## HAL

open science

# Global exact controllability in infinite time of Schrödinger equation 

Vahagn Nersesyan, Hayk Nersisyan

## To cite this version:

Vahagn Nersesyan, Hayk Nersisyan. Global exact controllability in infinite time of Schrödinger equation. 2010. hal-00489481v1

HAL Id: hal-00489481
https://hal.science/hal-00489481v1
Preprint submitted on 4 Jun 2010 (v1), last revised 5 Oct 2011 (v2)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire HAL, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# Global exact controllability in infinite time of Schrödinger equation 

Vahagn Nersesyan, Hayk Nersisyan ${ }^{\dagger}$


#### Abstract

In this paper, we study the problem of controllability of Schrödinger equation. We prove that the system is exactly controllable in infinite time to any position. The proof is based on an inverse mapping theorem for multivalued functions. We show also that the system is not exactly controllable in finite time in lower Sobolev spaces.

\section*{Contents} 1 IntroductionB 2 Controllability of linearized system ..... 42.1 Main result2.2 Proof of Theorem 2.6.9 2.3 Proof of Proposition 2.10 ..... 12 3 Controllability of nonlinear system ..... 14 3.1 Well-posedness of Schrödinger equation ..... 14 3.2 Exact controllability in infinite time ..... 17 3.3 Proof of Theorem 3.5 ..... 18 4 Non-controllability result ..... 20 4.1 Main result ..... 20 4.2 Some $\varepsilon$-entropy estimates ..... 21 4.3 Proof of Theorem 4.2 ..... 23 5 Appendix ..... 24 5.1 Genericity of Condition 2.5 ..... 24 5.2 Inverse mapping theorem for multifunctions ..... 26 References ..... 26

^[ *Laboratoire de Mathématiques, UMR CNRS 8100, Université de Versailles-Saint-Quentin-en-Yvelines, F-78035 Versailles, France, e-mail: Vahagn.Nersesyan@math.uvsq.fr ${ }^{\dagger}$ Laboratoire de Mathématiques, UMR CNRS 8088, Université de Cergy-Pontoise, F-95000 Cergy-Pontoise, France, e-mail: Hayk.Nersisyan@u-cergy.fr ]


## 1 Introduction

The paper is devoted to the study of the following controlled Schrödinger equation

$$
\begin{align*}
i \dot{z} & =-\Delta z+V(x) z+u(t) Q(x) z,  \tag{1.1}\\
\left.z\right|_{\partial D} & =0,  \tag{1.2}\\
z(0, x) & =z_{0}(x) . \tag{1.3}
\end{align*}
$$

We assume that space variable $x$ belongs to a rectangle $D \subset \mathbb{R}^{d}, V, Q \in$ $C^{\infty}(\bar{D}, \mathbb{R})$ is a given functions, $u$ is the control, and $z$ is the state. We prove that the linearization of this system is exactly controllable in Sobolev spaces in infinite time. Application of this result gives global exact controllability in infinite time in $H^{3}$ for $d=1$. We show also that the system is not exactly controllable in finite time in lower Sobolev spaces.

Let us recall some previous results on the controllability problem of Schrödinger equation. In [6], Beauchard proves an exact controllability result for the system with $d=1, D=(-1,1)$ and $Q(x)=x$ in $H^{7}$-neighborhoods of the eigefunctions. Beauchard and Coron [8] established later a partial global exact controllability result, showing that the system in question is also controlled between neighborhoods of eigenfunctions. Recently, Beauchard and Laurent [10] simplified the proof of [6] and generalized it to the case of the nonlinear equation. The proofs of [6, 8, 10] work also for the neighborhoods of finite linear combinations of eigenfunctions. In the case of infinite linear combinations, these arguments do not work, since the linearized system does not verify the property of spectral gap (even if the problem is 1-D), hence the Ingham inequality cannot be applied.

Chambrion et al. 12], Privat, Sigalotti 26], and Mason, Sigalotti 19] prove that (1.1), (1.2) is approximately controllable in $L^{2}$ generically with respect to function $Q$ and domain $D$. In [23, 22], the first author of this paper proves a stabilization result and a property of global approximate controllability to eigenstates for Schrödinger equation. Combination of these results with the local exact controllability property obtained by Beauchard [6] gives global exact controllability in finite time for $d=1$ in the spaces $H^{3+\varepsilon}, \varepsilon>0$. See also papers [27. 30, 3, 2, 1, 9] for controllability of finite-dimensional systems and papers [17. 18, 5, 31, 13, 20, 15] for controllability properties of various Schrödinger systems.

In this article, we study the properties of control system on the time half-line $\mathbb{R}_{+}$instead of a finite interval $[0, T]$, as in all above cited papers. We study the mapping, which associates to initial condition $z_{0}$ and control $u$ the $\omega$-limit set of the corresponding trajectory. We consider this mapping as a multivalued function in the phase space. We show that this multivalued function is differentiable with differential equal to the limit of the linearization of (1.1), (1.2), when time $t$ goes to infinity. Observing that the linearized system is controllable in infinite time at almost any point, we conclude the controllability of the nonlinear system (in the case $d=1$ ), using an inverse mapping theorem for multivalued functions
[21] by Nachi and Penot. Thus (1.1), (1.2) is exactly controllable near any point in the phase space, hence globally. The controllability of the linearized system is proved for any $d \geq 1$, but this result is not directly applicable to the study of the nonlinear system with $d \geq 2$. We have a loss of regularity: the solution of the nonlinear problem exists for more regular controls than the ones used to control the linear problem. A Nash-Moser type theorem for multivalued functions could be useful to solve this problem. This is a subject for further study.

To our knowledge, the inverse mapping theorem for multivalued functions was never used before in the theory of control of PDEs. Our proof does not rely on the particular asymptotics of the eigenvalues of Dirichlet Laplacian, so it is likely to work in other settings. Considering the problem in infinite time enables us to prove the controllability of the linearized system in the case of any space dimension $d \geq 1$, even when the gap condition is not verified for the eigenvalues (which is the case for $d \geq 3$ ).

In the second part of the paper, we study the problem of non-controllability for (1.1), (1.2) in finite time. We prove that the system is not exactly controllable in finite time in the spaces $H^{k}$ with $k \in(0, d)$. Let us recall that previously Ball, Marsden and Slemrod [4] and Turinici [29] have shown that the problem is not controllable in the space $H^{2}$. Our result is inspired by the paper 28] of Shirikyan, where the non-controllability of 2D Euler equation is established. More precisely, it is proved in [28] that, if the Euler system is controlled by finite dimensional external force, then the set of all reachable points in a given time $T>0$ cannot cover a ball in the phase space. Later this result was generalized by the second author of the present paper, in [24]: in the case of 3D Euler equation it is proved that the union of all sets of reachable points at all times $T>0$ also does not cover a ball.

Using ideas of Shirikyan, we prove that the image by the resolving operator of a ball in the space of controls has a Kolmogorov $\varepsilon$-entropy strictly less than that of a ball in the phase space $H^{k}$. This implies the non-controllability.

Acknowledgments. The authors would like to thank Armen Shirikyan for many fruitful conversations.

## Notation

In this paper, we use the following notation. Let

$$
\ell^{2}:=\left\{\left\{a_{j}\right\} \in \mathbb{C}^{\infty}:\left\|\left\{a_{j}\right\}\right\|_{\ell^{2}}^{2}=\sum_{j=1}^{+\infty}\left|a_{j}\right|^{2}<+\infty\right\}
$$

We denote by $H^{s}:=H^{s}(D)$ be the Sobolev space of order $s \geq 0$. Consider Schrödinger operator $-\Delta+V, V \in C^{\infty}(\bar{D}, \mathbb{R})$ with $\mathcal{D}(-\Delta+V):=H_{0}^{1} \cap H^{2}$. Let $\left\{\lambda_{j, V}\right\}$ and $\left\{e_{j, V}\right\}$ be the sets of eigenvalues and normalized eigenfunctions of this operator. Let $\langle\cdot, \cdot\rangle$ and $\|\cdot\|$ be the scalar product and the norm in the space $L^{2}$. Define the space $H_{(V)}^{s}:=D\left((-\Delta+V)^{\frac{s}{2}}\right)$ endowed with the
norm $\|\cdot\|_{s, V}=\left\|\left(\lambda_{j, V}\right)^{\frac{s}{2}}\left\langle\cdot, e_{j, V}\right\rangle\right\|_{\ell^{2}}$. When $D$ is the rectangle $(0,1)^{d}$ and $V\left(x_{1}, \ldots, x_{d}\right)=V_{1}\left(x_{1}\right)+\ldots+V_{d}\left(x_{d}\right), V_{k} \in C^{\infty}([0,1], \mathbb{R})$, the eigenvalues and eigenfunctions of $-\Delta+V$ on $D$ are of the form

$$
\begin{align*}
\lambda_{j_{1}, \ldots, j_{d}, V} & =\lambda_{j_{1}, V_{1}}+\ldots+\lambda_{j_{d}, V_{d}},  \tag{1.4}\\
e_{j_{1}, \ldots, j_{d}, V}\left(x_{1}, \ldots, x_{d}\right) & =e_{j_{1}, V_{1}}\left(x_{1}\right) \cdot \ldots \cdot e_{j_{d}, V_{d}}\left(x_{d}\right), \quad\left(x_{1}, \ldots, x_{d}\right) \in D, \tag{1.5}
\end{align*}
$$

where $\left\{\lambda_{j, V_{k}}\right\}$ and $\left\{e_{j, V_{k}}\right\}$ are the eigenvalues and eigenfunctions of operator $-\frac{\mathrm{d}^{2}}{\mathrm{~d} x^{2}}+V$ on $(0,1)$. Define the space

$$
\mathcal{V}:=\left\{z \in L^{2}:\|z\|_{\mathcal{V}}^{2}:=\sum_{j_{1}, \ldots, j_{d}=1}^{+\infty} \mid\left(\left.j_{1}^{3} \cdot \ldots \cdot j_{d}^{3}\left\langle z, e_{j_{1}, \ldots, j_{d}, V}\right\rangle\right|^{2}<+\infty\right\} .\right.
$$

Notice that, in the case $d=1$, the space $\mathcal{V}$ coincides with $H_{(V)}^{3}$. The eigenvalues and eigenfunctions of Dirichlet Laplacian on the interval $(0,1)$ are $\lambda_{k, 0}=k^{2} \pi^{2}$ and $e_{k, 0}(x)=\sqrt{2} \sin (k \pi x), x \in(0,1)$. It is well known that for any $V \in$ $L^{2}([0,1], \mathbb{R})$

$$
\begin{align*}
& \lambda_{k, V}=k^{2} \pi^{2}+\int_{0}^{1} V(x) \mathrm{d} x+r_{k}  \tag{1.6}\\
& \left\|e_{k, V}-e_{k, 0}\right\|_{L^{\infty}} \leq \frac{C}{k}  \tag{1.7}\\
& \left\|\frac{\mathrm{~d} e_{k, V}}{\mathrm{~d} x}-\frac{\mathrm{d} e_{k, 0}}{\mathrm{~d} x}\right\|_{L^{\infty}} \leq C \tag{1.8}
\end{align*}
$$

where $\sum_{k=1}^{+\infty} r_{k}^{2}<+\infty$ (e.g., see 25). For a Banach space $X$, we shall denote by $B_{X}(a, r)$ the open ball of radius $r>0$ centered at $a \in X$. For a set $A$, we write $2^{A}$ for the set consisting of all subsets of $A$. We denote by $C$ a constant whose value may change from line to line.

## 2 Controllability of linearized system

### 2.1 Main result

In this section, we suppose that $D$ is the rectangle $(0,1)^{d}, d \geq 1$ and $V\left(x_{1}, \ldots, x_{d}\right)$ $=V_{1}\left(x_{1}\right)+\ldots+V_{d}\left(x_{d}\right), V_{k} \in C^{\infty}([0,1], \mathbb{R})$. For any $\tilde{z} \in \mathcal{V}$, let $\mathcal{U}_{t}(\tilde{z}, 0)$ be the solution of (1.1)-(1.3) with $z_{0}=\tilde{z}$ and $u=0$. Clearly,

$$
\begin{equation*}
\mathcal{U}_{t}(\tilde{z}, 0)=\sum_{j=1}^{+\infty} e^{-i \lambda_{j, V} t}\left\langle\tilde{z}, e_{j, V}\right\rangle e_{j, V} \tag{2.1}
\end{equation*}
$$

Lemma 2.1. For any $\tilde{z} \in \mathcal{V}$, there is a sequence $T_{n} \rightarrow+\infty$ such that $\mathcal{U}_{T_{n}}(\tilde{z}, 0) \rightarrow \tilde{z}$ in $\mathcal{V}$.

Proof. The proof uses the following well known result (e.g., see 16]).

Lemma 2.2. For any $\varepsilon>0, N \geq 1$ and $\alpha_{j} \in \mathbb{R}, j=1, \ldots, N$, there is $k \in \mathbb{N}$ such that

$$
\sum_{j=1}^{N}\left|e^{i \alpha_{j} k}-1\right|<\varepsilon
$$

Applying this lemma, we see that for any $\varepsilon>0$ and for sufficiently large $N \geq 1$, we have

$$
\begin{aligned}
\left\|\mathcal{U}_{k}(\tilde{z}, 0)-\tilde{z}\right\|_{\mathcal{V}}^{2} \leq & \sum_{\lambda_{j_{1}, \ldots, j_{d}, V \leq N}}\left|e^{-i \lambda_{j_{1}}, \ldots, j_{d}, v k}-1\right|^{2} \mid j_{1}^{3} \cdot \ldots \cdot j_{d}^{3}\left\langle\tilde{z}, e_{\left.j_{1}, \ldots, j_{d}, V\right\rangle\left.\right|^{2}}\right. \\
& +2 \sum_{\left.\lambda_{j_{1}, \ldots, j_{d}, V}\right\rangle N}\left|j_{1}^{3} \cdot \ldots \cdot j_{d}^{3}\left\langle\tilde{z}, e_{\left.j_{1}, \ldots, j_{d}, V\right\rangle}\right\rangle\right|^{2} \leq \frac{\varepsilon}{2}+\frac{\varepsilon}{2}=\varepsilon
\end{aligned}
$$

for an appropriate choice of $k \in \mathbb{N}$. This proves Lemma 2.1.

