/2}^2}^2 \right\} \leq (\mathcal{L}g, g)_{L^2} \leq C_2 \|g\|_{H_{s+\gamma/2}^s}^2.$$

In the following, we will use the lower script Φ on the non-isotropic norm, and so use the notation $\|g\|_{\Phi}$ if we need to specify its dependence on the kinetic factor Φ . Notations J_1^{Φ}, J_2^{Φ} will be also used for the same purpose.

Part of the proof on the lower bound of the non-isotropic norm given in Proposition 2.2 is essentially due to the following equivalence relations.

Proposition 2.3. *Assume that the cross-section satisfies (1.2) with $0 < s < 1$ and $\gamma > -3$. Then we have*

$$\|g\|_{\Phi} \sim \|g\|_{\tilde{\Phi}}.$$

Concerning the dependence on the index γ in $\Phi_{\gamma} = |v - v_*|^{\gamma}$, we have

Proposition 2.4. *Assume that the cross-section satisfies (1.2) with $0 < s < 1$ and $\gamma > -3$. Then for any $\beta > -3$, we have*

$$\|g\|_{\Phi_{\gamma}} \sim \| \langle v \rangle^{(\gamma-\beta)/2} g \|_{\Phi_{\beta}}.$$

2.1. Bounds on the non-isotropic norm. This section is devoted to the proof of Proposition 2.2. We will often use the following elementary estimate stated in velocity dimension n , since it will be needed for both cases $n = 2$ and $n = 3$.

Lemma 2.5. *Let the velocity dimension be n , $n \in \mathbb{N}$, $\rho > 0, \delta \in \mathbb{R}$ and let $\mu_{\rho,\delta}(u) = \langle u \rangle^{\delta} e^{-\rho|u|^2}$ for $u \in \mathbb{R}^n$. If $\alpha > -n$ and $\beta \in \mathbb{R}$, then we have*

$$(2.3) \quad I_{\alpha,\beta}(u) = \int_{\mathbb{R}^n} |w|^{\alpha} \langle w \rangle^{\beta} \mu_{\rho,\delta}(w+u) dw \sim \langle u \rangle^{\alpha+\beta}.$$

Proof. Since we have

$$\langle u \rangle^{\beta} \langle u+w \rangle^{-|\beta|} \leq \langle w \rangle^{\beta} \leq \langle u \rangle^{\beta} \langle u+w \rangle^{|\beta|},$$

it suffices to show (2.3) with $\beta = 0$, by taking $\mu_{\rho,\delta \pm |\beta|}$ instead of $\mu_{\rho,\delta}$. Taking into account the fact that $\alpha > -n$, this estimate is obvious when $|u| \leq 1$. If $|u| \geq 1$, then we have

$$I_{\alpha,0}(u) \geq 4^{-|\alpha|} \langle u \rangle^{\alpha} \int_{\{|u+w| \leq 1/2\}} \mu_{\rho,\delta}(u+w) dw \gtrsim \langle u \rangle^{\alpha},$$

because $|u+w| \leq 1/2$ implies that $4^{-1} \langle u \rangle \leq |w| \leq 4 \langle u \rangle$. Noticing that $2|w| \geq \langle w \rangle$ if $|w| \geq 1$, we have

$$\begin{aligned} I_{\alpha,0}(u) &\leq \left(\max_{|w| \leq 1} \mu_{\rho,\delta}(u+w) \right) \int_{\{|w| \leq 1\}} |w|^{\alpha} dw + 2^{|\alpha|} \int_{\{|w| \geq 1\}} \langle w \rangle^{\alpha} \mu_{\rho,\delta}(u+w) dw \\ &\lesssim \left(\langle u \rangle^{|\delta|} e^{-\rho|u|^2/2} + \langle u \rangle^{\alpha} \int_{\mathbb{R}^n} \langle u+w \rangle^{\alpha} \mu_{\rho,\delta}(u+w) dw \right) \lesssim \langle u \rangle^{\alpha}. \end{aligned}$$

And this completes the proof of the lemma. \square

Recall from (2.1) that the non-isotropic norm contains two parts, denoted by J_1 and J_2 respectively. The estimation on each part will be given in the following subsections. We start with the estimation on J_2 because the analysis is easier.

Let us start with the following upper bound on J_2 .

Lemma 2.6. *Under the same assumptions as in Theorem 1.1, we have*

$$J_2 := \iiint b(\cos \theta) \Phi(|v - v_*|) g_*^2 (\sqrt{\mu'} - \sqrt{\mu})^2 dv dv_* d\sigma \lesssim \|g\|_{L^2_{s+\gamma/2}}^2.$$

Proof. Note that

$$\begin{aligned} J_2 &\leq 2 \iiint b \Phi(|v - v_*|) g_*^2 (\mu'^{1/4} - \mu^{1/4})^2 (\mu'^{1/2} + \mu^{1/2}) dv dv_* d\sigma \\ &\leq \iiint b |v' - v_*|^\gamma g_*^2 (\mu'^{1/4} - \mu^{1/4})^2 \mu'^{1/2} dv dv_* d\sigma \\ &\quad + \iiint b |v - v_*|^\gamma g_*^2 (\mu'^{1/4} - \mu^{1/4})^2 \mu^{1/2} dv dv_* d\sigma \\ &= F_1 + F_2. \end{aligned}$$

By the regular change of variables $v \rightarrow v'$, we have

$$\begin{aligned} F_1 &\lesssim \iint |v' - v_*|^\gamma \left(\int b(\cos \theta) \min(|v' - v_*|^2 \theta^2, 1) d\sigma \right) g_*^2 \mu'^{1/2} dv' dv_* \\ &\lesssim \int \left(\int |v' - v_*|^{\gamma+2s} \mu' dv' \right) g_*^2 dv_* \lesssim \|g\|_{L^2_{s+\gamma/2}}^2, \end{aligned}$$

where we have used Lemma 2.5 in the case $n = 3$ to get the last inequality. A direct estimation show that the same bound holds true for F_2 . And this completes the proof of the lemma. \square

Remark 2.7. *Note that the above lemma holds even if Φ is replaced by $\tilde{\Phi}$ by using Lemma 2.5.*

We now turn to the lower bound for J_2 .

Lemma 2.8. *Under the assumptions (1.2) with $\gamma > -3$, there exists a constant $C > 0$ such that*

$$J_2 := \iiint b(\cos \theta) \Phi(|v - v_*|) g_*^2 (\sqrt{\mu'} - \sqrt{\mu})^2 dv dv_* d\sigma \geq C \|g\|_{L^2_{s+\gamma/2}}^2.$$

Proof. We will apply the argument used in [39]. By shifting to the ω -representation,

$$v' = v - ((v - v_*) \cdot \omega) \omega \quad v_*' = v + ((v - v_*) \cdot \omega) \omega, \quad \omega \in \mathbb{S}^2,$$

in view of the change of variables $(v, v_*) \rightarrow (v_*, v)$, we get,

$$J_2 = 4 \iiint b(\cos \theta) \sin(\theta/2) \Phi(|v - v_*|) g_*^2 (\sqrt{\mu'_*} - \sqrt{\mu_*})^2 dv dv_* d\omega,$$

because $d\sigma = 4 \sin(\theta/2) d\omega$. Then, we use the Carleman representation. The idea of this representation is to replace the set of variables (v, v_*, ω) by the set (v, v', v'_*) . Here, $v, v' \in \mathbb{R}^3$ and $v'_* \in E_{vv'}$, where $E_{vv'}$ is the hyperplane passing through v and orthogonal to $v - v'$. By using the formula

$$dv_* d\omega = \frac{dv'_* dv'}{|v - v'|^2},$$

cf. page 347 of [39], and by taking the change of variables

$$(v, v', v'_*) \rightarrow (v, v + h, v + y),$$

with $h \in \mathbb{R}^3$ and $y \in E_h$, where E_h is the hyperplane orthogonal to h passing through the origin in \mathbb{R}^3 , we have

$$\begin{aligned} J_2 &\sim \int_{\mathbb{R}^3} \int_{\mathbb{R}_h^3} \int_{y \in E_h \cap \{|y| \geq |h|\}} \frac{|y|^{1+2s+\gamma}}{|h|^{1+2s}} g(v)^2 \\ &\quad \times (\sqrt{\mu(v+y)} - \sqrt{\mu(v+y+h)})^2 dv \frac{dh dy}{|h|^2}, \end{aligned}$$

because

$$\begin{aligned} |h| &= |v' - v| = |v'_* - v| \tan \frac{\theta}{2} = |y| \tan \frac{\theta}{2}, \quad \theta \in [0, \pi/2], \\ b(\cos \theta) \sin(\theta/2) \Phi(|v - v'_*|) &\sim \frac{|v'_* - v|^{1+2s+\gamma}}{|v - v'|^{1+2s}} \mathbf{1}_{\{|v'_* - v| \geq |v' - v|\}}. \end{aligned}$$

We decompose $v = v_1 + v_2$, where v_2 is the orthogonal projection of v on E_h . Since μ is invariant by rotation, we may assume $v = (0, 0, |v|)$ without loss of generality. By introducing the polar coordinates

$$h = (\rho \sin \vartheta \cos \phi, \rho \sin \vartheta \sin \phi, \rho \cos \vartheta), \quad \vartheta \in [0, \pi], \quad \phi \in [0, 2\pi], \quad \rho > 0,$$

we obtain $|v_1| = |v| |\cos \vartheta|$, $|v_1 + h| = ||v| \cos \vartheta + \rho|$ and $|v_2| = |v| \sin \vartheta$. Note that if $0 < \vartheta \leq \pi/2$, then

$$\begin{aligned} (\sqrt{\mu(v+y)} - \sqrt{\mu(v+y+h)})^2 &= \mu(v_2+y) (\sqrt{\mu(v_1)} - \sqrt{\mu(v_1+h)})^2 \\ &\geq \mu(v_2+y) \mu(v_1) (1 - e^{-\rho^2/4})^2 / (2\pi)^{3/2}. \end{aligned}$$

Therefore, we have for any $\delta > 0$

$$\begin{aligned} J_2 &\geq C \int_{\mathbb{R}_v^3} g(v)^2 \left\{ \int_{\mathbb{R}_h^3} \frac{(\sqrt{\mu(v_1)} - \sqrt{\mu(v_1+h)})^2}{|h|^{3+2s}} \right. \\ &\quad \times \left. \left(\int_{y \in E_h \cap \{|y| \geq |h|\}} |y|^{1+2s+\gamma} \mu(v_2+y) dy \right) dh \right\} dv \\ &\geq C \int_{\mathbb{R}_v^3} g(v)^2 \left\{ \int_{\pi/2-1/\langle v \rangle}^{\pi/2} \mu(v_1) \left(\int_0^\delta \frac{(1 - e^{-\rho^2/4})^2}{\rho^{1+2s}} \right. \right. \\ &\quad \times \left. \left(\int_{y \in E_h} |y|^{1+2s+\gamma} \mu(v_2+y) dy \right. \right. \\ &\quad \left. \left. - \int_{y \in E_h \cap \{|y| \leq \rho\}} |y|^{1+2s+\gamma} \mu(v_2+y) dy \right) d\rho \right\} \sin \vartheta d\vartheta \right\} dv. \end{aligned}$$

Since we have

$$\int_{y \in E_h \cap \{|y| \leq \rho\}} |y|^{1+2s+\gamma} \mu(v_2+y) dy \leq \delta^{2s} \int_{y \in E_h} |y|^{1+\gamma} \mu(v_2+y) dy, \quad \text{if } \rho \leq \delta,$$

and it follows from Lemma 2.5 in the case $n = 2$, that

$$\int_{y \in E_h} |y|^\beta \mu(v_2+y) dy \sim \langle v_2 \rangle^\beta \quad \text{if } \beta > -2,$$

there exist two constants $C_1, C_2 > 0$ such that if $\rho \leq \delta$, we have

$$\begin{aligned} \int_{y \in E_h} |y|^{1+2s+\gamma} \mu(v_2 + y) dy - \int_{y \in E_h \cap \{|y| \leq \rho\}} |y|^{1+2s+\gamma} \mu(v_2 + y) dy \\ \geq C_1 \langle v_2 \rangle^{1+2s+\gamma} - C_2 \delta^{2s} \langle v_2 \rangle^{1+\gamma}. \end{aligned}$$

Taking a sufficiently small $\delta > 0$ gives

$$\begin{aligned} J_2 &\geq C \int_{\mathbb{R}_v^3} g(v)^2 \left\{ \int_{\pi/2-1/\langle v \rangle}^{\pi/2} \mu(v_1) \right. \\ &\quad \left. \times \left(\int_0^\delta \frac{(1 - e^{-\rho^2/4})^2}{\rho^{1+2s}} d\rho \right) \langle v_2 \rangle^{1+2s+\gamma} \sin \vartheta d\vartheta \right\} dv \\ &\geq C_\delta \int_{\mathbb{R}_v^3} \langle v \rangle^{2s+\gamma} g(v)^2 \left\{ \int_{\pi/2-1/\langle v \rangle}^{\pi/2} e^{-|v|^2 \cos^2 \vartheta} \langle v \rangle d\vartheta \right\} dv \\ &\geq C_\delta \|g\|_{L_{s+\gamma/2}^2}^2. \end{aligned}$$

The proof of the lemma is now completed. \square

Remark 2.9. In the above proof, the factor $|y|^\gamma$ can be replaced by $\langle y \rangle^\gamma$, so that Lemma 2.8 is valid even if Φ is replaced by $\tilde{\Phi} = \langle v - v_* \rangle^\gamma$. By the above lemma together with Lemma 2.6 and the Remark after it, we can conclude

$$(2.4) \quad J_2^\Phi \sim \|g\|_{L_{s+\gamma/2}^2}^2 \sim J_2^{\tilde{\Phi}}.$$

We now turn to the estimation of the first term of the non-isotropic norm, that is, J_1 . We will firstly show that the singular behavior of the cross-section when $v = v_*$ can be smoothed out. This point is given by the following proposition.

Proposition 2.10. Under the same assumption as in Theorem 1.1, we have

$$J_1^\Phi + \|g\|_{L_{s+\gamma/2}^2}^2 \sim J_1^{\tilde{\Phi}} + \|g\|_{L_{s+\gamma/2}^2}^2.$$

Remark 2.11. This proposition is nothing but Proposition 2.3 by Remark 2.9.

Proof. By using similar arguments as in the proof of Lemma 2.8, it follows from the Carleman representation that

$$\begin{aligned} J_1^\Phi &\sim \int_{\mathbb{R}_v^3} \int_{\mathbb{R}_h^3} \int_{y \in E_h \cap \{|y| \geq |h|\}} \frac{|y|^{1+2s+\gamma}}{|h|^{1+2s}} \mu(v) (g(v+y) - g(v+y+h))^2 dv \frac{dh dy}{|h|^2} \\ &= \int_{\mathbb{R}_v^3} \int_{\mathbb{R}_h^3} \int_{y \in E_h \cap \{|y| \geq |h|\}} \frac{|y|^{1+2s+\gamma}}{|h|^{1+2s}} \mu(v+y) (g(v) - g(v+h))^2 dv \frac{dh dy}{|h|^2}, \end{aligned}$$

where the last equality is a direct consequence of the change of variables $(v+y, y) \rightarrow (v, -y)$.

Similarly, we have

$$J_1^{\tilde{\Phi}} \sim \int_{\mathbb{R}_v^3} \int_{\mathbb{R}_h^3} \int_{y \in E_h \cap \{|y| \geq |h|\}} \frac{|y|^{1+2s} \langle y \rangle^\gamma}{|h|^{1+2s}} \mu(v+y) (g(v+h) - g(v))^2 dv \frac{dh dy}{|h|^2}.$$

We claim that

$$(2.5) \quad \int_{\mathbb{R}_v^3} \int_{\mathbb{R}_h^3} \int_{y \in E_h \cap \{|y| \leq |h|\}} \frac{|y|^{1+2s+\gamma}}{|h|^{1+2s}} \mu(v+y) (g(v+h) - g(v))^2 dv \frac{dh dy}{|h|^2} \\ \lesssim \|g\|_{L_{s+\gamma/2}^2}^2,$$

$$(2.6) \quad \int_{\mathbb{R}_v^3} \int_{\mathbb{R}_h^3} \int_{y \in E_h \cap \{|y| \leq |h|\}} \frac{|y|^{1+2s} \langle y \rangle^\gamma}{|h|^{1+2s}} \mu(v+y) (g(v+h) - g(v))^2 dv \frac{dh dy}{|h|^2} \\ \lesssim \|g\|_{L_{s+\gamma/2}^2}^2.$$

Note carefully that the integration in these estimates is performed for "large" values of h .

Once we admit those estimates, to conclude the proof of the lemma, it suffices to show that

$$G(v, h) = \int_{y \in E_h} |y|^{1+2s+\gamma} \mu(v+y) dy \sim \int_{y \in E_h} |y|^{1+2s} \langle y \rangle^\gamma \mu(v+y) dy = \tilde{G}(v, h).$$

We decompose $v = v_1 + v_2$, where v_2 is the orthogonal projection of v on E_h . Then we have $\mu(v+y) = \mu(v_1) \mu(v_2+y)$, whence it follows from Lemma 2.5 together with $1+2s+\gamma > -2$ that

$$G(v, h) \sim \mu(v_1) \langle v_2 \rangle^{1+2s+\gamma} \sim \tilde{G}(v, h).$$

It remains to show (2.5) and (2.6). We write

$$\int_{\mathbb{R}_v^3} \int_{\mathbb{R}_h^3} \int_{y \in E_h \cap \{|y| \leq |h|\}} \frac{|y|^{1+2s+\gamma}}{|h|^{1+2s}} \mu(v+y) (g(v+h) - g(v))^2 dv \frac{dh dy}{|h|^2} \\ = \int_{\mathbb{R}_v^3} \int_{\{|h| \leq 1\}} \int_{y \in E_h \cap \{|y| \leq |h|\}} + \int_{\mathbb{R}_v^3} \int_{\{|h| \geq 1\}} \int_{y \in E_h \cap \{|y| \leq |h|\}} = A_1 + A_2.$$

Take a small $\delta > 0$ such that $\gamma - \delta > -3$. Then, in view of $1 + \gamma - \delta > -2$, we have

$$A_1 \leq \int_{\mathbb{R}_v^3} \int_{\{|h| \leq 1\}} \int_{y \in E_h} \frac{|y|^{1+\gamma-\delta}}{|h|^{1-\delta}} \mu(v+y) (g(v+h) - g(v))^2 dv \frac{dh dy}{|h|^2} \\ = \int_{\mathbb{R}_v^3} \mu(v_1) \int_{\{|h| \leq 1\}} \left(\int_{y \in E_h} |y|^{1+\gamma-\delta} \mu(v_2+y) dy \right) (g(v+h) - g(v))^2 \frac{dh}{|h|^{3-\delta}} dv \\ \lesssim \int_{\mathbb{R}_v^3} \mu(v_1) \langle v_2 \rangle^{1+\gamma-\delta} \int_{\{|h| \leq 1\}} (g(v+h) - g(v))^2 \frac{dh}{|h|^{3-\delta}} dv \\ \lesssim \int_{\mathbb{R}_v^3} \int_{\{|h| \leq 1\}} (\mu(v_1 - h) + \mu(v_1)) \langle v_2 \rangle^{1+\gamma-\delta} |g(v)|^2 \frac{dh}{|h|^{3-\delta}} dv,$$

where we have used the change of variables $v+h \rightarrow v$ for the factor $g(v+h)$. As in the proof of Lemma 2.8, by assuming $v = (0, 0, |v|)$, we introduce the polar coordinates

$$h = (\rho \sin \vartheta \cos \phi, \rho \sin \vartheta \sin \phi, \rho \cos \vartheta), \quad \vartheta \in [0, \pi], \quad \phi \in [0, 2\pi], \quad \rho > 0.$$

Since $|v_1| = |v| \cos \vartheta$, $|v_1 - h| = ||v| \cos \vartheta - \rho|$ and $|v_2| = |v| \sin \vartheta$, by using the change of variable $|v| \cos \vartheta = r$, we obtain

$$A_1 \lesssim \int_{\mathbb{R}_v^3} |g(v)|^2 \int_0^1 \frac{1}{\rho^{1-\delta}} \\ \times \left(\int_{-|v|}^{|v|} \frac{(1 + |v|^2 - r^2)^{(1+\gamma-\delta)/2}}{|v|} (e^{-|r-\rho|^2/2} + e^{-r^2/2}) dr \right) d\rho dv.$$

Similarly, if $1 + 2s - \delta > 1$, then we have

$$\begin{aligned} A_2 &\leq \int_{\mathbb{R}_v^3} \int_{\{|h|\geq 1\}} \int_{y \in E_h} \frac{|y|^{1+\gamma+2s-\delta}}{|h|^{1+2s-\delta}} \mu(v+y) (g(v+h) - g(v))^2 dv \frac{dh dy}{|h|^2} \\ &\lesssim \int_{\mathbb{R}_v^3} |g(v)|^2 \int_1^\infty \frac{1}{\rho^{1+2s-\delta}} \\ &\quad \times \left(\int_{-|v|}^{|v|} \frac{(1+|v|^2-r^2)^{(1+\gamma+2s-\delta)/2}}{|v|} (e^{-|r-\rho|^2/2} + e^{-r^2/2}) dr \right) d\rho dv. \end{aligned}$$

If $1 + \gamma + 2s - \delta \geq 0$, then

$$\begin{aligned} K(v, \rho) &= \int_{-|v|}^{|v|} \frac{(1+|v|^2-r^2)^{(1+\gamma+2s-\delta)/2}}{|v|} (e^{-|r-\rho|^2/2} + e^{-r^2/2}) dr \\ &\leq \langle v \rangle^{(\gamma+2s-\delta)/2} \int_{-|v|}^{|v|} (e^{-|r-\rho|^2/2} + e^{-r^2/2}) dr \lesssim \langle v \rangle^{\gamma+2s}, \end{aligned}$$

which shows

$$(2.7) \quad A_2 \lesssim \int_{\mathbb{R}_v^3} |g(v)|^2 \int_1^\infty \frac{K(v, \rho)}{\rho^{1+2s-\delta}} d\rho dv \lesssim \int \langle v \rangle^{\gamma+2s} |g(v)|^2 dv.$$

On the other hand, if $1 + \gamma + 2s - \delta < 0$ and $|v| \geq 3$, then

$$\begin{aligned} K(v, \rho) &\lesssim \int_0^{|v|} \frac{(|v|^2-r^2)^{(1+\gamma+2s-\delta)/2}}{|v|} (e^{-|r-\rho|^2/2} + 3e^{-r^2/2}) dr \\ &\lesssim |v|^{(-1+\gamma+2s-\delta)/2} \int_0^{|v|} (|v|-r)^{(1+\gamma+2s-\delta)/2} (e^{-|r-\rho|^2/2} + 3e^{-r^2/2}) dr \\ &\lesssim \langle v \rangle^{\gamma+2s} + |v|^{(-1+\gamma+2s-\delta)/2} \int_{|v|/2}^{|v|} (|v|-r)^{(1+\gamma+2s-\delta)/2} 3e^{-|r-\rho|^2/2} dr, \end{aligned}$$

because

$$\begin{aligned} \int_0^{|v|} (|v|-r)^{(1+\gamma+2s-\delta)/2} e^{-|r|^2/2} dr &\lesssim |v|^{(1+\gamma+2s-\delta)/2} \int_0^{|v|/2} e^{-|r|^2/2} dr \\ &\quad + e^{-|v|^2/8} \int_{|v|/2}^{|v|} (|v|-r)^{(1+\gamma+2s-\delta)/2} dr, \end{aligned}$$

where we have used that $(1 + \gamma + 2s - \delta)/2 > -1$ for small $\delta > 0$ that follows from the assumption $\gamma > -3$. We now consider

$$\begin{aligned} &\int_1^\infty \frac{d\rho}{\rho^{1+2s-\delta}} \int_{|v|/2}^{|v|} (|v|-r)^{(1+\gamma+2s-\delta)/2} e^{-|r-\rho|^2/2} dr \\ &\leq \int_{|v|/2}^{|v|} (|v|-r)^{(1+\gamma+2s-\delta)/2} \left(\int_{\{|r-\rho|\leq \sqrt{2\log|v|}\}} (|v|/3)^{-(1+2s-\delta)} d\rho \right) dr \\ &\quad + \int_{|v|/2}^{|v|} (|v|-r)^{(1+\gamma+2s-\delta)/2} \left(\int_{\{|r-\rho|\geq \sqrt{2\log|v|}\}} \frac{|v|^{-1} d\rho}{\rho^{1+2s-\delta}} \right) dr \\ &\lesssim (|v|^{(1+\gamma+2s-\delta)/2} \sqrt{2\log|v|} + |v|^{(1+\gamma+2s-\delta)/2}) \lesssim \langle v \rangle^{(1+\gamma+2s)/2}. \end{aligned}$$

Therefore, in the case when $1 + \gamma + 2s - \delta < 0$, we also have (2.7). Similarly, we have

$$A_1 \lesssim \int_{\mathbb{R}_+^3} |g(v)|^2 \int_0^1 \frac{K(v, \rho)}{\rho^{1-\delta}} d\rho dv \lesssim \int \langle v \rangle^{\gamma+2s} |g(v)|^2 dv,$$

which shows (2.5). The proof of (2.6) is similar, and thus the proof of the proposition is completed. \square

Lemma 2.12. *There exist constants $C_1, C_2 > 0$ such that*

$$(2.8) \quad J_1 \geq C_1 \|\langle v \rangle^{\gamma/2} g\|_{H^s}^2 - C_2 \|g\|_{L_{s+\gamma/2}^2}^2.$$

The same conclusion holds even with μ replaced by μ^ρ for any fixed $\rho > 0$.

Proof. It follows from Proposition 2.10 that

$$(2.9) \quad \begin{aligned} C(J_1^\Phi + \|g\|_{L_{s+\gamma/2}^2}^2) &\geq 2J_1^{\tilde{\Phi}} \\ &\geq \iiint b(\cos \theta) \frac{\mu_*}{\langle v_* \rangle^{|\gamma|}} (\langle v' \rangle^{\gamma/2} g' - \langle v \rangle^{\gamma/2} g)^2 d\sigma dv dv_* \\ &\quad - 2 \iiint b(\cos \theta) \frac{\mu_*}{\langle v_* \rangle^{|\gamma|}} (\langle v \rangle^{\gamma/2} - \langle v' \rangle^{\gamma/2})^2 |g|^2 d\sigma dv dv_* \\ &= A_1 + A_2, \end{aligned}$$

because $\tilde{\Phi}(|v - v_*|) \sim \langle v' - v_* \rangle^\gamma \geq \frac{\langle v' \rangle^\gamma}{\langle v_* \rangle^{|\gamma|}}$ and $2(a+b)^2 \geq a^2 - 2b^2$. Setting $\tilde{g} = \langle v \rangle^{\gamma/2} g$ and noting $C_\gamma \mu(v) \langle v \rangle^{-|\gamma|} \geq \mu(2v)$ for a $C_\gamma > 0$, as in Proposition 1 of [3], we have

$$\begin{aligned} C_\gamma A_1 &\geq \int_{\mathbb{R}^6} \int_{\mathbb{S}^2} b\left(\frac{v - v_*}{|v - v_*|} \cdot \sigma\right) \mu(2v_*) (\tilde{g}(v) - \tilde{g}(v'))^2 d\sigma dv_* dv \\ &= (4\pi)^{-3} \int_{\mathbb{R}^3} \int_{\mathbb{S}^2} b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) \left\{ \hat{\mu}(0) |\hat{\tilde{g}}(\xi)|^2 + \hat{\mu}(0) |\hat{\tilde{g}}(\xi^+)|^2 \right. \\ &\quad \left. - 2 \operatorname{Re} \hat{\mu}(\xi^-/2) \hat{\tilde{g}}(\xi^+) \bar{\hat{\tilde{g}}}(\xi) \right\} d\sigma d\xi \\ &\geq \frac{1}{2(4\pi)^3} \int_{\mathbb{R}^3} |\hat{\tilde{g}}(\xi)|^2 \left\{ \int_{\mathbb{S}^2} b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) (\hat{\mu}(0) - |\hat{\mu}(\xi^-/2)|) d\sigma \right\} d\xi. \end{aligned}$$

Since we have $\hat{\mu}(0) - |\hat{\mu}(\xi^-/2)| = c(1 - e^{-|\xi|^{-1/8}}) \geq c'|\xi|^{-2}$ if $|\xi| \leq 1$, in view of $|\xi|^{-2} = |\xi|^2 \sin^2 \theta/2 \geq |\xi|^2 (\theta/\pi)^2$, we obtain for $|\xi| \geq 1$

$$\begin{aligned} \int_{\mathbb{S}^2} b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) (\hat{\mu}(0) - |\hat{\mu}(\xi^-/2)|) d\sigma &\geq \int_{|\xi|(\theta/\pi) \leq 1} \sin \theta b(\cos \theta) |\xi|^2 (\theta/\pi)^2 d\theta \\ &\geq c'' K |\xi|^2 \int_0^{1/|\xi|} \theta^{-1-2s} \theta^2 d\theta \\ &= c'' K |\xi|^2 |\xi|^{2s-2} / (2-2s). \end{aligned}$$

Therefore, we have

$$(2.10) \quad \begin{aligned} A_1 &\geq C_1 \int_{|\xi| \geq 1} |\xi|^{2s} |\hat{\tilde{g}}(\xi)|^2 d\xi \geq C_1 2^{-2s} \int_{|\xi| \geq 1} (1 + |\xi|^2)^s |\hat{\tilde{g}}(\xi)|^2 d\xi \\ &\geq C_1 2^{-2s} \|\langle v \rangle^{\gamma/2} g\|_{H^s(\mathbb{R}_+^3)}^2 - C_1 \|\langle v \rangle^{\gamma/2} g\|_{L^2(\mathbb{R}_+^3)}^2. \end{aligned}$$

As for A_2 , we note that if $v_\tau = v' + \tau(v - v')$ for $\tau \in [0, 1]$, then

$$\langle v \rangle \leq \langle v - v_* \rangle + \langle v_* \rangle \leq \sqrt{2} \langle v_\tau - v_* \rangle + \langle v_* \rangle \leq (1 + \sqrt{2}) \langle v_\tau \rangle \langle v_* \rangle,$$

and $\langle v_\tau \rangle \leq (1 + \sqrt{2})\langle v \rangle \langle v_* \rangle$, which show $\langle v_\tau \rangle^\beta \leq C_\beta \langle v \rangle^\beta \langle v_* \rangle^{|\beta|}$ for any $\beta \in \mathbb{R}$. It follows that

$$\begin{aligned} \left| \langle v \rangle^{\gamma/2} - \langle v' \rangle^{\gamma/2} \right| &\leq C_\gamma \int_0^1 \langle v' + \tau(v - v') \rangle^{(\gamma/2-1)} d\tau |v - v_*| \theta \\ &\leq C'_\gamma \left(\langle v \rangle^{(\gamma/2-1)} \langle v_* \rangle^{|\gamma/2-1|} \right) \langle v - v_* \rangle \theta, \end{aligned}$$

and thus we have

$$\begin{aligned} A_2 &\leq C \iint \frac{\mu_*}{\langle v_* \rangle^{|\gamma|}} |g|^2 \left\{ \left(\langle v \rangle^{(\gamma-2)} \langle v_* \rangle^{|\gamma-2|} \right)^2 \left(\int_0^{\langle v-v_* \rangle^{-1}} \theta^{-1-2s} \langle v - v_* \rangle \theta \right)^2 d\theta \right. \\ &\quad \left. + \int_{\langle v-v_* \rangle^{-1}}^{\pi/2} \left(\langle v \rangle^\gamma + \langle v \rangle^\gamma \langle v_* \rangle^{|\gamma|} \right) \theta^{-1-2s} d\theta \right\} dv dv_* \\ &\leq C \iint \left(\langle v \rangle^{2s+\gamma} \langle v_* \rangle^{2s+\max(|\gamma-2|-|\gamma|, 0)} \right) \mu_* |g|^2 dv dv_* \leq C \|g\|_{L_{s+\gamma/2}^2}^2, \end{aligned}$$

which together with (2.10) yields the desired estimate (2.8). The last estimate of the lemma is obvious by replacing μ by μ^ρ in each step of the above arguments, so that the proof of the lemma is completed. \square

Lemma 2.8 together with Lemma 2.12 implies that we have the following lower bound on the non-isotropic norm,

$$\|g\|^2 \gtrsim \left(\|g\|_{H_{\gamma/2}^s}^2 + \|g\|_{L_{s+\gamma/2}^2}^2 \right).$$

Therefore, to complete the proof of Proposition 2.2, it remains to show

Lemma 2.13. *Let $\gamma > -3$. Then we have*

$$J_1 \lesssim \|g\|_{H_{s+\gamma/2}^s}^2 + \|g\|_{L_{s+\gamma/2}^2}^2.$$

The same conclusion holds even if μ in J_1 is replaced by μ^ρ for any fixed $\rho > 0$.

Proof. As for Lemma 2.12, it follows from Proposition 2.10 that, for suitable constants $C_1, C_2 > 0$, we have

$$\begin{aligned} (2.11) \quad C_1 J_1^\Phi - C_2 \|g\|_{L_{s+\gamma/2}^2}^2 &\leq J_1^{\tilde{\Phi}} \\ &\leq 2 \iiint b(\cos \theta) \mu_* \langle v_* \rangle^{|\gamma|} \left(\langle v' \rangle^{\gamma/2} g' - \langle v \rangle^{\gamma/2} g \right)^2 d\sigma dv dv_* \\ &\quad + 2 \iiint b(\cos \theta) \mu_* \langle v_* \rangle^{|\gamma|} \left(\langle v \rangle^{\gamma/2} - \langle v' \rangle^{\gamma/2} \right)^2 |g|^2 d\sigma dv dv_* \\ &= B_1 + B_2, \end{aligned}$$

because $\tilde{\Phi}(|v - v_*|) \sim \langle v' - v_* \rangle^\gamma \leq \langle v' \rangle^\gamma \langle v_* \rangle^{|\gamma|}$ and $(a + b)^2 \leq 2(a^2 + b^2)$.

By the same argument for A_2 in the proof of Lemma 2.12, we get $B_2 \lesssim \|g\|_{L_{s+\gamma/2}^2}^2$.

To estimate B_1 , we apply Theorem 2.1 of [6] about the upper bound on the collision operator in the Maxwellian molecule case. It follows from (2.1.9) of [6] with $(m, \alpha) = (-s, -s)$ that

$$\left| \left(Q^{\Phi_0}(F, G), G \right) \right| \lesssim \|F\|_{L_{s+2s}^1} \|G\|_{H_s^s}^2.$$

Since $2a(b-a) = -(b-a)^2 + (a^2 - b^2)$, we get

$$\begin{aligned} (\mathcal{Q}^{\Phi_0}(F, G), G) &= \iiint bF_*G(G' - G) \\ &= -\frac{1}{2} \iiint bF_*(G' - G)^2 + \frac{1}{2} \iint F_*(G'^2 - G^2), \end{aligned}$$

and therefore

$$\begin{aligned} \left| \iiint bF_*(G' - G)^2 \right| &\leq 2 \left| (\mathcal{Q}^{\Phi_0}(F, G), G) \right| + \left| \iint F_*(G'^2 - G^2) \right| \\ &\lesssim \|F\|_{L^1_{s_3}} \|G\|_{H^s_{s+\gamma/2}}^2 + \|F\|_{L^1} \|G\|_{L^2}^2, \end{aligned}$$

where we have used the cancellation lemma from [3] for the second term. Choosing $F = \mu \langle \nu \rangle^{|\gamma|}$ and $G = \langle \nu \rangle^{\gamma/2} g$, it follows that $B_1 \lesssim \|g\|_{H^s_{s+\gamma/2}}^2$, completing the proof of the lemma. \square

2.2. Equivalence to the linearized operator. We will now show that the Dirichlet form of the linearized collision operator is equivalent to the square of the non-isotropic norm, and therefore, the proof of Proposition 2.1 will be given. Let us note that for the bilinear operator $\Gamma(\cdot, \cdot)$, for suitable functions f, g , one has

$$\begin{aligned} (\Gamma(f, g), h)_{L^2} &= \iiint b(\cos \theta) \Phi(|v - v_*|) \sqrt{\mu_*} (f'_* g' - f_* g) h \\ &= \iiint b(\cos \theta) \Phi(|v - v_*|) \sqrt{\mu'_*} (f_* g - f'_* g') h', \end{aligned}$$

and by adding these two lines, it follows that

$$(2.12) \quad (\Gamma(f, g), h)_{L^2} = \frac{1}{2} \iiint b(\cos \theta) \Phi(|v - v_*|) (f'_* g' - f_* g) (\sqrt{\mu_*} h - \sqrt{\mu'_*} h').$$

The following lemma shows that \mathcal{L}_1 dominates \mathcal{L} .