This subsection is devoted to the study of the linearization of (1.1), (1.2) around the trajectory $\mathcal{U}_{t}(\tilde{z}, 0)$ :

$$
\begin{align*}
i \dot{z} & =-\Delta z+V(x) z+u(t) Q(x) \mathcal{U}_{t}(\tilde{z}, 0),  \tag{2.2}\\
\left.z\right|_{\partial D} & =0,  \tag{2.3}\\
z(0, x) & =z_{0} . \tag{2.4}
\end{align*}
$$

Let $S$ be the unit sphere in $L^{2}$. For $y \in S$, let $T_{y}$ be the tangent space to $S$ at $y \in S$ :

$$
T_{y}=\{z \in S: \operatorname{Re}\langle z, y\rangle=0\} .
$$

Lemma 2.3. For any $z_{0} \in T_{\tilde{z}} \cap H_{(0)}^{2}$ and $u \in L_{\text {loc }}^{1}\left(\mathbb{R}_{+}, \mathbb{R}\right)$, problem (2.2)-(2.4) has a unique solution $z \in C\left(\mathbb{R}_{+}, H_{(0)}^{2}\right)$. Furthermore, if $\mathcal{R}_{t}(\cdot, \cdot): T_{\tilde{z}} \cap H_{(0)}^{2} \times$ $L^{1}([0, t], \mathbb{R}) \rightarrow H_{(0)}^{2},\left(z_{0}, u\right) \rightarrow z(t)$ is the resolving operator of the problem, then
(i) $R_{t}\left(z_{0}, u\right) \in T_{\mathcal{U}_{t}(\tilde{z}, 0)}$ for any $t \geq 0$,
(ii) The operator $\mathcal{R}_{t}$ is linear continuous from $T_{\tilde{z}} \cap H_{(0)}^{2} \times L^{1}([0, t], \mathbb{R})$ to $H_{(0)}^{2}$.

Proof. The proof of existence and (ii) is standard (e.g., see 11]). To prove (i), notice that

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} t} \operatorname{Re}\left\langle R_{t}, \mathcal{U}_{t}\right\rangle & =\operatorname{Re}\left\langle\dot{R}_{t}, \mathcal{U}_{t}\right\rangle+\operatorname{Re}\left\langle R_{t}, \dot{\mathcal{U}}_{t}\right\rangle \\
& =\operatorname{Re}\left\langle i(\Delta-V) R_{t}-i u(t) Q(x) \mathcal{U}_{t}, \mathcal{U}_{t}\right\rangle+\operatorname{Re}\left\langle R_{t}, i(\Delta-V) \mathcal{U}_{t}\right\rangle \\
& =\operatorname{Re}\left\langle i(\Delta-V) R_{t}, \mathcal{U}_{t}\right\rangle+\operatorname{Re}\left\langle R_{t}, i(\Delta-V) \mathcal{U}_{t}\right\rangle=0
\end{aligned}
$$

Since $\operatorname{Re}\left\langle R_{0}, \mathcal{U}_{0}\right\rangle=\operatorname{Re}\left\langle z_{0}, \tilde{z}\right\rangle=0$, we get (i).

As (2.2)-(2.4) is a linear control problem, the controllability of system with $z_{0}=0$ is equivalent to that of with any $z_{0} \in T_{\tilde{z}}$. Henceforth, we take $z_{0}=0$ in (2.4). Let us rewrite this problem in the Duhamel form

$$
\begin{equation*}
z(t)=-i \int_{0}^{t} S(t-s) u(s) Q(x) \mathcal{U}_{s}(\tilde{z}, 0) \mathrm{d} s \tag{2.5}
\end{equation*}
$$

where $S(t)=e^{i t(\Delta-V)}$ is the free evolution. Using (2.1) and (2.5), we obtain

$$
\begin{equation*}
\left\langle z(t), e_{m, V}\right\rangle=-i \sum_{k=1}^{+\infty} e^{-i \lambda_{m, V} t}\left\langle\tilde{z}, e_{k, V}\right\rangle Q_{m k} \int_{0}^{t} e^{i \omega_{m k} s} u(s) \mathrm{d} s, m \geq 1 \tag{2.6}
\end{equation*}
$$

where $\omega_{m k}=\lambda_{m}-\lambda_{k}$ and $Q_{m k}:=\left\langle Q e_{m, V}, e_{k, V}\right\rangle$. Let $T_{n} \rightarrow+\infty$ be the sequence in Lemma 2.2. Then $e^{-i \lambda_{m, V} T_{n}} \rightarrow 1$ as $n \rightarrow+\infty$. Let us take $t=T_{n}$ in (2.6) and pass to the limit as $n \rightarrow+\infty$. For any $u \in L^{1}\left(\mathbb{R}_{+}, \mathbb{R}\right)$ the right-hand side has a limit. Equality (2.6) implies that the following limit exists

$$
R_{\infty}(0, u):=\lim _{n \rightarrow+\infty} z\left(T_{n}\right)=\lim _{n \rightarrow+\infty} \mathcal{R}_{T_{n}}(0, u)
$$

The choice of the sequence $T_{n}$ implies that

$$
\begin{equation*}
\left\langle R_{\infty}(0, u), e_{m, V}\right\rangle=-i \sum_{k=1}^{+\infty}\left\langle\tilde{z}, e_{k, V}\right\rangle Q_{m k} \int_{0}^{+\infty} e^{i \omega_{m k} s} u(s) \mathrm{d} s \tag{2.7}
\end{equation*}
$$

Moreover, $\mathcal{R}_{\infty}(0, u) \in T_{\tilde{z}}$. Indeed,

$$
\operatorname{Re}\left\langle\mathcal{R}_{\infty}(0, u), \tilde{z}\right\rangle=\lim _{n \rightarrow \infty} \operatorname{Re}\left\langle\mathcal{R}_{T_{n}}(0, u), \mathcal{U}_{T_{n}}(\tilde{z}, 0)\right\rangle=0
$$

by property (i).
For any $u \in L^{1}\left(\mathbb{R}_{+}, \mathbb{R}\right)$, denote by $\check{u}$ the inverse Fourier transform of the function obtained by extending $u$ as zero to $\mathbb{R}_{-}^{*}$ :

$$
\check{u}(\omega):=\int_{0}^{+\infty} e^{i \omega s} u(s) \mathrm{d} s .
$$

Define the following spaces

$$
\begin{aligned}
& \tilde{\ell}^{2}:=\left\{d=\left\{d_{m k}\right\}:\|d\|_{\tilde{\ell}^{2}}^{2}:=\left|d_{11}\right|^{2}+\sum_{m, k=1, m \neq k}^{+\infty}\left|d_{m k}\right|^{2}<+\infty, d_{m m}=d_{11}\right. \\
&\text { and } \left.d_{m k}=\bar{d}_{k m} \text { for all } m, k \geq 1\right\}, \\
& \mathcal{B}:=\left\{u \in L_{l o c}^{2}\left(\mathbb{R}_{+}, \mathbb{R}\right):\|u\|_{\mathcal{B}}^{2}:=\sum_{p=1}^{+\infty} p^{2}\|u\|_{L^{2}([p-1, p])}^{2}<+\infty\right\} . \\
& \mathcal{C}:=\left\{u \in L^{1}\left(\mathbb{R}_{+}, \mathbb{R}\right):\left\{\check{u}\left(\omega_{m k}\right)\right\} \in \tilde{\ell}^{2}\right\} .
\end{aligned}
$$

The set of admissible controls is the Banach space

$$
\Theta:=u \in \mathcal{B} \cap \mathcal{C} \cap H^{s}\left(\mathbb{R}_{+}, \mathbb{R}\right)
$$

endowed with the norm $\|u\|_{\Theta}:=\|u\|_{\mathcal{B}}+\|u\|_{L^{1}}+\left\|\left\{\check{u}\left(\omega_{m k}\right)\right\}\right\|_{\tilde{\ell}^{2}}+\|u\|_{H^{s}}$, where $s \geq 1$ is any fixed constant. Clearly, the space $\Theta$ is nontrivial.
Lemma 2.4. For any $\tilde{z} \in S \cap \mathcal{V}, R_{\infty}(0, \cdot)$ is linear continuous mapping from $\Theta$ to $T_{\tilde{z}} \cap \mathcal{V}$.

Proof. Step 1. Let us admit that for any $m_{j}, k_{j} \geq 1, j=1, \ldots, d$ we have

$$
\begin{equation*}
\left|\frac{\left(m_{1} \cdot \ldots \cdot m_{d}\right)^{3}}{\left(k_{1} \cdot \ldots \cdot k_{d}\right)^{3}}\left\langle Q e_{k_{1}, \ldots, k_{d}, V}, e_{m_{1}, \ldots, m_{d}, V}\right\rangle\right| \leq C \tag{2.8}
\end{equation*}
$$

Then (2.7), (2.8) and the Schwarz inequality imply that

$$
\begin{aligned}
& \left\|R_{\infty}(0, u)\right\|_{\mathcal{V}}^{2}=\sum_{m_{1}, \ldots, m_{d}=1}^{+\infty} \mid\left(\left.m_{1}^{3} \cdot \ldots \cdot m_{d}^{3}\left\langle R_{\infty}(0, u), e_{m_{1}, \ldots, m_{d}, V}\right\rangle\right|^{2}\right. \\
& \leq\|\tilde{z}\|_{\mathcal{V}}^{2} \sum_{k_{1}, m_{1}, \ldots, k_{d}, m_{d}=1}^{+\infty}\left|\frac{\left(m_{1} \cdot \ldots \cdot m_{d}\right)^{3}}{\left(k_{1} \cdot \ldots \cdot k_{d}\right)^{3}}\left\langle Q e_{k_{1}, \ldots, k_{d}, V}, e_{m_{1}, \ldots, m_{d}, V}\right\rangle \int_{0}^{+\infty} e^{i \omega_{m k} s} u(s) \mathrm{d} s\right|^{2} \\
& \leq C^{2}\|\tilde{z}\|_{\mathcal{V}}^{2}\|u\|_{\Theta}^{2}<+\infty .
\end{aligned}
$$

Step 2. Let us prove (2.8). To simplify notation, let us suppose that $d=2$; the proof of the general case is similar. Let $V\left(x_{1}, x_{2}\right)=V_{1}\left(x_{1}\right)+V_{2}\left(x_{2}\right)$. Integration by parts gives

$$
\begin{aligned}
\left\langle Q e_{k_{1}, k_{2}, V}, e_{m_{1}, m_{2}, V}\right\rangle= & \frac{1}{\lambda_{m_{1}, V_{1}}^{2}}\left\langle\left(-\frac{\partial^{2}}{\partial x_{1}^{2}}+V_{1}\right)\left(Q e_{k_{1}, k_{2}, V}\right),\left(-\frac{\partial^{2}}{\partial x_{1}^{2}}+V_{1}\right)\left(e_{m_{1}, m_{2}, V}\right)\right\rangle \\
= & \frac{1}{\lambda_{m_{1}, V_{1}}^{2}}\left(\left.\int_{0}^{1} \frac{\partial Q}{\partial x_{1}} \frac{\partial e_{k_{1}, k_{2}, V}}{\partial x_{1}} \frac{\partial e_{m_{1}, m_{2}, V}}{\partial x_{1}}\right|_{x_{1}=0} ^{x_{1}=1} \mathrm{~d} x_{2}\right. \\
& +\left\langle\frac{\partial}{\partial x_{1}}\left(-\frac{\partial^{2}}{\partial x_{1}^{2}}+V_{1}\right)\left(Q e_{k_{1}, k_{2}, V}\right), \frac{\partial e_{m_{1}, m_{2}, V}}{\partial x_{1}}\right\rangle \\
& \left.+\left\langle\left(-\frac{\partial^{2}}{\partial x_{1}^{2}}+V_{1}\right)\left(Q e_{k_{1}, k_{2}, V}\right), V_{1} e_{m_{1}, m_{2}, V}\right\rangle\right) \\
= & I_{1}+I_{2}+I_{3} .
\end{aligned}
$$

Again integrating by parts, we get

$$
\begin{aligned}
I_{1}= & \left.\frac{1}{\lambda_{m_{1}, V_{1}}^{2} \lambda_{m_{2}, V_{2}}^{2}} \int_{0}^{1} \frac{\partial Q}{\partial x_{1}} \frac{\partial e_{k_{1}, k_{2}, V}}{\partial x_{1}}\left(-\frac{\partial^{2}}{\partial x_{2}^{2}}+V_{2}\right) \frac{\partial e_{m_{1}, m_{2}, V}}{\partial x_{1}}\right|_{x_{1}=0} ^{x_{1}=1} \mathrm{~d} x_{2} \\
= & \frac{1}{\lambda_{m_{1}, V_{1}}^{2} \lambda_{m_{2}, V_{2}}^{2}}\left(\frac{\partial Q}{\partial x_{1}} \frac{\partial e_{k_{1}, k_{2}, V}}{\partial x_{1}} \frac{\partial^{2} e_{m_{1}, m_{2}, V}^{\partial x_{1}} \partial x_{2}}{\left.\left.\right|_{x_{1}=0} ^{x_{1}=1}\right|_{x_{2}=0} ^{x_{2}=1}}\right. \\
& +\int_{0}^{1}\left(\frac{\partial^{2} Q}{\partial x_{1} \partial x_{2}} \frac{\left.\partial e_{k_{1}, k_{2}, V}^{\partial x_{1}}+\frac{\partial Q}{\partial x_{1}} \frac{\partial^{2} e_{k_{1}, k_{2}, V}}{\partial x_{1} \partial x_{2}}\right)\left.\frac{\partial^{2} e_{m_{1}, m_{2}, V}}{\partial x_{1} \partial x_{2}}\right|_{x_{1}=0} ^{x_{1}=1} \mathrm{~d} x_{2}}{}\right. \\
& +\left.\int_{0}^{1} \frac{\partial Q}{\partial x_{1}} \frac{\partial e_{k_{1}, k_{2}, V}^{\partial x_{1}}}{V_{2}} \frac{\partial e_{m_{1}, m_{2}, V}}{\partial x_{1}}\right|_{x_{1}=0} ^{x_{1}=1} \mathrm{~d} x_{2} .
\end{aligned}
$$

In view of (1.4)-(1.8), this implies that

$$
\left|\frac{\left(m_{1} \cdot \ldots \cdot m_{d}\right)^{3}}{\left(k_{1} \cdot \ldots \cdot k_{d}\right)^{3}} I_{1}\right| \leq C .
$$

The terms $I_{2}, I_{3}$ are treated in the same way. We omit the details.

We prove the controllability of (2.2), (2.3) under below condition.
Condition 2.5. The functions $V, Q \in C^{\infty}(\bar{D}, \mathbb{R})$ are such that
(i) $\inf _{p_{1}, j_{1}, \ldots, p_{d}, j_{d} \geq 1}\left|\left(p_{1} j_{1} \cdot \ldots \cdot p_{d} j_{d}\right)^{3} Q_{p j}\right|>0, Q_{p j}:=\left\langle Q e_{p_{1}, \ldots, p_{d}, V}, e_{j_{1}, \ldots, j_{d}, V}\right\rangle$,
(ii) $\lambda_{i, V}-\lambda_{j, V} \neq \lambda_{p, V}-\lambda_{q, V}$ for all $i, j, p, q \geq 1$ such that $\{i, j\} \neq\{p, q\}$ and $i \neq j$.

See Appendix for the proof of genericity of this condition. Let us introduce the set

$$
\begin{aligned}
& \mathcal{E}:=\left\{z \in S: \exists p, q \geq 1, p \neq q, z=c_{p} e_{p, V}+c_{q} e_{q, V}\right. \\
& \left.\qquad\left|c_{p}\right|^{2}\left\langle Q e_{p, V}, e_{p, V}\right\rangle-\left|c_{q}\right|^{2}\left\langle Q e_{q, V}, e_{q, V}\right\rangle=0\right\} .
\end{aligned}
$$

The following result is proved in next subsection.
Theorem 2.6. Under Condition 2.5, for any $\tilde{z} \in S \cap \mathcal{V} \backslash \mathcal{E}$, the mapping $R_{\infty}(0, \cdot): \Theta \rightarrow T_{\tilde{z}} \cap \mathcal{V}$ admits a continuous right inverse, where the space $T_{\tilde{z}} \cap \mathcal{V}$ is endowed with the norm of $\mathcal{V}$. If $\tilde{z} \in S \cap \mathcal{V} \cap \mathcal{E}$, then $R_{\infty}(0, \cdot)$ is not invertible.
Remark 2.7. The invertibility of the mapping $R_{T}(0, \cdot)$ with finite $T>0$ and $\tilde{z}=e_{1}$ is studied by Beauchard et al. [7]. They prove that for space dimension $d \geq 3$ the mapping is not invertible. By Beauchard $[6], R_{T}$ is invertible in the case $d=1$ and $\tilde{z}=e_{1}$. The case $d=2$ is open to our knowledge.
Remark 2.8. Let us emphasize that the set $\left\{\omega_{m k}\right\}$ does not verify the gap condition (even in the case 1-D)

$$
\inf _{(m, k) \neq\left(m^{\prime}, k^{\prime}\right)}\left|\omega_{m k}-\omega_{m^{\prime} k^{\prime}}\right|>0
$$

Thus one cannot prove exact controllability in finite time using arguments based on the Ingham inequality.

Remark 2.9. The presence of the space $\mathcal{B}$ in the definition of $\Theta$ is motivated by the application to the nonlinear control system that we give in Section 3.

### 2.2 Proof of Theorem 2.6

The proof of the theorem is based on the following proposition, which is proved in next subsection.

Proposition 2.10. There is a linear continuous operator A from $\tilde{\ell}^{2}$ to $\Theta$ such that $\left\{A(d)\left(\omega_{m k}\right)\right\}=d$ for any $d \in \tilde{\ell}^{2}$.

The idea of the proof of Theorem 2.6 is to rewrite (2.7) in the form $d_{m k}=$ $\check{u}\left(\omega_{m k}\right)$ with $d=\left\{d_{m k}\right\} \in \tilde{\ell}^{2}$ and to apply the proposition. Let us take any $y \in T_{\tilde{z}} \cap \mathcal{V}$. Define

$$
d_{m k}:=\frac{i\left\langle y, e_{m}\right\rangle\left\langle e_{k}, \tilde{z}\right\rangle-i\left\langle e_{k}, y\right\rangle\left\langle\tilde{z}, e_{m}\right\rangle}{Q_{m k}}+C_{m k},
$$

where $C_{m k} \in \mathbb{C}$ and $e_{k}=e_{k, V}$. The fact that $\tilde{z} \in S$ implies

$$
\begin{aligned}
-i \sum_{k=1}^{+\infty}\left\langle\tilde{z}, e_{k}\right\rangle Q_{m k} d_{m k}= & \sum_{k=1}^{+\infty}\left\langle y, e_{m}\right\rangle\left|\left\langle\tilde{z}, e_{k}\right\rangle\right|^{2}-\sum_{k=1}^{+\infty}\left\langle e_{k}, y\right\rangle\left\langle\tilde{z}, e_{m}\right\rangle\left\langle\tilde{z}, e_{k}\right\rangle \\
& -i \sum_{k=1}^{+\infty}\left\langle\tilde{z}, e_{k}\right\rangle Q_{m k} C_{m k} \\
= & \left\langle y, e_{m}\right\rangle-\left\langle\tilde{z}, e_{m}\right\rangle\langle\tilde{z}, y\rangle-i \sum_{k=1}^{+\infty}\left\langle\tilde{z}, e_{k}\right\rangle Q_{m k} C_{m k} .
\end{aligned}
$$

By (2.7), we have $y=R_{\infty}(0, u)$, when

$$
\begin{equation*}
i \sum_{k=1}^{+\infty}\left\langle\tilde{z}, e_{k}\right\rangle Q_{m k} C_{m k}=-\left\langle\tilde{z}, e_{m}\right\rangle\langle\tilde{z}, y\rangle \tag{2.9}
\end{equation*}
$$

for all $m \geq 1$. Thus if we show that there are $C_{m k} \in \mathbb{C}$ such that (2.9) is verified and $d=\left\{d_{m k}\right\} \in \tilde{\ell}^{2}$, then the proof of the theorem will be completed, in view of Proposition 2.10. Notice that, under Condition 2.5, we have

$$
\sum_{m, k=1, m \neq k}^{+\infty}\left|\frac{\left\langle y, e_{m}\right\rangle\left\langle e_{k}, \tilde{z}\right\rangle}{Q_{m k}}\right|^{2} \leq C\|y\|_{\mathcal{V}}^{2}\|\tilde{z}\|_{\mathcal{V}}^{2}<+\infty
$$

Thus $\left\{d_{m k}\right\} \in \tilde{\ell}^{2}$, if $C_{m k} \in \mathbb{C}$ are such that

$$
\begin{align*}
& d_{m m}=\frac{i\left\langle y, e_{m}\right\rangle\left\langle e_{m}, \tilde{z}\right\rangle-i\left\langle e_{m}, y\right\rangle\left\langle\tilde{z}, e_{m}\right\rangle}{Q_{m m}}+C_{m m}=d_{0},  \tag{2.10}\\
& C_{m k}=\bar{C}_{k m},  \tag{2.11}\\
& \quad \sum_{m, k=1, m \neq k}^{+\infty}\left|C_{m k}\right|^{2}<+\infty, \tag{2.12}
\end{align*}
$$

where $d_{0} \in \mathbb{R}$. Let us show that, for an appropriate choice of $d_{0}$, there are $C_{m k}$ satisfying (2.9)-(2.12). Since $y \in T_{\tilde{z}}$, we have $\langle\tilde{z}, y\rangle=i \operatorname{Im}\langle\tilde{z}, y\rangle$. We can rewrite (2.9) and (2.10) in the following form

$$
\begin{align*}
& \sum_{k=1}^{+\infty}\left\langle\tilde{z}, e_{k}\right\rangle Q_{m k} C_{m k}=-\left\langle\tilde{z}, e_{m}\right\rangle \operatorname{Im}\langle\tilde{z}, y\rangle,  \tag{2.13}\\
& d_{m m}=\frac{-2 \operatorname{Im}\left(\left\langle y, e_{m}\right\rangle\left\langle e_{m}, \tilde{z}\right\rangle\right)}{Q_{m m}}+C_{m m}=d_{0} \tag{2.14}
\end{align*}
$$

Case 1. Let as suppose that $\tilde{z}=c e_{p}$, where $c \in \mathbb{C},|c|=1$ and $p \geq 1$. Then (2.11)-(2.14) is verified for $C_{m k}=0$, if $m \neq k$ and $C_{m m}$ defined by (2.14) with any $d_{0} \in \mathbb{R}$.

Case 2. Suppose $\tilde{z}=c_{p} e_{p}+c_{q} e_{q}$, where $c_{p}, c_{q} \in \mathbb{C},\left|c_{p}\right|^{2}+\left|c_{q}\right|^{2}=1$ and $p \neq q$. For any $m \geq 1$, define $C_{m m}$ by (2.14). If $m \neq p$, we set

$$
\begin{equation*}
C_{m p}:=\frac{-c_{m}\left(\operatorname{Im}\langle\tilde{z}, y\rangle+Q_{m m} C_{m m}\right)}{c_{p} Q_{m p}} \tag{2.15}
\end{equation*}
$$

where $c_{m}=0$, if $m \neq q$. Then all the equations in (2.13) are verified, excepted the case $m=p$. Let us show that, for an appropriate choice of $d_{0} \in \mathbb{R}$, this equation is also satisfied. Equation (2.13) for $m=p$ is

$$
c_{p} Q_{p p} C_{p p}+c_{q} Q_{p q} C_{p q}=-c_{p} \operatorname{Im}\langle\tilde{z}, y\rangle .
$$

Using (2.15) for $m=q\left(\right.$ taking $\left.C_{p q}=\bar{C}_{q p}\right)$ and (2.14) for $m=p$, we get

$$
\begin{aligned}
-c_{p} \operatorname{Im}\langle\tilde{z}, y\rangle= & c_{p} Q_{p p}\left(d_{0}+\frac{2 \operatorname{Im}\left(\left\langle y, e_{p}\right\rangle\left\langle e_{p}, \tilde{z}\right\rangle\right)}{Q_{p p}}\right) \\
& +c_{q} Q_{p q}\left(\frac{-c_{q}\left(\operatorname{Im}\langle\tilde{z}, y\rangle+Q_{q q} C_{q q}\right)}{c_{p} Q_{q p}}\right) \\
= & c_{p} Q_{p p}\left(d_{0}+\frac{2 \operatorname{Im}\left(\left\langle y, e_{p}\right\rangle\left\langle e_{p}, \tilde{z}\right\rangle\right)}{Q_{p p}}\right) \\
& +c_{q} Q_{p q}\left(\frac{-\bar{c}_{q} \operatorname{Im}\langle\tilde{z}, y\rangle}{\bar{c}_{p} Q_{q p}}\right) \\
& +c_{q} Q_{p q}\left(\frac{-\bar{c}_{q} Q_{q q} C_{q q}}{\bar{c}_{p} Q_{q p}}\right) .
\end{aligned}
$$

Now using (2.14) for $m=q$, we rewrite this equality in an equivalent form

$$
\left(\left|c_{p}\right|^{2} Q_{p p}-\left|c_{q}\right|^{2} Q_{q q}\right) d_{0}=A
$$

for some constant $A \in \mathbb{R}$. Thus if $\tilde{z}$ is such that $\left|c_{p}\right|^{2} Q_{p p}-\left|c_{q}\right|^{2} Q_{q q} \neq 0$, then we are able to find $C_{m k}$ satisfying (2.11)-(2.14). If $\left|c_{p}\right|^{2} Q_{p p}-\left|c_{q}\right|^{2} Q_{q q}=0$, then
linear system (2.2), (2.3) is not controllable, since for any $u \in \Theta$ and $t \geq 0$ we have

$$
\begin{aligned}
& \frac{\mathrm{d}}{\mathrm{~d} t} \operatorname{Im}\left\langle R_{t}(0, u), c_{p} e^{-i \lambda_{p} t} e_{p}-c_{q} e^{-i \lambda_{q} t} e_{q}\right\rangle \\
& =\operatorname{Im}\left\langle i(\Delta-V) R_{t}(0, u)-i u Q\left(c_{p} e^{-i \lambda_{p} t} e_{p}+c_{q} e^{-i \lambda_{q} t} e_{q}\right), c_{p} e^{-i \lambda_{p} t} e_{p}-c_{q} e^{-i \lambda_{q} t} e_{q}\right\rangle \\
& \quad+\operatorname{Im}\left\langle R_{t}(0, u), i(\Delta-V)\left(c_{p} e^{-i \lambda_{p} t} e_{p}-c_{q} e^{-i \lambda_{q} t} e_{q}\right)\right\rangle \\
& =\operatorname{Im}\left\langle-i u Q\left(c_{p} e^{-i \lambda_{p} t} e_{p}+c_{q} e^{-i \lambda_{q} t} e_{q}\right), c_{p} e^{-i \lambda_{p} t} e_{p}-c_{q} e^{-i \lambda_{q} t} e_{q}\right\rangle \\
& =-u\left(\left|c_{p}\right|^{2} Q_{p p}-\left|c_{q}\right|^{2} Q_{q q}\right)=0 .
\end{aligned}
$$

This non-controllability property is a remark of Beauchard and Coron [8].
Case 3. Here we suppose that $\tilde{z}=\sum_{j=1}^{+\infty} c_{j} e_{j}$ with $c_{p} c_{q} c_{r} \neq 0$, and $p, q, r$ are not equal to each other. If we define again $C_{m p}, m \neq p$ by (2.15) and follow the steps of case 2 , we get the following equation for $d_{0}$

$$
\left(\left|c_{p}\right|^{2} Q_{p p}-\sum_{m \neq p}\left|c_{m}\right|^{2} Q_{m m}\right) d_{0}=\tilde{A}
$$

for some constant $\tilde{A} \in \mathbb{R}$. This implies that for any $\tilde{z}$ such that $\left|c_{p}\right|^{2} Q_{p p}-$ $\sum_{m \neq p}\left|c_{m}\right|^{2} Q_{m m} \neq 0$, we can find $C_{m k}$ satisfying (2.11)-2.14). Let us suppose that

$$
\begin{equation*}
\left|c_{p}\right|^{2} Q_{p p}-\sum_{m \neq p}\left|c_{m}\right|^{2} Q_{m m}=0 \tag{2.16}
\end{equation*}
$$

In this case, we define $C_{m p}$ by (2.15) only for integers $m$ which are different from $p, q$ and $r$. Then all the equations in (2.13) are verified, exepted for $m=p, q, r$. We take any $C_{q p} \in \mathbb{C}$ and choose $C_{q r}$ and $C_{r p}$ such that

$$
\begin{align*}
& c_{p} Q_{r p} C_{r p}+c_{q} Q_{r q} C_{r q}+c_{p} Q_{r r} C_{r r}=-c_{r} \operatorname{Im}\langle\tilde{z}, y\rangle,  \tag{2.17}\\
& c_{p} Q_{q p} C_{q p}+c_{q} Q_{q q} C_{q q}+c_{r} Q_{q r} C_{q r}=-c_{q} \operatorname{Im}\langle\tilde{z}, y\rangle . \tag{2.18}
\end{align*}
$$

Replacing the value of $C_{q r}$ from (2.18) into (2.17), then the value of $C_{p r}$ from (2.17) into (2.13) with $m=p$, and using (2.11), we get the following equation for $d_{0}$

$$
\left(\left|c_{p}\right|^{2} Q_{p p}+\left|c_{q}\right|^{2} Q_{q q}-\sum_{m \neq p, q}\left|c_{m}\right|^{2} Q_{m m}\right) d_{0}=\tilde{\tilde{A}}
$$

for some constant $\tilde{\tilde{A}} \in \mathbb{R}$. Equality (2.16) implies that

$$
\left|c_{p}\right|^{2} Q_{p p}+\left|c_{q}\right|^{2} Q_{q q}-\sum_{m \neq p, q}\left|c_{m}\right|^{2} Q_{m m}=0
$$

if and only if $\left|c_{q}\right|^{2} Q_{q q}=0$, which is not the case: $c_{q} \neq 0, Q_{q q} \neq 0$. Thus solution $d_{0} \in \mathbb{R}$ exists, and the sequence $C_{m k}$ is constructed for any $\tilde{z} \notin \mathcal{E}$.