Lemma 2.14. *Under the conditions (1.2) on the cross-section with $0 < s < 1$ and $\gamma \in \mathbb{R}$, we have*

$$(2.13) \quad (\mathcal{L}_1 g, g)_{L^2} \geq \frac{1}{2} (\mathcal{L} g, g)_{L^2}.$$

Proof. By standard changes of variables, the following computations hold true

$$\begin{aligned} (\mathcal{L}_1 g, g)_{L^2} &= -(\Gamma(\sqrt{\mu}, g), g)_{L^2} \\ &= \frac{1}{2} \iiint B(|v - v_*|, \cos \theta) ((\mu'_*)^{1/2} g' - (\mu_*)^{1/2} g)^2 dv_* d\sigma dv \\ &= \frac{1}{2} \iiint B(|v - v_*|, \cos \theta) ((\mu')^{1/2} g'_* - (\mu)^{1/2} g_*)^2 dv_* d\sigma dv \\ &= \frac{1}{4} \iiint B(|v - v_*|, \cos \theta) \\ &\quad \times \left\{ ((\mu'_*)^{1/2} g' - (\mu_*)^{1/2} g)^2 + ((\mu')^{1/2} g'_* - (\mu)^{1/2} g_*)^2 \right\}, \end{aligned}$$

and

$$\begin{aligned}
(\mathcal{L}g, g)_{L^2} &= -(\Gamma(\sqrt{\mu}, g) + \Gamma(g, \sqrt{\mu}), g)_{L^2(\mathbb{R}_v^3)} \\
&= \iiint B(\mu_*^{1/2}g - (\mu'_*)^{1/2}g' + g_*(\mu)^{1/2} - g'_*(\mu')^{1/2})(\mu_*)^{1/2}g \\
&= \iiint B((\mu'_*)^{1/2}g' - (\mu_*)^{1/2}g + g'_*(\mu')^{1/2} - g_*(\mu)^{1/2})(\mu'_*)^{1/2}g' \\
&= \iiint B((\mu)^{1/2}g_* - (\mu')^{1/2}g'_* + g(\mu_*)^{1/2} - g'(\mu'_*)^{1/2})(\mu)^{1/2}g_* \\
&= \iiint B((\mu')^{1/2}g'_* - (\mu)^{1/2}g_* + g'(\mu'_*)^{1/2} - g(\mu_*)^{1/2})(\mu')^{1/2}g'_* \\
&= \frac{1}{4} \iiint B \{ (\mu_*)^{1/2}g - (\mu'_*)^{1/2}g' + (\mu)^{1/2}g_* - (\mu')^{1/2}g'_* \}^2.
\end{aligned}$$

Therefore, (2.13) follows from $(\alpha + \beta)^2 \leq 2(\alpha^2 + \beta^2)$ and the proof is completed. \square

Now for the term \mathcal{L}_2 , we have

Lemma 2.15. *One has*

$$|(\mathcal{L}_2g, h)_{L^2}| \lesssim \|\mu^{1/10^3}g\|_{L^2} \|\mu^{1/10^3}h\|_{L^2}.$$

Proof. It follows from (2.12) that

$$\begin{aligned}
(\mathcal{L}_2(g), h)_{L^2} &= -\frac{1}{2} \iiint B(g'_* \sqrt{\mu'} - g_* \sqrt{\mu})(\sqrt{\mu_*}h - \sqrt{\mu'_*}h') \\
&= (g, \mathcal{L}_2(h))_{L^2},
\end{aligned}$$

that is, \mathcal{L}_2 is symmetric. Hence it suffices to show the lemma in the case when $g = h$. Putting $G = \sqrt{\mu}g$, we have

$$\begin{aligned}
-\mathcal{L}_2g &= \mu^{-1/2}Q(G, \mu) \\
&= \mu^{-1/2} \iint b(\cos \theta) \Phi(|v - v_*|) G'_*(\mu' - \mu) dv_* d\sigma \\
&\quad + \sqrt{\mu} \iint b(\cos \theta) \Phi(|v - v_*|) (G'_* - G_*) dv_* d\sigma \\
&= I_1(v) + I_2(v).
\end{aligned}$$

Thanks to the cancellation lemma, we have $I_2(v) = \sqrt{\mu(v)}(S * G)(v)$ with $S(v) \sim |v|^\gamma$, whence we have

$$\begin{aligned}
(2.14) \quad |I_2, g)_{L^2}| &\lesssim \int_v \int_{v_*} |v - v_*|^\gamma \sqrt{\mu} \sqrt{\mu_*} |g| |g_*| dv dv_* \\
&\lesssim \int_v \int_{v_*} |v - v_*|^\gamma \{ (\mu_*^{1/4} \mu^{1/4} g)^2 + (\mu^{1/4} \mu_*^{1/4} g_*)^2 \} dv dv_* \\
&\lesssim \|\langle v \rangle^\gamma \mu^{1/4} g\|_{L^2}^2 \lesssim \|\mu^{1/8} g\|_{L^2}^2,
\end{aligned}$$

by means of Lemma 2.5.

Writing

$$\mu' - \mu = \sqrt{\mu'}(\sqrt{\mu'} - \sqrt{\mu}) + \sqrt{\mu}(\sqrt{\mu'} - \sqrt{\mu})$$

and using $\sqrt{\mu'\mu_*} = \sqrt{\mu\mu_*}$, we have

$$I_1(v) = \iint b(\cos \theta) \Phi(|v - v_*|) g'_*(\sqrt{\mu_*} + \sqrt{\mu'}) (\sqrt{\mu'} - \sqrt{\mu}) dv_* d\sigma.$$

Hence

$$\begin{aligned} (I_1, g)_{L^2} &= \iiint b(\cos \theta) \Phi(|v - v_*|) g'_*(\sqrt{\mu_*} - \sqrt{\mu'}) (\sqrt{\mu'} - \sqrt{\mu}) g dv dv_* d\sigma \\ &\quad + 2 \iiint b(\cos \theta) \Phi(|v - v_*|) G_*(\sqrt{\mu} - \sqrt{\mu'}) g' dv dv_* d\sigma \\ &= A_1 + A_2, \end{aligned}$$

where we have used the change of variables $(v, v_*) \rightarrow (v', v'_*)$ for the second term. We can write

$$\begin{aligned} A_1 &= \iiint b(\cos \theta) \Phi(|v - v_*|) (\mu_*^{1/4} - \mu_*'^{1/4}) (\mu'^{1/4} - \mu^{1/4}) g'_* g \\ &\quad \times (\mu_*^{1/4} + \mu_*'^{1/4}) (\mu'^{1/4} + \mu^{1/4}) d\sigma dv dv_*. \end{aligned}$$

Since we have

$$\begin{aligned} |v_*'|^2 &\leq (|v_*' - v'| + |v'|)^2 \leq (\sqrt{2}|v_* - v'| + |v'|)^2 \\ &\leq (\sqrt{2}|v_*| + (\sqrt{2} + 1)|v'|)^2 \leq 4|v_*|^2 + 2(\sqrt{2} + 1)^2|v'|^2, \end{aligned}$$

and in the same way, $|v|^2 \leq 4|v'|^2 + 2(\sqrt{2} + 1)^2|v_*|^2$, we get, by adding the two corresponding inequalities, that $\mu_*\mu' \leq (\mu_*'\mu)^{1/(10+4\sqrt{2})}$. Moreover, we have $\mu_*'\mu' = \mu_*\mu \leq (\mu_*'\mu)^{1/5}$ because $|v_*'|^2 \leq (|v_*' - v| + |v|)^2 \leq (|v_* - v| + |v|)^2 \leq 2|v_*|^2 + 8|v|^2$. Noticing that

$$\left| (\mu_*^{1/4} - \mu_*'^{1/4}) (\mu'^{1/4} - \mu^{1/4}) \right| \lesssim |v - v_*'|^2 \theta^2,$$

we get

$$\begin{aligned} (2.15) \quad |A_1| &\lesssim \iint |v - v_*'|^{\gamma+2} \left\{ \int_0^{\pi/2} \theta^{1-2s} d\theta \right\} (\mu_*'\mu)^{1/80} g'_* g dv dv_* \\ &\lesssim \iint |v - v_*'|^\gamma (\mu_*'\mu)^{1/160} g'_* g dv dv_* \lesssim \|\mu^{1/103} g\|_{L^2}^2, \end{aligned}$$

by an argument similar to the analysis of I_1 .

For A_2 , we use the regular change of variable $v \rightarrow v'$, and denote its inverse transformation by $v' \rightarrow \psi_\sigma(v') = v$. Then

$$\begin{aligned} A_2 &= 2 \iint \sqrt{\mu_*} g_* \left\{ \int_{\mathbb{S}^2} b\left(\frac{\psi_\sigma(v') - v_*}{|\psi_\sigma(v') - v_*|} \cdot \sigma\right) \Phi(|\psi_\sigma(v') - v_*|) \right. \\ &\quad \left. \times (\sqrt{\mu(\psi_\sigma(v'))} - \sqrt{\mu(v')}) \left| \frac{\partial(\psi_\sigma(v'))}{\partial(v')} \right| d\sigma \right\} g(v') dv_* dv'. \end{aligned}$$

It follows from the Taylor expansion that

$$\begin{aligned} \sqrt{\mu(\psi_\sigma(v'))} - \sqrt{\mu(v')} &= (\nabla \sqrt{\mu})(v') \cdot (\psi_\sigma(v') - v') \\ &\quad + \int_0^1 (1 - \tau) (\nabla^2 \sqrt{\mu})(v' + \tau(\psi_\sigma(v') - v')) (\psi_\sigma(v') - v')^2 d\tau. \end{aligned}$$

Note that the integral with respect to σ corresponding to the first order term vanishes, by means of the symmetry on \mathbb{S}^2 . Putting $v'_r = v' + \tau(\psi_\sigma(v') - v')$, we have $|v'|^2 \leq (|v' - v_*| + |v_*|)^2 \leq (|v'_r - v_*| + |v_*|)^2 \leq 2|v'_r|^2 + 8|v_*|^2$, so that

$$\left| \sqrt{\mu(v_*)} (\nabla^2 \sqrt{\mu})(v' + \tau(\psi_\sigma(v') - v')) \right| \lesssim (\mu(v_*) \mu(v'))^{1/12}.$$

Since $|\psi_\sigma(v') - v'| \lesssim |v' - v_*| \theta$, we have

$$\begin{aligned} |A_2| &\lesssim \iint \left\{ \int_0^{\pi/2} \theta^{1-2s} d\theta \right\} |v' - v_*|^{\gamma+2} (\mu_* \mu')^{1/12} |g_*| |g'| dv_* dv' \\ &\lesssim \iint |v' - v_*|^\gamma (\mu_* \mu')^{1/24} |g_*| |g'| dv_* dv' \lesssim \|\mu^{1/10^3} g\|_{L^2}^2. \end{aligned}$$

Together with (2.14) and (2.15), this yields the desired estimate and completes the proof of the lemma. \square

Let us note the following inequality between $(\mathcal{L}_1 g, g)_{L^2}$, corresponding to the first term of the linear operator, and the non-isotropic norm.

Proposition 2.16. *Let $\gamma > -3$. There exists a constant $C > 0$ such that*

$$(2.16) \quad \|g\|^2 \geq (\mathcal{L}_1 g, g)_{L^2} \geq \frac{1}{10} \|g\|^2 - C \|g\|_{L^2_{\gamma/2}}^2.$$

Proof. The equalities

$$\begin{aligned} 2(\mathcal{L}_1 g, g)_{L^2} &= -2(\Gamma(\sqrt{\mu}, g), g)_{L^2} \\ &= \iiint B \left((\mu'_*)^{1/2} g' - (\mu_*)^{1/2} g \right)^2 dv dv_* d\sigma \\ &= \iiint B \left((\mu'_*)^{1/2} (g' - g) + g((\mu'_*)^{1/2} - (\mu_*)^{1/2}) \right)^2 dv dv_* d\sigma, \end{aligned}$$

together with the inequality

$$2(a^2 + b^2) \geq (a + b)^2 \geq \frac{1}{2} a^2 - b^2$$

yields

$$(2.17) \quad \|g\|^2 \geq (\mathcal{L}_1 g, g)_{L^2} \geq \frac{1}{4} J_1 - \frac{1}{2} J_2 \geq \frac{1}{4} \|g\|^2 - \frac{3}{4} J_2.$$

It follows from the equality $(a + b)^2 = a^2 + b^2 + 2ab$ that

$$(2.18) \quad 2(\mathcal{L}_1 g, g)_{L^2} \geq J_2 - C \|g\|_{L^2_{\gamma/2}}^2,$$

which yields the desired estimate (2.16) together with (2.17).

Indeed, note that

$$\begin{aligned} &2(\mathcal{L}_1 g, g)_{L^2} \\ &= \iiint B \left((\mu'_*)^{1/2} (g' - g) + g((\mu'_*)^{1/2} - (\mu_*)^{1/2}) \right)^2 dv dv_* d\sigma \\ &= J_1 + J_2 + 2 \iiint B (g' - g) g (\mu'_*)^{1/2} ((\mu'_*)^{1/2} - (\mu_*)^{1/2}) dv dv_* d\sigma. \end{aligned}$$

Using the identity $2(\beta - \alpha)\alpha = \beta^2 - \alpha^2 - (\beta - \alpha)^2$, we have

$$\begin{aligned}
& 2(g' - g)g(\mu'_*)^{1/2}(\mu'_*)^{1/2} - (\mu_*)^{1/2}) \\
&= \frac{1}{2}(g'^2 - g^2 - (g' - g)^2)(\mu'_* - \mu_* + ((\mu'_*)^{1/2} - (\mu_*)^{1/2})^2) \\
&= -\frac{1}{2}(g' - g)^2((\mu'_*)^{1/2} - (\mu_*)^{1/2})^2 + \frac{1}{2}(g^2 - g'^2)(\mu_* - \mu'_*) \\
&\quad + \frac{1}{2}(g' - g)^2(\mu_* - \mu'_*) + \frac{1}{2}(g'^2 - g^2)((\mu'_*)^{1/2} - (\mu_*)^{1/2})^2 \\
&= I_1 + I_2 + I_3 + I_4.
\end{aligned}$$

Using the change of variables $(v', v'_*) \rightarrow (v, v_*)$, we see that

$$\left| \iiint B I_2 dv dv_* d\sigma \right| = \left| \iiint B \mu_* (g^2 - g'^2) dv dv_* d\sigma \right| \leq C \|g\|_{L^2_{\gamma/2}}^2,$$

by means of the cancellation lemma. Furthermore,

$$\begin{aligned}
\iiint B I_1 dv dv_* d\sigma &= -\frac{1}{2} \iiint B (\mu_* + \mu'_*) (g' - g)^2 dv dv_* d\sigma \\
&\quad + \iiint B (\mu_*)^{1/2} (\mu'_*)^{1/2} (g' - g)^2 dv dv_* d\sigma \geq -J_1,
\end{aligned}$$

where we have used the change of variables $(v', v'_*) \rightarrow (v, v_*)$. Thus, we obtain (2.18) because the integrals corresponding to the last two terms I_3 and I_4 vanish, ending the proof of the proposition. \square

End of the proof of Proposition 2.1: It follows from (2.16) and (2.13) that

$$\|g\|^2 \geq (\mathcal{L}_1 g, g)_{L^2} \geq \frac{1}{2} (\mathcal{L} g, g)_{L^2}.$$

On the other hand, note that $(\mathcal{L} g, g)_{L^2} = (\mathcal{L}(\mathbf{I} - \mathbf{P})g, (\mathbf{I} - \mathbf{P})g)_{L^2}$, from the very definition of the projection operator \mathbf{P} .

Thus, from Proposition 2.16 and Lemma 2.15, we get

$$\begin{aligned}
(\mathcal{L} g, g)_{L^2} &= (\mathcal{L}_1(\mathbf{I} - \mathbf{P})g, (\mathbf{I} - \mathbf{P})g)_{L^2} + (\mathcal{L}_2(\mathbf{I} - \mathbf{P})g, (\mathbf{I} - \mathbf{P})g)_{L^2} \\
&\geq \frac{1}{10} \|(\mathbf{I} - \mathbf{P})g\|^2 - C \|(\mathbf{I} - \mathbf{P})g\|_{L^2_{\gamma/2}}^2.
\end{aligned}$$

Since it is known from [31] that we have

$$(\mathcal{L} g, g)_{L^2} \geq C \|(\mathbf{I} - \mathbf{P})g\|_{L^2_{\gamma/2}}^2,$$

we get on the whole

$$\|(\mathbf{I} - \mathbf{P})g\|^2 \leq C (\mathcal{L} g, g)_{L^2}.$$

2.3. Non-isotropic norms with different kinetic factors. This subsection is devoted to the proof of Proposition 2.4. That is, we will show some equivalence relations between the non-isotropic norms with different kinetic factors and different weights.

For the proof, we introduce some further notations. Let $\rho > 0$, $\mu_\rho(v) = \mu(v)^\rho$, and set

$$J_{1,\rho}^{\Phi_\gamma}(g) = \iiint \Phi_\gamma(|v - v_*|) b(\cos \theta) \mu_{\rho,*} (g' - g)^2 dv dv_* d\sigma.$$

We simply write $J_1^{\Phi_\gamma}(g)$ if $\rho = 1$, and also introduce the notation $J_{2,\rho}^{\Phi_\gamma}(g)$ similarly with μ replaced by μ_ρ .

Then it follows from (2.4) and the change of variables $v \rightarrow v/\sqrt{\rho}$ that

$$J_{2,\rho}^{\Phi_\gamma}(g) \sim \|g\|_{L_{s+\gamma/2}^2}^2 = \|\langle v \rangle^{(\gamma-\beta)/2} g\|_{L_{s+\beta/2}^2}^2 \sim J_{2,\rho}^{\Phi_\beta}(\langle v \rangle^{(\gamma-\beta)/2} g).$$

By the last assertions of Lemmas 2.12 and 2.13, there exist constants $C_1, C_2 > 0$ such that

$$(2.19) \quad C_1 \|g\|_{H_{s+\gamma/2}^s}^2 \leq J_{1,\rho}^{\Phi_\gamma}(g) + \|g\|_{L_{s+\gamma/2}^2}^2 \leq C_2 \|g\|_{H_{s+\gamma/2}^s}^2.$$

Furthermore, it follows from (2.9), (2.11) and the proofs of Lemmas 2.12 and 2.13 that

$$J_{1,2}^{\Phi_0}(\langle v \rangle^{\gamma/2} g) \lesssim J_1^{\Phi_\gamma}(g) \lesssim J_{1,1/2}^{\Phi_0}(\langle v \rangle^{\gamma/2} g), \quad \text{modulo } \|g\|_{L_{s+\gamma/2}^2}^2,$$

because we have $C_1 \mu_2 \leq \mu \langle v \rangle^{\pm|\gamma|} \leq C_2 \mu_{1/2}$.

Therefore, to complete the proof of Proposition 2.3, it suffices to show that for any $\rho, \rho' > 0$

$$(2.20) \quad J_{1,\rho}^{\Phi_0}(g) \sim J_{1,\rho'}^{\Phi_0}(g), \quad \text{modulo } \|g\|_{L_s^2}^2.$$

In fact, note that

$$J_1^{\Phi_\gamma}(g) \sim J_{1,\rho}^{\Phi_0}(\langle v \rangle^{\gamma/2} g) \sim J_1^{\Phi_\beta}(\langle v \rangle^{(\gamma-\beta)/2} g), \quad \text{modulo } \|g\|_{L_{s+\gamma/2}^2}^2.$$

This equivalence looks quite obvious, however, for completeness, we shall give a proof. In fact, (2.20) is a direct consequence of the following lemma, by taking $f = \mu_{\rho'}$.

Lemma 2.17. *Assume that (1.2) holds with $0 < s < 1$. Then there exists a constant $C > 0$ such that*

$$\iiint b f_*^2(g' - g)^2 d\sigma dv dv_* \leq C \|f\|_{L_s^2}^2 (J_{1,\rho}^{\Phi_0}(g) + \|g\|_{L_s^2}^2).$$

Once the equivalence (2.20) has been established, we have

Corollary 2.18. *Assume that (1.2) holds with $0 < s < 1$. Then there exists a constant $C > 0$ such that*

$$\iiint b f_*^2(g' - g)^2 d\sigma dv dv_* \leq C \|f\|_{L_s^2}^2 \|g\|_{\Phi_0}^2.$$

Proof. It is enough to consider the case $\rho = 1$. As in the proof of Lemma 2.12, it follows from Proposition 2 of [3] that

$$(2.21) \quad \begin{aligned} J_1^{\Phi_0}(g) &= \iiint b(\cos \theta) \mu_*(g' - g)^2 dv_* d\sigma dv \\ &= \frac{1}{(2\pi)^3} \iint b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) (\widehat{\mu}(0) |\widehat{g}(\xi)|^2 + |\widehat{g}(\xi^+)|^2 \\ &\quad - 2\operatorname{Re} \widehat{\mu}(\xi^-) \widehat{g}(\xi^+) \overline{\widehat{g}(\xi)}) d\xi d\sigma \\ &= \frac{1}{(2\pi)^3} \iint b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) (\widehat{\mu}(0) |\widehat{g}(\xi) - \widehat{g}(\xi^+)|^2 \\ &\quad + 2\operatorname{Re} (\widehat{\mu}(0) - \widehat{\mu}(\xi^-)) \widehat{g}(\xi^+) \overline{\widehat{g}(\xi)}) d\xi d\sigma, \end{aligned}$$

and

$$\begin{aligned} A &= \iiint b(\cos \theta) f_*^2(g' - g)^2 dv_* d\sigma dv \\ &= \frac{1}{(2\pi)^3} \iint b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) (\widehat{f}^2(0) |\widehat{g}(\xi) - \widehat{g}(\xi^+)|^2 \\ &\quad + 2\operatorname{Re} (\widehat{f}^2(0) - \widehat{f}^2(\xi^-)) \widehat{g}(\xi^+) \overline{\widehat{g}(\xi)}) d\xi d\sigma. \end{aligned}$$

Since $\widehat{f^2}(0) = \|f\|_{L^2}^2$ and $\widehat{\mu}(0) = c_0 > 0$, we obtain

$$\begin{aligned} c_0 A &= c_0 \iiint b(\cos \theta) f_*^2(g' - g)^2 dv_* d\sigma dv \\ &= \|f\|_{L^2}^2 \mathcal{J}_1^{\Phi_0}(g) \\ &\quad - \frac{2}{(2\pi)^3} \|f\|_{L^2}^2 \iint b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) \operatorname{Re}\left(\widehat{\mu}(0) - \widehat{\mu}(\xi^-)\right) \widehat{g}(\xi^+) \overline{\widehat{g}(\xi)} d\xi d\sigma \\ &\quad + \frac{2c_0}{(2\pi)^3} \iint b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) \operatorname{Re}\left(\widehat{f^2}(0) - \widehat{f^2}(\xi^-)\right) \widehat{g}(\xi^+) \overline{\widehat{g}(\xi)} d\xi d\sigma \\ &= \|f\|_{L^2}^2 \mathcal{J}_1^{\Phi_0}(g) + A_1 + A_2. \end{aligned}$$

Write

$$\begin{aligned} A_2 &= \frac{2c_0}{(2\pi)^3} \left\{ \int |\widehat{g}(\xi)|^2 \left(\int b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) \operatorname{Re}\left(\widehat{f^2}(0) - \widehat{f^2}(\xi^-)\right) d\sigma \right) d\xi \right. \\ &\quad \left. + \iint b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) \operatorname{Re}\left(\widehat{f^2}(0) - \widehat{f^2}(\xi^-)\right) (\widehat{g}(\xi^+) - \widehat{g}(\xi)) \overline{\widehat{g}(\xi)} d\xi d\sigma \right\} \\ &= A_{2,1} + A_{2,2}. \end{aligned}$$

It follows from Cauchy-Schwarz's inequality that

$$\begin{aligned} |A_{2,2}| &\leq C \left(\iint b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) |\widehat{f^2}(0) - \widehat{f^2}(\xi^-)|^2 |\widehat{g}|^2(\xi) d\xi d\sigma \right)^{1/2} \\ &\quad \times \left(\iint b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) |\widehat{g}(\xi^+) - \widehat{g}(\xi)|^2 d\xi d\sigma \right)^{1/2} \\ &= B_1^{1/2} \times B_2^{1/2}. \end{aligned}$$

Since

$$|\widehat{f^2}(0) - \widehat{f^2}(\xi^-)| \leq \int f^2(v) |1 - e^{-iv \cdot \xi^-}| dv,$$

we have

$$\begin{aligned} B_1 &\leq C \iiint |\widehat{g}(\xi)|^2 f^2(v) f^2(w) \\ &\quad \times \left(\int b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) (|1 - e^{-iv \cdot \xi^-}|^2 + |1 - e^{-iw \cdot \xi^-}|^2) d\sigma \right) dv dw d\xi \\ &\leq C \|g\|_{H^s}^2 \|f\|_{L^2}^2 \|f\|_{L^2_s}^2, \end{aligned}$$

because

$$\begin{aligned} &\int b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) |1 - e^{-iv \cdot \xi^-}|^2 d\sigma \\ &\leq C \left(\int_0^{(\langle v \rangle \langle \xi \rangle)^{-1}} \theta^{-1-2s} (|v||\xi|)^2 \theta^2 d\theta + \int_{(\langle v \rangle \langle \xi \rangle)^{-1}}^{\pi/2} \theta^{-1-2s} d\theta \right) \\ &\leq C \langle v \rangle^{2s} \langle \xi \rangle^{2s}. \end{aligned}$$

Then we have $|A_{2,1}| \leq C \|g\|_{H^s}^2 \|f\|_{L^2_s}^2$ because

$$\begin{aligned} & \int b \left(\frac{\xi}{|\xi|} \cdot \sigma \right) \operatorname{Re} \left(\widehat{f^2}(0) - \widehat{f^2}(\xi^-) \right) d\sigma \\ &= \int f^2(v) \left(\int b \left(\frac{\xi}{|\xi|} \cdot \sigma \right) (1 - \cos(v \cdot \xi^-)) d\sigma \right) dv \\ &\leq C \langle \xi \rangle^{2s} \int f^2(v) \langle v \rangle^{2s} dv. \end{aligned}$$

Since $\widehat{\mu}(\xi)$ is real-valued, it follows that

$$\operatorname{Re} \left(\widehat{\mu}(0) - \widehat{\mu}(\xi^-) \right) \widehat{g}(\xi^+) \overline{\widehat{g}(\xi)} = \left(\int (1 - \cos(v \cdot \xi^-)) \mu(v) dv \right) \operatorname{Re} \widehat{g}(\xi^+) \overline{\widehat{g}(\xi)}.$$

Therefore, by using Cauchy-Schwarz's inequality and the change of variables $\xi \rightarrow \xi^+$ (see the proof of Lemma 2.8 in [7]), we obtain $|A_1| \leq C \|f\|_{L^2}^2 \|g\|_{H^s}^2$. Furthermore, it follows from (2.21) that

$$\begin{aligned} B_2 &= \iint b \left(\frac{\xi}{|\xi|} \cdot \sigma \right) |\widehat{g}(\xi) - \widehat{g}(\xi^+)|^2 d\xi d\sigma \\ &\leq C \left(J_1^{\Phi_0}(g) + \|g\|_{H^s}^2 \right), \end{aligned}$$

which yields $|A_{2,2}| \leq C \|f\|_{L^2} \|f\|_{L^2_s} \|g\|_{H^s} \left(J_1^{\Phi_0}(g) + \|g\|_{H^s}^2 \right)^{1/2}$. Hence

$$|A_2| \leq C \|f\|_{L^2_s}^2 \|g\|_{H^s} \left(J_1^{\Phi_0}(g) + \|g\|_{H^s}^2 \right)^{1/2}.$$

Finally, we have

$$A \leq C \|f\|_{L^2_s}^2 \|g\|_{H^s} \left(J_1^{\Phi_0}(g) + \|g\|_{H^s}^2 \right)^{1/2} \leq C \|f\|_{L^2_s}^2 \left(J_1^{\Phi_0}(g) + \|g\|_{L^2_s}^2 \right),$$

by means of (2.19) with $\gamma = 0$, completing the proof of the lemma. \square

3. ESTIMATES OF NONLINEAR COLLISION OPERATOR IN VELOCITY SPACE

In this section, we derive various estimates on the nonlinear collision operator. Even though we consider the soft potential case in this paper, some of the following estimates also hold for general case so that they will be used in Part II.

3.1. Upper bounds in general case. In this sub-section, we will establish various functional estimates which hold true under the more general assumption $0 < s < 1$ and $\gamma > -3$. In particular, all the results in this part are independent of the sign of $\gamma + 2s$.

Proposition 3.1. *For all $0 < s < 1$ and $\gamma > -3$, we have*

$$\left| (\Gamma(f, g), h)_{L^2} \right| \lesssim \|h\|_{\Phi_\gamma} \left\{ \|f\|_{L^\infty} \|g\|_{\Phi_\gamma} + \left(\|\nabla f\|_{L^\infty} + \|f\|_{L^\infty} \right) \|g\|_{L^2_{s+\gamma/2}}^2 \right\}.$$

Proof. Direct calculation gives

$$\begin{aligned} (\Gamma(f, g), h)_{L^2} &= \iiint b \Phi_\gamma \mu_*^{1/2} (f'_* g' - f_* g) h dv dv_* d\sigma \\ &= \frac{1}{2} \iiint (b \Phi_\gamma)^{1/2} \mu_*^{1/4} (f'_* g' - f_* g) (b \Phi_\gamma)^{1/2} (\mu_*^{1/4} h - \mu_*'^{1/4} h') \\ &\quad + \frac{1}{2} \iiint (b \Phi_\gamma)^{1/2} \mu_*^{1/4} (f'_* g' - f_* g) (b \Phi_\gamma)^{1/2} (\mu_*^{1/4} - \mu_*'^{1/4}) h. \end{aligned}$$

Noticing that

$$\mu_*^{1/4}h - \mu_*'^{1/4}h' = \mu_*'^{1/4}(h - h') + (\mu_*^{1/4} - \mu_*'^{1/4})h,$$

by using the Cauchy-Schwarz inequality, we have

$$\begin{aligned} |(\Gamma(f, g), h)| &\lesssim \left(\iiint b\Phi_\gamma \mu_*^{1/2} (f_*'g' - f_*g)^2 d\sigma dv dv_* \right)^{1/2} \|h\|_{\Phi_\gamma}, \\ &= A^{1/2} \|h\|_{\Phi_\gamma}, \end{aligned}$$

where we have used the fact that the non-isotropic norm is invariant by replacing μ by μ^ρ for any fixed $\rho > 0$ (see the previous section). We then estimate

$$\begin{aligned} A &\leq 3 \left(\iiint b\Phi_\gamma \mu_*^{1/4} (\mu_*^{1/8} f)_*' - (\mu_*^{1/8} f)_* \right)^2 g^2 d\sigma dv dv_* \\ &\quad + \iiint b\Phi_\gamma \mu_*^{1/8} (\mu_*^{1/8} f)_*'^2 (g' - g)^2 d\sigma dv dv_* \\ &\quad + \iiint b\Phi_\gamma \mu_*^{1/4} (\mu_*^{1/8} - \mu_*'^{1/8})^2 (f_*'g')^2 d\sigma dv dv_* \\ &= A_1 + A_2 + A_3. \end{aligned}$$

It is easy to see that

$$A_2 + A_3 \lesssim \|f\|_{L^\infty}^2 \|g\|_{\Phi_\gamma}^2.$$

Note that

$$\left| (\mu_*^{1/8} f)_*' - (\mu_*^{1/8} f)_* \right| \leq \min \left\{ (\|\nabla f\|_{L^\infty} + \|f\|_{L^\infty}) \theta |v - v_*|, \|f\|_{L^\infty} \right\}.$$

Then we have

$$\begin{aligned} A_1 &\lesssim (\|\nabla f\|_{L^\infty} + \|f\|_{L^\infty})^2 \iint \Phi_\gamma \left(\int b(\cos \theta) \min(\theta^2 |v - v_*|^2, 1) d\sigma \right) \mu_*^{1/4} g^2 dv dv_* \\ &\lesssim (\|\nabla f\|_{L^\infty} + \|f\|_{L^\infty})^2 \iint |v - v_*|^{\gamma+2s} \mu_*^{1/4} g^2 dv dv_* \lesssim (\|\nabla f\|_{L^\infty} + \|f\|_{L^\infty})^2 \|g\|_{s+\gamma/2}^2, \end{aligned}$$

where we have used $\gamma + 2s > -3$ and the fact that

$$\begin{aligned} \int b(\cos \theta) \min(\theta^2 |v - v_*|^2, 1) d\sigma &\leq |v - v_*|^2 \int_0^{\min(\pi/2, |v-v_*|^{-1})} \theta^{1-2s} d\theta \\ &\quad + \int_{\min(\pi/2, |v-v_*|^{-1})}^{\pi/2} \theta^{-1-2s} d\theta \lesssim |v - v_*|^{2s}. \end{aligned}$$

And this completes the proof of the proposition. \square

Lemma 3.2. *Let $\gamma \geq 0$. Assume that (1.2) holds with $0 < s < 1$. Then*

$$\iiint \Phi_\gamma(|v - v_*|) b f_*^2 (g' - g)^2 d\sigma dv dv_* \lesssim \|f\|_{L_{s+\gamma/2}^2}^2 \|g\|_{\Phi_\gamma}^2.$$

Proof. Since $\Phi_\gamma(|v - v_*|) \lesssim \langle v' \rangle^\gamma + \langle v_* \rangle^\gamma$, we have

$$\begin{aligned} & \iiint b(\cos \theta) \Phi(|v - v_*|) f_*^2 (g' - g)^2 d\sigma dv dv_* \\ & \lesssim \iiint b(\cos \theta) f_*^2 (\langle v' \rangle^{\gamma/2} g' - \langle v \rangle^{\gamma/2} g)^2 d\sigma dv dv_* \\ & \quad + \iiint b(\cos \theta) (\langle v_* \rangle^{\gamma/2} f_*)^2 (g' - g)^2 d\sigma dv dv_* \\ & \quad + \iiint b(\cos \theta) f_*^2 (\langle v \rangle^{\gamma/2} - \langle v' \rangle^{\gamma/2})^2 |g|^2 d\sigma dv dv_* \\ & = A_1 + A_2 + A_3. \end{aligned}$$

Noticing that

$$\begin{aligned} |\langle v \rangle^{\gamma/2} - \langle v' \rangle^{\gamma/2}| & \leq C_\gamma \int_0^1 \langle v' + \tau(v - v') \rangle^{(\gamma/2-1)^+} d\tau |v - v_*| \theta \\ & \leq C'_\gamma (\langle v \rangle^{(\gamma/2-1)^+} + \langle v_* \rangle^{(\gamma/2-1)^+}) \langle v - v_* \rangle \theta, \end{aligned}$$

we have

$$\begin{aligned} A_3 & \lesssim \iint f_*^2 |g|^2 \{ (\langle v \rangle^{(\gamma/2-1)^+} + \langle v_* \rangle^{(\gamma/2-1)^+})^2 \\ & \quad \times \left(\int_0^{\langle v-v_* \rangle^{-1}} \theta^{-1-2s} (\langle v - v_* \rangle \theta)^2 d\theta \right. \\ & \quad \left. + \int_{\langle v-v_* \rangle^{-1}}^{\pi/2} (\langle v \rangle^{\gamma/2} + \langle v_* \rangle^{\gamma/2})^2 \theta^{-1-2s} d\theta \right) dv dv_* \\ & \lesssim \iint (\langle v \rangle^{2s+\gamma} + \langle v_* \rangle^{2s+\gamma}) f_*^2 |g|^2 dv dv_* \\ & \lesssim (\|f\|_{L^2_{s+\gamma/2}}^2 \|g\|_{L^2}^2 + \|f\|_{L^2}^2 \|g\|_{L^2_{s+\gamma/2}}^2). \end{aligned}$$

Applying Corollary 2.18 to A_1 and A_2 , it follows that

$$\begin{aligned} A_1 + A_2 & \lesssim \|f\|_{L^2_s}^2 \| \langle v \rangle^{\gamma/2} g \|_{\Phi_0}^2 + \| \langle v \rangle^{\gamma/2} f \|_{L^2_s}^2 \|g\|_{\Phi_0}^2 \\ & \lesssim \|f\|_{L^2_{s+\gamma/2}}^2 \|g\|_{\Phi_\gamma}^2, \end{aligned}$$

where we have used Proposition 2.4 in the last inequality. \square

Proposition 3.3. *For all $0 < s < 1$ and $\gamma > -3$, one has*

$$\begin{aligned} |(\Gamma(f, g), h)_{L^2}| & \lesssim \|h\|_{\Phi_\gamma} \{ \|f\|_{L^2_{s+\gamma/2}} \|g\|_{\Phi_\gamma} + \|g\|_{L^2_{s+\gamma/2}} \|f\|_{\Phi_\gamma} \\ & \quad + \min(\|f\|_{L^2} \|g\|_{L^2_{s+\gamma/2}}, \|f\|_{L^2_{s+\gamma/2}} \|g\|_{L^2}) \\ & \quad + \|\mu^{1/40} f\|_{L^2} \|\mu^{1/60} g\|_{H^{\max(-\gamma/2, 1)}} + \|\mu^{1/40} f\|_{\Phi_\gamma} \|\mu^{1/40} g\|_{L^\infty}^2 \}. \end{aligned}$$

Proof. We will use the decomposition

$$\Phi_\gamma(z) = |z|^\gamma \mathbf{1}_{\{|z| \geq 1\}} + |z|^\gamma \mathbf{1}_{\{|z| \leq 1\}} = \Phi_A(z) + \Phi_B(z).$$

We denote by $\Gamma_A(\cdot, \cdot), \Gamma_B(\cdot, \cdot)$ the collision operators with the kinetic factor in the cross section given by Φ_A and Φ_B respectively. Similarly to the proof of Proposition 3.1, we

have

$$\begin{aligned} |(\Gamma_A(f, g), h)| &\lesssim \left(\iiint b \Phi_A \mu_*^{1/2} (f'_* g' - f_* g)^2 d\sigma dv dv_* \right)^{1/2} \|h\|_{\Phi_\gamma} \\ &= A^{1/2} \|h\|_{\Phi_\gamma}. \end{aligned}$$

Since $\Phi_A \leq 2^{|\gamma|} \tilde{\Phi}_\gamma$, we have

$$\begin{aligned} A &\lesssim \iiint b \tilde{\Phi}_\gamma \mu_*^{1/4} (\mu^{1/8} f'_* - (\mu^{1/8} f)_*)^2 g^2 d\sigma dv dv_* \\ &\quad + \iiint b \tilde{\Phi}_\gamma \mu_*^{1/8} ((\mu^{1/8} f)_*)^2 (g' - g)^2 d\sigma dv dv_* \\ &\quad + \iiint b \tilde{\Phi}_\gamma \mu_*^{1/4} (\mu_*^{1/8} - \mu_*'^{1/8})^2 (f'_* g')^2 d\sigma dv dv_* \\ &= A_1 + A_2 + A_3. \end{aligned}$$

In order to estimate A_1 , we make use of Lemma 3.2. Since $\tilde{\Phi}_\gamma(|v - v_*|) \mu_*^{1/4} \lesssim \langle v \rangle^\gamma$, we have by putting $f = \mu^{1/8} f$ and $g = \langle v \rangle^{\gamma/2} g$,

$$\begin{aligned} A_1 &\lesssim \iiint b (\langle v \rangle^{\gamma/2} g)^2 ((\mu^{1/8} f)'_* - (\mu^{1/8} f)_*)^2 d\sigma dv dv_* \\ &\lesssim \|\langle v \rangle^{\gamma/2} g\|_{L_s^2}^2 \|\mu^{1/8} f\|_{\Phi_0}^2 \lesssim \|g\|_{L_{s+\gamma/2}^2}^2 \|f\|_{\Phi_\gamma}^2. \end{aligned}$$

We decompose the estimation on A_2 as

$$\begin{aligned} A_2 &\lesssim \iiint b (\mu^{1/8} f'_*)^2 ((\langle v \rangle^{\gamma/2} g)' - (\langle v \rangle^{\gamma/2} g))^2 d\sigma dv dv_* \\ &\quad + \iiint b (\langle v \rangle^{\gamma/2} - \langle v' \rangle^{\gamma/2})^2 (\mu^{1/8} f'_*)^2 g'^2 d\sigma dv dv_* \\ &= A_{2,1} + A_{2,2}. \end{aligned}$$

Applying again Lemma 3.2 to $A_{2,1}$, we get

$$A_{2,1} \lesssim \|\mu^{1/8} f\|_{L_s^2}^2 \|\langle v \rangle^{\gamma/2} g\|_{\Phi_0}^2 \lesssim \|f\|_{L_{s+\gamma/2}^2}^2 \|g\|_{\Phi_\gamma}^2.$$

For $A_{2,2}$, we note that if $v_\tau = v' + \tau(v - v')$ for $\tau \in [0, 1]$, then

$$\langle v \rangle \leq \langle v - v_* \rangle + \langle v_* \rangle \leq \sqrt{2} \langle v_\tau - v_* \rangle + \langle v_* \rangle \leq (1 + \sqrt{2}) \langle v_\tau \rangle \langle v_* \rangle,$$

and $\langle v_\tau \rangle \leq (1 + \sqrt{2}) \langle v \rangle \langle v_* \rangle$. Thus, $\langle v_\tau \rangle^\beta \leq C_\beta \langle v \rangle^\beta \langle v_* \rangle^{\beta|}$ for any $\beta \in \mathbb{R}$. Since it follows that

$$\begin{aligned} |\langle v \rangle^{\gamma/2} - \langle v' \rangle^{\gamma/2}| &\leq C_\gamma \int_0^1 \langle v' + \tau(v - v') \rangle^{(\gamma/2-1)} d\tau |v - v_*| \theta \\ &\leq C'_\gamma (\langle v \rangle^{(\gamma/2-1)} \langle v_* \rangle^{|\gamma/2-1|}) \langle v - v_* \rangle \theta, \end{aligned}$$

we have, by using the change of variables $(v', v'_*) \rightarrow (v, v_*)$,

$$\begin{aligned} A_{2,2} &\lesssim \iint \frac{(\mu^{1/8} f'_*)^2}{\langle v_* \rangle^{|\gamma|}} |g|^2 \left\{ (\langle v \rangle^{(\gamma-2)} \langle v_* \rangle^{|\gamma-2|})^2 \left(\int_0^{\langle v-v_* \rangle^{-1}} \theta^{-1-2s} (\langle v - v_* \rangle \theta)^2 d\theta \right) \right. \\ &\quad \left. + \int_{\langle v-v_* \rangle^{-1}}^{\tau/2} (\langle v \rangle^\gamma + \langle v \rangle^\gamma \langle v_* \rangle^{|\gamma|}) \theta^{-1-2s} d\theta \right\} dv dv_* \\ &\lesssim \iint (\langle v \rangle^{2s+\gamma} \langle v_* \rangle^{2s+\max(|\gamma-2|-|\gamma|, 0)}) (\mu^{1/8} f'_*)^2 |g|^2 dv dv_* \lesssim \|\mu^{1/10} f\|_{L^2}^2 \|g\|_{L_{s+\gamma/2}^2}^2. \end{aligned}$$

Noticing that $(\mu_*^{1/8} - \mu_*'^{1/8})^2 \lesssim \min(|v - v_*|^2 \theta^2, 1)$, we have

$$\begin{aligned} A_3 &\lesssim \iint \tilde{\Phi}_\gamma \left(\int_{\mathbb{S}^2} b(\cos \theta) \min(|v - v_*|^2 \theta^2, 1) d\sigma \right) f_*^2 g^2 dv dv_* \\ &\lesssim \iint \langle v - v_* \rangle^{\gamma+2s} f_*^2 g^2 dv dv_* \\ &\lesssim \iint \langle v_* \rangle^{\gamma+2s} f_*^2 \langle v \rangle^{\gamma+2s} g^2 dv dv_* \lesssim \|f\|_{L^2_{s+\gamma/2}}^2 \|g\|_{L^2_{s+\gamma/2}}^2, \end{aligned}$$

when $\gamma + 2s \geq 0$ because $\langle v - v_* \rangle^{\gamma+2s} \leq \langle v_* \rangle^{\gamma+2s} \langle v \rangle^{\gamma+2s}$.

To consider the case $\gamma + 2s < 0$, we divide the space $\mathbb{R}_v^3 \times \mathbb{R}_{v_*}^3$ into three parts

$$\begin{aligned} U_1 &= \{|v - v_*| \leq |v_*|/8\}, \quad U_2 = \{|v - v_*| > |v_*|/8\} \cap \{|v_*| \leq 1\}, \\ U_3 &= \{|v - v_*| > |v_*|/8\} \cap \{|v_*| > 1\}. \end{aligned}$$

Then we have

$$\begin{aligned} \frac{1}{3} A_3 &= \iiint b \tilde{\Phi}_\gamma \mu_*'^{1/4} (\mu_*'^{1/8} - \mu_*^{1/8})^2 (f_* g)^2 d\sigma dv dv_* \\ &= \iint_{U_1} \int d\sigma dv dv_* + \iint_{U_2} \int d\sigma dv dv_* + \iint_{U_3} \int d\sigma dv dv_* \\ &= A_{3,1} + A_{3,2} + A_{3,3}. \end{aligned}$$

Since $|v' - v_*| \leq |v - v_*| \leq |v_*|/8$, we have $7|v_*|/8 \leq |v'|$, $|v| \leq 9|v_*|/8$ and $|v_*'|^2 = |v|^2 + |v_*|^2 - |v'|^2 \geq |v_*|^2/2$ in U_1 . Hence, in this region, we have $\mu_*'^{1/4} \leq C\mu_*^{1/8} \leq C(\mu_*\mu)^{1/20}$, and this to

$$A_{3,1} \lesssim \iint (\mu\mu_*)^{1/20} \langle v - v_* \rangle^{\gamma+2s} f_*^2 g^2 dv dv_* \lesssim \|f\|_{L^2_{s+\gamma/2}}^2 \|g\|_{L^2_{s+\gamma/2}}^2.$$

Furthermore, we have

$$A_{3,2} \lesssim \iint_{U_2} \langle v - v_* \rangle^{\gamma+2s} f_*^2 g^2 dv dv_* \lesssim \|f\|_{L^2_{s+\gamma/2}}^2 \|g\|_{L^2_{s+\gamma/2}}^2,$$

because $\langle v - v_* \rangle^{-1} \leq \langle v \rangle^{-1} \langle v_* \rangle^{-1} \langle v_* \rangle^2 \leq 2\langle v \rangle^{-1} \langle v_* \rangle^{-1}$ in U_2 . Since $\langle v - v_* \rangle^{-1} \leq 8|v_*|^{-1} \leq 16\langle v_* \rangle^{-1}$ in U_3 , we get

$$A_{3,3} \lesssim \|f\|_{L^2_{s+\gamma/2}}^2 \|g\|_{L^2}^2.$$

Therefore, we have, when $\gamma + 2s \leq 0$

$$A_3 \lesssim \|f\|_{L^2_{s+\gamma/2}}^2 \|g\|_{L^2}^2.$$

By considering another partition in \mathbb{R}_{v, v_*}^6 with v and v_* exchanged, the estimate $A_3 \lesssim \|f\|_{L^2}^2 \|g\|_{L^2_{s+\gamma/2}}^2$ holds, because $|v_*' - v| \leq |v_* - v| \leq |v|/8$ implies $7|v|/8 \leq |v_*'|$, $|v_*| \leq 9|v|/8$.

As a conclusion, we have in summary that

$$\begin{aligned} \left| (\Gamma_A(f, g), h) \right| &\lesssim \left\{ \|f\|_{L^2_{s+\gamma/2}} \|g\|_{\Phi_\gamma} + \|g\|_{L^2_{s+\gamma/2}} \|f\|_{\Phi_\gamma} \right. \\ &\quad \left. + \min(\|f\|_{L^2} \|g\|_{L^2_{s+\gamma/2}}, \|f\|_{L^2_{s+\gamma/2}} \|g\|_{L^2}) \right\} \|h\|_{\Phi_\gamma}. \end{aligned}$$

We now turn to Γ_B . For this, firstly, it holds that

$$\begin{aligned} \left| (\Gamma_B(f, g), h) \right| &\lesssim \left(\iiint b \Phi_B \mu_*^{1/2} (f_*' g' - f_* g)^2 d\sigma dv dv_* \right)^{1/2} \|h\|_{\Phi_\gamma} \\ &= B^{1/2} \|h\|_{\Phi_\gamma}. \end{aligned}$$

Since $|v - v_*| \leq 1$ implies $|v|^2 \leq 2 + 2|v_*|^2$ and then $\mu_* \leq e\mu^{1/2}$, we have

$$\begin{aligned} B &\lesssim \iiint_{\{|v-v_*| \leq 1\}} b\Phi_\gamma \mu_*^{1/10} \mu^{1/10} ((\mu^{1/8} f)'_* - (\mu^{1/8} f)_*)^2 g^2 d\sigma dv dv_* \\ &\quad + \iiint_{\{|v-v_*| \leq 1\}} b\Phi_\gamma \mu_*^{1/10} \mu^{1/10} ((\mu^{1/8} f)'_*)^2 (g' - g)^2 d\sigma dv dv_* \\ &\quad + \iiint_{\{|v-v_*| \leq 1\}} b\Phi_\gamma \mu_*^{1/10} \mu^{1/10} (\mu_*^{1/8} - \mu_*'^{1/8})^2 (f'_* g')^2 d\sigma dv dv_* \\ &= B_1 + B_2 + B_3. \end{aligned}$$

Obviously,

$$B_1 \lesssim \|\mu^{1/20} g\|_{L^\infty}^2 \|\mu^{1/8} f\|_{\Phi_\gamma} \lesssim \|\mu^{1/20} g\|_{L^\infty}^2 \|f\|_{\Phi_\gamma}.$$

Since $|\mu_*^{1/8} - \mu_*'^{1/8}| \lesssim |v - v_*| \theta$, we see that by the change of variables $(v', v'_*) \rightarrow (v, v_*)$

$$\begin{aligned} B_3 &\lesssim \iiint_{\{|v-v_*| \leq 1\}} b\Phi_\gamma \mu_*^{1/10} \mu^{1/10} (\mu_*'^{1/8} - \mu_*^{1/8})^2 (f_* g)^2 d\sigma dv dv_* \\ &\lesssim \iint_{\{|v-v_*| \leq 1\}} (\mu_*^{1/20} f_*)^2 (\mu^{1/20} g)^2 |v - v_*|^{\gamma+2} \left(\int b(\cos \theta) \theta^2 d\sigma \right) dv dv_* \\ &\lesssim \int (\mu_*^{1/20} f_*)^2 \left(\sup_{v_*} \int \frac{(\mu^{1/20} g)^2}{|v - v_*|^{(-\gamma-2)^+}} dv \right) dv_* \\ &\lesssim \|\mu^{1/20} f\|_{L^{s+\gamma/2}}^2 \|D_v|^{(-\gamma/2-1)^+} \mu^{1/20} g\|_{L^2}^2, \end{aligned}$$

where we have used the Hardy inequality if $\gamma + 2 < 0$, cf. [34]. If one writes

$$\begin{aligned} B_2 &\lesssim \iiint_{\{|v-v_*| \leq 1\}} b\Phi_\gamma \mu_*^{1/10} \mu^{1/20} ((\mu^{1/8} f)'_*)^2 ((\mu^{1/40} g)' - (\mu^{1/40} g))^2 d\sigma dv dv_* \\ &\quad + \iiint_{\{|v-v_*| \leq 1\}} b\Phi_\gamma \mu_*^{1/10} \mu^{1/20} ((\mu^{1/8} f)'_*)^2 g'^2 (\mu^{1/40} - \mu'^{1/40})^2 d\sigma dv dv_* \\ &= B_{2,1} + B_{2,2}, \end{aligned}$$

then the second term $B_{2,2}$ has a similar upper bound as B_3 . It remains to estimate

$$B_{2,1} = \iiint_{\{|v-v_*| \leq 1\}} b\Phi_\gamma \mu_*^{1/10} \mu^{1/20} ((\mu^{1/8} f)_*)^2 ((\mu^{1/40} g) - (\mu^{1/40} g)')^2 d\sigma dv dv_*,$$

by the change of variables $(v', v'_*) \rightarrow (v, v_*)$. By firstly putting $F = \mu^{1/8} f$ and $G = \mu^{1/40} g$, and denoting by $v_\tau = v' + \tau(v - v')$ for $\tau \in [0, 1]$, then by using

$$|G(v) - G(v')|^2 = \left| \int_0^1 \nabla G(v_\tau) \cdot (v - v') d\tau \right|^2 \leq |v - v_*|^2 (\sin^2 \theta/2) \left(\int_0^1 |\nabla G(v_\tau)|^2 d\tau \right),$$

we have

$$B_{2,1} \lesssim \int_0^1 \left\{ \iiint b |v - v_*|^{\gamma+2} \sin^2 \frac{\theta}{2} F_*^2 |\nabla G(v_\tau)|^2 dv dv_* d\sigma \right\} d\tau.$$

To estimate this term, we need the change of variables

$$v \rightarrow v_\tau = \frac{1+\tau}{2}v + \frac{1-\tau}{2}(v - v_*|\sigma + v_*).$$

The Jacobian of this transform is bounded from below uniformly in v_* , σ and τ , because

$$\begin{aligned} \left| \frac{\partial(v_\tau)}{\partial(v)} \right| &= \left| \det \left(\frac{1+\tau}{2} I + \frac{1-\tau}{2} \sigma \otimes \mathbf{k} \right) \right| \quad (\mathbf{k} = \frac{v-v_*}{|v-v_*|}) \\ &= \frac{(1+\tau)^3}{2^3} \left| 1 + \frac{1-\tau}{1+\tau} \mathbf{k} \cdot \sigma \right| = \frac{(1+\tau)^3}{2^3} \left| \frac{2\tau}{1+\tau} + 2 \frac{1-\tau}{1+\tau} \cos^2 \frac{\theta}{2} \right| \\ &\geq \frac{(1+\tau)^3}{2^3} \left| \frac{2\tau}{1+\tau} + \frac{1-\tau}{1+\tau} \right| = \frac{(1+\tau)^3}{2^3} \geq \frac{1}{2^3}. \end{aligned}$$

If we set $\tilde{b} = b(\mathbf{k} \cdot \sigma)(1 - \mathbf{k} \cdot \sigma)$, then we have $\int_{\mathbb{S}^2} \tilde{b} d\sigma < \infty$. Therefore,

$$\begin{aligned} B_{2,1} &\lesssim \int_0^1 \int F_*^2 \left\{ \int_{\mathbb{S}^2} \tilde{b} \int \frac{|\nabla G(v_\tau)|^2}{|v_\tau - v_*|^{-\gamma-2}} dv_\tau d\sigma \right\} dv_* d\tau \\ &\lesssim \int_0^1 \int \langle v_* \rangle^{(\gamma/2+1)^+} F_*^2 \left\{ \int_{\mathbb{S}^2} \tilde{b} \int \frac{\langle v_\tau \rangle^{(\gamma/2+1)^+} |\nabla G(v_\tau)|^2}{|v_\tau - v_*|^{(-\gamma-2)^+}} dv_\tau d\sigma \right\} dv_* d\tau \\ &\lesssim \|F\|_{L^2_{(\gamma/2+1)^+}}^2 \| |D|^{(-\gamma/2-1)^+} \nabla G \|_{L^2_{(\gamma/2+1)^+}}^2 \lesssim \|F\|_{L^2_{(\gamma/2+1)^+}}^2 \|G\|_{H^{\max(-\gamma/2, 1)}_{(\gamma/2+1)^+}}^2, \end{aligned}$$

where we have used $|v - v_*| \sim |v_\tau - v_*|$. Finally we obtain

$$B \lesssim \|\mu^{1/40} f\|_{\Phi_\gamma} \|\mu^{1/40} g\|_{L^\infty} + \|\mu^{1/40} f\|_{L^2} \|\mu^{1/60} g\|_{H^{\max(-\gamma/2, 1)}}.$$

This concludes the proof of Proposition 3.3. \square

Note that the above estimation is good enough for proving the local existence for the general case. However, the above upper bound related to B is given in Sobolev space with positive index and this can not be controlled by the non-isotropic norm. Hence, it is not sufficient for the proof of global existence.

3.2. A simple proof of Theorem 1.2 for $\gamma > -3/2$. We first give a simple proof of upper bound estimates on the Boltzmann nonlinear operator when $\gamma > -\frac{3}{2}$. We state it as

Proposition 3.4. *Assume that $0 < s < 1$ and $\gamma > -3/2$. Then*

$$\begin{aligned} \left| (\Gamma(f, g), h) \right| &\lesssim \left\{ \|f\|_{L^2_{s+\gamma/2}} \|g\|_{\Phi_\gamma} + \|g\|_{L^2_{s+\gamma/2}} \|f\|_{\Phi_\gamma} \right. \\ &\quad \left. + \min \left(\|f\|_{L^2} \|g\|_{L^2_{s+\gamma/2}}, \|f\|_{L^2_{s+\gamma/2}} \|g\|_{L^2} \right) \right\} \|h\|_{\Phi_\gamma}. \end{aligned}$$

Furthermore, together with $\gamma \geq -3s$, one has

$$\left| (\Gamma(f, g), h) \right| \lesssim \left\{ \|f\|_{L^2_{s+\gamma/2}} \|g\|_{\Phi_\gamma} + \|g\|_{L^2_{s+\gamma/2}} \|f\|_{\Phi_\gamma} \right\} \|h\|_{\Phi_\gamma}.$$

Let us note that the first statement deals with general values of $\gamma > -3/2$, that is not necessarily linked with the value of s . For the second statement, note that the condition $\gamma \geq -3s$ is always true in the physical cases mentioned above. Indeed recall here that $\gamma = 1 - 4s$, and that $0 < s < 1$. Therefore, we can conclude that together with the constraint $\gamma > -3/2$, the physical range $0 < s < 5/8$ is allowed.

Proof. **The case when $\gamma \geq 0$.** Note that

$$\begin{aligned}
(3.1) \quad (\Gamma(f, g), h)_{L^2} &= (\mu^{-1/2} Q(\mu^{1/2} f, \mu^{1/2} g), h)_{L^2} \\
&= \iiint \Phi_\gamma b(\cos \theta) \mu_*^{1/2} (f'_* g' - f_* g) h \\
&= \frac{1}{2} \iiint \Phi_\gamma b(\cos \theta) (f'_* g' - f_* g) (\mu_*^{1/2} h - \mu_*^{1/2} h') \\
&\leq \frac{1}{2} \left(\iiint \Phi_\gamma b(\cos \theta) (f'_* g' - f_* g)^2 \right)^{1/2} \\
&\quad \times \left(\iiint \Phi_\gamma b(\cos \theta) (\mu_*^{1/2} h - \mu_*^{1/2} h')^2 \right)^{1/2} \\
&\leq \frac{1}{2} A^{1/2} \times B^{1/2}.
\end{aligned}$$

For B , we have

$$\begin{aligned}
B &= \iiint \Phi_\gamma b(\cos \theta) \left((\mu'_*)^{1/2} (h' - h) + h (\mu'_*)^{1/2} - (\mu_*)^{1/2} \right)^2 \\
&\leq 2 \iiint \Phi_\gamma b(\cos \theta) \left\{ \mu'_* (h' - h)^2 + h^2 (\mu'_*)^{1/2} - (\mu_*)^{1/2} \right\}^2 \\
&\leq 2 \iiint \Phi_\gamma b(\cos \theta) \mu_* (h' - h)^2 + 2 \iiint \Phi_\gamma b(\cos \theta) h_*^2 (\mu'^{1/2} - \mu^{1/2})^2 \\
&= 2 \|h\|_{\Phi_\gamma}^2,
\end{aligned}$$

where we have used the change of variables $(v, v_*) \rightarrow (v', v'_*)$ for the first term and $(v, v_*) \rightarrow (v_*, v)$ for the second term. Similarly,

$$\begin{aligned}
A &= \iiint \Phi_\gamma b(\cos \theta) \left(f'_* (g' - g) + g (f'_* - f_*) \right)^2 \\
&\leq 2 \iiint \Phi_\gamma b(\cos \theta) \left\{ f_*'^2 (g' - g)^2 + g^2 (f'_* - f_*)^2 \right\} \\
&\leq 2 \iiint \Phi_\gamma b(\cos \theta) f_*'^2 (g' - g)^2 + 2 \iiint \Phi_\gamma b(\cos \theta) g_*^2 (f' - f)^2.
\end{aligned}$$

Then (3.2) implies that

$$A \lesssim \|f\|_{L_{s+\gamma/2}^2}^2 \|g\|_{\Phi_\gamma}^2 + \|g\|_{L_{s+\gamma/2}^2}^2 \|f\|_{\Phi_\gamma}^2,$$

which completes the proof in the case when $\gamma \geq 0$.

The case when $-3/2 < \gamma < 0$. As in Subsection 2.3, it is easy to check that for any fixed $\rho > 0$,

$$\begin{aligned}
(3.2) \quad \|g\|_{\Phi_\gamma}^2 &\sim J_{1,\rho}^{\Phi_\gamma}(g) + J_{2,\rho}^{\Phi_\gamma}(g) \sim J_{1,\rho}^{\Phi_\gamma}(g) + \|g\|_{L_{s+\gamma/2}^2}^2 \\
&\sim \iiint \frac{\Phi_{2\gamma}}{\Phi_\gamma} b_{\mu_{\rho,*}} (g' - g)^2 + \iiint \frac{\Phi_{2\gamma}}{\Phi_\gamma} b_{g_*^2} (\sqrt{\mu'_\rho} - \sqrt{\mu_\rho})^2,
\end{aligned}$$

where the assumption $2\gamma > -3$ is required for the existence of the above integral, and more precisely for

$$\int |v_*|^{2\gamma} \langle v_* \rangle^{2s-\gamma} \mu_\rho(v + v_*) dv_* \sim \langle v \rangle^{\gamma+2s}.$$

Instead of (3.1), we write

$$\begin{aligned} (\Gamma(f, g), h) &= \iiint b\Phi_\gamma \mu_*^{1/2} (f'_* g' - f_* g) h d v d v_* d \sigma \\ &= \frac{1}{2} \iiint (b\tilde{\Phi}_\gamma)^{1/2} (f'_* g' - f_* g) \left(b \frac{\Phi_{2\gamma}}{\tilde{\Phi}_\gamma}\right)^{1/2} \mu_*'^{1/4} (\mu_*^{1/4} h - \mu_*'^{1/4} h') \\ &\quad + \frac{1}{2} \iiint (b\tilde{\Phi}_\gamma)^{1/2} (f'_* g' - f_* g) \left(b \frac{\Phi_{2\gamma}}{\tilde{\Phi}_\gamma}\right)^{1/2} \mu_*^{1/4} (\mu_*^{1/4} - \mu_*'^{1/4}) h. \end{aligned}$$

Noticing that

$$\mu_*^{1/4} h - \mu_*'^{1/4} h' = \mu_*'^{1/4} (h - h') + (\mu_*^{1/4} - \mu_*'^{1/4}) h,$$

by Cauchy-Schwarz's inequality and (3.2), we have

$$\begin{aligned} |(\Gamma(f, g), h)| &\lesssim \left(\iiint b\tilde{\Phi}_\gamma \mu_*^{1/2} (f'_* g' - f_* g)^2 d\sigma d v d v_* \right)^{1/2} \|h\|_{\Phi_\gamma} \\ &= A^{1/2} \|h\|_{\Phi_\gamma}. \end{aligned}$$

We estimate

$$\begin{aligned} A &\leq 3 \left(\iiint b\tilde{\Phi}_\gamma \mu_*^{1/4} \left((\mu^{1/8} f)'_* - (\mu^{1/8} f)_* \right)^2 g^2 d\sigma d v d v_* \right. \\ &\quad + \iiint b\tilde{\Phi}_\gamma \mu_*^{1/8} \left((\mu^{1/8} f)'_* \right)^2 (g' - g)^2 d\sigma d v d v_* \\ &\quad \left. + \iiint b\tilde{\Phi}_\gamma \mu_*^{1/4} \left(\mu_*^{1/8} - \mu_*'^{1/8} \right)^2 (f'_* g')^2 d\sigma d v d v_* \right) \\ &= A_1 + A_2 + A_3. \end{aligned}$$

Since $\tilde{\Phi}_\gamma(|v - v_*|) \mu_*^{1/4} \lesssim \langle v \rangle^\gamma$, we have by means of Corollary 2.18

$$\begin{aligned} A_1 &\lesssim \iiint b \langle v \rangle^{\gamma/2} g^2 \left((\mu^{1/8} f)'_* - (\mu^{1/8} f)_* \right)^2 d\sigma d v d v_* \\ &\lesssim \|\langle v \rangle^{\gamma/2} g\|_{L^2_s}^2 \|\mu^{1/8} f\|_{\Phi_0}^2 \lesssim \|g\|_{L^2_{s+\gamma/2}}^2 \|f\|_{\Phi_\gamma}^2, \end{aligned}$$

where we have used Propositions 2.2 and 2.4 in the last inequality. As for A_2 , we decompose it as follows :

$$\begin{aligned} A_2 &\lesssim \iiint b \left((\mu^{1/8} f)'_* \right)^2 \left(\langle v \rangle^{\gamma/2} g' - \langle v \rangle^{\gamma/2} g \right)^2 d\sigma d v d v_* \\ &\quad + \iiint b \left(\langle v \rangle^{\gamma/2} - \langle v' \rangle^{\gamma/2} \right)^2 \left((\mu^{1/8} f)'_* \right)^2 g'^2 d\sigma d v d v_* \\ &= A_{2,1} + A_{2,2}. \end{aligned}$$

Apply Corollary 2.18 again to $A_{2,1}$. Then

$$A_{2,1} \lesssim \|\mu^{1/8} f\|_{L^2_s}^2 \|\langle v \rangle^{\gamma/2} g\|_{\Phi_0}^2 \lesssim \|f\|_{L^2_{s+\gamma/2}}^2 \|g\|_{\Phi_\gamma}^2.$$

The estimation for $A_{2,2}$ is the same as the one for A_2 in the proof of Lemma 2.12. By using the change of variables $(v', v'_*) \rightarrow (v, v_*)$, we obtain

$$A_{2,2} \lesssim \iint \left(\langle v \rangle^{2s+\gamma} \langle v_* \rangle^{2s+2} \right) (\mu^{1/8} f)_*^2 |g|^2 d v d v_* \lesssim \|\mu^{1/10} f\|_{L^2}^2 \|g\|_{L^2_{s+\gamma/2}}^2.$$

Noticing that $(\mu_*^{1/8} - \mu_*'^{1/8})^2 \lesssim \min(|v - v_*|^2 \theta^2, 1)$, we have

$$\begin{aligned} A_3 &\lesssim \iint \tilde{\Phi}_\gamma \left(\int_{\mathbb{S}^2} b(\cos \theta) \min(|v - v_*|^2 \theta^2, 1) d\sigma \right) f_*^2 g^2 dv dv_* \\ &\lesssim \iint \langle v - v_* \rangle^{\gamma+2s} f_*^2 g^2 dv dv_* \\ &\lesssim \iint \langle v_* \rangle^{\gamma+2s} f_*^2 \langle v \rangle^{\gamma+2s} g^2 dv dv_* \lesssim \|f\|_{L^2_{s+\gamma/2}}^2 \|g\|_{L^2_{s+\gamma/2}}^2, \end{aligned}$$

if $\gamma + 2s \geq 0$ because of $\langle v - v_* \rangle^{\gamma+2s} \leq \langle v_* \rangle^{\gamma+2s} \langle v \rangle^{\gamma+2s}$.

To consider the case $\gamma + 2s < 0$, we divide $\mathbb{R}_v^3 \times \mathbb{R}_{v_*}^3$ into three parts

$$\begin{aligned} U_1 &= \{|v - v_*| \leq |v_*|/8\}, \quad U_2 = \{|v - v_*| > |v_*|/8\} \cap \{|v_*| \leq 1\}, \\ U_3 &= \{|v - v_*| > |v_*|/8\} \cap \{|v_*| > 1\}. \end{aligned}$$

Then we have

$$\begin{aligned} \frac{1}{3} A_3 &= \iiint b \tilde{\Phi}_\gamma \mu_*'^{1/4} (\mu_*'^{1/8} - \mu_*^{1/8})^2 (f_* g)^2 d\sigma dv dv_* \\ &= \iint_{U_1} \int d\sigma dv dv_* + \iint_{U_2} \int d\sigma dv dv_* + \iint_{U_3} \int d\sigma dv dv_* \\ &= A_{3,1} + A_{3,2} + A_{3,3}. \end{aligned}$$

Since $|v' - v_*| \leq |v - v_*| \leq |v_*|/8$ implies $7|v_*|/8 \leq |v'|$, $|v| \leq 9|v_*|/8$ and $|v_*'|^2 = |v|^2 + |v_*|^2 - |v'|^2 \geq |v_*|^2/2$. Hence, we have $\mu_*'^{1/4} \leq C\mu_*^{1/8} \leq C(\mu_*\mu)^{1/20}$ on U_1 , which leads to

$$A_{3,1} \lesssim \iint (\mu\mu_*)^{1/20} \langle v - v_* \rangle^{\gamma+2s} f_*^2 g^2 dv dv_* \leq C \|f\|_{L^2_{s+\gamma/2}}^2 \|g\|_{L^2_{s+\gamma/2}}^2.$$

Furthermore, we have

$$A_{3,2} \lesssim \iint_{U_2} \langle v - v_* \rangle^{\gamma+2s} f_*^2 g^2 dv dv_* \lesssim \|f\|_{L^2_{s+\gamma/2}}^2 \|g\|_{L^2_{s+\gamma/2}}^2,$$

because $\langle v - v_* \rangle^{-1} \leq \langle v \rangle^{-1} \langle v_* \rangle^{-1} \langle v_* \rangle^2 \leq 2\langle v \rangle^{-1} \langle v_* \rangle^{-1}$ on U_2 . Since $\langle v - v_* \rangle^{-1} \leq 8|v_*|^{-1} \leq 16\langle v_* \rangle^{-1}$ on U_3 , we get

$$A_{3,3} \lesssim \|f\|_{L^2_{s+\gamma/2}}^2 \|g\|_{L^2}^2.$$

Therefore, we have in the case when $\gamma + 2s < 0$

$$A_3 \lesssim \|f\|_{L^2_{s+\gamma/2}}^2 \|g\|_{L^2}^2.$$

If one considers another partition in R_{v, v_*}^6 with v and v_* exchanged, then the estimate

$$A_3 \lesssim \|f\|_{L^2}^2 \|g\|_{L^2_{s+\gamma/2}}^2$$

holds, because $|v_*' - v| \leq |v_* - v| \leq |v|/8$ implies $7|v|/8 \leq |v_*'|$, $|v_*| \leq 9|v|/8$.

As a conclusion, when $\gamma > -3/2$ and $\gamma + 2s \leq 0$ we have

$$\begin{aligned} \left| (\Gamma(f, g), h) \right| &\lesssim \left\{ \|f\|_{L^2_{s+\gamma/2}} \|g\|_{\Phi_\gamma} + \|g\|_{L^2_{s+\gamma/2}} \|f\|_{\Phi_\gamma} \right. \\ &\quad \left. + \min \left(\|f\|_{L^2} \|g\|_{L^2_{s+\gamma/2}}, \|f\|_{L^2_{s+\gamma/2}} \|g\|_{L^2} \right) \right\} \|h\|_{\Phi_\gamma}. \end{aligned}$$

which concludes the proof of the first statement of Proposition 3.4.

The case $\gamma + 2s < 0$, $\gamma \geq -3s$. We go back to the definition of A_3 , that is (we have performed the usual change of variables)

$$\begin{aligned} A_3 &\sim \iiint b\tilde{\Phi}_\gamma \mu_*^{1/4} (\mu_*^{1/8} - \mu_*^{1/8})^2 f_*^2 g^2 d\sigma dv dv_* \\ &\lesssim \iiint b\tilde{\Phi}_\gamma (\mu_*^{1/8} - \mu_*^{1/8})^2 f_*^2 g^2 d\sigma dv dv_*. \end{aligned}$$

We estimate the spherical integral as usual, that is over the sets

$$|v'_* - v_*| \leq \frac{1}{2} \langle v_* \rangle \quad \text{and} \quad |v'_* - v_*| \geq \frac{1}{2} \langle v_* \rangle$$

It follows by Taylor formula that, on the first set (which is the singular part), one has, for another non important and non negative constant c

$$(\mu_*^{1/8} - \mu_*^{1/8})^2 \lesssim \theta^2 |v - v_*|^2 \mu_*^c.$$

On the other set, we just estimate the square by 1. Note that on the second set we have, $|v - v_*| \gtrsim \langle v_* \rangle$.

Then we find, by now standard computations, that

$$\begin{aligned} A_3 &\lesssim \iint \langle v - v_* \rangle^\gamma \mu_*^c |v - v_*|^{2s} \langle v_* \rangle^{2-2s} f_*^2 g^2 \\ &+ \iint \langle v - v_* \rangle^\gamma |v - v_*|^{2s} \langle v_* \rangle^{-2s} f_*^2 g^2 \\ &= \tilde{A}_{3,1} + \tilde{A}_{3,2} \end{aligned}$$

Now, for $\tilde{A}_{3,1}$, we write $\langle v - v_* \rangle^{\gamma+2s} \lesssim \langle v \rangle^{\gamma+2s} \langle v_* \rangle^{-\gamma-2s}$ and we see that we may absorb all the powers of $\langle v_* \rangle$ with the maxwellian, to get, for another non negative constant d

$$\tilde{A}_{3,1} \lesssim \|\mu^d f\|_{L^2}^2 \|g\|_{L^2_{\gamma/2+s}}^2.$$

For $\tilde{A}_{3,2}$, we write

$$\langle v - v_* \rangle^{\gamma+2s} \langle v_* \rangle^{-2s} \lesssim \langle v \rangle^{\gamma+2s} \langle v_* \rangle^{-\gamma-2s} \langle v_* \rangle^{-2s} = \langle v \rangle^{\gamma+2s} \langle v_* \rangle^{-\gamma-4s}.$$

Note that the power $-\gamma - 4s$ which enters the power over $\langle v_* \rangle$ can be written

$$-\gamma - 4s = -(\gamma + 2s) - 2s$$

the first term being positive. Of course $-\gamma - 4s \leq 0$ iff $\gamma \geq -4s$, and this is true since we have assumed that $\gamma \geq -3s$. Furthermore $\gamma + 4s \geq -\gamma - 2s$ again because $\gamma \geq -3s$. Therefore we obtained

$$\tilde{A}_{3,2} \lesssim \|f\|_{L^2_{\gamma/2+s}}^2 \|g\|_{L^2_{\gamma/2+s}}^2,$$

concluding the proof of the second statement. \square

Let us note that the proof of Proposition 3.4 gives the following corollary.

Corollary 3.5. *With the regularized potential, together with assumptions (1.2), $0 < s < 1$ and $\gamma \geq -3s$, one has*

$$\left| (\Gamma_{\tilde{\Phi}}(f, g), h) \right| \lesssim \left\{ \|f\|_{L^2_{s+\gamma/2}} \|g\|_{\Phi_\gamma} + \|g\|_{L^2_{s+\gamma/2}} \|f\|_{\Phi_\gamma} \right\} \|h\|_{\Phi_\gamma}.$$

Note carefully that in this last result, the constraint $\gamma > -3/2$ is removed, and we have retained the constraint $\gamma \geq -3s$, which is always true for physical cases, as we saw above.