### 2.3 Proof of Proposition 2.10

The construction of the operator $A$ is based on the following lemma.
Lemma 2.11. For any $d \in \tilde{\ell}^{2}$ and $\varepsilon>0$, there is $u \in B_{\Theta}(0, \varepsilon)$ such that $\left\{\check{u}\left(\omega_{m k}\right)\right\}=d$.
Proof of Proposition 2.10. Let $d^{n}$ be any orthonormal basis in $\tilde{\ell}^{2}$. Applying Lemma 2.11, we find a sequence $u_{n} \in B_{\Theta}\left(0, \frac{1}{n}\right)$ such that $\left\{\check{u}_{n}\left(\omega_{m k}\right)\right\}=d^{n}$. For any $d \in \tilde{\ell}^{2}$, there is $c \in \ell^{2}$ such that $d=\sum_{n=1}^{+\infty} c_{n} d^{n}$. Let us define $A$ in the following way

$$
A(d)=\sum_{n=1}^{+\infty} c_{n} u_{n}
$$

As $u_{n} \in B_{\Theta}\left(0, \frac{1}{n}\right)$, this sum converges in $\Theta$ :

$$
\|A(d)\|_{\Theta} \leq \sum_{n=1}^{+\infty}\left|c_{n}\right|\left\|u_{n}\right\|_{\Theta} \leq\left(\sum_{n=1}^{+\infty}\left|c_{n}\right|^{2}\right)^{\frac{1}{2}}\left(\sum_{n=1}^{+\infty}\left\|u_{n}\right\|_{\Theta}^{2}\right)^{\frac{1}{2}} \leq C\|d\|_{\tilde{\ell}^{2}}
$$

Thus $A: \tilde{\ell}^{2} \rightarrow \Theta$ is linear continuous and $\left\{A \check{(d)}\left(\omega_{m k}\right)\right\}=d$, by construction.
Proof of Lemma 2.11. Let us take any $d \in \tilde{\ell}^{2}$ and introduce the functional

$$
\mathcal{H}(u):=\left\|\left\{\check{u}\left(\omega_{m k}\right)\right\}-d\right\|_{\tilde{\ell}^{2}}^{2}=\left|\check{u}\left(\omega_{11}\right)-d_{11}\right|^{2}+\sum_{m, k=1, m \neq k}^{+\infty}\left|\check{u}\left(\omega_{m k}\right)-d_{m k}\right|^{2}
$$

defined on the space $\Theta$.
Step 1. First, let us show that there is $u_{0} \in \overline{B_{\Theta}(0, \varepsilon)}$ such that

$$
\begin{equation*}
\mathcal{H}\left(u_{0}\right)=\inf _{u \in B_{\Theta}(0, \varepsilon)} \mathcal{H}(u) . \tag{2.19}
\end{equation*}
$$

To this end, let $u_{n} \in \overline{B_{\Theta}(0, \varepsilon)}$ be an arbitrary minimizing sequence. Since $\mathcal{B} \cap H^{s}\left(\mathbb{R}_{+}, \mathbb{R}\right)$ is reflexive, without loss of generality, we can assume that there is $u_{0} \in \overline{B_{\mathcal{B} \cap H^{s}\left(\mathbb{R}_{+}, \mathbb{R}\right)}(0, \varepsilon)}$ such that $u_{n} \rightharpoonup u_{0}$ in $\mathcal{B} \cap H^{s}\left(\mathbb{R}_{+}, \mathbb{R}\right)$. Extracting a subsequence, if necessary, we can assume that $u_{n}(t) \rightarrow u_{0}(t)$ for almost any $t \in \mathbb{R}$. The Fatou lemma implies that

$$
\int_{0}^{+\infty}\left|u_{0}(s)\right| \mathrm{d} s \leq \liminf _{n \rightarrow \infty} \int_{0}^{+\infty}\left|u_{n}(s)\right| \mathrm{d} s \leq \varepsilon
$$

Again extracting a subsequence, one gets $\left\{\check{u}_{n}\left(\omega_{m k}\right)\right\} \rightharpoonup\left\{\check{u}_{0}\left(\omega_{m k}\right)\right\}$ in $\tilde{\ell}^{2}$ as $n \rightarrow+\infty$. This implies that $u_{0} \in \Theta$ and

$$
\mathcal{H}\left(u_{0}\right) \leq \inf _{u \in B_{\Theta}(0, \varepsilon)} \mathcal{H}(u)
$$

The fact that $u_{0} \in \overline{B_{\Theta}(0, \varepsilon)}$ follows from the Fatou lemma and lower weak semicontinuity of norms. Thus we have (2.19).

Step 2. To complete the proof, we need to show that $\mathcal{H}\left(u_{0}\right)=0$. Suppose, by contradiction, that $\mathcal{H}\left(u_{0}\right)>0$. As we shall see below, this implies that there is $v \in \overline{B_{\Theta}(0, \varepsilon)}$ such that

$$
\begin{equation*}
\left.\frac{\mathrm{d}}{\mathrm{~d} t} \mathcal{H}\left((1-t) u_{0}+t v\right)\right|_{t=0}<0 \tag{2.20}
\end{equation*}
$$

Since $(1-t) u_{0}+t v \in \overline{B_{\Theta}(0, \varepsilon)}$ for all $t \in[0,1]$, (2.20) is a contradiction to 2.19).
To construct such a function $v$, notice that the derivative is given explicitly by

$$
\begin{aligned}
\left.\frac{\mathrm{d}}{\mathrm{~d} t} \mathcal{H}\left((1-t) u_{0}+t v\right)\right|_{t=0} & =2 \operatorname{Re}\left[\left(\check{v}\left(\omega_{11}\right)-\check{u}_{0}\left(\omega_{11}\right)\right) \overline{\left(\check{u}_{0}\left(\omega_{11}\right)-d_{11}\right)}\right] \\
& +2 \sum_{m, k=1, m \neq k}^{+\infty} \operatorname{Re}\left[\left(\check{v}\left(\omega_{m k}\right)-\check{u}_{0}\left(\omega_{m k}\right)\right) \overline{\left(\check{u}_{0}\left(\omega_{m k}\right)-d_{m k}\right)}\right] .
\end{aligned}
$$

In view of this equality, the existence of $v$ follows immediately from the following lemma.

Lemma 2.12. The set

$$
U:=\left\{\left\{\check{u}\left(\omega_{m k}\right)\right\}: u \in B_{\Theta}(0, \varepsilon)\right\}
$$

is dense in $\tilde{\ell}^{2}$.
Proof. Suppose that $h \in \tilde{\ell}^{2}$ is orthogonal to $U$ : for any $u \in B_{\Theta}(0, \varepsilon)$ we have

$$
\check{u}\left(\omega_{11}\right) \overline{h_{11}}+\sum_{m, k=1, m \neq k}^{+\infty} \check{u}\left(\omega_{m k}\right) \overline{h_{m k}}=0 .
$$

Replacing in this equality $\check{u}\left(\omega_{m k}\right)$ by its integral representation, we get

$$
\sum_{m, k=1, m \neq k}^{+\infty} \int_{0}^{+\infty} e^{i \omega_{m k} s} u(s) \mathrm{d} s \overline{h_{m k}}=\int_{0}^{+\infty} u(s)\left(\sum_{m, k=1, m \neq k}^{+\infty} e^{i \omega_{m k} s} \overline{h_{m k}}\right) \mathrm{d} s=0
$$

As this equality holds for any $u \in B_{\Theta}(0, \varepsilon)$, we obtain

$$
e^{i \omega_{11} s} \overline{h_{11}}+\sum_{m=1, k, m \neq k}^{+\infty} e^{i \omega_{m k} s} \overline{h_{m k}}=0
$$

for any $s \geq 0$. Condition 2.5 (ii) and Lemma 3.10 in 22 imply that $h=0$. This proves that $U$ is dense.

## 3 Controllability of nonlinear system

### 3.1 Well-posedness of Schrödinger equation

In this section, we suppose that $d=1, D=(0,1)$. We consider the following Schrödinger equation

$$
\begin{align*}
i \dot{z} & =-\frac{\mathrm{d}^{2} z}{\mathrm{~d} x^{2}}+V(x) z+u(t) Q(x) z+v(t) Q(x) y,  \tag{3.1}\\
\left.z\right|_{\partial D} & =0  \tag{3.2}\\
z(0, x) & =z_{0}(x) . \tag{3.3}
\end{align*}
$$

See Proposition 2 in 10 for the proof of well-posedness of this system with $V=0$ in $H_{(0)}^{3}$. Here we give an estimate for the solution which is important for the study of the controllability property.
Proposition 3.1. For any $z_{0} \in H_{(V)}^{3}, u, v \in L^{1}\left(\mathbb{R}_{+}, \mathbb{R}\right)$ and $y \in C\left(\mathbb{R}_{+}, H_{(V)}^{3}\right)$, problem (3.1)-(3.5) has a unique solution $z \in C\left(\mathbb{R}_{+}, H_{(V)}^{3}\right)$. Furthermore, if $u, v \in L^{1}\left(\mathbb{R}_{+}, \mathbb{R}\right) \cap \mathcal{B}$, then there is a constant $C>0$ such that

$$
\begin{align*}
\sup _{t \in \mathbb{R}_{+}}\|z(t)\|_{3, V} \leq & C\left(\left\|z_{0}\right\|_{3, V}+\sup _{t \in \mathbb{R}_{+}}\|y(t)\|_{3, V}\left(\|v\|_{L^{1}\left(\mathbb{R}_{+}\right)}+\|v\|_{\mathcal{B}}\right)\right) \\
& \times \exp \left(C\left(\|u\|_{L^{1}\left(\mathbb{R}_{+}\right)}+1\right) \exp \left(\|u\|_{\mathcal{B}}^{2}\right)\right) \tag{3.4}
\end{align*}
$$

If $v=0$, then for all $t \geq 0$ we have

$$
\begin{equation*}
\|z(t)\|=\left\|z_{0}\right\| . \tag{3.5}
\end{equation*}
$$

Proof. Existence and uniqueness of solution are proved in Proposition 2 in 10. The proof of (3.4) is also uses some ideas of [10]. We give all the details for the sake of completeness.

Let us rewrite (3.1)-(3.3) in the Duhamel form

$$
\begin{equation*}
z(t)=S(t) z_{0}-i \int_{0}^{t} S(t-s)[u(s) Q z(s)+v(s) f(s)] \mathrm{d} s \tag{3.6}
\end{equation*}
$$

For any $u \in L^{1}\left(\mathbb{R}_{+}, \mathbb{R}\right) \cap \mathcal{B}$ and $z \in C\left(\mathbb{R}_{+}, H_{(V)}^{3}\right)$, we estimate the function

$$
G_{t}(z):=\int_{0}^{t} S(-s)(u(s) Q z(s)) \mathrm{d} s
$$

Integration by parts gives (we write $\lambda_{j}, e_{j}$ instead of $\lambda_{j, V}, e_{j, V}$ )

$$
\begin{aligned}
\left\langle Q z(s), e_{j}\right\rangle= & \frac{1}{\lambda_{j}}\left\langle\left(-\frac{\mathrm{d}^{2}}{\mathrm{~d} x^{2}}+V\right)(Q z), e_{j}\right\rangle \\
= & \frac{1}{\lambda_{j}^{2}}\left\langle\left(-\frac{\mathrm{d}^{2}}{\mathrm{~d} x^{2}}+V\right)(Q z),\left(-\frac{\mathrm{d}^{2}}{\mathrm{~d} x^{2}}+V\right) e_{j}\right\rangle=\left.\frac{1}{\lambda_{j}^{2}} \frac{\mathrm{~d}^{2}}{\mathrm{~d} x^{2}}(Q z) \frac{\mathrm{d}}{\mathrm{~d} x} e_{j}\right|_{x=0} ^{x=1} \\
& +\frac{1}{\lambda_{j}^{2}}\left(\left\langle V\left(-\frac{\mathrm{d}^{2}}{\mathrm{~d} x^{2}}+V\right)(Q z), e_{j}\right\rangle+\left\langle\frac{\mathrm{d}}{\mathrm{~d} x}\left(-\frac{\mathrm{d}^{2}}{\mathrm{~d} x^{2}}+V\right)(Q z), \frac{\mathrm{d}}{\mathrm{~d} x} e_{j}\right\rangle\right) \\
= & : I_{j}+J_{j} .
\end{aligned}
$$

Thus

$$
\begin{align*}
\left\|G_{t}(z)\right\|_{3, V}^{2} & =\sum_{j=1}^{+\infty}\left(j^{3} \int_{0}^{t} e^{i \lambda_{j} s} u(s)\left\langle Q z(s), e_{j}\right\rangle \mathrm{d} s\right)^{2} \\
& =\sum_{j=1}^{+\infty}\left(j^{3} \int_{0}^{t} e^{i \lambda_{j} s} u(s)\left(I_{j}+J_{j}\right) \mathrm{d} s\right)^{2} \tag{3.7}
\end{align*}
$$

Using (1.8), we get

$$
\left\langle\frac{\mathrm{d}}{\mathrm{~d} x}\left(-\frac{\mathrm{d}^{2}}{\mathrm{~d} x^{2}}+V\right) Q z, \frac{\mathrm{~d}}{\mathrm{~d} x} e_{j}\right\rangle=j \pi\left\langle\frac{\mathrm{~d}}{\mathrm{~d} x}\left(-\frac{\mathrm{d}^{2}}{\mathrm{~d} x^{2}}+V\right) Q z, \sqrt{2} \cos (j \pi x)\right\rangle+s_{j}(z),
$$

where $\left|s_{j}(z)\right| \leq C\|z\|_{3, V}$ for all $j \geq 1$. The definition of $J_{j}$, the fact that $\{\sqrt{2} \cos (j \pi x)\}$ is an orthonormal system in $L^{2}$, (1.6) and the Minkowski inequality yield

$$
\begin{equation*}
\sum_{j=1}^{+\infty}\left(j^{3} \int_{0}^{t} e^{i \lambda_{j} s} u(s) J_{j} \mathrm{~d} s\right)^{2} \leq C\left(\int_{0}^{t}|u(s)|\|z(s)\|_{3, V} \mathrm{~d} s\right)^{2} \tag{3.8}
\end{equation*}
$$

On the other hand, (1.8) implies that

$$
\left.\frac{\mathrm{d}^{2}}{\mathrm{~d} x^{2}}(Q z) \frac{\mathrm{d}}{\mathrm{~d} x} e_{j}\right|_{x=0} ^{x=1}=\left.j \pi \frac{\mathrm{~d}^{2}}{\mathrm{~d} x^{2}}(Q z) \sqrt{2} \cos (j \pi x)\right|_{x=0} ^{x=1}+\tilde{s}_{j}(z)=: j c_{j}(z)+\tilde{s}_{j}(z)
$$

where $\left|\tilde{s}_{j}\right| \leq C\|z\|_{3, V}$ for all $j \geq 1$ and $c_{j}(z)$ depends on the parity of $j$. Again applying the Minkowski inequality, we obtain

$$
\begin{equation*}
\sum_{j=1}^{+\infty}\left(j^{3} \int_{0}^{t} e^{i \lambda_{j} s} u(s) \tilde{s}_{j}(z) \mathrm{d} s\right)^{2} \leq C\left(\int_{0}^{t}|u(s)|\|z(s)\|_{3, V} \mathrm{~d} s\right)^{2} \tag{3.9}
\end{equation*}
$$

Without loss of generality, we can assume that $c(z):=c_{j}(z)$ does not depend on $j$. Thus we cannot conclude as in the case of $J_{j}$. Here we use the fact that $u \in \mathcal{B}$. Let $P \geq 1$ be an integer such that $P \leq t<P+1$. Using the Cauchy-Schwarz and the Ingham inequalities, we obtain

$$
\sum_{j=1}^{+\infty}\left(\int_{0}^{t} e^{i \lambda_{j} s} u(s) c(z) \mathrm{d} s\right)^{2}=\sum_{j=1}^{+\infty}\left(\left(\int_{P}^{t}+\sum_{p=1}^{P} \int_{p-1}^{p}\right) e^{i \lambda_{j} s} u(s) c(z) \mathrm{d} s\right)^{2}
$$

$$
\begin{aligned}
\leq & 2 \sum_{j=1}^{+\infty}\left(\int_{P}^{t} e^{i \lambda_{j} s} u(s) c(z) \mathrm{d} s\right)^{2} \\
& +2 \sum_{j=1}^{+\infty}\left(\sum_{p=1}^{P} \frac{1}{p^{2}}\right)\left(\sum_{p=1}^{P} p^{2}\left(\int_{p-1}^{p} e^{i \lambda_{j} s} u(s) c(z) \mathrm{d} s\right)^{2}\right) \\
\leq & C\|u(s) c(z)\|_{L^{2}([P, t])}^{2}+C \sum_{p=1}^{P} p^{2} \sum_{j=1}^{+\infty}\left(\int_{p-1}^{p} e^{i \lambda_{j} s} u(s) c(z) \mathrm{d} s\right)^{2} \\
\leq & C\|u(s) c(z)\|_{L^{2}([P, t])}^{2}+C \sum_{p=1}^{P} p^{2}\|u(s) c(z)\|_{L^{2}([p-1, p])}^{2} \\
\leq & C \int_{0}^{t} w(s)\|z(s)\|_{3, V}^{2} \mathrm{~d} s .
\end{aligned}
$$

where $w(s)=|u(s)|^{2} \chi_{[P, t]}(s)+\sum_{p=1}^{P} p^{2}|u(s)|^{2} \chi_{[p-1, p]}(s)$. Notice that

$$
\begin{equation*}
\int_{0}^{t} w(s) \mathrm{d} s \leq\|u\|_{\mathcal{B}}^{2} \quad \text { for all } t \geq 0 \tag{3.10}
\end{equation*}
$$

Combining (3.7)-(3.10), we get

$$
\begin{equation*}
\left\|G_{t}(z)\right\|_{3, V} \leq C\left(\int_{0}^{t} w(s)\|z(s)\|_{3, V}^{2} \mathrm{~d} s\right)^{\frac{1}{2}}+C \int_{0}^{t}|u(s)|\|z(s)\|_{3, V} \mathrm{~d} s \tag{3.11}
\end{equation*}
$$

The quantity

$$
\tilde{G}_{t}(f):=\int_{0}^{t} S(-s)(v(s) f(s)) \mathrm{d} s
$$

is estimated in a similar way

$$
\begin{align*}
\left\|\tilde{G}_{t}\right\|_{3, V} & \leq C\left(\int_{0}^{t} \tilde{w}(s)\|f(s)\|_{3, V}^{2} \mathrm{~d} s\right)^{\frac{1}{2}}+C \int_{0}^{t}|v(s)|\|f(s)\|_{3, V} \mathrm{~d} s \\
& \leq C\|f(s)\|_{3, V}\left(\|v\|_{L^{1}\left(\mathbb{R}_{+}\right)}+\|v\|_{\mathcal{B}}\right) \tag{3.12}
\end{align*}
$$

where $\tilde{w}(s)=|v(s)|^{2} \chi_{[P, t]}(s)+\sum_{p=1}^{P} p^{2}|v(s)|^{2} \chi_{[p-1, p]}(s)$. Let us prove (3.4).
From (3.6) and (3.11) we have

$$
\begin{aligned}
& \|z(t)\|_{3, V}^{2} \leq C\left(\left\|z_{0}\right\|_{3, V}^{2}+\left\|\tilde{G}_{t}\right\|_{3, V}^{2}+\left\|G_{t}\right\|_{3, V}^{2}\right) \\
& \leq C\left(\left\|z_{0}\right\|_{3, V}^{2}+\left\|\tilde{G}_{t}\right\|_{3, V}^{2}+\int_{0}^{t} w(s)\|z(s)\|_{3, V}^{2} \mathrm{~d} s+\left(\int_{0}^{t}|u(s)|\|z(s)\|_{3, V} \mathrm{~d} s\right)^{2}\right)
\end{aligned}
$$

The Gronwall inequality implies

$$
\begin{aligned}
&\|z(t)\|_{3, V}^{2} \leq C\left(\left\|z_{0}\right\|_{3, V}^{2}+\left\|\tilde{G}_{t}\right\|_{3, V}^{2}+\left(\int_{0}^{t}|u(s)|\|z(s)\|_{3, V} \mathrm{~d} s\right)^{2}\right) \\
& \times \exp \left(C \int_{0}^{t} w(s) \mathrm{d} s\right)
\end{aligned}
$$

Taking the square root of this inequality, using (3.10) and the Gronwall inequality, we obtain

$$
\begin{aligned}
\|z(t)\|_{3, V} \leq & C\left(\left\|z_{0}\right\|_{3, V}+\left\|\tilde{G}_{t}\right\|_{3, V}\right) \\
& \times \exp \left(C\left(\int_{0}^{t} w(s) \mathrm{d} s+\int_{0}^{t}|u(s)| \mathrm{d} s \exp \left(\int_{0}^{t} w(s) \mathrm{d} s\right)\right)\right) \\
\leq & C\left(\left\|z_{0}\right\|_{3, V}+\left\|\tilde{G}_{t}\right\|_{3, V}\right) \exp \left(C\left(\|u\|_{L^{1}\left(\mathbb{R}_{+}\right)}+1\right) \exp \left(\|u\|_{\mathcal{B}}^{2}\right)\right)
\end{aligned}
$$

In view of (3.12), this completes the proof of the proposition.
Remark 3.2. Let us notice that, one should not expect to have a well-posedness property in any Sobolev space $H^{k}$ with controls in $L^{1}$. Indeed, exact controllability property in $H^{3}$, proved by Beauchard and Laurent 10 in the case $d=1$, implies that the problem is not well posed in spaces $H^{3+\sigma}$ for any $\sigma>0$ (a point $z_{1} \in H^{3} \backslash H^{3+\sigma}$ would not be accessible from a point $z_{0} \in H^{3+\sigma}$ ). Schrödinger equation is well posed in higher Sobolev spaces, when control $u$ is more regular.

Corollary 3.3. Denote by $\mathcal{U}_{t}(\cdot, \cdot): H_{(V)}^{3} \times L^{1}\left(\mathbb{R}_{+}, \mathbb{R}\right) \cap \mathcal{B} \rightarrow H_{(V)}^{3}$ the resolving operator of (1.1), (1.2). Then $\mathcal{U}_{t}(\cdot, \cdot)$ is locally Lipschitz continuous, i.e., for any $\delta>0$ there is $C>0$ such that

$$
\begin{equation*}
\sup _{t \in \mathbb{R}_{+}}\left\|\mathcal{U}_{t}\left(z_{0}, u\right)-\mathcal{U}_{t}\left(z_{0}^{\prime}, u^{\prime}\right)\right\|_{3, V} \leq C\left\|\left(z_{0}, u\right)-\left(z_{0}^{\prime}, u^{\prime}\right)\right\|_{H_{(V)}^{3} \times L^{1}\left(\mathbb{R}_{+}, \mathbb{R}\right) \cap \mathcal{B}} \tag{3.13}
\end{equation*}
$$

for all $\left(z_{0}, u\right),\left(z_{0}^{\prime}, u^{\prime}\right) \in B_{H_{(V)}^{3} \times L^{1}\left(\mathbb{R}_{+}, \mathbb{R}\right) \cap \mathcal{B}}(0, \delta)$, where $L^{1}\left(\mathbb{R}_{+}, \mathbb{R}\right) \cap \mathcal{B}$ is endowed with the norm $\|\cdot\|_{L^{1}\left(\mathbb{R}_{+}, \mathbb{R}^{\prime}\right) \cap \mathcal{B}}:=\|\cdot\|_{L^{1}}+\|\cdot\|_{\mathcal{B}}$.
Proof. Notice that $z(t):=\mathcal{U}_{t}\left(z_{0}, u\right)-\mathcal{U}_{t}\left(z_{0}^{\prime}, u^{\prime}\right)$ is a solution of problem

$$
\begin{aligned}
i \dot{z} & =-\frac{\mathrm{d}^{2} z}{\mathrm{~d}^{2} x}+u(t) Q(x) z+\left(u(t)-u^{\prime}(t)\right) Q(x) \mathcal{U}_{t}\left(z_{0}^{\prime}, u^{\prime}\right), \\
\left.z\right|_{\partial D} & =0 \\
z(0, x) & =z_{0}(x)-z_{0}^{\prime}(x) .
\end{aligned}
$$

Applying Proposition 3.1, we get (3.13).

### 3.2 Exact controllability in infinite time

For any control $u \in \Theta$, problem (3.1), (3.2) is well-posed in Sobolev space $H_{(V)}^{3}$. Equality (3.5) implies that it suffices to consider the controllability properties of (3.1), (3.2) on the unit sphere $S$ in $L^{2}$. Let $\mathcal{U}_{\infty}\left(z_{0}, u\right)$ be the $H^{3}$-weak $\omega$ limit set of the trajectory corresponding to control $u \in \Theta$ and initial condition $z_{0} \in H_{(V)}^{3}$ :

$$
\begin{equation*}
\mathcal{U}_{\infty}\left(z_{0}, u\right):=\left\{z \in H_{(V)}^{3}: \mathcal{U}_{t_{n}}\left(z_{0}, u\right) \rightharpoonup z \text { in } H^{3} \text { for some } t_{n} \rightarrow+\infty\right\} \tag{3.14}
\end{equation*}
$$

By (3.4), $\mathcal{U}_{t}\left(z_{0}, u\right)$ is bounded in $H^{3}$, thus $\mathcal{U}_{\infty}\left(z_{0}, u\right)$ is non-empty.

Definition 3.4. We say that (3.1), (3.8) is exactly controllable in infinite time in subset $H \subset S$, if for any $z_{0}, z_{1} \in H$ there is a control $u \in \Theta$ such that $z_{1} \in \mathcal{U}_{\infty}\left(z_{0}, u\right)$.

Below theorem is one of the main results of this paper.
Theorem 3.5. Under Condition 2.5, for any $\tilde{z} \in S \cap H_{(V)}^{3}$ there is $\delta>0$ such that problem (3.1), (3.2) is exactly controllable in infinite time in $S \cap B_{H_{(V)}^{3}}(\tilde{z}, \delta)$.

See Section 3.3 for the proof. Working in higher Sobolev spaces, one can prove similar exact controllability results with more regular controls. For example:

Theorem 3.6. Under Condition 2.5, for any $\tilde{z} \in S \cap H_{(V)}^{3+\sigma}, \sigma \in(0,2]$ there is $\delta>0$ such that problem (3.1), (3.2) is exactly controllable in infinite time in $S \cap B_{H_{(V)}^{3+\sigma}}(\tilde{z}, \delta)$ with controls $u \in W^{1,1}\left(\mathbb{R}_{+}, \mathbb{R}\right) \cap H^{s}\left(\mathbb{R}_{+}, \mathbb{R}\right)$.