3.3. Proof of Theorem 1.2. The general case is long and will be divided into several steps. One of the key ingredients in the proof is to split the term $(\Gamma(f, g); h)$ into two parts. In this way, one part can be dealt with the method introduced in the previous subsection, while the other one will be analyzed by direct Fourier transform.

Lemma 3.6. *For any integer $k \geq 2$ we can write*

$$\begin{aligned} \mu_*^{1/2} &= (\mu^{a^1} - \mu_*^{a^1})^k \sum_{i=1}^{k+2} \alpha_i^2 \mu_*^{a_i^2} \mu^{b_i^2} + \sum_{i=1}^k \alpha_i^3 \mu_*^{a_i^3} \mu^{b_i^3} \\ &\stackrel{\text{def}}{=} \mu(v, v_*) + \sum_{i=1}^k \alpha_i^3 \mu_*^{a_i^3} \mu^{b_i^3}. \end{aligned}$$

In the above, α_i^j are real numbers for all i and j , and the other exponents are strictly positive, at the exception of $b_i^2 = 0$, with $b_i^3 > a_i^3$.

Proof. Differentiating $k-1$ times the identity $\sum_{j=0}^{2k} x^j = \frac{1-x^{2k+1}}{1-x}$ we have

$$\begin{aligned} \sum_{j=k-1}^{2k} \frac{j!}{(j-k+1)!} x^{j-k+1} &= \left(\frac{1}{1-x}\right)^{(k-1)} - \sum_{j=0}^{k-1} \binom{k-1}{j} \left(\frac{1}{1-x}\right)^{(j)} (x^{2k+1})^{(k-1-j)} \\ &= \frac{(k-1)!}{(1-x)^k} \left\{ 1 - \sum_{j=0}^{k-1} \frac{(2k+1)!}{j!} \left(\sum_{n=0}^{k-j-1} \frac{(-1)^n}{n!(2k-j-n+1)!} \right) x^{2k-j+1} \right\}. \end{aligned}$$

By setting $x = B/A$ and multiplying the above identity by $A^{2k+2}(1-B/A)^k/(k-1)!$, we obtain

$$\begin{aligned} A^{2k+2} &= (A-B)^k \sum_{j=k-1}^{2k} \frac{j!}{(j-k+1)!} A^{2k-j+1} B^{j-k+1} \\ &\quad + \sum_{j=0}^{k-1} \frac{(2k+1)!}{j!} \left(\sum_{n=0}^{k-j-1} \frac{(-1)^n}{n!(2k-j-n+1)!} \right) A^{j+1} B^{2k-j+1}, \end{aligned}$$

which gives the desired formula. \square

With the help of Lemma 3.6, we can analyze $(\Gamma(f, g); h)$ as follows. Write

$$(\Gamma(f, g), h) = (\Gamma_\mu(f, g), h) + (\Gamma_{\text{rest}}(f, g), h),$$

with

$$(3.3) \quad (\Gamma_\mu(f, g), h) = \iiint b\left(\frac{v-v_*}{|v-v_*|} \cdot \sigma\right) \Phi_\gamma(v-v_*) \mu(v, v_*) (f'_* g' - f_* g) h v d v_* d \sigma,$$

and $(\Gamma_{\text{rest}}(f, g); h)$ is a finite linear combination of terms in the form of

$$\begin{aligned} (3.4) \quad (\Gamma_{\text{mod}, i}(f, g), h) &= \iiint b\left(\frac{v-v_*}{|v-v_*|} \cdot \sigma\right) \Phi_\gamma(v-v_*) (f'_* g' - f_* g) \mu_*^{c_i} \mu^{d_i} h v d v_* d \sigma \\ &= (Q(\mu^{c_i} f, \mu^{c_i} g), \mu^{d_i} h), \end{aligned}$$

with $d_i > 0$, $c_i > 0$, for $1 \leq i \leq 3$.

The following two propositions give estimates on each of these scalar products, and all together imply Theorem 1.2.

Proposition 3.7. *For all $0 < s < 1$ and $\gamma > -3$, one has*

$$\begin{aligned} |(\Gamma_\mu(f, g); h)| &\lesssim \left\{ \|f\|_{L^2_{s+\gamma/2}} \|g\|_{\Phi_\gamma} + \|g\|_{L^2_{s+\gamma/2}} \|f\|_{\Phi_\gamma} \right. \\ &\quad \left. + \min\{\|f\|_{L^2} \|g\|_{L^2_{s+\gamma/2}}, \|f\|_{L^2_{s+\gamma/2}} \|g\|_{L^2}\} \right\} \|h\|_{\Phi_\gamma}. \end{aligned}$$

Proof. Since $\mu(v, v_*)$ is a finite sum of $(\mu^a - \mu_*^a)^k \mu_*^{2b} \mu^c$ with $a, b > 0$ and $c \geq 0$, by setting $H = \mu^c h$, we can write

$$\begin{aligned} (\Gamma_\mu(f, g), h) &= \iiint b \Phi_\gamma (\mu^a - \mu_*^a)^k \mu_*^b (f'_* g' - f_* g) \mu_*^b H d\sigma dv dv_* \\ &= \frac{1}{2} \iiint b \Phi_\gamma (\mu^a - \mu_*^a)^k \mu_*^b (f'_* g' - f_* g) (\mu_*^b H - \mu_*'^b H') d\sigma dv dv_* \\ &\quad + \frac{1}{2} \iiint b \Phi_\gamma (f'_* g' - f_* g) \left\{ (\mu'^a - \mu_*'^a)^k \mu_*'^b - (\mu^a - \mu_*^a)^k \mu_*^b \right\} \mu_*'^b H' d\sigma dv dv_*. \end{aligned}$$

By setting $\Phi_\gamma = \tilde{\Phi}_\gamma^{1/2} (\Phi_\gamma \tilde{\Phi}_\gamma^{-1/2})$, the Cauchy-Schwarz inequality gives

$$|(\Gamma_\mu(f, g), h)| \leq \frac{1}{2} A^{1/2} (D^{1/2} + E^{1/2}),$$

where

$$\begin{aligned} A &= \iiint b \tilde{\Phi}_\gamma \mu_*^{2b} (f'_* g' - f_* g)^2 d\sigma dv dv_*, \\ D &= \iiint b (\Phi_\gamma \tilde{\Phi}_\gamma^{-1/2} (\mu^a - \mu_*^a)^k)^2 (\mu_*^b H - \mu_*'^b H')^2 d\sigma dv dv_*, \\ E &= \iiint b (\Phi_\gamma \tilde{\Phi}_\gamma^{-1/2} \{ (\mu'^a - \mu_*'^a)^k \mu_*'^b - (\mu^a - \mu_*^a)^k \mu_*^b \})^2 H^2 d\sigma dv dv_*. \end{aligned}$$

It is easy to see $D + E \lesssim \|H\|_{\Phi_\gamma}^2$ because $(\Phi_\gamma \tilde{\Phi}_\gamma^{-1/2} (\mu^a - \mu_*^a)^k)^2 \lesssim \Phi_\gamma$, and

$$(\Phi_\gamma \tilde{\Phi}_\gamma^{-1/2} \{ (\mu'^a - \mu_*'^a)^k \mu_*'^b - (\mu^a - \mu_*^a)^k \mu_*^b \})^2 \lesssim \Phi_\gamma (\mu_*^b + \mu_*'^b) \min\{|v - v_*|^2 \theta^2, 1\}.$$

The estimation on A is just the same as the one in the proof of Proposition 3.3. And this completes the proof of the proposition. \square

To estimate the last term in (3.4), we have the following proposition.

Proposition 3.8. *For all $0 < s < 1$ and $\gamma > \max\{-3, -3/2 - 2s\}$, one has*

$$|(\Gamma_{mod,i}(f, g), h)| \lesssim \left\{ \|f \mu^{c_1}\|_{L^2_{s+\gamma/2}} \|g \mu^{c_2}\|_{\Phi_\gamma} + \|g \mu^{c_1}\|_{L^2_{s+\gamma/2}} \|f \mu^{c_2}\|_{\Phi_\gamma} \right\} \|h \mu^{d_1}\|_{\Phi_\gamma}.$$

For this, we consider

$$(\Gamma_{mod,i}(f, g), h) = (Q(\mu^{c_1} f, \mu^{c_2} g), \mu^{d_1} h).$$

We therefore consider $(Q(F, G), H)$ with

$$(3.5) \quad F = f \mu^{c_1}, \quad G = g \mu^{c_2}, \quad H = h \mu^{c_3},$$

for some positive constants c_1, c_2 and c_3 .

Let $0 \leq \phi(v) \leq 1$ be a smooth radial function with value 1 for v close to 0, and 0 for large values of v . Set

$$\Phi_\gamma(v) = \Phi_\gamma(v) \phi(v) + \Phi_\gamma(v) (1 - \phi(v)) = \Phi_c(v) + \Phi_{\bar{c}}(v).$$

And then correspondingly we can write

$$Q(F, G) = Q_c(F, G) + Q_{\bar{c}}(F, G),$$

where the kinetic factor in the collision operator is defined according to the decomposition respectively. To prove Proposition 3.8, it suffices to prove the following two lemmas, by taking $m = -s$ in the statements. The general form below for any real m will be needed in Part II.

Lemma 3.9. *For all $0 < s < 1$ and $\gamma > -3$, one has*

$$\left| (Q_{\bar{c}}(F, G), H) \right| \leq C \left\{ \|F\|_{L^2_{s+\gamma/2}} \|G\|_{\Phi_\gamma} + \|G\|_{L^2_{s+\gamma/2}} \|F\|_{\Phi_\gamma} \right\} \|H\|_{\Phi_\gamma}.$$

Proof. One has for some positive constant β

$$\begin{aligned} |(Q_{\bar{c}}(F, G); H)| &= \left| \iiint b \Phi_{\bar{c}}(v - v_*) \mu_*^\beta \mu^\beta [F'_* G' - F_* G] H dv dv_* d\sigma \right| \\ &= \frac{1}{2} \left| \iiint b \Phi_{\bar{c}}(v - v_*) [F'_* G' - F_* G] \mu_*^\beta \mu^\beta [H' - H] dv dv_* d\sigma \right| \\ &\lesssim A^{1/2} B^{1/2}, \end{aligned}$$

where

$$A = \iiint b \Phi_{\bar{c}}(v - v_*) [F'_* G' - F_* G]^2 \mu_*^\beta \mu^\beta dv dv_* d\sigma,$$

and

$$B = \iiint b \Phi_{\bar{c}}(v - v_*) \mu_*^\beta \mu^\beta [H' - H]^2 dv dv_* d\sigma.$$

B is clearly estimated from above by the dissipative norm of H , while for A , we note that $\Phi_{\bar{c}} \lesssim \tilde{\Phi}_\gamma$. The proof of Proposition 3.3 in Subsection 3.1 can then be applied to A . And this gives the desired estimate and then completes the proof of the lemma. \square

Next, let us note that, from the Appendix, $|\hat{\Phi}_c(\xi)| \lesssim \langle \xi \rangle^{-3-\gamma}$. We shall prove

Lemma 3.10. *Let $m \in \mathbb{R}$. For all $0 < s < 1$ and $\gamma > \max\{-3, -\frac{3}{2} - 2s\}$, one has*

$$|(Q_c(F, G); H)| \lesssim \left(\|F\|_{L^2} \|G\|_{H^{(m+2s)^+}} + \|F\|_{H^{s+(m+s)^+}} \|G\|_{L^2} + \|F\|_{H^s} \|G\|_{H^{(m+s)^+}} \right) \|H\|_{H^{-m}}.$$

It is important to note that even though the statement of Lemma 3.10 is not as sharp as the one of Lemma 3.9, by recalling (3.5), we have all the needed weights because we are dealing with functions of the form F , G and H that contain Gaussians. Hence, these two lemmas together imply Proposition 3.8.

For the proof of Lemma 3.10, first of all, by using the formula from the Appendix of [3], we have

$$\begin{aligned} \mathcal{F}(Q_c(F, G))(\xi) &= \iint_{\mathbb{R}_{\xi_*}^3 \times \mathbb{S}^2} b \left(\frac{\xi}{|\xi|} \cdot \sigma \right) \hat{\Phi}_c(\xi_*) \hat{F}(\xi^- + \xi_*) \hat{G}(\xi^+ - \xi_*) \\ &\quad - \iint_{\mathbb{R}_{\xi_*}^3 \times \mathbb{S}^2} b \left(\frac{\xi}{|\xi|} \cdot \sigma \right) \hat{\Phi}_c(\xi_*) \hat{F}(\xi_*) \hat{G}(\xi - \xi_*). \end{aligned}$$

We change variables in ξ_* in the first integral to obtain

$$\begin{aligned} (Q_c(F, G); H) &= \iiint b \left(\frac{\xi}{|\xi|} \cdot \sigma \right) [\hat{\Phi}_c(\xi_* - \xi^-) - \hat{\Phi}_c(\xi_*)] \hat{F}(\xi_*) \hat{G}(\xi - \xi_*) \overline{\hat{H}(\xi)} d\xi d\xi_* d\sigma \\ &= \iiint_{|\xi^-| \leq \frac{1}{2} \langle \xi_* \rangle} \dots d\xi d\xi_* d\sigma + \iiint_{|\xi^-| \geq \frac{1}{2} \langle \xi_* \rangle} \dots d\xi d\xi_* d\sigma \\ &= A_1(F, G, H) + A_2(F, G, H). \end{aligned}$$

$A_2(F, G, H)$ can be naturally decomposed into

$$\begin{aligned} A_2 &= \iiint b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) \mathbf{1}_{|\xi^-| \geq \frac{1}{2}\langle \xi_* \rangle} \hat{\Phi}_c(\xi_* - \xi^-) \hat{F}(\xi_*) \hat{G}(\xi - \xi_*) \overline{\hat{H}(\xi)} d\xi d\xi_* d\sigma \\ &\quad - \iiint b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) \mathbf{1}_{|\xi^-| \geq \frac{1}{2}\langle \xi_* \rangle} \hat{\Phi}_c(\xi_*) \hat{F}(\xi_*) \hat{G}(\xi - \xi_*) \overline{\hat{H}(\xi)} d\xi d\xi_* d\sigma \\ &= A_{2,1}(F, G, H) - A_{2,2}(F, G, H). \end{aligned}$$

For A_1 , we use the Taylor expansion of $\hat{\Phi}_c$ at order 2 to have

$$A_1 = A_{1,1}(F, G, H) + A_{1,2}(F, G, H)$$

where

$$A_{1,1} = \iiint b \xi^- \cdot (\nabla \hat{\Phi}_c)(\xi_*) \mathbf{1}_{|\xi^-| \leq \frac{1}{2}\langle \xi_* \rangle} \hat{F}(\xi_*) \hat{G}(\xi - \xi_*) \overline{\hat{H}(\xi)} d\xi d\xi_* d\sigma,$$

and $A_{1,2}(F, G, H)$ is the remaining term corresponding to the second order term in the Taylor expansion of $\hat{\Phi}_c$. The $A_{i,j}$ with $i, j = 1, 2$ are estimated by the following lemmas.

Lemma 3.11. *Let $m \in \mathbb{R}$. For all $\gamma > \max(-3, -\frac{3}{2} - 2s)$, one has*

$$|A_{1,j}| \lesssim (\|F\|_{L^2} \|G\|_{H^{(m+2s)^+}} + \|F\|_{H^{s+(m+s)^+}} \|G\|_{L^2} + \|F\|_{H^s} \|G\|_{H^{(m+s)^+}}) \|H\|_{H^{-m}}, \quad j = 1, 2.$$

Proof. We first consider $A_{1,1}$. By writing

$$\xi^- = \frac{|\xi|}{2} \left(\left(\frac{\xi}{|\xi|} \cdot \sigma \right) \frac{\xi}{|\xi|} - \sigma \right) + \left(1 - \left(\frac{\xi}{|\xi|} \cdot \sigma \right) \right) \frac{\xi}{2},$$

we see that the integral corresponding to the first term on the right hand side vanishes because of the symmetry on \mathbb{S}^2 . Hence, we have

$$A_{1,1} = \iint_{\mathbb{R}^6} K(\xi, \xi_*) \hat{F}(\xi_*) \hat{G}(\xi - \xi_*) \overline{\hat{H}(\xi)} d\xi d\xi_*,$$

where

$$K(\xi, \xi_*) = \int_{\mathbb{S}^2} b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) \left(1 - \left(\frac{\xi}{|\xi|} \cdot \sigma \right) \right) \frac{\xi}{2} \cdot (\nabla \hat{\Phi}_c)(\xi_*) \mathbf{1}_{|\xi^-| \leq \frac{1}{2}\langle \xi_* \rangle} d\sigma.$$

Note that $|\nabla \hat{\Phi}_c(\xi_*)| \lesssim \frac{1}{\langle \xi_* \rangle^{3+\gamma+1}}$, from the Appendix. If $\sqrt{2}|\xi| \leq \langle \xi_* \rangle$, then $|\xi^-| \leq \langle \xi_* \rangle/2$ and this implies that for $0 \leq \theta \leq \pi/2$,

$$\begin{aligned} |K(\xi, \xi_*)| &\lesssim \int_0^{\pi/2} \theta^{1-2s} d\theta \frac{\langle \xi \rangle}{\langle \xi_* \rangle^{3+\gamma+1}} \lesssim \frac{\langle \xi \rangle^s \langle \xi_* \rangle^s}{\langle \xi_* \rangle^{3+\gamma+2s}} \left(\frac{\langle \xi \rangle}{\langle \xi_* \rangle} \right)^{1-s} \\ &\lesssim \frac{\langle \xi \rangle^{m+s} \langle \xi_* \rangle^s}{\langle \xi_* \rangle^{3+\gamma+2s}} \langle \xi \rangle^{-m} \lesssim \langle \xi_* \rangle^s \frac{\langle \xi_* \rangle^{(m+s)^+} + \langle \xi - \xi_* \rangle^{(m+s)^+}}{\langle \xi_* \rangle^{3+\gamma+2s}} \langle \xi \rangle^{-m} \end{aligned}$$

On the other hand, if $\sqrt{2}|\xi| \geq \langle \xi_* \rangle$, then

$$\begin{aligned} |K(\xi, \xi_*)| &\lesssim \int_0^{\pi\langle \xi_* \rangle/(2|\xi|)} \theta^{1-2s} d\theta \frac{\langle \xi \rangle}{\langle \xi_* \rangle^{3+\gamma+1}} \lesssim \frac{\langle \xi \rangle^{2s}}{\langle \xi_* \rangle^{3+\gamma+2s}} \frac{\langle \xi_* \rangle}{\langle \xi \rangle} \\ &\lesssim \frac{(\langle \xi - \xi_* \rangle^{(m+2s)^+} + \langle \xi_* \rangle^{(m+2s)^+})}{\langle \xi_* \rangle^{3+\gamma+2s}} \langle \xi \rangle^{-m}. \end{aligned}$$

Since $\langle \xi_* \rangle^{-(3+\gamma+2s)} \in L^2$ when $\gamma > -3/2 - 2s$, we obtain the desired estimate for $A_{1,1}$.

Now we turn to $A_{1,2}(F, G, H)$, which comes from the second order term of the Taylor expansion. Note that

$$A_{1,2} = \iiint b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) \int_0^1 d\tau (\nabla^2 \hat{\Phi}_c)(\xi_* - \tau \xi^-) \cdot \xi^- \cdot \xi^- \hat{F}(\xi_*) \hat{G}(\xi - \xi_*) \bar{H}(\xi) d\sigma d\xi d\xi_*.$$

From the Appendix, we have

$$|(\nabla^2 \hat{\Phi}_c)(\xi_* - \tau \xi^-)| \lesssim \frac{1}{\langle \xi_* - \tau \xi^- \rangle^{3+\gamma+2}} \lesssim \frac{1}{\langle \xi_* \rangle^{3+\gamma+2}},$$

because $|\xi^-| \leq \langle \xi_* \rangle / 2$. Similar to $A_{1,1}$, we can obtain

$$|A_{1,2}| \lesssim \iint_{\mathbb{R}^6} \tilde{K}(\xi, \xi_*) \hat{F}(\xi_*) \hat{G}(\xi - \xi_*) \bar{H}(\xi) d\xi d\xi_*,$$

where $\tilde{K}(\xi, \xi_*)$ has the following upper bound

$$\tilde{K}(\xi, \xi_*) \lesssim \int_0^{\min(\pi/2, \pi \langle \xi_* \rangle / (2|\xi|))} \theta^{1-2s} d\theta \frac{\langle \xi \rangle^2}{\langle \xi_* \rangle^{3+\gamma+2}} \lesssim \frac{\langle \xi \rangle^{2s}}{\langle \xi_* \rangle^{3+\gamma+2s}},$$

which yields the desired estimate for $A_{1,2}$. And this completes the proof of the lemma. \square

Lemma 3.12. *Let $m \in \mathbb{R}$. For all $0 < s < 1$ and $\gamma > \max(-3, -\frac{3}{2} - 2s)$, one has*

$$|A_{2,1}| + |A_{2,2}| \lesssim (\|F\|_{L^2} \|G\|_{H^{(m+2s)^+}} + \|F\|_{H^{(m+2s)^+}} \|G\|_{L^2}) \|H\|_{H^{-m}}.$$

Proof. In view of the definition of $A_{2,2}$, Since we assume that $\theta \geq 1/2 |\xi_*| |\xi|^{-1}$, we also have $1/2 \langle \xi_* \rangle |\xi|^{-1} \leq \frac{\pi}{2}$, that is, $\langle \xi_* \rangle \leq |\xi|$. We can then directly compute the spherical integral appearing inside $A_{2,2}$ together with Φ by using the inequality

$$\frac{1}{\langle \xi_* \rangle^{3+\gamma}} \frac{\langle \xi \rangle^{2s}}{\langle \xi_* \rangle^{2s}} \lesssim \frac{1}{\langle \xi_* \rangle^{3+\gamma+2s}} \langle \xi \rangle^{-m} \left(\langle \xi_* \rangle^{(m+2s)^+} + \langle \xi - \xi_* \rangle^{(m+2s)^+} \right),$$

to obtain the estimate for $A_{2,2}$.

We now turn to

$$A_{2,1} = \iiint b \mathbf{1}_{|\xi^-| \geq \frac{1}{2} \langle \xi_* \rangle} \hat{\Phi}_c(\xi_* - \xi^-) \hat{F}(\xi_*) \hat{G}(\xi - \xi_*) \bar{H}(\xi) d\sigma d\xi d\xi_*.$$

Firstly, note that we can work on the set $|\xi_* \cdot \xi^-| \geq \frac{1}{2} |\xi^-|^2$. In fact, on the complementary of this set, we have $|\xi_* \cdot \xi^-| \leq \frac{1}{2} |\xi^-|^2$ so that $|\xi_* - \xi^-| \geq |\xi_*|$, and in this case, we can proceed in the same way as for $A_{2,2}$. Therefore, it suffices to estimate

$$\begin{aligned} A_{2,1,p} &= \iiint b \mathbf{1}_{|\xi^-| \geq \frac{1}{2} \langle \xi_* \rangle} \mathbf{1}_{|\xi_* \cdot \xi^-| \geq \frac{1}{2} |\xi^-|^2} \hat{\Phi}_c(\xi_* - \xi^-) \hat{F}(\xi_*) \hat{G}(\xi - \xi_*) \overline{\hat{H}(\xi)} d\sigma d\xi d\xi_* \\ &= \iiint b \frac{\mathbf{1}_{|\xi^-| \geq \frac{1}{2} \langle \xi_* \rangle} \mathbf{1}_{|\xi_* \cdot \xi^-| \geq \frac{1}{2} |\xi^-|^2}}{\left(\langle \xi_* \rangle^{(m+2s)^+} + \langle \xi - \xi_* \rangle^{(m+2s)^+} \right)} \hat{\Phi}_c(\xi_* - \xi^-) \\ &\quad \times \overline{\hat{H}(\xi)} \sum_{j=1}^2 \hat{F}_j(\xi_*) \hat{G}_j(\xi - \xi_*) d\sigma d\xi d\xi_*, \end{aligned}$$

where $\hat{F}_1 = \langle \xi \rangle^{(m+2s)^+} \hat{F}$, $\hat{G}_1 = \hat{G}$ and $\hat{F}_2 = \hat{F}$, $\hat{G}_2 = \langle \xi \rangle^{(m+2s)^+} \hat{G}$. On the set for the above integral, we have $\langle \xi_* - \xi^- \rangle^{2s} \lesssim \langle \xi_* \rangle^{2s}$, because $|\xi^-| \lesssim |\xi_*|$ that follows from $|\xi^-|^2 \leq 2|\xi_* \cdot \xi^-| \lesssim |\xi^-| |\xi_*|$. By the Cauchy-Schwarz inequality, we have

$$|A_{2,1,p}| \leq \sum_{j=1}^2 D^{1/2} D_j^{1/2},$$

where

$$D = \int_{\mathbb{R}^3} \left(\iint_{\mathbb{R}^3 \times \mathbb{S}^2} b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) \mathbf{1}_{|\xi^-| \geq \frac{1}{2}\langle \xi_* \rangle} \frac{|\hat{\Phi}_c(\xi_* - \xi^-)|^2}{\langle \xi \rangle^{2s+2m} \langle \xi_* - \xi^- \rangle^{2s}} d\sigma d\xi_* \right) |\hat{H}(\xi)|^2 d\xi$$

and

$$D_j = \iiint_{\mathbb{R}^6 \times \mathbb{S}^2} b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) \frac{\mathbf{1}_{|\xi^-| \geq \frac{1}{2}\langle \xi_* \rangle} \langle \xi \rangle^{2s+2m} \langle \xi_* \rangle^{2s}}{\left(\langle \xi_* \rangle^{(m+2s)^+} + \langle \xi - \xi_* \rangle^{(m+2s)^+}\right)^2} |\hat{F}_j(\xi_*)|^2 |\hat{G}_j(\xi - \xi_*)|^2 d\sigma d\xi_* d\xi.$$

Since $\int_{\mathbb{S}^2} b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) \mathbf{1}_{|\xi^-| \geq \frac{1}{2}\langle \xi_* \rangle} d\sigma \lesssim |\xi|^{2s} \langle \xi_* \rangle^{-2s}$, we obtain

$$D_j \lesssim \|F_j\|_{L^2}^2 \|G_j\|_{L^2}^2.$$

For D , we use the change of variables in ξ_* , $u = \xi_* - \xi^-$ to get

$$D = \int_{\mathbb{R}^3} \left(\iint_{\mathbb{R}^3 \times \mathbb{S}^2} b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) \mathbf{1}_{|\xi^-| \geq \frac{1}{2}\langle u + \xi^- \rangle} \frac{|\hat{\Phi}_c(u)|^2}{\langle \xi \rangle^{2s+2m} \langle u \rangle^{2s}} d\sigma du \right) |\hat{H}(\xi)|^2 d\xi.$$

By noting that $|\xi^-| \geq \frac{1}{2}\langle u + \xi^- \rangle$ implies $|\xi^-| \geq \langle u \rangle / \sqrt{10}$, we have

$$D \lesssim \int_{\mathbb{R}^6} \left(\frac{|\xi|}{\langle u \rangle} \right)^{2s} \frac{|\hat{\Phi}_c(u)|^2}{\langle \xi \rangle^{2s+2m} \langle u \rangle^{2s}} |\hat{H}(\xi)|^2 d\xi du \lesssim \|H\|_{H^{-m}}^2,$$

because $2(\gamma + 3) + 4s > 3$. And this completes the proof of the lemma. \square

3.4. Estimation of commutators. By using the arguments similar to those used in previous subsections, we now prove the following estimation on commutators.

Proposition 3.13. *Assume that $0 < s < 1$ and $\gamma > -3$. Then, for any $l \geq 1$, one has*

$$(3.6) \quad \begin{aligned} & \left| (W_l \Gamma(f, g) - \Gamma(f, W_l g), h) \right|_{L^2} \\ & \lesssim \left(\|f\|_{L^2_{s+\gamma/2}} \|W_{l-s} g\|_{L^2_{s+\gamma/2}} + \|g\|_{L^2_{s+\gamma/2}} \|W_{l-s} f\|_{L^2_{s+\gamma/2}} \right. \\ & \quad \left. + \min\{\|f\|_{L^2} \|W_{l-s} g\|_{L^2_{s+\gamma/2}}, \|f\|_{L^2_{s+\gamma/2}} \|W_{l-s} g\|_{L^2}\} \right. \\ & \quad \left. + \min\{\|g\|_{L^2} \|W_{l-s} f\|_{L^2_{s+\gamma/2}}, \|g\|_{L^2_{s+\gamma/2}} \|W_{l-s} f\|_{L^2}\} \right) \|h\|_{\Phi_\gamma}; \end{aligned}$$

and for any $0 < l < 1$, one has

$$(3.7) \quad \begin{aligned} & \left| (W_l \Gamma(f, g) - \Gamma(f, W_l g), h) \right|_{L^2} \\ & \lesssim \left(\|f\|_{L^2_{s+\gamma/2}} \|W_{l-s} g\|_{L^2_{s+\gamma/2}} + \|g\|_{L^2_{s+\gamma/2}} \|W_{l-s} f\|_{L^2_{s+\gamma/2}} \right. \\ & \quad \left. + \min\{\|f\|_{L^2} \|W_{l-s} g\|_{L^2_{s+\gamma/2}}, \|f\|_{L^2_{s+\gamma/2}} \|W_{l-s} g\|_{L^2}\} \right. \\ & \quad \left. + \min\{\|g\|_{L^2} \|W_{l-s} f\|_{L^2_{s+\gamma/2}}, \|g\|_{L^2_{s+\gamma/2}} \|W_{l-s} f\|_{L^2}\} \right) \|h\|_{\Phi_\gamma}. \end{aligned}$$

Remark 3.14. *Assume that $0 < s < 1$ and $\gamma > -3$. Then, for any $l \geq 0$, one has*

$$\begin{aligned} & \left| (\tilde{W}_l \Gamma(f, g) - \Gamma(f, \tilde{W}_l g), h) \right|_{L^2} \\ & \lesssim \left(\|f\|_{L^2_{s+\gamma/2}} \|\tilde{W}_l g\|_{L^2_{s+\gamma/2}} + \|g\|_{L^2_{s+\gamma/2}} \|\tilde{W}_l f\|_{L^2_{s+\gamma/2}} \right. \\ & \quad \left. + \min\{\|f\|_{L^2} \|\tilde{W}_l g\|_{L^2_{s+\gamma/2}}, \|f\|_{L^2_{s+\gamma/2}} \|\tilde{W}_l g\|_{L^2}\} \right. \\ & \quad \left. + \min\{\|g\|_{L^2} \|\tilde{W}_l f\|_{L^2_{s+\gamma/2}}, \|g\|_{L^2_{s+\gamma/2}} \|\tilde{W}_l f\|_{L^2}\} \right) \|h\|_{\Phi_\gamma}, \end{aligned}$$

where we use (3.6) if $l|\gamma/2 + s| \geq 1$ and (3.7) if $l|\gamma/2 + s| < 1$.

Proof. In view of the decomposition given for Γ , it is enough to consider (3.3) with

$$\mu(v, v_*) = (\mu^c - \mu_*^c)^2 \mu_*^a,$$

for some constants $a, c > 0$. Indeed, all the other terms have compensation by some Gaussian function so that any algebraic weight is not a problem. For this term, the commutator is then given by

$$\left| (W_l \Gamma_\mu(f, g) - \Gamma_\mu(f, W_l g), h)_{L^2} \right| = \iiint b \Phi_\gamma(v - v_*) \mu(v, v_*) f_*' g'(W_l - W_l') h d v d v_* d \sigma,$$

which can be written as

$$\begin{aligned} & \iiint b \Phi_\gamma(v - v_*) \mu(v, v_*) f_*' g'(W_l - W_l') h d v d v_* d \sigma \\ &= \iiint b \Phi_\gamma(v - v_*) \mu(v, v_*) f_*' g'(W_l - W_l') (h - h') d v d v_* d \sigma \\ &+ \iiint b \Phi_\gamma(v - v_*) [\mu(v', v_*') - \mu(v, v_*)] f_* g(W_l' - W_l) h d v d v_* d \sigma \\ &+ \iiint b \Phi_\gamma(v - v_*) \mu(v, v_*) f_* g(W_l' - W_l) h d v d v_* d \sigma \\ &= A + B + C, \end{aligned}$$

by the usual change of variables. For A , we use the Cauchy-Schwarz inequality to get

$$\begin{aligned} A &= \iiint b \Phi_\gamma(v - v_*) \mu(v, v_*) f_*' g'(W_l - W_l') [h - h'] d v d v_* d \sigma \\ &\lesssim \left(\iiint b \Phi_\gamma(v - v_*) (\mu^c - \mu_*^c)^4 \mu_*^a |f_*'|^2 |g'|^2 |W_l - W_l'|^2 d v d v_* d \sigma \right)^{1/2} \|h\|_{\Phi_\gamma} \\ &\lesssim U^{1/2} \|h\|_{\Phi_\gamma}, \end{aligned}$$

where

$$U = \iiint b \langle v - v_* \rangle^\gamma \mu_*^a |f_*'|^2 |g'|^2 |W_l' - W_l|^2 d v d v_* d \sigma.$$

If $l \geq 1$, by using the Taylor's formula, we have

$$|W_l' - W_l|^2 \lesssim \min\{\theta^2 |v - v_*|^2, (\langle v \rangle + \langle v_* \rangle)^2\} (\langle v_* \rangle^{l-1} + \langle v \rangle^{l-1})^2,$$

and then

$$\int_{\mathbb{S}^2} b |W_l' - W_l|^2 d \sigma \lesssim |v - v_*|^{2s} (\langle v \rangle^{2l-2s} + \langle v_* \rangle^{2l-2s}).$$

Then we note immediately that U is similar to the term A_3 in the proof of Proposition 3.3, because we have a Gaussian inside the definition of U . Taking into account the weights here gives

$$\begin{aligned} U &\lesssim \|f\|_{L^2_{s+\gamma/2}} \|W_{l-s} g\|_{L^2_{s+\gamma/2}} + \min\{\|f\|_{L^2} \|W_{l-s} g\|_{L^2_{s+\gamma/2}}, \|f\|_{L^2_{s+\gamma/2}} \|W_{l-s} g\|_{L^2}\} \\ &+ \|g\|_{L^2_{s+\gamma/2}} \|W_{l-s} f\|_{L^2_{s+\gamma/2}} + \min\{\|g\|_{L^2} \|W_{l-s} f\|_{L^2_{s+\gamma/2}}, \|g\|_{L^2_{s+\gamma/2}} \|W_{l-s} f\|_{L^2}\}. \end{aligned}$$

If $0 < l \leq 1$, then we note that

$$|W_l' - W_l|^2 \lesssim \langle v \rangle^{2l} + \langle v_* \rangle^{2l} \quad \text{and} \quad |W_l' - W_l|^2 \lesssim \theta^2 |v - v_*|^2,$$

so that

$$|W_l' - W_l|^2 \lesssim \min\{\theta^2 |v - v_*|^2, \langle v \rangle^{2l} + \langle v_* \rangle^{2l}\}.$$

Then we obtain

$$(3.8) \quad \int_{\mathbb{S}^2} b |W_l' - W_l|^2 d \sigma \lesssim |v - v_*|^{2s} (\langle v \rangle^{2(1-s)l} + \langle v_* \rangle^{2(1-s)l}),$$

and therefore the same argument gives

$$U \lesssim \|f\|_{L^2_{s+\gamma/2}} \|W_{l-s}g\|_{L^2_{s+\gamma/2}} + \min\{\|f\|_{L^2} \|W_{l-s}g\|_{L^2_{s+\gamma/2}}, \|f\|_{L^2_{s+\gamma/2}} \|W_{l-s}g\|_{L^2}\} \\ + \|g\|_{L^2_{s+\gamma/2}} \|W_{l-s}f\|_{L^2_{s+\gamma/2}} + \min\{\|g\|_{L^2} \|W_{l-s}f\|_{L^2_{s+\gamma/2}}, \|g\|_{L^2_{s+\gamma/2}} \|W_{l-s}f\|_{L^2}\},$$

which gives the final conclusion, for $\gamma > -3$. Terms B and C can be dealt with similarly so that we omit the details for brevity. And this completes the proof of the proposition. \square

4. FUNCTIONAL ESTIMATES IN FULL SPACE

In this section, we prove the estimations on the collision operators in some weighted function space of variables $(x, v) \in \mathbb{R}^6$. Together with the essential coercivity estimates proved in Section 2, we give coercivity results for the linear operator in some weighted spaces. These tools are crucial for the proofs of the existence results, both in the local and global cases. Recall the assumption $\gamma + 2s \leq 0$.