These local exact controllability properties imply the following global exact controllability result.

Theorem 3.7. Under Condition 2.5, problem (3.1), (3.2) is exactly controllable in infinite time in $S \cap H_{(V)}^{3}$ in the following sense: for any $z_{0} \in S \cap H_{(V)}^{3+\sigma}, \sigma \in$ $(0,2]$ and $z_{1} \in S \cap H_{(V)}^{3}$ there is a control $u \in L^{1}\left(\mathbb{R}_{+}, \mathbb{R}\right)$ such that $z_{1} \in$ $\mathcal{U}_{\infty}\left(z_{0}, u\right)$.

Proof. Let $\gamma:[0,1] \rightarrow S \cap H_{(V)}^{3}$ be any continuous function such that $\gamma(0)=z_{0}$, $\gamma(1)=z_{1}$ and $\gamma(s) \in H_{(V)}^{3+\sigma}$ for any $s \in[0,1)$. Using the compactness of the curve $\gamma$ and Theorem 3.6, we prove that there is a control $v$ and time $T>0$ such that $\mathcal{U}_{T}\left(z_{0}, v\right) \in B_{H_{(V)}^{3}}\left(z_{1}, \delta_{z_{1}}\right)$, where $\delta_{z_{1}}>0$ is the constant in Theorem 3.5 corresponding to $z_{1}$. This completes the proof.

Remark 3.8. We do not know if problem (1.1)-(1.3) is well posed in the space $\mathcal{V}$ for $d \geq 2$ with $\Theta$-controls. Well-posedness in $\mathcal{V}$ with $u \in \Theta$ would imply the controllability of the multidimensional problem. The nonlinear problem's solution is in $\mathcal{V}$ for more regular controls.

### 3.3 Proof of Theorem 3.5

The proof is based on an inverse mapping theorem for multivalued functions. We apply the inverse mapping theorem established by Nachi and Penot [21], which suits well to the setting of Schrödinger equation. For the reader's convenience, we recall the statement of their result in Appendix (see Theorem 5.3).

Let us first slightly modify the definition (3.14) of the set $\mathcal{U}_{\infty}\left(z_{0}, u\right)$. Let $T_{n} \rightarrow+\infty$ be the sequence defined in Section 2.1. Define

$$
\begin{equation*}
\mathcal{U}_{\infty}\left(z_{0}, u\right):=\left\{z \in H_{(V)}^{3}: \mathcal{U}_{T_{n_{k}}}\left(z_{0}, u\right) \rightharpoonup z \text { in } H^{3} \text { for some } n_{k} \rightarrow+\infty\right\} \tag{3.15}
\end{equation*}
$$

Consider the multivalued function

$$
\begin{array}{r}
\mathcal{U}_{\infty}(\cdot, \cdot): S \cap H_{(V)}^{3} \times \Theta \rightarrow 2^{S \cap H_{(V)}^{3}} \\
\left(z_{0}, u\right) \rightarrow \mathcal{U}_{\infty}\left(z_{0}, u\right) .
\end{array}
$$

Since the result of Nachi and Penot is stated in the case of Banach spaces, we cannot apply it directly to $\mathcal{U}_{\infty}$. Following Beauchard and Laurent 10, we project the system onto the tangent space $T_{\tilde{z}}$. We apply Theorem 5.3 to the following multivalued function

$$
\begin{aligned}
\tilde{\mathcal{U}}_{\infty}(\cdot, \cdot): T_{\tilde{z}} \cap H_{(V)}^{3} \times \Theta & \rightarrow 2^{T_{\tilde{z}} \cap H_{(V)}^{3}}, \\
\left(z_{0}, u\right) & \rightarrow P \mathcal{U}_{\infty}\left(P^{-1} z_{0}, u\right),
\end{aligned}
$$

where $P$ is the orthogonal projection in $L^{2}$ onto $T_{\tilde{z}}$, i.e., $P z=z-\operatorname{Re}\langle z, \tilde{z}\rangle \tilde{z}, z \in$ $L^{2}$. Notice that $P^{-1}: B_{T_{\tilde{z}}}(0, \delta) \rightarrow S$ is well defined for sufficiently small $\delta>0$. By the definition of $T_{n}$, we have $\lim _{n \rightarrow+\infty} \mathcal{U}_{T_{n}}(\tilde{z}, 0)=\tilde{z}$. Hence (3.15) implies that $\mathcal{U}_{\infty}(\tilde{z}, 0)=\tilde{z}$ and $\tilde{\mathcal{U}}_{\infty}(0,0)=\{0\}$. If we show that $\tilde{\mathcal{U}}_{\infty}$ is strictly differentiable at $\left(x_{0}, y_{0}\right)$ with $x_{0}=(0,0) \in T_{\tilde{z}} \cap H_{(V)}^{3} \times \Theta$ and $y_{0}=0 \in T_{\tilde{z}} \cap H_{(V)}^{3}$ (see Definition 5.2), and the derivative admits a right inverse, then Theorem 3.5 will be proved as a consequence of Theorem 5.3.

Proposition 3.9. The multifunction $\tilde{\mathcal{U}}_{\infty}$ is strictly differentiable at $((0,0), 0) \in$ $T_{\tilde{z}} \cap H_{(V)}^{3} \times \Theta \times T_{\tilde{z}} \cap H_{(V)}^{3}$ in the sense of Definition 5.7. Moreover, the differential is the mapping

$$
\begin{array}{r}
R_{\infty}(\cdot, \cdot): T_{\tilde{z}} \cap H_{(V)}^{3} \times \Theta \rightarrow T_{\tilde{z}} \cap H_{(V)}^{3}, \\
\left(z_{0}, u\right) \rightarrow R_{\infty}\left(z_{0}, u\right),
\end{array}
$$

where $R_{\infty}$ is defined in Section 2.1.

Proof of Theorem 3.5. Case 1. Let us suppose that $\tilde{z} \in S \cap H_{(V)}^{3} \backslash \mathcal{E}$. For any $\left(z_{0}, u\right) \in B_{T_{\tilde{z}} \cap H_{(V)}^{3} \times \Theta}(0, \delta)$, the set $\tilde{\mathcal{U}}_{\infty}\left(z_{0}, u\right)$ is closed and non-empty, if $\delta>0$ is sufficiently small. The mapping $R_{\infty}$ is invertible in view of Theorem 2.6. Thus Theorem 5.3 completes the proof.

Case 2. In the case $\tilde{z} \in S \cap H_{(V)}^{3} \cap \mathcal{E}$, the linearized system (2.2), (2.3) is not controllable, and $R_{\infty}$ is not invertible. Controllability near $\tilde{z}$ in finite time is proved by Beauchard and Coron [8]. They show that the linearized system is controllable up to codimension one. This implies that the nonlinear system is also controllable up to codimnsion one. The controllability in the missed directions is proved using the intermediate values theorem.

Proof of Proposition 3.8. It suffices to show that for any $\varepsilon>0$ there exists $\delta>0$ for which
$e\left(\tilde{\mathcal{U}}_{\infty}\left(z_{0}, u\right)-R_{\infty}\left(z_{0}, u\right), \tilde{\mathcal{U}}_{\infty}\left(z_{0}^{\prime}, u^{\prime}\right)-R_{\infty}\left(z_{0}^{\prime}, u^{\prime}\right)\right) \leq \varepsilon\left\|\left(z_{0}, u\right)-\left(z_{0}^{\prime}, u^{\prime}\right)\right\|_{T_{\tilde{z}} \cap H_{(V)}^{3} \times \Theta}$,
whenever $\left(z_{0}, u\right),\left(z_{0}^{\prime}, u^{\prime}\right) \in B_{T_{\tilde{z} \cap H^{\prime}}^{3} \times \Theta}((0,0), \delta)$. Here $e(\cdot, \cdot)$ stands for the Hausdorff distance (see Appendix for the definition). We show the following stronger estimate

$$
\begin{aligned}
& \sup _{t \in \mathbb{R}_{+}} \| \mathcal{U}_{t}\left(P^{-1} z_{0}, u\right)- R_{t}\left(z_{0}, u\right)-\mathcal{U}_{t}\left(P^{-1} z_{0}^{\prime}, u^{\prime}\right)+R_{t}\left(z_{0}^{\prime}, u^{\prime}\right) \|_{T_{\tilde{z}} \cap H_{(V)}^{3} \times T_{\tilde{z}} \cap H_{(V)}^{3}} \\
& \leq \varepsilon\left\|\left(z_{0}, u\right)-\left(z_{0}^{\prime}, u^{\prime}\right)\right\|_{T_{\tilde{z}} \cap H_{(V)}^{3} \times \Theta}
\end{aligned}
$$

To prove this, notice that the function

$$
y(t):=\mathcal{U}_{t}\left(P^{-1} z_{0}, u\right)-R_{t}\left(z_{0}, u\right)-\mathcal{U}_{t}\left(P^{-1} z_{0}^{\prime}, u^{\prime}\right)+R_{t}\left(z_{0}^{\prime}, u^{\prime}\right)
$$

is a solution of the problem

$$
\begin{aligned}
i \dot{y}=-\frac{\mathrm{d}^{2} y}{\mathrm{~d} x^{2}}+\left(u-u^{\prime}\right) Q\left(\mathcal{U}_{t}\left(P^{-1} z_{0}, u\right)\right. & \left.-\mathcal{U}_{t}(\tilde{z}, 0)\right) \\
& +u^{\prime} Q\left(\mathcal{U}_{t}\left(P^{-1} z_{0}, u\right)-\mathcal{U}_{t}\left(P^{-1} z_{0}^{\prime}, u\right)\right)
\end{aligned}
$$

$\left.y\right|_{\partial D}=0$,
$y(0, x)=P^{-1} z_{0}-z_{0}-P^{-1} z_{0}^{\prime}+z_{0}^{\prime}$.
We have

$$
\begin{equation*}
\|y(0)\|_{3, V} \leq \varepsilon\left\|z_{0}-z_{0}^{\prime}\right\|_{3, V} \tag{3.16}
\end{equation*}
$$

for any $z_{0}, z_{0}^{\prime} \in B_{T_{\tilde{z}} \cap H_{(V)}^{3}}(0, \delta)$ and for sufficiently small $\delta>0$. Using (3.4) (we use the version of the inequality with $v_{1} f_{1}+v_{2} f_{2}$ instead of $v f$ ), Corollary 3.3 and (3.16), we get

$$
\begin{aligned}
\sup _{t \in \mathbb{R}_{+}}\|\tilde{y}(t)\|_{3, V} \leq & C\left(\|y(0)\|_{3, V}+\sup _{t \in \mathbb{R}_{+}}\left\|\mathcal{U}_{t}\left(P^{-1} z_{0}, u\right)-\mathcal{U}_{t}(\tilde{z}, 0)\right\|_{3, V}\left\|u-u^{\prime}\right\|_{\Theta}\right. \\
& \left.+\sup _{t \in \mathbb{R}_{+}}\left\|\mathcal{U}_{t}\left(P^{-1} z_{0}, u\right)-\mathcal{U}_{t}\left(P^{-1} z_{0}^{\prime}, u\right)\right\|_{3, V}\left\|u^{\prime}\right\|_{\Theta}\right) \\
\leq & C\left(\|y(0)\|_{3, V}+\left(\left\|z_{0}\right\|_{3, V}+\left\|u^{\prime}\right\|_{\Theta}\right)\left\|u-u^{\prime}\right\|_{\Theta}+\left\|z_{0}-z_{0}^{\prime}\right\|_{3, V}\left\|u^{\prime}\right\|_{\Theta}\right) \\
\leq & \varepsilon\left\|\left(z_{0}, u\right)-\left(z_{0}^{\prime}, u^{\prime}\right)\right\|_{T_{\tilde{z}} \cap H_{(V)}^{3} \times \Theta}
\end{aligned}
$$

for sufficiently small $\delta$. This proves the proposition.

## 4 Non-controllability result

### 4.1 Main result

In this section, we study the problem of non-controllability of Schrödinger equation

$$
\begin{align*}
i \dot{z} & =-\Delta z+V(x) z+u(t) Q(x) z, \quad x \in D  \tag{4.1}\\
\left.z\right|_{\partial D} & =0  \tag{4.2}\\
z(0, x) & =z_{0}(x) \tag{4.3}
\end{align*}
$$

where $D \subset \mathbb{R}^{d}$ is a bounded domain with smooth boundary, $V, Q \in C^{\infty}(\bar{D}, \mathbb{R})$ are arbitrary given functions. The following lemma establishes the well-posedness of system (4.1)-(4.3) in the space $L^{2}$.

Lemma 4.1. For any $z_{0} \in L^{2}$ and for any $u \in L_{\text {loc }}^{1}\left(\mathbb{R}_{+}, \mathbb{R}\right)$, problem (4.1)(4.3) has a unique solution $z \in C\left(\mathbb{R}_{+}, L^{2}\right)$. Furthermore, the resolving operator $\mathcal{U}_{t}(\cdot, u): L^{2} \rightarrow L^{2}$ taking $z_{0}$ to $z(t)$ satisfies the relation

$$
\left\|\mathcal{U}_{t}\left(z_{0}, u\right)\right\|=\left\|z_{0}\right\|, \quad t \geq 0 .
$$

See 11] for the proof. Let us define the set of attainability of system (4.1), (4.2) from an initial point $z_{0} \in S$ :

$$
\begin{equation*}
\mathcal{A}\left(z_{0}\right):=\left\{\mathcal{U}_{t}\left(z_{0}, u\right): \text { for all } u \in W_{\mathrm{loc}}^{1,1}\left(\mathbb{R}_{+}, \mathbb{R}\right) \text { and } t \geq 0\right\} \tag{4.4}
\end{equation*}
$$

The following theorem is the main result of this section.
Theorem 4.2. For any constant $k \in(0, d)$, any initial condition $z_{0} \in S$ and any ball $B \subset H_{(V)}^{k}$, we have

$$
\mathcal{A}^{c}\left(z_{0}\right) \cap B \cap S \neq \varnothing .
$$

The proof of this theorem is an adaptation of ideas of Shirikyan 28 to the case of Schrödinger equation. Using a Hölder type estimate for the solution of the equation, we show that the image by the resolving operator $\mathcal{U}$ of a ball in the space of controls has a Kolmogorov $\varepsilon$-entropy strictly less than that of a ball $B$ in the phase space $H_{(V)}^{k}$. As we show, this implies the non-controllability.

### 4.2 Some $\varepsilon$-entropy estimates

Let $X$ be a Banach space. For any compact set $K \subset X$ and $\varepsilon>0$, we denote by $N_{\varepsilon}(K, X)$ the minimal number of sets of diameters $\leq 2 \varepsilon$ that are needed to cover $K$. The Kolmogorov $\varepsilon$-entropy of $K$ is defined as $H_{\varepsilon}(K, X)=\ln N_{\varepsilon}(K, X)$.

Let $Y$ be another Banach space and let $f: K \rightarrow Y$ be a Lipschitz continuous function:

$$
\begin{equation*}
\left\|f\left(u_{1}\right)-f\left(u_{2}\right)\right\|_{Y} \leq L\left\|u_{1}-u_{2}\right\|_{X}^{\theta} \tag{4.5}
\end{equation*}
$$

for any $u_{1}, u_{2} \in K$ and for some constants $L>0$ and $\theta \in(0,1)$. The following lemma follows immediately from the definition of $\varepsilon$-entropy (cf. Lemma 2.1 in (28).

Lemma 4.3. For any compact set $K \subset X$ and any function $f: K \rightarrow Y$ satisfying inequality (4.5), we have

$$
H_{\varepsilon}(f(K), Y) \leq H_{\left(\frac{\varepsilon}{L}\right)^{\frac{1}{\theta}}}(K, X) \text { for all } \varepsilon>0
$$

We also need the following two lemmas.

Lemma 4.4. For any $T>0$ and for any closed ball $B \subset W^{1,1}([0, T], \mathbb{R})$, there is a constant $C>0$ such that

$$
H_{\varepsilon}\left(B, L^{1}([0, T], \mathbb{R})\right) \leq \frac{C}{\varepsilon} \ln \frac{1}{\varepsilon}
$$

This is Proposition 2.3 in [28.
Lemma 4.5. For any $k>0$ and any closed ball $B:=\overline{B_{H_{(V)}^{k}}\left(z_{0}, r\right)}$ such that $B_{H_{(V)}^{k}}\left(z_{0}, r\right) \cap S \neq \varnothing$ there is a constant $C>0$ such that

$$
\begin{equation*}
H_{\varepsilon}\left(B \cap S, H^{k-1}\right) \geq C\left(\frac{1}{\varepsilon}\right)^{d} \tag{4.6}
\end{equation*}
$$

Proof. It is well known that

$$
\begin{equation*}
C_{1}\left(\frac{1}{\varepsilon}\right)^{d} \leq H_{\varepsilon}\left(B, H^{k-1}\right) \leq C_{2}\left(\frac{1}{\varepsilon}\right)^{d} \tag{4.7}
\end{equation*}
$$

for some constants $C_{1}, C_{2}>0$ (e.g., see 14). Consider the mapping

$$
\begin{gathered}
f:\left[\frac{1}{2}, \frac{3}{2}\right] \times B \cap S \rightarrow H^{k-1} \\
(s, z) \rightarrow s z
\end{gathered}
$$

The set $f\left(\left[\frac{1}{2}, \frac{3}{2}\right] \times B \cap S\right)$ has a non-empty interior, so there is a ball $\tilde{B}$ in $H^{k}$ such that

$$
\begin{equation*}
\tilde{B} \subset f\left(\left[\frac{1}{2}, \frac{3}{2}\right] \times B \cap S\right) \tag{4.8}
\end{equation*}
$$

Clearly,

$$
\left\|f\left(s_{1}, z_{1}\right)-f\left(s_{2}, z_{2}\right)\right\|_{k-1} \leq C\left(\left|s_{1}-s_{2}\right|+\left\|z_{1}-z_{2}\right\|_{k-1}\right)
$$

Using (4.8) and Lemma 4.3, we get

$$
\begin{aligned}
H_{\varepsilon}\left(\tilde{B}, H^{k-1}\right) & \leq H_{\varepsilon}\left(f\left(\left[\frac{1}{2}, \frac{3}{2}\right] \times B \cap S\right), H^{k-1}\right) \\
& \leq H_{\frac{\varepsilon}{C}}\left(\left[\frac{1}{2}, \frac{3}{2}\right] \times B \cap S, \mathbb{R} \times H^{k-1}\right) \\
& \leq H_{\frac{\varepsilon}{C}}\left(\left[\frac{1}{2}, \frac{3}{2}\right], \mathbb{R}\right)+H_{\frac{\varepsilon}{C}}\left(B \cap S, H^{k-1}\right) \\
& \leq C\left(\ln \frac{1}{\varepsilon}+H_{\varepsilon}\left(B \cap S, H^{k-1}\right)\right)
\end{aligned}
$$

Combining this with (4.7) for $\tilde{B}$, we obtain (4.6).

### 4.3 Proof of Theorem 4.2

Let us suppose, by contradiction, that there is $k \in(0, d)$, an initial point $z_{0} \in S$ and a ball $B \subset H_{(V)}^{k}$ such that

$$
\begin{equation*}
B \cap S \subset \mathcal{A}\left(z_{0}\right) \tag{4.9}
\end{equation*}
$$

where $\mathcal{A}$ is the set of attainability of system (4.1), (4.2) from the initial point $z_{0}$ defined by (4.4). Let us set

$$
\begin{aligned}
B_{m} & :=[0, m] \times B_{W^{1,1}([0, m], \mathbb{R})}(0, m), \\
\mathcal{U}\left(B_{m}\right) & :=\left\{\mathcal{U}_{t}\left(z_{0}, u\right): \text { for all }(t, u) \in B_{m}\right\} .
\end{aligned}
$$

We have

$$
\begin{align*}
\mathbb{R} \times W_{\text {loc }}^{1,1}\left(\mathbb{R}_{+}, \mathbb{R}\right) & =\bigcup_{m=1}^{\infty} B_{m} \\
\mathcal{A}\left(z_{0}\right) & =\bigcup_{m=1}^{\infty} \mathcal{U}\left(B_{m}\right) . \tag{4.10}
\end{align*}
$$

Combining (4.9), (4.10) and the Baire lemma, we see that there is a ball $Q \subset$ $H_{(V)}^{k}$ and an integer $m \geq 1$ such that $\mathcal{U}\left(B_{m}\right)$ is dense in $Q \cap S$ with respect to $H^{k}$-norm.

Step 1. Let us define the set

$$
\tilde{B}_{m}=\left\{(t, u) \in B_{m}: \text { such that } \mathcal{U}_{t}\left(z_{0}, u\right) \in Q\right\}
$$

Here we prove that $\tilde{B}_{m}$ is compact in $[0, m] \times L^{1}([0, m], \mathbb{R})$. Indeed, take any sequence $\left(t_{n}, u_{n}\right) \in \tilde{B}_{m}$. As $\left(t_{n}, u_{n}\right) \in B_{m}$ and $B_{m}$ is compact in $[0, m] \times$ $L^{1}([0, m], \mathbb{R})$, there is a sequence $n_{k} \rightarrow \infty$ and $\left(t_{0}, u_{0}\right) \in B_{m}$ such that

$$
\left|t_{n_{k}}-t_{0}\right|+\left\|u_{n_{k}}-u_{0}\right\|_{L^{1}([0, m], \mathbb{R})} \rightarrow 0, k \rightarrow \infty
$$

We need to show that $\left(t_{0}, u_{0}\right) \in \tilde{B}_{m}$. As $\mathcal{U}_{t_{n_{k}}}\left(z_{0}, u_{n_{k}}\right) \in Q$, there is $z \in Q$ such that $\mathcal{U}_{t_{n_{k}}}\left(z_{0}, u_{n_{k}}\right) \rightharpoonup z$ in $H^{k}$ (again extracting a subsequence, if necessary). On the other hand, Lemma 4.1 implies that $\mathcal{U}_{t_{n_{k}}}\left(z_{0}, u_{n_{k}}\right) \rightarrow \mathcal{U}_{t_{0}}\left(z_{0}, u_{0}\right)$ in $L^{2}$. Thus $\mathcal{U}_{t_{0}}\left(z_{0}, u_{0}\right)=z$ and $\left(t_{0}, u_{0}\right) \in \tilde{B}_{m}$. Thus $\tilde{B}_{m}$ is compact in $[0, m] \times L^{1}([0, m], \mathbb{R})$.

In particular, this implies that $\mathcal{U}\left(\tilde{B}_{m}\right)$ is compact in $L^{2}$, as an image of a compact set by a continuous mapping. On the other hand, $\mathcal{U}\left(\tilde{B}_{m}\right)$ is dense in the compact set $Q \cap S$ in $L^{2}$. Thus $Q \cap S=\mathcal{U}\left(\tilde{B}_{m}\right)$.

Step 2. Using standard arguments, one can show that we have

$$
\left\|\mathcal{U}_{t}\left(z_{0}, u\right)-\mathcal{U}_{t^{\prime}}\left(z_{0}, u^{\prime}\right)\right\| \leq C\left(\left|t-t^{\prime}\right|+\left\|u-u^{\prime}\right\|_{L^{1}([0, m], \mathbb{R})}\right)
$$

for any $(t, u),\left(t^{\prime}, u^{\prime}\right) \in \tilde{B}_{m}$, where $C>0$ is a constant not depending on $(t, u)$ and $\left(t^{\prime}, u^{\prime}\right)$. Combining this with the interpolation inequality

$$
\|z\|_{k-1} \leq C\|z\|^{\frac{1}{k}}\|z\|_{k}^{\frac{k-1}{k}}
$$

we get

$$
\left\|\mathcal{U}_{t}\left(z_{0}, u\right)-\mathcal{U}_{t^{\prime}}\left(z_{0}, u^{\prime}\right)\right\|_{k-1} \leq C\left(\left|t-t^{\prime}\right|^{\frac{1}{k}}+\left\|u-u^{\prime}\right\|_{L^{1}([0, m], \mathbb{R})}^{\frac{1}{k}}\right)
$$

for any $(t, u),\left(t^{\prime}, u^{\prime}\right) \in \tilde{B}_{m}$. Here we used the fact that $\mathcal{U}_{t}\left(z_{0}, u\right), \mathcal{U}_{t^{\prime}}\left(z_{0}, u^{\prime}\right) \in Q$. Appying Lemmas 4.3 and 4.4 and of the fact that $Q \cap S \subset \mathcal{U}\left(\tilde{B}_{m}\right)$, we obtain

$$
\begin{aligned}
H_{\varepsilon}\left(Q \cap S, H^{k-1}\right) & \leq H_{\varepsilon}\left(\mathcal{U}\left(\tilde{B}_{m}\right), H^{k-1}\right) \leq C H_{\varepsilon^{k}}\left(\tilde{B}_{m},[0, m] \times L^{1}([0, m], \mathbb{R})\right) \\
& \leq C H_{\varepsilon^{k}}\left(B_{m},[0, m] \times L^{1}([0, m], \mathbb{R})\right) \\
& \leq \frac{C}{\varepsilon^{k}} \ln \frac{1}{\varepsilon^{k}}
\end{aligned}
$$

This estimate contradicts Lemma 4.5 and proves the theorem.
Remark 4.6. The same proof works also in the case of Schrödinger equation with any finite number of controls:

$$
i \dot{z}=-\Delta z+V(x) z+u_{1}(t) Q_{1}(x) z+\ldots+u_{n}(t) Q_{n}(x) z
$$

where $n \geq 1$ is any integer, $Q_{j} \in C^{\infty}(\bar{D}, \mathbb{R})$ are arbitrary functions and $u_{j}$ are the controls $j=1, \ldots, n$.

## 5 Appendix

### 5.1 Genericity of Condition 2.5

Let us assume that $D=(0,1)^{d}$ and introduce the space

$$
\begin{aligned}
& \mathcal{G}:=\left\{V \in C^{\infty}(D, \mathbb{R}): V\left(x_{1}, \ldots, x_{d}\right)=V_{1}\left(x_{1}\right)+\ldots+V_{d}\left(x_{d}\right)\right. \\
&\text { for some } \left.V_{k} \in C^{\infty}([0,1], \mathbb{R}), k=1, \ldots, d\right\} .
\end{aligned}
$$

Then $\mathcal{G}$, endowed with the metric of $C^{\infty}(\bar{D}, \mathbb{R})$, is a closed subspace in $C^{\infty}(\bar{D}, \mathbb{R})$. By Lemma 3.12 in [22], the set $\mathcal{A}$ of all functions $V \in \mathcal{G}$ such that property (ii) in Condition 2.5 is verified is $G_{\delta}$ set (i.e., countable intersection of dense open sets). First let us prove genericity of property (i) in the case $d=1$.