Let $N \in \mathbb{N}$, $l \in \mathbb{R}$, we define the weighted function spaces

$$\mathcal{B}_l^N(\mathbb{R}^6) = \left\{ g \in \mathcal{S}'(\mathbb{R}^6_{x,v}); \|g\|_{\mathcal{B}_l^N(\mathbb{R}^6)}^2 = \sum_{|\alpha|+|\beta| \leq N} \int_{\mathbb{R}_x^3} \| |W_{l-|\beta|} \partial_\beta^\alpha g(x, \cdot) \|_{\Phi_\gamma}^2 dx < +\infty \right\}, \\ \tilde{\mathcal{B}}_l^N(\mathbb{R}^6) = \left\{ g \in \mathcal{S}'(\mathbb{R}^6); \|g\|_{\tilde{\mathcal{B}}_l^N(\mathbb{R}^6)}^2 = \sum_{|\alpha|+|\beta| \leq N} \int_{\mathbb{R}_x^3} \| |\tilde{W}_{l-|\beta|} \partial_\beta^\alpha g(x, \cdot) \|_{\Phi_\gamma}^2 dx < +\infty \right\},$$

and also

$$\mathcal{X}^N(\mathbb{R}^6) = \left\{ g \in \mathcal{S}'(\mathbb{R}^6); \|g\|_{\mathcal{X}^N(\mathbb{R}^6)}^2 = \sum_{|\alpha| \leq N} \int_{\mathbb{R}_x^3} \| \partial_x^\alpha g \|_{\Phi_\gamma}^2 dx < +\infty \right\}.$$

4.1. **Estimations without weight.** First of all, one has

Lemma 4.1. *For all $0 < s < 1$, $\gamma > -3$, and for any $\alpha, \beta \in \mathbb{N}^3$,*

$$(4.1) \quad \| \partial_\beta^\alpha \mathbf{P}g \|_{\mathcal{X}^0(\mathbb{R}^6)} + \| \mathbf{P}(\partial_\beta^\alpha g) \|_{\mathcal{X}^0(\mathbb{R}^6)} \leq C_\beta \| \partial_x^\alpha g \|_{L^2(\mathbb{R}^6)}^2,$$

$$(4.2) \quad \frac{\eta_0}{2} \|g\|_{\mathcal{X}^0(\mathbb{R}^6)}^2 - C \|g\|_{L^2(\mathbb{R}^6)}^2 \leq (\mathcal{L}g, g)_{L^2(\mathbb{R}^6)} \leq 2(\mathcal{L}_1g, g)_{L^2(\mathbb{R}^6)} \lesssim \|g\|_{\mathcal{X}^0(\mathbb{R}^6)}^2,$$

and

$$(4.3) \quad \|g\|_{L^2_{s+\frac{\gamma}{2}}(\mathbb{R}^6)}^2 + \|g\|_{L^2(\mathbb{R}_x^3, H^s_{\frac{\gamma}{2}}(\mathbb{R}_v^3))}^2 \lesssim \|g\|_{\mathcal{X}^0(\mathbb{R}^6)}^2 \lesssim \|g\|_{L^2(\mathbb{R}_x^3, H^s_{s+\frac{\gamma}{2}}(\mathbb{R}_v^3))}^2.$$

Proof. From [20], one has

$$\mathbf{P}g = (a_g(t, x) + v \cdot b_g(t, x) + |v|^2 c_g(t, x)) \mu^{1/2},$$

where

$$a_g(t, x) = \int_{\mathbb{R}^3} \left(2 - \frac{|v|^2}{2}\right) g(t, x, v) \mu^{1/2}(v) dv,$$

$$b_g(t, x) = \int_{\mathbb{R}^3} g(t, x, v) v \mu^{1/2}(v) dv,$$

and

$$c_g(t, x) = \int_{\mathbb{R}^3} \left(\frac{|v|^2}{6} - \frac{1}{2}\right) g(t, x, v) \mu^{1/2}(v) dv.$$

Thus, (4.1) can be obtained by using integration by parts. To get (4.2), we use the results from Section 2 to obtain

$$\begin{aligned} \|g\|_{\mathcal{X}^0(\mathbb{R}^6)}^2 &\gtrsim (\mathcal{L}g, g)_{L^2(\mathbb{R}^6)} \geq \eta_0 \|(\mathbf{I} - \mathbf{P})g\|_{\mathcal{X}^0(\mathbb{R}^6)}^2 \\ &\geq \frac{\eta_0}{2} \|g\|_{\mathcal{X}^0(\mathbb{R}^6)}^2 - C \|\mathbf{P}g\|_{\mathcal{X}^0(\mathbb{R}^6)}^2 \geq \frac{\eta_0}{2} \|g\|_{\mathcal{X}^0(\mathbb{R}^6)}^2 - C \|g\|_{L^2(\mathbb{R}^6)}^2. \end{aligned}$$

Finally, (4.3) follows directly from Section 2. \square

The following Lemma is an application of the Sobolev embedding for functions with values in a Hilbert space.

Lemma 4.2.

$$\sup_{x \in \mathbb{R}^3} \|f(x, \cdot)\|_{\Phi_\gamma} \lesssim \|f\|_{\mathcal{X}^2(\mathbb{R}^6)}.$$

Proof. It follows from the definition that

$$\begin{aligned} \left(\sup_{x \in \mathbb{R}^3} \|f(x, \cdot)\|_{\Phi_\gamma} \right)^2 &\leq \iiint B \mu_* \left(\sup_{x \in \mathbb{R}^3} (f(x, v') - f(x, v))^2 \right) dv dv_* d\sigma \\ &\quad + \iiint B \left(\sup_{x \in \mathbb{R}^3} f(x, v_*)^2 \right) (\sqrt{\mu'} - \sqrt{\mu})^2 dv dv_* d\sigma \\ &\lesssim \iiint B \mu_* \left(\sum_{|\alpha| \leq 2} \int (\partial_x^\alpha f(x, v') - \partial_x^\alpha f(x, v))^2 dx \right) dv dv_* d\sigma \\ &\quad + \iiint B \left(\sum_{|\alpha| \leq 2} \int f(x, v_*)^2 dx \right) (\sqrt{\mu'} - \sqrt{\mu})^2 dv dv_* d\sigma \\ &\lesssim \sum_{|\alpha| \leq 2} \int \|\partial_x^\alpha f(x, \cdot)\|_{\Phi_\gamma}^2 dx. \end{aligned}$$

\square

Proposition 4.3. *Under the assumption of Theorem 1.4, for any $N \geq 3$, we have, for all $\alpha \in \mathbb{N}^3$, $|\alpha| \leq N$,*

$$\begin{aligned} \left| \left(\partial_x^\alpha \Gamma(f, g), h \right)_{L^2(\mathbb{R}^6)} \right| &\lesssim \left\{ \|f\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \|g\|_{\mathcal{X}^N(\mathbb{R}^6)} \right. \\ &\quad \left. + \|f\|_{\mathcal{X}^N(\mathbb{R}^6)} \|g\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \right\} \|h\|_{\mathcal{X}^0(\mathbb{R}^6)}. \end{aligned}$$

Proof. Firstly, if $|\alpha_1| \leq N - 2$, we get from Theorem 1.2, Lemma 4.2, and usual Sobolev embedding, replacing the "min" term by the corresponding terms without the weights that

$$\begin{aligned} &\left| \left(\Gamma(\partial^{\alpha_1} f, \partial^{\alpha_2} g), h \right)_{L^2(\mathbb{R}^6)} \right| \\ &\lesssim \left(\int_{\mathbb{R}_x^3} (\|\partial^{\alpha_1} f\|_{L^2}^2 \|\partial^{\alpha_2} g\|_{\Phi_\gamma}^2 + \|\partial^{\alpha_1} f\|_{\Phi_\gamma}^2 \|\partial^{\alpha_2} g\|_{L^2}^2 + \|\partial^{\alpha_1} f\|_{L^2}^2 \|\partial^{\alpha_2} g\|_{L^2}^2) dx \right)^{1/2} \|h\|_{\mathcal{X}^0(\mathbb{R}^6)} \\ &\lesssim \left(\|f\|_{H^{|\alpha_1|+3/2+\epsilon}(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \|g\|_{\mathcal{X}^{|\alpha_2|}(\mathbb{R}^6)} + \|f\|_{\mathcal{X}^{|\alpha_1|+2}(\mathbb{R}^6)} \|g\|_{H^{|\alpha_2|}(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \right) \|h\|_{\mathcal{X}^0(\mathbb{R}^6)}. \end{aligned}$$

If $|\alpha_1| = N - 1$, N , then $|\alpha_2| + 2 \leq N$, we get in a similar way, again from Theorem 1.2 that

$$\begin{aligned} &\left| \left(\Gamma(\partial^{\alpha_1} f, \partial^{\alpha_2} g), h \right)_{L^2(\mathbb{R}^6)} \right| \\ &\lesssim \left(\int_{\mathbb{R}_x^3} (\|\partial^{\alpha_1} f\|_{L^2}^2 \|\partial^{\alpha_2} g\|_{\Phi_\gamma}^2 + \|\partial^{\alpha_1} f\|_{\Phi_\gamma}^2 \|\partial^{\alpha_2} g\|_{L^2}^2 + \|\partial^{\alpha_1} f\|_{L^2}^2 \|\partial^{\alpha_2} g\|_{L^2}^2) dx \right)^{1/2} \|h\|_{\mathcal{X}^0(\mathbb{R}^6)} \\ &\lesssim \left(\|f\|_{H^{|\alpha_1|}(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \|g\|_{\mathcal{X}^{|\alpha_2|+2}(\mathbb{R}^6)} + \|f\|_{\mathcal{X}^{|\alpha_1|}(\mathbb{R}^6)} \|g\|_{H^{|\alpha_2|+3/2+\epsilon}(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \right) \|h\|_{\mathcal{X}^0(\mathbb{R}^6)}. \end{aligned}$$

The proof of the proposition is then completed by recalling the Leibniz formula

$$\partial_x^\alpha \Gamma(f, g) = \sum_{\alpha_1 + \alpha_2 = \alpha} C_{\alpha_1, \alpha_2} \Gamma(\partial_x^{\alpha_1} f, \partial_x^{\alpha_2} g).$$

□

Remark 4.4. *The above proof shows that, for $|\alpha_1| < N$,*

$$\left| \left(\Gamma(\partial^{\alpha_1} f, \partial^{\alpha_2} g), h \right)_{L^2(\mathbb{R}^6)} \right| \lesssim \|f\|_{H^N(\mathbb{R}^3; L^2(\mathbb{R}^3))} \|h\|_{\mathcal{X}^0(\mathbb{R}^6)} \left(\|g\|_{\mathcal{X}^N(\mathbb{R}^6)} + \|g\|_{H^N(\mathbb{R}^3; L^2(\mathbb{R}^3))} \right).$$

Finally, the estimate on the linear operator \mathcal{L}_2 can be given as follows, which in fact can be deduced from Section 2.

Proposition 4.5. *For all $0 < s < 1$, $\gamma > -3$ and any $\alpha \in \mathbb{N}^3$, we have*

$$\left| \left(\partial_x^\alpha \mathcal{L}_2(f), h \right)_{L^2(\mathbb{R}^6)} \right| \leq C_{|\alpha|} \|f\|_{H^{|\alpha|}(\mathbb{R}^3; L^2(\mathbb{R}^3))} \|\mu^{1/10^3} h\|_{L^2(\mathbb{R}^6)}.$$

4.2. Estimation with weight. We now prove the following upper bound with weights.

Proposition 4.6. *For all $0 < s < 1$, $\gamma > -3$, and for any $N \geq 6$, $\ell \geq N$, $|\alpha| + |\beta| \leq N$, we have*

$$\begin{aligned} \left| \left(W_{\ell-|\beta|} \partial_\beta^\alpha \Gamma(f, g), h \right)_{L^2(\mathbb{R}^6)} \right| &\lesssim \|h\|_{\mathcal{B}_0^0(\mathbb{R}^6)} \left(\|f\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \|g\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \right. \\ &\quad \left. + \|f\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \|g\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)} + \|g\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \|f\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)} \right). \end{aligned}$$

Proof. By using Leibniz formula, we have

$$\begin{aligned} \left(W_{\ell-|\beta|} \partial_\beta^\alpha \Gamma(f, g), h \right) &= \sum C_{\alpha_2, \beta_2, \beta_3}^{\alpha_1, \beta_1} \left(W_{\ell-|\beta|} \mathcal{T}(\partial_{\beta_1}^{\alpha_1} f, \partial_{\beta_2}^{\alpha_2} g, \mu_{\beta_3}), h \right) \\ &= \sum C_{\alpha_2, \beta_2, \beta_3}^{\alpha_1, \beta_1} \left(W_{\ell-|\beta|} \mathcal{T}(\partial_{\beta_1}^{\alpha_1} f, \partial_{\beta_2}^{\alpha_2} g, \mu_{\beta_3}) - \mathcal{T}(\partial_{\beta_1}^{\alpha_1} f, W_{\ell-|\beta|} \partial_{\beta_2}^{\alpha_2} g, \mu_{\beta_3}), h \right) \\ &\quad + \sum C_{\alpha_2, \beta_2, \beta_3}^{\alpha_1, \beta_1} \left(\mathcal{T}(\partial_{\beta_1}^{\alpha_1} f, W_{\ell-|\beta|} \partial_{\beta_2}^{\alpha_2} g, \mu_{\beta_3}), h \right). \end{aligned}$$

Note that \mathcal{T} shares the same upper bound properties as Γ given in the previous propositions. In fact, by using Proposition 3.13

$$\begin{aligned} A &= \left| \left(W_{\ell-|\beta|} \mathcal{T}(\partial_{\beta_1}^{\alpha_1} f, \partial_{\beta_2}^{\alpha_2} g, \mu_{\beta_3}) - \mathcal{T}(\partial_{\beta_1}^{\alpha_1} f, W_{\ell-|\beta|} \partial_{\beta_2}^{\alpha_2} g, \mu_{\beta_3}), h \right)_{L^2(\mathbb{R}^6)} \right| \\ &\lesssim \|h\|_{\mathcal{B}_0^0(\mathbb{R}^6)} \left\{ \int_{\mathbb{R}^3} \left(\|W_{\ell-|\beta|+\gamma/2} \partial_{\beta_1}^{\alpha_1} f(x, \cdot)\|_{L^2(\mathbb{R}^3)}^2 \|\partial_{\beta_2}^{\alpha_2} g(x, \cdot)\|_{L^2_{\gamma/2+s}(\mathbb{R}^3)}^2 \right. \right. \\ &\quad \left. \left. + \|\partial_{\beta_1}^{\alpha_1} f(x, \cdot)\|_{L^2_{\gamma/2+s}(\mathbb{R}^3)}^2 \|W_{\ell-|\beta|+\gamma/2} \partial_{\beta_2}^{\alpha_2} g(x, \cdot)\|_{L^2(\mathbb{R}^3)}^2 \right) \right. \\ &\quad \left. + \min \left\{ \|\partial_{\beta_1}^{\alpha_1} f(x, \cdot)\|_{L^2(\mathbb{R}^3)}^2 \|W_{\ell-|\beta|+\gamma/2} \partial_{\beta_2}^{\alpha_2} g(x, \cdot)\|_{L^2(\mathbb{R}^3)}^2, \right. \right. \\ &\quad \left. \left. \|\partial_{\beta_1}^{\alpha_1} f(x, \cdot)\|_{L^2_{\gamma/2+s}(\mathbb{R}^3)}^2 \|W_{\ell-s} \partial_{\beta_2}^{\alpha_2} g(x, \cdot)\|_{L^2(\mathbb{R}^3)}^2 \right\} \right. \\ &\quad \left. + \min \left\{ \|\partial_{\beta_2}^{\alpha_2} g(x, \cdot)\|_{L^2(\mathbb{R}^3)}^2 \|W_{\ell-|\beta|+\gamma/2} \partial_{\beta_1}^{\alpha_1} f(x, \cdot)\|_{L^2(\mathbb{R}^3)}^2, \right. \right. \\ &\quad \left. \left. \|\partial_{\beta_2}^{\alpha_2} g(x, \cdot)\|_{L^2_{\gamma/2+s}(\mathbb{R}^3)}^2 \|W_{\ell-|\beta|-s} \partial_{\beta_1}^{\alpha_1} f(x, \cdot)\|_{L^2(\mathbb{R}^3)}^2 \right\} dx \right\}^{1/2}. \end{aligned}$$

For this, we divide the discussion into two cases.

Case 1: $|\alpha_1| + |\beta_1| \leq N - 2$. We have, by using $H^2(\mathbb{R}_x^3) \subset L^\infty(\mathbb{R}_x^3)$, and $\gamma + 2s \leq 0$ that

$$\begin{aligned} A &\lesssim \|h\|_{\mathcal{B}_0^0(\mathbb{R}^6)} \left(\|W_{\ell-|\beta|+\gamma/2} \nabla_x^2 \partial_{\beta_1}^{\alpha_1} f\|_{L^2(\mathbb{R}^6)} \|\partial_{\beta_2}^{\alpha_2} g\|_{L^2(\mathbb{R}^6)} \right. \\ &\quad \left. + \|\nabla_x^2 \partial_{\beta_1}^{\alpha_1} f\|_{L^2(\mathbb{R}^6)} \|W_{\ell-|\beta|+\gamma/2} \partial_{\beta_2}^{\alpha_2} g\|_{L^2(\mathbb{R}^6)} \right) \\ &\lesssim \|h\|_{\mathcal{B}_0^0(\mathbb{R}^6)} \|f\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \|g\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}. \end{aligned}$$

Case 2: $|\alpha_1| + |\beta_1| > N - 2$. Then $|\alpha_2| + |\beta_2| \leq 1$ and $|\alpha_2| + |\beta_2| + 5 \leq N$, and we have

$$\begin{aligned} A &\lesssim \|h\|_{\mathcal{B}_0^0(\mathbb{R}^6)} \left(\|W_{\ell-|\beta|+\gamma/2} \partial_{\beta_1}^{\alpha_1} f\|_{L^2(\mathbb{R}^6)} \|\nabla_x^2 \partial_{\beta_2}^{\alpha_2} g\|_{L^2(\mathbb{R}^6)} \right. \\ &\quad \left. + \|\partial_{\beta_1}^{\alpha_1} f\|_{L^2(\mathbb{R}^6)} \|W_{\ell-|\beta|+\gamma/2} \nabla_x^2 \partial_{\beta_2}^{\alpha_2} g\|_{L^2(\mathbb{R}^6)} \right) \\ &\lesssim \|h\|_{\mathcal{B}_0^0(\mathbb{R}^6)} \|f\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \|g\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}. \end{aligned}$$

Therefore, we get

$$(4.4) \quad \sum C_{\alpha_2, \beta_2, \beta_3}^{\alpha_1, \beta_1} \left| (W_{\ell-|\beta|} \mathcal{T}(\partial_{\beta_1}^{\alpha_1} f, \partial_{\beta_2}^{\alpha_2} g, \mu_{\beta_3}) - \mathcal{T}(\partial_{\beta_1}^{\alpha_1} f, W_{\ell-|\beta|} \partial_{\beta_2}^{\alpha_2} g, \mu_{\beta_3}), h) \right| \\ \leq C \|h\|_{\mathcal{B}_0^0(\mathbb{R}^6)} \|f\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \|g\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}.$$

Next, we deal with the following term and the discussion is also divided into several steps:

$$\sum C_{\alpha_2, \beta_2, \beta_3}^{\alpha_1, \beta_1} (\mathcal{T}(\partial_{\beta_1}^{\alpha_1} f, W_{\ell-|\beta|} \partial_{\beta_2}^{\alpha_2} g, \mu_{\beta_3}), h).$$

Case I: $|\alpha_1| + |\beta_1| \leq 1$. From Proposition 3.1, we get

$$\begin{aligned} B &= \left| (\mathcal{T}(\partial_{\beta_1}^{\alpha_1} f, W_{\ell-|\beta|} \partial_{\beta_2}^{\alpha_2} g, \mu_{\beta_3}), h)_{L^2(\mathbb{R}^6)} \right| \\ &\lesssim \|h\|_{\mathcal{B}_0^0} \left\{ \|f\|_{H^5(\mathbb{R}^6)} \|W_{\ell-|\beta|} \partial_{\beta_2}^{\alpha_2} g\|_{\mathcal{B}_0^0(\mathbb{R}^6)} + \|f\|_{H^6(\mathbb{R}^6)} \|W_{\ell-|\beta|} \partial_{\beta_2}^{\alpha_2} g\|_{L^2_{s+\gamma/2}(\mathbb{R}^6)} \right\} \\ &\lesssim \|h\|_{\mathcal{B}_0^0} \|f\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \|g\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)}. \end{aligned}$$

Case II: $2 \leq |\alpha_1| + |\beta_1| \leq 3$. Again from Proposition 3.1, one has

$$\begin{aligned} B &\lesssim \|h\|_{\mathcal{B}_0^0} \left\{ \|f\|_{H^5(\mathbb{R}^6)} \|W_{\ell-|\beta|} \nabla_x^2 \partial_{\beta_2}^{\alpha_2} g\|_{\mathcal{B}_0^0(\mathbb{R}^6)} + \|f\|_{H^6(\mathbb{R}^6)} \|W_{\ell-|\beta|} \nabla_x^2 \partial_{\beta_2}^{\alpha_2} g\|_{L^2_{s+\gamma/2}(\mathbb{R}^6)} \right\} \\ &\lesssim \|h\|_{\mathcal{B}_0^0} \|f\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \|g\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)}. \end{aligned}$$

Case III: $4 \leq |\alpha_1| + |\beta_1| \leq N - 1$. Then $|\alpha_2| + |\beta_2| \leq N - 4$, and from Proposition 3.3, we get

$$\begin{aligned} B &\lesssim \|h\|_{\mathcal{B}_0^0} \left\{ \|\partial_{\beta_1}^{\alpha_1} f\|_{L^2_{\gamma/2+s}(\mathbb{R}^6)} \|W_{\ell-|\beta|} \nabla_x^2 \partial_{\beta_2}^{\alpha_2} g\|_{\mathcal{B}_0^0(\mathbb{R}^6)} + \|\partial_{\beta_1}^{\alpha_1} f\|_{\mathcal{B}_0^0(\mathbb{R}^6)} \|W_{\ell-|\beta|} \nabla_x^2 \partial_{\beta_2}^{\alpha_2} g\|_{L^2_{s+\gamma/2}(\mathbb{R}^6)} \right\} \\ &\quad + \|\partial_{\beta_1}^{\alpha_1} f\|_{L^2(\mathbb{R}^6)} \|W_{\ell-|\beta|} \nabla_x^2 \partial_{\beta_2}^{\alpha_2} g\|_{L^2_{\gamma/2+s}(\mathbb{R}^6)} + \|\mu^{1/60} \partial_{\beta_1}^{\alpha_1} f\|_{L^2(\mathbb{R}^6)} \|\mu^{1/60} W_{\ell-|\beta|} \partial_{\beta_2}^{\alpha_2} g\|_{H^{3/2+\epsilon}(\mathbb{R}^3; H^{3/2-\epsilon}(\mathbb{R}^3))} \\ &\quad + \|\mu^{1/60} \partial_{\beta_1}^{\alpha_1} f\|_{H^s(\mathbb{R}^6)} \|\mu^{1/60} W_{\ell-|\beta|} \partial_{\beta_2}^{\alpha_2} g\|_{H^{3+\epsilon}(\mathbb{R}^6)} \lesssim \|h\|_{\mathcal{B}_0^0} \|f\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \|g\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}. \end{aligned}$$

Case IV: $|\alpha_1| + |\beta_1| = |\alpha| + |\beta| \leq N$. Again from Proposition 3.3, we have

$$\begin{aligned} B &= \left| (\Gamma(\partial_{\beta}^{\alpha} f, W_{\ell-|\beta|} g), h)_{L^2(\mathbb{R}^6)} \right| \lesssim \|h\|_{\mathcal{B}_0^0(\mathbb{R}^6)} \left\{ \|\partial_{\beta}^{\alpha} f\|_{L^2_{s+\gamma/2}(\mathbb{R}^6)} \|W_{\ell-|\beta|} \nabla_x^2 g\|_{\mathcal{B}_0^0(\mathbb{R}^6)} \right. \\ &\quad + \|W_{\ell-|\beta|} \nabla_x^2 g\|_{L^2_{s+\gamma/2}(\mathbb{R}^6)} \|\partial_{\beta}^{\alpha} f\|_{\mathcal{B}_0^0(\mathbb{R}^6)} + \|\partial_{\beta}^{\alpha} f\|_{L^2_{s+\gamma/2}(\mathbb{R}^6)} \|W_{\ell-|\beta|} \nabla_x^2 g\|_{L^2(\mathbb{R}^6)} \\ &\quad + \|\mu^{1/40} \partial_{\beta}^{\alpha} f\|_{L^2_{s+\gamma/2}(\mathbb{R}^6)} \|\mu^{1/40} W_{\ell-|\beta|} g\|_{H^{3/2+\epsilon}(\mathbb{R}^3; H^{\max(-\gamma/2, 1)}(\mathbb{R}^3))} \\ &\quad \left. + \|\mu^{1/40} \partial_{\beta}^{\alpha} f\|_{\mathcal{B}_0^0(\mathbb{R}^6)} \|\mu^{1/40} W_{\ell-|\beta|} g\|_{L^\infty(\mathbb{R}^6)} \right\} \\ &\lesssim \|h\|_{\mathcal{B}_0^0(\mathbb{R}^6)} \|g\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \|f\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)}. \end{aligned}$$

Finally, the proof of the proposition is completed. \square

4.3. Estimation with modified weight. In the sequel, we shall often use the inequalities

$$(4.5) \quad \|f\|_{L^2_{s+\gamma/2}} \leq \|f\|_{L^2}, \quad \|f\|_{L^2_{s+\gamma/2}} \leq \|f\|.$$

The first inequality is valid since we are assuming that $\gamma + 2s \leq 0$, while the second one comes from Section 2.

To obtain the global existence, the upper bound given in Proposition 4.6 for the general case is not enough because this bound can not be controlled by the coercivity estimate of the linearized operator that contains loss of weight in the case of soft potential. To

overcome this, we now prove the following upper bound under the conditions of Theorem 1.5.

Proposition 4.7. *Under the assumptions of Theorem 1.5 on the parameters γ and s , for any $N \geq 6, \ell \geq N, |\alpha| + |\beta| \leq N$, one has*

$$(4.6) \quad |(\tilde{W}_{l-|\beta|} \partial_\beta^\alpha \Gamma(f, g), \tilde{W}_{l-|\beta|} \partial_\beta^\alpha h)_{L^2(\mathbb{R}^6)}| \leq \left(\|f\|_{\tilde{\mathcal{H}}_l^N(\mathbb{R}^6)} \|g\|_{\tilde{\mathcal{H}}_l^N(\mathbb{R}^6)} + \|g\|_{\tilde{\mathcal{H}}_l^N(\mathbb{R}^6)} \|f\|_{\tilde{\mathcal{H}}_l^N(\mathbb{R}^6)} \right) \|h\|_{\tilde{\mathcal{H}}_l^N(\mathbb{R}^6)}.$$

Proof. First, notice that from Remark 3.14 and (4.5), we have for $\gamma > -3$,

$$(4.7) \quad |(\tilde{W}_l \Gamma(f, g) - \Gamma(f, \tilde{W}_l g), h)_{L^2(\mathbb{R}_+^3)}| \leq \left(\|f\|_{L^2(\mathbb{R}_+^3)} \|\tilde{W}_l g\|_{\Phi_\gamma} + \|\tilde{W}_l f\|_{\Phi_\gamma} \|g\|_{L^2(\mathbb{R}_+^3)} \right) \|h\|_{\Phi_\gamma}.$$

Recall the definition of \mathcal{T} to deduce that

$$\begin{aligned} |(\tilde{W}_l \partial_\beta \Gamma(f, g), \tilde{W}_l \partial_\beta h)_{L^2(\mathbb{R}_+^3)}| &\leq \sum_{\beta_1 + \beta_2 + \beta_3 = \beta} |(\tilde{W}_l \mathcal{T}(\partial_{\beta_1} f, \partial_{\beta_2} g, \partial_{\beta_3} \mu^{1/2}), \tilde{W}_l \partial_\beta h)_{L^2(\mathbb{R}_+^3)}| \\ &= \sum \left\{ |(\mathcal{T}(\partial_{\beta_1} f, \tilde{W}_l \partial_{\beta_2} g, \partial_{\beta_3} \mu^{1/2}), \tilde{W}_l \partial_\beta h)_{L^2(\mathbb{R}_+^3)}| \right. \\ &\quad \left. + |(\tilde{W}_l \mathcal{T}(\partial_{\beta_1} f, \partial_{\beta_2} g, \partial_{\beta_3} \mu^{1/2}) - \mathcal{T}(\partial_{\beta_1} f, \tilde{W}_l \partial_{\beta_2} g, \partial_{\beta_3} \mu^{1/2}), \tilde{W}_l \partial_\beta h)_{L^2(\mathbb{R}_+^3)}| \right\}. \end{aligned}$$

By using Theorem 1.2, we obtain

$$\begin{aligned} |(\Gamma(\partial_{\beta_1} f, \tilde{W}_l \partial_{\beta_2} g), \tilde{W}_l \partial_\beta h)_{L^2(\mathbb{R}_+^3)}| \\ \leq \left(\|\partial_{\beta_1} f\|_{L^2(\mathbb{R}_+^3)} \|\tilde{W}_l \partial_{\beta_2} g\|_{\Phi_\gamma} + \|\partial_{\beta_1} f\|_{\Phi_\gamma} \|\tilde{W}_l \partial_{\beta_2} g\|_{L^2(\mathbb{R}_+^3)} \right) \|\tilde{W}_l \partial_\beta h\|_{\Phi_\gamma}. \end{aligned}$$

Moreover, (4.7) implies that

$$\begin{aligned} |(\tilde{W}_l \Gamma(\partial_{\beta_1} f, \partial_{\beta_2} g) - \Gamma(\partial_{\beta_1} f, \tilde{W}_l \partial_{\beta_2} g), \tilde{W}_l \partial_\beta h)_{L^2(\mathbb{R}_+^3)}| \\ \leq \left(\|\partial_{\beta_1} f\|_{L^2(\mathbb{R}_+^3)} \|\tilde{W}_l \partial_{\beta_2} g\|_{\Phi_\gamma} + \|\tilde{W}_l \partial_{\beta_1} f\|_{\Phi_\gamma} \|\partial_{\beta_2} g\|_{L^2(\mathbb{R}_+^3)} \right) \|\tilde{W}_l \partial_\beta h\|_{\Phi_\gamma}. \end{aligned}$$

As a consequence,

$$(4.8) \quad |(\tilde{W}_l \partial_\beta \Gamma(f, g), \tilde{W}_l \partial_\beta h)_{L^2(\mathbb{R}_+^3)}| \leq \sum \left(\|\partial_{\beta_1} f\|_{L^2(\mathbb{R}_+^3)} \|\tilde{W}_l \partial_{\beta_2} g\|_{\Phi_\gamma} + \|\partial_{\beta_1} f\|_{\Phi_\gamma} \|\tilde{W}_l \partial_{\beta_2} g\|_{L^2(\mathbb{R}_+^3)} + \|\tilde{W}_l \partial_{\beta_1} f\|_{\Phi_\gamma} \|\partial_{\beta_2} g\|_{L^2(\mathbb{R}_+^3)} \right) \|\tilde{W}_l \partial_\beta h\|_{\Phi_\gamma}.$$

By the Leibniz rule in x variable, one has

$$\begin{aligned} |(\tilde{W}_{l-|\beta|} \partial_\beta^\alpha \Gamma(f, g), \tilde{W}_{l-|\beta|} \partial_\beta^\alpha h)_{L^2(\mathbb{R}^6)}| \\ \leq \sum_{\alpha_1 + \alpha_2 = \alpha} |(\tilde{W}_{l-|\beta|} \partial_\beta \Gamma(\partial^{\alpha_1} f, \partial^{\alpha_2} g), \tilde{W}_{l-|\beta|} \partial_\beta^\alpha h)_{L^2(\mathbb{R}^6)}| \\ \leq \sum \int_{\mathbb{R}^3} \left(\|\partial^{\alpha_1} f\|_{L^2(\mathbb{R}_+^3)} \|\tilde{W}_{l-|\beta|} \partial_{\beta_2}^{\alpha_2} g\|_{\Phi_\gamma} + \|\partial^{\alpha_1} f\|_{\Phi_\gamma} \|\tilde{W}_{l-|\beta|} \partial_{\beta_2}^{\alpha_2} g\|_{L^2(\mathbb{R}_+^3)} \right. \\ \quad \left. + \|\tilde{W}_{l-|\beta|} \partial_{\beta_1}^{\alpha_1} f\|_{\Phi_\gamma} \|\partial_{\beta_2}^{\alpha_2} g\|_{L^2(\mathbb{R}_+^3)} \right) \|\tilde{W}_{l-|\beta|} \partial_\beta^\alpha h\|_{\Phi_\gamma} dx \\ \leq \sum \int_{\mathbb{R}^3} \left(\|\tilde{W}_{l-\beta_1} \partial_{\beta_1}^{\alpha_1} f\|_{L^2(\mathbb{R}_+^3)} \|\tilde{W}_{l-\beta_2} \partial_{\beta_2}^{\alpha_2} g\|_{\Phi_\gamma} + \|\tilde{W}_{l-\beta_1} \partial_{\beta_1}^{\alpha_1} f\|_{\Phi_\gamma} \|\tilde{W}_{l-\beta_2} \partial_{\beta_2}^{\alpha_2} g\|_{L^2(\mathbb{R}_+^3)} \right. \\ \quad \left. + \|\tilde{W}_{l-\beta_1} \partial_{\beta_1}^{\alpha_1} f\|_{\Phi_\gamma} \|\tilde{W}_{l-\beta_2} \partial_{\beta_2}^{\alpha_2} g\|_{L^2(\mathbb{R}_+^3)} \right) \|\tilde{W}_{l-|\beta|} \partial_\beta^\alpha h\|_{\Phi_\gamma} dx \\ \leq \sum \int_{\mathbb{R}^3} \left(\|\tilde{W}_{l-\beta_1} \partial_{\beta_1}^{\alpha_1} f\|_{L^2(\mathbb{R}_+^3)} \|\tilde{W}_{l-\beta_2} \partial_{\beta_2}^{\alpha_2} g\|_{\Phi_\gamma} \right. \\ \quad \left. + \|\tilde{W}_{l-\beta_1} \partial_{\beta_1}^{\alpha_1} f\|_{\Phi_\gamma} \|\tilde{W}_{l-\beta_2} \partial_{\beta_2}^{\alpha_2} g\|_{L^2(\mathbb{R}_+^3)} \right) \|\tilde{W}_{l-|\beta|} \partial_\beta^\alpha h\|_{\Phi_\gamma} dx \\ = \sum (G_{\alpha_1, \beta_1, \beta_2}^1 + G_{\alpha_1, \beta_1, \beta_2}^2). \end{aligned}$$

Now these terms are discussed in the following two cases.

- When $|\alpha_1| + |\beta_1| \leq N/2$, we have

$$\begin{aligned} G_{\alpha_1, \beta_1, \beta_2}^1 &\leq \|\tilde{W}_{l-\beta_1} \partial_{\beta_1}^{\alpha_1} f\|_{L^\infty(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \int_{\mathbb{R}^3} \|\tilde{W}_{l-\beta_2} \partial_{\beta_2}^{\alpha_2} g\|_{\Phi_\gamma} \|\tilde{W}^{l-|\beta|} \partial_\beta^\alpha h\|_{\Phi_\gamma} dx \\ &\leq \|f\|_{\tilde{\mathcal{H}}_l^N(\mathbb{R}^6)} \|g\|_{\tilde{\mathcal{B}}_l^N(\mathbb{R}^6)} \|h\|_{\tilde{\mathcal{B}}_l^N(\mathbb{R}^6)}. \\ G_{\alpha_1, \beta_1, \beta_2}^2 &\leq \|\tilde{W}_{l-\beta_1} \partial_{\beta_1}^{\alpha_1} f\|_{\Phi_\gamma} \|L^\infty(\mathbb{R}_x^3)\| \int_{\mathbb{R}^3} \|\tilde{W}_{l-\beta_2} \partial_{\beta_2}^{\alpha_2} g\|_{L_v^2} \|\tilde{W}^{l-|\beta|} \partial_\beta^\alpha h\|_{\Phi_\gamma} dx \\ &\leq \|f\|_{\tilde{\mathcal{B}}_l^N(\mathbb{R}^6)} \|g\|_{\tilde{\mathcal{H}}_l^N(\mathbb{R}^6)} \|h\|_{\tilde{\mathcal{B}}_l^N(\mathbb{R}^6)}. \end{aligned}$$

- When $|\alpha_1| + |\beta_1| \geq N/2$, we have

$$\begin{aligned} G_{\alpha_1, \beta_1, \beta_2}^1 &\leq \|\tilde{W}_{l-\beta_2} \partial_{\beta_2}^{\alpha_2} g\|_{\Phi_\gamma} \|L^\infty(\mathbb{R}_x^3)\| \int_{\mathbb{R}^3} \|\tilde{W}_{l-\beta_1} \partial_{\beta_1}^{\alpha_1} f\|_{L_v^2} \|\tilde{W}^{l-|\beta|} \partial_\beta^\alpha h\|_{\Phi_\gamma} dx \\ &\leq \|f\|_{\tilde{\mathcal{H}}_l^N(\mathbb{R}^6)} \|g\|_{\tilde{\mathcal{B}}_l^N(\mathbb{R}^6)} \|h\|_{\tilde{\mathcal{B}}_l^N(\mathbb{R}^6)}. \\ G_{\alpha_1, \beta_1, \beta_2}^2 &\leq \|\tilde{W}_{l-\beta_2} \partial_{\beta_2}^{\alpha_2} g\|_{L_x^\infty(L_v^2)} \int_{\mathbb{R}^3} \|\tilde{W}_{l-\beta_1} \partial_{\beta_1}^{\alpha_1} f\|_{\Phi_\gamma} \|\tilde{W}^{l-|\beta|} \partial_\beta^\alpha h\|_{\Phi_\gamma} dx \\ &\leq \|f\|_{\tilde{\mathcal{B}}_l^N(\mathbb{R}^6)} \|g\|_{\tilde{\mathcal{H}}_l^N(\mathbb{R}^6)} \|h\|_{\tilde{\mathcal{B}}_l^N(\mathbb{R}^6)}. \end{aligned}$$

Here, we have used Lemma 4.2 to get

$$\|\|\|\tilde{W}_{l-\beta_1} \partial_{\beta_1}^{\alpha_1} f\|\|\|_{\Phi_\gamma} \|L^\infty(\mathbb{R}_x^3)\| \leq \|f\|_{\tilde{\mathcal{B}}_l^N(\mathbb{R}^6)},$$

for $|\alpha_1| + |\beta_1| \leq N/2$. Therefore, we complete the proof of the proposition. \square

4.4. Weighted coercivity of the linearized operator. We turn to the weighted lower estimates, more precisely, the lower bound for $(W_l \mathcal{L}g, W_l g)_{L^2(\mathbb{R}_v^3)}$. Let us recall that in Section 2 it was shown that, if $\gamma > -3$, then there exists a constant $C > 0$ such that

$$(4.9) \quad \|g\|_{\Phi_\gamma}^2 \geq (\mathcal{L}_1 g, g)_{L^2(\mathbb{R}_v^3)} \geq \eta_0 \|g\|_{\Phi_\gamma}^2 - C \|g\|_{L_{\gamma/2}^2(\mathbb{R}_v^3)}^2.$$

In the estimation on the weighted linearized collisional operator \mathcal{L}_1 , we need to consider the commutator estimate of the weight and the linearized operator. However, we can not apply the Proposition 3.13 directly because the error term then will have the weight of the order of $s + \gamma/2$. The purpose of the following proposition is to show that the error term coming from this commutator has the weight of order $\gamma/2$ only.

Proposition 4.8. *For all $0 < s < 1$, $\gamma > -3$, and for any $l \geq 0$, there exists a positive constant C such that*

$$(4.10) \quad (W_l \mathcal{L}_1 g, W_l g)_{L^2(\mathbb{R}_v^3)} \geq \frac{\eta_0}{2} \|W_l g\|_{\Phi_\gamma}^2 - C \|W_l g\|_{L_{\gamma/2}^2(\mathbb{R}_v^3)}^2.$$

Moreover, for any $\beta \in \mathbb{N}^3 \setminus \{0\}$, one has

$$\begin{aligned} (W_l \partial_\beta \mathcal{L}_1 g, W_l \partial_\beta g)_{L^2(\mathbb{R}_v^3)} &\geq \frac{\eta_0}{2} \|W_l \partial_\beta g\|_{\Phi_\gamma}^2 - C \|W_l \partial_\beta g\|_{L_{\gamma/2}^2(\mathbb{R}_v^3)}^2 \\ &\quad - C \left(\sum_{\beta_1 < \beta} \|W_l \partial_{\beta_1} g\|_{\Phi_\gamma} \right) \|W_l \partial_\beta g\|_{\Phi_\gamma}. \end{aligned}$$

Proof. According to (4.9), it is enough to show that

$$|(W_l \mathcal{L}_1 g, W_l g)_{L^2}| \leq C_\delta \|W_l g\|_{L_{\gamma/2}^2}^2 + \delta \|W_l g\|_{L_{s+\gamma/2}^2}^2,$$

where $\delta > 0$ can be arbitrarily small. By using the above expression, one has

$$\begin{aligned}
(W_l \mathcal{L}_1 g, W_l g)_{L^2} &= - \iiint b \Phi_\gamma (\mu_*')^{1/2} g' - (\mu_*)^{1/2} g \mu_*^{1/2} W_l W_l g dv_* d\sigma dv \\
&= - \iiint b \Phi_\gamma (\mu_*')^{1/2} W_l g' - (\mu_*)^{1/2} W_l g \mu_*^{1/2} W_l g dv_* d\sigma dv \\
&= - \iiint b \Phi_\gamma [(\mu_*')^{1/2} W_l' g' - (\mu_*)^{1/2} W_l g] + (\mu_*')^{1/2} g' (W_l - W_l') \mu_*^{1/2} W_l g \\
&= 1/2 \iiint b \Phi_\gamma (\mu_*')^{1/2} W_l' g' - (\mu_*)^{1/2} W_l g)^2 dv_* d\sigma dv \\
&\quad - \iiint b \Phi_\gamma (\mu_*')^{1/2} g' (W_l - W_l') \mu_*^{1/2} W_l g dv_* d\sigma dv \\
&= (\mathcal{L}_1(W_l g), (W_l g))_{L^2} + I,
\end{aligned}$$

where

$$I = ([W_l, \mathcal{L}_1]g, W_l g)_{L^2} = - \iiint b \Phi_\gamma (\mu_*')^{1/2} g' \mu_*^{1/2} W_l g (W_l - W_l') dv_* d\sigma dv.$$

Changing variables yields

$$I = - \iiint b \Phi_\gamma (\mu_*)^{1/2} g \mu_*^{1/2} W_l' g' (W_l' - W_l) dv_* d\sigma dv.$$

Adding the above two equations gives

$$2I = - \iiint b \Phi_\gamma (\mu_*)^{1/2} g \mu_*^{1/2} g' (W_l' - W_l)^2 dv_* d\sigma dv.$$

Then, by using the Cauchy-Schwarz inequality with respect to full variables, we find

$$|I| \lesssim \iiint b \Phi_\gamma (\mu_*)^{1/2} g^2 (W_l' - W_l)^2 dv_* d\sigma dv.$$

When $l \geq 1$, since

$$\begin{aligned}
\int_{\mathbb{S}^2} b (W_l' - W_l)^2 d\sigma &\lesssim |v - v_*|^{2s} [\langle v \rangle + \langle v_* \rangle]^{2l-2s} \\
&\lesssim \langle v \rangle^{2s} \langle v_* \rangle^{2s} (\langle v \rangle^{2l-2s} + \langle v_* \rangle^{2l-2s}) \\
&\lesssim \langle v_* \rangle^{2s} \langle v \rangle^{2l} + \langle v \rangle^{2s} \langle v_* \rangle^{2l} \lesssim \langle v_* \rangle^{2l} \langle v \rangle^{2l},
\end{aligned}$$

by using Section 2, we have

$$|I| \lesssim \|W_l g\|_{L_{\gamma/2}^2}^2.$$

When $0 < l \leq 1$, by (3.8), we have

$$\int_{\mathbb{S}^2} b |W_l' - W_l|^2 d\sigma \lesssim \langle v \rangle^{2l(1-s)} \langle v_* \rangle^{2l(1-s)} |v - v_*|^{2s}.$$

Consider

$$I = \iiint_{|v| \leq R} + \iiint_{|v| > R} = I_1 + I_2.$$

It is obvious that for any fixed R ,

$$|I_1| \lesssim \|W_l g\|_{L_{\gamma/2}^2}^2.$$

For I_2 , we split

$$I_2 = \iint_{|v| \geq R} \int_{\theta|v-v_*| \leq \frac{1}{2}|v|} + \iint_{|v| \geq R} \int_{\theta|v-v_*| \geq \frac{1}{2}|v|} = I_{2,1} + I_{2,2},$$

which are the singular part and the non-singular part respectively. For the singular part $I_{2,1}$, note that

$$\begin{aligned} \int_{\theta|v-v_*| \leq \frac{1}{2}|v|} b(\cos \theta) |W_l' - W_l|^2 d\sigma &\lesssim |v - v_*|^2 < v >^{2(l-1)} \frac{|v|^{2-2s}}{|v - v_*|^{2-2s}} \\ &\lesssim |v - v_*|^{2s} < v >^{2l-2s} \lesssim < v >^{2l} < v_* >^{2s}; \end{aligned}$$

while for the non singular part $I_{2,2}$, one has

$$\int_{\theta|v-v_*| \geq \frac{1}{2}|v|} b(\cos \theta) |W_l' - W_l|^2 d\sigma \lesssim |v - v_*|^{2s} < v >^{2l} < v_* >^{2l} |v|^{-2s}.$$

Therefore, we have

$$I_{2,1} \lesssim \|W_l g\|_{L_{\gamma/2}^2},$$

and

$$I_{2,2} \lesssim R^{-2s} \|W_l g\|_{L_{s+\gamma/2}^2}.$$

By taking R large enough, we complete the proof of (4.10).

Now we turn to the derivatives in v variable. For $\beta \in \mathbb{N}^3 \setminus \{0\}$, we have

$$\partial_\beta \mathcal{L}_1(g) = \mathcal{L}_1(\partial_\beta g) + \sum_{\beta_1 < \beta} C_{\beta_1, \beta_2} \mathcal{T}(\partial_{\beta_2} \mu, \partial_{\beta_1} g, \partial_{\beta_3} \mu).$$

By (4.10), we have

$$\begin{aligned} (W_l \partial_\beta \mathcal{L}_1(g), W_l \partial_\beta g) &= (W_l \mathcal{L}_1(\partial_\beta g), W_l \partial_\beta g) + \sum_{\beta_1 < \beta} C_{\beta_1, \beta_2} (W_l \mathcal{T}(\partial_{\beta_2} \mu, \partial_{\beta_1} g, \partial_{\beta_3} \mu), W_l \partial_\beta g) \\ &\geq \frac{\eta_0}{2} \|W_l \partial_\beta g\|_{\Phi_\gamma}^2 - C \|W_l \partial_\beta g\|_{L_{\gamma/2}^2} + II, \end{aligned}$$

where

$$II = \sum_{\beta_1 < \beta} C_{\beta_1, \beta_2} (W_l \mathcal{T}(\partial_{\beta_2} \mu, \partial_{\beta_1} g, \partial_{\beta_3} \mu); W_l \partial_\beta g).$$

Recall that the operator \mathcal{T} shares the same commutator properties as Γ . As in the proofs given in Section 2, the linearized operator $\mathcal{T}(\partial_{\beta_2} \mu, \partial_{\beta_1} g, \partial_{\beta_3} \mu)$ satisfies

$$|(W_l \mathcal{T}(\partial_{\beta_2} \mu, \partial_{\beta_1} g, \partial_{\beta_3} \mu); W_l \partial_\beta g)| \lesssim \|W_l \partial_{\beta_1} g\|_{\Phi_\gamma} \|W_l \partial_\beta g\|_{\Phi_\gamma}.$$

Hence

$$|II| \lesssim \left(\sum_{\beta_1 < \beta} \|W_l \partial_{\beta_1} g\|_{\Phi_\gamma} \right) \|W_l \partial_\beta g\|_{\Phi_\gamma}.$$

This completes the proof of the proposition. \square

5. LOCAL EXISTENCE

In the following two subsections, we prove Theorem 1.3 and the local existence of solutions in the function space considered in Theorem 1.4.

5.1. Classical solutions. We now proceed to the proof of Theorem 1.3. The restriction of soft potential $\gamma + 2s \leq 0$ will play an important role.

Consider the following Cauchy problem for a linear Boltzmann equation with a given function f ,

$$(5.1) \quad \partial_t g + v \cdot \nabla_x g + \mathcal{L}_1 g = \Gamma(f, g) - \mathcal{L}_2 f, \quad g|_{t=0} = g_0,$$

which is equivalent to the problem:

$$\partial_t G + v \cdot \nabla_x G = Q(F, G), \quad G|_{t=0} = G_0,$$

with $F = \mu + \sqrt{\mu} f$ and $G = \mu + \sqrt{\mu} g$. The proof is based on energy estimates in the functional space $\mathcal{H}_\ell^N(\mathbb{R}^6)$.

For $N \geq 6$, $\ell \geq N$ and $\alpha, \beta \in \mathbb{N}^3$, $|\alpha| + |\beta| \leq N$, taking

$$\varphi(t, x, v) = (-1)^{|\alpha|+|\beta|} \partial_\beta^\alpha W_{\ell-|\beta|}^2 \partial_\beta^\alpha g(t, x, v),$$

as a test function for equation (5.1), we get

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \|W_{\ell-|\beta|} \partial_\beta^\alpha g\|_{L^2(\mathbb{R}^6)}^2 + (W_{\ell-|\beta|} \partial_\beta^\alpha \mathcal{L}_1(g), W_{\ell-|\beta|} \partial_\beta^\alpha g)_{L^2(\mathbb{R}^6)} \\ & \quad + (W_{\ell-|\beta|} [\partial_\beta^\alpha, v] \cdot \nabla_x g, W_{\ell-|\beta|} \partial_\beta^\alpha g)_{L^2(\mathbb{R}^6)} \\ & = (W_{\ell-|\beta|} \partial_\beta^\alpha \Gamma(f, g), W_{\ell-|\beta|} \partial_\beta^\alpha g)_{L^2(\mathbb{R}^6)} - (W_{\ell-|\beta|} \partial_\beta^\alpha \mathcal{L}_2(f), W_{\ell-|\beta|} \partial_\beta^\alpha g)_{L^2(\mathbb{R}^6)}, \end{aligned}$$

where we have used the fact that

$$(v \cdot \nabla_x (W_{\ell-|\beta|} \partial_\beta^\alpha g), W_{\ell-|\beta|} \partial_\beta^\alpha g)_{L^2(\mathbb{R}^6)} = 0.$$

We have immediately

$$\left| (W_{\ell-|\beta|} [\partial_\beta^\alpha, v] \cdot \nabla_x g, W_{\ell-|\beta|} \partial_\beta^\alpha g)_{L^2(\mathbb{R}^6)} \right| \lesssim \|f\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \|g\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}.$$

By Proposition 4.5, one has

$$\left| (W_{\ell-|\beta|} \partial_\beta^\alpha \mathcal{L}_2(f), W_{\ell-|\beta|} \partial_\beta^\alpha g)_{L^2(\mathbb{R}^6)} \right| \lesssim \|g\|_{\mathcal{H}_\ell^{|\alpha|+|\beta|}(\mathbb{R}^6)}^2.$$

Now by using (4.4) with $f = \mu$, we have

$$\left| (W_{\ell-|\beta|} \partial_\beta^\alpha \mathcal{L}_1(g) - \mathcal{L}_1(W_{\ell-|\beta|} \partial_\beta^\alpha g), W_{\ell-|\beta|} \partial_\beta^\alpha g)_{L^2(\mathbb{R}^6)} \right| \leq C_\delta \|g\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}^2 + \delta \|g\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)}^2.$$

Applying Proposition 4.6, we get for any $|\alpha| + |\beta| \leq N$,

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \|W_{\ell-|\beta|} \partial_\beta^\alpha g\|_{L^2(\mathbb{R}^6)}^2 + (\mathcal{L}_1(W_{\ell-|\beta|} \partial_\beta^\alpha g), W_{\ell-|\beta|} \partial_\beta^\alpha g)_{L^2(\mathbb{R}^6)} \\ & \lesssim \|f\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \|g\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)}^2 + \|f\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)} \|g\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \|g\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)} \\ & \quad + \|g\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}^2 + \|f\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \|g\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)} + \delta \|g\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)}^2. \end{aligned}$$

By using now the coercivity estimate from Section 2 and Lemma 4.1, and by taking summation over $|\beta| \leq N$, the Cauchy-Schwarz inequality and soft potential assumption imply that

$$(5.2) \quad \begin{aligned} & \frac{d}{dt} \|g\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}^2 + \frac{\eta_0}{2} \|g\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)}^2 \leq C \left\{ \|f\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \|g\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)}^2 \right. \\ & \quad \left. + \|g\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \|f\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)}^2 + \|g\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}^2 + \|f\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}^2 \right\}. \end{aligned}$$

In conclusion, we are ready to prove the following proposition.

Proposition 5.1. *Assume that $0 < s < 1$, $\gamma > -3$ and let $N \geq 6, \ell \geq N$. Suppose that $g_0 \in \mathcal{H}_\ell^N(\mathbb{R}^6)$ and*

$$f \in L^\infty([0, T]; \mathcal{H}_\ell^N(\mathbb{R}^6)) \cap L^2([0, T]; \mathcal{B}_\ell^N(\mathbb{R}^6)).$$

If $g \in L^\infty([0, T]; \mathcal{H}_\ell^N(\mathbb{R}^6)) \cap L^2([0, T]; \mathcal{B}_\ell^N(\mathbb{R}^6))$ is a solution of the Cauchy problem (5.1), then there exists $\epsilon_0 > 0$ such that if

$$(5.3) \quad \|f\|_{L^\infty([0, T]; \mathcal{H}_\ell^N(\mathbb{R}^6))}^2 + \|f\|_{L^2([0, T]; \mathcal{B}_\ell^N(\mathbb{R}^6))}^2 \leq \epsilon_0^2,$$

we have

$$(5.4) \quad \|g\|_{L^\infty([0, T]; \mathcal{H}_\ell^N(\mathbb{R}^6))}^2 + \|g\|_{L^2([0, T]; \mathcal{B}_\ell^N(\mathbb{R}^6))}^2 \leq Ce^{CT} (\|g_0\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}^2 + \epsilon_0^2 T),$$

for a constant $C > 0$ depending only on N and ℓ .

Proof. From (5.2), we have, for $t \in]0, T[$,

$$\begin{aligned} \|g(t)\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}^2 + \frac{\eta_0}{2} e^{Ct} \int_0^t e^{-Cs} \|g(s)\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)}^2 ds &\leq e^{CT} \|g_0\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}^2 \\ &+ Ce^{CT} \left\{ \int_0^t e^{-Cs} \|f(s)\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \|g(s)\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)}^2 ds \right. \\ &\left. + \int_0^t e^{-Cs} \|g(s)\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}^2 \|f(s)\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)}^2 ds + \int_0^t e^{-Cs} \|f(s)\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}^2 ds \right\}. \end{aligned}$$

Then

$$\begin{aligned} \|g\|_{L^\infty([0, T]; \mathcal{H}_\ell^N(\mathbb{R}^6))}^2 + \frac{\eta_0}{2} \|g\|_{L^2([0, T]; \mathcal{B}_\ell^N(\mathbb{R}^6))}^2 &\leq e^{CT} \|g_0\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}^2 \\ &+ Ce^{CT} \left\{ \|f\|_{L^\infty([0, T]; \mathcal{H}_\ell^N(\mathbb{R}^6))} \|g\|_{L^2([0, T]; \mathcal{B}_\ell^N(\mathbb{R}^6))}^2 \right. \\ &\left. + \|g\|_{L^\infty([0, T]; \mathcal{H}_\ell^N(\mathbb{R}^6))}^2 \|f\|_{L^2([0, T]; \mathcal{B}_\ell^N(\mathbb{R}^6))}^2 + T \|f\|_{L^\infty([0, T]; \mathcal{H}_\ell^N(\mathbb{R}^6))}^2 \right\}. \end{aligned}$$

Hence, if we choose

$$Ce^{CT} \epsilon_0 \leq \frac{\eta_0}{4}, \quad Ce^{CT} \epsilon_0^2 \leq \frac{1}{2},$$

then (5.3) implies that

$$\begin{aligned} \frac{1}{2} \|g\|_{L^\infty([0, T]; \mathcal{H}_\ell^N(\mathbb{R}^6))}^2 + \frac{\eta_0}{4} \|g\|_{L^2([0, T]; \mathcal{B}_\ell^N(\mathbb{R}^6))}^2 \\ \leq e^{CT} \|g_0\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}^2 + Ce^{CT} T \epsilon_0^2. \end{aligned}$$

And this completes the proof of the proposition. \square

From the energy estimate (5.4), one can deduce the local existence as in [7], and we have proved the following precise version of Theorem 1.3.

Theorem 5.2. *Under the assumptions of Theorem 1.3, let $N \geq 6, \ell \geq N$. There exist $\epsilon_1, T > 0$ such that if $g_0 \in \mathcal{H}_\ell^N(\mathbb{R}^6)$ and*

$$\|g_0\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)} \leq \epsilon_1,$$

then the Cauchy problem (1.3) admits a solution

$$g \in L^\infty([0, T]; \mathcal{H}_\ell^N(\mathbb{R}^6)) \cap L^2([0, T]; \mathcal{B}_\ell^N(\mathbb{R}^6)).$$

Remark 5.3. *By using the Proposition 4.7, we can get the same results as Theorem 5.2 if we replace $\mathcal{H}_\ell^N(\mathbb{R}^6), \mathcal{B}_\ell^N(\mathbb{R}^6)$ by $\tilde{\mathcal{H}}_\ell^N(\mathbb{R}^6), \tilde{\mathcal{B}}_\ell^N(\mathbb{R}^6)$ respectively. In other words, Theorem 1.3 holds also in the function space $\tilde{\mathcal{H}}_\ell^N(\mathbb{R}^6)$.*

5.2. L^2 -solutions. Under some more restrictive conditions on the parameters γ and s , we can prove local existence of solutions with only differentiation in the x variable. That is, we will deduce the energy estimate for the equation (5.1) in the function space $H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))$. For $N \geq 3$ and $\beta \in \mathbb{N}^3$, $|\beta| \leq N$, by taking

$$\varphi(t, x, v) = (-1)^{|\beta|} \partial_x^\beta \partial_x^\beta g(t, x, v),$$

as a test function on $\mathbb{R}_x^3 \times \mathbb{R}_v^3$, we get

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \|\partial_x^\beta g\|_{L^2(\mathbb{R}^6)}^2 + (\partial_x^\beta \mathcal{L}_1(g), \partial_x^\beta g)_{L^2(\mathbb{R}^6)} \\ &= (\partial_x^\beta \Gamma(f, g), \partial_x^\beta g)_{L^2(\mathbb{R}^6)} - (\partial_x^\beta \mathcal{L}_2(f), \partial_x^\beta g)_{L^2(\mathbb{R}^6)}, \end{aligned}$$

where we have used the fact that $(v \cdot \nabla_x (\partial_x^\beta g), \partial_x^\beta g)_{L^2(\mathbb{R}^6)} = 0$.

Applying now Proposition 4.3, Proposition 4.5, we get for any $N \geq 3$ and $|\beta| \leq N$,

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \|\partial^\beta g\|_{L^2(\mathbb{R}^6)}^2 + (\mathcal{L}_1(\partial_x^\beta g), \partial_x^\beta g)_{L^2(\mathbb{R}^6)} \\ & \lesssim \|f\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \|g\|_{\mathcal{X}^N(\mathbb{R}^6)}^2 + \|f\|_{\mathcal{X}^N(\mathbb{R}^6)} \|g\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \|g\|_{\mathcal{X}^N(\mathbb{R}^6)} \\ & \quad + \|g\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))}^2 + \|f\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \|g\|_{\mathcal{X}^N(\mathbb{R}^6)} + \|g\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \|g\|_{\mathcal{X}^N(\mathbb{R}^6)}. \end{aligned}$$

By using now the coercivity estimate (4.2), and by taking summation over $|\beta| \leq N$, the Cauchy-Schwarz inequality leads to

$$(5.5) \quad \begin{aligned} & \frac{d}{dt} \|g\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))}^2 + \frac{\eta_0}{2} \|g\|_{\mathcal{X}^N(\mathbb{R}^6)}^2 \leq C \left\{ \|f\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \|g\|_{\mathcal{X}^N(\mathbb{R}^6)}^2 \right. \\ & \quad \left. + \|g\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))}^2 \|f\|_{\mathcal{X}^N(\mathbb{R}^6)}^2 + \|g\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))}^2 + \|f\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))}^2 \right\}. \end{aligned}$$

We are now ready to prove the following

Proposition 5.4. *Under the assumptions of Theorem 1.4, let $N \geq 3$, $g_0 \in H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))$ and*

$$f \in L^\infty([0, T]; H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))) \cap L^2([0, T]; \mathcal{X}^N(\mathbb{R}^6)).$$

If $g \in L^\infty([0, T]; H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))) \cap L^2([0, T]; \mathcal{X}^N(\mathbb{R}^6))$ is a solution of the Cauchy problem (5.1), then there exists $\epsilon_0 > 0$ such that if

$$(5.6) \quad \|f\|_{L^\infty([0, T]; H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3)))}^2 + \|f\|_{L^2([0, T]; \mathcal{X}^N(\mathbb{R}^6))}^2 \leq \epsilon_0^2,$$

we have

$$(5.7) \quad \|g\|_{L^\infty([0, T]; H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3)))}^2 + \|g\|_{L^2([0, T]; \mathcal{X}^N(\mathbb{R}^6))}^2 \leq C e^{CT} \left(\|g_0\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))}^2 + \epsilon_0^2 T \right),$$

for a constant $C > 0$ depending only on N .

Proof. From (5.5), we have, for $t \in]0, T[$,

$$\begin{aligned} & \|g(t)\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))}^2 + \frac{\eta_0}{2} e^{Ct} \int_0^t e^{-Cs} \|g(s)\|_{\mathcal{X}^N(\mathbb{R}^6)}^2 ds \leq e^{CT} \|g_0\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))}^2 \\ & \quad + C e^{CT} \left\{ \int_0^t e^{-Cs} \|f(s)\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \|g(s)\|_{\mathcal{X}^N(\mathbb{R}^6)}^2 ds \right. \\ & \quad \left. + \int_0^t e^{-Cs} \|g(s)\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))}^2 \|f(s)\|_{\mathcal{X}^N(\mathbb{R}^6)}^2 ds + \int_0^t e^{-Cs} \|f(s)\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))}^2 ds \right\}. \end{aligned}$$

Then

$$\begin{aligned} & \|g\|_{L^\infty([0,T];H^N(\mathbb{R}_x^3;L^2(\mathbb{R}_v^3)))}^2 + \frac{\eta_0}{2} \|g\|_{L^2([0,T];\mathcal{X}^N(\mathbb{R}^6))}^2 \leq e^{CT} \|g_0\|_{H^N(\mathbb{R}_x^3;L^2(\mathbb{R}_v^3))}^2 \\ & + C e^{CT} \left\{ \|f\|_{L^\infty([0,T];H^N(\mathbb{R}_x^3;L^2(\mathbb{R}_v^3)))} \|g\|_{L^2([0,T];\mathcal{X}^N(\mathbb{R}^6))}^2 \right. \\ & \left. + \|g\|_{L^\infty([0,T];H^N(\mathbb{R}_x^3;L^2(\mathbb{R}_v^3)))}^2 \|f\|_{L^2([0,T];\mathcal{X}^N(\mathbb{R}^6))}^2 + T \|f\|_{L^\infty([0,T];H^N(\mathbb{R}_x^3;L^2(\mathbb{R}_v^3)))}^2 \right\}. \end{aligned}$$

By choosing T small as in Proposition 5.1, we complete the proof. \square

As in [7], the energy estimate (5.7) yields

Theorem 5.5. *Under the assumptions of Theorem 1.4, for $N \geq 3$, there exist $\epsilon_1, T > 0$ such that if $g_0 \in H^N(\mathbb{R}_x^3;L^2(\mathbb{R}_v^3))$ and*

$$\|g_0\|_{H^N(\mathbb{R}_x^3;L^2(\mathbb{R}_v^3))} \leq \epsilon_1,$$

then the Cauchy problem (1.3) admits a solution

$$g \in L^\infty([0, T]; H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))) \cap L^2([0, T]; \mathcal{X}^N(\mathbb{R}^6)).$$

6. GLOBAL SOLUTIONS

We are now ready to prove the global existence of weak and classical solutions in the following two subsections.

6.1. L^2 -solutions. We now conclude for the global existence issue in Theorem 1.4. We already gave the macro-micro decomposition of solutions introduced in [20]:

$$\begin{aligned} g &= \mathbf{P}g + (\mathbf{I} - \mathbf{P})g = g_1 + g_2, \\ g_1 &= (a + v \cdot b + |v|^2 c) \sqrt{\mu}, \quad \mathcal{A} = (a, b, c). \end{aligned}$$

Notice that

$$\begin{aligned} \|g\|_{H^N(\mathbb{R}_x^3;L^2(\mathbb{R}_v^3))}^2 &\sim \|\mathcal{A}\|_{H^N(\mathbb{R}^3)}^2 + \|g_2\|_{H^N(\mathbb{R}_x^3;L^2(\mathbb{R}_v^3))}^2, \\ \|g\|_{\mathcal{X}^N(\mathbb{R}^6)}^2 &\sim \|\mathcal{A}\|_{H^N(\mathbb{R}^3)}^2 + \|g_2\|_{\mathcal{X}^N(\mathbb{R}^6)}^2. \end{aligned}$$

The temporal energy functional and dissipation integral of solutions are defined by

$$\begin{aligned} \mathcal{E}_N &= \|g\|_{H^N(\mathbb{R}_x^3;L^2(\mathbb{R}_v^3))}^2 = \|g_1\|_{H^N(\mathbb{R}_x^3;L^2(\mathbb{R}_v^3))}^2 + \|g_2\|_{H^N(\mathbb{R}_x^3;L^2(\mathbb{R}_v^3))}^2 \\ &\sim \|\mathcal{A}\|_{H^N(\mathbb{R}^3)}^2 + \|g_2\|_{H^N(\mathbb{R}_x^3;L^2(\mathbb{R}_v^3))}^2, \\ \mathcal{D}_N &= \|\nabla_x g_1\|_{H^{N-1}(\mathbb{R}_x^3;L^2(\mathbb{R}_v^3))}^2 + \|g_2\|_{\mathcal{X}^N(\mathbb{R}^6)}^2 \sim \|\nabla_x \mathcal{A}\|_{H^{N-1}(\mathbb{R}^3)}^2 + \|g_2\|_{\mathcal{X}^N(\mathbb{R}^6)}^2, \end{aligned}$$

respectively. Let $g = g(t, x, v)$ be a solution to

$$(6.1) \quad g_t + v \cdot \nabla_x g + \mathcal{L}g = \Gamma(g, g), \quad g|_{t=0} = g_0.$$

We start with the macroscopic energy estimate. It is well-known that the macroscopic component $g_1 = \mathbf{P}g \sim \mathcal{A} = (a, b, c)$, satisfies the following set of equations

$$(6.2) \quad \begin{cases} v_i |v|^2 \mu^{1/2} : & \nabla_x c = -\partial_t r_c + l_c + h_c, \\ v_i^2 \mu^{1/2} : & \partial_t c + \partial_i b_i = -\partial_t r_i + l_i + h_i, \\ v_i v_j \mu^{1/2} : & \partial_i b_j + \partial_j b_i = -\partial_t r_{ij} + l_{ij} + h_{ij}, \quad i \neq j, \\ v_i \mu^{1/2} : & \partial_t b_i + \partial_i a = -\partial_t r_{bi} + l_{bi} + h_{bi}, \\ \mu^{1/2} : & \partial_t a = -\partial_t r_a + l_a + h_a, \end{cases}$$

where

$$r = (g_2, e)_{L^2(\mathbb{R}_v^3)}, \quad l = -(v \cdot \nabla_x g_2 + \mathcal{L}g_2, e)_{L^2(\mathbb{R}_v^3)}, \quad h = (\Gamma(g, g), e)_{L^2(\mathbb{R}_v^3)},$$

stand for r_c, \dots, h_a , while

$$e \in \text{span}\{v_i|v|^2\mu^{1/2}, v_i^2\mu^{1/2}, v_iv_j\mu^{1/2}, v_i\mu^{1/2}, \mu^{1/2}\}.$$

Lemma 6.1. *Assume $N \geq 3$ and let $\partial^\alpha = \partial_x^\alpha$, $\alpha \in \mathbb{N}^3$, $|\alpha| \leq N$. Then,*

$$\|\partial^\alpha \mathcal{A}^2\|_{L^2(\mathbb{R}_x^3)} \leq \|\nabla_x \mathcal{A}\|_{H^{N-1}(\mathbb{R}_x^3)} \|\mathcal{A}\|_{H^{N-1}(\mathbb{R}_x^3)}.$$

Proof. Firstly, one has, for $|\alpha| = 0$

$$\|\mathcal{A}^2\|_{L^2(\mathbb{R}_x^3)} \leq \|\mathcal{A}\|_{L^6(\mathbb{R}_x^3)} \|\mathcal{A}\|_{L^3(\mathbb{R}_x^3)} \leq \|\nabla_x \mathcal{A}\|_{L^2(\mathbb{R}_x^3)} \|\mathcal{A}\|_{H^1(\mathbb{R}_x^3)}.$$

Also for $|\alpha| = 1$, we have

$$\begin{aligned} \|\partial \mathcal{A}^2\|_{L^2(\mathbb{R}_x^3)} &\leq \|\mathcal{A} \partial \mathcal{A}\|_{L^2(\mathbb{R}_x^3)} \leq \|\mathcal{A}\|_{L^\infty(\mathbb{R}_x^3)} \|\nabla_x \mathcal{A}\|_{L^2(\mathbb{R}_x^3)} \\ &\leq \|\mathcal{A}\|_{H^{N-1}(\mathbb{R}_x^3)} \|\nabla_x \mathcal{A}\|_{L^2(\mathbb{R}_x^3)}, \end{aligned}$$

and for $2 \leq |\alpha| \leq N$,

$$\begin{aligned} \|\partial^\alpha \mathcal{A}^2\|_{L^2(\mathbb{R}_x^3)} &\leq \sum_{k' \leq \frac{|\alpha|}{2}} \|\partial_x^{k'} \mathcal{A} \partial_x^{|\alpha|-k'} \mathcal{A}\|_{L^2(\mathbb{R}_x^3)} \\ &\leq \sum \|\partial_x^{k'} \mathcal{A}\|_{L^\infty(\mathbb{R}_x^3)} \|\partial_x^{|\alpha|-k'} \mathcal{A}\|_{L^2(\mathbb{R}_x^3)} \leq \|\mathcal{A}\|_{H^{N-1}(\mathbb{R}_x^3)} \|\nabla_x \mathcal{A}\|_{H^{N-1}(\mathbb{R}_x^3)}. \end{aligned}$$

This completes the proof of the lemma. \square

Lemma 6.2. *Assume $\gamma > -3$, $N \geq 3$. Let $\partial^\alpha = \partial_x^\alpha$, $\partial_i = \partial_{x_i}$, $|\alpha| \leq N-1$. Then, one has*

$$(6.3) \quad \|\partial_i \partial^\alpha r\|_{L^2(\mathbb{R}_x^3)} + \|\partial^\alpha l\|_{L^2(\mathbb{R}_x^3)} \leq \|g_2\|_{\mathcal{X}^N(\mathbb{R}^6)} \equiv A_1,$$

$$(6.4) \quad \begin{aligned} \|\partial^\alpha h\|_{L^2(\mathbb{R}_x^3)} &\leq \|\nabla_x \mathcal{A}\|_{H^{N-2}(\mathbb{R}_x^3)} \|\mathcal{A}\|_{H^{N-1}(\mathbb{R}_x^3)} \\ &\quad + \|\mathcal{A}\|_{H^{N-1}(\mathbb{R}_x^3)} \|g_2\|_{\mathcal{X}^N(\mathbb{R}^6)} + \|g_2\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \|g_2\|_{\mathcal{X}^N(\mathbb{R}^6)} \equiv A_2. \end{aligned}$$

Proof. (6.3) follows from

$$\begin{aligned} \|\partial_i \partial^\alpha r\|_{L^2(\mathbb{R}_x^3)} &= \|(\partial_i \partial^\alpha g_2, e)_{L^2(\mathbb{R}_v^3)}\|_{L^2(\mathbb{R}_x^3)} = \|(\tilde{W}^{-1} \partial_i \partial^\alpha g_2, \tilde{W} e)_{L^2(\mathbb{R}_v^3)}\|_{L^2(\mathbb{R}_x^3)} \\ &\leq \| \|\partial_i \partial^\alpha g_2\|_{L^2_{s+\gamma/2}(\mathbb{R}_x^3)} \| \leq \|\nabla g_2\|_{\mathcal{X}^{N-1}(\mathbb{R}^6)}, \end{aligned}$$

and

$$\begin{aligned} \|\partial^\alpha l\|_{L_x^2} &\leq \|(\tilde{W}^{-1} \nabla_x \partial^\alpha g_2, \tilde{W} v e)_{L^2(\mathbb{R}_v^3)}\|_{L^2(\mathbb{R}_x^3)} + \|(\tilde{W}^{-1} \partial^\alpha g_2, \tilde{W} \mathcal{L}^* e)_{L^2(\mathbb{R}_v^3)}\|_{L^2(\mathbb{R}_x^3)} \\ &\leq \|\nabla_x \partial^\alpha g_2\|_{L^2_{s+\gamma/2}(\mathbb{R}^6)} + \|\partial^\alpha g_2\|_{L^2_{s+\gamma/2}(\mathbb{R}^6)} \leq \|g_2\|_{\mathcal{X}^N(\mathbb{R}^6)}. \end{aligned}$$

Here, we have used (4.5).