Lemma 5.1. For any $V \in C^{\infty}([0,1], \mathbb{R})$, the set of functions $Q \in C^{\infty}([0,1], \mathbb{R})$ such that

$$
\begin{equation*}
\inf _{p, j \geq 1}\left|p^{3} j^{3}\left\langle Q e_{p, V}, e_{j, V}\right\rangle\right|>0 \tag{5.1}
\end{equation*}
$$

is dense in $C^{\infty}([0,1], \mathbb{R})$.
Proof. If $V=0$, then a straightforward calculation gives

$$
\left\langle x^{2} e_{p, 0}, e_{j, 0}\right\rangle= \begin{cases}\frac{(-1)^{p+j} 8 p j}{\pi^{2}\left(p^{2}-j^{2}\right)^{2}}, & \text { if } p \neq j, \\ \frac{2}{3}-\frac{1}{p^{2} \pi^{2}}, & \text { if } p=j,\end{cases}
$$

which implies (5.1) for $Q=x^{2}$ and $V=0$. In the general case, taking any $p \neq j$, we integrate by parts (we write $\lambda_{j}, e_{j}$ and $z^{\prime \prime}, z^{\prime}$ instead of $\lambda_{j, V}, e_{j, V}$ and $\frac{\mathrm{d}^{2} z}{\mathrm{~d} x^{2}}, \frac{\mathrm{~d} z}{\mathrm{~d} x}$, respectively)

$$
\begin{aligned}
\left\langle Q e_{p}, e_{j}\right\rangle & =\frac{1}{\lambda_{j}}\left\langle\left(-\frac{\mathrm{d}^{2}}{\mathrm{~d} x^{2}}+V\right)\left(Q e_{p}\right), e_{j}\right\rangle \\
& =\frac{1}{\lambda_{j}}\left(\left\langle-Q^{\prime \prime} e_{p}, e_{j}\right\rangle+\left\langle-Q^{\prime} e_{p}^{\prime}, e_{j}\right\rangle+\lambda_{p}\left\langle Q e_{p}, e_{j}\right\rangle\right) .
\end{aligned}
$$

This implies that

$$
\begin{equation*}
\left\langle Q e_{p}, e_{j}\right\rangle=-\frac{1}{\lambda_{j}-\lambda_{p}}\left(\left\langle Q^{\prime \prime} e_{p}, e_{j}\right\rangle+\left\langle Q^{\prime} e_{p}^{\prime}, e_{j}\right\rangle\right) . \tag{5.2}
\end{equation*}
$$

Again integrating by parts, we get

$$
\begin{align*}
\left\langle Q^{\prime} e_{p}^{\prime}, e_{j}\right\rangle & =\frac{1}{\lambda_{j}}\left\langle Q^{\prime} e_{p}^{\prime},\left(-\frac{\mathrm{d}^{2}}{\mathrm{~d} x^{2}}+V\right) e_{j}\right\rangle \\
& =-\left.\frac{1}{\lambda_{j}} Q^{\prime} e_{p}^{\prime} e_{j}^{\prime}\right|_{x=0} ^{x=1}+\frac{1}{\lambda_{j}}\left\langle\left(-\frac{\mathrm{d}^{2}}{\mathrm{~d} x^{2}}+V\right)\left(Q^{\prime} e_{p}^{\prime}, e_{j}\right\rangle\right. \tag{5.3}
\end{align*}
$$

Notice that

$$
\begin{aligned}
\left\langle\left(-\frac{\mathrm{d}^{2}}{\mathrm{~d} x^{2}}+V\right)\left(Q^{\prime} e_{p}^{\prime}\right), e_{j}\right\rangle= & \left\langle V Q^{\prime} e_{p}^{\prime}, e_{j}\right\rangle+\left\langle-Q^{\prime \prime \prime} e_{p}^{\prime}, e_{j}\right\rangle+\left\langle-Q^{\prime \prime} e_{p}^{\prime \prime}, e_{j}\right\rangle \\
& +\lambda_{p}\left\langle Q^{\prime} e_{p}^{\prime}, e_{j}\right\rangle-\left\langle Q^{\prime}\left(V e_{p}\right)^{\prime}, e_{j}\right\rangle
\end{aligned}
$$

Replacing this into (5.3), we get

$$
\begin{align*}
\left\langle Q^{\prime} e_{p}^{\prime}, e_{j}\right\rangle= & \frac{1}{\lambda_{j}-\lambda_{p}}\left(-\left.Q^{\prime} e_{p}^{\prime} e_{j}^{\prime}\right|_{x=0} ^{x=1}+\left\langle V Q^{\prime} e_{p}^{\prime}, e_{j}\right\rangle+\left\langle-Q^{\prime \prime \prime} e_{p}^{\prime}, e_{j}\right\rangle\right. \\
& \left.+\left\langle-Q^{\prime \prime} e_{p}^{\prime \prime}, e_{j}\right\rangle-\left\langle Q^{\prime}\left(V e_{p}\right)^{\prime}, e_{j}\right\rangle\right) . \tag{5.4}
\end{align*}
$$

Using (5.2) and (5.4) and the fact that

$$
\left\langle-Q^{\prime \prime} e_{p}^{\prime \prime}, e_{j}\right\rangle=-\left\langle Q^{\prime \prime} V e_{p}, e_{j}\right\rangle+\lambda_{p}\left\langle Q^{\prime \prime} e_{p}, e_{j}\right\rangle,
$$

we obtain

$$
\begin{aligned}
\left\langle Q e_{p}, e_{j}\right\rangle= & \left(-\frac{1}{\lambda_{j}-\lambda_{p}}\left\langle Q^{\prime \prime} e_{p}, e_{j}\right\rangle-\frac{\lambda_{p}}{\left(\lambda_{j}-\lambda_{p}\right)^{2}}\left\langle Q^{\prime \prime} e_{p}, e_{j}\right\rangle\right) \\
& -\frac{1}{\left(\lambda_{j}-\lambda_{p}\right)^{2}}\left(-\left.Q^{\prime} e_{p}^{\prime} e_{j}^{\prime}\right|_{x=0} ^{x=1}+\left\langle V Q^{\prime} e_{p}^{\prime}, e_{j}\right\rangle\right. \\
& \left.+\left\langle-Q^{\prime \prime \prime} e_{p}^{\prime}, e_{j}\right\rangle-\left\langle Q^{\prime \prime} V e_{p}, e_{j}\right\rangle-\left\langle Q^{\prime}\left(V e_{p}\right)^{\prime}, e_{j}\right\rangle\right) \\
= & I_{1}+I_{2} .
\end{aligned}
$$

Let $Q$ be such that $A:=\left.Q^{\prime}(x) \cos (p \pi x) \cos (j \pi x)\right|_{x=0} ^{x=1} \neq 0$. Clearly, this is verified for almost any $Q$, since $A$ depends only on the parity of $p$ and $j$. Let
us choose $Q$ such that $\left\langle Q e_{p}, e_{j}\right\rangle \neq 0$ for all $p, j \geq 1$; the set of such functions $Q$ is $G_{\delta}$, by Section 3.4 in (22]. Using the estimates (1.6)-(1.8), it is easy to see that $\inf _{p, j \geq 1, p \neq j}\left|p^{3} j^{3} I_{2}\right|>0$. Iterating the same arguments for $I_{1}$, we see that $\inf _{p, j \geq 1, p \neq j}\left|p^{3} j^{3}\left\langle Q e_{p, V}, e_{j, V}\right\rangle\right|>0$ for almost any polynomial $Q$.

If $p=j$, using (1.7), we get

$$
\left\langle Q e_{p}, e_{p}\right\rangle=2\left\langle Q, \sin ^{2}(p \pi x)\right\rangle+s_{p}
$$

where $s_{p} \rightarrow 0$. Thus

$$
\left\langle Q e_{p}, e_{p}\right\rangle=\langle Q, 1-\cos 2 p \pi x\rangle+s_{p}=\int_{0}^{1} Q \mathrm{~d} x-\langle Q, \cos 2 p \pi x\rangle+s_{p}
$$

Taking $Q$ such that $\int_{0}^{1} Q \mathrm{~d} x \neq 0$, we complete the proof of the lemma.

Take any functions $Q_{k} \in C^{\infty}([0,1], \mathbb{R}), k=1, \ldots, d$ in the dense set of Lemma 5.1 corresponding to $V_{k} \in C^{\infty}([0,1], \mathbb{R}), k=1, \ldots, d$. Then $Q\left(x_{1}, \ldots, x_{d}\right)$ $:=Q_{1}\left(x_{1}\right) \cdot \ldots \cdot Q_{d}\left(x_{d}\right)$ satisfies (i) with $V\left(x_{1}, \ldots, x_{d}\right):=V_{1}\left(x_{1}\right)+\ldots+V_{d}\left(x_{d}\right)$.

### 5.2 Inverse mapping theorem for multifunctions

In this section, we recall the statement of the inverse mapping theorem for multivalued functions or multifunctions. We refer the reader to the paper (21] by Nachi and Penot for details and for a review of the literature on this subject.

Let $X$ and $Y$ be Banach spaces. For any non-empty sets $C, D \subset X$, define the Hausdorff distance

$$
\begin{aligned}
& d(x, D)=\inf _{y \in D}\|x-y\|_{X}, \\
& e(C, D)=\sup _{x \in C} d(x, D) .
\end{aligned}
$$

We call a multifunction from $X$ to $Y$ any mapping $F$ from $X$ to $2^{Y}$.
Definition 5.2. A multifunction $F$ from an open set $X_{0} \subset X$ to $Y$ is said to be strictly differentiable at $\left(x_{0}, y_{0}\right)$ if there exists some continuous linear map $A: X \rightarrow Y$ such that for any $\varepsilon>0$ there exist $\beta, \delta>0$ for which

$$
e\left(F(x) \cap B_{Y}\left(y_{0}, \beta\right)-A(x), F\left(x^{\prime}\right)-A\left(x^{\prime}\right)\right) \leq \varepsilon\left\|x-x^{\prime}\right\|_{X}
$$

whenever $x, x^{\prime} \in B\left(x_{0}, \delta\right)$. The map $A$ is called a derivative of $F$ at $\left(x_{0}, y_{0}\right)$.
The following theorem is a generalization of the classical inverse function theorem to the case of multifunctions.

Theorem 5.3. Let $F$ be a multifunction from an open set $X_{0} \subset X$ to $Y$ with closed non-empty values. Suppose $F$ is strictly differentiable at $\left(x_{0}, y_{0}\right) \in$ $G r(F)$, and some derivative $A$ of $F$ at $\left(x_{0}, y_{0}\right)$ has a right inverse. Then for any neighborhood $U$ of $x_{0}$ there exists a neighborhood $V$ of $y_{0}$ such that $V \subset F(U)$.

See Theorem 22 in 21 for the proof.

## References

[1] A. Agrachev and T. Chambrion. An estimation of the controllability time for single-input systems on compact Lie groups. J. ESAIM Control Optim. Calc. Var., 12(3):409-441, 2006.
[2] F. Albertini and D. D'Alessandro. Notions of controllability for bilinear multilevel quantum systems. IEEE Transactions on Automatic Control, 48(8):1399-1403, 2003.
[3] C. Altafini. Controllability of quantum mechanical systems by root space decomposition of su(n). J. of Math. Phys., 43(5):2051-2062, 2002.
[4] J. M. Ball, J. E. Marsden, and M. Slemrod. Controllability for distributed bilinear systems. SIAM J. Control Optim., 20(4):575-597, 1982.
[5] L. Baudouin and J.-P. Puel. Uniqueness and stability in an inverse problem for the Schrödinger equation. Inverse Problems, 18(6):1537-1554, 2001.
[6] K. Beauchard. Local controllability of a 1D Schrödinger equation. J. Math. Pures et Appl., 84(7):851-956, 2005.
[7] K. Beauchard, Y. Chitour, D. Kateb, and R. Long. Spectral controllability of 2D and 3D linear Schrödinger equations. J. Funct. Anal., 256:3916-3976, 2009.
[8] K. Beauchard and J.-M. Coron. Controllability of a quantum particle in a moving potential well. J. Funct. Anal., 232(2):328-389, 2006.
[9] K. Beauchard, J.-M. Coron, M. Mirrahimi, and P. Rouchon. Implicit Lyapunov control of finite dimensional Schrödinger equations. Systems and Control Letters, 56(5):388-395, 2007.
[10] K. Beauchard and C. Laurent. Local controllability of linear and nonlinear Schrödinger equations with bilinear control. Preprint, 2010.
[11] T. Cazenave. Semilinear Schrödinger equations. Courant Lecture Notes in Mathematics, AMS, 10, 2003.
[12] T. Chambrion, P. Mason, M. Sigalotti, and U. Boscain. Controllability of the discrete-spectrum Schrödinger equation driven by an external field. Annales de l'IHP, Analyse non linéaire, 26(1):329-349, 2009.
[13] B. Dehman, P. Gérard, and G. Lebeau. Stabilization and control for the nonlinear Schrödinger equation on a compact surface. Math. Z., 254(4):729-749, 2006.
[14] D. E. Edmunds and H. Triebel. Function Spaces, Entropy Numbers, Differential Operators. Cambridge University Press, Cambridge, 1996.
[15] S. Ervedoza and J.-P. Puel. Approximate controllability for a system of Schrödinger equations modeling a single trapped ion. Annales de l'IHP, Analyse non linéaire, 26(6):2111-2136, 2009.
[16] A. Katok and B. Hasselblatt. Introduction to the modern theory of dynamical systems. Encyclopedia of Mathematics and its Applications, 54, Cambridge University Press, Cambridge, 1995.
[17] G. Lebeau. Contrôle de l'équation de Schrödinger. J. Math. Pures Appl., 71(3):267-291, 1992.
[18] E. Machtyngier and E. Zuazua. Stabilization of the Schrödinger equation. Portugaliae Matematica, 51(2):243-256, 1994.
[19] P. Mason and M. Sigalotti. Generic controllability properties for the bilinear Schrödinger equation. Comm. in PDE, 35(4):685-706, 2010.
[20] M. Mirrahimi. Lyapunov control of a particle in a finite quantum potential well. IEEE Conf. on Decision and Control, San Diego, 2006.
[21] K. Nachi and J.-P. Penot. Inversion of multifunctions and differential inclusions. Control and Cybernetics, 34(3):871-901, 2005.
[22] V. Nersesyan. Growth of Sobolev norms and controllability of the Schrödinger equation. Comm. Math. Phys., 290(1):371-387, 2009.
[23] V. Nersesyan. Global approximate controllability for Schrödinger equation in higher Sobolev norms and applications. Annales de l'IHP, Analyse non linéaire, 27(3):901-915, 2010.
[24] H. Nersisyan. Controllability of 3D incompressible Euler equations by a finite-dimensional external force. ESAIM: COCV, (DOI: 10.1051/cocv/2009017), 2009.
[25] J. Pöschel and E. Trubowitz. Inverse Spectral Theory. Academic Press, New York, 1987.
[26] Y. Privat and M. Sigalotti. The squares of Laplacian-Dirichlet eigenfunctions are generically linearly independent. Preprint, 2008.
[27] V. Ramakrishna, M. Salapaka, M. Dahleh, H. Rabitz, and A. Pierce. Controllability of molecular systems. Phys. Rev. A, 51(2):960966, 1995.
[28] A. Shirikyan. Euler equations are not exactly controllable by a finitedimensional external force. Physica D, 237:1317-1323, 2008.
[29] G. Turinici. On the controllability of bilinear quantum systems. Lecture Notes in Chem., 74, 2000.
[30] G. Turinici and H. Rabitz. Quantum wavefunction controllability. Chem. Phys., 267(1):1-9, 2001.
[31] E. Zuazua. Remarks on the controllability of the Schrödinger equation. CRM Proc. Lecture Notes, 33:193-211, 2003.