We prove (6.4) as follows. Firstly, set

$$\begin{aligned} H(f, g) &= (\Gamma(f, g), e)_{L^2(\mathbb{R}_v^3)} \\ &= \iiint b(\cos \theta) |v - v_*|^\gamma f_* g(\mu^{1/2})'_* e' - \mu^{1/2} e) dv dv_* d\sigma \\ &= \iiint b(\cos \theta) |v - v_*|^\gamma (\mu^{1/2} f)_* (\mu^{1/2} g)(q(v') - q(v)) dv dv_* d\sigma, \end{aligned}$$

where $q(v)$ is a polynomial. Of course, $h = H(g, g)$. We now write the Taylor expansion of the second order of the function $q(v)$ as

$$q(v) - q(v') = (\nabla q)(v) \cdot (v' - v) + \frac{1}{2} \int_0^1 \nabla^2 q(v + \tau(v' - v)) d\tau (v' - v)^2.$$

Setting $\mathbf{k} = \frac{v-v_*}{|v-v_*|}$, we recall

$$v' - v = \frac{1}{2}|v - v_*|(\sigma - (\sigma \cdot \mathbf{k})\mathbf{k}) + \frac{1}{2}((\sigma \cdot \mathbf{k}) - 1)(v - v_*),$$

and notice that by virtue of the symmetry

$$\int_{\mathbb{S}^2} b(\sigma \cdot \mathbf{k})(\sigma - (\sigma \cdot \mathbf{k})\mathbf{k})d\sigma = 0.$$

Therefore, we have

$$\begin{aligned} H(f, g) &= \frac{1}{2} \iint (\mu^{1/2} f)(\mu^{1/2} g)_* |v - v_*|^\gamma \left\{ \int_{\mathbb{S}} b(\cos \theta)(\cos \theta - 1)d\sigma \right\} (\nabla q)(v) \cdot (v - v_*) dv dv_* \\ &\quad + \frac{1}{2} \int_0^1 \left(\iiint (\mu^{1/2} f)(\mu^{1/2} g)_* |v - v_*|^\gamma b(\cos \theta) \nabla^2 q(v + \tau(v' - v)) d\tau (v' - v)^2 d\sigma dv dv_* \right) d\tau \\ &= H_1(f, g) + H_2(f, g). \end{aligned}$$

For H_1 , clearly, the integral in σ is bounded for $0 < s < 1$. By the Cauchy-Schwarz inequality, since $\gamma + 1 > -3$,

$$\begin{aligned} |H_1(f, g)| &\leq \iint (|\nabla q(v)| \mu^{1/2} |f|)(\mu^{1/2} |g|)_* |v - v_*|^{\gamma+1} dv dv_* \\ &\leq \iint \mu_*^{1/8} (\mu^{1/8} |f|)(\mu^{1/8} |g|)_* \mu^{1/8} |v - v_*|^{\gamma+1} dv_* dv \\ &\leq \left(\iint (\mu^{1/4} |g|^2)_* \mu^{1/4} |v - v_*|^{\gamma+1} dv_* dv \right)^{1/2} \left(\iint (\mu^{1/4} |f|^2) \mu_*^{1/4} |v - v_*|^{\gamma+1} dv_* dv \right)^{1/2} \\ &\leq \|\langle v \rangle^{(\gamma+1)/2} \mu^{1/8} f\|_{L_v^2} \|\langle v \rangle^{(\gamma+1)/2} \mu^{1/8} g\|_{L_v^2} \leq \|f\|_{L_v^2(\mathbb{R}_v^3)} \|g\|_{L_m^2(\mathbb{R}_v^3)}, \end{aligned}$$

for any $l, m \in \mathbb{R}$. Similarly, $\gamma + 2 > -3$ implies,

$$\begin{aligned} |H_2(f, g)| &\leq \iiint b(\cos \theta) \theta^2 (\mu^{1/2} |f|)(\mu^{1/2} |g|)_* |v - v_*|^{\gamma+2} (|\nabla^2 q(v)| + |\nabla^2 q(v')|) d\sigma dv dv_* \\ &\leq \|\langle v \rangle^{(\gamma+2)/2} \mu^{1/8} f\|_{L_v^2} \|\langle v \rangle^{(\gamma+2)/2} \mu^{1/8} g\|_{L_v^2} \leq \|f\|_{L_v^2(\mathbb{R}_v^3)} \|g\|_{L_m^2(\mathbb{R}_v^3)}. \end{aligned}$$

Combining these two estimates yields

$$(6.5) \quad |H(f, g)| \leq \|f\|_{L_v^2(\mathbb{R}_v^3)} \|g\|_{L_m^2(\mathbb{R}_v^3)}.$$

Now h is computed as follows.

$$h = H(g, g) = \sum_{i,j=1,2} H(g_i, g_j) = \sum_{i,j=1,2} H^{(ij)}.$$

Firstly, we have

$$H^{(11)} \sim \mathcal{A}^2 H(\varphi_k, \varphi_m),$$

where $\varphi_k \in \mathcal{N}$. Applying (6.5) for $l = m = 0$, we get by virtue of Lemma 6.1 and for $|\alpha| \leq N - 1$ that

$$\|\partial^\alpha H^{(11)}\|_{L^2(\mathbb{R}_x^3)} \leq \|\partial^\alpha \mathcal{A}^2\|_{L^2(\mathbb{R}_x^3)} \leq \|\nabla_x \mathcal{A}\|_{H^{N-2}(\mathbb{R}_x^3)} \|\mathcal{A}\|_{H^{N-1}(\mathbb{R}_x^3)};$$

while taking l, m to be 0 or $s + \gamma/2$ in (6.5) and by the Leibniz rule and by (4.5),

$$\begin{aligned}\|\partial^\alpha H^{(12)}\|_{L^2(\mathbb{R}_x^3)} &\leq \sum_{\alpha_1+\alpha_2=\alpha} \|\partial^{\alpha_1} \mathcal{A}\|_{L^2(\mathbb{R}_x^3)} \|\partial^{\alpha_2} g_2\|_{L^2_{s+\gamma/2}(\mathbb{R}_{x,v}^6)} \leq \|\mathcal{A}\|_{H^{N-1}(\mathbb{R}_x^3)} \|g_2\|_{\chi^N(\mathbb{R}^6)}, \\ \|\partial^\alpha H^{(21)}\|_{L^2(\mathbb{R}_x^3)} &\leq \sum_{\alpha_1+\alpha_2=\alpha} \|\partial^{\alpha_1} g_2\|_{L^2_{s+\gamma/2}(\mathbb{R}_{x,v}^6)} \|\partial^{\alpha_2} \mathcal{A}\|_{L^2(\mathbb{R}_x^3)} \leq \|g_2\|_{\chi^N(\mathbb{R}^6)} \|\mathcal{A}\|_{H^{N-1}(\mathbb{R}_x^3)}, \\ \|\partial^\alpha H^{(22)}\|_{L^2(\mathbb{R}_x^3)} &\leq \sum_{\alpha_1+\alpha_2=\alpha} \|\partial^{\alpha_1} g_2\|_{L^2(\mathbb{R}_{x,v}^6)}^2 \|\partial^{\alpha_2} g_2\|_{L^2_{s+\gamma/2}(\mathbb{R}_{x,v}^6)} \leq \|g_2\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \|g_2\|_{\chi^N(\mathbb{R}^6)}.\end{aligned}$$

Now the proof of Lemma 6.2 is completed. \square

Next, we shall prove

Lemma 6.3. *Assume $\gamma > -3$. Let $|\alpha| \leq N - 1$. Then,*

$$(6.6) \quad \|\nabla_x \partial^\alpha \mathcal{A}\|_{L^2(\mathbb{R}_x^3)}^2 \leq -\frac{d}{dt} \left\{ (\partial^\alpha r, \nabla_x \partial^\alpha (a, -b, c))_{L^2(\mathbb{R}_x^3)} + (\partial^\alpha b, \nabla_x \partial^\alpha a)_{L^2(\mathbb{R}_x^3)} \right\} \\ + \|g_2\|_{\chi^N(\mathbb{R}^6)}^2 + \|g_2\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \mathcal{D}_N.$$

Proof. (a) Estimate of $\nabla_x \partial^\alpha a$. Let A_1, A_2 be as in Lemma 6.2. From (6.2),

$$\begin{aligned}\|\nabla_x \partial^\alpha a\|_{L^2(\mathbb{R}_x^3)}^2 &= (\nabla_x \partial^\alpha a, \nabla_x \partial^\alpha a)_{L^2(\mathbb{R}_x^3)} \\ &= (\partial^\alpha (-\partial_t b - \partial_t r + l + h), \nabla_x \partial^\alpha a)_{L^2(\mathbb{R}_x^3)} \\ &\leq R_1 + C_\eta (A_1^2 + A_2^2) + \eta \|\nabla_x \partial^\alpha a\|_{L^2(\mathbb{R}_x^3)}^2.\end{aligned}$$

Here,

$$\begin{aligned}R_1 &= -(\partial^\alpha \partial_t b + \partial^\alpha \partial_t r, \nabla_x \partial^\alpha a)_{L^2(\mathbb{R}_x^3)} \\ &= -\frac{d}{dt} (\partial^\alpha (b + r), \nabla_x \partial^\alpha a)_{L^2(\mathbb{R}_x^3)} + (\nabla_x \partial^\alpha (b + r), \partial_t \partial^\alpha a)_{L^2(\mathbb{R}_x^3)} \\ &\leq -\frac{d}{dt} (\partial^\alpha (b + r), \nabla_x \partial^\alpha a)_{L^2(\mathbb{R}_x^3)} + C_\eta (\|\nabla_x \partial^\alpha b\|_{L^2(\mathbb{R}_x^3)}^2 + A_1^2) + \eta \|\partial_t \partial^\alpha a\|_{L^2(\mathbb{R}_x^3)}^2.\end{aligned}$$

(b) Estimate of $\nabla_x \partial^\alpha b$. From (6.2),

$$\begin{aligned}\Delta_x \partial^\alpha b_i + \partial_i^2 \partial^\alpha b_i &= \sum_{j \neq i} \partial_j \partial^\alpha (\partial_j b_i + \partial_i b_j) + \partial_i \partial^\alpha (2\partial_i b_i - \sum_{j \neq i} \partial_i b_j) \\ &= \partial_i \partial^\alpha (-\partial_t r + l + h), \\ \|\nabla_x \partial^\alpha b\|_{L^2(\mathbb{R}_x^3)}^2 + \|\partial_i \partial^\alpha b\|_{L^2(\mathbb{R}_x^3)}^2 &= -(\Delta_x \partial^\alpha b_i + \partial_i^2 \partial^\alpha b_i, \partial^\alpha b)_{L^2(\mathbb{R}_x^3)} = R_2 + R_3 + R_4,\end{aligned}$$

where

$$\begin{aligned}R_2 &= (\partial_t \partial^\alpha r, \partial_i \partial^\alpha b)_{L^2(\mathbb{R}_x^3)} = \frac{d}{dt} (\partial^\alpha r, \partial_i \partial^\alpha b)_{L^2(\mathbb{R}_x^3)} + (\partial_i \partial^\alpha r, \partial_t \partial^\alpha b)_{L^2(\mathbb{R}_x^3)} \\ &\leq \frac{d}{dt} (\partial^\alpha r, \partial_i \partial^\alpha b)_{L^2(\mathbb{R}_x^3)} + C_\eta A_1^2 + \eta \|\partial_t \partial^\alpha b\|_{L^2(\mathbb{R}_x^3)}^2, \\ R_3 &= -(\partial^\alpha l, \partial_i \partial^\alpha b)_{L^2(\mathbb{R}_x^3)} \leq C_\eta A_1^2 + \eta \|\partial_i \partial^\alpha b\|_{L^2(\mathbb{R}_x^3)}^2, \\ R_4 &= -(\partial^\alpha h, \partial_i \partial^\alpha b)_{L^2(\mathbb{R}_x^3)} \leq C_\eta A_2^2 + \eta \|\partial_i \partial^\alpha b\|_{L^2(\mathbb{R}_x^3)}^2.\end{aligned}$$

(c) Estimate of $\nabla_x \partial^\alpha c$. From (6.2),

$$\begin{aligned}\|\nabla_x \partial^\alpha c\|_{L^2(\mathbb{R}_x^3)}^2 &= (\nabla_x \partial^\alpha c, \nabla_x \partial^\alpha c)_{L^2(\mathbb{R}_x^3)} = (\partial^\alpha (-\partial_t r + l + h), \nabla_x \partial^\alpha c)_{L^2(\mathbb{R}_x^3)} \\ &\leq R_5 + C_\eta (A_1^2 + A_2^2) + \eta \|\nabla_x \partial^\alpha c\|_{L^2(\mathbb{R}_x^3)}^2,\end{aligned}$$

where

$$\begin{aligned} R_5 &= -(\partial^\alpha \partial_t r, \nabla_x \partial^\alpha c)_{L^2(\mathbb{R}_x^3)} = -\frac{d}{dt}(\partial^\alpha r, \nabla_x \partial^\alpha c)_{L^2(\mathbb{R}_x^3)} + (\nabla_x \partial^\alpha r, \partial_t \partial^\alpha c)_{L^2(\mathbb{R}_x^3)} \\ &\leq -\frac{d}{dt}(\partial^\alpha r, \nabla_x \partial^\alpha c)_{L^2(\mathbb{R}_x^3)} + C_\eta A_1^2 + \eta \|\partial_t \partial^\alpha c\|_{L^2(\mathbb{R}_x^3)}^2. \end{aligned}$$

(d) Estimate of $\partial_t \partial^\alpha(a, b, c)$.

$$\begin{aligned} \|\partial_t \partial^\alpha \mathcal{A}\|_{L^2(\mathbb{R}_x^3)} &= \|\partial^\alpha \partial_t \mathbf{P}g\|_{L^2(\mathbb{R}_{x,v}^6)} \\ &= \|\partial^\alpha \mathbf{P}(-v \cdot \nabla_x g - \mathcal{L}g + \Gamma(g, g))\|_{L^2(\mathbb{R}_{x,v}^6)} = \|\partial^\alpha \mathbf{P}(v \cdot \nabla_x g)\|_{L^2(\mathbb{R}_{x,v}^6)} \\ &\leq \|\nabla_x \partial^\alpha \mathcal{A}\|_{L^2(\mathbb{R}_x^3)} + \|\nabla_x \partial^\alpha \mathbf{P}(vw^{-1}wg_2)\|_{L^2(\mathbb{R}_{x,v}^6)} \\ &\leq \|\nabla_x \partial^\alpha(a, b, c)\|_{L^2(\mathbb{R}_x^3)} + \|\nabla_x \partial^\alpha g_2\|_{L^2_{s+\gamma/2}(\mathbb{R}_{x,v}^6)}. \end{aligned}$$

By combining the above estimates and taking $\eta > 0$ sufficiently small, we deduce

$$\begin{aligned} \|\nabla_x \partial^\alpha \mathcal{A}\|_{L^2(\mathbb{R}_x^3)}^2 &\leq -\frac{d}{dt} \left\{ (\partial^\alpha r, \nabla_x \partial^\alpha(a, -b, c))_{L^2(\mathbb{R}_x^3)} + (\partial^\alpha b, \nabla_x \partial^\alpha a)_{L^2(\mathbb{R}_x^3)} \right\} \\ &\quad + A_1^2 + A_2^2 + \eta \|\nabla_x \partial^\alpha g_2\|_{L^2_{s+\gamma/2}(\mathbb{R}_{x,v}^6)}. \end{aligned}$$

Finally, by choosing $|\alpha| \leq N-1$, and using Lemma 6.2, we obtain

$$A_1^2 + A_2^2 + \eta \|\nabla_x \partial^\alpha g_2\|_{L^2_{s+\gamma/2}(\mathbb{R}_{x,v}^6)}^2 \leq \mathcal{D}_N \|g\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} + \eta \|g\|_{\chi^N(\mathbb{R}^3)}^2,$$

which completes the proof of Lemma 6.3. \square

We now turn to the estimation on the microscopic component g_2 in the function space $H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))$. Actually, we shall establish

Lemma 6.4. *Under the assumptions of Theorem 1.4, for $N \geq 3$,*

$$(6.7) \quad \frac{d}{dt} \mathcal{E}_N + \|g_2\|_{\chi^N(\mathbb{R}^6)}^2 \lesssim \mathcal{E}_N^{1/2} \mathcal{D}_N.$$

Proof. We apply ∂_x^α to (6.1) and take the $L^2(\mathbb{R}^6)$ inner product with $\partial_x^\alpha g$. Since the inner product including $v \cdot \nabla_x g$ vanishes by integration by parts, we get

$$(6.8) \quad \frac{1}{2} \frac{d}{dt} \|\partial_x^\alpha g\|_N^2 + (\mathcal{L} \partial_x^\alpha g, \partial_x^\alpha g)_{L^2(\mathbb{R}^6)} = (\partial_x^\alpha \Gamma(g, g), \partial_x^\alpha g)_{L^2(\mathbb{R}^6)}.$$

In view of Section 2, we have, for all $\gamma > -3$,

$$\sum_{|\alpha| \leq N} (\mathcal{L} \partial_x^\alpha g, \partial_x^\alpha g)_{L^2(\mathbb{R}^6)} \geq \eta_0 \sum_{|\alpha| \leq N} \int_{\mathbb{R}_x^3} \|\partial_x^\alpha (\mathbf{I} - \mathbf{P})g\|^2 dx = \eta_0 \|g_2\|_{\chi^N(\mathbb{R}^6)}^2,$$

while we shall show below that for $|\alpha| \leq N$, $N \geq 3$,

$$(6.9) \quad \left| (\partial_x^\alpha \Gamma(g, g), \partial_x^\alpha g)_{L^2(\mathbb{R}^6)} \right| \lesssim \mathcal{E}_N^{1/2} \mathcal{D}_N.$$

Lemma 6.4 can then be concluded by plugging these two estimates into (6.8).

Proof of (6.9): Write

$$\begin{aligned} (\partial_x^\alpha \Gamma(g, g), \partial_x^\alpha g)_{L^2(\mathbb{R}^6)} &= J^{11} + J^{12} + J^{21} + J^{22}, \\ J^{ij} &= (\partial_x^\alpha \Gamma(g_i, g_j), \partial_x^\alpha g_2)_{L^2(\mathbb{R}^6)}. \end{aligned}$$

Estimation of J^{11} . We shall estimate

$$J^{11} \sim \int_{\mathbb{R}_x^3} (\partial_x^{\alpha_1} \mathcal{A})(\partial_x^{\alpha_2} \mathcal{A})(\Gamma(\varphi_k, \varphi_m), \partial_x^\alpha g_2)_{L^2(\mathbb{R}_v^3)} dx.$$

Firstly, for any function $\phi, \psi \in \mathcal{N}$, set

$$\Psi_1 = (\Gamma(\phi, \psi), h)_{L^2(\mathbb{R}^3)}.$$

We shall prove that for $\gamma > -3$,

$$(6.10) \quad |\Psi_1| \leq \|h\|_{L_m^2(\mathbb{R}^3)},$$

holds for any $m \in \mathbb{R}$.

Proof of (6.10). Notice that

$$\begin{aligned} \Psi_1 &= \iiint b(\cos \theta) |v - v_*|^\gamma (\mu_*)^{1/2} ((\phi_*)' \psi' - \phi_* \psi) h \, dv dv_* d\sigma \\ &= \iiint b(\cos \theta) |v - v_*|^\gamma (\mu_*)^{1/2} (\mu_*)^{1/2} (\mu)^{1/2} (q_*' r' - q_* r) h \, dv dv_* d\sigma, \end{aligned}$$

where $q = q(v)$ and $r = r(v)$ are some polynomials. First, write

$$q_*' r' - q_* r = (q_*' - q_*) (r' - r) + (q_*' - q_*) r + q_* (r' - r) = S_1 + S_2 + S_3,$$

and make a decomposition

$$\begin{aligned} \Psi_1 &= \sum_{i=1}^3 \iiint b(\cos \theta) |v - v_*|^\gamma (\mu_*)^{1/2} (\mu_*)^{1/2} (\mu)^{1/2} S_i h \, dv dv_* d\sigma \\ &= \Psi_1^1 + \Psi_1^2 + \Psi_1^3. \end{aligned}$$

Since $|S_1| \leq R_1(v, v_*) |v - v_*|^2 \theta^2$ where $R_1(v, v_*)$ is a polynomial of v, v_* , it holds that

$$\begin{aligned} |\Psi_1^1| &\lesssim \iint |v - v_*|^{\gamma+2} \left[\int_{\mathbb{S}^2} b(\cos \theta) \theta^2 d\sigma \right] (\mu_*)^{1/2} (\mu_*)^{1/2} (\mu)^{1/2} |R_1(v, v_*)| |h| \, dv dv_* \\ &\lesssim \int \langle v \rangle^{\gamma+2} \mu^{1/4} |h| \, dv \lesssim \|h\|_{L_m^2(\mathbb{R}^3)}, \end{aligned}$$

for any $m \in \mathbb{R}$.

On the other hand, the Taylor expansion of the second order gives

$$q(v_*') - q(v_*) = (\nabla q)(v_*) \cdot (v_*' - v_*) + \frac{1}{2} \int_0^1 \nabla^2 q(v_* + \tau(v_*' - v_*)) d\tau (v_*' - v_*)^2.$$

Since

$$\left| \int_0^1 \nabla^2 q(v_* + \tau(v_*' - v_*)) d\tau (v_*' - v_*)^2 \right| \lesssim |R_2(v, v_*)| |v - v_*|^2 \theta^2,$$

where R_2 is a polynomial of v, v_* , by symmetry, we have

$$\begin{aligned} |\Psi_1^2| &\lesssim \iint |v - v_*|^{\gamma+1} (\mu_*)^{1/2} (\mu_*)^{1/2} (\mu)^{1/2} (|\nabla q(v_*)| + |R_2(v, v_*)| |v - v_*|) |r(v) h| \, dv dv_* \\ &\lesssim \int (\langle v \rangle^{\gamma+1} + \langle v \rangle^{\gamma+2}) \mu^{1/4} |r(v) h| \, dv \lesssim \|h\|_{L_m^2(\mathbb{R}^3)}, \end{aligned}$$

for any $m \in \mathbb{R}$.

The estimation on Ψ_1^3 can be carried out exactly in the same way to have

$$|\Psi_1^3| \lesssim \int (\langle v \rangle^{\gamma+1} + \langle v \rangle^{\gamma+2}) \mu^{1/4} |h| \, dv \lesssim \|h\|_{L_m^2(\mathbb{R}^3)},$$

for any $m \in \mathbb{R}$. This completes the proof of (6.10).

Take $m = s + \gamma/2$ in (6.10) and use (4.5) to obtain $|\Psi_1| \lesssim \|h\|$. Set $h = \partial^\alpha g_2$. Now by the Sobolev embedding theorem, for $\alpha_1 + \alpha_2 = \alpha$, $1 \leq |\alpha| \leq N$, we have

$$\begin{aligned} |J^{11}| &\lesssim \int_{\mathbb{R}_x^3} |\partial_x^{\alpha_1} \mathcal{A}| |\partial_x^{\alpha_2} \mathcal{A}| \|\partial_x^\alpha g_2\|_{\Phi_\gamma} dx \\ &\lesssim \begin{cases} \|\partial_x^{\alpha_1} \mathcal{A}\|_{L_x^\infty} \|\partial_x^{\alpha_2} \mathcal{A}\|_{L_x^2} \|\partial_x^\alpha g_2\|_{\chi^0(\mathbb{R}^6)} \lesssim \|\mathcal{A}\|_{H_x^N} \|\nabla_x \mathcal{A}\|_{H_x^{N-1}} \|g_2\|_{\chi^N(\mathbb{R}^6)}, & |\alpha_1| \leq N-2, \\ \|\partial_x^{\alpha_1} \mathcal{A}\|_{L_x^2} \|\partial_x^{\alpha_2} \mathcal{A}\|_{L_x^\infty} \|\partial_x^\alpha g_2\|_{\chi^0(\mathbb{R}^6)} \lesssim \|\nabla_x \mathcal{A}\|_{H_x^{N-1}} \|\mathcal{A}\|_{H_x^N} \|g_2\|_{\chi^N(\mathbb{R}^6)}, & |\alpha_2| < 1, \end{cases} \\ &\lesssim \|\mathcal{A}\|_{H_x^N} (\|\nabla_x \mathcal{A}\|_{H_x^{N-1}}^2 + \|g_2\|_{\chi^N(\mathbb{R}^6)}^2) = \mathcal{E}_N^{1/2} \mathcal{D}_N, \end{aligned}$$

and for $|\alpha_1 + \alpha_2| = 0$, we have

$$\begin{aligned} \int_{\mathbb{R}_x^3} |\mathcal{A}|^2 \|g_2\| dx &\lesssim \|\mathcal{A}\|_{L_x^6} \|\mathcal{A}\|_{L_x^3} \|g_2\|_{\chi^N(\mathbb{R}^6)} \\ &\lesssim \|\nabla_x \mathcal{A}\|_{L_x^2} \|\mathcal{A}\|_{H_x^1} \|g_2\|_{\chi^N(\mathbb{R}^6)} \lesssim \mathcal{E}_N^{1/2} \mathcal{D}_N. \end{aligned}$$

Estimation of J^{12} : First, notice

$$J^{12} \sim \int_{\mathbb{R}_x^3} (\partial_x^{\alpha_1} \mathcal{A})(\Gamma(\varphi_k, \partial_x^{\alpha_2} g_2), \partial_x^\alpha g_2)_{L^2(\mathbb{R}^3)} dx.$$

For some $\phi \in \mathcal{N}$, set

$$\Psi_2 = (\Gamma(\phi, g), h)_{L^2(\mathbb{R}^3)}.$$

By virtue of Theorem 1.2, we get

$$\begin{aligned} |\Psi_2| &\lesssim \left(\|\phi\|_{L_{s+\gamma/2}^2(\mathbb{R}^3)} \|g\|_{\Phi_\gamma} + \|\phi\|_{\Phi_\gamma} \|g\|_{L_{s+\gamma/2}^2(\mathbb{R}^3)} \right. \\ &\quad \left. + \min\{\|\phi\|_{L^2(\mathbb{R}^3)} \|g\|_{L_{s+\gamma/2}^2(\mathbb{R}^3)}, \|g\|_{L^2} \|\phi\|_{L_{s+\gamma/2}^2(\mathbb{R}^3)}\} \right) \|h\|_{\Phi_\gamma} \\ &\lesssim \left(\|\phi\|_{L_{s+\gamma/2}^2(\mathbb{R}^3)} \|g\|_{\Phi_\gamma} + \|\phi\|_{\Phi_\gamma} \|g\|_{L_{s+\gamma/2}^2(\mathbb{R}^3)} + \|\phi\|_{L^2(\mathbb{R}^3)} \|g\|_{L_{s+\gamma/2}^2(\mathbb{R}^3)} \right) \|h\|_{\Phi_\gamma} \\ &\lesssim \left(\|g\|_{\Phi_\gamma} + \|g\|_{L_{s+\gamma/2}^2(\mathbb{R}^3)} \right) \|h\|_{\Phi_\gamma} \lesssim \|g\|_{\Phi_\gamma} \|h\|_{\Phi_\gamma}, \end{aligned}$$

where we have chosen the first factor in the min term.

Now we have

$$\begin{aligned} |J^{12}| &\lesssim \int_{\mathbb{R}_x^3} \|\partial_x^{\alpha_1} \mathcal{A}\| \|\partial_x^{\alpha_2} g_2\|_{\Phi_\gamma} \|\partial_x^\alpha g_2\|_{\Phi_\gamma} dx \\ &\lesssim \begin{cases} \|\partial_x^{\alpha_1} \mathcal{A}\|_{L_x^\infty} \int_{\mathbb{R}^3} \|\partial_x^{\alpha_2} g_2\|_{\Phi_\gamma} \|\partial_x^\alpha g_2\|_{\Phi_\gamma} dx \lesssim \|\mathcal{A}\|_{H^N(\mathbb{R}^3)} \|g_2\|_{\chi^N(\mathbb{R}^6)}^2, & |\alpha_1| \leq N-2, \\ \|\partial_x^{\alpha_2} g_2\|_{\Phi_\gamma} \|L_x^\infty \int_{\mathbb{R}^3} |\partial_x^{\alpha_1} \mathcal{A}| \|\partial_x^\alpha g_2\|_{\Phi_\gamma} dx \lesssim \|\mathcal{A}\|_{H^N(\mathbb{R}^3)} \|g_2\|_{\chi^N(\mathbb{R}^6)}^2, & |\alpha_2| \leq 1. \end{cases} \end{aligned}$$

Estimation of J^{21} : A similar argument applies to

$$J^{21} \sim \int_{\mathbb{R}_x^3} (\partial_x^{\alpha_2} \mathcal{A})(\Gamma(\partial_x^{\alpha_1} g_2, \varphi_k), \partial_x^\alpha g_2)_{L^2(\mathbb{R}^3)} dx.$$

In fact, for $\phi \in \mathcal{N}$, set

$$\Psi_3 = (\Gamma(f, \phi), h).$$

Again by virtue of Theorem 1.2 and taking the second factor in the min term, (4.5) gives

$$|\Psi_3| \lesssim \|f\|_{\Phi_\gamma} \|h\|_{\Phi_\gamma}.$$

Thus, proceeding as for J^{12} yields

$$|J^{21}| \lesssim \|\mathcal{A}\|_{H^N(\mathbb{R}^3)} \|g_2\|_{\chi^N(\mathbb{R}^6)}^2.$$

Estimation on J^{22} : It follows from the Leibniz rule that

$$|J^{22}| \lesssim \int |(\Gamma(\partial_x^{\alpha_1} g_2, \partial_x^{\alpha_2} g_2), \partial_x^{\alpha} g_2)_{L^2(\mathbb{R}^3)}| dx.$$

Different from the above, we now use Theorem 1.2 and (4.5) in the following way.

$$\begin{aligned} |(\Gamma(f, g), h)_{L^2(\mathbb{R}^3)}| &\lesssim \left(\|f\|_{L^2_{s+\gamma/2}(\mathbb{R}^3)} \|g\|_{\Phi_\gamma} + \|f\|_{\Phi_\gamma} \|g\|_{L^2_{s+\gamma/2}(\mathbb{R}^3)} \right. \\ &\quad \left. + \min \left\{ \|f\|_{L^2(\mathbb{R}^3)} \|g\|_{L^2_{s+\gamma/2}(\mathbb{R}^3)}, \|g\|_{L^2} \|f\|_{L^2_{s+\gamma/2}(\mathbb{R}^3)} \right\} \right) \|h\|_{\Phi_\gamma} \\ &\lesssim \left(\|f\|_{L^2(\mathbb{R}^3)} \|g\|_{\Phi_\gamma} + \|f\|_{\Phi_\gamma} \|g\|_{L^2(\mathbb{R}^3)} \right) \|h\|_{\Phi_\gamma}. \end{aligned}$$

This is valid for the assumptions imposed on γ and s from Theorem 1.4. Then,

$$|J^{22}| \lesssim \int \left(\|\partial_x^{\alpha_1} g_2\|_{L^2(\mathbb{R}^3)} \|\partial_x^{\alpha_2} g_2\|_{\Phi_\gamma} + \|\partial_x^{\alpha_1} g_2\|_{\Phi_\gamma} \|\partial_x^{\alpha_2} g_2\|_{L^2(\mathbb{R}^3)} \right) \|\partial_x^{\alpha} g_2\|_{\Phi_\gamma} dx.$$

Suppose $|\alpha_1| \leq N - 2$. Then, by the Sobolev embedding theorem,

$$\begin{aligned} |J^{22}| &\lesssim \|\partial_x^{\alpha_1} g_2\|_{L^\infty(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \int_{\mathbb{R}^3} \|\partial_x^{\alpha_2} g_2\|_{\Phi_\gamma} \|\partial_x^{\alpha} g_2\|_{\Phi_\gamma} dx \\ &\quad + \|\|\partial_x^{\alpha_1} g_2\|_{\Phi_\gamma}\|_{L^\infty(\mathbb{R}_x^3)} \int_{\mathbb{R}^3} \|\partial_x^{\alpha_2} g_2\|_{L^2(\mathbb{R}_v^3)} \|\partial_x^{\alpha} g_2\|_{\Phi_\gamma} dx \\ &\lesssim \|g_2\|_{H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))} \|g_2\|_{X^N(\mathbb{R}^3)}^2 \lesssim \mathcal{E}_N^{1/2} \mathcal{D}_N. \end{aligned}$$

Similarly, when $|\alpha_1| > N - 2$ then $|\alpha_2| \leq 1$, we get

$$\begin{aligned} |J^{22}| &\lesssim \|\partial_x^{\alpha_2} g_2\|_{L_x^\infty(L_v^2)} \int_{\mathbb{R}^3} \|\partial_x^{\alpha_1} g_2\|_{\Phi_\gamma} \|\partial_x^{\alpha} g_2\|_{\Phi_\gamma} dx \\ &\quad + \|\|\partial_x^{\alpha_2} g_2\|_{\Phi_\gamma}\|_{L_x^\infty} \int_{\mathbb{R}^3} \|\partial_x^{\alpha_1} g_2\|_{L_v^2} \|\partial_x^{\alpha} g_2\|_{\Phi_\gamma} dx \\ &\lesssim \mathcal{E}_N^{1/2} \mathcal{D}_N. \end{aligned}$$

Now, combining the above estimates yields the estimate (6.9) and this completes the proof of the Lemma 6.4. \square

Taking a suitable linear combination of (6.6) and (6.7) gives the

Proposition 6.5. (Global Energy Estimate without Weight) *Under the assumptions of Theorem 1.4, for $N \geq 3$, there exists a constant $C > 0$ such that*

$$\frac{d}{dt} \mathcal{E}_N + \mathcal{D}_N \leq C \mathcal{E}_N^{1/2} \mathcal{D}_N$$

holds as far as g exists.

This proposition assures that a usual continuation argument of local solutions can be carried out under the smallness assumption of initial data. Thus, we established the existence of global solutions in the space $H^N(\mathbb{R}_x^3; L^2(\mathbb{R}_v^3))$.

6.2. Classical solutions. We now turn to the energy estimates involving also v derivatives of solutions. To close this type of energy estimate, we then need to use the weighted norms in the v variable, cf. also Guo [19], with the weight function \tilde{W} . Recall that we assume $s + \gamma/2 \leq 0$. Set

$$\begin{aligned} \mathcal{E}_{N,\ell} &= \mathcal{E}_N + \|g\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}^2 \sim \|g_1\|_{\mathcal{H}^N(\mathbb{R}^6)}^2 + \|g\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}^2 \sim \|\mathcal{A}\|_{H^N(\mathbb{R}^3)}^2 + \|g_2\|_{\mathcal{H}_\ell^N(\mathbb{R}^6)}^2, \\ \mathcal{D}_{N,\ell} &= \mathcal{D}_N + \|g_2\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)}^2 \sim \|\nabla g_1\|_{\mathcal{H}^{N-1}(\mathbb{R}^3)}^2 + \|g_2\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)}^2 \sim \|\nabla_x \mathcal{A}\|_{H^{N-1}(\mathbb{R}^3)}^2 + \|g_2\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)}^2. \end{aligned}$$

Recall

$$\partial_\beta^\alpha = \partial_x^\alpha \partial_v^\beta, \quad |\alpha| + |\beta| \leq N, \quad \beta \neq 0, \quad N \geq 6,$$

and apply $\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha$ to (6.1) to get

$$\begin{aligned} & \partial_t (\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2) + v \cdot \nabla_x (\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2) + \mathcal{L}_1 (\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2) \\ &= \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha \Gamma(g, g) + [v \cdot \nabla_x, \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha] g_2 - \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha [\mathbf{P}, v \cdot \nabla] g \\ &+ [\mathcal{L}_1, \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha] g_2 - \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha \mathcal{L}_2(g_2). \end{aligned}$$

Then take the $L^2(\mathbb{R}^6)$ inner product with $\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2$ to get

$$(6.11) \quad \frac{1}{2} \frac{d}{dt} \|\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2\|_{L^2(\mathbb{R}^6)}^2 + D \leq K.$$

Here, D is a dissipation rate given by

$$D = (\mathcal{L}_1 (\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2), \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2)_{L^2(\mathbb{R}^6)}.$$

Due to the coercivity inequality from Section 2, which holds true for $\gamma > -3$, we get

$$\begin{aligned} D &\geq \eta_0 \int_{\mathbb{R}^3} \|(\mathbf{I} - \mathbf{P}) \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2\|^2 dx \\ &\geq \eta_0 \| \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2 \|_{L^2(\mathbb{R}_x^3)}^2 - C \| \partial_x^\alpha g_2 \|_{L^2_{s+\gamma/2}(\mathbb{R}^6)}^2, \end{aligned}$$

where we have used, with $\psi \in \mathcal{N}$ and $m \in \mathbb{N}$,

$$\begin{aligned} \| \mathbf{P} \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2 \|^2 &= \| \partial_x^\alpha (\psi, \tilde{W}_{\ell-|\beta|} \partial_v^\beta g_2)_{L^2(\mathbb{R}_v^3)} \psi \|^2 \\ &= |(\tilde{\psi}, \tilde{W}^{-m} \partial_x^\alpha g_2)_{L_v^2}| \| \psi \|^2 \lesssim \| \partial_x^\alpha g_2 \|_{L^2_{-m}(\mathbb{R}_{x,v}^6)}^2. \end{aligned}$$

Note that we will use the above estimate later by choosing $m = -|s + \gamma/2|$. On the other hand, K is given by

$$\begin{aligned} K &= (\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha \Gamma(g, g), \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2)_{L^2(\mathbb{R}^6)} + ([v \cdot \nabla_x, \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha] g_2, \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2)_{L^2(\mathbb{R}^6)} \\ &- (\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha [\mathbf{P}, v \cdot \nabla] g, \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2)_{L^2(\mathbb{R}^6)} + ([\mathcal{L}_1, \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha] g_2, \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2)_{L^2(\mathbb{R}^6)} \\ &- (\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha \mathcal{L}_2(g_2), \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2)_{L^2(\mathbb{R}^6)} \\ &= K_1 + K_2 + K_3 + K_4 + K_5. \end{aligned}$$

(I) Estimation of K_1 : First, we show that

$$(6.12) \quad |K_1| \lesssim \mathcal{E}_{N,\ell}^{1/2} \mathcal{D}_{N,\ell}.$$

For the proof, write

$$\begin{aligned} K_1 &= \sum_{i,j=1}^2 (\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha \Gamma(g_i, g_j), \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2)_{L^2(\mathbb{R}^6)} \\ &= K_{111} + K_{112} + K_{121} + K_{122}. \end{aligned}$$

(1) Estimation on K_{111} : Proceeding as in the computation for Ψ_1 in (6.10), we get for $\gamma > -3$,

$$|(\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha \Gamma(\varphi_k, \varphi_m), \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2)_{L^2(\mathbb{R}_v^3)}| \lesssim \| \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2 \|_{L^2(\mathbb{R}_v^3)},$$

which leads to

$$\begin{aligned} K_{111} &\sim \sum_{\alpha_1+\alpha_2=\alpha} \int_{\mathbb{R}^3} |(\partial_x^{\alpha_1} \mathcal{A})(\partial_x^{\alpha_2} \mathcal{A})| \|\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2\|_{L^2(\mathbb{R}^3)} dx \\ &\lesssim \sum_{\alpha_1+\alpha_2=\alpha} \|(\partial_x^{\alpha_1} \mathcal{A})(\partial_x^{\alpha_2} \mathcal{A})\|_{L^2(\mathbb{R}^3)} \|g_2\|_{\tilde{\mathcal{B}}_\ell^N(\mathbb{R}^6)} \\ &\lesssim \|\mathcal{A}\|_{H_x^N} (\|\nabla_x \mathcal{A}\|_{H_x^{N-1}}^2 + \|g_2\|_{\tilde{\mathcal{B}}_\ell^N(\mathbb{R}^6)}^2) \lesssim \mathcal{E}_{N,\ell}^{1/2} \mathcal{D}_{N,\ell}. \end{aligned}$$

Here, we used that for $\alpha_1 = \alpha_2 = 0$,

$$\|(\mathcal{A})^2\|_{L^2(\mathbb{R}^3)} \lesssim \|\mathcal{A}\|_{L^3(\mathbb{R}^3)} \|\mathcal{A}\|_{L^6(\mathbb{R}^3)} \lesssim \|\mathcal{A}\|_{H^1(\mathbb{R}^3)} \|\nabla \mathcal{A}\|_{L^2(\mathbb{R}^3)}.$$

(2) Estimation on K_{112} : Since $g_1 \sim \mathcal{A}\varphi$, Proposition 4.7 implies

$$\begin{aligned} |K_{112}| &\lesssim \sum_{\substack{\alpha_1+\alpha_2=\alpha \\ \beta_1+\beta_2\leq\beta}} \int_{\mathbb{R}^3} \|\tilde{W}_{\ell-|\beta|} \partial_{\beta_1}^{\alpha_1} g_1\|_{\Phi_\gamma} \|\tilde{W}_{\ell-|\beta|} \partial_{\beta_2}^{\alpha_2} g_2\|_{\Phi_\gamma} \|\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2\|_{\Phi_\gamma} dx \\ &\lesssim \sum_{\substack{\alpha_1+\alpha_2=\alpha \\ \beta_1+\beta_2\leq\beta}} \int_{\mathbb{R}^3} |\partial^{\alpha_1} \mathcal{A}| \|\tilde{W}_{\ell-|\beta|} \partial_{\beta_2}^{\alpha_2} g_2\|_{\Phi_\gamma} \|\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2\|_{\Phi_\gamma} dx = \sum_{\substack{\alpha_1+\alpha_2=\alpha \\ \beta_1+\beta_2\leq\beta}} L_{\beta_1, \beta_2}^{\alpha_1, \alpha_2}. \end{aligned}$$

We have, for $|\alpha_1| < N/2$

$$\begin{aligned} L_{\beta_1, \beta_2}^{\alpha_1, \alpha_2} &\lesssim \|\partial^{\alpha_1} \mathcal{A}\|_{L^\infty(\mathbb{R}^3)} \int_{\mathbb{R}^3} \|\tilde{W}^{\ell-|\beta_2|} \partial_{\beta_2}^{\alpha_2} g_2\|_{\Phi_\gamma} \|\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2\|_{\Phi_\gamma} dx \\ &\lesssim \|\partial^{\alpha_1} \mathcal{A}\|_{H^{3/2+\epsilon}(\mathbb{R}^3)} \|g_2\|_{\tilde{\mathcal{B}}_\ell^N(\mathbb{R}^6)}^2 \lesssim \|\mathcal{A}\|_{H^N(\mathbb{R}^3)} \|g_2\|_{\tilde{\mathcal{B}}_\ell^N(\mathbb{R}^6)}^2; \end{aligned}$$

while for $|\alpha_2| \leq N/2$

$$\begin{aligned} L_{\beta_1, \beta_2}^{\alpha_1, \alpha_2} &\lesssim \|\tilde{W}^{\ell-|\beta_2|} \partial_{\beta_2}^{\alpha_2} g_2\|_{\Phi_\gamma} \|_{L^\infty(\mathbb{R}^3)} \int_{\mathbb{R}^3} |\partial^{\alpha_1} \mathcal{A}| \|\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2\|_{\Phi_\gamma} dx \\ &\lesssim \|\tilde{W}^{\ell-|\beta_2|} \partial_{\beta_2}^{\alpha_2} g_2\|_{\Phi_\gamma} \|_{H^{3/2+\epsilon}(\mathbb{R}^3)} \|\partial^{\alpha_1} \mathcal{A}\|_{L^2(\mathbb{R}^3)} \|g_2\|_{\tilde{\mathcal{B}}_\ell^N(\mathbb{R}^6)} \\ &\lesssim \|\mathcal{A}\|_{H^N(\mathbb{R}^3)} \|g_2\|_{\tilde{\mathcal{B}}_\ell^N(\mathbb{R}^6)}^2. \end{aligned}$$

Consequently,

$$|K_{112}| \lesssim \mathcal{E}_{N,\ell}^{1/2} \mathcal{D}_{N,\ell}.$$

(3) Estimation on K_{121} : As for K_{112} , we get

$$\begin{aligned} |K_{121}| &\leq \sum_{\substack{\alpha_1+\alpha_2=\alpha \\ \beta_1+\beta_2\leq\beta}} \int_{\mathbb{R}^3} |\partial^{\alpha_2} \mathcal{A}| \|\tilde{W}_{\ell-|\beta|} \partial_{\beta_1}^{\alpha_1} g_2\|_{\Phi_\gamma} \|\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2\|_{\Phi_\gamma} dx \\ &\leq \mathcal{E}_{N,\ell}^{1/2} \mathcal{D}_{N,\ell}. \end{aligned}$$

(4) Estimation on K_{122} : We shall re-use (4.6) in the form,

$$\begin{aligned}
|K_{122}| &= |(\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha \Gamma(g_2, g_2), \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2)_{L^2(\mathbb{R}^6)}| \\
&\lesssim \sum_{\alpha_1+\alpha_2=\alpha} |(\tilde{W}_{\ell-|\beta|} \partial_\beta \Gamma(\partial^{\alpha_1} g_2, \partial^{\alpha_2} g_2), \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2)_{L^2(\mathbb{R}^6)}| \\
&\lesssim \sum_{\substack{\alpha_1+\alpha_2=\alpha \\ \beta_1+\beta_2\leq\beta}} \int_{\mathbb{R}^3} \left(\|\tilde{W}_{\ell-|\beta|} \partial_{\beta_1}^{\alpha_1} g_2\|_{L^2_{s+\gamma/2}(\mathbb{R}^3)} \|\tilde{W}_{\ell-|\beta|} \partial_{\beta_2}^{\alpha_2} g_2\|_{\Phi_\gamma} \right. \\
&\quad \left. + \|\tilde{W}_{\ell-|\beta|} \partial_{\beta_1}^{\alpha_1} g_2\|_{\Phi_\gamma} \|\tilde{W}_{\ell-|\beta|} \partial_{\beta_2}^{\alpha_2} g_2\|_{L^2_{s+\gamma/2}(\mathbb{R}^3)} \right) \|\tilde{W}_{\ell-|\beta|} \partial_\beta g_2\|_{\Phi_\gamma} dx \\
&\lesssim \sum_{\substack{\alpha_1+\alpha_2=\alpha \\ \beta_1+\beta_2\leq\beta}} \int_{\mathbb{R}^3} \left(\|\tilde{W}^{\ell-|\beta_1|} \partial_{\beta_1}^{\alpha_1} g_2\|_{L^2(\mathbb{R}^3)} \|\tilde{W}^{\ell-|\beta_2|} \partial_{\beta_2}^{\alpha_2} g_2\|_{\Phi_\gamma} \right. \\
&\quad \left. + \|\tilde{W}^{\ell-|\beta_1|} \partial_{\beta_1}^{\alpha_1} g_2\|_{\Phi_\gamma} \|\tilde{W}^{\ell-|\beta_2|} \partial_{\beta_2}^{\alpha_2} g_2\|_{L^2(\mathbb{R}^3)} \right) \|\tilde{W}_{\ell-|\beta|} \partial_\beta g_2\|_{\Phi_\gamma} dx.
\end{aligned}$$

Suppose $|\alpha_1| \leq N-2$. Then,

$$\begin{aligned}
&\int_{\mathbb{R}^3} \|\tilde{W}^{\ell-|\beta_1|} \partial_{\beta_1}^{\alpha_1} g_2\|_{L^2(\mathbb{R}^3)} \|\tilde{W}^{\ell-|\beta_2|} \partial_{\beta_2}^{\alpha_2} g_2\|_{\Phi_\gamma} \|\tilde{W}_{\ell-|\beta|} \partial_\beta g_2\|_{\Phi_\gamma} dx \\
&\lesssim \|\tilde{W}^{\ell-|\beta_1|} \partial_{\beta_1}^{\alpha_1} g_2\|_{L^\infty(\mathbb{R}^3; L^2(\mathbb{R}^3))} \int_{\mathbb{R}^3} \|\tilde{W}^{\ell-|\beta_2|} \partial_{\beta_2}^{\alpha_2} g_2\|_{\Phi_\gamma} \|\tilde{W}_{\ell-|\beta|} \partial_\beta g_2\|_{\Phi_\gamma} dx \\
&\lesssim \|g_2\|_{\tilde{\mathcal{H}}^N(\mathbb{R}^6)} \|g_2\|_{\tilde{\mathcal{B}}^N(\mathbb{R}^6)}^2.
\end{aligned}$$

The case $|\alpha_2| \leq 1$ can be computed similarly, we finally conclude

$$|K_{122}| \lesssim \|g_2\|_{\tilde{\mathcal{H}}^N(\mathbb{R}^6)} \|g_2\|_{\tilde{\mathcal{B}}^N(\mathbb{R}^6)}^2 \lesssim \mathcal{E}_{N,\ell}^{1/2} \mathcal{D}_{N,\ell},$$

and therefore the estimate (6.12) holds.

(II) Estimation of K_2 : On the other hand, we have, for $|\alpha| + |\beta| \leq N, \beta \neq 0$,

$$\begin{aligned}
|K_2| &= |([v \cdot \nabla_x, \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha] g_2, \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2)_{L^2(\mathbb{R}^6)}| \\
&\lesssim \|\tilde{W}^{\ell-(|\beta|-1)-1/2} \partial_x^{\alpha+1} \partial_v^{\beta-1} g_2\|_{L^2(\mathbb{R}^6)} \|\tilde{W}^{\ell-|\beta|-1/2} \partial_\beta^\alpha g_2\|_{L^2(\mathbb{R}^6)} \\
&\lesssim C_\delta \|\tilde{W}^{\ell-(|\beta|-1)-1/2} \partial_x^{\alpha+1} \partial_v^{\beta-1} g_2\|_{L^2(\mathbb{R}^6)}^2 + \delta \|\tilde{W}^{\ell-|\beta|-1/2} \partial_\beta^\alpha g_2\|_{L^2(\mathbb{R}^6)}^2 \\
&= C_\delta \|\tilde{W}^{\ell-(|\beta|-1)} \partial_x^{\alpha+1} \partial_v^{\beta-1} g_2\|_{L^2_{s+\gamma/2}(\mathbb{R}^6)}^2 + \delta \|\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2\|_{L^2_{s+\gamma/2}(\mathbb{R}^6)}^2 \\
&\lesssim C_\delta \|\tilde{W}^{\ell-(|\beta|-1)} \partial_x^{\alpha+1} \partial_v^{\beta-1} g_2\|_{\Phi_\gamma} \|g_2\|_{L^2(\mathbb{R}^3)}^2 + \delta \|g_2\|_{\tilde{\mathcal{B}}^N(\mathbb{R}^6)}^2.
\end{aligned}$$

(III) Estimation of K_3 : Again we assume $\beta \neq 0$ or $|\alpha| \leq N-1$.

$$\begin{aligned}
|K_3| &= |(\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha [\mathbf{P}, v \cdot \nabla] g, \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2)_{L^2(\mathbb{R}^6)}| \\
&\lesssim |(\partial_\beta \{ \tilde{W}^{2(\ell-|\beta|)} \partial_\beta^\alpha [\mathbf{P}, v \cdot \nabla] g \}, \partial^\alpha g_2)_{L^2(\mathbb{R}^6)}| \\
&\lesssim \|\nabla_x \partial^\alpha g_1\|_{L^2(\mathbb{R}^3_x)}^2 + \delta \|\tilde{W}^{-1} \nabla_x \partial^\alpha g_2\|_{L^2(\mathbb{R}^6_{x,v})}^2 + \delta \|\tilde{W}^{-1} \partial^\alpha g_2\|_{L^2(\mathbb{R}^6_{x,v})}^2 \\
&\lesssim \|\nabla_x \partial^\alpha \mathcal{A}\|_{L^2(\mathbb{R}^3_x)}^2 + \delta \|\nabla_x \partial^\alpha g_2\|_{\mathcal{B}_0^0} + \delta \|\partial^\alpha g_2\|_{\mathcal{B}_0^0} \\
&\lesssim \|\nabla_x \mathcal{A}\|_{H^{N-1}(\mathbb{R}^3_x)}^2 + \delta \|g_2\|_{\tilde{\mathcal{B}}^N(\mathbb{R}^6)}^2,
\end{aligned}$$

where \mathcal{D}_N is the dissipation integral with only x derivatives.

(IV) Estimation of K_4 : The main ingredients of the estimation are the commutator estimates I and II established in the proof of Proposition 4.8 that are valid for $\gamma > -3$. We

re-produce them here.

$$\begin{aligned}
|I| &= |([\tilde{W}_l, \mathcal{L}_1]g, \tilde{W}_l g)_{L^2(\mathbb{R}^3)}| \lesssim \|\tilde{W}_l g\|_{L^2_{\gamma/2}(\mathbb{R}^3)}^2 \\
|II| &= |([\tilde{W}_l[\partial_\beta, \mathcal{L}_1]g, \tilde{W}_l \partial_\beta g)_{L^2(\mathbb{R}^3)}| \\
&\lesssim \sum_{\beta_1+\beta_2+\beta_3=\beta, \beta_2 \neq 0} |(\tilde{W}_l \mathcal{T}(\partial_{\beta_1} \mu^{1/2}, \partial_{\beta_2} g, \partial_{\beta_3} \mu^{1/2}); \tilde{W}_l \partial_\beta g)_{L^2(\mathbb{R}^3)}| \\
&\lesssim \left(\sum_{\beta_1+\beta_2=\beta, \beta_2 \neq 0} \|\tilde{W}_l \partial_{\beta_1} g\|_{\Phi_\gamma} \right) \|\tilde{W}_l \partial_\beta g\|_{\Phi_\gamma}.
\end{aligned}$$

We also need the interpolation inequality

$$(6.13) \quad \|\tilde{W}^l \partial_\beta h\|_{L^2_{\gamma/2}} \lesssim C_\delta \|\tilde{W}^l \partial_{|\beta|-1} h\|_{L^2_{\gamma/2}} + \delta \|\tilde{W}^l \partial_\beta h\|_{H^s_{\gamma/2}} \lesssim C_\delta \|\tilde{W}^l \partial_{\beta-1} h\|_{L^2} + \delta \|\tilde{W}^l \partial_\beta h\|_{\Phi_\gamma}.$$

We shall prove

$$(6.14) \quad |K_4| \lesssim C_\delta \|g_2\|_{\tilde{\mathcal{B}}_t^{N-1}(\mathbb{R}^6)}^2 + \delta \|g_2\|_{\tilde{\mathcal{B}}_t^N(\mathbb{R}^6)}^2.$$

To this end, first, notice that

$$\begin{aligned}
K_4 &= ([\tilde{W}_{\ell-|\beta|} \partial_\beta, \mathcal{L}_1] \partial^\alpha g_2, \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2)_{L^2(\mathbb{R}^6)} \\
&= ([\tilde{W}_{\ell-|\beta|}, \mathcal{L}_1] \partial_\beta^\alpha g_2, \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2)_{L^2(\mathbb{R}^6)} \\
&\quad + (\tilde{W}_{\ell-|\beta|} [\partial_\beta, \mathcal{L}_1] \partial^\alpha g_2, \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2)_{L^2(\mathbb{R}^6)} = K_{41} + K_{42}.
\end{aligned}$$

Then, by virtue of the estimate for $|I|$ and the interpolation inequality (6.13),

$$\begin{aligned}
|K_{41}| &\lesssim \|\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2\|_{L^2_{\gamma/2}(\mathbb{R}^6)}^2 \\
&\lesssim C_\delta \|\tilde{W}_{\ell-|\beta|} \partial_{\beta-1}^\alpha g_2\|_{L^2(\mathbb{R}^3; L^2_{\gamma/2}(\mathbb{R}^3))}^2 + \delta \|\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2\|_{L^2(\mathbb{R}^3; H^s_{\gamma/2}(\mathbb{R}^3))}^2 \\
&\lesssim C_\delta \|\tilde{W}^{\ell-(|\beta|-1)} \partial_{\beta-1}^\alpha g_2\|_{L^2(\mathbb{R}^3; L^2_{s+\gamma/2}(\mathbb{R}^3))}^2 + \delta \|\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2\|_{\Phi_\gamma}^2 \\
&\lesssim C_\delta \|g_2\|_{\tilde{\mathcal{B}}_t^{N-1}(\mathbb{R}^6)}^2 + \delta \|g_2\|_{\tilde{\mathcal{B}}_t^N(\mathbb{R}^6)}^2.
\end{aligned}$$

On the other hand, the estimate for $|II|$ leads to

$$\begin{aligned}
|K_{42}| &\lesssim \sum_{\beta_1+\beta_2=\beta, \beta_2 \neq 0} \int_{\mathbb{R}^3} \|\tilde{W}^{\ell-|\beta|} \partial_{\beta_1}^\alpha g_2\|_{\Phi_\gamma} \|\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2\| dx \\
&\lesssim \|g_2\|_{\tilde{\mathcal{B}}_t^{N-1}(\mathbb{R}^6)} \|g_2\|_{\tilde{\mathcal{B}}_t^N(\mathbb{R}^6)} \\
&\lesssim C_\delta \|g_2\|_{\tilde{\mathcal{B}}_t^{N-1}(\mathbb{R}^6)}^2 + \delta \|g_2\|_{\tilde{\mathcal{B}}_t^N(\mathbb{R}^6)}^2.
\end{aligned}$$

This completes the proof of (6.14).

(V) Estimation of K_5 : Further, by Proposition 4.5 that holds for $\gamma > -3$, we can proceed as in the computation for K_{41} to obtain

$$\begin{aligned}
|K_5| &\lesssim \|\mu^{\min(a,b)} \tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2\|_{L^2(\mathbb{R}^6)}^2 \lesssim \|\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2\|_{L^2_{\gamma/2}(\mathbb{R}^6)}^2 \\
&\lesssim C_\delta \|\tilde{W}^{\ell-(|\beta|-1)} \partial_{\beta-1}^\alpha g_2\|_{L^2(\mathbb{R}^6)}^2 + \delta \|\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2\|_{\Phi_\gamma}^2 \\
&\lesssim C_\delta \|g_2\|_{\tilde{\mathcal{B}}_t^{N-1}(\mathbb{R}^6)}^2 + \delta \|g_2\|_{\tilde{\mathcal{B}}_t^N(\mathbb{R}^6)}^2.
\end{aligned}$$

Conclusion: Plug the above estimates into (6.11) to deduce, for $|\alpha + \beta| \leq N$, $|\beta| \geq 1$,

$$\begin{aligned}
(6.15) \quad & \frac{d}{dt} \left(\|\partial_\beta^\alpha g_2\|_{L_{\ell-|\beta|}^2(\mathbb{R}^6)} + (\partial^\alpha r, \nabla_x \partial^\alpha (a, -b, c))_{L^2(\mathbb{R}_x^3)} + (\partial^\alpha b, \nabla_x \partial^\alpha a)_{L^2(\mathbb{R}_x^3)} \right) \\
& + \frac{1}{2} \| \|\tilde{W}_{\ell-|\beta|} \partial_\beta^\alpha g_2\|_{\Phi_\gamma} \|_{L^2(\mathbb{R}_x^3)}^2 \\
& \lesssim \|\partial^\alpha g_2\|_{L_{s+\gamma/2}^2(\mathbb{R}^6)}^2 + \mathcal{E}_{N,\ell}^{1/2} \mathcal{D}_{N,\ell} \\
& + \| \|\tilde{W}^{\ell-(|\beta|-1)} \partial_x^{\alpha+1} \partial_v^{\beta-1} g_2\|_{\Phi_\gamma} \|_{L^2(\mathbb{R}_x^3)}^2 + \delta \|g_2\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)}^2 + \mathcal{D}_N \\
& + \|g_2\|_{\mathcal{B}_\ell^{N-1}(\mathbb{R}^6)}^2 + \delta \|g_2\|_{\mathcal{B}_\ell^N(\mathbb{R}^6)}^2.
\end{aligned}$$

We can then make a suitable linear combination of (6.3), (6.7), and (6.15) to deduce the following energy estimate.

Proposition 6.6. (Global Energy Estimate with Weight) *Under the assumptions of Theorem 1.5, for $N \geq 6$, $\ell \geq N$,*

$$\frac{d}{dt} \mathcal{E}_{N,\ell} + \mathcal{D}_{N,\ell} \lesssim \mathcal{E}_{N,\ell}^{1/2} \mathcal{D}_{N,\ell}$$

holds as far as g exists.

We can now conclude in a standard way that the global classical solutions exist for small initial data in the weighted space $\widetilde{\mathcal{H}}_\ell^N$, and this completes the proof of Theorem 1.5.

7. APPENDIX

Let us recall that $\Phi_\gamma = |\nu|^\gamma$. Let ϕ be a smooth, positive radial function that takes value 1 for small value and 0 for large value of $|\nu|$. Set $\Phi_c(\nu) = |\nu|^\gamma \phi(\nu)$. We shall show the following

Lemma 7.1. *Assume $\gamma > -3$. Then, for all integer k , one has*

$$|D^k \hat{\Phi}_c(\xi)| \lesssim \langle \xi \rangle^{-3-\gamma-k}, \text{ for all } \xi \in \mathbb{R}^3.$$

Proof. Since Φ_c is bounded and compactly supported, clearly, for any integer k , $|D^k \hat{\Phi}_c(\xi)| \leq 1$ so we can only consider the case when $|\xi| \gg 1$.

We first consider the case : $-3 < \gamma < 0$. We use the fact that the Fourier transform of $|\nu|^\gamma$ is (up to constant) $|\xi|^{-3-\gamma}$, see Page 243 of [34].

Let $\psi = \psi(\xi)$ a smooth positive function supported on $|\xi| \leq 1$, and is equal to 1 for on $|\xi| \leq 1/2$. Write

$$\hat{\Phi}_c(\xi) = \int_\eta \frac{1}{|\xi - \eta|^{3+\gamma}} \hat{\phi}(\eta) d\eta = J_1 + J_2,$$

where

$$J_1 = \int_\eta \frac{1}{|\xi - \eta|^{3+\gamma}} \psi(\xi - \eta) \hat{\phi}(\eta) d\eta, \quad \text{and} \quad J_2 = \int_\eta \frac{1}{|\xi - \eta|^{3+\gamma}} [1 - \psi(\xi - \eta)] \hat{\phi}(\eta) d\eta.$$

For J_1 , the support is on $|\xi - \eta| \leq 1$. This means that $|\eta| \geq |\xi| - 1 \geq c|\xi|$, for some constant c and because we have assumed $|\xi| \gg 1$. Then, we can use the decay of $\hat{\phi}$ to get, for any m positive

$$J_1 \lesssim \int_{\eta, |\xi - \eta| \leq 1} \frac{1}{|\xi - \eta|^{3+\gamma}} \langle \eta \rangle^{-m} d\eta \lesssim \langle \xi \rangle^{-m}.$$

For J_2 , the integration is over $|\xi - \eta| \geq 1/2$. So we can replace $|\xi - \eta|$ by $\langle \xi - \eta \rangle$ to get

$$J_2 = \int_{\eta, |\xi - \eta| \geq 1/2} \frac{1}{\langle \xi - \eta \rangle^{3+\gamma}} \langle \eta \rangle^{-m} d\eta.$$

Choose $m = M + 3 + \gamma$ with M large enough. Then

$$\begin{aligned} J_2 &= \int_{\eta, |\xi - \eta| \geq 1/2} \frac{1}{\langle \xi - \eta \rangle^{3+\gamma}} \langle \eta \rangle^{-M} \langle \eta \rangle^{-3-\gamma} d\eta \\ &\lesssim \langle \xi \rangle^{-3-\gamma} \int_{\eta, |\xi - \eta| \geq 1/2} \langle \eta \rangle^{-M} d\eta. \end{aligned}$$

Thus, we have shown that $|\hat{\Phi}_c(\xi)| \lesssim \langle \xi \rangle^{-3-\gamma}$, which proves the Lemma in the case when $k = 0$.

This proof works well for derivatives. For example, consider the case when $k = 1$. First note that

$$\nabla \hat{\Phi}_c(\xi) = \int_{\eta} \frac{1}{|\xi - \eta|^{3+\gamma}} \nabla \hat{\phi}(\eta) d\eta = K_1 + K_2,$$

where

$$K_1 = \int_{\eta} \frac{1}{|\xi - \eta|^{3+\gamma}} \psi(\xi - \eta) \nabla \hat{\phi}(\eta) d\eta \text{ and } K_2 = \int_{\eta} \frac{1}{|\xi - \eta|^{3+\gamma}} [1 - \psi(\xi - \eta)] \nabla \hat{\phi}(\eta) d\eta.$$

K_1 is estimated directly as for J_1 , with all the decay.

For K_2 , integration by parts gives

$$\begin{aligned} K_2 &= - \int_{\eta} \nabla \left[\frac{1}{|\xi - \eta|^{3+\gamma}} \right] [1 - \psi(\xi - \eta)] \hat{\phi}(\eta) d\eta \\ &\quad + \int_{\eta} \frac{1}{|\xi - \eta|^{3+\gamma}} \nabla \psi(\xi - \eta) \nabla \hat{\phi}(\eta) d\eta. \end{aligned}$$

Here, the first term has the good decay in $-3 - \gamma - 1$, while the second one has all the decay.

We now consider the case $\gamma \geq 0$. Of course, for $\gamma = 0$, the result is clear, because then $\hat{\Phi}_c$ is in \mathcal{S} .

For $2 > \gamma > 0$ we have

$$\begin{aligned} |v|^\gamma \varphi(|v|) &= \int \left(-\Delta_\xi e^{iv \cdot \xi} \right) \mathcal{F}_{v \rightarrow \xi} \left(|v|^{\gamma-2} \varphi(v) \right) (\xi) d\xi / (2\pi^3) \\ &= - \int e^{iv \cdot \xi} \Delta_\xi \mathcal{F}_{v \rightarrow \xi} \left(|v|^{\gamma-2} \varphi(v) \right) (\xi) d\xi / (2\pi^3), \end{aligned}$$

which gives

$$\left| \partial_\xi^\alpha \mathcal{F} \left(|v|^\gamma \varphi(|v|) \right) (\xi) \right| = \left| \partial_\xi^\alpha \Delta_\xi \mathcal{F} \left(|v|^{\gamma-2} \varphi(v) \right) (\xi) \right| \leq C_\alpha \langle \xi \rangle^{-3-\gamma-|\alpha|},$$

by using the previous negative case since $\gamma - 2 < 0$. The remaining cases are similar and this completes the proof of the lemma. \square

Acknowledgements : The research of the first author was supported in part by the Zhiyuan foundation and Shanghai Jiao Tong University. The research of the second author was supported by Grant-in-Aid for Scientific Research No.22540187, Japan Society of the Promotion of Science. The last author's research was supported by the General Research Fund of Hong Kong, CityU No.103109, and the Lou Jia Shan Scholarship programme of Wuhan University. The authors would like to thank the financial supports from City University of Hong Kong, Kyoto University, Rouen University and Wuhan University for their visits.

REFERENCES

- [1] R. Alexandre, Some solutions of the Boltzmann equation without angular cutoff, *J. Stat. Physics*, **104** (2001) 327–358.
- [2] R. Alexandre, A review of Boltzmann equation with singular kernels. *Kinetic and related models*, **2-4** (2009) 551–646.
- [3] R. Alexandre, L. Desvillettes, C. Villani and B. Wennberg, Entropy dissipation and long-range interactions, *Arch. Rational Mech. Anal.*, **152** (2000) 327-355.
- [4] R. Alexandre, M. ElSafadi, Littlewood-Paley decomposition and regularity issues in Boltzmann homogeneous equations. I. Non cutoff and Maxwell cases. *Math. Models Methods Appl. Sci.* **15** (2005) 907–920.
- [5] R. Alexandre, Y. Morimoto, S. Ukai, C.-J. Xu and T. Yang, Uncertainty principle and kinetic equations, *J. Funct. Anal.*, **255** (2008) 2013-2066.
- [6] R. Alexandre, Y. Morimoto, S. Ukai, C.-J. Xu and T. Yang, Regularizing effect and local existence for non-cutoff Boltzmann equation, *Arch. Rational Mech. Anal.*, **198** (2010), 39-123.
- [7] R. Alexandre, Y. Morimoto, S. Ukai, C.-J. Xu and T. Yang, Global existence and full regularity of the Boltzmann equation without angular cutoff, to appear in *Comm. Math. Phys.*
- [8] R. Alexandre, Y. Morimoto, S. Ukai, C.-J. Xu and T. Yang, Global well-posedness theory for the spatially inhomogeneous Boltzmann equation without angular cutoff, *C. R. Math. Acad. Sci. Paris, Ser. I*, **348** (2010) 867-871.
- [9] R. Alexandre, Y. Morimoto, S. Ukai, C.-J. Xu and T. Yang, Boltzmann equation without angular cutoff in the whole space : II, global existence for hard potential, preprint HAL, <http://hal.archives-ouvertes.fr/hal-00510633/fr/>
- [10] R. Alexandre, Y. Morimoto, S. Ukai, C.-J. Xu and T. Yang, Qualitative properties of solutions to the Boltzmann equation without angular cutoff, preprint.
- [11] R. Alexandre and C. Villani, On the Boltzmann equation for long-range interaction, *Comm. Pure Appl. Math.* **55** (2002) 30–70.
- [12] C. Cercignani, *The Boltzmann equation and its applications*, Applied mathematical sciences, **67**, Springer-Verlag, 1988.
- [13] L. Desvillettes, About the regularization properties of the non cut-off Kac equation, *Comm. Math. Phys.*, **168** (1995) 417–440.
- [14] L. Desvillettes, Regularization for the non Cutoff 2D Radially Symmetric Boltzmann Equation with a Velocity Dependant Cross Section, *Transport Theory and Statistical Physics*, vol. 25, n. 3-5 (1996), 383–394
- [15] L. Desvillettes, Regularization Properties of the 2-Dimensional Non Radially Symmetric Non Cutoff Spatially Homogeneous Boltzmann Equation for Maxwellian Molecules, *Transport Theory and Statistical Physics*, vol. 26, n. 3 (1997) 341–357.
- [16] L. Desvillettes and B. Wennberg, Smoothness of the solution of the spatially homogeneous Boltzmann equation without cutoff. *Comm. Partial Differential Equations*, **29-1-2** (2004) 133–155.
- [17] R. J. DiPerna and P. L. Lions, On the Cauchy problem for Boltzmann equations: global existence and weak stability. *Ann. Math.*, **130** (1989), 321-366.
- [18] H. Grad, *Asymptotic Theory of the Boltzmann Equation II*, Rarefied Gas Dynamics, J. A. Laurmann, Ed. **Vol. 1**, Academic Press, New York, 1963, 26–59.
- [19] Y. Guo, The Landau equation in a periodic box. *Comm. Math. Phys.*, **231** (2002) 391–434.
- [20] Y. Guo, The Boltzmann equation in the whole space. *Indiana Univ. Maths. J.*, **53-4** (2004) 1081–1094.
- [21] P.-T. Gressman, R.-M. Strain, Global strong solutions of the Boltzmann equation without angular cut-off. Preprint arXiv:0912.0888v1.pdf
- [22] P.-T. Gressman, R.-M. Strain, Global classical solutions of the Boltzmann equation with long-range interactions and soft potentials. Preprint arXiv:1002.3639v1.pdf
- [23] P.-T. Gressman, R.-M. Strain, Global classical solutions of the Boltzmann equation with long-range interactions. *Proc. Nat. Acad. Sci.*, **107** (2010) 5744–5749.
- [24] Z. H. Huo, Y. Morimoto, S. Ukai and T. Yang, Regularity of solutions for spatially homogeneous Boltzmann equation without Angular cutoff. *Kinetic and Related Models*, **1** (2008) 453-489.
- [25] P. L. Lions, Compactness in Boltzmann’s equation via Fourier integral operators and applications, I, II, and III. *J. Math. Kyoto Univ.*, **34** (1994), 391–427, 429–461, 539–584.
- [26] P. L. Lions, Regularity and compactness for Boltzmann collision operator without angular cut-off. *C. R. Acad. Sci. Paris Series I*, **326** (1998), 37–41.
- [27] T.-P. Liu and S.-H. Yu, Micro-macro decompositions and positivity of shock profiles. *Comm. Math. Phys.*, **246** (2004), no. 1, 133–179.
- [28] T.-P. Liu, T. Yang and S.-H. Yu, Energy method for Boltzmann equation. *Phys. D*, **188** (2004), 178-192.

- [29] Y. Morimoto and S. Ukai, Gevrey smoothing effect of solutions for spatially homogeneous nonlinear Boltzmann equation without angular cutoff, *J. Pseudo-Differ. Oper. Appl.* **1** (2010), 139-159.
- [30] Y. Morimoto, S. Ukai, C.-J. Xu and T. Yang, Regularity of solutions to the spatially homogeneous Boltzmann equation without angular cutoff. *Discrete and Continuous Dynamical Systems - Series A* **24** 1 (2009), 187-212.
- [31] C. Mouhot, Explicit coercivity estimates for the linearized Boltzmann and Landau operators. *Comm. Part. Diff. Equations*, vol. 31, (2006) 1321-1348.
- [32] C. Mouhot and R. M. Strain, Spectral gap and coercivity estimates for linearized Boltzmann collision operators without angular cutoff, *J. Math. Pures Appl.* (9) **87** (2007), no. 5, 515–535.
- [33] Y. P. Pao, Boltzmann collision operator with inverse power intermolecular potential, I, II. *Commun. Pure Appl. Math.*, **27** (1974), 407–428, 559–581.
- [34] M.-E. Taylor, *Partial Differential Equations*, **Vol. I**, Springer, New York .
- [35] S. Ukai, On the existence of global solutions of mixed problem for non-linear Boltzmann equation. *Proc. Japan Acad.* , **50** (1974), 179–184.
- [36] S. Ukai, Les solutions globales de l'équation de Boltzmann dans l'espace tout entier et dans le demi-espace, *C. R. Acad. Sci. Paris Ser. A-B* **282** (1976), no. 6, 317–A320.
- [37] S. Ukai, Local solutions in Gevrey classes to the nonlinear Boltzmann equation without cutoff, *Japan J. Appl. Math.*, **1-1** (1984) 141–156.
- [38] S. Ukai, Solutions of the Boltzmann equation, *Pattern and Waves – Qualitative Analysis of Nonlinear Differential Equations* (eds. M.Mimura and T.Nishida), Studies of Mathematics and Its Applications **18**, pp37-96, Kinokuniya-North-Holland, Tokyo, 1986.
- [39] C. Villani, Regularity estimates via entropy dissipation for the spatially homogeneous Boltzmann equation, *Rev. Mat. Iberoamericana*, **15-2** (1999) 335–352.
- [40] C. Villani, *A review of mathematical topics in collisional kinetic theory*. Handbook of Fluid Mechanics. Ed. S. Friedlander, D.Serre, 2002.

R. ALEXANDRE, DEPARTMENT OF MATHEMATICS, SHANGHAI JIAO TONG UNIVERSITY
SHANGHAI, 200240, P. R. CHINA

AND

IRENAV RESEARCH INSTITUTE, FRENCH NAVAL ACADEMY BREST-LANVÉOC 29290, FRANCE
E-mail address: radjesvarane.alexandre@ecole-navale.fr

Y. MORIMOTO, GRADUATE SCHOOL OF HUMAN AND ENVIRONMENTAL STUDIES, KYOTO UNIVERSITY
KYOTO, 606-8501, JAPAN

E-mail address: morimoto@math.h.kyoto-u.ac.jp

S. UKAI, 17-26 IWASAKI-CHO, HODOGAYA-KU, YOKOHAMA 240-0015, JAPAN

E-mail address: ukai@kurims.kyoto-u.ac.jp

C.-J. XU, SCHOOL OF MATHEMATICS, WUHAN UNIVERSITY 430072, WUHAN, P. R. CHINA

AND

UNIVERSITÉ DE ROUEN, UMR 6085-CNRS, MATHÉMATIQUES
AVENUE DE L'UNIVERSITÉ, BP.12, 76801 SAINT ÉTIENNE DU ROUVRAY, FRANCE

E-mail address: Chao-Jiang.Xu@univ-rouen.fr

T. YANG, DEPARTMENT OF MATHEMATICS, CITY UNIVERSITY OF HONG KONG, HONG KONG, P. R. CHINA

AND

SCHOOL OF MATHEMATICS, WUHAN UNIVERSITY 430072, WUHAN, P. R. CHINA

E-mail address: matyang@cityu.edu.hk